

Anomalous transport in anomaly free currents from symmetry breaking



Karl Landsteiner

Based on:

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In collaboration with:

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[International Conference on Symmetry Breaking Phenomena in Quantum Field Theory](#)

15-18 May 2026, Hefei, China

Outline

- Anomaly induced transport
- Holography
- Holographic model with anomalies and symmetry breaking
- Discussion and Outlook

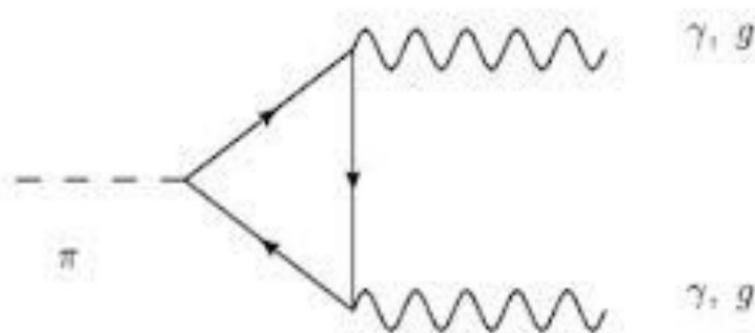
Anomaly induced transport

- Chiral anomalies

Conflict between two core concepts of theoretical physics:

Symmetry \longleftrightarrow **Quantum mechanics**

- Pion decay into photons or gravitons



$$D_\mu J_5^\mu = \frac{1}{8\pi^2} F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{1}{192\pi^2} R_{\alpha\beta\mu\nu} \tilde{R}^{\alpha\beta\mu\nu}$$

[Adler],[Bell,Jackiw] 1969
[Kimura] 1969

Anomaly induced transport

- Chiral anomalies

$$\nabla_{\mu} J_a^{\mu} = \frac{d_{abc}}{16\pi^2} F_{\mu\nu}^a \tilde{F}_a^{\mu\nu} + \frac{b_a}{348\pi^2} R_{\mu\nu}^{\alpha\beta} \tilde{R}_{\alpha\beta}^{\mu\nu}$$

- Anomaly coefficients:

$$d_{abc} = \text{str}(T_a T_b T_c)_r - \text{str}(T_a T_b T_c)_\ell$$

$$b_a = \text{str}(T_a)_r - \text{str}(T_a)_\ell$$

Anomaly induced transport

- Chiral Magnetic and Vortical effects

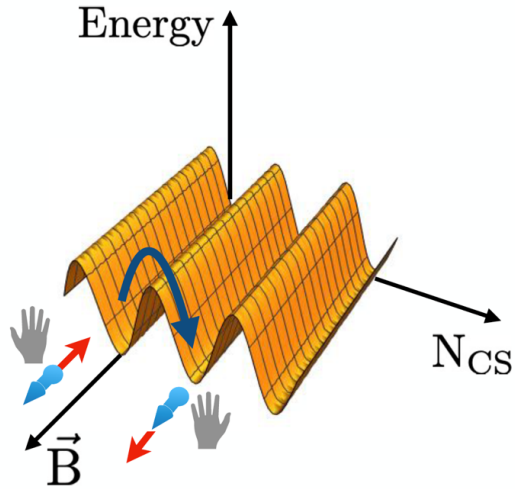
$$\vec{J}_A = d_{ABC} \frac{\mu_B}{4\pi^2} \vec{B}_C + \left(d_{ABC} \frac{\mu_B \mu_C}{4\pi^2} + b_A \frac{T^2}{12} \right) \vec{\Omega}$$

$$\vec{J}_\epsilon = \left(d_{ABC} \frac{\mu_B \mu_C}{8\pi^2} + b_A \frac{T^2}{24} \right) \vec{B}_A + \left(d_{ABC} \frac{\mu_A \mu_B \mu_C}{6\pi^2} + b_A \frac{\mu_A T^2}{6} \right) \vec{\Omega}$$

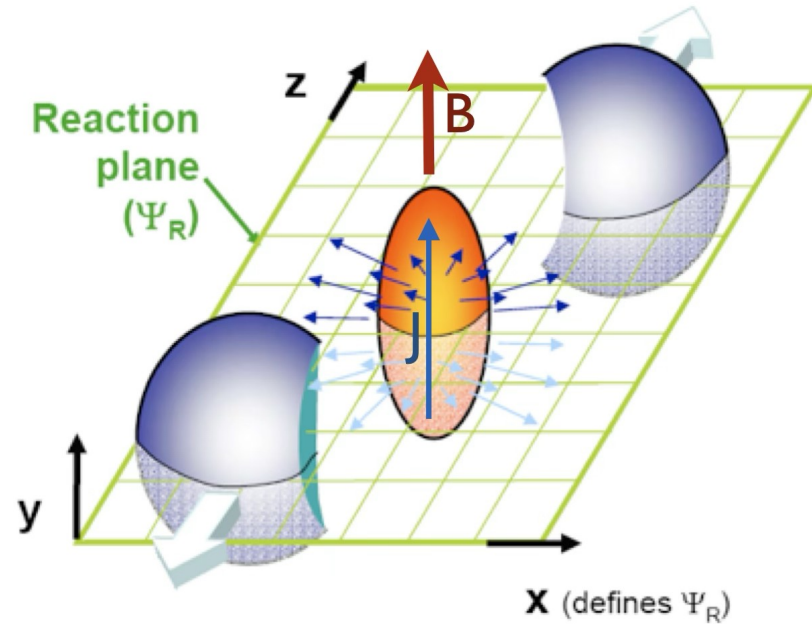
[Vilenkin]	1980
[Aleksseev, Fröhlich, Chaianov]	.
[Giovannini, Shaposhnikov]	.
[Fukushima, Kharzeev, Warringa]	.
[Newman]	.
[Son, Surowka]	.
[Newman, Oz]	.
[Erdmenger, Haack, Kaminski, Yarom]	.
[Banerjee, Bhattacharya, Bhattacharyya, Dutta, Loganayagam, Surowka]	.
[Megias, Landsteiner, Pena-Benitez]	.
[Megias, Melgar, Landsteiner, Pena-Benitez]	2012

Anomaly induced transport

- Realizations: Heavy Ion Collisions ?



$$\partial_\mu J_5^\mu = \frac{g^2 N_f}{16\pi^2} (G_{\mu\nu} \tilde{G}^{\mu\nu})$$



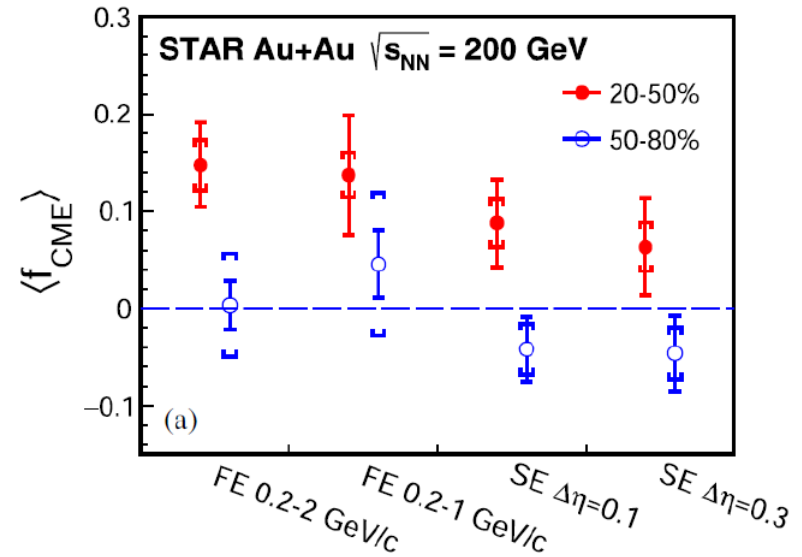
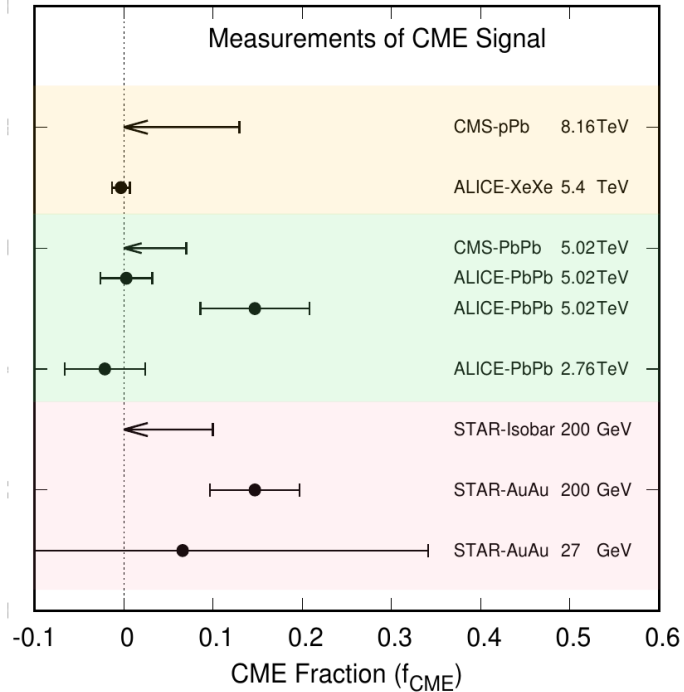
$$\vec{J} = \frac{\mu_5}{2\pi^2} \vec{B}$$

- Analogy to electroweak baryogenesis!
- Zhakharov conditions!

McLerran, Kharazeev, Warringa] 2008
[Fukushima, Kharazeev, Warringa] 2008

Anomaly induced transport

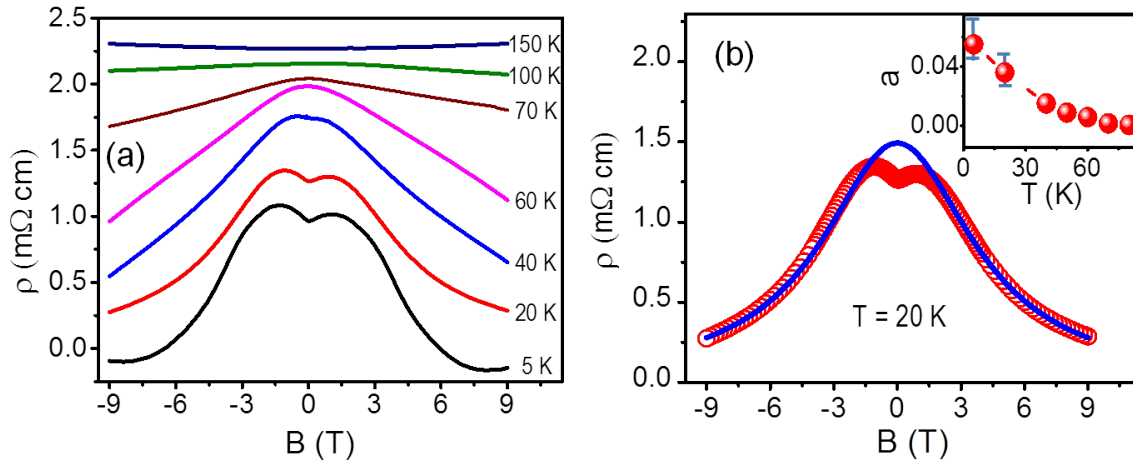
- Realizations: Heavy Ion Collisions ?



Anomaly induced transport

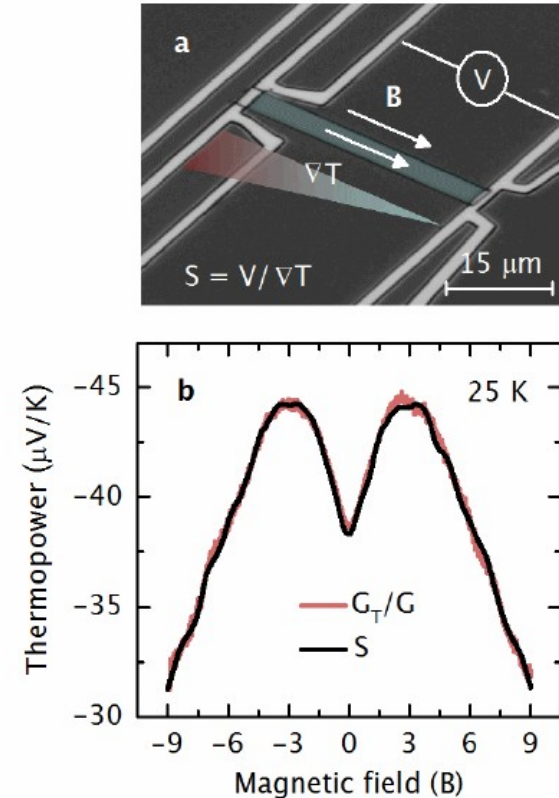
- **Realizations: Weyl semi-metals !**

CME in ZrTe₅



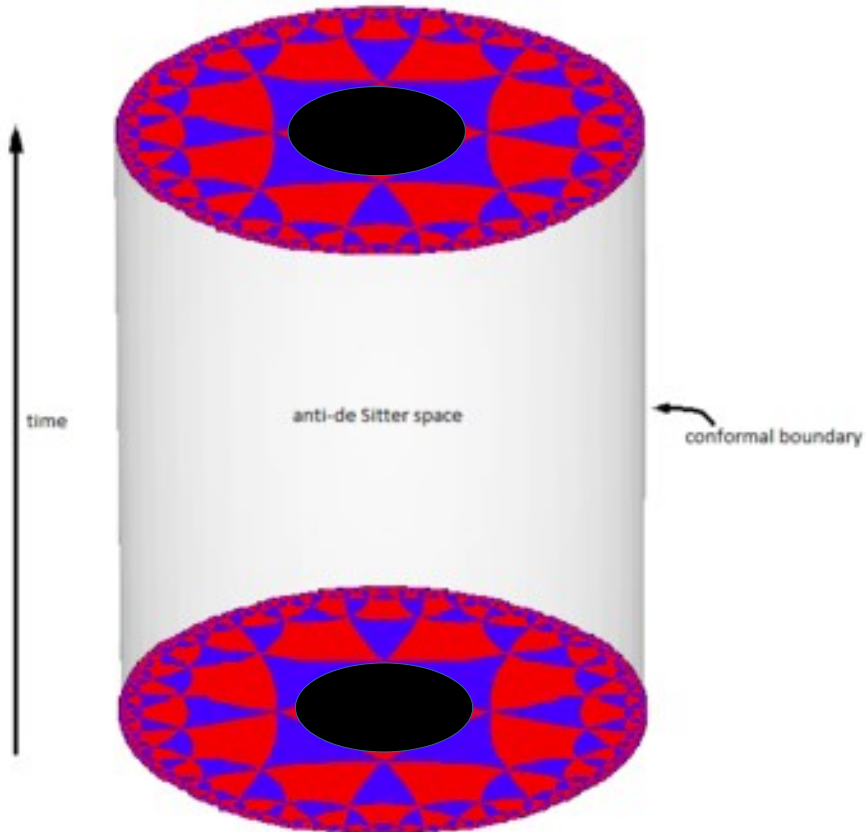
[Li, Kharzeev, Zhang, Huang, Pletikoscic, Fedorov, Zhong, Schneeloch, Gu, Valla] 2015

CthME in NbP



[Gooth, Niemann, Meng, Grushin, Landsteiner, Gotsmann, Menges, Schmidt, Shekhar, Suß, Huhne, Rellinghaus, Felser, Yan, Nielsch] 2017

(The value of) Holography



Gravity in asymptotically AdS = QFT

Holographic Dictionary

Metric

Energy Momentum
Tensor

Gauge field

Conserved current =
symmetry

Scalar field

Scalar operator

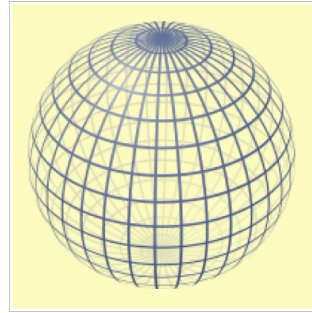
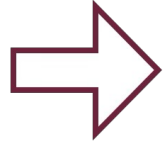
Boundary value

Coupling

Black Hole

Temperature

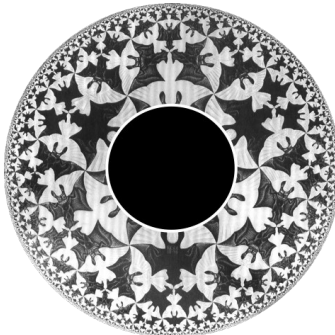
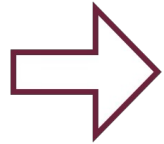
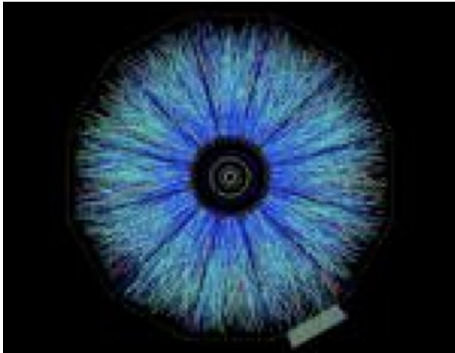
(The value of) Holography



"AdS is the hyperbolic cow of sQGP"

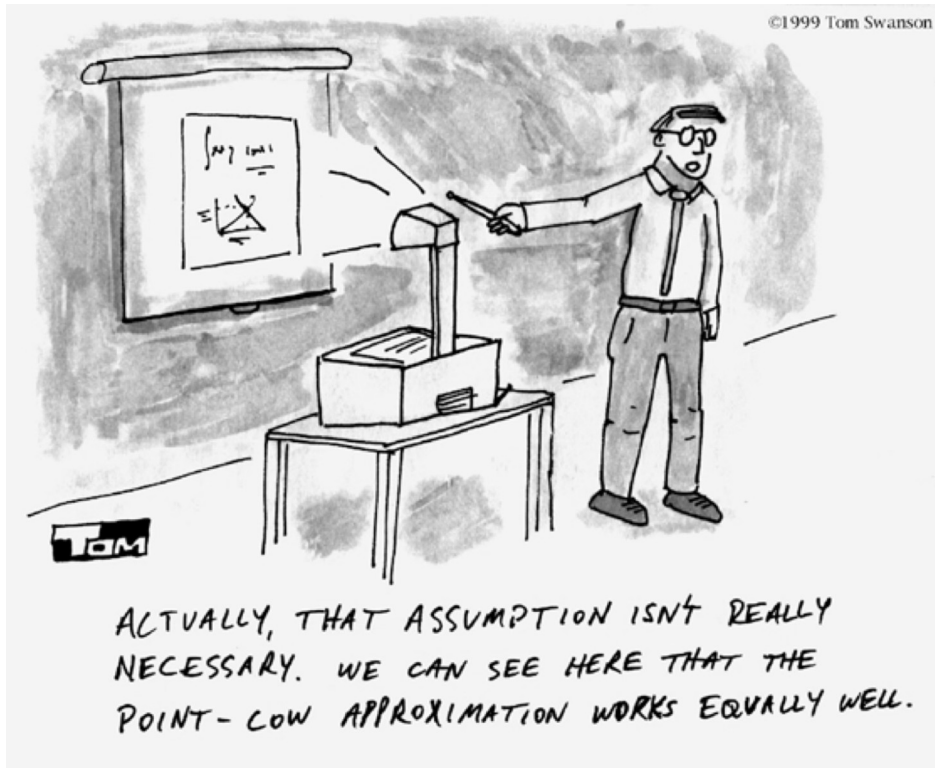
$$\frac{\eta}{s} = \frac{1}{4\pi}$$

Son, Starinets, Policastro, Kovtun ...



(The value of) Holography

Anomalies are better than that:



- CVE thought to be impossible

$$\vec{J} = \frac{\mu\mu_5}{2\pi^2} \vec{\Omega}$$

- Naively gravitational anomaly 4-th order in derivatives

$$\vec{J}_5 = \frac{T^2}{12\pi} \vec{\Omega}$$

- Strict equilibrium = smooth Euclidean section vanishing CME

$$A_5^0 = \mu_5$$

$$\vec{J} = \frac{\mu_5 - A_5^0}{2\pi^2} \vec{B}$$

Anomalous transport and symmetry breaking

Holography is a discovery tool for (anomalous) transport phenomena!

Goal: what happens if an anomalous and a non-anomalous symmetry are broken by a **common** symmetry breaking parameter?

Model:

$$S_{\text{kin}} = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \left(R + 2\Lambda - \frac{1}{4}F_V^2 - \frac{1}{4}F_A^2 - \frac{1}{4}F_W^2 \right)$$
$$S_{\text{CS}} = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \epsilon^{mnpqr} A_m \left(\frac{\kappa}{3} F_{A,np} F_{A,qr} + \kappa F_{V,np} F_{V,qr} + R_{np}^{ab} R_{abqr} \right)$$
$$S_{\partial} = \frac{1}{8\pi G_5} \int \sqrt{h} \left(K - 4n_m \epsilon^{mnpqr} A_n K_{pl} D_q K_r^l \right)$$

Matching to axial anomalies of Dirac fermions:

$$\frac{\kappa}{16\pi G_5} = -\frac{N_f}{16\pi^2} \quad \frac{\lambda}{16\pi G_5} = -\frac{N_f}{384\pi^2}$$

Anomalous transport and symmetry breaking

- Symmetries:
- J_V ... Vector - consistent current is anomaly free
covariant current has anomaly
 - J_A ... Axial – axial anomaly
 - J_W ... “Anomaly free” current
 - T Energy-momentum tensor (similar to J_V)

Symmetry breaking: $S_\phi = \frac{1}{16\pi G_5} \int \sqrt{-g} (-|D\phi|^2 - m^2\phi^2)$ $D_n\phi = (\partial_n - i(A_n - W_n))\phi$

$$m^2 = -3 \quad \phi(r) \approx \frac{M}{r} + \dots + \frac{\mathcal{O}}{r^3} + \dots$$

QFT intuition: $S_M = \int d^4x M \bar{\Psi} \Psi$

Anomalous transport and symmetry breaking

- Symmetries:
- J_V ... Vector - consistent current is anomaly free
covariant current has anomaly
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 - J_W ... Anomaly free current
 - T energy momentum tensor, as J_V

Symmetry breaking: $S_\phi = \frac{1}{16\pi G_5} \int \sqrt{-g} (-|D\phi|^2 - m^2\phi^2)$ $D_n\phi = (\partial_n - i(A_n - W_n))\phi$

$$m^2 = -3 \quad \phi(r) \approx \frac{M}{r} + \dots + \frac{\mathcal{O}}{r^3} + \dots$$

QFT interpretation: $S_M = \int d^4x M \bar{\Psi}\Psi$

Anomalous transport and symmetry breaking

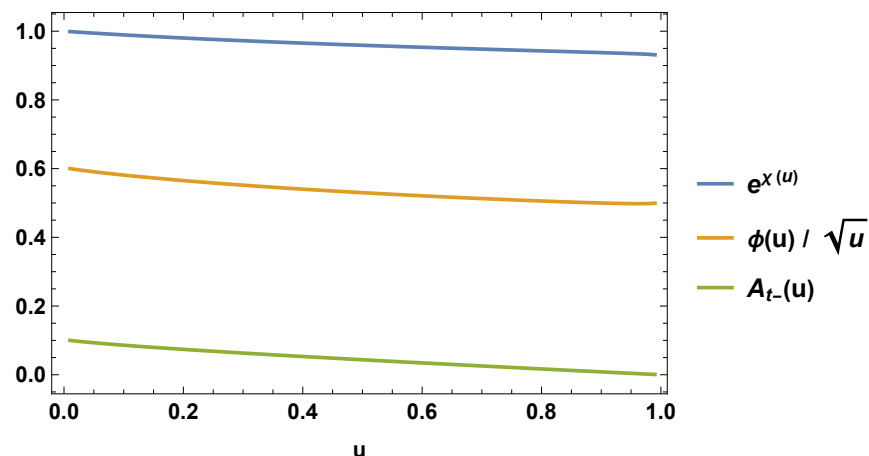
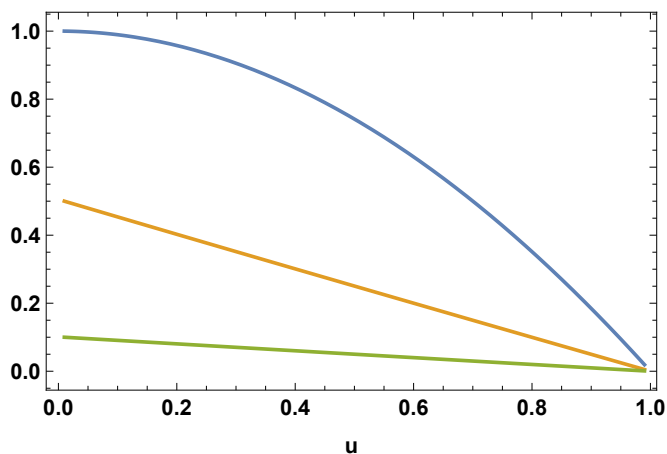
Ansatz: $ds^2 = \frac{1}{u^2} \left(-f(u)dt^2 + e^{\chi(u)} \frac{du^2}{f(u)} + d\vec{x}^2 \right) \quad \phi = \phi(u), \quad \lim_{u \rightarrow 0} \frac{\phi(u)}{\sqrt{u}} = M$

$$S = S_t(u)dt, \quad \lim_{u \rightarrow 0} S_t(u) = \mu_S, \quad S \in \{V, A, W\}$$

M=0: $f(u) = 1 - (1 + \sum_S Q_S^2)u^2 + \sum_S Q_S^2 u^3 \quad S_t(u) = \mu_S(1 - u)$

$$\pi T = 1 - \frac{\sum_S \mu_S^2}{6} \quad Q_S = \frac{\mu_S}{\sqrt{3}}$$

M≠0: numerical solutions $M = 0.7T, \quad \mu_V = 0.12T, \quad \mu_A = 0.35T, \quad \mu_W = 0.23T$



Anomalous transport and symmetry breaking

Response and Kubo formulas:

$$\begin{aligned}\vec{J}_v &= \sigma_{vv}^B \vec{B}_v + \sigma_{va}^B \vec{B}_a + \sigma_{vw}^B \vec{B}_w + \sigma_v^V \vec{\Omega}, \\ \vec{J}_a &= \sigma_{av}^B \vec{B}_v + \sigma_{aa}^B \vec{B}_a + \sigma_{aw}^B \vec{B}_w + \sigma_a^V \vec{\Omega}, \\ \vec{J}_w &= \sigma_{wv}^B \vec{B}_v + \sigma_{wa}^B \vec{B}_a + \sigma_{ww}^B \vec{B}_w + \sigma_w^V \vec{\Omega}, \\ \vec{J}_\varepsilon &= \sigma_{\varepsilon v}^B \vec{B}_v + \sigma_{\varepsilon a}^B \vec{B}_a + \sigma_{\varepsilon w}^B \vec{B}_w + \sigma_\varepsilon^V \vec{\Omega}.\end{aligned}$$

$$\sigma_{sb}^B = \lim_{p_k \rightarrow 0} \frac{i}{2p_k} \sum_{i,j} \epsilon_{ijk} \langle J_s^i J_b^j \rangle |_{\omega=0} \quad \sigma_s^V = \lim_{p_k \rightarrow 0} \frac{i}{2p_k} \sum_{i,j} \epsilon_{ijk} \langle J_s^i J_\varepsilon^j \rangle |_{\omega=0}, \quad (s \in \{v, a, w, \varepsilon\}, b \in \{v, a, w\})$$

M=0:

$$\begin{aligned}\sigma_{sv}^B(0) &= \frac{1}{2\pi^2} \begin{cases} \mu_a, & (s=v) \\ \mu_v, & (s=a) \\ \mu_v \mu_a, & (s=\varepsilon) \end{cases}, \\ \sigma_{sa}^B(0) &= \frac{1}{2\pi^2} \begin{cases} \mu_v, & (s=v) \\ \mu_a, & (s=a) \\ \frac{\mu_v^2 + \mu_a^2}{2} + \frac{\pi^2}{6} T^2, & (s=\varepsilon) \end{cases}, \\ \sigma_s^V(0) &= \frac{1}{2\pi^2} \begin{cases} \mu_v \mu_a, & (s=v) \\ \frac{\mu_v^2 + \mu_a^2}{2} + \frac{\pi^2}{6} T^2, & (s=a) \\ \frac{\mu_a}{3} (3\mu_v^2 + \mu_a^2 + \pi^2 T^2), & (s=\varepsilon) \end{cases}.\end{aligned}$$

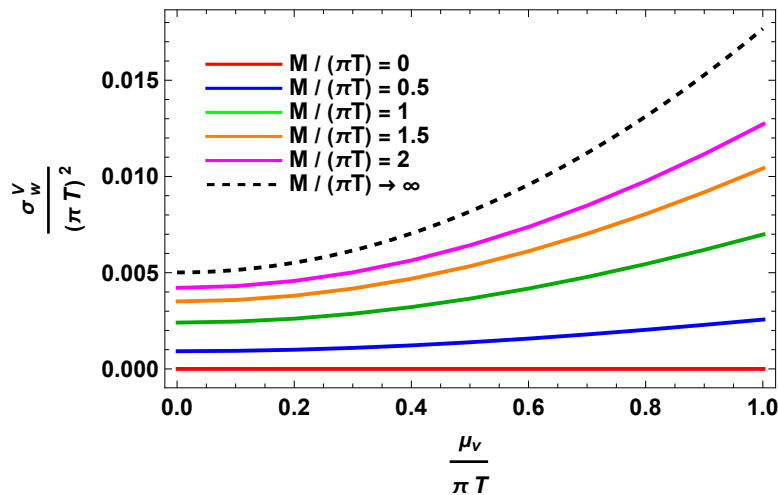
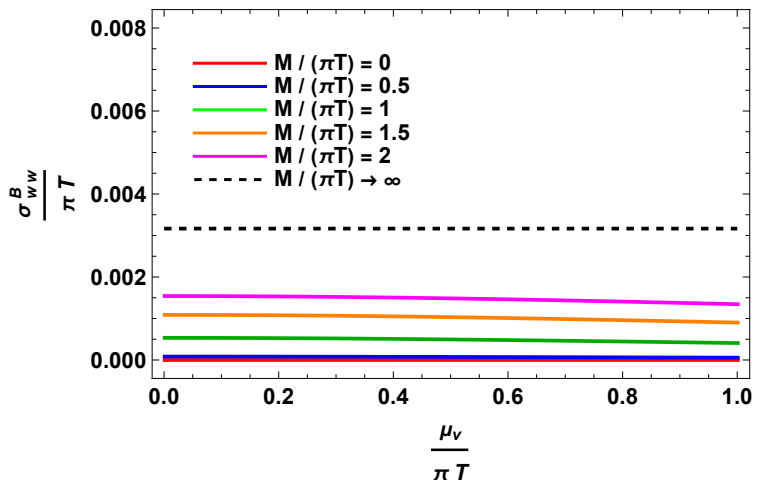
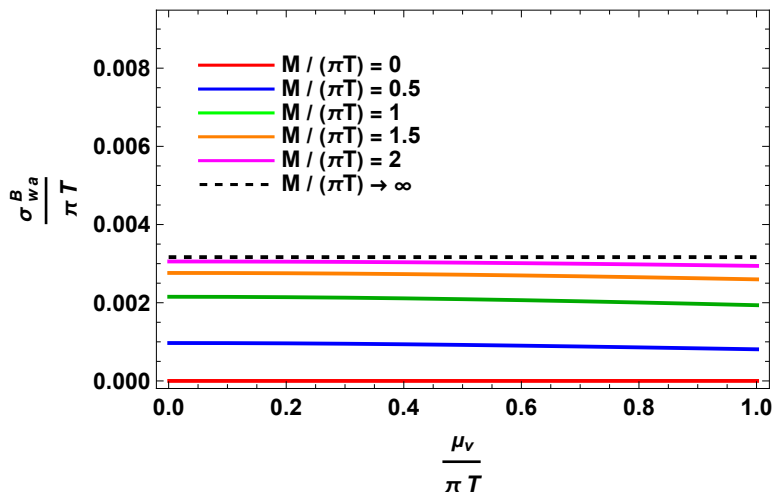
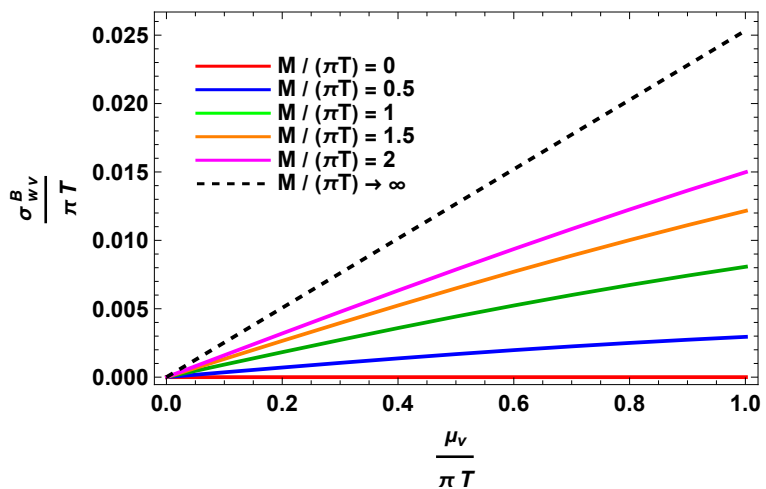
All conductivities related to non-anomalous symmetry vanish!

$$\sigma_{wx} = 0$$

Anomalous transport and symmetry breaking

$M \neq 0$:

$$\vec{J}_w = \sigma_{wv}^B \vec{B}_v + \sigma_{wa}^B \vec{B}_a + \sigma_{ww}^B \vec{B}_w + \sigma_w^V \vec{\Omega}$$



Anomalous transport and symmetry breaking

Chiral magnetic conductivities vs symmetry breaking

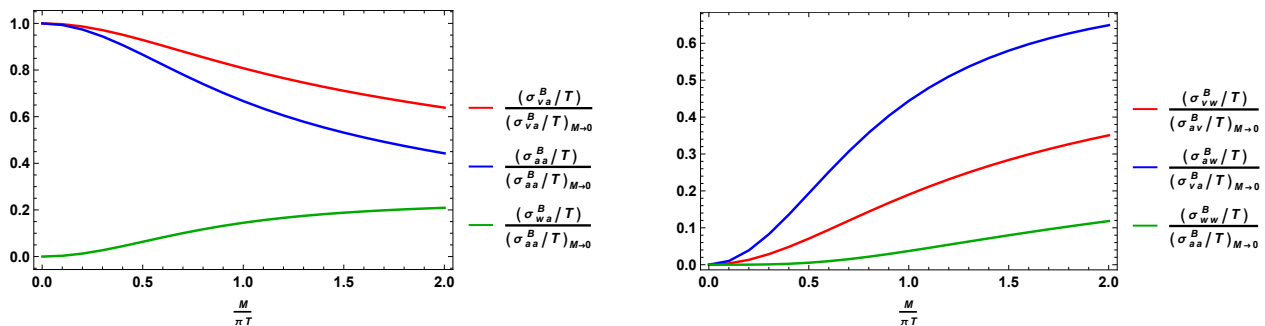


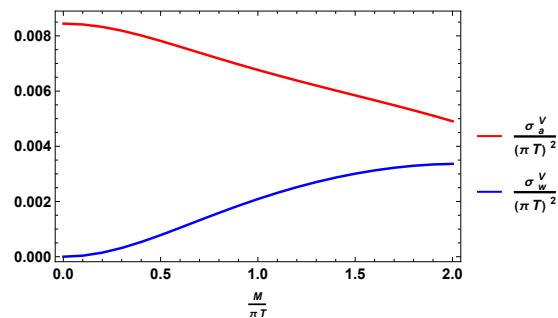
Figure 7. Plots of the chiral magnetic conductivities σ_{sa}^B (left panel) and σ_{sw}^B (right panel), where

Vortical conductivities: temperature part

$$\mu_v = \mu_a = \mu_w = 0$$

$$\vec{J}_a = \sigma_a^V \vec{\Omega}$$

$$\vec{J}_w = \sigma_w^V \vec{\Omega}$$



Anomalous transport and symmetry breaking

The covariant derivative $D_\mu \phi = \partial_\mu \phi - i(A_\mu - W_\mu)\phi = \partial_\mu \phi - iA_\mu^- \phi$ $A_\mu^\pm = A_\mu \pm W_\mu$

This implies for the conductivities $\sigma_{AS} = \frac{1}{2}(\sigma_{+S} + \sigma_{-S})$ $\sigma_{WS} = \frac{1}{2}(\sigma_{+S} - \sigma_{-S})$

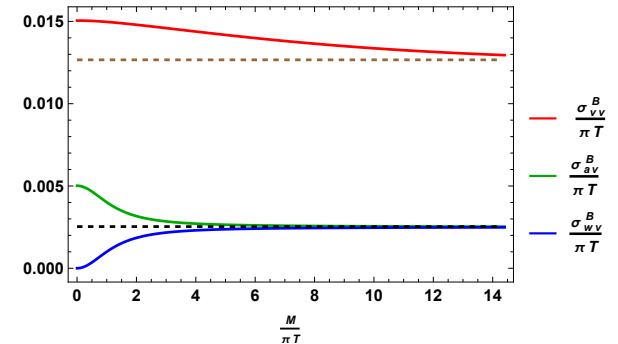
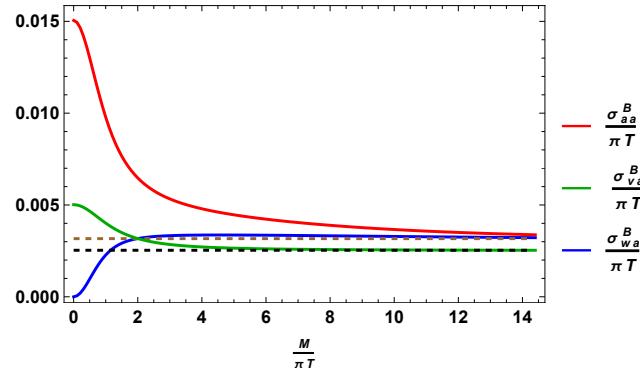
The $M \rightarrow \infty$ limit the W-field symmetry is maximally broken: $\sigma_{-S} = 0$

$$\sigma_{WV} = \frac{\mu_V}{4\pi^2}$$

$$\sigma_{WA} = \frac{\mu_A + \mu_W}{16\pi^2}$$

$$\sigma_{WW} = \frac{\mu_A + \mu_W}{16\pi^2}$$

$$\sigma_W^V = \frac{\mu_V^2 + (\mu_A + \mu_W)^2/4}{8\pi^2} + \frac{\pi^2 T^2}{24}$$



Discussion and Outlook

- Anomalies give rise to Chiral “transport” effects
- Conductivities are determined by anomaly coefficients
- Surprisingly non-anomalous currents can participate if symmetries are broken
- Realization in terms of holographic model explicit symmetry breaking
- What happens for spontaneous symmetry breaking? [work in progress]
- Weak coupling QFT calculation
- Real world applications?
 - TaAs: Weyl semi-metal with 24 Weyl nodes,
 - Possible complicated symmetry breaking patterns

谢谢

