

NLO QCD Corrections to J/ψ Inclusive Production in Photon-Photon Collision

Based on arXiv: 1608.06231

陈自强
陈龙斌
乔从丰

中国科学院大学 物理学院

2016-11-03

Outline

- Introduction
- J/ψ Inclusive Production in Photon-Photon Collision
- NLO QCD Corrections to $\gamma\gamma \rightarrow J/\psi + c\bar{c}$
- $J/\psi + ggg$ Final State Subprocess
- Summary

Introduction

Introduction

- In non-relativistic QCD(NRQCD), heavy quarkonium is treated as a nonrelativistic system. Then we get a hierarchy of energy scales:

$$(M_Q v^2)^2 \ll (M_Q v)^2 \ll M_Q^2$$

Introduction

- In non-relativistic QCD(NRQCD), heavy quarkonium is treated as a nonrelativistic system. Then we get a hierarchy of energy scales:

$$(M_Q v^2)^2 \ll (M_Q v)^2 \ll M_Q^2$$

- Based on this hierarchy, the quarkonium production and decay amplitudes can be factorized into short- and long-distance sectors:

$$d\sigma(H + X) = \sum d\hat{\sigma}(Q\bar{Q} + X)\langle\mathcal{O}_H\rangle$$

Introduction

- In non-relativistic QCD(NRQCD), heavy quarkonium is treated as a nonrelativistic system. Then we get a hierarchy of energy scales:

$$(M_Q v^2)^2 \ll (M_Q v)^2 \ll M_Q^2$$

- Based on this hierarchy, the quarkonium production and decay amplitudes can be factorized into short- and long-distance sectors:

$$d\sigma(H + X) = \sum d\hat{\sigma}(Q\bar{Q} + X) \langle \mathcal{O}_H \rangle$$

- NRQCD factorization model is not so intuitive as color-singlet(CS) model for the introduction of color-octet(CO) mechanism($Q\bar{Q}$ can be CO state).

Introduction

- In non-relativistic QCD(NRQCD), heavy quarkonium is treated as a nonrelativistic system. Then we get a hierarchy of energy scales:

$$(M_Q v^2)^2 \ll (M_Q v)^2 \ll M_Q^2$$

- Based on this hierarchy, the quarkonium production and decay amplitudes can be factorized into short- and long-distance sectors:

$$d\sigma(H + X) = \sum d\hat{\sigma}(Q\bar{Q} + X) \langle \mathcal{O}_H \rangle$$

- NRQCD factorization model is not so intuitive as color-singlet(CS) model for the introduction of color-octet(CO) mechanism($Q\bar{Q}$ can be CO state).
- The test of NRQCD factorization is a exigent task in quarkonium physics. And the study of J/ψ inclusive production through $\gamma\gamma$ collision at LEP II unravel a rather confusing pattern.

J/ψ Inclusive Production in Photon-Photon Collision

- Experimental Data
 - LEP II $\gamma\gamma \rightarrow J/\psi + X$ Data
- Theoretical Calculations
 - Three Classes of Subprocesses
 - The LO Calculations
 - The NLO Corrections
 - $J/\psi + c\bar{c}$ Final State Subprocess
 - Full NLO Corrections
 - About Our Work

LEP II $\gamma\gamma \rightarrow J/\psi + X$ Data

In 2001, the DELPHI Collaboration presented preliminary data on the J/ψ inclusive cross section in photon-photon collision ($e^+e^- \rightarrow e^+e^- J/\psi + X$) at LEP II.

LEP II $\gamma\gamma \rightarrow J/\psi + X$ Data

In 2001, the DELPHI Collaboration presented preliminary data on the J/ψ inclusive cross section in photon-photon collision ($e^+e^- \rightarrow e^+e^- J/\psi + X$) at LEP II.

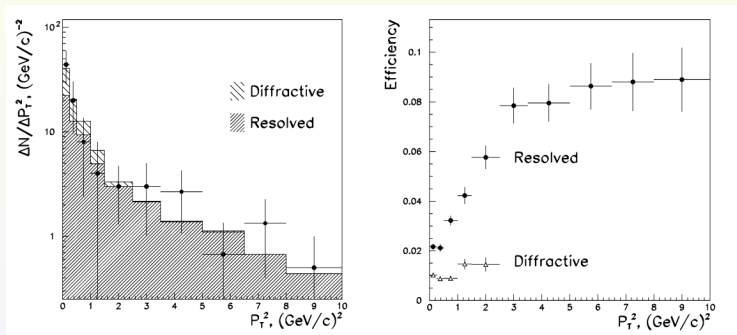


Figure: “DELPHI Collaboration, Phys. Lett. B 565(2003)76-86”.

The requirement of at least 4 reconstructed tracks suppress the $J/\psi + \gamma$ final state.

Three Classes of Subprocesses

There are 3 classes of subprocesses for $\gamma\gamma \rightarrow J/\psi + X$:

Three Classes of Subprocesses

There are 3 classes of subprocesses for $\gamma\gamma \rightarrow J/\psi + X$:

i direct process:

$$\gamma\gamma \rightarrow c\bar{c} \left[{}^3S_1^{(8)} \right] g,$$

$$\gamma\gamma \rightarrow c\bar{c} \left[{}^3S_1^{(1)} \right] \gamma \quad (\text{suppressed at at LEP II})$$

Three Classes of Subprocesses

There are 3 classes of subprocesses for $\gamma\gamma \rightarrow J/\psi + X$:

i direct process:

$$\gamma\gamma \rightarrow c\bar{c} \left[{}^3S_1^{(8)} \right] g,$$

$$\gamma\gamma \rightarrow c\bar{c} \left[{}^3S_1^{(1)} \right] \gamma \quad (\text{suppressed at at LEP II})$$

ii single-resolved process:

$$\gamma g \rightarrow c\bar{c} \left[{}^3S_1^{(1)} \right] g, c\bar{c}[8]g,$$

$$\gamma q \rightarrow c\bar{c}[8]q$$

Three Classes of Subprocesses

There are 3 classes of subprocesses for $\gamma\gamma \rightarrow J/\psi + X$:

i direct process:

$$\gamma\gamma \rightarrow c\bar{c} \left[{}^3S_1^{(8)} \right] g,$$

$$\gamma\gamma \rightarrow c\bar{c} \left[{}^3S_1^{(1)} \right] \gamma \quad (\text{suppressed at at LEP II})$$

ii single-resolved process:

$$\gamma g \rightarrow c\bar{c} \left[{}^3S_1^{(1)} \right] g, c\bar{c}[8]g,$$

$$\gamma q \rightarrow c\bar{c}[8]q$$

iii double-resolved process:

$$gg \rightarrow c\bar{c} \left[{}^3S_1^{(1)} \right] g, c\bar{c} \left[{}^3P_J^{(1)} \right] g, c\bar{c}[8]g,$$

$$gq \rightarrow c\bar{c} \left[{}^3P_J^{(1)} \right] q, c\bar{c}[8]q,$$

$$q\bar{q} \rightarrow c\bar{c} \left[{}^3S_1^{(8)} \right]$$

The LO Calculations

The LO calculations for direct process was finished early[1]. And full LO calculations was finished soon after the presentation of DELPHI data[2].

The NLO Corrections

Before the global NLO analysis performed, there are several NLO corrections for some subprocess:

The NLO Corrections

Before the global NLO analysis performed, there are several NLO corrections for some subprocess:

- 1 M.Klasen, B.A.Kniehl *et al.*, Nucl. Phys. B 713(2005)487-521:

NLO corrections to $\gamma\gamma \rightarrow Q\bar{Q} + g$ (TESLA)

The NLO Corrections

Before the global NLO analysis performed, there are several NLO corrections for some subprocess:

- 1 M.Klasen, B.A.Kniehl *et al.*, Nucl. Phys. B 713(2005)487-521:

NLO corrections to $\gamma\gamma \rightarrow Q\bar{Q} + g$ (TESLA)

- 2 M.Klasen, B.A.Kniehl *et al.*, Phys. Rev. D 71, 014016 (2005):

NLO corrections to $\gamma\gamma \rightarrow Q\bar{Q} + \gamma$ (TESLA)

The NLO Corrections

Before the global NLO analysis performed, there are several NLO corrections for some subprocess:

- 1 M.Klasen, B.A.Kniehl *et al.*, Nucl. Phys. B 713(2005)487-521:

NLO corrections to $\gamma\gamma \rightarrow Q\bar{Q} + g$ (TESLA)

- 2 M.Klasen, B.A.Kniehl *et al.*, Phys. Rev. D 71, 014016 (2005):

NLO corrections to $\gamma\gamma \rightarrow Q\bar{Q} + \gamma$ (TESLA)

- 3 M.Kramer, Nucl. Phys. B 459(1996)3-50:

NLO corrections to $\gamma g \rightarrow Q\bar{Q} + g$ (HERA)

The NLO Corrections

Before the global NLO analysis performed, there are several NLO corrections for some subprocess:

- 1 M.Klasen, B.A.Kniehl *et al.*, Nucl. Phys. B 713(2005)487-521:

NLO corrections to $\gamma\gamma \rightarrow Q\bar{Q} + g$ (TESLA)

- 2 M.Klasen, B.A.Kniehl *et al.*, Phys. Rev. D 71, 014016 (2005):

NLO corrections to $\gamma\gamma \rightarrow Q\bar{Q} + \gamma$ (TESLA)

- 3 M.Kramer, Nucl. Phys. B 459(1996)3-50:

NLO corrections to $\gamma g \rightarrow Q\bar{Q} + g$ (HERA)

- 4

The NLO Corrections

Before the global NLO analysis performed, there are several NLO corrections for some subprocess:

- 1 M.Klasen, B.A.Kniehl *et al.*, Nucl. Phys. B 713(2005)487-521:

NLO corrections to $\gamma\gamma \rightarrow Q\bar{Q} + g$ (TESLA)

- 2 M.Klasen, B.A.Kniehl *et al.*, Phys. Rev. D 71, 014016 (2005):

NLO corrections to $\gamma\gamma \rightarrow Q\bar{Q} + \gamma$ (TESLA)

- 3 M.Kramer, Nucl. Phys. B 459(1996)3-50:

NLO corrections to $\gamma g \rightarrow Q\bar{Q} + g$ (HERA)

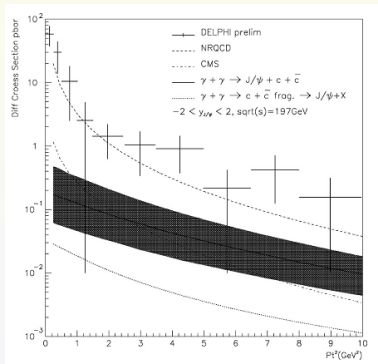
- 4

Although none of these analyses is performed in LEP II condition, people infer that similar corrections should not change the former conclusion, and shift their focus to a new subprocess: $J/\psi + c\bar{c}$ final state subprocess.

$J/\psi + c\bar{c}$ Final State Subprocess

The direct production of $J/\psi + c\bar{c}$, namely $\gamma\gamma \rightarrow J/\psi + c\bar{c}$ process, prove to be the dominant CS process[1].

A comparison is also made between pure fragmentation result and full calculation in [1]. It is found that this process cannot be mimicked by simple fragmentation scheme.



After including this process, the prediction of CS is doubled, but still insufficient to explain experiment data.

[1] C.F. Qiao and J.X. Wang, Phys. Rev. D 69, 014015(2004).

$J/\psi + c\bar{c}$ Final State Subprocess

Further studies:

$J/\psi + c\bar{c}$ Final State Subprocess

Further studies:

1 R. Li and K.-T. Chao, *Phys. Rev. D* **79**, 114020(2009):

$$\gamma g \rightarrow J/\psi + c\bar{c}, q\bar{q} \rightarrow J/\psi + c\bar{c}, gg \rightarrow J/\psi + c\bar{c},$$

$$\gamma\gamma \rightarrow J/\psi + c\bar{c}$$

$J/\psi + c\bar{c}$ Final State Subprocess

Further studies:

- 1 R. Li and K.-T. Chao, Phys. Rev. D 79, 114020(2009):

$$\gamma g \rightarrow J/\psi + c\bar{c}, q\bar{q} \rightarrow J/\psi + c\bar{c}, gg \rightarrow J/\psi + c\bar{c},$$

$$\gamma\gamma \rightarrow J/\psi + c\bar{c}$$

- 2 G. Chen, X.-G. Wu, *et al.* Phys. Rev. D 90, 034004(2014),

Z. Sun, X.-G. Wu and H.-F. Zhang, Phys. Rev. D 92, 074021(2015):

$$\gamma\gamma \rightarrow [Q\bar{Q}] + Q\bar{Q} \text{ at ILC}$$

$J/\psi + c\bar{c}$ Final State Subprocess

Further studies:

- 1 R. Li and K.-T. Chao, Phys. Rev. D 79, 114020(2009):

$$\gamma g \rightarrow J/\psi + c\bar{c}, q\bar{q} \rightarrow J/\psi + c\bar{c}, gg \rightarrow J/\psi + c\bar{c},$$

$$\gamma\gamma \rightarrow J/\psi + c\bar{c}$$

- 2 G. Chen, X.-G. Wu, *et al.* Phys. Rev. D 90, 034004(2014),

Z. Sun, X.-G. Wu and H.-F. Zhang, Phys. Rev. D 92, 074021(2015):

$$\gamma\gamma \rightarrow [Q\bar{Q}] + Q\bar{Q} \text{ at ILC}$$

- 3 M. Klasen and J.P.Lansberg, Nucl. Phys. B 179-180 (2008) 226-231:

photon-photon collisions at LHC

$J/\psi + c\bar{c}$ Final State Subprocess

Further studies:

- 1 R. Li and K.-T. Chao, Phys. Rev. D 79, 114020(2009):

$$\gamma g \rightarrow J/\psi + c\bar{c}, q\bar{q} \rightarrow J/\psi + c\bar{c}, gg \rightarrow J/\psi + c\bar{c},$$

$$\gamma\gamma \rightarrow J/\psi + c\bar{c}$$

- 2 G. Chen, X.-G. Wu, *et al.* Phys. Rev. D 90, 034004(2014),

Z. Sun, X.-G. Wu and H.-F. Zhang, Phys. Rev. D 92, 074021(2015):

$$\gamma\gamma \rightarrow [Q\bar{Q}] + Q\bar{Q} \text{ at ILC}$$

- 3 M. Klasen and J.P.Lansberg, Nucl. Phys. B 179-180 (2008) 226-231:

photon-photon collisions at LHC

- 4

$J/\psi + c\bar{c}$ Final State Subprocess

Further studies:

- 1 R. Li and K.-T. Chao, *Phys. Rev. D* **79**, 114020(2009):

$$\gamma g \rightarrow J/\psi + c\bar{c}, q\bar{q} \rightarrow J/\psi + c\bar{c}, gg \rightarrow J/\psi + c\bar{c},$$

$$\gamma\gamma \rightarrow J/\psi + c\bar{c}$$

- 2 G. Chen, X.-G. Wu, *et al.* *Phys. Rev. D* **90**, 034004(2014),

Z. Sun, X.-G. Wu and H.-F. Zhang, *Phys. Rev. D* **92**, 074021(2015):

$$\gamma\gamma \rightarrow [Q\bar{Q}] + Q\bar{Q} \text{ at ILC}$$

- 3 M. Klasen and J.P.Lansberg, *Nucl. Phys. B* **179-180** (2008) 226-231:

photon-photon collisions at LHC

- 4

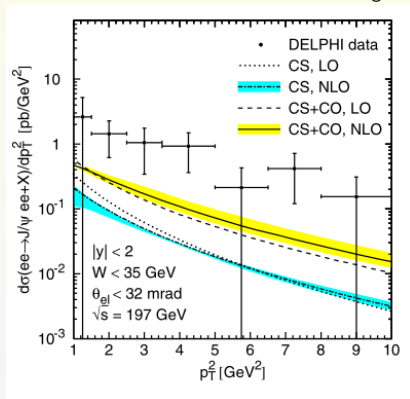
These researches indicate that the contributions of single-resolved and double-resolved are significantly less important than the direct one with in the CS prescription. Their contributions do little to reduce the gap between CS prediction and experiment data.

Full NLO analysis

Kniehl (one of the author of full LO analysis) et al. then carried out a full NLO NRQCD analysis and found the situation become more enigmatic[1].

Full NLO analysis

Kniehl (one of the author of full LO analysis) et al. then carried out a full NLO NRQCD analysis and found the situation become more enigmatic[1].



Full NLO analysis

Kniehl (one of the author of full LO analysis) et al. then carried out a full NLO NRQCD analysis and found the situation become more enigmatic[1].

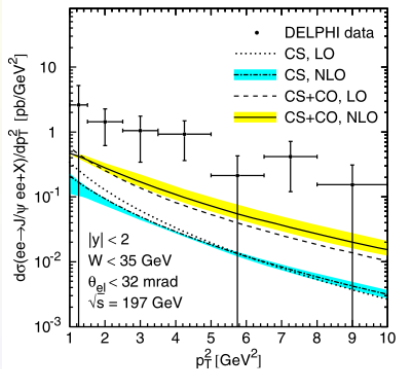


TABLE I. NLO fit results for the J/ψ CO LDMEs.

$\langle \mathcal{O}^{J/\psi}({}^1S_0^{[8]}) \rangle$	$(4.97 \pm 0.44) \times 10^{-2} \text{ GeV}^3$
$\langle \mathcal{O}^{J/\psi}({}^3S_1^{[8]}) \rangle$	$(2.24 \pm 0.59) \times 10^{-3} \text{ GeV}^3$
$\langle \mathcal{O}^{J/\psi}({}^3P_0^{[8]}) \rangle$	$(-1.61 \pm 0.20) \times 10^{-2} \text{ GeV}^5$

Owing to the negative value of $\langle \mathcal{O}^{J/\psi}({}^3P_0^{[8]}) \rangle$, the DELPHI data cannot be reproduced even with the CO contributions.

Full NLO analysis

Kniehl (one of the author of full LO analysis) et al. then carried out a full NLO NRQCD analysis and found the situation become more enigmatic[1].

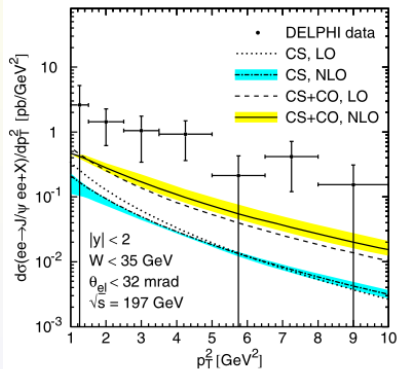


TABLE I. NLO fit results for the J/ψ CO LDMEs.

$\langle \mathcal{O}^{J/\psi}({}^1S_0^{[8]}) \rangle$	$(4.97 \pm 0.44) \times 10^{-2} \text{ GeV}^3$
$\langle \mathcal{O}^{J/\psi}({}^3S_1^{[8]}) \rangle$	$(2.24 \pm 0.59) \times 10^{-3} \text{ GeV}^3$
$\langle \mathcal{O}^{J/\psi}({}^3P_0^{[8]}) \rangle$	$(-1.61 \pm 0.20) \times 10^{-2} \text{ GeV}^5$

Owing to the negative value of $\langle \mathcal{O}^{J/\psi}({}^3P_0^{[8]}) \rangle$, the DELPHI data cannot be reproduced even with the CO contributions.

Kniehl et al. started to emphasise that the data sample only consisted of 16 events above $p_t > 1\text{GeV}$.

About our work

About our work

- However, before we give the definite conclusion, we should stretch our calculation as far as we can. To this aim, we calculate the tedious NLO QCD corrections to $\gamma\gamma \rightarrow J/\psi + c\bar{c}$ process.

About our work

- However, before we give the definite conclusion, we should stretch our calculation as far as we can. To this aim, we calculate the tedious NLO QCD corrections to $\gamma\gamma \rightarrow J/\psi + c\bar{c}$ process.
- This process deserves our attention, because
 - i It's the dominant CS process in the J/ψ inclusive production through $\gamma\gamma$ collision.
 - ii $J/\psi + c\bar{c}$ is the experimentally distinguishable final state.

About our work

- However, before we give the definite conclusion, we should stretch our calculation as far as we can. To this aim, we calculate the tedious NLO QCD corrections to $\gamma\gamma \rightarrow J/\psi + c\bar{c}$ process.
- This process deserves our attention, because
 - i It's the dominant CS process in the J/ψ inclusive production through $\gamma\gamma$ collision.
 - ii $J/\psi + c\bar{c}$ is the experimentally distinguishable final state.
- This is the first true NLO calculation of 2 to 3 inclusive process in heavy quarkonium production.

NLO QCD Corrections to $\gamma\gamma \rightarrow J/\psi + c\bar{c}$

- Formalism and Calculation
- Numerical Results

Formalism and Calculation

Formalism and Calculation

- The overall differential cross section can be written as

$$d\sigma = \int dx_1 dx_2 f_\gamma(x_1) f_\gamma(x_2) d\hat{\sigma}(\gamma\gamma \rightarrow c\bar{c}[^3S_1] + c\bar{c}) \langle \mathcal{O}^{J/\psi} (^3S_1) \rangle$$

Formalism and Calculation

- The overall differential cross section can be written as

$$d\sigma = \int dx_1 dx_2 f_\gamma(x_1) f_\gamma(x_2) d\hat{\sigma}(\gamma\gamma \rightarrow c\bar{c}[^3S_1] + c\bar{c}) \langle \mathcal{O}^{J/\psi} (^3S_1) \rangle$$

- To NLO calculation, the cross section is

$$d\hat{\sigma}(\gamma\gamma \rightarrow c\bar{c}[^3S_1] + c\bar{c}) = d\hat{\sigma}_{born} + d\hat{\sigma}_{virtual} + d\hat{\sigma}_{real} + O(\alpha^2 \alpha_s^4)$$

Formalism and Calculation

- The overall differential cross section can be written as

$$d\sigma = \int dx_1 dx_2 f_\gamma(x_1) f_\gamma(x_2) d\hat{\sigma}(\gamma\gamma \rightarrow c\bar{c}[^3S_1] + c\bar{c}) \langle \mathcal{O}^{J/\psi} (^3S_1) \rangle$$

- To NLO calculation, the cross section is

$$d\hat{\sigma}(\gamma\gamma \rightarrow c\bar{c}[^3S_1] + c\bar{c}) = d\hat{\sigma}_{born} + d\hat{\sigma}_{virtual} + d\hat{\sigma}_{real} + O(\alpha^2 \alpha_s^4)$$

- The quarkonium spin projection operator is

$$v(P/2)\bar{u}(P/2) = \frac{1}{4\sqrt{2}E(E+2m_c)} \left(\frac{\not{P}}{2} - m_c\right) \not{\epsilon}_S^* (\not{P} + 2E) \left(\frac{\not{P}}{2} + m_c\right)$$

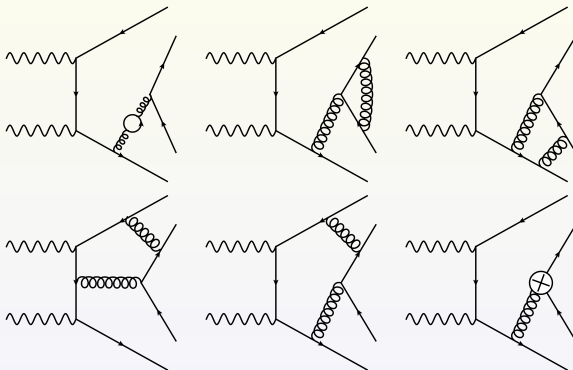
Formalism and Calculation

- Virtual corrections (582 Diagrams)

Formalism and Calculation

■ Virtual corrections (582 Diagrams)

Selfenergies, Triangles, Boxes, Pentagons, Hexagons and Counter terms



Masses of Feynman integrals are evaluated.

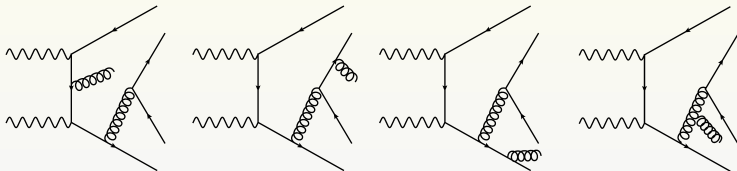
Formalism and Calculation

- Real corrections (200 Diagrams)

Formalism and Calculation

■ Real corrections (200 Diagrams)

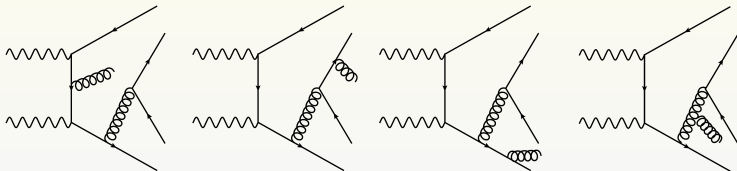
$$\gamma\gamma \rightarrow J/\psi + c\bar{c} + g$$



Formalism and Calculation

■ Real corrections (200 Diagrams)

$$\gamma\gamma \rightarrow J/\psi + c\bar{c} + g$$



According to B. W. Harris and J. F. Owens, “Two cutoff phase space slicing method”,
 Phys. Rev. D 65, 094032 (2002),

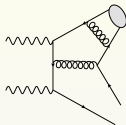
$$d\sigma_{real} = d\sigma_{soft}^{IR} \Big|_{p_g^0 < \delta \frac{\sqrt{s}}{2}} + d\sigma_{hard}^{IR-free} \Big|_{p_g^0 > \delta \frac{\sqrt{s}}{2}}$$

Formalism and Calculation

- Cancellation of singularities

Formalism and Calculation

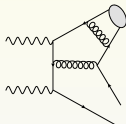
- Cancellation of singularities
 - i Coulombic singularities are attributed to NRQCD long-distance matrix elements



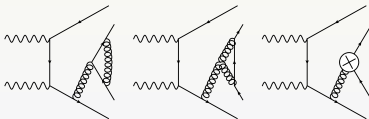
Formalism and Calculation

■ Cancellation of singularities

- i Coulombic singularities are attributed to NRQCD long-distance matrix elements



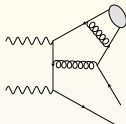
- ii Ultraviolet singularities are canceled by renormalization



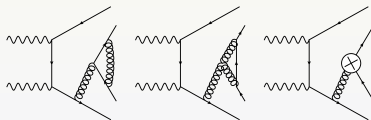
Formalism and Calculation

■ Cancellation of singularities

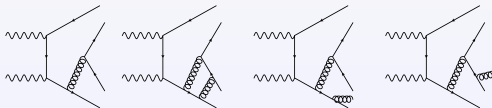
- i Coulombic singularities are attributed to NRQCD long-distance matrix elements



- ii Ultraviolet singularities are canceled by renormalization



- iii Infrared singularities involved in virtual corrections and real emissions are canceled each other



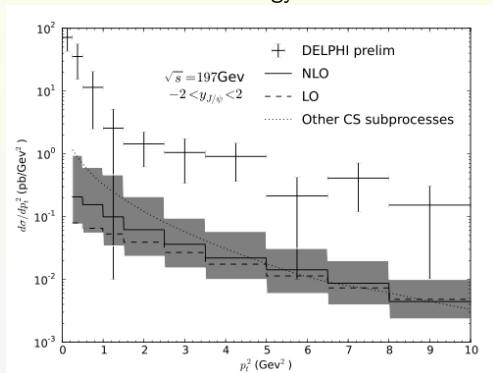
Numerical Results

Numerical Results

- LEP II collider energy

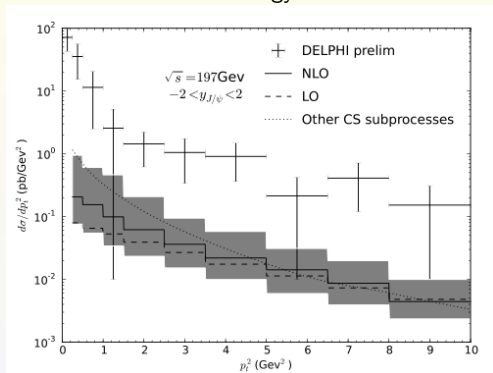
Numerical Results

■ LEP II collider energy



Numerical Results

■ LEP II collider energy

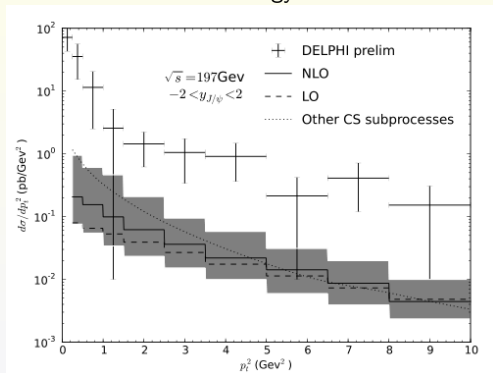


The shaded band is the NLO result of $\gamma\gamma \rightarrow J/\psi + c\bar{c}$ process with its upper bound obtained at $r = 0.5$, $m_c = 1.4 \text{ GeV}$ and lower bound at $r = 2$, $m_c = 1.6 \text{ GeV}$. The solid and dashed lines represent the NLO and LO results with $r = 1$, $m_c = 1.5 \text{ GeV}$.

$$\mu = r\sqrt{4m_c^2 + p_t^2}$$

Numerical Results

■ LEP II collider energy



The shaded band is the NLO result of $\gamma\gamma \rightarrow J/\psi + c\bar{c}$ process with its upper bound obtained at $r = 0.5$, $m_c = 1.4 \text{ GeV}$ and lower bound at $r = 2$, $m_c = 1.6 \text{ GeV}$. The solid and dashed lines represent the NLO and LO results with $r = 1$, $m_c = 1.5 \text{ GeV}$.

$$\mu = r\sqrt{4m_c^2 + p_t^2}$$

The NLO correction is moderate but not big enough to remove the huge discrepancy between theoretical prediction and experimental observation.

Numerical Results

- LEP II collider energy

While integrated over the range $1 \leq p_t^2 \leq 10\text{GeV}^2$,

$\sigma(\text{pb})$	$m_c = 1.4\text{GeV}$	$m_c = 1.5\text{GeV}$	$m_c = 1.6\text{GeV}$
$r = 0.5$	0.766(0.436)	0.459(0.283)	0.299(0.187)
$r = 1$	0.363(0.236)	0.227(0.156)	0.152(0.105)
$r = 2$	0.216(0.152)	0.138(0.101)	0.093(0.069)

Table: NLO(LO) results of total cross sections with different renormalization scale and charm quark mass. The K factor of central value is about 1.46.

Numerical Results

- LEP II collider energy

While integrated over the range $1 \leq p_t^2 \leq 10\text{GeV}^2$,

$\sigma(\text{pb})$	$m_c = 1.4\text{GeV}$	$m_c = 1.5\text{GeV}$	$m_c = 1.6\text{GeV}$
$r = 0.5$	0.766(0.436)	0.459(0.283)	0.299(0.187)
$r = 1$	0.363(0.236)	0.227(0.156)	0.152(0.105)
$r = 2$	0.216(0.152)	0.138(0.101)	0.093(0.069)

Table: NLO(LO) results of total cross sections with different renormalization scale and charm quark mass. The K factor of central value is about 1.46.

The result of DELPHI and other CS processes read $(6.4 \pm 2.0)\text{pb}$ and $0.39_{-0.09}^{+0.16}\text{pb}$, respectively. And the DELPHI data can not be reproduced by CS model (according to [1] NRQCD neither) after including our corrections.

[1] M.Butenschoen and B.A.Kniehl, Phys. Rev. D 84, 051501(2011).

Numerical Results

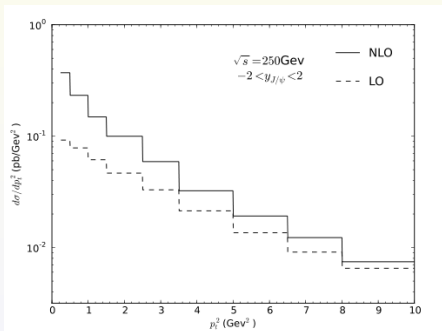
- CEPC collider energy

In the future, the e^+e^- collider CEPC will run at $\sqrt{s} = 250\text{GeV}$,

Numerical Results

■ CEPC collider energy

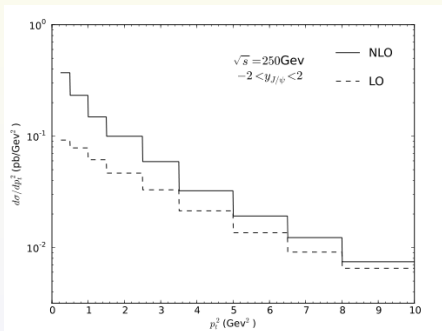
In the future, the e^+e^- collider CEPC will run at $\sqrt{s} = 250\text{GeV}$,



Numerical Results

■ CEPC collider energy

In the future, the e^+e^- collider CEPC will run at $\sqrt{s} = 250\text{GeV}$,

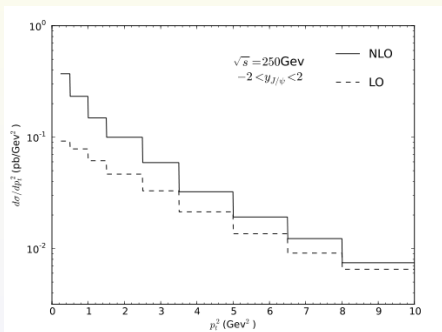


The p_t^2 distribution of $J/\psi + c\bar{c}$ production through $\gamma\gamma$ collision at CEPC. The solid and dashed lines represent the NLO and LO results with $r = 1$, $m_c = 1.5\text{GeV}$.

Numerical Results

■ CEPC collider energy

In the future, the e^+e^- collider CEPC will run at $\sqrt{s} = 250\text{GeV}$,



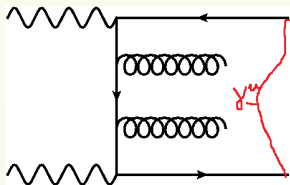
The p_t^2 distribution of $J/\psi + c\bar{c}$ production through $\gamma\gamma$ collision at CEPC. The solid and dashed lines represent the NLO and LO results with $r = 1$, $m_c = 1.5\text{GeV}$. Integrated over $p_t^2 \geq 1\text{GeV}^2$, the total NLO(LO) cross section is $0.432\text{pb}(0.245\text{pb})$. The K factor is about 1.76.

$J/\psi + ggg$ Final State Subprocess

- $\gamma\gamma \rightarrow J/\psi + gg$ Process and Furry's Theorem
- Feynman Diagrams
- Numerical Results

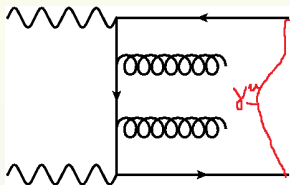
$\gamma\gamma \rightarrow J/\psi + gg$ Process and Furry's Theorem

Process $\gamma\gamma \rightarrow J/\psi + gg$ is of the same order as $\gamma\gamma \rightarrow J/\psi + c\bar{c}$. But this process is forbidden by Furry's theorem based on charge conjugation invariance.

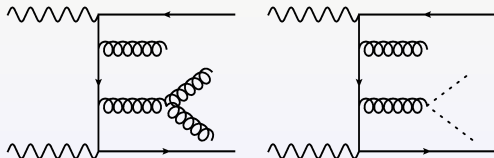


$\gamma\gamma \rightarrow J/\psi + gg$ Process and Furry's Theorem

Process $\gamma\gamma \rightarrow J/\psi + gg$ is of the same order as $\gamma\gamma \rightarrow J/\psi + c\bar{c}$. But this process is forbidden by Furry's theorem based on charge conjugation invariance.



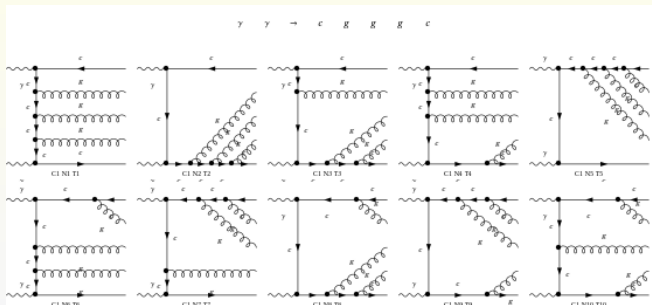
The same as



This was also verified by explicit calculation.

Feynman Diagrams

There are 120 Feynman diagrams for $\gamma\gamma \rightarrow J/\psi + ggg$.



Other diagrams can be obtained by exchange gluons and photons.

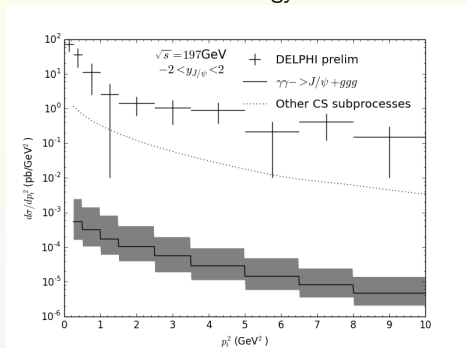
Numerical Results

Numerical Results

- LEP II collider energy

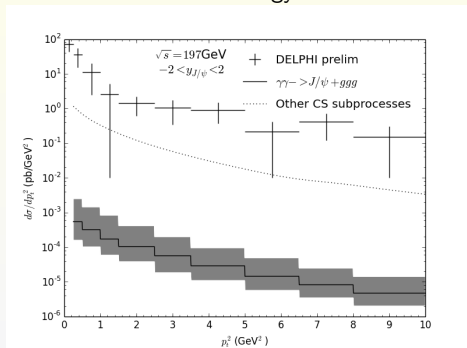
Numerical Results

■ LEP II collider energy



Numerical Results

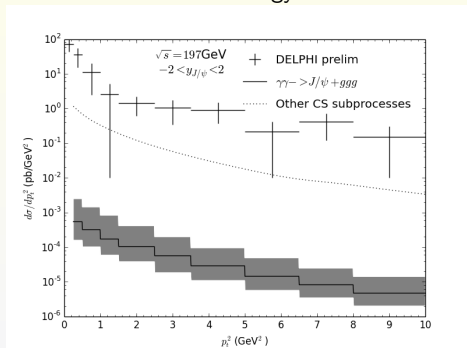
■ LEP II collider energy



The shaded band is the cross section of $\gamma\gamma \rightarrow J/\psi + ggg$ process with its upper bound obtained at $r = 0.5$, $m_c = 1.4 \text{ GeV}$ and lower bound at $r = 2$, $m_c = 1.6 \text{ GeV}$. The solid lines represent the results with $r = 1$, $m_c = 1.5 \text{ GeV}$.

Numerical Results

LEP II collider energy



The shaded band is the cross section of $\gamma\gamma \rightarrow J/\psi + ggg$ process with its upper bound obtained at $r = 0.5$, $m_c = 1.4 \text{ GeV}$ and lower bound at $r = 2$, $m_c = 1.6 \text{ GeV}$. The solid lines represent the results with $r = 1$, $m_c = 1.5 \text{ GeV}$.

Integrated over the range $1 \leq p_t^2 \leq 10 \text{ GeV}^2$,

$\sigma(\text{fb})$	$m_c = 1.4 \text{ GeV}$	$m_c = 1.5 \text{ GeV}$	$m_c = 1.6 \text{ GeV}$
$r = 0.5$	1.32	0.82	0.54
$r = 1$	0.52	0.33	0.22
$r = 2$	0.26	0.17	0.12

Numerical Results

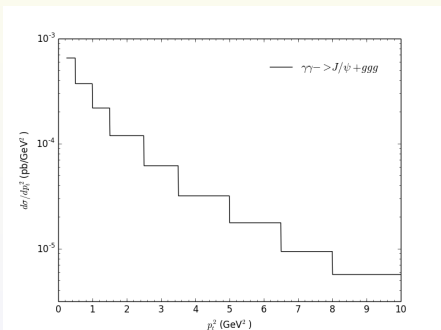
- CEPC collider energy

At the e^+e^- collider CEPC,

Numerical Results

- CEPC collider energy

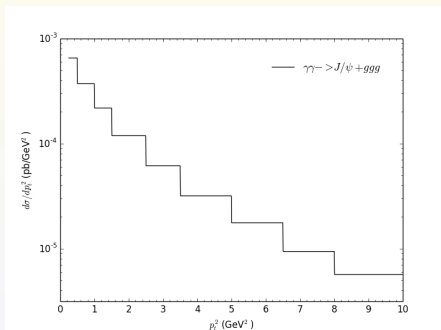
At the e^+e^- collider CEPC,



Numerical Results

■ CEPC collider energy

At the e^+e^- collider CEPC,



The p_t^2 distribution of $J/\psi + ggg$ production through $\gamma\gamma$ collision at CEPC. The solid lines represent the cross section with $r = 1$, $m_c = 1.5\text{GeV}$. Integrated over $p_t^2 \geq 1\text{GeV}^2$, the total cross section is 0.39fb.

Summary

Summary

- We have calculated the NLO QCD corrections to $\gamma\gamma \rightarrow J/\psi + c\bar{c}$ process. This is the first truly NLO calculation of 2 to 3 inclusive process for heavy quarkonium production.

Summary

- We have calculated the NLO QCD corrections to $\gamma\gamma \rightarrow J/\psi + c\bar{c}$ process. This is the first truly NLO calculation of 2 to 3 inclusive process for heavy quarkonium production.
- Numerical results shows that the cross section of this process at LEP II can be enhanced with a K factor of about 1.46. And the discrepancy between theoretical prediction and experiment data reduced accordingly.

Summary

- We have calculated the NLO QCD corrections to $\gamma\gamma \rightarrow J/\psi + c\bar{c}$ process. This is the first truly NLO calculation of 2 to 3 inclusive process for heavy quarkonium production.
- Numerical results shows that the cross section of this process at LEP II can be enhanced with a K factor of about 1.46. And the discrepancy between theoretical prediction and experiment data reduced accordingly.
- As a experimentally distinguishable final state process, It is meaningful to predict its cross section at CEPC. Our results indicate that The NLO corrections is more significant there. And the K factor can reach up to 1.76.

Summary

- We have calculated the NLO QCD corrections to $\gamma\gamma \rightarrow J/\psi + c\bar{c}$ process. This is the first truly NLO calculation of 2 to 3 inclusive process for heavy quarkonium production.
- Numerical results shows that the cross section of this process at LEP II can be enhanced with a K factor of about 1.46. And the discrepancy between theoretical prediction and experiment data reduced accordingly.
- As a experimentally distinguishable final state process, It is meaningful to predict its cross section at CEPC. Our results indicate that The NLO corrections is more significant there. And the K factor can reach up to 1.76.
- According to Furry's theorem, CS process $\gamma\gamma \rightarrow J/\psi + gg$ is forbidden. The cross section of $\gamma\gamma \rightarrow J/\psi + ggg$ is relatively small.