

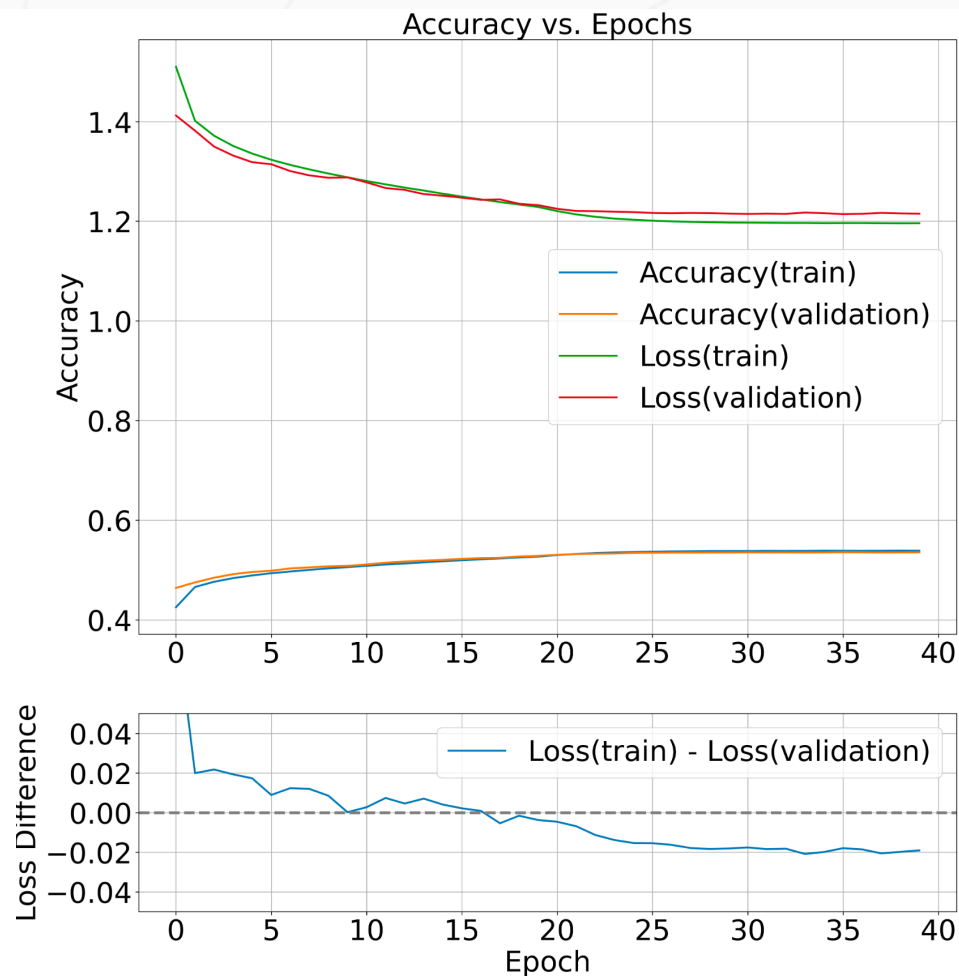
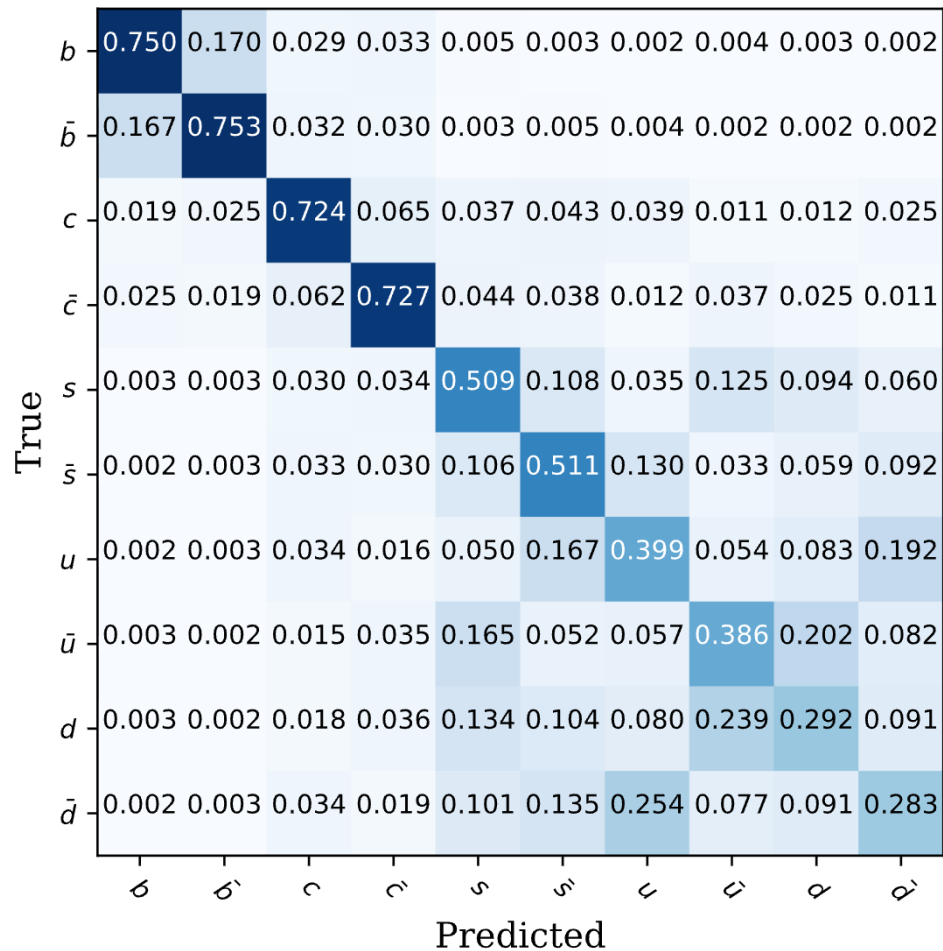


Update of AFB measurement @ $Z \rightarrow q\bar{q}$

Libo Liao

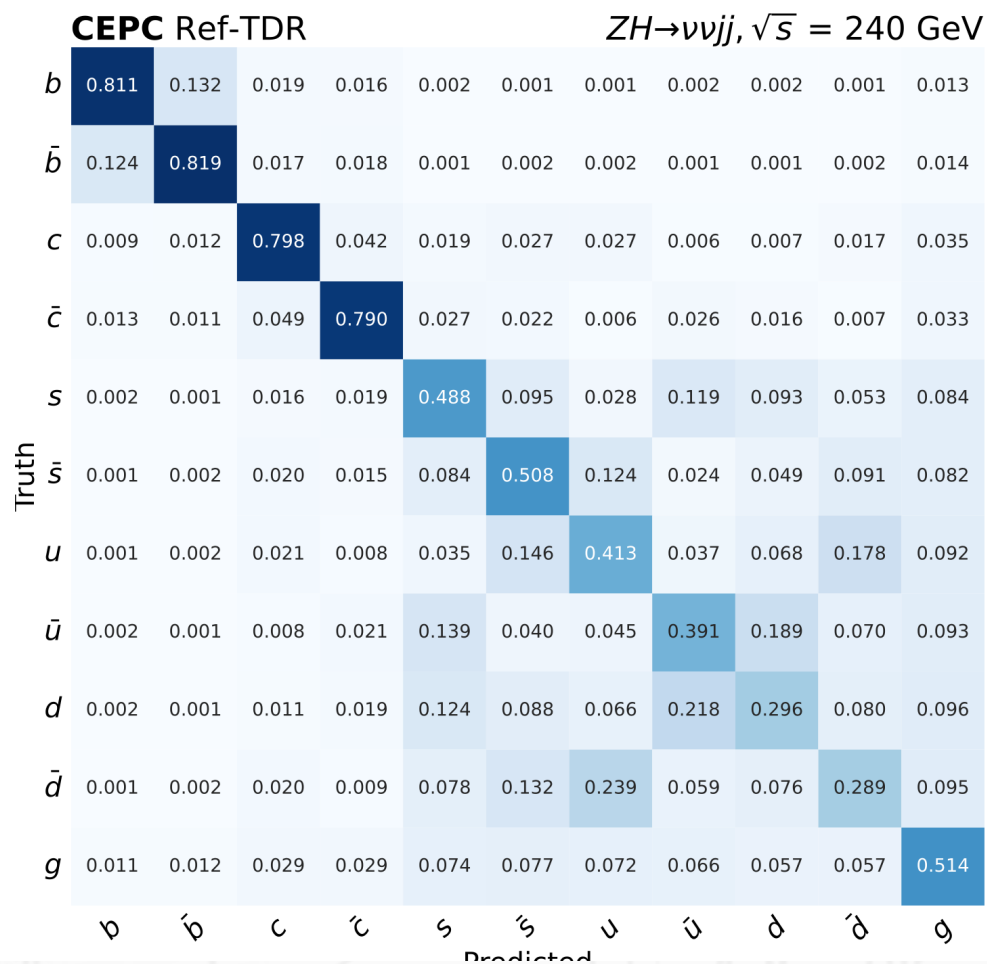
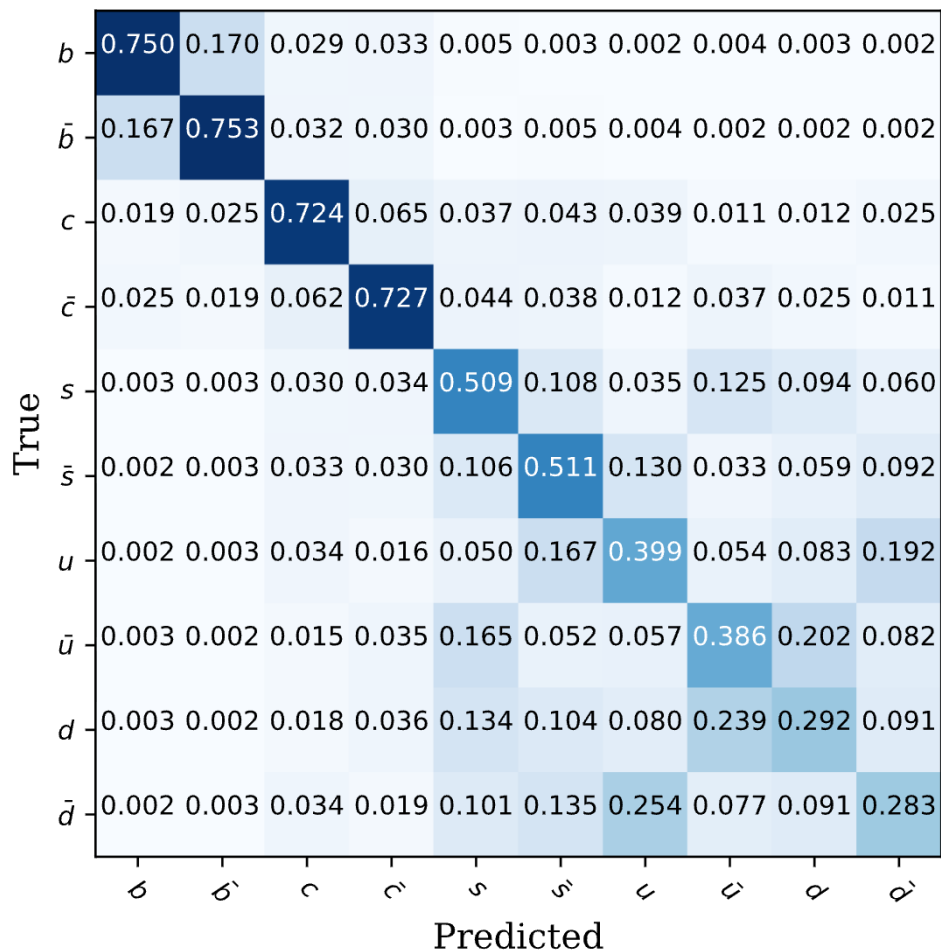
Gang LI, Weimin Song, Jin Zhang, Jiabao Gong, Jiarong Li, Yilun Wang, and Zhaoling Zhang

Transformer - zqq



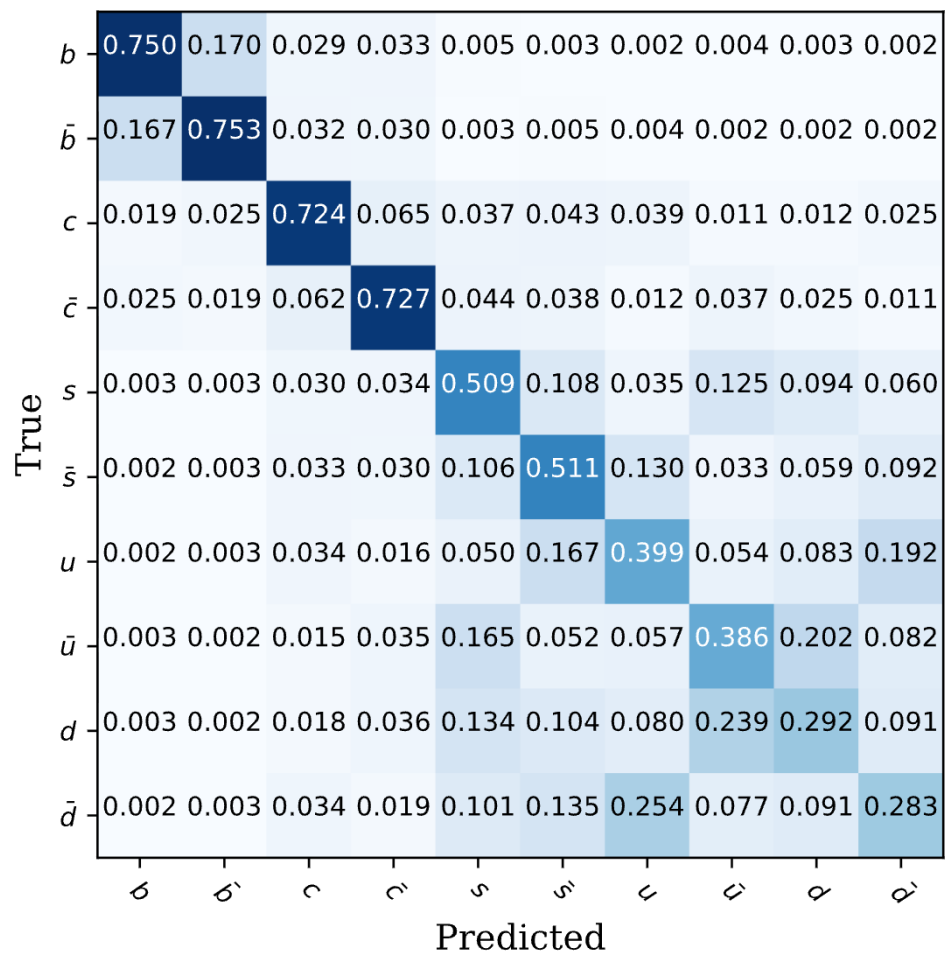
- stable training
- no obvious over-fitting

Compared to TDR



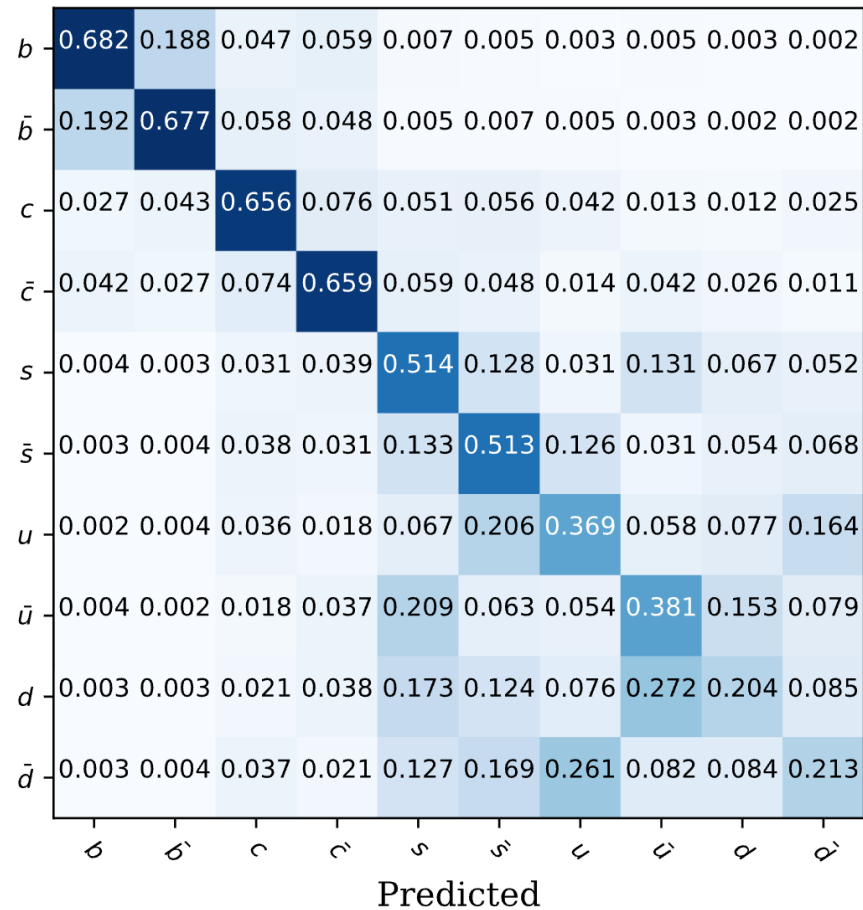
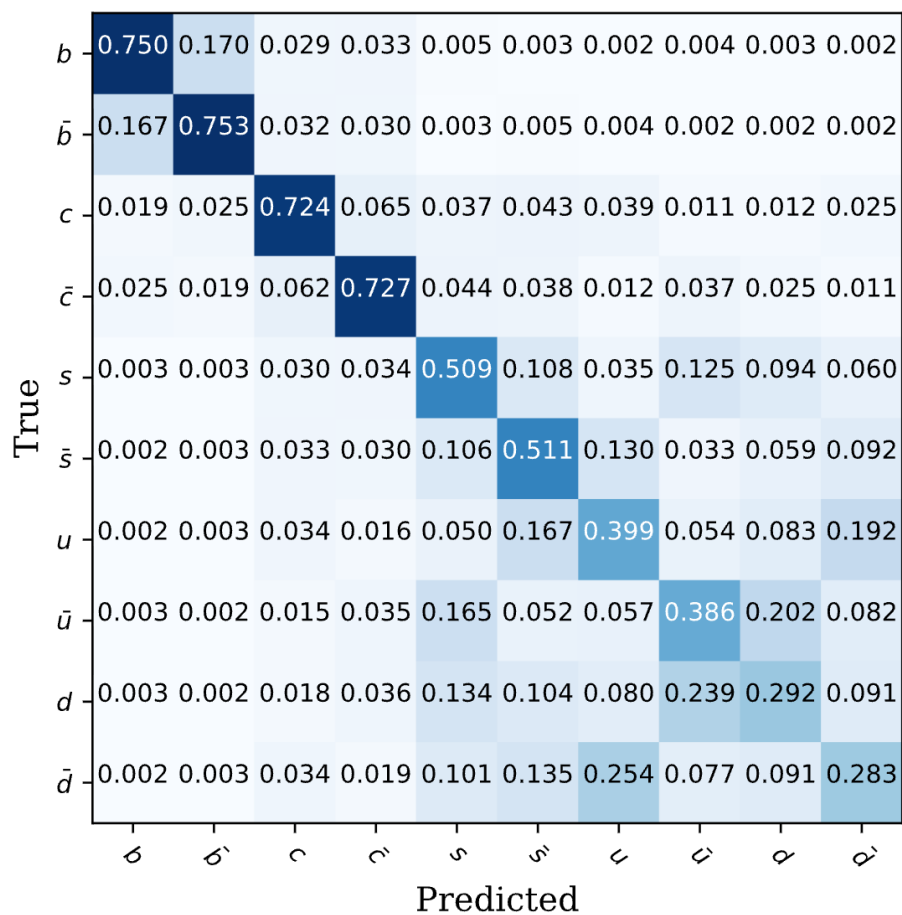
- comparable misclassification distributions.
- different samples, Based on Z pole and ZH

Compared to Z->ss analysis



□ comparable results in each elements

Compared to ParticleNet



□ slight improved in b, c identification

Update physics results

AFB (1e-3)	PN	ParT
b	0.013	0.012
c	0.023	0.022
s	0.019	0.017

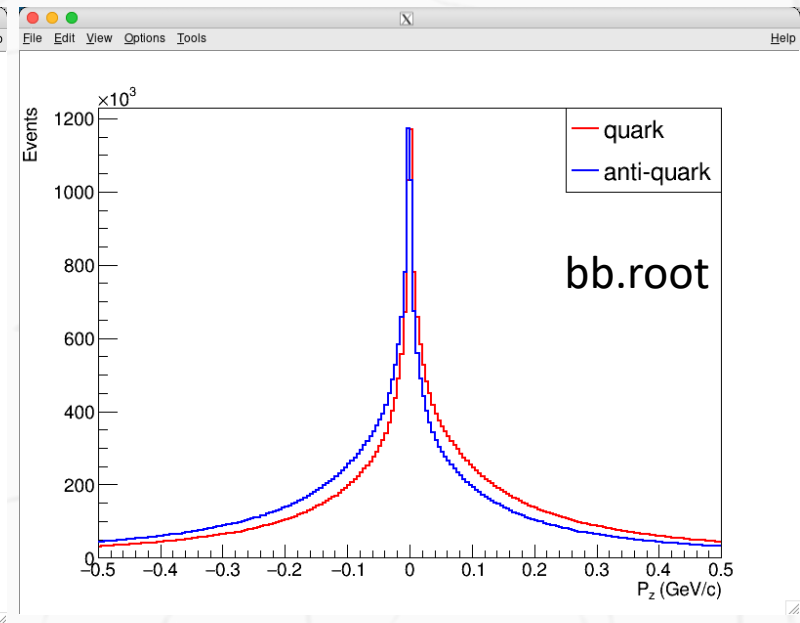
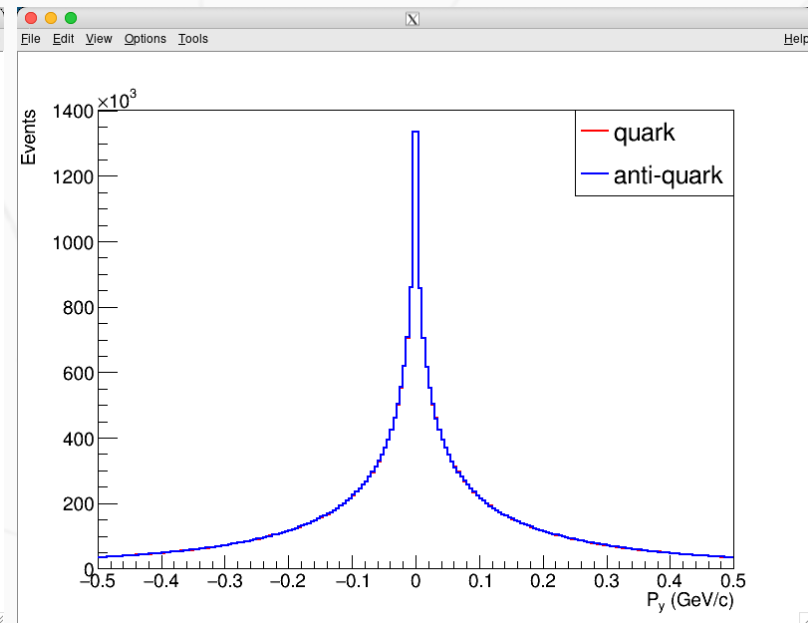
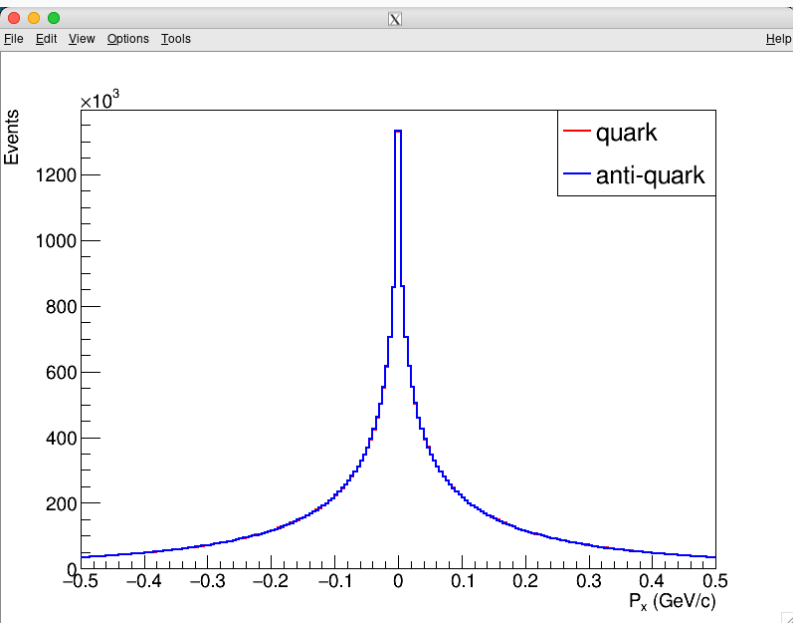
sinW (1e-6)	fast sim	PN	ParT
b	1.0	1.1	1.0
c	1.5	1.7	1.6
s	1.4	1.5	1.4

Pe	0.5		0.6		0.7		0.8	
sinW(1e-6)	PN	ParT	PN	ParT	PN	ParT	PN	ParT
b	0.57	0.54	0.56	0.53	0.54	0.51	0.52	0.50
c	0.81	0.72	0.79	0.71	0.77	0.69	0.75	0.67
s	0.93	0.88	0.92	0.87	0.91	0.86	0.89	0.84

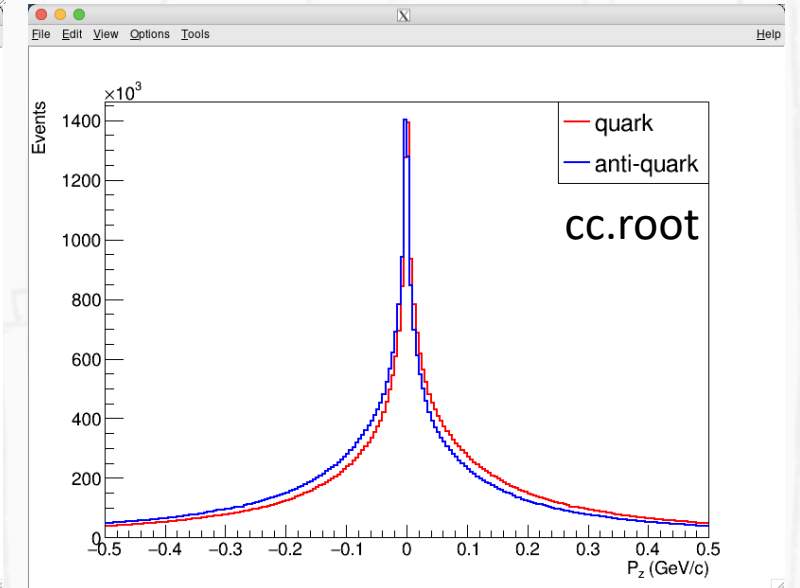
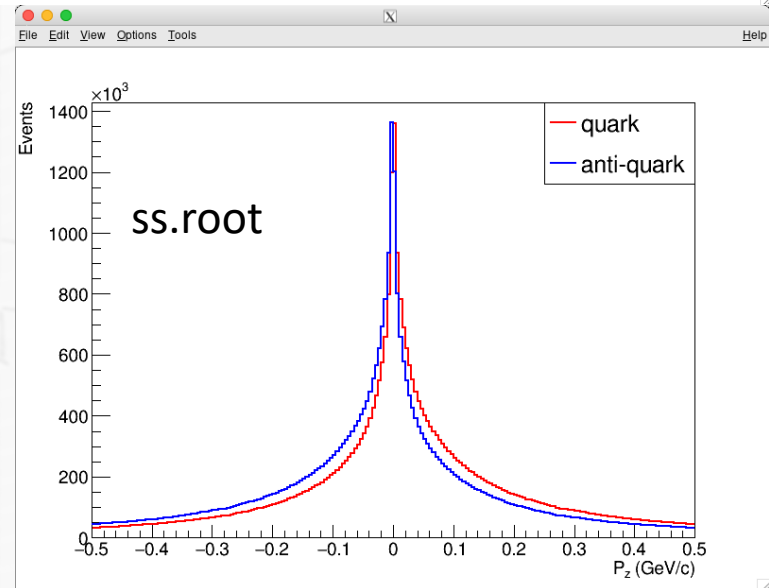
- ▣ Comparable results
- ▣ Dominant by statistics

BACKUP

unforeseen issue



The momentum is not consistent in P_z

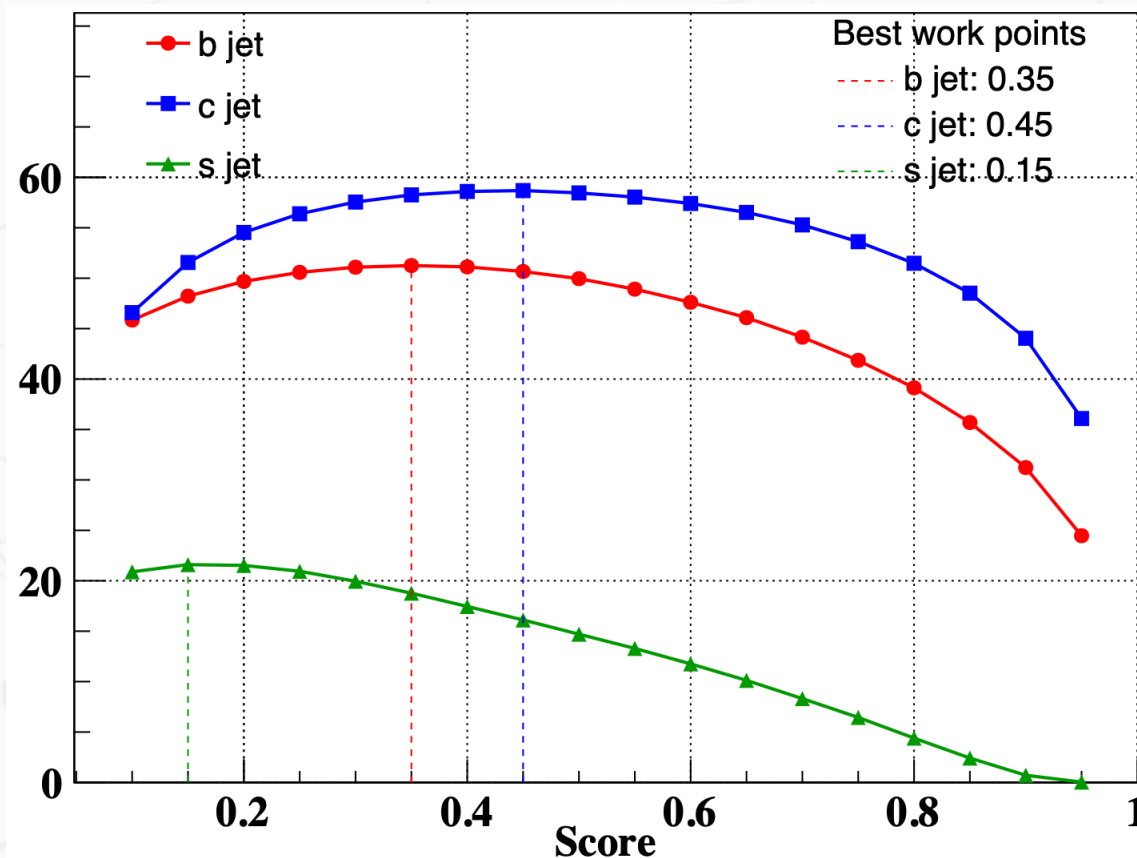


Determination of working points

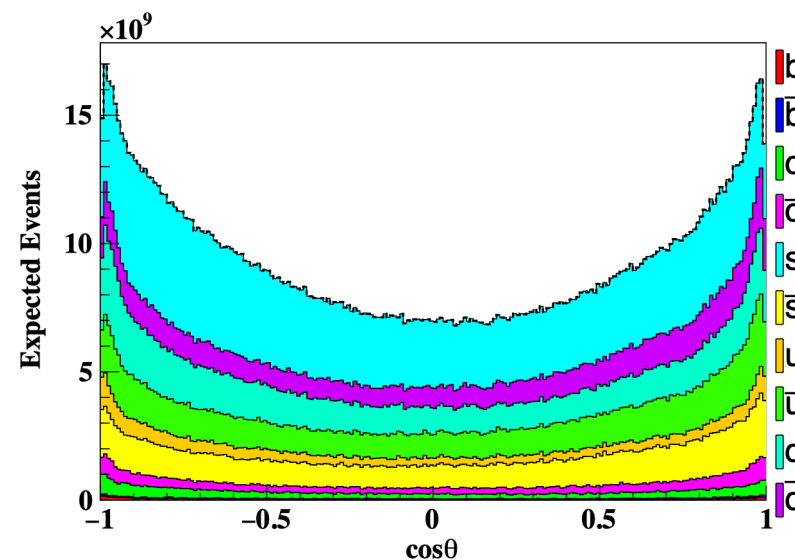
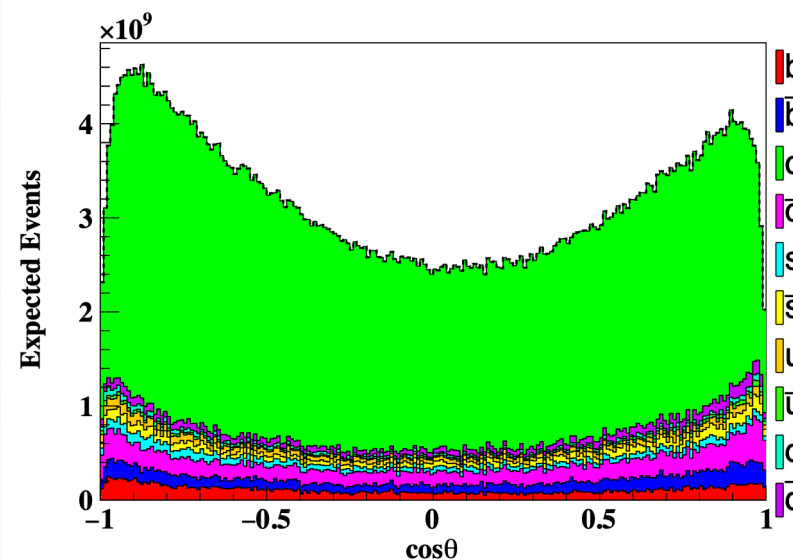
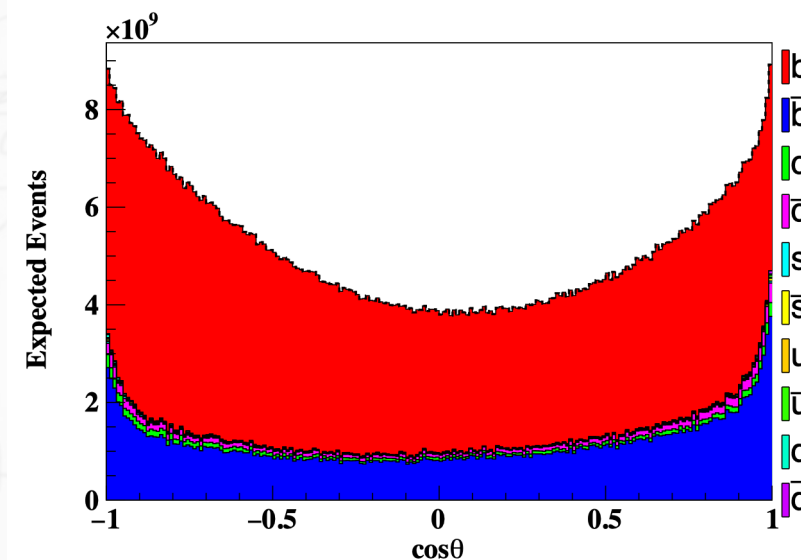
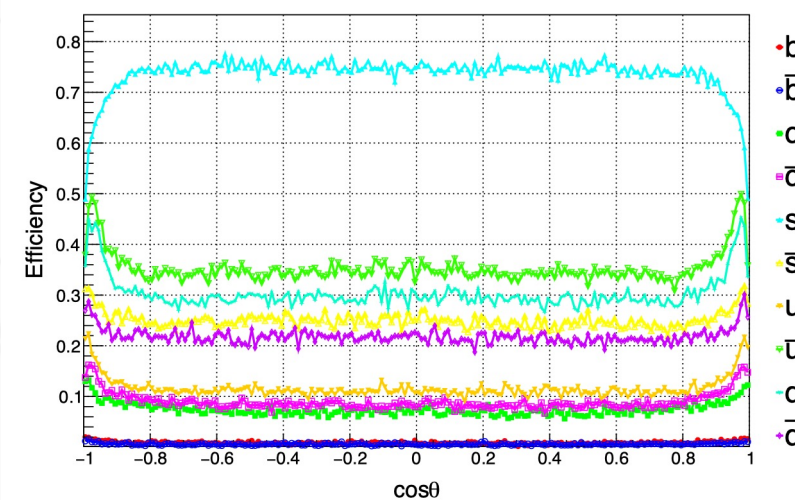
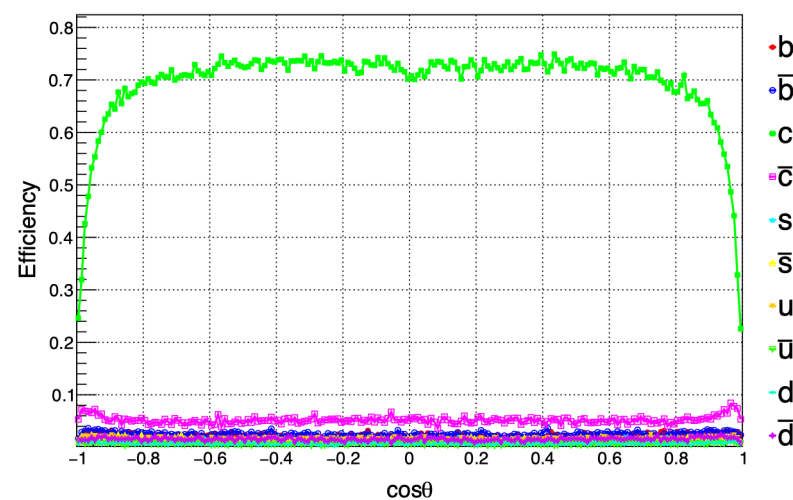
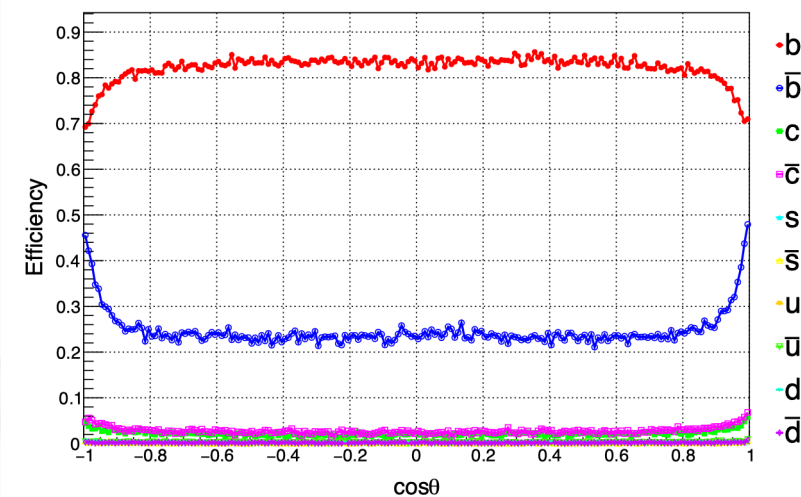
- The working point is chosen to maximize the overall jet charge identification performance, weighted by the production cross sections

$$\begin{aligned}\text{FoM} &= \sum_i \sigma_i \times \epsilon_i \times \rho_i \times \left(\frac{1}{\sigma_i} \frac{\partial \sigma_i}{\partial A_{\text{FB}}} \right)^2 \\ &= \sum_i \epsilon_i \times \rho_i \times \frac{1}{\sigma_i} \left(\frac{8}{3} \cos \theta_i \right)^2,\end{aligned}$$

- The results are consistent with expectation.
 - ◆ The lower threshold for s-jets reflects the greater challenge in distinguishing strange quarks due to the absence of a heavy quark decay vertex.



Efficiency curves



Measurement of A_{FB}

- The expected number of events for reconstructed flavor j (5 flavor) in bin i (200) is:

$$N_{i,j}^{\text{exp}} = \mathcal{L} \sum_k \sigma_{i,j} \epsilon_{i,j,k},$$

$$\sigma_{i,j} = (1 + \cos^2 \theta_i + \frac{8}{3} A_{FB,j} \cos \theta_i) \Delta,$$

- A_{FB} could be extracted via a binned maximum likelihood fit

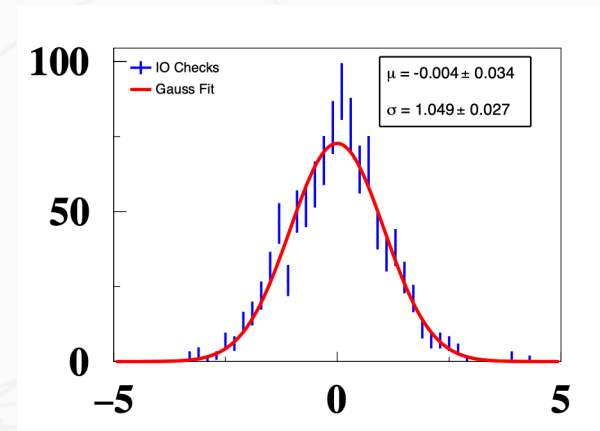
$$\mathcal{L}(A_{FB,j}, N_j^{\text{prod}}) = \prod_i \frac{\mu_{i,j}^{n_{i,j}} e^{-\mu_{i,j}}}{n_{i,j}!},$$

- To estimate the statistical precision, and joint flavor and charge identification

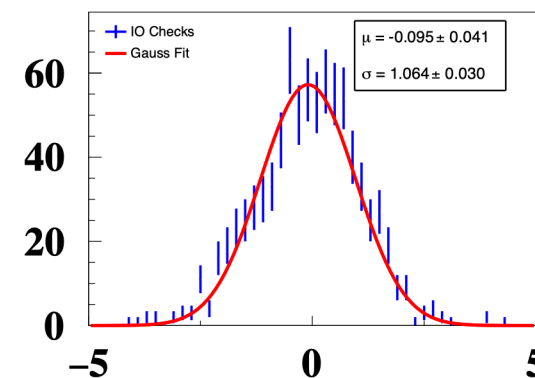
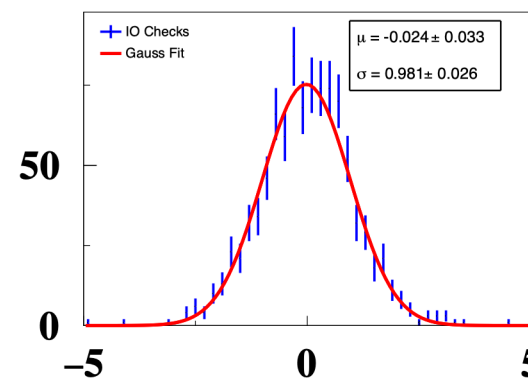
◆ Generate 1000 toy Monte Carlo samples.

- The results are listed below:

	This work (10^{-3})	PDG values(10^{-3})
ΔA_{FB}^b	0.013	16
ΔA_{FB}^c	0.021	50
ΔA_{FB}^s	0.018	148



$$\text{Pull} = (O - I) / \text{error}$$





Measurement of $\sin^2 \theta_W^{eff}$

- The published sensitivities of $\sin^2 \theta_{eff}^{lept}$ are at the 10^{-3} level
 - ◆ The sensitivity obtained from $A_{FB}^{0,b}$ is better than that from $A_{FB}^{0,l}$, because the cross section for quark final states is larger than that for leptonic decays.

	$\sin^2 \theta_{eff}^{lept}$	$\Delta \sin^2 \theta_{eff}^{lept} (10^{-3})$
LEP, $A_{FB}^{0,b}$	0.23221 ± 0.00029	1.3
LEP, $A_{FB}^{0,l}$	0.23099 ± 0.00053	2.3
SLC, A_{LR}	0.23098 ± 0.00026	1.2
LEP+SLC	0.23153 ± 0.00016	0.70

- To extract $\sin^2 \theta_W^{eff}$ from the simulated CEPC data, a similar binned extended maximum likelihood is performed
- 1000 toy Monte Carlo samples, pull distributions are also standard
- The results (10^{-6}) are listed
 - ◆ Improvements in both statistical precision and analysis methods

	b	c	s
$\Delta \sin^2 \theta_{eff}^{lept}$	1.0	1.5	1.4



Extension to polarization scenario

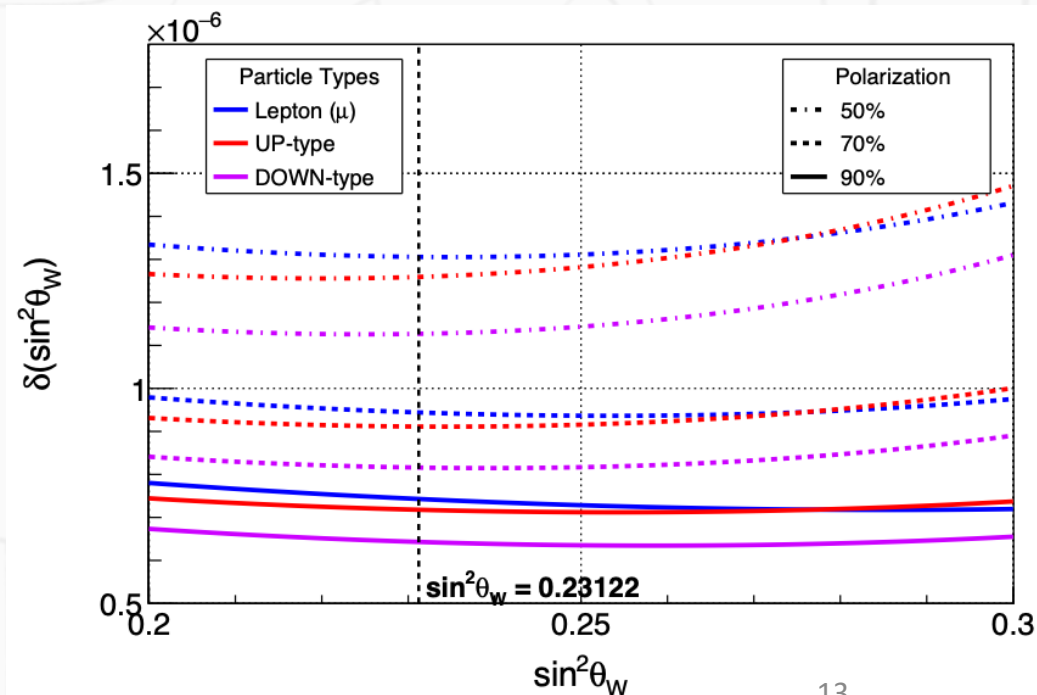
- The power of beam polarization for electroweak precision measurements was convincingly demonstrated by the SLC experiment
 - ◆ With less statistics using a polarized electron beam
 - ◆ A precision comparable to — and in some respects exceeding

	$\sin^2 \theta_{\text{eff}}^{\text{lept}}$	$\Delta \sin^2 \theta_{\text{eff}}^{\text{lept}} (10^{-3})$
LEP, $A_{FB}^{0,b}$	0.23221 ± 0.00029	1.3
LEP, $A_{FB}^{0,l}$	0.23099 ± 0.00053	2.3
SLC, A_{LR}	0.23098 ± 0.00026	1.2
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- The statistical advantage of the polarized measurement could be estimated using Fisher Information
 - The factor $\sqrt{2}$ accounts for the fact that the total integrated luminosity is split between the left-handed and right-handed polarization states.

$$\frac{\Delta(\sin^2 \theta_W^{\text{eff}})_{\text{pol}}}{\Delta(\sin^2 \theta_W^{\text{eff}})_{\text{unpol}}} = \frac{\sqrt{2}}{P_e} \cdot \frac{|\partial A_{FB} / \partial \sin^2 \theta_W^{\text{eff}}|}{|\partial A_{LR} / \partial \sin^2 \theta_W^{\text{eff}}|},$$

- The expected precision of $\sin^2 \theta_W^{\text{eff}}$ measurements using the Fisher information method



Extension to polarization scenario

- The expected bin-by-bin number of left(right)-handed events

$$N_{L(R),i,j}^{\text{exp}} = \mathcal{L} \Delta \sum_k \sigma_{L(R),i,j} \epsilon_{i,j,k},$$

- $\sin^2 \theta_W^{\text{eff}}$ could be extracted via a binned maximum likelihood fit

$$\mathcal{L}(\sin^2 \theta_W^{\text{eff}}, N_{L,j}^{\text{prod}}, N_{R,j}^{\text{prod}}) = \prod_i \frac{\mu_{L,i,j}^{n_{L,i,j}} e^{-\mu_{L,i,j}}}{n_{L,i,j}!} \times \frac{\mu_{R,i,j}^{n_{R,i,j}} e^{-\mu_{R,i,j}}}{n_{R,i,j}!},$$

- 1000 toy Monte Carlo samples, pull distributions are also standard
- The results (10^{-6}) are listed

	$P_e = 0.5$	$P_e = 0.6$	$P_e = 0.7$	$P_e = 0.8$
b	0.55	0.54	0.52	0.51
c	0.71	0.69	0.68	0.65
s	0.83	0.82	0.81	0.79