# A unitary coupled-channel approach to $J/\psi$ radiative decay to pseudoscalar pairs

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#### **Outline**

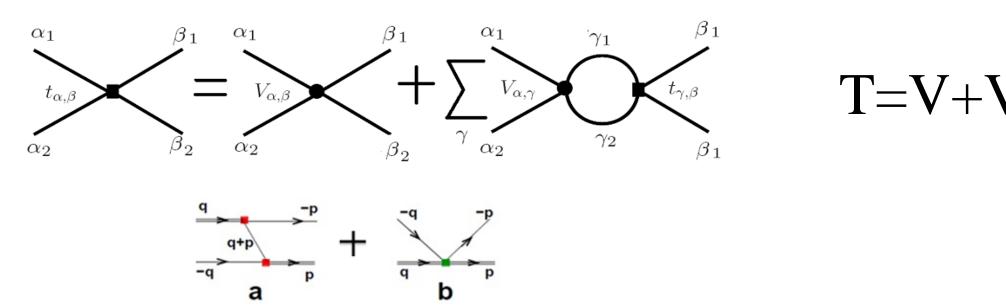
- Background
- Coupled-Channel model for  $J/\psi \rightarrow \gamma PP$
- Toy model
- Summary





#### Background

• Coupled Channel Approach: It describes interactions between different channels or pathways that a quantum system can take. Different possible states or interactions (channels) are considered simultaneously.









#### Background

- Why Coupled Channel Approach?
- Non-perturbative effect / Loop contribution:  $\frac{1}{1-x} \sim 1 + x + x^2 + \cdots$
- **Unitary:** |**S**|=**1**
- Unified: Connection between different processes, conjoint analysis
- Resonance description: BW:  $\frac{1}{s-m^2+i\Gamma m}$  VS Single channel:  $\frac{1}{E-m_0-Re(\Sigma(E))+i\Gamma(E)/2}$
- •

$$T=V+VGT$$



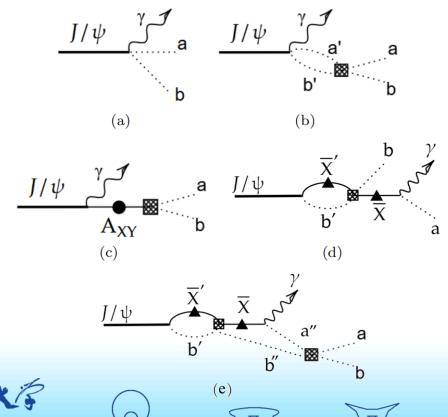




 $J/\psi \rightarrow \gamma PP$ : Three body final state, but electromagnetic interaction is much smaller than strong interaction, thus  $\gamma P$  re-scattering is negligible.

$$J/\psi \rightarrow \gamma PP$$
:

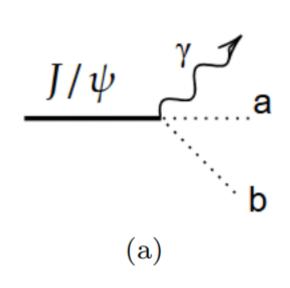
- a. bare coupling
- b. PP re-scattering
- $c. X \rightarrow PP$
- d.  $\bar{X}P$  re-scattering and then  $\bar{X} \to \gamma P$
- e. d plus PP re-scattering







 $J/\psi \rightarrow \gamma PP$ : a. bare coupling



$$\mathcal{M}_{J/\psi \to \gamma \alpha}^{(a)} = \sum_{L_{1}, L_{2}, L_{1}^{z}, L_{2}^{z}, S, S^{z}} \langle 1S_{\gamma}^{z}, L_{2}L_{2}^{z} | SS^{z} \rangle$$

$$\times \langle SS^{z}, L_{1}L_{1}^{z} | 1S_{J/\psi}^{z} \rangle$$

$$\times u_{J/\psi \to \gamma \alpha}^{L_{1}L_{2}S} Y_{L_{1}L_{1}^{z}} (-\hat{q}_{\gamma}) f_{(J/\psi, \gamma)}^{L_{1}} (q_{\gamma})$$

$$\times Y_{L_{2}L_{2}^{z}} (\hat{q}_{\alpha}) f_{\alpha}^{L_{2}} (q_{\alpha}),$$

Lorentz structure

$$\epsilon_{\mu\nu\alpha\beta}\epsilon^{\mu}_{J/\psi}\epsilon^{*\nu}_{\gamma}P^{\beta}_{(J/\psi)}\tilde{t}^{(1)\alpha}$$

$$f_{\alpha}^{L}(k) = \frac{(1 + \frac{k^{2}}{\Lambda^{2}})^{-2-L/2} \cdot (k/m_{0})^{L}}{\sqrt{E_{\alpha_{1}}(k)E_{\alpha_{2}}(k)}}$$

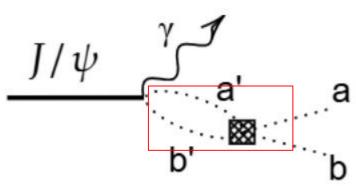








 $J/\psi \rightarrow \gamma PP$ : b. PP re-scattering



$$\mathcal{M}_{J/\psi \to \ \gamma\alpha}^{(b)} = \sum_{L_1,L_2,L_1^z,L_2^z,S,S^z} Y_{L_1L_1^z}(-\hat{q}_{\gamma}) f_{(J/\psi,\gamma)}^{L_1}(q_{\gamma})$$

$$\times \langle 1S_{\gamma}^z, L_2L_2^z | SS^z \rangle \langle SS^z, L_1L_1^z | 1S_{J/\psi}^z \rangle$$

$$\times \sum_{\beta} u_{J/\psi \to \ \gamma\beta}^{L_1L_2S} M_{\beta\beta}^{L_2}(E) \tilde{t}_{\beta\alpha}^{L_2}(E) f_{\alpha}^{L_2}(q_{\alpha}) Y_{L_2L_2^z}(\hat{q}_{\alpha})$$
**h**

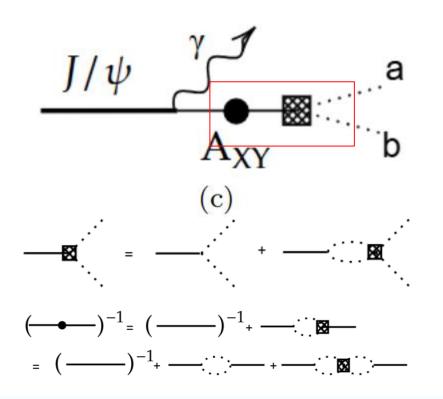








 $J/\psi \rightarrow \gamma PP$ : c.  $X \rightarrow PP$ 



$$\mathcal{M}_{J/\psi \to \gamma \alpha}^{(c)} = \sum_{L_{1}, L_{2}, L_{1}^{z}, L_{2}^{z}, S, S^{z}} Y_{L_{1}L_{1}^{z}}(-\hat{q}_{\gamma}) f_{(J/\psi, \gamma)}^{L_{1}}(q_{\gamma})$$

$$\times \langle 1S_{\gamma}^{z}, L_{2}L_{2}^{z} | SS^{z} \rangle \langle SS^{z}, L_{1}L_{1}^{z} | 1S_{J/\psi}^{z} \rangle$$

$$\times \sum_{XY} u_{J/\psi \to \gamma X}^{L_{1}L_{2}S} A_{XY}(E) \mathcal{G}_{Y\alpha}^{L_{2}}(q_{\alpha}, E) Y_{L_{2}L_{2}^{z}}(\hat{q}_{\alpha})$$

Additional Vertex for  $J/\psi \rightarrow \gamma X$ 

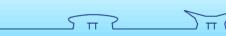
$$\mathcal{G}_{X\alpha}^{L}(k;E) = G_{X\alpha}^{L}(k) + \sum_{\gamma} g_{X\gamma}(E)\tilde{t}_{\gamma\alpha}^{L}(E)f_{\alpha}(k)$$

$$(A^{-1})_{XY}(E) = \delta_{XY}(E - m_X) - \Sigma_{XY}^0(E) - \Sigma_{XY}^I(E)$$

$$\Sigma_{XY}^{0}(E) = \sum_{\gamma} \int dq q^{2} \frac{G_{X\gamma}(q) G_{Y\gamma}(q)}{E - \omega_{\gamma}(k) + i\epsilon}, \quad \Sigma_{XY}^{\mathbf{I}}(E) = \sum_{\alpha, \beta} g_{X\alpha}(E) \tilde{t}_{\alpha\beta}^{L}(E) g_{Y\beta}(E)$$



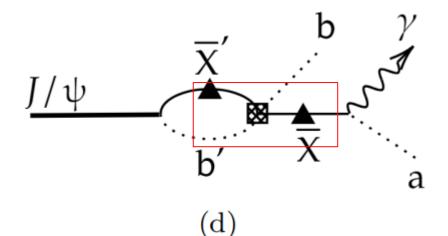






 $I/\psi \rightarrow \gamma PP$ : d.  $\overline{X}P$  re-scattering and then  $\overline{X} \rightarrow \gamma P$ 

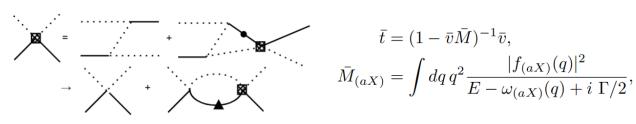
Additional Vertex for  $J/\psi \rightarrow \bar{Y}c$ 



$$\mathcal{M}_{J/\psi \to \gamma \alpha}^{(d)} = \sum_{L_1, L_1^z} \sum_{\bar{Y}c, S_{\bar{X}}^z} Y_{L_1 L_1^z}^* (\hat{q}_c) \langle S_{\bar{Y}} S_{Y}^z, L_1 L_1^z | 1 S_{J/\psi}^z \rangle$$

$$\times u_{J/\psi \to \bar{Y}c}^{L_1} \bar{M}_{\bar{Y}c} (m_{J/\psi})$$

$$\times \sum_{L_2, L_2^z} \sum_{\bar{X}, S_{\bar{X}}^z} \langle S_{\bar{X}} S_{\bar{X}}^z, L_2 L_2^z | 1 S_{J/\psi}^z \rangle Y_{L_2 L_2^z} (\hat{q}_b)$$



$$\bar{t} = (1 - \bar{v}\bar{M})^{-1}\bar{v},$$

$$\bar{M}_{(aX)} = \int dq \, q^2 \frac{|f_{(aX)}(q)|^2}{E - \omega_{(aX)}(q) + i \, \Gamma/2}$$

$$\times \sum_{L_3L_3^z} \langle 1S_{\gamma}^z, L_3L_3^z | S_{\bar{X}}S_{\bar{X}}^z \rangle Y_{L_3L_3^z}(-\hat{q}_{\gamma})$$

$$\times \frac{f_{\bar{X}\to\gamma a}^{L_3}(q_a)}{E(q_a) - m_{\bar{X}} + i \ \Gamma/2},$$

 $\times \bar{t}_{\bar{Y}_c\bar{X}_b}(m_{J/\psi})f^{L_1}_{\bar{\mathbf{Y}}_b}(q_b)$ 

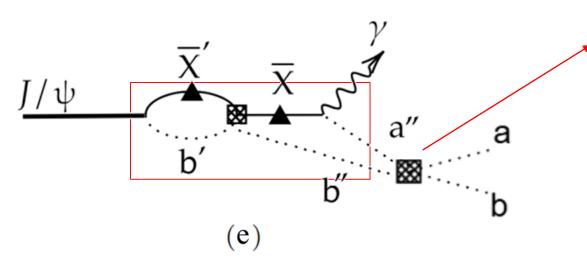


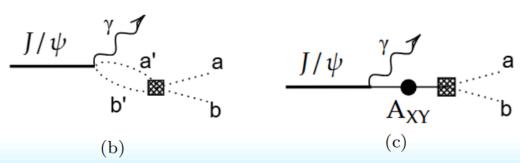






 $J/\psi \rightarrow \gamma PP$ : e. d plus PP re-scattering





#### **Full Amplitude**



$$\tilde{T}_{\alpha\beta}^{L}\left(k,k';E\right) = \tilde{V}_{\alpha\beta}^{L}\left(k,k';E\right) + \sum_{\gamma} \int dq \, q^{2} \frac{\tilde{V}_{\alpha\gamma}^{L}(k,q;E)\tilde{T}_{\gamma\beta}^{L}\left(q,k';E\right)}{E - \omega_{\gamma}(q) + i\epsilon}$$

$$\tilde{V}_{\alpha\beta}^{L}(k,k',E) = V_{\alpha\beta}^{L}(k,k') + \sum_{X} \frac{G_{X\alpha}^{L}(k)G_{X\beta}^{L}(k')}{E - m_{X} + i\epsilon}$$

$$\tilde{T}^{L} = t^{L} + T^{L}$$

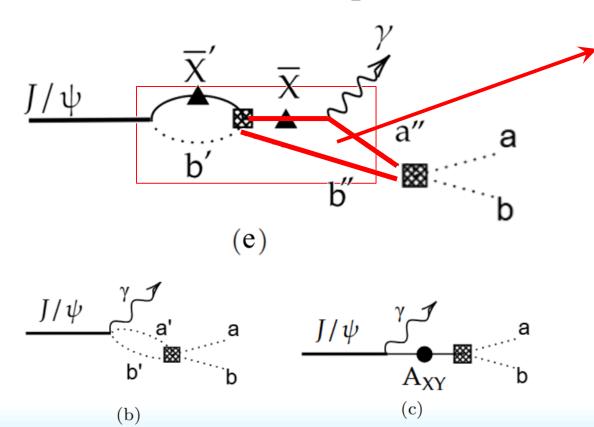








 $J/\psi \rightarrow \gamma PP$ : e. d plus PP re-scattering



#### **Triangle Singularity**

We should consider it additionally, here we do not express explicitly.





#### Toy model

 $J/\psi \to \gamma K_S K_S$  and  $J/\psi \to \gamma \pi^0 \pi^0$ 

Two 0<sup>+</sup> bare states,  $f_0 \sim$  m = 1.221 GeV and  $f_0' \sim$  m = 1.451 GeV One 1<sup>+</sup> bare state for  $\gamma K_S$  channel,  $K_1 \sim m = 1.403$  GeV and  $\Gamma = 0.173$  GeV

$R$ $\{L\}$	$f_0$ $\{0\}$	$f_{0}^{'} \{1\}$
$\overline{m_R}$	1.221	1.457
$g_{(R,KK)}$	0.810	-0.710
$g_{(R,\pi\pi)}$	-0.740	0.910
$u_{(J/\psi,\gamma R)}$	0.300	0.500
$\gamma PP$		
$\overline{v_{KK,KK}}$	0.980	-
$v_{KK,\pi\pi}$	0.870	-
$v_{\pi\pi,\pi\pi}$	0.980	-
$u_{(J/\psi,\gamma KK)}$	0.800	-
$u_{(J/\psi,\gamma\pi\pi)}$	0.800	-

$\overline{\mathbb{R}\{L\}}$	$K_1 \{0\}$	
$\overline{m_R}$	1.403	
$\Gamma_R$	0.174	-
$\bar{u}_{(J/\psi,KR)}$	1.400	-
$\bar{v}_{(KR,KR)}$	1.000	-
$\bar{g}_{(R,\gamma K)}$	1.000	_

$\overline{\mathbb{R}\{L\}}$	$M_{pole}(\mathrm{GeV})$	RS
$f_0\{0\}$	1.225 - 0.010 i	(uu)
$f_{0}^{'}\{0\}$	1.457 - 0.007 i	(uu)
$K_1\{0\}$	1.403 - 0.087 i	-









#### Toy model

$R\{L\}$	$M_{pole}(\mathrm{GeV})$	RS
$f_0\{0\}$	1.225 - 0.010 i	(uu)
$f_0^{'}\{0\}$	1.457 - 0.007 i	(uu)
$K_1\{0\}$	1.403 - 0.087 i	-

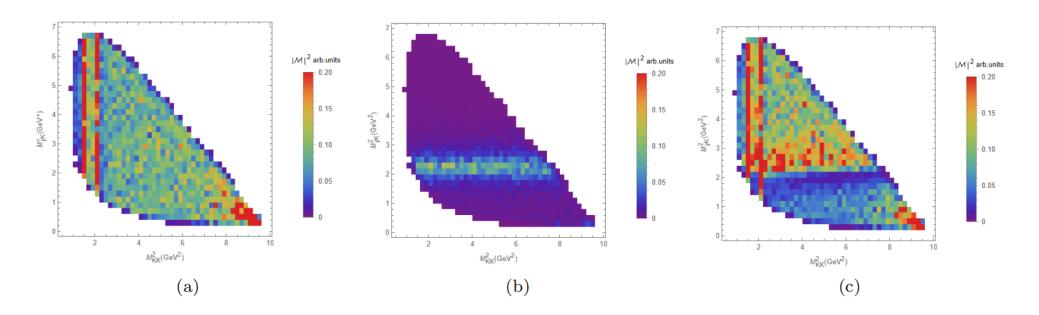
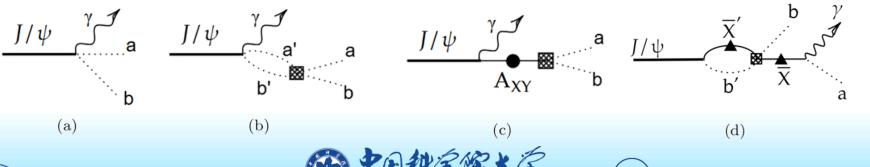


FIG. 5: The Dalitz plot on  $M^2(KK)$  and  $M^2(\gamma K)$  for different amplitudes' square.(a) Tree-level, PP-rescattering and Bare states;(b) XP-rescattering; (c) All amplitudes.



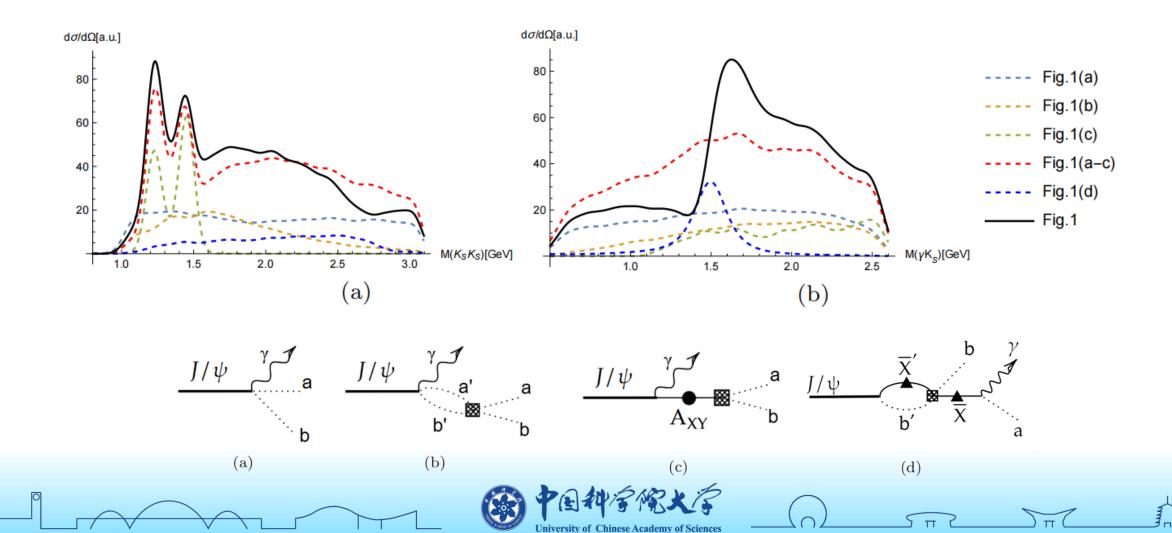






#### Toy model

$R\{L\}$	$M_{pole}(\mathrm{GeV})$	RS
$f_0\{0\}$	1.225 - 0.010 i	$\overline{(uu)}$
$f_0^{'}\{0\}$	1.457 - 0.007 i	(uu)
$K_1\{0\}$	1.403 - 0.087 i	_



#### Summary

- We give a coupled-channel approach for  $J/\psi \to \gamma PP$
- We also give numerical calculation based on a toy model
- Such formalism will be used in the BESIII analysis in future.







# Thanks very much!



