



New axion misalignment mechanism effect

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History and progress

QCD Axion

Strong CP problem

Peccei-Quinn mechanism

Axion

$$d_n = 5.2 \times 10^{-16} \bar{\theta} e \cdot \text{cm}$$

$$d_n \leq 10^{-26} e \cdot \text{cm}$$

Pseudo-scalar particle

[Preskill, Wise, Wilczek (1983)]

[Abbott, Sikivie (1983)]

[Dine, Fischler (1983)]

Misalignment mechanism : $\phi_0 \neq 0, \dot{\phi} = 0$

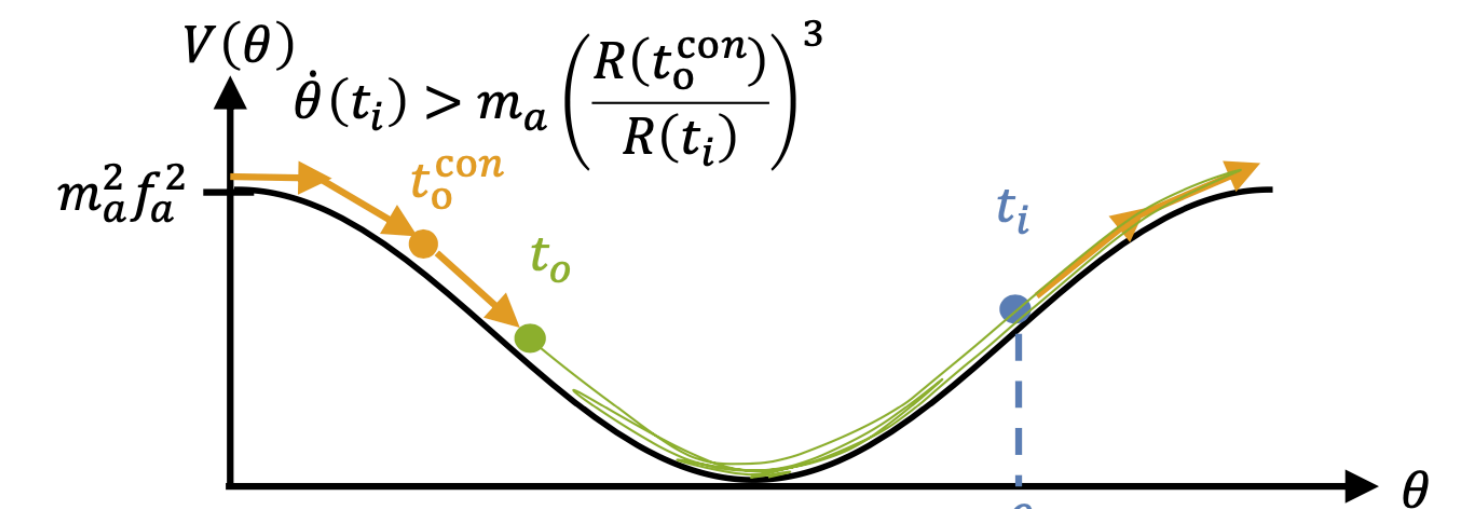
$$\ddot{\phi} + 3H\dot{\phi} + \frac{\Lambda^4(T)}{f^2}\phi = 0$$

$$\rho_{\text{DM}} \sim m_\phi \left[\frac{a(T_{\text{osc}})}{a_0} \right]^3 \left[\frac{\Lambda(T_{\text{osc}})^4 \theta_i^2}{m_\phi(T_{\text{osc}})} \right]$$

Kinetic Misalignment mechanism & Axiogenesis

Kinetic Misalignment mechanism : $\phi_0 \neq 0, \dot{\phi} \neq 0$

[Co, Hall, Harigaya (2019)]
[Chang, Cui (2019)]



Non-zero Peccei-Quinn number \rightarrow Axiogenesis!

Strong sphaleron

$$n_{PQ} = S^2 \dot{\phi}$$

EWsphaleron

Quake chiral asymmetry

EWsphaleron

Baryon asymmetry

PRL, 124, 111602

History and progress

Trapped Misalignment mechanism : initial $\theta \sim \pi$

- *L. Di Luzio, B. Gavela, P. Quilez, A. Ringwald, JCAP 10 (2021) 001*

Hilltop Misalignment mechanism : initial $\theta \sim \pi$

- *Raymond T. Co, Eric Gonzalez, Keisuke Harigaya, JHEP 05 (2019) 163*

Acoustic Misalignment mechanism: Fluctuations of θ

- *A. Bodas, R.T. Co, A. Ghalsasi, K. Harigaya and L.T. Wang, JHEP 08(2025)131*

Thermal Misalignment mechanism: new thermal potential term

- *B. Batell, A. Ghalsasi, Phys. Rev. D 107(2023)L091701*
- *Y. Zhang, PRL, 132(2024)8, 081003*

Outline

In this talk, we focus on the axion misalignment mechanism in the presence of exotic axion mass generation mechanism, and in the presence of primordial magnetic field.

- **Scenario-2 : Axion misalignment mechanism with specific ALP mass generation mechanism**

Wei Chao, Mingjie Jin, Hai-jun Li, Ying-quan Peng and Yue Wang, Phys.Rev.D 109(2024)115027

- **Scenario-3 : Axion misalignment mechanism in the presence of primordial magnetic field**

Wei Chao, Chang-jie Dai, arXiv: [2604.02012](https://arxiv.org/abs/2604.02012)

ALP mass via the type-II seesaw mechanism

Type-II seesaw + spontaneous breaking $U(1)_L$ symmetry

$$V(S, \Phi, \Delta) = V(\Phi, \Delta) - \mu_S^2(S^\dagger S) + \lambda_6(S^\dagger S)^2$$

LNV term!

$$+ \lambda_7(S^\dagger S)(\Phi^\dagger \Phi) + \lambda_8(S^\dagger S)\text{Tr}(\Delta^\dagger \Delta) + \mu\Phi^T i\tau_2 \Delta^\dagger \Phi + \lambda S\Phi^T i\tau_2 \Delta^\dagger \Phi + \text{h.c.},$$

$$\Phi = \begin{pmatrix} \phi^+ \\ \frac{v_\phi + \phi + i\chi}{\sqrt{2}} \end{pmatrix}$$

$$\Delta = \begin{pmatrix} \frac{\Delta^+}{\sqrt{2}} & \Delta^{++} \\ \frac{v_\Delta + \delta + i\xi}{\sqrt{2}} & \frac{\Delta^+}{\sqrt{2}} \end{pmatrix}$$

$$S = \frac{v_s + \tilde{s} + i\tilde{a}}{\sqrt{2}}$$

\tilde{a} : Majoron

Yukawa Interaction

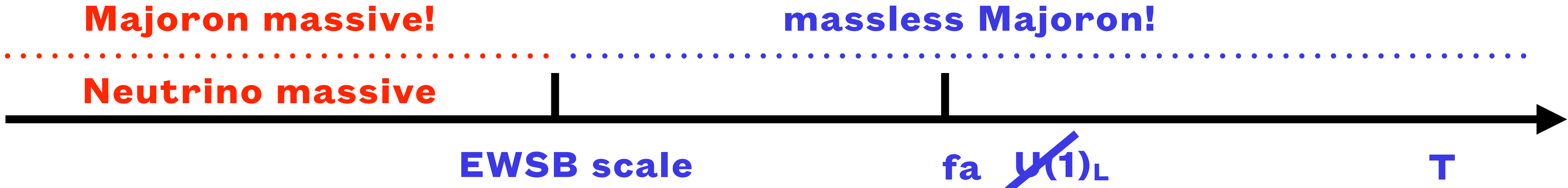
$$-\mathcal{L}_\Delta = Y_{\alpha\beta} \overline{\ell_L^{\alpha C}} i\sigma^2 \Delta \ell_L^\beta + \text{h.c.}$$

Key term:

$$\mu\Phi^T i\sigma^2 \Delta \Phi + \text{h.c.}$$

ALP mass via the type-II seesaw mechanism

Sequential breaking of various symmetries



Temperature of Universe

$$(m_\nu)_{\alpha\beta} = y_{\alpha\beta} v_\Delta / \sqrt{2}.$$

$$m_a^2 = \frac{\sqrt{2} \mu v_\phi^2 v_\Delta (v_\phi^2 + 4v_\Delta^2)}{2v_\phi^2 (v_\Delta^2 + v_s^2) + 8v_\Delta^2 v_s^2} \approx \frac{\mu v_\phi^2 v_\Delta}{\sqrt{2} v_s^2},$$

ALP masses

Majoron mass should arise from cosine like potential!

$$\left. \begin{aligned} \ell_L &\rightarrow e^{-\frac{ia}{2f}} \ell_L & S &\rightarrow e^{+\frac{ia}{f}} S \\ E_R &\rightarrow e^{-\frac{ia}{2f}} E_R & \Delta &\rightarrow e^{-\frac{ia}{f}} \Delta \\ H &\rightarrow H \end{aligned} \right\}$$

$$-\mathcal{L}_{\text{int}} \supset \mu e^{i\frac{a}{f_a}} \Phi^T i\tau_2 \Delta^\dagger \Phi + \text{h.c.}..$$

**After electroweak
symmetry breaking**

$$-\mathcal{L}_{\text{int}} \supset \frac{\mu v_\Phi^2 v_\Delta}{\sqrt{2}} \cos\left(\frac{a}{f_a}\right)$$

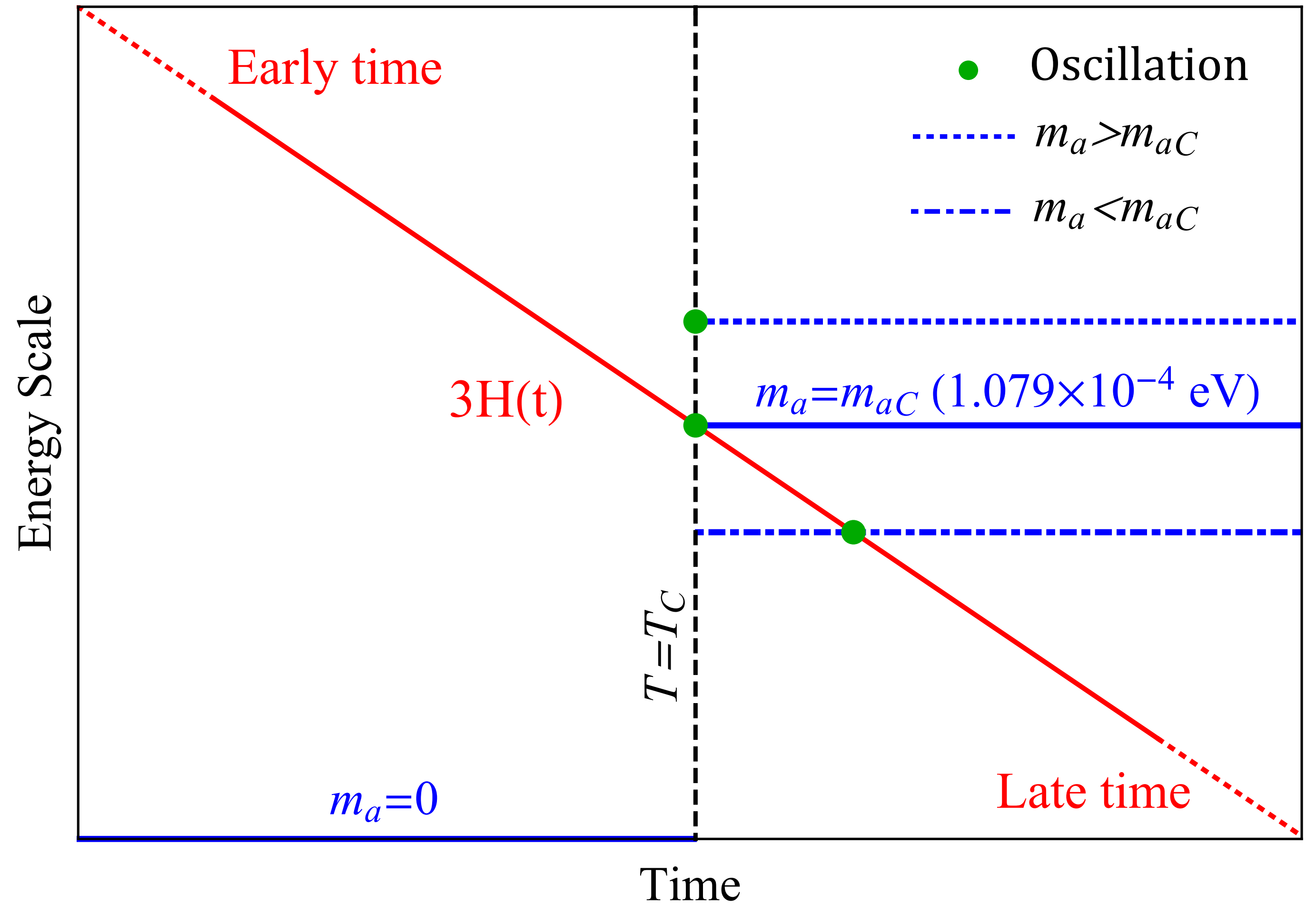
Non-zero Majoron mass

ALP oscillation time is relevant to EW scale !

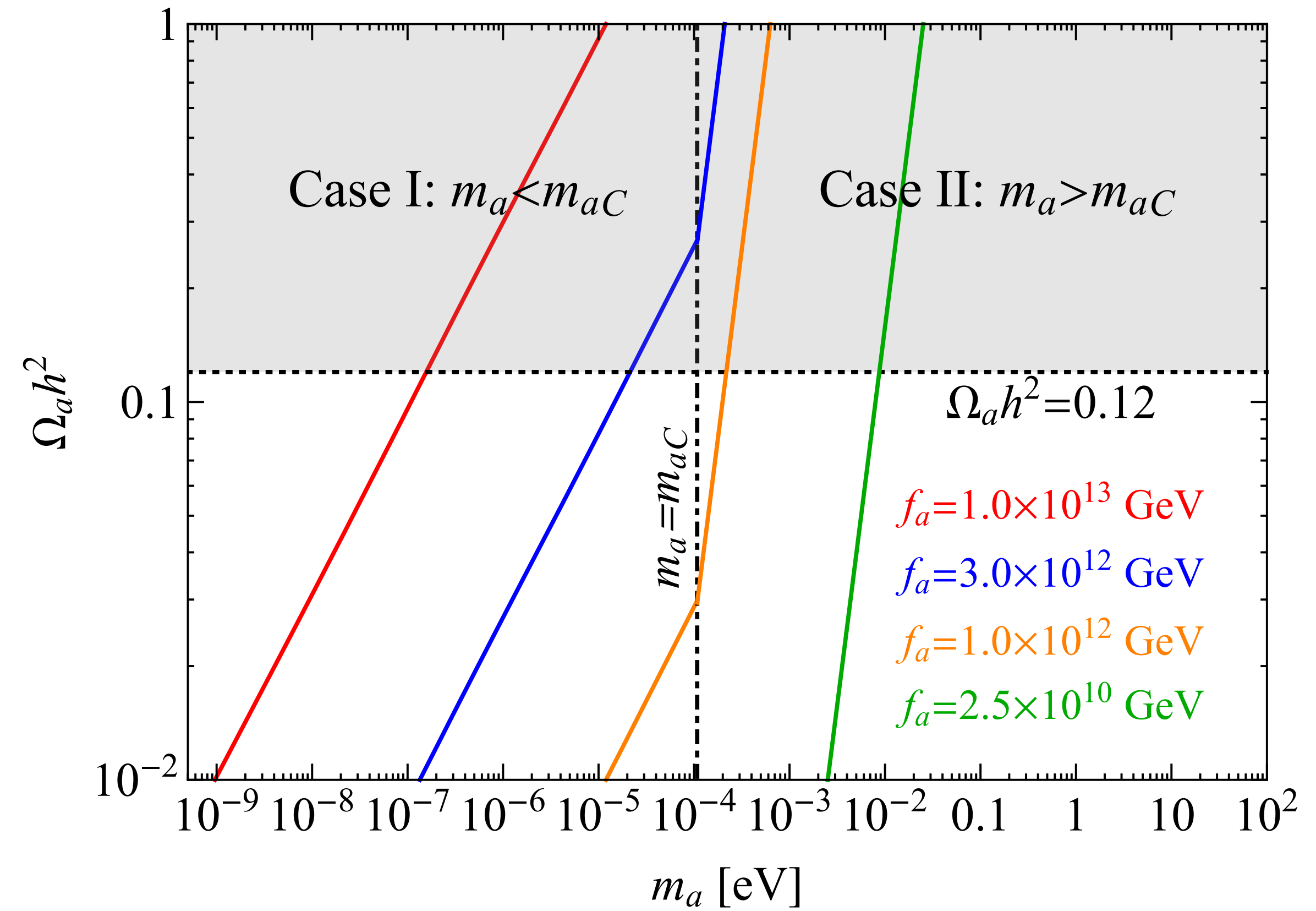
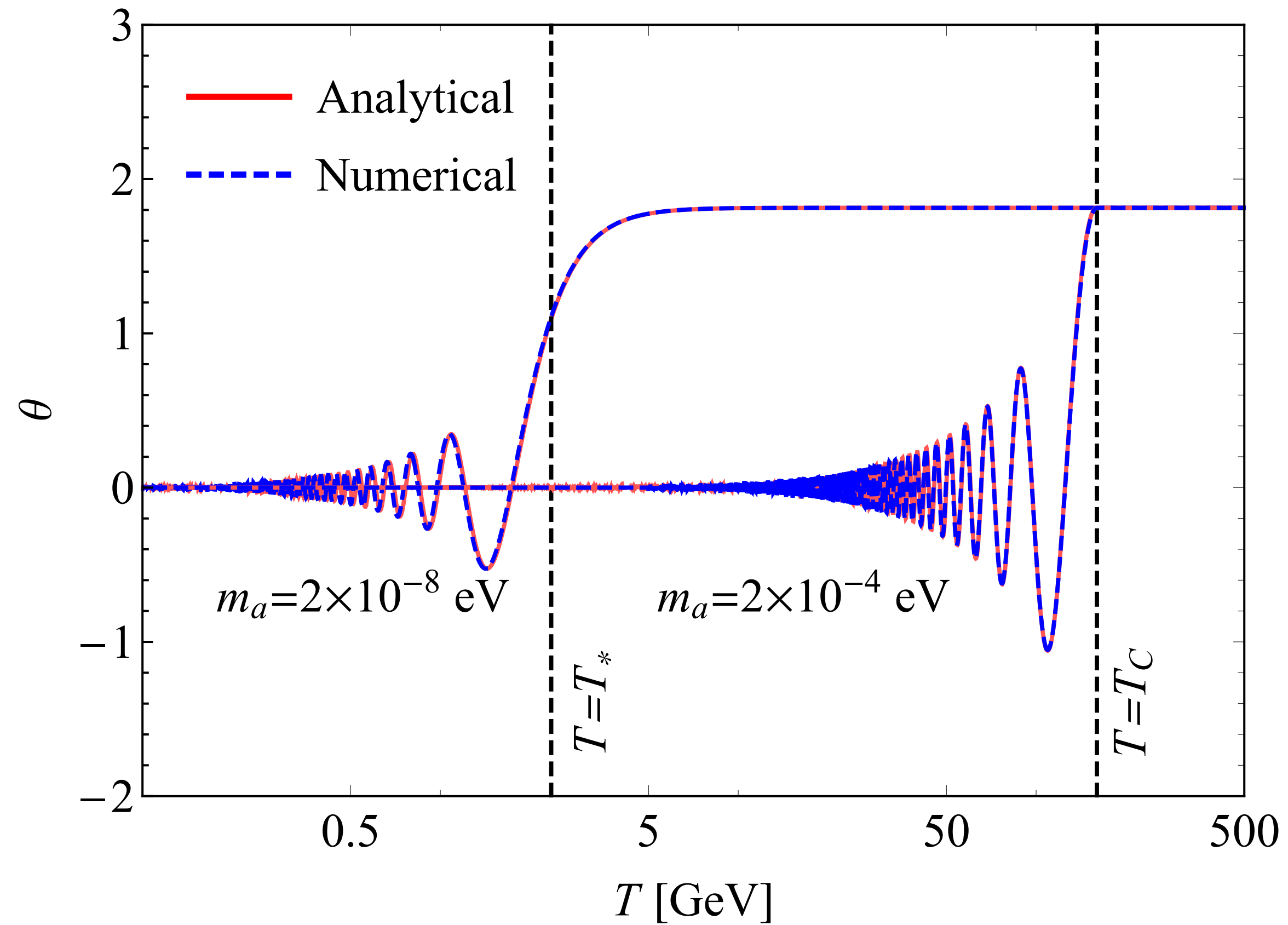
$$m_a^2(T) = \begin{cases} \frac{\mu v_\phi^2(T) v_\Delta(T)}{\sqrt{2} f_a^2}, & T \leq T_C \\ 0, & T > T_C \end{cases}$$

$$T_{\text{osc}} = \begin{cases} T_*, & m_a < m_{aC} \\ T_C, & m_a \geq m_{aC} \end{cases}$$

$$m_{aC} = 1.079 \times 10^{-4} \text{ eV}$$



ALP Relic Density



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Magnetic field in the early universe

Mechanisms for the production of primordial magnetic field

- PMF can be generated during the pseudo-scalar inflation, ϕ , with the Chern-Simons interaction, $\frac{\phi}{4\pi f_a} F \widetilde{F}$.
- PMF can be produced from the axion oscillations.
- PMF can be generated from phase transitions in the early universe.
- PMF can be generated from graviton.

Redshifted as $(1+z)^{-2}$

Axion dynamics in the presence of the PMF

Action

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{2} (\partial_\mu \phi) (\partial^\mu \phi) - V(\phi) - \frac{g_{\phi\gamma}}{4} \phi F_{\mu\nu} \widetilde{F}^{\mu\nu} \right]$$

Equation of motion

$$\partial_\eta^2 \phi + 2\mathcal{H} \partial_\eta \phi + a^2 m_\phi^2 \phi = g_{\phi\gamma} a^{-2} \mathbf{E}^* \cdot \mathbf{B}^*$$

Helicity

$$h^* = \lim_{V^* \rightarrow \infty} \frac{1}{V^*} \int d^3x \mathbf{A}^* \cdot \mathbf{B}^* ,$$

Evolution EQ

$$\frac{\partial h^*}{\partial \eta} = \lim_{V^* \rightarrow \infty} \frac{1}{V^*} \int d^3x (-2 \mathbf{E}^* \cdot \mathbf{B}^*) .$$

Crucial for solving
EOM of the axion!

Chiral Magnetic Effect

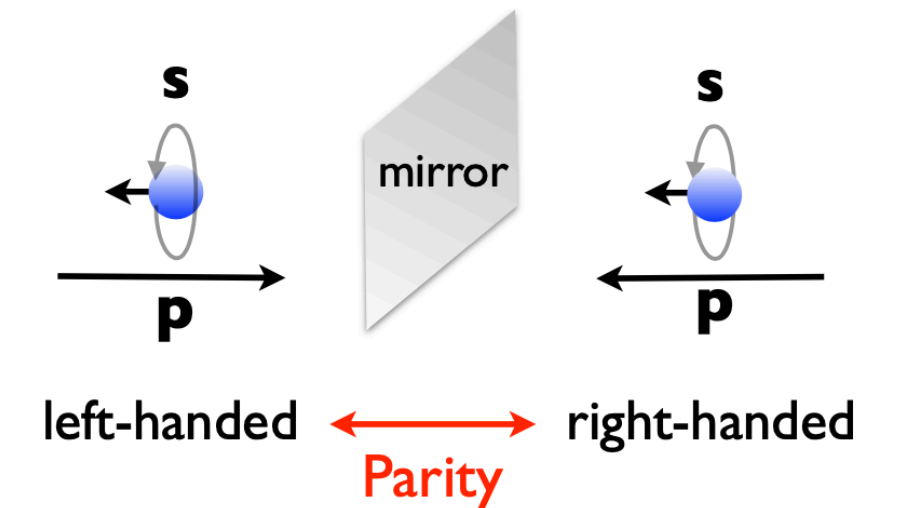
CME: When a medium with a chiral asymmetry is exposed to a magnetic field there is an induced electric current.

- Electric current induced by magnetic field in usual media? **NO!**

$$P : j \rightarrow -j \quad E \rightarrow -E \quad B \rightarrow B$$

Chirality of fermions

In the presence of a parity odd quantity, electric current can be induced!

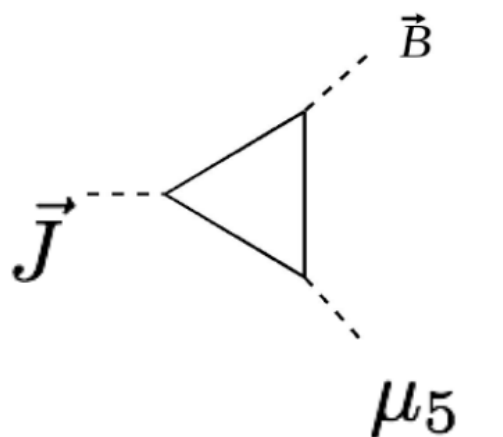


Chiral chemical potential can be generated due to topological transition.

$$\mu_5 \neq 0$$

In the presence of magnetic field, electric current is generated

$$\partial_\mu j^\mu = \frac{e^2}{32\pi^2} F \widetilde{F} \rightarrow \vec{J} = \frac{e^2}{2\pi^2} \mu_5 \vec{B}$$



Chiral Magnetic Effect

@finite T: Lagrangian/Action for hyper magnetic field in the presence of chemical potential

$$-\mathcal{L} \sim \frac{1}{4} Y_{\mu\nu} Y^{\mu\nu} + C_{EW} n_{CS}^{EW} + C_Y N_{CS}^Y + j^\mu A_\mu + \dots$$

$$\mu_5 \sim \sum_{i=1}^3 \left[-2\mu_{R_i} + \mu_{L_i} - \frac{2}{3}\mu_{d_i} - \frac{8}{3}\mu_{u_{R_i}} + \frac{1}{3}\mu_{Q_i} \right] \quad N_{CS}^Y = \frac{g'^2}{32\pi^2} 2A \cdot B_Y$$

• Maxwell equations

$$\frac{\partial E}{\partial \eta} - \nabla \times B + J = 0 \quad \frac{\partial B}{\partial \eta} + \nabla \times E = 0 \quad \nabla \cdot E = \rho \quad \nabla \cdot B = 0$$

$$J = \sigma(E + v \times B) + \frac{2\alpha}{\pi} \mu_5 B - g_{\phi\gamma} B \partial_\eta \phi$$

The generalized Ohm's law

EOM for helicity and the energy density

• EOM

$$\frac{\partial h^*}{\partial \eta} = \lim_{V^* \rightarrow \infty} \frac{1}{V^*} \int d^3x \frac{2}{\sigma^*} (\nabla^2 A^*) \cdot B^* + \frac{1}{\sigma^*} \left(\frac{8\alpha}{\pi} \mu_5^* - 4g_{\phi\gamma} \partial_\eta \phi \right) \rho_B^*,$$

$$\frac{\partial \rho_B^*}{\partial \eta} = \lim_{V^* \rightarrow \infty} -\frac{1}{V^*} \frac{1}{\sigma^*} \int d^3x \mathbf{B}^* \cdot \left[\nabla^2 \mathbf{B}^* + \nabla \times \left(-\frac{2\alpha}{\pi} \mu_5^* \mathbf{B}^* + g_{\phi\gamma} \mathbf{B}^* \partial_\eta \phi \right) \right]$$

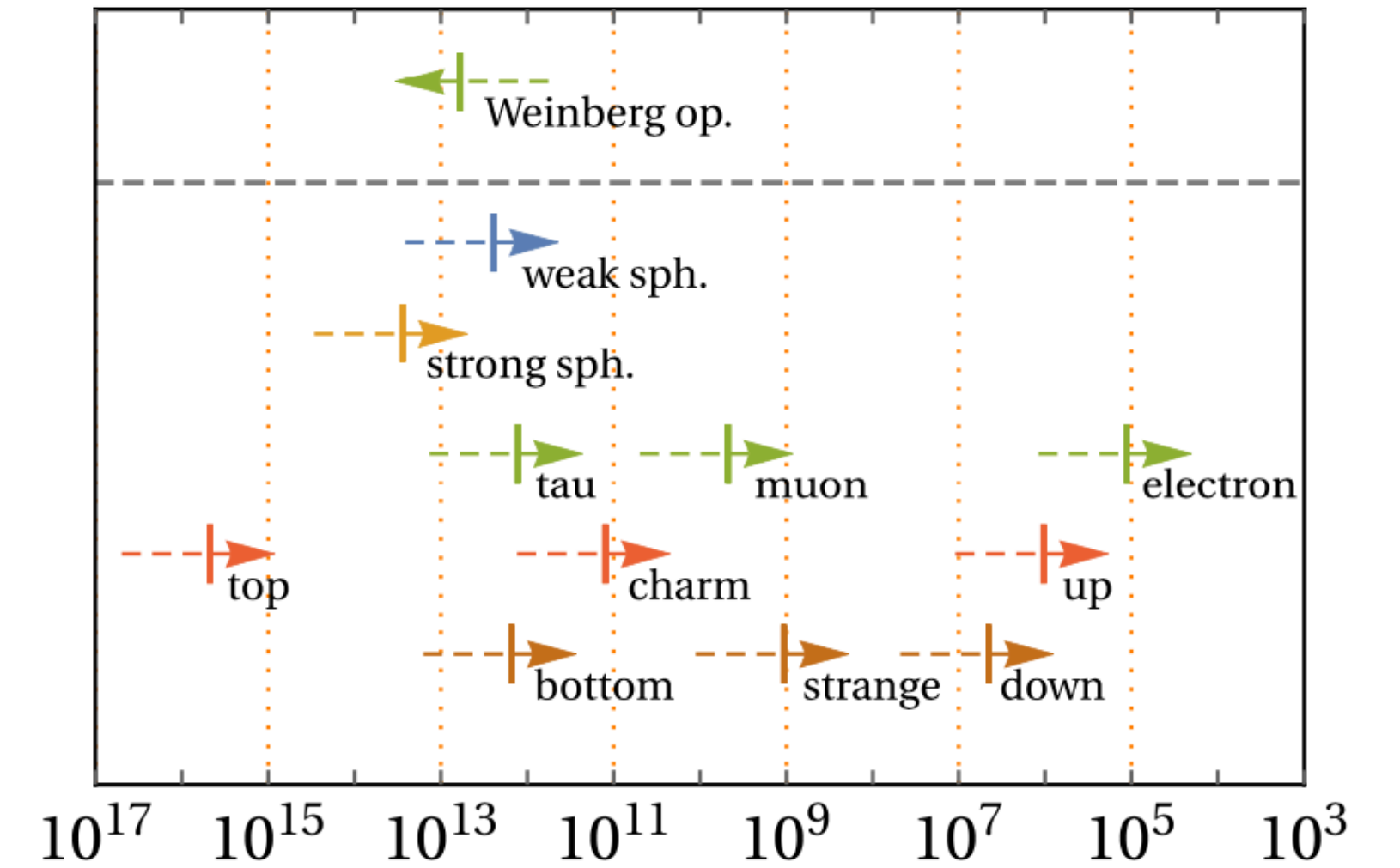
• Further define

$$h = \int \frac{d^3x}{V} A \cdot B = \int dk h_k \quad \rho_B = \int \frac{d^3x}{V} \frac{1}{2} B^2 = \int dk \rho_k$$

$$\frac{\partial h_k^*}{\partial \eta} = -\frac{2(k^*)^2}{\sigma^*} h_k^* + \frac{1}{\sigma^*} \left(\frac{8\alpha}{\pi} \mu_5^* - 4g_{\phi\gamma} \partial_\eta \phi \right) \rho_{B,k}^* \quad \frac{\partial \rho_{B,k}^*}{\partial \eta} = -\frac{2(k^*)^2}{\sigma^*} \rho_{B,k}^* + \frac{1}{\sigma^*} \left(\frac{2\alpha}{\pi} \mu_5^* - g_{\phi\gamma} \partial_\eta \phi \right) (k^*)^2 h_k^*$$

EOM for chemical potentials

Interaction	Weinberg	WS	SS	Y_e	Y_μ	Y_τ
Γ_α/T^4	$\kappa_W \frac{m_\nu^2 T^2}{v_{EW}^4}$	$\frac{1}{2} \kappa_{WS} \alpha_2^5$	$\frac{1}{2} \kappa_{SS} \alpha_3^5$	$\kappa_{Y_e} y_e^2$	$\kappa_{Y_\mu} y_\mu^2$	$\kappa_{Y_\tau} y_\tau^2$
T_α [GeV]	6.0×10^{12}	2.5×10^{12}	2.8×10^{13}	1.1×10^5	4.7×10^9	1.3×10^{12}
Interaction	Y_u	Y_c	Y_t	Y_d	Y_s	Y_b
Γ_α/T^4	$\kappa_{Y_u} y_u^2$	$\kappa_{Y_c} y_c^2$	$\kappa_{Y_t} y_t^2$	$\kappa_{Y_d} y_d^2$	$\kappa_{Y_s} y_s^2$	$\kappa_{Y_b} y_b^2$
T_α [GeV]	1.0×10^6	1.2×10^{11}	4.7×10^{15}	4.5×10^6	1.1×10^9	1.5×10^{12}



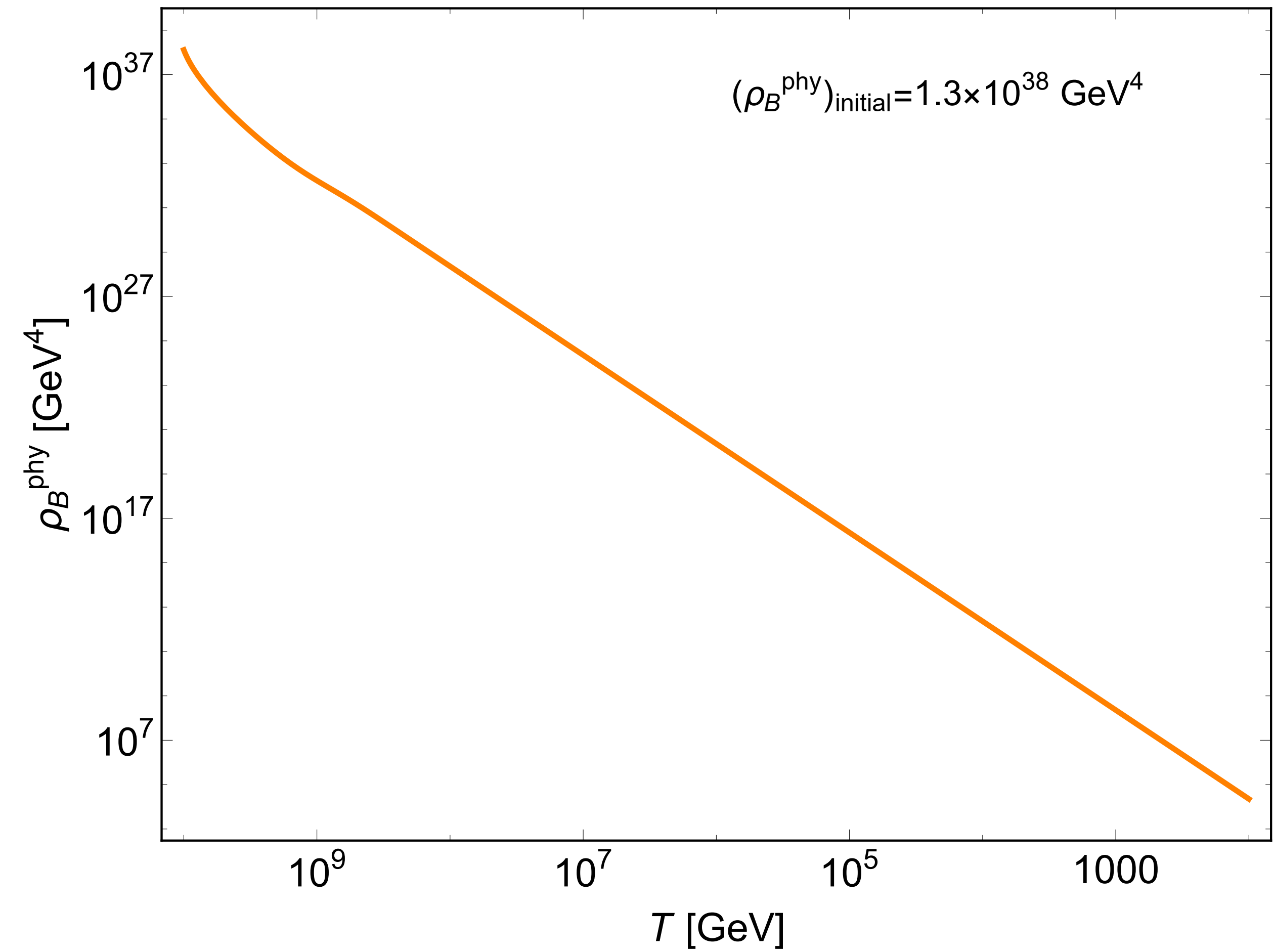
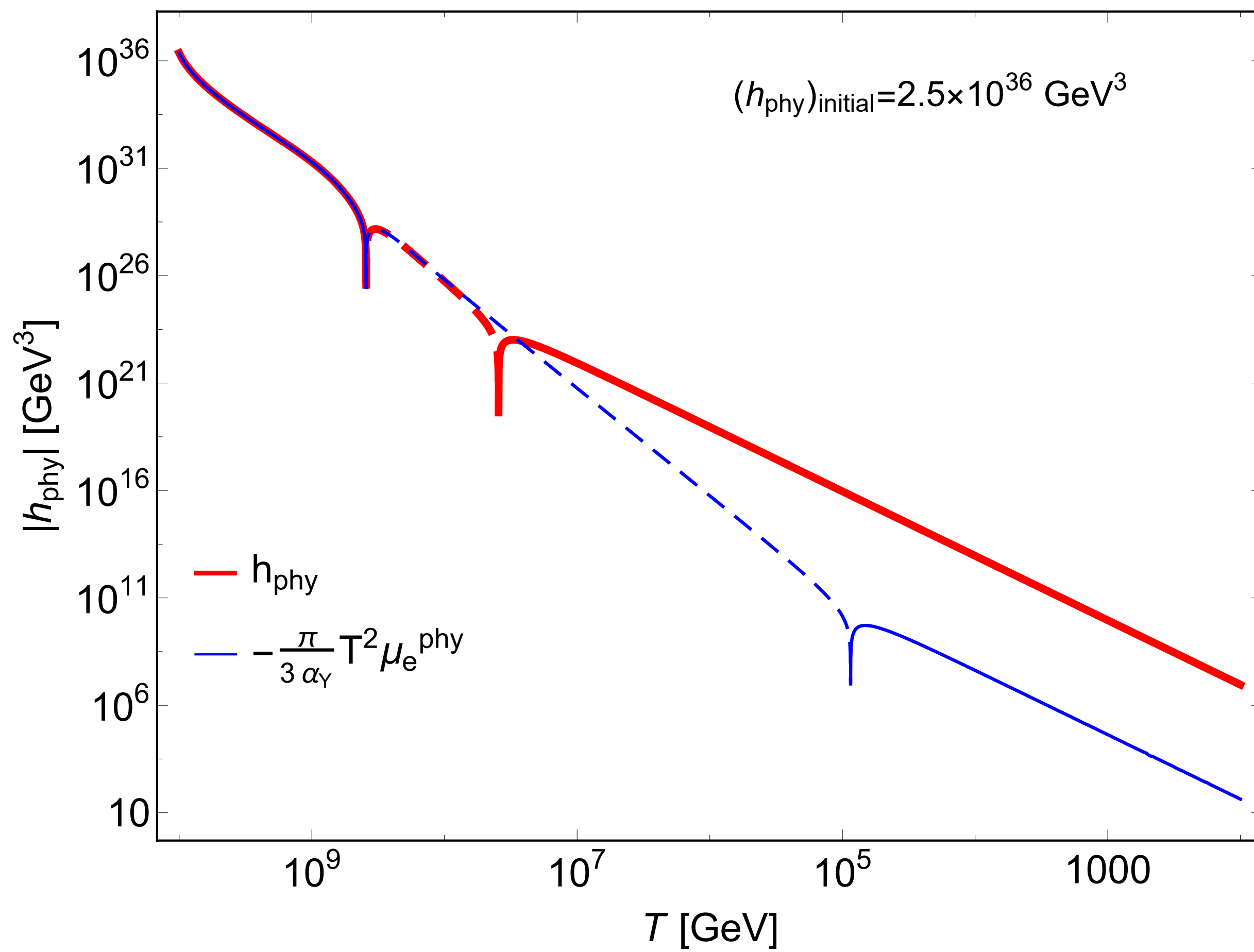
Valerie Domck, Yohei Ema, Kyohei Mukai and Masaki Yamada, JHEP 08(2020)096

- **Toy scenario: Assuming all interactions except right-handed electron are in thermal equilibrium**

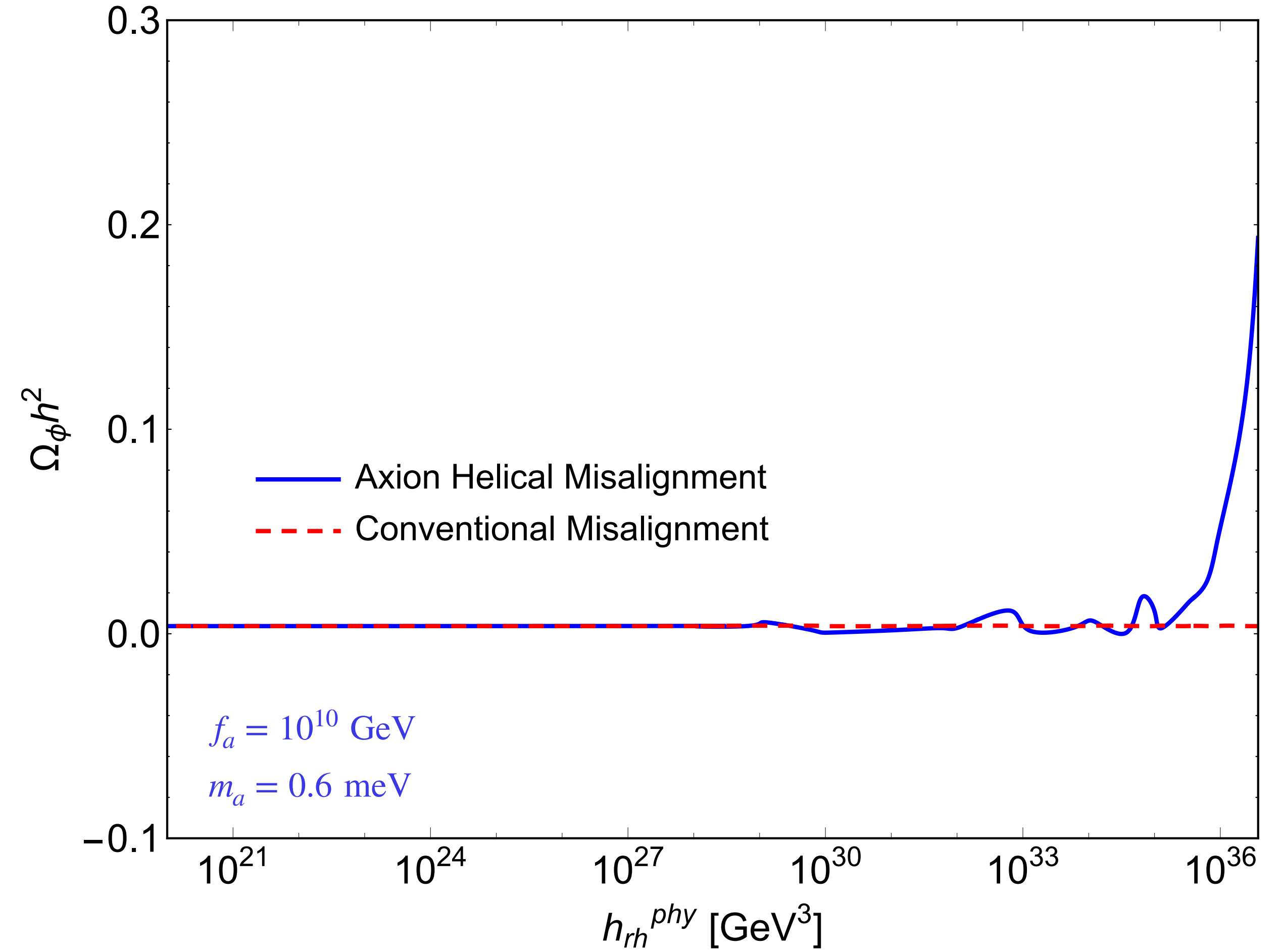
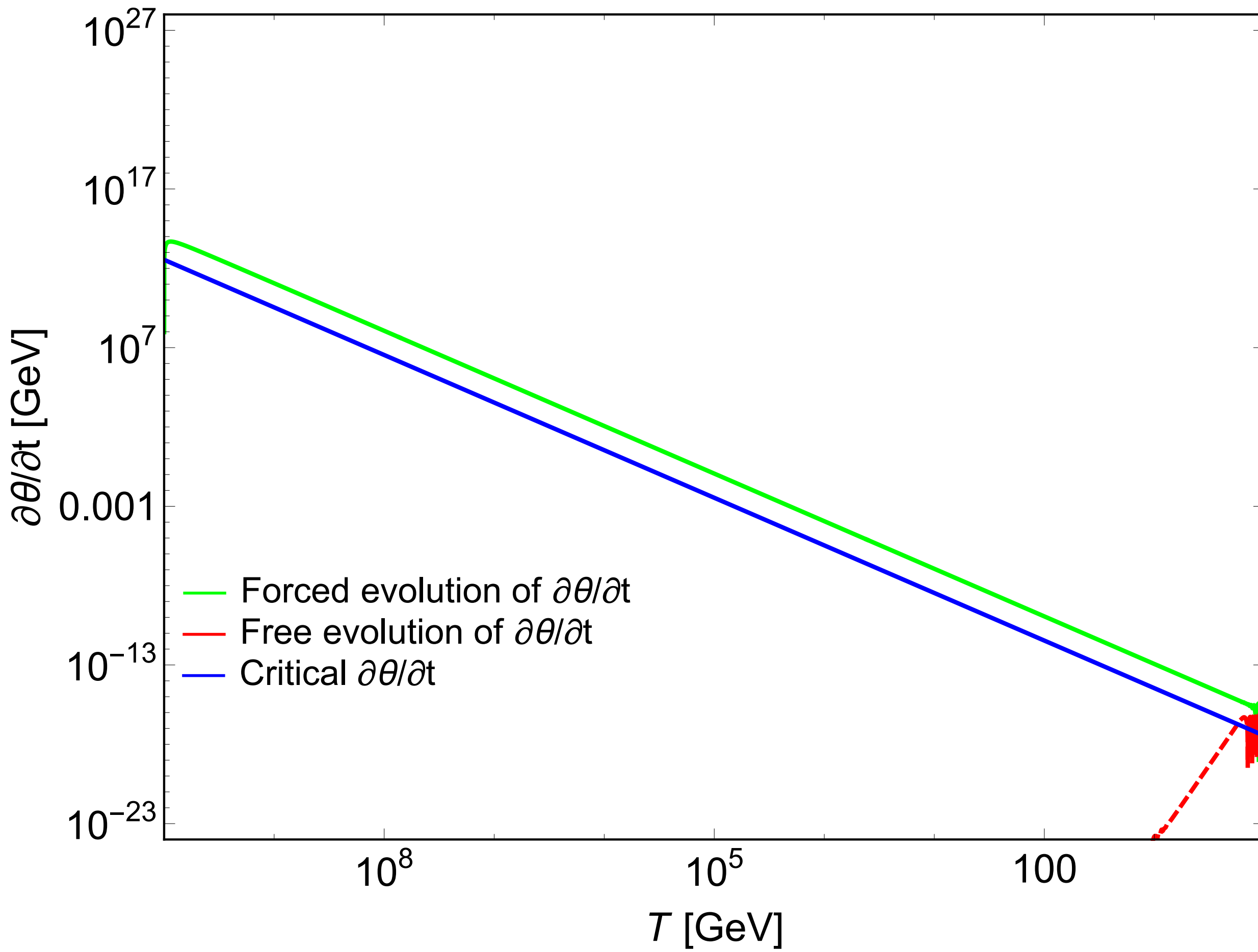
$$\mu_5^* = \frac{711}{481} \mu_e^*,$$

$$\frac{\partial \mu_e^*}{\partial \eta} = - \frac{3}{(T^*)^2} \frac{\alpha}{\pi} \frac{\partial h}{\partial \eta} - \Gamma_{Y_e} \frac{711}{481} \mu_e^*$$

Numerical results



Numerical results



Conclusion

Axion misalignment mechanism is crucial for evaluation of the axion relic abundance. There are still several effects need to be clarified for this mechanism:

- Axion misalignment mechanism is relevant to mass generation mechanism of the axion.
- Axion misalignment mechanism is relevant to the evolution of the PMF and Chiral asymmetries.

Thank you for your attention!