



第八届全国重味物理和量子色动力学研讨会 (重庆)

2026年4月24日-28日

New Perspectives on Precision Nucleon Tomography

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Phys. Rev. Lett. 136 (2026) 151902

Phys. Rev. Lett. 136 (2026) 021901

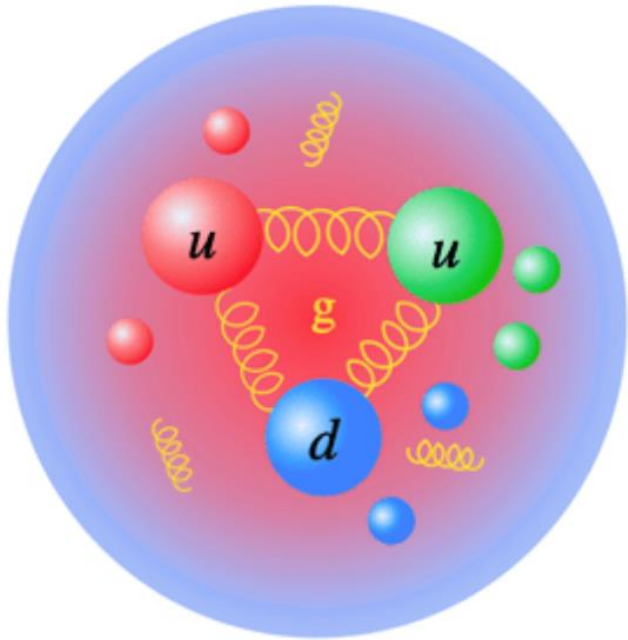
Phys. Rev. Lett. 135 (2025) 171903

JHEP 11 (2025) 082

2602.20249, 2602.22972, 2510.13951

$$\mathcal{L}_{\text{QCD}} = \bar{\psi}_i (i(\gamma^\mu D_\mu)_{ij} - m \delta_{ij}) \psi_j - \frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu}$$

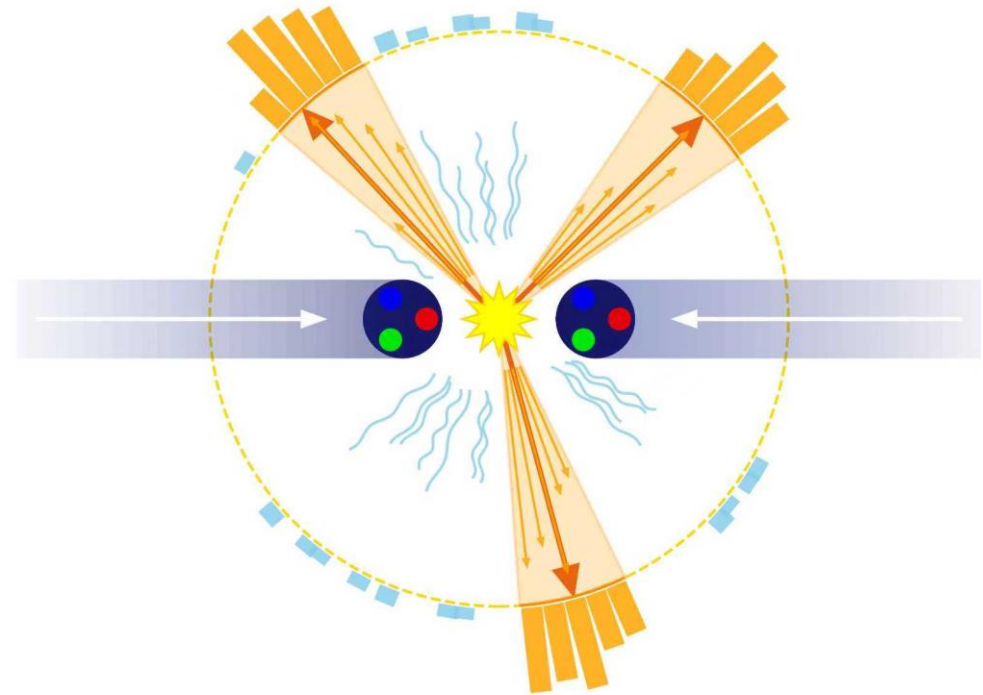
强子结构



色禁闭
渐近自由
红外奇点

.....

喷注产生



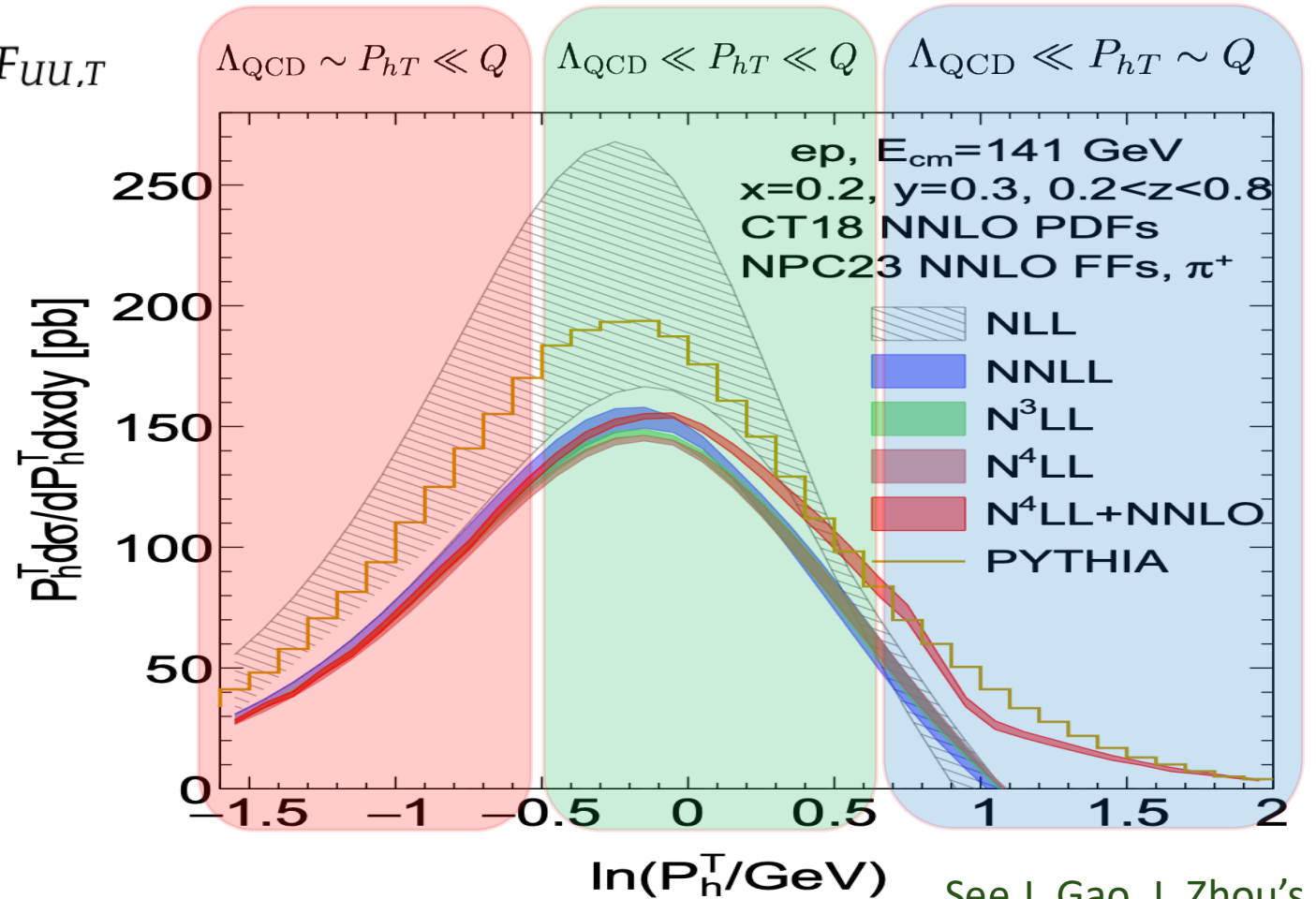
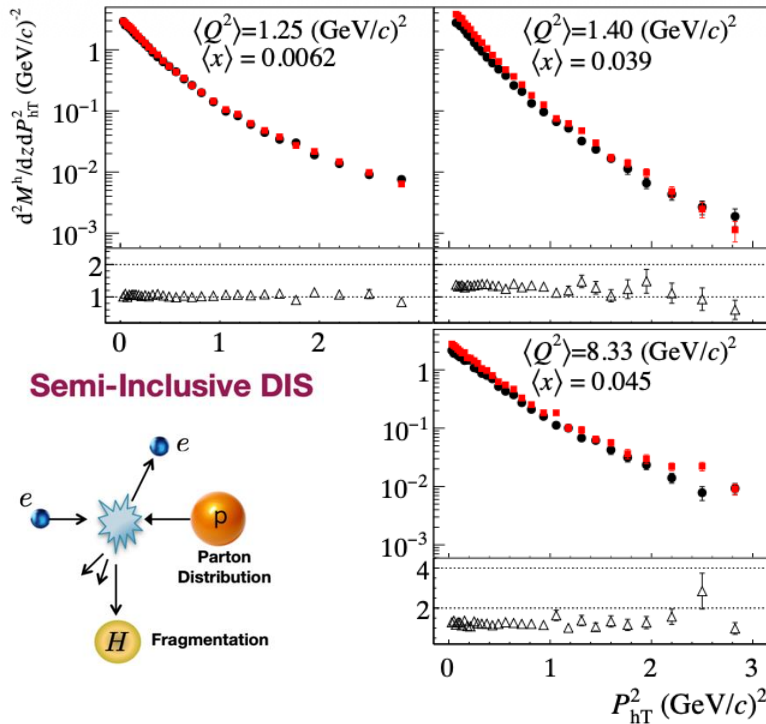


Quark TMDs in e-p collisions

One of the most fundamental measurements in SIDIS related to TMDs is the study of the unpolarized P_{hT} differential cross section

$$\frac{d^4\sigma_{\text{SIDIS}}}{dx dy dz_h dP_{hT}} = 4\pi P_{hT} \frac{\alpha_{em}^2}{x y Q^2} \left(1 - y + \frac{1}{2}y^2\right) F_{UU,T}$$

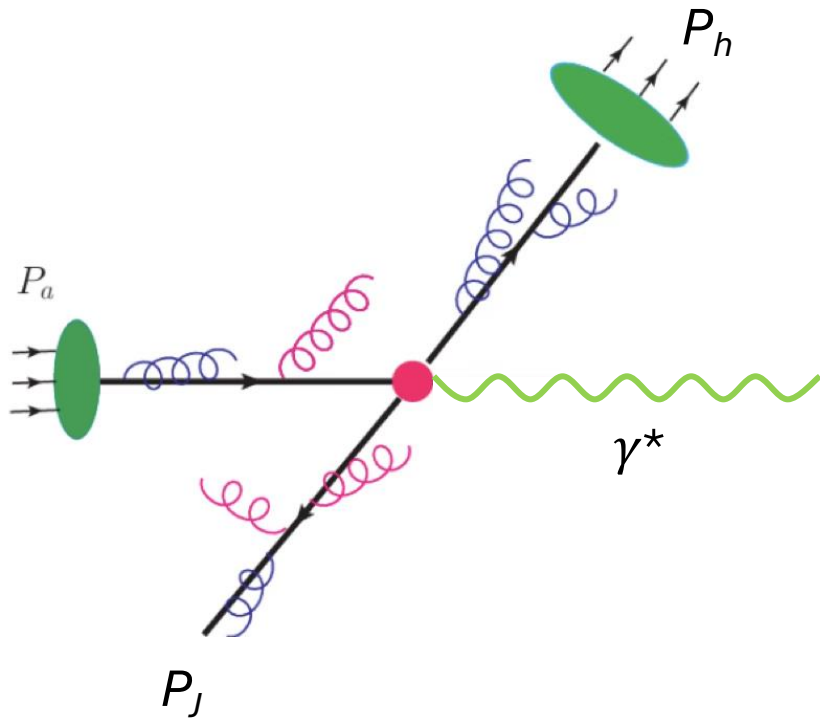
Gao, Li, Zhu, Zhu '26



See J. Gao, J. Zhou's talks

NNLO QCD corrections for P_{hT} distribution

Hadron + jets @ NNLO



- A key challenge is handling infrared (IR) divergences, which must cancel between real and virtual configurations.

$$\sigma_{\text{tot}} = \left| \begin{array}{c} e^+e^- \rightarrow q\bar{q} \\ \text{tree} + \text{loop} \end{array} \right|^2 + \left| \begin{array}{c} e^+e^- \rightarrow q\bar{q}g \\ \text{tree} + \text{tree} \end{array} \right|^2$$

- **FMNLO** framework for hadron production at NLO Liu, Shen, Zhou, Gao '23
- **Transverse momentum (qT) subtraction** for hadron production at NNLO Gao, Li, Zhu, Zhu '26
- **qT subtraction works well for colorless or colored (but massive) system**
 - **N3LO Drell-Yan** Chen, Gehrmann, Glover, Huss, Yang, Zhu '23
 - **N3LO di-photon** Czakon, Eschment, Generet, Poncelet '26
- **For processes with jets, qT is unsuitable because qT=0 for radiation emitted inside jets.**

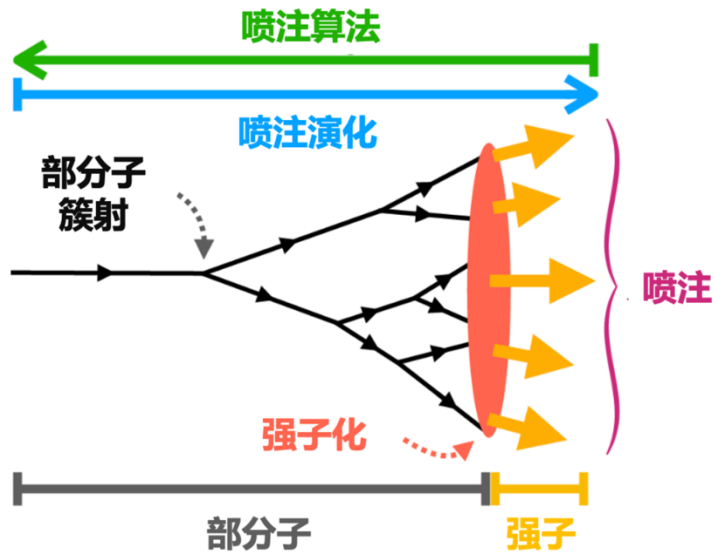


Winner-take-all qT subtraction for jets

Fu(符荣峻), Rahn, DYS, Waalewijn, Wu PRL135 (2025) 171903

- **Winner-take-all recombination scheme** Bertolini, Chan, Thaler '13

Also see S. Y. Wei's talk



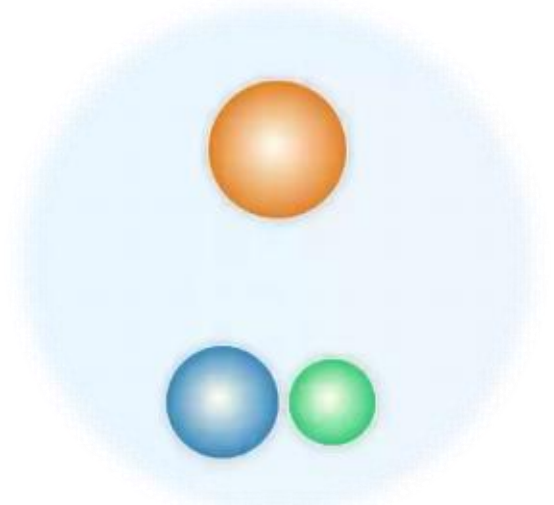
“赢者通吃” 动量重组方案

$$p_{t,r} = p_{t,i} + p_{t,j}$$

$$\phi_r = (w_i \phi_i + w_j \phi_j) / (w_i + w_j)$$

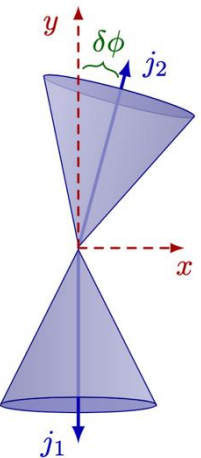
$$y_r = (w_i y_i + w_j y_j) / (w_i + w_j)$$

$$w_i = (p_{t,i})^n \quad n \rightarrow \infty$$



- **Factorization theorem is significantly simplified** Chien, Rahn, DYS, Waalewijn & Wu '22 JHEP + Schrignder '21 PLB

- **NNLL resummation for $\delta\phi$ @ ep & eA** Fang(方申), Ke, DYS, Terry '24 JHEP
- **N3LL + NNLO resummation for $\delta\phi$ @ ep** Fang(方申), Gao(高梅森), Li, DYS '25 JHEP





Winner-take-all qT subtraction for jets

Fu(符荣峻), Rahn, DYS, Waalewijn, Wu PRL135 (2025) 171903

- We propose two ways of extending qT subtraction to processes with jets

The key ingredient is the use of a recoil-free WTA jet axis!

- Azimuthal decorrelation: $\delta\phi$ or q_x

- Suitable for processes that are **planar at LO**: $pp \rightarrow 2 \text{ jets}, \gamma^* p \rightarrow jj, e^+e^- \rightarrow 3 \text{ jets}, \text{ etc.};$

- **Standard TMD PDFs and soft functions: known at N3LO** Luo, Yang, Zhu, Zhu, '19; Ebert, Mistlberger, Vita '20 ...

- **TMD jet functions: partially known at NNLO**; Reyes, Scimemi, Waalewijn, Zoppi '19; Bell, Brune, Das, Wald '23; Fang, Gao, Li, DYS '24

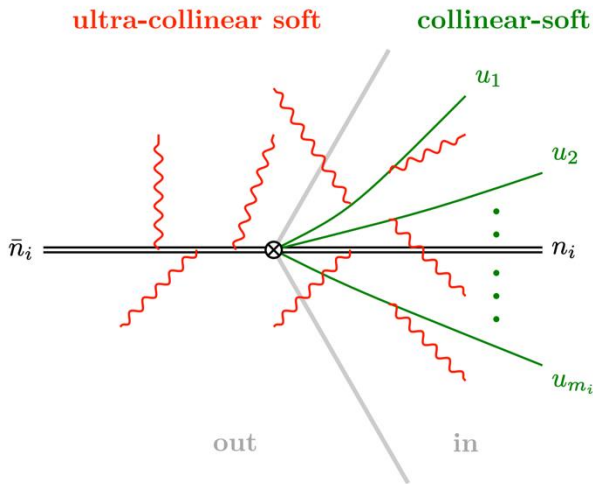
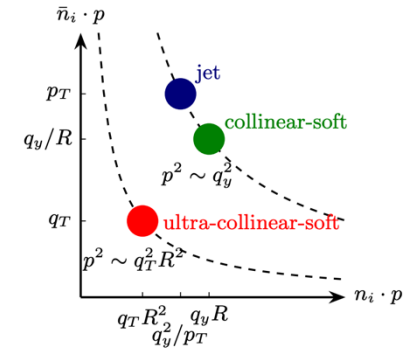
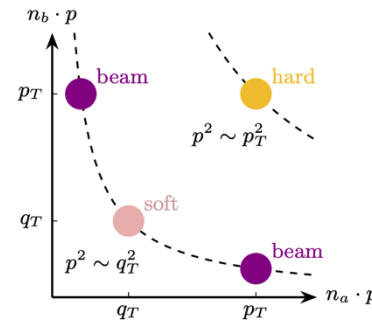
- Total transverse momentum q_T

- Suitable for **LO non-planar process**, such as $pp \rightarrow V/H + 2 \text{ jets}, pp \rightarrow 3 \text{ jets}, \text{ etc.}$

- Soft function is new: Outside the jet (b_T), Inside the jet (b_x)

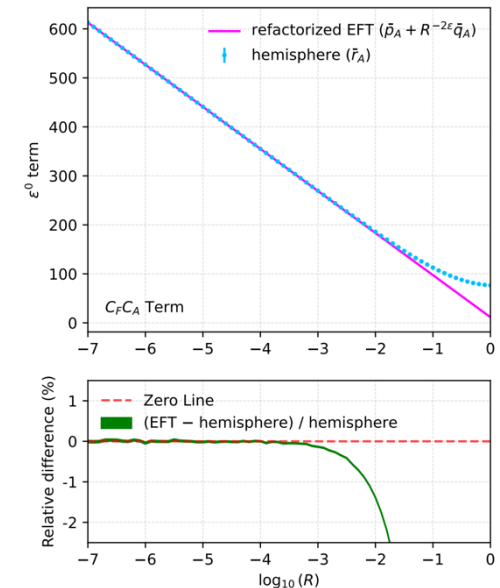
➤ In small R limit, soft function refactorizes into simpler global and (**ultra**)-**collinear-soft** contributions

- **soft:** $p_s^\mu \sim q_T (1, 1, 1)$,
- **n_i -collinear-soft:** $p_{cs}^\mu \sim |q_y|/R (R^2, 1, R)_{n_i \bar{n}_i}$,
- **n_i -ultra-collinear-soft:** $p_{ucs}^\mu \sim q_T (R^2, 1, R)_{n_i \bar{n}_i}$,



Two-loop collinear-soft function

$$\begin{aligned}
 S_i^{\text{hemi},(2)}(b_\perp, b^+) &= \left(\frac{\mu|b_\perp|}{b_0}\right)^{4\epsilon} \left(\frac{\nu|b_\perp|R_i}{b_0}\right)^\eta \omega^2 \left[\omega^2 \left(\frac{\nu|b_\perp|R_i}{b_0}\right)^\eta C_i^2 \frac{h_{\text{in}}^2}{2} \right. \\
 &\quad \left. + C_i C_A (h_A + v_A^{\text{in}}) + C_i n_f T_F h_f \right] \\
 &+ \left(\frac{i b^+ \mu}{b_0 R_i}\right)^{4\epsilon} \left[C_i^2 \frac{s_{\text{out}}^2}{2} + C_i C_A (g_A + v_A^{\text{out}}) + C_i n_f T_F g_f \right] \\
 &+ \left(\frac{\mu|b_\perp|}{b_0}\right)^{2\epsilon} \left(\frac{i b^+ \mu}{b_0 R_i}\right)^{2\epsilon} (C_i C_A r_A(r) + C_i n_f T_F r_f(r)) \\
 &+ \left(\frac{\mu|b_\perp|}{b_0}\right)^{2\epsilon} \left(\frac{i b^+ \mu}{b_0 R_i}\right)^{2\epsilon} \left(\frac{\nu|b_\perp|R_i}{b_0}\right)^\eta \omega^2 C_i^2 h_{\text{in}} s_{\text{out}}.
 \end{aligned}$$



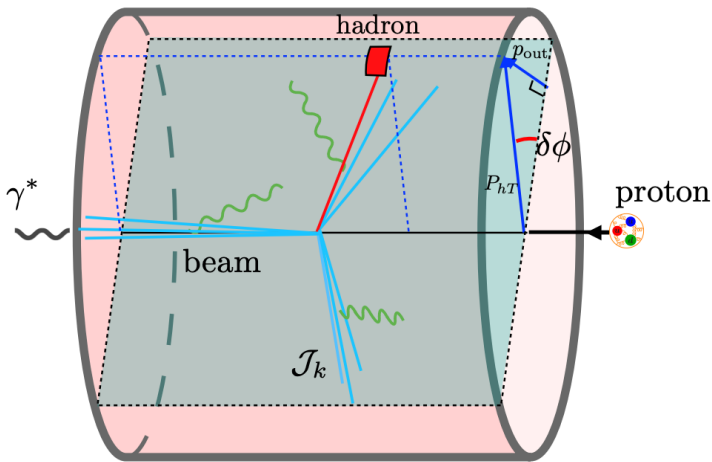
Multi-Wilson lines: Becher, Neubert, Rothen, DYS '16 PRL

NNLO QCD corrections for P_{hT} distribution

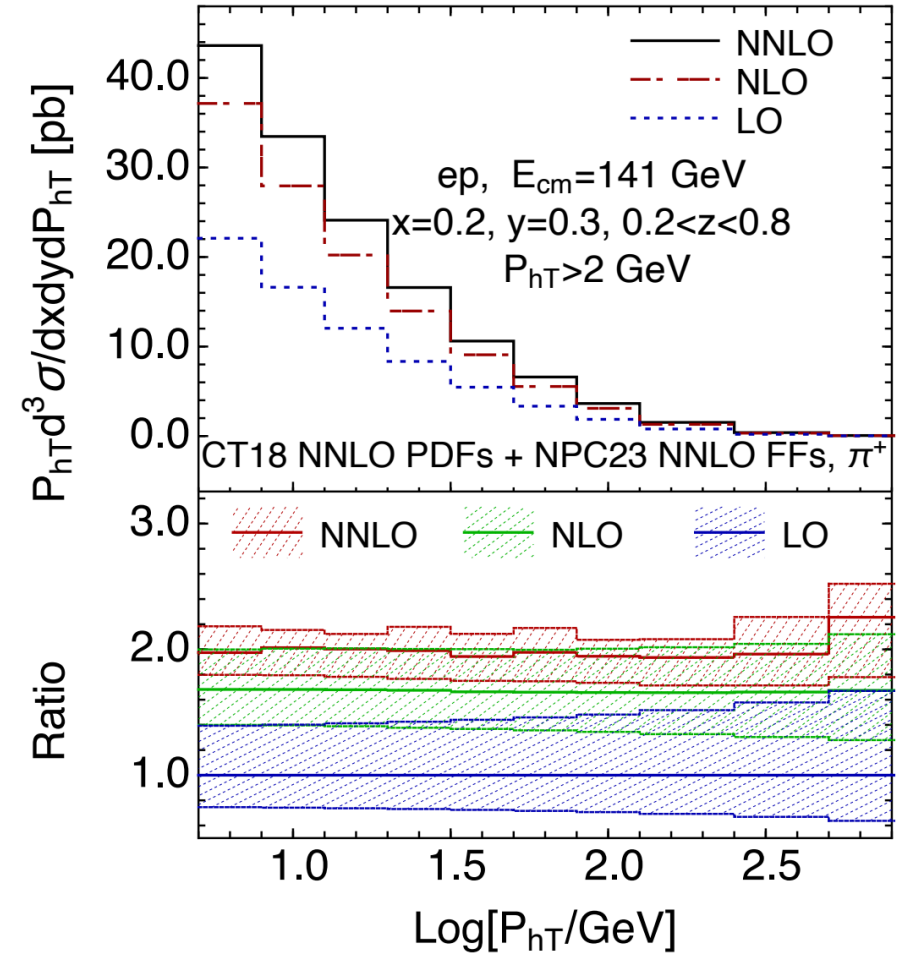
Dong(董亮), Fang(方申), Gao, DYS, Li, Zhu (朱雨蛟) 2602.22972

- We present the **first** complete calculation of hadron production in DIS at finite transverse momentum to NNLO

$$\frac{d\sigma}{d\mathcal{O}} = \int_0^{\delta\phi^{\text{cut}}} d\delta\phi \frac{d\sigma}{d\delta\phi d\mathcal{O}} + \int_{\delta\phi^{\text{cut}}}^{\delta\phi^{\text{max}}} d\delta\phi \frac{d\sigma}{d\delta\phi d\mathcal{O}}$$

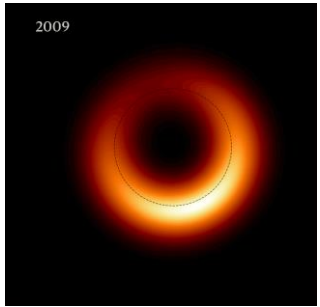


$$\begin{aligned} \frac{d\sigma_{\text{LP}}}{dx dy dz d^2\vec{P}_{hT} dp_{\text{out}}} &= \int \frac{db}{2\pi} e^{ip_{\text{out}}b/\zeta} \\ &\times \sum_{ijk} \int d\xi H_{ei \rightarrow ejk}(Q, \xi, \zeta) \\ &\times \mathcal{B}_{i/p}(\xi, b) \mathcal{D}_{h/j}(\zeta, b) \mathcal{J}_k(b) S_{ijk}(b) \end{aligned}$$

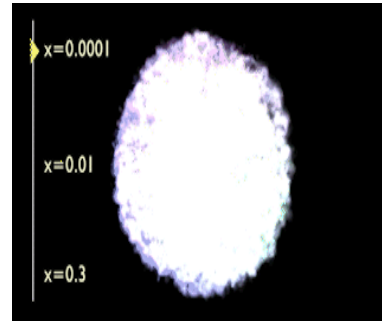
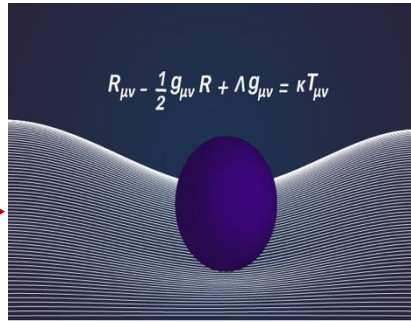


- **First N3LO SIDIS** Dong(董亮), Fang(方申), Gao, DYS, Li, Zhu (朱雨蛟) Zhu 2603.29673

➤ **核子成像** Also see J. Zhou's talk

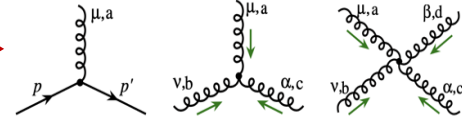


黑洞成像将深化人类对引力的理解



核子成像将深化人类对强力的理解

$$-\frac{1}{4}G_{\mu\nu,a}^2[A] + \sum_f \bar{\psi}_f (iD_\mu[A]\gamma^\mu - m_f) \psi_f$$

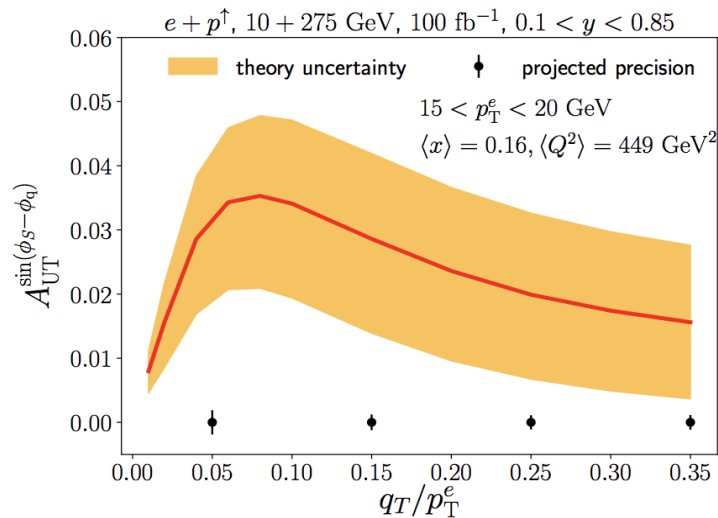
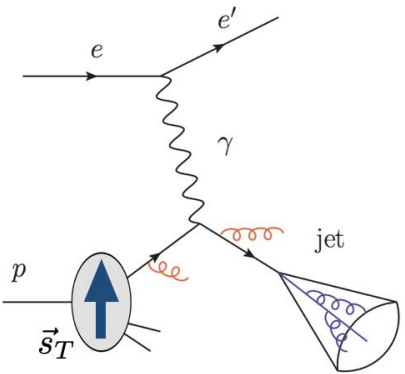


➤ **喷注及其精细结构为抽取核子分布函数提供了新的探针**

➤ **与传统SIDIS观测量相比可以有效分离出末态的非微扰信息**

$$p(p_A, S_A) + e(p_B, \lambda_e) \rightarrow [\text{jet}(p_C) h(z_h, \mathbf{j}_\perp)] + e(p_D) + X$$

PDF	\tilde{f}_1	$\tilde{f}_{1T}^{\perp(1)}$	\tilde{g}_{1L}	$\tilde{g}_{1T}^{(1)}$
JFF				
\mathcal{D}_1	1	$\sin(\phi_q - \phi_{S_A})$	1	$\cos(\phi_q - \phi_{S_A})$
\mathcal{D}_{1T}^\perp	$\sin(\hat{\phi}_h - \hat{\phi}_{S_h})$	$\sin(\hat{\phi}_h - \hat{\phi}_{S_h}) \sin(\phi_q - \phi_{S_A})$	$\sin(\hat{\phi}_h - \hat{\phi}_{S_h})$	$\sin(\hat{\phi}_h - \hat{\phi}_{S_h}) \cos(\phi_q - \phi_{S_A})$
\mathcal{G}_{1L}	1	$\sin(\phi_q - \phi_{S_A})$	1	$\cos(\phi_q - \phi_{S_A})$
\mathcal{G}_{1T}	$\cos(\hat{\phi}_h - \hat{\phi}_{S_h})$	$\cos(\hat{\phi}_h - \hat{\phi}_{S_h}) \sin(\phi_{S_A} - \phi_q)$	$\cos(\hat{\phi}_h - \hat{\phi}_{S_h})$	$\cos(\hat{\phi}_h - \hat{\phi}_{S_h}) \cos(\phi_q - \phi_{S_A})$
PDF	$\tilde{h}_1^{\perp(1)}$	$\tilde{h}_{1L}^{\perp(1)}$	\tilde{h}_1	$\tilde{h}_{1T}^{\perp(2)}$
JFF				
\mathcal{H}_1^\perp	$\cos(\hat{\phi}_h - \phi_q)$	$\sin(\phi_q - \hat{\phi}_h)$	$\sin(\phi_{S_A} - \hat{\phi}_h)$	$\sin(2\phi_q - \hat{\phi}_h - \phi_{S_A})$
\mathcal{H}_{1L}^\perp	$\sin(\hat{\phi}_h - \phi_q)$	$\cos(\phi_q - \hat{\phi}_h)$	$\cos(\phi_{S_A} - \hat{\phi}_h)$	$\cos(2\phi_q - \hat{\phi}_h - \phi_{S_A})$
\mathcal{H}_1	$\sin(\hat{\phi}_{S_h} - \phi_q)$	$\cos(\phi_q - \hat{\phi}_{S_h})$	$\cos(\phi_{S_A} - \hat{\phi}_{S_h})$	$\cos(2\phi_q - \hat{\phi}_{S_h} - \phi_{S_A})$
\mathcal{H}_{1T}^\perp	$\sin(2\hat{\phi}_h - \hat{\phi}_{S_h} - \phi_q)$	$\cos(2\hat{\phi}_h - \hat{\phi}_{S_h} - \phi_q)$	$\cos(2\hat{\phi}_h - \hat{\phi}_{S_h} - \phi_{S_A})$	$\cos(2\hat{\phi}_h - \hat{\phi}_{S_h} + \phi_{S_A} - 2\phi_q)$



Kang, Lee, DYS, Zhao '21

Liu, Ringer, Vogelsang, Yuan 19
Arratia, Kang, Prokudin, Ringer 19

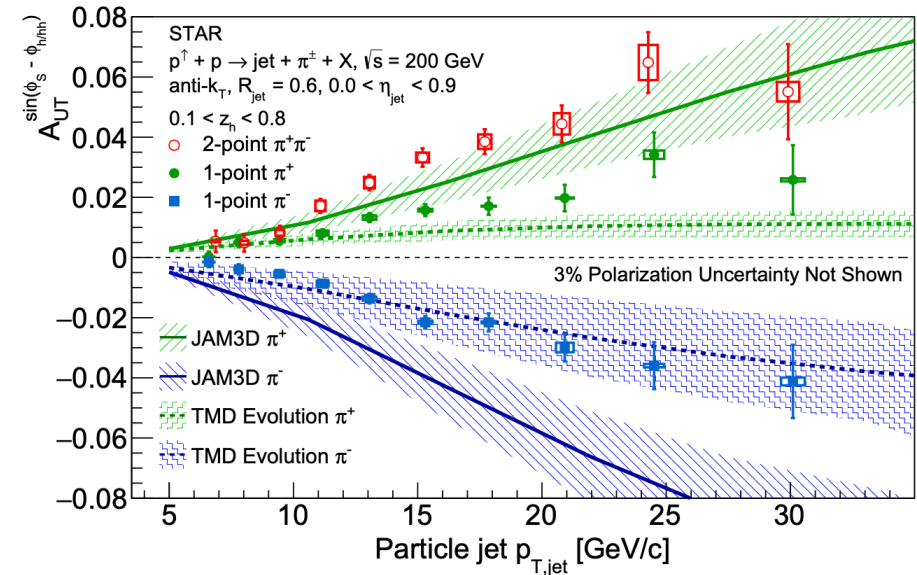
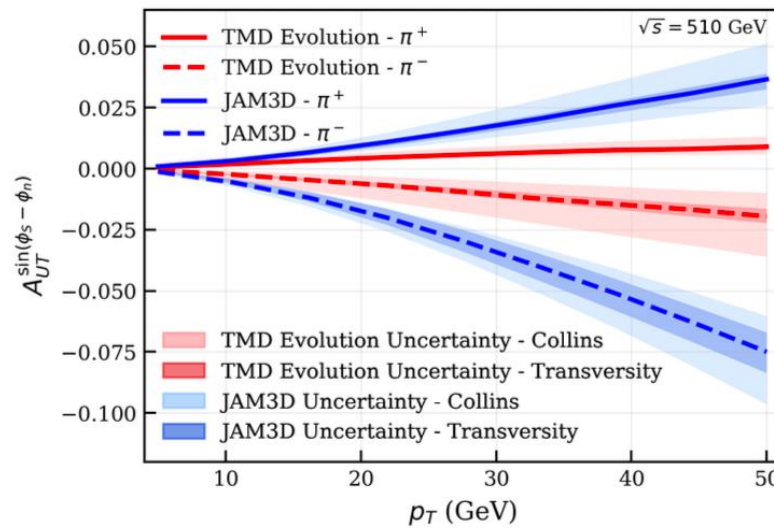
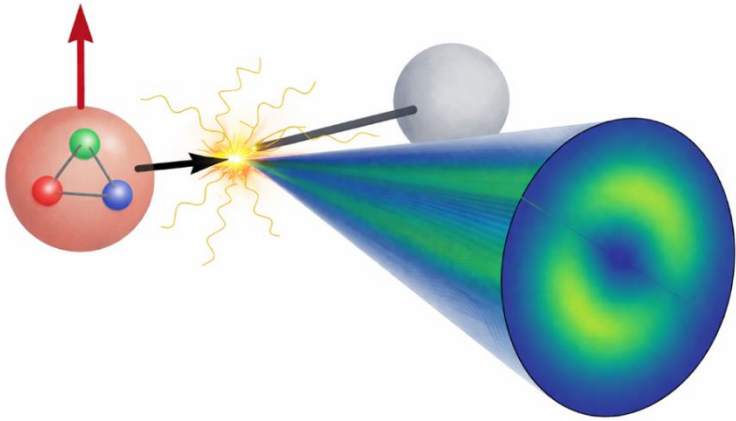
Gao(高梅森), Kang, Li(李婉辰), DYS PRL 136 (2026) 151902

- Transversity distribution h_1^q is a fundamental, yet less known, parton distribution.
- Its chiral-odd nature prevents access via inclusive DIS
- Nucleon tensor charge δq is essential for BSM search, Lattice QCD benchmarks
- We introduce the one-point energy correlator (OPEC) to probe the nucleon transversity

$$\Delta^q(z, \hat{n}) = \sum_X \sum_{i \in J} \langle \Omega | \bar{\chi}_n \delta_{Q, P_n} \delta^{(2)}(\hat{n} - \hat{n}_i) | JX \rangle \frac{E_i}{E_J} \langle JX | \chi | \Omega \rangle$$

Also see H. X. Zhu & X. H. Liu's talks

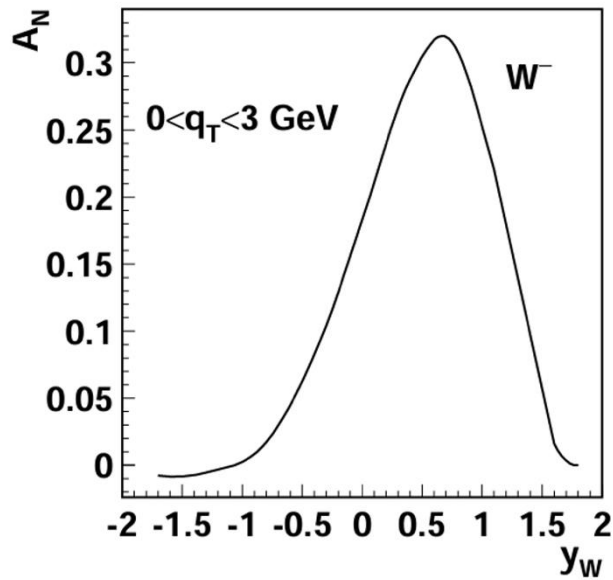
RHIC-STAR 2604.15543





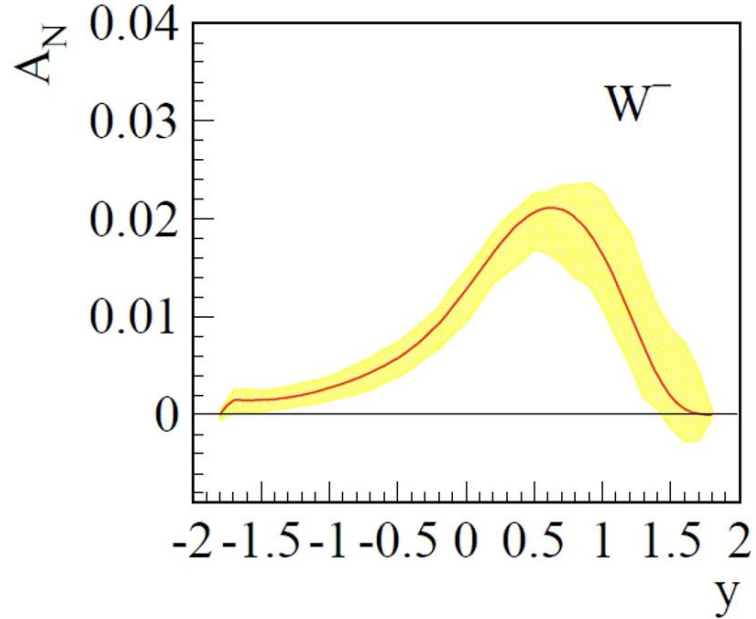
Spin asymmetry and QCD evolution

Without evolution



Kang & Qiu '09

With evolution



Echevarria, Idilbi, Kang, Vitev '14

- Soft and collinear radiation destroy superposition
- Soft radiation decohere the directions of the momenta and the gauge representations of the particles
- Collinear radiation decohere the energies and the spin

Carney, Chaurette, Neuenfeld, Semenov '17

Neill, Waalewijn '19

S.J. Lin, M.J. Liu, DYS, S.Y. Wei '25

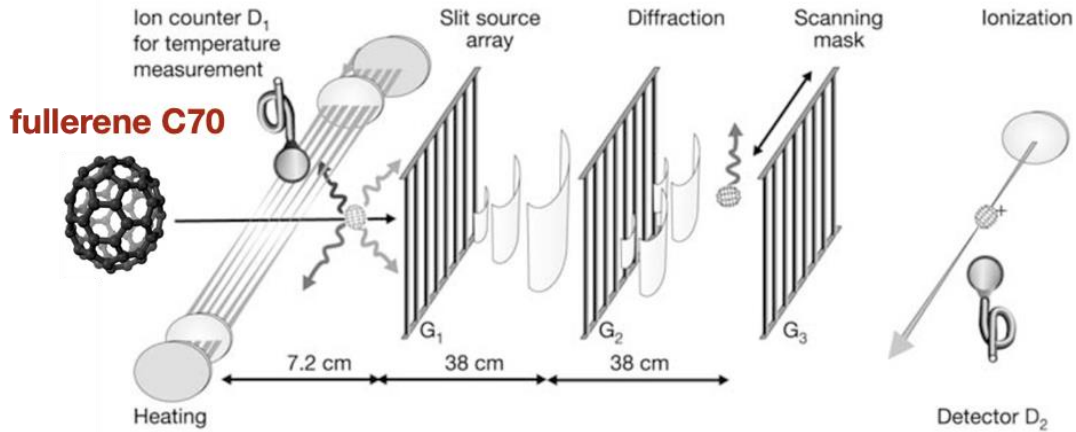
J.Y. Gu, S.J. Lin, DYS, S.X. Yang, L.T Wang '25

... ..

Evolution reduces the asymmetry

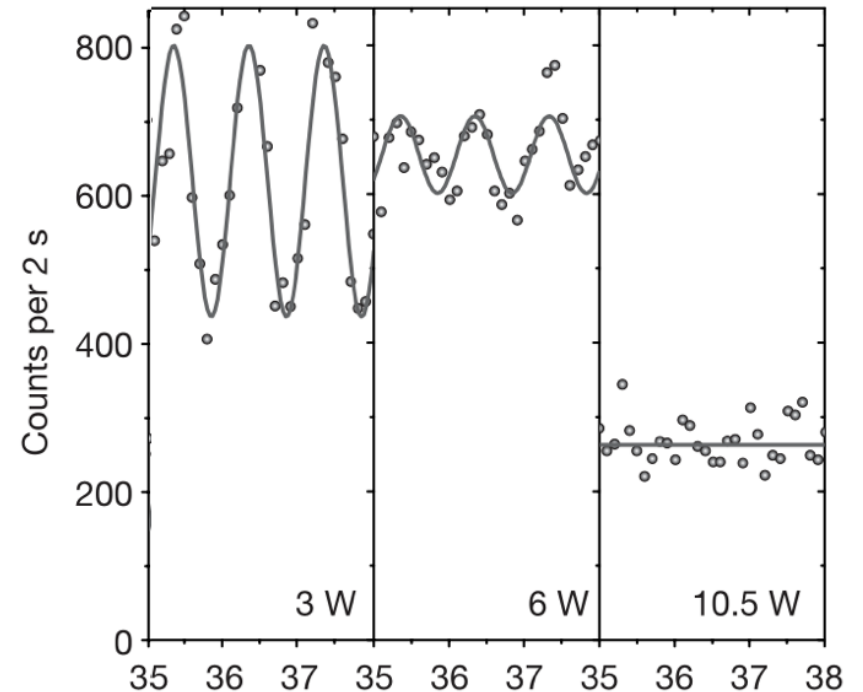
热光子辐射引起的退相干

Hackermuller, Hornberger, Brezger, Zeilinger, Arndt Nature '04



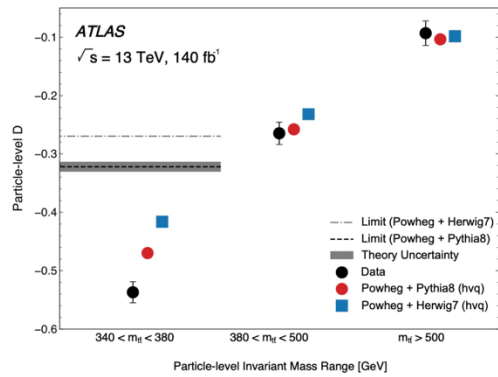
- 塔尔博特-劳厄干涉仪
- 在进入干涉仪之前，C70分子会穿过一束激光以提高温度
- 温度调节范围从低于1000K（低温）至3000K（高温）

- 低温：
 - 分子辐射可忽略不计
 - 结果：观测到清晰的干涉条纹
- 高温：
 - 分子发射携带路径信息的热光子
 - 结果：干涉条纹可见度降至0%

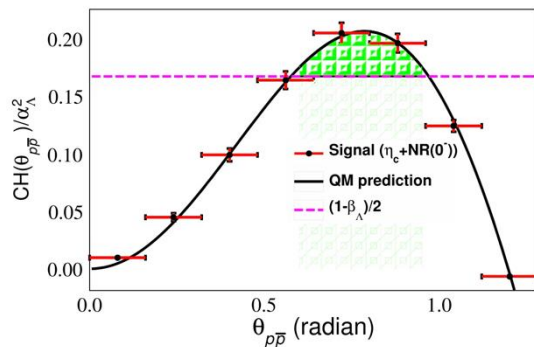


辐射可以导致退相干

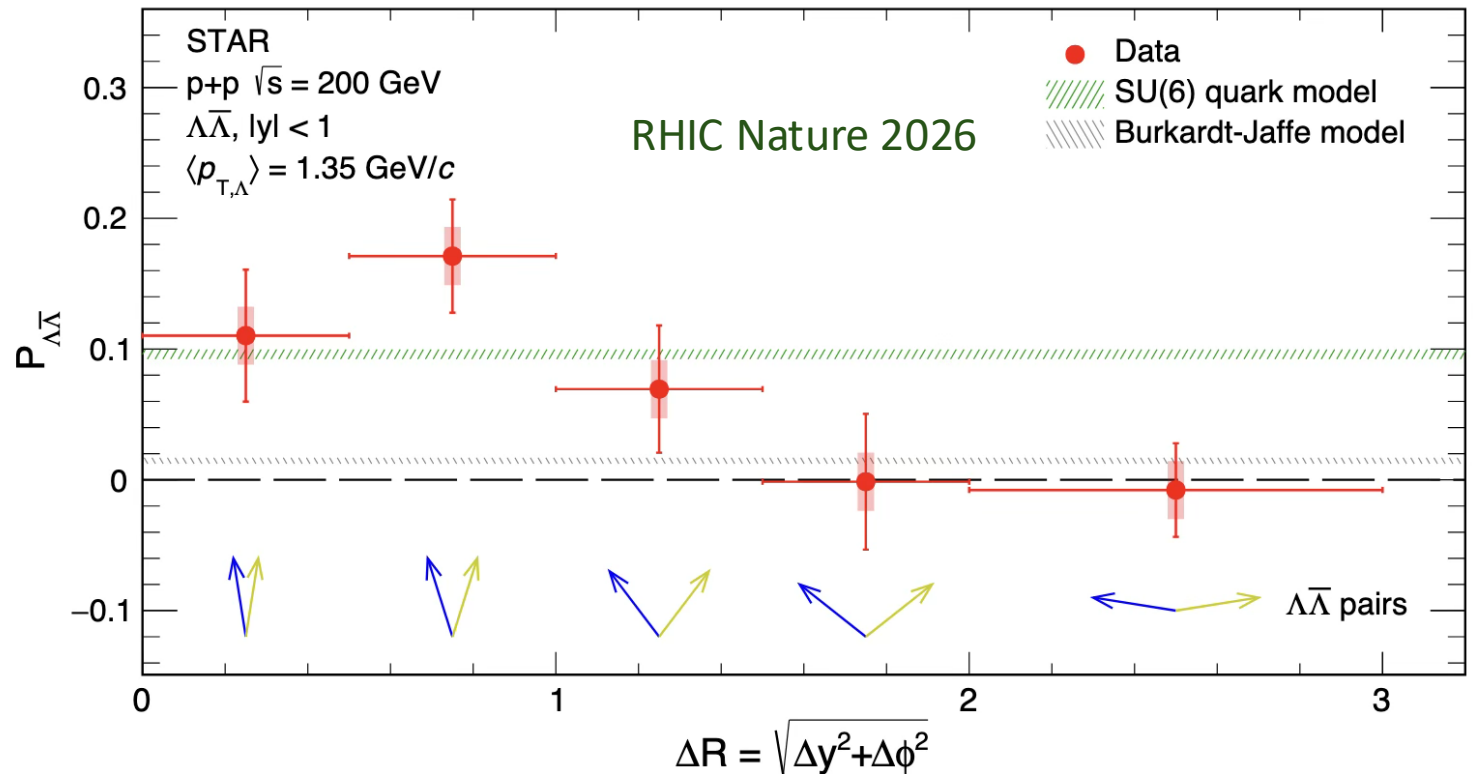
The study of quantum information in high-energy collider physics is rapidly transitioning from a theoretical curiosity to an experimental reality.



ATLAS Nature 2024



“Measuring spin correlation between quarks during QCD confinement”



BESIII Nature Communications 2025

Effective field theory for decoherence

S.J. Lin(林士佳), M.J. Liu(刘铭钧), DYS, S.Y. Wei 25 JHEP; J.Y. Gu, S.J. Lin(林士佳), DYS, S.X. Yang(杨斯翔), L.T Wang 25

- **A systematic framework calculating decoherence in high energy collisions by combing SCET with the formalism of open quantum systems.**
- **The renormalization group evolution constitutes a quantum channel, where the RG flow parameter, rather than time, drives a Markovian loss of quantum coherence.**

➤ **E.g. Applications in QED processes** $e^+e^- \rightarrow l^+l^- + n\gamma$

➤ **Lindblad equation and jump operators**

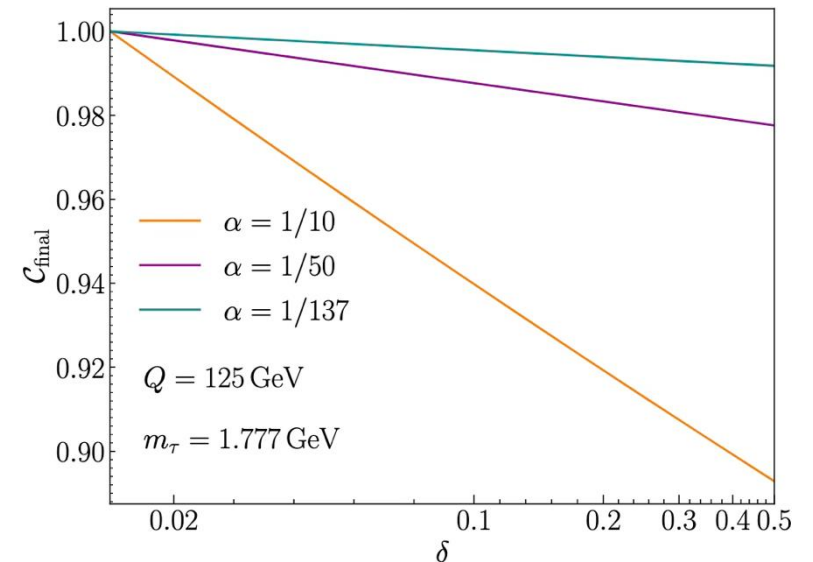
$$\frac{d\hat{\rho}_{\text{eff}}}{dt} = -\frac{\alpha}{2\pi}\hat{\rho}_{\text{eff}} + \frac{\alpha}{4\pi} [(\hat{\sigma}_3 \otimes \mathbb{I}) \hat{\rho}_{\text{eff}} (\hat{\sigma}_3 \otimes \mathbb{I}) + (\mathbb{I} \otimes \hat{\sigma}_3) \hat{\rho}_{\text{eff}} (\mathbb{I} \otimes \hat{\sigma}_3)]$$

$$t \equiv \log(Q\delta/\mu). \quad \hat{L}_1 = \sqrt{\alpha/4\pi} \hat{\sigma}_3 \otimes \mathbb{I} \quad \hat{L}_2 = \sqrt{\alpha/4\pi} \mathbb{I} \otimes \hat{\sigma}_3$$

➤ **Each “jump” corresponds to an unresolved collinear photon emission from either of the fermion legs, which induces a stochastic **phase-flip** quantum channel**

➤ **Reduction of the concurrence**

$$\mathcal{C}(t) = \mathcal{C}(0)e^{-\frac{\alpha}{\pi}t}$$

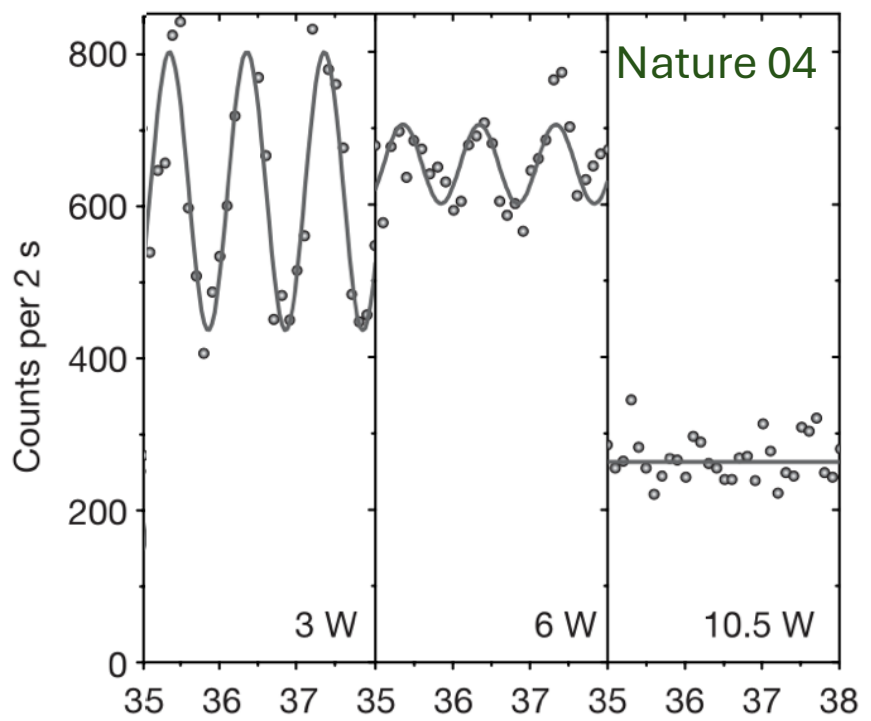


**decoherence = RG flow
anomalous dimensions in RGE
determine the information loss**



Nucleon tomography

- TMDs encode the quantum correlations and the motion of quarks and gluons
- However, at high energies, initial state radiation and environmental noise, causing distortion of these signals.
- We introduce a time scale τ_0 into a tomographic cross-section
- We demonstrate the dependence of the cross-section on the time scale τ_0 and its asymmetry

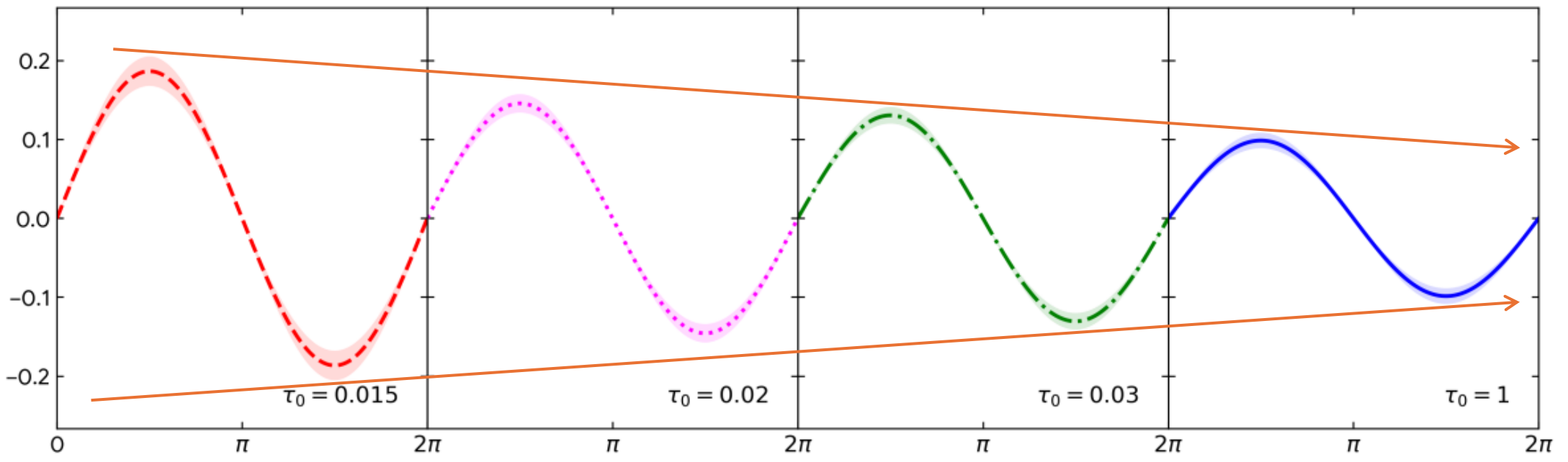


木硕), DYS, Zhou PRL136 (2026)021901

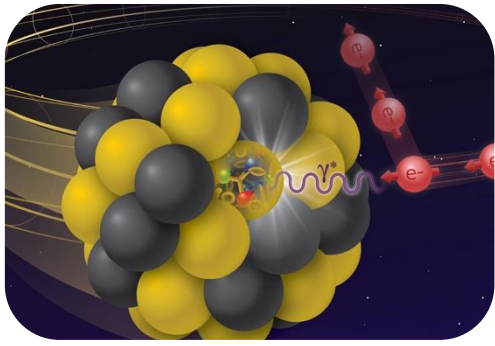
$$\frac{\vec{s}_T}{S} = F_{UU} + \sin(\phi_s - \phi_q) F_{UT}^{\sin(\phi_s - \phi_q)}$$

↓
↓

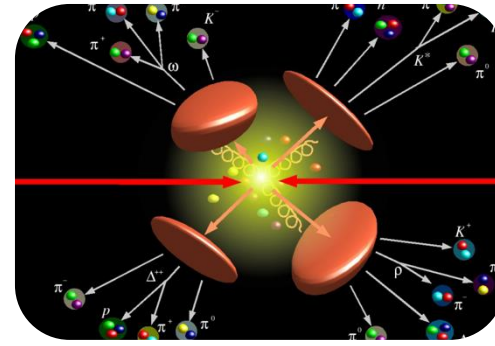
$$f_q(x, k_T) \quad \frac{1}{M} \epsilon_{\alpha\beta} s_T^\alpha k_T^\beta f_{1T}^{\perp q}(x, k_T)$$



QCD对撞机物理研究从“总截面”的宏观成功，迈向“微分与极化”的微观精密时代



核子结构与自旋物理



喷注物理与精确计算



量子信息的视角

谢谢
