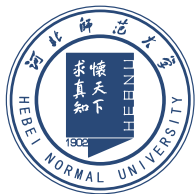


Evaluation of Muon $g-2$ from $\tau \rightarrow \pi\pi\nu$ Data

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- ① Current status on muon $g-2$
- ② The electromagnetic correction $G_{EM}(t)$
- ③ Updated evaluation of a_μ from τ data
- ④ Summary

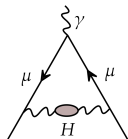
Current status on muon g-2

- The **Leading-order hadronic vacuum polarization** (LO HVP) contributes

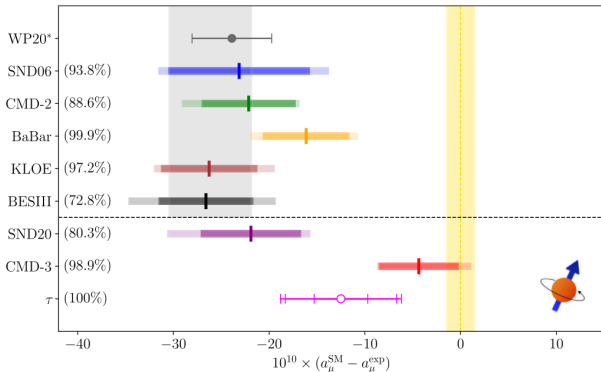
[Gourdin and Rafael, NPB'69.] [R. Aliberti et al., Phys. Rept'25]

$$a_{\mu}^{\text{HVP,LO}} = \frac{1}{4\pi^3} \int_{4m_{\pi}^2}^{\infty} dt K(t) \sigma_{e^+e^- \rightarrow \text{hadrons}}^0(t)$$

e^+e^- hadrons cross sections $\Rightarrow \pi\pi > 70\%$



- Data-driven $e^+e^- \rightarrow \pi^+\pi^-$ inputs show a sizable spread.



Alternative way to address HVP from $\pi\pi$

- $e^+e^- \rightarrow \pi^+\pi^-$ cross section

$$a_\mu^{\text{HVP,LO}} = \frac{1}{4\pi^3} \int_{4m_\pi^2}^{\infty} dt K(t) \sigma_{e^+e^- \rightarrow \text{hadrons}}^0(t)$$

$$\sigma_{e^+e^- \rightarrow \pi^+\pi^-}^0 = \frac{\pi\alpha^2}{3t} \beta_{\pi^+\pi^-}^3(t) \left| F_{\pi\pi}^{(0)}(t) \right|^2$$

- $\pi\pi$ invariant-mass distribution of τ data

$$\frac{d\Gamma_{\tau\pi\pi[\gamma]}}{dt} = \frac{G_F^2 |V_{ud}|^2 m_\tau^3 S_{\text{EW}}}{384\pi^3} \left(1 - \frac{t}{m_\tau^2}\right)^2 \left(1 + \frac{2t}{m_\tau^2}\right) \beta_{\pi^-\pi^0}^3 \left| F_{\pi\pi}^{(-)}(t) \right|^2 G_{\text{EM}}(t)$$

- Conserved Vector Current (CVC)**

$$\left. \begin{array}{l} e^+e^- \rightarrow \pi^+\pi^- : \quad \nu_0(t) \sim \beta_{\pi^+\pi^-}^3(t) \left| F_{\pi\pi}^{(0)}(t) \right|^2 \\ \tau^- \rightarrow \pi^-\pi^0 \nu_\tau : \quad \nu_-(t) \sim \beta_{\pi^-\pi^0}^3(t) \left| F_{\pi\pi}^{(-)}(t) \right|^2 \end{array} \right\} \Rightarrow \nu_0(t) = \nu_-(t) \text{ (CVC)}$$

$$\sigma_{e^+e^- \rightarrow \pi^+\pi^-}^0 = \left[\frac{K_\sigma(t)}{K_\Gamma(t)} \frac{d\Gamma_{\tau\pi\pi[\gamma]}}{dt} \right] \times \left[\frac{R_{\text{IB}}(t)}{S_{\text{EW}}} \right]$$

$$K_\sigma(t) = \frac{\pi\alpha^2}{3t}, \quad K_\Gamma(t) = \frac{G_F^2 |V_{ud}|^2 m_\tau^3}{384\pi^3} \left(1 - \frac{t}{m_\tau^2}\right)^2 \left(1 + 2\frac{t}{m_\tau^2}\right).$$

Current status on muon g-2

- Key problem: isospin breaking (IB) effects

$$a_{\mu}^{\text{HVP,LO}} = \frac{1}{4\pi^3} \int_{4m_{\pi}^2}^{\infty} dt K(t) \sigma_{e^{+}e^{-} \rightarrow \text{hadrons}}^0(t)$$
$$\sigma_{e^{+}e^{-} \rightarrow \pi^{+}\pi^{-}}^0 = \left[\frac{K_{\sigma}(t)}{K_{\Gamma}(t)} \frac{d\Gamma_{\tau\pi\pi[\gamma]}}{dt} \right] \times \left[\frac{R_{\text{IB}}(t)}{S_{\text{EW}}} \right]$$

- Correction term

$$R_{\text{IB}}(t) = \frac{\text{FSR}(t) \beta_{\pi^{+}\pi^{-}}^3(t) \left| F_{\pi\pi}^{(0)}(t) \right|^2}{G_{\text{EM}}(t) \beta_{\pi^{-}\pi^0}^3(t) \left| F_{\pi\pi}^{(-)}(t) \right|^2}$$

- $\pi\pi$ form factors ($F_{\pi\pi}^{(0,-)}(t)$)
- kinematical factors ($\beta_{+,0}$)
- final-state radiation (FSR)
- electroweak short-distance correction (S_{EW})
- **long-distance electromagnetic correction (G_{EM})**

The electromagnetic correction function $G_{EM}(t)$

- The EM correction function $G_{EM}(t)$ for the $\tau \rightarrow \pi\pi\nu_\tau$

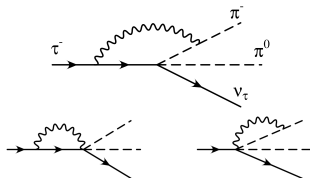
$$\frac{d\Gamma_{\tau\pi\pi[\gamma]}}{dt} = \frac{G_F^2 |V_{ud}|^2 m_\tau^3 S_{EW}}{384\pi^3} \left(1 - \frac{t}{m_\tau^2}\right)^2 \left(1 + \frac{2t}{m_\tau^2}\right) \beta_{\pi^-\pi^0}^3 \left|F_{\pi\pi}^{(-)}(t)\right|^2 G_{EM}(t)$$

- Non-radiative two-pion tau decay

$$\frac{d\Gamma_{\tau\pi\pi}^{(0)}}{dt} = \frac{G_F^2 |V_{ud}|^2 m_\tau^3 S_{EW}}{384\pi^3} \left(1 - \frac{t}{m_\tau^2}\right)^2 \left(1 + \frac{2t}{m_\tau^2}\right) \beta_{\pi^-\pi^0}^3 \left|F_{\pi\pi}^{(-)}(t)\right|^2$$

virtual photon (loops)

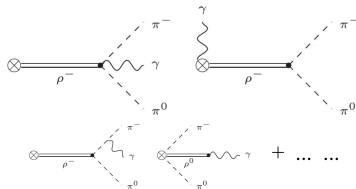
[Cirigliano, et al., PLB' 01]



$$\frac{d\Gamma_{\tau\pi\pi[\gamma]}}{dt} = \frac{d\Gamma_{\tau\pi\pi}^{(0)}}{dt} G_{EM}(t) \Rightarrow G_{EM}(t) = 1 + G_{EM}^{(v)}(t) + G_{EM}^{(r)}(t)$$

real photon

$(\tau \rightarrow \pi\pi\nu_\tau\gamma)$



The electromagnetic correction function $G_{EM}(t)$

- Amplitude of $\tau \rightarrow \pi\pi\nu\tau\gamma$

$$\mathcal{M} = eG_F V_{ud}^* \varepsilon^\mu(k)^* \left\{ F_\nu \bar{u}(q) \gamma^\nu (1 - \gamma_5) (m_\tau + \not{P} - \not{k}) \gamma_\mu u(P) + (V_{\mu\nu} - A_{\mu\nu}) \bar{u}(q) \gamma^\nu (1 - \gamma_5) u(P) \right\}$$

- The hadronic tensor amplitude

$$V_{\mu\nu} = F_{\pi\pi}^{(-)}(u) \frac{p_{1\mu}}{p_1 \cdot k} (p_1 + k - p_2)_\nu - F_{\pi\pi}^{(-)}(u) g_{\mu\nu} + \frac{F_{\pi\pi}^{(-)}(u) - F_{\pi\pi}^{(-)}(t)}{(p_1 + p_2) \cdot k} (p_1 + p_2)_\mu (p_2 - p_1)_\nu + v_1 (g_{\mu\nu} p_1 \cdot k - p_{1\mu} k_\nu) + v_2 (g_{\mu\nu} p_2 \cdot k - p_{2\mu} k_\nu) + v_3 (p_{1\mu} p_2 \cdot k - p_{2\mu} p_1 \cdot k) p_{1\nu} + v_4 (p_{1\mu} p_2 \cdot k - p_{2\mu} p_1 \cdot k) (p_1 + p_2 + k)_\nu$$

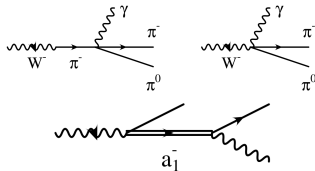
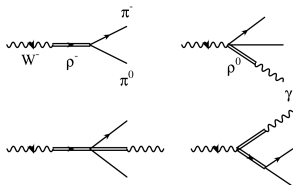
$$A_{\mu\nu} = i \left(a_1 \epsilon_{\mu\nu\rho\sigma} p_1^\rho k^\sigma + a_2 \epsilon_{\mu\nu\rho\sigma} p_2^\rho k^\sigma + a_3 p_{1\nu} \epsilon_{\mu\rho\beta\sigma} k^\rho p_1^\beta p_2^\sigma + a_4 p_{2\nu} \epsilon_{\mu\rho\beta\sigma} k^\rho p_1^\beta p_2^\sigma \right)$$

$$F_\nu = (p_2 - p_1)_\nu F_{\pi\pi}^{(-)}(t) / (2P \cdot k)$$

- Form Factor $F_{\pi\pi}^{(-)}(t)$ and $v_i(a_i)$**

Minimal resonance chiral theory contributions (CEN)

[V. Cirigliano, G. Ecker, and H. Neufeld, JHEP'02]

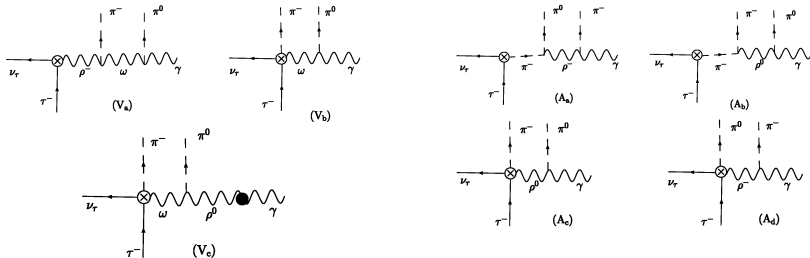


The electromagnetic correction function $G_{EM}(t)$

• Contributions from VVP and VJP operators in resonance chiral theory

[P. D. Ruiz-Femenia, et al., JHEP'03] [C.Chen, C.-G.Duan and Z.-H.Guo, JHEP'22]

$$\begin{aligned}
 \mathcal{L}_{VVP} &= d_1 \epsilon_{\mu\nu\rho\sigma} \langle \{V^{\mu\nu}, V^{\rho\alpha}\} \nabla_\alpha u^\sigma \rangle + id_2 \epsilon_{\mu\nu\rho\sigma} \langle \{V^{\mu\nu}, V^{\rho\sigma}\} \chi_- \rangle \\
 &\quad + d_3 \epsilon_{\mu\nu\rho\sigma} \langle \{ \nabla_\alpha V^{\mu\nu}, V^{\rho\alpha} \} u^\sigma \rangle + d_4 \epsilon_{\mu\nu\rho\sigma} \langle \{ \nabla^\sigma V^{\mu\nu}, V^{\rho\alpha} \} u_\alpha \rangle, \\
 \mathcal{L}_{VJP} &= \frac{c_1}{M_V} \epsilon_{\mu\nu\rho\sigma} \langle \{V^{\mu\nu}, f_+^{\rho\alpha}\} \nabla_\alpha u^\sigma \rangle + \frac{c_2}{M_V} \epsilon_{\mu\nu\rho\sigma} \langle \{V^{\mu\alpha}, f_+^{\rho\sigma}\} \nabla_\alpha u^\nu \rangle \\
 &\quad + \frac{ic_3}{M_V} \epsilon_{\mu\nu\rho\sigma} \langle \{V^{\mu\nu}, f_+^{\rho\sigma}\} \chi_- \rangle + \frac{ic_4}{M_V} \epsilon_{\mu\nu\rho\sigma} \langle V^{\mu\nu} [f_-^{\rho\sigma}, \chi_+] \rangle \\
 &\quad + \frac{c_5}{M_V} \epsilon_{\mu\nu\rho\sigma} \langle \{ \nabla_\alpha V^{\mu\nu}, f_+^{\rho\alpha} \} u^\sigma \rangle + \frac{c_6}{M_V} \epsilon_{\mu\nu\rho\sigma} \langle \{ \nabla_\alpha V^{\mu\alpha}, f_+^{\rho\sigma} \} u^\nu \rangle \\
 &\quad + \frac{c_7}{M_V} \epsilon_{\mu\nu\rho\sigma} \langle \{ \nabla^\sigma V^{\mu\nu}, f_+^{\rho\alpha} \} u_\alpha \rangle.
 \end{aligned}$$



- **High energy constraints** to the resonance couplings

$$c_1 + 4c_3 = 0 \quad c_1 - c_2 + c_5 = 0 \quad c_5 - c_6 = \frac{N_C M_V}{64\sqrt{2}\pi^2} F_V$$

$$d_1 + 8d_2 = -\frac{N_C M_V^2}{(8\pi F_V)^2} + \frac{F^2}{4F_V^2} \quad d_3 = -\frac{N_C M_V^2}{(8\pi F_V)^2} + \frac{F^2}{8F_V^2}$$

- Determination of low-energy coupling constant d_4

[C.Chen, C.-G.Duan and Z.-H.Guo, JHEP'22]

$$\omega \rightarrow \pi^0 \pi^0 \gamma$$

Scalar

- Conditions for cancelling ultraviolet divergence:

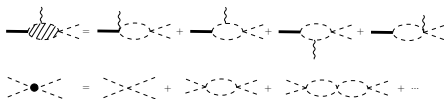
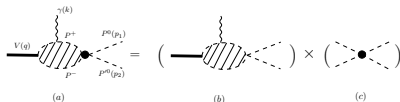
$$F_A = F_\pi, \quad F_V = \sqrt{2}F_\pi, \quad G_V = F_\pi/\sqrt{2}$$

- Determine d_4 from the radiative decay $\omega \rightarrow \pi^0 \pi^0 \gamma$

$$\Gamma_{\omega \rightarrow \pi^0 \pi^0 \gamma}^{\text{Exp}} = (5.8 \pm 1.0) \times 10^{-4} \text{ MeV}$$

$$d_4 = -0.42 \pm 0.07 \quad \text{Sol-A}$$

$$d_4 = 1.01 \pm 0.07 \quad \text{Sol-B}$$



- We are left with a parameter free theoretical amplitude for the $\tau \rightarrow \pi\pi\nu_\tau\gamma$ process

The electromagnetic correction function $G_{EM}(t)$

- Results for G_{EM}

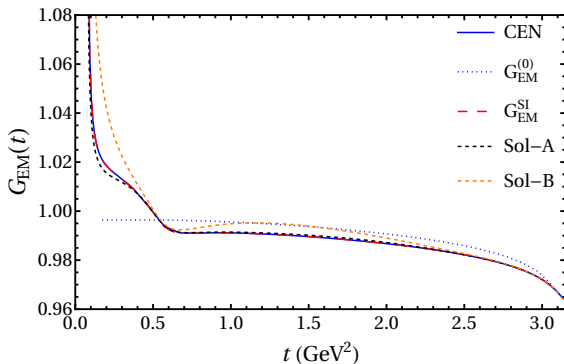
$$G_{EM}(t) = 1 + G_{EM}^{(v)}(t) + G_{EM}^{(r)}(t)$$

- real-photon correction function

$$G_{EM}^{(r)}(t) = \frac{1}{\frac{d\Gamma_{\tau\pi\pi}^{(0)}}{dt}} \frac{\pi^2}{32(2\pi)^{12}m_\tau^2} \int_{S_{\pi\pi\nu}^-}^{S_{\pi\pi\nu}^+} dS_{\pi\pi\nu} \int_{S_{\nu\gamma}^-}^{S_{\nu\gamma}^+} dS_{\nu\gamma} \int_{S_{\pi\nu}^-}^{S_{\pi\nu}^+} dS_{\pi\nu} \int_{S_{\pi\nu\gamma}^-}^{S_{\pi\nu\gamma}^+} \frac{dS_{\pi\nu\gamma}}{\sqrt{-\Delta_4}} |\mathcal{M}|_{\tau \rightarrow \pi\pi\nu\tau\gamma}^2$$

- virtual correction

$$G_{EM}^{(v)}(t) = \frac{12 \int_{u_{\min}}^{u_{\max}} D(t,u) f_{\text{loop}}^{\text{elm}}(u, M_\gamma) du}{m_\tau^6 \left(1 - \frac{4m_\pi^2}{t}\right)^{3/2} \left(1 - \frac{t}{m_\tau^2}\right)^2 \left(1 + \frac{2t}{m_\tau^2}\right)}$$



Updated evaluation of a_μ from τ data

The **relative shifts** of a_μ caused by the IB corrections in the $\pi\pi$ channel

$$\Delta a_\mu^{\text{HVP,LO}} = \frac{1}{4\pi^3} \int_{4m_\pi^2}^{t_{max}} dt K(t) \left[\frac{K_\sigma(t)}{K_\Gamma(t)} \frac{d\Gamma_{\tau\pi\pi[\gamma]}}{dt} \right] \times \left(\frac{R_{\text{IB}}(t)}{S_{\text{EW}}} - 1 \right).$$

$\Delta a_\mu^{\text{HVP,LO}}[\pi\pi]$ from $G_{\text{EM}}(t)$ in units of 10^{-11}						
$[t_{min}, t_{max}]$	Sol-A	Sol-B	CEN	MR[$\mathcal{O}(p^4)$]	MR[$\mathcal{O}(p^6)$]	
$[4m_\pi^2, 1 \text{ GeV}^2]$	-7.1	-44.9	-10.6	-10.4	-15.9	-63.2 ± 16.5
$[4m_\pi^2, 2 \text{ GeV}^2]$	-6.4	-44.5	-9.8	-9.6	-15.2	-58.1 ± 12.2
$[4m_\pi^2, 3 \text{ GeV}^2]$	-6.3	-44.4	-9.7	-9.5	-15.1	-67.8 ± 17.5
$[4m_\pi^2, m_\tau^2]$	-6.3	-44.4	-9.7	-9.5	-15.1	-64.9 ± 13.4

- The dominant contributions arise from the **low-energy** regions.

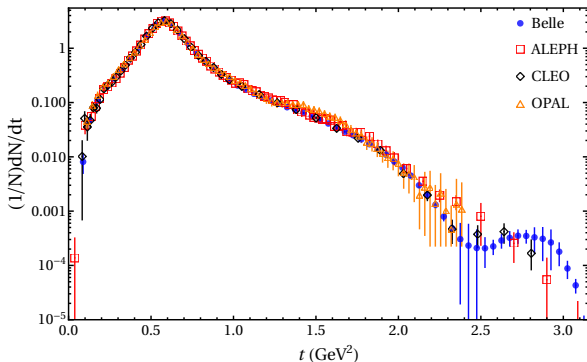
$\Delta a_\mu^{\text{HVP,LO}}[\pi\pi]$ from various IB sources in units of 10^{-11}					
Parameters	t_{max}/GeV^2	S_{EW}	$\beta_{\pi^+\pi^-}^3/\beta_{\pi^-\pi^0}^3$	$ F_{\pi\pi}^{(0)}/F_{\pi\pi}^{(-)} ^2$	FSR
Sol-A	1.0	-112.7	-74.1	76.8	42.7
	2.0	-114.6	-74.3	76.7	43.3
	3.0	-114.7	-74.3	76.7	43.3
	m_τ^2	-114.7	-74.3	76.7	43.3
Sol-B	1.0	-113.6	-76.9	77.1	43.2
	2.0	-115.4	-77.1	77.0	43.8
	3.0	-115.5	-77.1	77.0	43.8
	m_τ^2	-115.5	-77.1	77.0	43.8

Updated evaluation of a_μ from τ data

- Evaluation of $e^+e^- \rightarrow \pi^+\pi^-$ cross section by using tau data with IB corrections

$$\sigma_{e^+e^- \rightarrow \pi^+\pi^-}^0(t) = \frac{2\pi\alpha^2 m_\tau^2}{3|V_{ud}|^2 t} \frac{1}{\left(1 - \frac{t}{m_\tau^2}\right)^2 \left(1 + \frac{2t}{m_\tau^2}\right)} \frac{\mathcal{B}_{\tau\pi\pi}}{\mathcal{B}_{\tau e}} \frac{1}{N_{\tau\pi\pi}} \frac{dN_{\tau\pi\pi}}{dt} \frac{R_{\text{IB}}(t)}{S_{\text{EW}}}$$

- Input two-pion tau decay data



Updated evaluation of a_μ from τ data

- LO HVP

$$a_\mu^{\text{HVP,LO}|\pi\pi,\tau\text{data}} = \frac{1}{4\pi^3} \int_{4m_\pi^2}^{\infty} dt K(t) \sigma_{e^+e^- \rightarrow \pi^+\pi^-}^0(t)$$

Parameters	Experiments	$a_\mu^{\text{HVP,LO} \pi\pi,\tau\text{data}}$
Sol-A	Belle	$516.7 \pm 2.1 \pm 7.9 \pm 2.2$
	ALEPH	$513.3 \pm 4.3 \pm 2.8 \pm 2.1$
	CLEO	$516.9 \pm 3.2 \pm 8.8 \pm 2.2$
	OPAL	$527.2 \pm 9.8 \pm 6.8 \pm 2.1$
Sol-B	Belle	$513.4 \pm 2.0 \pm 7.9 \pm 2.2$
	ALEPH	$510.2 \pm 4.2 \pm 2.8 \pm 2.1$
	CLEO	$513.7 \pm 3.2 \pm 8.8 \pm 2.2$
	OPAL	$523.5 \pm 9.5 \pm 6.8 \pm 2.1$

- The first error is from the uncertainties of the experimental invariant-mass spectra.
- The second error arises from the branching ratios of $\mathcal{B}_{\tau\pi\pi}$ and $\mathcal{B}_{\tau e}$.
- The third error comes from the IB correction function R_{IB} .

- We take **Sol-A** as the baseline solution

$$a_\mu^{\text{HVP,LO}|\pi\pi,\tau\text{data, Belle}} = 516.7 \pm 2.1_{\text{Spec}} \pm 7.9_{\text{BR}} \pm 2.2_{\text{IB}} \pm 3.3_{\text{Sys}}$$

$$a_\mu^{\text{HVP,LO}|\pi\pi,\tau\text{data, ALEPH}} = 513.3 \pm 4.3_{\text{Spec}} \pm 2.8_{\text{BR}} \pm 2.1_{\text{IB}} \pm 3.2_{\text{Sys}}$$

$$a_\mu^{\text{HVP,LO}|\pi\pi,\tau\text{data, CLEO}} = 516.9 \pm 3.2_{\text{Spec}} \pm 8.9_{\text{BR}} \pm 2.2_{\text{IB}} \pm 3.3_{\text{Sys}}$$

$$a_\mu^{\text{HVP,LO}|\pi\pi,\tau\text{data, OPAL}} = 527.2 \pm 9.8_{\text{Spec}} \pm 6.8_{\text{BR}} \pm 2.1_{\text{IB}} \pm 3.7_{\text{Sys}}$$

- The average of the central values weighted

$$10^{10} \cdot a_\mu^{\text{HVP,LO}|\pi\pi,\tau\text{data}} = 516.0 \pm 2.9_{\text{Spec+BR}} \pm 4.0_{\text{IB+Sys}}$$

- **Full LO HVP contribution** (including $\pi^+\pi^-\pi^0$, $2\pi^+2\pi^-$, ...)

$$10^{10} \cdot a_\mu^{\text{HVP,LO}|\tau\text{data}} = 702.1 \pm 2.9_{\text{Spec+BR}} \pm 4.0_{\text{IB+Sys}} \pm 2.1_{\text{Others}}$$

- The deviation between **SM** and **BNL+FNAL** experimental result

$$\Delta a_\mu \equiv a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = (14.9 \pm 5.6) \times 10^{-10} \Rightarrow 2.7\sigma$$

Summary

- The real-photon electromagnetic correction to the $\tau \rightarrow \pi\pi\nu_\tau$ process has been calculated.
- By taking both vector and scalar contributions into account in $\omega \rightarrow \pi^0\pi^0\gamma$, we updated the determination of the key coupling d_4 .
- Using the experimental $\tau \rightarrow \pi\pi\nu_\tau$ data from Belle, ALEPH, CLEO and OPAL to calculate $a_\mu^{\text{HVP,LO}}|_{\pi\pi,\tau\text{data}}$.
- The resulting deviation of a_μ between our estimation and the updated world average lies at the level of 2.7σ .

Thanks!