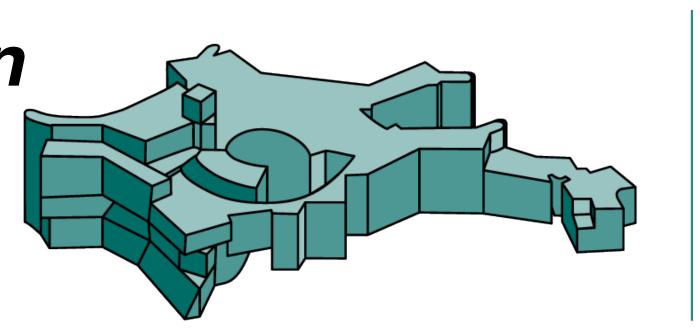
Reference: EK, Nature Rev. Phys. 4, 452 (2022)

A Tantalizing Hint of Cosmological Parity Violation in the Polarized Light of the CMB

Does the Universe distinguish between left and right?





Eiichiro Komatsu (Max Planck Institute for Astrophysics)

"Gravitational Waves, Cosmology and Fundamental Physics", MAX-PLANCK-INSTITUT Hangzhou Institute for Advanced Study, Nov. 1, 2025

Overarching Theme

There are many ideas, but how can we make progress?

- The current cosmological model (ACDM) requires new physics beyond the standard model of elementary particles and fields.
 - What is dark matter (CDM)? => CDM, WDM, FDM, ...
 - What is dark energy (/I)? => Dynamical field, modified gravity, quantum gravity, ...

New in cosmology!

Violation of parity symmetry may hold the answer to these fundamental questions.

Reference: nature reviews physics

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Review Article | Published: 18 May 2022

New physics from the polarized light of the cosmic microwave background

Eiichiro Komatsu

Key Words:

Cosmic Microwave

Background (CMB)

- Polarization
- **Parity Symmetry**

Nature Reviews Physics 4, 452-469 (2022) Cite this article

1. Parity

Probing Parity Symmetry

Definition

- Parity transformation = Inversion of all spatial coordinates
 - $(x, y, z) \rightarrow (-x, -y, -z)$
- Parity symmetry in physics states:
 - The laws of physics are invariant under inversion of all spatial coordinates.
- Violation of parity symmetry = The laws of physics are **not** invariant under...

 We ask, "When we observe a certain phenomenon in nature, do we also observe its mirror image with equal probability?"

Parity and Rotation

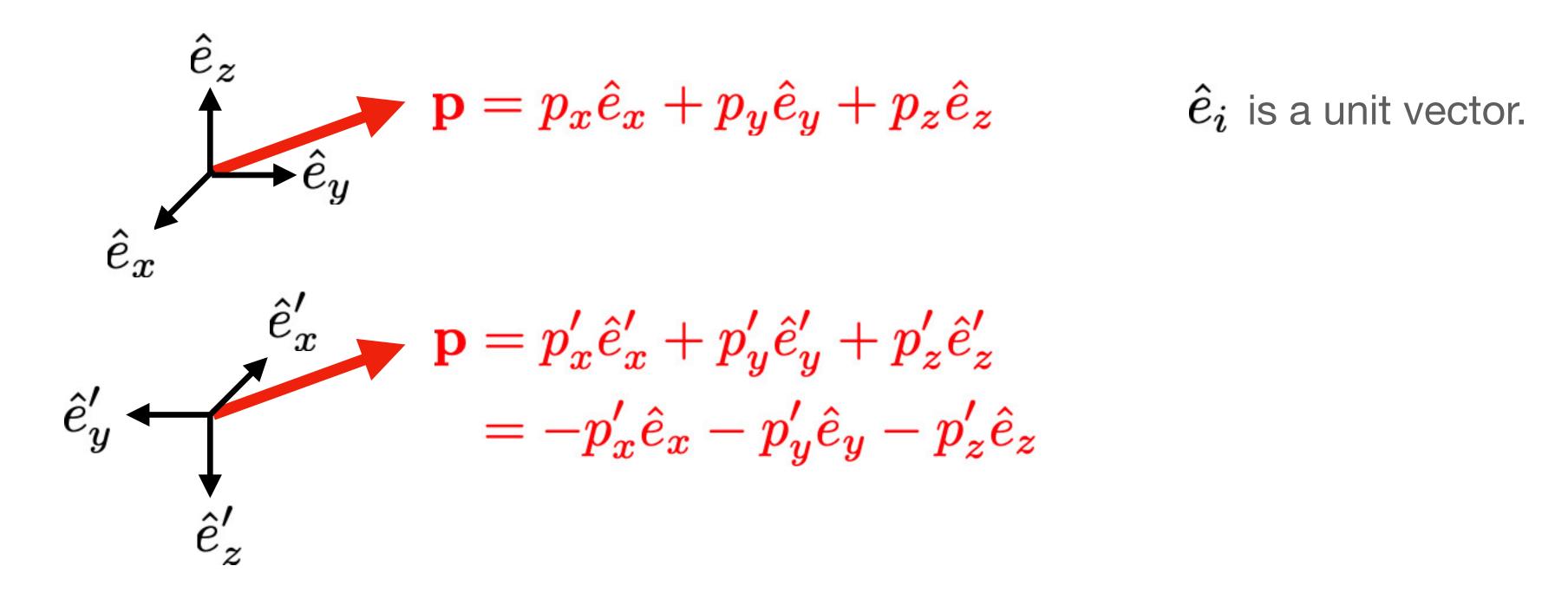
- Parity transformation (x -> -x) and 3d rotation (x -> Rx) are different.
 - R is a continuous transformation and the determinant of R is det(R) = +1.
 - Parity is a discrete transformation and the determinant is -1, as

$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} \rightarrow \begin{pmatrix} -x \\ -y \\ -z \end{pmatrix} = \begin{pmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix}$$

2. Pseudovector, Pseudoscalar

Parity Transformation: Vector

E.g., momentum, electric field



- p is the same vector, written using two different basis vectors.
- Therefore, **p**'s components are transformed as $(p_x', p_y', p_z') = (-p_x, -p_y, -p_z)$

Parity Transformation: Pseudovector

E.g., angular momentum, magnetic field

- Orbital angular momentum, $\mathbf{L} = \mathbf{r} \times \mathbf{p}$, is a pseudovector. Its components do not change under parity transformation: $(L_x', L_y', L_z') = (L_x, L_y, L_z)$
 - Both $\mathbf{r} = (X, Y, Z)$ and $\mathbf{p} = (p_x, p_y, p_z)$ are vectors whose components change sign. Thus, their products do not change, e.g.,

$$L_x' = Y'p_z' - Z'p_y'$$
 $= (-Y)(-p_z) - (-Z)(-p_y)$
 $= L_x$

Parity Transformation: Pseudoscalar How to test parity symmetry?

- A dot product of a vector and a pseudovector is a pseudoscalar.
 - Like a scalar, a pseudoscalar is invariant under rotation.
 - But, a pseudoscalar changes sign under parity transformation.
- Experimental test of parity symmetry: Construct a pseudoscalar and see if the average value is zero. If not, the system violates parity symmetry!
 - Example: a dot product of particle A's momentum and particle B's angular momentum: $\mathbf{p}_A \cdot \mathbf{L}_B$. Measure this and average over many trials. Does the average vanish, $\langle \mathbf{p}_A \cdot \mathbf{L}_B \rangle = 0$?

3. Discovery of Parity Violation in β-decay (weak interaction)

Experimental Test of Parity Conservation in Beta Decay*

C. S. Wu, Columbia University, New York, New York

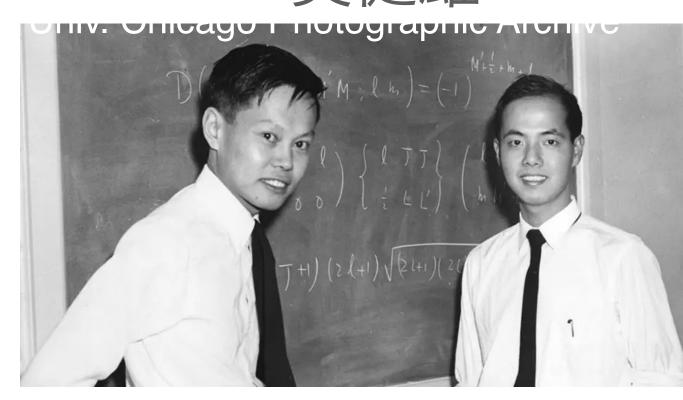
AND

E. Ambler, R. W. Hayward, D. D. Hoppes, and R. P. Hudson, National Bureau of Standards, Washington, D. C. (Received January 15, 1957)

TN a recent paper on the question of parity in weak interactions, Lee and Yang critically surveyed the experimental information concerning this question and reached the conclusion that there is no existing evidence either to support or to refute parity conservation in weak interactions. They proposed a number of experiments on beta decays and hyperon and meson decays which would provide the necessary evidence for parity conservation or nonconservation. In beta decay, one could measure the angular distribution of the electrons coming from beta decays of polarized nuclei. If an asymmetry in the



Chien-Shiung Wu 吳健雄

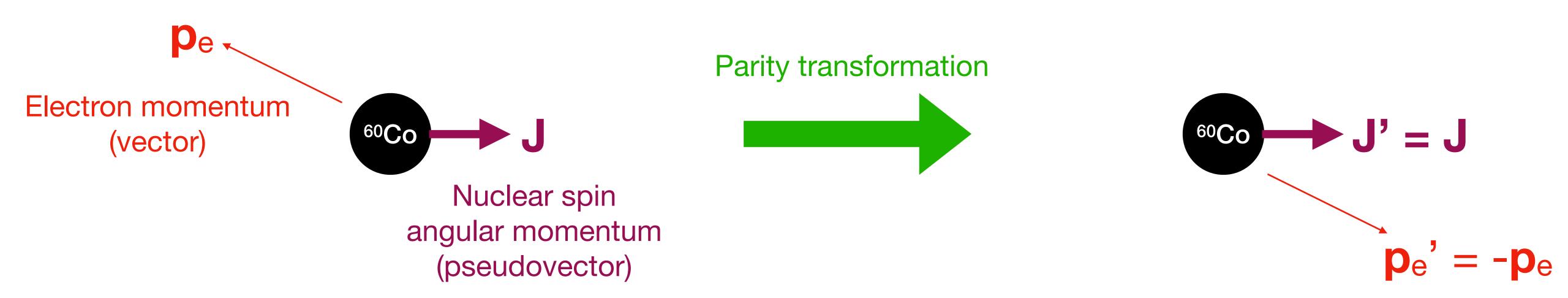


Chen-Ning Yang 楊振寧

Tsung-Dao Lee 李政道

The Wu Experiment of \beta-decay

 60 Co -> 60 Ni + e⁻ + $\overline{\nu}_e$ + 2 Y



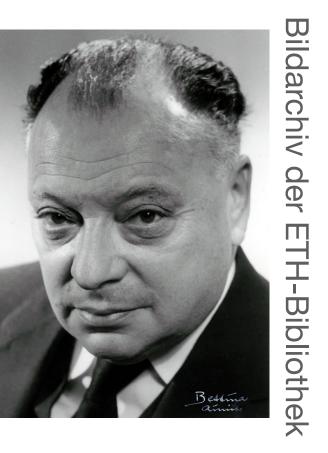
- Electrons must be emitted with equal probability in all directions relative to J, if parity symmetry is respected in β-decay.
 - This was not observed: $\langle \mathbf{p}_e \cdot \mathbf{J} \rangle \neq 0$. Parity symmetry is violated in β -decay!

"This Month in Physics History", APS News, October 2022

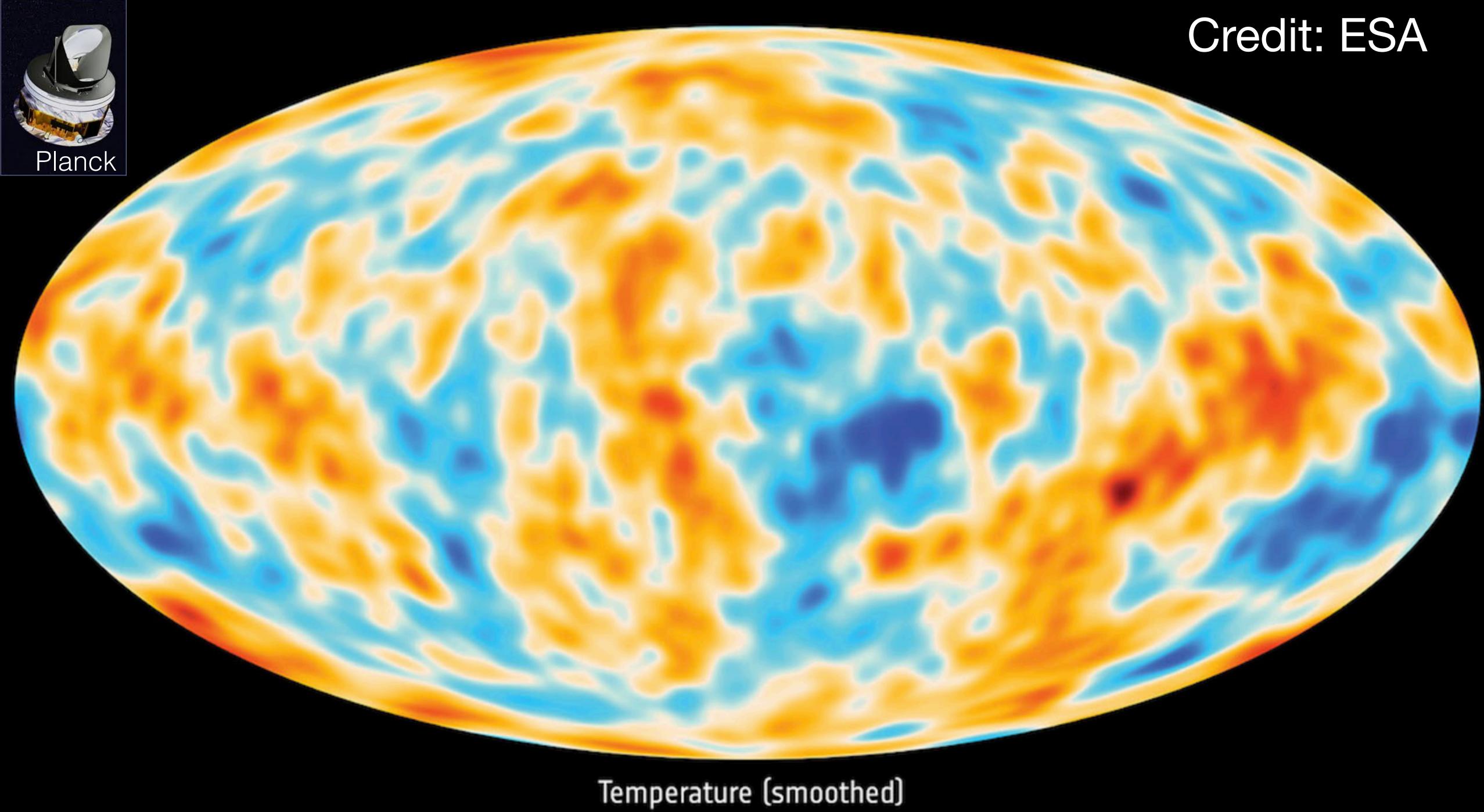
Initial reaction

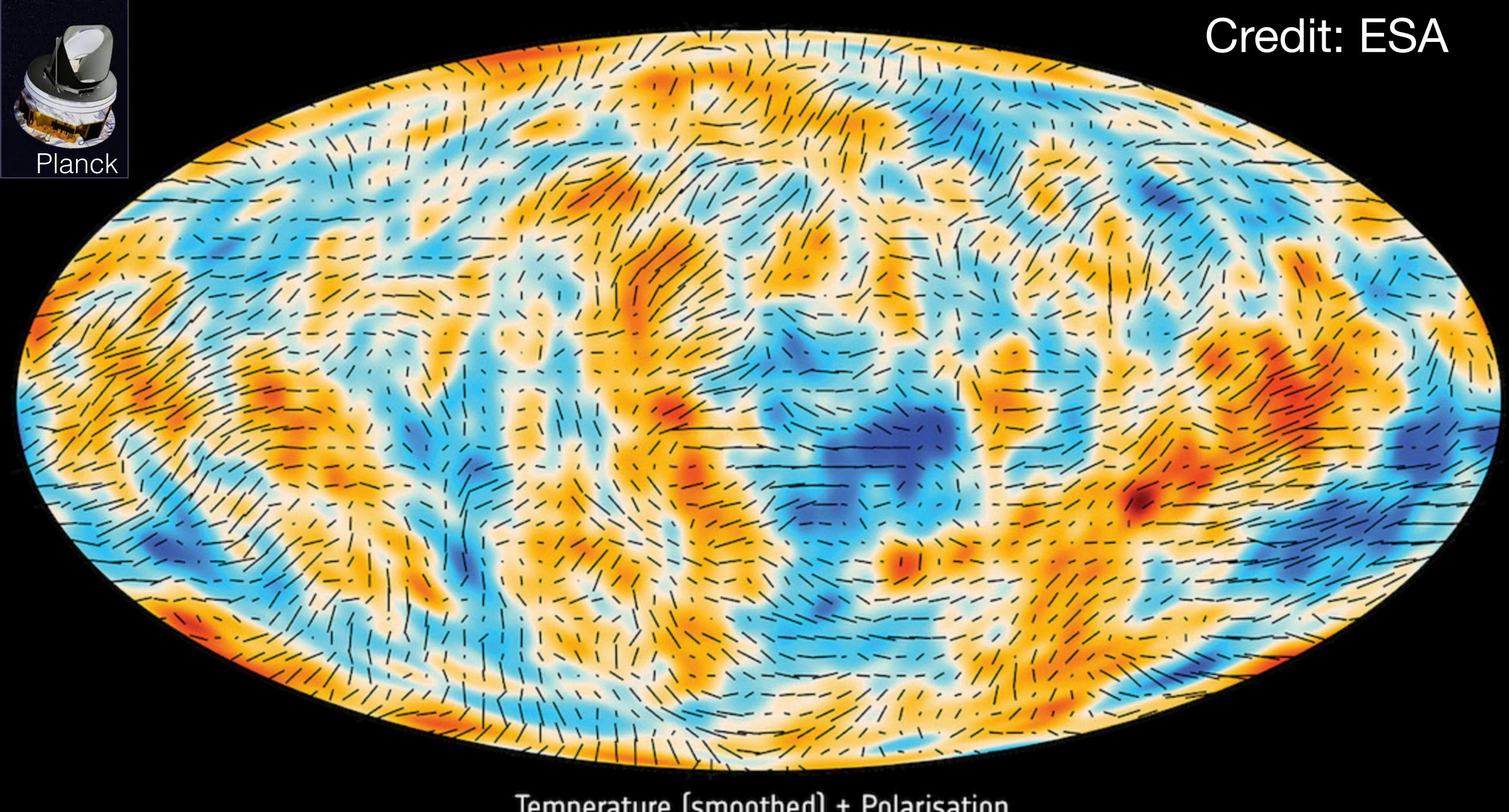
Many physicists did not believe it initially.

- To Lee and Yang's theoretical paper on parity violation in β-decay:
 - Wolfgang Pauli said, "Ich glaube aber nicht, daß der Herrgott ein schwacher Linkshänder ist" (I do not believe that the Lord is a weak left-hander).
- To Wu's discovery paper:
 - Wolfgang Pauli said, "Sehr aufregend. Wie sicher ist die Nachricht?" (Very exciting. How sure is this news?)
- This was shocking news. The weak interaction distinguishes between left and right!
- In this talk we ask, "Does the Universe distinguish between left and right?"

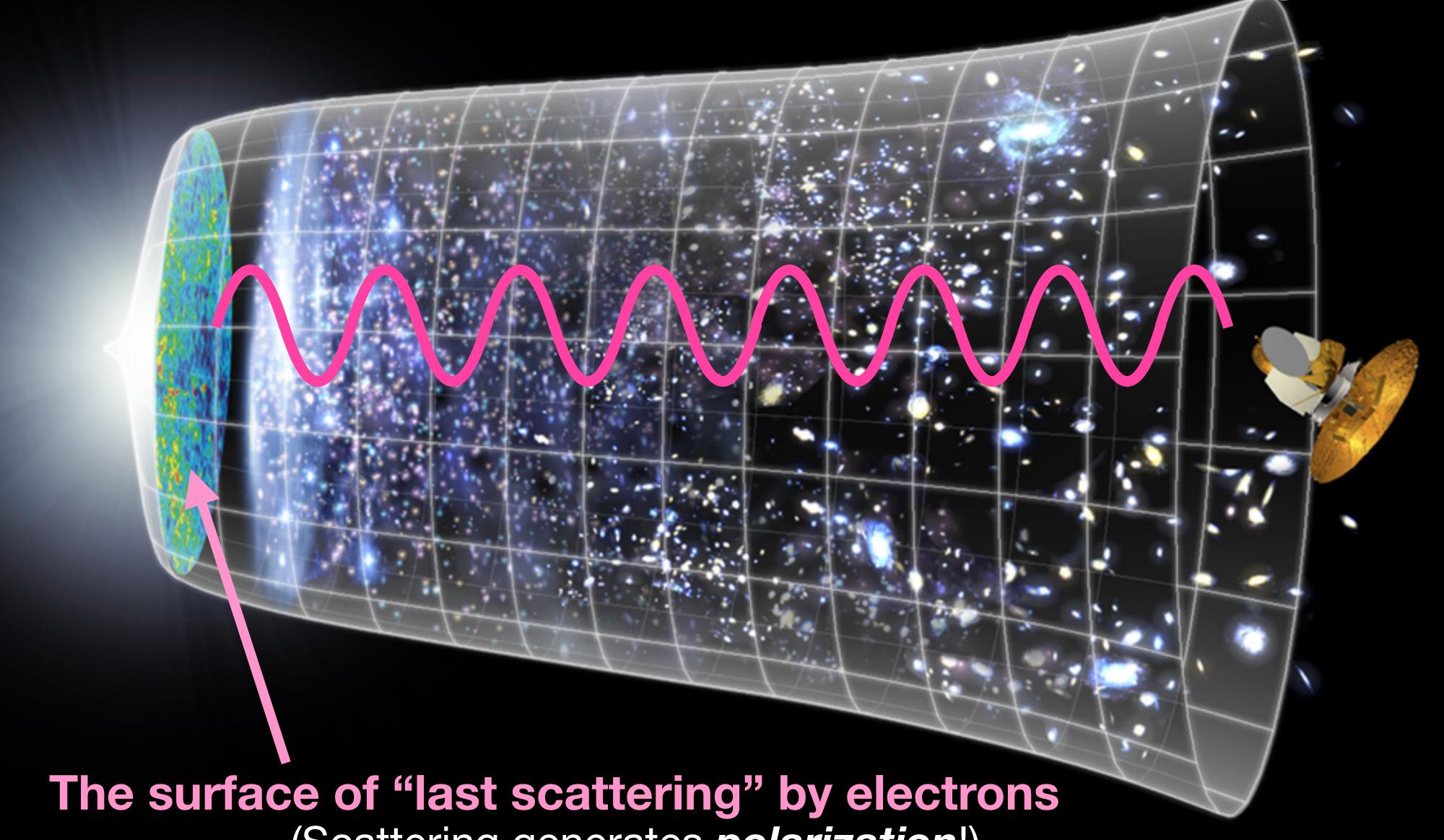


4. Parity Violation in the Cosmic Microwave Background (CMB)





CMB photons have traveled for a very long time!

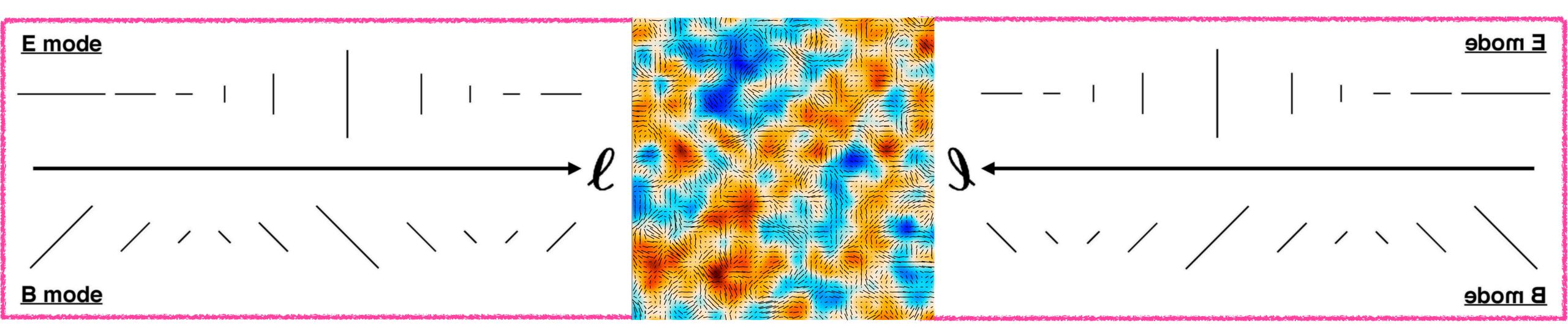


(Scattering generates polarization!)

Credit: WMAP Science Team

Pseudoscalar: EB correlation

• The observed pattern of the CMB polarization can be decomposed into eigenstates of parity, called "E modes" and "B modes".

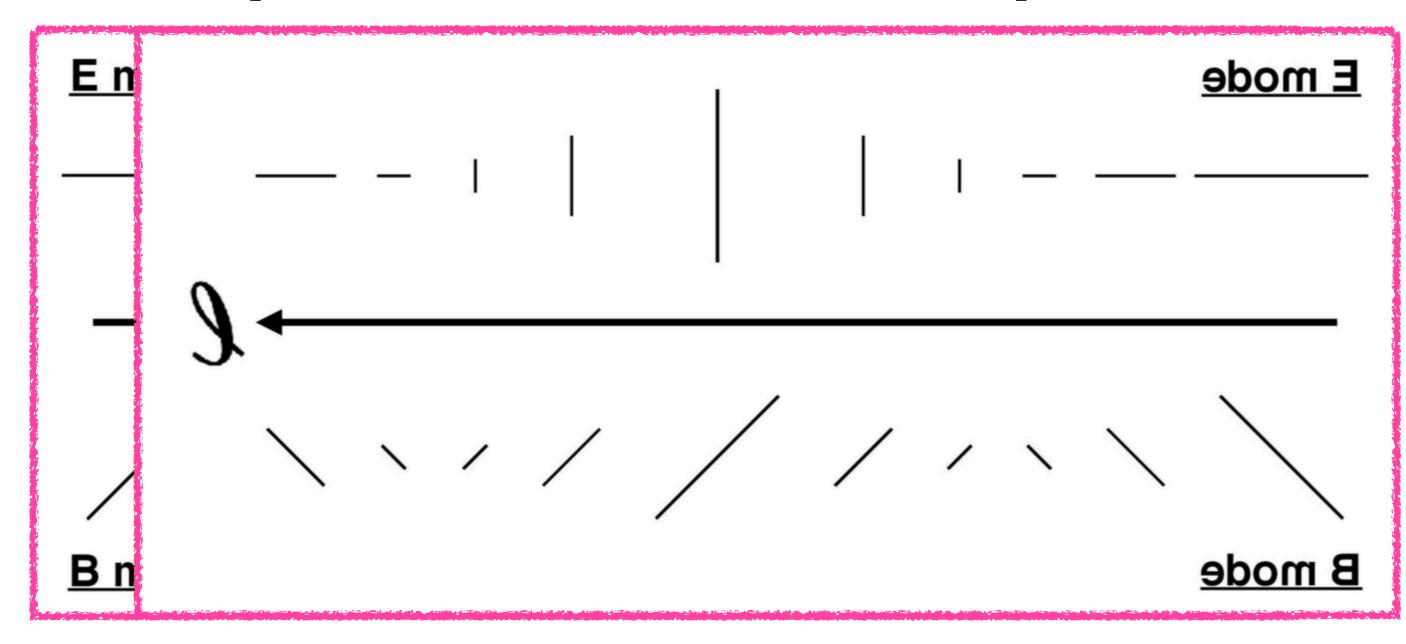


- E and B modes are transformed differently under the parity transformation. Therefore, the product of the two, the "EB correlation", is a pseudoscalar.
- The full-sky average of the EB correlation must vanish (to within the measurement uncertainty), if there is no parity violation!

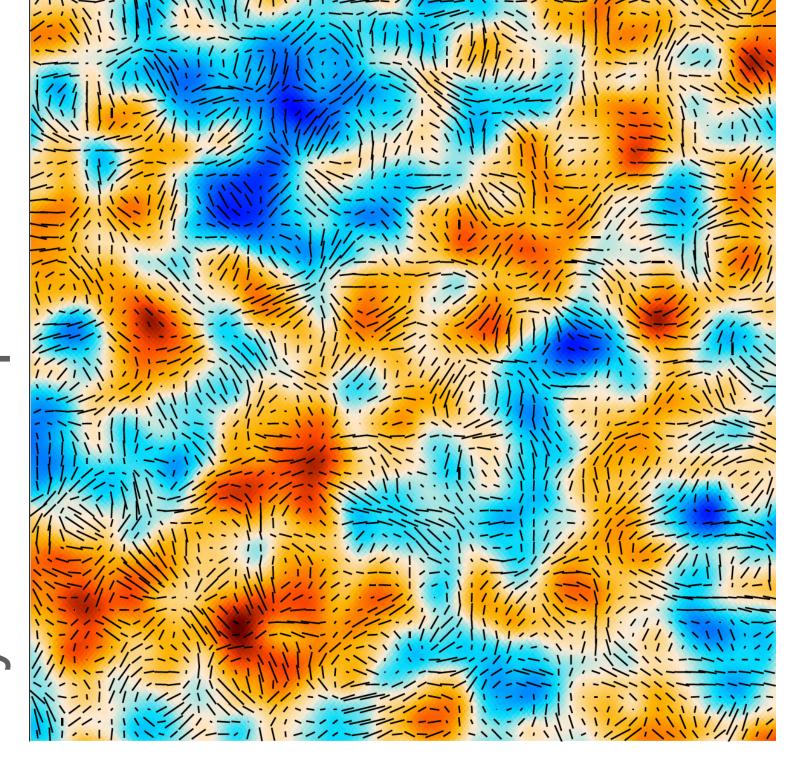
Zaldarriaga, Seljak (1997); Kamionkowski, Kosowsky, Stebbins (1997)

Parity eigenstates: E and B modes

Concept defined in Fourier space



This map is dominated by E-mode polarization



$$\langle E_{\ell} E_{\ell'}^* \rangle = (2\pi)^2 \delta_D^{(2)} (\ell - \ell') C_{\ell}^{EE}$$

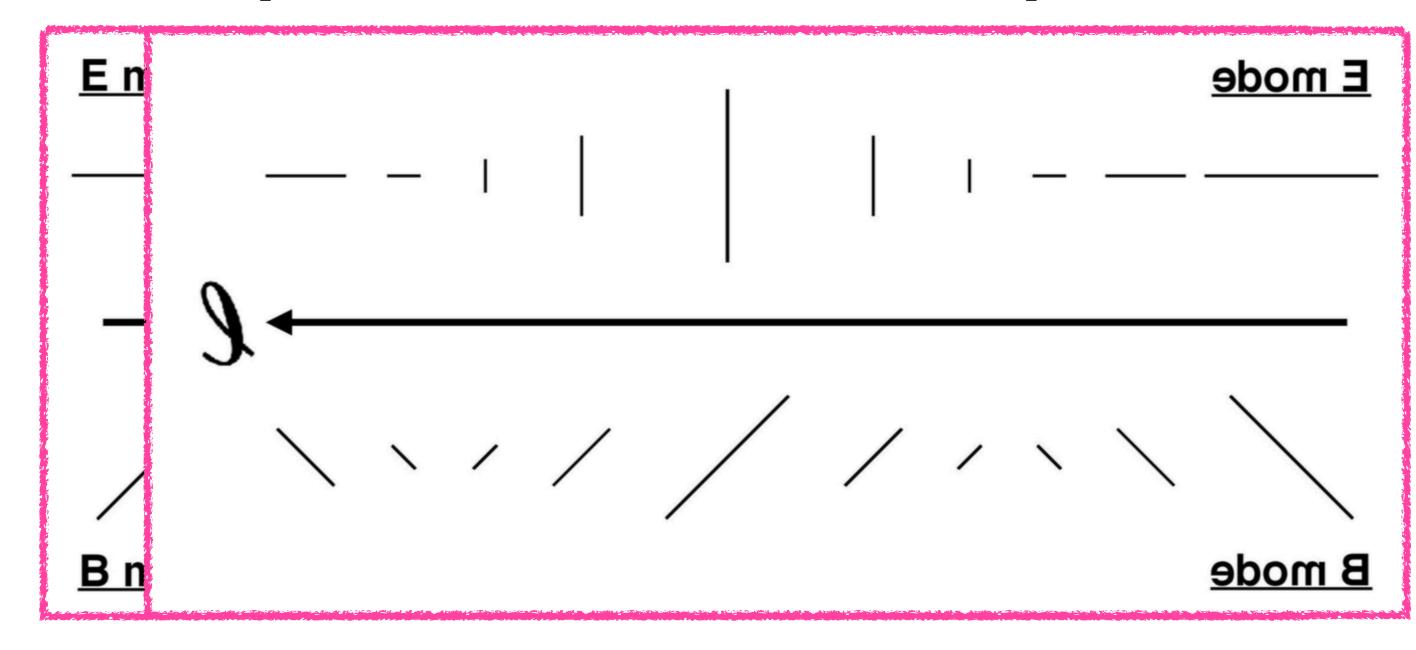
$$\langle B_{\ell} B_{\ell'}^* \rangle = (2\pi)^2 \delta_D^{(2)} (\ell - \ell') C_{\ell}^{BB}$$

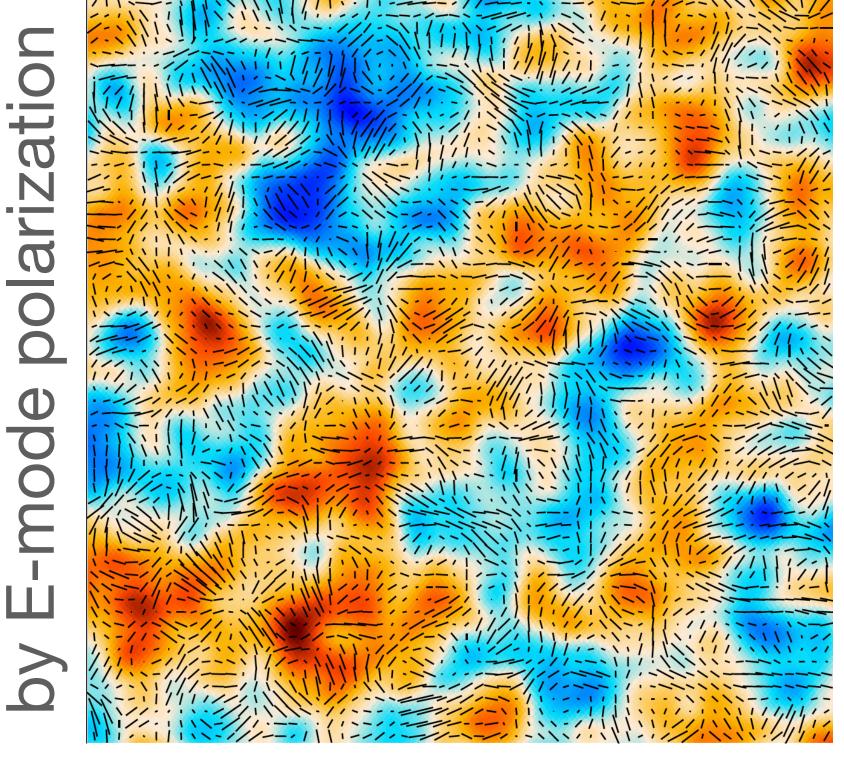
$$\langle T_{\ell} E_{\ell'}^* \rangle = \langle T_{\ell}^* E_{\ell'} \rangle = (2\pi)^2 \delta_D^{(2)} (\ell - \ell') C_{\ell}^{TE}$$

These are scalars and insensitive to parity violation.

Parity eigenstates: E and B modes

Concept defined in Fourier space





$$\langle E_{\ell}E_{\ell'}^* \rangle = (2\pi)^2 \delta_D^{(2)}(\ell - \ell')C^{EE}$$
 The other combinations, and ,

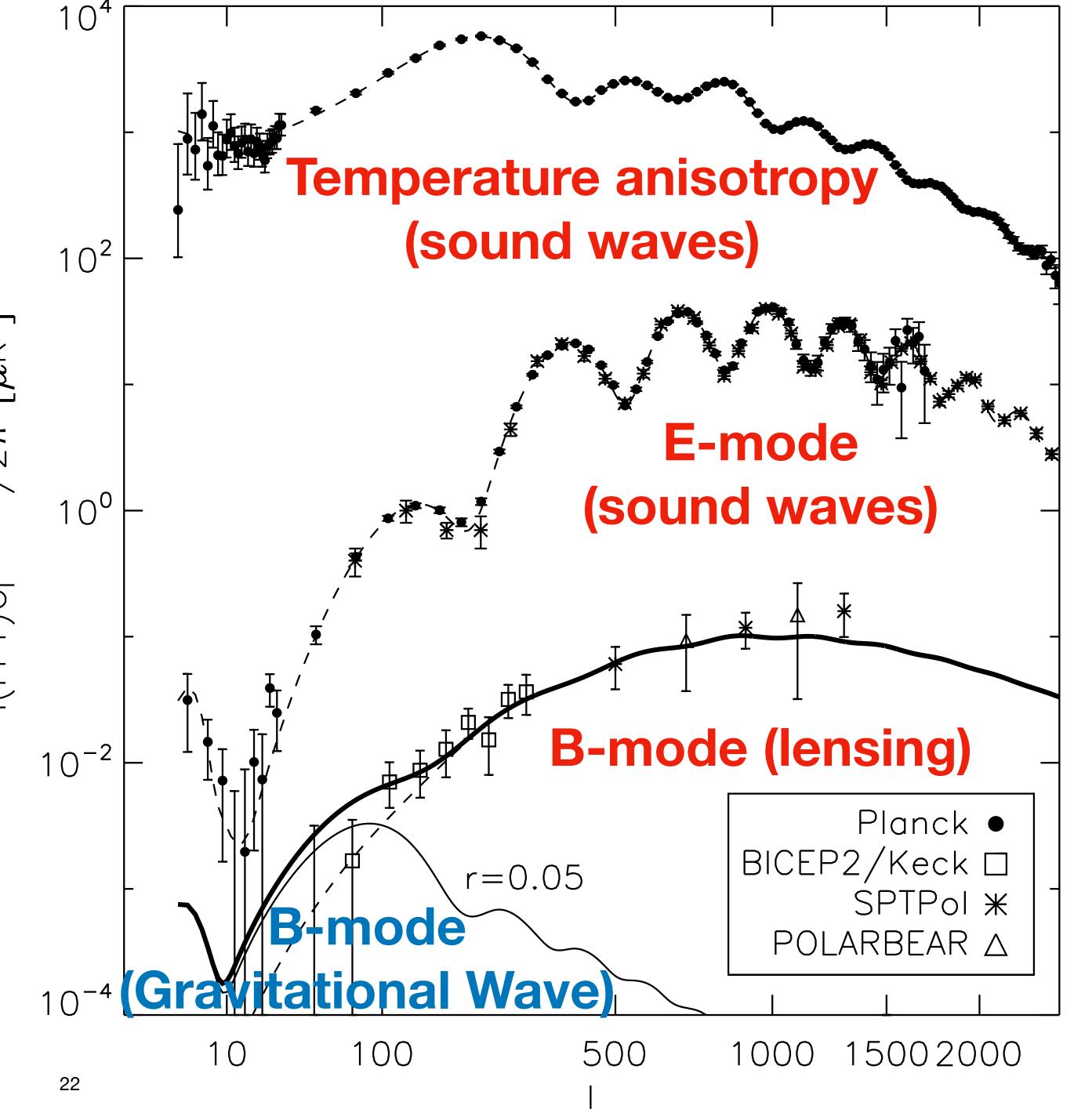
 $\langle B_{\ell}B_{\ell'}^*\rangle = (2\pi)^2 \delta_D^{(2)}(\ell$ are pseudoscalars and sensitive to parity violation!

$$\langle T_{\ell} E_{\ell'}^* \rangle = \langle T_{\ell}^* E_{\ell'} \rangle = (2\pi)^2 \delta_D^{(2)} (\ell - \ell') C_{\ell}^{TL}$$

CMB Power Spectra

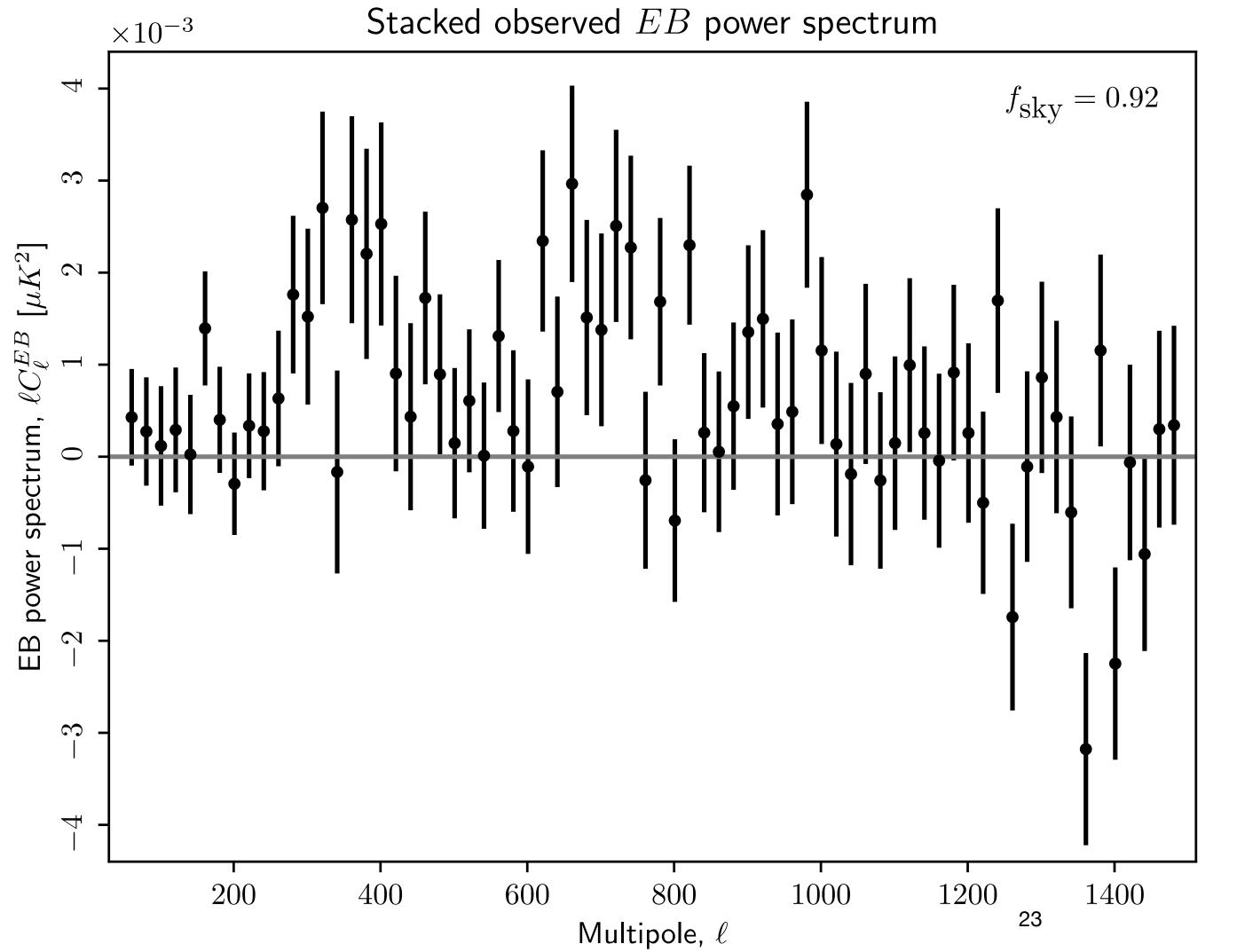
Progress over 30 years

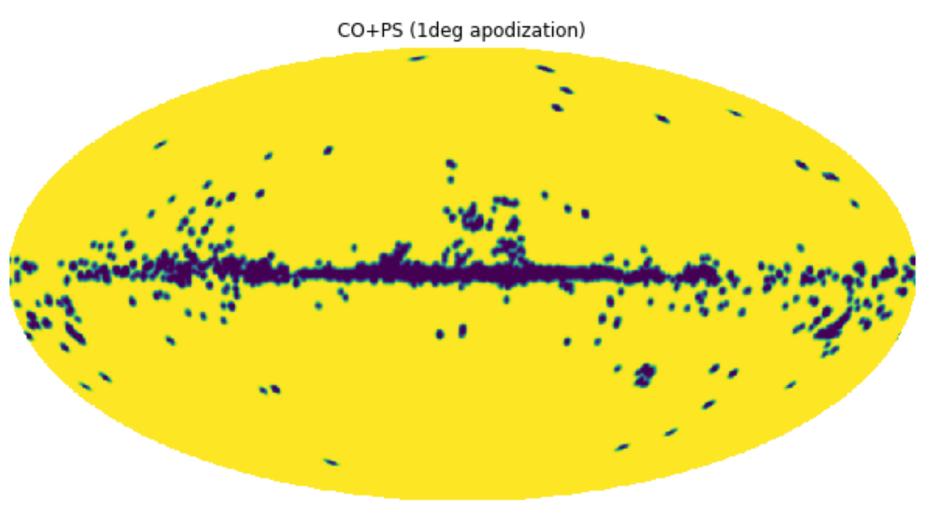
- This is the typical figure seen in talks and lectures on the CMB.
 - The temperature and the E- and B-mode polarization power spectra are well measured.
- Parity violation appears in the TB and EB power spectra, not shown here.



This is the EB power spectrum (WMAP+Planck)

Nearly full-sky data (92% of the sky)

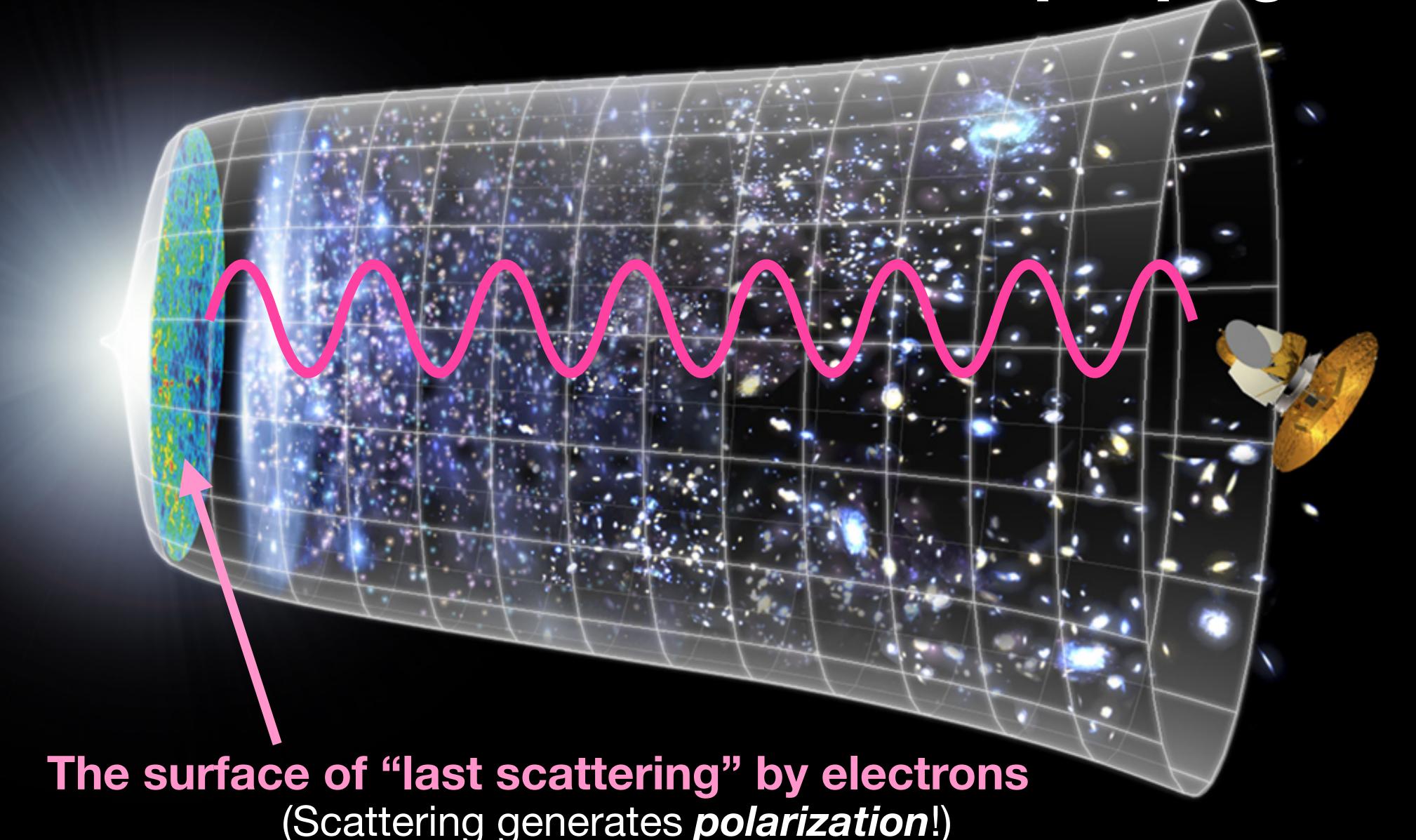




- $\chi^2 = 125.5$ for DOF=72
 - Unambiguous signal of something!

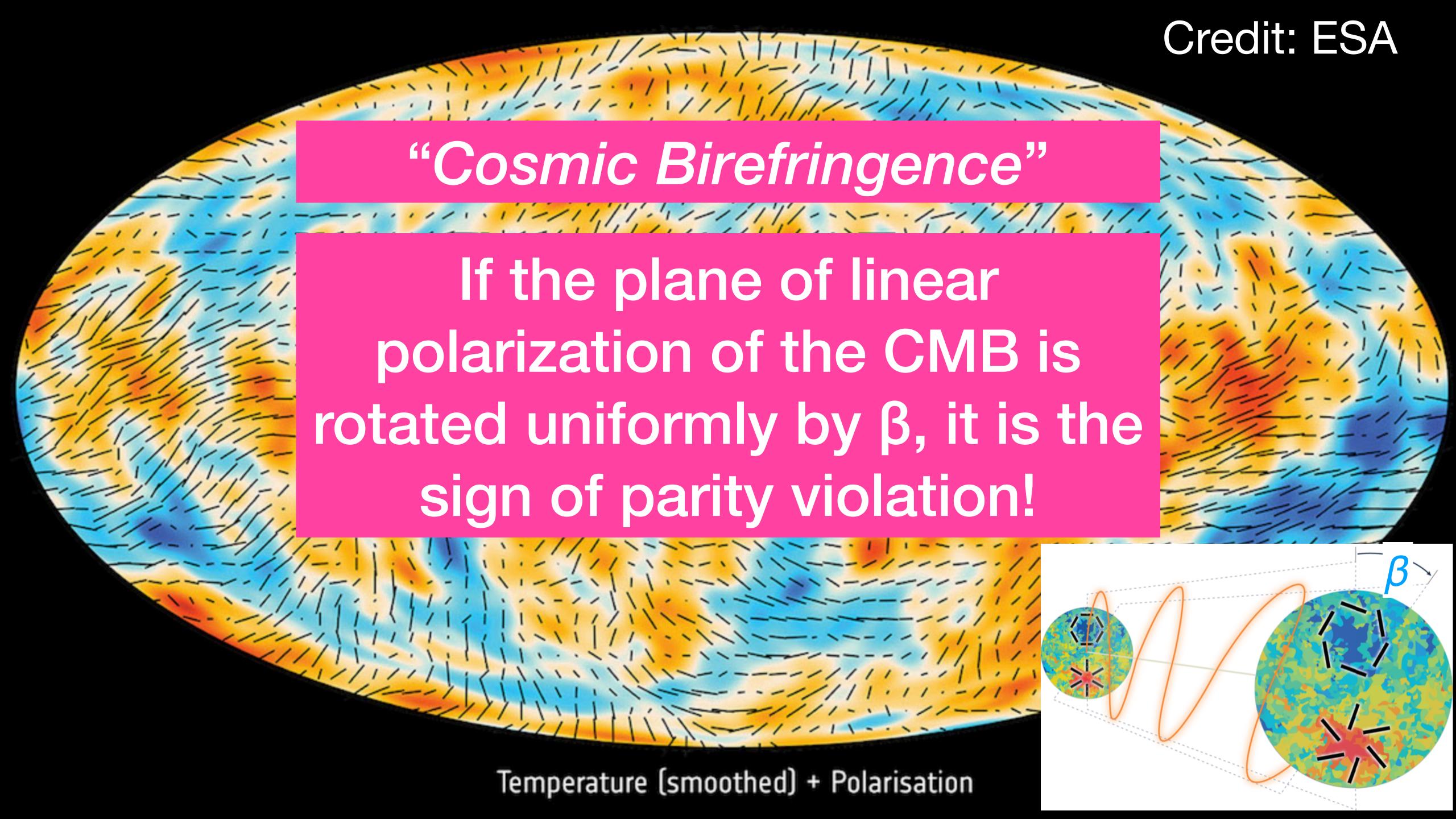
5. Cosmic Birefringence: Rotation of the Plane of Linear Polarization

How does the EM wave of the CMB propagate?



Credit: WMAP Science Team

How does the EM wave of the CMB propagate?



Lue, Wang, Kamionkowski (1999); Feng et al. (2005, 2006)

E-B mixing by rotation of the plane of linear polarization

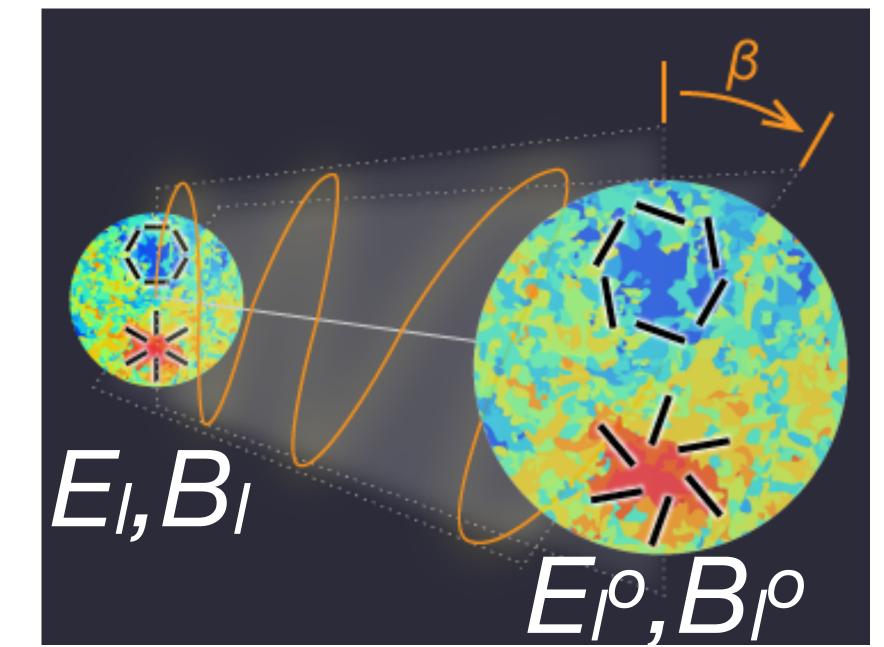
• Observed E- and B-mode polarization, E_I° and B_I°, are related to those before rotation as

$$E_\ell^{\rm o} \pm iB_\ell^{\rm o} = (E_\ell \pm iB_\ell)e^{\pm 2i\beta}$$

which gives

$$E_{\ell}^{o} = E_{\ell} \cos(2\beta) - B_{\ell} \sin(2\beta)$$

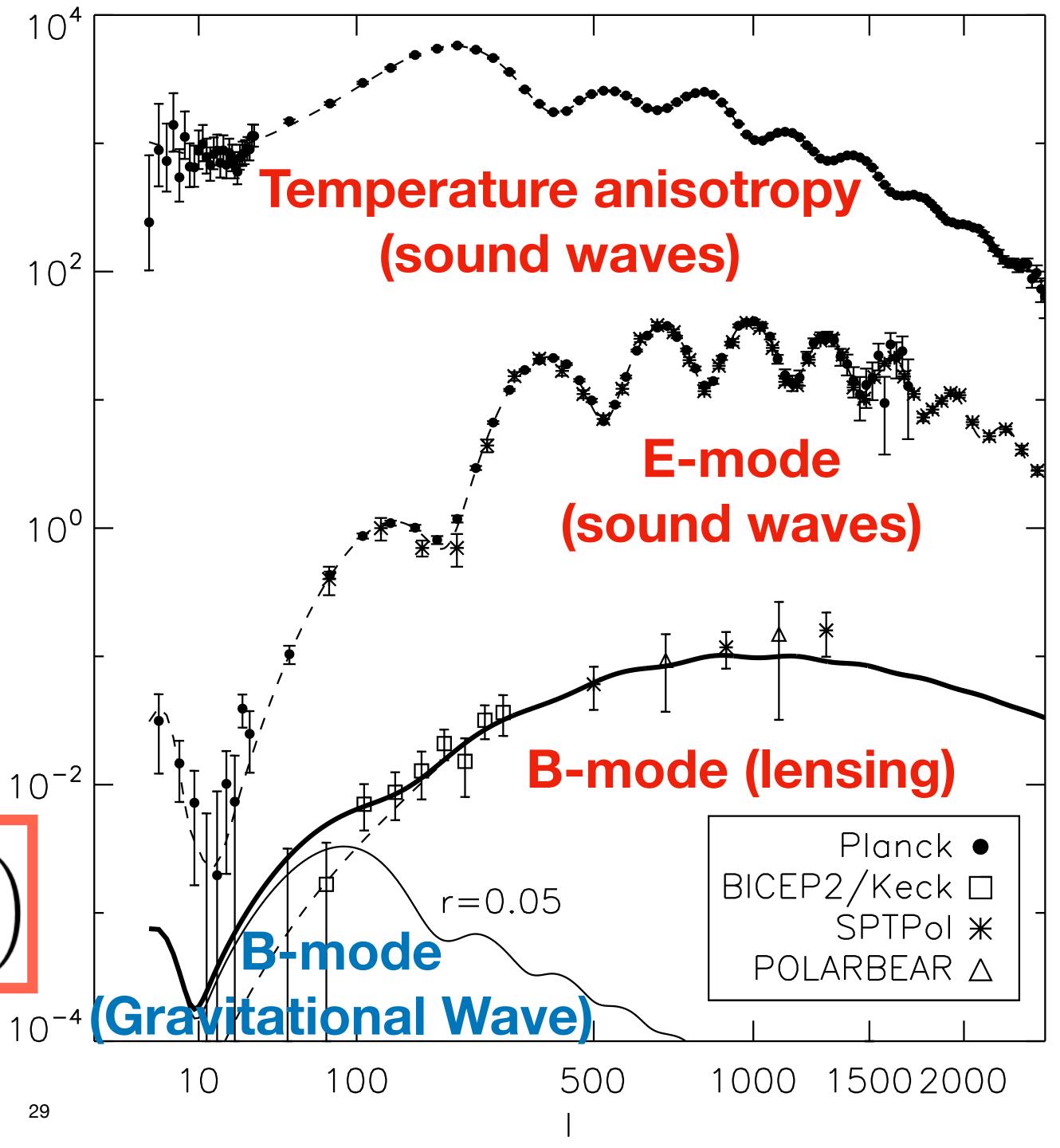
$$B_{\ell}^{o} = E_{\ell} \sin(2\beta) + B_{\ell} \cos(2\beta)$$



CMB Power Spectra

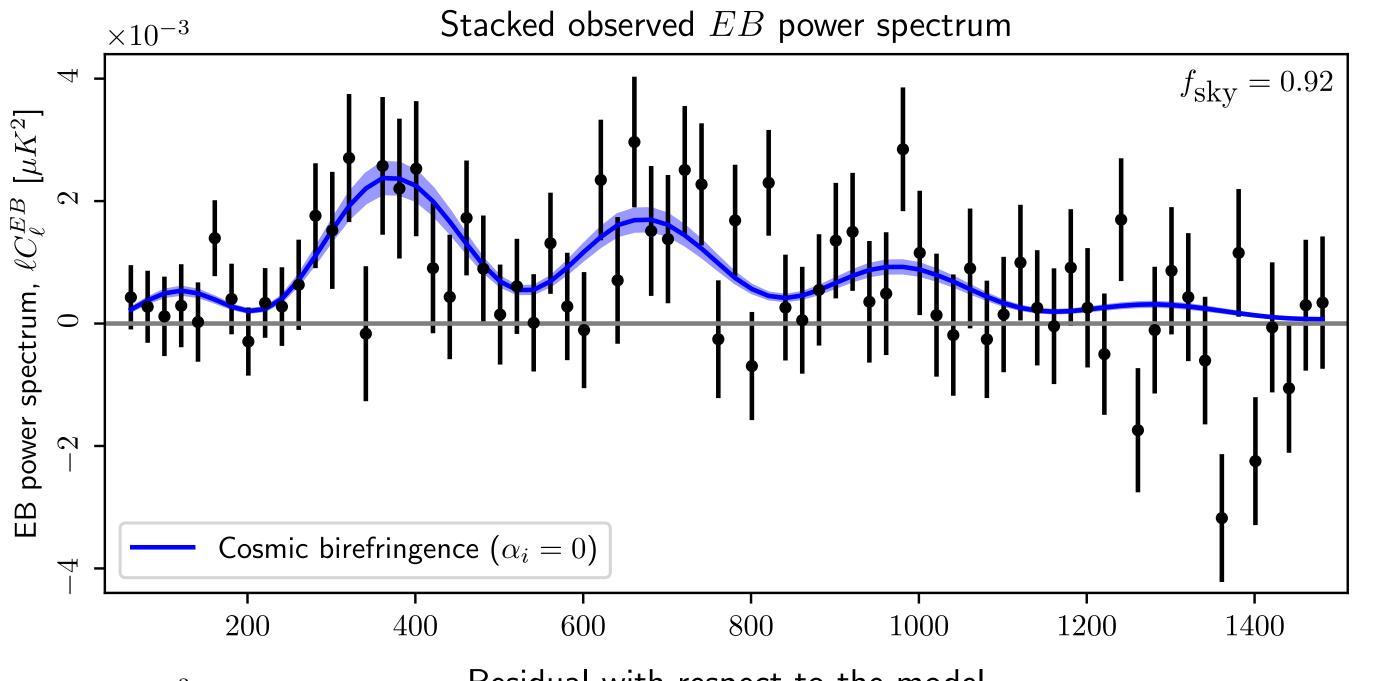
- Rotation of the plane of linear polarization **mixes** E and B modes.
- Therefore, the EB correlation will be given by the difference between the EE and BB correlations.
- Observed EE is much greater than BB. We expect EB to look like EE!

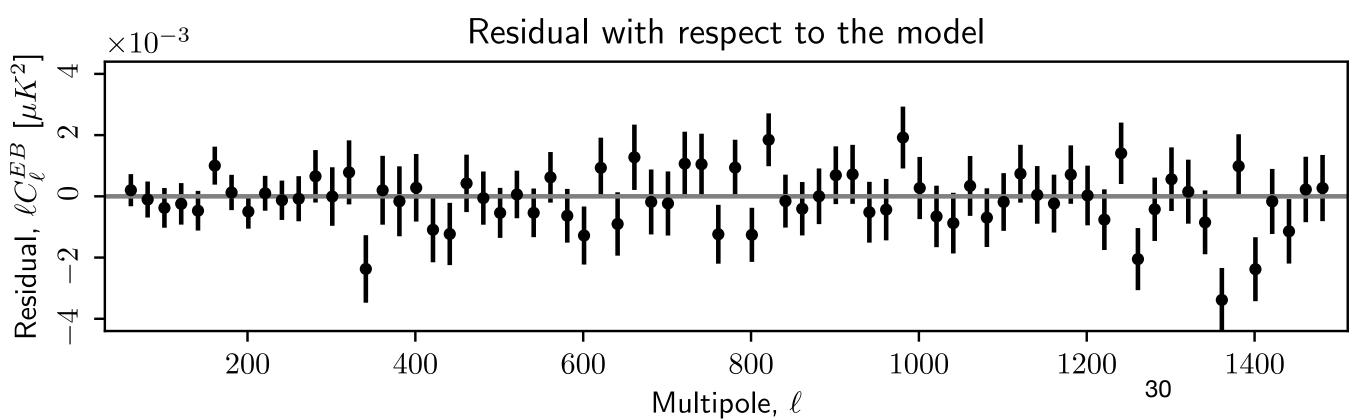
$$C_{\ell}^{EB,o} = \frac{\tan(4\beta)}{2} \left(C_{\ell}^{EE,o} - C_{\ell}^{BB,o} \right)$$

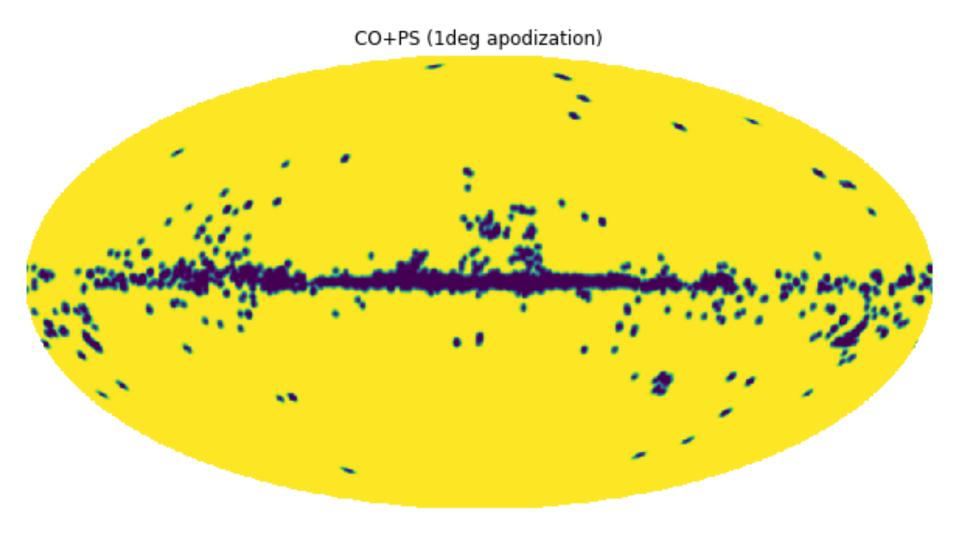


Cosmic Birefringence fits well(?) $C_{\ell}^{EB,o} = \frac{\tan(4\beta)}{2}$

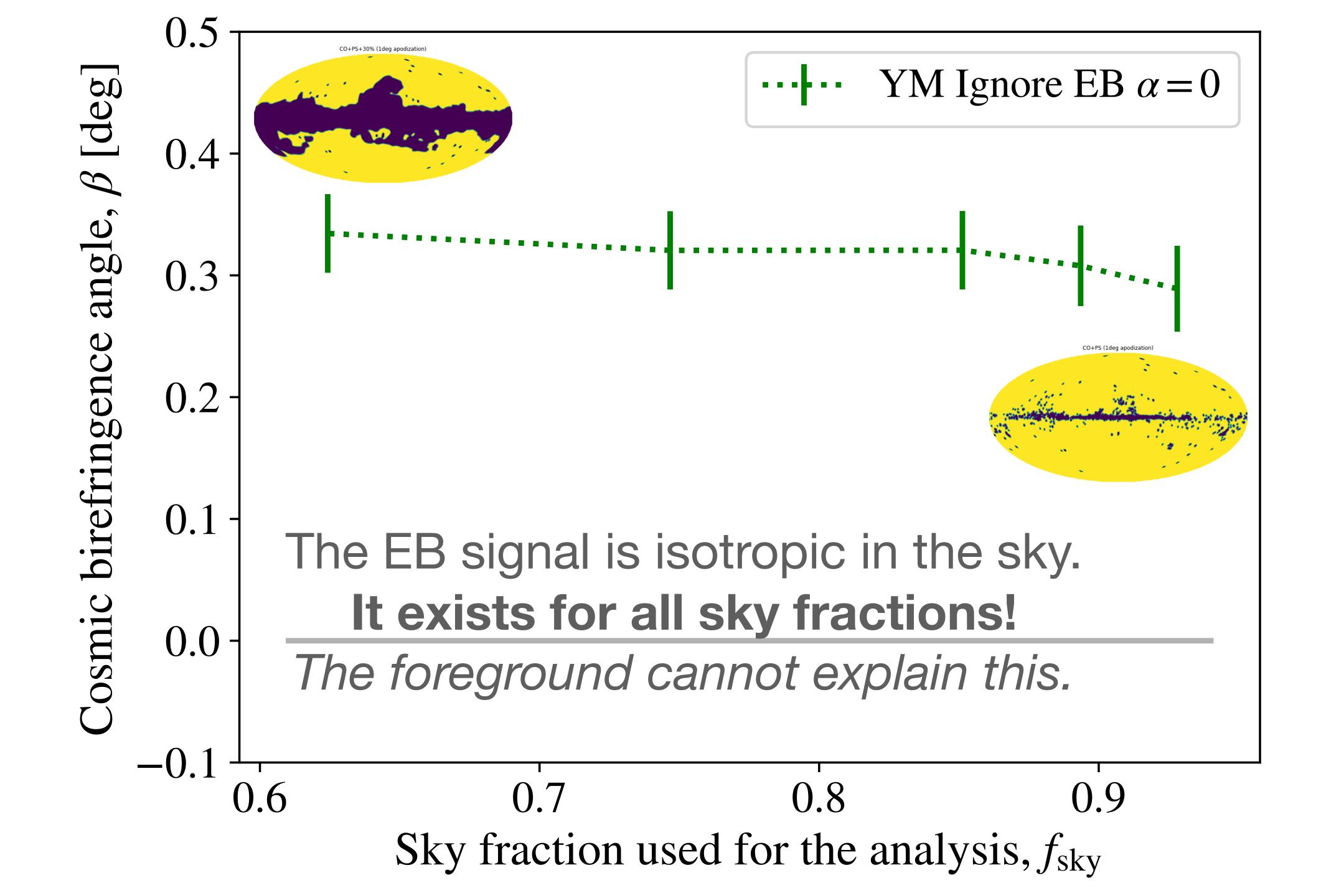
Nearly full-sky data (92% of the sky)







- $\beta = 0.288 \pm 0.032 \text{ deg}$
- $\chi^2 = 66.1$ for DOF=71
 - Good fit! 9σ detection?



The Biggest Problem: Miscalibration of detectors

Wu et al. (2009); Miller, Shimon, Keating (2009); EK et al. (2011)

Impact of miscalibration of polarization angles

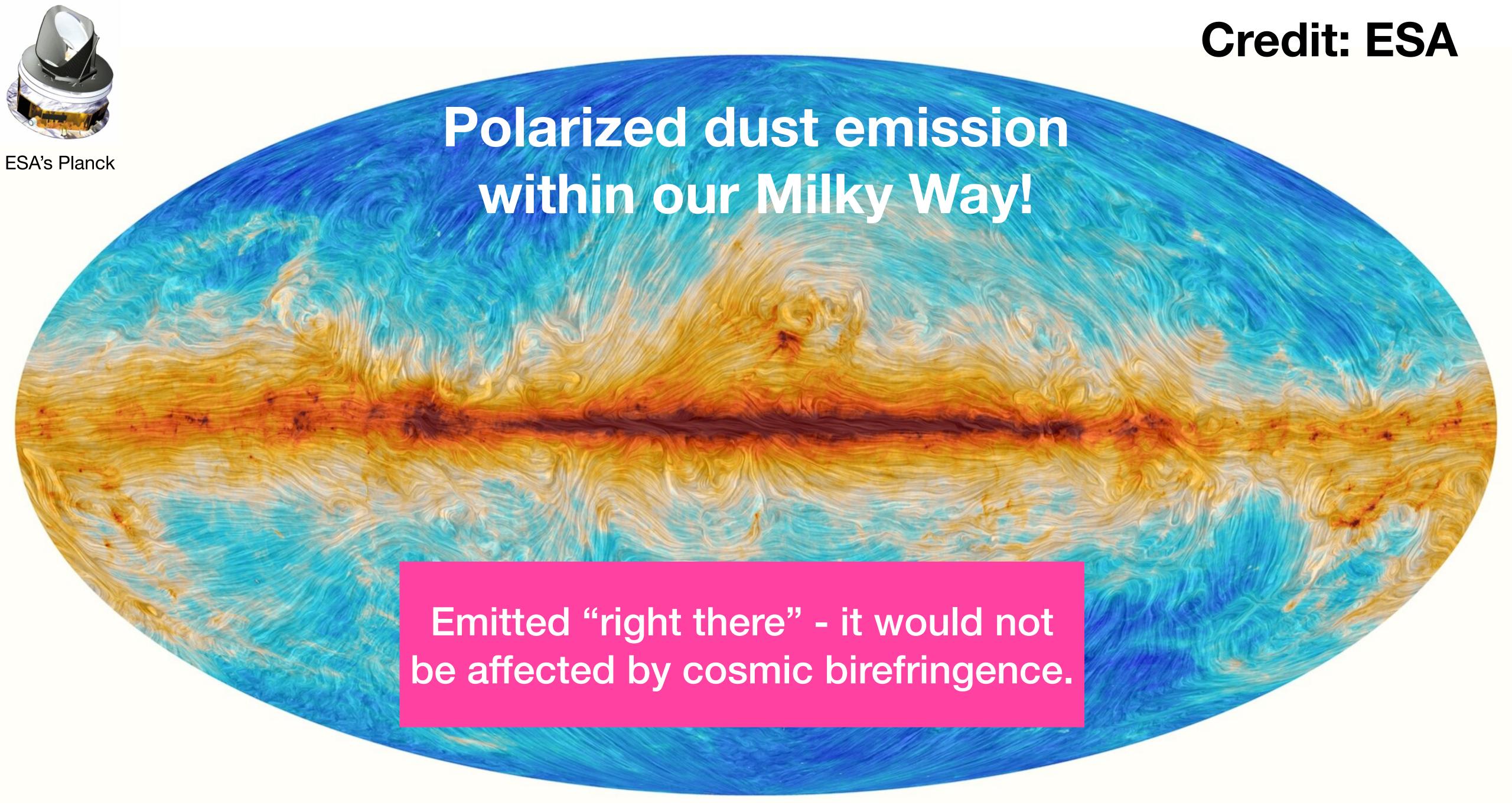
Cosmic or Instrumental?



- Is the plane of linear polarization rotated by the genuine cosmic birefringence effect, or simply because the polarization-sensitive directions of the detectors are rotated with respect to the sky coordinates (and we did not know it)?
- If the detectors are rotated by α , it seems that we can measure only the SUM $\alpha+\beta$.

Minami et al. (2019); Minami, EK (2020)

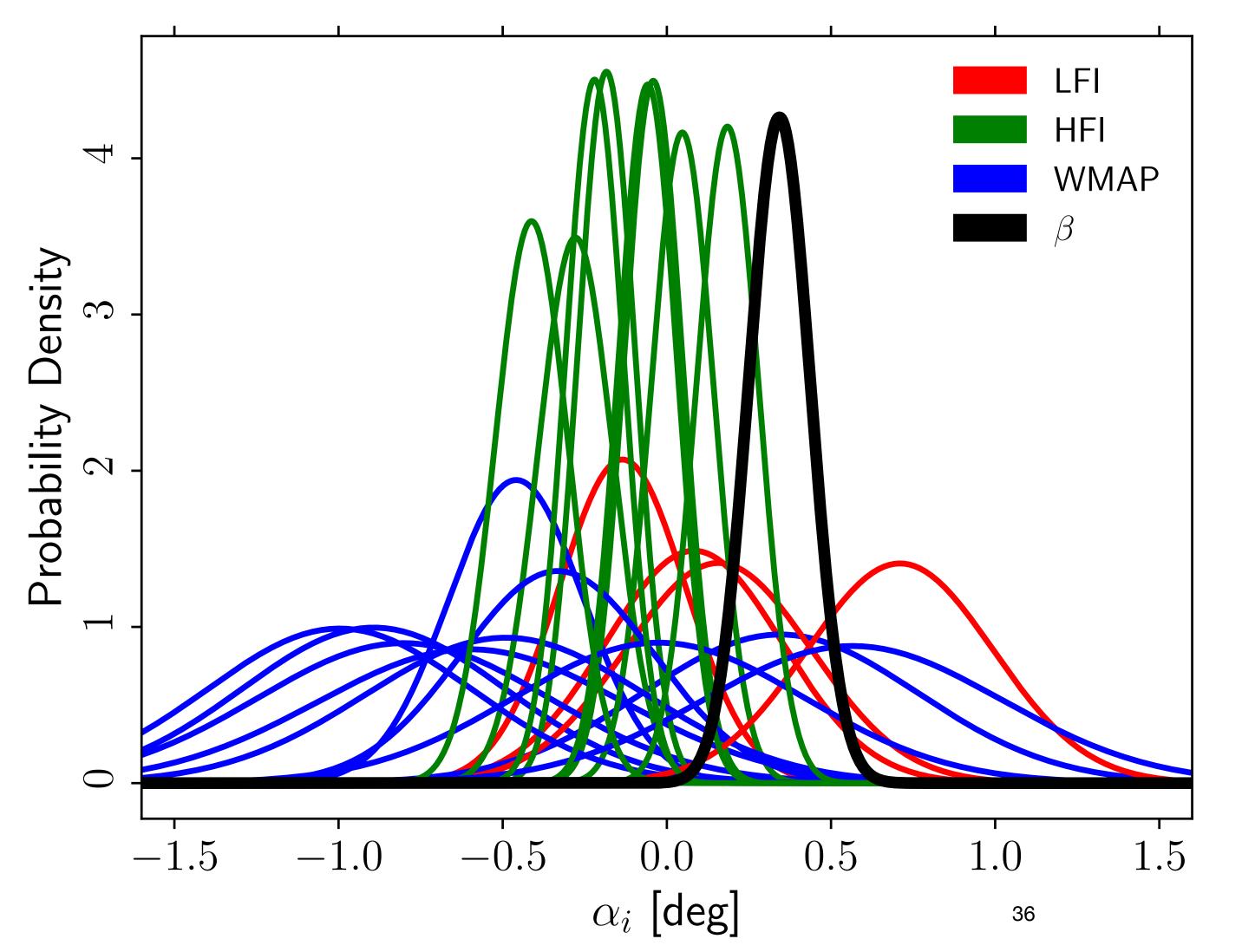
The Key Idea: The polarized Galactic foreground emission as a calibrator



Directions of the magnetic field inferred from polarization of the thermal dust emission in the Milky Way

Miscalibration angles (WMAP and Planck)

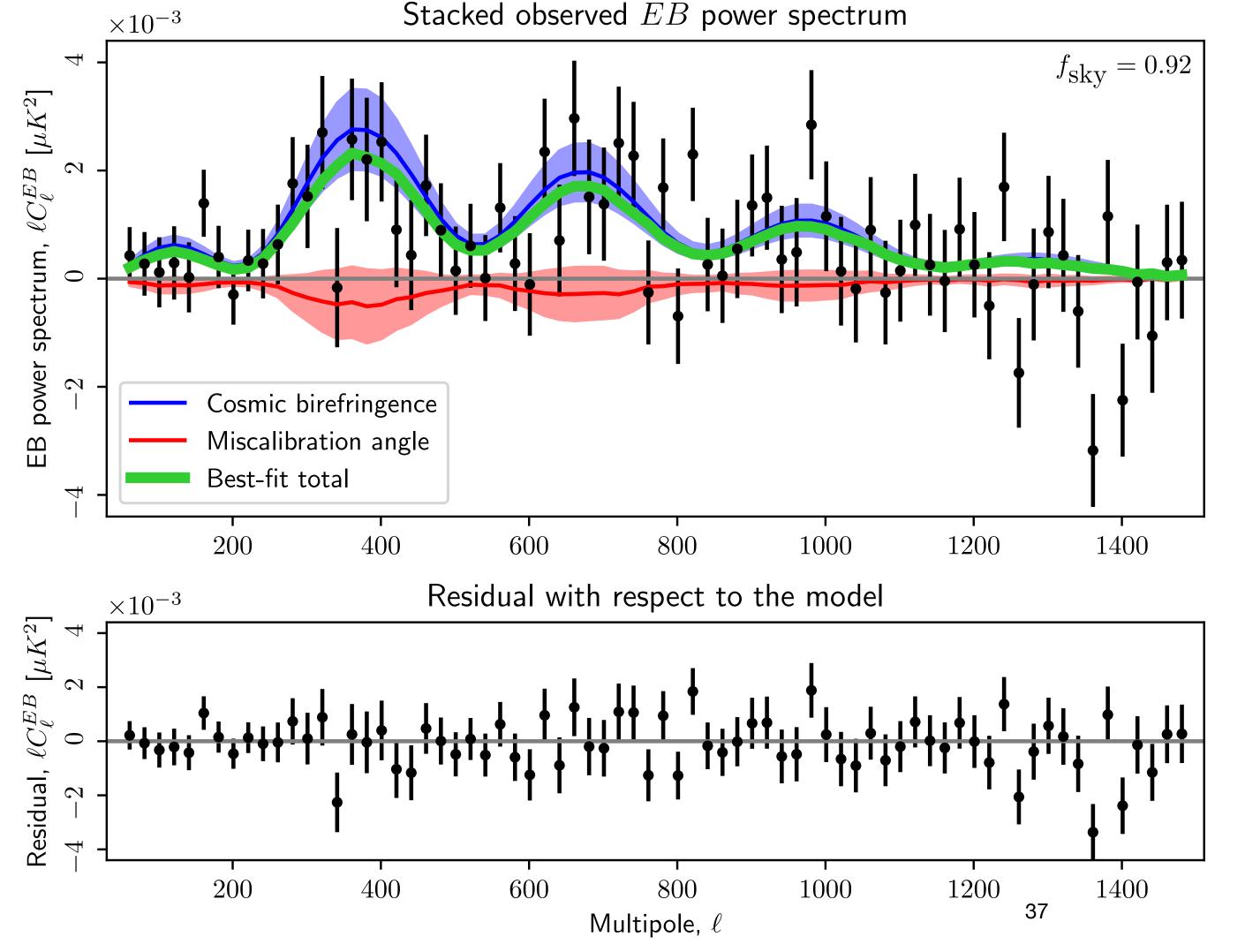
Nearly full-sky data (92% of the sky)

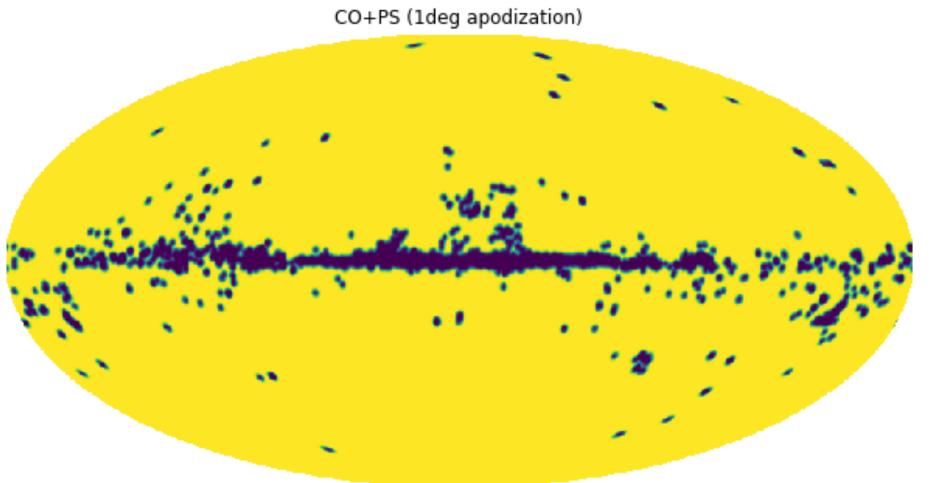


- The angles are all over the place, and are well within the quoted calibration uncertainty of instruments.
 - 1.5 deg for WMAP
 - 1 deg for Planck
- They cancel!
 - The power of adding independent datasets.

Cosmic Birefringence fits well (WMAP+Planck)

Nearly full-sky data (92% of the sky)

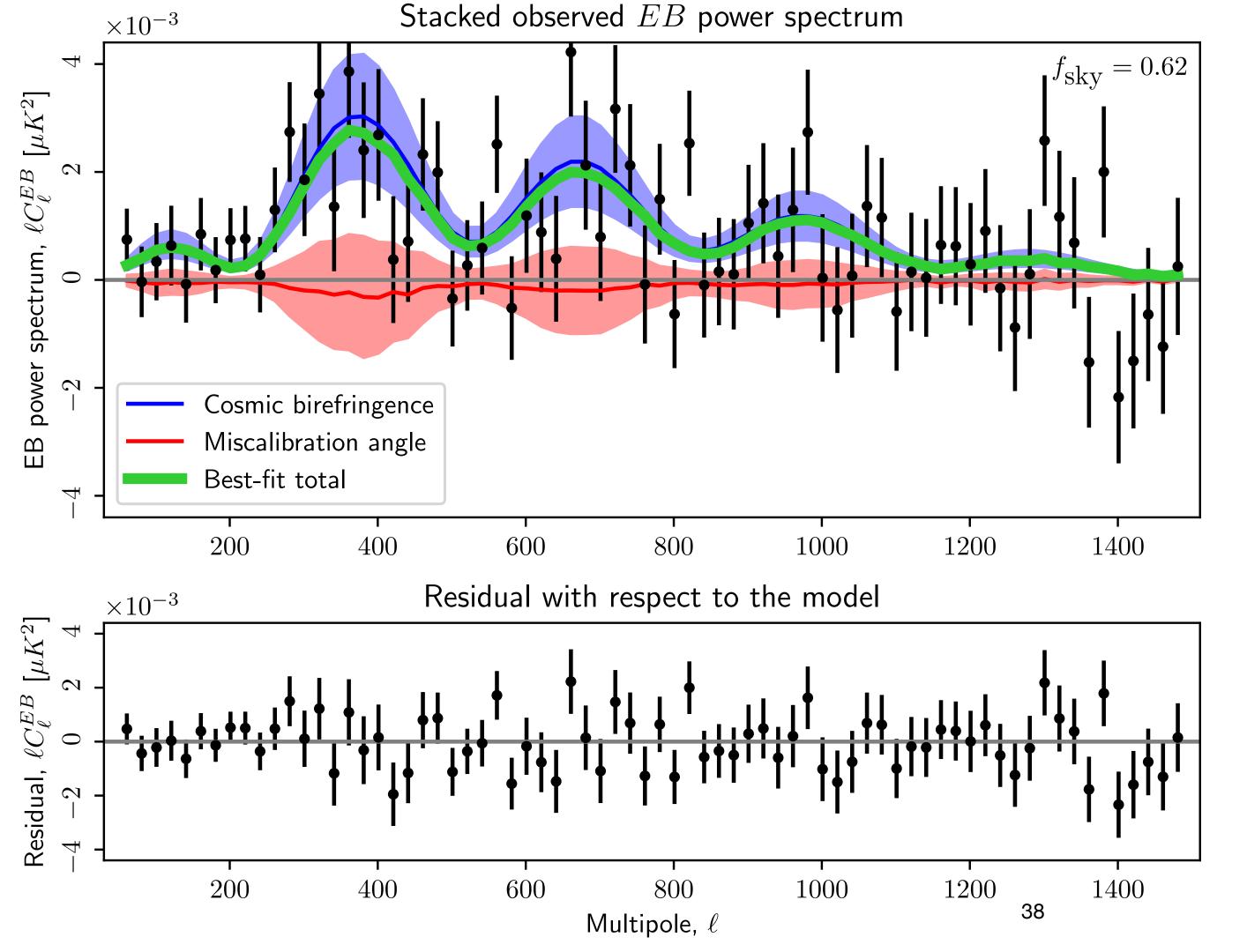


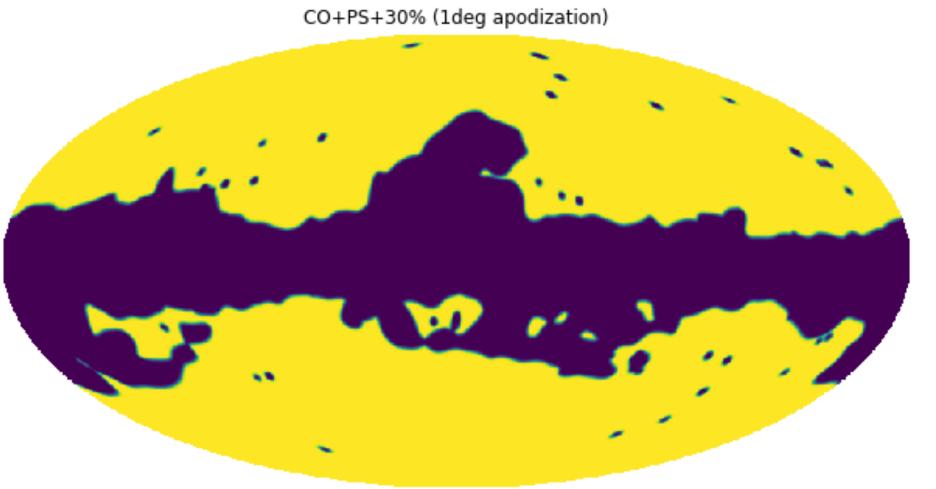


- Miscalibration angles make only small contributions thanks to the cancellation.
- $\beta = 0.34 \pm 0.09 \text{ deg}$
- $\chi^2 = 65.3$

Cosmic Birefringence fits well (WMAP+Planck)

Robust against the Galactic mask (62% of the sky)

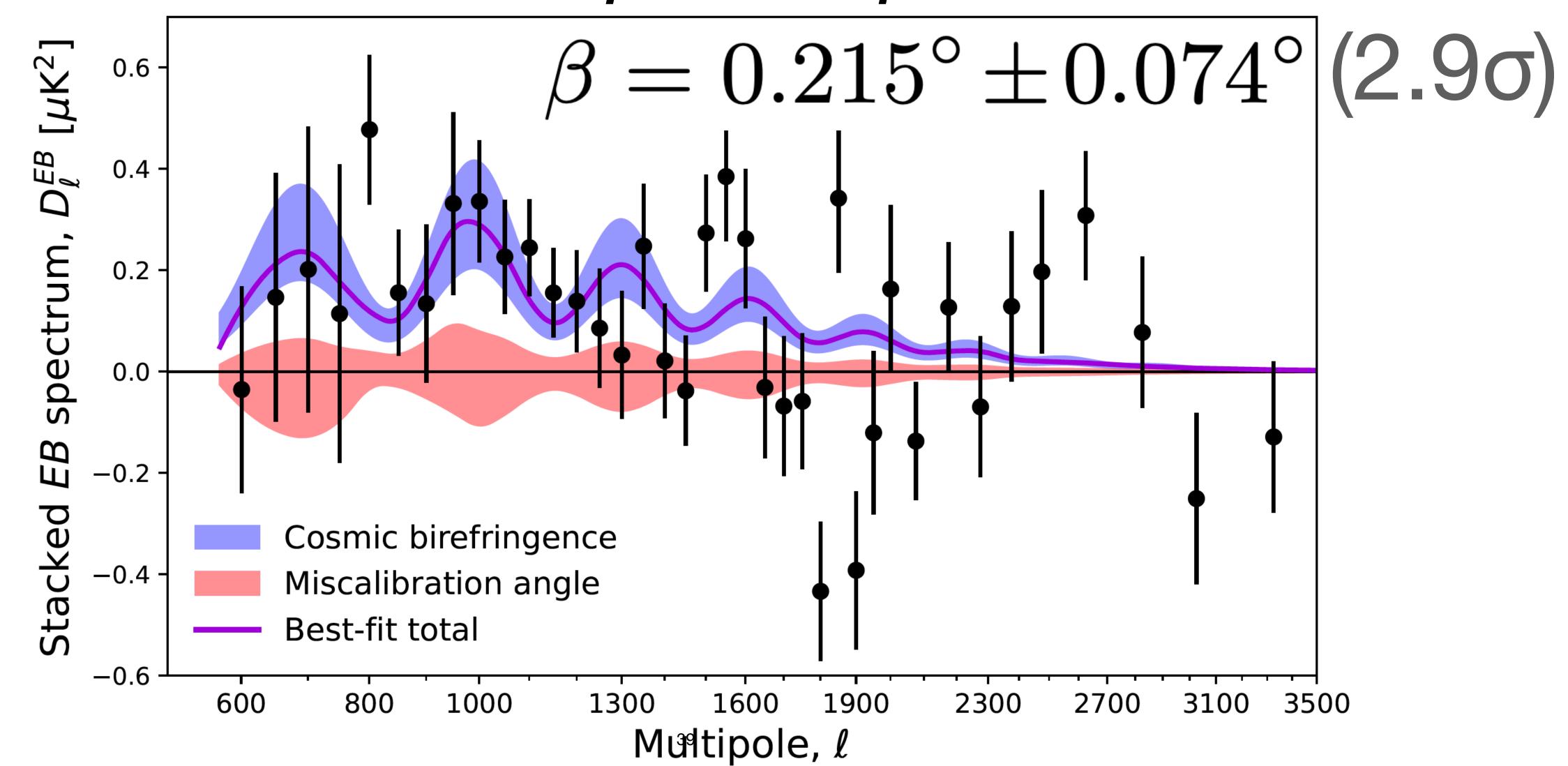


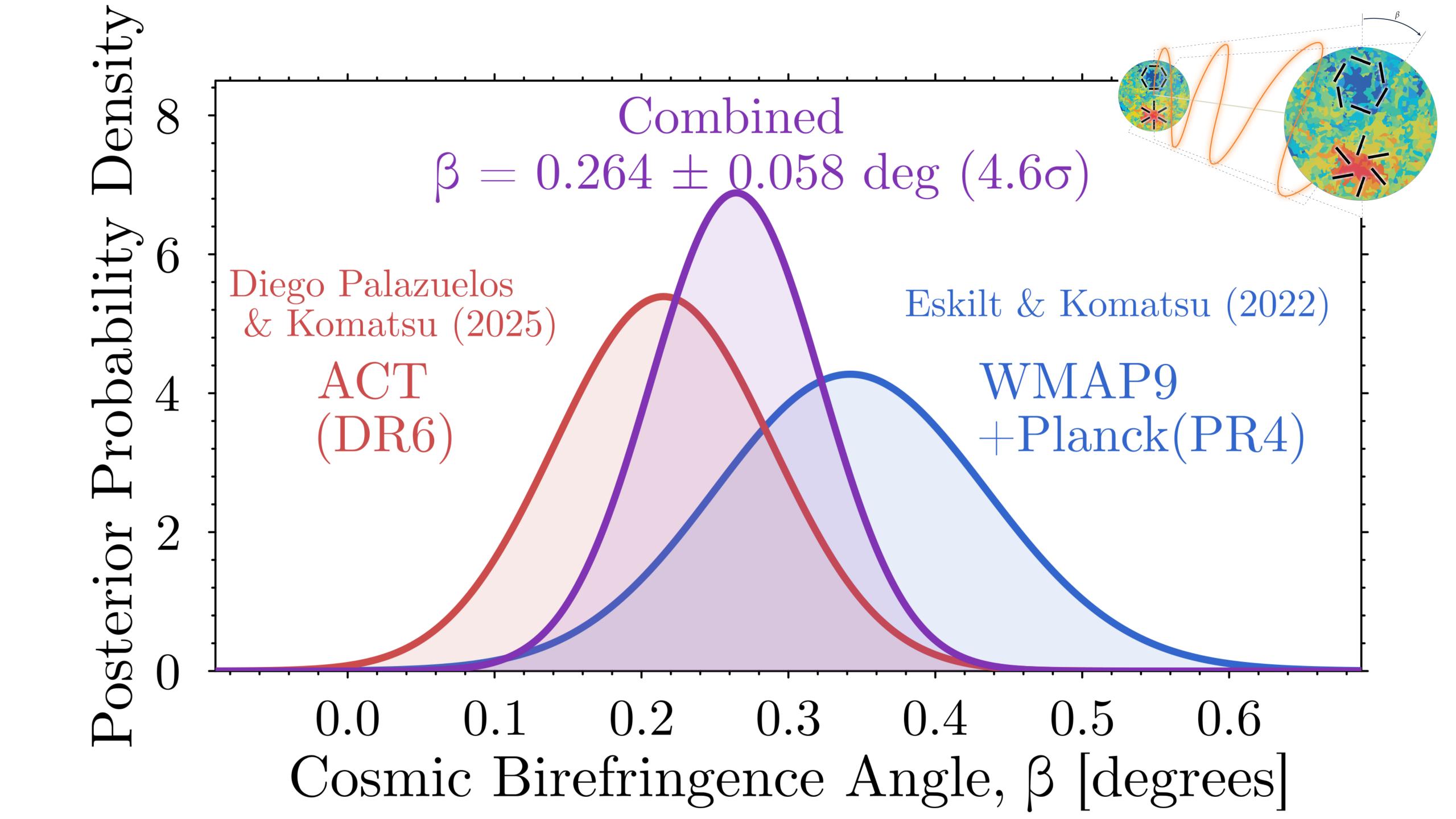


- Miscalibration angles make only small contributions thanks to the cancellation.
- $\beta = 0.37 \pm 0.14 \deg$
- $\chi^2 = 65.8$

Atacama Cosmology Telescope (DR6)

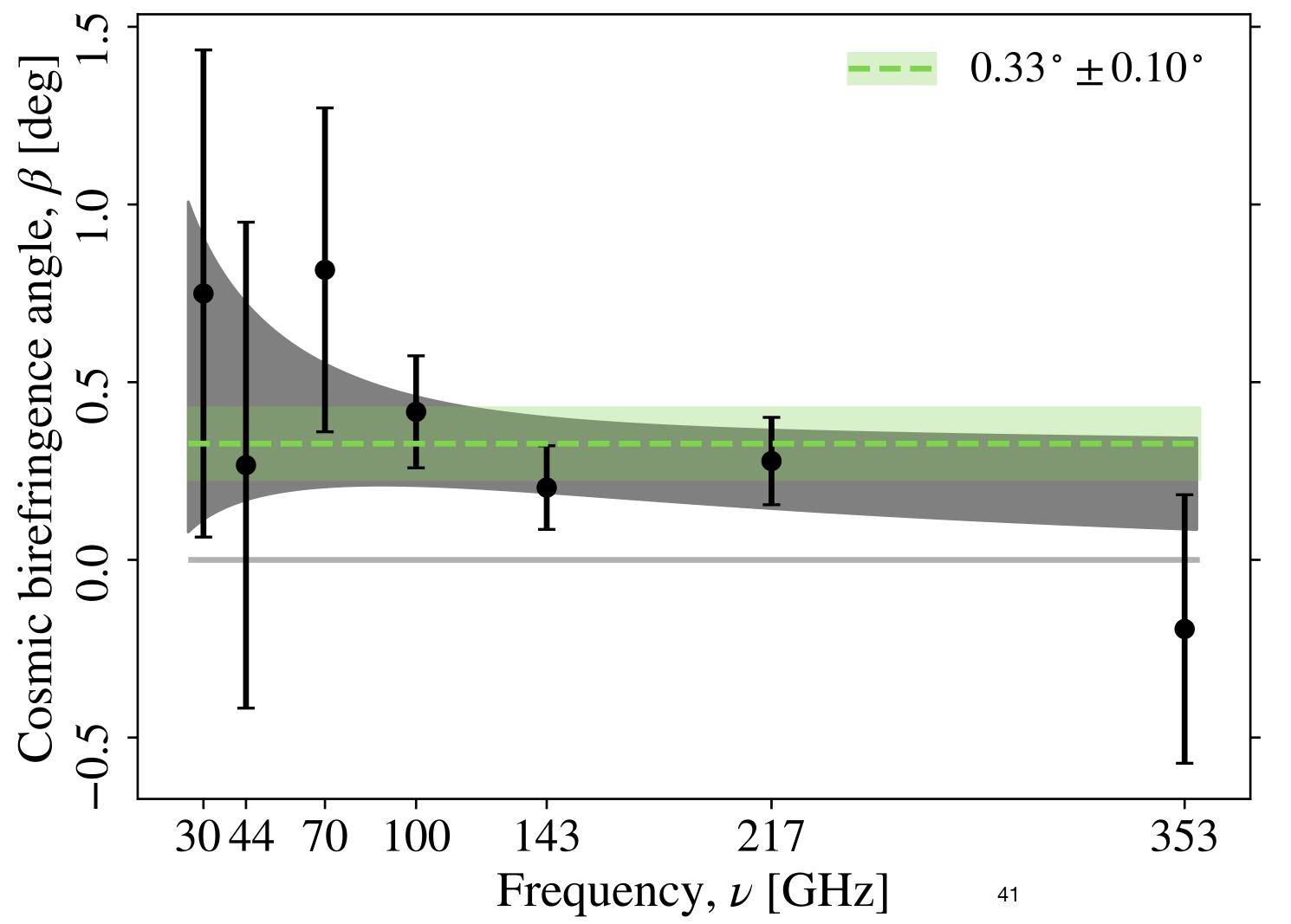
We need to measure it in independent experiments -> here it is!





No frequency dependence is found

It is not due to Faraday rotation.



- Light traveling in a uniform magnetic field also experiences a rotation of the plane of linear polarization, called "Faraday rotation". However, the rotation angle depends on the frequency, as $\beta(\nu) \propto \nu^{-2}$.
- No evidence for frequency dependence is found!
 - For $\beta \propto \nu^n$, $n = -0.20^{+0.41}_{-0.39}$ (68% CL)
 - Faraday rotation (n = -2) is disfavoured.

What are we worried about now?

"Unknown Unknowns"

WMAP+Planck

 The biggest worry: Unknown systematics in the Planck HFI at 353 GHz, since our results depend crucially on it.

ACT

- The biggest worry: The model for the optics of the ACT telescope and instruments may not capture all the systematics.
- The way forward: We will need another independent measurement, using an artificial polarization source. This will remove the dependence on any models.
 - BICEP3 (Cornelison et al., arXiv:2410.12089) and the Simons Observatory (Murata et al., arXiv:2309.02035) are doing exactly that. The final word is coming soon!

6. Implications for Dark Matter and Dark Energy

Imagine that space is filled with a pseudoscalar field coupled to photons via the CS term.

Scalar field DM/DE coupled to photons

DM = Dark Matter; DE = Dark Energy

$$I = \int d^4x \sqrt{-g} \left[-\frac{1}{2} (\partial \chi)^2 - V(\chi) - \frac{1}{4} F^2 - \frac{\alpha}{4f} \chi F \tilde{F} \right]$$

- χ is a neutral pseudoscalar field (spin 0).
- The scalar field sector: DM or DR.
- Why consider χ as a good DM/DE candidate?
 - Why not? We have an example in the Standard Model: a neutral pion.
 - We expect $\alpha \simeq \alpha_{\rm EM} \simeq 10^{-2}$ and $f < M_{\rm Pl} \simeq 2.4 \times 10^{18}$ GeV.
- χ can be composed of fermions like a pion, or a fundamental pseudoscalar like an "axion" field.

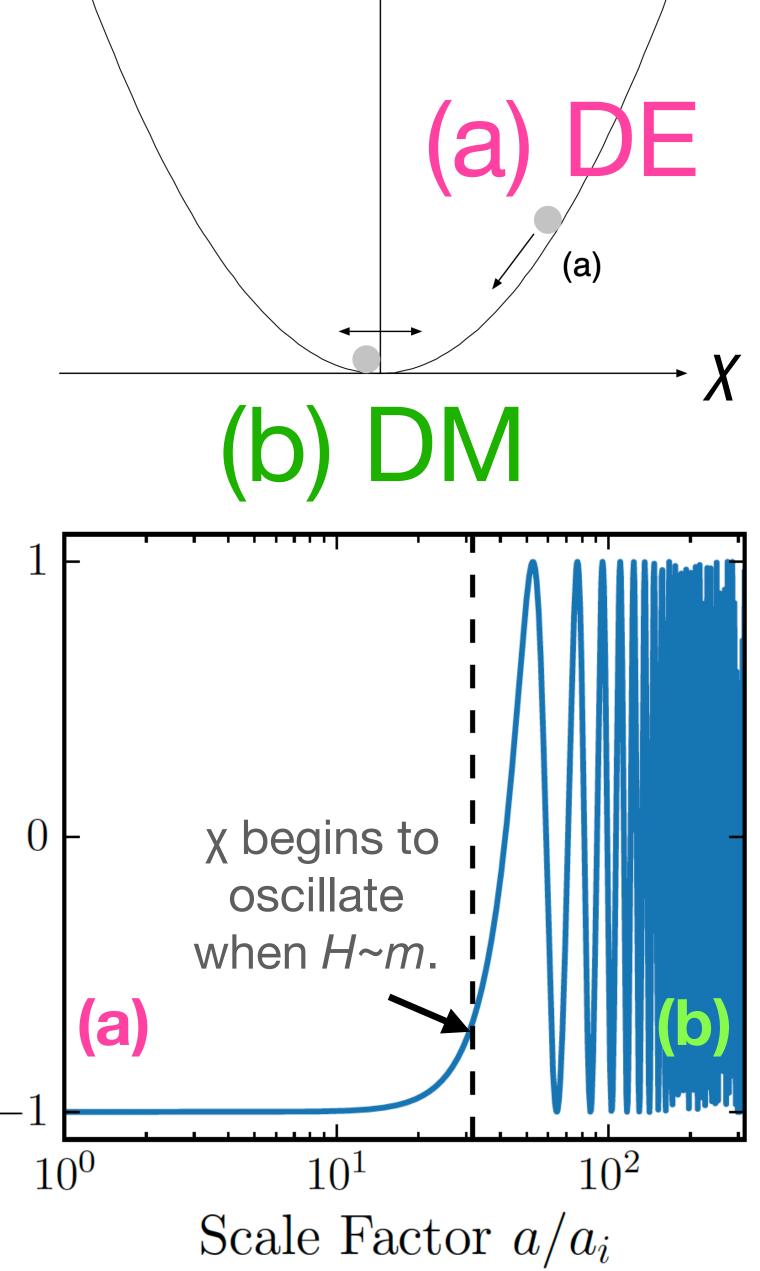
Distinction between DE and DM

How small is its mass? Example of $V(\chi) = m^2 \chi^2/2$

• The useful criterion is the equation of state parameter, w.

$$w = \frac{P}{\rho} = \frac{\langle \dot{\chi}^2 \rangle - m^2 \langle \chi^2 \rangle}{\langle \dot{\chi}^2 \rangle + m^2 \langle \chi^2 \rangle}$$

- $w \simeq -1$: Dark Energy (DE)
 - $m \lesssim H_0 \simeq 10^{-33} \text{ eV}$
- $w \simeq 0$: Dark Matter (DM)
 - $m \gtrsim H_0$



 $i^2\chi^2)/(\dot{\chi}^2$

 $V(\chi)$

Imagine that space is filled with a pseudoscalar field coupled to photons via the CS term.

Scalar field DM/DE coupled to photons

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$$I = \int d^4x \sqrt{-g} \left[-\frac{1}{2} (\partial \chi)^2 - V(\chi) - \frac{1}{4} F^2 - \frac{\alpha}{4f} \chi F \tilde{F} \right]$$

The electromagnetic sector: Coupling to photons via the "Chern-Simons term".

$$F^2 \equiv F_{\mu\nu}F^{\mu\nu} = 2({f B}\cdot{f B}-{f E}\cdot{f E})$$
 This is a scalar and is invariant under parity transformation.

$$F\tilde{F} \equiv F_{\mu\nu}\tilde{F}^{\mu\nu} = -4\mathbf{B} \cdot \mathbf{E}$$

This is a *pseudoscalar* and changes sign under parity transformation!

Ni (1977); Sikivie (1983); Turner, Widrow (1987) Carroll, Field, Jackiw (1990)

FF in the action

Chern-Simons term

• Consider
$$I_{\text{CS}} = -\frac{1}{4} \alpha \int d^4 x \; \theta F \tilde{F}$$
 with $F \tilde{F} = 2 \partial_{\mu} (A_{\nu} \tilde{F}^{\mu \nu})$
• α : a dimensionless constant
• θ : a dimensionless (pseudo)scalar field

We wrote $\theta = \frac{\chi}{f}$

Why Chern-Simons Terms?

Published Superature of the resiment Terms of the state of the resiment Terms of the r

- θ : a dimensionless (pseudo)scalar field
- This is not a surface term! Integration by parts gives

$$I_{\rm CS} = \frac{1}{2} \alpha \int d^4 x \ (\partial_{\mu} \theta) A_{\nu} \tilde{F}^{\mu\nu}$$

Jim Simons in 2023

https://einstein-chair.github.io/simons2023/

• This is a special case of the so-called Chern-Simons term(*), $p_{\mu}A_{\nu}F^{\mu\nu}$ (*) Strictly speaking, Chern-Simons 3-form for an Abelian gauge field is $A_
u ilde{F}^{0
u} = A_i e^{ijk} \partial_i A_
u$ with $p_\mu = \partial_\mu heta$ Adler (1969); Bell, Jackiw (1969); Fujikawa (1979)

Is there a known example of this term in particle physics?

Yes, a pion. The ABJ anomaly!



- A pion is a composite meson composed of a quark and an antiquark.
 - A neutral pion, π^0 , is composed of either $u\bar{u}$ or $d\bar{d}$, and is a pseudoscalar.

(Chinowsky & Steinberger, 1954)

- π⁰ is coupled to photons via l_{CS} where
 - $\theta = \pi^0 / f_{\pi}$ with $f_{\pi} \sim 184$ MeV (pion decay constant)
 - $\alpha = 2\alpha_{EM}N_c/(3\pi)$ with $N_c = 3$ (the number of quark colors) and $\alpha_{EM} \sim 1/137$ (EM fine structure constant)
- π⁰ decays into 2 photons via this term, which has been observed. So, this
 possibility is not completely crazy!

Multiple discoveries of $I_{\rm CS}=-\frac{1}{4}\alpha\int d^4x\;\theta F\tilde{F}$ The presence of this term is well motivated.

- This electromagnetic coupling term has been discovered at least 4 times in the past.
 - 1969: The ABJ anomaly [Adler, Bell, Jackiw]
 - 1974: Chern-Simons 3-form [Chern, Simons]
 - Promoted to 4-dimensional theory in 1990 [Carroll, Field, Jackiw]
 - 1977: Equivalence Principle [Ni (倪維斗)]
 - 1983: Axion electrodynamics [Sikivie]

Implications

DM = Dark Matter; DE = Dark Energy

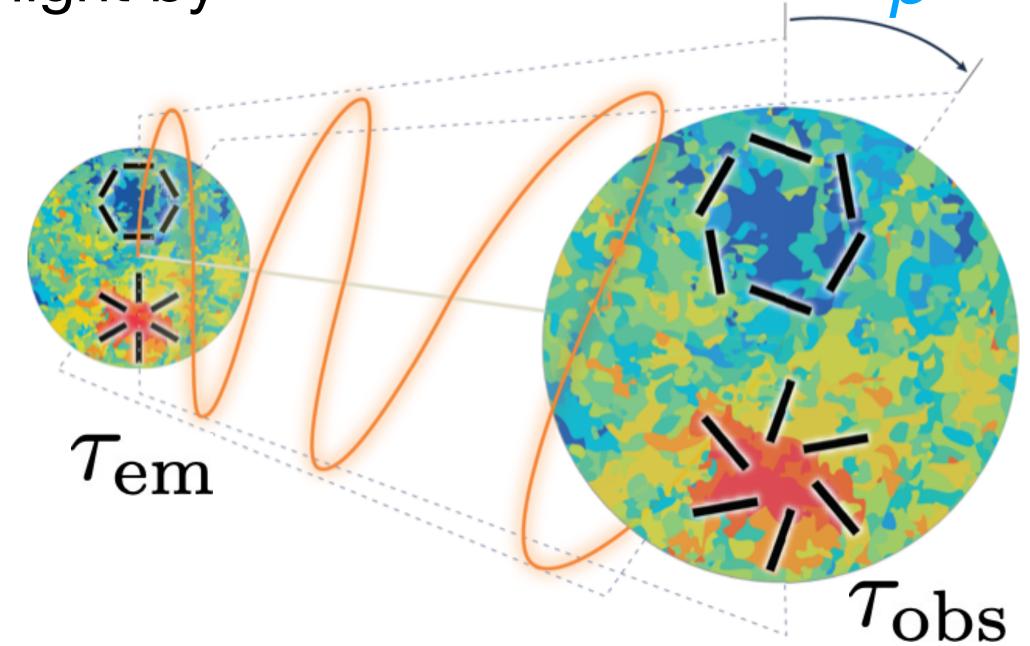
This term exists for a pion.

What if DM/DE is "pion-like particle"?

$$I = \int d^4x \sqrt{-g} \left[-\frac{1}{2} (\partial \chi)^2 - V(\chi) - \frac{1}{4} F^2 - \frac{\alpha}{4f} \chi F \tilde{F} \right]$$

• This rotates the plane of linear polarization of light by

$$eta = -\int_{ au_{ ext{em}}}^{ au_{ ext{obs}}} d au \left(\omega_{+} - \omega_{-}
ight) \ = rac{lpha}{2f} \left[\chi(au_{ ext{obs}}) - \chi(au_{ ext{em}})
ight]^{ au_{ ext{em}}}$$



Implications

DM = Dark Matter; DE = Dark Energy

This term exists for a pion.

What if DM/DE is "pion-like particle"?

$$I = \int d^4x \sqrt{-g} \left[-\frac{1}{2} (\partial \chi)^2 - V(\chi) - \frac{1}{4} F^2 - \frac{\alpha}{4f} \chi F \tilde{F} \right]$$

• The measured angle, β, implies that the field has evolved by

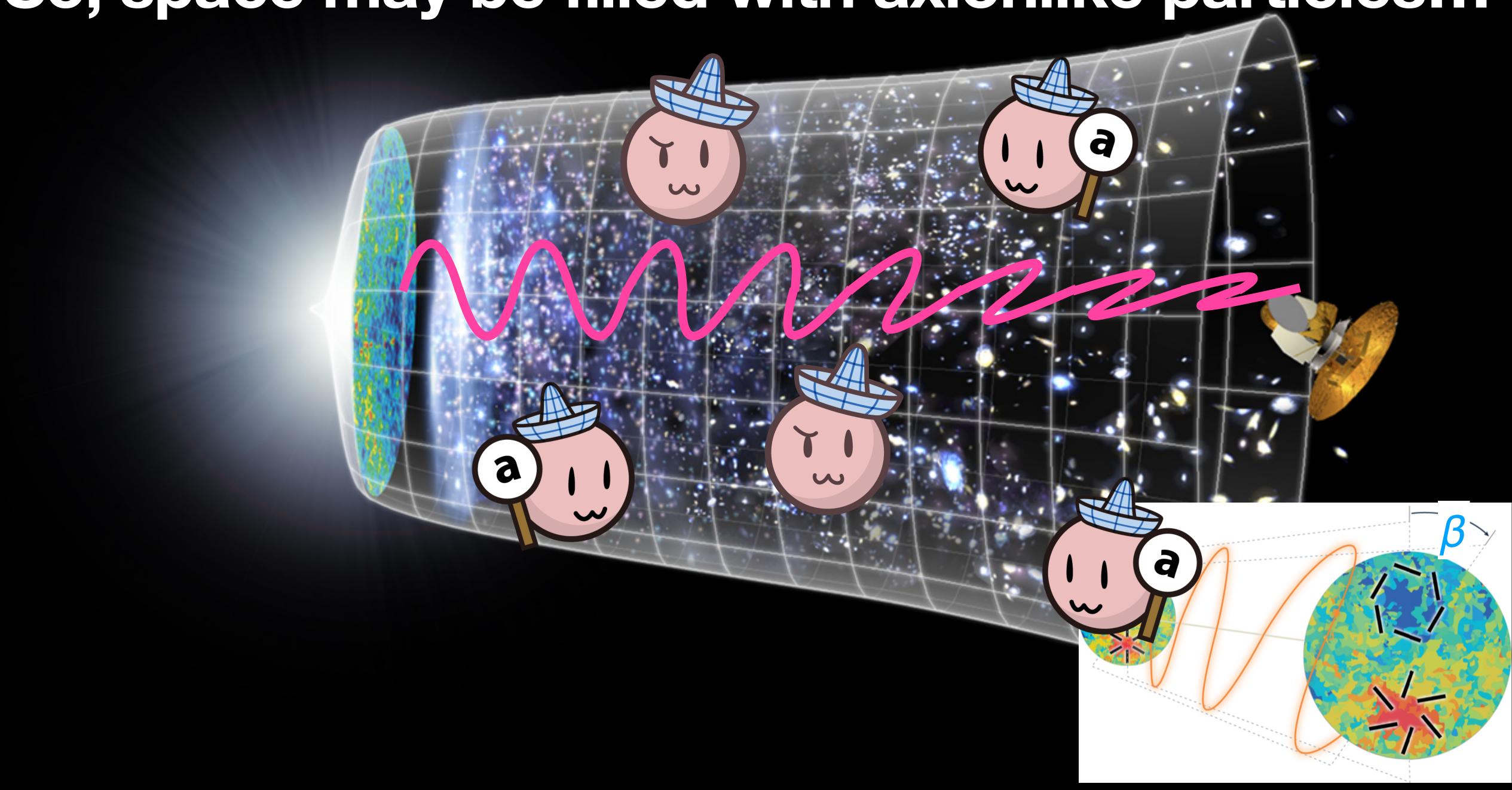
$$\Delta \chi = \chi(\tau_{\rm obs}) - \chi(\tau_{\rm em}) \simeq \frac{10^{-2}}{\alpha} f$$

Axionlike particle?

- If it is due to DE: this measurement rules out DE being a cosmological constant.
- If it is due to DM: at least a fraction of DM violates parity symmetry.



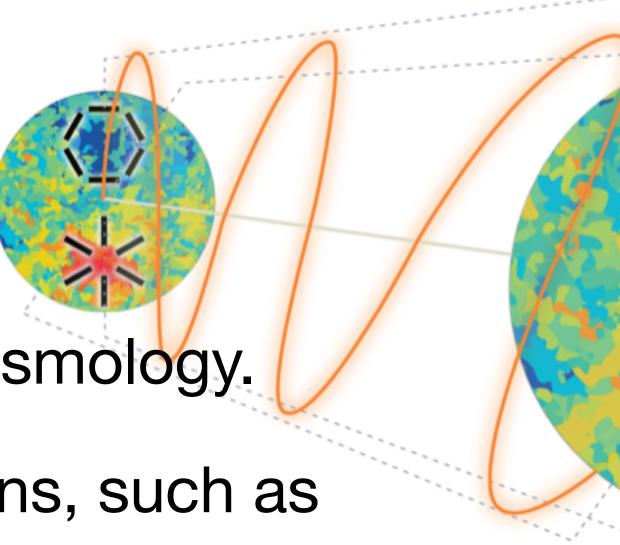
So, space may be filled with axionlike particles...



Summary

Let's find new physics!

- Violation of parity symmetry is a new topic in cosmology.
- It may hold the answers to fundamental questions, such as
 - What is Dark Matter?
 - What is Dark Energy?
- Since parity is violated in the weak interaction, it seems natural to expect that parity is also violated in the Dark Sector.
 - 4σ hint of Cosmic Birefringence: Space may be filled with parity-violating DM and DE fields?
- What else should we be looking? New and exciting research topics.



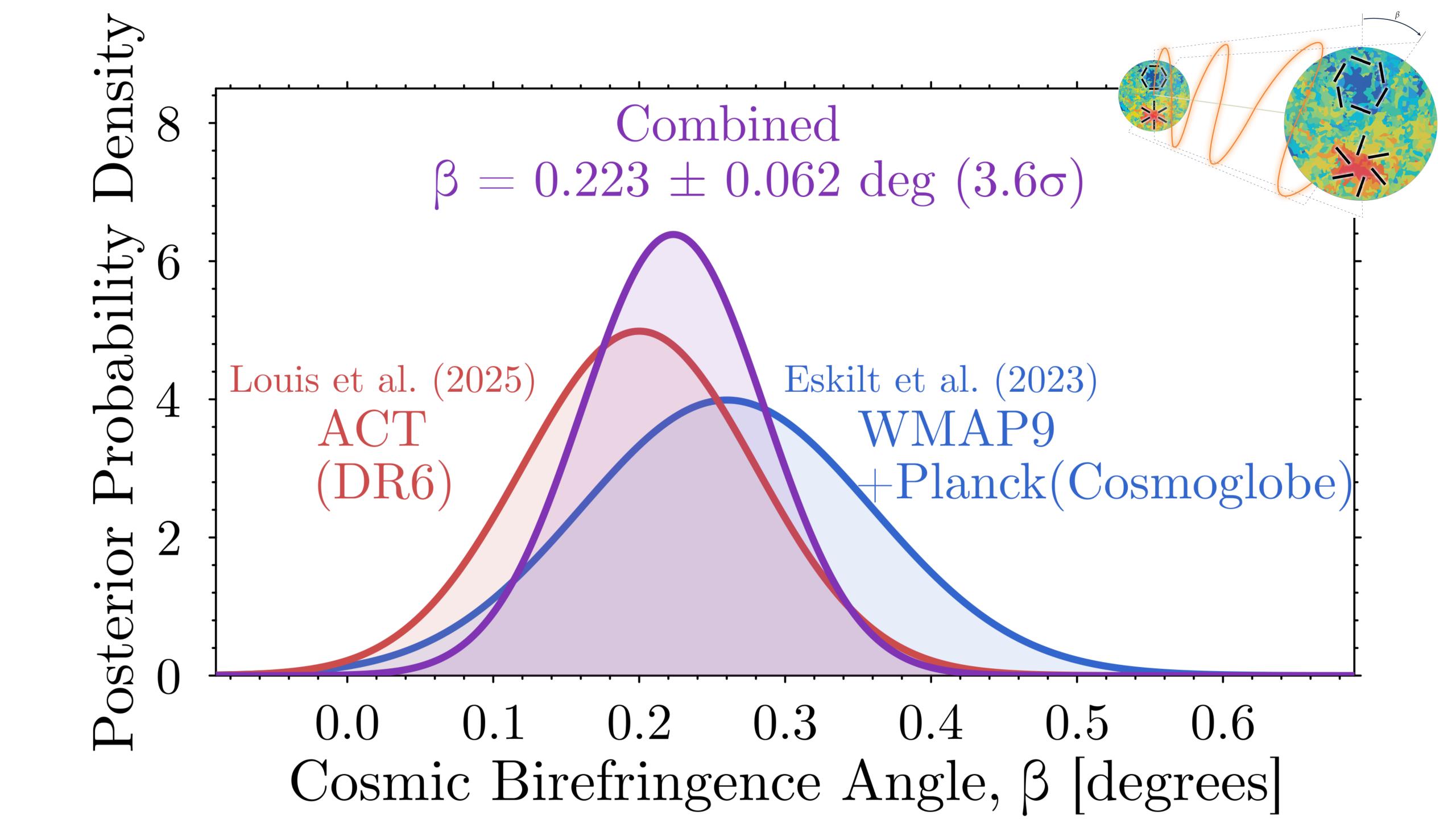
Back up slides

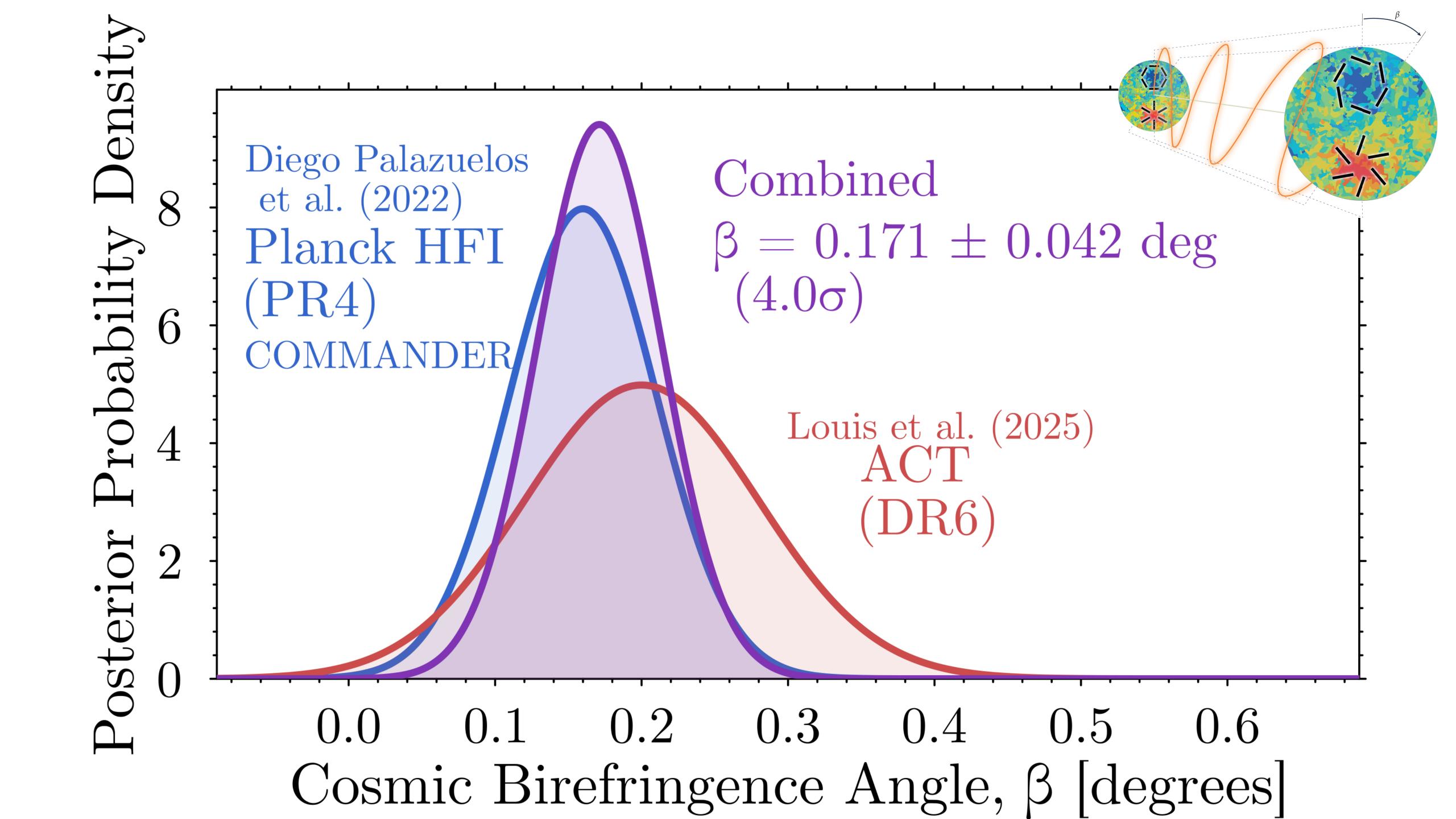
Diego-Palazuelos et al. (2022, 2023); Eskilt et al. [Cosmoglobe] (2023)

Is \(\beta \) caused by non-cosmological effects?

We need to measure it in independent experiments.

- The known instrumental effects of the WMAP and Planck missions are shown to have negligible effects on β .
 - However, we can never rule out **unknown** instrumental effects... We need to measure β in independent experiments.
- The polarized Galactic foreground emission was used to calibrate the instrumental polarization angles, α. The intrinsic EB correlations of the Galactic foreground emission (polarized dust and synchrotron emission) could affect the results.
 - We need to measure β without relying on the foreground by calibrating α well, e.g., Murata et al. (Simons Observatory), arXiv:2309.02035; Murphy et al. (ACT), arXiv:2403.00763; Cornelison et al. (BICEP3), arXiv:2410.12089; Ritacco et al. (COSMOCal), arXiv:2405.12135.





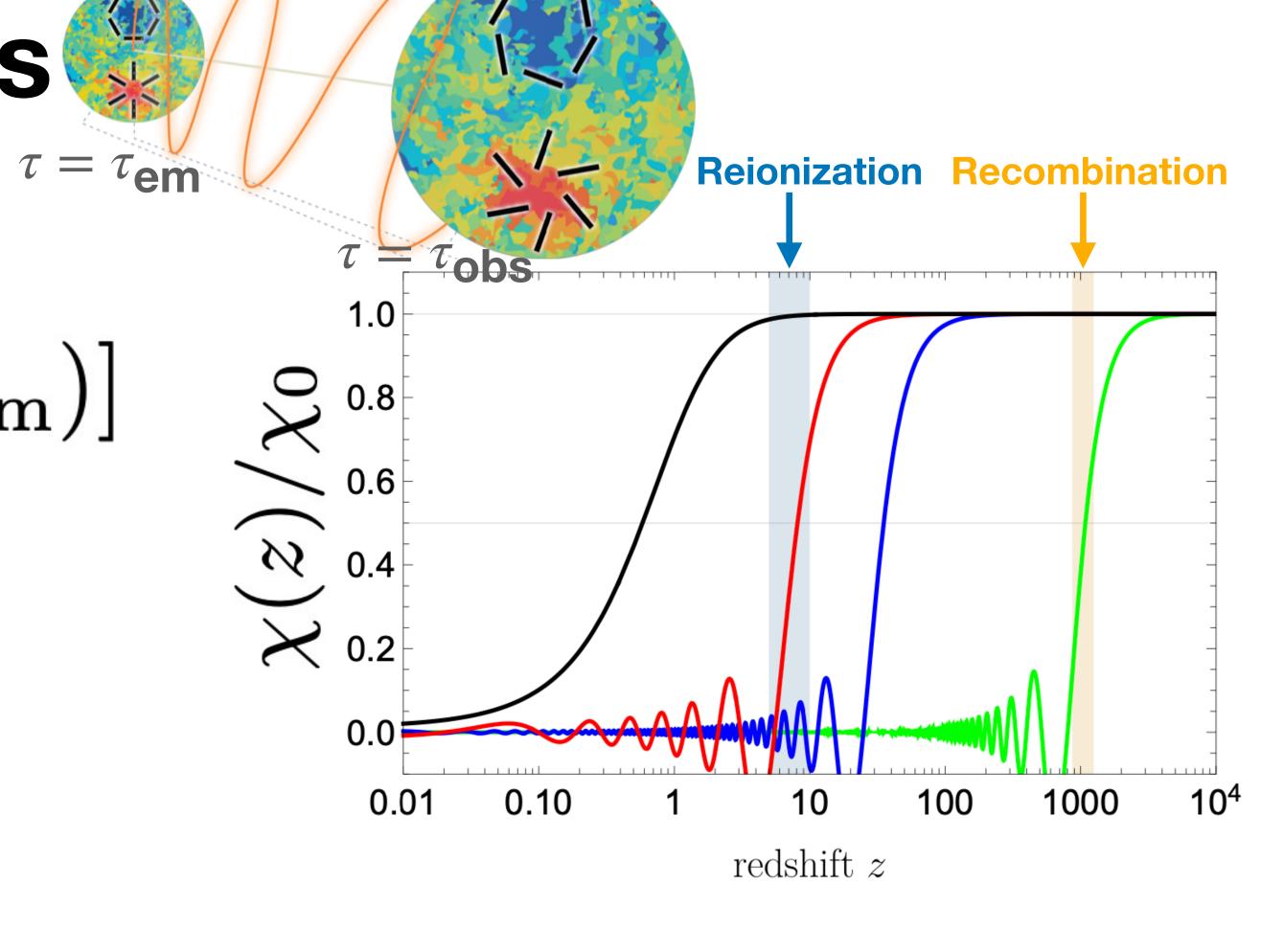
Is cosmic birefringence due to dark matter or dark energy?

Nakatsuka, Namikawa, EK (2022)

How to measure mass

$$\beta = +\frac{\alpha}{2f} \left[\chi(\tau_{\rm obs}) - \chi(\tau_{\rm em}) \right]$$

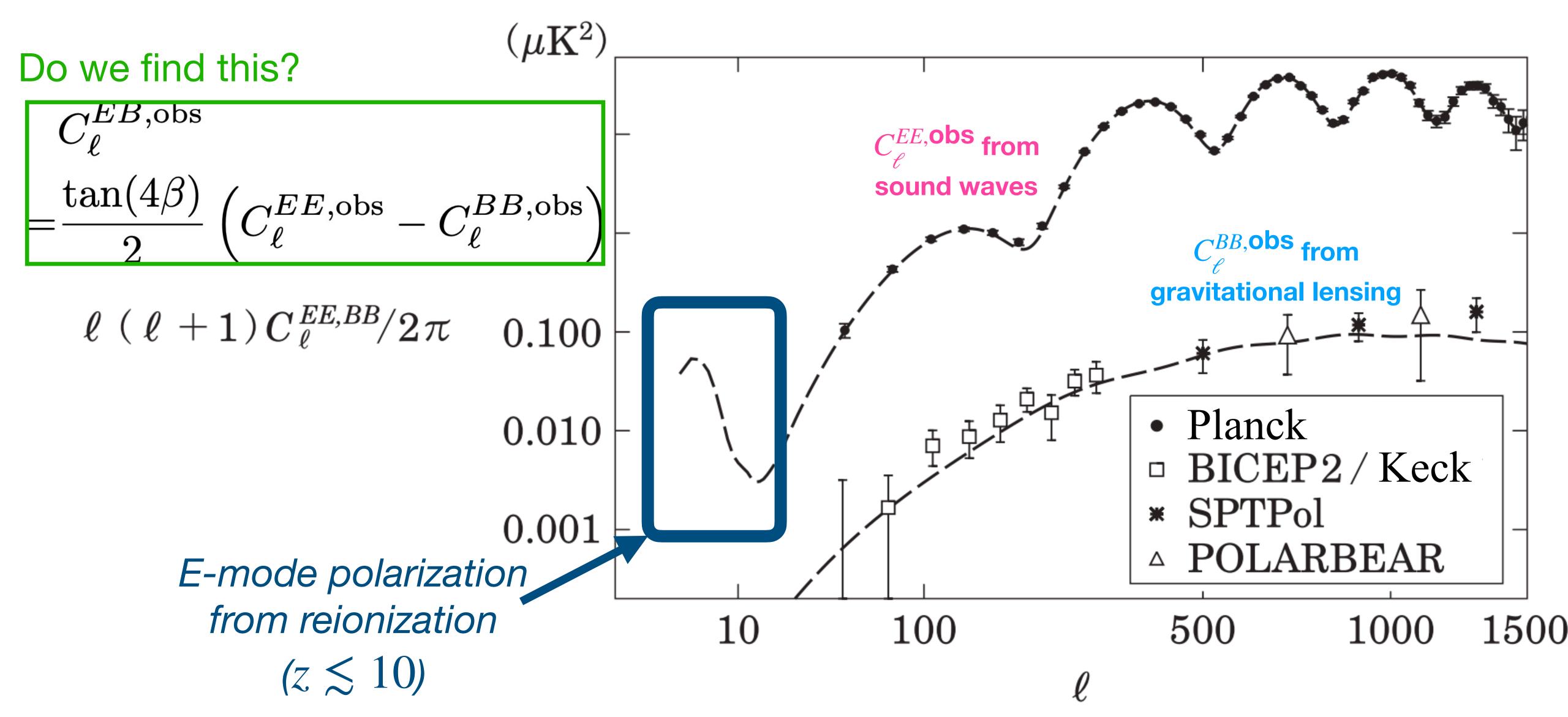
- There are 2 epochs when the CMB polarization was produced.
 - z~1100: Recombination
 - z~10: Reionization
- β from these 2 epochs can be different!



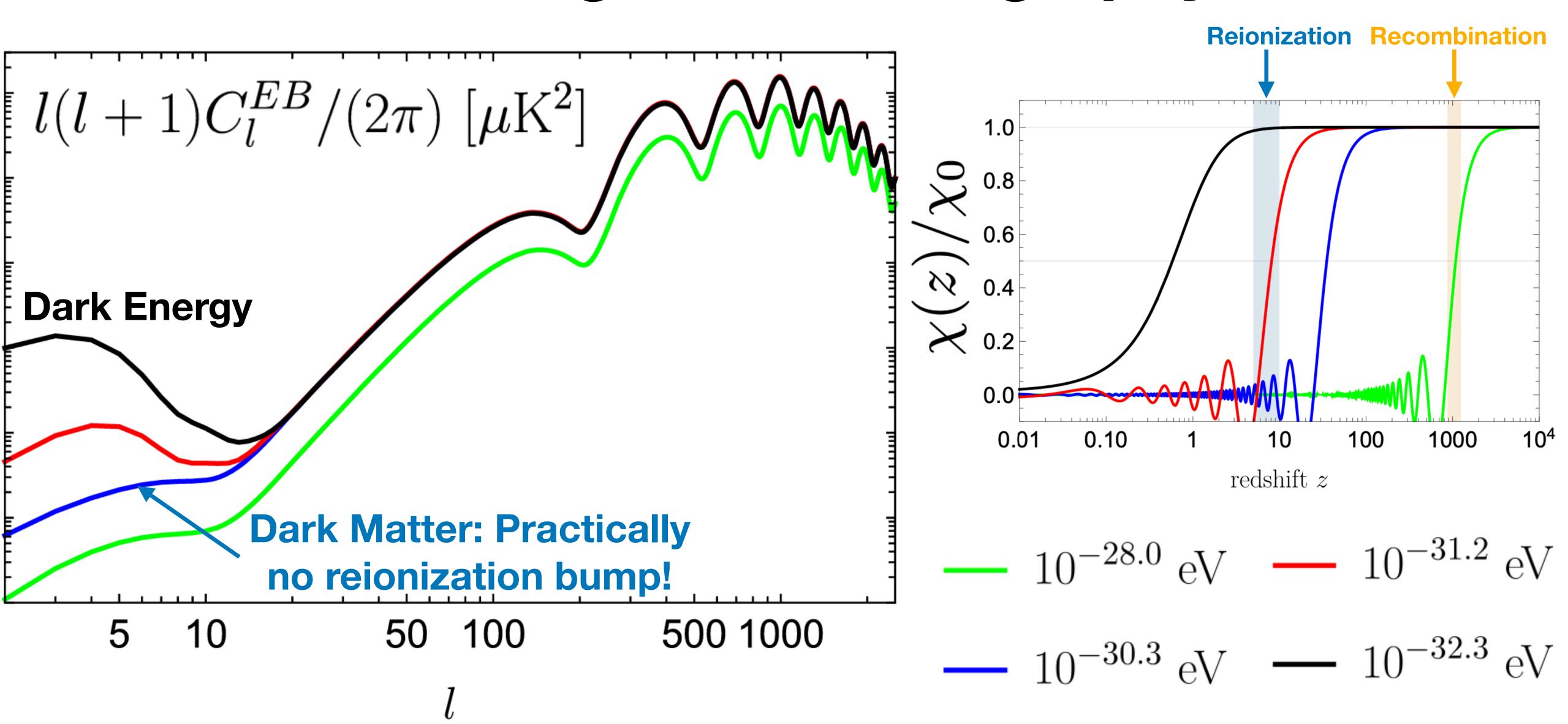
$$-- 10^{-28.0} \text{ eV} -- 10^{-31.2} \text{ eV}$$

$$-10^{-30.3} \text{ eV} - 10^{-32.3} \text{ eV}$$

"Reionization bump" at low multipoles ($\ell \lesssim 10$)



Cosmic Birefringence "Tomography"



Time-dependent Rotation Angle

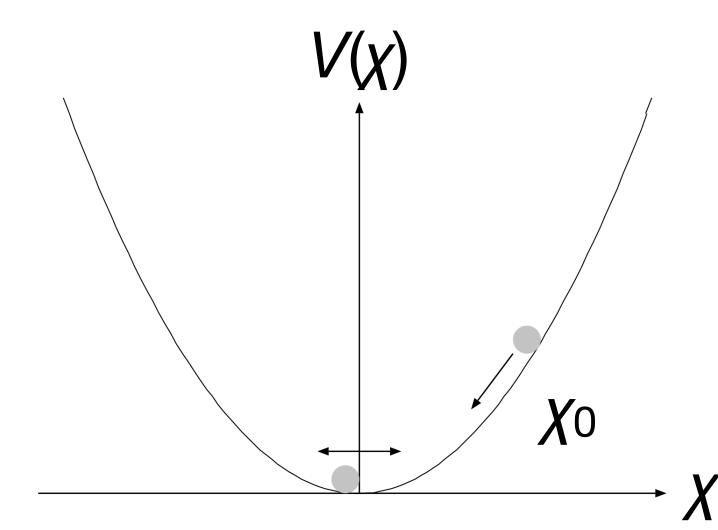
Pressure of a massive free scalar field

$$V(x) = m^2x^2/2$$
 (c=1 and \hbar =1)

- To simplify notation, we will omit the overline, $\overline{\chi}(t) \rightarrow \chi(t)$.
- The equation of motion for a massive free scalar field is

$$\ddot{\chi} + m^2 \chi = 0$$

• The solution with $\chi(0) = \chi_0$ and $\dot{\chi}(0) = 0$ is



$$\chi(t) = \chi_0 \cos(mt)$$
 Oscillations with the period $T = 2\pi/m$.

$$P = 0$$

Obata, Fujita, Michimura (2018); Fedderke, Graham, Rajendran (2019)

CDM-induced parity violation in EM waves

"Time-domain cosmology"

The Chern-Simons interaction between photons and CDM gives

$$\ddot{A}_{\pm} + \left(k^2 \mp rac{k lpha \dot{\chi}}{f}
ight) A_{\pm} = 0$$
 $\ddot{A}_{\pm} + \left[k^2 \pm rac{k lpha \chi_0 m}{f} rac{\sin(mt)}{f}
ight] A_{\pm} = 0$ This is a human for parity violation in EM waves $\simeq 1~{
m day} \left(rac{5 imes 10^{-20}~{
m eV}}{m}
ight)$