

Long-lived particle searches at FASER and future machines

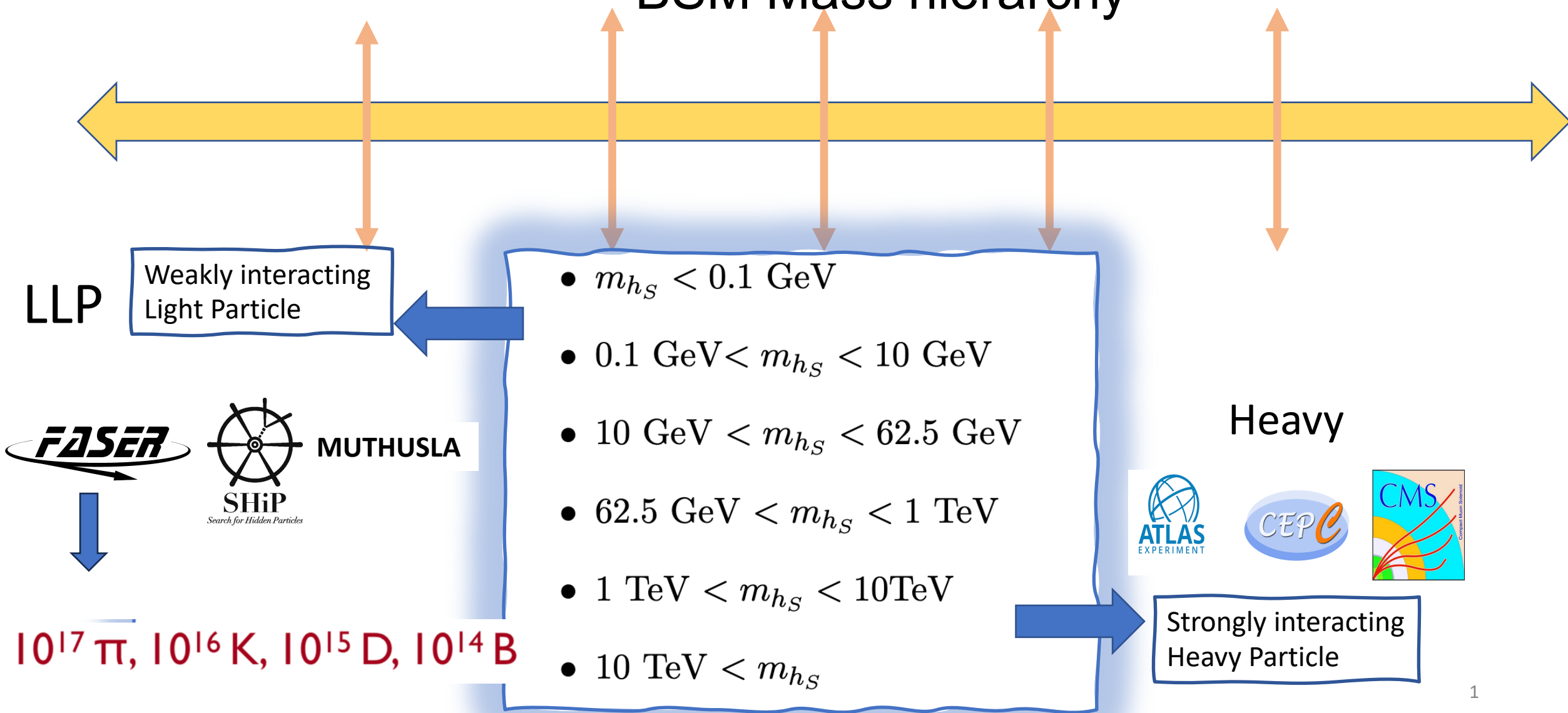
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NPhIS 2026 Nanjing

[2508.12309](#) (XY Qi, HY Song, WS)
[2512.24376](#) (S.Su, WS, JM. Yang, PX. Zhu, R. Zhu)
2606.XXXXX (YZ. Shi, WS, etc.)

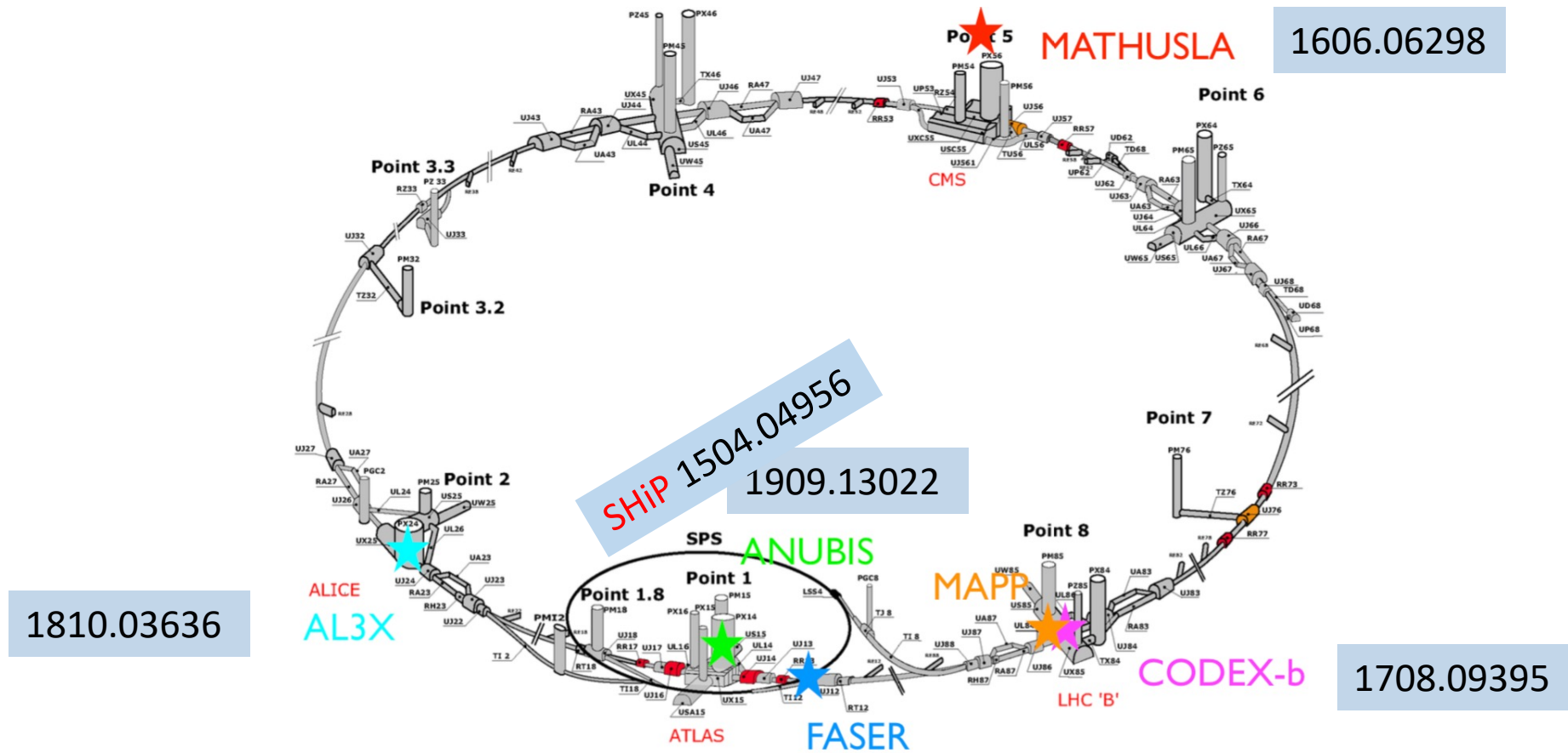


Motivation: LLP

BSM Mass hierarchy



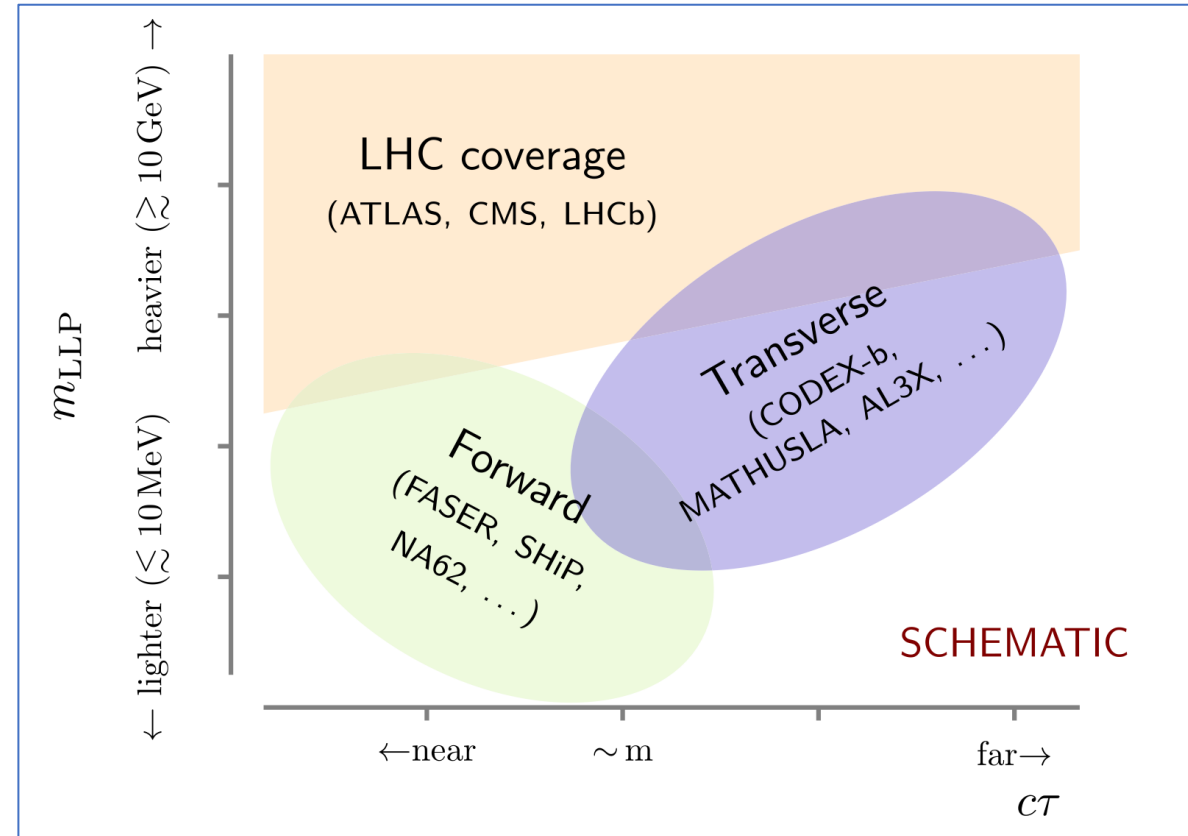
Proposals for LLP



Light LLP searches

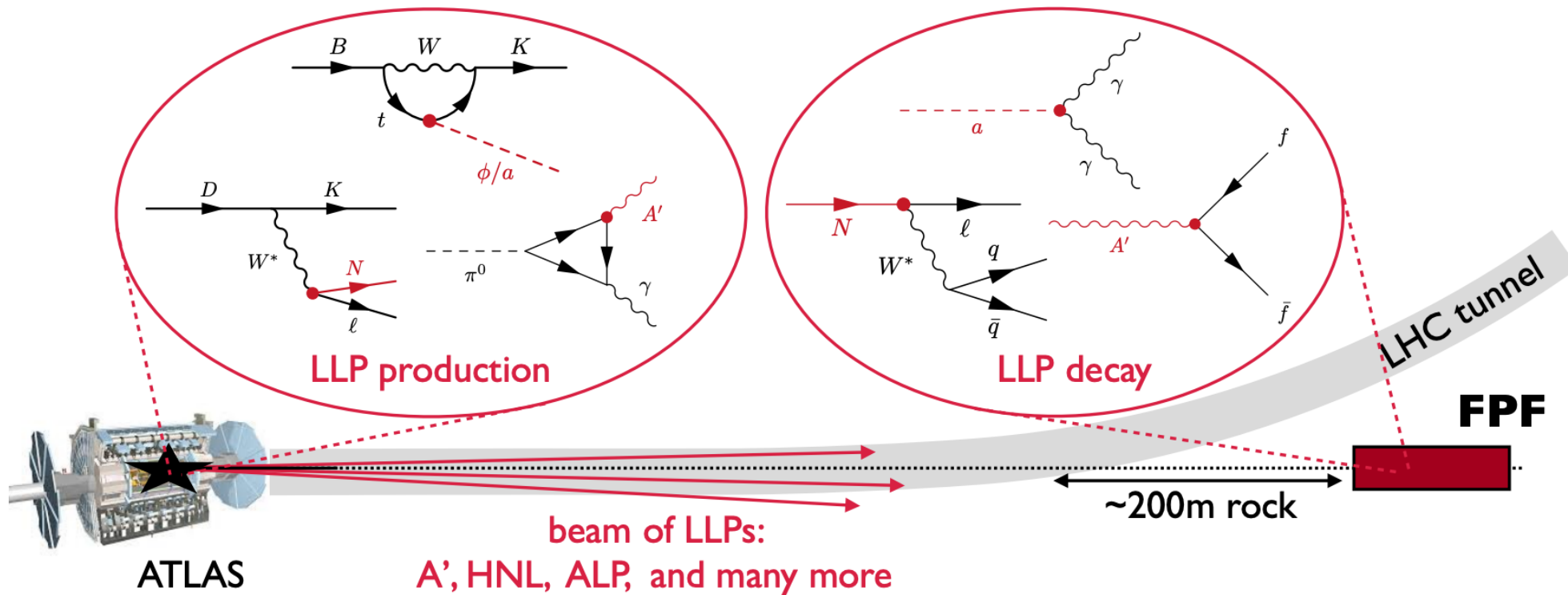
Eur. Phys. J. C (2020) 80:1177

1. ATLAS and CMS: Heavy LLPs ($m_{\text{LLP}} \gtrsim 10 \text{ GeV}$) for all lifetimes ($c\tau \lesssim 10^7 \text{ m}$).
2. LHCb: Short to medium lifetimes ($c\tau \lesssim 1 \text{ m}$) for light LLPs ($0.1 \text{ GeV} \lesssim m_{\text{LLP}} \lesssim 10 \text{ GeV}$).
3. Forward/beam dump detectors (FASER, NA62, SHiP): Medium to long lifetime regime ($0.1 \lesssim c\tau \lesssim 10^7 \text{ m}$) for light LLPs ($m_{\text{LLP}} \lesssim \text{few GeV}$), for low $\sqrt{s}$ production channels.
4. Shielded, transversely displaced detectors (MATHUSLA, CODEX-b, AL3X): Relatively light LLPs² ($m_{\text{LLP}} \lesssim 10\text{--}100 \text{ GeV}$) in the long lifetime regime ($1 \lesssim c\tau \lesssim 10^7 \text{ m}$), and high $\sqrt{s}$ production channels.



Eg: Production at FASER

many hadrons: 10^{17} π , 10^{16} K , 10^{15} D , 10^{14} B with $E \sim \text{TeV}$



$pp \rightarrow \text{LLP} + X$, LLP travels ~ 480 m, $\text{LLP} \rightarrow \text{charged tracks} + X$

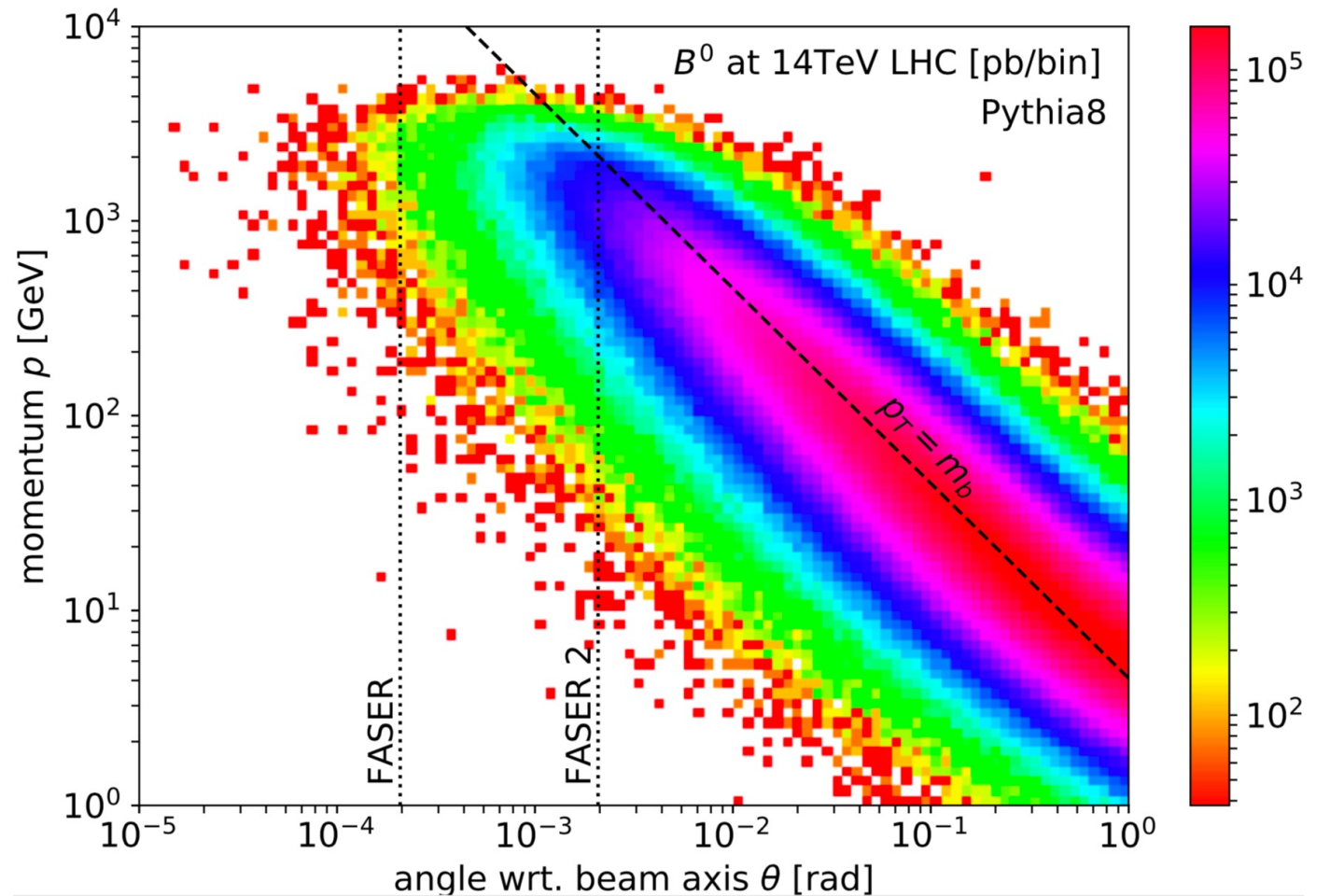
Production

B^0

boosted mesons highly collimated

$$p \cdot \theta = p_T \sim \text{Mb}$$

FASER: radius $R = 10$ cm,
length $D = 1.5$ m,
FASER 2: radius $R = 1$ m,
length $D = 5$ m.



Production: CP even scalar

Main contribution

$$K^\pm: 493.677 \pm 0.016 \text{ MeV/}$$

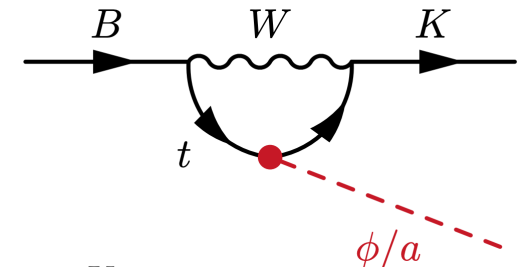
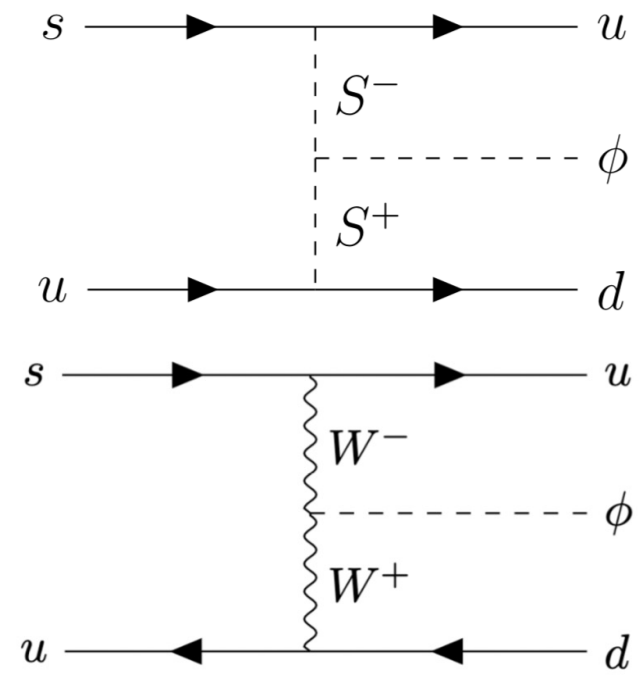
$$K^0: 497.611 \pm 0.013 \text{ MeV/}$$

$$\text{Br}(K^\pm \rightarrow \pi^\pm \phi) = \frac{1}{\Gamma_{K^\pm}} \frac{2p_\phi^0}{m_{K^\pm}} \frac{|\mathcal{M}|^2}{16\pi m_{K^\pm}},$$

$$\mathcal{M}(K^\pm \rightarrow \pi^\pm \phi) = G_F^{1/2} 2^{1/4} \xi_\phi^W \left[\frac{7\lambda(m_{K^\pm}^2 + m_{\pi^\pm}^2 - m_\phi^2)}{18} - \frac{7Am_{K^\pm}^2}{9} \right] + \frac{\xi_\phi^{ds}}{2v} m_s \frac{m_{K^\pm}^2 - m_{\pi^\pm}^2}{m_s - m_d} f_0^{K^\pm \pi^\pm}(q^2)$$

$$\frac{\text{Br}(B \rightarrow X_s \phi)}{\text{Br}(B \rightarrow X_c e \nu)} = \frac{\Gamma(b \rightarrow s \phi)}{\Gamma(b \rightarrow c e \nu)} = \frac{12\pi^2 v^2}{m_b^2} \left(1 - \frac{m_\phi^2}{m_b^2}\right)^2 \frac{1}{f(m_c^2/m_b^2)} \left| \frac{\xi_\phi^{bs}}{V_{cb}} \right|^2$$

b: 4.18 GeV, B: around 5.3 GeV,



Decay: CP even scalar

$$\begin{aligned} H &\rightarrow \pi\pi \\ H &\rightarrow KK \\ H &\rightarrow \pi\pi\pi\pi \end{aligned}$$

Hadronic decays into pions and kaons for

$$m_\phi < 2 \text{ GeV}$$

$$\begin{aligned} \mathcal{L} &\supset \frac{\Phi}{v} \left(\xi_\Phi^g \frac{\alpha_s}{12\pi} G_{\mu\nu}^a G^{a\mu\nu} - \xi_\Phi^u m_u \bar{u}u - \xi_\Phi^d m_d \bar{d}d - \xi_\Phi^s m_s \bar{s}s \right) \\ &= -\frac{\Phi}{v} \left\{ \xi_\Phi^g \left[\frac{2}{27} \Theta_\mu^\mu - \frac{2}{27} (m_u \bar{u}u + m_d \bar{d}d + m_s \bar{s}s) \right] + (\xi_\Phi^u m_u \bar{u}u + \xi_\Phi^d m_d \bar{d}d + \xi_\Phi^s m_s \bar{s}s) \right\} \end{aligned}$$

$$\Theta_\mu^\mu = -\frac{9\alpha_s}{8\pi} G_{\mu\nu}^a G^{a\mu\nu} + m_u \bar{u}u + m_d \bar{d}d + m_s \bar{s}s.$$

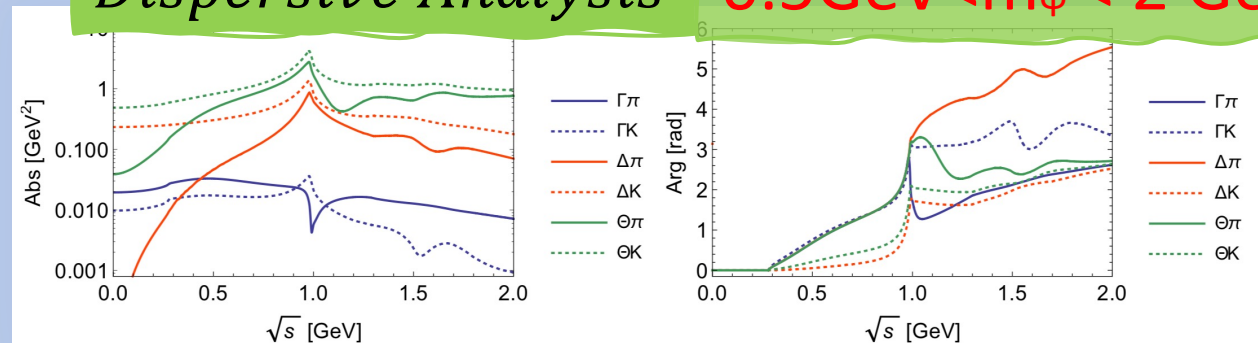
$$\Gamma_\pi = \langle \pi\pi | m_u \bar{u}u + m_d \bar{d}d | 0 \rangle, \quad \Delta_\pi = \langle \pi\pi | m_s \bar{s}s | 0 \rangle, \quad \Theta_\pi = \langle \pi\pi | \Theta_\mu^\mu | 0 \rangle$$

$$m_\phi < 0.5 \text{ GeV}$$

Leading order chiral
perturbation theory

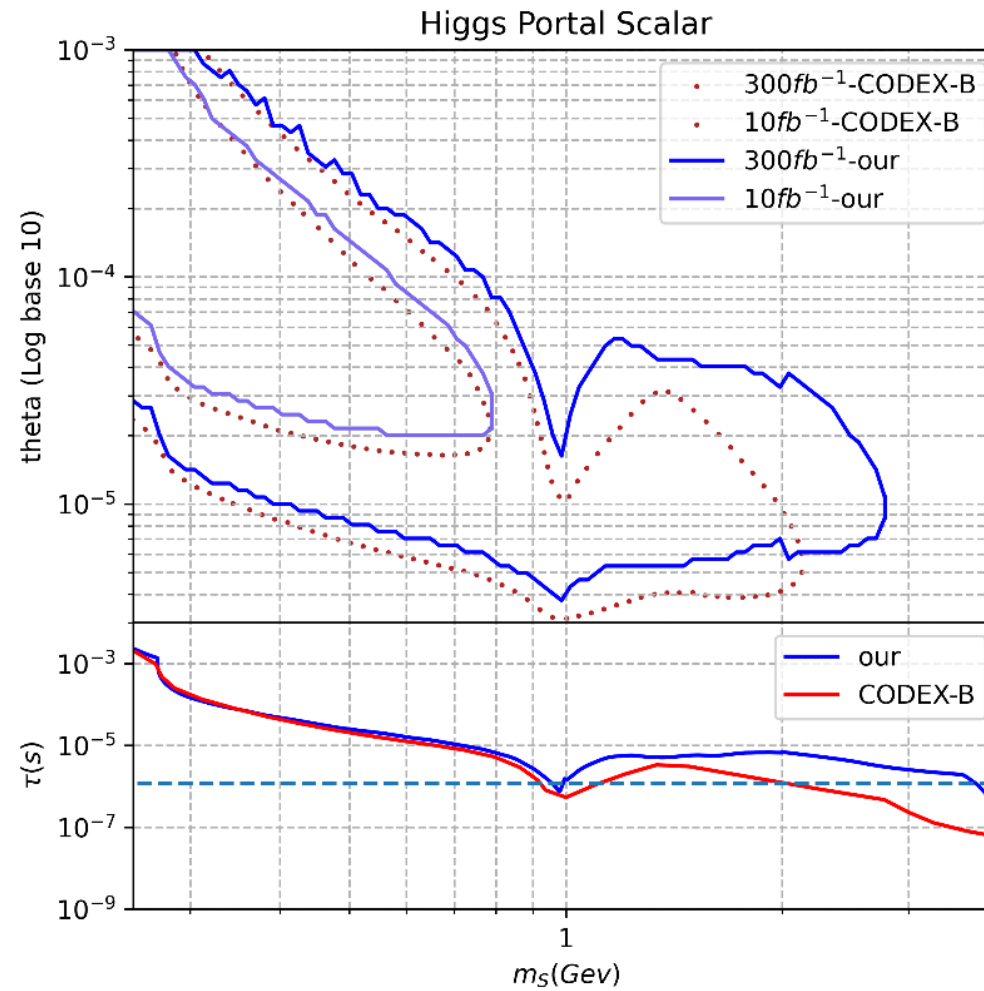
Dispersive Analysis

$$0.5 \text{ GeV} < m_\phi < 2 \text{ GeV}$$



Comparison

- SM+singlet

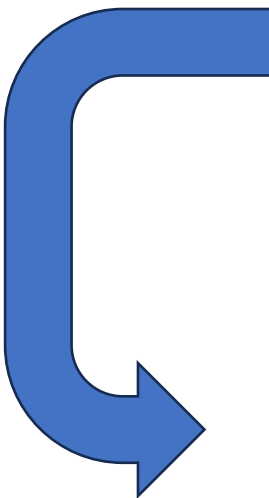


https://github.com/shiggs90/Light_scalar_decay.git.

Complete Light Higgs

All constraints @ Type-I 2HDM

$$\cos(\beta - \alpha) \simeq \frac{1}{\tan \beta}, \quad m_A \in (54, 600) \text{ GeV}, \quad m_{H^\pm} \sim m_A, \quad \lambda v^2 = 0.$$

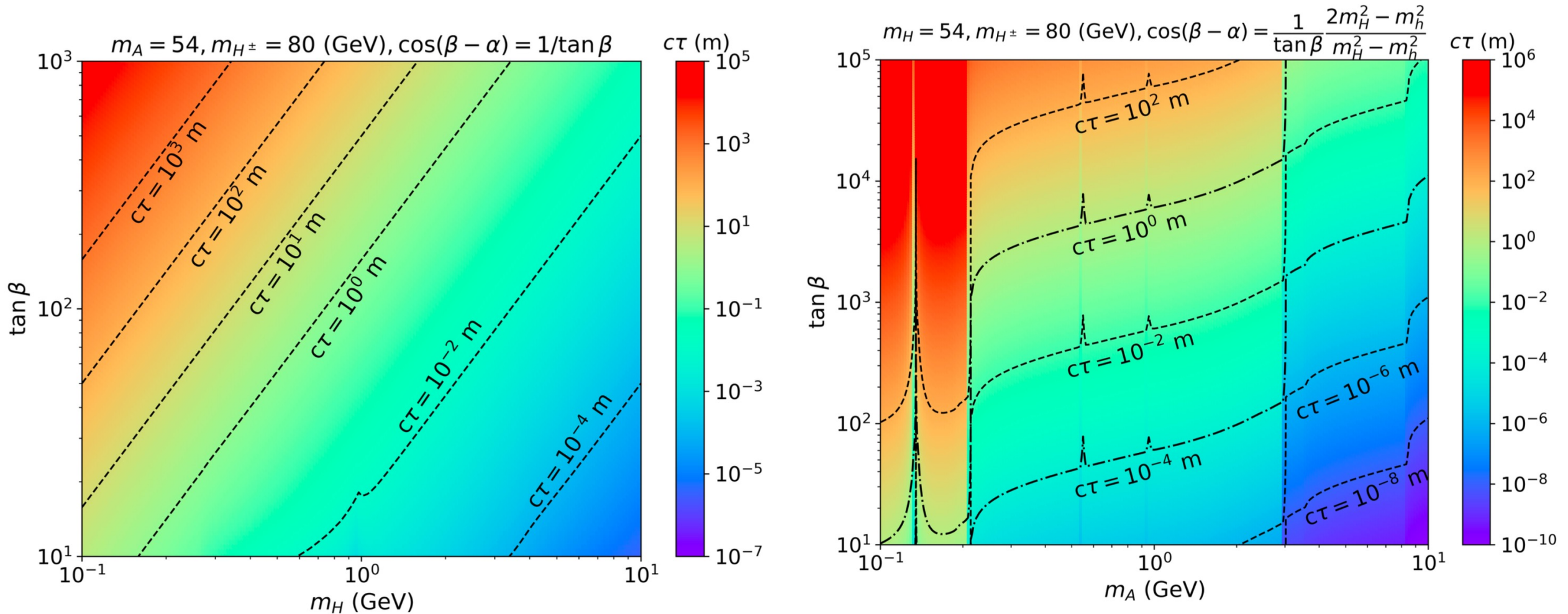

$$\cos(\beta - \alpha) \simeq \frac{1}{\tan \beta} \frac{2m_H^2 - m_h^2}{m_H^2 - m_h^2}, \quad m_H \in (54, m_h) \text{ GeV}, \quad m_{H^\pm} \sim m_A, \quad \lambda v^2 = 0.$$

$$\cos(\beta - \alpha) = 0, \quad m_H \in (86.9, 90) \text{ GeV}, \quad m_{H^\pm} \sim m_H, \quad \lambda v^2 = 0,$$

$$\cos(\beta - \alpha) = -\frac{1}{\tan \beta}, \quad m_H \in (101.2, 103) \text{ GeV}, \quad m_{H^\pm} \sim m_H, \quad \lambda v^2 = 0.$$

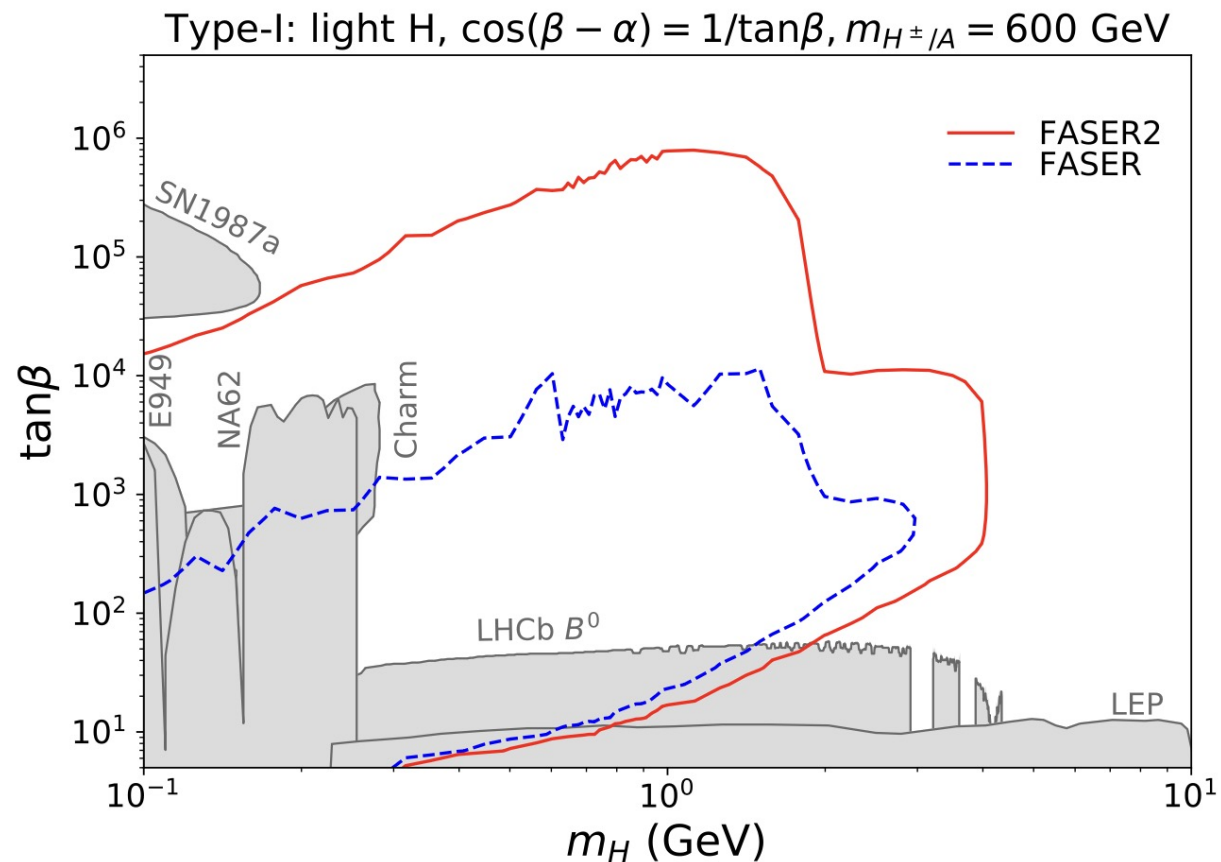
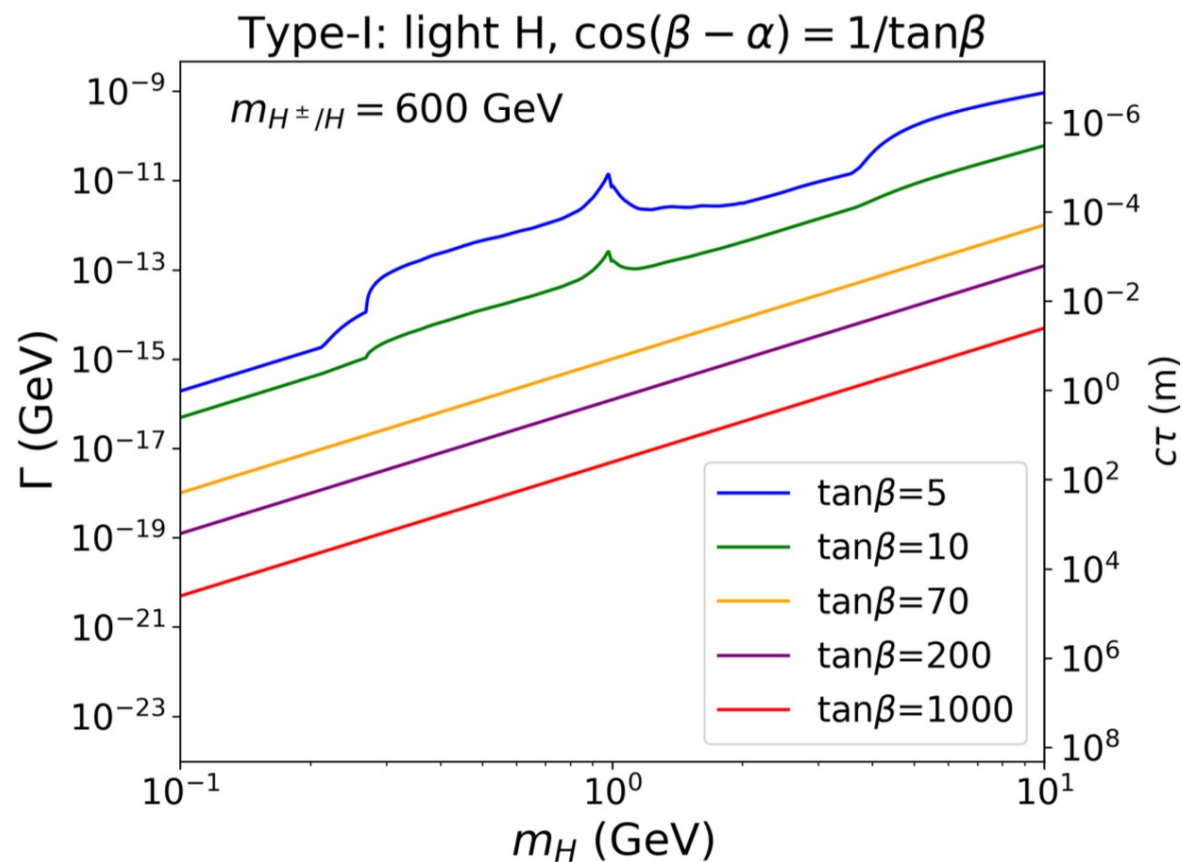
Complete Light LLP

@ Type-I 2HDM



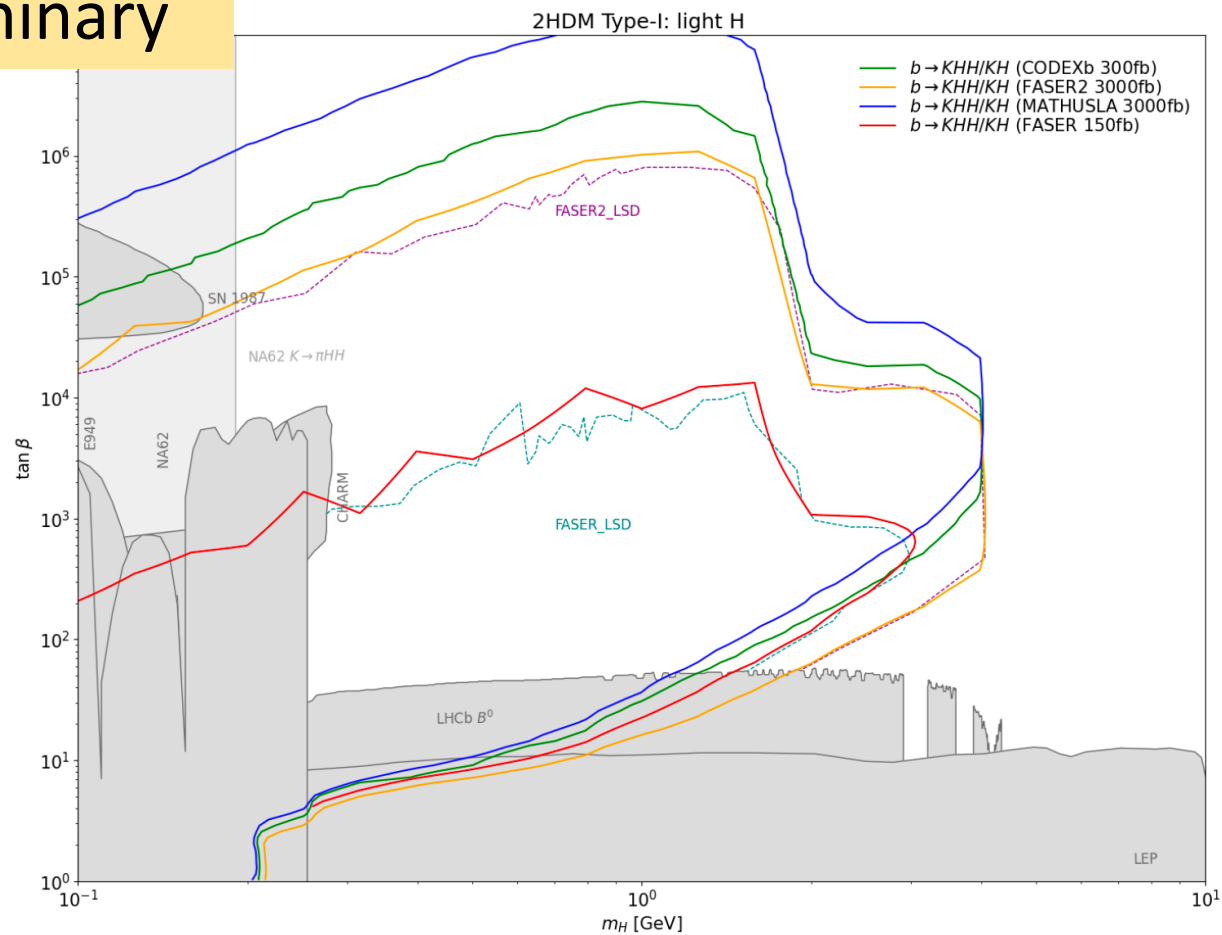
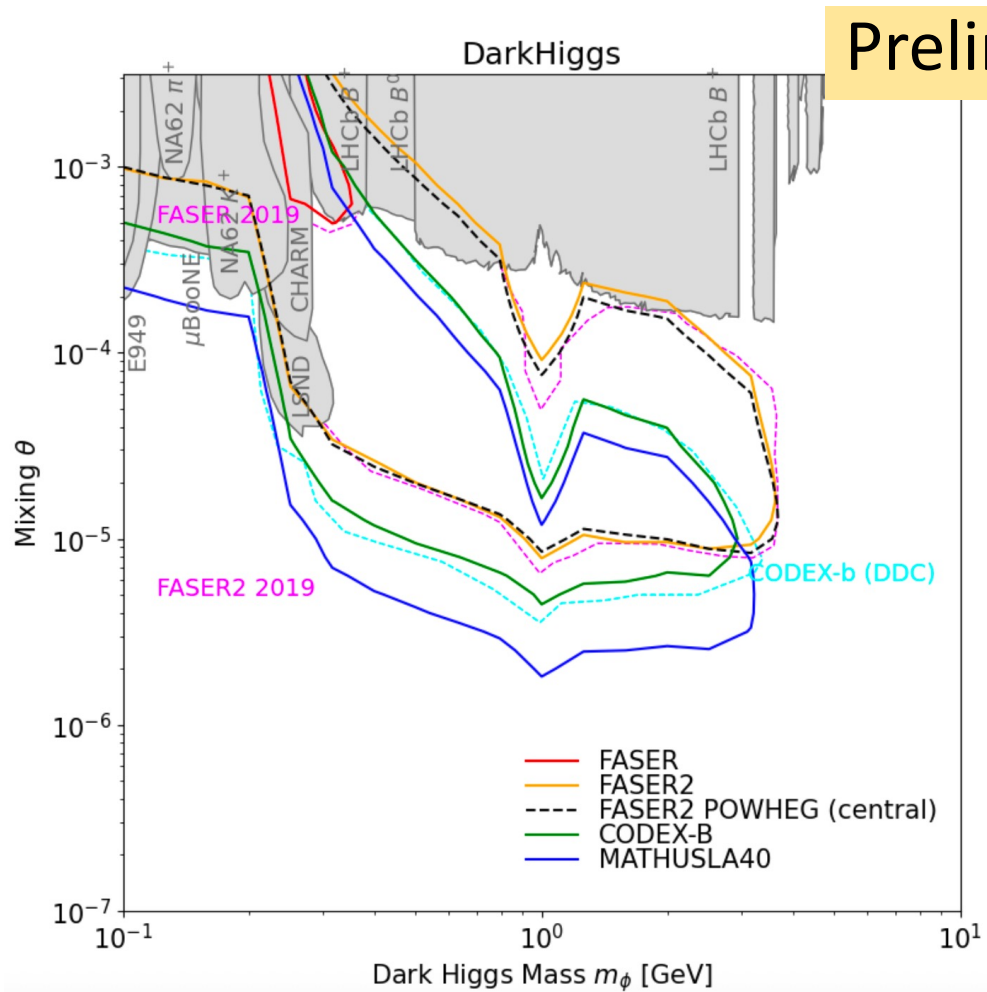
Results: CP even BM

@ Type-I 2HDM

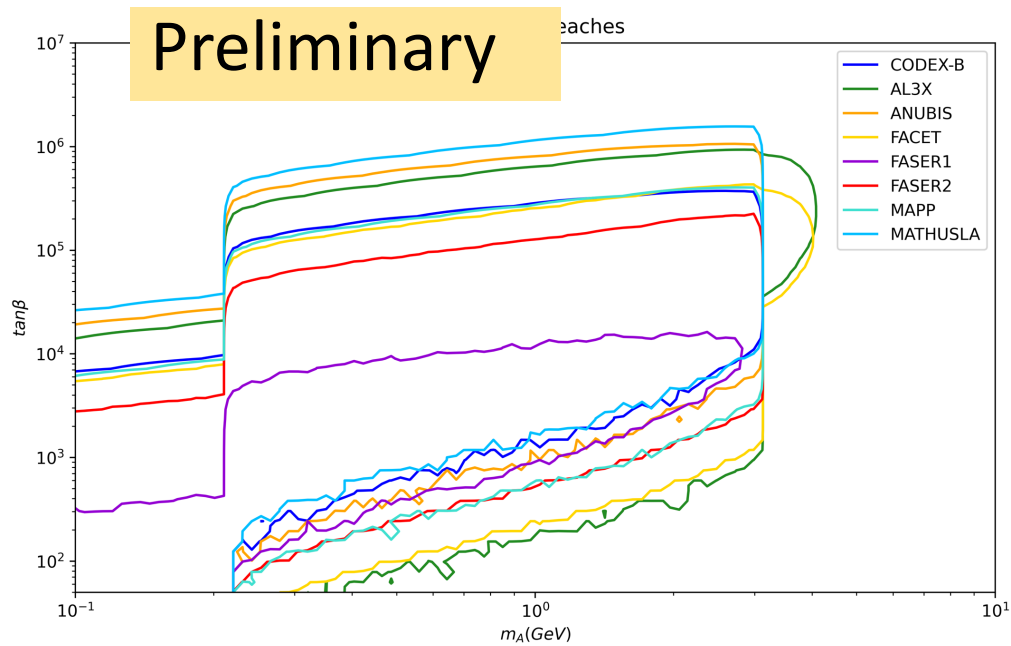


Future machines

Experiment	Geometry	Luminosity	Energy	Color in Figure
MATHUSLA40	Box Max=[20, 93, 110], Box Min=[-20, 81, 70]	$3\ ab^{-1}$	14TeV	Blue
CODEX-b	Box Max=[36, 3, 15], Box Min=[26, -7, 5]	$300\ fb^{-1}$	14TeV	Green
FASER	Cylinder L=480m, h=1.5m, r=0.1m	$150\ fb^{-1}$	14TeV	Red
FASER2	Cylinder L=480m, h=5m, r=1m	$3\ ab^{-1}$	14TeV	Orange



Case study

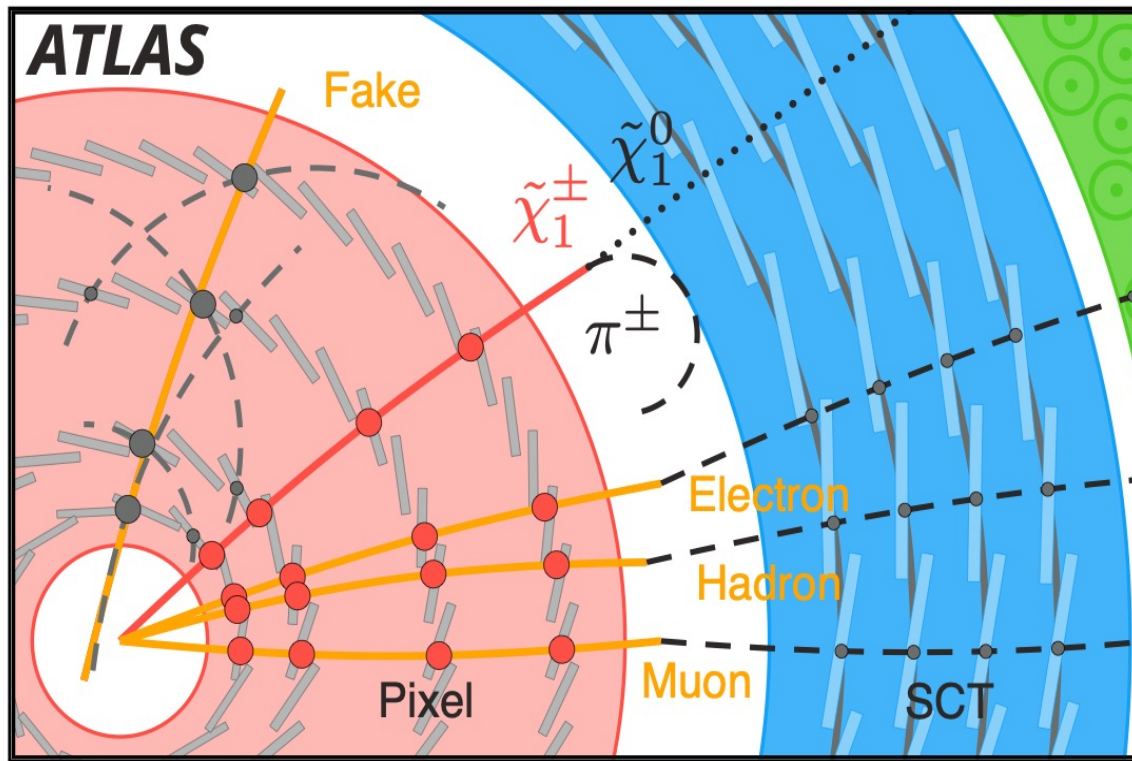


higher luminosity
helps to reach the
weaker coupling
region.

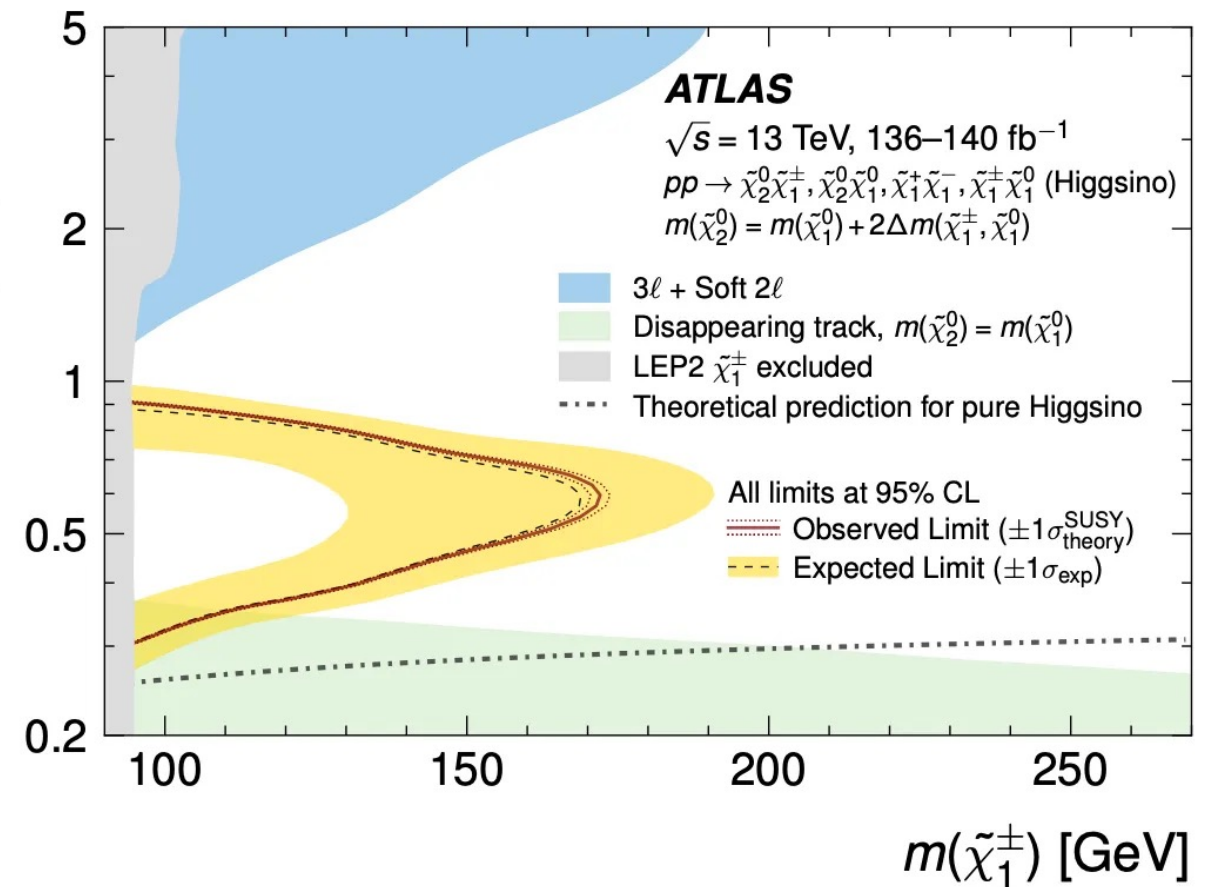
A larger detector,
especially the radius
helps to extend the
reach in m_A .

Q: possible to go beyond M_b @ FASER?

eg: Long-lived charginos $M_{\text{LLP}} > M_b$

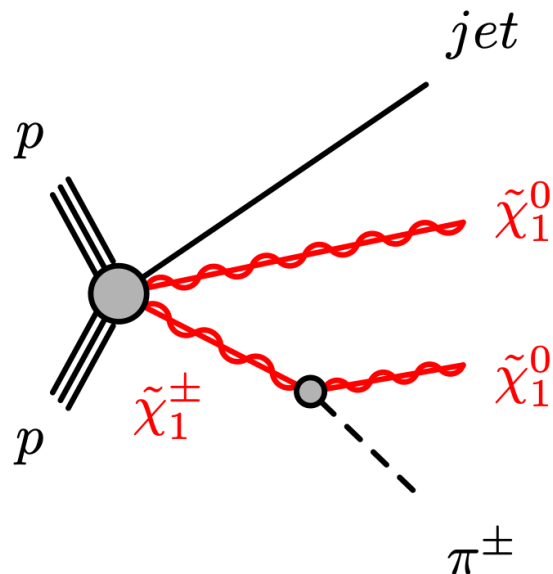


PhysRevLett.132.221801



$$M_{\text{LLP}} > M_b$$

Both the two methods,
disappearing and displaced
tracks, are focused on
charged LLP



some assumptions

Disappearing track, $m(\tilde{\chi}_2^0) = m(\tilde{\chi}_1^0)$

Displaced track :

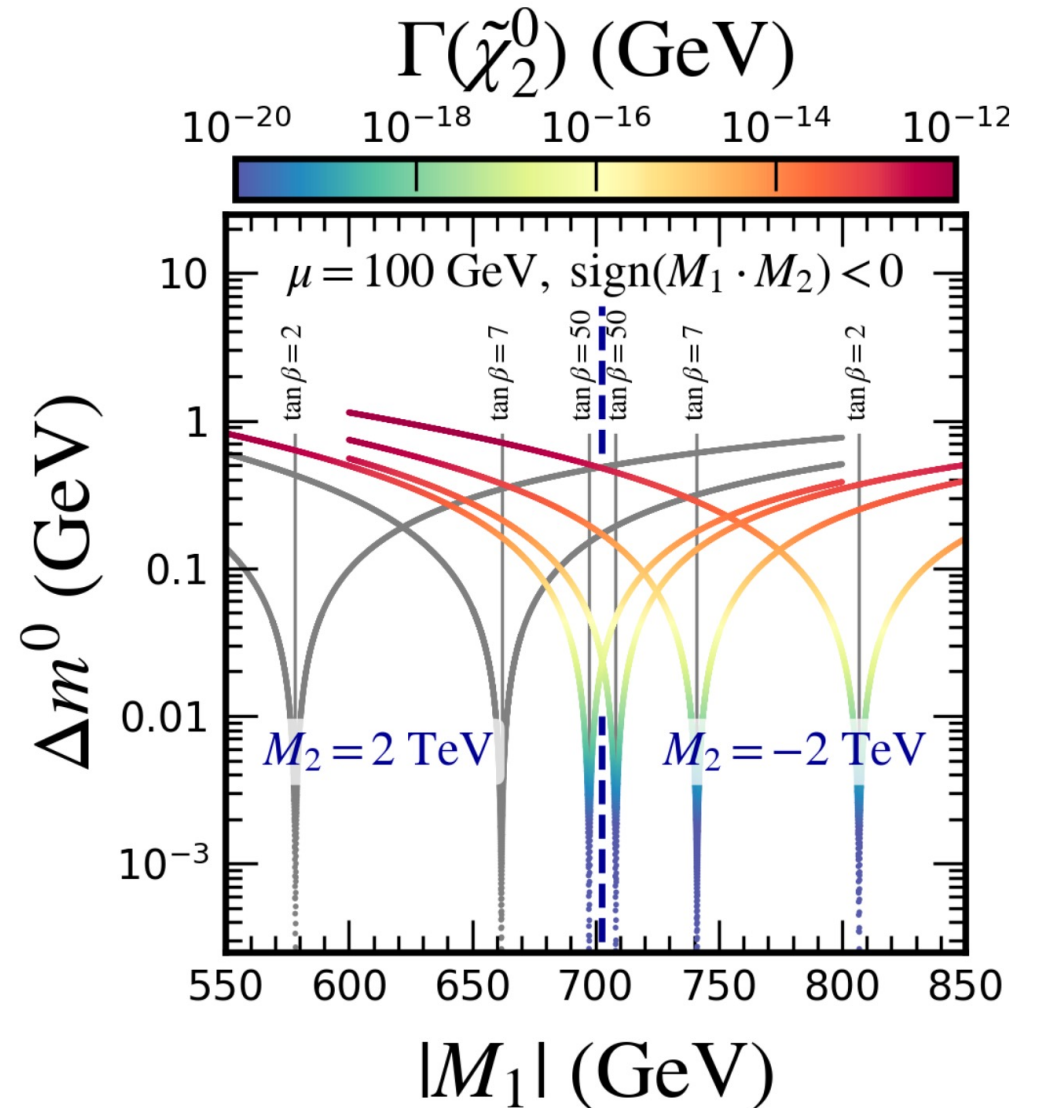
$$m(\tilde{\chi}_2^0) = m(\tilde{\chi}_1^0) + 2\Delta m(\tilde{\chi}_1^\pm, \tilde{\chi}_1^0)$$

$M_{\text{LLP}} > M_b$: Long-lived Higgsino

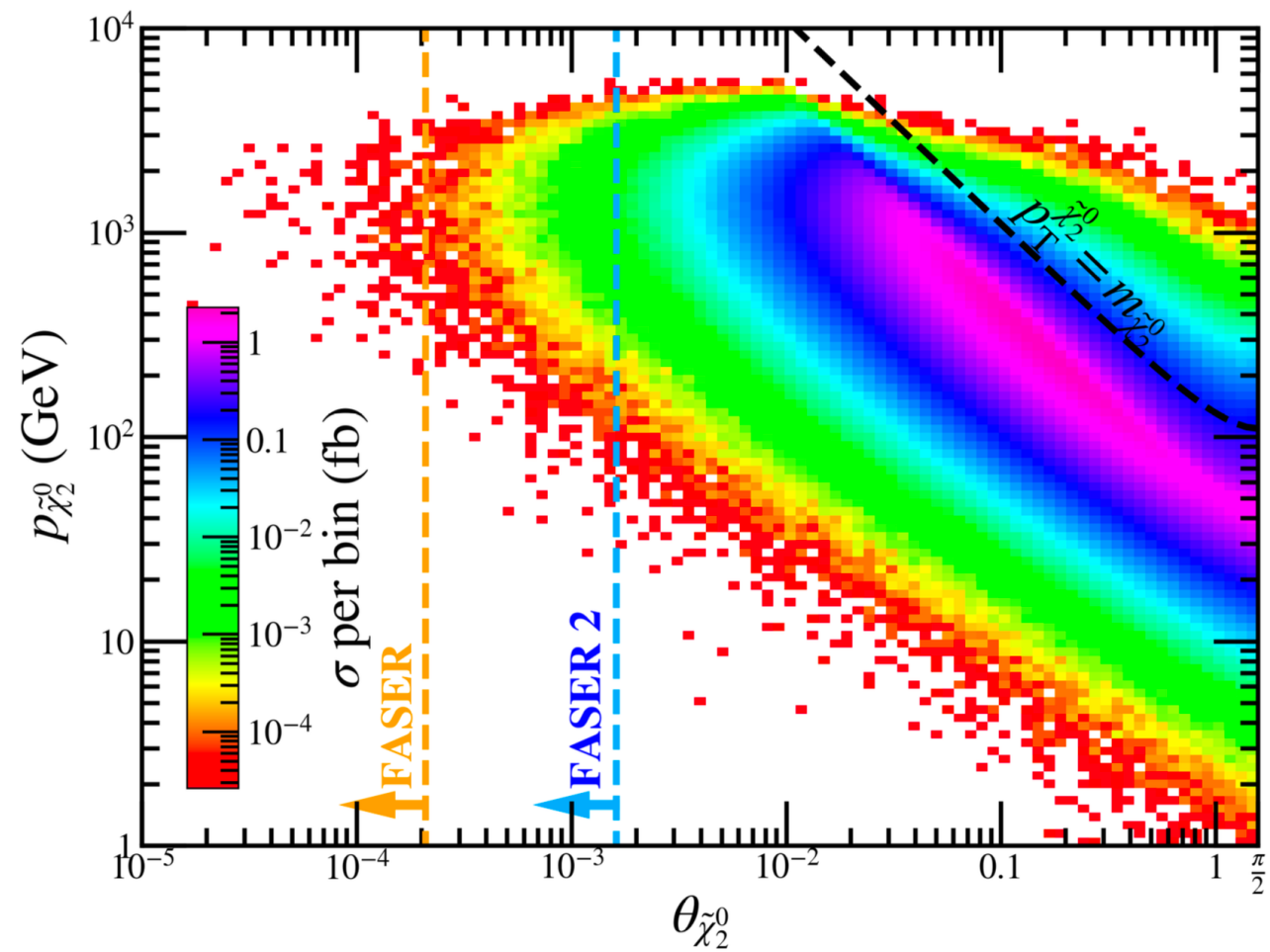
Higgsino-like neutralino @ MSSM

$$\begin{aligned}\tilde{\chi}_2^0 &\rightarrow \tilde{\chi}_1^0 + \gamma, \\ &\rightarrow \tilde{\chi}_1^0 + Z^* \rightarrow \tilde{\chi}_1^0 + f\bar{f}.\end{aligned}$$

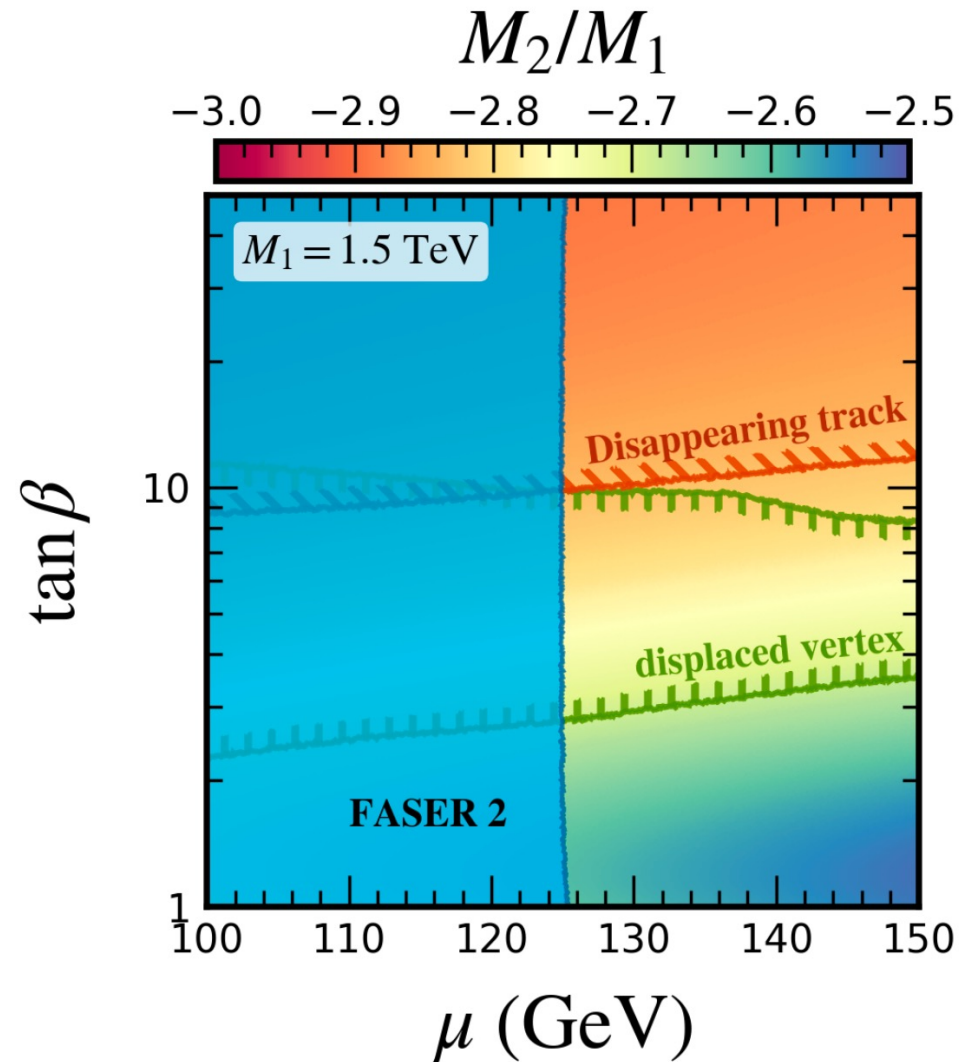
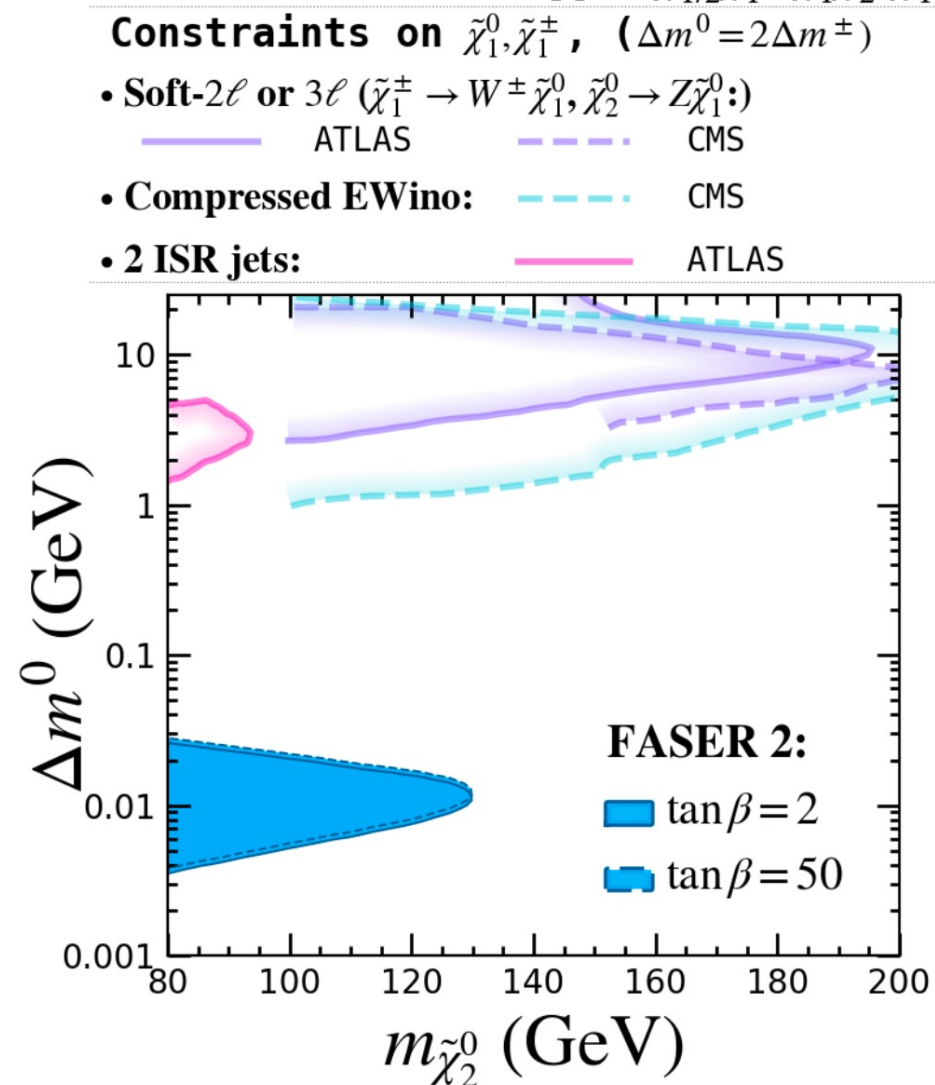
$$\Gamma(\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 \gamma) = \frac{g_{\tilde{\chi}_2^0 \tilde{\chi}_1^0 \gamma}^2}{8\pi} \frac{(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{\chi}_1^0}^2)^3}{m_{\tilde{\chi}_2^0}^5},$$



Production



FASER for $\mathcal{O}(100)$ GeV LLP



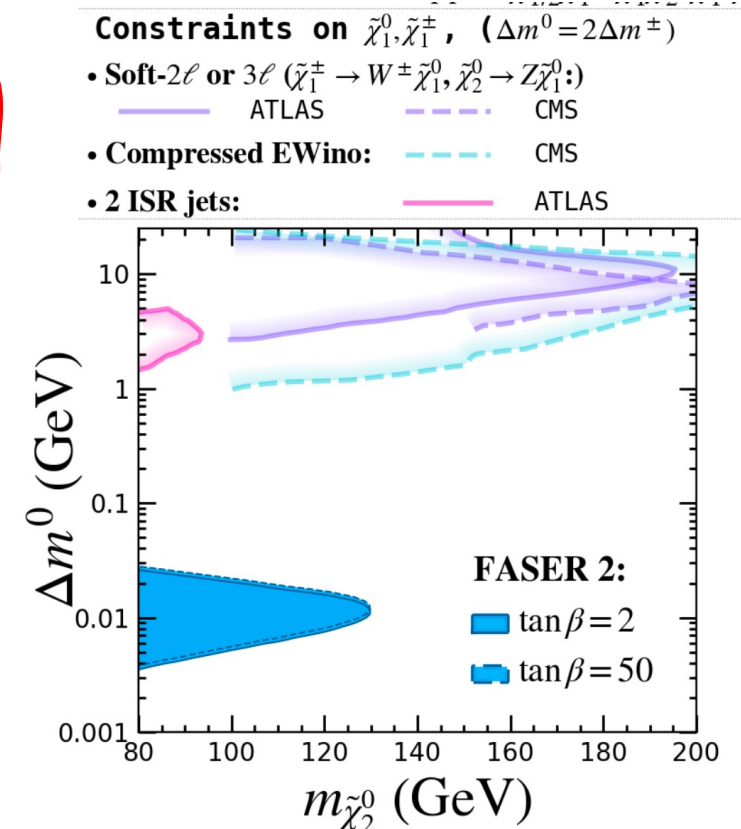
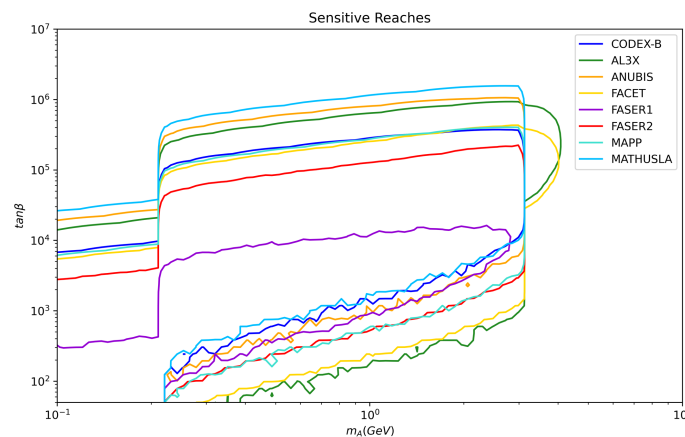
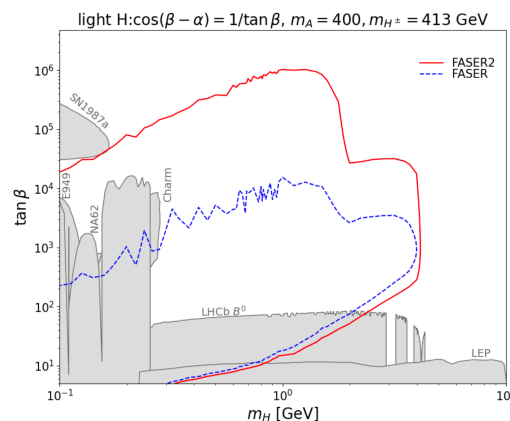
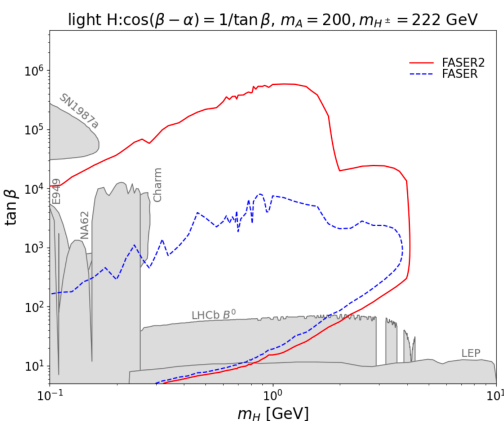
Conclusion

- Light scalar decay: GeV scale
- Complete LLP at Type-I 2HDM
- Various facilities
- $M_{\text{LLP}} > M_B$

https://github.com/shiggs90/Light_scalar_decay.git.

$$\cos(\beta - \alpha) \simeq \frac{1}{\tan \beta}, \quad m_A \in (54, 600) \text{ GeV}, \quad m_{H^\pm} \sim m_A, \quad \lambda v^2 = 0.$$

$$\cos(\beta - \alpha) \simeq \frac{1}{\tan \beta} \frac{2m_H^2 - m_h^2}{m_H^2 - m_h^2}, \quad m_H \in (54, m_h) \text{ GeV}, \quad m_{H^\pm} \sim m_A, \quad \lambda v^2 = 0.$$



Thanks !

Case study: 2HDM

	ξ_H^u	ξ_H^d	ξ_H^ℓ	ξ_A^u	ξ_A^d	ξ_A^ℓ
Type-I	$\cot \beta$	$\cot \beta$	$\cot \beta$	$\cot \beta$	$-\cot \beta$	$-\cot \beta$
Type-II	$\cot \beta$	$-\tan \beta$	$-\tan \beta$	$\cot \beta$	$\tan \beta$	$\tan \beta$
Type-L	$\cot \beta$	$\cot \beta$	$-\tan \beta$	$\cot \beta$	$-\cot \beta$	$\tan \beta$
Type-F	$\cot \beta$	$-\tan \beta$	$\cot \beta$	$\cot \beta$	$\tan \beta$	$-\cot \beta$

● Two Higgs Doublet Model

$$\Phi_i = \begin{pmatrix} \phi_i^+ \\ (v_i + \phi_i^0 + iG_i)/\sqrt{2} \end{pmatrix}$$

$$v_u^2 + v_d^2 = v^2 = (246\text{GeV})^2$$

$$\tan \beta = v_u/v_d$$

	ϕ_1	ϕ_2
Type I	u,d,l	
Type II	u	d,l
lepton-specific	u,d	l
flipped	u,l	d

$$\begin{pmatrix} H^0 \\ h^0 \end{pmatrix} = \begin{pmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{pmatrix} \begin{pmatrix} \phi_1^0 \\ \phi_2^0 \end{pmatrix}, \quad \begin{aligned} A &= -G_1 \sin \beta + G_2 \cos \beta \\ H^\pm &= -\phi_1^\pm \sin \beta + \phi_2^\pm \cos \beta \end{aligned}$$

● Parameters (CP-conserving, Flavor Limit, Z_2 Symmetry)

$$m_{11}^2, m_{22}^2, \lambda_1, \lambda_2, \lambda_3, \lambda_4, \lambda_5$$



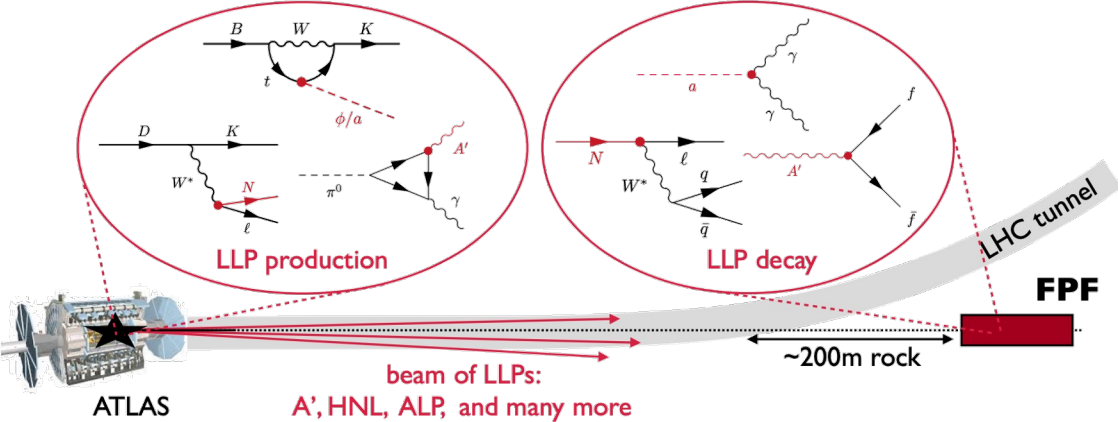
$$v, \tan \beta, \alpha, m_h, m_H, m_A, m_{H^\pm}$$

Soft Z_2 symmetry breaking: m_{12}^2

246 GeV

125. GeV

Decay



Well studied

?

Scale-ind

Scale > 2/3 GeV

$A \rightarrow \gamma\gamma$	$H \rightarrow \gamma\gamma$
$A \rightarrow e^+e^-$	$H \rightarrow e^+e^-$
$A \rightarrow \mu^+\mu^-$	$H \rightarrow \mu^+\mu^-$
$A \rightarrow \tau^+\tau^-$	$H \rightarrow \tau^+\tau^-$
$A \rightarrow q\bar{q}$	$H \rightarrow c\bar{c}$
$A \rightarrow gg$	$H \rightarrow s\bar{s}$
	$H \rightarrow gg$

$H \rightarrow \pi\pi$
$H \rightarrow KK$
$H \rightarrow \pi\pi\pi\pi$

Chiral Perturbativity...

arXiv:1809.01876
arXiv:1612.06538

- $A \rightarrow \pi\pi\pi$
- $A \rightarrow \eta\pi\pi$
- $A \rightarrow \eta'\pi\pi$
- $A \rightarrow \eta\eta\pi$
- $A \rightarrow KK\pi$
- $A \rightarrow \gamma\pi\pi$
- $A \rightarrow \eta\eta'\pi$
- $A \rightarrow \eta'\eta'\pi$
- $A \rightarrow \eta\eta\eta$
- $A \rightarrow \eta\eta\eta'$
- $A \rightarrow \eta\eta'\eta'$
- $A \rightarrow \eta'\eta'\eta'$
- $A \rightarrow \eta KK$
- $A \rightarrow \eta' KK$

Production: CP even scalar

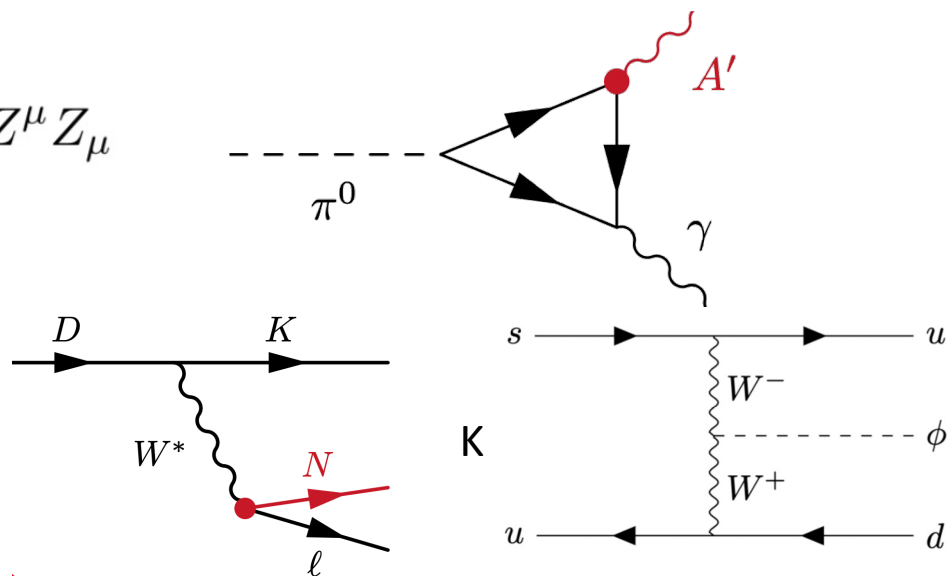
$$\xi_\phi^g = \sum_{f \in q} \frac{3}{2} \xi_\phi^i \mathcal{A}_{1/2}^\phi(\tau_f^\phi),$$

$$\xi_\phi^\gamma = \sum_{f \in q, \ell} N_c Q_f^2 \xi_\phi^f \mathcal{A}_{1/2}^\phi(\tau_f^\phi) + \xi_\phi^V \mathcal{A}_1^\phi(\tau_W^\phi)$$

$$\mathcal{L} = -\frac{1}{2} m_\phi^2 \phi^2 - \sum_f \xi_\phi^f \frac{m_f}{v} \phi \bar{f} f + \xi_\phi^W \frac{2m_W^2}{v} \phi W^{\mu+} W_\mu^- + \xi_\phi^Z \frac{m_Z^2}{v} \phi Z^\mu Z_\mu \\ + \xi_\phi^g \frac{\alpha_s}{12\pi v} \phi G_{\mu\nu}^a G^{a\mu\nu} + \xi_\phi^\gamma \frac{\alpha}{4\pi v} \phi F_{\mu\nu} F^{\mu\nu}.$$

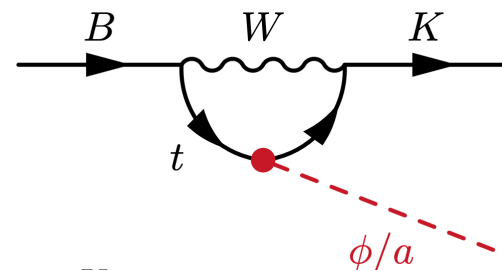
$$K \rightarrow \pi\phi, \eta^{(\prime)} \rightarrow \pi\phi, D \rightarrow X_u\phi, B \rightarrow X_s\phi$$

$$\pi^+ \rightarrow \ell\nu\phi \quad K^+ \rightarrow \ell\nu\phi \quad \Upsilon \rightarrow \phi\gamma$$



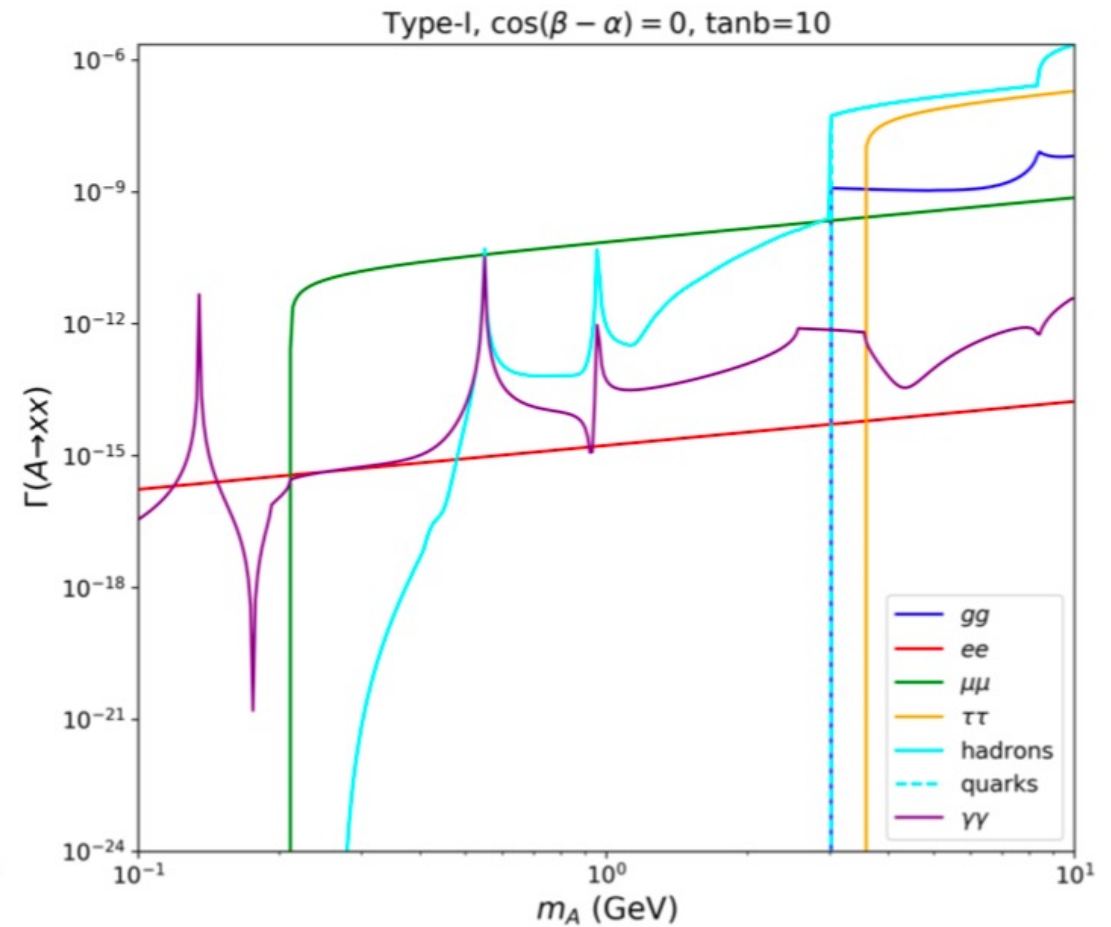
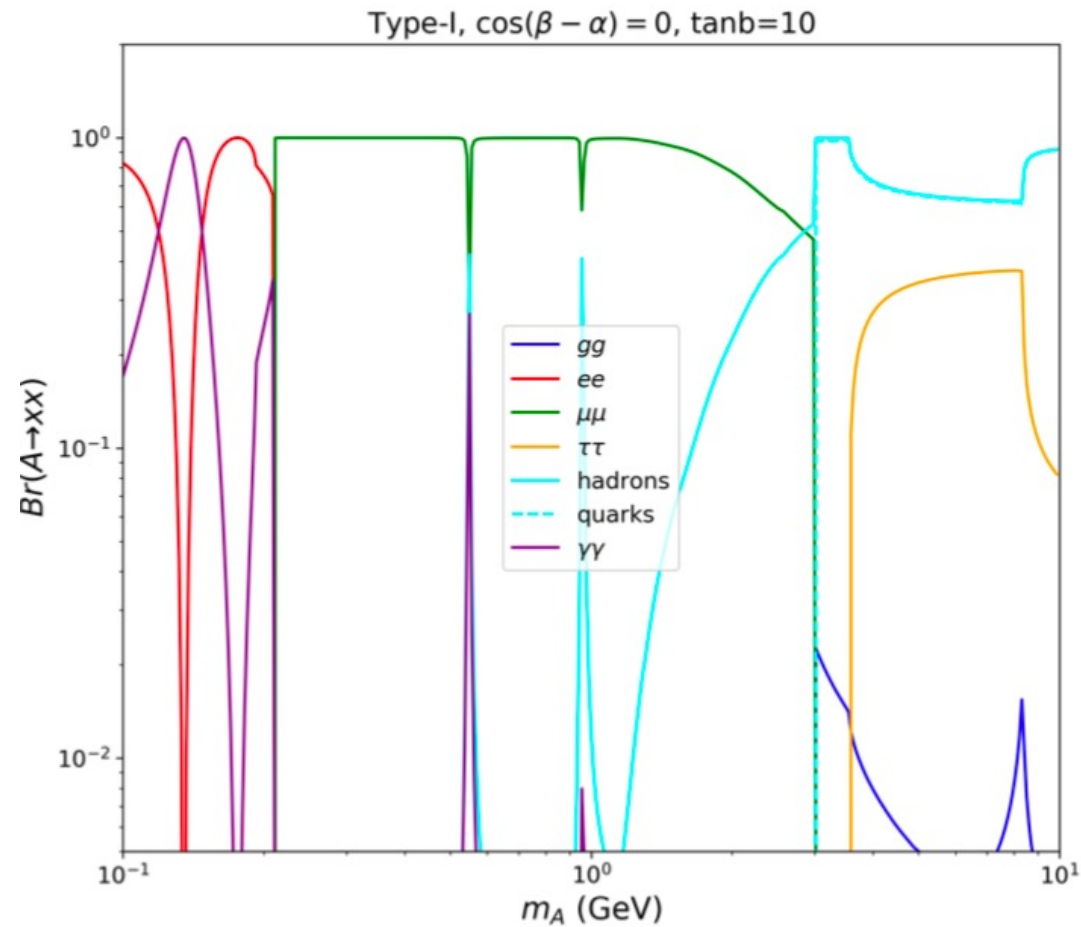
$$\frac{\text{Br}(B \rightarrow X_s\phi)}{\text{Br}(B \rightarrow X_c e \nu)} = \frac{\Gamma(b \rightarrow s\phi)}{\Gamma(b \rightarrow c e \nu)} = \frac{12\pi^2 v^2}{m_b^2} \left(1 - \frac{m_\phi^2}{m_b^2}\right)^2 \frac{1}{f(m_c^2/m_b^2)} \left| \frac{\xi_\phi^{bs}}{V_{cb}} \right|^2$$

$$\text{Br}(K^\pm \rightarrow \pi^\pm \phi) = \frac{1}{\Gamma_{K^\pm}} \frac{2p_\phi^0}{m_{K^\pm}} \frac{|\mathcal{M}|^2}{16\pi m_{K^\pm}},$$



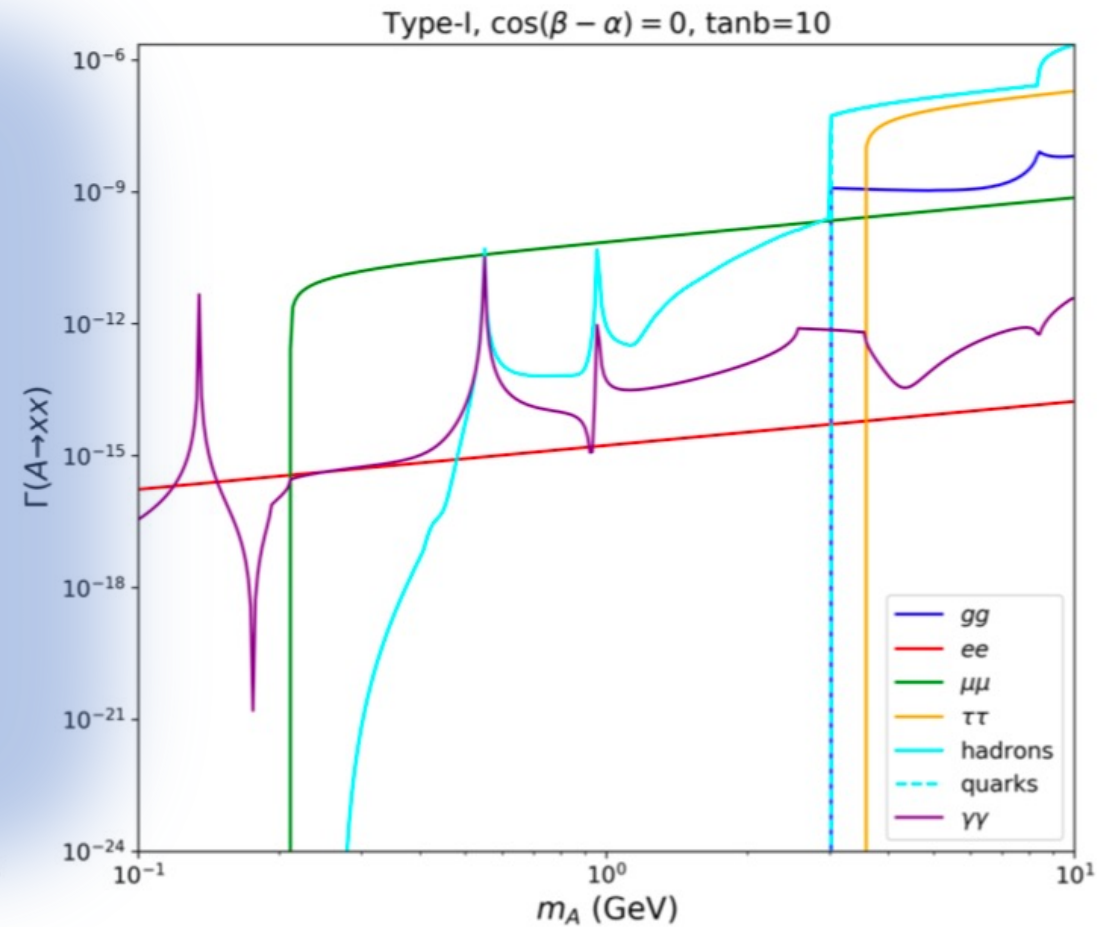
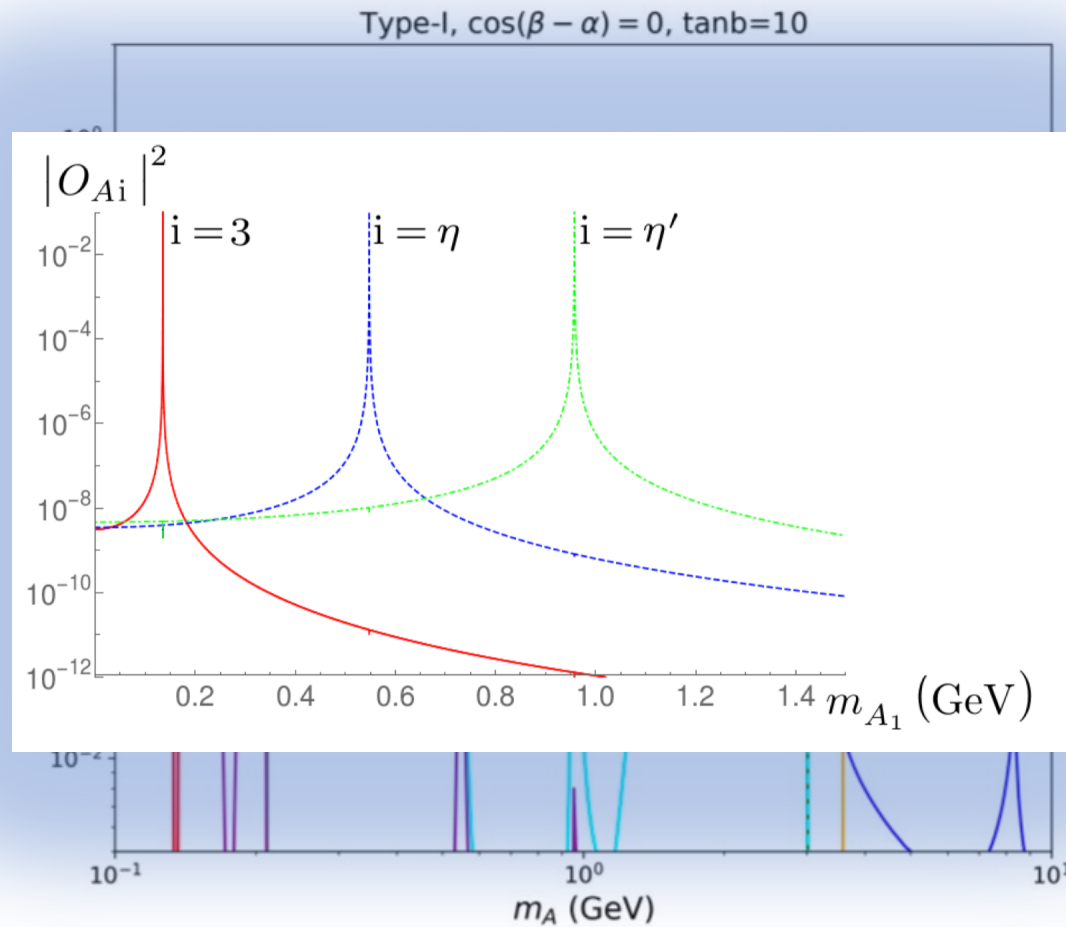
Results: CP odd

$$\xi_A^f|_{\cos(\beta-\alpha)=0} = 1/\tan\beta$$



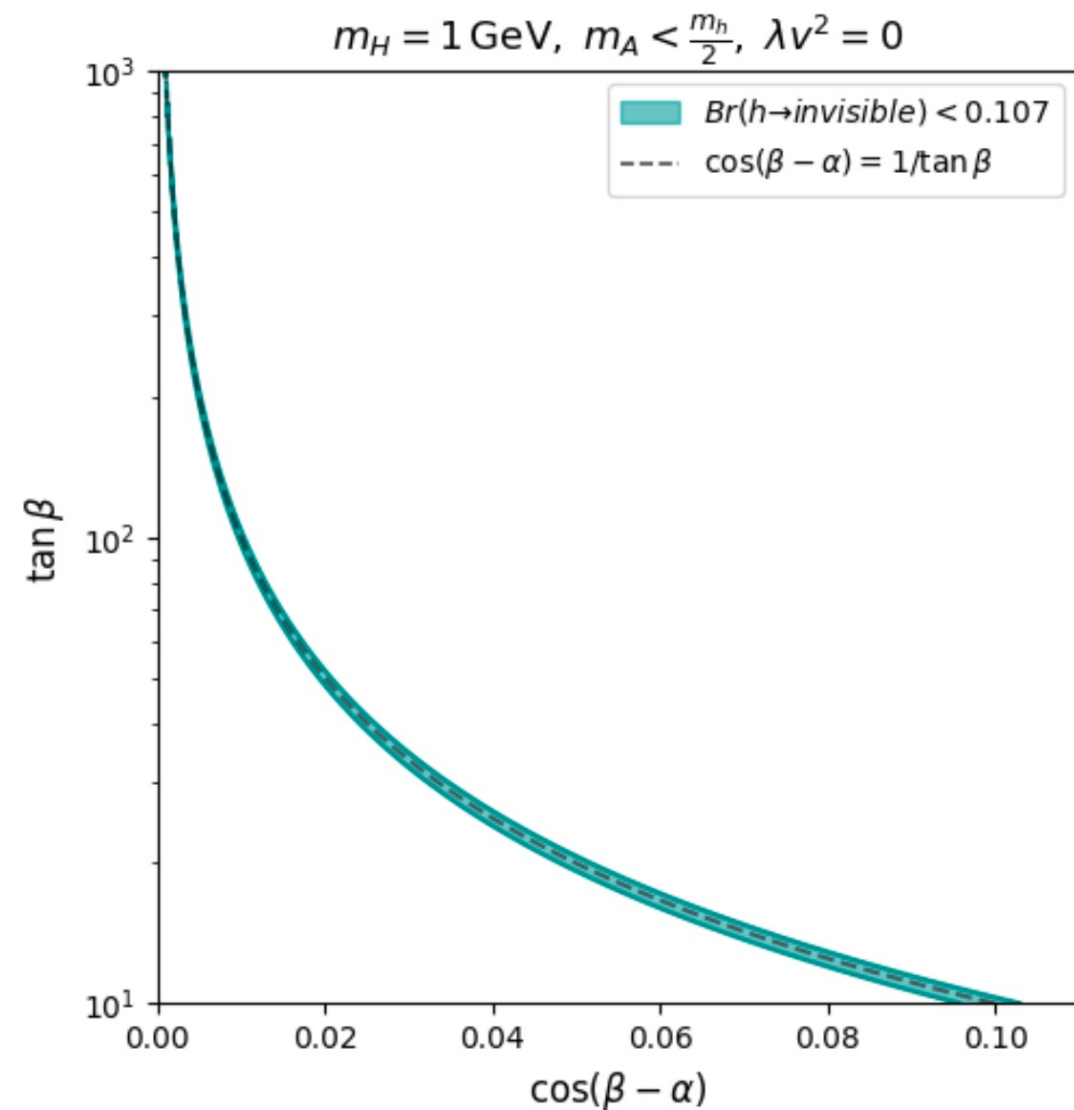
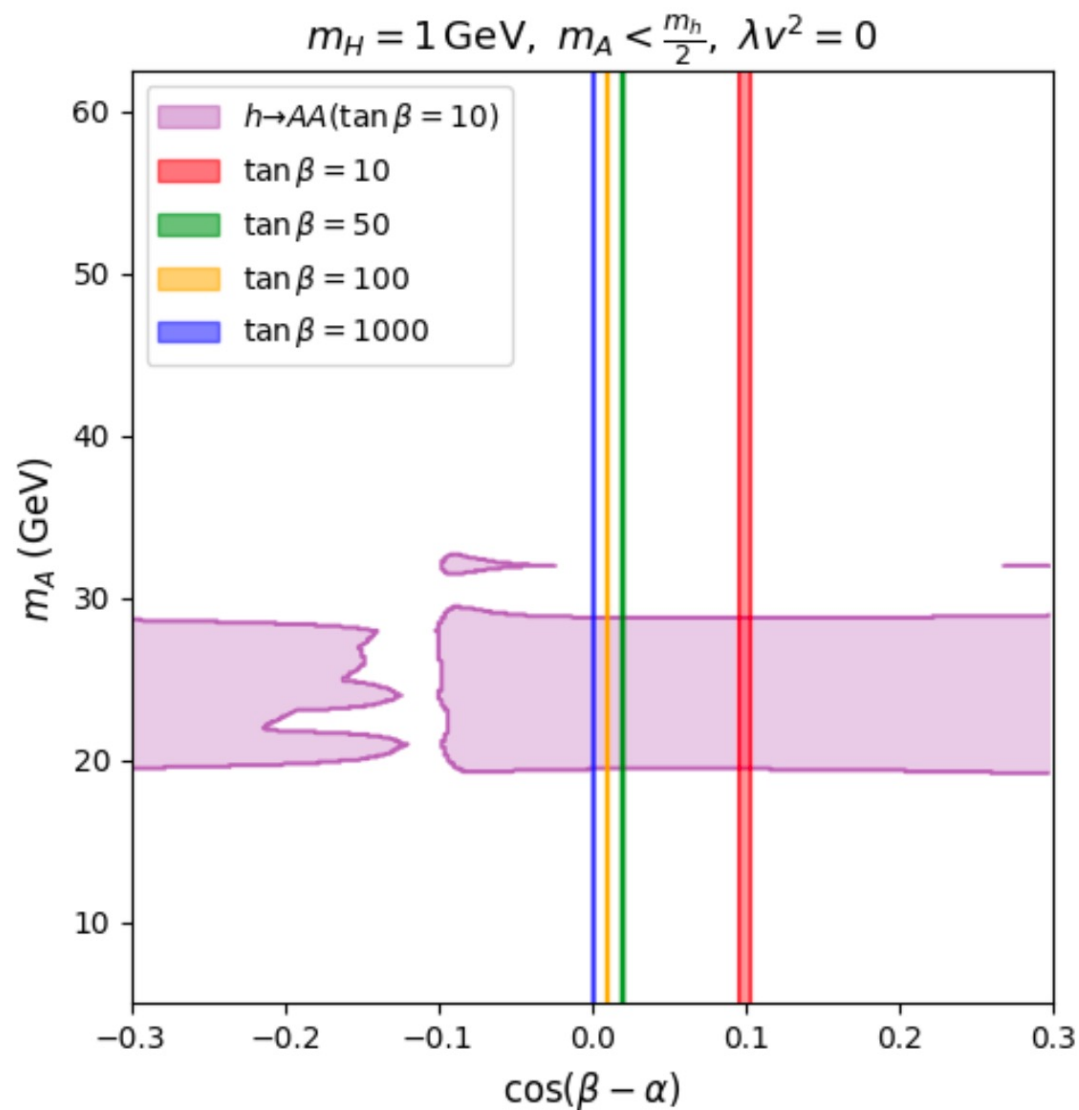
Results: CP odd

$$\xi_A^f|_{\cos(\beta-\alpha)=0} = 1/\tan\beta$$



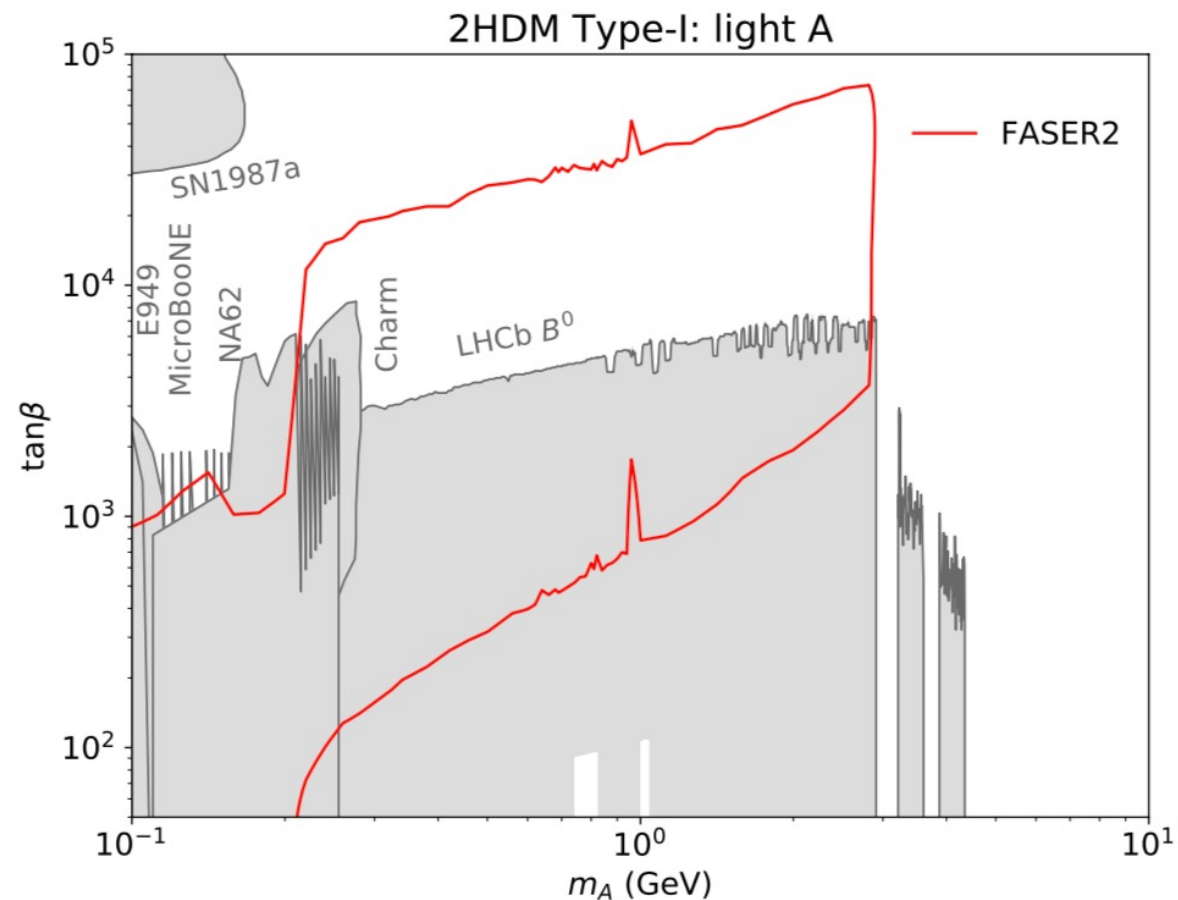
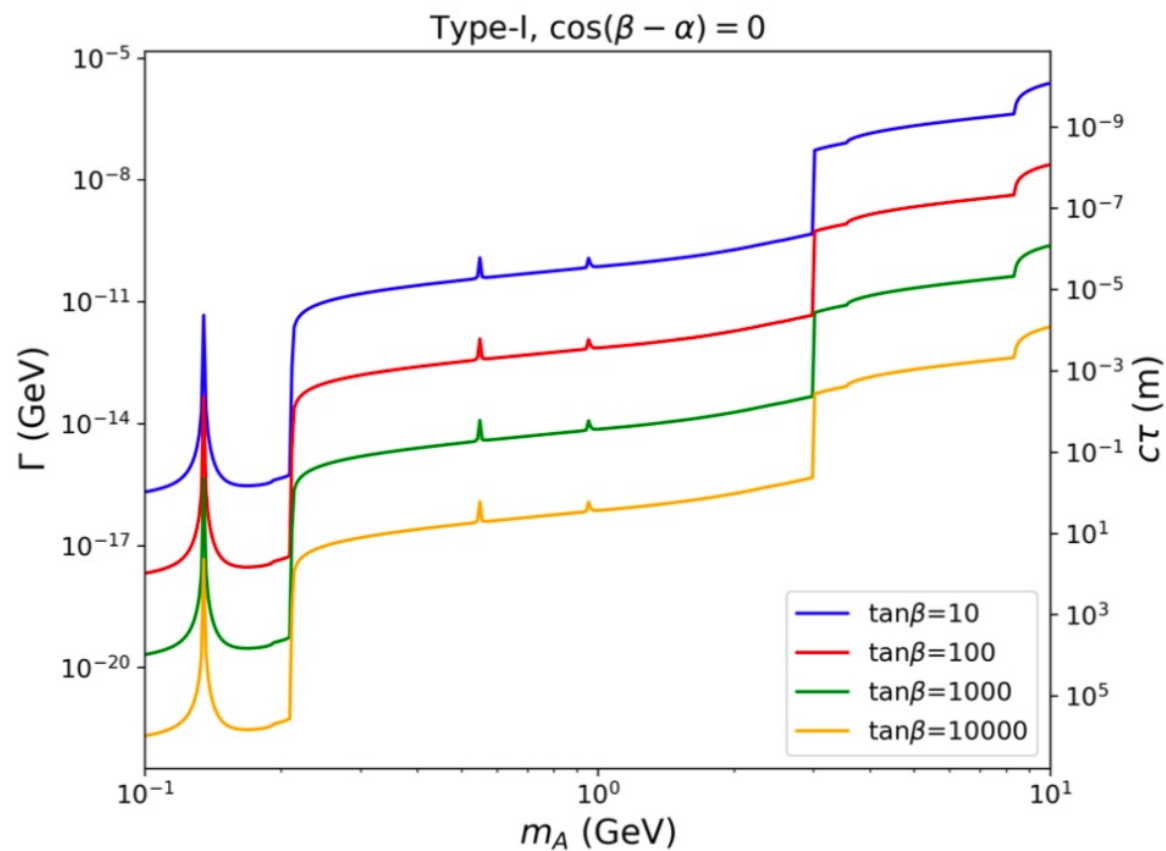
Light H

$$\cos(\beta - \alpha) \simeq \frac{1}{\tan \beta}, \quad m_A \in (54, 600) \text{ GeV}, \quad m_{H^\pm} \sim m_A, \quad \lambda v^2 = 0.$$

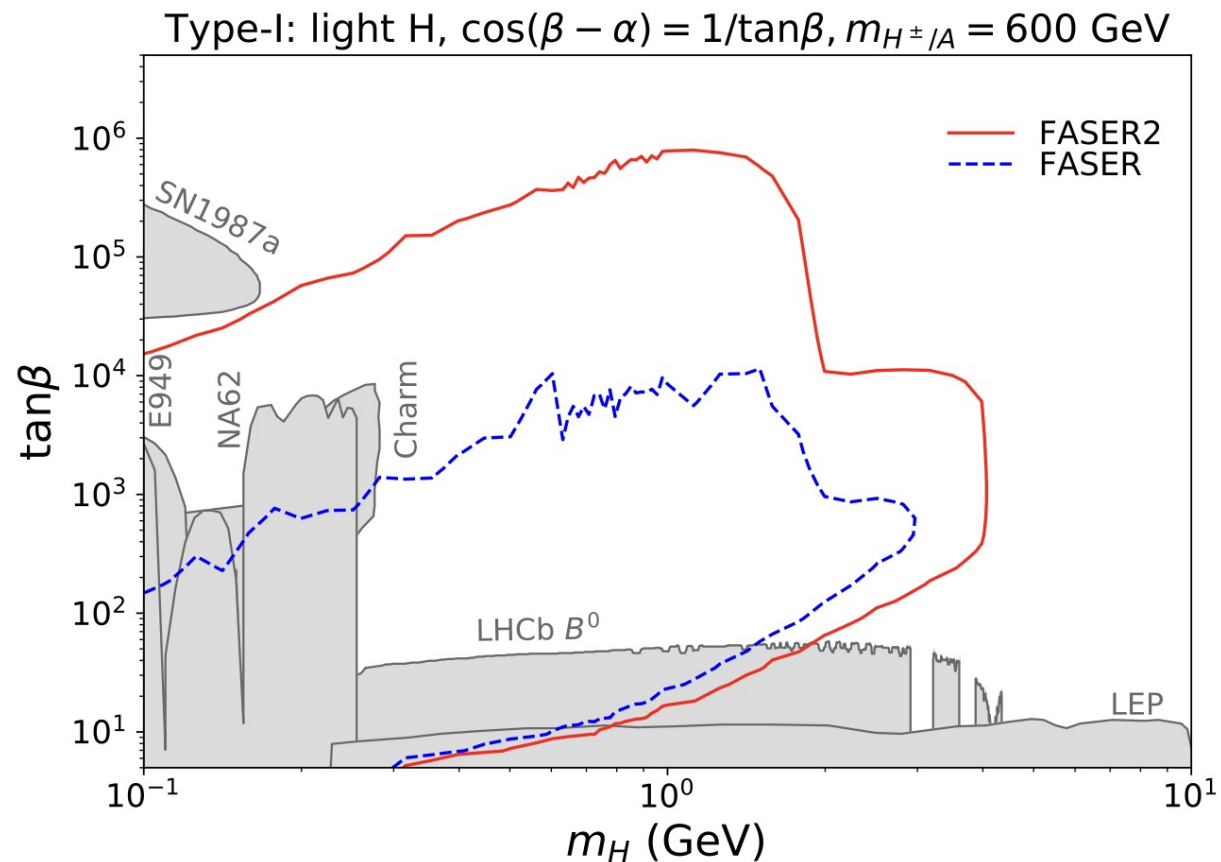
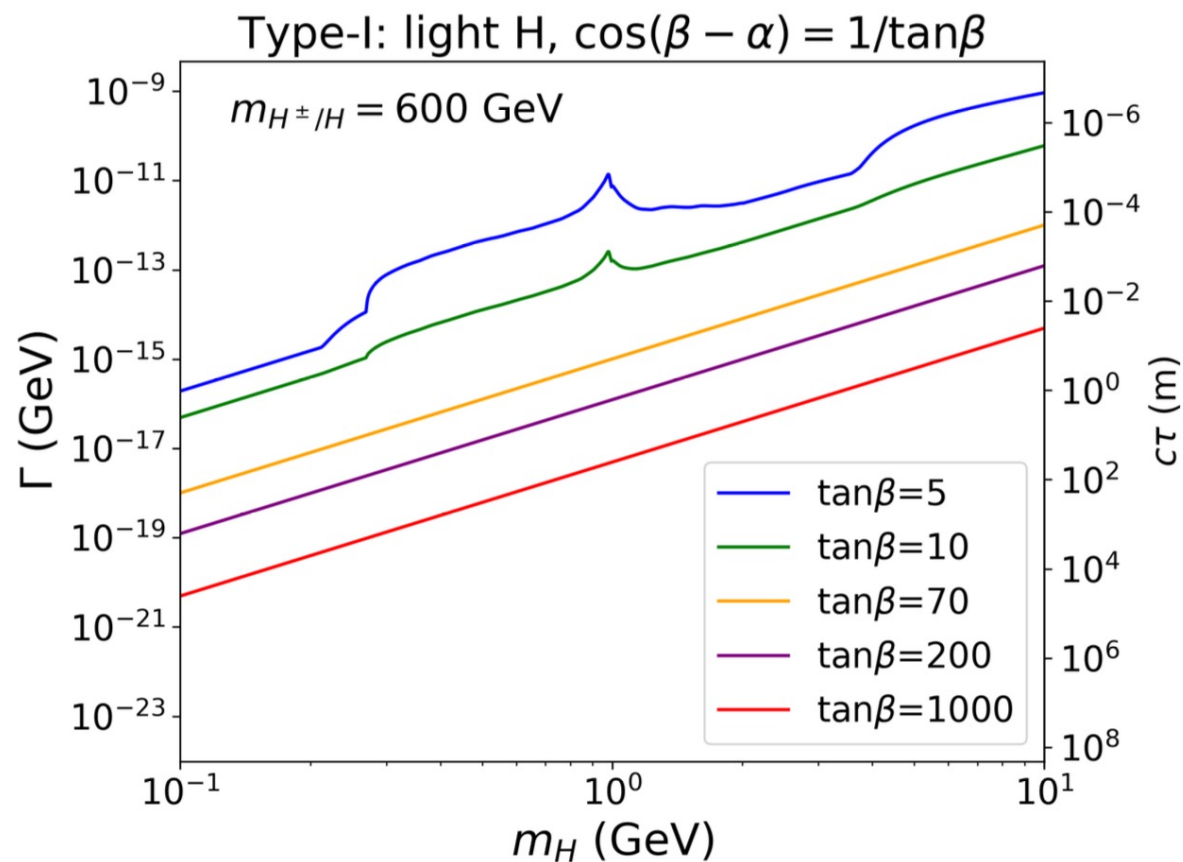


Results: CP odd

$$\xi_A^f|_{\cos(\beta-\alpha)=0} = 1/\tan\beta$$



Results: CP even BM



Decay: CP even scalar

$$m_\phi < 2 \text{ GeV}$$

$$\Gamma_\pi = \langle \pi\pi | m_u \bar{u}u + m_d \bar{d}d | 0 \rangle, \quad \Delta_\pi = \langle \pi\pi | m_s \bar{s}s | 0 \rangle, \quad \Theta_\pi = \langle \pi\pi | \Theta_\mu^\mu | 0 \rangle$$

$$H \rightarrow \pi\pi$$

$$H \rightarrow KK$$

$$H \rightarrow \pi\pi\pi\pi$$

Leading order chiral
perturbation theory

$$\Gamma_\pi^0 = m_\pi^2,$$

$$\Delta_\pi^0 = 0,$$

$$\Theta_\pi^0 = s + 2m_\pi^2$$

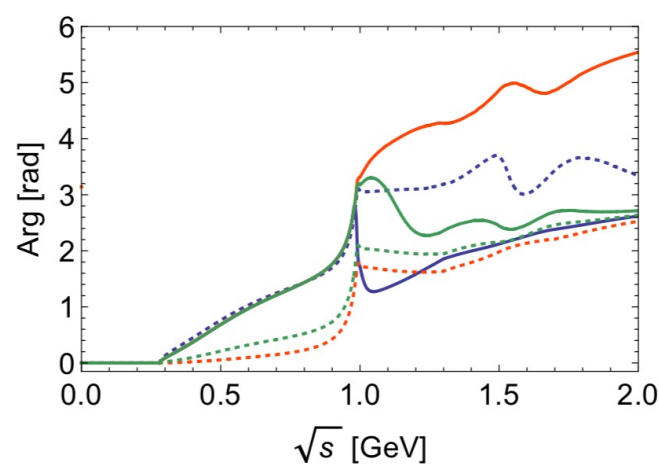
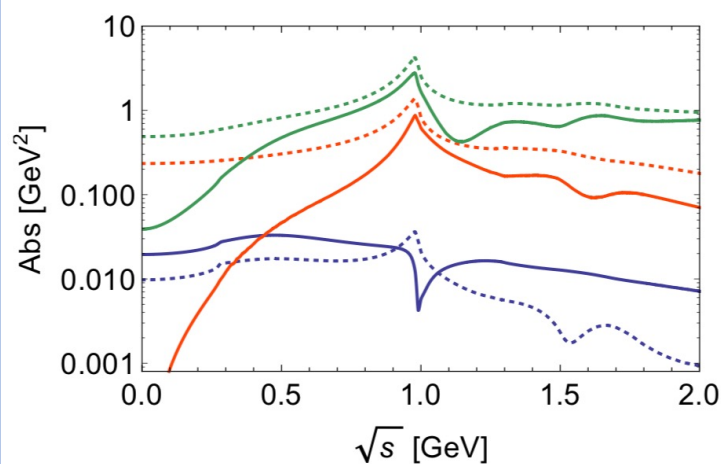
$$\Gamma_K^0 = \frac{1}{2}m_\pi^2,$$

$$\Delta_K^0 = m_K^2 - \frac{1}{2}m_\pi^2,$$

$$\Theta_K^0 = s + 2m_K^2$$

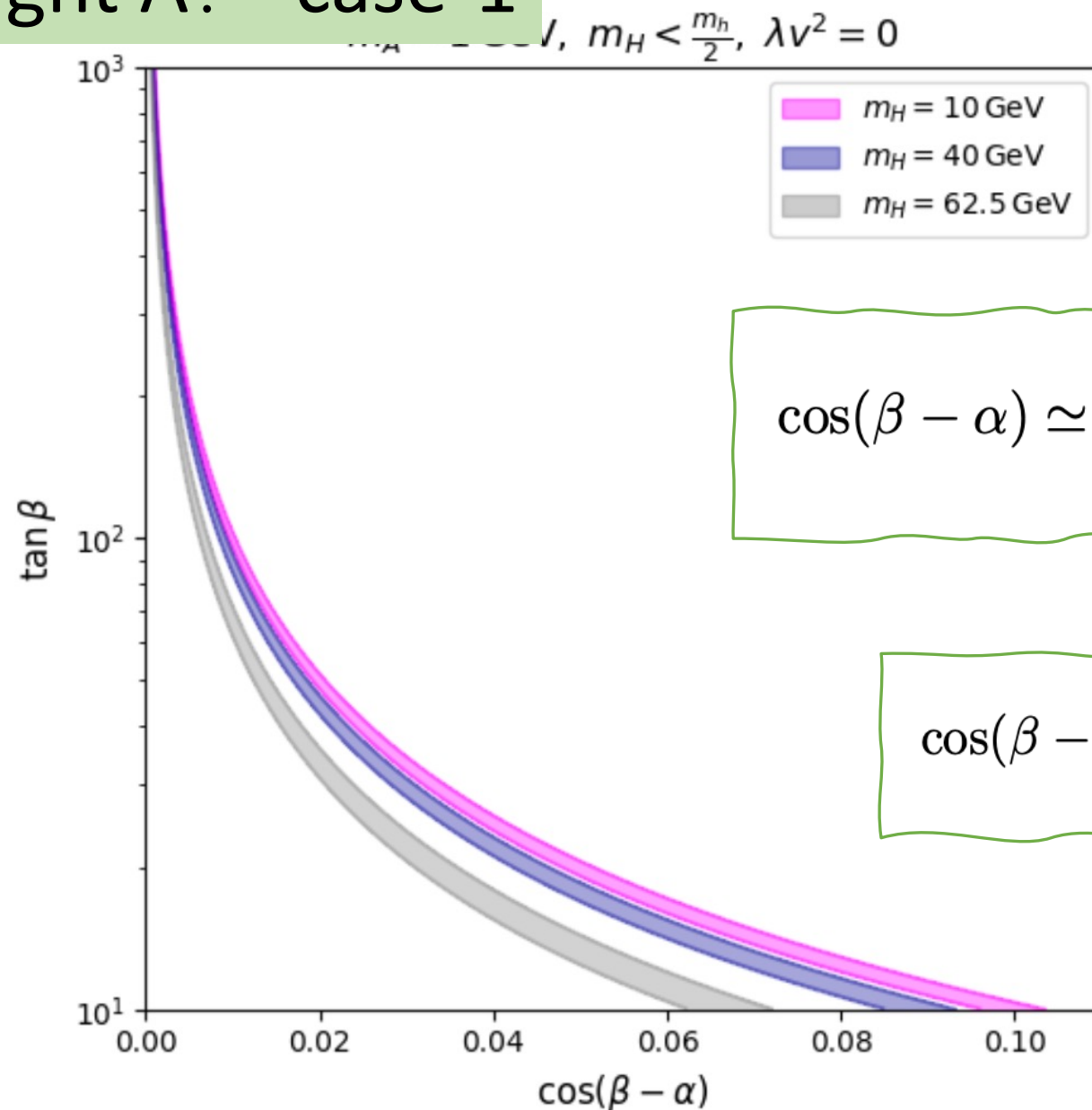
$$m_\phi < 0.5 \text{ GeV}$$

Dispersive Analysis



$$0.5 \text{ GeV} < m_\phi < 2 \text{ GeV}$$

Light A: case-1



$$\cos(\beta - \alpha) \simeq \frac{1}{\tan \beta} \frac{2m_H^2 - m_h^2}{m_H^2 - m_h^2}, \quad m_H \in \left(5, \frac{m_h}{2}\right) \text{ GeV.}$$



$$\cos(\beta - \alpha) \simeq \frac{1}{\tan \beta}, \quad m_H \in (5, 25) \text{ GeV.}$$