

Decoherence in high energy collisions as renormalization group flow

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Introduction

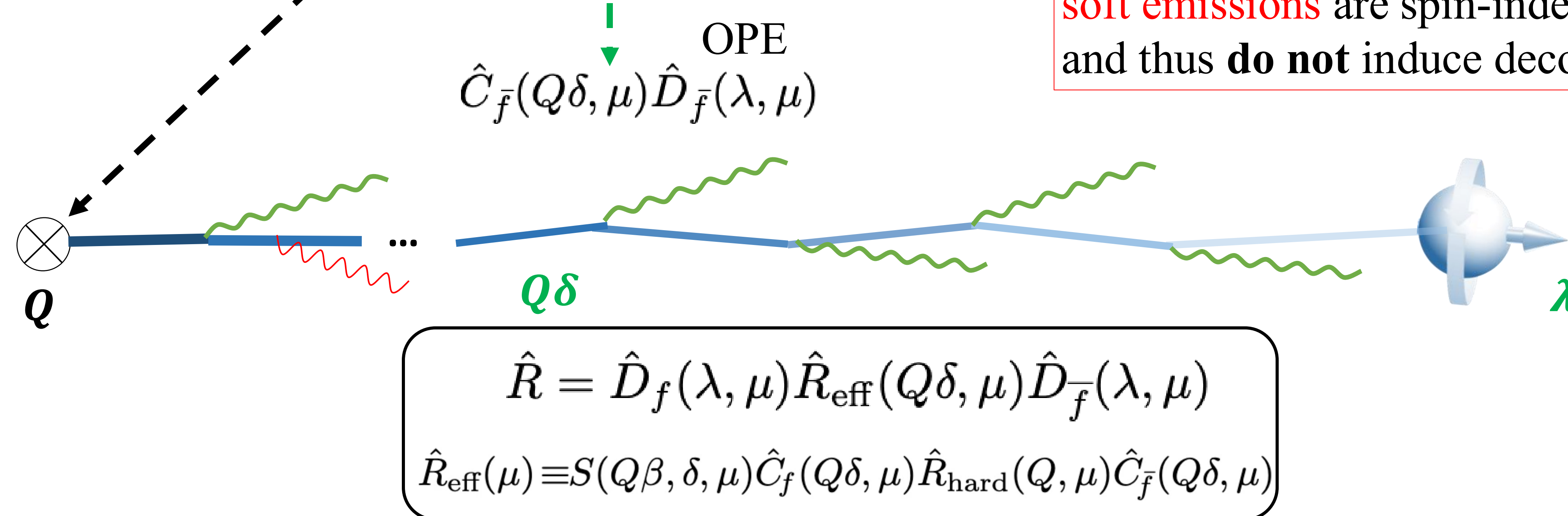
- Observations such as top-quark **spin entanglement** demonstrate quantum effects at high energies^[1].
- However, the quark pairs are **open quantum systems**, as unresolved soft & collinear photon/gluon radiation inevitably causes **decoherence**.

SCET Factorization

- Using SCET factorization theorems, the production density matrix of the fermion pair is factorized:

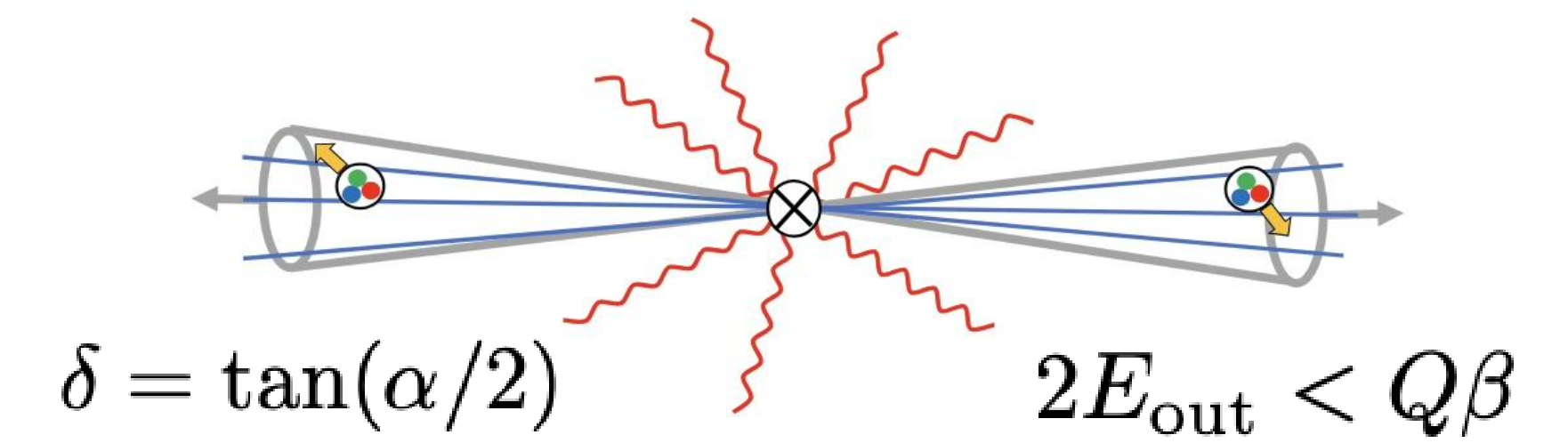
$$\hat{R} = \hat{J}_f(Q\delta, \lambda, \mu) \hat{R}_{\text{hard}}(Q, \mu) \hat{J}_{\bar{f}}(Q\delta, \lambda, \mu) S(Q\beta, \delta, \mu)$$

At leading order, **soft emissions** are spin-independent and thus **do not** induce decoherence^[4-5].



Effective Field Theory for Decoherence

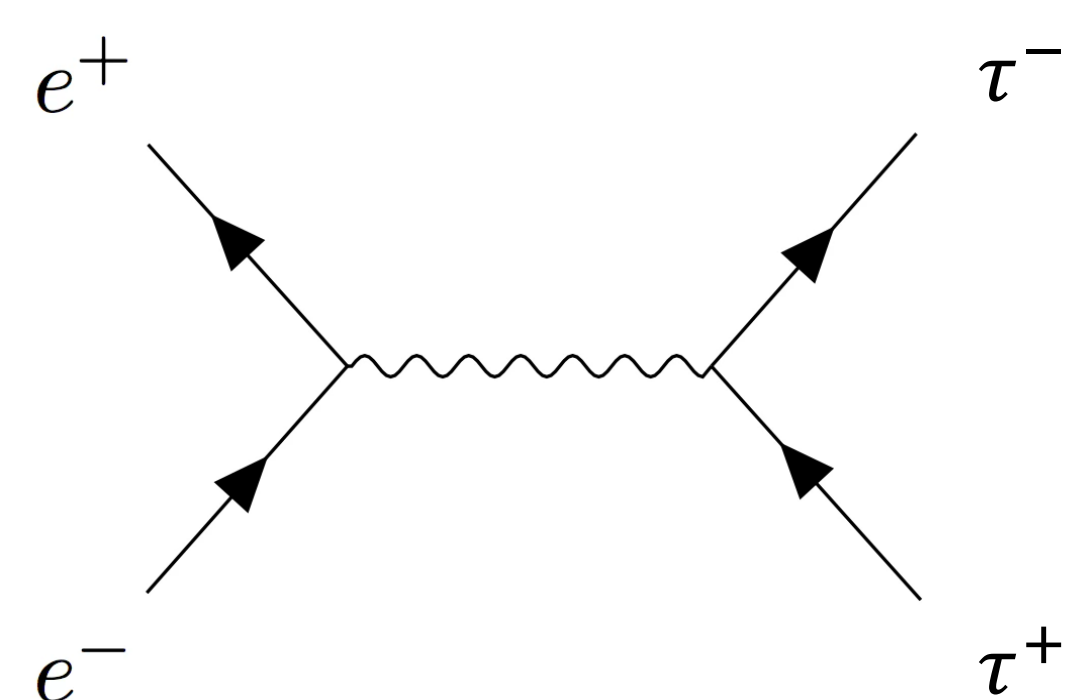
- Radiation should be considered **unresolved** if either **soft** or **collinear**.
- We introduce the energy and angular resolution parameters, similar to **Sterman-Weinberg cone jet** definition^[2].
- We apply **Soft Collinear Effective Theory (SCET)**^[3].



At scale Q: Production

- The fermion pair is produced at an energy scale Q much higher than an IR scale λ .
- The spin state of two spin-1/2 particles can be characterized by a two-qubit density operator $\hat{\rho}_{\text{hard}} = \frac{\hat{R}_{\text{hard}}}{\text{Tr}[\hat{R}_{\text{hard}}]}$.

- eg. $e^+ + e^- \rightarrow \tau^- + \tau^+$



$$\hat{\rho}_{\text{hard}}(Q, \mu) = \frac{1}{4} \left(\hat{I}_2 \otimes \hat{I}_2 + \frac{\sin^2 \theta}{1 + \cos^2 \theta} \hat{\sigma}_1 \otimes \hat{\sigma}_1 + \frac{\sin^2 \theta}{1 + \cos^2 \theta} \hat{\sigma}_2 \otimes \hat{\sigma}_2 - \hat{\sigma}_3 \otimes \hat{\sigma}_3 \right)$$

From $Q\delta$ to λ : Collinear Emission

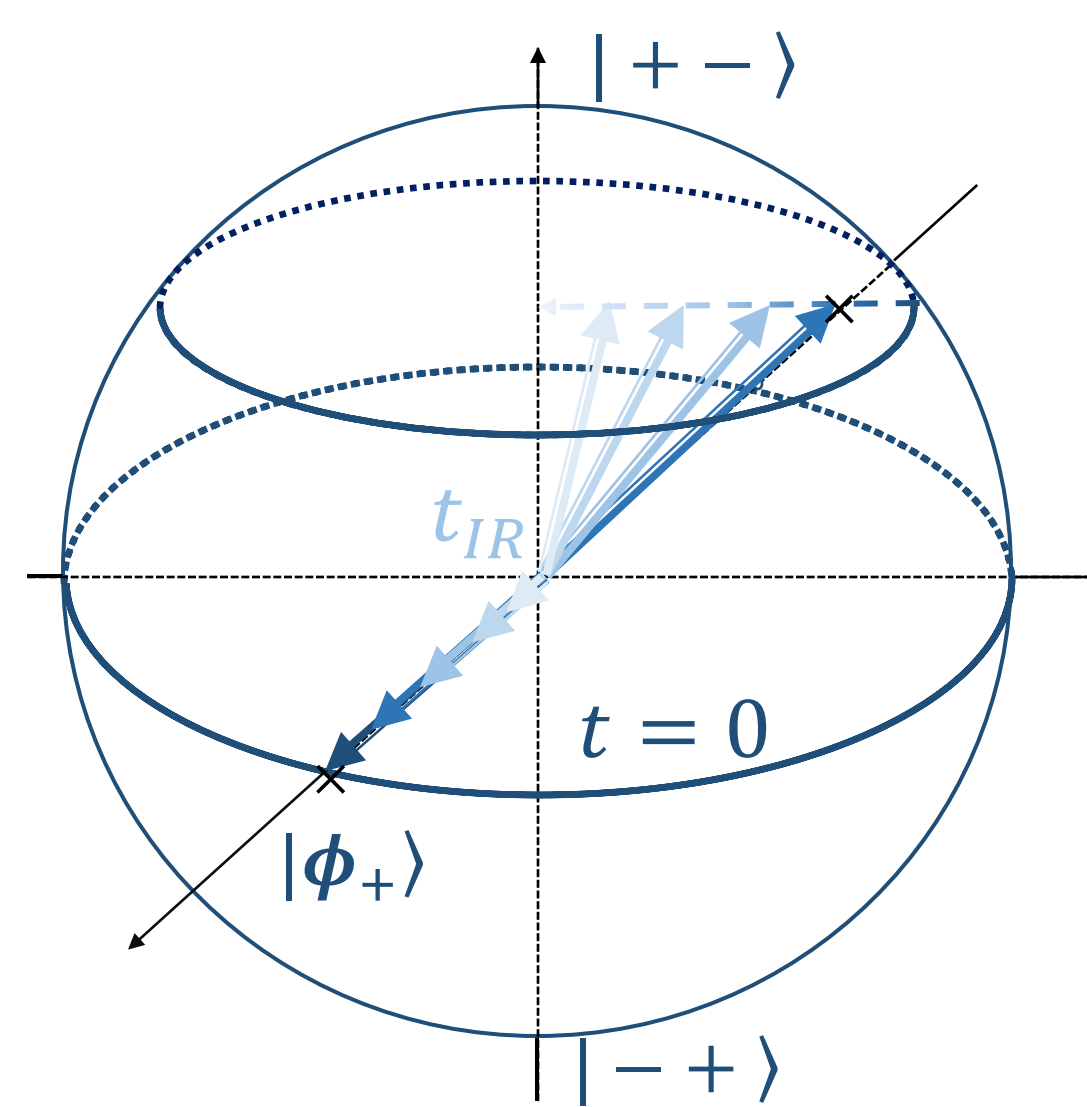
- From scale $Q\delta$ to λ , the fermion pair emits collinear photons.
- Solution to the Renormalization Group Equation(RGE)

$$\hat{R}_{\text{eff}}(t) = \hat{U}_f(t, 0) \hat{R}_{\text{eff}}(0) \hat{U}_{\bar{f}}(t, 0) \quad t \equiv \log(Q\delta/\mu)$$
- The evolution functions: $U^{\mathcal{P}}(t, 0) = \exp\left(\int_0^t dt \gamma^{\mathcal{P}}\right)$, $\gamma^{\mathcal{P}}$ as anomalous dimension.

- The **RGE** encapsulates decoherence from collinear radiation, whose **Kraus representation** can be determined.
- **Anomalous dimensions** determine the information loss.

- eg. collinear photon emission \leftrightarrow **phase-flip** channel

$$U^U = U^L = 1, \quad U^T = \exp\left(-\frac{\alpha}{2\pi} t\right) \leftrightarrow \hat{K}_{(i,j)} = \hat{K}_i^{\ell-} \otimes \hat{K}_j^{\ell+}$$



$$\hat{K}_0^{\ell-} = \hat{K}_0^{\ell+} = \sqrt{1-p^2} \mathbb{I},$$

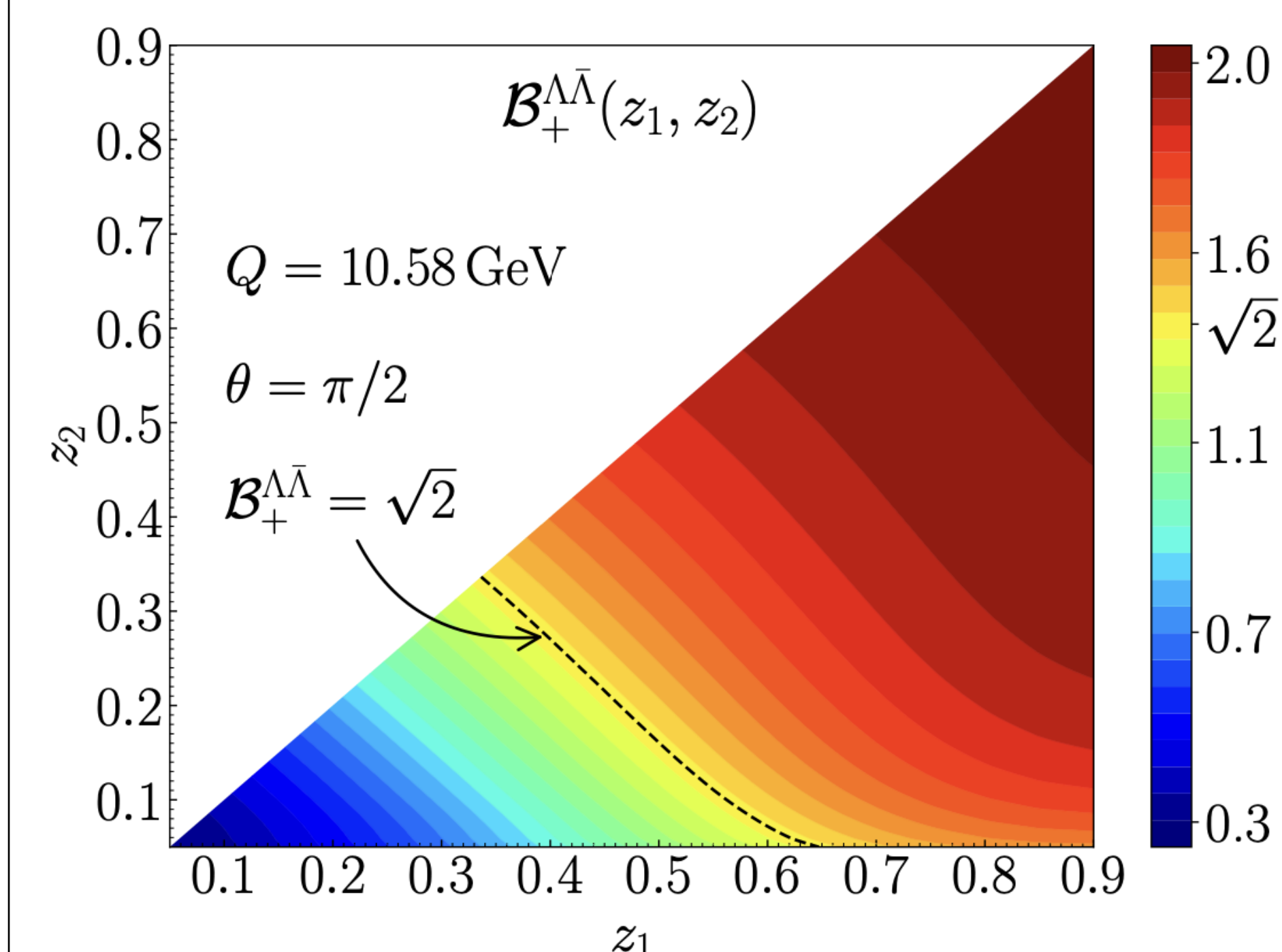
$$\hat{K}_1^{\ell-} = \hat{K}_1^{\ell+} = p \hat{\sigma}_3, \quad p = \sqrt{\frac{1}{2} \left[1 - \exp\left(-\frac{\alpha}{2\pi} t\right) \right]}$$

- Final-state concurrence for general cases

$$\mathcal{C}_{\text{final}} \leq \mathcal{C}(0) \left(\frac{Q\delta}{\lambda} \right)^{-\frac{\alpha}{\pi}}$$

At scale λ : Measurement

- When the state is evolved to scale λ , the final stage of the process is **the projection of the evolved spin state onto definite experimental outcome**.
- eg. Defining spin measurement operators as $\hat{M}_f(\mathcal{S}_f) \equiv \hat{D}_f \hat{P}_f$ to contain final state projection and the spin measurement.



(arXiv:2507.15387, decoherence in QCD)

Summary

- Developed a systematic framework for computing spin decoherence from final-state radiations in high-energy processes by unifying SCET with open quantum system formalism.
- Demonstrated that renormalization group evolution constitutes a quantum channel, where the RG flow parameter drives a Markovian loss of quantum coherence.
- Framework extends naturally to QCD, including subleading- power soft effects, etc.

Reference

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