

Searching for dark matter at future colliders

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BeiHang, 2016.09.02

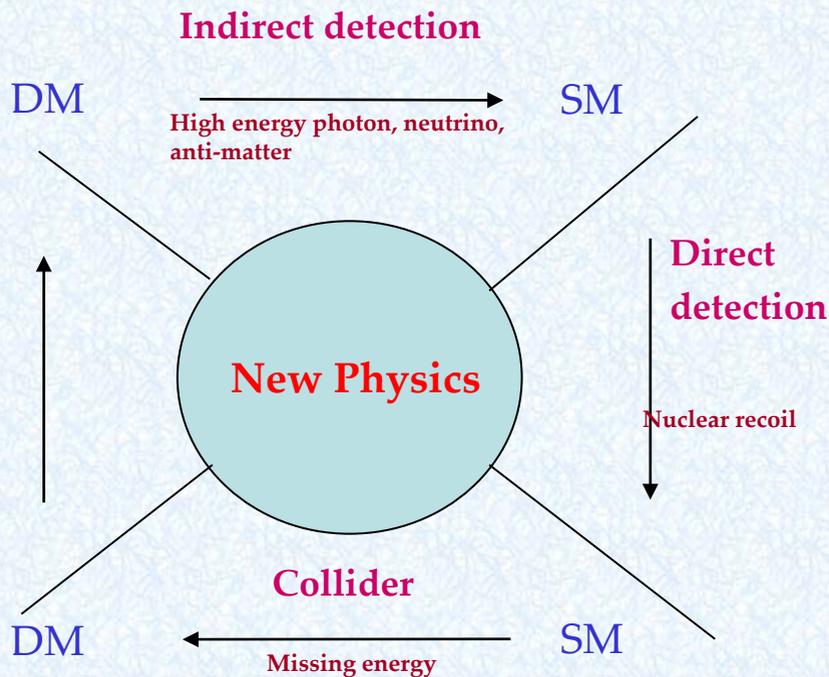


Outline

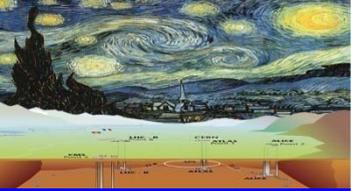
- ⊕ **Introduction: DM searches and DM collider searches**
- ⊕ **Direct searches at e^+e^- colliders**
- ⊕ **Indirect searches at e^+e^- colliders**
- ⊕ **Mass measurements of DM**
- ⊕ **Searches at future pp colliders**



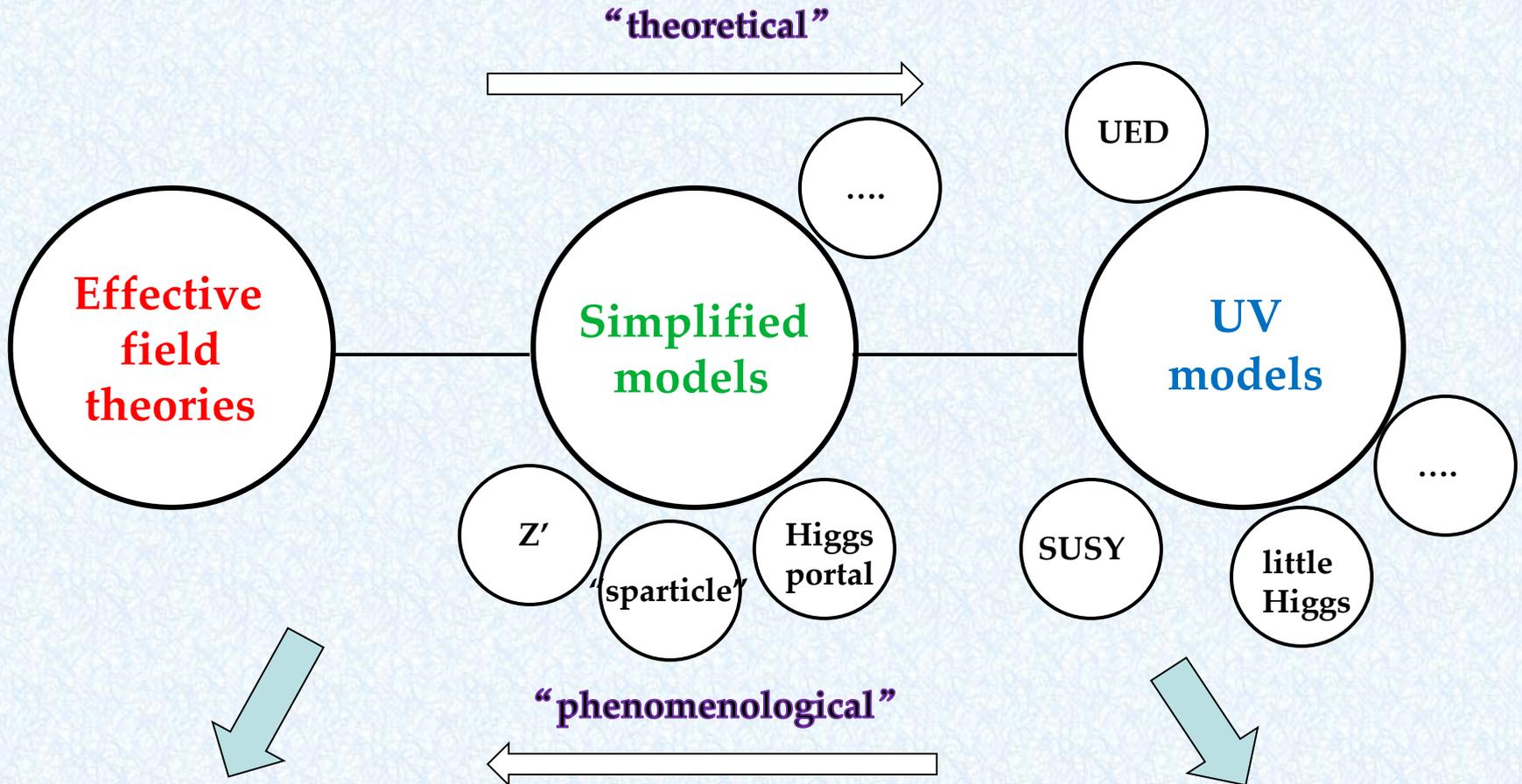
DM searches



- ⊕ We still do not know what DM is
- ⊕ Neutral, stable, and weakly interacting
- ⊕ $\Omega_{\text{DM}} h^2 \sim 0.12$
thermal freeze out? $\sigma v \sim 3 \times 10^{-26} \text{cm}^3/\text{s}$
 $\sigma \sim g^4/m^2$, $m < 1.8 \text{TeV}$ ($g^2/0.3$)
- ⊕ **Weakly Interacting Massive Particle !**
- ⊕ Missing energy at colliders
(transverse missing energy at hadron colliders)
- ⊕ Complementarity of three detections



Theoretical approaches



Mono-X

Direct production, "model-independent"

Multi final states +ME

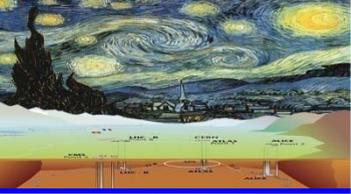
Cascade decay, "model-dependent"

Other collider signatures and constraints



Possible studies at e^+e^- colliders

- ‡ **Disadvantage: Low CM energy;**
Difficult to directly discover heavy new particles
- ‡ **Advantage: No large QCD background; precise beam energy; polarized beams;**
Can accurately measure the mass, spin, and other quantum numbers
Complementary to hadron colliders
- ‡ **Search for light DM $< \sim 100 \text{ GeV}$**
Search for interactions between the DM and electrons/EW gauge bosons/Higgs
- ‡ **Probe relevant particles in new physics models (t-channel annihilation mediator, DM in the EW multiplet...)**
Indirect search through loop effects, e.g. in Higgs measurements, EW precision measurements



Possible studies at e^+e^- colliders

- ‡ Precise measurement rather than discovery

Full missing energy can be obtained

Accurately measure the mass, spin, and other quantum numbers of DM

- ‡ If a missing particle χ is discovered at colliders, can it make up all the DM component in the Universe ?

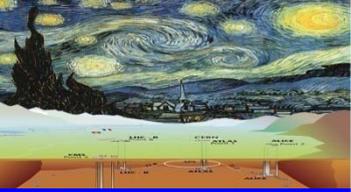
Check the thermal relic density of χ by using measured DM mass and coupling coefficients.

Probe the DM production mechanism and cosmological effects of DM

$\Omega_\chi h^2 = 0.12$, strong support for the discovery of DM

$\Omega_\chi h^2 < 0.12$, multi-DM particles ? Non-thermal production?

$\Omega_\chi h^2 > 0.12$, some other unexpected annihilation channels? Resonant annihilation?
co-annihilation?

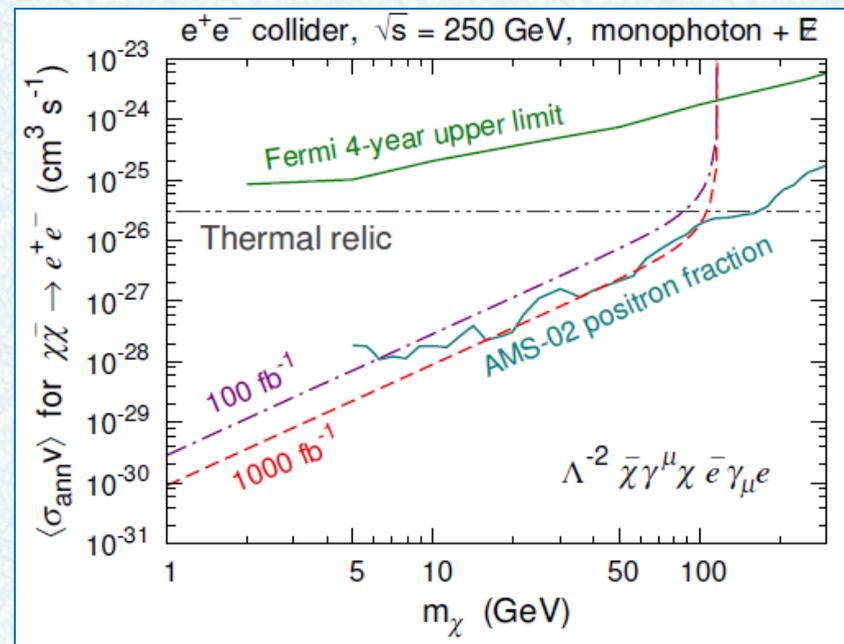
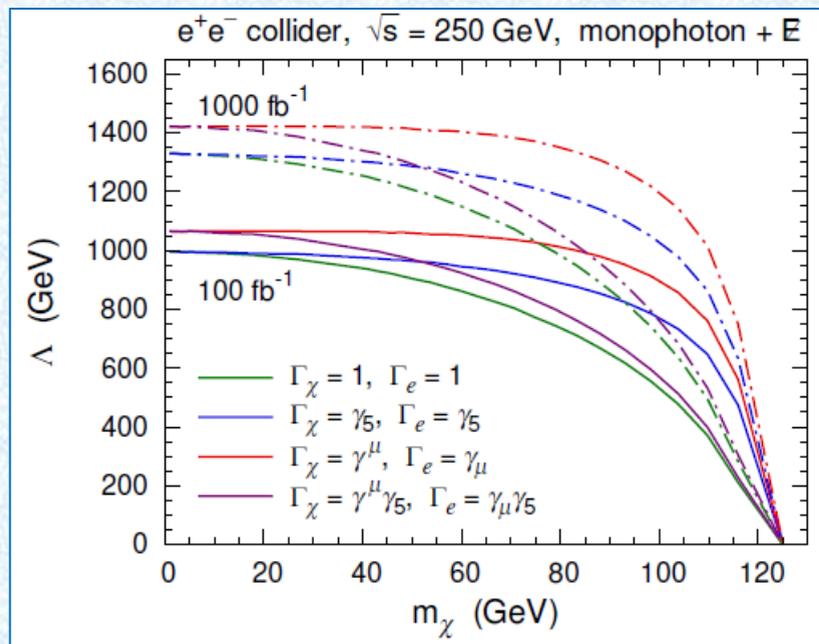


Results for DM interactions with electron/positron

⊕ Effective Field Theory: heavy mediator can be integrated out leaving a scale Λ

$$\mathcal{O}_e = \frac{1}{\Lambda^2} \bar{\chi} \Gamma_\chi \chi \bar{e} \Gamma_e e \quad \Gamma_\chi, \Gamma_e \in \{1, \gamma_5, \gamma^\mu, \gamma^\mu \gamma_5, \sigma^{\mu\nu}\}$$

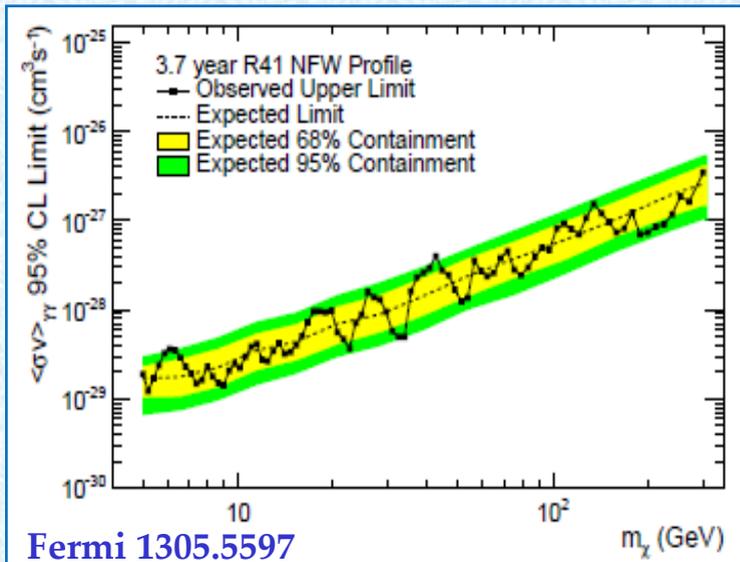
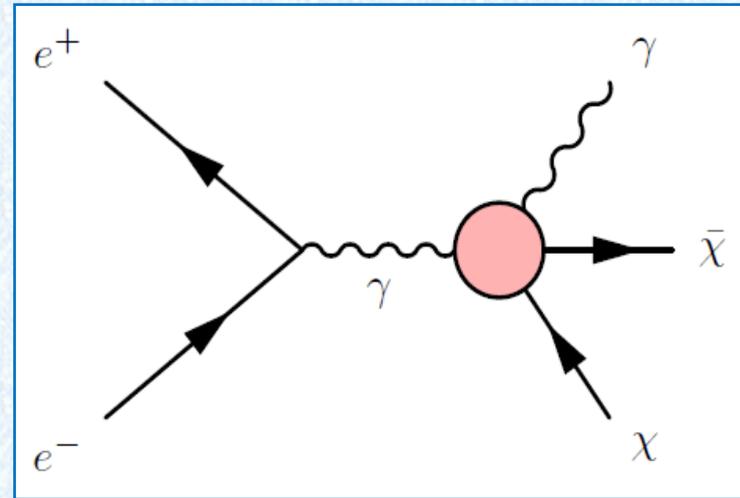
Search for mono-photon signals from initial state radiations





Gamma-ray line and mono-photon

- ⊕ Gamma-ray line is a critical evidence of DM annihilation/decay
- ⊕ ~3-4 σ signals reported by some groups but not confirmed by the Fermi-LAT collaboration.
- ⊕ Search for corresponding mono-photon signals at future e^+e^- colliders



- ⊕ Effective operator

Yu, Yan, PFY, 1307.5740

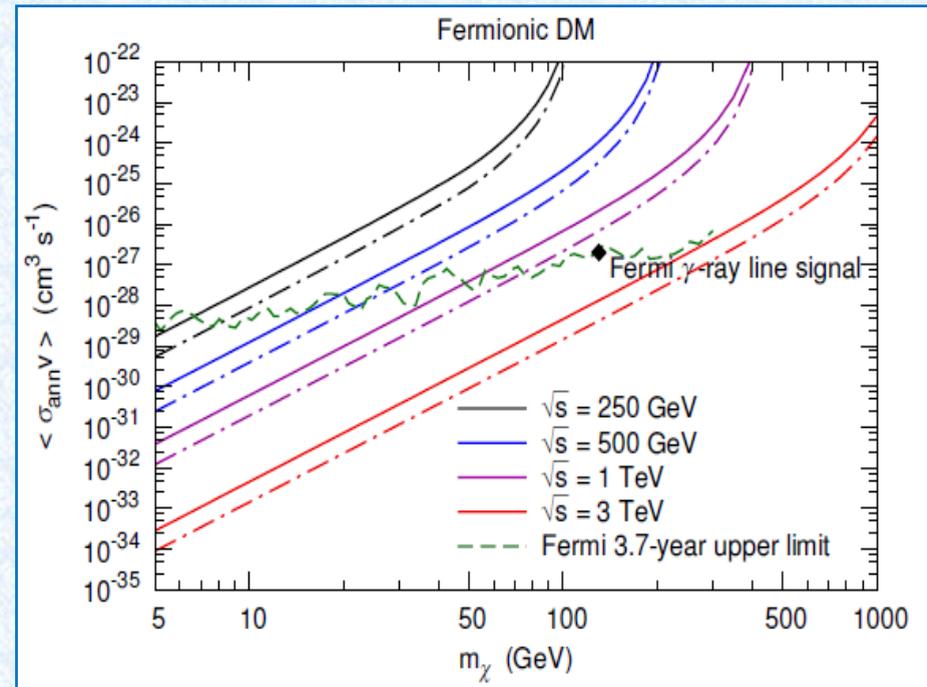
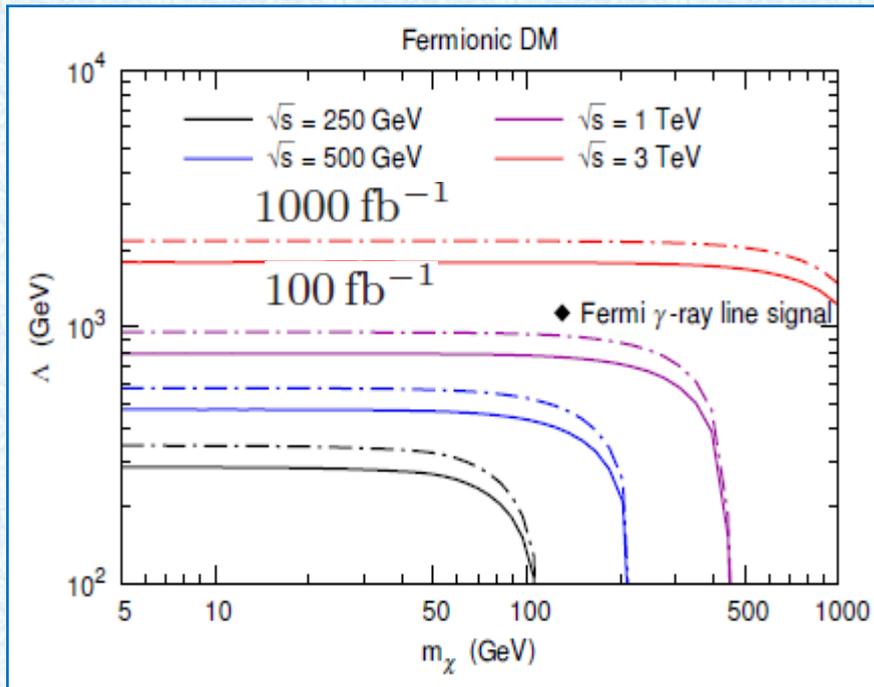
$$\mathcal{O}_F = \frac{1}{\Lambda^3} \bar{\chi} i\gamma_5 \chi F_{\mu\nu} \tilde{F}^{\mu\nu}$$

Cut scale for a ~100 GeV DM and a detectable cross section is ~TeV

$$\langle\sigma_{\text{ann}}v\rangle_{\chi\bar{\chi}\rightarrow 2\gamma} \simeq \frac{4m_\chi^4}{\pi\Lambda^6} = 10^{-27} \text{ cm}^3 \text{ s}^{-1} \left(\frac{m_\chi}{130 \text{ GeV}}\right)^4 \left(\frac{1272 \text{ GeV}}{\Lambda}\right)^6$$

Sensitivities

- ⊕ Consider possible e^+e^- colliders with CM energies of 250, 500, 1000 and 3000 GeV
- ⊕ 3σ sensitivities for mass scale and annihilation cross section
- ⊕ Require large luminosities

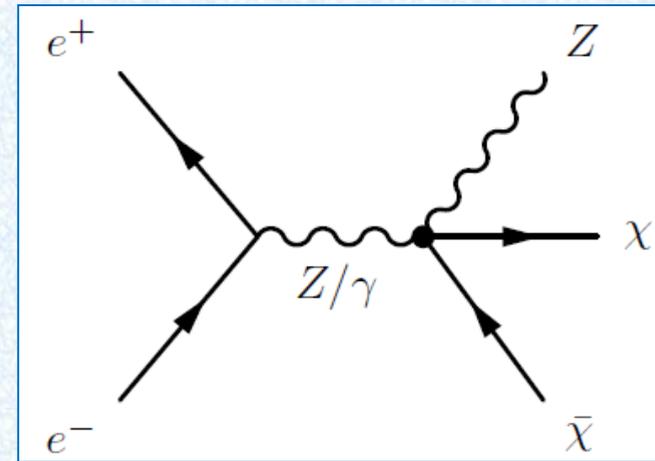




Mono-Z signals

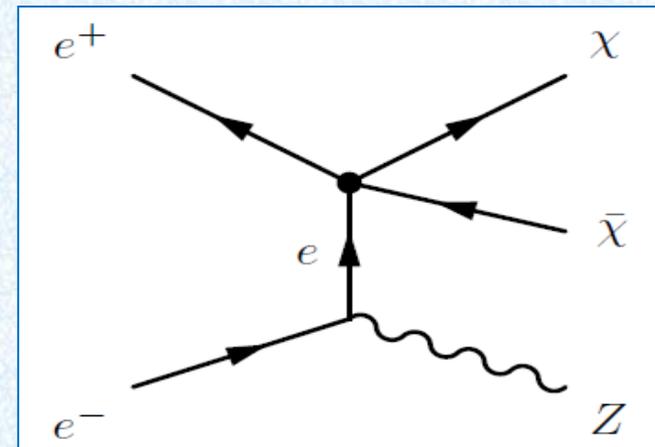
- ⊕ DM can interact with both the photon and Z boson
- ⊕ Consider effective operators

$$\begin{aligned} \mathcal{O}_{F1} &= \frac{1}{\Lambda_1^3} \bar{\chi} \chi B_{\mu\nu} B^{\mu\nu} + \frac{1}{\Lambda_2^3} \bar{\chi} \chi W_{\mu\nu}^a W^{a\mu\nu} \\ &\supset \bar{\chi} \chi (G_{ZZ} Z_{\mu\nu} Z^{\mu\nu} + G_{AZ} A_{\mu\nu} Z^{\mu\nu}) \\ \mathcal{O}_{F2} &= \frac{1}{\Lambda_1^3} \bar{\chi} i\gamma_5 \chi B_{\mu\nu} \tilde{B}^{\mu\nu} + \frac{1}{\Lambda_2^3} \bar{\chi} i\gamma_5 \chi W_{\mu\nu}^a \tilde{W}^{a\mu\nu} \\ &\supset \bar{\chi} i\gamma_5 \chi (G_{ZZ} Z_{\mu\nu} \tilde{Z}^{\mu\nu} + G_{AZ} A_{\mu\nu} \tilde{Z}^{\mu\nu}) \\ \mathcal{O}_{FH} &= \frac{1}{\Lambda^3} \bar{\chi} \chi (D_\mu H)^\dagger D_\mu H \rightarrow \frac{m_Z^2}{2\Lambda^3} \bar{\chi} \chi Z_\mu Z^\mu \end{aligned}$$



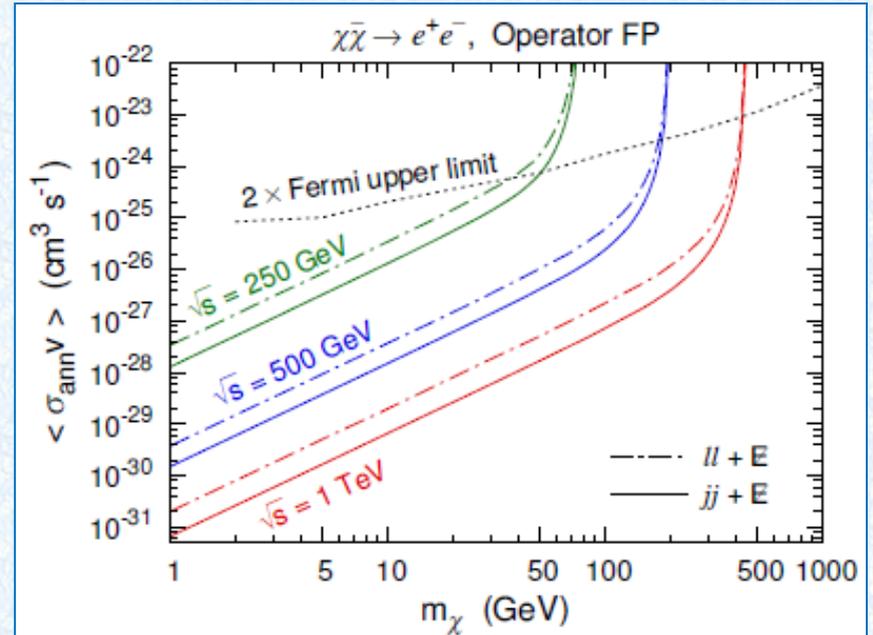
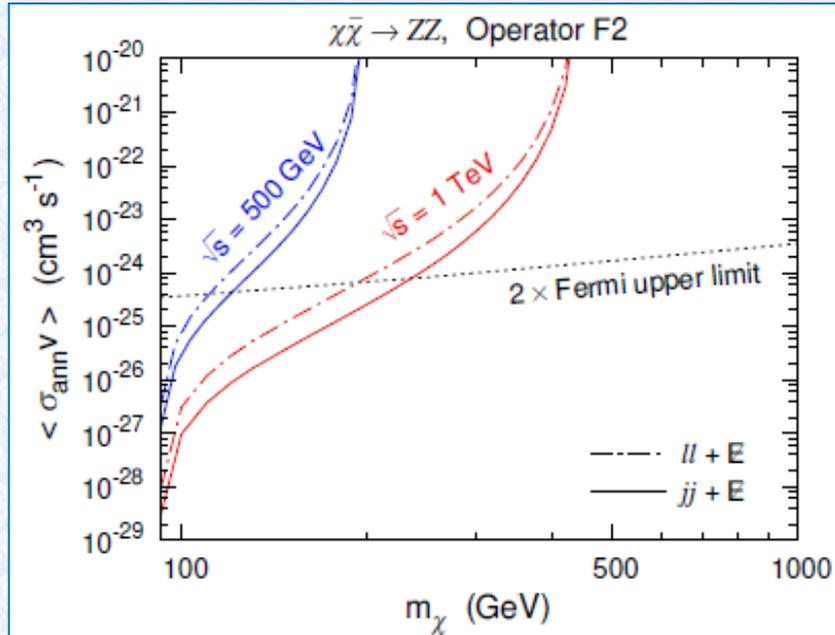
- ⊕ Z boson can also come from initial state radiation

$$\begin{aligned} \mathcal{O}_{FP} &= \frac{1}{\Lambda^2} \bar{\chi} \gamma_5 \chi \bar{e} \gamma_5 e, \\ \mathcal{O}_{FA} &= \frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \gamma_5 \chi \bar{e} \gamma_\mu \gamma_5 e \end{aligned}$$



Sensitivities

- ⊕ 3 σ sensitivities for interactions between the DM and gauge bosons/electrons
- ⊕ Assume: 1000 fb⁻¹ of data; $\Lambda = \Lambda_1 = \Lambda_2$
- ⊕ Compare with the limits from Fermi-LAT dwarf galaxy observations





Light DM coupled to SM mediators

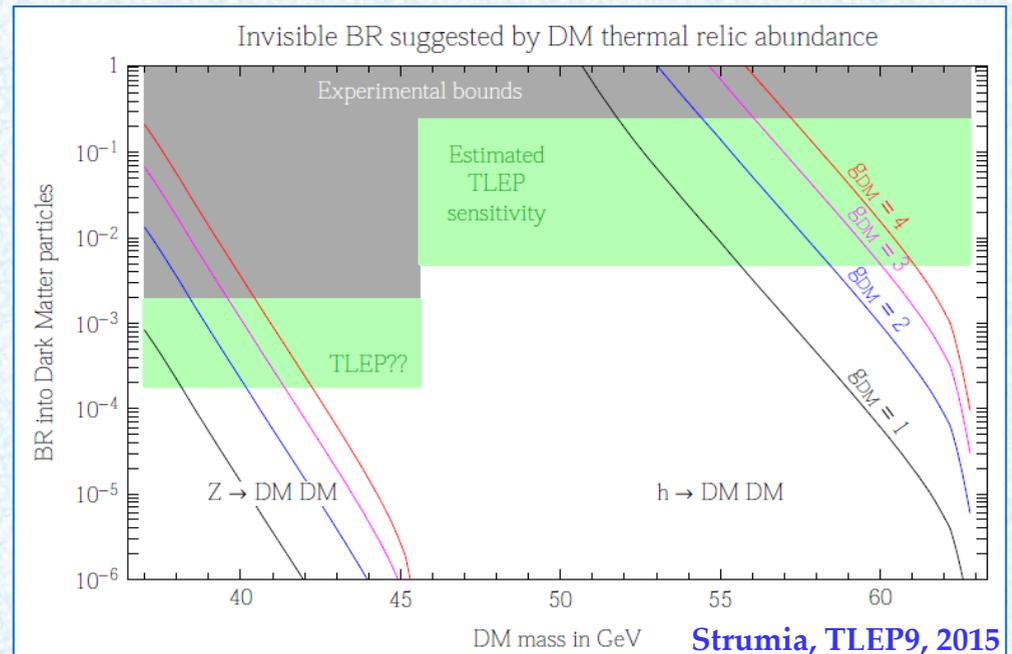
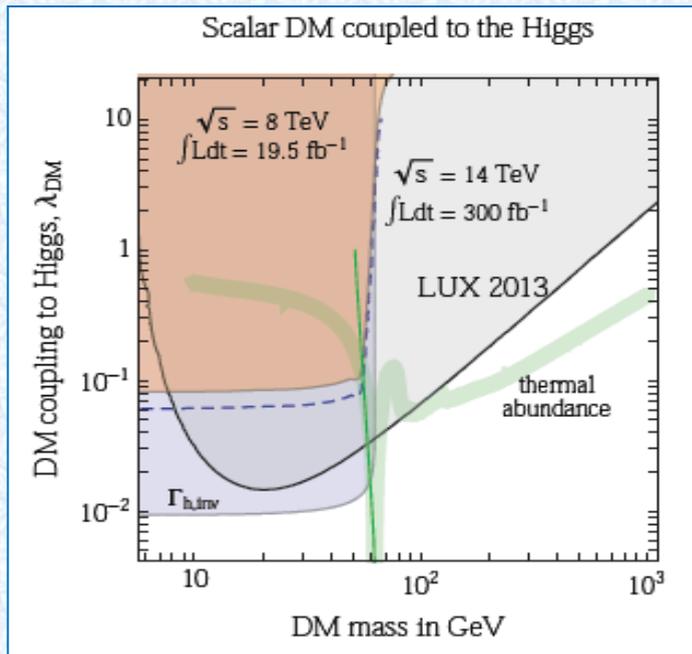
- Assume that DM couples to the SM particles through SM mediators

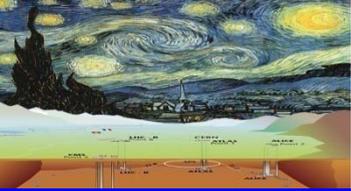
$$\mathcal{L} = -hJ_h, \quad J_h = \frac{1}{\sqrt{2}} \left[\sum_f y_f \bar{f}f + \bar{\psi}_{\text{DM}}(y_{\text{DM}} + iy_{\text{DM}}^P \gamma_5)\psi_{\text{DM}} + \frac{\lambda_{\text{DM}} v}{2} s_{\text{DM}}^2 \right]$$

$$\mathcal{L} = -Z_\mu J_Z^\mu, \quad J_Z^\mu = \frac{g_2}{\cos \theta_W} \left[\sum_f [\bar{f} \gamma_\mu (g_V^f + \gamma_5 g_A^f) f] + \sum_s g_s [s^* (i\partial_\mu s) - (i\partial_\mu s^*) s] \right]$$

- Searches for invisible Higgs/Z decays are useful to probe light DM

See also H.Zhang's talk on the Higgs exotic decay at CEPC



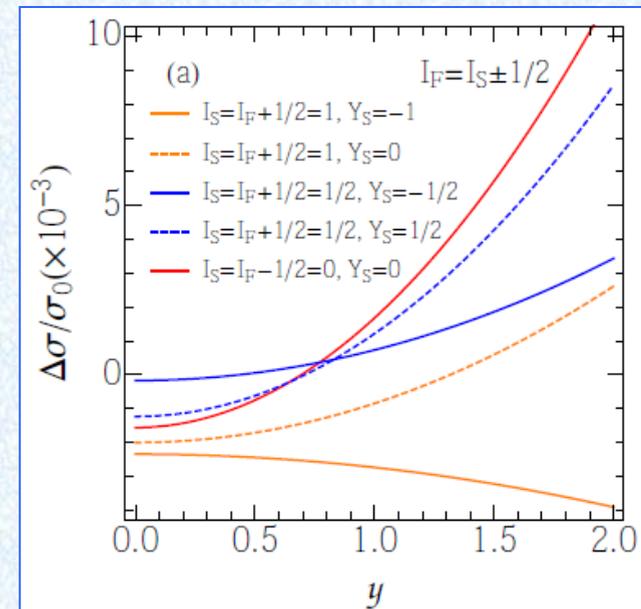
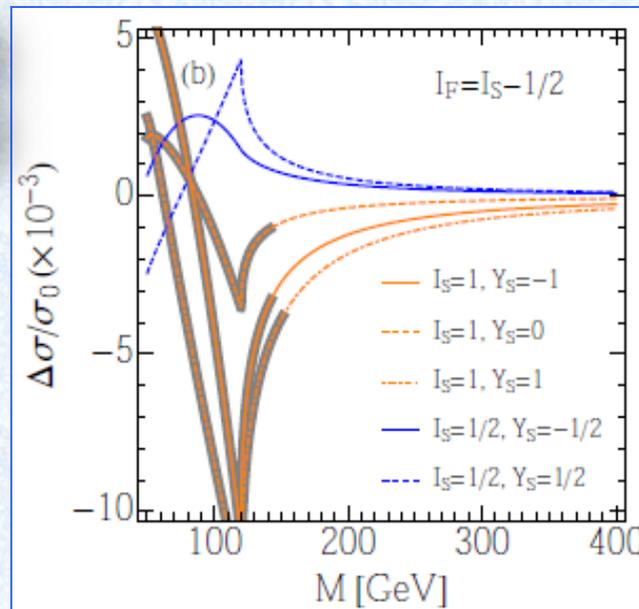
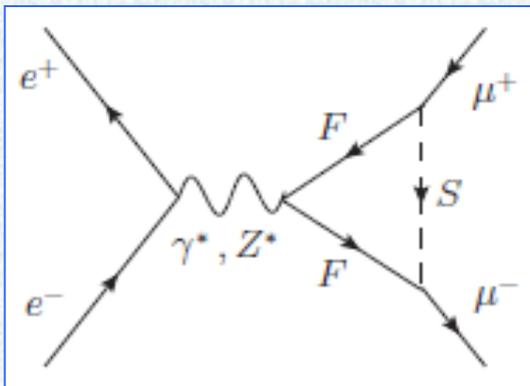


Indirect searches at colliders

- ⊕ There would be some exotic particles carrying EW charges in DM models. These particles might affect SM observations via high order effects
- ⊕ May be required by the DM relic density, e.g. t-channel mediator for annihilation, coannihilation
- ⊕ Indirect searches for new physics. Depend on detailed models
- ⊕ For instance, calculate the loop correction to the $e^+e^- \rightarrow \mu^+\mu^-$ from a leptophilic DM models with a vector-like fermion F (mediator) and a scalar S (DM)

$$\mathcal{L}_{\text{Yuk}} = yC_{ijk}S^i\bar{\mu}_L^kF^j + h.c.$$

$$\mathcal{L}_{\text{Yuk}} = yC_{ij}S^i\bar{\mu}_RF^j + h.c.$$





Test from Higgs measurements

- DM embedded in EW multiplets may interact with Higgs
For instance, a singlet-doublet Fermionic DM model

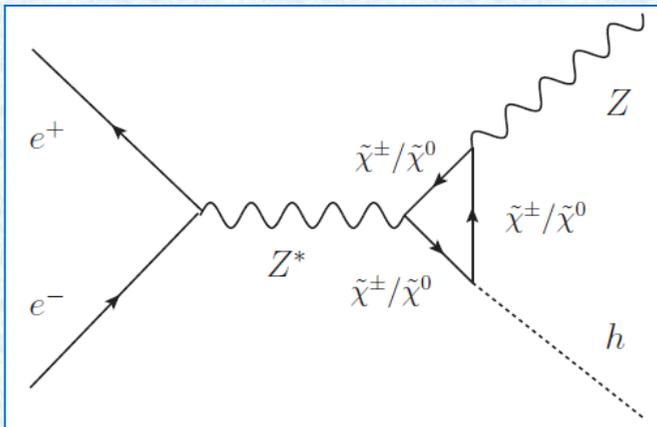
e.g. D'Eramo, 0705.4493;
Cohen et al, 1109.2604

$$\mathcal{L}_S = iS^\dagger \bar{\sigma}^\mu \partial_\mu S - \frac{1}{2}(m_S SS + \text{h.c.}), \quad \mathcal{L}_D = iD_{1i}^\dagger \bar{\sigma}^\mu D_{\mu} D_1^i + iD_{2i}^\dagger \bar{\sigma}^\mu D_{\mu} D_2^i - (m_D \varepsilon_{ij} D_1^i D_2^j + \text{h.c.})$$

$$\mathcal{L}_{\text{HSD}} = y_1 H_i S D_1^i - y_2 H_i^\dagger S D_2^i + \text{h.c.}$$

DM could also be embedded in triplet, quadruplet

- Could affect the Higgs production process at e^+e^- colliders



- DM and other exotic neutral/charged particles may be directly produced at hadron colliders

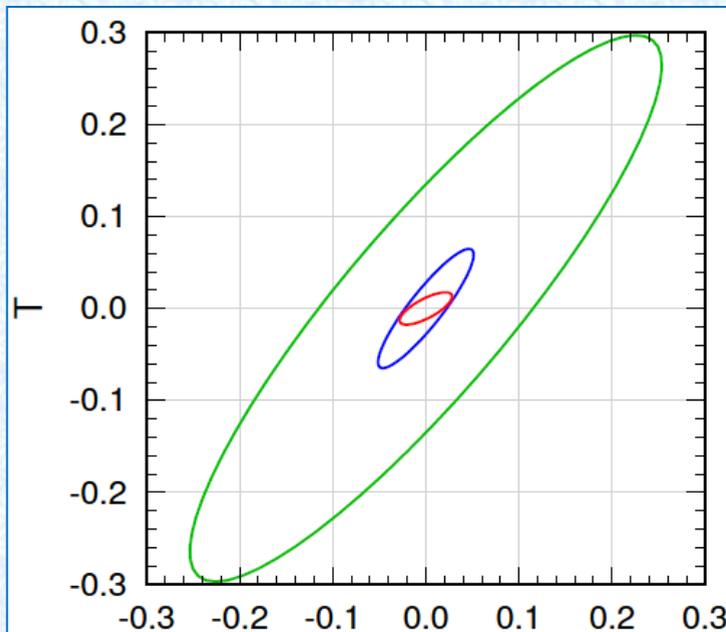


Test from EW precision measurements

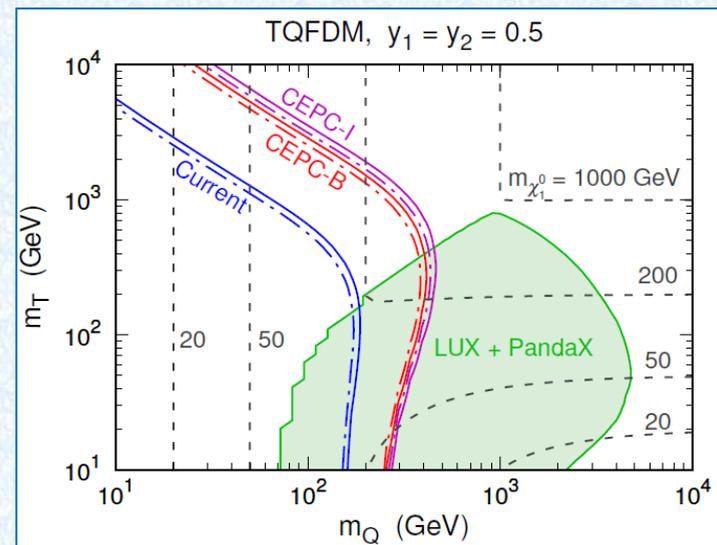
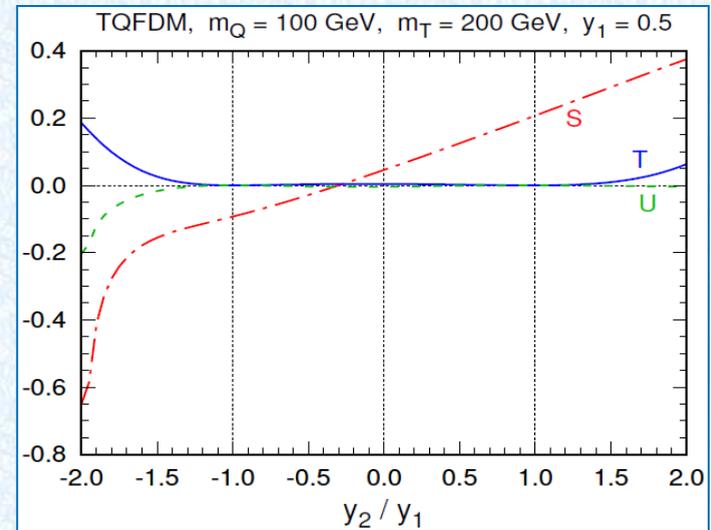
- ⊕ CEPC measurements would improve the precisions of EW parameters

See Fan, Reece, Wang, 1411.1054

- ⊕ Can set constraints on DM (new physics) models



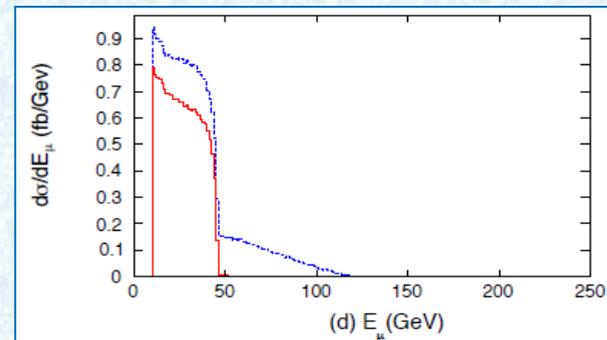
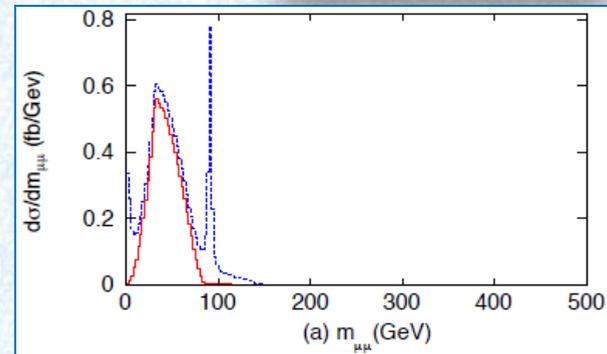
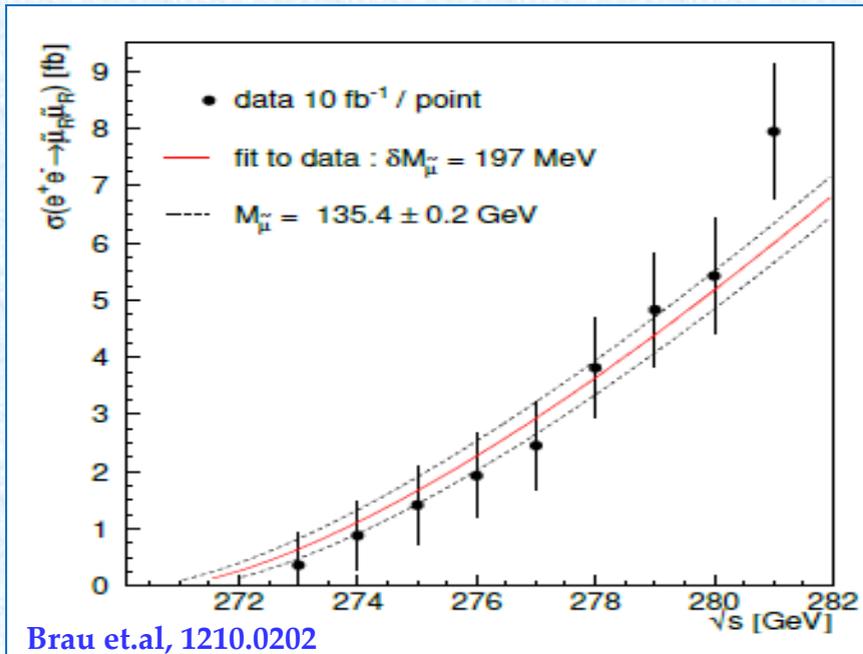
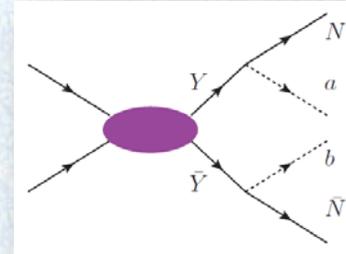
Cai, Yu, et al, in progress



Determination of DM mass at e^+e^- colliders



- ⊕ Depend on the topology of DM production process
- ⊕ Consider a process $e^+ + e^- \rightarrow Y + \bar{Y}$ with $Y \rightarrow a(p_a) + N(k_1)$, $\bar{Y} \rightarrow b(p_b) + \bar{N}(k_2)$
An example is the slepton pair production in SUSY models
- ⊕ Determine the mass of Y through a threshold scan
- ⊕ Use kinematic variables to determine (m_N, m_Y)
e.g. visible particle energy endpoint, recoil mass, angle of visible particles....





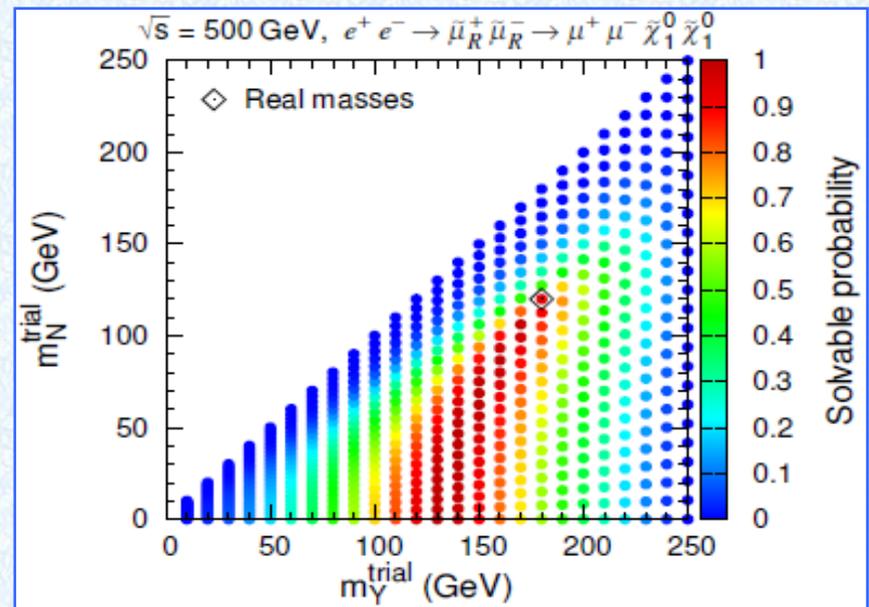
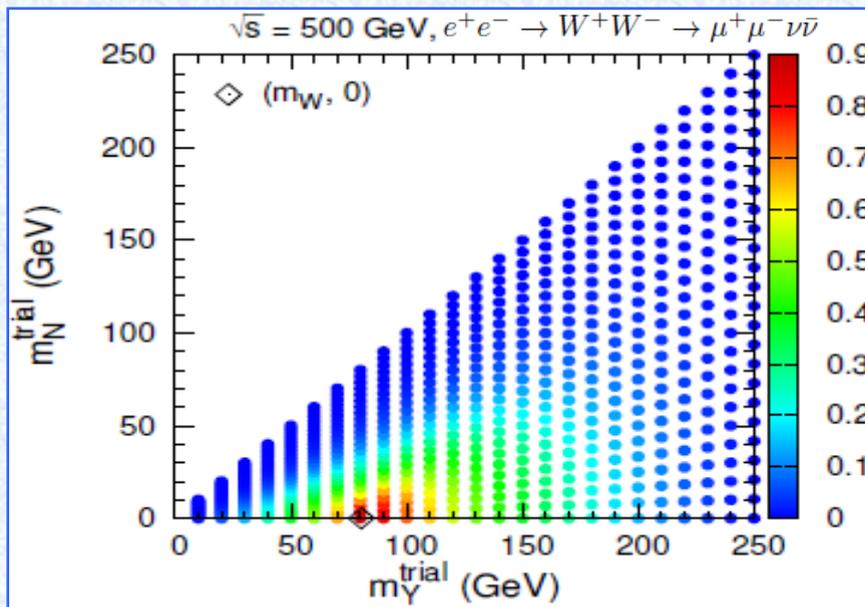
A method of DM mass determination

- Consider a process $e^+ + e^- \rightarrow Y + \bar{Y}$ with $Y \rightarrow a(p_a) + N(k_1)$, $\bar{Y} \rightarrow b(p_b) + \bar{N}(k_2)$
 For the known initial four-momentum and given particle masses, we have 8 Eqs and 8 variables

$$q^\mu = p_a^\mu + p_b^\mu + k_1^\mu + k_2^\mu, \quad k_1^2 = k_2^2 = m_N^2, \quad (p_a + k_1)^2 = (p_b + k_2)^2 = m_Y^2$$

This is a solvable problem at the e^+e^- collider

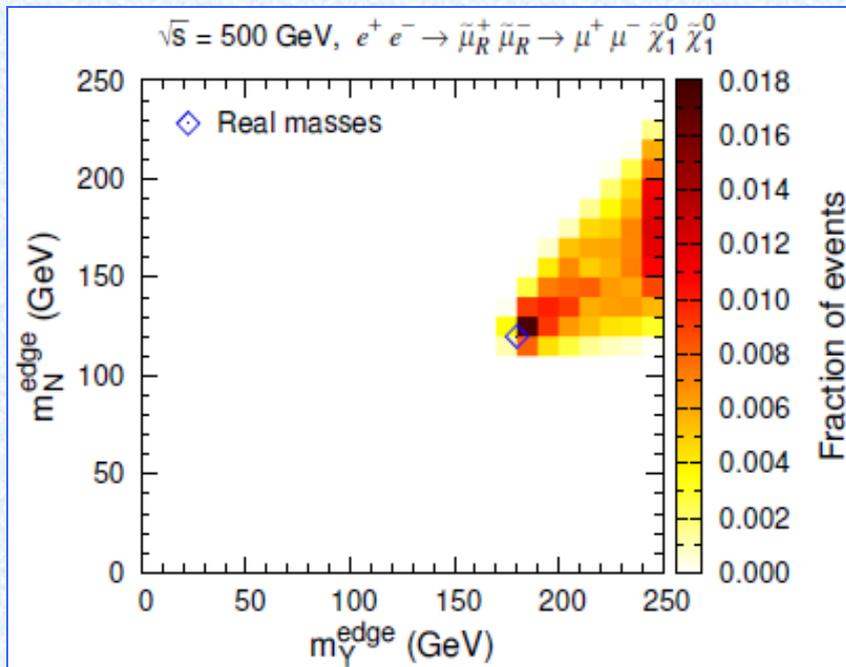
- For a wrong set of (m_N, m_Y) , above Eqs might not have solutions
 Scan the trial m_N - m_Y plane and find realistic values by solving the Eqs



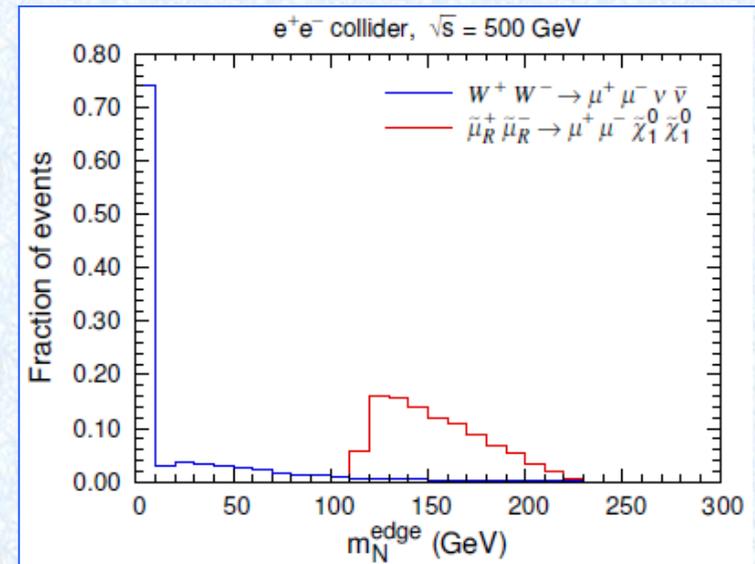
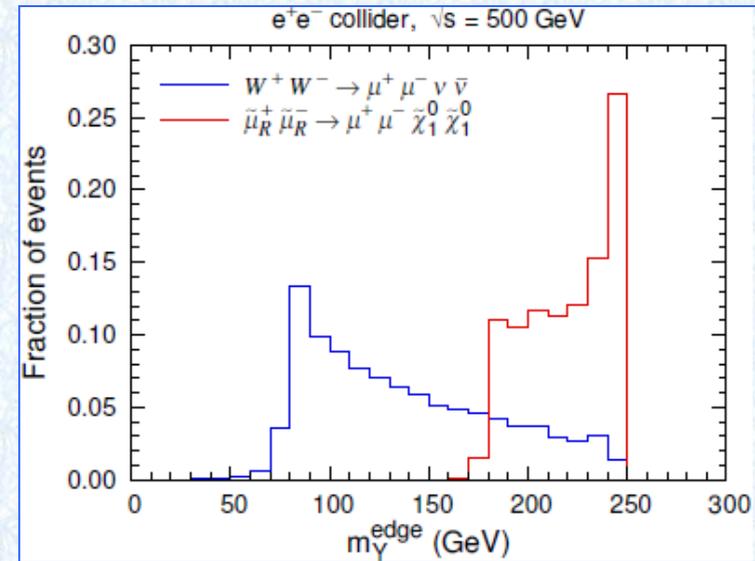


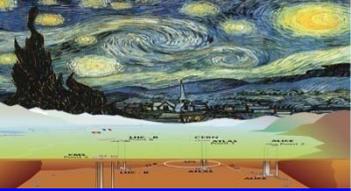
A method of DM mass determination

- ⊕ For every event, obtain a savable point with the largest distance from the origin in the trial m_Y - m_N plane. The coordinates are defined as $(m_Y^{\text{edge}}, m_N^{\text{edge}})$
- ⊕ Extract (m_Y, m_N) from distributions of $(m_Y^{\text{edge}}, m_N^{\text{edge}})$



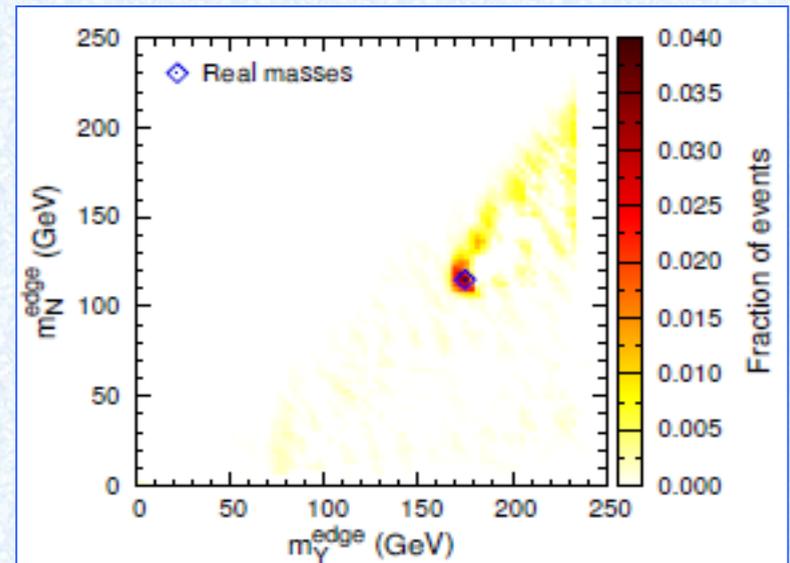
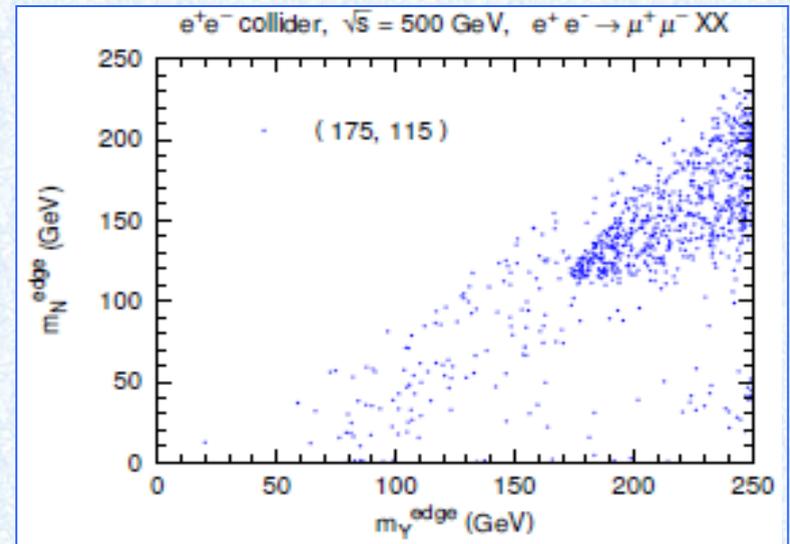
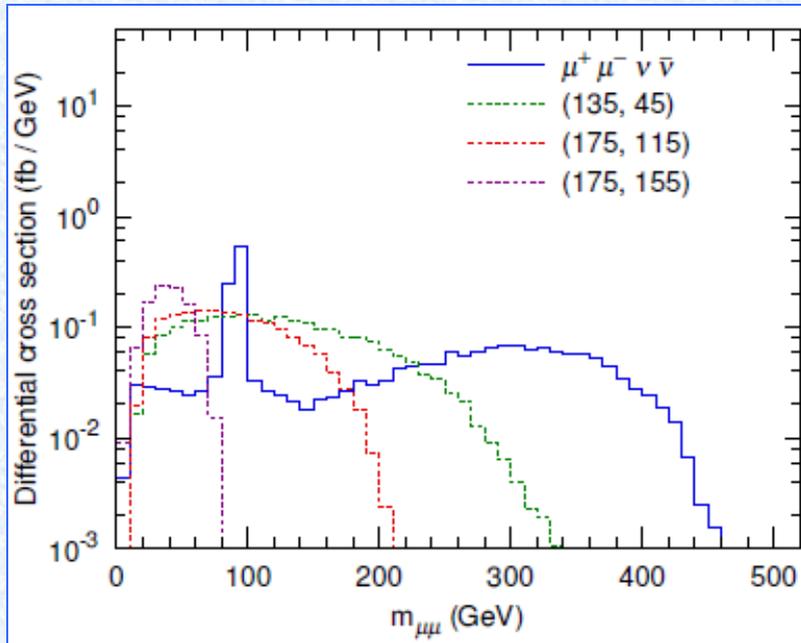
See also Harlamd-Lang et al, 1202.0047

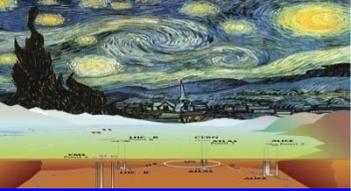




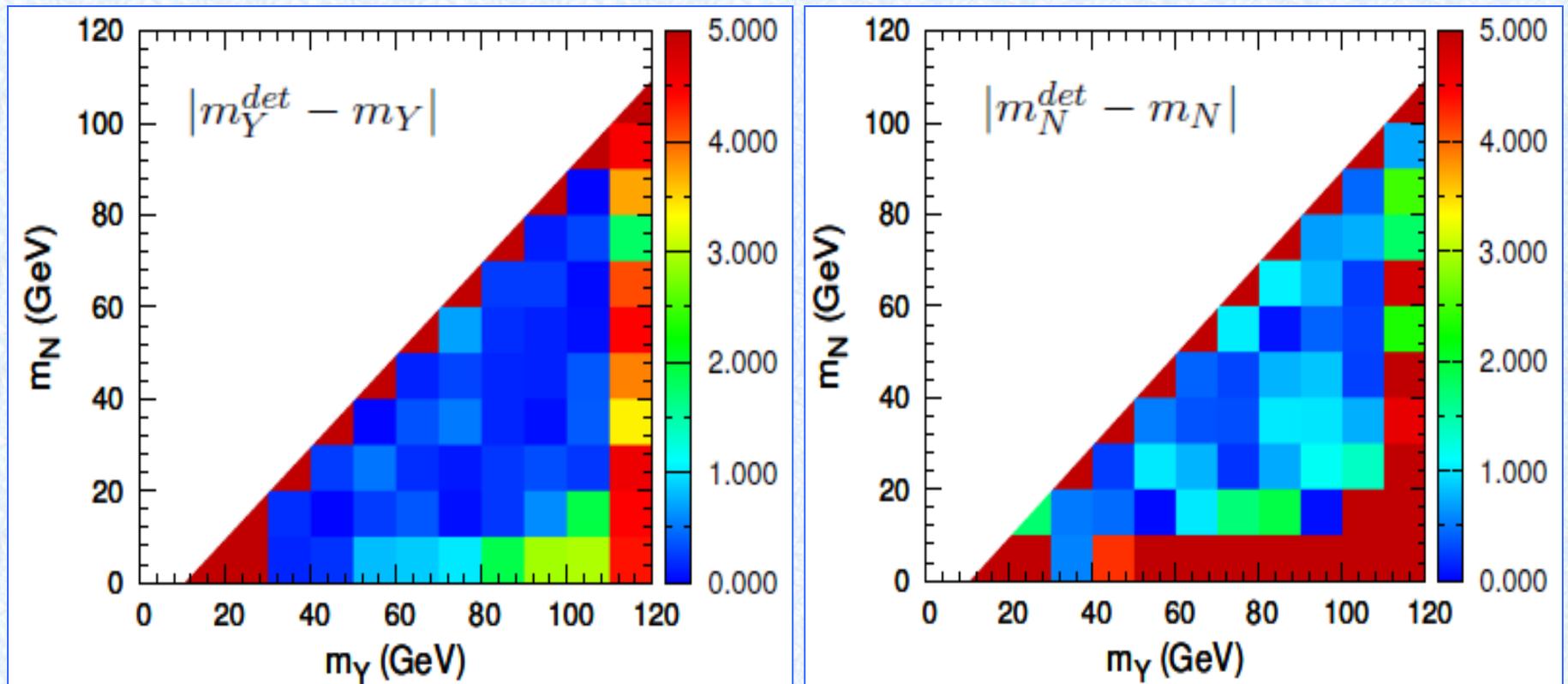
A method of DM mass determination

- ⊕ Include detector effects in the simulation
- ⊕ Consider the backgrounds from SM $\mu^+\mu^-$ +MET
- ⊕ Find realistic values of (m_N, m_Y)

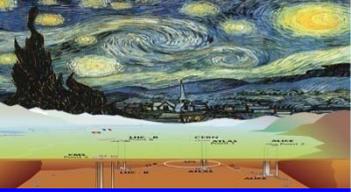




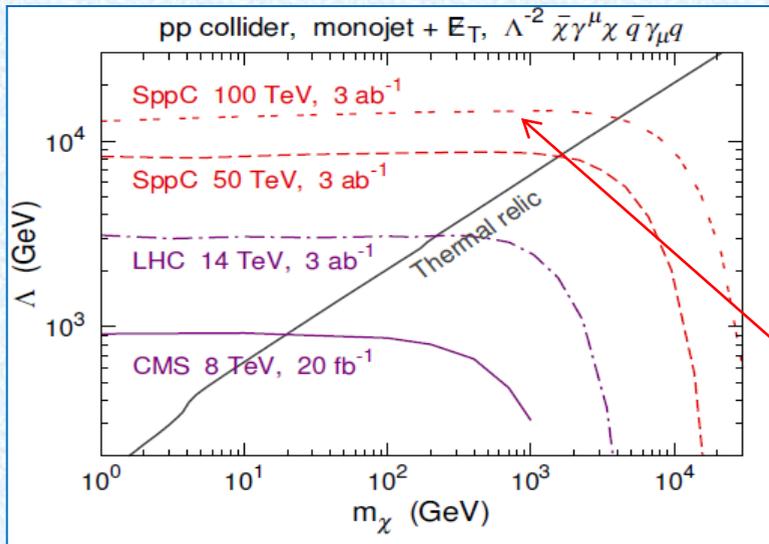
A method of DM mass determination



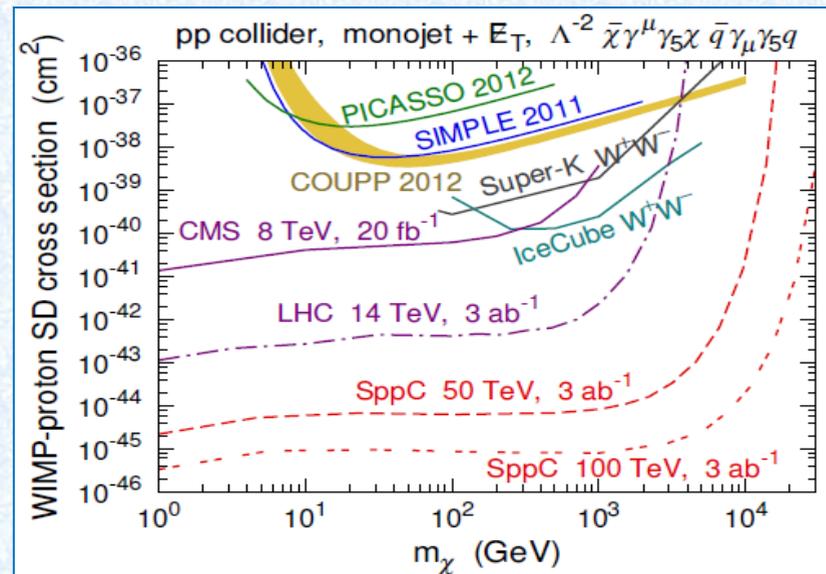
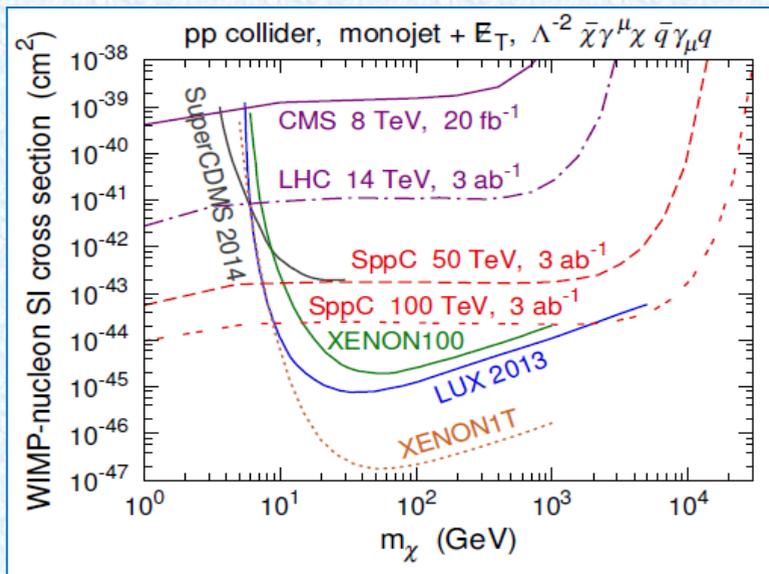
⊕ Deviations of determined masses at a CM energy of 250 GeV with 200 fb⁻¹ of data



Searching for DM at the SPPC



- Advantage: large energy and luminosity for the DM production
- Disadvantage: large backgrounds
- ideal for probing light DM and spin-dependent interactions
- If $Q \gg m_{\text{med}}$, EFT limit is not valid
mediator can be directly produced





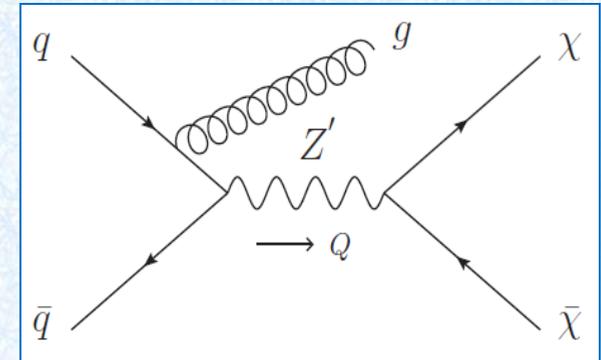
From EFT to simplified model

- ⊕ (minimal) simplified models including Z'

$$\mathcal{L}_{\text{FV}} = \sum_q g_q Z'_\mu \bar{q} \gamma^\mu q + g_\chi Z'_\mu \bar{\chi} \gamma^\mu \chi$$

$$\mathcal{L}_{\text{FA}} = \sum_q g_q Z'_\mu \bar{q} \gamma^\mu \gamma_5 q + g_\chi Z'_\mu \bar{\chi} \gamma^\mu \gamma_5 \chi$$

$$\mathcal{L}_{\text{SV}} = \sum_q g_q Z'_\mu \bar{q} \gamma^\mu q + i g_\chi Z'_\mu [\chi^* \partial^\mu \chi - (\partial^\mu \chi^*) \chi]$$

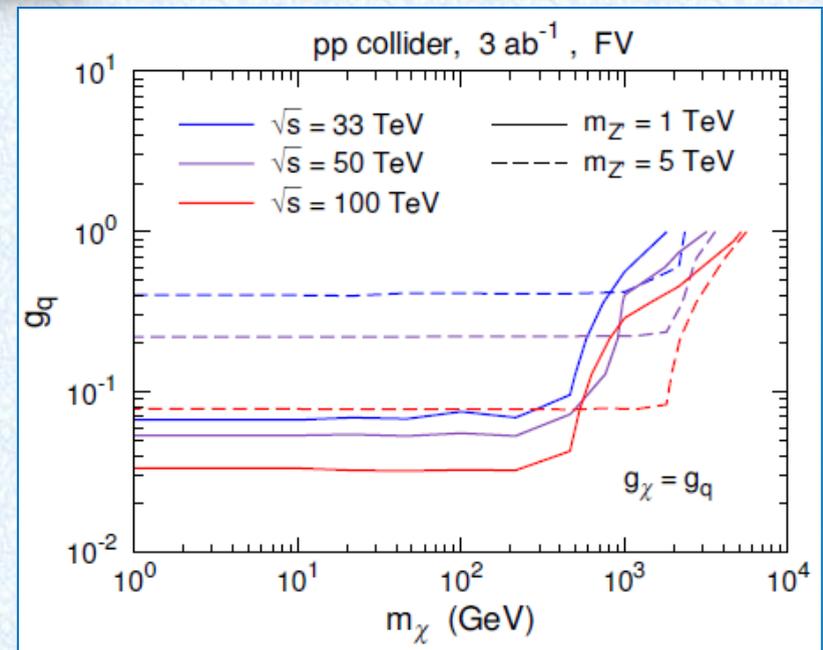


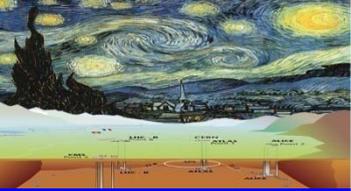
- ⊕ Require $\Gamma_{Z'} < m_{Z'}$

- ⊕ For simplicity, assume $g_q = g_\chi$

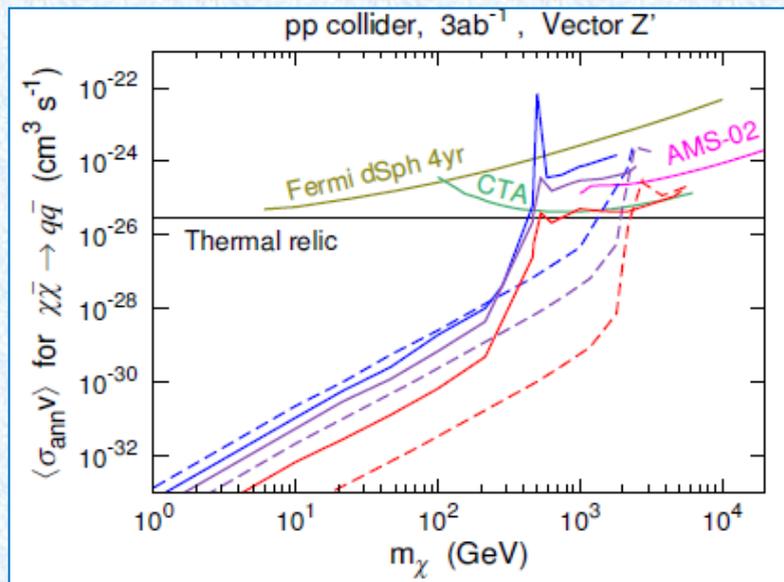
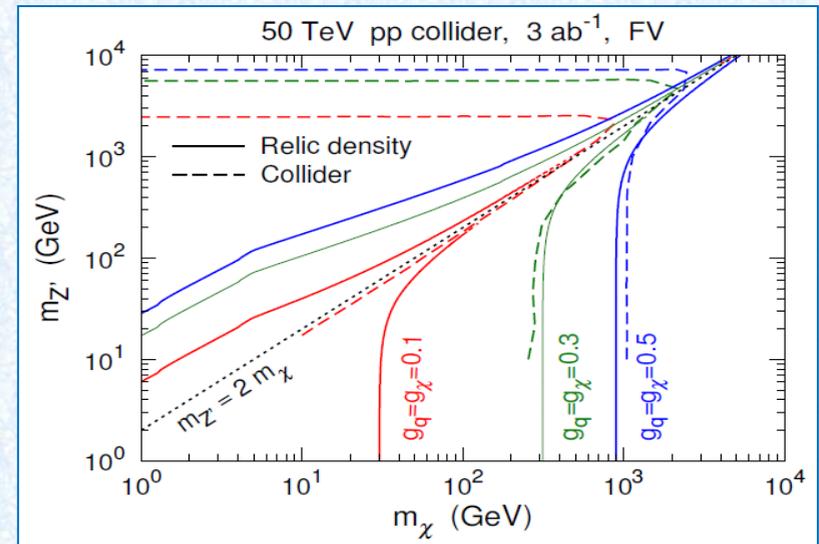
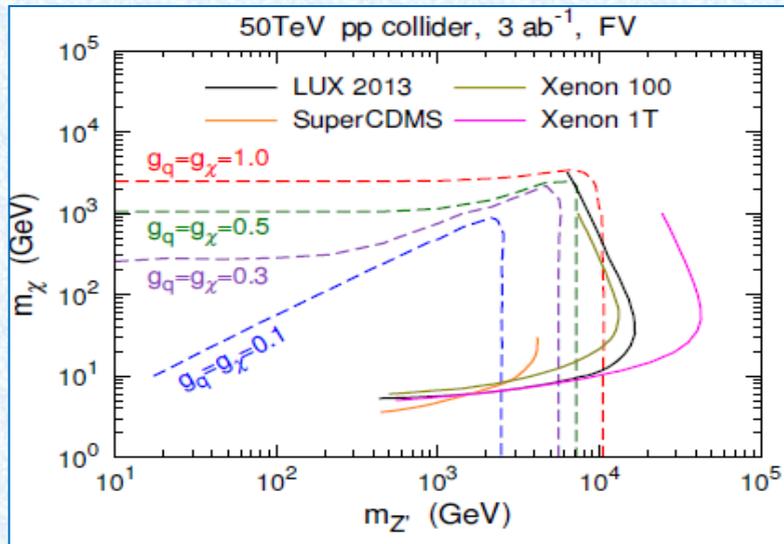
g_q can be limited by di-jet resonance searches.

In the resonance region, Z' width would affect the DM production rate

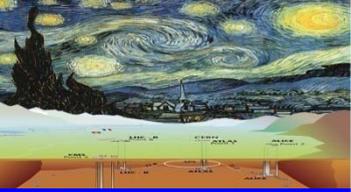




Sensitivities of SppC



- ⊕ Compare with the limits from the direct detection and indirect detection
- ⊕ Compare with the parameter regions for the correct thermal DM relic density
- ⊕ Direct detection limits are derived for $g=0.5$



“Sparticle” simplified model

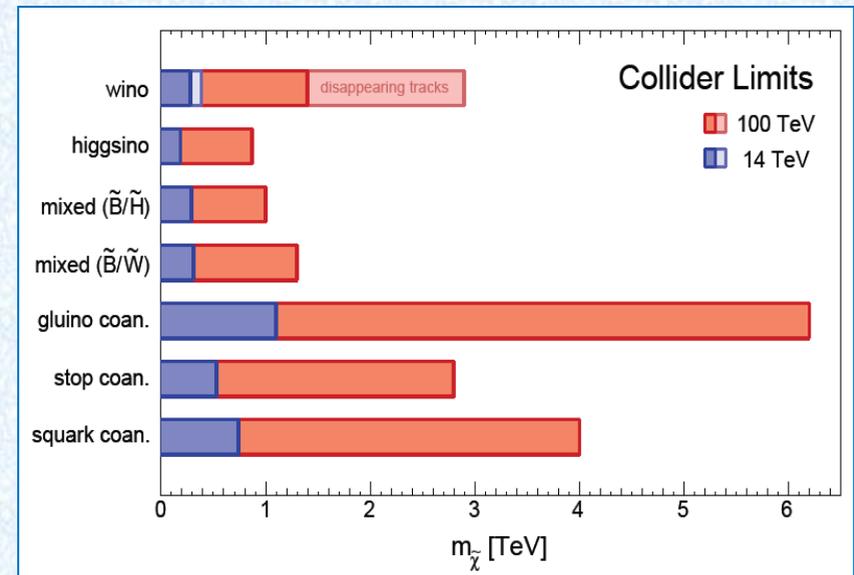
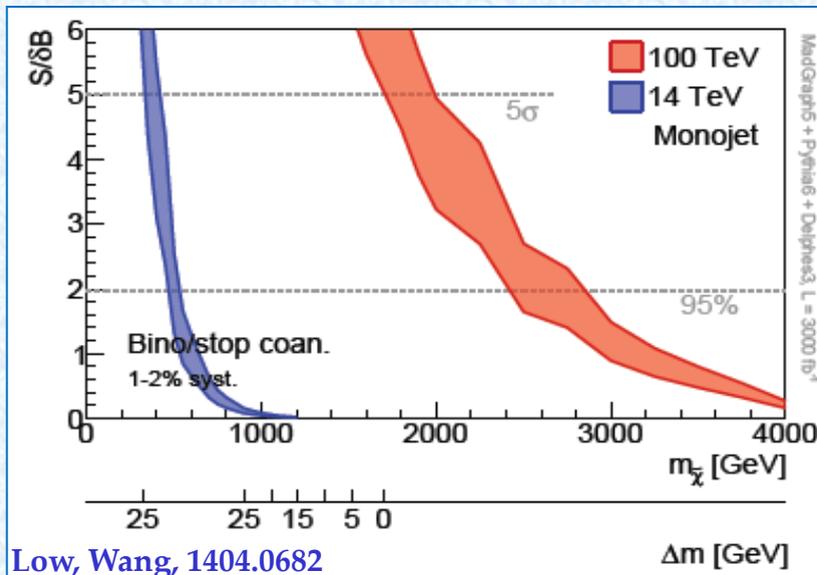
- ⊕ Consider a Lagrangian similar to a SUSY model

$$\Delta L = M_1 \tilde{B} \tilde{B} + M_2 \tilde{W} \tilde{W} + \mu \tilde{H}_u \tilde{H}_d + \sqrt{2} \kappa_1 h^\dagger \tilde{W} \tilde{H}_u + \sqrt{2} \kappa_2 h \tilde{W} \tilde{H}_d + \frac{\kappa'_1}{\sqrt{2}} h^\dagger \tilde{B} \tilde{H}_u + \frac{\kappa'_2}{\sqrt{2}} h \tilde{B} \tilde{H}_d$$

- ⊕ DM may be a pure electroweak state or a mixture

Arkani et. al, 1511.06495

- ⊕ Search strategies are also similar to those for the SUSY: mono-jet, soft leptons, disappearing tracks





Summary

- ⊕ Difficult to directly detect DM signals at e^+e^- colliders
- ⊕ Colliders are ideal tools for probing the light DM
- ⊕ e^+e^- colliders are suitable to measure the DM mass
- ⊕ Complementary to direct and indirect detections



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Thank you !