

SM predictions for electroweak precision observables

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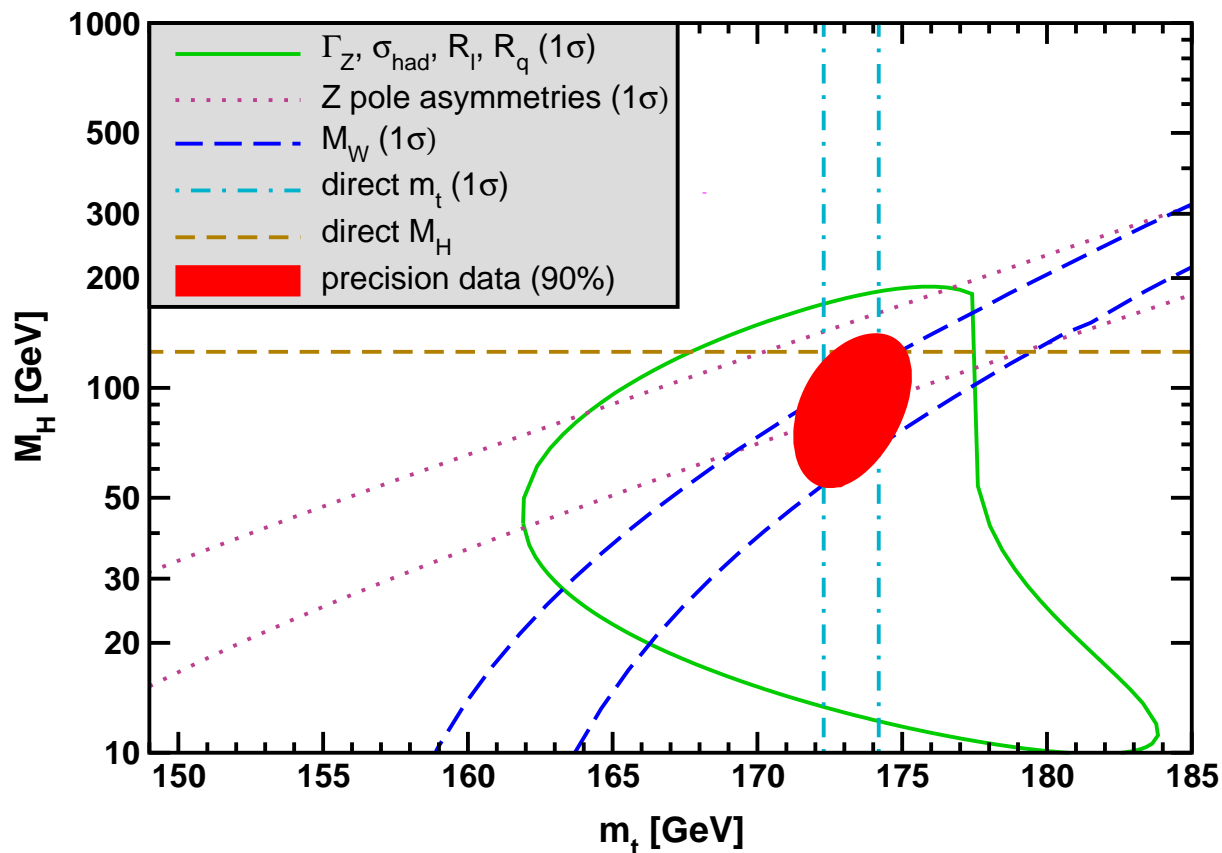
Workshop on Physics at the CEPC,
10-12 August 2015

- 1. Introduction**
- 2. Electroweak precision observables**
- 3. Current status of SM loop results**
- 4. Future projections**

Standard Model after Higgs discovery:

- Good agreement between measured mass and indirect prediction
- Very good agreement over large number of observables

Erler '13



Direct measurements:

$$M_H = 125.6 \pm 0.4 \text{ GeV}$$

$$m_t = 173.24 \pm 0.95 \text{ GeV}$$

Indirect prediction:

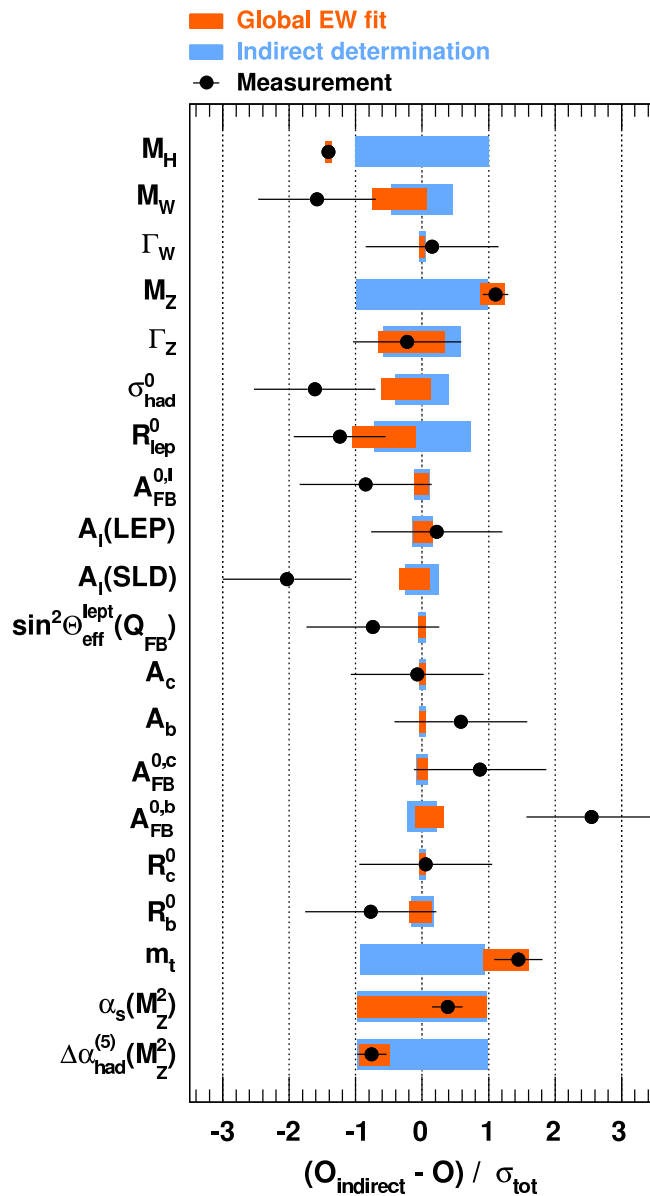
$$M_H = 123.7 \pm 2.3 \text{ GeV}$$

(with LHC BRs)

$$M_H = 89^{+22}_{-18} \text{ GeV}$$

(w/o LHC data)

$$m_t = 177.0 \pm 2.1 \text{ GeV}$$

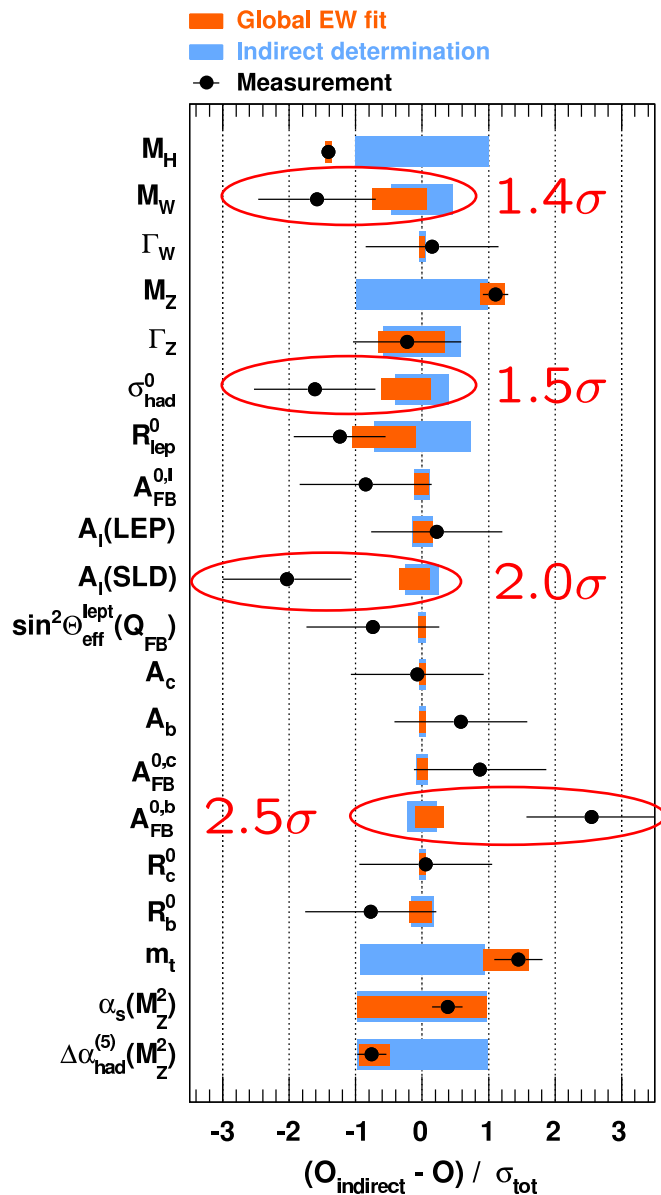


Surprisingly good agreement:

$$\chi^2/\text{d.o.f.} = 18.1/14 \quad (p = 20\%)$$

Most quantities measured with
1%–0.1% precision

GFitter coll. '14



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$$\chi^2/\text{d.o.f.} = 18.1/14 \quad (p = 20\%)$$

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A few interesting deviations:

$$M_W \quad (\sim 1.4\sigma)$$

$$\sigma_{had}^0 \quad (\sim 1.5\sigma)$$

$$A_l(\text{SLD}) \quad (\sim 2\sigma)$$

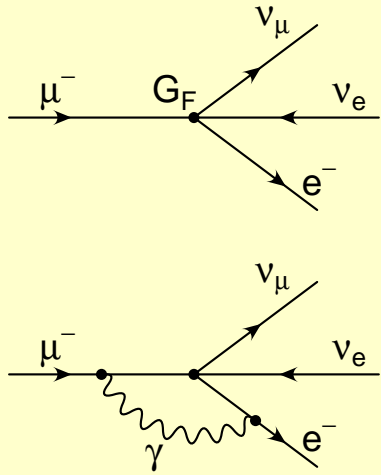
$$A_{FB}^b \quad (\sim 2.5\sigma)$$

$$(g_\mu - 2) \quad (\gtrsim 3\sigma)$$

GFitter coll. '14

W mass

μ decay in Fermi Model



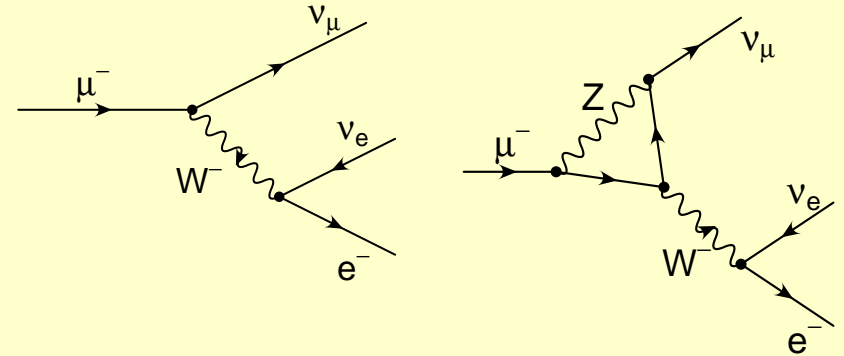
$$\Gamma_\mu = \frac{G_F^2 m_\mu^5}{192\pi^3} F\left(\frac{m_e^2}{m_\mu^2}\right) (1 + \Delta q)$$

← QED corr.
(2-loop)

Ritbergen, Stuart '98

Pak, Czarnecki '08

μ decay in Standard Model



$$\frac{G_F^2}{\sqrt{2}} = \frac{e^2}{8s_w^2 M_W^2} (1 + \Delta r)$$

electroweak corrections

$e^+e^- \rightarrow f\bar{f}$ for $\sqrt{s} \sim m_Z$:

$$\sigma = \mathcal{R}_{\text{ini}} \left[12\pi \frac{\Gamma_{ee}\Gamma_{ff}}{(s - m_Z^2)^2 - m_Z^2\Gamma_Z^2} (1 + \delta X) (1 - \mathcal{P}_e \mathcal{A}_e) + \sigma_{\text{non-res}} \right],$$

$$\Gamma_{ff} = \mathcal{R}_V^f g_{Vf}^2 + \mathcal{R}_A^f g_{Af}^2, \quad \Gamma_Z = \sum_f \Gamma_{ff},$$

$$\mathcal{A}_f = 2 \frac{g_{Vf}/g_{Af}}{1 + (g_{Vf}/g_{Af})^2} = \frac{1 - 4|Q_f| \sin^2 \theta_{\text{eff}}^f}{1 - 4|Q_f| \sin^2 \theta_{\text{eff}}^f + 8(|Q_f| \sin^2 \theta_{\text{eff}}^f)^2}.$$

$$A_{\text{FB}}^f = \frac{3}{4} \mathcal{A}_e \mathcal{A}_f \quad A_{\text{LR}} = \mathcal{A}_e$$

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QED/QCD corrections on ext. fermions

Chetyrkin, Kataev, Tkachov '79

Dine, Sphirstein '79

Celmaster, Gonsalves '80

Gorishnii, Kataev, Larin '88,91

Chetyrkin, Kühn '90

Surguladze, Samuel '91

Kataev '92

Chetyrkin '93

etc...

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$$A_{\text{FB}}^f = \frac{3}{4} \mathcal{A}_e \mathcal{A}_f \quad A_{\text{LR}} = \mathcal{A}_e$$

additional initial-state QED corrections

Kuraev, Fadin '85

Berends, Burgers, v. Neerven '88

Kniehl, Krawczyk, Kühn, Stuart '88

Beenakker, Berends, v. Neerven '89

Bardin et al. '89,91

Montagna, Nicosini, Piccinini '97

etc...

$e^+e^- \rightarrow f\bar{f}$ for $\sqrt{s} \sim m_Z$:

$$\sigma = \mathcal{R}_{\text{ini}} \left[12\pi \frac{\Gamma_{ee}\Gamma_{ff}}{(s - m_Z^2)^2 - m_Z^2\Gamma_Z^2} (1 + \delta X) (1 - \mathcal{P}_e\mathcal{A}_e) + \sigma_{\text{non-res}} \right],$$

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$$A_{\text{FB}}^f = \frac{3}{4} \mathcal{A}_e \mathcal{A}_f \quad A_{\text{LR}} = \mathcal{A}_e$$

electroweak corrections

Correction term first at NNLO:

$$\delta X_{(2)} = -(\text{Im } \Sigma'_{Z(1)})^2 - 2\bar{\Gamma}_Z \bar{M}_Z \text{Im } \Sigma''_{Z(1)}$$

Grassi, Kniehl, Sirlin '01

Freitas '13

Test of running $\overline{\text{MS}}$ weak mixing angle $\sin^2 \bar{\theta}(\mu)$

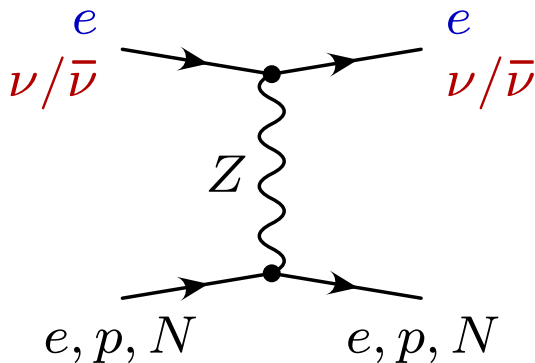
- Polarized ee , ep , ed scattering

$(Q_W(e), Q_W(p), \text{eDIS})$

E158 '05; Qweak '13; JLab Hall A '13

- $\nu N/\bar{\nu}N$ scattering

NuTeV '02

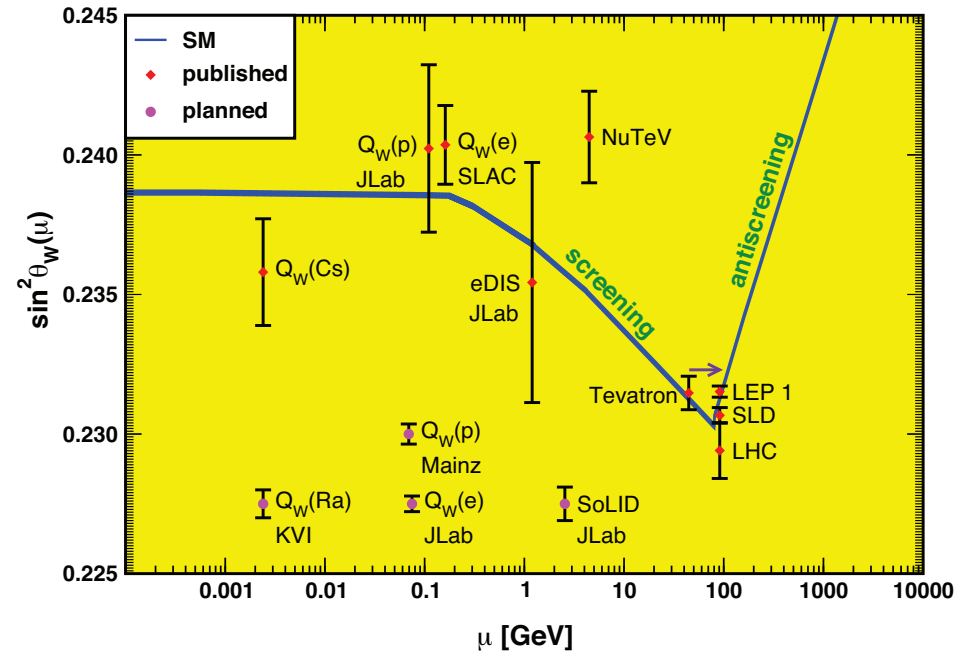


- Atomic parity violation

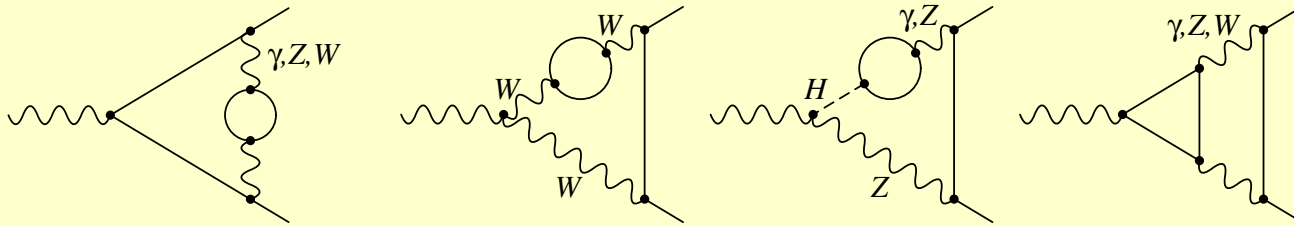
$(Q_W(^{133}\text{Cs}))$ Wood et al. '97

Guéna, Lintz, Bouchiat '05

Erlar '14



Known corrections to Δr , $\sin^2 \theta_{\text{eff}}^f$, g_{Vf} , g_{Af} :



- Complete NNLO corrections (Δr , $\sin^2 \theta_{\text{eff}}^l$) Freitas, Hollik, Walter, Weiglein '00
 Awramik, Czakon '02; Onishchenko, Veretin '02
 Awramik, Czakon, Freitas, Weiglein '04; Awramik, Czakon, Freitas '06
 Hollik, Meier, Uccirati '05,07; Degrandi, Gambino, Giardino '14
- “Fermionic” NNLO corrections (g_{Vf} , g_{Af}) Czarnecki, Kühn '96
 Harlander, Seidensticker, Steinhauser '98
 Freitas '13,14
- Partial 3/4-loop corrections to ρ/T -parameter
 $\mathcal{O}(\alpha_t \alpha_s^2)$, $\mathcal{O}(\alpha_t^2 \alpha_s)$, $\mathcal{O}(\alpha_t \alpha_s^3)$ Chetyrkin, Kühn, Steinhauser '95
 Faisst, Kühn, Seidensticker, Veretin '03
 Boughezal, Tausk, v. d. Bij '05
 Schröder, Steinhauser '05; Chetyrkin et al. '06
 Boughezal, Czakon '06

$$(\alpha_t \equiv \frac{y_t^2}{4\pi})$$

	Experiment	Theory error	Main source
M_W	80.385 ± 0.015 MeV	4 MeV	$\alpha^3, \alpha^2\alpha_s$
Γ_Z	2495.2 ± 2.3 MeV	0.5 MeV	$\alpha_{\text{bos}}^2, \alpha^3, \alpha^2\alpha_s, \alpha\alpha_s^2$
σ_{had}^0	41540 ± 37 pb	6 pb	$\alpha_{\text{bos}}^2, \alpha^3, \alpha^2\alpha_s$
$R_b \equiv \Gamma_Z^b / \Gamma_Z^{\text{had}}$	0.21629 ± 0.00066	0.00015	$\alpha_{\text{bos}}^2, \alpha^3, \alpha^2\alpha_s$
$\sin^2 \theta_{\text{eff}}^l$	0.23153 ± 0.00016	4.5×10^{-5}	$\alpha^3, \alpha^2\alpha_s$

Methods for theory error estimates:

- Parametric factors, *i. e.* factors of α, N_c, N_f, \dots
- Geometric progression, *e. g.* $\frac{\mathcal{O}(\alpha^3)}{\mathcal{O}(\alpha^2)} \sim \frac{\mathcal{O}(\alpha^2)}{\mathcal{O}(\alpha)}$
- Renormalization scale dependence (often underestimates error)
- Renormalization scheme dependence (may underestimate error)

Use of $\overline{\text{MS}}$ renormalization for m_t reduces h.o. QCD corrections of $\mathcal{O}(\alpha_t \alpha_s^n)$:

loops (n+1)	$\Delta\rho_{(n)}^{\overline{\text{MS}}} / \left(\frac{3G_F \overline{m}_t^2}{8\sqrt{2}\pi^2} \right)$	$\Delta\rho_{(n)}^{\text{OS}} / \left(\frac{3G_F m_t^2}{8\sqrt{2}\pi^2} \right)$	
2	$-0.193 \left(\frac{\alpha_s}{\pi} \right)$	$-3.970 \left(\frac{\alpha_s}{\pi} \right)$	Djouadi, Verzegnassi '87 Kniehl '90
3	$-2.860 \left(\frac{\alpha_s}{\pi} \right)^2$	$-14.59 \left(\frac{\alpha_s}{\pi} \right)^2$	Avdeev, Fleischer, et al. '94 Chetyrkin, Kühn, Steinhauser '95
4	$-1.680 \left(\frac{\alpha_s}{\pi} \right)^3$	$-93.15 \left(\frac{\alpha_s}{\pi} \right)^3$	Schröder, Steinhauser '05 Chetyrkin, Faisst, Kühn, et al. '06 Boughezal, Czakon '06

No clear pattern of this kind known for $\mathcal{O}(\alpha^n)$

→ Only few results available that allow direct comparison

e.g. Faisst, Kühn, Seidensticker, Veretin '03

ILC: $\sqrt{s} \approx M_Z$ with 30 fb^{-1} , $\sqrt{s} \approx 2M_W$ with 100 fb^{-1}
FCC-ee: $\sqrt{s} \approx M_Z$ with $4 \times 3000 \text{ fb}^{-1}$, $\sqrt{s} \approx 2M_W$ with $4 \times 3000 \text{ fb}^{-1}$
CEPC: $\sqrt{s} \approx M_Z$ with $100\text{--}1000 \text{ fb}^{-1}$, $\sqrt{s} \approx 2M_W$ with 100 fb^{-1}

	Current exp.	ILC	FCC-ee	CEPC	Current perturb.
M_W [MeV]	15	3–5	~ 1	3–5	4
Γ_Z [MeV]	2.3	~ 1	~ 0.1	~ 0.5	0.5
R_b [10^{-5}]	66	15	$\lesssim 5$	17	15
$\sin^2 \theta_{\text{eff}}^l$ [10^{-5}]	16	1.3	0.3	~ 3	4.5

→ Existing theoretical calculations adequate for LEP/SLC/LHC,
but not future e^+e^- machines!

	ILC	CEPC	perturb. error with 3-loop [†]	Param. error ILC*	Param. error CEPC**
M_W [MeV]	3–5	3–5	1	2.6	2.1
Γ_Z [MeV]	~ 1	~ 0.5	$\lesssim 0.2$	0.5	0.09
R_b [10^{-5}]	15	17	5–10	< 1	< 1
$\sin^2 \theta_{\text{eff}}^\ell$ [10^{-5}]	1.3	3	1.5	2	2

[†] **Theory scenario:** $\mathcal{O}(\alpha\alpha_S^2)$, $\mathcal{O}(N_f\alpha^2\alpha_S)$, $\mathcal{O}(N_f^2\alpha^2\alpha_S)$
 (N_f^n = at least n closed fermion loops)

Parametric inputs:

* **ILC:** $\delta m_t = 100$ MeV, $\delta\alpha_S = 0.001$, $\delta M_Z = 2.1$ MeV

****CEPC:** $\delta m_t = 600$ MeV, $\delta\alpha_S = 0.0001$, $\delta M_Z = 0.5$ MeV

also: $\delta(\Delta\alpha) = 5 \times 10^{-5}$

■ Subtraction of QED radiation contributions

→ Known to $\mathcal{O}(\alpha^2)$, $\mathcal{O}(\alpha^3 L^3)$ for **ISR**,
 $\mathcal{O}(\alpha^2)$ for **FSR** and $\mathcal{O}(\alpha^2 L^2)$ for **A_{FB}**

$$(L = \log \frac{s}{m_e^2})$$

Berends, Burgers, v.Neerven '88

Kniehl, Krawczyk, Kühn, Stuart '88

Beenakker, Berends, v.Neerven '89

Skrzypek '92; Montagna, Nicrosini, Piccinini '97

→ $\mathcal{O}(0.1\%)$ uncertainty on σ_Z , A_{FB}

→ Improvement needed for ILC/FCC-ee

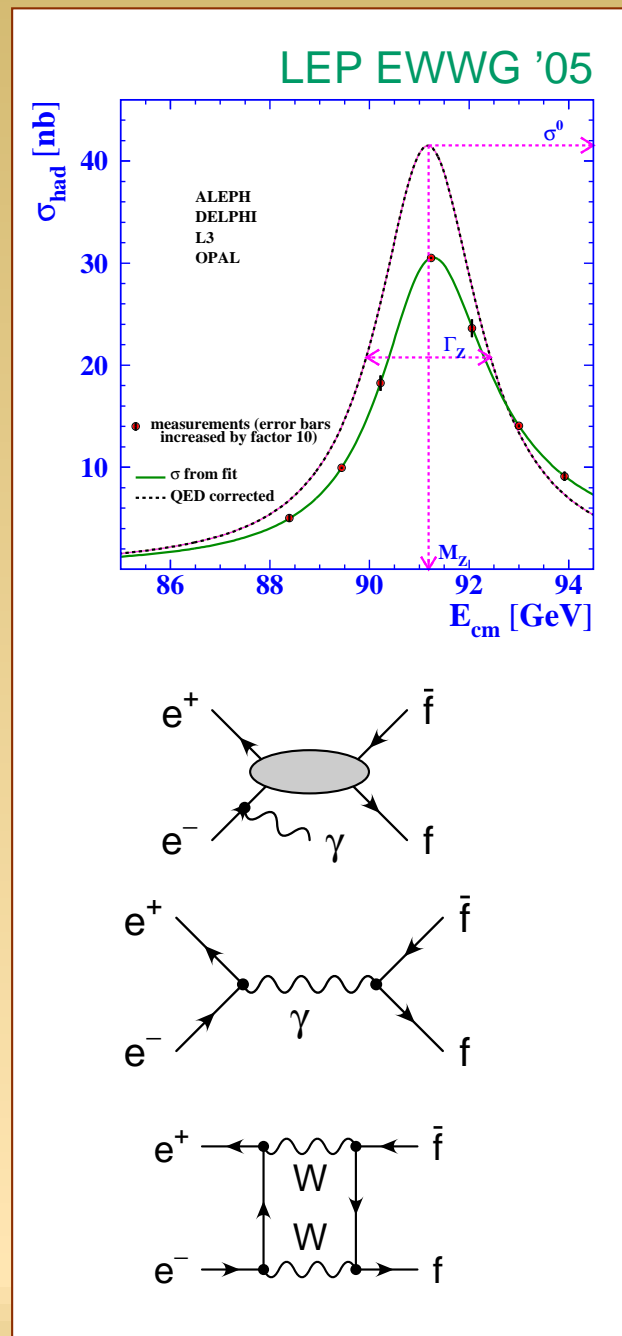
■ Subtraction of non-resonant γ -exchange, γ -Z interf., box contributions, Bhabha scattering

see, e.g., Bardin, Grünewald, Passarino '99

→ $\mathcal{O}(0.01\%)$ uncertainty within SM

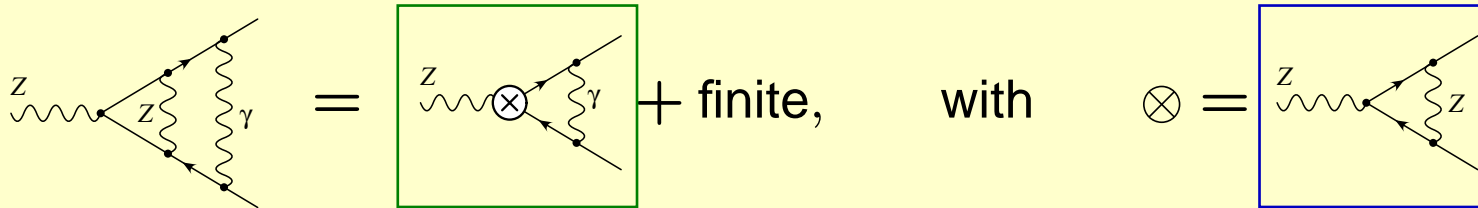
(improvements may be needed)

→ Sensitivity to some NP beyond EWPO

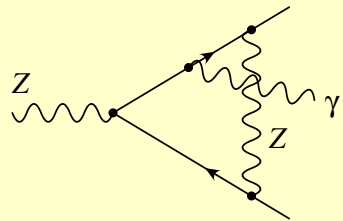


Factorization of massive and QED/QCD FSR:

$$\Gamma_{ff} \propto \mathcal{R}_V^f g_{Vf}^2 + \mathcal{R}_A^f g_{Af}^2$$



Additional non-factorizable contributions, e.g.



→ Known at $\mathcal{O}(\alpha\alpha_s)$ Czarnecki, Kühn '96
 Harlander, Seidensticker, Steinhauser '98

→ Currently not known at $\mathcal{O}(\alpha^2)$ and beyond

→ $\mathcal{O}(0.01\%)$ uncertainty on Γ_Z, σ_Z , maybe larger for A_b
 (improvements may be needed)

Full SM corrections at >2 -loop:

- Large number of diagrams and tensor integrals, $\mathcal{O}(1000) - \mathcal{O}(10000)$
 - Use of computer algebra tools

- Many different scales (masses and ext. momenta)
 - In general not possible analytically
 - Numerical methods must be automizable, stable, fastly converging
 - Need procedure for isolating divergent pieces

- Useful for diagrams with up to two scales
(*e. g.* M_W & m_t or M_W & M_Z)
- Reduce to master integrals with integration-by-parts and Lorentz-invariance identities
Chetyrkin, Tkachov '81; Gehrmann, Remiddi '00; Laporta '00; ...
- Evaluate master integrals with differential equations or Mellin-Barnes representations
Kotikov '91; Remiddi '97; Smirnov '00,01; ...

Current status:

Single-scale problems: $Z f \bar{f}$ QED/QCD vertex corrections up to 4-loop

Gorishnii, Kataev, Larin '88,91; Chetyrkin, Kühn, Kwiatkowski '96
Baikov, Chetyrkin, Kühn, Rittinger '12

Two-scale problems: $Z f \bar{f}$ electroweak 2-loop vertex diagrams with $m_f = 0$

Awramik, Czakon, Freitas, Weiglein '04

Extendability: Possible, but much work needed

- Exploit large mass ratios, *e. g.* $M_Z^2/m_t^2 \approx 1/4$
- Fast numerical evaluation

Current status:

Two-scale problems: $\mathcal{O}(\alpha\alpha_s^n)$ corr. to $\Delta\rho$, Δr , ...

→ Several expansion terms up to 3-loop, leading term up to 4-loop

Djouadi, Verzegnassi '87; Bardin, Chizhov '88

Chetyrkin, Kühn, Steinhauser '95

Faisst, Kühn, Seidensticker, Veretin '03

...

Up to three-scale problems: $Zf\bar{f}$ ew. 2-loop vertex corrections

Barbieri et al. '92,93; Fleischer, Tarasov, Jegerlehner '93,95

Degrassi, Gambino, Sirlin '97; Awramik, Czakon, Freitas, Weiglein '04

Extendability: Promising, mostly limited by computing/algorithmic power

General form of Feynman integral:

$$I = \int_0^1 dx_1 \dots dx_n \delta(1 - \sum_i x_i) \frac{N(x_i)}{D(x_i)^{r+\varepsilon}}$$

→ Can be integrated numerically (if finite)

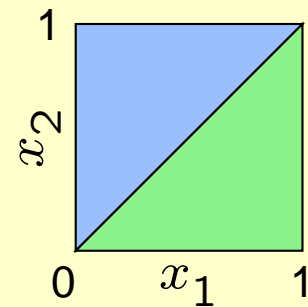
Alternatives: Integration in momentum space, Mellin-Barnes space

Treatment of divergencies:

- **Sector decomposition:** Sub-divide integration space such that divergent terms factorize

Binoth, Heinrich '00,03

- **Subtraction terms:** Remove divergencies with simple terms that can be integrated analytically



Nagy, Soper '03

Becker, Reuschle, Weinzierl '10; Freitas '12

- Automizable, but computing intensive
- Internal thresholds reduce numerical convergence

Current status:

Several 2-loop applications with many scales

Anastasiou et al., Petriello et al., Borowka et al.,

Individual 3-loop integrals

Extendability: Likely, but more work needed

- **Current SM predictions** for electroweak precision observables under good control (compared to experimental uncertainties)
- **LHC** will provide independent results for $\sin^2 \theta_{\text{eff}}$ and M_W , but overall precision not substantially improved
- **ILC/CEPC/FCC-ee** with $\sqrt{s} \sim M_Z$ will reduce exp. error of some EWPO by $\mathcal{O}(10)$
 - 3-loop (and maybe some 4-loop) corrections needed!
- **Asymptotic expansion and numerical integration techniques** are promising but more work needed
- **Open questions** in evaluation of theory errors, resummation and optimal choice of inputs

Backup slides

■ Geometric perturbative series

$$\alpha_t = \alpha m_t^2$$

$$\mathcal{O}(\alpha^3) - \mathcal{O}(\alpha_t^3) \sim \frac{\mathcal{O}(\alpha^2) - \mathcal{O}(\alpha_t^2)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha^2) \sim 0.26 \text{ MeV}$$

$$\mathcal{O}(\alpha^2 \alpha_s) - \mathcal{O}(\alpha_t^2 \alpha_s) \sim \frac{\mathcal{O}(\alpha^2) - \mathcal{O}(\alpha_t^2)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha \alpha_s) \sim 0.30 \text{ MeV}$$

$$\mathcal{O}(\alpha \alpha_s^2) - \mathcal{O}(\alpha_t \alpha_s^2) \sim \frac{\mathcal{O}(\alpha \alpha_s) - \mathcal{O}(\alpha_t \alpha_s)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha \alpha_s) \sim 0.23 \text{ MeV}$$

$$\mathcal{O}(\alpha \alpha_s^3) - \mathcal{O}(\alpha_t \alpha_s^3) \sim \frac{\mathcal{O}(\alpha \alpha_s) - \mathcal{O}(\alpha_t \alpha_s)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha \alpha_s^2) \sim 0.035 \text{ MeV}$$

$$\mathcal{O}(\alpha_{\text{bos}}^2) \sim \mathcal{O}(\alpha_{\text{bos}})^2 \sim 0.1 \text{ MeV}$$

■ Parametric prefactors: $\mathcal{O}(\alpha_{\text{bos}}^2) \sim \Gamma_Z \alpha^2 \sim 0.1 \text{ MeV}$

$$\mathcal{O}(\alpha \alpha_s^2) - \mathcal{O}(\alpha_t \alpha_s^2) \sim \frac{\alpha n_{|q}}{\pi} \alpha_s^2 \sim 0.29 \text{ MeV}$$

Total: $\delta \Gamma_Z \approx 0.5 \text{ MeV}$

■ Renormalization scheme dependence:

- a) Uncertainty of $\mathcal{O}(\alpha^2)$ corrections beyond leading $\alpha^2 m_t^4$ and $\alpha^2 m_t^2$ from comparison of $\overline{\text{MS}}$ and OS schemes: Degrassi, Gambino, Sirlin '96

$$\delta M_W \sim 2 \text{ MeV} \quad (\text{for } M_H \sim 100 \text{ GeV})$$

Actual remaining $\mathcal{O}(\alpha^2)$ corrections: Freitas, Hollik, Walter, Weiglein '00

$$\delta M_W \sim 3 \text{ MeV} \quad (\text{for } M_H \sim 100 \text{ GeV})$$

- b) Estimate of missing $\mathcal{O}(\alpha^3)$ corrections from comparison of $\overline{\text{MS}}$ and OS results:

Awramik, Czakon, Freitas, Weiglein '03

Degrassi, Gambino, Giardino '14

$$\delta M_W \sim 4 \dots 5 \text{ MeV} \quad (\text{after accounting for } \mathcal{O}(\alpha_t \alpha_s^3) \text{ corrections})$$

→ Saturates previous δM_W estimate!

Note: Differences in (implicitly) resummed higher-order contributions

Parametrization of perturbation series: α vs. G_F ?

G_F can resum some leading one-loop terms

$$\Delta\alpha \equiv 1 - \frac{\alpha(0)}{\alpha(M_Z)} \approx 0.059 \qquad \Delta\rho = \frac{3\alpha}{16\pi s^2 c^2} \frac{m_t^2}{M_Z^2}$$

But: Strong cancellations between $\Delta\alpha$ and $\Delta\rho$ terms beyond one-loop:

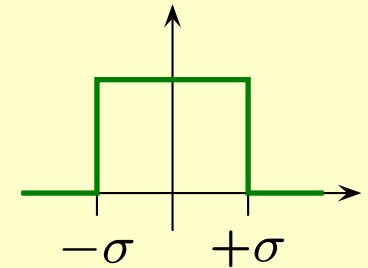
$$\begin{aligned} \Delta r_{\text{res}}^{(3)} &= (\Delta\alpha)^3 - 3(\Delta\alpha)^2 \left(\frac{c^2}{s^2} \Delta\rho\right) + 6(\Delta\alpha) \left(\frac{c^2}{s^2} \Delta\rho\right)^2 - 5 \left(\frac{c^2}{s^2} \Delta\rho\right)^3 \\ &\approx (2.05 \quad -3.40 \quad +3.74 \quad -1.72) \times 10^{-4} \\ &= 0.68 \times 10^{-4} \end{aligned}$$

→ Not *the* numerically leading contribution anymore

- Add theory errors from each source **linearly**:

Idea: each value within error range is equally likely

→ Use flat prior in global fits



- Add theory errors from each source **quadratically**:

Idea: different error sources are uncorrelated

→ Use Gaussian prior in global fits (central limit theorem)

