

Automated One-Loop Computation in Quarkonium Process Within NRQCD Framework

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Outline

- General steps for the one-loop computation in quarkonium physics within the NRQCD framework and basic tools used in each step
 - Generating Feynman diagrams and amplitudes FeynArts, QGraf ...
 - Dirac- and Color-algebra simplification FeynCalc, FORM, FormLinkFeynCalc ...
 - Partial fraction and IBP reduction to Master Integrals (MI)
\$ Apart, FIRE ...
 - Post-processing final results
- Specific process in the examples
$$e^+ + e^- \rightarrow \gamma^* \rightarrow J/\psi + \text{S-, P- or D-wave}$$
- Summary

Generating Feynman Diagrams and Amplitudes

- Generating Feynman diagrams
 - Replacing hadrons with on-shell partonic states

$$J/\psi \rightarrow c(p_1)\bar{c}(p_2)$$

$$J/\psi \rightarrow c(p_1)\bar{c}(p_2)g(k)$$

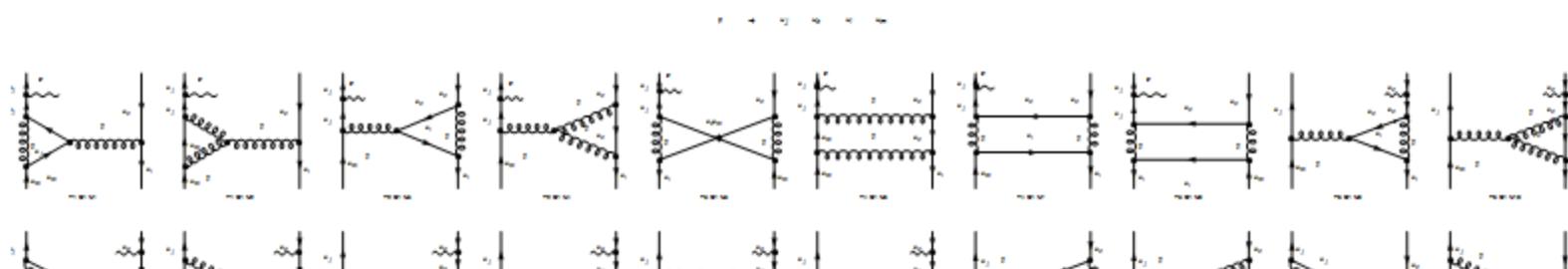
- FeynArts – Mathematica 9
<http://www.feynarts.com>

- QGraf – Fortran 77
<http://cfif.ist.utl.pt/~paulo/qgraf.html>

- Generating Amplitude for each diagram
 - Standard Model
 - Adjusting for later processing
 - Export to files

Generating Feynman Diagrams and Amplitudes

```
<< FeynArts`FeynArts`  
  
$ExcludeTopologies[V4onExt] = FreeQ[Cases[#, Propagator[External][__]], Vertex[4]] &;  
  
top = CreateTopologies[1, 1 → 4, ExcludeTopologies → {WF Corrections, Tadpoles, V4onExt}];  
  
tmp = InsertFields[top, {V[1]} → {F[3], -F[3], F[3], -F[3]}, Model → "SMQCD",  
  ExcludeParticles → {V[1], V[2], V[3], V[4], S[_], F[4]}, InsertionLevel → {Classes}];  
  
tmp1 =  
  DiagramSelect[tmp,  
    (FreeQ[List @@ #2, {___, PatternSequence[Propagator[Outgoing][_, v_, Field[2]], ___,  
      Propagator[Outgoing][_, v_, Field[3]]], ___}]) &];  
tmp2 =  
  DiagramSelect[tmp1,  
    (FreeQ[List @@ #2, {___, PatternSequence[Propagator[Outgoing][_, v_, Field[4]], ___,  
      Propagator[Outgoing][_, v_, Field[5]]], ___}]) &];  
all = tmp2;  
  
all = DiagramDelete[all, 3 ... 4, 13 ... 14, 25 ... 26, 33 ... 34, 42 ... 43];  
  
Paint[all, ColumnsXRows → {10, 7}];
```



$$e^+ + e^- \rightarrow \gamma^* \rightarrow J/\psi + \eta_c$$

Generating Feynman Diagrams and Amplitudes

```

FAToFC[d_] := Module[{temp},
  temp = CreateFeynAmp[d, PreFactor + 1] /. {FourMomentum[Incoming, 1] → p3 + p4, FourMomentum[Outgoing, 1] →  $\frac{p_3^3}{2} + q_3$ ,
    FourMomentum[Outgoing, 2] →  $\frac{p_3^3}{2} - q_3$ , FourMomentum[Outgoing, 3] →  $\frac{p_4^4}{2} + q_4$ , FourMomentum[Outgoing, 4] →  $\frac{p_4^4}{2} - q_4$ ,
    MQU[Index[g_, i_]] → mc, MQD[Index[g_, i_]] → mc, EL → e, GS → Gstrong};
  temp = (Part[#, 3] & ) /@ (List @@ ToFA1Conventions[temp]);
  temp = temp /. {DiracMatrix[li_].ChiralityProjector[1] → DiracMatrix[li]/2,
    DiracMatrix[li_].ChiralityProjector[-1] → DiracMatrix[li]/2};
  temp = temp /. {exp_DiracTrace → FactorSUNT[exp]};
  temp = temp /. {IndexDelta[Index[x1_, y1_], Index[x2_, y2_]] → 1, SUNT[c_, _, _] → SUNT[c], SumOver[_] → 1,
    PolarizationVector[_] → 1};
  Return[temp];
];

Amp = FAToFC[all];

Amp = Amp /. fm : DiracSpinor[ $\frac{p_4^4}{2} + q_4$ , mc] . (mx__ /; FreeQ[{mx}, DiracSpinor]). DiracSpinor[q3 -  $\frac{p_3^3}{2}$ , mc] → $LineA1(Dot[mx]);
Amp = Amp /. fm : DiracSpinor[ $\frac{p_3^3}{2} + q_3$ , mc] . (mx__ /; FreeQ[{mx}, DiracSpinor]). DiracSpinor[q4 -  $\frac{p_4^4}{2}$ , mc] → $LineA2(Dot[mx]);
Amp = Amp /. fm : DiracSpinor[ $\frac{p_4^4}{2} + q_4$ , mc] . (mx__ /; FreeQ[{mx}, DiracSpinor]). DiracSpinor[q4 -  $\frac{p_4^4}{2}$ , mc] → $LineB1(Dot[mx]);
Amp = Amp /. fm : DiracSpinor[ $\frac{p_3^3}{2} + q_3$ , mc] . (mx__ /; FreeQ[{mx}, DiracSpinor]). DiracSpinor[q3 -  $\frac{p_3^3}{2}$ , mc] → $LineB2(Dot[mx]);

SetDirectory[ToFileName[NotebookDirectory[], "dmp"]];
Export["Amp.m", Amp];
ResetDirectory[];
]

```

$$e^+ + e^- \rightarrow \gamma^* \rightarrow J/\psi + \eta_c$$

Dirac- and Color-Algebra Simplification

- Covariant spin projectors Bodwin, Petrelli (2002) & Braaten, Lee (2003)

$$v(\bar{p})\bar{u}(p) \rightarrow \frac{1}{4\sqrt{2}E(E+m_c)}(\not{p}-m_c)[\gamma_5,\not{\epsilon}^*](\not{P}+2E)(\not{p}+m_c) \otimes \left[\frac{1}{\sqrt{N_c}}, T^a \right]$$

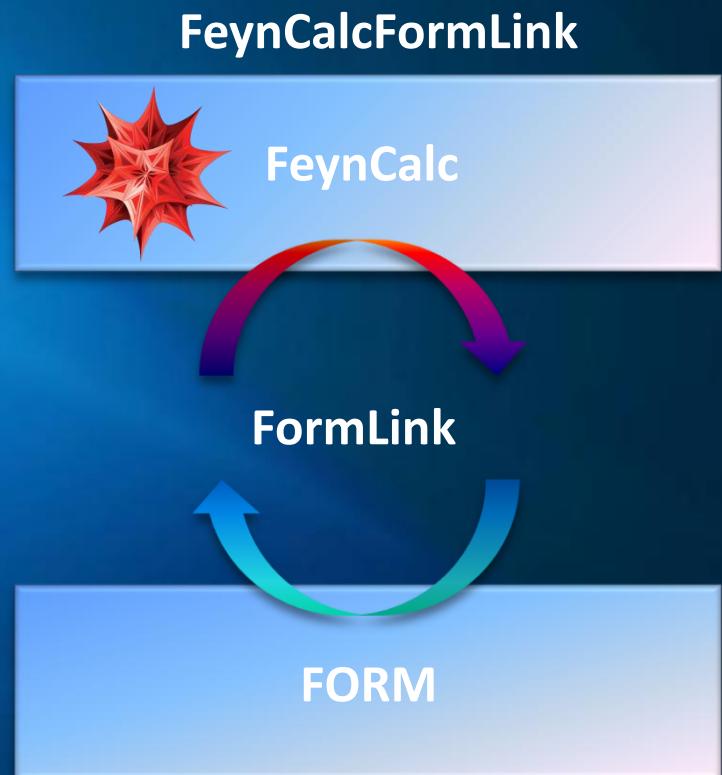
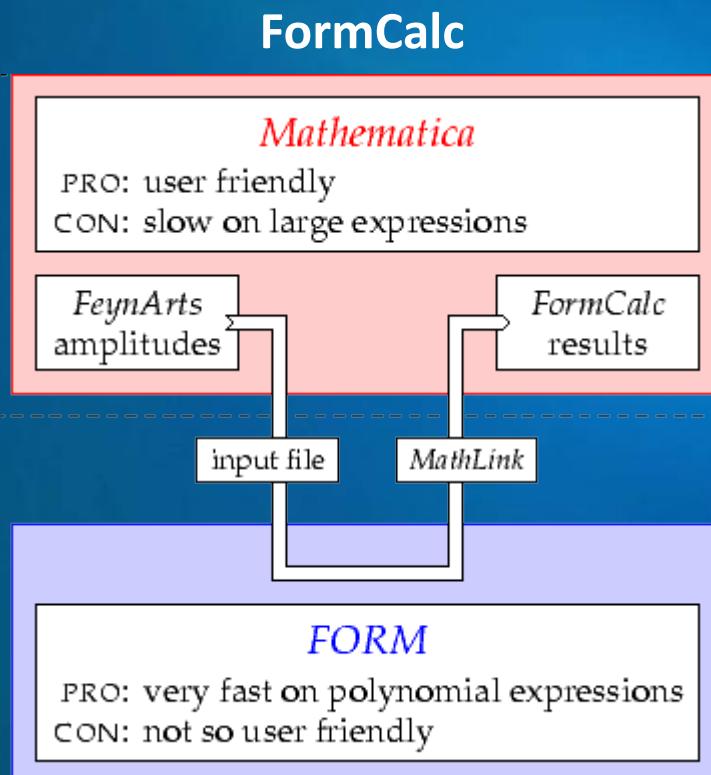
- Performing Dirac- and Color-algebra

- FORM – written in C
<http://www.nikhef.nl/~form/>
- FeynCalc – Mathematica 9
<http://www.feyncalc.org/>
- FeynCalcFormLink – Combine FORM & FeynCalc
<http://www.feyncalc.org/formlink/>

{
 FormCalc : Exchange data through Input & Output files
 FormLink : Exchange data through Piping

Introduction to FeynCalcFormLink

- Communication Between Processes
 - Input & Output Files
 - Piping



Introduction to FeynCalcFormLink

- Basic ideas of FormLink

- FormLink : create two unnamed pipes: r#, w#
- FormLink : form -pipe r#, w# init
- PID in w# : FORM → FormLink
- PID,PID in r# : FORM ← FormLink
- Form start running init.frm

- Content of init.frm

```
Off Statistics;  
#ifndef 'PIPES_'  
#message "No pipes found";  
.end;  
#endif  
#if ('PIPES_' <= 0)  
#message "No pipes found";  
.end;  
#endif
```

```
#procedure put(fmt, mexp)  
#toexternal 'fmt', 'mexp'  
#toexternal "#THE-END-MARK#"  
#endprocedure  
#setexternal 'PIPE1_';  
#toexternal "OK"  
#fromexternal  
.end
```

Introduction to FeynCalcFormLink

• Interactive FORM

```
In[1]:= << FormLink`  
  
In[2]:= FormStart[];  
  
In[3]:= FormWrite["Symbol a,b,c"];  
  
In[4]:= FormWrite["L exp=(a+b+c)^3;"];  
  
In[5]:= FormWrite[".sort;"];  
  
In[6]:= FormWrite["#call put(\"%E\", exp);"];  
  
In[7]:= FormWrite[".sort;"];  
  
In[8]:= TraditionalForm@ToExpression@FormRead[]  
  
Out[8]//TraditionalForm=  
a3 + 3 a2 b + 3 a2 c + 3 a b2 + 6 a b c + 3 a c2 + b3 + 3 b2 c + 3 b c2 + c3  
  
In[9]:= FormWrite["L exp2 = exp^2;"];  
  
In[10]:= FormWrite[".sort;"];  
  
In[11]:= FormWrite["#call put(\"%E\", exp2);"];  
  
In[12]:= FormWrite[".sort;"];  
  
In[13]:= TraditionalForm@ToExpression@FormRead[]  
  
Out[13]//TraditionalForm=  
a6 + 6 a5 b + 6 a5 c + 15 a4 b2 + 30 a4 b c + 15 a4 c2 + 20 a3 b3 + 60 a3 b2 c + 60 a3 b c2 + 20 a3 c3 + 15 a2 b4 + 60 a2 b3 c + 90 a2 b2 c2 +  
60 a2 b c3 + 15 a2 c4 + 6 a5 b + 30 a5 c + 60 a4 b2 c + 60 a4 b3 c + 30 a4 b c2 + 60 a4 c3 + 6 a5 c + b6 + 6 b5 c + 15 b4 c2 + 20 b3 c3 + 15 b2 c4 + 6 b c5 + c6|  
9
```

Introduction to FeynCalcFormLink

• A Simple Trace with FeynCalc

```
In[1]:= << HighEnergyPhysics`fc`  
  
In[2]:= Tr[GS[p1, p2, p3, p4, p5, p6]]  
  
Out[2]= 4 (p1·p6 p2·p5 p3·p4 - p1·p5 p2·p6 p3·p4 + p1·p2 p3·p4 p5·p6 - p1·p6 p2·p4 p3·p5 + p1·p4 p2·p6 p3·p5 +  
p1·p5 p2·p4 p3·p6 - p1·p4 p2·p5 p3·p6 + p1·p6 p2·p3 p4·p5 - p1·p3 p2·p6 p4·p5 + p1·p2 p3·p6 p4·p5 -  
p1·p5 p2·p3 p4·p6 + p1·p3 p2·p5 p4·p6 - p1·p2 p3·p5 p4·p6 + p1·p4 p2·p3 p5·p6 - p1·p3 p2·p4 p5·p6)
```

• Do it with FeynCalcFormLink

```
In[8]:= ? FormLink
```

FormLink[" form statements "] runs the form
statements in Form and returns the result to Mathematica.

If only one Local assignment is present the return value is that result. If
more Form Local variables are present, a list of results is returned.

```
In[9]:= ? FeynCalcFormLink
```

FeynCalcFormLink[expr] translates the FeynCalc expression expr to FORM,
calculates it, pipes it back to Mathematica and translates it to FeynCalc syntax.

Introduction to FeynCalcFormLink

Do it with FeynCalcFormLink

```
In[1]:= << FeynCalcFormLink`  
  
In[2]:= exp = DiracTrace[GS[p1, p2, p3, p4, p5, p6]]  
  
Out[2]=  $\text{tr}((\gamma \cdot p1).(\gamma \cdot p2).(\gamma \cdot p3).(\gamma \cdot p4).(\gamma \cdot p5).(\gamma \cdot p6))$   
  
In[3]:= FeynCalcFormLink[exp]  
  
Vectors p1,p2,p3,p4,p5,p6;  
Format Mathematica;  
L resFL = (g_(1,p1)*g_(1,p2)*g_(1,p3)*g_(1,p4)*g_(1,p5)*g_(1,p6));  
trace4,1;  
contract 0;  
.sort;  
#call put("%E", resFL)  
#fromexternal  
  
Piping the script to FORM and running FORM  
Time needed by FORM : 0.016 seconds. FORM finished. Got the result back to Mathematica as a string.  
Start translation to Mathematica / FeynCalc syntax  
Total wall clock time used: 1.31 seconds. Translation to Mathematica and FeynCalc finished.  
  
Out[3]= 4 p1 · p6 p2 · p5 p3 · p4 - 4 p1 · p5 p2 · p6 p3 · p4 + 4 p1 · p2 p3 · p4 p5 · p6 - 4 p1 · p6 p2 · p4 p3 · p5 + 4 p1 · p4 p2 · p6 p3 · p5 +  
4 p1 · p5 p2 · p4 p3 · p6 - 4 p1 · p4 p2 · p5 p3 · p6 + 4 p1 · p6 p2 · p3 p4 · p5 - 4 p1 · p3 p2 · p6 p4 · p5 + 4 p1 · p2 p3 · p6 p4 · p5 -  
4 p1 · p5 p2 · p3 p4 · p6 + 4 p1 · p3 p2 · p5 p4 · p6 - 4 p1 · p2 p3 · p5 p4 · p6 + 4 p1 · p4 p2 · p3 p5 · p6 - 4 p1 · p3 p2 · p4 p5 · p6
```

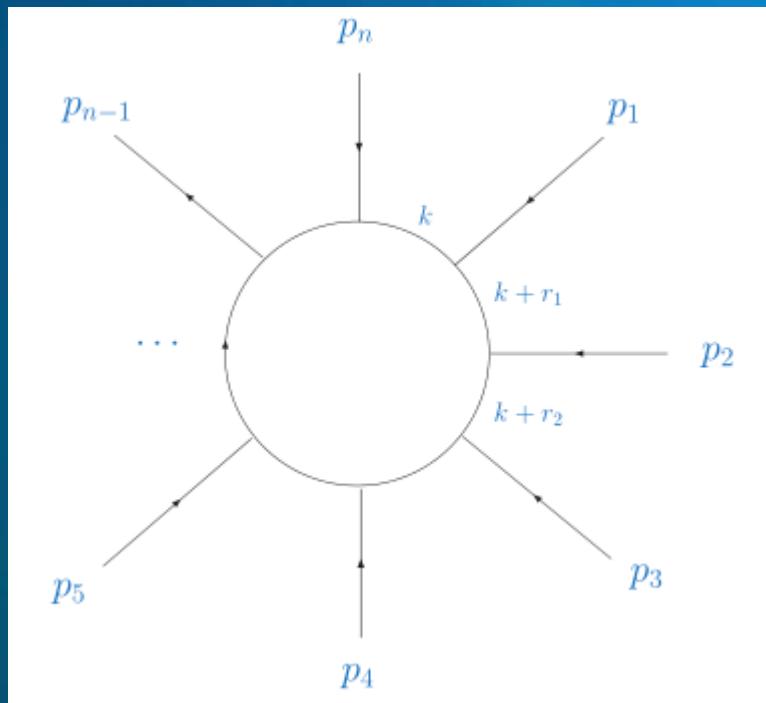
The Method of Region Expansion in NRQCD

- Expanding the relative momentum q Before or After loop integration?
 - Generally, we expand q after performing the loop integration, and then project the S-, P- or D-waves for the quarkonium.
 - We can also expand q before performing the loop integration, as long as only the hard region concerned, according to the method of region expansion [Benek, Smirnov (1998)].
- Hard region
 - If the NRQCD factorization is valid, it will be safe to use the method of region expand to compute the short-distance coefficients, which correspond to the hard region.

Passarino-Veltman Reduction

- Generic one-loop integral

$$\mathcal{T}^{\mu_1 \dots \mu_p} \equiv \frac{(2\pi\mu)^{4-d}}{i\pi^2} \int d^d k \frac{k^{\mu_1} \dots k^{\mu_p}}{D_0 D_1 D_2 \dots D_{n-1}}$$



$$D_i = (k + r_i)^2 - m_i^2 + i\varepsilon$$

$$r_i = \sum_{k=1}^i p_k, \quad i = 1, \dots, n-1$$

$$r_0 = \sum_{k=1}^n p_k = 0$$

$$r_{ij} = r_i - r_j$$

Passarino-Veltman Reduction

- Generic n-Point Tensor integrals

$$B^\mu(r_{10}^2, m_0^2, m_1^2) = \frac{(2\pi\mu)^{4-d}}{i\pi^2} \int d^d k \ k^\mu \prod_{i=0}^1 \frac{1}{(k+r_i)^2 - m_i^2}$$

$$B^{\mu\nu}(r_{10}^2, m_0^2, m_1^2) = \frac{(2\pi\mu)^{4-d}}{i\pi^2} \int d^d k \ k^\mu k^\nu \prod_{i=0}^1 \frac{1}{(k+r_i)^2 - m_i^2}$$

$$C^\mu(r_{10}^2, r_{12}^2, r_{20}^2, m_0^2, m_1^2, m_2^2) = \frac{(2\pi\mu)^{4-d}}{i\pi^2} \int d^d k \ k^\mu \prod_{i=0}^2 \frac{1}{(k+r_i)^2 - m_i^2}$$

$$C^{\mu\nu}(r_{10}^2, r_{12}^2, r_{20}^2, m_0^2, m_1^2, m_2^2) = \frac{(2\pi\mu)^{4-d}}{i\pi^2} \int d^d k \ k^\mu k^\nu \prod_{i=0}^2 \frac{1}{(k+r_i)^2 - m_i^2}$$

$$C^{\mu\nu\rho}(r_{10}^2, r_{12}^2, r_{20}^2, m_0^2, m_1^2, m_2^2) = \frac{(2\pi\mu)^{4-d}}{i\pi^2} \int d^d k \ k^\mu k^\nu k^\rho \prod_{i=0}^2 \frac{1}{(k+r_i)^2 - m_i^2}$$

$$D^\mu(r_{10}^2, r_{12}^2, r_{23}^2, r_{30}^2, r_{20}^2, r_{13}^2, m_0^2, m_1^2, m_2^2, m_3^2) = \frac{(2\pi\mu)^{4-d}}{i\pi^2} \int d^d k \ k^\mu \prod_{i=0}^3 \frac{1}{(k+r_i)^2 - m_i^2}$$

$$D^{\mu\nu}(r_{10}^2, r_{12}^2, r_{23}^2, r_{30}^2, r_{20}^2, r_{13}^2, m_0^2, m_1^2, m_2^2, m_3^2) = \frac{(2\pi\mu)^{4-d}}{i\pi^2} \int d^d k \ k^\mu k^\nu \prod_{i=0}^3 \frac{1}{(k+r_i)^2 - m_i^2}$$

$$D^{\mu\nu\rho}(r_{10}^2, r_{12}^2, r_{23}^2, r_{30}^2, r_{20}^2, r_{13}^2, m_0^2, m_1^2, m_2^2, m_3^2) = \frac{(2\pi\mu)^{4-d}}{i\pi^2} \int d^d k \ k^\mu k^\nu k^\rho \prod_{i=0}^3 \frac{1}{(k+r_i)^2 - m_i^2}$$

$$D^{\mu\nu\rho\sigma}(r_{10}^2, r_{12}^2, r_{23}^2, r_{30}^2, r_{20}^2, r_{13}^2, m_0^2, m_1^2, m_2^2, m_3^2) = \frac{(2\pi\mu)^{4-d}}{i\pi^2} \int d^d k \ k^\mu k^\nu k^\rho k^\sigma \prod_{i=0}^3 \frac{1}{(k+r_i)^2 - m_i^2}$$

Passarino-Veltman Reduction

- Generic n-Point Scalar integrals

$$A_0(m_0^2) = \frac{(2\pi\mu)^{4-d}}{i\pi^2} \int d^d k \frac{1}{k^2 - m_0^2}$$

$$B_0(r_{10}^2, m_0^2, m_1^2) = \frac{(2\pi\mu)^{4-d}}{i\pi^2} \int d^d k \prod_{i=0}^1 \frac{1}{(k + r_i)^2 - m_i^2}$$

$$C_0(r_{10}^2, r_{12}^2, r_{20}^2, m_0^2, m_1^2, m_2^2) = \frac{(2\pi\mu)^{4-d}}{i\pi^2} \int d^d k \prod_{i=0}^2 \frac{1}{(k + r_i)^2 - m_i^2}$$

$$D_0(r_{10}^2, r_{12}^2, r_{23}^2, r_{30}^2, r_{20}^2, r_{13}^2, m_0^2, m_1^2, m_2^2, m_3^2) = \frac{(2\pi\mu)^{4-d}}{i\pi^2} \int d^d k \prod_{i=0}^3 \frac{1}{(k + r_i)^2 - m_i^2}$$

- Pa-Ve Reduction

$$\begin{aligned} D^{\mu\nu\rho\sigma} &= (g^{\mu\nu}g^{\rho\sigma} + g^{\mu\rho}g^{\nu\sigma} + g^{\mu\sigma}g^{\nu\rho})D_{0000} \\ &= + \sum_{i,j}^3 (g^{\mu\nu}r_i^\rho r_j^\sigma + g^{\nu\rho}r_i^\mu r_j^\sigma + g^{\mu\rho}r_i^\nu r_j^\sigma + g^{\mu\sigma}r_i^\nu r_j^\rho + g^{\nu\sigma}r_i^\mu r_j^\rho + g^{\rho\sigma}r_i^\mu r_j^\nu)D_{00ij} \\ &= + \sum_{i,j,k,l=1}^3 r_\mu^i r_j^\nu r_k^\rho r_l^\sigma D_{ijkl} \end{aligned}$$

Passarino-Veltman Reduction

- Passarino-Veltman functions in FeynCalc

```
In[1]:= << HighEnergyPhysics`fc`  
  
In[2]:= PaVe[0, 0, 0, 0, {r10^2, r12^2, r23^2, r30^2, r20^2, r13^2}, {m0^2, m1^2, m2^2, m3^2}]  
Out[2]= D0 0 0 0  
  
In[3]:= PaVe[0, 1, 2, 3, {r10^2, r12^2, r23^2, r30^2, r20^2, r13^2}, {m0^2, m1^2, m2^2, m3^2}]  
Out[3]= D0 1 2 3  
  
In[4]:= PaVe[0, 0, 0, 0, {1, 2, 3, 4, 5, 6}, {1, 1, 1, 1}] // PaVeReduce  
Out[4]= - $\frac{135 C_0(1, 2, 5, 1, 1, 1)}{2401} - \frac{5751 C_0(1, 4, 6, 1, 1, 1)}{192080} - \frac{88691 C_0(2, 3, 6, 1, 1, 1)}{2650704} -$   
 $\frac{5755 C_0(3, 4, 5, 1, 1, 1)}{633864} + \frac{1587 D_0(1, 2, 3, 4, 5, 6, 1, 1, 1, 1)}{38416} + \frac{51 B_0(1, 1, 1)}{1960} +$   
 $\frac{907 B_0(2, 1, 1)}{27048} + \frac{1025 B_0(3, 1, 1)}{99176} + \frac{347 B_0(4, 1, 1)}{32340} - \frac{50 B_0(5, 1, 1)}{1617} - \frac{181 B_0(6, 1, 1)}{22540} + \frac{5}{72}$ 
```

- Break-down for Gram Determinant = 0
 - Happens when we expand the relative momentum q before loop integration, due to taking the derivative over q.

Integrate-By-Part (IBP) Reduction

- Integrate By Part (IBP)

[JHEP 0810, 107 (2008)]

$$F(a_1, \dots, a_n) = \int \dots \int \frac{d^d k_1 \cdots d^d k_h}{E_1^{a_1} \cdots E_n^{a_n}}$$

where k_i , $i = 1, \dots, h$, are loop momenta and the denominators E_r are either quadratic or linear with respect to the loop momenta k_i of the graph. Irreducible polynomials in the numerator can be represented as denominators raised to negative powers.

- Basic idea:

$$\int \dots \int d^d k_1 d^d k_2 \cdots \frac{\partial}{\partial k_i} \left[\frac{p_j}{E_1^{a_1} \cdots E_n^{a_n}} \right] = 0$$

- List of equations:

$$\sum \alpha_i F(a_1 + b_{i,1}, \dots, a_n + b_{i,n}) = 0$$

- Master Integrals

- Propagators should be independent!

Partial Fraction on Propagators

- Apart in Mathematica

```
In[1]:= Apart[1/((x - a) (x - b) (x - c))]  
Out[1]= 1/(a - b) (b - c) (x - b) - 1/(a - c) (c - b) (x - c) + 1/(a - b) (a - c) (x - a)  
  
In[2]:= Apart[1/((x + a) (y + 2 b) (3 x + 4 y))]  
Out[2]= 1/(a + x) (3 x - 8 b) (2 b + y) - 4/(a + x) (3 x - 8 b) (3 x + 4 y)
```

- Do it with \$Apart function

```
In[1]:= << CalcExt`Apart`  
  
In[2]:= $Apart[1/((x - a) (x - b) (x - c)), {x}]  
Out[2]= -(1/(a - x))/(a - b) (a - c) + (1/(b - x))/(a - b) (b - c) + (1/(c - x))/(a - c) (c - b)  
  
In[3]:= $Apart[1/((x + a) (y + 2 b) (3 x + 4 y)), {x, y}]  
Out[3]= -(1/(a + x) (2 b + y))/(3 a + 8 b) + 4 (1/(a + x) (3 x + 4 y))/(3 a + 8 b) + 3 (1/(2 b + y) (3 x + 4 y))/(3 a + 8 b)
```

Partial Fraction on Propagators

- What does Linear Independent mean?
 - V_x : vector space spanned by n linear-independent vectors: $\{x_i\}$ over the coefficient field \mathcal{F} .
 - $V_x^* \equiv V_x \oplus \mathcal{F}$
 - $\{e_i = v_i + f_i\}$ are linear independent if and only if $\{v_i\}$ are linear independent
 - $\sum f_i e_i = f \quad (\text{not all } f_i = 0)$
- Product Operator: inherited from n -variate polynormial
- What does \$Apart do? [Comput.Phys.Commun.183,2158(2012)]

$$\prod_{i=1}^N e_i^{n_i} = \sum_j f_j \prod_{i=1}^{N_j} e_{k_{ji}}^{n_{ji}}$$

where $0 < k_{ji} < N$, and $\{e_{k_{ji}}\}$ are lienar independent.

Partial Fraction on Propagators

- A simply physical example:

$$\exp = \frac{(k \cdot p_1)(k \cdot p_2)}{k^2[(k + p_1)^2 - m^2][(k + p_2)^2 - m^2]}$$

Taking k^2 , $k \cdot p_1$ and $k \cdot p_2$ as indepdent vectors.

- Partial Fraction on it:

```
In[1]:= << CalcExt` Apart`  
  
In[2]:= << HighEnergyPhysics` fc`  
  
In[3]:= exp =  $\frac{\text{SP}[\mathbf{k}, \mathbf{p1}] \text{SP}[\mathbf{k}, \mathbf{p2}]}{\text{SP}[\mathbf{k}] (\text{SP}[\mathbf{k} + \mathbf{p1}] - m^2) (\text{SP}[\mathbf{k} + \mathbf{p2}] - m^2)^2}$  // FCI // ScalarProductExpand;  
  
In[4]:= xs = FCI /@ {SP[k], SP[k, p1], SP[k, p2]}  
  
Out[4]= {k2, k · p1, k · p2}  
  
In[5]:= $Apart[exp, xs]  
  
Out[5]= 
$$-\frac{1}{4} \left( \frac{1}{(m^2 - k^2 - 2k · p2 - p2^2)^2} \right) + \frac{1}{4} (m^2 - p1^2) \left( \frac{1}{(m^2 - k^2 - 2k · p1 - p1^2)(m^2 - k^2 - 2k · p2 - p2^2)^2} \right) + \frac{1}{4} (m^2 - p2^2) \left( \frac{1}{k^2 (-m^2 + k^2 + 2k · p2 + p2^2)^2} \right) +$$
  

$$\frac{1}{2} \left( \frac{k · p2}{(m^2 - k^2 - 2k · p1 - p1^2)(-m^2 + k^2 + 2k · p2 + p2^2)^2} \right) + \frac{1}{4} (m^2 - p1^2)(m^2 - p2^2) \left( \frac{1}{k^2 (-m^2 + k^2 + 2k · p1 + p1^2)(-m^2 + k^2 + 2k · p2 + p2^2)^2} \right) +$$
  

$$\frac{1}{4} \left( \frac{1}{k^2 (-m^2 + k^2 + 2k · p2 + p2^2)} \right) + \frac{1}{4} (m^2 - p1^2) \left( \frac{1}{k^2 (-m^2 + k^2 + 2k · p1 + p1^2)(-m^2 + k^2 + 2k · p2 + p2^2)} \right)$$

```

Integrate-By-Part (IBP) Reduction

- Public codes performing IBP reduction
 - AIR – Maple JHEP **0407** (2004) 046
 - FIRE – Mathematica/C++ JHEP **0810**, 107 (2008)
 - Reduze – C++ Comput. Phys. Commun. 181, 1293 (2010)
 - LiteRed – Mathematica arXiv:1212.2685 (2012)
 - Many other private codes
 - FIRE – Feynman Integral Reduction

```
In[5]:= $Apart[exp, xs]
Out[5]= -\frac{1}{4} \left\| \frac{1}{(m^2 - k^2 - 2k \cdot p2 - p2^2)^2} \right\| + \frac{1}{4} (m^2 - p1^2) \left\| \frac{1}{(m^2 - k^2 - 2k \cdot p1 - p1^2)(m^2 - k^2 - 2k \cdot p2 - p2^2)^2} \right\| + \frac{1}{4} (m^2 - p2^2) \left\| \frac{1}{k^2 (-m^2 + k^2 + 2k \cdot p2 + p2^2)^2} \right\| +
```

$$\frac{1}{2} \left\| \frac{k \cdot p2}{(m^2 - k^2 - 2k \cdot p1 - p1^2)(-m^2 + k^2 + 2k \cdot p2 + p2^2)^2} \right\| + \frac{1}{4} (m^2 - p1^2) (m^2 - p2^2) \left\| \frac{1}{k^2 (-m^2 + k^2 + 2k \cdot p1 + p1^2)(-m^2 + k^2 + 2k \cdot p2 + p2^2)^2} \right\| +$$

$$\frac{1}{4} \left\| \frac{1}{k^2 (-m^2 + k^2 + 2k \cdot p2 + p2^2)} \right\| + \frac{1}{4} (m^2 - p1^2) \left\| \frac{1}{k^2 (-m^2 + k^2 + 2k \cdot p1 + p1^2)(-m^2 + k^2 + 2k \cdot p2 + p2^2)} \right\|$$

Integrate-By-Part (IBP) Reduction

• FIRE – Feynman Integral Reduction

$$F[l, m, n] = \int \frac{d^4 k}{(2\pi)^4} \frac{(k \cdot p_2)^{-l}}{(m^2 - k^2 - 2k \cdot p_1 - p_1^2)^m (-m^2 + k^2 + 2k \cdot p_2 + p_2^2)^n}$$

• Basic usage of FIRE

```
In[1]:= << HighEnergyPhysics`fc`  
  
In[2]:= << FIRE`  
  
In[3]:= Replacement = {p1^2 → m^2, p2^2 → m^2, p1.p2 → SP[p1, p2]};  
Internal = {k};  
External = {p1, p2};  
Propagators = {k.p2, -2 k.p1 - k^2 + m^2 - p1^2, 2 k.p2 + k^2 - m^2 + p2^2};  
PrepareIBP[];  
startinglist = {IBP[k, k], IBP[k, p1], IBP[k, p2]} /. Replacement;  
Prepare[];  
  
In[10]:= Burn[];  
  
In[11]:= F[-1, 1, 2]  
Out[11]= 
$$\frac{(d-2) G(\{0, 0, 1\})}{8(m^2 - p1 \cdot p2)} + \frac{(d-2) G(\{0, 1, 0\})}{8(m^2 - p1 \cdot p2)} + \frac{1}{4} (4-d) G(\{0, 1, 1\})$$

```

Final Result

- After Partial Fraction and IBP Reduction:

$$\int \frac{d^4 k}{(2\pi)^4} \frac{(k \cdot p_1)(k \cdot p_2)}{k^2[(k + p_1)^2 - m^2][(k + p_2) - m^2]^2}$$



Assuming $p_1^2 = p_2^2 = m^2$:

$$\begin{aligned} & \text{Out[27]= } \frac{(D-2) \left\| \frac{1}{-2 k \cdot p1 - k^2} \right\|}{16 (m^2 - p1 \cdot p2)} + \frac{(D-2) \left\| \frac{1}{2 k \cdot p2 + k^2} \right\|}{16 (m^2 - p1 \cdot p2)} + \\ & \quad \frac{1}{8} (4-D) \left\| \frac{1}{(-2 k \cdot p1 - k^2) (2 k \cdot p2 + k^2)} \right\| + \frac{1}{4} \left\| \frac{1}{k^2 (2 k \cdot p2 + k^2)} \right\| \end{aligned}$$

where $\left\| \cdots \right\|$ is defined as:

$$\left\| \exp \right\| = \int \frac{d^4 k}{(2\pi)^4} \exp$$

Using Apart

```
In[2]:= << CalcExt`Apart`
```

```
In[3]:= exp =  $\frac{SP[k, p1] SP[k, p2]}{SP[k] (SP[k + p1] - m^2) (SP[k + p2] - m^2)^2}$  // FCI // ScalarProductExpand;
```

```
In[4]:= xs = FCI /@ {SP[k], SP[k, p1], SP[k, p2]}
```

```
Out[4]= {k2, k · p1, k · p2}
```

```
In[5]:= ApartReasult = $Apart[exp, xs]
```

```
Out[5]= 
$$-\frac{1}{4} \left( \frac{1}{(m^2 - k^2 - 2k · p2 - p2^2)^2} \right) + \frac{1}{4} (m^2 - p1^2) \left( \frac{1}{(m^2 - k^2 - 2k · p1 - p1^2)(m^2 - k^2 - 2k · p2 - p2^2)^2} \right) + \frac{1}{4} (m^2 - p2^2) \left( \frac{1}{k^2 (-m^2 + k^2 + 2k · p2 + p2^2)^2} \right) +$$


$$\frac{1}{2} \left( \frac{k · p2}{(m^2 - k^2 - 2k · p1 - p1^2)(-m^2 + k^2 + 2k · p2 + p2^2)^2} \right) + \frac{1}{4} (m^2 - p1^2)(m^2 - p2^2) \left( \frac{1}{k^2 (-m^2 + k^2 + 2k · p1 + p1^2)(-m^2 + k^2 + 2k · p2 + p2^2)^2} \right) +$$


$$\frac{1}{4} \left( \frac{1}{k^2 (-m^2 + k^2 + 2k · p2 + p2^2)} \right) + \frac{1}{4} (m^2 - p1^2) \left( \frac{1}{k^2 (-m^2 + k^2 + 2k · p1 + p1^2)(-m^2 + k^2 + 2k · p2 + p2^2)} \right)$$

```

```
In[6]:= ApartReasult = $ApartComplete[ApartReasult, xs];
```

Generate All Propagators

```
In[7]:= PropagatorsList = Union @@ ((\$ApartUnion[#] /. \$ApartIR[_, __, x_, __] :> x) & /@ Flatten[{ApartReasult}])
```

```
Out[7]= 
$$\begin{pmatrix} k^2 & -m^2 + k^2 + 2k · p1 + p1^2 & -m^2 + k^2 + 2k · p2 + p2^2 \\ k^2 & -m^2 + k^2 + 2k · p2 + p2^2 & k · p1 \\ k · p2 & m^2 - k^2 - 2k · p1 - p1^2 & -m^2 + k^2 + 2k · p2 + p2^2 \\ m^2 - k^2 - 2k · p1 - p1^2 & m^2 - k^2 - 2k · p2 - p2^2 & k^2 \\ m^2 - k^2 - 2k · p2 - p2^2 & k^2 & k · p1 \end{pmatrix}$$

```

```
In[8]:= For[i = 1, i ≤ Length[PropagatorsList], i++, PNS[PropagatorsList[[i]]] = i];
```

```
In[9]:= SPropagatorsList = PropagatorsList /. {FCI@SP[k] → k2, FCI@SP[k, p1] → k · p1, FCI@SP[k, p2] → k · p2};
```

Setup FIRE

```
In[10]:= << FIRE`
```

Init Parameters

```
In[11]:= Replacement = {p1^2 → m^2, p2^2 → m^2, p1 p2 → SP[p1, p2]} // FCI;
d = D;
```

Prepare

```
In[13]:= For[$i = 1, $i ≤ Length[SPropagatorsList], $i ++,
  ClearIBP[];
  Internal = {k};
  External = {p1, p2};
  Propagators = {k p2, -2 k p1 - k^2 + m^2 - p1^2, 2 k p2 + k^2 - m^2 + p2^2};
  PrepareIBP[];
  startinglist = {IBP[k, k], IBP[k, p1], IBP[k, p2]} /. Replacement;
  Prepare[$i];
];
ClearIBP[];
```

Burn

```
In[15]:= Burn[];
```

Results

```
In[16]:= Clear[IntF];
FireReduce[exp_] := exp /. $ApartIR[_,_,{x_,n_}] :> F[PNS[x], -n] /. G[pn_, ns_] :> IntF[PropagatorsList[[pn]], -ns];
IntF[x_List, n_List] := IntF[Times @@ x^n, q1];
IntF /: MakeBoxes[IntF[exp_, ___], TraditionalForm] := RowBox[{"", MakeBoxes[exp, TraditionalForm], ""}];
IntFCollect[exp_] := Collect[exp, _IntF];
```

```
In[21]:= FireReduce[ApartReasult] /. {FCI@SP[p1] :> m^2, FCI@SP[p2] :> m^2}
```

$$\text{Out[21]} = \frac{1}{2} \left(\frac{(D-2) \left\| \frac{1}{-2k \cdot p1 - k^2} \right\|}{8(m^2 - p1 \cdot p2)} + \frac{(D-2) \left\| \frac{1}{2k \cdot p2 + k^2} \right\|}{8(m^2 - p1 \cdot p2)} + \frac{1}{4} (4-D) \left\| \frac{1}{(-2k \cdot p1 - k^2)(2k \cdot p2 + k^2)} \right\| \right) + \frac{1}{4} \left\| \frac{1}{k^2 (2k \cdot p2 + k^2)} \right\|$$

```
In[22]:= IntFCollect[%]
```

$$\text{Out[22]} = \frac{(D-2) \left\| \frac{1}{-2k \cdot p1 - k^2} \right\|}{16(m^2 - p1 \cdot p2)} + \frac{(D-2) \left\| \frac{1}{2k \cdot p2 + k^2} \right\|}{16(m^2 - p1 \cdot p2)} + \frac{1}{8} (4-D) \left\| \frac{1}{(-2k \cdot p1 - k^2)(2k \cdot p2 + k^2)} \right\| + \frac{1}{4} \left\| \frac{1}{k^2 (2k \cdot p2 + k^2)} \right\|$$

Summary

- Using FeynArts for Feynman diagrams and amplitudes generation.
- Using FeynCalc/FormLink for Dirac- and Color-algebra simplification.
- Using \$Apart and FIRE for partial fraction and Integrate-By-Part reduction respectively.
- Solving Master Integrals by any other mean and substituting them to get the final result.
- Post-processing the final result in last previous step.

Thank You!

TID

```
In[1]:= << HighEnergyPhysics`fc`
```

```
In[2]:= << CalcExt`TID`
```

```
In[3]:= FVD[k, μ] f[SP[p1], SP[p2], SP[p, p2]] // FCI
```

```
Out[3]=  $k^\mu f(p_1^2, p_2^2, p \cdot p_2)$ 
```

```
In[4]:= $TID[FVD[k, μ] f[SP[p1], SP[p2], SP[p, p2]] // FCI, k, {p1, p2}]
```

```
Out[4]=  $\left( \frac{p_2^\mu k \cdot p_1 p_1 \cdot p_2 - p_1^2 p_2^\mu k \cdot p_2}{p_1 \cdot p_2^2 - p_1^2 p_2^2} + \frac{p_1^\mu k \cdot p_2 p_1 \cdot p_2 - p_2^2 p_1^\mu k \cdot p_1}{p_1 \cdot p_2^2 - p_1^2 p_2^2} \right) f(p_1^2, p_2^2, p \cdot p_2)$ 
```

```
In[5]:= FVD[k, μ] FVD[k, ν] f[SP[p1], SP[p2], SP[p, p2]] // FCI
```

```
Out[5]=  $k^\mu k^\nu f(p_1^2, p_2^2, p \cdot p_2)$ 
```

```
In[6]:= $TID[FVD[k, μ] FVD[k, ν] f[SP[p1], SP[p2], SP[p, p2]] // FCI, k, {p1, p2}]
```

```
Out[6]=  $f(p_1^2, p_2^2, p \cdot p_2)$ 
```

$$\left(\frac{g^{\mu\nu} (p_2^2 k \cdot p_1^2 - 2k \cdot p_2 p_1 \cdot p_2 k \cdot p_1 + k^2 p_1 \cdot p_2^2 + k \cdot p_2^2 p_1^2 - k^2 p_1^2 p_2^2)}{(D-2)(p_1 \cdot p_2^2 - p_1^2 p_2^2)} + \frac{1}{(D-2)(p_1 \cdot p_2^2 - p_1^2 p_2^2)^2} p_2^\mu p_2^\nu (D p_1 \cdot p_2^2 k \cdot p_1^2 - 2 p_1 \cdot p_2^2 k \cdot p_1^2 + p_1^2 p_2^2 k \cdot p_1^2 - 2 D k \cdot p_2 p_1^2 p_1 \cdot p_2 k \cdot p_1 + 2 k \cdot p_2 p_1^2 p_1 \cdot p_2 k \cdot p_1 + D k \cdot p_2^2 p_1^{22} - k \cdot p_2^2 p_1^{22} + k^2 p_1^2 p_1 \cdot p_2^2 - k^2 p_1^{22} p_2^2) + \frac{1}{(D-2)(p_1 \cdot p_2^2 - p_1^2 p_2^2)^2} p_2^\mu p_1^\nu (-k^2 p_1 \cdot p_2^3 + D k \cdot p_1 k \cdot p_2 p_1 \cdot p_2^2 - D k \cdot p_2^2 p_1^2 p_1 \cdot p_2 + k \cdot p_2^2 p_1^2 p_1 \cdot p_2 - D k \cdot p_1^2 p_2^2 p_1 \cdot p_2 + k \cdot p_1^2 p_2^2 p_1 \cdot p_2 + k^2 p_1^2 p_1 \cdot p_2^2 p_1 \cdot p_2 + D k \cdot p_1 k \cdot p_2 p_1^2 p_2^2 - 2 k \cdot p_1 k \cdot p_2 p_1^2 p_2^2) + \frac{1}{(D-2)(p_1 \cdot p_2^2 - p_1^2 p_2^2)^2} p_1^\mu p_2^\nu (-k^2 p_1 \cdot p_2^3 + D k \cdot p_1 k \cdot p_2 p_1 \cdot p_2^2 - D k \cdot p_2^2 p_1^2 p_1 \cdot p_2 + k \cdot p_2^2 p_1^2 p_1 \cdot p_2 - D k \cdot p_1^2 p_2^2 p_1 \cdot p_2 + k \cdot p_1^2 p_2^2 p_1 \cdot p_2 + k^2 p_1^2 p_1 \cdot p_2^2 p_1 \cdot p_2 + D k \cdot p_1 k \cdot p_2 p_1^2 p_2^2 - 2 k \cdot p_1 k \cdot p_2 p_1^2 p_2^2) + \frac{1}{(D-2)(p_1 \cdot p_2^2 - p_1^2 p_2^2)^2} p_1^\mu p_1^\nu (D p_1 \cdot p_2^2 k \cdot p_2^2 - 2 p_1 \cdot p_2^2 k \cdot p_2^2 + p_1^2 p_2^2 k \cdot p_2^2 - 2 D k \cdot p_1 p_1 \cdot p_2 p_2^2 k \cdot p_2 + 2 k \cdot p_1 p_1 \cdot p_2 p_2^2 k \cdot p_2 + D k \cdot p_1^2 p_2^{22} - k \cdot p_1^2 p_2^{22} - k^2 p_1^2 p_2^{22} + k^2 p_1 \cdot p_2^2 p_2^2) \right)$$