



中国物理学会高能物理分会
第十二届全国会员代表大会暨学术年会

Baryon–baryon and meson–baryon interactions in Lorentz invariant ChEFT

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In collaboration with

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Ulf-G. Meißner (Bonn)



UNIVERSITÄT  BONN

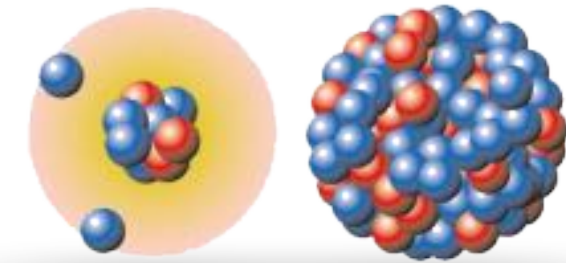
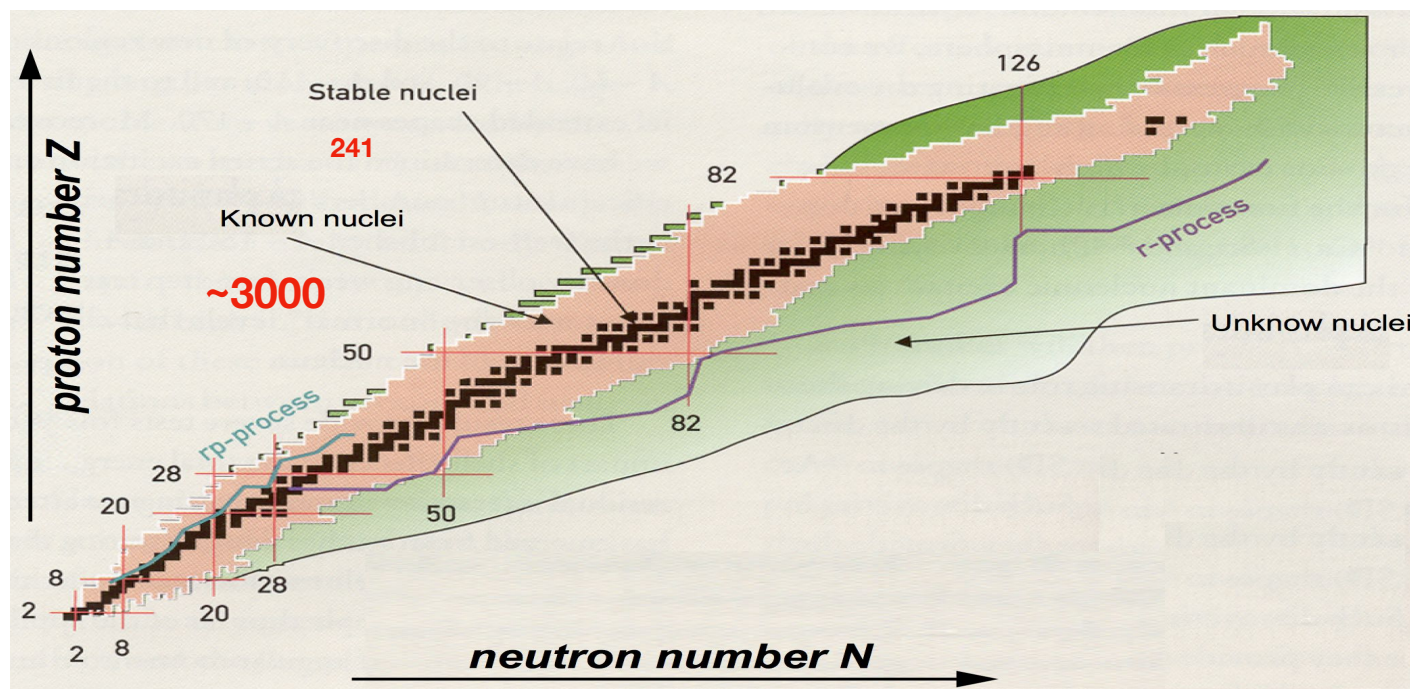
OUTLINE

- ▣ Introduction
- Theoretical framework
- Results and discussion
- Summary and outlook

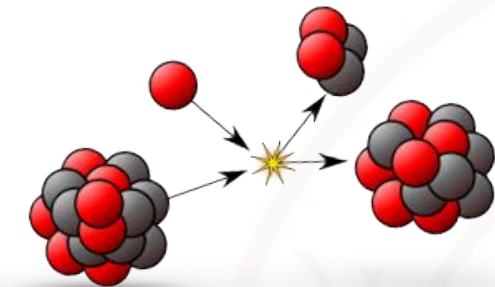


(Hyper)Nuclear Chart

□ Nuclear Chart (2D — N, Z)



Nuclear structure



Nuclear reaction

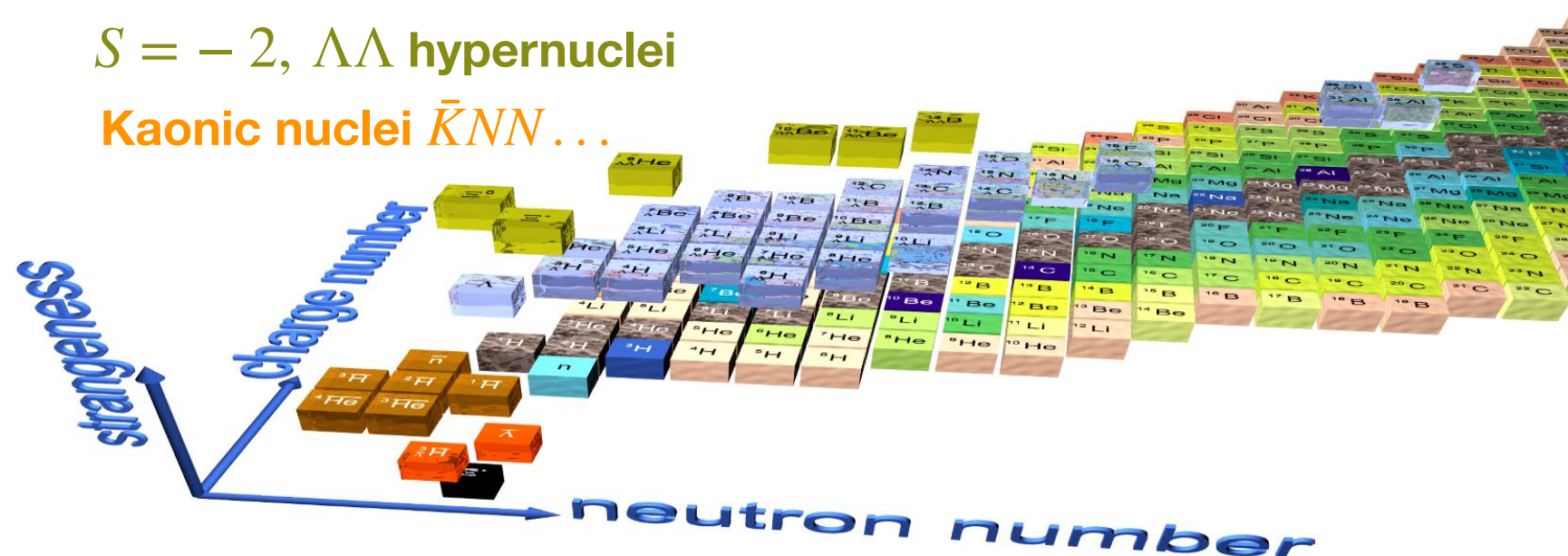
- Detailed understanding of the strong **nuclear force (NN)** is essential!

□ Hypernuclear Chart (3D — N, Z, Strangeness)

$S = -1$, Λ , Ξ hypernuclei

$S = -2$, $\Lambda\Lambda$ hypernuclei

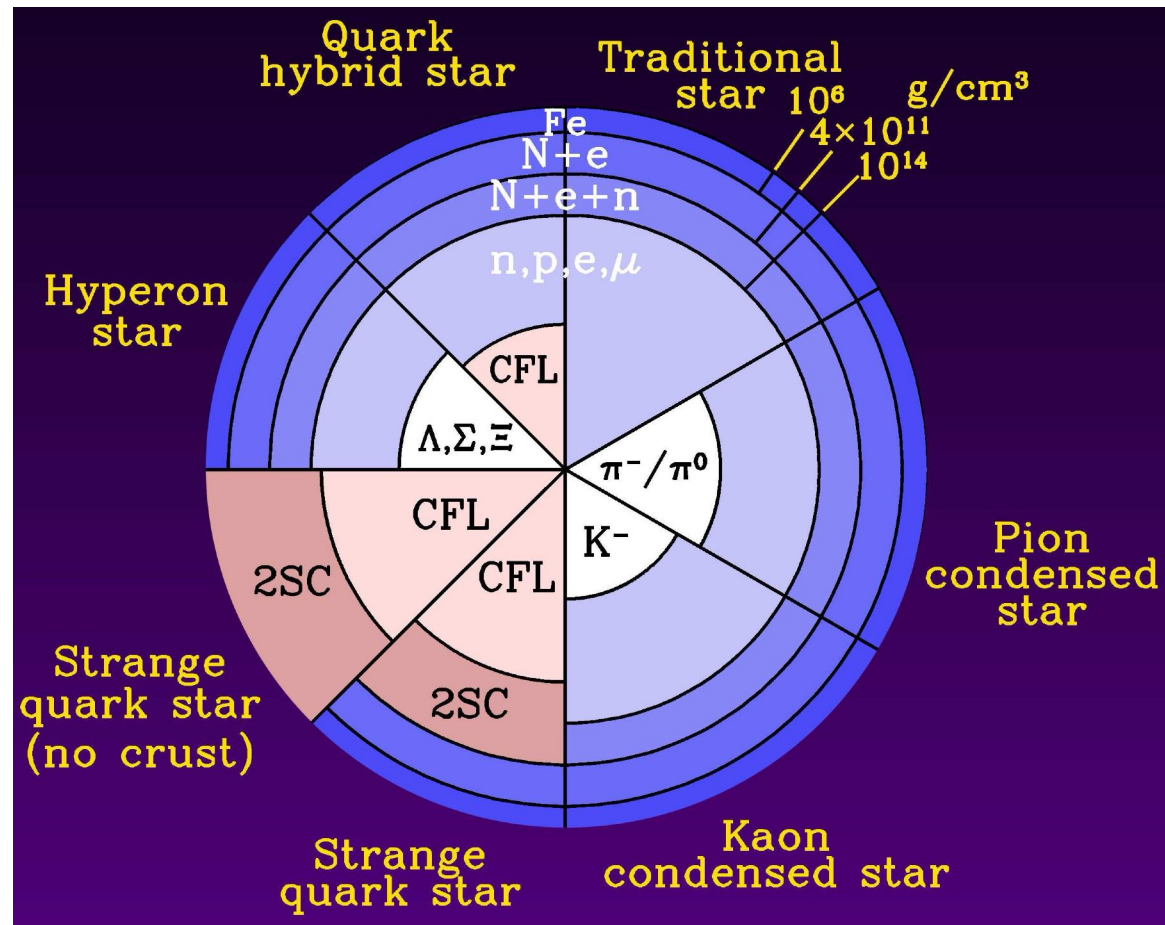
Kaonic nuclei $\bar{K}NN$...



- Hyperon-nucleon (**YN**)
- Hyperon-hyperon (**YY**)
- Antikaon-nucleon (**$\bar{K}N$**) interactions

Neutron star

Unknown interior

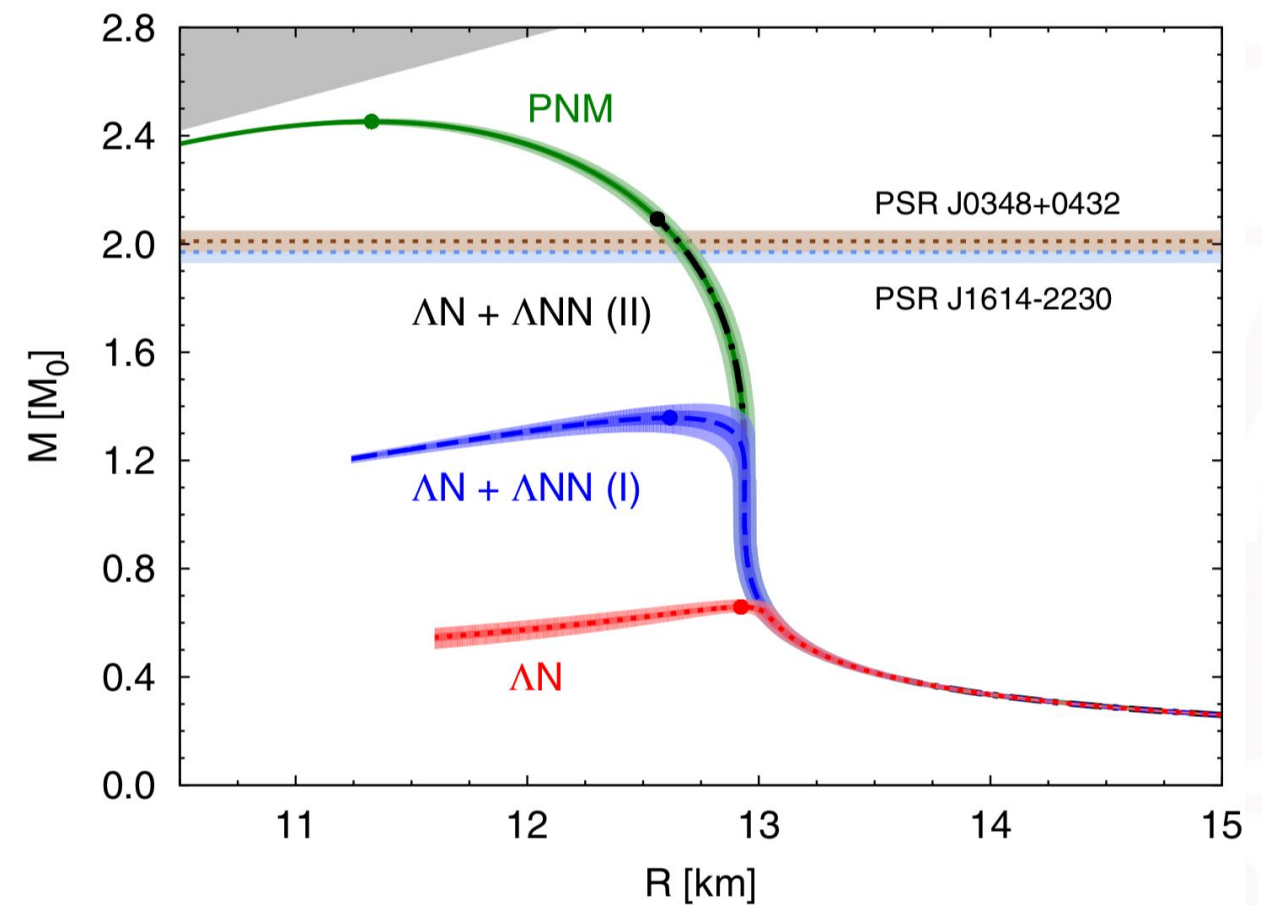


• Inner Core:

- ✓ Hyperon star ?
- ✓ Kaon-condensed star ?
- ✓ Strange quark star ?

Equation of state

• Hyperon puzzle

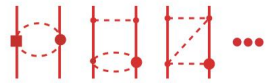


D.Lonardon et al., PRL 114 (2015) 092301

Precise knowledge of YN and $\bar{K}N$ interactions needed

BB and MB scatterings in ChEFT

- BB interactions extensively studied in **non-relativistic** framework

	2NF	3NF	4NF
LO (Q^0)	 <i>S. Weinberg, PLB 1990, NPB1990</i>	—	—
NLO (Q^2)	 <i>U. van Kolck et al., PLB1992, PRL1994 N. Kaiser et al., NPA1997</i>	—	—
N ² LO (Q^3)	 <i>U. van Kolck et al., PRC1994 E. Epelbaum et al., NPA1998, 2000</i>	 <i>U. van Kolck et al., PRC1994</i>	—
N ³ LO (Q^4)	 <i>R. Machleidt et al., PRC2003 E. Epelbaum et al., NPA2005</i>	 <i>S. Ishikwas et al., PRC2007 V. Bernard et al., PRC2007</i>	 <i>E. Epelbaum, PLB2006, EPJA2007</i>
N ⁴ LO (Q^5)	 <i>R. Machleidt et al., PRC2015 E. Epelbaum et al., PRL2015</i>	 <i>H. Krebs et al., PRC2012,2013</i>	

	Two-baryon force	Three-baryon force	Four-baryon force
LO (Q^0)	 <i>S=-1,-2: H. Polinder et al., NPA2006,PLB2007 S=-3,-4: J. Haidenbauer & U.-G. Meißner PLB2010</i>	—	—
NLO (Q^2)	 <i>S=-1: J. Haidenbauer et al., NPA2013, EPJA2020 S=-2: J. Haidenbauer et al., NPA2016</i>	—	—
N ² LO (Q^3)	 <i>S=-1: J. Haidenbauer et al., arXiv:2301.00722</i>	 <i>ΔNN: S. Petschauer et al., PRC2016, NPA2017</i>	—

- Solve Lippmann-Schwinger equation

$$T(\mathbf{p}', \mathbf{p}) = V(\mathbf{p}', \mathbf{p}) + \int \frac{d^3k}{(2\pi)^3} V(\mathbf{p}', \mathbf{k}) \frac{m_N}{p^2 - k^2 + i\epsilon} T(\mathbf{k}, \mathbf{p})$$

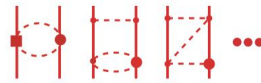
- **Renormalization issue is still under debate!**

Nuclear Forces for Precision Nuclear Physics: A Collection of Perspectives

Few-Body Syst (2022) 63:67

BB and MB scatterings in ChEFT

- BB interactions extensively studied in **non-relativistic** framework

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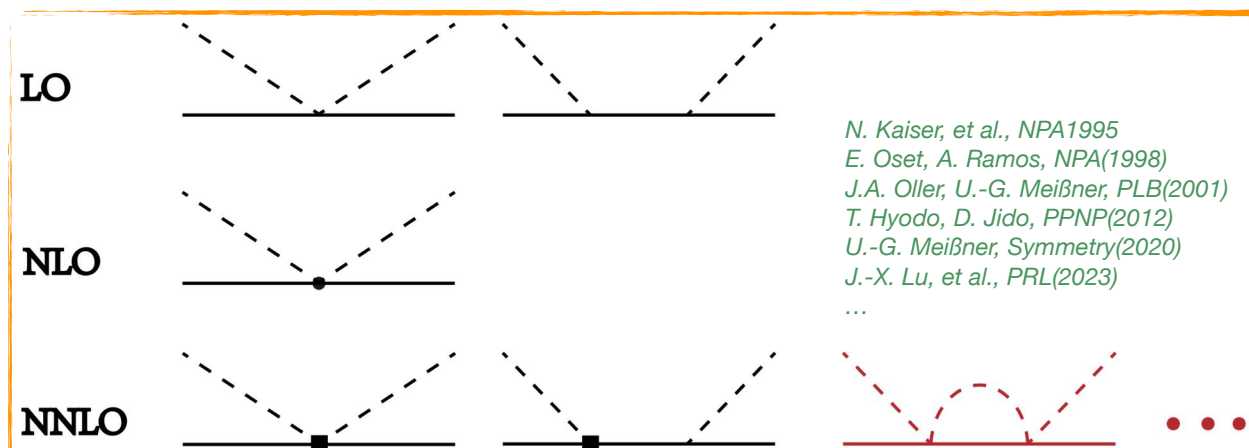
- Solve Lippmann-Schwinger equation

$$T(p', p) = V(p', p) + \int \frac{d^3k}{(2\pi)^3} V(p', k) \frac{m_N}{p^2 - k^2 + i\epsilon} T(k, p)$$

- **Renormalization issue is still under debate!**

Nuclear Forces for Precision Nuclear Physics: A Collection of Perspectives
Few-Body Syst (2022) 63:67

- MB interactions and **resonances** (usually) studied in **covariant** framework



- Solve Bethe-Salpeter equation

$$T(p', p) = V(p', p) + i \int \frac{d^4k}{(2\pi)^4} V(p', k) G(k) T(k, p)$$

- On-shell factorization
- Model dependence: Finite cutoff or subtraction constants

In this work

□ We tentatively proposed a unified framework for baryon-baryon and meson-baryon interactions within time-ordered perturbation theory (TOPT) using the Lorentz invariant chiral Lagrangians

- Formulate NN and YN interactions at leading order
V. Baru, E. Epelbaum, J. Gegelia, XLR*, Phys. Lett. B 798,134987 (2019)
XLR, E.Epelbaum, J.Gegelia, Phys. Rev. C 101,034001 (2020)
- Extend NN interaction to **next-to-next-to-leading order**
XLR, E. Epelbaum, J. Gegelia, Phys. Rev. C 106, 034001 (2022)
XLR, E. Epelbaum, J. Gegelia, Phys. Rev. C 113, 034002(2026)
- Investigate $\bar{K}N$, KN interactions and $\Lambda(1405)$ at LO/NLO
XLR, E. Epelbaum, J. Gegelia and U.-G. Meißner, Eur. Phys. J. C80 (2020) 406
XLR, E. Epelbaum, J. Gegelia and U.-G. Meißner, Eur. Phys. J. C81 (2021) 582
XLR, Phys. Lett. B 855 (2024) 138802
XLR, Phys. Rev. D 113 (2026) 114027
XLR, et al., in preparation (2026)

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Time-ordered perturbation theory

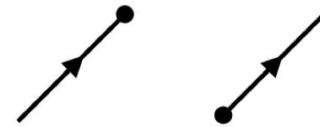
- Re-express the Feynman integral in a form that **makes the connection with on-mass-shell (off-energy shell) state explicit**.
- This form is called **TOPT (or old-fashioned perturbation theory)**

▶ External lines

Spin 0 boson (in, out)



Spin 1/2 fermion (in, out)



1

$u(\mathbf{p}), \quad \bar{u}(\mathbf{p}')$

▶ Internal lines

Spin 0 (anti-)boson



Spin 1/2 fermion



anti-fermion



$$\frac{1}{2\epsilon_q}$$

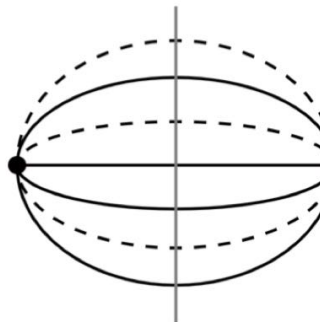
$$\epsilon_q \equiv \sqrt{\mathbf{q}^2 + M^2}$$

$$\frac{m}{\omega_p} \sum u(\mathbf{p})\bar{u}(\mathbf{p}) \quad \omega_p \equiv \sqrt{\mathbf{p}^2 + m^2}$$

$$\frac{m}{\omega_p} \sum u(\mathbf{p})\bar{u}(\mathbf{p}) - \gamma_0$$

▶ Intermediate state

A set of lines between two vertices



$$\frac{1}{E - \sum_i \omega_{p_i} - \sum_j \epsilon_{q_j} + i\epsilon}$$

▶ Interaction vertices: the standard Feynman rules

- Take care of zeroth components of integration momenta

✓ particle $p^0 \rightarrow \omega(p, m)$

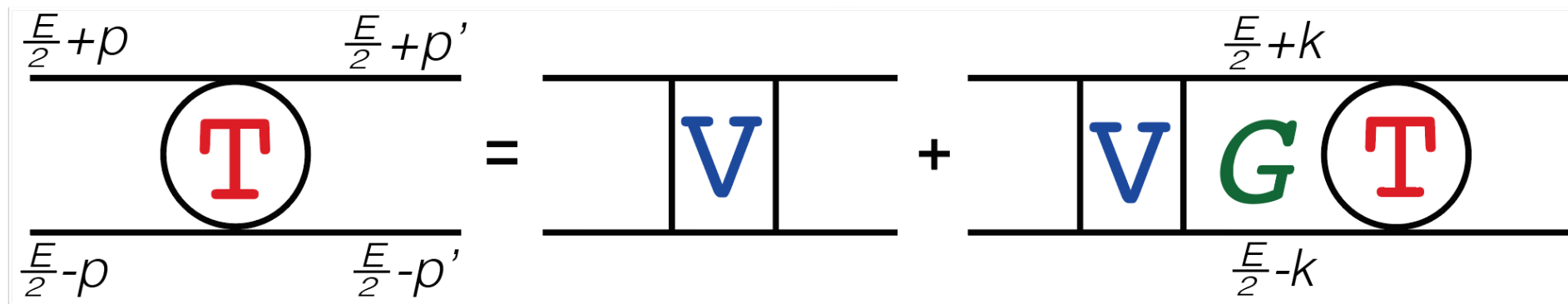
✓ antiparticle $p^0 \rightarrow -\omega(p, m)$

XLR, PoS(CD2021)007



Baryon-baryon scattering in TOPT

□ Baryon-baryon scattering amplitude



- **Potential V** : sum up the two-baryon **irreducible time-ordered diagrams**
 - ✓ Employ **the Weinberg power counting** to perturbatively calculate potential

$$V = \begin{pmatrix} V_{NN,NN} & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & V_{\Lambda N,\Lambda N} & V_{\Lambda N,\Sigma N} & 0 & 0 & 0 & 0 \\ 0 & V_{\Sigma N,\Lambda N} & V_{\Sigma N,\Sigma N} & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & V_{\Lambda\Lambda,\Lambda\Lambda} & V_{\Lambda\Lambda,\Xi N} & V_{\Lambda\Lambda,\Sigma\Sigma} & V_{\Lambda\Lambda,\Sigma\Lambda} \\ 0 & 0 & 0 & V_{\Xi N,\Lambda\Lambda} & V_{\Xi N,\Xi N} & V_{\Xi N,\Sigma\Sigma} & V_{\Xi N,\Sigma\Lambda} \\ 0 & 0 & 0 & V_{\Sigma\Sigma,\Lambda\Lambda} & V_{\Sigma\Sigma,\Xi N} & V_{\Sigma\Sigma,\Sigma\Sigma} & V_{\Sigma\Sigma,\Sigma\Lambda} \\ 0 & 0 & 0 & V_{\Sigma\Lambda,\Lambda\Lambda} & V_{\Sigma\Lambda,\Xi N} & V_{\Sigma\Lambda,\Sigma\Sigma} & V_{\Sigma\Lambda,\Sigma\Lambda} \end{pmatrix} \begin{matrix} \mathbf{S}=0 \\ \mathbf{S}=-1 \\ \mathbf{S}=-2 \end{matrix}$$

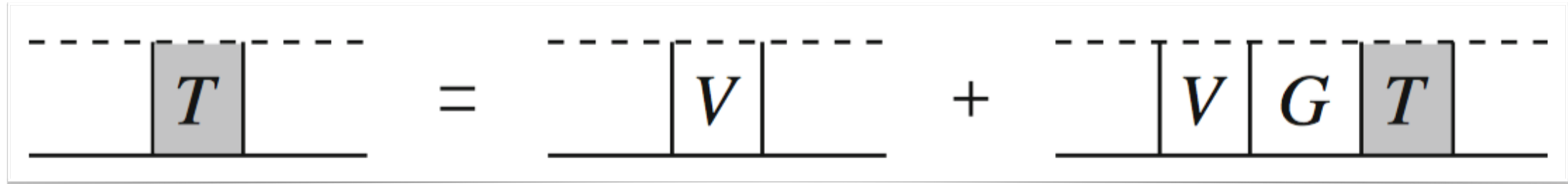
- **Two-baryon Green function (TOPT rules)**

$$G_{ij}^{BB}(E) = \frac{m_i m_j}{\omega(k, m_i) \omega(k, m_j)} \frac{1}{E - \omega(k, m_i) - \omega(k, m_j) + i\epsilon}$$

- ✓ This is the generalized Kadyshevsky propagator of NN scattering *V. Kadyshevsky, NPB (1968)*
- ✓ **SELF-CONSISTENTLY obtained** in TOPT

Meson-baryon scattering in TOPT

□ Meson-baryon scattering amplitude



- Interaction kernel/ **potential V**
 - ✓ **Define:** sum up the one-meson & one-baryon **irreducible time-ordered diagrams**
 - ✓ **Power counting:** Q/Λ_χ **systematic ordering of all graphs**
- Meson-baryon Green functions in TOPT:

$$G_{MB}(E) = \frac{m}{2\omega(k, M) \omega(k, m)} \frac{1}{E - \omega(k, M) - \omega(k, m) + i\epsilon}$$

- Coupled-channel integral equation for T-matrix

$$T_{M_j B_j, M_i B_i}(\mathbf{p}', \mathbf{p}; E) = V_{M_j B_j, M_i B_i}(\mathbf{p}', \mathbf{p}; E) + \sum_{MB} \int \frac{d^3 \mathbf{k}}{(2\pi)^3} V_{M_j B_j, MB}(\mathbf{p}', \mathbf{k}; E) G_{MB}(E) T_{MB, M_i B_i}(\mathbf{k}, \mathbf{p}; E)$$

OUTLINE

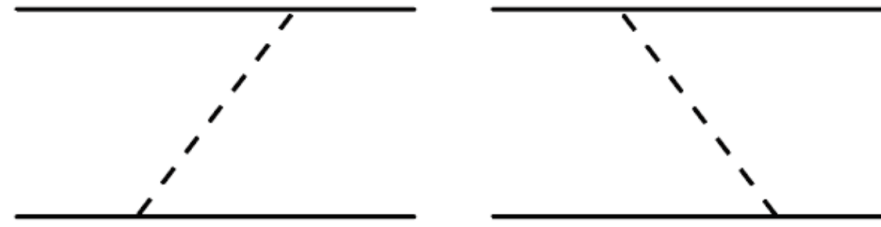
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NN and YN interactions at LO

□ LO potential

Weinberg power counting

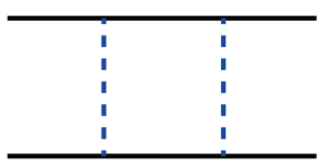


- NN interaction **XLR, et al., 1905.02116**

$$V_{LO,C} = (C_S + C_V) - (C_{AV} - 2C_T) \boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2$$

$$V_{\text{OPE}} = -\frac{g_A^2}{4f_\pi^2} \boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2 \frac{4m_N^2}{\omega(q, M_\pi) (m_N + \omega(p, m_N)) (m_N + \omega(p', m_N))} \\ \times \frac{\boldsymbol{\sigma}_1 \cdot \mathbf{q} \boldsymbol{\sigma}_2 \cdot \mathbf{q}}{\omega(p, m_N) + \omega(p', m_N) + \omega(q, M_\pi) - E - i\epsilon}$$

- **Milder UV behaviour than that of the non-relativistic OPEP**



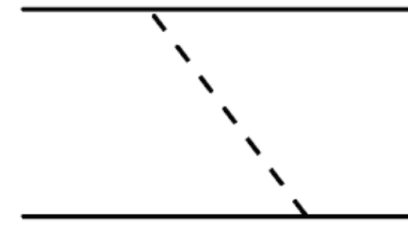
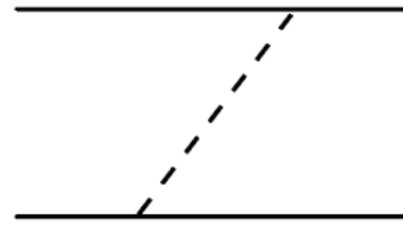
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NN and YN interactions at LO

LO potential

Weinberg power counting

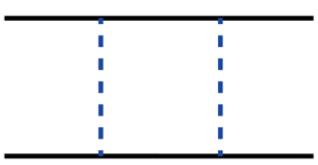


- NN interaction **XLR, et al., 1905.02116**

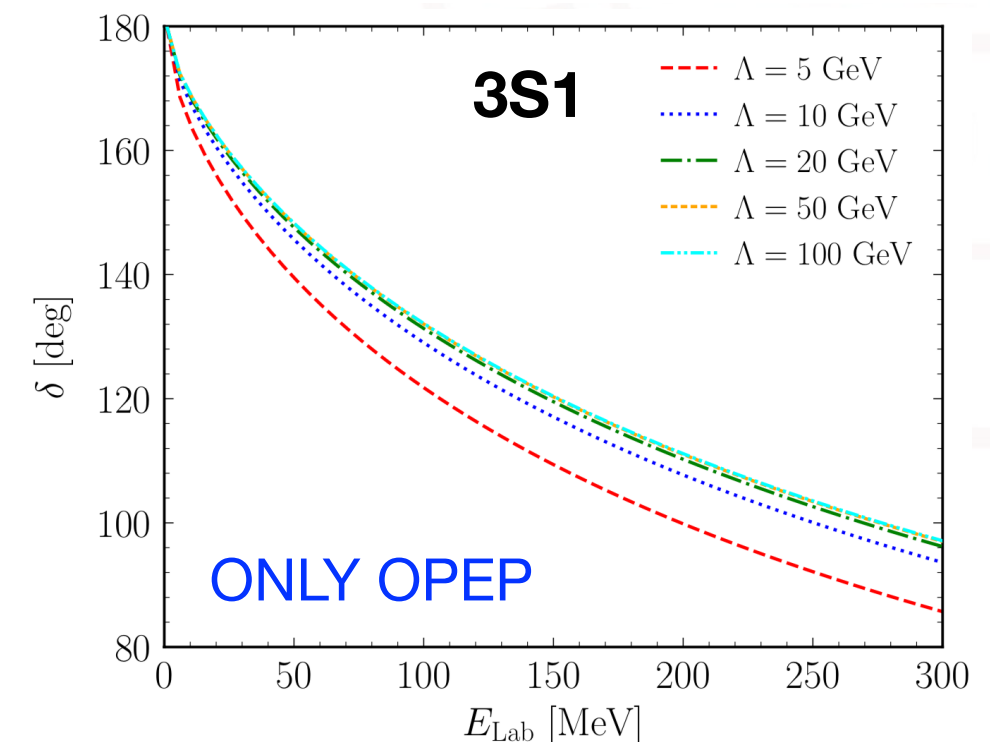
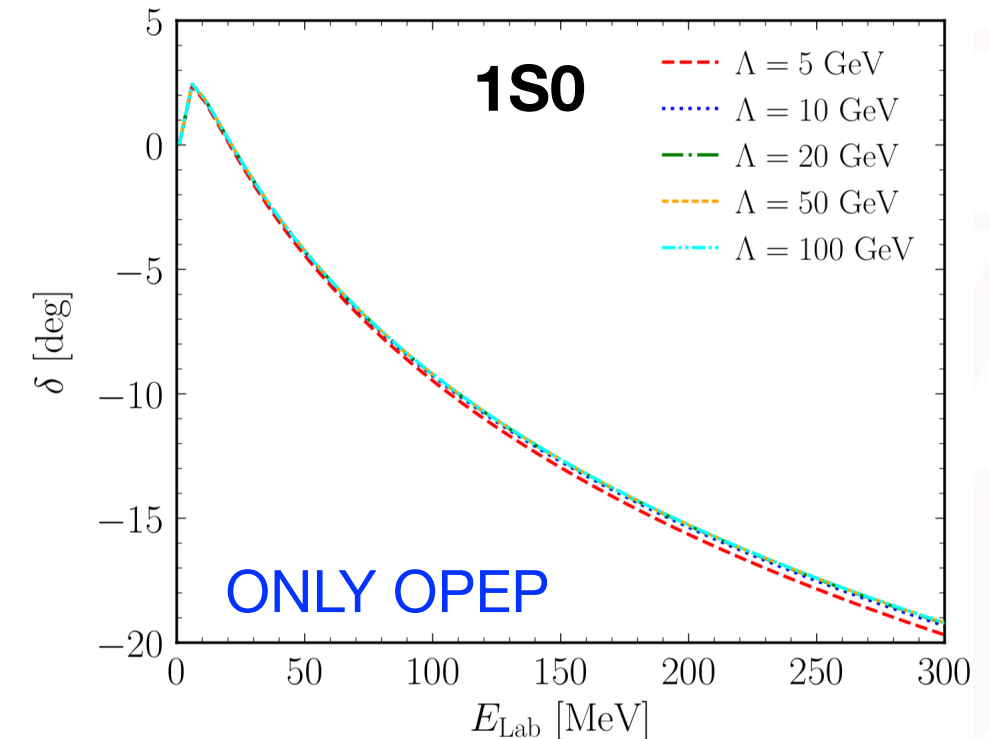
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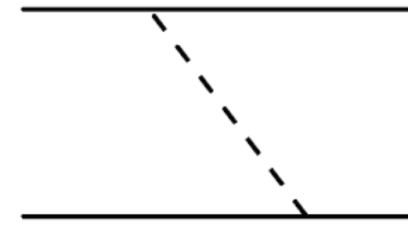
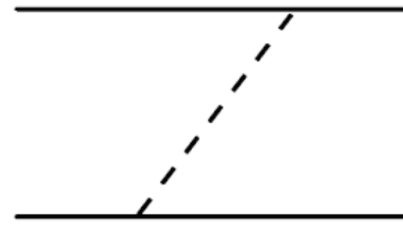
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NN and YN interactions at LO

LO potential

Weinberg power counting

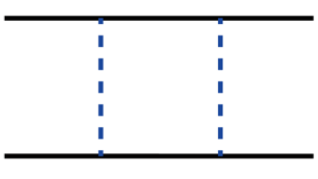


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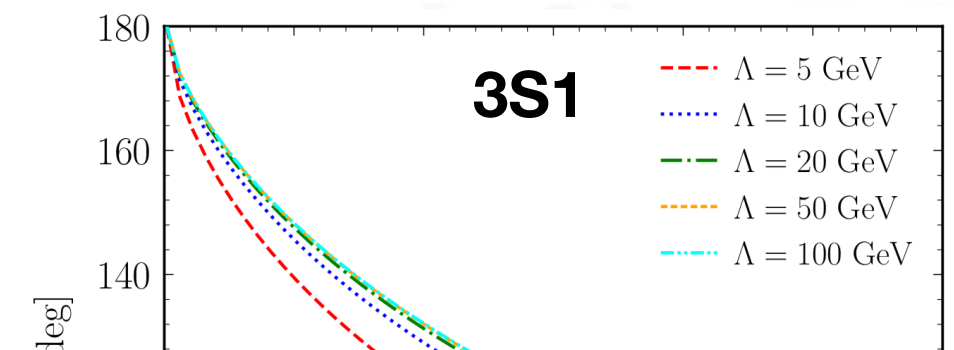
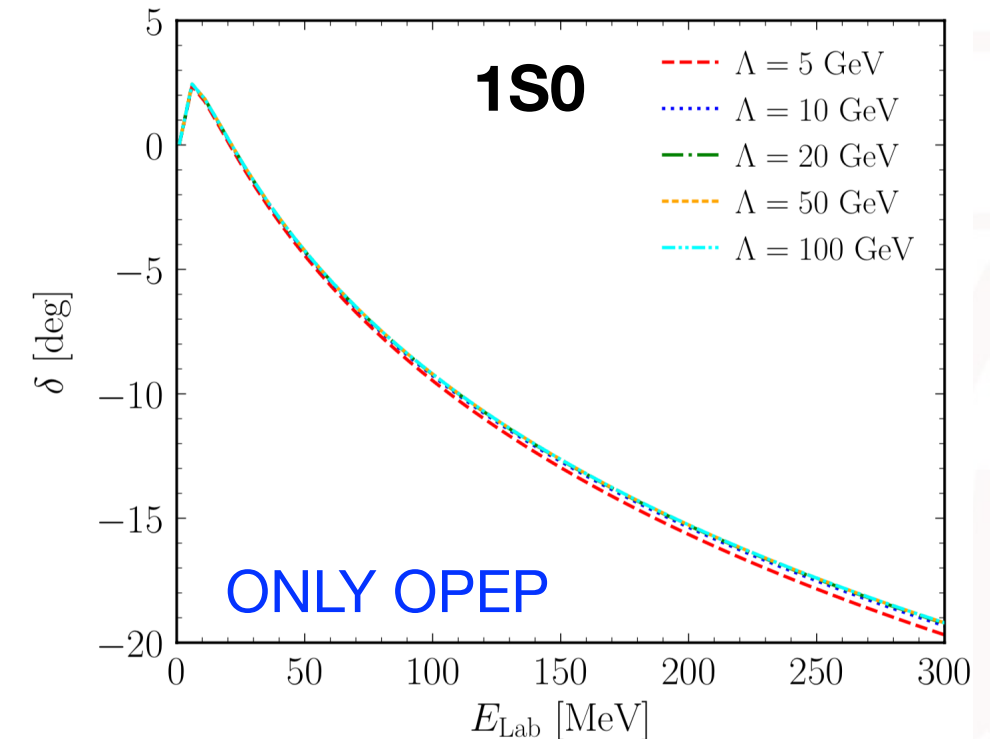
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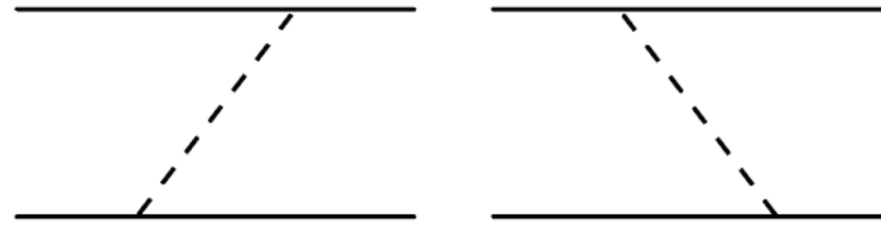
→ **Our LO potential is renormalizable!**

- **Unique solutions for all partial waves, no limit-cycle behavior**
- **Avoid finite-cutoff artefacts** inherent to the conventional NR framework

NN and YN interactions at LO

LO potential

Weinberg power counting



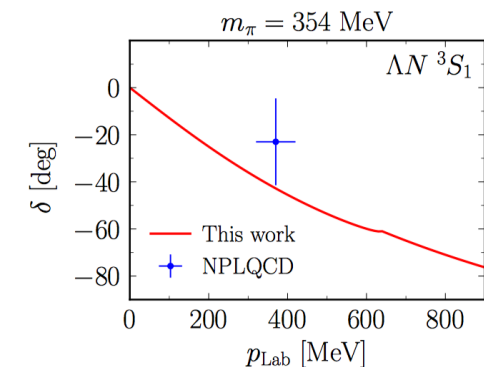
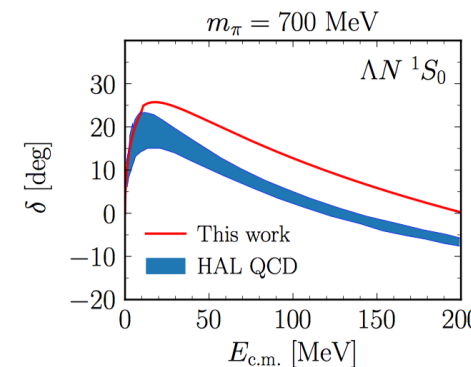
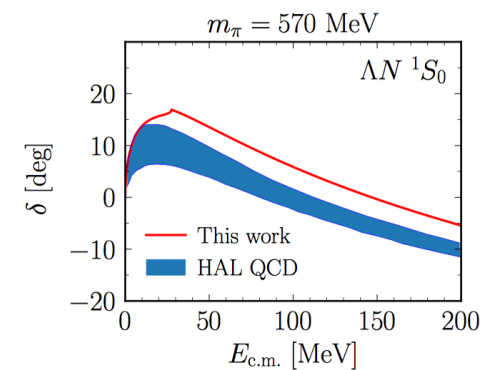
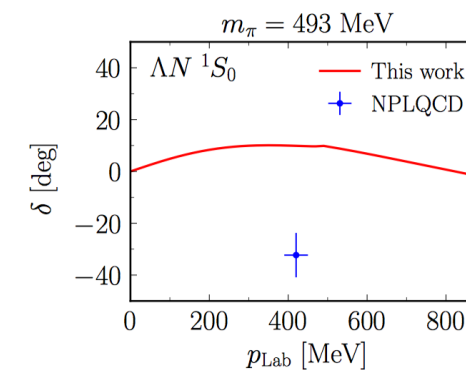
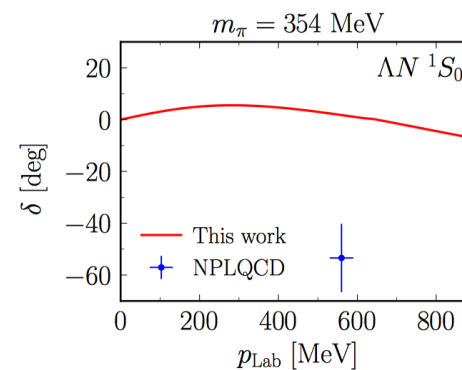
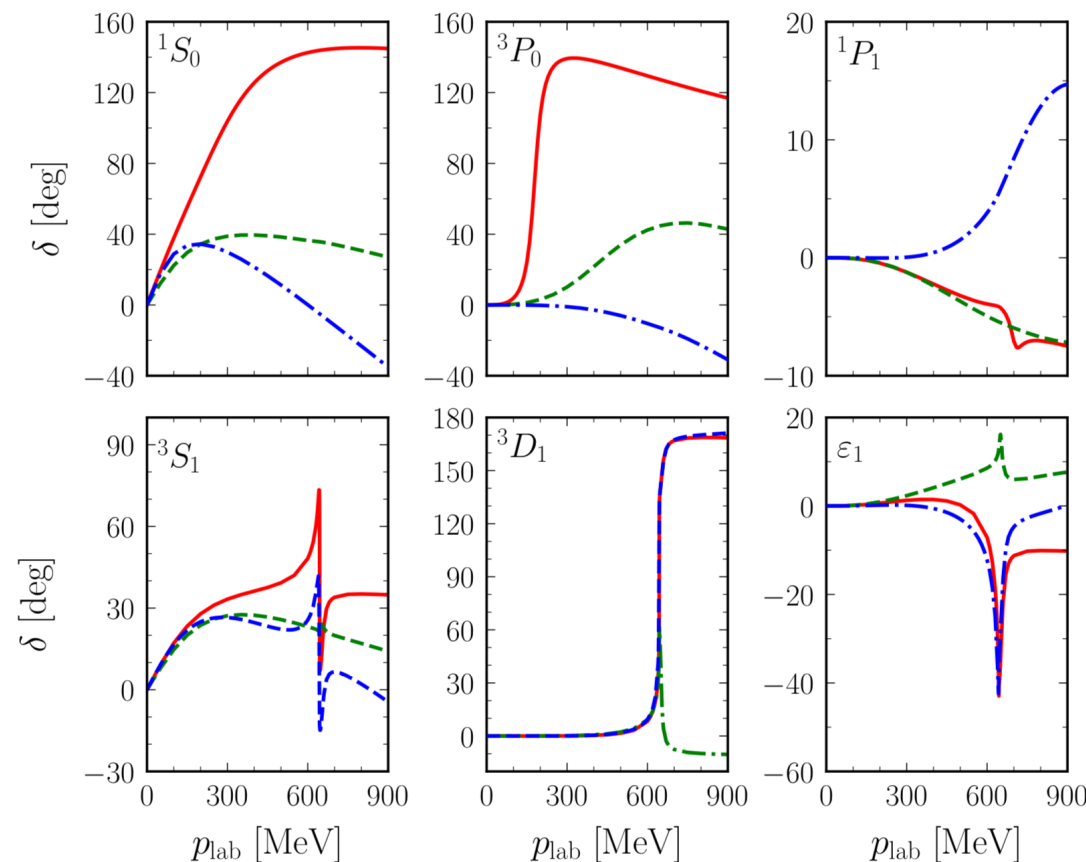
- Extend the renormalizable framework to the ΛN - ΣN interaction

$$V_{\text{LO}} = V_{\text{C}} + V_{\text{OME}} = \begin{pmatrix} V_{\Lambda N, \Lambda N}, & V_{\Lambda N, \Sigma N} \\ V_{\Sigma N, \Lambda N}, & V_{\Sigma N, \Sigma N} \end{pmatrix}$$

XLR, et al., 1911.05616

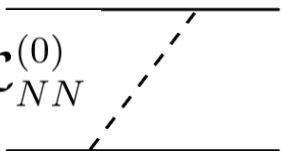
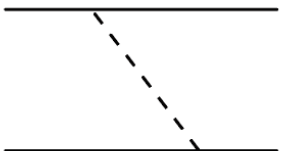
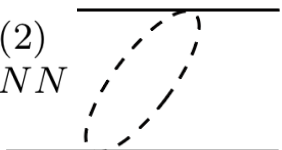
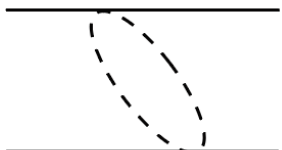
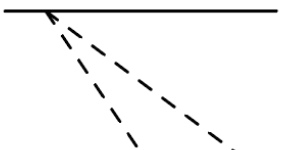
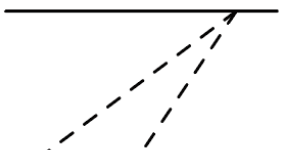
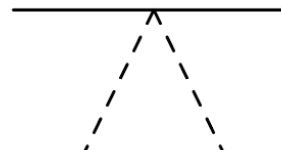
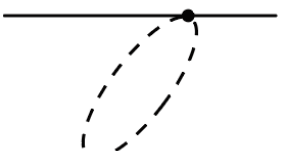
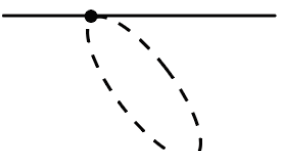
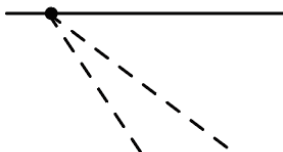
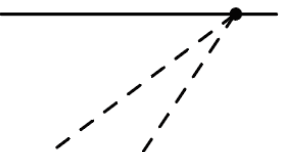
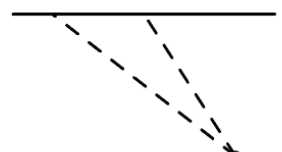
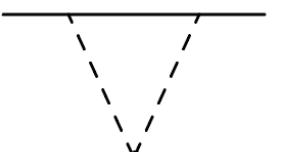
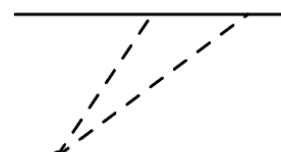
✓ 3 unknown LECs for each partial waves ($1S_0$ and $3S_1$)

✓ Carry out the renormalization



NN interaction up to NNLO

Two-nucleon irreducible time-ordered diagrams

LO		$\mathcal{L}_{NN}^{(0)}$			$\mathcal{L}_{\pi N}^{(1)} = \bar{\Psi}_N \left\{ i\not{D} - m_N + \frac{1}{2}g_A \psi \gamma^5 \right\} \Psi_N$		
NLO		$\mathcal{L}_{NN}^{(2)}$					
NNLO			$\mathcal{L}_{\pi N}^{(2)} = \bar{\Psi}_N \left\{ c_1 \langle \chi + \rangle - \frac{c_2}{4m_N^2} \langle u^\mu u^\nu \rangle (D_\mu D_\nu + \text{h.c.}) + \frac{c_3}{2} \langle u^\mu u_\mu \rangle - \frac{c_4}{4} \gamma^\mu \gamma^\nu [u_\mu, u_\nu] \right\} \Psi_N$				
							$c_1 = -0.74, c_2 = 1.81, c_3 = -3.61, c_4 = 2.17 \text{ GeV}^{-1}$ <i>D. Siemens, et al., PLB(2017)</i>

XLR, E. Epelbaum, J. Gegelia, Phys. Rev. C 106, 034001 (2022)

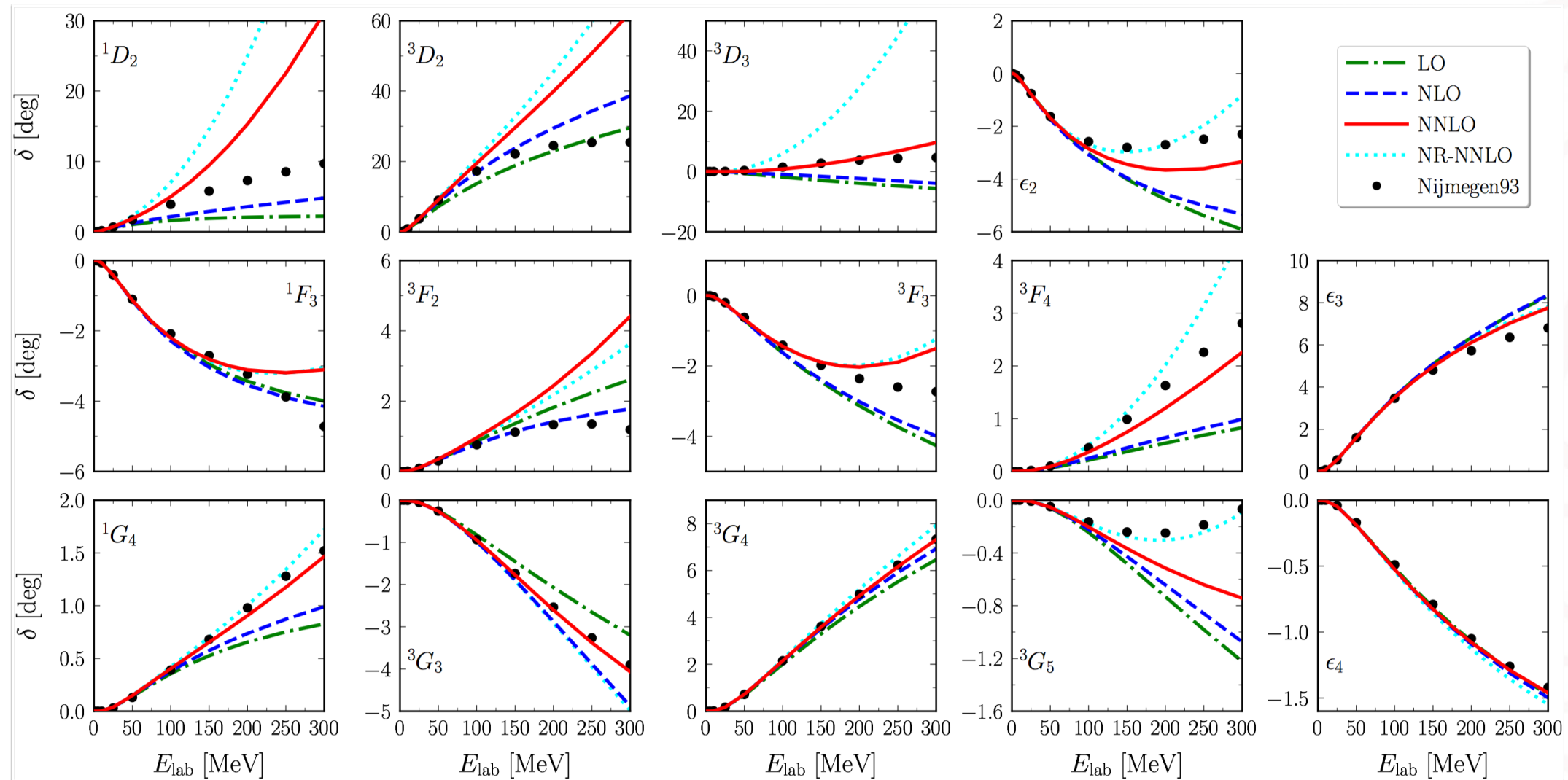
XLR, E. Epelbaum, J. Gegelia, Phys. Rev. C 113, 034002 (2026)

NN phase shifts at NLO and NNLO

On-shell T-matrix under the Born approximation

$$T(p', p) = V_{\text{OPE}}(p', p) + V_{2\pi, \text{irr}}^{(2)}(p', p) + V_{2\pi, \text{irr}}^{(3)}(p', p) + V_{\text{OPE}} G V_{\text{OPE}}$$

Prediction: phase shifts of D, F, G waves



✓ **Improve the description** of D waves; globally similar results for F, G waves

- 3G_5 : non-rel. result is accidental, c_i/m_N effect (N⁴LO) is large *D. Entem, et al., PRC 91, 014002 (2015)*

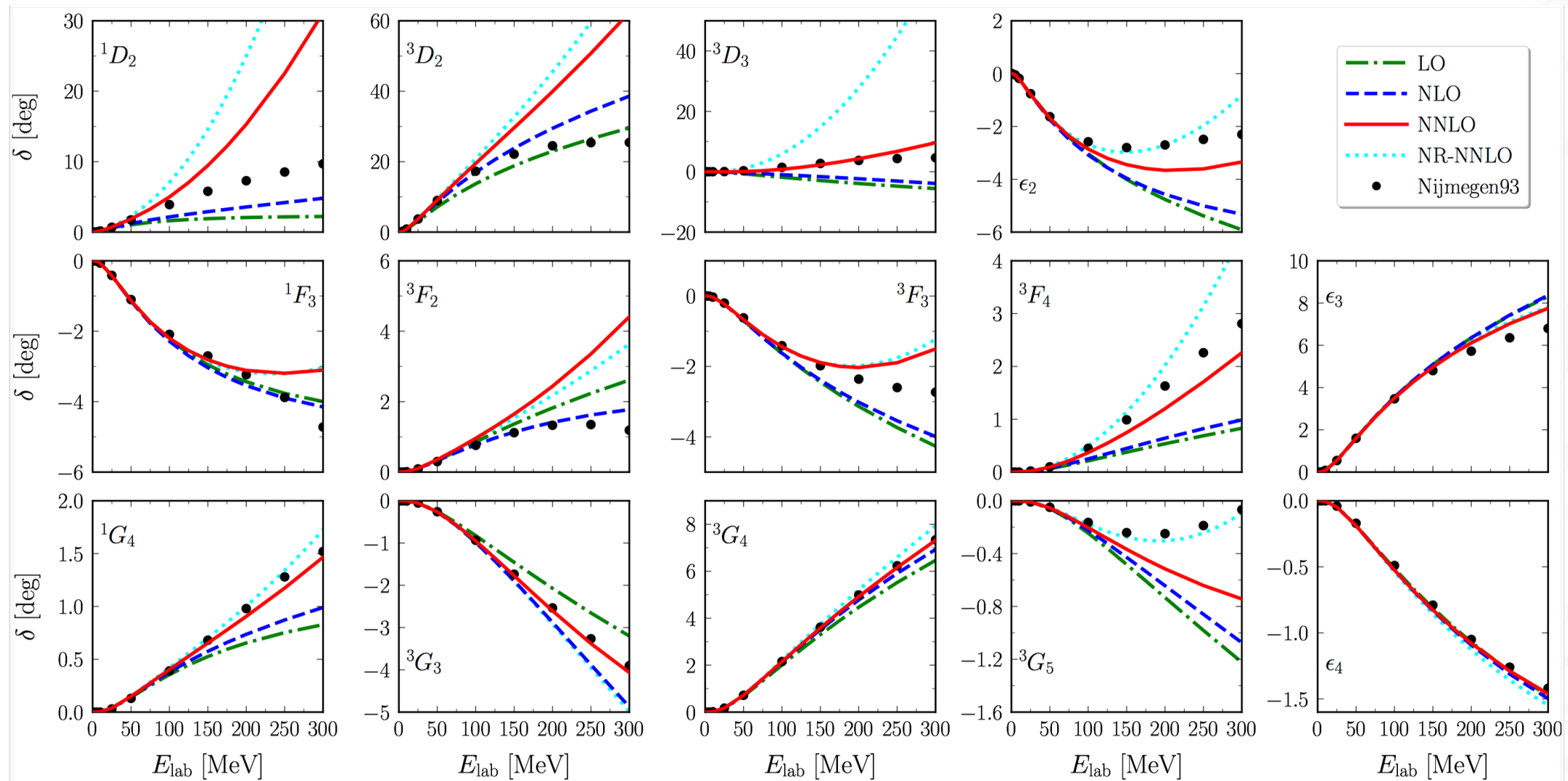
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XLR, E. Epelbaum, J. Gegelia, Phys. Rev. C 106, 034001 (2022)

NN phase shifts at NLO and NNLO

Partial wave T-matrix to describe S, P waves

- $V_{\text{NLO, NNLO}}$ **non-perturbatively** iterated in the Kadyshevsky equation

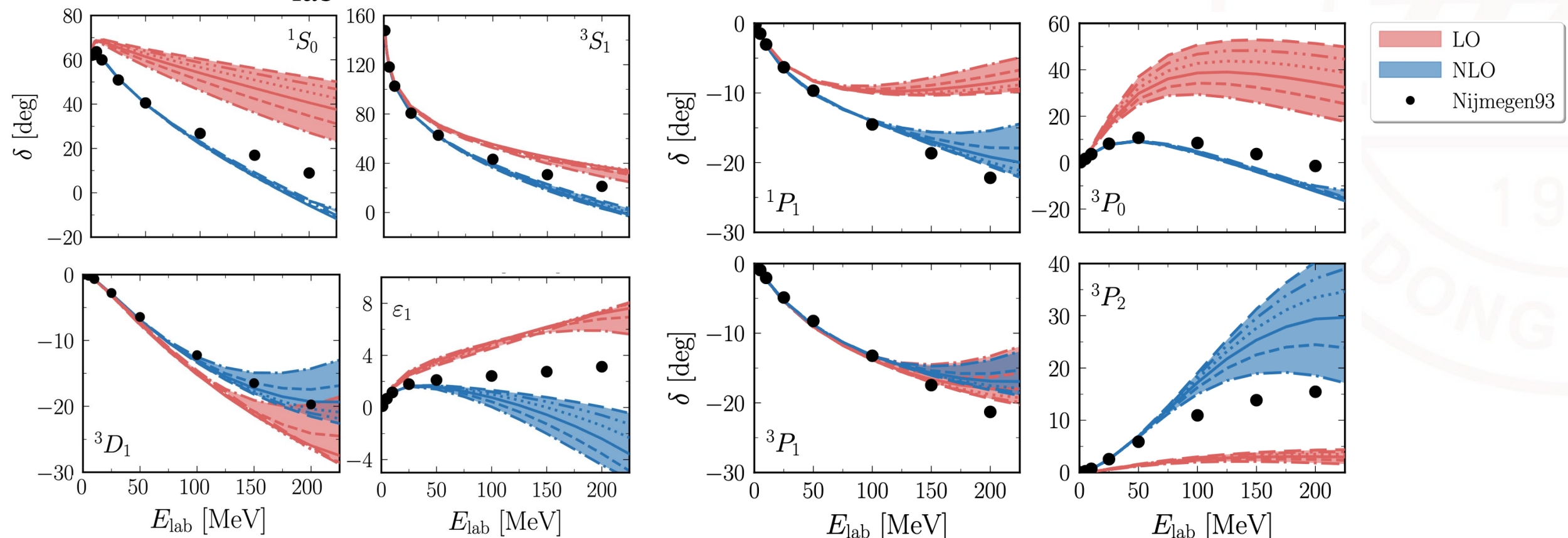
$$T_{ll'}^{sj}(p', p) = V_{ll'}^{sj}(p', p) + \sum_{l''} \int \frac{d^3k}{(2\pi)^3} V_{ll''}^{sj}(p', k) \frac{m_N^2}{2(k^2 + m_N^2)} \frac{1}{\sqrt{p^2 + m_N^2} - \sqrt{k^2 + m_N^2} + i\epsilon} T_{l''l'}^{sj}(k, p)$$

- Exponential regulator

$$F(p) = \exp(-p^{2n}/\Lambda^{2n}), \text{ with } n = 2 \quad \Lambda = 400, \dots, 650 \text{ MeV}$$

Phase shifts

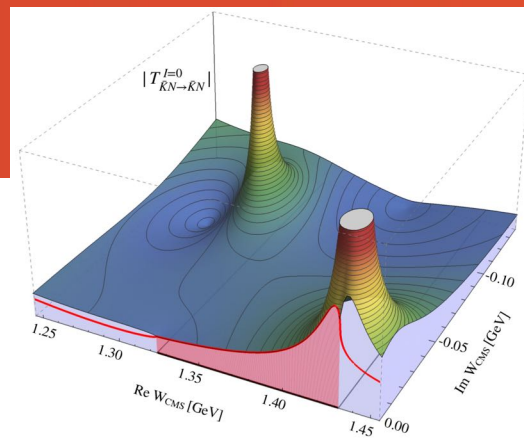
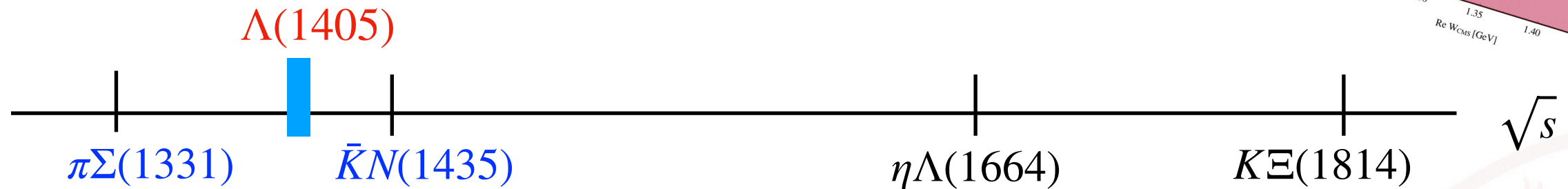
- Fit **NPWA** ($E_{\text{lab}} \leq 50 \text{ MeV}$): **np** scattering Nijmegen 93



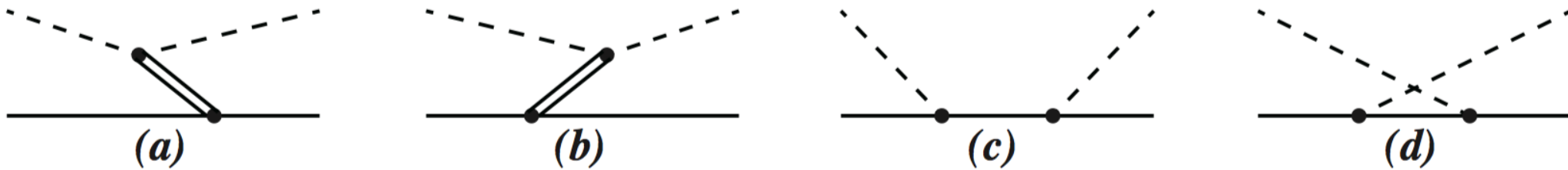
XLR, E. Epelbaum, J. Gegelia, PRC 113, 034002 (2026)

$\bar{K}N$ interaction at LO

- Four coupled channels $\bar{K}N$, $\pi\Sigma$, $\eta\Lambda$, $K\Xi$



- Focus on the S-wave potential



- Vector mesons included as explicit degrees of freedom

✓ VME potential instead of the Weinberg-Tomozawa term

Weinberg:1968de, Borasoy:1995ds

✓ Improve the ultraviolet behaviour without changing the low-energy physics

- Solve coupled-channel scattering equation

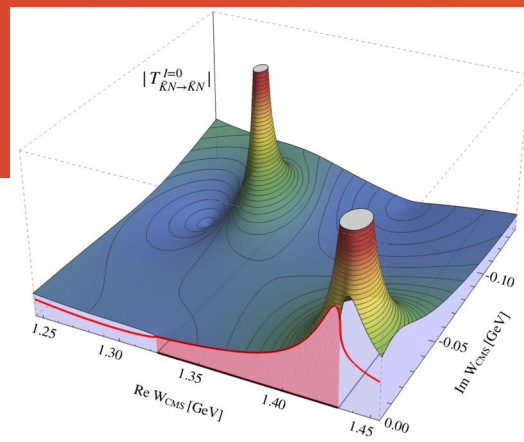
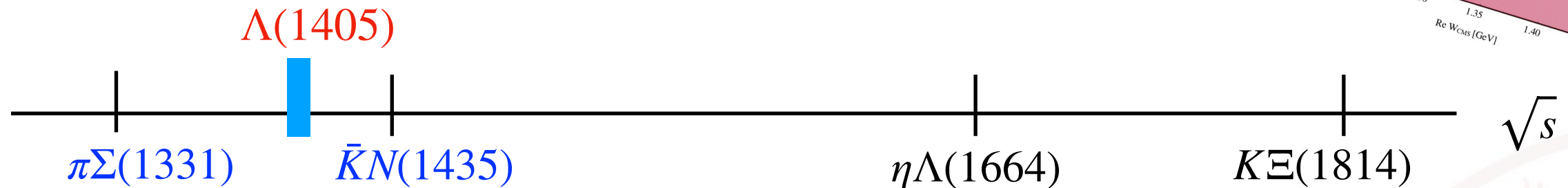
XLR, et al. EPJC 80 (2020) 406

✓ Taking into account the off-shell effects of potential

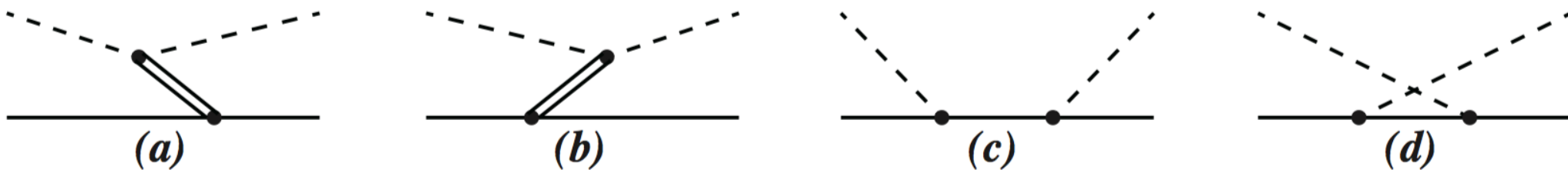
✓ Using subtractive renormalization to obtain the renormalized T-matrix

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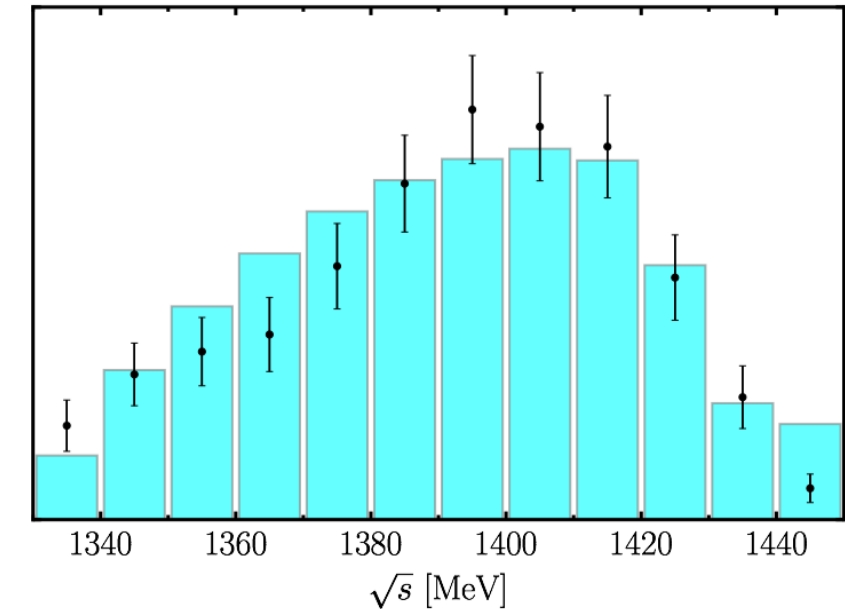
No free parameters needed to be fitted!

$\bar{K}N$ interaction in TOPT

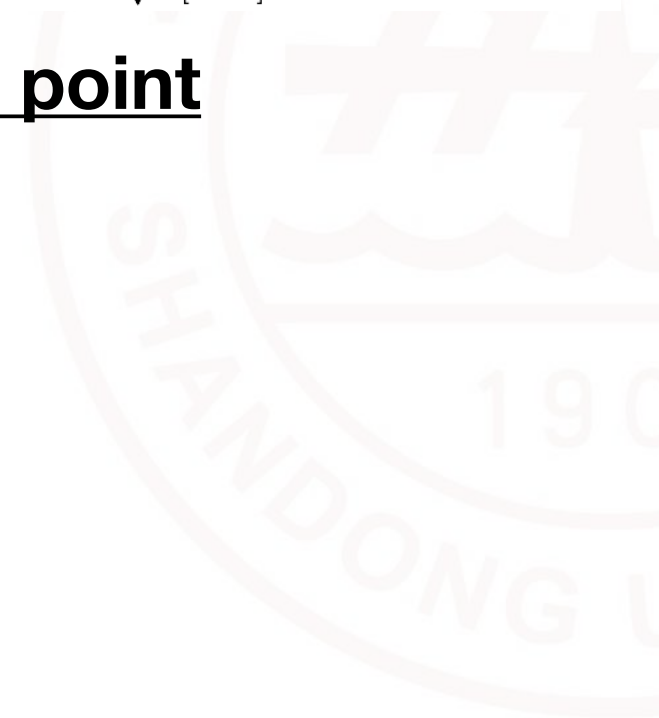
- Two pole positions of $\Lambda(1405)$ **XLR, et al. EPJC81 (2021) 582**

		lower pole	higher pole
This work	$F_0 = F_\pi$	$1337.7 - i 79.1$	$1430.9 - i 8.0$
(LO)	$F_0 = 103.4$	$1348.2 - i 120.2$	$1436.3 - i 0.7$
NLO →	<i>Y. Ikeda, NPA(2012)</i>	$1381^{+18}_{-6} - i 81^{+19}_{-8}$	$1424^{+7}_{-23} - i 26^{+3}_{-14}$
	<i>Z.-H. Guo, PRC(2013)-Fit II</i>	$1388^{+9}_{-9} - i 114^{+24}_{-25}$	$1421^{+3}_{-2} - i 19^{+8}_{-5}$
	<i>M. Mai, EPJA(2015)-sol-2</i>	$1330^{+4}_{-5} - i 56^{+17}_{-11}$	$1434^{+2}_{-2} - i 10^{+2}_{-1}$
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$\pi\Sigma$ invariant mass spectrum



- Consistent with the NLO study by M. Mai at physical point

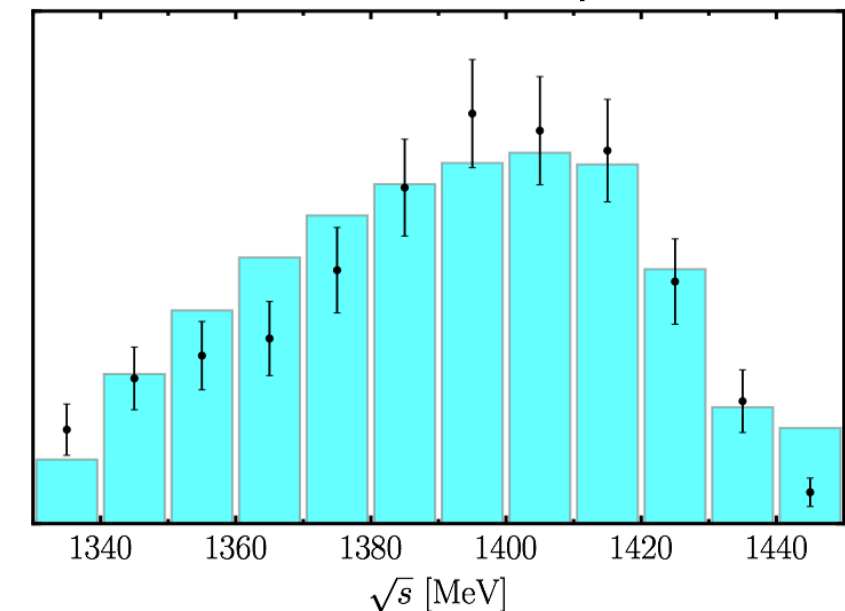


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$\pi\Sigma$ invariant mass spectrum



- Consistent with the NLO study by M. Mai at physical point

- Extend to the unphysical quark mass region $m_\pi \approx 200$ MeV, $m_K \approx 487$ MeV

- Consistent with the BaSc results **PRL 132, 051901 (2024)**

	BaSc [PRL2024]	This work			
$\Lambda(1405)$	z_R [MeV]	z_R [MeV]	$g_{\pi\Sigma}$	$g_{\bar{K}N}$	$ g_{\pi\Sigma} / g_{\bar{K}N} $
Lower pole	1392(18)	1387.14	$0.021 + i1.87$	$0.017 + i1.55$	1.21
Higher pole	1455(21) - i11.5(6.0)	1469.86 - i4.71	$0.038 + i0.98$	$1.51 - i1.22$	0.50

XLR, Phys. Lett. B 855 (2024) 138802

Meson–baryon interaction @ NLO

□ Maintain the scattering T-matrix renormalizable

- Take LO potential non-perturbatively
- Include higher order corrections perturbatively

□ Up to next-to-leading order

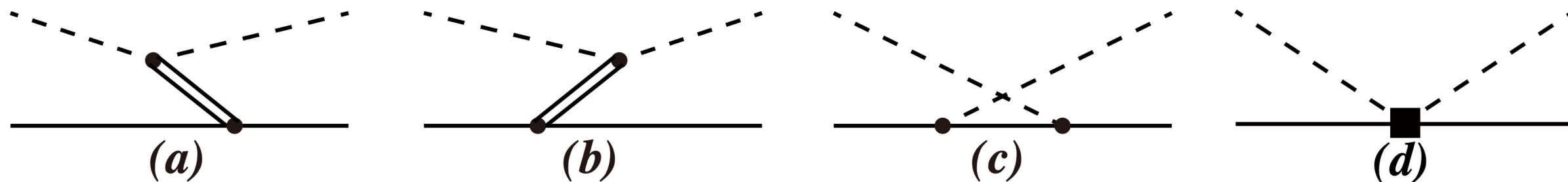
- Potential: $V = V_{\text{LO}} + V_{\text{NLO}}$
- T-matrix: $T = T_{\text{LO}} + T_{\text{NLO}}$

$$T_{\text{LO}} = V_{\text{LO}} + V_{\text{LO}}GT_{\text{LO}} \quad (\text{non-perturbative})$$

$$T_{\text{NLO}} = V_{\text{NLO}} + T_{\text{LO}}GV_{\text{NLO}} + V_{\text{NLO}}GT_{\text{LO}} + T_{\text{LO}}GV_{\text{NLO}}GT_{\text{LO}}$$

- Use **the subtractive renormalization** to remove divergent terms and power-counting breaking terms

KN interaction up to NLO



Chiral effective Lagrangians at NLO

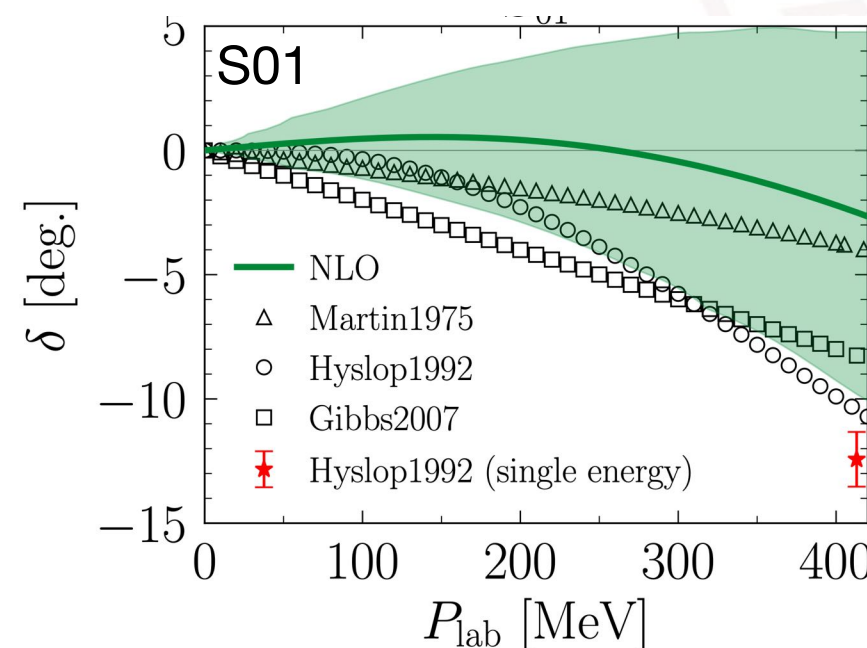
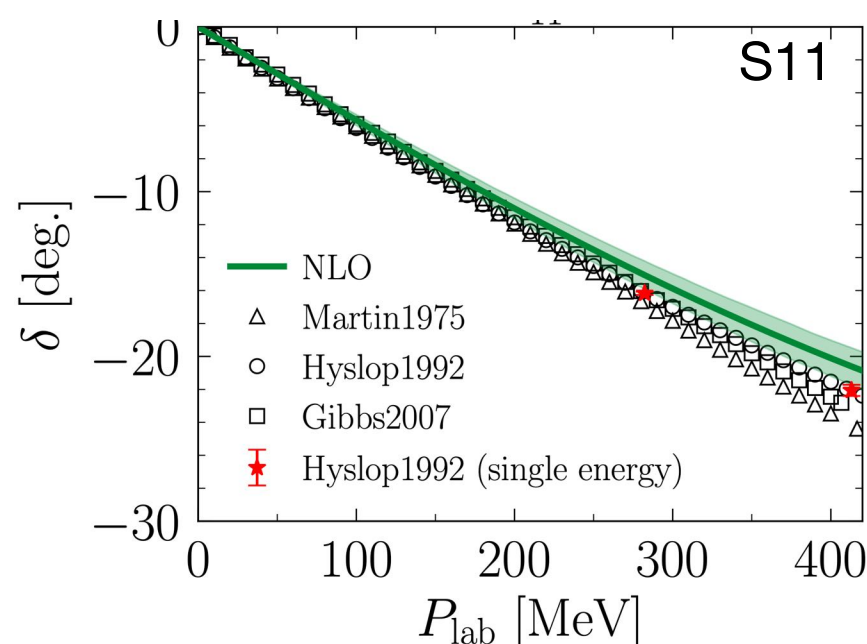
For S-wave scattering

$$\begin{aligned} \mathcal{L}_{\text{NLO}} = & b_0 \langle \chi_+ \rangle \langle \bar{B}B \rangle + b_D \langle \bar{B} \{ \chi_+, B \} \rangle + b_F \langle \bar{B} [\chi_+, B] \rangle \\ & + d_1 \langle \bar{B} \{ u_\mu, [u^\mu, B] \} \rangle + d_2 \langle \bar{B} [u_\mu, [u^\mu, B]] \rangle \\ & + d_3 \langle \bar{B} u_\mu \rangle \langle u^\mu B \rangle + d_4 \langle \bar{B}B \rangle \langle u^\mu u_\mu \rangle, \end{aligned}$$

Four parameters at NLO

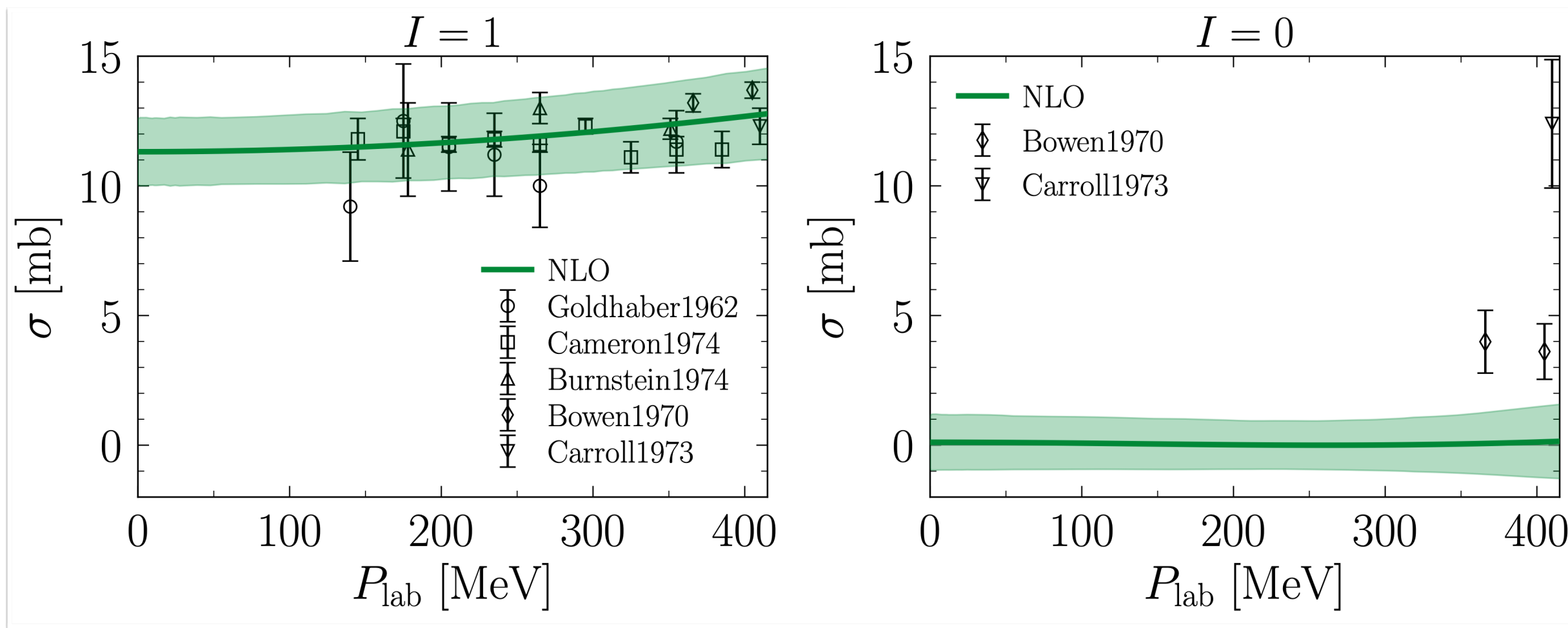
- Fixed by RQCD's $b_{0,D,F}$
- Determined by S_{11} and S_{01} scattering lengths

Prediction of s-wave phase shifts XLR, Phys. Rev. D 113 (2026) 114027

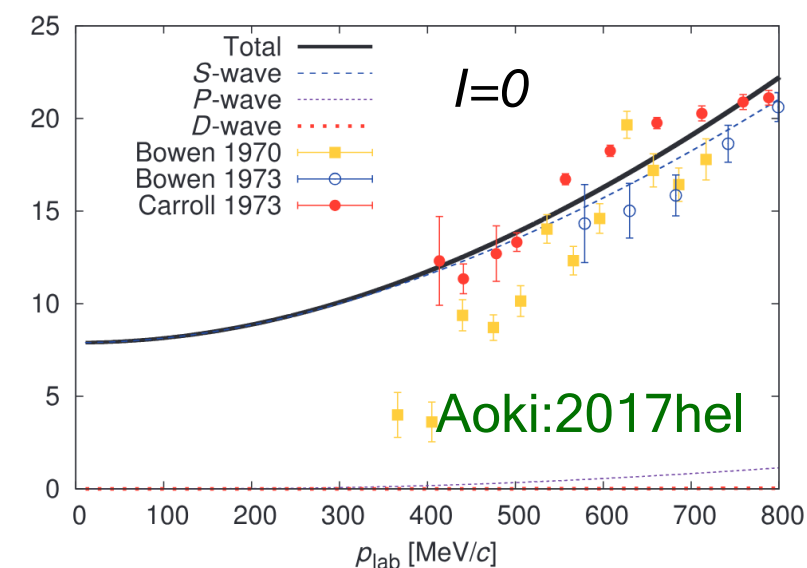


KN interaction up to NLO

■ S-wave contributions to the total cross sections



- $I=1$ channel: **agrees well** with experimental data and slowly increases with P_{lab}
- $I=0$ channel: NLO s-wave contribution small
✓ P-wave contribution dominant !



XLR, Phys. Rev. D 113 (2026) 114027

Summary

- ❑ Tentatively proposed **a unified framework** to formulate the baryon-baryon and meson-baryon interactions
 - based on time-ordered perturbation theory
 - using covariant chiral effective Lagrangians
- ❑ **Baryon-baryon interactions**
 - Obtained **non-singular LO BB potential**
 - ✓ Avoid finite-cutoff artefacts and take cutoff $\Lambda \rightarrow \infty$
 - Formulated the chiral NN potential up to **NNLO**
 - ✓ Achieved a rather reasonable description of phase shifts
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Thank you for your attention!