

Instrumentation of Particle Physics Experiments

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Instrumentation

- Instrumentation is a very broad topic, which is everywhere in daily life
 - Oscilloscope, signal generator, digital multimeter ...
 - Smartphone – Pocket Geiger for iPhone after Fukushima nuclear disaster
 - Commercial, scientific, military ...
 - **Wikipedia: Instrumentation** is defined as the art and science of measurement and control of process variables within a production or manufacturing area
- This lecture will focus on the instrumentation of particle physics experiments
 - Detector & Electronics



Organization of Lectures

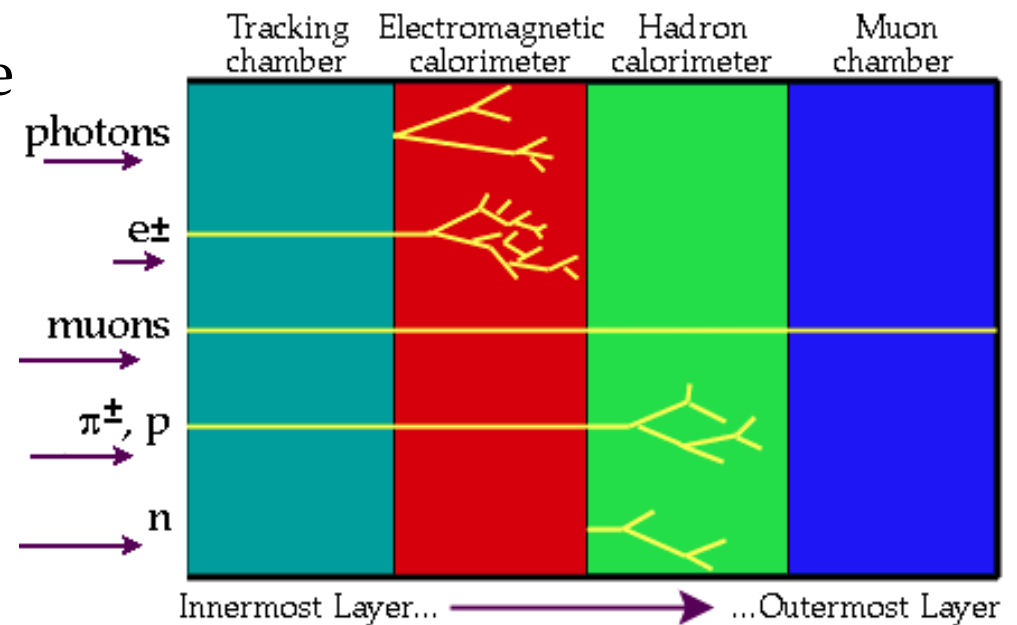
- Lecture 1
 - Calorimetry and Development of Noble Liquid Calorimeter
- Lecture 2
 - Readout Electronics of ATLAS LAr Calorimeters
- Lecture 3
 - ATLAS LAr Readout Electronics Upgrade
- Lecture 4
 - Accelerator Based Neutrino Experiments and Cold Electronics Development

Lecture 1 – Calorimetry

- Overview of Calorimetry
 - Electromagnetic Calorimeter
 - Hadronic Calorimeter
 - New Calorimeter Development
- Noble Liquid Calorimeter
 - Invention of LAr Calorimeter
 - R806 @ CERN ISR
 - HELIOS @ CERN SPS
 - NA48 @ CERN SPS
 - Accordion – From RD3 to GEM @ SSC and ATLAS @ LHC

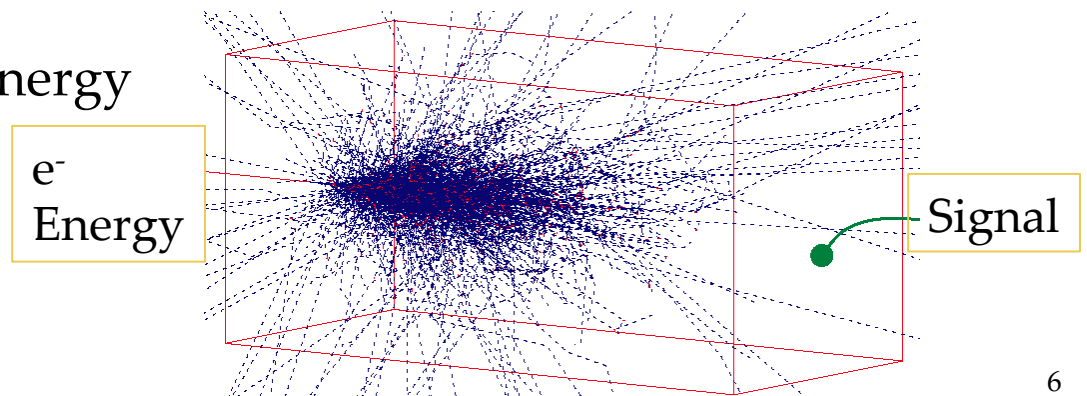
Calorimetry

- A typical particle physics detector
- Particles characteristics are measured through different type of detectors and identified from specific behaviors due to their interaction with matter
- Calorimeters are used to detect
 - γ , e
 - Jets (q, g)
 - Missing energy (e.g. ν)



Calorimetry

- Calorimeter is the detector for energy measurement via total absorption of particles
 - Most calorimeters are position sensitive to measure energy depositions depending on their location
 - Calorimeter is a “destructive” method. Energy and particle get absorbed
- Principle of operation
 - Incoming particle initiates particle shower
 - Energy deposited in form of: heat, ionization, excitation of atoms, Cherenkov light ...
 - Signal \sim total deposited energy
 - Signal collection



Calorimetry

- Energy vs. momentum measurement

- Calorimeter $\frac{\sigma_E}{E} \sim \frac{1}{\sqrt{E}}$

- Gas Detector $\frac{\sigma_p}{p} \sim p$

- ATLAS:

$$\frac{\sigma_E}{E} \approx \frac{0.1}{\sqrt{E}}, \quad \frac{\sigma_E}{E} = 1\% @ 100 GeV \quad \frac{\sigma_p}{p} \approx 5 \cdot 10^{-4} \cdot p_t, \quad \frac{\sigma_p}{p} = 5\% @ 100 GeV$$

- At very high energies one has to switch to calorimeters because their resolution improves while those of a magnetic spectrometer decreases with E

- Shower depth

- Calorimeter

- Shower depth nearly energy independent, compact calorimeter is possible

- Magnetic spectrometer

- Detector size has to grow quadratically to maintain resolution

Calorimetry

- Calorimeter features

- Calorimeters can be built as 4π detectors, i.e. they can detect particles over almost the full solid angle
 - Compactness: dimension necessary to containment is proportional to $\ln E$
- Calorimeters can provide fast timing signals (1 to 10 ns); can be used for triggering
- Calorimeters can measure the energy of both, charged and neutral particles, if they interact via electromagnetic or strong forces, e.g. $\gamma, \pi^0, K^0, \dots$
- Segmentation in depth allows separation of hadrons (p, n, π^\pm), from particles which only interact electromagnetically (γ, e)
 - Measure of position, direction & particle id on topological basis

Rules of Thumbs – EM Shower

- Radiation length $X_0 = \frac{180A}{Z^2} \frac{g}{cm^2}$
- Critical energy $E_c = \frac{550MeV}{Z}$
- Shower maximum $t_{max} = \ln \frac{E}{E_c} - \begin{cases} 1.0 & e^- \text{ induced shower} \\ 0.5 & \gamma \text{ induced shower} \end{cases}$
- Longitudinal energy containment

$$L(95\%) = t_{max} + 0.08Z + 9.6 [X_0]$$

- Transverse energy containment

$$R(90\%) = R_M$$

$$R(95\%) = 2R_M$$

EM Calorimeter – Homogeneous

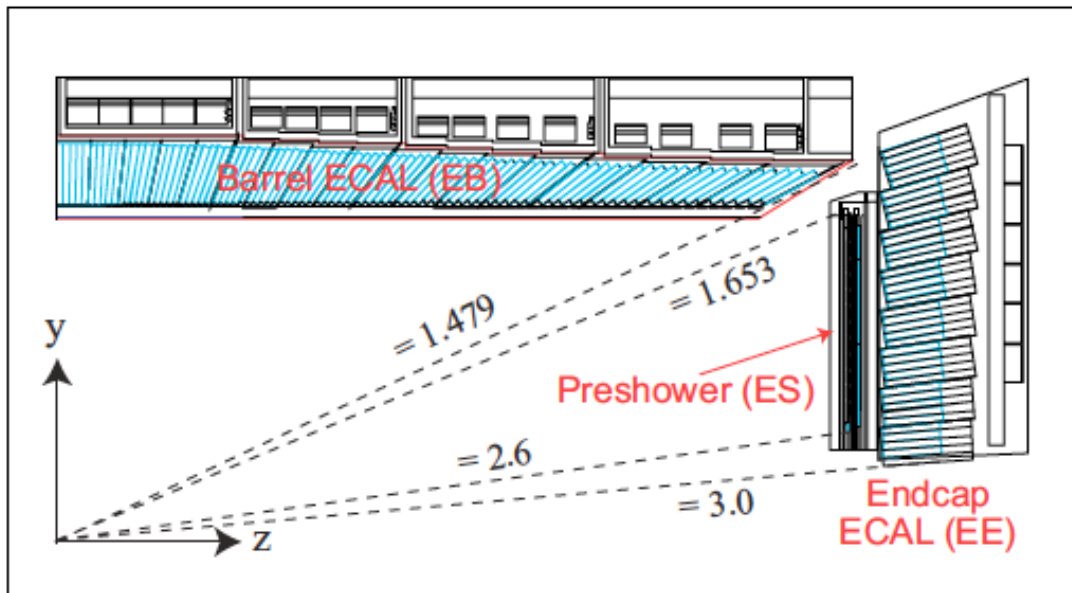
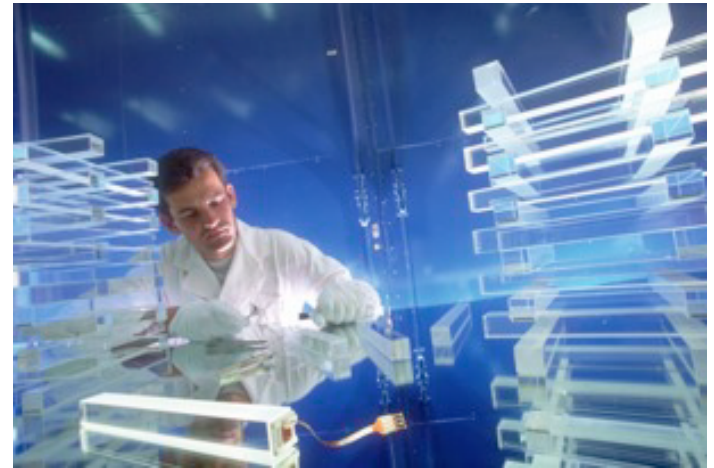
- Homogeneous calorimeters: all the energy is deposited in the active medium.
 - Absorber is active medium as well

Signal	Material
Scintillation light	BGO, BaF ₂ , CeF ₃ , ...
Cherenkov light	Lead Glass
Ionization signal	Liquid noble gases (Ar, Kr, Xe)

- Pros
 - Excellent energy resolution
- Cons
 - Expensive
 - No information on longitudinal shower shape

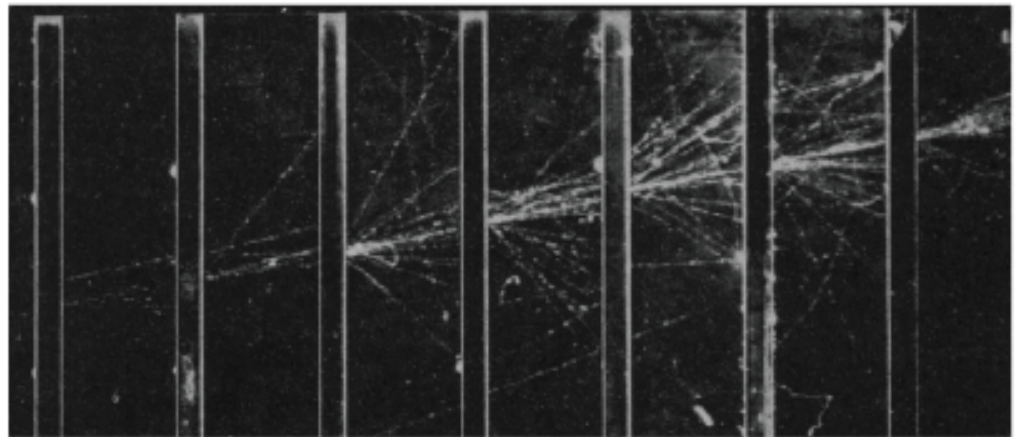
EM Calorimeter – Homogeneous

- CMS crystal calorimeter
 - Scintillator: PbWO_4
 - Photosensor: APDs
 - Number of crystals: $\sim 70,000$
 - Light output: 4.5 photons/MeV



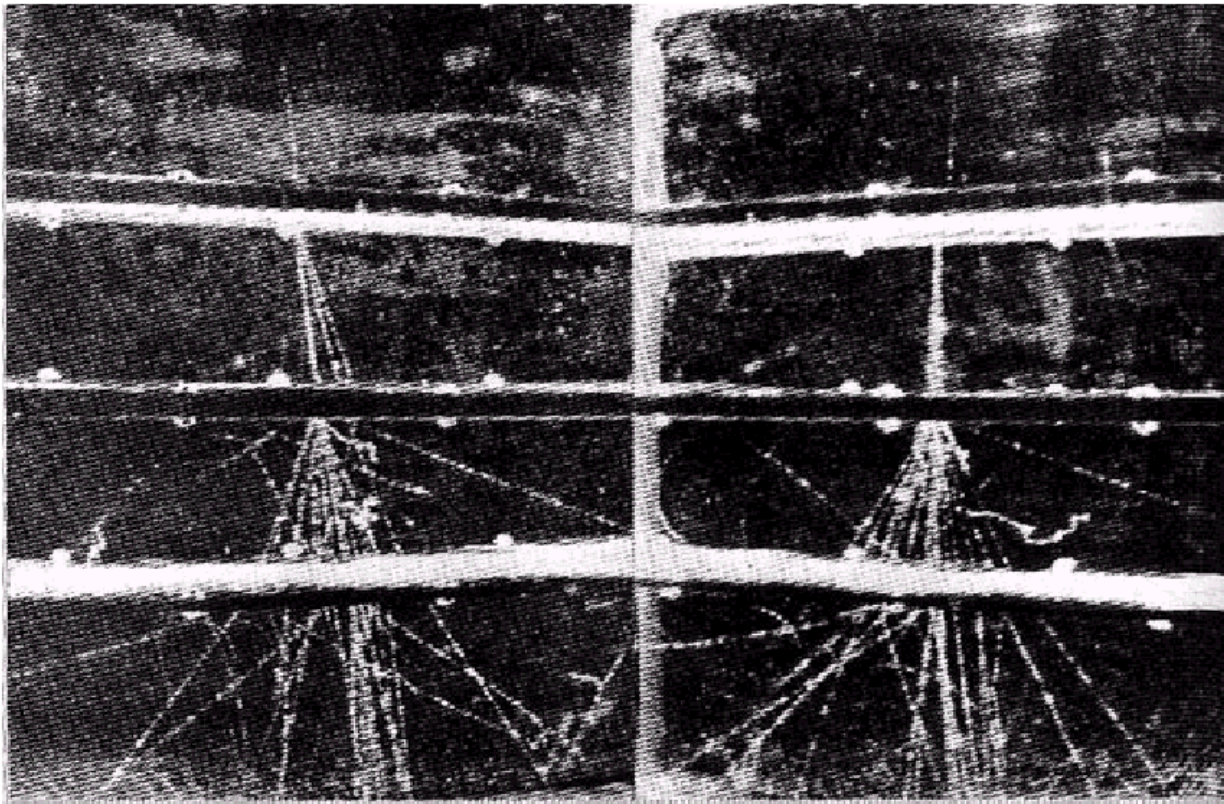
EM Calorimeter – Sampling

- Sampling calorimeters: shower is sampled by layers of active medium (low Z) alternated with dense absorber (high Z) material
 - Absorber is shower generator, active layers negligible in the shower development
- Absorber materials
 - Iron (Fe), Lead (Pb), Uranium (U)
- Active materials
 - Plastic scintillator
 - Silicon detectors
 - Liquid ionization chamber
 - Gas detector



EM Calorimeter – Sampling

- Cloud chamber photograph of EM shower developing in lead plates
 - Thickness from top down 1.1, 1.1, 0.13 X_0
 - Exposed to cosmic radiation



EM Calorimeter – Sampling

■ Pros

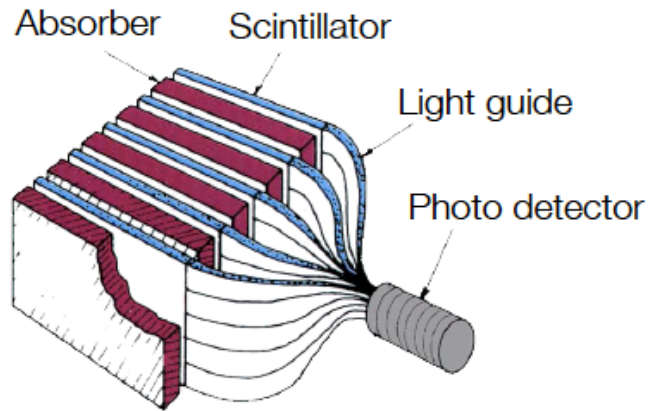
- By separating passive and active layers the different layer materials can be optimally adapted to the corresponding requirements
- By freely choosing high-density material for the absorbers one can build very compact calorimeters
- Sampling calorimeters are simpler with more passive material and thus cheaper than homogeneous calorimeters
- Detailed shower shaper information is available

■ Cons

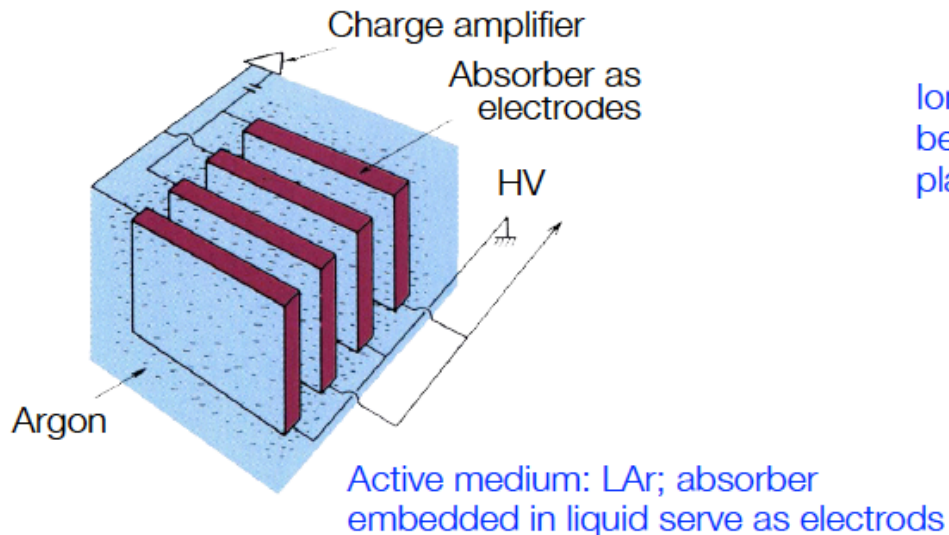
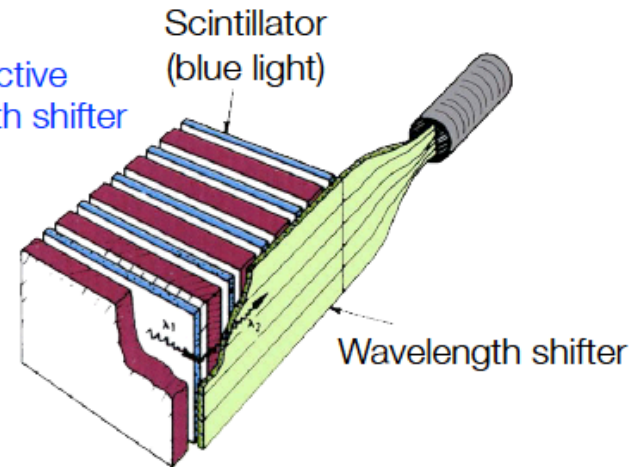
- Only part of the deposited particle energy is actually detected in the active layers; typically a few percent [for gas detectors even only $\sim 10^{-5}$]
- Due to this sampling-fluctuations typically result in a reduced energy resolution for sampling calorimeters

Sampling Calorimeters – Possible Setup

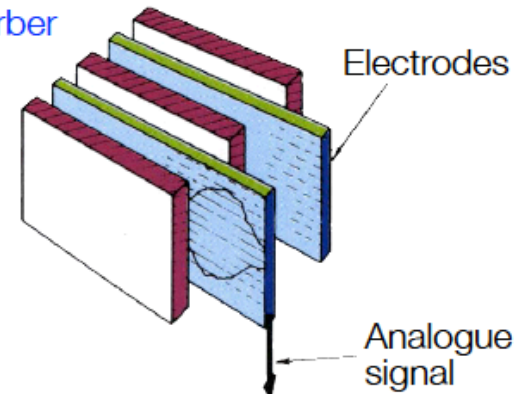
Scintillators as active layer;
signal readout via photo multipliers



Scintillators as active layer; wave length shifter to convert light

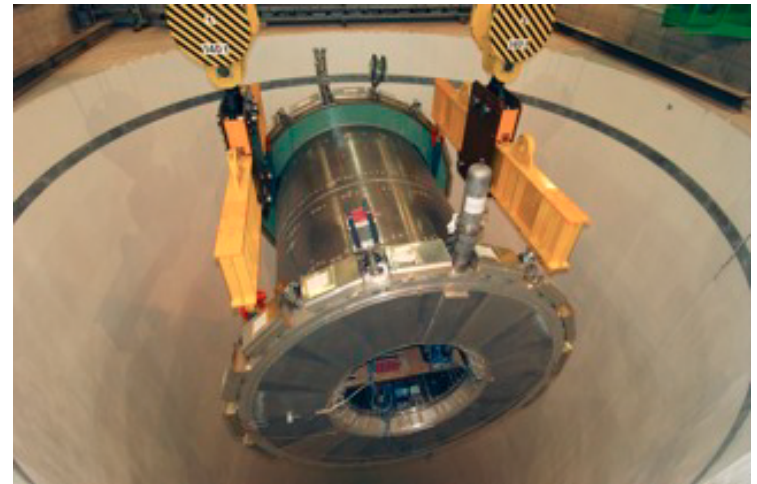
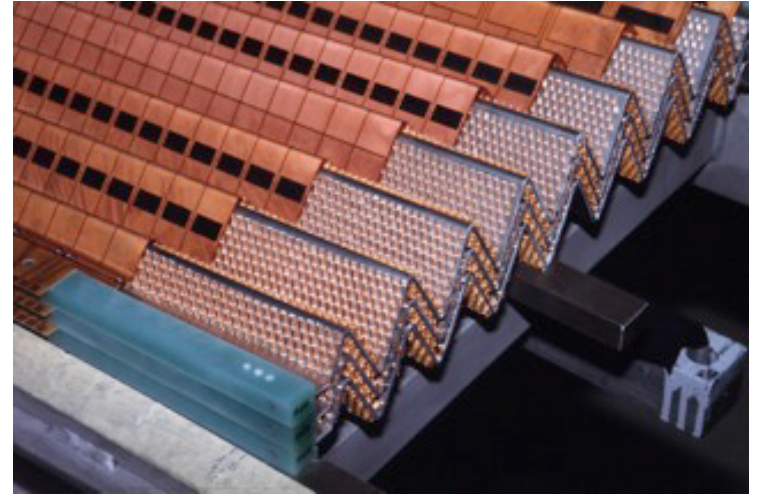
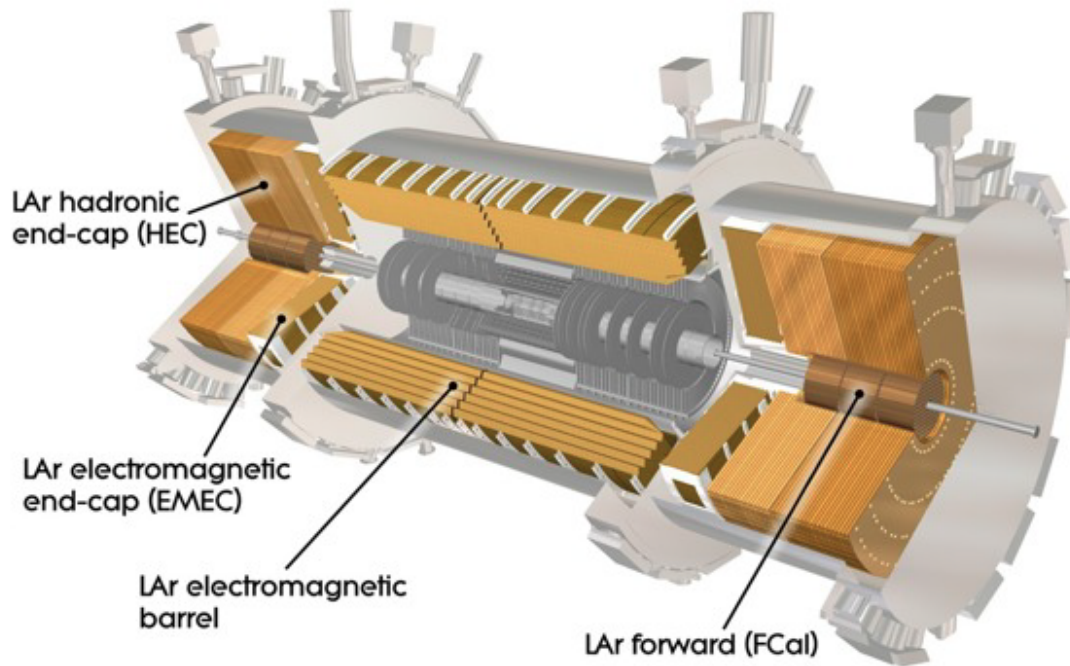


Ionization chambers between absorber plates



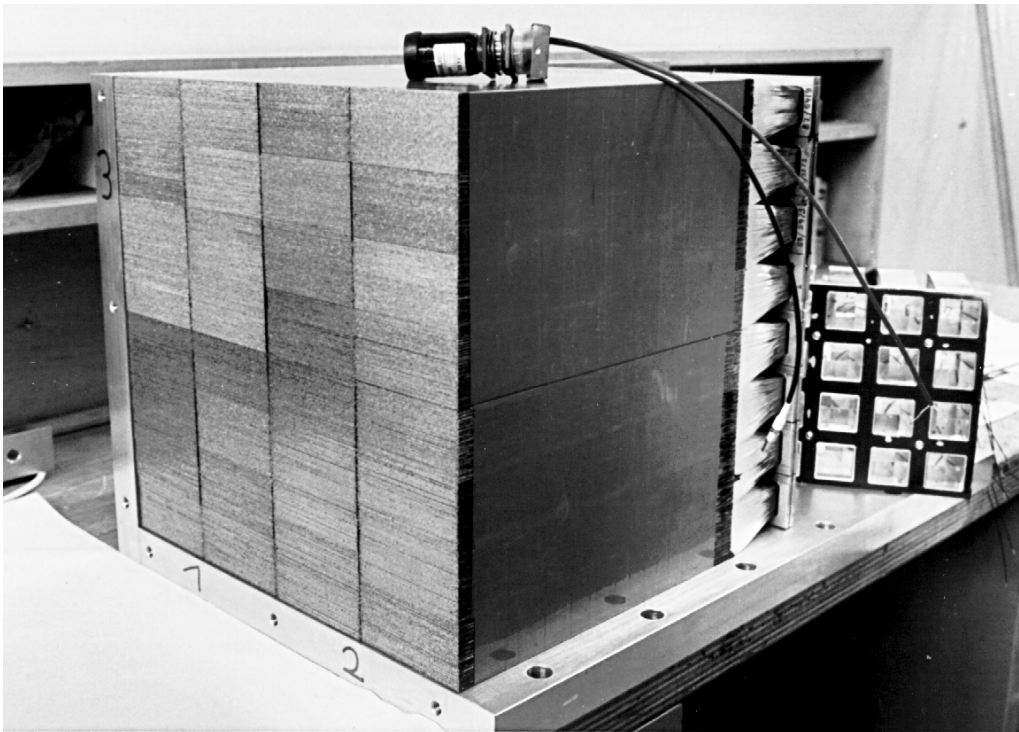
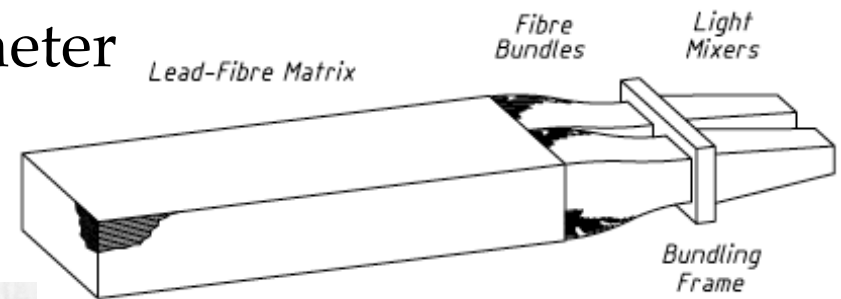
EM Calorimeter – Sampling

- ATLAS Liquid Argon Calorimeter

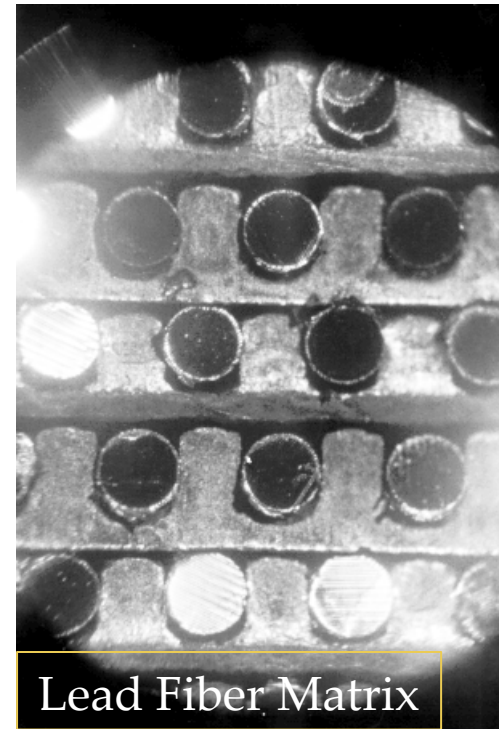


EM Calorimeter – Sampling

- H1 SpaCal: Spaghetti Calorimeter



4 SpaCal Supermodules



Lead Fiber Matrix

Homogeneous vs. Sampling Calorimeter

- Resolution of typical EM calorimeter [E is in GeV]

Technology (Experiment)	Depth	Energy resolution	Date
NaI(Tl) (Crystal Ball)	$20X_0$	$2.7\%/E^{1/4}$	1983
Bi ₄ Ge ₃ O ₁₂ (BGO) (L3)	$22X_0$	$2\%/\sqrt{E} \oplus 0.7\%$	1993
CsI (KTeV)	$27X_0$	$2\%/\sqrt{E} \oplus 0.45\%$	1996
CsI(Tl) (BaBar)	$16-18X_0$	$2.3\%/E^{1/4} \oplus 1.4\%$	1999
CsI(Tl) (BELLE)	$16X_0$	1.7% for $E_\gamma > 3.5$ GeV	1998
PbWO ₄ (PWO) (CMS)	$25X_0$	$3\%/\sqrt{E} \oplus 0.5\% \oplus 0.2/E$	1997
Lead glass (OPAL)	$20.5X_0$	$5\%/\sqrt{E}$	1990
Liquid Kr (NA48)	$27X_0$	$3.2\%/\sqrt{E} \oplus 0.42\% \oplus 0.09/E$	1998
Scintillator/depleted U (ZEUS)	$20-30X_0$	$18\%/\sqrt{E}$	1988
Scintillator/Pb (CDF)	$18X_0$	$13.5\%/\sqrt{E}$	1988
Scintillator fiber/Pb spaghetti (KLOE)	$15X_0$	$5.7\%/\sqrt{E} \oplus 0.6\%$	1995
Liquid Ar/Pb (NA31)	$27X_0$	$7.5\%/\sqrt{E} \oplus 0.5\% \oplus 0.1/E$	1988
Liquid Ar/Pb (SLD)	$21X_0$	$8\%/\sqrt{E}$	1993
Liquid Ar/Pb (H1)	$20-30X_0$	$12\%/\sqrt{E} \oplus 1\%$	1998
Liquid Ar/depl. U (DØ)	$20.5X_0$	$16\%/\sqrt{E} \oplus 0.3\% \oplus 0.3/E$	1993
Liquid Ar/Pb accordion (ATLAS)	$25X_0$	$10\%/\sqrt{E} \oplus 0.4\% \oplus 0.3/E$	1996

Homogeneous

Sampling

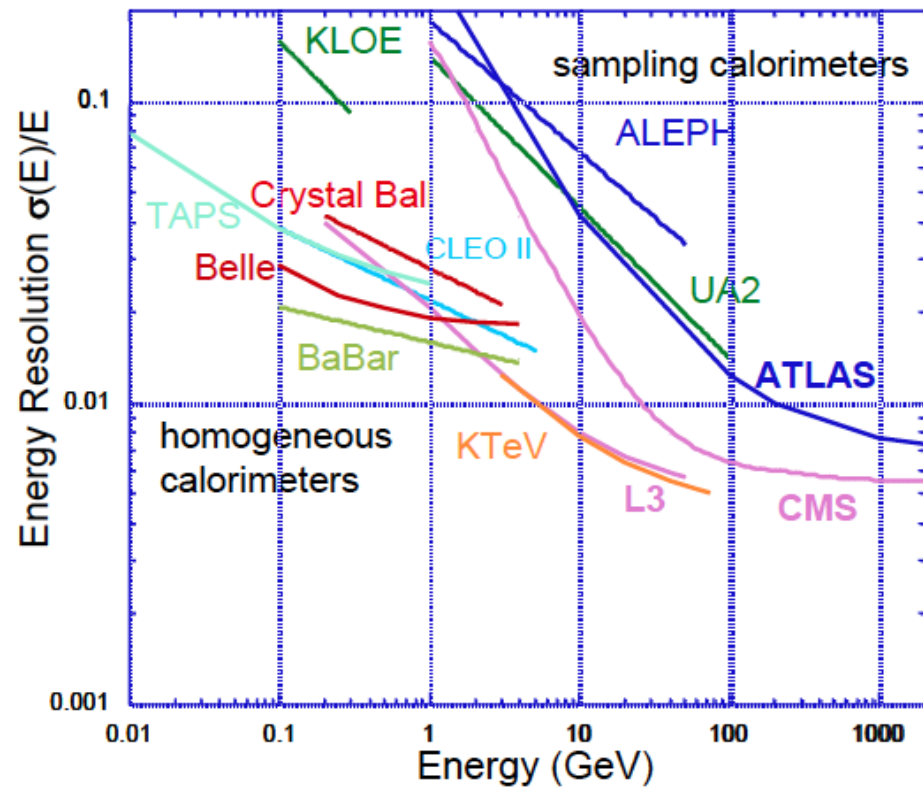
Homogeneous vs. Sampling Calorimeter

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Bi ₄ Ge ₃ O ₁₂ (BGO) (L3)	22X ₀	2%/√E ⊕ 0.7%	1993
CsI (KTeV)	27X ₀	2%/√E ⊕ 0.45%	1996
CsI(Tl) (BaBar)	16–18X ₀	2.3%/E ^{1/4} ⊕ 1.4%	1999
CsI(Tl) (BELLE)	16X ₀	1.7% for E _γ > 3.5 GeV	1998
PbWO ₄ (PWO) (CMS)	25X ₀	3%/√E ⊕ 0.5% ⊕ 0.2/E	1997
Lead glass (OPAL)	20.5X ₀	5%/√E	1990
Liquid Kr (NA48)	27X ₀	3.2%/√E ⊕ 0.42% ⊕ 0.09/E	1998
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Liquid Ar/Pb (NA31)	27X ₀	7.5%/√E ⊕ 0.5% ⊕ 0.1/E	1988
Liquid Ar/Pb (SLD)	21X ₀	8%/√E	1993
Liquid Ar/Pb (H1)	20–30X ₀	12%/√E ⊕ 1%	1998
Liquid Ar/depl. U (DØ)	20.5X ₀	16%/√E ⊕ 0.3% ⊕ 0.3/E	1993
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Homogeneous

Sampling



Energy Resolution

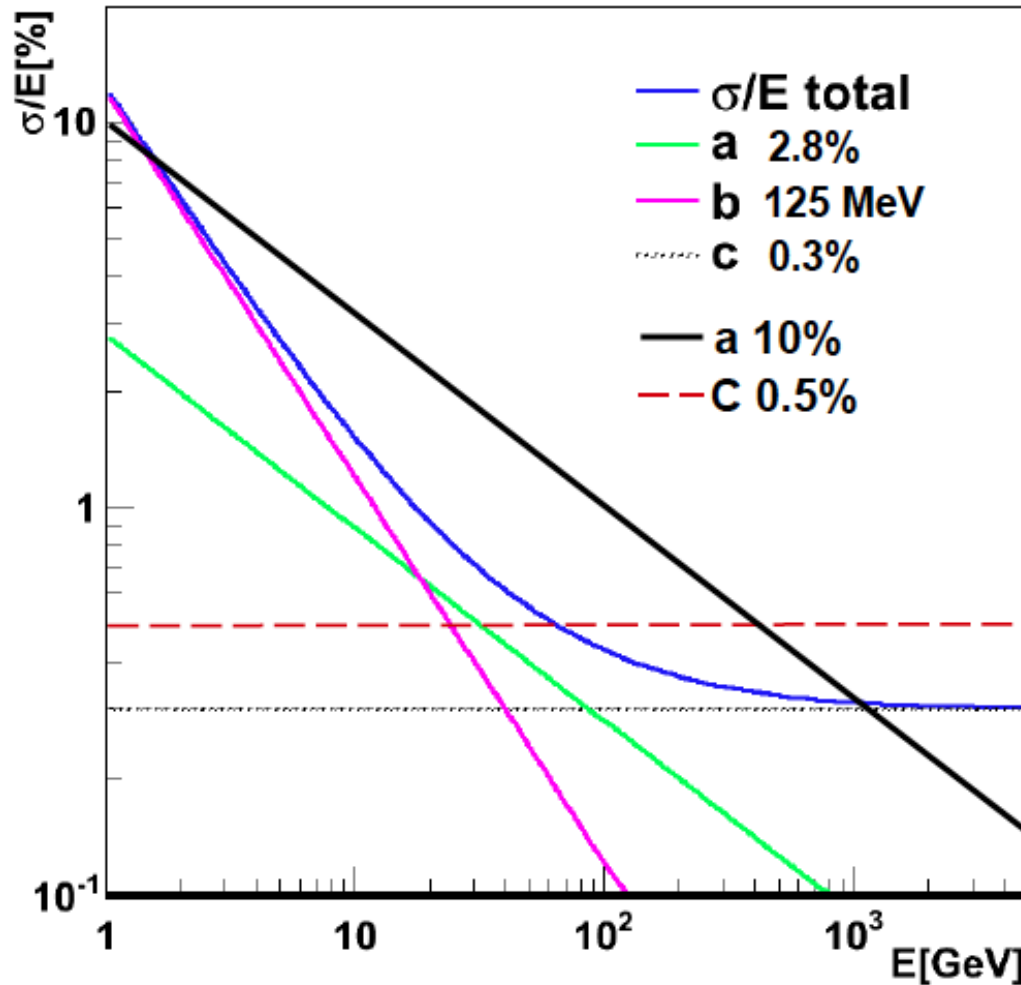
- Energy resolution of a calorimeter can be parameterized as

- Note the quadratic sum

$$\frac{\sigma_E}{E} = \frac{a}{\sqrt{E}} \oplus \frac{b}{E} \oplus c$$

- **a**: the stochastic term accounts for any kind of Poisson-like fluctuations
 - Natural merit of homogeneous calorimeters
 - Several contributions add to the “intrinsic one”
- **b**: the noise term responsible for degradation of low energy resolution
 - Mainly the energy equivalent of the electronic noise
 - Contribution from pileup: the fluctuation of energy entering the measurement area from sources other than the primary particle
- **c**: the constant term dominates at high energy
 - Its relevance is strictly connected to the small value of **a**
 - It is mostly dominated by the stability of calibration
 - Contributions from energy leakage, non-uniformity of signal generation and/or collection, loss of energy in dead materials

When do we have to worry about c



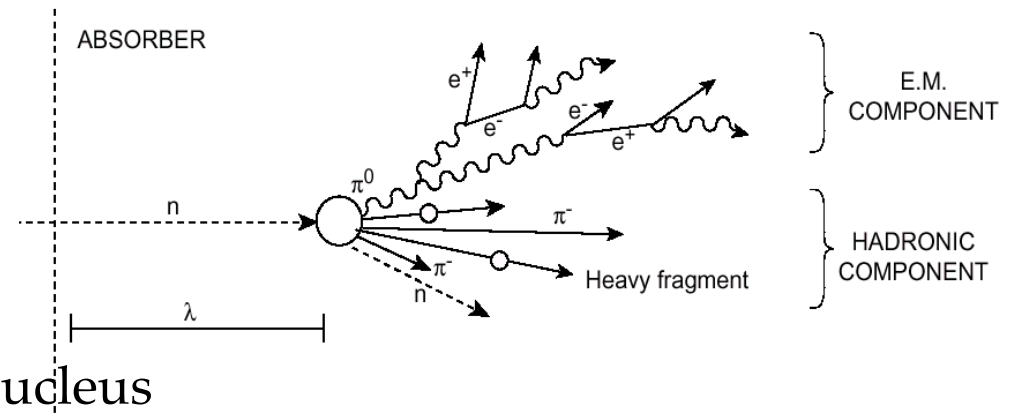
Hadronic Shower

- Hadronic interaction

- Elastic: $p + \text{Nucleus} \rightarrow p + \text{Nucleus}$
- Inelastic: $p + \text{Nucleus} \rightarrow \pi^+ + \pi^- + \pi^0 + \text{Nucleus}^*$

- Shower development

- $p + \text{Nucleus} \rightarrow \text{Pions} + N^* + \dots$
- Secondary particles undergo further elastic collisions until they fall below pion production threshold, $E \sim 2m_\pi = 0.28 \text{ GeV}$
- Sequential decays
 - $\pi^0 \rightarrow \gamma\gamma$ yields electromagnetic shower
 - Fission fragments $\rightarrow \beta$ -decay, γ -decay
 - Neutron capture \rightarrow fission
 - Spallation ...



Hadronic Shower

- Hadronic interaction length

$$\lambda_{\text{int}} = \frac{1}{\sigma_{\text{tot}} \cdot n} = \frac{A}{\sigma_{pp} A^{2/3} \cdot N_A \rho} \sim A^{1/3} \approx 35 \frac{\text{g}}{\text{cm}^2} \cdot A^{1/3} \quad N(x) = N_0 \exp\left(-\frac{x}{\lambda_{\text{int}}}\right)$$

Some numerical values for materials typical used in hadron calorimeters

- Interaction length characterizes both, longitudinal and transverse profile of hadronic showers

$$\left. \begin{array}{l} X_0 \sim \frac{A}{Z^2} \\ \lambda_{\text{int}} \sim A^{1/3} \end{array} \right\} \rightarrow \frac{\lambda_{\text{int}}}{X_0} \sim A^{4/3}$$

- Typical longitudinal size [95% containment]: $6 \dots 9 \lambda_{\text{int}}$
 - EM: 15-20 X_0
- Typical transverse size [95% containment]: $1 \lambda_{\text{int}}$
 - EM: $2 R_M$, compact
- Hadronic calorimeter needs more depth than electromagnetic calorimeter

	λ_{int} [cm]	X_0 [cm]
Szint.	79.4	42.2
LAr	83.7	14.0
Fe	16.8	1.76
Pb	17.1	0.56
U	10.5	0.32
C	38.1	18.8

Hadronic Shower

- Hadronic interaction length

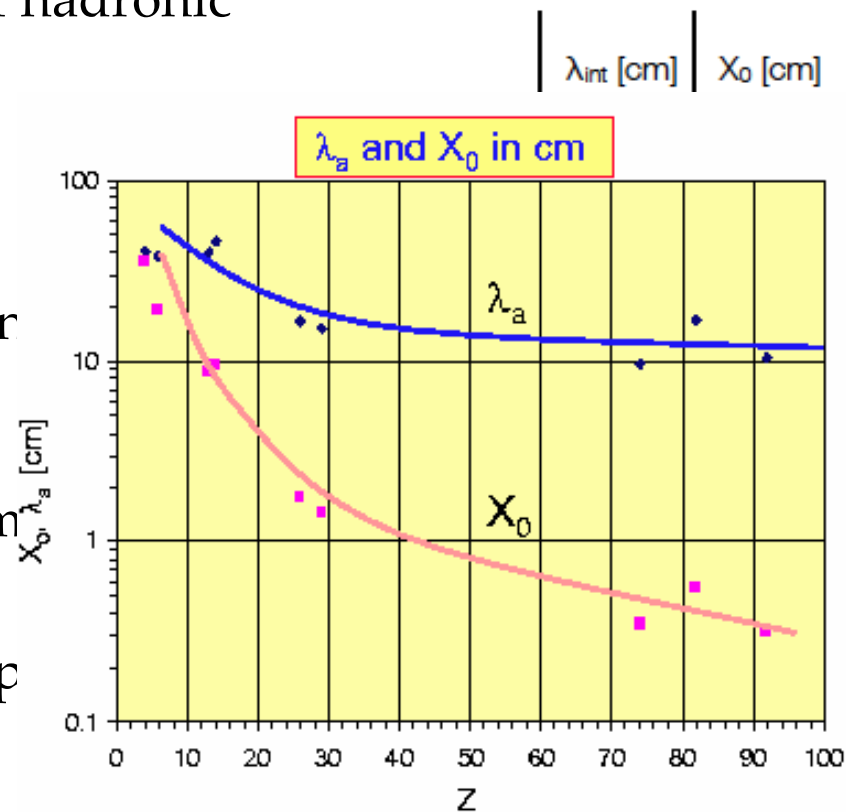
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- Typical longitudinal size [95% contain ... $9 \lambda_{\text{int}}$
 - EM: 15-20 X_0
- Typical transverse size [95% containm
 - EM: 2 R_M , compact
- Hadronic calorimeter needs more dep electromagnetic calorimeter

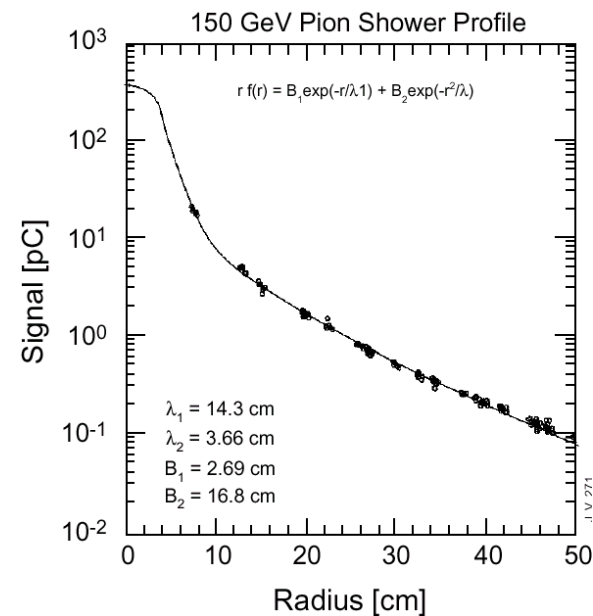
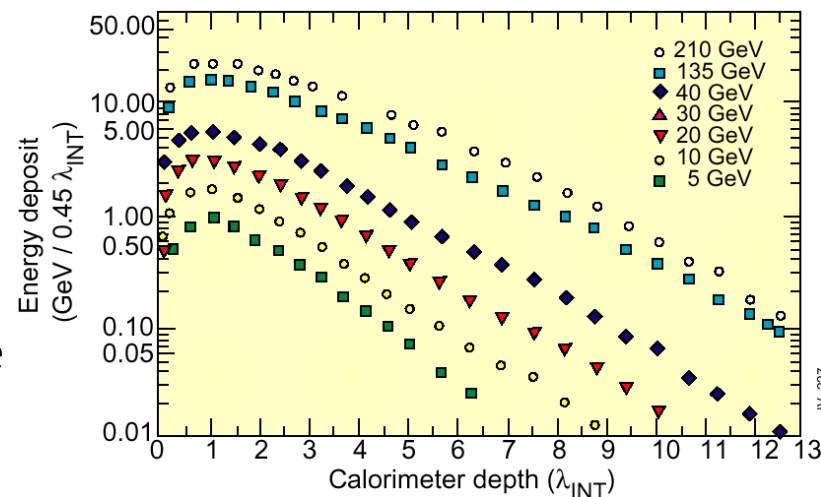


Hadronic Shower

- Longitudinal
 - Sharp peak from π^0 from the first interaction
 - Gradual extinction with typical scale
 - Need to sample

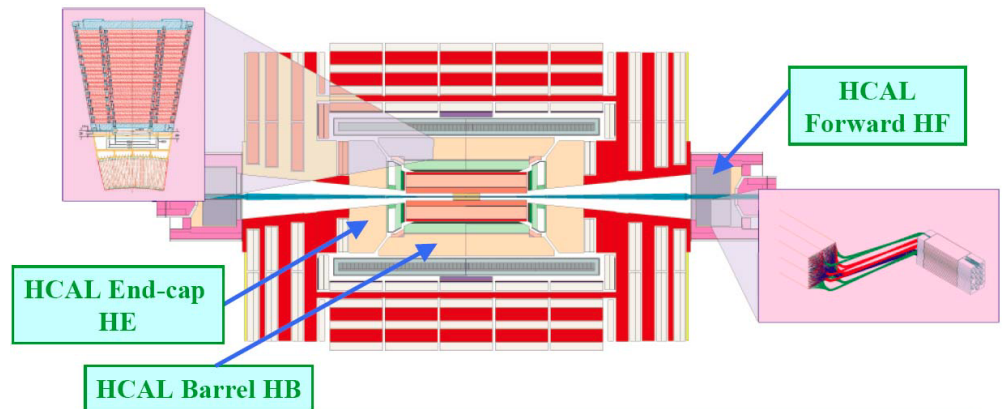
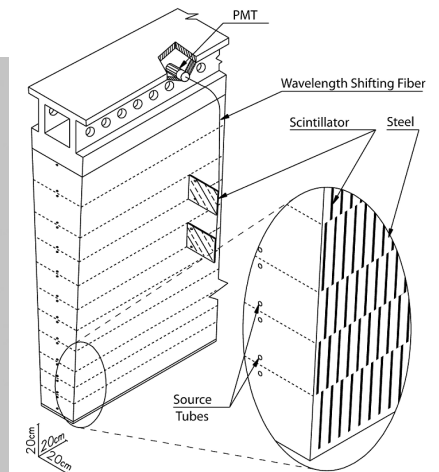
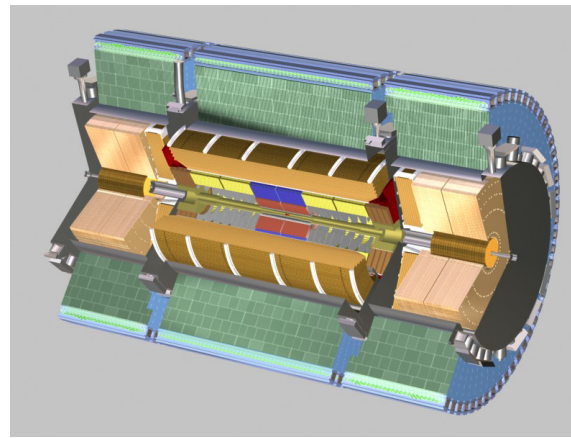
- Lateral
 - Average p_t secondaries ~ 300 MeV
 - Typical transverse scale λ_{int} for 95% E containment
 - Dense core due to π^0

WA78 : 5.4λ of 10mm U / 5mm Scint + 8λ of 25mm Fe / 5mm Scint



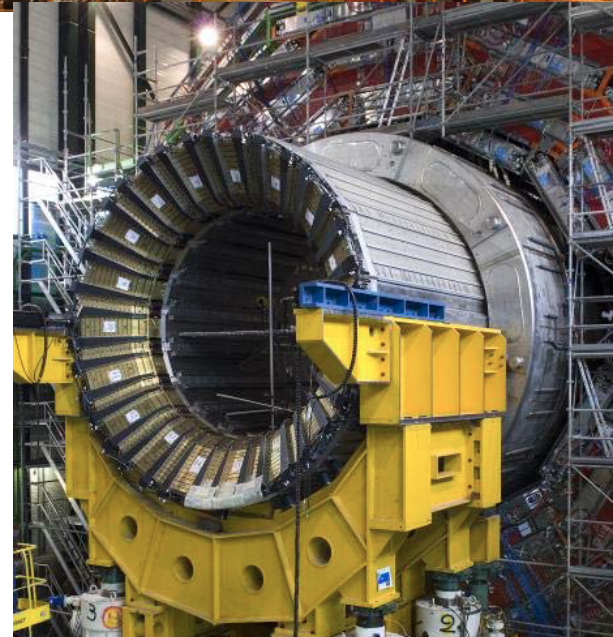
Hadronic Calorimeter

- Most common realization – Sampling
- ATLAS Tile Calorimeter
 - Iron/plastic scintillator sampling calorimeter
- CMS HCAL
 - Barrel and Endcap are brass/scintillator sampling calorimeter
 - Forward is steel/quartz fibers sampling calorimeter



Hadronic Calorimeter

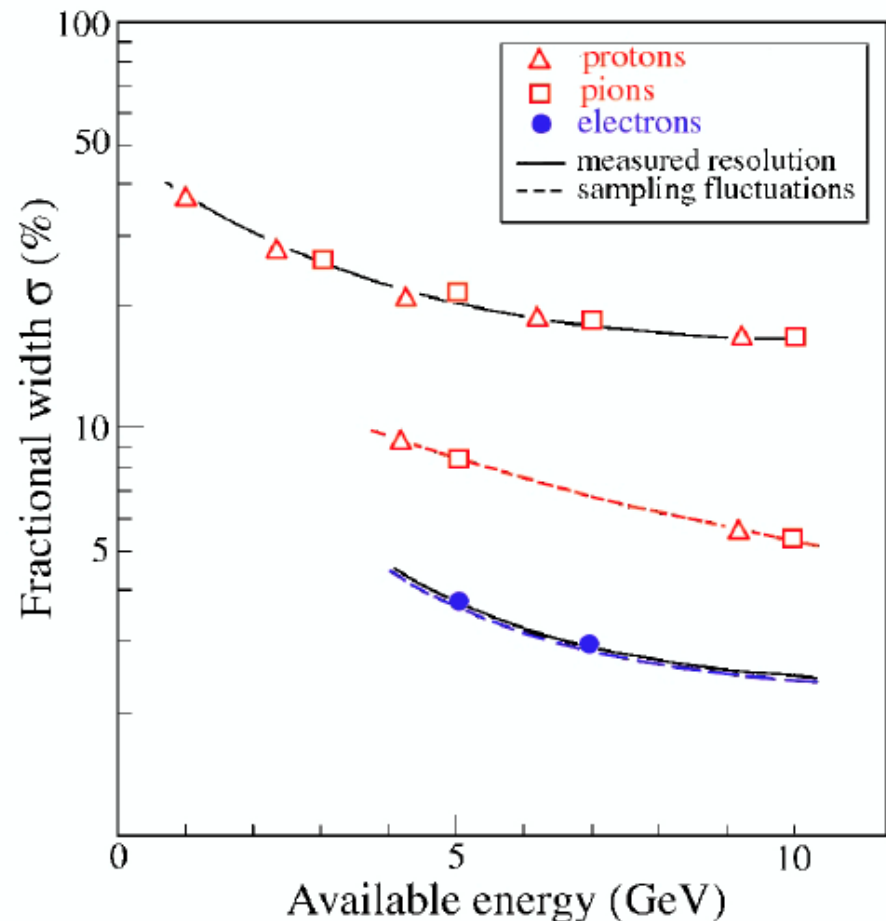
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Energy Resolution

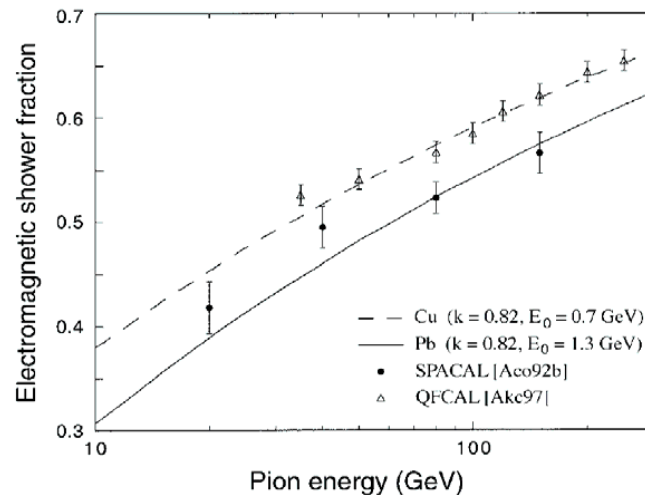
$$\frac{\sigma_E}{E} = \frac{A}{\sqrt{E}} \oplus \frac{B}{E} \oplus C$$

- A – typical 0.5 – 1.0 [record: 0.35]
 - Leakage fluctuations, sampling fluctuation, fluctuation of EM fraction, nuclear excitations, fission, binding energy fluctuations, heavily ionizing particles
- B – typical few%
 - Sampling fraction variations, electronic noise
- C – typical 0.03 – 0.05
 - Inhomogeneous shower leakage

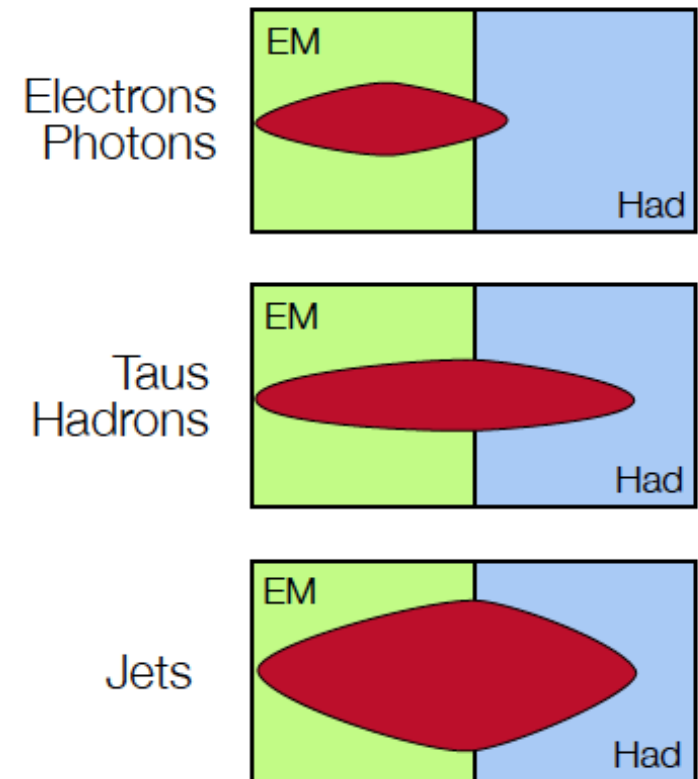


Hadronic Calorimeter

- Typical calorimeter has two components
 - Electromagnetic (EM)
 - Hadronic (Had)
- Hadronic energy measured in both parts of calorimeter
 - Needs careful consideration of different response
 - A priori e and h in a calorimeter give a different response, e.g. $e/h > 1$
 - The fluctuations in the fraction of energy deposited by e and h limits resolution
 - Moreover in average this fraction is energy dependent

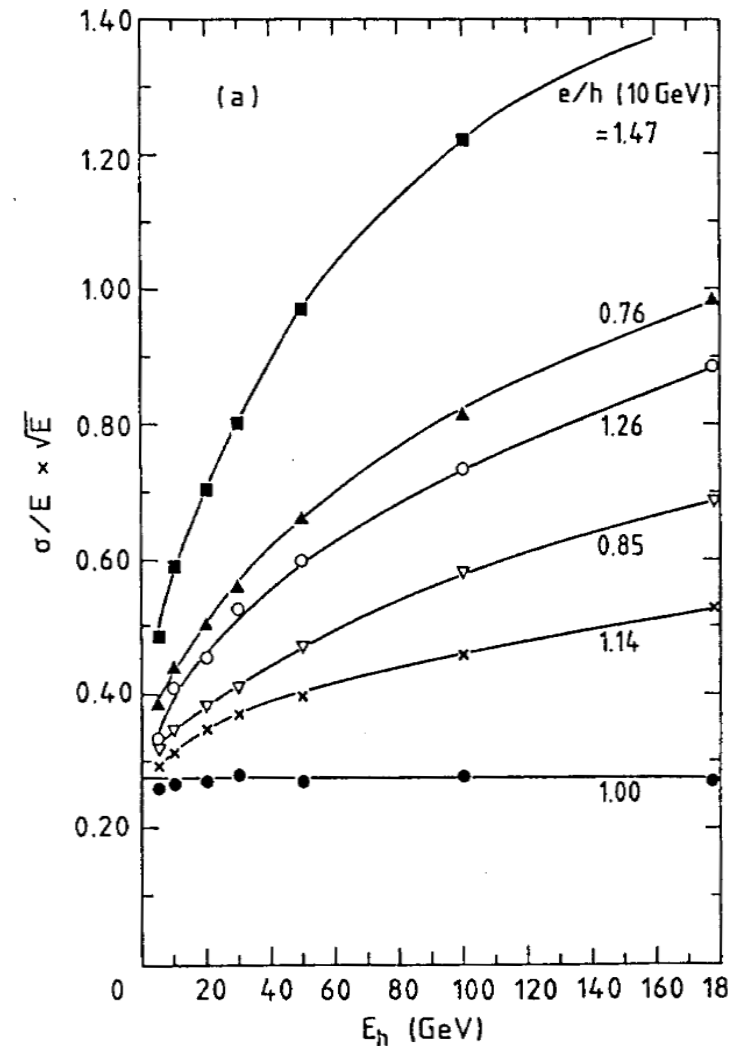


Schematic of a typical HEP calorimeter



Hadronic Calorimeter

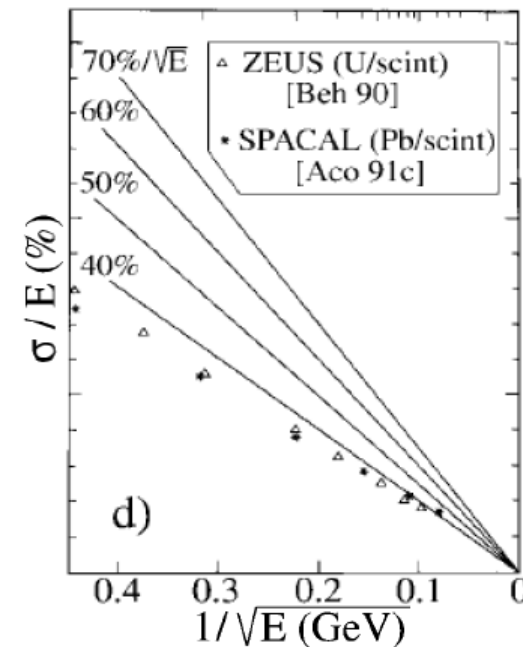
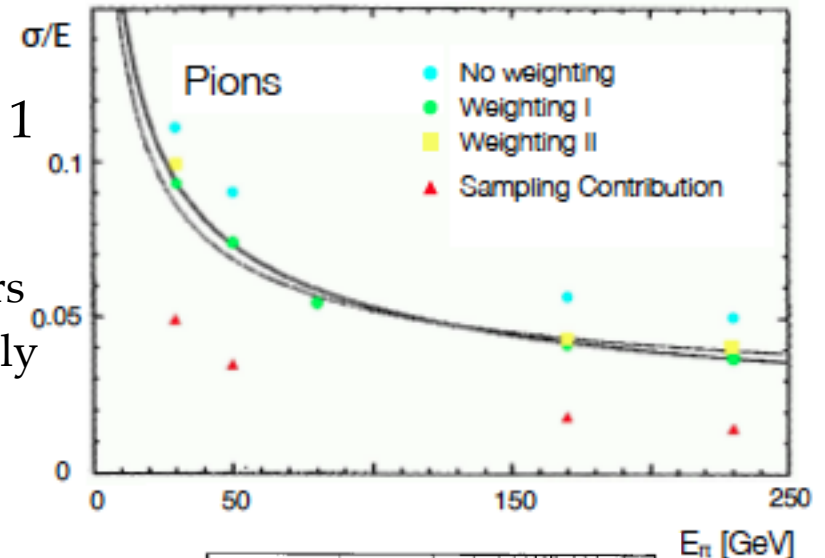
- e/h ratio
 - Response of calorimeters very different to EM and Had energy deposits
 - Usually higher weight for EM component, e.g. $e/h > 1$
- $e/h \neq 1$ leads to
 - Non-uniform energy response due to fluctuations in f_{em}
 - Non-linear behavior
 - Worsening of resolution
 - Deviation from $1/\sqrt{E}$ resolution dependence



Hadronic Calorimeter

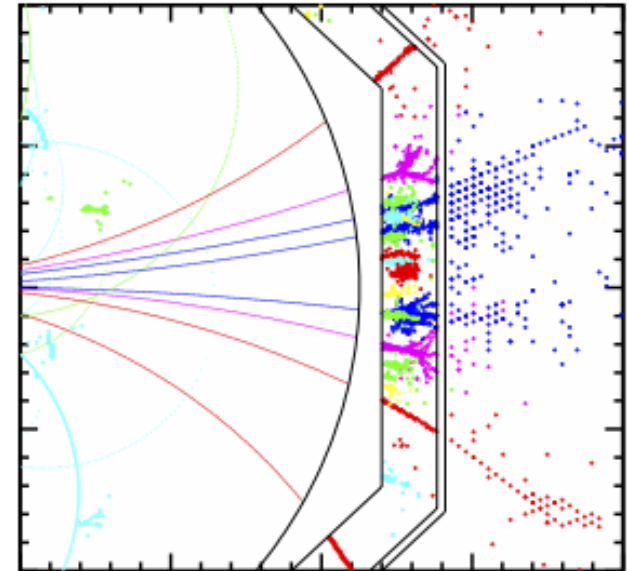
- Compensation is important to get $e/h = 1$
- Software compensation
 - Segmentation of calorimeter into cells/layers
 - Identification of cells/layers with particularly large energy
 - Give small weight to layers with larger energy density
- Hardware compensation
 - Choose suitable hardware parameters to either increase h/mip or decrease e/mip
 - Suppress EM component with high Z absorber
 - Enhance h production through fission and spallation
 - Enhance response to n using active materials hydrogen rich

Energy resolution of LAr calorimeter with and without weighting ...



New Calorimeter Development

- PFA – Particle Flow Approach
 - Imaging calorimeter for Linear Collider
 - Highly granular calorimeter can identify the energy deposit of every visible particle (10^7 to 10^{12} readout channels)
 - The energy of charged hadrons will be measured by the tracking detectors
 - The energy of photons will be measured by the electromagnetic calorimeter
 - The hadronic calorimeter is then used only to measure the energy of neutral hadrons
 - The reduced dependence on the hadronic calorimeter will lead to an unprecedented jet energy resolution ($\sim 25\%$)



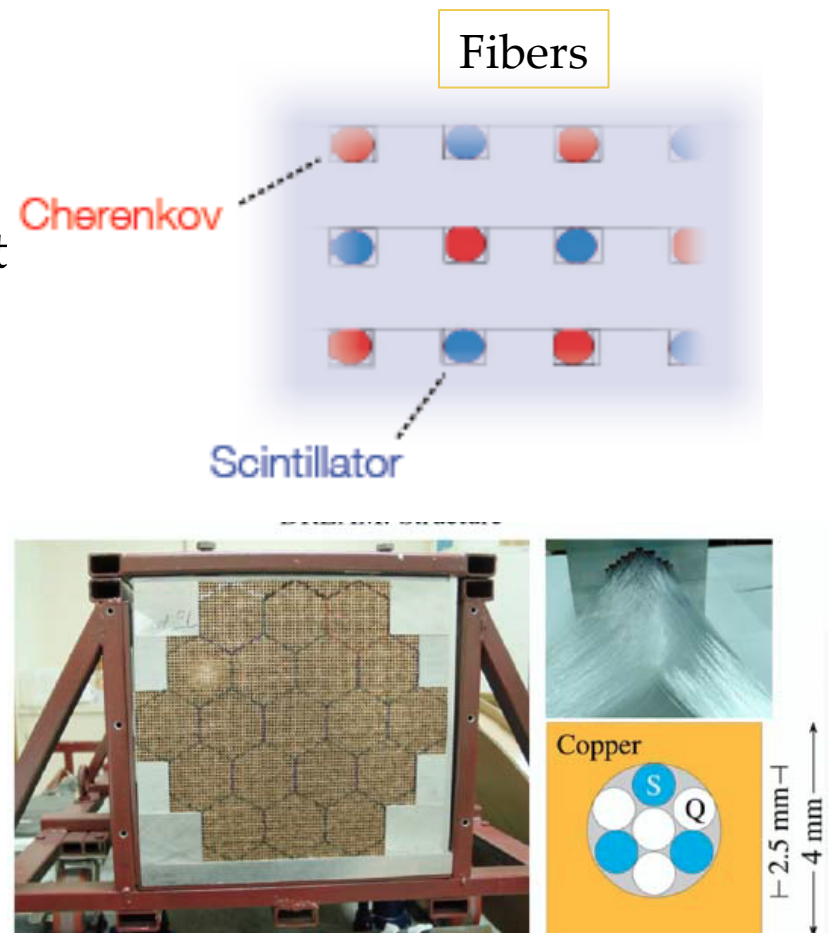
Component	Detector	Fraction	Part. resolution	Jet Energy Res.
Charged (X^\pm)	Tracker	60%	$10^{-4} E_x$	negligible
Photons (γ)	ECAL	30%	$0.1/\sqrt{E_\gamma}$	$.06/\sqrt{E_{jet}}$
Neutral Hadrons (h)	E/HCAL	10%	$0.5/\sqrt{E_{had}}$	$.16/\sqrt{E_{jet}}$

New Calorimeter Development

- Dual Readout Calorimetry
 - Cherenkov assisted Hadron Calorimetry
 - EM: clear fibers for Cherenkov light to sample EM part of the shower [$E_{\text{Cherenkov}}$]
 - Charged: scintillation fibers to sample all components [$E_{\text{Ionization}}$]

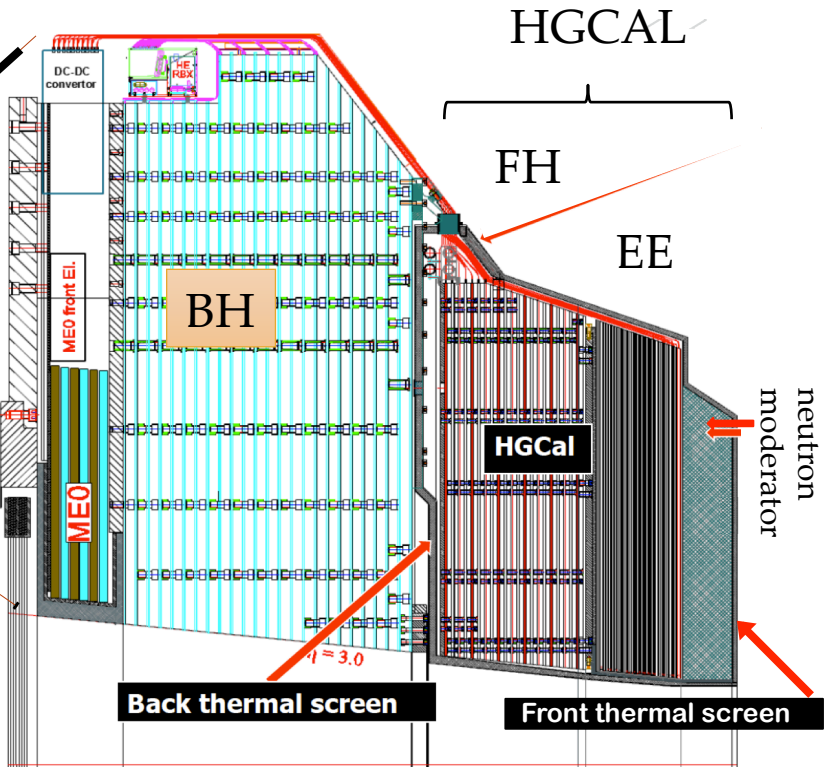
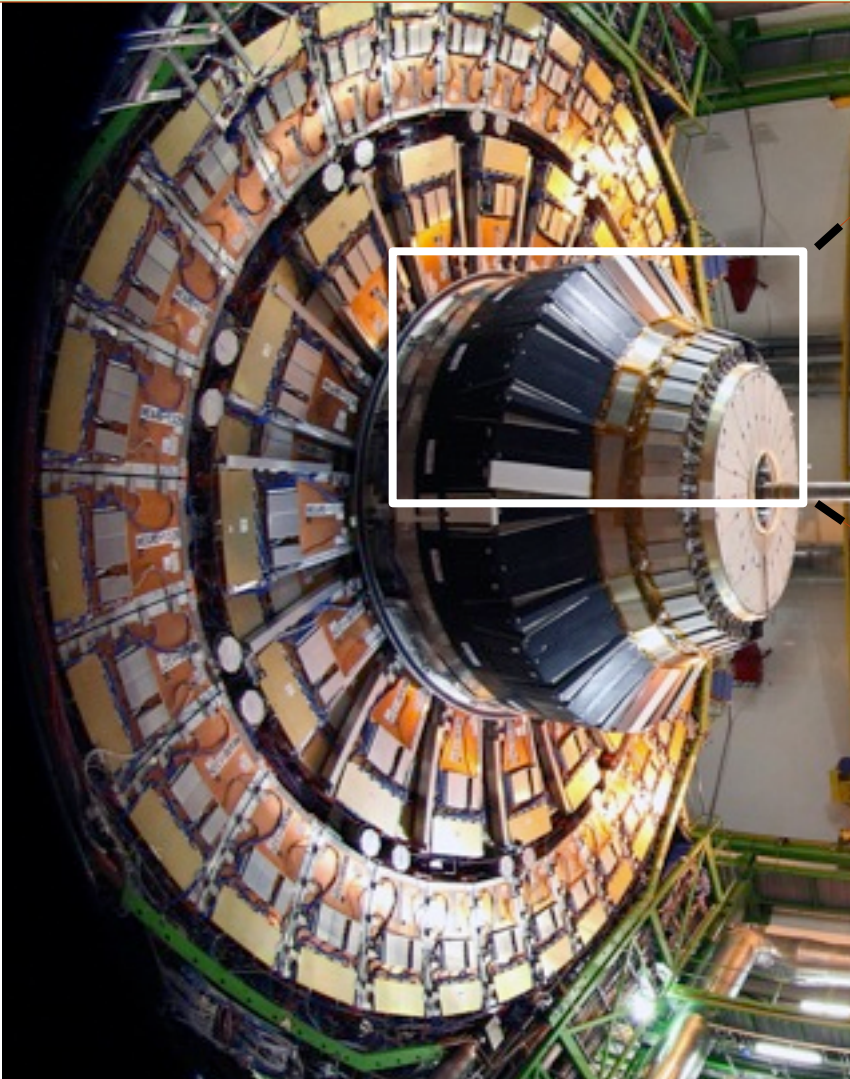
$$\frac{E_{em}}{E_{tot}} = \frac{E_{Cherenkov}}{E_{Ionization}}$$

- RD52 – DREAM (Dual-REAdout Method)
 - Measure f_{em} event by event to improve hadronic energy resolution



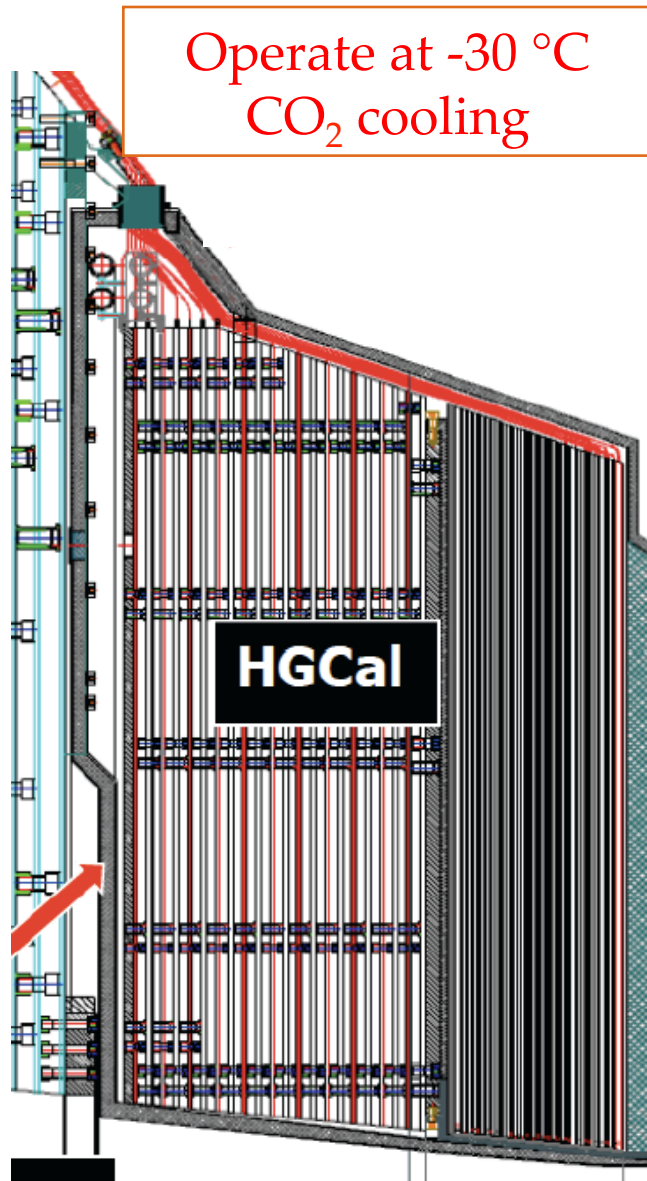
Endcap Calorimeter for HL-LHC: High Granularity Cal

Integrated sampling Silicon ECAL+HCAL and Backing Calorimeters



EE	Cu-W / Si	$26 X_0$ (1.5λ)
FH	Brass / Si	3.5λ
BH	Brass / scint. tiles	5λ

Endcap Calorimeter for HL-LHC: HGCal



Operate at $-30\text{ }^{\circ}\text{C}$
 CO_2 cooling

Si/W-ECAL Section ($\Sigma_{\text{depth}} > 25X_0, 1.5\lambda$)

$10 \times 0.65X_0$

$10 \times 0.88X_0$

$8 \times 1.26X_0$

Si/Brass Front HCAL (FH) Section ($\Sigma_{\text{depth}} > 3.5\lambda$)

$12 \times 0.3\lambda$

Scint/Brass Backing HCAL (BH) Section ($\Sigma_{\text{depth}} > 5\lambda$)

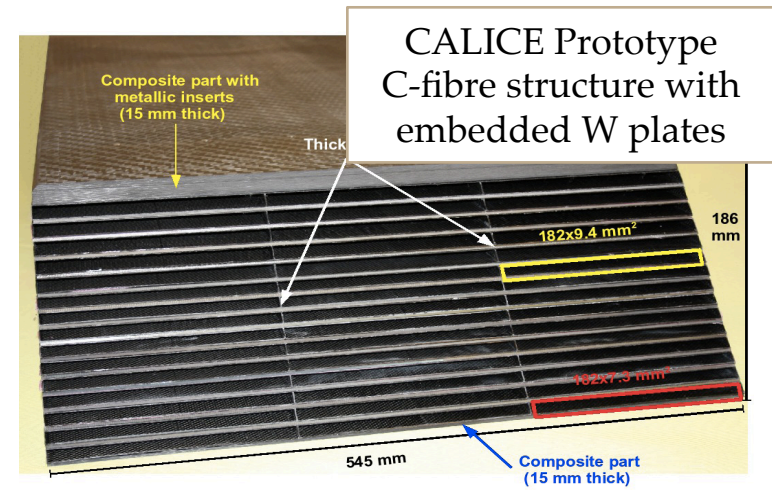
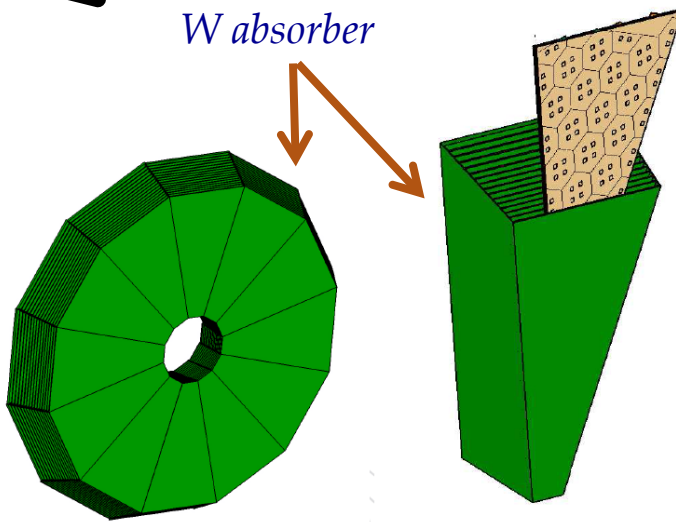
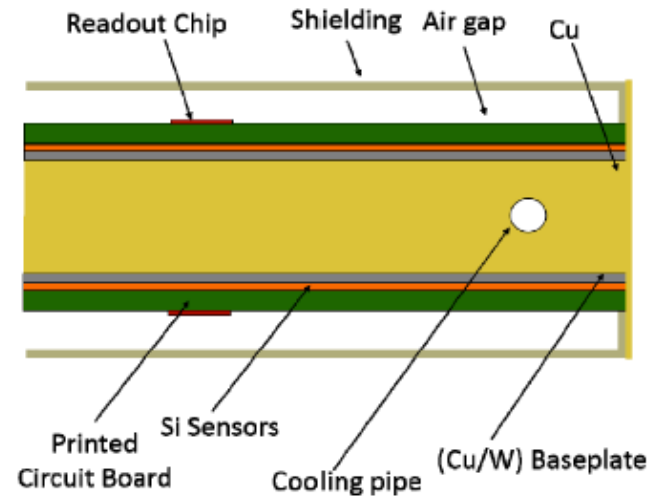
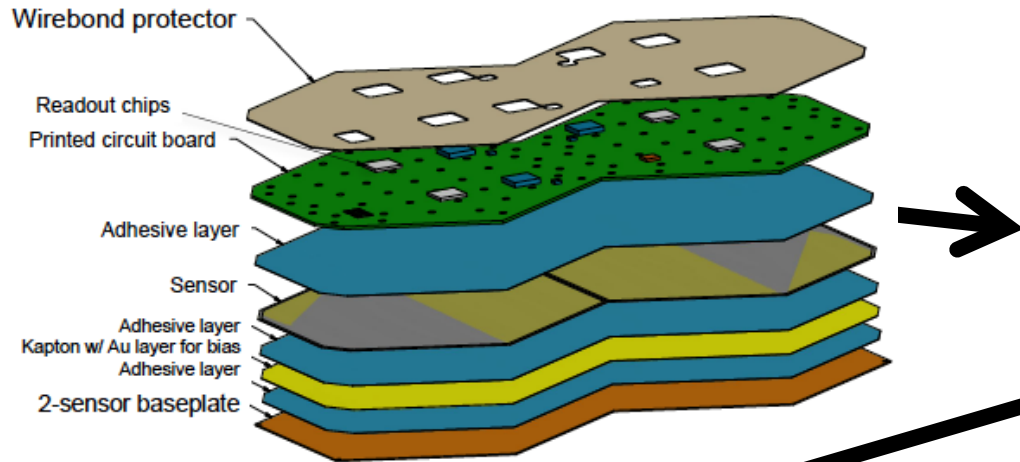
$12 \times 0.45\lambda$

Total Depth $> 10\lambda$

Table 3.2: Parameters of the EE and FH.

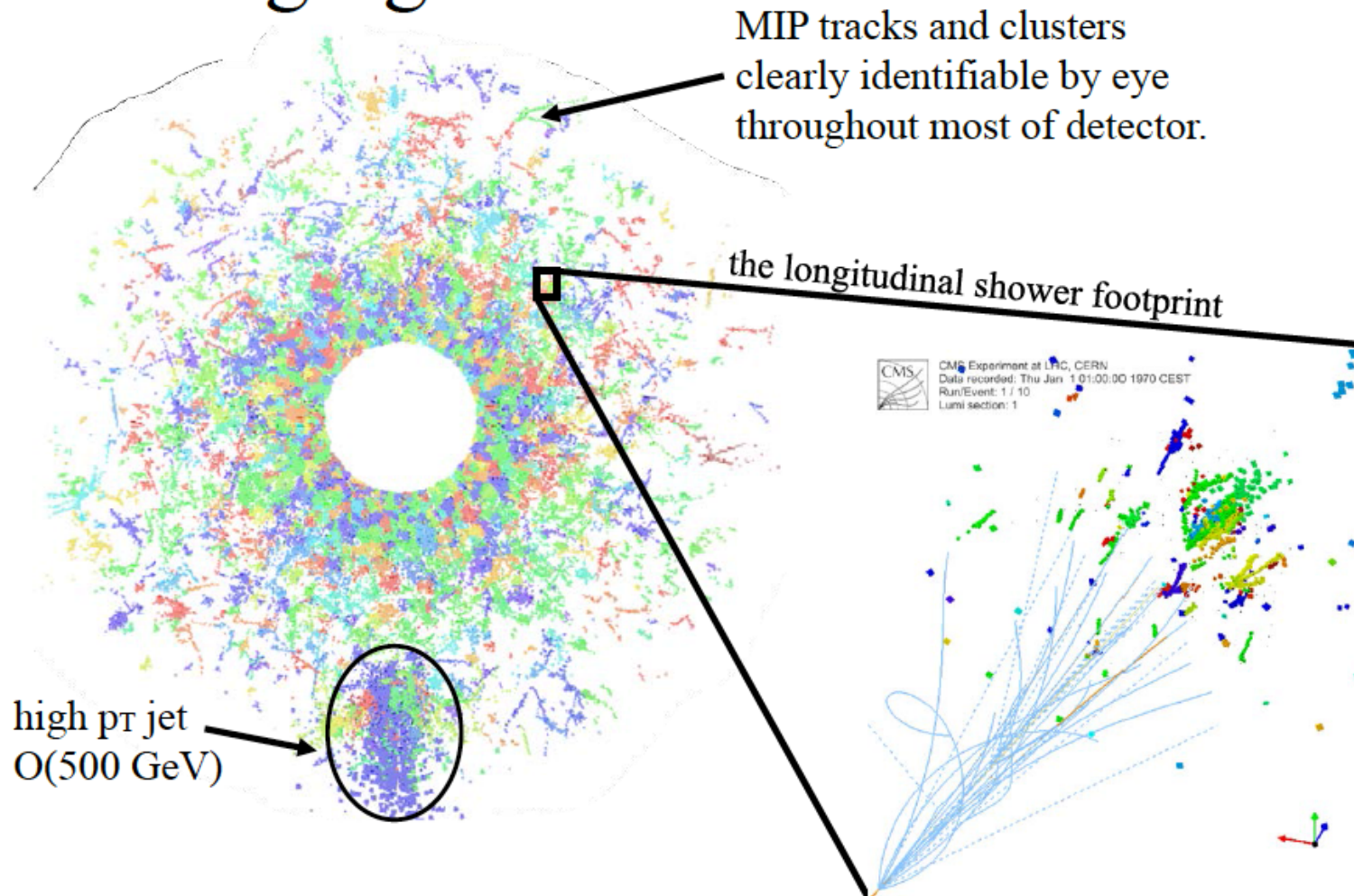
	EE	FH	Total
Area of silicon (m^2)	380	209	589
Channels	4.3M	1.8M	6.1M
Detector modules	13.9k	7.6k	21.5k
Weight (one endcap) (tonnes)	16.2	36.5	52.7
Number of Si planes	28	12	40

Endcap Calorimeter for HL-LHC: HGCal



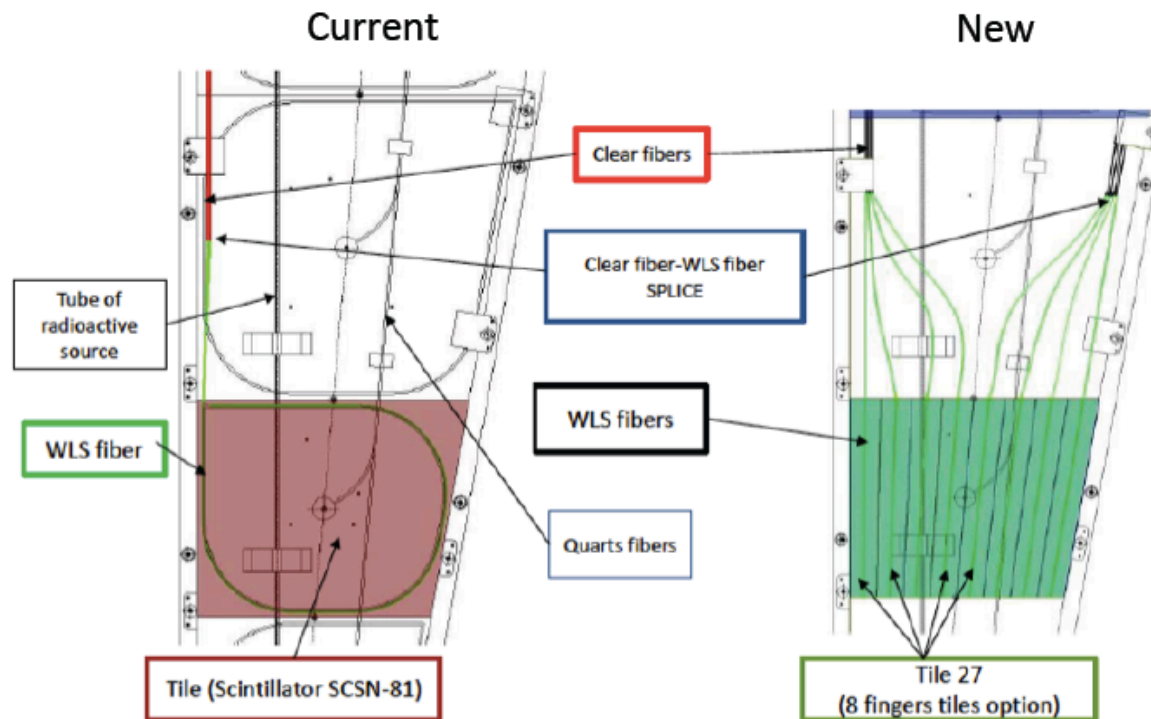
Net effect: extends tracking into calorimeter!

Imaging Showers with the HGC



Endcap Calorimeter: Backing Hadronic

- Improvement of current HE tiles for ≈ 5 Mrad tolerance
 - Doubly-doped plastic scintillator x 2 light after irradiation
 - Finger tile design - shorter light path



Synergy with Barrel HCAL upgrade – same scintillator, different geometry depending on location (i.e. expected radiation doses)

- And also increased granularity $\approx x 2$ in Φ & $x 1.3$ in η

The Invention of Liquid Argon Calorimeter

- By Bill Willis and Veljko Radeka in 1974

NUCLEAR INSTRUMENTS AND METHODS 120 (1974) 221-236; © NORTH-HOLLAND PUBLISHING CO.

LIQUID-ARGON IONIZATION CHAMBERS AS TOTAL-ABSORPTION DETECTORS*

W. J. WILLIS†

Department of Physics, Yale University, New Haven, Connecticut 06520, U.S.A.

and

V. RADEKA

Instrumentation Division, Brookhaven National Laboratory, Upton, New York 11973, U.S.A.

Received 14 May 1974

A new detector for the measurement of energy by total absorption, based on the use of multiple-plate ion chambers, is described. The use of liquid argon as the working medium and optimized readout results in an electronic noise contribution to the resolution of less than 0.1 GeV, in a large detector. The use of thin plates,

0.1 radiation length, ensures that sampling fluctuations are small. The technique allows absolute calibration and very good gain stability. Tests on a detector large enough to absorb a high-energy electromagnetic shower are described, where the energy resolution is limited by the residual sampling fluctuations.

1. Principles and limitations of calorimetric detectors

If a high-energy particle enters a sufficiently large block of matter, all of its energy will be transformed into ionization and eventually into heat, with certain important exceptions. Thus, detectors relying on total

for incident electrons or photons. Thus, the first effect above can be appreciable for incident hadrons of low energy ($\lesssim 2$ GeV), but is always rather small for electrons. Effect (ii) comes mainly from positive pions which come to rest, which then convert about 135 MeV

The Invention of LAr Calorimeter

- By Bill Willis and Veljko Radeka in 1974

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$$\frac{\sigma_E}{E} = \frac{a}{\sqrt{E}} \oplus \frac{b}{E} \oplus c$$

Received 14 May 1974

A new detector for the measurement of energy by total absorption, based on the use of multiple-plate ion chambers, is described. The use of liquid argon as the working medium and optimized readout results in an electronic noise contribution to the resolution of less than 0.1 GeV, in a large detector. The use of thin plates,

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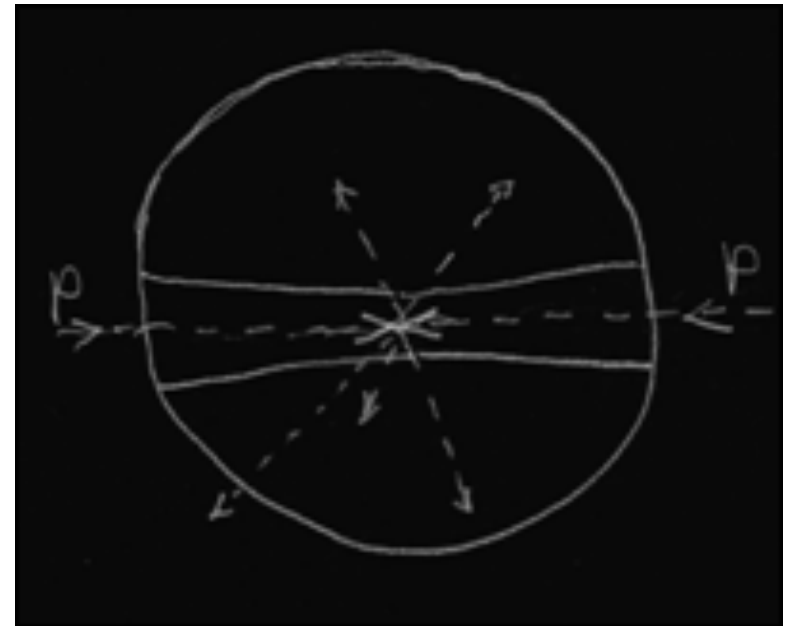
1. Principles and limitations of calorimetric detectors

If a high-energy particle enters a sufficiently large block of matter, all of its energy will be transformed into ionization and eventually into heat, with certain important exceptions. Thus, detectors relying on total

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How LAr Calorimeter was Invented

- Bill Willis, a professor at Yale University, was in sabbatical at BNL 1972-1973
- Bill met Veljko Radeka, a scientist/electronics engineer at Instrumentation Division of BNL, in October 1972
- Impactometer
 - A circle representing a sphere with two openings at the poles, through which beams of particles from opposite directions would collide in the center of the sphere
 - Measure the energy of *almost all* the debris that comes from such collisions
 - The idea was motivated by the high luminosity challenge of the future colliders. It was to be demonstrated at the CERN ISR



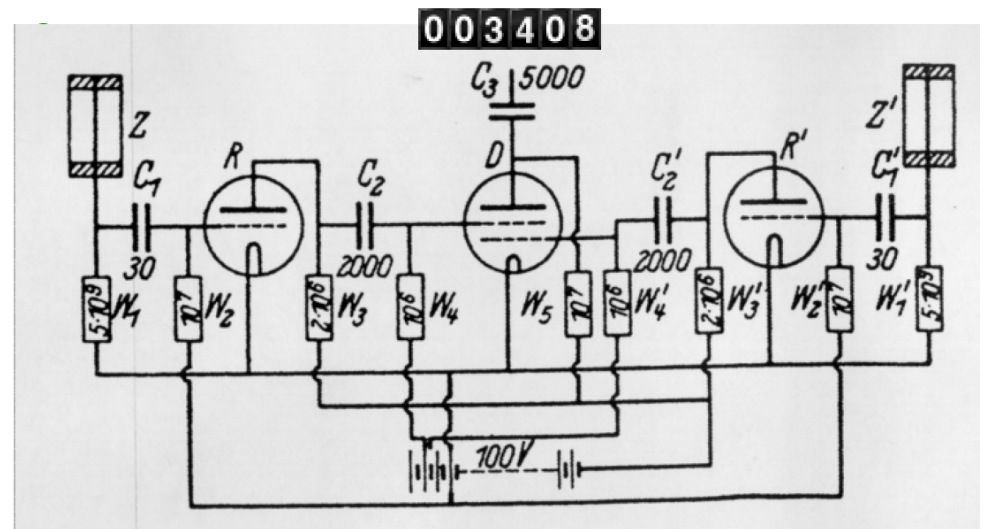
How to realize an “Impactometer”

- Bill Willis *“After working with a calorimeter made with steel and scintillator plates, I knew one could make a calorimeter covering all solid angle, but what about very good energy resolution and fine spatial resolution? I built a 1m long, LA(r) device and measured light, but it did not seem attractive. Detecting charge would be better, if charge gain isn’t needed. “*
- What: “Just a simple ionization chamber?”
- Detecting charge from ionization with high precision was not a strong trait of “counter detectors”, aka “electronic detectors” in particle physics at that time ...

Electronic detectors in particle physics then...

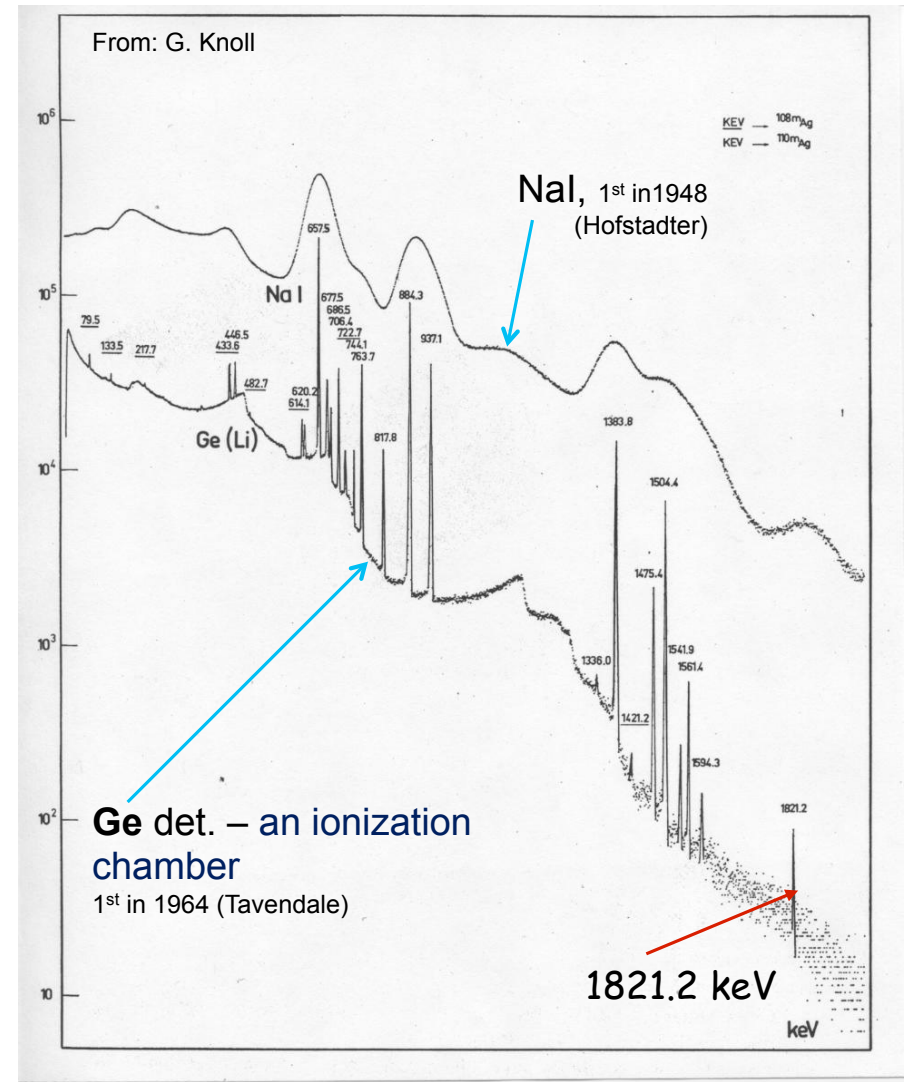
- Example: Coincidence and timing circuits to measure angular distributions and correlations
- **1960-ies: nanosecond** (“*millimicrosecond*”) resolution was achieved using *tunnel diodes* and new transistors
- Scintillators and photomultipliers had replaced Geiger-Miller tubes
- Spark chambers were producing Nobel prizes
- Radeka looked elsewhere how to read LAr ionization chambers...

‘Zur Vereinfachung von Koinzidenzzählungen’
(‘On the simplification of coincidence counting’)
W. Bothe, November 1929

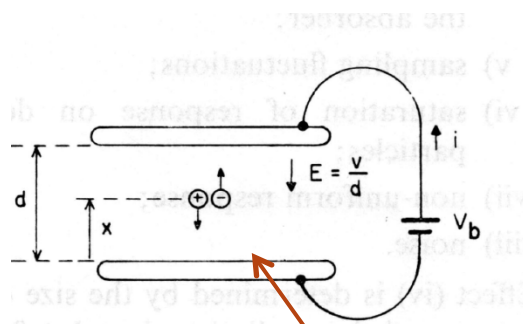


Detectors in nuclear physics

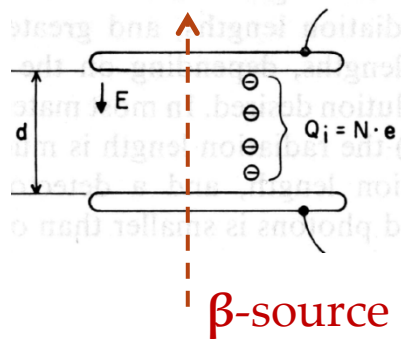
- Germanium vs. Sodium Iodide for gamma-ray spectrometry
- High precision ionization charge measurements, high precision calibration and low noise electronics were developed for germanium detectors in ~1965-1970
 - For gamma ray energy resolution of ~0.1% in the ~0.1 to 10 MeV range
 - This provided the basis for later use of such techniques in particle physics



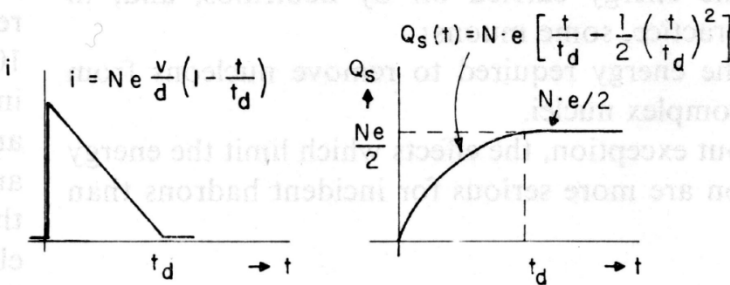
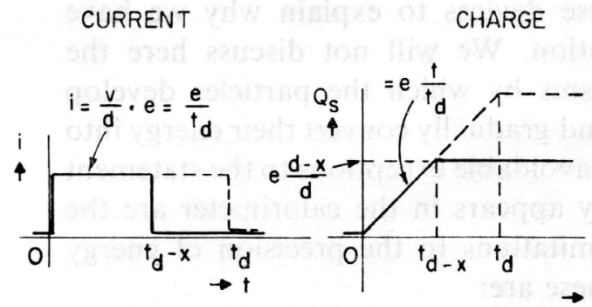
Dec 1972 – Jan 1973: Charge Collection and Drift Velocity Studies in LAr



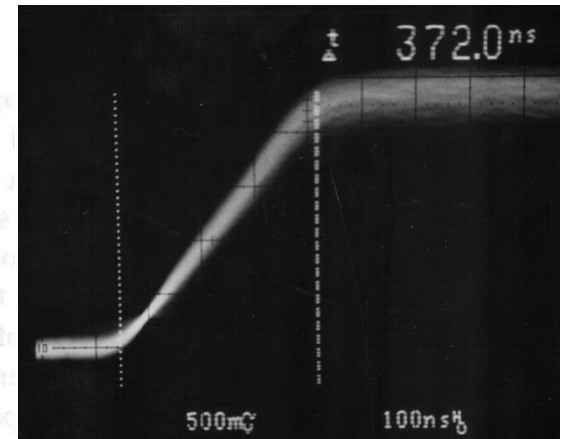
Electroplated α - source



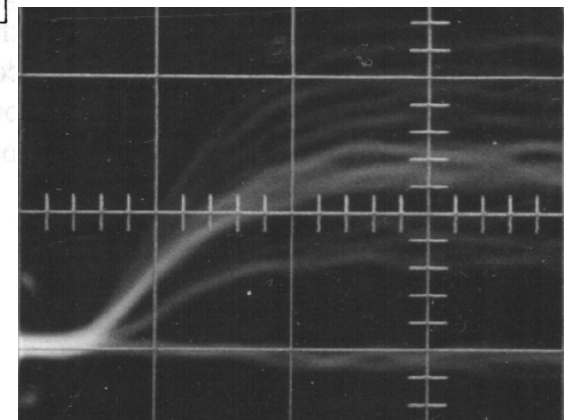
β -source



Only electrons induce a signal in the time of interest $< 1 \mu\text{s}$; mobility $\mu_{\text{electron}} \gg \mu_{\text{ion}}$

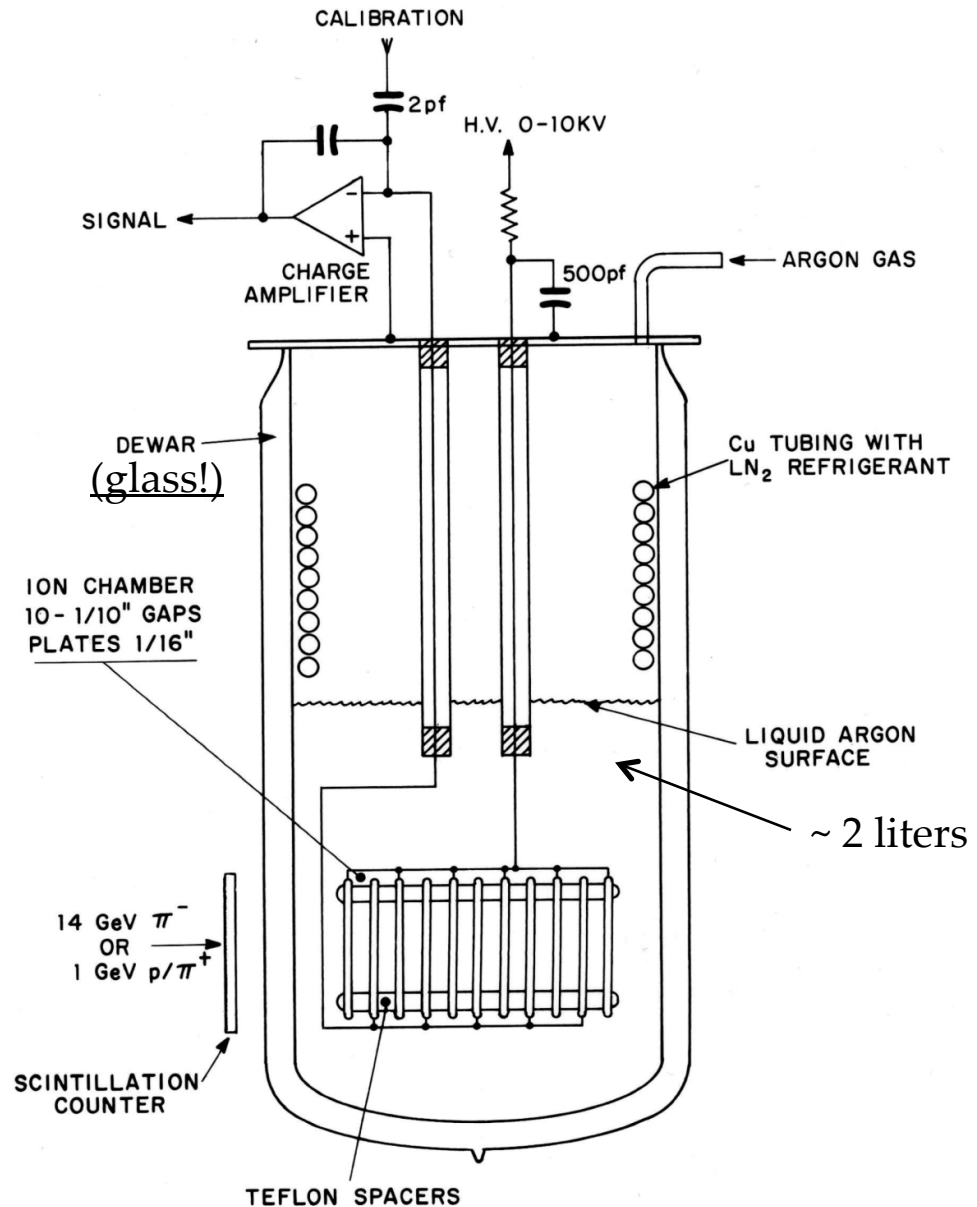


2 mm electrode spacing

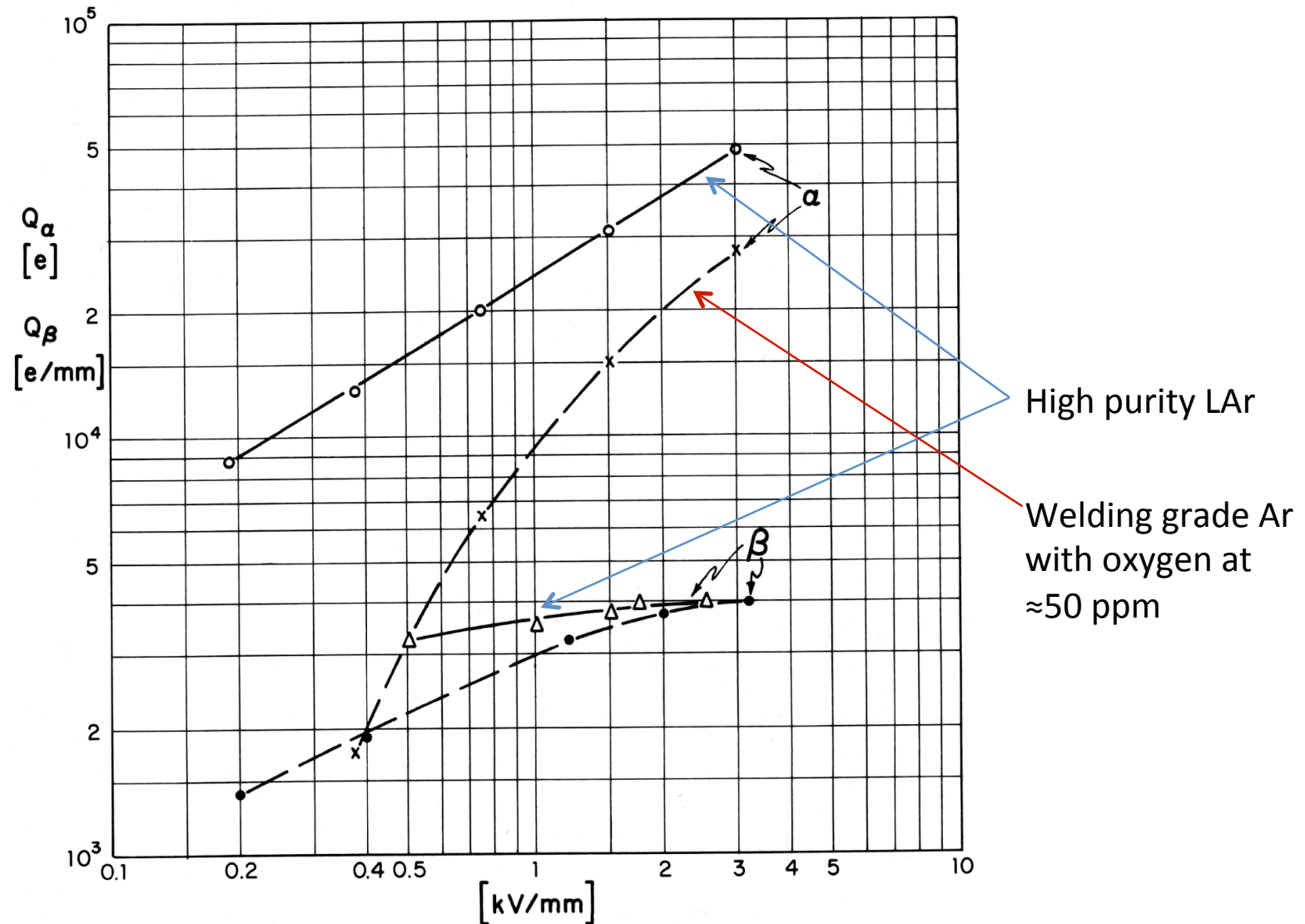


Dec 1972 – Jan 1973: Small test cells to study charge collection

Charge calibration
~0.1%

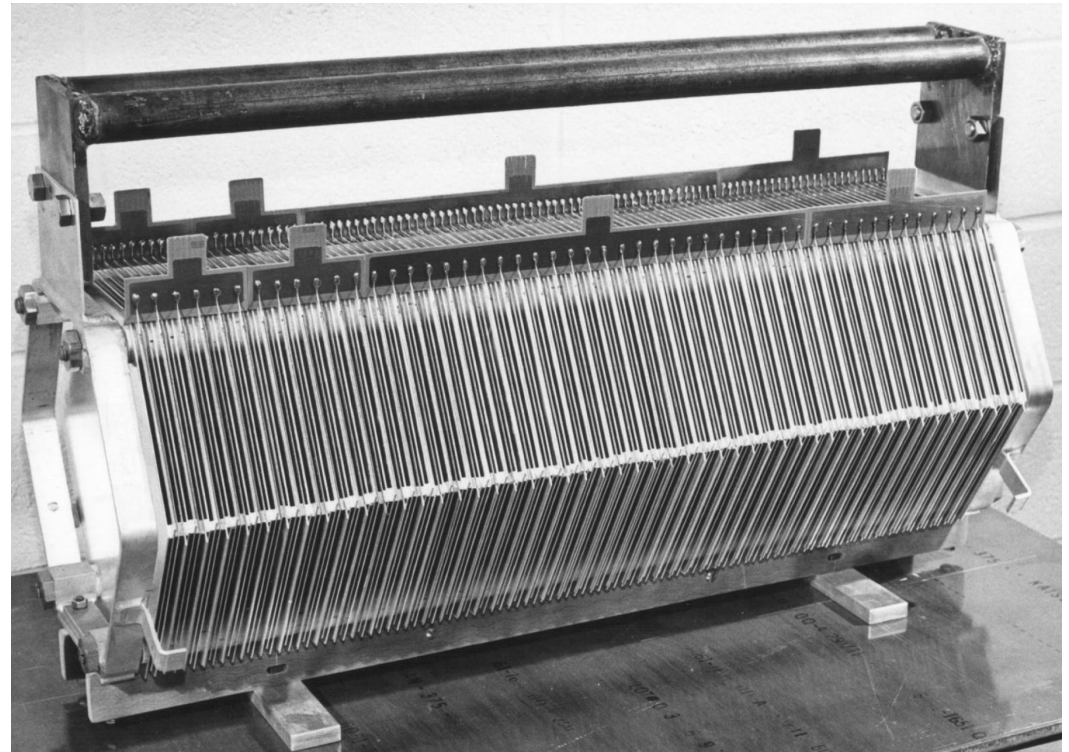
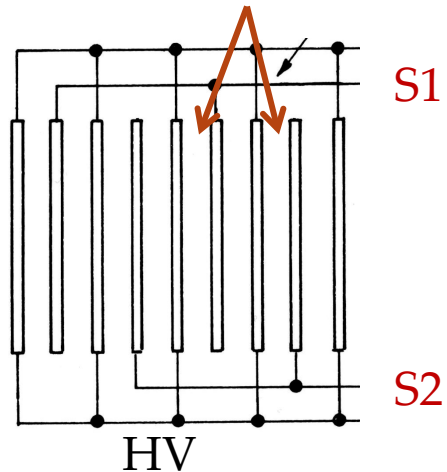


Measured charge vs electric field for 5.5 MeV α and MIPs

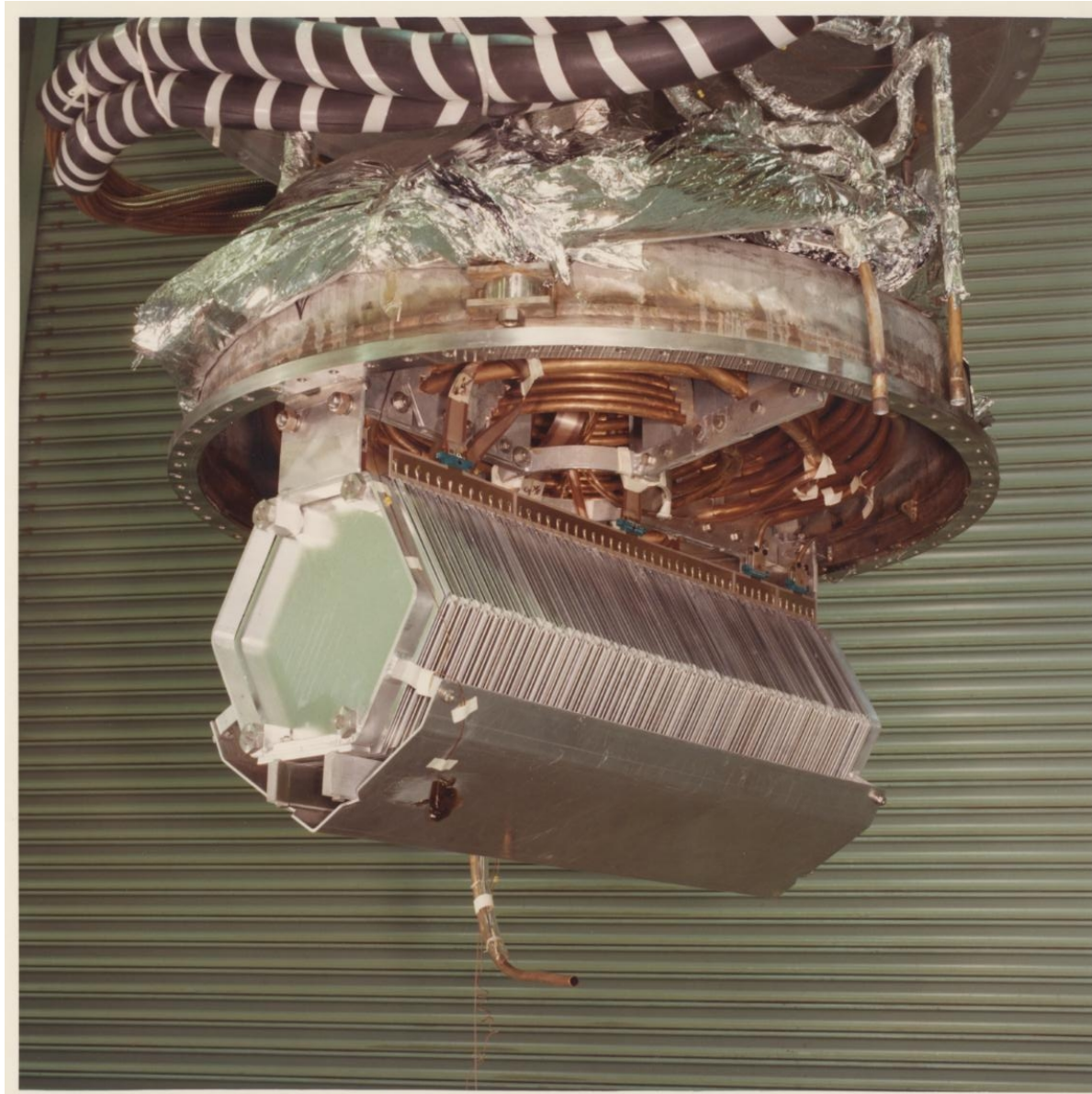


Jan-Apr 1973: Design and build the first LAr sampling electromagnetic calorimeter

- 20 radiation lengths long, 200 steel 1.6 mm plates with 2 mm LAr gaps; $\Delta E/\Delta x \approx (11.6 + 2.1)$ MeV/cm, or $0.1 X_0$
- Charge from ion chamber with ~ 100 nF electrode capacitance brought out on low impedance strip transmission lines
- Interleaved readout from alternate gaps for studies of sampling fluctuations



April 1973: Ready for cooling and beam tests at AGS



May-June 1973: The 1st LAr EM Cal. – Spectra and Difference Signals

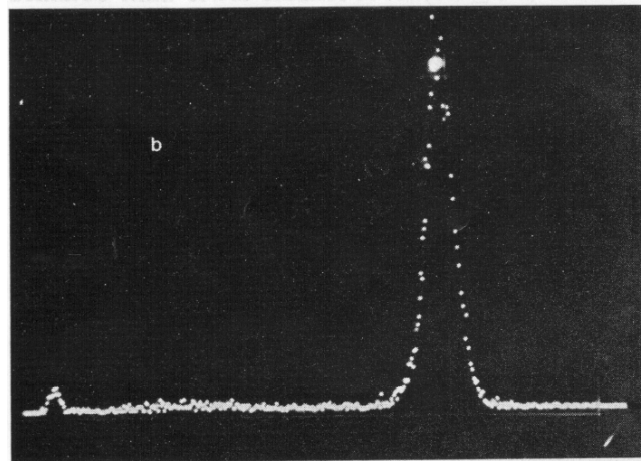
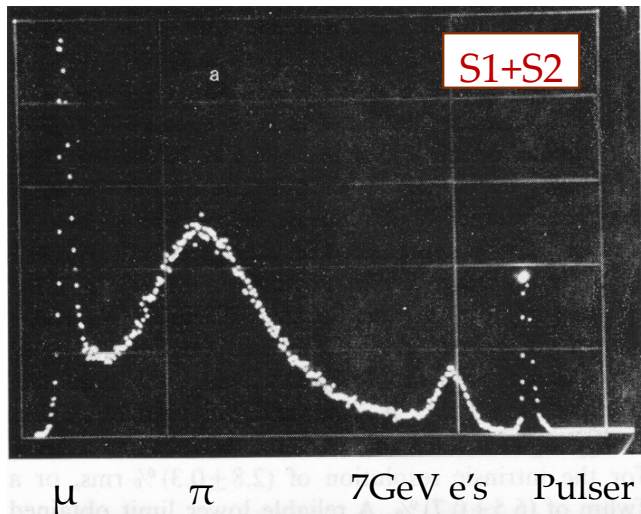


Fig. 12. Charge spectrum with large test chamber for 7 GeV/c negative beam. (a) peaks from left to right: muons, π 's, electrons, calibration pulser; (b) electrons enhanced with a Cherenkov detector.

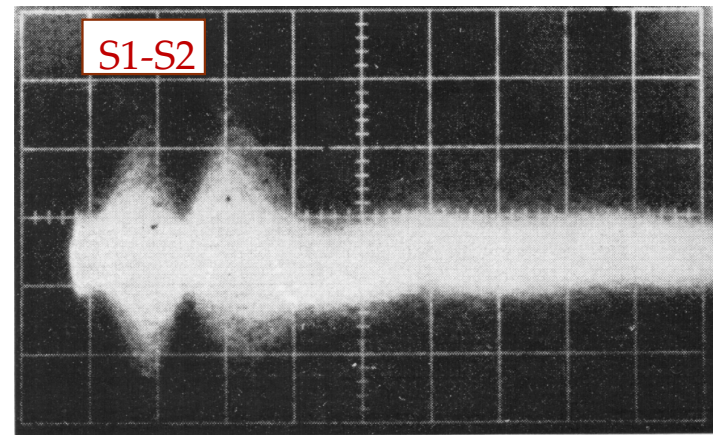


Fig. 15. Difference signals from two interleaved chambers for 7 GeV electrons (see text).

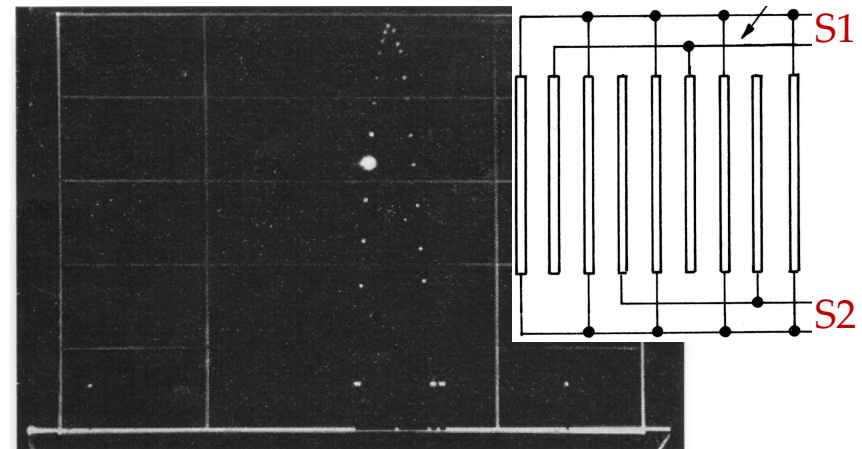


Fig. 16. Distribution of difference signals from two interleaved chambers with 7 GeV electrons. The three-decade logarithmic display shows a Gaussian distribution.

R806

EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH

M.H. BLEWETT

ISR * 16/11/73

ISRC/73/33

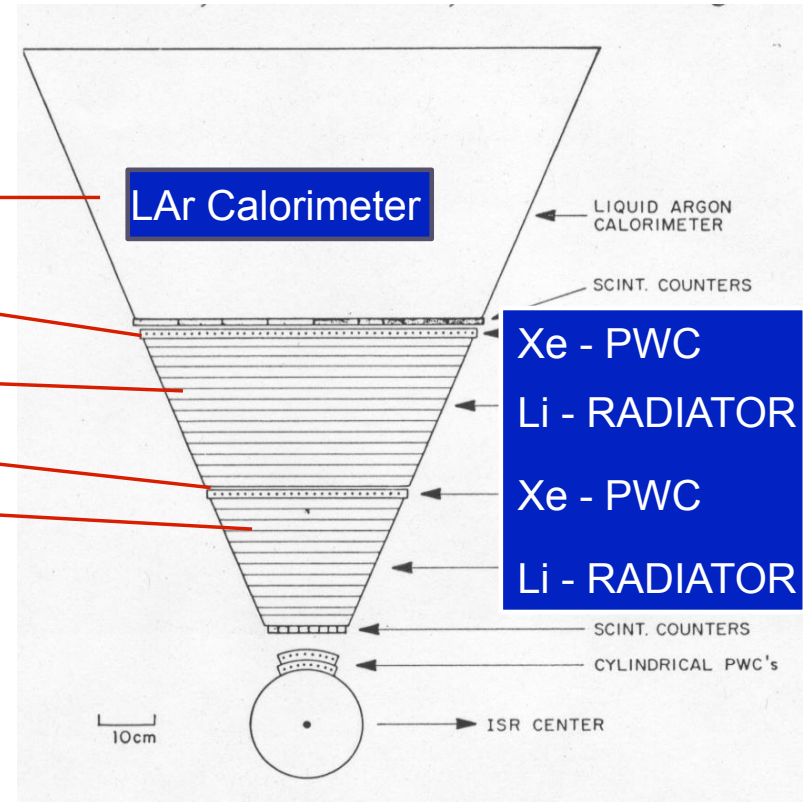
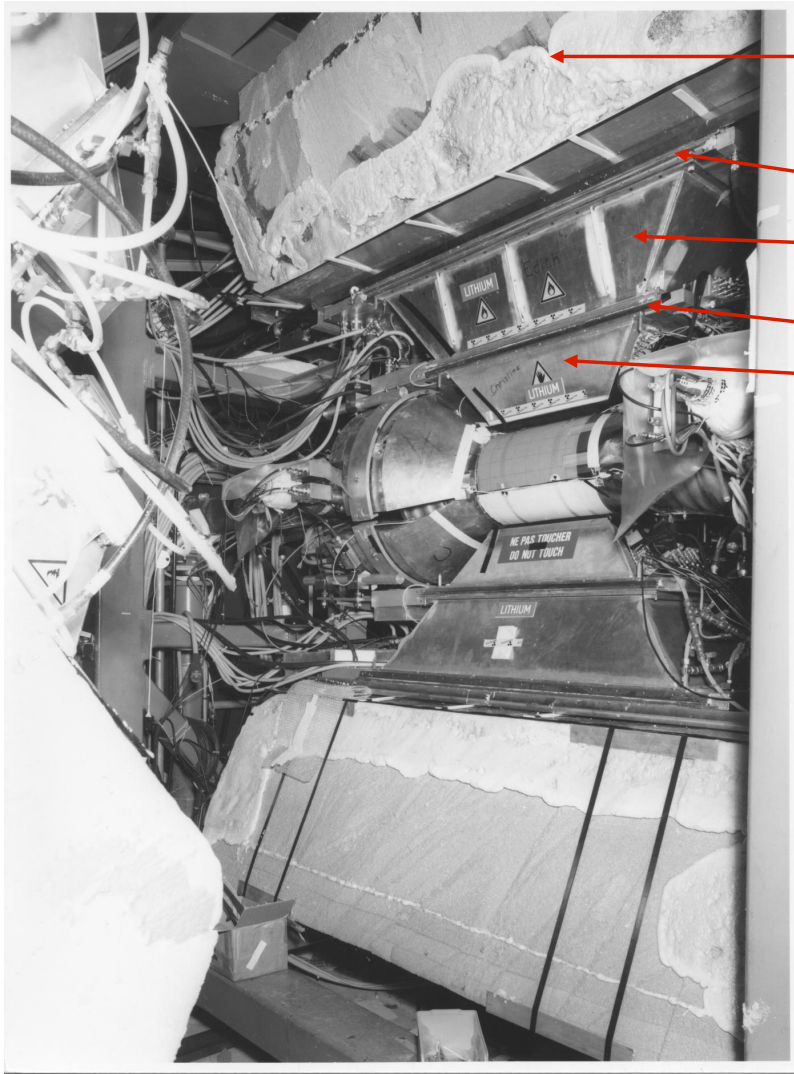
STUDY OF LARGE TRANSVERSE MOMENTUM PHENOMENA

Brookhaven, CERN, Saclay, Syracuse, Yale Collaboration

INTRODUCTION

Three of the most exciting questions in high-energy physics are: (i) the existence of massive particles¹⁾ (ii) the production of massive virtual photons²⁾ and (iii) the properties of "jets" at high transverse momentum³⁾ These questions are all related from an experimental point of view, as well as theoretically. For instance, we have in mind the detec-

R806



R806: work in progress – early days



Understanding the physics limitation of Hadron Calorimetry

- Distinguish instrumental effects, such as sampling fluctuations from intrinsic effects
- It's not the sampling which limits performance
- What matters is the response difference to electron and hadrons
- Uranium is the answer

Willis' notes on Uranium compensation

ENERGY RELEASE IN
 ^{238}U FISSION

	E_f MeV
FRAGMENTS	16.7
NEW TRONS	5.5
INITIAL γ 's	7.5
γ 's / β	8.4
β	8.9
$\bar{\nu}$	11.9
TOTAL	206.

} "VISIBLE"
13 MeV

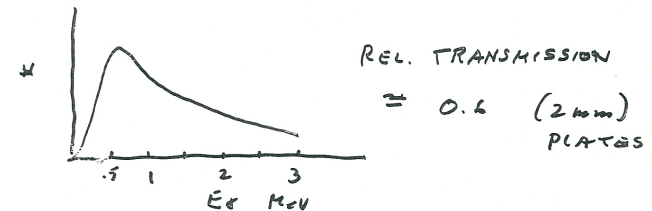
FROM ORNL-3940 LIVINGSTON ET AL
 1 n FOR 16 MeV, ^{238}U
 35 MeV Pb
 n energy mult. factor 1.87

5. COMPARISON:

Pb	U
$n \rightarrow \gamma \rightarrow 5 \text{ MeV}$	$n \rightarrow .71 \times 13 = 9.23 \text{ MeV}$
	$+ 3.3 \times 5 = 16.5$
	<hr style="width: 50px; margin-left: 0;"/>
	25.73 MeV

DIFFERENCE IN $\gamma = 20.7 \text{ MeV}$

6. BUT SPECTRUM IS SOFTER
 THAN SHOWER SPECTRUM:



7. DETECTED EXCESS γ IN ^{238}U / PRIMARY
 NEUT. = $0.6 \times 20.7 = 12.4 \text{ MeV}$ IN γ

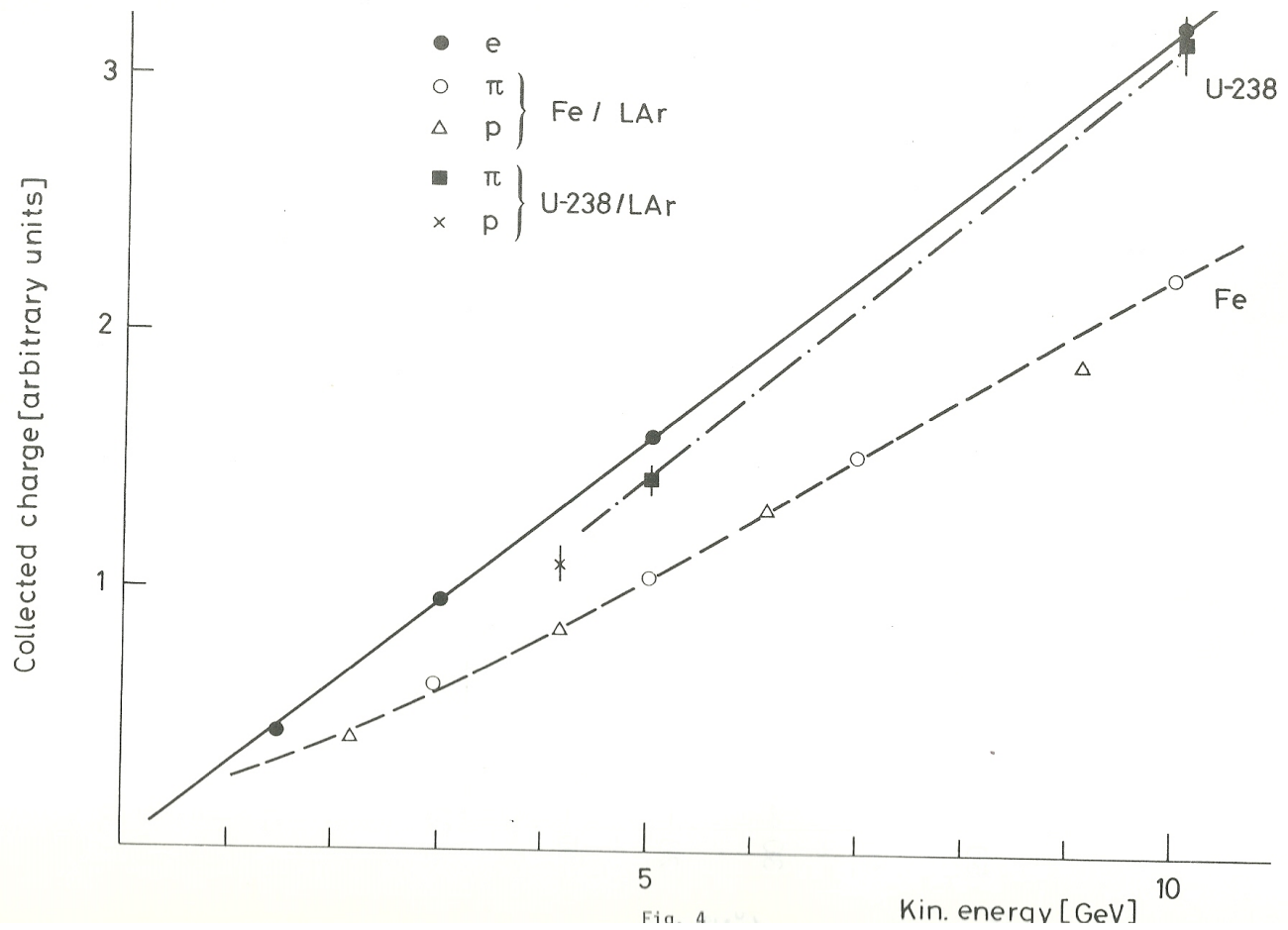
8. OR, PER GeV INCIDENT / GeV γ
 $\Rightarrow 29 \times 12.4 = 360 \text{ MeV}$

9. ADD TO $\sim 700 \text{ MeV}$ "HADRONIC VIS." (FE)
 $= 1.06 \text{ GeV}$ HURRAN!

PHYSICS LIMITATIONS ON CALORIMETRY

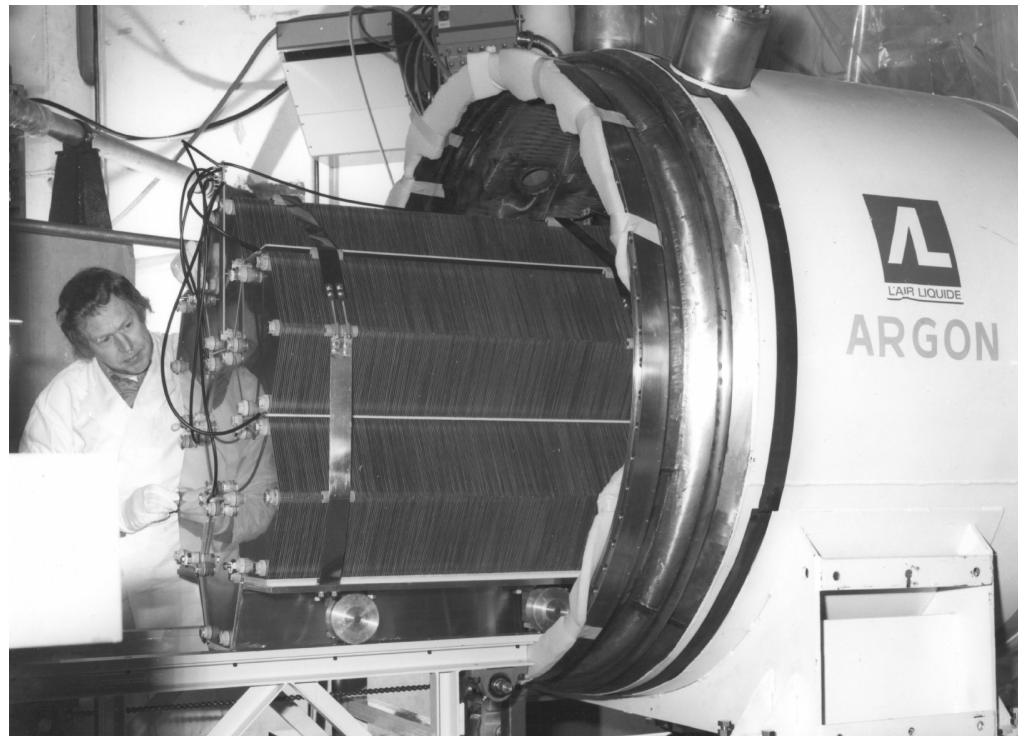
(Summary of a Contribution presented to the FNAL Workshop
on Calorimeters, May 9/10, 1975)

C.W. Fabjan and W.J. Willis
CERN, Geneva, Switzerland



1974-1976 LAr “Impactometer”

- Depleted Uranium as absorber in the center hexagon, surrounded by six hexagons with steel plates
- Studies of hadron showers
- Opening of the e/h response compensation issue
 - The contribution of fission to the hadron signal was not as high as expected
 - There was more to the compensation mechanism
- Uranium was later used in HELIOS, ZEUS and D0 calorimeters



NA34 – HELIOS

- 1983 – The new horizon was to convert the SPS into an Ion accelerator to study nuclear matter under extreme conditions

LETTER OF INTENT TO THE SPSC

SPS EXPERIMENTS WITH PRECISE FULL ANGLE ENERGY MEASUREMENTS

C.W. Fabjan - D. Lissauer - W.J. Willis

1. Introduction
2. Origin of Leptons, Neutrinos and Neutrino-like Objects
3. Deeply Inelastic Reactions of Hadrons with Nuclei
4. Interaction of Nuclear Beams with Heavy Nuclei
5. The New Spectrometer and its Performance.

NA34 – HELIOS

- HELIOS – High Energy Lepton and Ion Spectrometer
- Also an experiment in sociology: particle physicists meeting the nuclear physicists
- BNL-CERN-Heidelberg-Lund-McGill-Montreal-Lebedev-Novosibirsk-Pittsburgh-Saclay-Syracuse-Tel Aviv
- 1983-1988: at the CERN SPS, HELIOS single-lepton detector with a uranium-LAr em+hadron calorimeter in p-p reactions and 4π showers in *heavy ion* experiment
 - Radeka introduced cryogenic electronics!

PROPOSAL

STUDY OF HIGH ENERGY DENSITIES OVER EXTENDED NUCLEAR VOLUMES VIA NUCLEUS-NUCLEUS COLLISIONS AT THE SPS

NA34 Collaboration

H. Gordon, T. Ludlam, V. Polychronakos, D.C. Rahn, I. Stumer, C. Woody
Brookhaven National Laboratory

T. Åkesson, H. Atherton, H. Breuer, C.W. Fabjan, U. Goerlach
S. Katsanevas, U. Mjornmark, J. Schukraft, W.J. Willis
CERN

P. Glässel, A. Pfaiffer, J. Soltani, H.J. Specht ^{*)}
Heidelberg University

N.J. DiGiacomo, P.L. McGaughey, W.E. Sondheim, J.W. Sunier
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V. Sidorov
Novosibirsk Institute of Nuclear Physics

W. Cleland, J. Thompson
Pittsburgh University

A. Gaidot, F. Gibrat, C.W. London, J.P. Pansart
CEN DPhPE Saclay

D. Bettoni, M. Goldberg, N. Horwitz, G.C. Moneti, L. Olsen
Syracuse University

G. Benary, S. Dagan, D. Lissauer, Y. Oren
Tel Aviv University

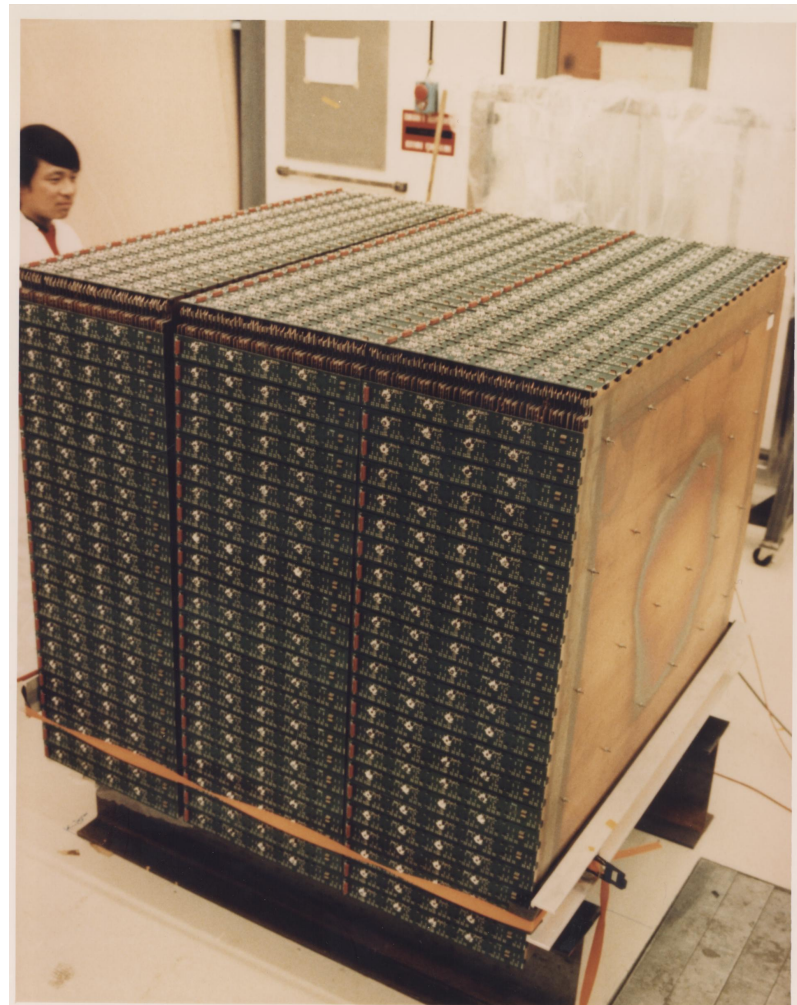
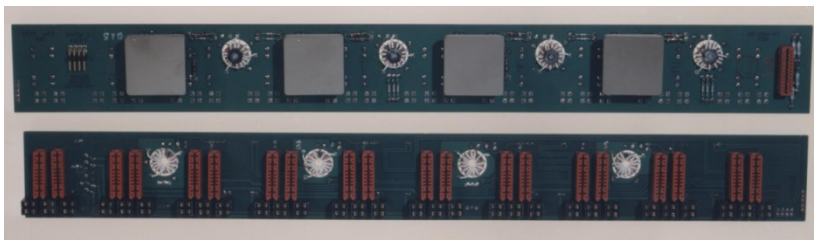
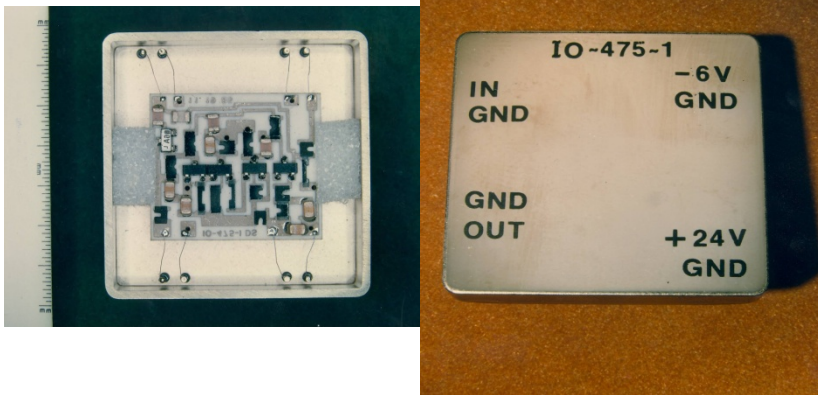
ABSTRACT

We propose to examine in detail the characteristics of ultra-relativistic nucleus-nucleus interactions using ^{16}O beams of 200 GeV/A from the SPS. The experiment combines 4π calorimeter coverage with measurements of inclusive particle spectra, two-particle correlations, low- and high-mass lepton pairs and photons. A multi-wire active target allows maximum interaction rates with a minimum of secondary interactions.

^{*)} Spokesman

HELIOS Uranium-LAr: Hadronic Module

- 576 x and y strips with interleaved readout and preamplifiers operated in LAr

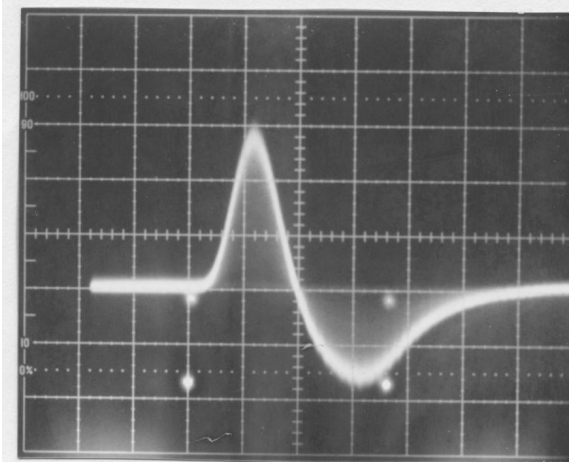


HELIOS Uranium-LAr EM+Hadron Calorimeter ready for cooling

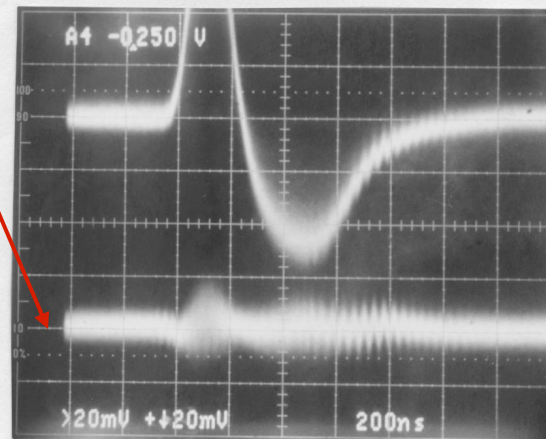


HELIOS Uranium-LAr: Sum response to 200 GeV protons (576 signal channels)

- Difference signal from alternate LAr gaps \rightarrow **sampling fluctuations**
- Energy resolution (single particles): $\sim 21\%/\sqrt{E}$



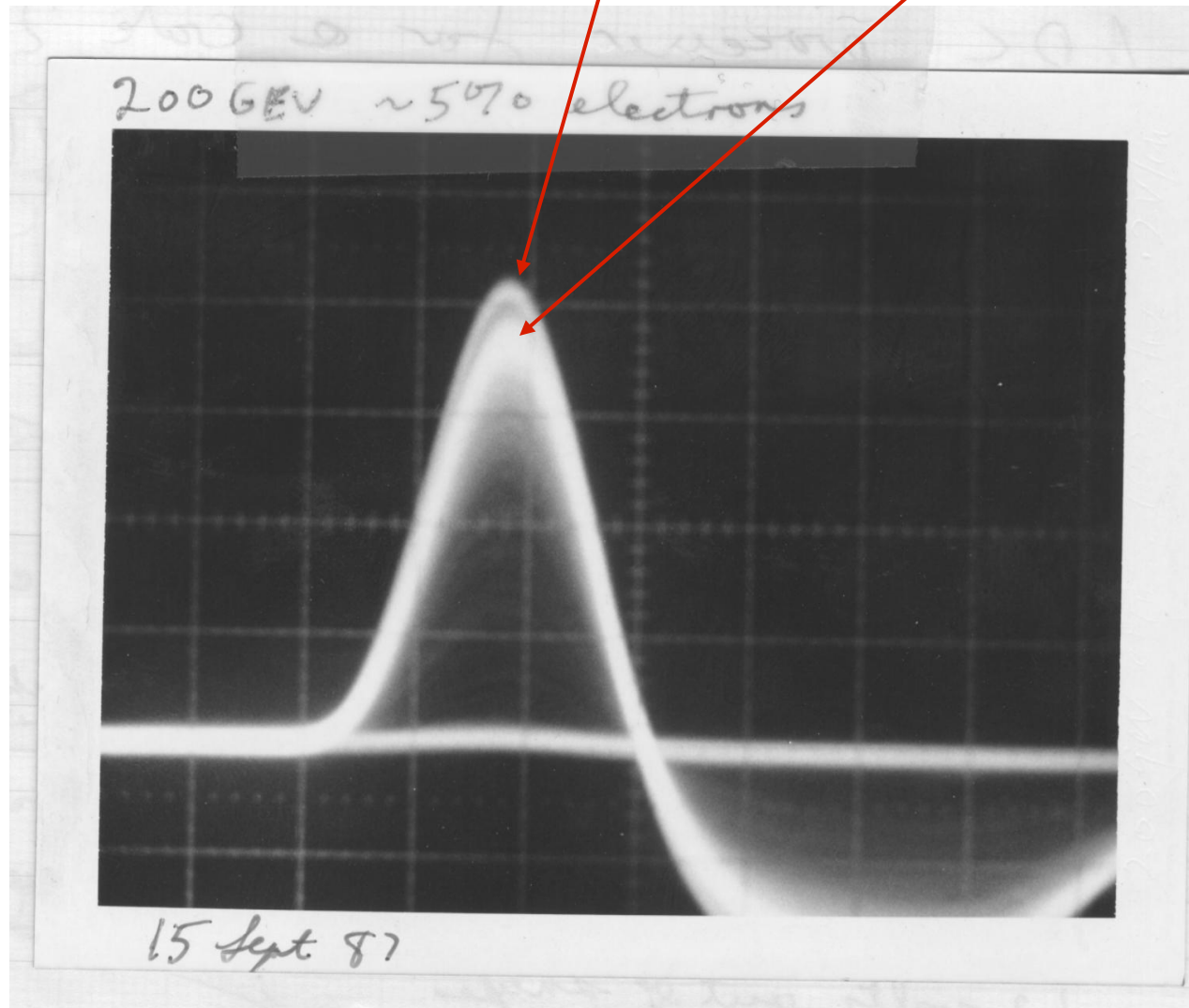
200 GeV hadrons
 $E_x + E_{x'} + E_y + E_{y'}$
(576 channels, 1.8 m^3)
Rate: $\sim 4 \times 10^5 \text{ sec}^{-1}$



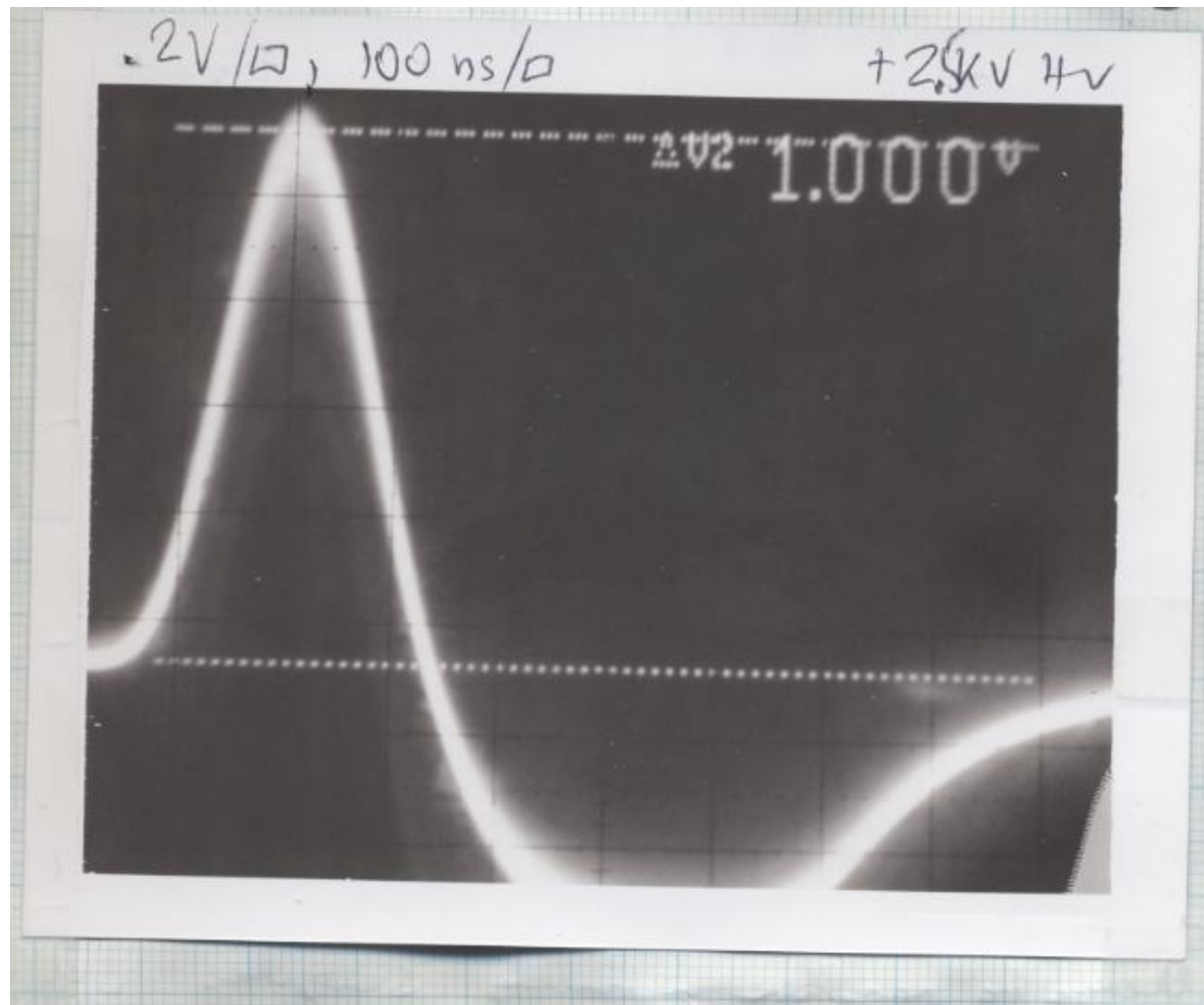
$\leftarrow E_y$ (5 divisions peak amplitude)

$\leftarrow E_y - E_{y'}$

HELIOS Uranium-LAr: Electrons and Pions ($e/\pi \approx 1.15$)

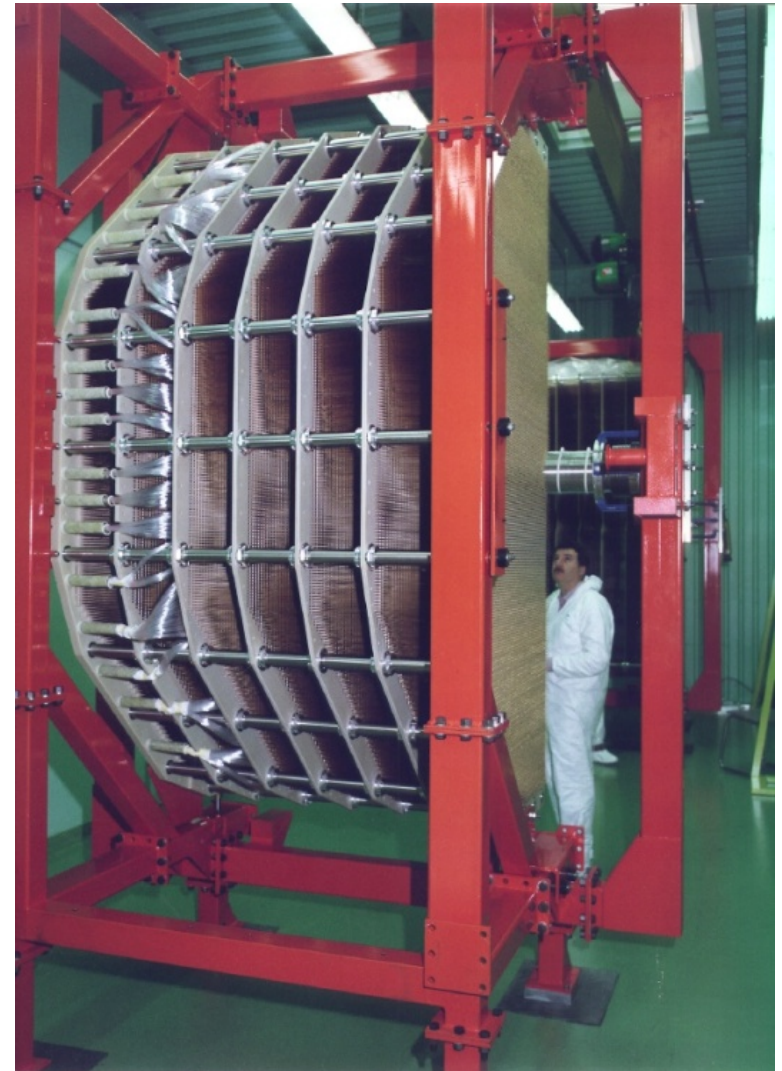
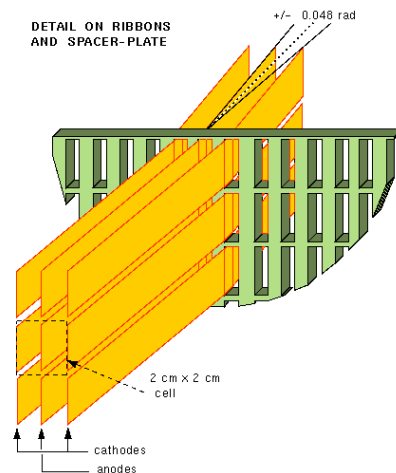


HELIOS Uranium-LAr: Sulphur ions 200 GeV/nucleon 6.4 TeV total



NA48

- NA48 is an experiment launched in the early 1990s to study direct CP violation in the neutral kaon decay system at the North Area of the SPS at CERN
 - Homogeneous calorimeter
 - NA62 is the successor and still running
- Liquid Krypton calorimeter with silicon JFET preamplifiers operated in LKr, total 13,212 channels
 - Continuation of cryogenic electronics



Accordion

- 1990 – RD3 proposal
- R&D proposal – Liquid Argon Calorimetry with LHC-performance specifications
- LAr technology for collider
- Pros
 - Reasonable cost
 - Good energy resolution
 - Good stability
 - Accurate inter-calibration
 - Intrinsically radiation resistant
- Cons
 - Relatively slow: 400 ns for 2 mm gap
 - Long connections to electronics → large C → large noise
 - Large dead space due to readout electronics

CERN LIBRARIES, GENEVA
CERN/DRDC/90-31
DRDC/P5
13th August, 1990

CERN
BIBLIOTHEQUE
SCP
CERN-DRDC
90-31

R&D PROPOSAL

LIQUID ARGON CALORIMETRY WITH
LHC-PERFORMANCE SPECIFICATIONS

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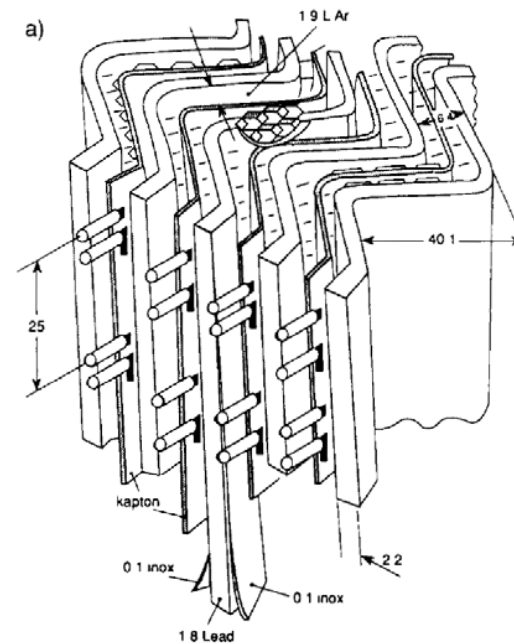
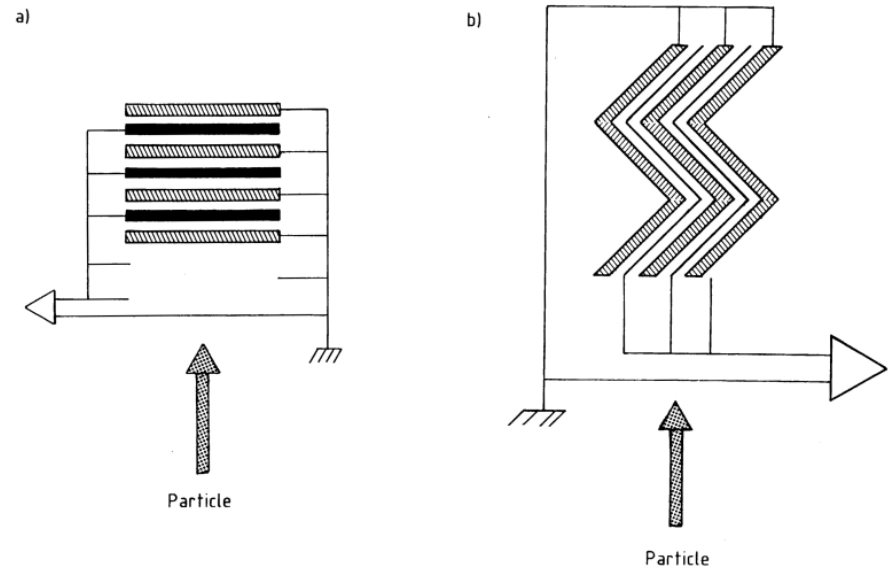
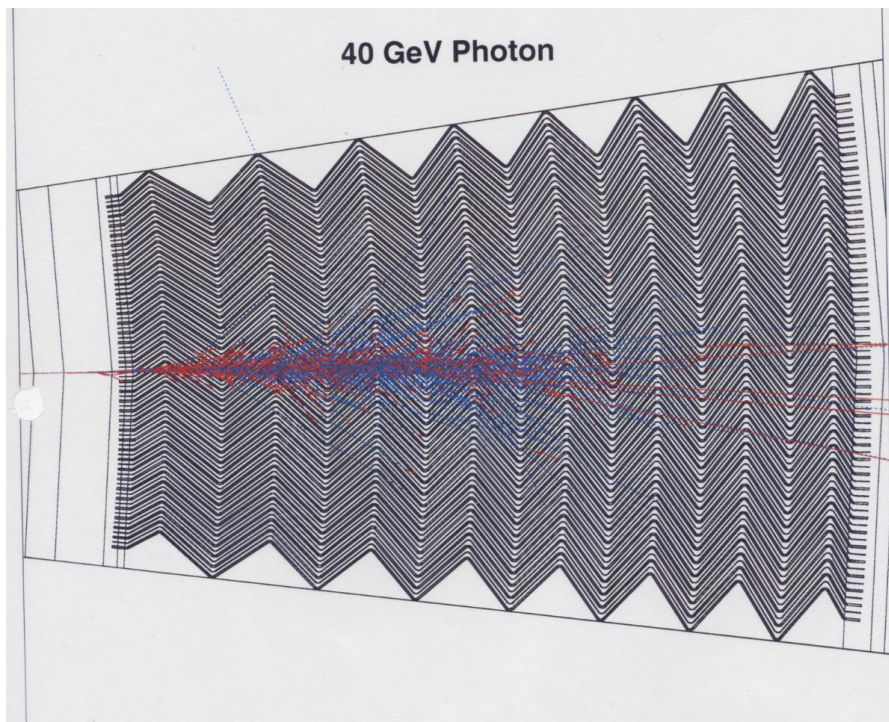
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2004

Accordion

- 1990/RD3 – Daniel Fournier introduced accordion geometry
- From RD3 to GEM/SSC and ATLAS/LHC: LKr and LAr accordion EM Calorimetry



GEM – Gammas, Electrons and Muons

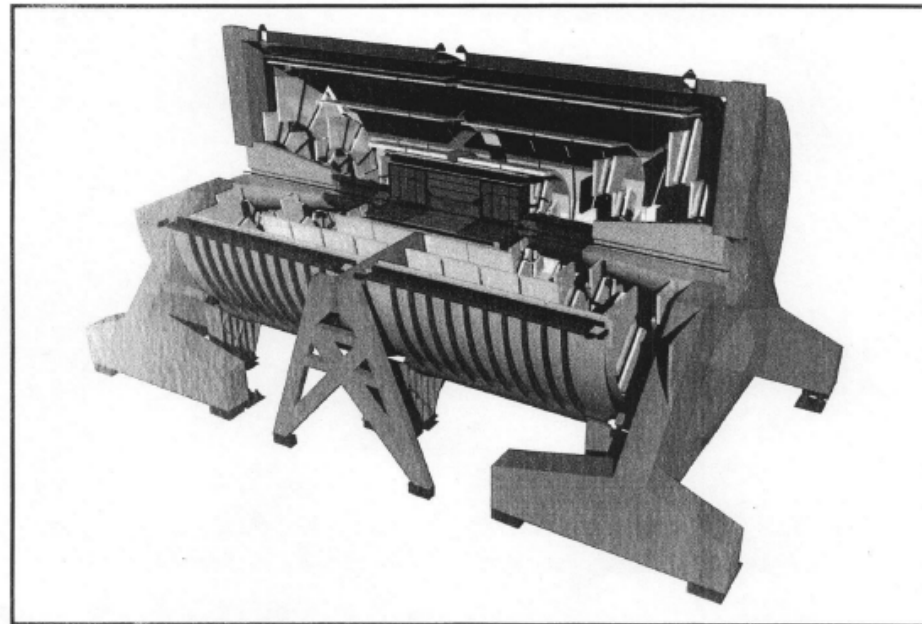


GEM



GEM-TN-93-262
SSCL-SR-1219

Technical Design Report



April 30, 1993

Consider Liquid Krypton for GEM detector

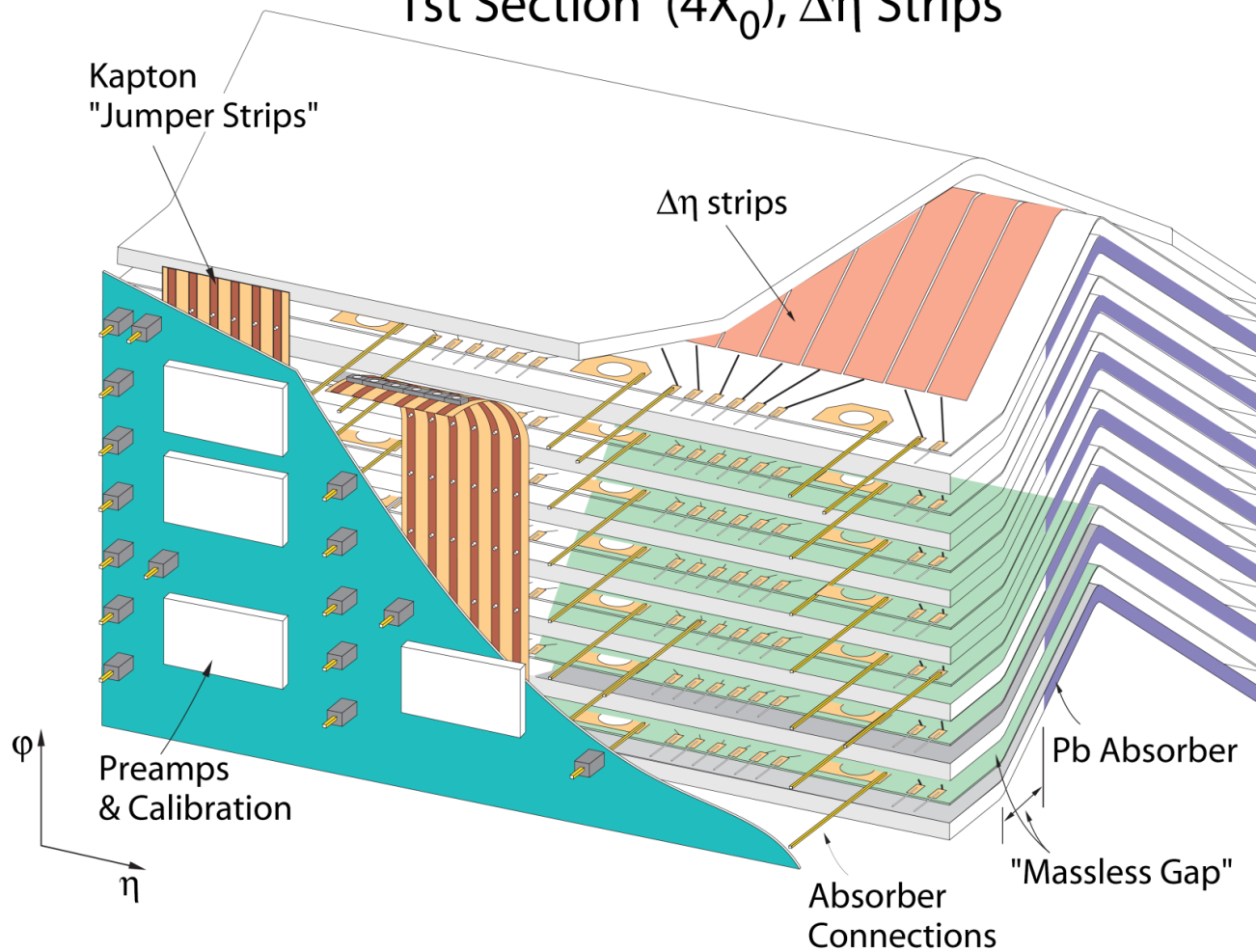
	LAr	LKr	LXe
X_o [cm]	14.0	<u>4.72</u>	2.77
R_M [cm]	10.1	4.8	4.1
W_i [eV / pair] at 10 kV/cm	25.1	<u>19.1</u>	
dE/dx [MeV / cm]	2.11	<u>3.45</u>	
v_d [cm / μ s]	- " -	4.4	3.83
l for $25X_o$ [m]		3.5	1.2

Peak induced current in a sampling calorimeter:

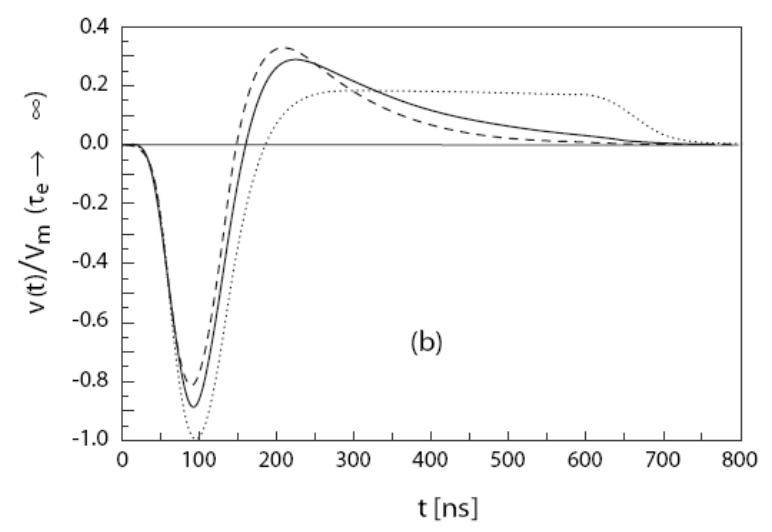
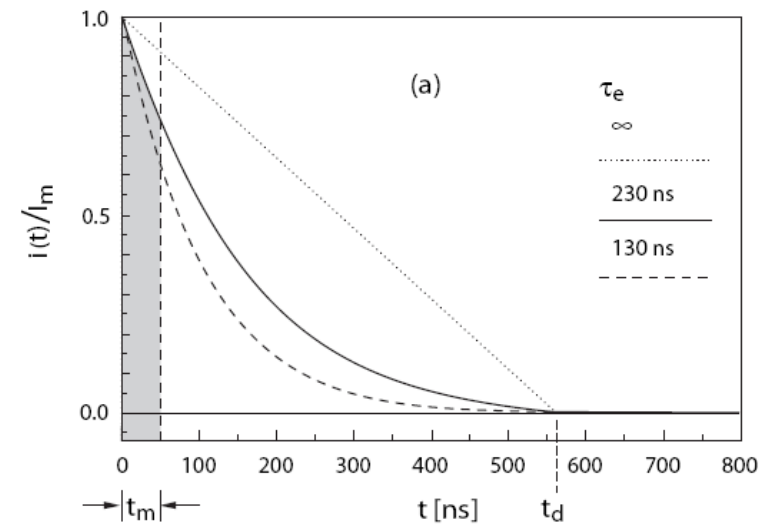
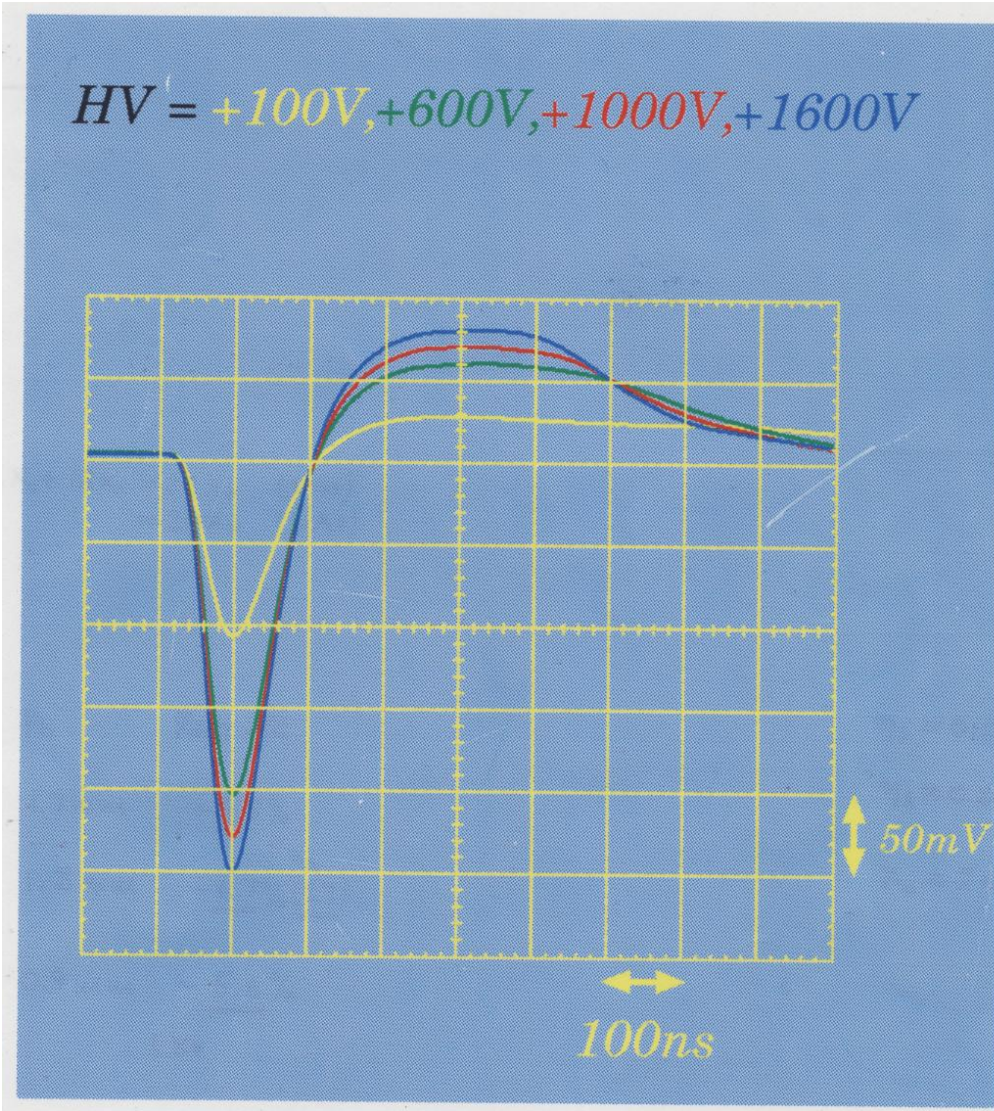
$$I_m / E \propto \frac{(dE/dx)}{(dE/dx)_{abs} + (dE/dx)} \cdot \frac{v_d}{W_i} \cdot ("e/\mu")$$

Tests of a new design with LKr and LAr at AGS in 1992 and at CERN SPS in 1994

EM Accordion Electrodes 1st Section ($4X_0$), $\Delta\eta$ Strips



LKr Accordion EM tests: Waveforms (shower sums)



LAr and LKr Accordion EM: Results

- Energy Resolution:**

		Pb absorber [mm]	$\frac{\sigma_E}{E} \sqrt{E}$ [%]
CERN RD3	LAr	1.8	9.6
BNL 1992	LAr	1.3	7.7
"	LKr	1.3	6.7
"	LAr	2 x 0.8	6.3

CERN 1994 (by BNL et al.)	LKr	1mm	$\frac{5.6\%}{\sqrt{E}} \oplus 0.2$
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- Timing Resolution:**

(single tower)

$$\frac{1.16 \text{ ns}}{E} \oplus 73 \text{ ps}$$

- Pointing Resolution:**

$$\frac{32.2}{\sqrt{E}} \text{ mrad}$$

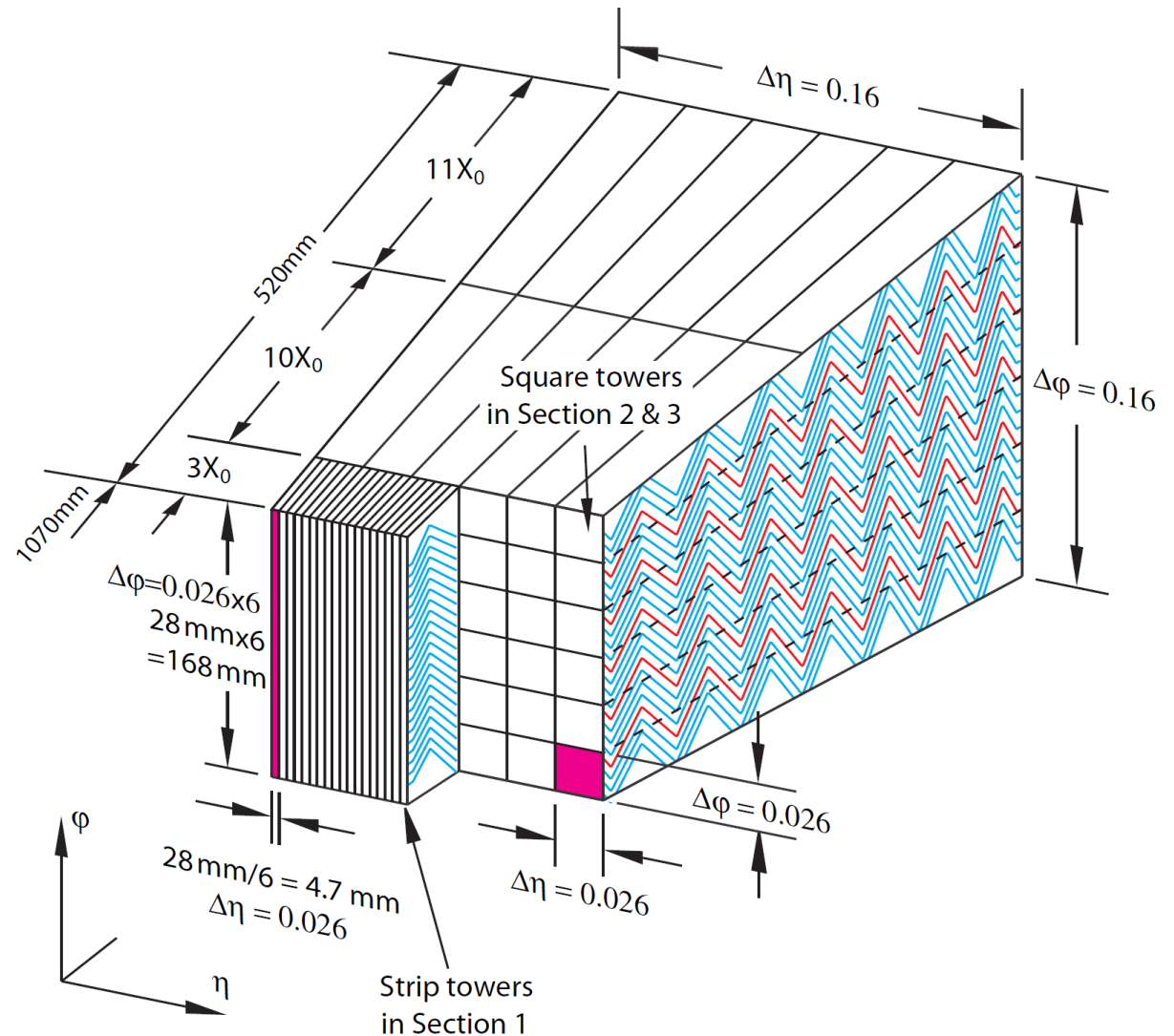
- Peak Current LKr / LAr:**

(for equal length cal.)

$$\approx 2$$

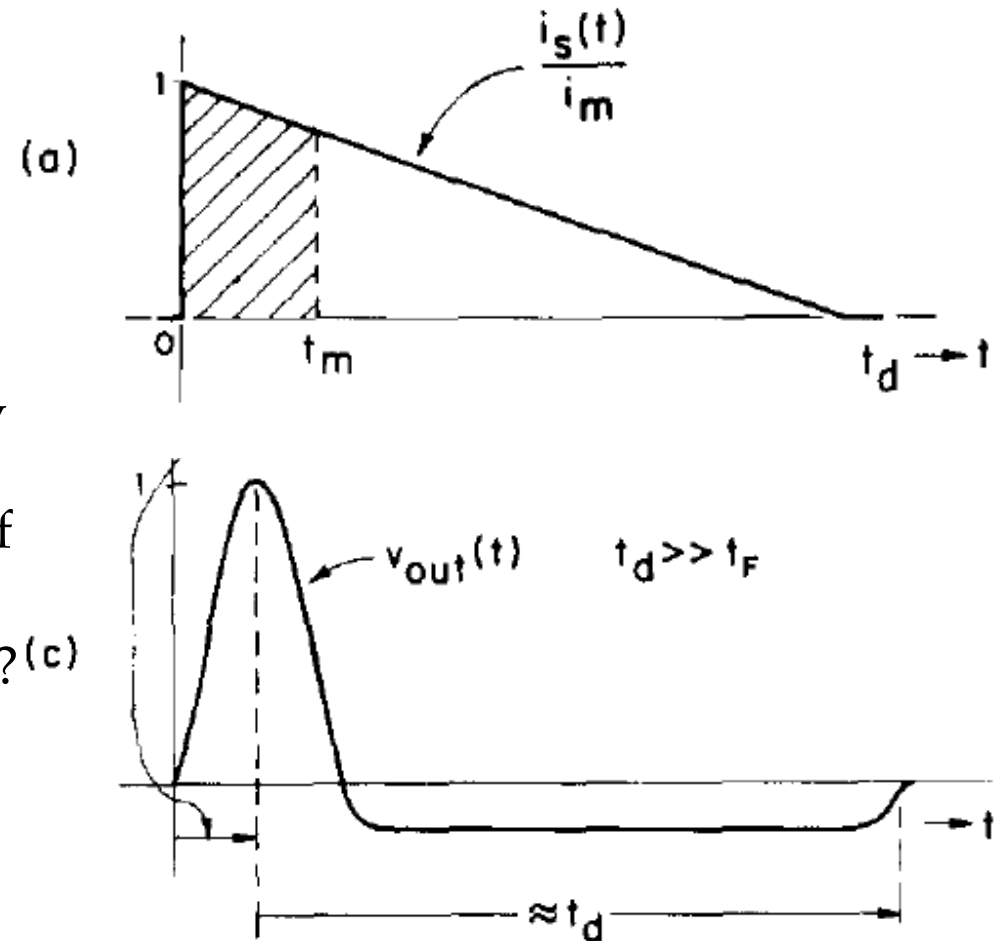
EM electrode design with fine granularity intended for GEM and later adapted for ATLAS

- Used in ATLAS with a different longitudinal subdivision



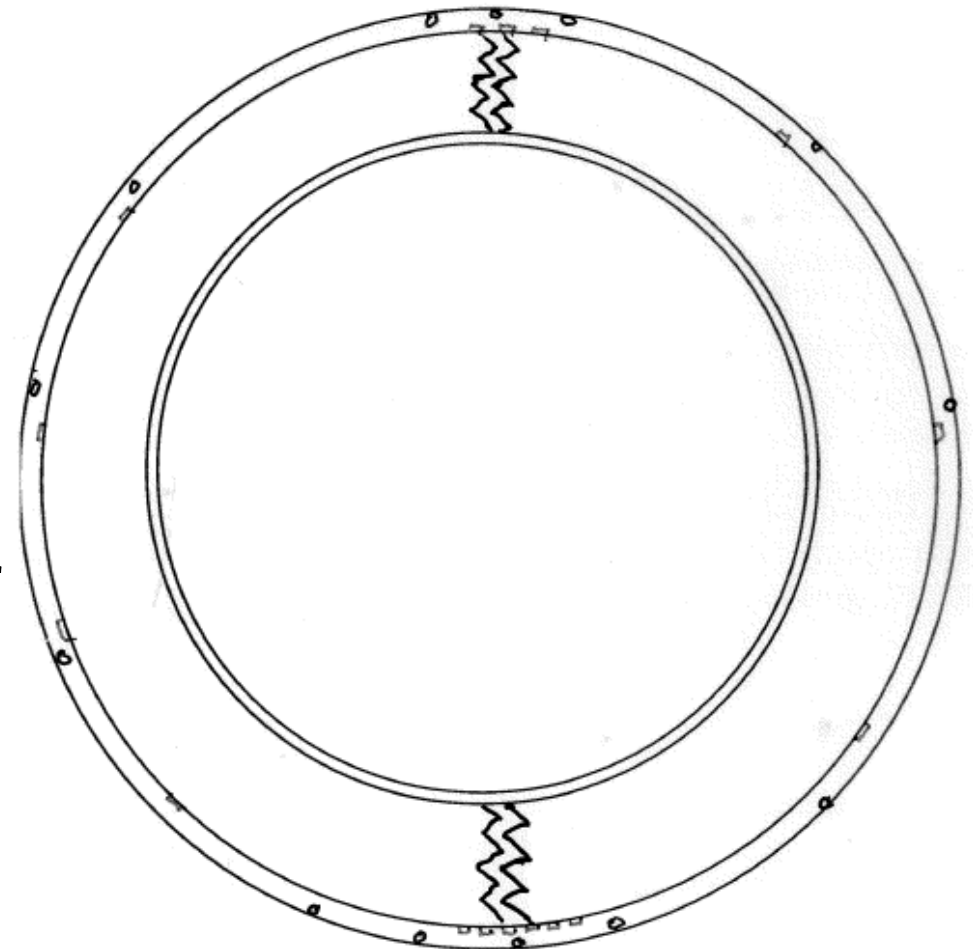
ATLAS LAr Calorimeter

- Two remaining problems
- Problem #1: LAr maybe too slow for LHC?
 - 400 ns (2 mm gap) \leftrightarrow collision rate = 25 ns @ LHC
 - Induced current rises quickly \rightarrow take only the front part using bipolar shaping time of $\tau < 50$ ns
 - How to handle negative tails? ^(c) \rightarrow pileup may compensate. This is an advantage to restore baseline in high rate operation with high pileup



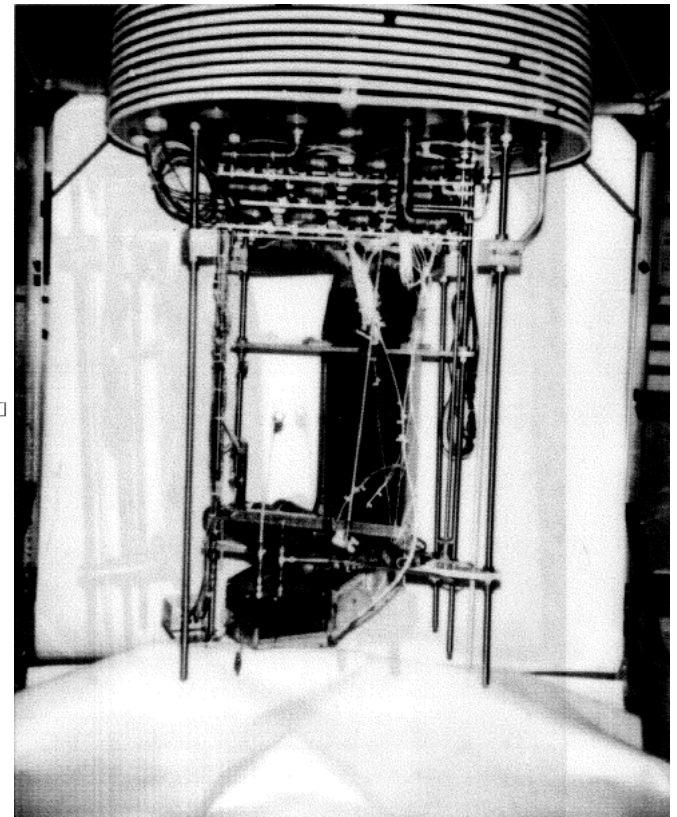
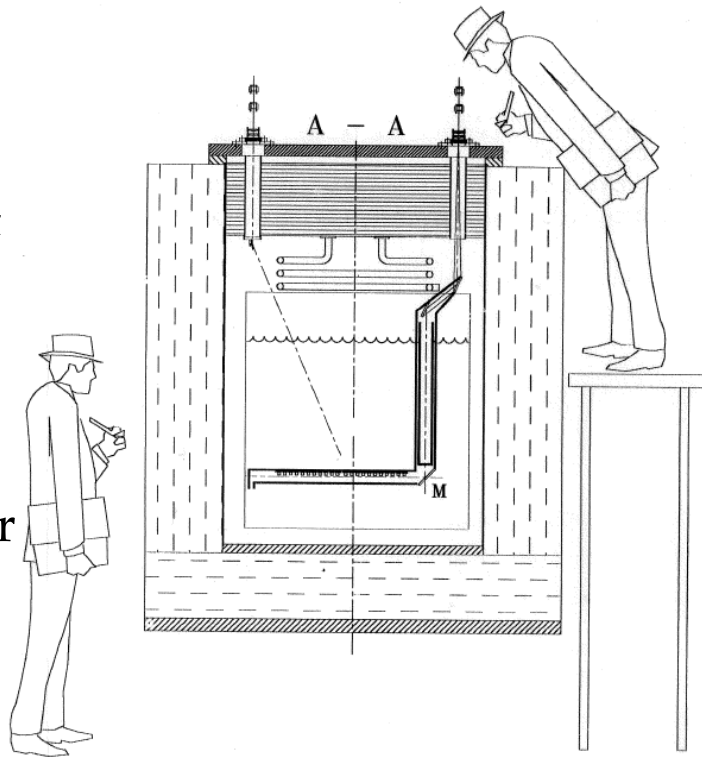
ATLAS LAr Calorimeter

- Problem #2: Preamplifiers be in the calorimeter?
 - With preamplifiers directly mounted, no additional cables are needed and the most favorable configuration for high speed, low noise and small cross-talk is reached
 - The ATLAS barrel calorimeter will have a problem which is not permitting convective cooling of the preamps



Cooling Tests at BNL

- Cooling tubes located at the top of the preamps were unable to stop the boiling
- Signal variation with respect to temperature is at least 2% per degree
- The temperature variations over different regions of the calorimeter would lead to large corresponding variations in the signal outputs



ATLAS LAr Calorimeter

- Problem #2: Preamplifiers be in the calorimeter?
 - It was found that cables do not contribute to noise for fast shaping with accurately matched line-termination
 - ATLAS decided to place all preamps outside the cryostats for EM calorimeter, because of maintenance, pile-up effects ...

Nuclear Instruments and Methods in Physics Research A330 (1993) 228-242
North-Holland

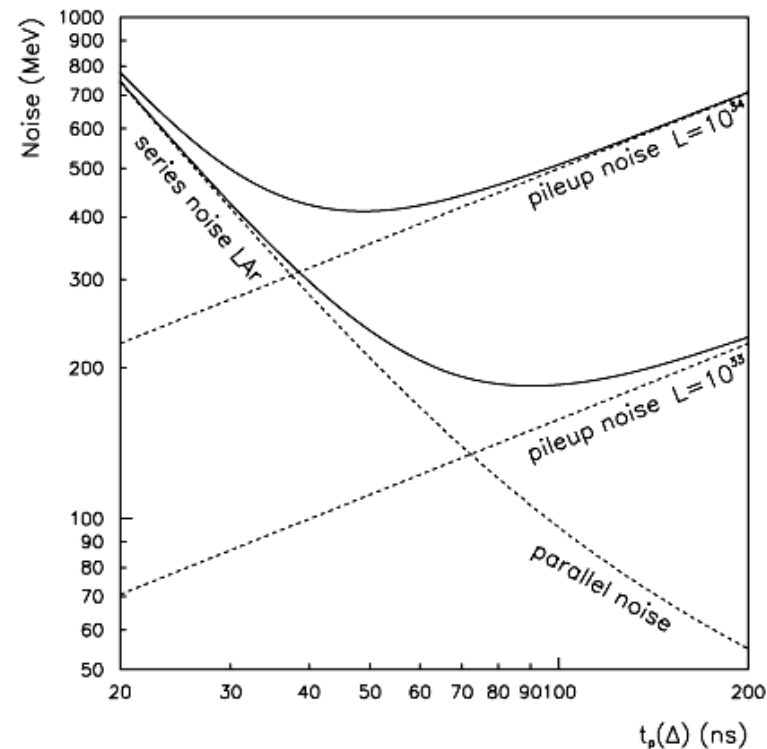
Transmission line connections between detector and front end electronics in liquid argon calorimetry *

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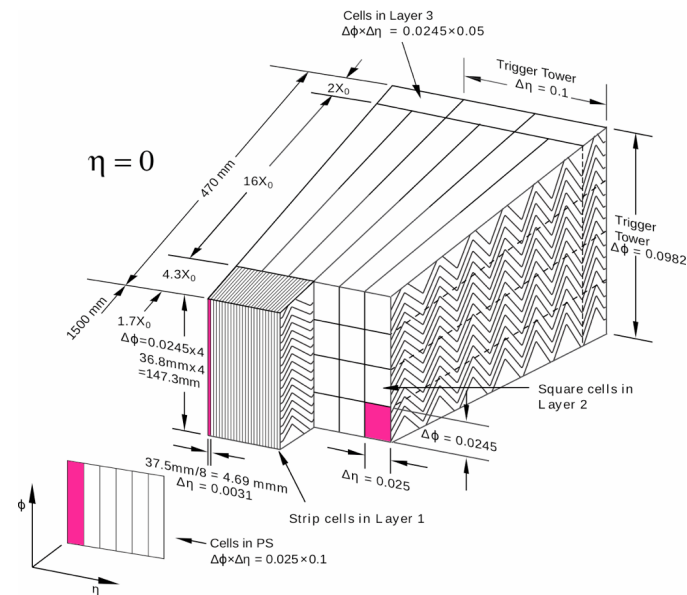
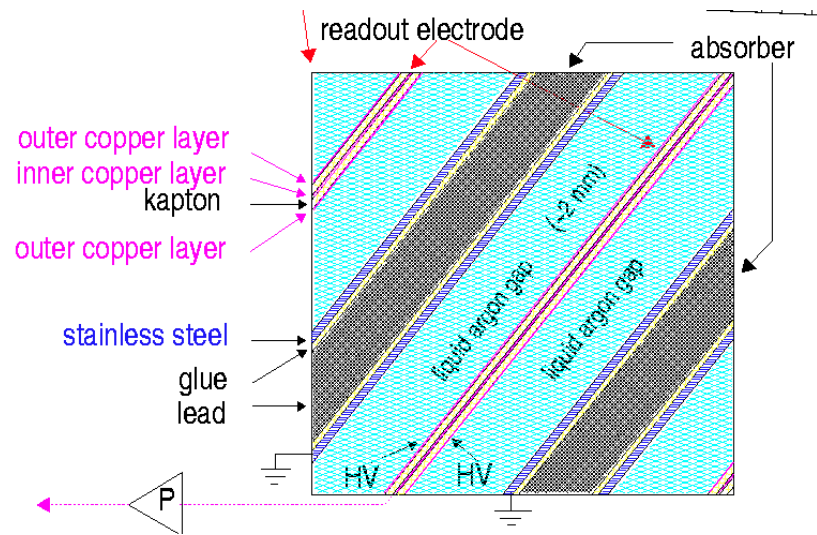
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Received 9 August 1992 and in revised form 8 January 1993



Design of ATLAS LAr EM Calorimeter

- Acceptance: $|\eta| < 2.4$ (full),
2.4~3.2 (coarse)
- Fine segmentations
 - Depth: 1+3 layers
 - Transverse: 0.025x0.1 (presampler)
 - 0.003125x0.1 (front)
 - 0.025x0.025 (middle)
 - 0.025 x 0.05 (back)
- Thickness
 - 26 X_0 for barrel
 - 28 X_0 for endcap

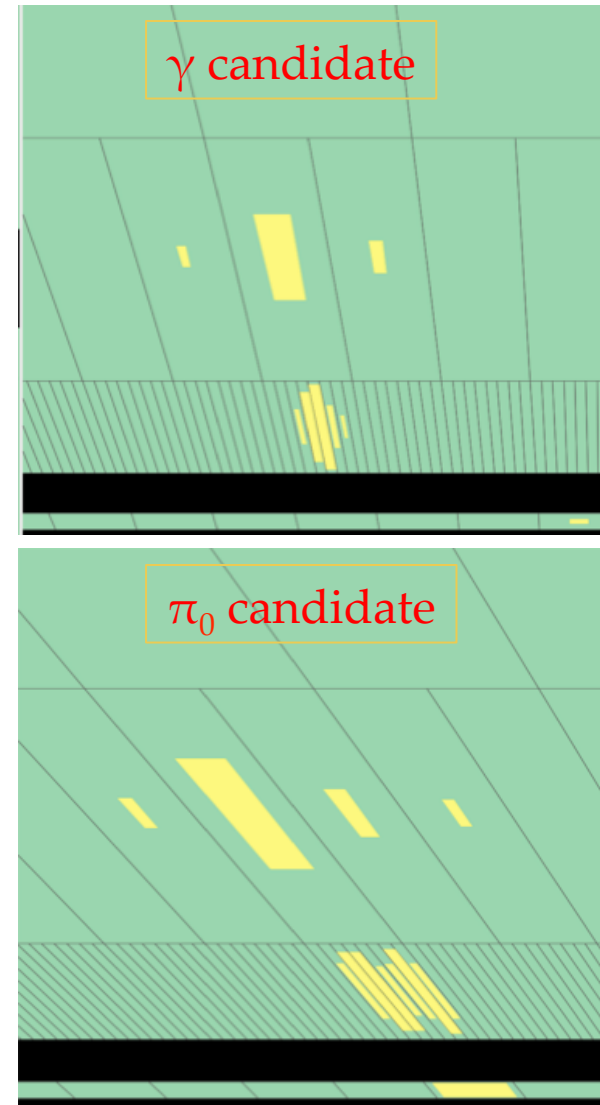


Design of ATLAS LAr EM Calorimeter

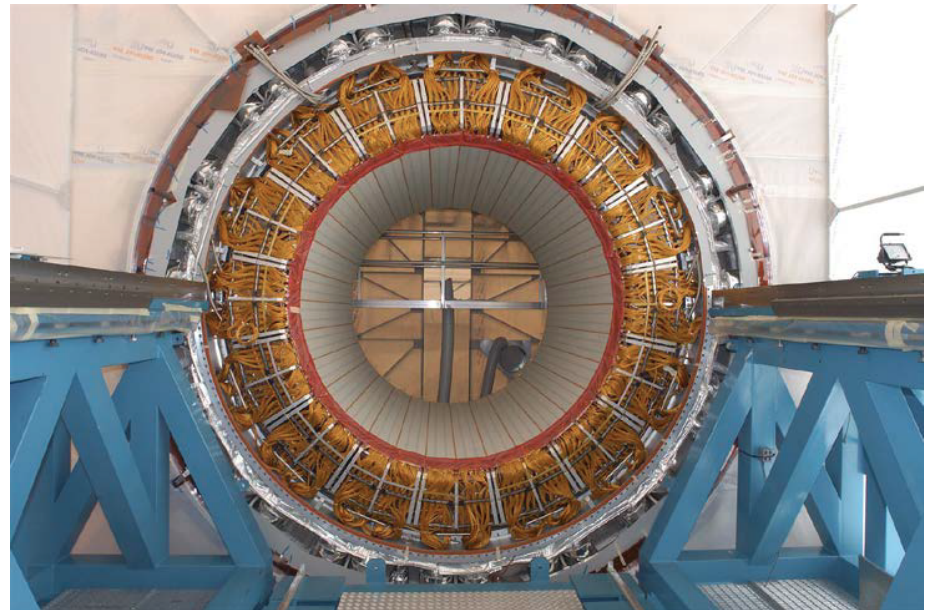
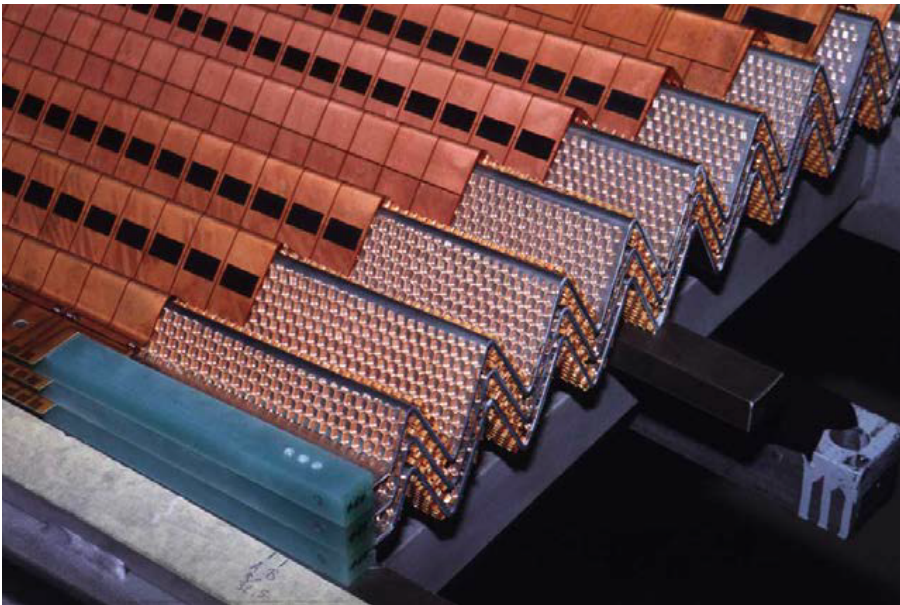
- Dynamic range
 - 35 MeV ~ 3 TeV → readout electronics 16-bit accuracy
- Total readout channels: 182,468
- Energy resolution (E in [GeV])

$$\frac{\sigma_E}{E} = \frac{10\%}{\sqrt{E}} \oplus \frac{0.4}{E} \oplus 0.7\%$$

- Angle resolution: 40 mrad/ \sqrt{E}



ATLAS LAr Calorimeter – From Construction



ATLAS LAr Calorimeter – To Higgs Discovery

