

# Predicting the $\sin \phi_s$ transverse single-spin asymmetry of pion production at an electron ion collider

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# OUTLINE

1. Introduction

2. Formalism

3. Numerical Calculation

4. Conclusion

# INTRODUCTION

## ◆ Transversity $h_1^q(x)$ :

- Fundamental observables to encode nucleon structure
- Chiral-odd (difficult to probe)

To manifest the effect of  $h_1^q(x)$  in a process, another chiral-odd function is needed to couple with transversity to ensure chirality conservation.

## ◆ Some approaches to extract transversity function:

- (TMD) factorization frame work in SIDIS ——Collins function  $H^\perp$ .
- Collinear factorization formalism in SIDIS——dihadron fragmentation function.
- In the Drell-Yan process——the antiquark transversity.

# INTRODUCTION

- ◆ An alternative approach to access transversity in SIDIS—  
—the twist-3 chiral-odd fragmentation function  $\tilde{H}(z)$   
serve as a “spin analyzer”.

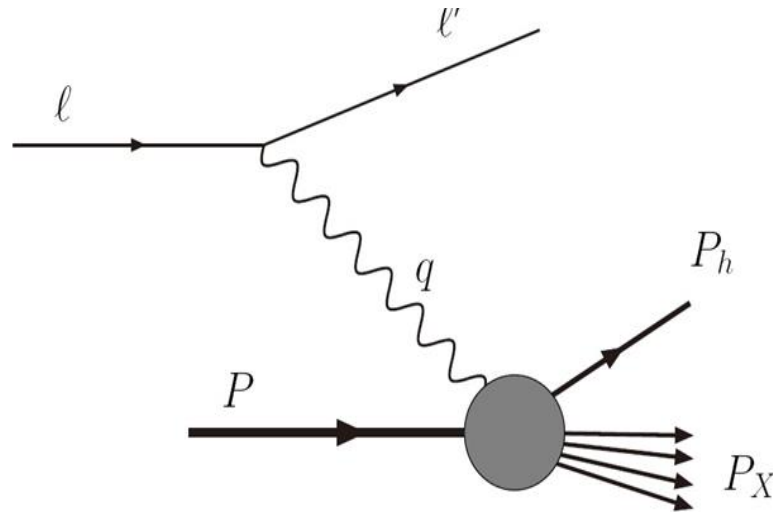
Motivation: In collinear picture, only the coupling between the function  $\tilde{H}(z)$  and the transversity remains as a contribution to the  $\sin \phi_S$  azimuthal modulation in the leptonproduction of a single hadron off a transversely polarized nucleon.

- ◆ The advantage of this approach is that the transverse momentum of the final state hadron is not necessarily to be measured, in contrast to the Collins effect.

# FORMALISM

- ◆ The process under study is the pion electroproduction off a transversely polarized proton target

$$e(\ell) + p^\uparrow(P) \rightarrow e(\ell') + \pi(P_h) + X(P_X),$$

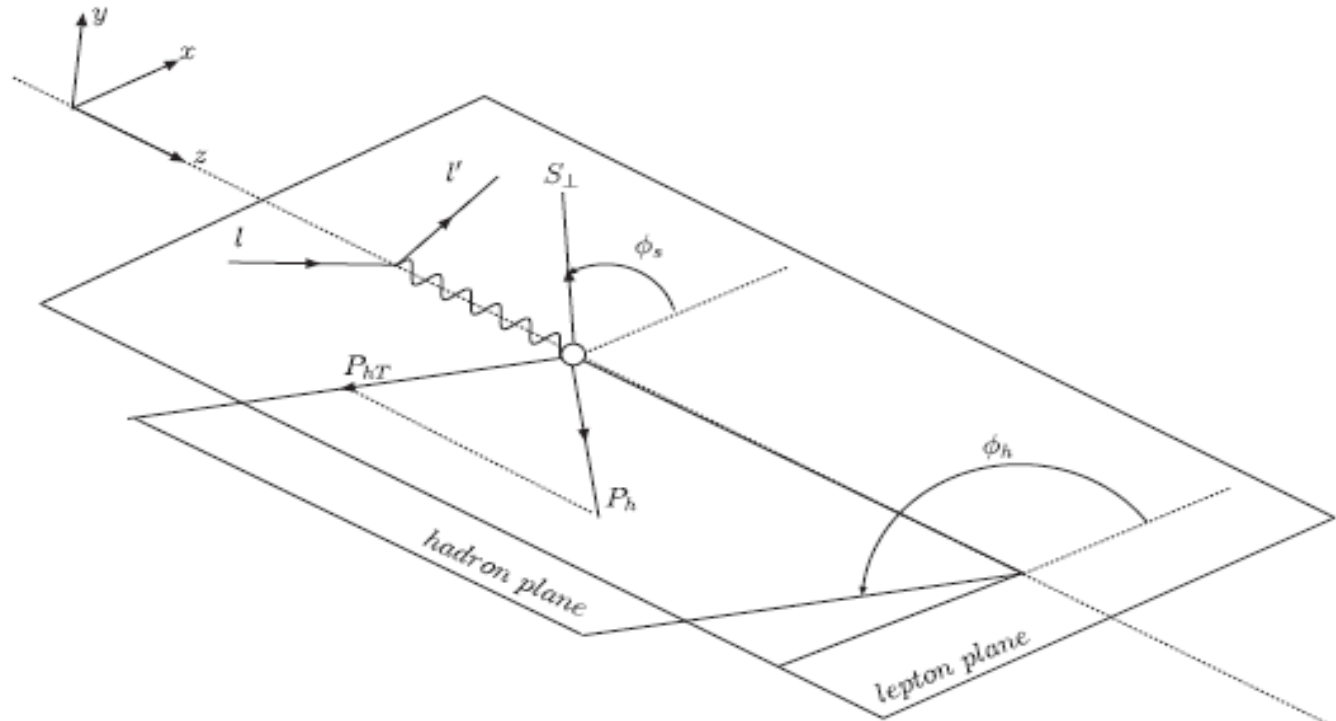


- ◆ The invariants used to express the differential cross section

$$x = \frac{Q^2}{2P \cdot q}, \quad y = \frac{P \cdot q}{P \cdot l}, \quad z = \frac{P \cdot P_h}{P \cdot q},$$
$$\gamma = \frac{2Mx}{Q}, \quad Q^2 = -q^2, \quad s = (P + l)^2.$$

# FORMALISM

- ◆ The reference frame of the process under study is



# FORMALISM

- ◆ Up to twist-3 level, the sixfold  $(x, y, z, \phi_h, \phi_S$  and  $P_{hT}$ ) differential cross section in SIDIS with a transversely polarized target has the general form

$$\begin{aligned} \frac{d^6\sigma}{dx dy dz d\phi_h d\phi_S dP_{hT}^2} = & \frac{\alpha^2}{xy Q^2} \frac{y^2}{2(1-\varepsilon)} \left( 1 + \frac{\gamma^2}{2x} \right) \\ & \times \sqrt{2\varepsilon(1+\varepsilon)} \{ \sin\phi_S F_{UT}^{\sin\phi_S}(x, z, P_T) \\ & + \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)}(x, z, P_T) \\ & + \text{leading twist terms} \}, \end{aligned} \quad (3)$$

# FORMALISM

- ◆ Perform the integration over the  $P_{hT}$ , the differential cross section turns to the form

$$\frac{d^4\sigma}{dx dy dz d\phi_s} = \frac{2\alpha^2}{xyQ^2} \frac{y^2}{2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \times \sqrt{2\epsilon(1+\epsilon)} \sin\phi_s F_{UT}^{\sin\phi_s}(x, z).$$

- ◆ The x-dependent and z-dependent asymmetry are defined as

$$A_{UT}^{\sin\phi_s}(x) = \frac{\int dy \int dz \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \sqrt{2\epsilon(1+\epsilon)} F_{UT}^{\sin\phi_s}(x, z)}{\int dy \int dz \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x}\right) F_{UU}(x, z)},$$

$$A_{UT}^{\sin\phi_s}(z) = \frac{\int dx \int dy \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \sqrt{2\epsilon(1+\epsilon)} F_{UT}^{\sin\phi_s}(x, z)}{\int dx \int dy \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x}\right) F_{UU}(x, z)}.$$



# FORMALISM

- ◆  $F_{UT}^{\sin \phi_s}(x, z)$  and  $F_{UU}(x, z)$  are the collinear counterpart of the original structure function

$$F_{UT}^{\sin \phi_s}(x, z) = -x \sum_q e_q^2 \frac{2M_h}{Q} h_1^q(x) \frac{\tilde{H}^q(z)}{z},$$

$$F_{UU}(x, z) = x \sum_q e_q^2 f_1^q(x) D_1^q(z),$$

- ◆ Twist-3 fragmentation function

$$\tilde{H}^{h/q}(z) = 2z^3 \int_z^\infty \frac{dz_1}{z_1^2} \frac{1}{\frac{1}{z} - \frac{1}{z_1}} \hat{H}_{FU}^{h/q, \mathcal{S}}(z, z_1).$$

- ◆ At the energy scale  $Q^2 = 1 \text{ GeV}^2$ ,  $\hat{H}_{FU}^{h/q, \mathcal{S}}(z, z_1)$  has been extracted as

$$\frac{\hat{H}_{FU}^{\pi^+/(u, \bar{d}), \mathcal{S}}(z, z_1)}{D_1^{\pi^+/(u, \bar{d})}(z) D_1^{\pi^+/(u, \bar{d})}(z/z_1)} = \frac{N_{\text{fav}}}{2I_{\text{fav}} J_{\text{fav}}} z^{\alpha_{\text{fav}}}(z/z_1)^{\alpha'_{\text{fav}}} \times (1-z)^{\beta_{\text{fav}}}(1-z/z_1)^{\beta'_{\text{fav}}},$$

# FORMALISM

- ◆ The  $\pi^-$  fragmentation functions may be fixed through charge conjugation

$$\hat{H}_{FU}^{\pi^-/(d,\bar{u}),\mathcal{S}}(z, z_1) = \hat{H}_{FU}^{\pi^+/(u,\bar{d}),\mathcal{S}}(z, z_1)$$

$$\hat{H}_{FU}^{\pi^-/(u,\bar{d}),\mathcal{S}}(z, z_1) = \hat{H}_{FU}^{\pi^+/(d,\bar{u}),\mathcal{S}}(z, z_1)$$

- ◆ The  $\pi^0$  fragmentation functions are given by the average of the fragmentation functions for  $\pi^+$  and  $\pi^-$ .
- ◆ At the initial scale  $Q^2 = 2.14 \text{ GeV}^2$ , we adopt the standard parametrization for the transversity

$$h_1^q(x) = \frac{1}{2} \mathcal{N}_q^T(x) [f_1^q(x) + g_1^q(x)],$$

$$\mathcal{N}_q^T(x) = N_q^T x^\alpha (1 - \beta)^\beta \frac{(\alpha + \beta)^{\alpha + \beta}}{\alpha^\alpha \beta^\beta}.$$

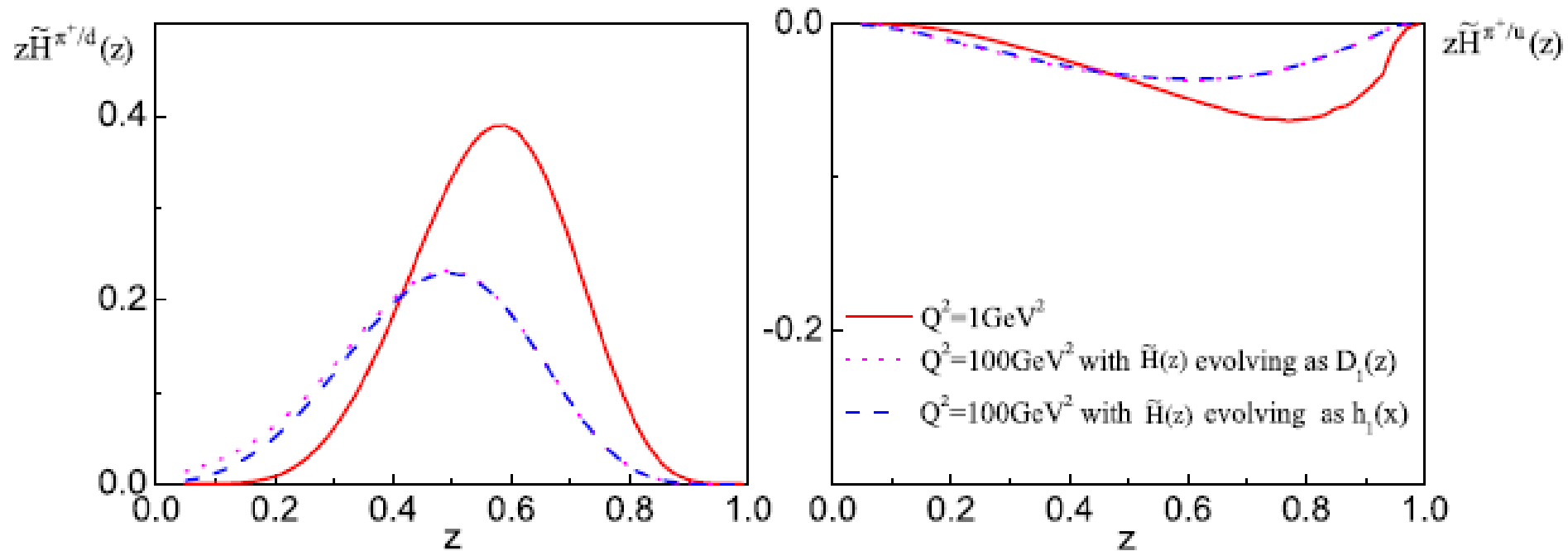
# NUMERICAL CALCULATION

- ◆ Kinematics at EIC covers a wide range of  $Q \longrightarrow$  QCD evolution of the transversity and fragmentation functions  $\tilde{H} \longrightarrow$  implement LO QCD evolution
- ◆ Adopt two different choices to evolve  $\tilde{H}$ :
  - Evolve as the fragmentation function  $D_1$ .  
K. Kanazawa, Y. Koike, A. Metz, and D. Pitonyak, Phys. Rev. D 89, 111501 (2014)
  - Evolve as the transversity  $h_1(x)$ .  
Motivation:  $\tilde{H}$  is a chiral-odd fragmentation function.

# NUMERICAL CALCULATION

- ◆ Result of  $z\tilde{H}^{\pi^+/d}$  (left panel)  $z\tilde{H}^{\pi^+/u}$  (right panel) at the initial scale  $Q^2 = 1 \text{ GeV}^2$  (solid lines) and the evolved results at  $Q^2 = 100 \text{ GeV}^2$

Dotted lines: evolving as  $D_1$ , dashed lines: evolving as  $h_1$



# NUMERICAL CALCULATION

## ◆ The kinematical region available at EIC

$$Q^2 > 1 \text{ GeV}^2, \quad 0.001 < x < 0.4, \quad 0.01 < y < 0.95,$$

$$0.2 < z < 0.8, \quad \sqrt{s} = 45 \text{ GeV}, \quad W > 5 \text{ GeV},$$

$$W^2 = (P + q)^2 \approx \frac{1-x}{x} Q^2.$$

A. Accardi et al., arXiv:1212.1701

## ◆ Parametrization we adopt:

- unpolarized distribution  $f_1^q(x)$

M. Glück, E. Reya, and A. Vogt, Eur. Phys. J. C 5, 461 (1998)

- helicity distribution  $g_1^q(x)$

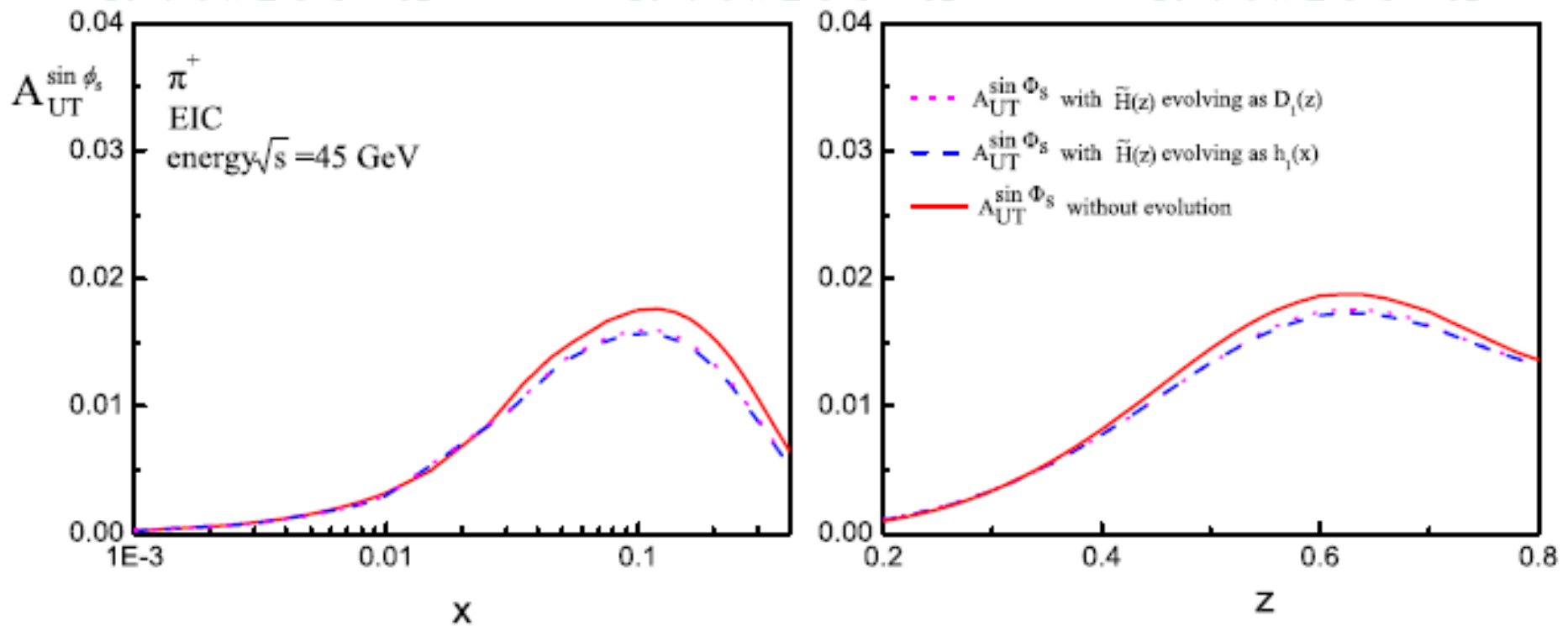
M. Glück, E. Reya, M. Stratmann, and W. Vogelsang, Phys. Rev. D 63, 094005 (2001)

- unpolarized integrated fragmentation function  $D_1^q(z)$

D. de Florian, R. Sassot, and M. Stratmann, Phys. Rev. D 75, 114010 (2007): LO set

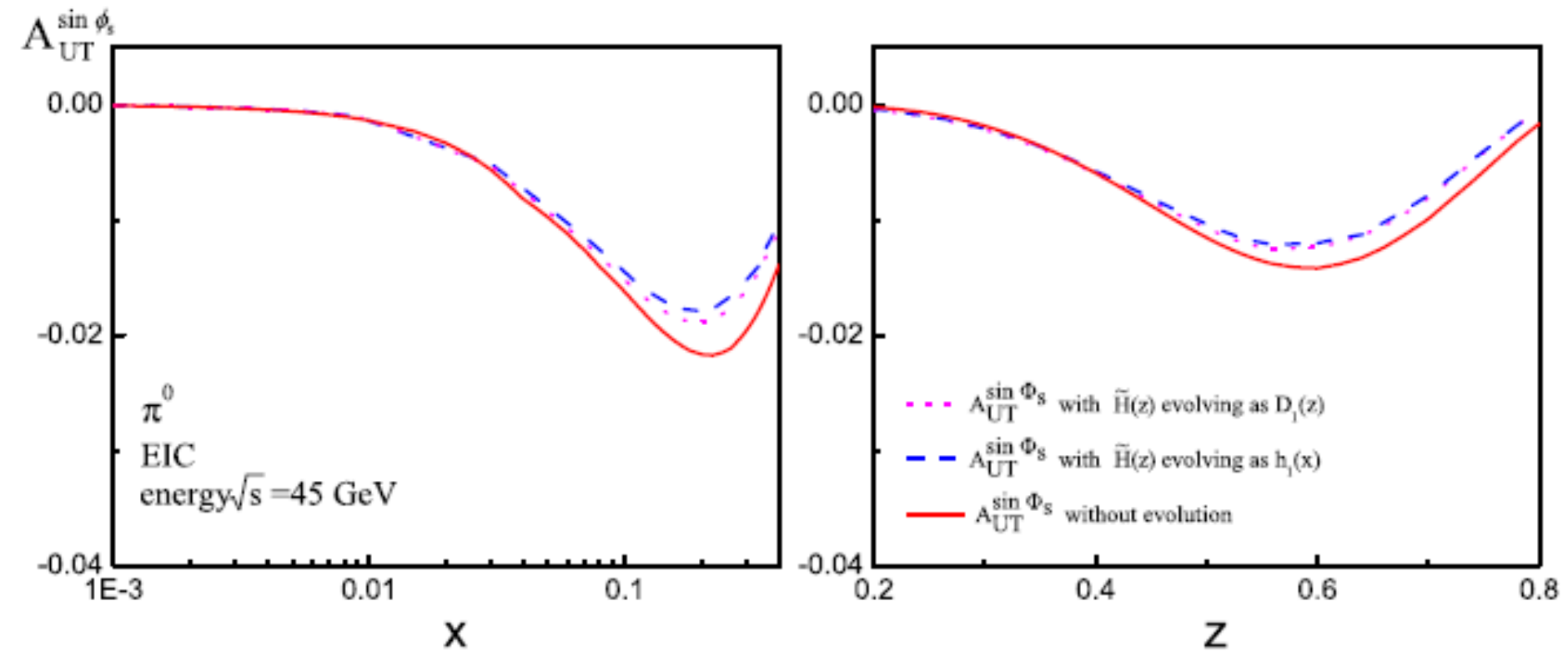
# NUMERICAL CALCULATION

- ◆ Transverse SSA  $\sin \phi_S$  of  $\pi^+$  production in SIDIS at EIC for  $\sqrt{s}=45$  GeV. The left panel shows the x-dependent asymmetry, while the right one shows the z-dependent asymmetry.



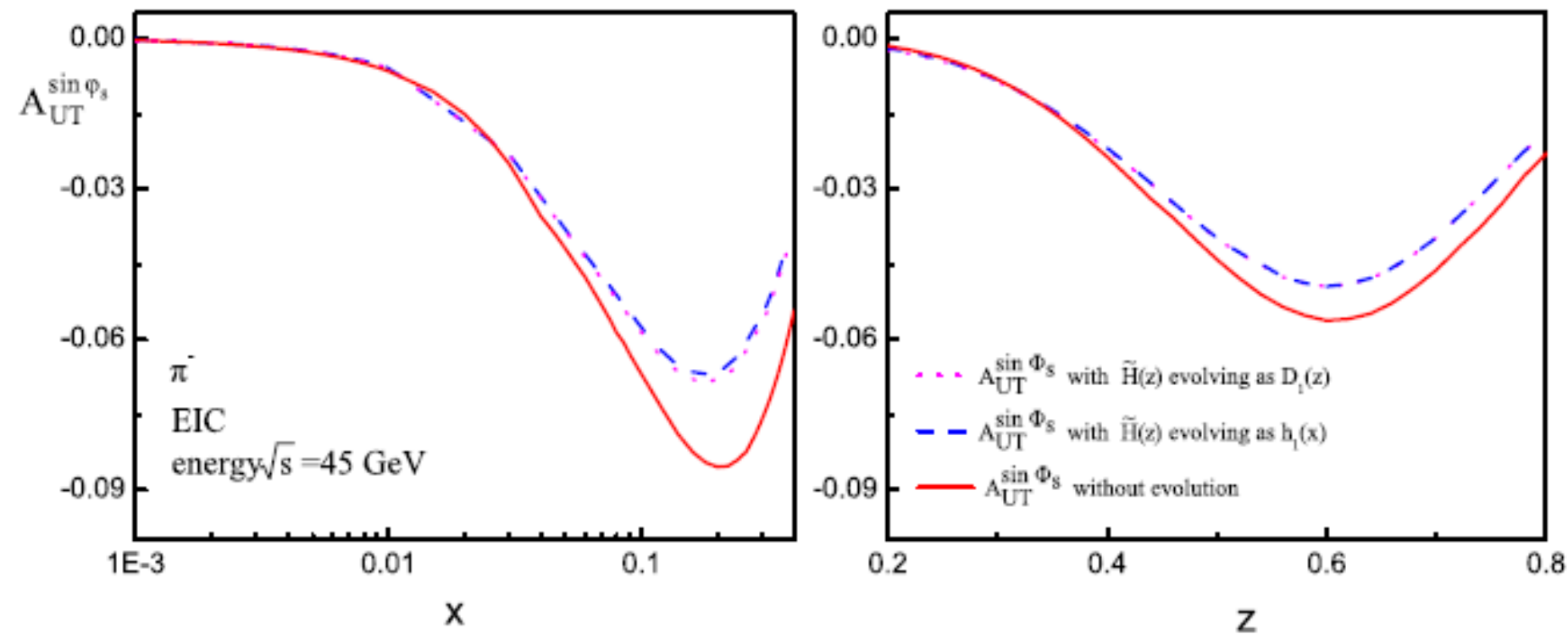
# NUMERICAL CALCULATION

- ◆ Similarly, transverse SSA  $\sin \phi_S$  of  $\pi^0$  production.



# NUMERICAL CALCULATION

- ◆ Similarly, transverse SSA  $\sin \phi_S$  of  $\pi^-$  production.





# CONCLUSION

- ◆ The numerical prediction shows that the asymmetries for the charged and neutral pions are all sizable, about several percent.
- ◆ It is quite promising that the  $\sin \phi_S$  asymmetries of meson production in SIDIS could be measured at the kinematics of EIC.
- ◆ The inclusion of the evolution effect may be important for the interpretation of future experimental data.



# Thanks for listening !

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