

The study of D_{sJ} at Belle and Belle II

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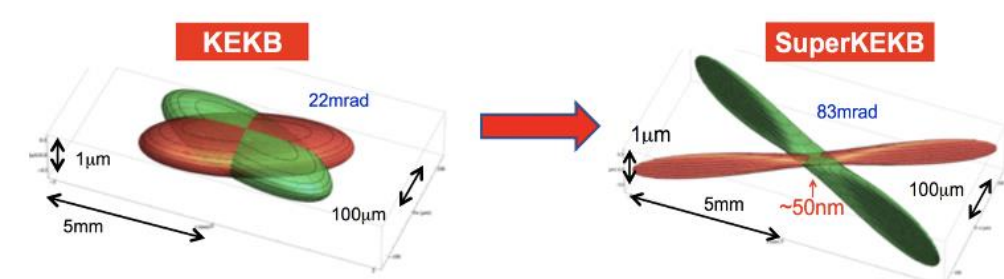
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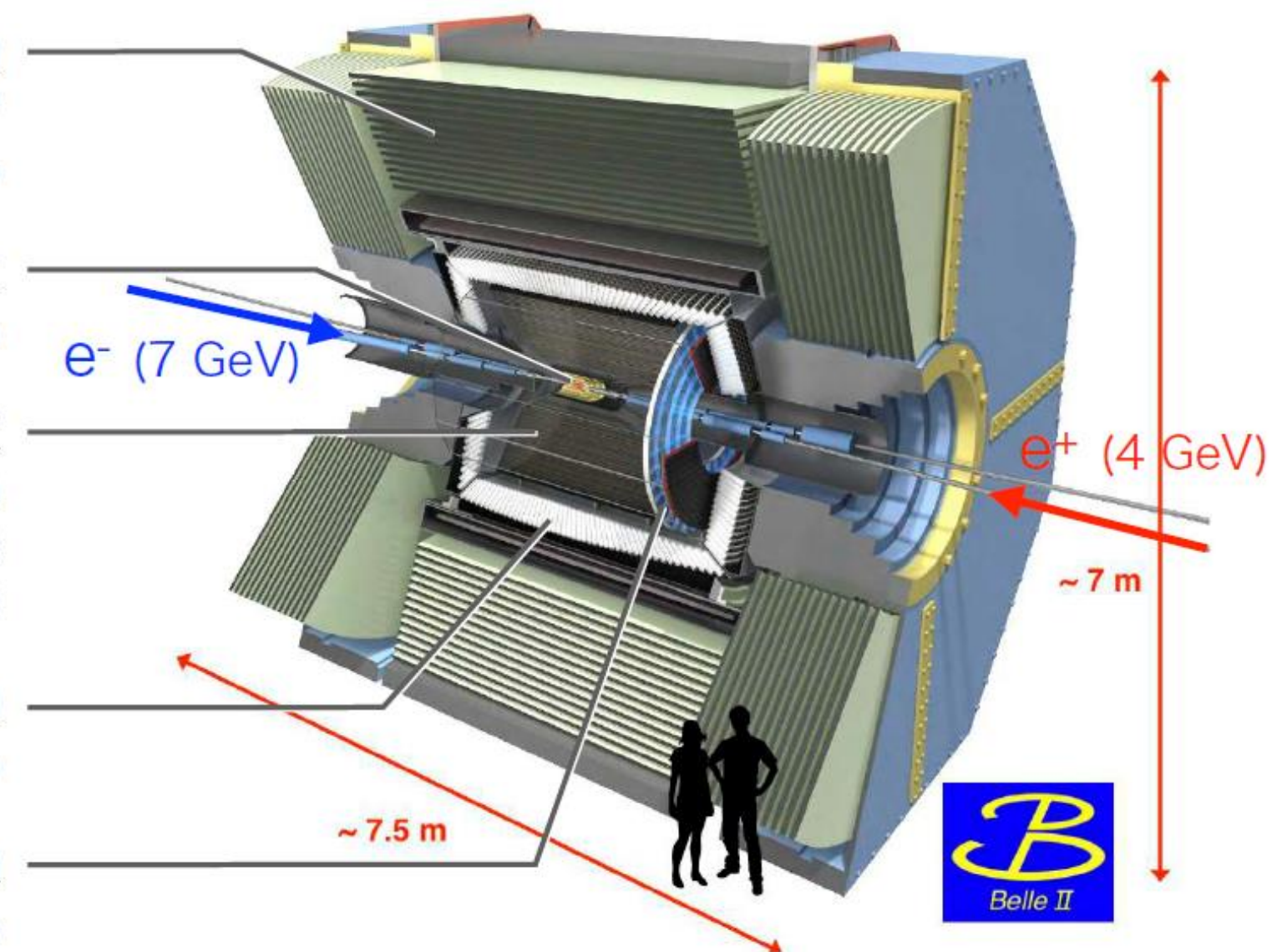
$$e^+ e^- \rightarrow c \bar{c}$$

D_{sJ} physics

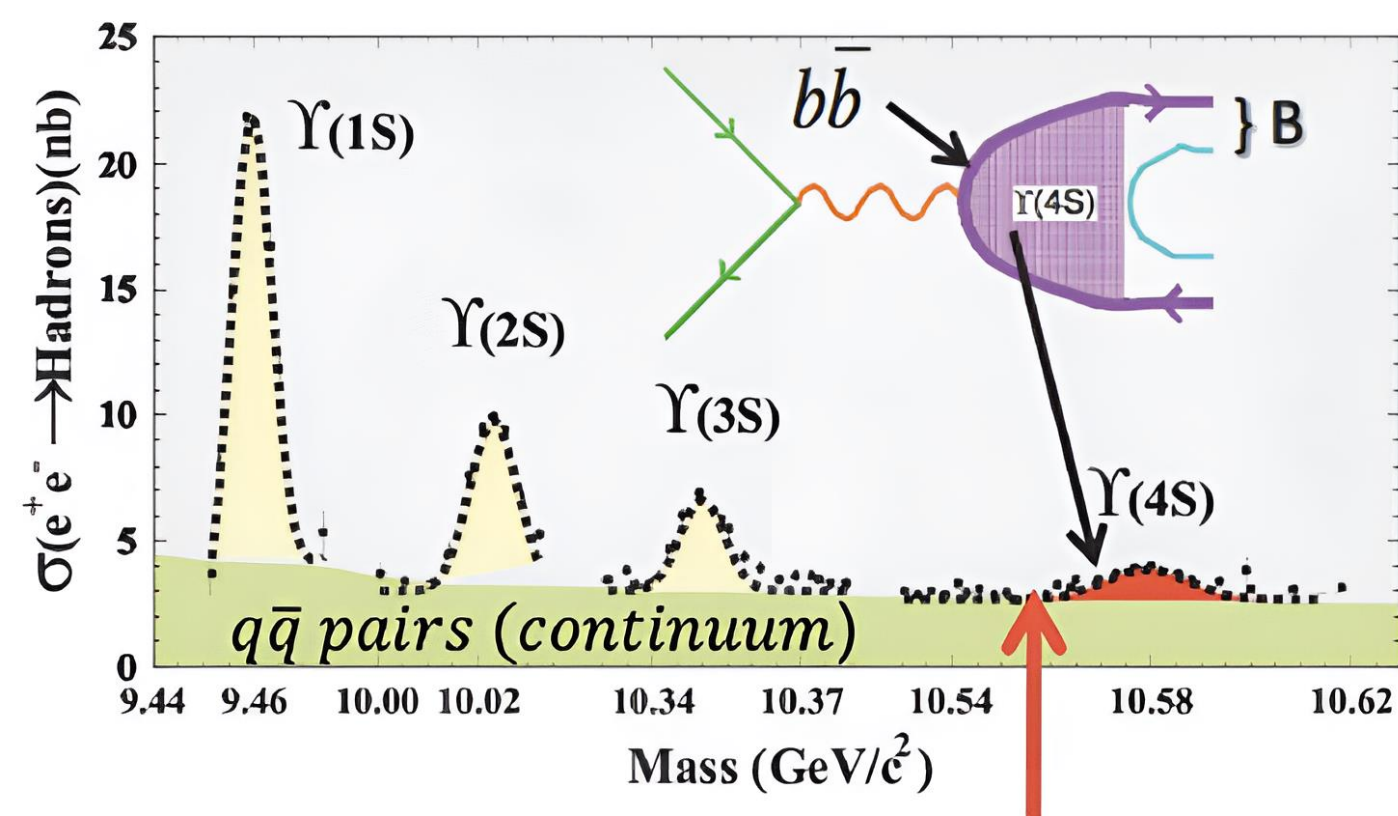
- Extract $|V_{cs}|$ via (semi)leptonic decay $\rightarrow$ BESIII
- Amplitude analyses $\rightarrow$ BESIII
- Spectroscopy $\rightarrow$ Belle and LHCb
-



- KL and μ Detection
KL⁰ p resolution: 15 MeV
 μ identification eff.: ~90%
- Vertex Detector
vertex resolution: 15 μ m
- Central Drift Chamber
spatial resolution: 100 μ m
dE/dx resolution: 5%
P_T resolution: 0.4%
- EM Calorimeter, CsI(Tl)
energy resolution: 1.6%-4%
- Particle Identification
K eff.: 90%, fake π rate: 5%



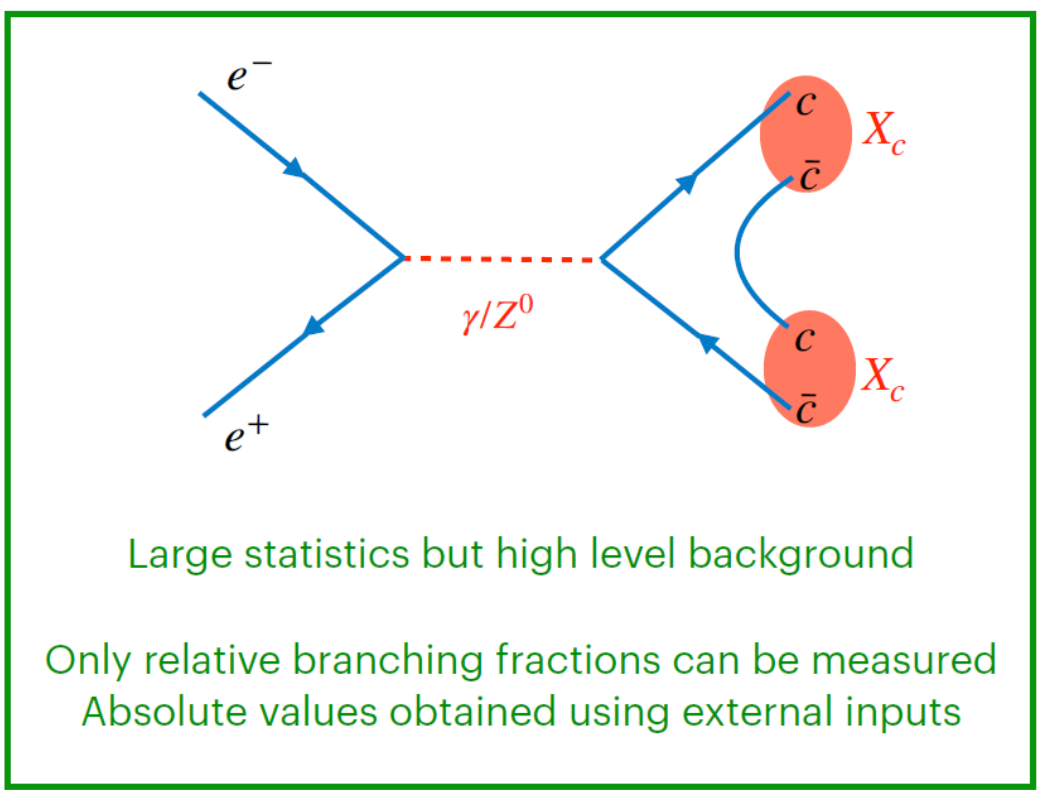
Two processes of D_{sJ} generation at Belle (II) :



$B\bar{B}$ thres.

$e^+e^- \rightarrow$	Cross section [nb]
$\Upsilon(4S)$	1.05 ± 0.10
$c\bar{c}$	1.30
$s\bar{s}$	0.38
$u\bar{u}$	1.61
$d\bar{d}$	0.40
$\tau^+\tau^-(\gamma)$	0.919
$\mu^+\mu^-(\gamma)$	1.148
$e^+e^-(\gamma)$	300 ± 3

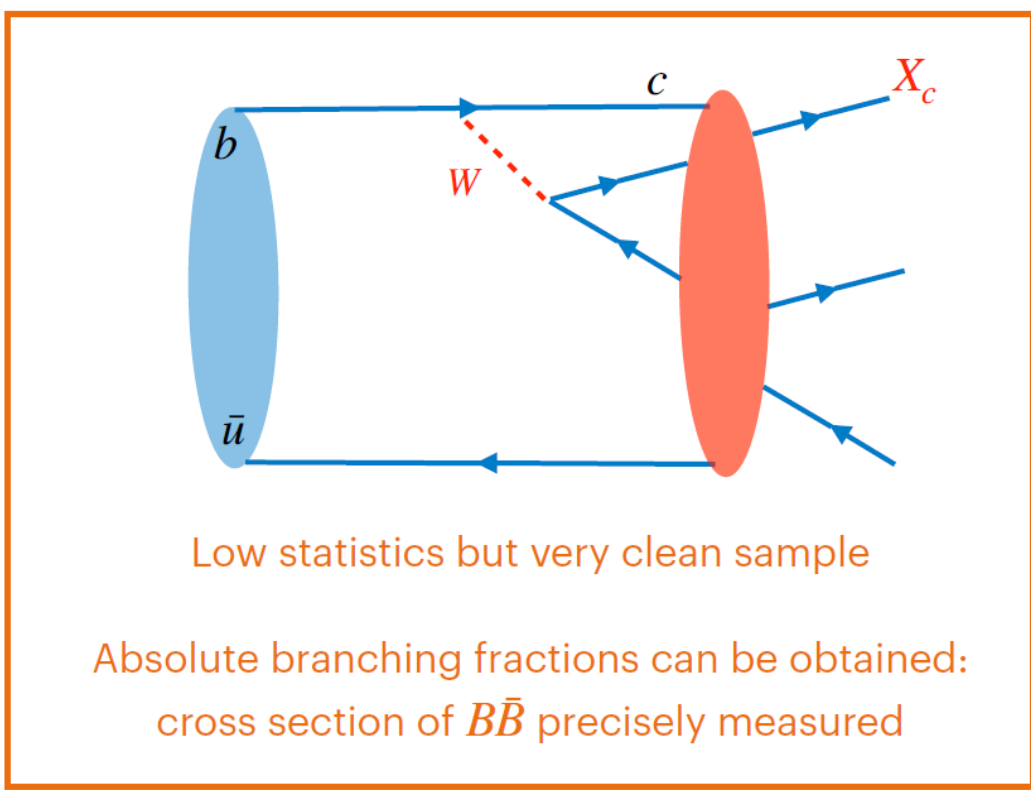
Two charmed hadrons produced from continuum



Large statistics but high level background

Only relative branching fractions can be measured
Absolute values obtained using external inputs

One or more charmed hadrons produced in B decays



Low statistics but very clean sample

Absolute branching fractions can be obtained:
cross section of $B\bar{B}$ precisely measured

Spectrum of D_{sJ}

Xian-Hui Zhong 2110.05024

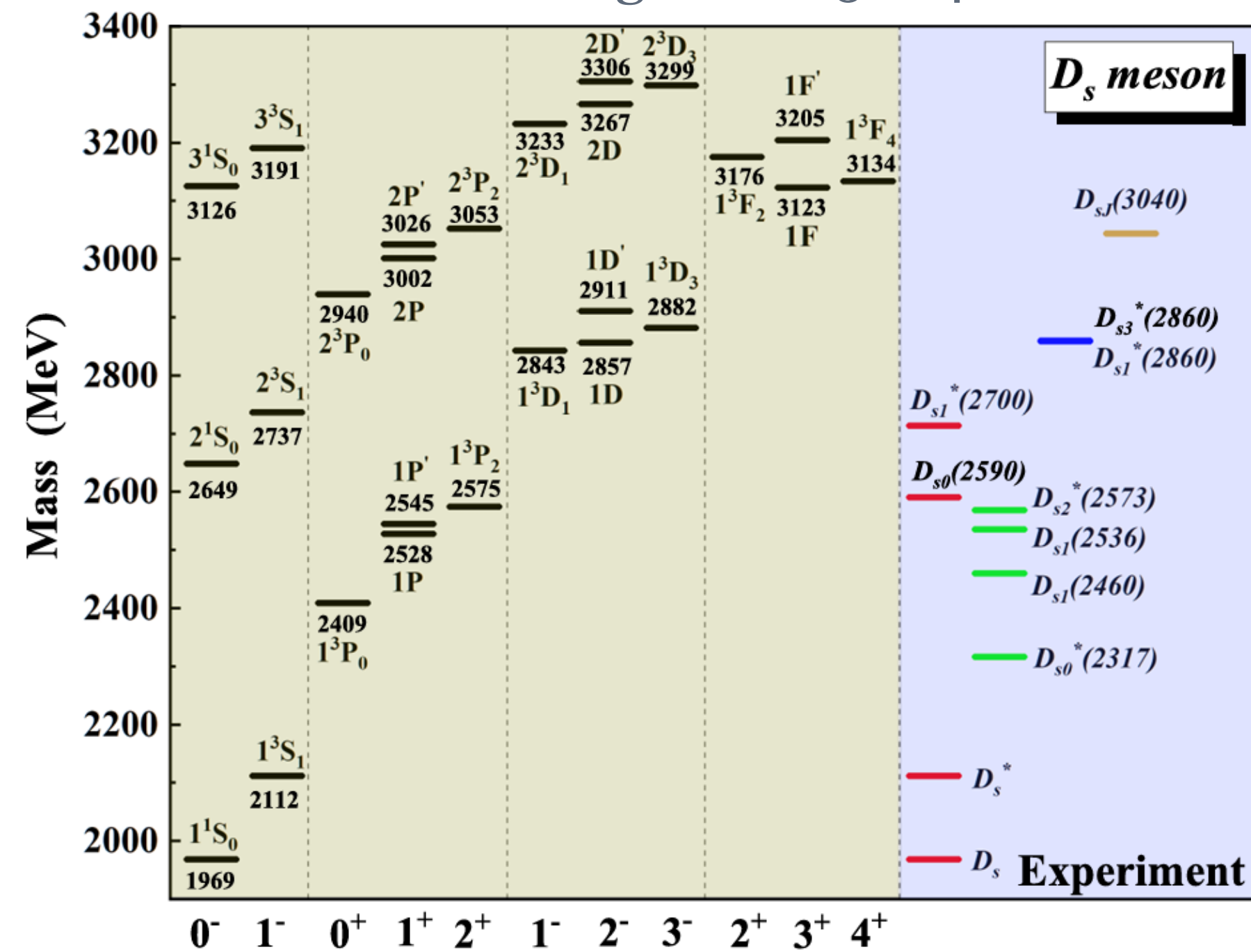
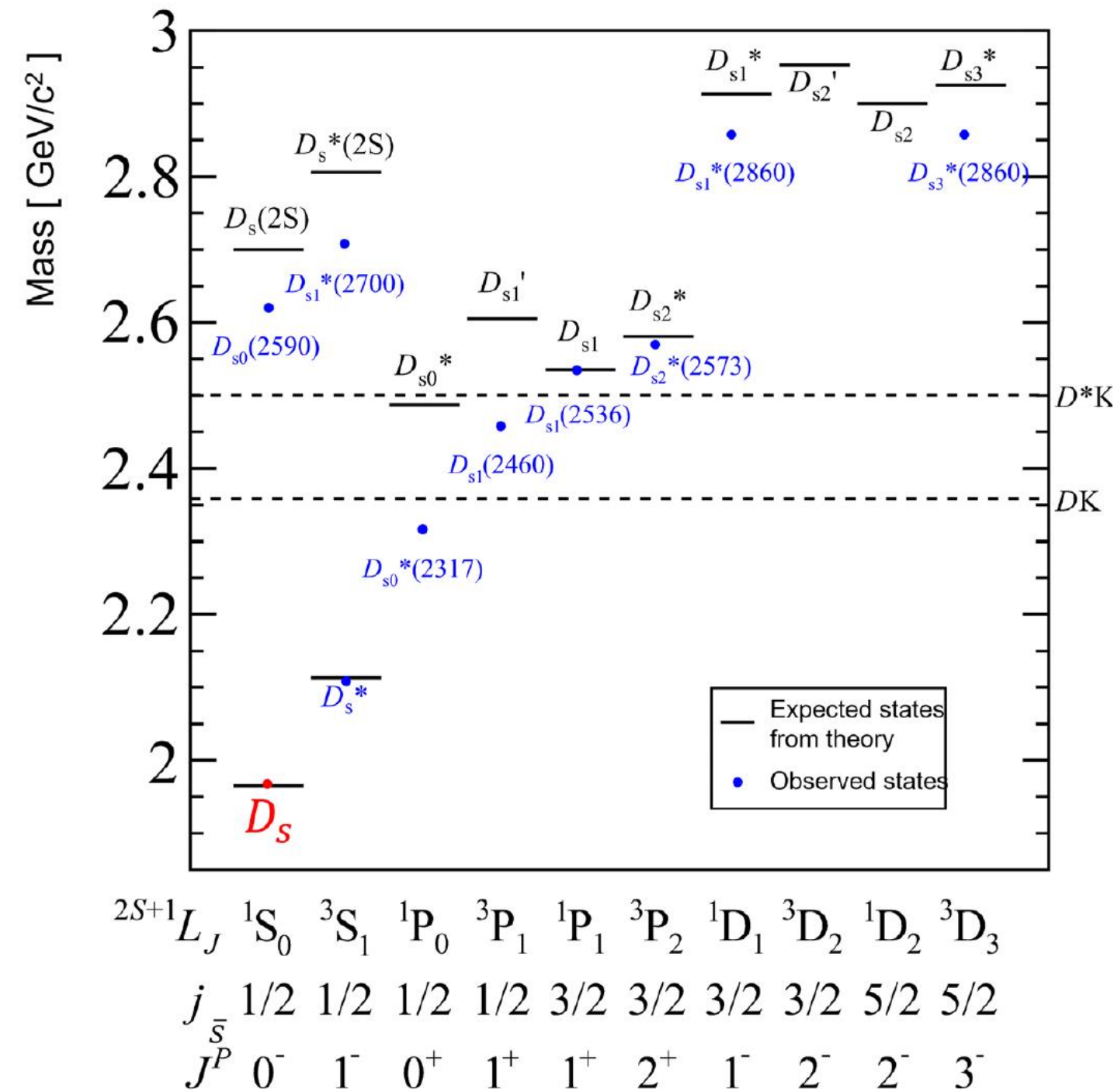


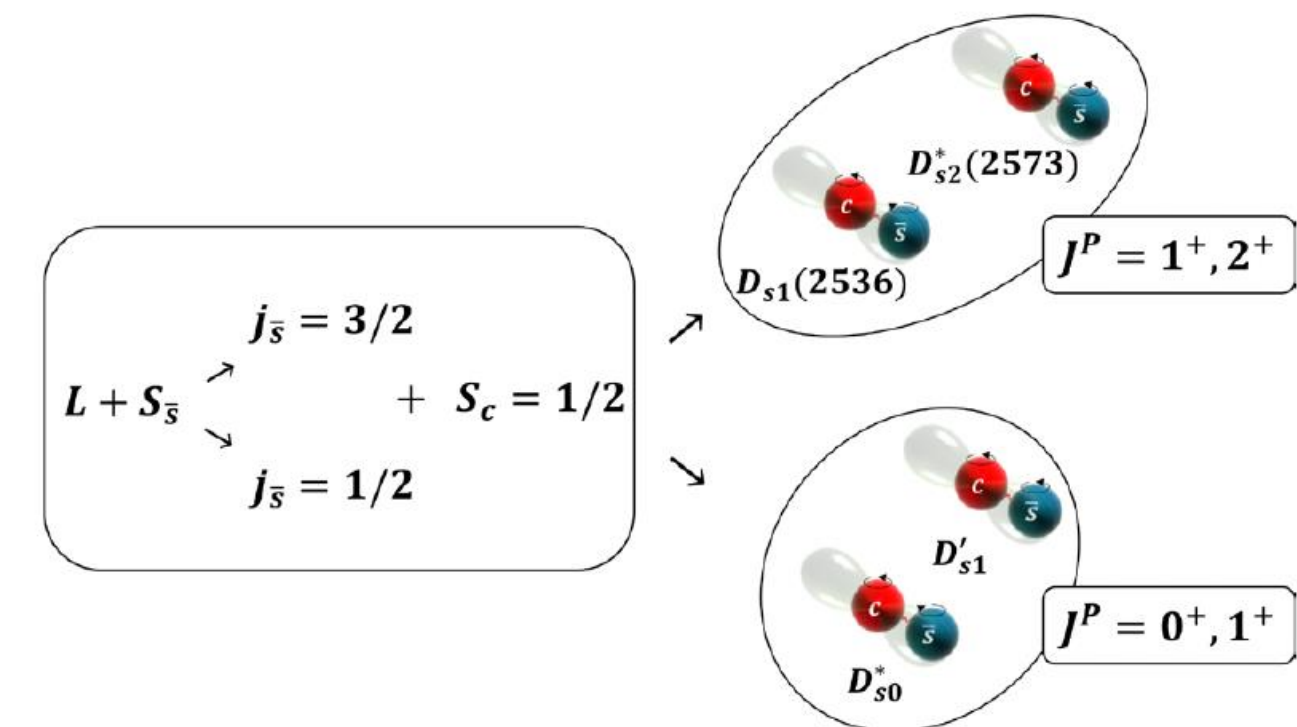
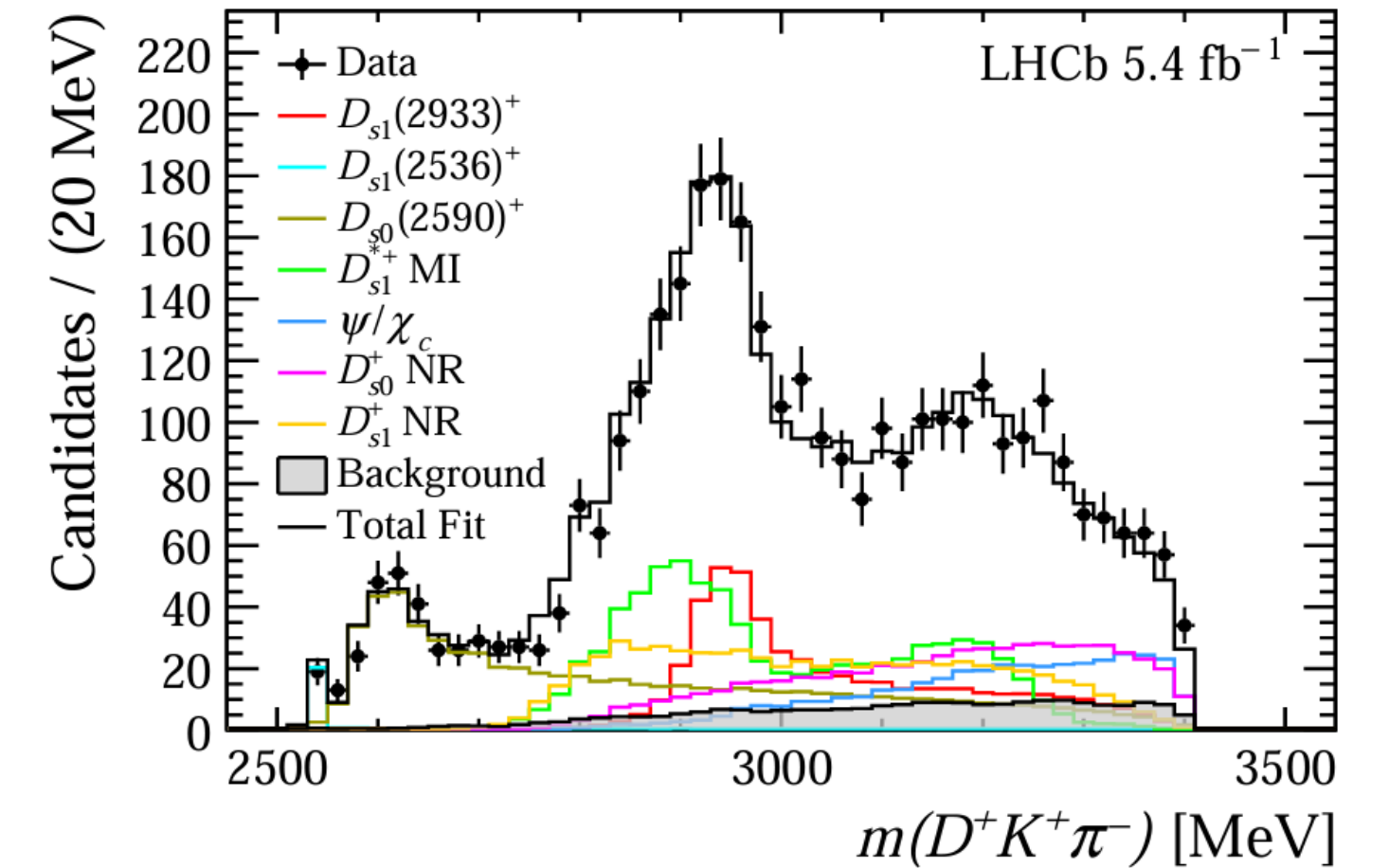
FIG. 2: Predicted charmed-strange meson mass spectrum compared with the observations.

- low-lying status are well known
- High excited states less know: mixing states?
- Possible exotic particle.
- There is still missing D_s

QCD@Work 2022



2604.21257 $D_{s1}(2933)^+, J^P = 1^+$



$D_{sJ}@Belle$ (II)

Production:

- $B \rightarrow \bar{D} D_{s0}^{*+}$ (2317): [Phys.Rev.D 91 \(2015\) 9, 092011](#)
 - measured the branching fraction
 - Search for the isospin partners of D_{s0}^{*+} (2317)

$$\frac{\mathcal{B}(B^+ \rightarrow \bar{D}^0 D_{s0}^{*+})}{\mathcal{B}(B^+ \rightarrow \bar{D}^0 D_s^+)} = 0.081_{-0.021}^{+0.026}$$

$$\frac{\mathcal{B}(B^0 \rightarrow D^- D_{s0}^{*+})}{\mathcal{B}(B^0 \rightarrow D^- D_s^+)} = 0.13_{-0.05}^{+0.06}$$

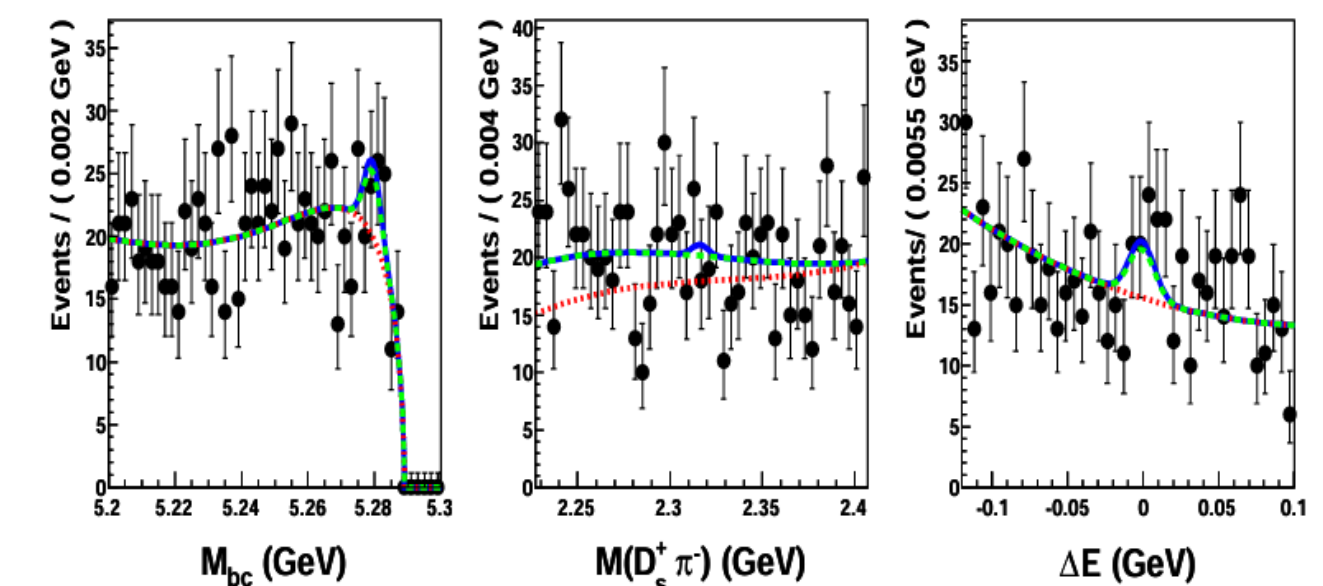
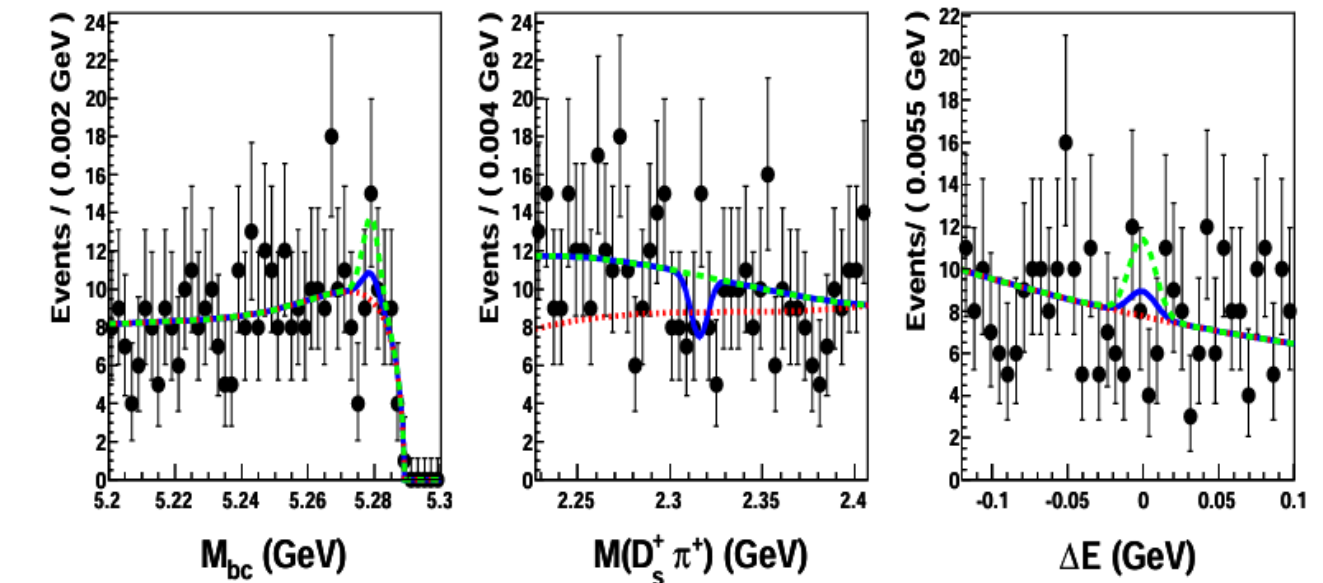
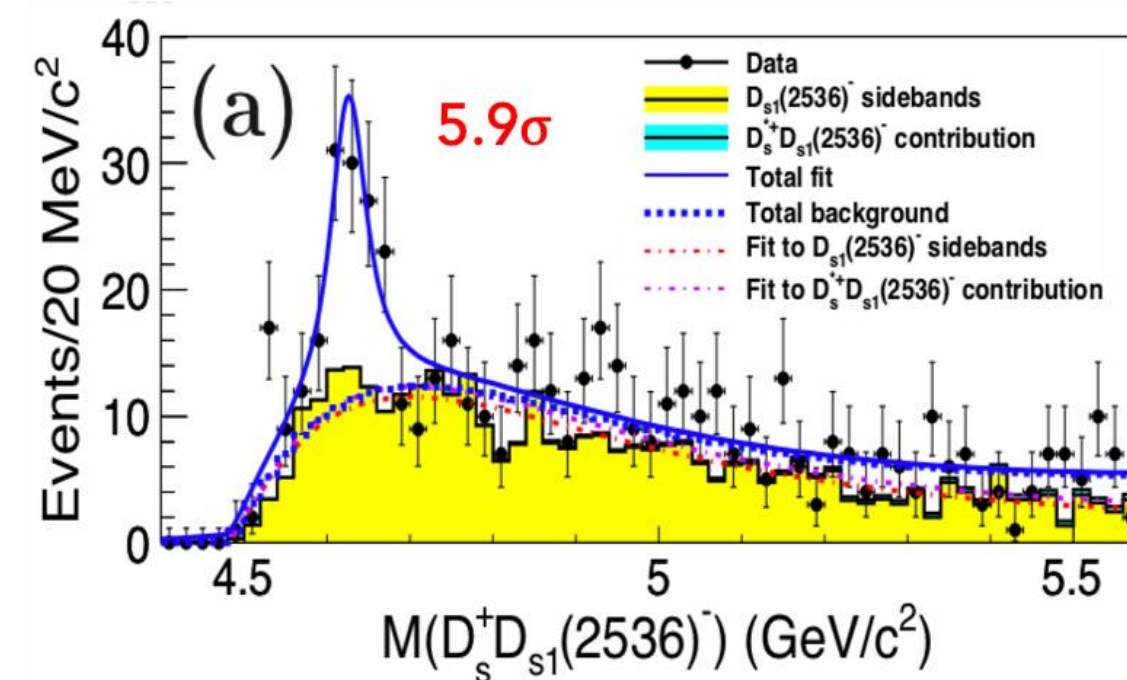
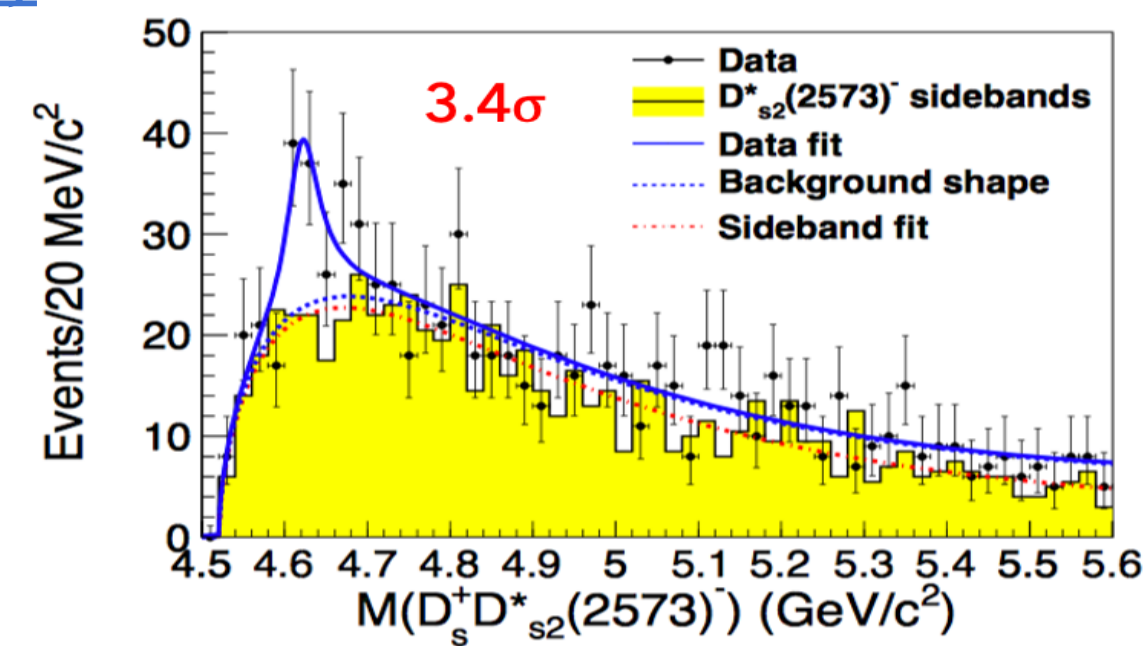
below expectations for a purely $\bar{c}s$ state

- $e^+ e^- \rightarrow D_s^+ D_{s1}(2536)^- / D_s^+ D_{s2}^*(2573)^-$:

[PRD 101, 091101\(R\)](#)

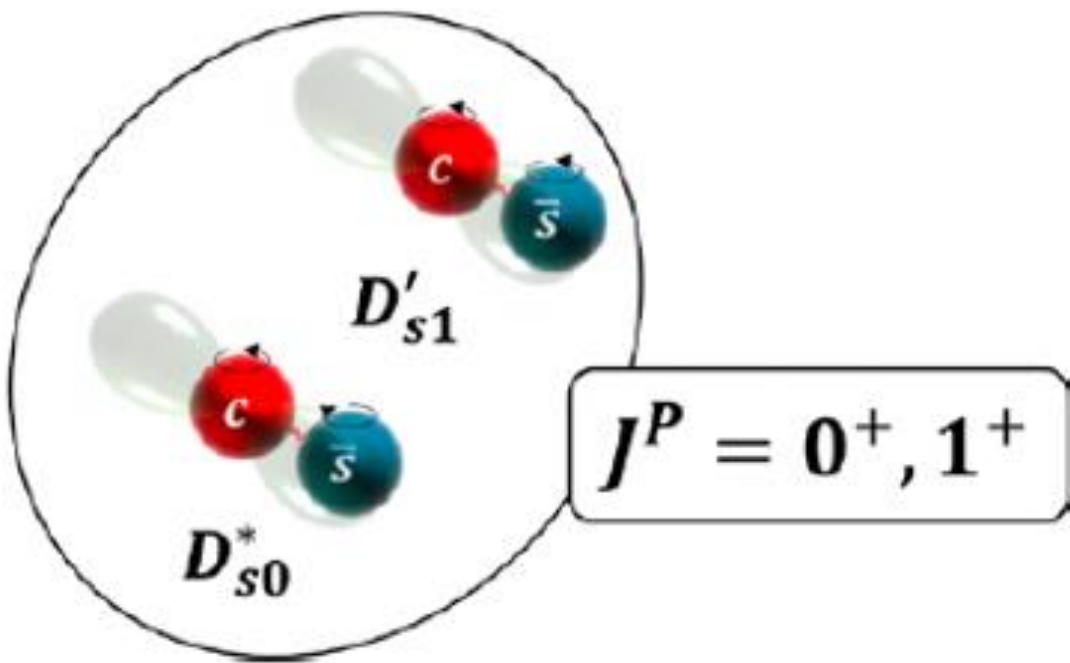
[PRD 100, 111103\(R\)](#)

Experiment	Mass (MeV)	Width (MeV)
Belle, $\Lambda_c^+ \Lambda_c^-$	$4634_{-7}^{+8} {}_{-8}^{+5}$	$92_{-24}^{+40} {}_{-21}^{+10}$
Belle, $\pi^+ \pi^- \psi(2S)$	$4652 \pm 10 \pm 8$	$68 \pm 11 \pm 1$
BaBar, $\pi^+ \pi^- \psi(2S)$	$4669 \pm 21 \pm 3$	$104 \pm 48 \pm 10$
Belle, $D_s^+ D_{s1}(2536)^-$	$4626_{-7}^{+7} \pm 1$	$49.8_{-12}^{+14} \pm 4$
Belle, $D_s^+ D_{s2}^*(2573)^-$	$4620_{-8}^{+9} \pm 3$	$47.0_{-15}^{+32} \pm 5$



$D_{s0}(2317)^+$ and $D_{s1}(2460)^+$

- Mass of $D_{s0}(2317)^+$ and $D_{s1}(2460)^+$ are much lower than the quark model predictions of the lowest $c\bar{s}$ mesons with the corresponding J^P quantum number
 - modifying the $c\bar{s}$ quark model
 - D^*K hadronic molecule
 - compact tetraquarks
 - chiral partners of the ground state D_s^* meson
- Electromagnetic transitions can be well described by QED in the quark level
 - more comprehensively reflect the inner structure than nonperturbative strong decay
- **Width of them can not be measured currently.**
 - The width of $D_s(2317)^+$ is unknown: $\Gamma(D_s(2317)^+) < 3.8 \text{ MeV} @95\% \text{ CL}$.



Partial decay widths:
unique in discriminating between various models

$D_{s0}^*(2317)^\pm$ DECAY MODES

$D_{s0}^*(2317)^-$ modes are charge conjugates of modes below.

Mode	Fraction (Γ_i / Γ)	Scale Factor/ Conf. Level	$P(\text{MeV}/c)$	
$\Gamma_1 \quad D_s^+ \pi^0$	$(100^{+0}_{-20}) \%$		298	✓
$\Gamma_2 \quad D_s^+ \gamma$	$<5 \%$	CL=90%	323	✓
$\Gamma_3 \quad D_s^*(2112)^+ \gamma$	$<6 \%$	CL=90%		✓
$\Gamma_4 \quad D_s^+ \gamma \gamma$	$<18 \%$	CL=95%	323	✓
$\Gamma_5 \quad D_s^*(2112)^+ \pi^0$	$<11 \%$	CL=90%		✓
$\Gamma_6 \quad D_s^+ \pi^+ \pi^-$	$<4 \times 10^{-3}$	CL=90%	194	✓
$\Gamma_7 \quad D_s^+ \pi^0 \pi^0$	not seen		205	✓

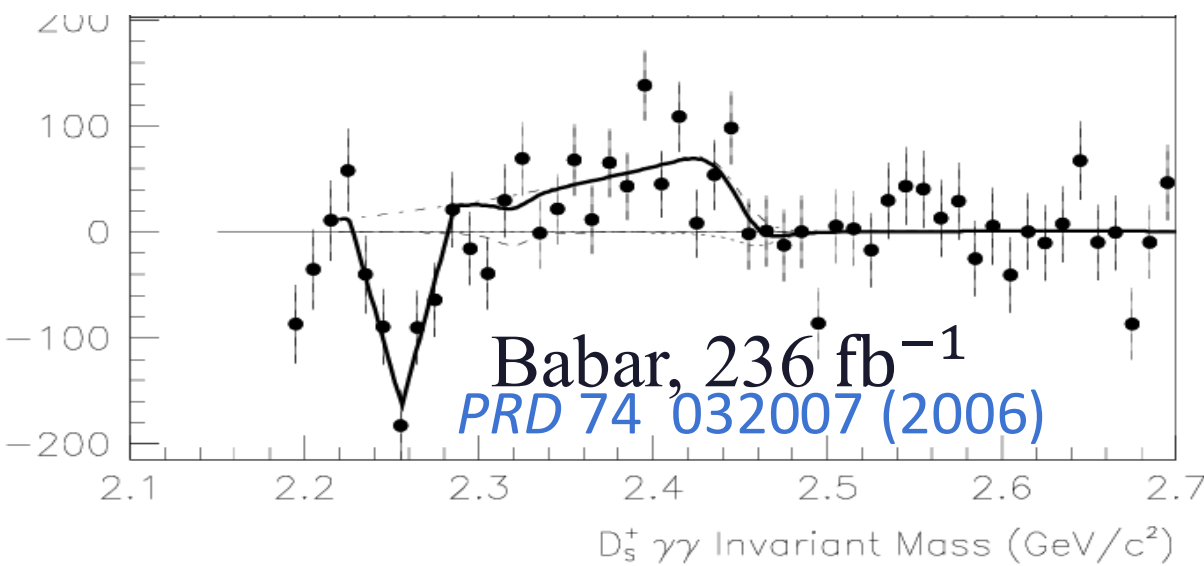
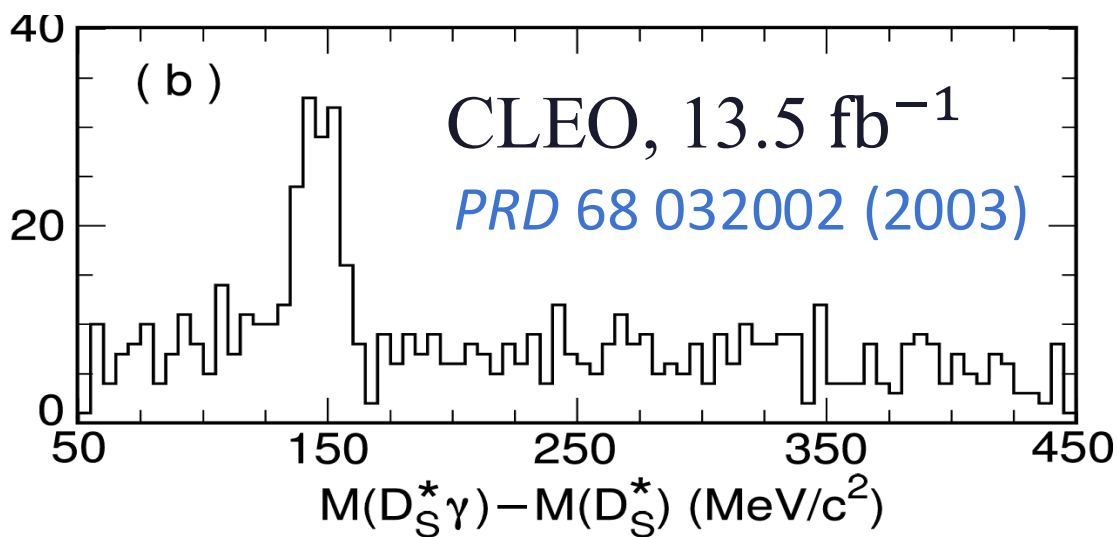
$D_{s1}(2460)^+$ DECAY MODES

$D_{s1}(2460)^-$ modes are charge conjugates of the modes below.

Mode	Fraction (Γ_i / Γ)	Scale Factor/ Conf. Level	$P(\text{MeV}/c)$	
$\Gamma_1 \quad D_s^{*+} \pi^0$	$(48 \pm 11) \%$		297	✓
$\Gamma_2 \quad D_{s+} \gamma$	$(18 \pm 4) \%$		442	✓
$\Gamma_3 \quad D_{s+} \pi^+ \pi^-$	$(4.3 \pm 1.3) \%$	S=1.1	363	✓
$\Gamma_4 \quad D_s^{*+} \gamma$	$<8 \%$	CL=90%	323	✓
$\Gamma_5 \quad D_{s0}^*(2317)^+ \gamma$	$(3.7^{+5.0}_{-2.4}) \%$		138	✓

Radiative decay of $D_{s0}(2317)^+$ and $D_{s1}(2460)^+$

- Precise absolute branching fraction is unknown: $\mathcal{B}(D_s(2317)^+ \rightarrow D_s^+\pi^0) = (100^{+0}_{-20})\%$
- The radiative decay of $D_{s0}(2317)^+$ has not been founded.
 - $\mathcal{R} < 0.059$ @ 90% CL from CLEO, $N(D_s(2317)^+ \rightarrow D_s^{*+}\gamma) = -6.5 \pm 5.2$



$D_{s0}(2317)^+$ Decay mode	J^p transition	Possible wave
$D_s^+\gamma$	$0^+ \rightarrow 0^-1^-$	Forbidden
$D_s^{*+}\gamma$	$0^+ \rightarrow 1^-1^-$	$L = 0, 2$ ★

Eg. $\frac{\Gamma(D_{s1}(2460)^+ \rightarrow D_s^{*+}\gamma)}{\Gamma(D_{s1}(2460)^+ \rightarrow D_s^+\gamma)} = \frac{4.2}{7.3} \tan^2 \theta$

PRD72, 094004 (2005)

Radiative decay:

- $\Delta J = 0, \pm 1$ (But not $0 \rightarrow 0$)
- sign of P $\begin{cases} \text{change: E1 transition} \\ \text{unchange: M1 transition} \end{cases}$

$D_{s1}(2460)^+$ Decay mode	J^p transition	Transition
$D_s^+\gamma$	$1^+ \rightarrow 0^-1^-$	E1
$D_s^{*+}\gamma$	$1^+ \rightarrow 1^-1^-$	
$D_{s0}^{*+}(2317)^+\gamma$	$1^+ \rightarrow 0^+1^-$	M1

- $D_{s1}(2460)^+ \rightarrow D_s^+\gamma$ is observed but $D_s^{*+}\gamma$ not
 - $D_{s1}(2460)^+$ state has the components with both $\mathbf{S} = 0$ (1P_1 state) and $\mathbf{S} = 1$ (3P_1 state)

$$D_{s1}(2460)^+ = \sin\theta |^1P_1\rangle + \cos\theta |^3P_1\rangle$$
 - For E1 transition, $\Delta S = 0$, So, it is useful to **measure $D_s^+\gamma$ and $D_s^{*+}\gamma$ decay ratio**

Radiative decay of $D_{s0}(2317)^+$ and $D_{s1}(2460)^+$

Table I: The theoretical predictions of the $\mathcal{B}(D_{s0}^*(2317)^+ \rightarrow D_s^{*+}\gamma)/\mathcal{B}(D_{s0}^*(2317)^+ \rightarrow D_s^+\pi^0)$ ($\mathcal{R}$)

Reference	$D_{s0}^*(2317)^+$ Structure	Approach	Radiative Decay Width (keV)	Hadronic Decay Width (keV)	$\mathcal{R}(\%)$
Godfrey [40]	$c\bar{s}$ state	quark model	1.9	~ 10	19
Ke [41]		light front quark mode	1.71 ± 2.9	-	$> (0.45 \pm 0.11)$
Bardeen [18]		chiral theory	1.74	21.5	8.1
Faessler [42]	molecule state	effective Lagrangian	$0.55 \sim 1.41$	79.6 ± 33	$0.5 \sim 3.0$
Fu [44]		heavy quark flavor symmetry	3.7 ± 0.3	132 ± 7	2.8 ± 0.9
Lutz [45]		chiral theory	1.94 or 6.47	140	1.35 or 4.57

The branching fraction ratio of $\mathcal{R} = \frac{\mathcal{B}(D_s(2317)^+ \rightarrow D_s^{*+}\gamma)}{\mathcal{B}(D_s(2317)^+ \rightarrow D_s^+\pi^0)}$ and $\mathcal{R} = \frac{\mathcal{B}(D_s(2460)^+ \rightarrow D_s^{*+}\gamma)}{\mathcal{B}(D_s(2460)^+ \rightarrow D_s^+\gamma)}$ are the probe of the $D_{s0}(2317)^+$ and $D_{s1}(2460)^+$ internal structure

Table I: Summary of theoretical predictions for the radiative partial widths of $D_{s1}(2460)^+ \rightarrow D_s^{*+}\gamma$ and $D_{s1}(2460)^+ \rightarrow D_s^+\gamma$, and for their ratio, under different interpretations of the $D_{s1}(2460)^+$. All partial widths are given in keV.

References	Interpretations	$\Gamma(D_{s1}^+ \rightarrow D_s^{*+}\gamma)$	$\Gamma(D_{s1}^+ \rightarrow D_s^+\gamma)$	$\frac{\Gamma(D_{s1}^+ \rightarrow D_s^{*+}\gamma)}{\Gamma(D_{s1}^+ \rightarrow D_s^+\gamma)}$
Godfrey [13]	$c\bar{s}$	5.5	6.2	0.89
Bardeen <i>et al.</i> [14]	$c\bar{s}$	4.66	5.08	0.92
Ferrandes [15]	$c\bar{s}$	1.5	3.3	0.45
Radford <i>et al.</i> [18]	$c\bar{s}$	11.0	8.3	1.33
Green <i>et al.</i> [21]	$c\bar{s}$	5.40	15.2	0.36
Bondar <i>et al.</i> [40]	$c\bar{s}$	104	297	0.35
Colangelo <i>et al.</i> [41]	$c\bar{s}$	0.6–1.1	19–29	0.021–0.058
Close <i>et al.</i> [42]	$c\bar{s}$	$4.2 \sin^2 \theta_s$	$7.3 \cos^2 \theta_s$	$(4.2/7.3) \tan^2 \theta_s$
	D^*K	~ 30	~ 30	~ 1
Fu <i>et al.</i> [43]	D^*K	13 ± 2	42 ± 4	0.31 ± 0.06
Zhang <i>et al.</i> [44]	D^*K	1.65/1.60/1.17	5.58/15.36/16.26	0.30/0.10/0.07
	$c\bar{s} + D^*K$	1.74/2.19/2.75	0.98/6.41/6.99	1.78/0.34/0.39
Lutz <i>et al.</i> [45]	hadrogenesis	21.8/12.47	44.50/45.14	0.49/0.28

Radiative decay of $D_{s0}(2317)^+$ and $D_{s1}(2460)^+$

PRL 136, 241901
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Decay chain: $e^+e^- \rightarrow c\bar{c} \rightarrow D_{s0}(2317)^+ + \text{anythings}$

$D_{s0}(2317)^+ \rightarrow D_s^+[\phi\pi, K^*K] \pi^0$ **Reference channel**

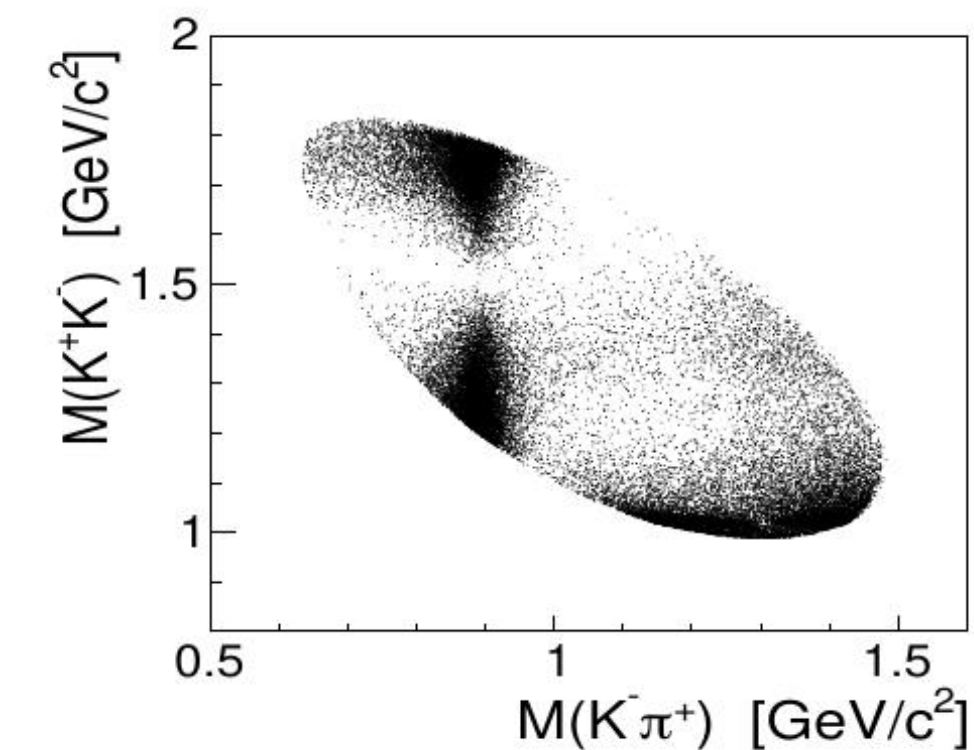
$\rightarrow D_s^{*+}[D_s^+\gamma]\gamma$ **Signal Mode**

Datasets:

- All Belle data: 980 fb⁻¹
- All Belle II data: 428 fb⁻¹

Signal MC:

- $e^+e^- \rightarrow c\bar{c}$: Pythia
- $D_s(2317)^+ \rightarrow D_s^+\pi^0, D_s^{*+}\gamma$: PHSP (S-Wave)
- $D_s^{*+} \rightarrow D_s^+\gamma$: VSP
- $D_s^+ \rightarrow KK\pi$: Dalitz



Pre-selection:

Photon:

- π^0 veto: the invariant mass of any combination of one signal γ_{sig} with one γ_{roe} in rest of event **should not** in the region of $|M(\gamma_{sig}\gamma_{roe}) - m(\pi^0)| < 15 \text{ MeV}/c^2$

Radiative decay of $D_{s0}(2317)^+$

The correction of x_p distribution: get the correct efficiency

$$xp = \frac{p_{D_s^* \gamma}^*}{\sqrt{E_{beam}^2 - M_{D_s^* \gamma}^2}}$$

Dots:

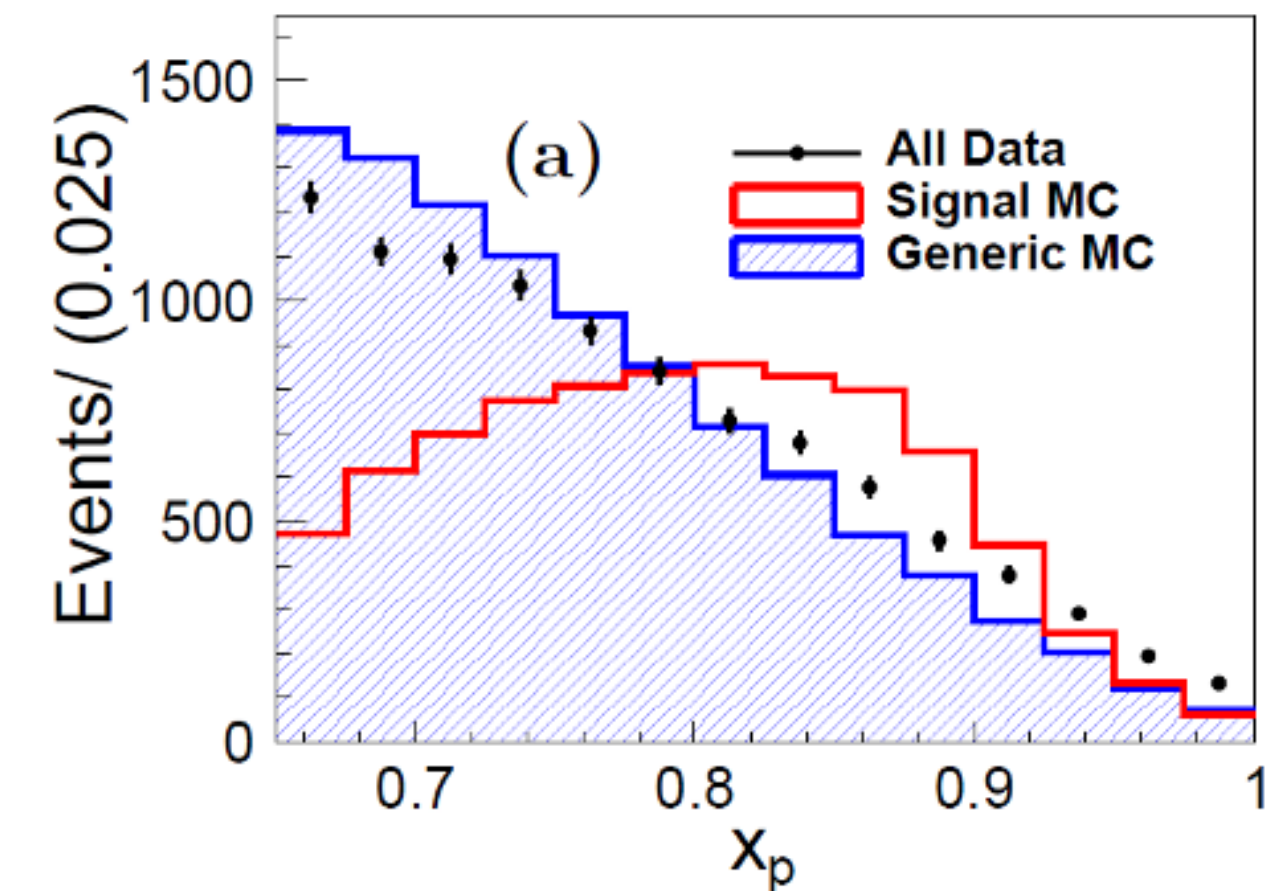
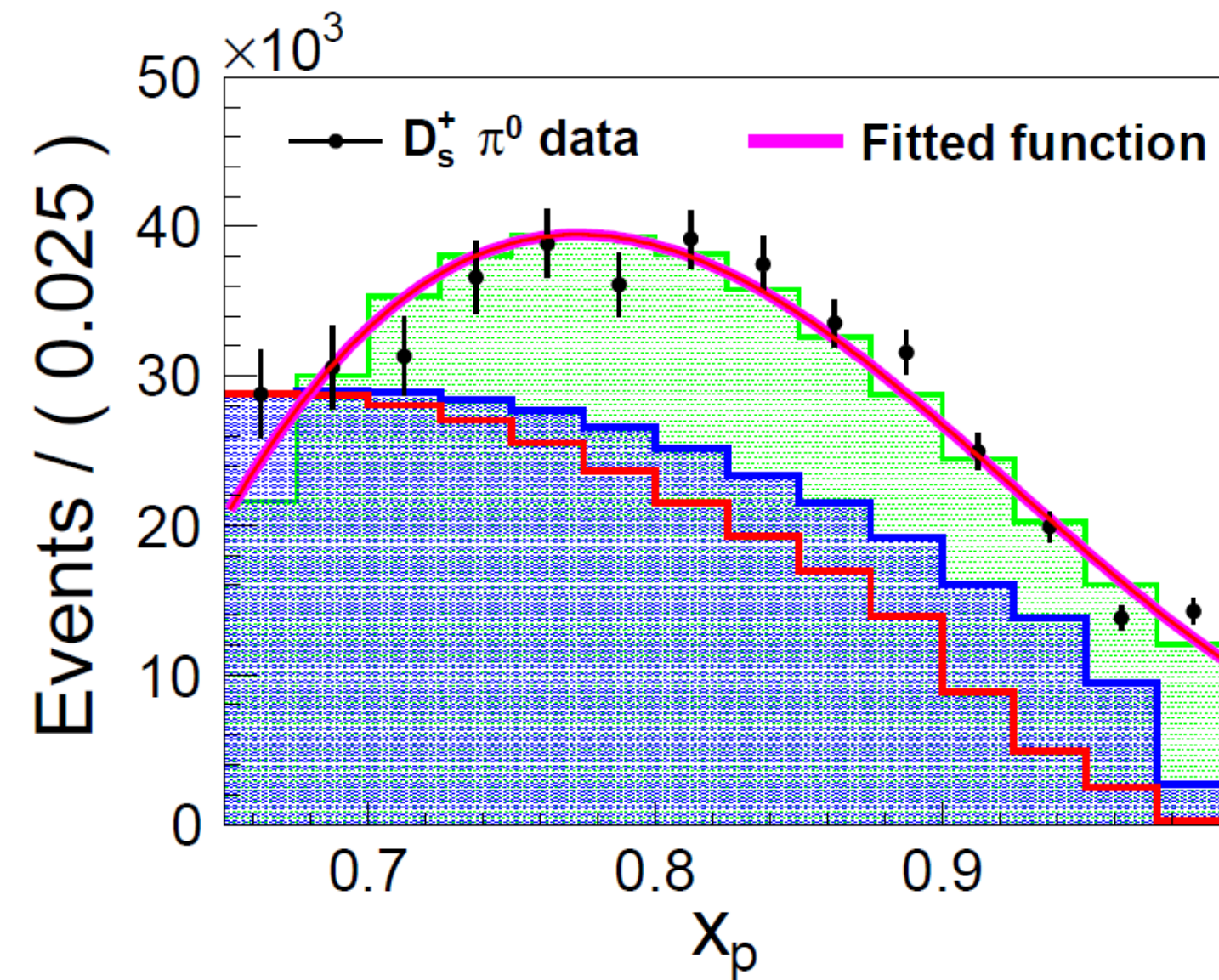
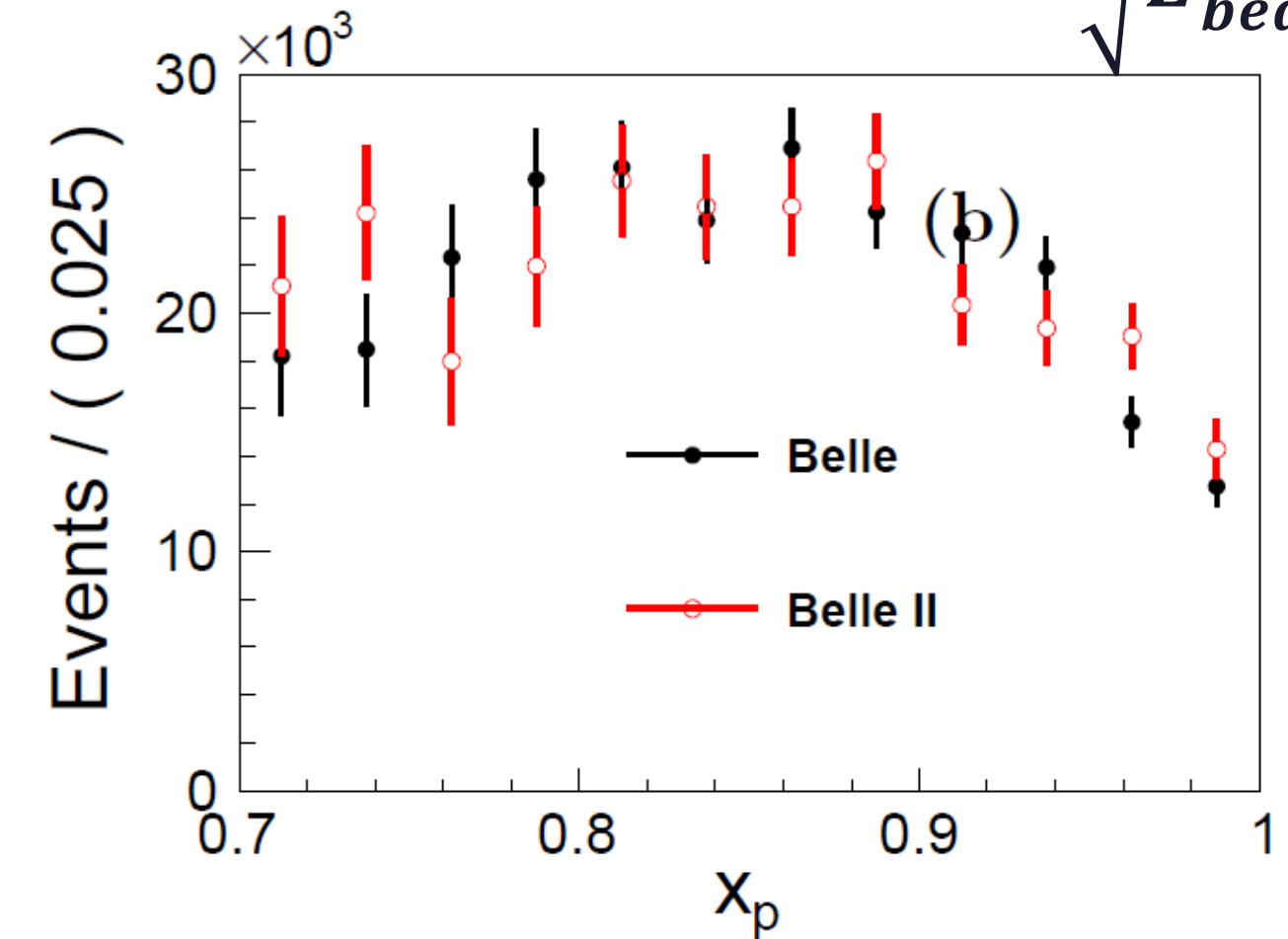
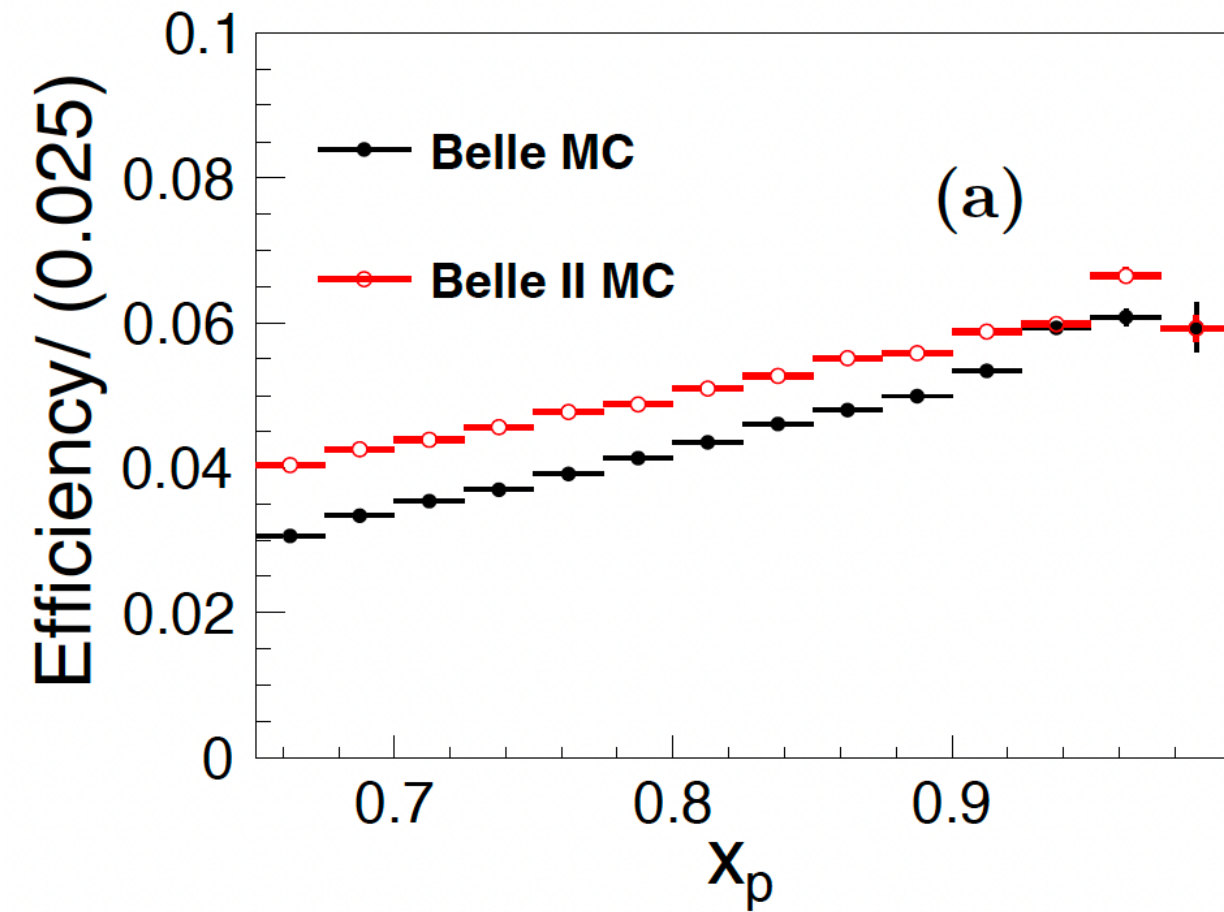
Efficiency-corrected from
combined Belle and Belle II data

Histograms:

Simulated distributions from

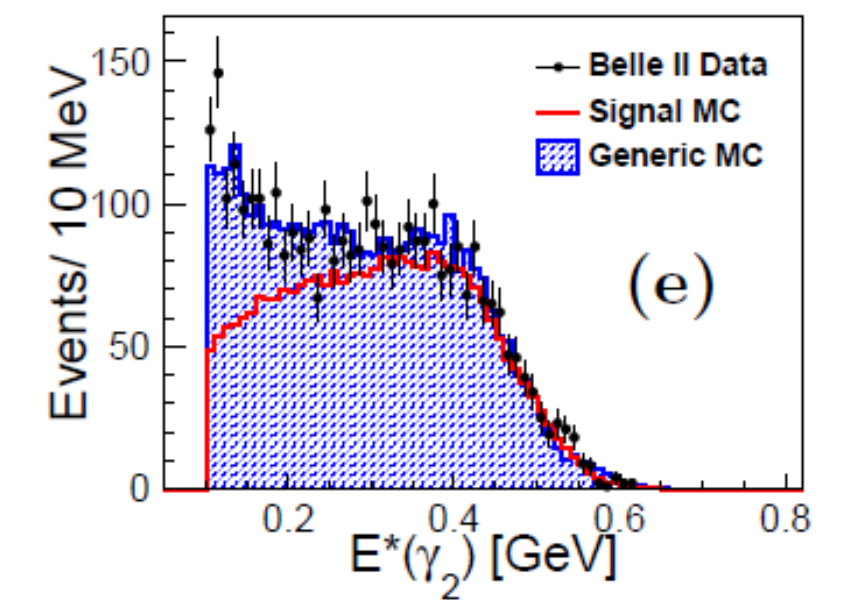
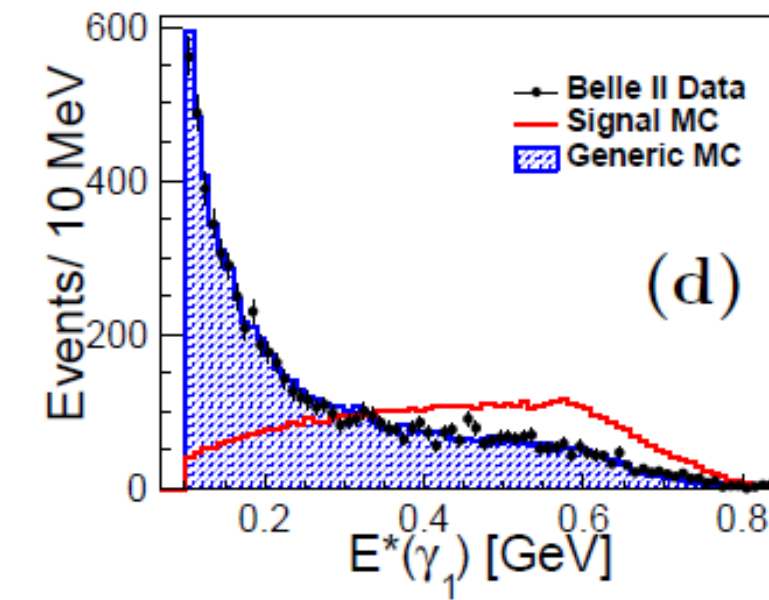
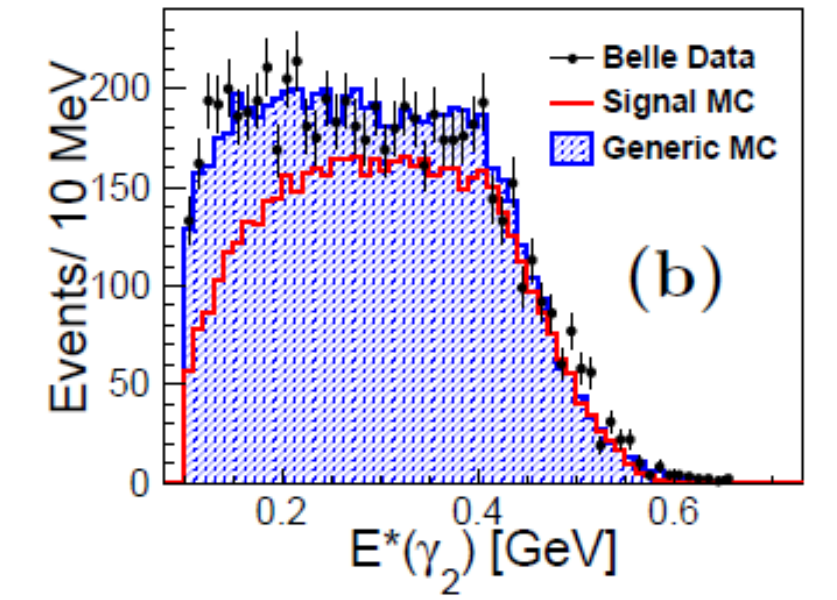
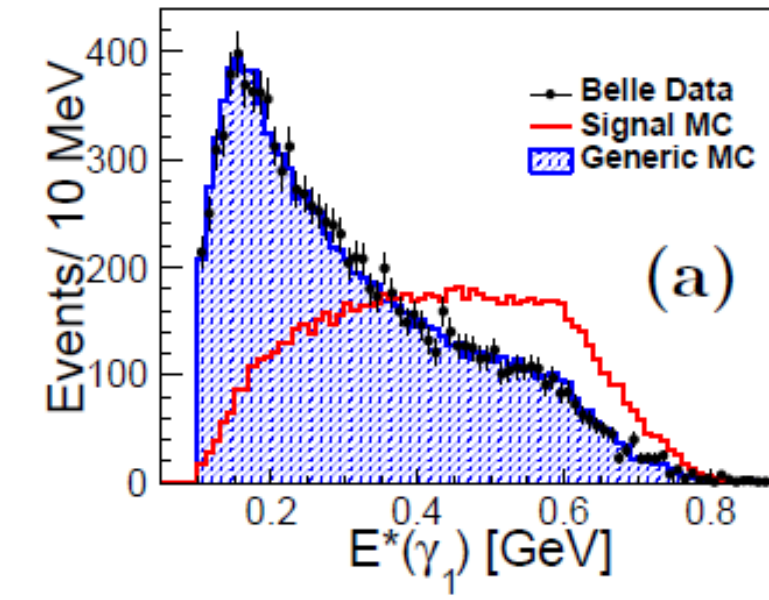
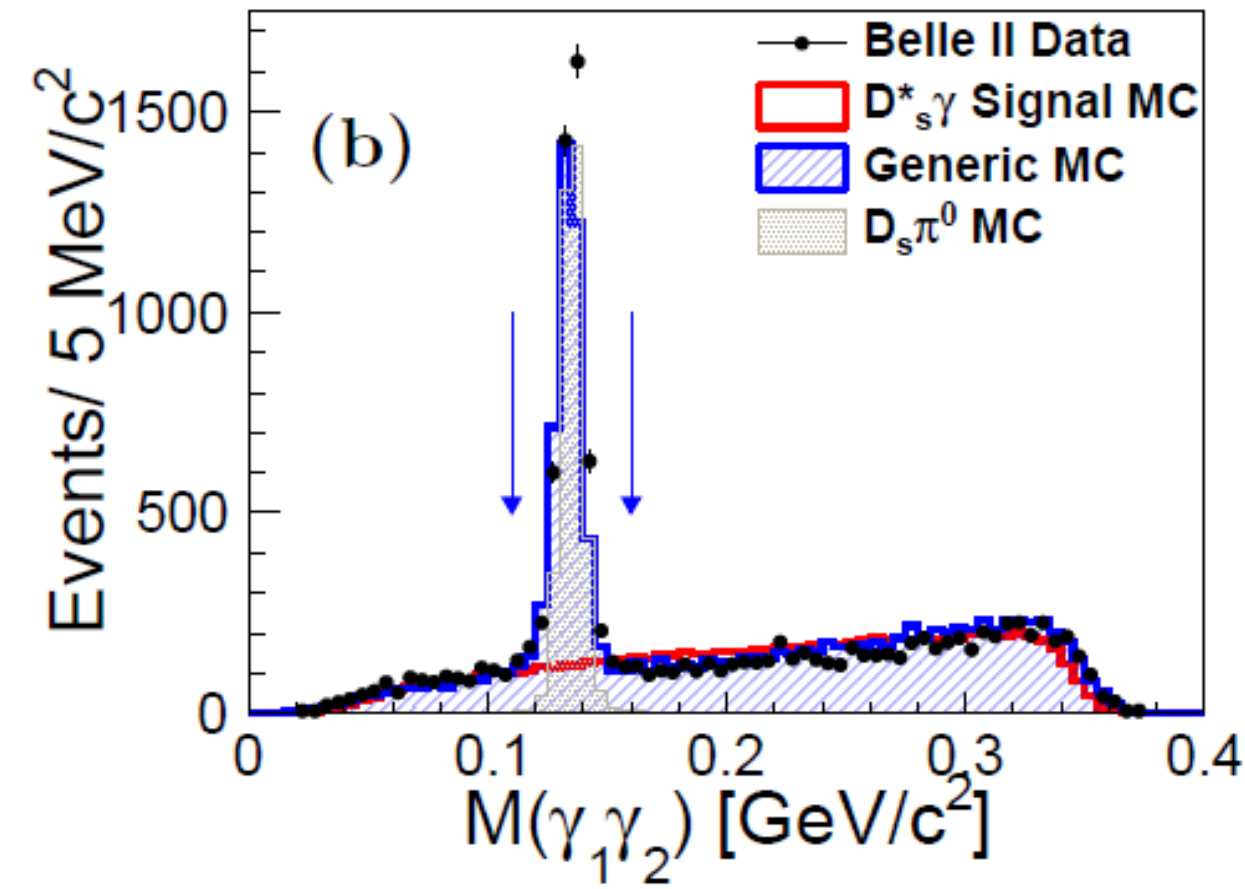
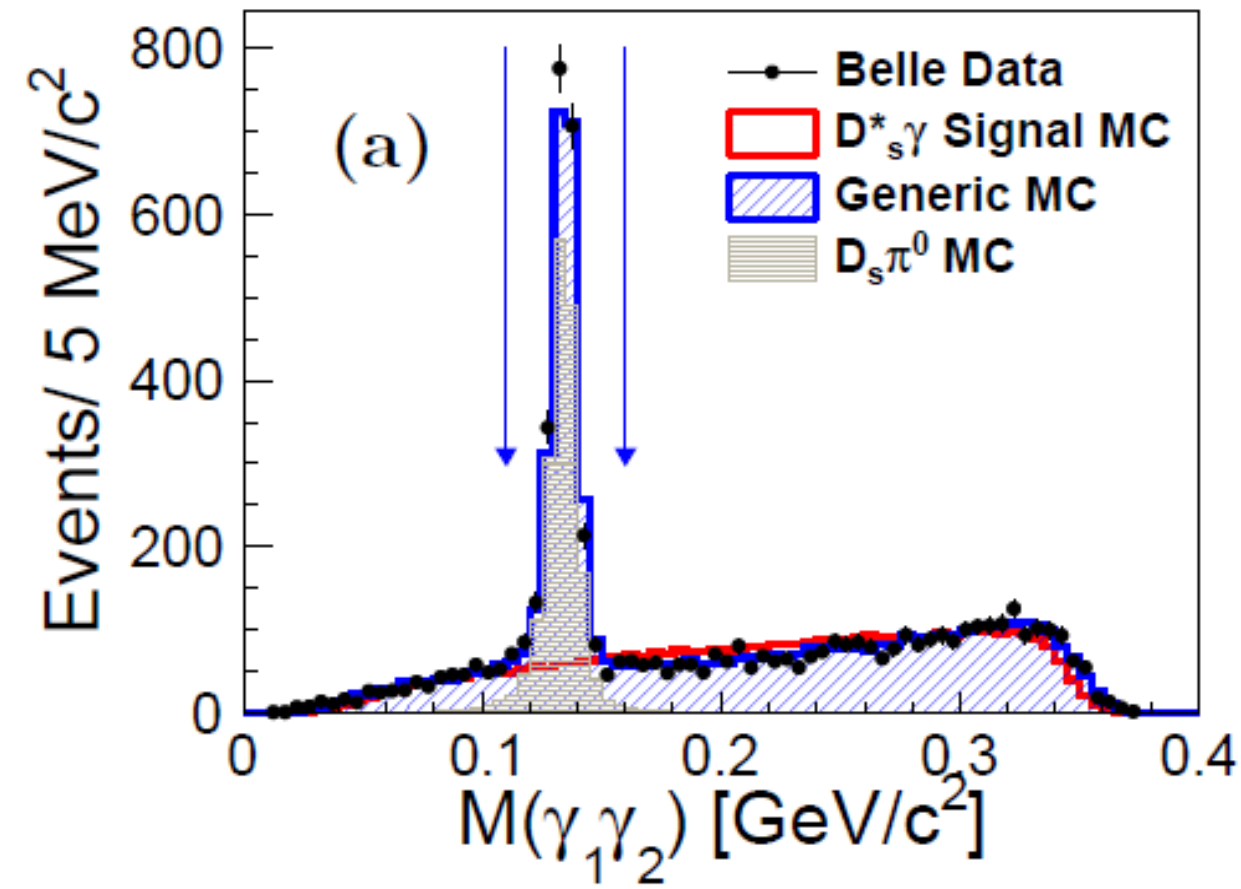
- ◆ Original Belle MC samples
- ◆ Original Belle II MC samples
- ◆ Reweighted MC samples

Only background from $e^+e^- \rightarrow q\bar{q}$
Apply cut: $x_p > 0.7$



Radiative decay of $D_{s0}(2317)^+$

γ selection: $D_s(2317)^+ \rightarrow [D_s^{*+} \rightarrow D_s^+ \gamma_2] \gamma_1$



Veto events in the region:

$$0.100 < |M(\gamma_1\gamma_2)| < 0.160 \text{ GeV}/c^2$$

To suppress the contributions from the $D_s(2317)^+ \rightarrow D_s^+\pi^0$ decay

Apply:

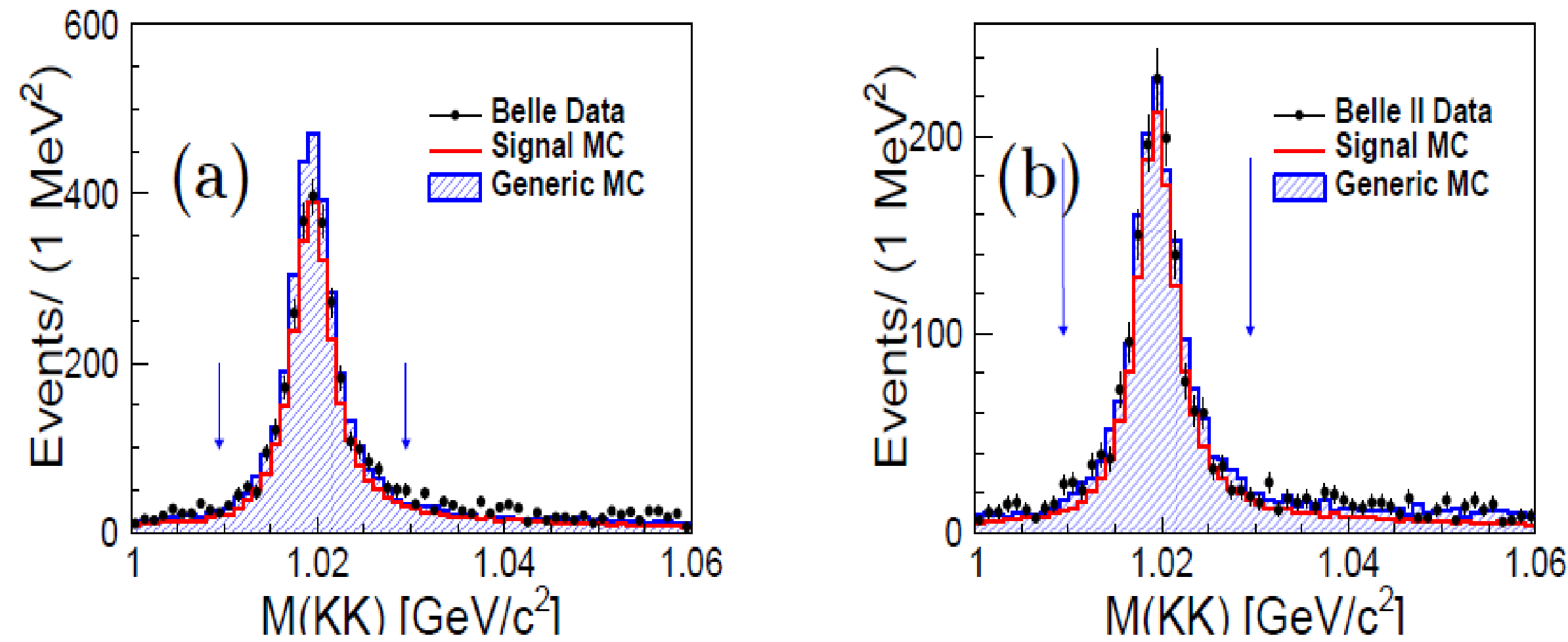
$$E^*(\gamma_1) > 0.22 \text{ GeV}$$

$$E^*(\gamma_2) > 0.10 \text{ GeV}$$

Radiative decay of $D_{s0}(2317)^+$

PRL 136, 241901
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ϕ selection:



Apply:

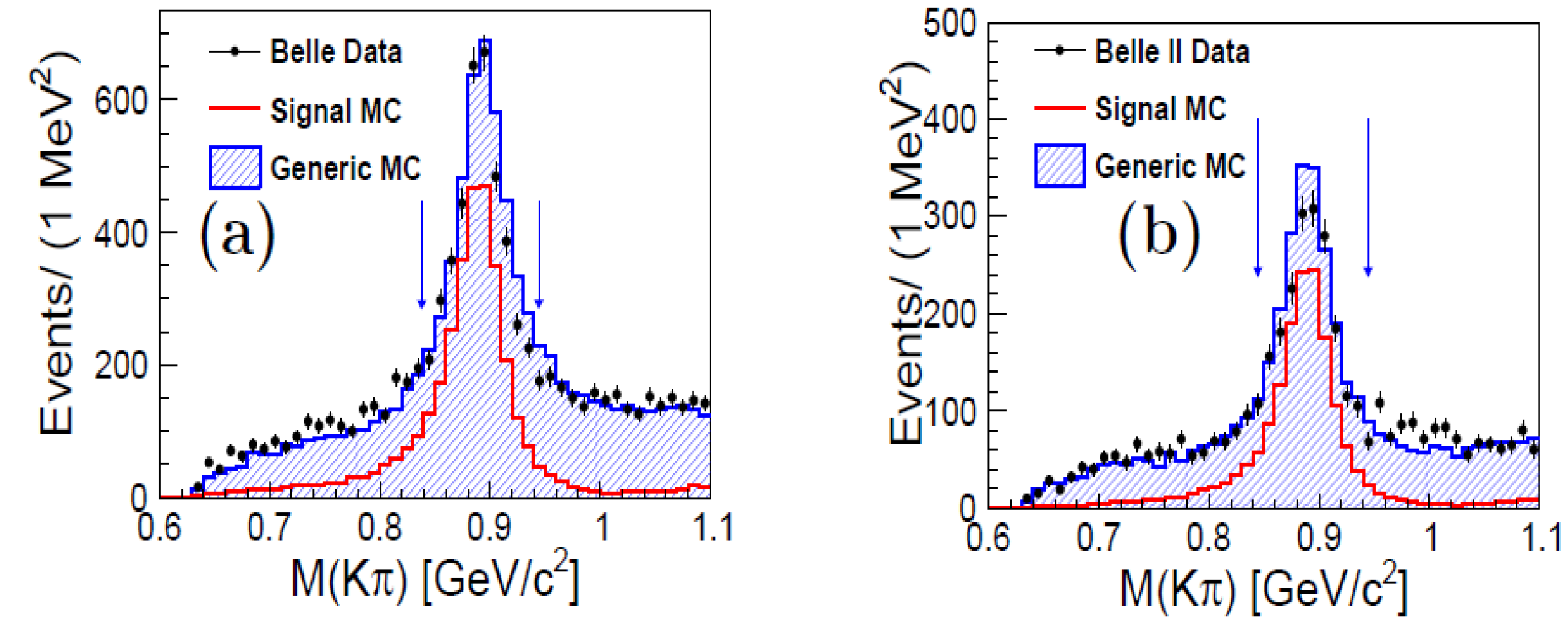
$$|M(K^+K^-) - m(\phi)| < 10 \text{ MeV}/c^2$$

D_s^+ selection:

Mass window :

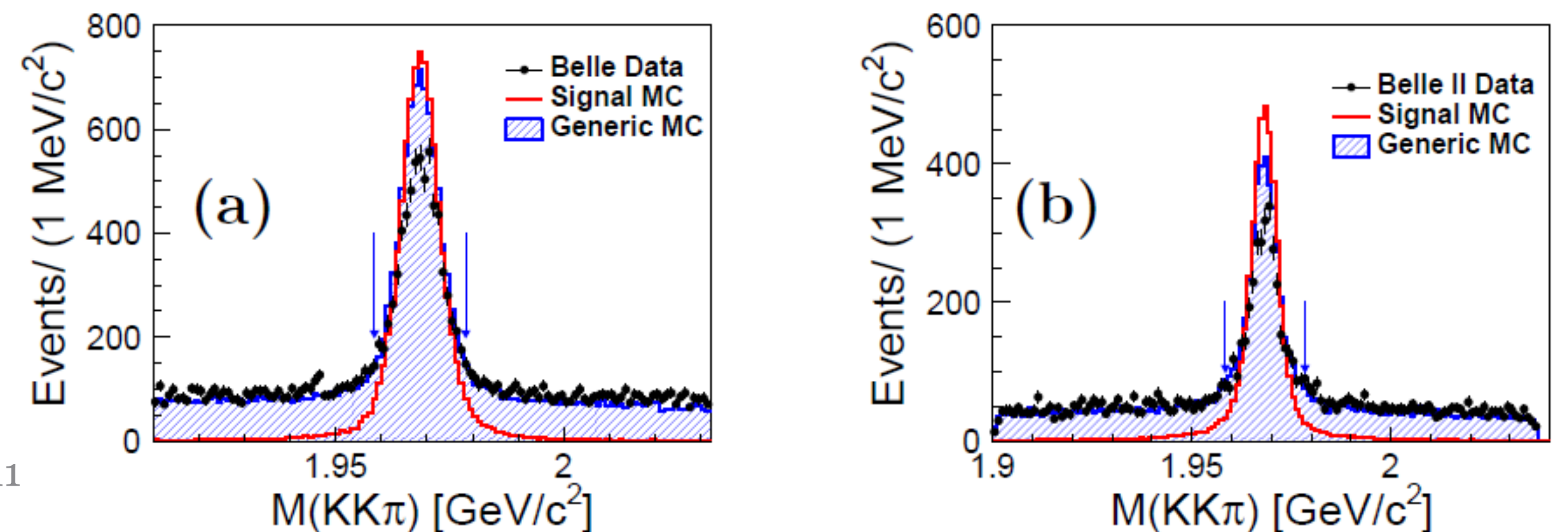
$$|M(K^+K^-\pi^+) - m(D_s^+)| < 10 \text{ MeV}/c^2$$

$\bar{K}^{*0}$ selection:



Apply:

$$|M(K\pi) - m(\bar{K}^{*})| < 50 \text{ MeV}/c^2$$

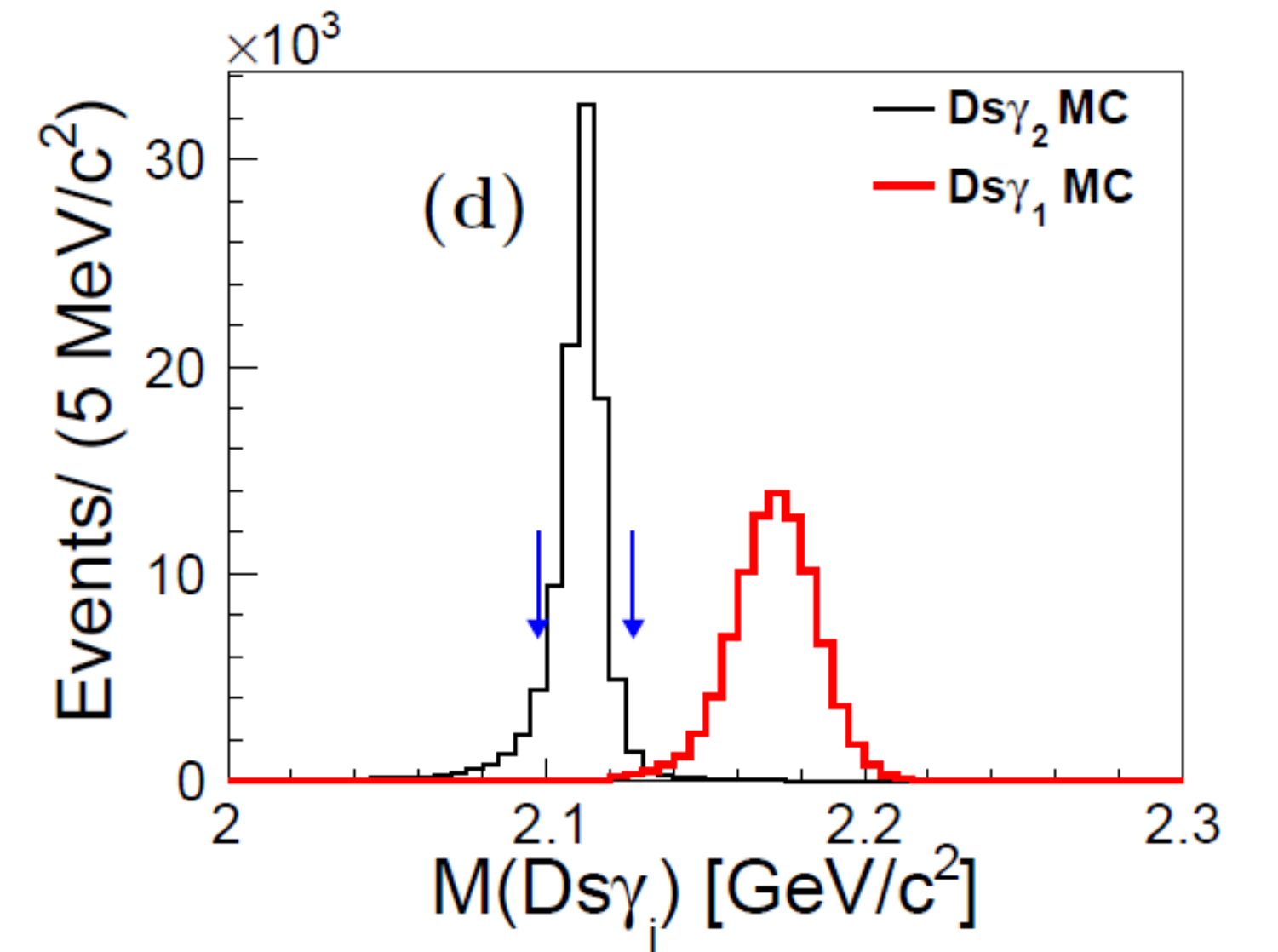
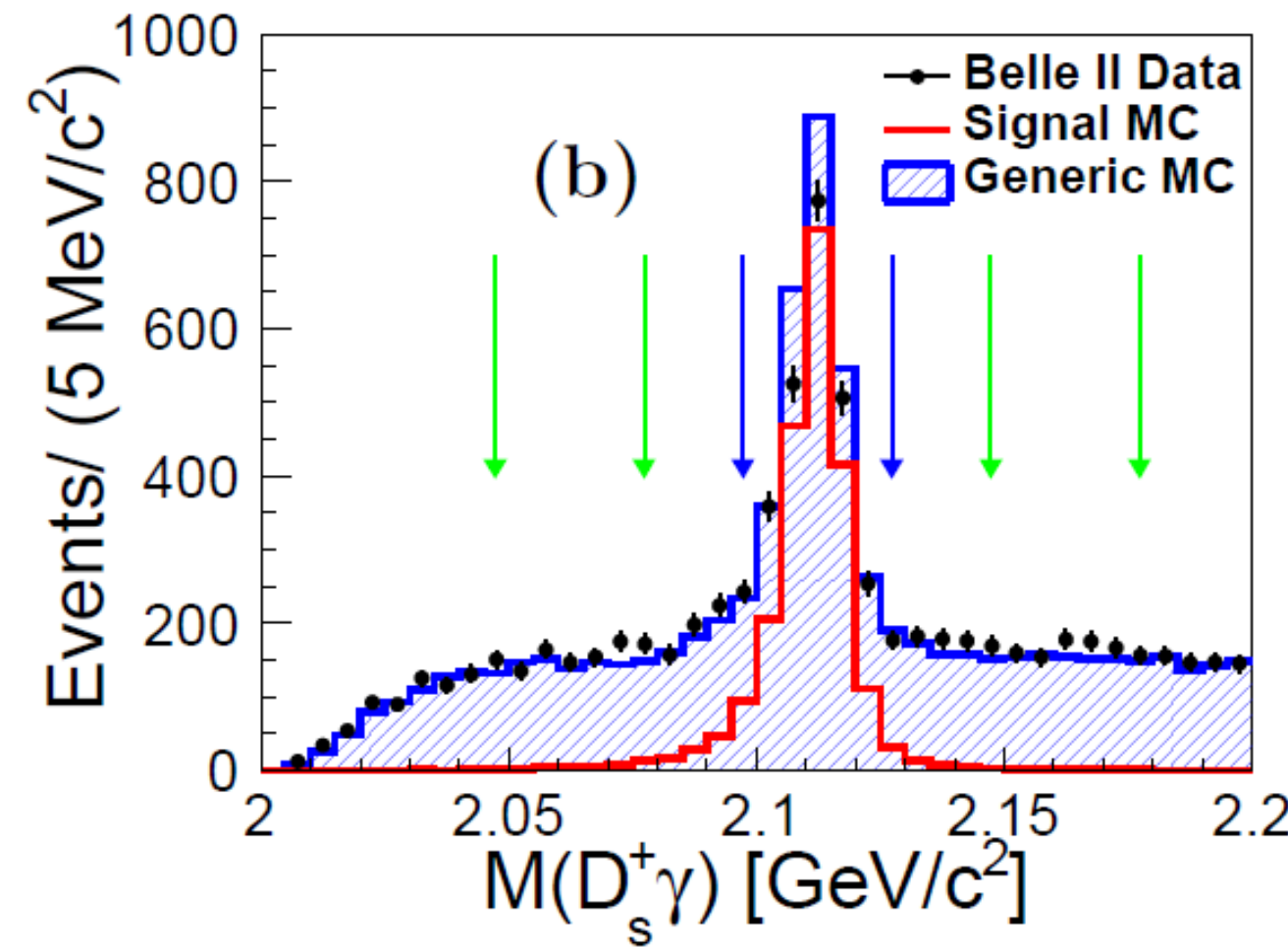
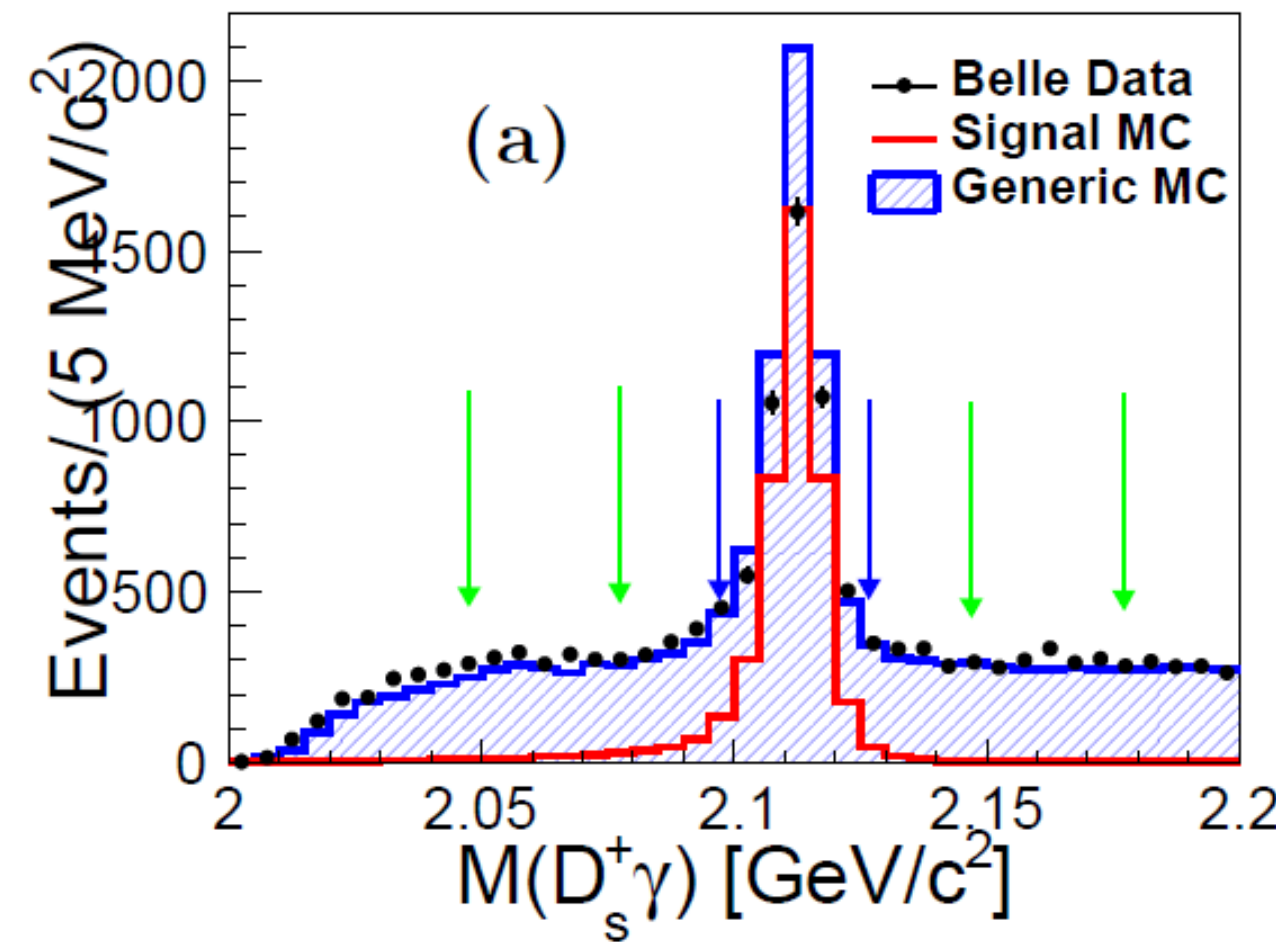


D_s^{*+} selection:

$$D_s(2317)^+ \rightarrow [D_s^{*+} \rightarrow D_s^+ \gamma_2] \gamma_1$$

To reduce the resolution of $m(D_s^+)$:

$$M(D_s^+ \gamma) = M^{rec}(D_s^+ \gamma) - M(K^+ K^- \pi^+) + m(D_s^+)$$



Signal window:

$$|M(D_s^+ \gamma) - m(D_s^{*+})| < 0.015 \text{ GeV}/c^2$$

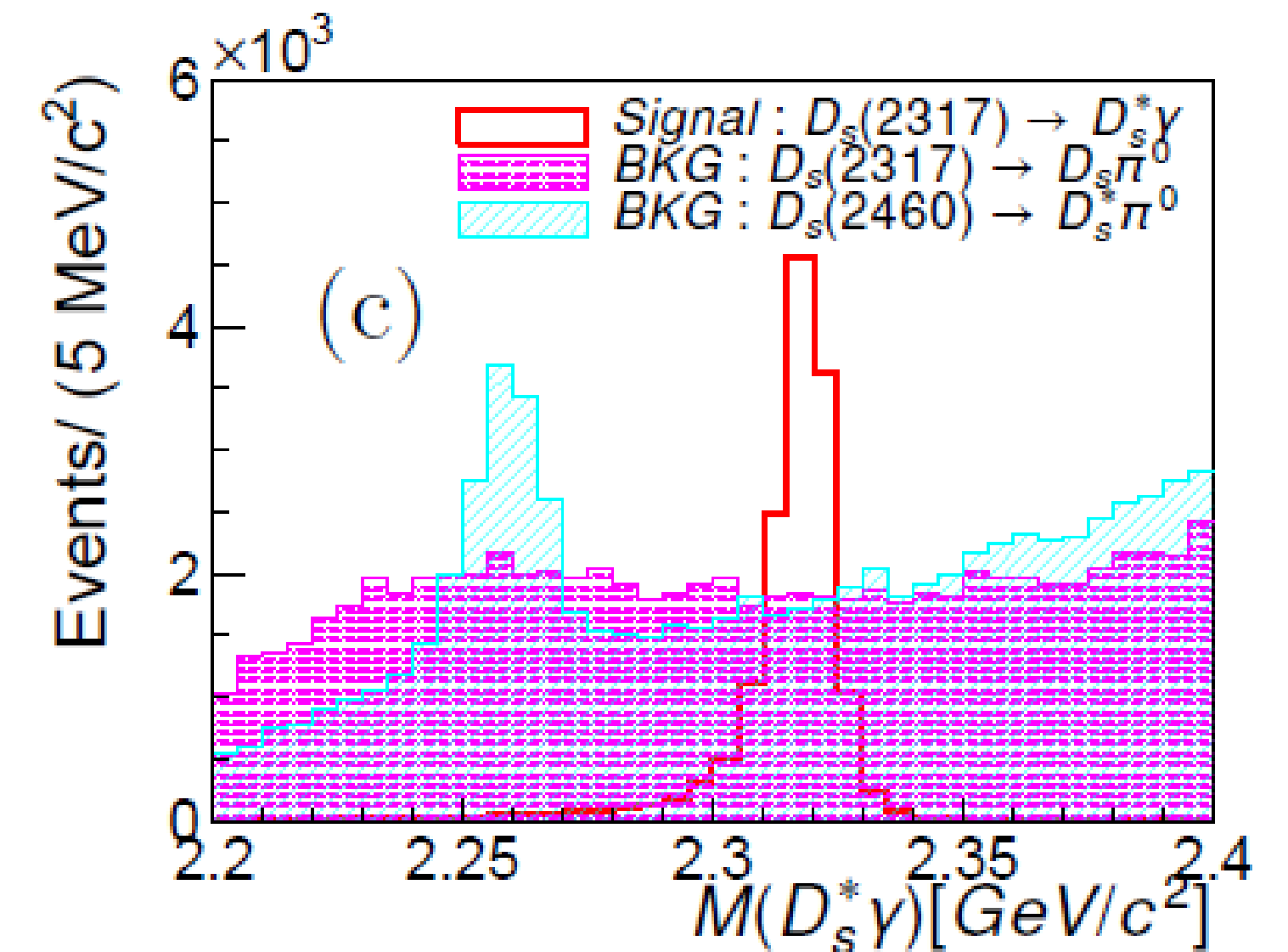
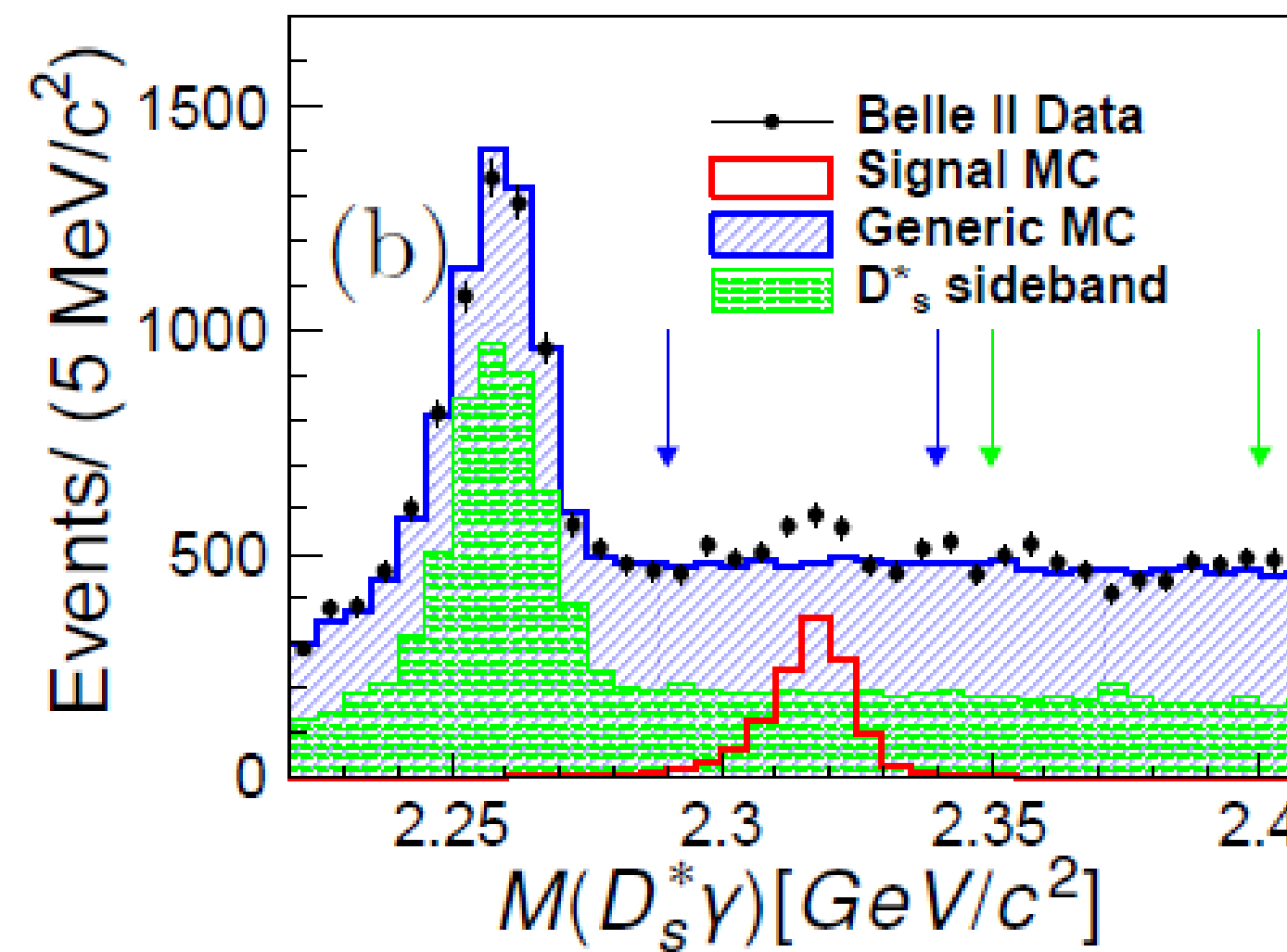
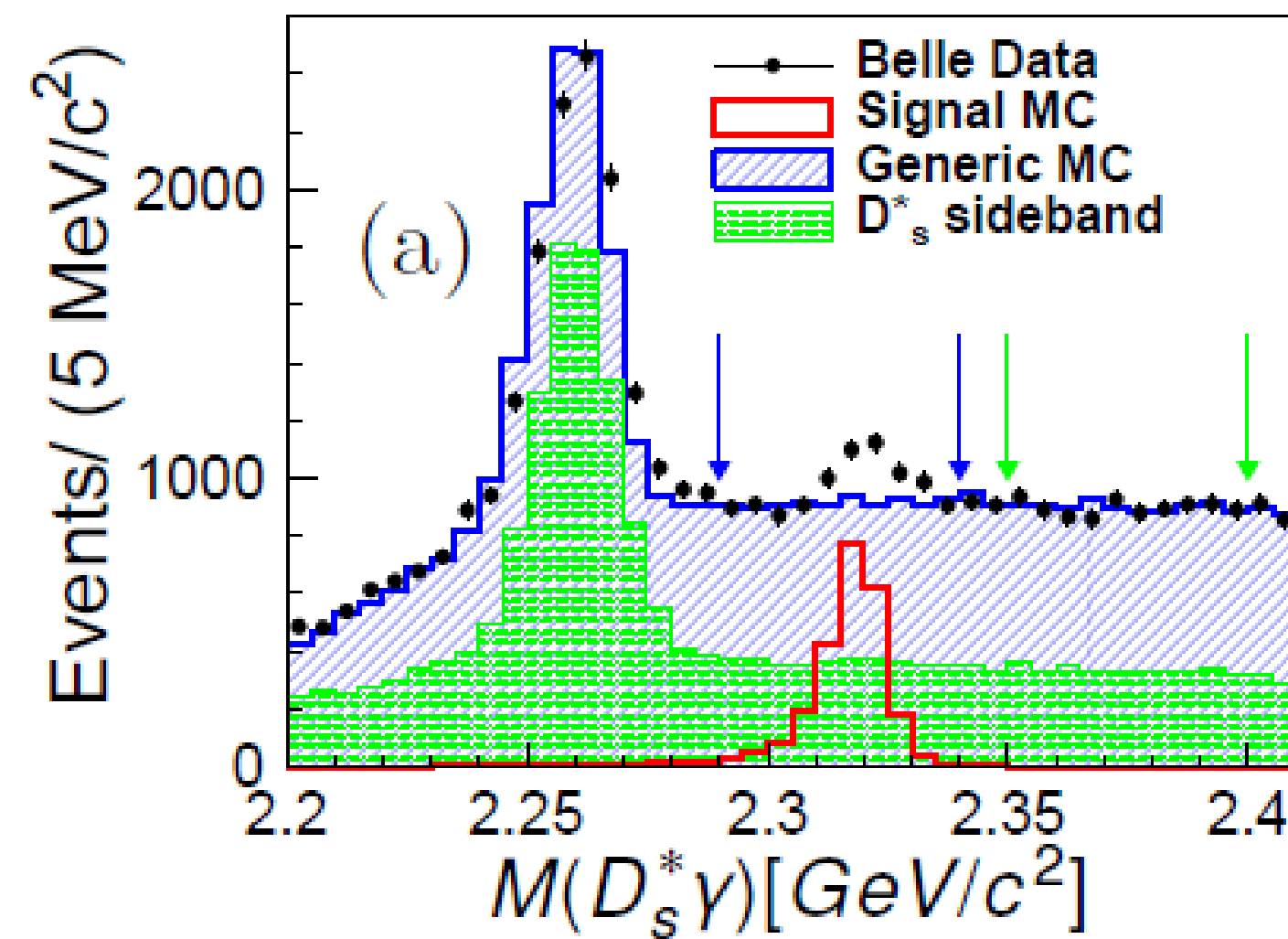
Sideband:

$$|M(D_s^+ \gamma) - m(D_s^{*+}) \pm 0.05| < 0.015 \text{ GeV}/c^2$$

Radiative decay of $D_{s0}(2317)^+$

$$D_s(2317)^+ \rightarrow [D_s^{*+} \rightarrow D_s^+ \gamma_2] \gamma_1$$

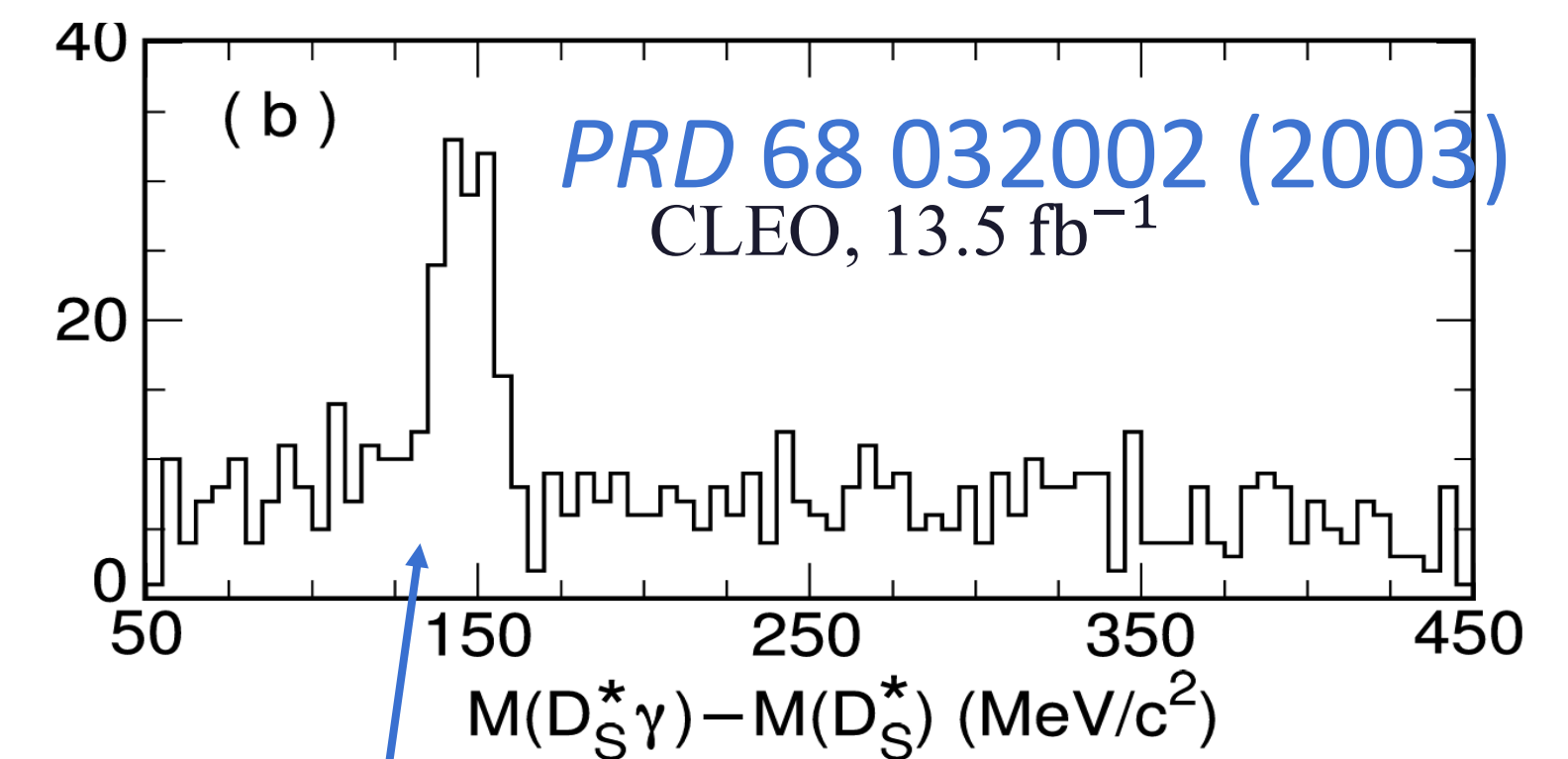
$D_s(2317)$ distributions:



To reduce the resolution of $M(D_s^{*+})$:

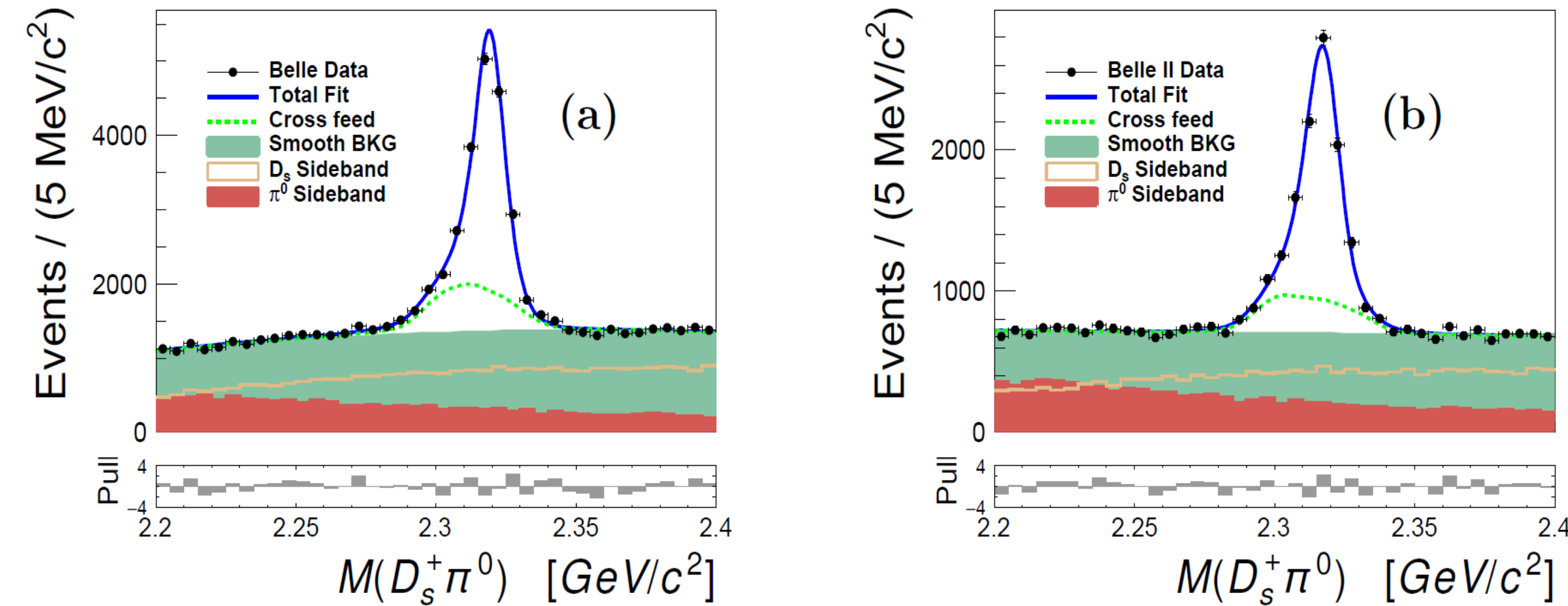
$$M(D_s^{*+} \gamma) = M^{rec}(D_s^{*+} \gamma_1) - M^{rec}(D_s^+ \gamma_2) + m(D_s^{*+})$$

- **Clear $D_s(2317)$ signal peaks**
- **Flat background:** from both D_s^* sideband and generic MC
- Peaks around $2.26 \text{ GeV}/c^2$ are from $D_s^* + \gamma_{\text{free}}$ (same as CLEO)



The same conclusion holds for the $D_s^* \gamma$ final state, shown in Fig. 3(b), where we combine selected D_s^* candidates with photons of energy above 150 MeV. The peak in the $\Delta M(D_s^* \gamma)$ spectrum in Fig. 3(b) near 150 MeV/c² is due to real $D_s^{*+} \rightarrow D_s^+ \gamma$ decays in which a random photon has been combined with the D_s^* candidate to form the D_s^* candidate, and the actual photon from this transition is combined with this system to form the D_{sJ}^* candidate. There is no sign of any structure in this spectrum near 205 MeV/c² where a signal from $D_s^* (2317)$

Radiative decay of $D_{s0}(2317)^+$



Separate fit to extract $D_s(2460)^+ \rightarrow D_s \pi^0$ yields

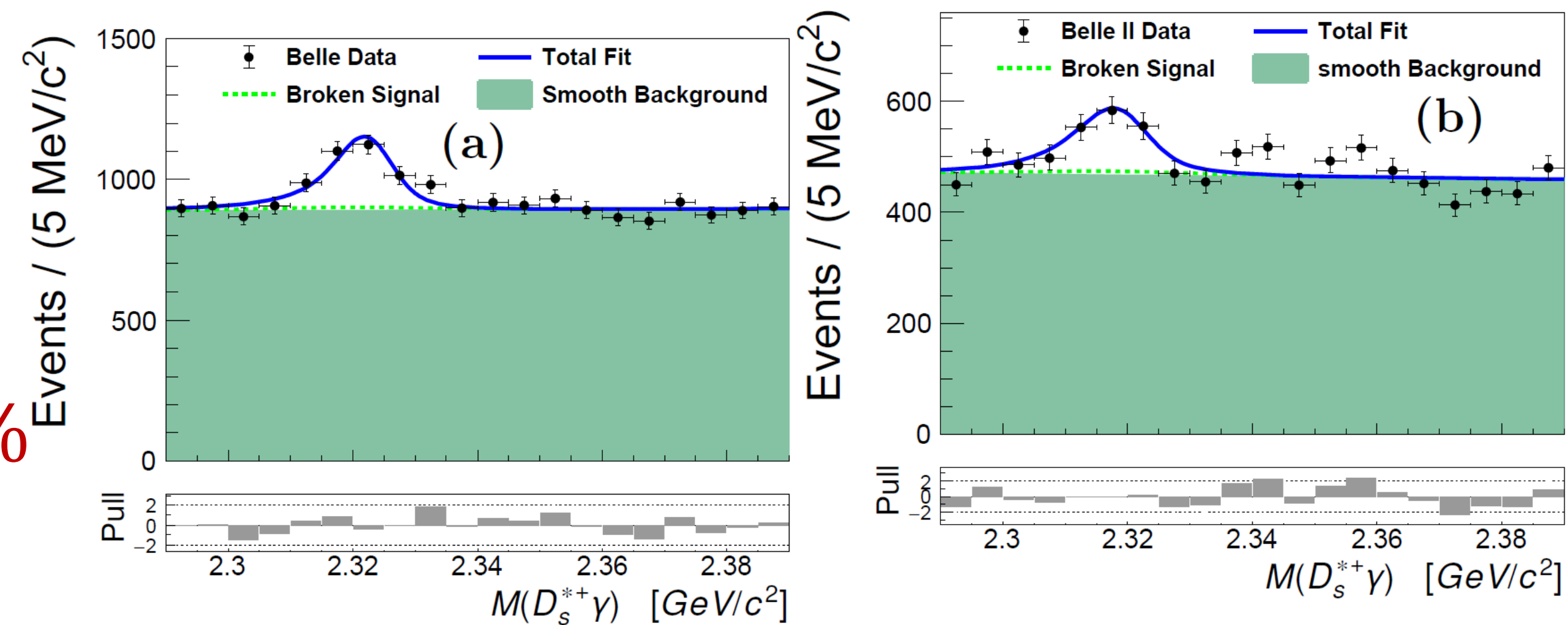
- Signal: Crystal Ball convoluted with triple-Gaussian
- Broken Signal: asymmetry Gaussian
- Smooth Background: 2nd order polynomial
- cross feed: $D_s(2460)^+ \rightarrow D_s^{*+} \pi^0$, (i.e. missing one photon)

Simultaneous fit to extract $\mathcal{R}$

$$\bullet N_{\text{exp}}(D_s^{*+} \gamma) = \mathcal{R} \frac{N_{\text{exp}}^{\text{fit}}(D_s^+ \pi^0) \varepsilon_{\text{exp}}(D_s^{*+} \gamma)}{\varepsilon_{\text{exp}}(D_s^+ \pi^0)},$$

exp represents Belle or Belle II.

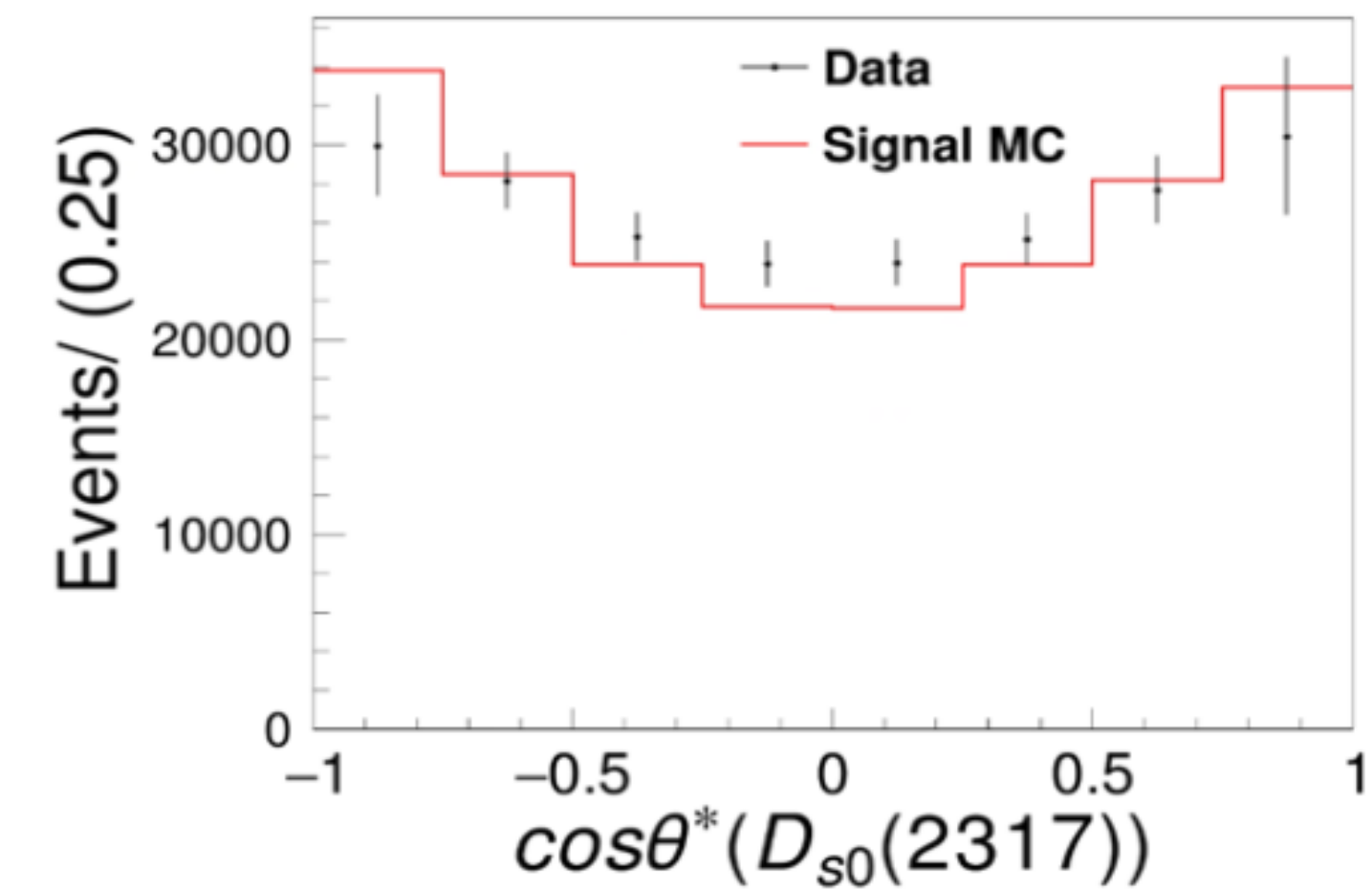
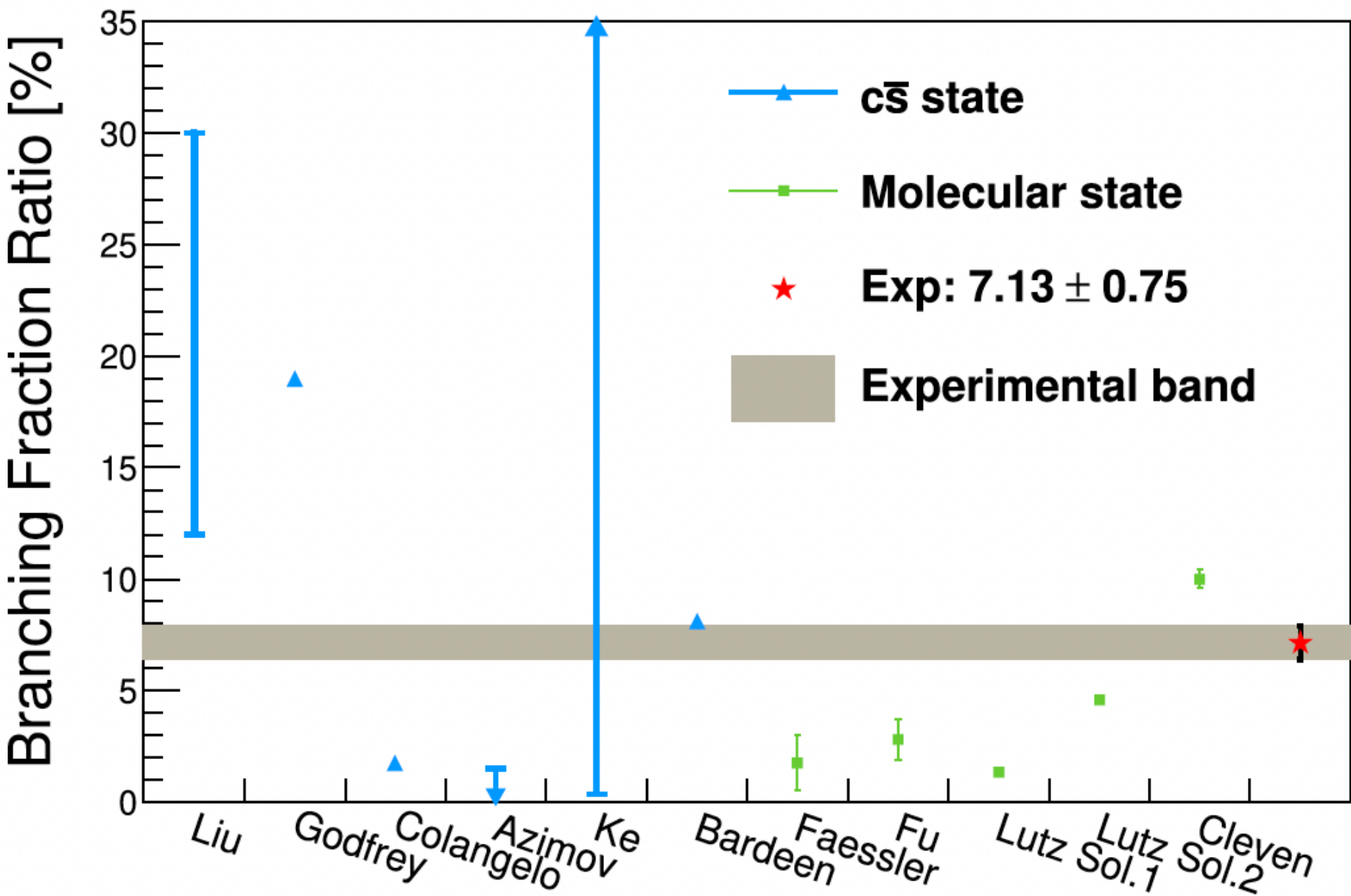
- $\mathcal{R} = [7.13 \pm 0.70(\text{stat.}) \pm 0.23(\text{syst.})]\%$
- Significance: 10.1σ , with syst. Uncertainty included.



Radiative decay of $D_{s0}(2317)^+$

Systematic uncertainty

Source	$D_s^+ \pi^0$	$D_s^{*+} \gamma$
Fit region and background pdf	0.8	1.3
Fixed pdf parameters	0.7	2.5
Cross-feed or broken signal	0.6	0.7
x_p reweighting		0.5
$\cos \theta_{\gamma_1 \gamma_2}$ reweighting		1.7
MC sample size		0.5
Sum		3.6



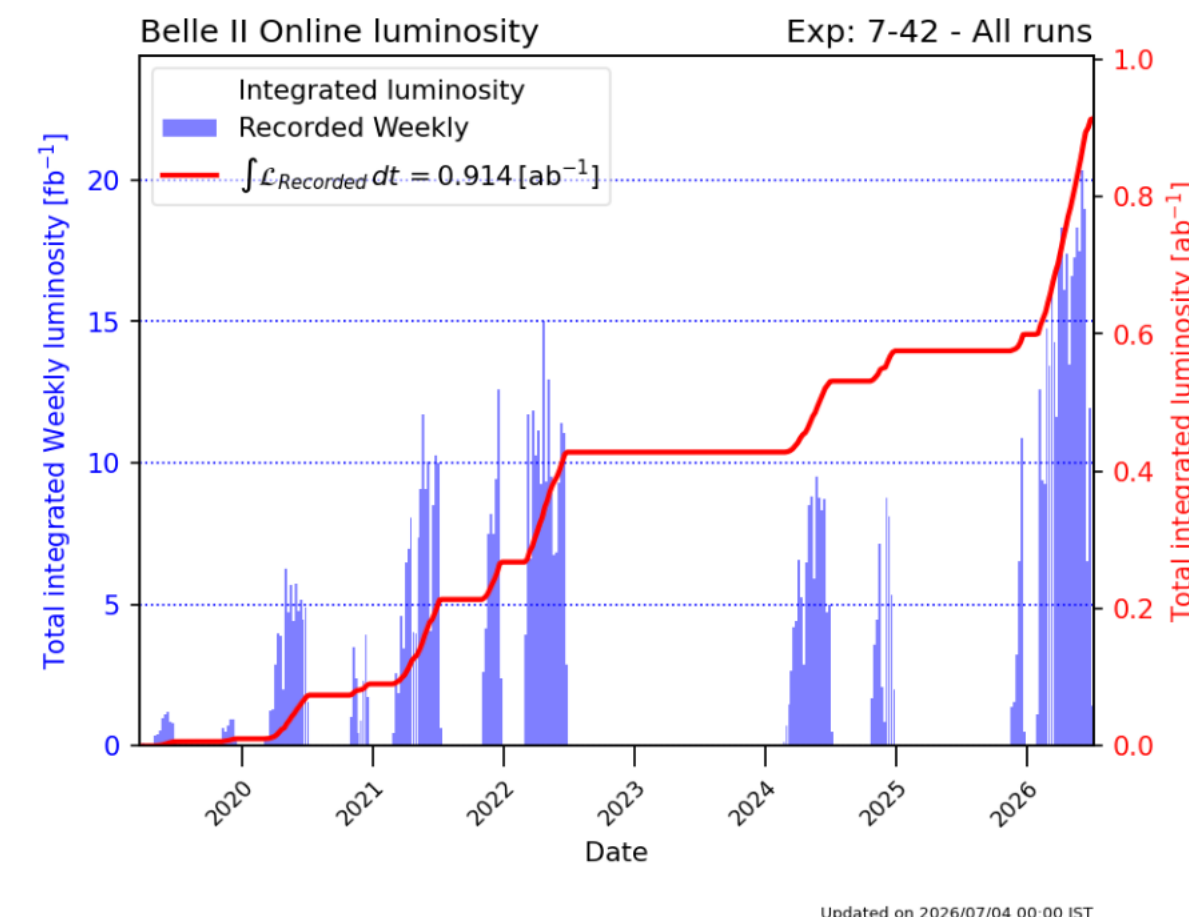
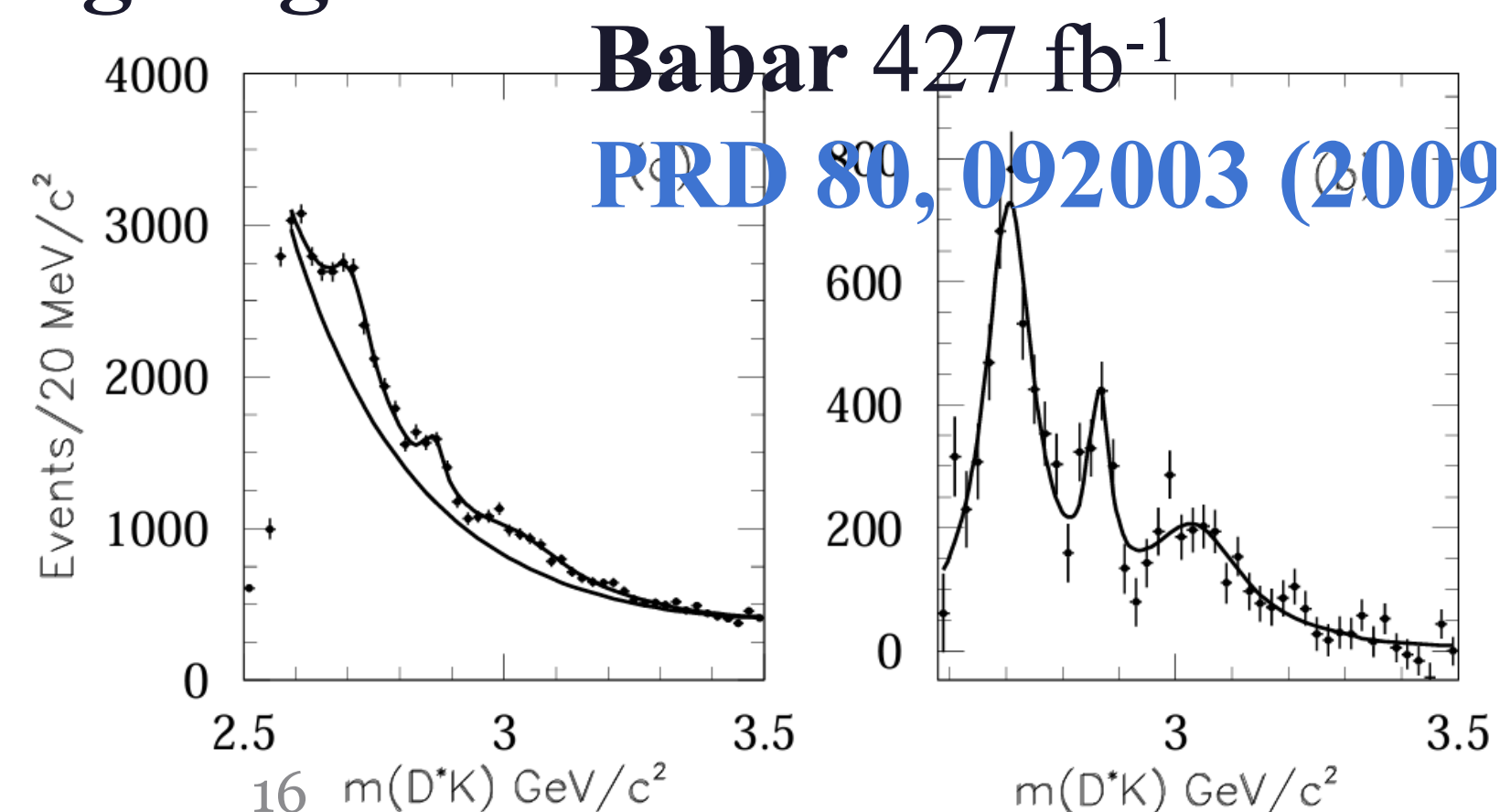
Reweight MC as:

$I(\cos \theta) \propto 1 + \cos \theta_{\gamma\gamma}$, where $\theta_{\gamma\gamma}$ is the angle between the two photons in the D_s^{*+} rest frame.

θ^* : polar angle of the $D_{s0}(2317)^+$ in center-of-mass system

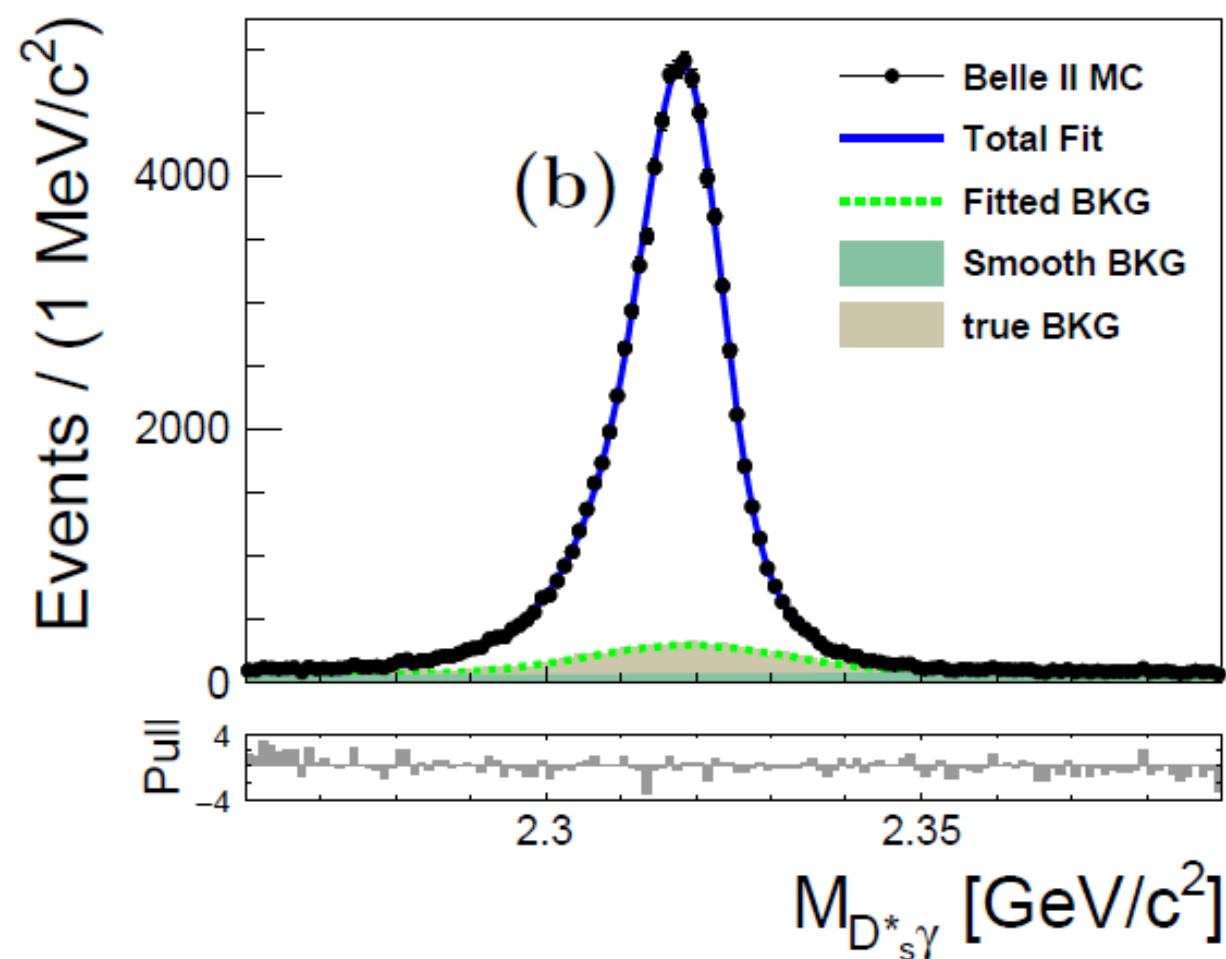
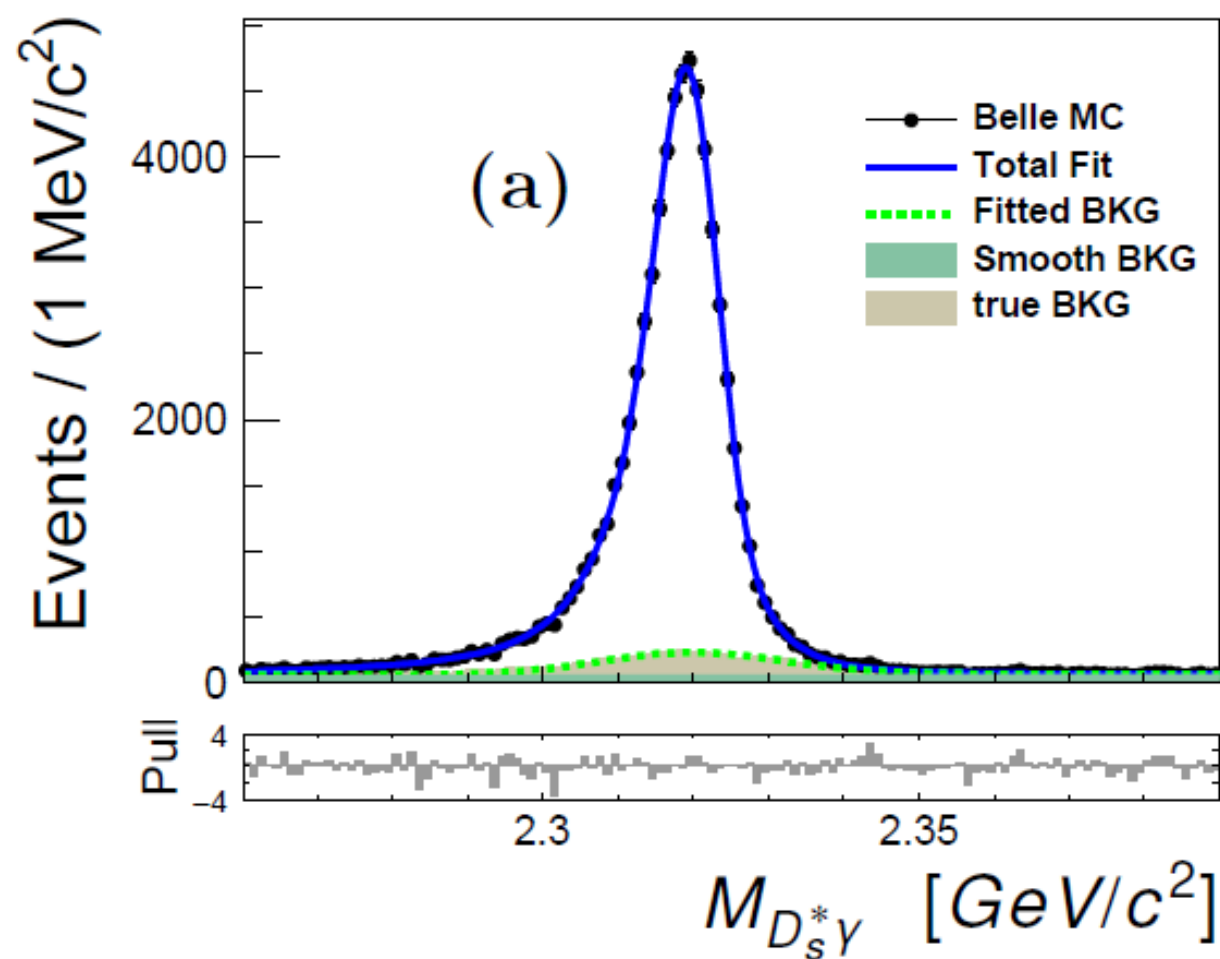
Summary

- Using all Belle data (983 fb⁻¹) and Belle II data (427 fb⁻¹ +xxx)
- Observed $D_{s0}(2317)^+ \rightarrow D_s^{*+} \gamma$ decay for the first time with **10.1 σ**
- $\mathcal{R} = \frac{\mathcal{B}(D_s(2317)^+ \rightarrow D_s^{*+} \gamma)}{\mathcal{B}(D_s(2317)^+ \rightarrow D_s^+ \pi^0)} = [7.14 \pm 0.70(stat.) \pm 0.23(syst.)]\%$
 - uncertainty dominated by Statistical error of $D_s(2317)^+ \rightarrow D_s^{*+} \gamma$
 - higher than the molecular assumption for $D_{s0}(2317)^+$
 - possible interpretation: $D_{s0}(2317)^+$ could be the mixture state of pure $c\bar{s}$ and molecular state
- Analysis of $D_{s1}(22460)^+ \rightarrow D_s^{*+} \gamma$ is ongoing.
- More analysis can be done:
 - $XYZ \rightarrow D_{sJ} XXX$
 - Higher excited states



Radiative decay of $D_{s0}(2317)^+$

$D_s(2317)$ MC distributions



- Signal: Crystal Ball convoluted with triple-Gaussian
- Broken Signal: asymmetry Gaussian
- Smooth Background: 2nd order polynomial
- All parameters free to float in the fits

Experiment s	Belle	Belle II
N^{fit}	$(6.44 \pm 0.02) \times 10^4$	$(8.00 \pm 0.03) \times 10^4$
Truth yields	64,508	80,154

The IO checks are OK for the signal MC samples

Systematic Uncertainty Study

- The systematic uncertainties due to D_s and **photon** selection, and cut on x_p are cancelled

- x_p reweighting:

- Change the order of the polynomial used to fit x_p distribution
- Redo the reweighting to get the new detection efficiency ϵ_{exp}
- Fit $\mathcal{R}$ again

$$N_{\text{exp}}(D_s^{*+}\gamma) = \mathcal{R} \frac{N_{\text{exp}}^{\text{fit}}(D_s^+\pi^0)\epsilon_{\text{exp}}(D_s^{*+}\gamma)}{\epsilon_{\text{exp}}(D_s^+\pi^0)}$$

Uncertainties on $\mathcal{R}$:
Difference in Fit $\mathcal{R}$

- $D_{s0}(2317) \rightarrow D_s\pi^0$ mode:

In the fits to $M(D_s^+\pi^0)$ spectra, By

1. Changing the fit region, order of background polynomial
 2. Enlarging the widths of Gaussians by 1σ (gotten from fits to signal MC)
 3. Using the shapes of $M(D_s\pi)$ from $D_s(2460)^+ \rightarrow D_s^{*+}\pi^0$ MC directly
- Differences in the fitted signal yields: Uncertainties ($\Delta_{\text{exp}}^{N(D_s^+\pi^0)}$) of $N_{\text{exp}}^{\text{fit}}(D_s^+\pi^0)$

Then, in the simultaneous fit to $M(D_s^{*+}\gamma)$ spectra,

By varying the $N_{\text{exp}}^{\text{fit}}(D_s^+\pi^0)$ based on Gaussian($N_{\text{exp}}^{\text{fit}}(D_s^+\pi^0)$, $\Delta_{\text{exp}}^{N(D_s^+\pi^0)}$) in the fits

Widths of the distributions of fitted values of $\mathcal{R}$:
Uncertainties on $\mathcal{R}$:

Systematic Uncertainty Study

- $D_{s0}(2317) \rightarrow D_s^* \gamma$ mode:

In the simultaneous fit to $M(D_s^{*+} \gamma)$ spectra, by:

1. Changing the fit region, order of background polynomial
2. Enlarging the widths of Gaussian by 1σ (gotten from fits to signal MC)
3. Changing resolutions of broken signal by 2σ
4. Changing ratios of broken signal to signal yields by 2σ

➤ Fit $\mathcal{R}$ again

Uncertainties on $\mathcal{R}$:
Difference in Fit $\mathcal{R}$

Source	$D_s^+ \pi^0$	$D_s^{*+} \gamma$
Fit region and background pdf	0.8	1.3
Fixed PDF parameters	0.7	2.5
Cross-feed or broken signal	0.6	0.7
x_p reweighting	0.5	
MC statistic	0.5	
Sum	3.2	

**Assuming all systematic
uncertainties are
independent**