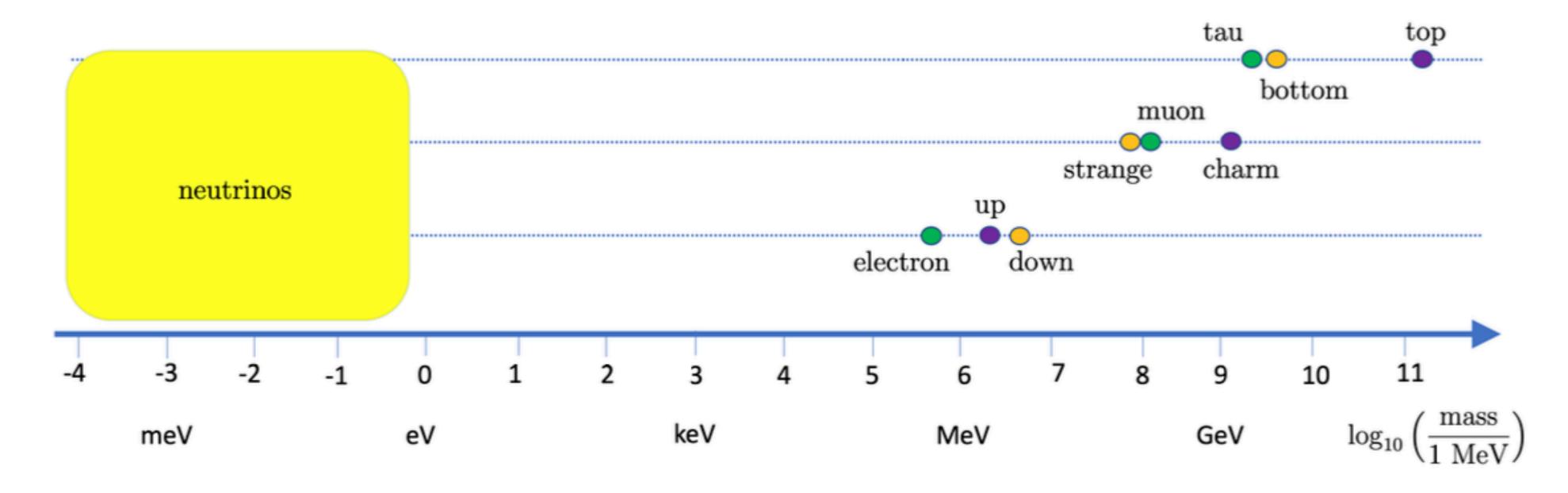
Theory of "toponium"

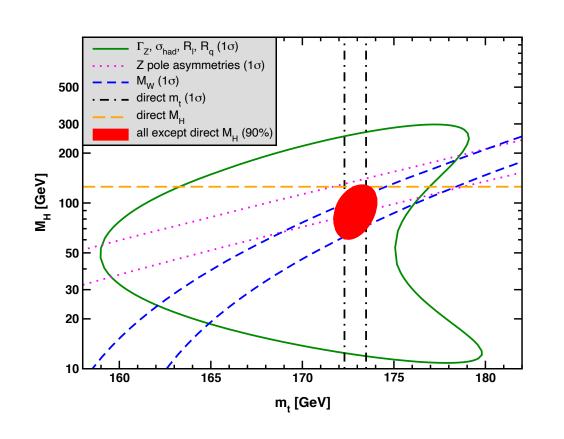
Li Lin Yang Zhejiang University



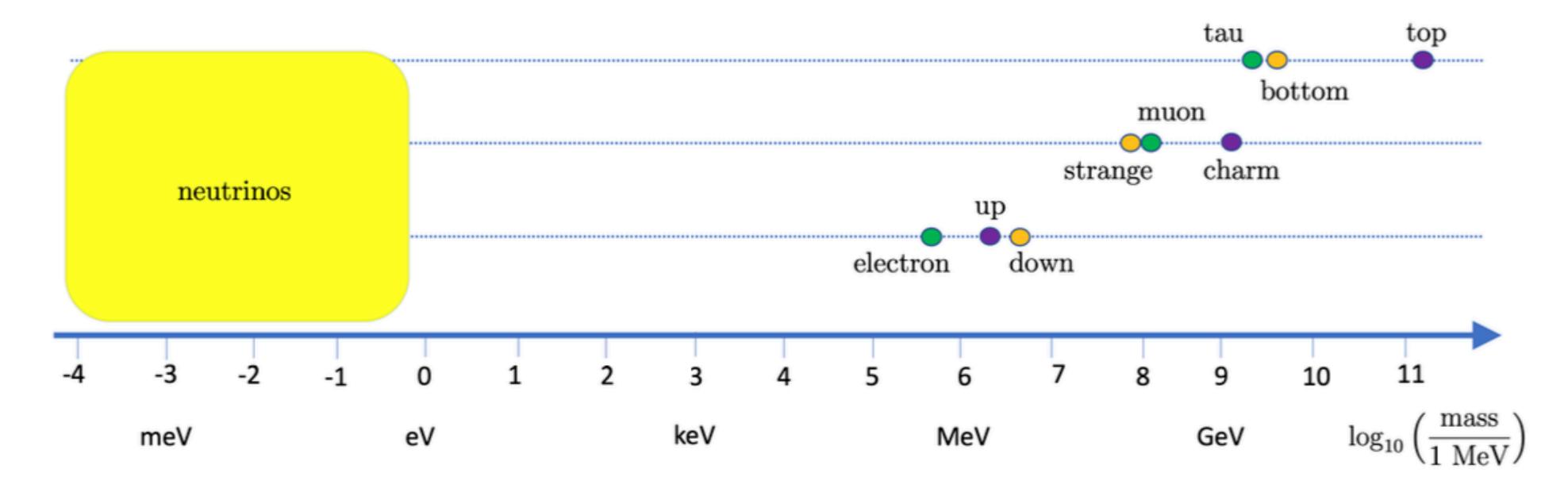
The top quark



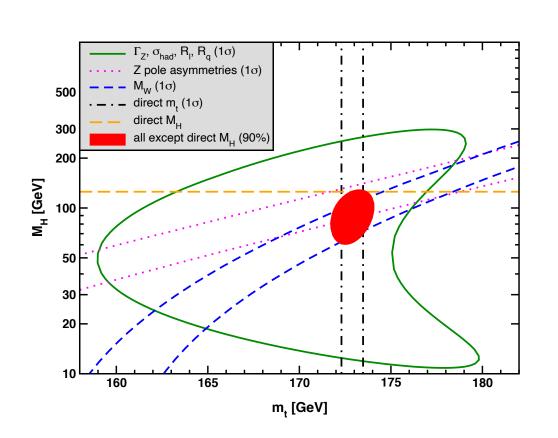
- ➤ The heaviest fundamental particle
- ➤ Large Yukawa coupling with the Higgs field
- ➤ Top quark mass: a crucial parameter of the SM
- ➤ Decay before hadronization: "bare quark"



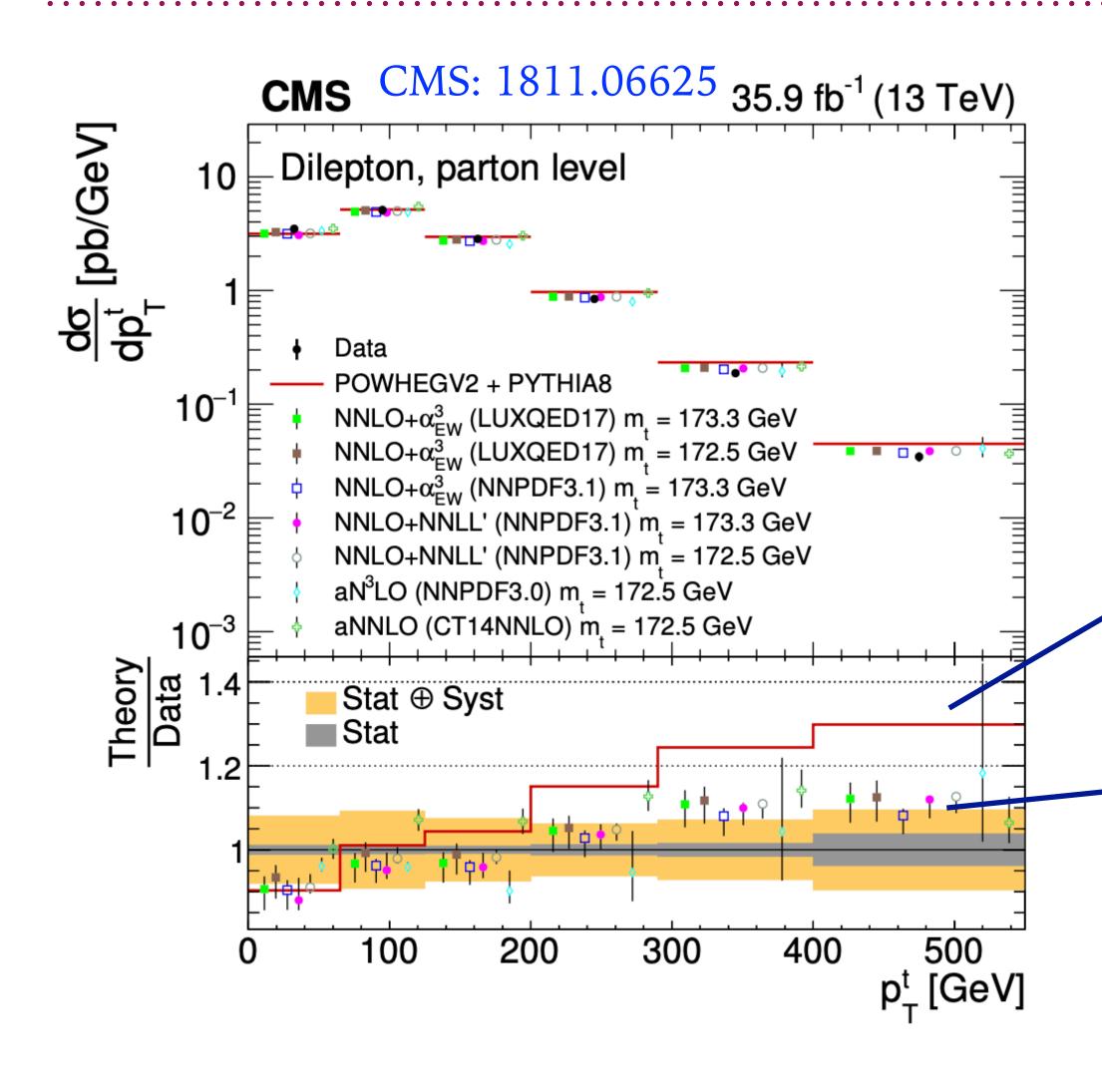
The top quark



- ➤ The heaviest fundamental particle
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Top quark pair production



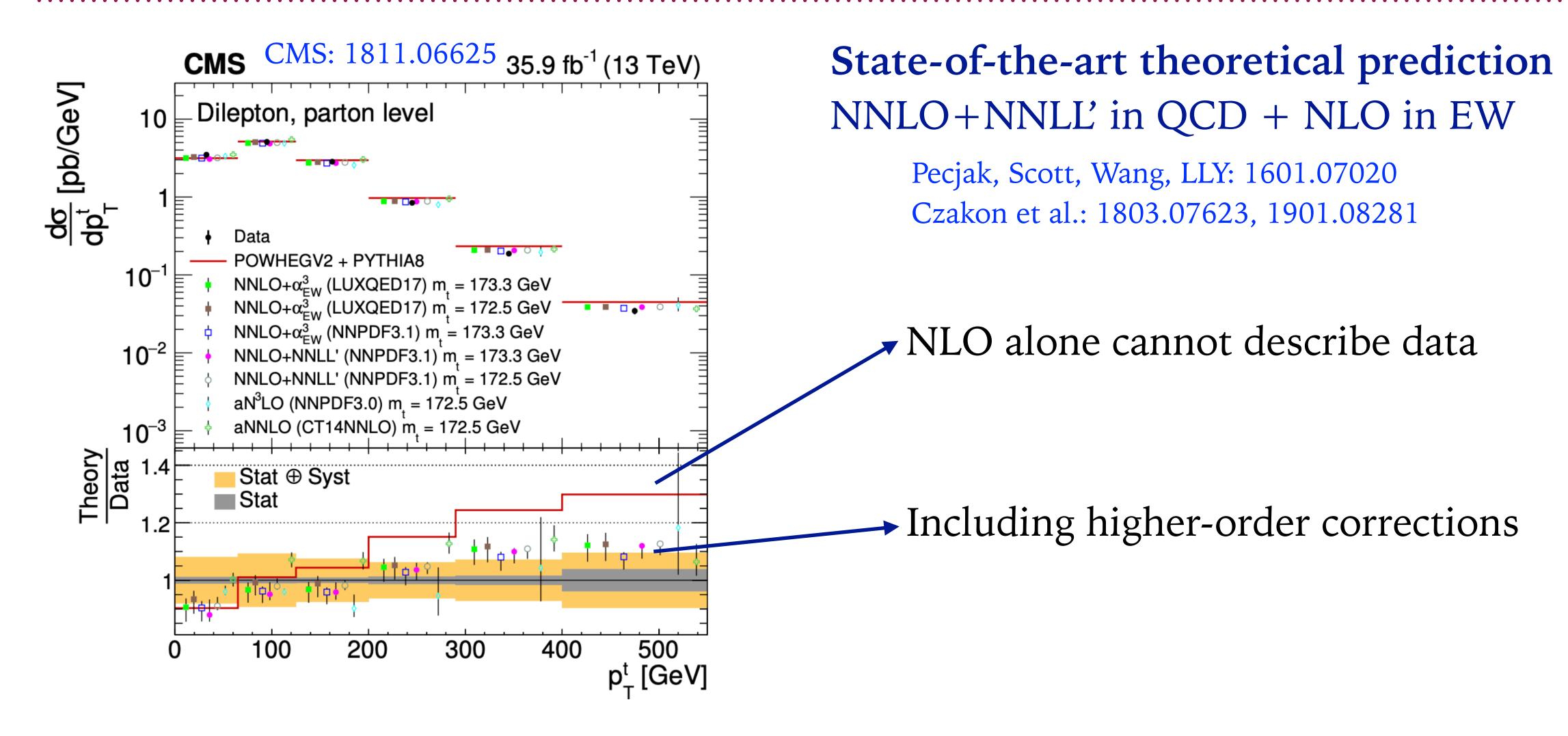
State-of-the-art theoretical prediction NNLO+NNLL' in QCD + NLO in EW

Pecjak, Scott, Wang, LLY: 1601.07020 Czakon et al.: 1803.07623, 1901.08281

NLO alone cannot describe data

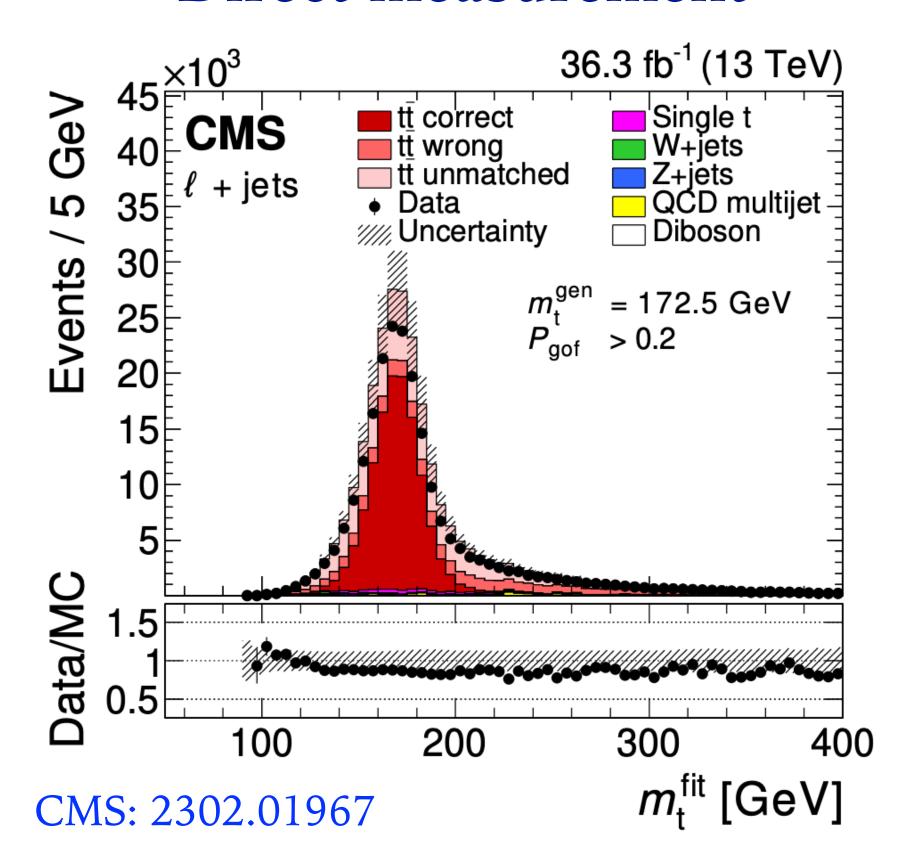
Including higher-order corrections

Top quark pair production



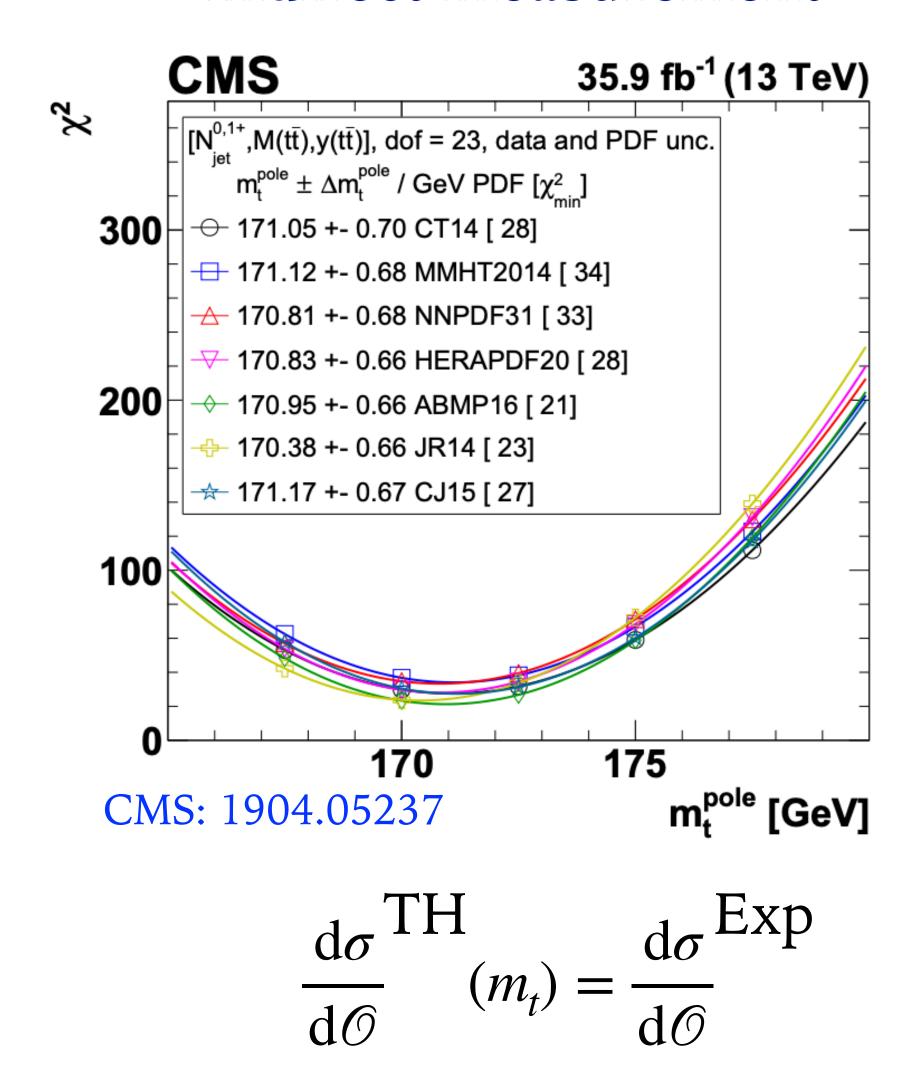
With such precision data and predictions, one would like to extract m_t from them...

Direct measurement



Reconstruct the invariant mass of decay products

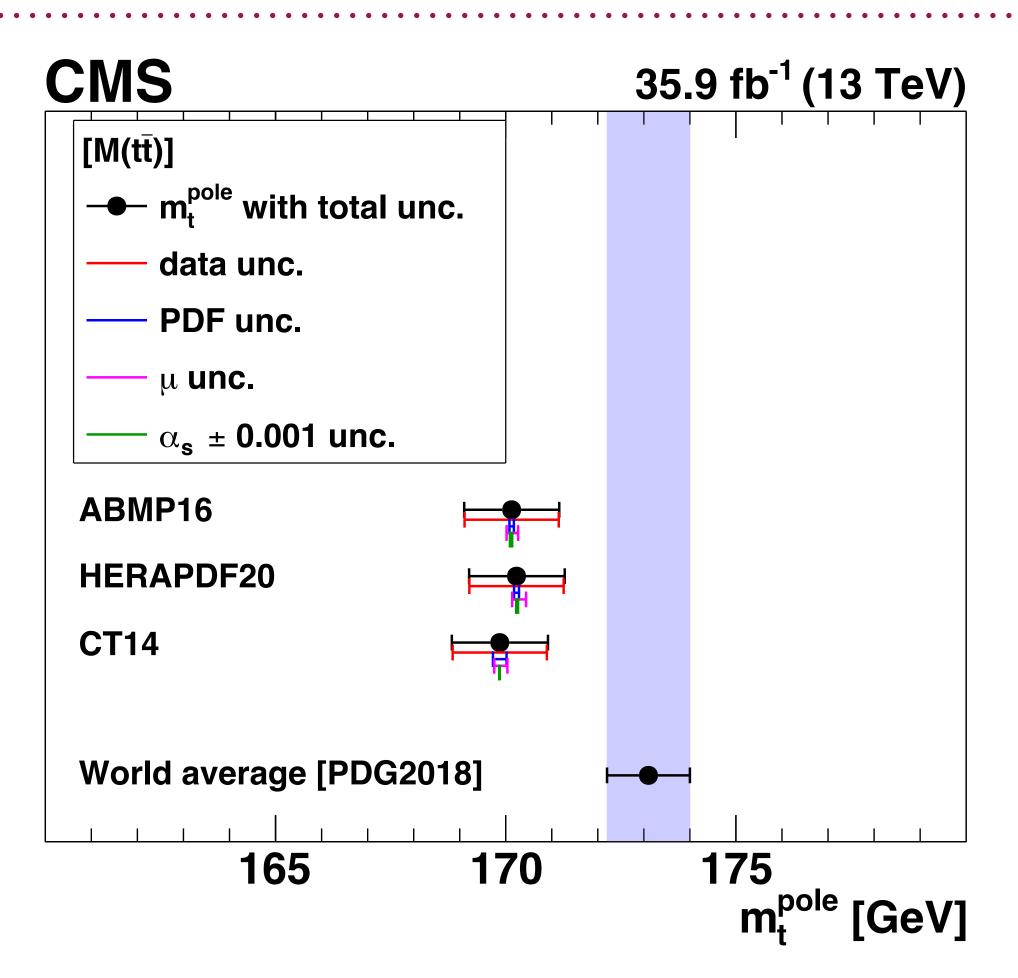
Indirect measurement



VS

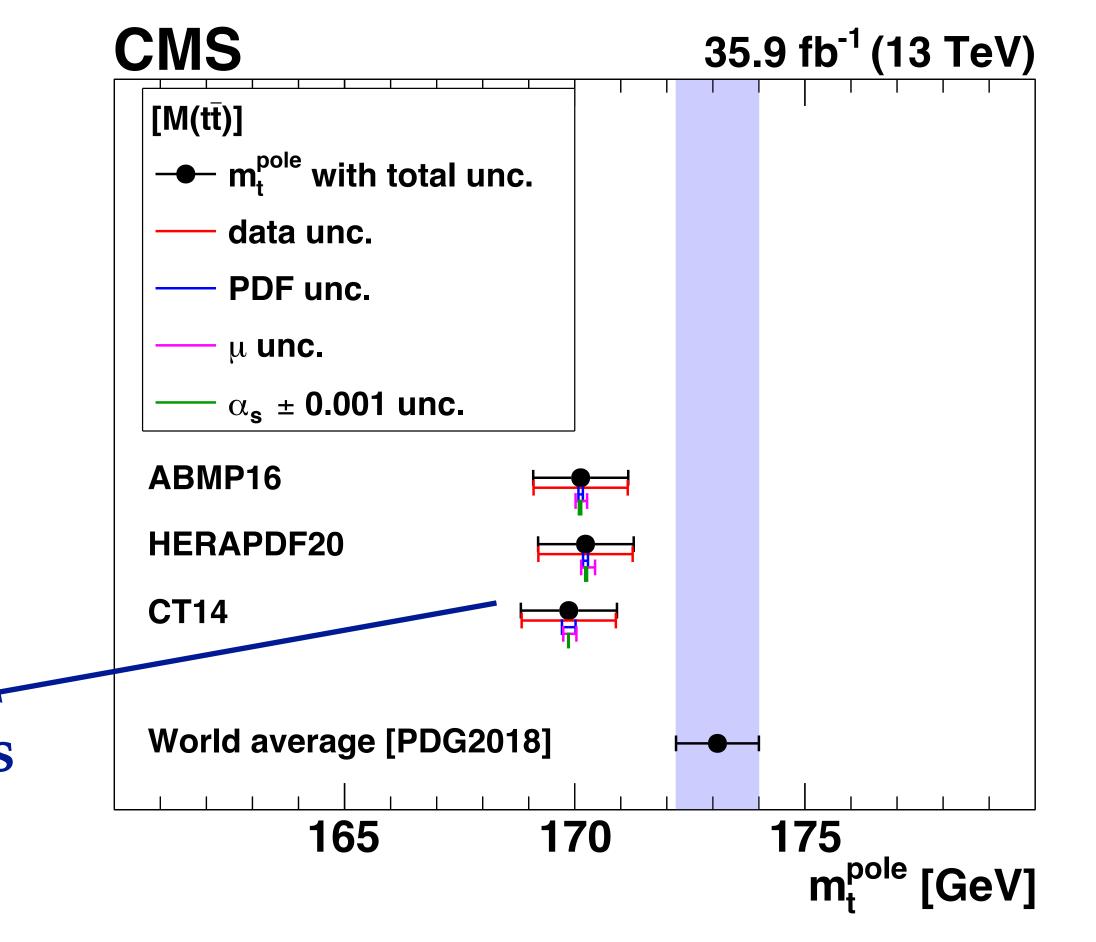
ATLAS: 1905.02302

CMS: 1904.05237

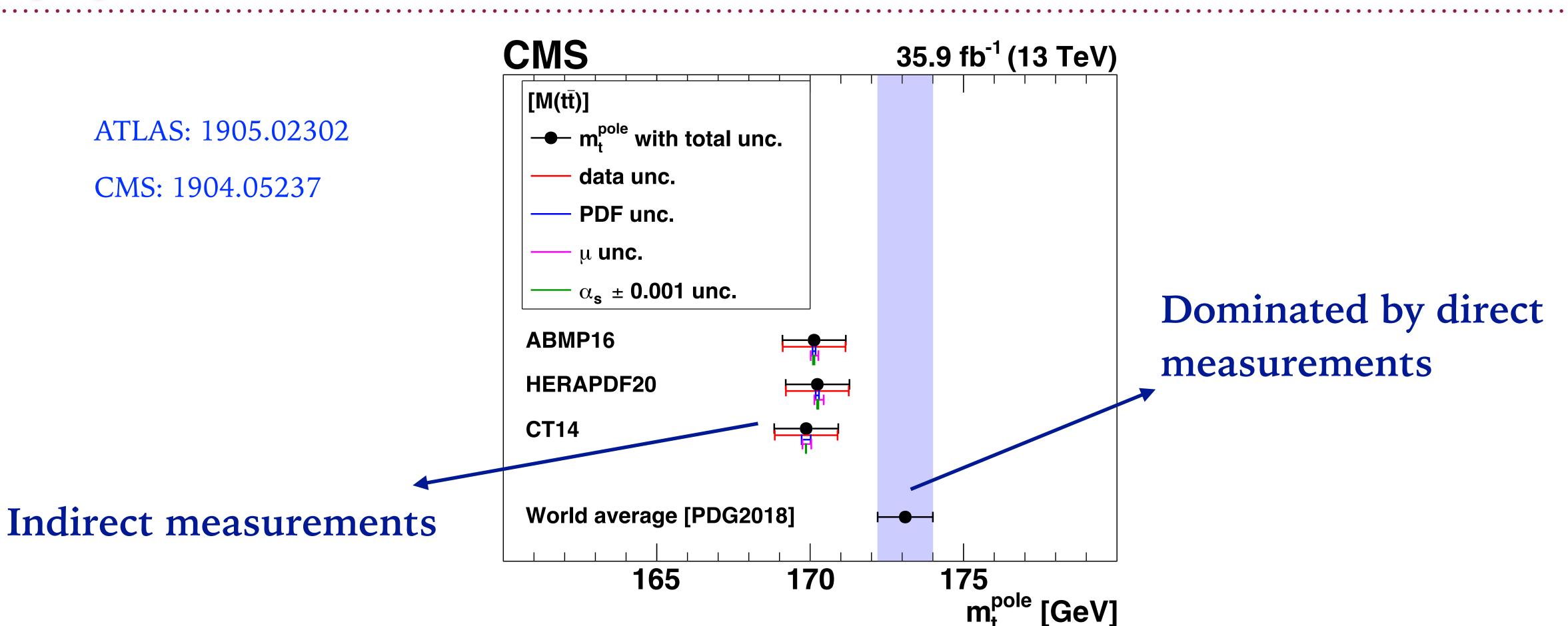


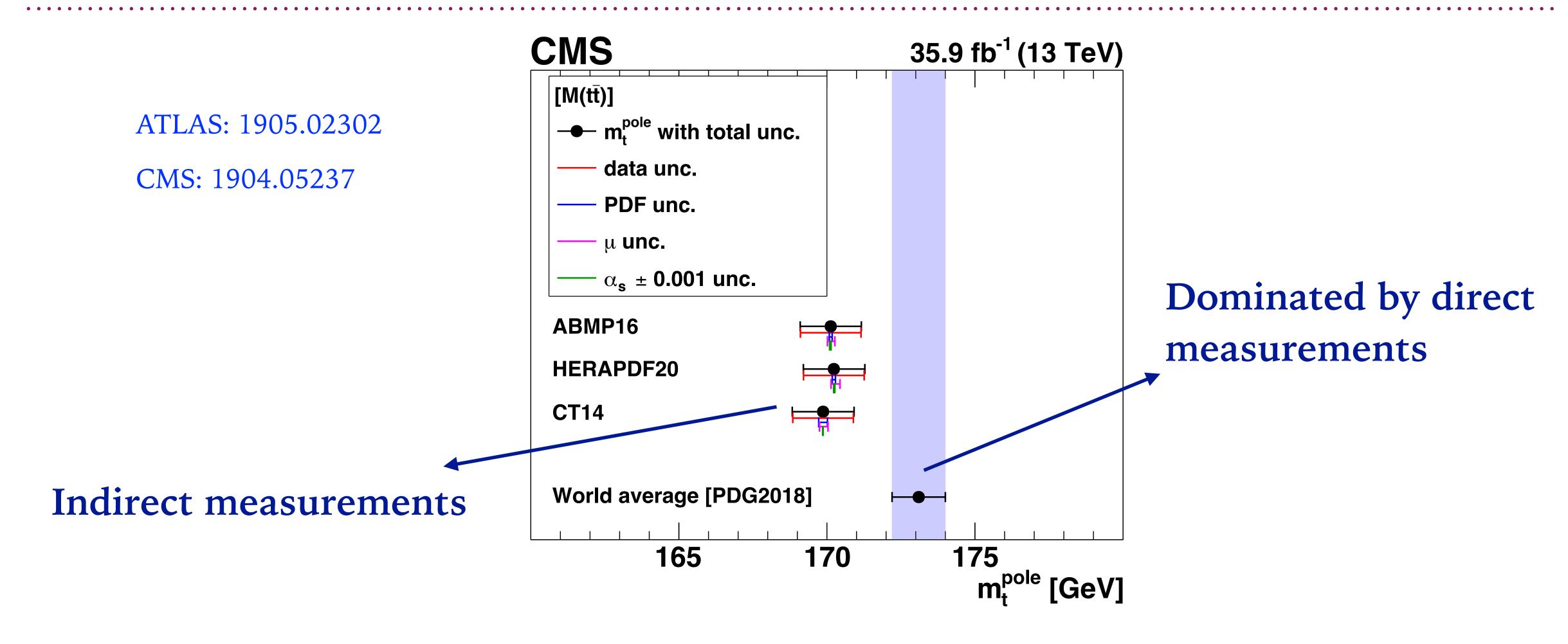


CMS: 1904.05237

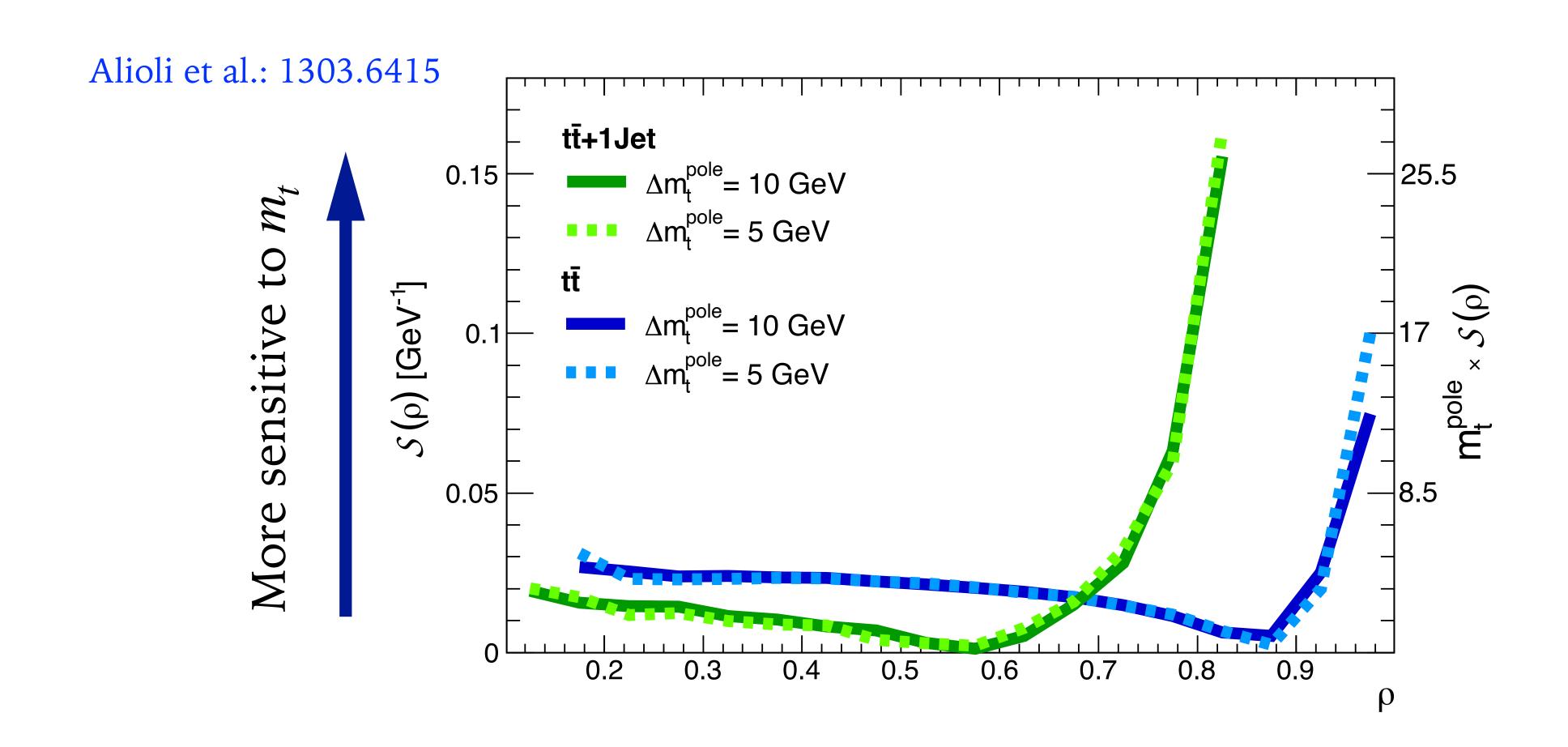


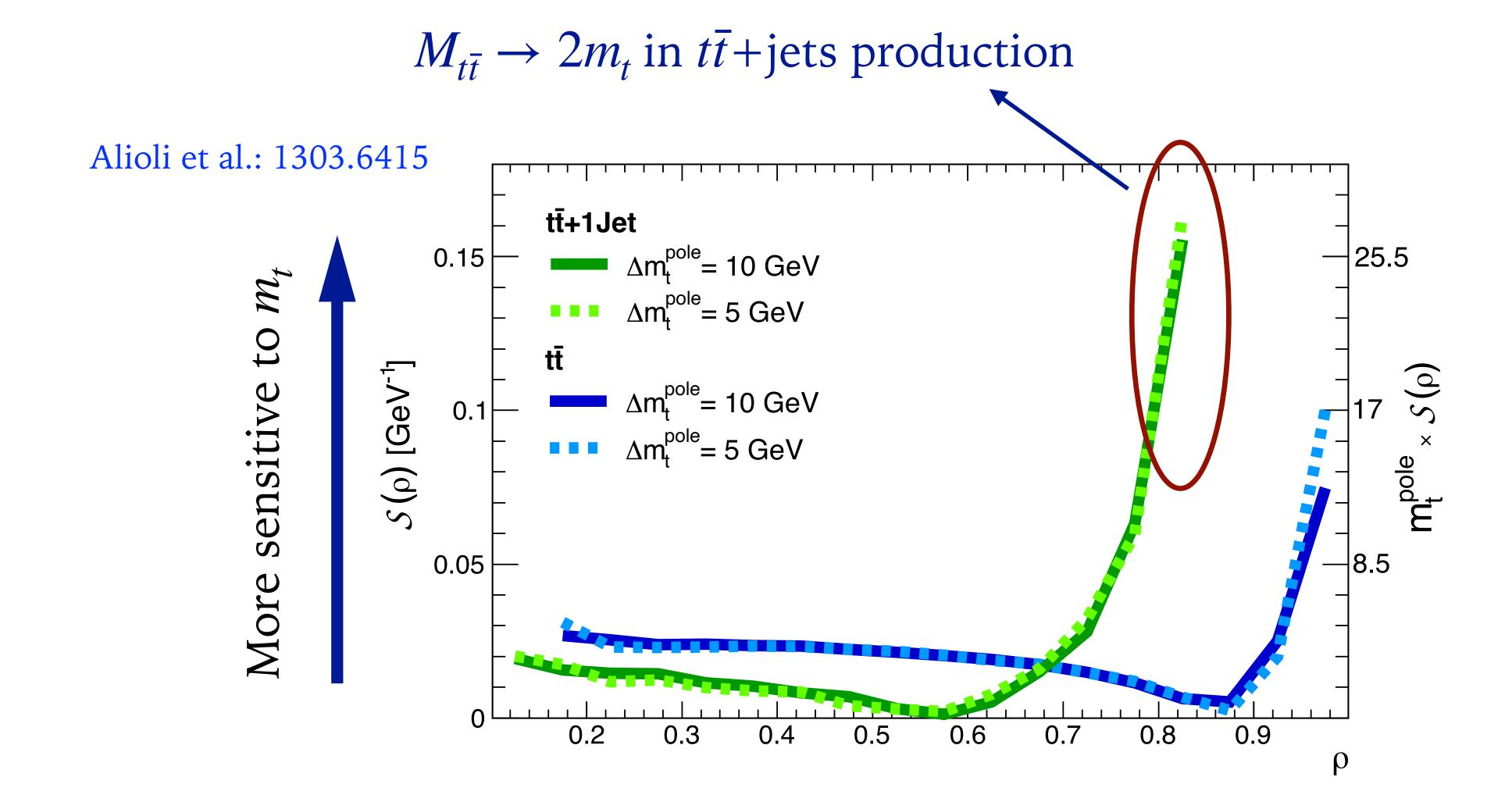
Indirect measurements

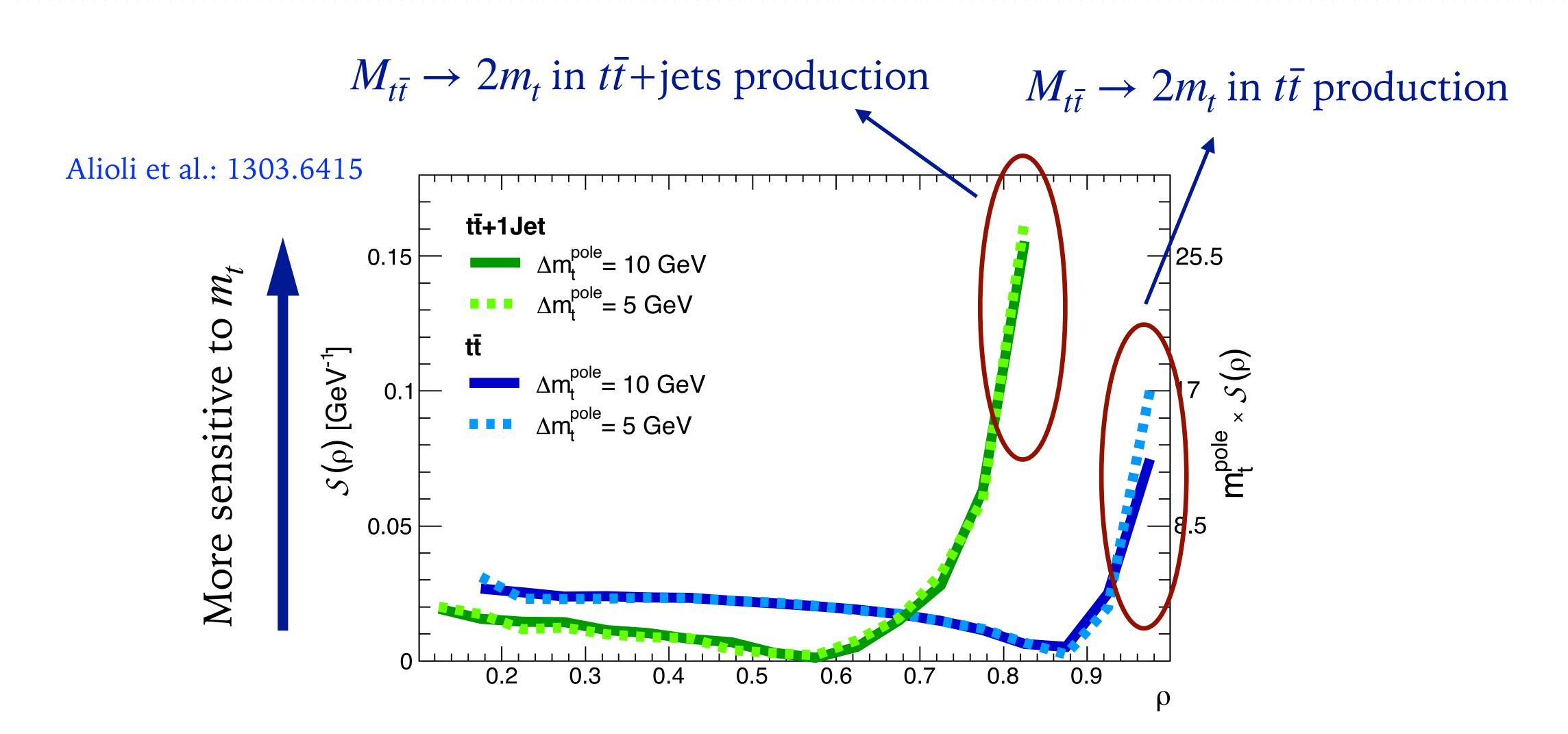


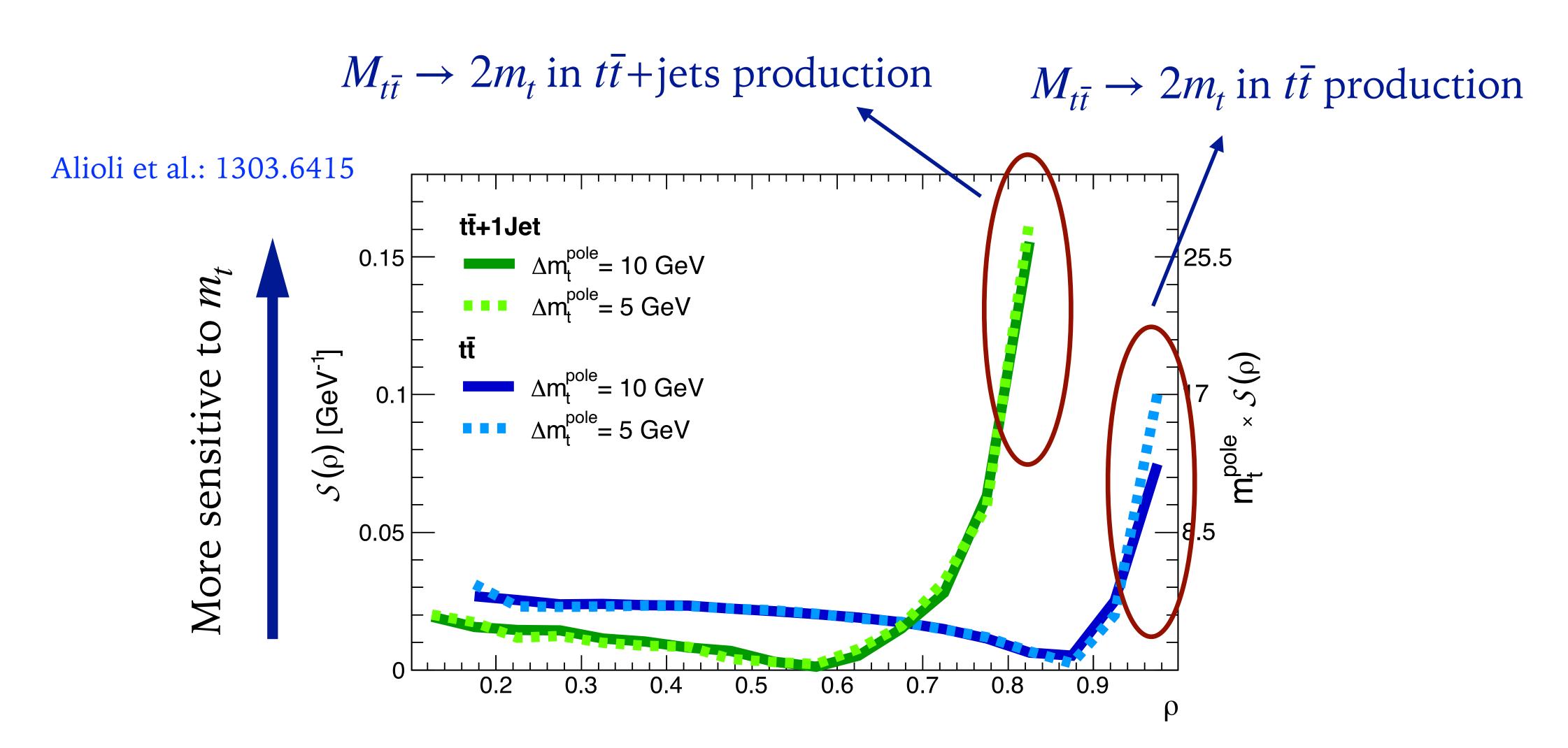


Tension! Why?

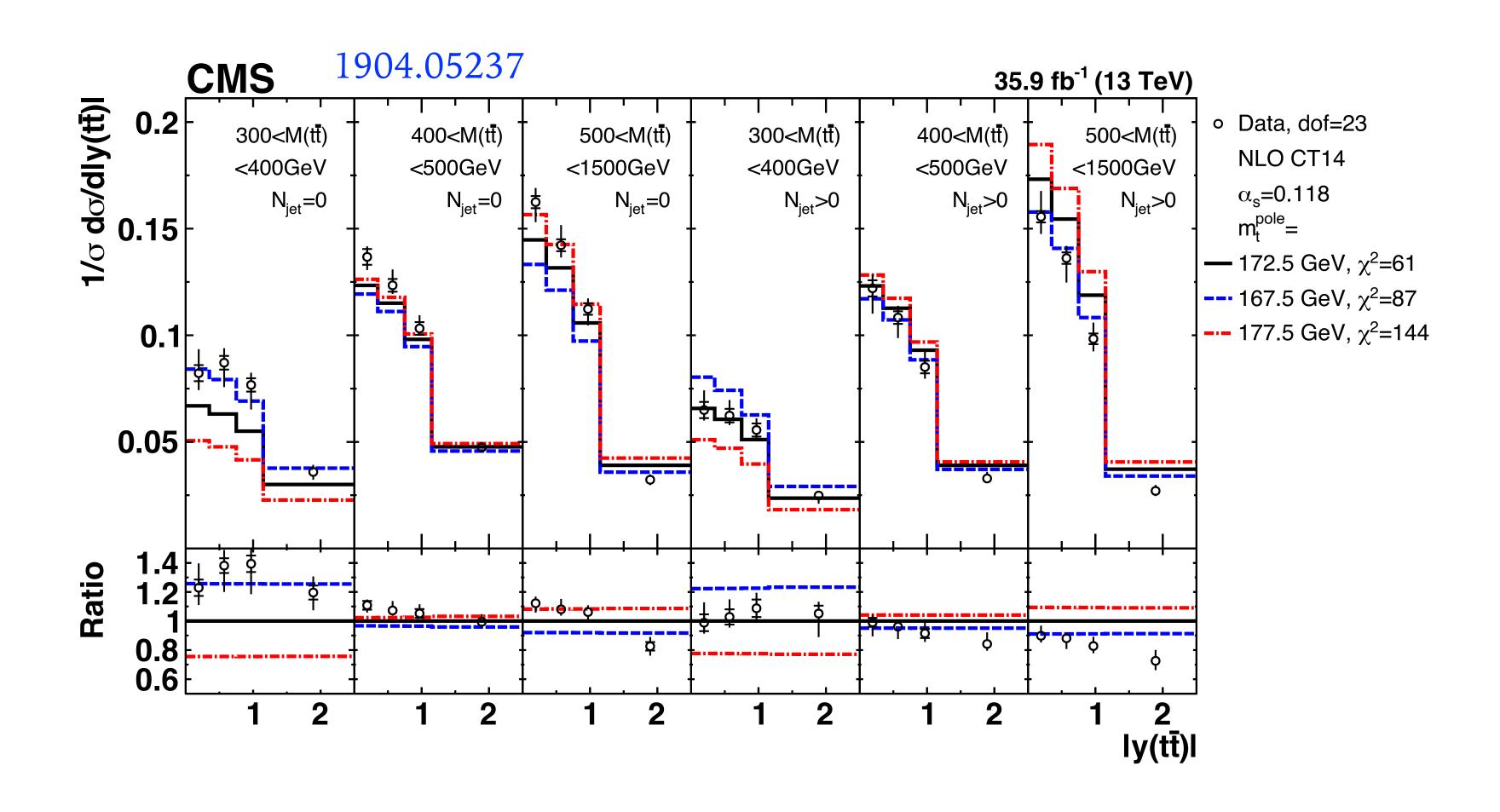


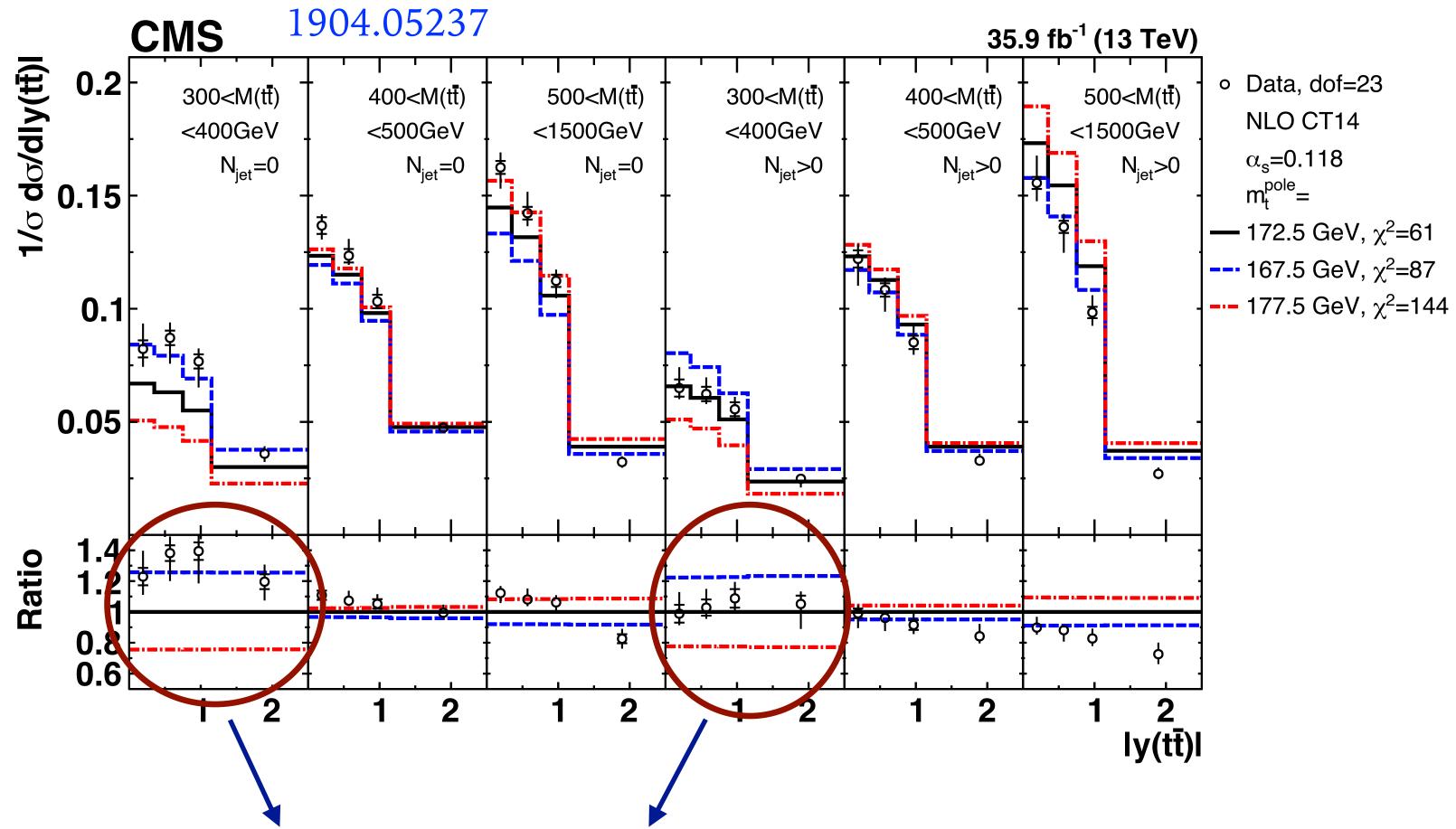




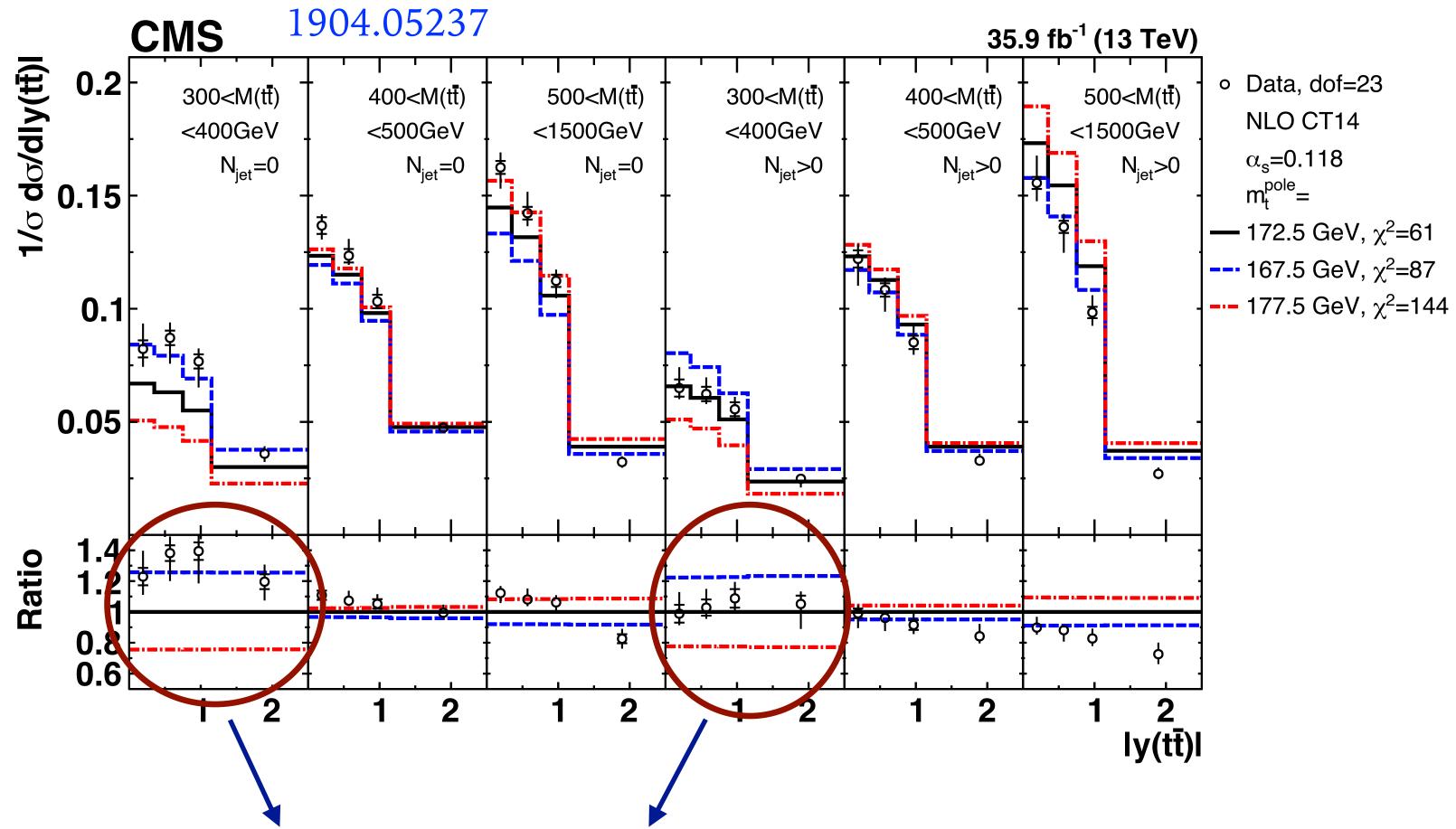


It is well-known that threshold region is most sensitive to m_t



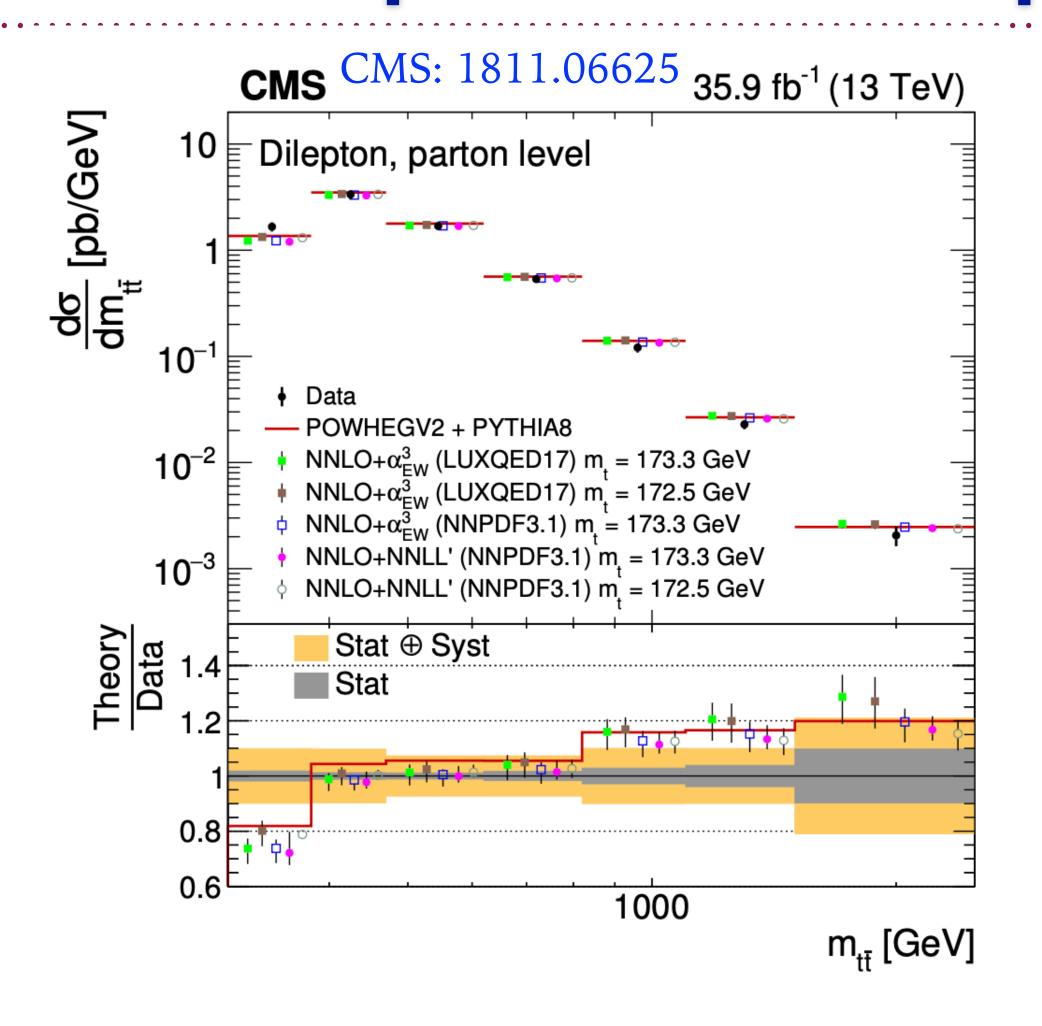


Data also shows that most of the mass-sensitivity comes from the low $M_{t\bar{t}}$ region

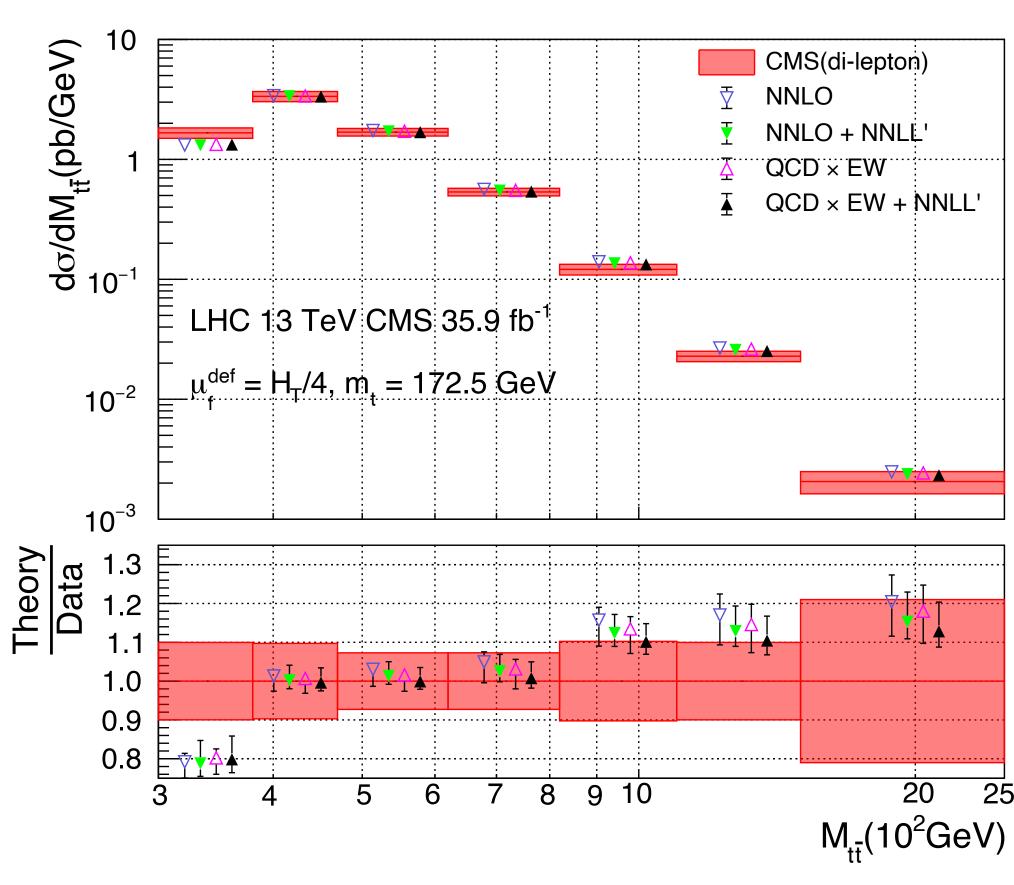


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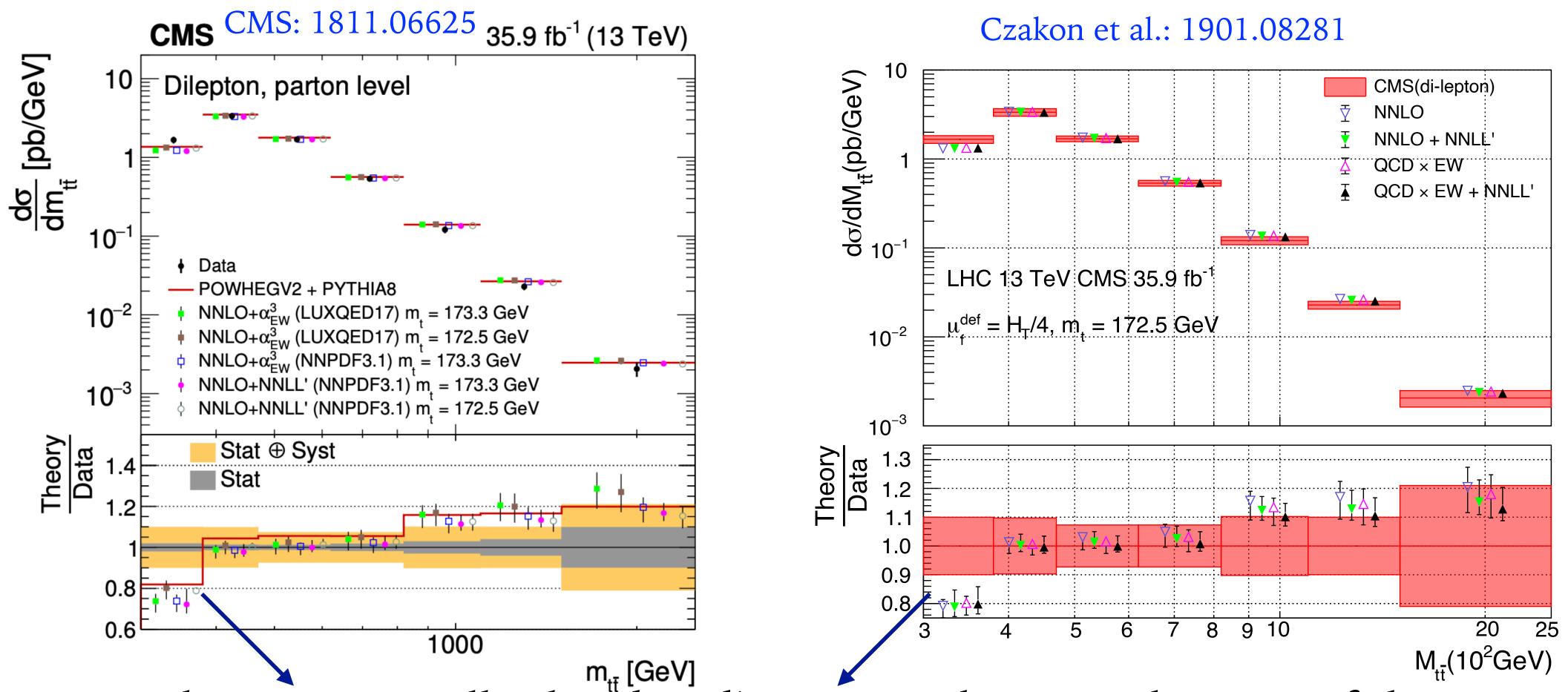
Theoretical predictions vs experimental data



Czakon et al.: 1901.08281

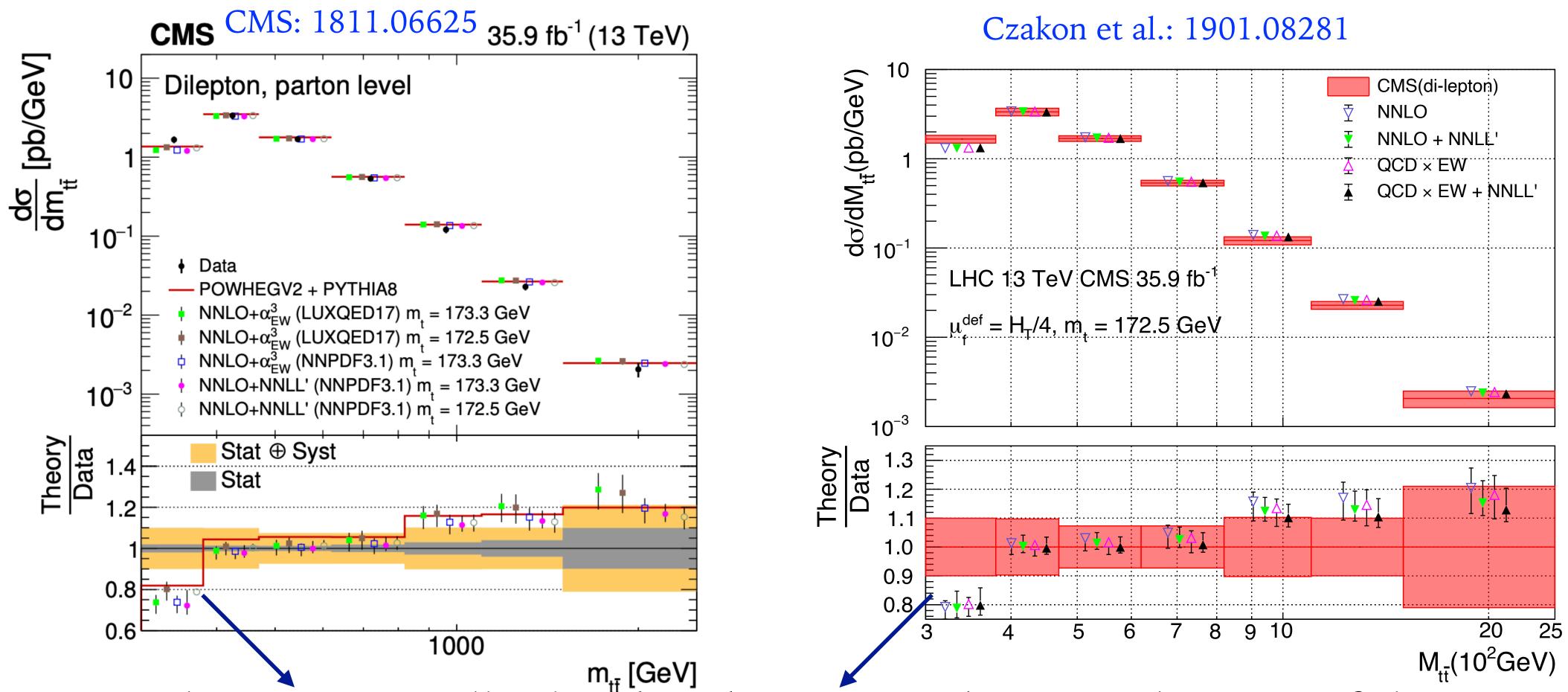


Theoretical predictions vs experimental data

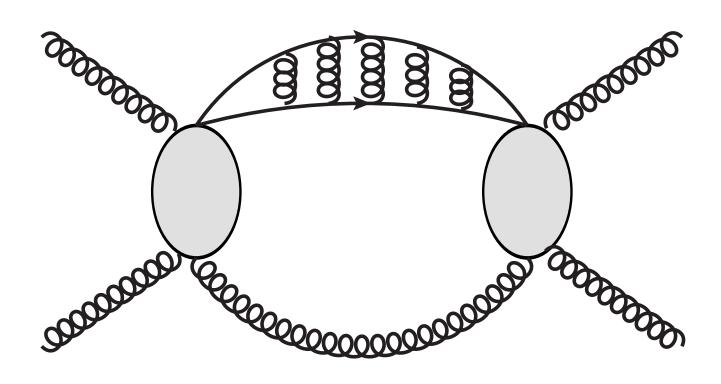


There was actually already a discrepancy between the state-of-the-art theoretical predictions and experimental data in the threshold region!

Theoretical predictions vs experimental data



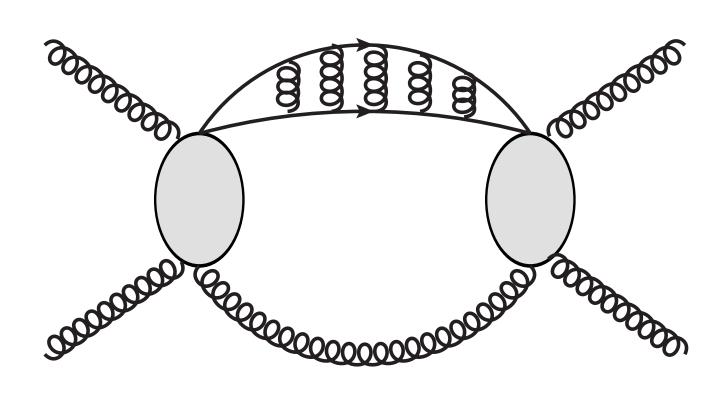
There was actually already a discrepancy between the state-of-the-art theoretical predictions and experimental data in the threshold region!



$$\frac{\mathrm{d}\sigma}{\mathrm{d}M_{t\bar{t}}\mathrm{d}\Theta} \sim \beta \left(1 + \frac{\alpha_s}{\beta} + \frac{\alpha_s^2}{\beta^2} + \frac{\alpha_s^3}{\beta^3} + \cdots \right)$$

$$\beta = \sqrt{1 - \frac{4m_t^2}{M_{t\bar{t}}^2}} \qquad \beta \sim \alpha_s \ll 1$$

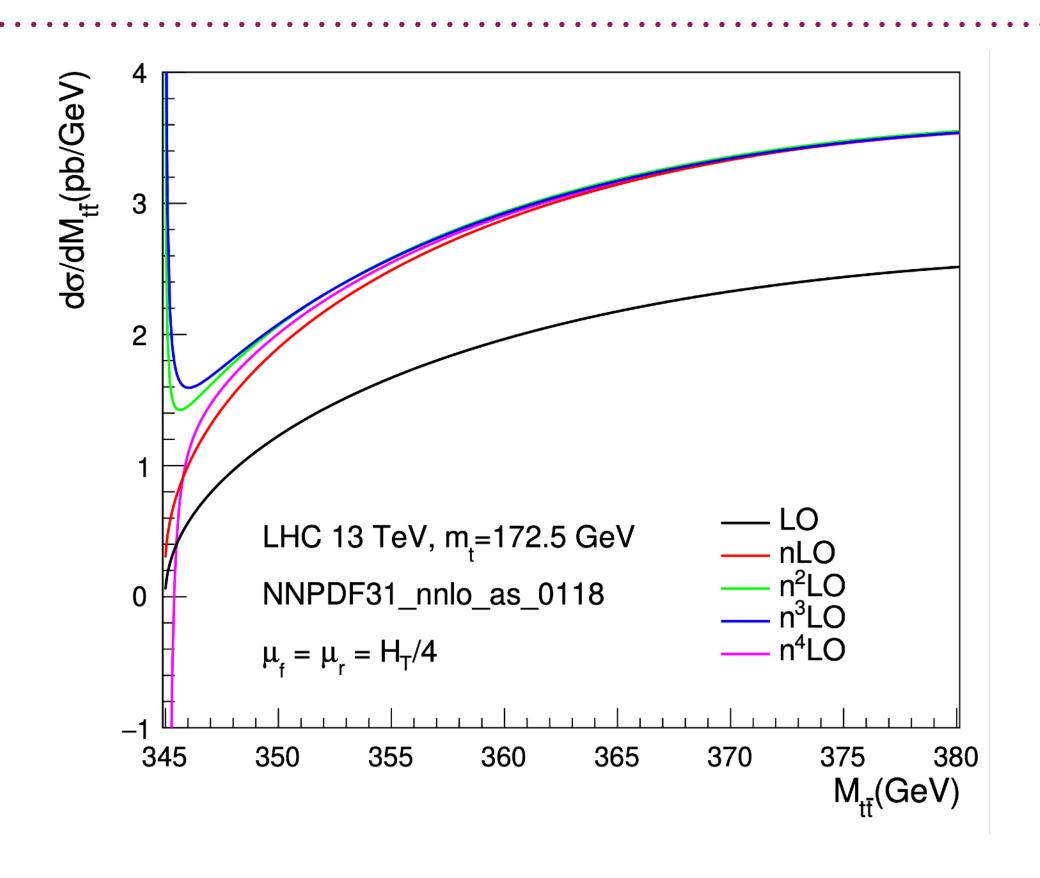
3-velocity of the top quark in the $t\bar{t}$ center-of-mass frame

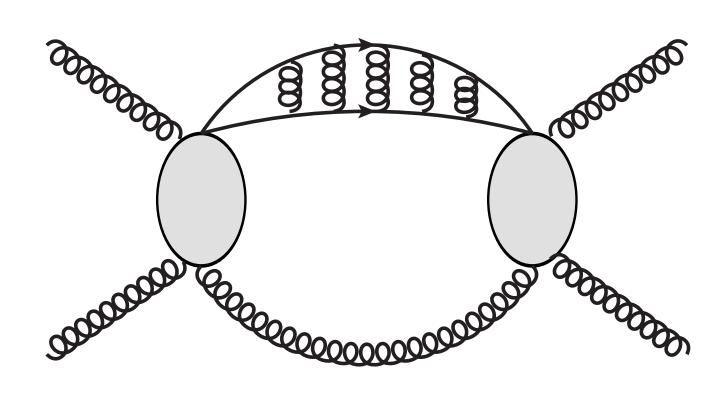


$$\frac{\mathrm{d}\sigma}{\mathrm{d}M_{t\bar{t}}\mathrm{d}\Theta} \sim \beta \left(1 + \frac{\alpha_s}{\beta} + \frac{\alpha_s^2}{\beta^2} + \frac{\alpha_s^3}{\beta^3} + \cdots \right)$$

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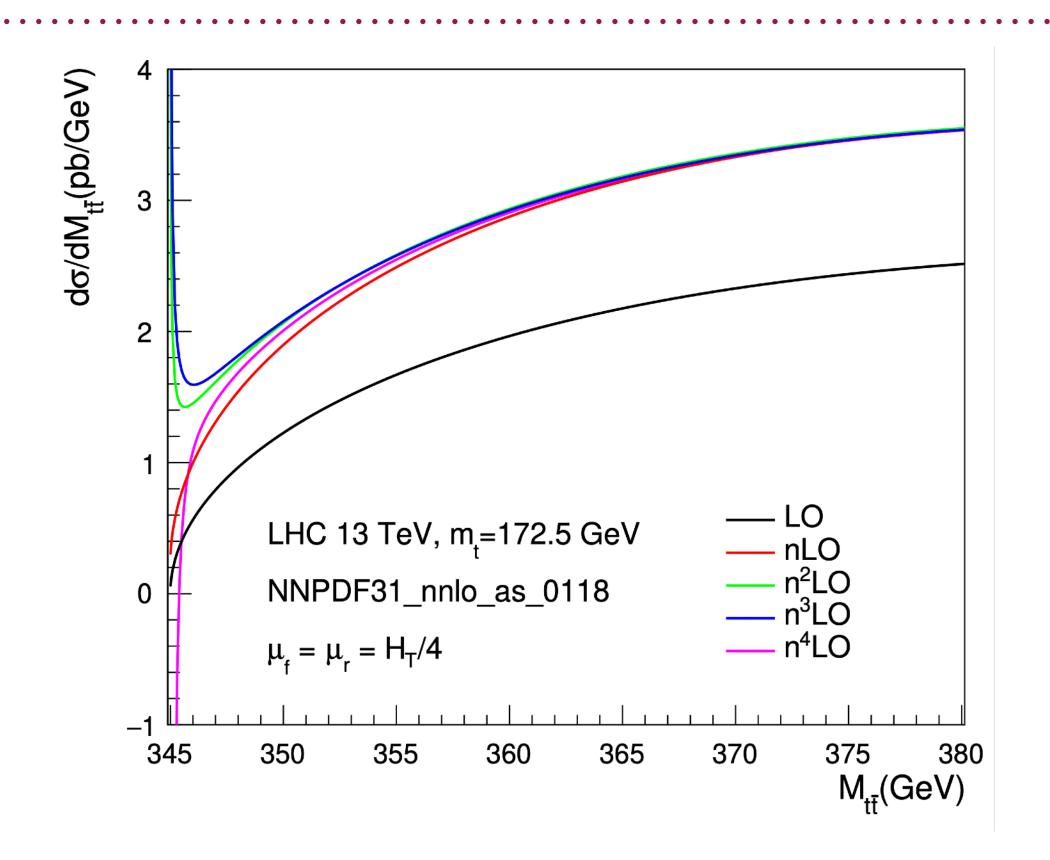




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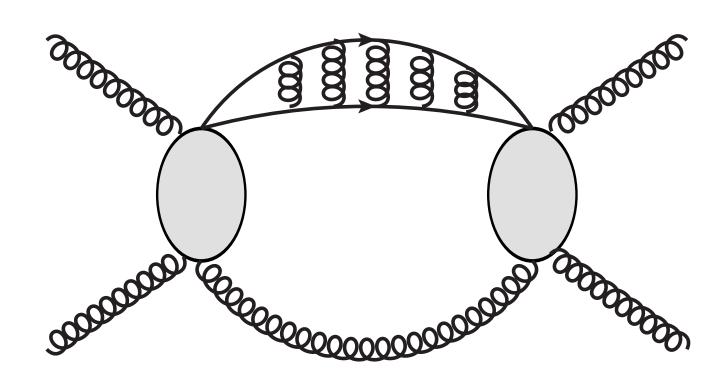
3-velocity of the top quark in the $t\bar{t}$ center-of-mass frame



- ➤ No big deal at NLO
- ➤ Integrable at NNLO

 $\mathrm{d}M_{t\bar{t}} \sim \beta \mathrm{d}\beta$

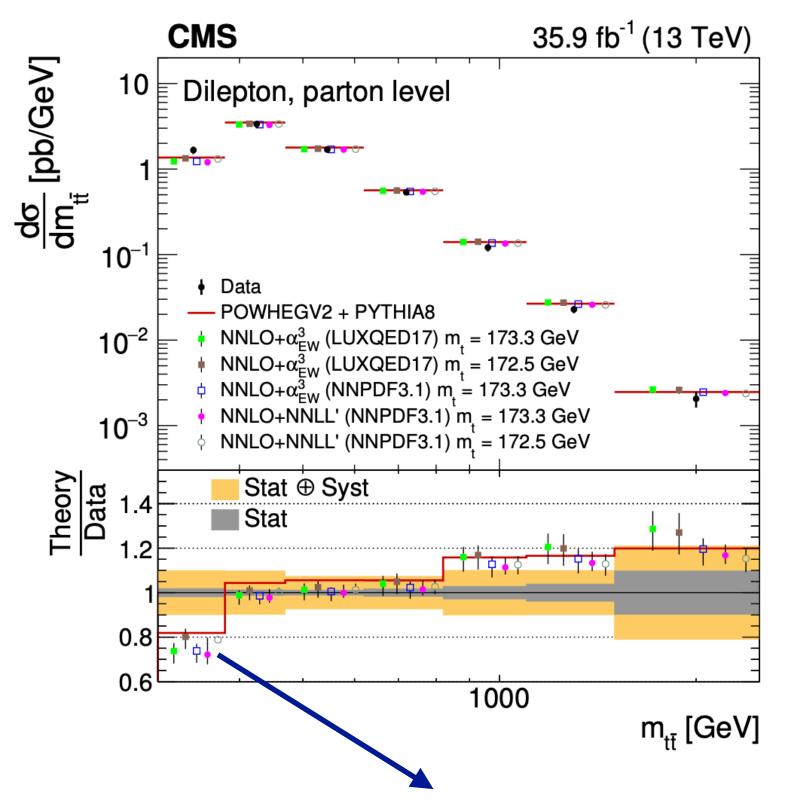
Subtleties beyond NNLO



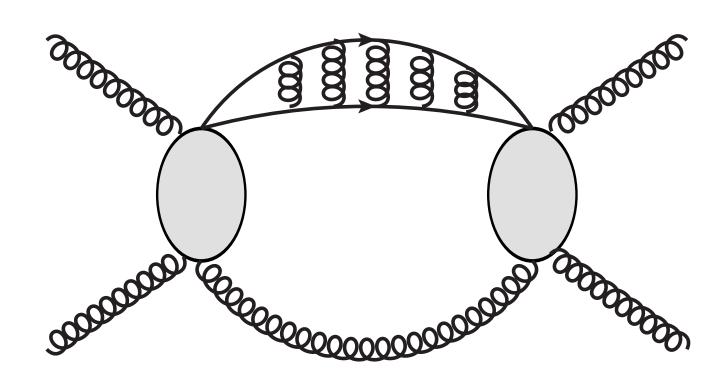
$$\frac{\mathrm{d}\sigma}{\mathrm{d}M_{t\bar{t}}\mathrm{d}\Theta} \sim \beta \left(1 + \frac{\alpha_s}{\beta} + \frac{\alpha_s^2}{\beta^2} + \frac{\alpha_s^3}{\beta^3} + \cdots \right)$$

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3-velocity of the top quark in the $t\bar{t}$ center-of-mass frame



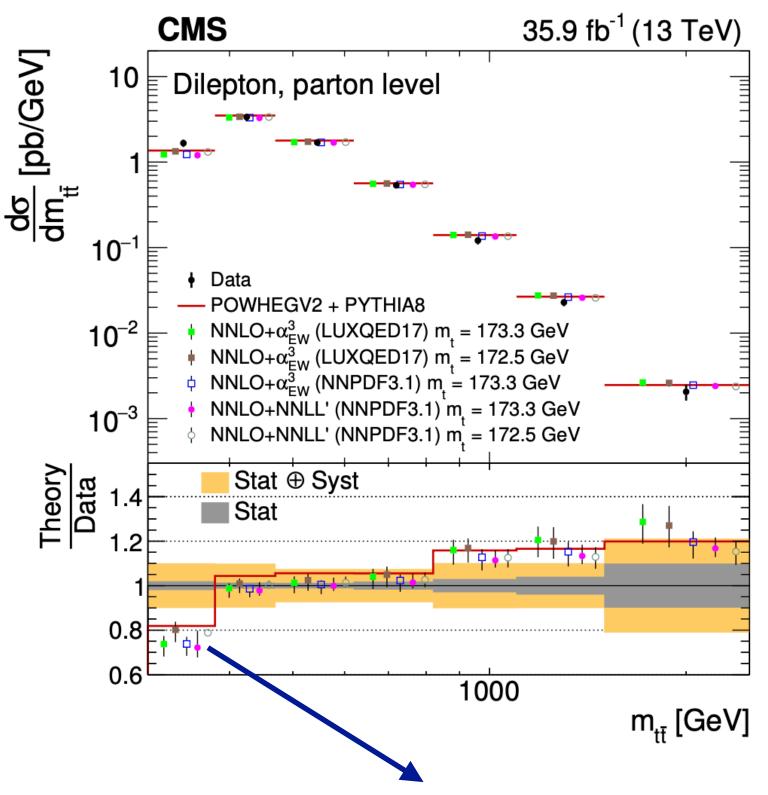
It's not "new" effects — already included in the NNLO+NNLL' calculations, but only up to α_s^2



$$\frac{\mathrm{d}\sigma}{\mathrm{d}M_{t\bar{t}}\mathrm{d}\Theta} \sim \beta \left(1 + \frac{\alpha_s}{\beta} + \frac{\alpha_s^2}{\beta^2} + \frac{\alpha_s^3}{\beta^3} + \cdots \right)$$

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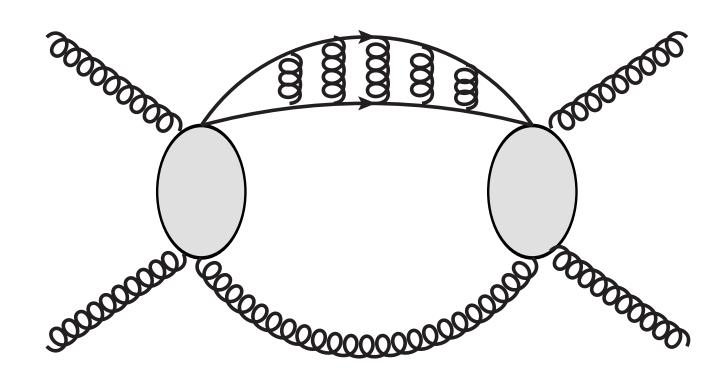
3-velocity of the top quark in the $t\bar{t}$ center-of-mass frame



It's not "new" effects — already included in the NNLO+NNLL' calculations, but only up to α_s^2

Higher order terms are important!

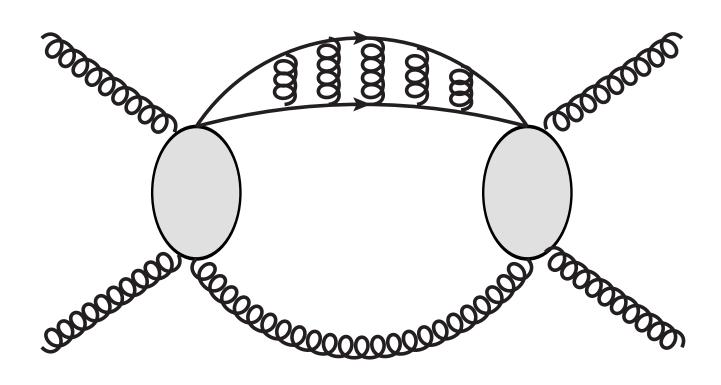
The "toponium"



$$\frac{\mathrm{d}\sigma}{\mathrm{d}M_{t\bar{t}}\mathrm{d}\Theta} \sim \beta \left(1 + \frac{\alpha_s}{\beta} + \frac{\alpha_s^2}{\beta^2} + \frac{\alpha_s^3}{\beta^3} + \cdots \right)$$

- ➤ Similar effects well-studied in quarkonium systems (hence the name "toponium")
- ➤ However, there are important differences!

The "toponium"



$$\frac{\mathrm{d}\sigma}{\mathrm{d}M_{t\bar{t}}\mathrm{d}\Theta} \sim \beta \left(1 + \frac{\alpha_s}{\beta} + \frac{\alpha_s^2}{\beta^2} + \frac{\alpha_s^3}{\beta^3} + \cdots \right)$$

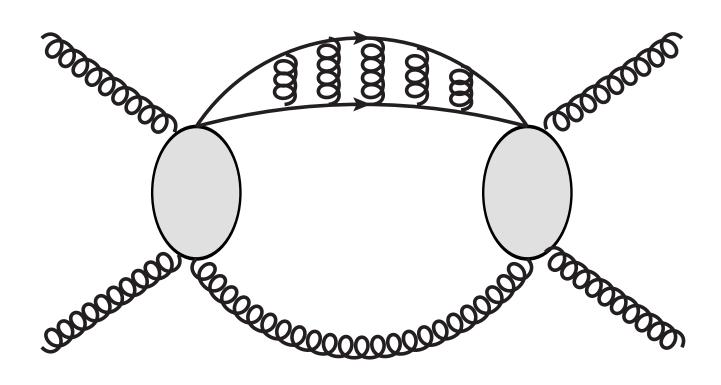
- ➤ Similar effects well-studied in quarkonium systems (hence the name "toponium")
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 η_c, η_b Lifetime $\sim 10^{-22}\,\mathrm{s}$ $J/\psi, \Upsilon$ Lifetime $\sim 10^{-20}\,\mathrm{s}$ Decay mainly through strong interactions (annihilation)

 $t\bar{t}$ Lifetime $\sim 10^{-25}$ s

Decay mainly through weak interactions (individually)

The "toponium"



$$\frac{\mathrm{d}\sigma}{\mathrm{d}M_{t\bar{t}}\mathrm{d}\Theta} \sim \beta \left(1 + \frac{\alpha_s}{\beta} + \frac{\alpha_s^2}{\beta^2} + \frac{\alpha_s^3}{\beta^3} + \cdots \right)$$

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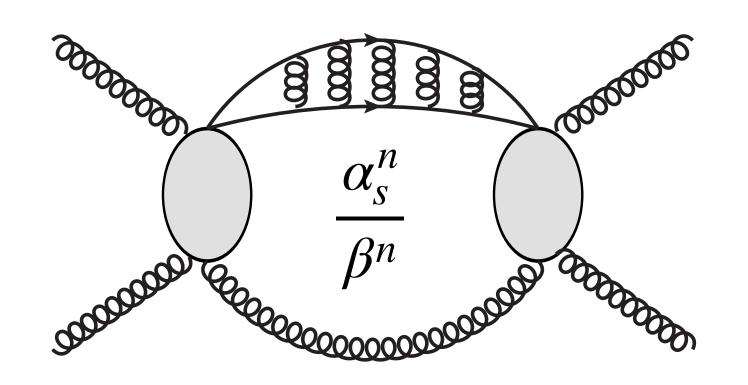
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 $t\bar{t}$ Lifetime $\sim 10^{-25}$ s

Decay mainly through weak interactions (individually)

"quasi-bound state"



See also:

Fadin et al.: 1987~1990

Kiyo et al.: 0812.0919

Sumino and Yokoya: 1007.0075

Beneke et al.: 1109.1536

Some further works:

Fuks et al.: 2102.11281, 2411.18962,

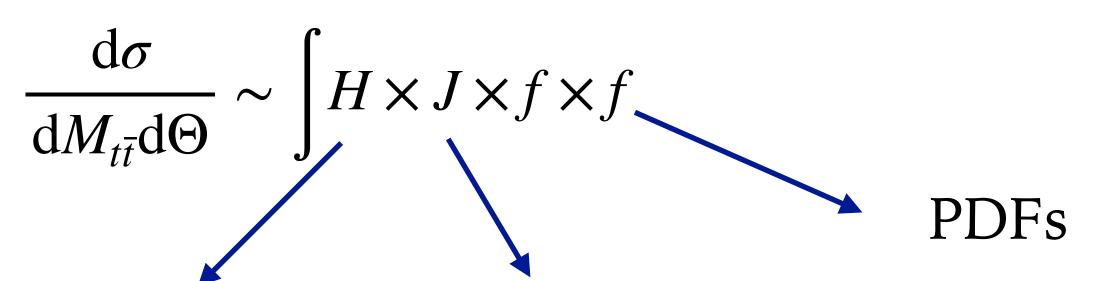
2509.03596

Garzelli et al.: 2412.16685

Wang et al.: 2411.17955

Fu et al.: 2412.11254, 2504.12634

Factorization in the $\beta \sim \alpha_s \ll 1$ limit



Hard functions

Potential functions in NRQCD

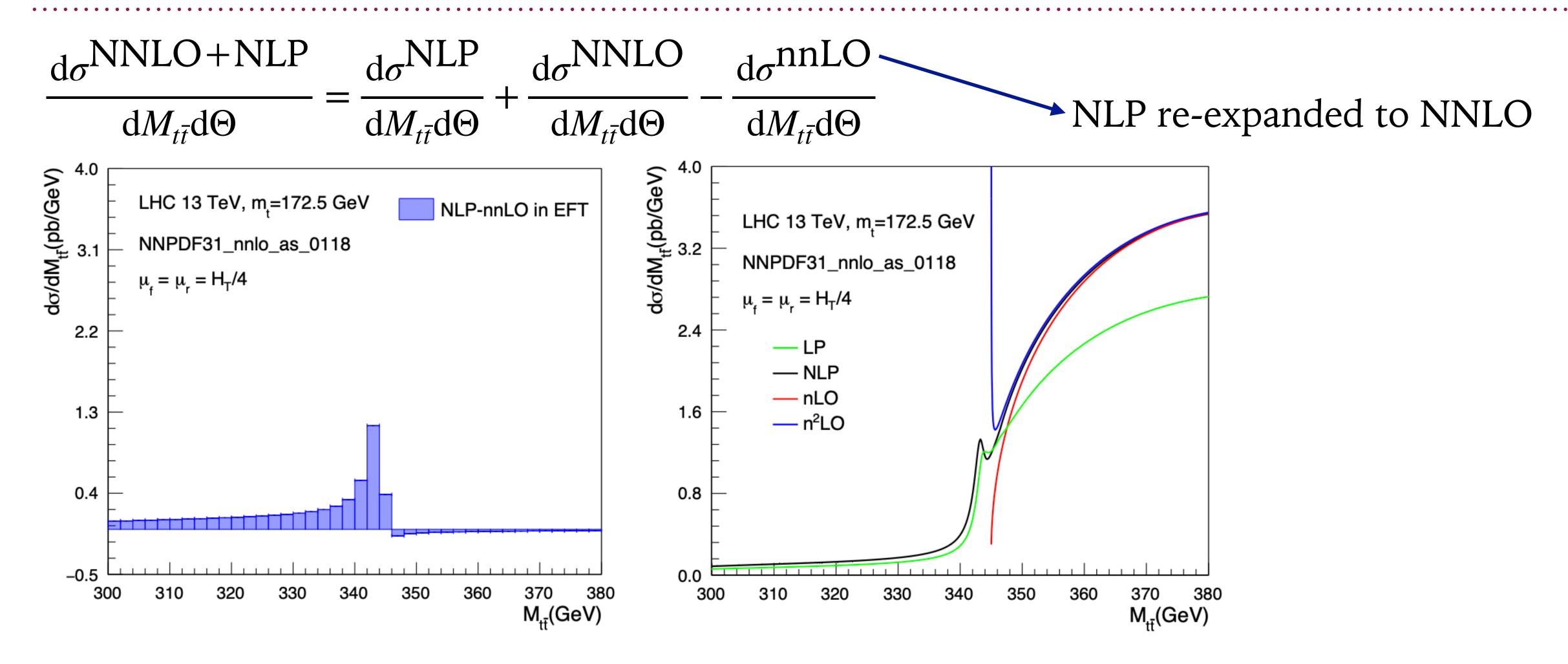
NNLO+NLP resummation

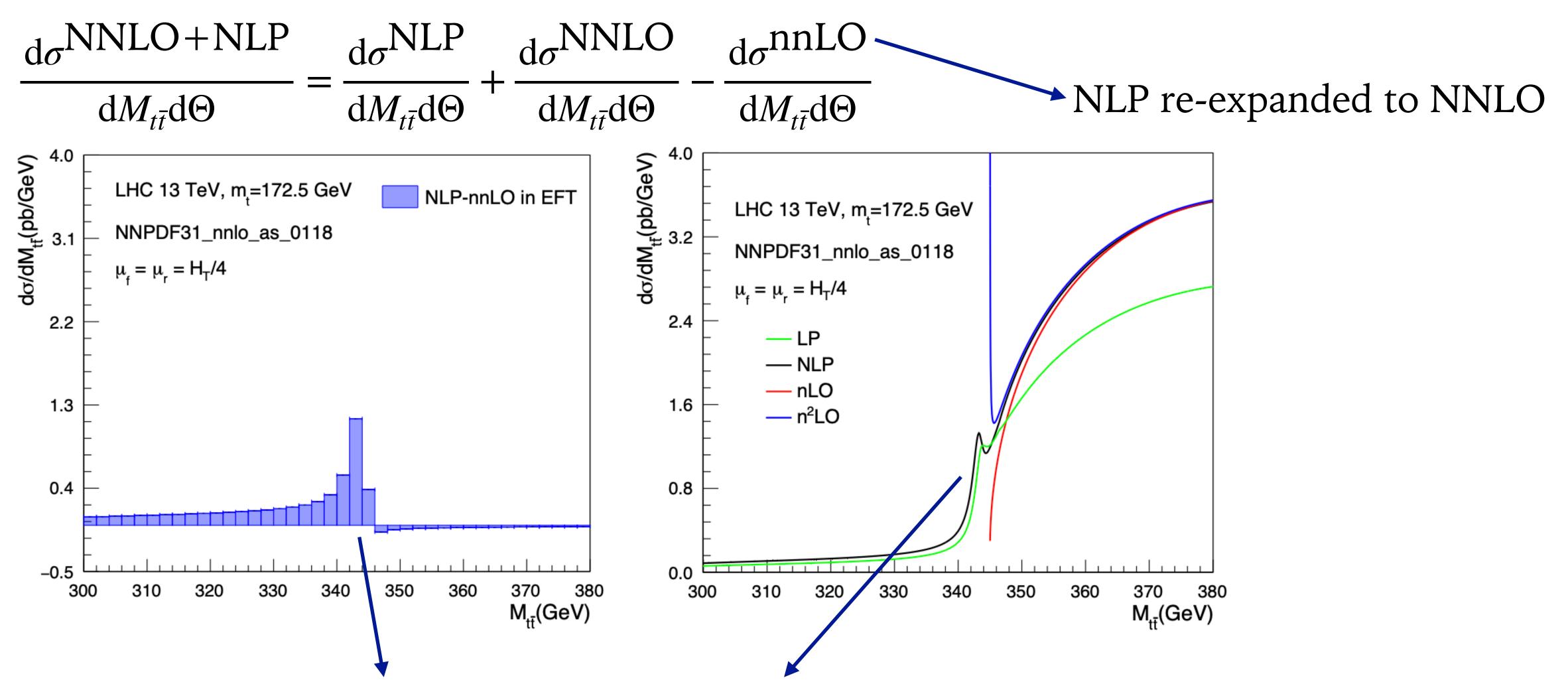
- ➤ Calculated NLO hard functions relevant for double differential distributions and dynamic scales
- ➤ Incorporated potential functions at NLP
- ➤ Combined with NNLO fixed-order results by matching
- First to point out this effect in Run 2 data

Ju, Wang, Wang, Xu, Xu, LLY: 1908.02179, 2004.03088

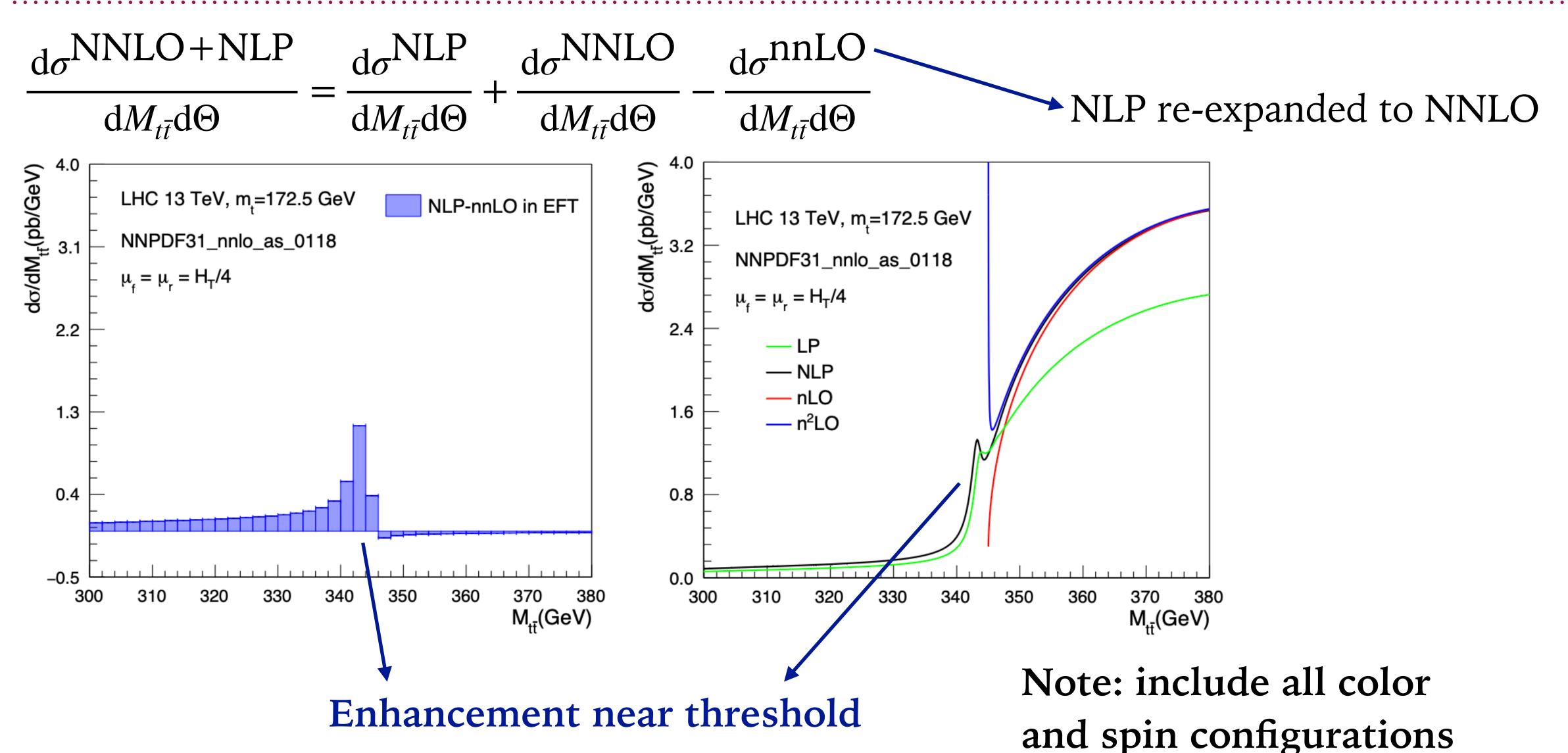
$$\frac{\mathrm{d}\sigma^{\mathrm{NNLO} + \mathrm{NLP}}}{\mathrm{d}M_{t\bar{t}}\mathrm{d}\Theta} = \frac{\mathrm{d}\sigma^{\mathrm{NLP}}}{\mathrm{d}M_{t\bar{t}}\mathrm{d}\Theta} + \frac{\mathrm{d}\sigma^{\mathrm{NNLO}}}{\mathrm{d}M_{t\bar{t}}\mathrm{d}\Theta} - \frac{\mathrm{d}\sigma^{\mathrm{nnLO}}}{\mathrm{d}M_{t\bar{t}}\mathrm{d}\Theta}$$

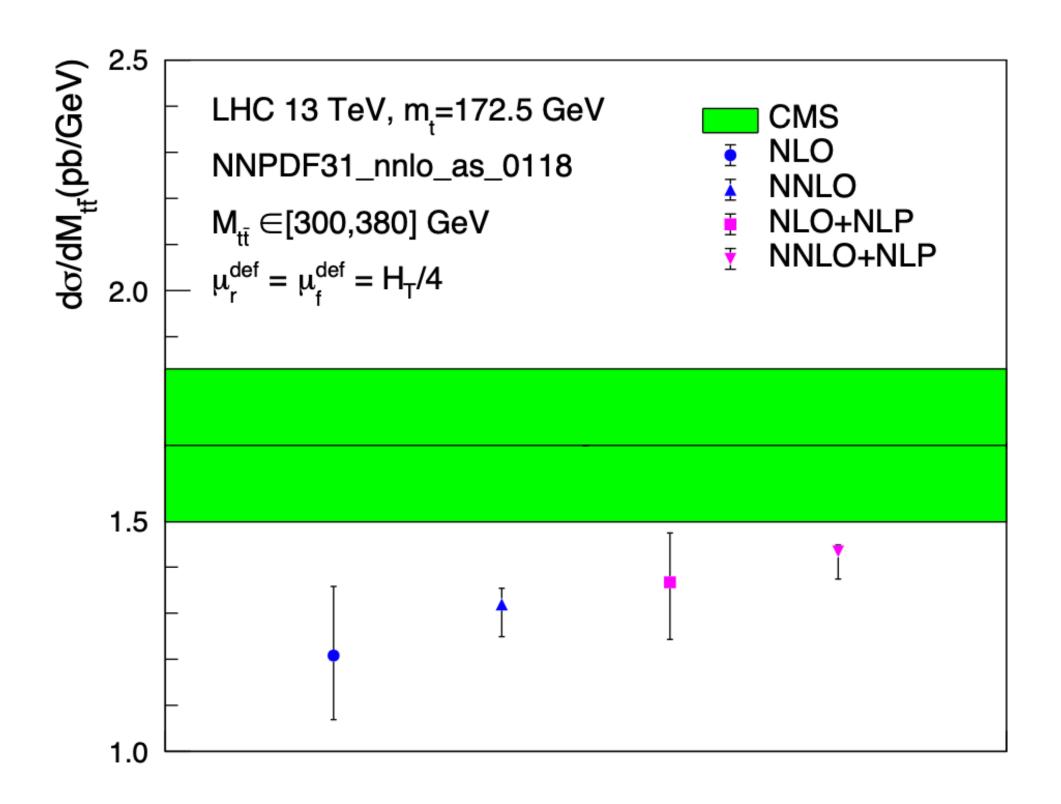
→ NLP re-expanded to NNLO



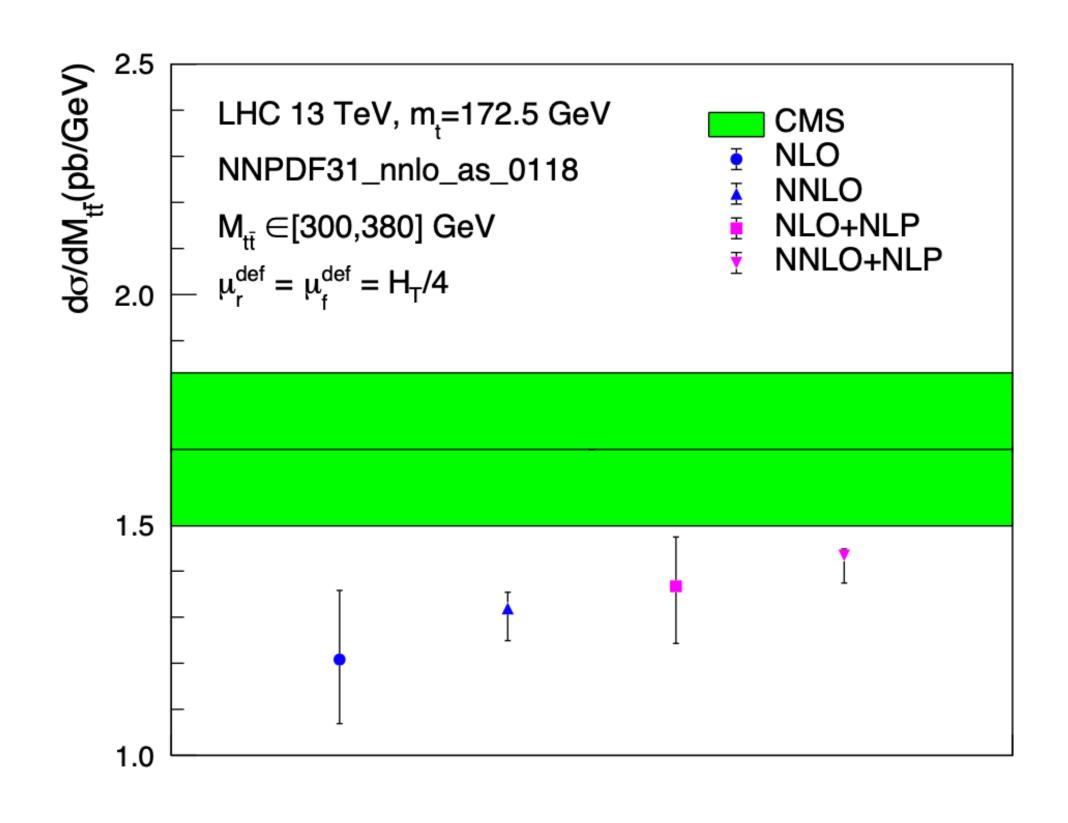


Enhancement near threshold





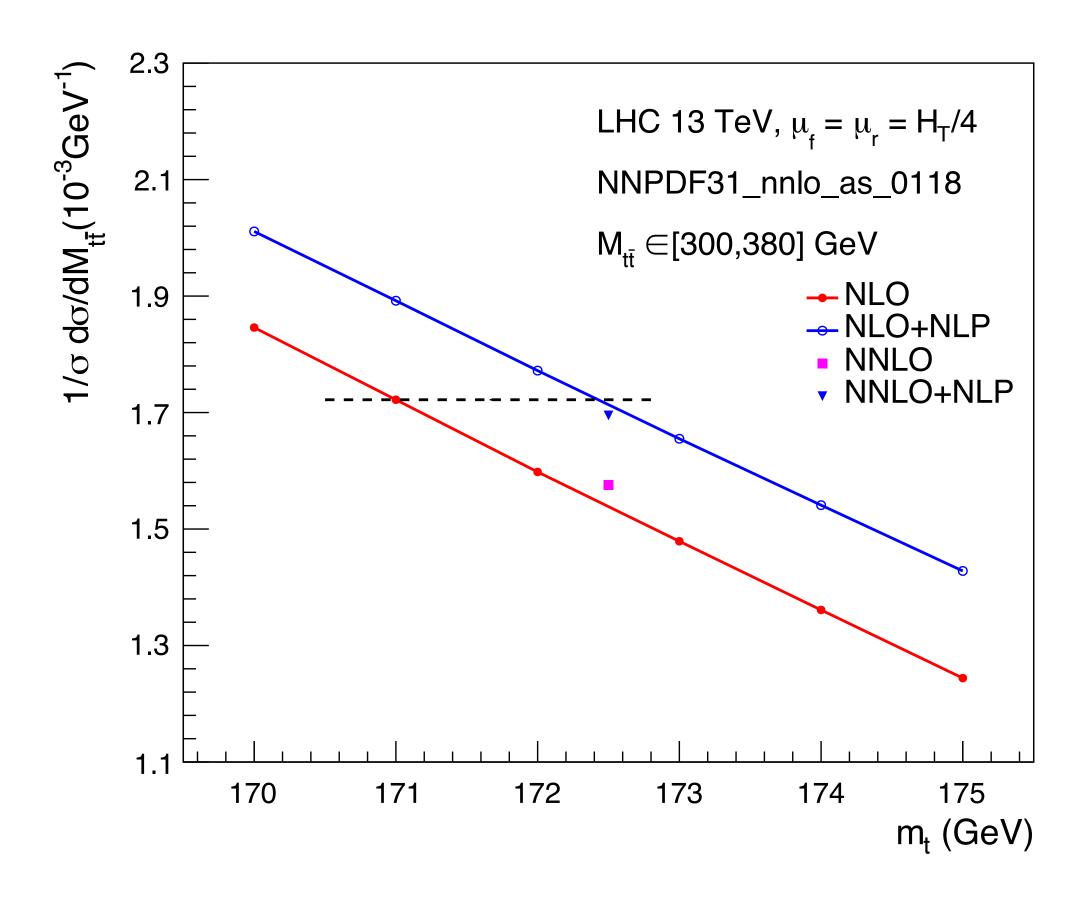
Enhancement ~ 10 pb

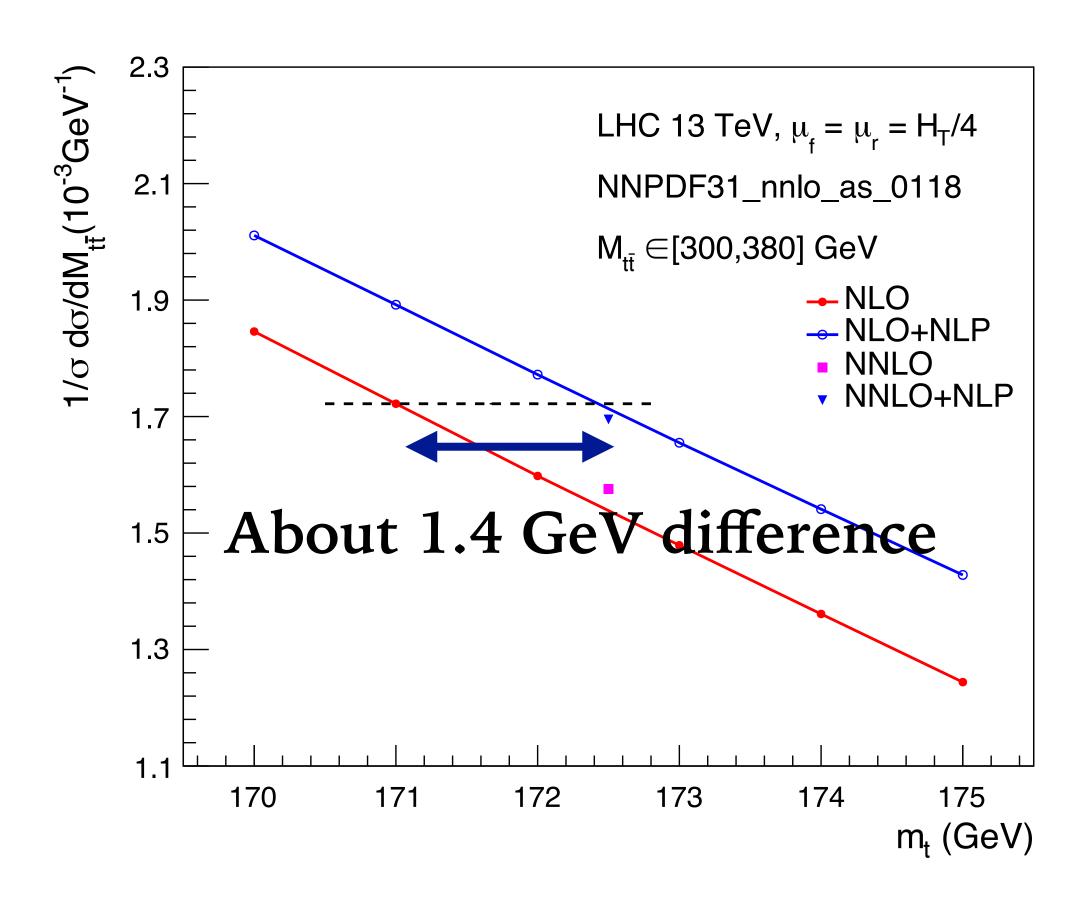


1/5 ds/dY ™ 0.20 LHC 13 TeV, m₌ 172.5 GeV NNPDF31_nnlo_as_0118 $\mu_r^{\text{def}} = H_T/4$ CMS NLO 0.10 NLO+NLP 0.05 M_{,₹} ∈ [300,400] GeV 0.00 Ratio 1.2 1.0 8.0 0.6

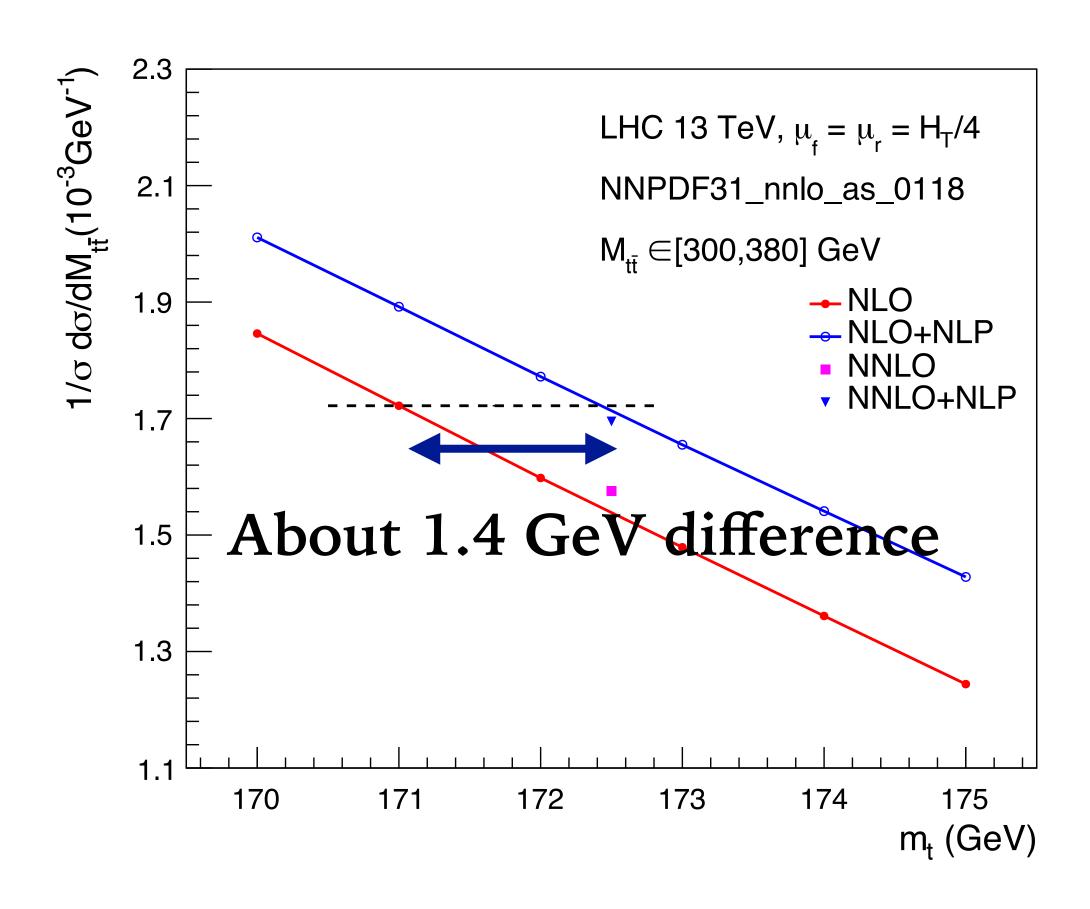
Enhancement ~ 10 pb

Similar effects in double distributions



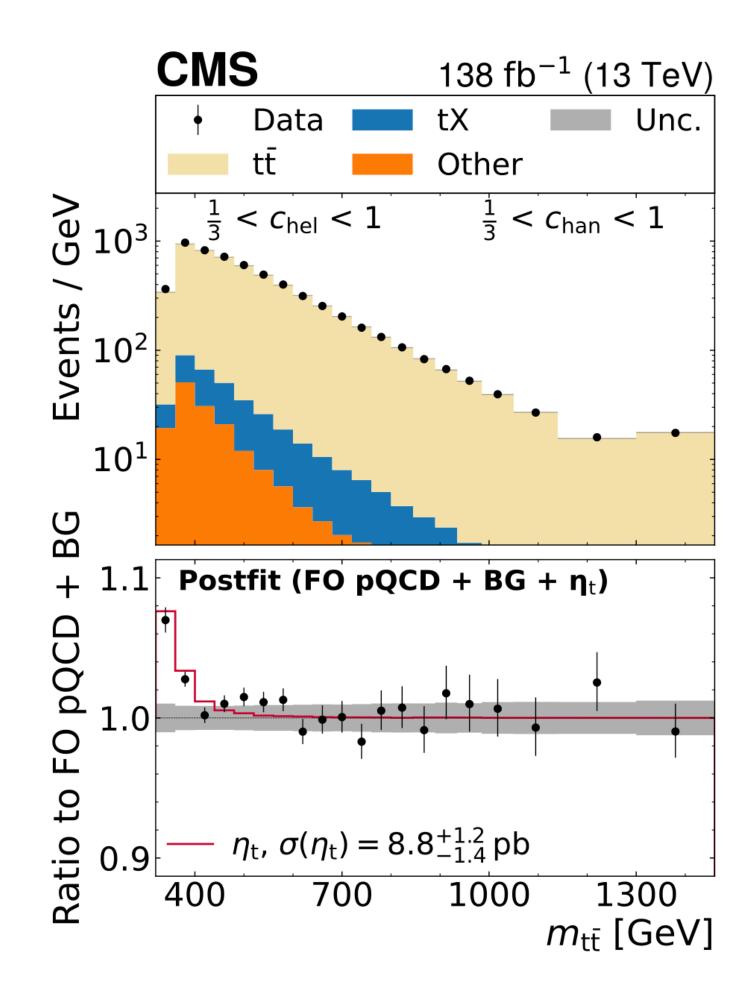


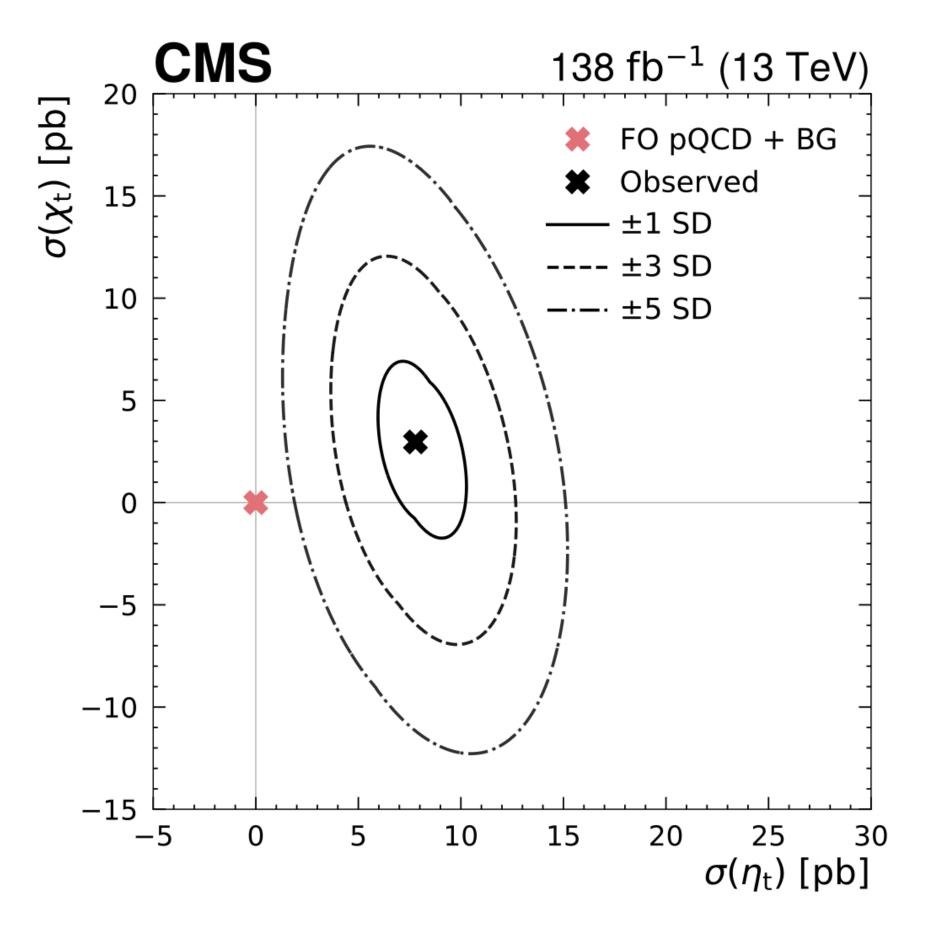
Resummation of non-relativistic enhancement



Account for most of the tension between direct and indirect measurements!

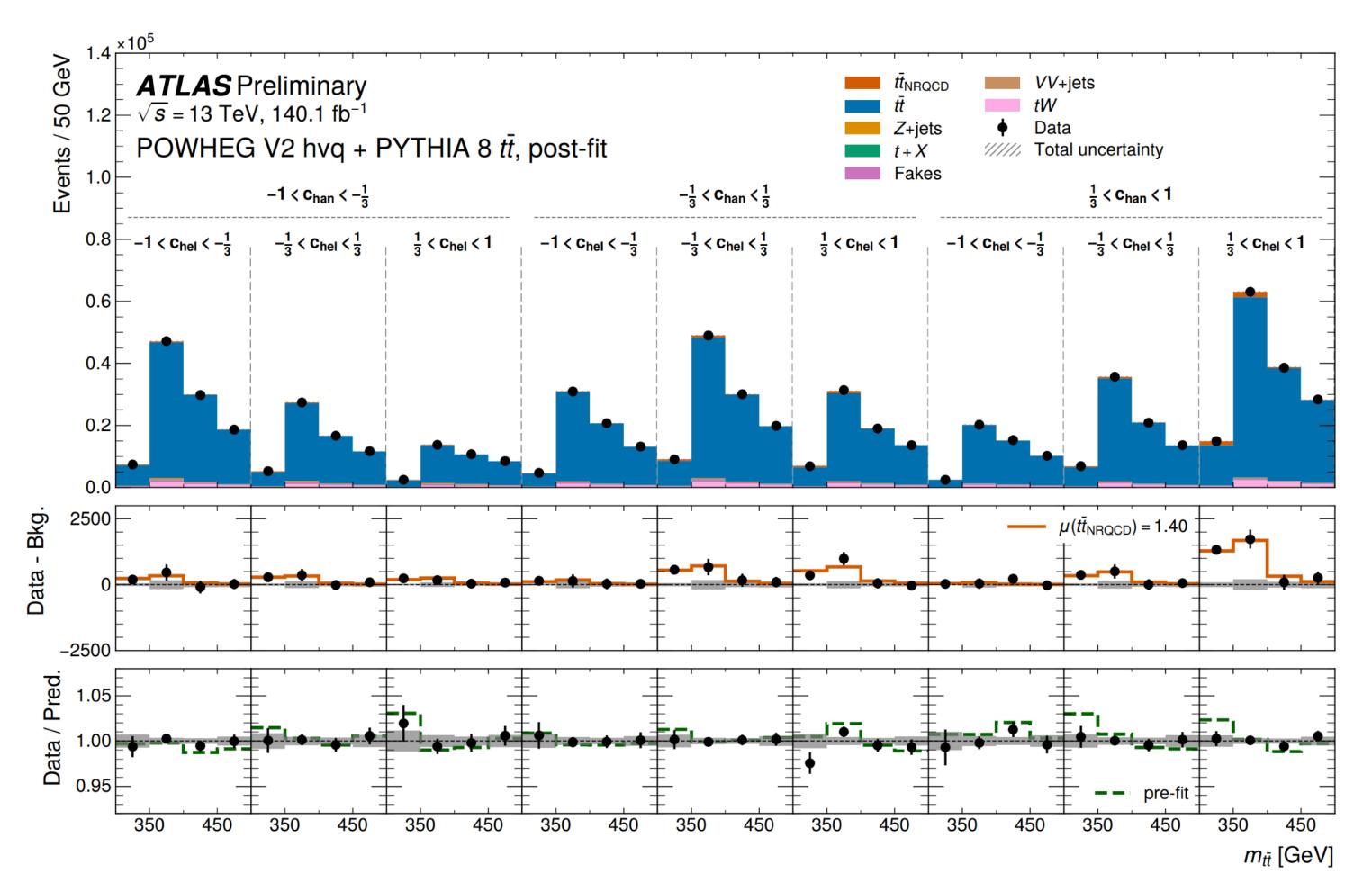
Recent CMS measurement





CMS: 2503.22382

Recent ATLAS measurement



See talk by Hai-Feng Li

ATLAS-CONF-2025-008

Recent works often focus on color-singlet pseudo-scalar configuration " η_t "

Signatures of toponium formation in LHC run 2 data

2102.11281

Benjamin Fuks, 1, 2, * Kaoru Hagiwara, 3, † Kai Ma, 4, 5, 6, ‡ and Ya-Juan Zheng^{7, §}

¹Sorbonne Université, CNRS, Laboratoire de Physique Théorique et Hautes Énergies, LPTHE, F-75005 Paris, France

²Institut Universitaire de France, 103 boulevard Saint-Michel, 75005 Paris, France

³KEK Theory Center, Tsukuba 305-0801, Japan

⁴School of Fundamental Physics and Mathematical Science,

Hangzhou Institute for Advanced Study, UCAS, Hangzhou 310024, Zhejiang, China

⁵International Centre for Theoretical Physics Asia-Pacific, Beijing/Hangzhou, China

⁶Department of Physics, Shaanxi University of Technology, Hanzhong 723000, Shaanxi, China

⁷Department of Physics and Astronomy, University of Kansas, Lawrence, KS 66045, U.S.A.

Simulating toponium formation signals at the LHC

2411.18962

Benjamin Fuks ^{[D] a,1}, Kaoru Hagiwara^{b,2}, Kai Ma ^{[D] c,3} and Ya-Juan Zheng ^{[D] d,4}

Toponium: Implementation of a toponium model in FeynRules

2504.12634

Jing-Hang Fu^a, Yu-Jie Zhang^a, Guang-Zhi Xu^b, Kui-Yong Liu^{c,*}

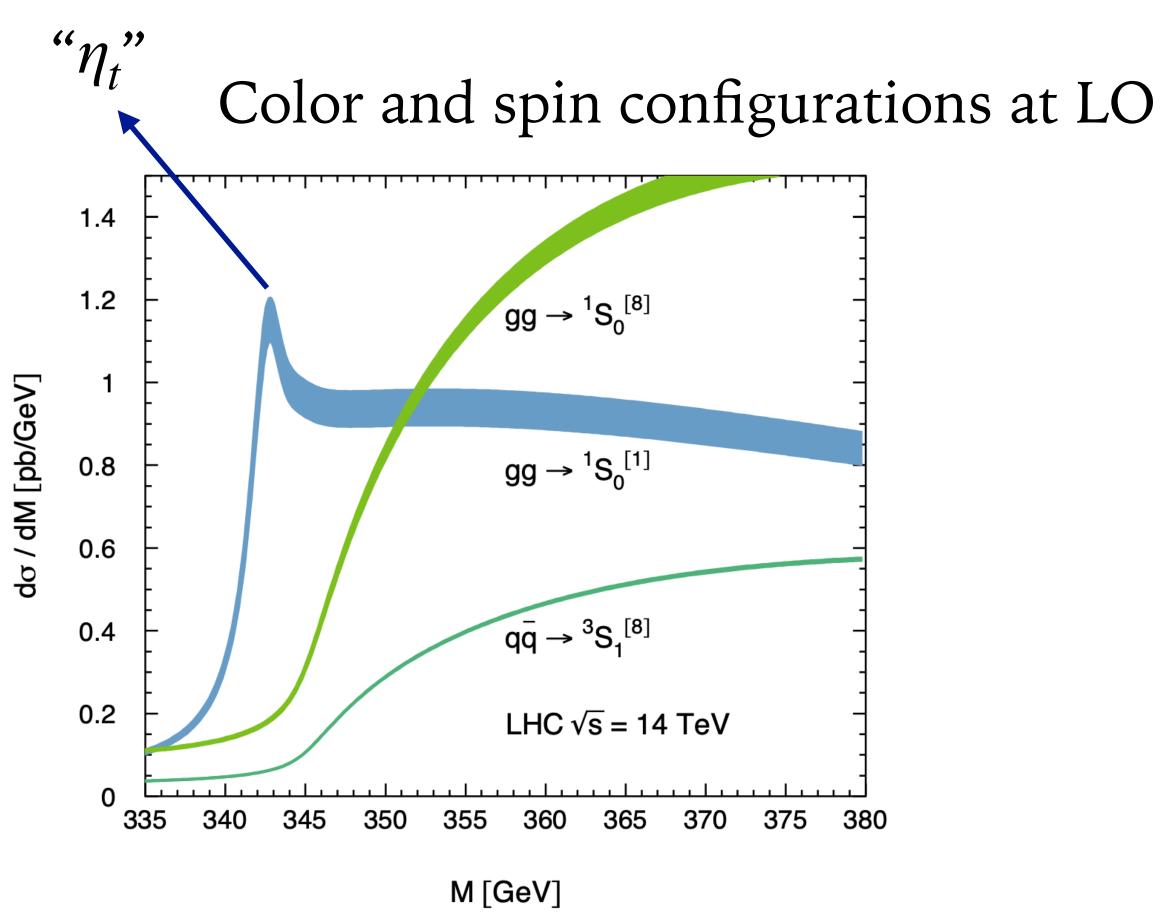
^aSchool of Physics, Beihang University, Beijing, China.
 ^bSchool of Physics, Liaoning University, Shenyang, China.
 ^cSchool of Physics, Liaoning Normal University, Dalian, China.

¹ Laboratoire de Physique Théorique et Hautes Énergies (LPTHE), UMR 7589, Sorbonne Université et CNRS, 4 place Jussieu, 75252 Paris Cedex 05, France

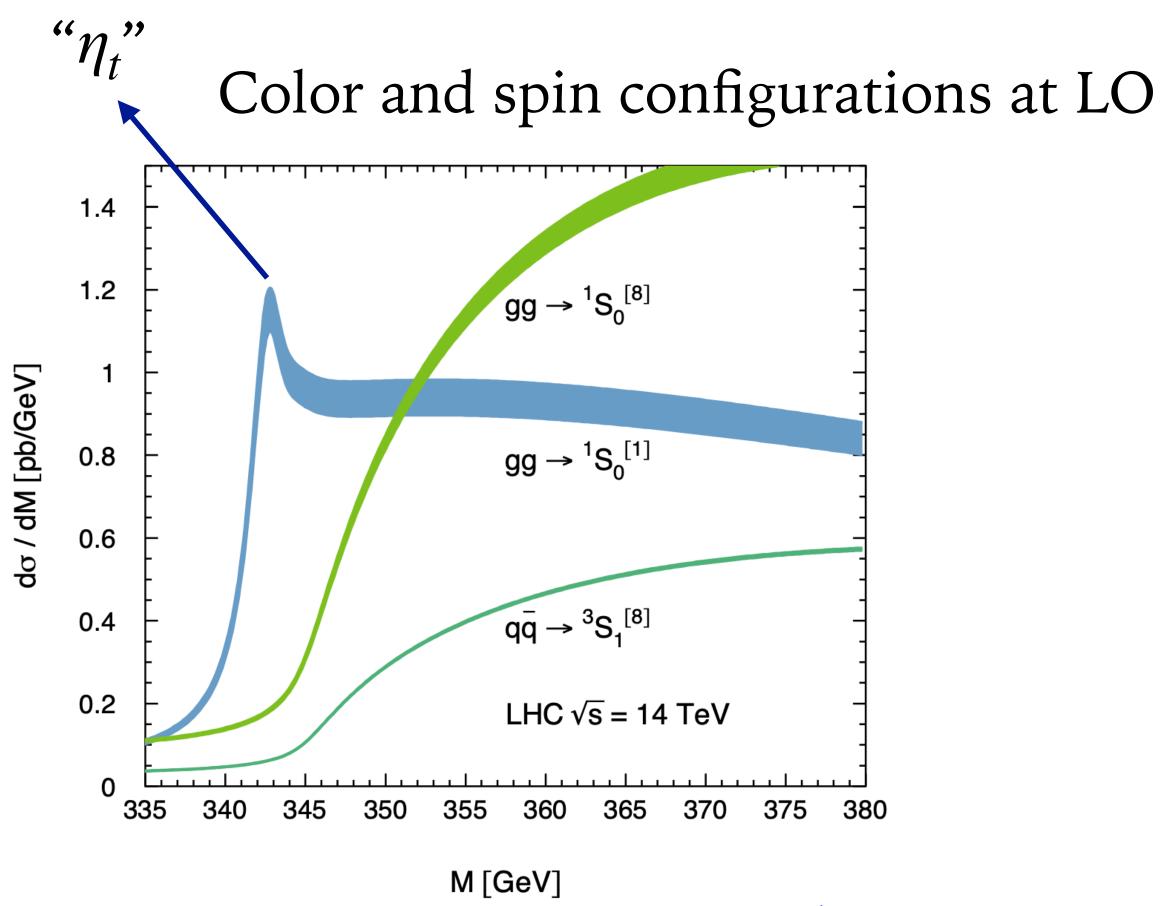
² KEK Theory Center, Tsukuba 305-0801, Japan

³ Faculty of Science, Xi'an University of Architecture and Technology, Xi'an, 710055, China

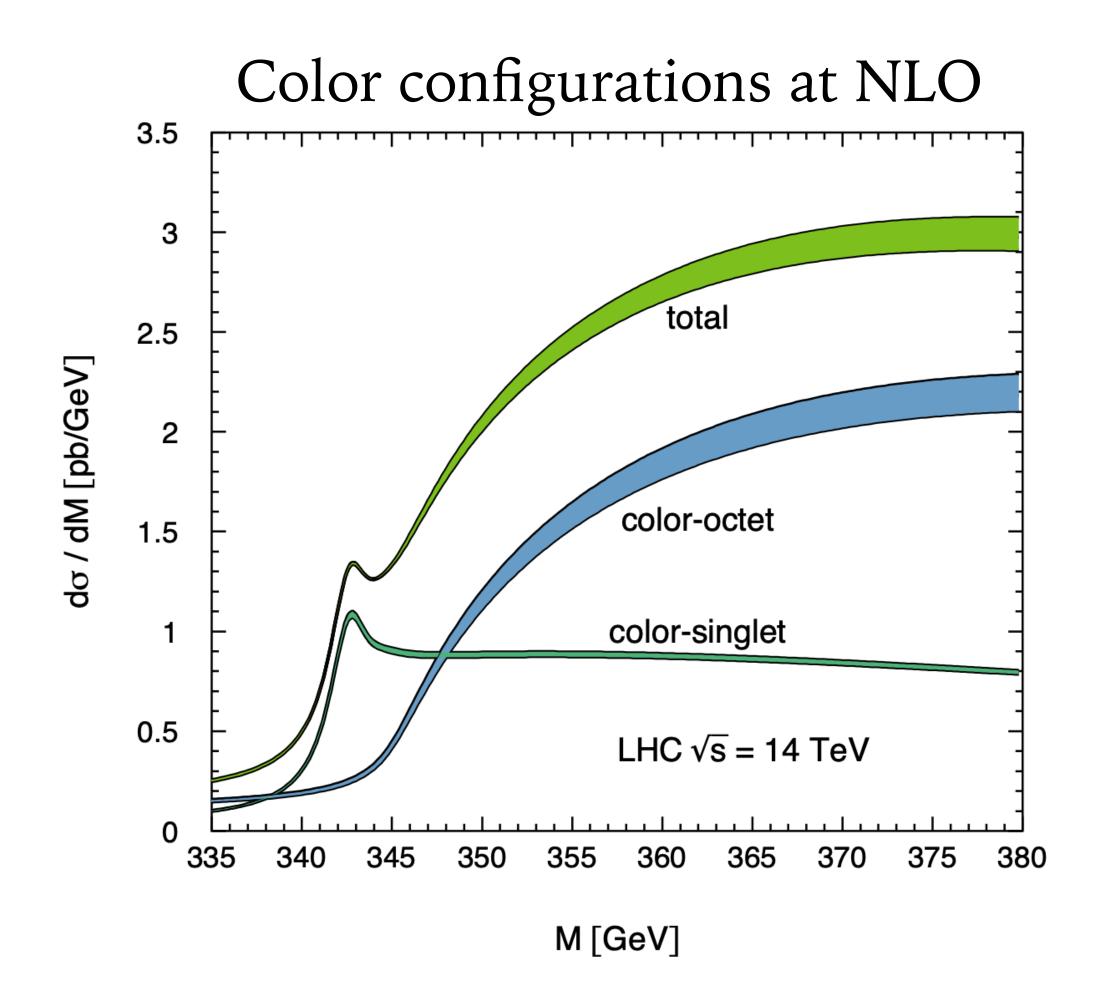
⁴ Faculty of Education, Iwate University, Morioka, Iwate 020-8550, Japan

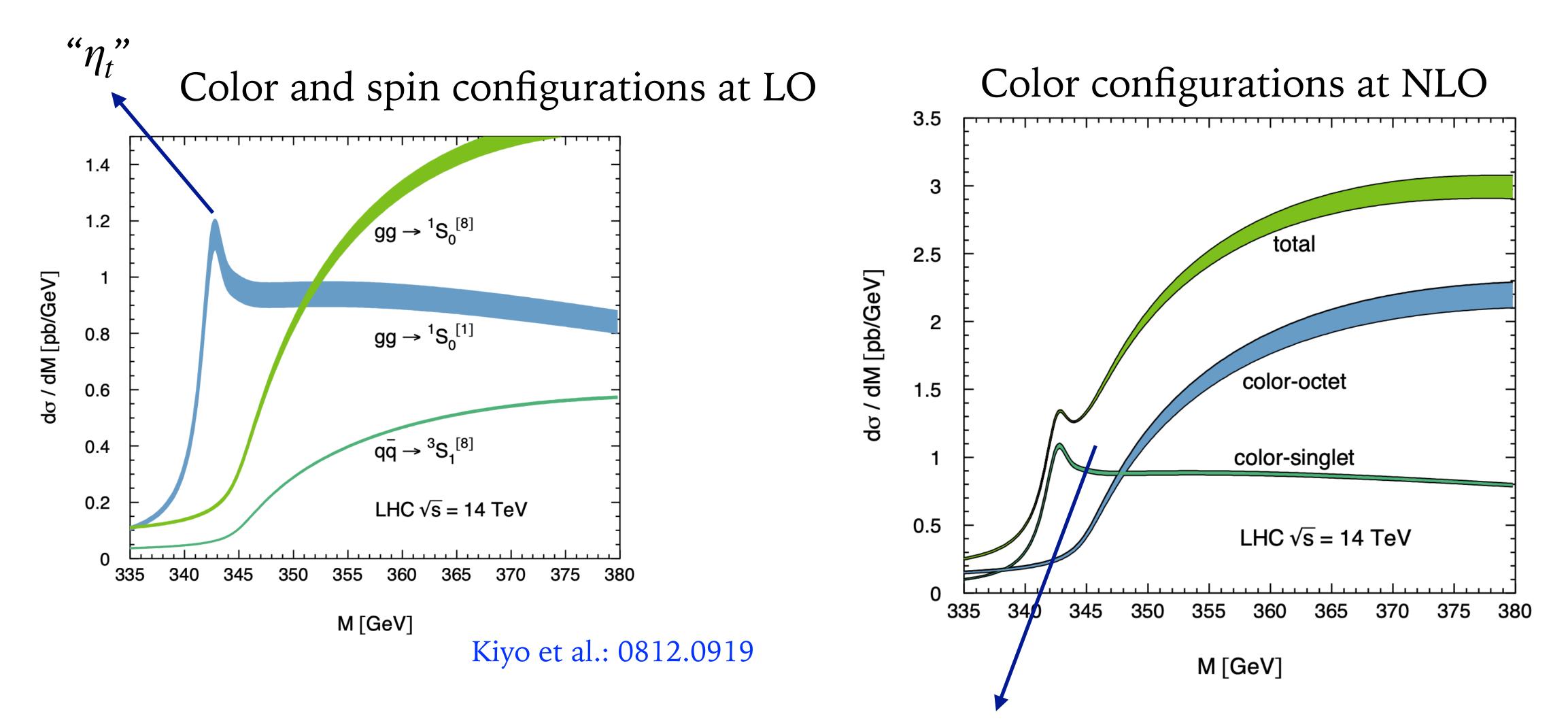


Kiyo et al.: 0812.0919

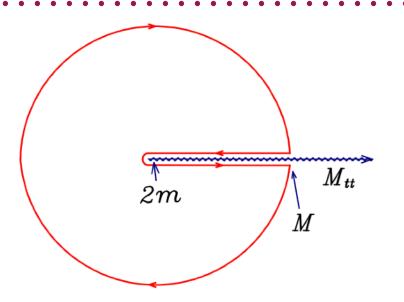








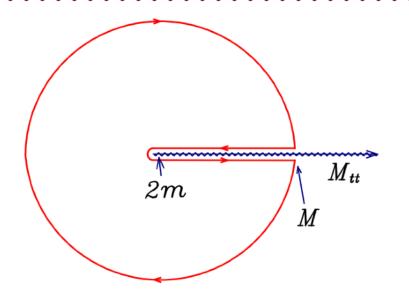
Color-octet contributions are non-negligible



Typical value of β depends on the upper edge of the first bin

Nason et al.: 2505.00096

See also Beneke et al.: 1606.02434



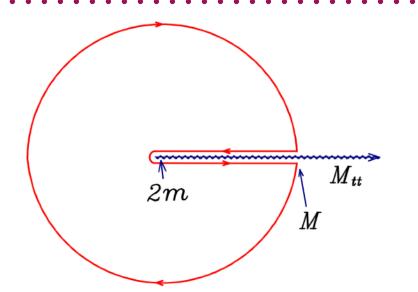
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Nason et al.: 2505.00096

See also Beneke et al.: 1606.02434

Fixed-order expansion is enough instead of a full resummation if $\beta \gg \alpha_s$

$$\frac{\mathrm{d}\sigma}{\mathrm{d}M_{t\bar{t}}\mathrm{d}\Theta} \sim \int H \times J \times f \times f \sim \beta \left(1 + \frac{\alpha_s}{\beta} + \frac{\alpha_s^2}{\beta^2} + \frac{\alpha_s^3}{\beta^3} + \cdots \right)$$



Typical value of β depends on the upper edge of the first bin

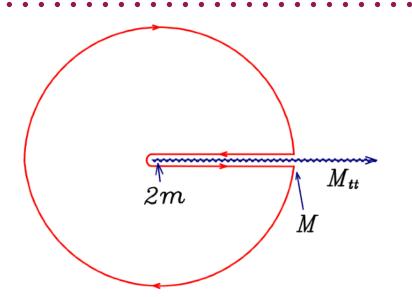
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$$M_{t\bar{t}}$$
 (GeV) β 4000.503800.423600.283500.16



Typical value of β depends on the upper edge of the first bin

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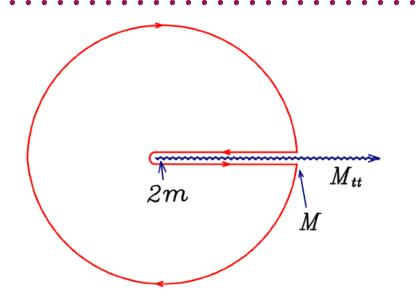
 $M_{t\bar{t}}$ (GeV) β 400 0.50
380 0.42

350

360 0.28

0.16

Resummation is needed for $M_{t\bar{t}} \lesssim 350$ GeV!



Typical value of β depends on the upper edge of the first bin

Nason et al.: 2505.00096

See also Beneke et al.: 1606.02434

Fixed-order expansion is enough instead of a full resummation if $\beta \gg \alpha_s$

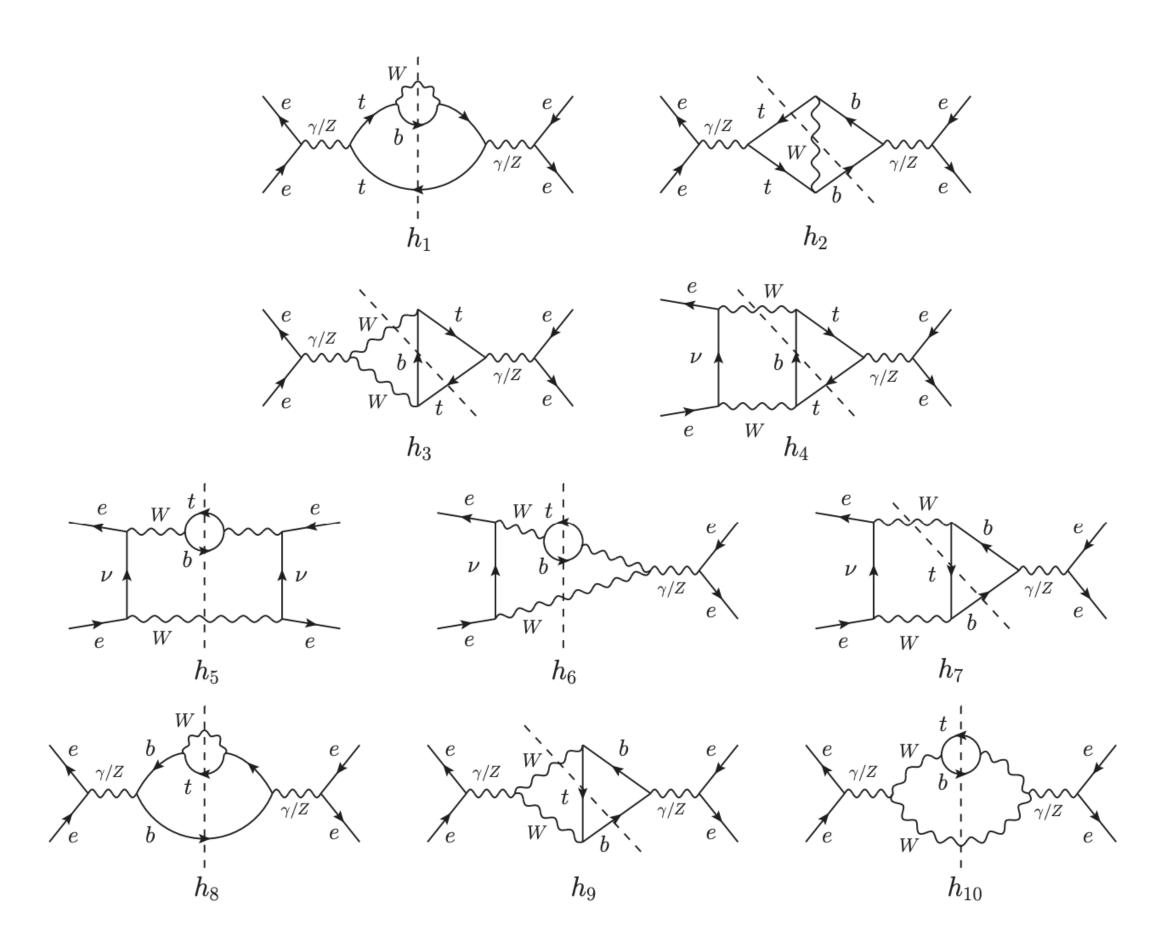
$$\frac{\mathrm{d}\sigma}{\mathrm{d}M_{t\bar{t}}\mathrm{d}\Theta} \sim \int H \times J \times f \times f \sim \beta \left(1 + \frac{\alpha_s}{\beta} + \frac{\alpha_s^2}{\beta^2} + \frac{\alpha_s^3}{\beta^3} + \cdots \right)$$

 $M_{t\bar{t}} ext{ (GeV)} \qquad \beta \ 400 \qquad 0.50 \ 380 \qquad 0.42 \ 360 \qquad 0.28 \ 350 \qquad 0.16 \ \$

Resummation is needed for $M_{t\bar{t}} \lesssim 350$ GeV!

Note also: $M_{t\bar{t}}\beta^2$ is still a perturbative scale

$$E = M_{t\bar{t}} - 2m_t + i\Gamma_t \sim \frac{1}{2}M_{t\bar{t}}\beta^2 + i\Gamma_t$$



Top quark decay and non-resonant effects?

Existing studies only focus on lepton colliders

See, e.g., Beneke et al.: 1711.10429, Bach et al.: 1712.02220

Ongoing: extending to hadron colliders

Summary

- ➤ Near-threshold production of top quark pairs is phenomenologically important
 - ➤ Indirect top quark mass measurements from differential cross sections
 - > Spin correlations and quantum entanglement
- ➤ Indication of non-relativistic enhancement in LHC Run 2 data
 - ➤ Performed NLP resummation combined with NNLO fixed-order results for differential distributions
 - ➤ Enhancement ~ 10 pb, bridging the discrepancy between indirect and direct top quark mass measurements
- ➤ Recent observations by both CMS and ATLAS collaborations

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