

DARK PION DARK MATTER AND PHENOMENOLOGY

Lingfeng Li
李凌风

Mar.5, 2026

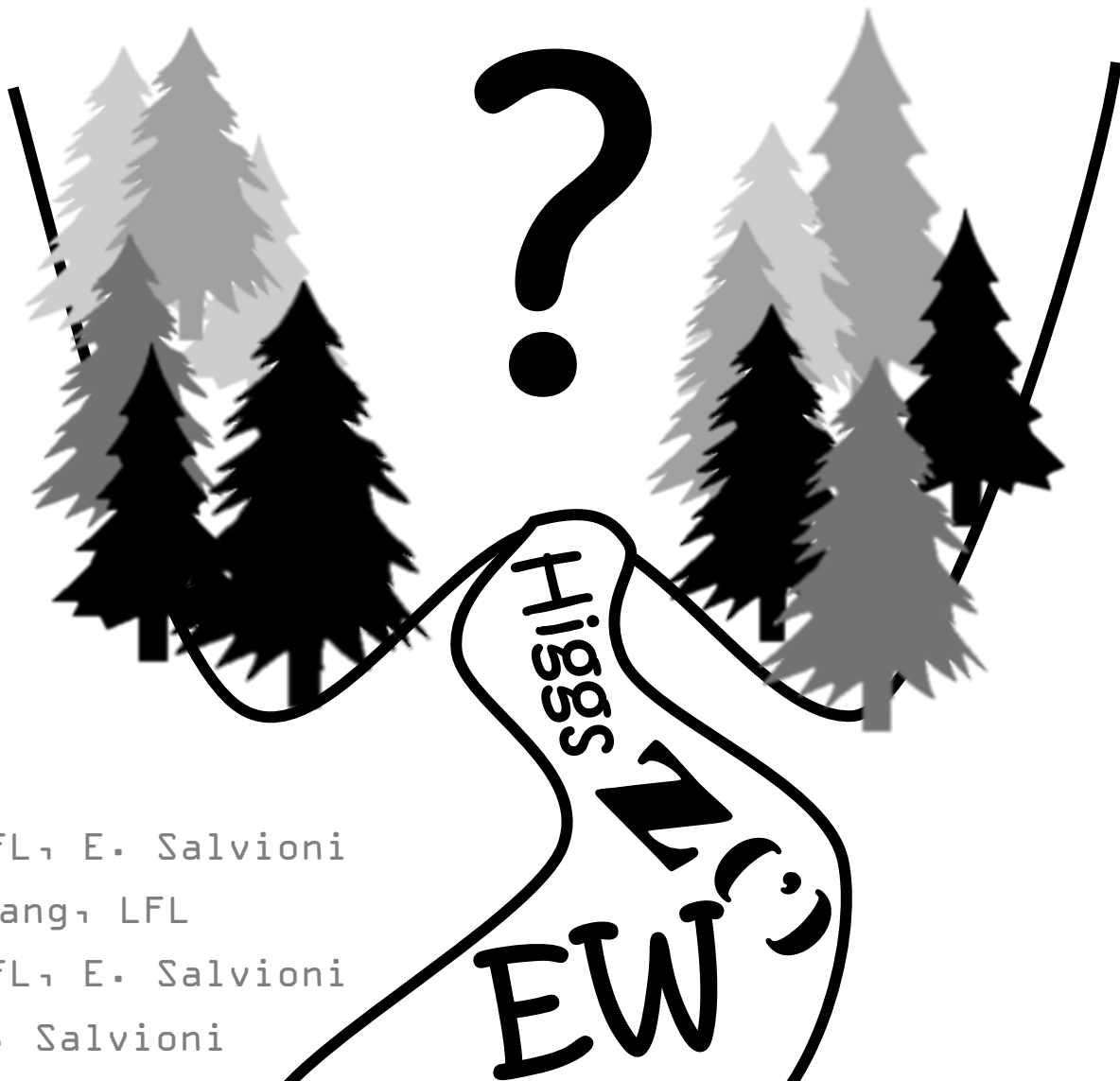


26xx.abcd H-C. Cheng, X. Jiang, LFL, E. Salvioni

2408.13304 H-C. Cheng, X. Jiang, LFL

2401.08785 H-C. Cheng, X. Jiang, LFL, E. Salvioni

2110.10691 H-C. Cheng, LFL, E. Salvioni



The overview of this talk

The Electroweak Portal...

COSMIC

- The legacy of SM as a tool
- “Irrelevant” portal EFT
- EW precision tests
- Flavor portal

To the Hidden Valley ...

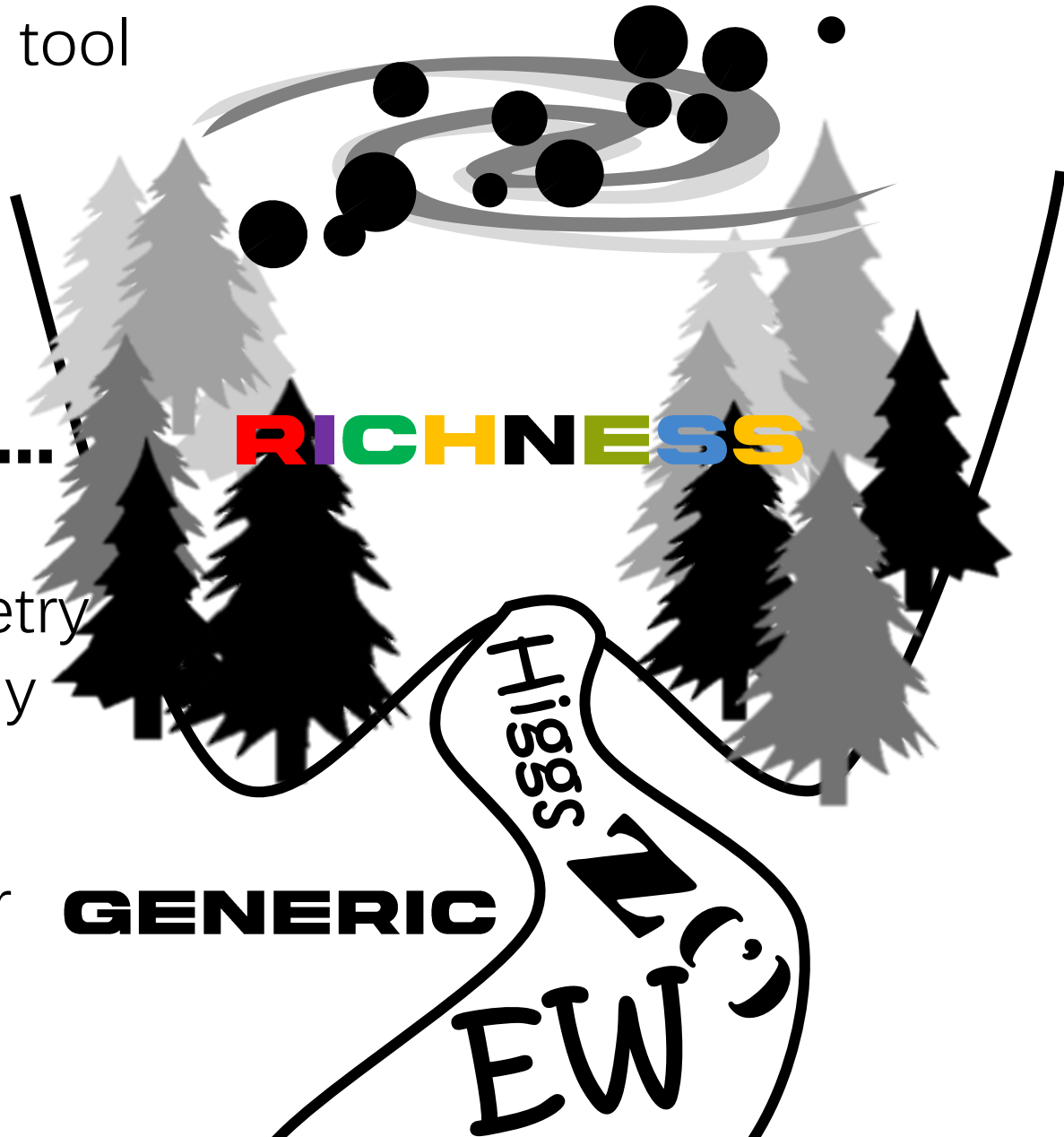
RICHNESS

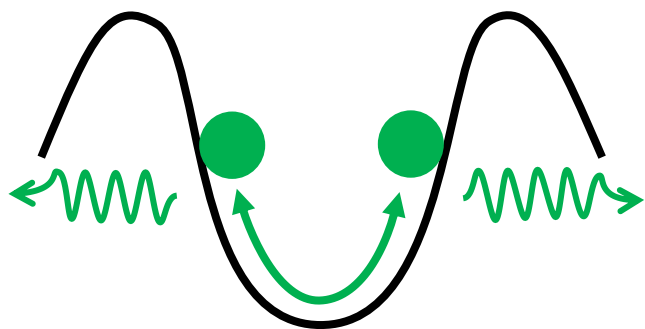
- Long-lived particles
- Dark hadrons & symmetry
- UV& IR Phenomenology

...and Beyond

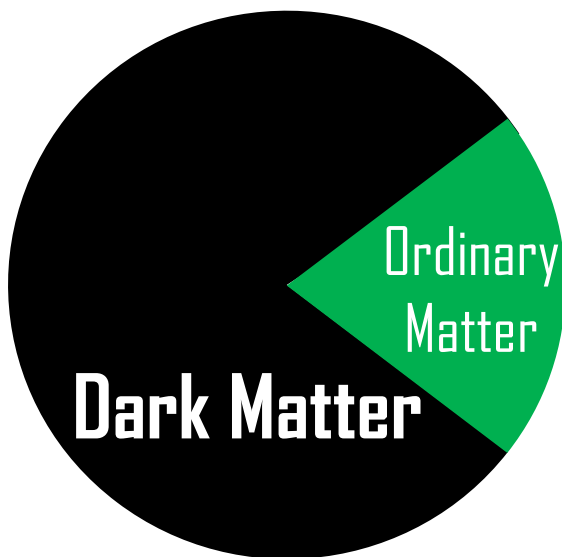
- Composite Dark Matter
- Cosmological Probes

GENERIC





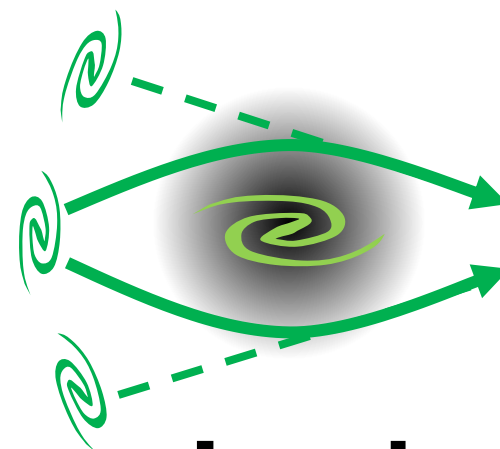
**CMB Power
Spectra**



**Large Scale
Structure**



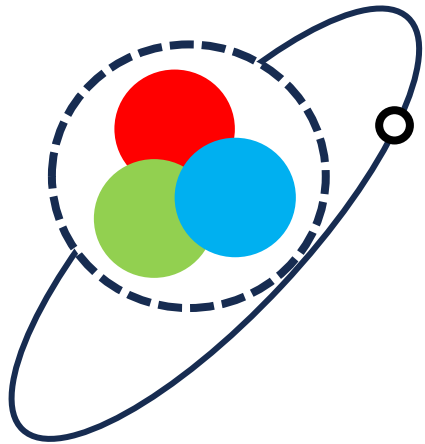
**Galaxy
Formation**



Lensing

Motivation of Compositeness

"Baryons" are Composite



Hierarchy
Asymmetry

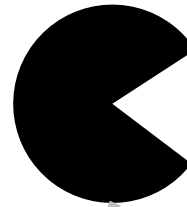
GUT

FOPT

CFT

UV Model Building

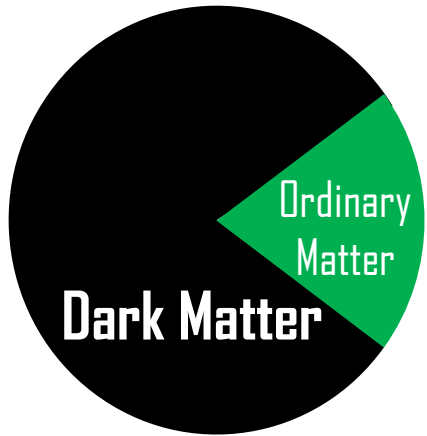
"Minimal" Setup (?)



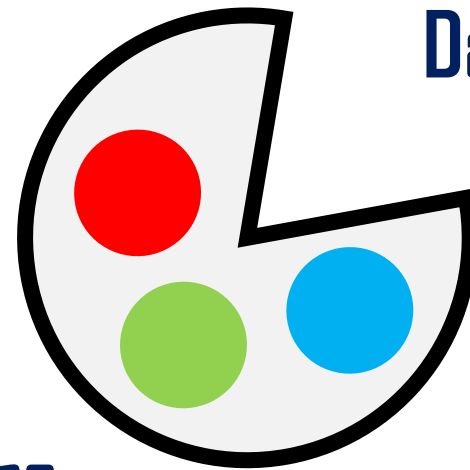
G
only?



Unique Dynamics & Phenomena



Composite DM Candidates

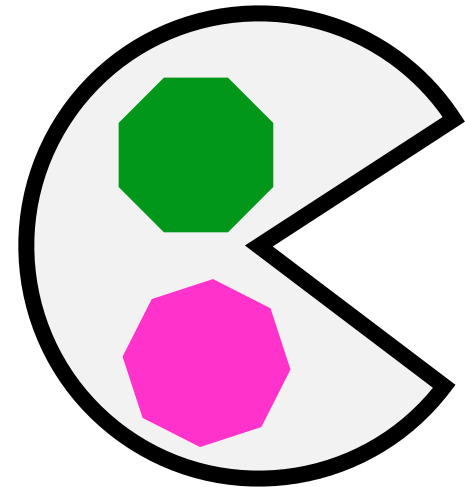
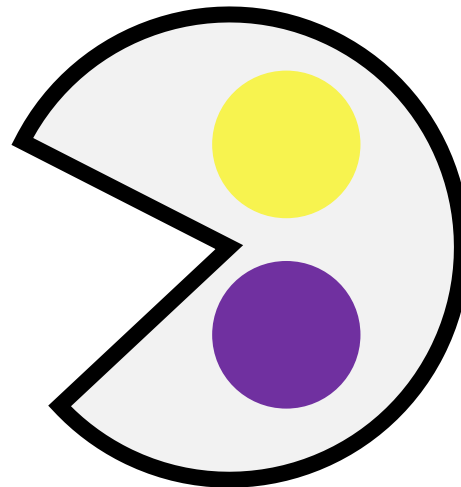


Dark Baryons



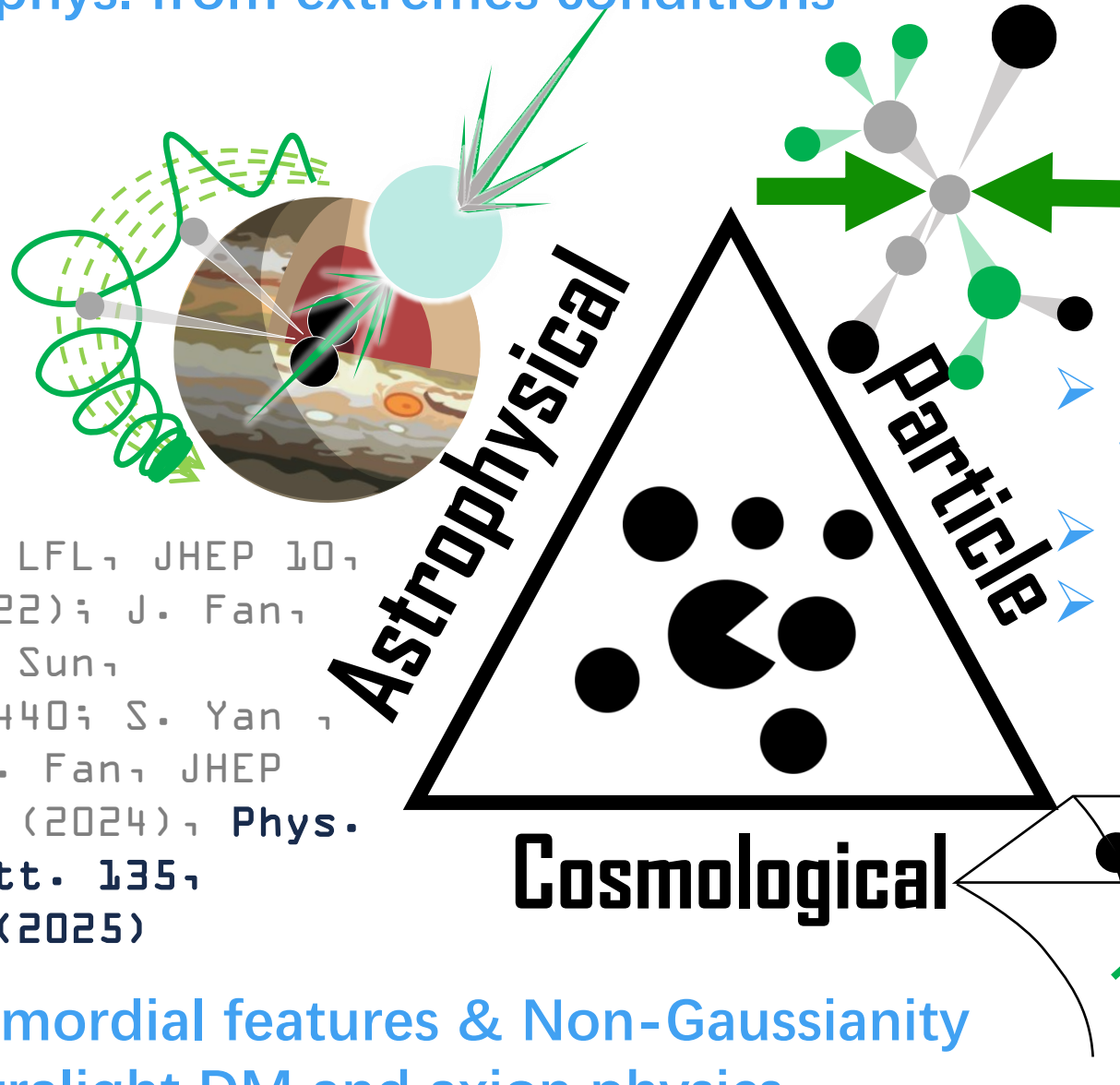
Dark R-hadrons

Dark Mesons



Dark Blueballs

- Planetary science implications
- New phys. from extremes conditions



J. Fan, LFL, JHEP 10, 186 (2022); J. Fan, LFL, C. Sun, 2501.12440; S. Yan, LFL, J. Fan, JHEP 06, 028 (2024), Phys. Rev. Lett. 135, 231801 (2025)

- Primordial features & Non-Gaussianity
- Ultralight DM and axion physics

LFL, S. Lu, Y. Wang, S. Zhou, JHEP 07, 231 (2020); L. Fong, LDL, T. Liu, H. Lu, T. Qiu, H. Tye, JCAP 08, 057 (2021); Y. Bao, J. Fan, LFL, Phys. Rev. Lett. 130, no.24, 241001 (2023); X. Chen, J. Fan, LFL, JHEP 12, 197 (2023)

S. Chai, J. Gu, LFL JHEP 05, 292 (2024); T.S. M. Ho, X. H. Jiang, T. H. Kwok, LFL, T. Liu, Phys. Rev. D 109, no.9, 093004 (2024);

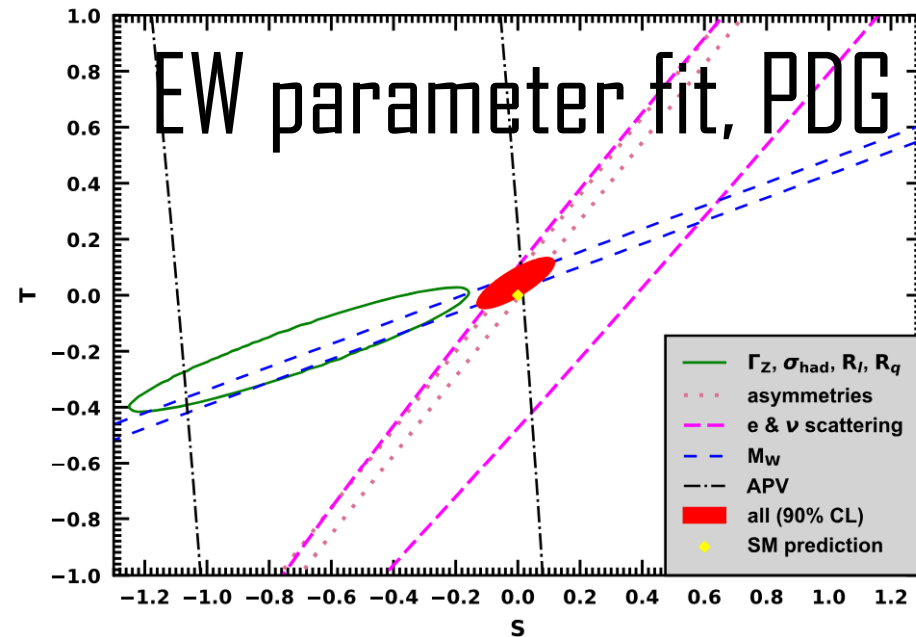
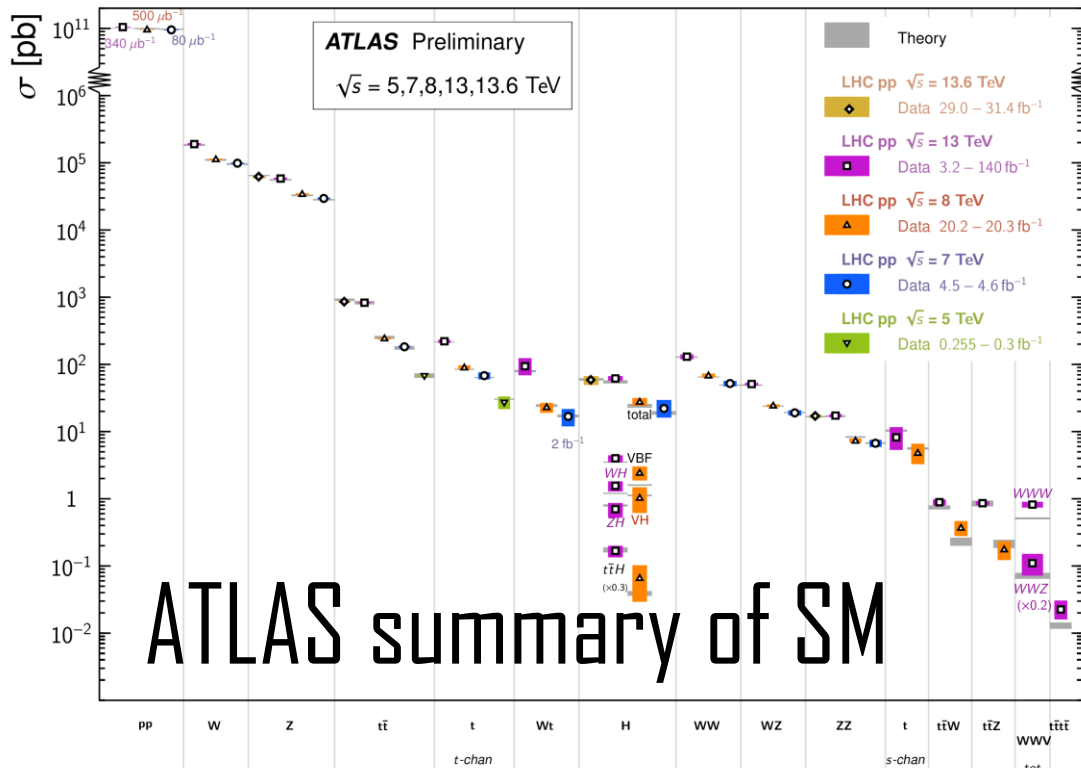
- Pheno of Hidden Valley
- Flavor tests of SM
- Machine learning developments

SM Legacy as Tools

- SM is “complete” and known with reasonable precision

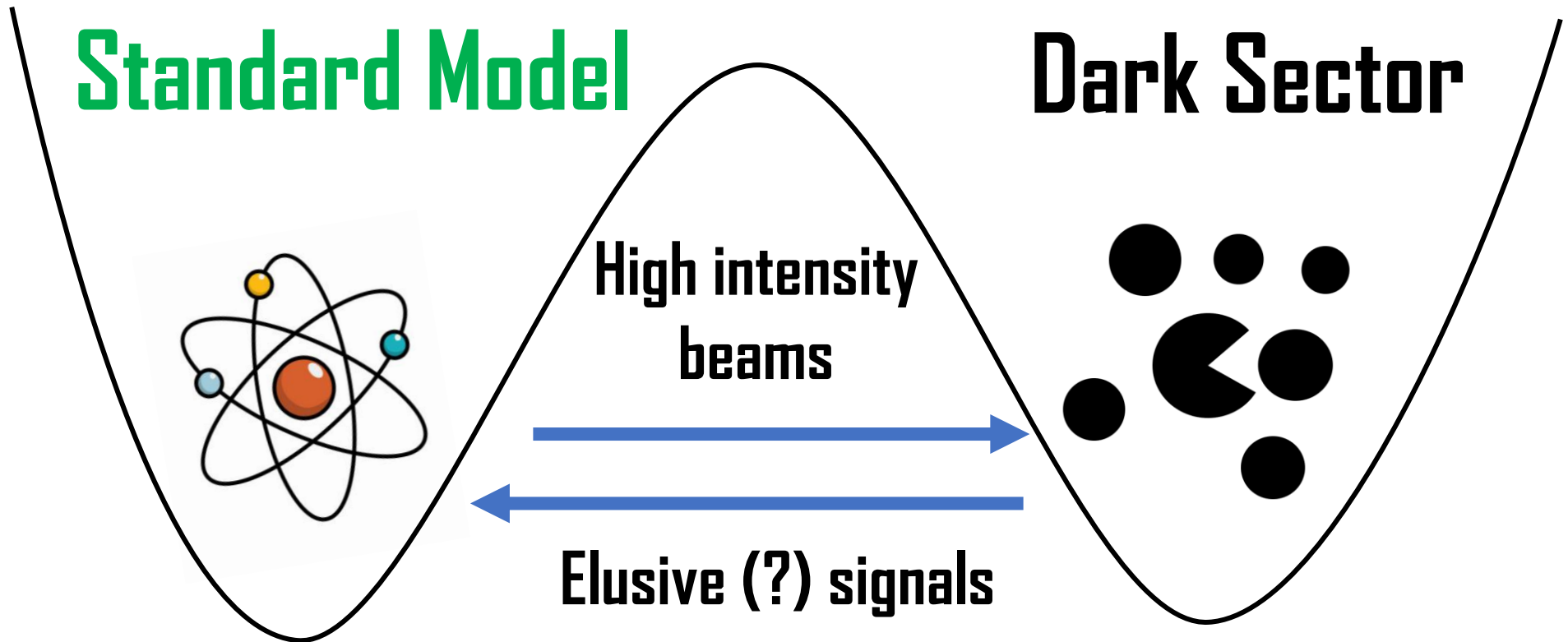
Standard Model Total Production Cross Section Measurements

Status: October 2023



- What can we learn from all these?

The Irrelevant Portal



$\epsilon \mathcal{O}_{\text{SM}} J_{\text{Dark}}$ **Relevant portal:** small couplings to keep the valley hidden

$\frac{1}{\Lambda^n} \mathcal{O}_{\text{SM}} J_{\text{Dark}}$ **Irrelevant portal:** Hidden behind EFT

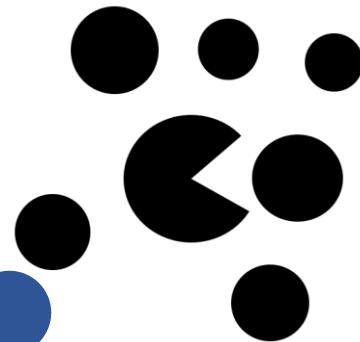
The Electroweak Portal

$$\frac{c}{M^2} (iH^\dagger \overleftrightarrow{D}_\mu H) (\bar{\psi} \gamma^\mu \psi)$$



h

Z

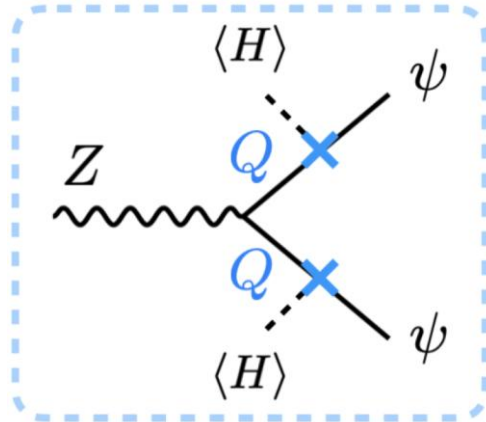


EW bosons stand for a scale of 246 GeV, close enough to TeV+

>TeV Scale $\approx$ Dark Enough

The EW Portal: $\frac{c}{M^2} (iH^\dagger \overleftrightarrow{D}_\mu H) (\bar{\psi} \gamma^\mu \psi)$ from Two UV Setups

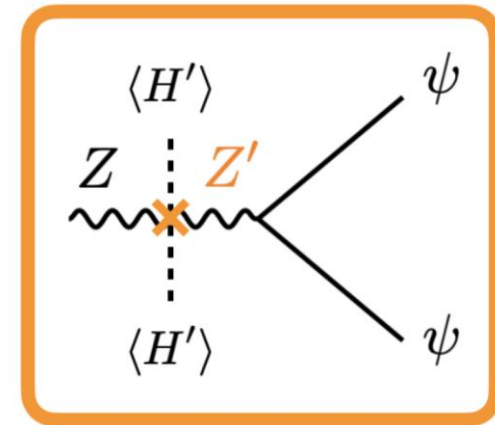
$$\mathcal{L}_{\text{UV}} \supset \bar{Q} Y \psi H \longrightarrow \frac{c}{M^2} \sim \frac{Y^2}{M_Q^2}$$



Heavy fermion doublet mediation

H-C. Cheng, LFL, E. Salvioni, 2110.10691

$$\mathcal{L}_{\text{UV}} \supset \delta \hat{M}^2 Z^\mu Z'_\mu \longrightarrow \frac{c}{M^2} \sim \frac{g_D^2 \delta \hat{M}^2}{m_Z^2 m_{Z'}^2}$$



Dark Z' mediation: a Z' heavier or lighter(!) than Z with mass mixing

H-C. Cheng, X. Jiang, LFL, E. Salvioni, 2401.08785

See also Y. C. San, M. Perelstein, and P. Tanedo, 2022

Exotic BR of Z

A key phenomenon for EW portal models

As the doublet mediation benchmark:

$$\text{BR}(Z \rightarrow \psi' \bar{\psi}') \approx 1.8 \times 10^{-4} \left(\frac{N_d \text{Tr}(\mathbf{Y} \mathbf{Y}^\dagger \mathbf{Y} \mathbf{Y}^\dagger)}{3} \right) \left(\frac{1 \text{ TeV}}{M} \right)^4$$

For the Z' mediation :

$$\text{BR}(Z \rightarrow \psi' \bar{\psi}') \approx 1.4 \times 10^{-4} \left(\frac{N_d}{3} \right) \left(\frac{\xi}{0.01} \right)^2 g_D^2 \sum_i (x_{Li}^2 + x_{Ri}^2)$$

$$\text{The effective mixing angle } \xi \sim \frac{\delta M^2}{M_Z^2 - M_{Z'}^2}$$

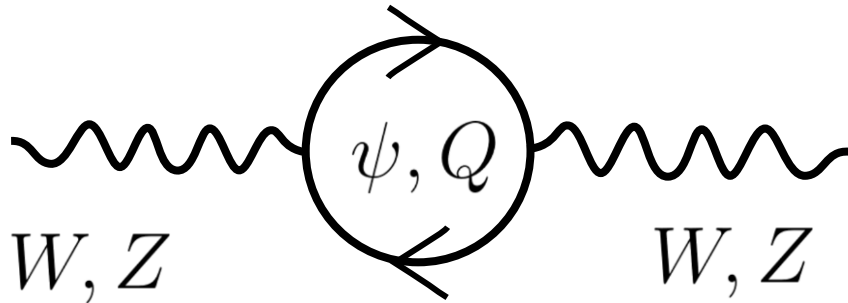
Exotic BR of h is not large ($< \sim 10^{-4}$), the constraints are weaker than those from Z decays ($\sim 10^3$ more Z than h at LHC)

EWPT Corrections

M.E. Peskin, T. Takeuchi, 1990

M.E. Peskin, T. Takeuchi, 1992

Doublet mediation



Leading by EW loops with mixing

L. Lavoura, J. P. Silva, 1993;

C. Anastasiou, E. Furlan, J.

Santiago, 2009

$$S \simeq \frac{2N_d}{45\pi M^2} \left(\frac{5Y^2 v^2}{2} - 3m_Z^2(1 - 2s_W^2 + 2s_W^4) \right)$$

$$T \simeq \frac{N_d y_{11}^4 v^4}{48\pi s_W^2 m_W^2 M^2}$$

Z' mediation



Leading by LEP Z pole observables defined with:

$$\mathcal{L}_{Zf\bar{f}} = -\frac{\bar{Z}e}{s_W c_W} \bar{f} \gamma^\mu (T_{Lf}^3 - s_*^2 Q_f) f Z_\mu$$

B. Holdom, 1991; G. Altarelli

et. al, 1991

$$S = \frac{4s_W^2}{\alpha} \left(\frac{c_W^2}{s_W} \xi t_\chi - c_W^2 \xi^2 \right)$$

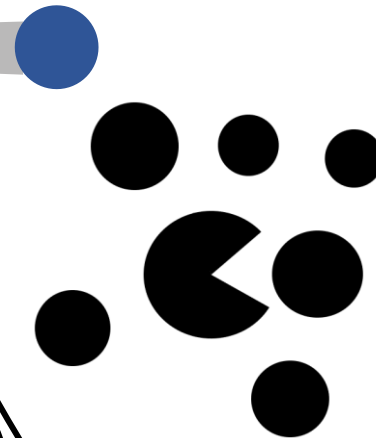
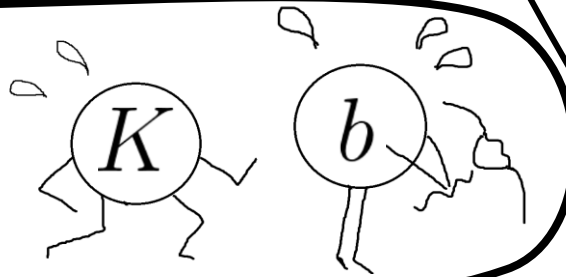
$$T = \frac{1}{\alpha} \left(2s_W \xi t_\chi + \xi^2 \left(\frac{M_{Z'}^2}{M_Z^2} - 2 \right) \right)$$

The Flavor Portal

$$\frac{c}{M^2} (iH^\dagger \overleftrightarrow{D}_\mu H) (\bar{\psi} \gamma^\mu \psi) \quad \text{Tree level, meson decay and dark shower}$$



$h \ Z$



$$\frac{\kappa}{M^2} (\bar{d}_i \gamma^\mu d_j) (\bar{\psi} \gamma_\mu \psi)$$

W. Altmannshofer, A. Crivellin, H. Haigh, G. Inguglia and J. Martin Camalich, 2311.14629; J. Martin Camalich, M. Pospelov, P. N. H. Vuong, R. Ziegler and J. Zupan, 2002.04623; M. K. Gaillard, M. B. Gavela, R. Houtz, P. Quilez and R. Del Rey, 1805.06465; T. Li, M.A. Schmidt and M. Yuan, 2506.08342.....

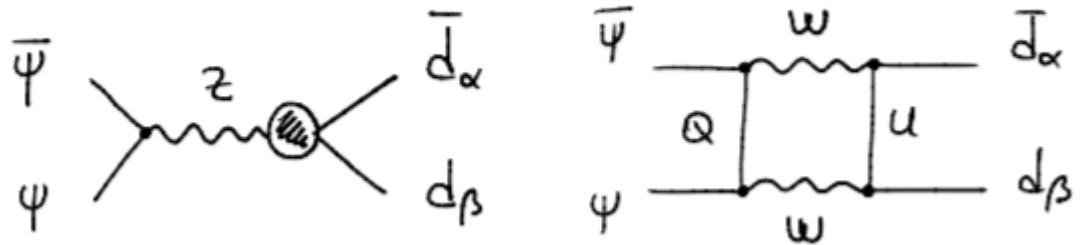
Flavor Portal from EW only

$\Gamma_{B,K,\dots}$ are suppressed by $(M_W)^{-4}$ in SM.

Inami, Lim 1980

X.G. He, J. Tandean,

G. Valencia, 0909.3638



The four-fermion EFT respecting MFV

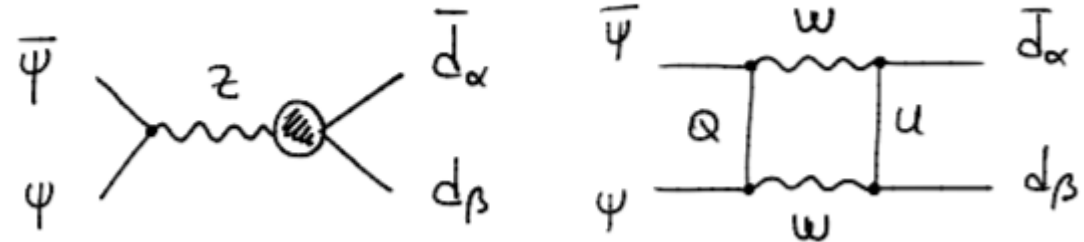
$$-\frac{c}{M^2} \frac{g^2}{128\pi^2} \underline{J_D^\mu \bar{d}_j \gamma_\mu P_L d_i} \sum_{q \in u,c,t} \underline{V_{qj}^* V_{qi} \mathcal{K}_q}$$

Flavor Portal from EW only

Inami, Lim 1980

X.G. He, J. Tandean,

G. Valencia, 0909.3638



In heavy doublet mediation:

$$\mathcal{K}_q \equiv 3 + x_q \left(\log \frac{\Lambda_{UV}^2}{m_q^2} - 2 \right) \quad x_q \equiv m_q^2 / M_W^2$$

In Z' mediation (the mZ' terms cancel out):

Typically $O(10)$

$$\mathcal{K}_q \equiv x_q \log \frac{\Lambda_{UV}^2}{M_W^2} + \frac{-7x_q + x_q^2}{2(1 - x_q)} - \frac{4x_q - 2x_q^2 + x_q^3}{(1 - x_q)^2} \log x_q$$

Does it constrain anything?

We can barely see the SM Z & W in $bs\nu\nu$ and $sd\nu\nu$
(for now...)

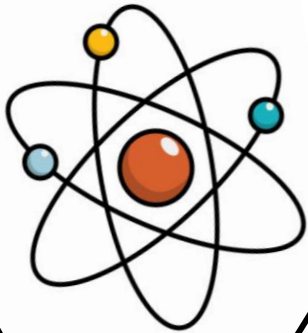
Γ_{39}	$\pi^+ \nu \bar{\nu}$	$(1.14^{+0.40}_{-0.33}) \times 10^{-10}$
Γ_{40}	$\pi^+ \pi^0 \nu \bar{\nu}$	$< 4.3 \times 10^{-5}$
Γ_{582}	$\pi^+ \nu \bar{\nu}$	$< 1.4 \times 10^{-5}$
Γ_{588}	$K^+ \bar{\nu} \nu$	$< 1.6 \times 10^{-5}$
Γ_{589}	$\rho^+ \nu \bar{\nu}$	$< 3.0 \times 10^{-5}$

Want to Kick Things Up?

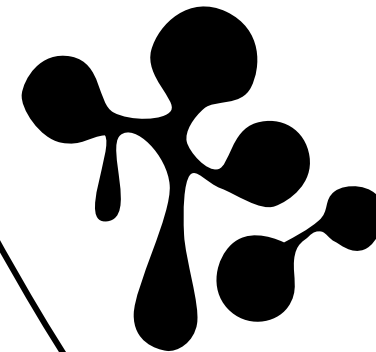


Hidden Valley Models

Standard Model



Confining Dark Sector



High intensity
beams

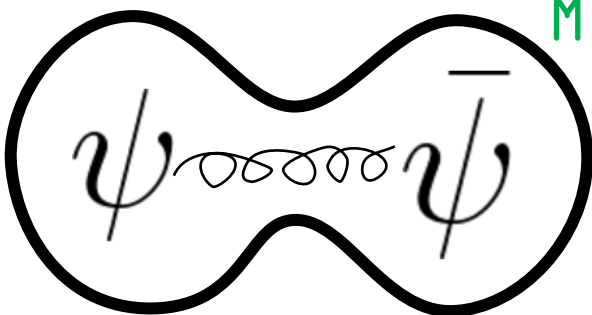


long-lived signals

M. J. Strassler and K. M. Zurek, 0604261

$$\Gamma \sim \text{Tr}[Y^\dagger Y \sigma_{\text{hadron}}] \text{ or } \text{Tr}[X \sigma_{\text{hadron}}]$$

Symmetries Yes

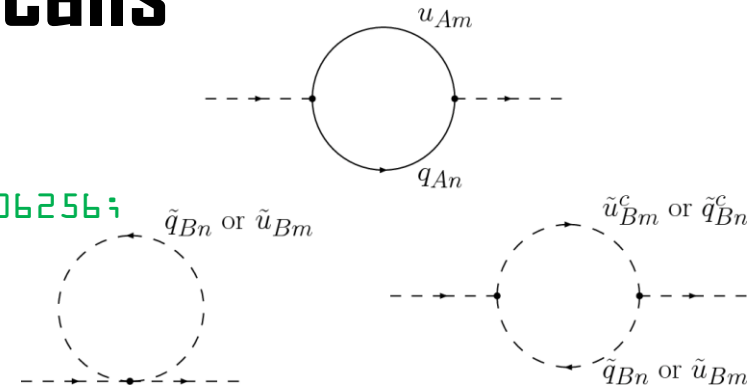


The Dark Force for Confinement



➤ Neutral naturalness calls
for a non-QCD color

Z. Chacko, H.S. Goh, and R. Harnik, 0506256;
G. Burdman, Z. Chacko, H.S. Goh and
R. Harnik, 0609152;
H-C. Cheng, LL, E. Salvioni, and
C. Verhaaren, 1803.03561

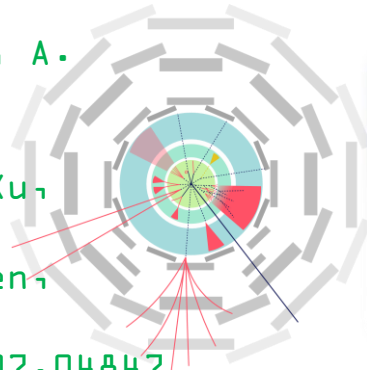


➤ Makes good dark
matter candidate



Y. Hochberg, E. Kuflik,
H. Murayama, T. Volansky,
J. Wacker, 1411.3727;
A. Katz, E. Salvioni, and
B. B. Shakya, 2006.15148; s

P. Schwaller, D. Stolarski, A.
Weiler, 1502.05409;
CMS, 1810.10069;
S. Knapen, J. Shelton, D. Xu,
2103.01238;
S. Born, R. Karur, S. Knapen,
J. Shelton, 2303.04167;
J. Carroasco, J. Zivita, 2307.04847

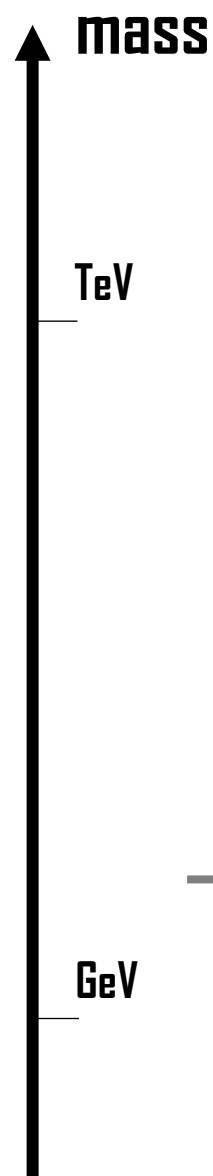


➤ Rich LLP
phenomenology

➤ Duality ? ALPs ? Baryogenesis?

The Cartoon of Dark Spectrum

SM



Dark QCD

$Q_{L,R}$

Dark Force

$\Rightarrow \psi_1 \quad \psi_2 \Leftarrow$



Z'

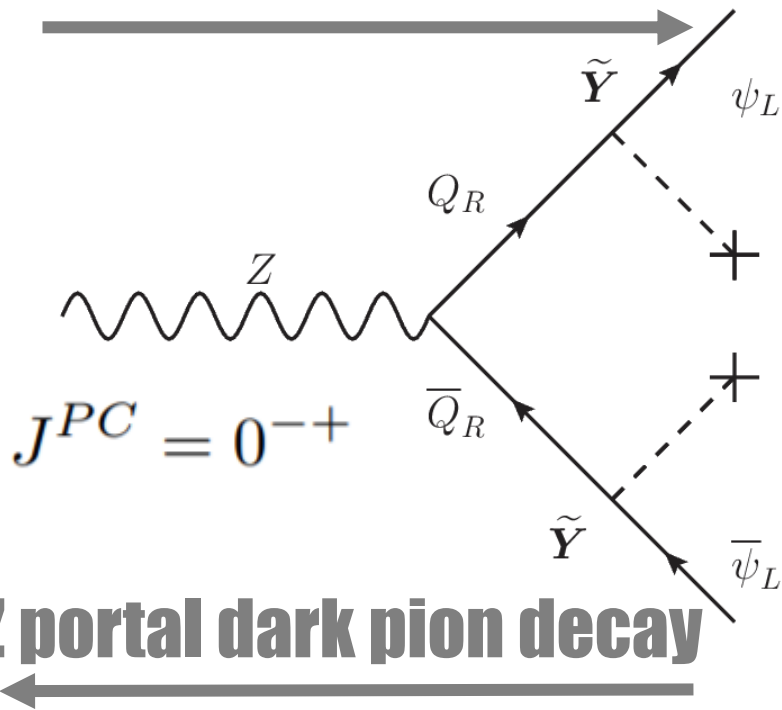
Dark Hadrons



$\psi_{L,R}$

Dark Pions

Z portal dark pion production



Dark pions rearrange into **CP eigenstates** (like K_S and K_L in the SM)

The π_1 and π_3 decay via Z portal, ALP-like (axion-like-particle) with **effective ALP decay constants**:

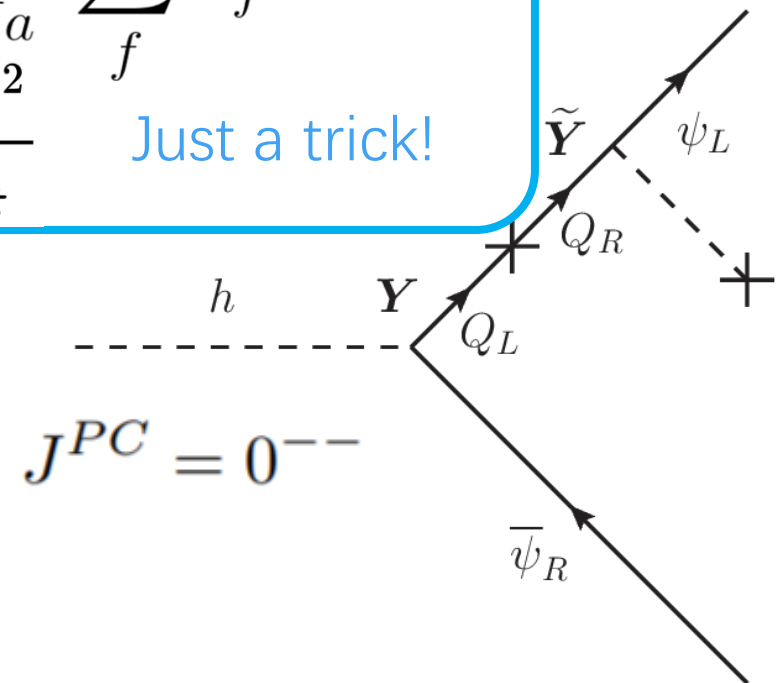
$$\mathcal{L}_a \supset -\frac{\partial_\mu a}{f_a} \sum_f T_f^3 \bar{f} \gamma^\mu \gamma^5 f$$

$$f_a \propto \frac{M^2}{f_{\hat{\pi}}} \quad \text{Just a trick!}$$

Z portal dark pion decay

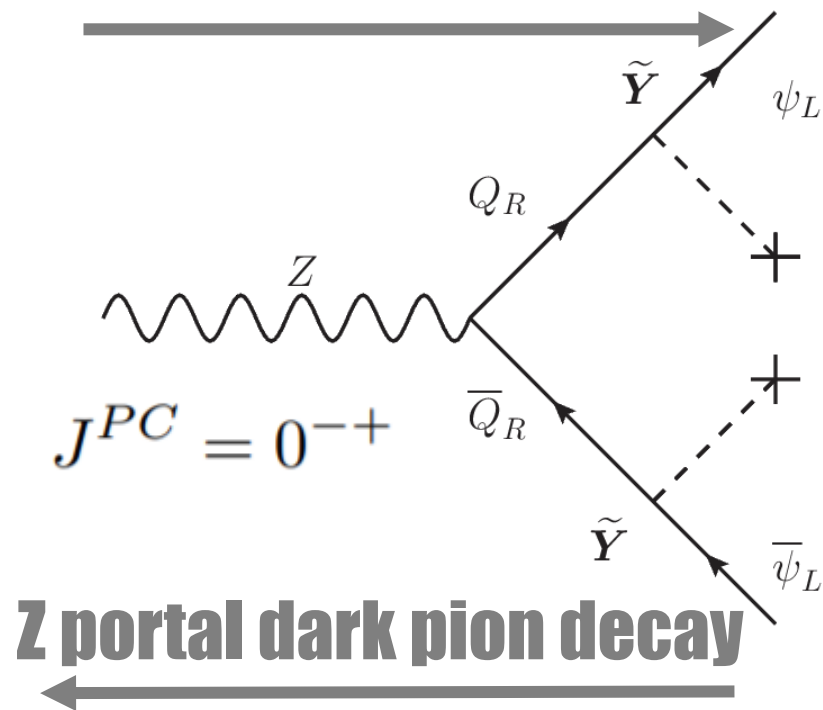
The π_2 mix with the Higgs since it's CP-even, with **mixing angle**:

$$\mathcal{L}_s \supset -\sin \theta_s \frac{m_f}{v} s \bar{f} f, \quad \theta_s \lesssim 10^{-6}$$



Dark Pions

Z portal dark pion production



Dark pions rearrange into **CP eigenstates** (like K_S and K_L in the SM)

The π_1 and π_3 decay via Z portal, ALP-like (axion-like-particle) with **effective ALP decay constants**:

$$\mathcal{L}_a \supset -\frac{\partial_\mu a}{f_a} \sum_f T_f^3 \bar{f} \gamma^\mu \gamma^5 f$$

$$f_a \propto \frac{M^2}{f_{\hat{\pi}}} \quad \text{Just a trick!}$$

$f_{\hat{\pi}}$ The dark meson decay constant, O(confinement scale)

f_a The effective (composite) ALP decay constant, O(PeV) or above

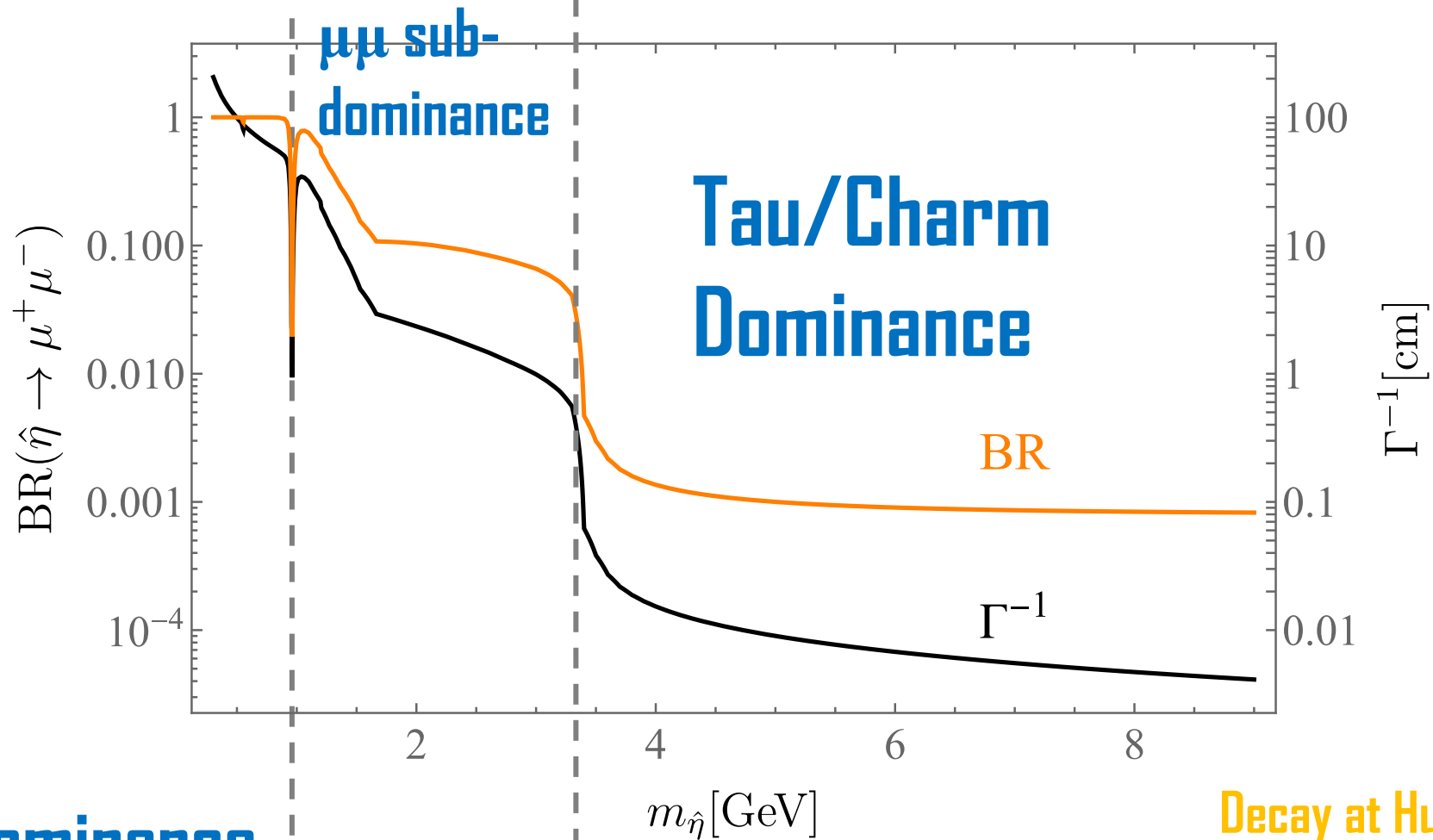
Convenient for the matching with existing ALP searches

Dark (Flavor) Symmetry

- Dark isospin symmetry: identical mass and couplings $\Rightarrow$ All dark pions stable, DM?
- Dark U(1) symmetry: mass and charge matrix diagonal $\Rightarrow$ dark pions carrying non-zero U(1) global charge are stable, DM?
- Dark CP-symmetry: all entries above are real $\Rightarrow$ No mixing between CP-odd and –even species
-

Dark Pion Decays (ALP-Like)

Recent Update: R. Balkin, T. Coren, Y. Soreq and
M. Williams, 2506.15637



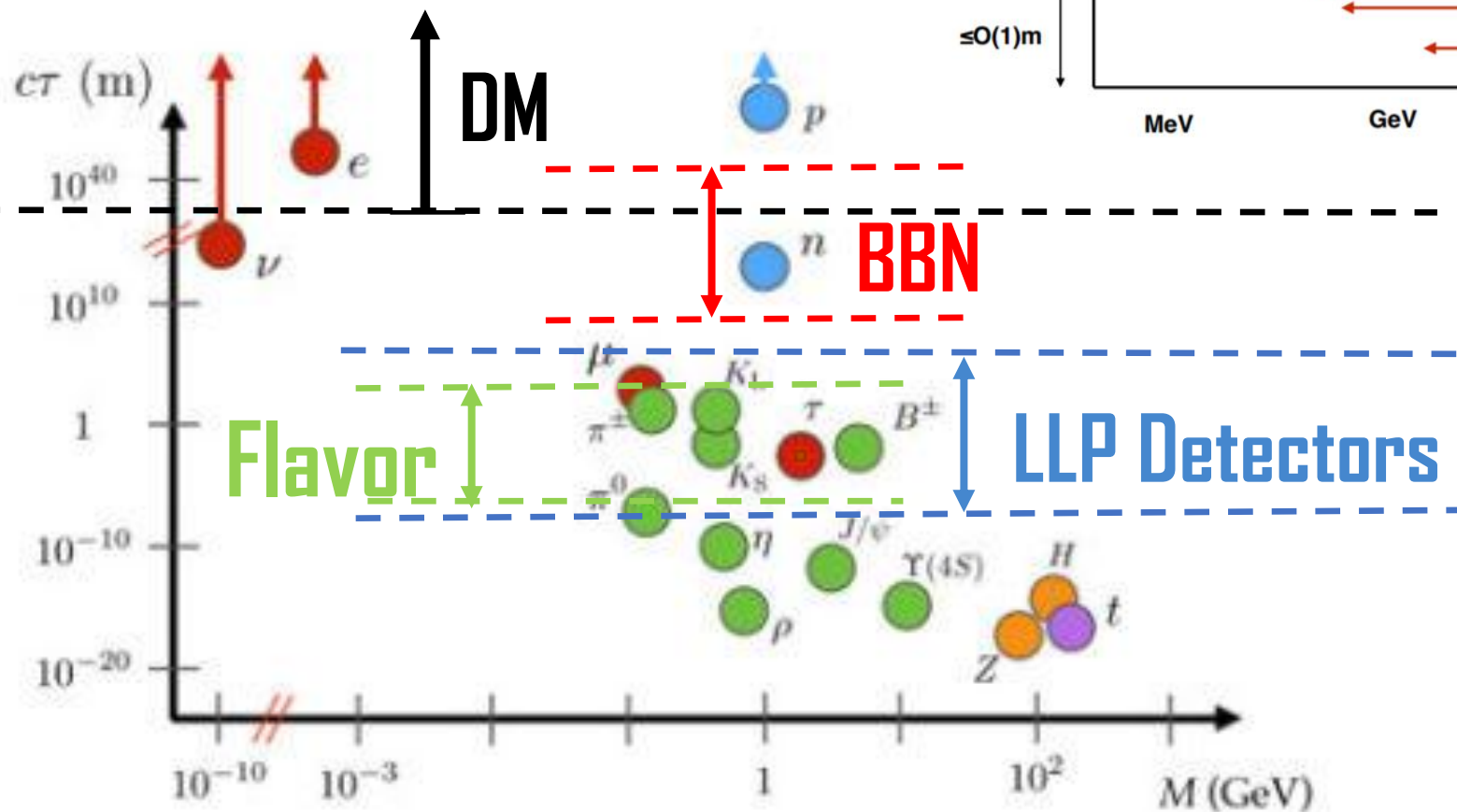
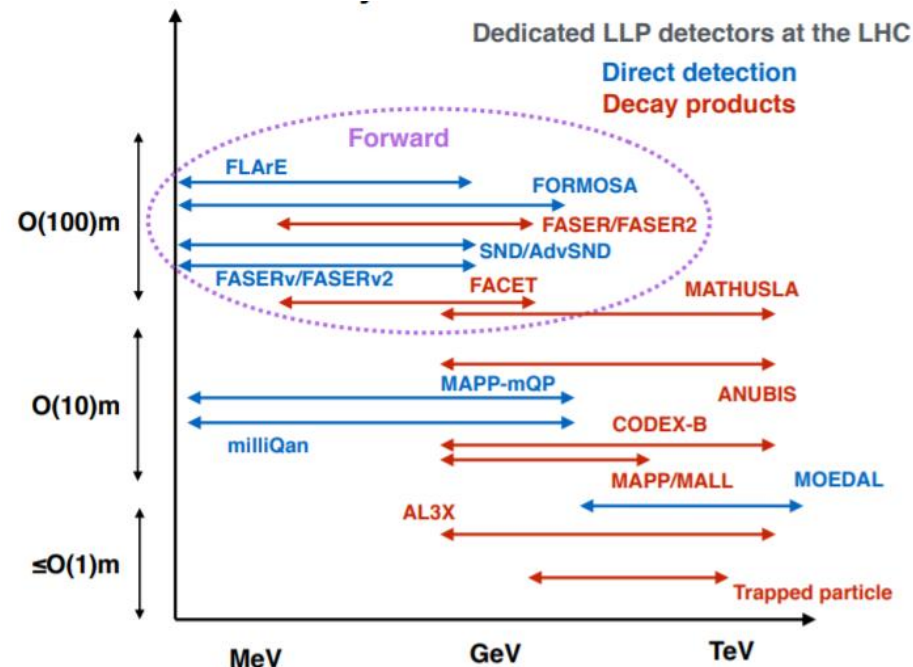
See also:

D. Aloni, Y. Soreq and M. Williams, 1811.03474; M. Ovchinnikov and A.~Zaporozhchenko, 2501.04525; Y. Bai, T. K. Chen, J. Liu and X. Ma, 2505.24822

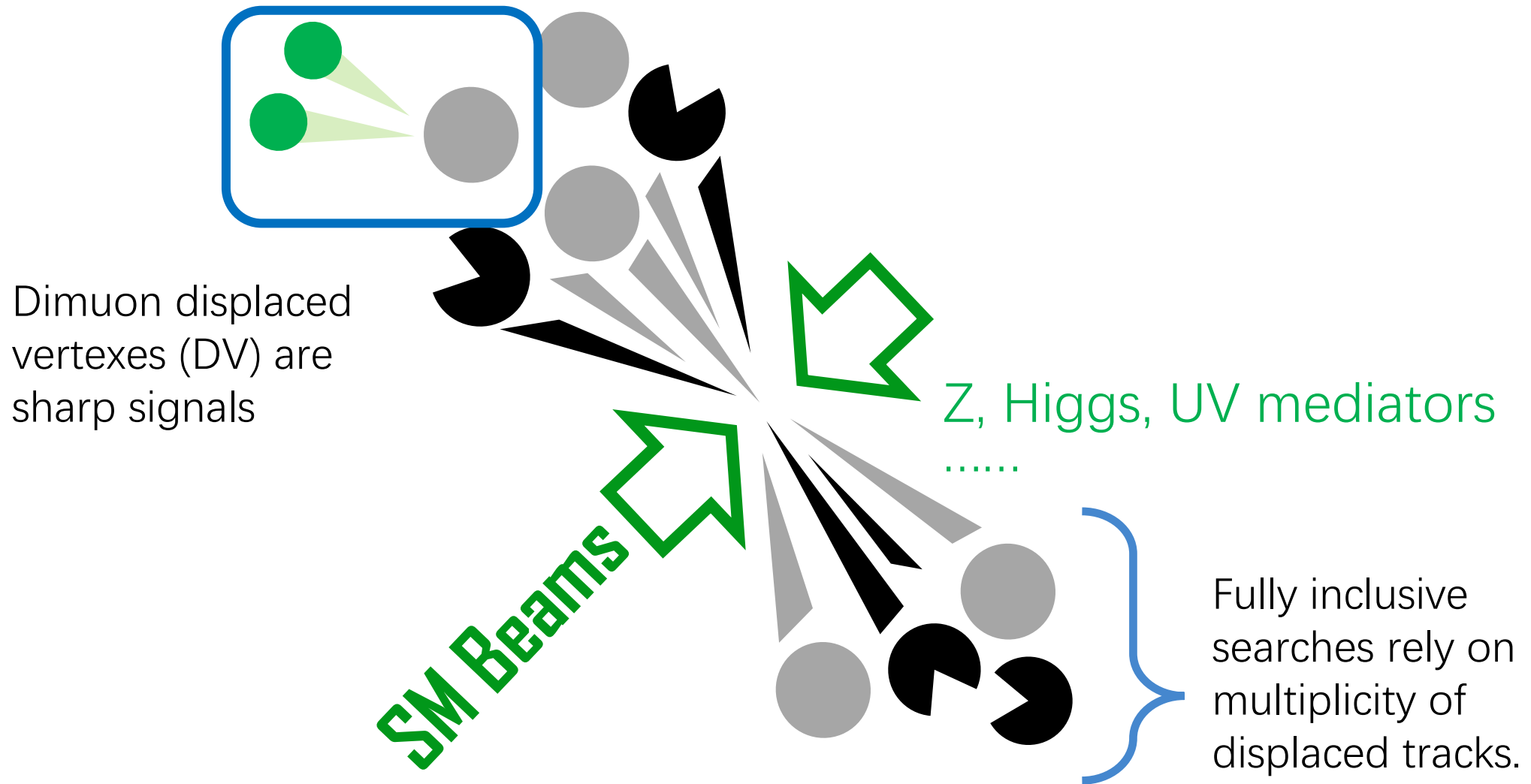
Long-Lived Particles



They are more common and fundamental than it appears



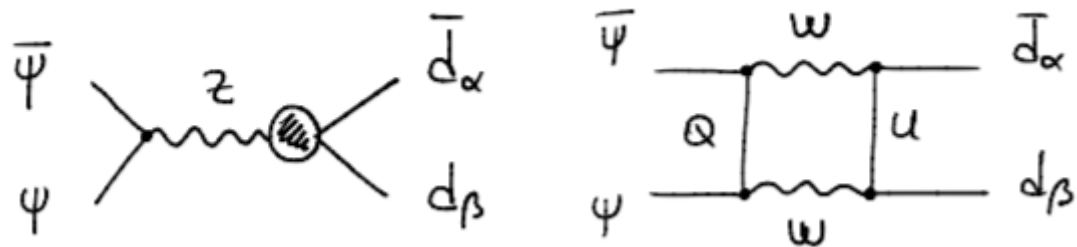
Phenomenology at (Semi-)UV



If any dark pion is an LLP $\rightarrow$ emerging energy outside the innermost detector

P. Schwaller, D. Stolarski, A. Weiler, 1502.05409; CMS, 1810.10069; S. Knapen, J. Shelton, D. Xu, 2103.01238; . Born, R. Karur, S. Knapen, J. Shelton, 2303.04167; J. Carroasco, J. Zurita, 2307.04847

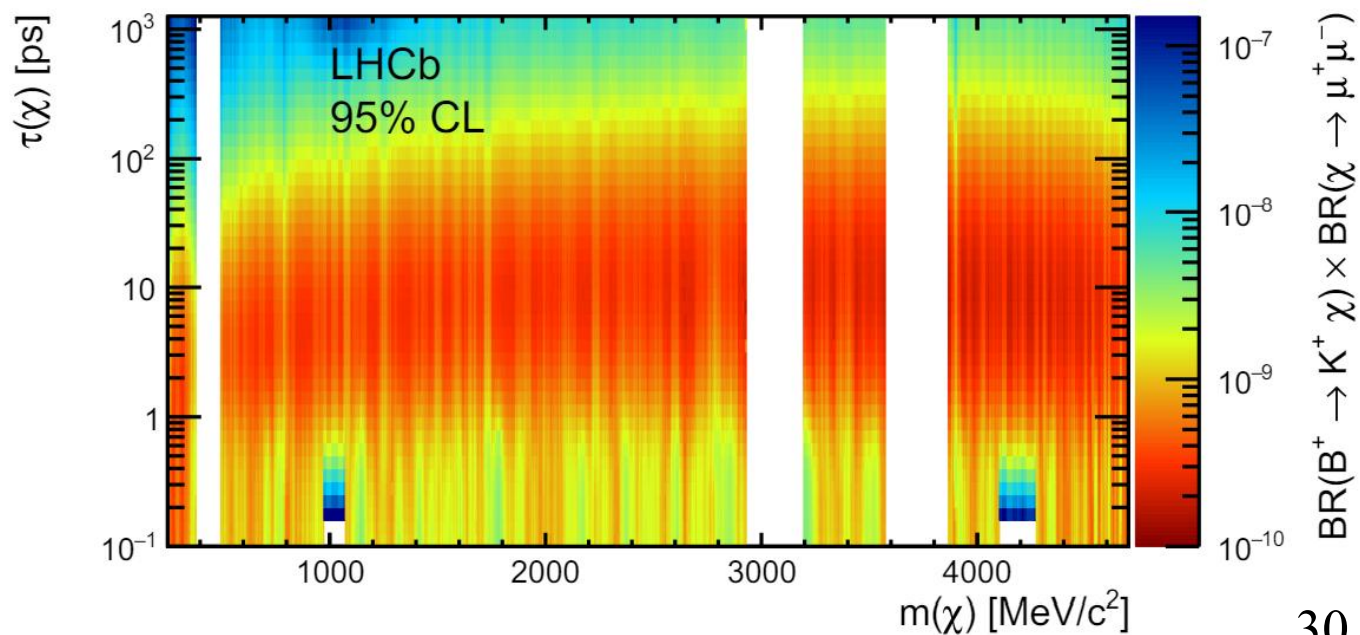
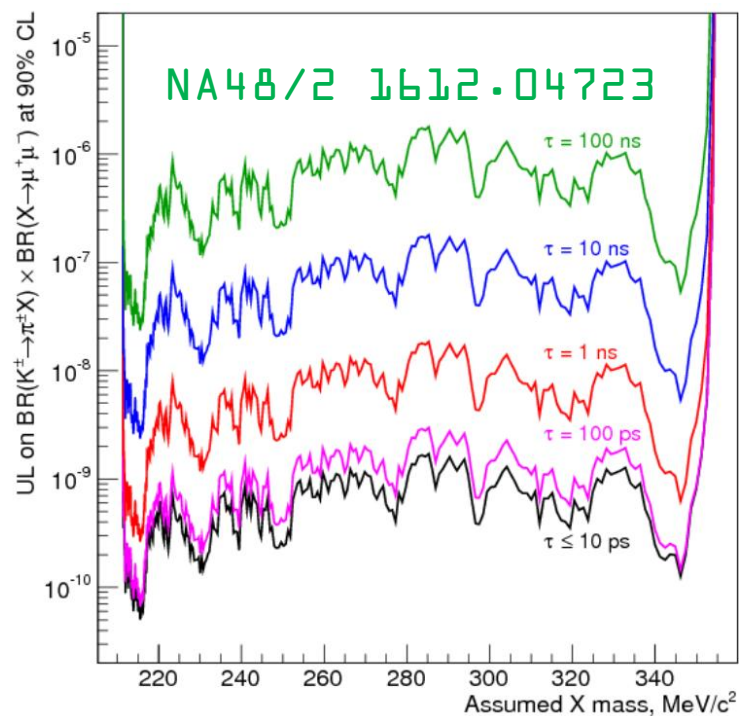
IR: Flavor Portal to LLP



➤ 5-10 times larger BR for inclusive X_s decays

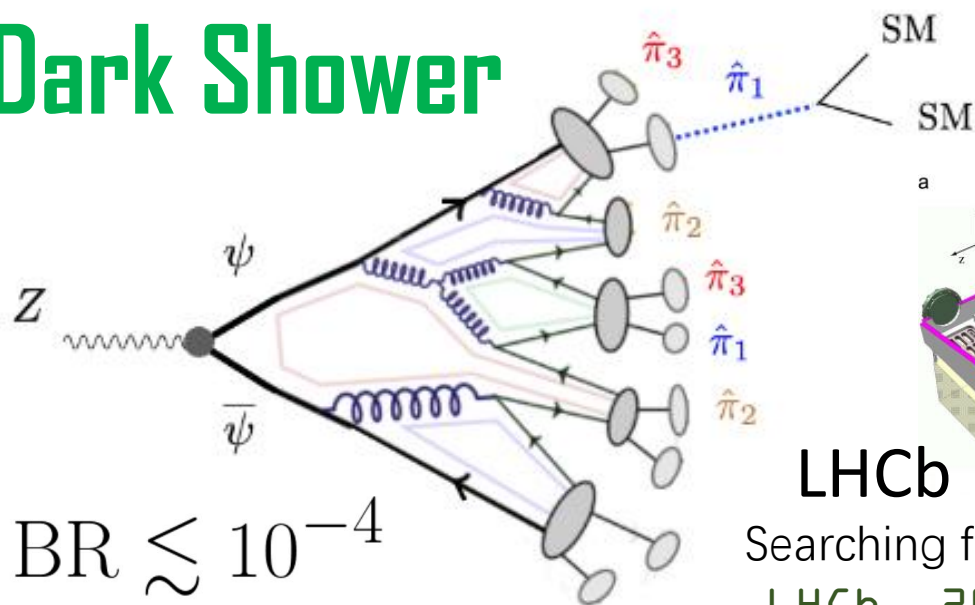
$$\text{BR}(B^{+,0} \rightarrow \{K^+ \hat{\pi}, K^{*0} \hat{\pi}\}) \approx \{0.92, 1.1\} \times 10^{-8} \left(\frac{1 \text{ PeV}}{f_a} \right)^2 \left(\frac{\mathcal{K}_t}{10} \right)^2 \{ \lambda_{BK\hat{\pi}}^{1/2}, \lambda_{BK^*\hat{\pi}}^{3/2} \}$$

$$\text{BR}(K^+ \rightarrow \pi^+ \hat{\pi}) \approx 3.9 \times 10^{-11} \left(\frac{1 \text{ PeV}}{f_a} \right)^2 \left(\frac{\mathcal{K}_t}{10} \right)^2 \lambda_{K\pi\hat{\pi}}^{1/2}$$



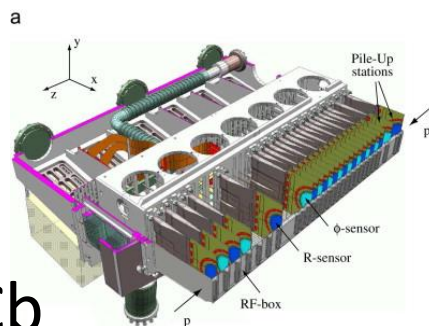
LLP Phenomenology

Dark Shower

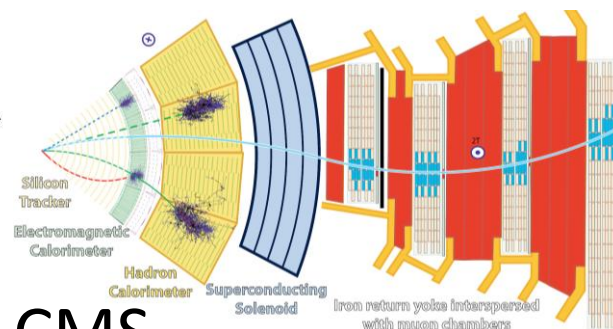


Signal: Muon rich jet-like structure with long-lived tracks + MET

LHCb

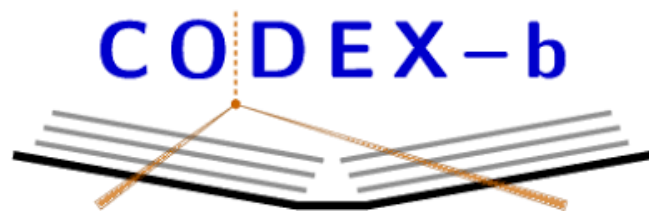


CMS

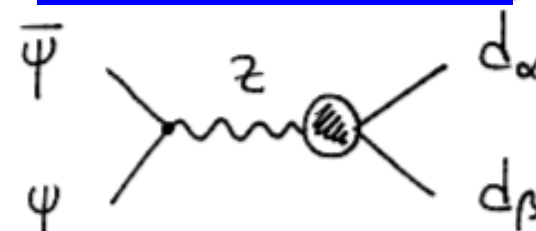


Searching for displaced dimuon resonances (one or more)

LHCb, 2007.03923; CMS, 2112.13769

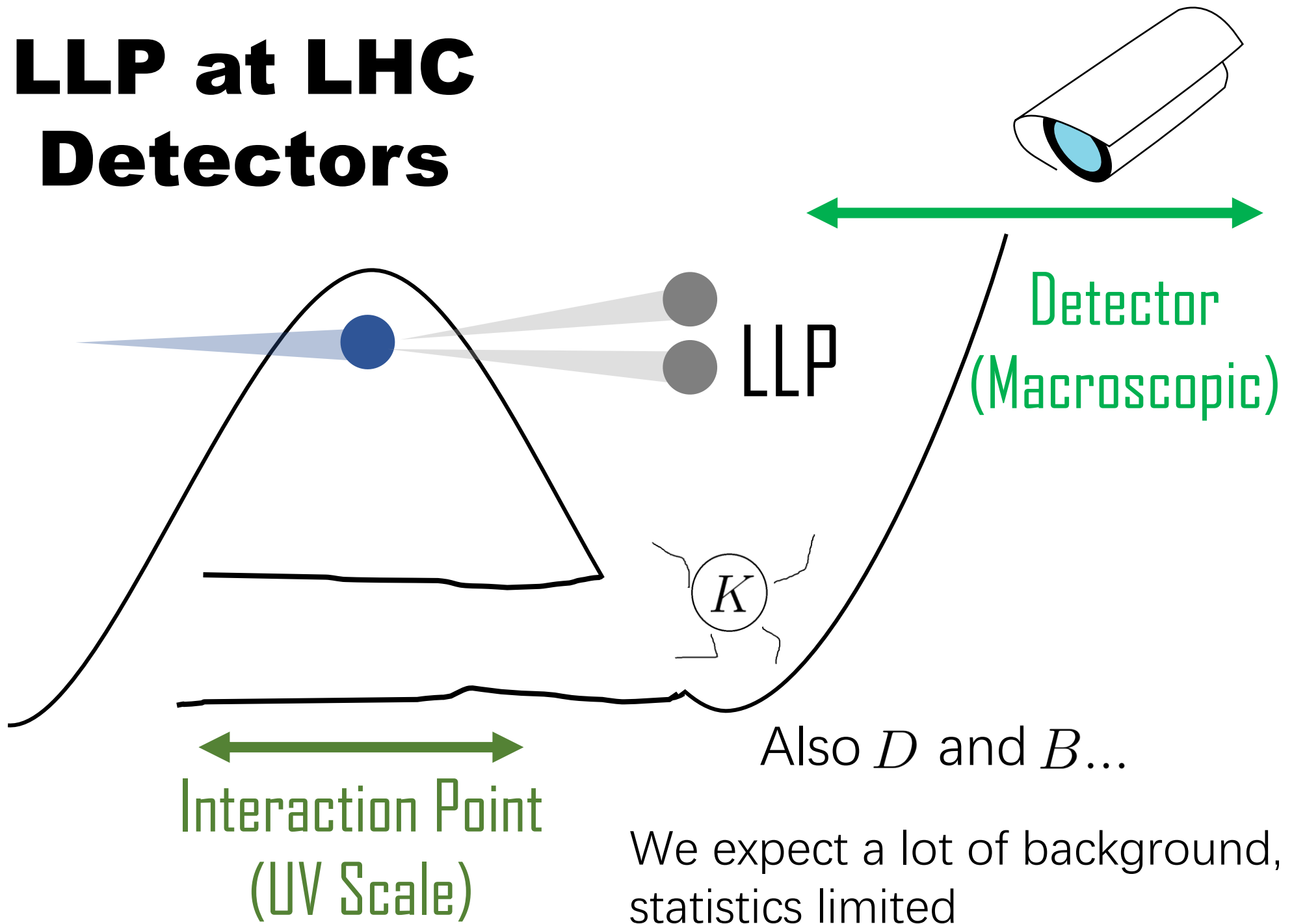


$$\text{BR} \approx 10^{-8} \left(\frac{1 \text{ PeV}}{f_a} \right)^2$$

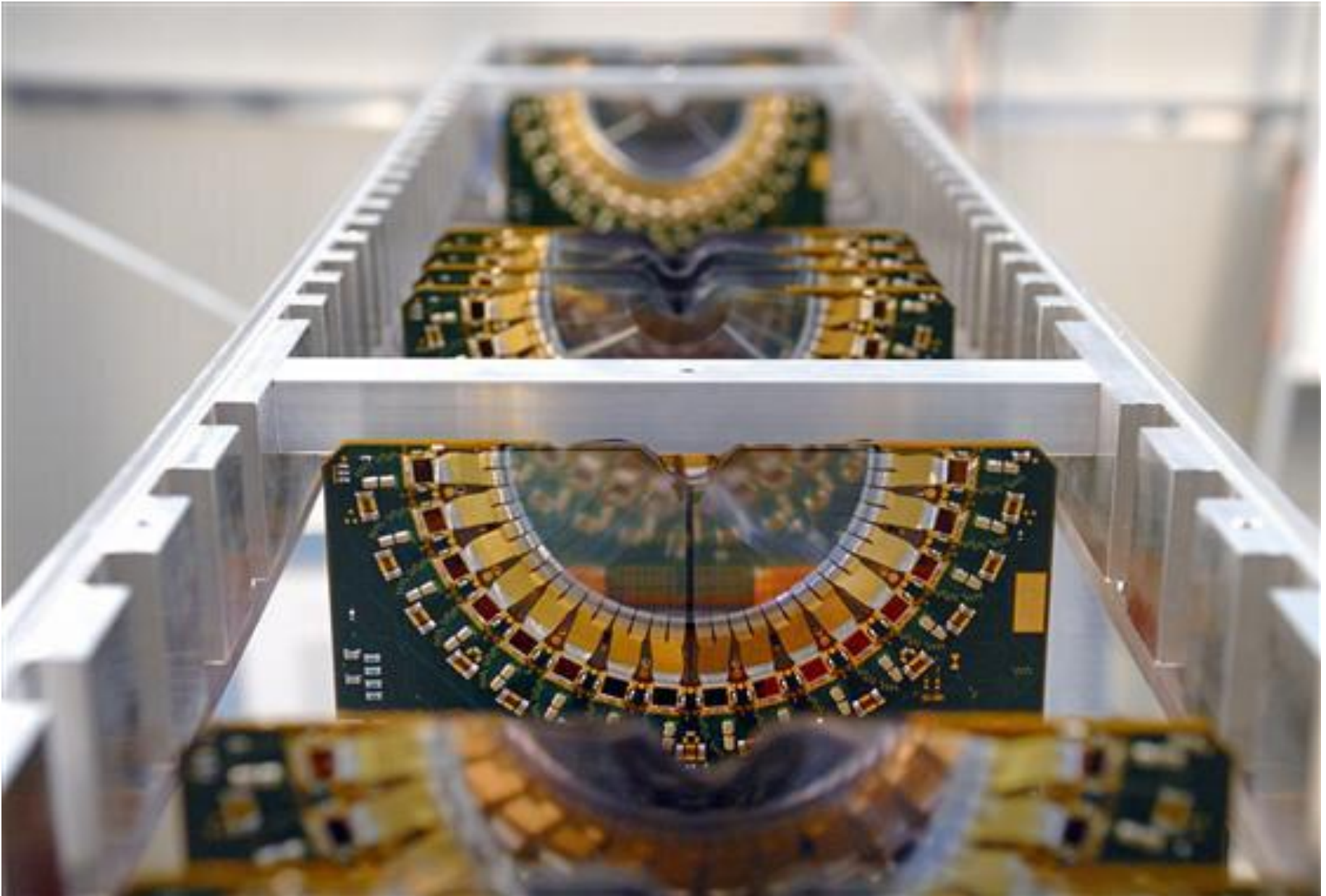


FCNC Decays

LLP at LHC Detectors

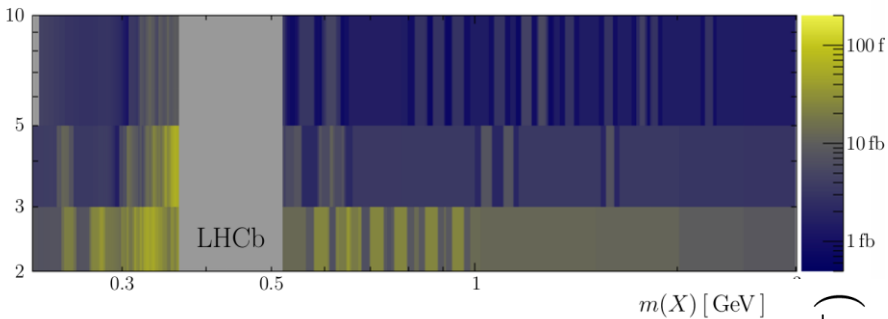


LHCb Vertex Locator (VELO)



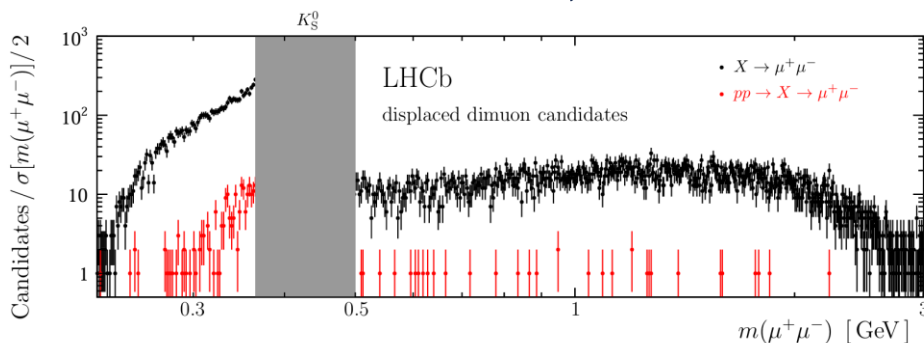
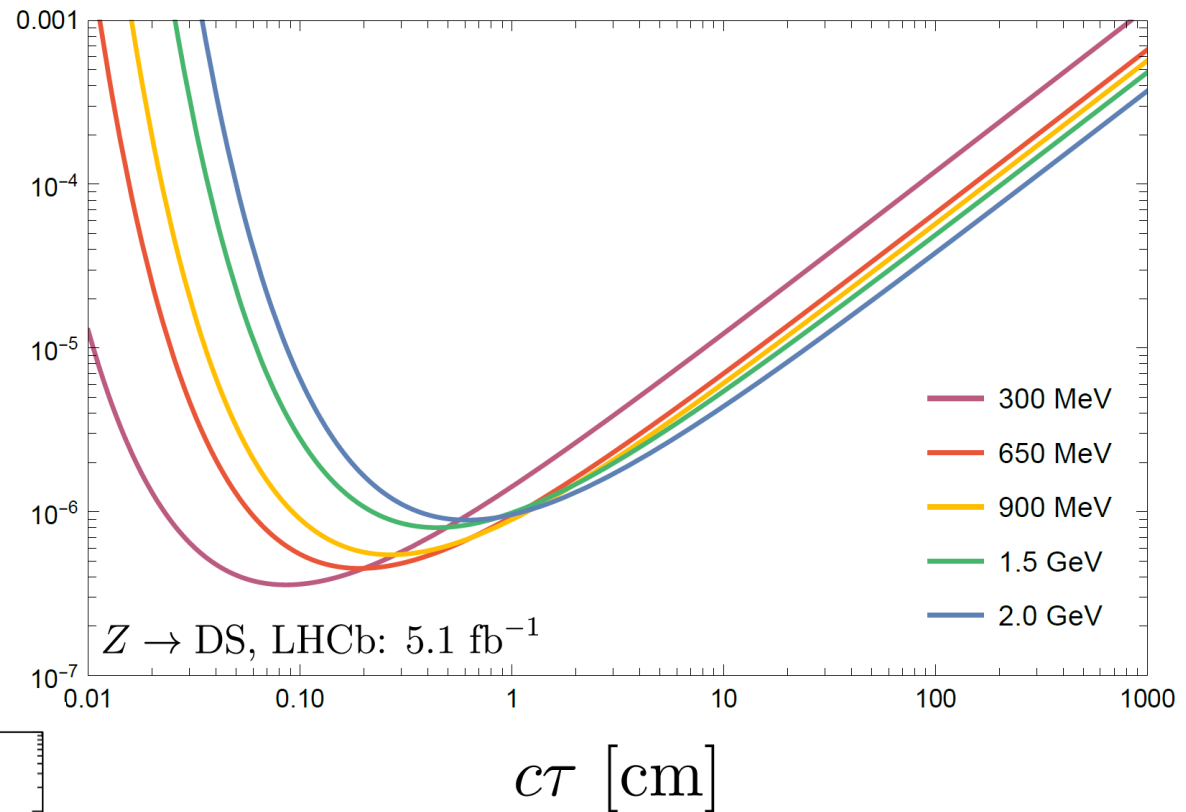
LHCb: “Trigger-free” Searches

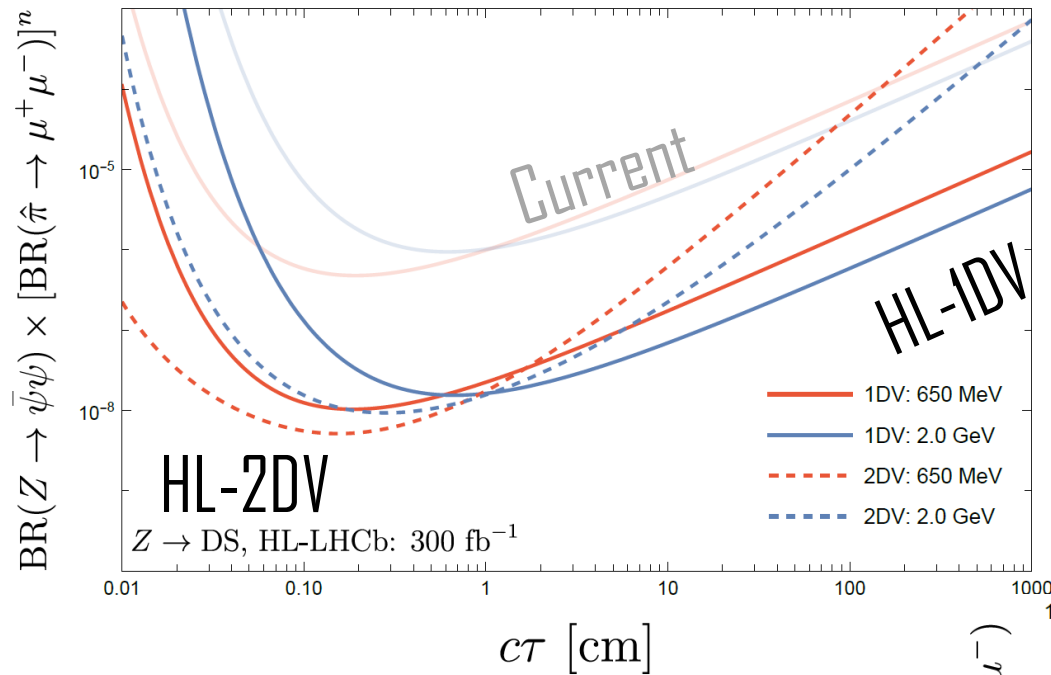
See [LHCb, 2007.03923](#)



Matching with multiple ways to calibrate the projection

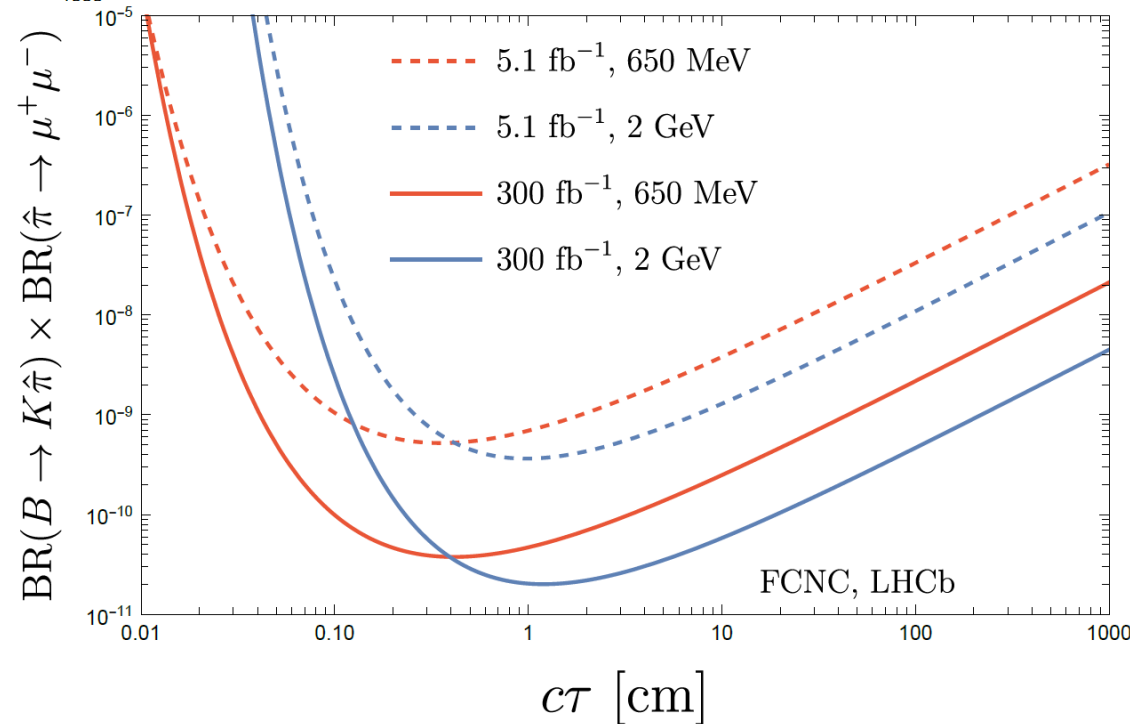
$$\text{BR}(Z \rightarrow \bar{\psi}\psi) \times \text{BR}(\hat{\pi} \rightarrow \mu^+ \mu^-)$$



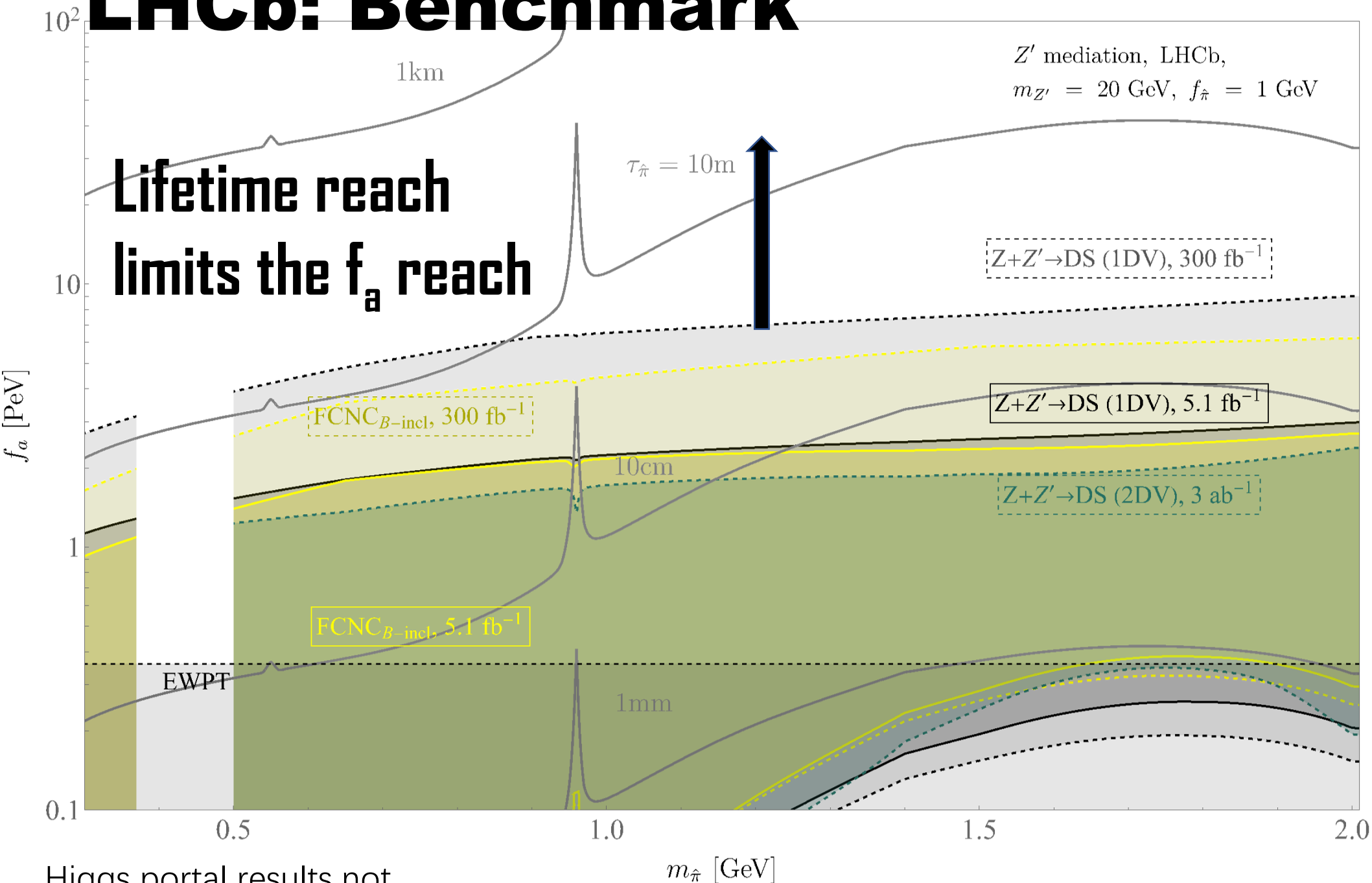


➤ Extended to the HL-LHC era, improve by $\sim \text{O}(10^2)$

➤ Also sensitive to FCNC, unsuppressed B production in the forward region



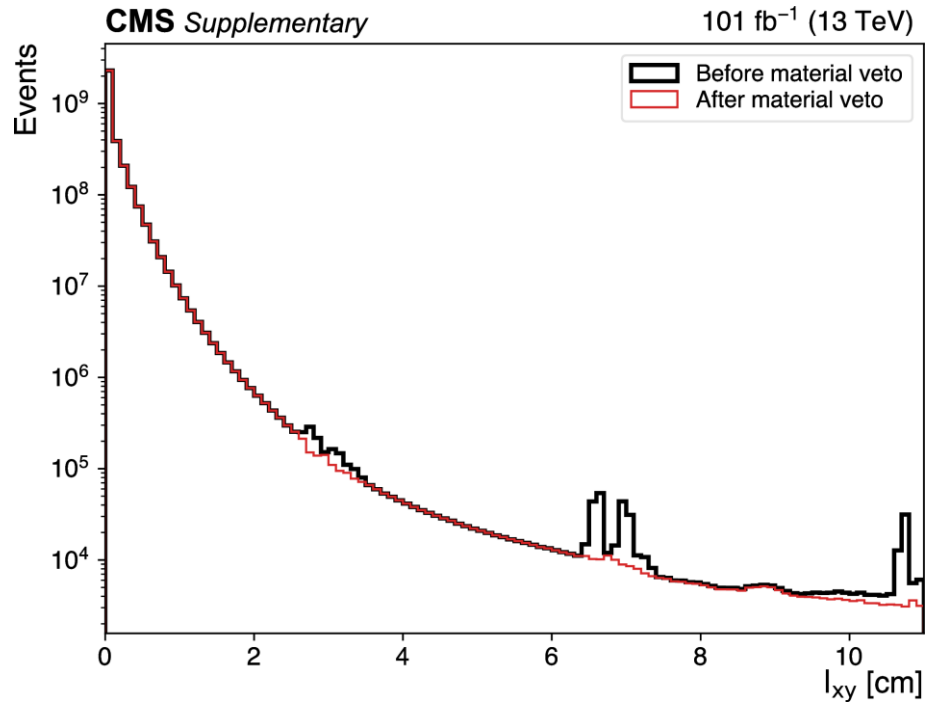
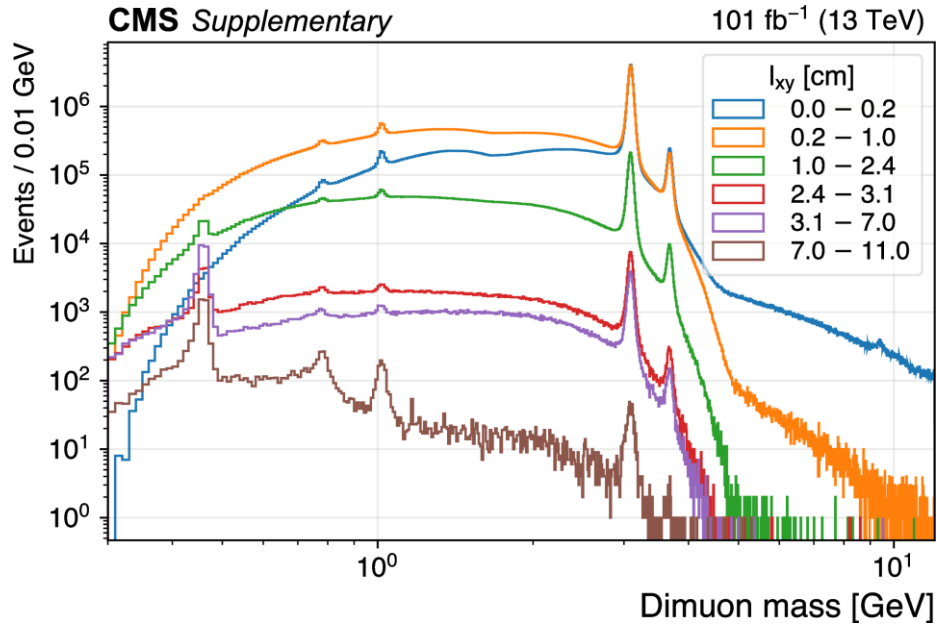
LHCb: Benchmark



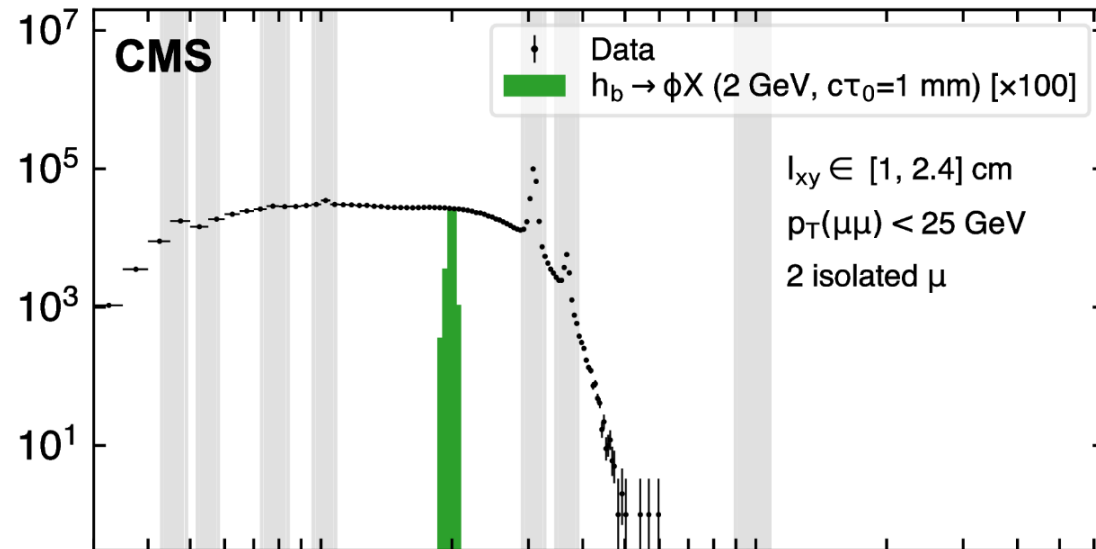
Higgs portal results not shown here

CMS: Dimuon Scouting

CMS, 2112.13769



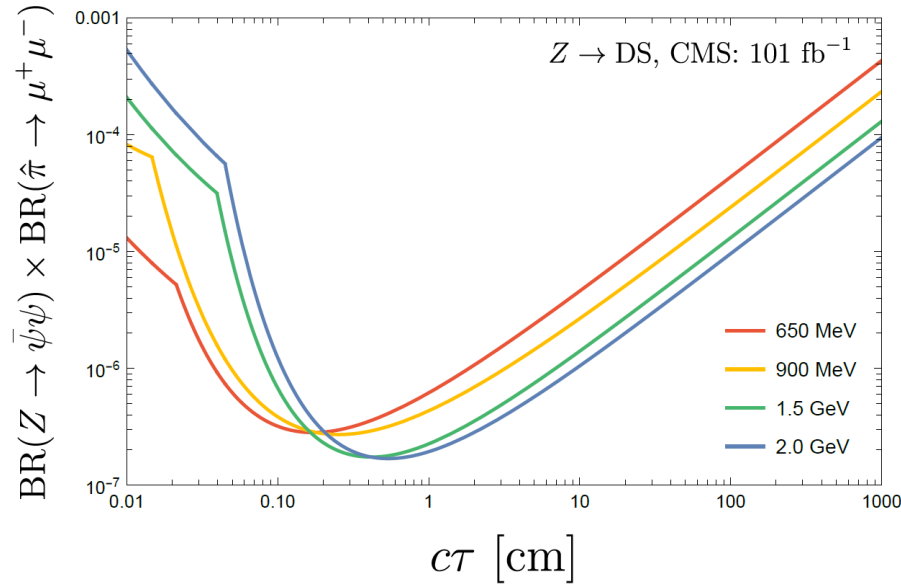
- Drop-out detailed info to boost dimuon trigger rate, useful for low mass resonances
- Not bkg free



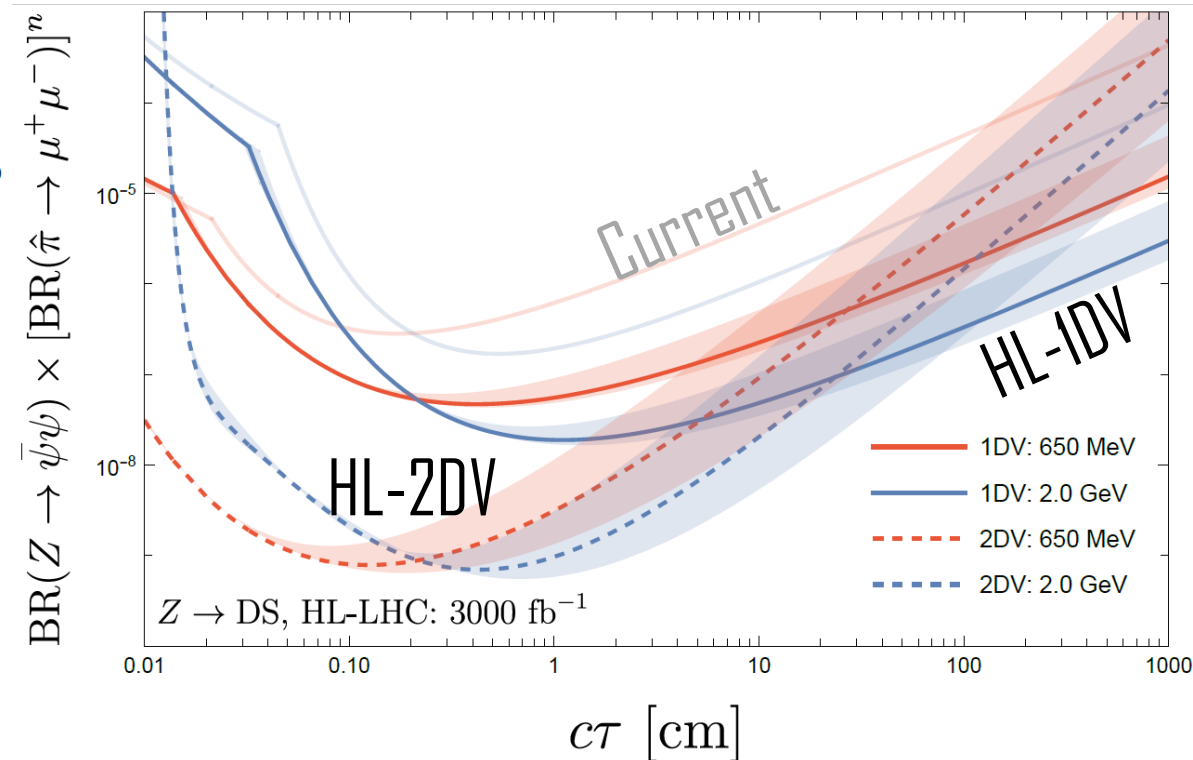
- Applies to Dark shower or FCNC

CMS: Model-Independent Limits

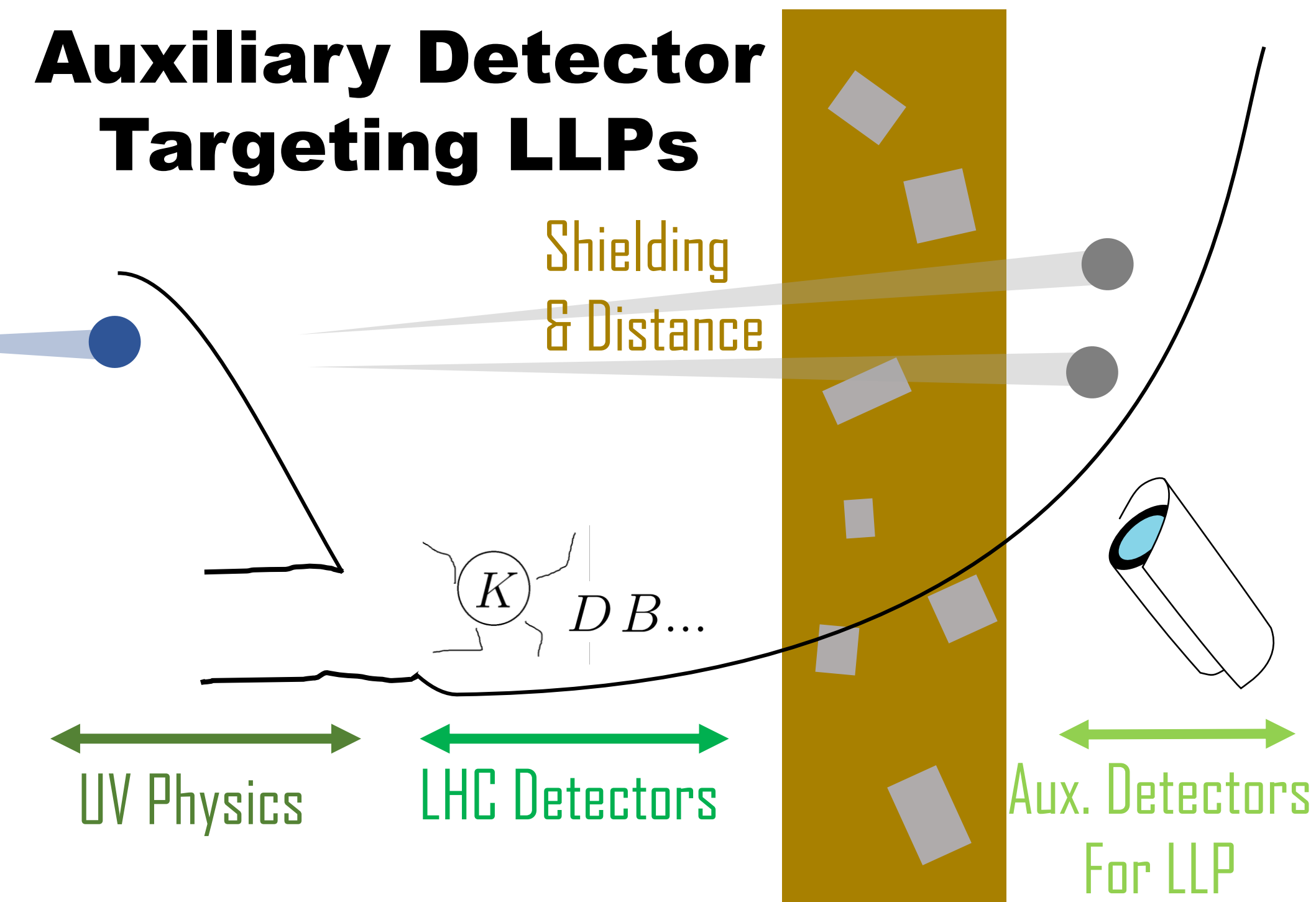
- Current limit with <11 cm transverse displacement
- 30 times luminosity, detector/algorithm upgrades.....
- ↘ HL-LHC limits



- At HL-LHC, high efficiency to do 2DV search: two dimuon DV with the same mass

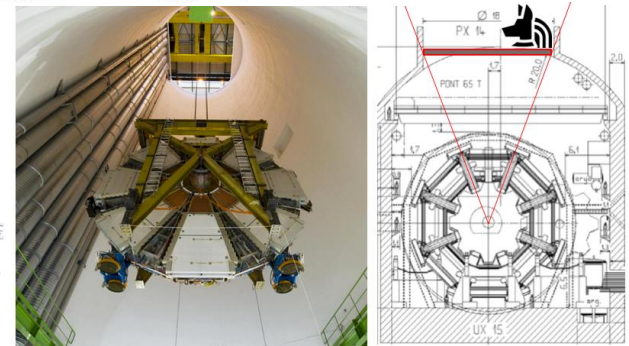
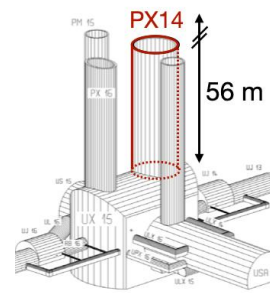
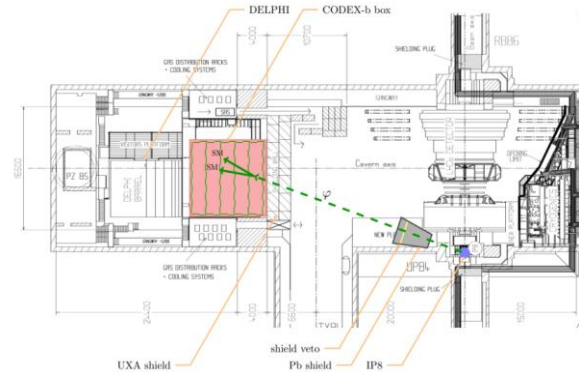
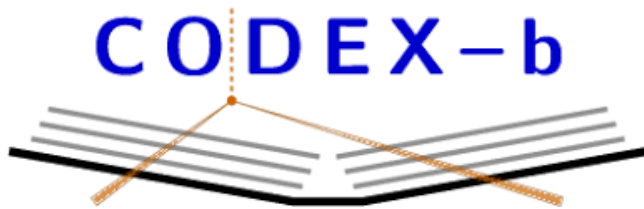
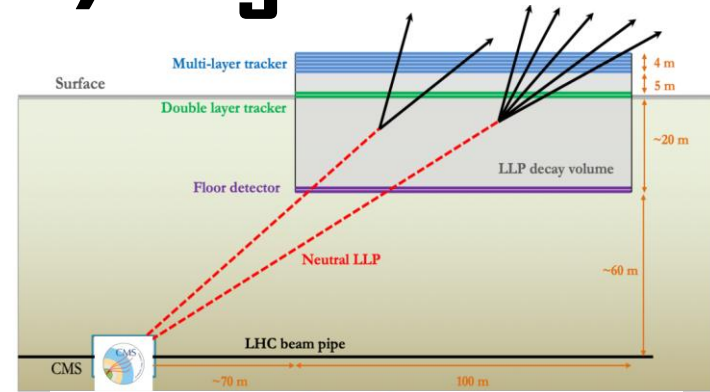


Auxiliary Detector Targeting LLPs

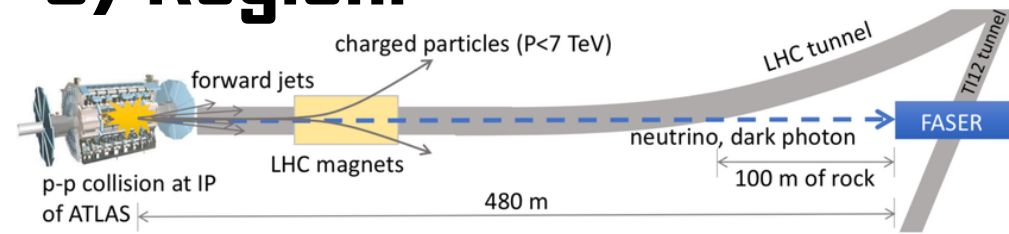


Central ($|\eta| < 5$) Region:

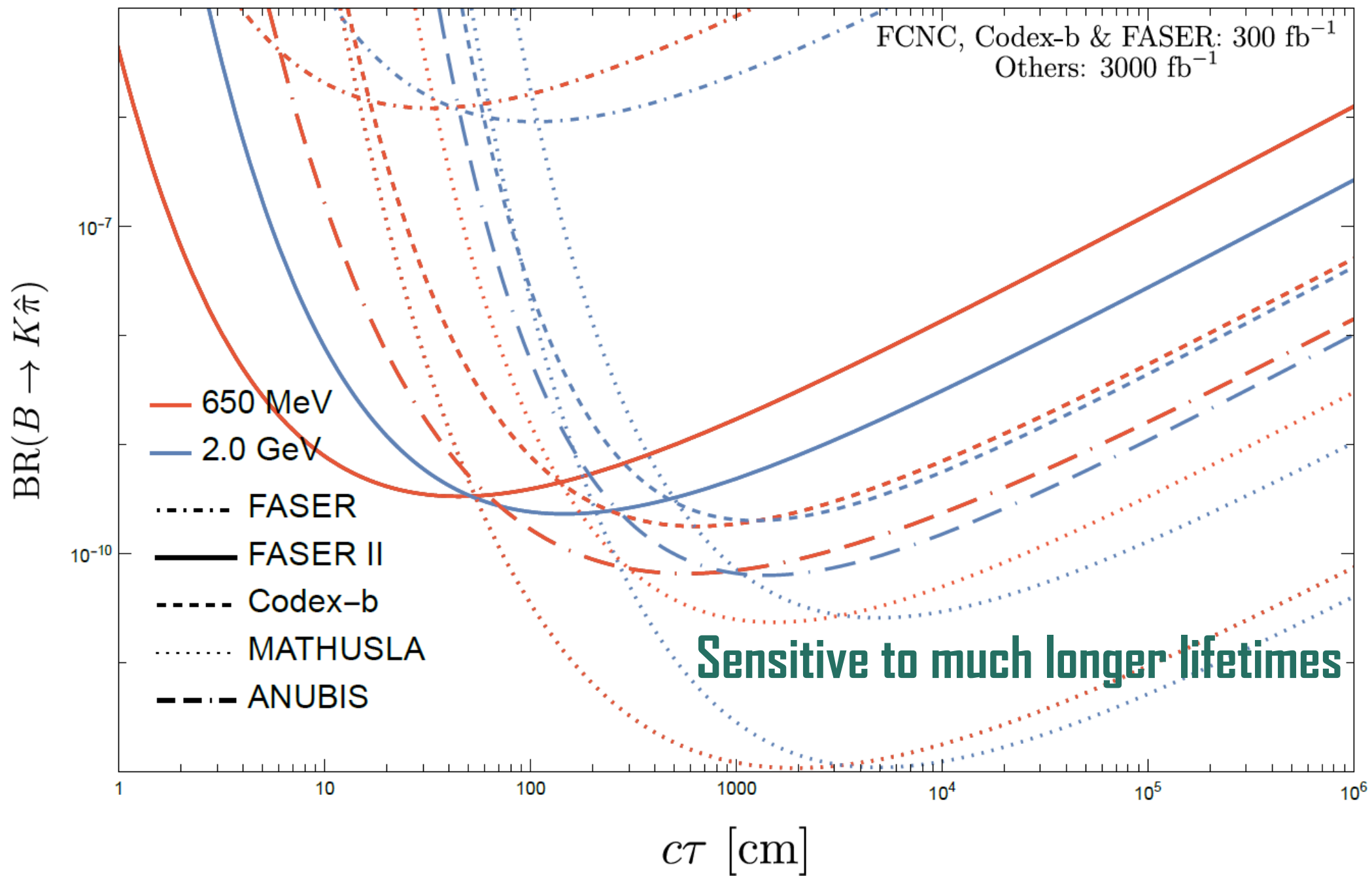
MATIS A

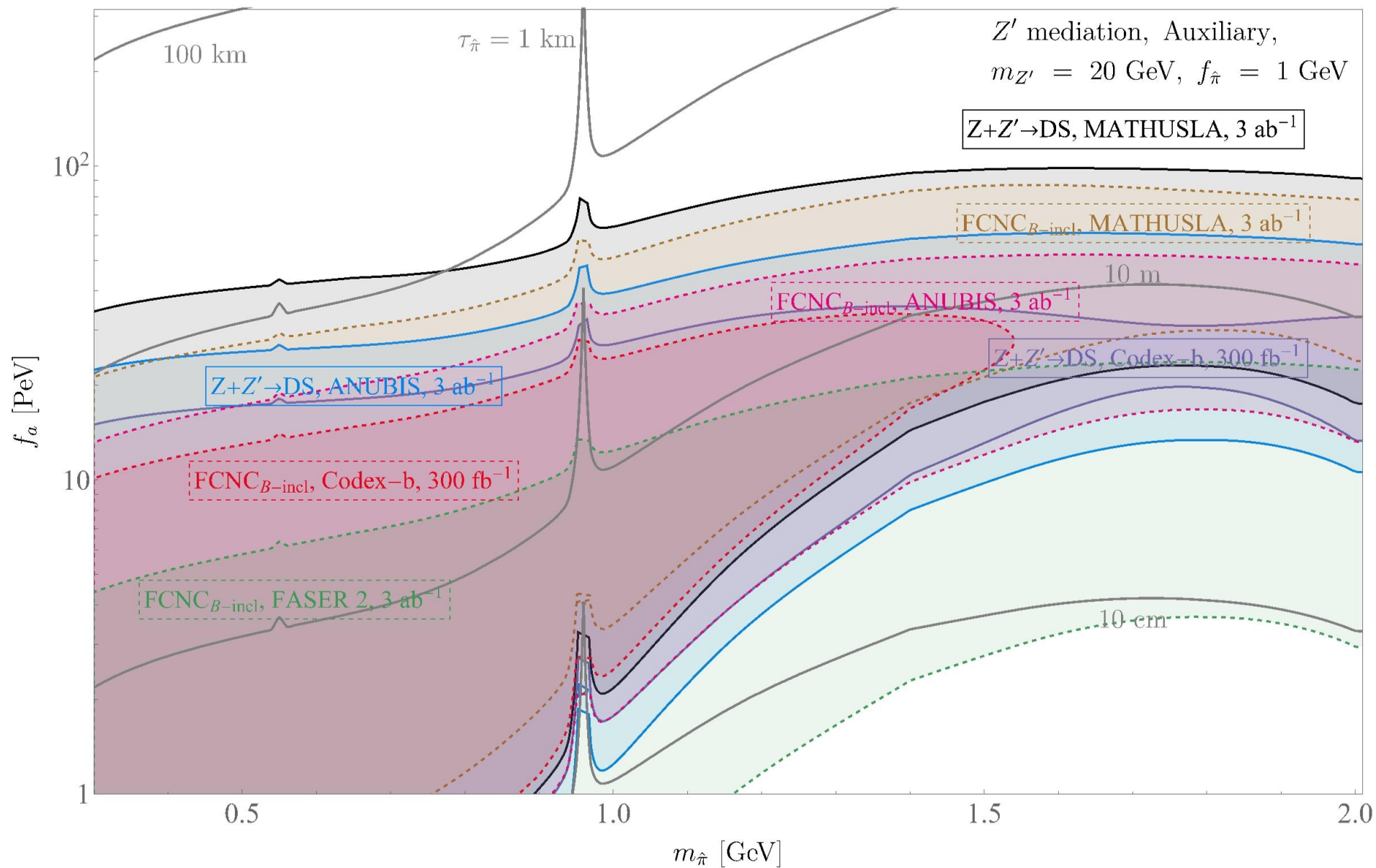


Forward ($|\eta| > 5$) Region:



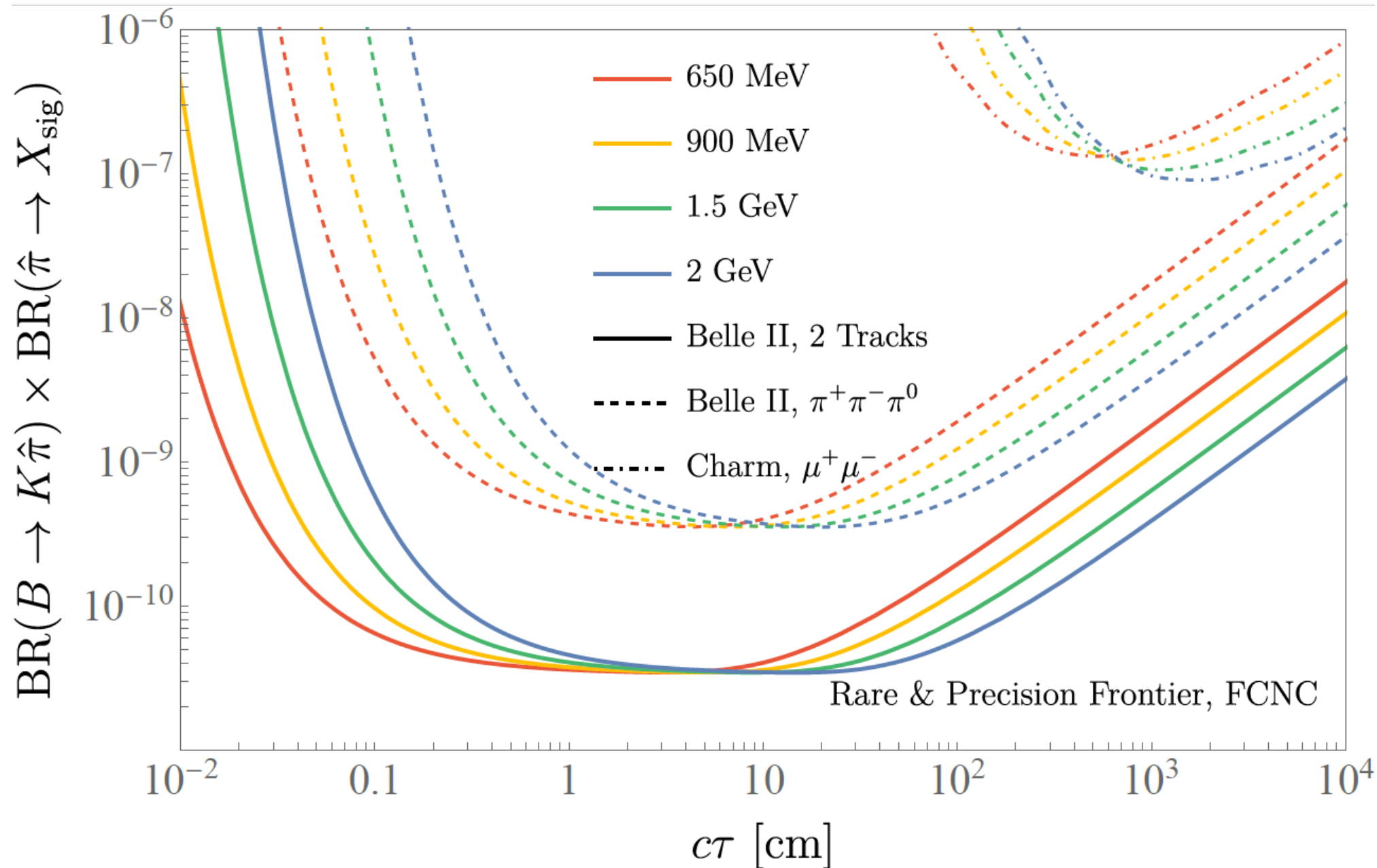
Auxiliary LLP Detectors

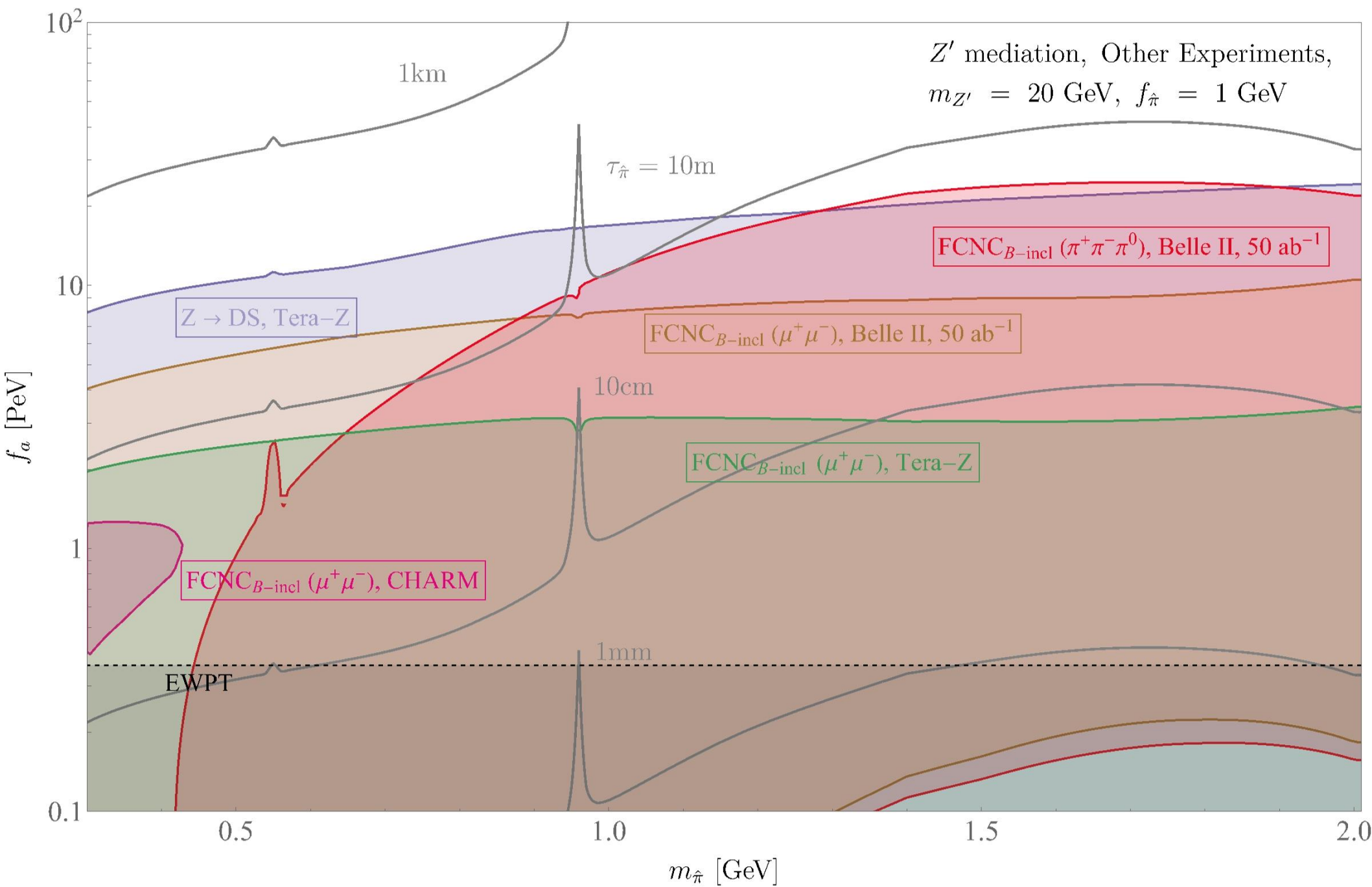




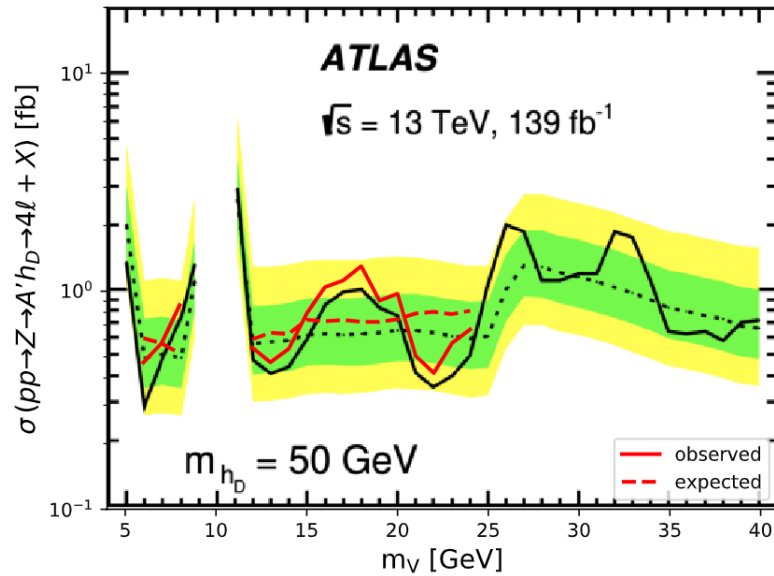
Corresponding f_a range increase up to $O(100)$ PeV

Rare/Precision/Intensity Frontier



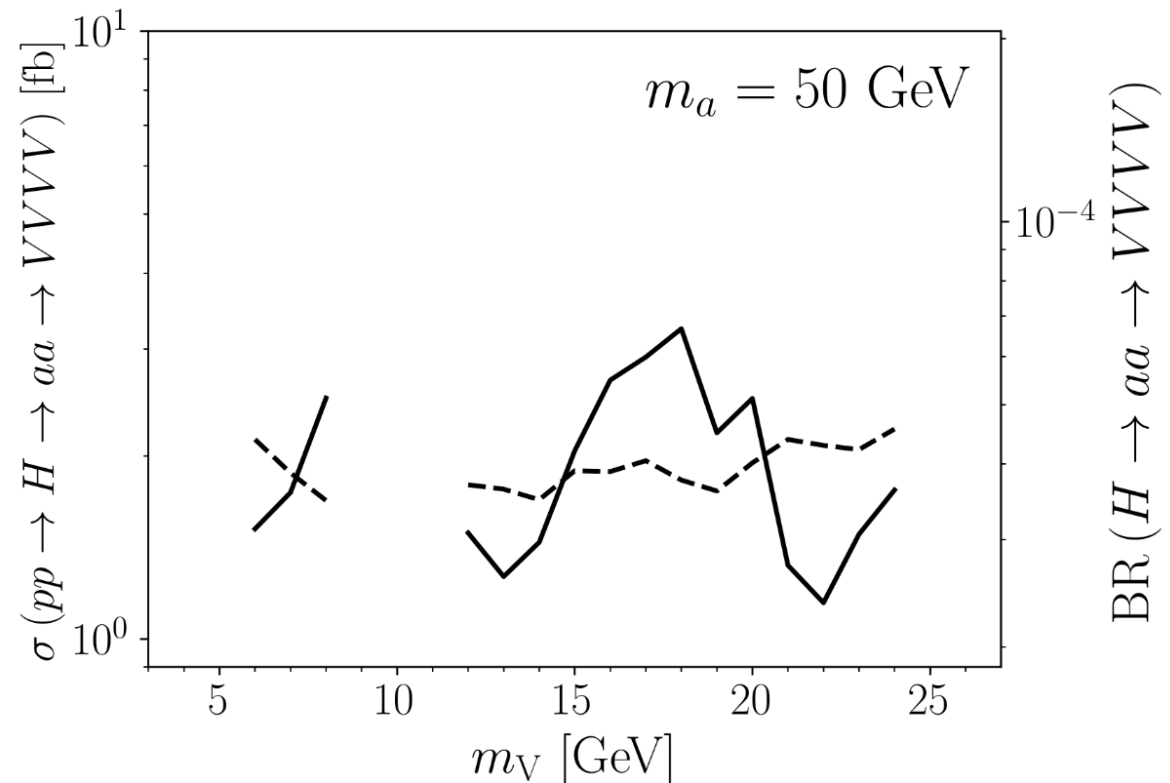
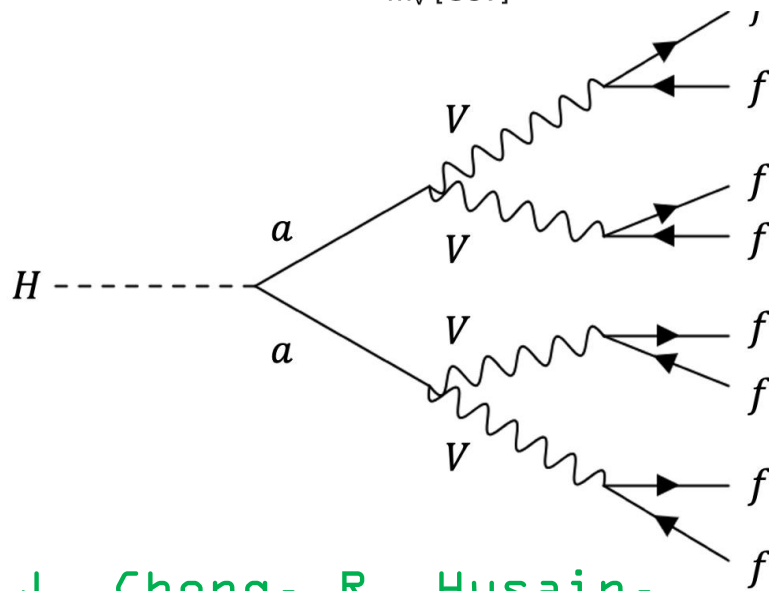


Prompt Dark Vector from DS



- Exotic Higgs or Z decays
- Recast from ATLAS $Z \rightarrow 4(6)$ lepton exotic searches

ATLAS 2306.07413



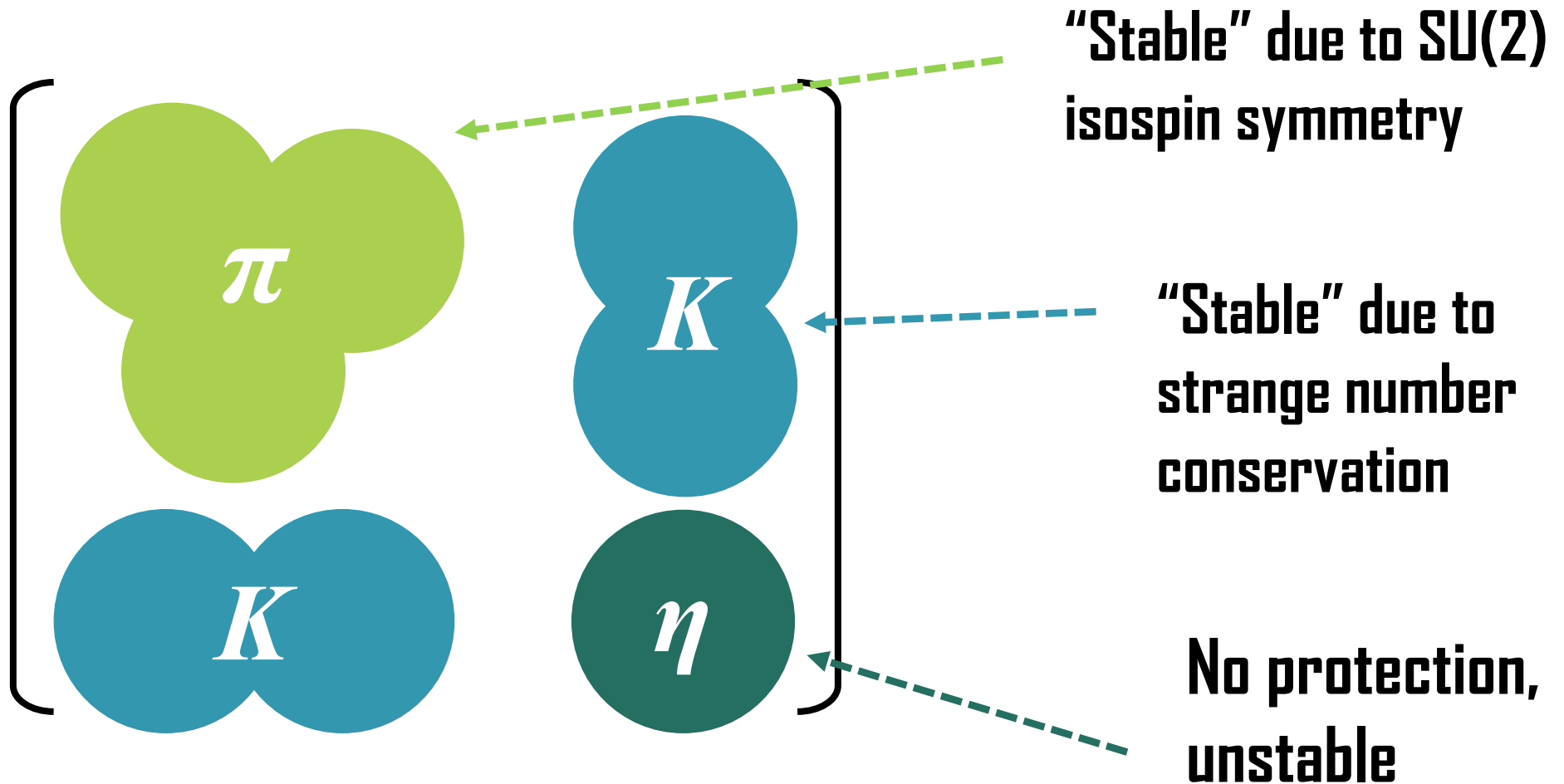
J. Cheng, R. Husain,
 LFL, M. Strassler,
 2412.14452 + in prep.

Time to Find out THE Dark but Colorful



GENERIC✓
RICHNESS✓

Back in the SM



Thermal Mechanisms



The messenger: heavy meson that decays, contact with the SM



The DM meson, non-interacting



The intermediate, exists in many models



SIMP-like



Pairwise



Scattering

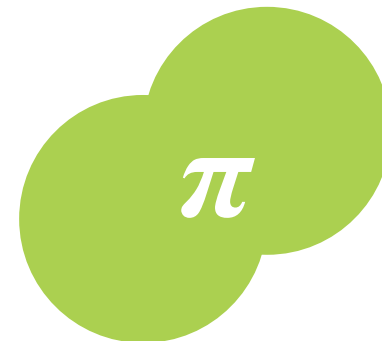
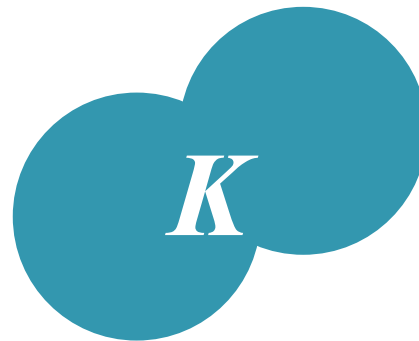
Y. Hochberg, E.
Kuflik, H.
Murayama, T.
Volansky, J.
Wacker, 1411.3727,
A. Katz, E.
Salvioni, and B.
Shakya,
2006.15148.....

The Model

Gauge group: $SO(N_c)$ $N_c \geq 4$

$SU(3)/SO(3) \Rightarrow$ 5 Goldstone bosons

Two π , two K and one η



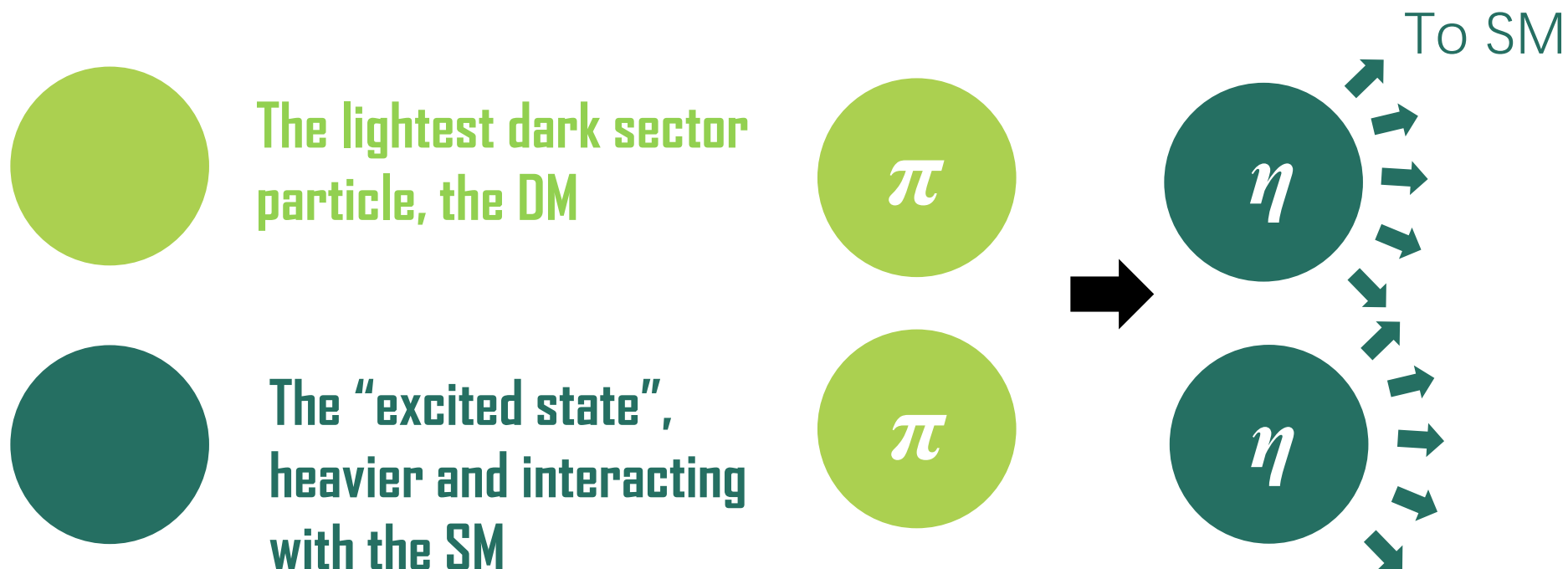
$$m_\eta = (1 + \Delta_\eta) m_\pi > m_K = (1 + \Delta_K) m_\pi > m_\pi$$

Leading ChPT Lagrangian (Low Temp)

$$\begin{aligned}
 \mathcal{L}_2 \supset & \frac{1}{24f^2} \left\{ m_{\hat{\pi}}^2 \left[\hat{\eta}_0^2 + 2\hat{\pi}_+\hat{\pi}_- + 2\hat{K}_{+1/2}\hat{K}_{-1/2} \right]^2 + 2(m_{\hat{K}}^2 - m_{\hat{\pi}}^2) \left[\frac{8}{9}\hat{\eta}_0^4 + 3\hat{\eta}_0^2\hat{K}_{+1/2}\hat{K}_{-1/2} \right. \right. \\
 & \left. \left. - \sqrt{\frac{2}{3}}(\hat{\eta}_0\hat{K}_{+1/2}^2\hat{\pi}_- + \hat{\eta}_0\hat{K}_{-1/2}^2\hat{\pi}_+) + 2\hat{K}_{+1/2}^2\hat{K}_{-1/2}^2 + 2\hat{K}_{+1/2}\hat{K}_{-1/2}\hat{\pi}_+\hat{\pi}_- \right] \right\} \\
 & + \frac{1}{24f^2} \left\{ 4(\partial\hat{\pi}_+)^2\hat{\pi}_-^2 + 4\hat{\pi}_+^2(\partial\hat{\pi}_-)^2 - 8\partial\hat{\pi}_+\hat{\pi}_+\partial\hat{\pi}_-\hat{\pi}_- + (\partial\hat{K}_{+1/2})^2\hat{K}_{-1/2}^2 + \hat{K}_{+1/2}^2(\partial\hat{K}_{-1/2})^2 \right. \\
 & - 2\partial\hat{K}_{+1/2}\hat{K}_{+1/2}\partial\hat{K}_{-1/2}\hat{K}_{-1/2} + 8\hat{K}_{+1/2}\partial\hat{K}_{-1/2}\hat{\pi}_+\partial\hat{\pi}_- + 8\partial\hat{K}_{+1/2}\hat{K}_{-1/2}\partial\hat{\pi}_+\hat{\pi}_- \\
 & - 4\hat{K}_{+1/2}\hat{K}_{-1/2}\partial\hat{\pi}_+\partial\hat{\pi}_- - 4\hat{K}_{+1/2}\partial\hat{K}_{-1/2}\partial\hat{\pi}_+\hat{\pi}_- - 4\partial\hat{K}_{+1/2}\hat{K}_{-1/2}\hat{\pi}_+\partial\hat{\pi}_- \\
 & - 4\partial\hat{K}_{+1/2}\partial\hat{K}_{-1/2}\hat{\pi}_+\hat{\pi}_- - 6\hat{\eta}_0^2\partial\hat{K}_{+1/2}\partial\hat{K}_{-1/2} - 6(\partial\hat{\eta}_0)^2\hat{K}_{+1/2}\hat{K}_{-1/2} \\
 & + 6\hat{\eta}_0\partial\hat{\eta}_0\partial\hat{K}_{+1/2}\hat{K}_{-1/2} + 6\hat{\eta}_0\partial\hat{\eta}_0\hat{K}_{+1/2}\partial\hat{K}_{-1/2} + 2\sqrt{6} \left[\partial\hat{\eta}_0\partial\hat{K}_{-1/2}\hat{K}_{-1/2}\hat{\pi}_+ \right. \\
 & + \partial\hat{\eta}_0\partial\hat{K}_{+1/2}\hat{K}_{+1/2}\hat{\pi}_- - \partial\hat{\eta}_0\hat{K}_{+1/2}^2\partial\hat{\pi}_- - \partial\hat{\eta}_0\hat{K}_{-1/2}^2\partial\hat{\pi}_+ - \hat{\eta}_0(\partial\hat{K}_{-1/2})^2\hat{\pi}_+ \\
 & \left. \left. - \hat{\eta}_0(\partial\hat{K}_{+1/2})^2\hat{\pi}_- + \hat{\eta}_0\partial\hat{K}_{-1/2}\hat{K}_{-1/2}\partial\hat{\pi}_+ + \hat{\eta}_0\partial\hat{K}_{+1/2}\hat{K}_{+1/2}\partial\hat{\pi}_- \right] \right\}.
 \end{aligned}$$

**Mass & 4-point
interaction terms**

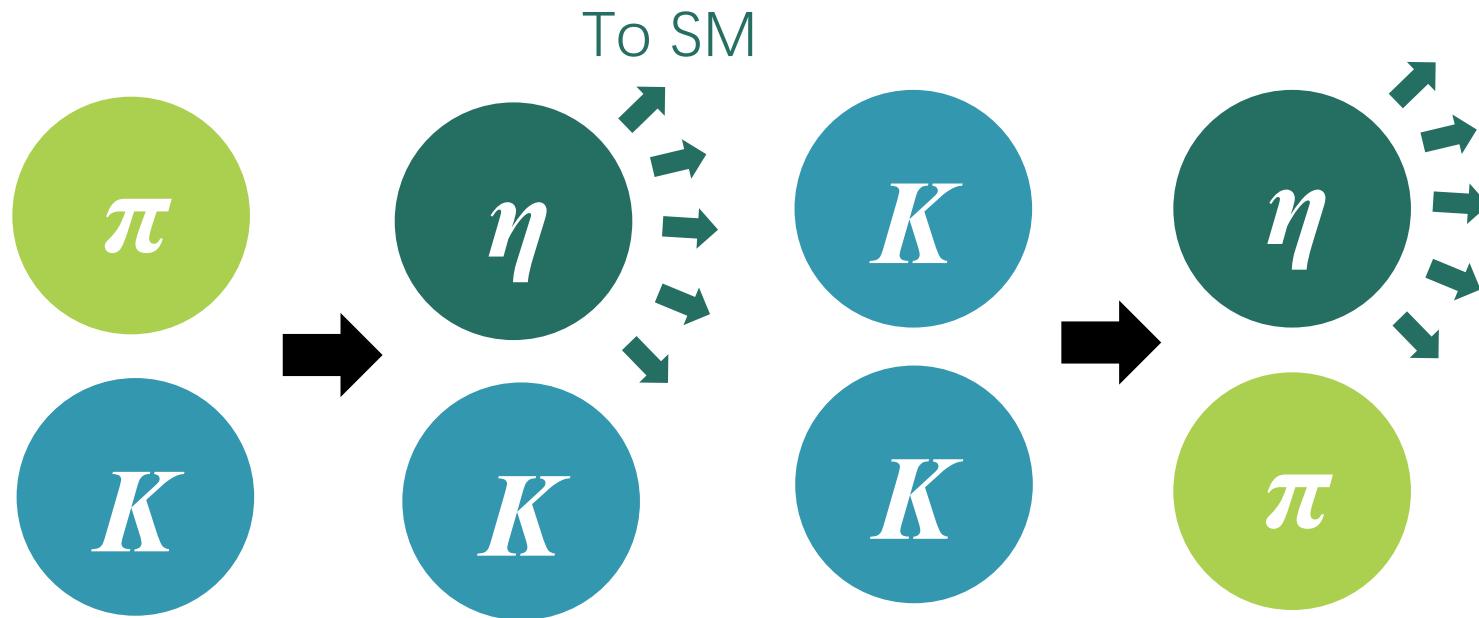
Forbidden Composite DM



**Only energetic particle can go through
the mass gap, exponential sensitivity**

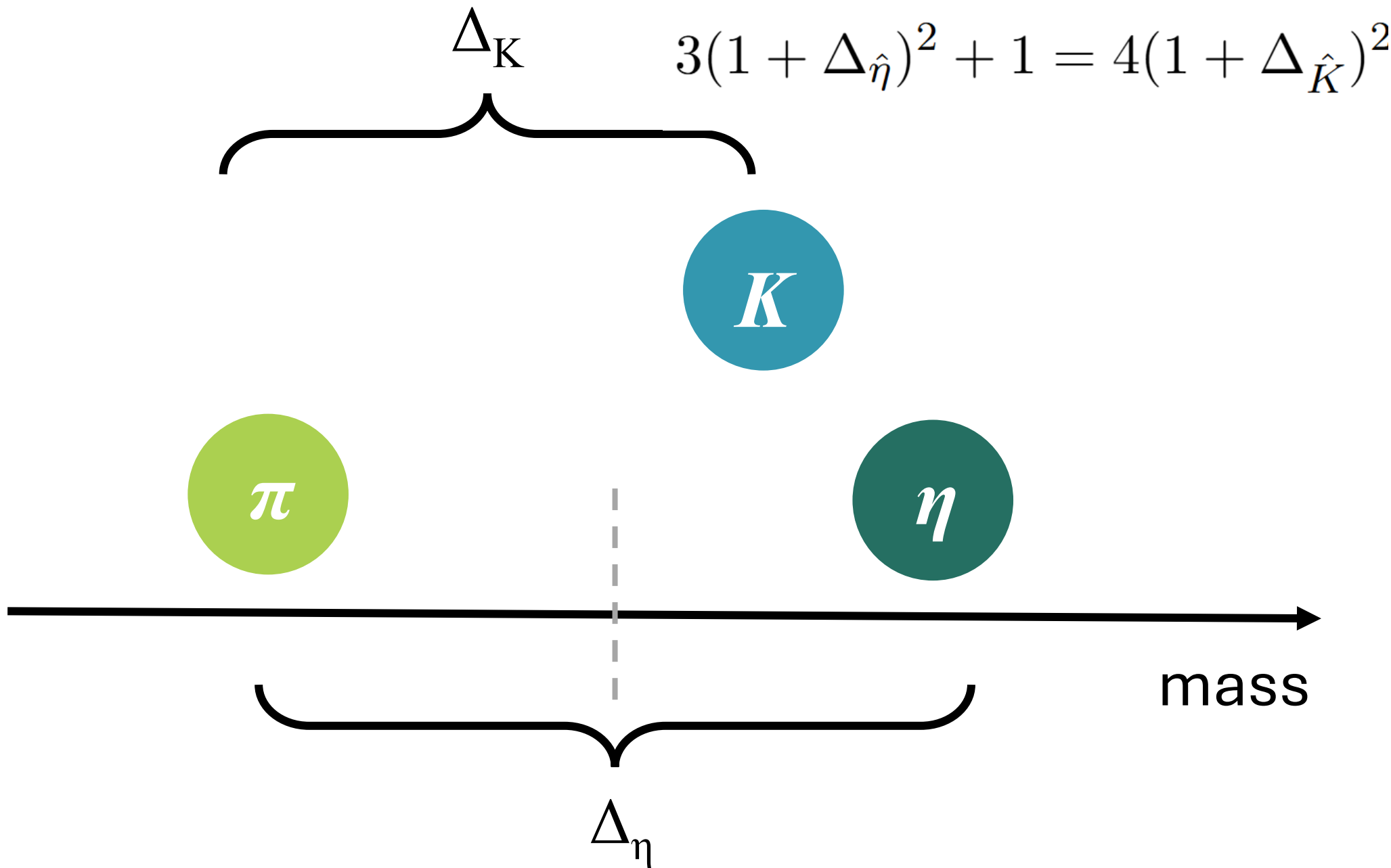
R.T. D'Agnolo and J.T. Ruderman, 1505.07107; LFL, Y. Tsai, 1901.0993; E. Bernreuther, F. Kahlhoefer, M. Kramer and P. Tunney, 1907.04346.....

Important Alternative Processes



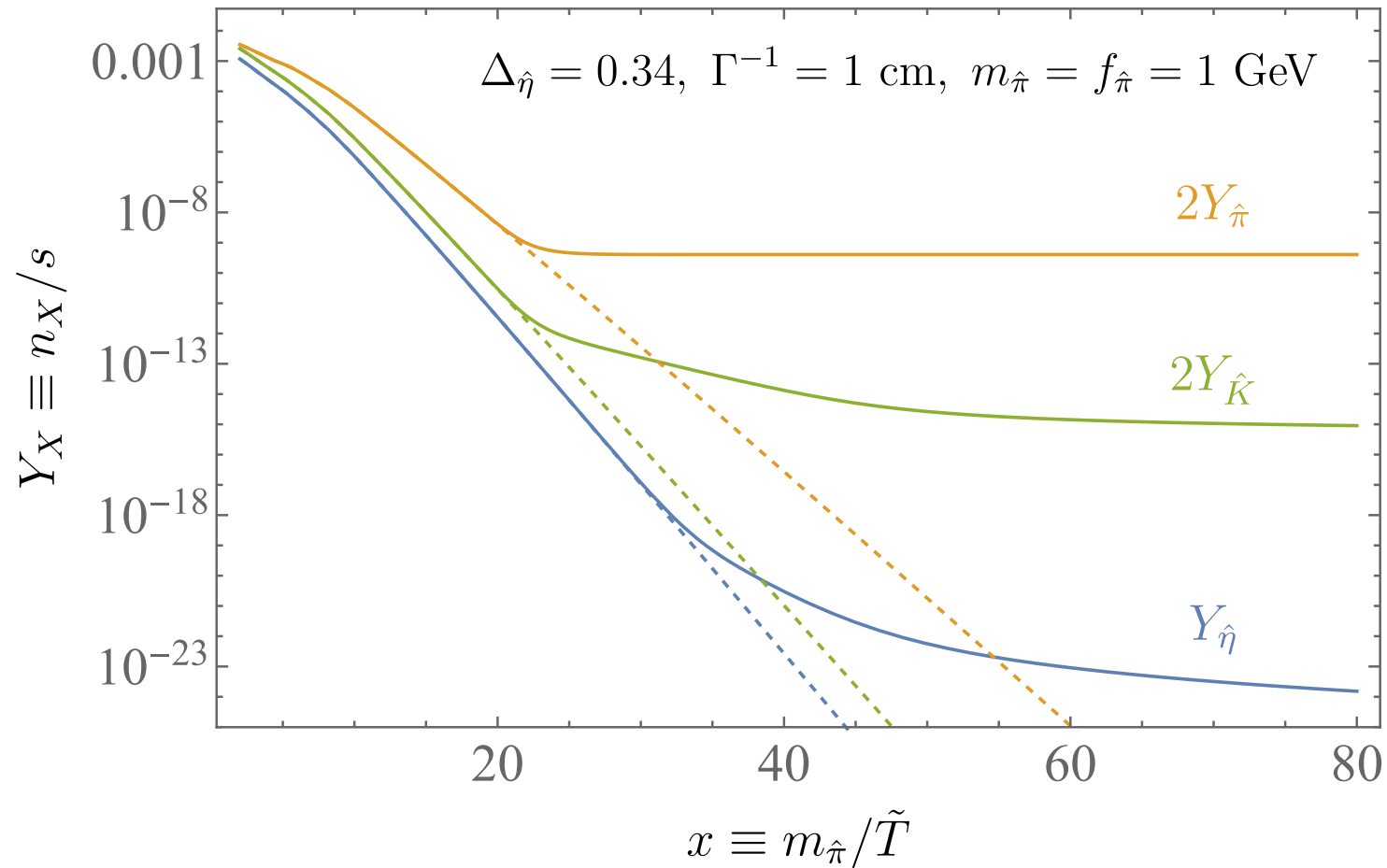
K assisted conversion, low Boltzmann suppression

K annihilation, gives indirect detection signals

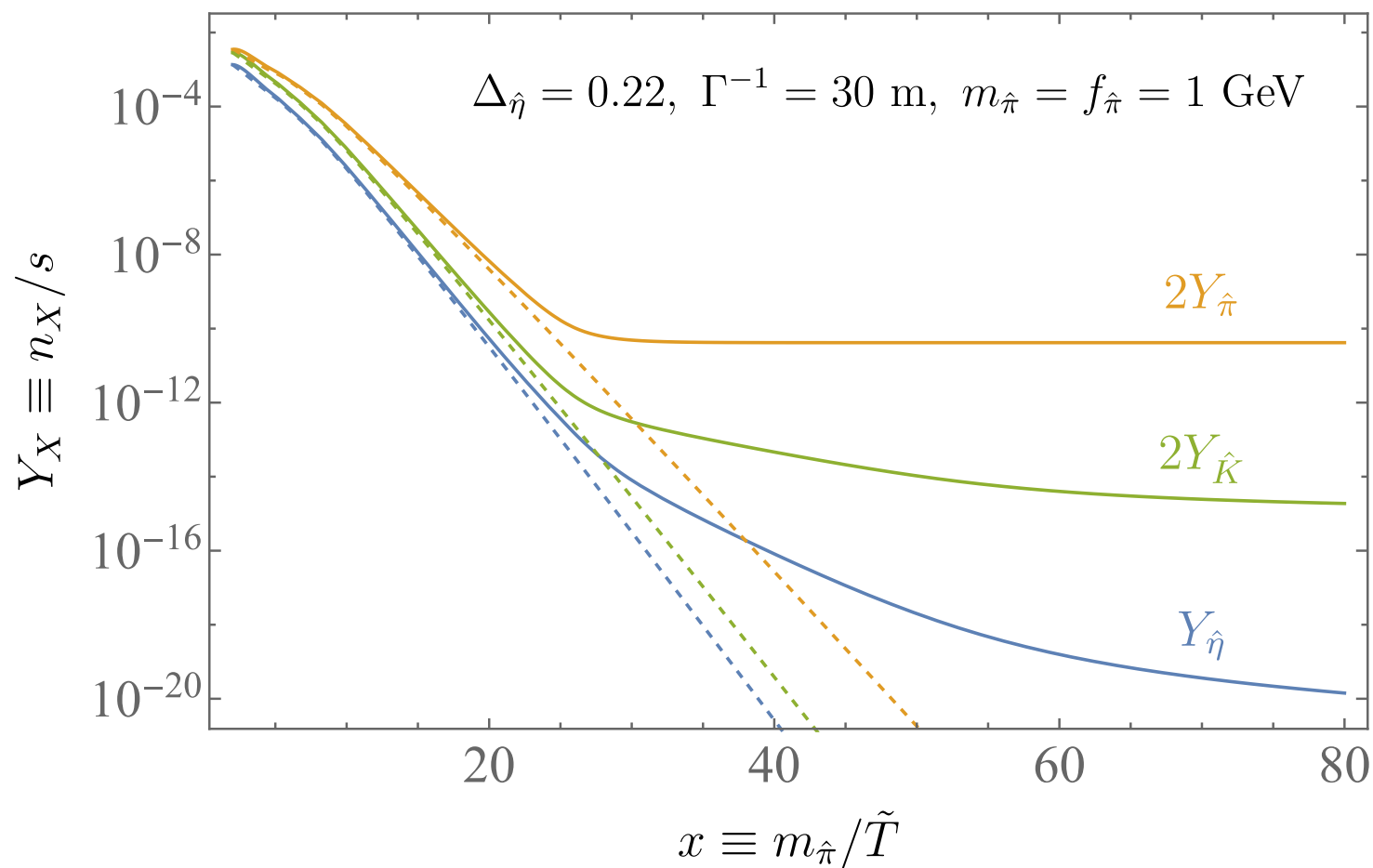


KK tp $\pi\eta$ always onshell and unsuppressed

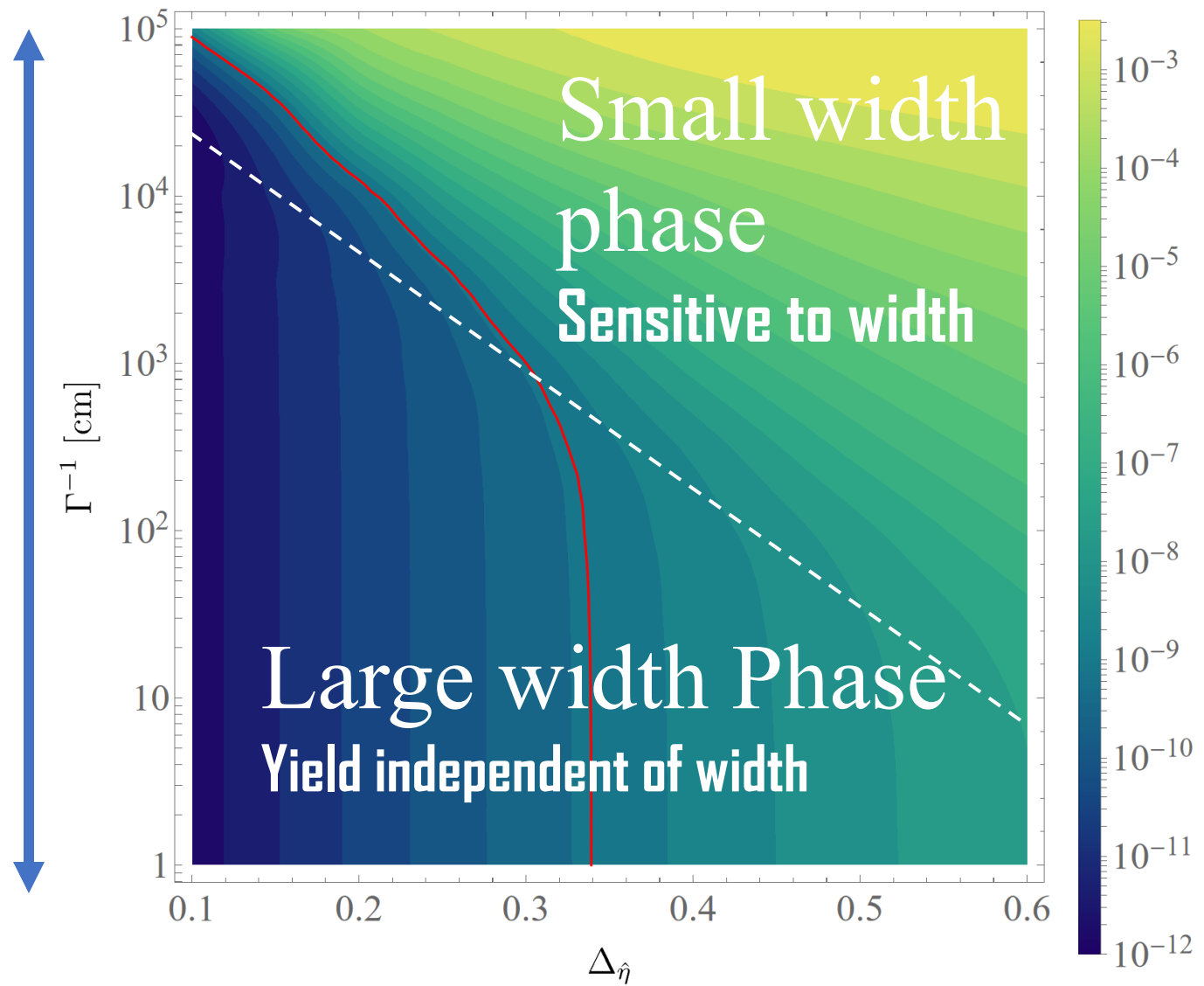
Large Decay Width Limit: Freezeout Controlled by Conversion

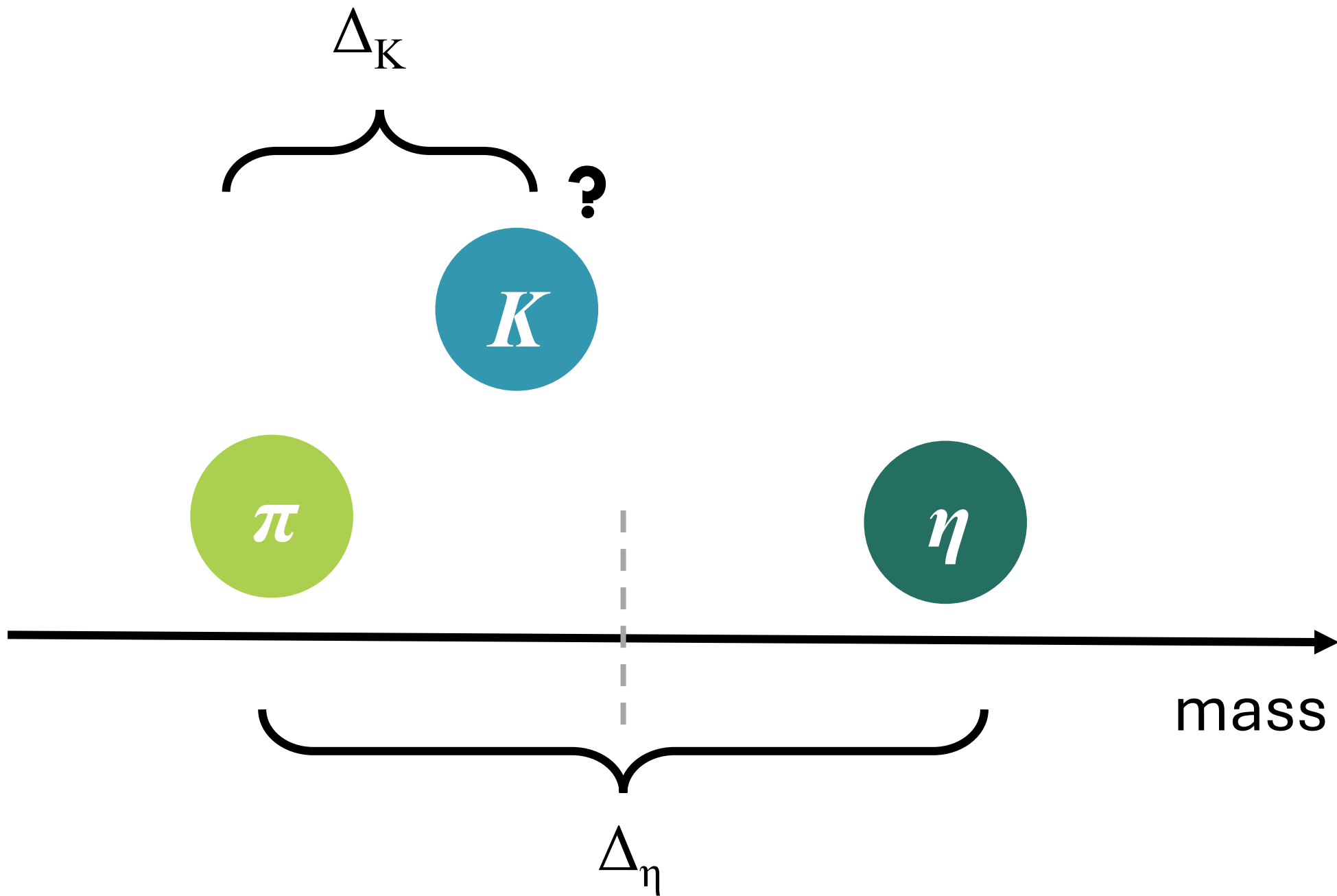


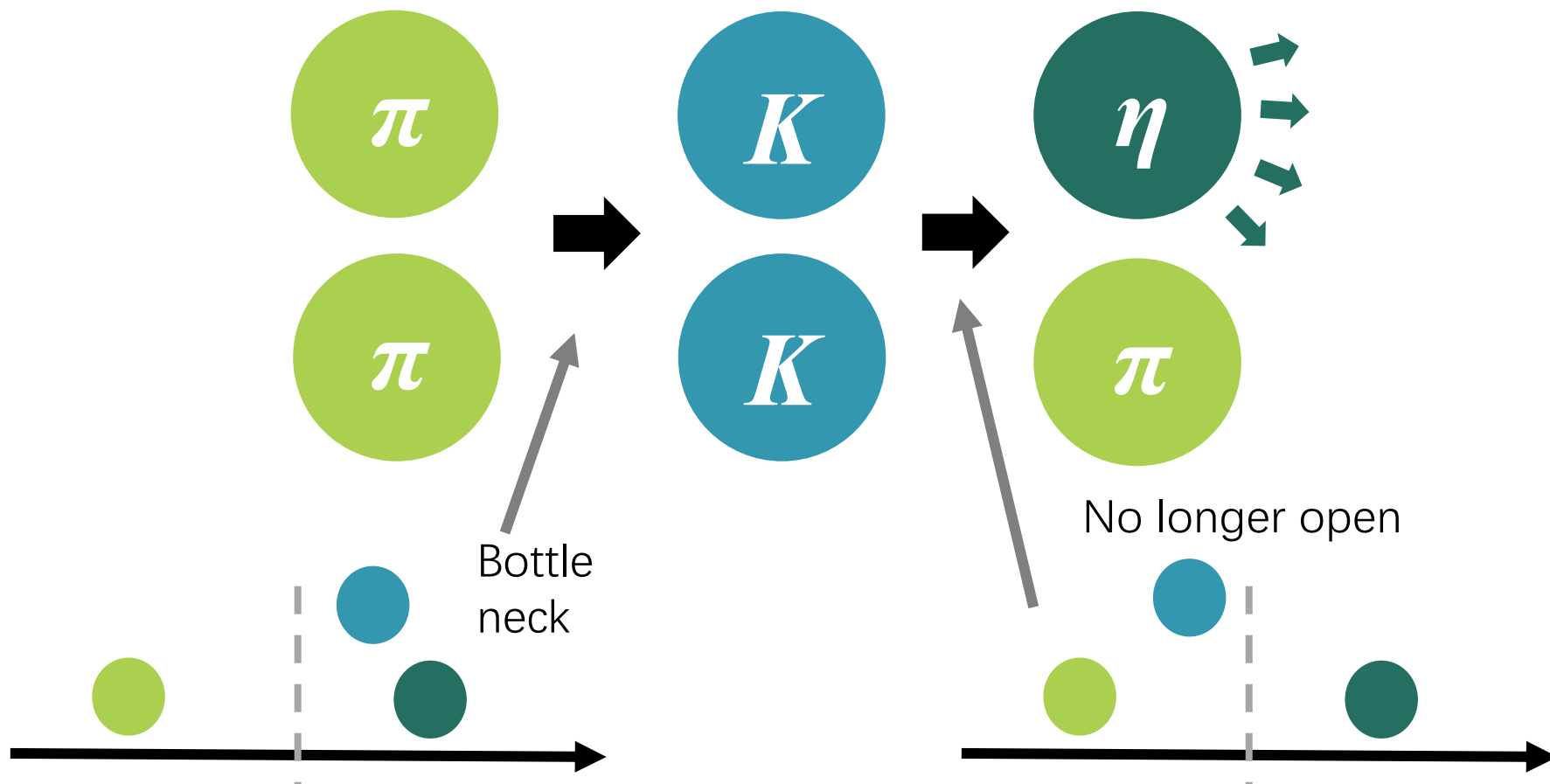
Small Decay Width Limit: Freezeout Controlled by the Decay



Decay at
Human/Lab
Length Scales

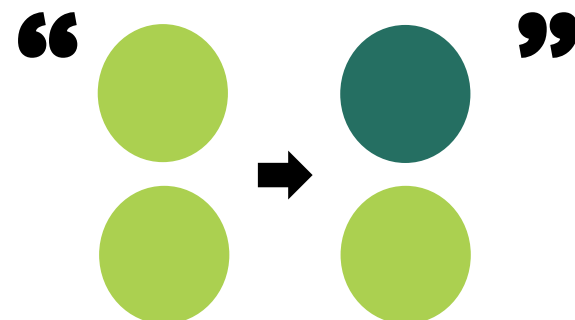
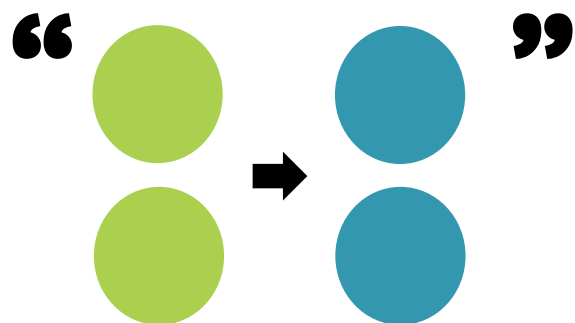






Effectively $\pi\pi$ to KK

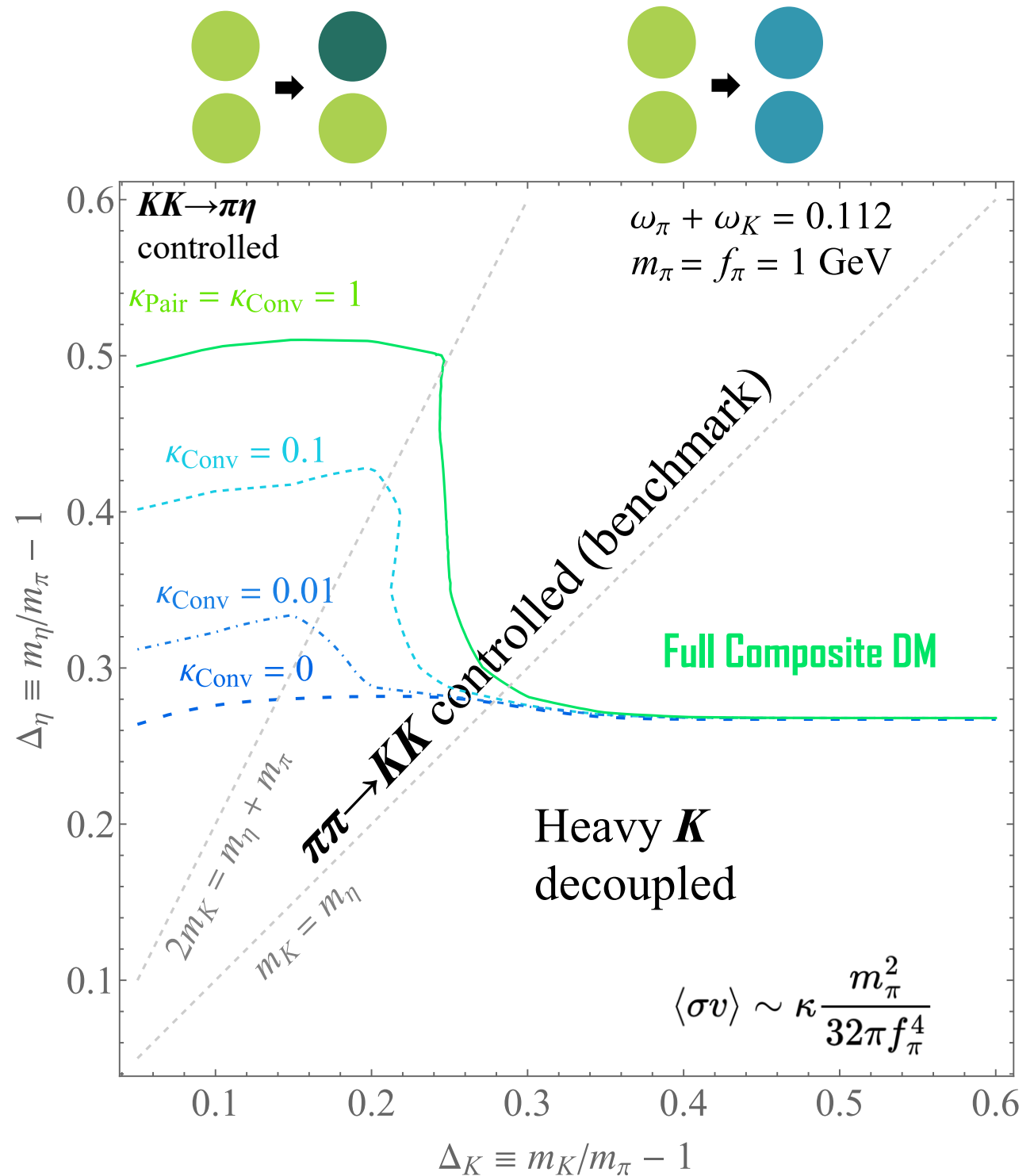
Effectively $\pi\pi$ to $\pi\eta$

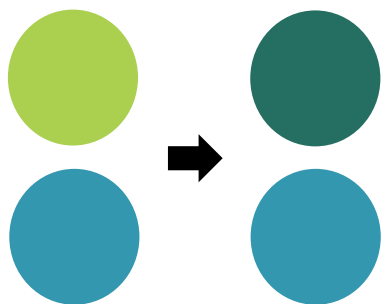


Extended Phases

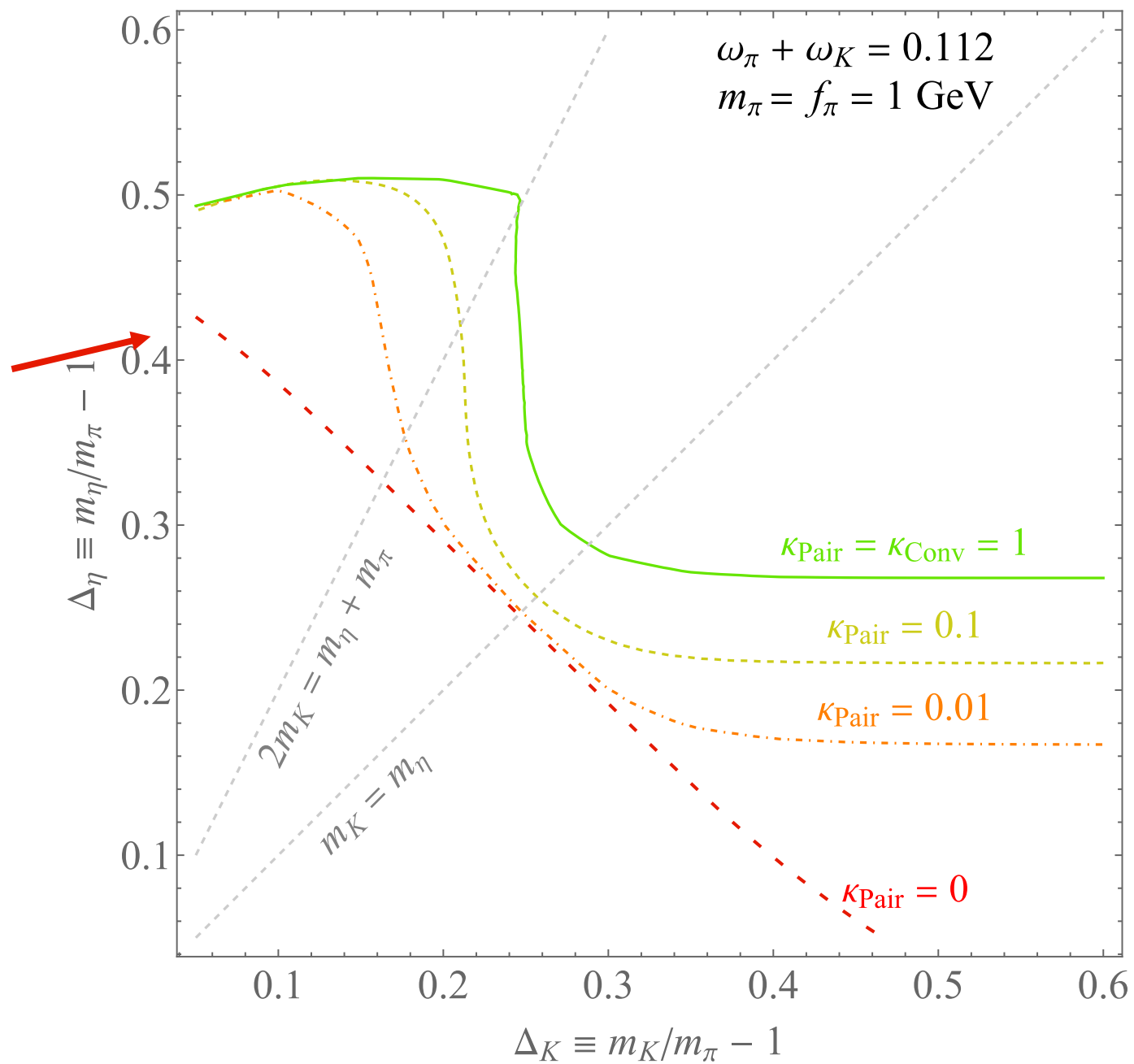
At least 3 extra dynamic phases recognized

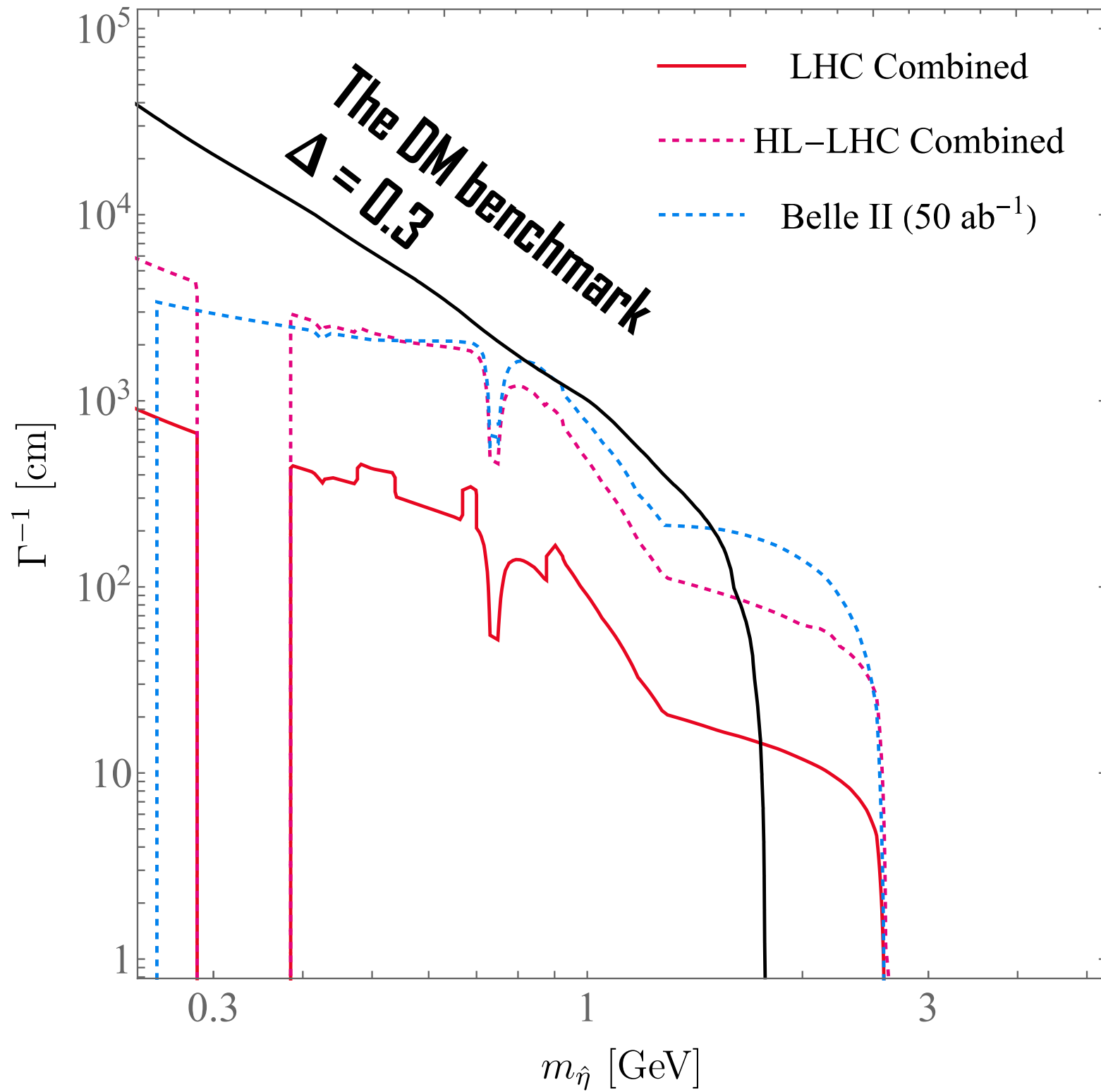
* Assuming all processes having similar cross sections, not necessarily true for realistic models

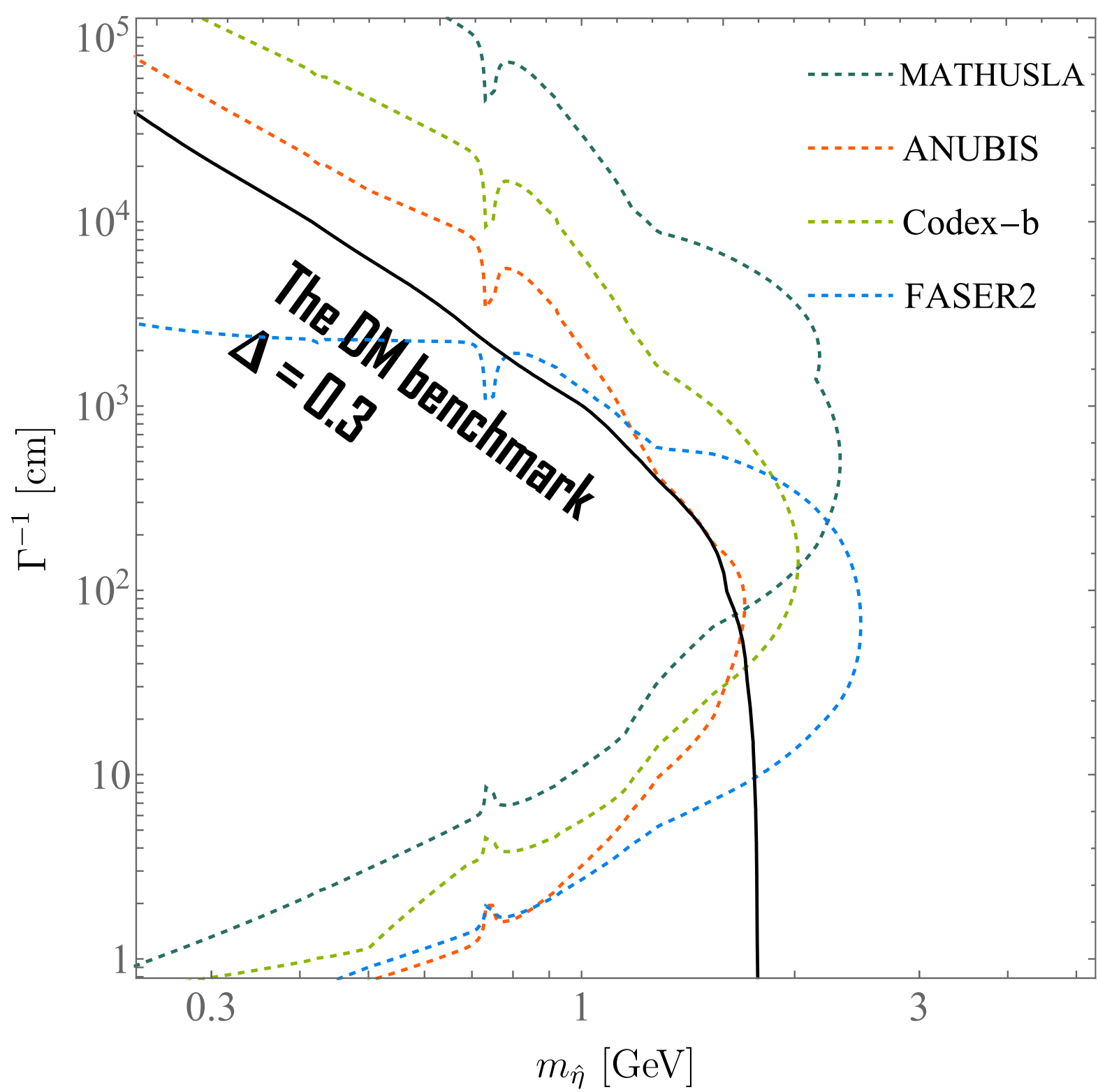




**Full $\pi K \rightarrow \eta K$
dominance**



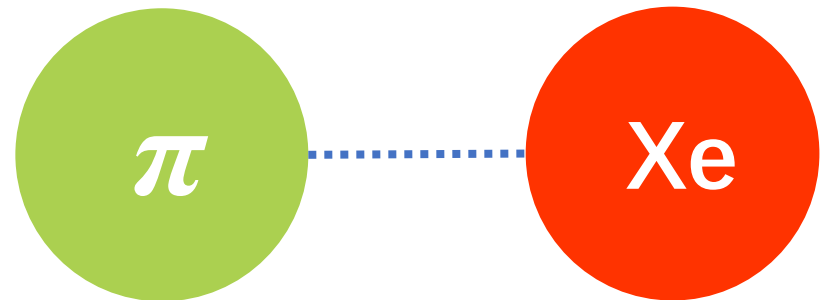




Direct Detection Prospects

The leading Z exchange between DM and SM are forbidden by the same symmetry that stabilizes the DM, the remaining Higgs exchange:

$$\mathcal{L} \supset 4B \frac{y_1 \tilde{y}_1}{M_{Q_1}} \pi_+ \pi_- |H|^2 ,$$



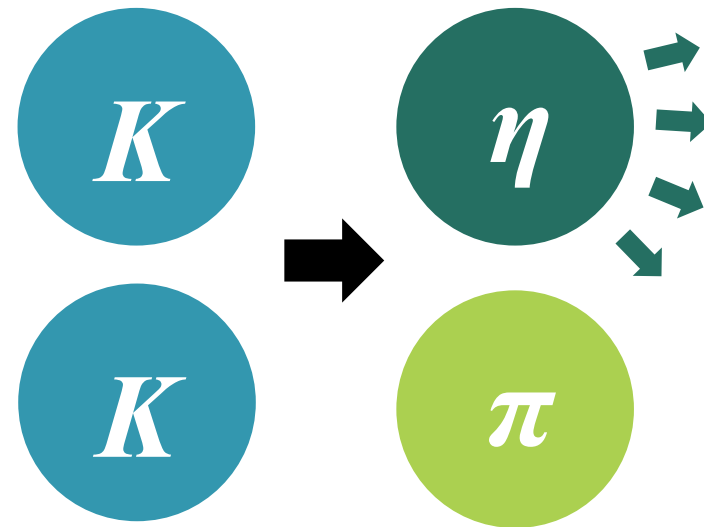
Nucleon
scattering Xsec:

$$\simeq 2.5 \times 10^{-45} \text{ cm}^2 b^2 \left(\frac{m_N}{m_\pi + m_N} \right)^2 \left(\frac{f}{\text{GeV}} \right)^2 \left(\frac{y_1 \tilde{y}_1}{10^{-2}} \right)^2 \left(\frac{\text{TeV}}{M_{Q_1}} \right)^2$$

When $m > 8 \text{ GeV}$, could be above the neutrino fog

Indirect Detection Prospects

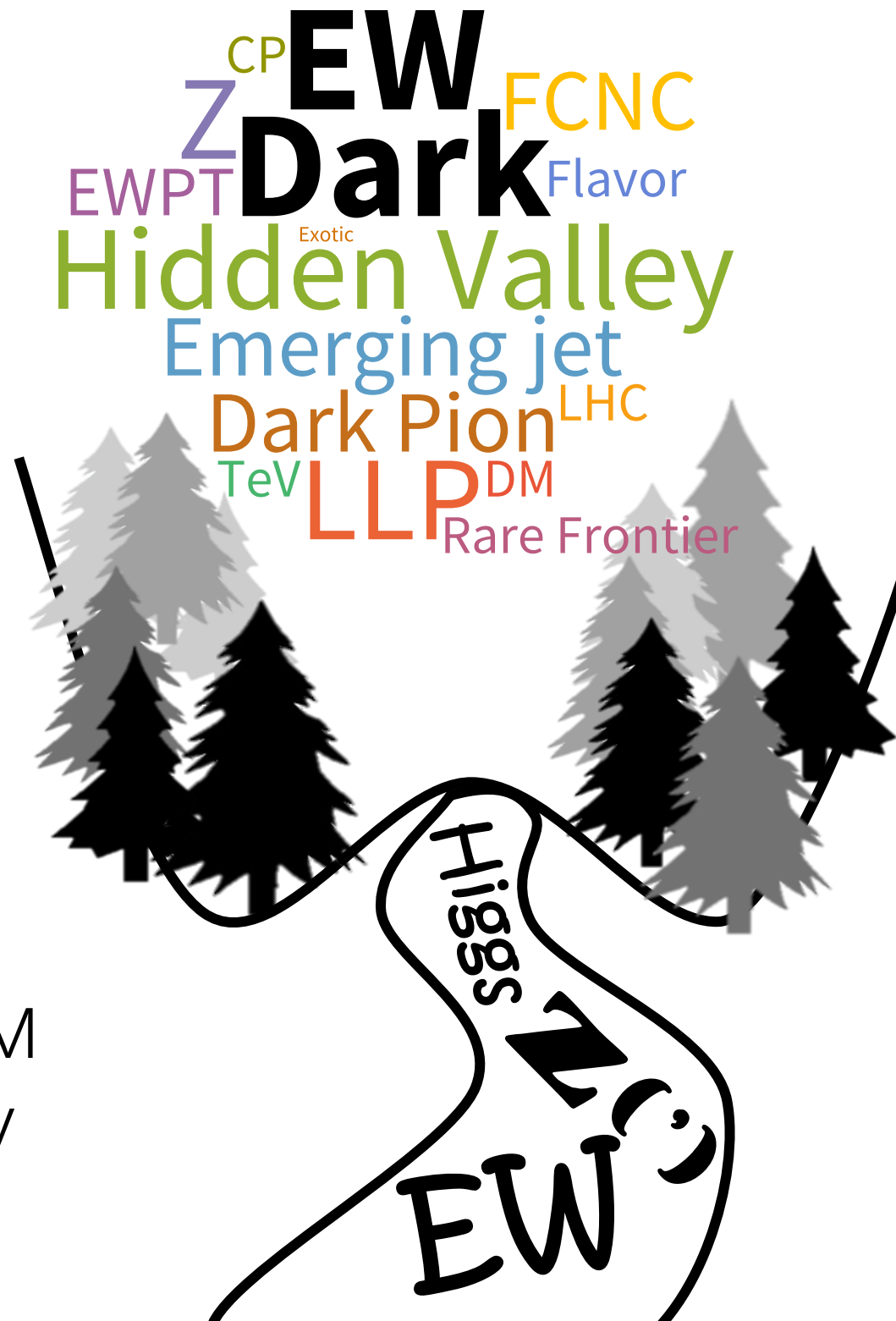
- As forbidden dark matter, usually free from indirect detection
- The remaining secondary Kaon DM can annihilate to SM, strongly depending on how many K remain



The most stringent CMB spectral distortion limit constraints the highly degenerate case ($\Delta < 1\%$)

Summary

- EW portal to the dark is a generic feature to be searched
- Dark confinement gives rise to LLP with good motivations
- At GeV scale, current and future limits are projected
- May lead to composite DM and non-trivial cosmology



Backup Slides

Lagrangian ($\mathbf{Z'}$ Mediation)

$$\mathcal{L}_{\text{mix}} = -\frac{\sin \chi}{2} \hat{Z}'_{\mu\nu} \hat{B}^{\mu\nu} + \delta \hat{M}^2 \hat{Z}'^\mu \hat{Z}_\mu - \kappa \Phi^* \Phi H^\dagger H$$

$$\begin{pmatrix} A_\mu \\ Z_\mu \\ Z'_\mu \end{pmatrix} = L^{-1} \begin{pmatrix} \hat{A}_\mu \\ \hat{Z}_\mu \\ \hat{Z}'_\mu \end{pmatrix} \quad L = \begin{pmatrix} 1 & -\hat{c}_W \sin \xi \tan \chi & -\hat{c}_W \cos \xi \tan \chi \\ 0 & \cos \xi + \hat{s}_W \sin \xi \tan \chi & -\sin \xi + \hat{s}_W \cos \xi \tan \chi \\ 0 & \sin \xi / \cos \chi & \cos \xi / \cos \chi \end{pmatrix}$$

$$\tan 2\xi = \frac{-2 \cos \chi (\delta \hat{M}^2 + \hat{M}_Z^2 \hat{s}_W \sin \chi)}{\hat{M}_{Z'}^2 - \hat{M}_Z^2 \cos^2 \chi + \hat{M}_Z^2 \hat{s}_W^2 \sin^2 \chi + 2\delta \hat{M}^2 \hat{s}_W \sin \chi}$$

Lagrangian (Fermion Mediation)

$$- \mathcal{L}_{\text{UV}} = \overline{Q}_L \mathbf{Y} \psi_R H + \overline{Q}_R \tilde{\mathbf{Y}} \psi_L H + \overline{Q}_L \mathbf{M} Q_R + \overline{\psi}_L \boldsymbol{\omega} \psi_R + \text{h.c.}$$

$$\begin{aligned} \mathcal{L}_{\text{EFT}} = & \frac{1}{2} \overline{\psi}_R \mathbf{Y}^\dagger \mathbf{M}^{-2} \mathbf{Y} \left[|H|^2 i \not{D} + i \gamma^\mu H^\dagger D_\mu H \right] \psi_R + \text{h.c.} \\ & + \frac{1}{2} \overline{\psi}_L \tilde{\mathbf{Y}}^\dagger \mathbf{M}^{-2} \tilde{\mathbf{Y}} \left[|H|^2 i \not{D} + i \gamma^\mu H^\dagger D_\mu H \right] \psi_L + \text{h.c.} \\ & - \overline{\psi}_L \boldsymbol{\omega} \psi_R + \overline{\psi}_L \tilde{\mathbf{Y}}^\dagger \mathbf{M}^{-1} \mathbf{Y} \psi_R |H|^2 + \text{h.c.} , \end{aligned}$$

Z Pole EWPT (Z' Mediation)

$$\mathcal{L}_{Zf\bar{f}} = -\frac{\bar{Z}e}{s_W c_W} \bar{f} \gamma^\mu (T_{Lf}^3 - s_*^2 Q_f) f Z_\mu$$

$$\bar{Z} = 1 + \frac{\alpha T}{2}, \quad s_*^2 = s_W^2 + \frac{\alpha}{c_W^2 - s_W^2} \left(\frac{S}{4} - s_W^2 c_W^2 T \right)$$

$$S = \frac{4s_W^2}{\alpha} \left(\frac{c_W^2}{s_W} \xi t_\chi - c_W^2 \xi^2 \right)$$

$$T = \frac{1}{\alpha} \left(2s_W \xi t_\chi + \xi^2 \left(\frac{M_{Z'}^2}{M_Z^2} - 2 \right) \right)$$

$$U = \frac{4s_W^2}{\alpha} c_W^2 \xi^2$$

Flashback I: Accidental SUSY with Neutral Naturalness

$$W'_{Z_3} = y_t(Q_A H u_A^c + Q_B H u_B^c + Q_C H u_C^c) + \omega(u'_B u_B^c + u'_C u_C^c) + M(Q_B Q_B'^c + Q_C Q_C'^c)$$

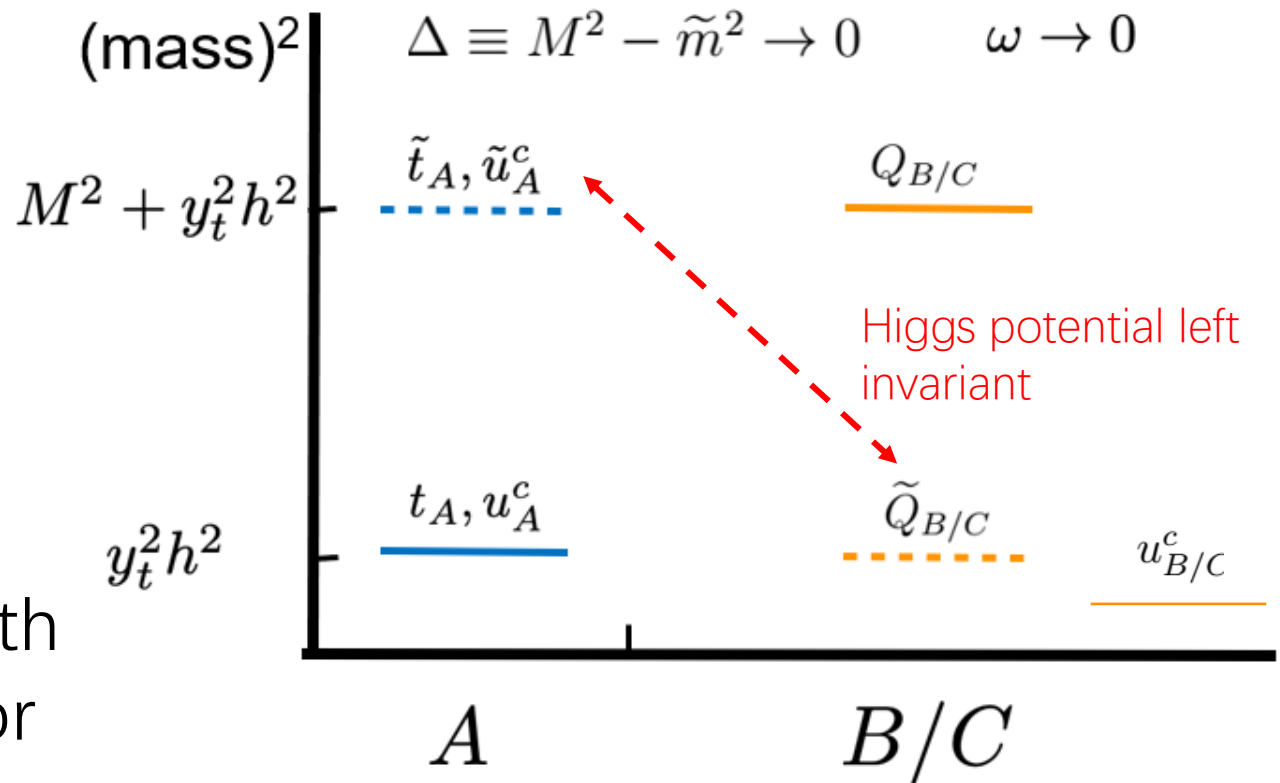
$$V'_s = \tilde{m}^2(|\tilde{Q}_A|^2 + |\tilde{u}_A^c|^2) - \tilde{m}^2(|\tilde{Q}_B|^2 + |\tilde{Q}_C|^2)$$

$Q_A \ u_A^c$ SM top sector

$$Q_{B,C} = \begin{pmatrix} t_{B,C} \\ b_{B,C} \end{pmatrix} \sim 2_{-1/2}$$

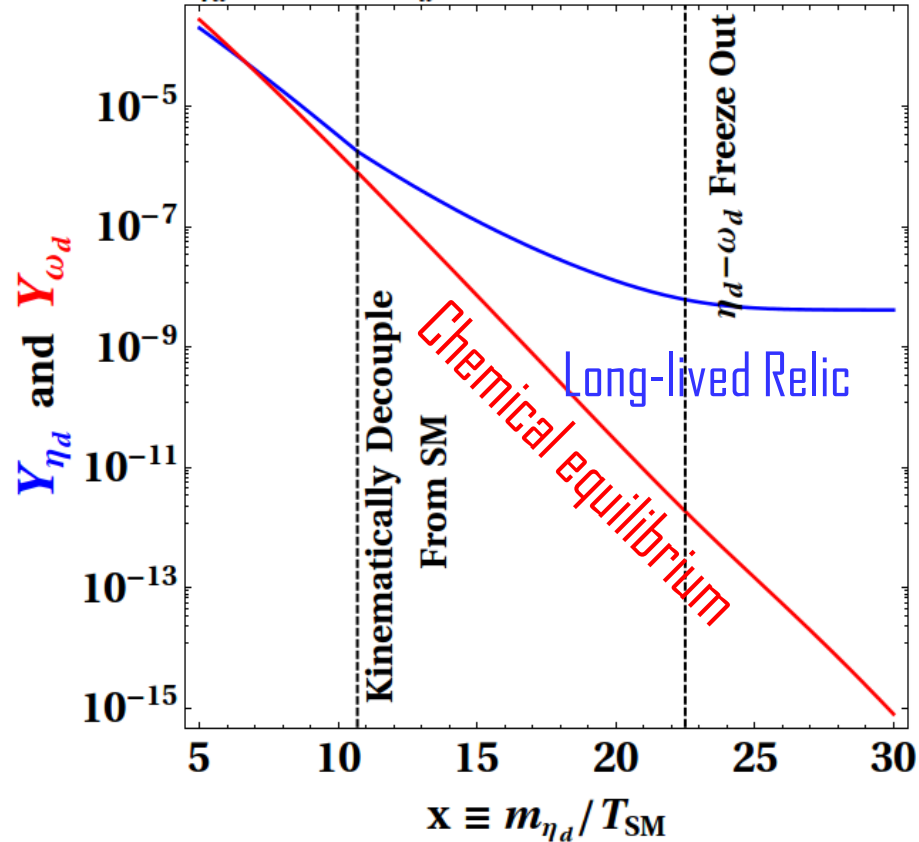
$$u_{B/C}^c \sim 1_0$$

➤ Only compatible with
the LHC if dark color
exists



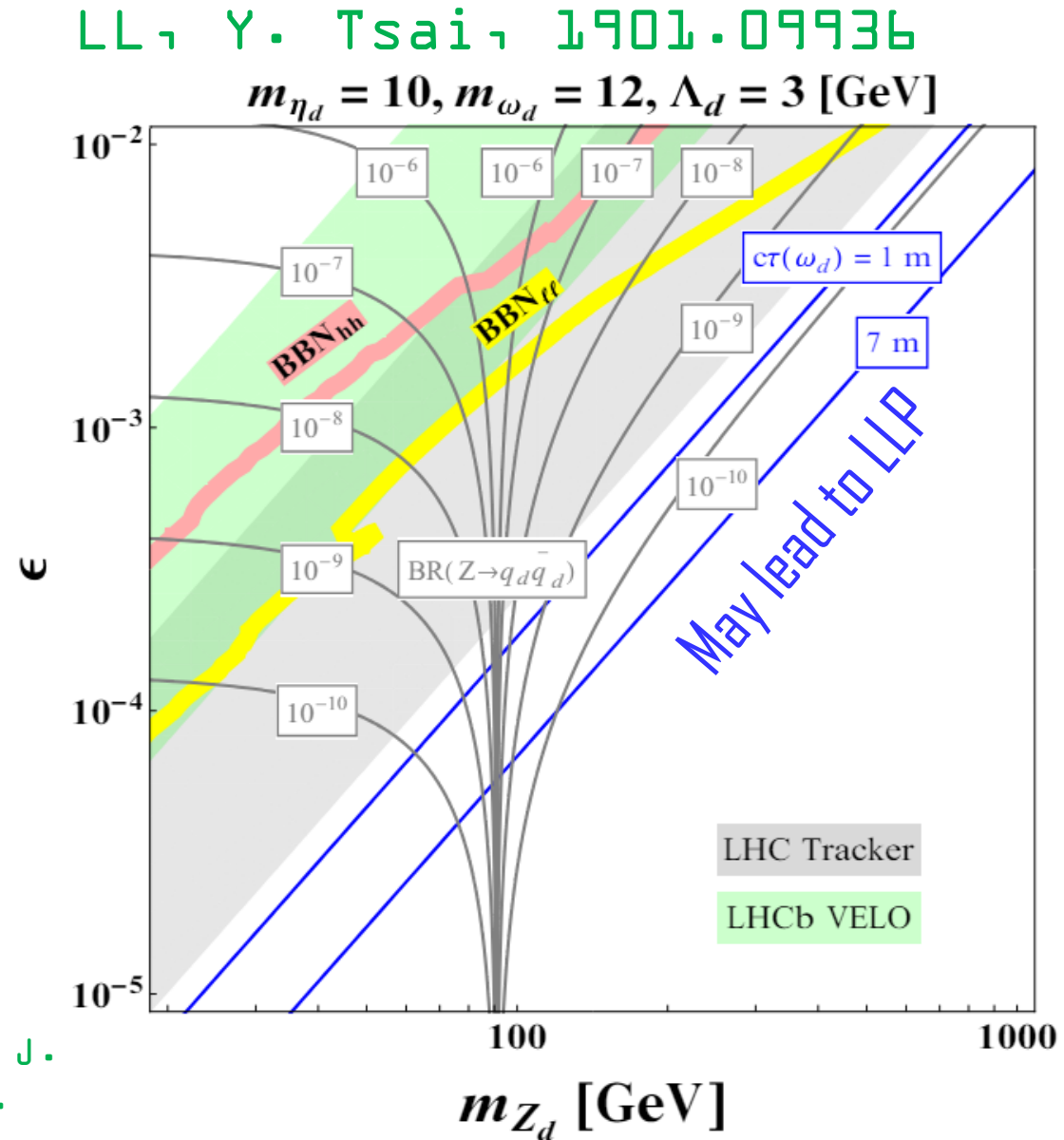
Flashback II: Composite DM

$$m_{\eta_d} = 10, m_{\omega_d} = 12 \text{ GeV}, c\tau(\omega_d) = 1 \text{ cm}$$



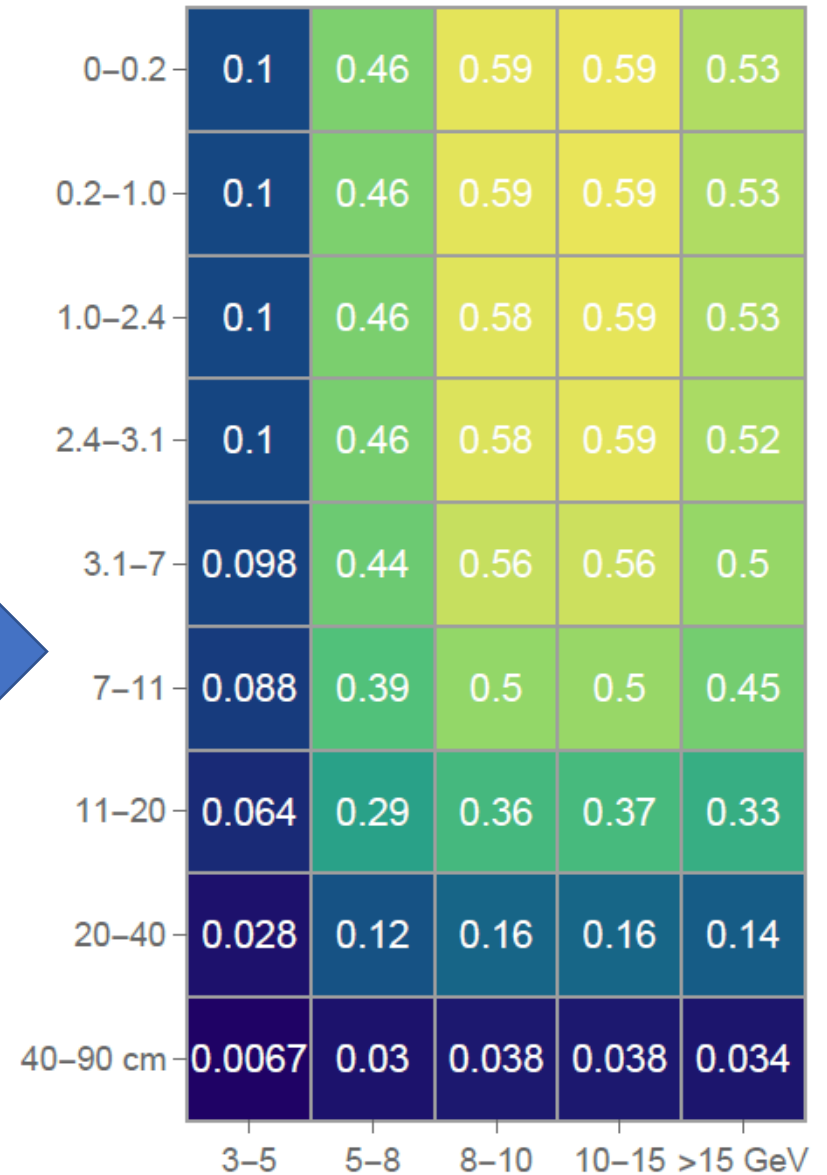
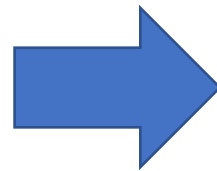
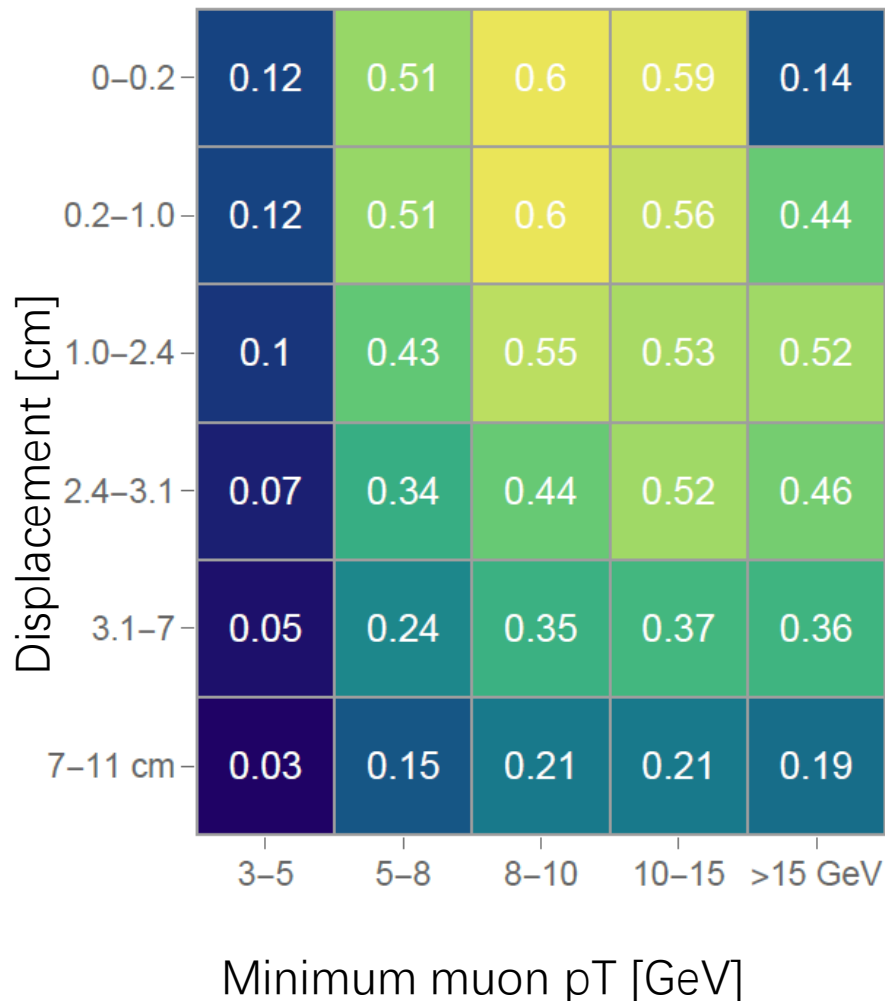
E. Bernreuther, F. Kahlhoefer, M. Kramer and P. Tunney, 1907.04346; S. Koren and R. McGehee, 1908.03559; Y. Hochberg, E. Kuflik, T. Volansky and J. G. Wacker, 1402.5143; Y. Hochberg, E. Kuflik and H. Murayama, 1512.07917

 LL, Y. Tsai, 1901.09936



Displaced Vertex Efficiency for (HL-)LHC

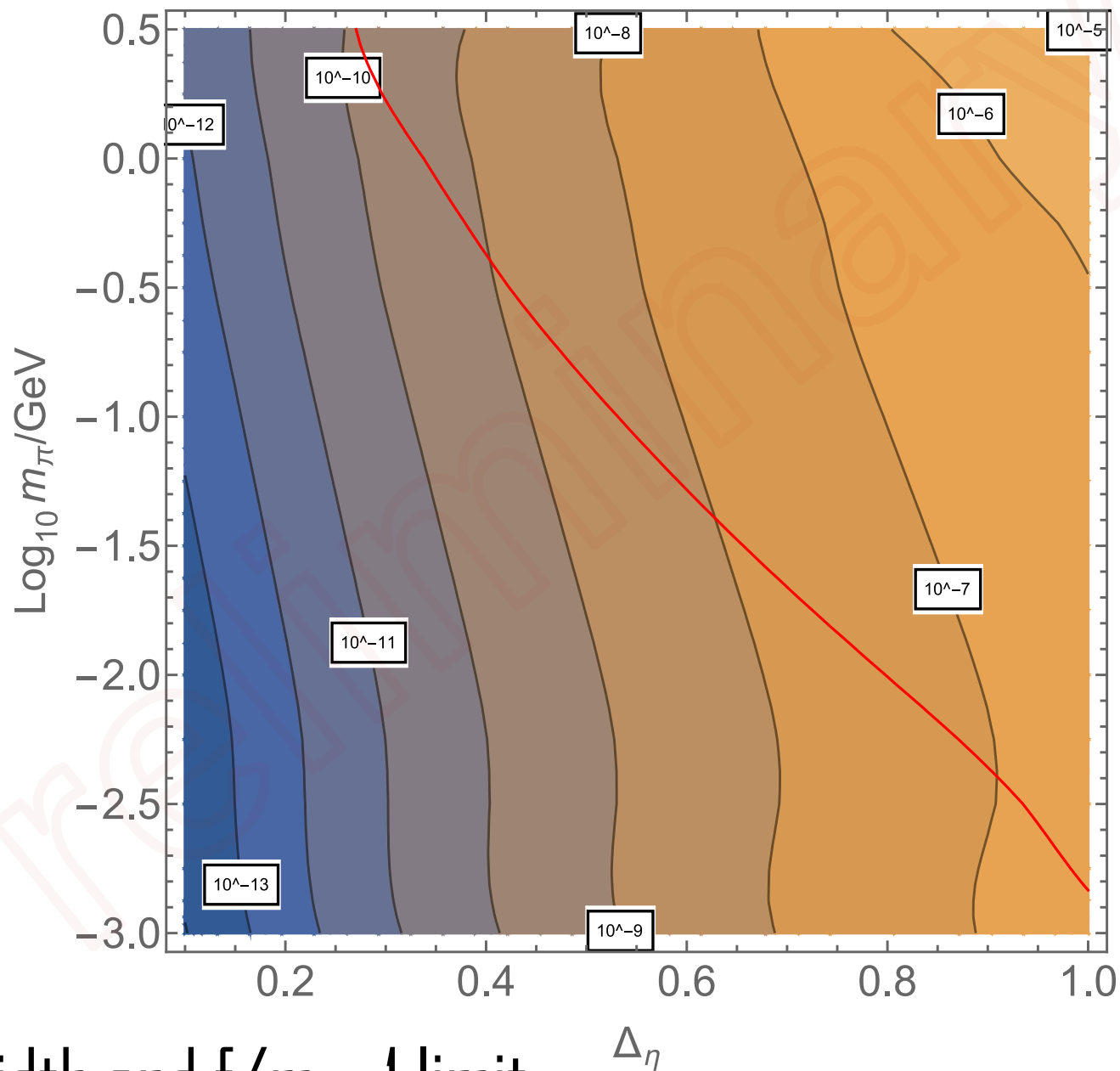
Current CMS:



Extended Higgs Sector

$$H \sim \mathbf{2}_{1/2,0}, H' \sim \mathbf{2}_{1/2,x_\Phi}, \Phi \sim \mathbf{1}_{0,x_\Phi}$$

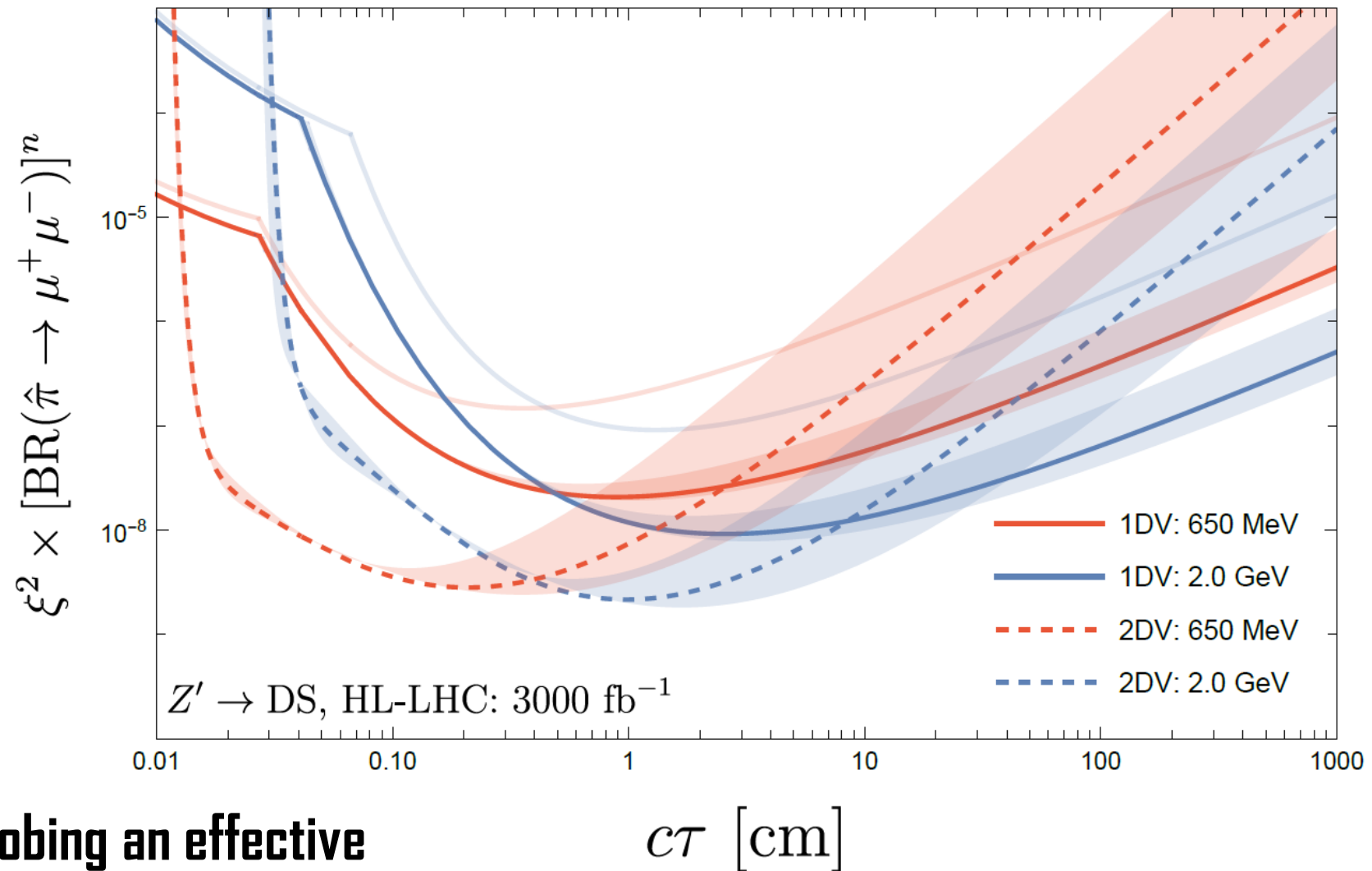
$$V = -\frac{m_H^2}{2}h^2 + \frac{\lambda_H}{4}h^4 + \frac{m_{H'}^2}{2}h'^2 + \frac{\lambda_{H'}}{4}h'^4 - m_\Phi^2\varphi^2 + \lambda_\Phi\varphi^4 - \mu hh'\varphi$$



In large width and $f/m = 1$ limit

CMS: Results on a Light Z'

Efficiency drop due to soft muons (can't reach the muon chamber of $p_T < 3$ GeV)



Still probing an effective
mixing of 10^{-3} - 10^{-4}

Dark Pion Decays (ALP-Like)

Applies to arbitrary
flavor couplings

See also:

D. Aloni, Y. Soreq and M.
Williams, 1811.03474

A.1 $a \rightarrow \gamma\gamma$

A.2 $a \rightarrow \pi^+\pi^-\gamma$

A.3 $a \rightarrow \pi^+\pi^-\pi^0$

A.4 $a \rightarrow 3\pi^0$

A.5 $a \rightarrow \pi^0\pi^0\eta, \pi^+\pi^-$

A.6 $a \rightarrow \pi^0\pi^0\eta', \pi^+\pi^-$

A.7 $a \rightarrow \eta\eta\pi^0$

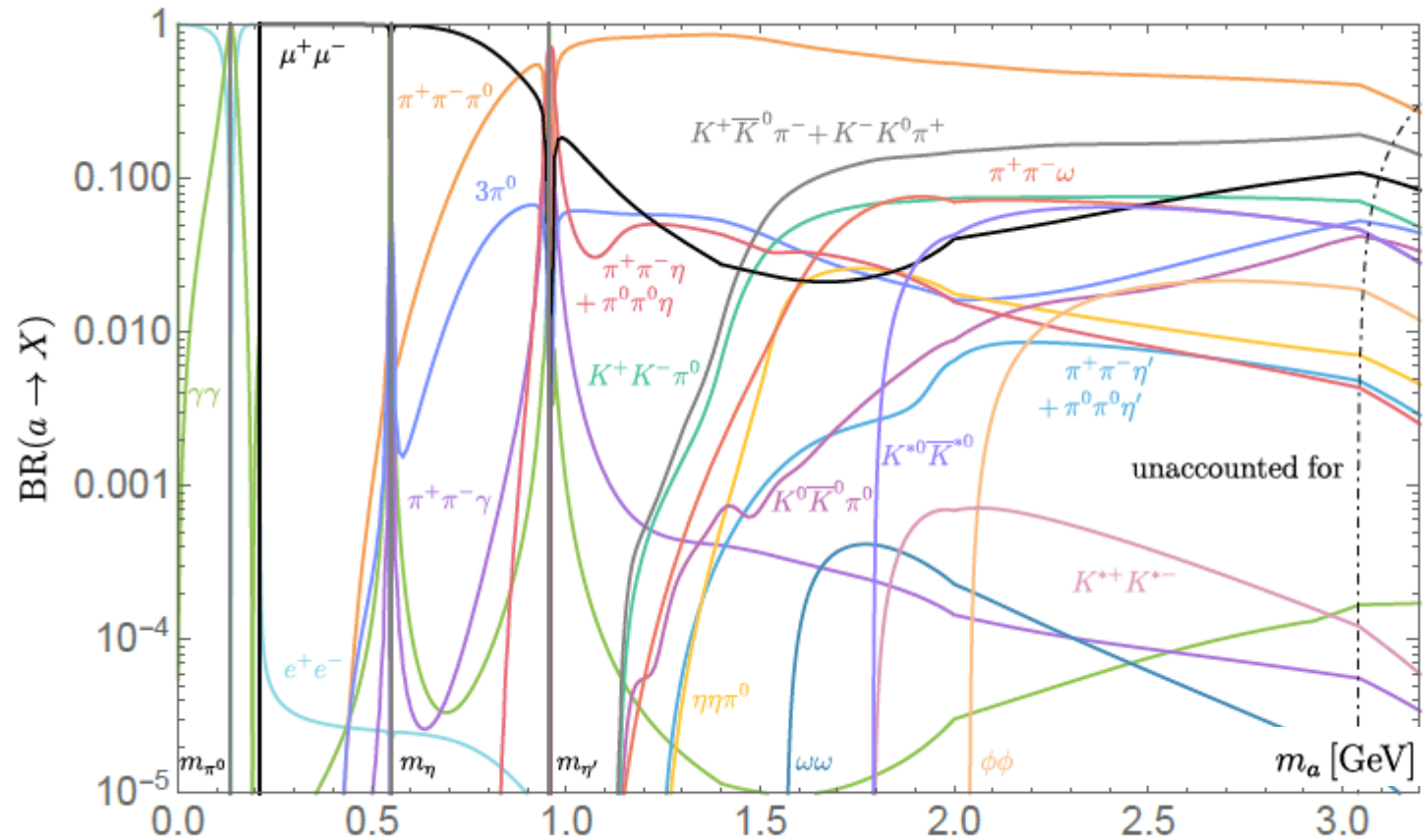
A.8 $a \rightarrow K^0\bar{K}^0\pi^0$

A.9 $a \rightarrow K^+K^-\pi^0$

A.10 $a \rightarrow K^+\bar{K}^0\pi^-, K^-K^0\pi^+$

A.11 $a \rightarrow \omega\omega, \phi\phi, K^{*+}K^{*-}, K^{*0}\bar{K}^{*0}$

A.12 $a \rightarrow \pi^+\pi^-\omega$



$$\Gamma_{(\hat{\pi} \rightarrow \mu^+ \mu^-)}^{-1} \simeq 45 \text{ cm} \left(\frac{f_a}{1 \text{ PeV}} \right)^2 \left(\frac{1 \text{ GeV}}{m_{\hat{\pi}}} \right) \sqrt{\frac{m_{\hat{\pi}}^2}{m_{\hat{\pi}}^2 - 4m_{\mu}^2}}$$

