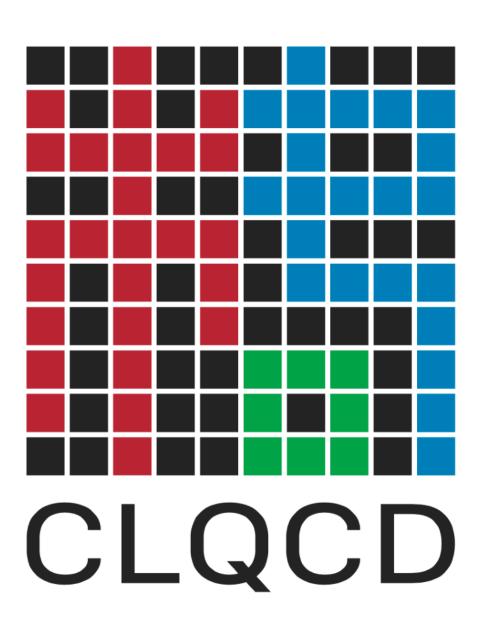
QED correction of the hadron masses with heavy quark

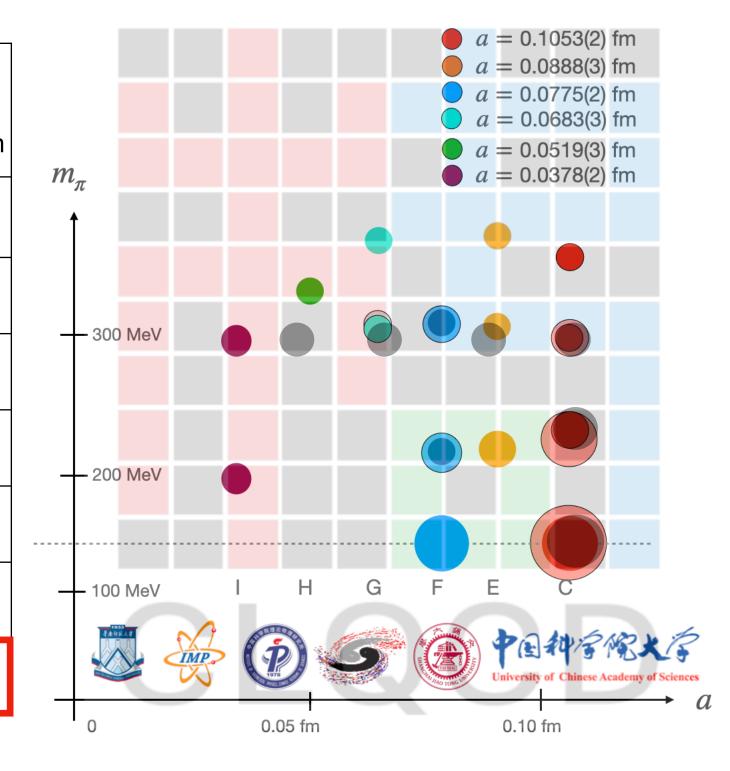


Yi-Bo Yang For CLQCD collaboration

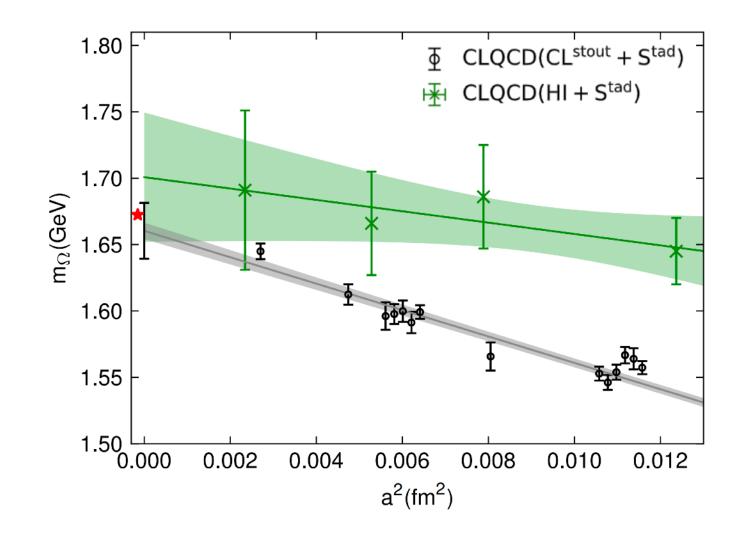




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	Country/ Region	Smallest lattice spacing	No. of physical point ensembles	Largest spacial size	No. of fermion discretization
MILC	US	0.03 fm	5	5.8 fm	1
RBC	US	0.06 fm	3	5.5 fm	1
BMW	EN	0.05 fm	15	10 fm	2
CLS	EN	0.04 fm	2	5.5 fm	1
ETM	EN	0.05 fm	5	6.3 fm	1
PACS	JP	0.06 fm	3	10 fm	1
CLQCD	CN	0.04 fm	3	6.7 fm	2

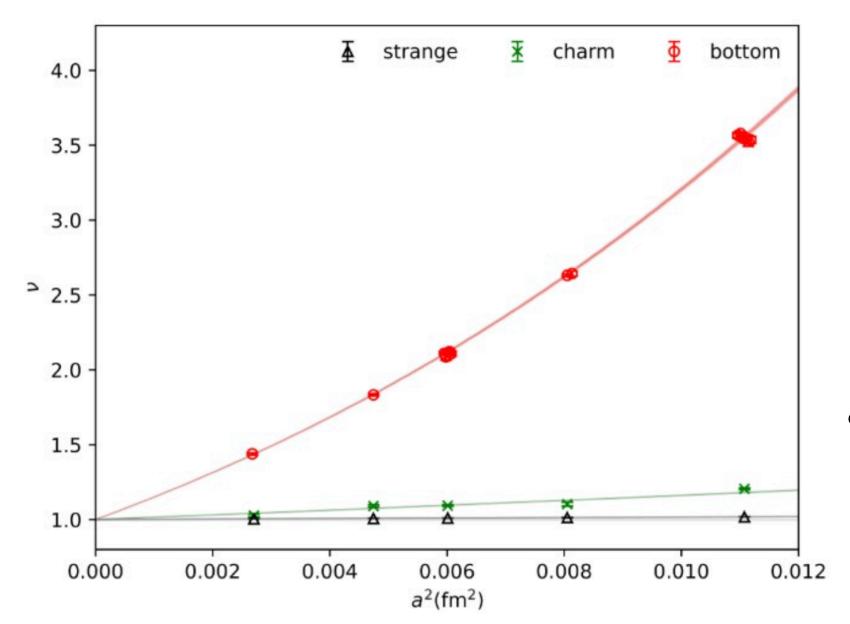


- The first
 ensemble set
 from China which
 can control most
 of the systematic
 uncertainties;
- Unique advantage on finite volume studies.

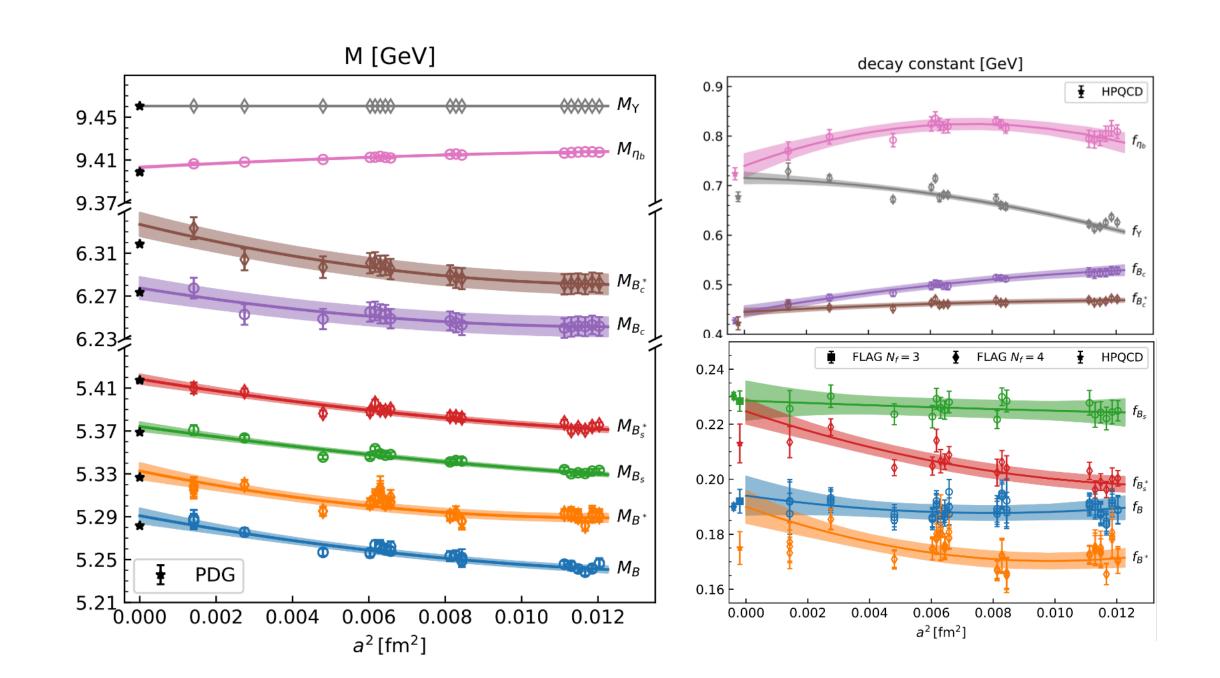


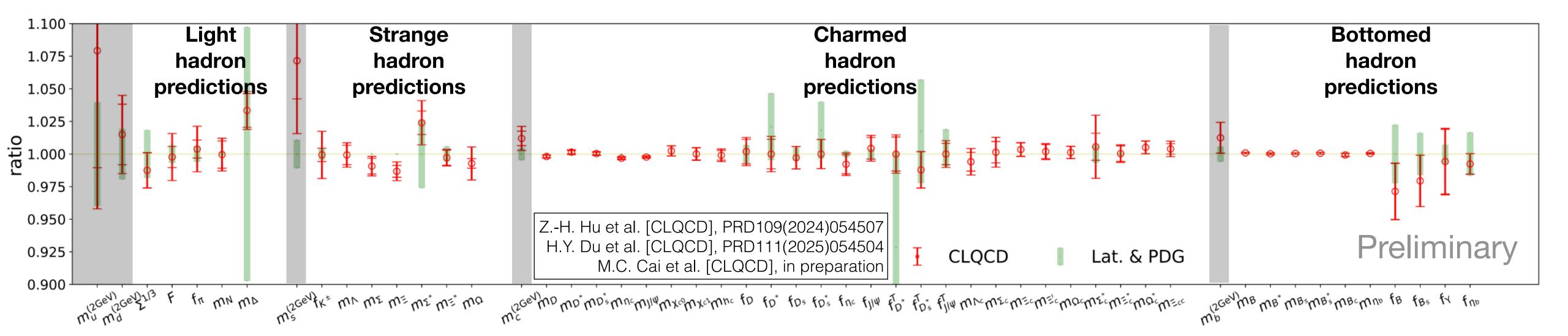
- New ensembles (HI + S^{tad}) with 2+1+1 flavor HISQ fermion can provide proper estimate of the charm sea effects;
- Compared to the current 2+1 flavor Clover fermion ensembles $(CL^{stout} + S^{tad})$, the discretization errors are also suppressed in kinds of the cases.

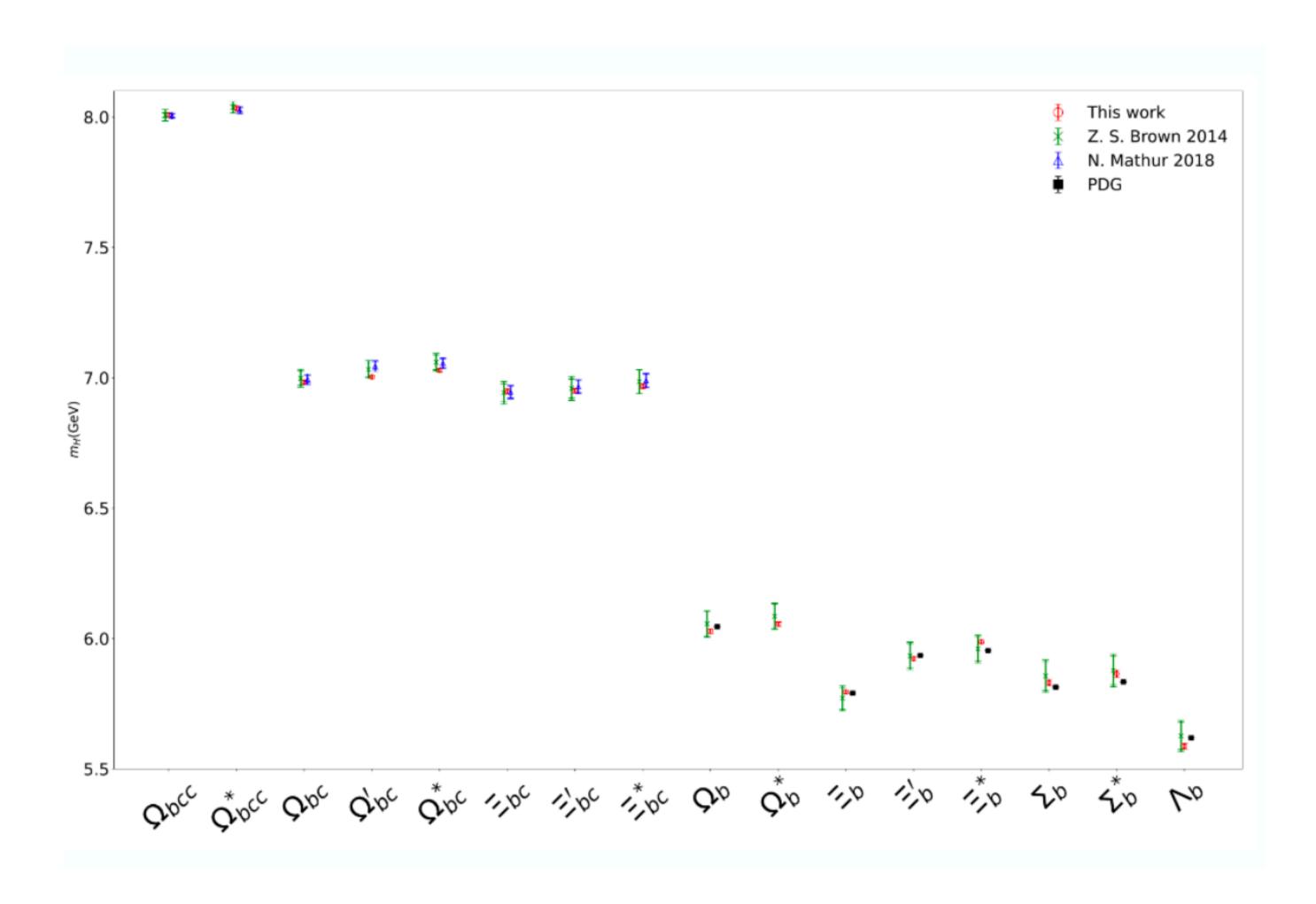
Bottom physics

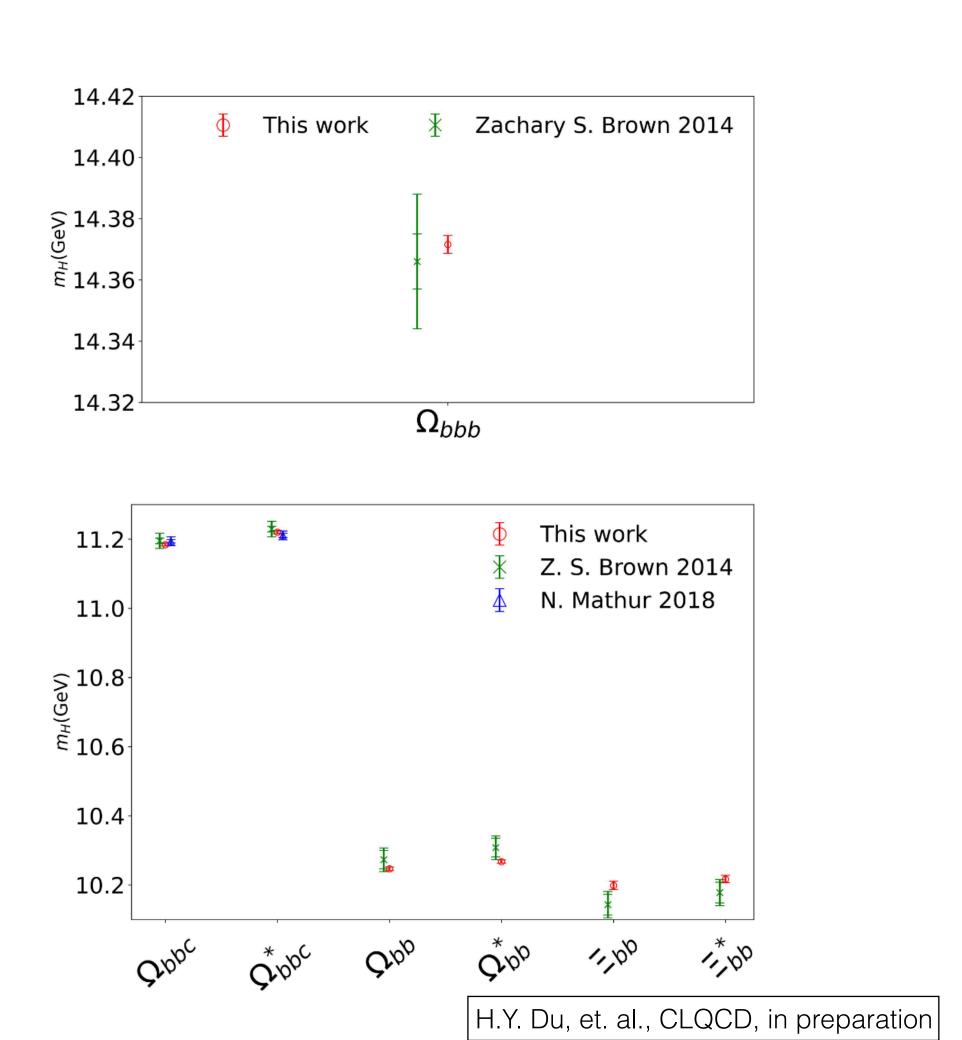


- Absorb most of the discretization errors of the bottom quark into an rescale factor along the temporal direction;
- And developed the corresponding non-perturbative renormalization.

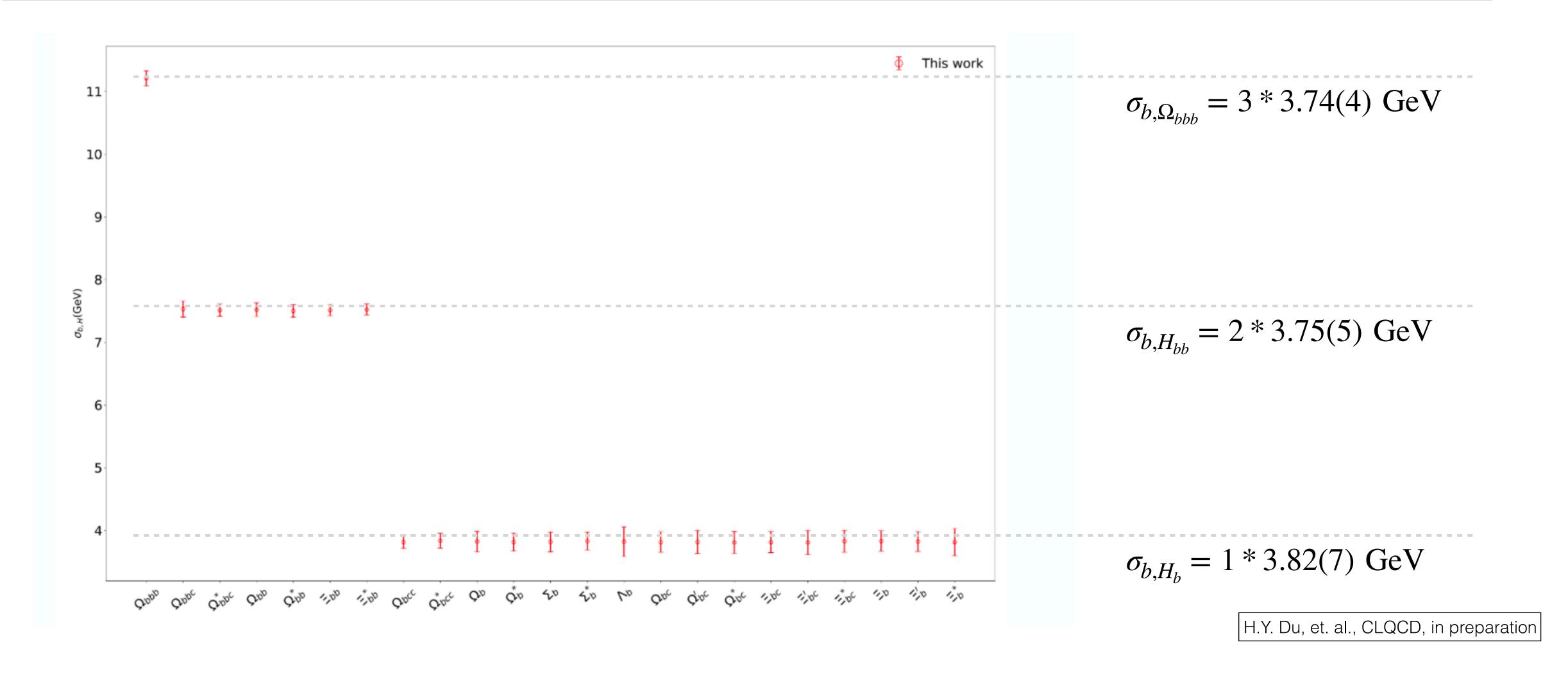








- Our prediction of the bottomed baryon masses avoid the dependences on additional nonrelativistic QCD parameters;
- And much more precise the previous lattice calculations.



- The $\sigma_{b,H}=m_b\langle \bar{b}b\rangle_H$ is the scale and scheme independent bottom quark mass contribution to the baryon mass;
- "Intrinsic scale" to make $\langle \bar{b}b \rangle = 1$ would be around $\overline{\rm MS}$ 7.6 GeV.

• The hadron mass and matrix elements in the real world require the full QCD+QED calculation. But since the lattice calculation can only reach an $\mathcal{O}(1\%)$ ($\mathcal{O}(0.1\%)$) in the heavy quark case) precision, one can expand the prediction in term of the polynomial of α and also $\delta_{\rm ISB} \equiv (m_d - m_u)/\Lambda_{\rm OCD}$:

$$\mathcal{M}^{\text{QCD+QED}} = \mathcal{M}^{\text{isoQCD}} + \alpha \mathcal{M}^{(0,1)} + \delta_{\text{ISB}} \mathcal{M}^{(1,0)} + \mathcal{O}(\alpha^2, \alpha \delta_{\text{ISB}}, \delta_{\text{ISB}}^2).$$

- Naive power counting suggests that both ISB and QED corrections are 1%;
- There are kinds of known results for the ISB and QED corrections:

$$m_n - m_p = 2.52_{\text{ISB}}(29) \text{ MeV} - 1.00_{\text{QED}}(16) \text{ MeV},$$

$$m_{D^+} - m_{D^0} = 2.54_{\text{ISB}}(13) \text{ MeV} + 2.14_{\text{QED}}(13) \text{ MeV}, \qquad m_{B^+} - m_{B^0} = -1.88_{\text{ISB}}(60) \text{ MeV} + 1.58_{\text{QED}}(24) \text{ MeV}.$$

BMWc, Science 347(2015)1452

M. Rowe, R. Zwicky, JHEP(2023)089

- ISB effect is only 0.1% for the charmed hadron, how about the QED effect?
- How to understand the sizable QED correction for the charged pion mass in the chiral limit?

$$m_{\pi^+}|_{m_q \to 0} = 0_{m_q \bar{q}q} + 0_{\mathcal{O}(\alpha_s)G^2} + 32_{\mathcal{O}(\alpha)} \text{ MeV}$$
 given $(m_{\pi^+}^2 - m_{\pi^0}^2)|_{m_q \to 0} \simeq 1000 \text{ MeV}^2$

• The QCD+QED calculation can be done under the quenched QED approximation using QED_L for the valence fermion:

$$U_{\mu}^{\text{QCD+QED}} = U_{\mu}^{\text{QCD}} e^{-iee_q A_{\mu}}, A_{\mu}(x) = \int \frac{\mathrm{d}^4 p}{(2\pi)^4} e^{-ip \cdot x} A_{\mu}(p), P_{A_{\mu}(p)}|_{\vec{p} \neq 0} \propto e^{-\frac{1}{2N_V} \hat{p}^2 A_{\mu}^2(p)}.$$

• The QED finite volume correction (FVC) of hadron masses using QED_L is independent of the hadron structure until $\mathcal{O}(1/(m_H L)^3)$:

$$\delta_{\text{QED-FVC}} m_H = e_H^2 \frac{c_1}{L} (1 + \frac{2}{m_H L} + \mathcal{O}(\frac{1}{(m_H L)^2})).$$

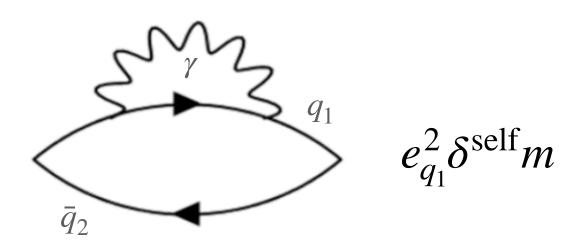
BMWc, Science 347(2015)1452

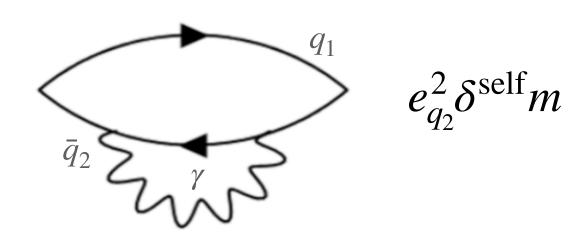
Thus the QED-FVC of neutral particles using QED_L are highly suppressed, likes that using QED_∞ .

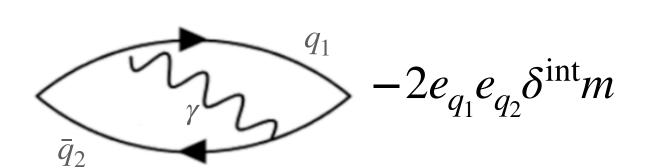
ullet One can further improve QED_L by enlarging the weights of the near-zero momentum modes:

$$\delta_{\text{QED-FVC}}^{\text{inf.imp.}} m_H = \mathcal{O}(\frac{e_H^2}{(m_H L)^2}), \; P_{A_\mu(p)}^{\text{inf.imp.}} \mid_{\vec{p} \neq 0} \propto e^{-\frac{1 + 1.4856 \delta(p^2 - \frac{4\pi^2}{L^2})}{2N_V} \hat{p}^2 A_\mu^2(p)}.$$
 Z. Davoudi et.al., PRD99(2019)034510

• One can define the neutral pion uncorrected (NPU) scheme by tuning the bare mass $m_q^b(e_q)$ of the quark with a QED charge e_q , to ensure the neutral iso-vector pseudoscalar meson mass to be the same as that using the QED-neutral quark with bare mass $m_q^b(0)$:







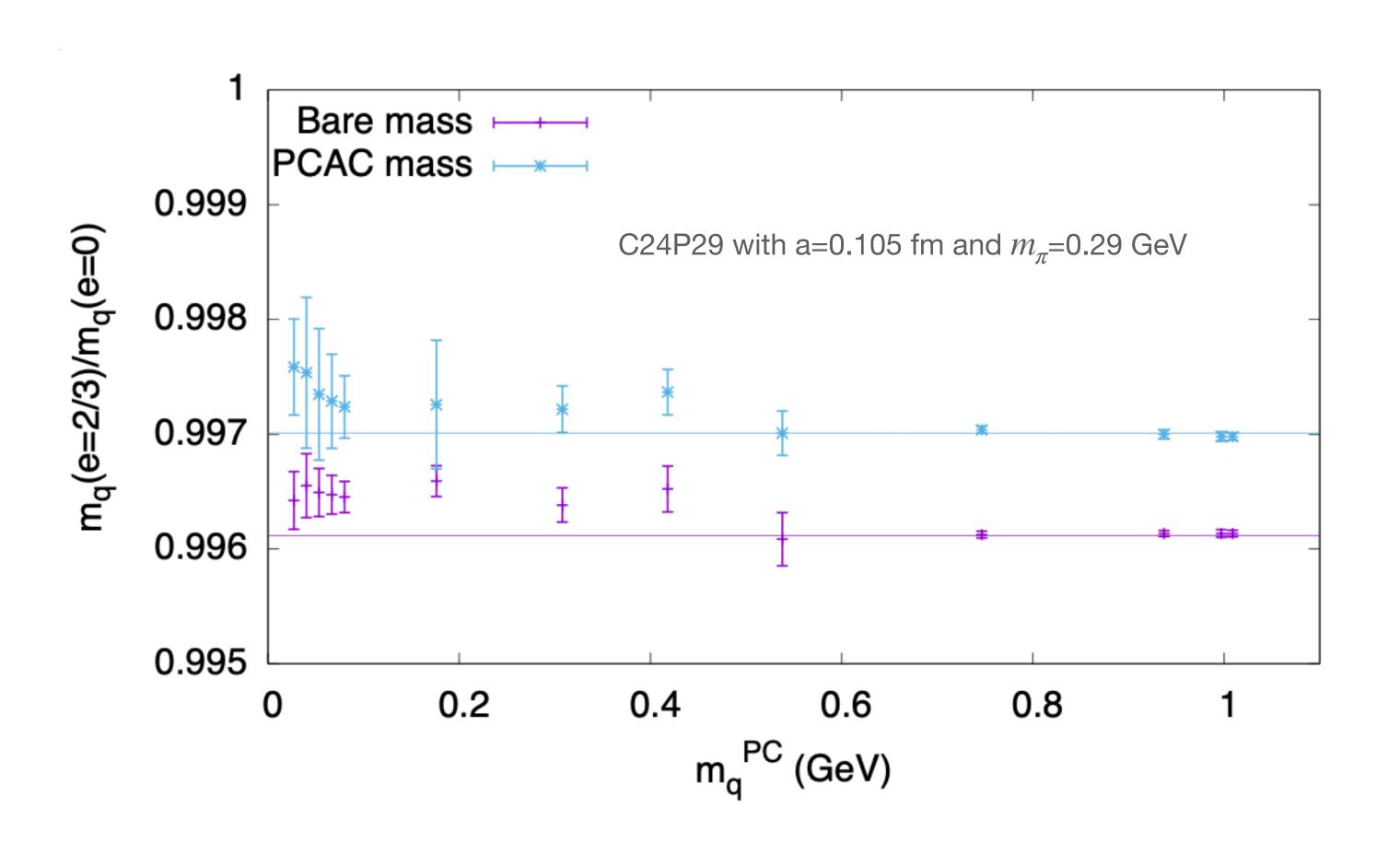
$$m_{\eta_q}^{\text{QCD+QED}}(m_q^b(e_q)) = m_{\eta_q}^{\text{QCD}}(m_q^b(0)).$$

• Using the QED quark diagram decomposition, we have $(\delta^{\text{self}} m = m^{\text{QCD+QED}} - m^{\text{QCD}})$:

$$\begin{split} \delta m_{\pi^0} &= \frac{5}{18} \delta^{\text{self}} m_{\pi^0} - \frac{5}{18} \delta^{\text{int}} m_{\pi^0} + \mathcal{O}(\alpha \alpha_s^2, \alpha^2), \\ \delta m_{\pi^+} &= \frac{5}{18} \delta^{\text{self}} m_{\pi^+} + \frac{4}{18} \delta^{\text{int}} m_{\pi^+} + \mathcal{O}(\alpha \alpha_s^2, \alpha^2), \\ \delta m_{\eta_c} &= \frac{4}{9} \delta^{\text{self}} m_{\eta_c} - \frac{4}{9} \delta^{\text{int}} m_{\eta_c} + \mathcal{O}(\alpha \alpha_s^2, \alpha^2) \\ \delta m_{\eta_b} &= \frac{1}{9} \delta^{\text{self}} m_{\eta_b} - \frac{1}{9} \delta^{\text{int}} m_{\eta_b} + \mathcal{O}(\alpha \alpha_s^2, \alpha^2) \end{split}$$

- $\delta^{\mathrm{self}} m_{\eta_q}$ requires from the QED UV renormalization and a matching condition.
 - The NPU scheme defines $\delta^{\rm self} m_{\eta_q} = \delta^{\rm int} m_{\eta_q} \text{, and then}$ $\delta m_{\pi^+} = \frac{1}{2} \delta^{\rm int} m_{\pi^+}$ $+ \mathcal{O}(\alpha \alpha_s^2, \alpha^2)$

Quark mass renormalization



- For the u-type quarks:
- 1. PCAC quark mass is changed by 0.30(5)%,
- 2. Bare quark mass $m_q^b m_q^{\rm crti}$ is changed by 0.39(5)%,

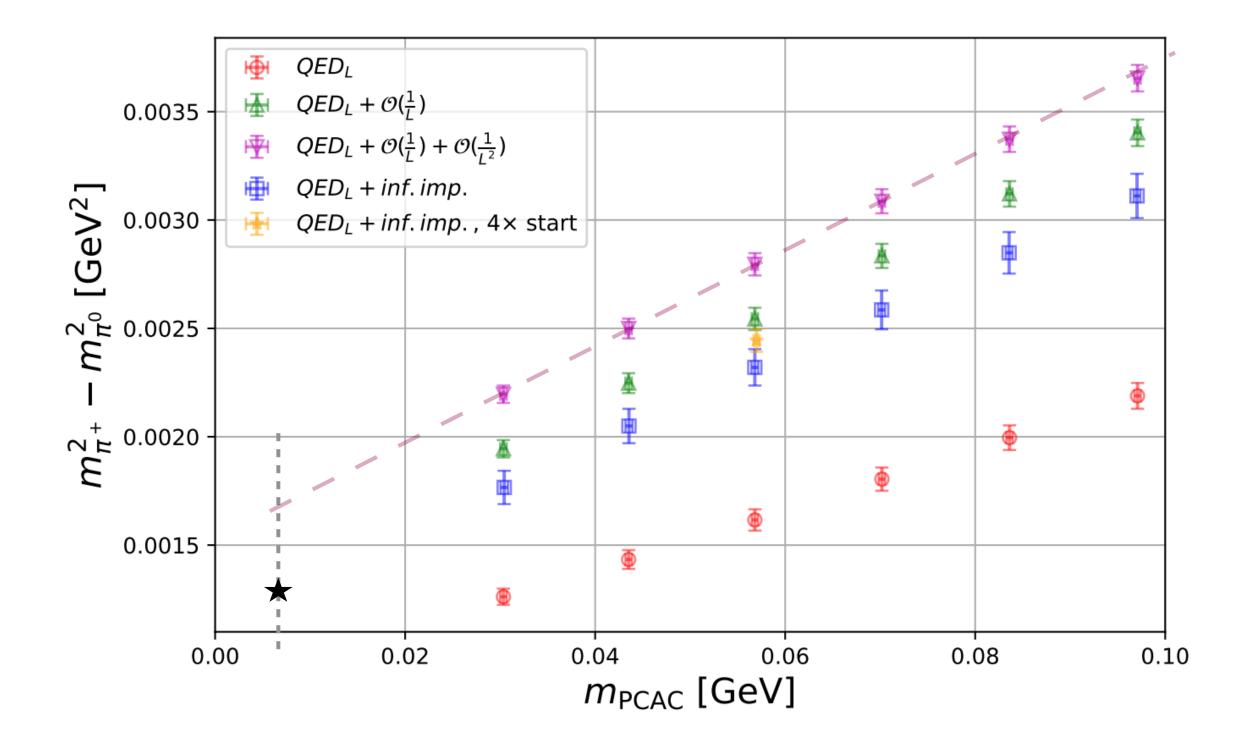
using the NPU scheme, with their difference coming from the additive chiral symmetry breaking of the clover fermion;

- The correction would be quark mass independent and more statistics is necessary to verify it.
- The perturbative calculation shows that the QED UV scale dependence is 0.12% from a=0.105 fm to a=0.052 fm.
- That of the d-type quarks will be suppressed by a factor of 4.

$$\delta_{\text{QED-FVE}} m_H^2 = e_H^2 m_H \frac{2c_1}{L} (1 + \frac{2}{m_H L} + \mathcal{O}(\frac{1}{(m_H L)^2}))$$

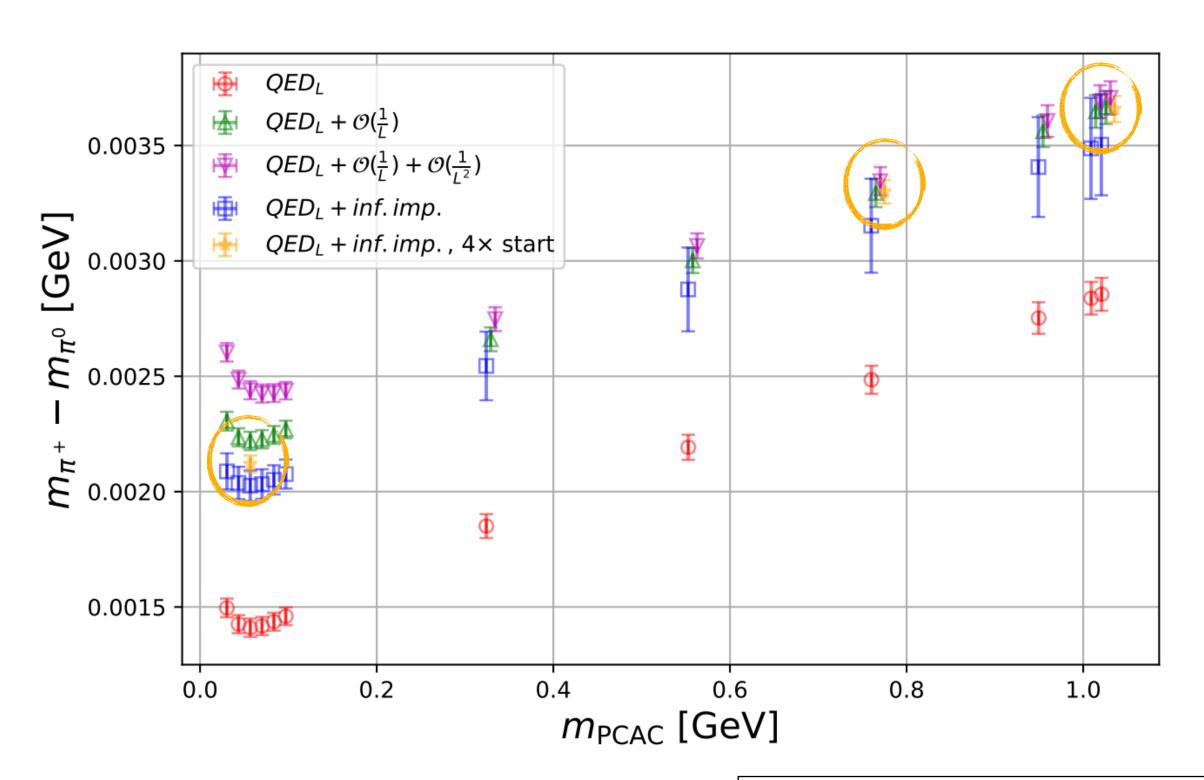
$$\uparrow \qquad \uparrow \qquad \uparrow$$

$$\mathcal{O}(\frac{1}{L}) \qquad \mathcal{O}(\frac{1}{L^2})$$



C24P29 with a=0.105 fm and m_{π} =0.29 GeV

- Ignoring the iso-spin breaking (ISB) correction, the mass difference between $m_{\pi^+}[\bar{q}(e_q)\gamma_5 q(e_q)]$ and $m_{\pi^0}[\bar{q}(-e_q)\gamma_5 q(e_q)]$ is a pure QED correction;
- With the NNLO QED-FVC, we have $m_{\pi^+}^2 m_{\pi^0}^2 = 1.5 \times 10^3 \ \mathrm{MeV^2}$ after the chiral extrapolation of the valence quark mass, at a=0.105 fm and m_{π}^{sea} =0.29 GeV, which is not far away from the **physical value**.
- The inferred improved QED_L result can include the $\mathcal{O}(1/L)$ FVC automatically and closes to the NLO QED_L result.



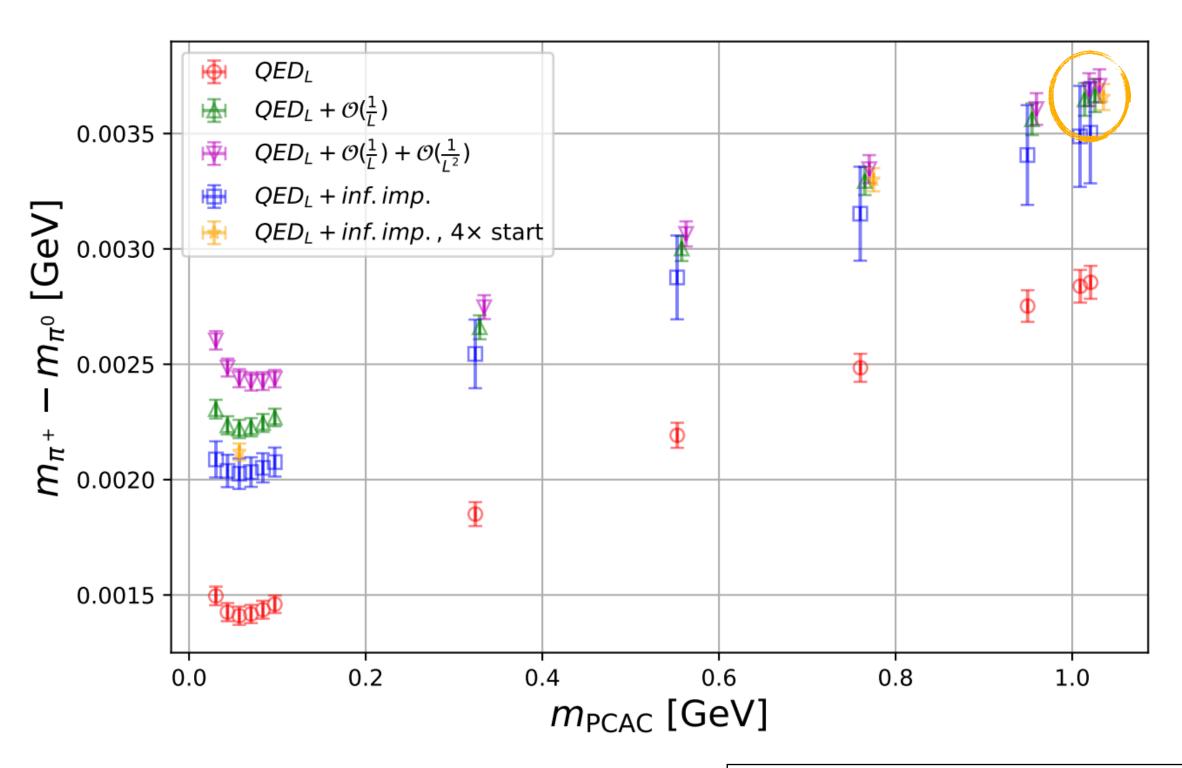
Y.Y. Liu et al. [CLQCD], in preparation

C24P29 with a=0.105 fm and m_π =0.29 GeV

- With heavier quark mass, the NNLO QED-FVC becomes negligible and then only the NLO QED-FVE matters;
- The inferred improved ${\rm QED}_L$ result becomes closer to the NLO QED-FVE result with heavier quark mass, while the statistical uncertainty is larger;
- The agreement becomes even better after the statistics of the inferred improved QED_L result is improved.

QED corrections

Impact on the charm physics



Y.Y. Liu et al. [CLQCD], in preparation

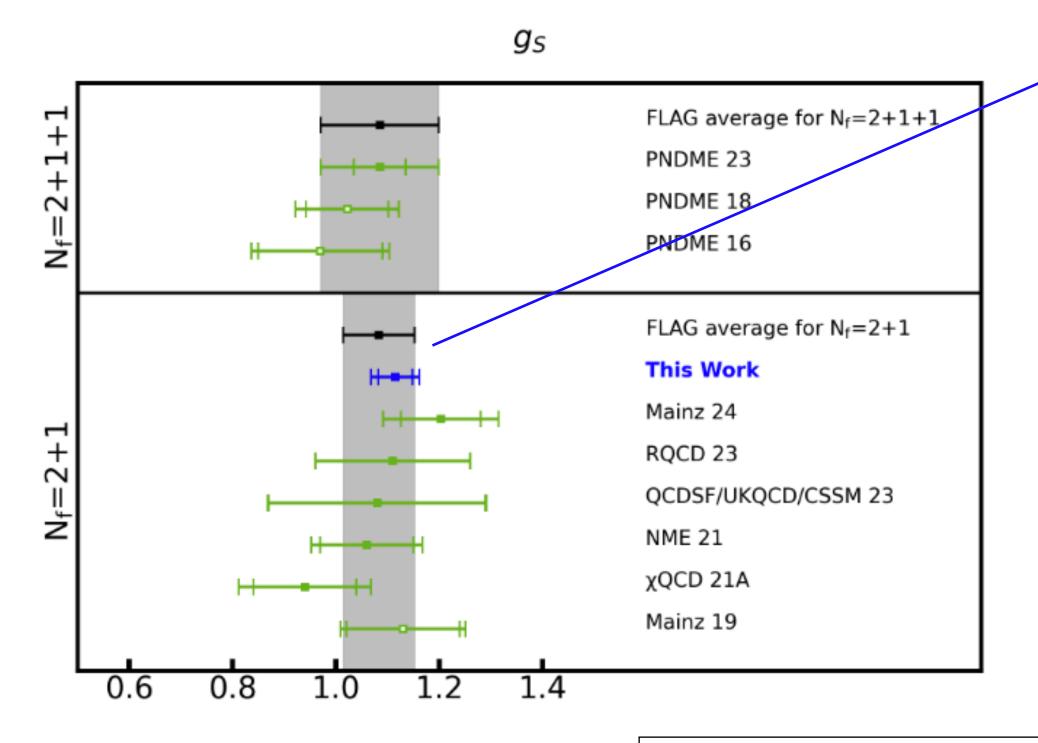
C24P29 with a=0.105 fm and m_{π} =0.29 GeV

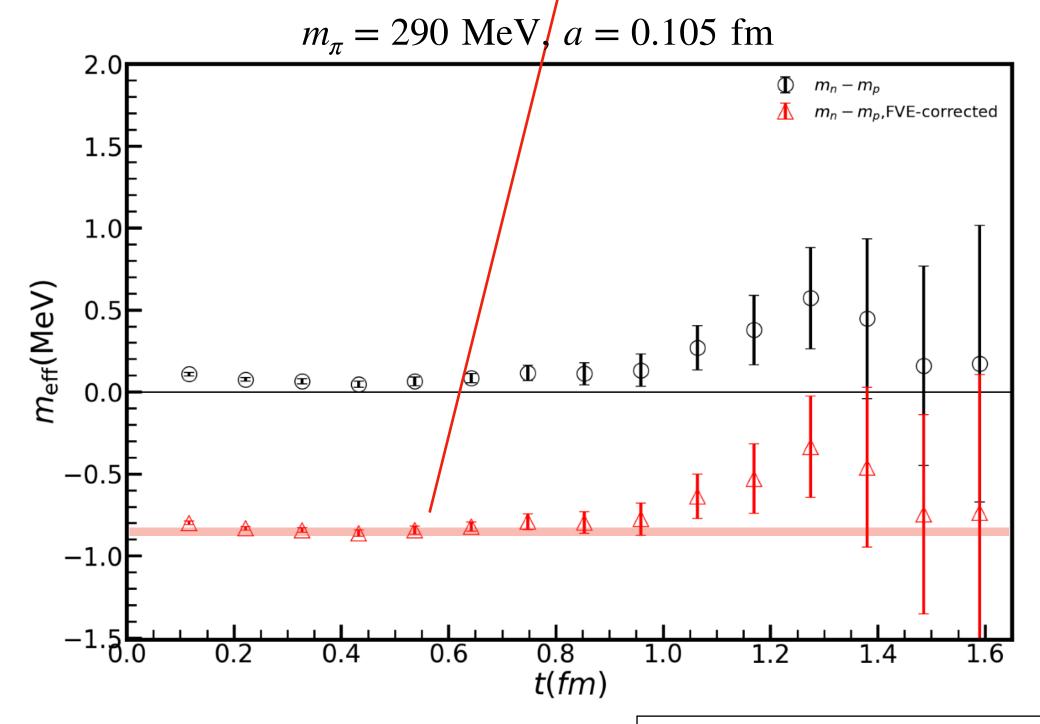
- For the charm quark:
- 1. QED interaction correction $-3.7*2\times(2/3)^2\simeq-3.2(1)$ MeV of m_{η_c} is also similar with -3.0(1) MeV obtained in 2009.07667.
- 2. QED self energy correction will be $3.7 \times (2/3)^2 \simeq 1.6(1)$ MeV which is 0.15% of the charm quark mass;
- 3. Combining the QED interaction correction $-4.7(5)e_ce_l$ MeV of m_{D^0} obtained in 2009.07667 and ignoring the light quark self energy correction, we have $\delta_{\rm QED}m_{D^0}=-0.5(2)$ MeV and $\delta_{\rm QED}m_{D^{+(s)}}=2.6(2)$ MeV which also agree with those from literature.

QCD+QED prediction

Proton-neutron mass difference:

$$\begin{split} m_n - m_p &= m_u (\frac{\partial m_n}{\partial m_u} - \frac{\partial m_p}{\partial m_u}) + m_d (\frac{\partial m_n}{\partial m_d} - \frac{\partial m_p}{\partial m_d}) + \delta^{\text{QED}} m_p^{\text{isoQCD}} = (m_d - m_u) g_S^{u-d} + \delta^{\text{QED}} m_p^{\text{isoQCD}} \\ &= (2.35(12) \text{ MeV})_{m_d - m_u} * 1.12(5)_{g_s} - 0.85(2)(?) \text{ MeV}_{\text{QED}} \\ &= 1.78(17)(?) \text{ MeV} \,. \end{split}$$

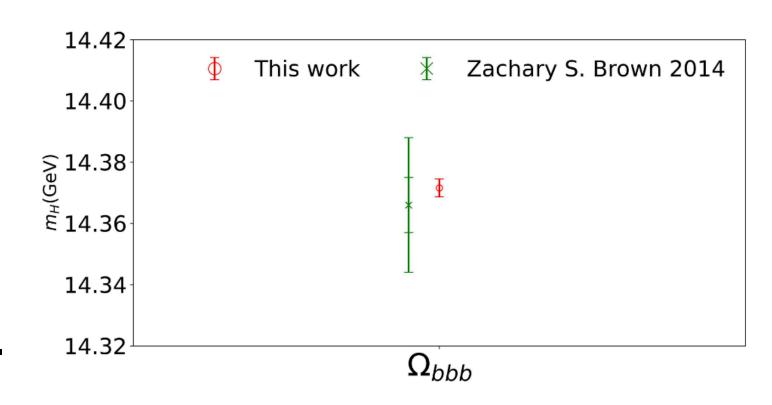


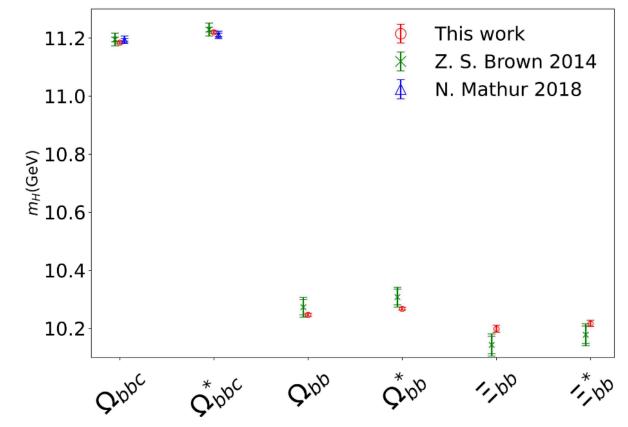


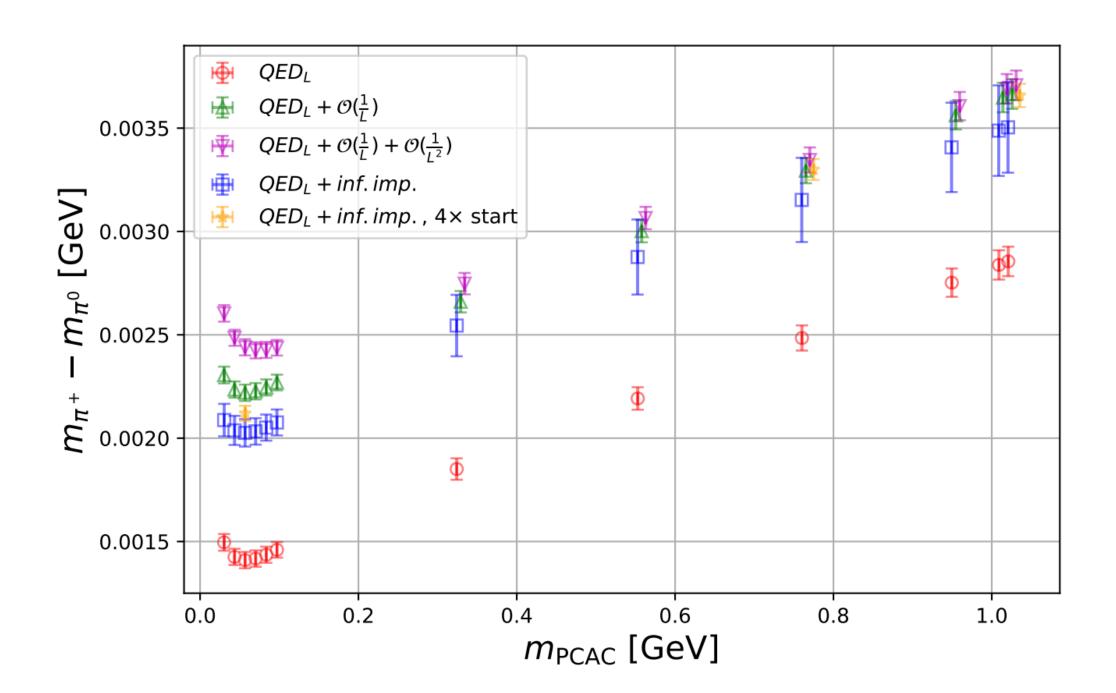
Current status

Summary

- CLQCD ensembles can now provide high precision hadron matrix elements with heavy quark;
- And also control all the systematic uncertainties of the iso-symmetric QCD.







- The QCD+QED simulation on the CLQCD ensemble C24P29 (a=0.105 fm and m_π =0.29 GeV) agrees with that in the literature reasonably;
- The QED corrections of the charm and bottom quark masses would be around 3-4 MeV.
- Full QCD+QED prediction of baryon masses from CLQCD are in progress.