

RTA-Resummed Stochastic Diffusion Beyond White Noise: Markovization and W_n Cumulant Dynamics



Navid Abbasi

With **Mei Huang, Dirk Rischke, Pei Zhang**



1. QCD CEP

- Landau-Ginzburg picture in thermal equilibrium

$$\Omega = \int d^3x \left[\frac{1}{2} (\nabla \sigma)^2 + \frac{m_\sigma^2}{2} \sigma^2 + \frac{\lambda_3}{3} \sigma^3 + \frac{\lambda_4}{4} \sigma^4 + \dots \right]. \quad \xi = m_\sigma^{-1}.$$

$$\begin{aligned} \kappa_2 = \langle \sigma_0^2 \rangle &= \frac{T}{V} \xi^2; & \kappa_3 = \langle \sigma_0^3 \rangle &= \frac{2\lambda_3 T}{V} \xi^6; \\ \kappa_4 = \langle \sigma_0^4 \rangle_c &\equiv \langle \sigma_0^4 \rangle - \langle \sigma_0^2 \rangle^2 = \frac{6T}{V} [2(\lambda_3 \xi)^2 - \lambda_4] \xi^8. \end{aligned}$$

[Stephanov PRL (2009)]

- Critical point $\equiv \xi \rightarrow \infty$
- Higher cumulants scale with higher powers of ξ
- non-Gaussian fluctuations (specially κ_4) are most sensitive to CEP

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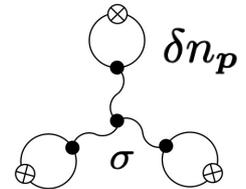
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- Contribution to particle correlators

$$\begin{aligned} \mathcal{C}_3 &= \langle \delta n_{\mathbf{p}_1} \delta n_{\mathbf{p}_2} \delta n_{\mathbf{p}_3} \rangle_\sigma \\ &\sim \kappa_3 G_{\sigma \delta n}(\mathbf{p}_1) G_{\sigma \delta n}(\mathbf{p}_2) G_{\sigma \delta n}(\mathbf{p}_3) \end{aligned}$$



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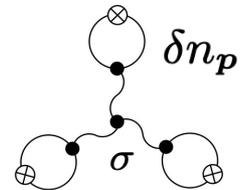
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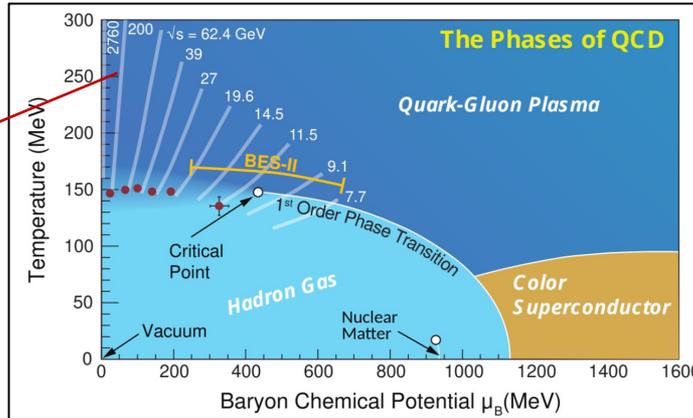
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Fluctuations of the order parameter (baryon density fluctuations)

The fluctuations of the net proton number (cumulants)

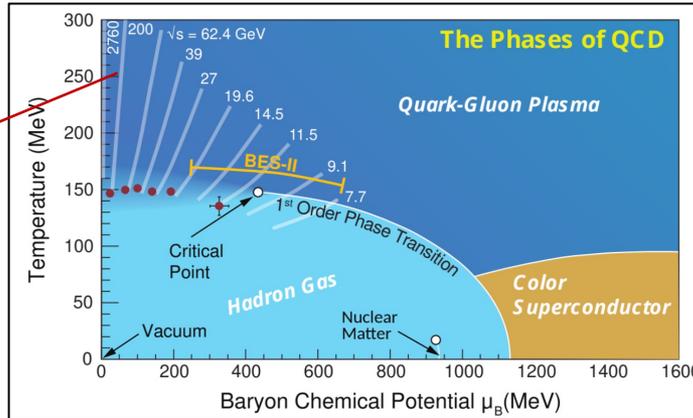
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QGP droplet trajectory

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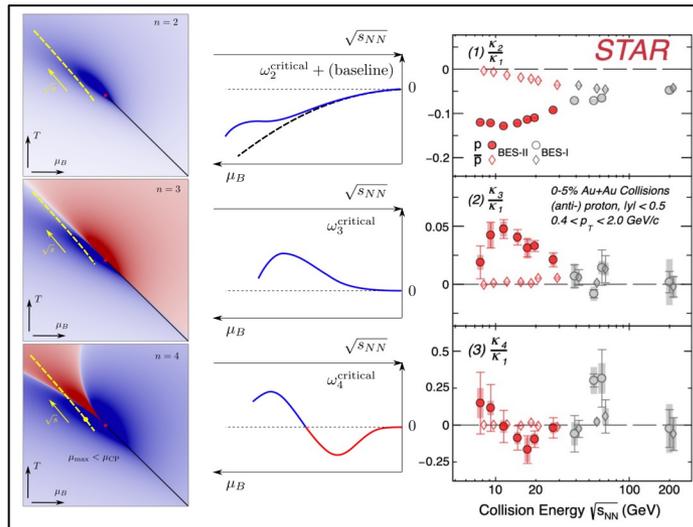


QGP droplet trajectory

- Experiment measures net-proton cumulants:

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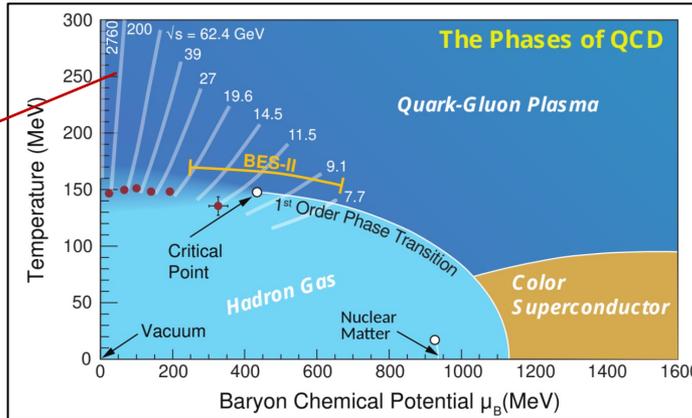
Non-monotonic behavior of the cumulants:



$$\omega_n \equiv K_n / K_1$$

[Stephanov 2024]

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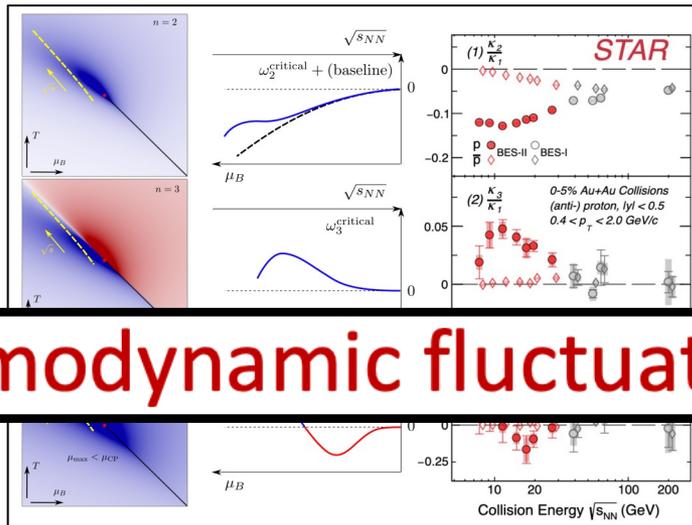


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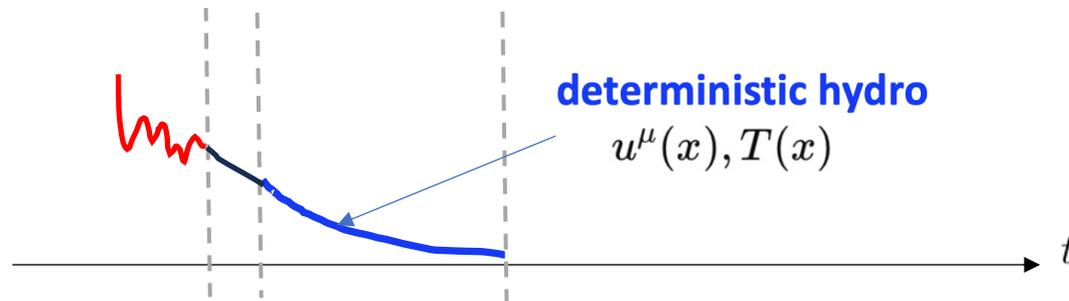
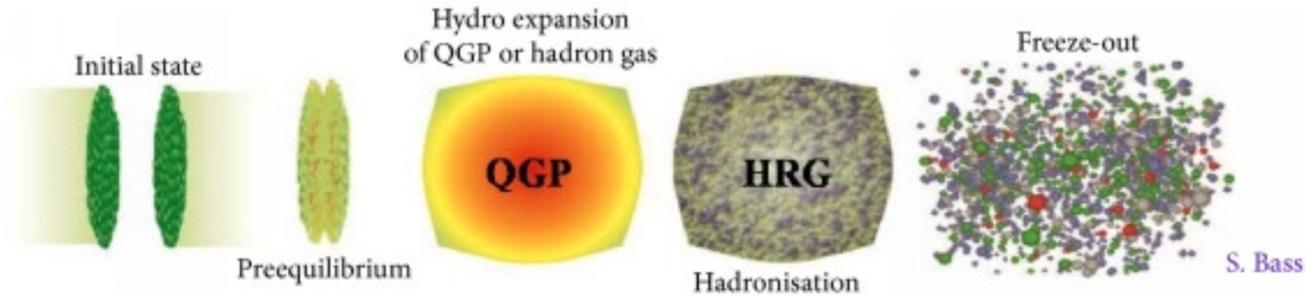


thermodynamic fluctuations!

[Stephanov 2024]

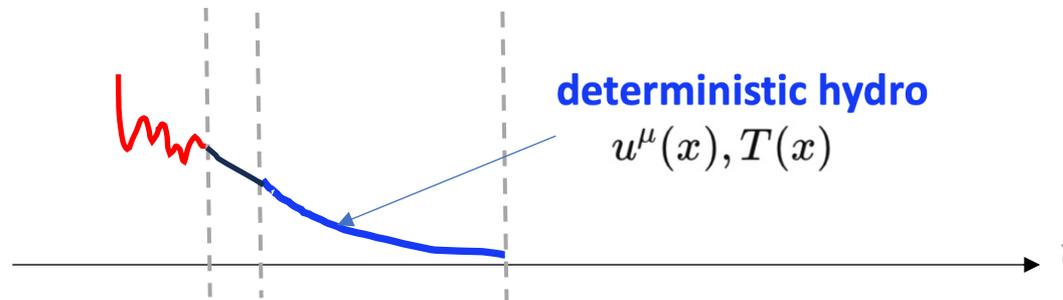
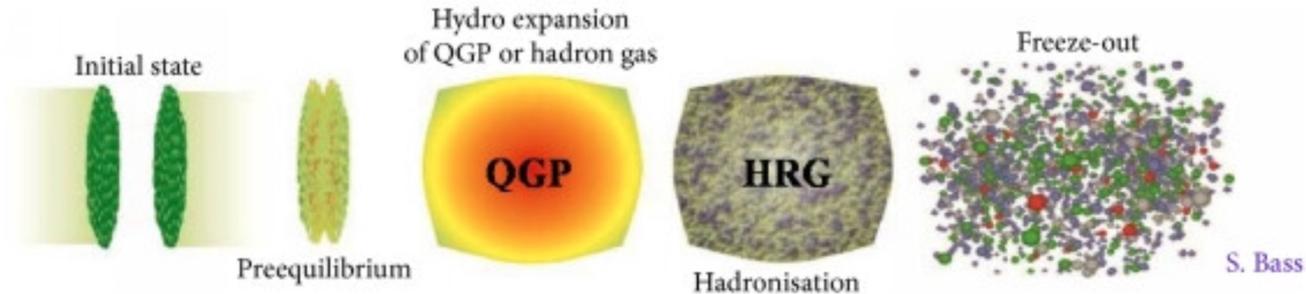
Theory needs **dynamical evolution + acceptance mapping**

3. QGP dynamics and Fluctuations



- Initial state Fluctuations $\rightarrow v_n$ [Heinz, Snellings Ann.Rev.Nucl.Part.Sci. (2013)]
- Sampling at freeze-out \rightarrow additional event-by-event fluctuations [Bzdak, Koch, Shen PRC (2012)]

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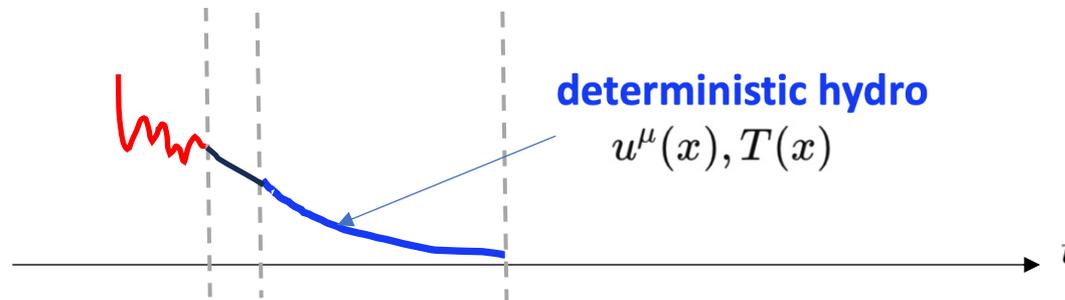
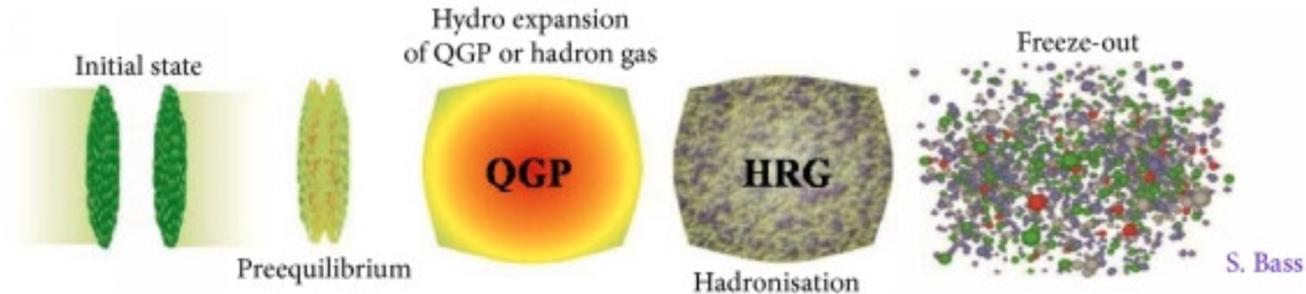
(local equilibrium + white noise)

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[Landau-Lifshitz Vol 9]

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what happens near CEP?

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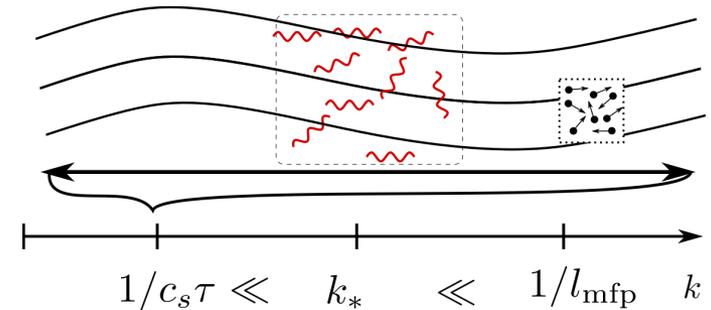
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- ❑ Fluctuations of $\sim k^*$ are out of equilibrium:

- $\gamma_\eta k_*^2 \sim \theta$, $k_* \sim \left(\frac{\theta}{\gamma_\eta}\right)^{1/2}$ separates equilibrated / non-equilibrated modes.
- Wigner function for $q \sim k_*$ evolves locally:

$$W_2(t, \mathbf{x}; \mathbf{q}) = \int_{\mathbf{q}} e^{-i\mathbf{d}\cdot\mathbf{y}} \langle \delta\phi(t, \mathbf{x} - \frac{\mathbf{y}}{2}) \delta\phi(t, \mathbf{x} + \frac{\mathbf{y}}{2}) \rangle$$



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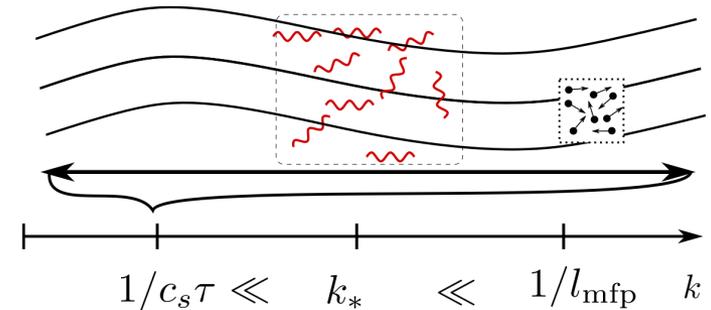
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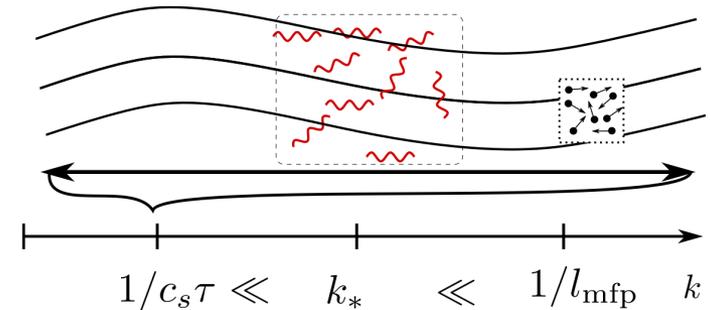
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- Evolution of W_2 in a Bjorken flow

[Akumatsu, Mazeliauskas, Teaney PRC (2017)]

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- Evolution of W_2 in relativistic hydro with a critical point

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5. Deterministic W_n from stochastic dynamics

In the diffusion theory, response is **instantaneous**:

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$$\partial_t \check{n} = -\nabla \cdot \check{\mathbf{J}}$$

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with the **white noise**:

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Evolution equations:

$$\begin{aligned} \partial_t W_2(\mathbf{q}_1) &= - \left[\gamma \mathbf{q}_1^2 W_2(\mathbf{q}_2) + \lambda \mathbf{q}_1 \cdot \mathbf{q}_2 \right]_{\text{P}\mathbf{q}_1 \mathbf{q}_2} \\ \partial_t W_3(\mathbf{q}_1, \mathbf{q}_2) &= - \left[\frac{1}{2} \gamma \mathbf{q}_1^2 W_3(\mathbf{q}_2, \mathbf{q}_3) + \frac{1}{2} \gamma' \mathbf{q}_1^2 W_2(\mathbf{q}_2) W_2(\mathbf{q}_3) + \lambda' \mathbf{q}_1 \cdot \mathbf{q}_2 W_2(\mathbf{q}_3) \right]_{\text{P}\mathbf{q}_1 \mathbf{q}_2 \mathbf{q}_3} \end{aligned}$$

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This motivates going beyond white noise.

6. W_n and cumulants

Evolution equations should be constructed for

- The set of slow fields near CEP: j 's = (m, p)
- Coupled to an expanding background

$$u \cdot \partial G_{mm}^c(x_1, x_2) = 2 \left[- (y_1 \cdot \partial u) \cdot \frac{\partial}{\partial x_1} G_{mm}^c(x_1, x_2) + L_{m,m}(x_1) G_{mm}^c(x_1, x_2) + Q_{mm}(x_1, x_2) \right]_{12} \quad (3.15a)$$

$$u \cdot \partial G_{mmm}^c(x_1, x_2, x_3) = 3 \left[- (y_1 \cdot \partial u) \cdot \frac{\partial}{\partial x_1} G_{mmm}^c(x_1, x_2, x_3) + L_{m,m}(x_1) G_{mmm}^c(x_1, x_2, x_3) \right. \\ \left. + L_{m,p}(x_1) G_{pmm}^c(x_1, x_2, x_3) + L_{m,mm}(x_1) G_{mm}^c(x_1, x_2) G_{mm}^c(x_1, x_3) \right. \\ \left. + 2Q_{mm,m}(x_1, x_2) G_{mm}^c(x_1, x_3) \right]_{123} \quad (3.15b)$$

$$u \cdot \partial G_{mmmm}^c(x_1, x_2, x_3, x_4) = 4 \left[- (y_1 \cdot \partial u) \cdot \frac{\partial}{\partial x_1} G_{mmmm}^c(x_1, x_2, x_3, x_4) + L_{m,m}(x_1) G_{mmmm}^c(x_1, x_2, x_3, x_4) \right. \\ \left. + L_{m,p}(x_1) G_{pmmm}^c(x_1, x_2, x_3, x_4) + 3L_{m,mm}(x_1) G_{mm}^c(x_1, x_2) G_{mm}^c(x_1, x_3, x_4) \right. \\ \left. + 3L_{m,mp}(x_1) G_{mm}^c(x_1, x_2) G_{pmm}^c(x_1, x_3, x_4) \right. \\ \left. + L_{m,mmm}(x_1) G_{mm}^c(x_1, x_2) G_{mm}^c(x_1, x_3) G_{mm}^c(x_1, x_4) \right. \\ \left. + 3Q_{mm,p}(x_1, x_2) G_{pmm}^c(x_1, x_3, x_4) + 3Q_{mm,m}(x_1, x_2) G_{mm}^c(x_1, x_3, x_4) \right. \\ \left. + 3Q_{mm,mm}(x_1, x_2) G_{mm}^c(x_1, x_3) G_{mm}^c(x_1, x_4) \right]_{1234} \quad (3.15c)$$

[An, Basar, Stephanov, Yee PRC (2023)]

the next step is to map W_n onto the particle multiplicity cumulant C_n

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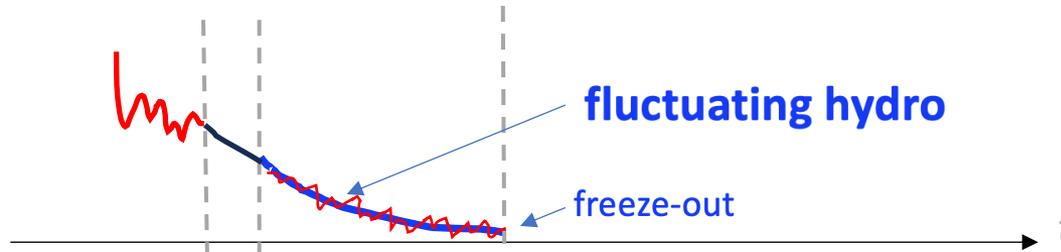
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[An, Basar, Stephanov, Yee PRC (2023)]

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- Coupled to an expanding background

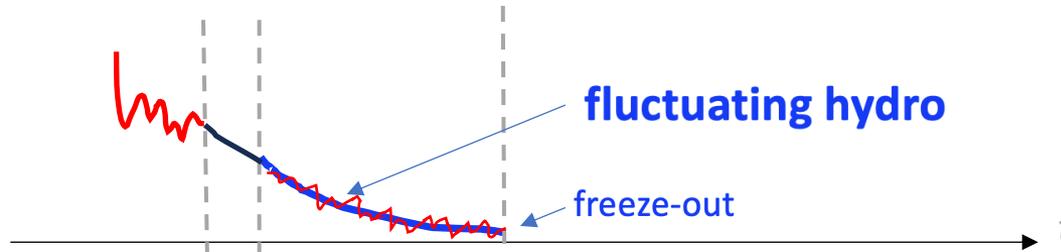
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- Sudden freeze-out time + acceptance filter $\rightarrow C_n$

$$C_m(\Delta y) = \int \left[\prod_{i=1}^m \frac{dq_i}{2\pi} A(q_i) \right] (2\pi) \delta \left(\sum_{i=1}^m q_i \right) W_m(q_1, \dots, q_m; t_f).$$

$$A(q) = e^{iqy_0} \frac{2 \sin(q\Delta y/2)}{q}$$

↑ rapidity
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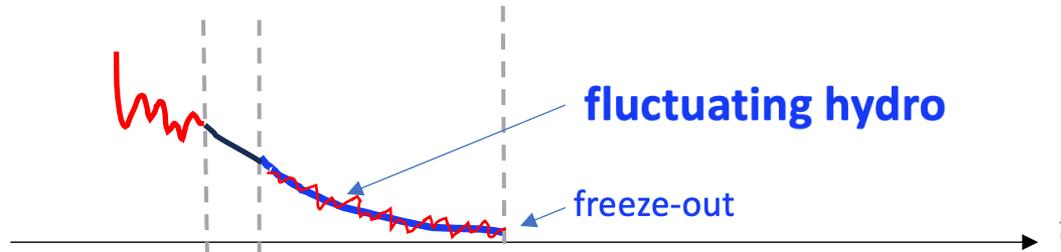
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We aim to explore **the effect of colored noise** on the **evolution of W_n** , and consequently on **the cumulants**, in an RTA setup

7. Noise in RTA diffusion

- Integrating RTA gives a resummed theory of diffusion

[NA, Rischke JHEP (2025)]

$$p^\mu \partial_\mu f(x, \mathbf{p}) = \frac{p^\alpha u_\alpha}{\tau} [f(x, \mathbf{p}) - f^{(0)}(x, \mathbf{p})] \quad \xrightarrow{\mathbf{D} = \partial_t + \mathbf{v} \cdot \nabla} \quad \boxed{\int_\Omega \frac{\mathbf{D}}{1 + \tau \mathbf{D}} n(x) = 0}$$

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Derivative expansion: $\partial_t n + \nabla_i \left[-\frac{\tau}{3} \nabla_i n + \frac{\tau^3}{45} \nabla_i \nabla^2 n - \frac{2\tau^5}{945} \nabla_i \nabla^4 n + \dots \right] = 0$

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$$\langle \zeta(\omega, k) \zeta(\omega', k') \rangle = (2\pi)^{d+1} \delta(\omega + \omega') \delta(k + k') \mathcal{N}(\omega, k) \frac{|1 - L_p|^2}{\tau^2} \chi T \tau \left(\frac{L_p}{1 - L_p} + \frac{L_p^*}{1 - L_p^*} \right)$$

noise is colored

$$L_p = \frac{1}{2i\tau k} \ln \left(\frac{\frac{i}{\tau} + \omega - k}{\frac{i}{\tau} + \omega + k} \right)$$

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RTA-Resummed Stochastic Diffusion Beyond White Noise:

Markovization and W_n Cumulant Dynamics

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we use the “**Markovian embedding technique**”

[Gardiner 2004]

8. Markovian embedding

non-Markovian

$$\left\langle \frac{\mathcal{D}}{1 + \tau \mathcal{D}} n_k(t) \right\rangle_{\mu} = \zeta_k(t)$$

$$\langle f \rangle_{\mu} \equiv \frac{1}{2} \int_{-1}^1 d\mu f(\mu) \quad \mathcal{D} \equiv \partial_t + ik\mu$$

→
Markovianization by introducing

a hidden variable:

$$\mu = \cos \theta$$

Markovian

$$(1 + \tau \mathcal{D}) g_k(\mu, t) = n_k(t) + \tau \xi_k(\mu, t)$$

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- $\xi_k(\mu, t)$ is chosen to be white.
- Integrating out $\mu = \cos \theta \rightarrow$ the original non-Markovian dynamics:

$$\langle \xi(\mu, t) \xi(\mu', t') \rangle = \frac{2\chi T}{\tau} [2\delta(\mu - \mu') - 1] \delta(t - t').$$

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The goal is to approximately solve $g_k(\mu, t)$

9. Hierarchy of Markovian embeddings

- The key point is:

$$\mu P_\ell = \frac{\ell+1}{2\ell+1} P_{\ell+1} + \frac{\ell}{2\ell+1} P_{\ell-1}$$

- The projected equation onto ℓ

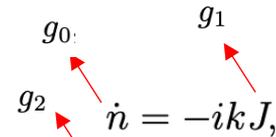
$$\tau \dot{g}_\ell + g_\ell + i\tau k \left(\frac{\ell}{2\ell+1} g_{\ell-1} + \frac{\ell+1}{2\ell+1} g_{\ell+1} \right) = \delta_{\ell 0} n + \tau \xi_\ell$$

- Hierarchy of the equations

$$\ell = 0.$$

$$\ell = 1.$$

$$\ell = 2.$$

$$\vdots$$


$$\tau \dot{J} + J + i\tau k \left(\frac{1}{3} n + \frac{2}{3} Q \right) = \eta_1,$$

$$\tau \dot{Q} + Q + i\tau k \left(\frac{2}{5} J \right) = \eta_2.$$

$$\langle \eta_1(t) \eta_1^*(t') \rangle = \frac{2\chi T \tau}{3} \delta(t - t'),$$

$$\langle \eta_2(t) \eta_2^*(t') \rangle = \frac{2\chi T \tau}{5} \delta(t - t'),$$

Truncation at $\ell = \ell_{\max}$ gives a finite Markovian system.

- Evolution equation for W_2 (truncated at $\ell_{\max} = 1$)

$$W_2 \equiv \langle n_k n_{-k} \rangle, \quad P \equiv -i \langle J_k n_{-k} \rangle, \quad V \equiv \langle J_k J_{-k} \rangle$$

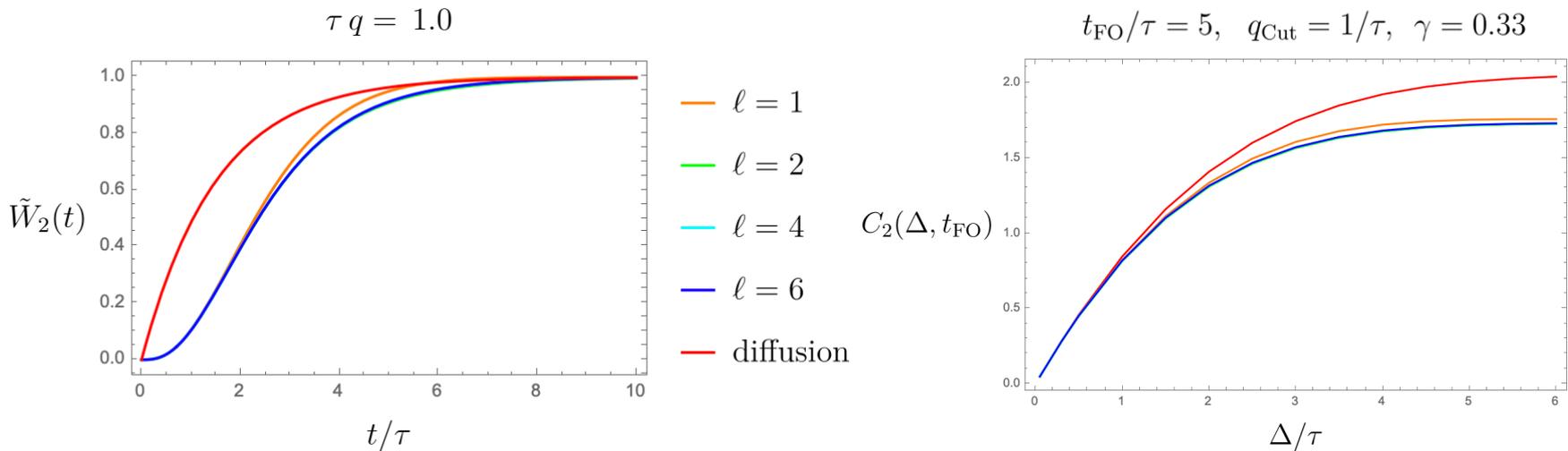
$$\dot{W}_2 = 2kP$$

$$\dot{P} = -\frac{1}{\tau} P - \frac{k}{3} W_2 + kV$$

$$\dot{V} = -\frac{2}{\tau} V - \frac{2k}{3} P + \frac{2\chi T}{3\tau}$$

10. Numerical results: C_2

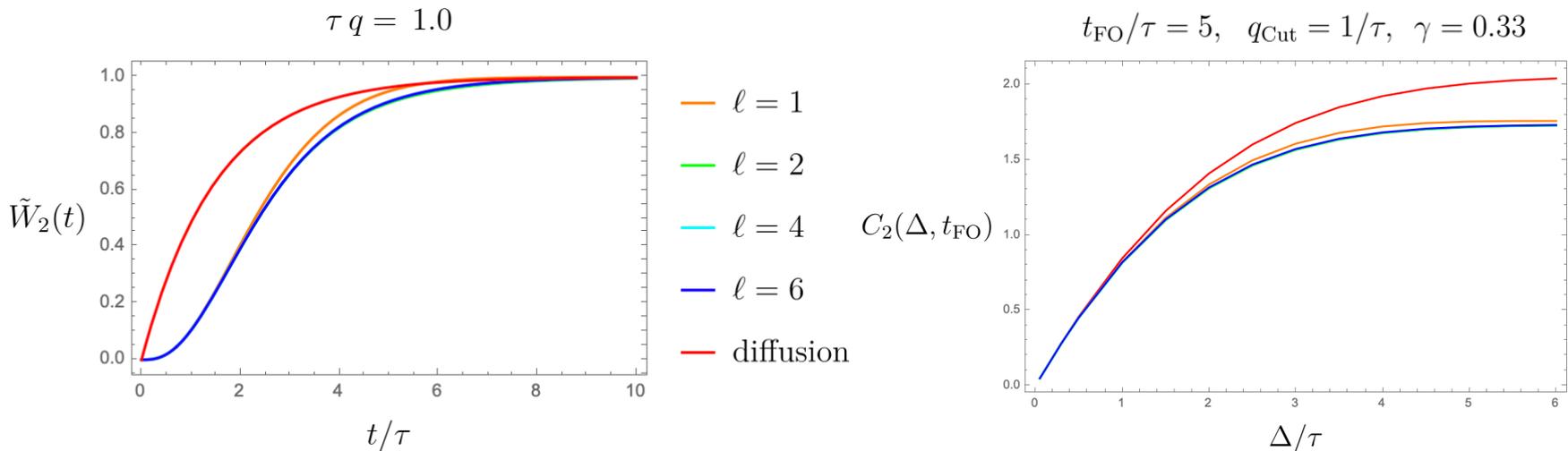
diffusion vs. RTA truncated at $\ell_{max} = 1, 2, 4, 6$



- Memory \rightarrow lower equilibration of high q -modes
- W_2 rises more slowly \rightarrow suppressed C_2
- Stable for $\ell_{max} \gtrsim 4 \rightarrow$ baseline benchmark for C_3 and C_4

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diffusion vs. RTA truncated at $\ell_{max} = 1, 2, 4, 6$

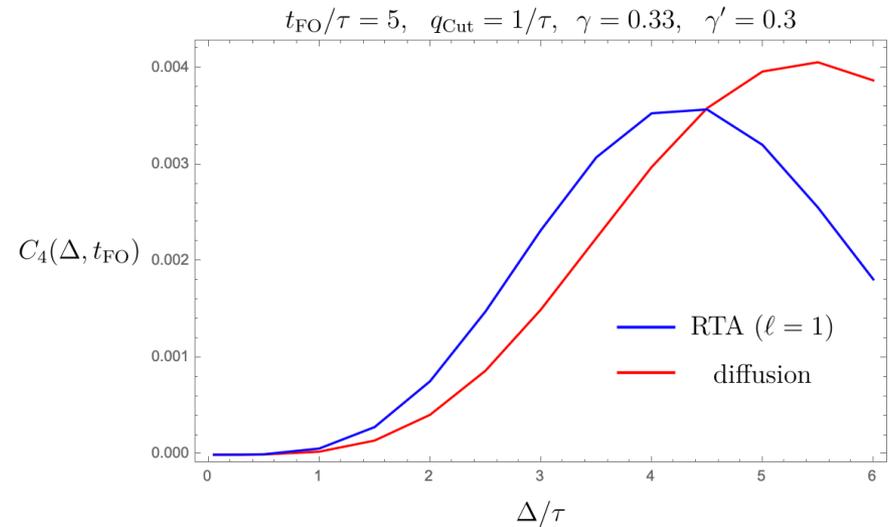
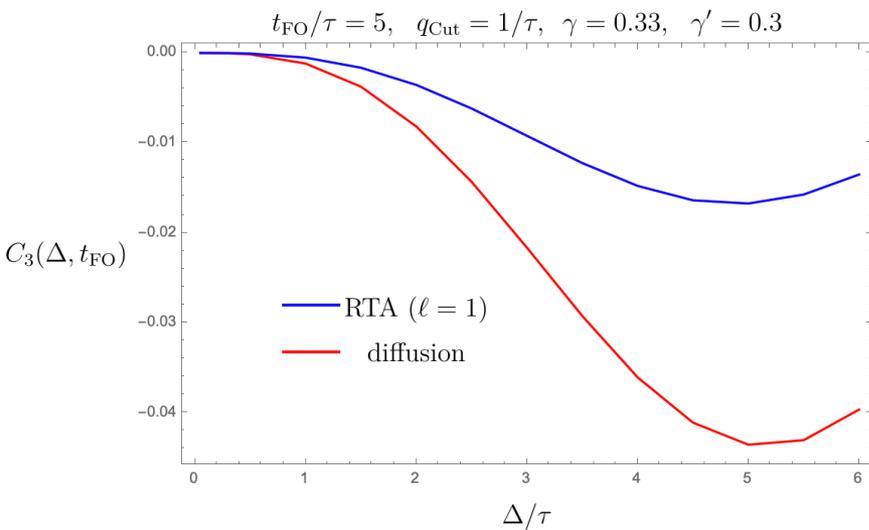


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“ C_2 suppression from memory”

11. Numerical results: C_3 and C_4

- Currently shown for $\ell_{max} = 1$



- C_3 : the same trend as C_2
- C_4 : ordering changes with Δ
→ C_4 is more sensitive to nonlinear source terms

colored-noise memory effects can change
the **magnitude** and also **the shape of higher cumulants**
→ Important for quantitative comparisons to data

12. Conclusion and outlook

1. Memory isn't a small detail!

- it's required by FDT once relaxation is finite.
- It mainly affects high- k ,
- so it shows up first at small rapidity windows.

2. We have done a **microscopic derivation** of the colored noise, tied to a transport model.

3. We developed a concrete, computable scheme for W_n hierarchy.

3. The next step is to couple the equations to a **expanding background** and **critical EoS**.

4. Final goal: a fully fluctuating RTA + expanding background + critical EoS + slow mode

Thank you for your attention!