



Search for the Strange Dibaryons with Baryon-Baryon Correlations at STAR

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I. Motivation

II. Femtoscopy

III. Analysis Details

IV. Lednicky-Lyuboshitz Model

V. Results

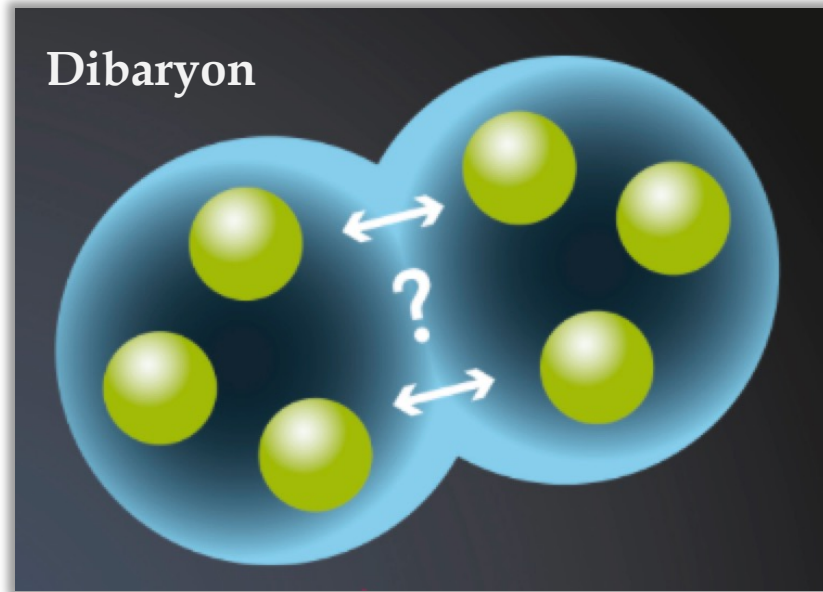
- $|S| = 2$: p - Ξ^- Correlation
- $|S| = 3$: p - Ω^- Correlation

VI. Spin Femtoscopy Method

VII. Summary

Motivation

(Strange) Dibaryons have never been found experimentally



Particle	Mass (MeV)	Quark composition	Decay mode
f_0	980	$q\bar{q}s\bar{s}$	$\pi\pi$
a_0	980	$q\bar{q}s\bar{s}$	$\pi\eta$
$K(1460)$	1460	$q\bar{q}q\bar{s}$	$K\pi\pi$
$\Lambda(1405)$	1405	$qqqs\bar{q}$	$\pi\Sigma$
$\Theta^+(1530)$	1530	$qqqq\bar{s}$	KN
H	2245	$uuddss$	$\Lambda\Lambda$
$N\Omega$	2573	$qqqsss$	$\Lambda\Xi$
$\Xi\Xi$	2627	$qqssss$	$\Lambda\Xi$
$\Omega\Omega$	3228	$ssssss$	$\Lambda K^- + \Lambda K^-$

Sungtae Cho, et al. (ExHIC), Phys. Rev. C 84, 064910 (2011)

➤ The possible formation channels:

$$(|S|=2) \text{ Dibaryon} \Leftrightarrow p + \Xi^-$$

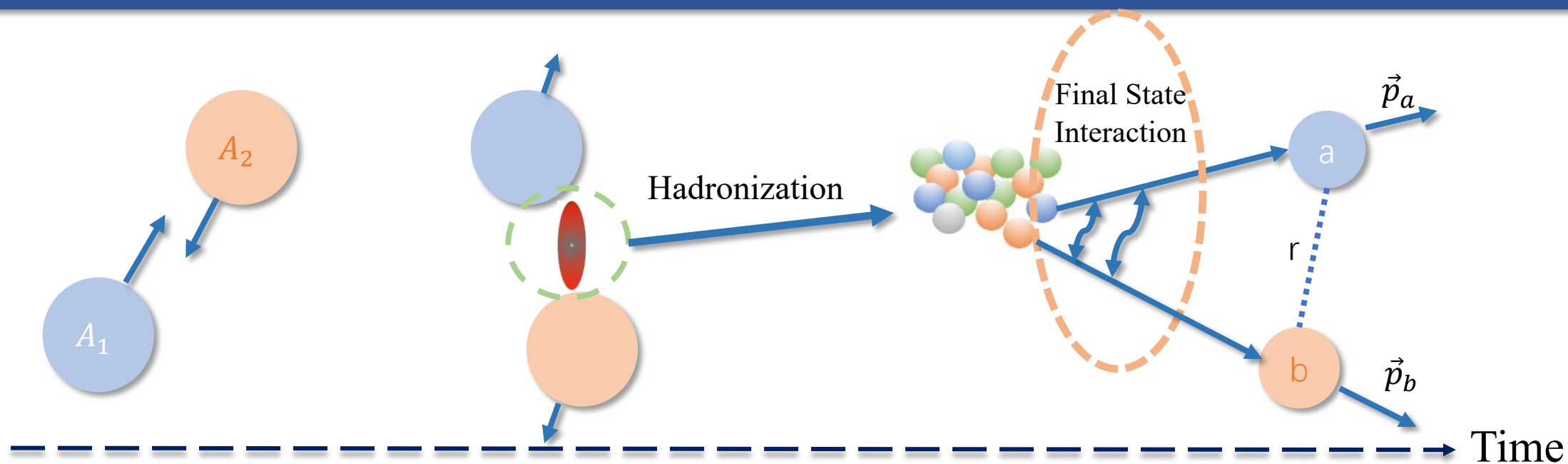
$$(|S|=2) \text{ Dibaryon} \Leftrightarrow \Lambda + \Lambda$$

$$(|S|=3) \text{ Dibaryon} \Leftrightarrow p + \Omega^-$$

➤ Hyperon-Nucleon (Y-N) and Hyperon-Hyperon (Y-Y) interactions provide important information to constrain the Equation-of-State and help to understand the inner structure of compact stars

More experimental measurements are needed !

Femtoscscopy



➤ Two-particle correlation function:

$$C(k^*) = \frac{\text{Model}}{\text{Experimental}} = \frac{\int S(\vec{r}) |\Psi(\vec{k}^*, \vec{r})|^2 d^3\vec{r}}{\frac{N_{\text{same}}(k^*)}{N_{\text{mixed}}(k^*)}}$$

$S(\vec{r})$: Source function

$\Psi(\vec{k}^*, \vec{r})$: Pair wave function

$k^* = \frac{1}{2} |\vec{p}_a - \vec{p}_b|$, relative momentum

$\vec{r}$: relative distance

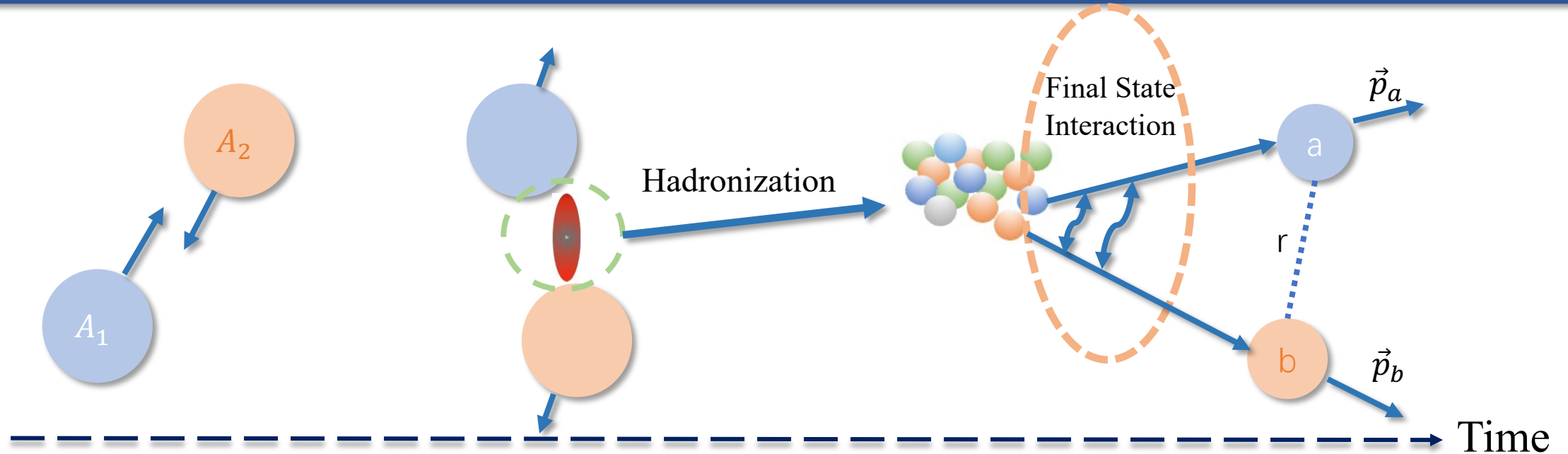
Depends on ...

- Emission source
- Strong interaction
- Coulomb interaction

R. Lednicky, et al. Sov.J.Nucl.Phys.35(1982)770

L. Michael, et al. Ann.Rev.Nucl.Part.Sci. 55 (2005) 357-402

Femtoscscopy



➤ Two-particle correlation function:

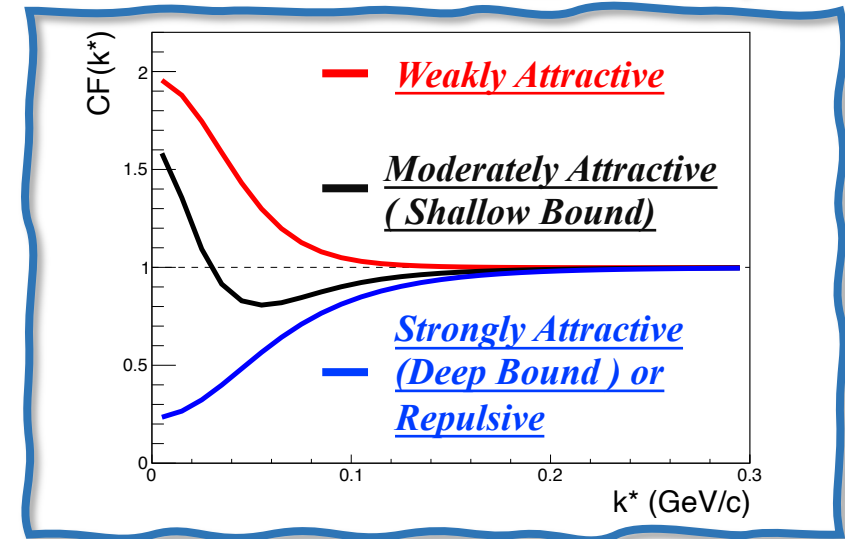
$$C(k^*) = \frac{\text{Model}}{\text{Experimental}} = \frac{\int S(\vec{r}) |\Psi(\vec{k}^*, \vec{r})|^2 d^3\vec{r}}{N_{\text{same}}(k^*) / N_{\text{mixed}}(k^*)}$$

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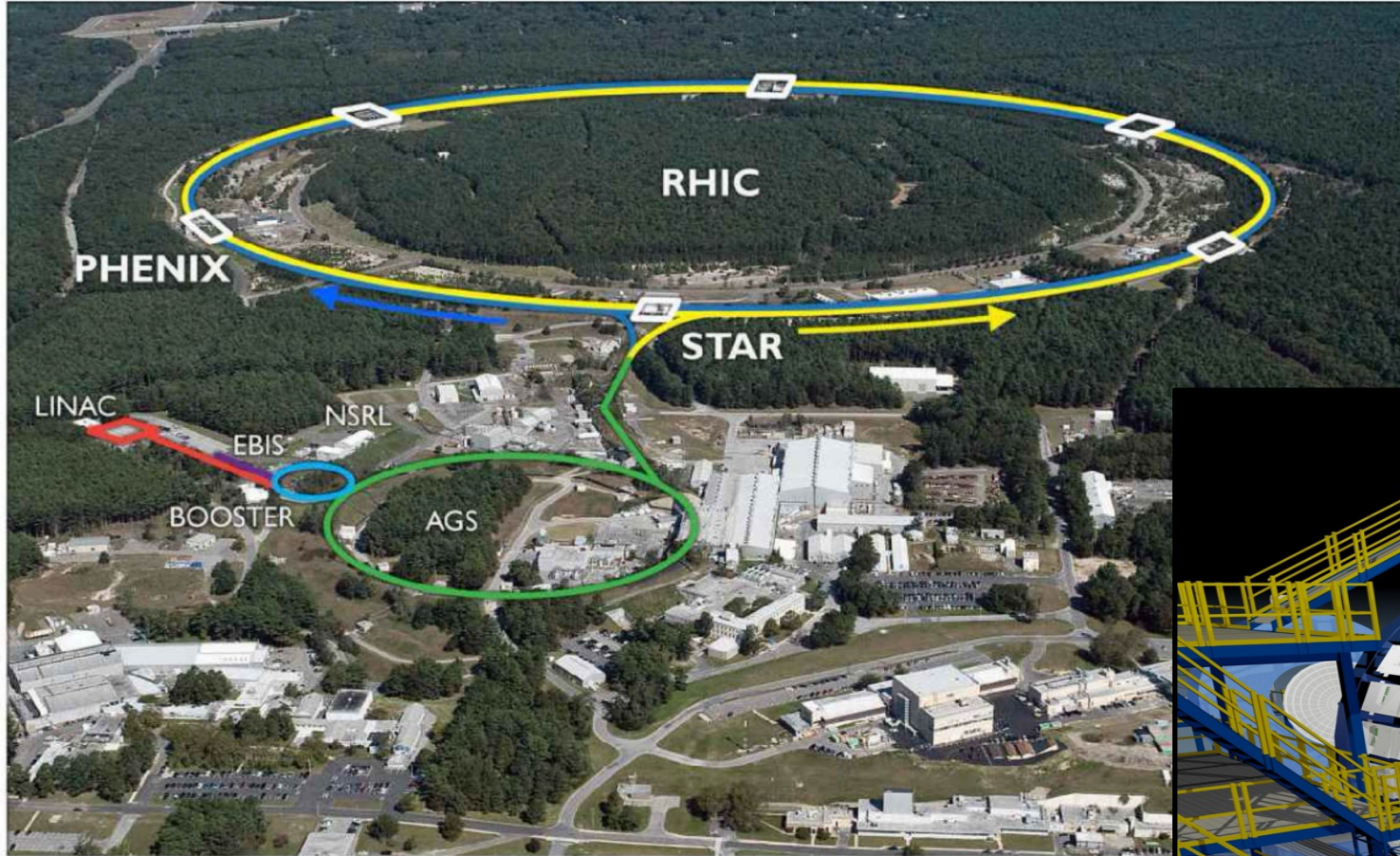


R. Lednicky, et al. Sov.J.Nucl.Phys.35(1982)770

L. Michael, et al. Ann.Rev.Nucl.Part.Sci. 55 (2005) 357-402

Zhi-Wei Liu, et al., Phys. Rev. D 107, 074019 (2023)

RHIC-STAR Experiment



Relativistic Heavy Ion Collider (RHIC)

Brookhaven National Laboratory, Upton

→ p+p, Au+Au, Zr+Zr, Ru+Ru, d+Au...

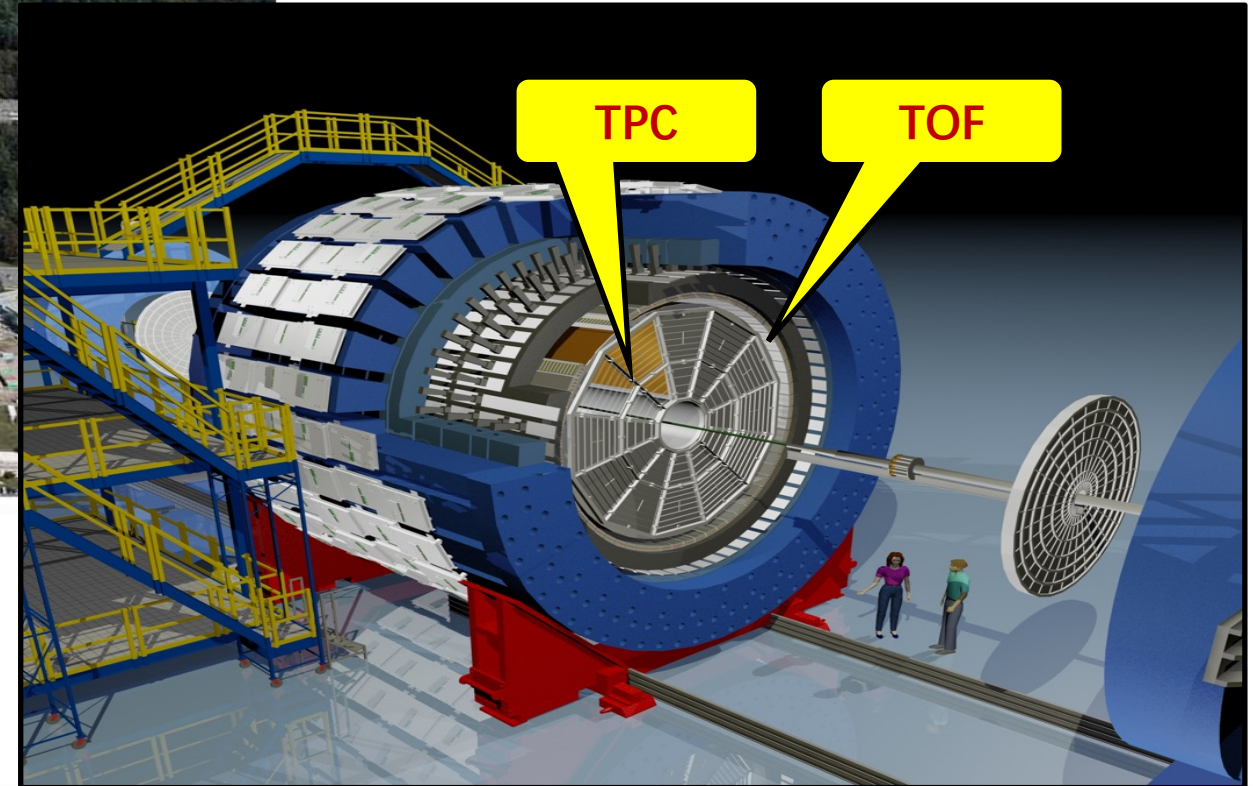
→ Beam Energy Scan Program I, II (Au+Au)

$$\sqrt{s_{NN}} = 3.0 - 200 \text{ GeV}$$

The Solenoidal Tracker At RHIC (STAR)

→ Excellent particle identification

→ Large, uniform acceptance at Mid-rapidity



Lednický-Lyuboshitz Model

$$CF(k^*) = \int d^3r^* S(r^*) |\psi(r^*, k^*)|^2 = \frac{N_{same}(k^*)}{N_{mixed}(k^*)}$$

➤ Formalism with Lednický-Lyuboshitz (L-L) model:

- ❑ Only consider s-wave
- ❑ Effective range expansion for $\psi(r^*, k^*)$
- ❑ Static and spherical Gaussian source assumed

R. Lednický and V. L. Lyuboshitz, Sov. J. Nucl. Phys. 35, 770 (1982)

➤ Wave function:

- ❑ $\psi(r^*, k^*) = e^{-ik^*r^*} + f(k^*) \frac{e^{ik^*r^*}}{r^*}$

➤ Scattering amplitude:

- ❑ $f(k^*) = [\frac{1}{f_0} + \frac{1}{2}d_0k^{*2} - ik^*]^{-1}$ (No Coulomb)
 - ❑ $f(k^*) = [\frac{1}{f_0} + \frac{1}{2}d_0k^{*2} - \frac{2}{a_c}h(\eta) - ik^*A_c(\eta)]^{-1}$ (Include Coulomb)
- a_c : Bohr radius $\eta = (k^*a_c)^{-1}$
 A_c, h : Coulomb interaction factor

L-L Fitting by
Bayesian method

➤ Physics quantity:

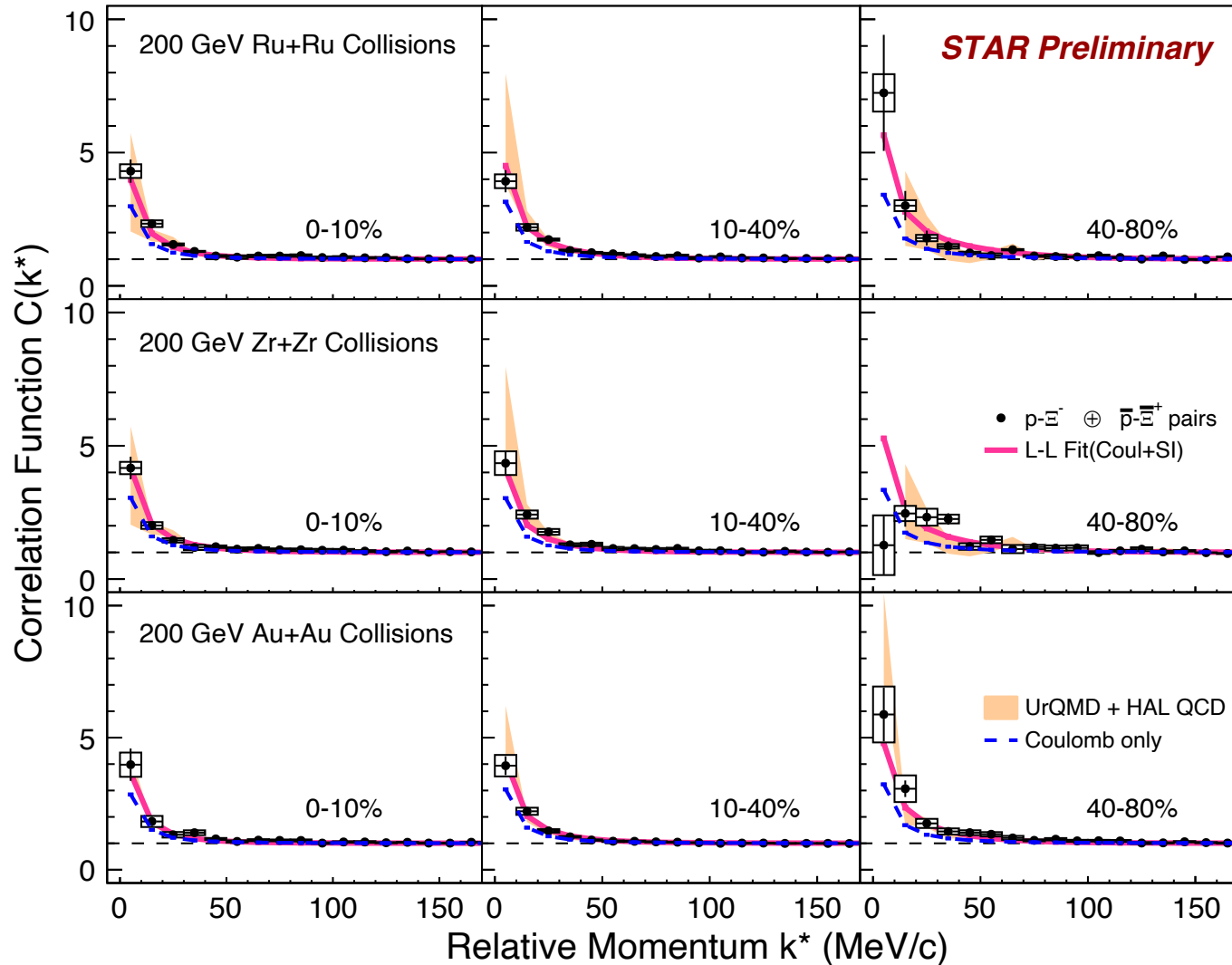
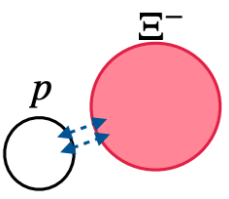
- ❑ f_0 : Scattering Length
- ❑ d_0 : Effective Range
- ❑ R_G : Spherical Gaussian Source Size

$f_0 > 0$: Attractive Interaction

$f_0 < 0$: Repulsive Interaction or Bound State

Bayesian method: https://github.com/chunshen1987/bayesian_analysis
Mäntysaari, H., et al., Phys. Lett. B 833 (2022) 137348

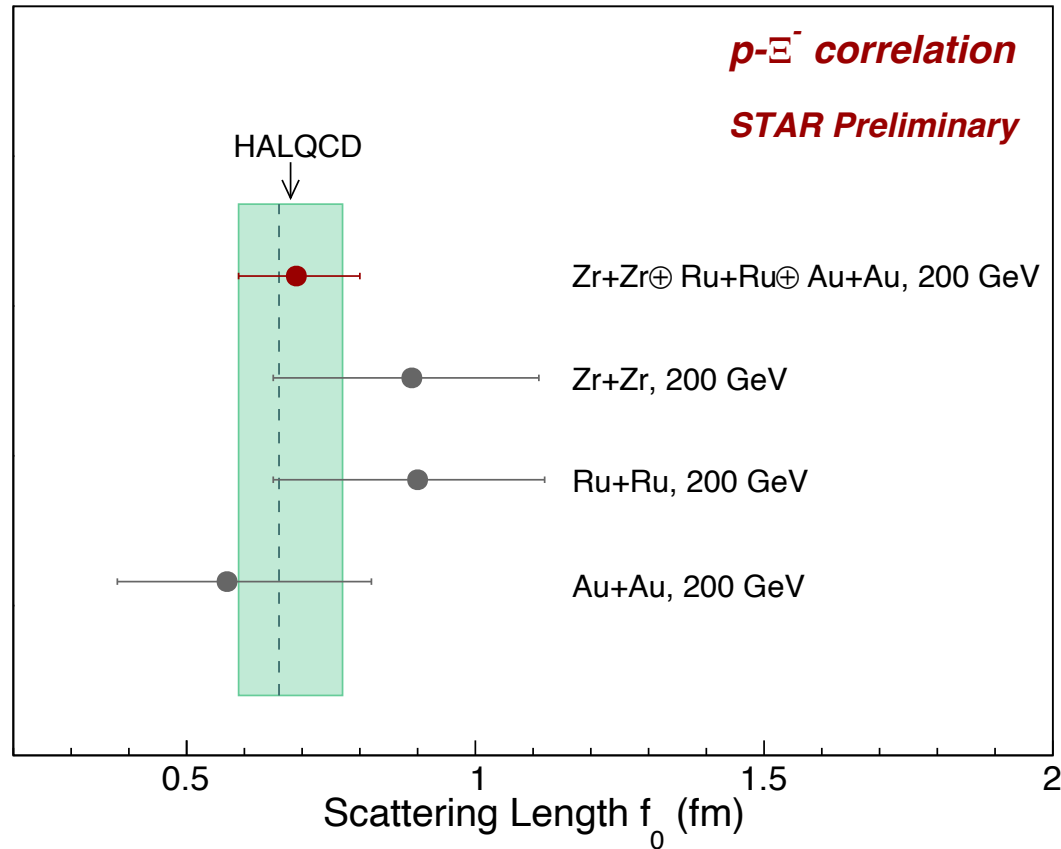
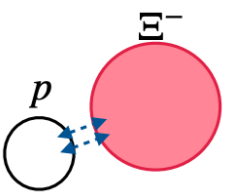
p- Ξ^- Correlation Function ($|S|=2$)



- I. Measure $p\text{-}\Xi^- \oplus \bar{p}\text{-}\bar{\Xi}^+$ CFs at $\sqrt{s_{NN}} = 200$ GeV in Au+Au and Isobar collisions
- II. CFs show enhancement at low k^*
- III. Simultaneously fit with L-L function for different centralities in each collision system to extract R_G , f_0 and d_0 by Bayesian method
- IV. UrQMD + HAL QCD model is consistent with data
 - ⊙ Particle phase space provided by UrQMD
 - ⊙ Interaction potential provided by HAL QCD

Y. Kamiya, et al., Phys. Rev. C 105, 014915 (2022)

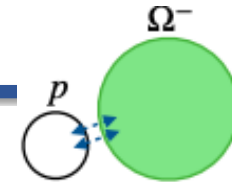
p- Ξ^- Interaction Parameters ($|S|=2$)



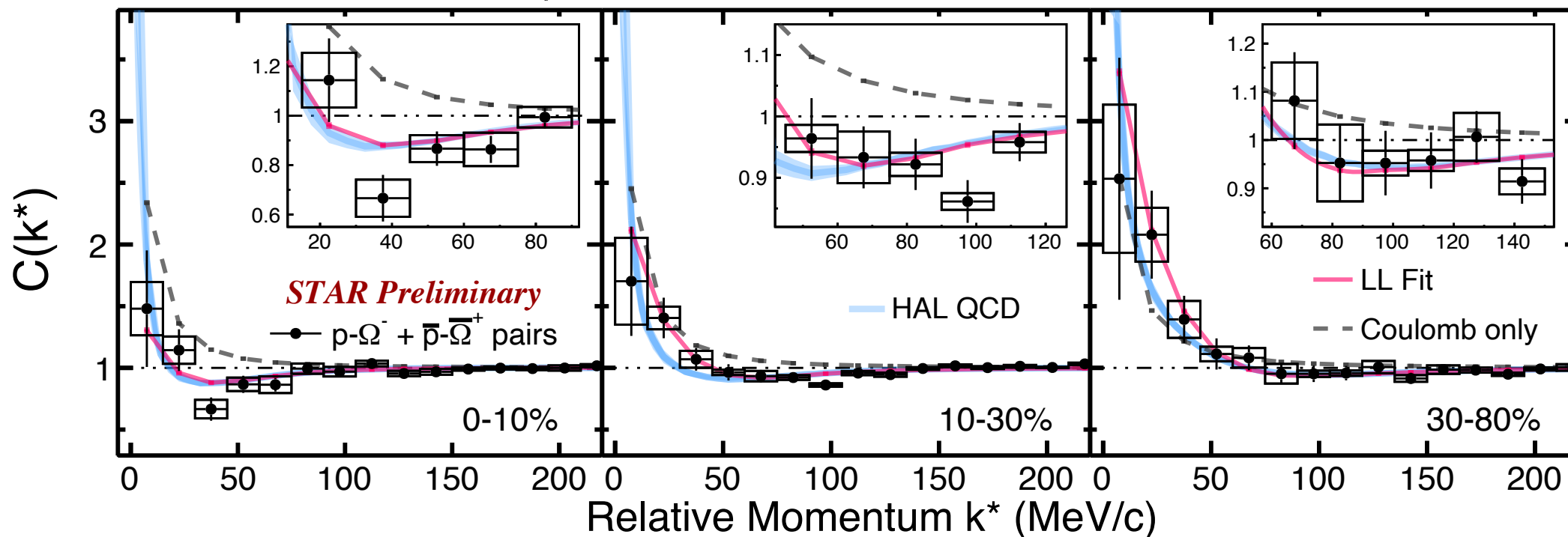
- I. Simultaneously fit with L-L function for different centralities in each collision system to extract R_G , f_0 and d_0 by Bayesian method
- II. First experimental constraints of strong interaction parameters in p- Ξ^- pairs in heavy-ion collisions
- III. Extracted spin averaged scattering length:
 $f_0 = 0.7^{+0.1}_{-0.1}$ fm (stat.+sys.)
 - Weakly attractive interaction
 - Consistent with HAL QCD prediction

Y. Kamiya, et al., Phys. Rev. C 105, 014915 (2022)

p- Ω^- Correlation Function ($|S|=3$)



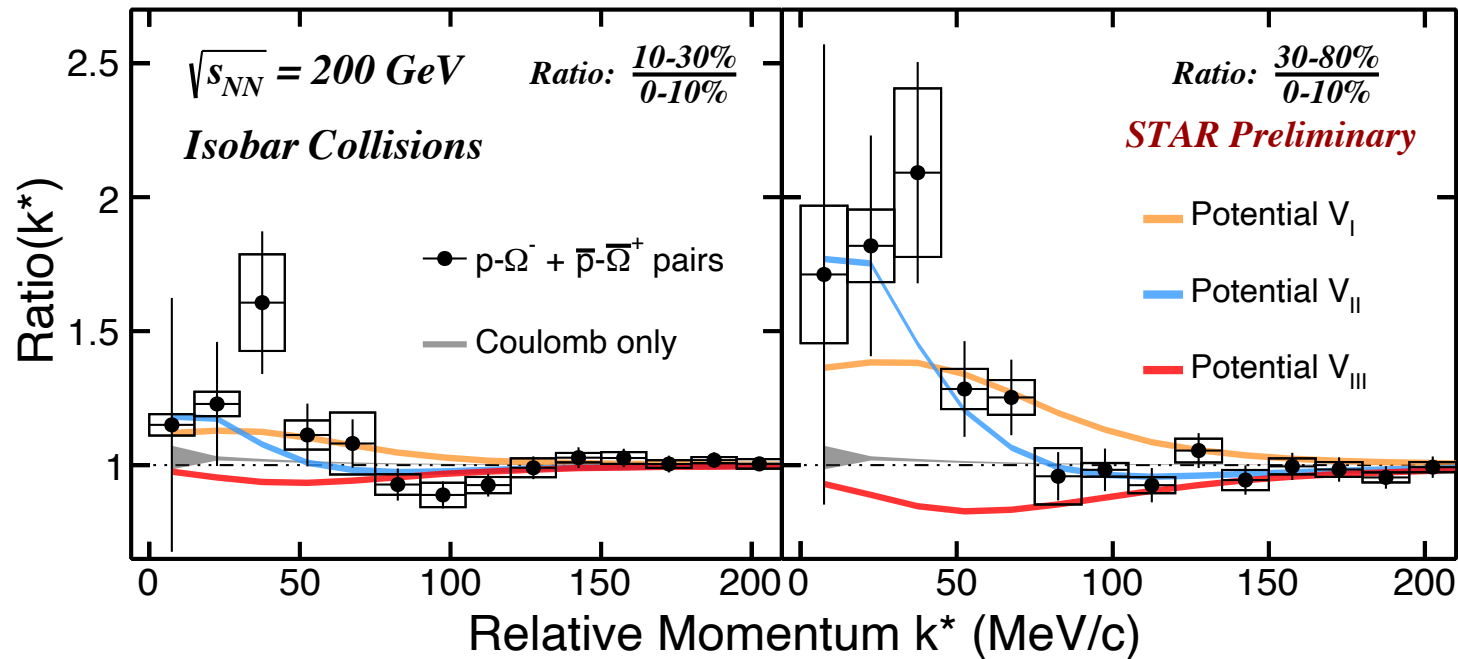
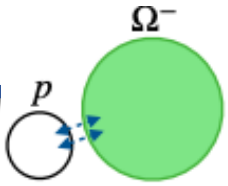
$\sqrt{s_{NN}} = 200 \text{ GeV}$ Isobar Collisions



- I. Precise measurements of p- $\Omega^- \oplus \bar{p}$ - $\bar{\Omega}^+$ correlation functions in Isobar collisions
 - ▣ CFs show enhancement at low k^* $\rightarrow$ mainly due to Coulomb attraction interaction
 - ▣ CFs show depletion at $k^* \sim 30$ -100 MeV/c $\rightarrow$ mainly due to the strong interaction
- II. Simultaneously fit with L-L function for 3 centralities to extract R_G , f_0 and d_0 by Bayesian method
- III. CFs obtained by HAL QCD theory with extracted R_G by L-L model is consistent with the data

Takumi Iritani, et al. (HAL QCD), Phys. Lett. B792 (2019)

p- Ω^- CF Ratio ($|S|=3$)



I. By taking CF ratio, Coulomb effect can be largely canceled

II. CF Ratio shows enhancement at low k^* and depletion around $k^* \sim 100$ MeV/c

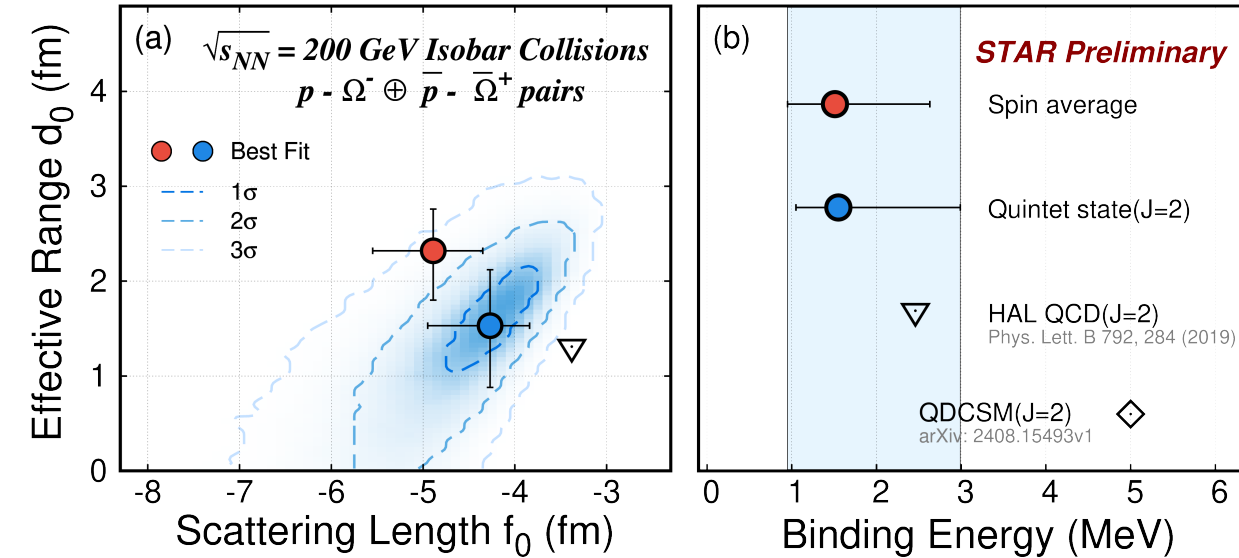
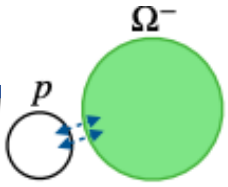
□ Due to the presence of shallow bound state

III. The potential V_{II} , with a p-value of 0.812, provides a better description of the data

Type	f_0 (fm)	d_0 (fm)	BE (MeV)	χ^2/ndf	p-value	
V_I [1]	1.12	1.16	--	48.2/29	0.014	No Bound
V_{II} [2]	-3.38	1.31	2.15	22.2/29	0.812	Shallow Bound
V_{III} [1]	-1.29	0.65	26.9	58.7/29	0.001	Deeply Bound

[1] Kenji Morita, et al., Phys. Rev. C 94, 031901 (2016)
 [2] Kenji Morita, et al., Phys. Rev. C 101, 015201 (2020)

p- Ω^- Interaction Parameters ($|S|=3$)



	Spin ave.	Quintet	HAL QCD
f_0 (fm)	$-4.9^{+0.5}_{-0.7}$	$-4.3^{+0.4}_{-0.7}$	-3.4
d_0 (fm)	$2.3^{+0.4}_{-0.5}$	$1.5^{+0.5}_{-0.7}$	1.3
BE (MeV)	$1.5^{+1.1}_{-0.6}$	$1.6^{+1.4}_{-0.5}$	2.3

Kenji Morita, et al., Phys. Rev. C 101, 015201 (2020)

- I. First experimental constraints in heavy-ion collisions of strong interaction parameters in p- Ω^- pair
- II. Extracted **negative f_0** ($|f_0| > 2d_0$) by Spin average method and Quintet method
 - First experimental evidence of Strange Dibaryon
- III. Calculate Binding Energy (BE) via Betha formula:

$$\text{Reduced mass: } m_{p\Omega} = \frac{m_p m_\Omega}{m_p + m_\Omega}$$

$$BE_{p\Omega} = \frac{1}{2m_{p\Omega}d_0^2} \left(1 - \sqrt{1 + \frac{2d_0}{f_0}} \right)^2$$

- Calculated BE are consistent with HAL QCD prediction

H-dibaryon Search($S = -2, J = 0$)

• Search for the H-dibaryon

The H-dibaryon is an $S=-2, I=0, J=0$ state with quark content $uuddss$, first predicted by Robert L. Jaffe in 1977

R. L. Jaffe, Phys. Rev. Lett. 38, 617 (1977)

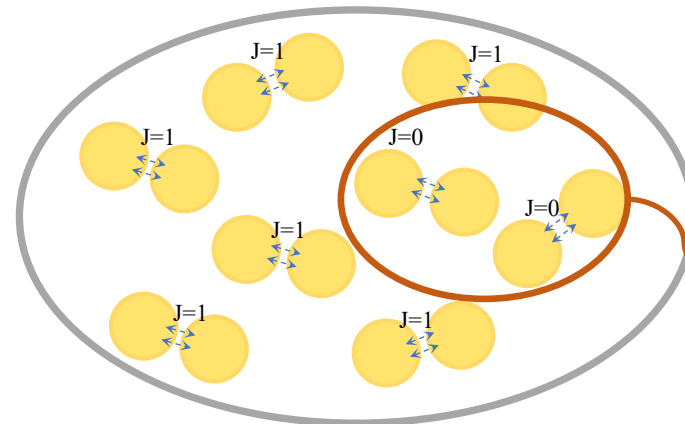
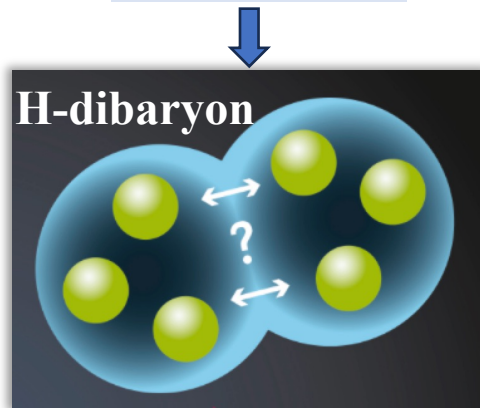
$$\text{H-dibaryon} \Leftrightarrow p + \Xi^-$$

$$\text{H-dibaryon} \Leftrightarrow \Lambda + \Lambda$$

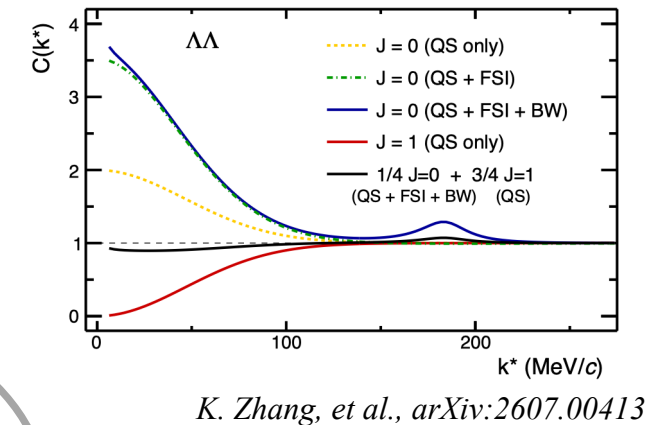
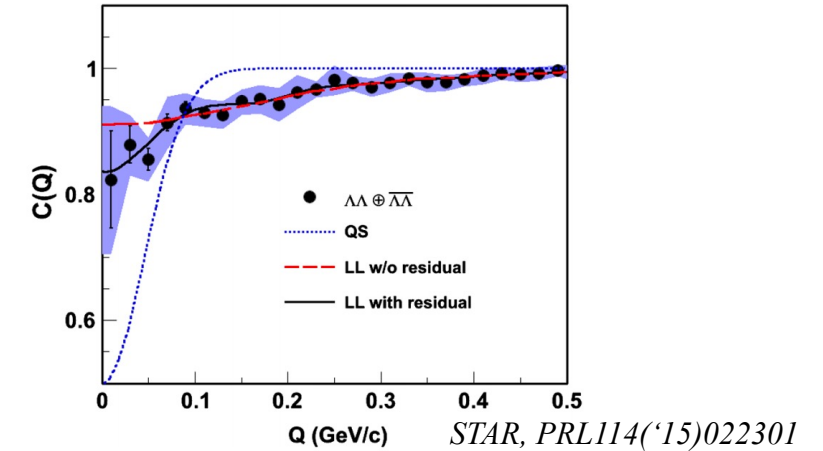
In $\Lambda\Lambda$ and $p-\Xi^-$ pairs, there are two spin states:

$$C_{p-\Xi} = \frac{1}{4} C_{J=0, \text{singlet}} + \frac{3}{4} C_{J=1, \text{triplet}}$$

$$C_{\Lambda-\Lambda} = \frac{1}{4} C_{J=0, \text{singlet}} + \frac{3}{4} C_{J=1, \text{triplet}}$$



If the spin state ($J=0$) can be selected, a signal of the H-dibaryon may become observable.



Spin Femtoscopy Method

- $\Lambda\Lambda$ spin correlation:**

$$\frac{dN}{d\cos\theta^*} \propto 1 + \alpha_\Lambda^2 P_{\Lambda\Lambda} \cos\theta^* \quad P_{\Lambda\Lambda} = \frac{1}{3} \langle \sigma_1 \cdot \sigma_2 \rangle$$

$$\alpha_\Lambda = 0.746$$

- Different spin states:**

$$J=0, P_{\Lambda\Lambda} = -1 \quad \frac{dN}{d\cos\theta^*} \propto 1 - \alpha_\Lambda^2 \cos\theta^*$$

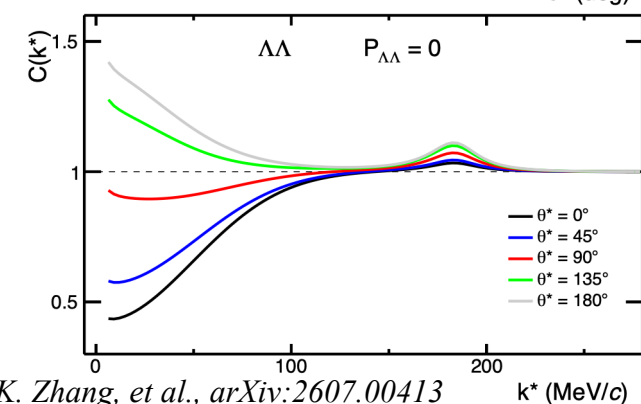
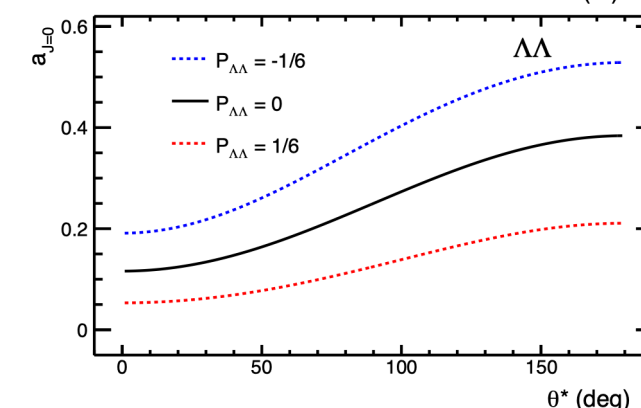
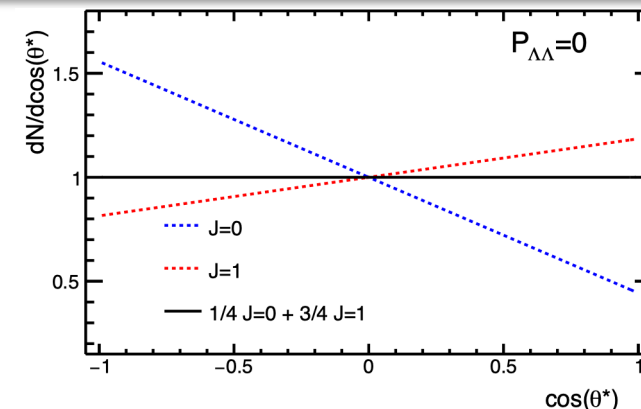
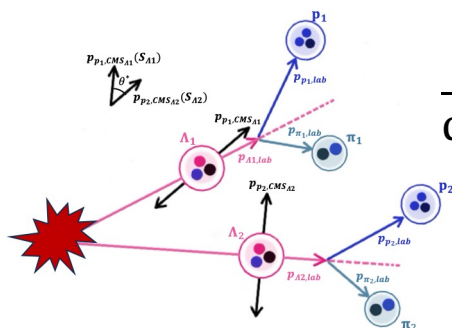
$$J=1, P_{\Lambda\Lambda} = 1/3 \quad \frac{dN}{d\cos\theta^*} \propto 1 + \frac{1}{3} \alpha_\Lambda^2 \cos\theta^*$$

$$\begin{aligned} \frac{dN}{d\cos\theta^*} &\propto w_0(1 - \alpha_\Lambda^2 \cos\theta^*) + (1 - w_0)(1 + \frac{1}{3} \alpha_\Lambda^2 \cos\theta^*) \\ &= 1 + \alpha_\Lambda^2 \left(\frac{1 - 4w_0}{3} \right) \cos\theta^* \end{aligned}$$

w_0 is the total spin $J=0$ state weight

- spin weight dependence on θ^* :**

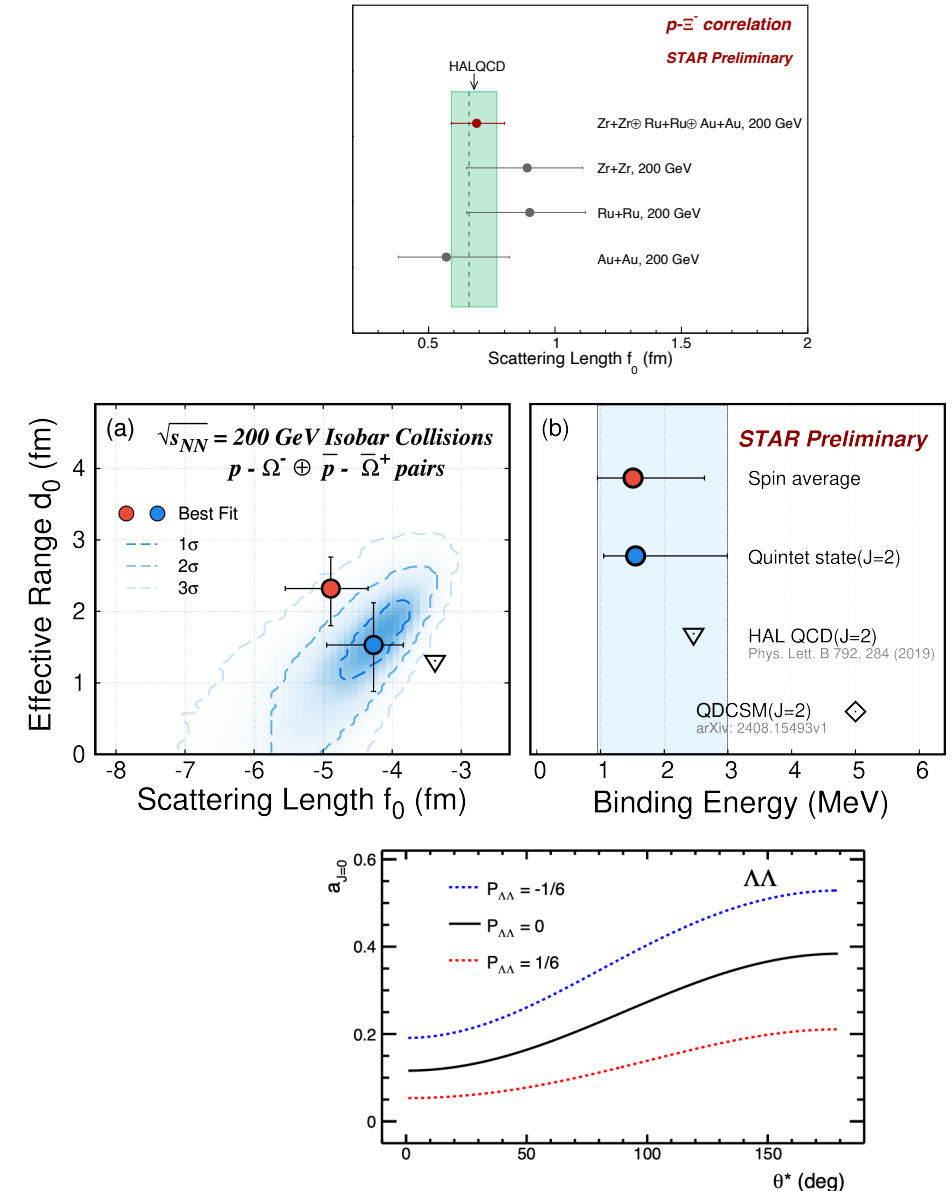
$$a_{J=0}(\theta^*) = \frac{w_0(1 - \alpha_\Lambda^2 \cos\theta^*)}{w_0(1 - \alpha_\Lambda^2 \cos\theta^*) + (1 - w_0)(1 + \frac{1}{3} \alpha_\Lambda^2 \cos\theta^*)}$$



K. Zhang, et al., arXiv:2607.00413

Summary

- I. Precision measurements of $p\text{-}\Xi^- \oplus \bar{p}\text{-}\bar{\Xi}^+$ and $p\text{-}\Omega^- \oplus \bar{p}\text{-}\bar{\Omega}^+$ correlation function at STAR
- II. Extracted strong interaction parameters using L-L model by Bayesian method
 - $p\text{-}\Xi^- \oplus \bar{p}\text{-}\bar{\Xi}^+$ pairs: $f_0 > 0$
 - $p\text{-}\Omega^- \oplus \bar{p}\text{-}\bar{\Omega}^+$ pairs: $f_0 < 0 \rightarrow$ Bound state
- III. Extracted Binding Energy $BE = 1.6_{-0.5}^{+1.4}$ MeV in $p\text{-}\Omega^- \oplus \bar{p}\text{-}\bar{\Omega}^+$ pair consistent with HAL QCD prediction
- IV. Spin femtoscopy provides a novel approach to select spin states and enhance the sensitivity to H-dibaryon searches.





Thank you for your attention!