

Neutrinos in Cosmology: Some recent developments

Marco Drewes
Université catholique de Louvain

NOPP 2026
IHEP Beijing, China

Towards a precision calculation of N_{eff} in the Standard Model

Marco Drewes
Université catholique de Louvain

NOPP 2026
IHEP Beijing, China

N_{eff} in the Standard Model

- N_{eff} is the “effective number of neutrino species”; parameterises expansion rate during BBN and CMB decoupling

$$\left. \frac{\rho_\nu}{\rho_\gamma} \right|_{T/m_e \rightarrow 0} \equiv \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{\text{eff}}^{\text{SM}}$$

- In the SM it is defined as

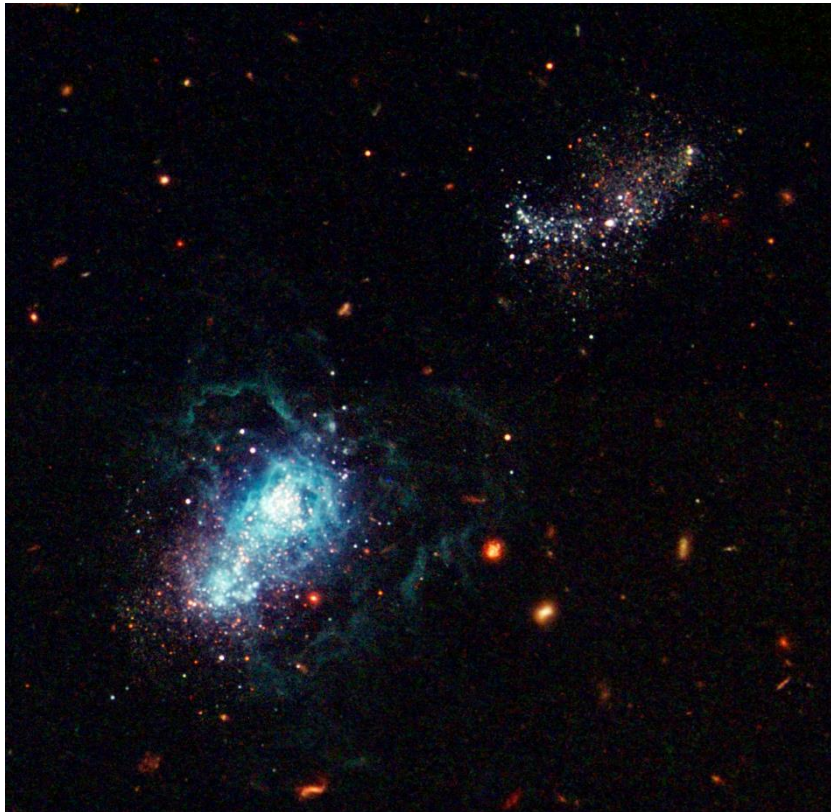
- $N_{\text{eff}} = 3$ in the SM if

- i) primordial plasma is ideal gas,
- ii) neutrinos decouple instantaneously,
- iii) neutrinos decoupled at temperature $T \gg m_e$

None of this is really true, leading to deviations from $N_{\text{eff}} = 3$

- BSM phenomena also lead to deviations from $N_{\text{eff}} = 3$ (extra light particles, modified expansion history, non-standard neutrino interactions...), making N_{eff} a powerful probe of BSM physics
- One can measure N_{eff} with percent accuracy

Observational Perspectives



Big Bang Nucleosynthesis

- Recently the primordial Helium abundance has been measured with unprecedented precision

Aver et al 2601.22238

- This has — combined with various CMB and LSS measurement — yielded the currently strongest bound on $N_{eff} = 2.990 \pm 0.070$

Goldstein/Hill 2603.13226

Cosmic Microwave Background

- Planck alone imposes $N_{eff} = 2.99 \pm 0.17$

Planck 1807.06209

- Simons Observatory is foreseen to achieve ± 0.045

Simons Observatory 2503.00636



N_{eff} in the Standard Model

- N_{eff} is the “effective number of neutrino species”; parameterises expansion rate during BBN and CMB decoupling

$$\left. \frac{\rho_\nu}{\rho_\gamma} \right|_{T/m_e \rightarrow 0} \equiv \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{\text{eff}}^{\text{SM}}$$

- In the SM it is defined as

- $N_{\text{eff}} = 3$ in the SM if
 - i) primordial plasma is ideal gas,
 - ii) neutrinos decouple instantaneously,
 - iii) neutrinos decoupled at temperature $T \gg m_e$

None of this is really true, leading to deviations from $N_{\text{eff}} = 3$

- BSM phenomena also lead to deviations from $N_{\text{eff}} = 3$ (extra light particles, modified expansion history, non-standard neutrino interactions...), making N_{eff} a powerful probe of BSM physics
- One can measure N_{eff} with percent accuracy

- **Status 2020 (current state-of-the-art value):**

$$N_{\text{eff}}^{\text{SM}} = 3.0440 \pm 0.0002$$

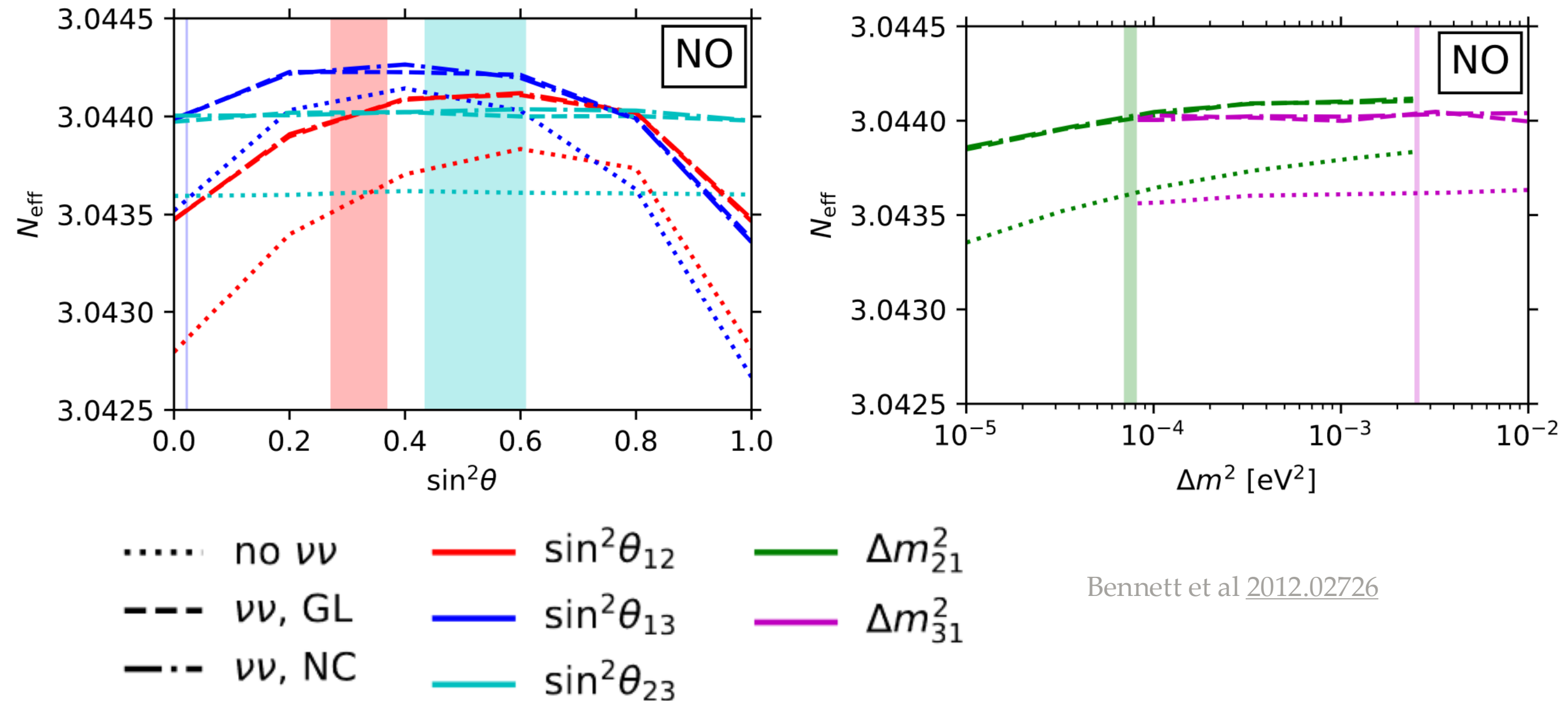
Akita/Yamaguchi al [2005.07047](#)

Froustey/Pitrou [2008.01074](#)

Bennett et al [2012.02726](#)

(used by PDG, in CAMB and CLASS codes, major collaborations like DES, DESI...)

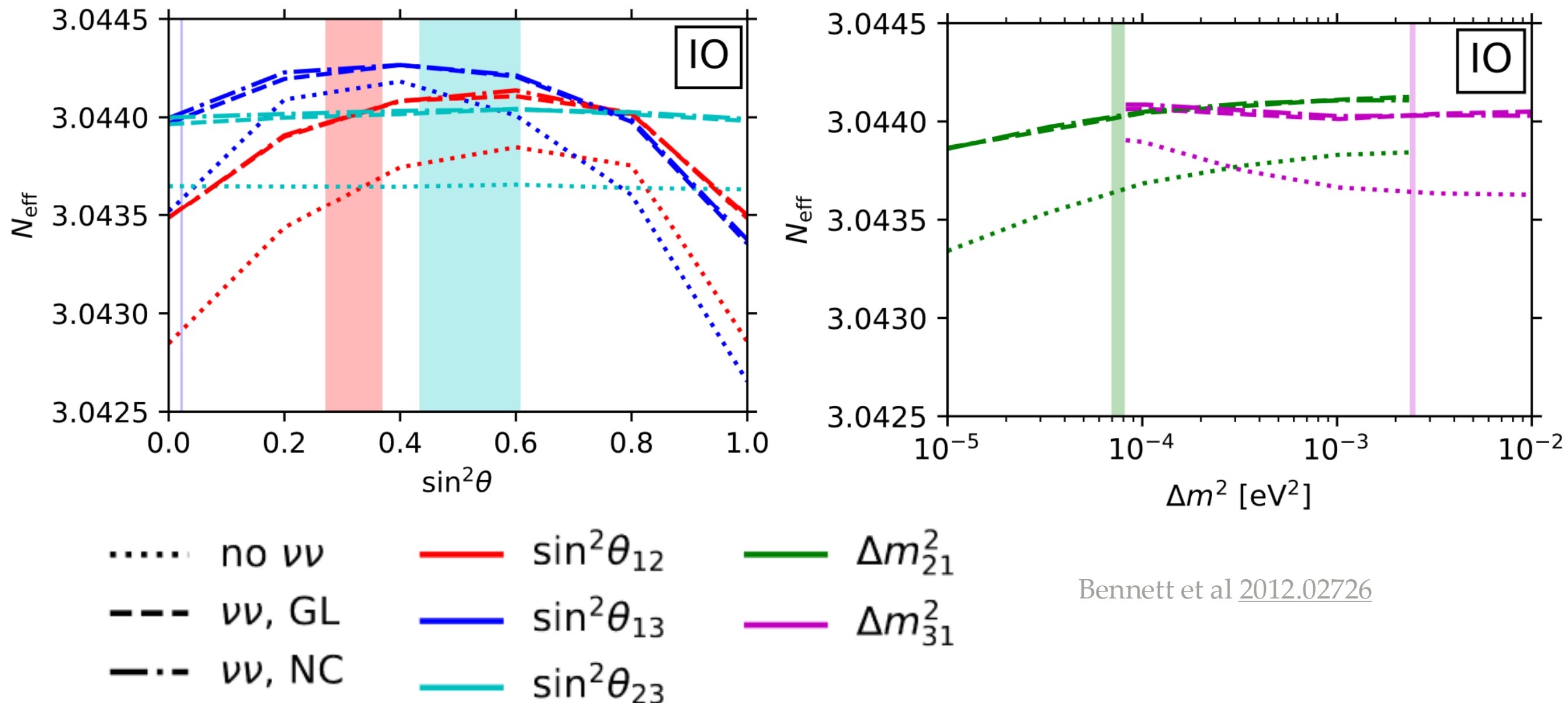
Experimental Uncertainties



Bennett et al [2012.02726](#)

- Largest uncertainty from experiment: mixing angles
- Uncertainty from neutrino oscillation data is sub-dominant compared to numerical error (Gauss-Laguerre vs Newton-Cotes quadrature)

Experimental Uncertainties



Bennett et al [2012.02726](#)

- Largest uncertainty from experiment: mixing angles
- Uncertainty from neutrino oscillation data is sub-dominant compared to numerical error (Gauss-Laguerre vs Newton-Cotes quadrature)

Impact of NLO QED Corrections

- N_{eff} in the SM is found by solving

$$(d/dt)\rho_{\text{tot}} + 3H(\rho_{\text{tot}} + P_{\text{tot}}) = 0, \quad \partial_t \varrho - pH \partial_p = -i[\mathbb{H}, \varrho] + \mathcal{I}[\varrho],$$

- $\mathcal{O}(e^3)$ correction to equation of state is more important than neutrino oscillations!

Standard-model corrections to $N_{\text{eff}}^{\text{SM}}$	Leading-digit contribution
m_e/T_d correction	+0.04
$\mathcal{O}(e^2)$ FTQED correction to the QED EoS	+0.01
Non-instantaneous decoupling+spectral distortion	-0.005
$\mathcal{O}(e^3)$ FTQED correction to the QED EoS	-0.001
Flavour oscillations	+0.0005
Type (a) FTQED corrections to the weak rates	$\lesssim 10^{-4}$

Impact of NLO QED Corrections

- N_{eff} in the SM is found by solving

$$(d/dt)\rho_{\text{tot}} + 3H(\rho_{\text{tot}} + P_{\text{tot}}) = 0, \quad \partial_t \varrho - pH \partial_p = -i[\mathbb{H}, \varrho] + \mathcal{I}[\varrho],$$

- $\mathcal{O}(e^3)$ correction to equation of state is more important than neutrino oscillations!

Standard-model corrections to $N_{\text{eff}}^{\text{SM}}$	Leading-digit contribution
m_e/T_d correction	+0.04
$\mathcal{O}(e^2)$ FTQED correction to the QED EoS	+0.01
Non-instantaneous decoupling+spectral distortion	-0.005
$\mathcal{O}(e^3)$ FTQED correction to the QED EoS	-0.001
Flavour oscillations	+0.0005
Type (a) FTQED corrections to the weak rates	$\lesssim 10^{-4}$

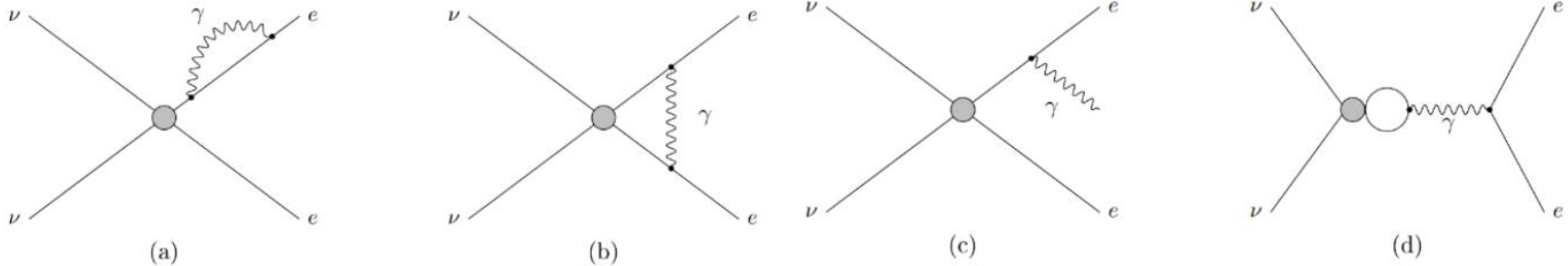
- The result $N_{\text{eff}} = 3.0440 \pm 0.0002$ obtained in Bennett et al 2012.02726 consistently included finite temperature QED corrections to the equation of state at NLO...

$$\ln Z^{(2)} = -\frac{1}{2} \text{ (diagram: circle with wavy line) } \quad \ln Z^{(3)} = \frac{1}{2} \left[\frac{1}{2} \text{ (diagram: circle with 2 wavy lines) } - \frac{1}{3} \text{ (diagram: circle with 3 wavy lines) } + \frac{1}{4} \text{ (diagram: circle with 4 wavy lines) } + \dots \right]$$

- ...but only roughly estimates NLO corrections to the collision term
- Cielo et al 2306.05460 found that NLO corrections to collision term lead to $N_{\text{eff}} = 3.043$

NLO Corrections to Collision Term

QED Corrections to Collision Term



Processes in Fig.1 included	Quantum statistics	Hadronic corrections	Non-linearities in QKE	Electron mass in NLO terms	Flavour dependence	First principles derivation
(a) roughly estimated	partially (fermions)	no	yes	no	partially	no
(a)-(c) partially	partially	no	no	partially	no	no
all	yes	partially	no	no	yes	yes
(d)	yes	no	no	yes	yes	yes

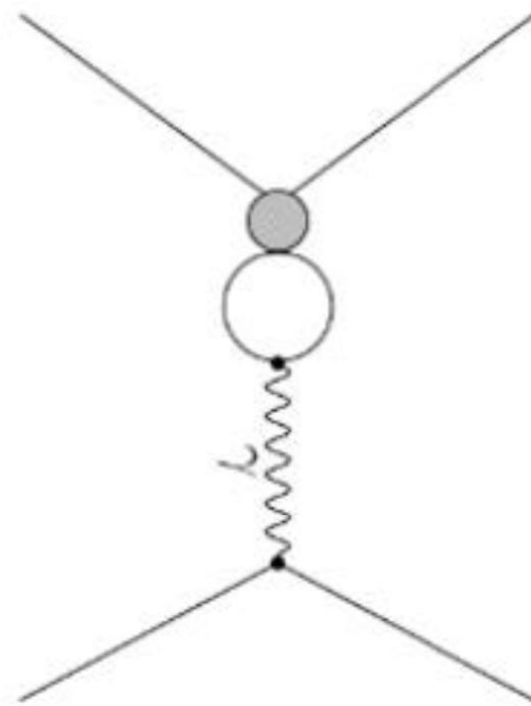
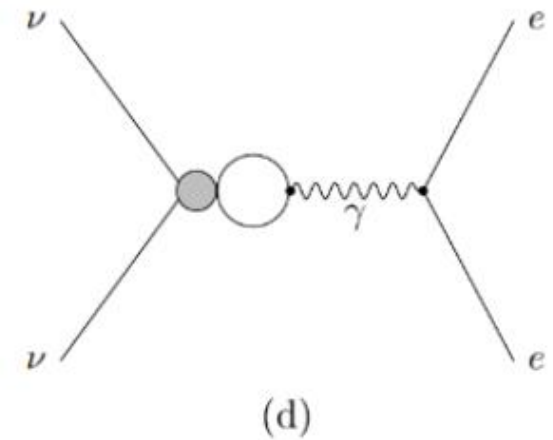
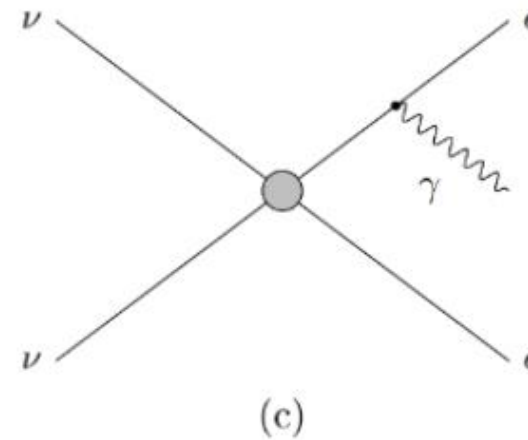
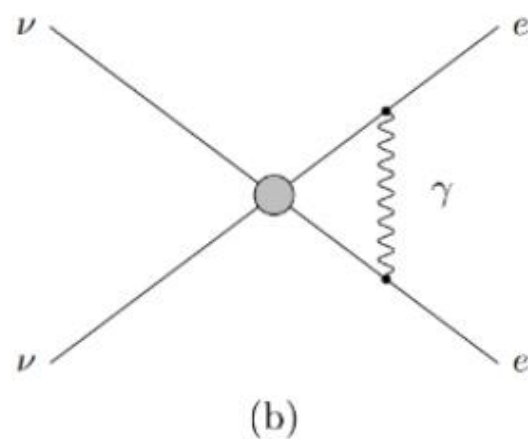
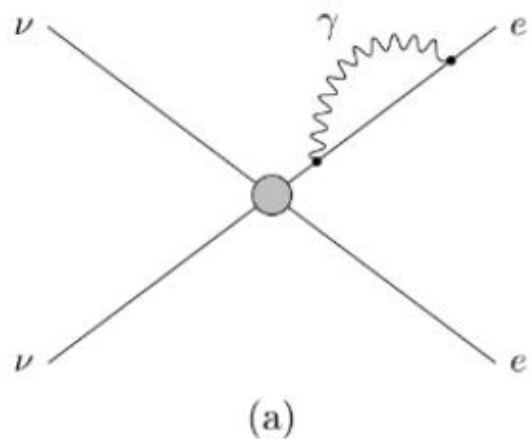
Bennett et al
[2012.02726](#)

Cielo et al
[2306.05460](#)

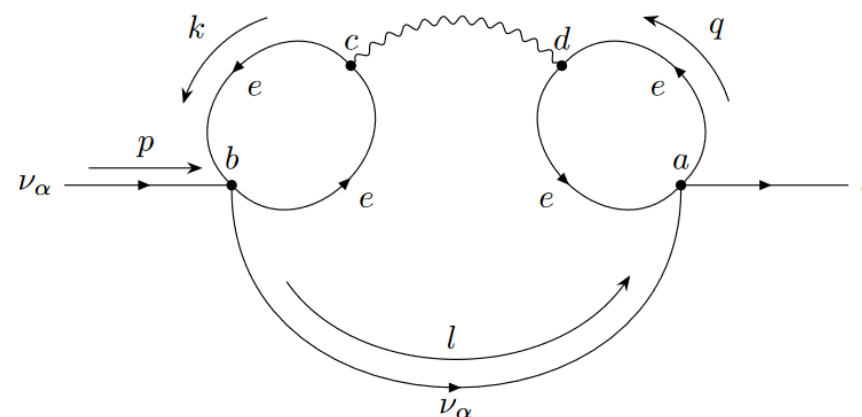
Jackson/Laine
[2312.07015](#), [2412.03958](#)

MaD et al [2402.18481](#)

QED Corrections to Collision Term



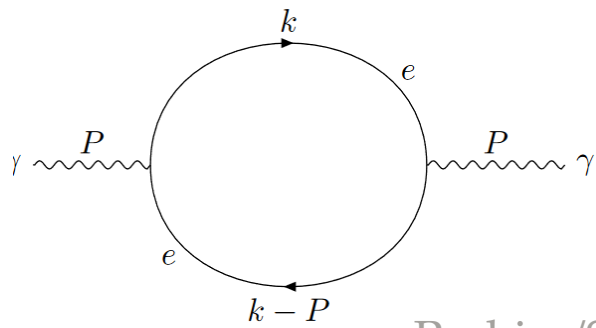
- t-channel version of diagram (d) is IR divergent in the t-channel,
- Hence we focussed on (d) in MaD et al 2402.18481
- In Closed-Time-Path formalism appears as a cut of



- requires usage of resummed finite temperature photon propagator

- But: annihilation channels transfer more energy than scatterings, see Jackson/Laine 2412.03958

Resummed Photon Propagator



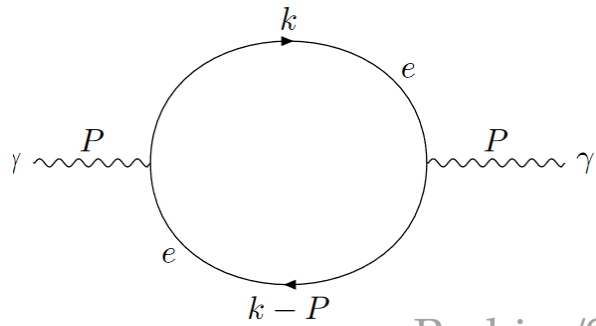
$$\Pi_{ab}^{\mu\nu}(P) = (-1)^{a+b} i e^2 \int \frac{d^4 k}{(2\pi)^4} \text{Tr} \left[i S_e^{ab}(k) \gamma^\mu i S_e^{ba}(k - P) \gamma^\nu \right]$$

Peshier/Schertler/Thoma et [9708434](#)

Thoma/Carrington [9708363](#)

MaD et al [2402.18481](#)

Resummed Photon Propagator



$$\Pi_{ab}^{\mu\nu}(P) = (-1)^{a+b} i e^2 \int \frac{d^4 k}{(2\pi)^4} \text{Tr} \left[i S_e^{ab}(k) \gamma^\mu i S_e^{ba}(k-P) \gamma^\nu \right]$$

Peshier/Schertler/Thoma et [9708434](#)

Thoma/Carrington [9708363](#)

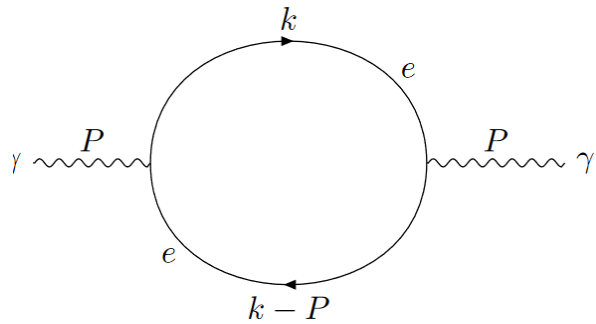
MaD et al [2402.18481](#)

$$\begin{aligned} \text{Re } \Pi_{11,T \neq 0}^T &= \frac{\alpha_{\text{em}}}{\pi |\mathbf{P}|^3} \int_{m_e}^{\infty} d\omega \left[2\ell_2(\omega, P) P_0 P^2 \omega - 4|\mathbf{k}||\mathbf{P}|(P_0^2 + |\mathbf{P}|^2) \right. \\ &\quad \left. + \ell_1(\omega, P)(|\mathbf{P}|^2 P^2 + 2P_0^2 \omega^2 - 2|\mathbf{k}|^2 |\mathbf{P}|^2 + \frac{1}{2} P^4) \right] f_F(\omega) \\ &\stackrel{\text{HTL}}{\approx} -\frac{3}{2} m_\gamma^2 \frac{P_0^2}{|\mathbf{P}|^2} \left[1 - \left(1 - \frac{|\mathbf{P}|^2}{P_0^2} \right) \frac{P_0}{2|\mathbf{P}|} \ln \left| \frac{P_0 + |\mathbf{P}|}{P_0 - |\mathbf{P}|} \right| \right] \longrightarrow 0, \quad \text{for } P_0 = 0 \end{aligned}$$

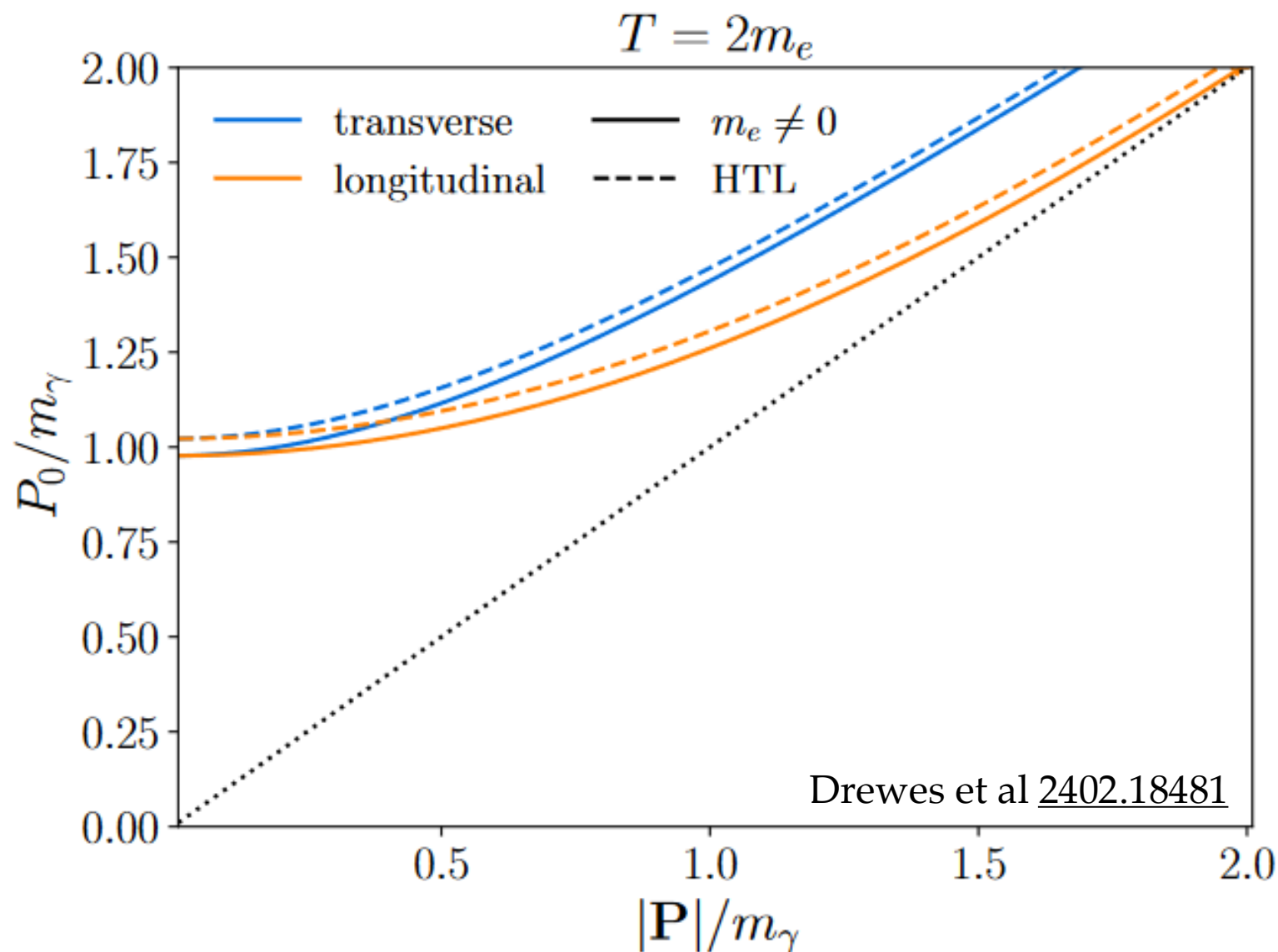
$$\begin{aligned} \text{Re } \Pi_{11,T \neq 0}^L &= \frac{\alpha_{\text{em}} P^2}{\pi |\mathbf{P}|^3} \int_{m_e}^{\infty} d\omega \left[8|\mathbf{k}||\mathbf{P}| - \ell_1(\omega, P)(P^2 + 4\omega^2) - 4\ell_2(\omega, P) P_0 \omega \right] f_F(\omega) \\ &\stackrel{\text{HTL}}{\approx} -3m_\gamma^2 \left(1 - \frac{P_0^2}{|\mathbf{P}|^2} \right) \left[1 - \frac{P_0}{2|\mathbf{P}|} \ln \left| \frac{P_0 + |\mathbf{P}|}{P_0 - |\mathbf{P}|} \right| \right] \longrightarrow -3m_\gamma^2 \quad \text{for } P_0 = 0, \end{aligned}$$

Resummation: momentum-dependent photon mass and longitudinal photon mode

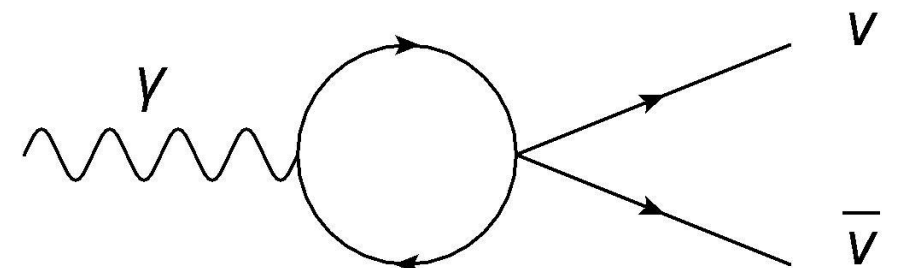
Neutrino Decay



$$\Pi_{ab}^{\mu\nu}(P) = (-1)^{a+b} i e^2 \int \frac{d^4 k}{(2\pi)^4} \text{Tr} \left[i S_e^{ab}(k) \gamma^\mu i S_e^{ba}(k-P) \gamma^\nu \right]$$



Gives rise to well-known plasmon process that is relevant for White Dwarf cooling



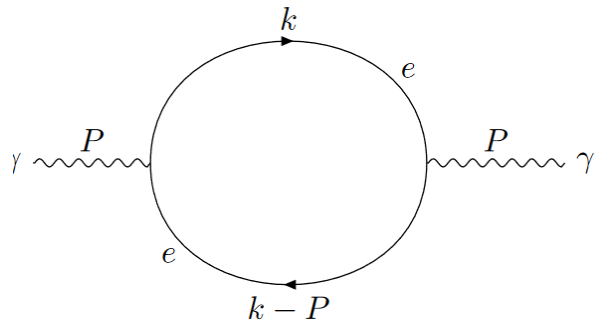
Braaten/Segel [9302213](#)

Haft/Raffelt/Weiss [9309014](#)

MaD et al [2109.06158](#)

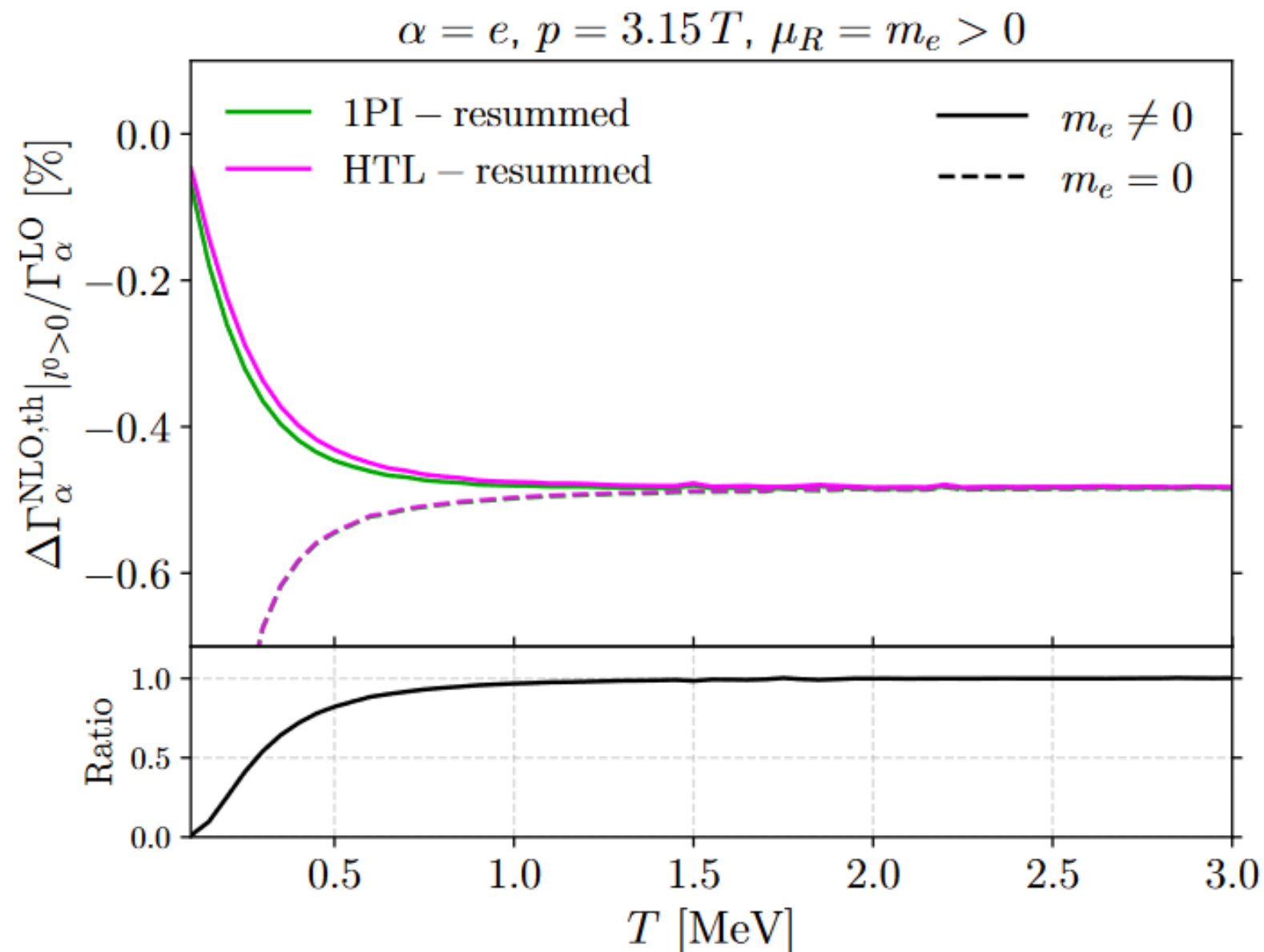
Impact on N_{eff} is negligible

Correction to Neutrino Scattering



$$\Pi_{ab}^{\mu\nu}(P) = (-1)^{a+b} i e^2 \int \frac{d^4 k}{(2\pi)^4} \text{Tr} \left[i S_e^{ab}(k) \gamma^\mu i S_e^{ba}(k-P) \gamma^\nu \right]$$

MaD et al [2402.18481](#)



Impact on N_{eff}

- In principle: need to solve density matrix equation for all neutrino modes
- Large set of coupled equations, highly non-linear due to neutrino self-interactions (the Γ are functionals of q)
- We use two simpler methods to estimate impact on N_{eff} :
 - entropy conservation
 - solve density matrix equation in “damping approximation”

Method	$\delta N_{\text{eff}}/N_{\text{eff}}^{\text{LO}}$	$\delta N_{\text{eff}}^{m_e=0}/N_{\text{eff}}^{\text{LO}}$
Entropy conservation (common decoupling)	-1.1×10^{-5}	-1.2×10^{-5}
Entropy conservation (flavour-dependent decoupling)	-5.4×10^{-6}	-6.1×10^{-6}
Boltzmann damping approximation (mean momentum)	-7.8×10^{-6}	-1.6×10^{-5}
Boltzmann damping approximation (full momentum)	-7.9×10^{-6}	-2.6×10^{-5}

- Ultimately: full density matrix needs to be solved with all contributions
- Our results indicate that **NLO corrections to collision integral do not modify N_{eff} considerably** , indicating that $N_{\text{eff}} = 3.044$ in the SM

MaD et al [2402.18481](#)

- Independent computations by Jackson and Laine including all annihilation channels come to the same conclusion

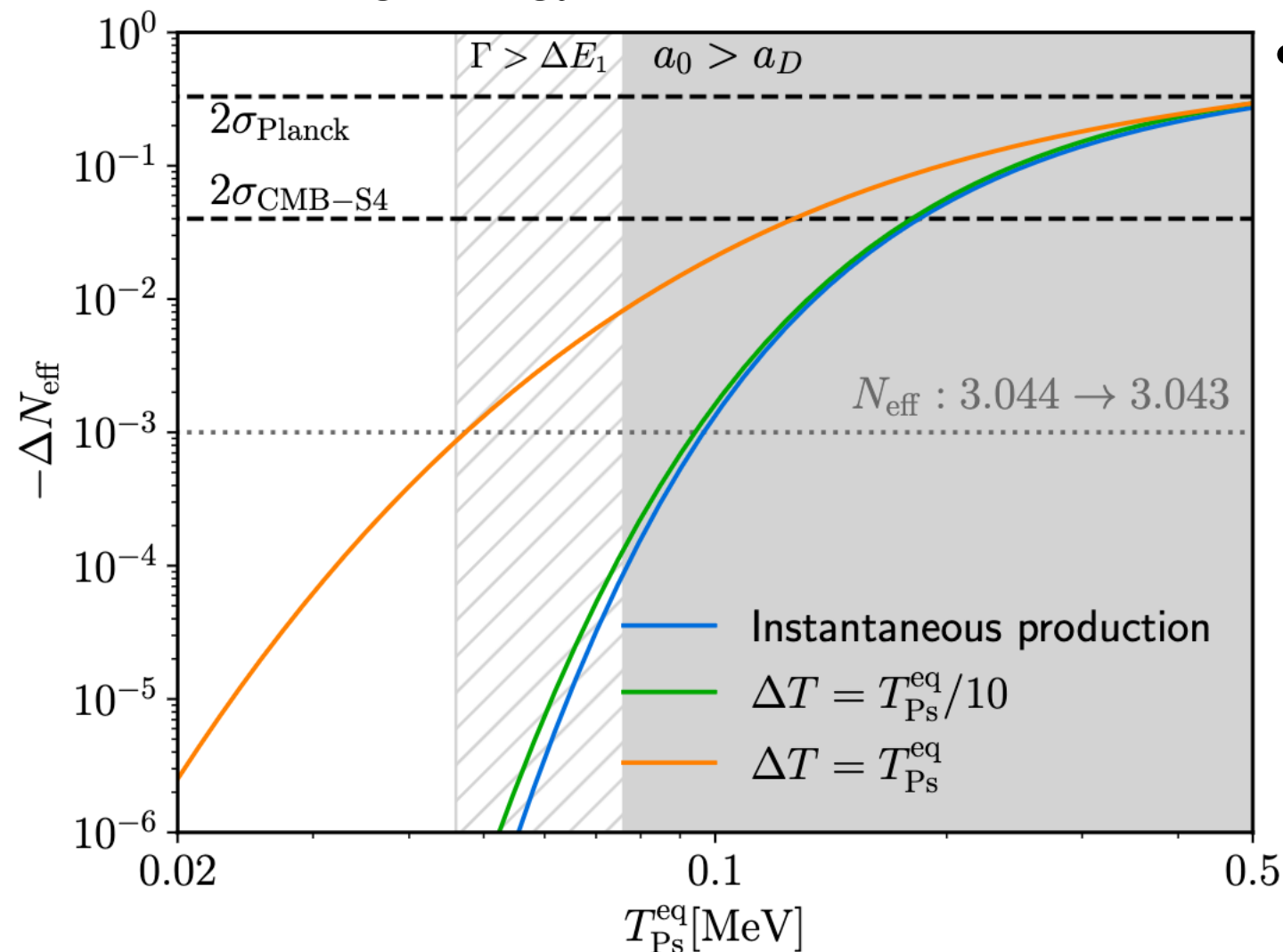
Jackson/Laine [2312.07015](#), [2412.03958](#)

Positronium Formation

Bound State Formation

- Transient population of positronium affects QED equation of state

Approach I: Estimate impact of entropy change from adding bound states
(assuming energy conservation)



- Bound state properties: heuristic description by finite T effective potential

$$V(r) = -\alpha m_D - \frac{\alpha}{r} e^{-m_D r} - i\alpha T \phi(m_D r)$$

$$\phi(y) = 2 \int_0^{+\infty} du \frac{u}{(u^2 + 1)^2} \left(1 - \frac{\sin(uy)}{uy} \right)$$

- Formation occurs when Yukawa potential accommodates bound states that are separated from the continuum by a gap exceeding their thermal width

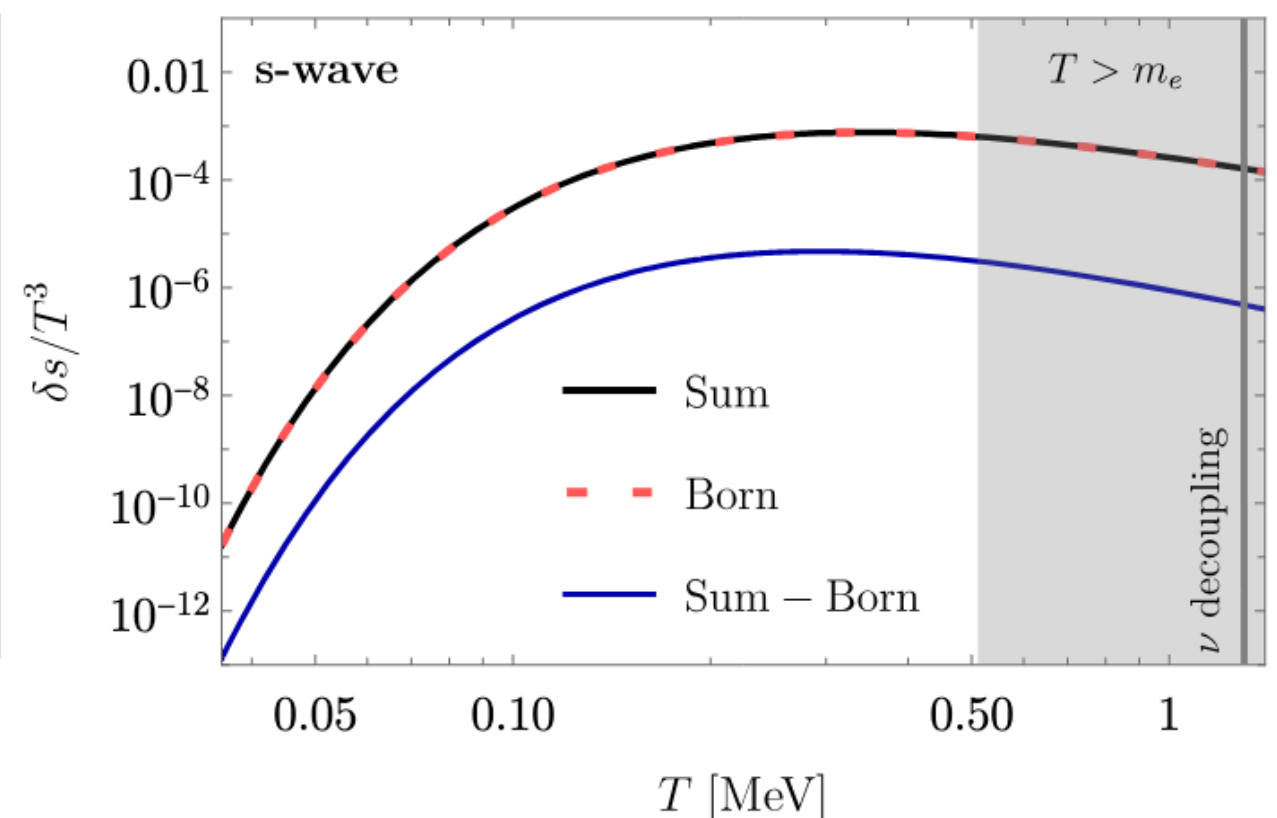
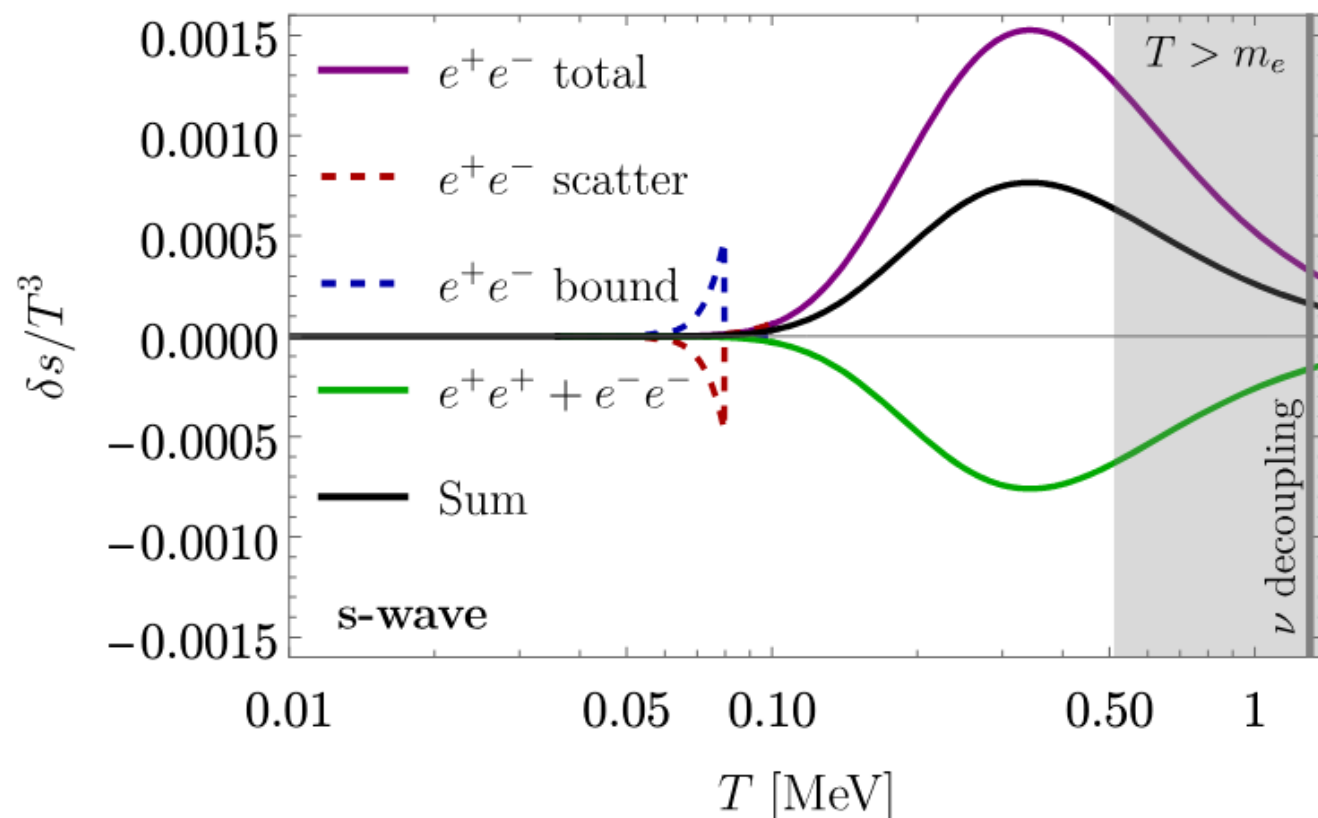
- Magnitude of effect strongly depends on when positronium forms

Bound State Formation

- Transient population of positronium affects QED equation of state

Approach II: Quantify deviation from ideal gas assuming equilibrium
(assuming energy and entropy conservation)

- For $T \gg m_e$: system is perturbative, use results from [Bennett et al 1911.04504](#)
- For $T \ll m_e$: non-perturbative treatment with Beth-Uhlenbeck formula

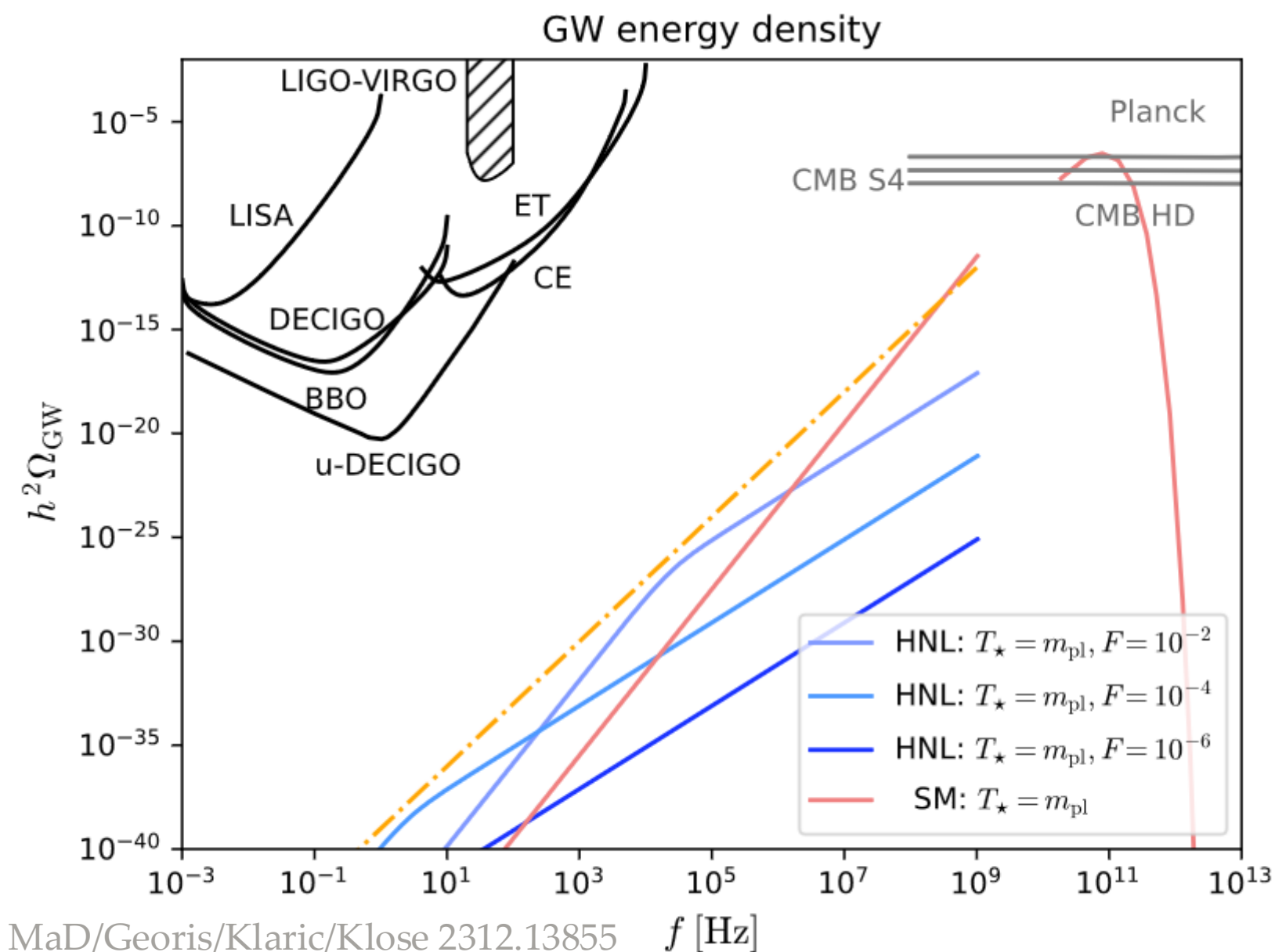


- If the system remains in local equilibrium, impact on N_{eff} is below 10^{-6}

GW Contribution

GWs as a Big Bang Thermometer

- Any plasma in thermal equilibrium emits a black body spectrum of GWs, the graviton equivalent to the CMB and CvB Ghiglieri/Laine 1504.02569
- This irreducible GW background contributes to N_{eff} in the SM, but depends on T_{re}



Ringwald/Schütte-Engel/Tamarit 2011.04731
 Ghiglieri/Jackson/Laine/Zhu 2004.11392
 MaD/Georis/Klaric/Klose 2312.13855

- Direct detection is extremely difficult
- Contribution to N_{eff} imposes the first “direct” constraints on the reheating temperature T_{re}

The Need for Speed
(if you wanna go BSM)

Speeding up N_{eff} Computations

- Suppose the conclusions from Jackson/Laine 2312.07015 & 2412.03958 and MaD et al 2402.18481 are correct and NLO corrections from the collision integral are negligible, so that $N_{eff} = 3.044$ in the SM... then, what next?
- Precision computation sampling the v-momentum distributions in the non-linear collision term are numerically demanding... not suitable to scan parameter spaces of BSM models
- Using moments (instead of sampling momenta) massively reduces the computation time and only leads to errors around 0.04%, recovering $N_{eff} = 3.044$; a public code is provided in Escudero/Jackson/Laine/Sandner 2511.04747
- Direct Simulation Monte Carlo methods can further accelerate the computations, Ihnatenko/Ovchinnikov 2508.08379 recover $N_{eff} = 3.0439 \pm 0.0006$ in the SM (compared to $N_{eff} = 3.0440 \pm 0.0002$ in Bennett et al 2012.02726)

Summary & Conclusions

- In 2020 a new benchmark value $N_{eff} = 3.044$ in the SM was established that includes neutrino oscillations and NLO corrections to the QED EoS
Akita/Yamaguchi et al 2005.07047, Froustey/Pitrou 2008.01074, Bennett et al 2012.02726
- NLO corrections to the collision term have been studied by different authors, suggesting $N_{eff} = 3.043$ in Cielo et al 2306.05460
 $N_{eff} = 3.044$ in Jackson/Laine 2312.07015 & 2412.03958 and MaD et al 2402.18481
The discrepancy has not (yet) been resolved, my personal working assumption remains $N_{eff} = 3.0440 \pm 0.0002$ based on Bennett et al 2012.02726
- Bound state formation is unlikely to have a sizeable impact Binder et al 2411.14091
- GW contribution to N_{eff} in the SM is detectable for very high reheating temperature
Ghiglieri et al 1504.02569 & 2004.11392, Ringwald et al 2011.04731, MaD et al 2312.13855
- Beyond the SM increasing speed may be more important than increasing precision; e.g.
using moments instead of momenta Escudero/Jackson/Laine/Sandner 2511.04747
using Direct Simulation Monte Carlo Ihnatenko/Ovchynnikov 2508.08379