

Bottom baryon mass and sigma term on lattice QCD

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Hadron Mass Decomposition

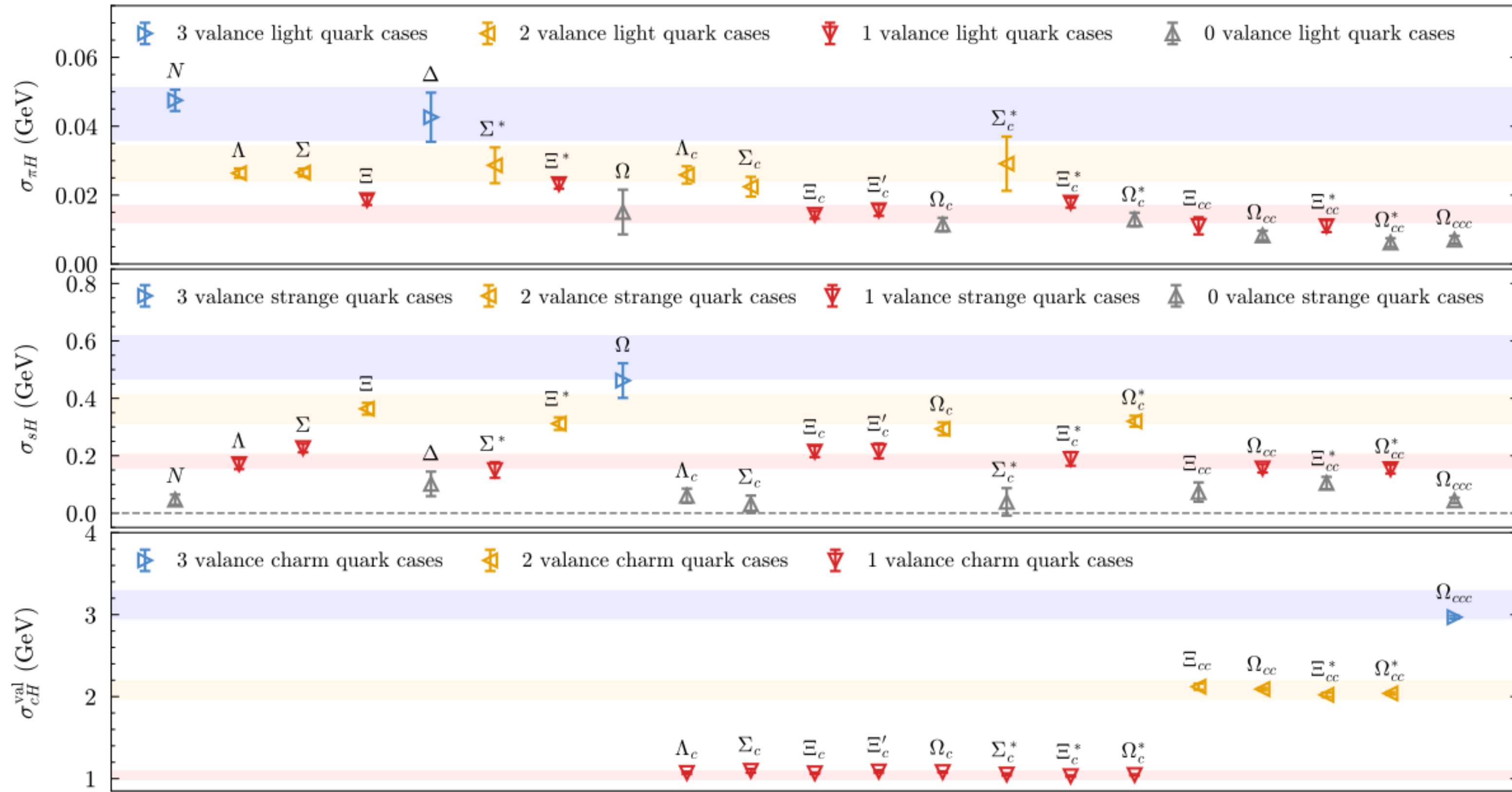
- Hadron mass decomposition: $m_H = \sum_q \sigma_{q,H} + \gamma_m \sum_q \sigma_{q,H} + \frac{\beta(\alpha_s)}{2\alpha_s} \langle G^{\mu\nu} G_{\mu\nu} \rangle_H$
- Sigma term: quark flavor contribution $\sigma_{q,H} \equiv m_q \langle \bar{q}q \rangle_H$, where $\langle \bar{q}q \rangle_H$ with $\langle O \rangle_H \equiv \langle H | O | H \rangle / \langle H | H \rangle$ can be understood as the condensate of quark and antiquark within the hadron H .
- Scale violation of quark mass dimension $\gamma_m = \frac{2}{\pi} \alpha_s + O(\alpha_s^2)$
- Gluon trace anomaly part: $\frac{\beta(\alpha_s)}{2\alpha_s} = (-\frac{11}{8\pi} + \frac{N_f}{12\pi}) \alpha_s + O(\alpha_s^2)$
- Extract sigma term: Feynman-Hellman theorem $\sigma_{q,H}^{\text{FH}} = m_q^{\text{PC}} \frac{\partial m_H}{\partial m_q^{\text{PC}}}$ or directly 3-pt calculation.

Hadron Mass Decomposition

Light and charmed baryon case

- Hadron mass decomposition:

$$m_H = \sum_q \sigma_{q,H} + \gamma_m \sum_q \sigma_{q,H} + \frac{\beta(\alpha_s)}{2\alpha_s} \langle G^{\mu\nu} G_{\mu\nu} \rangle_H$$
- Flavor-dependent enhancement for quark contribution: 4-8 for light quark, 2-3 for strange quark and 1.2-1.3 for charm quark.
- What about bottom case?



Bolun Hu, Hai-Yang Du, et.al, 2411.18402

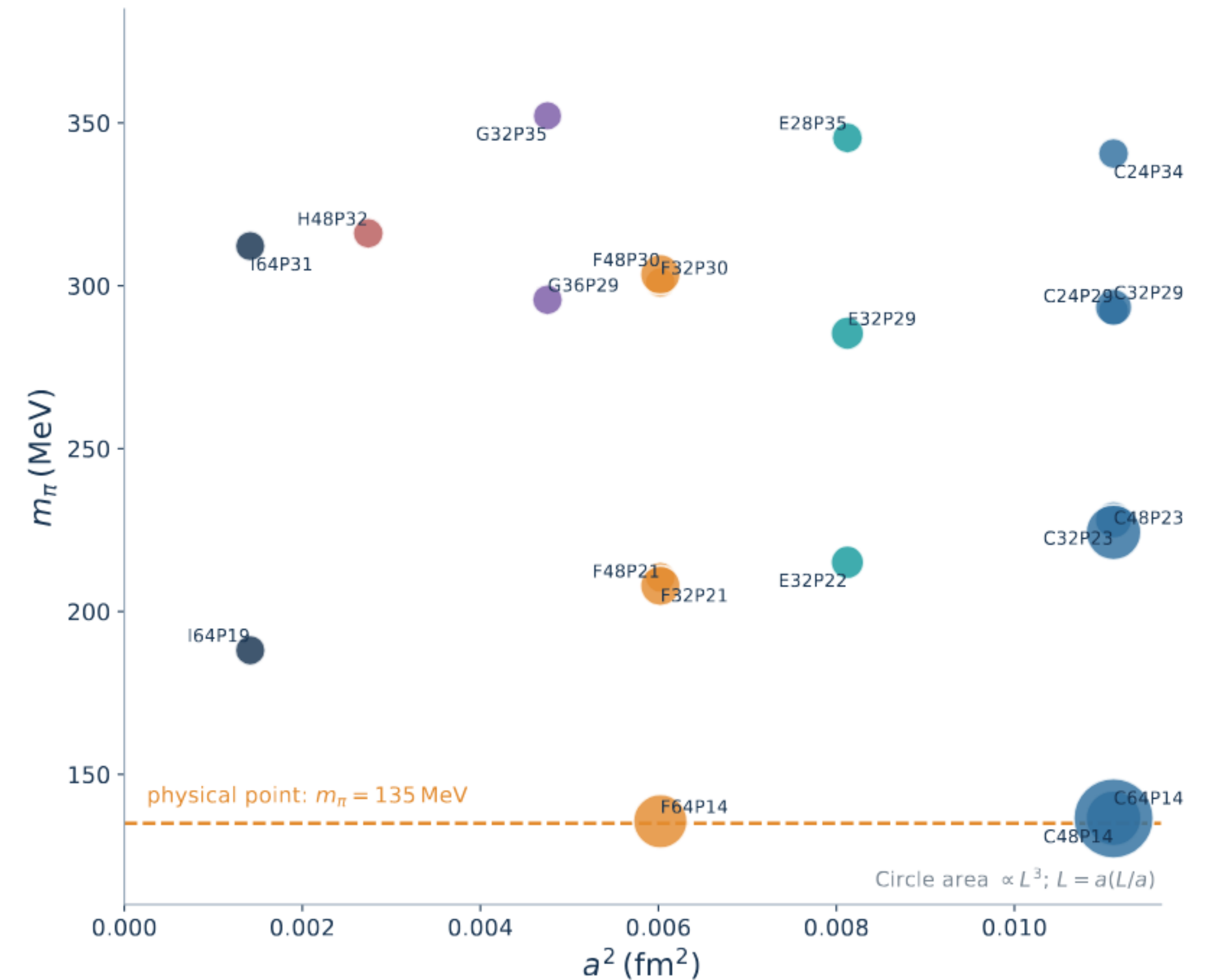
Ensemble Information

CLQCD clover ensembles

- CLQCD Collaboration ensembles: 2+1 clover action, 6 lattice spacing, several sea π and η_s mass and different volume.
- More than 4000 configurations.
- Chiral and continuum extrapolation.
- Simulate light, strange and charm quark very well.

Z.-C. Hu et al. (CLQCD), *Phys.Rev.D*, 109, 054507 (2024).

H-Y. Du et al. (CLQCD), *Phys.Rev.D*, 111, 054504 (2025).



Bottom Quark on Lattice QCD

Anisotropic clover action

- Heavy quark discretization error of bottom quark: order $(m_b a)^2$ and even $(m_b a)^4$, need finer lattice spacing.
- Extrapolation from lighter masses, Heavy Quark Effective Theory (HQET), Non-Relativistic QCD (NRQCD).

- Physical bottom quark on lattice QCD: Anisotropic clover action for bottom quark

$$S_Q = a^4 \sum_x \bar{Q} \left[m_Q + \gamma_4 \nabla_4 - \frac{a}{2} \nabla_4^2 + \nu \sum_{i=1}^3 (\gamma_i \nabla_i - \frac{a}{2} \nabla_i^2) - c_E \frac{a}{2} \sum_{i=1}^3 \sigma_{0i} F_{0i} - c_B \frac{a}{4} \sum_{i,j=1}^3 \sigma_{ij} F_{ij} \right] Q$$

Z. S. Brown, et.al *Phys.Rev.D*, 90(9) 094507.

L. Liu, et.al *Phys.Rev.D*, 81 094507

- Anisotropic valence bottom quark on isotropic gauge ensembles.

- Clover coefficients $c_E = \frac{1 + \nu}{2u_0^3}$, $c_B = \frac{\nu}{u_0^3}$. Heavy quark anisotropy ν .

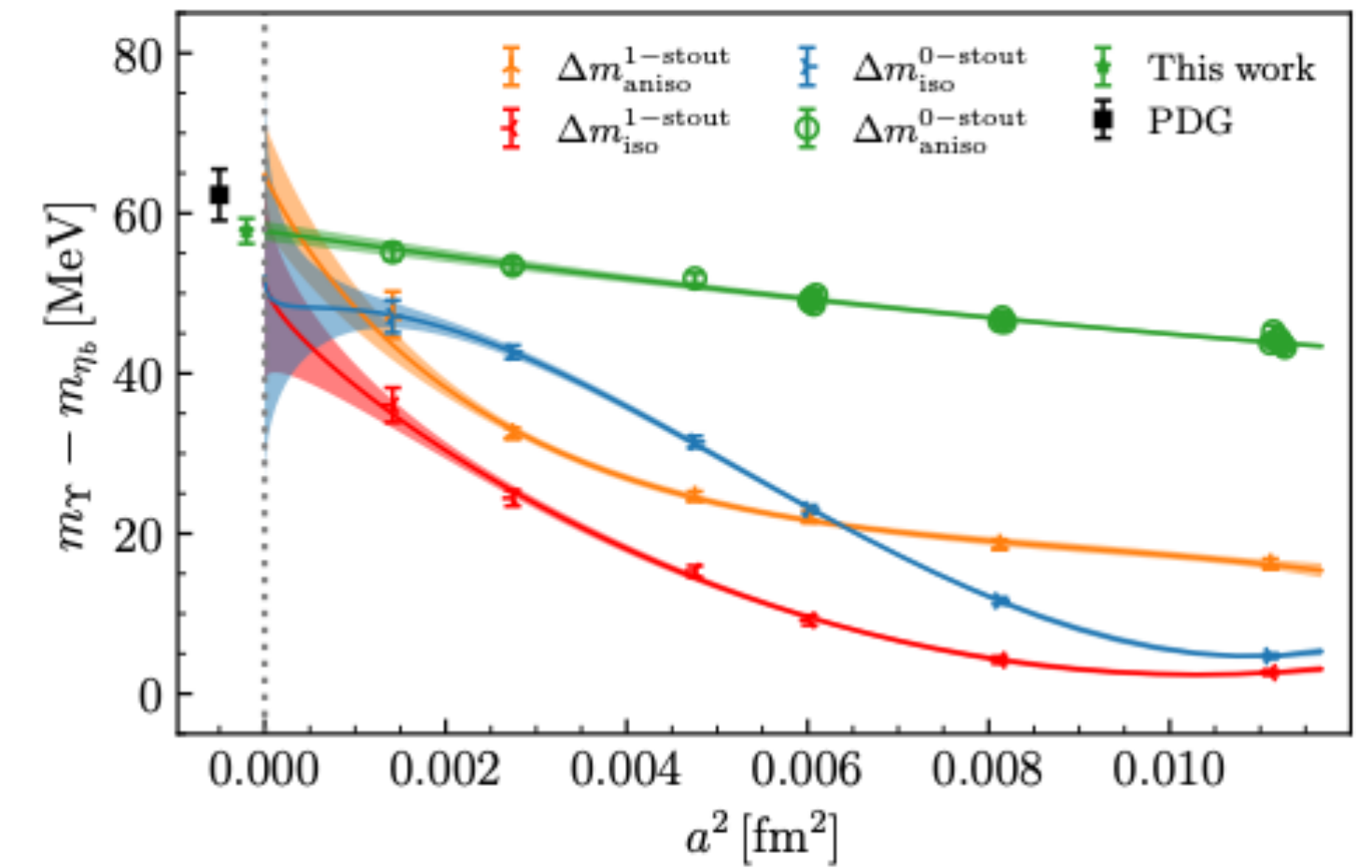
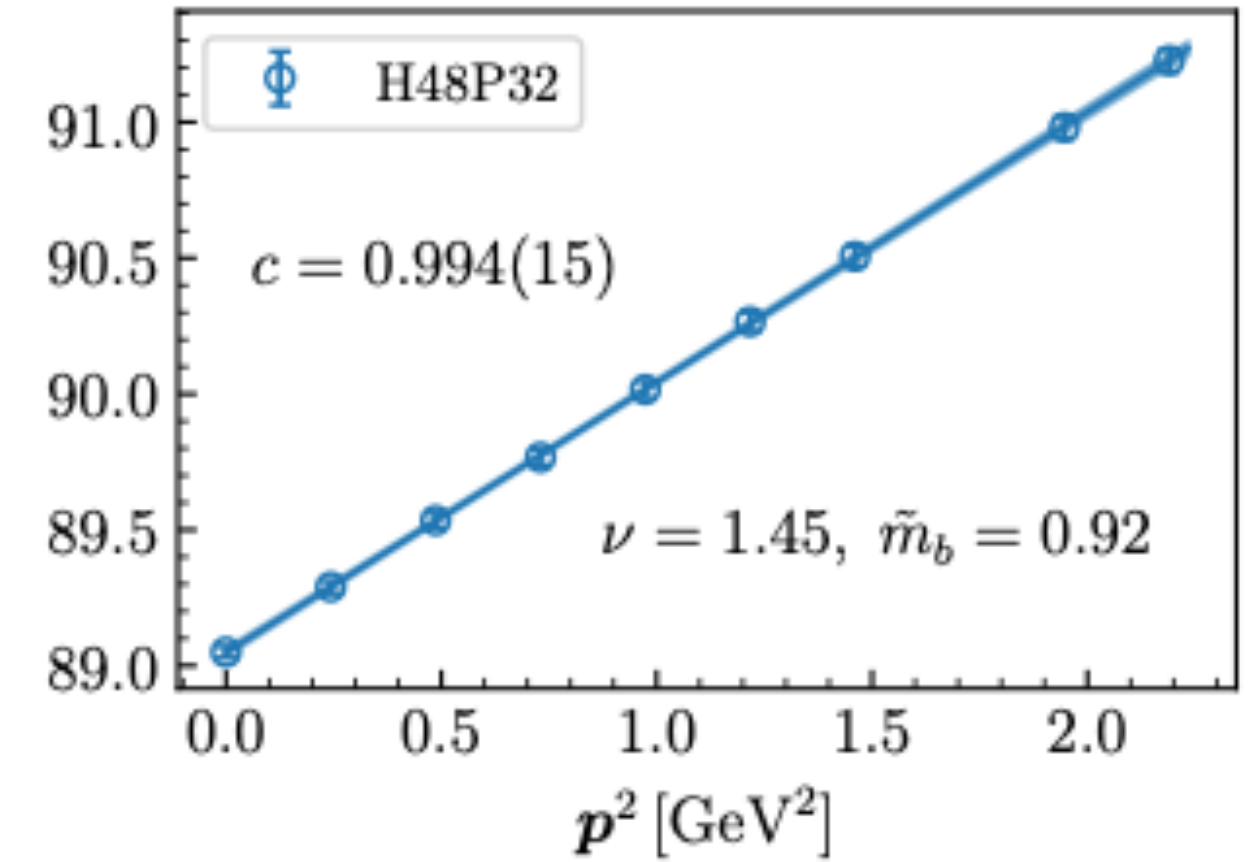
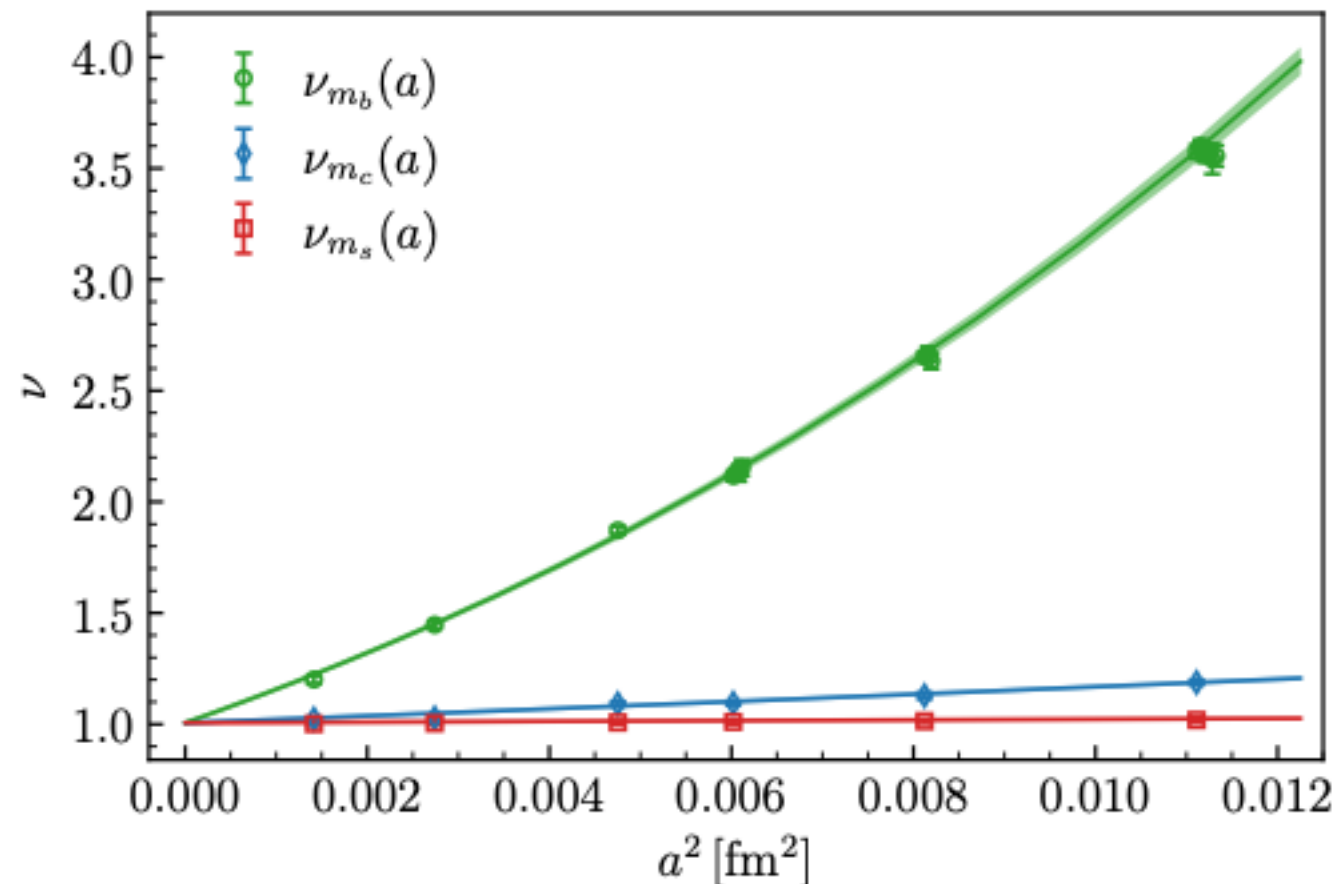
P. Chen, *Phys.Rev.D*, 64 034509

- Turning the values of m_Q and ν to satisfy dispersion relation and physical Υ (1S vector bottomonium) mass:
 $E_V(\vec{p})^2 = m_V^2 + c^2 p^2$, $c = 1$, $m_V = m_\Upsilon = 9460.40(10)$ MeV.

Bottom Quark on Lattice QCD

Parameters turning

- Turning the values of m_Q and ν to satisfy dispersion relation and physical Υ (1S vector bottomonium) mass:
 $E_V(\vec{p})^2 = m_V^2 + c^2 p^2, c = 1, m_V = m_\Upsilon = 9460.40(10) \text{ MeV}$
- Bottomonium hyperfine splitting with 0-step stout-link smearing and anisotropic bottom action: $\Delta_{\text{HFS}}^{bb} \equiv m_\Upsilon - m_{\eta_b} = 57.8(1.2)(1.0) \text{ MeV}$ agree with PDG well.
- ν tends to 1 in continuum.
- Empirical formula: $\nu = \frac{\text{Sinh}(cm_V a)}{cm_V a} \times [1 + c_1(m_\pi^2 - m_{\pi,phy}^2) + c_2(m_{\eta_s}^2 - m_{\eta_s,phy}^2)]$

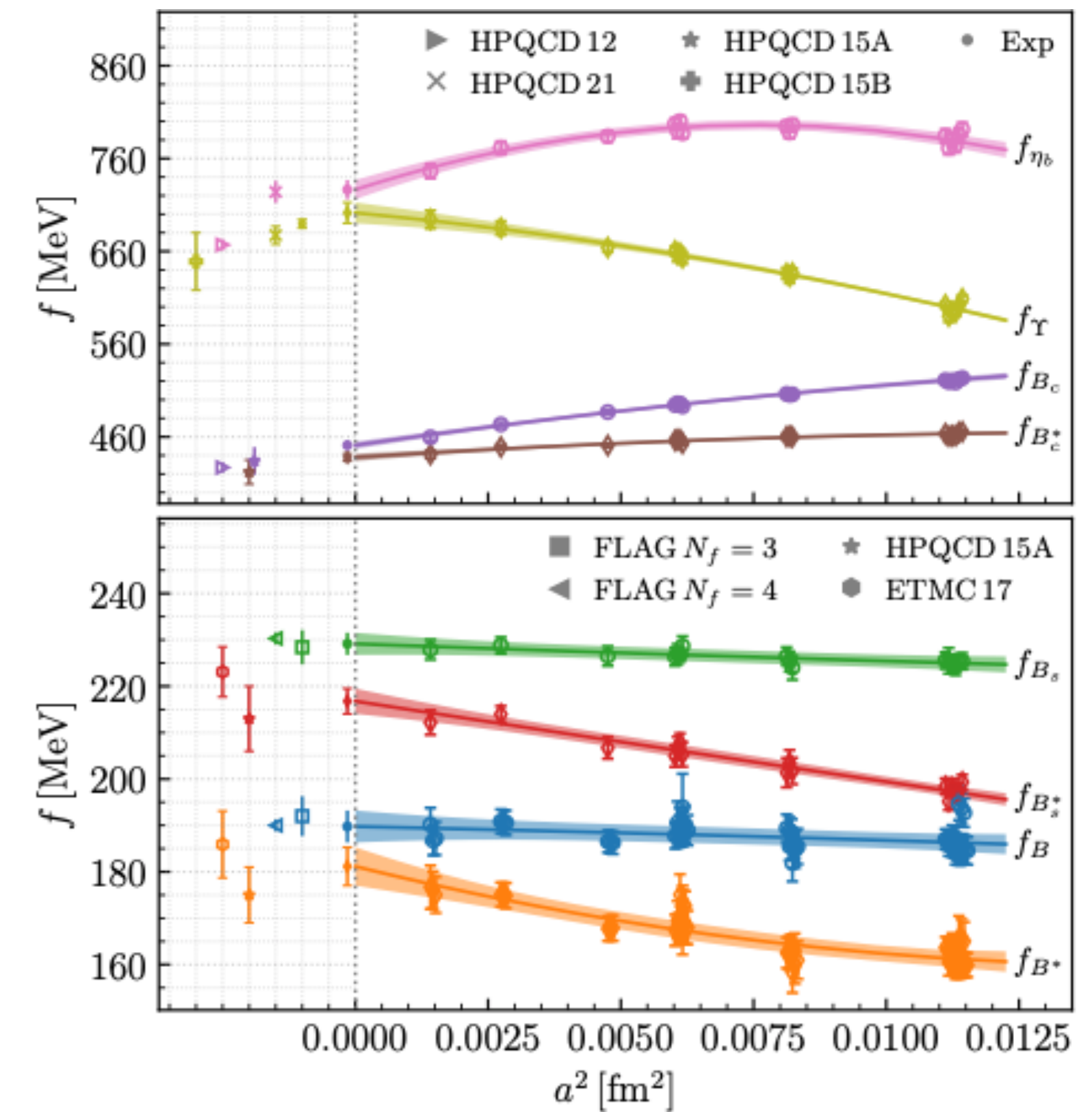
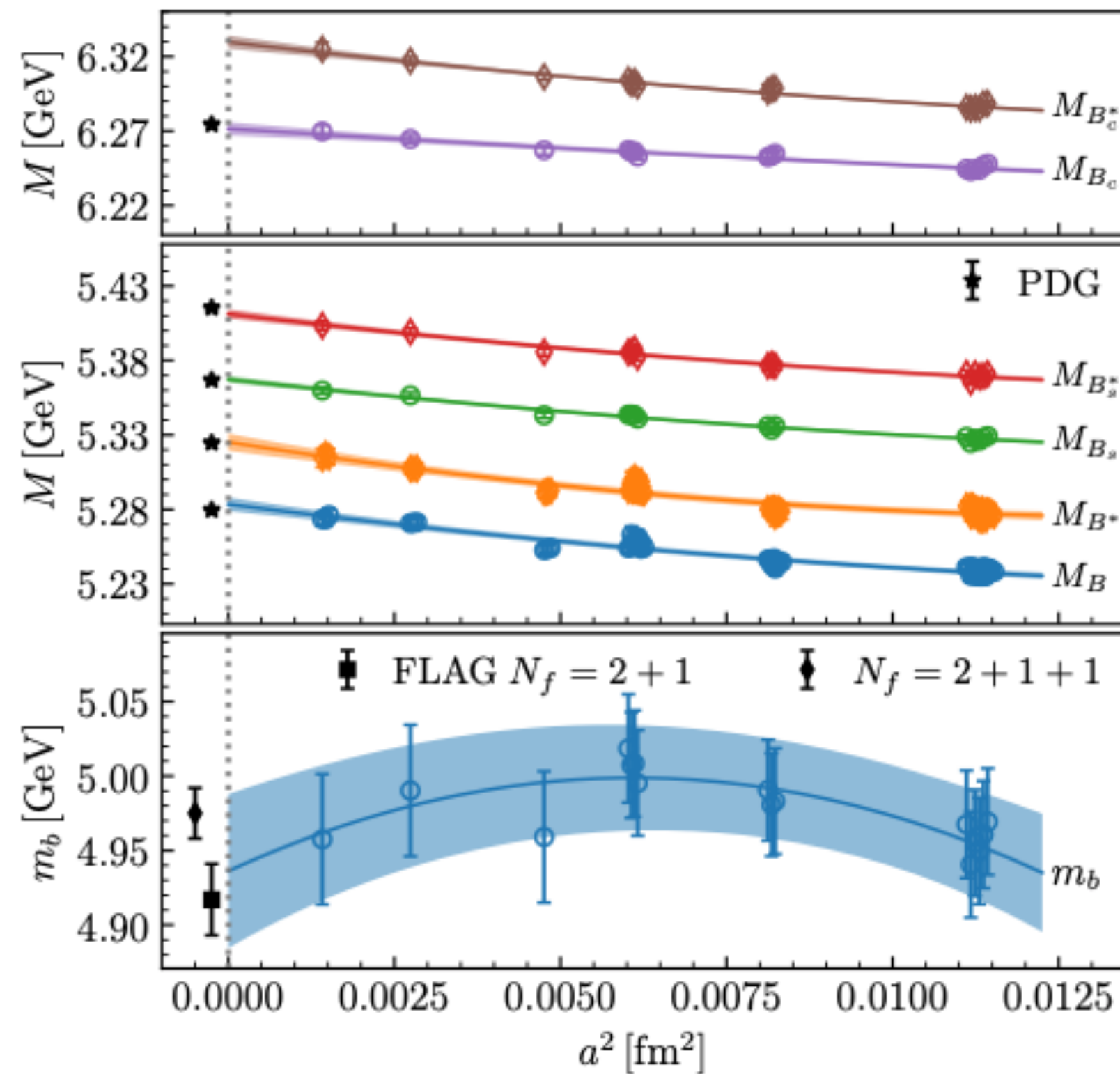


Mengchu Cai, Hai-Yang Du, et.al, 2603.01846

Bottom Quark on Lattice QCD

Meson results

- Bottom quark mass at $\overline{\text{MS}}$ 2 GeV.
- Open bottom meson mass.
- Decay constants of bottom meson.



Mengchu Cai, Hai-Yang Du, et.al, 2603.01846

Bottom Baryon

Interpolating field

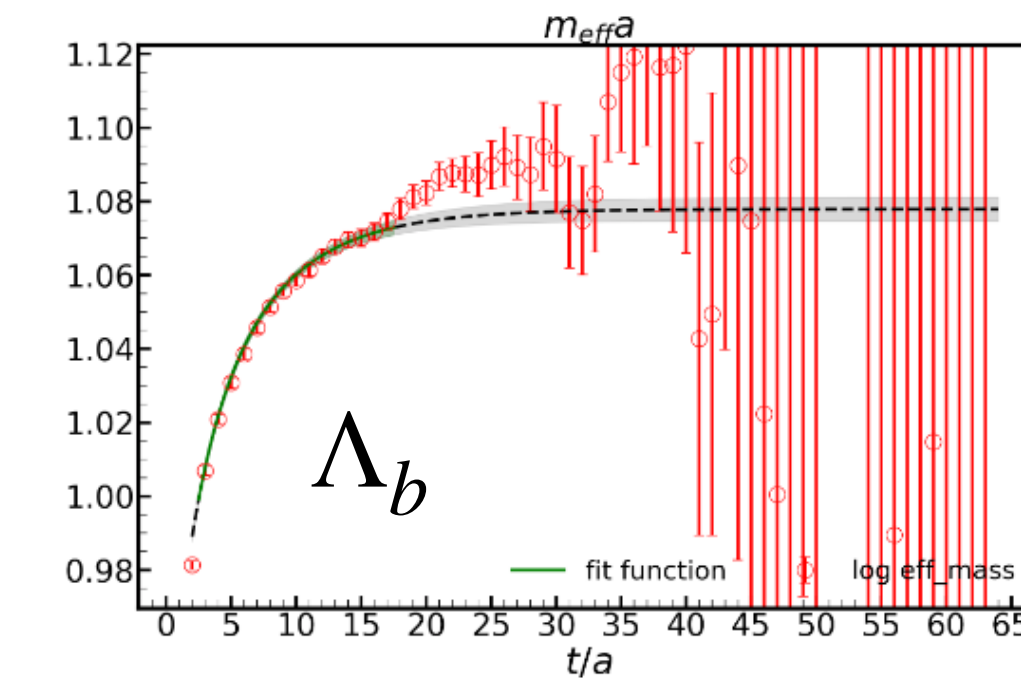
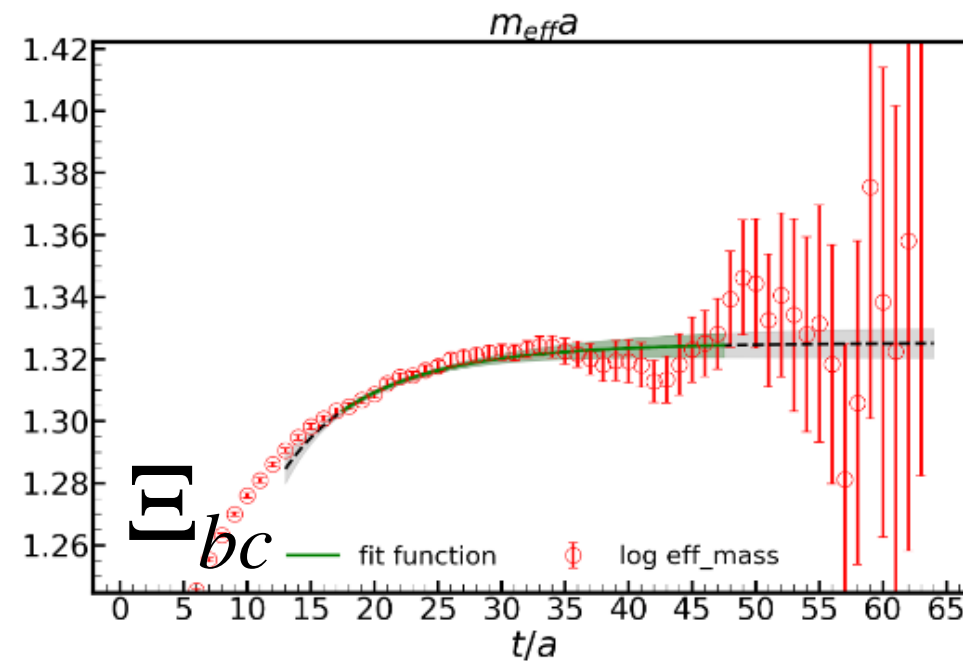
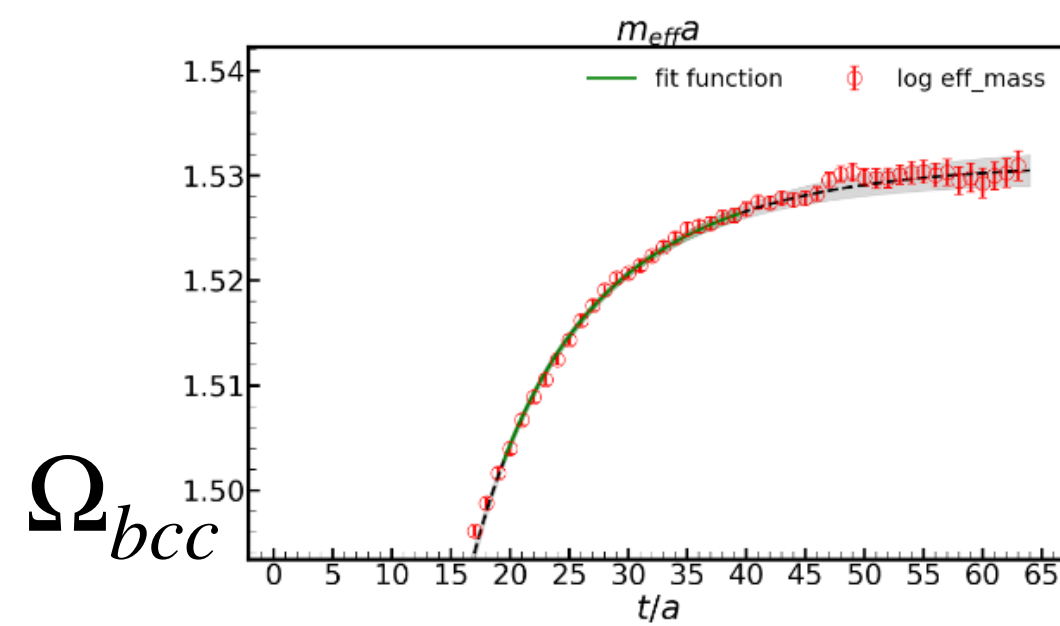
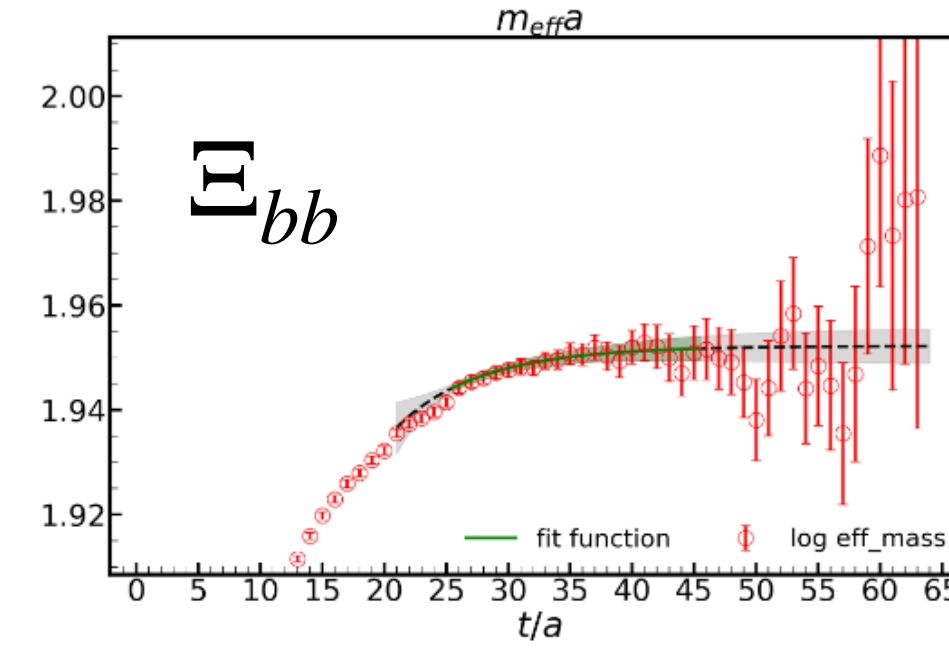
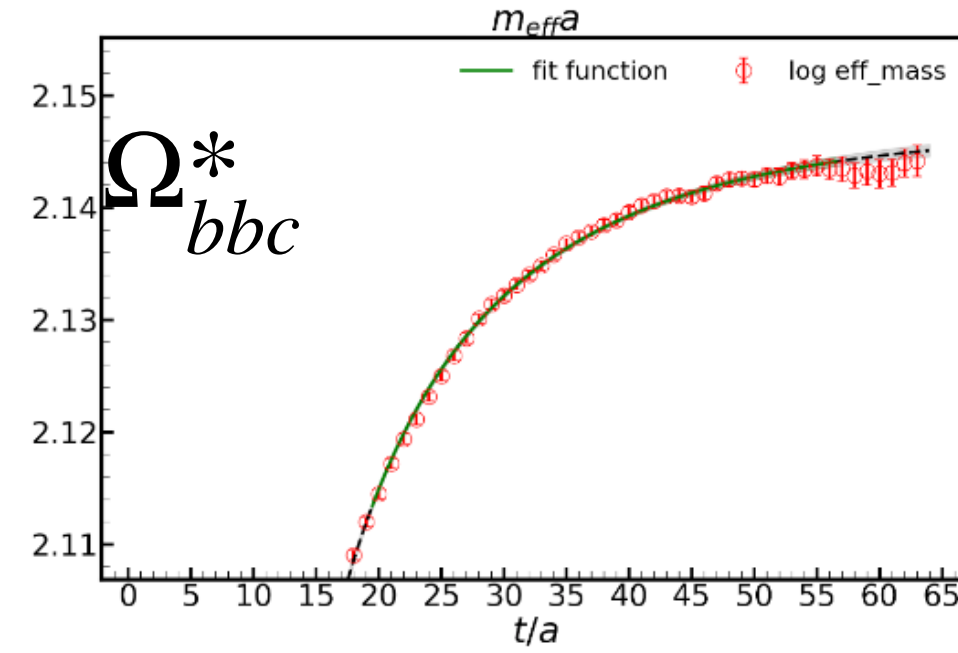
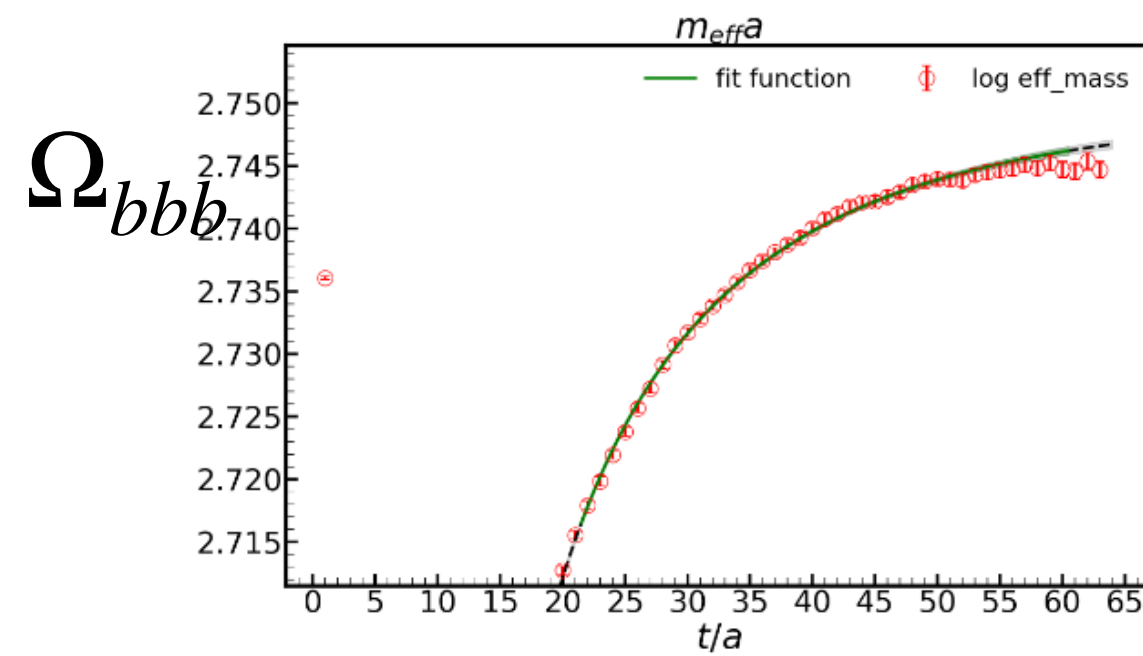
- Bottom baryon in quark model: $SU(5)_F$ symmetry.
- Interpolating field on lattice QCD:
 - Spin-1/2, Σ -type: $\Sigma : \epsilon^{abc} P^+ (Q_a^T C \gamma^5 q_b) q_c$. For Σ_Q , q is u or d and for Ω_Q , q is s .
 - Spin-1/2, Ξ' -type: $\Xi' : \frac{1}{\sqrt{2}} \epsilon^{abc} P^+ [(q_a^T C \gamma^5 Q_b) q'_c + (q_a'^T C \gamma^5 Q_b) q_c]$ where q is u or d , $q' = s$.
 - Spin-1/2, Λ -type: octet lambda $\Lambda_o : \frac{1}{\sqrt{6}} \epsilon^{abc} P^+ [2(q_a^T C \gamma^5 q'_b) Q_c + (q_a^T C \gamma^5 Q_b) q'_c - (q_a'^T C \gamma^5 Q_b) q_c]$ where $q = u, q' = d$ for Λ_Q , and $q = u, q' = s$ for Ξ_Q .
 - Spin-3/2, Σ^* : $\frac{1}{\sqrt{3}} \epsilon^{abc} P^+ [2(Q_a^T C \gamma^\mu q_b) q_c + (q_a^T C \gamma^\mu q_b) Q_c]$ $\Xi_Q^* : \frac{2}{\sqrt{3}} \epsilon^{abc} P^+ [(q_a^T C \gamma^\mu q'_b) Q_c + (q_a'^T C \gamma^\mu Q_b) q_c + (Q_a^T C \gamma^\mu q_b) q'_c]$
 - Interchange the role of light and heavy fields for double heavy quark.

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Bottom Baryon

Two state fit for baryon mass

- $m_H = c_0 e^{-m_H t} (1 + c_1 e^{-Et})$ I64P30



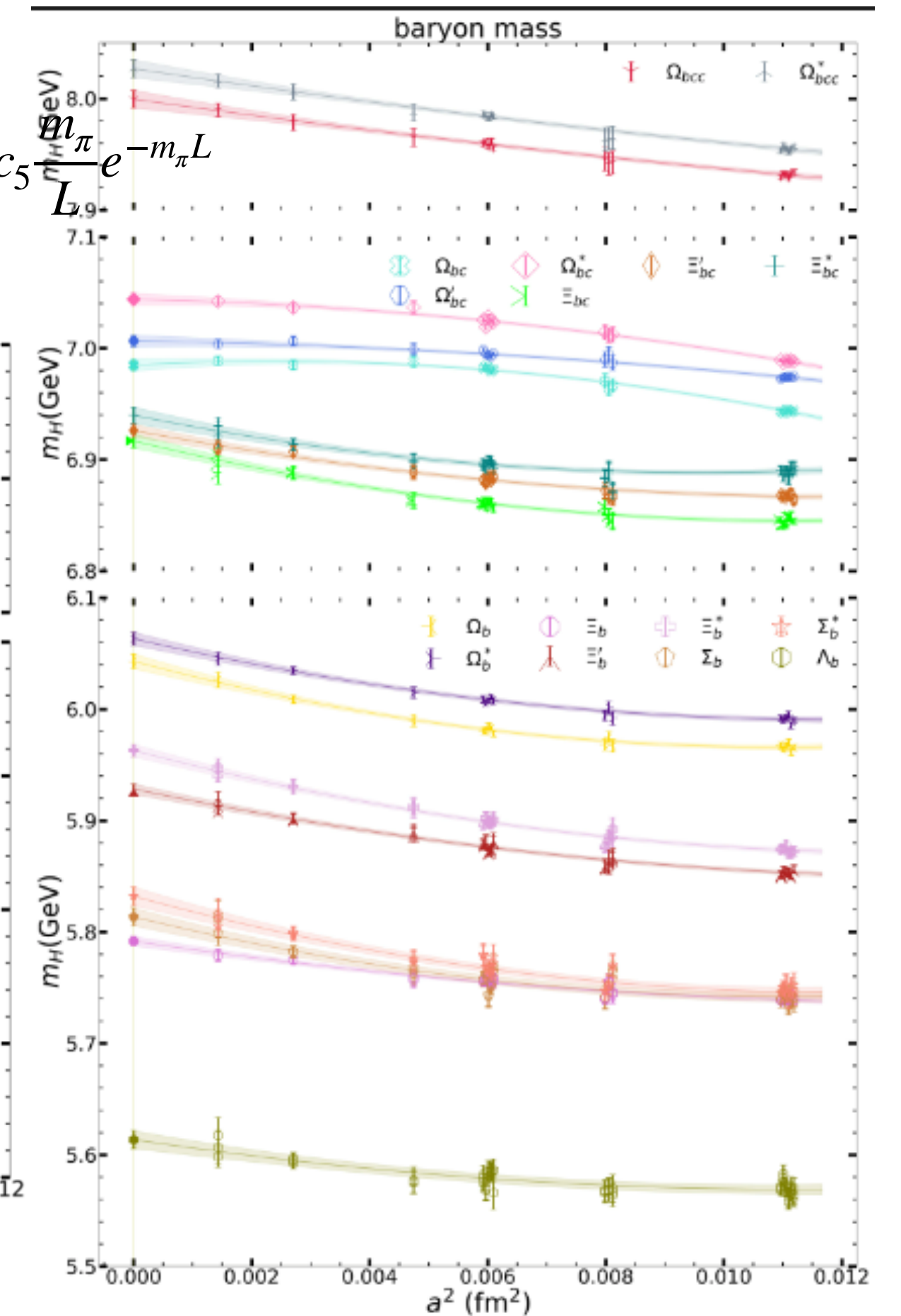
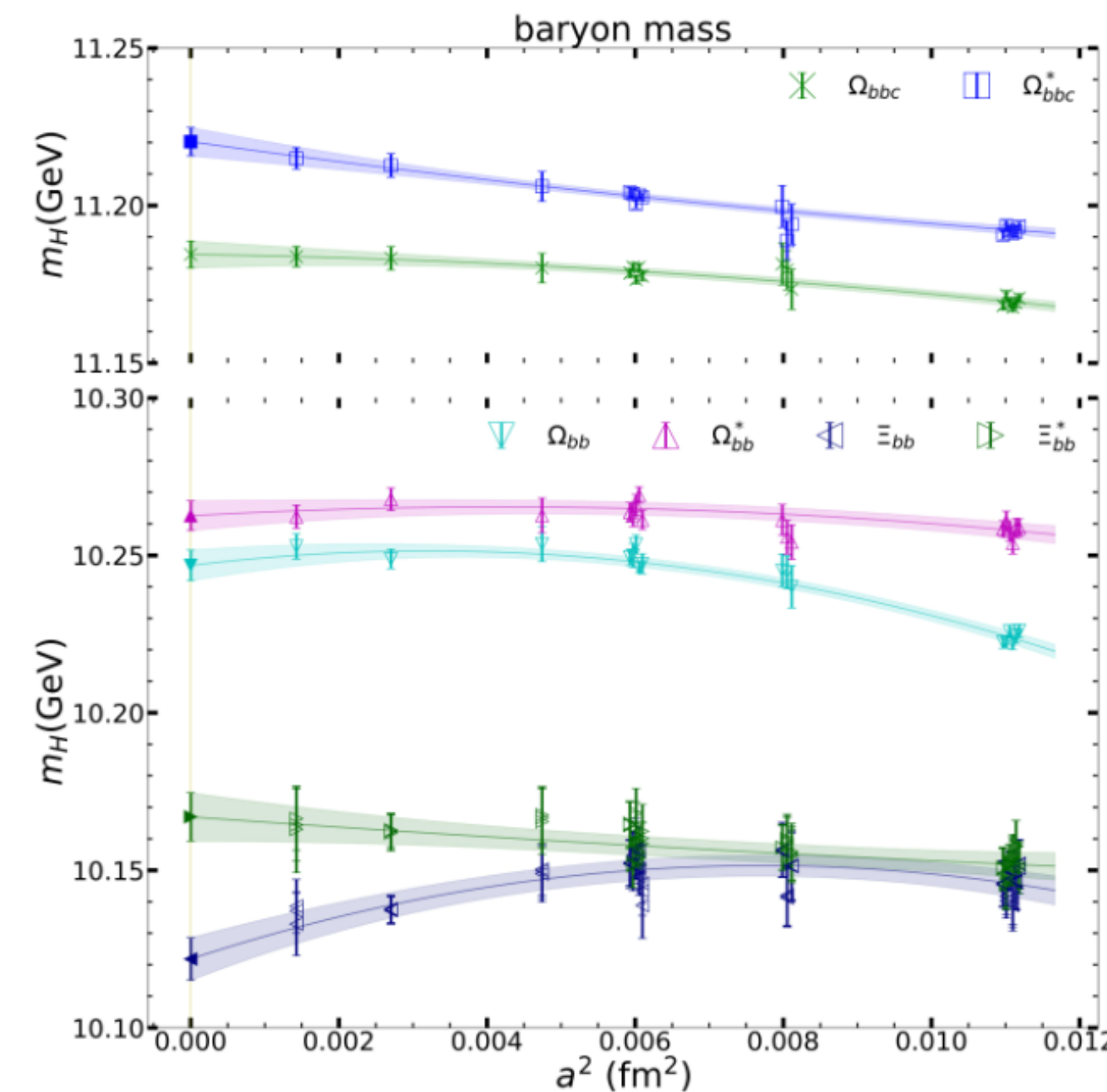
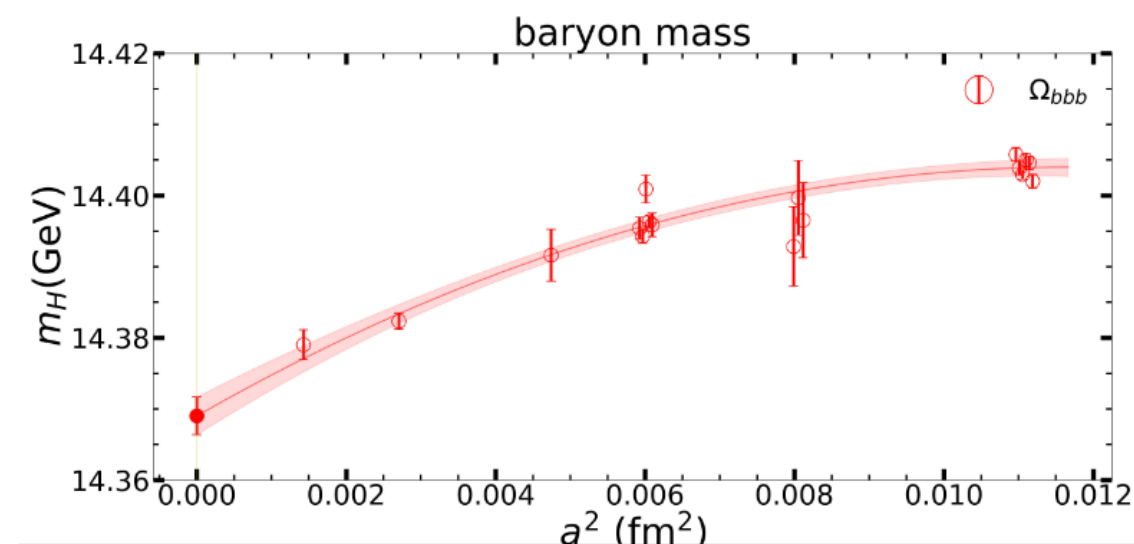
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Bottom Baryon

CLQCD preliminary

Global fit on baryon mass

- Global fit: $m_H(m_{\pi,\text{val}}, m_{\pi,\text{sea}}, m_{\eta_s,\text{sea}}, a) = m_H^{\text{phy}} + c_1(m_{\pi,\text{val}}^2 + m_{\pi,\text{sea}}^2 - m_{\pi,\text{phy}}^2) + c_2(m_{\eta_s,\text{sea}}^2 - m_{\eta_s,\text{phy}}^2) + c_3a^2 + c_4a^4 + c_5 \frac{m_\pi}{L} e^{-m_\pi L}$
- The joint fit interpolates valence, sea and strange masses to the physical point.
- Smooth lattice-spacing dependence enables a controlled continuum extrapolation.
- Consistent trends across channels support the preliminary mass determination.



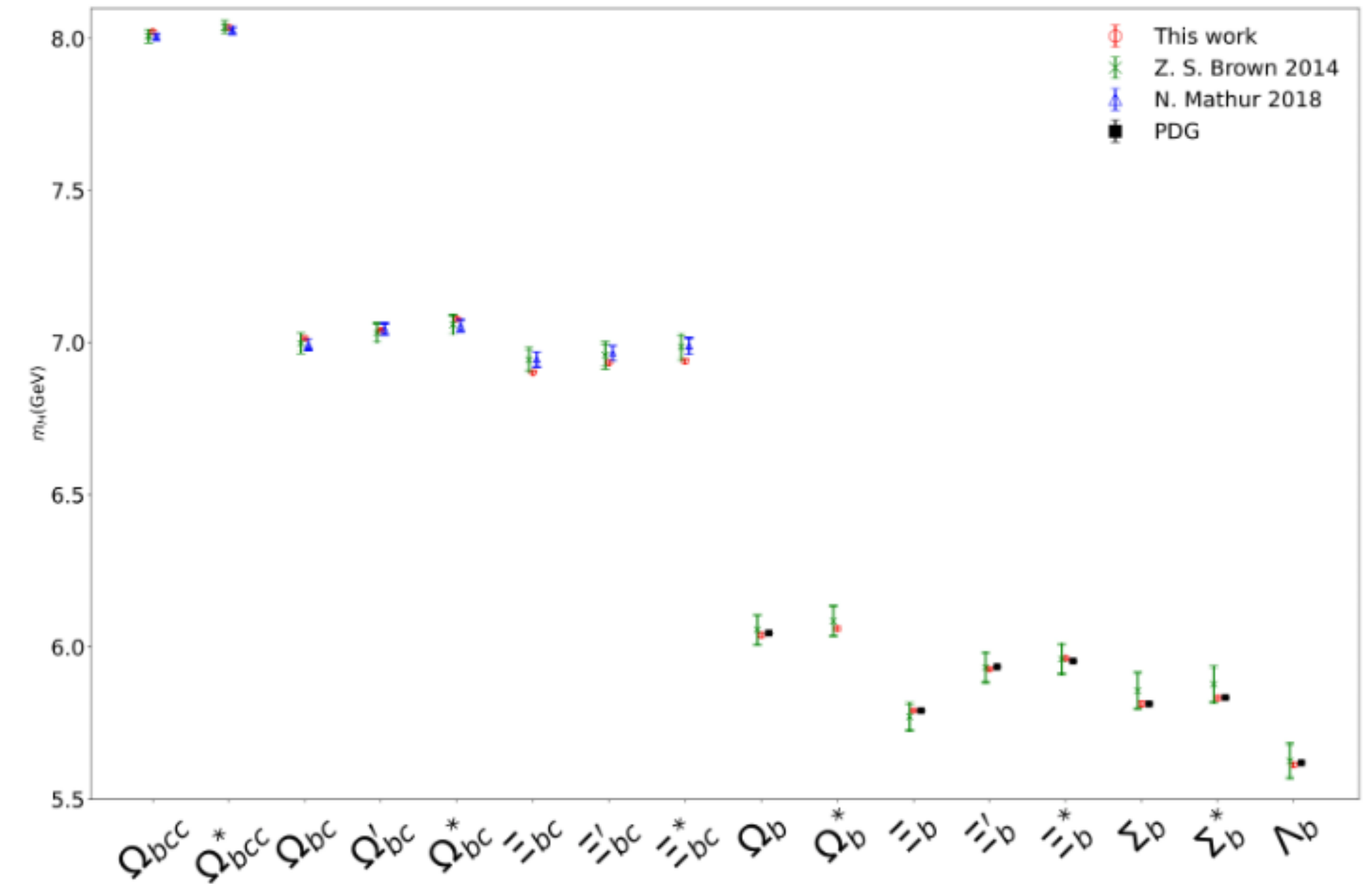
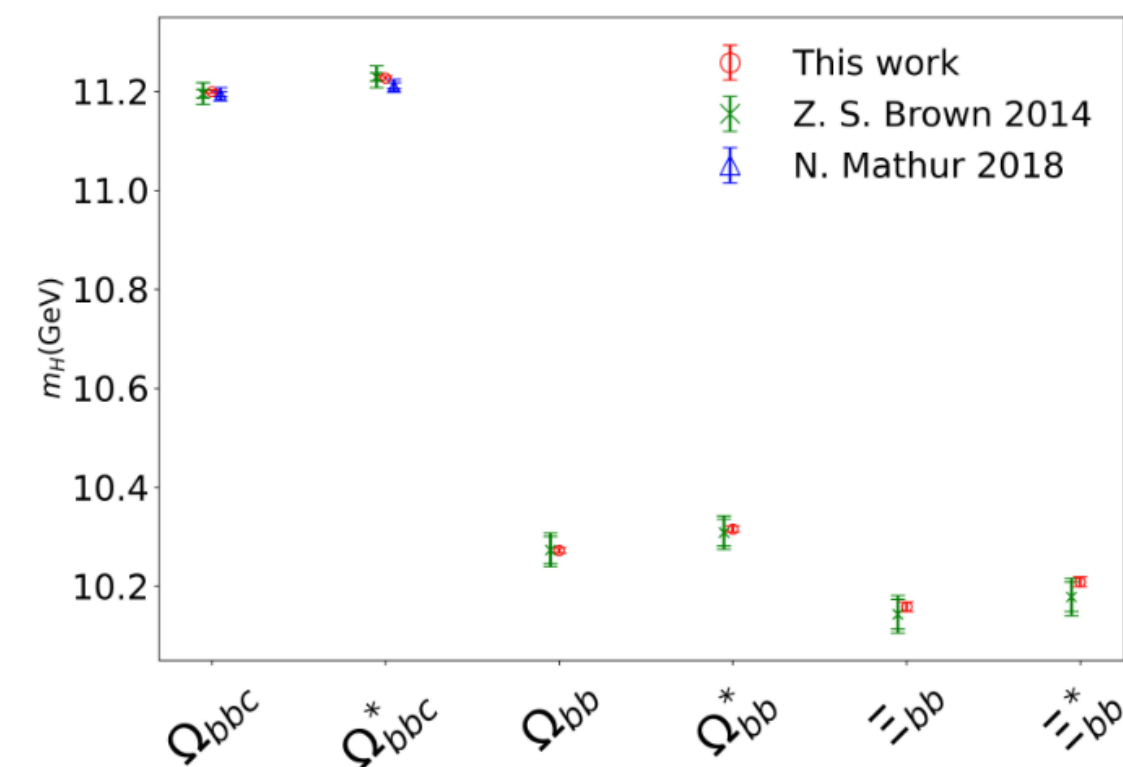
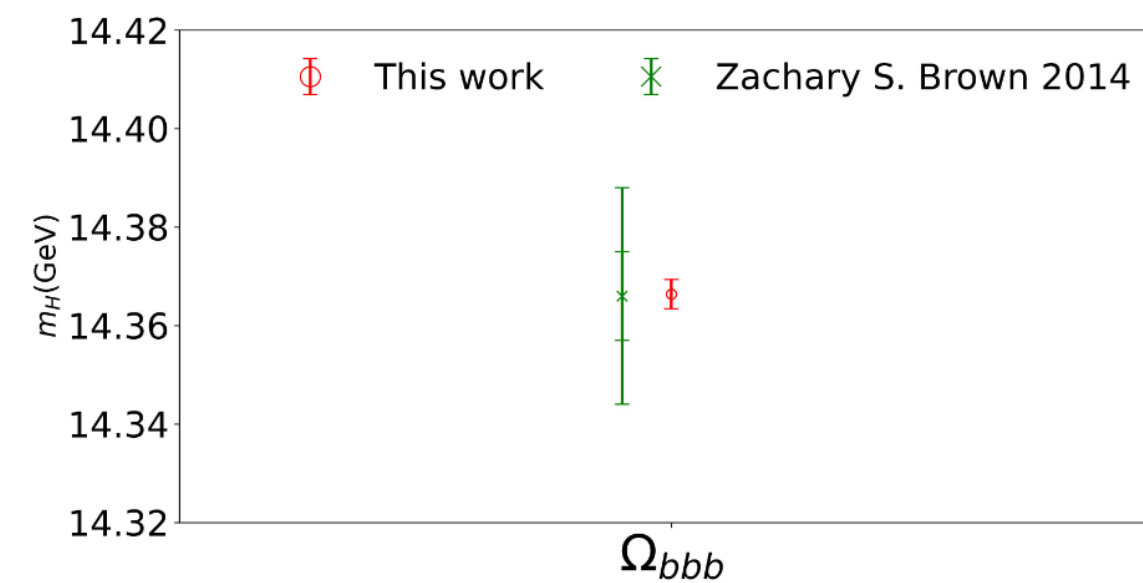
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Bottom Baryon

CLQCD preliminary

Bottom baryon mass results

- For the experimentally known single-bottom states, this work is consistent with the PDG values within the quoted uncertainties.
- The results are also compatible with previous lattice calculations, and our results have higher precision.
- The agreement across single-, double- and mixed-heavy channels supports the control of the physical-bottom tuning and continuum effects. The unobserved states therefore provide useful lattice predictions.



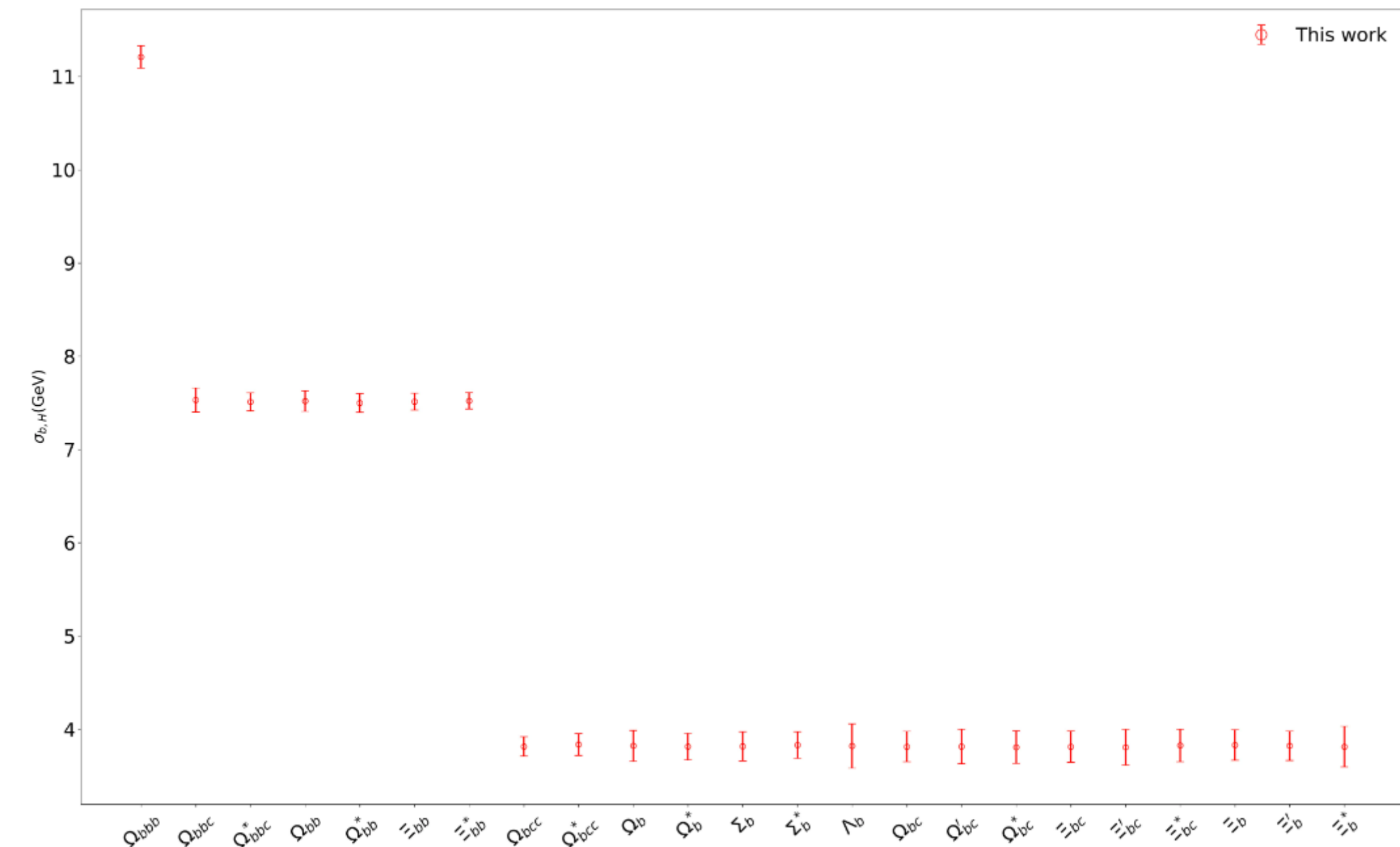
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Bottom Baryon

CLQCD preliminary

Sigma term calculation: valence bottom, charm and strange

- Extract valence bottom, charm and strange sigma term using Feynman-Hellman theorem $\sigma_{q,H}^{\text{FH}} = m_q^{\text{PC}} \frac{\partial m_H}{\partial m_q^{\text{PC}}}$ on each ensemble.
- Global fit: sea mass and lattice spacing dependence
 $\sigma_{b,H}(m_{\pi,\text{val}}, m_{\pi,\text{sea}}, m_{\eta_s,\text{sea}}, a) = \sigma_{b,H}^{\text{phy}} + c_1(m_{\pi,\text{sea}}^2 - m_{\pi,\text{phy}}^2) + c_2(m_{\eta_s,\text{sea}}^2 - m_{\eta_s,\text{phy}}^2) + c_3 a^2 + c_4 a^4 + c_5 e^{-m_\pi L}$
- Valence bottom contribution: $\sigma_{b,H}/n_b$: $\sim 3.74(4)$ GeV for thiple bottom quarks (Ω_{bbb}), $\sim 3.75(5)$ GeV for double bottom quarks and $\sim 3.82(17)$ GeV for single bottom quarks



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Bottom Baryon

Sigma term calculation: sea light and strange

- Sea mass contribution from quark mass: sea mass dependence of m_s^{val} , m_c^{val} and m_b^{val} lead to correction of sea light and sea strange sigma term.

$$\begin{aligned}\delta m_H &= \delta m_l^{\text{sea}} \langle \bar{l}l \rangle_H^{\text{sea}} + \delta m_s^{\text{sea}} \langle \bar{s}s \rangle_H^{\text{sea}} + \delta m_s^{\text{val}} \langle \bar{s}s \rangle_H^{\text{val}} + \delta m_c^{\text{val}} \langle \bar{c}c \rangle_H^{\text{val}} + \delta m_b^{\text{val}} \langle \bar{b}b \rangle_H^{\text{val}} \\ &= \delta m_l^{\text{sea}} (\langle \bar{l}l \rangle_H^{\text{sea}} + \frac{\partial m_s^{\text{val}}}{\partial m_l^{\text{sea}}} \langle \bar{s}s \rangle_H^{\text{val}} + \frac{\partial m_c^{\text{val}}}{\partial m_l^{\text{sea}}} \langle \bar{c}c \rangle_H^{\text{val}} + \frac{\partial m_b^{\text{val}}}{\partial m_l^{\text{sea}}} \langle \bar{b}b \rangle_H^{\text{val}}) + \delta m_s^{\text{sea}} (\langle \bar{s}s \rangle_H^{\text{sea}} + \frac{\partial m_s^{\text{val}}}{\partial m_s^{\text{sea}}} \langle \bar{s}s \rangle_H^{\text{val}} + \frac{\partial m_c^{\text{val}}}{\partial m_s^{\text{sea}}} \langle \bar{c}c \rangle_H^{\text{val}} + \frac{\partial m_b^{\text{val}}}{\partial m_s^{\text{sea}}} \langle \bar{b}b \rangle_H^{\text{val}})\end{aligned}$$

$$\begin{aligned}\sigma_{\pi,H} &= \sigma_{\pi,H}^{\text{val}} + \sigma_{\pi,H}^{\text{sea}} \\ \sigma_{\pi,H}^{\text{val}} &= m_l^{\text{phy}} \frac{\partial m_H(m_\pi^{\text{val}}, m_\pi^{\text{phy}}, m_{\eta_c}^{\text{phy}}, 0, 0)}{\partial m_l^{\text{val}}} \Big|_{m_l^{\text{val}}=m_l^{\text{phy}}} = (m_\pi^{\text{phy}})^2 \frac{\partial m_H(m_\pi^{\text{val}}, m_\pi^{\text{phy}}, m_{\eta_c}^{\text{phy}}, 0, 0)}{\partial (m_\pi^{\text{val}})^2} \Big|_{m_\pi^{\text{val}}=m_\pi^{\text{phy}}} \\ \sigma_{\pi,H}^{\text{sea}} &= m_l^{\text{phy}} \left[\frac{\partial m_H(m_\pi^{\text{phy}}, m_\pi^{\text{sea}}, m_{\eta_c}^{\text{phy}}, 0, 0)}{\partial m_l^{\text{sea}}} \Big|_{m_l^{\text{sea}}} - \frac{\partial m_s^{\text{val}}}{\partial m_l^{\text{sea}}} \langle \bar{s}s \rangle_H^{\text{val}} - \frac{\partial m_c^{\text{val}}}{\partial m_l^{\text{sea}}} \langle \bar{c}c \rangle_H^{\text{val}} - \frac{\partial m_b^{\text{val}}}{\partial m_l^{\text{sea}}} \langle \bar{b}b \rangle_H^{\text{val}} \right] \\ &= (m_\pi^{\text{phy}})^2 \left[\frac{\partial m_H(m_\pi^{\text{val}}, m_\pi^{\text{phy}}, m_{\eta_c}^{\text{phy}}, 0, 0)}{\partial (m_\pi^{\text{val}})^2} \Big|_{m_\pi^{\text{sea}}} - \frac{\partial m_s^{\text{val}}}{\partial (m_\pi^{\text{sea}})^2} \langle \bar{s}s \rangle_H^{\text{val}} - \frac{\partial m_c^{\text{val}}}{\partial (m_\pi^{\text{sea}})^2} \langle \bar{c}c \rangle_H^{\text{val}} - \frac{\partial m_b^{\text{val}}}{\partial (m_\pi^{\text{sea}})^2} \langle \bar{b}b \rangle_H^{\text{val}} \right] \\ \sigma_{s,H} &= \sigma_{s,H}^{\text{val}} + \sigma_{s,H}^{\text{sea}} \\ \sigma_{s,H}^{\text{val}} &= m_s^{\text{phy}} \frac{\partial m_H(m_\pi^{\text{sea}}, m_\pi^{\text{sea}}, m_{\eta_c}^{\text{phy}}, 0, 0)}{\partial m_s^{\text{val}}} \Big|_{m_s^{\text{val}}=m_s^{\text{phy}}} = (m_{\eta_c}^{\text{phy}})^2 \frac{\partial m_H(m_\pi^{\text{sea}}, m_\pi^{\text{sea}}, m_{\eta_c}^{\text{phy}}, 0, 0)}{\partial (m_{\eta_c}^{\text{val}})^2} \Big|_{m_{\eta_c}^{\text{val}}=m_{\eta_c}^{\text{phy}}} \\ \sigma_{s,H}^{\text{sea}} &= m_s^{\text{phy}} \left[\frac{\partial m_H(m_\pi^{\text{phy}}, m_\pi^{\text{phy}}, m_{\eta_c}^{\text{sea}}, 0, 0)}{\partial m_s^{\text{sea}}} \Big|_{m_s^{\text{sea}}} - \frac{\partial m_s^{\text{val}}}{\partial m_s^{\text{sea}}} \langle \bar{s}s \rangle_H^{\text{val}} - \frac{\partial m_c^{\text{val}}}{\partial m_s^{\text{sea}}} \langle \bar{c}c \rangle_H^{\text{val}} - \frac{\partial m_b^{\text{val}}}{\partial m_s^{\text{sea}}} \langle \bar{b}b \rangle_H^{\text{val}} \right] \\ &= (m_{\eta_c}^{\text{phy}})^2 \left[\frac{\partial m_H(m_\pi^{\text{phy}}, m_\pi^{\text{phy}}, m_{\eta_c}^{\text{sea}}, 0, 0)}{\partial (m_{\eta_c}^{\text{val}})^2} \Big|_{m_{\eta_c}^{\text{sea}}} - \frac{\partial m_s^{\text{val}}}{\partial (m_{\eta_c}^{\text{sea}})^2} \langle \bar{s}s \rangle_H^{\text{val}} - \frac{\partial m_c^{\text{val}}}{\partial (m_{\eta_c}^{\text{sea}})^2} \langle \bar{c}c \rangle_H^{\text{val}} - \frac{\partial m_b^{\text{val}}}{\partial (m_{\eta_c}^{\text{sea}})^2} \langle \bar{b}b \rangle_H^{\text{val}} \right] \\ \sigma_{c,H} &= \sigma_{c,H}^{\text{val}} \\ \sigma_{b,H} &= \sigma_{b,H}^{\text{val}}\end{aligned}$$

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Bottom Baryon

CLQCD preliminary

Sigma term contribution to baryon mass

- Total sigma term: sum over all flavor contributions, including valence and sea.
- Valence bottom, charm and strange quark sigma term approximately proportional to the number of quark in bottom baryon.
- Sea strange sigma term contribution: ~ 30 MeV
- Sea light sigma term contribution: ~ 10 MeV
- Total sigma term is the main component of the bottom baryon mass.
- Remain part: trace anomaly.

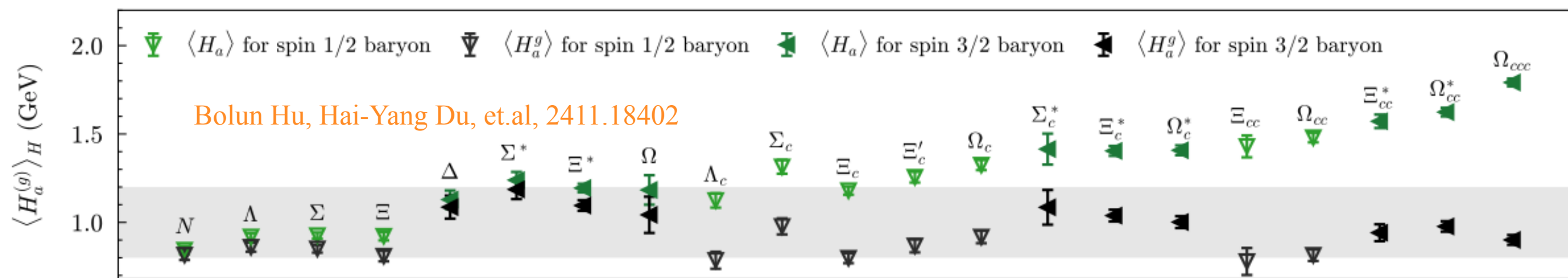
	m_H	σ_H	$\sigma_{b,H}$	$\sigma_{c,H}$	$\sigma_{s,H}^{\text{val}}$	$\sigma_{s,H}^{\text{sea}}$	$\sigma_{s,H}$	$\sigma_{\pi,H}^{\text{val}}$	$\sigma_{\pi,H}^{\text{sea}}$	$\sigma_{\pi,H}$
Ω_{bbb}	14.3664(30)	11.061(18)	11.023(16)	0	0	0.0212(72)	0.0212(72)	0	0.01588(38)	0.01588(38)
Ω_{bbc}	11.1991(30)	8.442(22)	7.446(20)	0.9584(65)	0	0.0254(55)	0.0254(55)	0	0.01136(28)	0.01136(28)
Ω_{bb^*c}	11.2270(36)	8.450(32)	7.442(28)	0.962(12)	0	0.0353(61)	0.0353(61)	0	0.01132(32)	0.01132(32)
Ω_{bb}	10.2720(55)	7.657(33)	7.510(30)	0	0.097(11)	0.037(10)	0.134(16)	0	0.01229(50)	0.01229(50)
Ω_{bb}^*	10.3156(62)	7.632(42)	7.485(39)	0	0.092(13)	0.042(11)	0.134(18)	0	0.01256(56)	0.01256(56)
Ξ_{bb}	10.158(10)	7.540(85)	7.481(83)	0	0	0.044(18)	0.044(18)	0.00467(38)	0.01065(97)	0.01531(92)
Ξ_{bb}^*	10.2086(98)	7.571(76)	7.520(74)	0	0	0.033(19)	0.033(19)	0.00423(32)	0.0134(10)	0.0176(10)
Ω_{bcc}	8.0228(48)	5.784(50)	3.782(47)	1.951(13)	0	0.0425(52)	0.0425(52)	0	0.00747(29)	0.00747(29)
Ω_{bcc}^*	8.0359(54)	5.739(62)	3.750(59)	1.926(14)	0	0.0560(65)	0.0560(65)	0	0.00789(35)	0.00789(35)
Ω_{bc}	7.0149(43)	5.017(47)	3.806(42)	1.041(17)	0.1077(82)	0.0541(67)	0.162(11)	0	0.00860(33)	0.00860(33)
Ω_{bc}'	7.0403(46)	4.976(39)	3.788(35)	1.030(12)	0.0971(94)	0.0522(69)	0.149(12)	0	0.00890(35)	0.00890(35)
Ω_{bc}^*	7.0777(52)	4.971(33)	3.789(29)	1.014(10)	0.1034(75)	0.0552(83)	0.159(11)	0	0.00938(42)	0.00938(42)
Ξ_{bc}	6.9023(68)	4.927(95)	3.804(91)	1.059(23)	0	0.051(11)	0.051(11)	0.00356(24)	0.00955(56)	0.01311(59)
Ξ_{bc}'	6.9348(67)	4.873(82)	3.780(78)	1.029(21)	0	0.053(11)	0.053(11)	0.00346(26)	0.00827(58)	0.01174(58)
Ξ_{bc}^*	6.9565(71)	4.813(90)	3.749(85)	1.002(23)	0	0.049(11)	0.049(11)	0.00356(25)	0.00925(60)	0.01281(62)
Ω_b	6.0390(79)	4.106(39)	3.809(35)	0	0.229(16)	0.058(13)	0.286(21)	0	0.01053(65)	0.01053(65)
Ω_b^*	6.0608(79)	4.112(31)	3.798(27)	0	0.233(11)	0.070(13)	0.303(18)	0	0.01112(65)	0.01112(65)
Ξ_b	5.7802(78)	3.98(10)	3.758(99)	0	0.190(18)	0.018(14)	0.207(23)	0.00568(18)	0.01045(62)	0.01614(65)
Ξ_b'	5.9332(79)	4.020(87)	3.813(85)	0	0.137(12)	0.054(12)	0.191(18)	0.00375(23)	0.01275(64)	0.01650(63)
Ξ_b^*	5.9561(78)	3.972(95)	3.781(93)	0	0.123(15)	0.050(12)	0.174(20)	0.00328(29)	0.01387(68)	0.01714(64)
Σ_b	5.813(11)	3.856(81)	3.803(78)	0	0	0.030(20)	0.030(20)	0.00747(43)	0.0151(10)	0.0226(10)
Σ_b^*	5.828(12)	3.796(95)	3.760(92)	0	0	0.012(22)	0.012(22)	0.00702(49)	0.0165(11)	0.0236(11)
Λ_b	5.625(12)	3.89(11)	3.83(11)	0	0	0.034(22)	0.034(22)	0.01513(52)	0.0079(12)	0.0230(12)

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Bottom Baryon

Trace anomaly contribution

- Hadron mass decomposition: $m_H = \sum_q \sigma_{q,H} + \gamma_m \sum_q \sigma_{q,H} + \frac{\beta(\alpha_s)}{2\alpha_s} \langle G^{\mu\nu} G_{\mu\nu} \rangle_H$
- Trace anomaly: $\langle H_a \rangle_H = m_H^{\text{phy}} - \sigma_H$, $\langle H_a^g \rangle_H = m_H^{\text{phy}} - (1 + \gamma_m) \sigma_H$. Scale dependence for both quark and gluon part.
- $\gamma_m = 0.295$ at $\overline{\text{MS}}$ 2 GeV: $\langle H_a \rangle_H \sim 14.4$ GeV for Ω_{bbb} , ~ 9.7 GeV for Ξ_{bb} , ~ 4.9 GeV for Σ_b .
 $\langle H_a^g \rangle_H$ at $\overline{\text{MS}}$ 2 GeV: ~ 0 GeV for Ω_{bbb} , ~ 0.4 GeV for Ξ_{bb} , ~ 0.9 GeV for Σ_b .
- Light and charmed baryon: gluon trace anomaly ~ 1 GeV at $\overline{\text{MS}}$ 2 GeV.

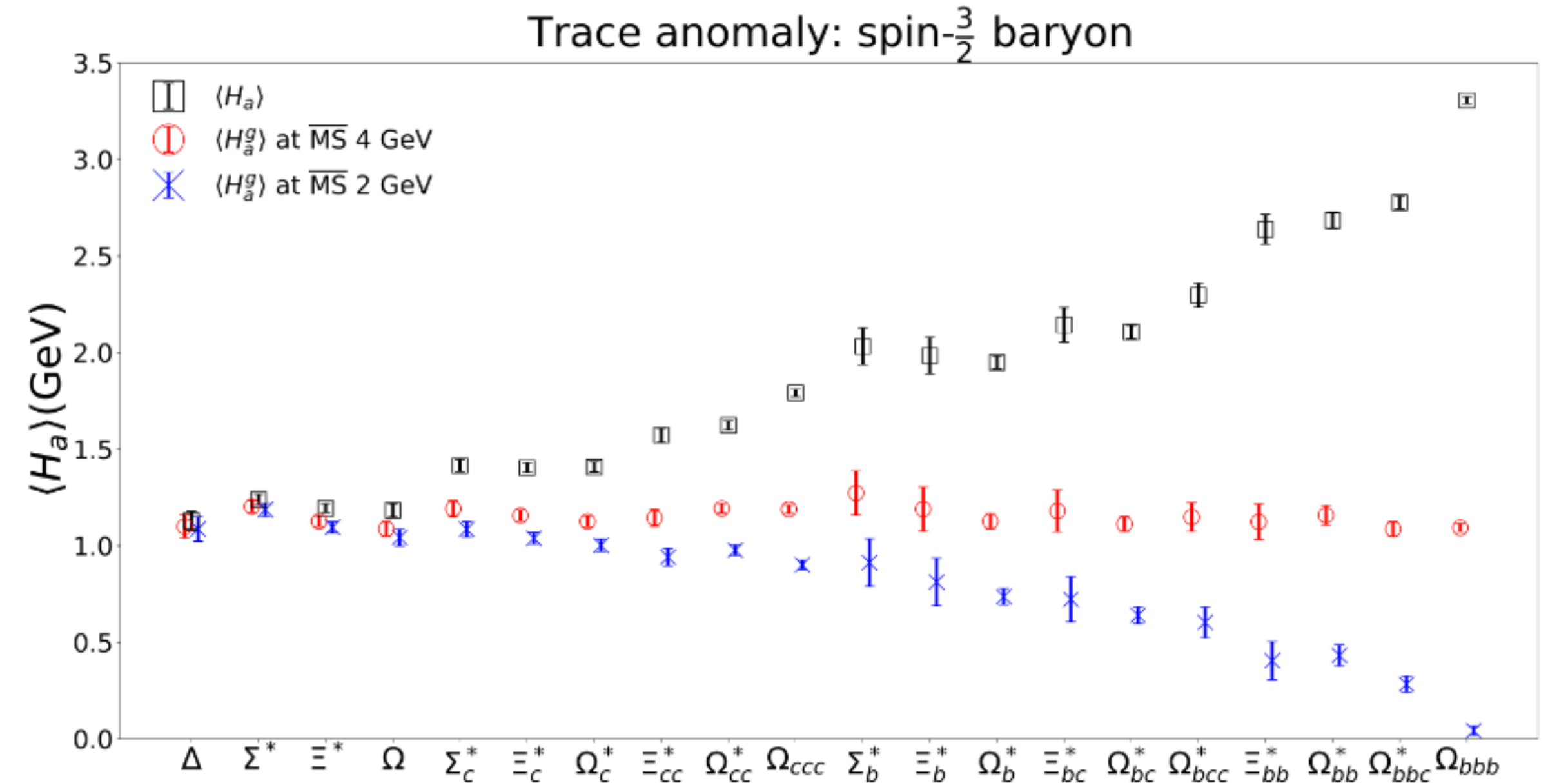
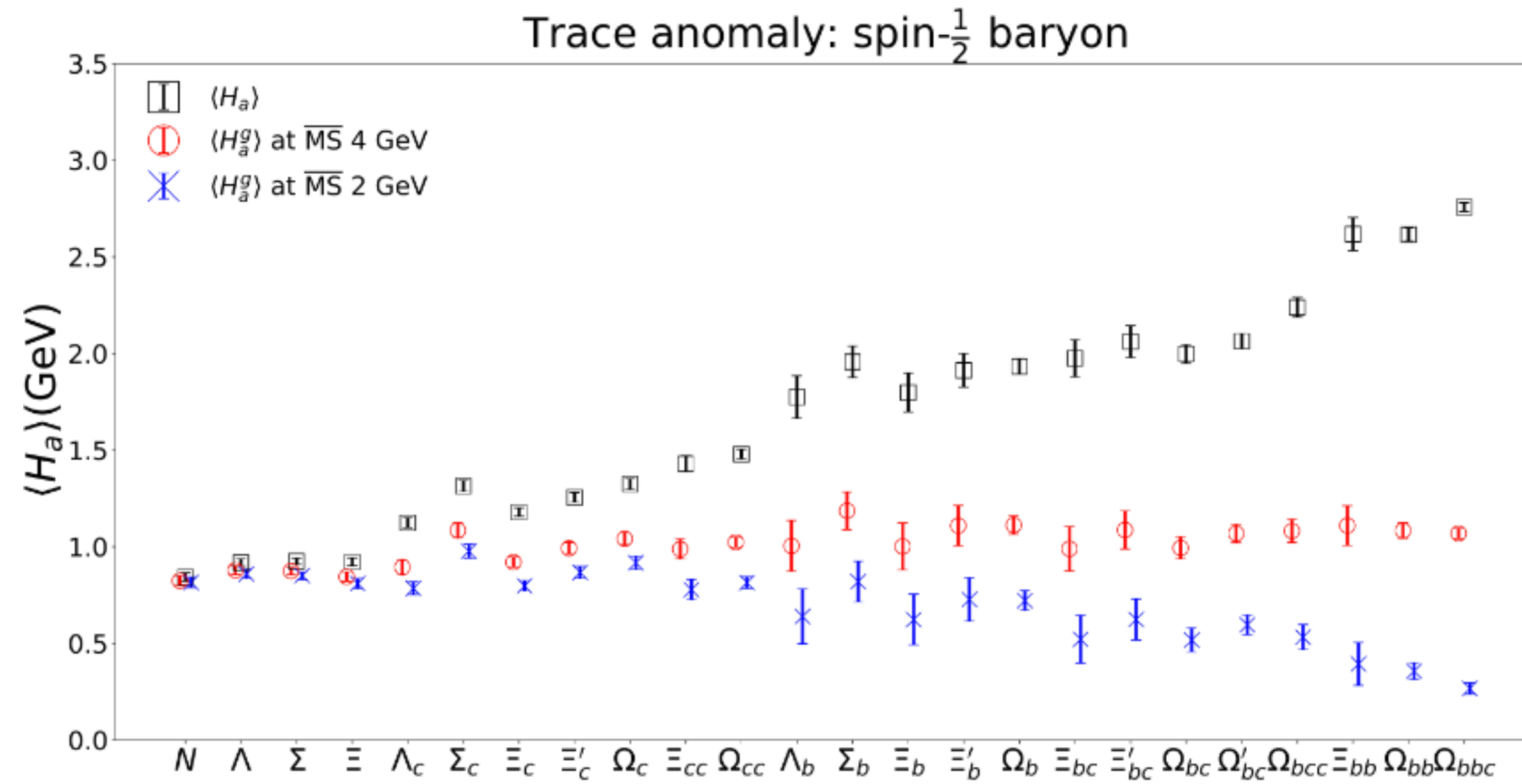


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Bottom Baryon

Scale dependence of trace anomaly

- $m_H = \sum_q \sigma_{q,H} + \gamma_m \sum_q \sigma_{q,H} + \frac{\beta(\alpha_s)}{2\alpha_s} \langle G^{\mu\nu} G_{\mu\nu} \rangle_H$, $\langle H_a \rangle_H = m_H^{\text{phy}} - \sigma_H$, $\langle H_a^g \rangle_H = m_H^{\text{phy}} - (1 + \gamma_m) \sigma_H$.
- $\gamma_m = 0.2$ at $\overline{\text{MS}}$ 4 GeV:



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Summary

- Anisotropic clover action for heavy quark: relativistic action for bottom quark, with discretization error better controlled.
- Bottom meson mass have good agreement with PDG, and given high precision decay constants.
- High precision lattice results on bottom baryon, and predicted the baryon mass with experimentally undiscovered.
- Sigma term contributions to bottom baryon mass.
- Trace anomaly contributions to bottom baryon mass, and its scale dependence.