## Automated EW Corrections with MG5_aMC@NLO (and Top Quark)

mainly based on: arXiv:1407.0823, arXiv:1504.03446, arXiv:1507.05640 and preliminary work for 100 TeV processes


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IHEP, Beijing, China
Miniworkshop

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20-11-2015
$$

## OUTLINE

## Automation of the EW corrections in MG5_aMC@NLO

- Status and technical aspects


## NLO QCD and EW corrections to $t \bar{t} V$

- Completely automated results at $8,13,100 \mathrm{TeV}$ and in a boosted regime


## EW corrections at 100 TeV

- Sudakov logs at high scales
- Effects from the Photon PDF


## Conclusions and Outlook

## Automation of NLO corrections in Madgraph5_aMC@NLO

## What do we mean with automation of EW corrections?

The possibility of calculating QCD and EW corrections for SM processes (matched to shower effects) with a process-independent approach.

```
generate process [QCD]
output process_QCD
```

```
generate process [QCD EW]
output process_QCD_EW
```

The automation of NLO QCD has been achieved, but we need higher precision to match the experimental accuracy at the LHC and future colliders.

- NNLO QCD automation is out of our theoretical capabilities at the moment.
- NLO EW corrections are of the same order $\left(\alpha_{s}^{2} \sim \alpha\right)$, the Sudakov logarithms can enhance their size. NLO QCD and EW corrections can be automated.


## Automation of NLO corrections in Madgraph5_aMC@NLO

The complete automation has already been achieved for QCD.


## Automation of NLO corrections in Madgraph5_aMC@NLO

The complete automation for $\mathbf{Q C D}+\mathbf{E W}$ is in progress.


## Amplitudes and matrix elements

| NLO UFO models: | - SM-alpha $(\mathrm{mZ})$ | $(\mathrm{EW}+\mathrm{QCD}$, Weak+QCD) |
| :--- | :--- | :--- |
| (UV CT, R2) | $-\mathrm{SM}-\mathrm{G} \mu$ | $(\mathrm{EW}+\mathrm{QCD}$, Weak+QCD) |

Weak $=$ EW without photonics corrections (to be used when gauge invariant). The matrix element calculation is completely automated. Example: $t \bar{t} V$ NLO orders of $t \bar{t} V$



## Amplitudes and matrix elements

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## Processes with only final-state massive particles

The generation of EW-QCD loops, real emission of gluons, quarks and photons is completely automated.

FKS IR counterterms completely automated. Also for photons in the initial state.

Heavy Boson Radiation (HBR)

$$
\begin{aligned}
& p p \rightarrow t \bar{t} H+V \\
& V=H, W, Z
\end{aligned}
$$

Formally of order $\alpha_{S}^{2} \alpha^{2}$

## Pheno studies

NLO purely Weak and QCD corrections to $t \bar{t} H$ production have been produced "assembling by hand" the FKS counterterms.

Frixione, Hirschi, DP, Shao, Zaro '14

Now, for the complete NLO QCD and EW corrections, with photons in the initial state, we need to type:

```
define p = p b b~ a
generate p p > t t~ h [QCD QED]
output ttbarh_QCD_QED
```



In this talk I present results for NLO QCD and EW corrections to
$t \bar{t} V \quad V=H, W, Z$

## Structure of NLO EW-QCD corrections



## Structure of NLO EW-QCD corrections



## Structure of NLO EW-QCD corrections



LO

NLO


## Structure of NLO EW-QCD corrections



LO

NLO


## Structure of NLO EW-QCD corrections



LO

NLO


## Structure of NLO EW-QCD corrections



## $t \bar{t} V$ production: numerical results

Alpha(mZ)-scheme, NNPDF2.3_QED, $\quad \mu=\frac{H_{T}}{2}, \quad \frac{1}{2} \mu \leq \mu_{R}, \mu_{F} \leq 2 \mu$

## Contributions

$\operatorname{HBR}\left(p p \rightarrow t \bar{t} V+V^{\prime}\right)$ is of the same order of NLO EW.

The Photon PDF (with large uncertainties) enters in LO EW and NLO EW.


## NLO QCD NLO EW

```
define p = p b b~ a
generate p p > t t~ h [QCD QED]
output ttbarh_QCD_QED
```


## $t \bar{t} V$ production: numerical results

Alpha(mZ)-scheme, $\quad$ NNPDF2.3_QED, $\quad \mu=\frac{H_{T}}{2}, \quad \frac{1}{2} \mu \leq \mu_{R}, \mu_{F} \leq 2 \mu$

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## NLO QCD NLO EW

## Boosted regime

$p_{T}(t) \geq 200 \mathrm{GeV}, \quad p_{T}(\bar{t}) \geq 200 \mathrm{GeV}, \quad p_{T}(H) \geq 200 \mathrm{GeV}$
S/B increases for boosted tops and Higgs.
Plehn, Salam, Spannowsky '10
Sudakov logs are relevant in these regions!


## Numerical results



## Scale variation

(NLO QCD+EW) PDF var.

Frixione, Hirschi, DP, Shao, Zaro '15

## Numerical results

| $10^{-1}$ | Itz production at the 13 TeV LHC | $t \bar{t} Z: \delta(\%)$ | 8 TeV | 13 TeV | 100 TeV |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | NLO QCD | $43.2{ }_{-15.9}^{+12.8}$ | $45.9_{-15.5}^{+13.2}\left(40.2_{-15.0}^{+11.1}\right)$ | $50.4_{-10.9}^{+11.4}$ |
|  |  | LO EW | $0.5 \pm 0.9$ | $0.0 \pm 0.7(2.1 \pm 1.6)$ | $-1.1 \pm 0.2$ |
|  |  | LO EW no $\gamma$ | $-0.8 \pm 0.1$ | $-1.1 \pm 0.0(-0.3 \pm 0.0)$ | $-1.6 \pm 0.0$ |
|  |  | NLO EW | $-3.3 \pm 0.3$ | $-3.8 \pm 0.2(-11.1 \pm 0.5)$ | $-5.2 \pm 0.1$ |
| $10^{-2}$ |  | - NLO EW no $\gamma$ | $-3.7 \pm 0.1$ | $-4.1 \pm 0.1(-11.5 \pm 0.3)$ | $-5.4 \pm 0.0$ |
|  |  | 免 HBR | 0.95 | 0.96 (2.13) | 0.85 |

(Boosted regime in brackets)

## Scale variation

(NLO QCD+EW) PDF var.

## $t \bar{t} Z$

Frixione, Hirschi, DP, Shao, Zaro '15

## Numerical results



| $t \bar{t} W^{+}: \delta(\%)$ | 8 TeV |  | 13 TeV |  | 100 TeV |
| :---: | :---: | :---: | :---: | :---: | :---: |
| NLO QCD | $40.8_{-12.3}^{+11.2}$ | $50.1_{-13.5}^{+14.2}\left(59.7_{-17.7}^{+18.9}\right)$ | $156.4_{-35.0}^{+38.3}$ |  |  |
| LO EW | 0 | 0 | 0 |  |  |
| LO EW no $\gamma$ | 0 | 0 | 0 |  |  |
| NLO EW | $-6.9 \pm 0.2$ | $-7.7 \pm 0.2(-19.2 \pm 0.7)$ | $-9.3 \pm 0.2$ |  |  |
| NLO EW no $\gamma$ | $-7.1 \pm 0.2$ | $-8.0 \pm 0.2(-20.0 \pm 0.5)$ | $-9.6 \pm 0.1$ |  |  |
| HBR | 2.41 | $3.88(7.41)$ | 21.52 |  |  |

(Boosted regime in brackets)
Scale variation
(NLO QCD+EW) PDF var.

## Numerical results



Frixione, Hirschi, DP, Shao, Zaro '15

## 100 TeV

We want to estimate the typical size of EW corrections at 100 TeV by looking at similar distributions in different "simple processes":
transverse momenta and invariant masses in:
WZ, HW, t̄t
At the same time, we want to also estimate the impact of the photon PDF (central value and errors) at 100 TeV .

## Disclaimer

We want to raise possible issues related to NLO EW corrections at 100 TeV . No final answers or definitive statements will be given!

WZ

## WZ

## differential plots


photon PDF NO

photon PDF YES

## WZ

## differential plots



photon PDF NO

# WZ what's going on? 



$$
\begin{gathered}
d \sigma^{d g \rightarrow W^{-} Z_{u}}=c_{W Z}^{d} d \sigma_{L}^{d g \rightarrow Z d} \frac{\alpha}{2 \pi} \log ^{2}\left[\frac{\left(p_{T}^{Z}\right)^{2}}{M_{W}^{2}}\right] \\
\mathrm{Zj}+ \\
\text { soft and collinear W } \\
d \sigma^{d \gamma \rightarrow W^{-}-Z u}=\frac{c_{L, d}^{2} c_{W Z}^{d}}{a_{W}^{2}} d \sigma_{L}^{d \gamma \rightarrow W^{-} u} \frac{\alpha}{2 \pi} \log ^{2}\left[\frac{\left(p_{T}^{W-}\right)^{2}}{M_{Z}^{2}}\right]
\end{gathered}
$$

See also Baglio, Ninh, Weber '13

- The large growth for high pt in NLO QCD corrections and photon-quark in NLO EW have similar origins: the same cuts may suppress both effects.
- The photon couples to the W, originating new t-channel configurations that enhance the relative size of photon-quark contributions in NLO EW. NLO QCD corrections do not exhibit similar features.


## WZ

cumulative plots

photon PDF NO

photon PDF YES

## WZ <br> cumulative plots



HW

## HW

## differential plots


photon PDF NO

photon PDF YES

## HW

## differential plots


photon PDF NO

photon PDF YES

## HW <br> cumulative plots




## HW <br> cumulative plots




# HW what's going on? 



At variance with VV processes, $\log \wedge 2$ enhancements in pt distributions are not present at NLO QCD or EW; there is not photon-quark or gluon-quark to Higgs + jet at tree level that can factorize W or Z soft and collinear emission.

However, HW is produced only via s-channel at LO, while the photonquark initial state introduces $t$-channels via W and photon interaction, which are much less suppressed for high $\mathrm{m}(\mathrm{HW})$ and lead to huge (100-600) K-factors.

Heavy Boson Radiation (HBR) is of the same perturbative order and the same numerical size of "genuine" NLO EW corrections.

## HW

## what if additional cuts are applied?



Cuts on H and W pseudo rapidities strongly reduce the photon-quark contribution at high $\mathrm{m}(\mathrm{HW})$, without affecting LO and NLO QCD results. Besides HW, cuts can in general strongly affect size of radiative corrections: Which cuts and on which particles at 100 TeV ?

## top-quark pair production

$t \bar{t}$
cumulative plots

photon PDF NO

photon PDF YES
$t \bar{t}$
cumulative plots

photon PDF NO

photon PDF YES

## CONCLUSION

The automation of NLO EW and QCD corrections in MadGraph5_aMC@NLO is in progress. NLO QCD and EW corrections to the $t \bar{t} V$ processes have been calculated in a completely automated approach. It is in general possible with massive final states.

NLO EW corrections at high energies can involve large negative contributions from Sudakov logs, which point to the necessity of going beyond NLO and resumming the logs.
The contribution from photon-induced processes can be huge and with very large uncertainties. However, this does not necessarily apply to all processes and distributions. Cuts may be essential for a realistic prediction.

## OUTLOOK

- Complete the automation of EW+QCD corrections for all processes
- Match NLO EW corrections to shower effects
- More pheno studies, playing with the tools


## EXTRA SLIDES

## Energy dependence




Maltoni, DP, Tsinikos '15

## Distributions: representative results at fixed order




Maltoni, DP, Tsinikos '15

## Distributions: representative results at fixed order




Maltoni, DP, Tsinikos '15

## Distributions: representative results at fixed order



Central Asymmetries

| $13 \mathrm{TeV} A_{c}[\%]$ | $t \bar{t} H$ | $t \bar{t} Z$ | $t \bar{t} W^{ \pm}$ | $t \bar{t} \gamma$ |
| :---: | :---: | :---: | :---: | :---: |
| LO | - | $-0.12_{-0.01}^{+0.01}{ }_{-0.02}^{+0.01} \pm 0.03$ | - | $-3.93_{-0.23}^{+0.26}{ }_{-0.11}^{+0.14} \pm 0.03$ |
| NLO | $1.00_{-0.20}^{+0.30}{ }_{-0.04}^{+0.06} \pm 0.02$ | $0.85_{-0.17}^{+0.25}{ }_{-0.05}^{+0.06} \pm 0.03$ | $2.90_{-0.47}^{+0.67}$ | ${ }_{-0.07}^{+0.06} \pm 0.07$ |

## Why do we care about photons in the proton? 2 representative examples:

High-mass Drell-Yan, $\mathrm{Z} \rightarrow \mathrm{e}^{+} \mathrm{e}^{-}$

| $t \bar{t}$ | Process | $\sigma_{\text {tot }}$ without cuts [pb] |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  |  | Born | correction |  |
|  | $u \bar{u}$ | 34.25 | -1.41 | NLO QED |
|  | $d \bar{d}$ | 21.61 | -0.228 |  |
|  | $s \bar{s}$ | 4.682 | -0.0410 |  |
|  | $c \bar{c}$ | 2.075 | -0.0762 |  |
|  | - gg | 407.8 | 2.08 |  |
|  | $g \gamma$ |  | 4.45 |  |
|  | - pp | 470.4 | 4.78 |  |

Integrated hadronic cross section for $t \bar{t}$ production at the LHC, at NLO QED

Hollik, Kollar '07
MRST2004QED



Carrazza '14
NNPDF2.3QED
LO

NLO


Set-up and photon-PDF perturbative orders
$G_{\mu}$ scheme, $\quad \quad \mathrm{NNPDF} 2.3 \_\mathrm{QED}, \quad \mu=\frac{H_{T}}{2}, \quad \frac{1}{2} \mu \leq \mu_{R}, \mu_{F} \leq 2 \mu$
WW

## Rapidity distributions: unboosted vs. boosted



## Transverse momentum distributions: unboosted vs. boosted



## Comparison between different schemes

$$
m_{W}=80.385 \mathrm{GeV}, \quad m_{Z}=91.188 \mathrm{GeV}
$$

$$
\alpha\left(m_{Z}\right) \text { scheme } \quad \frac{1}{\alpha\left(m_{Z}\right)}=128.93
$$

$$
G_{\mu} \text { scheme } G_{\mu}=1.16639 \cdot 10^{-5} \quad \longrightarrow \frac{1}{\alpha}=132.23
$$



Table 11: Comparison between results in the $\alpha\left(m_{Z}\right)$ and $G_{\mu}$ scheme, at 13 TeV .

## Why Weak corrections to $t \bar{t} H$ production?

We calculated NLO corrections of mixed QCD-Weak origin, ignoring QED effects. We compared them to NLO QCD corrections.

## Phenomenology motivations

Electroweak corrections are in general small. However, the Sudakov logarithms $\alpha_{W} \ln ^{2} s / M_{W}^{2}$ can enhance their size. They originate only from Weak corrections

The cross section of $t \bar{t} H$ depends directly on $\lambda_{t \bar{t} H}^{2}$. At NLO, only Weak corrections introduce a dependence on other Higgs couplings.

## Automation of NLO corrections

Without QED (photons), the structure of IR singularities is simpler $t \bar{t} H$ was the first pheno study of EW corrections in the MG5_aMC@NLO framework.

## Numerical results weak corrections

## Inclusive rates

(Boosted regime in brackets)
NLO corrections

| $\delta_{\mathrm{NLO}}(\%)$ | 8 TeV | 13 TeV | 100 TeV |
| :--- | :---: | :---: | :---: |
| QCD | $+25.6_{-11.8}^{+6.2}\left(+19.6_{-11.0}^{+3.7}\right)$ | $+29.3_{-11.6}^{+7.4}\left(+23.9_{-11.2}^{+5.4}\right)$ | $+40.4_{-11.6}^{+9.9}\left(+39.1_{-10.4}^{+9.7}\right)$ |
| weak | $-1.2(-8.3)$ | $-1.8(-8.2)$ | $-3.0(-7.8)$ |

Heavy Boson Radiation

| $\delta_{\text {HBR }}(\%)$ | 8 TeV | 13 TeV |  | 100 TeV |
| :--- | :---: | :---: | :---: | :---: |
| $W$ | $+0.42(+0.74)$ | $+0.37(+0.70)$ | $+0.14(+0.22)$ |  |
| $Z$ | $+0.29(+0.56)$ | $+0.34(+0.68)$ | $+0.51(+0.95)$ |  |
| $H$ | $+0.17(+0.43)$ | $+0.19(+0.48)$ | $+0.25(+0.53)$ |  |
| sum | $+0.88(+1.73)$ | $+0.90(+1.86)$ | $+0.90(+1.70)$ |  |

Partial compensation of Sudakov logs

NLO weak subchannels

| $\delta_{\mathrm{NLO}}(\%)$ | 8 TeV |  | 13 TeV |
| :--- | :---: | :---: | :---: |
| $g g$ | $-0.67(-2.9)$ | $-1.12(-4.0)$ | $-2.64(-6.8)$ |
| $u \bar{u}$ | $-0.01(-3.2)$ | $-0.15(-2.3)$ | $-0.10(-0.5)$ |
| $d \bar{d}$ | $-0.55(-2.2)$ | $-0.52(-1.9)$ | $-0.23(-0.5)$ |


| $\sigma[\mathrm{fb}]$ | 8 TeV | 13 TeV | 14 TeV | 25 TeV | 33 TeV | 50 TeV | 100 TeV |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $t \bar{t} Z Z$ | $0.502_{-8.6 \%}^{+2.9 \%}+2.7 \%$ | $2.12_{-8.6 \%}^{+3.8 \%}{ }_{-1.8 \%}^{+1.9 \%}$ | $2.59_{-8.7 \%}^{+4.3 \%}{ }_{-1.8 \%}^{+1.8 \%}$ | $11.1_{-9.1 \%}^{+6.9 \%}{ }_{-1.4 \%}^{+1.2 \%}$ | $21.1_{-9.4 \%}^{+8.1 \%}{ }_{-1.3 \%}^{+1.1 \%}$ | $51.6_{-9.8 \%}^{+9.9 \%}{ }_{-1.1 \%}^{+0.9 \%}$ | $204_{-9.9 \%}^{+11.3 \%}{ }_{-1.0 \%}^{+0.8 \%}$ |
| $t \bar{t} W^{+} W^{-}[4 \mathrm{f}]$ | $2.67{ }_{-11.1 \%}^{+6.2 \%}{ }_{-2.7 \%}^{+2.9 \%}$ | $11.8{ }_{-11.2 \%}^{+8.3 \%}{ }_{-2.4 \%}^{+2.3 \%}$ | $14.4{ }_{-12.8 \%}^{+12.2 \% ~}{ }_{-2.9 \%}^{+2.6 \%}$ | $66.6_{-10.8 \%}^{+9.5 \%}{ }_{-2.0 \%}^{+1.6 \%}$ | $130_{-10.8 \%}^{+10.2 \% ~}{ }_{-1.8 \%}$ | $327_{-10.6 \%}^{+10.9 \%}{ }_{-1.6 \%}^{+1.3 \%}$ | $1336_{-9.9 \%}^{+10.3 \%}{ }_{-1.3 \%}^{+1.0 \%}$ |
| $t \bar{t} \gamma \gamma$ | $2.77_{-10.5 \%}^{+6.4 \%}{ }_{-1.5 \%}^{+1.9 \%}$ | $10.3_{-13.3 \%}^{+13.9 \%}+1.3 \%$ | $12_{-12.6 \%}^{+12.5 \%}+1.2 \%$ | $44.8_{-13.5 \%}^{+15.7 \%}{ }_{-0.9 \%}^{+0.9 \%}$ | $78.2_{-13.6 \%}^{+16.4 \%}{ }_{-0.9 \%}^{+0.8 \%}$ | $184_{-14.7 \%}^{+19.2 \%}{ }_{-0.9 \%}^{+0.8 \%}$ | $624_{-13.4 \%}^{+15.5 \%}{ }_{-1.0 \%}^{+0.7 \%}$ |
| $t \bar{t} W^{ \pm} Z$ | $1.13_{-9.8 \%}^{+5.8 \%}+3.1 \%$ | $4.16_{-10.7 \%}^{+9.8 \%}{ }_{-1.6 \%}^{+2.2 \%}$ | $4.96_{-10.8 \%}^{+10.4 \%}+2.1 \%$ | $17.8_{-12.6 \%}^{+15.1 \%}{ }_{-1.1 \%}^{+1.5 \%}$ | $30.2_{-14.1 \%}^{+18.3 \%}{ }_{-0.9 \%}^{+1.2 \%}$ | $66_{-14.3 \%}^{+18.9 \%}{ }_{-0.8 \%}^{+1.1 \%}$ | $210_{-15.8 \%}^{+21.6 \%}{ }_{-0.8 \%}^{+1.0 \%}$ |
| $t \bar{t} Z \gamma$ | $1.39_{-11.2 \%}^{+6.9 \%}{ }_{-2.2 \%}^{+2.5 \%}$ | $5.77_{-12.1 \%}^{+10.5 \%}+1.8 \%$ | $6.95_{-12.1 \%}^{+10.7 \% ~+1.8 \%}$ | $29.9_{-12.4 \%}^{+12.9 \%}+1.3 \%$ | $56.5{ }_{-12.2 \%}^{+13.2 \%}{ }_{-1.4 \%}$ | $138_{-12.0 \%}^{+13.7 \%}{ }_{-1.1 \%}^{+1.0 \%}$ | $533{ }_{-11.1 \%}^{+13.3 \% ~}{ }_{-1.0 \%}^{+0.8 \%}$ |
| $t \bar{t} W^{ \pm} \gamma$ | $2.01_{-10.5 \%}^{+7.9 \%} \begin{gathered} +2.6 \% \\ -1.8 \% \end{gathered}$ | $6.73_{-11.6 \%}^{+12.0 \%}{ }_{-1.4 \%}^{+1.8 \%}$ | $7.99_{-11.9 \%}^{+12.8 \%}+1.7 \%$ | $27.6_{-14.4 \%}^{+18.7 \%}{ }_{-0.9 \%}^{+1.2 \%}$ | $46.3_{-15.1 \%}^{+20.2 \%}{ }_{-0.8 \%}^{+1.1 \%}$ | $98.4_{-15.9 \%}^{+21.9 \%}{ }_{-0.7 \%}^{+1.0 \%}$ | $318_{-17.7 \%}^{+22.5 \%}{ }_{-0.7 \%}^{+1.0 \%}$ |
| $t \bar{t} t \bar{t}$ | $1.71_{-26.2 \%}^{+24.9 \%}{ }_{-8.4 \%}^{+7.9 \%}$ | $13.3_{-25.3 \%}^{+25.8 \%}{ }_{-6.6 \%}^{+5.8 \%}$ | $17.8_{-25.4 \%}^{+26.6 \%}{ }_{-6.4 \%}^{+5.5 \%}$ | $130_{-24.3 \%}^{+26.7 \%}+3.8 \%$ | $297_{-23.3 \%}^{+25.5 \%}{ }_{-3.9 \%}^{+3.1 \%}$ | $929_{-22.4 \%}^{+24.9 \%}{ }_{-3.0 \%}^{+2.4 \%}$ | $4934_{-21.3 \%}^{+25.0 \%}{ }_{-2.1 \%}^{+1.7 \%}$ |
| $\sigma[\mathrm{pb}]$ | 8 TeV | 13 TeV | 14 TeV | 25 TeV | 33 TeV | 50 TeV | 100 TeV |
| $t \bar{t} Z$ | $0.226_{-11.9 \%}^{+9.0 \%}{ }_{-3.0 \%}^{+2.6 \%}$ | $0.874_{-11.7 \%}^{+10.3 \%}+2.0 \%$ | $1.057_{-11.7 \%}^{+10.4 \%}{ }_{-2.4 \%}^{+1.9 \%}$ | $4.224_{-11.0 \%}^{+11.0 \%}{ }_{-1.8 \%}^{+1.5 \%}$ | $7.735_{-10.8 \%}^{+11.2 \%}{ }_{-1.5 \%}^{+1.3 \%}$ | $18_{-10.2 \%}^{+11.1 \%}{ }_{-1.3 \%}^{+1.1 \%}$ | $64.17_{-11.0 \%}^{+11.1 \%}{ }_{-1.2 \%}^{+0.9 \%}$ |
| $t \bar{t} W^{ \pm}$ | $0.23_{-10.6 \%}^{+9.6 \%}{ }_{-1.7 \%}^{+2.3 \%}$ | $0.645_{-11.6 \%}^{+13.0 \%}{ }_{-1.3 \%}^{+1.7 \%}$ | $0.745_{-11.8 \%}^{+13.5 \%}+1.6 \%$ | $2.188_{-13.2 \%}^{+17.0 \%}{ }_{-0.9 \%}^{+1.3 \%}$ | $3.534_{-13.7 \%}^{+18.1 \%}{ }_{-0.8 \%}^{+1.2 \%}$ | $7.03_{-14.3 \%}^{+19.2 \%}{ }_{-0.8 \%}^{+1.1 \%}$ | $20.55_{-18.1 \%}^{+21.5 \%}{ }_{-0.8 \%}^{+1.1 \%}$ |
| $t \bar{t} \gamma$ | $0.788_{-13.5 \%}^{+12.7 \%}+2.1 \%$ | $2.746_{-13.5 \%}^{+14.2 \%}+1.6 \%$ | $3.26_{-13.4 \%}^{+14.2 \%}{ }_{-1.9 \%}^{+1.6 \%}$ | $11.77_{-12.7 \%}^{+14.5 \%}+1.2 \%$ | $20.84_{-12.5 \%}^{+14.9 \%}{ }_{-1.3 \%}^{+1.1 \%}$ | $45.68_{-11.7 \%}^{+14.2 \%}{ }_{-1.2 \%}^{+1.0 \%}$ | $152.6_{-13.7 \%}^{+14.3 \%}{ }_{-1.2 \%}^{+0.9 \%}$ |
| $t \bar{t} H$ | $0.136_{-9.1 \%}^{+3.3 \%}+2.8 \%$ | $0.522_{-9.4 \%}^{+6.0 \%}+2.1 \%$ | $0.631_{-9.4 \%}^{+6.3 \%}{ }_{-2.5 \%}^{+2.0 \%}$ | $2.505_{-9.4 \%}^{+8.3 \%}{ }_{-1.9 \%}^{+1.6 \%}$ | $4.567_{-9.2 \%}^{+8.8 \%}{ }_{-1.7 \%}^{+1.4 \%}$ | $10.55_{-9.0 \%}^{+9.5 \%}{ }_{-1.4 \%}^{+1.2 \%}$ | $37.65_{-9.8 \%}^{+10.0 \%} \quad{ }_{-1.3 \%}^{+1.0 \%}$ |

## differential plots


photon PDF NO


## differential plots


photon PDF NO

photon PDF YES

