Angular analysis of $B^0 \to K^{*0} \mu^+ \mu^-$ at large recoil within the Aligned Two-Higgs-Doublet Model

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1. Motivation



"clean" observables at large recoil

$$P_5' = \frac{S_5}{2\sqrt{-S_2^s S_2^c}}, \quad P_2 = \frac{S_6^s}{8S_2^s}$$

[S. Descotes-Genon, et al.: 1303.5794]



[Belle: 1604.04042]

2. The aligned two-Higgs doublet model (A2HDM)

In so-called "Higgs basis"

$$\Phi_1 = \begin{bmatrix} G^+ \\ \frac{1}{\sqrt{2}} (v + S_1 + iG^0) \end{bmatrix}, \qquad \Phi_2 = \begin{bmatrix} H^+ \\ \frac{1}{\sqrt{2}} (S_2 + iS_3) \end{bmatrix}$$

Five physical degrees of freedom: two charged fields $H^{\pm}(x)$ and three neutral fields $\varphi_i^0(x) = \{h(x), H(x), A(x)\} = \mathcal{R}_{ij}S_j$.

The Yukawa Lagrangian

$$\mathcal{L}_{Y} = -\frac{\sqrt{2}}{v} \left[\bar{Q}'_{L} (M'_{d} \Phi_{1} + Y'_{d} \Phi_{2}) d'_{R} + \bar{Q}'_{L} (M'_{u} \tilde{\Phi}_{1} + Y'_{u} \tilde{\Phi}_{2}) u'_{R} + \bar{L}'_{L} (M'_{\ell} \Phi_{1} + Y'_{\ell} \Phi_{2}) \ell'_{R} \right] + \frac{1}{2} \left[\bar{Q}'_{L} (M'_{d} \Phi_{1} + Y'_{d} \Phi_{2}) d'_{R} + \bar{Q}'_{L} (M'_{u} \tilde{\Phi}_{1} + Y'_{u} \tilde{\Phi}_{2}) u'_{R} + \bar{L}'_{L} (M'_{\ell} \Phi_{1} + Y'_{\ell} \Phi_{2}) \ell'_{R} \right] + \frac{1}{2} \left[\bar{Q}'_{L} (M'_{d} \Phi_{1} + Y'_{d} \Phi_{2}) d'_{R} + \bar{Q}'_{L} (M'_{u} \tilde{\Phi}_{1} + Y'_{u} \tilde{\Phi}_{2}) u'_{R} + \bar{L}'_{L} (M'_{\ell} \Phi_{1} + Y'_{\ell} \Phi_{2}) \ell'_{R} \right] \right]$$

Aligned hypothesis [A. Pich and P. Tuzón: 0908.1554]

$$Y_{d,\ell} = \varsigma_{d,\ell} M_{d,\ell} , \qquad Y_u = \varsigma_u^* M_u ,$$

Charged Higgs sector

$$\mathcal{L}_{H^{\pm}} = -\frac{\sqrt{2}}{v} H^{+} \left\{ \bar{u} \left[\varsigma_{d} V_{\text{CKM}} M_{d} P_{R} - \varsigma_{u} M_{u}^{\dagger} V_{\text{CKM}} P_{L} \right] d + \varsigma_{\ell} \bar{\nu} M_{\ell} P_{R} \ell \right\} + \text{h.c.} ,$$

NFC models (based on discrete \mathcal{Z}_2 symmetries)

Model	$(\varsigma_d, \varsigma_u, \varsigma_\ell)$
Type I	$(\coteta,\coteta,\coteta)$
Type II	$(-\tan\beta,\cot\beta,-\tan\beta)$
Type X (lepton-specific)	$(\cot\beta,\cot\beta,-\tan\beta)$
Type Y (flipped)	$(-\tan\beta,\cot\beta,\cot\beta)$
Inert	$(0,\!0,\!0)$

3. Effective Hamiltonian

Effective Hamiltonian for $b \to s\mu^+\mu^-$ transitions

$$\mathcal{H}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_i \left(C_i O_i + C_i' O_i' \right) \,,$$

radiative and semileptonic operators

$$\begin{split} O_7 &= \frac{e}{16\pi^2} \bar{m}_b \left(\bar{s} \sigma^{\mu\nu} P_R b \right) F_{\mu\nu} , \qquad O_7' &= \frac{e}{16\pi^2} \bar{m}_b \left(\bar{s} \sigma^{\mu\nu} P_L b \right) F_{\mu\nu} , \\ O_9 &= \frac{e^2}{16\pi^2} \left(\bar{s} \gamma^{\mu} P_L b \right) \left(\bar{\mu} \gamma_{\mu} \mu \right) , \qquad O_9' &= \frac{e^2}{16\pi^2} \left(\bar{s} \gamma^{\mu} P_R b \right) \left(\bar{\mu} \gamma_{\mu} \mu \right) , \\ O_{10} &= \frac{e^2}{16\pi^2} \left(\bar{s} \gamma^{\mu} R_L b \right) \left(\bar{\mu} \gamma_{\mu} \gamma_5 \mu \right) , \qquad O_{10}' &= \frac{e^2}{16\pi^2} \left(\bar{s} \gamma^{\mu} P_R b \right) \left(\bar{\mu} \gamma_{\mu} \gamma_5 \mu \right) , \end{split}$$

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Wilson coefficients in the A2HDM



Result



4. Numerical results

 Restrict C₇^{NP} the experimental combined result [Y. Amhis, et al.(HFAG) Collaboration: 1412.7515]

$$\mathcal{B}[b \to s\gamma]_{\text{exp}} = (3.43 \pm 0.21 \pm 0.07) \times 10^{-4}.$$

The SM prediction at the NNLO in QCD is [M. Misiak, et al.: 1503.01789].

$$\mathcal{B}[b \to s\gamma]_{\rm SM} = (3.36 \pm 0.23) \times 10^{-4}.$$

The $C_7^{\rm H^{\pm}}$ and $C_8^{\rm H^{\pm}}$ fulfill
 $-0.0634 \le C_7^{\rm H^{\pm}} + 0.242 C_8^{\rm H^{\pm}} \le 0.0464.$

(2) Global fit results for NP contributions on Wilson coefficients C^(I)_{9,10} [S. Descotes-Genon, et al.: 1510.04239; S. Meinel and D. van Dyk: 1603.02974]

$$\begin{split} -2.2 \, &\leq C_9^{\rm NP} \,\leq\, 2.5\,, \quad -0.5 \,\leq C_{10}^{\rm NP} \,\leq\, 4.2\,, \\ -1.3 \,\leq C_9^{\prime \rm NP} \,\leq\, 3.7\,, \quad -1.0 \,\leq C_{10}^{\prime \rm NP} \,\leq\, 3.1\,, \end{split}$$



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• $|\varsigma_d| \le 212, 476, \text{ and } 622;$

• $|\varsigma_u| \le 0.506, 0.763$, and 0.990; corresponding to $M_{H^{\pm}} = 80, 300$, and 500 GeV, respectively.

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Numerical results

The correlations between each Wilson coefficients are shown



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Numerical results



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	\mathbf{SM}	A2HDM-case 1	A2HDM-case 2
$q_0^2(P_2)$	$3.43_{-0.32}^{+0.33}$	(3.02, 3.90)	(3.02, 4.79)
$q_0^2(P_5')$	$2.02^{+0.19}_{-0.15}$	(1.77, 2.32)	(1.79, 4.85)

Table: The zero-crossing points of P_2 (nonzero one) and P'_5 .

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 $M_{H^{\pm}} \geq 432 \,\mathrm{GeV}$ in types-II,Y

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5. Conclusions

- We presented a complete one-loop calculation of the $C_{7,9,10}^{(\prime)H^{\pm}}$ in A2HDM.
- We obtain the allowed region in the $\varsigma_u \varsigma_d$ plane. The constraints are $|\varsigma_u| \leq 0.506 \ (0.990)$ and $|\varsigma_d| \leq 212 \ (622)$ for $M_{H^{\pm}} = 80 \ (500)$ GeV.
- In A2HDM, there are only two cases for Wilson coefficients of NP

$$\begin{aligned} & \operatorname*{case} 1: C_7^{\mathrm{H}^{\pm}}, \, C_9^{\mathrm{H}^{\pm}}, \, C_{10}^{\mathrm{H}^{\pm}} & \mathrm{with} \quad C_{9,10}^{\prime \mathrm{H}^{\pm}} \simeq 0 \,, \\ & \operatorname*{case} 2: C_7^{\mathrm{H}^{\pm}}, \, C_7^{\prime \mathrm{H}^{\pm}}, \, C_{10}^{\prime \mathrm{H}^{\pm}} & \mathrm{with} \quad C_{9,10}^{\mathrm{H}^{\pm}} \simeq 0 \,; \end{aligned}$$

There are only a small impact on P_2 and P'_5 in case 1. The case 2 can supply a bigger P'_5 than SM at large recoil (in low q^2), to conform with the experimental results.

• The charged-Higgs mass should be greater than 432 GeV in types-II,Y.

Thank You!

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