

Bottomonium Spectroscopy at Belle



April 22- 26, 2013, IHEP, Beijing

Saurabh Sandilya

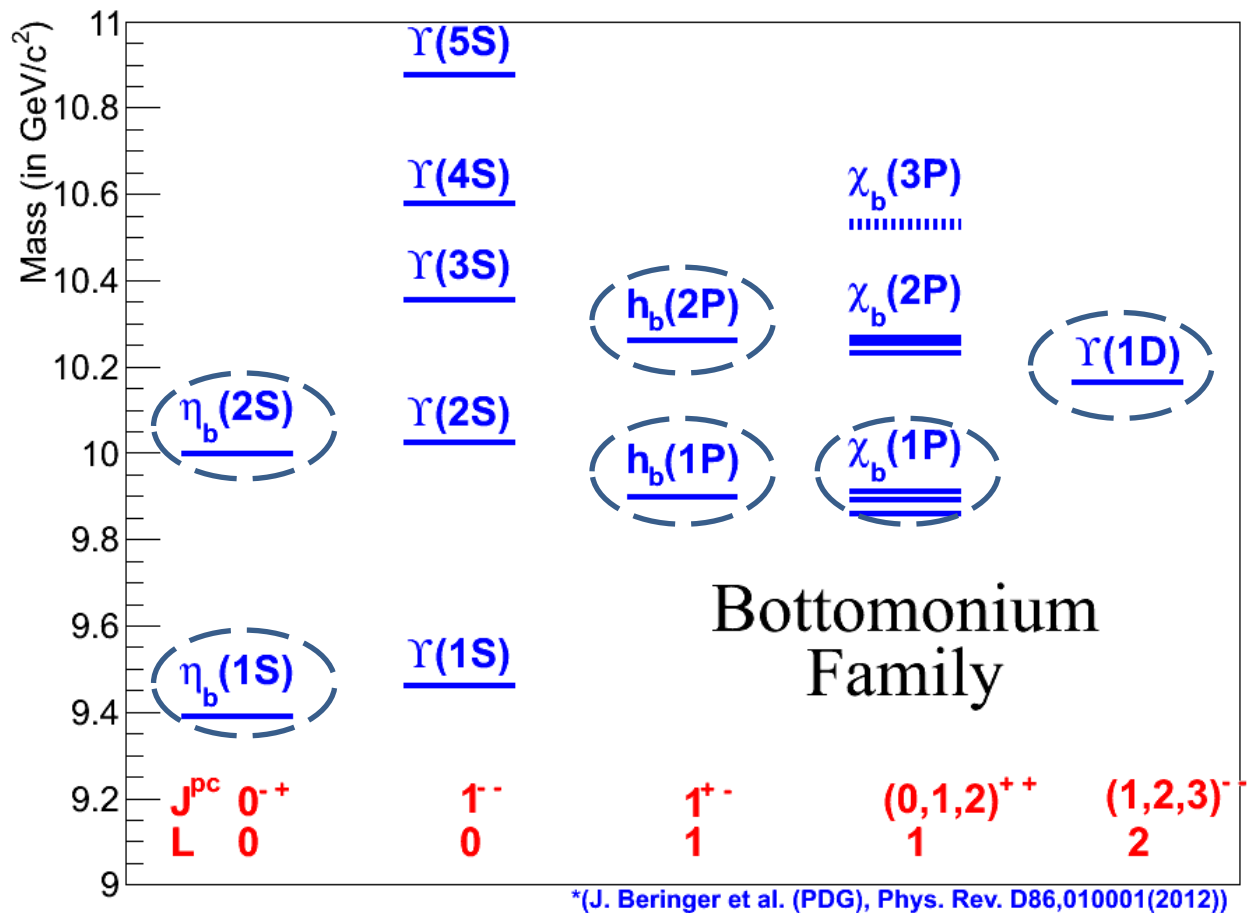
TIFR, Mumbai


(On behalf of the Belle Collaboration)

Outline

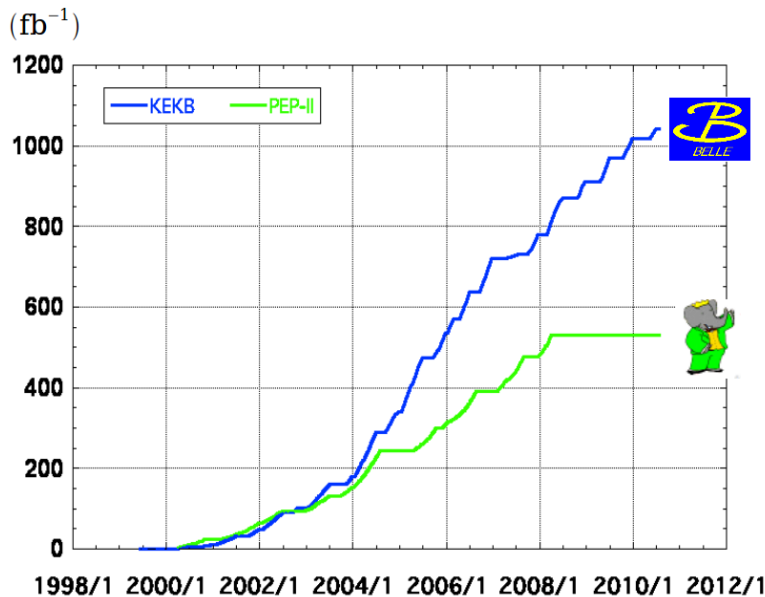
- $h_b(nP) \rightarrow \eta_b(mS)\gamma$
- Radiative $\Upsilon(2S)$ decay 
- $\Upsilon(1D)$ mass 
- R_b Scan
- Summary

Introduction



- My talk is a tour of some of the bottomonium states, indicated by  studied using the unique data set of Belle

Belle datasets



> 1 ab⁻¹
On resonance:
 Y(5S): 121 fb⁻¹
 Y(4S): 711 fb⁻¹
 Y(3S): 3 fb⁻¹
 Y(2S): 25 fb⁻¹
 Y(1S): 6 fb⁻¹
Off reson./scan:
 ~ 100 fb⁻¹

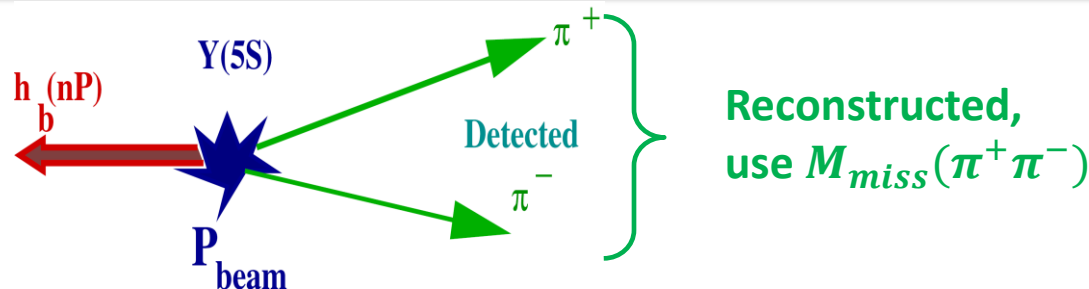
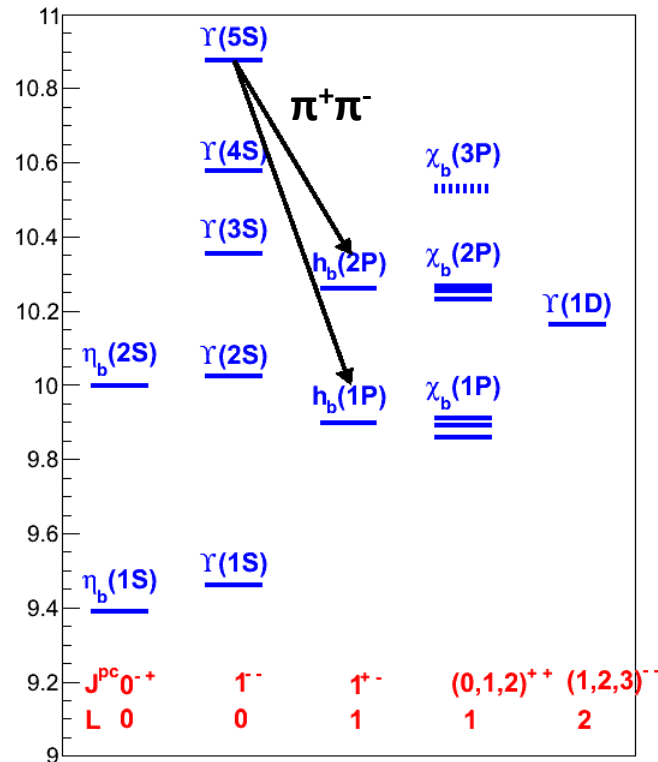
~ 550 fb⁻¹
On resonance:
 Y(4S): 433 fb⁻¹
 Y(3S): 30 fb⁻¹
 Y(2S): 14 fb⁻¹
Off resonance:
 ~ 54 fb⁻¹

Resonance	Data Collected @ Belle (June 30, 2010)
Y(1S)	6 fb ⁻¹ Largest
Y(2S)	24.7 fb ⁻¹ Largest
Y(3S)	3 fb ⁻¹
Y(4S)	711 fb ⁻¹ Largest
Y(5S)	121 fb ⁻¹ Largest

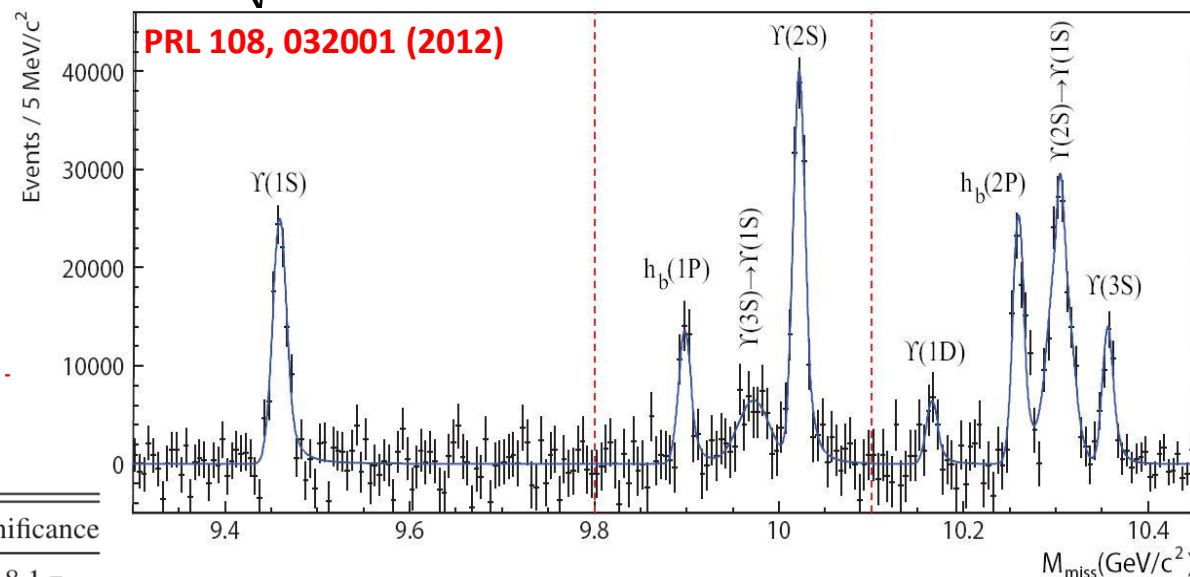
- Belle has collected most of its data at the Y(4S) resonance
 - To study various aspects of B-meson decays
 - Most celebrated result has been “CP violation in B-meson decays”
- e^+e^- colliders produce a particularly clean environment to study properties of the Υ states.
- The entire collision energy of the initial e^+e^- system turns into the Υ rest mass.

Observation of $h_b(nP)$

Looking for $h_b(nP)$
 $Y(5S) \rightarrow h_b \pi^+ \pi^-$ reconstruction



$$M(h_b) = \sqrt{(E_{\text{CM}} - E_{\pi^+\pi^-}^*)^2 - E_{\pi^+\pi^-}^{*2}} \equiv M_{\text{miss}}(\pi^+\pi^-)$$



$$\Delta M_{HF}(1P) = +0.8 \pm 1.1 \text{ MeV}$$

$$\Delta M_{HF}(2P) = +0.5 \pm 1.2 \text{ MeV}$$

consistent with zero, as expected.

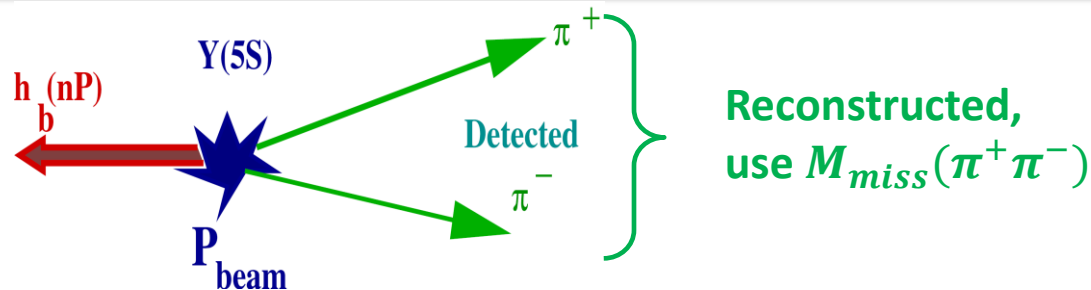
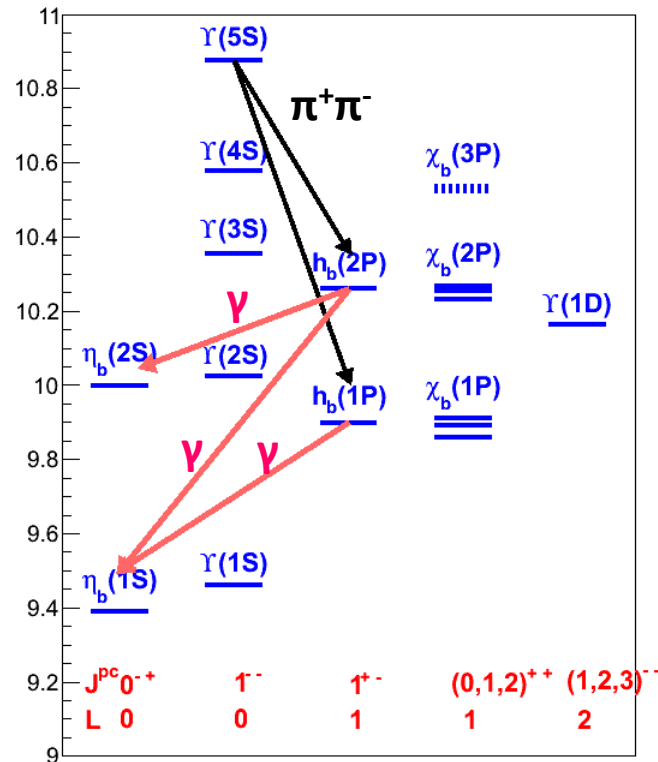
large $h_b(1,2P)$ production rate $\Rightarrow Z_b^+$ Discovery

See R. Mizuk's talk, on 26th April

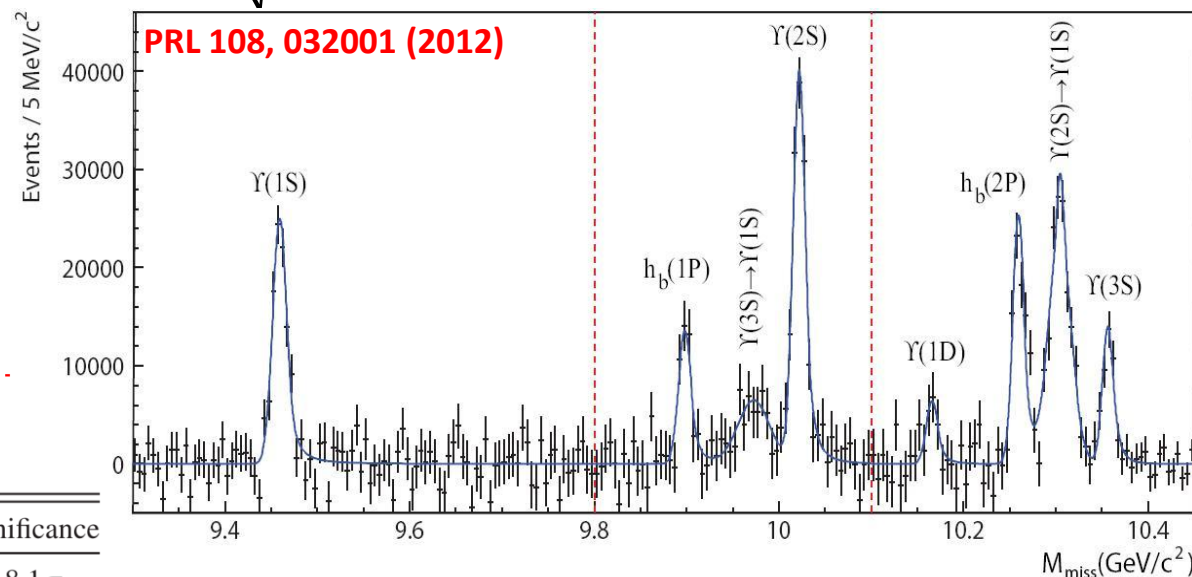
	Yield, 10^3	Mass, MeV/c^2	Significance
$Y(1S)$	$104.9 \pm 5.8 \pm 3.0$	$9459.4 \pm 0.5 \pm 1.0$	18.1σ
$h_b(1P)$	$50.0 \pm 7.8^{+4.5}_{-9.1}$	$9898.2^{+1.1+1.0}_{-1.0-1.1}$	6.1σ
$3S \rightarrow 1S$	55 ± 19	9973.01	2.9σ
$Y(2S)$	$143.7 \pm 8.7 \pm 6.8$	$10022.2 \pm 0.4 \pm 1.0$	17.1σ
$Y(1D)$	22.4 ± 7.8	10166.1 ± 2.6	2.4σ
$h_b(2P)$	$83.9 \pm 6.8^{+23.0}_{-10.0}$	$10259.8 \pm 0.6^{+1.4}_{-1.0}$	12.3σ
$2S \rightarrow 1S$	$151.3 \pm 9.7^{+9.0}_{-20.0}$	$10304.6 \pm 0.6 \pm 1.0$	15.7σ
$Y(3S)$	$45.5 \pm 5.2 \pm 5.1$	$10356.7 \pm 0.9 \pm 1.1$	8.5σ

Observation of $h_b(nP)$

Looking for $h_b(nP)$
 $Y(5S) \rightarrow h_b \pi^+ \pi^-$ reconstruction



$$M(h_b) = \sqrt{(E_{CM} - E_{\pi^+ \pi^-}^*)^2 - E_{\pi^+ \pi^-}^{*2}} \equiv M_{miss}(\pi^+ \pi^-)$$

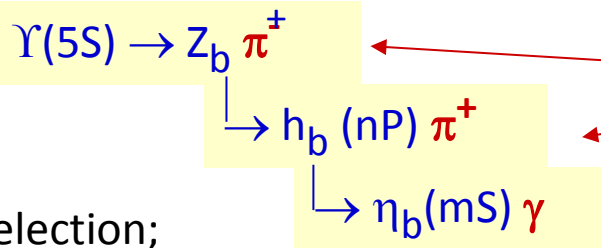


High yield of $h_b(nP)$ opens new perspective to study $\eta_b(mS)$!!

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Observation of $h_b(nP) \rightarrow \eta_b(mS)\gamma$

Decay chain



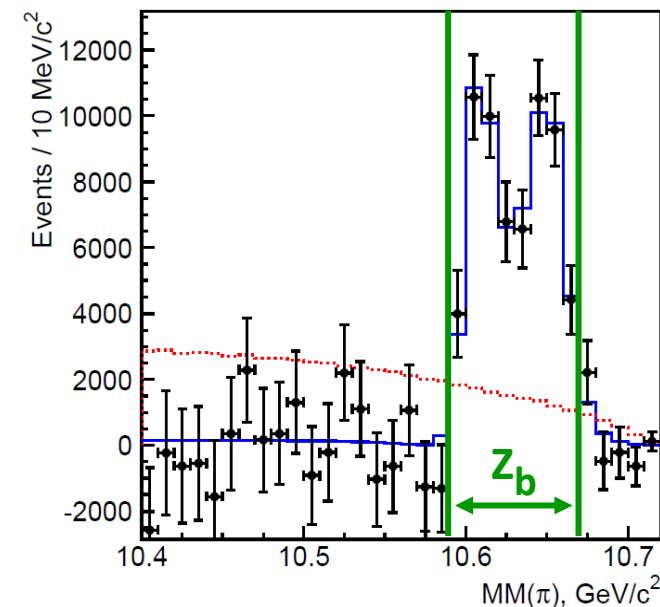
reconstruct

Use missing mass to identify signals

Hadronic event selection;
continuum suppression using event shape;
 π^0 veto.

Require intermediate Z_b : $10.59 < MM(\pi) < 10.67 \text{ GeV}$

$$\Delta M_{\text{miss}}(\pi^+\pi^-\gamma) \equiv M_{\text{miss}}(\pi^+\pi^-\gamma) - M_{\text{miss}}(\pi^+\pi^-) + M[h_b]$$

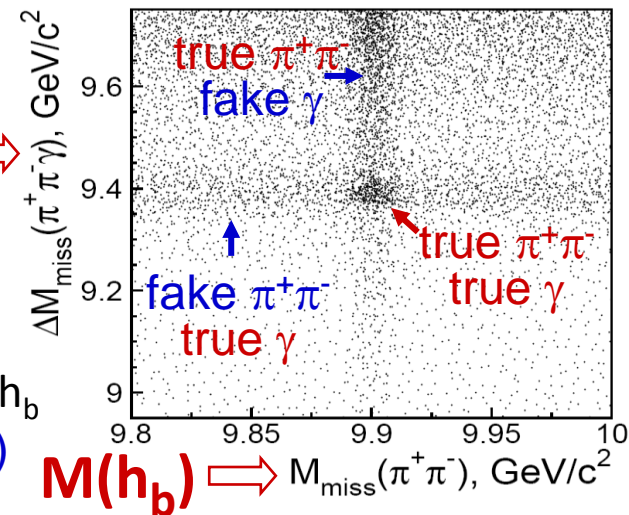


Approach:

$$M(\eta_b) \Rightarrow$$

Signal is a cluster in the 2D plane and we determine the h_b yield in bins of $\Delta M_{\text{miss}}(\pi^+\pi^-\gamma)$

- rectangular bg bands
- no correlation
MC simulation

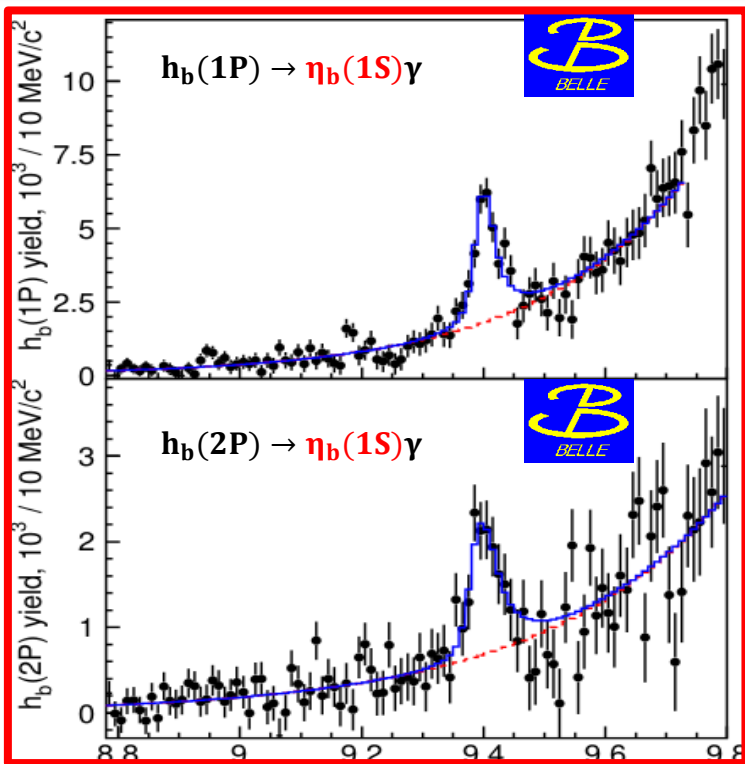


133.4 fb⁻¹ Data used at $\Upsilon(5S)$ resonance

fit $M_{\text{miss}}(\pi^+\pi^-)$ spectra
in $\Delta M_{\text{miss}}(\pi^+\pi^-\gamma)$ bins

Re-Discovery of $\eta_b(1S)$ at Belle

PRL 109,232002 (2012)



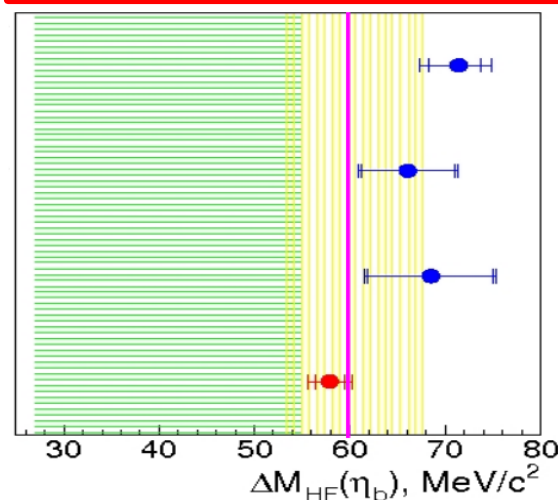
$$m_{\eta_b(1S)} = 9402.4 \pm 1.5 \pm 1.8 \text{ MeV}/c^2$$

First measurement of $\Gamma = 10.8_{-3.7}^{+4.0} {}_{-2.0}^{+4.5} \text{ MeV}/c^2$

$$\text{B.F. } [h_b(1P) \rightarrow \eta_b(1S)\gamma] = (49.2 \pm 5.7 {}_{-3.3}^{+5.6})\%$$

$$\text{B.F. } [h_b(2P) \rightarrow \eta_b(1S)\gamma] = (22.3 \pm 5.7 {}_{-3.3}^{+5.6})\%$$

More precise than PDG 2012 (avg)[9391.0 \pm 2.8 MeV],
decreases tension with theory



BaBar $\Upsilon(3S)$

BaBar $\Upsilon(2S)$

PNRQCD@NLL

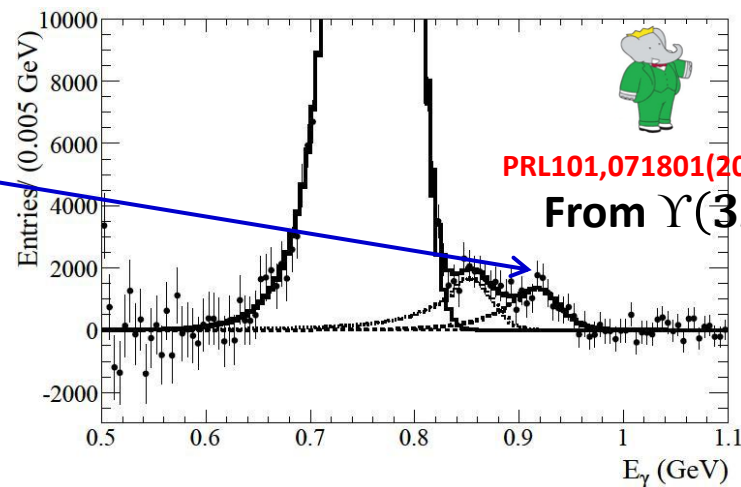
CLEO PRL92,242001(2004)

Lattice QCD

Belle PRD82,114502(2010)

Godfrey-Isgur

PRD32,189(1985)



PRL101,071801(2008)
From $\Upsilon(3S)$

First evidence of $\eta_b(2S)$ at Belle

PRL 109,232002 (2012)

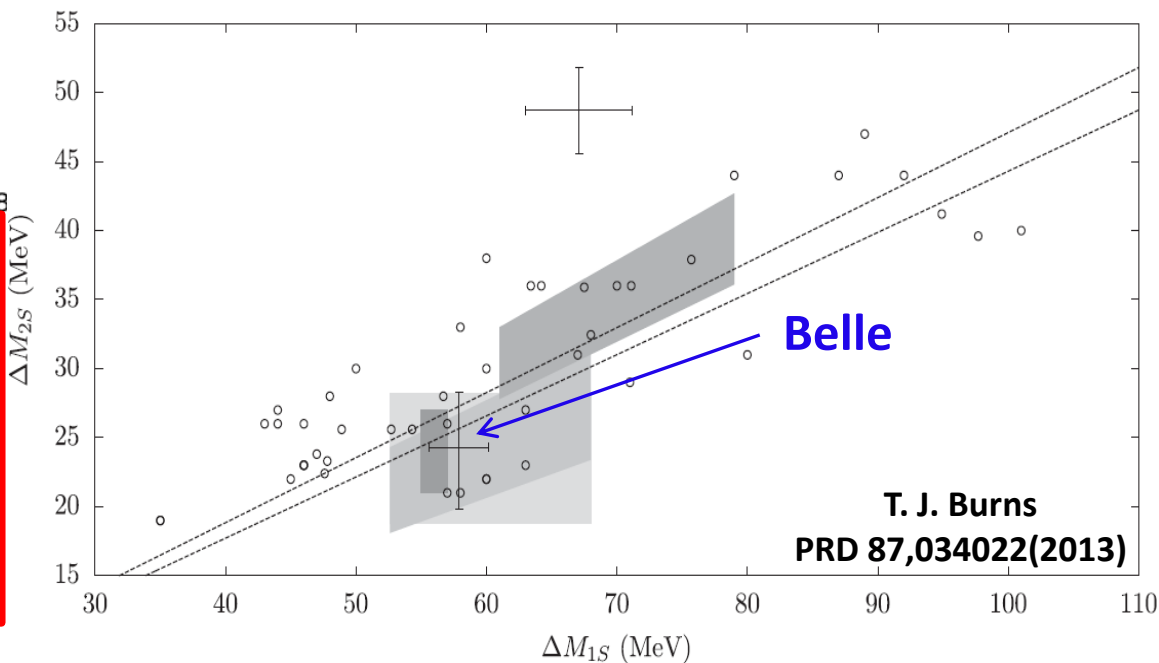
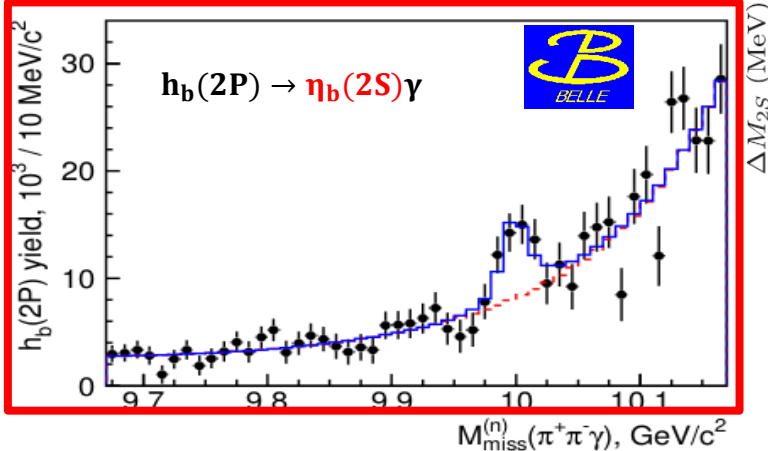
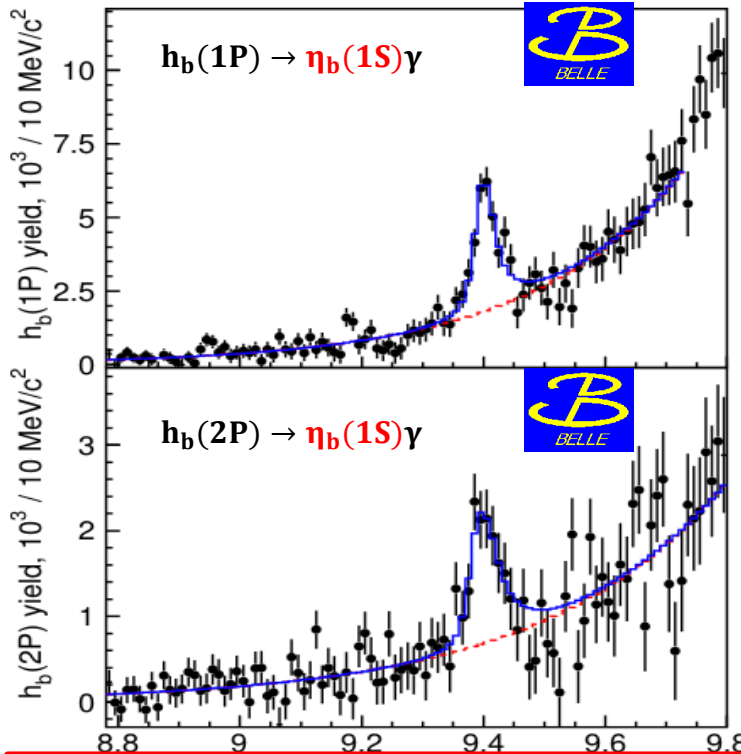
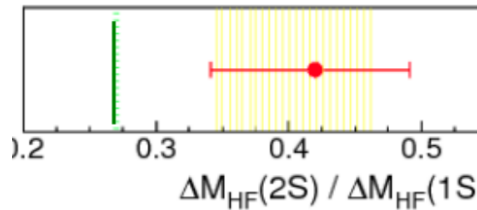
$$m_{\eta_b(2S)} = 9999.0 \pm 3.5^{+2.8}_{-1.9} \text{ MeV}/c^2$$

with significance (including sys.) at 4.2σ

$$\text{B.F. } [h_b(2P) \rightarrow \eta_b(2S)\gamma] = (47.5 \pm 10.5^{+6.8}_{-7.7})\%$$

$$\Delta M_{\text{HF}}(2S) = 24.3^{+4.0}_{-4.5} \text{ MeV}/c^2$$

Belle



First evidence of $\eta_b(2S)$ at Belle

PRL 109,232002 (2012)

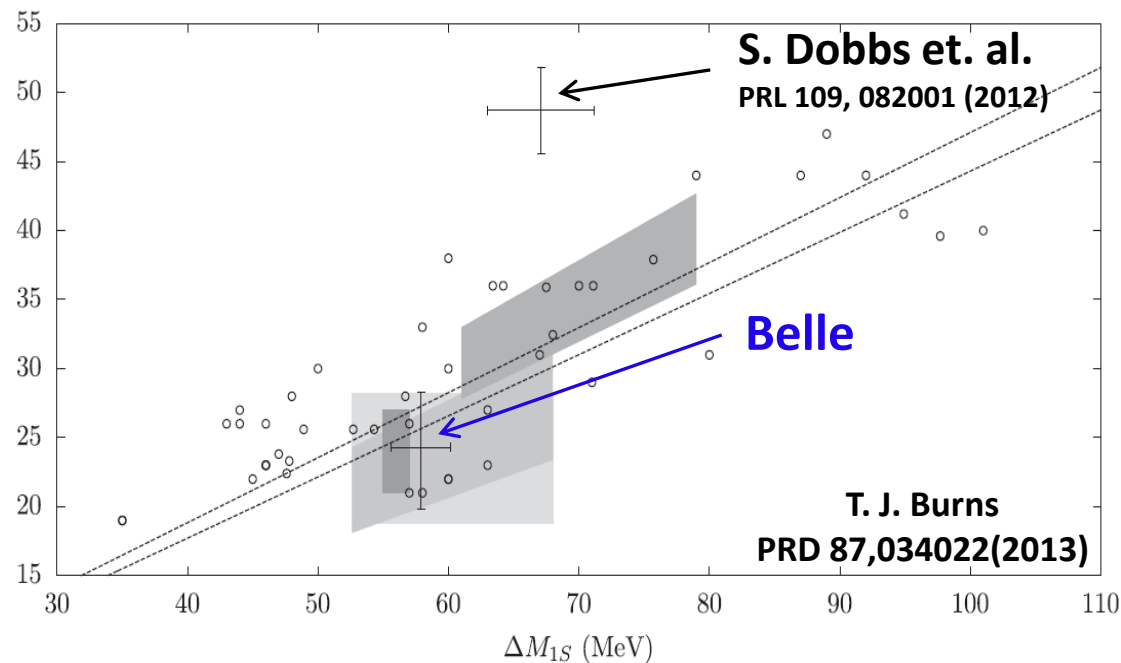
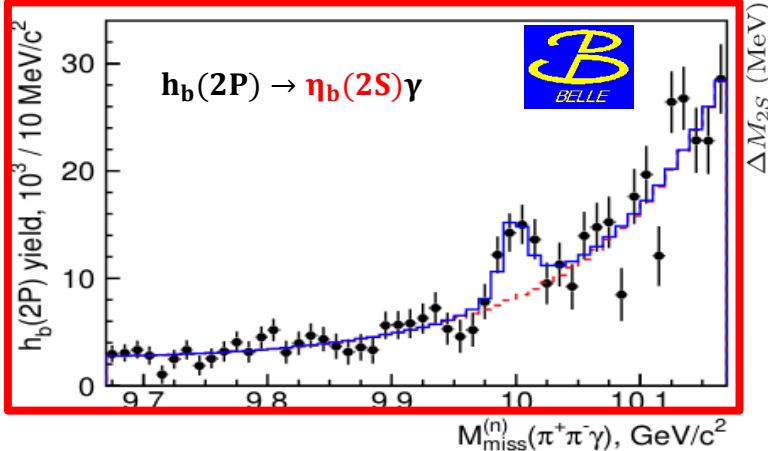
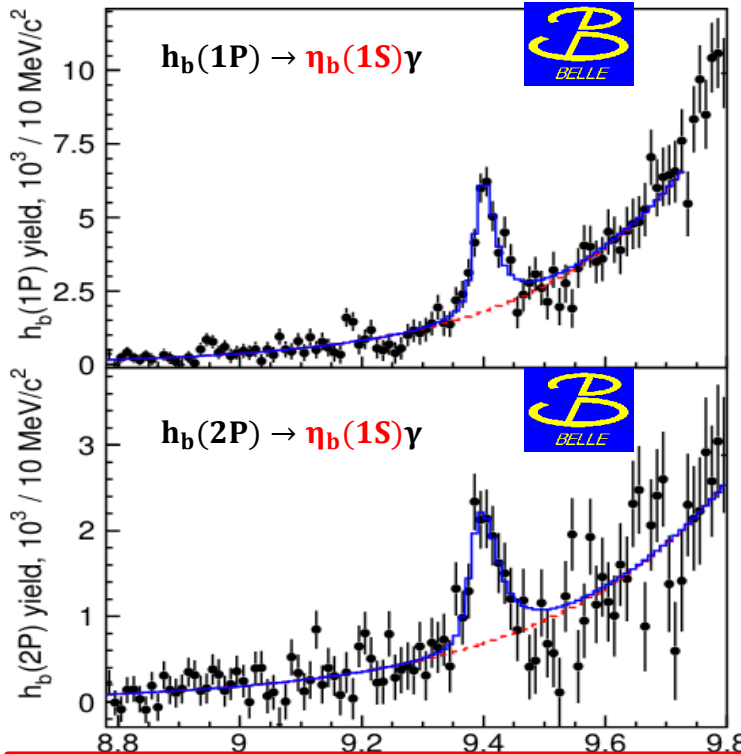
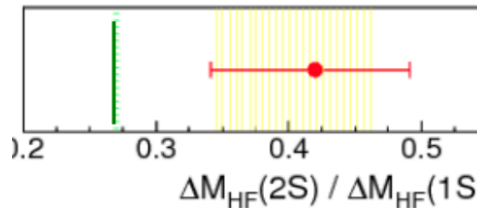
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Belle



$\eta_b(2S)$ claim based on CLEO-c data

Observation of the $\eta_b(2S)$ Meson in $Y(2S) \rightarrow \gamma\eta_b(2S)$, $\eta_b(2S) \rightarrow$ Hadrons and Confirmation of the $\eta_b(1S)$ Meson

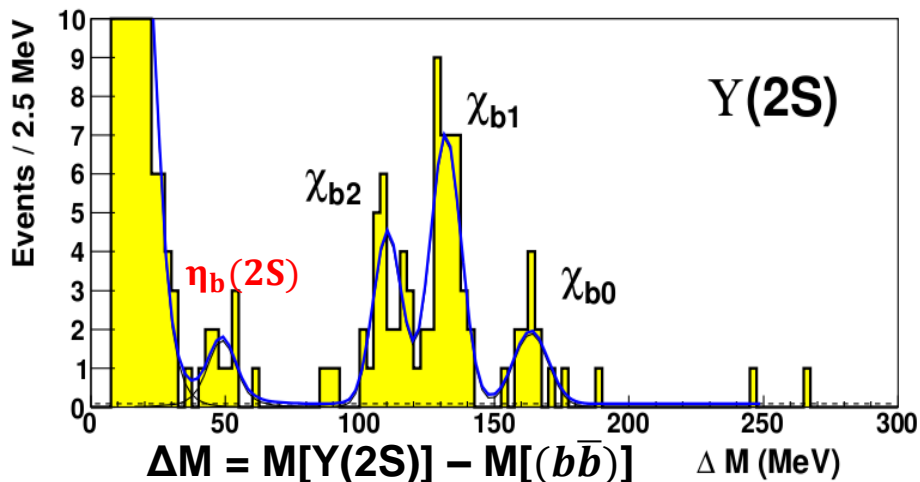
PRL 109, 082001 (2012)

S. Dobbs, Z. Metreveli, A. Tomaradze, T. Xiao, and Kamal K. Seth

Northwestern University, Evanston, Illinois 60208, USA

(Received 18 April 2012; published 24 August 2012)

The data for 9.3 million $Y(2S)$ and 20.9 million $Y(1S)$ taken with the CLEO III detector have been used to study the radiative population of states identified by their decay into 26 different exclusive hadronic final states. In the $Y(2S)$ decays, an enhancement is observed at a $\sim 5\sigma$ level at a mass of $9974.6 \pm 2.3(\text{stat}) \pm 2.1(\text{syst})$ MeV. It is attributed to $\eta_b(2S)$ and corresponds to the $Y(2S)$ hyperfine splitting of $48.7 \pm 2.3(\text{stat}) \pm 2.1(\text{syst})$ MeV. In the $Y(1S)$ decays, the identification of $\eta_b(1S)$ is confirmed at a $\sim 3\sigma$ level with $M[\eta_b(1S)]$ in agreement with its known value.



The measurement is carried out in 26 exclusive decays of the $\eta_b(2S)$ into charged hadrons.

$$\mathcal{B}_1 \times \mathcal{B}_2 \equiv \mathcal{B}_1[Y(nS) \rightarrow \gamma\eta_b(nS)] \times \sum_{i=1}^{26} \mathcal{B}_{2i}[\eta_b(nS) \rightarrow h_i]$$

$$\mathcal{B}_1 \times \mathcal{B}_2(\eta_b(2S)) = (46.2 \pm_{-14.2}^{+29.7} \pm 10.6) \times 10^{-6}$$

Reminder: Our $\Delta M_{\text{HF}}(2S) = 24.3_{-4.5}^{+4.0}$ MeV/c²

S. Dobbs's $\eta_b(2S)$ signal is not consistent with theory as well as our measurement

$\Upsilon(2S) \rightarrow \gamma (b\bar{b})$

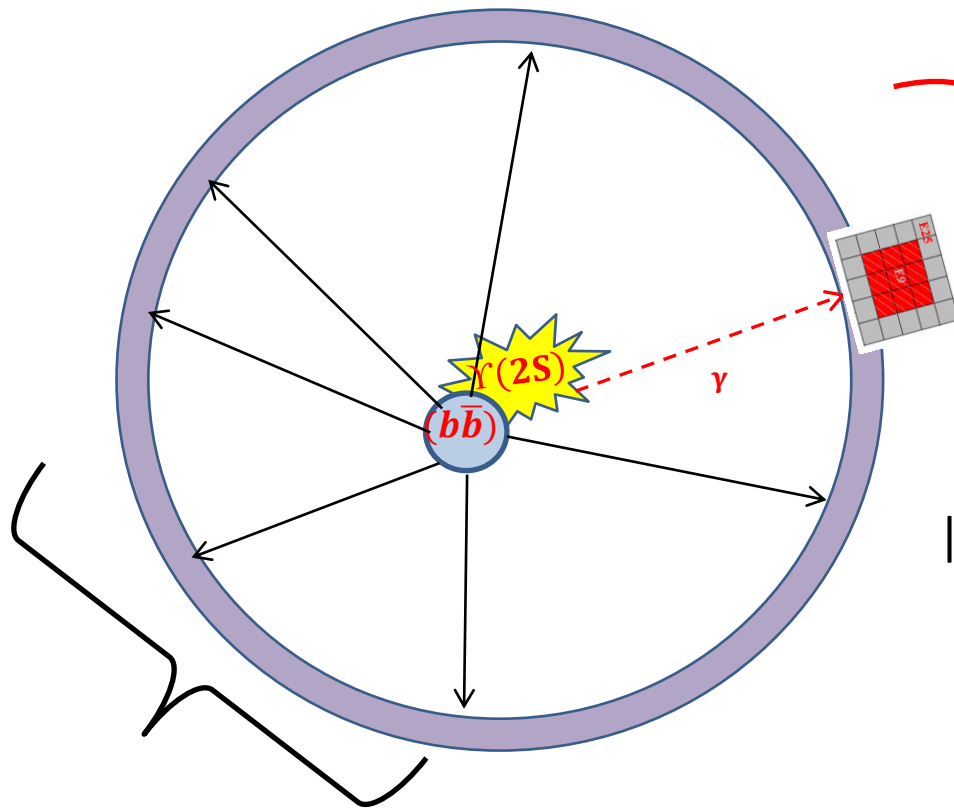
- Study is performed using the 25 fb^{-1} data ($157.8 \times 10^6 \Upsilon(2S)$ events) .
- ~17 times more data than CLEO-c's $\Upsilon(2S)$ sample**
- We study $\Upsilon(2S) \rightarrow \gamma (b\bar{b})$; where $(b\bar{b})$ decays hadronically (same 26 exclusive hadronic final states as mentioned in S. Dobbs et. al.)

$$\begin{aligned} \mathbf{x}_1 : & 2(\pi^+\pi^-), 3(\pi^+\pi^-), 4(\pi^+\pi^-), 5(\pi^+\pi^-), \\ & K^+K^-\pi^+\pi^-, K^+K^-2(\pi^+\pi^-), K^+K^-3(\pi^+\pi^-), \\ & K^+K^-4(\pi^+\pi^-), 2(K^+K^-), 2(K^+K^-)\pi^+\pi^-, \\ & 2(K^+K^-)2(\pi^+\pi^-), 2(K^+K^-)3(\pi^+\pi^-), p\bar{p}\pi^+\pi^-, \\ & p\bar{p}2(\pi^+\pi^-), p\bar{p}3(\pi^+\pi^-), p\bar{p}4(\pi^+\pi^-), p\bar{p}K^+K^-\pi^+\pi^-, \\ & p\bar{p}K^+K^-2(\pi^+\pi^-), p\bar{p}K^+K^-3(\pi^+\pi^-), K_S^0K^\pm\pi^\mp, \\ & K_S^0K^\pm\pi^\mp\pi^+\pi^-, K_S^0K^\pm\pi^\mp2(\pi^+\pi^-), K_S^0K^\pm\pi^\mp3(\pi^+\pi^-), \\ & 2K_S^0\pi^+\pi^-, 2K_S^02(\pi^+\pi^-), 2K_S^03(\pi^+\pi^-). \end{aligned}$$
- Following decay channels are good control samples

$$\Upsilon(2S) \rightarrow \gamma \chi_{bJ} \quad (J = 0, 1, 2)$$

and χ_{bJ} can decay to the hadronic modes (comprising charged pions, kaons, protons and K_S mesons)
- Off-resonance $\Upsilon(4S)$ data [89.5 fb^{-1} ~4 times larger than our $\Upsilon(2S)$ data] used for background shape study.

$$\Upsilon(2S) \rightarrow \gamma (b\bar{b})$$



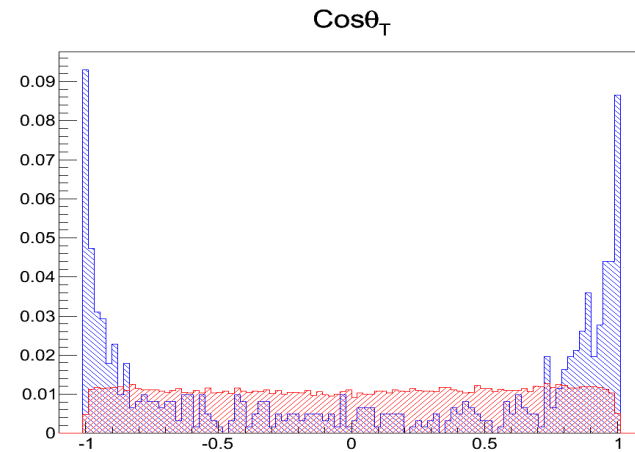
γ - selection

- Isolated cluster
- Energy of gamma > 22 MeV
- $E9/E25 > 0.85$
- Exclude endcaps

$|\cos\theta_T| < 0.8$: Continuum - Suppression

$(b\bar{b})$ reconstruction:

- Impact parameter cuts
- Number of charged tracks
- Particle identification (pion, kaon, proton)



Simple and Straightforward !!

$\Upsilon(2S) \rightarrow \gamma (b\bar{b})$

- More variables exploited to suppress backgrounds :
 - ΔE
 $E_{\Upsilon(2S)}^* - E_{CM}$
should peak around 0.
[$\Delta E > -0.04$ GeV & $\Delta E < 0.05$ GeV]
 - $P_{\Upsilon(2S)}^*$
momentum of the $\Upsilon(2S)$ candidate in the center-of-mass.
should peak around 0.
[$P_{\Upsilon(2S)}^* < 0.03$ GeV/c]
 - $\theta_{\gamma(bb)}$
Angle between γ candidate and $(b\bar{b})$ in the CM Frame.
should peak around 180° .
[$\theta_{\gamma(bb)} > 150^\circ$]
- Cut values obtained from optimization (assuming S. Dobbs et. al. B. F.)
- Multiple Candidates found at this stage is 8-10% .
- Energy-Momentum constrained kinematic fit (4C) is used to improve the resolution as well as for the best candidate selection.

$\Upsilon(2S) \rightarrow \gamma (b\bar{b})$

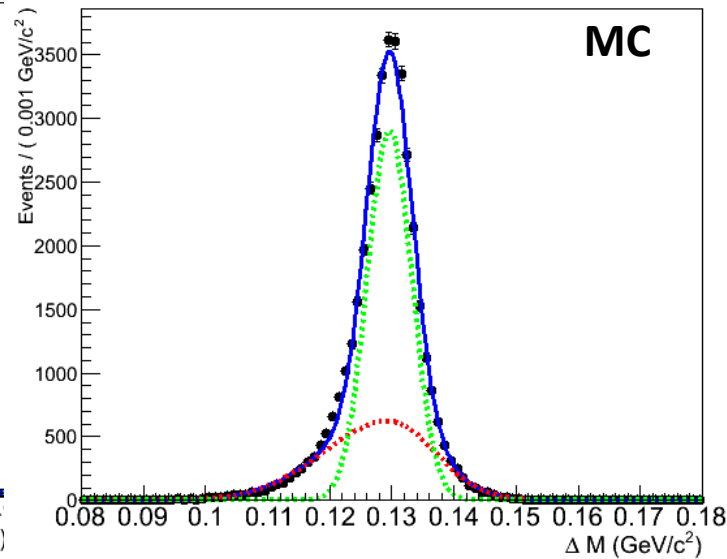
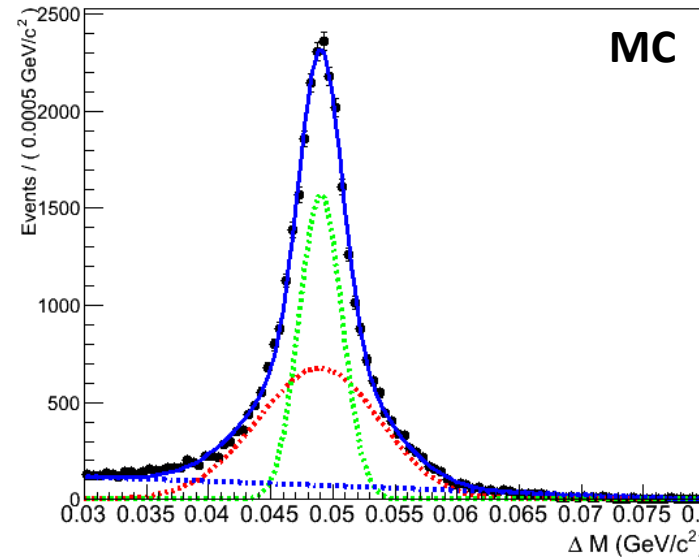
$\eta_b(2S)$: Signal MC

$\chi_{b1}(1P)$: Signal MC

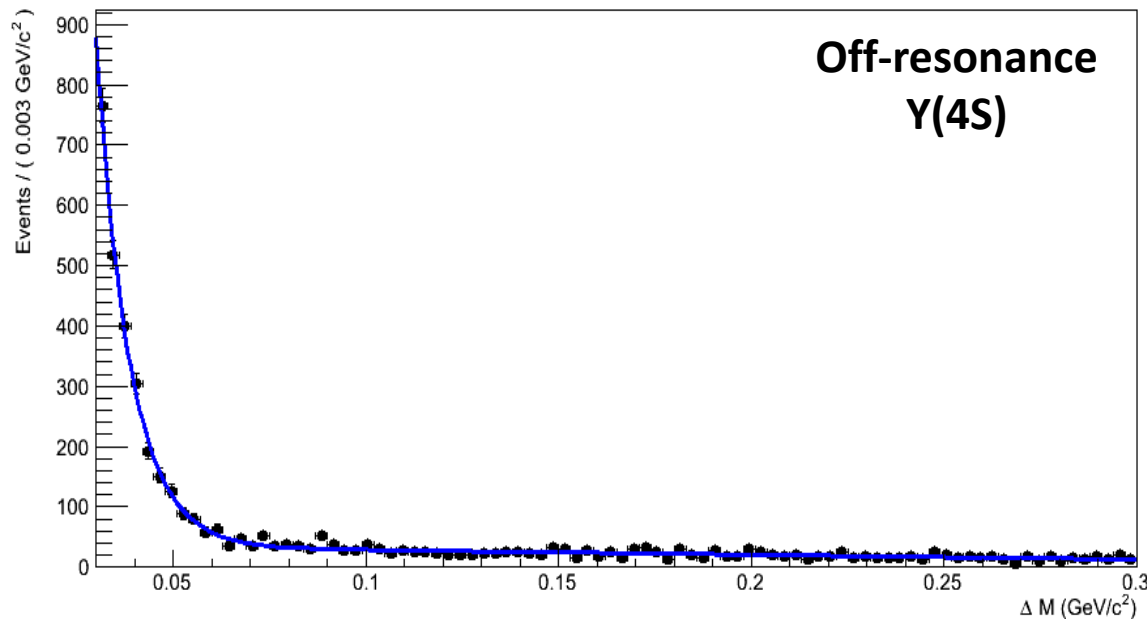
MC

MC

Signal MC is fitted with double gauss [Bifur gauss+gauss]



$\Upsilon(4S)$ off-resonance $\eta_b(2S)$ and χ_{b1} region



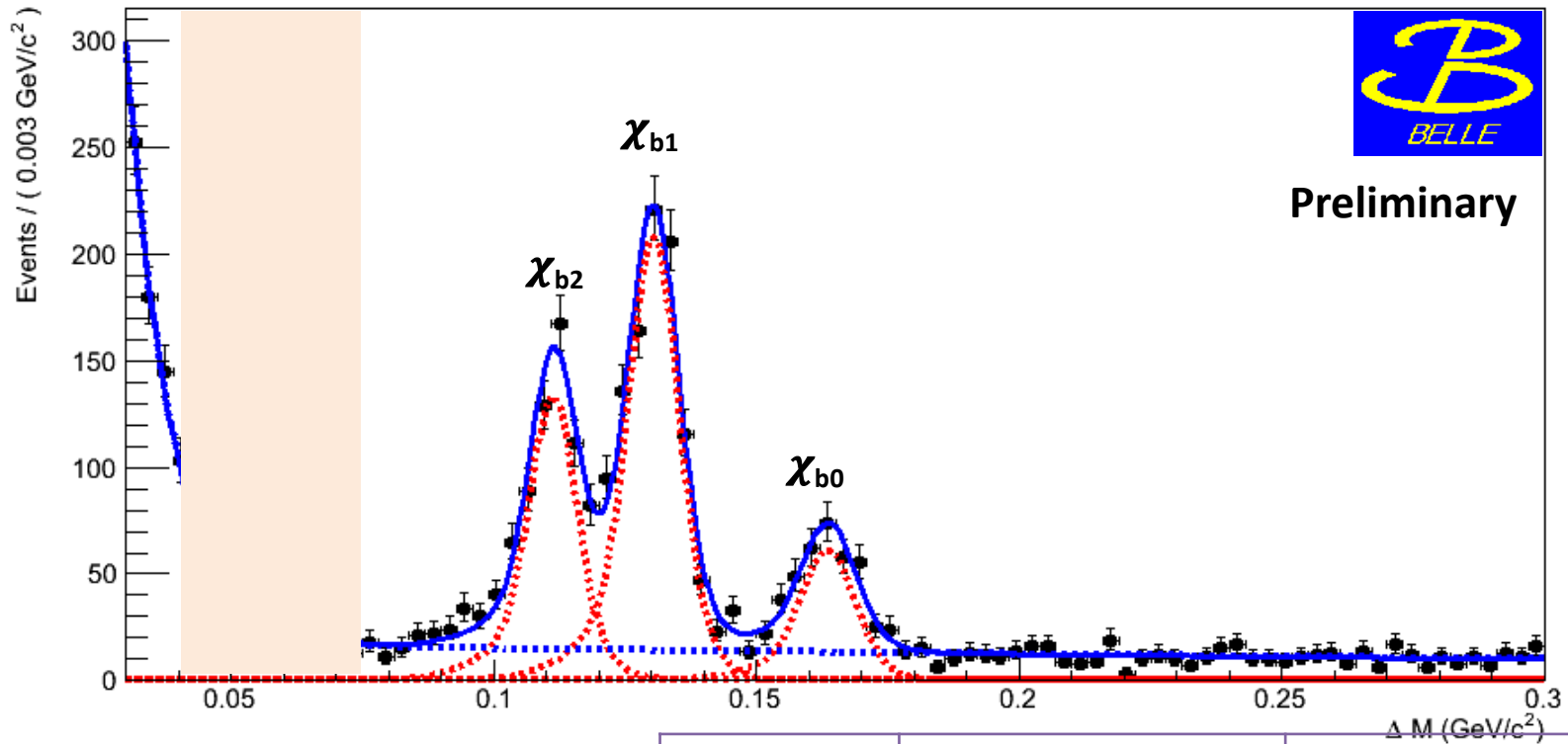
Off-resonance
 $\Upsilon(4S)$

The background at low energy is mainly coming from beam-background, which is exponential in nature and has long tail.

(To demonstrate this, we fitted $\Upsilon(4S)$ [89.5 fb⁻¹] off-resonance)
[exponential+chebyshev pol.]

$\Upsilon(2S) \rightarrow \gamma (b\bar{b})$

ΔM for $\eta_b(2S)$ and χ_{bJ} region

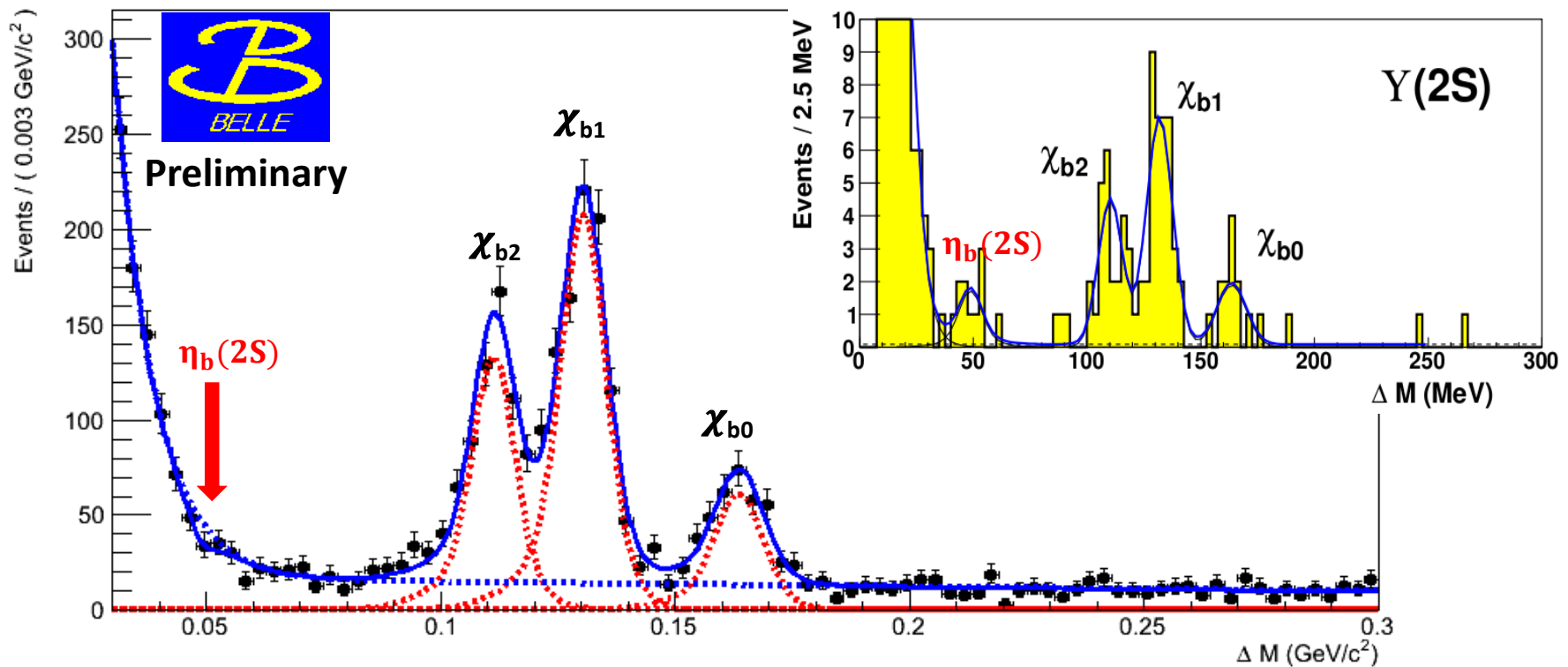


Large statistics available in our sample for χ_{bJ} (300-950 candidates) allows to determine precisely the χ_{bJ} masses. (with an accuracy competitive with PDG 2012).

	Mass (MeV/c ²)	Mass PDG (MeV/c ²)
$\chi_{b0}(1P)$	9859.63 ± 0.49	$9859.42 \pm 0.42 \pm 0.31$
$\chi_{b1}(1P)$	9892.83 ± 0.23	$9892.78 \pm 0.26 \pm 0.31$
$\chi_{b2}(1P)$	9912.00 ± 0.34	$9912.21 \pm 0.26 \pm 0.31$

$\Upsilon(2S) \rightarrow \gamma (b\bar{b})$

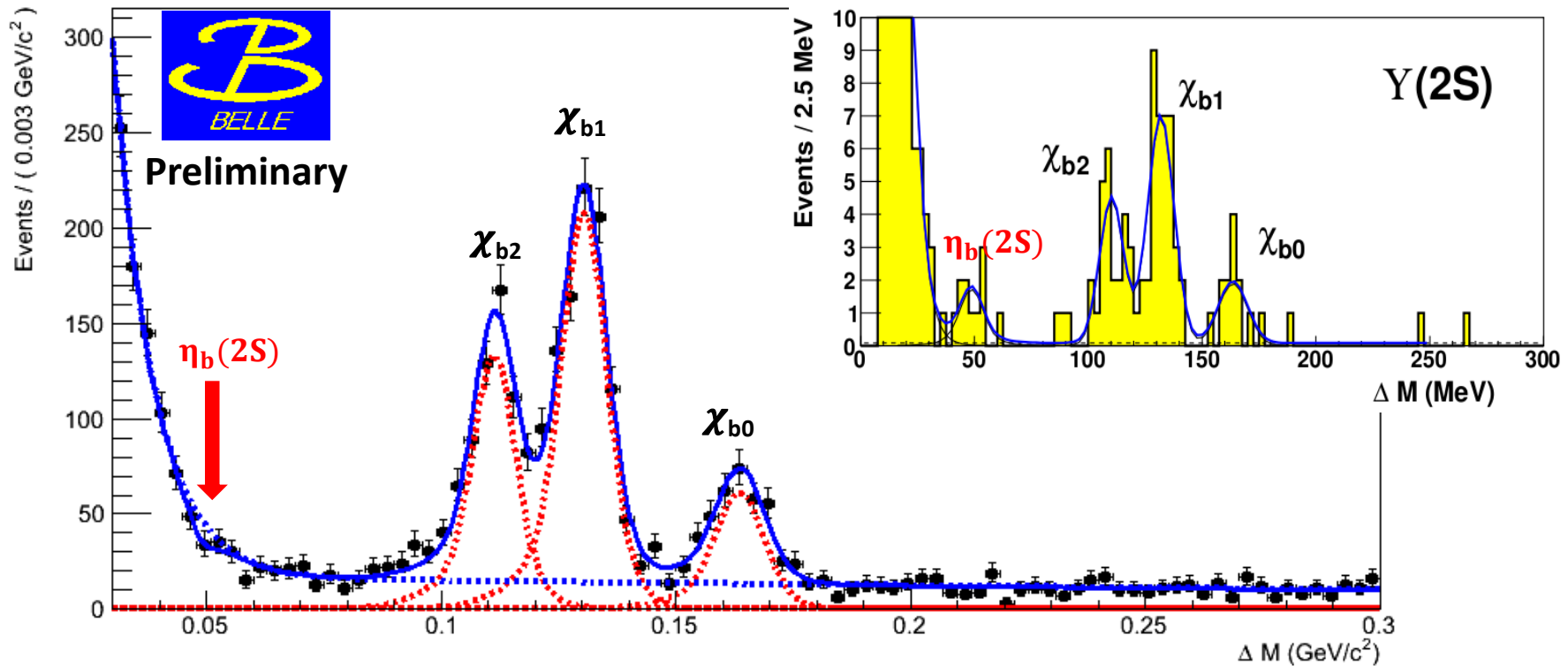
ΔM for $\eta_b(2S)$ and χ_{bJ} region



We did not find any signal corresponding to S. Dobbs et. al. $\eta_b(2S)$ candidate !!
 $N[\eta_b(2S)]$ candidate = -29.6 ± 19 (negative signal yield).

$\Upsilon(2S) \rightarrow \gamma (b\bar{b})$

ΔM for $\eta_b(2S)$ and χ_{bJ} region



We did not find any signal corresponding to S. Dobbs et. al. $\eta_b(2S)$ candidate !!

$$\mathcal{B}_1 \times \mathcal{B}_2 \equiv \mathcal{B}_1[\Upsilon(nS) \rightarrow \gamma \eta_b(nS)] \times \sum_{i=1}^{26} \mathcal{B}_{2i}[\eta_b(nS) \rightarrow h_i]$$

The Upper Limit on the B.F. (90% C.L.) $< 5.1 \times 10^{-6}$ (with sys.)

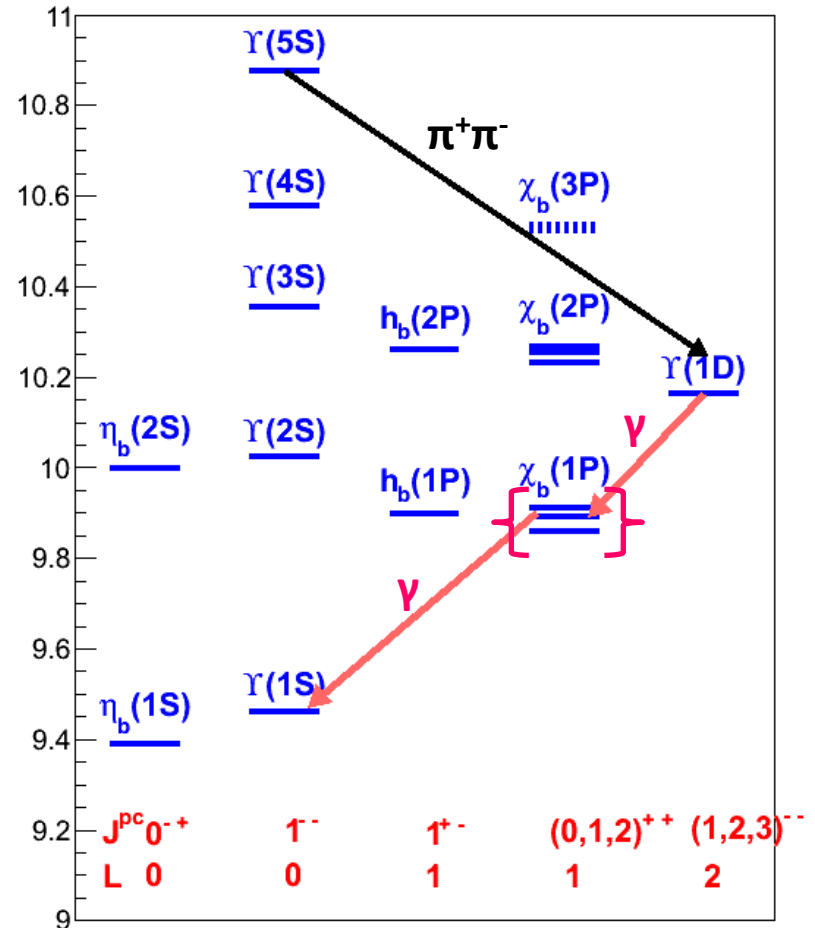
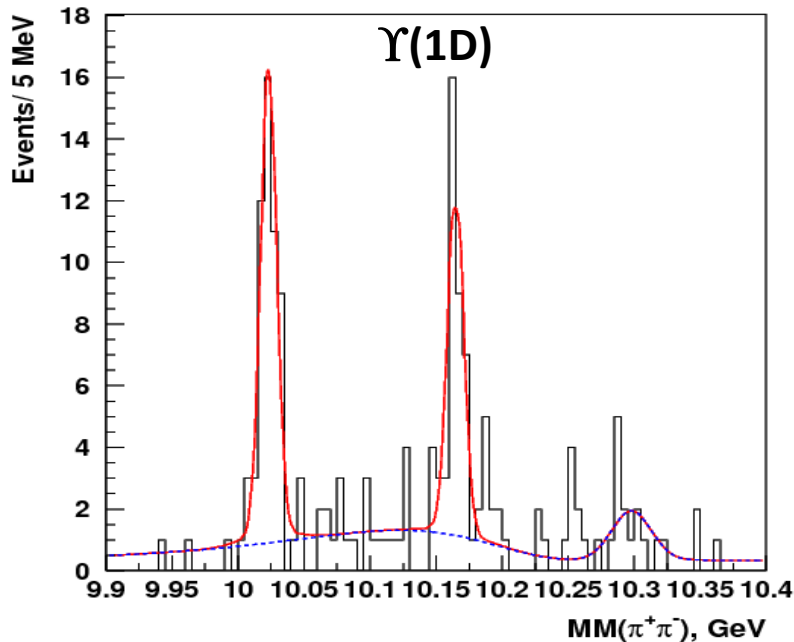
Reminder: $(46.2 \pm_{-14.2}^{+29.7} \pm 10.6) \times 10^{-6}$ S. Dobbs et. al.

$\Upsilon(1D)$ Mass

$$\Upsilon(5S) \rightarrow \Upsilon(1D)\pi^+\pi^-$$

$$\searrow \chi_{bJ}(1P)\gamma$$

$$\searrow \Upsilon(1S)\gamma$$



$$m_{\Upsilon(1D)} = 10164.7 \pm 1.4 \pm 1.0 \text{ MeV}/c^2$$

$$m_{\Upsilon(2S)} = 10023.2 \pm 1.0 \text{ MeV}/c^2 \quad \text{coincides with PDG value } 10023.26 \pm 0.31 \text{ MeV}/c^2$$

Babar [PRD82(2010)111102]: $m_{\Upsilon(1D)} = 10164.5 \pm 0.8 \pm 0.5 \text{ MeV}/c^2$

Cleo [PRD70(2004)032001]: $m_{\Upsilon(1D)} = 10161.1 \pm 0.6 \pm 1.6 \text{ MeV}/c^2$

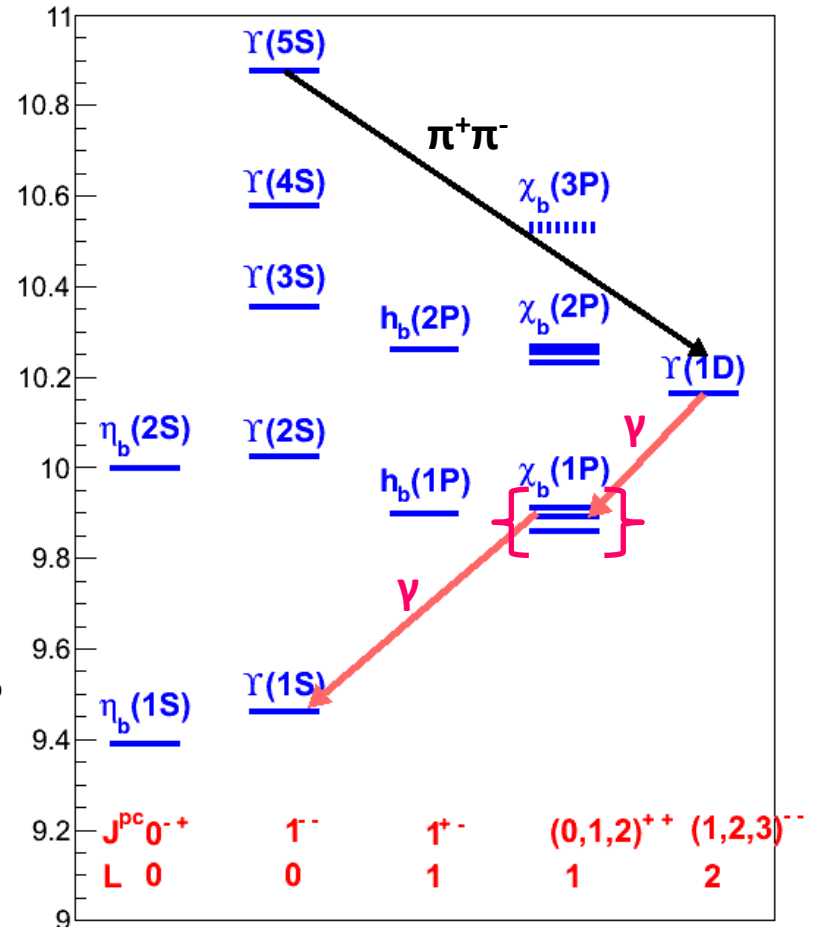
$\Upsilon(1D)$ Mass

$$\Upsilon(5S) \rightarrow \Upsilon(1D)\pi^+\pi^-$$

$$\searrow \chi_{bJ}(1P)\gamma$$

$$\searrow \Upsilon(1S)\gamma$$

- Three $\Upsilon(1D)$ states are predicted by theory
 $L=2, S=1 \Rightarrow J=1, 2$ and 3
- We assume production of $\Upsilon_1(1D)$, $\Upsilon_2(1D)$ and $\Upsilon_3(1D)$ proportional to $(2J+1)$ i.e. 3:5:7.
- Use B.F.s of $\Upsilon_1(1D) \rightarrow \chi_{b0}\gamma$, $\Upsilon_2(1D) \rightarrow \chi_{b1}\gamma$ and $\Upsilon_3(1D) \rightarrow \chi_{b2}\gamma$ from Kwong, Rosner
PRD 38, 279 (1998)
- $\mathcal{B}(\chi_{b0,1,2} \rightarrow \Upsilon(1S)\gamma) = 1.76\%, 33.9\%$ and 19.1% respectively from PDG-2012.
- $N \Upsilon_1(1D) : N \Upsilon_2(1D) : N \Upsilon_3(1D) = 10\% : 49\% : 41\%$
- Assuming only $\Upsilon_2(1D)$ and $\Upsilon_3(1D)$ contribute, (conservative assumption) we can fit the distribution to two peaks with fixed relative yields.



Splitting between $J=2$ and $J=3$ is $\Delta M < 10$ MeV at 90% CL (with sys.)

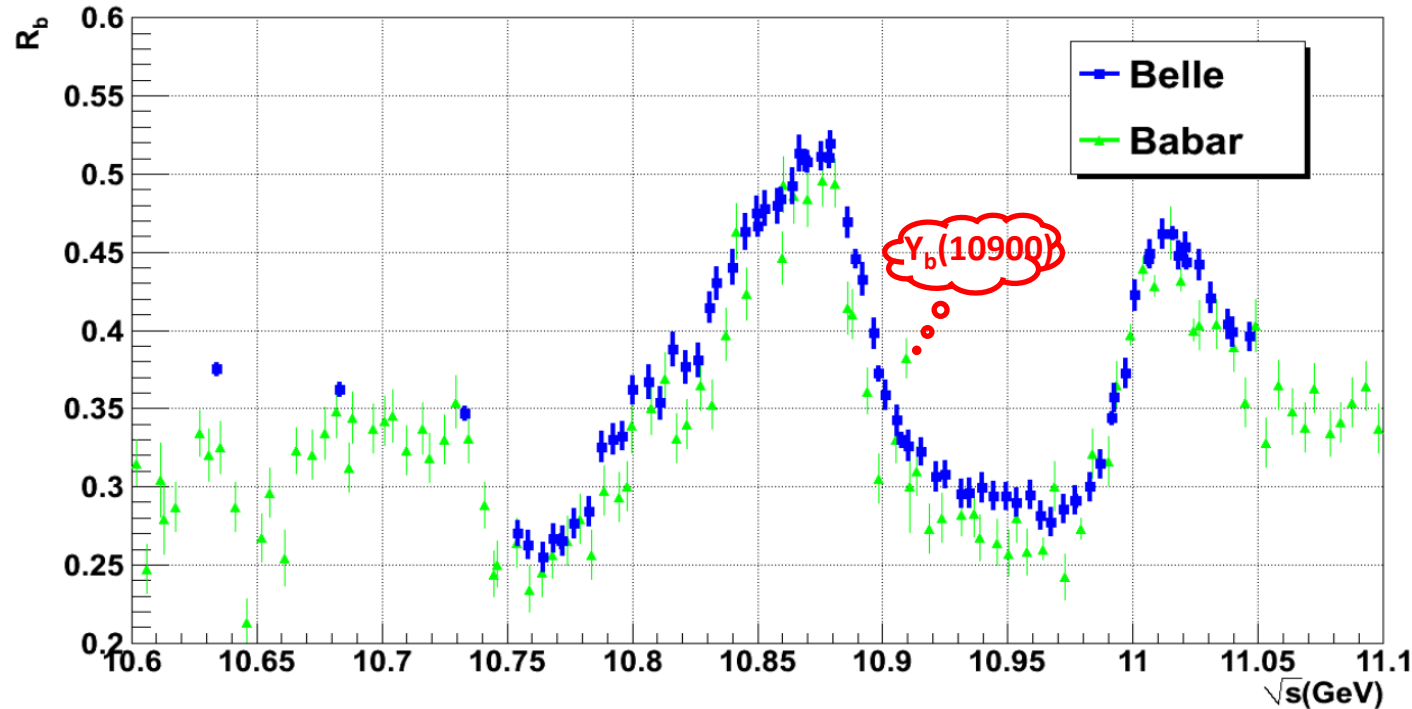
Potential model expectations: 4-11MeV

Search for Ali's $Y_b(10900)$

By definition R_b is given by:

$$R_b(s) \equiv \frac{\sigma_{b\bar{b}(\gamma)(s)}}{\sigma_{\mu\mu}^0(s)}$$

61 points from 10.750 GeV to 11.050 GeV with a step 5 MeV
around 50pb^{-1} for each energy point



- Better statistical errors, but covers a smaller energy range compared to Babar
- R_b is slightly higher by 0.0185
- No Ali's $Y_b(10900)$ (Phys. Lett. B 684, 28-39 2010)

Summary

- $h_b(nP) \rightarrow \eta_b(mS)\gamma$

$$m_{\eta_b(1S)} = 9402.4 \pm 1.5 \pm 1.8 \text{ MeV}/c^2 \quad \Gamma = 10.8_{-3.7}^{+4.0} {}_{-2.0}^{+4.5} \text{ MeV}/c^2$$

$$m_{\eta_b(2S)} = 9999.0 \pm 3.5_{-1.9}^{+2.8} \text{ MeV}/c^2$$

- $\Upsilon(2S) \rightarrow \gamma(b\bar{b})$

no signal found similar to S. Dobbs et. al. [attributed to $\eta_b(2S)$] 

Upper Limit on the B.F. (90% C.L.) $< 5.1 \times 10^{-6}$

- $\Upsilon(1D)$ 

$$m_{\Upsilon(1D)} = 10164.7 \pm 1.4 \pm 1.0 \text{ MeV}/c^2$$

Splitting between J=2 and J=3, $\Delta M < 10 \text{ MeV}$ at 90% CL.

- No Ali's $Y_b(10900)$ found in R_b scan.



Backup

S. Dobbs signal MC embedded

ΔM for $\eta_b(2S)$ and χ_{bJ} region

