

Review of polarization in heavy ion collisions at STAR

Jinhui Chen

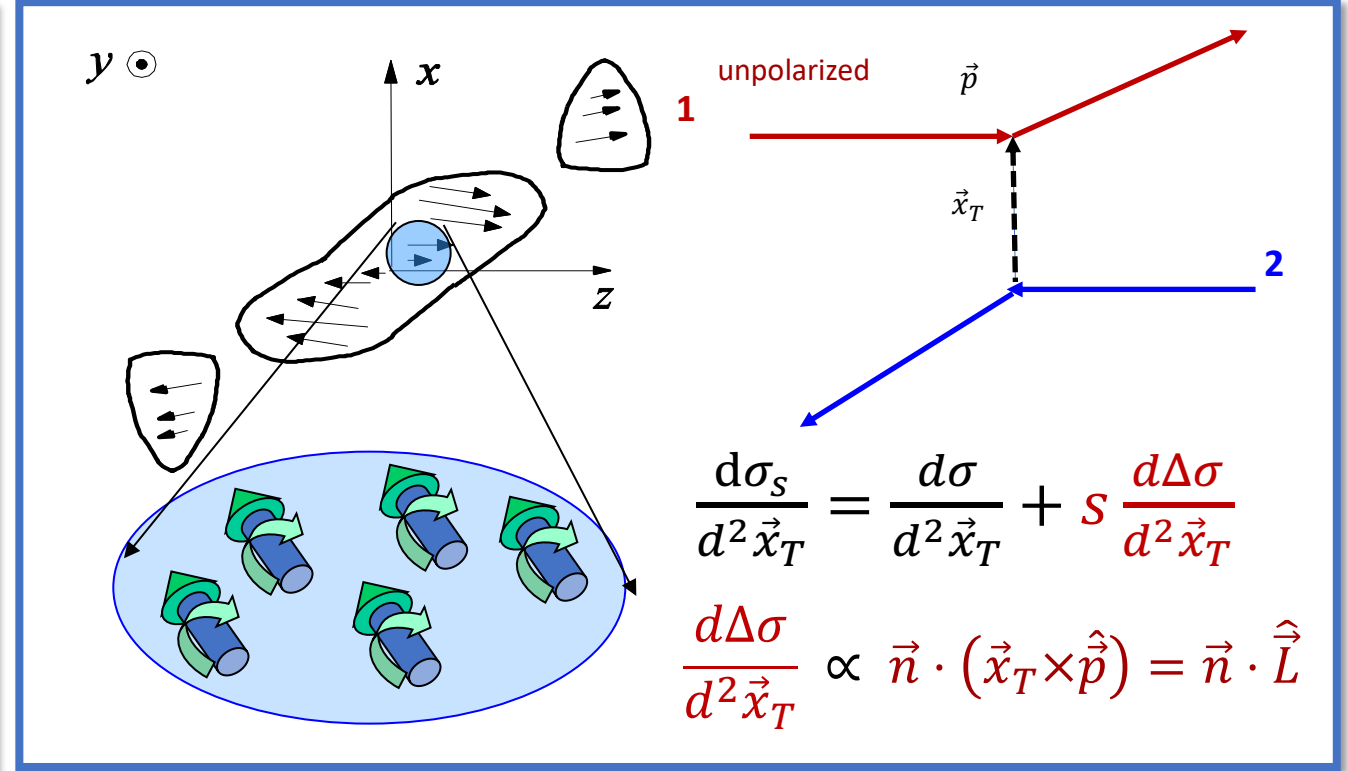
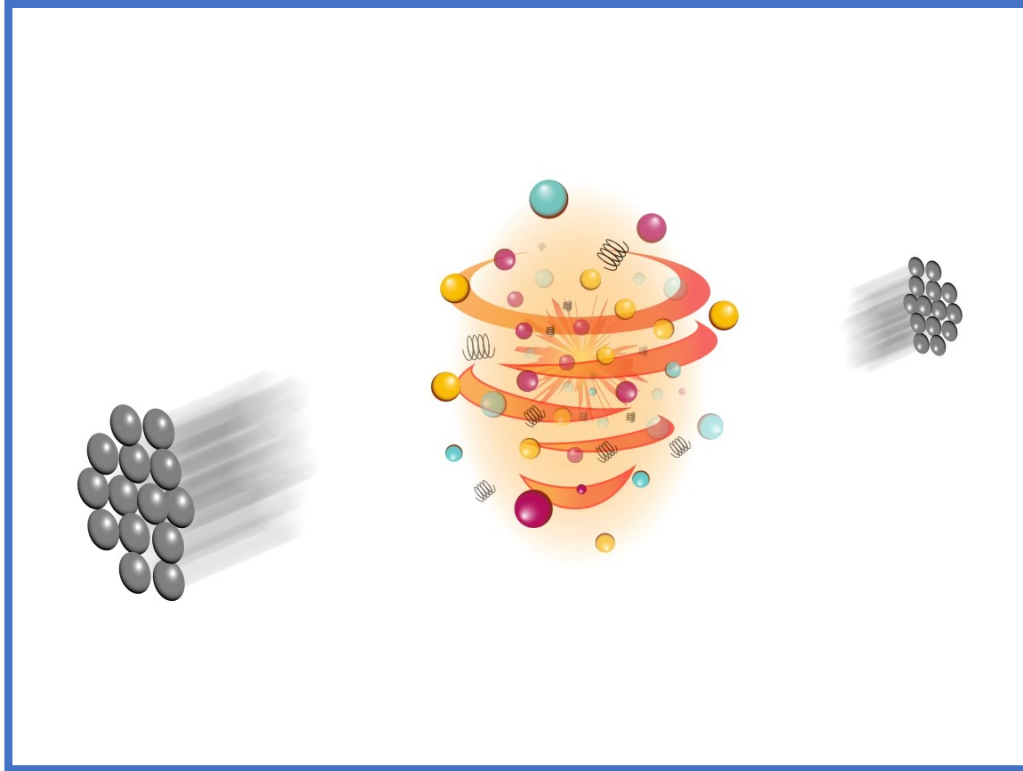
Fudan University

“The 1st Conference on Strong-Interaction Spin Physics”

Qingdao, 2026/7/26-31

New phenomenon: global pol. in HIC

Liang and Wang Phys. Rev. Lett. **94**, 102301 (2005); Phys. Lett. B **629**, 20 (2025)



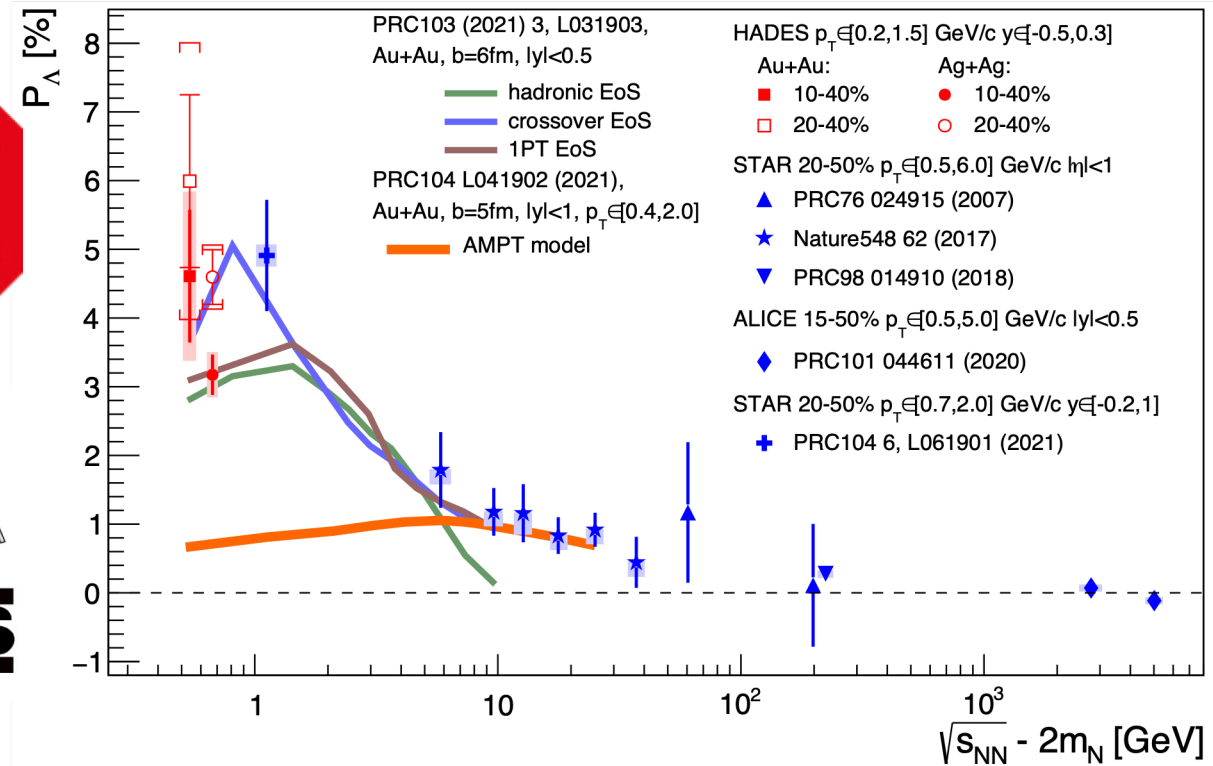
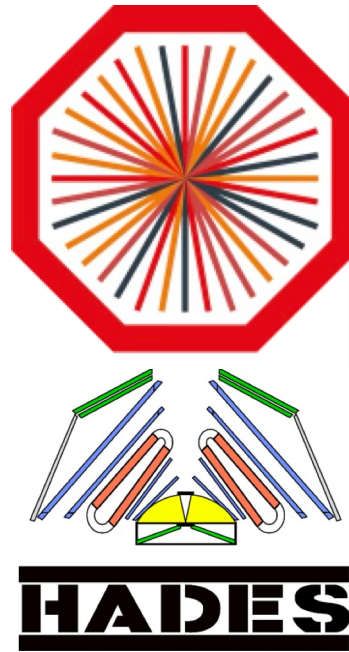
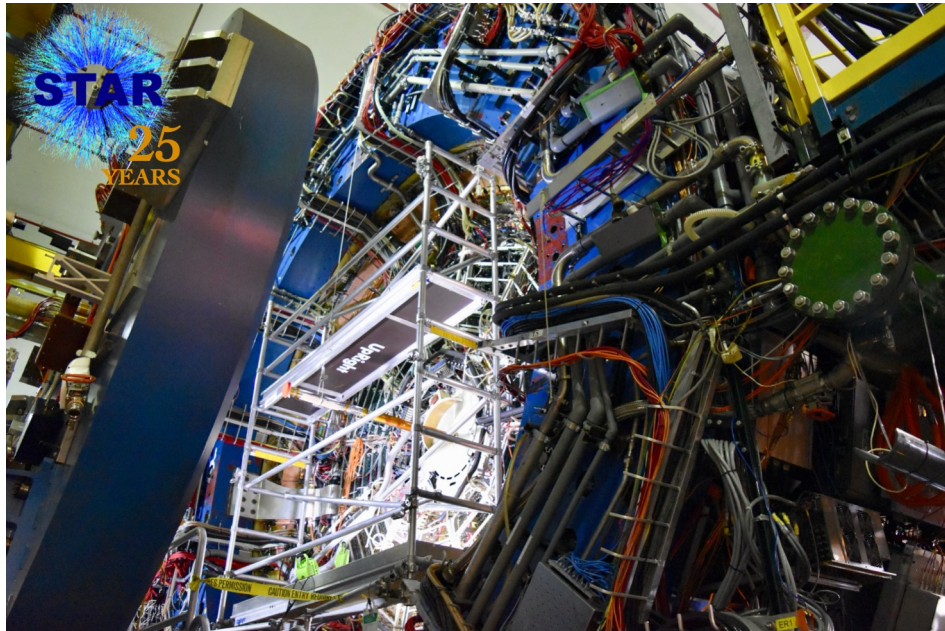
- Produced partons are shown to have large local relative OAM in noncentral heavy-ion collisions.
- Parton scattering with given relative OAM is shown to polarize quarks along the same direction due to spin-orbit coupling.
- Such global quark polarization has many measurable consequences in high-energy heavy-ion collisions.

Measurable consequences of Λ polarization

- The idea has yielded an active 20-year collective efforts in experimental and theoretical side.

Liang and Wang, PRL (2005) Inspire 693, STAR Col., Nature (2017) Inspire 1081

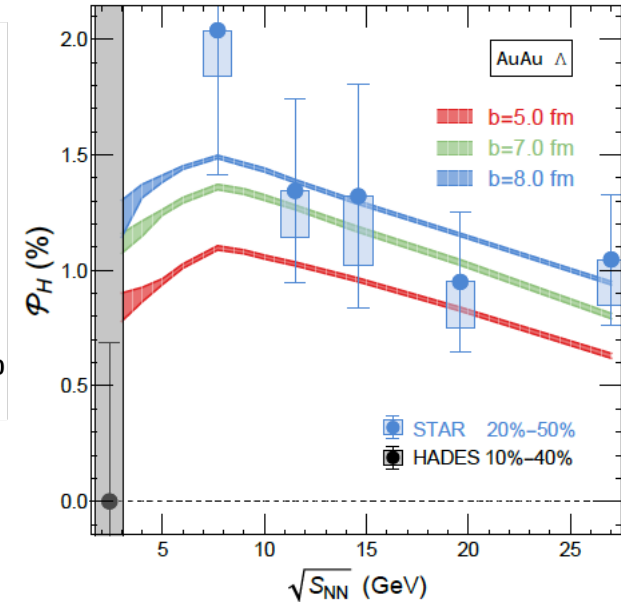
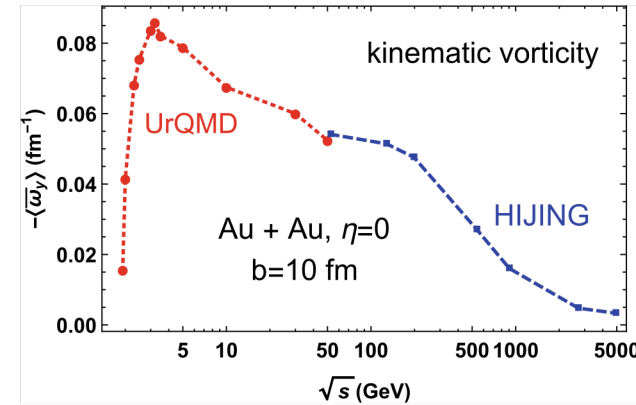
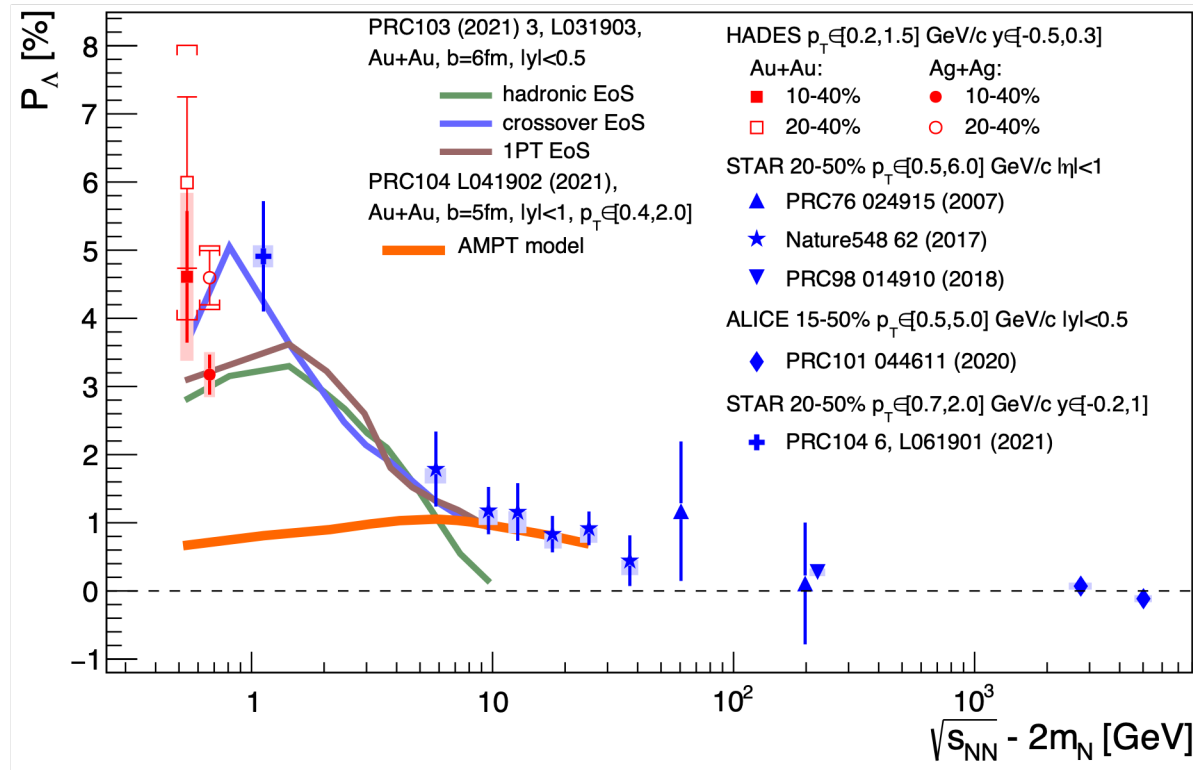
$$\frac{dN}{d \cos \theta^*} = \frac{1}{2} (1 + \alpha_H |\mathcal{P}_H| \cos \theta^*) \quad \overline{\mathcal{P}}_H \equiv \langle \mathcal{P}_H \cdot \hat{j}_{\text{sys}} \rangle = \frac{8}{\pi \alpha_H} \frac{\langle \cos(\phi_p^* - \phi_{j_{\text{sys}}}) \rangle}{R_{\text{EP}}^{(1)}}$$



A collection of data show global polarization in broad energy range, and identify the most vortical fluid ever created.

Λ polarization: energy dept. or independent

HADES Col. Phys. Lett. B **835**, 137506 (2022)



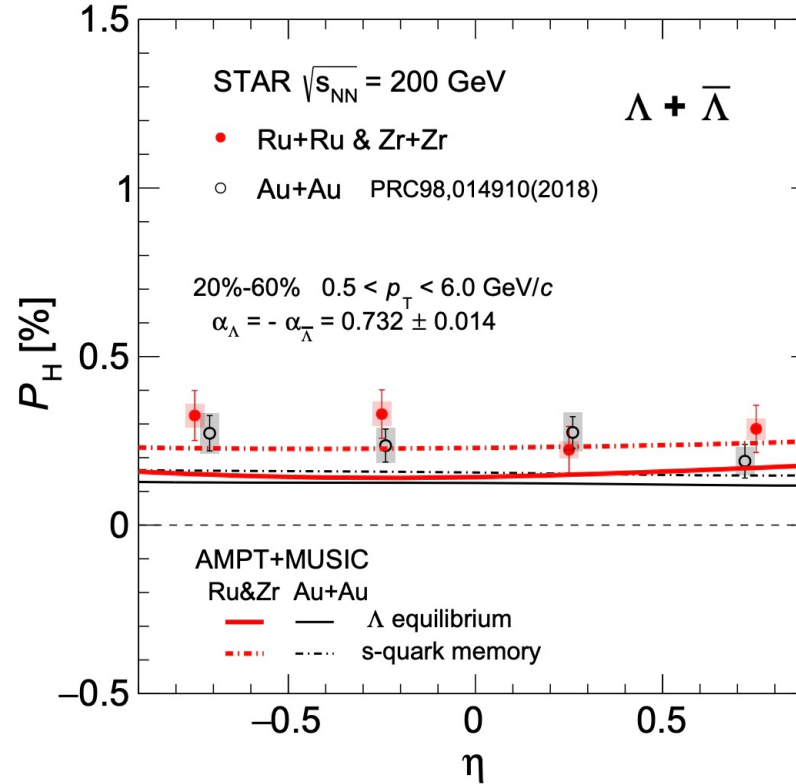
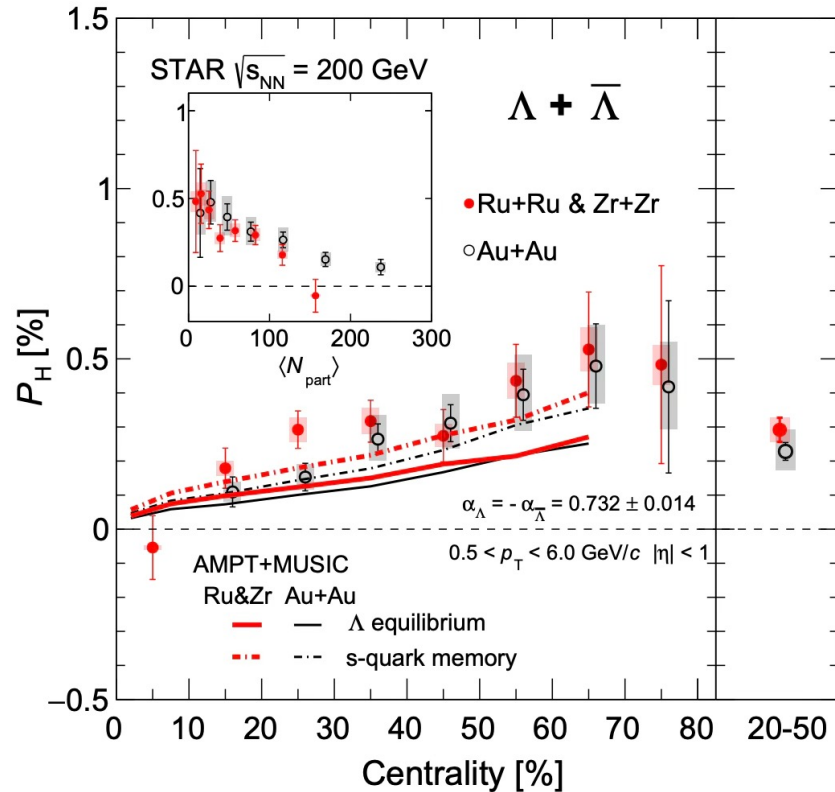
Deng, Huang, Ma, Zhang Phys. Rev. C **101**, 064908 (2020)
 Guo, Liao, Wang, Xing, Zhang Phys. Rev. C **104**, L041902 (2021)
 Wu, Yi, Qin, Pu, Phys. Rev. C **105**, 064909 (2022)
 (HADES data, Springer Proc. Phys. 250, 435 (2020))

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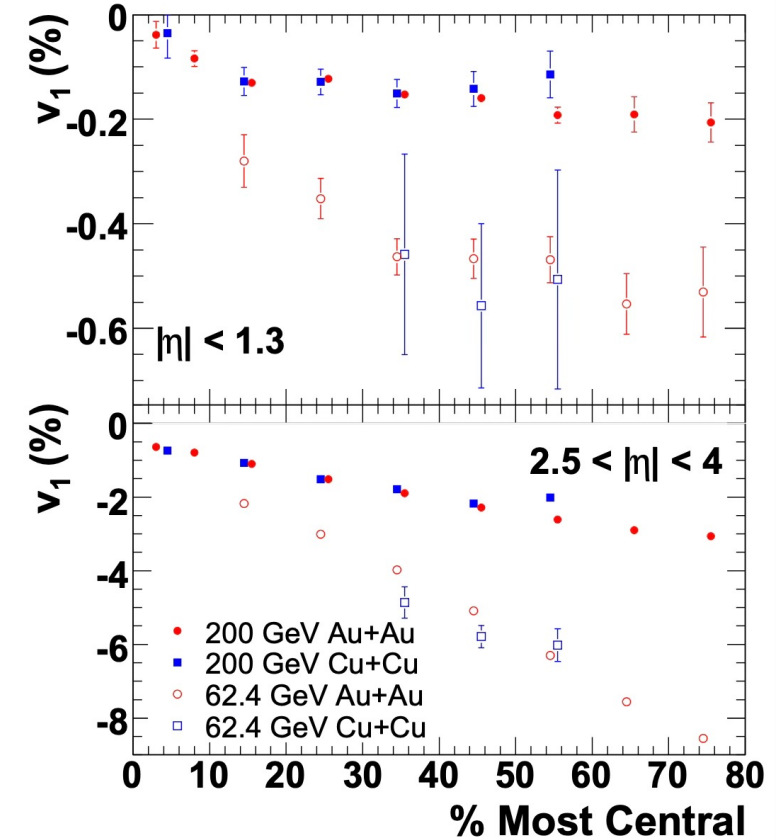
- Lambda polarization at 1-10 GeV presents strong energy dependence, which was described by the fluid dynamics, while transport model AMPT described the BES-I data was lower than data at 1 GeV.
- How to understand the underlying physics, is there any turn over point at lower energies?

Λ polarization: system size dept. or independent

STAR Col., Phys. Lett. B **870**, 139891 (2025)

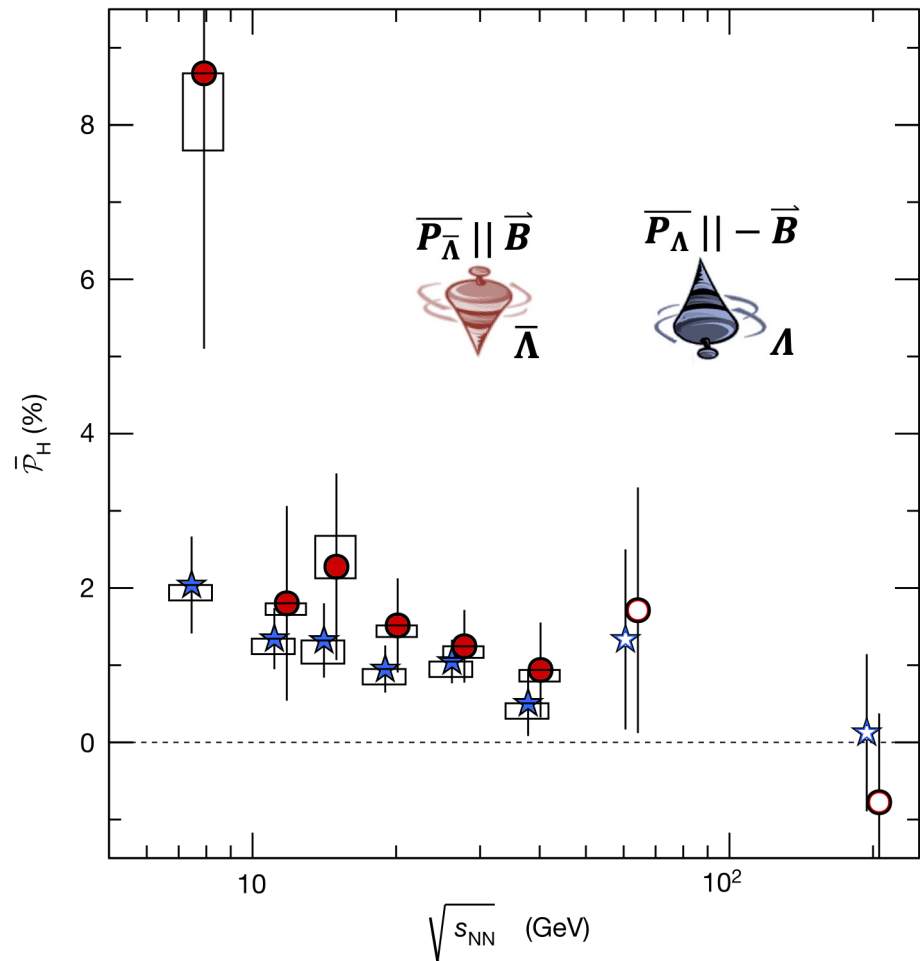


STAR Col., Phys. Rev. Lett. **101**, 252301 (2008)

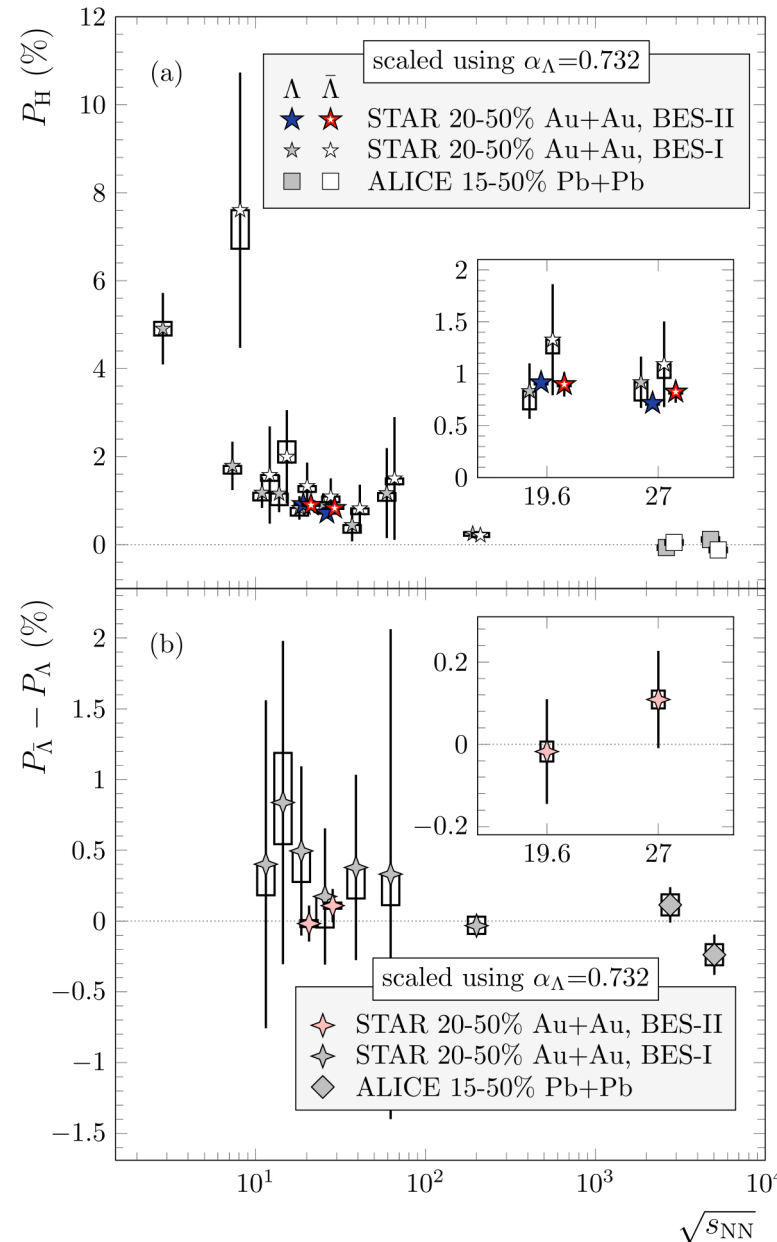


- System size independence?
- Both polarization and directed flow show system size independence
- physics or coincidence?

Measurable consequences of Λ polarization splitting



$$|B| \approx \frac{T_s |P_{\bar{\Lambda}} - P_{\Lambda}|}{2|\mu_{\Lambda}|}$$



STAR Col. Nature **548**, 62 (2017);
STAR Col., Phys. Rev. C **108**, 014910 (2023)

- High statistics BES-II data, 19.6 and 27 GeV
 - -0.018 ± 0.127 (stat.) $\pm$ 0.024 (syst.)
 - $B < 9.4 \times 10^{12}$ T
- 0.109 ± 0.118 (stat.) $\pm$ 0.022 (syst.)
- $B < 1.4 \times 10^{13}$ T

Λ polarization splitting: B-field vs. meson?

Csernai, Kapusta, Welle, Phys. Rev. C **99**, 021901(R) (2019)

- Interactions of baryons with mesons can result in a splitting of the polarization of Lambdas in high energy heavy ion collisions.

$$H_{\text{spin-orbit}}^V = \frac{g_{V\Lambda}}{2m_\Lambda^2} \frac{1}{r} \frac{\partial V_0}{\partial r} \mathbf{S} \cdot \mathbf{L}, \quad (5)$$

$$H_{\text{spin-orbit}}^\sigma = \frac{g_{\sigma\Lambda}}{2m_\Lambda^2} \mathbf{S} \cdot \nabla \sigma \times \mathbf{p}, \quad (6)$$

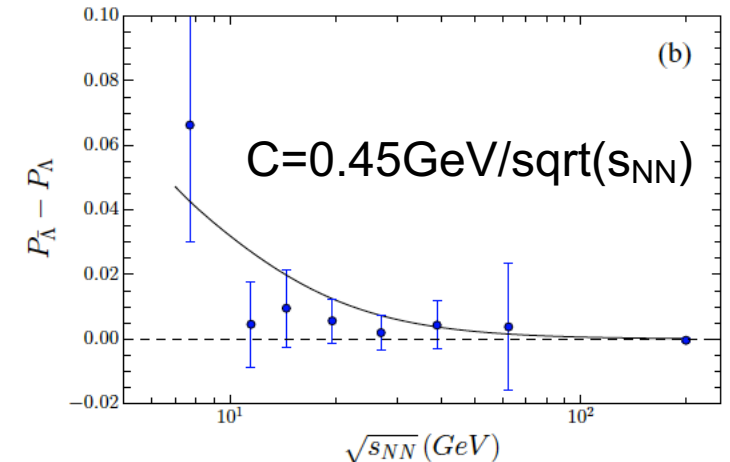
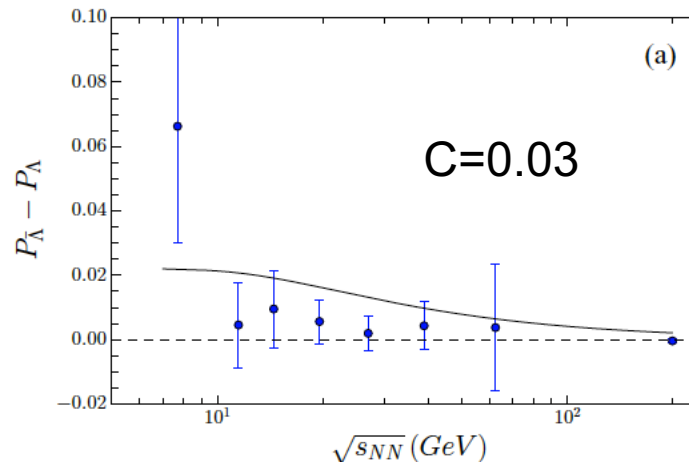
‘It is the **phi meson** which is the primary origin of the polarization difference, which should be no surprise because of the matching of the strange quark spin content of these hadrons.’

- The interaction is orders of magnitude larger than the one arising from electromagnetic fields.

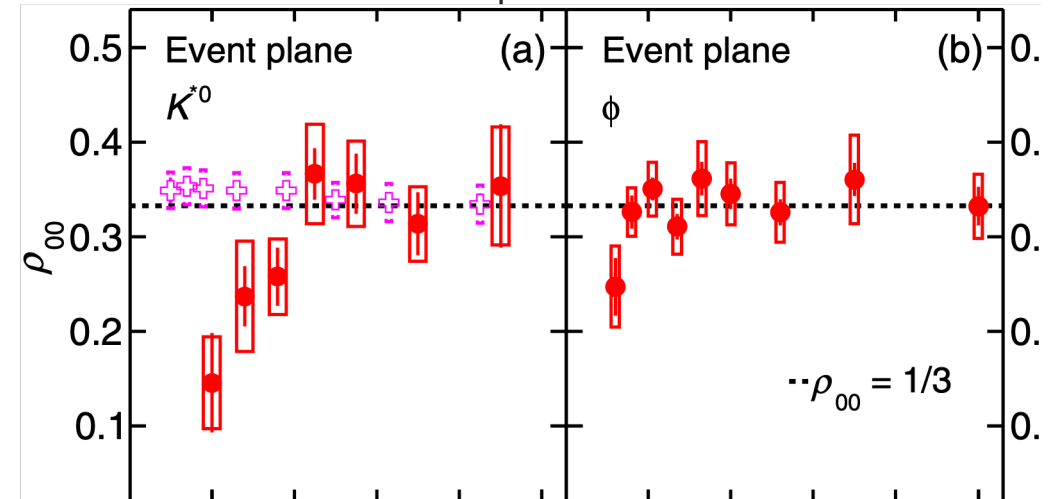
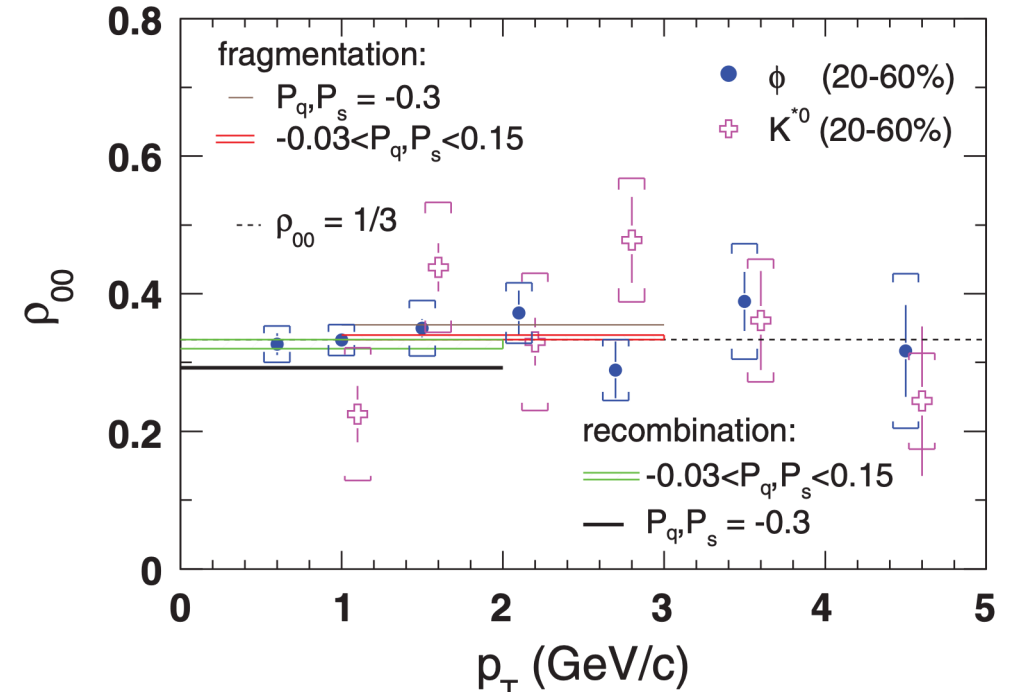
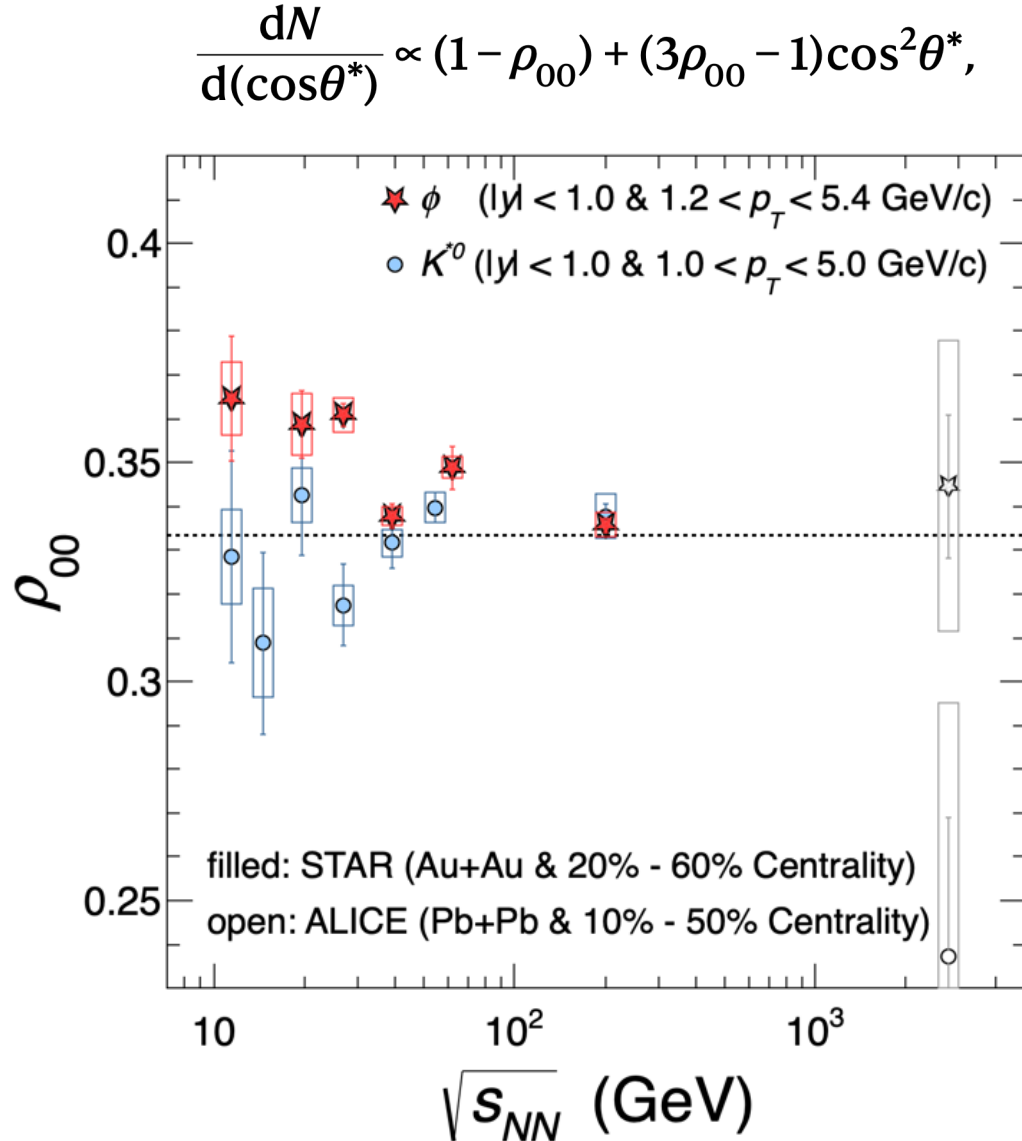
$$B = \frac{eb\sigma_E}{8\pi t^2} \exp(-b^2\sigma_E/4t), \quad P_\gamma \sim 10^{-6}$$

$$P_{\bar{\Lambda}} - P_{\Lambda} = C \left(\frac{n_B(t_{\text{ch}})}{0.15/\text{fm}^3} \right) \left(\frac{140 \text{ MeV}}{T(t_{\text{ch}})} \right).$$

‘The results are sensitive to the space-time evolution of the baryon current.’

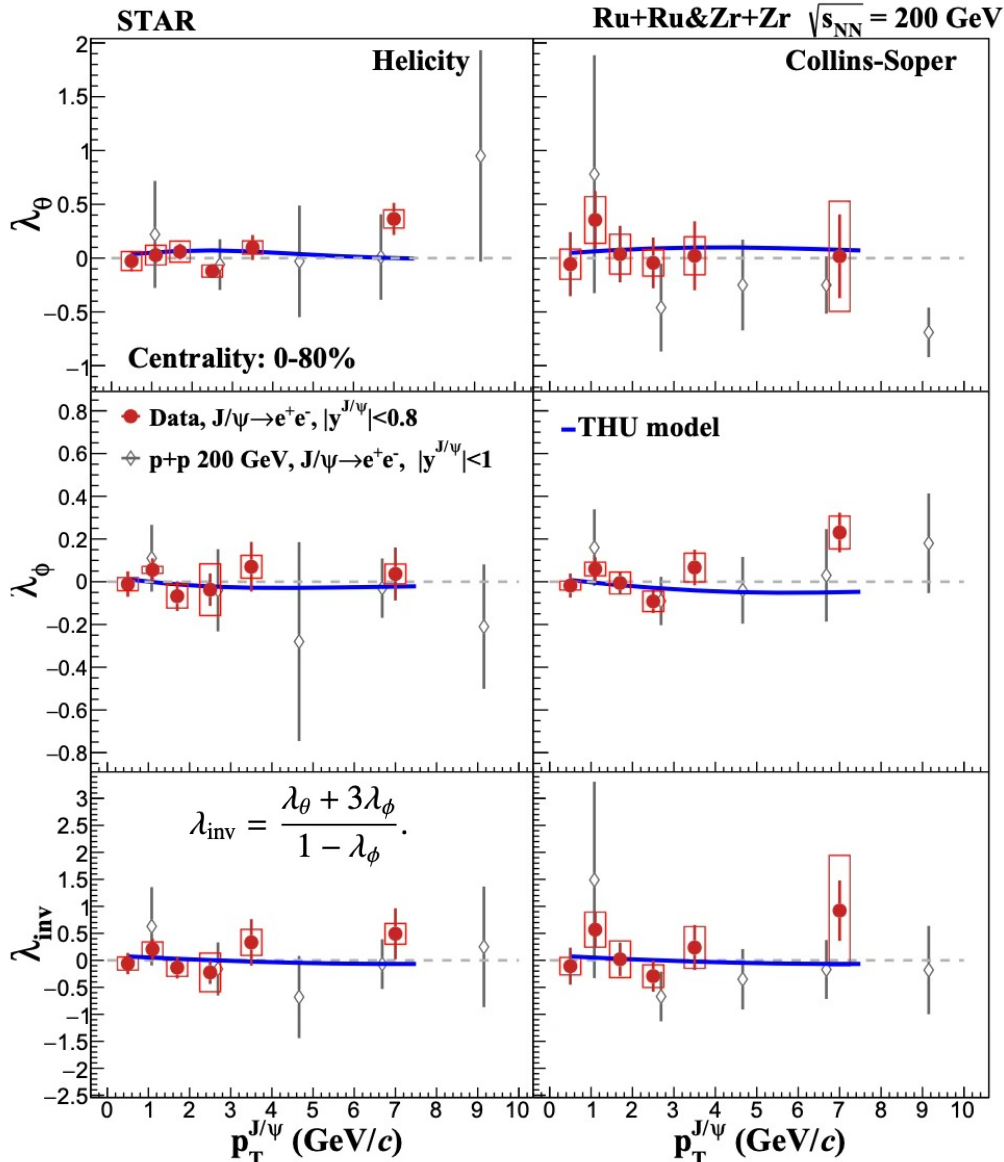


Measurable consequences of vector meson spin alignment



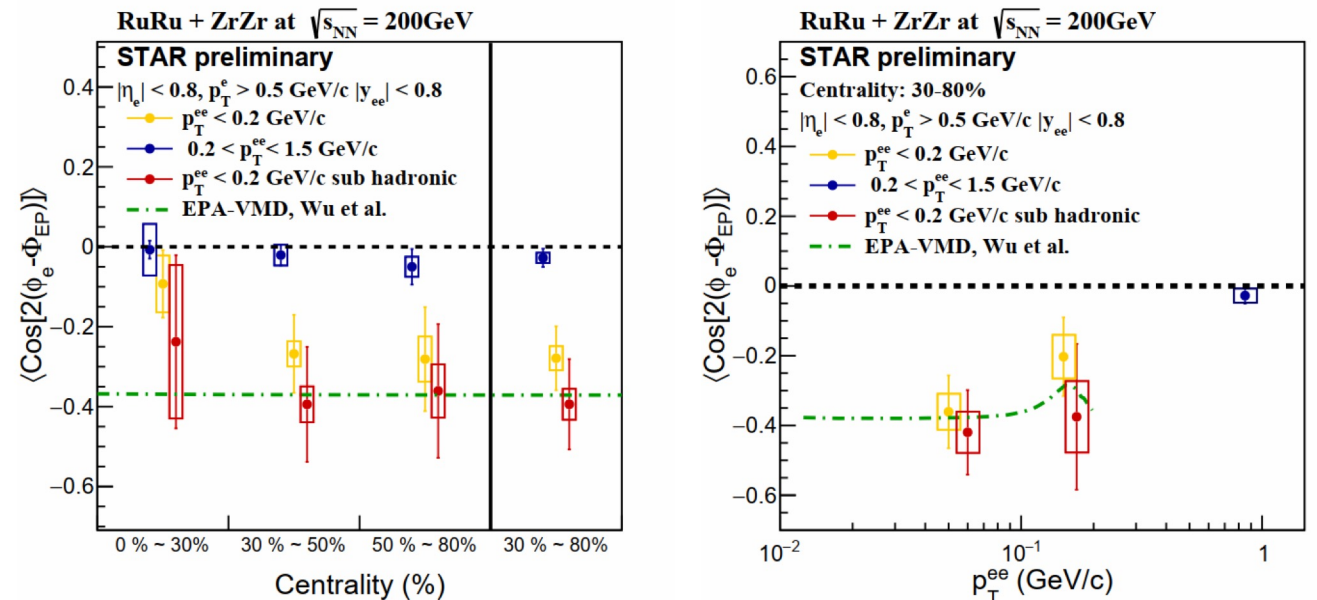
Spin alignment of charm mesons

STAR Col., arXiv:2604.07005

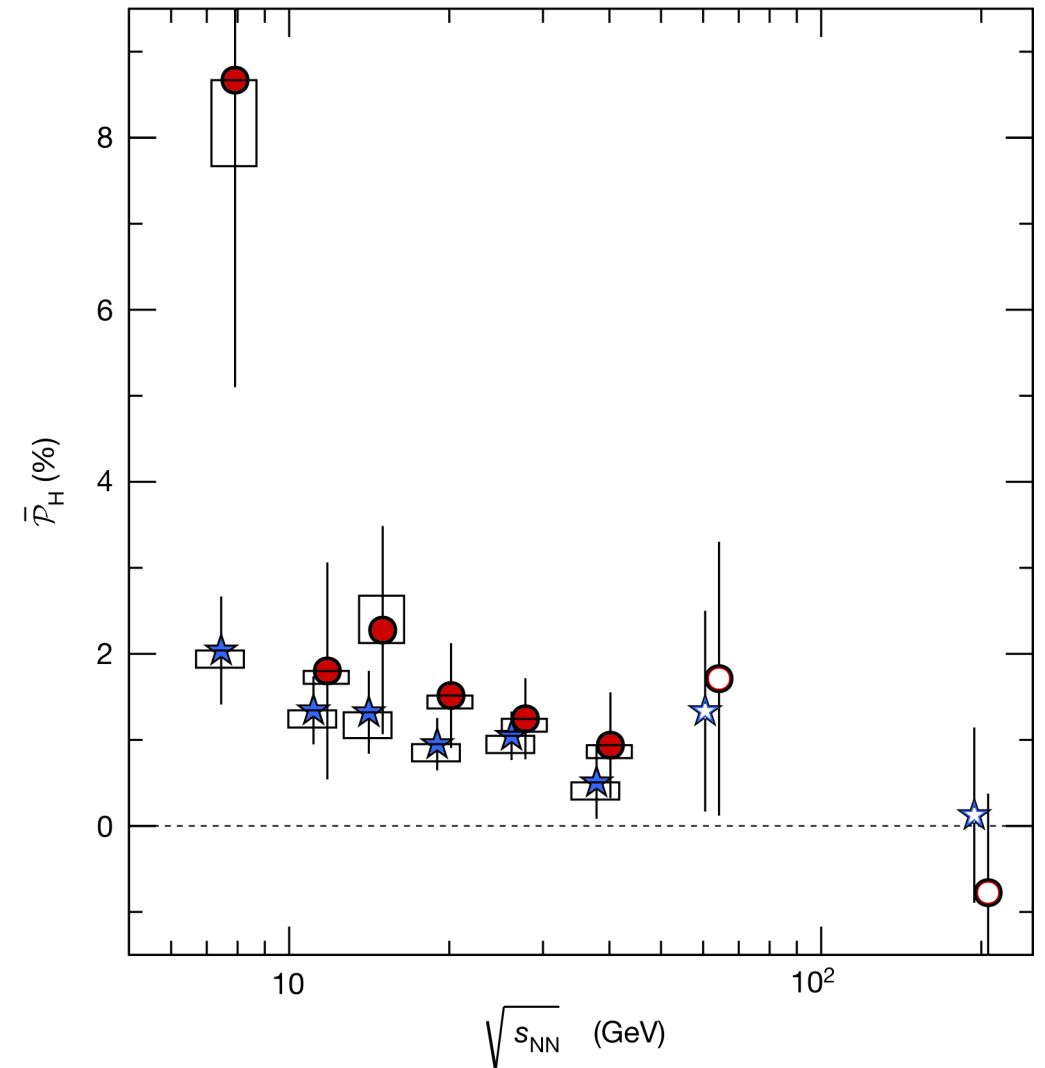
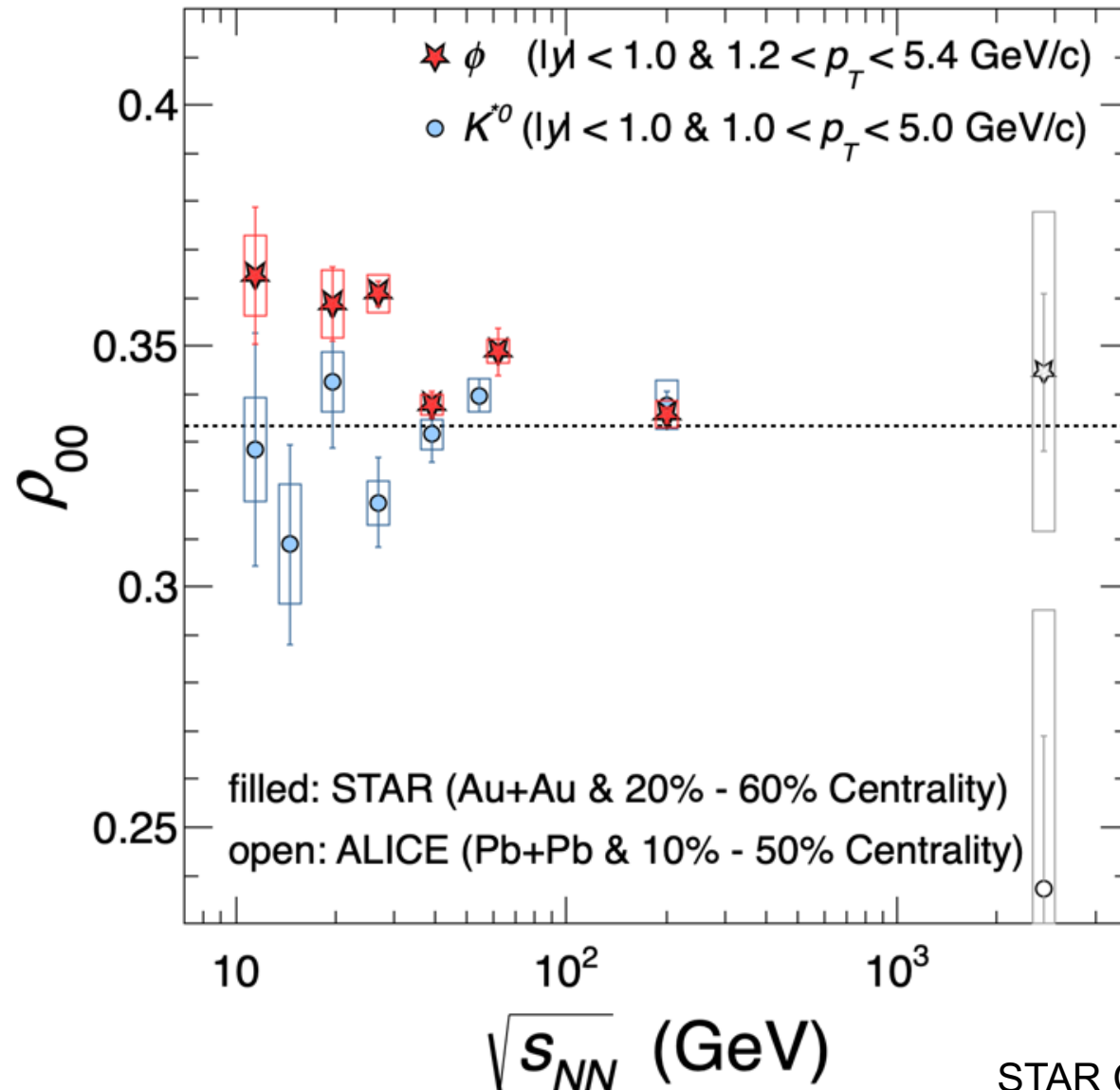


- STAR measured the J/ψ polarization in HIC and pp collisions
 - Polarization are consistent with zero within uncertainties.
 - Results in A+A collisions and p+p collisions are similar.
- From hadronic interaction to photon-induced dominant
Photon induced J/ψ present evidence of decay anisotropy
 - accessing impact parameter direction with photon-induced process

Kaiyang Wang, QM25



Large ϕ -meson alignment vs. Small Λ pol.



STAR Col. Nature **548**, 62 (2017), Nature **614**, 244 (2023)

Large ϕ -meson alignment vs. Small Λ pol. (cont.)

Liang, Chirality 2023

$$\left. \begin{aligned} \left| \rho_{00}^V - \frac{1}{3} \right| &\gg P_\Lambda^2 \sim P_q^2 \\ \rho_{00}^V - \frac{1}{3} &\sim \langle P_q P_{\bar{q}} \rangle \end{aligned} \right\} \text{The STAR data show that: } \langle P_q P_{\bar{q}} \rangle \neq \langle P_q \rangle \langle P_{\bar{q}} \rangle \quad \langle P_q P_{\bar{q}} \rangle \gg \langle P_q \rangle \langle P_{\bar{q}} \rangle$$

By studying P_H , we study the **average** of quark polarization P_q ;
by studying ρ_{00}^V , we study the **correlation** between P_q and $P_{\bar{q}}$.

How to separate long range or local correlations

$$C_{NN}^{H_i \bar{H}_j} \equiv \frac{N_{H_i \bar{H}_j}^{\uparrow\uparrow} + N_{H_i \bar{H}_j}^{\downarrow\downarrow} - N_{H_i \bar{H}_j}^{\uparrow\downarrow} - N_{H_i \bar{H}_j}^{\downarrow\uparrow}}{N_{H_i \bar{H}_j}^{\uparrow\uparrow} + N_{H_i \bar{H}_j}^{\downarrow\downarrow} + N_{H_i \bar{H}_j}^{\uparrow\downarrow} + N_{H_i \bar{H}_j}^{\downarrow\uparrow}}$$

sensitive to the **long range** correlation

$$\rho_{10}^V = \frac{P_{qz}(1 + P_{\bar{q}y}) + (1 + P_{qy})P_{\bar{q}z} - iP_{qx}(1 + P_{\bar{q}y}) - i(1 + P_{qy})P_{\bar{q}x}}{\sqrt{2}(3 + \vec{P}_q \cdot \vec{P}_{\bar{q}})}$$

$$\rho_{0-1}^V = \frac{P_{qz}(1 - P_{\bar{q}y}) + (1 - P_{qy})P_{\bar{q}z} - iP_{qx}(1 - P_{\bar{q}y}) - i(1 - P_{qy})P_{\bar{q}x}}{\sqrt{2}(3 + \vec{P}_q \cdot \vec{P}_{\bar{q}})}$$

$$\rho_{1-1}^V = \frac{P_{qz}P_{\bar{q}z} - P_{qx}P_{\bar{q}x} + i(P_{qx}P_{\bar{q}y} + P_{qy}P_{\bar{q}x})}{3 + \vec{P}_q \cdot \vec{P}_{\bar{q}}}$$

They should be sensitive to the **local** correlations.

Spin-spin correlation

Lv, Yu, Liang, Wang, Wang, PRD **109**, 114003 (2024); Zhang, Lv, Yu, Liang, PRD **110**, 074019 (2024)

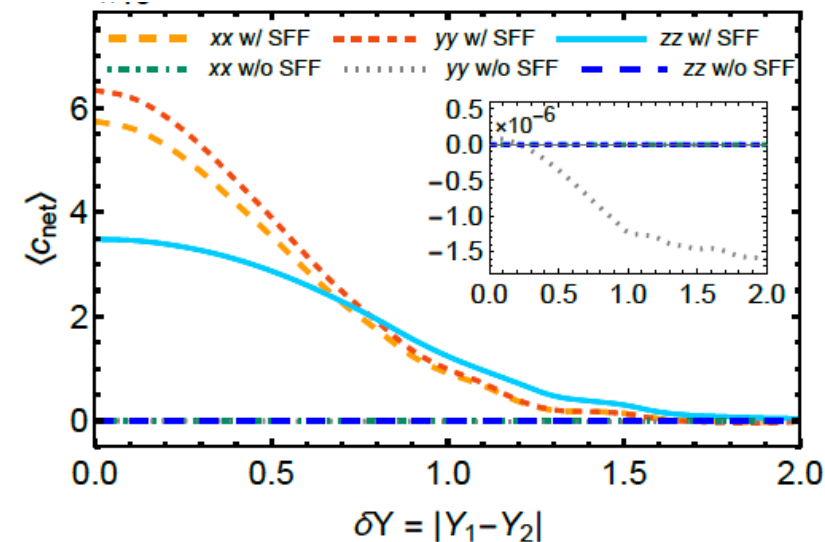
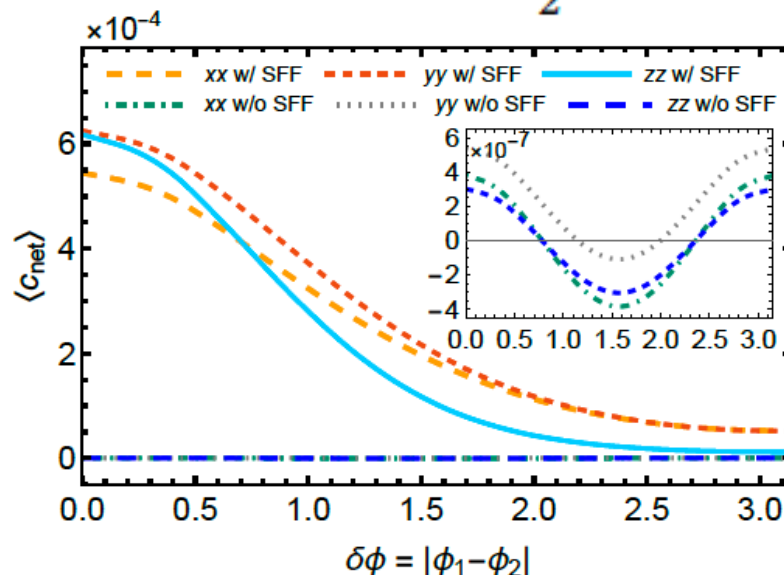
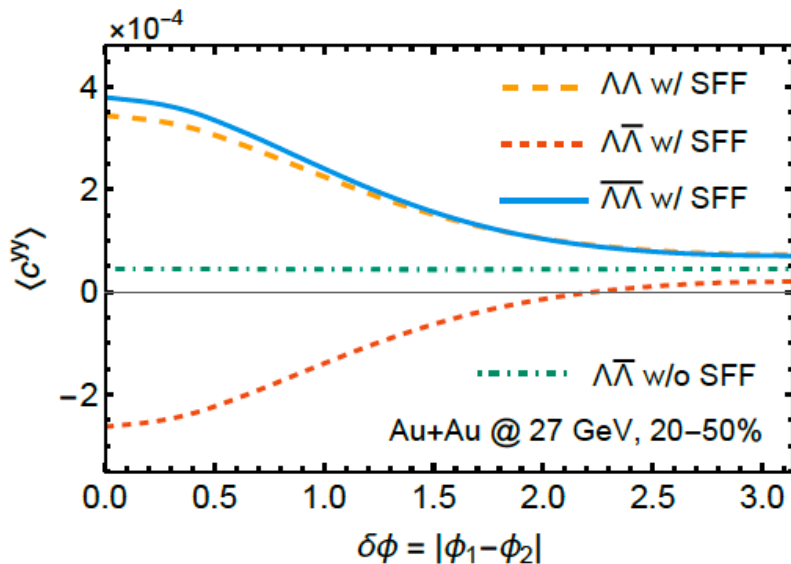
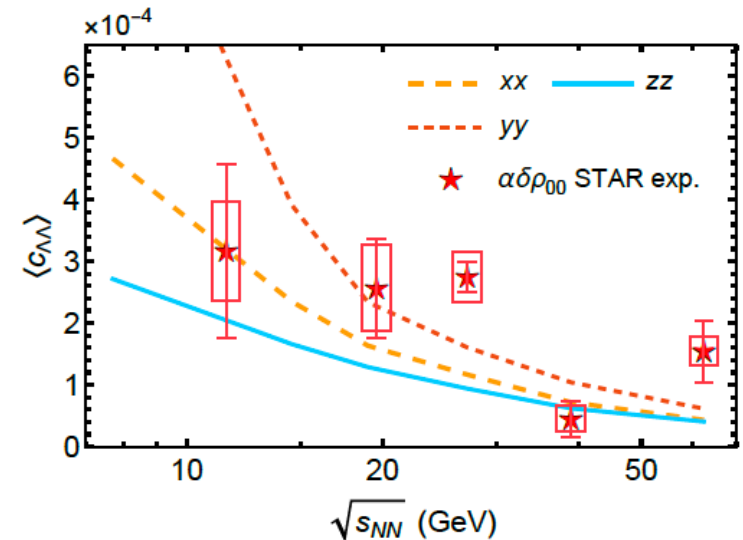
Sheng, Wu, Rischke, Wang, Phys. Rev. Lett. **136**, 082301 (2026)...

$$C_{12}^{\mu\nu}(p_1, p_2) \equiv \langle P_1^\mu(x_1, p_1) P_2^\nu(x_2, p_2) \rangle, \quad (1)$$

$$C_{12}^{\mu\nu}(p_1, p_2) = \frac{1}{N_{\text{event}}} \sum_{\text{event}} \times \frac{\int d\Sigma(x_1) \cdot p_1 \int d\Sigma(x_2) \cdot p_2 P_1^\mu P_2^\nu f_1 f_2}{\left[\int d\Sigma(x_1) \cdot p_1 f_1 \right] \left[\int d\Sigma(x_2) \cdot p_2 f_2 \right]} \quad (8)$$

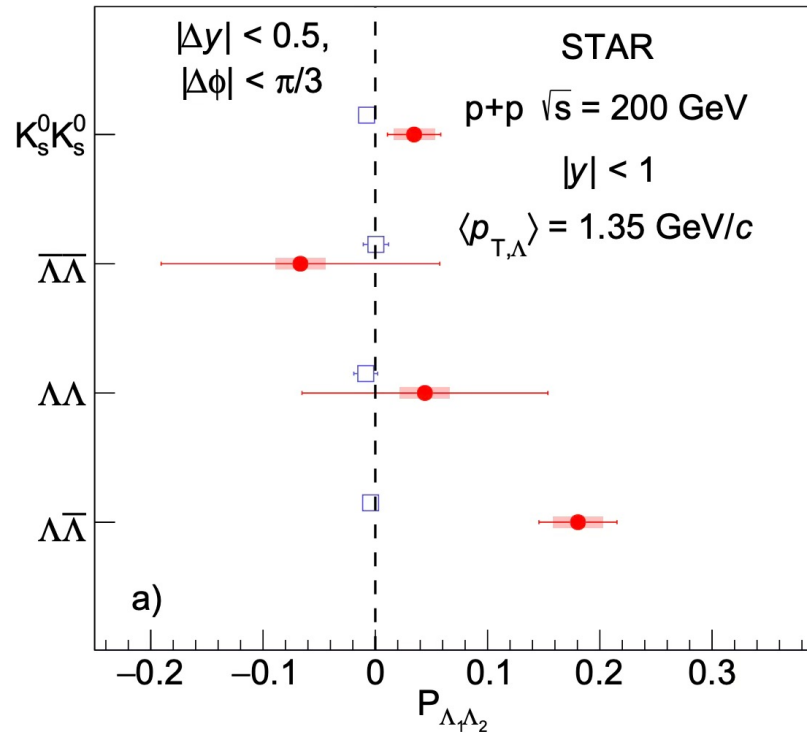
‘A key unique feature of the hyperon spin correlation caused by the SFF is that the spin correlation of $\Lambda\Lambda$ and anti- Λ anti- Λ has the opposite sign as the Λ anti- Λ spin correlation.’

$$c_{\text{net}}^{ab} \equiv \frac{1}{2}(c_{\Lambda\Lambda}^{ab} + c_{\bar{\Lambda}\bar{\Lambda}}^{ab}) - c_{\Lambda\bar{\Lambda}}^{ab}, \quad (12)$$

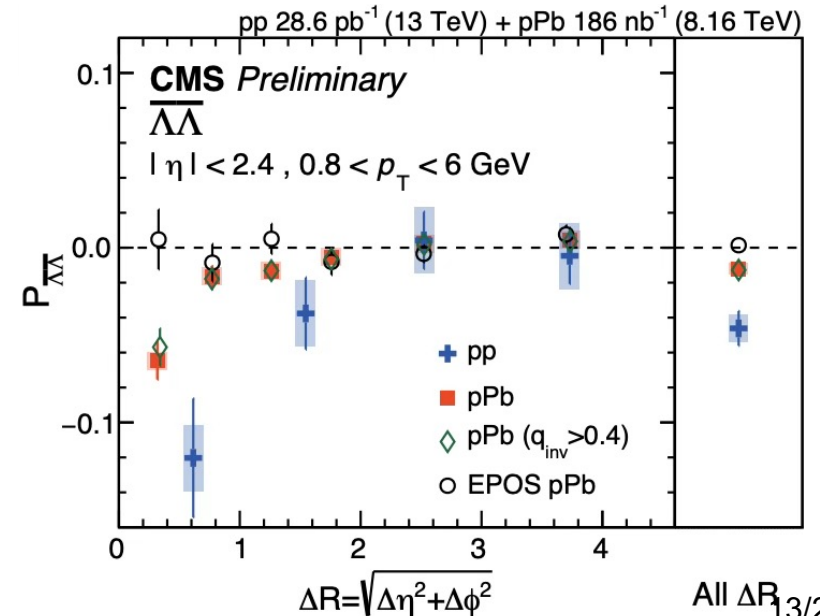
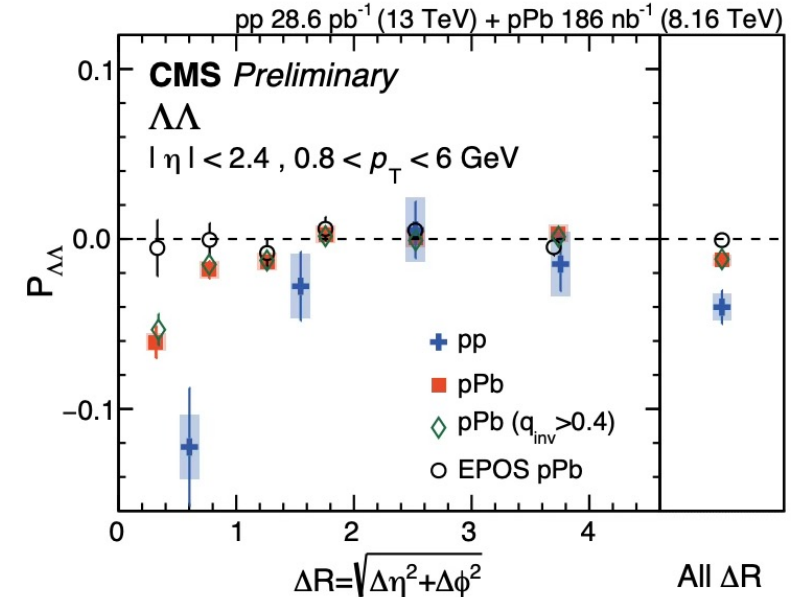
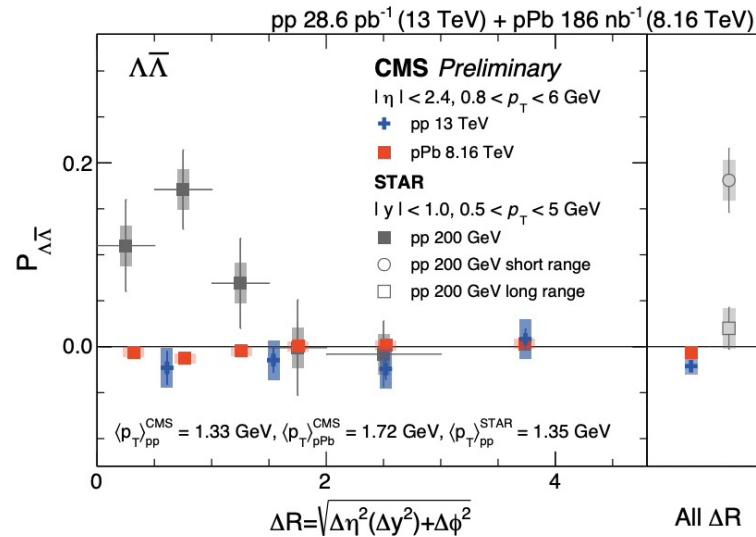


Measure the spin-spin correlation

$$\frac{dN}{d\cos(\theta^*)} \sim 1 + \alpha_1 \alpha_2 P_{\Lambda_1 \Lambda_2} \cos(\theta^*)$$



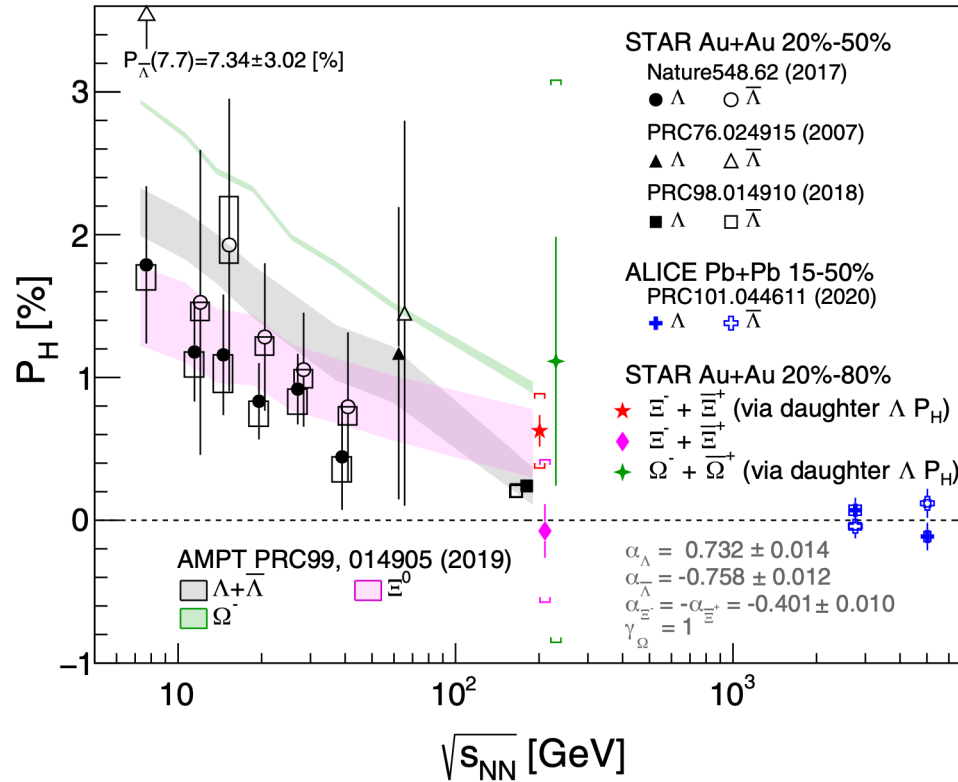
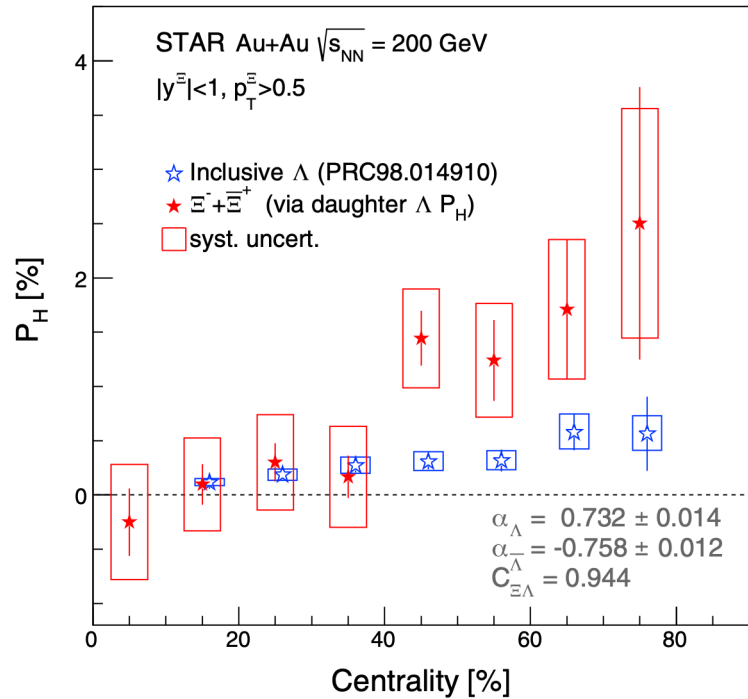
Jieke Wang DIS 2026



- Lambda spin correlation has been measured in small systems.
 - p+p 200 GeV: $P_{OS} > 0$, $P_{SS} \sim 0$
 - p+p 13 TeV and p+Pb 8.16 TeV: $P_{OS} \sim 0$, $P_{SS} < 0$

Measurements of spin 3/2 baryon

STAR Col. Phys. Rev. Lett. **126**, 162301 (2021)



- Xi data confirm the global polarization picture based on the system fluid vorticity
- New data agree with AMPT model prediction

• spin 1/2

$$\mathbf{P}_{\Lambda}^* = \frac{(\alpha_{\Xi} + \mathbf{P}_{\Xi}^* \cdot \hat{\mathbf{p}}_{\Lambda}^*)\hat{\mathbf{p}}_{\Lambda}^* + \beta_{\Xi}\mathbf{P}_{\Xi}^* \times \hat{\mathbf{p}}_{\Lambda}^* + \gamma_{\Xi}\hat{\mathbf{p}}_{\Lambda}^* \times (\mathbf{P}_{\Xi}^* \times \hat{\mathbf{p}}_{\Lambda}^*)}{1 + \alpha_{\Xi}\mathbf{P}_{\Xi}^* \cdot \hat{\mathbf{p}}_{\Lambda}^*}$$

$$\mathbf{P}_{\Lambda}^* = C_{\Xi-\Lambda}\mathbf{P}_{\Xi}^* = \frac{1}{3}(1 + 2\gamma_{\Xi})\mathbf{P}_{\Xi}^* \quad C_{\Xi-\Lambda} = \frac{1}{3}(1 + 2 \times 0.916) = +0.944$$

T.D. Lee and C.-N. Yang, Phys. Rev. **108**, 1645 (1957)

• spin 3/2

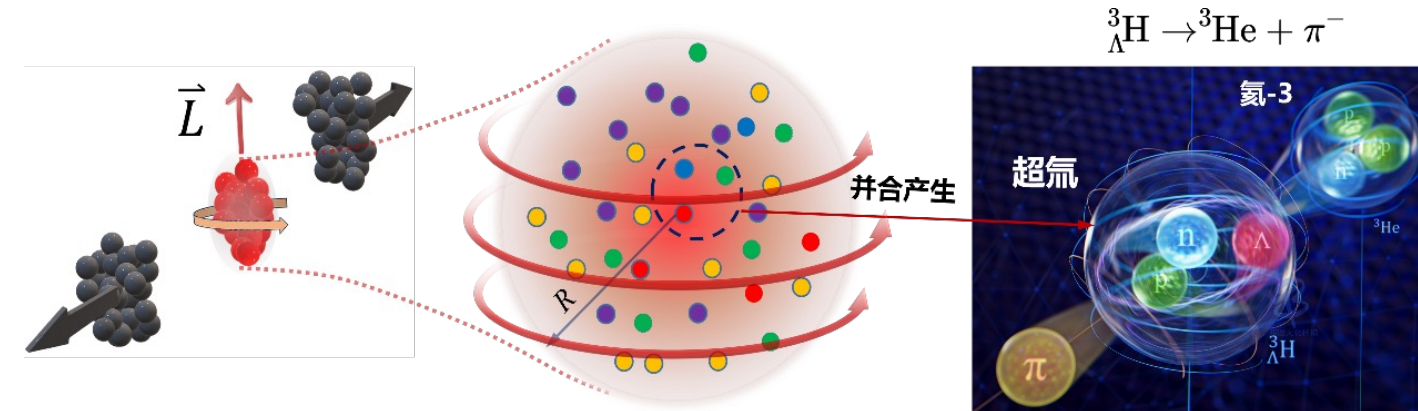
$$\mathbf{P}_{\Lambda}^* = C_{\Omega-\Lambda}\mathbf{P}_{\Omega}^* = \frac{1}{5}(1 + 4\gamma_{\Omega})\mathbf{P}_{\Omega}^*$$

$$C_{\Omega-\Lambda} \approx 1 \quad C_{\Omega-\Lambda} \approx -0.6$$

*different component?

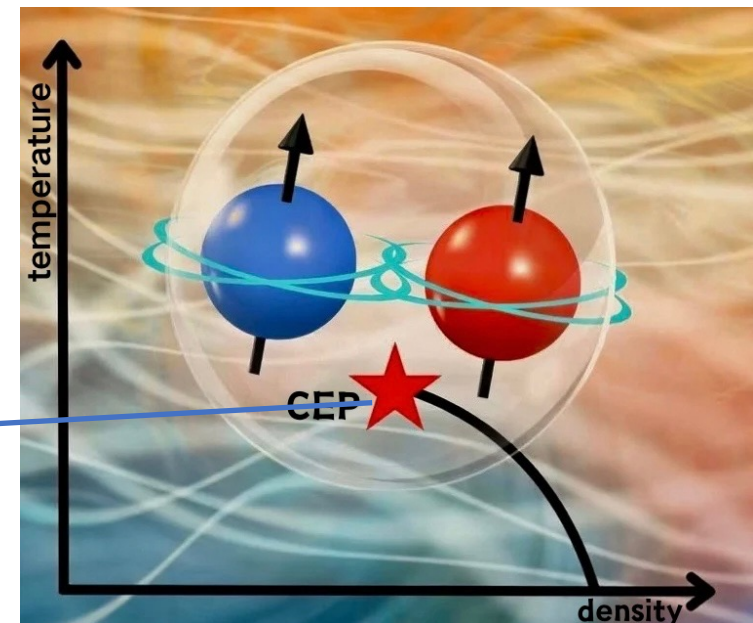
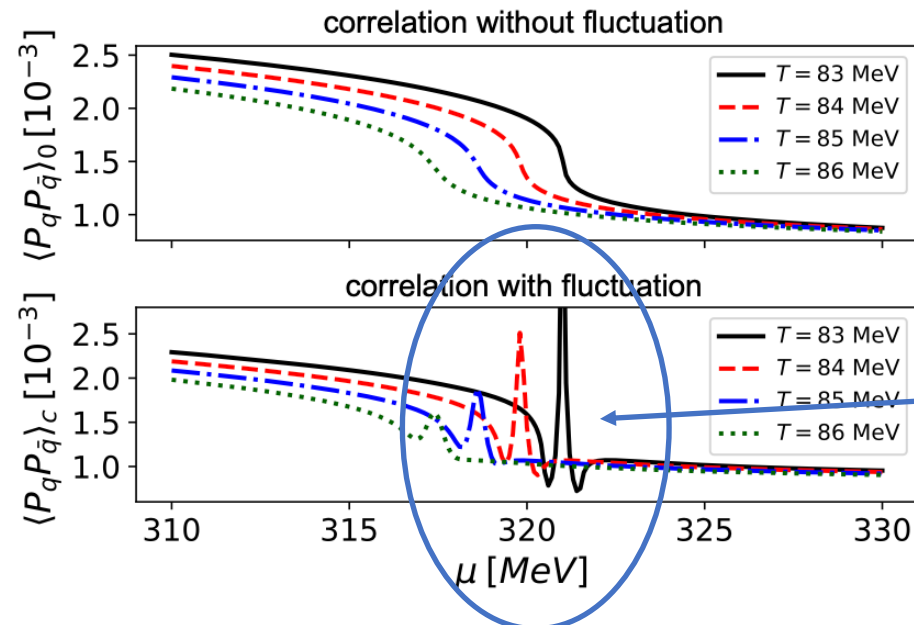
Spin polarization of light cluster and correlation

Sun, Liu, Zheng, Chen, Ko and Ma, Phys. Rev. Lett. **134**, 022301 (2025), Phys. Rev. C **113**, 054906 (2026), Phys. Rev. Res. **8**, 033057 (2026)...

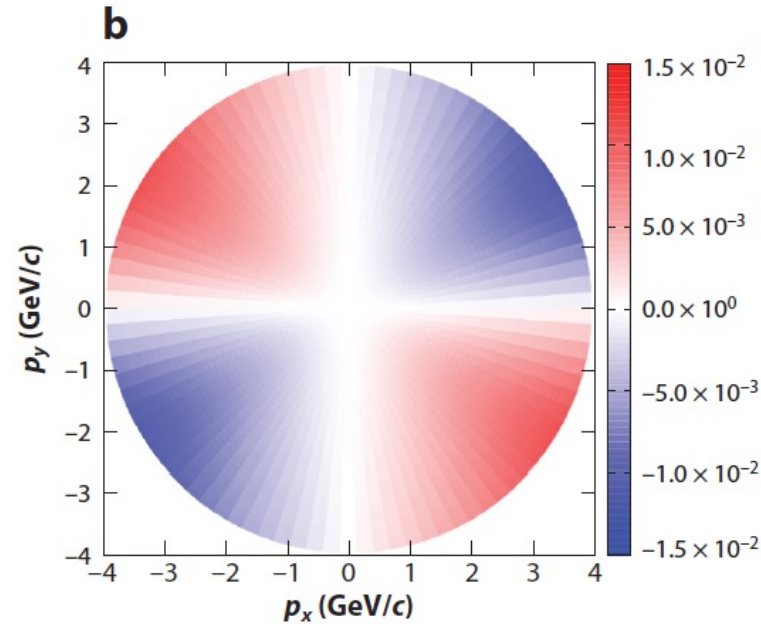
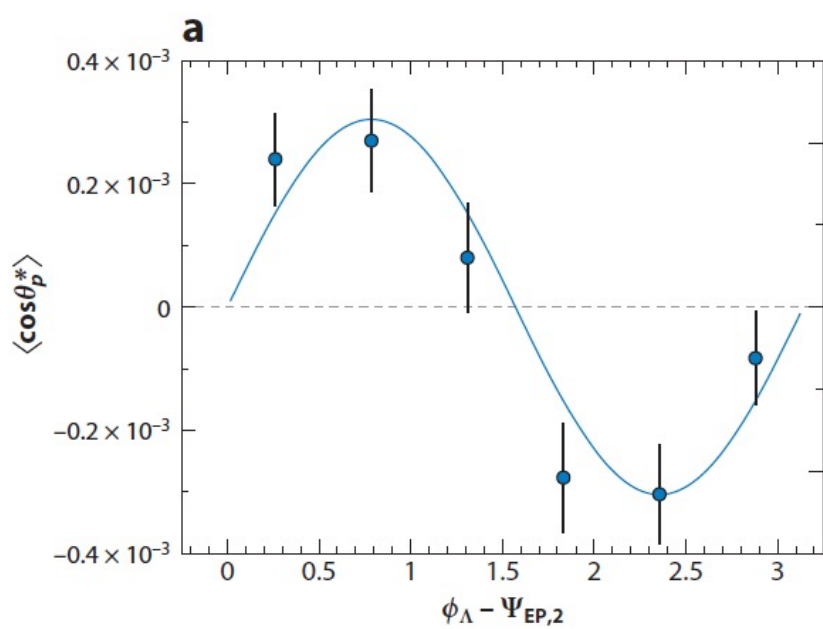


J^π	Structure	Decay mode	$dN/(\sin\theta^* d\theta^*)$
$\frac{1}{2}^+$	$\Lambda(\frac{1}{2}^+) - np(1^+)$	${}^3_\Lambda\text{H} \rightarrow \pi^- + {}^3\text{He}$	$\frac{1}{2}[1 - (1/2.58)\alpha_\Lambda \mathcal{P}_\Lambda \cos\theta^*]$
$\frac{1}{2}^+$	$\Lambda(\frac{1}{2}^+) - np(0^+)$	${}^3_\Lambda\text{H} \rightarrow \pi^- + {}^3\text{He}$	$\frac{1}{2}(1 + \alpha_\Lambda \mathcal{P}_\Lambda \cos\theta^*)$
$\frac{3}{2}^+$	$\Lambda(\frac{1}{2}^+) - np(1^+)$	${}^3_\Lambda\text{H} \rightarrow \pi^- + {}^3\text{He}$	$\frac{1}{2}\left(1 - \mathcal{P}_\Lambda^2(3\cos^2\theta^* - 1)\right)$
$\frac{1}{2}^-$	$\bar{\Lambda}(\frac{1}{2}^-) - \bar{n}\bar{p}(1^+)$	${}^3_{\bar{\Lambda}}\bar{\text{H}} \rightarrow \pi^+ + {}^3\bar{\text{He}}$	$\frac{1}{2}[1 - (1/2.58)\alpha_{\bar{\Lambda}} \mathcal{P}_{\bar{\Lambda}} \cos\theta^*]$
$\frac{1}{2}^-$	$\bar{\Lambda}(\frac{1}{2}^-) - \bar{n}\bar{p}(0^+)$	${}^3_{\bar{\Lambda}}\bar{\text{H}} \rightarrow \pi^+ + {}^3\bar{\text{He}}$	$\frac{1}{2}(1 + \alpha_{\bar{\Lambda}} \mathcal{P}_{\bar{\Lambda}} \cos\theta^*)$
$\frac{3}{2}^-$	$\bar{\Lambda}(\frac{1}{2}^-) - \bar{n}\bar{p}(1^+)$	${}^3_{\bar{\Lambda}}\bar{\text{H}} \rightarrow \pi^+ + {}^3\bar{\text{He}}$	$\frac{1}{2}\left(1 - \mathcal{P}_{\bar{\Lambda}}^2(3\cos^2\theta^* - 1)\right)$

Chen, Fu, Huang and Ma, Phys. Rev. Lett. **135**, 032302 (2025)

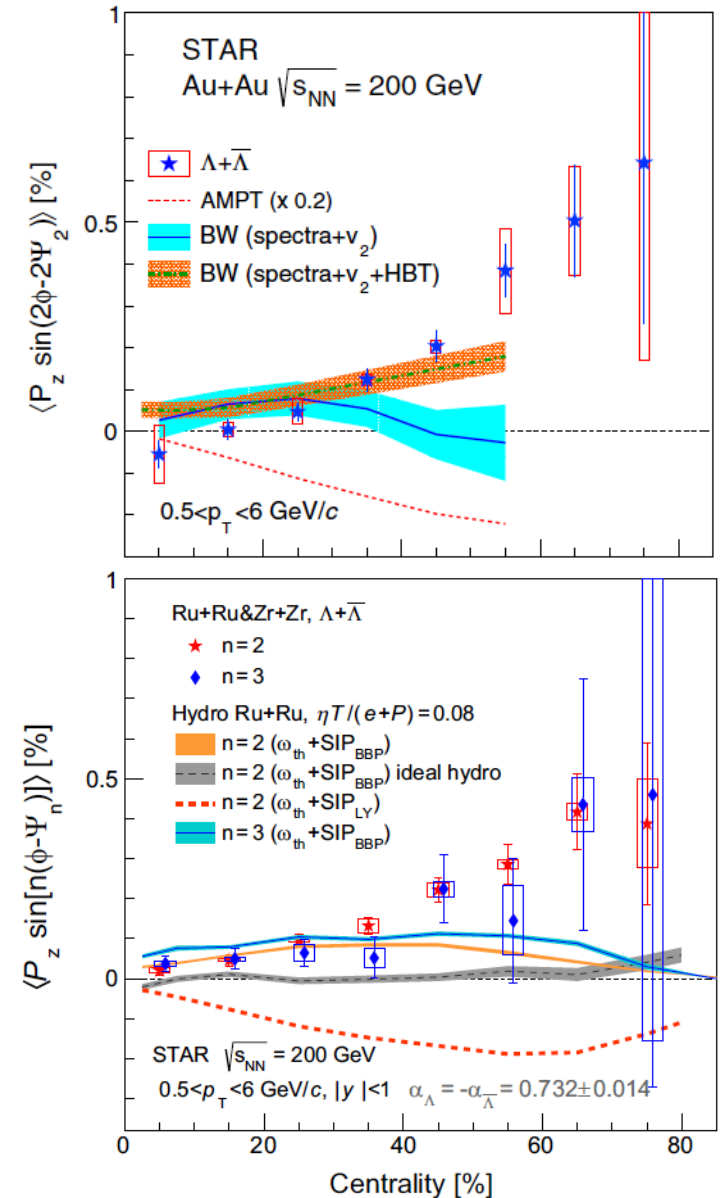


Λ polarization along beam direction



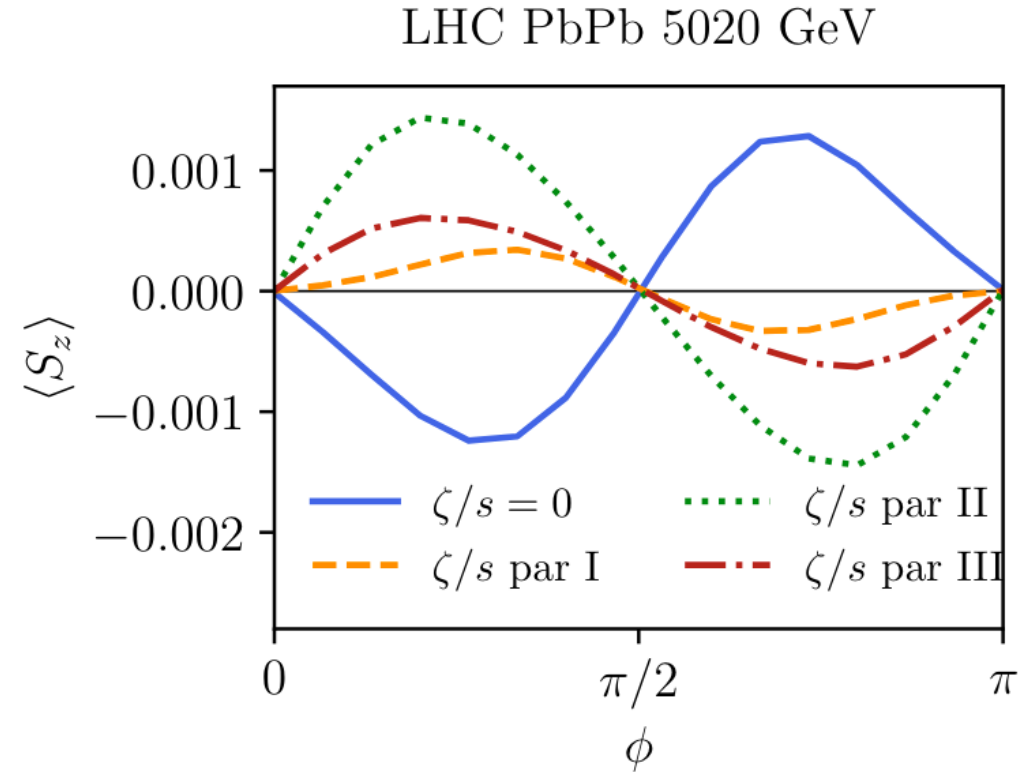
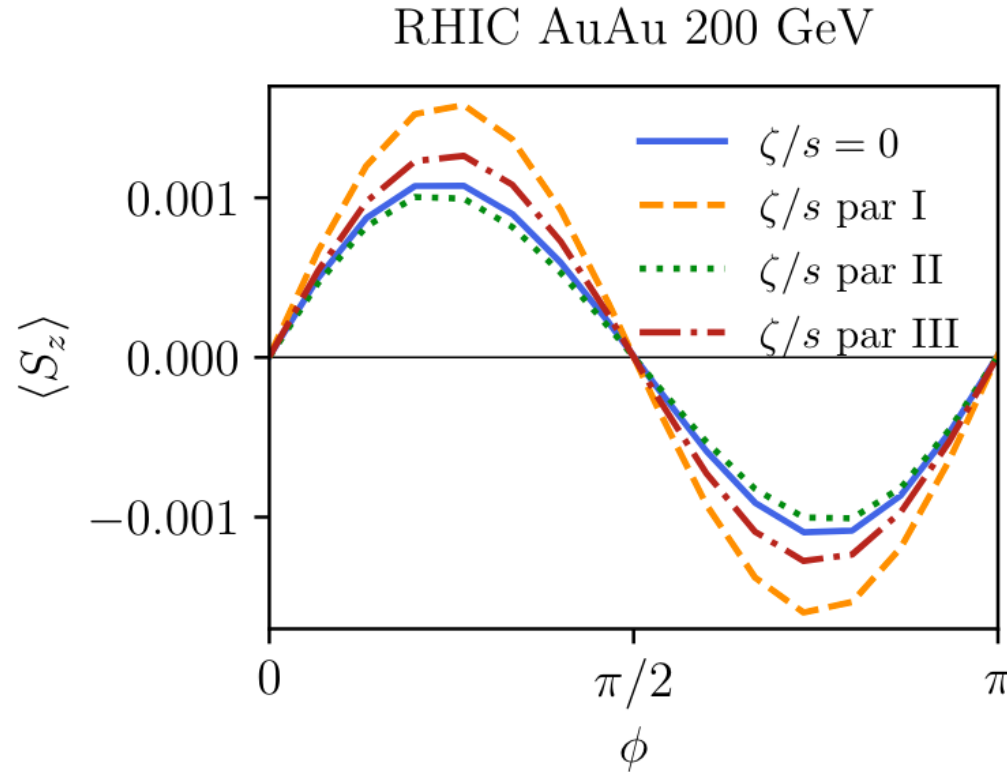
STAR Col., Phys. Rev. Lett. **123**, 132301 (2019), Phys. Rev. Lett. **131**, 202301 (2023),
ALICE Col., Phys. Rev. Lett. **128**, 172005 (2022), CMS Col., Phys. Rev. Lett. **135**, 132301 (2025).
Becattini, Lisa, Annu. Rev. Nucl. Part. Sci. 70, 395 (2020)...

- A quadrupole modulation of the polarization along the beam direction is observed. It is qualitatively consistent with the expectation from the vorticity component along the beam direction due to collective flow.
- ‘Spin sign puzzle’: data present a positive sine modulation, while hydrodynamics and transport model calculations predict a negative sign.



Polarization is sensitive to bulk viscosity

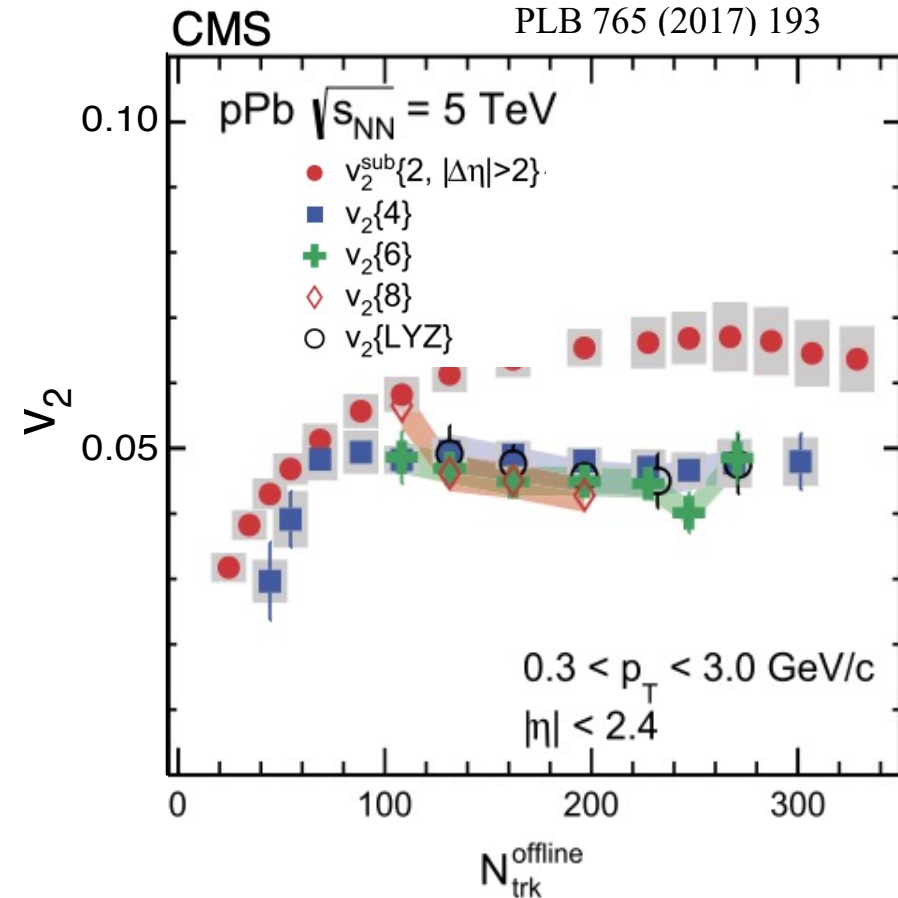
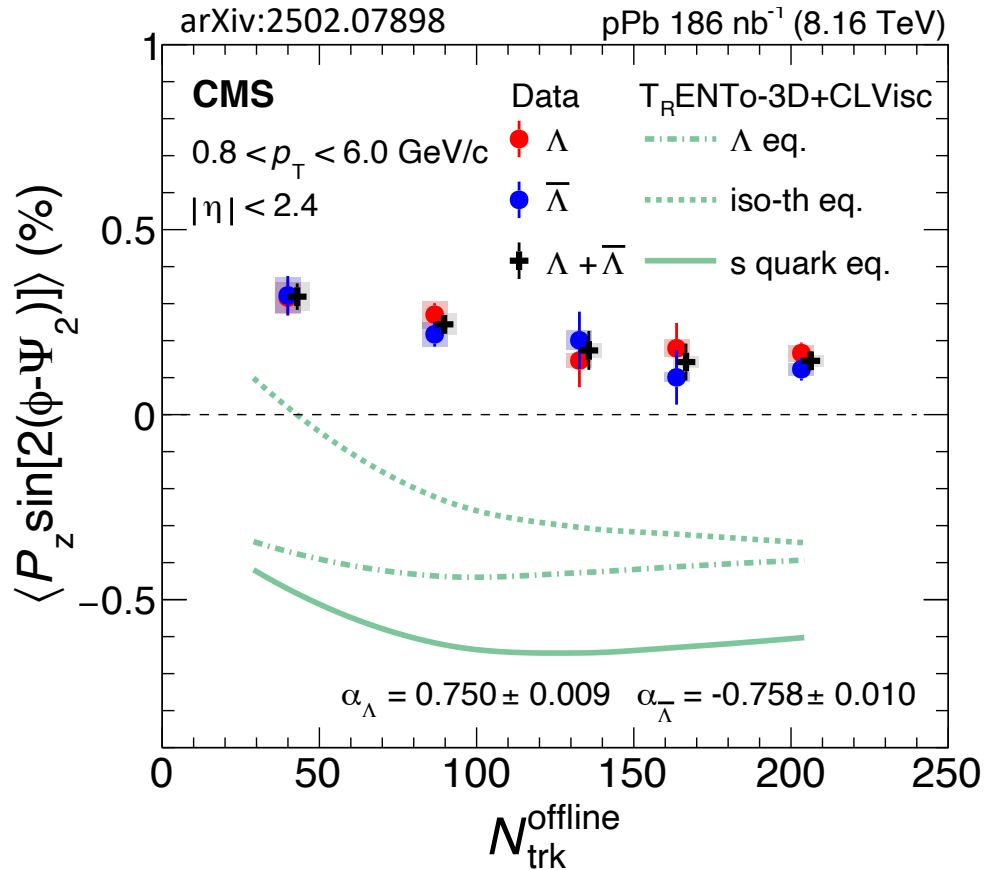
Palermo, Grossi, Karpenko, Becattini, Eur. Phys. J. C **84**, 920 (2024)



- The bulk viscosity significantly affects local polarization
- Once bulk viscosity is implemented, the sign of polarization is flipping in LHC energy (LHC more sensitive)
- Challenging to accurately extract the viscosity from $\langle S_z \rangle$

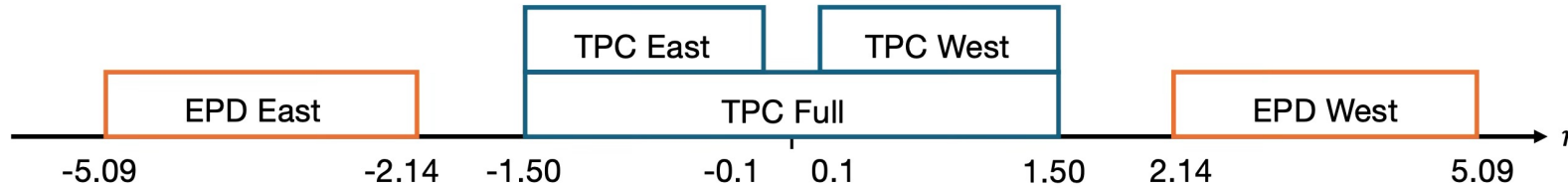
Measurements at LHC energies

CMS Col. Phys. Rev. Lett. **135**, 132301 (2025)

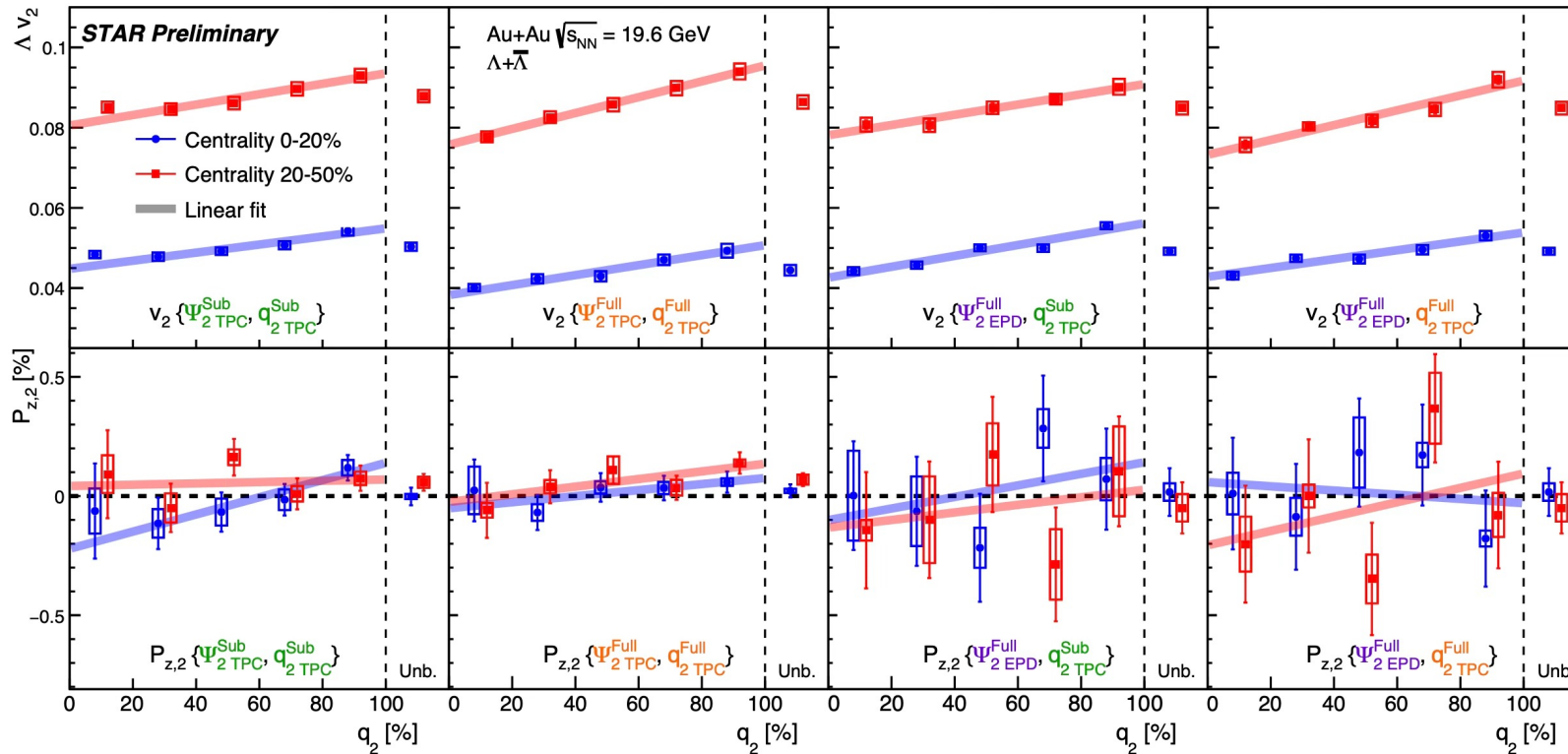


- Local polarization was observed in small system. Hydro calculations is not consistent with data
- Polarization increases monotonically toward 0 multiplicity, not consistent with the trend of v_2 .
- In large system, apply ESE to investigate any local polarization dependence? In small system, flow + initial state effect?

Local polarization with ESE



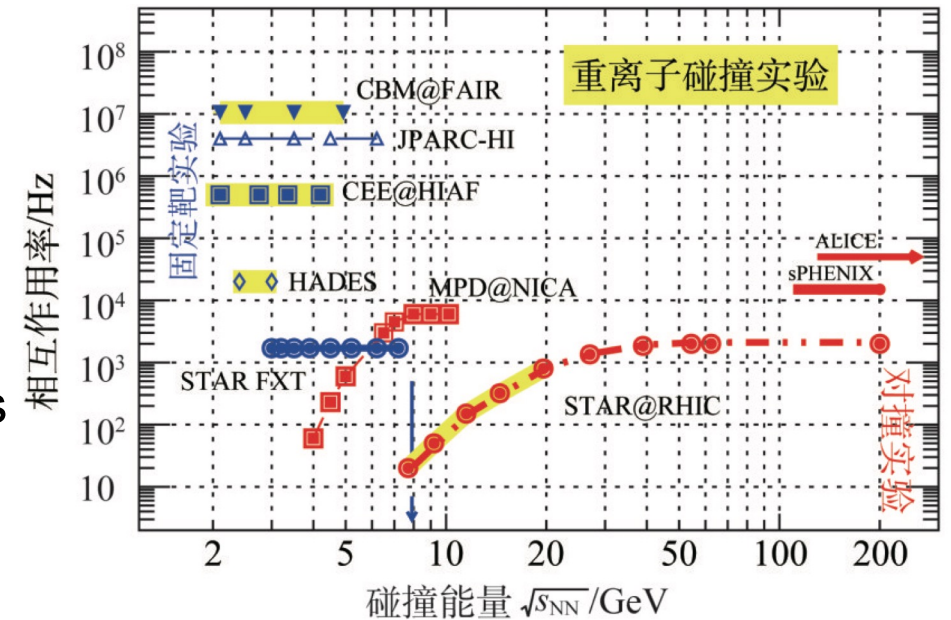
Taiki Kondo SQM 2026, Zhiyi, Diyu...



- Selecting different initial geometry and elliptic flow via event shape engineering at a given centrality.
- Making flow the only variable and study its relationship with local polarization.
- No significant v_2 dependence of $P_{z,2}$ in Au+Au 19.6 GeV has been observed.

Summary and Outlook

- Spin polarization opens a new avenue of heavy-ion collisions physics
- Global hyperon polarization is observed with the order of a few percent. It represents a measure of the average value of the global quark polarization in the system
- Global vector meson spin alignment is observed with a surprisingly large pattern for ϕ -meson. It represents a local fluctuation/correlation between quarks polarization
- Local polarization show a profound structure, which is sensitive to the fluctuation of vorticity and bulk properties
- Measurements of spin-spin correlations are progressing well, rich physics to be explored



Thank you for your attentions!