An Updated Analysis on Lepton Flavor Violating Radiative Decays in Non-sterile Electroweakscale Right-handed Neutrino Model

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- · 2. The Model
 - 3. The calculation
 - Process $I_i \rightarrow I_j + \gamma$
 - Magnetic Dipole Moment and Electric Dipole Moment
 - 4. Analysis and discussion
 - 5. Summary

• The Standard Model can explain most of the experimental results. However, there are many undetermined parameters and issues \rightarrow BSM

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- Neutrino oscillations = Neutral lepton flavour violation.
 - → Why not have charged lepton flavour violation (cLFV) ?



- If observed:
 - Probe the origin of lepton mixing
 - Probe the origin of new physics

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$$B(\mu \to e\gamma) = \frac{3\alpha}{32\pi} \Big| \sum_{l} (V_{MNS})^*_{\mu_l} (V_{MNS})_{el} \frac{m_{\nu_l}^2}{M_W^2} \Big|^2 ~~ \sim O(10^{-54})$$

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Could there be other mechanisms for LFV?



Dominant diagram for $l_i \rightarrow l_j + \gamma$

Such processes have been discussed in a generic fashion in **EW vR Model** (*arXiv:0711.0733 [hep-ph], P. L. B.* 659, 585 (2008))

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The Model

Non-sterile Electroweak-scale Right-handed Neutrino Model (The EW VR Model) [P.Q Hung 2007]

The Model

The EW vR Model [PQ Hung 2007]

What is it?

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Model in which right-handed neutrinos have Majorana masses of the order of Λ_{EW} naturally.



Gauge group



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 $SU(3)_{C} \times SU(2) \times U(1)_{Y}$



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Model Content



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Leptons

 $I_L = \left(\begin{array}{c} \nu_L \\ e_L \end{array}\right)$

e_R



Gauge group

 $SU(3)_{C} \times SU(2) \times U(1)_{Y}$

Model Content

Leptons

 $I_L = \begin{pmatrix} \nu_L \\ e_I \end{pmatrix} \longleftrightarrow I_R^M = \begin{pmatrix} \nu_R \\ e_P^M \end{pmatrix},$ e_I^M e_R



Quarks

 $\begin{pmatrix} u_L \\ d_L \end{pmatrix}$ $u_R, \ d_R$

 $q_L =$



Quarks



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$$q_{L} = \begin{pmatrix} u_{L} \\ d_{L} \end{pmatrix} \longleftrightarrow q_{R}^{M} = \begin{pmatrix} u_{R}^{M} \\ d_{R}^{M} \end{pmatrix},$$
$$u_{R}, d_{R} \longleftrightarrow u_{L}^{M}, d_{L}^{M}$$

Mirror particles are totally different from the SM particles!



Higgs Sector:

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$$L_{Y_l} = g_l \,\overline{l}_L \,\Phi \,e_R + h.c.$$
$$L_{Y_q} = g_q \,\overline{q}_L \,\Phi \,u_R + h.c.$$
$$\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}, \ \langle \phi^0 \rangle = \frac{v_2}{\sqrt{2}}$$

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$$\begin{aligned} \tilde{\chi} &= (3, Y/2 = 1) \\ \tilde{\chi} &= \frac{1}{\sqrt{2}} \, \vec{\tau} . \vec{\chi} = \begin{pmatrix} \frac{1}{\sqrt{2}} \chi^+ & \chi^{++} \\ \chi^0 & -\frac{1}{\sqrt{2}} \chi^+ \end{pmatrix} \end{aligned}$$

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$$\mathcal{L}_{M} = g_{M} \left(l_{R}^{M,T} \sigma_{2} \right) \left(i \tau_{2} \tilde{\chi} \right) l_{R}^{M}$$

$$= g_{M} \nu_{R}^{T} \sigma_{2} \nu_{R} \chi^{0} - \frac{1}{\sqrt{2}} \nu_{R}^{T} \sigma_{2} e_{R}^{M} \chi^{+} + \dots$$

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$$L_{M} = g_{M} \left(l_{R}^{M,T} \sigma_{2} \right) \left(i \tau_{2} \tilde{\chi} \right) l_{R}^{M} \\ = g_{M} \nu_{R}^{T} \sigma_{2} \nu_{R} \chi^{0} - \frac{1}{\sqrt{2}} \nu_{R}^{T} \sigma_{2} e_{R}^{M} \chi^{+} + \dots$$

If $\langle \chi^0 \rangle = v_M \sim \Lambda_{EW} \rightarrow Majorana mass M_R = g_M v_M$

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In order to restore the Custodial Symmetry (ρ = 1) (Chanowitz, Golden and Georgi, Machacek)

$$\langle \chi^0 \rangle = \langle \xi^0 \rangle = v_M$$
Higgs Sector: One Singlet (in original Model)

$$\phi_{S}(1, Y/2 = 0)$$

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 $L_S = g_{Sl} \bar{l}_L \phi_S l_R^M + h.c.$ $= g_{Sl} (\bar{\nu}_L \nu_R + \bar{e}_L e_R^M) \phi_S + h.c..$

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= $g_{Sl} (\bar{\nu}_{L} \nu_{R} + \bar{e}_{L} e_{R}^{M}) \phi_{S} + h.c.$

With $\langle \Phi_{S} \rangle = V_{S,} \rightarrow \text{gives the Dirac mass } m^{D} = g_{S'} V_{S}$

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The Yukawa interactions

 $L_S = \bar{l}_L^0 \left(g_{0S} \phi_{0S} + g_{1S} \tilde{\phi}_S + g_{2S} \tilde{\phi}_S \right) l_R^{M,0} + H.c. ,$

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Obtain the neutrino mass matrix

$$M_{\nu}^{D} = \begin{pmatrix} g_{0S}v_{0} & g_{1S}v_{3} & g_{2S}v_{2} \\ g_{2S}v_{3} & g_{0S}v_{0} & g_{1S}v_{1} \\ g_{1S}v_{2} & g_{2S}v_{1} & g_{0S}v_{0} \end{pmatrix}$$

Some works done on this Model:

EW precision

V. Hoang, P. Q. Hung and A. S. Kamat, *Nucl. Phys. B* 877, 190 (2013) [arXiv:1303.0428 [hep-ph]].

 Implications of the 125-GeV SM-like scalar: Dr Jekyll (SM-like) and Mr Hyde (very different from SM) V. Hoang, P. Q. Hung and A. S. Kamat, arXiv:1412.0343 [hep-ph] (To appear in Nuclear Physics B).

Signals of mirror fermions (Paper in preparation)
 P.Q. Hung, Trinh Le (UVA); Nandi, Chakdar, Gosh
 (Oklahoma State University).

On neutrino and charged lepton masses and mixings: A view from the electroweak-scale righthanded neutrino model P. Q. Hung and T. Le, arXiv:1501.02538 [hep-ph].

The calculation

$$\begin{split} & l_{j} \rightarrow l_{j} + \gamma \\ & \text{The Yukawa interaction can be cast into the following form:} \\ & \mathcal{L}_{\mathrm{Charged},S}^{A_{4}} = -\sum_{k=0}^{3}\sum_{i,m=1}^{3} \left(\overline{l_{Li}}\mathcal{U}_{im}^{Lk}l_{Rm}^{M} + \overline{l_{Ri}}\mathcal{U}_{im}^{Rk}l_{Lm}^{M}\right)\phi_{Sk} + \mathrm{H.c.} \end{split}$$

$$I_{i} \rightarrow I_{j} + \gamma$$
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Where

$$\mathcal{U}_{im}^{Lk} = \left(U_{PMNS}^{\dagger} \cdot M^{k} \cdot U_{PMNS}^{lM} \right)_{im}$$

$$\mathcal{U}_{im}^{Rk} = \left(U_{PMNS}^{\dagger} \cdot M^{k} \cdot U_{PMNS}^{lM} \right)$$

$$\mathcal{U}_{im}^{R\,\kappa} = \left(U_{\rm PMNS}^{\prime\,\tau} \cdot M^{\prime\,\kappa} \cdot U_{\rm PMNS}^{\prime\,\mu} \right)_{im}$$

100 A

$$I_i \rightarrow I_j + \gamma$$

The Lorentz-invariant amplitude for the process $I_i(p)$ $\rightarrow I_j(p') + \gamma(q)$ is

$$\mathcal{M}\left(l_{i}^{-} \to l_{j}^{-} \gamma\right) = \epsilon_{\mu}^{*}(q) \bar{u}_{j}(p') \left\{ i\sigma^{\mu\nu}q_{\nu} \left[C_{L}^{ij}P_{L} + C_{R}^{ij}P_{R}\right] \right\} u_{i}(p)$$

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The Lorentz-invariant amplitude for the process $I_i(p)$ $\rightarrow I_i(p') + \gamma(q)$ is $\mathcal{M}\left(l_{i}^{-} \to l_{j}^{-} \gamma\right) = \epsilon_{\mu}^{*}(q) \bar{u}_{j}(p') \left\{ i\sigma^{\mu\nu} q_{\nu} \left[C_{L}^{ij} P_{L} + C_{R}^{ij} P_{R} \right] \right\} u_{i}(p)$ Where $P_{R,L} = 1/2 (1 \pm \gamma^5)$ and $C_L^{ij} = +\frac{e}{16\pi^2} \sum_{l=0}^3 \sum_{j=1}^3 \left\{ \frac{1}{m_{lM}^2} \left[m_i \mathcal{U}_{jm}^{Rk} \left(\mathcal{U}_{im}^{Rk} \right)^* + m_j \mathcal{U}_{jm}^{Lk} \left(\mathcal{U}_{im}^{Lk} \right)^* \right] \mathcal{I} \left(\frac{m_{\phi_{Sk}}^2}{m_{lM}^2} \right) \right\}$ $+\frac{1}{m_{l_m^M}}\mathcal{U}_{jm}^{R\,k}\left(\mathcal{U}_{im}^{L\,k}\right)^*\mathcal{J}\left(\frac{m_{\phi_{Sk}}^2}{m_{I_M}^2}\right)\right\}$ $C_R^{ij} = +\frac{e}{16\pi^2} \sum_{k=0}^3 \sum_{m=1}^3 \left\{ \frac{1}{m_{l_m}^2} \left[m_i \mathcal{U}_{jm}^{L\,k} \left(\mathcal{U}_{im}^{L\,k} \right)^* + m_j \mathcal{U}_{jm}^{R\,k} \left(\mathcal{U}_{im}^{R\,k} \right)^* \right] \mathcal{I} \left(\frac{m_{\phi_{Sk}}^2}{m_{l_m}^2} \right) \right\}$ $+\frac{1}{m_{l_{m}}}\mathcal{U}_{jm}^{L\,k}\left(\mathcal{U}_{im}^{R\,k}\right)^{*}\mathcal{J}\left(\frac{m_{\phi_{Sk}}^{2}}{m_{l_{M}}^{2}}\right)\right\}$

$$I_i \rightarrow I_j + \gamma$$

The partial width

$$\Gamma(l_i \to l_j \gamma) = \frac{1}{16\pi} m_{l_i}^3 \left(1 - \frac{m_{l_j}^2}{m_{l_i}^2} \right)^3 \left(|C_L^{ij}|^2 + |C_R^{ij}|^2 \right)$$

The magnetic dipole moment anomaly for
lepton
$$l_i$$

$$\Delta a_{l_i} = \frac{2m_{l_i}}{e} \left(\frac{C_L^{ii} + C_R^{ii}}{2} \right)$$
$$= +\frac{1}{16\pi^2} \left\{ \sum_{k=0}^3 \sum_{m=1}^3 2 \left(|\mathcal{U}_{im}^{L\,k}|^2 + |\mathcal{U}_{im}^{R\,k}|^2 \right) \frac{m_{l_i}^2}{m_{l_m}^2} \mathcal{I} \left(\frac{m_{\phi_{Sk}}^2}{m_{l_m}^2} \right) \right\}$$
$$+ \sum_{k=0}^3 \sum_{m=1}^3 \operatorname{Re} \left(\mathcal{U}_{im}^{L\,k} \left(\mathcal{U}_{im}^{R\,k} \right)^* \right) \frac{m_{l_i}}{m_{l_m}} \mathcal{J} \left(\frac{m_{\phi_{Sk}}^2}{m_{l_m}^2} \right) \right\}$$

The electric dipole moment for lepton I_i

$$d_{l_i} = \frac{i}{2} \left(C_L^{ii} - C_R^{ii} \right) ,$$

$$= + \frac{e}{16\pi^2} \sum_{k=0}^3 \sum_{m=1}^3 \frac{1}{m_{l_m^M}} \operatorname{Im} \left(\mathcal{U}_{im}^{L\,k} \left(\mathcal{U}_{im}^{R\,k} \right)^* \right) \mathcal{J} \left(\frac{m_{\phi_{Sk}}^2}{m_{l_m^M}^2} \right)$$

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• For the masses of the Higgs singlets, we take

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• For the masses of the mirror lepton, we take

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• We assume all the Yukawa couplings $g_{_{0S}}$, $g_{_{1S}}$, $g_{_{2S}}$, $g'_{_{S0}}$, $g'_{_{S1}}$, and $g'_{_{S2}}$ to be all real and we also take $g_{_{0S}} = g'_{_{S0}}$, $g_{_{1S}} = g'_{_{S1}}$, $g_{_{S2}} = (g'_{_{S1}})^*$

For the three unknown mixing matrices, we consider 2 scenarios

- Scenario 1

$$U_{\rm PMNS}^{l^M} = U_{\rm PMNS}' = U_{\rm PMNS}'^{l^M} = U_{CW}^{\dagger}$$

- Scenario 2

$$U_{\rm PMNS}^{l^M} = U_{\rm PMNS}' = U_{\rm PMNS}'^{l^M} = U_{\rm PMNS}$$

Where the Cabibbo-Wolfenstein matrix

$$U_{CW} = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 1 & 1 \\ 1 & \omega & \omega^2 \\ 1 & \omega^2 & \omega \end{pmatrix}$$





• Limit on B($\mu \rightarrow e\gamma$) from MEG experiment:

B(μ → e γ) ≤ 5.7 × 10⁻¹³ (90 C.L.)[MEG, 2013] B(μ → e γ) ~ 4 × 10⁻¹⁴ [Projected Sensitivity]

• Δa_{μ} from E821 experiment

$$\Delta a_{\mu} \equiv a_{\mu}^{\exp} - a_{\mu}^{SM} = 288(63)(49) \times 10^{-11}$$









Scenario 2









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- 3. However, for Scenario 2, there is slightly different between normal and inverted cases. These differences depend on these couplings, for $g1S \ge 0.5 \times g0S$, these differences diminish.
- 4.The sensitivity of the couplings in the B($\mu \rightarrow e + \gamma$) has been weakened by one to two order of magnitudes for scenario 2, while for the muon anomalous magnetic moment it stays more or less the same.


Contour plots of Log10B($\mu \rightarrow e + \gamma$) on the (Log10 g_{0S} ;M_{mirror}) plane for normal (left panel) and inverted (right panel) hierarchy in Scenario 1 (red) and 2 (blue) with $g_{0S} = g'_{0S}$ and $g_{1S} = g'_{1S} = 10^{-2} g_{0S}$.



Contour plots of Log10B($\mu \rightarrow e + \gamma$) on the (Log10 g_{0S} ;M_{mirror}) plane for normal (left panel) and inverted (right panel) hierarchy in Scenario 1 (red) and 2 (blue) with $g_{0S} = g'_{0S}$ and $g_{1S} = g'_{1S} = 0.5 g_{0S}$.



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By the recent and future expectation data at MEG we figured out some interesting constraints on the model

