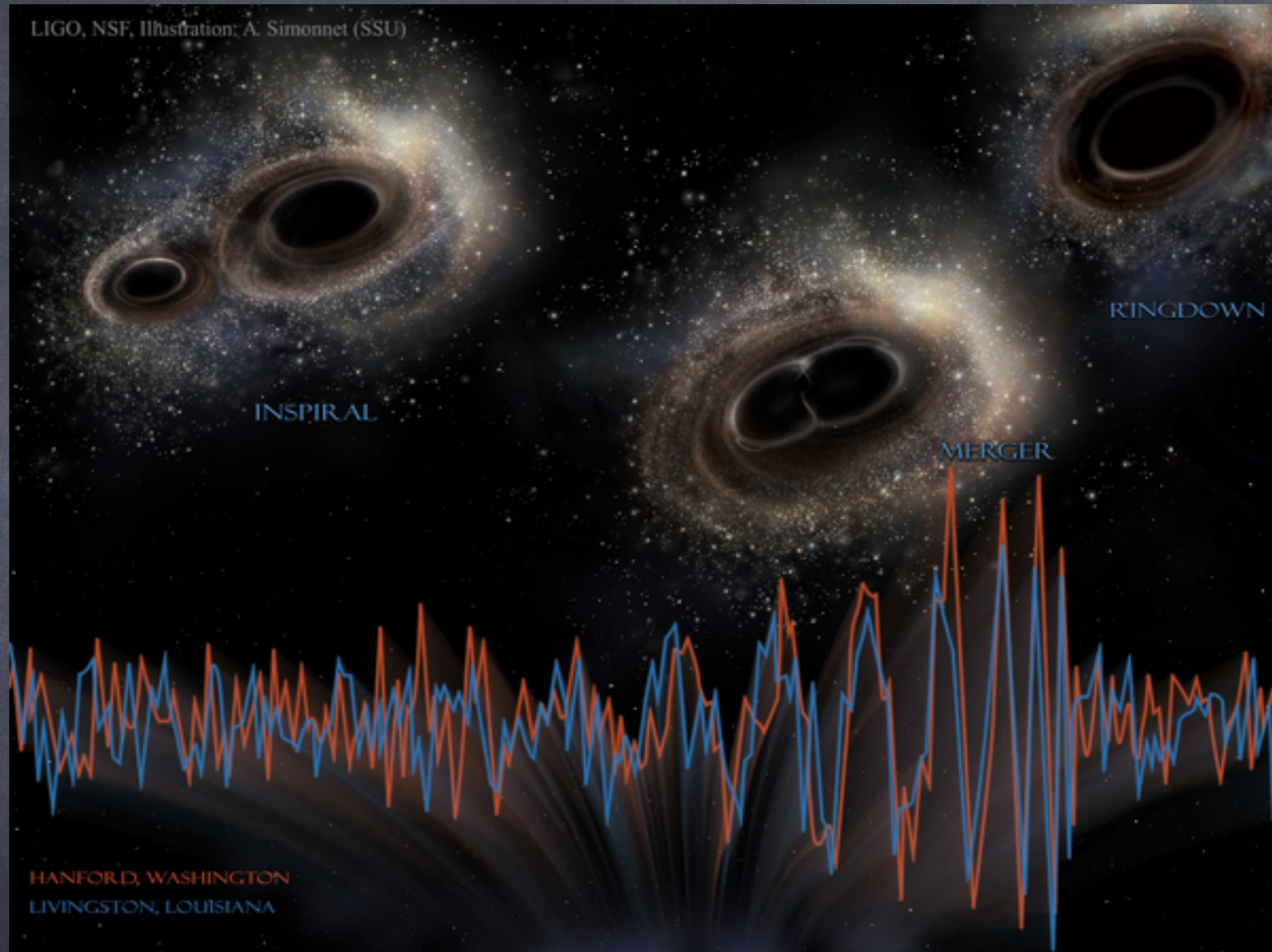


Prominence: A discriminator of  
gravitational wave signals

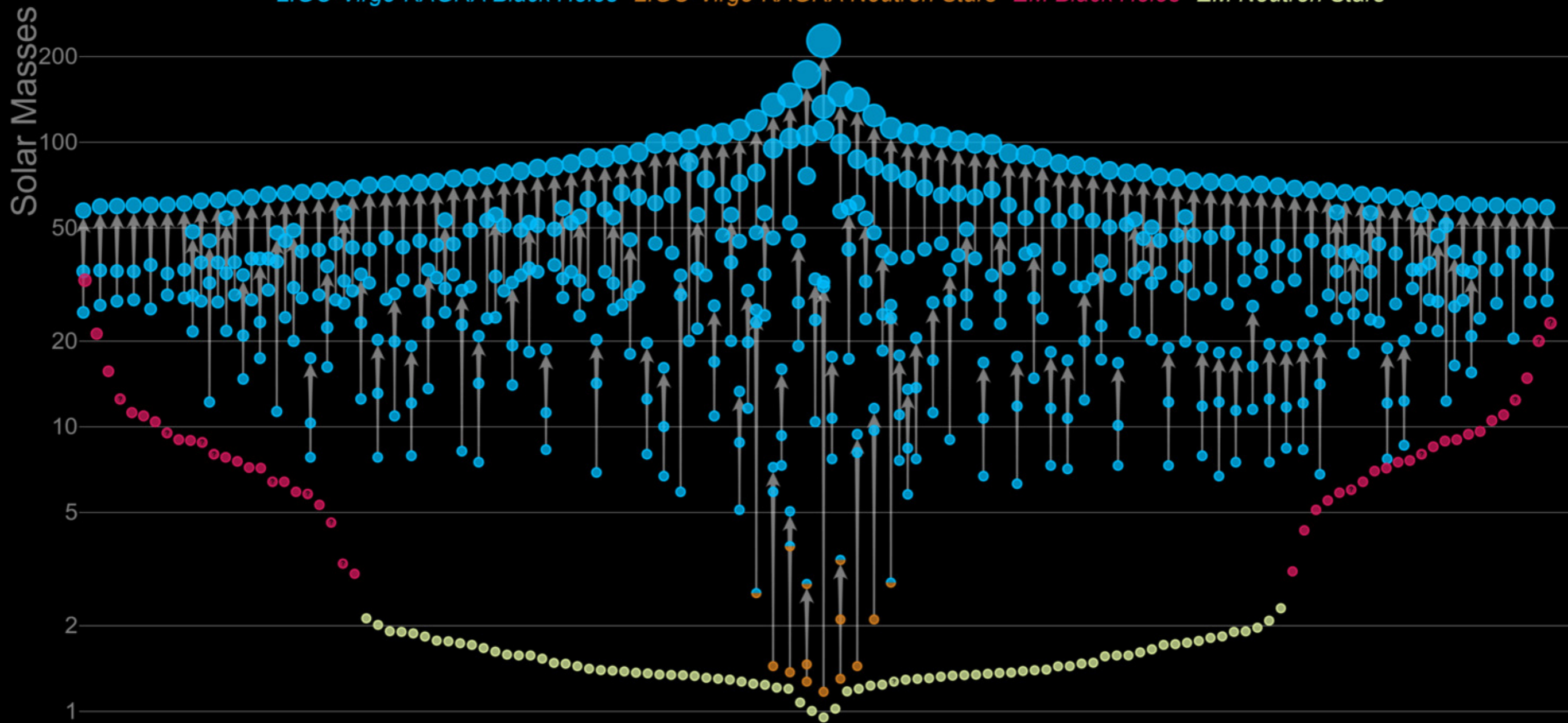
# Dawn of gravitational wave astronomy



- Detection of GW150914 on September 14, 2015 by LIGO
- First observation of heavy black hole binaries

# Masses in the Stellar Graveyard as of 09/25

*LIGO-Virgo-KAGRA Black Holes* *LIGO-Virgo-KAGRA Neutron Stars* *EM Black Holes* *EM Neutron Stars*



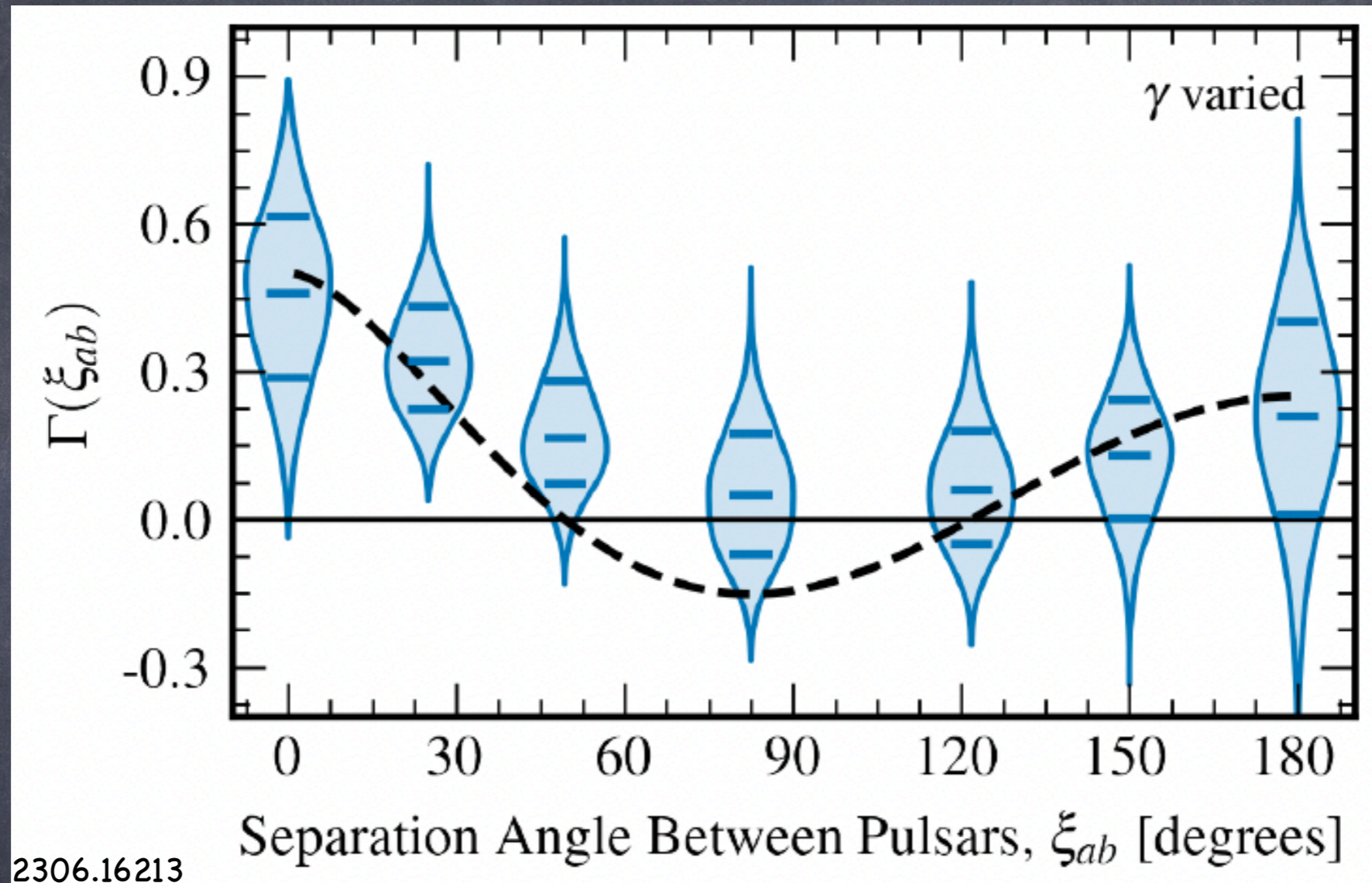
# Pulsar Timing Arrays



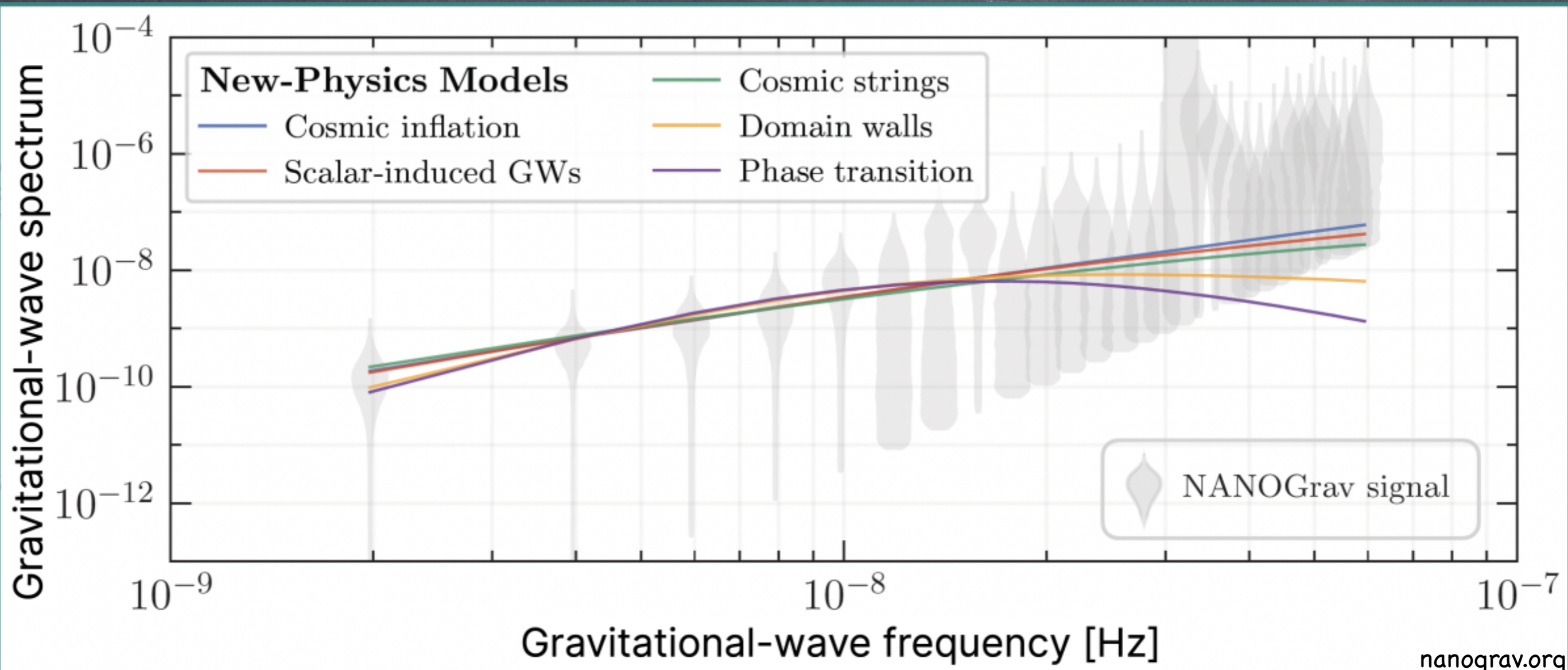
- Galactic size GW detector made of millisecond pulsar array
- Precise rotation periods and radio pulses from poles of pulsars make them ultra-precise clocks

- GW perturbs spacetime between the pulsar and Earth and changes the time of arrival (TOA) of pulses
- Measure residuals in TOA:  $R^a = TOA_{measured}^a - TOA_{model}^a$
- Cross-correlate timing residuals of pairs of pulsars separated by angle  $\xi_{ab}$
- Sensitive to frequencies between  $1/(\text{total observation time})$  and  $1/\text{cadence}$  i.e., [1/years, 1/weeks]

# Hellings-Downs curve

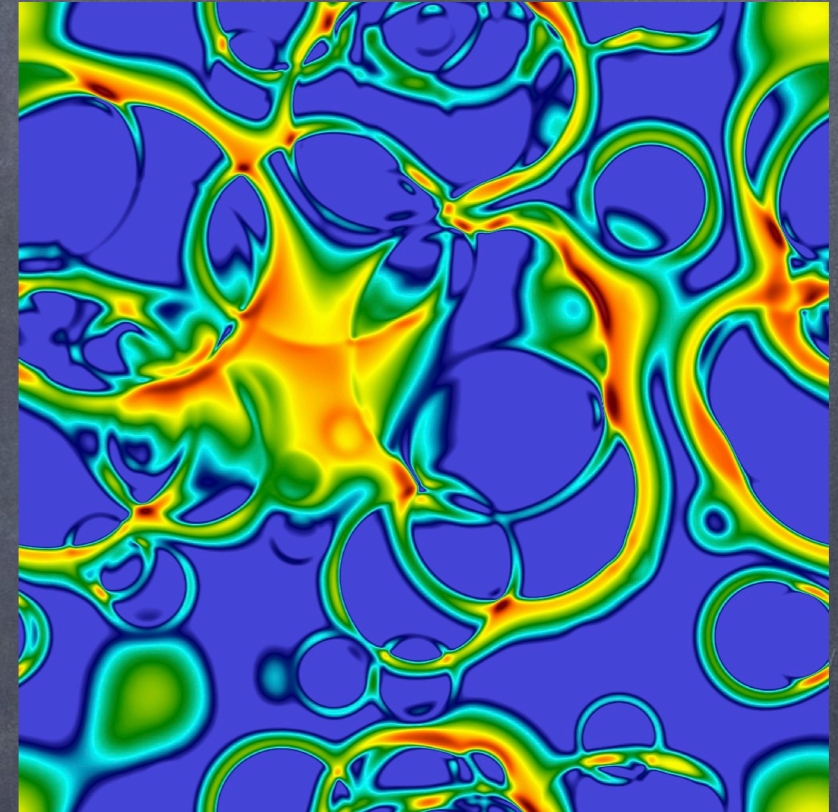


- PTAs see  $\sim 3\sigma$  quadrupole correlation of timing residuals
- Smoking gun signal of stochastic GW background



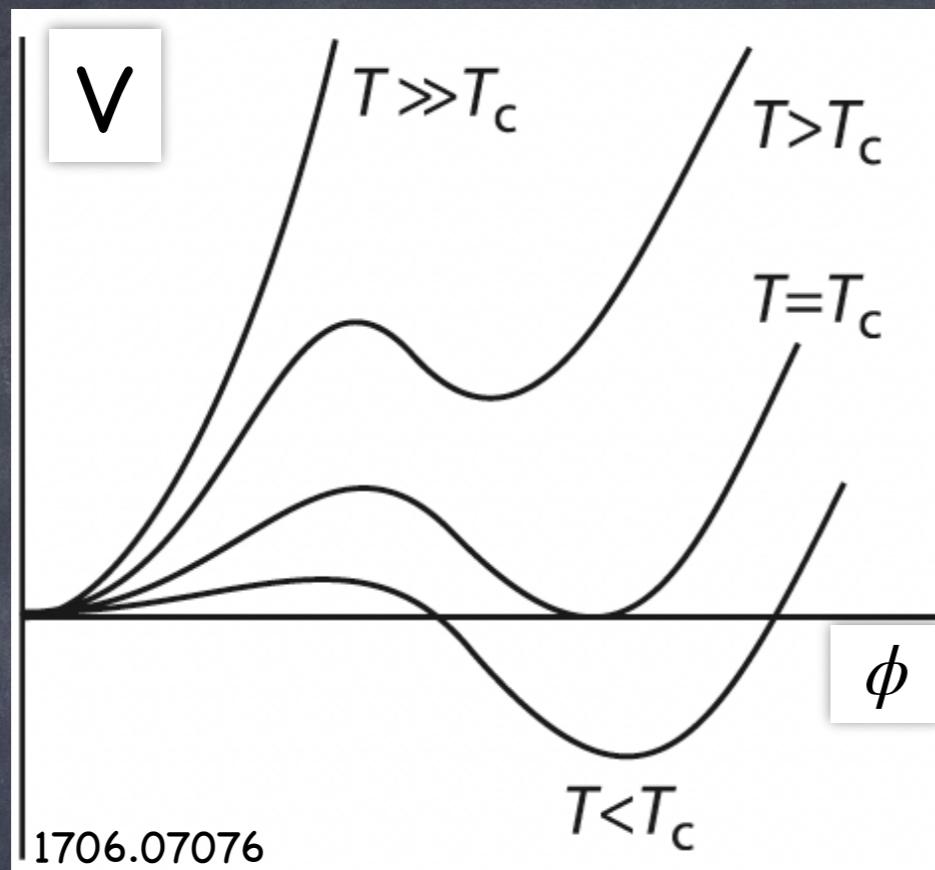
# Sources of gravitational waves

# Phase transitions

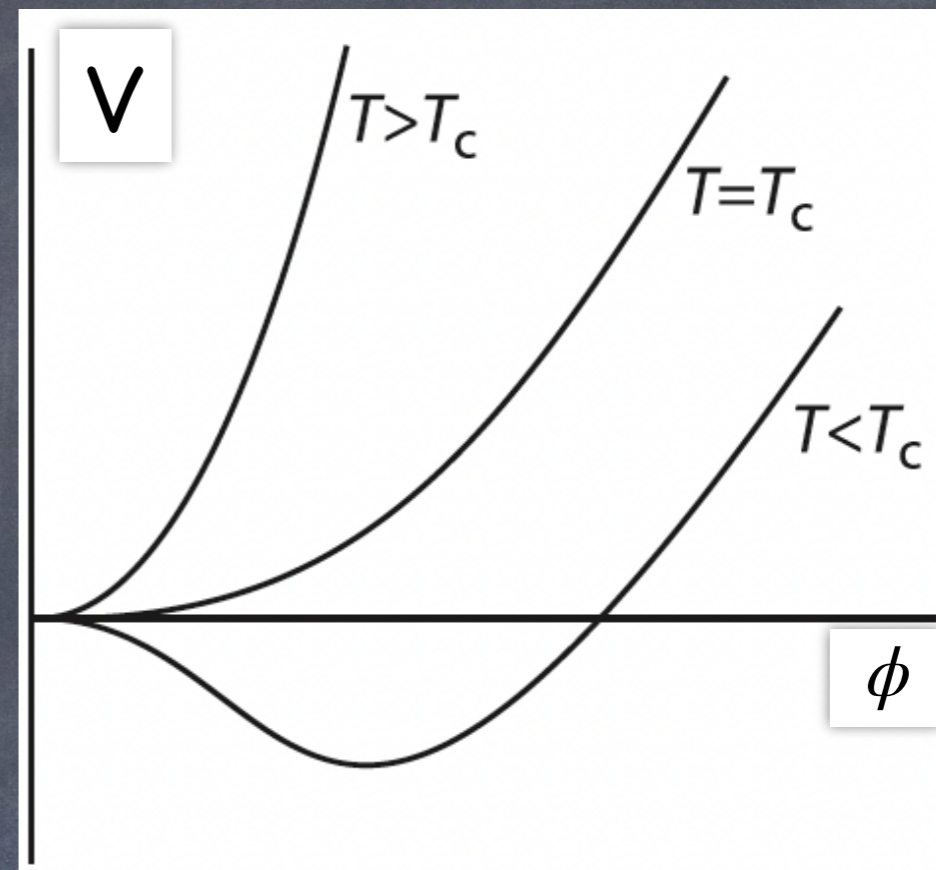


- Change in physical state (phase) of a system
- Change in vacuum structure as Universe cools
- Caused by spontaneous breaking of symmetries

## First order PT



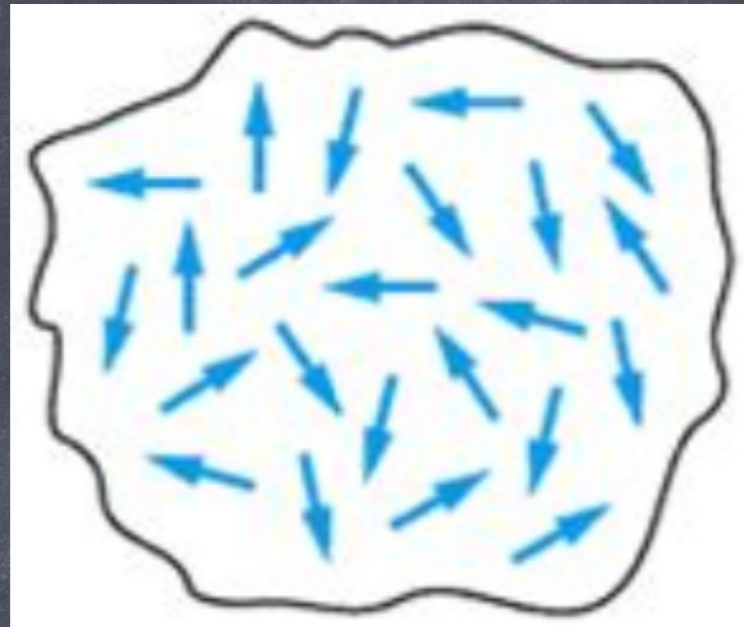
## Second order PT



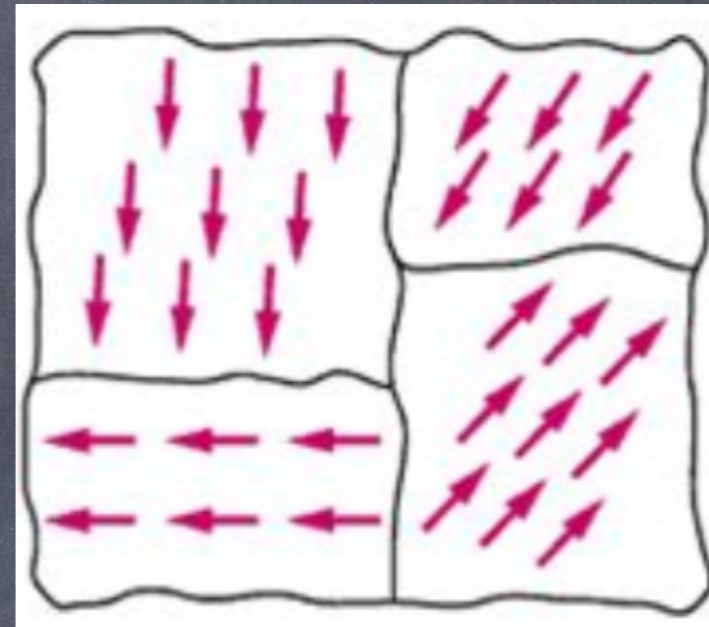
- FOPT associated with potential barrier and involves latent heat
- Thermal jump/quantum tunneling leads to nucleation of true vacuum bubbles
- Bubbles expand, transfer energy to plasma, collide thereby producing GWs

# Cosmic defects

Above  $T_c$

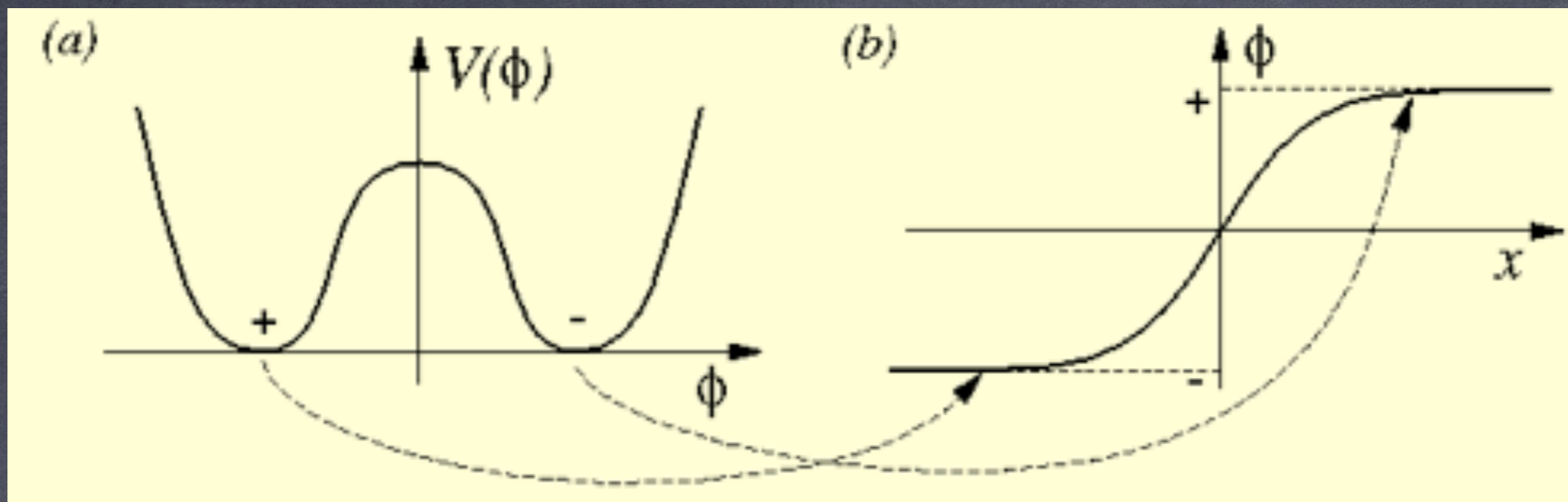


Below  $T_c$

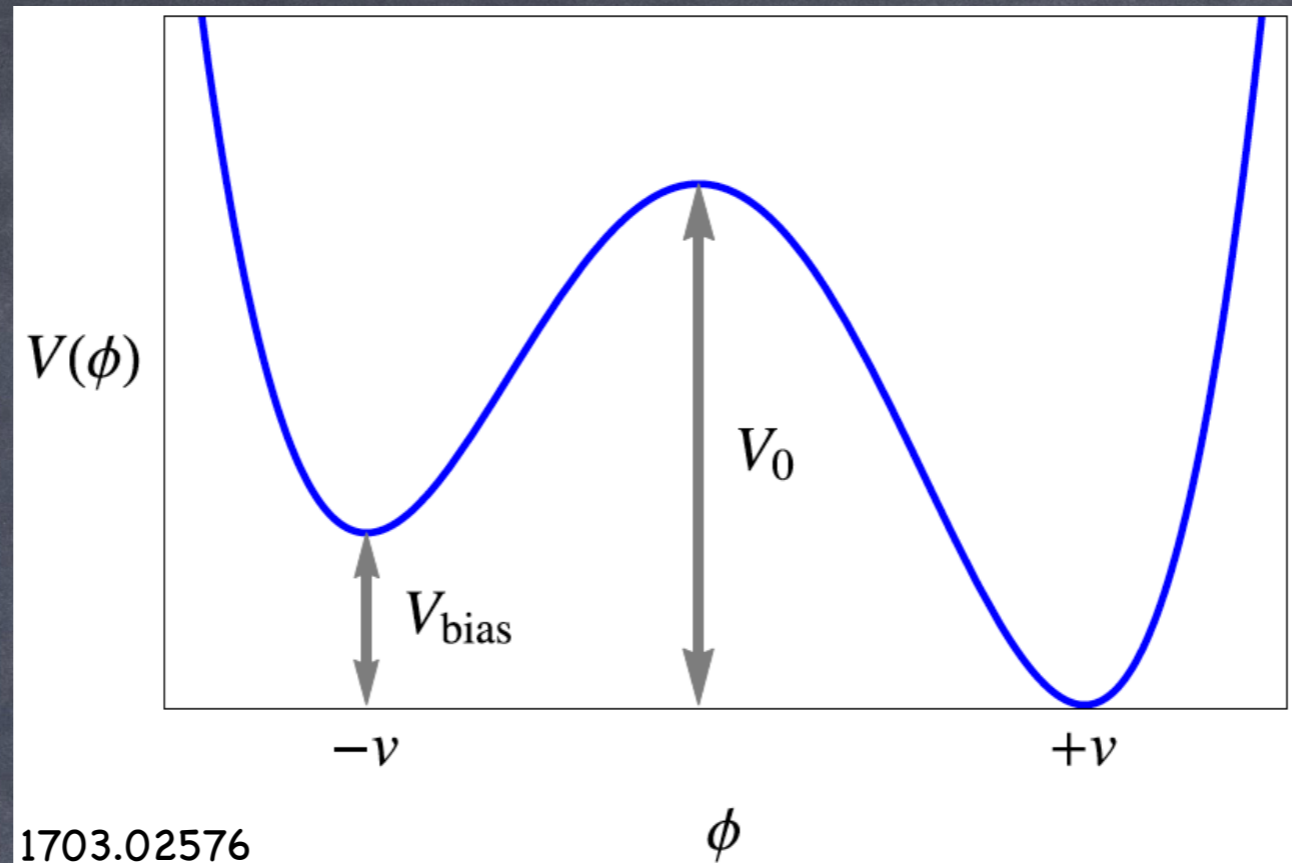


- As temperature is lowered below Curie temperature paramagnetic material becomes ferromagnetic
- Internal spin symmetry spontaneously broken
- Magnetic domains separated by stable domain walls

# Cosmic domain walls

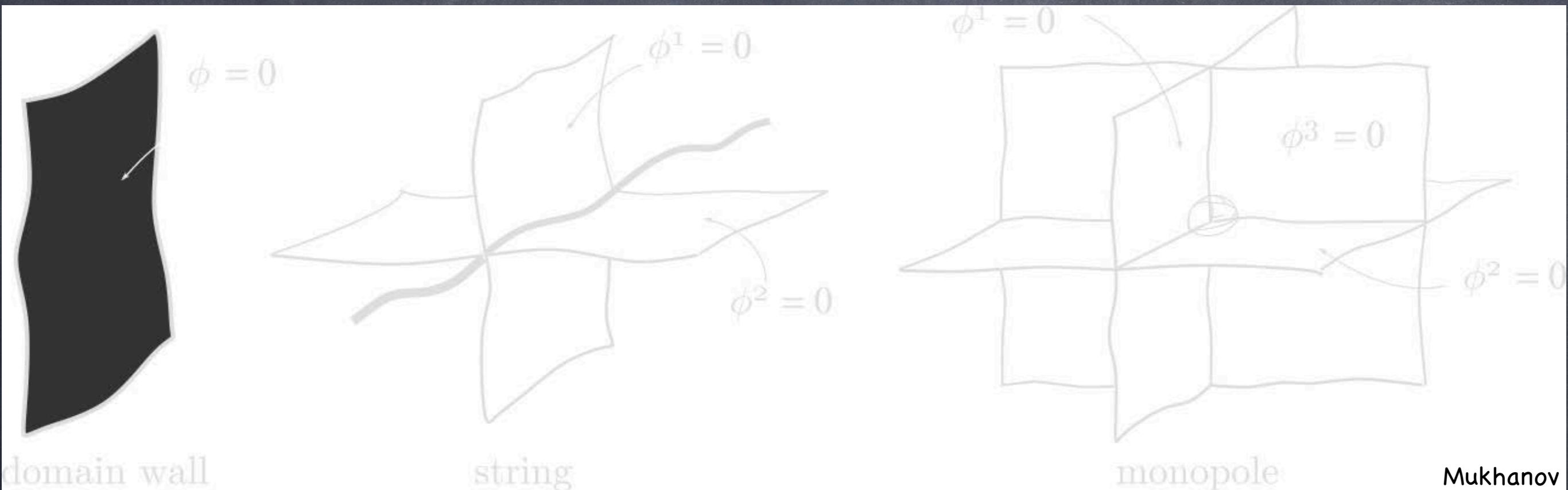


- Spontaneous breaking of  $\phi \rightarrow -\phi$  symmetry
- + and - domains separated by domain wall which has surface energy density (called surface tension)
- Symmetry unbroken at center of wall  $\phi = 0$
- Stable domain walls can dominate the energy density of the Universe since  $\rho_{walls} \sim 1/t \sim H \implies \Omega_{walls} \sim 1/H \sim t$



- Solution to domain wall problem is to break the symmetry
- Energy difference produces pressure  $\sim V_{\text{bias}}$  on the walls
- Collapsing domain walls produce GWs

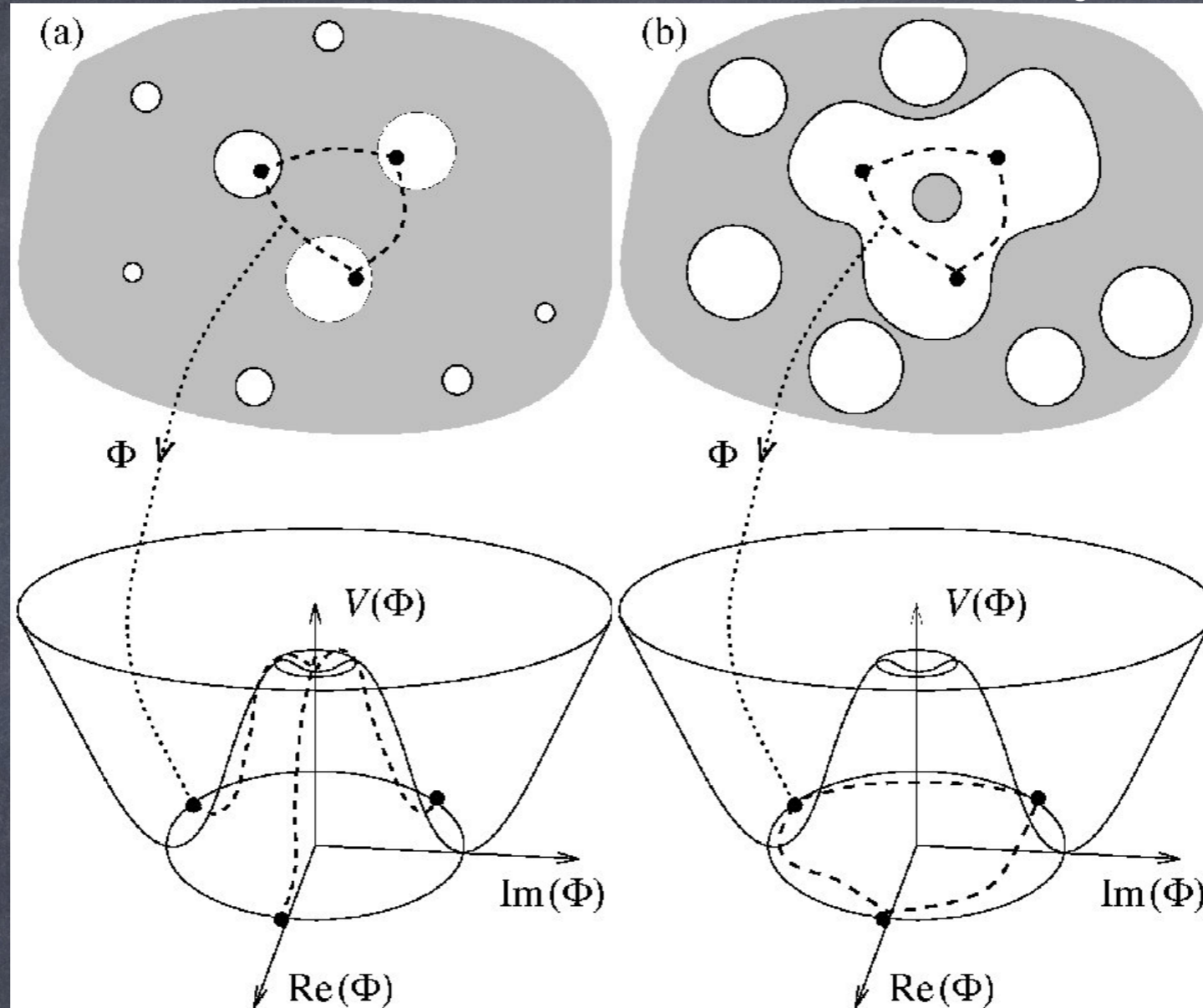
# Other cosmic defects



Spontaneous symmetry breaking in a higher dimensional field space

# Cosmic strings

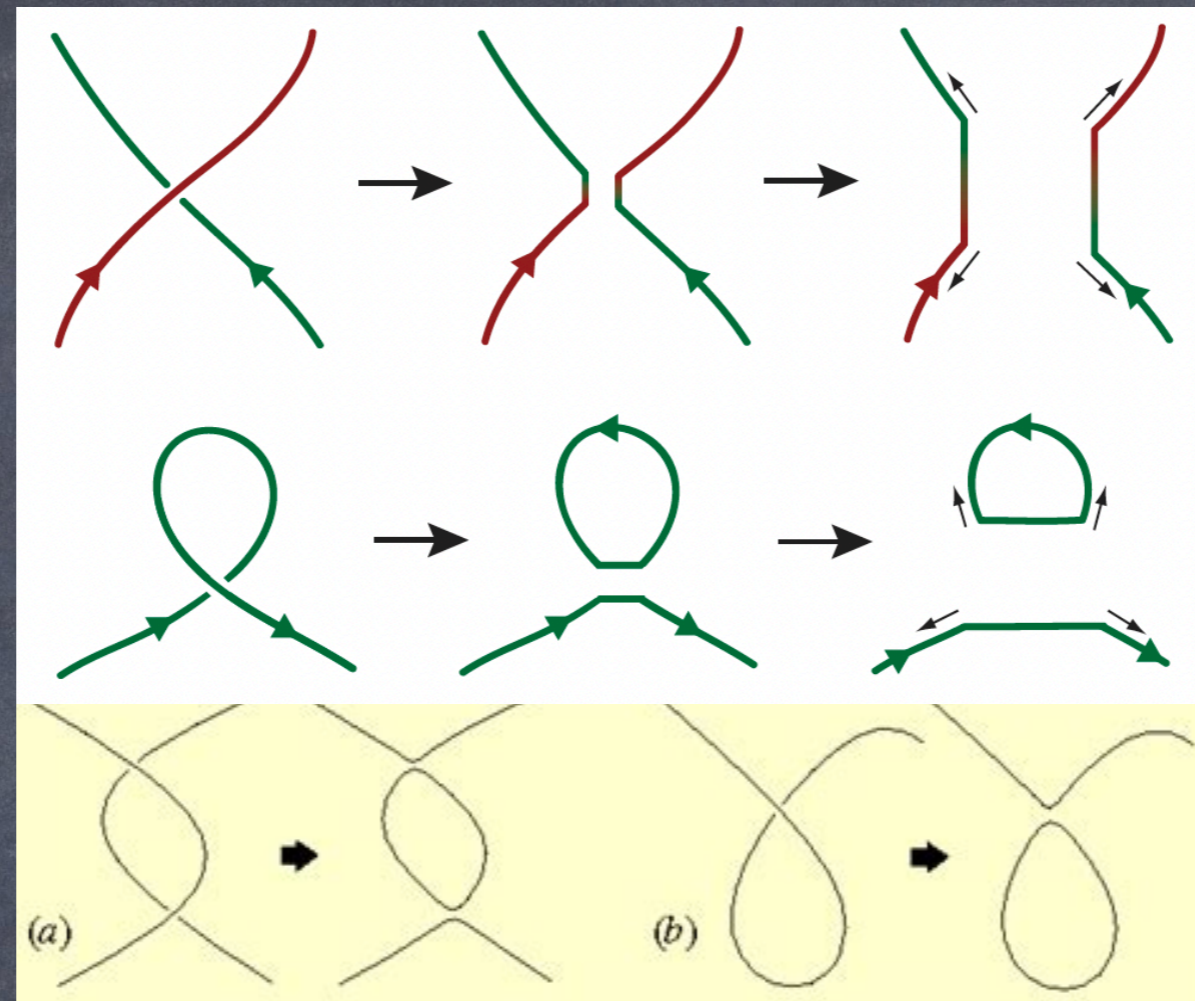
Jeong, Smoot



- Can have field theoretic origin, e.g., from spontaneous  $U(1)$  symmetry breaking
- Can be fundamental strings of superstring theory stretched to cosmic scales

# String interactions

0812.4020



- Strings collide and reconnect (with segments exchanged) with intercommutation probability
- Produce loops by self-interactions or pair-wise interactions
- Oscillating loops lose energy mainly by GW emission

## Scalar-induced GWs

- At linear order in perturbation theory, scalar and tensor modes evolve independently
- At second order these modes become coupled
- Large first-order scalar perturbations generate second-order tensor perturbations
- These tensor perturbations appear as GWs when scalar modes reenter the horizon at the end of inflation

Discriminating between GW sources

# Signal-to-Noise Ratio

$$\text{SNR} = \sqrt{T_{\text{obs}} \int_{f_{\text{min}}}^{f_{\text{max}}} df \frac{\Omega^2(f)}{\Omega_{\text{Sens}}^2(f)}}$$

- Used in most theoretical studies
- Spectral information of signal and background washed out because SNR is an integrated measure
- Would like observable that utilizes the shape of the signal to discriminate between sources

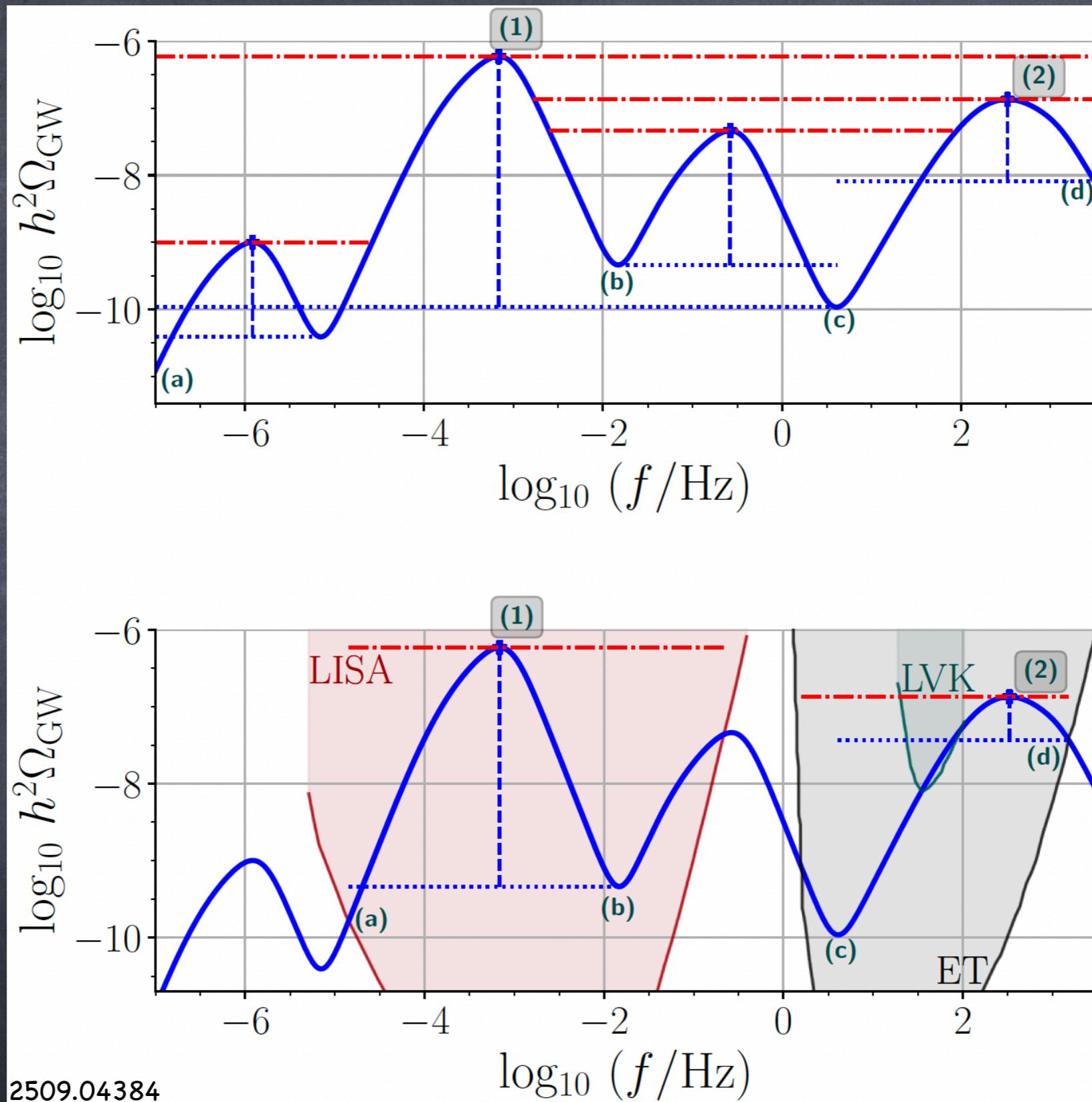
Prominence is a measure of how much a peak rises above the surrounding terrain

Four eight-thousanders



Lhotse has greater elevation than Makalu, but is much less prominent

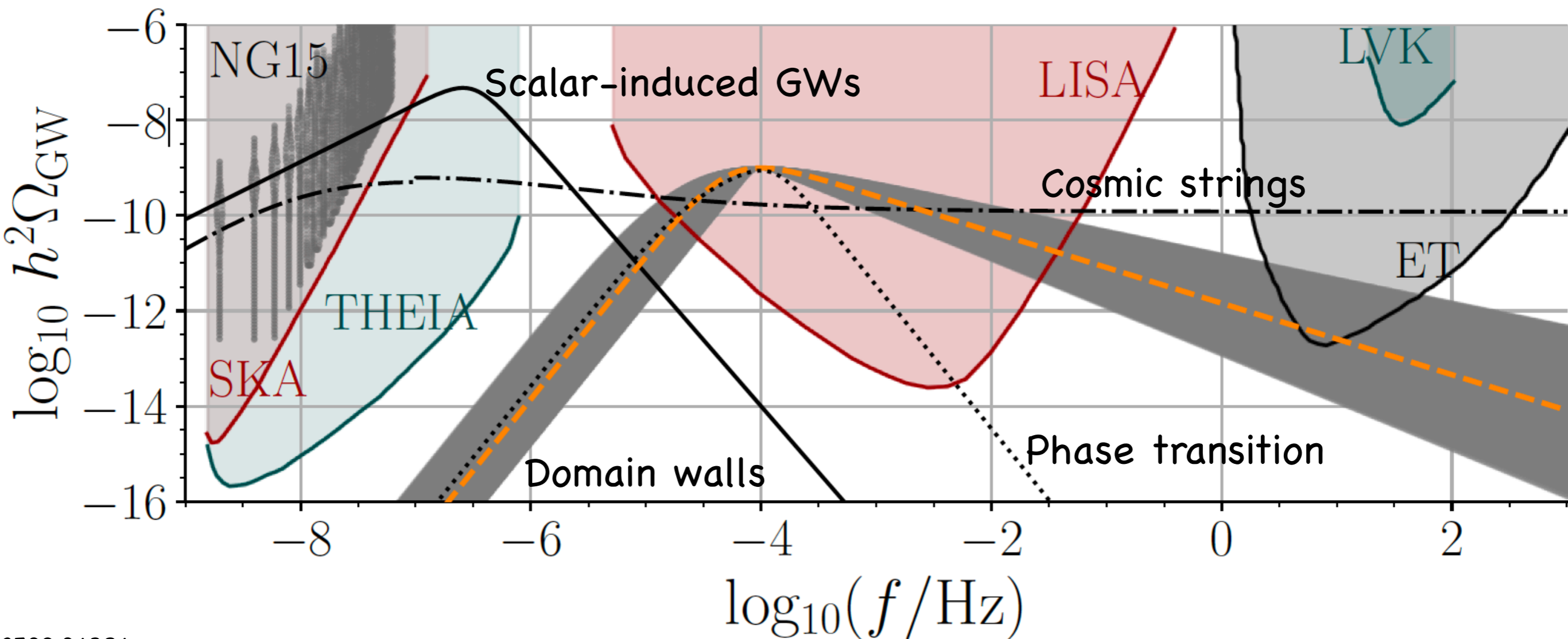
# Prominence



2509.04384

$$\mathcal{P}_i = \log_{10} \left[ h^2 \Omega_{\text{GW}}^{\text{peak}} \right]_i - \log_{10} \left[ h^2 \Omega_{\text{GW}}^{\text{base}} \right]_i$$

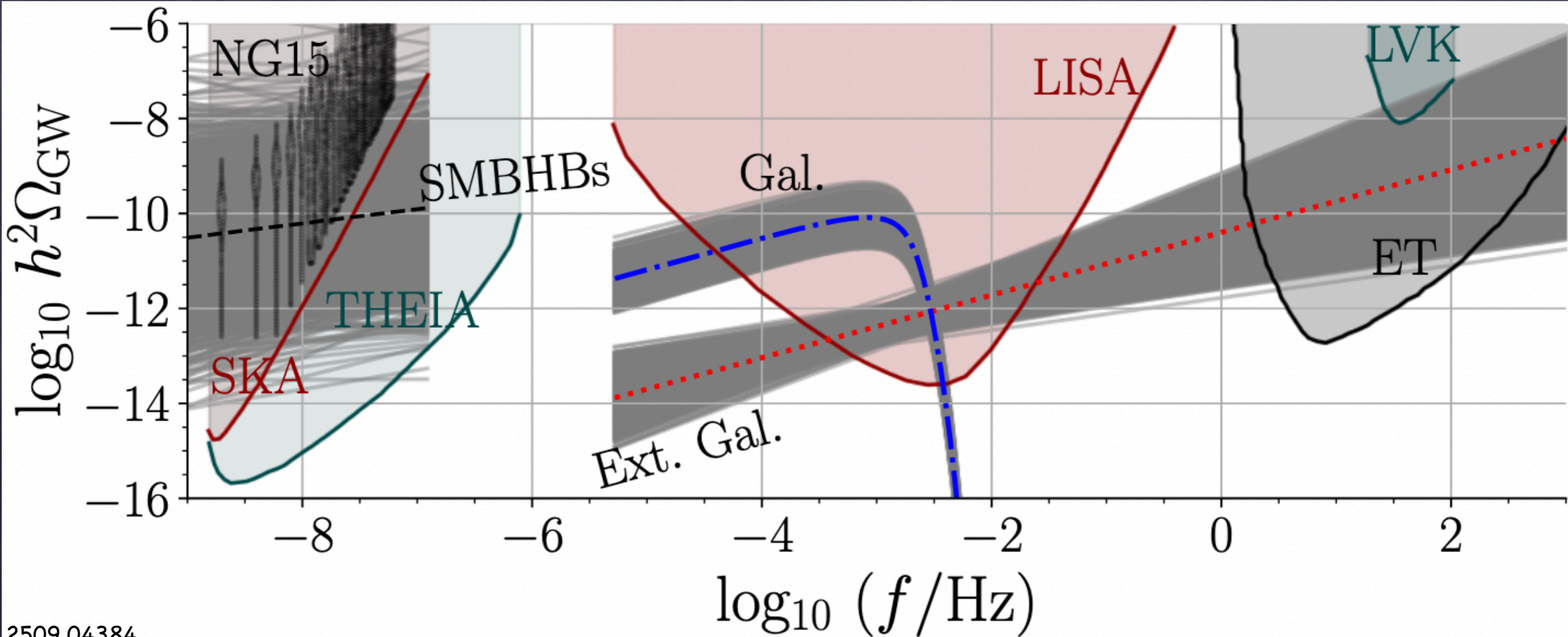
# GW signals



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- DW spectrum is broader than FOPT spectrum
- DW and FOPT spectra are similar below peak frequency
- CS spectrum in interferometer band is almost flat
- SIGWs only relevant in PTA band

# Backgrounds

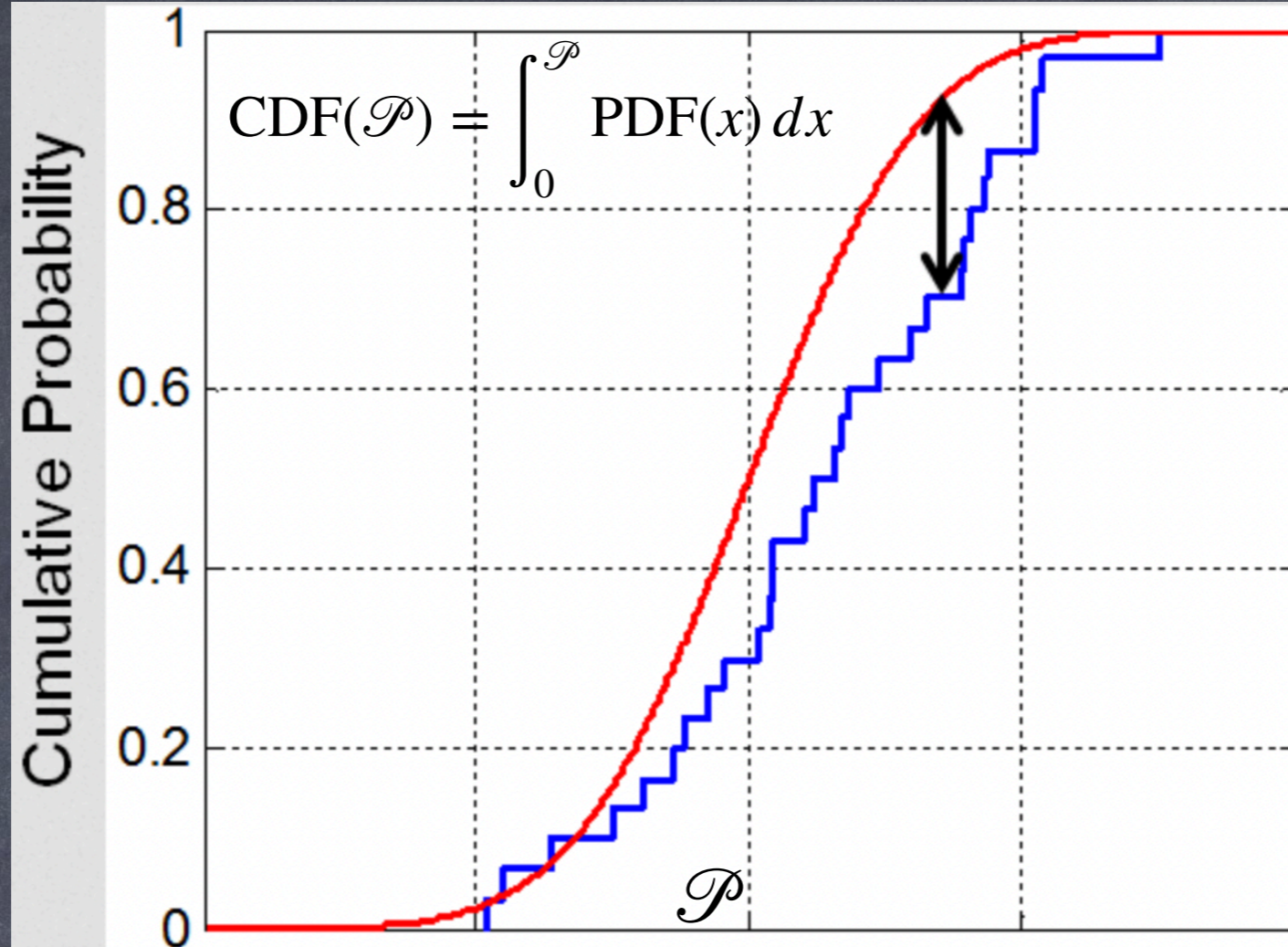


2509.04384

- SMBHB spectrum extends to higher frequencies and falls off above 1 Hz
- Interferometer bkg templates subtract individually resolvable binaries
- We analyze PTA and interferometer bands separately

## Assumptions and procedure to illustrate how $\mathcal{P}$ works ...

- Observation time: 4 years
- Detection threshold: SNR = 10
- Spectrum must have at least one peak
- CS contribution in interferometer band is subdominant to FOPT and DW peaks
- Peak amplitude and frequency of FOPT and DW spectra are identical (most conservative case)
- Simulate datasets with backgrounds varied within ranges determined by simulations
- For each peak construct PDF of  $\mathcal{P}$



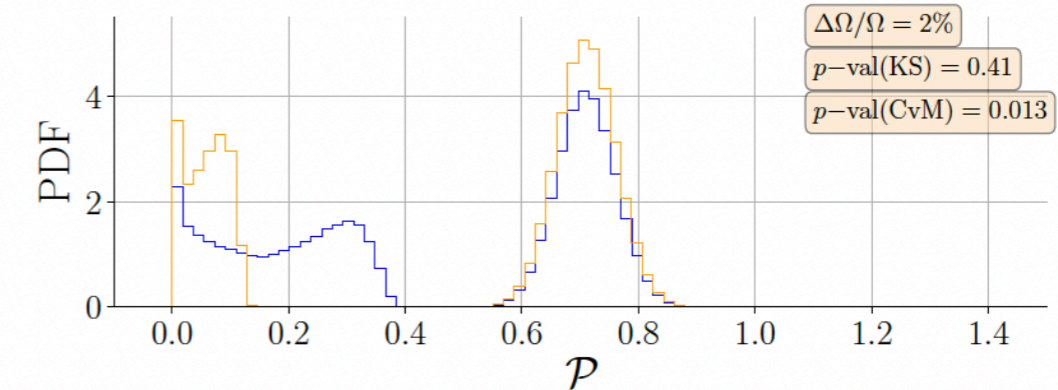
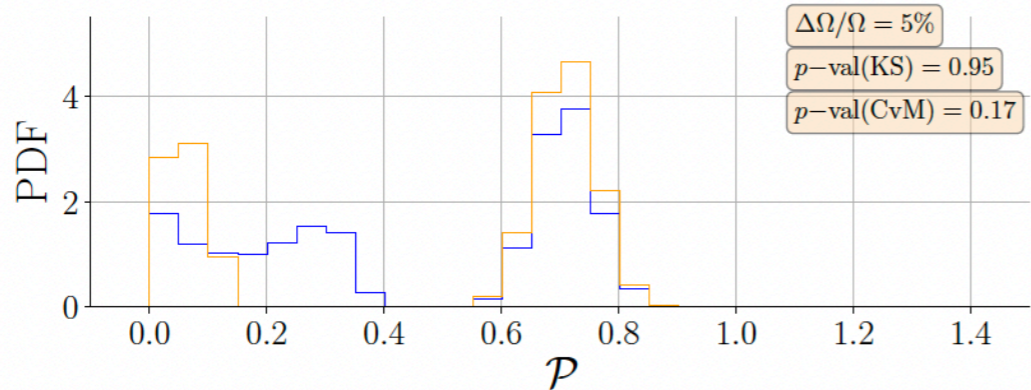
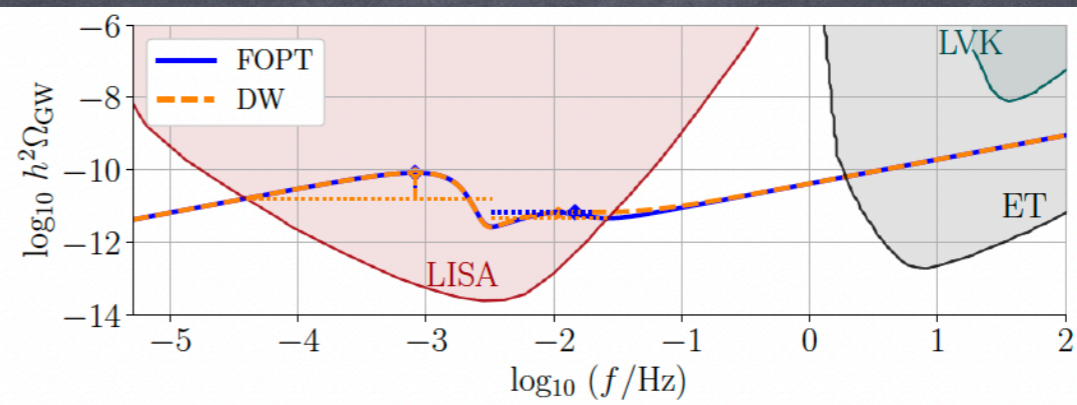
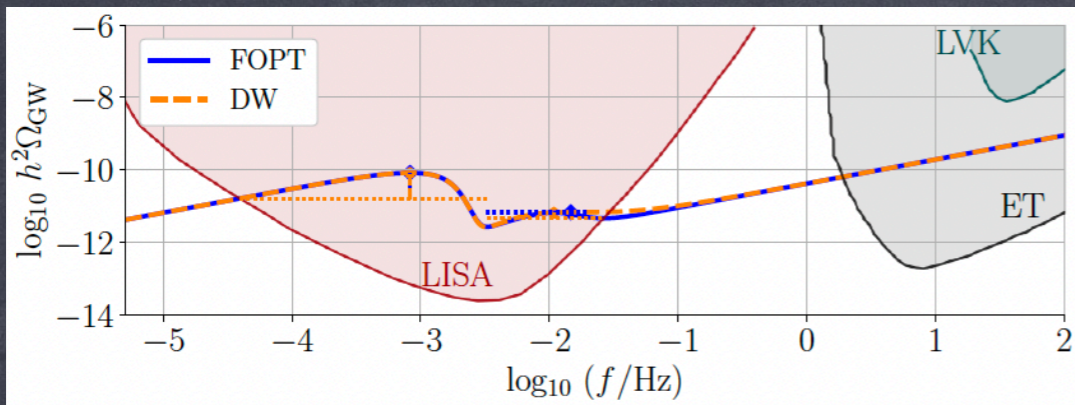
Kolmogorov-Smirnov test:

$$\text{KS} = \max_{\mathcal{P}} \left| \text{CDF}_1(\mathcal{P}) - \text{CDF}_2(\mathcal{P}) \right|$$

Cramer-von Mises test (measures the aggregate difference in CDFs):

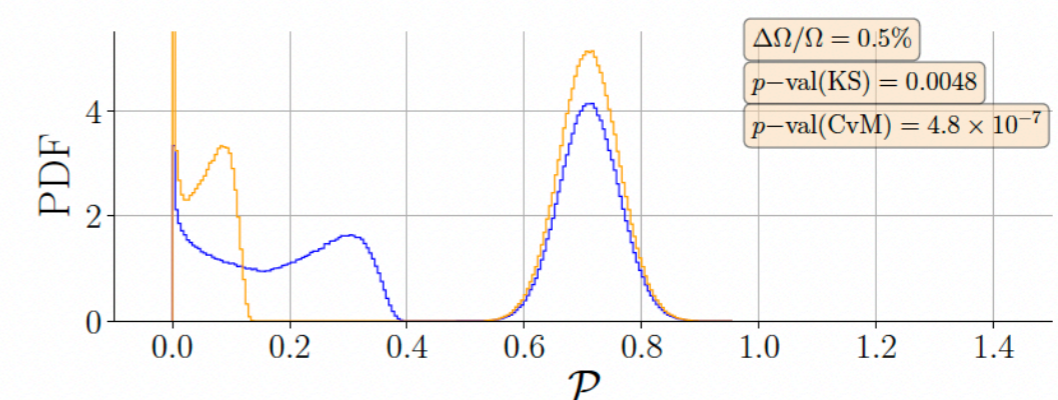
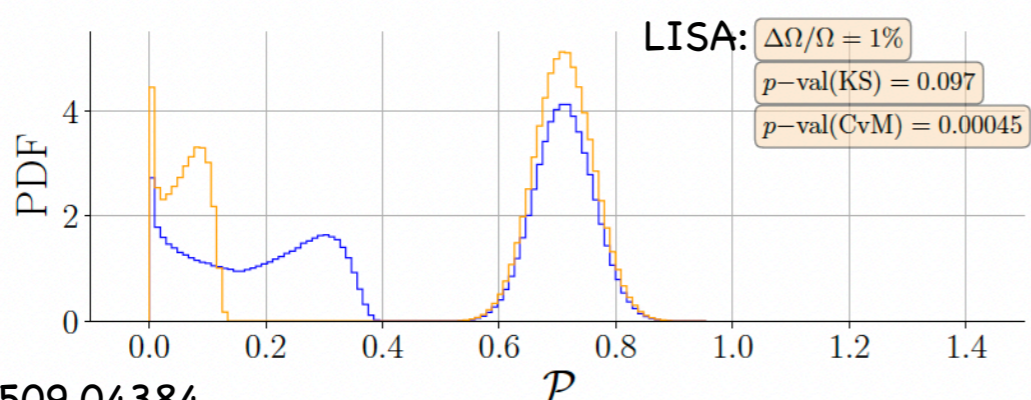
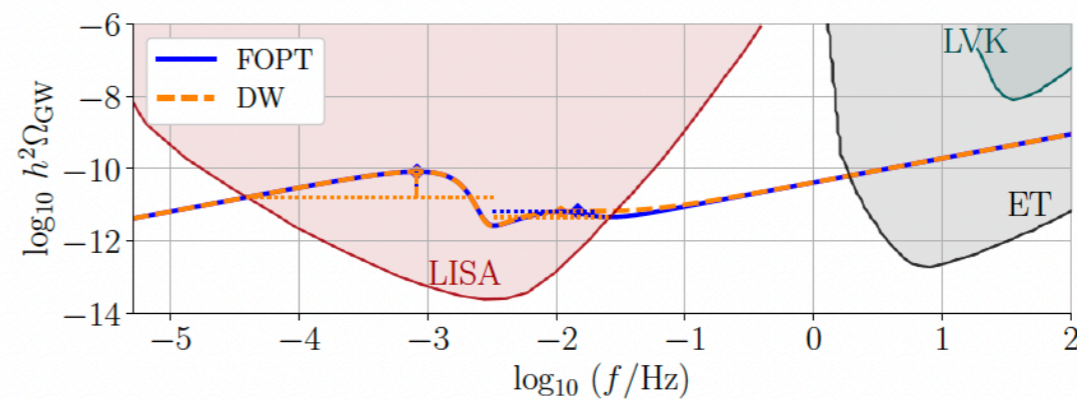
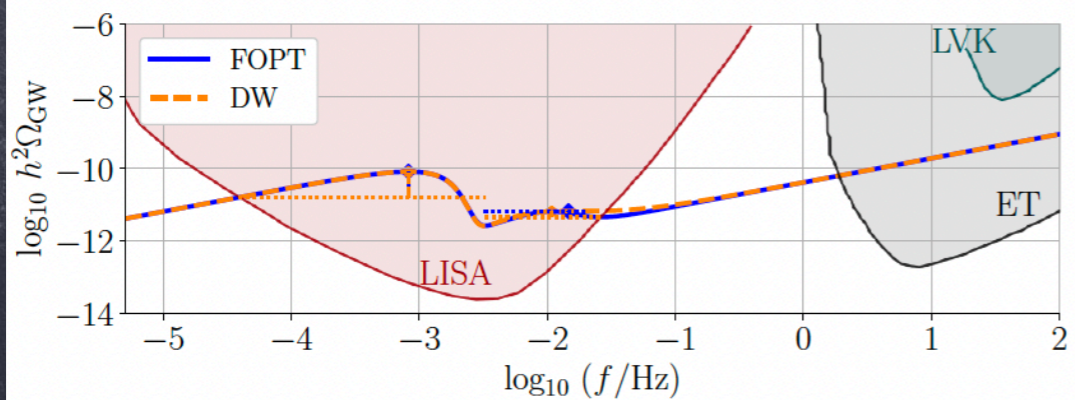
$$\text{CvM} = \int_0^{\infty} \left[ \text{CDF}_1(\mathcal{P}) - \text{CDF}_2(\mathcal{P}) \right]^2 d\text{CDF}_1(\mathcal{P})$$

# Impact of amplitude uncertainties



(a)

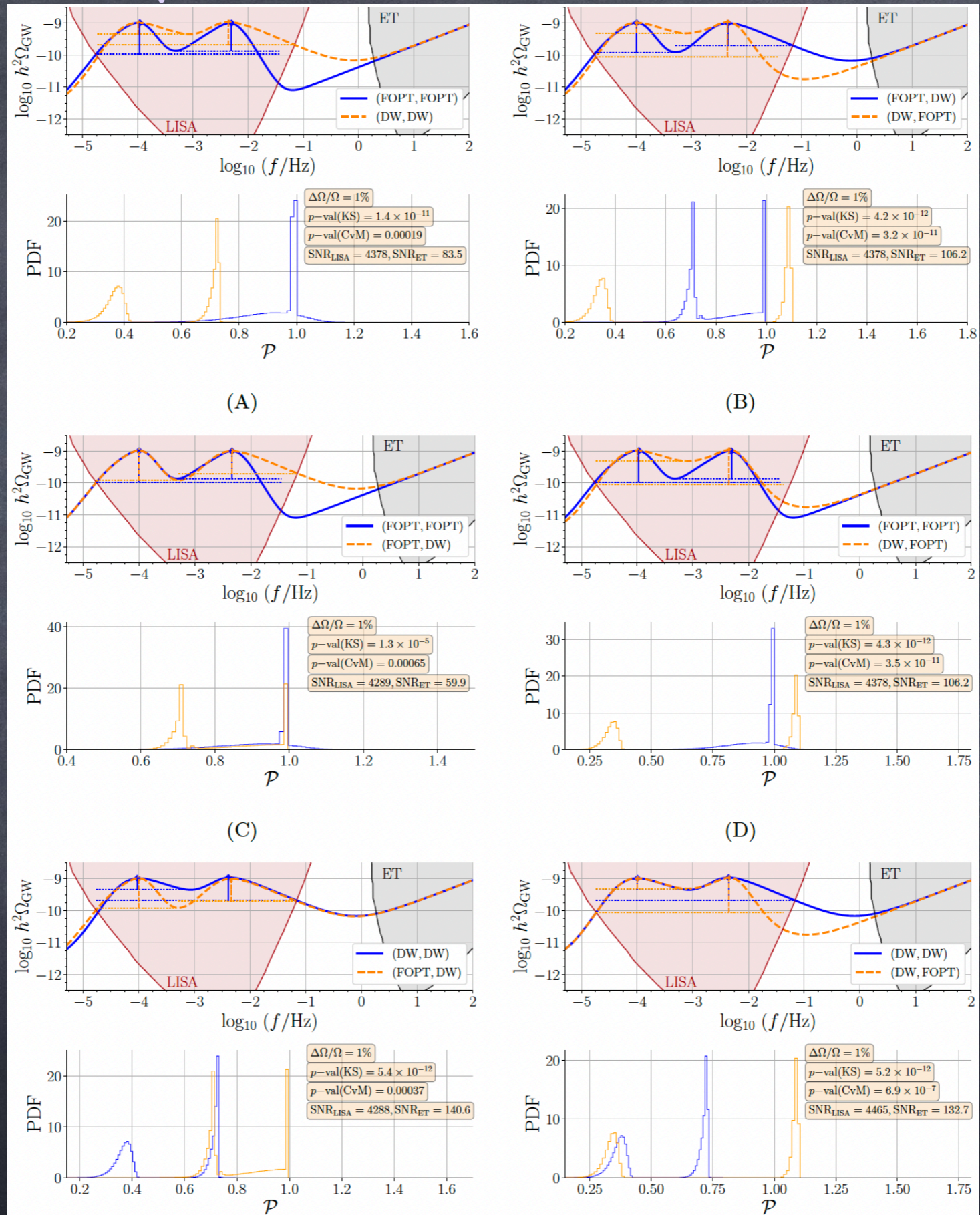
(b)



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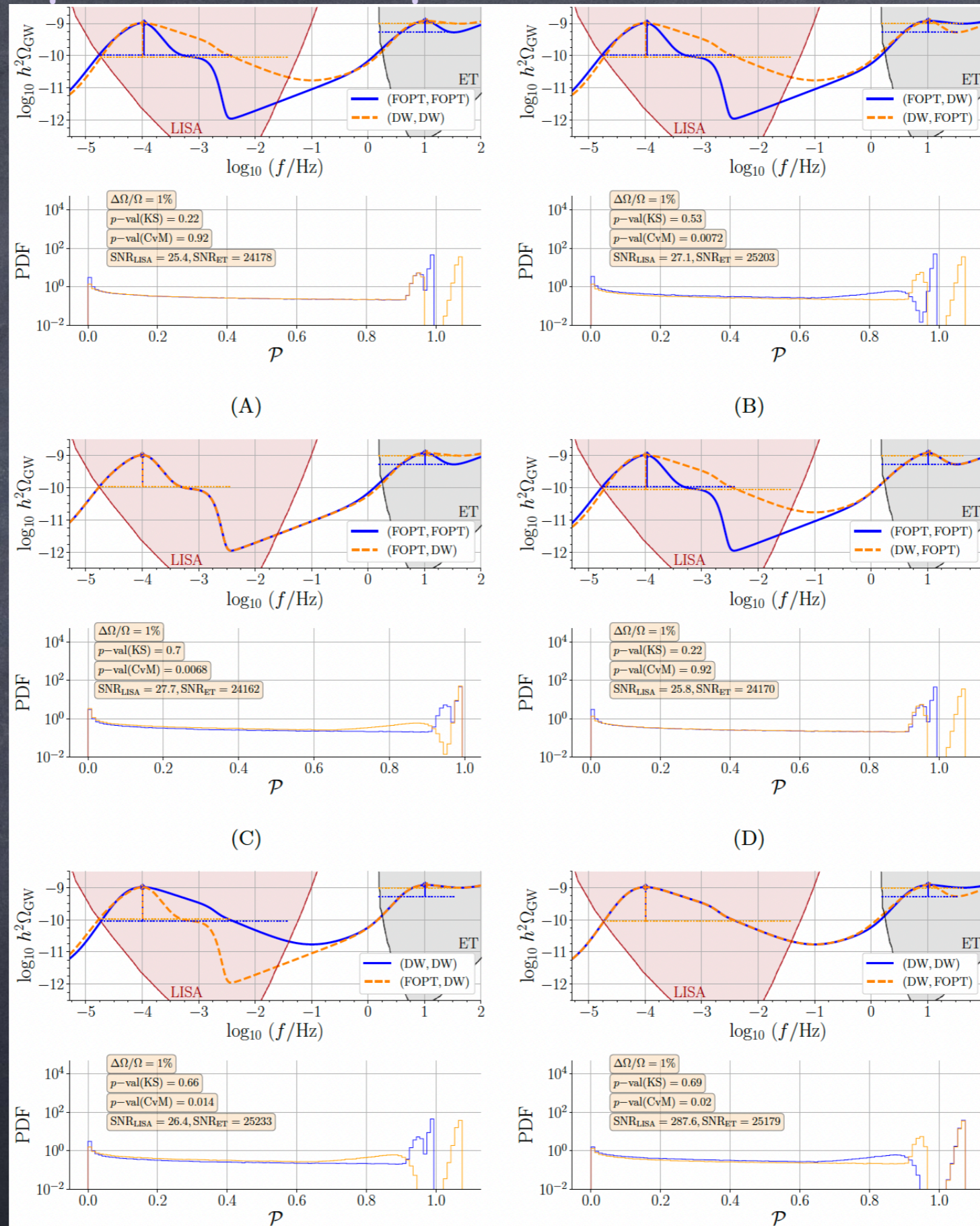
Bin size determined by measurement uncertainty in  $\Omega_{\text{GW}}$

# Two-peaked spectra in the LISA band

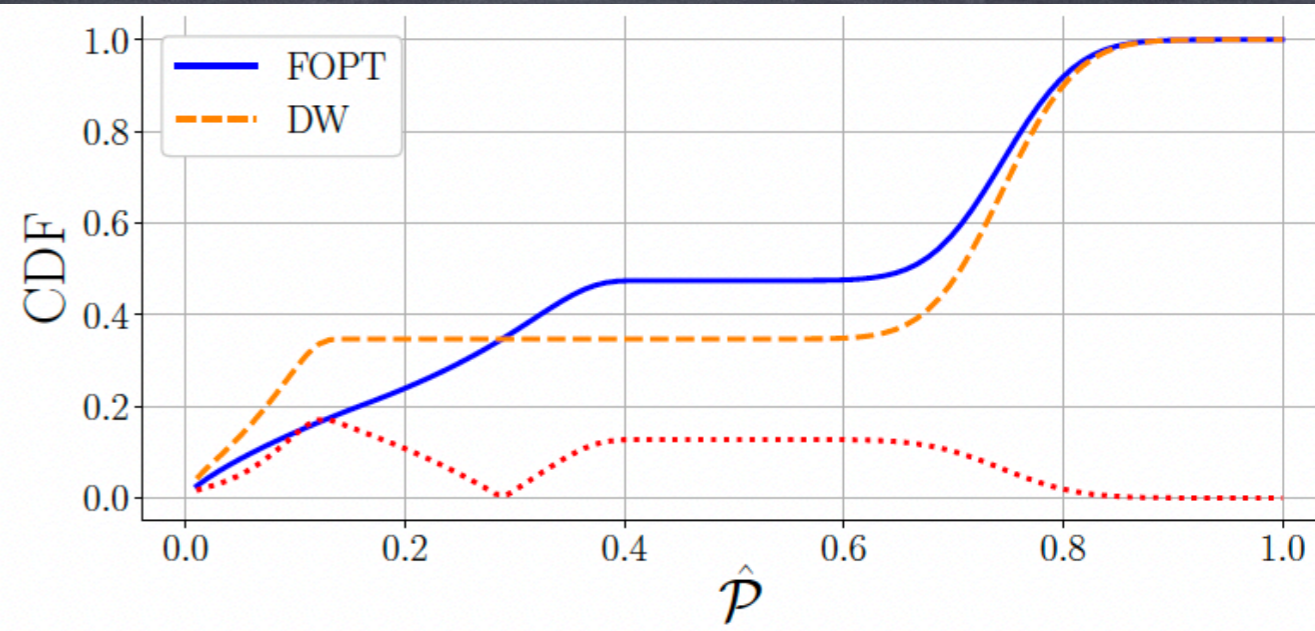
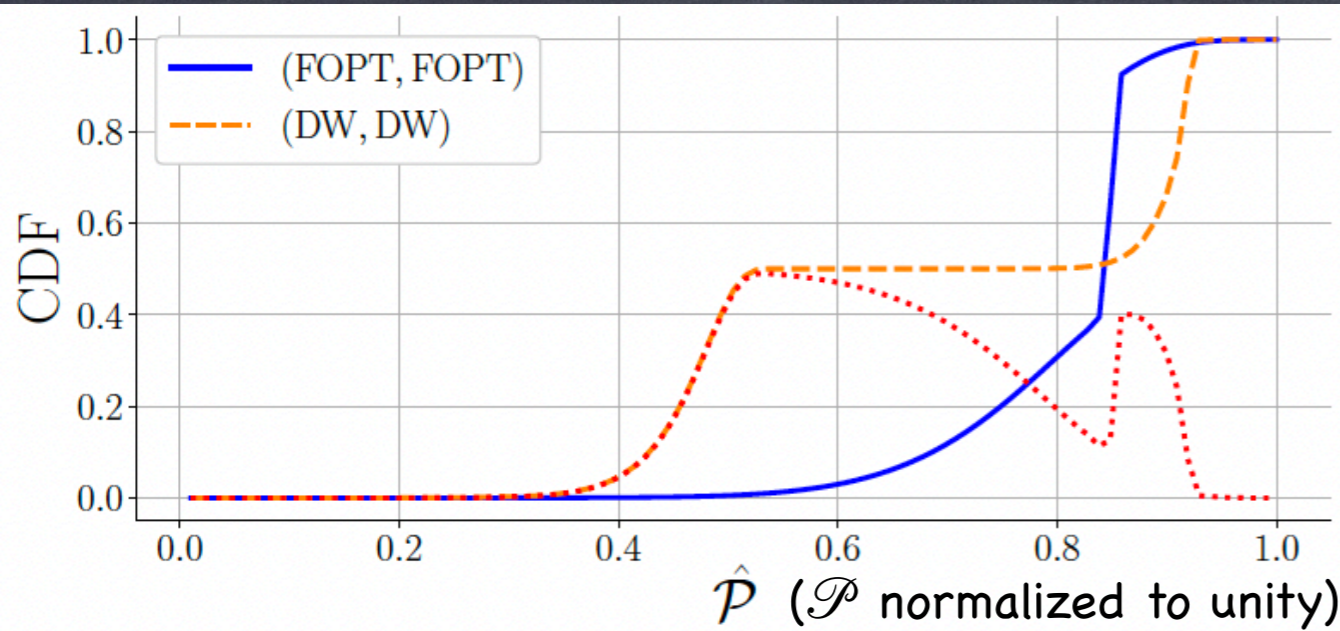


$>3.4\sigma$  discrimination with CvM test. Much more sensitivity with KS test

# Two-peaked spectra with peaks in the LISA and ET bands

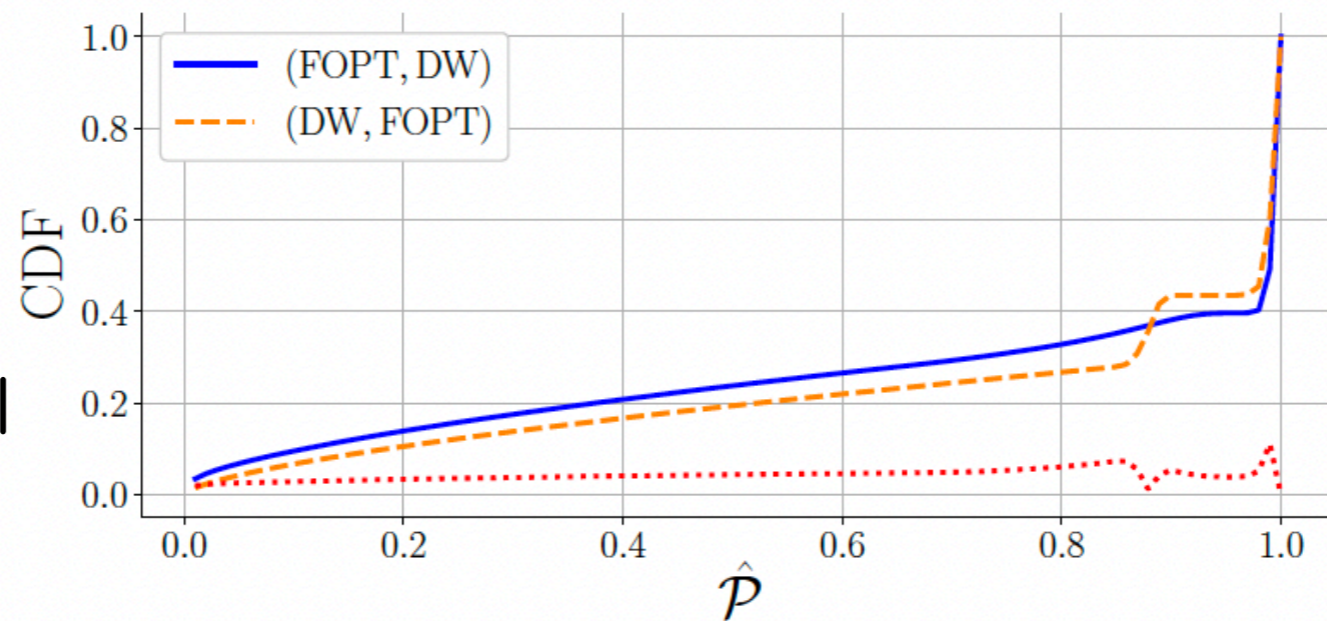


Poor discrimination because of large extragalactic bkg in the ET band



(a)  $p\text{-val(KS)} = 1.4 \times 10^{-11}$ ,  $p\text{-val(CvM)} = 0.00019$ .

(b)  $p\text{-val(KS)} = 0.097$ ,  $p\text{-val(CvM)} = 0.00045$ .



(c)  $p\text{-val(KS)} = 0.53$ ,  $p\text{-val(CvM)} = 0.0072$ .

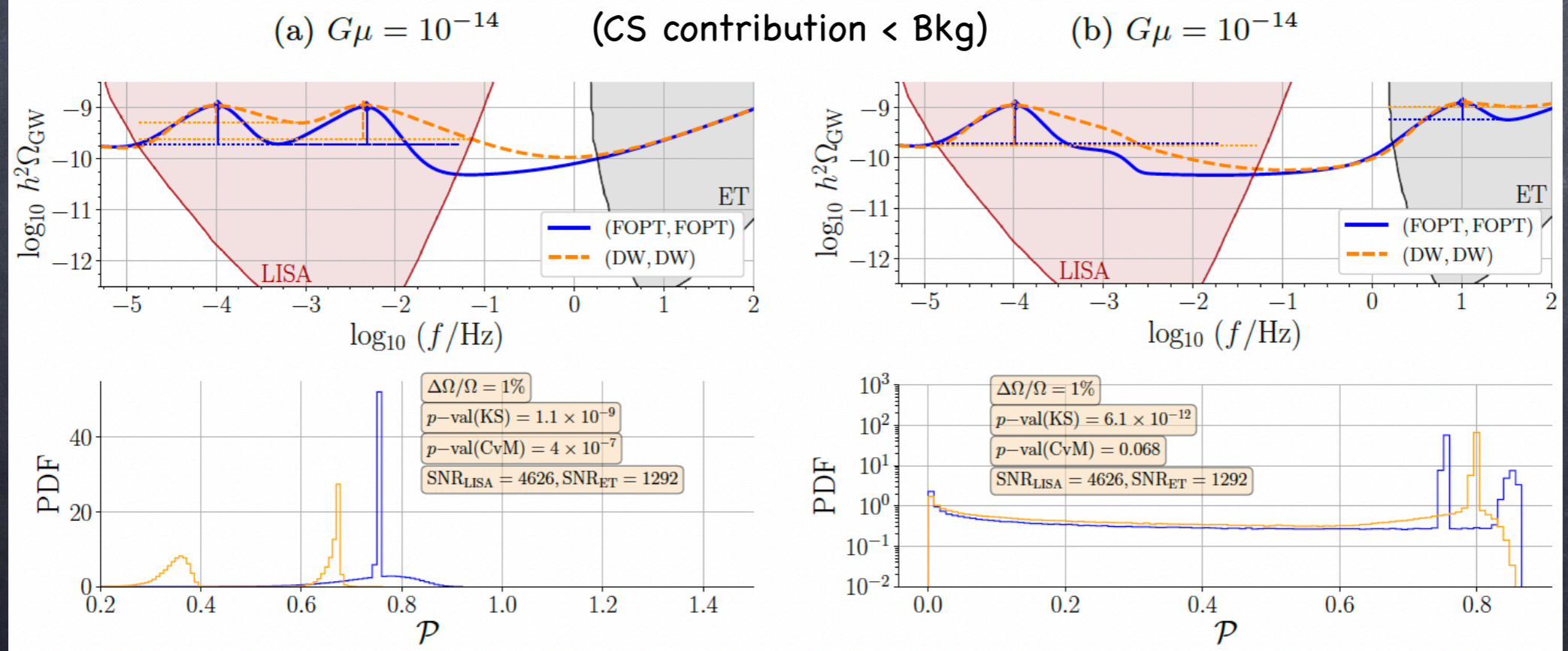
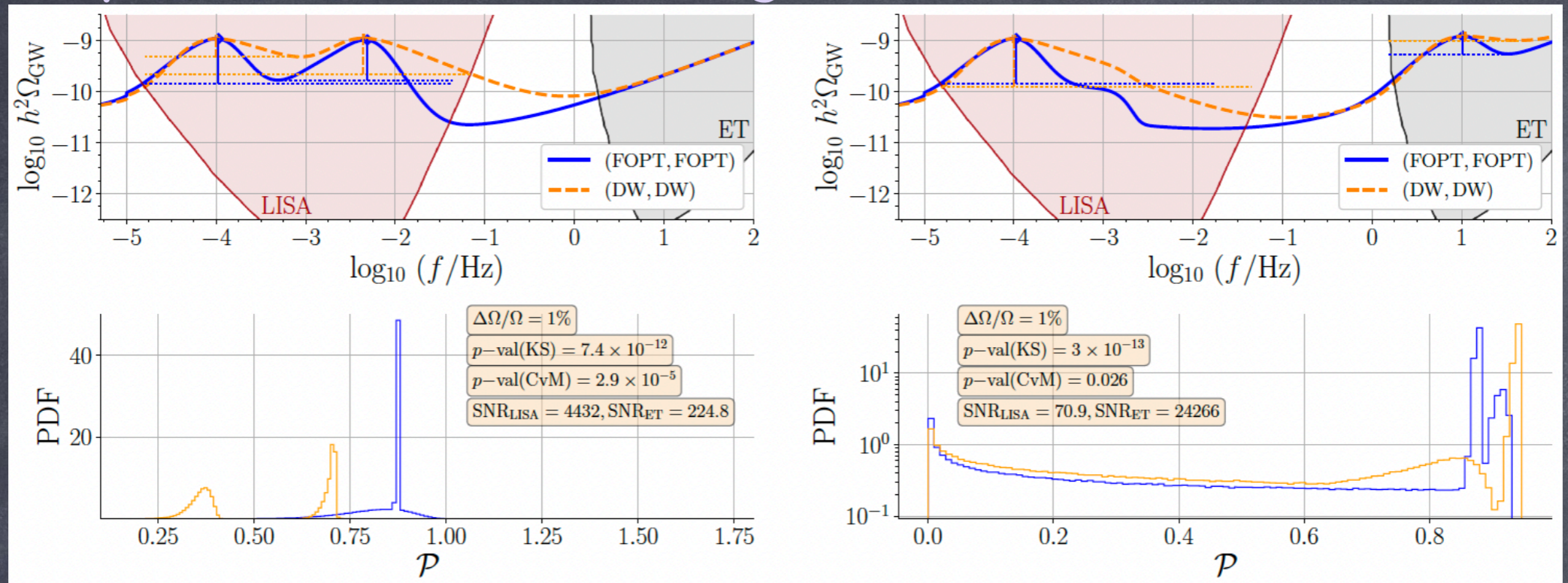
dotted = |solid - dashed|

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CvM is usually more conservative than KS, but can give smaller p-values if the CDFs

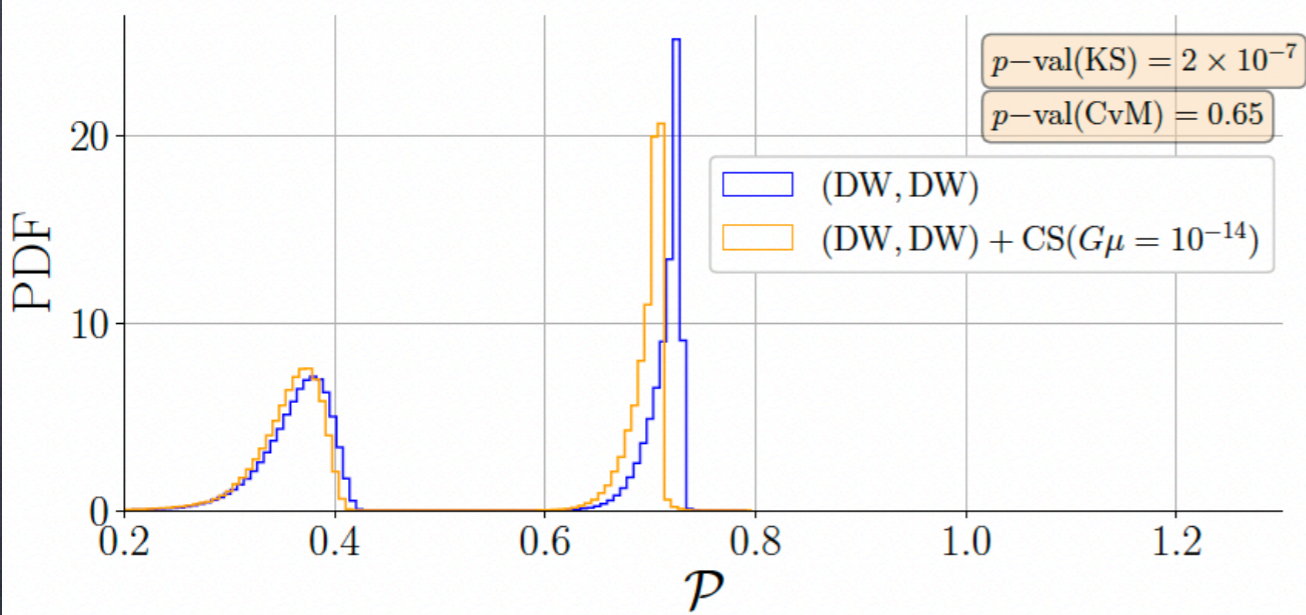
- 1) do not have sharp features
- 2) exhibit a systematic shift or persistent differences

# Impact of cosmic strings

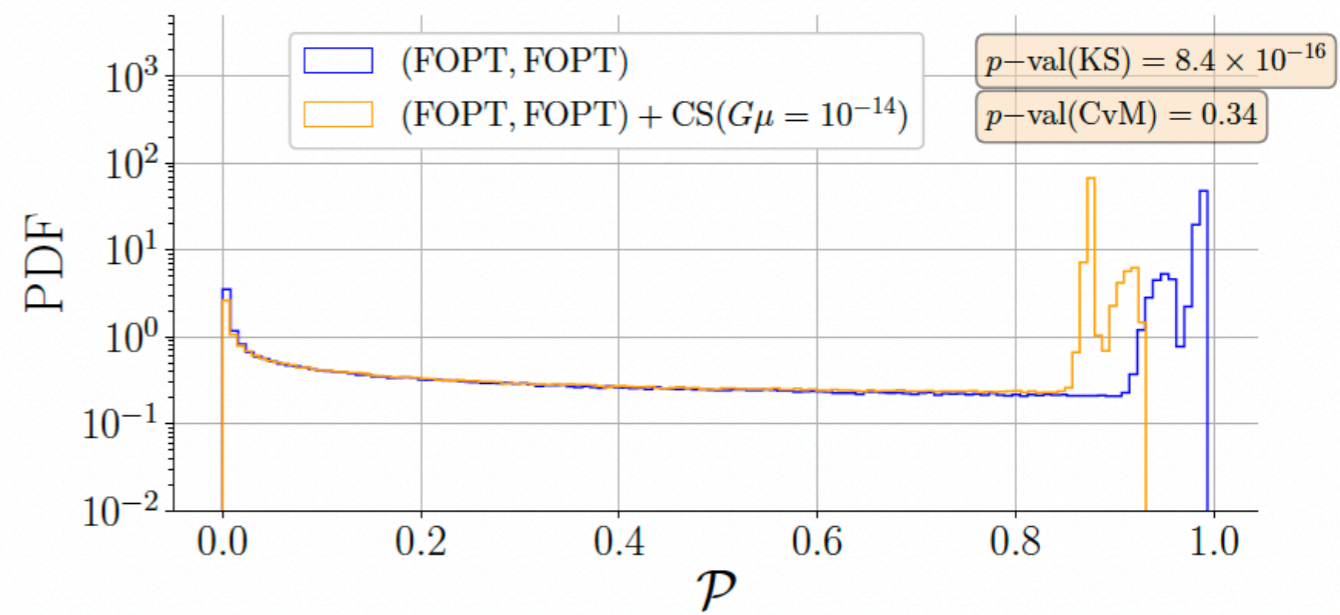


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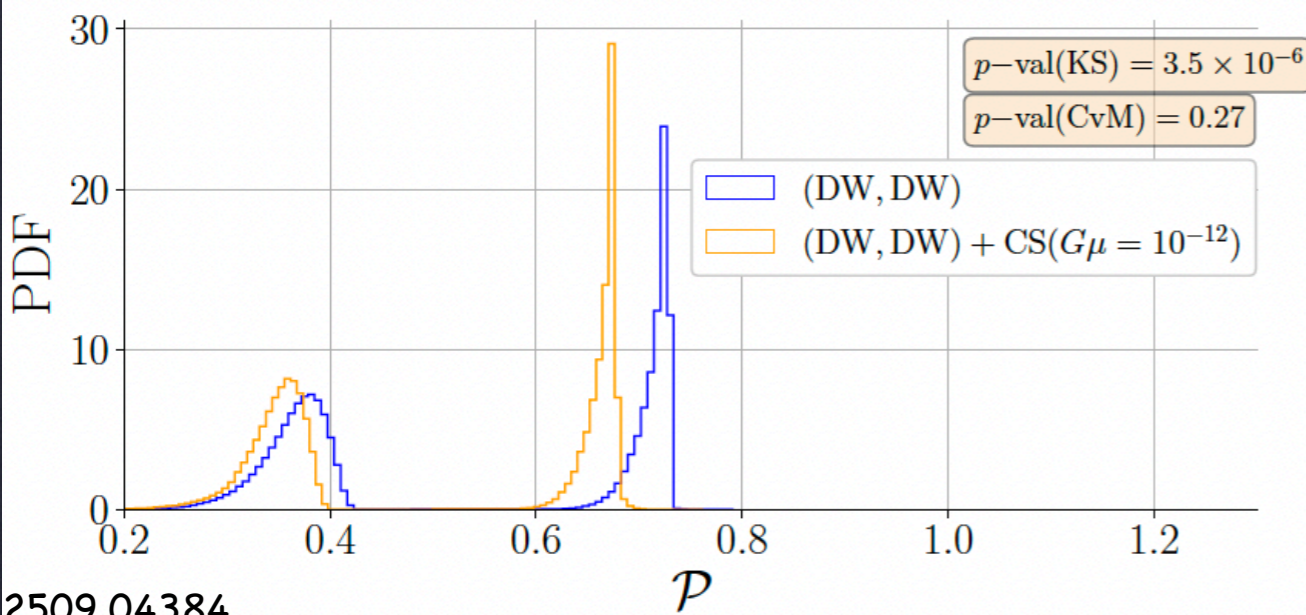
CSs can enhance the discrimination power of  $\mathcal{P}$



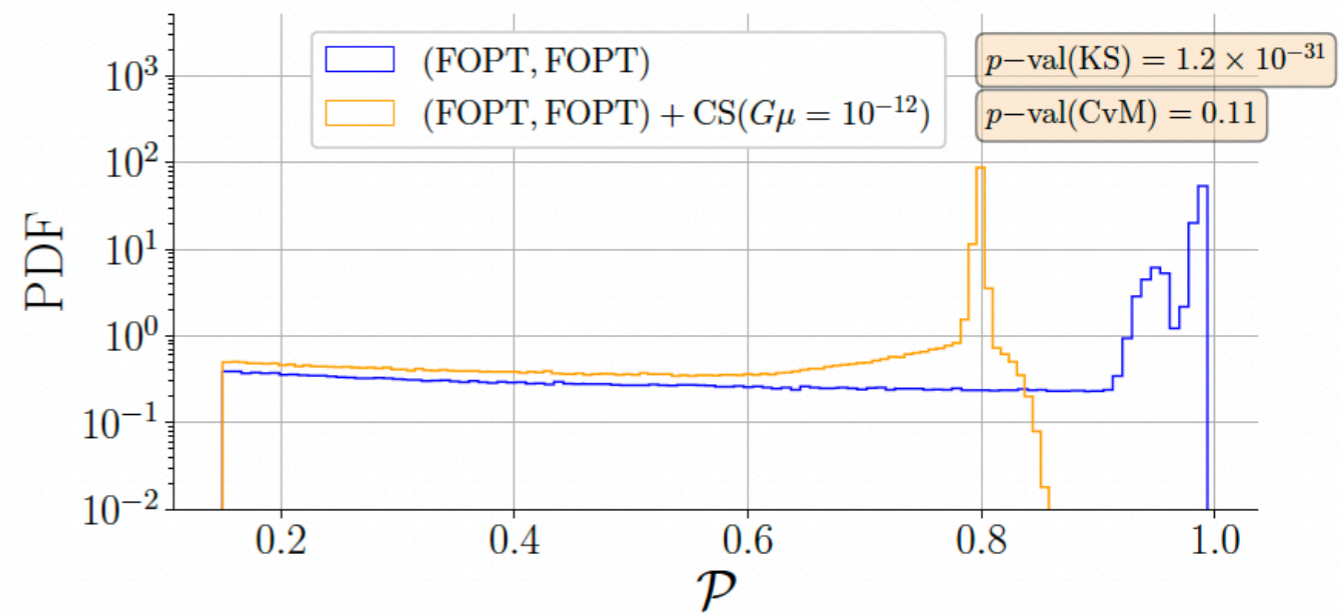
(a)



(b)



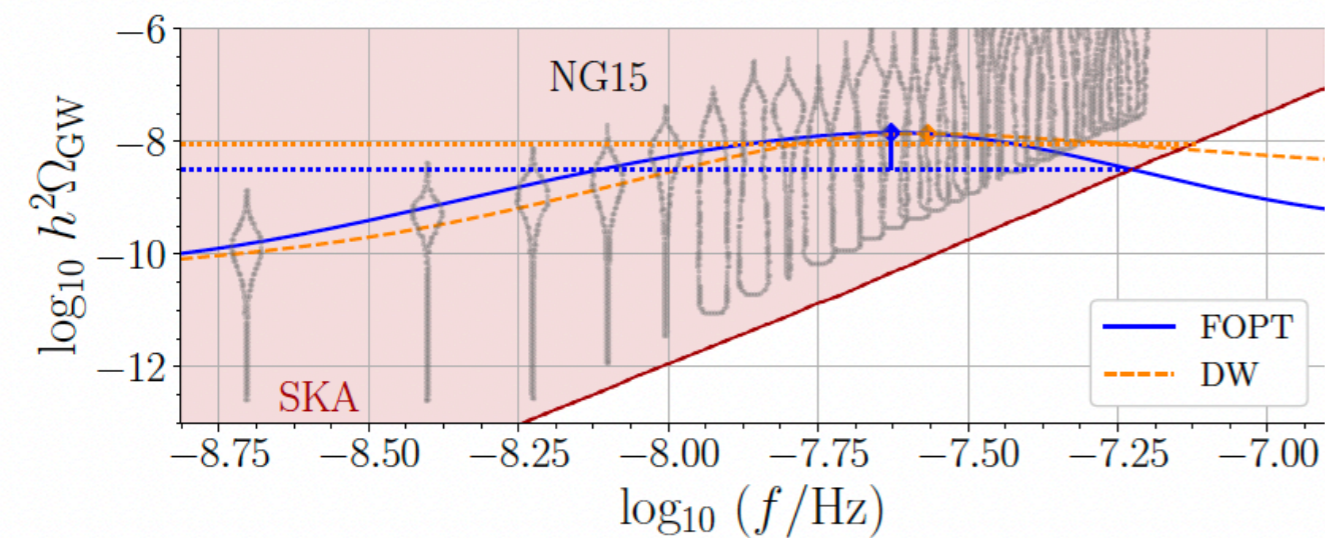
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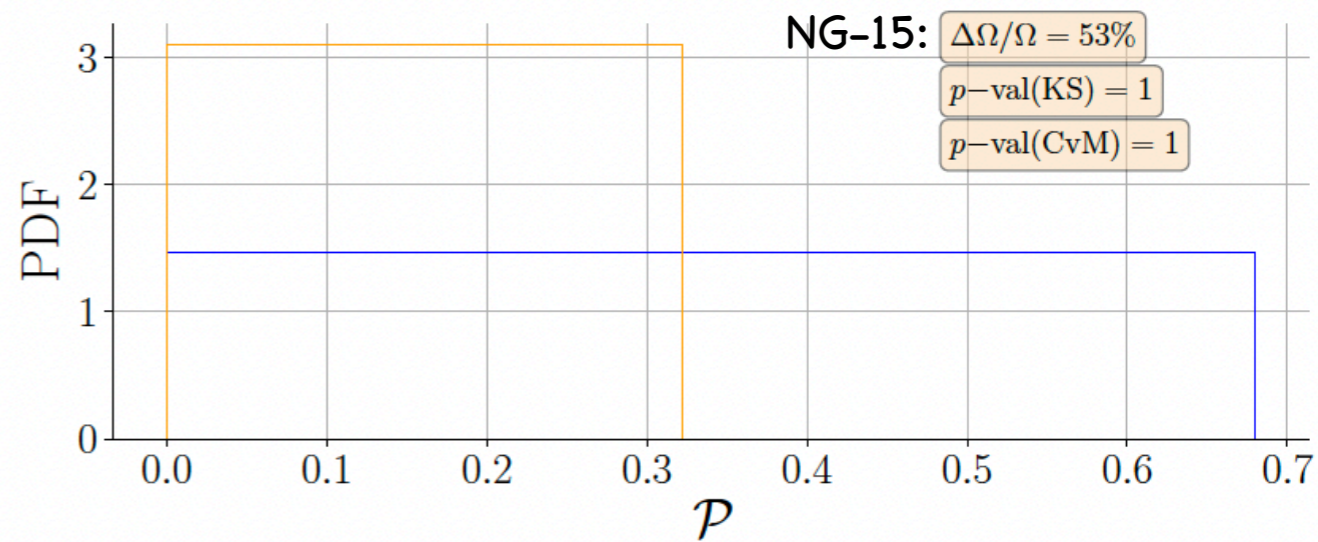
CSs can bury the bkg and produce large local differences giving small  $p$ -values for the KS test

Bigger effect for FOPTs than DWs because spectrum is narrower

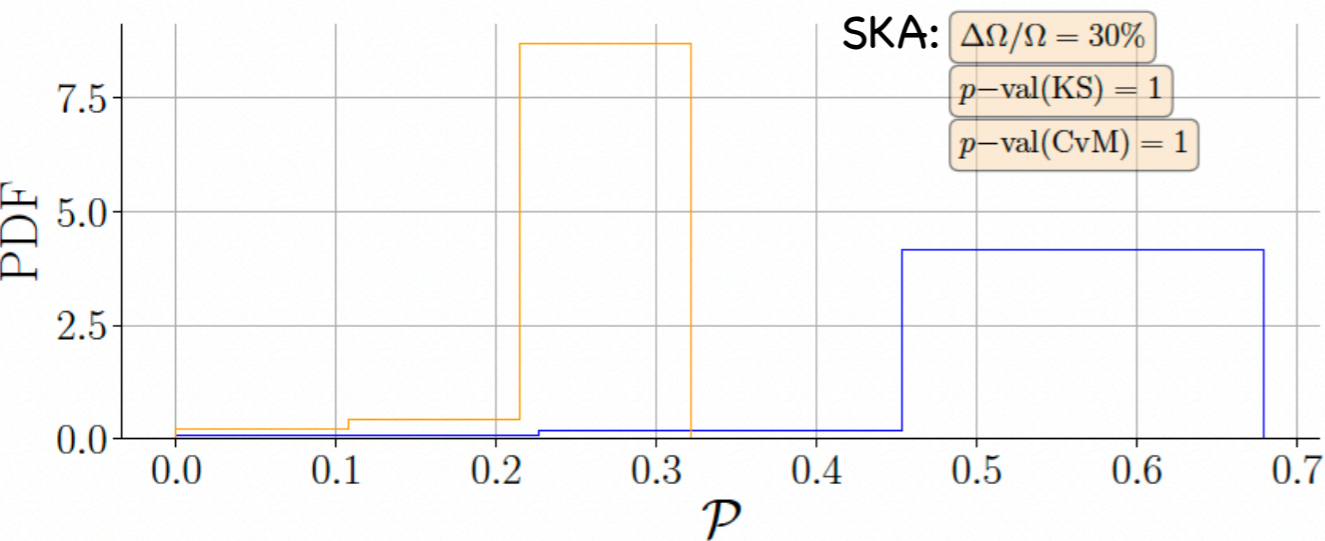
# Signals in the PTA band



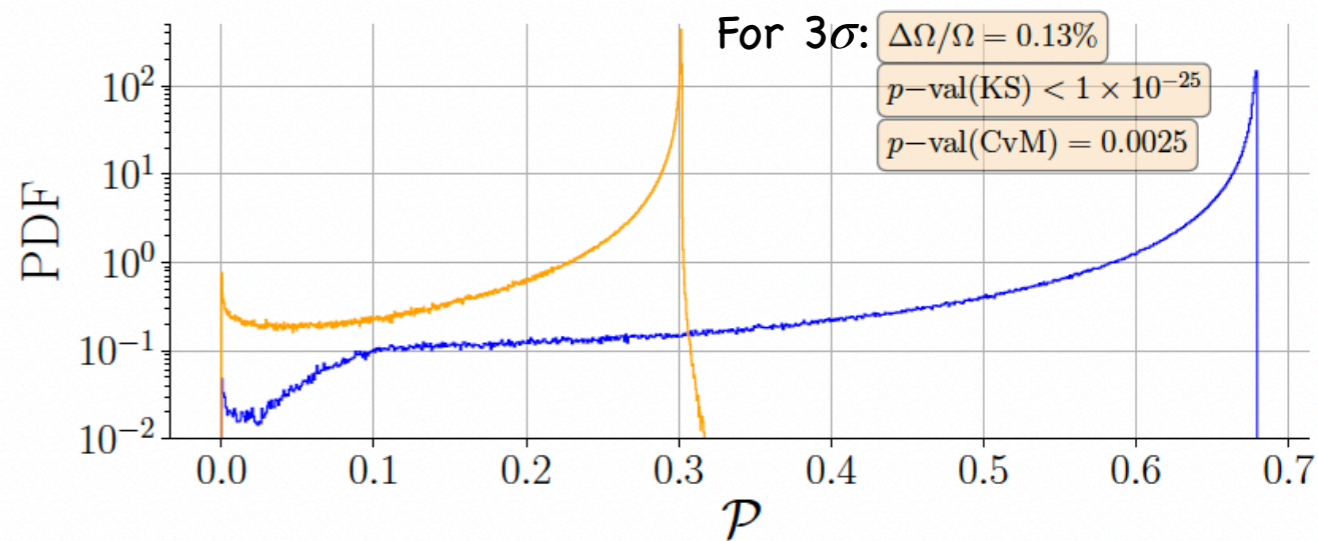
(a)



(b)



(c)

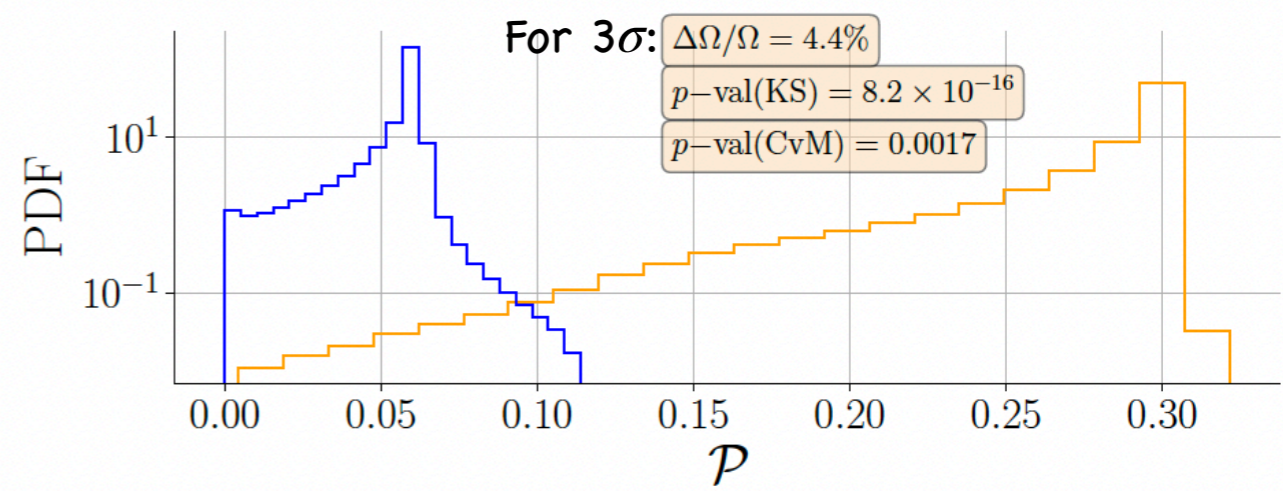
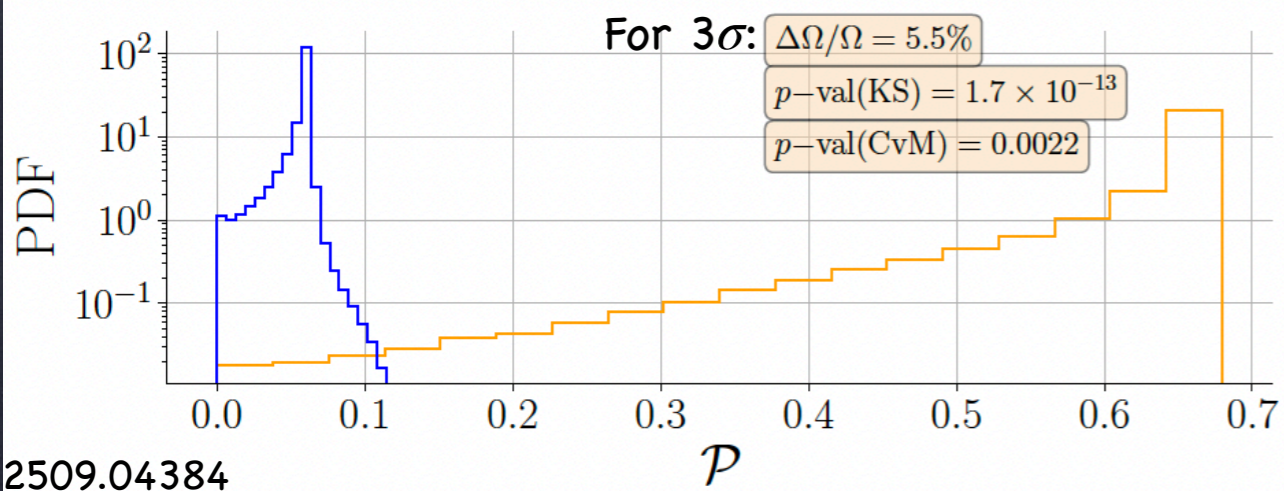
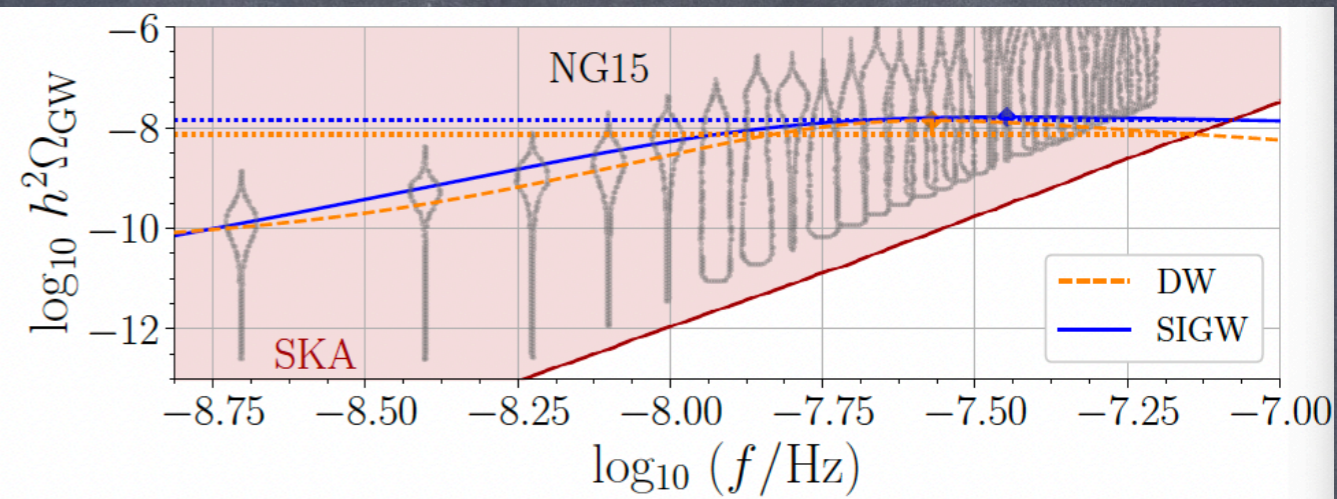
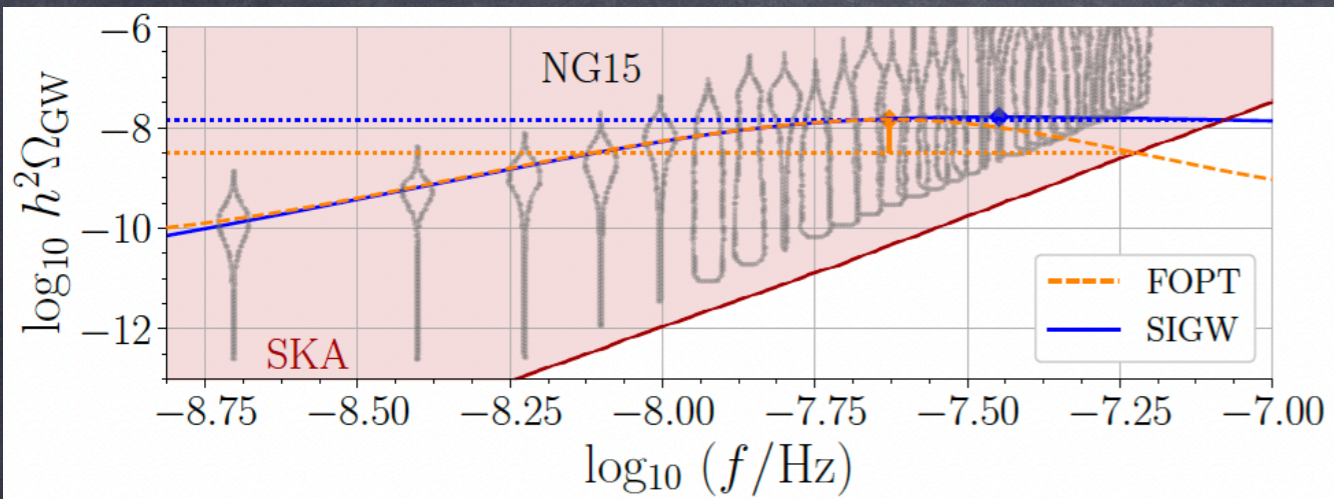


(d)

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FOPT, DW, SIGW parameters set by fitting to NANOGrav-15 data

Gauge CSs do not fit these data well



Easier to discriminate between FOPT/SIGW, DW/SIGW than FOPT/DW

Prominence in real-world settings

# Prominence vs SNR-based signal discrimination

$$\Omega_i = \Omega_{\text{SIG}_i} + \Omega_{\text{Bkg}}$$

$$\langle \Omega_1 | \Omega_2 \rangle = T_{\text{obs}} \int_{f_{\text{min}}}^{f_{\text{max}}} df \frac{\Omega_1(f) \Omega_2(f)}{\Omega_{\text{Sens}}^2(f)}$$

If  $\Omega_1 = \Omega_2 \equiv \Omega$ ,

$$\text{SNR} = \sqrt{\langle \Omega | \Omega \rangle} = \sqrt{T_{\text{obs}} \int_{f_{\text{min}}}^{f_{\text{max}}} df \frac{\Omega^2(f)}{\Omega_{\text{Sens}}^2(f)}}$$

Signal-space statistic:

$$\chi^2 = \langle \Omega_1 - \Omega_2 | \Omega_1 - \Omega_2 \rangle$$

$$\text{Mismatch : } \mathcal{M}_{12} \equiv 1 - \frac{\langle \Omega_1 | \Omega_2 \rangle}{\text{SNR}_1 \text{SNR}_2}$$

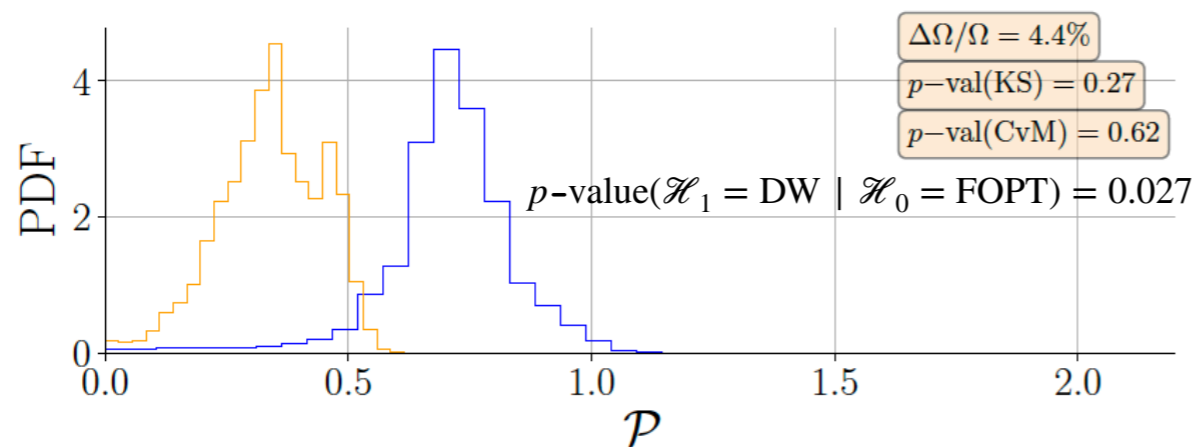
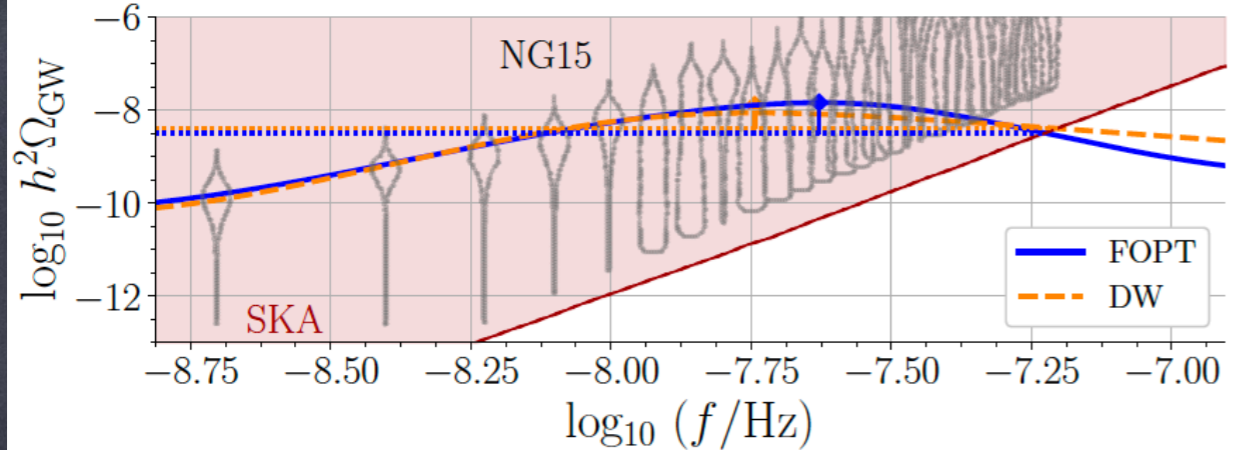
$$\therefore \chi^2 = (\text{SNR}_1 - \text{SNR}_2)^2 + 2\mathcal{M}_{12} \text{SNR}_1 \text{SNR}_2$$

For identical signals,  $\mathcal{M}_{12} = 0 \implies \chi^2 = 0$

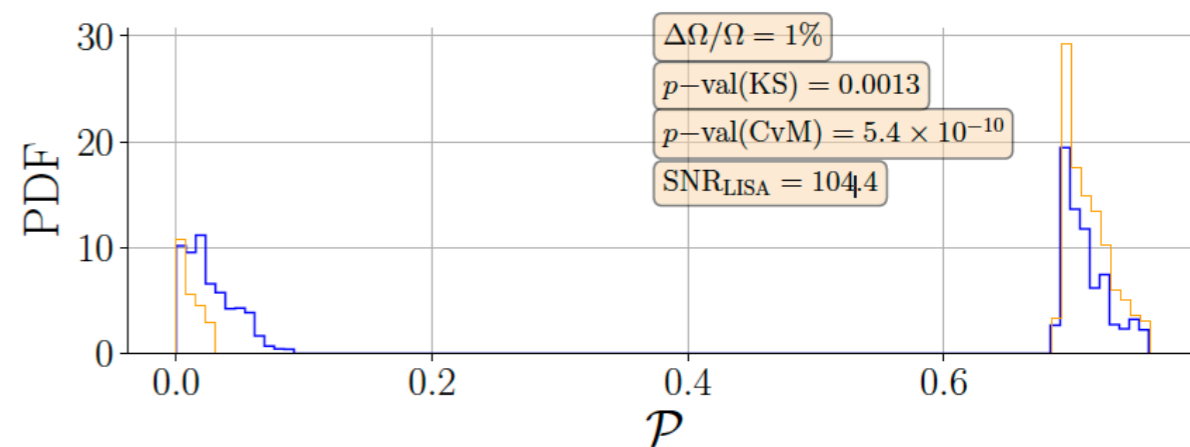
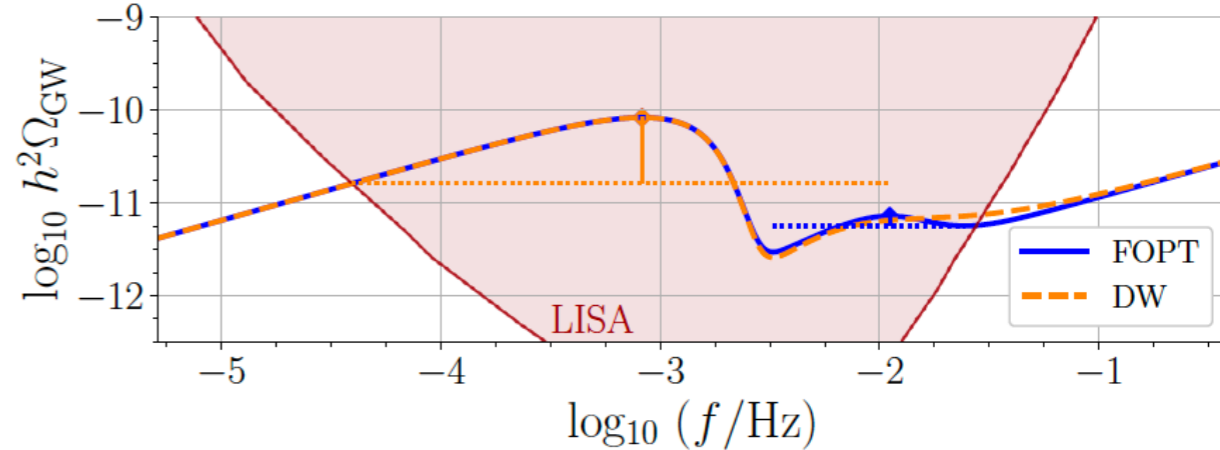
For no signal overlap,  $\mathcal{M}_{12} = 1 \implies \chi^2 = \text{SNR}_1^2 + \text{SNR}_2^2$

- Fix  $\Omega_1$  (simulating experimental spectrum)
- Fit with  $\Omega_2$  by varying  $\Omega_{\text{SIG}_2}$  and  $\Omega_{\text{Bkg}}$
- Scan over parameters in table for simulated source and hypothetical source
- Select points for each source within  $2\sigma$  as per SNR
- Produce CDFs for  $\mathcal{P}$  and apply KS and CvM tests

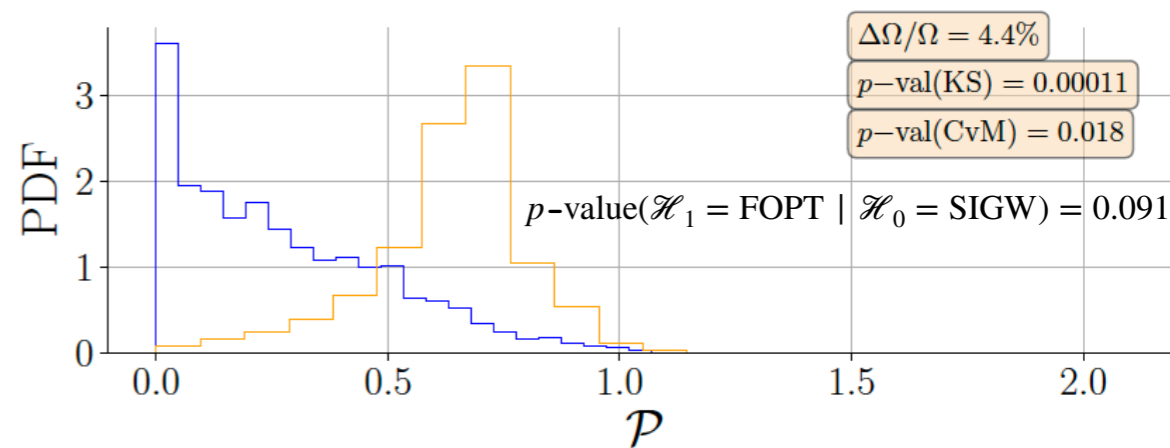
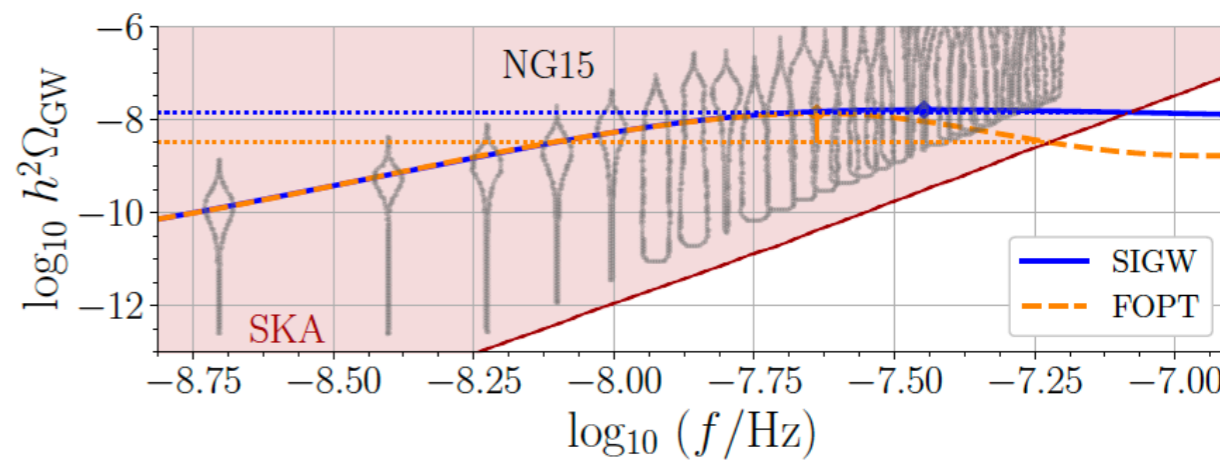
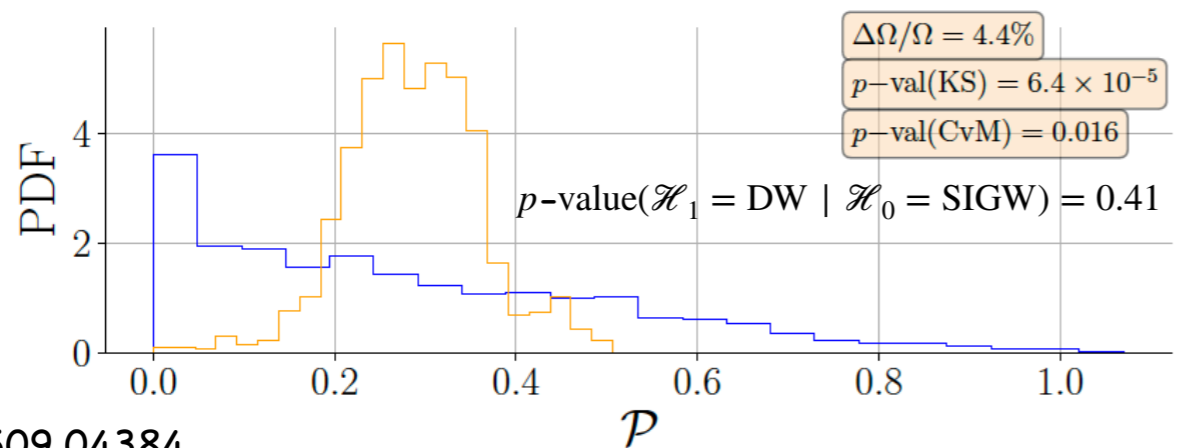
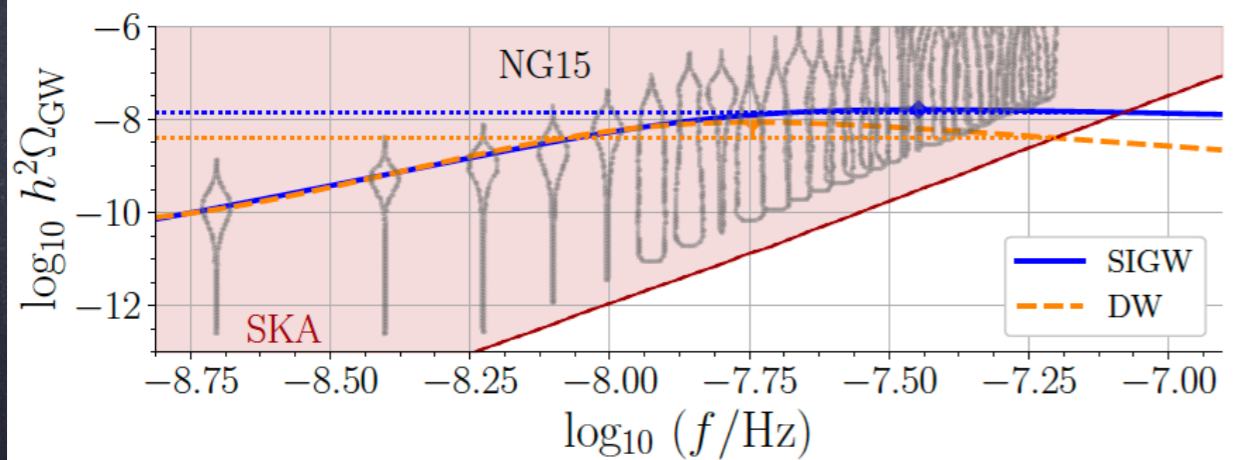
Simulated source	Hypothesis	Signal parameters	Background parameters	$\chi^2$ threshold
FOPT (PTA)	DW	$V_{\text{bias}}, \mathcal{E}, b, c$	$A_{\text{BHB}}, \gamma_{\text{BHB}}$	9.70
FOPT (PTA)	FOPT	$\alpha, \beta/H(T_*), T_*$	$A_{\text{BHB}}, \gamma_{\text{BHB}}$	8.02
FOPT (LISA)	DW	$V_{\text{bias}}, \mathcal{E}, b, c$	$h^2\Omega_{\text{Gal}}, h^2\Omega_{\text{Ext}}, n_s, \alpha_{\text{Ext}}$	9.70
FOPT (LISA)	FOPT	$\alpha, \beta/H(T_*), T_*$	$h^2\Omega_{\text{Gal}}, h^2\Omega_{\text{Ext}}, n_s, \alpha_{\text{Ext}}$	8.02
SIGW (PTA)	SIGW	$\alpha_{\text{SIG}}, \beta_{\text{SIG}}, A, f_c$	$A_{\text{BHB}}, \gamma_{\text{BHB}}$	9.70
SIGW (PTA)	DW	$V_{\text{bias}}, \mathcal{E}, b, c$	$A_{\text{BHB}}, \gamma_{\text{BHB}}$	9.70
SIGW (PTA)	FOPT	$\alpha, \beta/H(T_*), T_*$	$A_{\text{BHB}}, \gamma_{\text{BHB}}$	8.02



(a)



(b)



# Summary

- Prominence is sensitive not only to the peak amplitude of the GW spectrum, but also its shape
- $\mathcal{P}$  is especially useful when signals are similar and have relatively small SNR
- Amplitude uncertainty at PTAs needs to be at the sub-percent to percent level for  $\mathcal{P}$  to distinguish signals at  $3\sigma$
- Uncertainty at LISA and ET will be small enough ( $\sim 1\%$ ) that  $\mathcal{P}$  can play a central role in signal discrimination
- Signals that cannot be distinguished at  $2\sigma$  using SNR can be distinguished using  $\mathcal{P}$  at much  $>3\sigma$  in some cases