



Heavy to light form factors from LQCD

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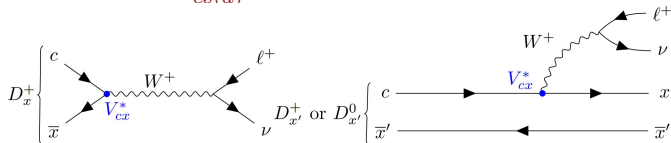
第22届全国重味物理与CP破坏研讨会，2025年10月24-28日，北京

Charmed high-precision frontier

"**November Revolution**"— The discovery of J/ψ particle in 1974, greatly facilitated the establishment of the Standard Model.



- Precise test for Standard Model and search for new physics
 - High-precision measurements(BESIII,etc) and theoretical calculations
 - CKM matrix element $V_{cs(d)}$



- Test the interplay of pert and non-pert QCD
 - intermediate energy scale

Charmed high-precision frontier

- FLAG2024:

$$f_+^{D \rightarrow \pi}(0) = 0.6296(50)$$

$$f_+^{D \rightarrow K}(0) = 0.7430(27)$$

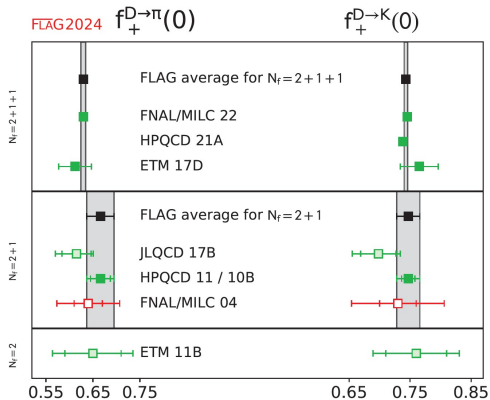
- HFLAV2023

$$|V_{cs}| = 0.9639 \pm 0.0044_{\text{EXP}} \pm 0.0032_{\text{LQCD}}$$

$$|V_{cd}| = 0.2265 \pm 0.0029_{\text{EXP}} \pm 0.0018_{\text{LQCD}}$$

- QED correction needs to be investigated

@杨一玻 talk on QED correction



$P \rightarrow V$ semileptonic decay

- For $P' \rightarrow P$:

$$\frac{d\Gamma(P' \rightarrow P e \nu_e)}{dq^2 d\cos\theta_l} = \frac{G_F^2}{32\pi^3} |\vec{p}|^3 |V_{xy}|^2 \sin^2\theta_l |f_+(q^2)|^2$$

for example, $D \rightarrow \pi$ and $D \rightarrow K$

- For $P \rightarrow V$:

$$\frac{d\Gamma(P \rightarrow V e \nu_e)}{dq^2 d\chi d\cos\theta_\ell d\cos\theta_V} = \frac{G_F^2 |V_{xy}|^2 |\vec{p}| q^2}{12(2\pi)^4 M^2} W(q^2, \theta_V, \theta_\ell, \chi)$$

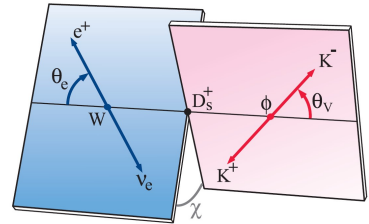
- **Independent** determination of the CKM matrix element
- **Stronger test** with abundant polarization information

$P \rightarrow V$ decay amplitude

$$\begin{aligned}
 W(\theta_K, \theta_\ell, \chi) = & \frac{9}{32} (1 + \cos^2 \theta_\ell) \sin^2 \theta_K (|H_+|^2 + |H_-|^2) + \frac{9}{8} \sin^2 \theta_\ell \cos^2 \theta_K |H_0|^2 \\
 & + \frac{9}{16} \cos \theta_\ell \sin^2 \theta_K (|H_+|^2 - |H_-|^2) - \frac{9}{16} \sin^2 \theta_\ell \sin^2 \theta_K \cos 2\chi (H_+ H_-) \\
 & + \frac{9}{16} \sin \theta_\ell \sin 2\theta_K \cos \chi (H_+ H_0 - H_- H_0) \\
 & + \frac{9}{32} \sin 2\theta_\ell \sin 2\theta_K \cos \chi (H_+ H_0 + H_- H_0) \\
 & + \frac{m_\ell^2}{2q^2} \left[\frac{9}{4} \cos^2 \theta_K |H_t|^2 - \frac{9}{2} \cos \theta_\ell \cos^2 \theta_K (H_0 H_t) + \frac{9}{4} \cos^2 \theta_\ell \cos^2 \theta_K |H_0|^2 \right. \\
 & + \frac{9}{16} \sin^2 \theta_\ell \sin^2 \theta_K (|H_+|^2 + |H_-|^2) + \frac{9}{8} \sin^2 \theta_\ell \sin^2 \theta_K \cos 2\chi (H_+ H_-) \\
 & + \frac{9}{8} \sin \theta_\ell \sin 2\theta_K \cos \chi (H_+ H_t + H_- H_t) \\
 & \left. - \frac{9}{16} \sin 2\theta_\ell \sin 2\theta_K \cos \chi (H_+ H_0 + H_- H_0) \right]
 \end{aligned}$$

and the helicity amplitudes are

$$\begin{aligned}
 H_\pm(q^2) &= (M + m) A_1(q^2) \mp \frac{2M|\vec{p}|}{M + m} V(q^2) \\
 H_0(q^2) &= \frac{1}{2M\sqrt{q^2}} \times \left[(M^2 - m^2 - q^2) (M + m) A_1(q^2) - 4 \frac{M^2 |\vec{p}|^2}{M + m} A_2(q^2) \right] \\
 H_t(q^2) &= \frac{2M|\vec{p}|}{\sqrt{q^2}} A_0(q^2)
 \end{aligned}$$



$P \rightarrow V$ matrix element

- For $P' \rightarrow P$:

$$\langle P(p) | V_\mu | H(p') \rangle = f_+(q^2) \left[(p' + p)_\mu - \frac{M^2 - m^2}{q^2} q_\mu \right] + f_0(q^2) \frac{M^2 - m^2}{q^2} q_\mu$$

- For $P \rightarrow V$:

$$\begin{aligned} \langle h_V(p) | V_\mu - A_\mu | H(p') \rangle &= \varepsilon_\nu^* \epsilon_{\mu\nu\alpha\beta} p'_\alpha p_\beta \frac{2V}{m+M} + (M+m) \varepsilon_\mu^* A_1 \\ &+ \frac{\varepsilon^* \cdot q}{M+m} (p+p')_\mu A_2 - 2m \frac{\varepsilon^* \cdot q}{q^2} q_\mu (A_0 - A_3) \end{aligned}$$

- Two key physical quantities:

$$r_V \equiv V(0)/A_1(0), \quad r_2 \equiv A_2(0)/A_1(0)$$

Charmed $P \rightarrow V$ decay

- Experiments:

Channels	r_V	r_2	Refs
$D^0 \rightarrow K^*(892)$			
(μ)	$1.37 \pm 0.09 \pm 0.03$	$0.76 \pm 0.06 \pm 0.02$	PRL134,011803(2025)
(μ)	$1.46 \pm 0.11 \pm 0.04$	$0.71 \pm 0.08 \pm 0.03$	PRL135,111803(2025)
(e)	$1.48 \pm 0.05 \pm 0.02$	$0.70 \pm 0.04 \pm 0.02$	JHEP03,197(2025)
Average	$1.456 \pm 0.040 \pm 0.016$	$0.715 \pm 0.031 \pm 0.014$	
$D_s \rightarrow K^{*,0} e \nu_e$	$1.67 \pm 0.34 \pm 0.16$	$0.77 \pm 0.28 \pm 0.07$	PRL122,061801(2019)
$D_s \rightarrow \phi$			
(e)	$1.807 \pm 0.046 \pm 0.065$	$0.816 \pm 0.036 \pm 0.030$	PRD78,051101(2008)
(μ)	$1.58 \pm 0.17 \pm 0.02$	$0.71 \pm 0.14 \pm 0.02$	JHEP12,072(2023)
$D^0 \rightarrow \rho^- e^+ \nu_e$	$1.695 \pm 0.083 \pm 0.051$	$0.845 \pm 0.056 \pm 0.039$	PRL122,062001(2019)

- We target on $D \rightarrow K^*$ and $D_s \rightarrow \phi$ using lattice QCD
 - Both experimental precision $\sim 3\%$

$$D_s \rightarrow \phi \text{ \& } D \rightarrow K^*(892)$$

$$K^*(892) \quad I(J^P) = 1/2(1^-)$$

$K^*(892)$ T-Matrix Pole $\sqrt{s}$

$(890 \pm 14) - i(26 \pm 6) \text{ MeV}$

$K^*(892)$ DECAY MODES

Mode	Fraction (Γ_i / Γ)	Scale Factor/ Conf. Level	P(MeV/c)
$\Gamma_1 \quad K\pi$	$\sim 100 \%$		289

$$\phi(1020) \quad I^G(J^{PC}) = 0^-(1^{--})$$

Expand/Collapse All

$\phi(1020)$ MASS

$1019.460 \pm 0.016 \text{ MeV}$



$\phi(1020)$ WIDTH

$4.249 \pm 0.013 \text{ MeV (S = 1.1)}$



- Γ_{K^*} is not so small, $K^* \rightarrow K\pi$ may have visible effect
- Γ_ϕ is only $\sim 4 \text{ MeV}$, a reasonable approximation to regard ϕ as stable

This work: $D_s \rightarrow \phi l \nu_l$

- Previous lattice calculation: [HPQCD,PRD90,074506\(2014\)](#)
 - **Three ensembles** with two coarse lattice spacings, ~ 0.09 fm and ~ 0.12 fm
 - Extrapolated to physical limit by z -expansion

$$F(z) = \sum_{n=0}^3 B_n^F [1 + C_n^F a^2 + D_n^F a^4 + E_n^F (m_l/m_{s,\text{phys}})] z^n$$

leading to $r_V = 1.72(21)$, $r_2 = 0.74(12)$

- Experiment:

	r_V	r_2	Refs
BaBar(e)	$1.807 \pm 0.046 \pm 0.065$	$0.816 \pm 0.036 \pm 0.030$	PRD78,051101(2008)
BESIII(μ)	$1.58 \pm 0.17 \pm 0.02$	$0.71 \pm 0.14 \pm 0.02$	JHEP12,072(2023)

- **A more precise and systematic lattice study is essential**

Extraction of the form factor: scalar function method

- Matrix element parameterization

$$\langle h_\nu / J_\nu^{\text{em}}(p) | J_\mu^{\text{em}/W}(0) | H(p') \rangle = \sum_j \mathcal{Q}_{\mu\nu}^j(p, p') F_j(q^2)$$

- $P \rightarrow 2\gamma$: $\pi_0/\eta/\eta'/\eta_c \rightarrow 2\gamma$
- $V \rightarrow \gamma P$: $J/\psi \rightarrow \gamma\eta_c, D_s^* \rightarrow \gamma D_s, D^* \rightarrow \gamma D$
- $P \rightarrow P' l \nu_l$, $P \rightarrow V l \nu_l$
- $P \rightarrow \gamma l \nu_l, P \rightarrow l' \bar{l}' l \nu_l$

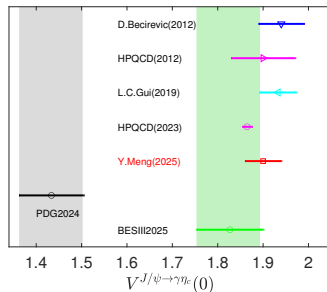
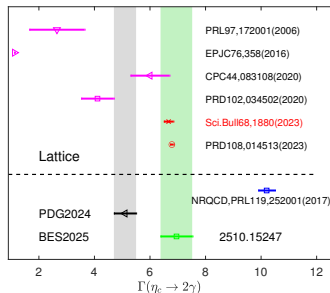
- Constructing scalar functions

$$\begin{aligned} \tilde{I}_i(q^2) &\equiv \mathcal{Q}_{\mu\nu}^i \cdot \langle h_\nu(p) | J_\mu^{\text{em}/W}(0) | H(p') \rangle \\ &= \sum_i \omega_{ij}(p, p') F_j(q^2), \quad \omega_{ij} \equiv \mathcal{Q}_{\mu\nu}^i \cdot \mathcal{Q}_{\mu\nu}^j \end{aligned}$$

- Solving the linear equation

$$F = \omega^{-1} \cdot \tilde{I}$$

Scalar function method: applications



- $\Gamma(\eta_c \rightarrow 2\gamma) = 6.67(16)(6)$ keV, $\Gamma(J/\psi \rightarrow \gamma\eta_c) = 2.30(10)$ keV
Y.M et al, Sci Bull 68, 1880 (2023) Y.M et al, PRD 111, 014508 (2025)
- A first step toward resolving the long-standing charmonium-decay puzzle
 - We first find **2.9 σ tension and PDG-fit is misleading**, arXiv:2109.09381
 - H-P.Wang and C-Z.Yuan, New puzzle in charmonium decays, arXiv:2112.08584
 - HPQCD(2023) and BESIII(2025)

Scalar function method: applications

- The most precise value $\Gamma(D_s^* \rightarrow \gamma D_s) = 0.0549(54)$ keV
 - Precision $\sim 10\%$, much improved than $0.066(26)$ keV by HPQCD PRL112,212002(2014)
 - Combined with experiment: $f_{D_s^*}|V_{cs}| = (190.5^{+55.1}_{-41.7_{\text{stat.}}} \pm 9.1_{\text{syst.exp}} \pm 8.7_{\text{syst.latt}})$ MeV

Y.M et al, PRD109,074511(2024)

- First lattice study on J/ψ semileptonic decay

channels	Upper limit/Br	Refs
$J/\psi \rightarrow D_s e \nu_e$	4.9×10^{-5}	PLB639,418(2006)
$J/\psi \rightarrow D_s e \nu_e$	1.3×10^{-6}	PRD90,112014(2014)
$J/\psi \rightarrow D e \nu_e$	7.1×10^{-8}	JHEP06,157(2021)
$J/\psi \rightarrow D \mu \nu_\mu$	5.6×10^{-7}	JHEP01,126(2024)
$J/\psi \rightarrow D_s e \nu_e$	$1.90(8) \times 10^{-10}$	this work
$J/\psi \rightarrow D_s \mu \nu_\mu$	$1.84(8) \times 10^{-10}$	this work
$J/\psi \rightarrow D e \nu_e$	$1.21(11) \times 10^{-11}$	this work
$J/\psi \rightarrow D \mu \nu_\mu$	$1.18(11) \times 10^{-11}$	this work

Y.M et al,PRD110,074510(2024)

Scalar function method: $D_s \rightarrow \phi l \nu_l$

- $\langle \phi_\sigma(\vec{p}) | J_\mu^W(0) | D_s(p') \rangle = \frac{F_0(q^2)}{Mm} \epsilon_{\mu\sigma\alpha\beta} p'_\alpha p_\beta + F_1(q^2) \delta_{\mu\sigma} + \frac{F_2(q^2)}{Mm} p_\mu p'_\sigma + \frac{F_3(q^2)}{M^2} p'_\mu p'_\sigma$

- Constructing various scalar functions $\tilde{I}_i(q^2)$

$$F_0 = \frac{m}{2} \tilde{I}_0, \quad F_1 = \frac{1}{2} \tilde{I}_1 + \frac{1}{2} \tilde{I}_2 - \frac{m^2}{2M^2} \tilde{I}_3$$

$$F_2 = \frac{mE}{2(E^2 - m^2)} \tilde{I}_1 + \frac{mE^2 + 2m^3}{2(E^3 - Em^2)} \tilde{I}_2 - \frac{3Em^3}{2M^2(E^2 - m^2)} \tilde{I}_3$$

$$F_3 = -\frac{m^2}{2(E^2 - m^2)} \tilde{I}_1 - \frac{3m^2}{2(E^2 - m^2)} \tilde{I}_2 + \frac{3m^4}{2M^2(E^2 - m^2)} \tilde{I}_3$$

- Matching to V, A_0, A_1, A_2

$$V = \frac{(m+M)}{2mM} F_0, \quad A_0 = \frac{F_1}{2m} + \frac{m^2 - M^2 + Q^2}{4m^2 M} F_2 + \frac{m^2 - M^2 - Q^2}{4mM^2} F_3$$

$$A_1 = \frac{F_1}{M+m}, \quad A_2 = \frac{M+m}{2mM^2} (MF_2 + mF_3)$$

Lattice setup

Configuration	Volume	a/fm	m_π/MeV	$N_{\text{conf}} \times N_t \times N_s$
C24P29	$24^3 \times 72$	0.10524(05)	292.3(1.0)	$450 \times 72 \times 2$
C32P29	$32^3 \times 64$	0.10524(05)	293.1(0.8)	$397 \times 64 \times 2$
C32P23	$32^3 \times 64$	0.10524(05)	227.9(1.2)	$333 \times 64 \times 3$
F32P30	$32^3 \times 96$	0.07753(03)	300.4(1.2)	$360 \times 96 \times 2$
F48P21	$48^3 \times 96$	0.07753(03)	207.5(1.1)	$241 \times 48 \times 4$
G36P29	$36^3 \times 108$	0.06887(12)	297.2(0.9)	$300 \times 54 \times 2$
F48P32	$48^3 \times 144$	0.05199(08)	316.6(1.0)	$300 \times 72 \times 2$

- Four lattice spacings $[0.052 - 0.105]$ fm $\rightarrow$ continuum limit
- $m_\pi = 208 - 317$ MeV $\rightarrow$ physical pion masses
- Same lattice spacings and pion mass: C24P29 vs C32P29 $\rightarrow$ finite-volume effect

Mass and decay constant

- CLQCD, PRD111, 054504 (2025)

$$m_{D_s} = 1966.7(1.5) \text{ MeV}$$

$$f_{D_s} = 247.0(2.1) \text{ MeV}$$

- This work 2510.14478

$$m_\phi = 1021.1(7.6) \text{ MeV}$$

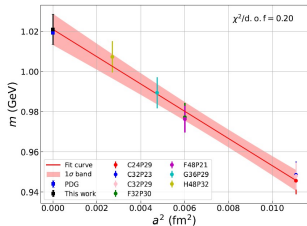
$$f_\phi = 246.2(4.1) \text{ MeV}$$

much improved than

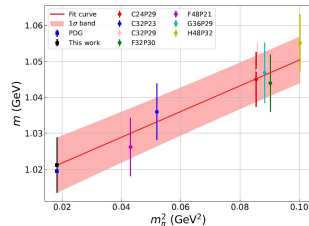
$$m_\phi = 1032(16) \text{ MeV}$$

$$f_\phi = 241(18) \text{ MeV}$$

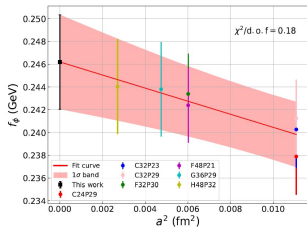
HPQCD, PRD90, 074506 (2014)



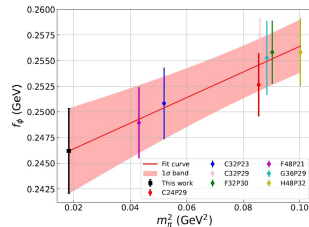
(i) ϕ mass at $m_\pi = 0.135 \text{ GeV}$.



(ii) ϕ mass at $a = 0.0 \text{ fm}$.



(iii) ϕ decay constant at $m_\pi = 0.135 \text{ GeV}$.



(iv) ϕ decay constant at $a = 0.0 \text{ fm}$.

$D_s \rightarrow \phi$ form factors

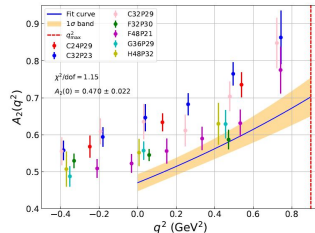
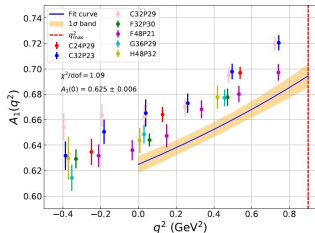
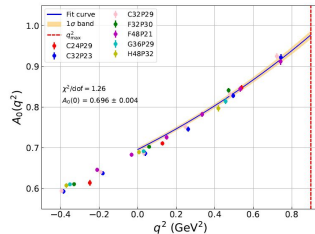
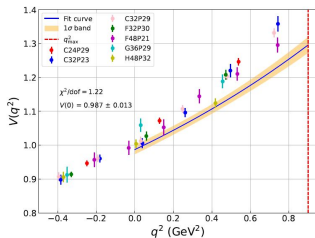
- z -expansion:

$$V/A(q^2) = \frac{1}{1 - q^2/m_{V/A}^2} \sum_{i=0}^2 (c_i + d_i a^2) \times \left[1 + f_i (m_\pi^2 - m_{\pi, \text{phys}}^2) \right] z^i$$

- Test other schemes

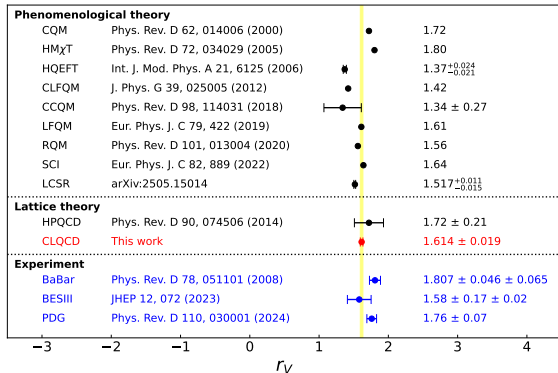
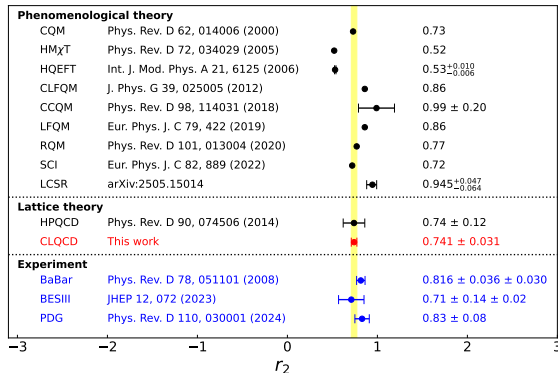
- Single pole
- Modified pole
- Phase moment form

D.L.Yao et al, PRD101,034014(2020)



$D_s \rightarrow \phi$ form factors

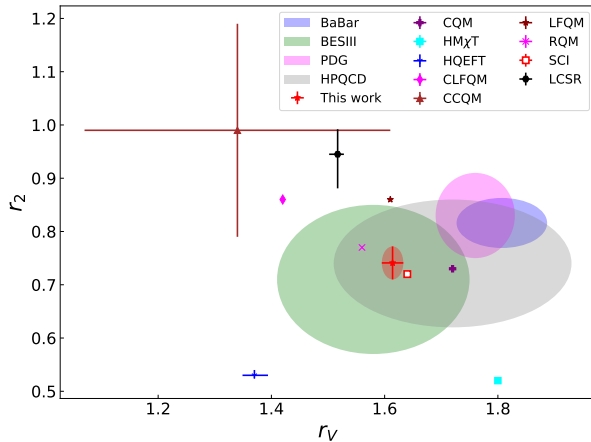
$$r_2 \equiv A_2(0)/A_1(0), \quad r_V \equiv V(0)/A_1(0)$$



G.-F.Fan, Y.M, C.Liu, Z.-F Liu, T.-X Wang, T.-H Shen, K.-Z Zhang, and L.Zhang, arXiv:2510.14478

- Precision improved by up to an order of magnitude

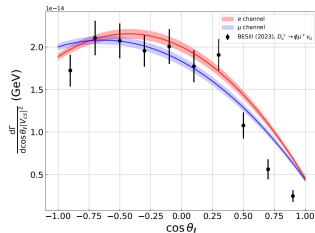
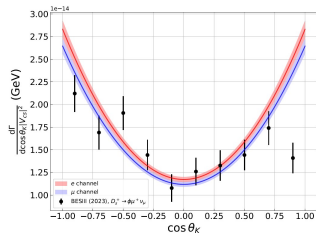
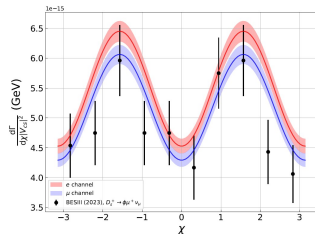
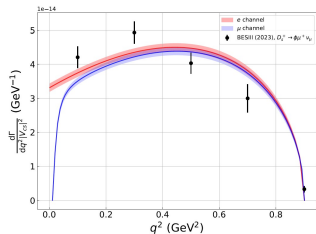
$D_s \rightarrow \phi$ form factors



- The PDG may be misleading

Differential decay width

$\mathcal{B}(D_s \rightarrow \phi \ell \nu_\ell) \times 10^{-2}$	μ channel	e channel
This work	2.351(67)	2.493(73)
BaBar [7]	—	2.61(17)
CLEO [8]	—	2.14(19)
BESIII (2018) [9]	1.94(54)	2.26(46)
BESIII (2023) [10]	2.25(11)	—
PDG [39]	2.24(11)	2.34(12)

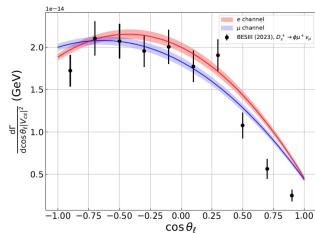
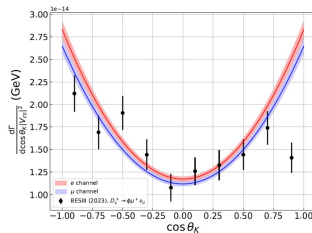
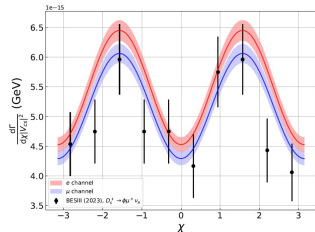
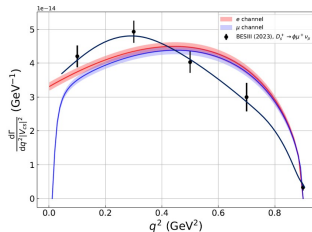


- Discrepancy in q^2 -distribution compared to BESIII measurement on $D_s \rightarrow \phi \mu \nu_\mu$

Differential decay width

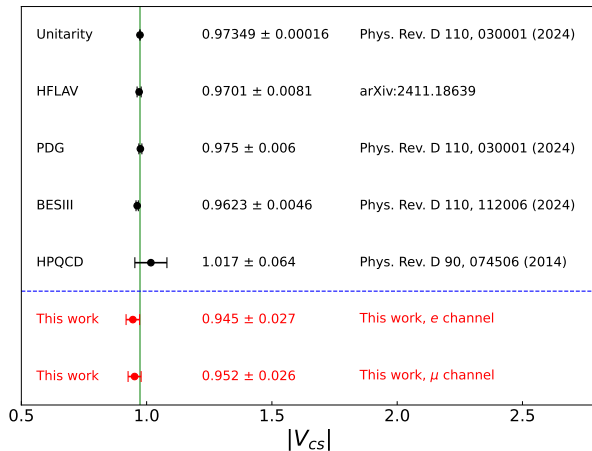
$\mathcal{B}(D_s \rightarrow \phi \ell \nu_\ell) \times 10^{-2}$	μ channel	e channel
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- $D_s \rightarrow \phi \ell \nu_\ell$ is needed



- Discrepancy in q^2 -distribution compared to BESIII measurement on $D_s \rightarrow \phi \mu \nu_\mu$

Extraction of $|V_{cs}|$



Conclusion and outlook

- Conclusion

- Scalar function method for $P \rightarrow V$ semileptonic decay
- Systematic study on $D_s \rightarrow \phi l \nu_l$ using seven CLQCD ensembles
- Most precise form factors at physical point after continuum and chiral limit

- Outlook

- **Unstable decay** of final vector particle, $P \rightarrow V(\rightarrow P_1 P_2) l \nu_l$
- Charmed meson semileptonic decay: $D \rightarrow K^*$, $D \rightarrow \rho$, etc

Thank you for attention!