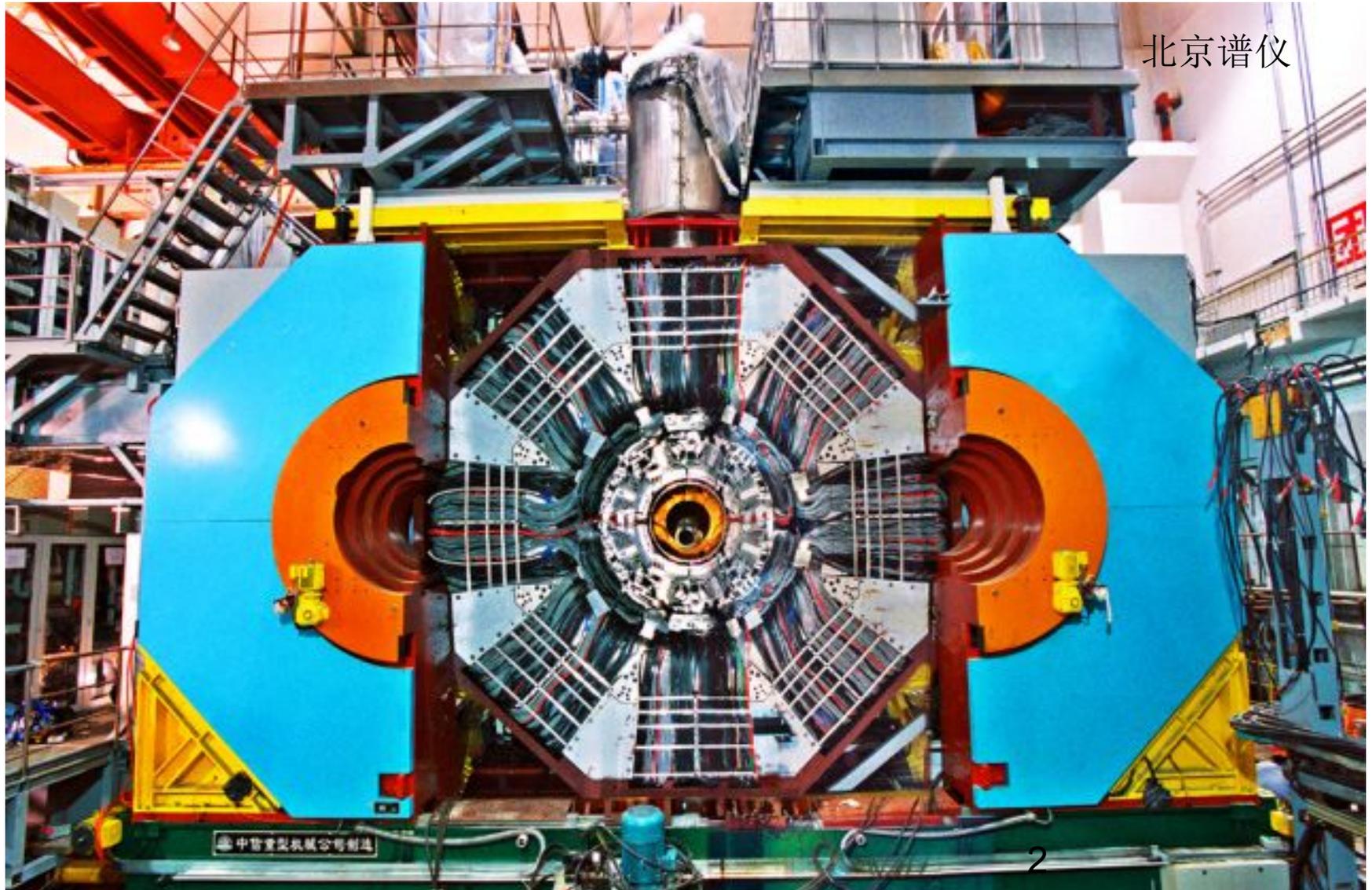


Particle Identification

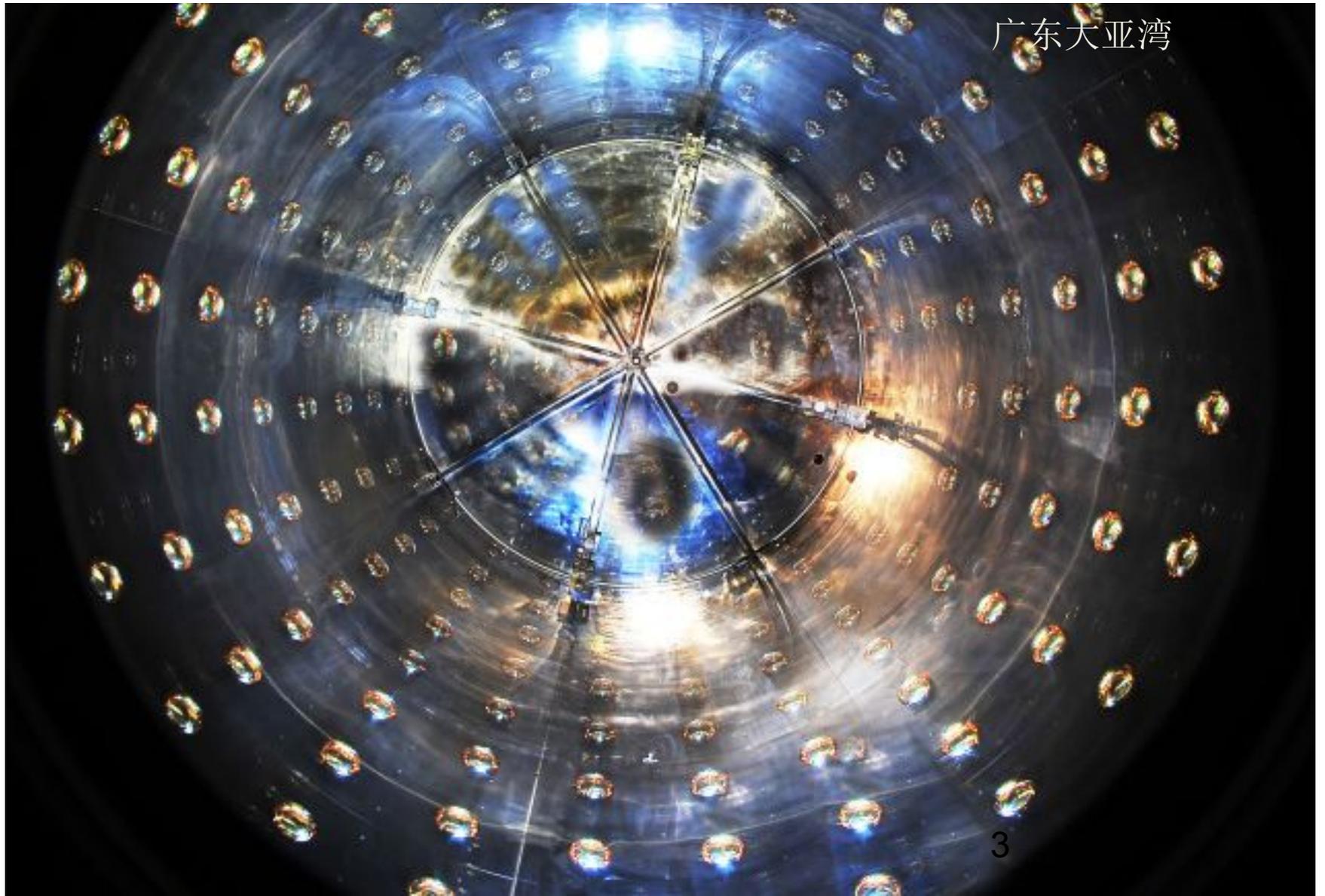
Huaqiao ZHANG (IHEP)

Some materials from Roger Forty's lecture

Example of high energy physics detectors



Example of high energy physics detectors

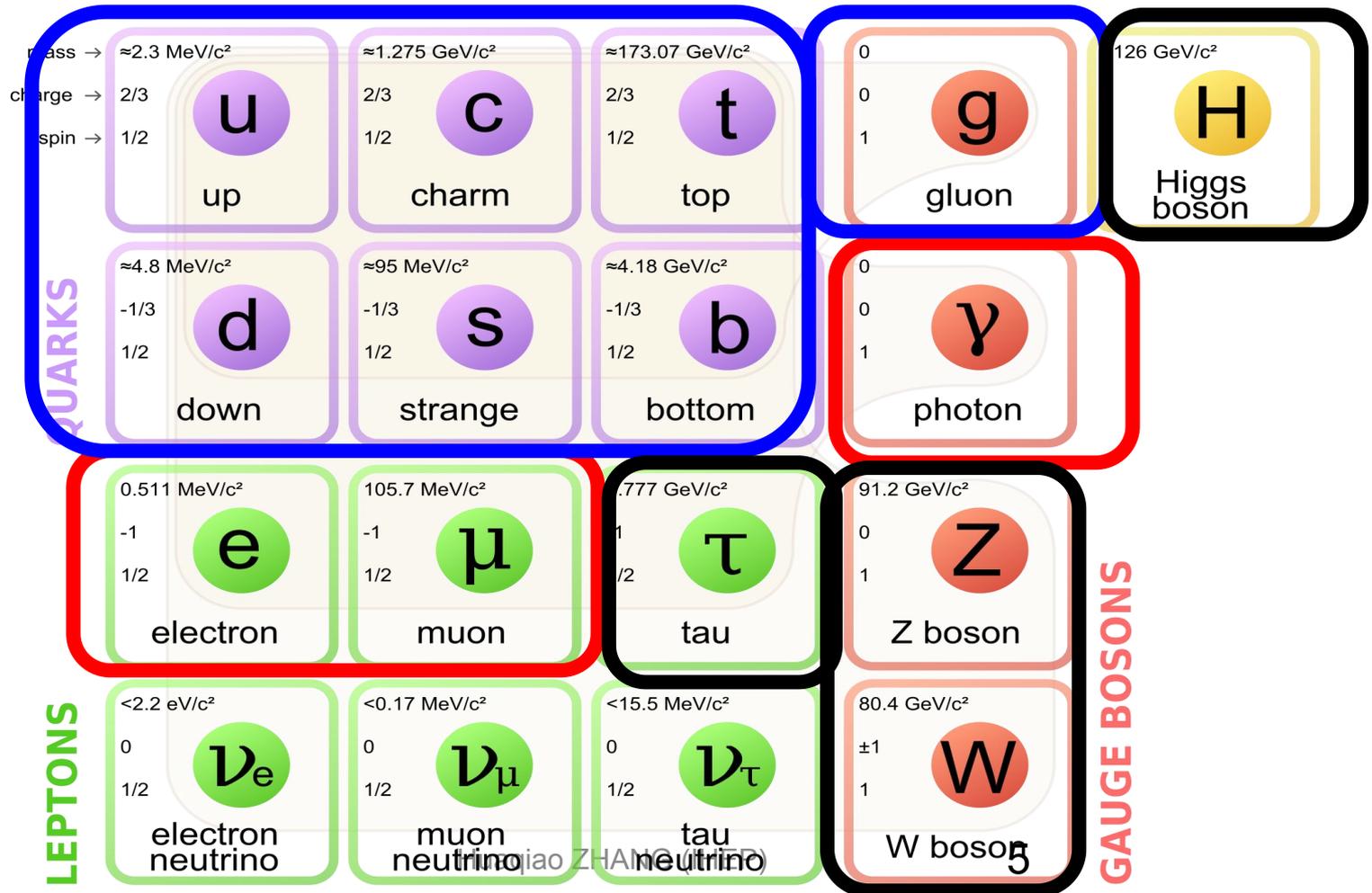


Example of high energy physics detectors



What particle to identify

- Long life time (Stable)
- Interaction with detectors => signals



Hadrons

- Instead of making do with jet reconstruction, often the physics under study requires the identification of *individual* hadrons
- Most are unstable, and decay into a few long-lived particles:

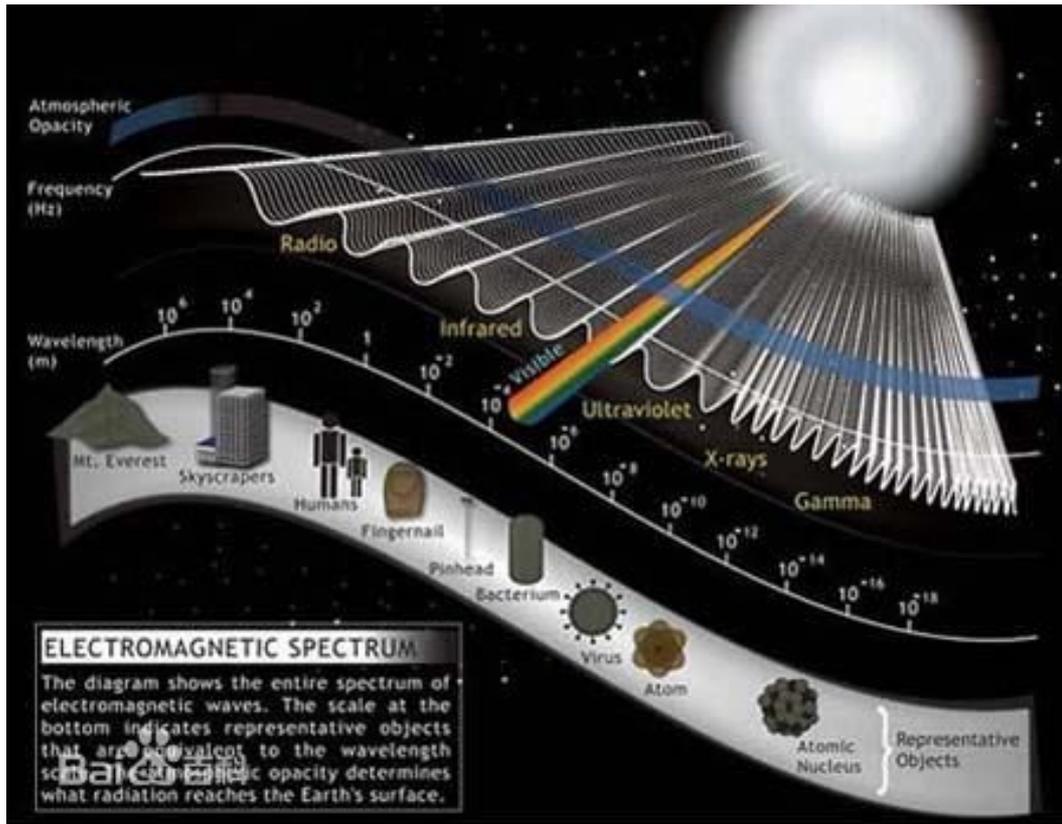
Particle	m [MeV]	Quarks	Main decay	Lifetime	$c\tau$ [cm]
π^\pm	140	$u\bar{d}$	$\mu\nu_\mu$	2.6×10^{-8} s	780
K^\pm	494	$u\bar{s}$	$\mu\nu_\mu, \pi\pi^0$	1.2×10^{-8} s	370
K_S^0	498	$d\bar{s}$	$\pi\pi$	0.9×10^{-10} s	2.7
K_L^0	498	$d\bar{s}$	$\pi\pi\pi, \pi l\nu$	5×10^{-8} s	1550
p	938	uud	stable	$> 10^{25}$ years	
n	940	udd	$p e \nu_e$	890 s	2.7×10^{13}
Λ	1116	uds	$p\pi$	2.6×10^{-10} s	7.9

Outline

- General principle for particle identification
- Examples of High Energy physics detector
 - CMS detector
 - Photons/Electrons ID
 - Muons ID
 - Taus/jets ID
 - Others ID

例子：光子的探测

- 每个人都有的光子探测器



电磁波：收音机
可见光：眼睛
红外，紫外探测器等

高能物理探测的光子能量：
keV, MeV, GeV, even TeV

General principle for particle identification

- Use interactions

	电磁相互作用	弱相互作用	强相互作用	引力相互作用
源	电荷	弱荷	色荷	质量
作用强度	$\frac{e^2}{\hbar c} \simeq \frac{1}{137}$	$(\frac{m_p c}{\hbar})^2 \frac{G}{\hbar c} \sim 10^{-5}$	$\frac{g^2}{\hbar c} \simeq 10$	$\frac{Gm_p^2}{\hbar c} \sim 10^{-38}$
力程 (m)	长程 ∞	短程 $\sim 10^{-18}$	短程 10^{-15}	长程 ∞
传递者	光子 γ	中间玻色子 W^\pm, Z^0	胶子 g	引力子(尚未发现) G
自旋宇称	1^-	$1^?$	1^-	2^+
质量(GeV)	$M_\gamma = 0$	$M_{W^\pm} = 80.4$ $M_{Z^0} = 91.2$	$M_g = 0$	$M_G = 0$
理论描述	电弱统一理论 (EW)		量子色动力学 (QCD)	几何动力学 (广义相对论)

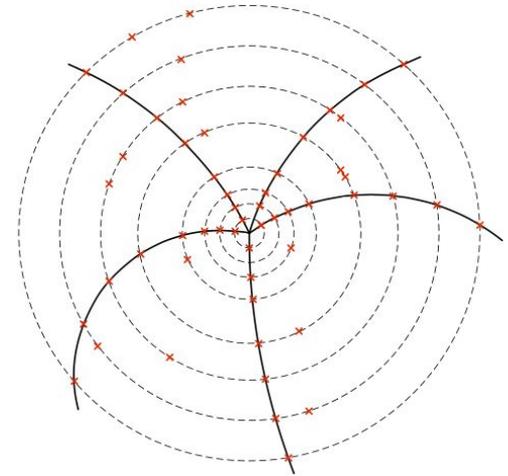
表 1.1: 四种相互作用的性质

How to Identify particles?

- The signature of these particles
 - Interactions with material
 - Masses
 - Lifetime
 - C,P,S... quantum numbers
- Methodology
 - For charged particles:
 - Measurement of their mass
 - Interaction with absorbers
 - For neutral and stable particle
 - Interaction with absorbers
 - For unstable particles
 - Study the decay products

Charged particles ID

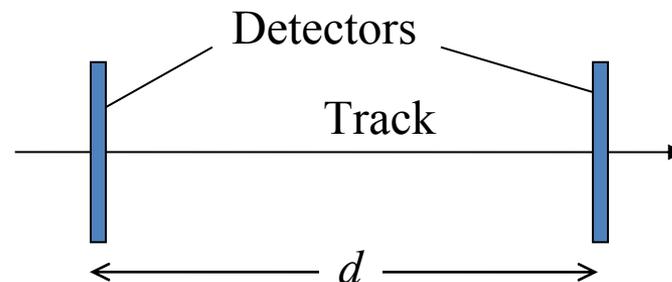
- **Everyone has charge**
 - *Determine their **mass** to ID particles*
- **$p = \gamma m v$**
 - momentum is measured by the check the curving in the magnet field of tracking system: $pT = q B R$
 - Then need to know the **velocity**



- Velocity measured by 4 main processes :
 1. **Time Of Flight (TOF)** of the particles over a fixed distance
 2. Alternatively one can look at the detail of their interaction with matter
The main source of energy loss is via **ionization** (dE/dx)
 3. If the velocity of the particle changes compared to the local speed of light it will radiate photons, detected as **Transition radiation**
 4. If a particle travels at *greater* than the local speed of light, it will radiate **Cherenkov radiation**

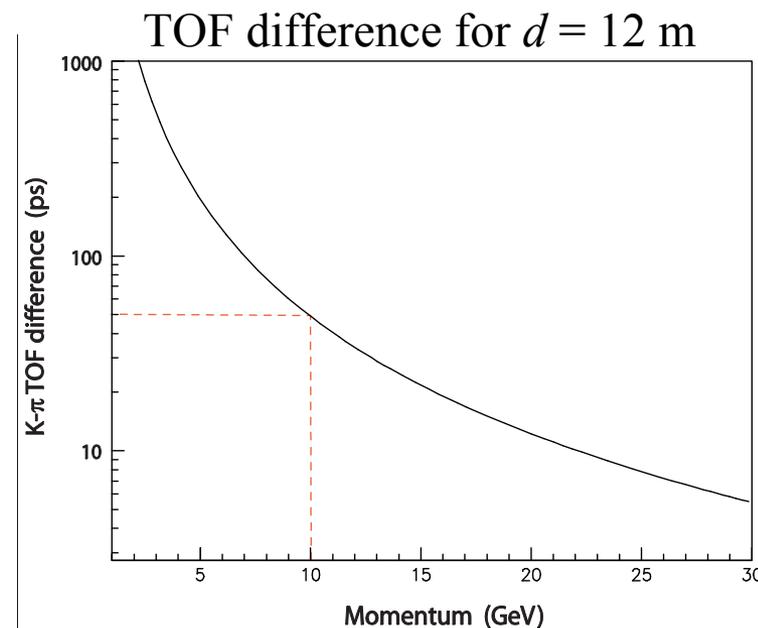
Time Of Flight

- Simple concept: measure the time difference between two detector planes
 $\beta = d / c\Delta t$
- At high energy, particle speeds are relativistic, closely approaching to c
- For a 10 GeV K , the time to travel 12 m is 40.05 ns, whereas for a π it would be 40.00 ns, so the difference is only 50 ps



- Moderr
 have resolution $\sigma_t \sim 10$ ns, fast enough for the LHC (bunch crossings 25 ns apart) but need $\sigma_t < 1$ ns to do useful TOF
- TOF gives good ID at low momentum
 Very precise timing required for $p > 5$ GeV

$$\delta t = L \left(\frac{1}{v_1} - \frac{1}{v_2} \right) \approx \frac{Lc}{2p^2} (m_1^2 - m_2^2)$$

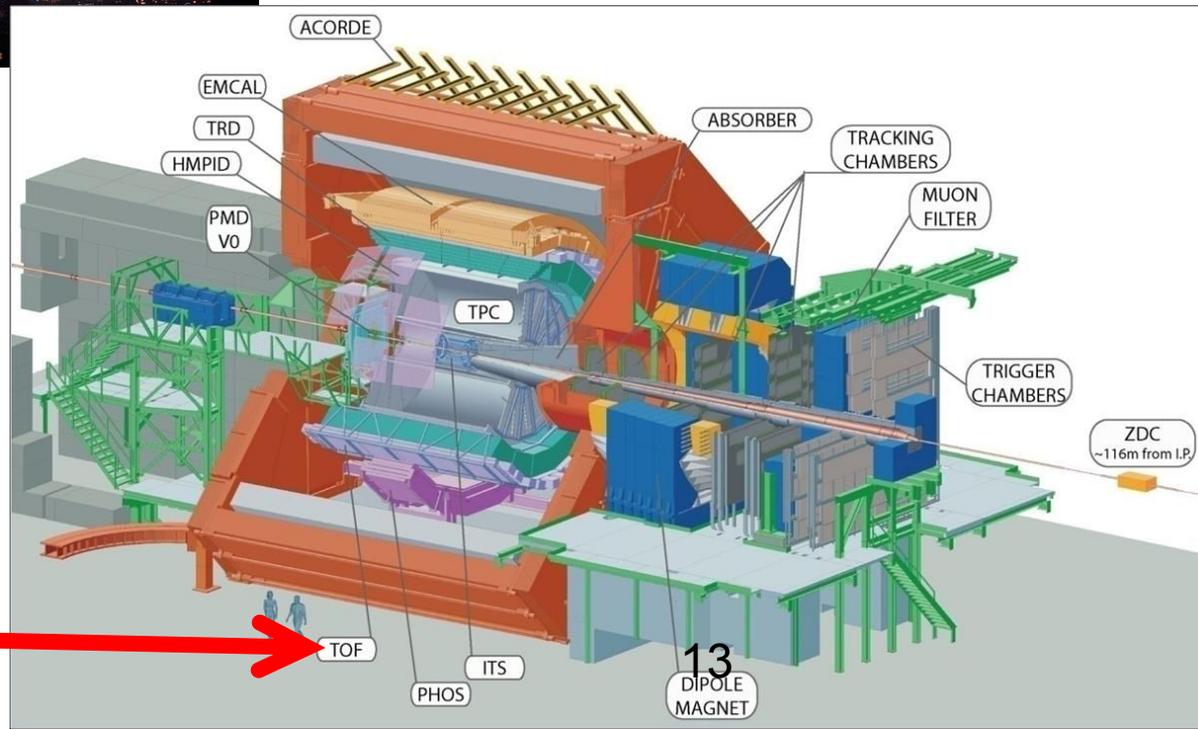


Example usage



对于远处的闪电
先看到闪电：光波（光子）
后听到雷声：声波（声子）

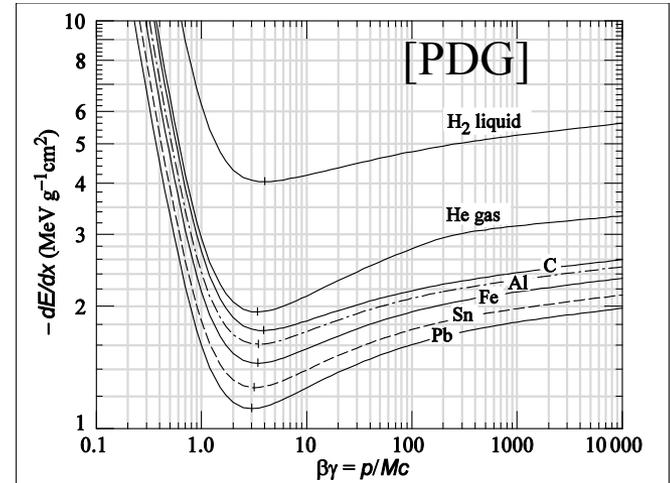
对于太近的闪电，人耳可能无法
区分闪电和雷声的先后



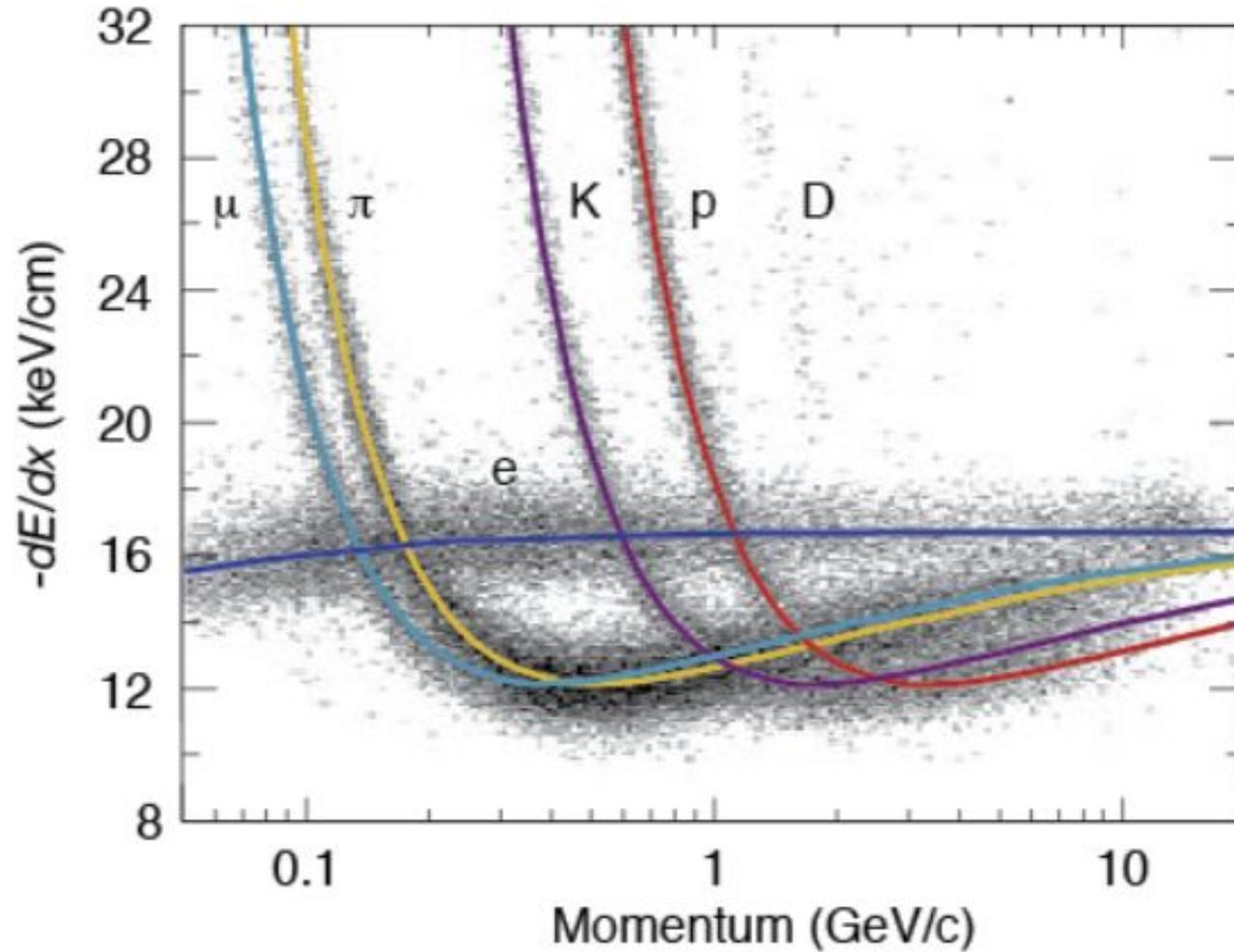
TOF在ALICE实验上的运用
2015-8-14

Ionization

- Charged particles passing through matter can knock out electrons from atoms of the medium: *ionization*
- Energy loss described by the Bethe-Bloch formula, which gives the universal velocity dependence: $dE/dx \propto \log(\beta^2 \gamma^2) / \beta^2$
- This can be used to identify particles, particularly at low momentum where dE/dx varies rapidly
- *Advantage:* uses existing detectors needed for tracking (but requires the accurate measurement of the charge)
- *Note:* these techniques all provide signals for charged *leptons* e, μ as well as π, K, p
But $m_\mu \approx m_\pi$, so they are not well separated (dedicated detectors do a better job)

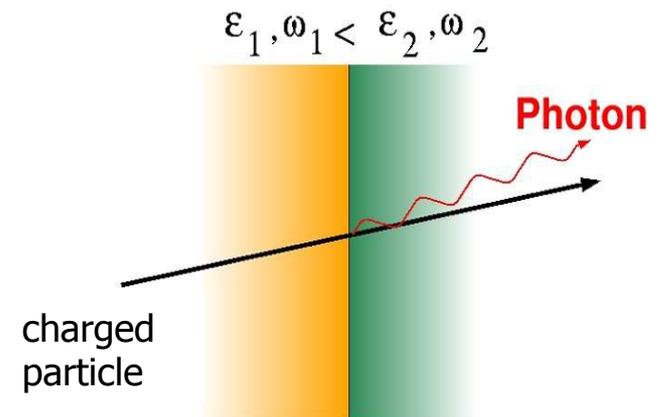


Examples



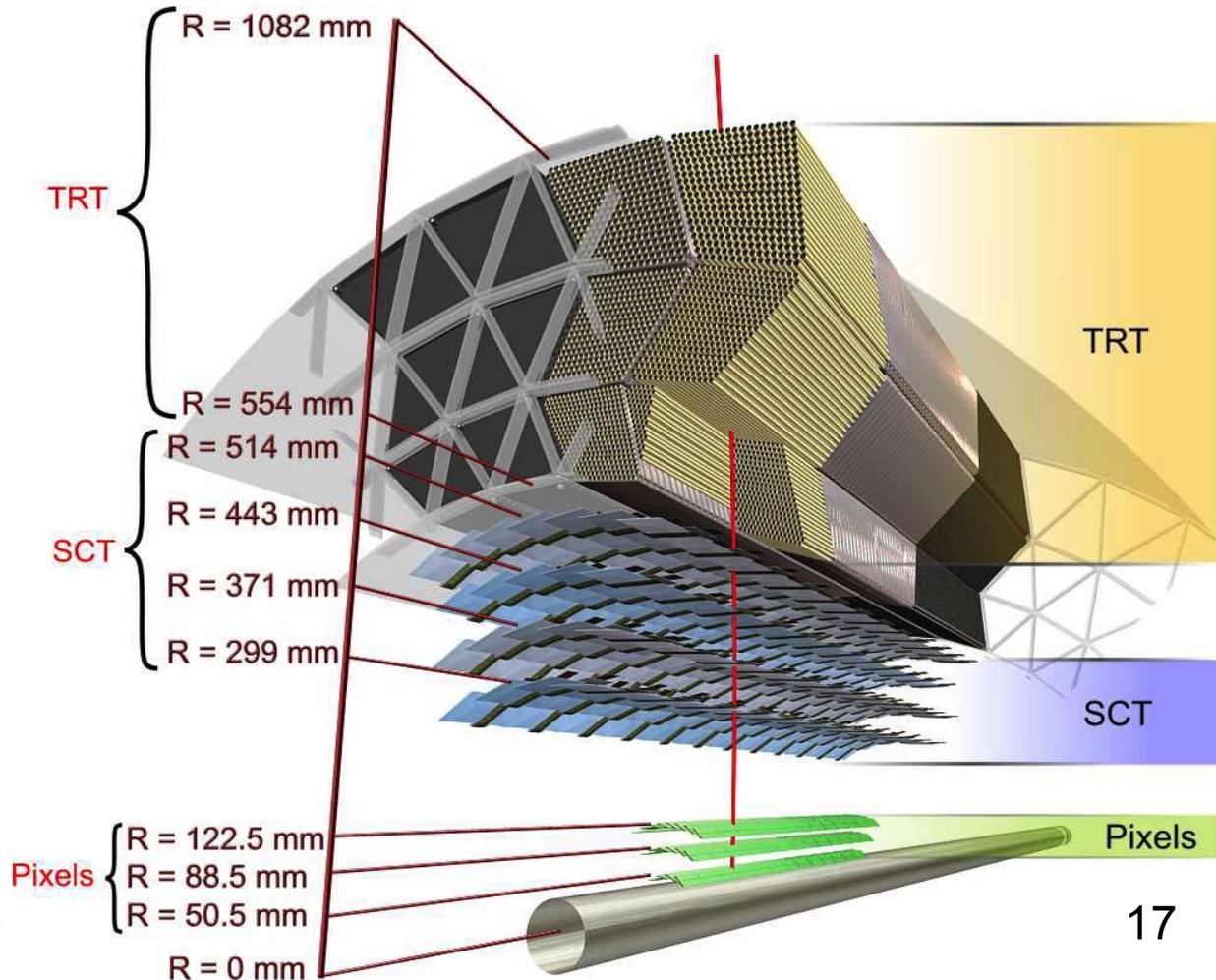
Transition radiation

- Local speed of light in a medium with refractive index n is $c_p = c/n$
- If its relative velocity v/c_p changes, a particle will radiate photons:
 1. Change of direction \mathbf{v} (in magnetic field) → Synchrotron radiation
 2. Change of $|\mathbf{v}|$ (passing through matter) → Bremsstrahlung radiation
 3. Change of refractive index n of medium → Transition radiation
- Transition radiation is emitted whenever a relativistic charged particle traverses the border between two media with different dielectric constants ($n \sim \sqrt{\epsilon}$)
- The energy emitted is proportional to the boost γ of the particle
 - Particularly useful for electron ID
 - Can also be used for hadrons at high energy



Example usage: ATLAS TRT

- Transition Radiation Tracker
 - Electron has faster velocity than other hadrons for same momentum



Cherenkov light

- Named after the Russian scientist P. Cherenkov who was the first to study the effect in depth (he won the Nobel Prize for it in 1958)
- From Relativity, nothing can go faster than the speed of light c (in vacuum)
- However, due to the refractive index n of a material, a particle *can* go faster than the *local* speed of light in the medium $c_p = c/n$
- This is analogous to the bow wave of a boat travelling over water or the sonic boom of an aeroplane travelling faster than the speed of sound



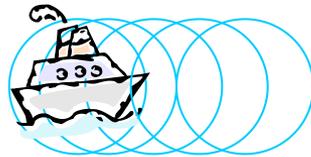
Propagating waves

- A stationary boat bobbing up and down on a lake, producing waves



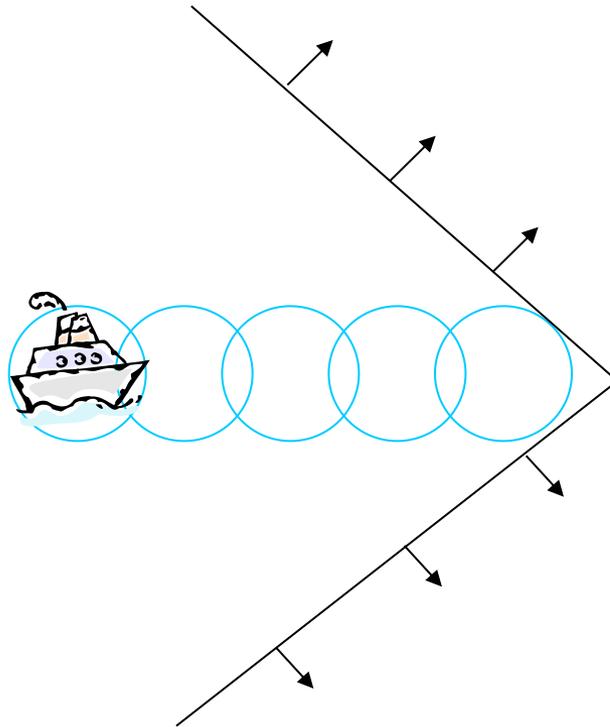
Propagating waves

- Now the boat starts to move, but slower than the waves
- No coherent wavefront is formed



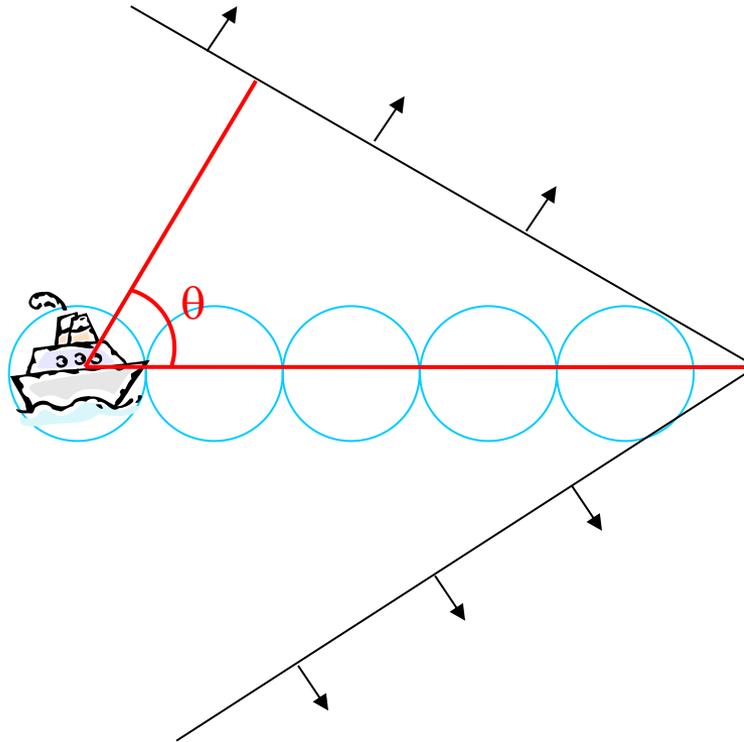
Propagating waves

- Next the boat moves faster than the waves
- A coherent wavefront is formed



Propagating waves

- Finally the boat moves even faster
- The angle of the coherent wavefront changes



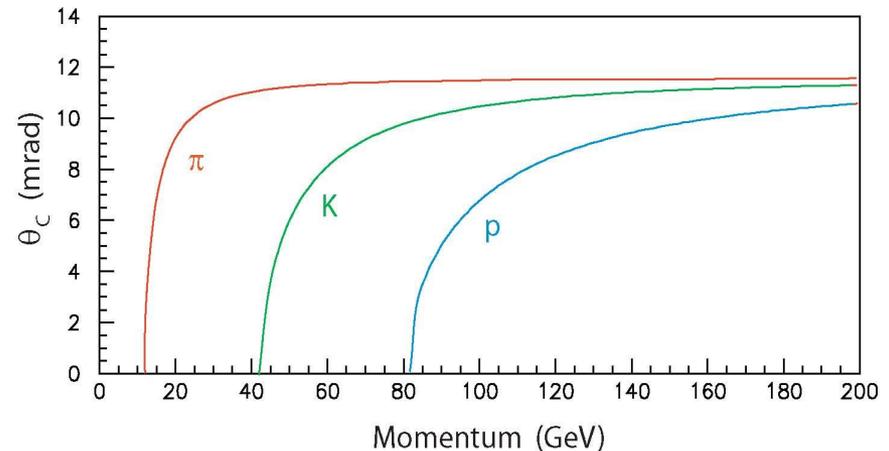
$$\cos \theta = \frac{v_{\text{wave}}}{v_{\text{boat}}}$$

Speed calculation

- Using this construction, we can determine (roughly) the boat speed:
 $\theta \approx 70^\circ$, $v_{\text{wave}} = 2 \text{ km/h}$ on water
 $\rightarrow v_{\text{boat}} = v_{\text{wave}} / \cos \theta \approx 6 \text{ km/h}$
- Cherenkov light is produced when charged particle ($v_{\text{boat}} = \beta c$) goes faster than the speed of light ($v_{\text{wave}} = c/n$)
 $\rightarrow \cos \theta_C = 1 / \beta n$
- Produced in three dimensions, so the wavefront forms a *cone* of light around the particle direction
- Measuring the opening angle of cone
 \rightarrow particle velocity can be determined

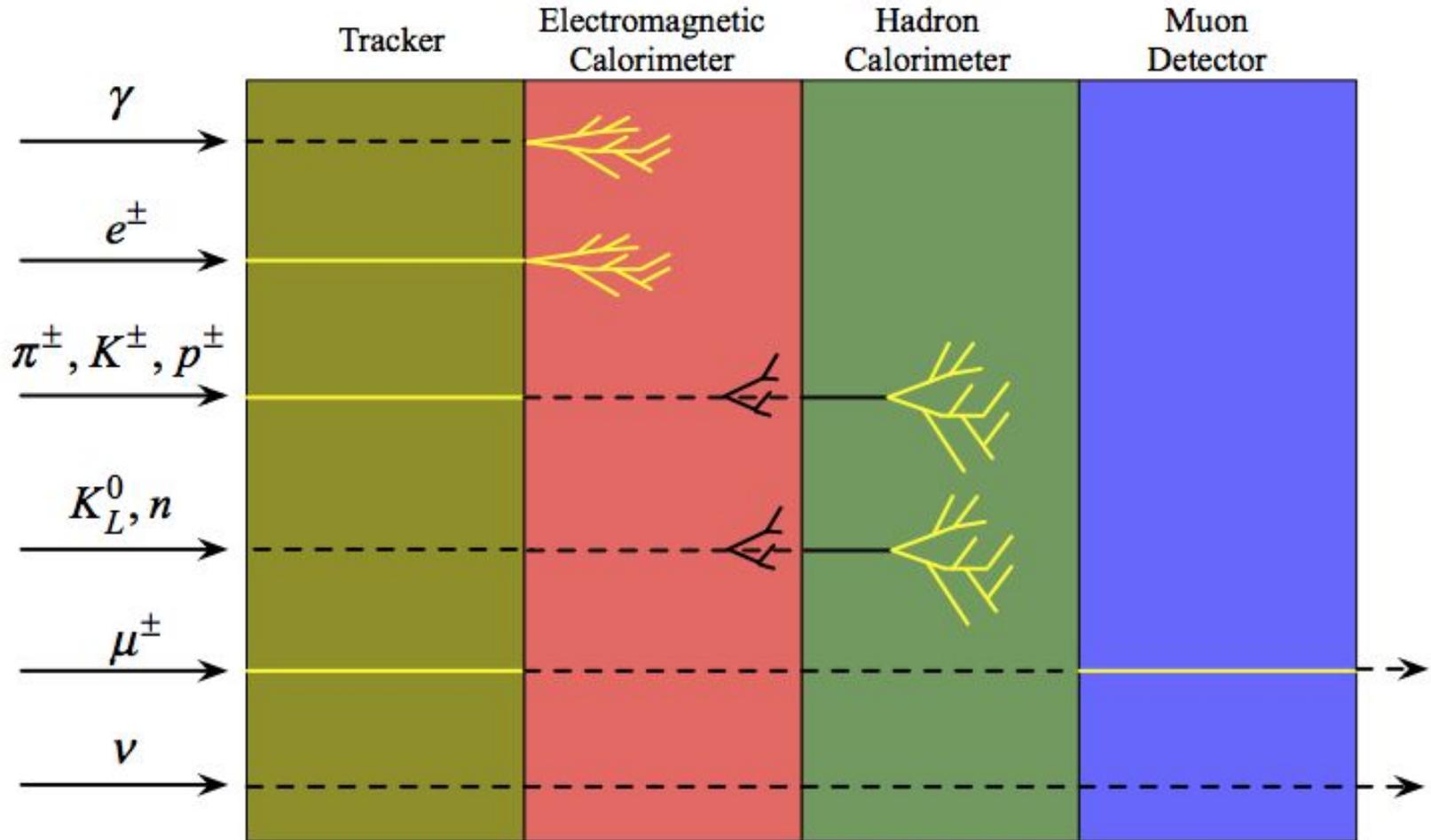


For Ne gas ($n = 1.000067$)

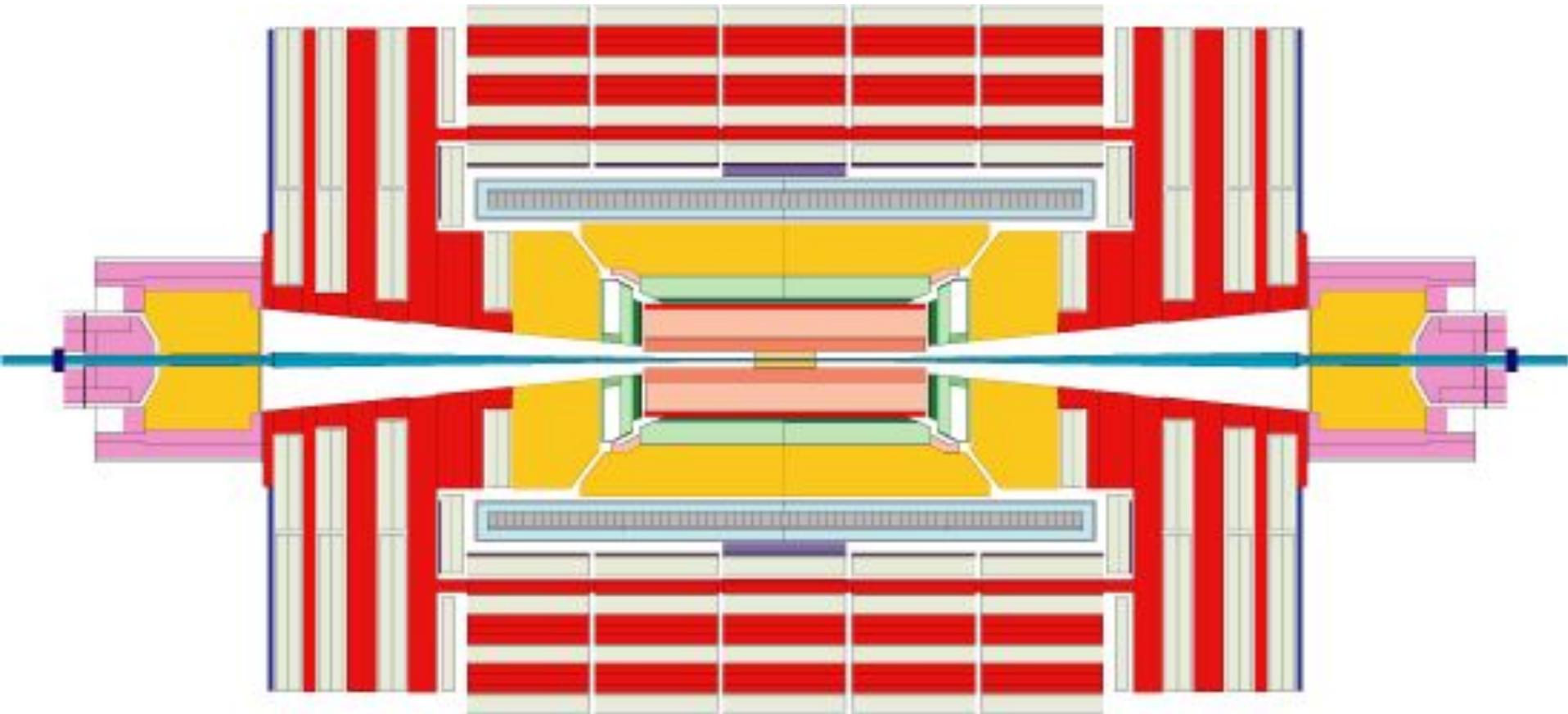


Stable particle identification

Detector Model

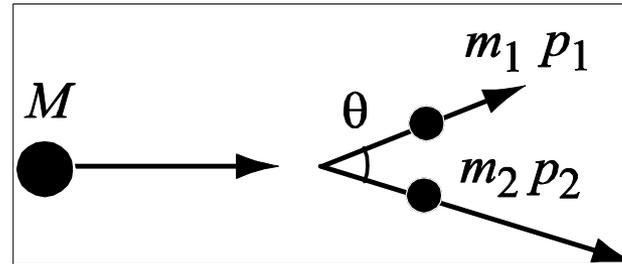


Examples: CMS detector



Unstable particle ID: two body decay

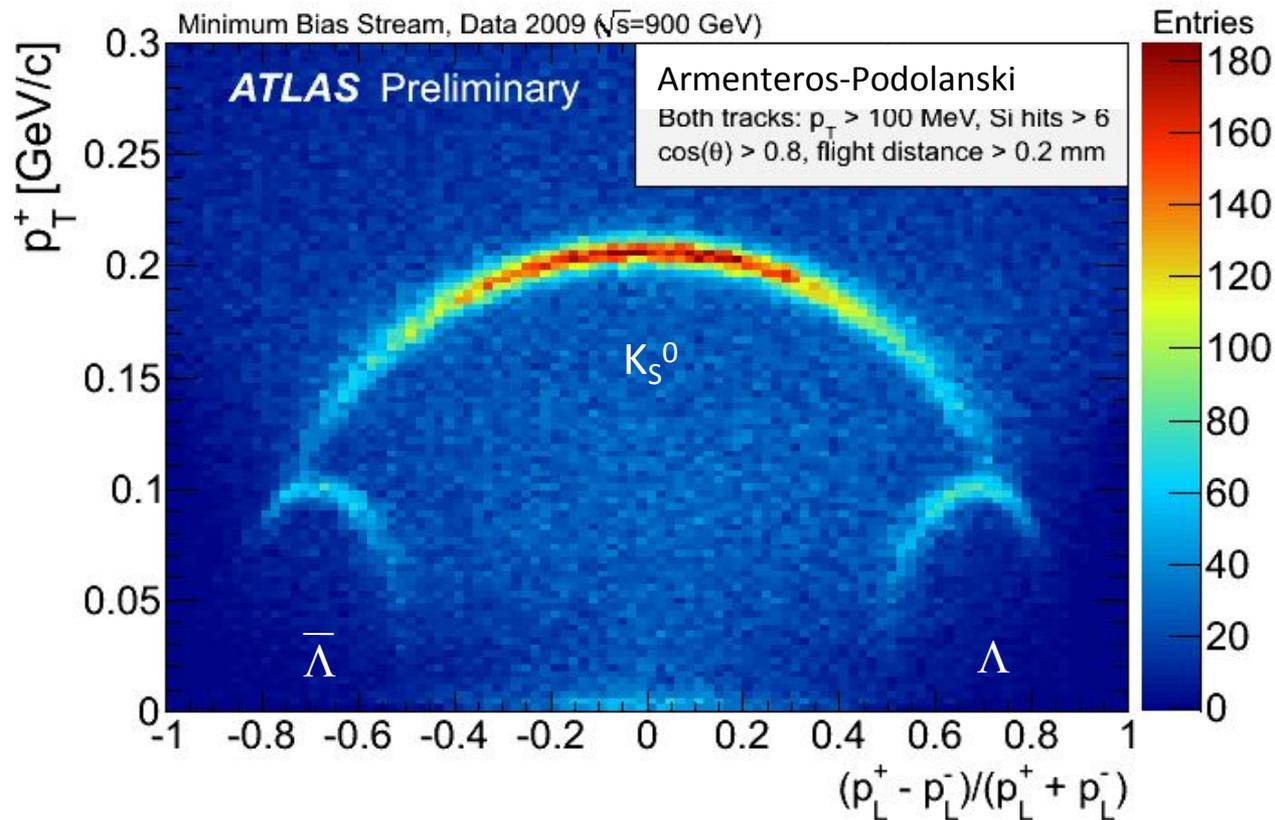
- $E^2 = p^2 + m^2$
 - The full expression is $E^2 = p^2c^2 + m^2c^4$ but factors of c are often dropped
- Consider a particle that decays to give two daughter particles:



- The *invariant mass* of the two particles from the decay:
- $M^2 = m_1^2 + m_2^2 + 2(E_1E_2 - p_1p_2 \cos\theta)$
- \rightarrow to reconstruct the parent mass a precise knowledge of the momentum and the angle θ of decay products is needed, from the tracking system, as well as their particle type, which determines their masses m_1 and m_2

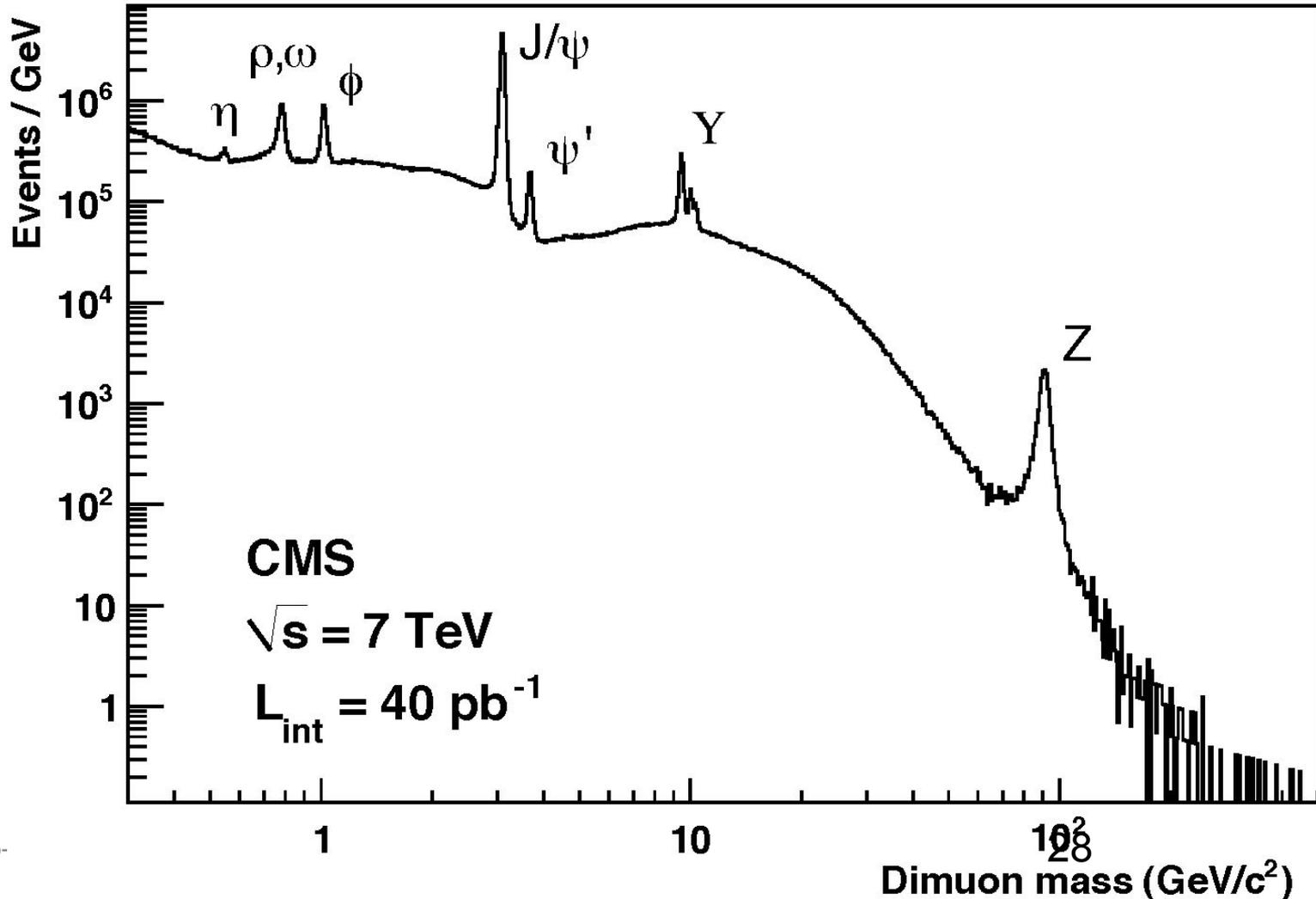
V^0 2-prong decay reconstruction

- V^0 s can be reconstructed from the kinematics of their positively and negatively charged decay products, without needing to identify the π or p

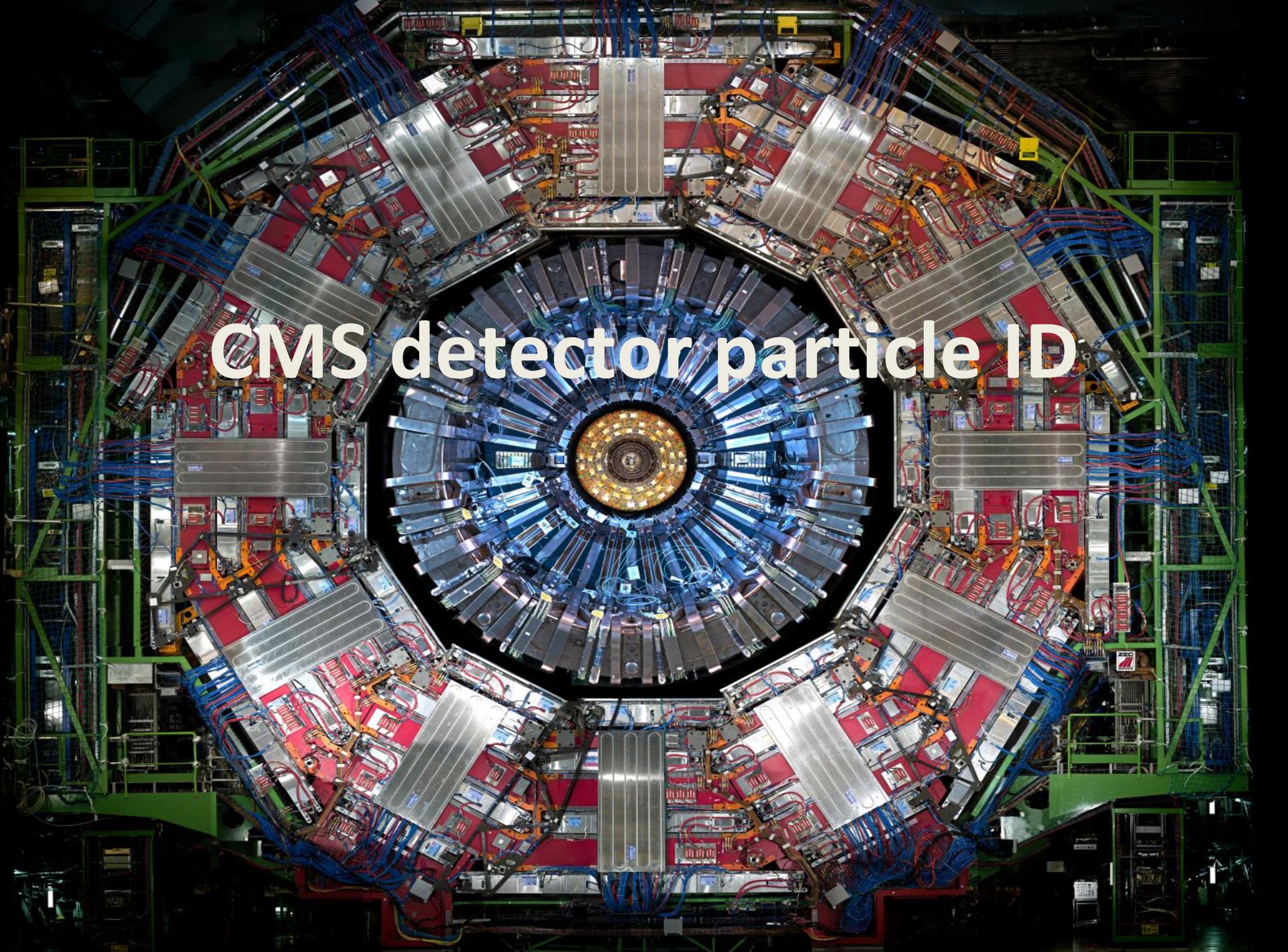


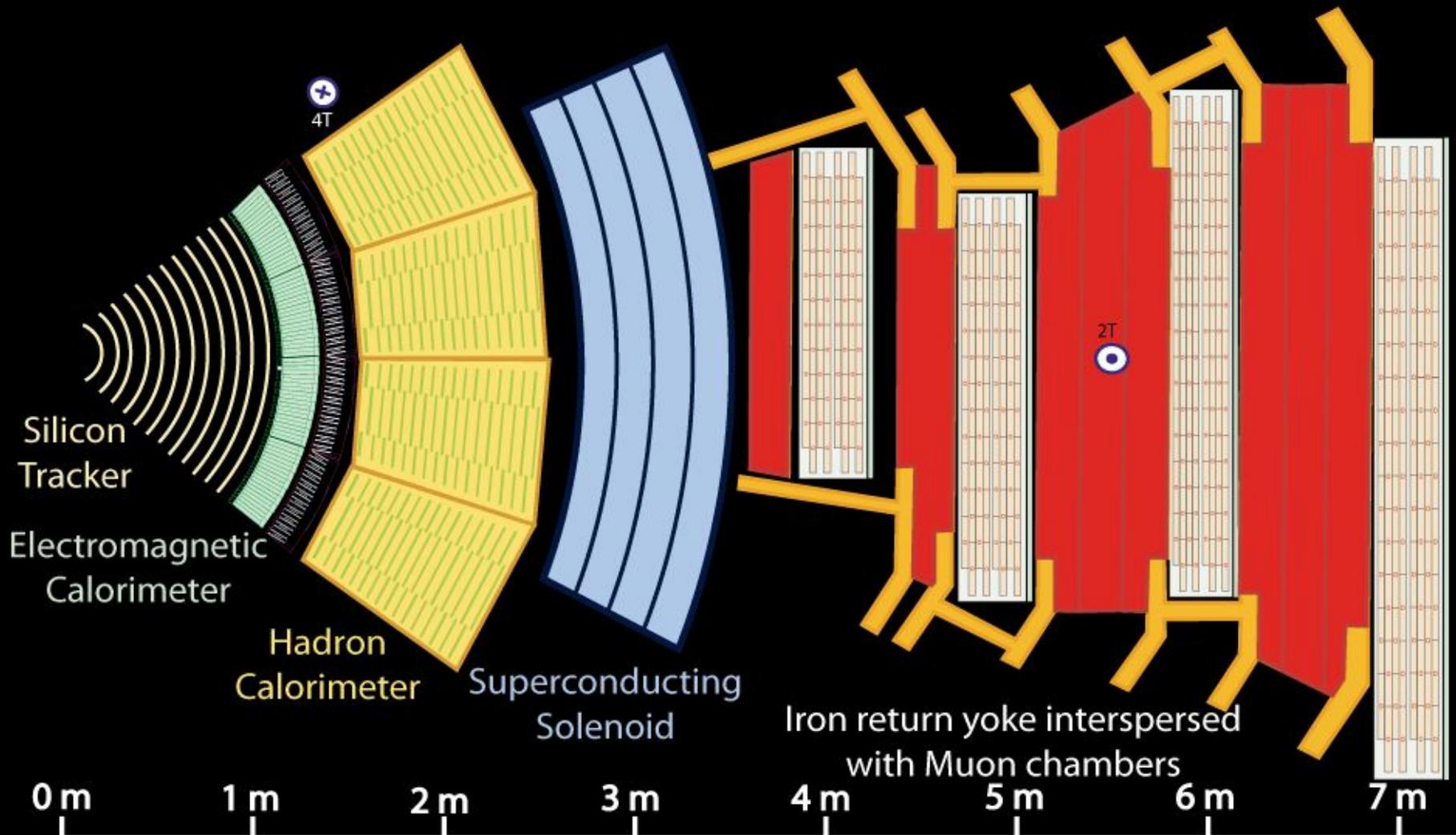
Mass reconstruction

- $M^2 = (E_1 + E_2 + \dots + E_n)^2 - (p_1 + p_2 + \dots + p_n)^2$



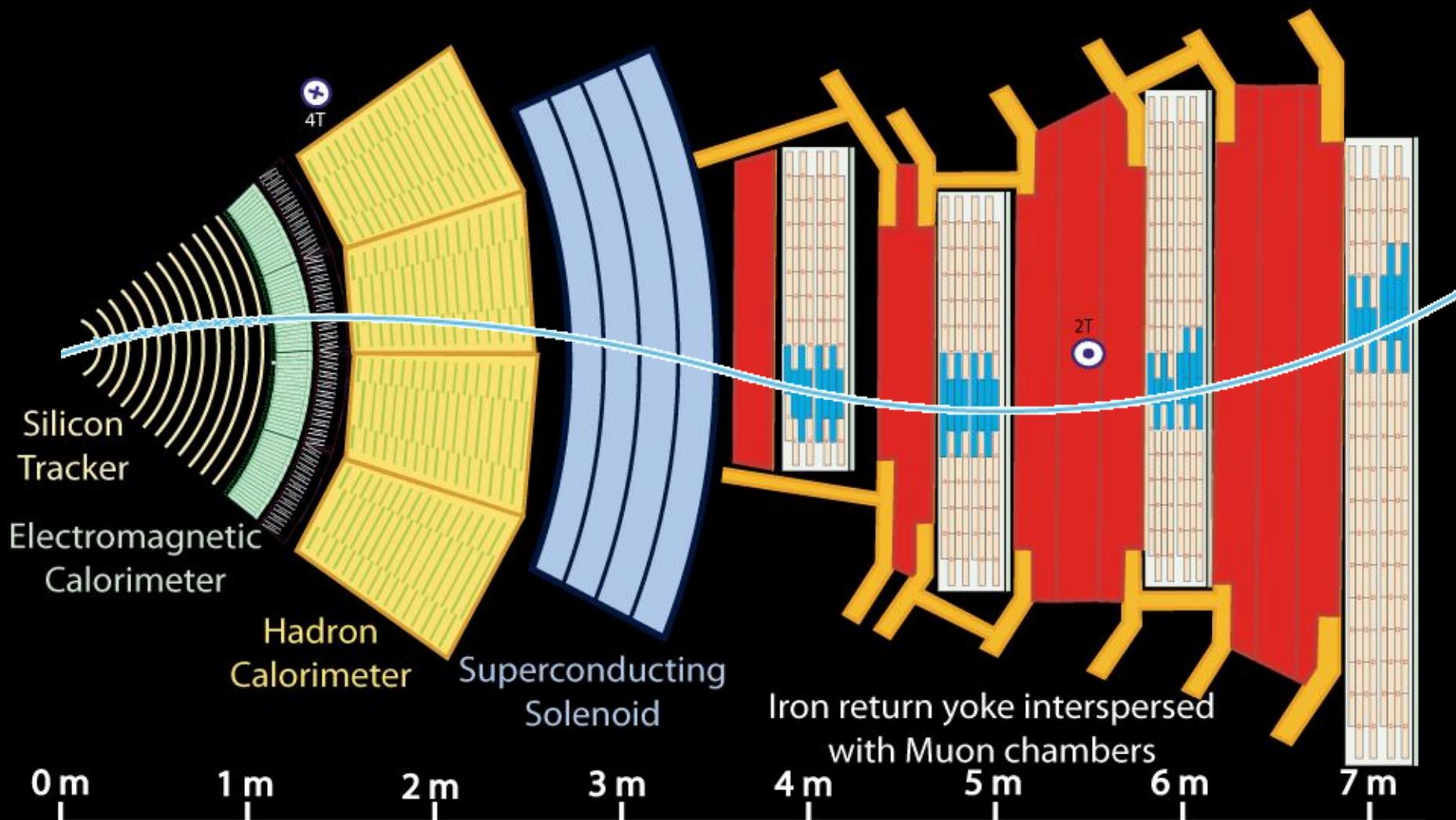
CMS detector particle ID





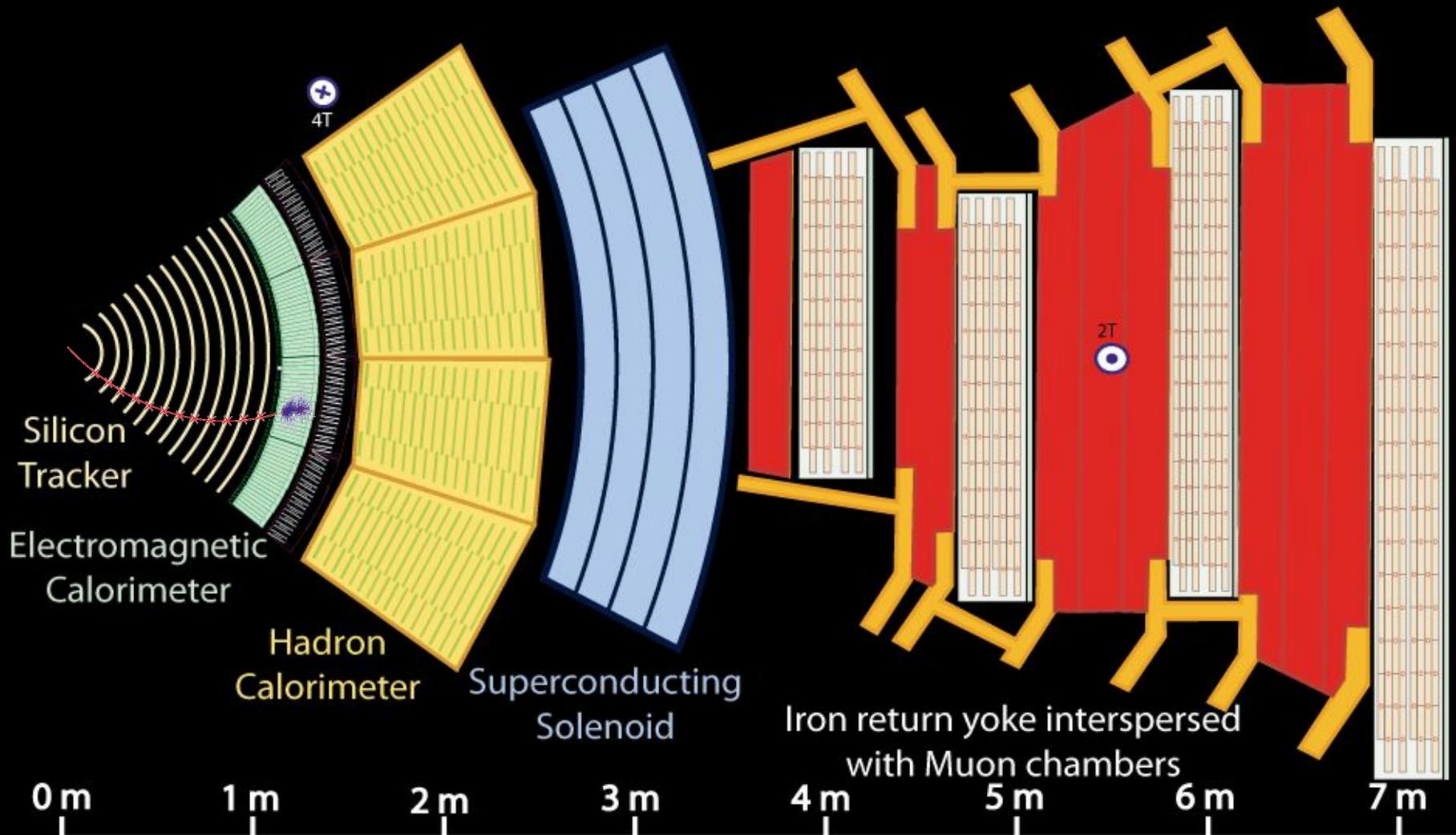
Key:

- Muon
- Electron
- Charged Hadron (e.g. Pion)
- - - Neutral Hadron (e.g. Neutron)
- - - Photon



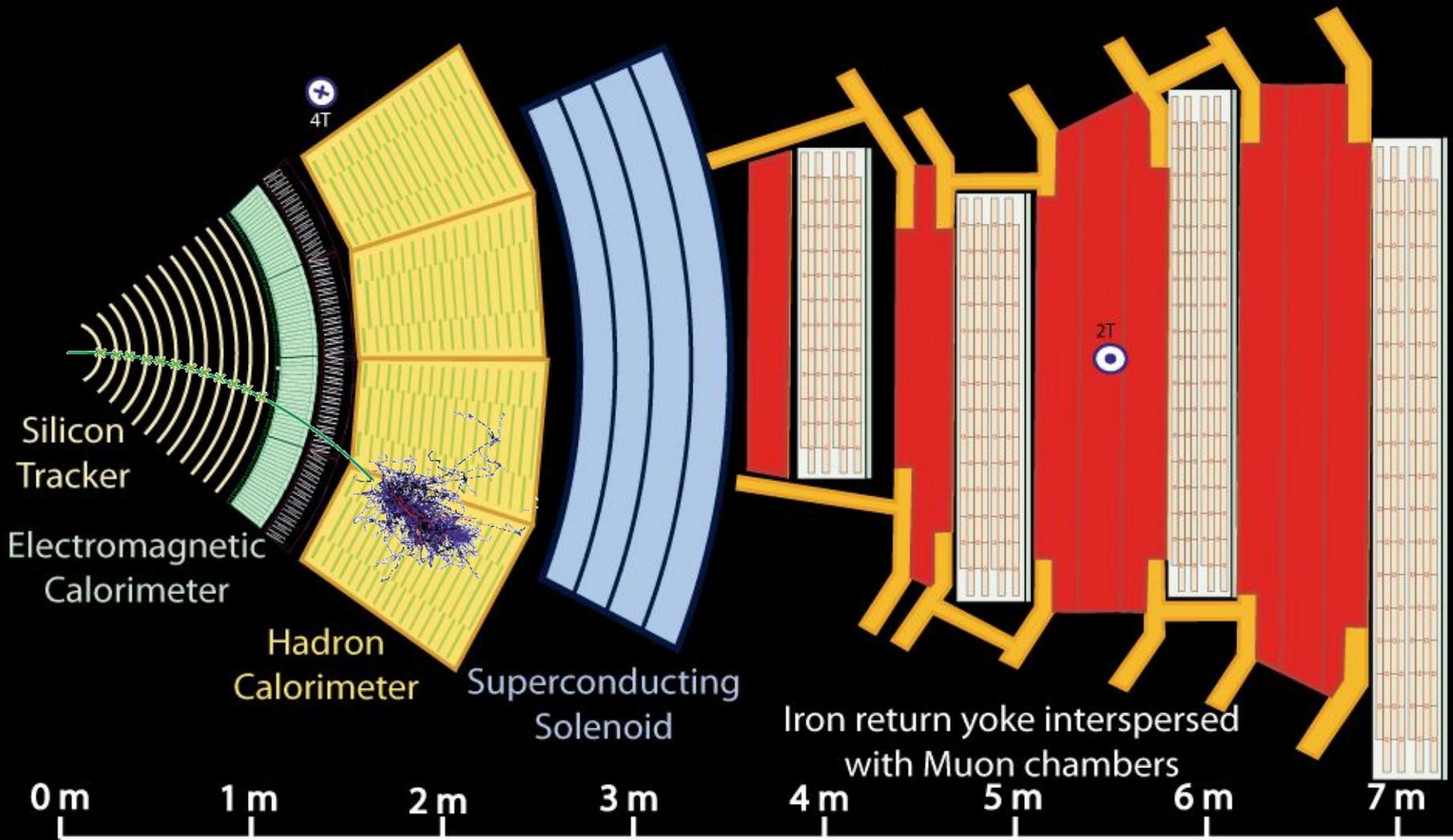
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- Muon
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Key:

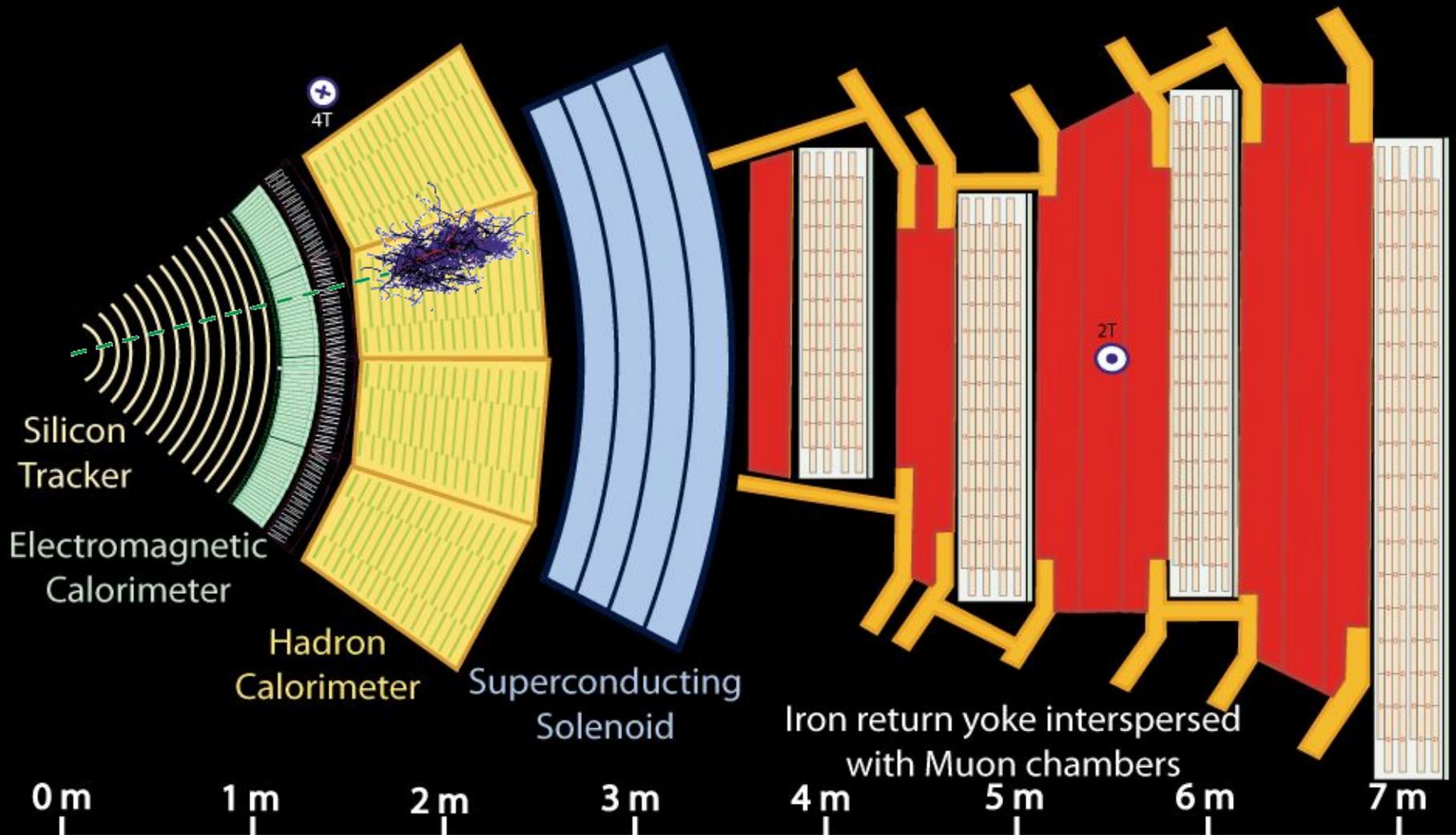
— Muon

— Electron

— Charged Hadron (e.g. Pion)

- - - Neutral Hadron (e.g. Neutron)

- - - Photon



Key:

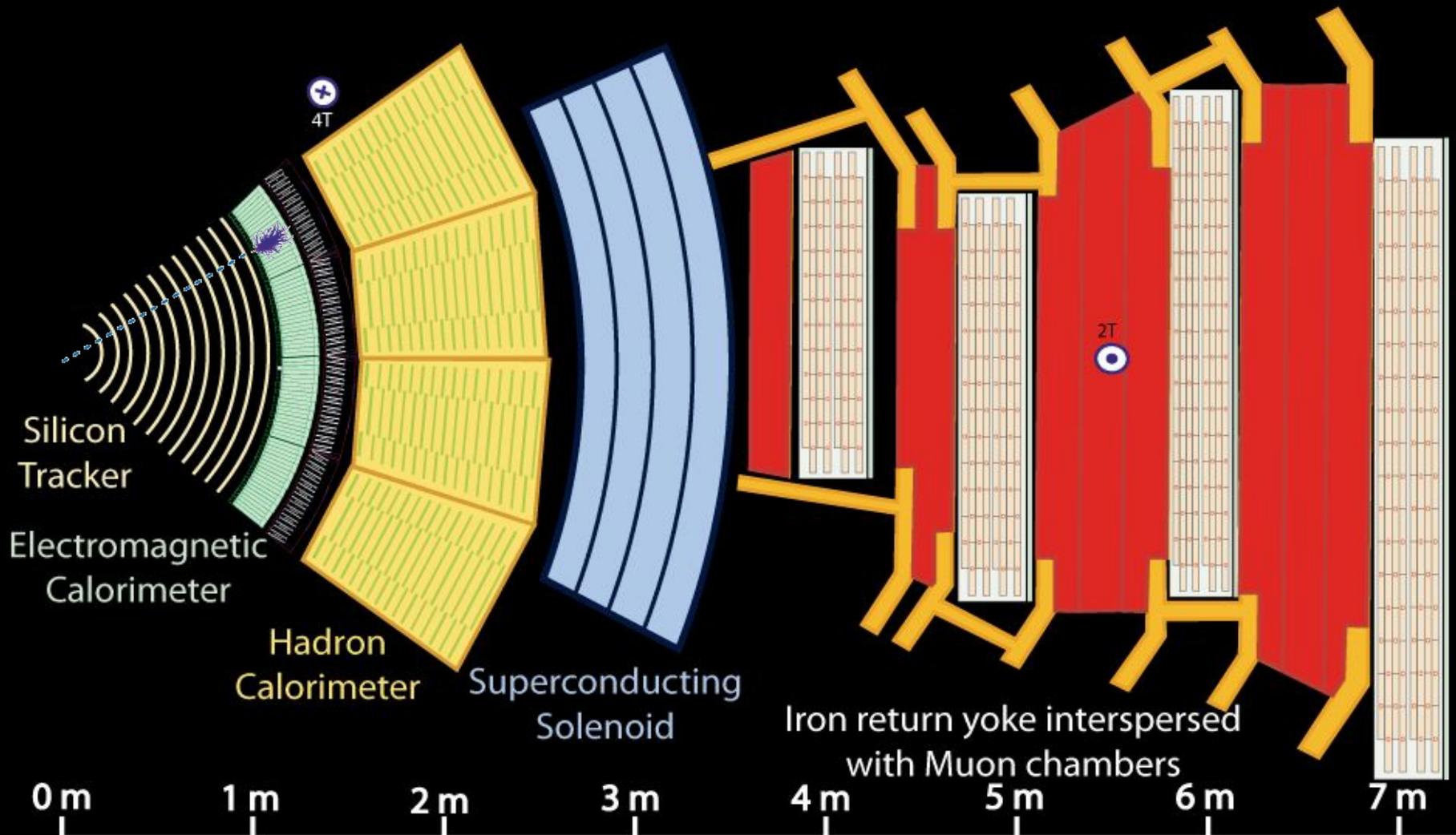
— Muon

— Electron

— Charged Hadron (e.g. Pion)

- - - Neutral Hadron (e.g. Neutron)

- - - Photon



Key:

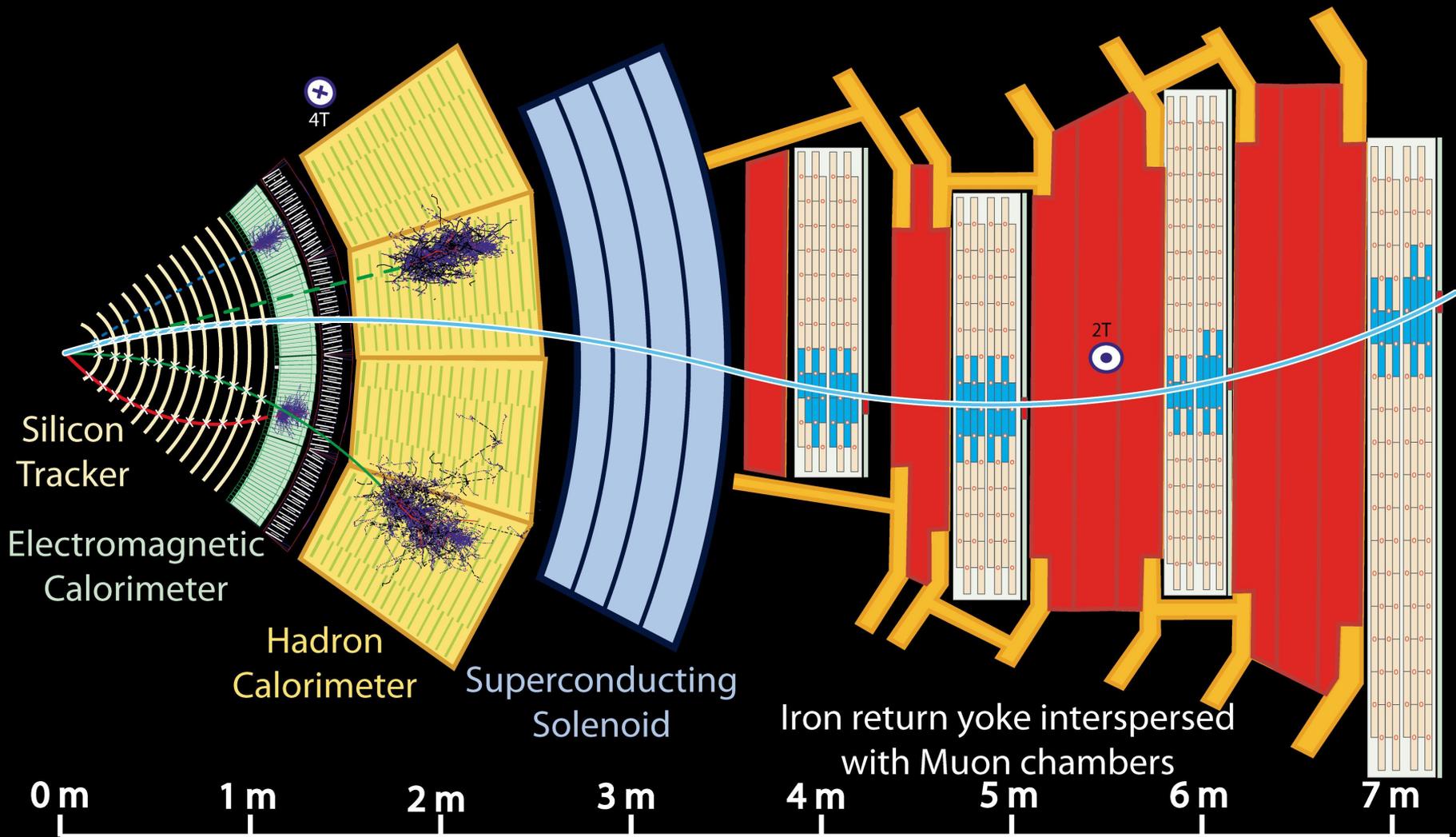
— Muon

— Electron

— Charged Hadron (e.g. Pion)

- - - Neutral Hadron (e.g. Neutron)

- - - Photon

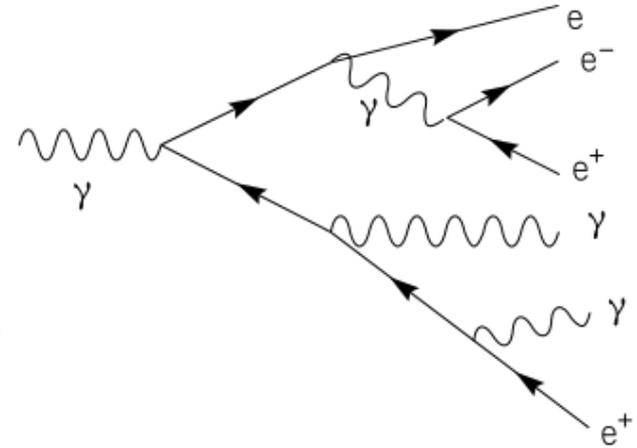


Key:

- Muon
- Electron
- Charged Hadron (e.g. Pion)
- - - Neutral Hadron (e.g. Neutron)
- - - Photon

e/γ identification (1)

- When incident on matter at high energy, photons convert to e^+e^- pairs
Since the electrons (and positrons) produce more photons by **Bremsstrahlung**, a shower develops of e^\pm and γ , until the energy of the incident particle has been used up
 - **Special shower profile**



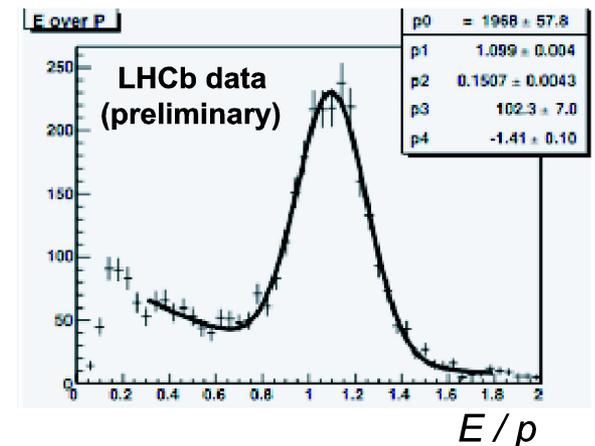
- Such showers are similar for electrons and photons

Distinguished by the existence (or not)

of a track associated to the shower

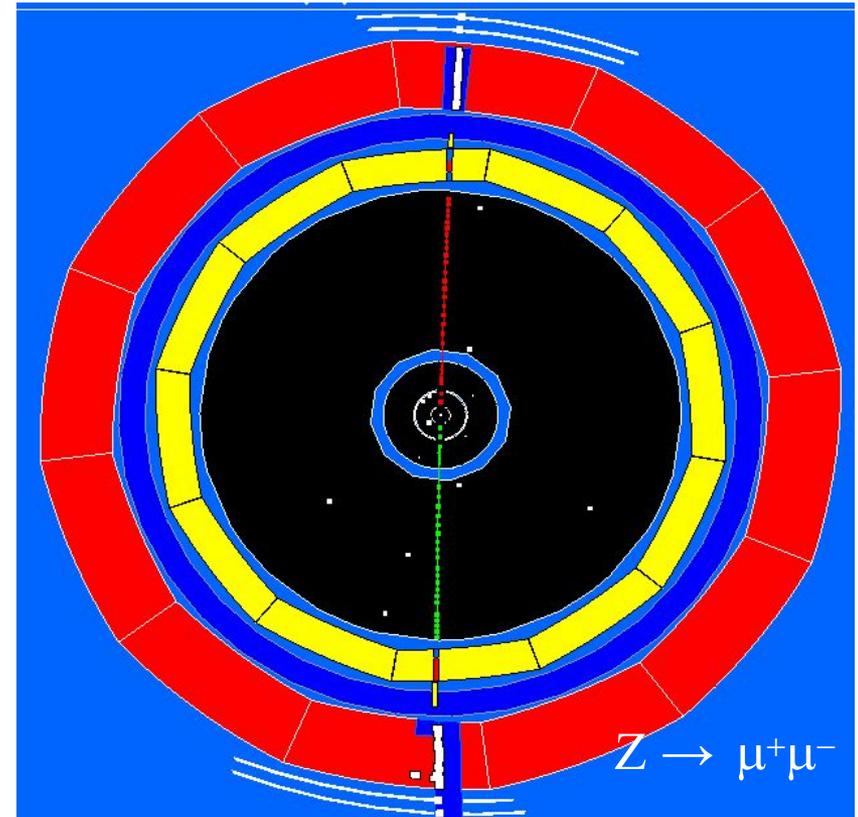
e/γ identification (2)

- Radiation length X_0 = mean distance to reduce energy by $1/e$
eg $X_0 = 1.76$ cm for Fe, so these *electromagnetic* showers are compact
- Distinguish electron/photon by the existence (or not) of a track associated to the shower
- For the electron, E (energy measured in EM calorimeter) and p (momentum from tracker) should be equal: $E/p = 1$
- Not the case for other charged particles



Muons

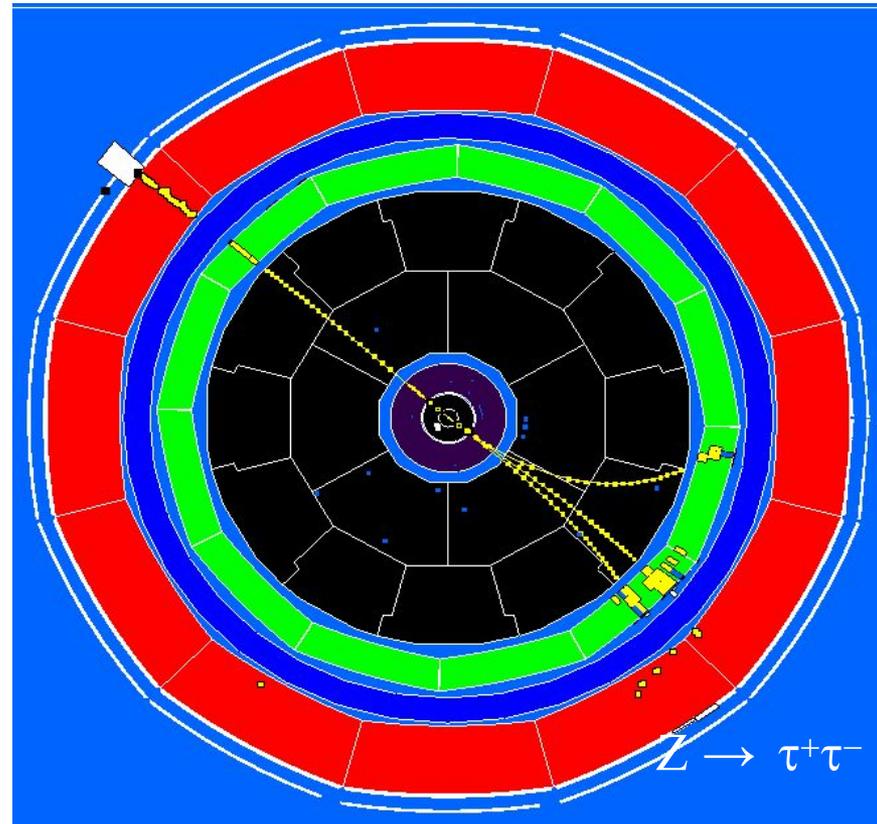
- Muons act like heavier versions of the electron, with mass 105.7 MeV
- They decay to electrons $\mu^- \rightarrow e^- \nu_e \nu_\mu$ with (proper) lifetime $\tau_\mu = 2.2 \mu\text{s}$
- Distance they travel (on average) before decay: $d = \beta\gamma c\tau_\mu$
where velocity $\beta = v/c$
boost $\gamma = E/m = 1/\sqrt{1-\beta^2}$
- So a 10 GeV muon flies ~ 60 km before decay
 \gg detector size
- effectively stable
- Since mass is large, Bremsstrahlung radiation is small, and as a lepton it does not feel the strong interaction



most penetrating charged particle

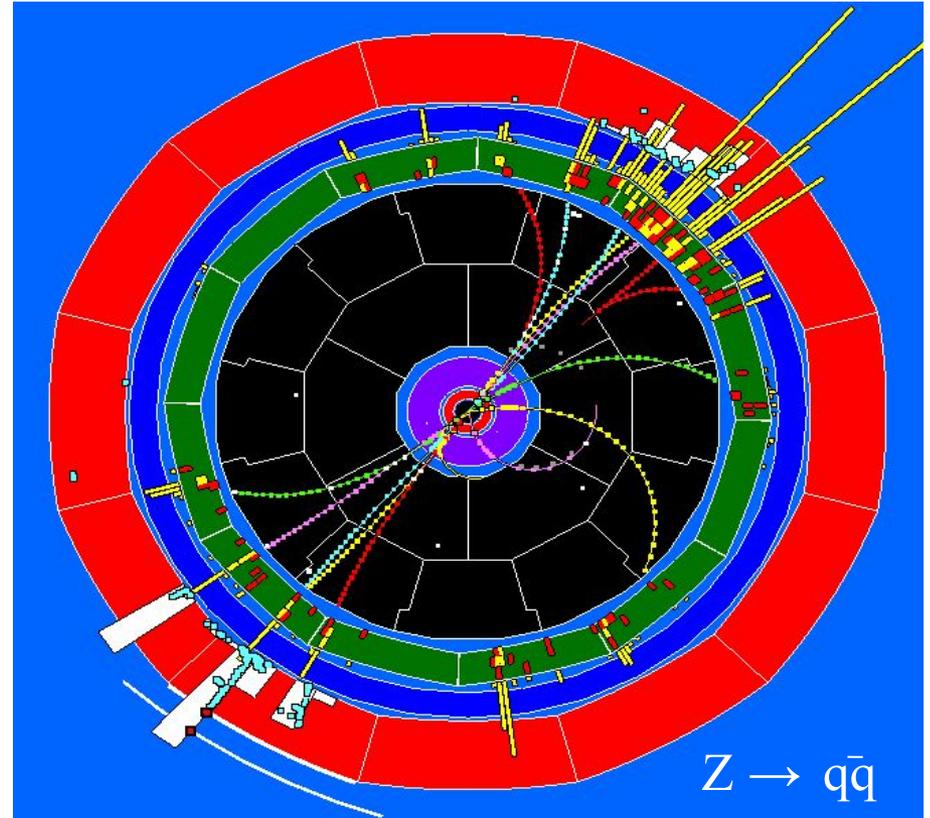
Tau leptons

- Taus are heavier still, $m_\tau = 1.78 \text{ GeV}$
- Heavy enough that can decay to many final states: $\tau^- \rightarrow \mu^- \nu_\mu \nu_\tau$, $\pi^- \nu_\tau$, $\pi^- \pi^0 \nu_\tau$, $\pi^- \pi^+ \pi^- \nu_\tau$, ...
- Lifetime $\tau_\tau = 0.29 \text{ ps}$ ($\text{ps} = 10^{-12} \text{ s}$)
so a 10 GeV tau flies $\sim 0.5 \text{ mm}$
- This is typically too short to be seen directly in the detectors
 - Instead the decay products are seen
- Accurate vertex detectors can detect that they do not come exactly from the interaction point



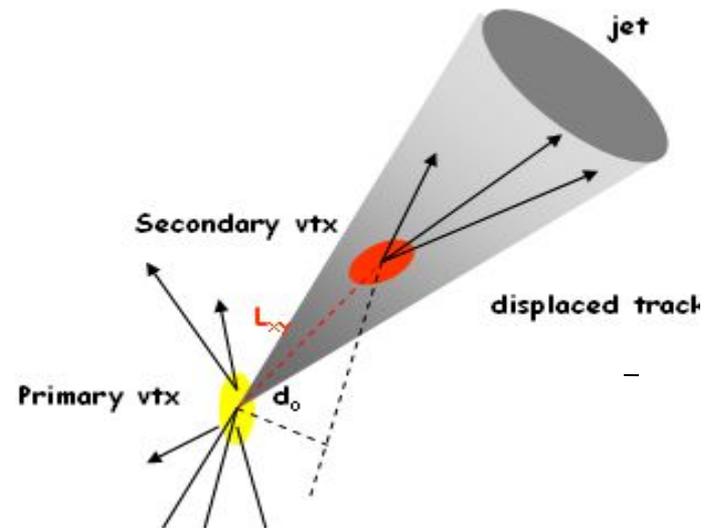
Quarks

- Quarks feel the strong interaction, mediated by gluons
- Not seen in the detector, due to confinement property of QCD
- Instead, they *hadronize* into mesons (qq) or baryons (qqq)
- At high energy $\gg m_q$ initial quark (or gluon) produces a “jet” of hadrons
- Gluon and quark jets are difficult to distinguish: gluon jets tend to be wider, and have a softer particle spectrum



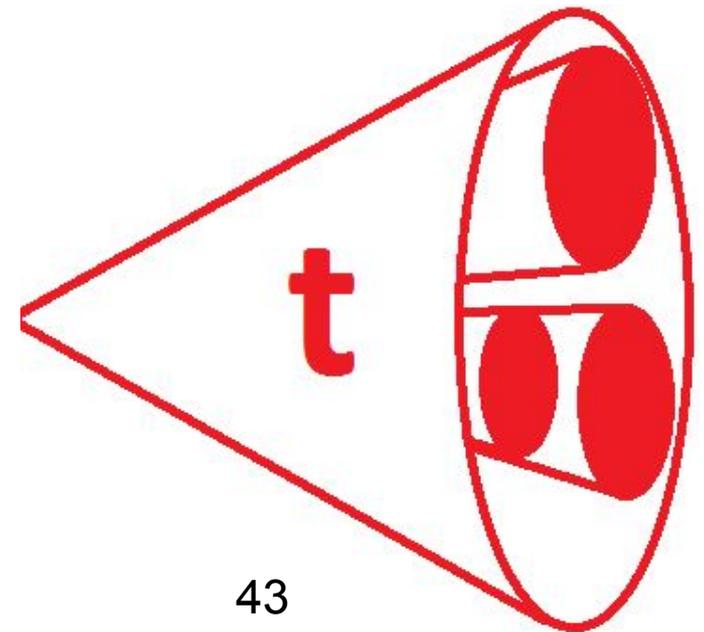
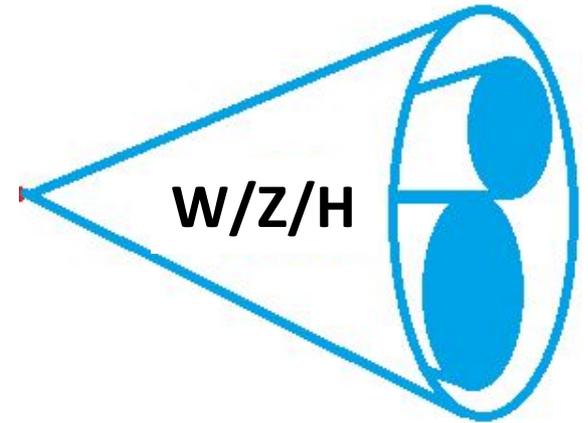
Jet reconstruction

- **Jets are reconstructed by summing up the particles assigned to the jet**
- Typically check within a cone size around the direction of a “seed” particle, or by iteratively adding up pairs of particles that give the lowest invariant mass
- Different quark flavours can be separated (at least statistically) by looking for displaced tracks from b- and c-hadron decays
- The jet properties can be used to approximate the quark or gluon
- Special alg. Used to ID b-jets: life time



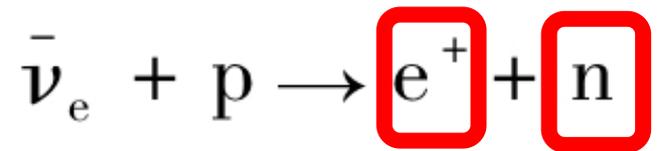
Boost objects

- Highly boosted $W/Z/H/t$...
 - Typical $P_t > \sim 500$ GeV
 - Form a big fat jet
 - With several energy centers
 - Invariant mass
 - b-jets inside
- Very interesting for higher energy colliders ($> \sim 13$ TeV)



Neutrinos

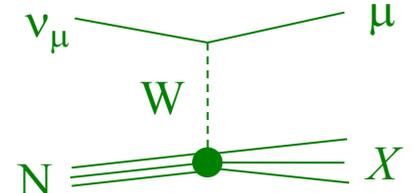
- Neutral (i.e. no track) and only weak interaction
 - pass through matter easily
- Interaction length $\lambda_{\text{int}} = A / (\rho\sigma N_A)$, cross section $\sigma \sim 10^{-38} \text{ cm}^2 \times E [\text{GeV}]$
 - a 10 GeV neutrino can pass through > million km of rock
- Neutrinos are usually detected in HEP experiments through *missing energy* (applying E conservation to rest of the event, usually in transverse plane E_T)
- Dedicated experiments can also identify neutrinos



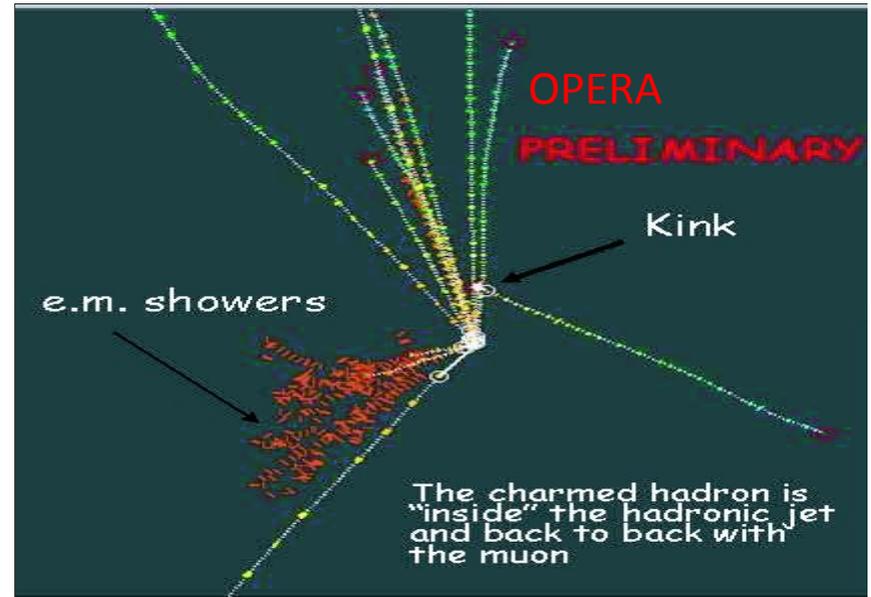
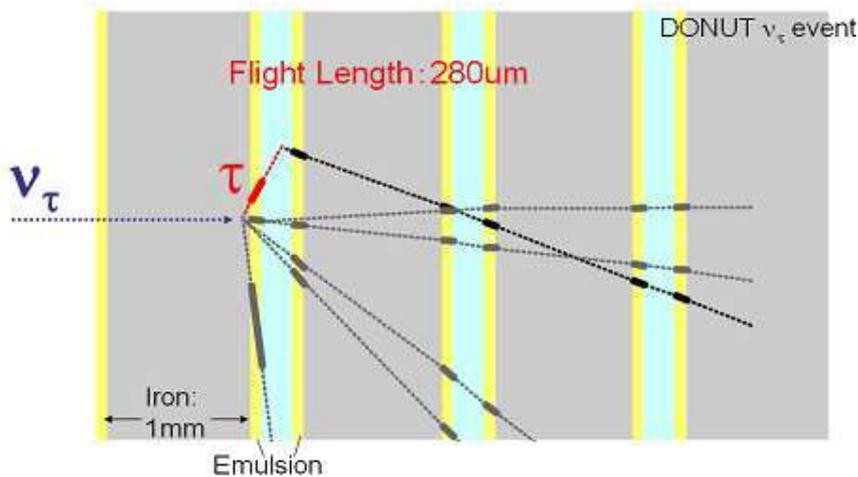
良好的屏蔽去掉其他粒子信号的可能
正负电子湮灭+液闪捕获中子: 光子
光子探测器阵列探测到光子信号

Neutrino flavours

- Can even determine the neutrino flavour (ν_e, ν_μ, ν_τ) from their charged-current interaction: $\nu_\mu N \rightarrow \mu^- X$, etc
- OPERA searches for ν_τ created by neutrino oscillation from a ν_μ beam (sent 730 km from CERN to Italy)
- Tau decay seen as track kink in a high precision emulsion detector, interleaved with lead sheets to provide the high mass of the target



First observation of ν_τ (DONUT)



Summary

- General principle for particle ID
- Example of how CMS identify particles
 - Photons/Electrons/Muons/Taus/jets/boost objects
- Enjoy

