## Interpretation of Positive Parity, J/ $\Psi-\varphi$ Resonances at LHCb

Luciano Maiani, Roma University and INFN Roma LHCb Workshop, CERN, Oct. 13, 2016 Jiao Tong University, Shanghai, Nov. 4, 2016

## 1. Old and new structures observed by LHCb arXiv:1606.07895

## Results of fit: m(J/ $\psi \phi)$



- 4 visible structures fit with BW amplitudes

28 Recontres de Blois, June 2, 2016

- Four structures
- positive parity, $\mathrm{J}=0$ and 1 , positive charge conjugation
- X(4140) seen previously by CDF, D0, CMS and by BELLE


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## Results of fit: m(J/ $\psi \phi)$



28 Recontres de Blois, June 2, 2016


- JP also measured all with $>4 \sigma$ significances

| Particle | JP | Significance | Mass <br> (MeV) | $\Gamma$ $(\mathrm{MeV})$ | Fit Fraction (\%) |
| :---: | :---: | :---: | :---: | :---: | :---: |
| X(4140) | $1^{+}$ | 8.4 \% | $4146.5 \pm 4.5{ }_{-28}^{4+6}$ | $83 \pm 2{ }_{-14}^{1+1}$ | $13.0 \pm 3.2{ }_{-20}^{1+88}$ |
| X(4274) | $1^{+}$ | 6.0 o | $4273.3 \pm 8.3_{-3.6}^{+17.2}$ | $56 \pm 11_{-11}^{+8}$ | $7.1 \pm 2.5{ }_{-24}^{135}$ |
| X(4500) | $0^{+}$ | 6.1 \% | $4506 \pm 11_{-15}^{+12}$ | $92 \pm 2{ }_{-20}^{2+1}$ | $6.6 \pm 2.4$ |
| X(4700) | $0^{+}$ | 5.6 \% | $4704 \pm 10_{-24}^{1+14}$ | $120 \pm 3{ }_{-33}^{1+32}$ | $12 \pm 5$ |
| NR | $0^{+}$ | $6.4 \sigma$ |  |  | $46 \pm 11$ |

- Four structures
- positive parity, $\mathrm{J}=0$ and 1 , positive charge conjugation
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We suggest to fit the structures in two tetraquark multiplets, $S$-wave ground state and the first radial excitation, with composition $[c s][\bar{c} \bar{s}]$. L. Maiani, A. Polosa, V. Riquer, PRD 94 (2016) 054026 With the previously identified $[c q][\bar{c} \bar{q}](q=u, d)$ multiplet, the new resonances would make a step towards a full nonet of S-wave tetraquarks made by c c-bar with a pair of light ( $u, d$, s) quarks.

## 2. Forces between colored objects (Han-Nambu)

- Interaction proportional to: $\frac{g^{2}}{q^{2}} \times 2\langle R| \mathbf{T} \cdot \mathbf{T}|R\rangle=\frac{g^{2}}{q^{2}}\left[C_{2}(R)-C_{2}(A)-C_{2}(B)\right]$
- quark-antiquark
- singlet: attractive (-8/3)
- octet: repulsive (+1/3)
- quark-quark
- three bar: attractive (-4/3)

$$
C_{2}(R)=\mathbf{T}_{R} \cdot \mathbf{T}_{R}
$$

| R | $\mathrm{C}_{2}(\mathrm{R})$ |
| :---: | :---: |
| $\mathbf{1}$ | 0 |
| $\mathbf{8}$ | 3 |
| $\mathbf{3}, \overline{\mathbf{3}}$ | $4 / 3$ |
| $\mathbf{6}$ | $10 / 3$ |

- six: repulsive (+2/3)
- quark-antiquark pairs bind in color singlet mesons;
- diquarks bind:
to another quark, to make a color singlet baryon to an antidiquark, to make a color singlet tetraquark.


## Constituent quark model

Old days: Sakharov\&Zeldovich
In QCD: De Rujula, Georgi and Glashow, PRL 38 (1977) 317
Revisited \& applied to tetraquarks: L.M., A. Polosa, V. Riquer, F, Piccinini, PRD 71 (2005) 014028

- color string forces produce an overall spin-independent potential that confines quarks inside a definite volume (bag), with some wave functions $\psi\left(\mathrm{X}_{1}, \mathrm{X}_{2}, \mathrm{X}_{3}\right)$
- residual quark-quark or quark-antiquark interactions are color-magnetic, spinspin, forces of the form

$$
H_{i j}=\frac{g^{2}}{m_{1} m_{2}}\left(T_{1} \cdot T_{2}\right) \mathbf{s}_{i} \cdot \mathbf{s}_{j} \delta^{(3)}\left(x_{1}-x_{2}\right)
$$

- T are the color charges, g the color coupling, the form is derived from the non relativistic limit of QCD
- if $\mathrm{i}, \mathrm{j}$ are in a color representation R , the formula simplifies to ( CF is a color factor similar to the one introduced for color interactions)

$$
\begin{aligned}
& \mathbf{H}_{i j}=2 \kappa_{i j} \mathbf{s}_{i} \cdot \mathbf{s}_{j} \\
& \kappa_{i j}=\mathrm{CF}(\mathrm{R}) \times \frac{g^{2}}{m_{1} m_{2}}|\psi(0)|^{2}
\end{aligned}
$$

$$
H=\sum_{i} m_{i}+\sum_{i<j} 2 \kappa_{i j} \mathbf{s}_{j} \cdot \mathbf{s}_{j} \quad \kappa_{i j}=\mathrm{CF}(\mathrm{R}) \times \frac{g^{2}}{m_{1} m_{2}}|\psi(0)|^{2}
$$

the Hamiltonian can be developed to first order in the small mass differences, like $\mathrm{m}_{\mathrm{s}}-\mathrm{m}_{\mathrm{u}, \mathrm{d}}$,

- there is also a first order contribution from the spin-spin interaction, which is very crucial for baryons
- one usually assumes that the wave function overlap is the same in all mesons and all baryons, but this is dubious in the case of hadrons with a very heavy quark, c or b .
- Works well for mesons and baryons (too well?)
- Few parameters: $\mathrm{m}_{\mathrm{u}}, \mathrm{m}_{\mathrm{d}}, \mathrm{m}_{\mathrm{s}}, \mathrm{m}_{\mathrm{c}}, \mathrm{m}_{\mathrm{s}}, \kappa_{\mathrm{ij}}$
- same values of masses, within $\pm 30 \mathrm{MeV}$, reproduce masses of different hadrons
- $\kappa_{\mathrm{ij}}$ scale approx. like $1 / \mathrm{m}_{\mathrm{i}} \mathrm{m}_{\mathrm{j}}$ and not far from scaling with color factors


## MESONS



## + similar multiplet of vector mesons



Spin 3/2 baryons

- Discovered by CLEO


## Data from PdG live

| MESONS | $q$ | $s$ | $c$ | $b$ |
| :---: | :---: | :---: | :---: | :---: |
| $\operatorname{mass}(\mathrm{MeV})$ | 308 | 484 | 1664 | 5005 |


| BARYONS | $q$ | $s$ | $c$ | $b$ |
| :---: | :---: | :---: | :---: | :---: |
| $\operatorname{mass}(\mathrm{MeV})$ | 362 | 540 | 1710 | 5044 |


| MESONS | $q \bar{q}$ | $q \bar{s}$ | $s \bar{s}$ | $c \bar{q}$ | $c \bar{s}$ | $b \bar{q}$ | $b \bar{s}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\left(\kappa_{i j}\right)_{1}(\mathrm{MeV})$ | 318 | 200 | 103 | 69 | 72 | 23 | 24 |
| $\left(\kappa_{i j}\right)_{1} m_{i} m_{j} / \Lambda_{Q C D}^{3}$ | 1.9 | 1.9 | 1.5 | 2.3 | 3.7 | 0.7 | 1.2 |


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| $\left(\kappa_{i j}\right)_{\overline{3}}(\mathrm{MeV})$ | 98 | 59 | 23 | 15 | 50 | 2.5 | 38 |
| $\left(\kappa_{i j}\right)_{\overline{3}} m_{i} m_{j} / \Lambda_{Q C D}^{3}$ | 0.82 | 0.74 | 0.43 | 0.57 | 2.7 | 0.29 | 6.6 |


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## 3.Tetraquark constituent picture of unexpected quarkonia

L.Maiani, F.Piccinini, A.D.Polosa and V.Riquer, Phys. Rev. D 71 (2005) 014028

- $\mathrm{I}=1,0$
- S-wave: positive parity
- total spin of each diquark, $S=1,0 \quad[c q]_{S=0,1}\left[\bar{c} \bar{q}^{\prime}\right]_{\bar{S}=0,1}$
- neutral states may be mixtures of isotriplet and isosinglet
$\bar{c} \uparrow$
- mass splitting described by

$$
\mathbf{H}=2 m_{(\text {diquark })}+\sum_{i<j} 2 \kappa_{i j}\left(\mathbf{s}_{i} \cdot \mathbf{s}_{j}\right)
$$

The S -wave, $\mathrm{J}^{\mathrm{P}}=1{ }^{+}$charmonium tetraquarks
in the $|S, \bar{S}\rangle_{J}$ basis we have the following states

$$
\begin{aligned}
& \mathrm{J}^{\mathrm{P}}=0^{+} \quad C=+\quad X_{0}=|0,0\rangle_{0}, \quad X_{0}^{\prime}=|1,1\rangle_{0} \\
& \mathrm{~J}^{\mathrm{P}}=1^{+} \quad C=+\quad X_{1}=\frac{1}{\sqrt{2}}\left(|1,0\rangle_{1}+|0,1\rangle_{1}\right) \\
& \mathrm{J}^{\mathrm{P}}=1^{+} \quad C=-\quad Z=\frac{1}{\sqrt{2}}\left(|1,0\rangle_{1}-|0,1\rangle_{1}\right), Z^{\prime}=|1,1\rangle_{1} \\
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& \text { diagonalize } \mathrm{H} \\
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# Can we extrapolate spin-spin couplings from mesons and baryons? 

## Phenomenology confirms

- q-qbar interaction in tetra quarks is not directly deducible from the meson spectrum, since the q-qbar pair in tetra is in a superposition of color octet and singlet.
- using one-gluon dominance, we estimate $\quad \kappa_{q \bar{q}, \text { tetra }}=\frac{1}{4}\left(\kappa_{q \bar{q}}\right)_{1}$
- if we take $\left(\kappa_{q \bar{q}}\right)_{1}$ from mesons, we obtain a value of $\kappa_{q \bar{q}}$, tetra $\gg\left(\kappa_{c q}\right)_{\overline{3}}$ from baryons and this leads to the wrong ordering of the $Z(3900)$ and $Z(4030)$ with respect to $\mathrm{X}(3872)$ :

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\begin{aligned}
& X(3872): \mathcal{C}=+, S_{c \bar{c}}=S_{q \bar{q}}=1 \\
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- $\mathrm{Z}\left(\mathrm{S}_{\mathrm{q} q \mathrm{qbar}}=1\right)=\mathrm{Z}_{\mathrm{deg}}$ is $\sim$ degenerate with $\mathrm{X}: \mathrm{Z}_{\mathrm{deg}}=\mathrm{Z}(3900)$
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## The new ansatz

- Mass ordering is explained by the hypohesis that the dominant spin-spin interactions in tetraquarks are those inside the diquark or the antidiquark.
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Assuming: $\quad \begin{aligned} X(3872) & =X_{1}=\frac{1}{\sqrt{2}}\left(|1,0\rangle_{1}+|0,1\rangle_{1}\right) \\ Z(3900) & =Z=\frac{1}{\sqrt{2}}\left(|1,0\rangle_{1}-|0,1\rangle_{1}\right), Z(4030)=Z^{\prime}=|1,1\rangle_{1}\end{aligned}$
$X(3872)$ and $Z(3900)$ have one spin 1, while $Z(4030)$ has two spins 1 and therefore it is heavier.

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- With this hypothesis, H is diagonal in the diquark spin basis (used before) and the Hamiltonian simply counts the number of diquarks with spin 1

Assuming: $\quad \begin{aligned} X(3872) & =X_{1}=\frac{1}{\sqrt{2}}\left(|1,0\rangle_{1}+|0,1\rangle_{1}\right) \\ Z(3900) & =Z=\frac{1}{\sqrt{2}}\left(|1,0\rangle_{1}-|0,1\rangle_{1}\right), Z(4030)=Z^{\prime}=|1,1\rangle_{1}\end{aligned}$
$X(3872)$ and $Z(3900)$ have one spin 1 , while $Z(4030)$ has two spins 1 and therefore it is heavier.

- A simple explanation of the dominance of inter-diquark interaction could be that diquarks and antidiquarks are at such relative distance in the hadron, as to suppress the overlap probability, unlike what happens, e.g., in the usual baryons


## Parameters

- The spectrum of 1S ground states is characterised by two quantities:
- the diquark mass, $\mathrm{m}_{[\mathrm{cq}]}$
- the spin-spin interaction inside the diquark or the antidiquark, $\kappa_{\mathrm{cq}}$.
- The first radially excited, 2 S , states are shifted up by a common quantity, the radial excitation energy, $\Delta \mathrm{E}_{\mathrm{r}}$ expected to be mildly dependent on the diquark mass: $\mathrm{E}_{\mathrm{r}}(\mathrm{cq}) \sim$ $\sim \mathrm{E}_{\mathrm{r}}$ (cs)


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& m_{[c q]}=1980 \mathrm{MeV} \\
& \kappa_{c q}=67 \mathrm{MeV}
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$$
\Delta E_{r}(c q)=530 \mathrm{MeV}
$$




## 4. J/ $\Psi-\varphi$ structures and S-wave tetraquarks


below threshold

- $\mathbf{0}^{++}$?
- $2^{++}$?
- 2 unresolved, almost degenerate lines with $0^{++}+2^{++}$??


## 4. J/Ч- $\varphi$ structures and S-wave tetraquarks



## 4. J/Ч- $\varphi$ structures and S-wave tetraquarks

$$
\begin{aligned}
& \overline{\mathrm{X}(4700)} \quad \bullet \quad \Delta m=m_{c s}-m_{c q}=129 \mathrm{MeV} \text {; } \\
& \kappa_{s c}=50 \mathrm{MeV}\left(\kappa_{q c}=67 \mathrm{MeV}\right) \\
& \text { radial excit. }=460 \mathrm{MeV} \\
& {[Z(4430)-Z(3900)=530 \mathrm{MeV}]} \\
& \text { NOTE : } \\
& X(4140)-X(3872) \sim 270 \mathrm{MeV} \text {; } \\
& \phi(1020)-\rho(770) \sim 244 \mathrm{MeV} \\
& X(4274) \text { cannot be } 1^{++} \\
& \text {- } 2 \text { unresolved, almost degenerate } \\
& \text { Decay modes of } \mathbf{J}^{\mathbf{P}}=\mathbf{1}^{+}, \mathrm{C}=-1 \text { : } \\
& s_{c \bar{c}}=1: J / \Psi+\eta, \chi_{c}+\eta(P-\text { wave }) \\
& s_{c \bar{c}}=0: \eta_{c}+\phi, h_{c}+\phi(P-\text { wave })
\end{aligned}
$$

## Summary Table

| radial | particle | $J^{P C}$ | input | predicted | notes |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $1 S$ | $X_{0}$ | $0^{++}$ |  | 4040 | below the $J / \Psi \phi$ threshold |
| $1 S$ | $X$ | $1^{++}$ | 4140 |  | -- |
| $1 S$ | $X^{(1)}$ | $1^{+-}$ |  | 4140 | decays in $\chi_{c} \phi ?(M \geq 4435)$ |
| $1 S$ | $X^{(2)}$ | $1^{+-}$ |  | 4240 | decays in $\eta_{c} \phi ?$ |
| $1 S$ | $X_{0}^{\prime}$ | $0^{++}$ |  | 4240 | part of 4274 structure? |
| $1 S$ | $X_{2}^{\prime}$ | $2^{++}$ |  | 4240 | part of 4274 structure? |
| $2 S$ | $X_{0}$ | $0^{++}$ | 4500 |  | -- |
| $2 S$ | $X$ | $1^{++}$ |  | 4600 | -- |
| $2 S$ | $X^{(1)}$ | $1^{+-}$ |  | 4600 | $S_{c \bar{c}}=1$ decays in $\chi_{c} \phi ?$ |
| $2 S$ | $X^{(2)}$ | $1^{+-}$ |  | 4700 | $S_{c \bar{c}}=0$ decays in $\eta_{c}(2 S) \phi ?$ |
| $2 S$ | $X_{0}^{\prime}$ | $0^{++}$ | 4700 |  | -- |
| $2 S$ | $X_{2}^{\prime}$ | $2^{++}$ |  | 4700 | decays in $J / \Psi \phi, \psi^{\prime} \phi ?$ |

Table 1: Input and predicted masses for $1 S$ and $2 S c s$ tetraquarks.

## 5. Variations on the theme

- J/Ч- $\varphi$ spectrum obtained with meson\&baryon spin-spin parameters does not fit with experiment
N.V. Drenska, R. Faccini and A. D. Polosa, Phys. Rev. D 79 (2009) 077502
- QCD sum rules with tetra quark currents tried with some success and support $X(4500)$ and $X(4700)$ to be higher excitations, radial or D-wave
Z. G. Wang, arXiv:1607.00701 [hep-ph];
- flavour $\operatorname{SU}(3)$ nonet including $\mathrm{J} / \Psi-\varphi$ has been considered in:
R. Zhu, Phys Rev. D 94 (2016) 054009
- diquarks in color 6 have been considered by several authors
J. Wu et al., arXiv:1608.07900 [hep-ph]
- if at all bound, tetraquarks made by color 6 diquarks would double the spectrum
- an option if $\mathrm{X}(4270)$ turns out to be a pure $1^{++}$resonance?
- basic masses of diquark in color 3 and 6 must be different: $X(4270)-X(4140)$ is not due only to spin-spin interactions and will be essentially incalculable.


## what about the strange members of the nonet?

- We expect strangeness $= \pm 1$ tetra quarks: $X_{\bar{s}}=[c q][\bar{c} \bar{s}] ; X_{s}=[c s][\bar{c} \bar{q}]$
- partners of X(4140) should decay in: $J / \Psi+K^{*} / \bar{K}^{*} \rightarrow \mu^{+} \mu^{-}+\pi+K_{S}$
- while partners of $\mathrm{C}=-1$ states decay in: $\quad J / \Psi+K / \bar{K} \rightarrow \mu^{+} \mu^{-}+K_{S}$
- Mass can be estimated at: $M\left(X_{s}\right) \sim \frac{4140+3872}{2} \sim 4006$

$$
\left[M(J / \Psi)+M\left(K^{*}\right) \sim 4000\right]
$$

- are they visible at LHCb/BELLE/BES III?


## 6.Conclusions

- J/Ч- $\varphi$ resonances go well with simple, $S$-wave, tetraquarks...
- ....except for the puzzling $1^{++}$duplication of X(4140) and X(4270)
- model predicts two almost degenerate states around 4240: could X(4270) be a complex of two unresolved lines with $\mathrm{J}^{\mathrm{PC}}$ $=0^{++}, 2^{++} \quad$ ???????????????????????
- if duplication of $1^{++}$is confirmed, diquarks with color 6 could be an option, but....are there other similar duplications?
- could the $\mathrm{X}(3940)$ be $1^{++}$as well ??
- Experimental clarification is essential for progress

