(Electroweak) Pion and photon emission in a (chiral) effective field theory for nuclei (and beyond)

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For NuFact 2013 at IHEP, Beijing, China

# Outline

- Motivation: MiniBooNE low energy excess
- QHD EFT frame work: a quick but important introduction
- Pion and photon neutrinoproduction from nucleon
- Incoherent and coherent productions from nucleus, reaction kernel modification (s channel) in the medium
- MiniBooNE neutral current (NC) photon production and excess events: a short report
- Summary

#### MiniBooNE





#### 800 tons mineral oil (CH2)



Introduction to quantum hadrondynamics effective field theory (QHD EFT)

# • NN interactions (relativistic field theory since 1970):



B. Serot and J. Walecka, Adv. Nucl. Phys. **16**, 1 (1986)

- NN interactions (relativistic field theory)
- Mean-field approximation (RMF): for nuclear matter and mid-heavy nuclei; meson fields develop expectation values; nucleon spinorbital coupling...

Beyond (local) Fermi gas (LFG)

$$h_{\text{LS,T}} = \left[\frac{1}{4\overline{M}^2} \frac{1}{r} \left(\frac{d\Phi}{dr} + \frac{dW}{dr}\right) + \frac{f_v}{2M\overline{M}} \frac{1}{r} \frac{dW}{dr} + O(v^4)\right] \boldsymbol{\sigma} \cdot \boldsymbol{L}$$

- NN interactions (relativistic field theory)
- Mean-field approximation (RMF): for nuclear matter and mid--heavy nuclei; boson fields develop expectation values; nucleon spin-orbital coupling...
- Symmetries:

$$SU(2)_L \otimes SU(2)_R \otimes U(1)_R$$

- Lorentz; C, P, T, and the breakings
- Chiral symmetry and it's spontaneous breaking → isospin symmetry; pion dynamics; and electroweak currents (CVC and PCAC)

B. Serot and X.Z., arXiv:1011.5913; Advances in QFT (InTech, 2012) (arXiv:1110.2760) 6

- NN interactions (relativistic field theory)
- Mean-field approximation (RMF)
- Symmetries
- QHD can be used to study hadron behavior in medium: Delta modification; pion and nucleon optical potentials. A new development of loop calculations in QHD EFT.

Y. Hu, J. McIntire, and B. Serot (NPA 794:187, 2007)

#### Calibration at nucleon level: pion production



B. Serot and X.Z., Phys. Rev. C 86, 015501 (2012) (arXiv:1206.3812)

### Current form factors



 $F_2^{V,md} = \frac{1}{2} \left( 2\lambda^{(1)} - \frac{f_{\rho}g_{\rho}}{g_{\nu}} \frac{q^2}{q^2 - m_{\rho}^2} \right)$ 

$$\begin{split} \langle N, B | V_{\mu}^{i} | N, A \rangle &= \langle B | \frac{\tau^{i}}{2} | A \rangle \overline{u}_{f} \left( \gamma_{\mu} + 2\delta F_{1}^{V,md} \frac{q^{2} \gamma_{\mu} - \not{q} q_{\mu}}{q^{2}} + 2F_{2}^{V,md} \frac{\sigma_{\mu\nu} i q^{\nu}}{2M} \right) u_{i} \equiv \langle B | \frac{\tau^{i}}{2} | A \rangle \overline{u}_{f} \Gamma_{V\mu}(q) u_{i}, \\ \langle N, B; \pi, j, k_{\pi} | A_{\mu}^{i} | N, A \rangle &= -\frac{\epsilon^{i}_{jk}}{f_{\pi}} \langle B | \frac{\tau^{k}}{2} | A \rangle \overline{u}_{f} \gamma^{\nu} u_{i} \left[ g_{\mu\nu} + 2\delta F_{1}^{V,md} ((q - k_{\pi})^{2}) \frac{q \cdot (q - k_{\pi})g_{\mu\nu} - (q - k_{\pi})_{\mu} q_{\nu}}{(q - k_{\pi})^{2}} \right] \\ &- \frac{\epsilon^{i}_{jk}}{f_{\pi}} \langle B | \frac{\tau^{k}}{2} | A \rangle \overline{u}_{f} \frac{\sigma_{\mu\nu} i q^{\nu}}{2M} u_{i} \left[ 2\lambda^{(1)} + 2\delta F_{2}^{V,md} ((q - k_{\pi})^{2}) \frac{q \cdot (q - k_{\pi})}{(q - k_{\pi})^{2}} \right] \\ &= \frac{\epsilon^{i}_{jk}}{f_{\pi}} \langle B | \frac{\tau^{k}}{2} | A \rangle \overline{u}_{f} \Gamma_{A\pi\mu}(q, k_{\pi}) u_{i}. \qquad F_{1}^{V,md} = \frac{1}{2} \left( 1 + \frac{\beta^{(1)}}{M^{2}} q^{2} - \frac{g_{\rho}}{g_{\nu}} \frac{q^{2}}{q^{2} - m_{\rho}^{2}} \right) \end{split}$$

VMD. This relation will be used in high energy extrapolation.

#### Current form factors

 $\begin{aligned} \langle \Delta, a, p_{\Delta} | V^{i\mu}(A^{i\mu}) | N, A, p_N \rangle \\ &\equiv T_a^{\dagger i A} \, \overline{u}_{\Delta \alpha}(p_{\Delta}) \, \Gamma_{V(A)}^{\alpha \mu}(q) \, u_N(p_N). \end{aligned}$ 

$$\begin{split} \Gamma^{\alpha\mu}_A &= -h_A \left( g^{\alpha\mu} - \frac{q^{\alpha}q^{\mu}}{q^2 - m_{\pi}^2} \right) \\ &+ \frac{2d_{2\Delta}}{M^2} \left( q^{\alpha}q^{\mu} - g^{\alpha\mu}q^2 \right) \end{split}$$

$$-\frac{2d_{4\Delta}}{M}(q^{\alpha}\gamma^{\mu}-g^{\alpha\mu}\not{q})-\frac{4d_{7\Delta}}{M^{2}}q^{\alpha}\sigma^{\mu\nu}iq_{\mu}$$



$$h_A(q^2) \equiv h_A + h_{\Delta a_1} \frac{q^2}{q^2 - m_{a_1}^2}$$

$$d_{i\Delta}(q^2) \equiv d_{i\Delta} + d_{i\Delta a_1} \frac{q^2}{q^2 - m_{a_1}^2}$$

#### Calibration at nucleon level: NC photon prod.



# **Results?**

$$\frac{c_1}{M^2} \overline{N} \gamma^{\mu} N \operatorname{Tr}(\widetilde{a}^{\nu} \overline{F}^{(+)}_{\mu\nu}), \quad \frac{e_1}{M^2} \overline{N} \gamma^{\mu} \widetilde{a}^{\nu} N \overline{f}_{s\mu\nu}$$

Already in the previous lagrangian. Also related to electro(photo) pion prod.

R. J. Hill, Phys. Rev. D 81, 013008 (2010)



K. Graczyk, D. Kiełczewska, P. Przewłocki, and J. Sobczyk, PRD 80, 093001 (2009).
E. Hern´andez, J. Nieves, and M. Valverde, PRD 76, 033005 (2007).
G. M. Radecky et al., PRD 25, 1161 (1982); T. Kitagaki et al., PRD 34, 2554 (1986).



#### **Power counting of the calculation**



# Incoherent pion and photon production from Nucleus

Final nucleon wave function, final state interaction (FSI), optical potential

Initial nucleon (shell) wave function



# Incoherent pion and photon production from Nucleus



X.Z. and B. Serot, Phys. Rev. C 86, 035502 (2012) (arXiv: 1206.6324)

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### Delta dynamics in nuclear medium

#### • Self energy: real part $\rightarrow$ spin-orbital coupling in nucleus

$$\mathcal{L}_{\Delta;\pi,\rho,V,\phi} = \frac{-i}{2} \overline{\Delta}^a_\mu \left\{ \sigma^{\mu\nu} , \left( i \ \widetilde{\partial} - h_\rho \ \phi - h_v \ V - m + h_s \phi \right) \right\}^b_a \Delta_{b\nu}$$

$$p_{\Delta}^{0} = h_{v} \langle V^{0} \rangle + \sqrt{m^{*2} + \vec{p}_{\Delta}^{2}}$$

$$\equiv h_{v} \langle V^{0} \rangle + p_{\Delta}^{*0}$$

$$= h_{v} \langle V^{0} \rangle + \sqrt{m^{*2} + \vec{p}_{\Delta}^{*2}},$$

$$m^{*} \equiv m - h_{s} \langle \phi \rangle.$$

$$h_{\Delta} = \frac{1}{3} \left[ \frac{1}{2\overline{m}^{2} r} \frac{d}{dr} \left( h_{s} \langle \phi \rangle + h_{v} \langle V^{0} \rangle \right) - \frac{\widetilde{f}_{v}}{m\overline{m} r} \frac{d}{dr} \left( h_{v} \langle V^{0} \rangle \right) \right] \vec{S} \cdot \vec{L}$$

$$\equiv \alpha(r) \vec{S} \cdot \vec{L}.$$

$$17$$



Y. Horikawa, M. Thies, and F. Lenz, NPA **345**, 386 (1980). S. X. Nakamura, T. Sato, T.-S. H. Lee, B. Szczerbinska, and K. Kubodera, PRC **81**, 035502 (2010).

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#### Delta dynamics in nuclear medium

Self energy: real part → spin-orbital coupling in nucleus
 Self energy: imaginary part; collision broadening

$$\Gamma_{\Delta} = \Gamma_{\pi} + \Gamma_{\rm sp} ,$$
  
$$\Gamma_{\rm sp} = V_0 \times \frac{\rho(r)}{\rho(0)}$$

$$V_0 \approx 80 \,\,\mathrm{MeV}$$

E. Oset and L. Salcedo, NPA **468**, 631 (1987)

### • Check: **incoherent** electro-production of pion from C12. $\Gamma_{\Delta} \rightarrow 120 \text{ MeV} + 40 \text{ MeV}$



E. Oset and L. Salcedo, NPA 468, 631 (1987), P. Barreau et al., NPA 402, 515 (1983).

# • Check: **incoherent** electro-production of pion from C12.



T. W. Donnelly (private communication).

# Incoherent neutrinoproduction of pion from C12



### Incoh. neutrinoprod. of pion from C12



## Incoh. neutrinoprod. of photon from C12



# Coherent production of pion



# Coherent production of pion

#### • "Optimal" approximation (factorization):

$$\begin{aligned} \frac{1}{m_A} \langle A, \pi(\vec{k}_{\pi}) | J_{had}^{\mu} | A \rangle \\ \approx \begin{cases} \int_A d\vec{r} e^{i(\vec{q} - \vec{k}_{\pi}) \cdot \vec{r}} \langle J_{had}^{\mu}(\vec{q}, \vec{k}_{\pi}, \vec{r}) \rangle & \text{PW}, \\ \int_A d\vec{r} e^{i(\vec{q} - \vec{k}_{\pi}) \cdot \vec{r}} e^{-i\int_z^{\infty} \frac{\Pi(\rho, l)}{2|\vec{k}_{\pi}|} dl} \langle J_{had}^{\mu}(\vec{q}, \vec{k}_{\pi}, \vec{r}) \rangle & \text{DW}. \end{cases} \\ & \langle J_{had}^{\mu}(\vec{q}, \vec{k}_{\pi}, \vec{r}) \rangle \approx \rho_n(\vec{r}) \frac{1}{2} \sum_{s_z} \frac{1}{p_{ni}^{*0}} \langle n, s_z, \frac{\vec{q} - \vec{k}_{\pi}}{2} | J_{had}^{\mu}(\vec{q}, \vec{k}_{\pi}) | n, s_z, \frac{\vec{k}_{\pi} - \vec{q}}{2} \rangle \\ & + \rho_p(\vec{r}) \frac{1}{2} \sum_{s_z} \frac{1}{p_{ni}^{*0}} \langle p, s_z, \frac{\vec{q} - \vec{k}_{\pi}}{2} | J_{had}^{\mu}(\vec{q}, \vec{k}_{\pi}) | p, s_z, \frac{\vec{k}_{\pi} - \vec{q}}{2} \rangle . \end{aligned}$$

PRC 86,

#### • Check: photo-production of pions from C12.



M. Schmitz, Ph.D. thesis, Johannes Gutenberg Universit<sup>®</sup> at Mainz, 1996. W. Peters, H. Lenske, and U. Mosel, NPA **640**, 89 (1998).

# Coh. neutrinoprod. of pion from C12



Dr. Geralyn Zeller (private communication).

### Coh. neutrinoprod. of photon from C12



# MiniBooNE NC photon

X.Z. and B. Serot, PLB 719, 409 (2013) (arXiv: 1210.3210)

Extrapolation of previous results to higher energy



$$\frac{c_1}{M^2}\overline{N}\gamma^{\mu}N\operatorname{Tr}(\widetilde{a}^{\nu}\overline{F}^{(+)}_{\mu\nu}), \quad \frac{e_1}{M^2}\overline{N}\gamma^{\mu}\widetilde{a}^{\nu}N\overline{f}_{s\mu\nu}$$



#### MiniBooNE NC photon

Coherent production



#### MiniBooNE NC photon events

$E_{QE}(GeV)$	[0.2, 0.3]	[0.3, 0.475]	[0.475, 1.25]
coh	1.5 (2.9)	6.0 (9.2)	2.1 (8.0)
inc	12.0 (14.1)	25.5 (31.1)	12.6 (23.2)
Н	4.1 (4.4)	10.6 (11.6)	4.6 (6.3)
Total	17.6 (21.4)	42.1 (51.9)	19.3 (37.5)
MiniBN	19.5	47.3	19.4
Excess	$42.6 \pm 25.3$	$82.2 \pm 23.3$	$21.5\pm34.9$
$E_{QE}(GeV)$	[0.2, 0.3]	[0.3, 0.475]	[0.475, 1.25]
coh	1.0 (2.2)	3.1 (5.5)	0.87 (5.4)
inc	4.5 (5.3)	10.0 (12.2)	4.0 (10.2)
Н	1.3 (1.6)	3.6 (4.3)	1.1 (2.4)
Total	6.8 (9.1)	16.7 (22.0)	6.0 (18.0)
MiniBN	8.8	16.9	6.8
Excess	$34.6 \pm 13.6$	$23.5 \pm 13.4$	$20.2\pm22.8$

# Summary

- QHD EFT→ a unified frame work for studying nuclear structure, EW response, nucleon, pion and Delta behavior in the medium
- Study EW pion and photon prod. in the QHD EFT, by using "LFG" (incoh.) and optimal factorization (coh.) approx.
- Extrapolate the EFT results to the 1—2 GeV lepton energy region; the kernel is from the EFT calculation.
- Calculate NC photons at MiniBooNE: the low energy excess can not be fully explained as NC photons.
- Interesting things to be done: a systematic study of Delta and pion dynamics in QHD EFT; go beyond "LFG"; treat the pion and photon prods., and quasi elastic scattering in this EFT.

# Back up

#### Where Are the Pions?

• For nuclear equation of state (EOS), 1- and 2-loop calculations (including pion) are done by Y. Hu, J. McIntire, and B. Serot (NPA 794:187, 2007); Infrared Regularization.



# Spin-3/2 Particle in EFT

• Redundant degrees of freedom in Rarita-Schwinger representation ( $\psi^{\mu}$ ) do NOT show up.

$$S_{F} = (S_{F}^{0(\frac{3}{2})} + S_{F}^{0(\frac{3}{2}\perp)}) + (S_{F}^{0(\frac{3}{2})} + S_{F}^{0(\frac{3}{2}\perp)})(\Sigma^{(\frac{3}{2})} + \Sigma^{(\frac{3}{2}\perp)})(S_{F}^{0(\frac{3}{2})} + S_{F}^{0(\frac{3}{2}\perp)}) + \cdots$$

$$= S_{F}^{0(\frac{3}{2})} + S_{F}^{0(\frac{3}{2})}\Sigma^{(\frac{3}{2})}S_{F}^{0(\frac{3}{2})} + \cdots$$
This can be generalized
$$= S_{F}^{0(\frac{3}{2}\perp)} + S_{F}^{0(\frac{3}{2}\perp)}\Sigma^{(\frac{3}{2}\perp)}S_{F}^{0(\frac{3}{2}\perp)} + \cdots$$
This con be generalized
to other spins

Related work: V. Pascalutsa, PRD 58: 096002, 1998; V. P and D. Phillips, PRC 67: 055202, 2003; H. Krebs, E. Epelbaum, and U. Meissner, PRC 80: 028201, 2009; PLB 683: 222, 2010

# Spin-3/2 Particle in EFT

- Redundant degrees of freedom in Rarita-Schwinger representation (  $\psi^{\mu}$  ) do NOT show up.
- Off-shell couplings:  $\gamma_{\mu}\psi^{\mu}$ ,  $\partial_{\mu}\psi^{\mu}$ ,  $\overline{\psi}^{\mu}\gamma_{\mu}$ , and  $\partial_{\mu}\overline{\psi}^{\mu}$





### Kernel



(e) is small
(f) are the anomalous terms with FFs included.

$$\frac{c_1}{M^2} \overline{N} \gamma^{\mu} N \operatorname{Tr}(\widetilde{a}^{\nu} \overline{F}^{(+)}_{\mu\nu}), \quad \frac{e_1}{M^2} \overline{N} \gamma^{\mu} \widetilde{a}^{\nu} N \overline{f}_{s\mu\nu}$$

# Benchmarks: pion prod.



FFs included in an EFT inspired way: CVC and PCAC.
 K-R, and anomalous diagrams included.

- NN interactions (relativistic field theory)
- Mesons nonlinear interactions



*Three body force* 

• Chiral symmetry in QCD:  $SU(2)_L \otimes SU(2)_R \otimes U(1)_B$ 

$$\mathcal{L} = \mathcal{L}_{QCD} + \overline{q}\gamma_{\mu}(\mathbf{v}^{\mu} + B\mathbf{v}^{\mu}_{(s)} + \gamma_{5}\mathbf{a}^{\mu})q - \overline{q}(\mathbf{s} - i\gamma_{5}\mathbf{p})q$$

$$q_{LA} \to \exp\left[-i\frac{\theta(x)}{3}\right] \left(\exp\left[-i\theta_{Li}(x)\frac{\tau^{i}}{2}\right]\right)_{A}^{B} q_{LB} \equiv \exp\left[-i\frac{\theta(x)}{3}\right] (L)_{A}^{B} q_{LB},$$
$$q_{R} \to \exp\left[-i\frac{\theta(x)}{3}\right] \exp\left[-i\theta_{Ri}(x)\frac{\tau^{i}}{2}\right] q_{R} \equiv \exp\left[-i\frac{\theta(x)}{3}\right] Rq_{R},$$

B. Serot and X.Z., arXiv:1011.5913; Advances in QFT (InTech, 2012) (arXiv:1110.2760) 41

• Chiral symmetry in QCD:

$$q_{LA} \to \exp\left[-i\frac{\theta(x)}{3}\right] \left(\exp\left[-i\theta_{Li}(x)\frac{\tau^{i}}{2}\right]\right)_{A}^{B} q_{LB} \equiv \exp\left[-i\frac{\theta(x)}{3}\right] (L)_{A}^{B} q_{LB}$$
$$q_{R} \to \exp\left[-i\frac{\theta(x)}{3}\right] \exp\left[-i\theta_{Ri}(x)\frac{\tau^{i}}{2}\right] q_{R} \equiv \exp\left[-i\frac{\theta(x)}{3}\right] Rq_{R},$$

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Its nonlinear realization at low energy EFT:

$$\begin{split} U &\equiv \exp\left[2i\frac{\pi_i(x)}{f\pi}t^i\right] \to LUR^{\dagger}, \\ \xi &\equiv \sqrt{U} = \exp\left[i\frac{\pi_i}{f\pi}t^i\right] \to L\xi h^{\dagger} = h\,\xi R^{\dagger}, \\ \widetilde{v}_{\mu} &\equiv \frac{-i}{2}[\xi^{\dagger}(\partial_{\mu} - il_{\mu})\xi + \xi(\partial_{\mu} - ir_{\mu})\xi^{\dagger}] \equiv \widetilde{v}_{i\mu}t^i \to h\,\widetilde{v}_{\mu}h^{\dagger} - ih\,\partial_{\mu}h^{\dagger}, \\ \widetilde{a}_{\mu} &\equiv \frac{-i}{2}[\xi^{\dagger}(\partial_{\mu} - il_{\mu})\xi - \xi(\partial_{\mu} - ir_{\mu})\xi^{\dagger}] \equiv \widetilde{a}_{i\mu}t^i \to h\,\widetilde{a}_{\mu}h^{\dagger}, \end{split}$$

- Chiral symmetry in QCD:
- Its nonlinear realization at low energy EFT:

$$\begin{split} (\widetilde{\partial}_{\mu}\psi)_{\alpha} &\equiv (\partial_{\mu} + i\,\widetilde{v}_{\mu} - i\mathsf{v}_{(s)\mu}B)^{\beta}_{\alpha}\psi_{\beta} \to \exp\left[-i\theta(x)B\right]h^{\beta}_{\alpha}(\widetilde{\partial}_{\mu}\psi)_{\beta}, \\ \widetilde{v}_{\mu\nu} &\equiv -i[\widetilde{a}_{\mu},\,\widetilde{a}_{\nu}] \to h\,\widetilde{v}_{\mu\nu}h^{\dagger}, \\ F^{(+)}_{\mu\nu} &\equiv \xi^{\dagger}f_{L\mu\nu}\,\xi + \xi f_{R\mu\nu}\,\xi^{\dagger} \to hF^{(+)}_{\mu\nu}h^{\dagger}, \\ F^{(-)}_{\mu\nu} &\equiv \xi^{\dagger}f_{L\mu\nu}\,\xi - \xi f_{R\mu\nu}\,\xi^{\dagger} \to hF^{(-)}_{\mu\nu}h^{\dagger}, \end{split}$$

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#### • The lagrangian, baryon section:

$$\mathcal{L}_{N(\hat{\nu}\leqslant3)} = \overline{N} \Big( i\gamma^{\mu} [\widetilde{\partial}_{\mu} + ig_{\rho}\rho_{\mu} + ig_{v}V_{\mu}] + g_{A}\gamma^{\mu}\gamma^{5} \widetilde{a}_{\mu} - M + g_{s}\phi \Big) N$$
  
$$- \frac{f_{\rho}g_{\rho}}{4M} \overline{N}\rho_{\mu\nu}\sigma^{\mu\nu}N - \frac{f_{v}g_{v}}{4M} \overline{N}V_{\mu\nu}\sigma^{\mu\nu}N - \frac{\kappa_{\pi}}{M} \overline{N}\widetilde{v}_{\mu\nu}\sigma^{\mu\nu}N + \frac{4\beta_{\pi}}{M} \overline{N}N \operatorname{Tr}(\widetilde{a}_{\mu}\widetilde{a}^{\mu}) \Big)$$

$$\mathcal{L}_{\Delta} = \frac{-i}{2} \overline{\Delta}^{a}_{\mu} \{ \sigma^{\mu\nu} , (i \,\widetilde{\partial} - h_{\rho} \, \not\!\!{\rho} - h_{v} \, \not\!\!{V} - m + h_{s} \phi) \}^{b}_{a} \Delta_{b\nu} + \widetilde{h}_{A} \overline{\Delta}^{a}_{\mu} \, \widetilde{a}^{b}_{a} \gamma^{5} \Delta^{\mu}_{b}$$

$$\mathcal{L}_{\Delta,N,\pi} = h_A \overline{\Delta}^{a\mu} T_a^{\dagger iA} \widetilde{a}_{i\mu} N_A + \text{c.c.},$$

•

• The lagrangian, baryon section.

• The lagrangian, meson section:

$$\mathcal{L}_{\text{meson}(\hat{\nu} \leqslant 4)} = \frac{1}{2} \partial_{\mu} \phi \, \partial^{\mu} \phi + \frac{1}{4} f_{\pi}^{2} \operatorname{Tr}[\widetilde{\partial}_{\mu} U(\widetilde{\partial}^{\mu} U)^{\dagger}] \\ + \frac{1}{4} f_{\pi}^{2} m_{\pi}^{2} \operatorname{Tr}(U + U^{\dagger} - 2) \\ - \frac{1}{2} \operatorname{Tr}(\rho_{\mu\nu} \rho^{\mu\nu}) - \frac{1}{4} V^{\mu\nu} V_{\mu\nu} \\ + \frac{1}{2g_{\gamma}} \left( \operatorname{Tr}(F^{(+)\mu\nu} \rho_{\mu\nu}) + \frac{1}{3} f_{s}^{\mu\nu} V_{\mu\nu} \right) \right)$$

Vector meson dominance (VMD)

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#### • Electroweak (EW) interactions:

$$\mathcal{L}_{\Delta,N,bg} = \frac{ic_{1\Delta}}{M} \overline{\Delta}^{a}_{\mu} \gamma_{\nu} \gamma^{5} T_{a}^{\dagger iA} F_{i}^{(+)\mu\nu} N_{A} + \frac{ic_{3\Delta}}{M^{2}} \overline{\Delta}^{a}_{\mu} i\gamma^{5} T_{a}^{\dagger iA} (\widetilde{\partial}_{\nu} F^{(+)\mu\nu})_{i} N_{A} + \frac{c_{6\Delta}}{M^{2}} \overline{\Delta}^{a}_{\lambda} \sigma_{\mu\nu} T_{a}^{\dagger iA} (\widetilde{\partial}^{\lambda} \overline{F}^{(+)\mu\nu})_{i} N_{A} - \frac{id_{4\Delta}}{M} \overline{\Delta}^{a}_{\mu} \gamma_{\nu} T_{a}^{\dagger iA} F_{i}^{(-)\mu\nu} N_{A} - \frac{id_{7\Delta}}{M^{2}} \overline{\Delta}^{a}_{\lambda} \sigma_{\mu\nu} T_{a}^{\dagger iA} (\widetilde{\partial}^{\lambda} F^{(-)\mu\nu})_{i} N_{A} + \text{c.c.}$$
$$\mathcal{L}_{\Delta,N,\rho} = \frac{ic_{1\Delta\rho}}{M} \overline{\Delta}^{a}_{\mu} \gamma_{\nu} \gamma^{5} T_{a}^{\dagger iA} \rho_{i}^{\mu\nu} N_{A} + \frac{ic_{3\Delta\rho}}{M^{2}} \overline{\Delta}^{a}_{\mu} i\gamma^{5} T_{a}^{\dagger iA} (\widetilde{\partial}_{\nu} \rho^{\mu\nu})_{i} N_{A} + \frac{c_{6\Delta\rho}}{M^{2}} \overline{\Delta}^{a}_{\lambda} \sigma_{\mu\nu} T_{a}^{\dagger iA} (\widetilde{\partial}^{\lambda} \overline{\rho}^{\mu\nu})_{i} N_{A} + \text{c.c.}$$

#### • Chiral symmetry in QCD:

$$SU(2)_L \otimes SU(2)_R \otimes U(1)_B$$

$$\mathcal{L} = \mathcal{L}_{QCD} + \overline{q}\gamma_{\mu}(\mathbf{v}^{\mu} + B\mathbf{v}^{\mu}_{(s)} + \gamma_{5}\mathbf{a}^{\mu})q - \overline{q}(\mathbf{s} - i\gamma_{5}\mathbf{p})q$$

B. Serot and X.Z., arXiv:1011.5913; Advances in QFT (InTech, 2012) (arXiv:1110.2760) 47

• Chiral symmetry in QCD

#### • The lagrangian, baryon section:

$$\mathcal{L}_{N(\hat{\nu}\leqslant3)} = \overline{N} \Big( i\gamma^{\mu} [\widetilde{\partial}_{\mu} + ig_{\rho}\rho_{\mu} + ig_{v}V_{\mu}] + g_{A}\gamma^{\mu}\gamma^{5} \widetilde{a}_{\mu} - M + g_{s}\phi \Big) N$$
  
$$- \frac{f_{\rho}g_{\rho}}{4M} \overline{N}\rho_{\mu\nu}\sigma^{\mu\nu}N - \frac{f_{v}g_{v}}{4M} \overline{N}V_{\mu\nu}\sigma^{\mu\nu}N - \frac{\kappa_{\pi}}{M} \overline{N}\widetilde{v}_{\mu\nu}\sigma^{\mu\nu}N + \frac{4\beta_{\pi}}{M} \overline{N}N \operatorname{Tr}(\widetilde{a}_{\mu}\widetilde{a}^{\mu}) \Big)$$

$$\mathcal{L}_{\Delta} = \frac{-i}{2} \overline{\Delta}^{a}_{\mu} \{ \sigma^{\mu\nu} , (i \,\widetilde{\partial} - h_{\rho} \, \not\!\!{\rho} - h_{v} \, \not\!\!{V} - m + h_{s} \phi) \}^{b}_{a} \Delta_{b\nu} + \widetilde{h}_{A} \overline{\Delta}^{a}_{\mu} \, \widetilde{a}^{b}_{a} \gamma^{5} \Delta^{\mu}_{b}$$

$$\mathcal{L}_{\Delta,N,\pi} = h_A \overline{\Delta}^{a\mu} T_a^{\dagger iA} \widetilde{a}_{i\mu} N_A + \text{c.c.},$$

- Chiral symmetry in QCD
- The lagrangian, baryon section.

#### • The lagrangian, meson section:

$$\mathcal{L}_{\text{meson}(\hat{\nu} \leqslant 4)} = \frac{1}{2} \partial_{\mu} \phi \, \partial^{\mu} \phi + \frac{1}{4} f_{\pi}^{2} \operatorname{Tr}[\widetilde{\partial}_{\mu} U(\widetilde{\partial}^{\mu} U)^{\dagger}] \\ + \frac{1}{4} f_{\pi}^{2} m_{\pi}^{2} \operatorname{Tr}(U + U^{\dagger} - 2) \\ - \frac{1}{2} \operatorname{Tr}(\rho_{\mu\nu} \rho^{\mu\nu}) - \frac{1}{4} V^{\mu\nu} V_{\mu\nu} \\ + \frac{1}{2g_{\gamma}} \left( \operatorname{Tr}(F^{(+)\mu\nu} \rho_{\mu\nu}) + \frac{1}{3} f_{s}^{\mu\nu} V_{\mu\nu} \right) \right)$$

Vector meson dominance (VMD)

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. . .

#### • Electroweak (EW) interactions:

$$\begin{split} l_{\mu} &= -e \, \frac{\tau^{0}}{2} \, A_{\mu} + \frac{g}{\cos \theta_{w}} \sin^{2} \theta_{w} \, \frac{\tau^{0}}{2} \, Z_{\mu} \\ &- \frac{g}{\cos \theta_{w}} \frac{\tau^{0}}{2} \, Z_{\mu} - g V_{ud} \, \left( W_{\mu}^{+1} \, \frac{\tau_{+1}}{2} + W_{\mu}^{-1} \frac{\tau_{-1}}{2} \right) \, , \\ r_{\mu} &= -e \, \frac{\tau^{0}}{2} \, A_{\mu} + \frac{g}{\cos \theta_{w}} \sin^{2} \theta_{w} \, \frac{\tau^{0}}{2} \, Z_{\mu} \, , \\ v_{(s)\mu} &= -e \, \frac{1}{2} \, A_{\mu} + \frac{g}{\cos \theta_{w}} \sin^{2} \theta_{w} \, \frac{1}{2} \, Z_{\mu} \, . \end{split}$$

# A quick look (recap)

- Chiral symmetry
- The lagrangian
- Electroweak (EW) interactions (CVC and PCAC)

## Benchmarks



### Benchmarks



### Benchmarks

