



Mellin Amplitudes for Energy Correlators

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New Frontiers of Quantum Field and Gravity

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Based on work in progress with Ian Mout, Di Wu, Anastasia Volovich

Asymptotic observables for Quantum Field Theory:

correlation function of detector operators

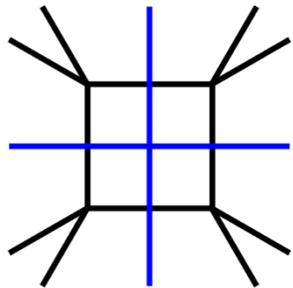
$$\langle \Psi | E(n_1) E(n_2) \dots E(n_N) | \Psi \rangle \quad \mathcal{E}(n) = \int_{-\infty}^{+\infty} du \lim_{r \rightarrow \infty} r^2 T_{0i}(t = u + r, r \vec{n}) n^i$$

integrated correlation functions || infinite sum over squared amplitudes

[Belisky, Hohenegger, Korchemsky, Sokatchev, Zhiboedov; Henn, Sokatchev, Yan, Zhiboedov; Dixon, Luo, Shtabovenko, Yang, Zhu; He, Jiang, Yang, Zhang, ...]

-- insights into higher point data very limited

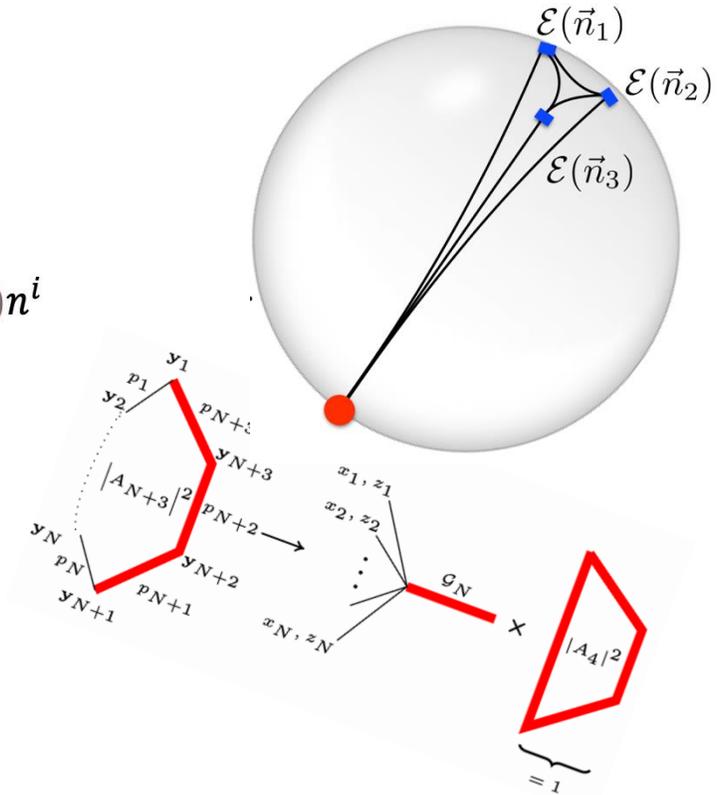
Mellin space as a new language for the multi-point asymptotic observables



$$\langle \phi_{\Delta_1}(x_1) \dots \phi_{\Delta_n}(x_n) \rangle = \int [d\delta_{ij}] M(\delta_{ij}) \prod_{i < j} \Gamma(\delta_{ij}) x_{ij}^{-\delta_{ij}}$$

Mellin formalism developed in the context of AdS/CFT correlation functions [Mack, Fitzpatrick, Kaplan, Penedones, Raju, van Rees; Nandan, Volovich, Wen]

Powerful in conformally invariant setting : conformal correlator/dual-conformal feynman integrals [Nandan, Paulos, Spradlina, Volovich]



Why Mellin space?

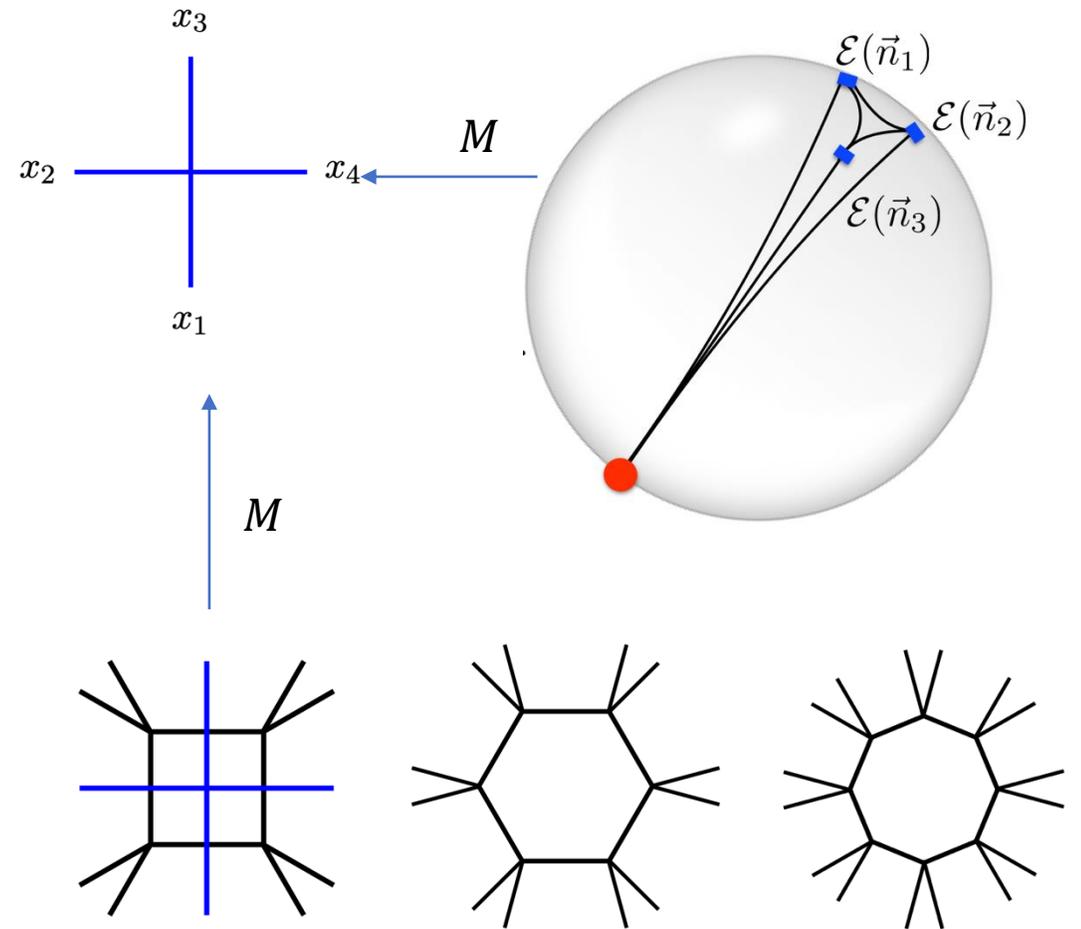
Physics of correlators transparent, allowing conformal block decomposition.

simple integro-differential relations between various types of multi-loop integrals [Drummond, Henn, Trnka; Dixon, Drummond, Henn; Ferro], complementary to standard IBPs
mapping onto basic building blocks: n-point star integral, as volumes of simplices in hyperbolic space [Schnetz, Goncharov, Bourjaily, Gardi, McLeod, Vergu, Ren, Spradlina, Vergu, Volovich...]

Profit from extensive studies on star integrals, extending knowledge on the unknowns

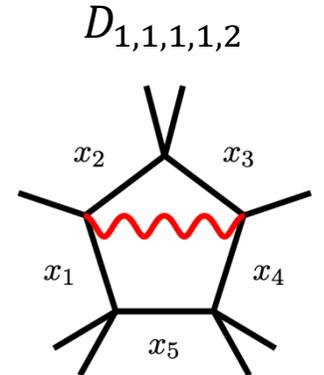
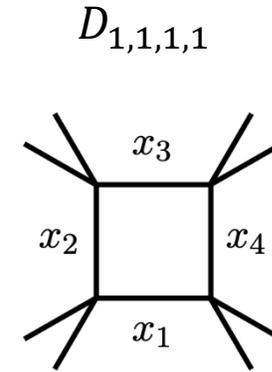
Symbology for physical observables:

What can we learn from physical constraints on the observables?



Preliminary: star formula

$$I_n \equiv \pi^{-h} \int d^d Q \prod_{i=1}^n \frac{\Gamma(\Delta_i)}{(-P_i \cdot Q)^{\Delta_i}} \quad := D_{\Delta_1, \dots, \Delta_n}$$



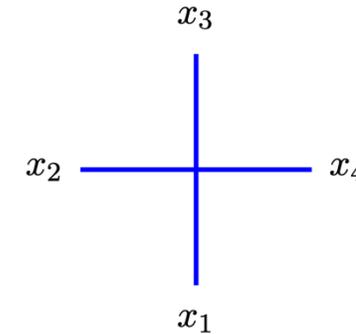
$$I_n = 2 \int_0^{+\infty} \prod_i \frac{dt_i}{t_i} t_i^{\Delta_i} \exp \left(\sum_{i<j} t_i t_j P_{ij} \right)$$

Conformal constraints

$$\delta_{ii} = -\Delta_i, \quad \sum_j \delta_{ij} = 0.$$

$$= \frac{1}{2} \int d\delta_{ij} \prod_{i<j} \Gamma(\delta_{ij}) P_{ij}^{-\delta_{ij}}$$

$$\delta_{ij} = k_i \cdot k_j$$



$$\sum_i k_i = 0.$$

$$k_i^2 = -\Delta_i$$

of $\Gamma(\delta_{ij})$: $\frac{n(n-1)}{2}$ # of independent δ'_{ij} s : $\frac{n(n-3)}{2}$

For N -point energy correlators, in collinear limit :

of (multiplicatively) independent kinematic variables: $\frac{N(N-1)}{2} - 1 = \frac{(N+1)(N-2)}{2}$ same as $(N+1)$ -gon

Preliminary: N-point EC in the multi-collinear limit

N -fold iterated integral over $1 \rightarrow N$ collinear splitting function

Toy example: ϕ^4 -theory

$$|\text{Split}|_{1 \rightarrow 3}^2 = \frac{1}{s_{123}^2} \quad s_{I \dots J} := \sum_{(i,j) \in [I,J]} x_i x_j |z_{ij}|^2$$

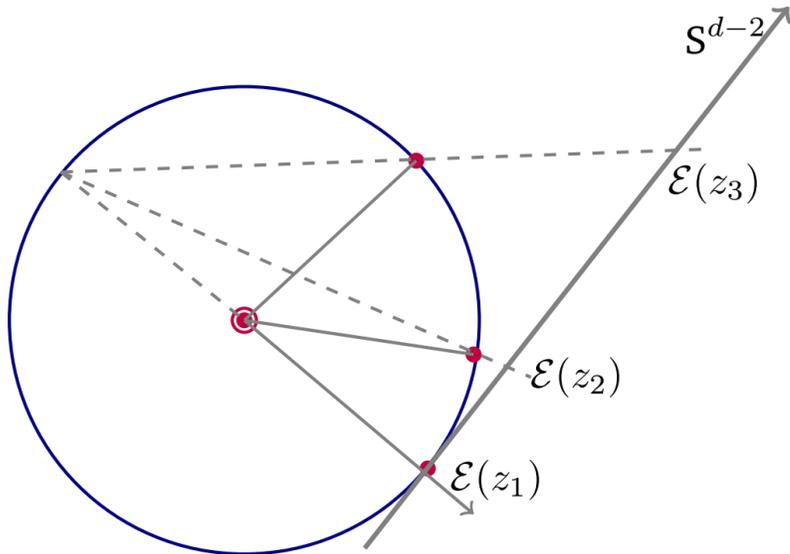
$$G_3 := |z_{12} z_{23}|^2 \int_0^1 dx_1 dx_2 dx_3 \delta(1 - x_1 - x_2 - x_3) \frac{x_1^{-J_1-1} x_2^{-J_2-1} x_3^{-J_3-1}}{s_{123}^2}$$

For energy correlators: $J_1 = J_2 = J_3 = 1 - d = -3$

Introduce an auxiliary point P_{N+1} in embedding space

$$P_i = (z_i^\mu, |z_i|^2, 1), \quad P_{N+1} = (0, 1, 0).$$

$$|z_{ij}|^2 = \frac{P_i \cdot P_j}{(P_i \cdot P_4)(P_4 \cdot P_j)} \quad i, j = 1, 2, 3$$



Writes into projective integral $x_i \rightarrow x_i P_{i4}$

$$\frac{P_{12}P_{23}}{P_{14}P_{24}P_{24}P_{34}} \times \prod_{i=1}^3 P_{i4}^{-J_i} \int \frac{d^3x}{\text{GL}(1)} \frac{\prod_i x_i^{-J_i-1}}{\left(\sum_{i,j=1}^3 x_i x_j P_{ij}\right)^2 \left(\sum_i x_i P_{i4}\right)^{-J-4}} \quad J := J_1 + J_2 + J_3$$

Introduce x_{N+1} , promote linear singularity into quadratic

$$\int d^4x \prod_{i=1}^3 x_i^{-J_i-1} x_4^{-J-4-1} \exp \left[- \sum_{i,j=1}^4 x_i x_j P_{ij} \right] := D_{-J_1, -J_2, -J_3, -J-4}$$

$$D_{3,3,3,5} \rightarrow? \int \circ \hat{\partial}^k [D_{1,1,1,1}]$$

Kinematics mapping: $u = \frac{P_{12}P_{24}}{P_{13}P_{24}} \rightarrow \frac{|z_{12}|^2}{|z_{13}|^2}$

$$v = \frac{P_{23}P_{14}}{P_{13}P_{24}} \rightarrow \frac{|z_{23}|^2}{|z_{13}|^2}$$

Solving conformal constraints

$$\delta_{14} + \delta_{24} + \delta_{34} = 3$$

$$\delta_{12} + \delta_{23} + \delta_{24} = 3$$

$$\delta_{13} + \delta_{23} + \delta_{34} = 3$$

$$\delta_{14} + \delta_{24} + \delta_{34} = 5$$

Shift mellin variables and
map onto the scalar box

$$\delta_{14} \rightarrow \delta_{14} + 1, \delta_{24} \rightarrow \delta_{24} + 1, \\ \delta_{24} \rightarrow \delta_{24} + 2, \delta_{13} \rightarrow \delta_{13} + 1$$

$$\delta_{14} + \delta_{24} + \delta_{34} = 1$$

$$\delta_{12} + \delta_{23} + \delta_{24} = 1$$

$$\delta_{13} + \delta_{23} + \delta_{34} = 1$$

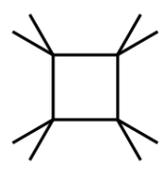
$$\delta_{14} + \delta_{24} + \delta_{34} = 1$$

$$M[D_{1,1,1,1}] = \prod \Gamma(\delta_{ij}) = \Gamma^2(\delta_u) \Gamma^2(\delta_u) \Gamma^2(1 - \delta_u - \delta_v)$$

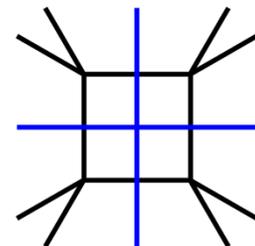
$$\delta_{12} = \delta_u, \delta_{23} = \delta_v, \delta_{13} \\ = 1 - \delta_u - \delta_v$$

$$M[D_{3,3,3,5}] \rightarrow \delta_u \delta_v (1 - \delta_v - \delta_u)^2 (2 - \delta_u - \delta_v) M[D_{1,1,1,1}]$$

$$\delta_u \rightarrow u \partial_u, \delta_v \rightarrow v \partial_v$$



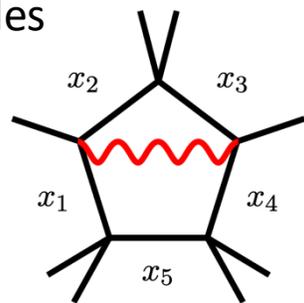
$$= P_{13}P_{24} \oint d\delta_{ij} \prod_{i<j}^4 \Gamma(\delta_{ij}) P_{ij}^{-\delta_{ij}} \implies M = 1$$



2 variables

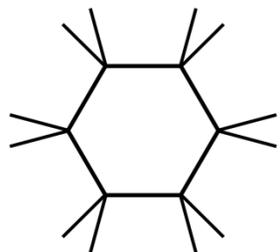
$$G_3^{\phi^4} \implies M = \left(\frac{s}{2} + 1\right) \left(\frac{u}{2} + 1\right) \left(\frac{t}{2} + 1\right)^2 \left(\frac{t}{2} + \frac{3}{2}\right)$$

5 variables



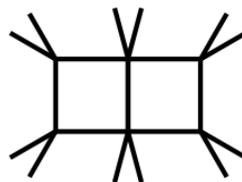
$$I_n^1 = \frac{1}{\pi^2} \int d^4Q (-Q \cdot Y) \prod_{i=1}^n \frac{1}{(-P_i \cdot Q)} \implies M = 1$$

9 variables



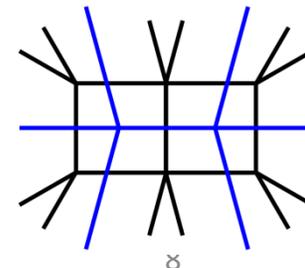
\implies

$$M = 1$$



\implies

$$M = \frac{1}{1 - s_{123}}$$



EC in N=4 super Yang Mills Theory

$$E^N C \stackrel{\text{coll.}}{=} \frac{1}{|z_{12} \cdots z_{N-1N}|^2} \int \frac{d^N x}{\text{GL}(1)} (x_1 + \cdots + x_N)^{-N} \mathcal{G}_N$$

$$+ \text{perm}(z_1, \cdots, z_N)$$

$$s_{I \dots J} := \sum_{(i,j) \in [I,J]} x_i x_j |z_{ij}|^2$$

$$N = 3: \quad \mathcal{G}_3 = \frac{s_{12} x_{123}}{s_{123} x_{12}} + \frac{s_{23} x_{123}}{s_{123} x_{23}} + \frac{x_1 x_3}{x_{12} x_{23}}$$

$$x_{I \dots J} := x_I + \cdots + x_J$$

Introduce a deforming parameter t

$$\frac{s_{12}}{s_{123} x_{12} x_{123}^2} \longrightarrow \int d^4 x dt \prod_{i=1}^3 x_1 x_2 x_4^2 \exp \left[- \sum_{i,j=3}^4 x_i x_j P_{ij} + x_4 t (P_{14} + P_{24}) + x_4 (P_{14} + P_{24} + P_{34}) \right]$$

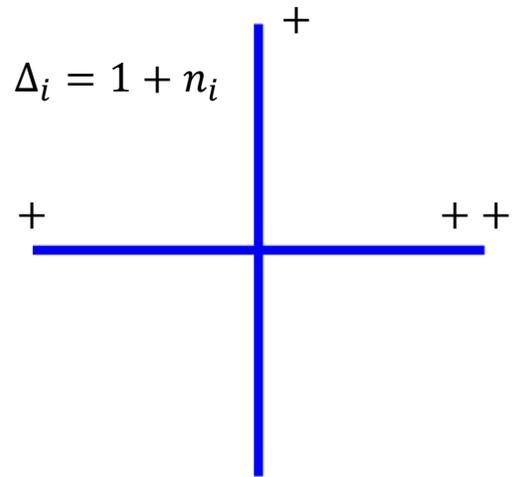
Euler type integral over t : $\rightarrow \frac{1}{\delta_{14} + \delta_{24} - 1}$

$$\frac{s_{12}}{s_{123} x_{12} x_{123}^2} \xrightarrow{M} P_{12} \times \frac{1}{\delta_{14} + \delta_{24} - 1} \times M[D_{2,2,1,3}]$$

$$\delta_{14} \rightarrow \delta_{14} + 1, \delta_{24} \rightarrow \delta_{24} + 1$$

$$\prod \Gamma(\delta_{ij})$$

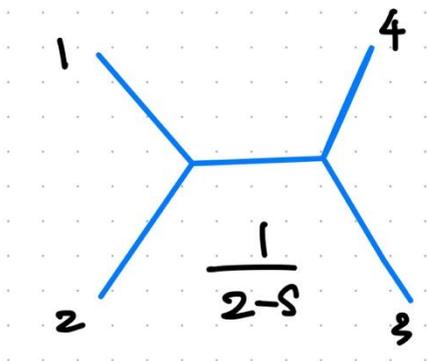
$$\sum_{j \neq i} k_i \cdot k_j = -k_i^2 = \Delta_i = 1 + n_i$$



$$M[D_{2,2,1,3}] \rightarrow \frac{1}{P_{14}P_{24}} \times \delta_{14}\delta_{24} \times M[D_{1,1,1,1}] \quad M[D_{1,1,1,1}] = \frac{1}{P_{13}P_{24}}$$

$$\frac{s_{12}}{s_{123} x_{12} x_{123}^2} \xrightarrow{M} \frac{1}{P_{14}P_{24}P_{34}} \left[\frac{P_{12}P_{34}}{P_{13}P_{24}} \right] \times \frac{\delta_{14}\delta_{24}}{\delta_{14} + \delta_{24} + 1}$$

Solving conformal constraints



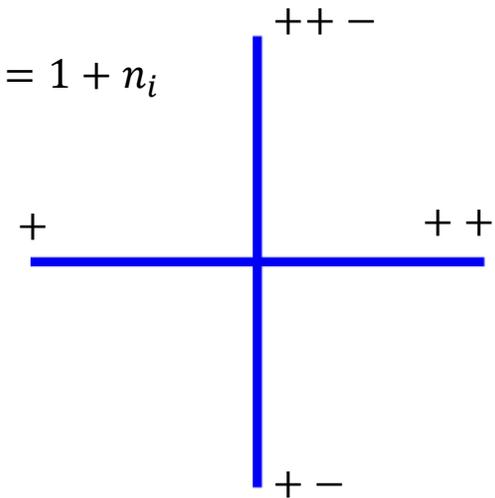
$$[u] \times \frac{\delta_v(1 - \delta_u - \delta_v)}{2 - \delta_u}$$

$$\delta_{12} = \delta_u, \delta_{23} = \delta_v, \delta_{13} = 1 - \delta_u - \delta_v$$

$$\delta_{14} \rightarrow \delta_{14} + 1, \delta_{34} \rightarrow \delta_{34} + 1, \delta_{12} \rightarrow \delta_{12} + 1, \\ \delta_{13} \rightarrow \delta_{13} - 1$$

$$\prod \Gamma(\delta_{ij})$$

$$\sum_{j \neq i} k_i \cdot k_j = -k_i^2 = \Delta_i = 1 + n_i$$

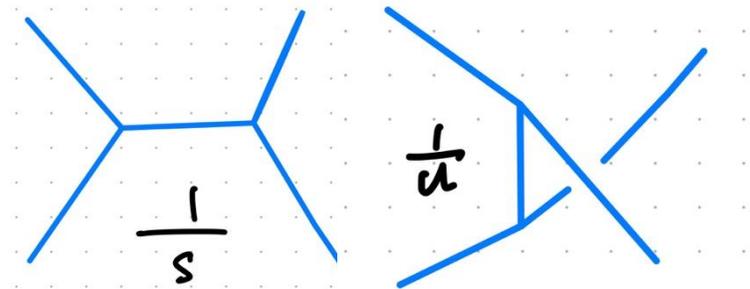


$$M[D_{2,2,1,3}] \rightarrow \frac{P_{13}}{P_{12}P_{14}P_{34}} \times \frac{\delta_{12}\delta_{14}\delta_{34}}{\delta_{13} - 1} \times M[D_{1,1,1,1}]$$

$$\frac{s_{12}}{s_{123} x_{12} x_{123}^2} \xrightarrow{M} \frac{1}{P_{14}P_{24}P_{34}} \times \frac{\delta_{12}\delta_{14}\delta_{34}}{(\delta_{14} + \delta_{24})(\delta_{13} - 1)}$$

Solving conformal constraints

$$-\frac{\delta_u^2 \delta_v}{(1 - \delta_u)(\delta_u + \delta_v)}$$



Contour shift: constructing UT integrals

$$[u] \times \frac{\delta_v(1 - \delta_u - \delta_v)}{2 - \delta_u}$$

$\delta_u \rightarrow \delta_u + 1$

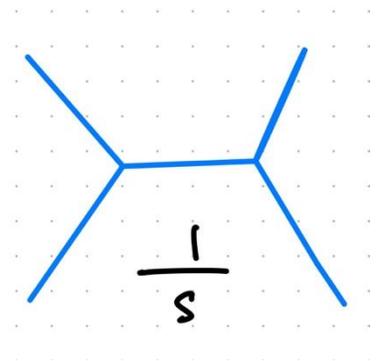
$$-\frac{\delta_u^2 \delta_v}{(1 - \delta_u)(\delta_u + \delta_v)}$$

$$\frac{\Gamma(\delta_u)^2 \Gamma(\delta_v)^2 \Gamma(1 - \delta_u - \delta_v)^2}{v^{-\delta_v} u^{-\delta_u}}$$

$$-\frac{\delta_u^2 \delta_v}{(1 - \delta_u)(\delta_u + \delta_v)} = -\frac{\delta_u \delta_v}{1 - \delta_u} + \frac{\delta_v^2 \delta_u}{(1 - \delta_u)(\delta_u + \delta_v)}$$

$\delta_v \rightarrow \delta_v - 1$

$$-[1] \times \frac{\delta_u \delta_v}{1 - \delta_u} - [v] \times \frac{\delta_u(1 - \delta_u - \delta_v)}{1 - \delta_u}$$



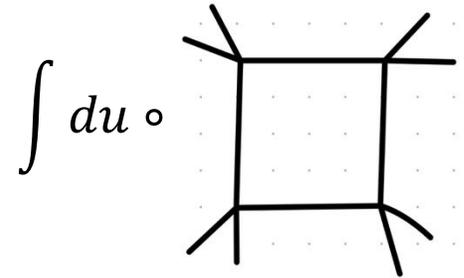
$E^3 C \xrightarrow{M}$

$$-[1] \times \left(\frac{\delta_u \delta_v}{1 - \delta_u} + \frac{\delta_u \delta_v}{1 - \delta_v} \right) - [v] \times \frac{\delta_u (1 - \delta_u - \delta_v)}{1 - \delta_u} - [u] \times \frac{\delta_v (1 - \delta_u - \delta_v)}{1 - \delta_v}$$

+ boundary constants

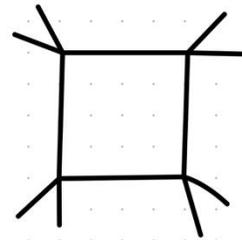
$$\hat{\partial}_u := u \partial_u, \hat{\partial}_v := v \partial_v$$

$$\mathcal{I}_{box'} := \int_0^1 dt I_4(ut, v) \xrightarrow{M} \frac{1}{1 - \delta_u}$$

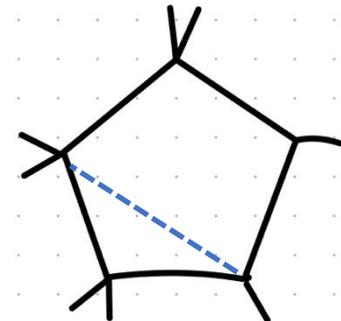


$\int du \circ$

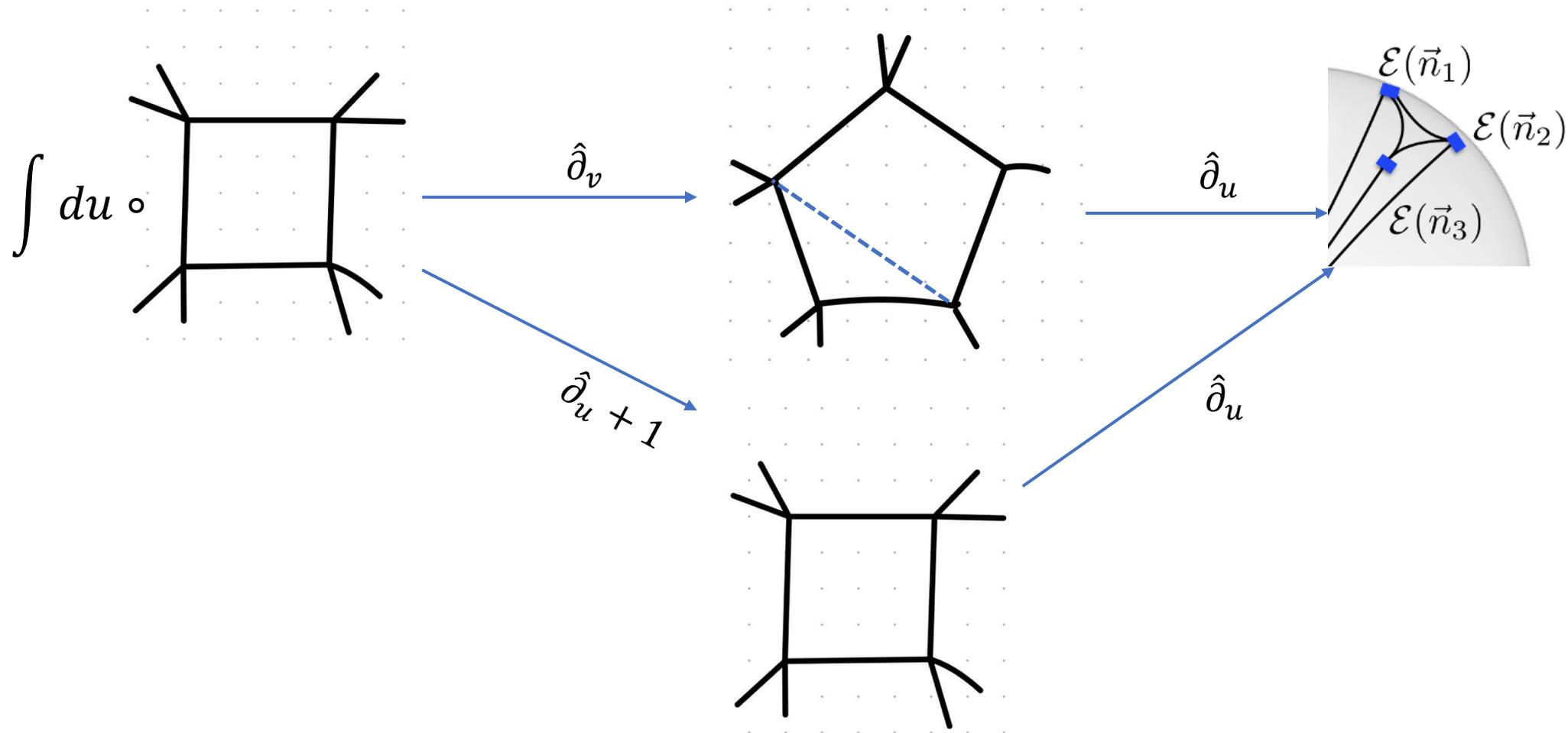
$$[\hat{\partial}_u + 1] \mathcal{I}_{box'} = I_4 \xrightarrow{M} 1$$



$$\hat{\partial}_v \mathcal{I}_{box'} = g_5^1 \xrightarrow{M} \frac{\delta_v}{\delta_u - 1} \xrightarrow{M^{-1}} \frac{x_1}{s_{123} x_{12} x_{123}}$$



E3C given by second derivatives acting on a transcendental weight 3 integral



$$\mathcal{S}[\mathcal{I}_{box'}] =$$

$$|1 - z|^2 \otimes (1 - |1 - z|^2) \otimes |1 - z|^2 - \frac{1}{2} |1 - z|^2 \otimes |z|^2 \otimes |1 - z|^2 - \frac{1}{2} |z|^2 \otimes |1 - z|^2 \otimes |1 - z|^2$$

$$+ \frac{1}{2} |z|^2 \otimes \frac{1 - z}{1 - \bar{z}} \otimes \frac{1 - z}{1 - \bar{z}} - \frac{1}{2} |1 - z|^2 \otimes \frac{z}{\bar{z}} \otimes \frac{1 - z}{1 - \bar{z}}$$

First entry $\otimes u, \otimes v$

Last entry $\otimes 1 - z, \otimes 1 - \bar{z}$

Determines the LS of weight-2 functions

Key input for bootstrapping the observable in QCD

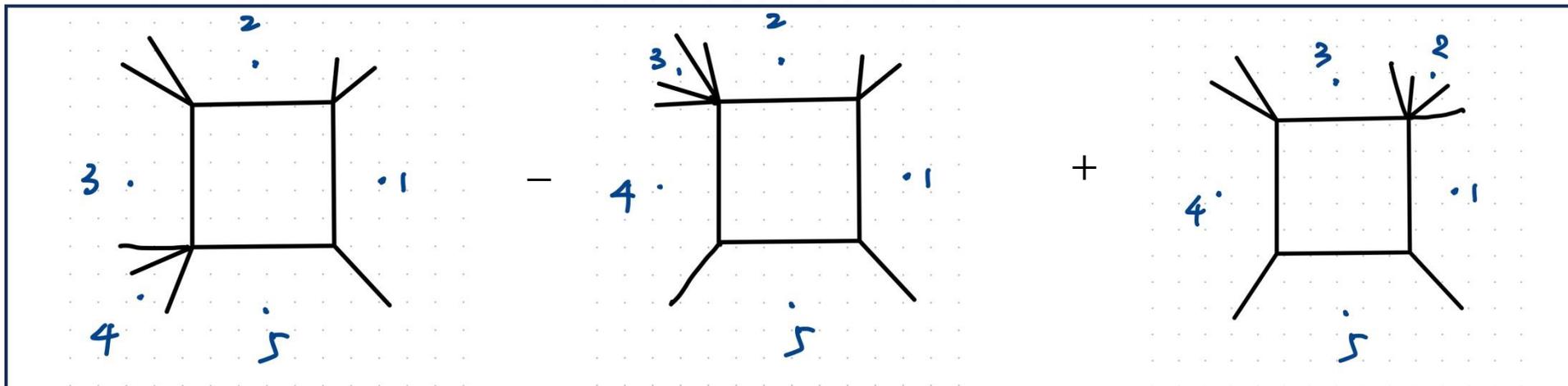
[Gong, Pokraka, Yan, Zhang]

$$\left(\begin{array}{cccc} \frac{1}{2} d \log v & \frac{1}{2} d \log \frac{1-u+v+\sqrt{\Delta}}{1-u+v-\sqrt{\Delta}} & 0 & 0 \\ d \log \frac{(1-v)^2}{u} & -d \log v & 0 & 0 \\ d \log \frac{1+u-v+\sqrt{\Delta}}{1+u-v-\sqrt{\Delta}} & d \log \frac{1-u+v-\sqrt{\Delta}}{1-u+v+\sqrt{\Delta}} & 0 & 0 \\ & & d \log v & \\ & & d \log u & \\ & & 0 & \end{array} \right) \begin{array}{l} B_1 \\ B_2 \end{array}$$

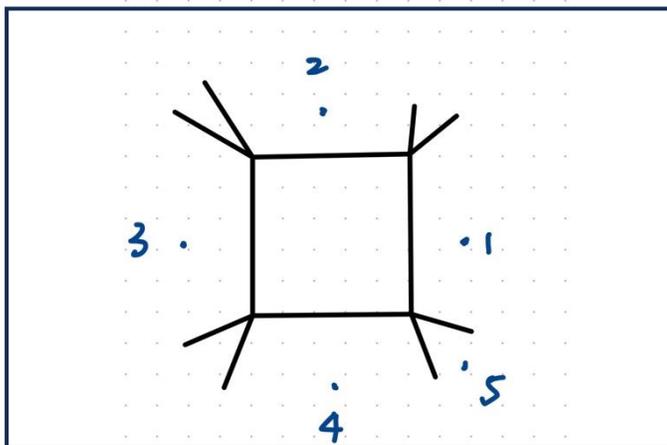
B_i : UT master integrals

$Li_2(1 - v) - \frac{1}{2} \ln v \ln u$ (D_+ - function [1912.11050])

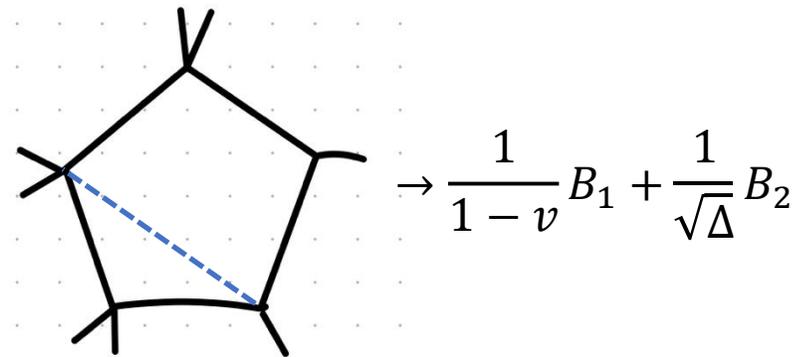
B_1



B_2



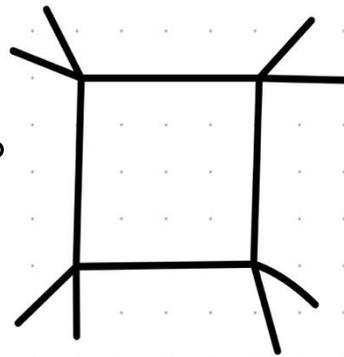
$\Phi_2(z, \bar{z})$ (D_- - function [1912.11050])



$$\begin{aligned}
G(z) = & \frac{(1 + |z|^2 + |1 - z|^2)}{2|z|^2|1 - z|^2} (1 + \zeta_2) + \frac{(-1 + |z|^2 + |z|^4 - |z|^6 - |1 - z|^4 - |z|^2|1 - z|^4 + 2|1 - z|^6)}{2|z|^2|1 - z|^2(z - \bar{z})^2} \log|1 - z|^2 \\
& + \frac{(-1 - |z|^4 + 2|z|^6 + |1 - z|^2 - |z|^4|1 - z|^2 + |1 - z|^4 - |1 - z|^6)}{2|z|^2|1 - z|^2(z - \bar{z})^2} \log|z|^2 \\
& + \frac{|z|^4 - 1}{2|z|^2|1 - z|^4} D_2^+(z) + \frac{|1 - z|^4 - 1}{2|z|^4|1 - z|^2} D_2^+(1 - z) + \frac{(|z|^2 - |1 - z|^2)(|z|^2 + |1 - z|^2)}{2|z|^2|1 - z|^2} D_2^+\left(\frac{z}{z - 1}\right) \\
& + \frac{2iD_2^-(z)}{2|1 - z|^4|z|^4(z - \bar{z})^3} p_3(|z|^2, |1 - z|^2)
\end{aligned}$$

$$\begin{aligned}
p_3(|z|^2, |1 - z|^2) = & (-1 + |z|^2 - |1 - z|^2)(1 + |z|^2 - |1 - z|^2)(-1 + |z|^2 + |1 - z|^2) \\
& \times \left[(-1 + |1 - z|^2)^2 |1 - z|^2 + |z|^6 (1 + |1 - z|^2) \right. \\
& \left. - 2|z|^4 (1 + |1 - z|^2)^2 + |z|^2 (1 + |1 - z|^2) (1 + (-5 + |1 - z|^2) |1 - z|^2) \right].
\end{aligned}$$

$$\hat{\partial}_u [-\hat{\partial}_v + v(1 + \hat{\partial}_u + \hat{\partial}_v)] \circ \int du \circ$$

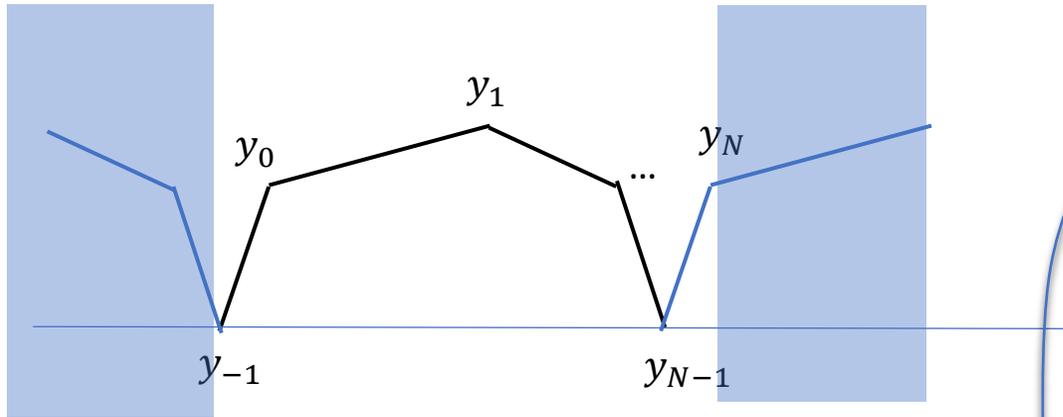


+ perms.

Moving onto Higher
Point EC

The multi-collinear limit of $|F_{N+1}|^2$ can be written in terms of dual conformal invariants including $N + 2$ points: $[y_{-1}, \dots, y_N]$

$$p_i = y_i - y_{i-1}, p_{i+N+1} := p_i, s_{i+1, \dots, k} = y_{ik}^2$$



1 \rightarrow N splitting function with
 $N = 3, 4$ (NMHV/MHV):

$$(a, b, c, d) := \frac{y_{ab}^2 y_{cd}^2}{y_{ac}^2 y_{bd}^2}$$

$$N = 3: \quad (-1, 1, 2, 4) + (-1, 3, 2, 0) + (3, 1, 0, 4)$$

In the N -particle collinear limit,

$$y_{iN}^2 \rightarrow x_{i+1} + \dots + x_{N-1}$$

$$y_{-1i}^2 \rightarrow x_1 + \dots + x_i$$

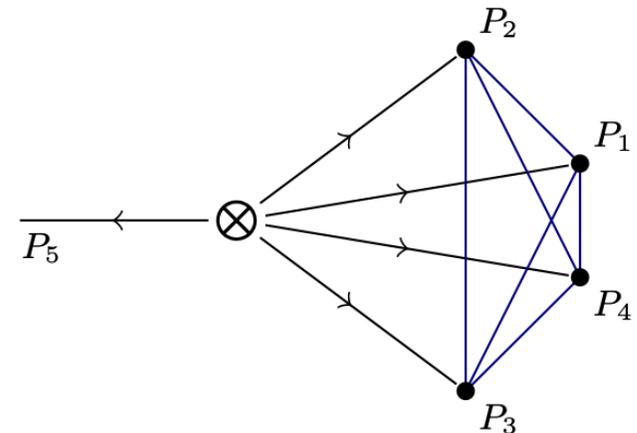
$N = 4$:

$$\begin{aligned} & -1 + (-1, 1, 2, 5) + (-1, 2, 3, 5) + (-1, 4, 3, 0) + (4, 1, 0, 5) + (-1, 3, 2, 0) + (4, 2, 1, 5) \\ & + (0, 4, 3, 1) + (0, 4, 3, 1)(-1, 1, 3, 5) + (-1, 4, 3, 1)(3, 1, 0, 5) \\ & + (-1, 4, 2, 0)(0, 2, 3, 5) + (-1, 1, 2, 4)(4, 2, 0, 5) + (-1, 3, 2, 0)(4, 2, 0, 5) + (-1, 4, 2, 0)(4, 2, 1, 5) \\ & + (-1, 4, 3, 0)(-1, 1, 2, 4) + (4, 1, 0, 5)(0, 2, 3, 5) + (-1, 4, 3, 1)(-1, 4, 2, 0) + (3, 1, 0, 5)(4, 2, 0, 5) \\ & + (-1, 4, 3, 1)(-1, 1, 2, 5) + (3, 1, 0, 5)(-1, 2, 3, 5) + (-1, 1, 2, 4)(-1, 1, 3, 5) + (0, 2, 3, 5)(-1, 1, 3, 5) \end{aligned}$$

Four – point correlator

$$G_4 := \frac{|F_5|^2}{|F_{5,MHV}|^2}$$

$$s_{I\dots J} := \sum_{(i,j) \in [I,J]} x_i x_j |z_{ij}|^2$$

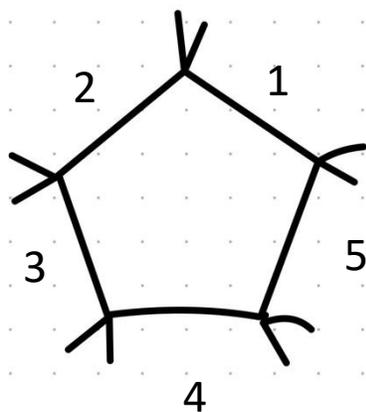


$$(-1, 4, 3, 0) = \frac{s_{123} x_{1234}}{s_{1234} x_{123}}$$

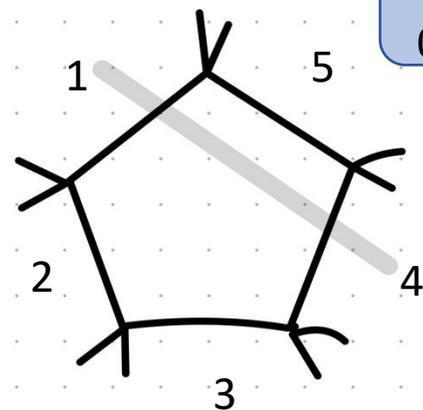
$$(-1, 4, 2, 0)(0, 2, 3, 5) = \frac{s_{12} x_{1234} x_4}{s_{1234} s_{123} x_{12} x_{34}}$$

$$(-1, 4, 3, 0)(-1, 1, 2, 4) = \frac{s_{123} s_{34} x_{1234} x_1}{s_{1234} s_{234} x_{12} x_{123}}$$

$$(0, 4, 3, 1)(-1, 1, 3, 5) = \frac{s_{1234} s_{23} x_1 x_4}{s_{123} s_{234} x_{123} x_{234}}$$

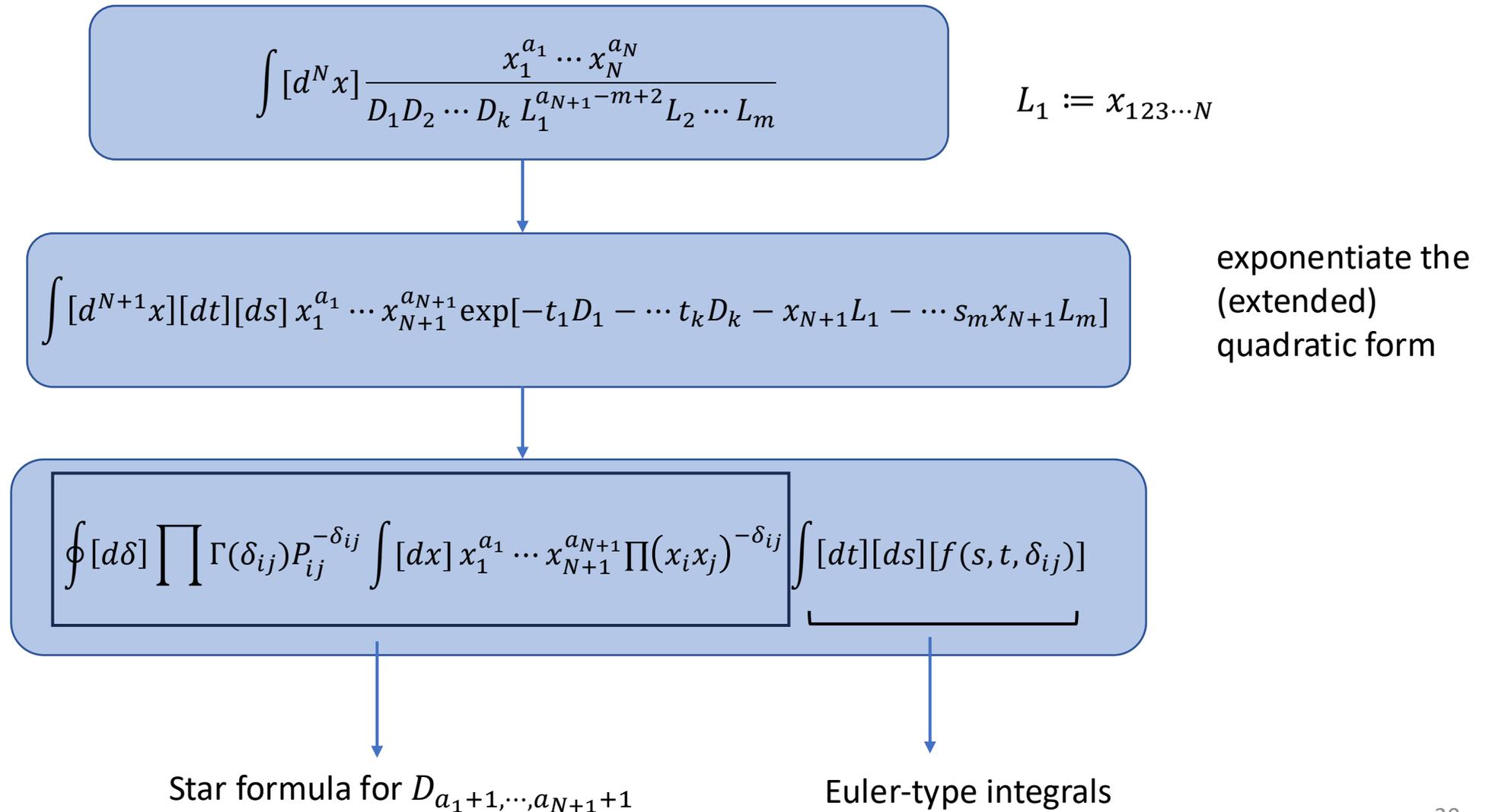


Singularities depend on 5 angles : five mass chiral pentagon as matching kernel



Singularities depend on 4 angles: impose $P_{14} = 0$ in the matching kernel

Rules for Mellin transform



Resolving overlapping singularities

$$\delta_{\sigma_I} := \sum_{(i,j) \in \sigma_I} \delta_{ij},$$

$$\sigma_I := \{(i,j) \mid P_{ij} \text{ appears in } D_I, \text{ but not in others}\}$$

Euler-type integral for a k -propagator topology:

$$\int_0^\infty [d^k t] t_1^{-\delta_{\sigma_1}} t_2^{-\delta_{\sigma_2}} \dots (t_1 + t_2)^{-\delta_{\sigma_{12}}} \dots (t_1 + t_2 + t_3)^{-\delta_{\sigma_{123}}} \dots$$

compatible poles: $\frac{1}{s_{123\dots N} s_{23\dots N\dots}}, \frac{1}{x_{123\dots N} x_{23\dots N\dots}}$

$$t_1 < t_1 + t_2 < t_1 + t_2 + t_3 \dots$$

$$\delta\left(\sum_{\sigma} \delta_{\sigma} - k\right) \times \frac{1}{\delta_{\sigma_{12}} - n_{12}} \times \frac{1}{\delta_{\sigma_{123}} - n_{123}} \times \dots$$

non-compatible poles, e.g. $\frac{1}{s_{123} s_{234}}$:

$$\int [dt_1 dt_2] t_1^{-\delta_{12} - \delta_{13}} t_2^{-\delta_{24} - \delta_{34}} (t_1 + t_2)^{-\delta_{23}} = \delta\left(\sum_{(ij) \neq (14)} \delta_{ij} - 2\right) \frac{\Gamma(1 - \delta_{12} - \delta_{13}) \Gamma(1 - \delta_{24} - \delta_{34})}{\Gamma(\delta_{23})}$$

Mapping onto Feynman Integrals

Matching Kernel: $N+1$ -point Star Integral (in even dimension)

*Search for simpler integrals (to map onto)
(e.g. higher-point star integrals with restricted kinematics)*

Generalized Euler Operator:

Differential Operator $\hat{\partial}_i := u_i \partial_{u_i}$

Integral Operator $(\hat{\partial}_i + n_i)^{-1} \rightarrow \int dt t^{n_i-1} I_{N+1}(u_i t)$

$$\frac{s_{234} x_{1234}}{s_{1234} x_{234}} + \frac{s_{12} s_{234} x_{1234} x_4}{s_{123} s_{1234} x_{234} x_{34}} + \frac{s_{23} s_{34} x_{1234}^2}{s_{123} s_{1234} x_{234} x_{34}}$$

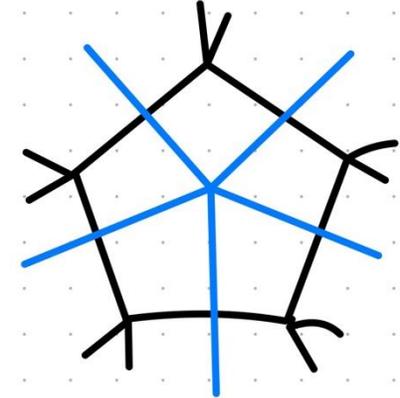
$$+ \frac{s_{12}^2 x_{1234} x_4}{s_{123} s_{1234} x_{12} x_{34}} + \frac{s_{12} s_{34} x_{123} x_{1234}}{s_{123} s_{1234} x_{12} x_{34}} + (1 \leftrightarrow 4, 2 \leftrightarrow 3)$$

M

$$\frac{1}{-c[1] + c[3] + c[4]}$$

$$\left(\frac{c[3] c[4]^2}{2(-2 + c[4])} - \frac{c[1] \times c[2] \times c[4]}{(-1 + c[4])(-1 + c[2] - c[5])} - \frac{c[4](-c[1] + c[4] + c[3] \times c[4]) c[5]}{2(-2 + c[4])(-2 + c[2] - c[5])} + \frac{c[2] \times c[4] \times c[5]}{2(-1 + c[2] - c[5])} - \frac{c[2] \times c[4](-1 + c[2] + c[3] - c[5]) c[5]}{2(-2 + c[2] - c[5])(1 + c[2] - c[5])} + \frac{c[3] \times c[4] c[5]^2}{2(-2 + c[2] - c[5])(1 + c[2] - c[5])} \right)$$

$$+(c[3] \leftrightarrow c[4], c[2] \leftrightarrow c[5])$$



$$c_5 := \delta_{12}, \quad c_2 := \delta_{34},$$

$$c_3 := \delta_{123}, \quad c_4 := \delta_{234},$$

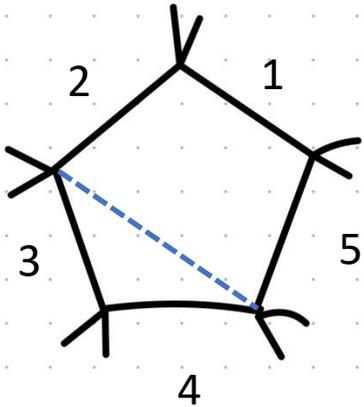
$$c_1 := \delta_{23}$$

$$c_3 + c_4 - c_1 = \delta_{14} - 1, \quad c_5 - c_2 = \delta_{35} + \delta_{45} - 1,$$

Amplitude side: matching kernels

Five-mass chiral pentagon in 4d $I_n^1 = \frac{1}{\pi^2} \int d^4 Q (-Q \cdot Y) \prod_{i=1}^n \frac{1}{(-P_i \cdot Q)} \quad Y \cdot P_{1,2,3,4} = 0$

$$I_5 = \frac{(-Y \cdot P_5)}{P_{25} P_{14} P_{35}} \oint \prod_{i=1}^5 \left(\frac{dc_i}{(2\pi i)} u_i^{c_i-1} \Gamma(1 - c_i) \right) \Gamma(c_1 + c_2 - c_4 - 1) \\ \times \Gamma(c_3 + c_4 - c_1) \Gamma(c_2 + c_3 - c_5) \Gamma(c_1 + c_5 - c_3 - 1) \Gamma(c_4 + c_5 - c_2)$$



$$u_1 = \frac{P_{14} P_{23}}{P_{13} P_{24}}, \quad u_2 = \frac{P_{25} P_{34}}{P_{24} P_{35}}, \quad u_3 = \frac{P_{13} P_{45}}{P_{14} P_{35}}, \quad u_4 = \frac{P_{15} P_{24}}{P_{14} P_{25}}, \quad u_5 = \frac{P_{12} P_{35}}{P_{13} P_{25}}$$

$$u_i = u_{i,i+2} = \frac{P_{i,i+3} P_{i+1,i+2}}{P_{i,i+2} P_{i+1,i+3}}, \quad i = 1, \dots, 5.$$

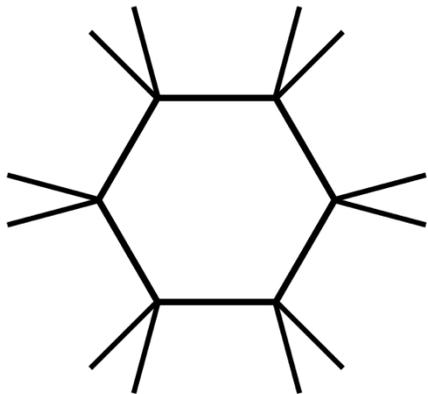
Reduction for a Hexagon with restricted kinematics: $P_{16} = P_{56} = 0, \frac{P_{62}}{P_{63}} = \frac{P_{52}}{P_{53}}, \frac{P_{64}}{P_{63}} = \frac{P_{54}}{P_{53}}$

$$\hat{I}_6 \equiv (P_{14}P_{25}P_{36}) I_6 = \oint \prod_{i=1}^9 \frac{dc_i}{2\pi i} u_i^{-c_i} \left(\prod_{i=4}^9 \Gamma(c_i) \right) \Gamma(c_1 - c_5 - c_6) \Gamma(c_2 - c_6 - c_7) \\ \times \Gamma(1 - c_1 - c_3 + c_5 + c_8) \Gamma(c_3 - c_7 - c_8) \Gamma(c_1 - c_8 - c_9) \Gamma(1 - c_1 - c_2 + c_6 + c_9) \\ \times \Gamma(1 - c_2 - c_3 + c_4 + c_7) \Gamma(c_3 - c_4 - c_5) \Gamma(c_2 - c_4 - c_9),$$

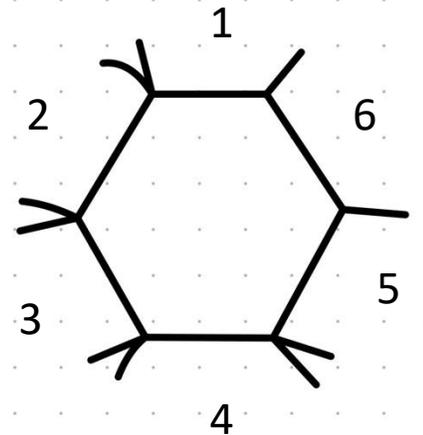
$$u_{ij} \equiv \frac{P_{i,j+1} P_{i+1,j}}{P_{ij} P_{i+1,j+1}}.$$

$$u_i = u_{i,i+3}, \quad i = 1, 2, 3,$$

$$u_{i+3} = u_{i+1,i+5}, \quad i = 1, \dots, 6$$

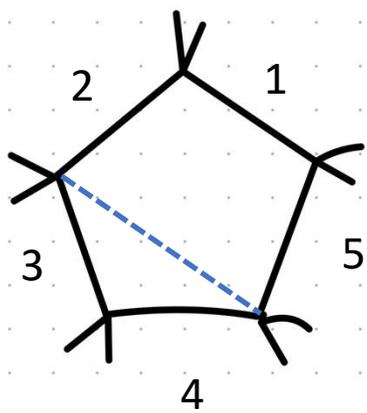


Set $u_7 = u_8 = 0, u_2 = u_9 = 1$
Apply Bannes Lemma to
integrate out a_2, a_9

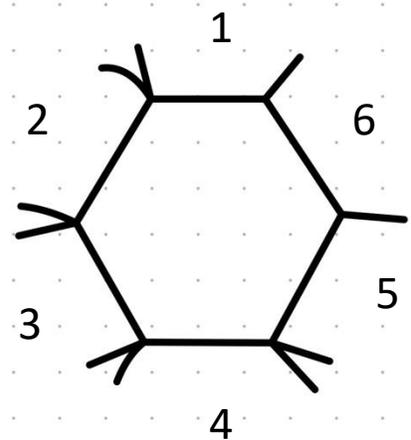


M
normalize by I_5^1

$$\frac{1}{1 - c_4}$$

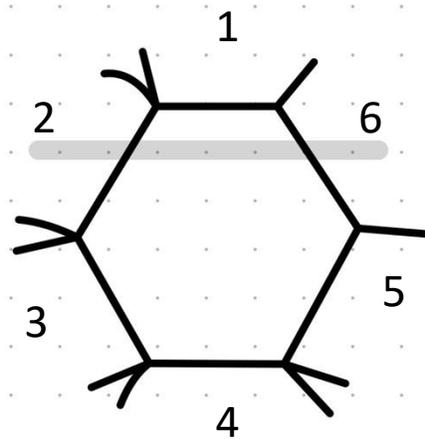


$$M \longrightarrow 1$$



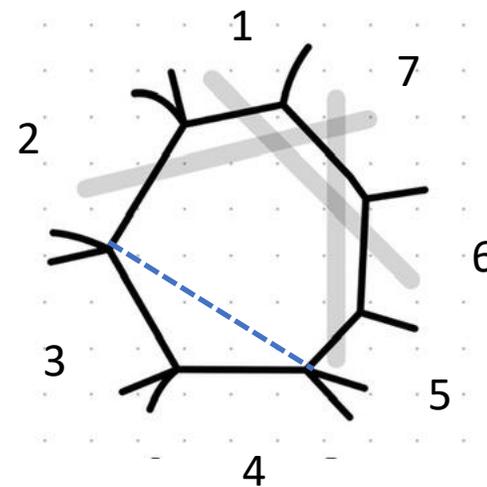
Hex_A

$$M \longrightarrow \frac{1}{1 - \delta_{234}}$$



Hex_B

$$M \longrightarrow \frac{1}{\delta_{12} - \delta_{34}}$$



$$M \longrightarrow \frac{\delta_{123}}{(1 - \delta_{234})(\delta_{12} - \delta_{34} + 1)}$$

Hep_A $D_{1,1,1,2,1,1,1}$

Deformed Hexagon

$$\int \frac{dt}{(1+t)} \text{Hex}_{A,B}((1+t)u_3) \quad \text{Each contains } >3000 \text{ terms in its symbol}$$

First 2 entries:

$$\{-1 + w, w, -1 + wb, wb, -1 + z, z, -w + z, -1 + zb, zb, -wb + zb, \boxed{-1 + zzb}\}$$

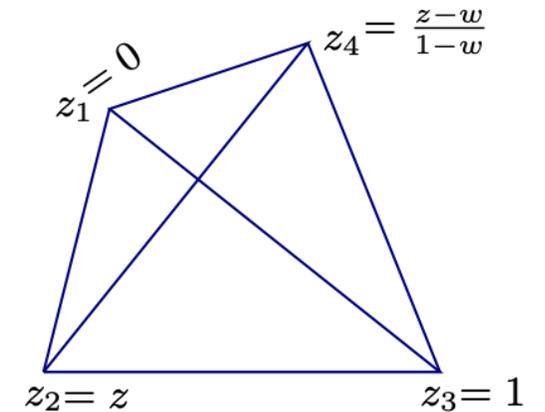
3rd entry \ first 2 :

$$\{z - zb, -wbz + wzb, \boxed{-wbz + wzb, -wb + wzb, -w + wbz, 1 - w - wb + wzb, 1 - w - wb + wz, -1 + w + wb - wbz - wzb + zzb}\}$$

Last entry \ first 3 :

$$\{1 - wwb, 1 - wbz, wzb - zb, 1 - wzb, z - wzb, 1 - z - zb + wzb, 1 - z + wzb - zb, 1 - wwb - z + wzb - zb + wzb\}$$

$$z \leftrightarrow w, \bar{z} \leftrightarrow \bar{w}$$



Summary

- Systematic studies on deformed star integrals
 - extra structures in the symbol alphabets,
 - universal patterns independent of supersymmetry?
- Interpretation on the Mellin amplitudes
 - Physical v.s. spurious poles in Mellin space; connection to symbol letters
 - Last entry: physical origin of singularities in the LS
- How to build differential equations systematically (without/with minimum IBP)?
- Higher loop @ Lower point : generalize with dimensional regularization