

Lattice studies of multiply charmed baryons

Gunnar Bali

(Regensburg)



QWG13, IHEP Beijing, 23.4.2013



Outline

- Motivation
- Simulation parameters
- Multiply charmed baryons
- Summary

What I will (not) cover

I concentrate on recent results of

- SFB-TR55 (Regensburg-Wuppertal): [P Pérez-Rubio, 1302.5774](#) and [preliminary](#) $n_f = 2 + 1$ 2-HEX and SLiNC
2-HEX Charmonium: → [Craig McNeile](#), Thursday 10:20
- Hadron Spectrum Collaboration: [M Padmanath et al, preliminary](#), $n_f = 2 + 1$ anisotropic clover-Wilson
Charmonia: → [Christopher Thomas](#), earlier today.
- Recent results not covered in detail:
 - [L Liu et al 0909.3294](#) clover-Wilson and DW on $n_f = 2 + 1$ asqtad
 - [C Alexandrou et al \[ETMC\] 1205.6856](#) $n_f = 2$ Twisted Mass, 2 a 's
 - [R Briceño et al 1207.3536](#) clover on $n_f = 2 + 1 + 1$ HISQ, 3 a 's
 - [S Basak et al \[ILGTI\] 1211.6277](#) overlap on $n_f = 2 + 1 + 1$ HISQ
 - [Y Namekawa et al \[PACS-CS\] 1301.4743](#) $n_f = 2 + 1$ NP clover-Wilson

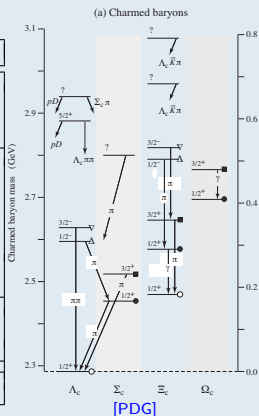
Singly charmed baryons

Bar.	M(MeV)	(qqq)	I(J ^P)	St.	Bar.	M(MeV)	(qqq)	I(J ^P)	St.
Λ_c^+	2286		$0(1/2^+)$	****	Ξ'^+	2575	(usc)	$1/2(1/2^+)$	***
	2595		$0(1/2^-)$	***	Ξ'^0	2578	(dsc)	$1/2(1/2^+)$	***
	2625	(udc)	$0(3/2^+)$	***		2645		$1/2(3/2^+)$	***
	2880		$0(5/2^+)$	***		2790	(usc),	$0(?^?)$	***
	2940		$0(?^?)$	***	Ξ_c	2815		$1/2(3/2^-)$	***
Σ_c	2455	(uuc),	$1(1/2^+)$	****		2980	(dsc)	$1(?^?)$	***
	2520	(udc),	$1(3/2^+)$	***		3080		$1/2(?^?)$	***
	2800	(ddc)	$1(?^?)$	***					
Ξ_c^+	2468	(usc)	$1/2(1/2^+)$	***	Ω_c	2695	(ssc)	$0(1/2^+)$	***
						2770		$0(3/2^+)$	***
Ξ_c^0	2470	(dsc)	$1/2(1/2^+)$	***	Ξ_{cc}^+	3519	(dcc)	$?(?^?)$	*

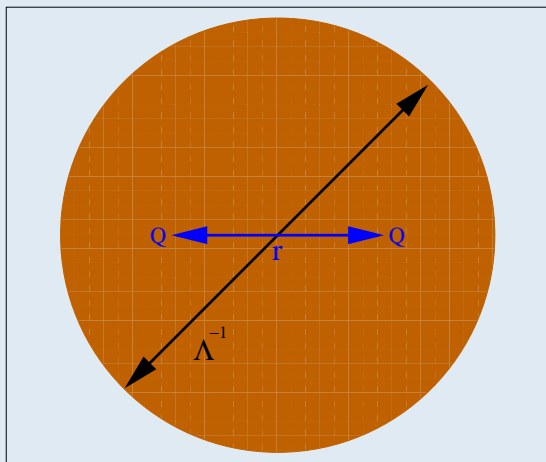
$$m_{\Sigma_c^{*+}} - m_{\Sigma_c^+} = (64.6 \pm 2.3) \text{ MeV}$$

$$m_{\Xi_c'^{*+}} - m_{\Xi_c'^+} = (70.9 \pm 3.4) \text{ MeV}$$

$$m_{\Omega_c^*} - m_{\Omega_c^+} = (70.8 \pm 1.5) \text{ MeV}$$

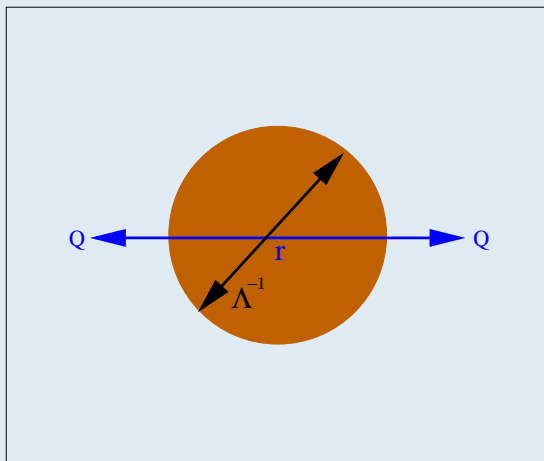


Doubly charmed baryons



HQET picture for $r \ll \Lambda^{-1}$.

Doubly charmed baryons

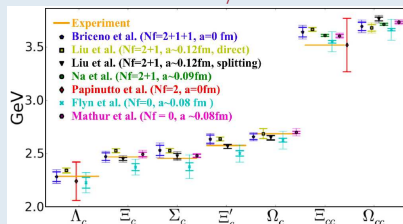


$r \gg \Lambda$: NRQCD/Charmonium-alike?

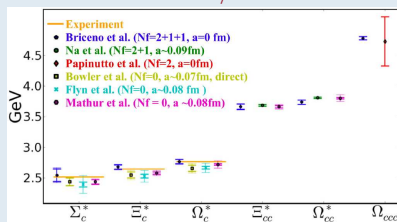
Previous lattice calculations

[R A Briceño, D Bolten, H-W Lin 11]

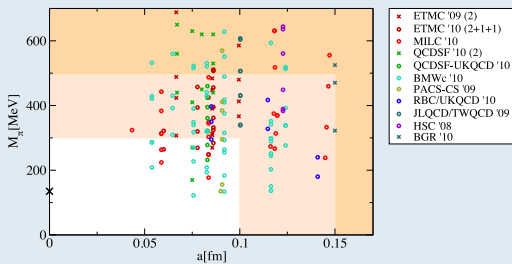
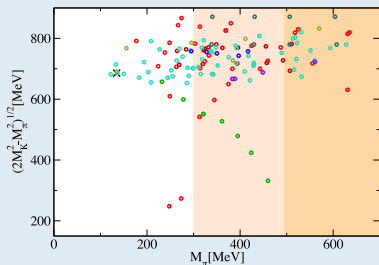
$$J^P = 1/2^+$$



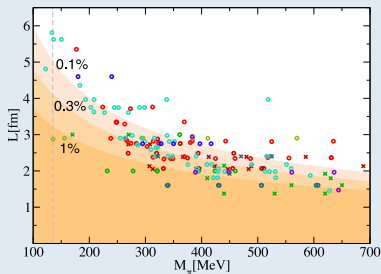
$$J^P = 3/2^+$$



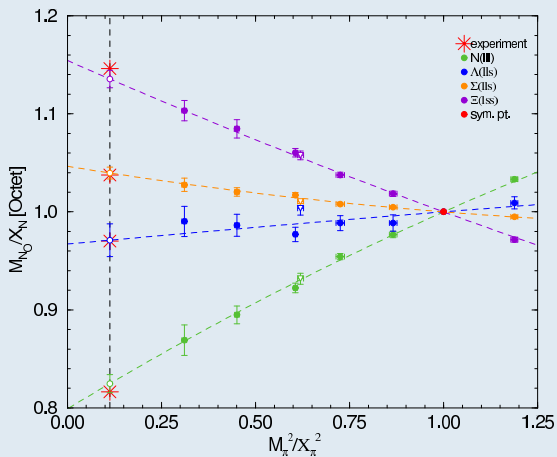
Landscape of current lattice simulations



Figures taken from [C Hoelbling 1102.0410](#)

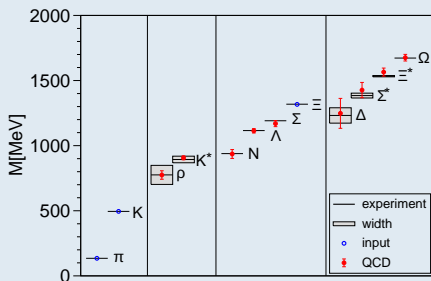
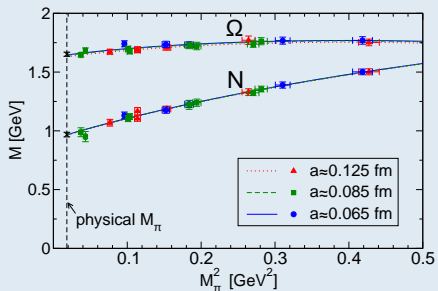


QCDSF: $m_u + m_d + m_s \propto m_\pi^2 + 2m_K^2 \approx \text{const.} = \text{physical}$



QCDSF + UKQCD: W Bietenholz et al. 10

BMW-c: chiral extrapolation and the continuum limit



S Dürer et al. 09/10

Analyzed configurations, $N_f = 2 + 1$

- BMW-c: 2-HEX action**

β	Volume	a (fm)	M_π (MeV)	no.	β	Volume	a (fm)	M_π (MeV)	no.
3.8	$32^3 \times 64$	0.054	205	454	3.7	$32^3 \times 64$	0.065	282	390
3.8	$48^3 \times 64$	0.054	193	317	3.7	$48^3 \times 64$	0.065	282	248
3.8	$64^3 \times 144$	0.054	117	220	3.7	$64^3 \times 64$	0.065	205	187
3.5	$24^3 \times 64$	0.092	199	415	3.7	$48^3 \times 64$	0.065	261	250
3.5	$32^3 \times 64$	0.092	109	260	3.61	$32^3 \times 48$	0.077	391	390
3.5	$32^3 \times 64$	0.092	167	256	3.61	$64^3 \times 72$	0.077	94	120
3.5	$24^3 \times 64$	0.092	340	371	3.61	$48^3 \times 48$	0.077	97	265
3.5	$32^3 \times 64$	0.092	213	254	3.61	$32^3 \times 64$	0.077	90	517
3.5	$64^3 \times 64$	0.092	165	133	3.61	$32^3 \times 64$	0.077	146	340

Analyzed configurations, $N_f = 2 + 1$

- QCDSF: SLiNC action**

* $m_u + m_s + m_d$ kept constant ($2M_\pi^2 + M_K^2 \approx \text{const.}$)

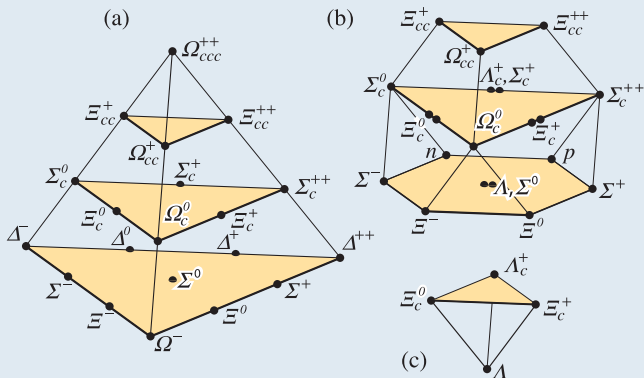
β	Volume	a (fm)	M_π (MeV)	M_K (MeV)
5.50	$24^3 \times 48$	0.0795(3)	442	442
5.50	$24^3 \times 48$	0.0795(3)	412	458
5.50	$24^3 \times 48$	0.0795(3)	375	471
5.50	$24^3 \times 48$	0.0795(3)	348	484
5.50	$32^3 \times 64$	0.0795(3)	432	432
5.50	$32^3 \times 64$	0.0795(3)	332	474
5.50	$32^3 \times 64$	0.0795(3)	283	490
5.50	$48^3 \times 96$	0.0795(3)	220	500

- HadSpect (Padmanath et al): anisotropic clover-Wilson**

$a_s \approx 0.12$ fm, $a_t \approx 0.035$ fm,

Volumes: $16^3 \times 128$, $24^3 \times 128$, $M_\pi \approx 390$ MeV.

- SU(4) representations



- Flavour symmetry is not respected but
- simplest way to see which baryons should exist.

- SU(4): $4 \otimes 4 \otimes 4 = 20 \oplus 20 \oplus 20 \oplus \bar{4}$

$$\square \otimes \square \otimes \square = \begin{array}{|c|c|c|} \hline \square & \square & \square \\ \hline \end{array} \oplus \begin{array}{|c|c|} \hline \square & \square \\ \hline \end{array} \oplus \begin{array}{|c|c|} \hline \square & \square \\ \hline \end{array} \oplus \begin{array}{|c|} \hline \square \\ \hline \end{array}$$

Interpolating fields: SU(4) reps

- SU(4) 20-PLET CONTAINING SU(3) OCTETS 

- * **N-alike:** $\Omega_{cc}^+, \Xi_{cc}^+, \Xi_{cc}^{++}$.

$$\mathcal{O}_\gamma^P(x) = \epsilon^{abc} [q_1^a(x)^T (C\gamma_5) q_2^b(x)] q_{2\gamma}^c(x).$$

- SU(4) 20-PLET CONTAINING SU(3) DECUPLET 

- * Δ^{++} -**alike:** Ω_{ccc}^{++} .

$$\mathcal{O}_\gamma^{\Delta^{++}} = \epsilon^{abc} (q_1^{aT} (C\gamma_\mu) q_1^b) q_{1\gamma}^c$$

- * Σ^{*-} -**alike:** $\Xi_{cc}^{*+}, \Xi_{cc}^{*++}, \Omega_{cc}^{*+}$.

$$\mathcal{O}_\gamma^{\Sigma^{*-}} = \epsilon^{abc} \{ 2 (q_1^{aT} (C\gamma_\mu) q_2^b) q_{2\gamma}^c + (q_2^{aT} (C\gamma_\mu) q_2^b) q_{1\gamma}^c \}$$

- CORRELATORS

$$C(t, \mathbf{p} = 0) = T_{\bar{\gamma}\gamma} \sum_{\mathbf{x}} \langle \mathcal{O}_\gamma(x) \bar{\mathcal{O}}_{\bar{\gamma}}(0) \rangle,$$

where T_γ is a parity and spin projection matrix.

Interpolating fields: HQET

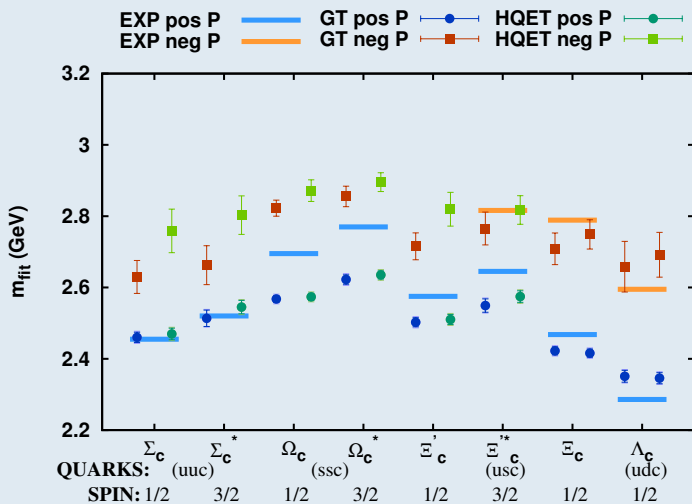
cc diquark: $\mathbf{3} \times \mathbf{3} = \bar{\mathbf{3}} + \mathbf{6}$
 diquark spin s_d , parity π_d .

$$J^P = \frac{1}{2}^+ \quad J^P = \frac{3}{2}^+$$

(S)	(I)	$s_d^{\pi_d}$	(qq)q	\mathcal{O}	Name	Name
(0)	($\frac{1}{2}$)	(1) ⁺	(cc)u	$\mathcal{O}_\mu = \epsilon_{abc}(c^{aT} C \gamma_\mu c^b)u^c$	Ξ_{cc}	Ξ_{cc}^*
(-1)	(0)	(1) ⁺	(cc)s	$\mathcal{O}_\mu = \epsilon_{abc}(c^{aT} C \gamma_\mu c^b)s^c$	Ω_{cc}	Ω_{cc}^*

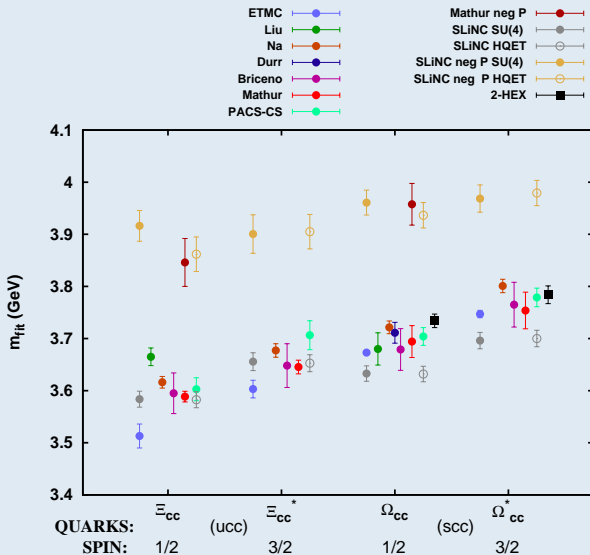
- All \mathcal{O}_x 's require projections $T_{\bar{\gamma}\gamma}$.
- Correlators, $C_{\mu\nu}(t)$ from $\mathcal{O}_\mu^{(t)}$ need projections into the desired $J = 1/2, 3/2$.
 Projectors $P_{\mu\nu}^{3/2}, P_{\mu\nu}^{1/2}$ used.

Singly charmed baryons: SLiNC preliminary



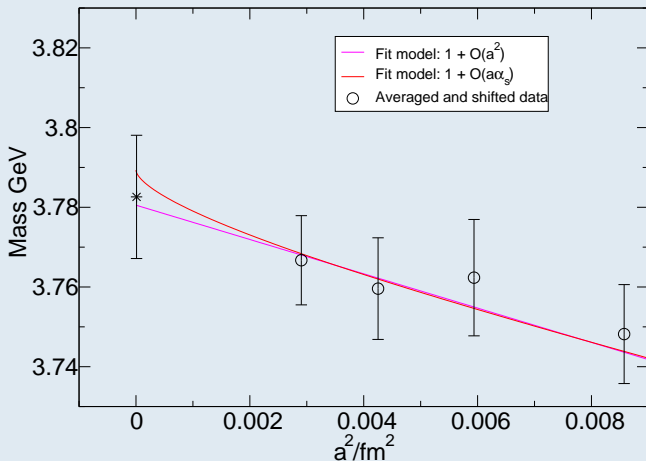
$m_s/m_u \approx 2.9 \Rightarrow$ light quarks too heavy.

cc baryons: other groups, SLiNC + 2-HEX preliminary

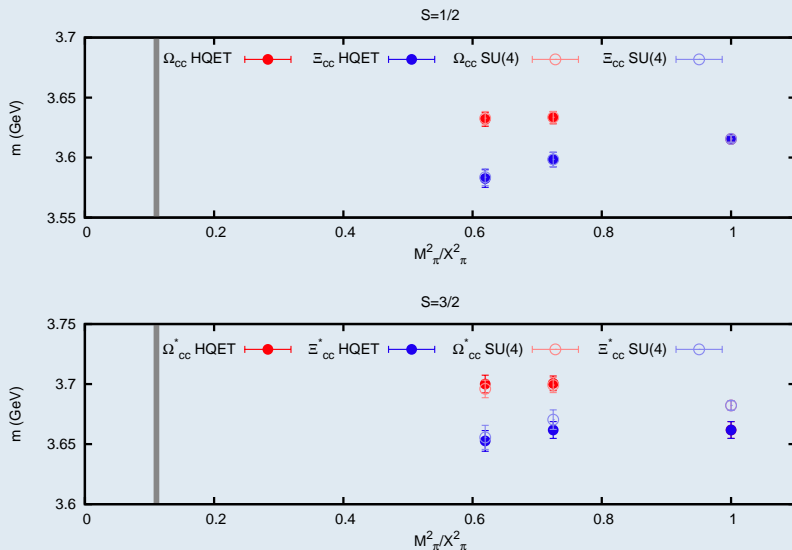


2-HEX continuum limit extrapolation: preliminary!

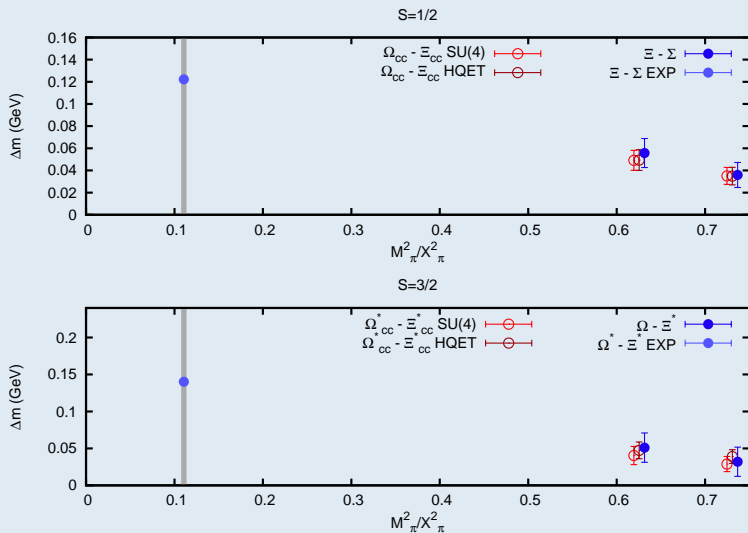
Spin 3/2 ccs baryon



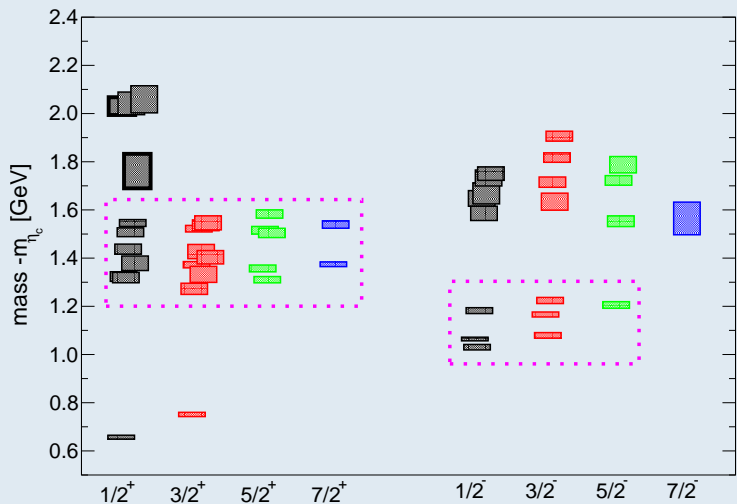
Mass dependence of ccs - ccu splittings: SLiNC preliminary



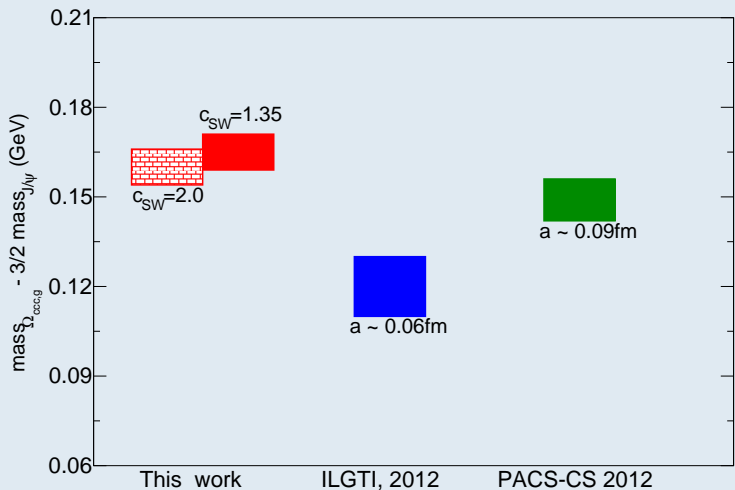
ccs-ccu splittings II: SLiNC preliminary



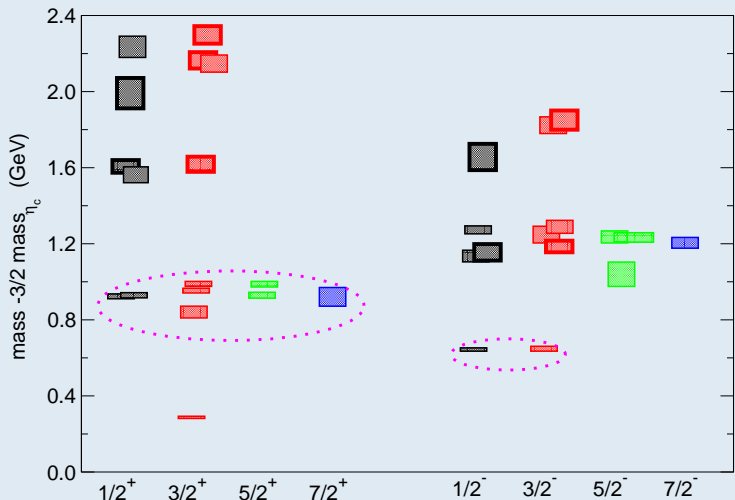
ccu baryons: Padmanath et al, preliminary



ccc ground state: Padmanath et al, preliminary



Triply charmed baryons: Padmanath et al, preliminary



Summary and outlook

- Hidden and open charm states are narrower and experimentally cleaner than many light quark resonances. Theoretically, the heavy quark limit provides guidance.
- Particularly interesting are the baryonic analogues of charmonia. Unfortunately, these states are still awaiting to be discovered.
- In the doubly charmed baryon sector we aim to establish if the spectrum resembles that of D , D_s “heavy mesons”: $(cc)u$, $(cc)s$ (HQET picture) — or that of $c(cu)$, $c(cs)$ “charmonia”.
- Positive parity ccs ground states (Ω_{cc} and Ω_{cc}^*) have been interpolated to the physical point and the continuum limit taken. Lattice studies of other states are of a more qualitative nature.
- The tip of the pyramid — even less discoverable: Ω_{ccc}^{++} .