

Multiprocess imaging of nuclear modifications on parton distributions in proton-nucleus collisions

Meng-Quan Yang

Collaborated with Peng Ru, Ben-Wei Zhang

M.-Q. Yang, P. Ru and B.-W. Zhang, Phys. Rev. D 112 (2025) no.7, 074008

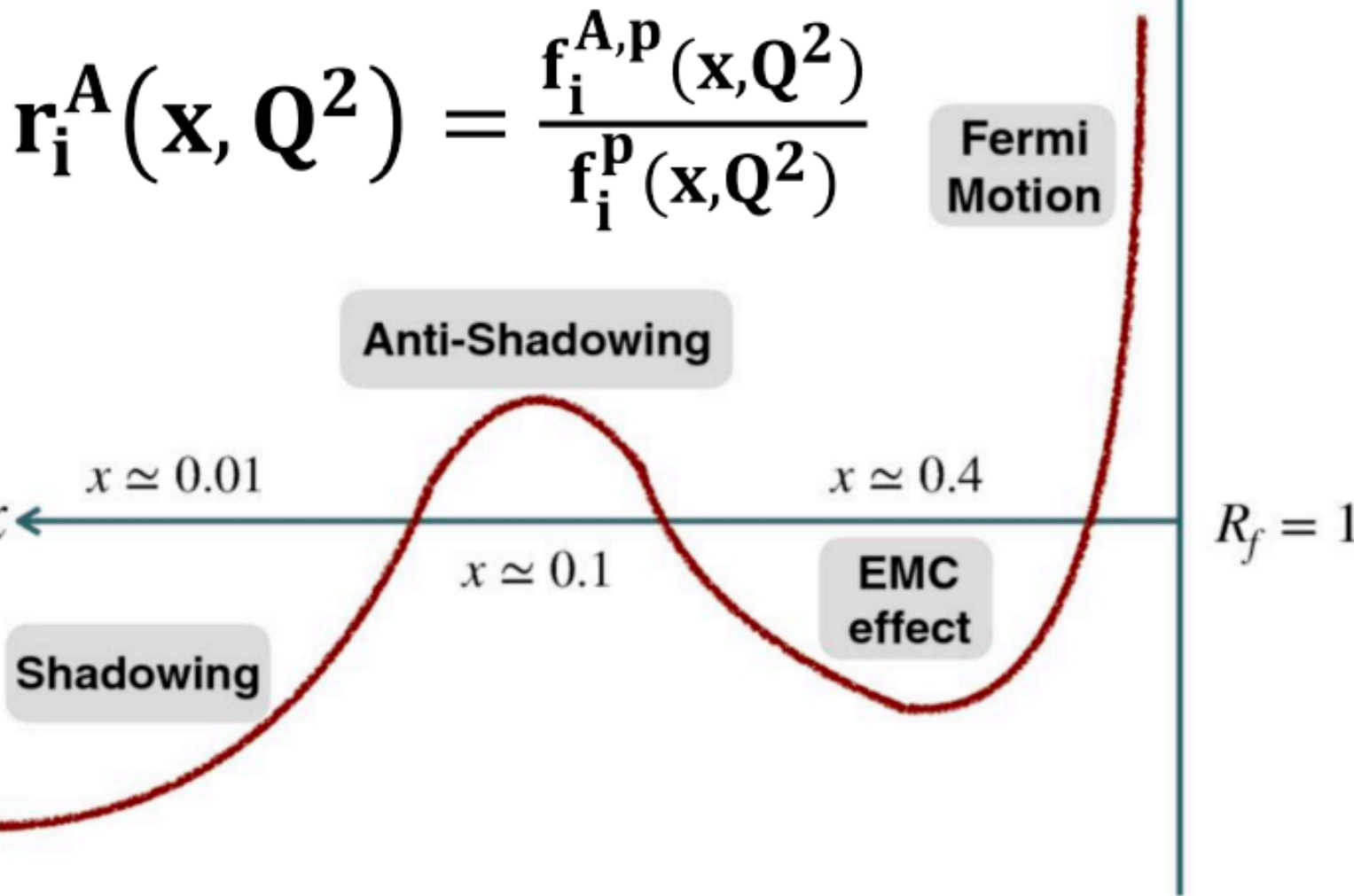
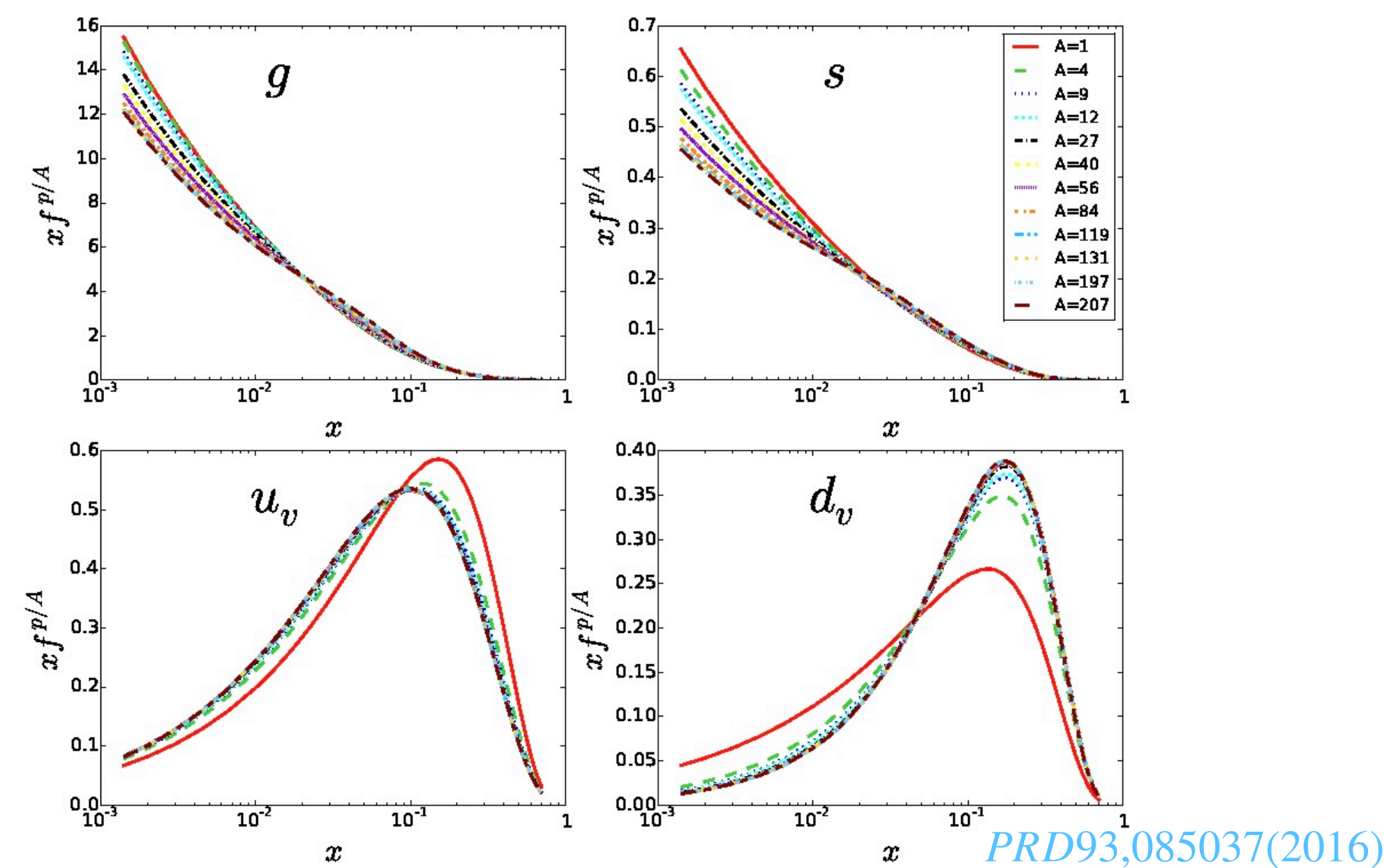
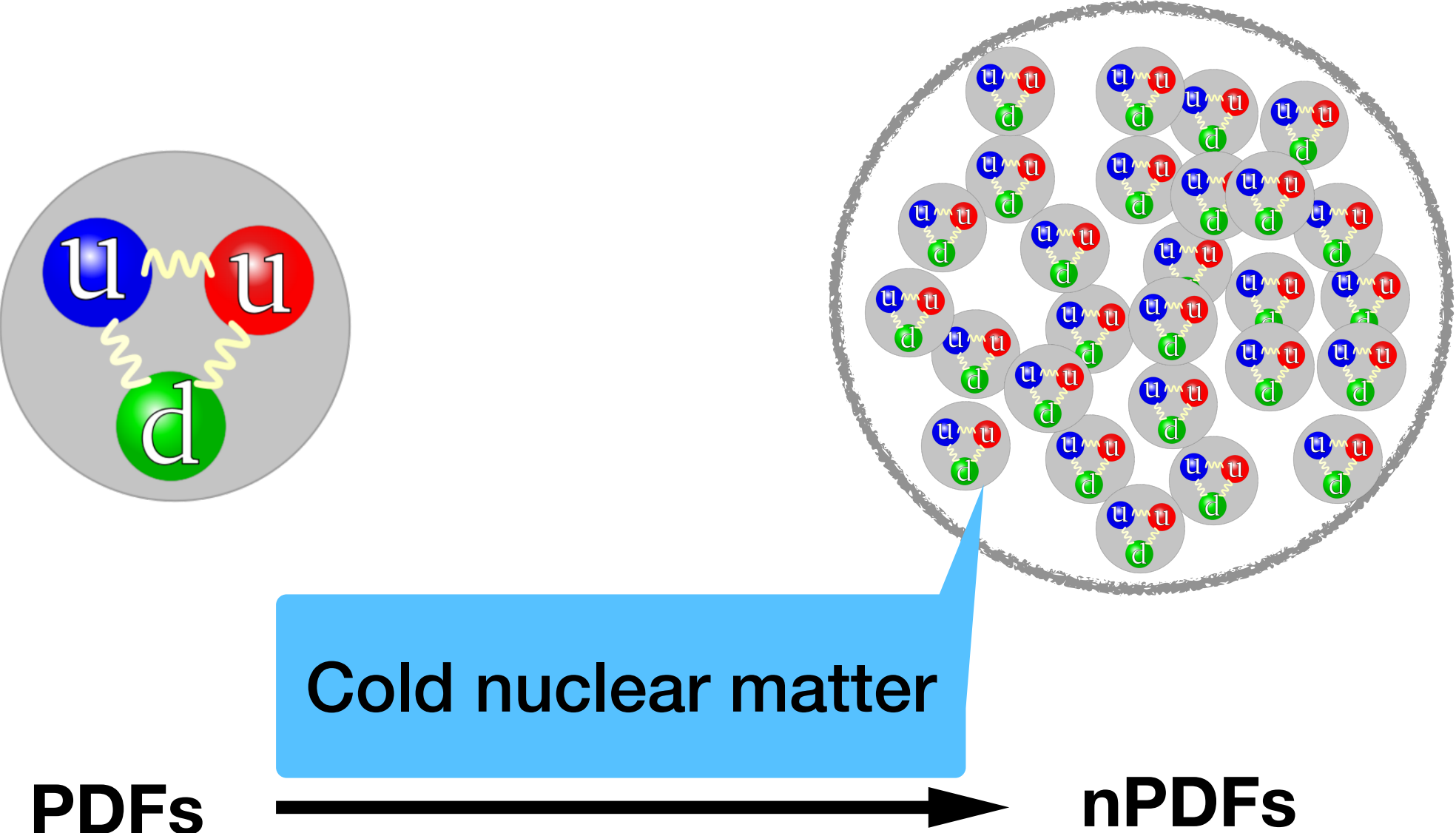
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Outline

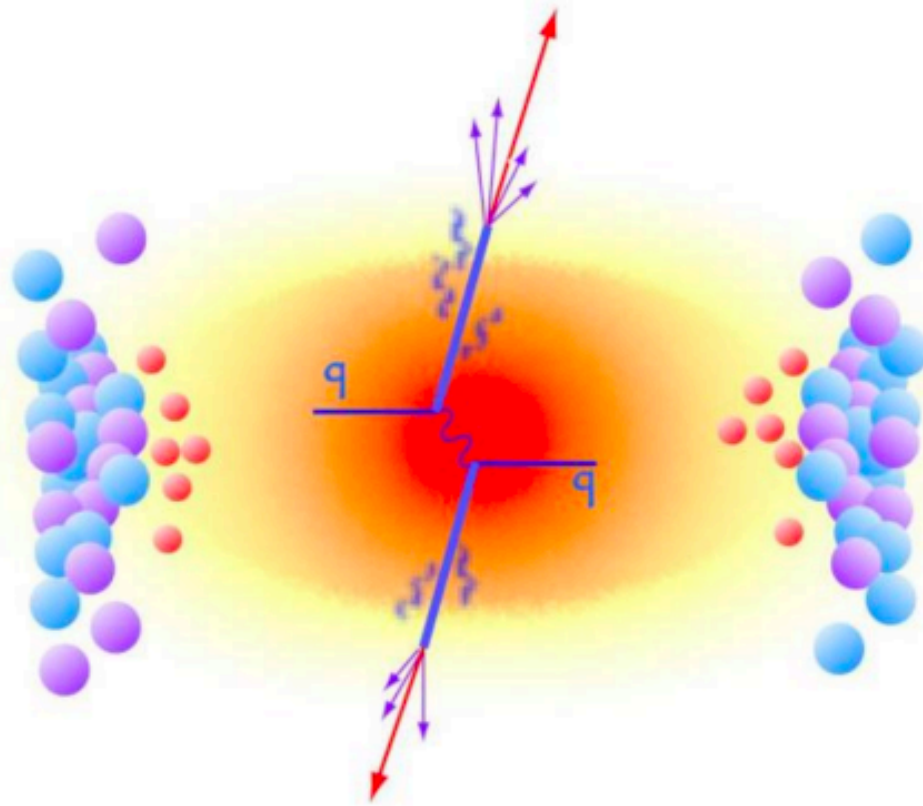


- Introduction
- Imaging $r_i^A(x, Q^2)$ in pA collisions
- Summary

Introduction



Nuclear modification ratios of collinear PDFs.



An essential baseline for disentangling final-state nuclear matter effects probed by hard particles.

Introduction



Eur. Phys. J. C (2017) 77:163

Table 1 The data sets used in the EPPS16 analysis, listed in the order of growing nuclear mass number. The number of data points and their contribution to χ^2 counts only those data points that fall within the kinematic cuts explained in the EPS09 analysis are marked with

Experiment	Observable	Collisions	Data points
SLAC E139	DIS	$e^- \text{He}(4), e^- \text{D}$	21
CERN NMC 95, re	DIS	$\mu^- \text{He}(4), \mu^- \text{D}$	16
CERN NMC 95	DIS	$\mu^- \text{Li}(6), \mu^- \text{D}$	15
CERN NMC 95, Q^2 dep	DIS	$\mu^- \text{Li}(6), \mu^- \text{D}$	153
SLAC E139	DIS	$e^- \text{Be}(9), e^- \text{D}$	20
CERN NMC 96	DIS	$\mu^- \text{Be}(9), \mu^- \text{C}$	15
SLAC E139	DIS	$e^- \text{C}(12), e^- \text{D}$	7
CERN NMC 95	DIS	$\mu^- \text{C}(12), \mu^- \text{D}$	15
CERN NMC 95, Q^2 dep	DIS	$\mu^- \text{C}(12), \mu^- \text{D}$	165
CERN NMC 95, re	DIS	$\mu^- \text{C}(12), \mu^- \text{D}$	16
CERN NMC 95, re	DIS	$\mu^- \text{C}(12), \mu^- \text{Li}(6)$	20
FNAL E772	DY	$p\text{C}(12), p\text{D}$	9
SLAC E139	DIS	$e^- \text{Al}(27), e^- \text{D}$	20
CERN NMC 96	DIS	$\mu^- \text{Al}(27), \mu^- \text{C}(12)$	15
SLAC E139	DIS	$e^- \text{Ca}(40), e^- \text{D}$	7
FNAL E772	DY	$p\text{Ca}(40), p\text{D}$	9
CERN NMC 95, re	DIS	$\mu^- \text{Ca}(40), \mu^- \text{D}$	15
CERN NMC 95, re	DIS	$\mu^- \text{Ca}(40), \mu^- \text{Li}(6)$	20
CERN NMC 95	DIS	$\mu^- \text{Ca}(40), \mu^- \text{C}(12)$	15
SLAC E139	DIS	$e^- \text{Fe}(56), e^- \text{D}$	26
FNAL E772	DY	$e^- \text{Fe}(56), e^- \text{D}$	9
CERN NMC 96	DIS	$\mu^- \text{Fe}(56), \mu^- \text{C}(12)$	15
FNAL E866	DY	$p\text{Fe}(56), p\text{Be}(9)$	28
CERN EMC	DIS	$\mu^- \text{Cu}(64), \mu^- \text{D}$	19
SLAC E139	DIS	$e^- \text{Ag}(108), e^- \text{D}$	7
CERN NMC 96	DIS	$\mu^- \text{Sn}(117), \mu^- \text{C}(12)$	15
CERN NMC 96, Q^2 dep	DIS	$\mu^- \text{Sn}(117), \mu^- \text{C}(12)$	144
FNAL E772	DY	$p\text{W}(184), p\text{D}$	9
FNAL E866	DY	$p\text{W}(184), p\text{Be}(9)$	28
CERN NA10 ^a	DY	$\pi^- \text{W}(184), \pi^- \text{D}$	10
FNAL E615 ^a	DY	$\pi^+ \text{W}(184), \pi^- \text{W}(184)$	11
CERN NA3 ^a	DY	$\pi^- \text{Pt}(195), \pi^- \text{H}$	7
SLAC E139	DIS	$e^- \text{Au}(197), e^- \text{D}$	21
RHIC PHENIX	π^0	$d\text{Au}(197), pp$	20
CERN NMC 96	DIS	$\mu^- \text{Pb}(207), \mu^- \text{C}(12)$	15
CERN CMS ^a	W^\pm	$p\text{Pb}(208)$	10
CERN CMS ^a	Z	$p\text{Pb}(208)$	6
CERN ATLAS ^a	Z	$p\text{Pb}(208)$	7
CERN CMS ^a	dijet	$p\text{Pb}(208)$	7
CERN CHORUS ^a	DIS	$\nu\text{Pb}(208), \bar{\nu}\text{Pb}(208)$	824
Total			1811

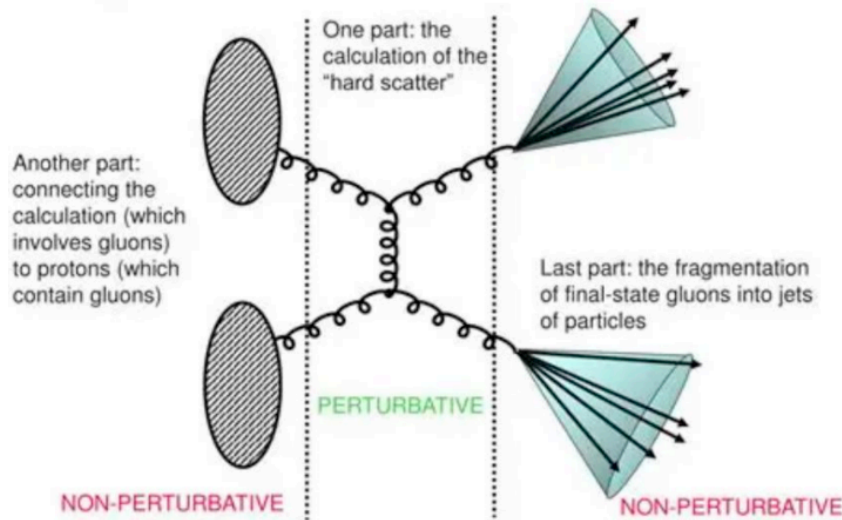
Both PDFs and their nuclear modification ratios rely on the global QCD analyses of diverse experimental data.

$$r_i^A(x, Q^2) = \frac{f_i^{A,p}(x, Q^2)}{f_i^p(x, Q^2)}$$

Challenge in global analyses:

For a realistic observable, the dependencies on x , Q^2 , and i are intricately convoluted in calculations with collinear factorization in perturbative QCD.

$$d\sigma = \sum_{abd} f_{a/A}(x_a, \mu_F) \otimes f_{b/B}(x_b, \mu_F) \otimes d\hat{\sigma}_{ab \rightarrow cd}(x_a P_A, x_b P_B, \mu_F, \mu_R) \otimes D_{c \rightarrow h}(z_h, \mu_F)$$



Introduction



Challenge in global analyses:

An analogy to solving a system of equations

Lower degree of variable mixing results in faster solving.

$$\begin{cases} a_{11}x_1 + a_{12}x_2 + \cdots + a_{1n}x_n = O_1 \\ a_{21}x_1 + a_{22}x_2 + \cdots + a_{2n}x_n = O_2 \\ \vdots \\ a_{n1}x_1 + a_{n2}x_2 + \cdots + a_{nn}x_n = O_n \end{cases}$$



$$\{x_1, x_2, \cdots, x_n\}$$

$$\begin{cases} a_{11}x_1 = O_1 \\ a_{21}x_1 + a_{22}x_2 = O_2 \\ \vdots \\ a_{nn}x_n = O_n \end{cases}$$



$$\{x_1, x_2, \cdots, x_n\}$$

Introduction



Challenge in global analyses:

An analogy to solving a system of equations

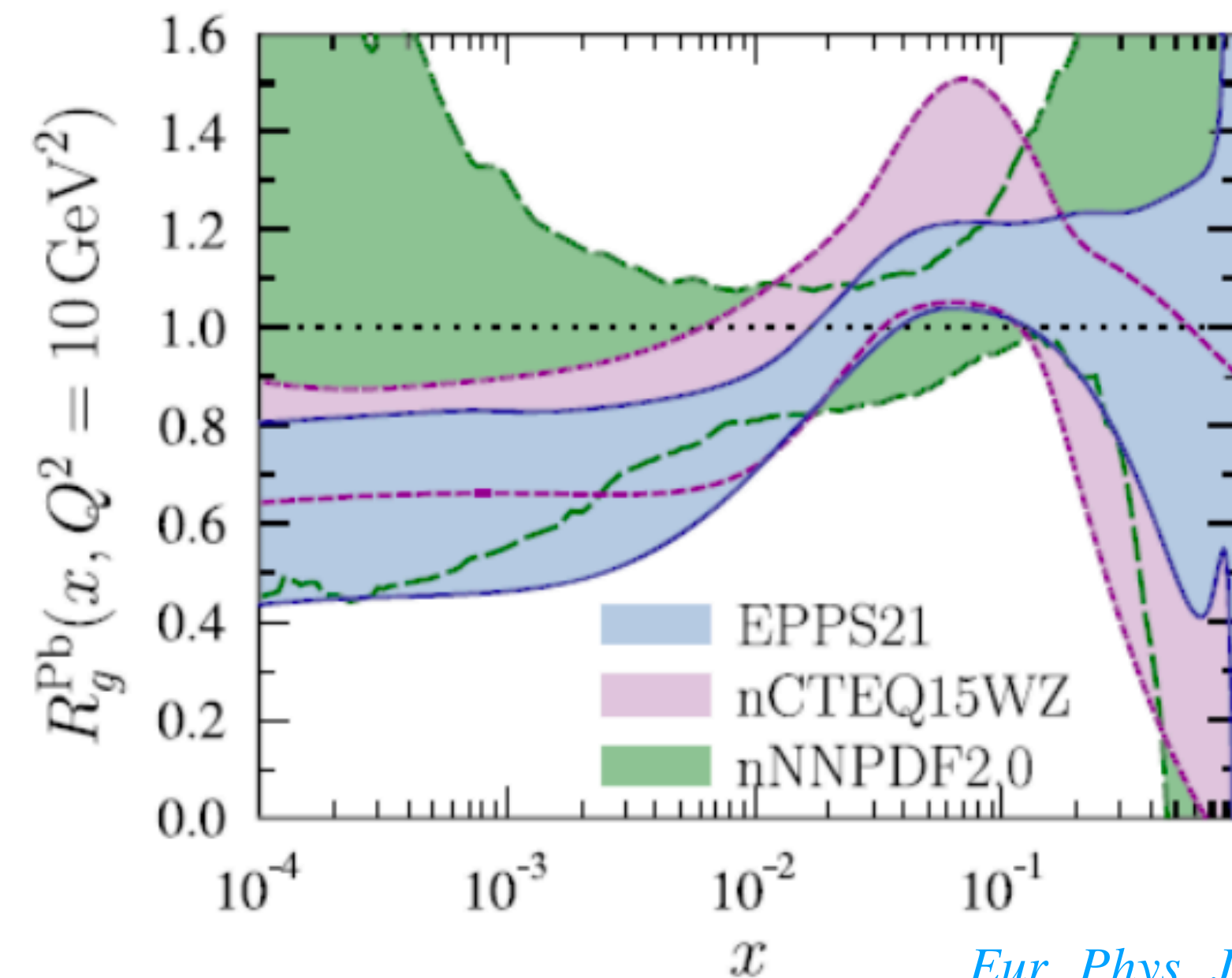
$$\begin{cases} a_{11}x_1 + a_{12}x_2 + \cdots + a_{1n}x_n = O_1 \\ a_{21}x_1 + a_{22}x_2 + \cdots + a_{2n}x_n = O_2 \\ \vdots \\ a_{n1}x_1 + a_{n2}x_2 + \cdots + a_{nn}x_n = O_n \end{cases}$$



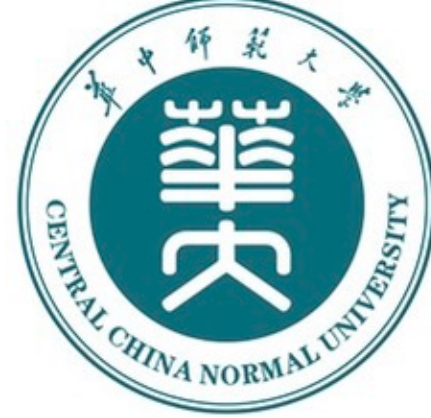
$$\{x_1, x_2, \cdots, x_n\}$$

Challenge from mixing contributions of variables

1. Overfitting.
2. Parameters degeneracy.
3. Complicated uncertainty propagation.
4. Slow and unstable convergence.

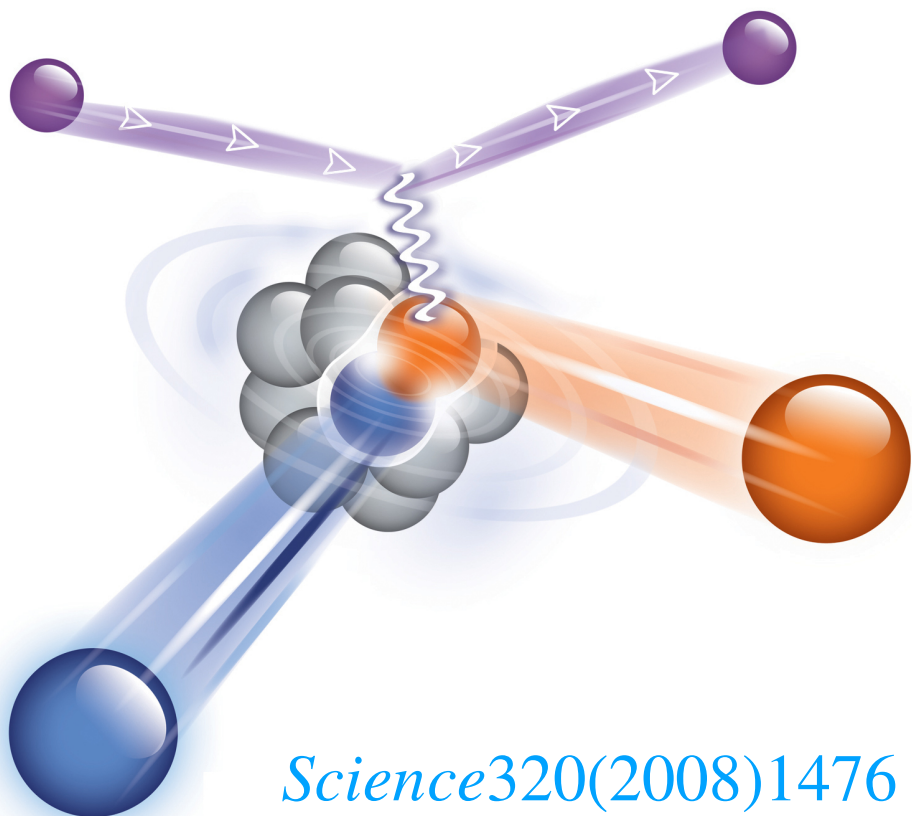


Introduction



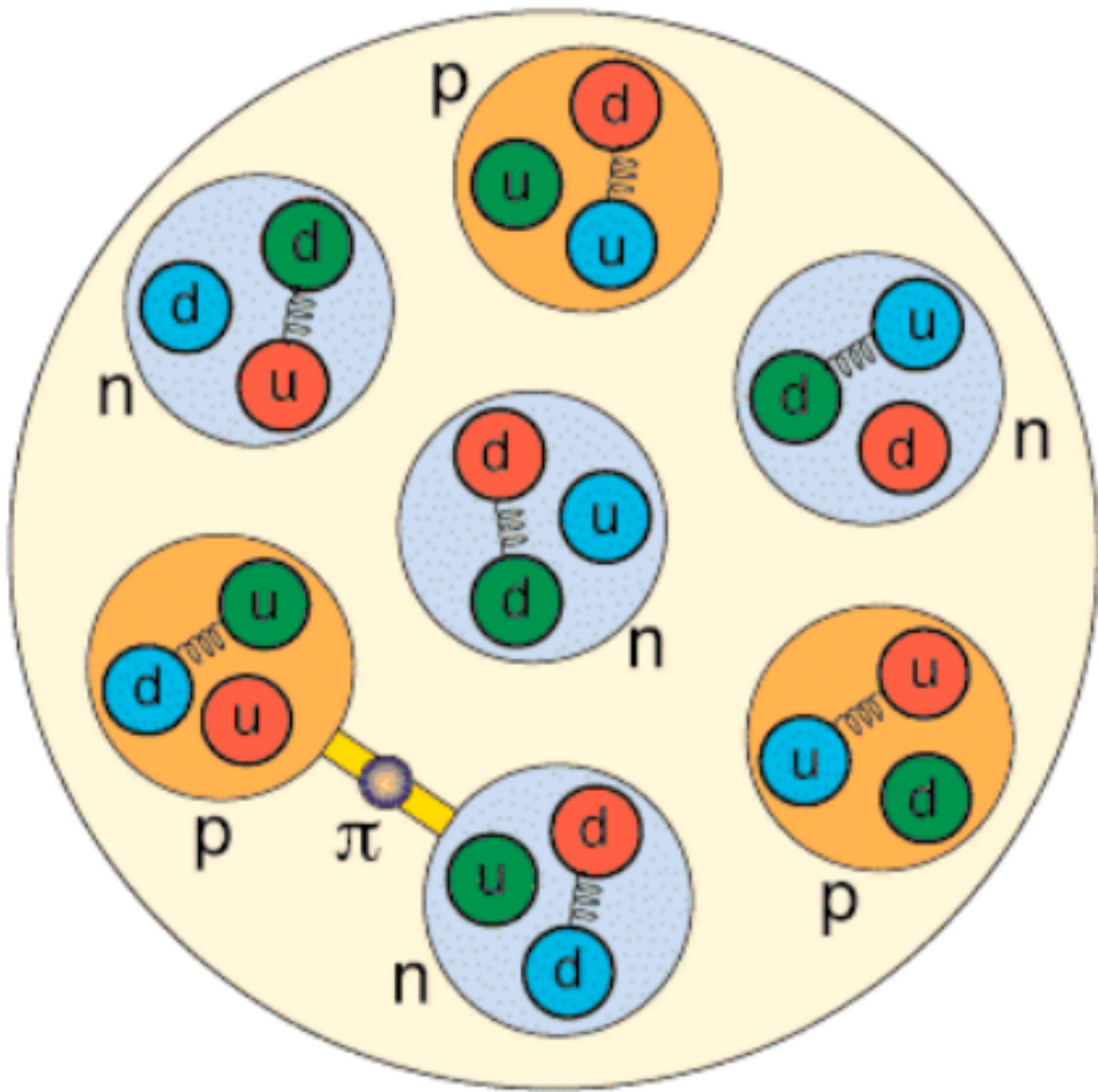
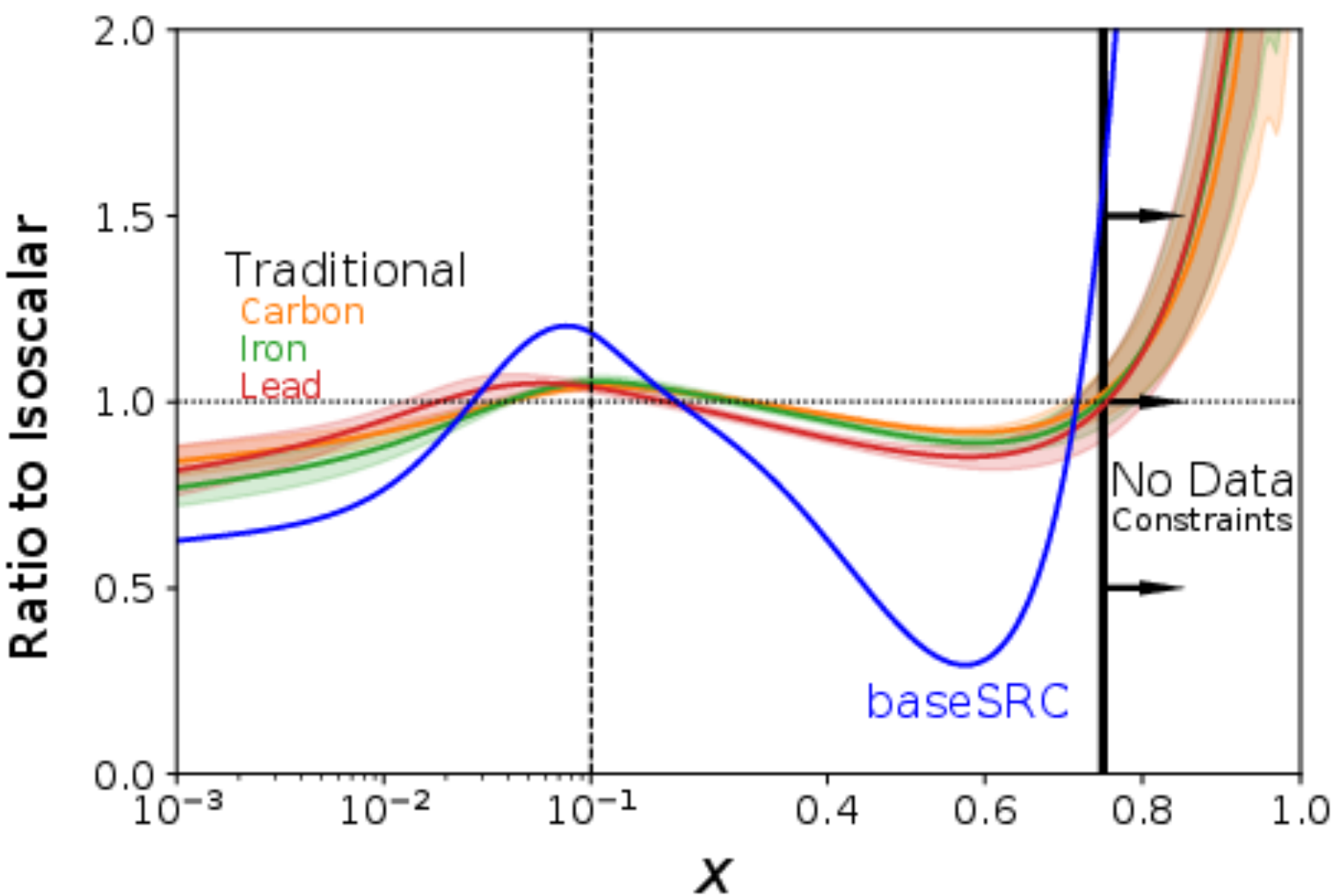
Insights from nuclear binding dynamics are useful.

Short-Range Correlations



*Science*320(2008)1476

*Phys. Rev. Lett.*133(2024)15,152502



Meson exchange current
Off-shell corrections
Coherent nuclear shadowing
Fermi motion

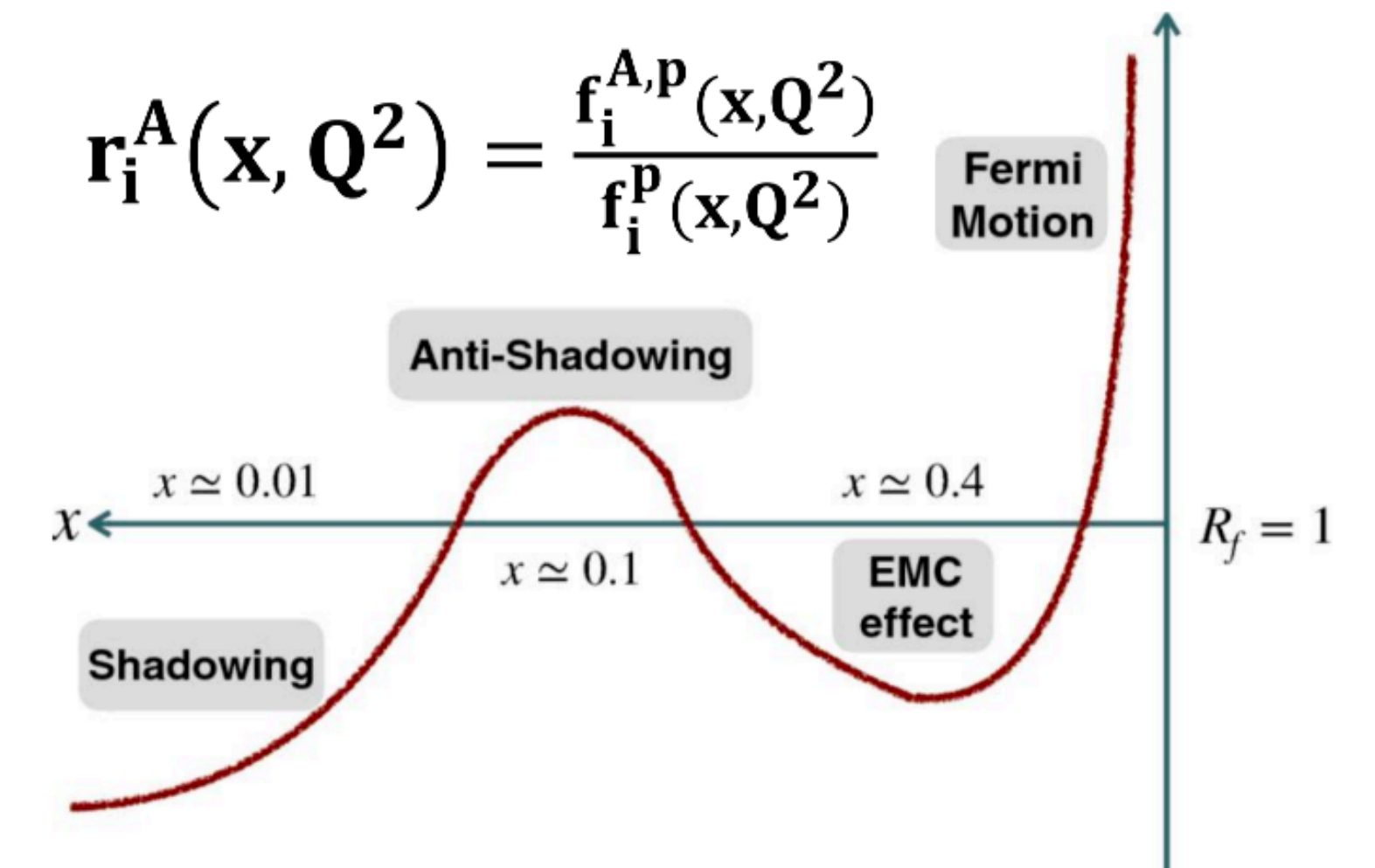
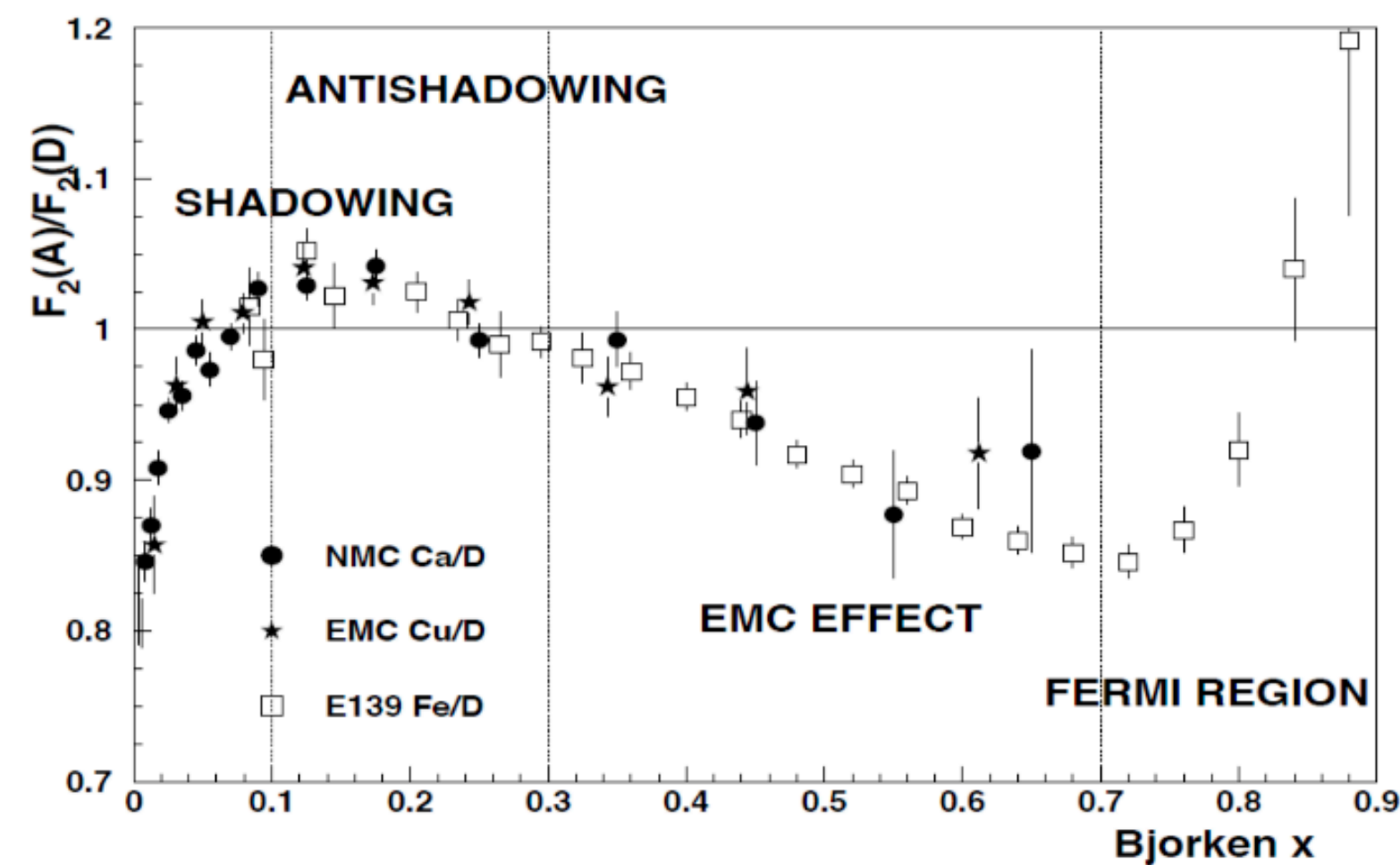
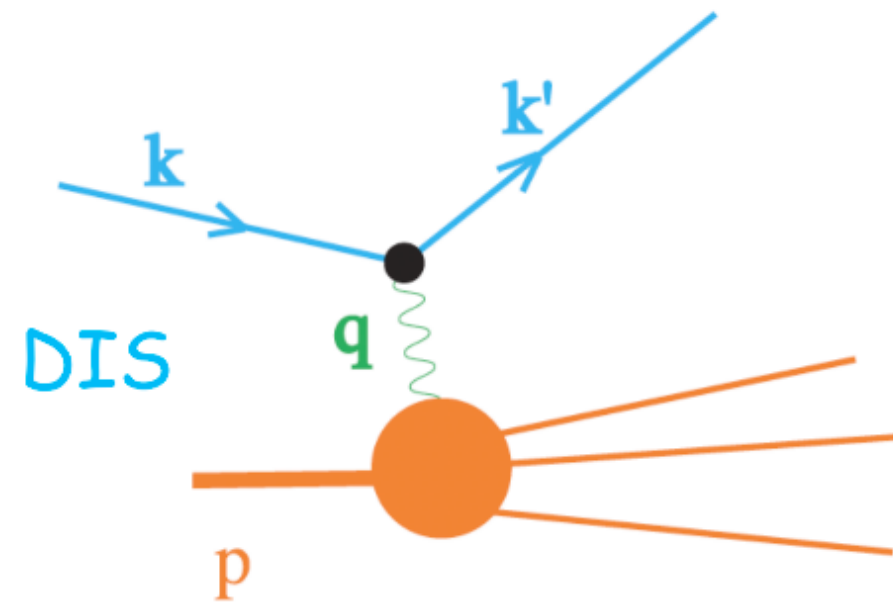
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*Phys. Rev. C*90(2014)4,045204

Introduction



Data in DIS serve as an effective image of $r_i^A(x, Q^2)$



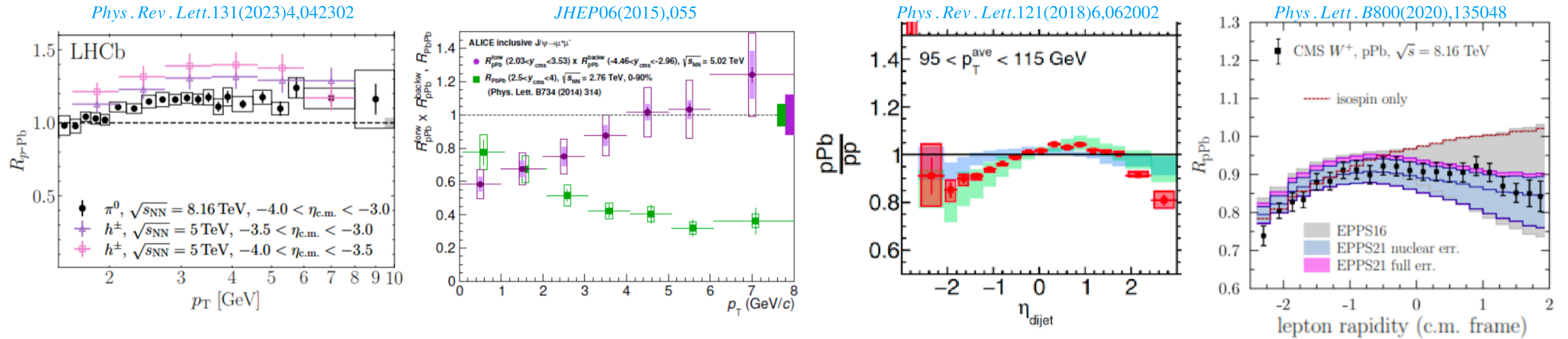
Measurement of such observable with low degree of variable mixing:

1. inspire parametrizations of modifications
2. improve the efficiency of global analysis.

Motivation



Status for pA collisions at the LHC



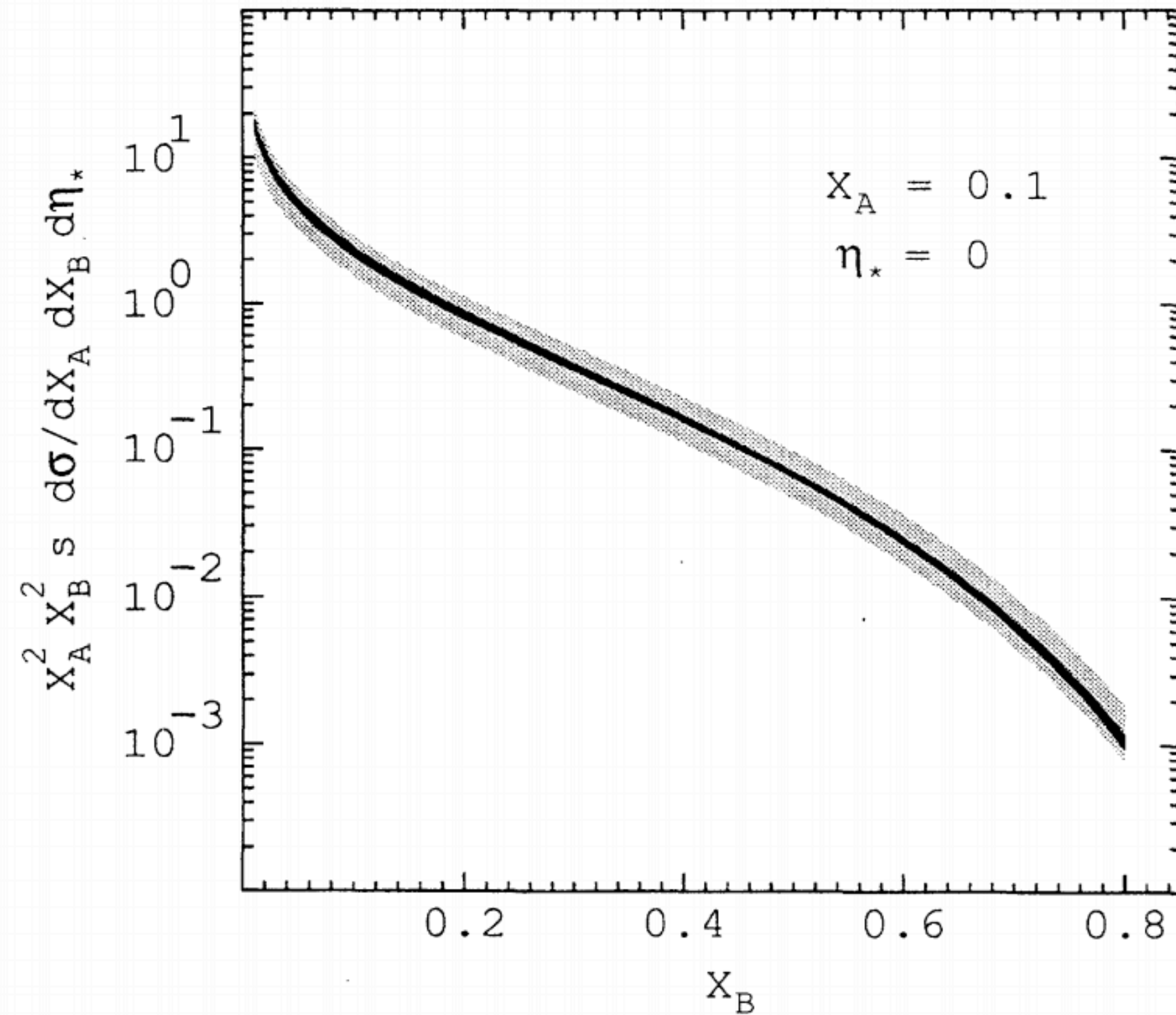
However, the imaging of $r_i^A(x, Q^2)$ achieved in DIS has not been replicated for most LHC processes.

Can we image $r_i^A(x, Q^2)$ at LHC?

Construction of the image of $r_i^A(x, Q^2)$ in pA



Triple differential dijet cross section

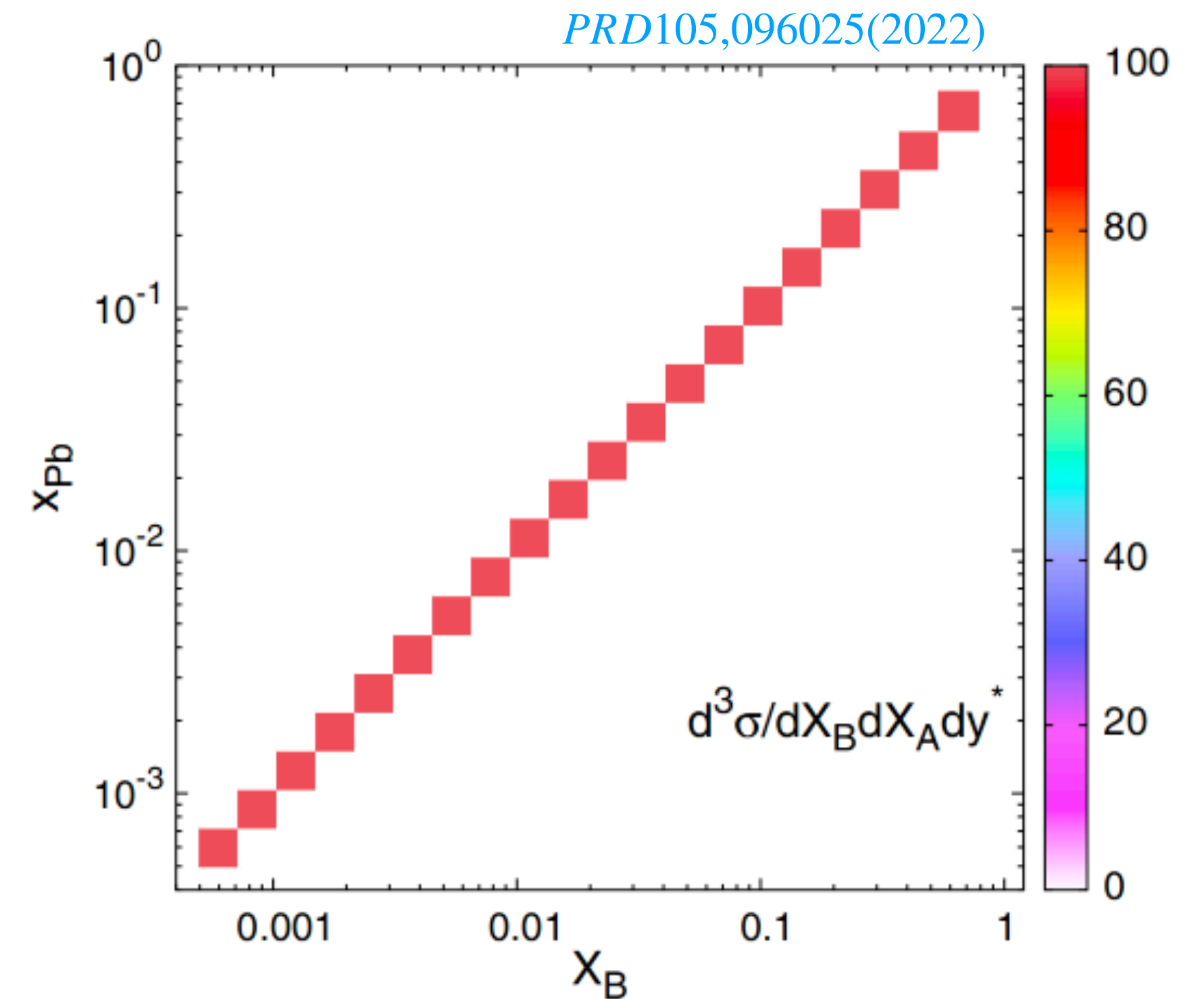


Phys. Rev. Lett. 74(1995), 5182 – 5185

$$V^{(3)} = \{X_B, X_A, y^*\}$$

$$X_A = \sum_{n \in \text{dijets}} \frac{E_{Tn}}{\sqrt{s}} e^{+y_n}$$

$$X_B = \sum_{n \in \text{dijets}} \frac{E_{Tn}}{\sqrt{s}} e^{-y_n}$$

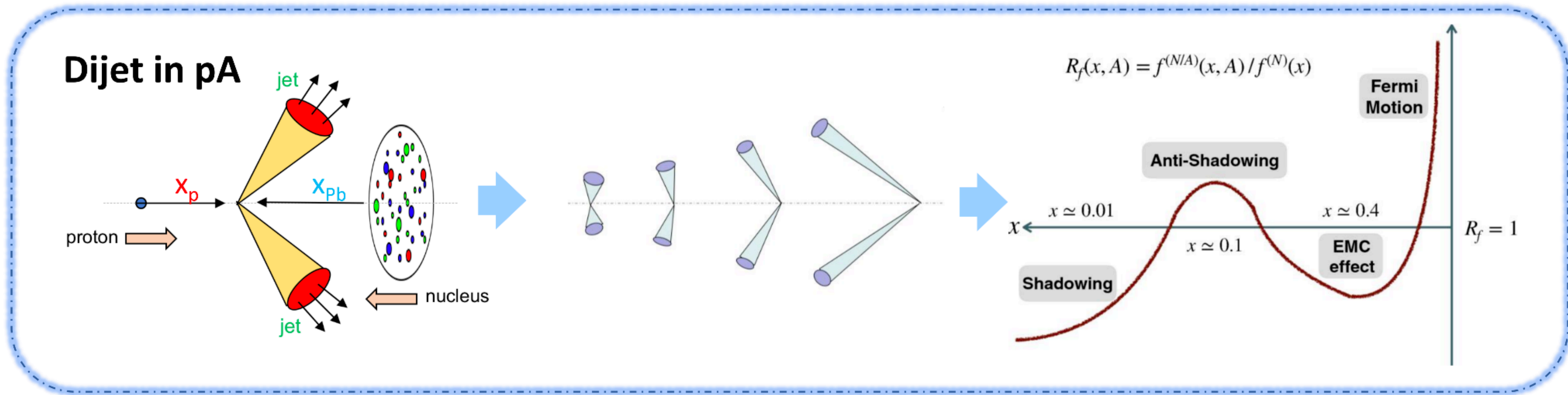


PRD 105, 096025 (2022)

LO: $X_A = x_p$ control the probe

$X_B = x_{Pb}$ scan the target

Construction of the image of $r_i^A(x, Q^2)$ in pA



Shen, Ru, Zhang, PRD 105, 096025 (2022)

More advantages in pPb collisions

$$R_{pA}(v_1, v_2, v_3) \approx \frac{\sum_{a,b} f_a^p(x_a, \mu^2) f_b^A(x_b, \mu^2) H_{ab}(v_1, v_2, v_3)}{\sum_{a,b} f_a^p(x_a, \mu^2) f_b^p(x_b, \mu^2) H_{ab}(v_1, v_2, v_3)}$$

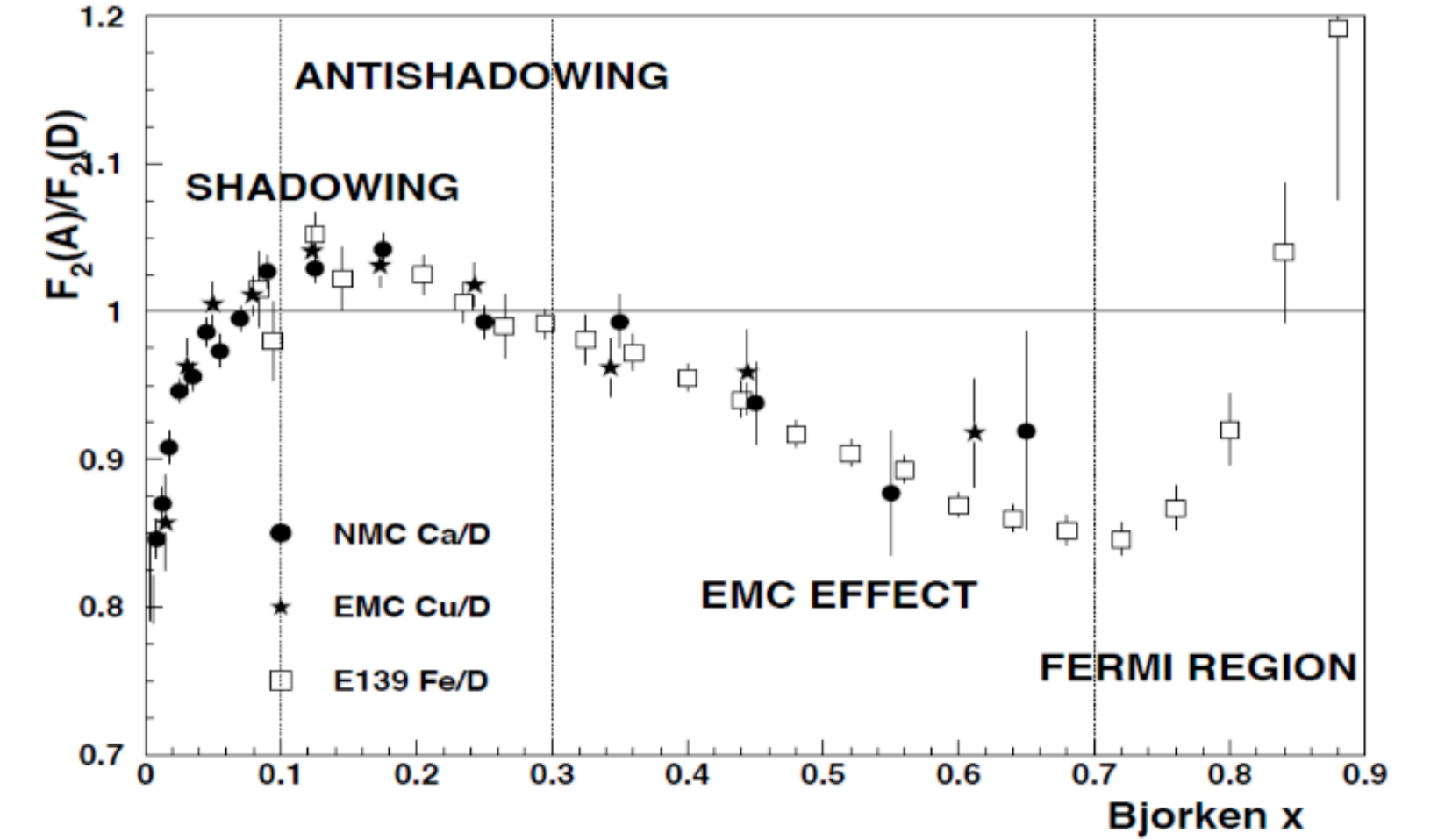
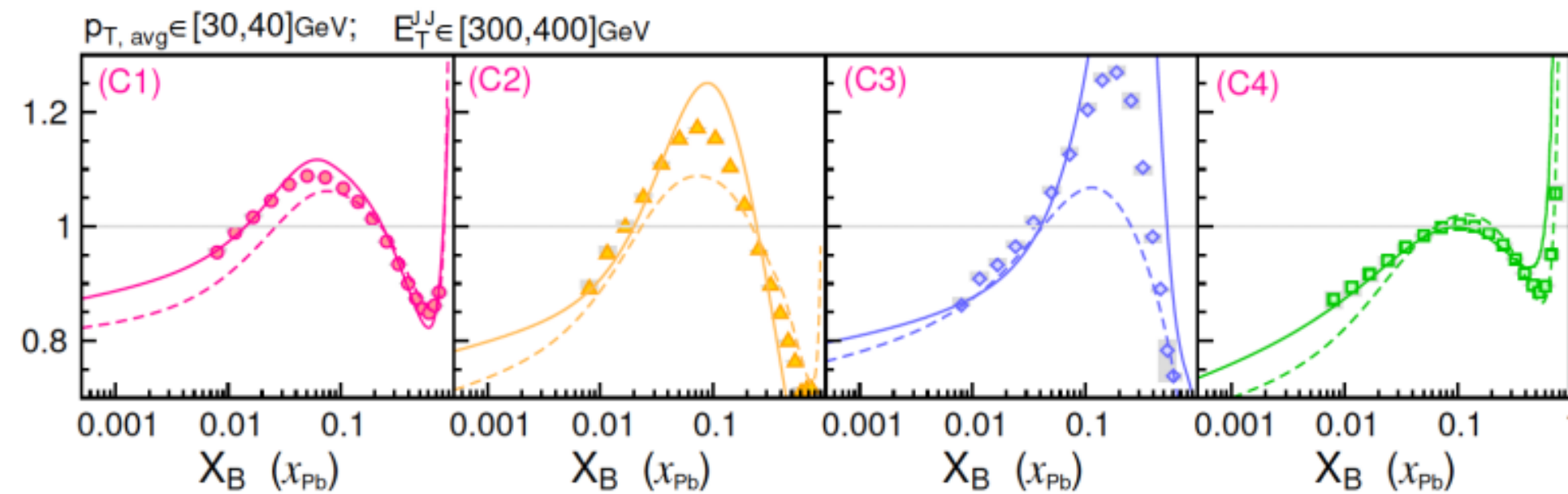
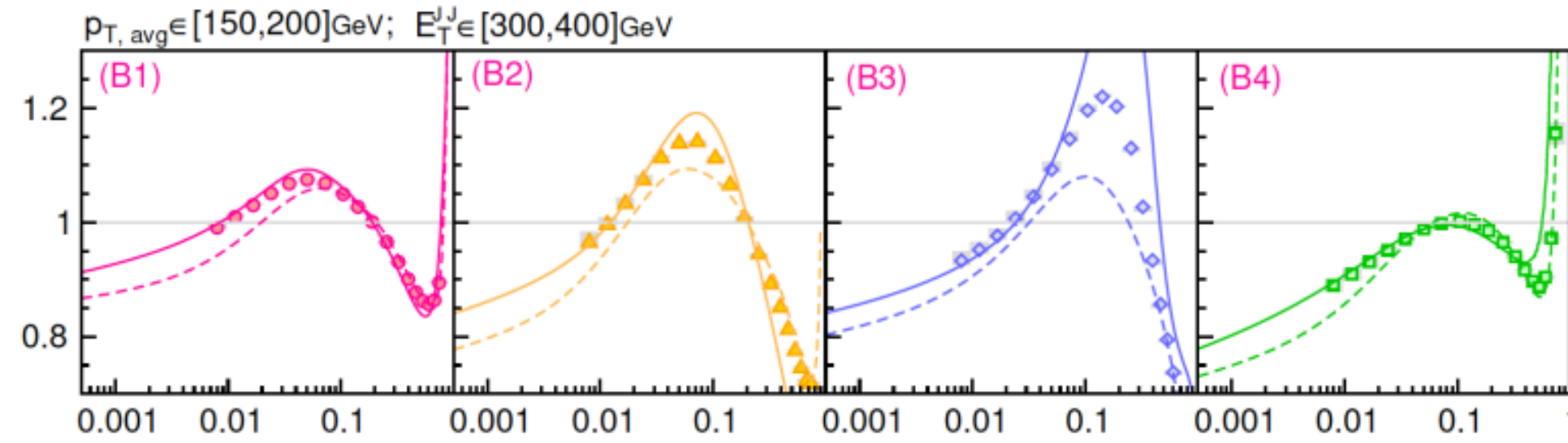
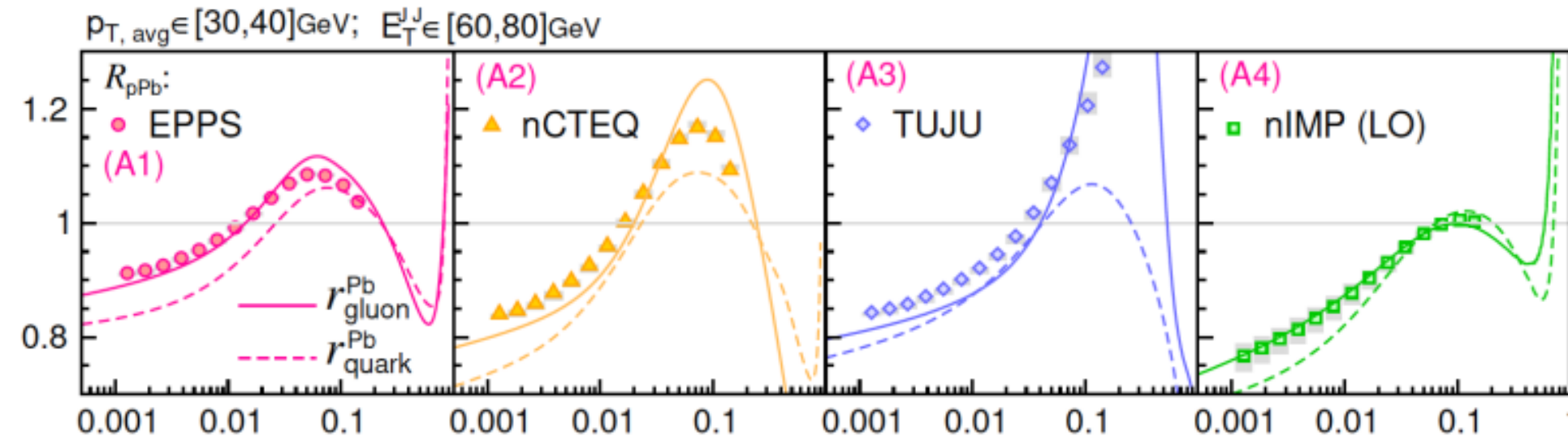
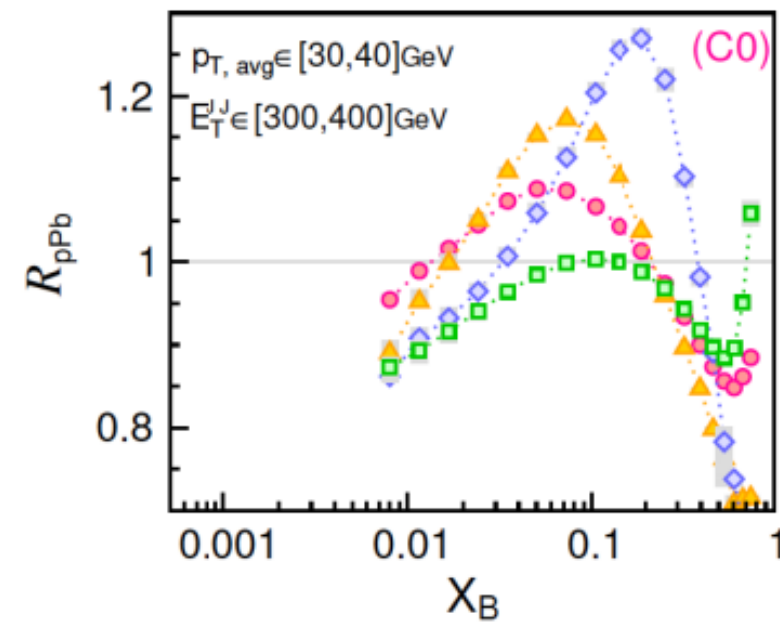
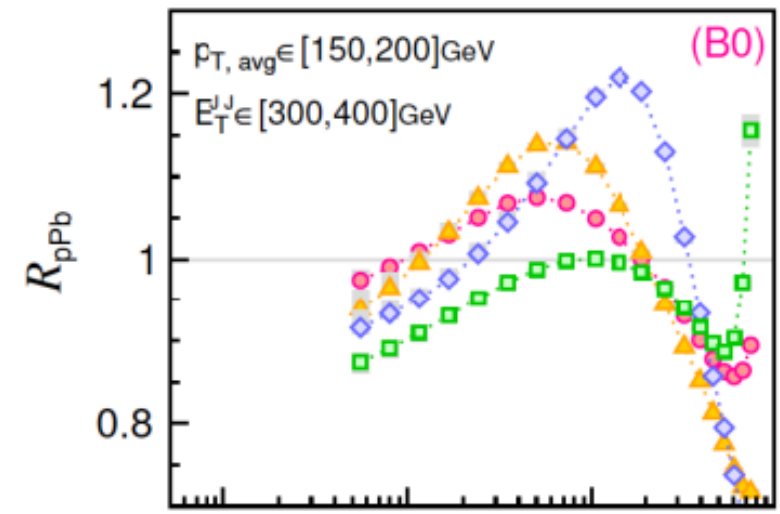
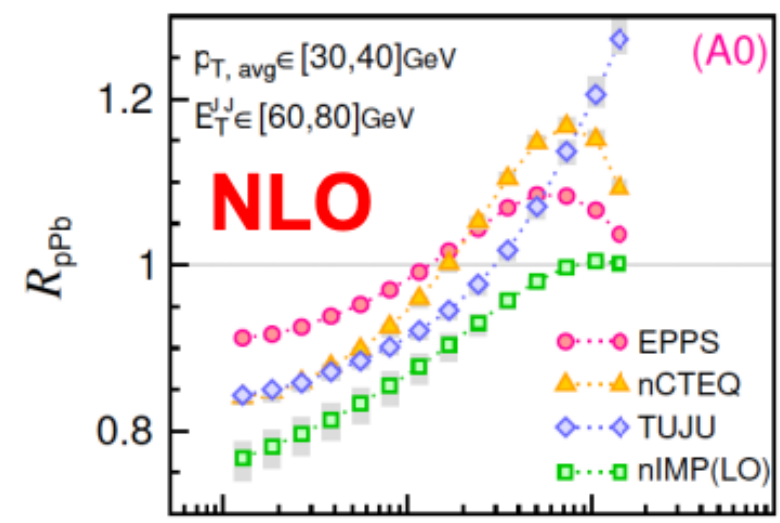
$r_i^A(x, Q^2)$

1. Uncertainties from proton PDFs reduced.
2. Uncertainties from high-order corrections reduced.
3. Incoming from one side.
4. Experimental uncertainties also reduced.

Construction of the image of $r_i^A(x, Q^2)$ in pA



Extended to cases with fixed probing scales. $V^{(3)} = \{X_B, E^{JJ}, p_{T,avg}\}$



1. Scanning of $r_i^A(x, Q^2)$ in pA, verified at NLO.
2. An analogy to the image in DIS.
3. However, parton flavors are still mixed.

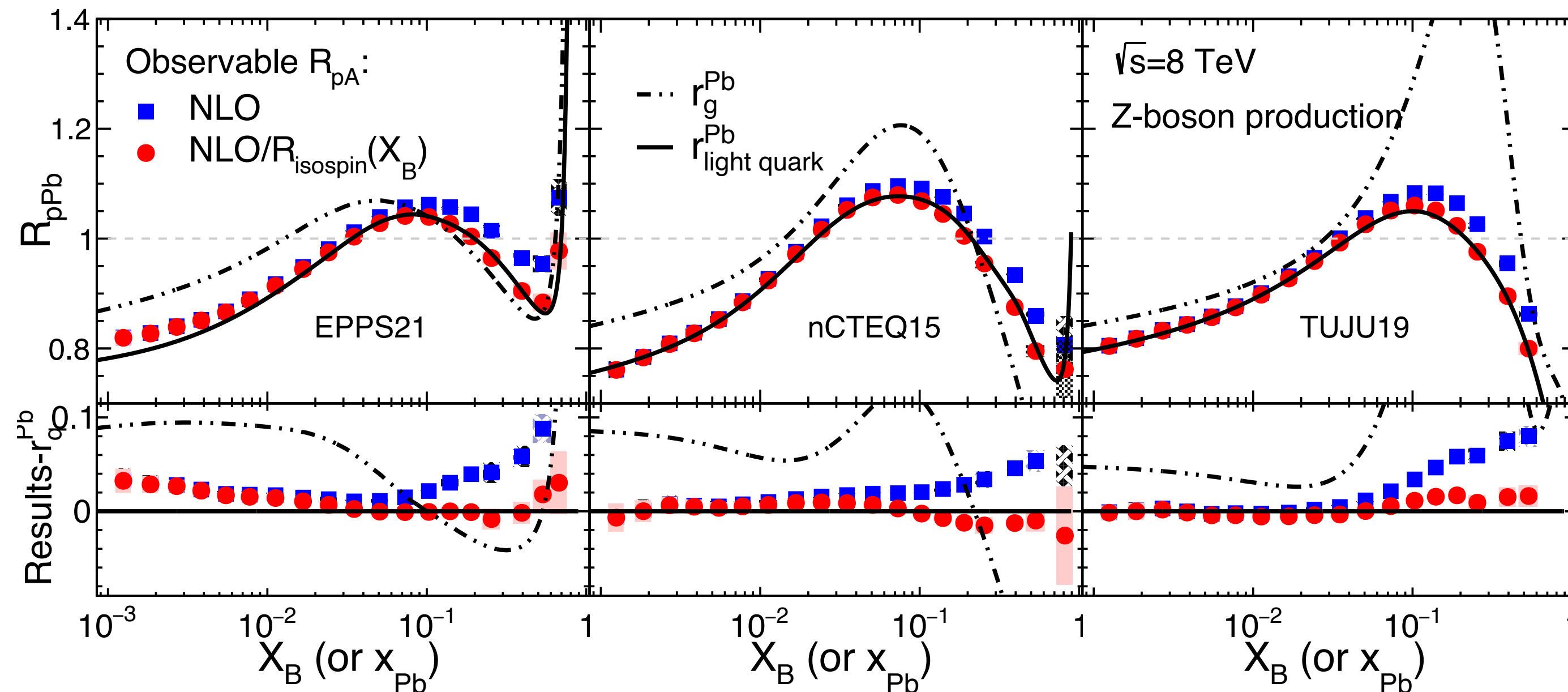
Construction of the image of $r_i^A(x, Q^2)$ in pA



Can parton flavors be further separated?

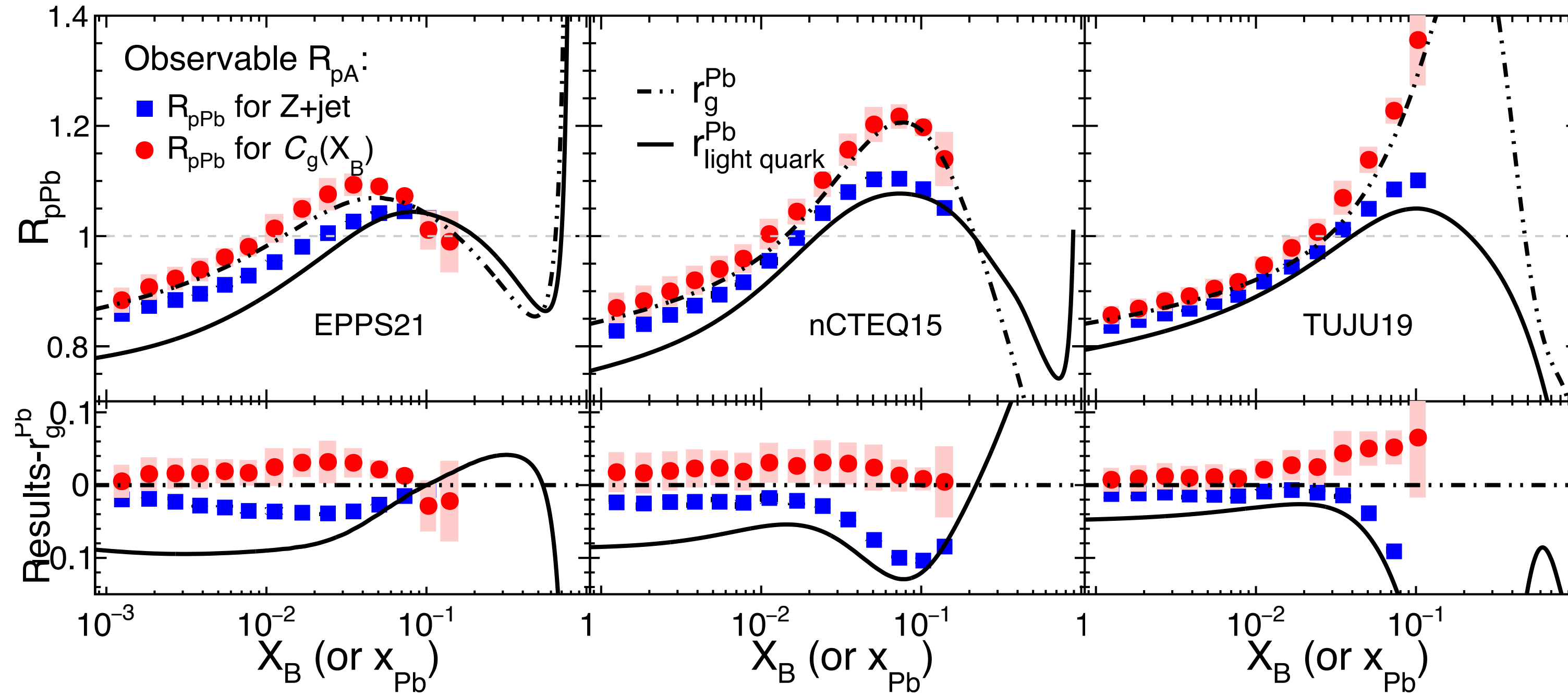
Multiprocess imaging based on the scanning with (x, Q^2)

Process	Partonic subprocess at LO
Z-boson	$q + \bar{q} \rightarrow Z \rightarrow l^+ l^-$
Z + jet	$q + \bar{q} \rightarrow Z + g$ $q(\bar{q}) + g \rightarrow Z + q(\bar{q})$
Z + c-jet	$c(\bar{c}) + g \rightarrow Z + c(\bar{c})$



R_{pA} for Z-boson productions nicely images the $r_i^A(x, Q^2)$ for light quarks!

Construction of the image of $r_i^A(x, Q^2)$ in pA



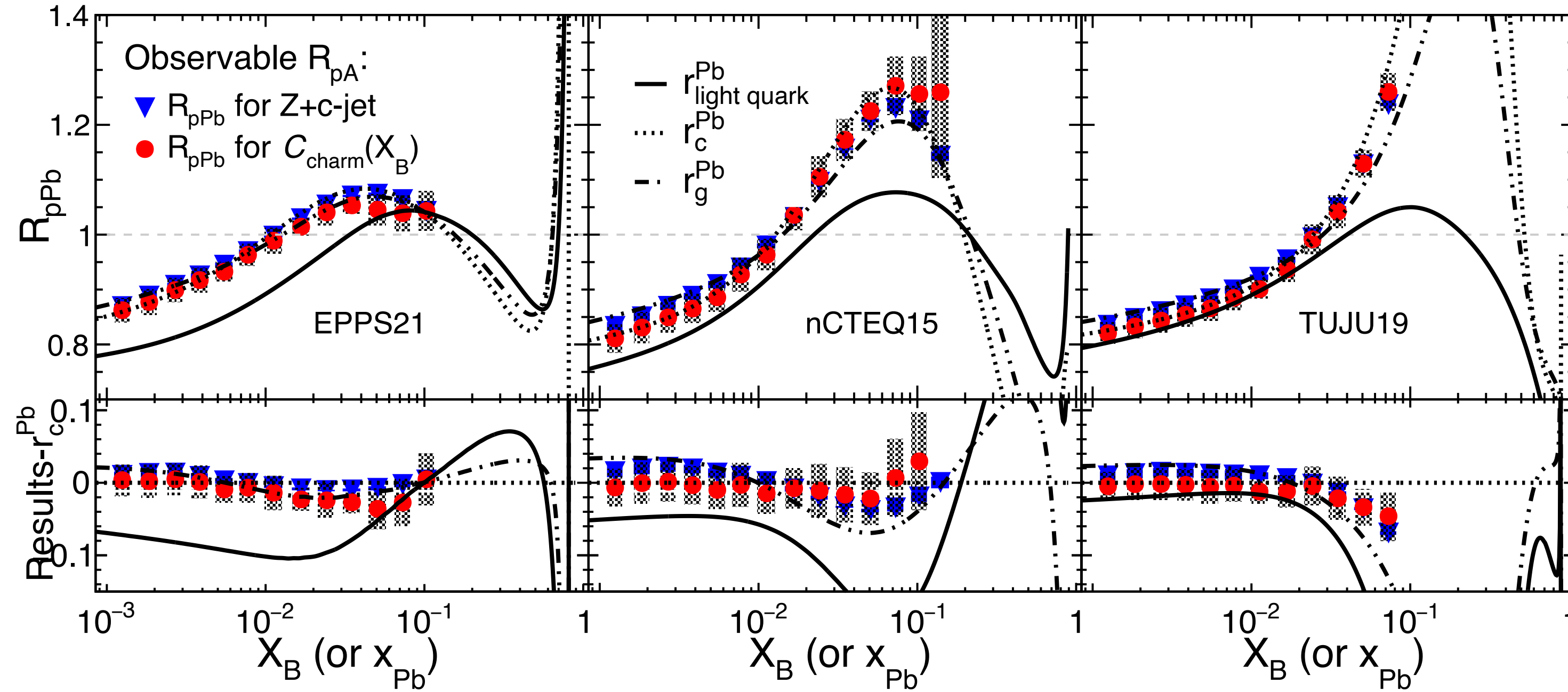
R_{pA} for Z+jet productions
images the mixture of $r_i^A(x, Q^2)$
for light quarks and gluons.

A combined observable with Z-boson and Z+jet productions:

$$C_g(X_B) = \kappa_1(X_B) \times d\sigma^{Z+\text{jet}}(X_B) - d\sigma^Z(X_B), \quad \kappa_1(X_B) = \frac{d\sigma^Z(X_B)}{[d\sigma^{Z+\text{jet}}(X_B)]_{\text{nuclear quark}}}$$

which suppress the effects from nuclear quarks and nicely images the $r_i^A(x, Q^2)$ for gluons.

Construction of the image of $r_i^A(x, Q^2)$ in pA



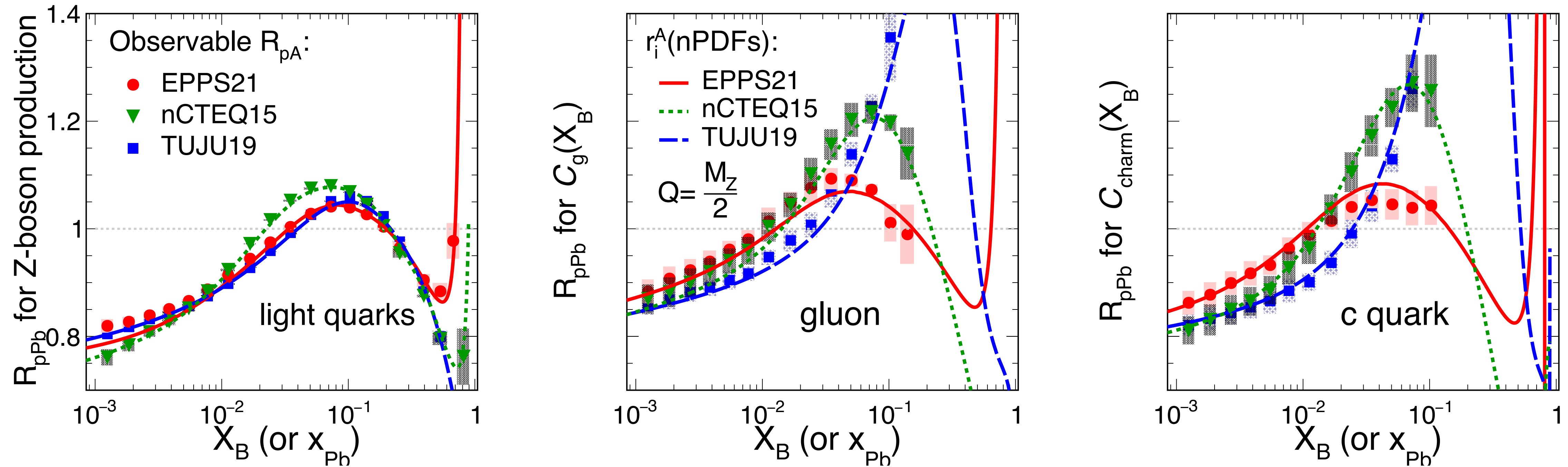
R_{pA} for Z+c-jet productions
images the mixture of $r_i^A(x, Q^2)$
for charm quarks and gluons.

A combined observable with Z-boson, Z+jet and Z+c-jet productions:

$$C_{\text{charm}}(X_B) = \kappa_2(X_B) \times d\sigma^{Z+c\text{-jet}}(X_B) - C_g(X_B), \quad \kappa_2(X_B) = \frac{C_g(X_B)}{[d\sigma^{Z+c\text{-jet}}(X_B)]_{\text{nuclear gluon}}}$$

which nicely images the $r_i^A(x, Q^2)$ for charm quarks.

Construction of the image of $r_i^A(x, Q^2)$ in pA



The results show the probability to separately image the $r_i^A(x, Q^2)$ for certain parton flavors at LHC.

Summary



- An imaging methodology is developed to optimize the future measurements at LHC.

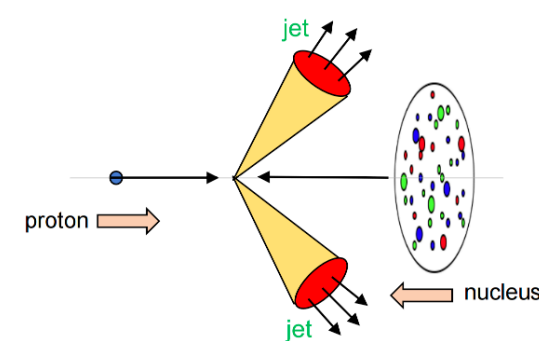
With better disentangled contributions of (x, Q^2, i) for $r_i^A(x, Q^2)$, facilitate more efficient global analysis .

- Neither PDFs nor $r_i^A(x, Q^2)$ is directly measurable.

Compared to traditional observables, the proposed imaging observables provide somewhat preprocessed data with theoretical guidance.

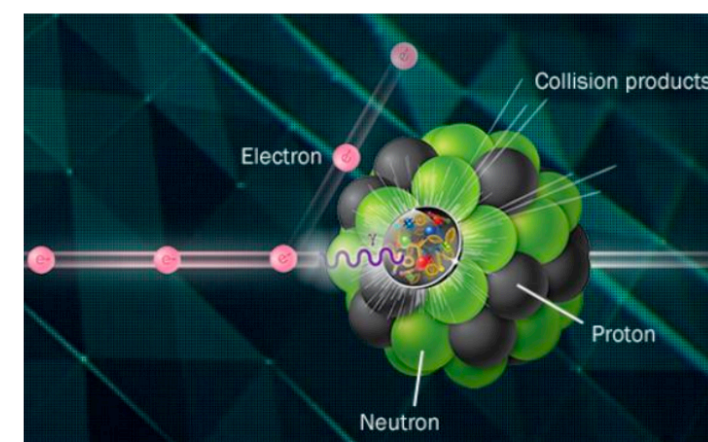
- Future applications:

Other processes at the LHC



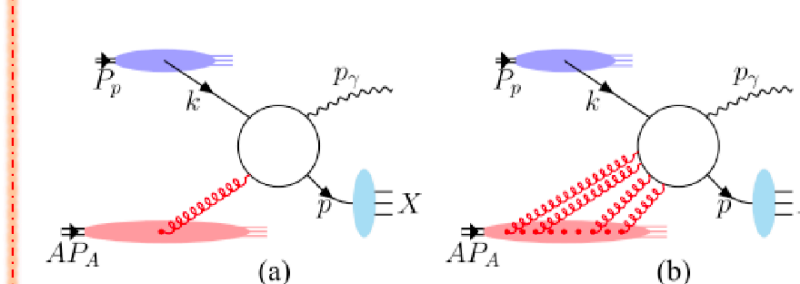
Dijet, Drell-Yan, γ -jet,

Combined study of LHC and EicC/EIC



Systematic comparison
spanning from pA to eA

Nuclear effects at higher-twist



Fu, Kang, Salazar, Wang, Xing
PRL 135, 032301 (2025)