Work @ CERN

Analysis Strategy on Acceptance, Trigger & Reconstruction efficiency Inclusive $b \rightarrow J/\psi X$, $J/\psi \rightarrow \mu \mu$

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Work @ CERN

- CSC trivial tasks: small gears on cable and installation in Carven P5, etc. (Feb. – May.)
- □ CSC shift for Beijing IHEP
 - Common commissioning and test: CRUZET 2 (Jun.-Jul.)
 - CRUZET 3,4(Aug. –Sep.)
 - CRAFT 1-4 (Oct. Nov.)
- □ Inclus b analysis in B physics group

Analysis Strategy

Outline

Acceptance, Trigger & Reconstruction efficiency

- Motivation
- □ Strategy in anal Inclus b & B+ production @ CDF
- □ Strategy for Jpsi production @ CMS
- □ Inclus b analysis @ CMS
- To do list

Acceptance, Trigger & Reconstruction efficiency

- Acceptance: the detector geometric and kinematic acceptances. Determinateed from Monte Carlo simulation
- Efficiency: Triggered vs. unTriggered
 & reconstructed vs. nonreconstructed
- Muon: hits , Segments, Stubs, Macthing

PHYSICAL REVIEW D 71, 032001 (2005) CDF: Acceptance, Trigger & Reconstruction efficiency I

$$\begin{split} \mathcal{A} \left(p_T, y \right) &= \frac{N^{\text{rec}} \left(p_T(J/\psi), |y(J/\psi)| < 0.6 \right)}{N^{\text{gen}} (p_T'(J/\psi), |y'(J/\psi)| < 0.6)} \\ \mathcal{A}' &= \frac{N^{\text{rec}} (|y|_{\text{gen}} > 0.6, |y|_{\text{rec}} < 0.6)}{N^{\text{gen}} (|y|_{\text{gen}} < 0.6)} \end{split}$$

The GEANT simulation is validated by comparing the resulting distributions of various kinematic quantities such as , pT, the track-stub matching distance, and the *z* vertex distribution in reconstructed data and reconstructed Monte Carlo events. Differences in the data and Monte Carlo distributions are used to estimate the systematic uncertainties on the modeling of the CDF detector geometry in the simulation.

For the level 1 di-muon trigger efficiency, used *Jspi* events that were taken with a high-pT single-muon trigger. At level 1, this trigger requires a muon with pT greater than 4.0 GeV/c. In level 3, a *Jpsi* is reconstructed using the triggered high-pT muon and a second muon which is not required to pass the level 1 requirements. This second muon is then used to measure the level 1 single-muon efficiency. The denominator of the efficiency measurement is the number of *Jpsi* reconstructed using the level 3 track and muon information.

$$\epsilon_{L1}^{\mu}(p_T^{\mu}) = E \cdot \text{freq}\left(\frac{A - 1/p_T}{R}\right), \qquad \epsilon_{L1}^{J/\psi}(p_T^{J/\psi}) = \epsilon_{L1}^{\mu}(p_T^{\mu_1}) \cdot \epsilon_{L1}^{\mu}(p_T^{\mu_2})$$

The level 3 reconstruction efficiency is the difference between the on-line and the offline tracking efficiencies. A fast tracking algorithm is used for pattern recognition in the COT in level 3. In the off-line reconstruction a more accurate tracking algorithm is combined with the result of the level 3 algorithm to give a higher overall COT-tracking efficiency.

$$\epsilon^{\mu}_{L3/Offline} = 0.997 \pm 0.001 (stat) \pm 0.002 (syst).$$

PHYSICAL REVIEW D 71, 032001 (2005) CDF: Acceptance, Trigger & Reconstruction efficiency II

TABLE II. Summar	y or J/ψ reconstruction efficiencies.
J/ψ selection	Efficiency
Level 3 muon reconstruction COT off-line tracking	$\epsilon_{L3} = 0.997 \pm 0.001 \pm 0.002$ $\epsilon_{COT} = 0.9961 \pm 0.0002^{+0.0034}_{-0.0034}$
Muon off-line reconstruction Muon z_0 position less than ± 90 cm	$\epsilon_{CMU} = 0.986 \pm 0.003 \pm 0.010$ $\epsilon_{z_0} = 0.9943 \pm 0.0016$
Dimuon z_0 separation less than 5 cm	$\epsilon_{\Delta z_0} = 1.0 \pm 0.001$
Total reconstruction	$\boldsymbol{\epsilon}_{\mathrm{rec}} = \boldsymbol{\epsilon}_{\mathrm{L3}}^2 \cdot \boldsymbol{\epsilon}_{\mathrm{COT}}^2 \cdot \boldsymbol{\epsilon}_{\mathrm{CMU}}^2 \cdot \boldsymbol{\epsilon}_{z_0} \cdot \boldsymbol{\epsilon}_{\Delta z_0} = 95.5\% \pm 2.7\%$

the track-stub matching

 $\epsilon_{\chi^2} = (1.0018 \pm 0.0003) - (0.0024 \pm 0.0001) p_T^{\mu}$

An event-by-event weighting to determine the Jpsi yield

$$1/w_i = \epsilon_{L1}(p_T^{\mu 1}) \cdot \epsilon_{L1}(p_T^{\mu 2}) \cdot \epsilon_{\chi^2}(p_T^{\mu 1}) \cdot \epsilon_{\chi^2}(p_T^{\mu 2}) \cdot \mathcal{A}(p_T^{J/\psi}, y^{J/\psi}).$$

The Jpsi differential cross section calculated as follows

$$\frac{d\sigma}{dp_T} \cdot \operatorname{Br}(J/\psi \to \mu \mu) = \frac{N(p_T)_{\text{corrected}} \cdot (1 - \mathcal{A}')}{\epsilon_{\text{rec}} \cdot \int \mathcal{L} dt \cdot \Delta p_T}$$
$$\sigma_i(H_b) = \frac{\sigma_i(\text{raw})}{f_{\sigma}^i} = \frac{\sum_{j=1}^N w_{ij}\sigma_j(J/\psi)}{f_{\sigma}^j}$$

PHYSICAL REVIEW D 75, 012010 (2007) CDF: Acceptance, Trigger & Reconstruction efficiency III

TABLE I. Detector acceptance, \mathcal{A} , as a function of the $B^{\pm} p_T$. The acceptance \mathcal{A}_{corr} includes corrections evaluated using the data. The average $\langle p_T \rangle$ is the value at which the theoretical differential cross section [1] equals the integrated cross section in each momentum bin divided by the bin width.

p_T range (GeV/c)	$\left< p_T \right> ({\rm GeV}/c)$	A (%)	$\mathcal{A}_{\mathrm{corr}}$ (%)
6-9	7.37	1.545	1.780 ± 0.045
9-12	10.38	3.824	4.405 ± 0.111
12-15	13.39	5.966	6.872 ± 0.173
15-25	19.10	8.819	10.16 ± 0.25
≥ 25		12.516	14.42 ± 0.36

The bin width pT and Acorr, the geometric and kinematic acceptance that includes trigger and tracking efficiencies measured with the data, are listed in Table I.

TABLE II. Summary of efficiencies for reconstructing B^{\pm} candidates in the data and the simulation. The last column indicates the corrections applied to the simulated acceptance and used to derive \mathcal{A}_{corr} in Table I.

Source	Data	Simulation	Con.
COT tracking	(0.996 ± 0.006) ³	(0.998 ± 0.002) ³	1.00 ± 0.02
Kaon interaction			1.000 ± 0.003
CMU acc. and eff.	$(0.6251 \pm 0.0047)^2$	$(0.6439 \pm 0.0004)^2$	0.942 ± 0.014
L1 CMU primitives	$(0.9276 \pm 0.0005)^2$	$(0.8369 \pm 0.0004)^2$	1.228 ± 0.002
L1 eff.	0.9879 ± 0.0009	0.9868	1.0011 ± 0.0009
L2 eff.	0.9948 ± 0.0001	0.9939	1.0009 ± 0.0001
L3 eff.	$(0.997 \pm 0.002)^2$	1	0.994 ± 0.004
Total	0.328 ± 0.008	0.283 ± 0.002	1.152 ± 0.029

In the simulation, the efficiencies of the L1 and L2 triggers are 0.9868 and 0.9939, respectively. By studying *Jpsi* candidates acquired with the CMUP*pT*4 trigger, the L1

efficiency is measured to be 0.9879 \pm 0.0009, and that of the

L2 trigger 0.9948 \pm 0.0001. The L3 trigger is not simulated. The L3 trigger efficiency is dominated by differences between the online and offline reconstruction code efficiency.5 The relative L3 efficiency for reconstructing a single muon identified

by the offline code has been measured to be 0.997 \pm 0.002 [26]. The reconstruction efficiencies in the data and in the simulation are summarized in Table II.

The B+ differential cross section calculated as follows

$$\frac{d\sigma(B^+)}{dp_T} = \frac{N/2}{\Delta p_T \times \mathcal{L} \times \mathcal{A}_{\text{corr}} \times BR}$$

CMS AN-2007/023

Zongchang' Jpsi production @ CMS

Ngen is the *total* true number of *Jpsi*'s present, i.e. *before applying the generator dimuon filter*

$$\epsilon_{offline}^{J/\psi}(p_T^{J/\psi},\eta_{J/\psi},\theta_{J/\psi}) = \epsilon_{\mu_1}(p_T^{\mu_1},\eta_{\mu_1}) \times \epsilon_{\mu_2}(p_T^{\mu_2},\eta_{\mu_2})$$

 μ

$$\epsilon_{trig}(p_T^{J/\psi}, \eta_{J/\psi}) = \frac{N_{trig}(p_T^{J/\psi}, \eta_{J/\psi})}{N_{of\,flinereco}(p_T^{J/\psi}, \eta_{J/\psi})}$$

tit I/ol

$$\frac{d\sigma}{dp_T}(J/\psi) \cdot Br(J/\psi \to \mu^+\mu^-) = \frac{N_{J/\psi}^{fit}}{\int Ldt \cdot A \cdot \lambda_{trigger}^{corr} \cdot \lambda_{reco}^{corr} \cdot \Delta p_T} \quad (p_T^\mu, \eta^\mu) = \frac{\epsilon_{trigger}^{data}(p_T^\mu, \eta^\mu)}{\epsilon_{trigger}^{MC}(p_T^\mu, \eta^\mu)}$$

• $N_{J/\psi}^{fit}$ is the number of reconstructed J/ψ candidates in a given p_T bin. This is obtained by fitting the J/ψ mass spectrum with a linear background and double gaussian signal hypothesis, as was explained in Sec. 4.2.3.

- A is the total efficiency for triggering and offline reconstructing the J/ψ events, as extracted from Monte Carlo. It will be explained in Sec. 5.1.
- λ^{corr}_{trigger} and λ^{corr}_{reco} are correction factors to the trigger and offline efficiencies, respectively, as measured in
 data compared to the Monte Carlo simulation, and will be addressed in Sec.5.2.1.
- ∫ Ldt is the integrated luminosity.
- 2009 Δp_T is the size of the p_T bin.

PHYSICAL REVIEW D 71, 032001 (2005)

J/ \V & b-hadron production

Inv. Mass of J/ψ vs diff. pT bin



$J \neq \psi \& b$ -hadron production

CDF II's Results

PHYSICAL REVIEW D 71, 032001 (2005)

•PPbar energy: 1.96 TeV with 39.7 pb^{-1.}

iμμ) nb/(GeV/c)

do/dp+(J/ψ)*Br(J/ψ-



FIG. 9. Differential cross-section distribution of J/ψ events from the decays of b hadrons as a function of J/ψ transverse momentum integrated over the rapidity range |y| < 0.6. The crosses with error bars are the data with systematic and statistical uncertainties added including correlated uncertainties. The solid line is the central theoretical values using the FONLL calculations outlined in [41]; the dashed line is the theoretical uncertainty.



FIG. 10 (color online). The inclusive J/ψ cross section as $\frac{1}{100}$ function of $J/\psi p_T$ integrated over the rapidity range |y| < 0.6 is plotted as points with error bars where all uncertainties hav the been added. The hatched histogram indicates the contribution to the cross-hatched histogram is the contribution from decays of *l* hadrons. 1

$$\sigma[p\overline{p} \rightarrow J/\psi X, |y(J/\psi)| < 0.6]$$

 $= 4.08 \pm 0.02 (\text{stat})^{+0.36}_{-0.33} (\text{syst}) \ \mu \text{b}.$ $\sigma[p\bar{p} \to H_b, H_b \to J/\psi, p_T(J/\psi) > 1.25 \text{ GeV}/c, |y(J/\psi)| < 0.6] = 0.330 \pm 0.005 (\text{stat})^{+0.036}_{-0.033} (\text{syst}) \ \mu \text{b}.$ $\sigma[p\bar{p} \to J/\psi_p X, \ p_T(J/\psi) > 1.25 \text{ GeV}/c, |y(J/\psi)| < 0.6] = 2.86 \pm 0.01 (\text{stat})^{+0.34}_{-0.45} (\text{syst}) \ \mu \text{b}.$

$$\sigma(p\bar{p} \rightarrow H_b X, |y| < 0.6) = 17.6 \pm 0.4(\text{stat})^{+2.5}_{-2.3}(\text{syst}) \ \mu b$$



FIG. 11. Differential cross-section distribution of *b*-hadron production as a function of *b*-hadron transverse momenta. The crosses with error bars are the data with systematic and statistical uncertainties added, including correlated uncertainties. The solid line is the central theoretical values using the FONLL calculations outlined in [41]; the dashed line is the theoretical uncertainty.



In Figs. 9 and 11, we compare our measurement to a QCD calculation of the *b*-hadron cross section by Cacciari *et al.* [41]. This calculation uses a fixed-order approach with a next-to-leading-log resummation and a new technique to extract the *b*-hadron fragmentation function from LEP data [20,41]. The single *b*-hadron cross section from this FONLL calculation using the CTEQ6M parton distribution functions [43] is $\sigma_{(lyl<0.6)}^{\text{FONLL}} = 16.8^{+7.0}_{-5.0} \,\mu\text{b}$ which is in good agreement with our measurement of $17.6 \pm 0.4(\text{stat})^{+2.5}_{-2.5}(\text{syst}) \,\mu\text{b}.$

We also compare this result to the QCD calculation described in Ref. [19]. This calculation employs a factorization scheme where the mass of the quark is considered negligible and a different treatment of the *b*-hadron fragmentation function is used. The cross-section calculation in [19] is repeated using $\sqrt{s} = 1960 \text{ GeV}/c$ and the MRST2001 parton distribution functions [44]. The central value of the calculated cross section integrated over the rapidity range |y| < 0.6 and $p_T(J/\psi) > 5.0 \text{ GeV}/c$ is $\sigma(p\bar{p} \rightarrow H_b X, |y| < 0.6) \cdot \text{Br}(H_b \rightarrow J/\psi X) \cdot \text{Br}(J/\psi \rightarrow \mu\mu) = 3.2 \text{ nb [45] which is in good agreement with our result of <math>3.06 \pm 0.04(\text{stat}) \pm 0.22(\text{syst})$ nb.

A more complete discussion of the changes in QCD calculations can be found in Refs. [18,20,41]. Updated determinations of proton parton densities and bottom quark fragmentation functions have brought the QCD calculations into better agreement with the CDF measurements of the total *b*-hadron cross section and the *b*-hadron p_T distribution.

a fixed-order approach with a next-to-leading-log resummation

 $\tau_B = 1.526 \pm 0.034 \text{ (stat)} \pm 0.035 \text{ (syst) ps}$

Inclus b @ CMS

- □ First measurement of b-production cross section at sqrt(s) = 10 and 14 TeV
- □ Test of QCD calculations
- Essential measurement of background for many other processes, validation of b-tagging

Inclusive J/ ψ cross section $\sigma[pp \rightarrow J/\psi X]$

- **J**/ ψ from b;
- Prompt J/ ψ: directly produced, prompt decay of heavier charmonium ¹p₀ x_{c0}, ³p₁ x_{c1}, ³p₂ x_{c2}, etc.

to measure: $\sigma[pp \rightarrow H_b, H_b \rightarrow J/\psi X]$

□ Inclusive b production cross section

(another approach based on bJets Andreev, et al. (CMS NOTE 2006/120)

from the fraction of inclusive $b \rightarrow J/\psi X$ to deduce the total inclusive b production

 $\begin{array}{ll} \mathsf{Br}(\mathsf{H}_{\mathsf{b}} \to \mathsf{J}/\psi \mathsf{X} \ \mathtt{)} = 1.16\% \pm 0.10\% & \sigma[pp \to H_b X] \cdot Br(H_b \to J/\psi X) \\ \mathsf{Br}(\mathsf{J}/\psi \to \mu \ \mu \ \mathtt{)} = 5.93\% \pm 0.06\% & = \sigma[pp \to H_b, H_b \to J/\psi X] \end{array}$

- □ The further study on inclusive b lifetime
 - calibrate resolution function
 - understand calibration & alignment



• The inclusive b-> J/ ψ cross-section is calculated by

$$\frac{d\sigma}{dp_T} \cdot \operatorname{Br}(J/\psi \to \mu^+ \mu^-) = f_b(p_T) \cdot \frac{N_{Reco}(p_T) \cdot (1 - A')}{\int L dt \cdot A \cdot \varepsilon_{Trig} \cdot \varepsilon_{Offline} \cdot \Delta p_T}$$

- **1.** $\int L dt$: the integral luminosity
- **2.** f_b : fraction for J/ Ψ from b
- 3. ΔP_T : the size of the P_T bin.
- 4. N_{Reco} : the number of reconstructed J/ ψ signals
- 5. A, A': the acceptance and relative acceptance
- 6. ϵ_{Trig} : trigger efficiency
- 7. $\epsilon_{Offline}$: off-line reconstruction efficiency



 Acceptances include the detector geometric and kinematic acceptances, Can be obtained by Monte Carlo simulation. is treated and defined as:

$$A(p_T^{J/\psi}, \eta^{J/\psi}) = \frac{N^{Rec}(p_T^{J/\psi}, |\eta^{J/\psi}| < 2.4)}{N^{Gen}(p_T^{J/\psi}, |\eta^{J/\psi}| < 2.4)}$$

pT: 9-10GeV



Offline p-jpsi: 11545 b: 4288

1. a finite experimental resolution on each measurement

$$F_{I}(t) = \exp(-t/\tau)$$

$$F_R(t) = \exp(-t/\tau) \otimes G(t, \mu, \sigma)$$

= $\int dt' \exp(-t'/\tau) G(t-t', \mu, \sigma)$

$$F_E(t, dt) = \exp(-t/\tau) \otimes G(t, \mu, dt)$$

 $F_E(t, dt) = \exp(-t/\tau) \otimes G(t, \mu, \mathbf{s} \cdot \mathbf{dt})$

2009-01-16

Unbinned combined MLH fit & analysis method: 3pb-1 "data"







- □ The differential *b*-hadron cross section vs. pT(H_b) is extracted from the measured differential ones of H_b ->J/ΨX
- □ Distortions between pT distribution of b hadrons and J/Ψs from them

Unfolding Method: Bayes'

Unfolding on M.C. data



3pb-1: M.C. vs "Data" inclusive b differential production Cross section



Fast Sim:

ProductionWinter2009

<u>TWiki</u> > C<u>MS Web</u> > <u>GeneratorMain</u> > <u>GeneratorProduction</u>

> GeneratorProduction2009 > ProductionWinter2009

The 1G Winter09 pp@10TeV Fast Sim Production for Physics with CMSSW 22x

•Min bias: 100 Mevt with Pythia

•QCD jets: 600 Mevt with Madgraph+Pythia

•tt + Jets: 10 Mevt with Madgraph)

•t+ jets: 3 Mevt with Madgraph

•Photon+jets: 25 Mevt with Madgraph

•Z/W + jets: 50 Mevt with Madgraph

•Enriched electrons/muons (mostly from b-decays) 25Mevt with Pythia+Madgraph

•Enriched photons ? 10 Mevt with Pythia+Madgraph

•bbbar 50 Mevt with Pythia+Madgraph

•Onia 5 Mevt with Pythia

•120 Mevt of additional requests

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To do List

- Solve the discrepancy (data vs. M.C.) of CSA07 data anal(1 or 2 weeks)
- Start CMSSW_2_1_12 anal on Summer08 data
 - Efficiency: Accept., Trig & reco. M.C. & T.P.
 - Comprehensive anal method
 - □ Ctau efficiency
 - Scale factor
 - Unfolding method
- □ Fast Sim Production: QCD @Beijing T3?
- □ Prepare PAS and Note draft.

backups

Pseudo proper decay length



- pseudo proper decay length distribution
- □ Measure the I distribution of the background under the J/ Ψ by studying the $\mu + \mu$ - mass sidebands of the J/ Ψ
- Fit the distribution to the sum of background, direct (zero-lifetime) and B decay (non-zero lifetime) Contributions and extract the lifetime

Unfolding methods I

- Bin-to-bin correction: no into account migrations a bin to the others; neglect correlation between adjacent bins.
- The matrix method: solve the problem of migrations; singular problem; statistical fluctuations; results unstable.
- Regularized unfolding: satisfactory results but technical complications; only with one dimension

Unfolding Method II: Bayes'

A Multidimensional unfolding method based on Bayes' theorem by G.D'Agostini, Nucl. Instr. Meth. A362 (1995) 487-498. -- Model independent method

$$P(\mathbf{C}_i | \mathbf{E}_j) = \frac{P(\mathbf{E}_j | \mathbf{C}_i) P_0(\mathbf{C}_i)}{\sum_{l=1}^{n_{\mathbf{C}}} P(\mathbf{E}_j | \mathbf{C}_l) P_0(\mathbf{C}_l)}.$$
$$\hat{n}(\mathbf{C}_i)|_{\text{obs}} = \sum_{j=1}^{n_{\mathbf{E}}} n(\mathbf{E}_j) P(\mathbf{C}_i | \mathbf{E}_j)$$

$$\hat{n}(\mathbf{C}_i) = \frac{1}{\epsilon_i} \sum_{j=1}^{n_{\mathbf{E}}} n(\mathbf{E}_j) P(\mathbf{C}_i | \mathbf{E}_j) \quad \epsilon_i \neq 0$$

Ci: cause in i-th bin_: Ej: effect in j-th bin P(Ci/Ej): corelation matrix for Ej to Ci

$$\hat{N}_{\text{true}} = \sum_{i=1}^{n_{\text{C}}} \hat{n}(C_i),$$
$$\hat{P}(C_i) \equiv P(C_i | n(E)) = \frac{\hat{n}(C_i)}{\hat{N}_{\text{true}}},$$
$$\hat{\epsilon} = \frac{N_{\text{obs}}}{\hat{N}}.$$

the unfolding can be performed through the following steps:

1) choose the initial distribution of $P_0(C)$ from the best knowledge of the process under study, and hence the initial expected number of events $n_0(C_i) = P_0(C_i)N_{obs}$; in case of complete ignorance, $P_0(C)$ will be just a uniform distribution: $P_0(C_i) = 1/n_C$;

2) calculate $\hat{n}(C)$ and $\hat{P}(C)$;

3) make a χ^2 comparison between $\hat{n}(C)$ and $n_0(C)$;

4) replace $P_0(C)$ by $\hat{P}(C)$, and $n_0(C)$ by $\hat{n}(C)$, and start again; if, after the second iteration the value of χ^2 is "small enough", stop the iteration; otherwise go to step 2. Some criteria about the optimum number of iterations will be discussed later.

Unbinned combined MLH fit & analysis method: 3pb-1 "data" PT: 5-40 GeV/c

Jpsi: prompt, b hadrons, and inclusive production
 abstraction for b fraction

2) abstraction for b fraction





