

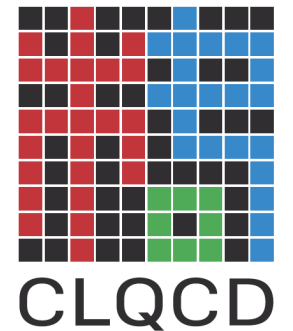
2026.07.28 The 1st Conference on Strong-Interaction Spin Physics

Intrinsic quark spin in octet baryons from 2+1 flavour LQCD at the physical point

Speaker: Zhi-Cheng Hu(胡志成)

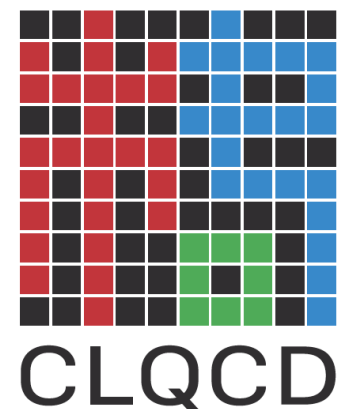


Institute of Theoretical Physics, CAS
Institute of Modern Physics, CAS



Collaborators: Liuming Liu, Peng Sun, Yi-Bo Yang, Long Chen

- ◆ Chapter 1: Background
- ◆ Chapter 2: Excited-State Contaminations on LQCD
- ◆ Chapter 3: Numerical result
- ◆ Chapter 4: Summary



◆ Chapter 1: Background

◆ Chapter 2: Excited-State Contaminations on LQCD

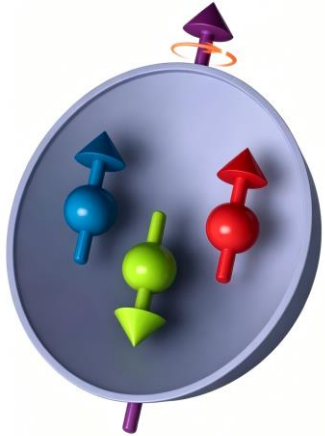
◆ Chapter 3: Numerical result

◆ Chapter 4: Summary



1. Background: The Origin of Proton Spin

Quark model



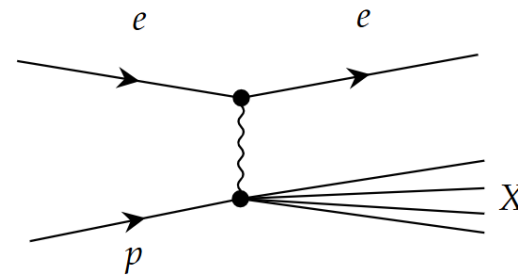
1980s

- ◆ Proton spin comes from sum of quarks.

$$S_p = \sum_q S_q = \frac{1}{2} \Sigma.$$

Proton spin crisis

DIS : $e + p \rightarrow e + X$



➤ 1989: $\frac{1}{2} \Sigma = 12(17)\% S_p$

EMC, PLB **206**,364(1988)

➤ Now: $\frac{1}{2} \Sigma \sim 30\% S_p$

Rept.Prog.Phys. **82** (2019) 076201

Spin Sum rules

- Jaffe-Manohar:

$$\frac{\hbar}{2} = \frac{1}{2} \Sigma + \Delta G + \ell_q + \ell_g$$

R.L. Jaffe et.al, NPB **337** (1990) 509-546

- Ji's sum rule:

$$\frac{\hbar}{2} = \frac{1}{2} \Sigma + L_q + J_g$$

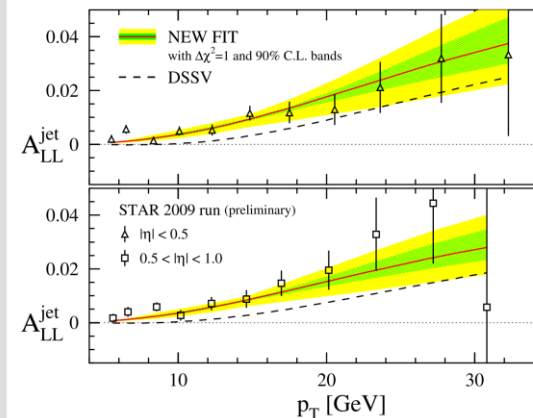
X.D. Ji, PRL **78**, 610(1997)

Constrain Gluon Spin

$p + p \rightarrow \text{jet} + X$

$$\int_{0.05}^1 dx \Delta g(x) = 0.20(7)$$

D.D. Florian, et.al, PRL **113**, 012001 (2014)
STAR, PRL **115**, 092002 (2015)

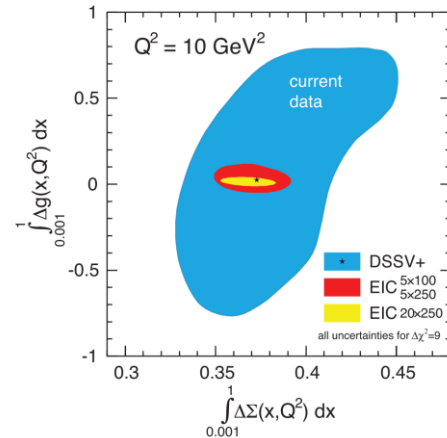
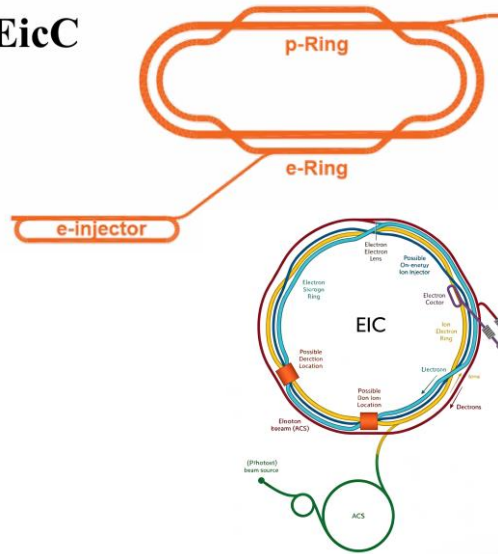


1. Background: The Origin of Proton Spin

- EIC physics in the future: precision on the total quark/gluon helicity contributions to the **proton spin**

$$\text{DIS} : \vec{e} + \vec{p} \rightarrow \vec{e} + X$$

EicC



NPA 1026 (2022) 122447
Front. Phys. 16, 64701 (2021)

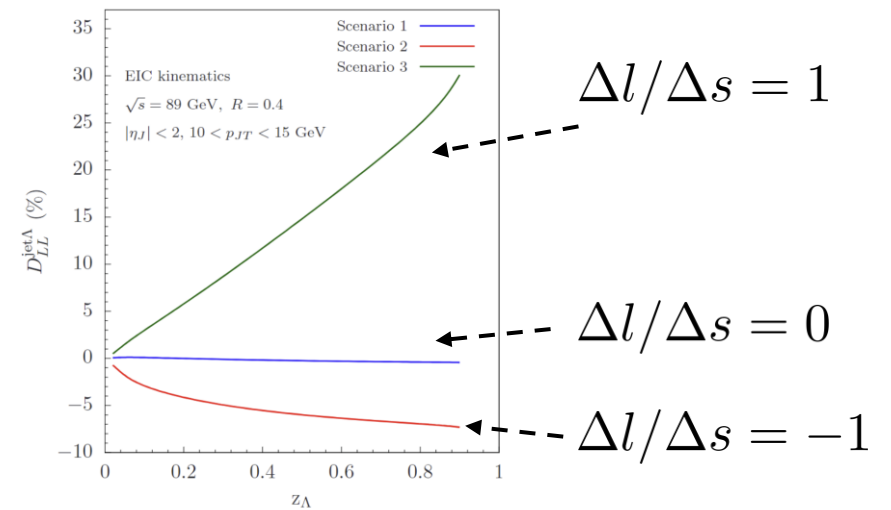
Higher precision for proton

- Λ spin in EIC:

$$\vec{p} + \vec{e} \rightarrow \text{jet}(\Lambda) + X$$

Λ polarization asymmetry:

$$D_{LL}^{\text{jet}\Lambda} = \frac{d\Delta\sigma^{\vec{p}\vec{p} \rightarrow (\text{jet}\vec{\Lambda}X)}}{dp_{JT}d\eta_J dZ_\Lambda} / \frac{d\sigma^{pp \rightarrow (\text{jet}\Lambda X)}}{dp_{JT}d\eta_J dZ_\Lambda}.$$



Z. B. Kang, PLB 809, 135756(2020)

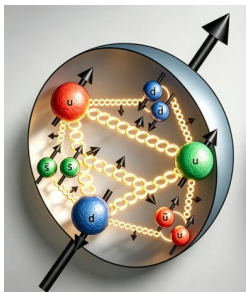
More baryons under study

1. Background: The Origin of Proton Spin



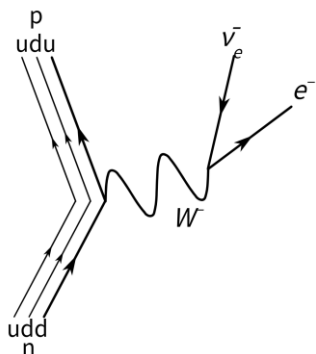
- Hadron matrix elements contain rich information about the hadronic structure;

- Proton spin puzzle

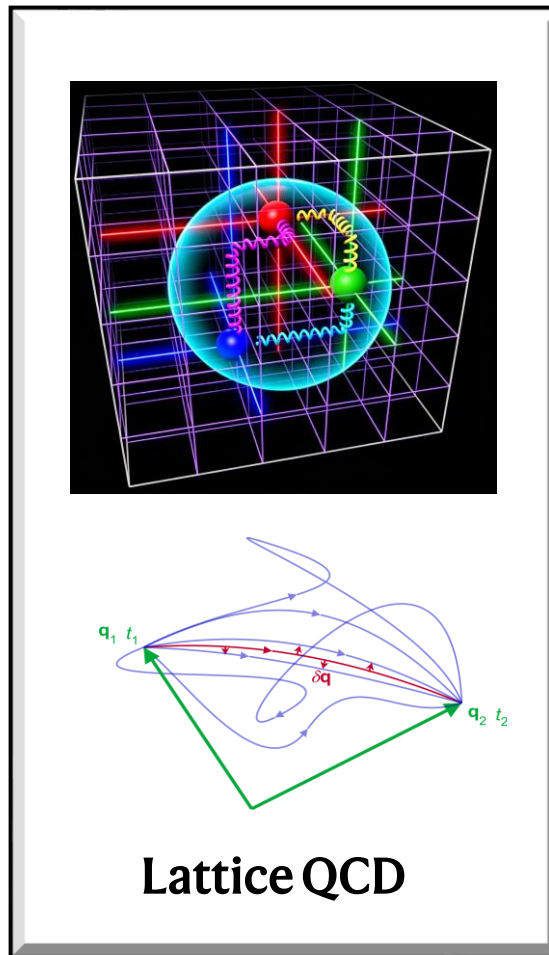


EMC, PLB 206,364(1988)
X.D. Ji, PRL 78, 610(1997)

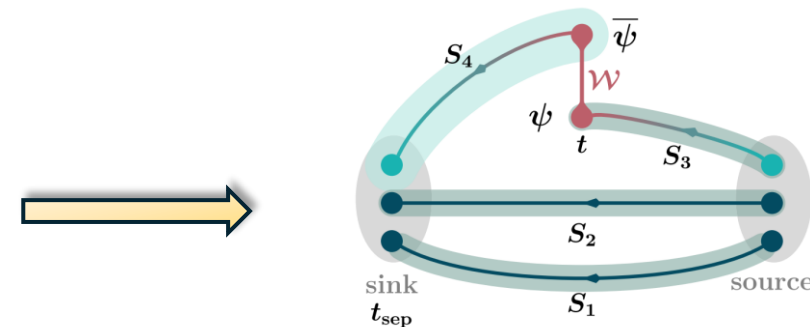
- Neutron beta decay



B. Markisch; et.al; PRL. 122, 242501 (2019)
CalLAT, Nature 558, 91(2018);



- Parton Structure (PDFs)



X.D. Ji, PRL 110, 262002(2013)
LPC, PRL 131.261901(2023)

- ✓ Lattice QCD provides a **first-principle, non-perturbative** description of the strong interaction.

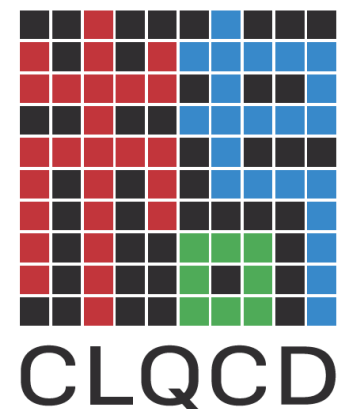
K. G. Wilson, PRD. 10, 8 (1974);

◆ Chapter 1: Background

◆ **Chapter 2: Excited-State Contaminations on LQCD**

◆ Chapter 3: Numerical result

◆ Chapter 4: Summary



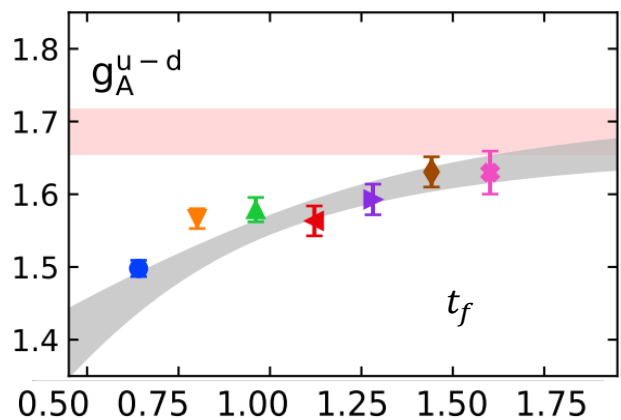
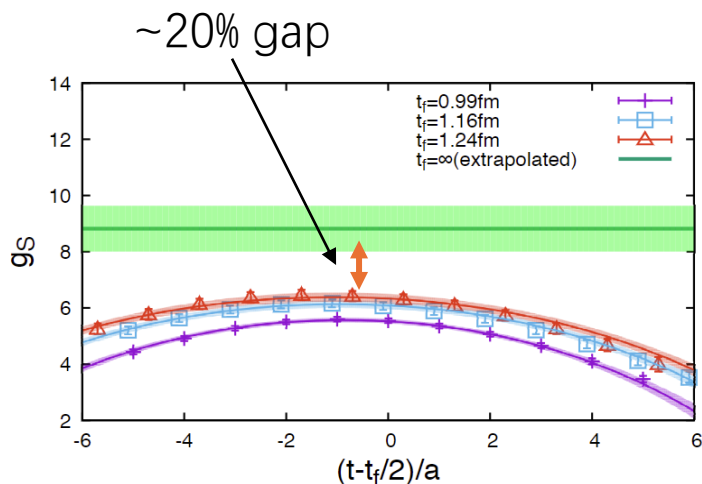
2. Excited-State Contaminations in LQCD



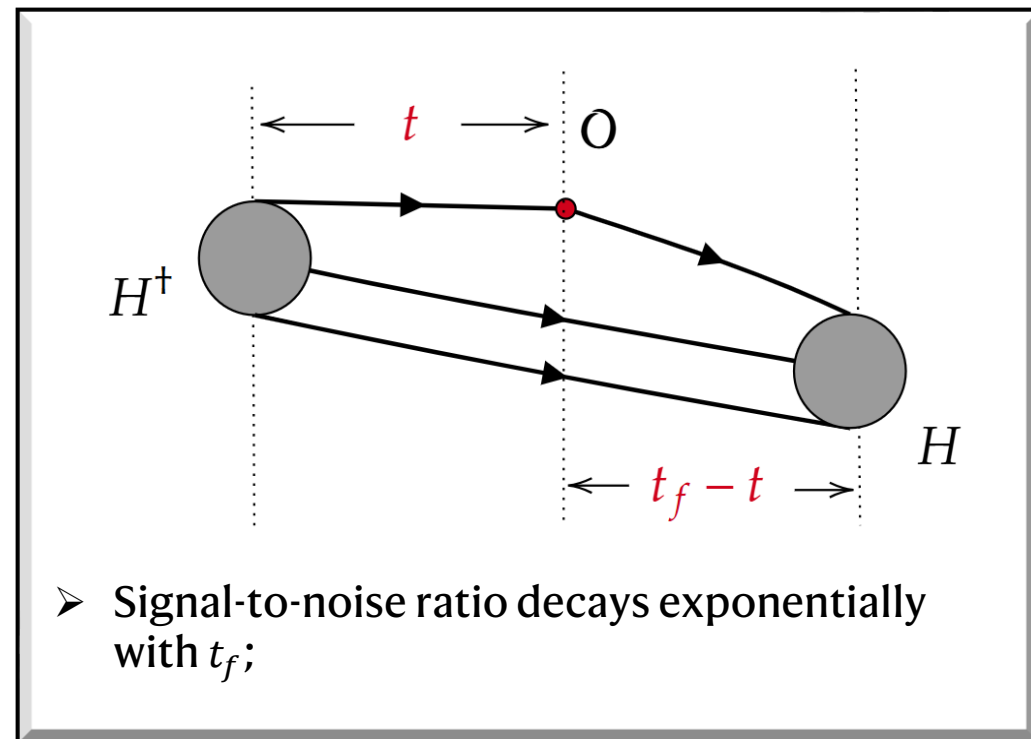
1. Imaginary-time evolution presents a serious challenge to the signal-to-noise ratio;

2. The excited-state contamination of the baryon matrix elements is serious, which will introduce systematic errors;

$$\langle H(t_f) \mathcal{O}(t) H^\dagger(t=0) \rangle = \exp(-E_0 t_f) \left(\langle 0 | \mathcal{O} | 0 \rangle + \sum_{i>0} \langle i | \mathcal{O} | i \rangle e^{-\Delta_i t_f} + \sum_{i \neq 0; j \neq 0} \langle i | \mathcal{O} | j \rangle (e^{-\Delta_i(t_f-t) - \Delta_j t} + e^{-\Delta_j(t_f-t) - \Delta_i t}) \right)$$



Y.B. Yang, et al, PRD **94**(2016)5, 054503
ETMC, PRD **102**(2020)5, 054517



2. Excited-State Contaminations in LQCD



- ChPT guides the very low-lying excited-states contamination (<300 MeV) from $N\pi$ and $N\pi\pi$ in 3-pts:

- $N\pi$ ESC in N 2-pt and 3-pt(nucleon charge)

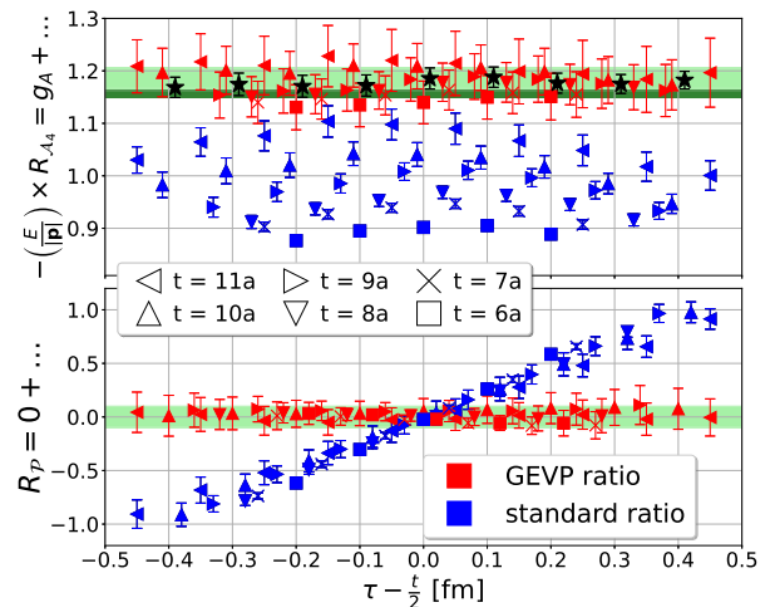
B. Tiburzi, PRD 91 (2015) 094510

- $N\pi$ ESC in N 3-pt with $J^\mu = \bar{q}\gamma_5 q$ can be very large;

O. Bär, PRD 100 (2019) 5, 054507

- $N\pi$ and $N\pi\pi$ in N 3-pts with $J^\mu = \bar{q} I q$ can be very large;

R. Gupta, et al. PRD 99 (2019) 5, 054506



$$\langle \mathcal{N} | \bar{q} \gamma^5 q(Q^2) | \mathcal{N} \rangle \quad \langle \mathcal{N} | \bar{q} \gamma^5 \gamma^4 q(Q^2) | \mathcal{N} \rangle$$

L. Barca, PRD 112(2025) 9, L091503

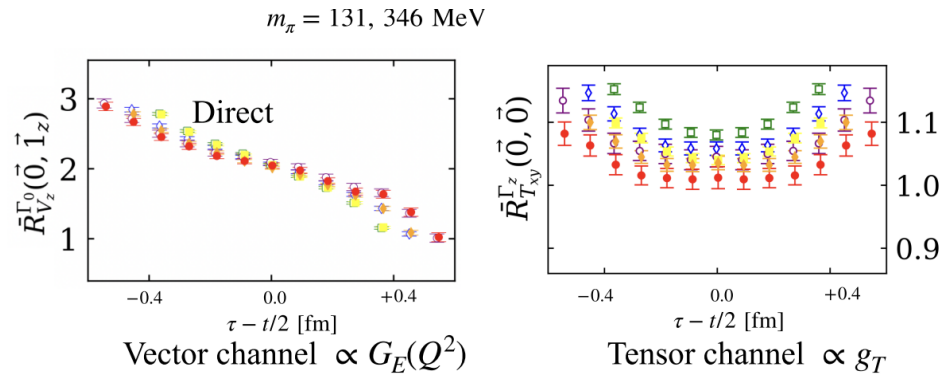
- ✓ After GEVP with $N\pi$, the convergence behaviors have been improved significantly for **axial vector** and **pseudoscalar matrix** element in LQCD;

- ◆ $N\pi$ is not general !!!

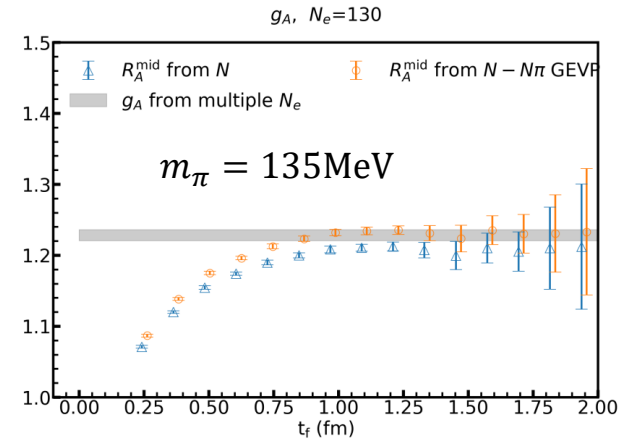
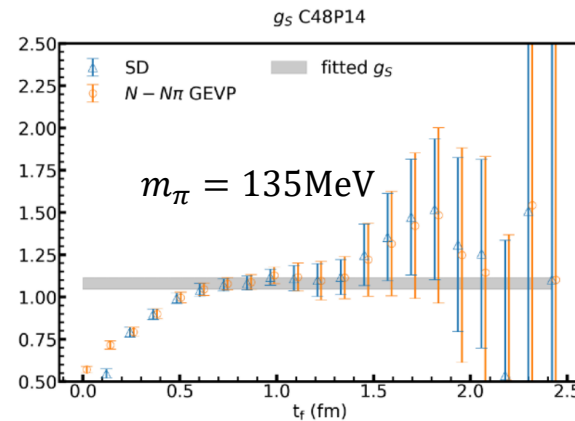
2. Current enhanced Excited-State contaminations



- $N\pi$ states **is not import** for nucleon charge at rest framework;

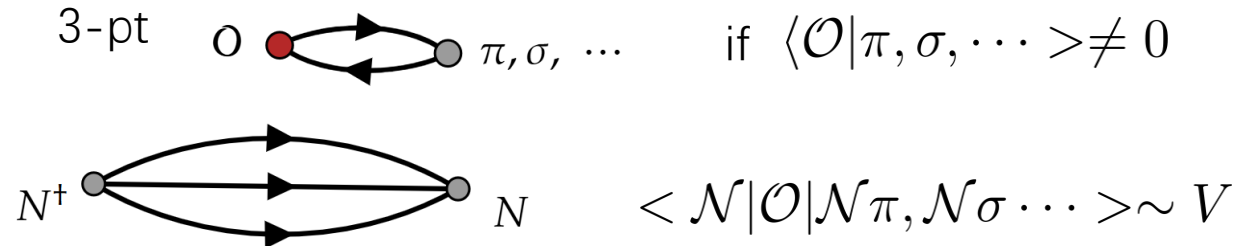
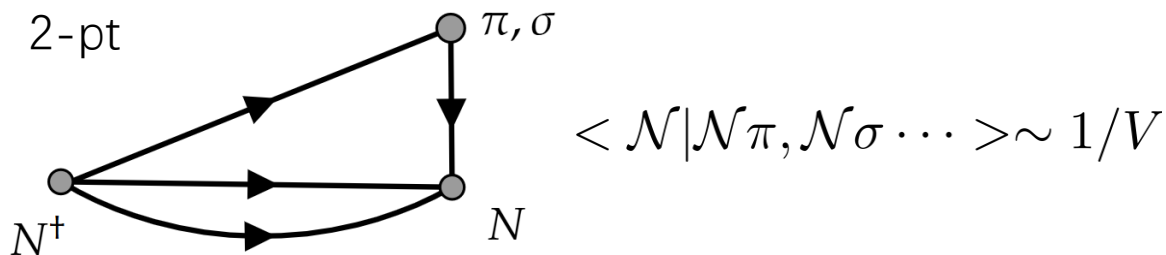


ETMC, *PRD 110 (2024) 9, 094514*



CLQCD, *arXiv:2511.02326*

- **Current enhanced Excited-State contamination;** *Lorenzo Barcal, PRD 112 (2025) 9, L091503*

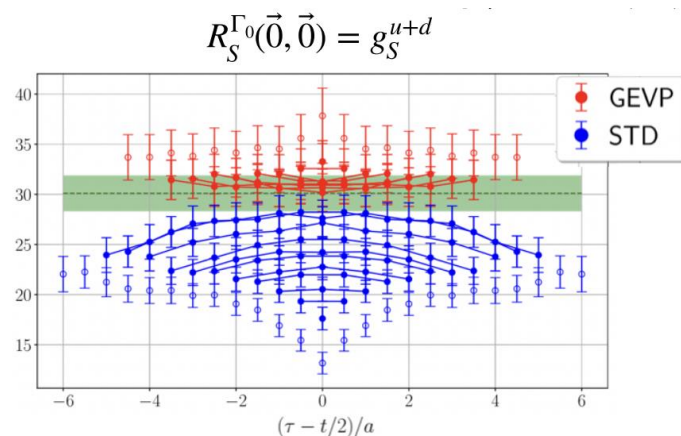


- After volume enhancement, nucleon matrix elements get **O(1)** excited state contamination;

2. Current enhanced Excited-State contaminations

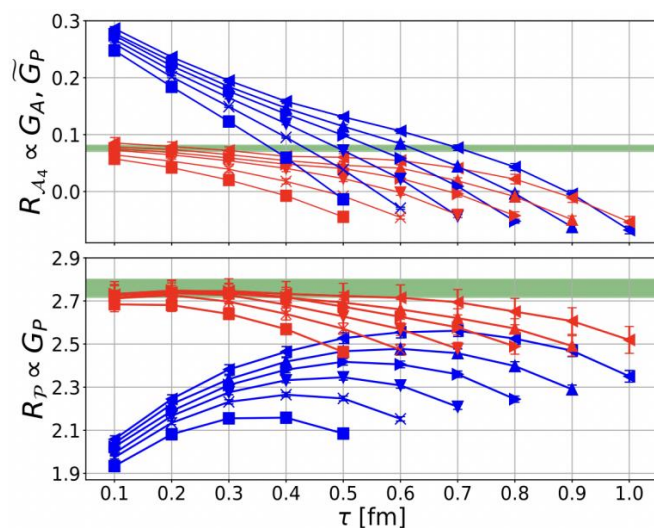


g_s v.s. $N\sigma$



Lorenzo Barcal, PRD 111 (2025) 3, L031505

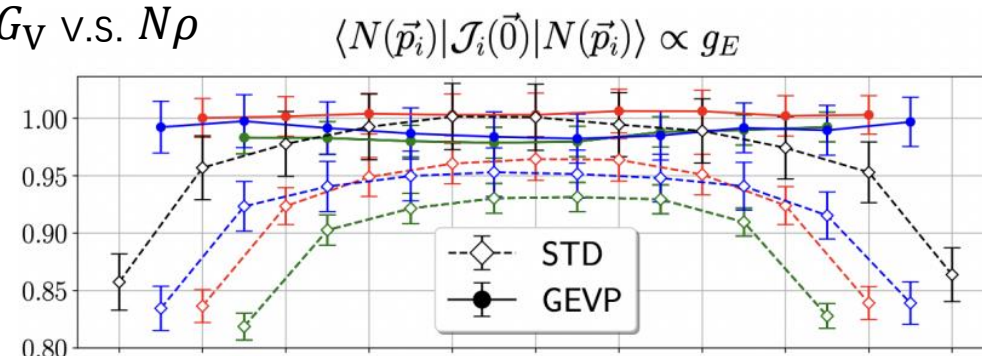
G_p v.s. $N\pi$



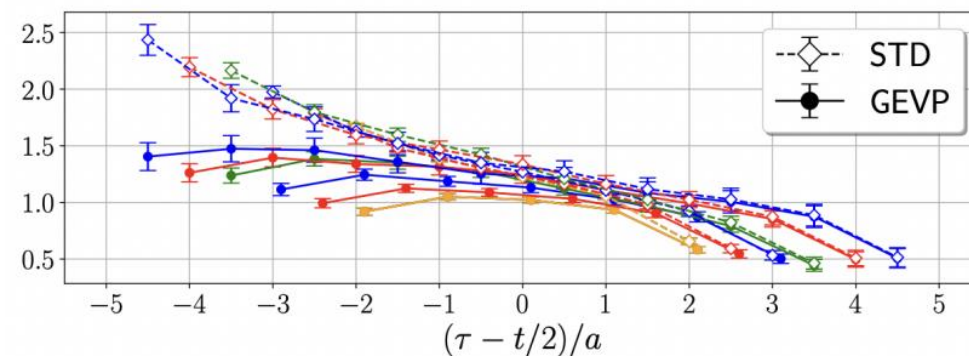
Lorenzo Barcal, PRD 107 (2023) 5, L051505

✓ GEVP removes leading order excited contaminations;

G_V v.s. $N\rho$



$\langle N(\vec{p}_i) | \mathcal{J}_i(\vec{0}) | N(\vec{p}_i) \rangle \propto g_E$



Lorenzo Barcal, PRD 112 (2025) 9, L091503

◆ Chapter 1: Background

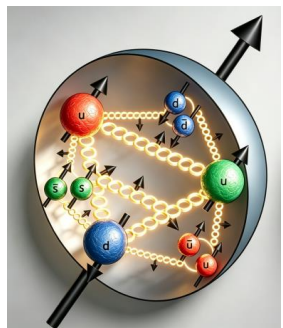
◆ Chapter 2: Excited-State Contaminations on LQCD

◆ **Chapter 3: Numerical result**

◆ Chapter 4: Summary



3. Application on nucleon matrix element: GEVP



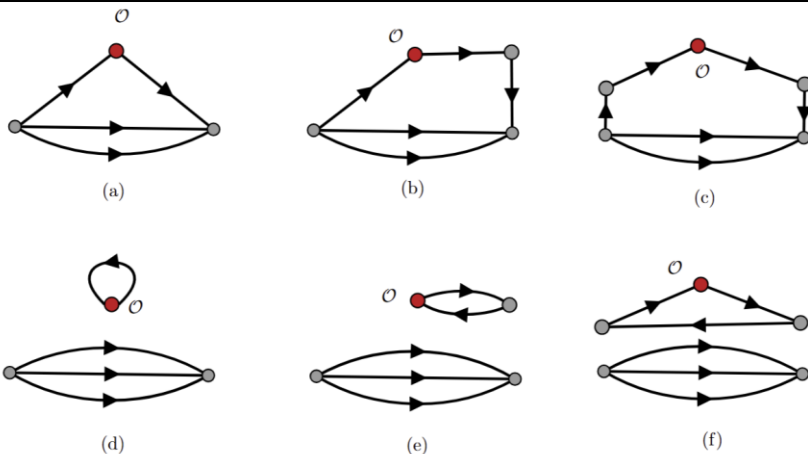
● Proton spin puzzle

$$\frac{\hbar}{2} = \frac{1}{2} \sum_q \Delta q + L_q + J_g$$

X.D. Ji, PRL 78, 610(1997)

- Quark spin can be extracted from the local matrix element g_A^{u+d+s} ;

3x3 GEVP based on the operator basis $(\mathcal{N}, \mathcal{N}f_1, \mathcal{N}f'_1)$



- ✓ Major excited-state contamination comes from $\mathcal{N}f_1$ scattering state for g_A^{u+d+s} ;

$f_1(1285)$ width 23MeV

$$\mathcal{N} = P^+ u(uC\gamma^5 d)$$

$$f_1 = \frac{1}{\sqrt{3}} (\bar{u}\gamma_5\gamma^i u + \bar{d}\gamma_5\gamma^i d + \bar{s}\gamma_5\gamma^i s)$$

$$f'_1 = \frac{1}{\sqrt{6}} (\bar{u}\gamma_5\gamma^i u + \bar{d}\gamma_5\gamma^i d - 2\bar{s}\gamma_5\gamma^i s)$$

- ✓ The **s-wave** scattering of nucleons and $f_1(1^+)$ mesons can couple to the nucleon $(1/2^+)$ representation.

- ✓ The Full GEVP analysis need **$\mathcal{O}(1000)$ Wick Construction**;

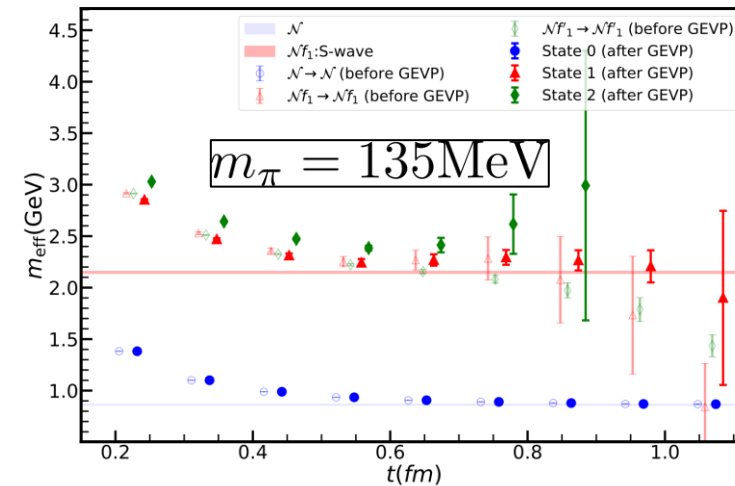
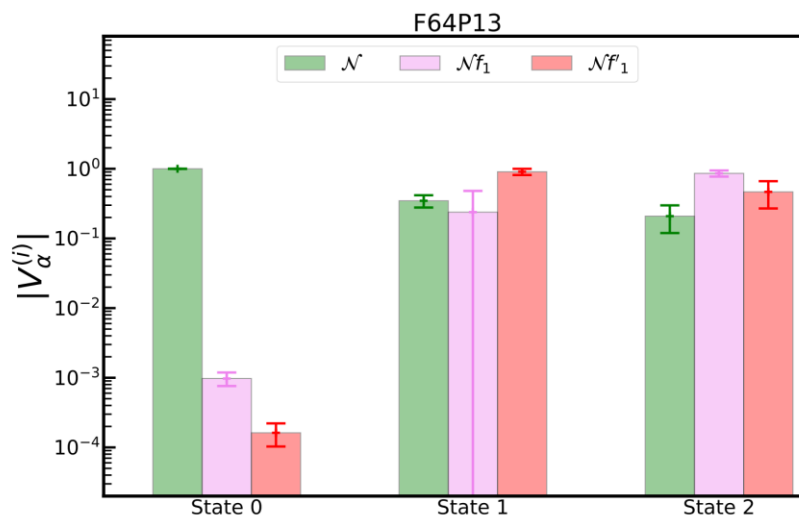
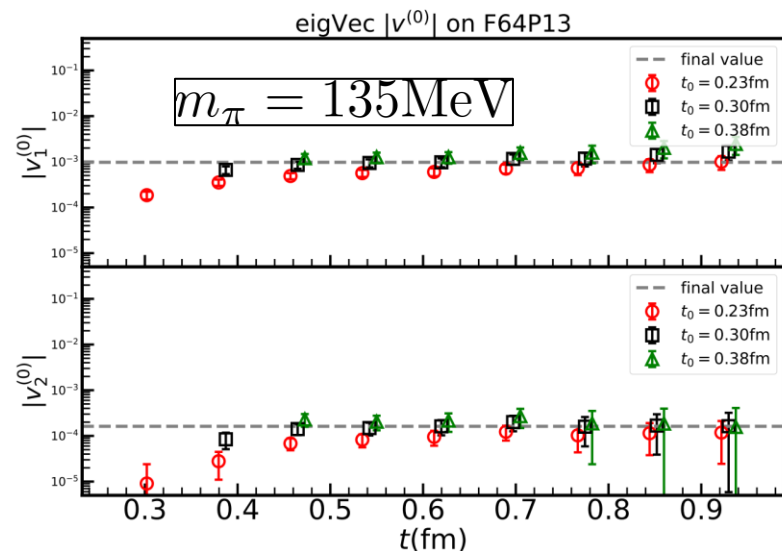
- ✓ “Blending” method realizes **all-to-all** propagator.

Z.C. Hu [CLQCD] arXiv: 2505.01719

3. Application on nucleon matrix element: GEVP



$$C_{2,\alpha\beta}(t)V_{\beta}^n(t,t_0) = \lambda^n(t,t_0)C_{2,\alpha\beta}(t_0)V_{\beta}^n(t,t_0),$$



- ✓ GEVP vector can be uniquely determined by solving the GEVP equation and by the imaginary-time evolution;
- ✓ After applying the GEVP, a state with an energy near the no-interaction Nf_1 threshold is obtained;

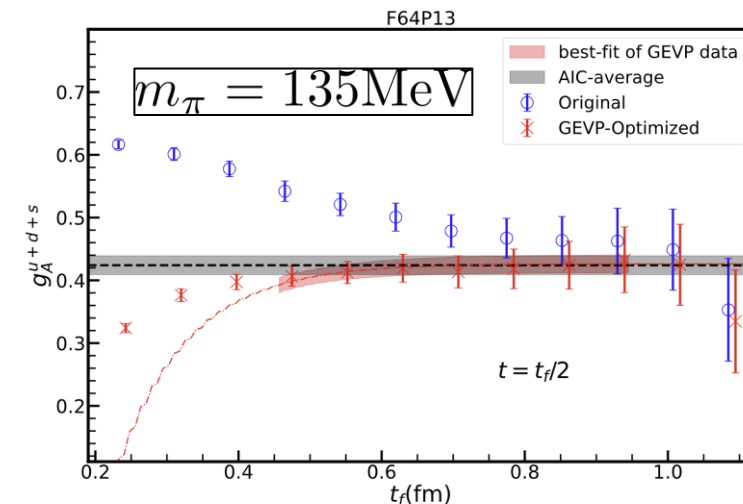
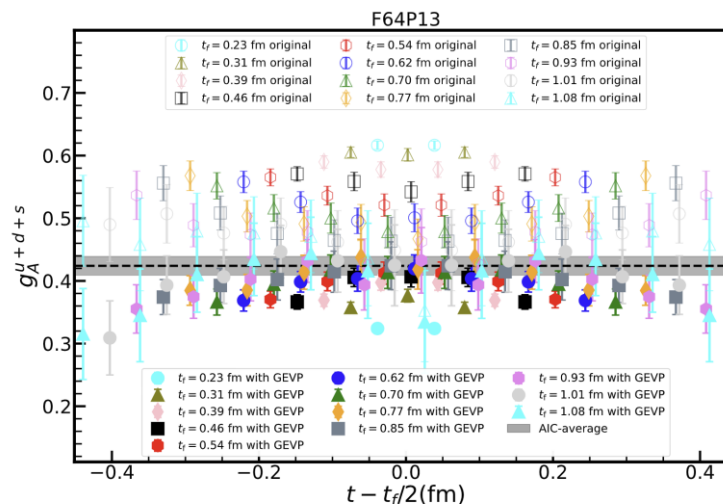
3. Application on nucleon matrix element: GEVP



✓ Applying GEVP to 3-p function, Convergent behaviors are largely improved after GEVP;

$$C_{3,\alpha\beta}(t_f, t) \equiv \langle X_\alpha(t_f) A_\Sigma(t) X_\beta^\dagger(0) \rangle,$$

$$\tilde{C}_{3,00}(t_f, t) \equiv V_\alpha^0(t, t_0) C_{3,\alpha\beta}(t_f, t) V_\beta^0(t, t_0),$$



✓ Even though the overlap is only about 0.1%, it plays a very effective role in the 3-p process;

✓ The result with AIC on fit windows is $g_A^{u+d+s}(a = 0.077\text{fm}) = 0.417(17)$, achieving world leading precision at physical point directly.

	Ensemble	L	T	a(fm)	m_π	n_{cfg}	g_A^s	g_A^{u+d+s}	Inversion	Inv.to 4% g_A^{u+d+s}
PNDME	a09m310	64	96	0.09	138	847	-0.047(27)	—	8.8 M	—
ETM17	cA2.09.48	48	96	0.09	135	2136	-0.021(5)	0.402(35)	7.2 M	34.1 M
ETM19	cB211.072.64	64	128	0.08	139	750	-0.046(7)	0.382(31)	1.37 M	5.63 M
This Work	F64P13	64	128	0.078	134	46	-0.0409(37)	0.416(17)	0.39 M	0.39 M

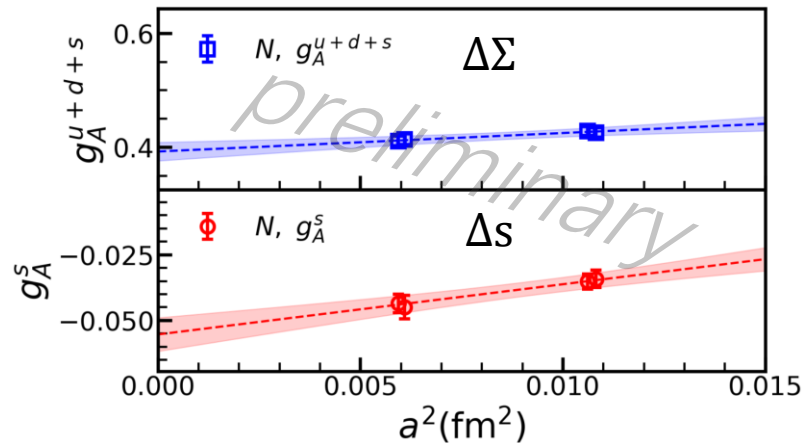
3. Quark spin at continue limit



➤ Continuum extrapolation based 4 ensembles($a \rightarrow 0$)

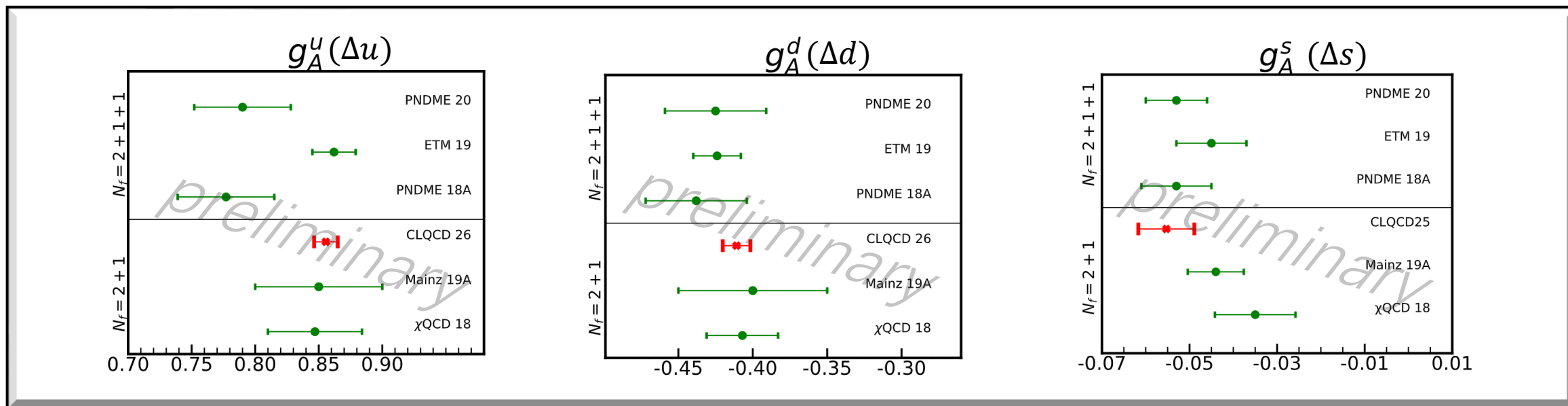
	a [fm]	$L^3 \times T$	m_π [MeV]	$m_\pi L$	N_{cfg}
C32P29	0.105	$32^3 \times 64$	292.4(1.1)	5.01	160
C48P13	0.105	$48^3 \times 96$	135.5(1.6)	3.47	53
F48P30	0.077	$48^3 \times 96$	303.4(0.9)	5.72	40
F64P13	0.077	$64^3 \times 128$	134.1(1.5)	3.37	43

$$g_A(a, m_\pi) = g_A^{\text{phys}} + c_1(m_\pi^2 - m_{\pi, \text{phys}}^2) + c_2 a^2.$$



$$\frac{1}{2} \Delta \Sigma = 39(2)\% S_p$$

$$\frac{1}{2} \Delta s = -5.5(6)\% S_p$$



3. Quark spin on Octet baryons



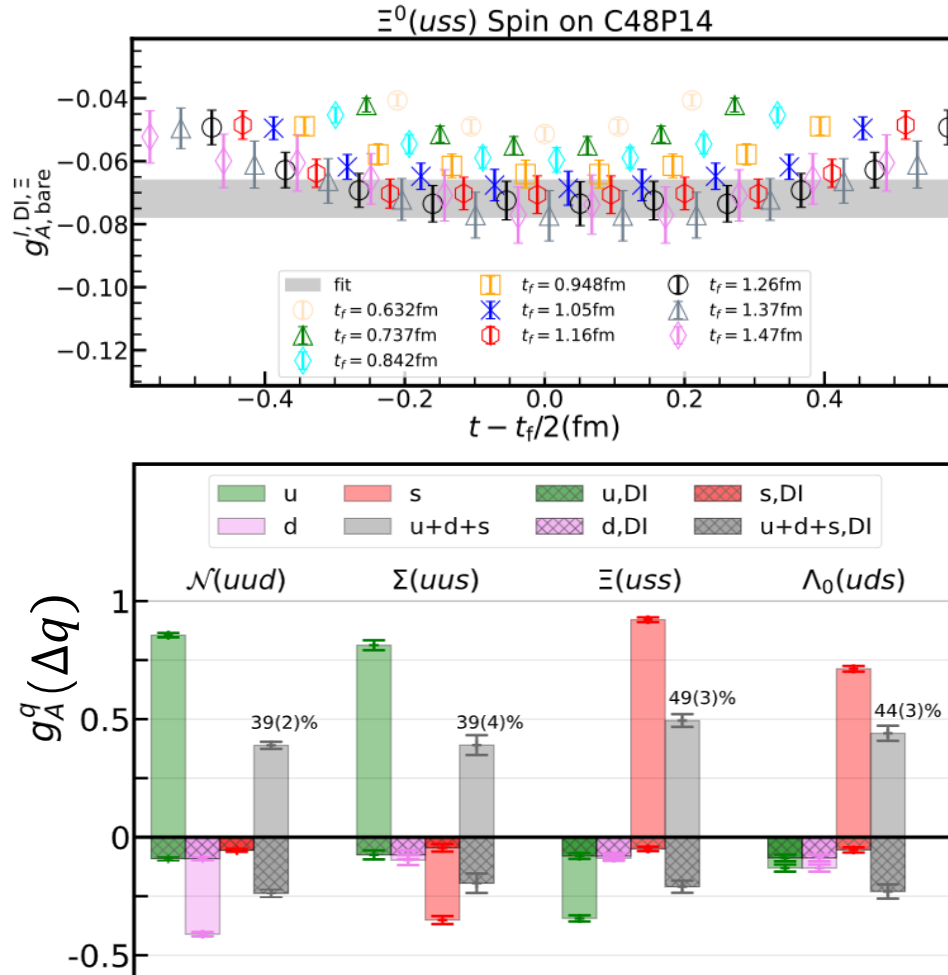
✓ For the first time, the quark spin of the hyperon was obtained at the physical point;

◆ The contribution of the sea quarks is significantly non-zero;

1. $\mathcal{N}(uud)$ with $\Delta s = -5.53(64)\%S_P$;
2. $\Sigma^+(uus)$ with $\Delta d = -10(2)\%S_\Sigma$;
3. $\Xi^0(uss)$ with $\Delta d = -9(1)\%S_\Xi$;

◆ $\Delta u, \Delta d$ on Λ_0 are significantly non-zero too;

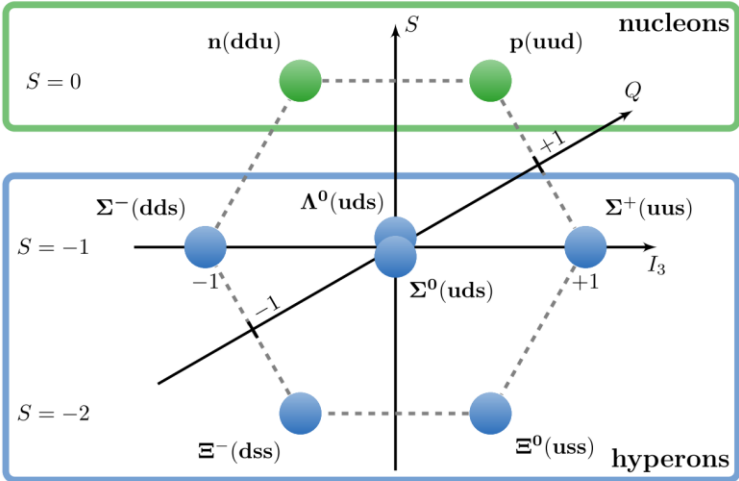
$$\Delta u = \Delta d = -13(2)\%S_{\Lambda_0}$$



3. Quark spin in octet baryons and SU(3)-flavour symmetry



➤ Quark Spin could be extrapolated to Hyperon
From nucleon data based in SU(3)_f symmetry

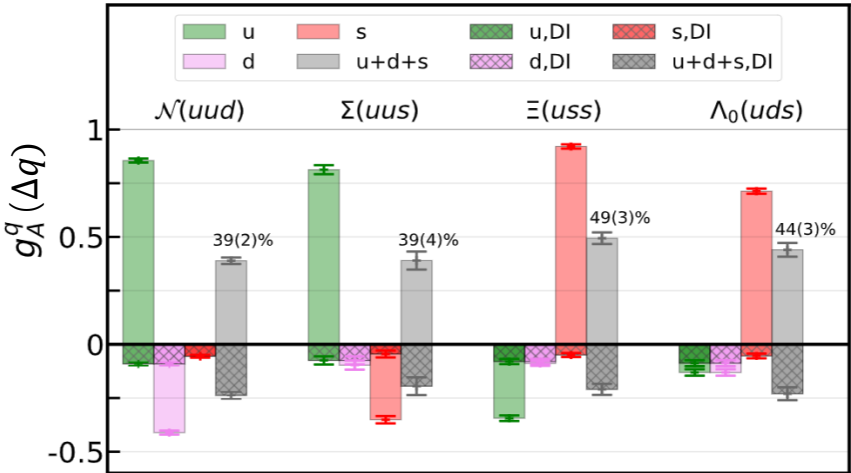


SU(3)_f symmetry

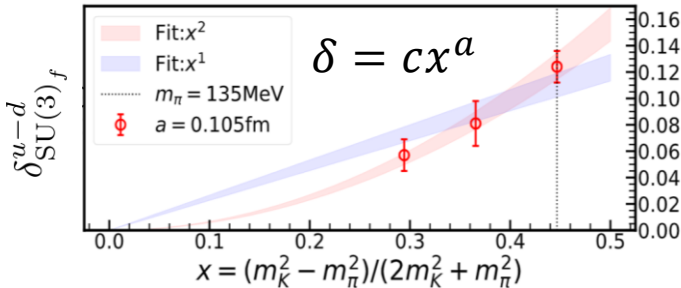
- Isospin vector:
$$g_{A,p}^{u-d} + g_{A,\Xi^0}^{u-d} - g_{A,\Sigma^+}^{u-d} = 0;$$
- Flavour singlet:
$$g_{A,N}^{u+d+s} = g_{A,\Xi}^{u+d+s} = g_{A,\Sigma}^{u+d+s} = g_{A,\Lambda}^{u+d+s};$$

Baryon	ΔU	ΔD	ΔS
p	0.83 ± 0.03	-0.43 ± 0.03	-0.10 ± 0.03
n	-0.43 ± 0.03	0.83 ± 0.03	-0.10 ± 0.03
Σ^+	0.83 ± 0.03	-0.10 ± 0.03	-0.43 ± 0.03
Σ^-	-0.10 ± 0.03	0.83 ± 0.03	-0.43 ± 0.03
Ξ^0	-0.43 ± 0.03	-0.10 ± 0.03	0.83 ± 0.03
Ξ^-	-0.10 ± 0.03	-0.43 ± 0.03	0.83 ± 0.03
Σ^0	0.37 ± 0.03	0.37 ± 0.03	-0.43 ± 0.03
Λ	-0.17 ± 0.03	-0.17 ± 0.03	0.63 ± 0.03

D. Ashery; et.al; PLB 469,263-269 (1999)



➤ $(m_s - m_l)$ drives the breaking of SU(3)_f symmetry;



$$\delta_{\text{SU}(3)}^{u-d} = \frac{g_{A(u-d)}^{\mathcal{N}} - g_{A(u-d)}^{\Sigma^+} + g_{A(u-d)}^{\Xi^0}}{g_A^{\text{exp}}} = 11(1)\%$$

$$\delta_{\text{SU}(3)}^{u+d+s} = \frac{g_{A,\Xi}^{u+d+s} - g_{A,\mathcal{N}}^{u+d+s}}{(g_{A,\Xi}^{u+d+s} + g_{A,\mathcal{N}}^{u+d+s})/2} = 24(7)\%,$$

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◆ Chapter 3: Numerical result

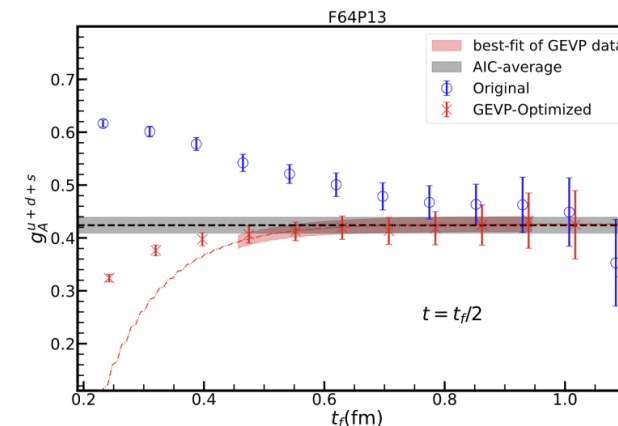
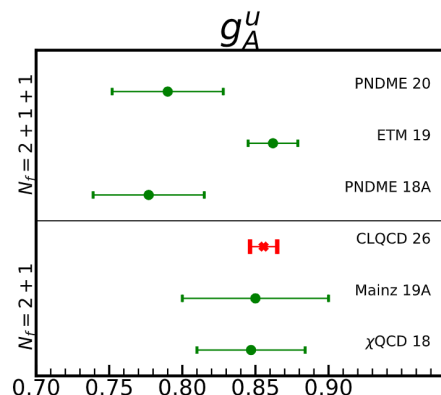
◆ Chapter 4: Summary



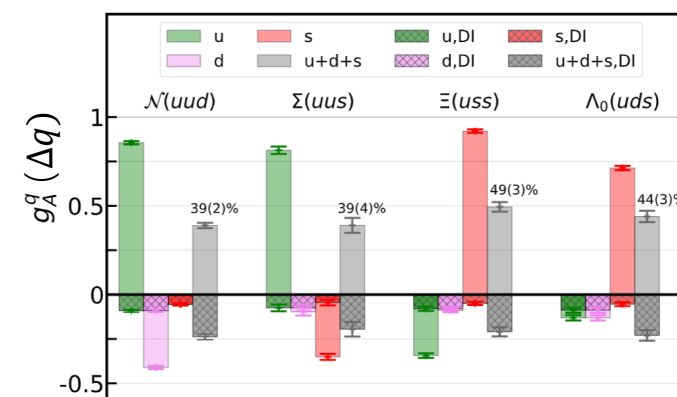
4. Summary



1. We completed GEVP-analysis based on nucleon and f_1 scattering state to improve convergence behaviors of 3-pt effectively based on the idea of “Current enhanced excited state contaminations”;



2. We achieve world leading precision of quark spin $\Delta\Sigma$ based on the “Blending” method and the GEVP analysis.



3. We obtained the quark spin on octet baryons systematically;

Thank you!