

Xiong-Hui Cao 曹雄辉

XHC, F.-K. Guo, Q.-Z. Li, D.-L. Yao, Nature Commun. 16 (2025) 6979;

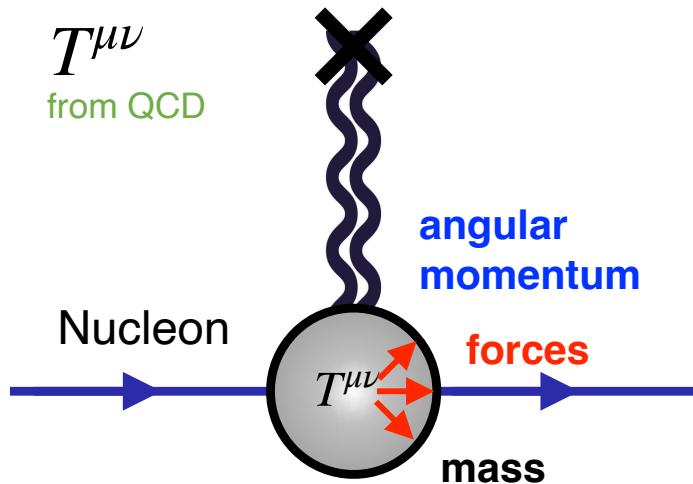
XHC, F.-K. Guo, Q.-Z. Li, B.-W. Wu, D.-L. Yao, Eur. Phys. J. ST 235 (2026) 9, 2121



第十二届高能物理大会暨第十四届晨光杯评选

2026年7月15日，广州，华南师范大学

Gravitational structure of nucleons



- Gravity couples to matter through energy-momentum tensor (EMT) $T^{\mu\nu}$

study gravitational forces within hadrons



study mechanical properties of hadrons



- Definition:

Forces inside hadron, **mechanical picture of quark confinement?**

$$\langle N(p') | T^{\mu\nu} | N(p) \rangle = \frac{1}{4m_N} \bar{u}(p') \left[A(t) P^\mu P^\nu + J(t) \left(i P^{\{\mu} \sigma^{\nu\} \rho} \Delta_\rho \right) + D(t) \left(\Delta^\mu \Delta^\nu - g^{\mu\nu} \Delta^2 \right) \right] u(p)$$

Kobzarev, Okun (1962)
Pagels (1966)

Mass normalization:

$$m_N = \int d^3r T_{00}(r)$$

$$A(0) = 1$$

Spin normalization:

$$J^i = \epsilon_{ijk} \int d^3r r_j T_{0k}(r)$$

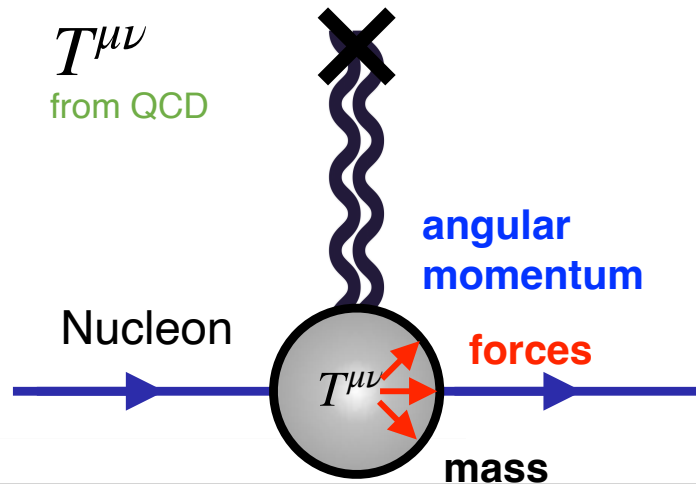
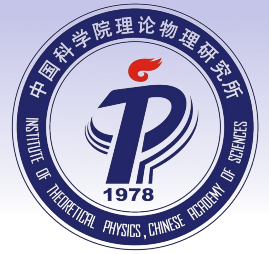
$$J(0) = 1/2$$

D-term: $D \equiv D(0)$

$$D = -\frac{m_N}{2} \int d^3r \left(r_i r_j - \frac{1}{3} \delta_{ij} \right) T_{ij}(r)$$

$$D = ?$$

Gravitational structure of nucleons



**D-term as the
“last unknown global property”**

Polyakov, and Schweitzer, (2018)

em:	$\partial_\mu J_{\text{em}}^\mu = 0$	$\langle N' J_{\text{em}}^\mu N \rangle \longrightarrow$	$Q = 1.602176487(40) \times 10^{-19} \text{C}$
			$\mu = 2.792847356(23) \mu_N$

weak:	PCAC	$\langle N' J_{\text{weak}}^\mu N \rangle \longrightarrow$	$g_A = 1.2694(28)$
			$g_p = 8.06(55)$

gravity:	$\partial_\mu T_{\text{grav}}^{\mu\nu} = 0$	$\langle N' T_{\text{grav}}^{\mu\nu} N \rangle \longrightarrow$	$m = 938.272013(23) \text{ MeV}/c^2$
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$J = \frac{1}{2}$
$D = ?$

D -term form factor from experiments

- All GFFs can be measured from DVCS through sum rules derived for GPDs Ji, PRL (1997); Ji, Melnitchouk and Song, PRD (1997)

- π^0 : $\gamma^* \gamma \rightarrow \pi^0 \pi^0$ in $e^+ e^-$ Belle data: PRD 93, 032003 (2016)

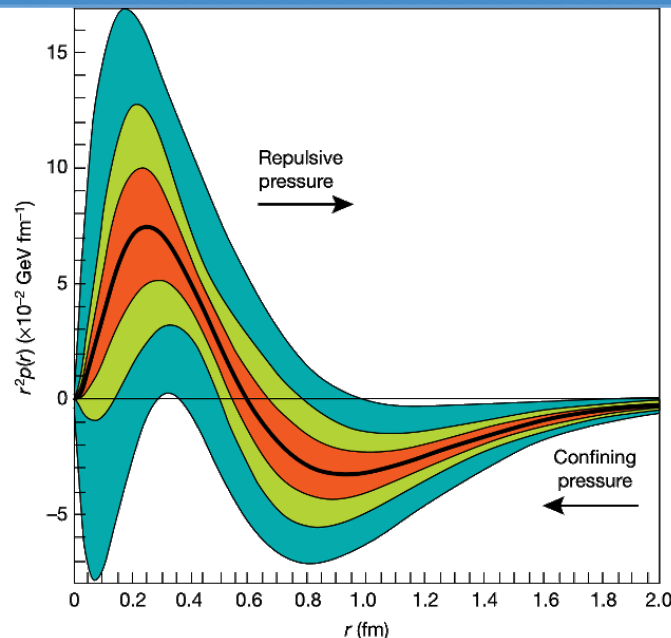
$$D_{\pi^0}^Q \simeq -0.7 \text{ at } \langle Q^2 \rangle = 16.6 \text{ GeV}^2$$

Kumano, Song, Teryaev, PRD (2018)

Chiral symmetry: total $D_{\pi^0} \simeq -1$ (gluons contribute the rest)

- Proton: Burkert, Elouadrhiri, Girod, Nature 557, 396 (2018)

“Pressure” distribution inside proton using DVCS data

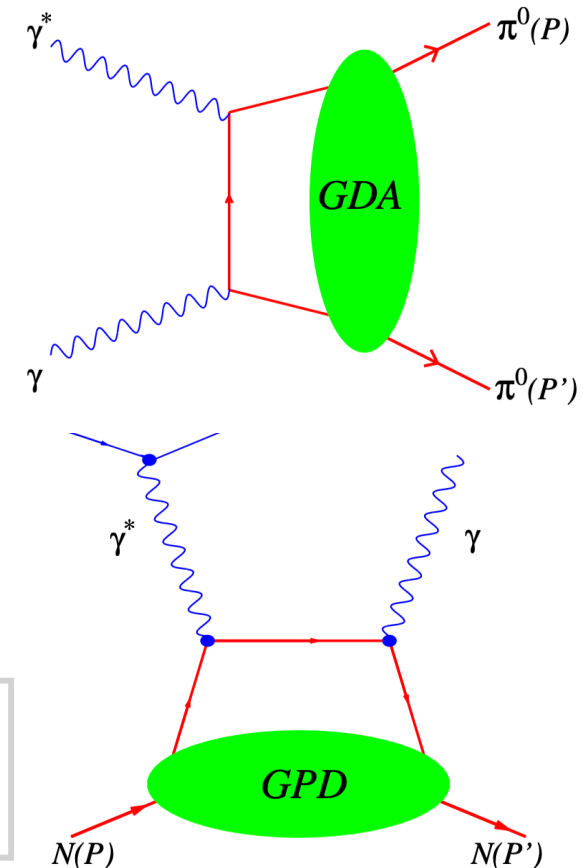


Subtraction constant in dispersion relations for DVCS amplitudes

$C_{\mathcal{H}}(t, \mu^2) \rightarrow D^Q(t, \mu^2)$ scale-dependent

model-dependent

Explore t range at EIC & EicC



Scalar radius puzzle

- Trace FF:

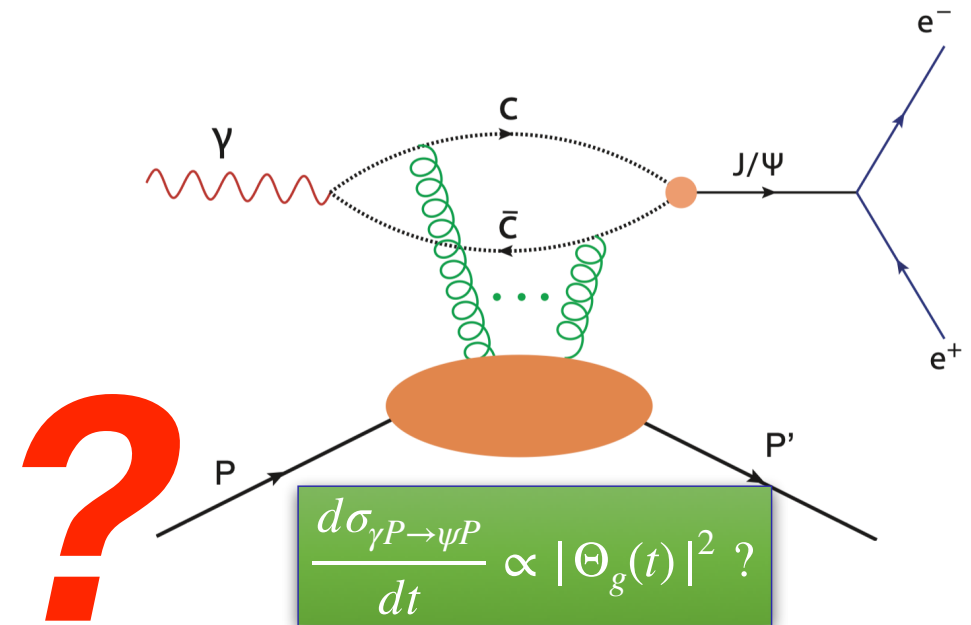
$$\langle N(p') | T^\mu_\mu | N(p) \rangle = m_N \bar{u}(p') \left[A(t) - \frac{t}{4m_N^2} [A(t) - 2J(t) + 3D(t)] \right] u(p) \equiv \bar{u}(p') \Theta(t) u(p)$$

- Trace anomaly in QCD: $T^\mu_\mu = \frac{\beta(g)}{2g} F^{a,\mu\nu} F^a_{\mu\nu} + (1 + \gamma_m) \sum_q m_q \bar{\psi}_q \psi_q$

- Scalar radius: $\Theta(t) = m_N (1 + t \langle r_\Theta^2 \rangle / 6 + \dots)$

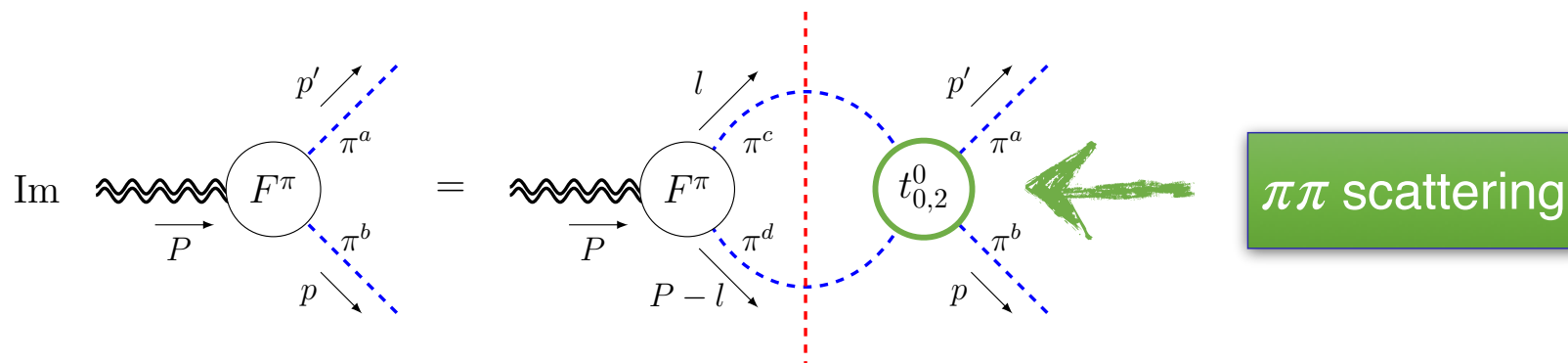
It could be extracted from the threshold photoproduction of the vector-meson, e.g., J/ψ , and the fit result is ~ 0.55 fm
Kharzeev, PRD (2021).....

Recently, two LQCD calculations at near physical quark mass result in a large scalar radius, ~ 1 fm
Hackett et al., PRL (2024); Wang et al. [χQCD], PRD (2024)



Unitarity relation for pion GFFs

Partial-wave unitarity



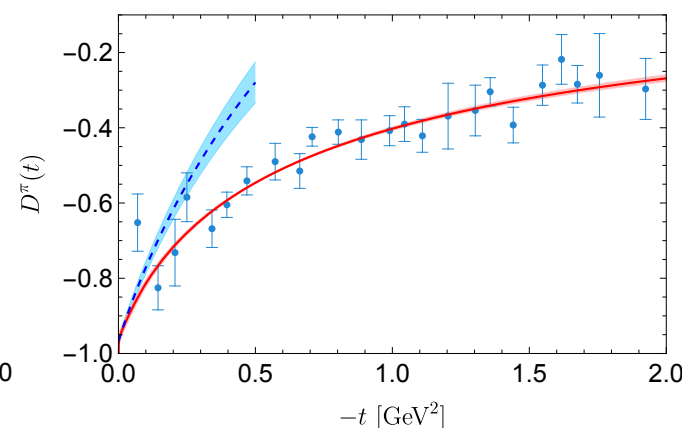
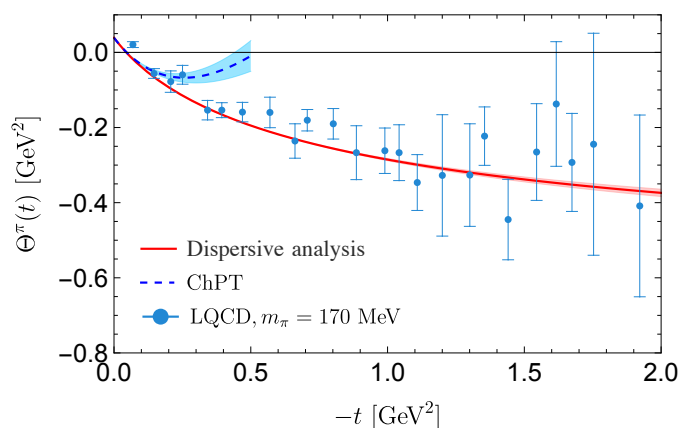
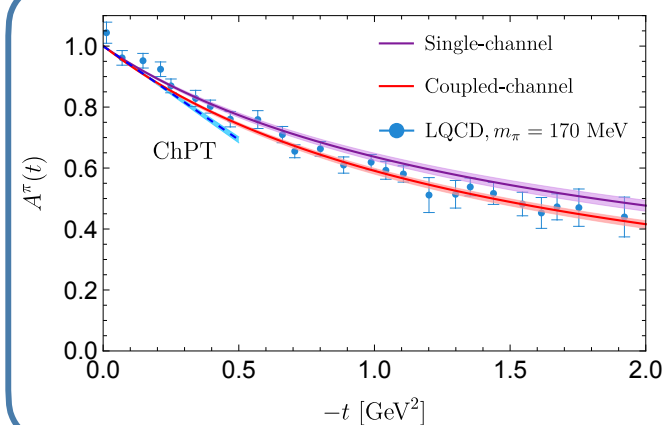
$$\text{disc } A^\pi(t) = 2iA^\pi(t)\sin\delta_2^0(t)e^{-i\delta_2^0(t)}\theta(t-t_\pi)$$

Omnes problem

Pion

Prediction, Not fit

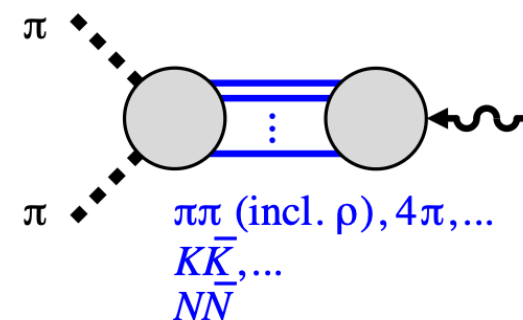
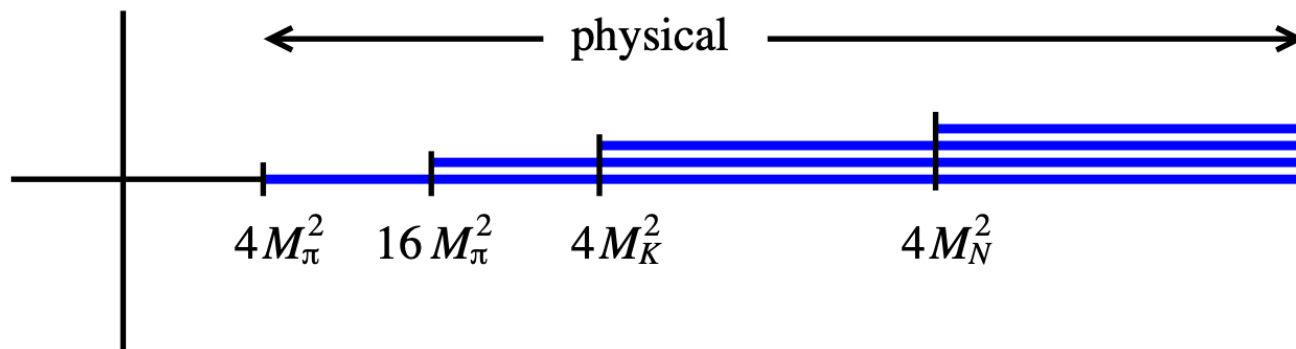
LQCD ($m_\pi \simeq 170$ MeV): [Hackett et al., PRD \(2023\)](#)



Unitarity relation for nucleon GFFs

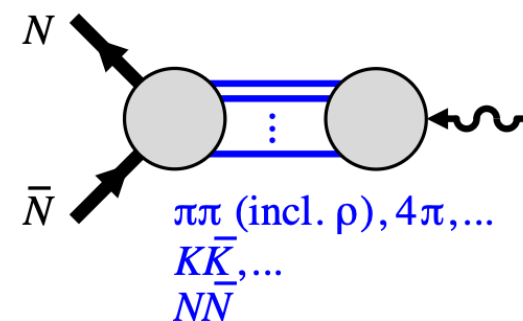
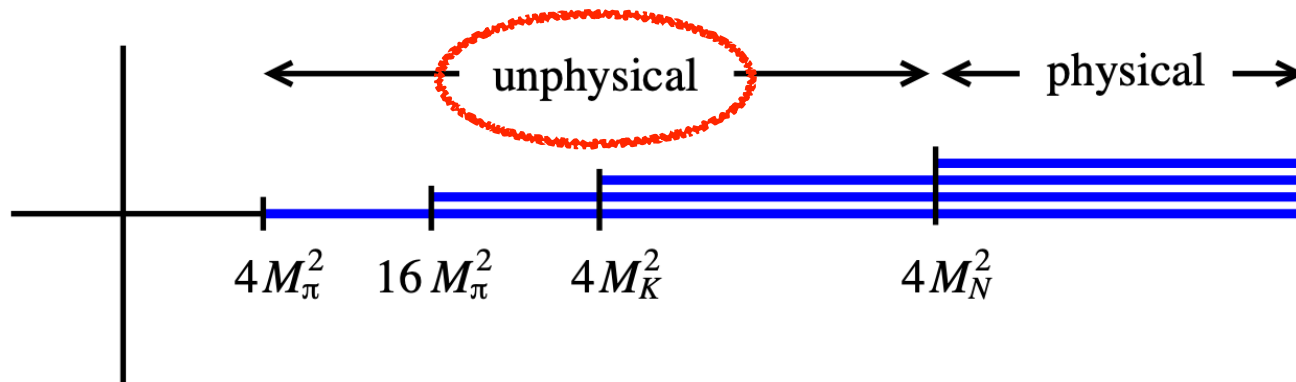


Pion GFFs



Nucleon GFFs

Analytic continuation!



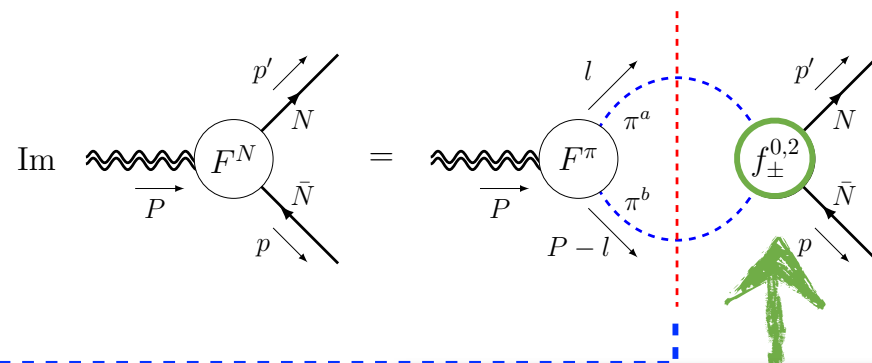
Unitarity relation for nucleon GFFs



Discontinuity

$$\text{Disc} \left\langle N(p') \bar{N}(p) \left| \hat{T}^{\mu\nu}(0) \right| 0 \right\rangle$$

$$\propto \sum_n \langle N(p') \bar{N}(p) | n \rangle \langle n | \hat{T}^{\mu\nu}(0) | 0 \rangle^* \delta^4(p + p' - p_n)$$



$\pi\pi(K\bar{K}) \rightarrow N\bar{N}$ scattering

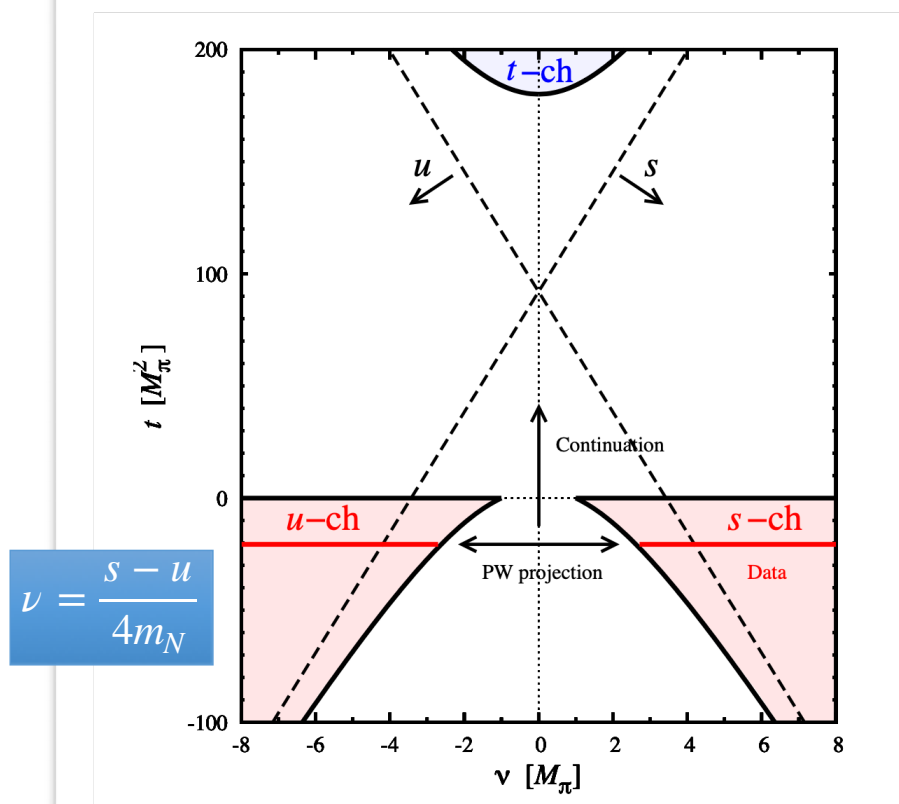
- Idea: Construct $\pi\pi \rightarrow N\bar{N}$ PWAs $\Gamma_i(t)$ from $\pi N \rightarrow \pi N$ scattering data using amplitude analysis techniques

Analytically continue PWAs to region $t > 4M_\pi^2$ via Roy-Steiner eq.

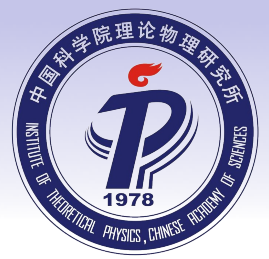
Ditsche, et al., JHEP (2012);

Hoferichter, et al., Phys. Rept. (2016);

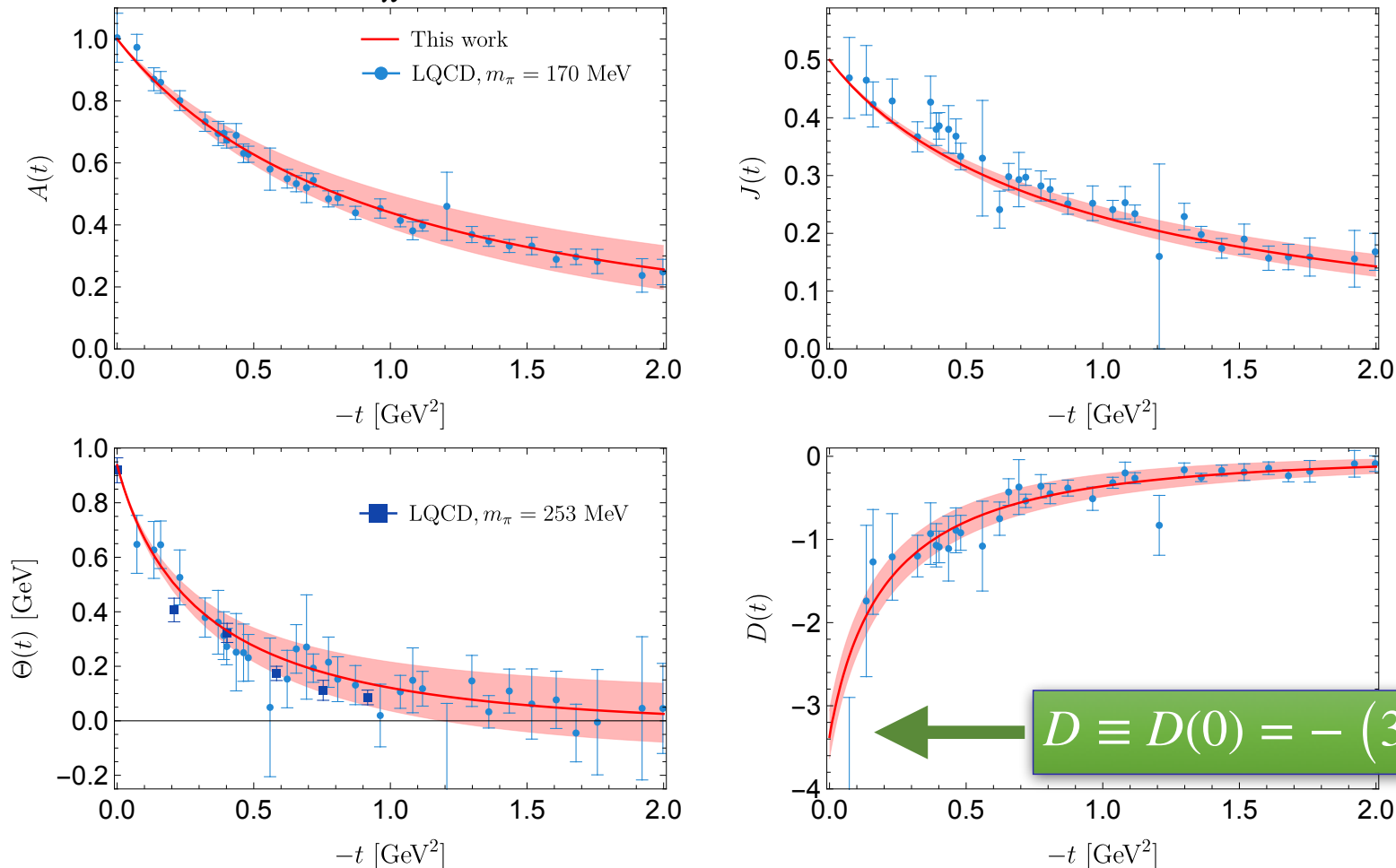
XHC, Q.-Z. Li & H.-Q. Zheng, JHEP (2022)



Nucleon GFFs



● Predictions: LQCD ($m_\pi \simeq 170$ MeV): [Hackett et al., PRL \(2024\)](#)



LQCD ($m_\pi \simeq 253$ MeV): [Wang et al. \[\$\chi\$ QCD\], PRD \(2024\)](#)

- ☑ GFFs be **model-independently** predicted using data-driven dispersion relations + ChPT
- ☑ Not a fit!

Physics of the scalar (trace) radius

- The scalar (trace) radius is larger than the mass radius as long as $D < 0$

$$T_{\mu}^{\mu}$$

$$\sqrt{\langle r_{\Theta}^2 \rangle} = 0.97^{+0.03}_{-0.03} \text{ fm}$$



“the largest radius
in strong interaction”

$$> \sqrt{\langle r_{C,p}^2 \rangle} = 0.84 \text{ fm}$$

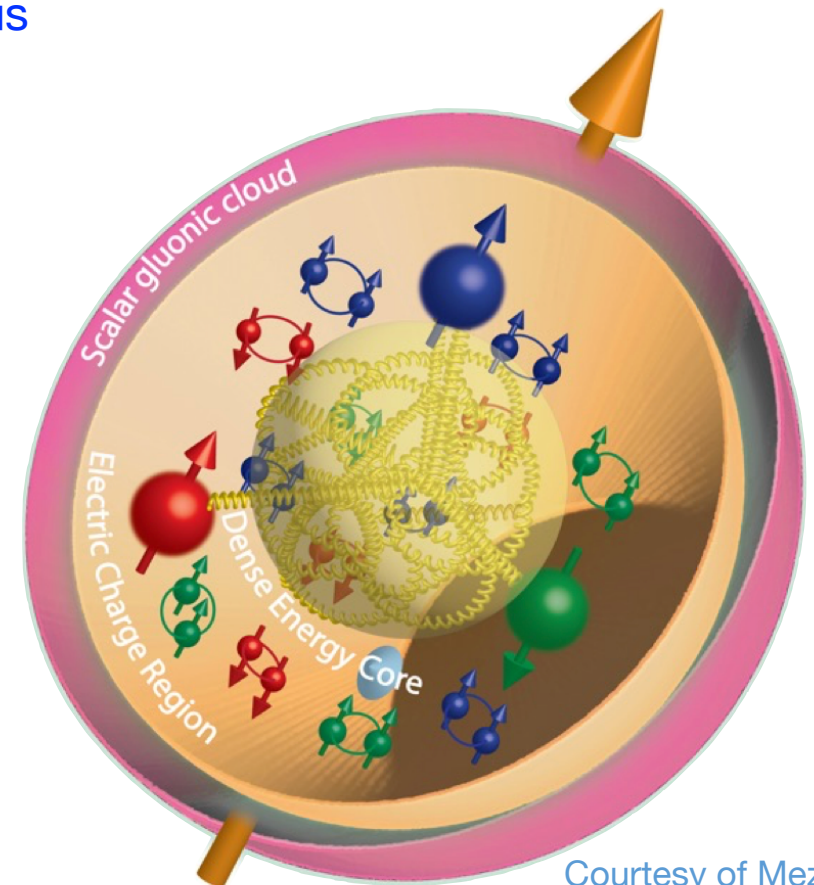
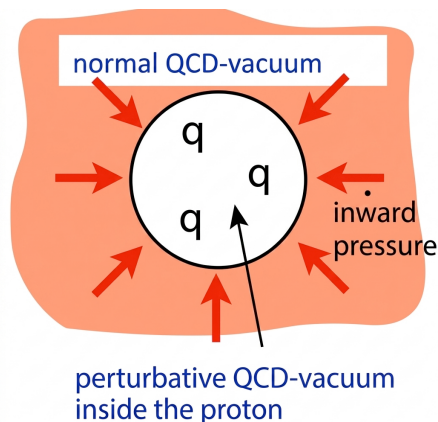
proton charge radius

$$T^{00}$$

$$> \sqrt{\langle r_{\text{Mass}}^2 \rangle} = 0.70^{+0.03}_{-0.04} \text{ fm}$$

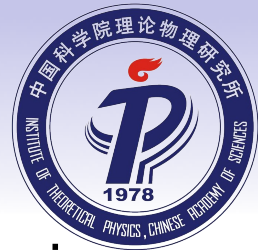
- In the MIT bag model, scalar radius may be considered as the bag (confinement) radius $\sim 1 \text{ fm}$ (physical boundary of confinement)

Ji, Front.Phys.(Beijing) (2021)



Courtesy of Meiziani

Summary



☑ The pion, kaon and nucleon GFFs are precisely determined using dispersive method with some inputs.

☑ Nucleon static properties:

$$D = -3.38^{+0.34}_{-0.35}, \quad \sqrt{\langle r_{\Theta}^2 \rangle} = 0.97^{+0.03}_{-0.03} \text{fm} > \sqrt{\langle r_{C,p}^2 \rangle} = 0.84 \text{fm} > \sqrt{\langle r_{\text{Mass}}^2 \rangle} = 0.70^{+0.03}_{-0.04} \text{fm}$$

proton charge radius

☑ Foundation for further analyses, interpretations, etc.

Thank you for your attention!

Nucleon GFFs: D-term & radii

● Various radii in the Breit frame

Radius of the scalar trace density

$$\langle r_{\Theta}^2 \rangle = \frac{6\dot{\Theta}(0)}{m_N} = 6\dot{A}(0) - \frac{9D}{2m_N^2}$$

Radius of the mass (energy) density

$$\langle r_{\text{Mass}}^2 \rangle = 6\dot{A}(0) - \frac{3D}{2m_N^2}$$

Mechanical radius Polyakov, PLB (2003)
Polyakov, & Schweitzer, JMPA (2018)

$$\langle r_{\text{Mech}}^2 \rangle = \frac{6D}{\int_{-\infty}^0 dt D(t)}$$

Radius of the angular momentum (spin) density $J(t) + \frac{2}{3}t \frac{dJ(t)}{dt}$

$$\langle r_J^2 \rangle = 20J'(0)$$

Polyakov, PLB (2003)
Lorcé et al., PLB (2018)

Quantity	Result	Error budget
D-term	$-3.38^{+0.34}_{-0.35}$	$+(0.18)_{\text{ChPT}}(0.12)_{\text{PWA}}(0.26)_{\text{eff}}$ $-(0.16)_{\text{ChPT}}(0.12)_{\text{PWA}}(0.29)_{\text{eff}}$
$\sqrt{\langle r_{\Theta}^2 \rangle}$	$0.97^{+0.03}_{-0.03}$	$+(0.01)_{\text{ChPT}}(0.01)_{\text{PWA}}(0.03)_{\text{eff}}$ $-(0.02)_{\text{ChPT}}(0.01)_{\text{PWA}}(0.26)_{\text{eff}}$
$\sqrt{\langle r_{\text{Mass}}^2 \rangle}$	$0.70^{+0.03}_{-0.04}$	$+(0.02)_{\text{ChPT}}(0.01)_{\text{PWA}}(0.02)_{\text{eff}}$ $+(0.02)_{\text{ChPT}}(0.01)_{\text{PWA}}(0.03)_{\text{eff}}$
$\sqrt{\langle r_{\text{Mech}}^2 \rangle}$	$0.72^{+0.09}_{-0.08}$	$+(0.02)_{\text{ChPT}}(0.00)_{\text{PWA}}(0.09)_{\text{eff}}$ $-(0.03)_{\text{ChPT}}(0.01)_{\text{PWA}}(0.07)_{\text{eff}}$
$\sqrt{\langle r_J^2 \rangle}$	$0.70^{+0.02}_{-0.02}$	$+(0.01)_{\text{ChPT}}(0.01)_{\text{PWA}}(0.01)_{\text{eff}}$ $-(0.01)_{\text{ChPT}}(0.00)_{\text{PWA}}(0.02)_{\text{eff}}$

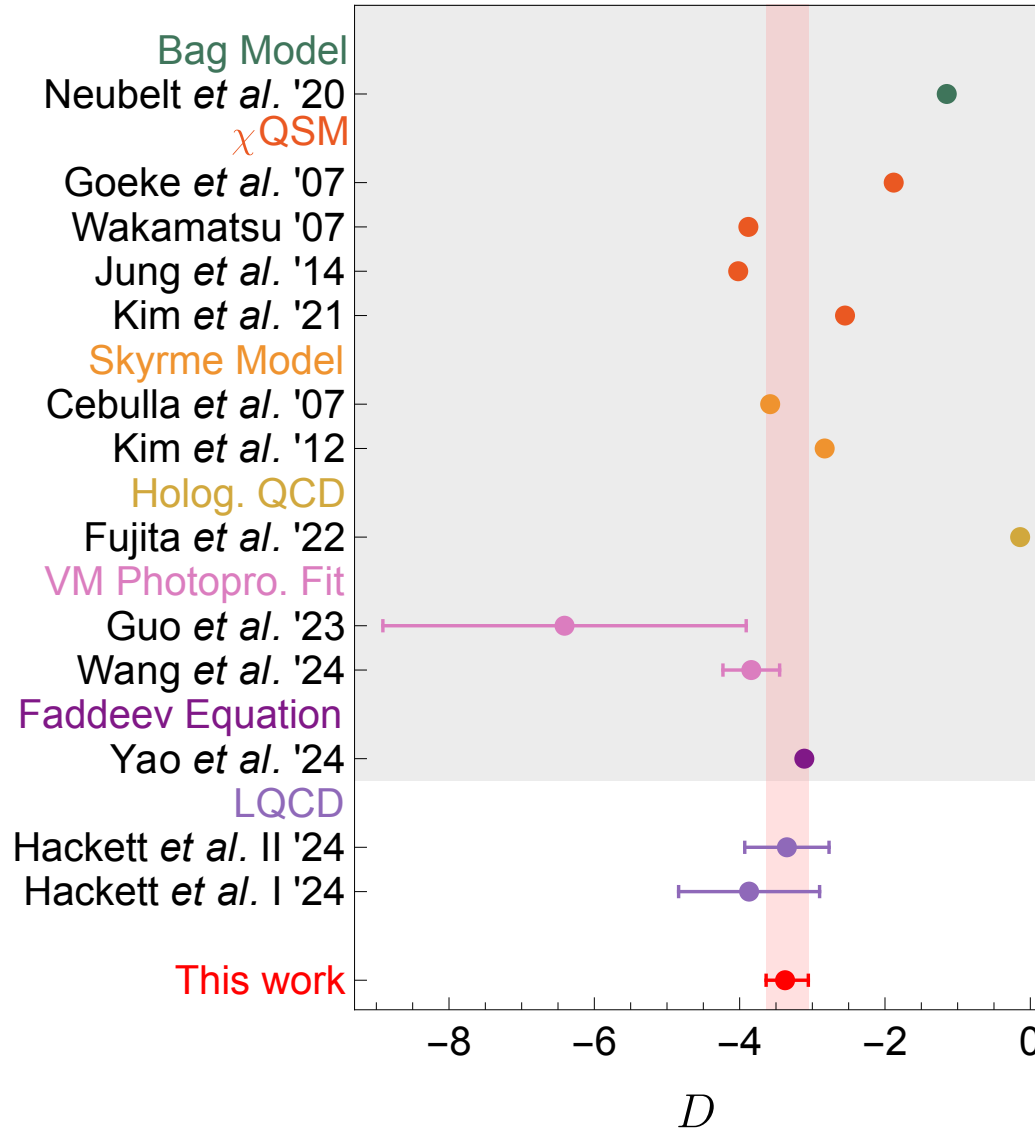
ChPT: NLO ChPT inputs
pwa: $\pi\pi/K\bar{K} \rightarrow N\bar{N}$
eff: effective poles m_S, m_D

Nucleon D-term



- Nucleon D-term $D \equiv D(0)$:

$$D = - \left(3.38^{+0.34}_{-0.35} \right)$$



- Positivity bound:

$$D < -0.2$$

Gegelia & Polyakov, PLB (2021)

“The stability condition”

$$\text{normal force } \frac{2}{3}s(r) + p(r) \geq 0$$

$$\widetilde{D}(r) = \int \frac{d^3\Delta}{(2\pi)^3} e^{-i\Delta r} D(-\Delta^2) < 0$$

Comparison with other works

