

Charged Lepton Flavor Violations

CLFV

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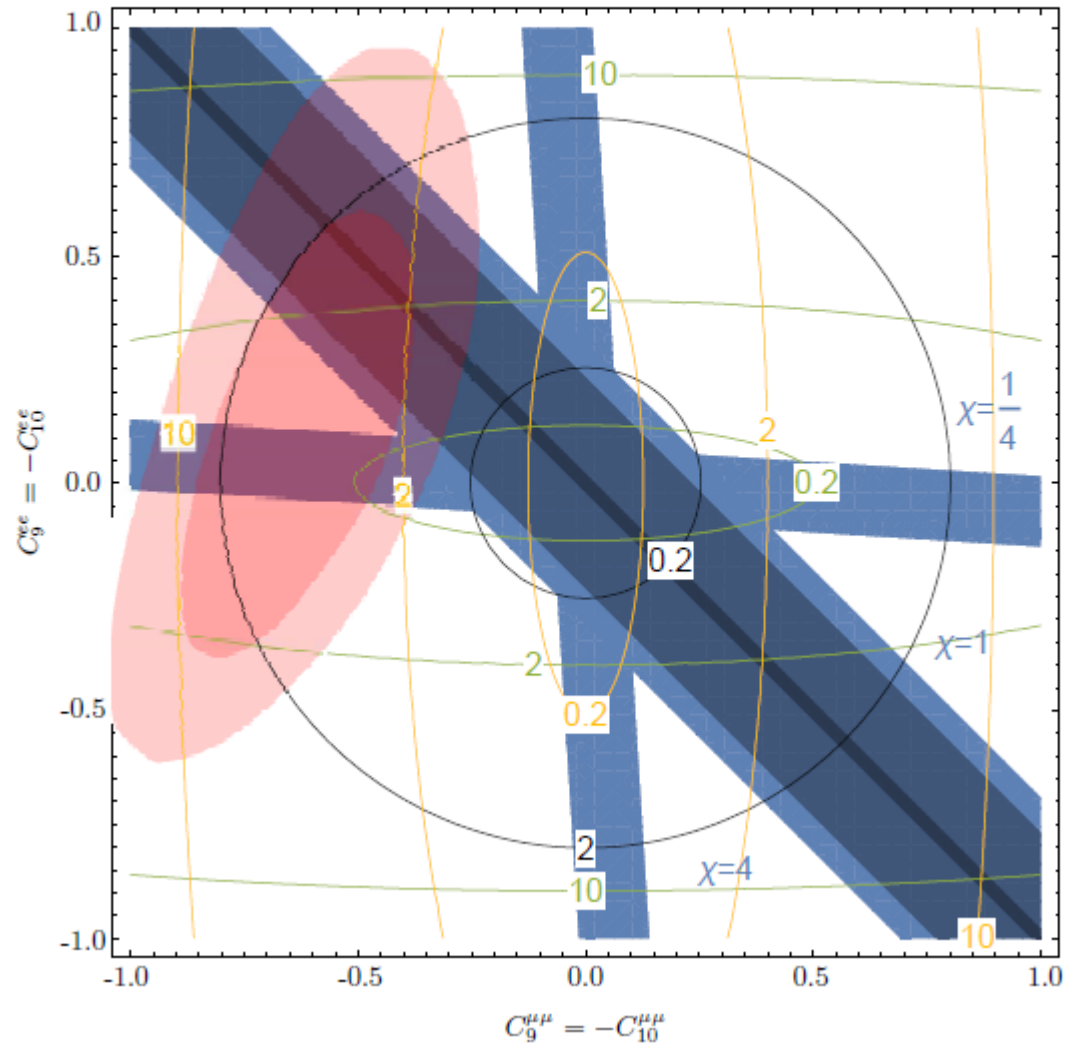
Most stringent LFV upper limits

Process	Current Limit	Next Generation exp
$\tau \rightarrow \mu\eta$	BR < 6.5 E-8 BaBar	10 ⁻⁹ - 10 ⁻¹⁰ (Belle II)
$\tau \rightarrow \mu\gamma$	BR < 6.8 E-8 BaBar	
$\tau \rightarrow \mu\mu\mu$	BR < 3.2 E-8 Belle	
$\tau \rightarrow eee$	BR < 3.6 E-8 Belle	
$K_L \rightarrow e\mu$	BR < 4.7 E-12 BNL	NA62 might improve by O(10)
$K^+ \rightarrow \pi^+e^-\mu^+$	BR < 1.3 E-11 BNL	
$B^0 \rightarrow e\mu$	BR < 7.8 E-8 LHCb	Belle II - LHCb
$B^+ \rightarrow K^+e\mu$	BR < 9.1 E-8 BaBar	
$\mu^+ \rightarrow e^+\gamma$	BR < 4.2 E-13 MEG@PSI	10 ⁻¹⁴ (MEG@PSI)
$\mu^+ \rightarrow e^+e^+e^-$	BR < 1.0 E-12 SINDRUM@PSI	10 ⁻¹⁶ (PSI)
$\mu N \rightarrow eN$	R _{μe} < 7.0 E-13 SINDRUM@PSI	10 ⁻¹⁷ (Mu2e, COMET)

R-parity conserved SUSY: meson decays strongly suppressed by μ constraints
(not true if R –parity is violated): Belayev et al., 2000

The recent LHCb results on possible LFU violations could be a sign of new physics giving rise to LFV: A. Crivellin et al., 2017 (LQ model)

- $b \rightarrow s\mu^+\mu^-$ (1σ)
 - $\text{Br}[\mu \rightarrow e\gamma] < 4.2 \cdot 10^{-13}$ with Φ_3
 - $\text{Br}[\mu \rightarrow e\gamma] < 4.2 \cdot 10^{-13}$ with V_1^μ
 - $\text{Br}[\mu \rightarrow e\gamma] < 4.2 \cdot 10^{-13}$ with V_3^μ
- $b \rightarrow s\mu^+\mu^-$ (2σ)
 - $10^8 \text{Br}[B \rightarrow K\mu^\pm e^\mp]$ with $\gamma = 1/2$
 - $10^8 \text{Br}[B \rightarrow K\mu^\pm e^\mp]$ with $\gamma = 1$
 - $10^8 \text{Br}[B \rightarrow K\mu^\pm e^\mp]$ with $\gamma = 2$



Search for LFV decays of the Higgs boson at LHC

(Phys. Lett. B 763(2016) 472; CMS-PAS-HIG-17-001)

$B(H \rightarrow \mu\tau) < 0.25\%$; $B(H \rightarrow e\tau) < 0.61\%$; $B(H \rightarrow e\mu) < 0.035\%$

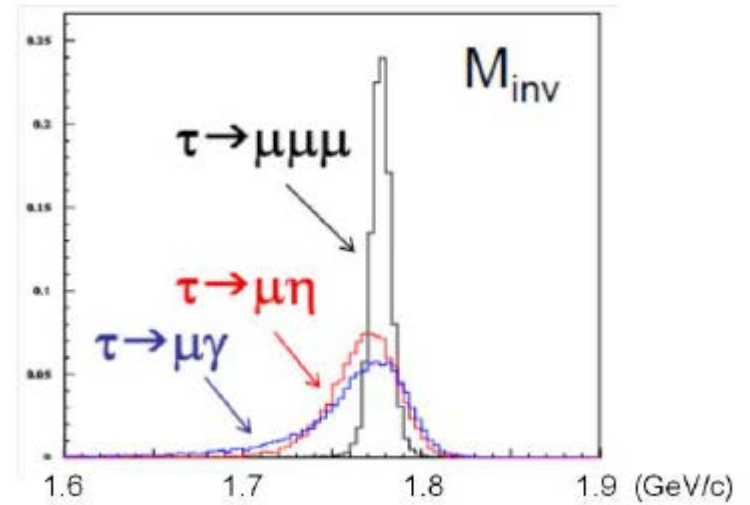
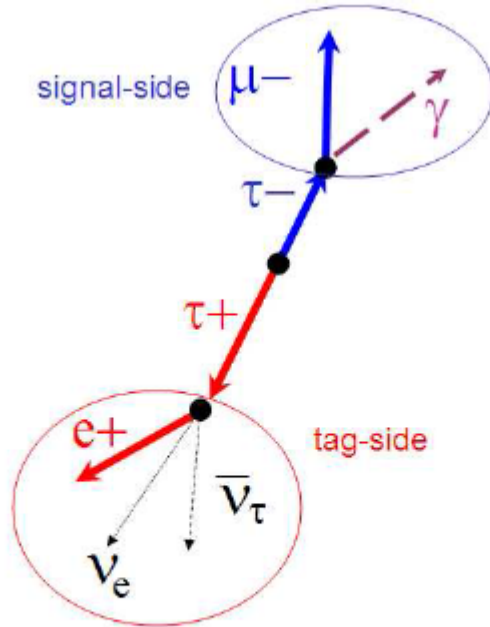
Higgs induced CLFV: parameterized with the following lagrangian

$$-\mathcal{L} \supset (m_e)_i \bar{e}_{Li} e_{Ri} + (Y_e^h)_{ij} \bar{e}_{Li} e_{Rj} h + h.c.$$

Process	Coupling	Bound	
$h \rightarrow \mu e$	$\sqrt{ Y_{\mu e}^h ^2 + Y_{e\mu}^h ^2}$	$< 5.4 \times 10^{-4}$	LHC
$\mu \rightarrow e\gamma$	$\sqrt{ Y_{\mu e}^h ^2 + Y_{e\mu}^h ^2}$	$< 2.1 \times 10^{-6}$	
$\mu \rightarrow eee$	$\sqrt{ Y_{\mu e}^h ^2 + Y_{e\mu}^h ^2}$	$\lesssim 3.1 \times 10^{-5}$	
$\mu Ti \rightarrow e Ti$	$\sqrt{ Y_{\mu e}^h ^2 + Y_{e\mu}^h ^2}$	$< 1.2 \times 10^{-5}$	
$h \rightarrow \tau e$	$\sqrt{ Y_{\tau e}^h ^2 + Y_{e\tau}^h ^2}$	$< 2.3 \times 10^{-3}$	LHC
$\tau \rightarrow e\gamma$	$\sqrt{ Y_{\tau e}^h ^2 + Y_{e\tau}^h ^2}$	< 0.014	
$\tau \rightarrow eee$	$\sqrt{ Y_{\tau e}^h ^2 + Y_{e\tau}^h ^2}$	$\lesssim 0.12$	
$h \rightarrow \tau\mu$	$\sqrt{ Y_{\tau\mu}^h ^2 + Y_{\mu\tau}^h ^2}$	$< 1.4 \times 10^{-3}$	LHC
$\tau \rightarrow \mu\gamma$	$\sqrt{ Y_{\tau\mu}^h ^2 + Y_{\mu\tau}^h ^2}$	< 0.016	
$\tau \rightarrow \mu\mu\mu$	$\sqrt{ Y_{\tau\mu}^h ^2 + Y_{\mu\tau}^h ^2}$	$\lesssim 0.25$	

Leptons: τ

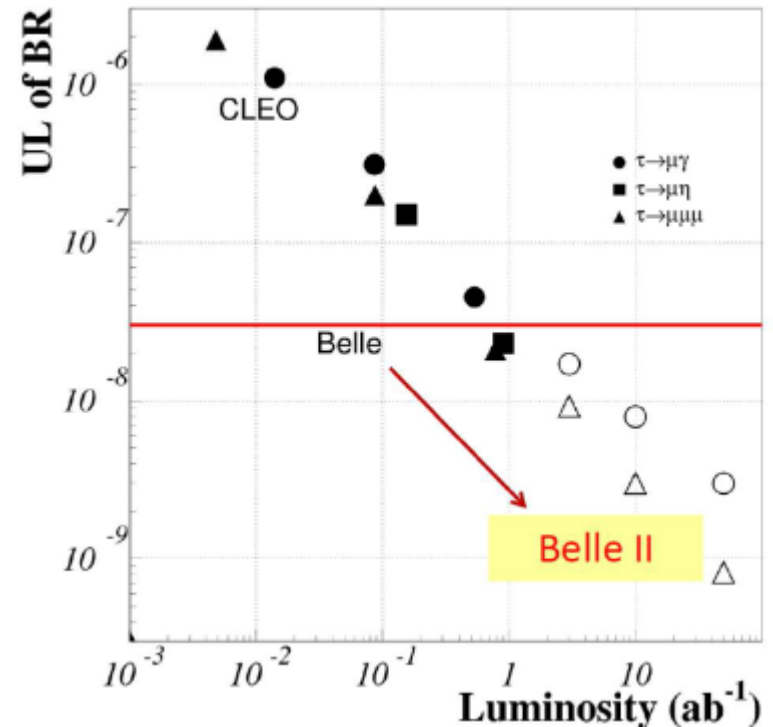
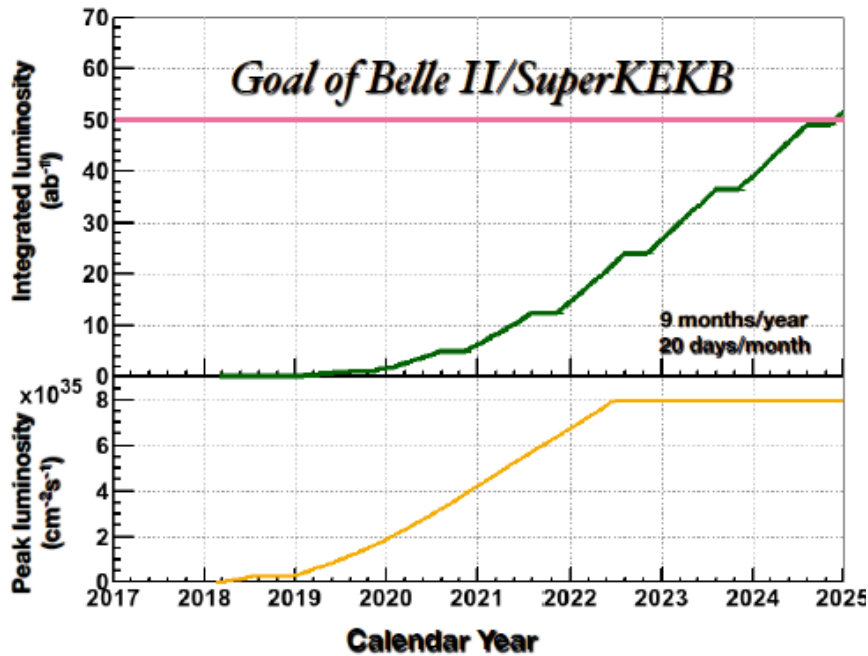
τ cLFV sensitivity (BelleII): 1



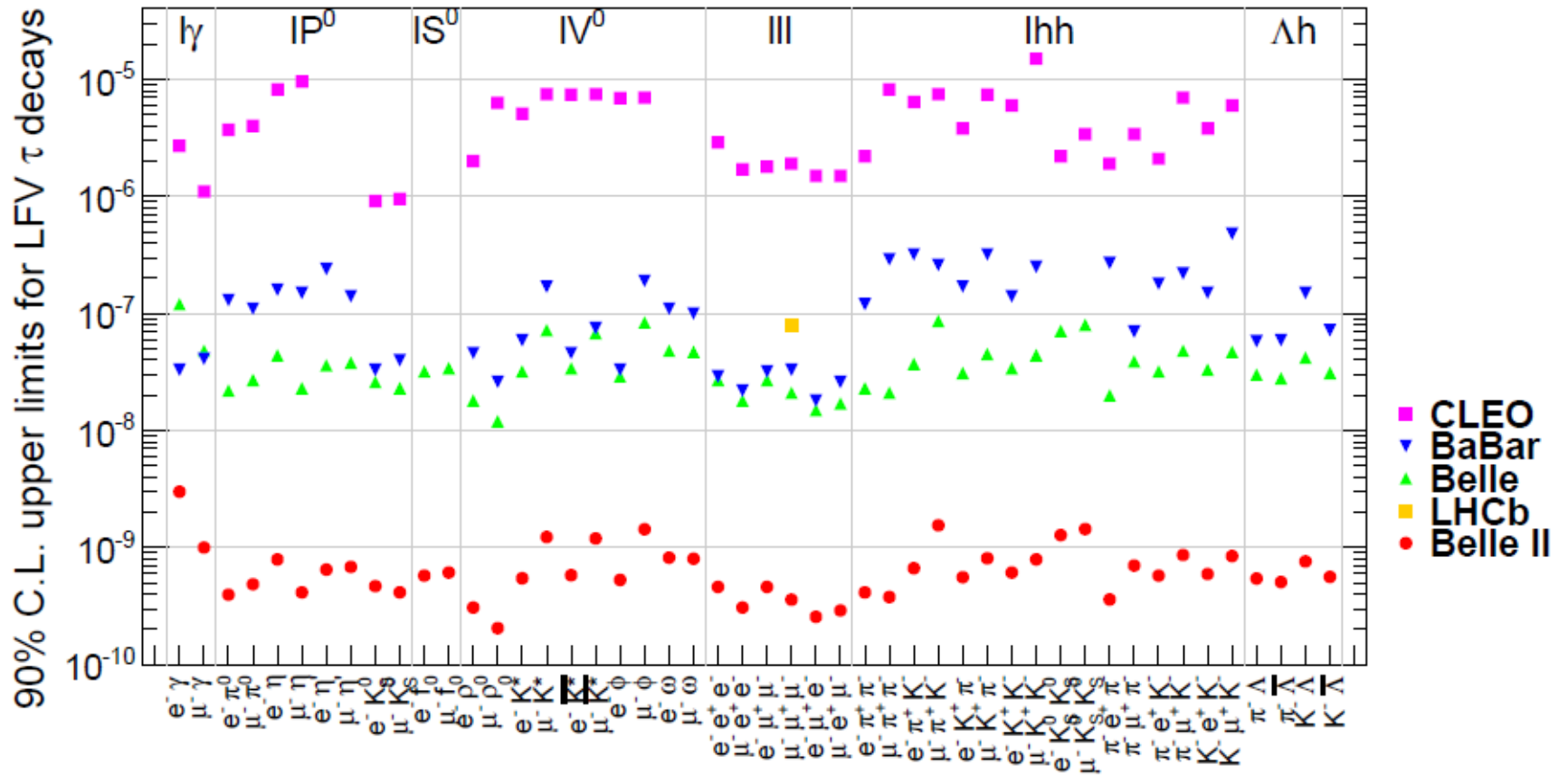
$$m_{\text{inv}} = \sqrt{E_{\mu\gamma}^2 - p_{\mu\gamma}^2} \quad \Delta E = E_{\mu\gamma}^{\text{CM}} - E_{\text{beam}}^{\text{CM}}$$

τ cLFV sensitivity (BelleII): 2

- Luminosity of 50 ab^{-1} ($5 \times 10^5 \tau \bar{\tau}$) reachable by 2025
- Golden channels ($\tau \rightarrow 3\mu$) might be non background limited

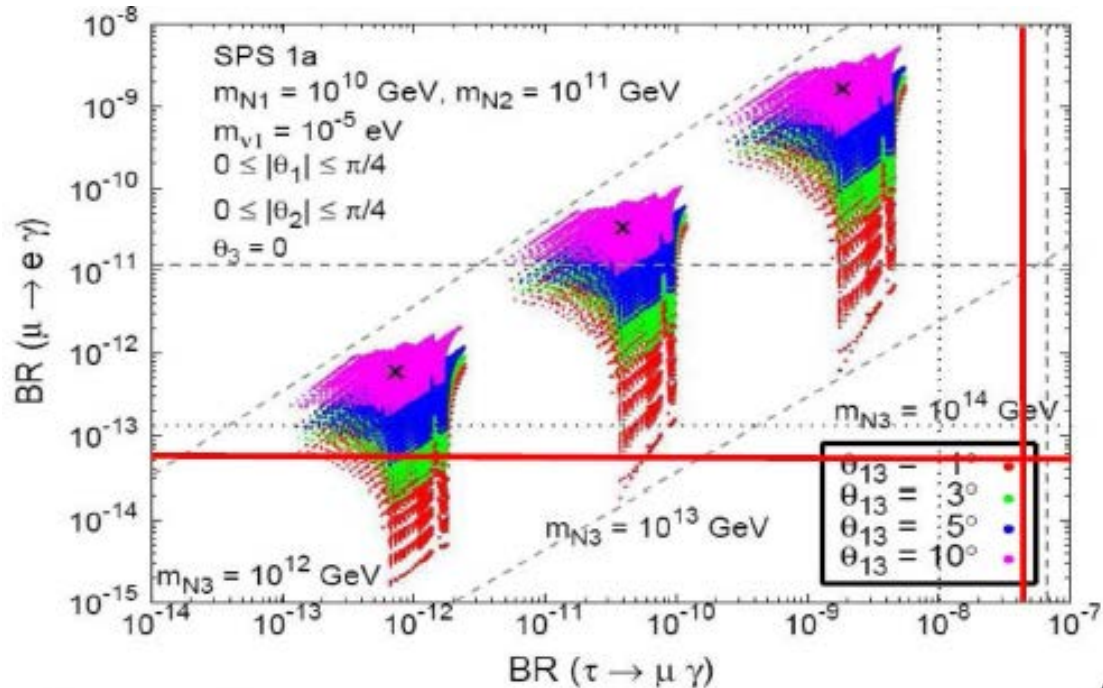


The foreseen improvements on all the CLFV decays



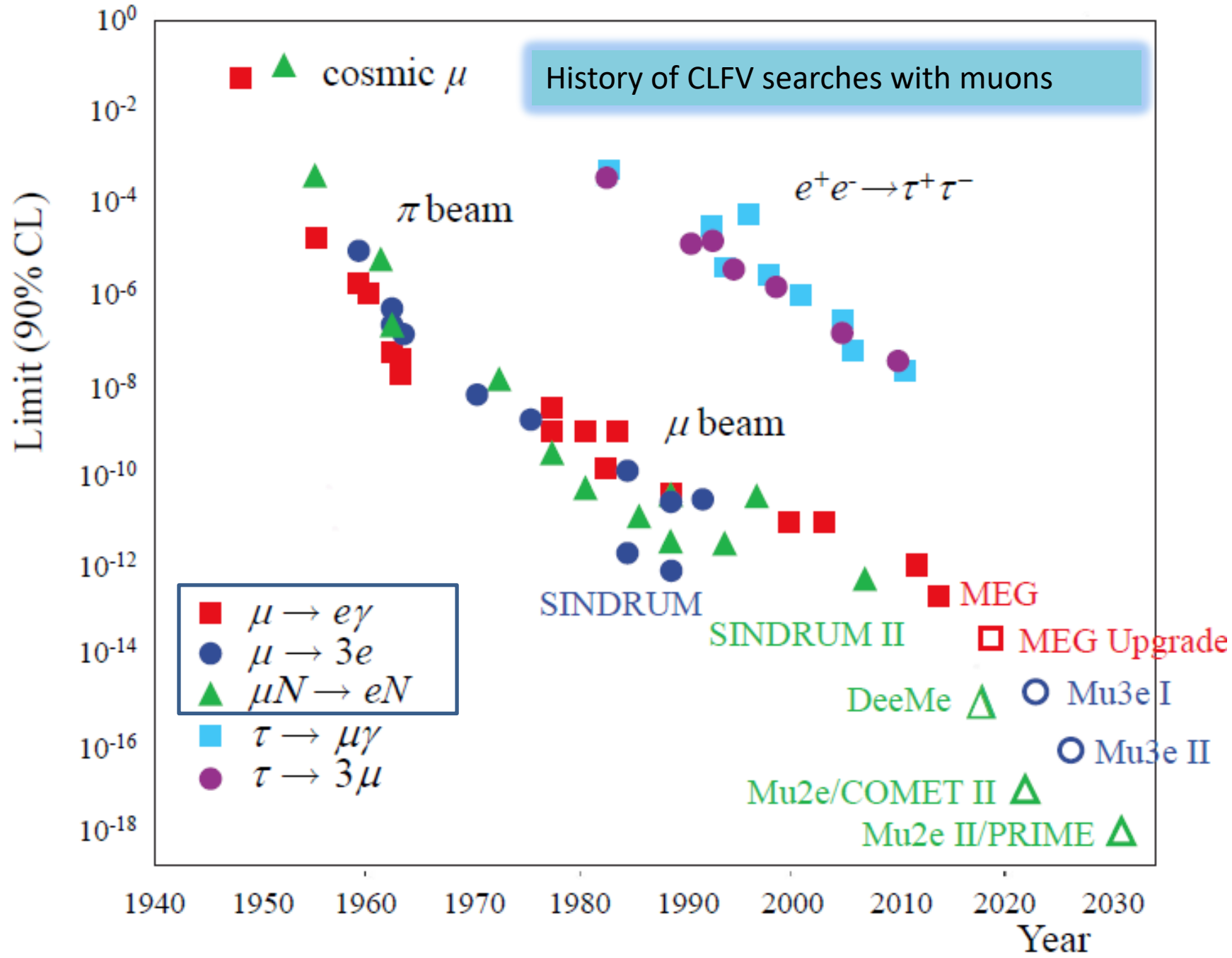
The sensitivity of the τ channels wrt the μ channels again depends very much on the model considered

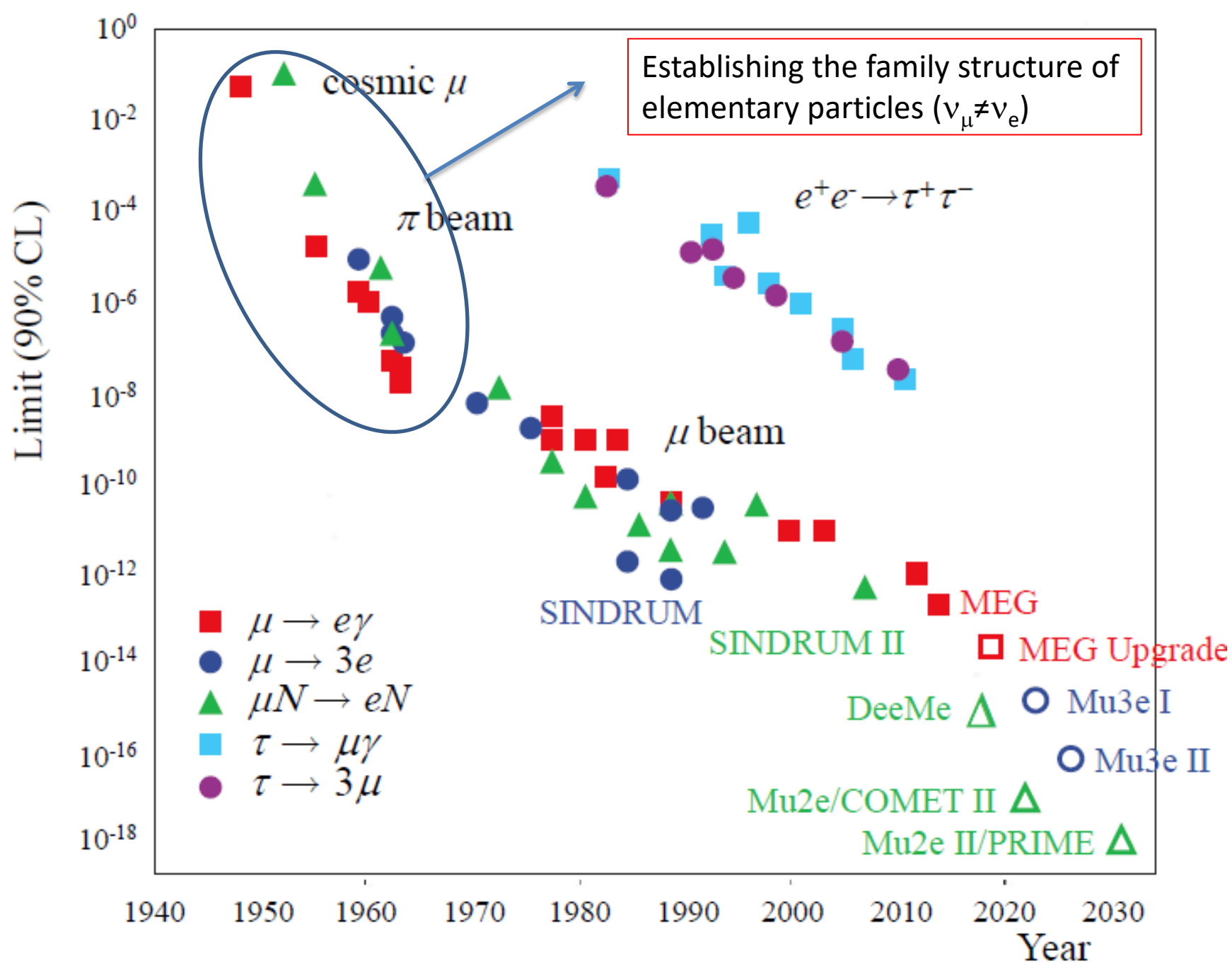
Antusch et al., 2006

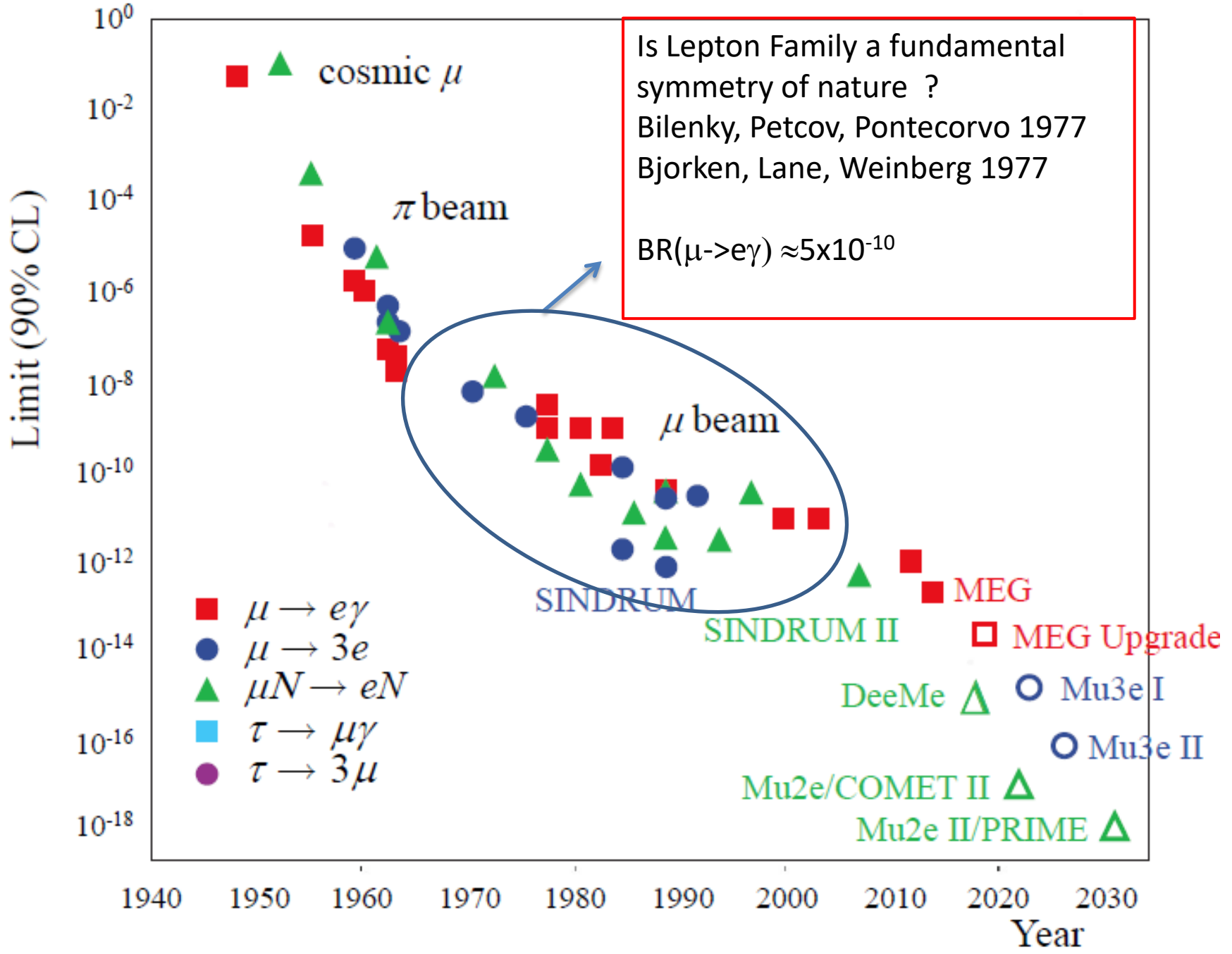


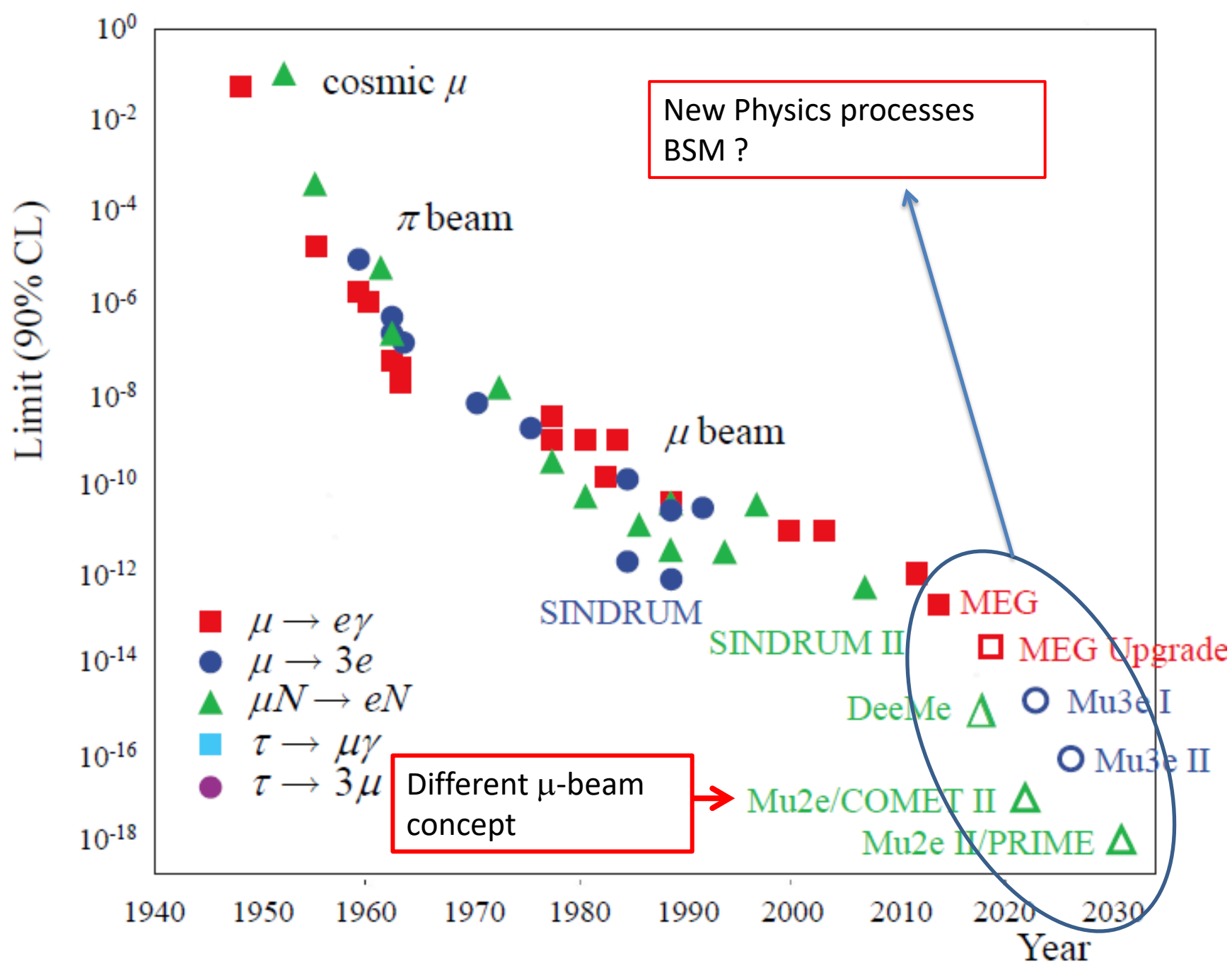
Bu the enhancement of the τ decay can be as large as 10^4 : Barbieri et al., 1994

History of CLFV searches with muons









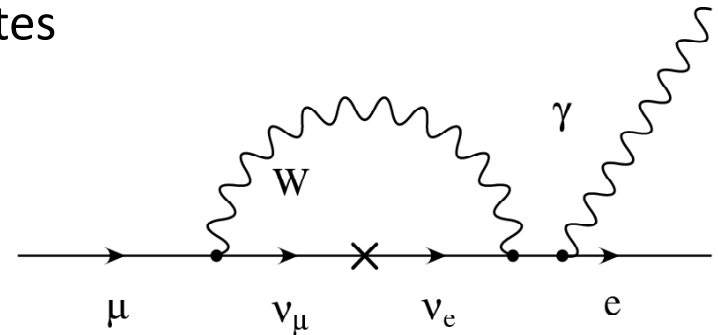
Why are these processes so sensitive to BSM ?

First: CLFV in the SM are NOT observable

In the pre- ν oscillations SM cLFV amplitudes are 0 due to the fact that neutrino masses are 0

But even the introduction of conventional neutrino masses to account for oscillations gives negligible rates

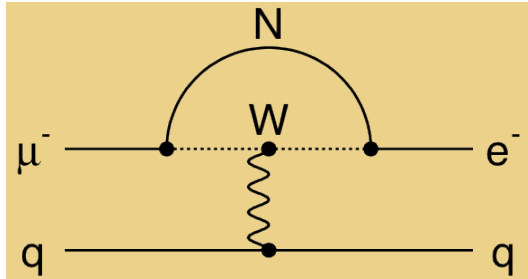
$\mu \rightarrow e \gamma$ rate in the standard model
after neutrino oscillations



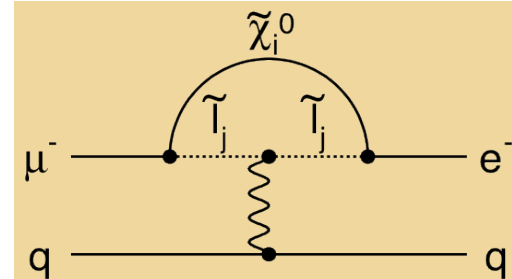
$$BR \sim \left(\frac{\delta m_\nu^2}{M_w^2} \right)^2 ; \delta m_\nu^2 \sim 10^{-5} eV^2 \Rightarrow BR \sim 10^{-54}$$

CLFV are VERY clean channels for new physics searches

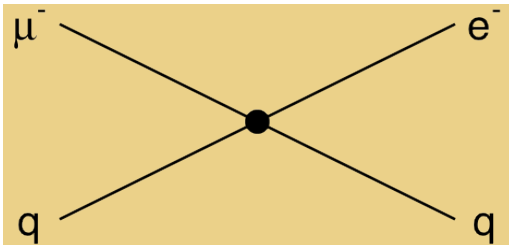
Second: as soon as one starts adding terms to the SM Lagrangian:



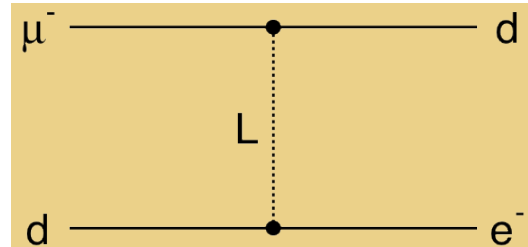
Heavy Neutrinos
(see saw)



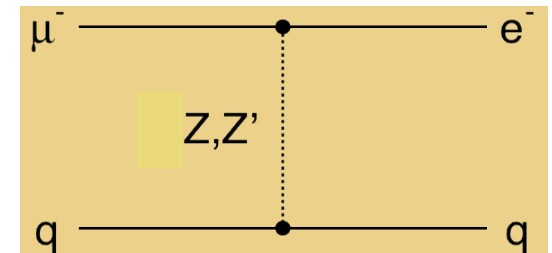
Supersymmetry



Compositeness



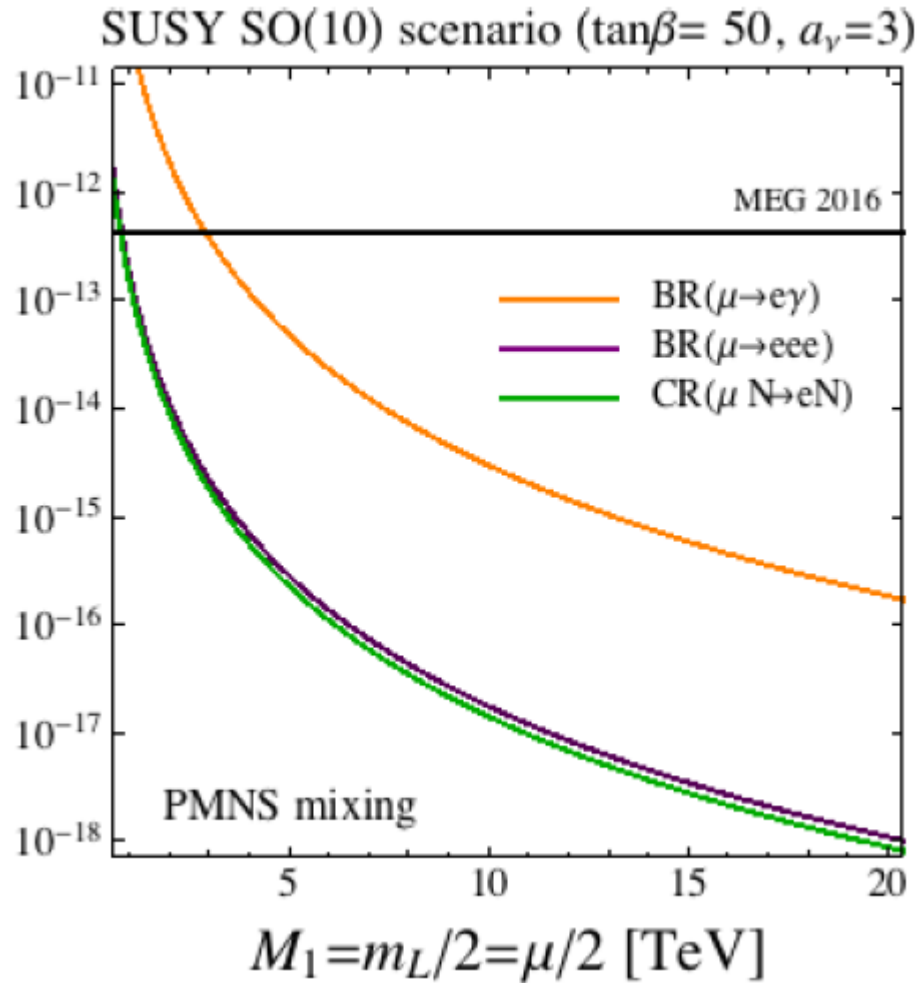
Leptoquarks



New Heavy Bosons /
Anomalous Couplings

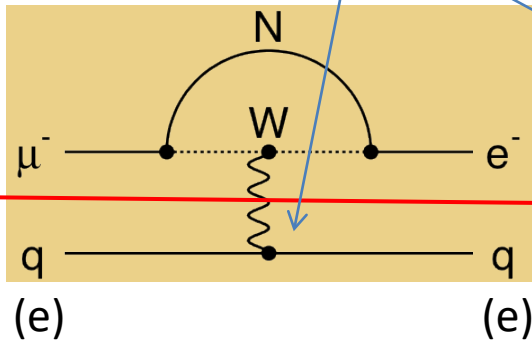
Charged Lepton flavor violating amplitudes arise at measurable levels !!

An example...

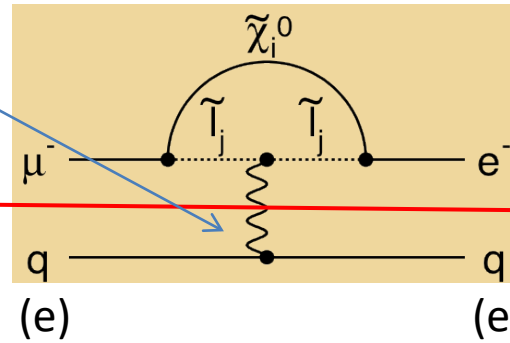


SUSY seesaw SO(10) with PMNS slepton mixing; Calibbi, Signorelli 2017 and references therein.

In case model prefer dipole operator (Supersymmetry) $\mu \rightarrow e \gamma$ is enhanced by roughly a factor $1/\alpha$ wrt $\mu \rightarrow e$ conversion or $\mu \rightarrow 3e$

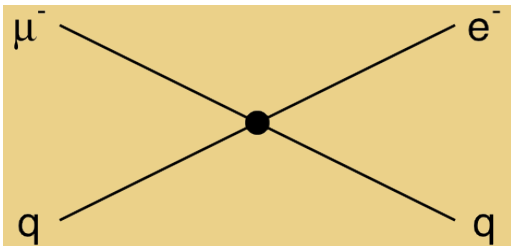


Heavy Neutrinos
(see saw)

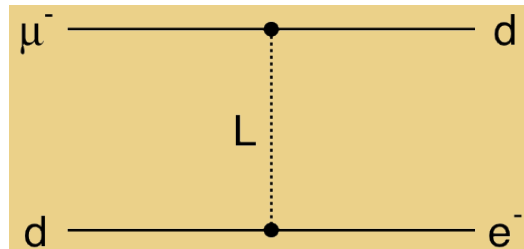


Supersymmetry

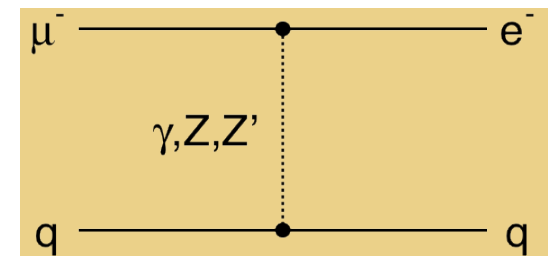
$\mu \rightarrow e \gamma$



Compositeness



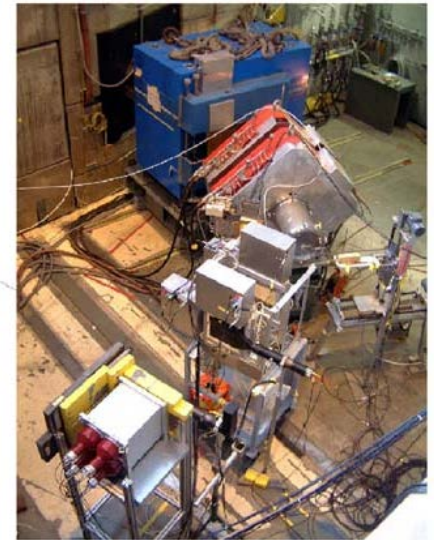
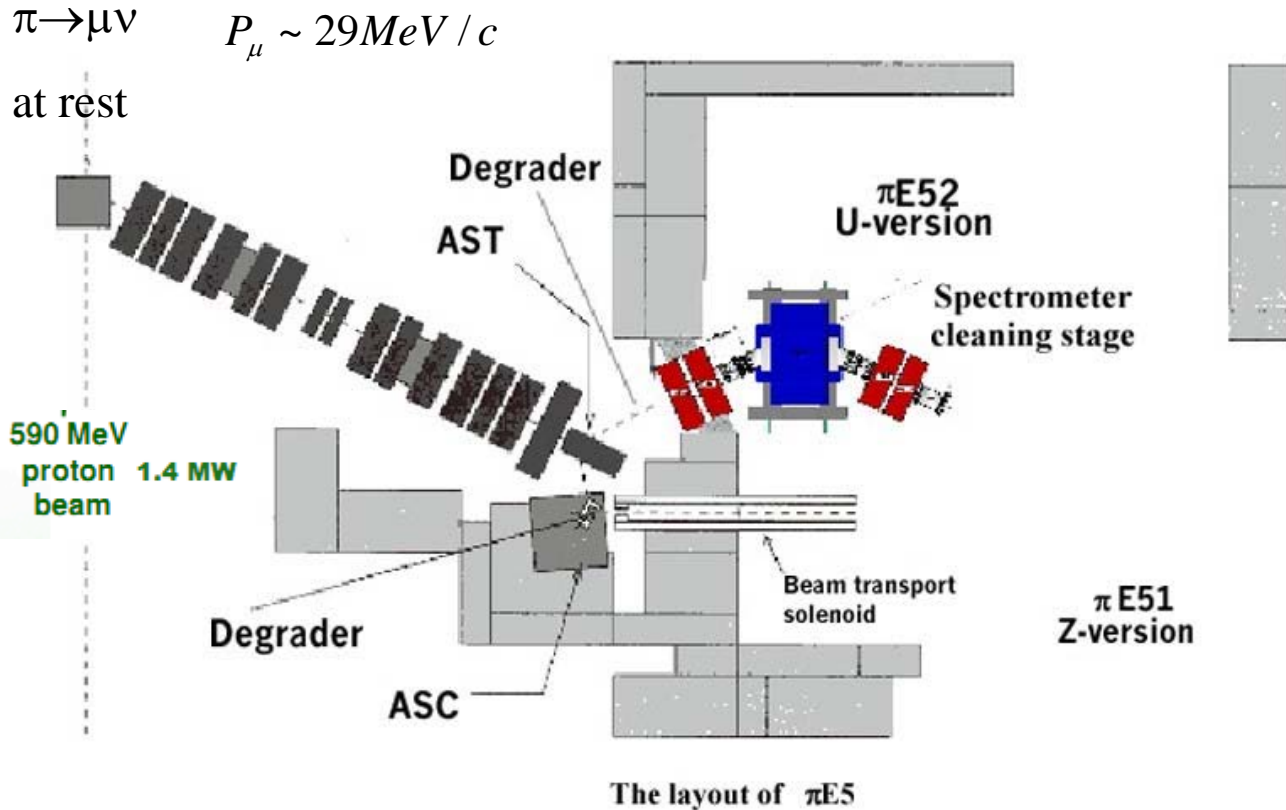
Leptoquarks



New Heavy Bosons /
Anomalous Couplings

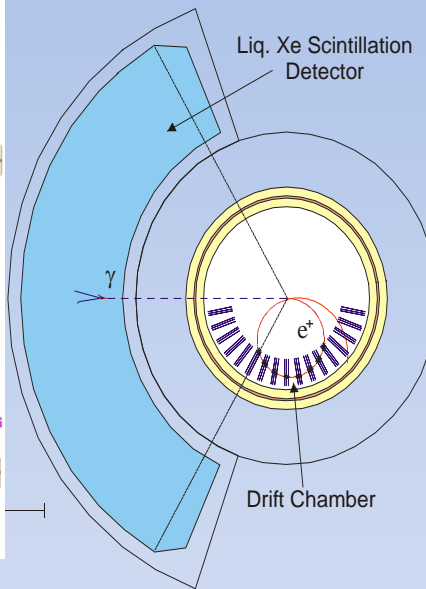
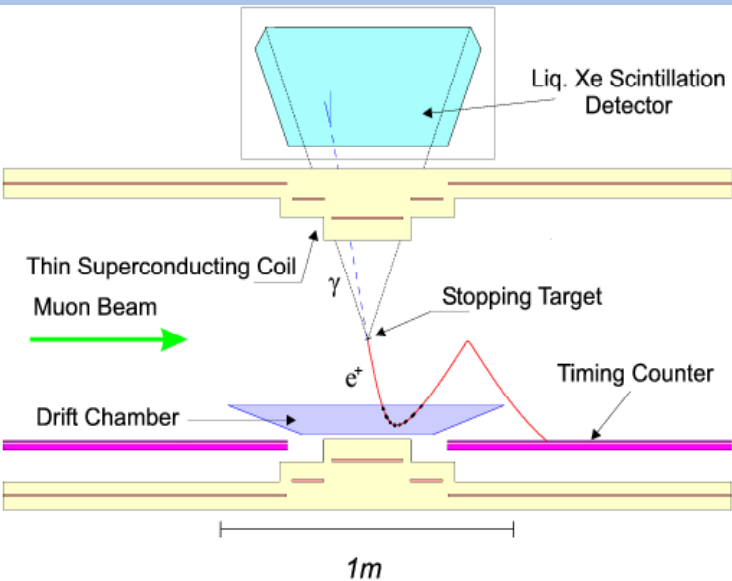
In other cases $\mu \rightarrow e$ conversion or $\mu \rightarrow 3e$ dominate \rightarrow Complementarity + possibility of distinguishing among different models

The PSI surface muon beam



- $10^8 \mu/s$ from the decay at rest of π^+ on the target surface (the μ range is approx. $.1 \text{ gr/cm}^3$): almost completely polarized
- It is possible to focalize and stop the μ beam in a thin target

Used by the MEG experiment...

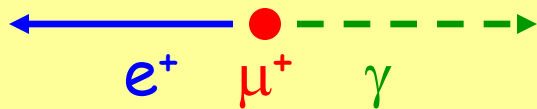


1. Stopped beam of $3 \cdot 10^7 \mu$ /sec in a $150 \mu\text{m}$ target
2. Solenoid spectrometer & drift chambers for e^+ momentum
3. Scintillation counters for e^+ timing
4. Liquid Xenon calorimeter for γ detection (scintillation)

Signal and background

signal

$$\mu \rightarrow e \gamma$$



$$\theta_{e\gamma} = 180^\circ$$

$$E_e = E_\gamma = 52.8 \text{ MeV}$$

$$T_e = T_\gamma$$

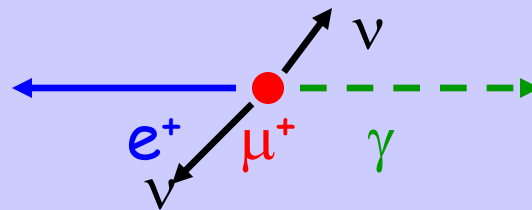
background

accidental

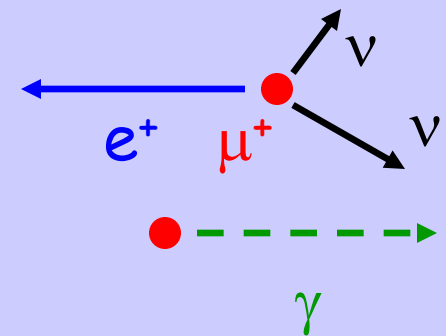
$$\mu \rightarrow e \nu \nu$$

physical

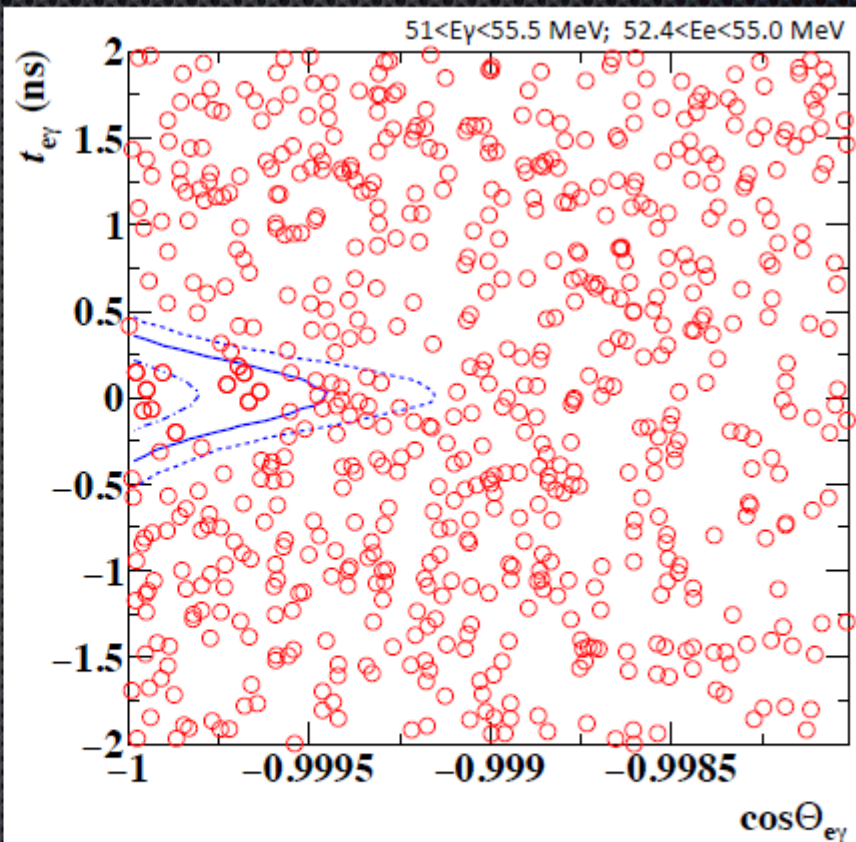
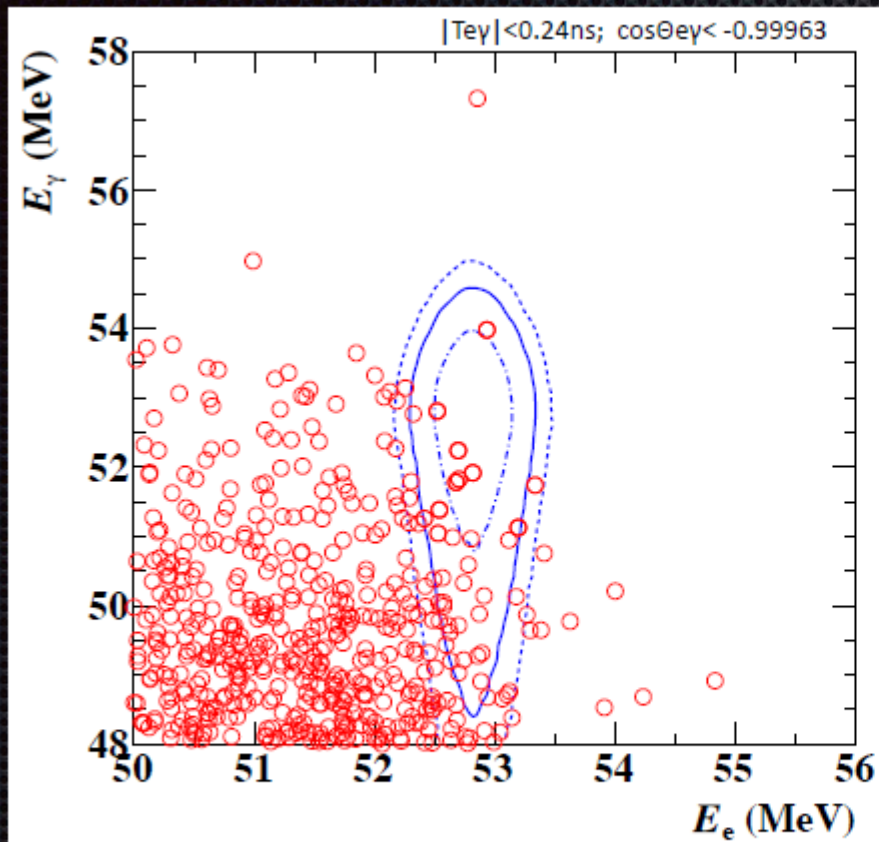
$$\mu \rightarrow e \gamma \nu \nu$$



$$\left\{ \begin{array}{l} \mu \rightarrow e \gamma \nu \nu \\ ee \rightarrow \gamma \gamma \\ eZ \rightarrow eZ \gamma \end{array} \right.$$

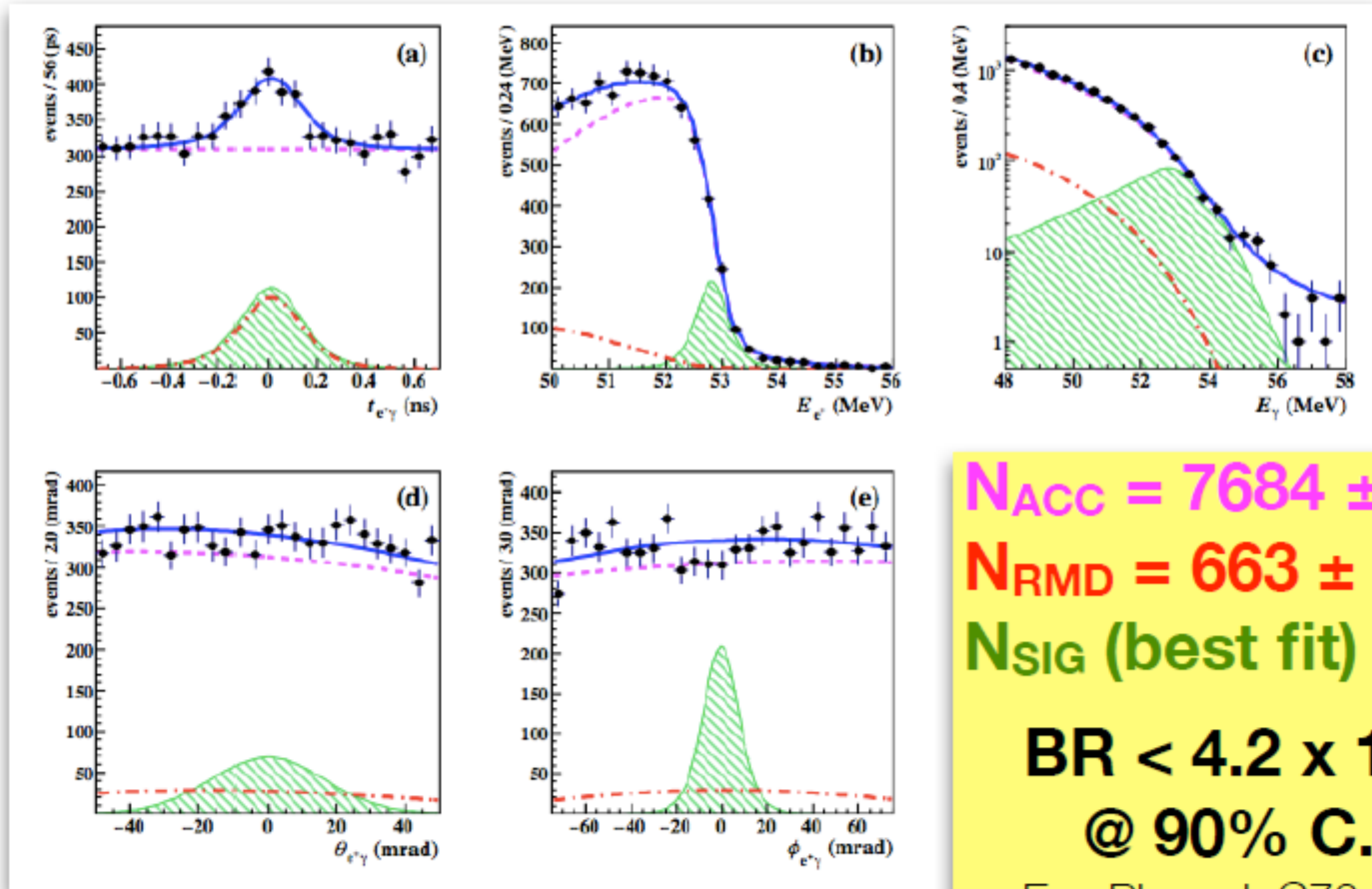


The interesting 4D region split in 2 bidimensional plots



signal contours of 1, 1.64, 2σ are shown

A more quantitative comparison...



$N_{\text{Acc}} = 7684 \pm 103$
 $N_{\text{RMD}} = 663 \pm 59$
 $N_{\text{SIG}} (\text{best fit}) = -2.2$

 $\text{BR} < 4.2 \times 10^{-13}$
@ 90% C.L.
 Eur. Phys. J. C76 (2016)

Magnified signal (BR = 4×10^{-11})

The relative angle is split into zenith and azimuth

The MEG II detector

Liquid Xenon Gamma-ray Detector

COBRA
Superconducting
Magnet

Gamma ray

better uniformity w/
VUV-sensitive
 $12 \times 12 \text{ mm}^2$ SiPM

x2 resolution everywhere

full available
intensity
 $7 \times 10^7 / \text{s}$

Muon

Drift Chamber
single-volume $\text{He:iC}_4\text{H}_{10}$
small stereo cells

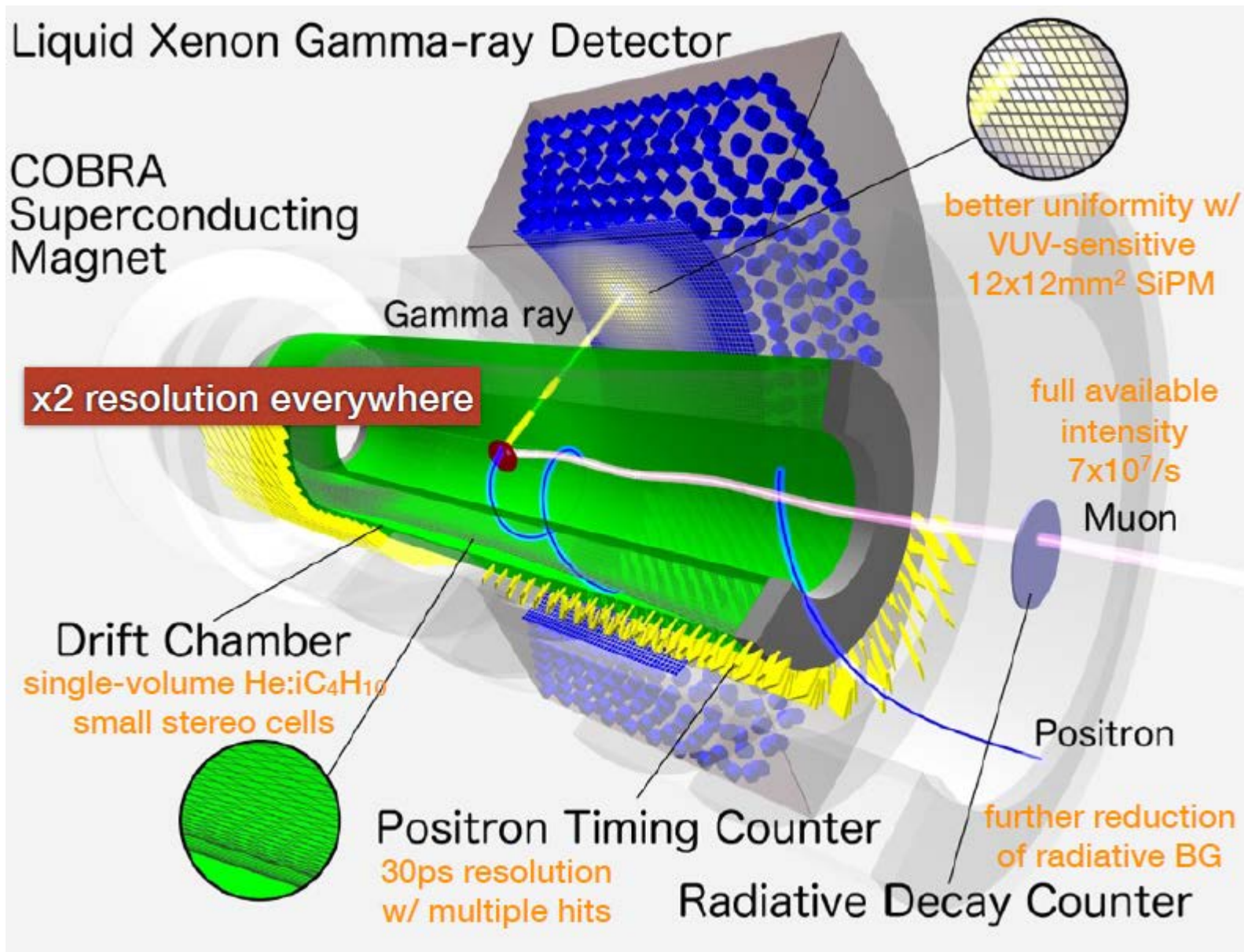
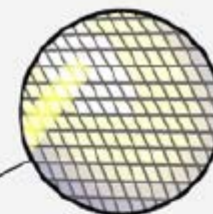
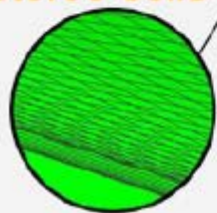
Positron

Positron Timing Counter

30ps resolution
w/ multiple hits

further reduction
of radiative BG

Radiative Decay Counter



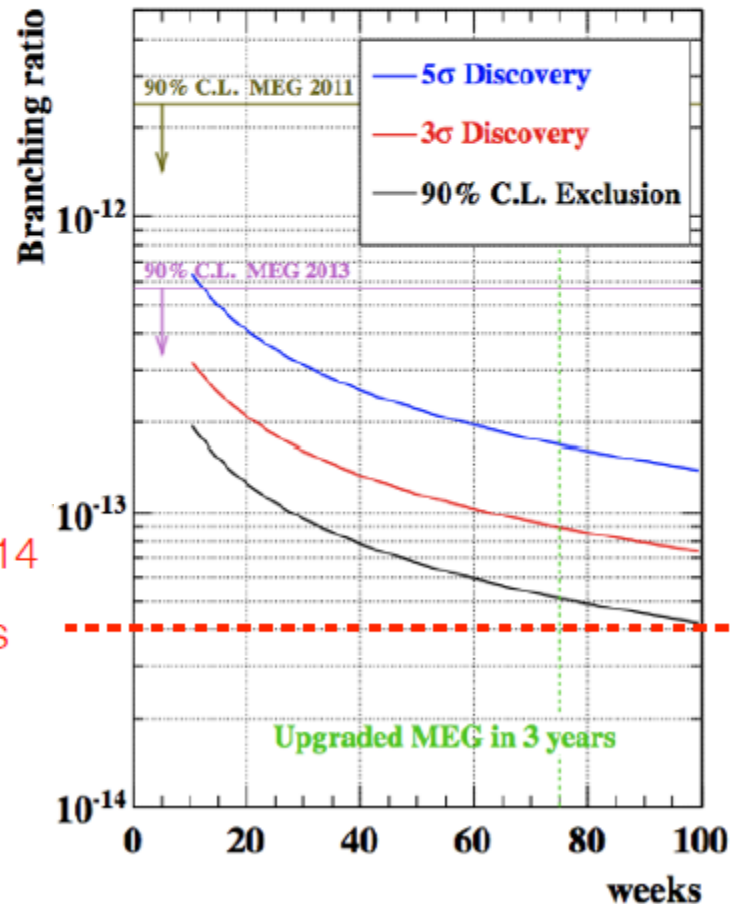
For a final sensitivity of 4×10^{-14}

- MEG-II is expected to start taking data with the full detector next year
- x10 improvement in sensitivity w.r.t. MEG

TABLE VIII: Resolution (Gaussian σ) and efficiencies for MEG upgrade

PDF parameters	Present MEG	Upgrade scenario
$\sigma_{E_{e^-}}$ (keV)	380	110
$e^+ \sigma_\theta$ (mrad)	9	5
$e^+ \sigma_\phi$ (mrad)	11	5
$e^+ \sigma_z / \sigma_y(\text{core})$ (mm)	2.0/1.0	1.2/0.7
$\frac{\sigma_{E_\gamma}}{E_\gamma}$ (%) $w > 2$ cm	1.6	1.0
γ position at LXe $\sigma_{(x,y)} - \sigma_w$ (mm)	4	2
γ - e^+ timing (ps)	120	80
Efficiency (%)		
trigger	≈ 99	≈ 99
γ reconstruction	60	60
e^+ reconstruction	40	95
event selection	80	85

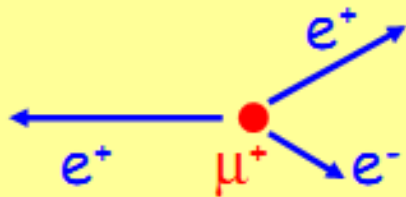
4×10^{-14}
in 4 years



The same μ beam will be used by $\mu \rightarrow 3e$

signal

$\mu \rightarrow eee$



$$\Sigma p_i = 0$$

Vertex

$$\Sigma E_e = m_\mu$$

$$te^+ = te^+ = te^-$$

background

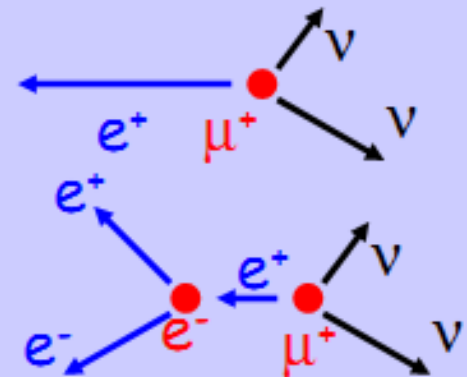
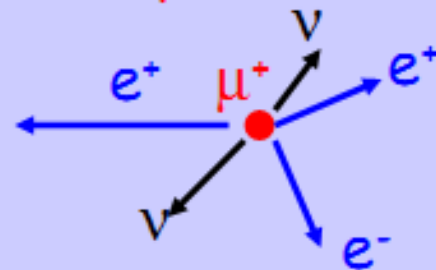
accidental

$\mu \rightarrow e \nu \nu$

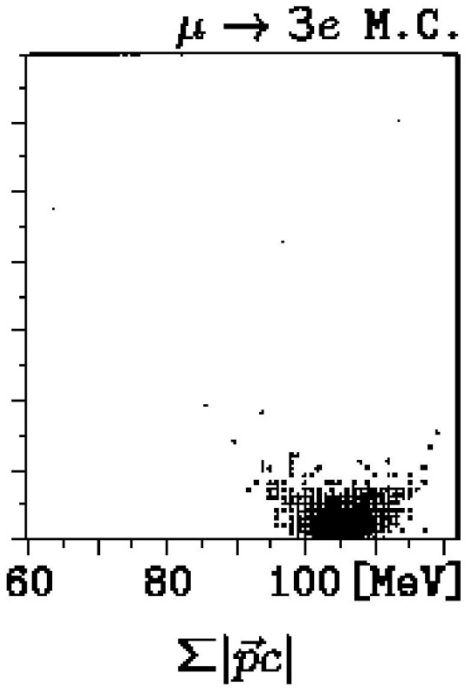
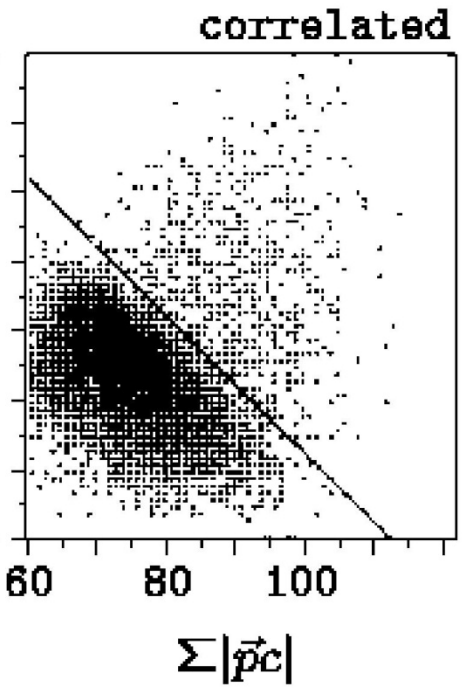
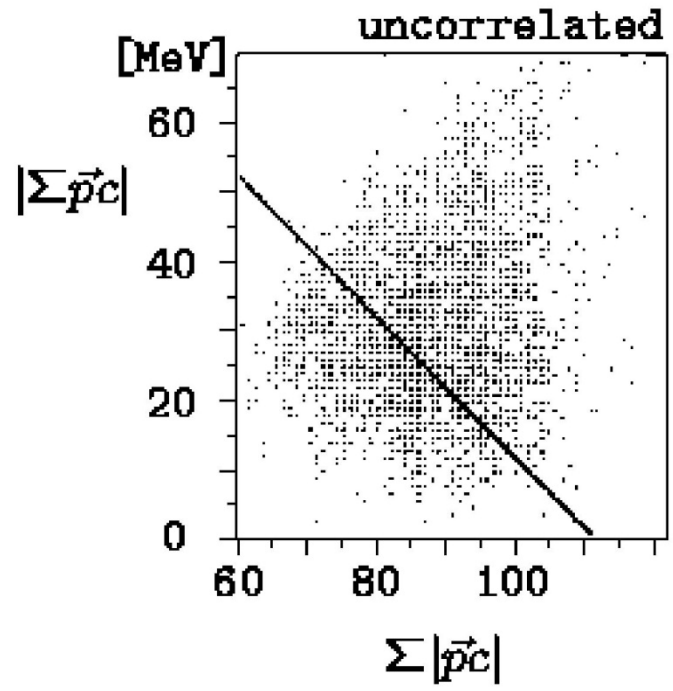
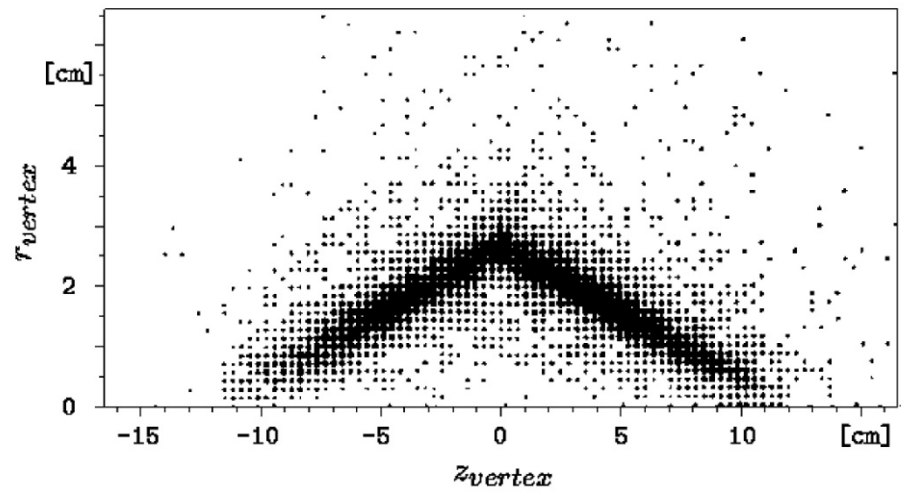
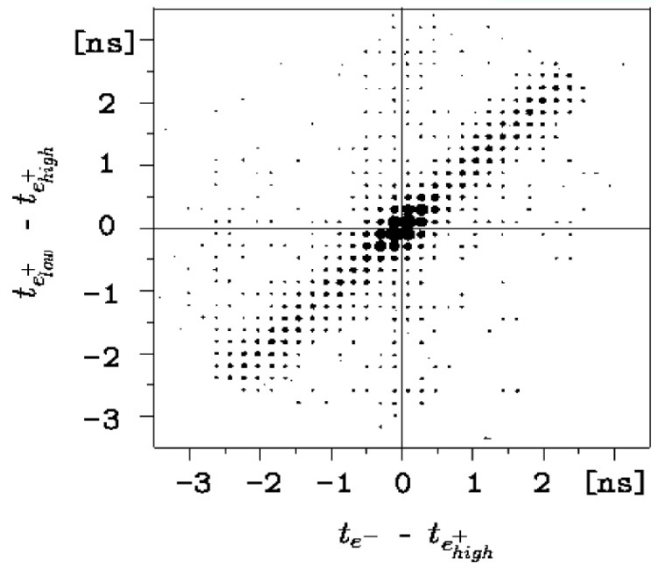
$\left\{ \begin{array}{l} \mu \rightarrow e \nu \nu \\ e^+e^- \rightarrow e^+e^- \end{array} \right.$

correlated
(prompt)

$\mu \rightarrow e e e \nu \nu$



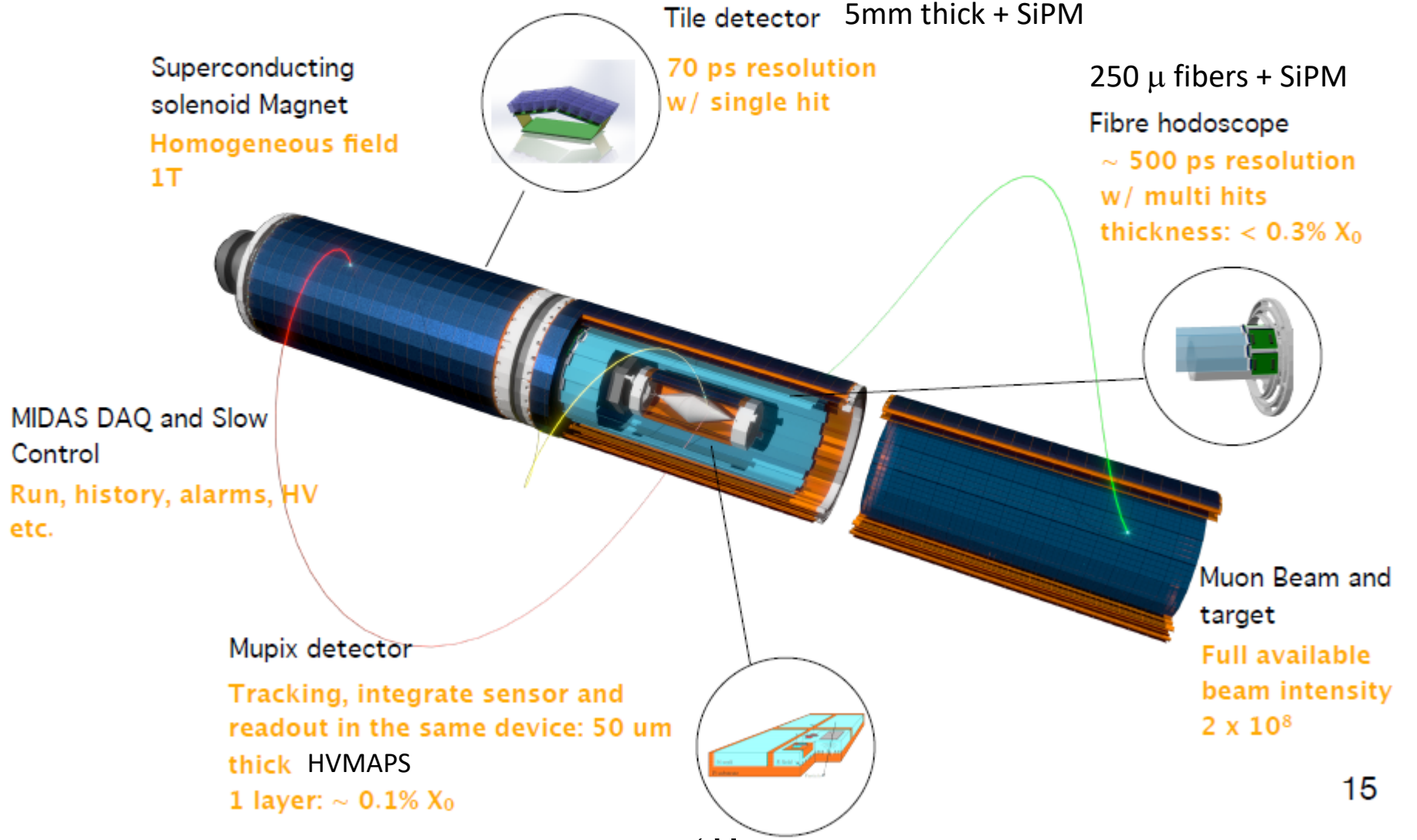
Data from SINDRUM I



2013-5	2015-7	2017	2018-20
Design	Construction	Eng Run	Run

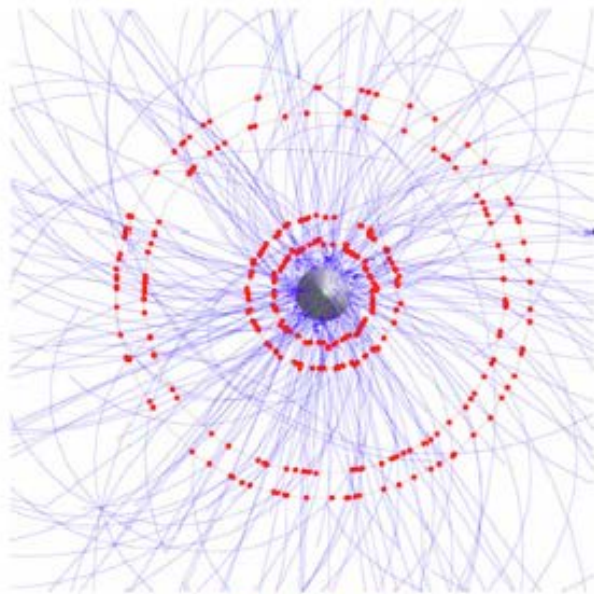
The detector concept

Sensitivity phase I >2020 $\sim 10^{-15}$
 (Final sensitivity phase II [202x] $\sim 10^{-16}$
 needs 10^9 muons/s)

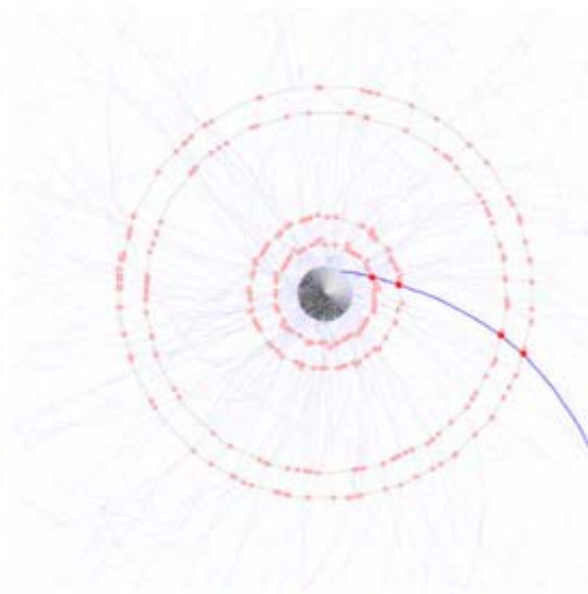


Precise timing is critical to reduce the accidental background

- Scintillating fibers (SciFi) $O(1 \text{ ns})$, full detection efficiency ($>99\%$)
- Scintillating tiles $O(100 \text{ ps})$, full detection efficiency ($>99\%$)



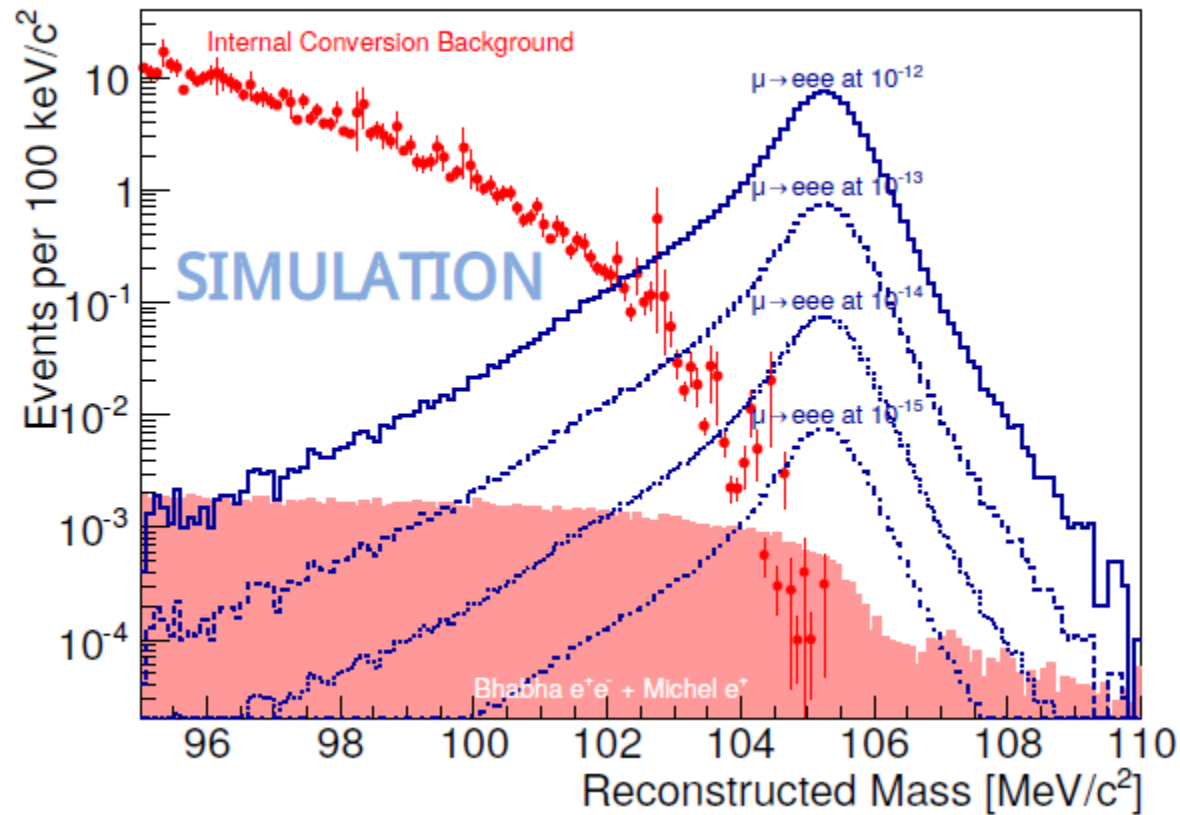
Pixels: $O(50 \text{ ns})$



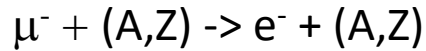
Scintillating fibres $O(1 \text{ ns})$;
Scintillating tiles $O(100 \text{ ps})$

Reconstruction simulation

Mu3e: $1 \cdot 10^{15}$ μ on Target; Rate 10^8 μ/s



$\mu \rightarrow e$ conversion



Signature: single monoenergetic e^-

$$E = m_\mu - B_\mu : 105 \text{ MeV in Al}$$

Lifetime $\approx 1 \mu\text{s}$

No accidental backgrounds!

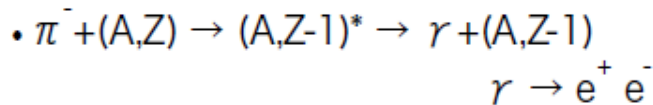
- Physics backgrounds

- Muon Decay in Orbit (DIO)

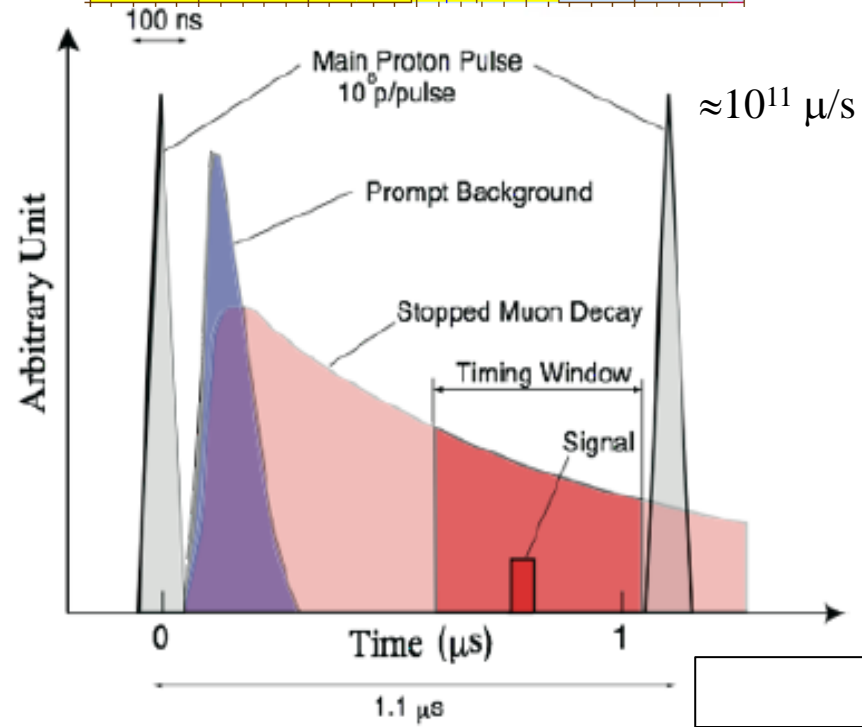
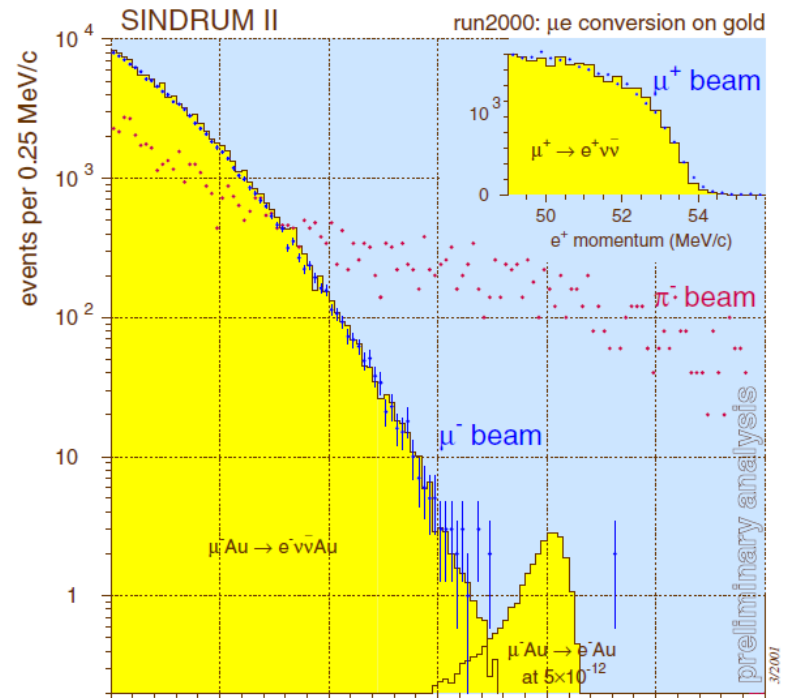
- $E_e > 102.5 \text{ MeV}$ (BR: 10^{-14})

- $E_e > 103.5 \text{ MeV}$ (BR: 10^{-16})

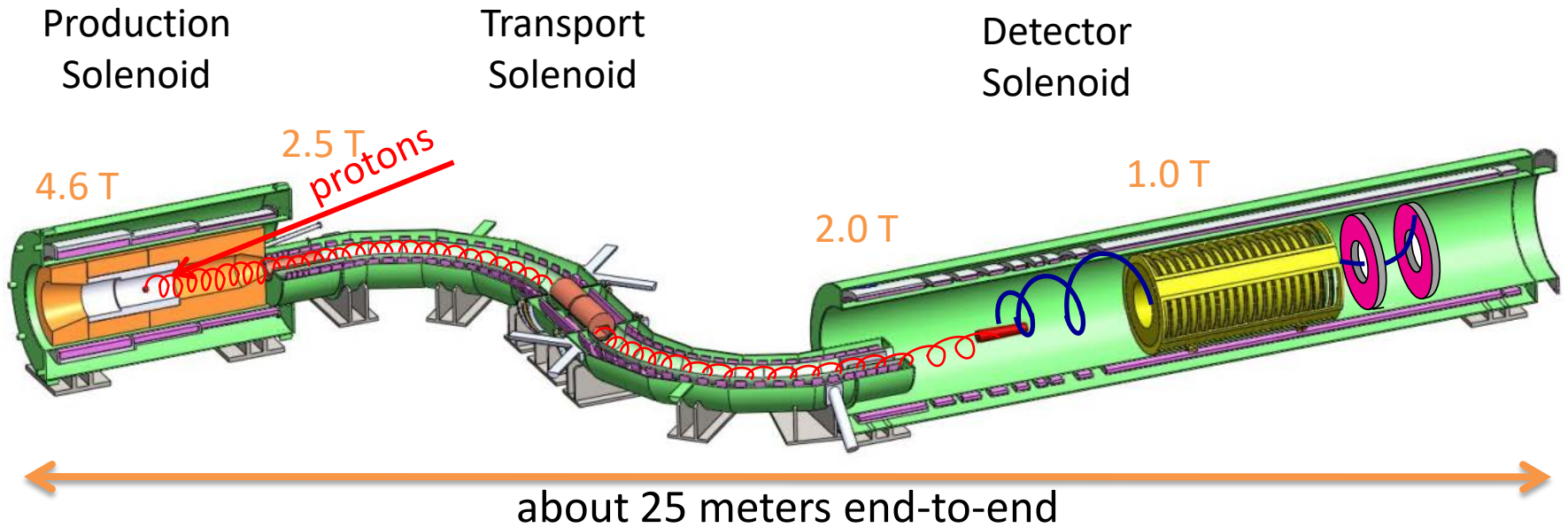
- Beam Pion Capture



$$R_{\text{ext}} = \frac{\text{number of proton between pulses}}{\text{number of proton in a pulse}}$$



Mu2e beam line elements



Production Solenoid:

8 GeV protons interact with a tungsten target to produce μ^- (from π^- decay)

Transport Solenoid:

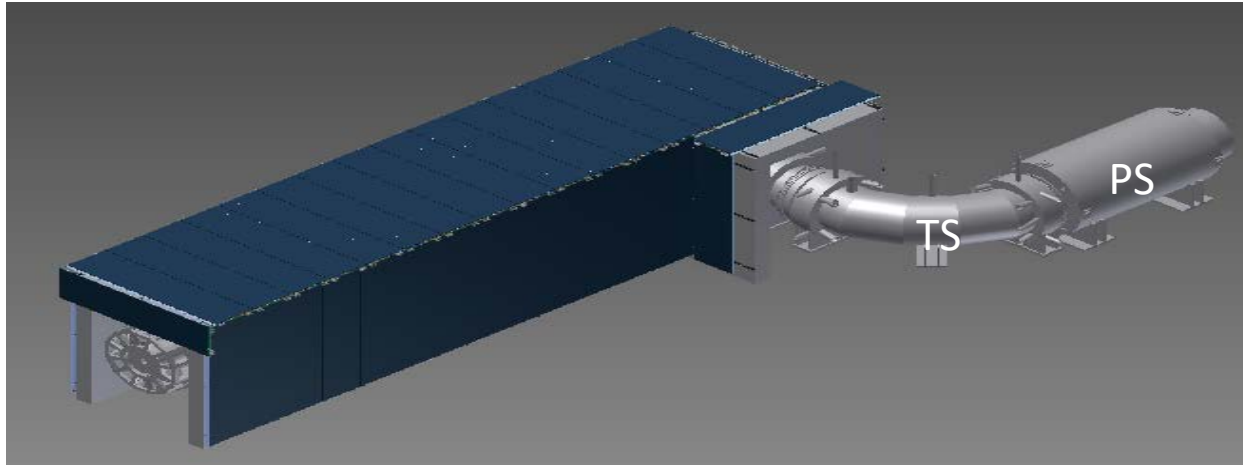
Captures π^- and subsequent μ^- ; momentum- and sign-selects beam

Detector Solenoid:

Upstream – Al. stopping target, Downstream – tracker, calorimeter

(not shown – cosmic ray veto system, extinction monitor, target monitor)

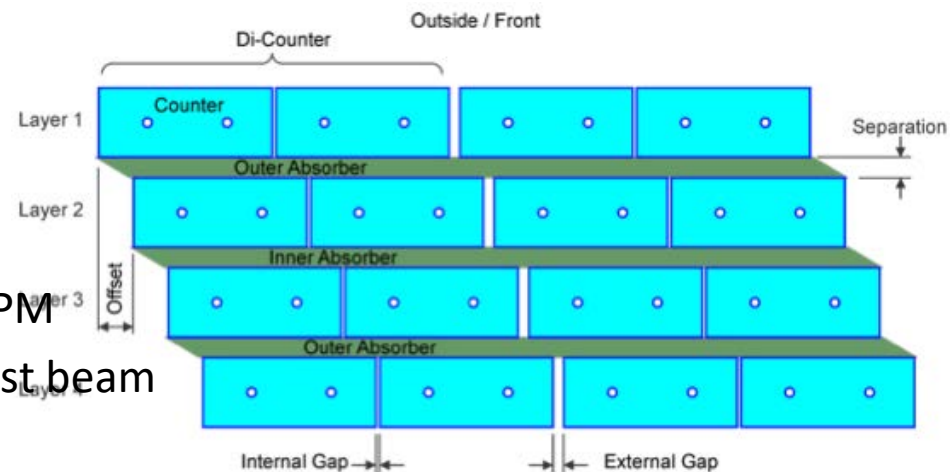
Mu2e Cosmic-Ray Veto



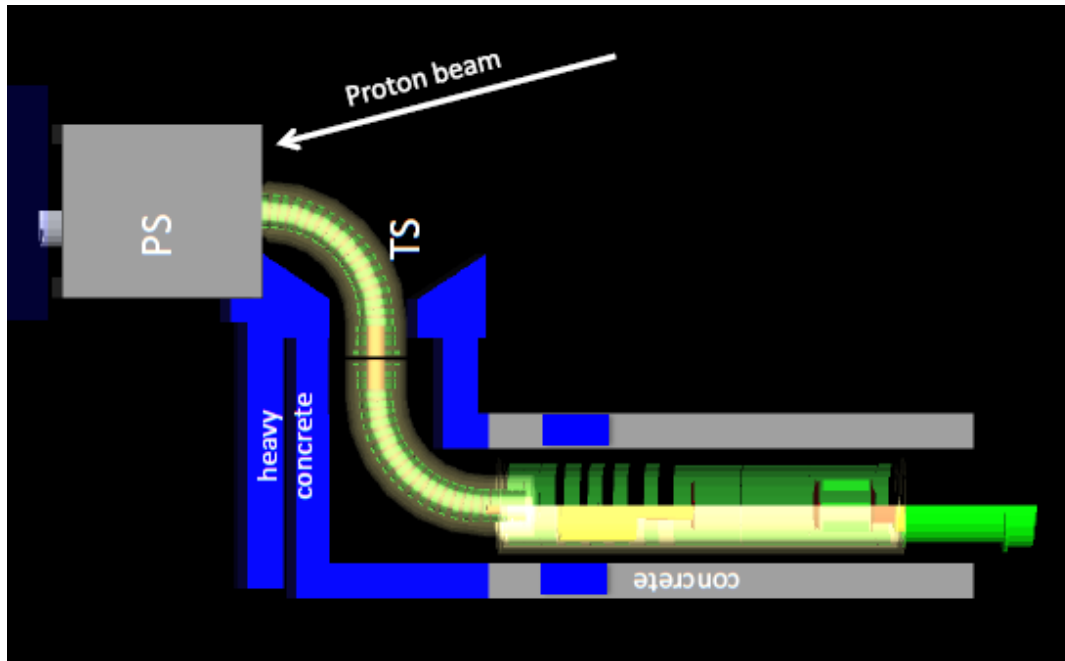
Without the veto system, ~ 1 cosmic-ray induced background event per day

- 4 overlapping layers of scintillator

- Each bar is $5 \times 2 \times \sim 450 \text{ cm}^3$
- 2 WLS fibers / bar
- Read-out both ends of each fiber with SiPM
- Have achieved $e > 99.4\%$ (per layer) in test beam

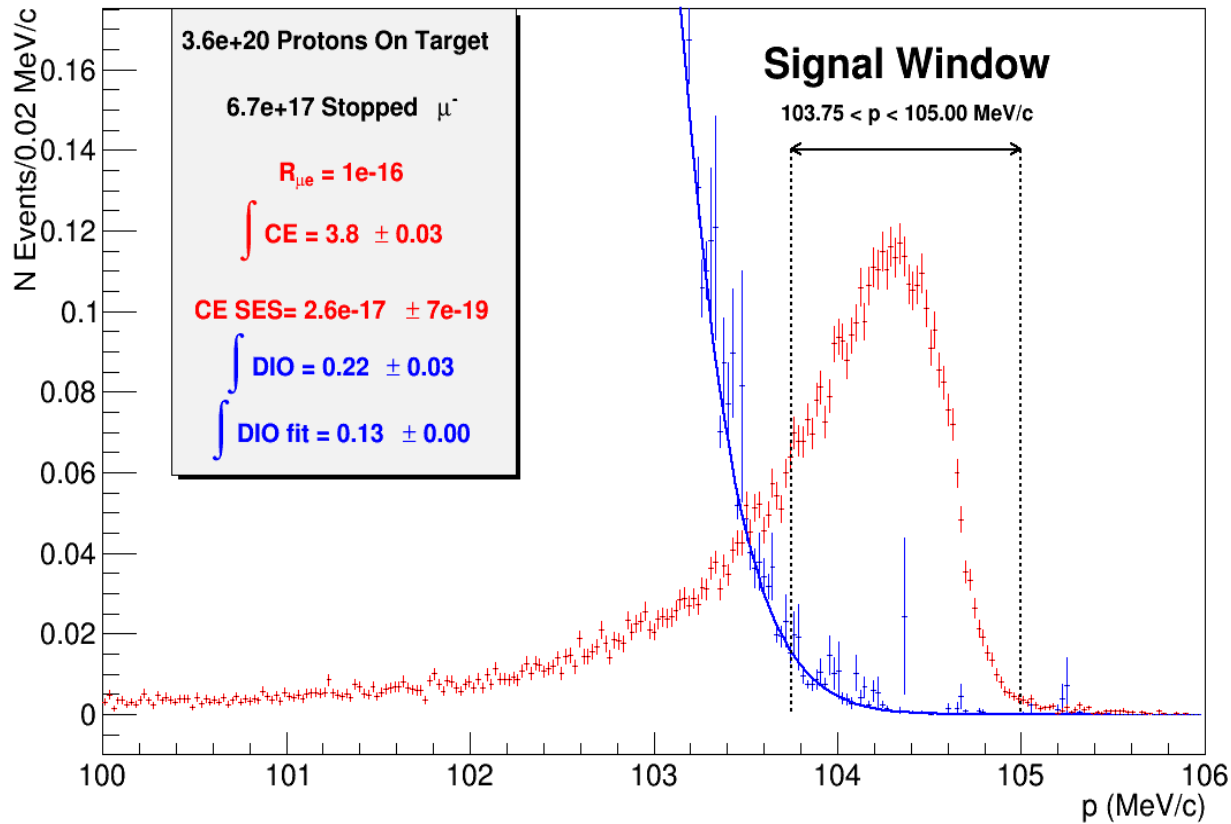


Plus a neutrons shielding of the detectors' area



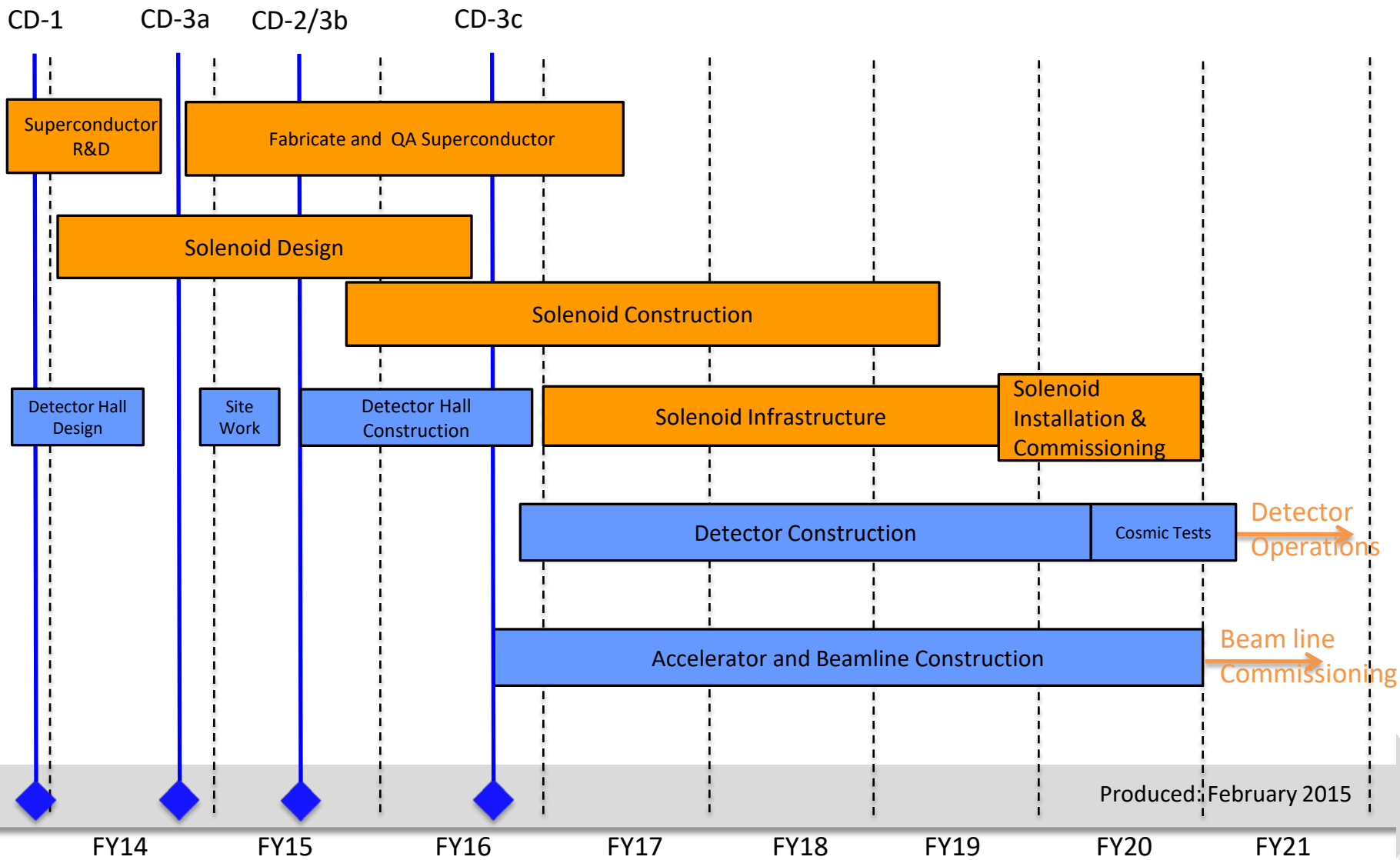
...and joint effort Mu2e and Comet to study the proton emission due to muon capture in Al which may significantly contribute to the detectors hit rate: **ALCap experiment at PSI**

Detailed simulation of the electron momentum at the signal region

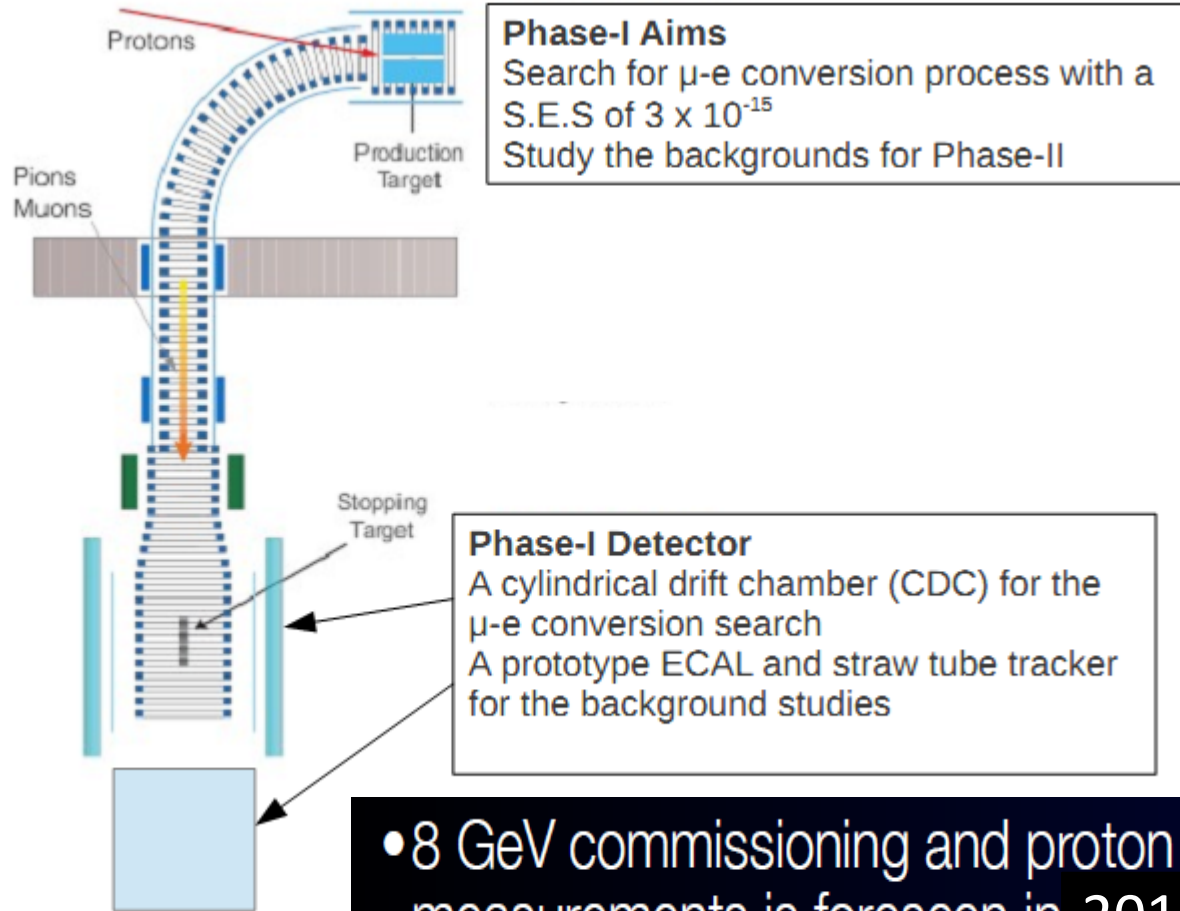


- Single-event-sensitivity = 2.6×10^{-17}
(SES goal 2.4×10^{-17})
- Total background < 0.5 events

Mu2e Schedule

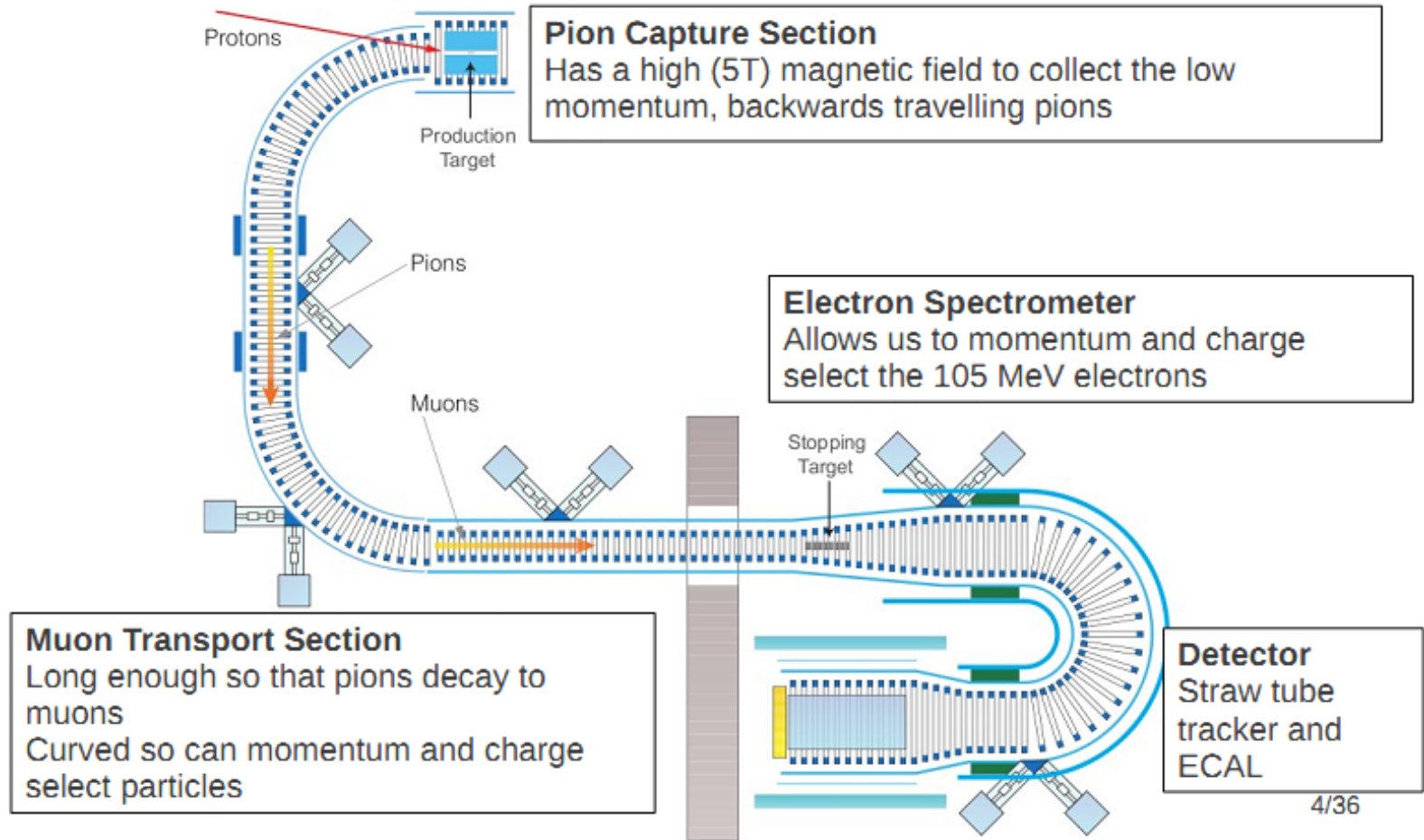


COMET Phase I



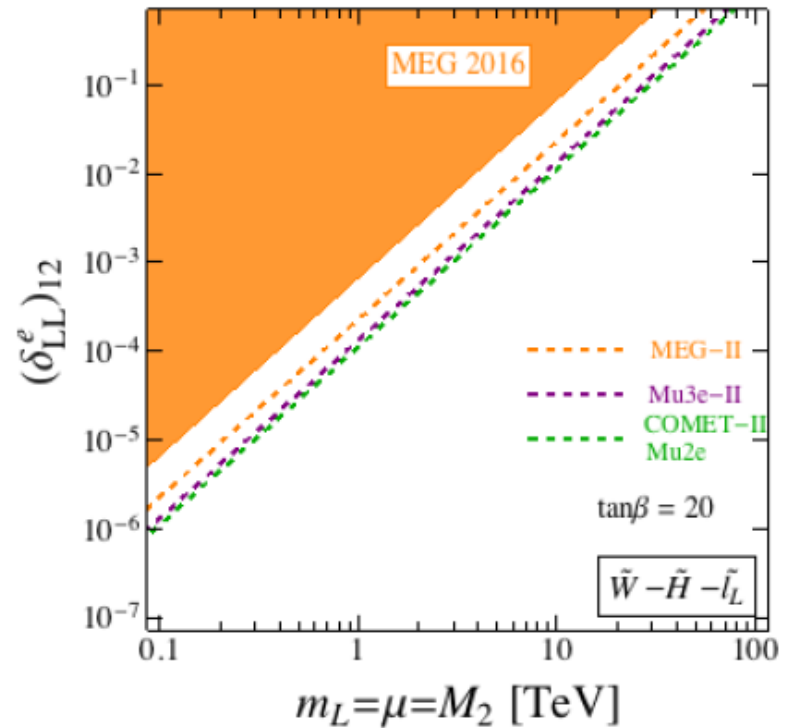
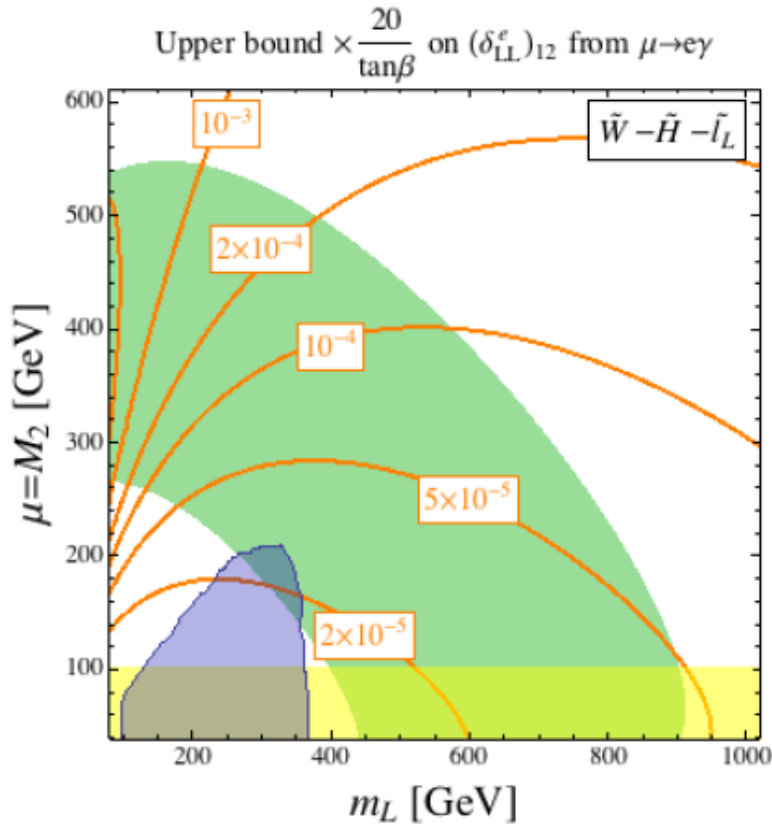
COMET Phase II

Starting in 2020
Measurement in 2022
S.E.S = 3×10^{-17}



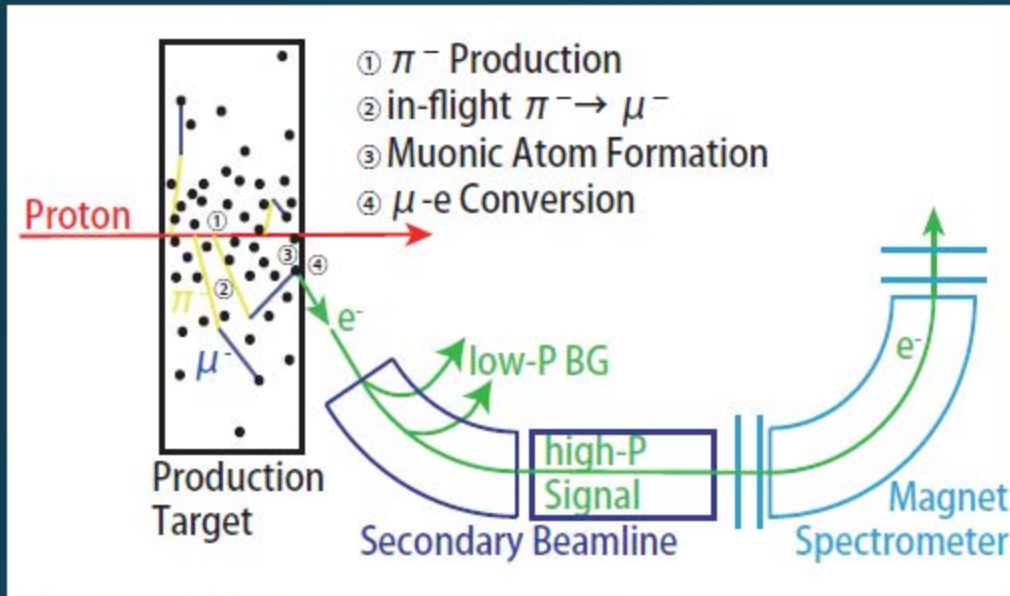
Comparison with SUSY searches at LHC

$$(\delta_{LL})_{ij} = \frac{(\Delta_{LL})_{ij}}{\sqrt{(\tilde{m}_L^2)_{ii}(\tilde{m}_L^2)_{jj}}}$$



Calibbi, Signorelli, NC 2017

DeeMe experiment



$\mu N \rightarrow e N$
signal electron

- single
- mono energetic
- delayed

The signal electron is identified by measuring their momentum

Start with Carbon target

- Lifetime of muonic atom $\sim 2 \mu\text{s}$
- Energy of electron from μ -e conversion = 105 MeV
- Single event sensitivity (1 year = 2×10^7 sec)
 - 1×10^{-13}
 - 2.5×10^{-14} (4 years)



- 2×10^{-14}
- 5×10^{-15} (4 years)

Summary

CLFV: excellent probe for search for physics Beyond the Standard Model

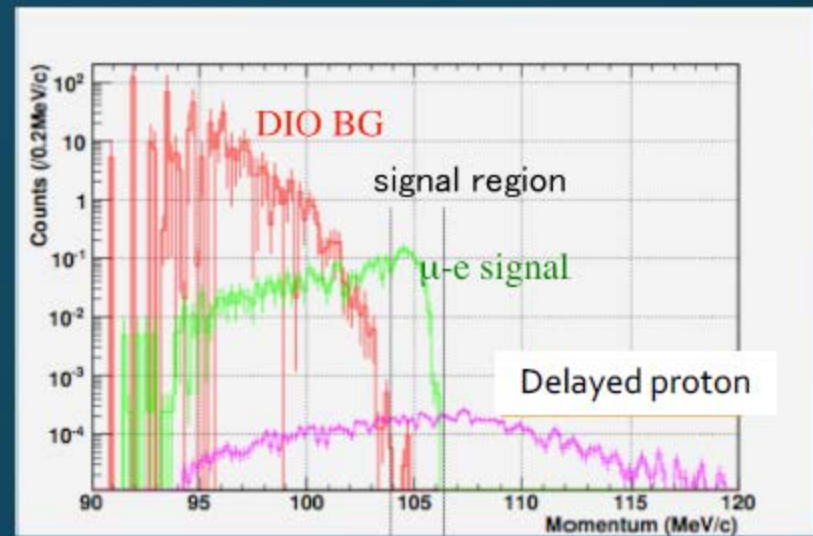
Complementarity among the different channels enabling to check the different New Physics models

Complementarity of intensity (and precision) frontier wrt to the high energy frontier

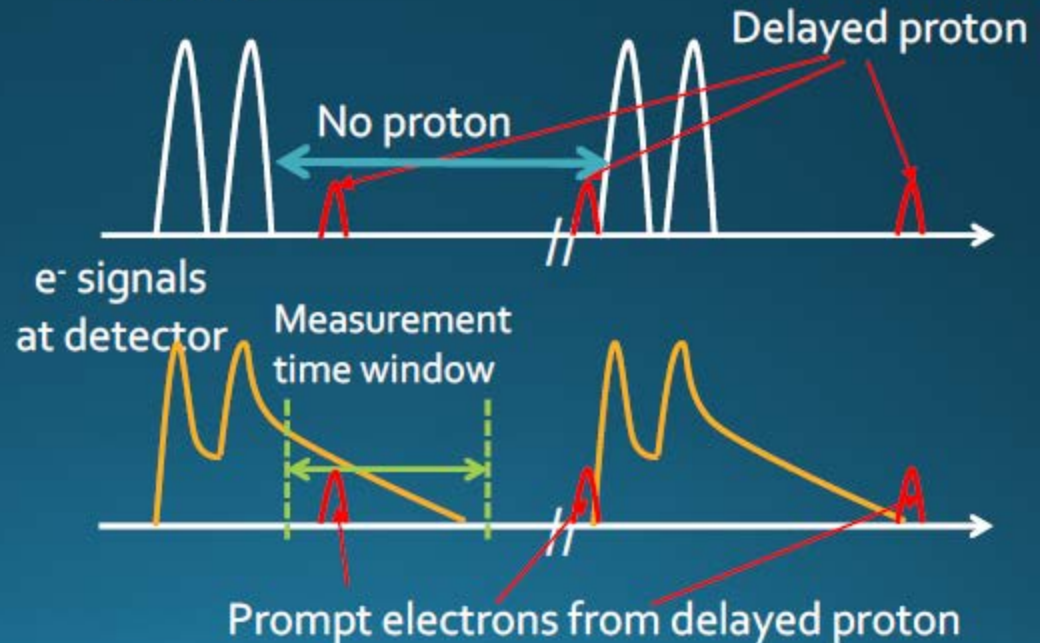
We would need some luck, soon or later...

Background

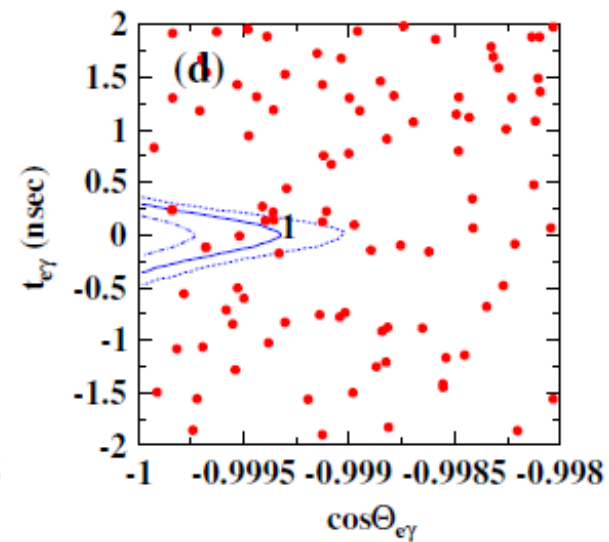
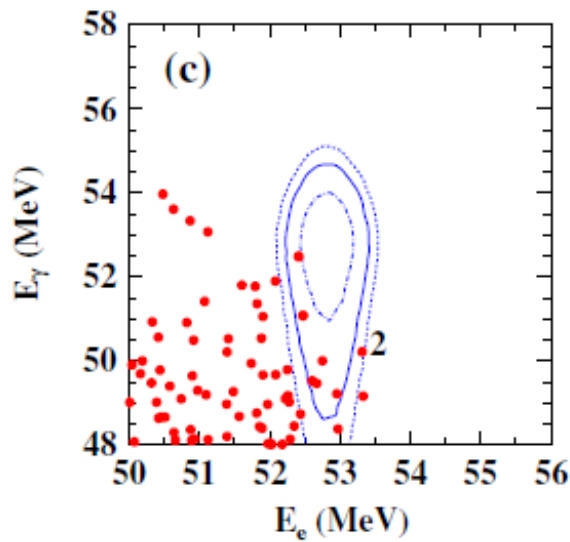
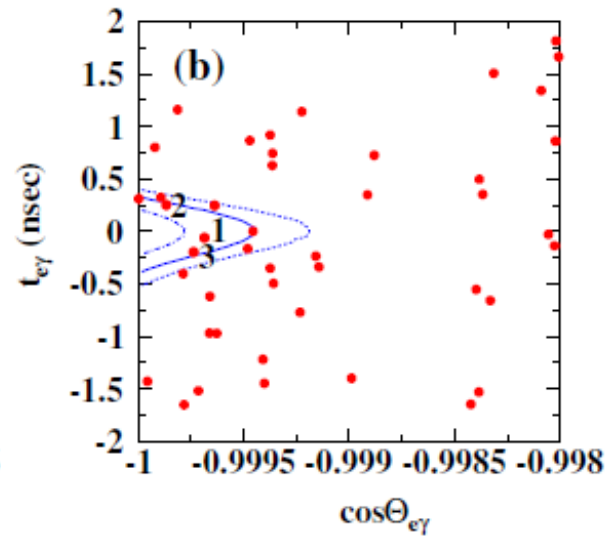
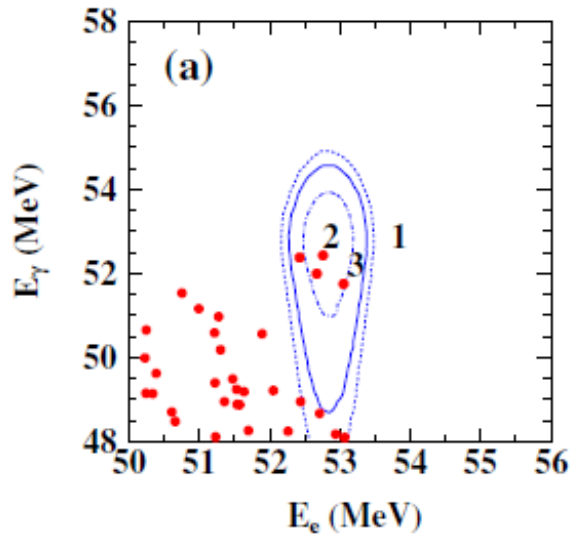
- Decay In Orbit 0.015 in the signal region
- Delayed proton < 0.027 (zero in principle)
Synchrotron is empty at the delayed timing by fast extraction
- Cosmic-ray induced
 $e: < 0.018, \mu: < 0.001$
duty factor = $1/20000$
Horizontal tracking direction
- Anti-Proton zero
beam energy (=3 GeV)
 $< \bar{p}$ production threshold



Primary Proton
from RCS



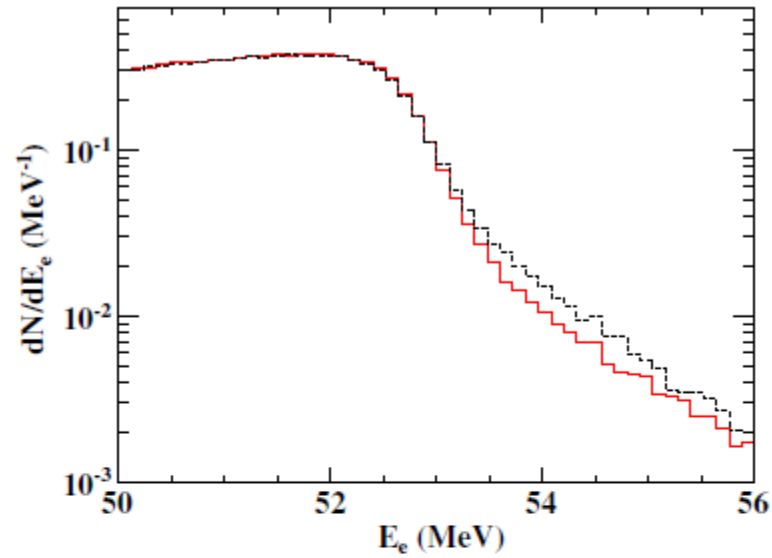
Aiming to start DeeMe experiment from 2017!



**2011
paper**

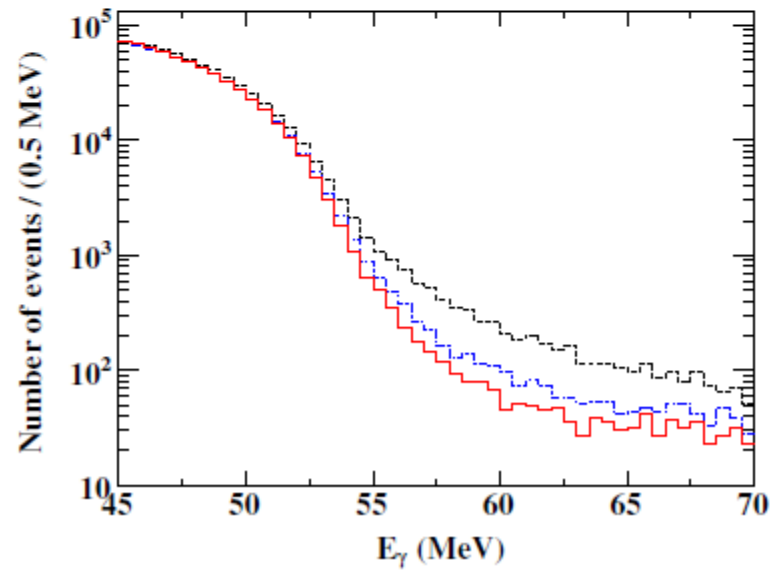
Data set	\mathcal{B}_{fit}	LL	UL
2009	3.2×10^{-12}	1.7×10^{-13}	9.6×10^{-12}
2010	-9.9×10^{-13}	...	1.7×10^{-12}
2009 and 2010	-1.5×10^{-13}	...	2.4×10^{-12}

New Kalman filter



2013

Better pile-up rejection

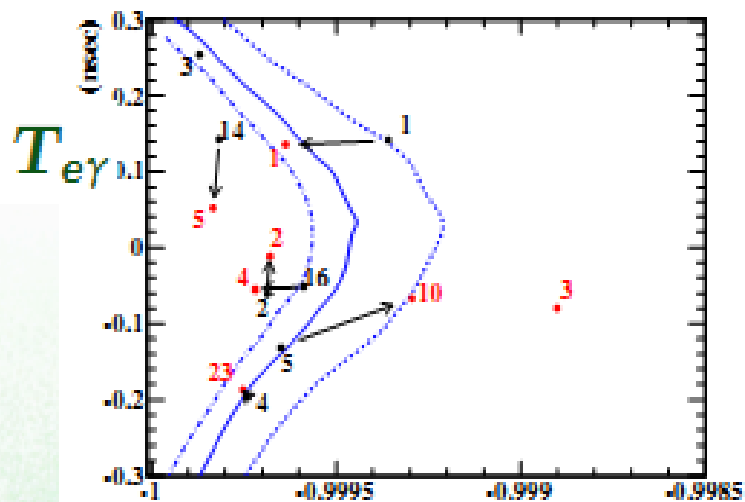
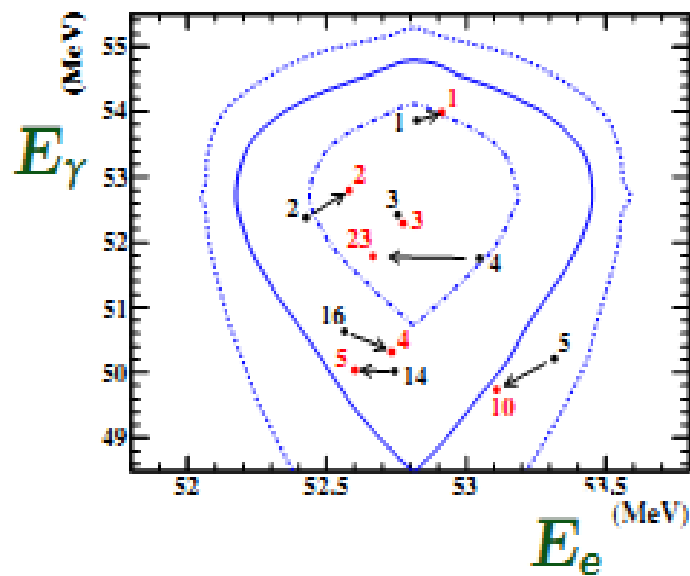


Comparison with previous analysis

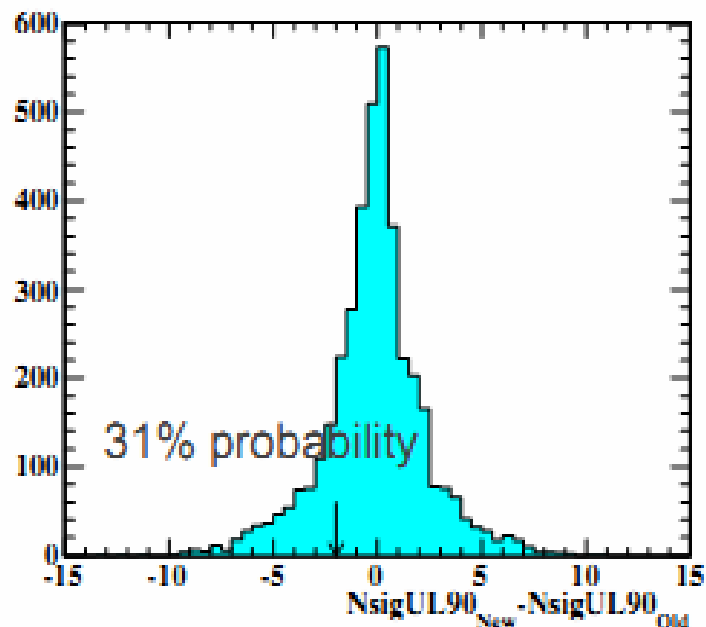


- Previous analysis
- New analysis

Numbers : rank in each dataset

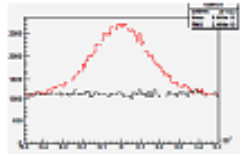
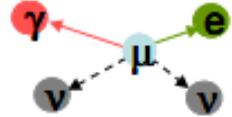


Change of UL by modifications of reconstruction algorithms. (MC)



- High ranked events are stable
- Differences of observables by modifications of reconstruction algorithms are smaller than resolutions.

μ radiative decay

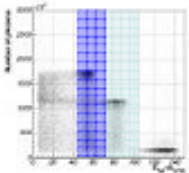


Lower beam intensity $< 10^7$ is necessary to reduce pile-ups

Better σ_{T1} makes it possible to take data with higher beam intensity

A few days ~ 1 week to get enough statistics

$\pi^0 \rightarrow \gamma\gamma$



$\pi + p \rightarrow \pi^0 + n$
 $\pi^0 \rightarrow \gamma\gamma$ (55MeV, 83MeV)
 $\pi + p \rightarrow \gamma + n$ (129MeV)

10 days to scan all volume precisely

(faster scan possible with less points)

LH₂ target



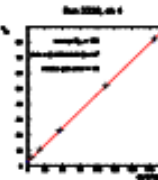
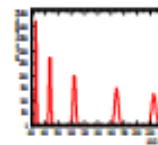
Laser

(rough) relative timing calib.

$< 2-3$ nsec



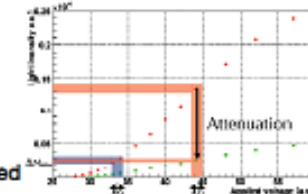
LED



PMT Gain

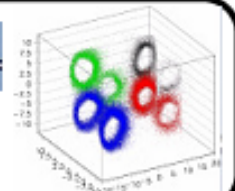
Higher V with light att.

Can be repeated frequently



Xenon Calibration

alpha

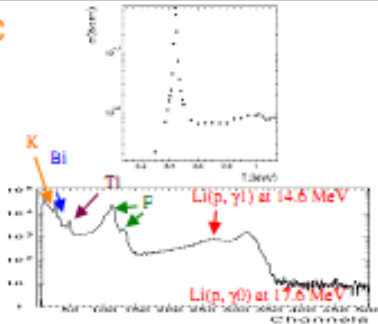


PMT QE & Att. L

Cold GXe

LXe

Proton Acc



Li(p,γ)Be

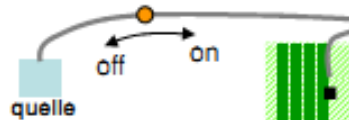
LiF target at COBRA center

17.6MeV γ

~daily calib.

Can be used also for initial setup

Nickel γ Generator



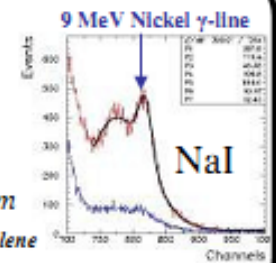
Illuminate Xe from the back

Source (Cf) transferred by comp air \rightarrow on/off

3 cm 20 cm

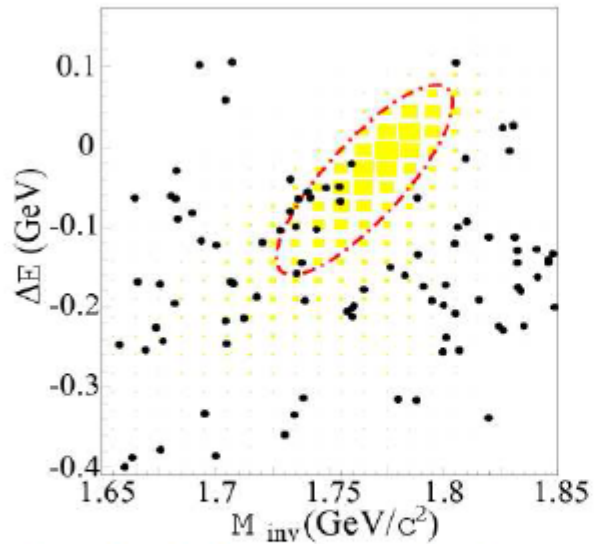
Polyethylene

0.25 cm Nickel plate



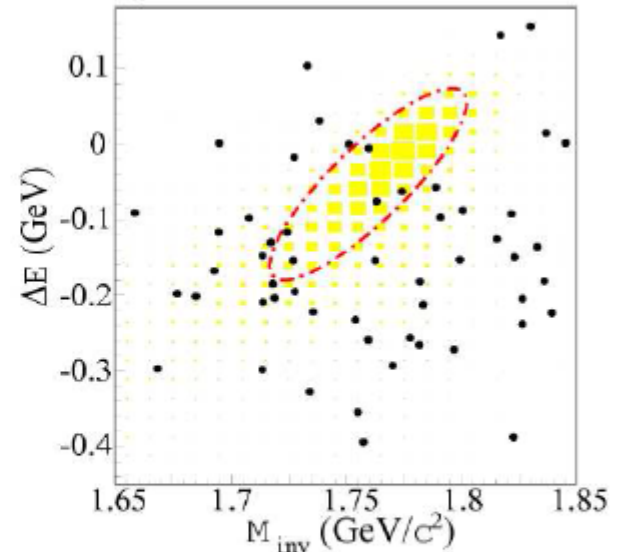
Belle

$\tau \rightarrow \mu \gamma$



– $\text{Br} < 4.5 \times 10^{-8}$ at 90% C.L.

$\tau \rightarrow e \gamma$



– $\text{Br} < 1.2 \times 10^{-7}$ at 90% C.L.

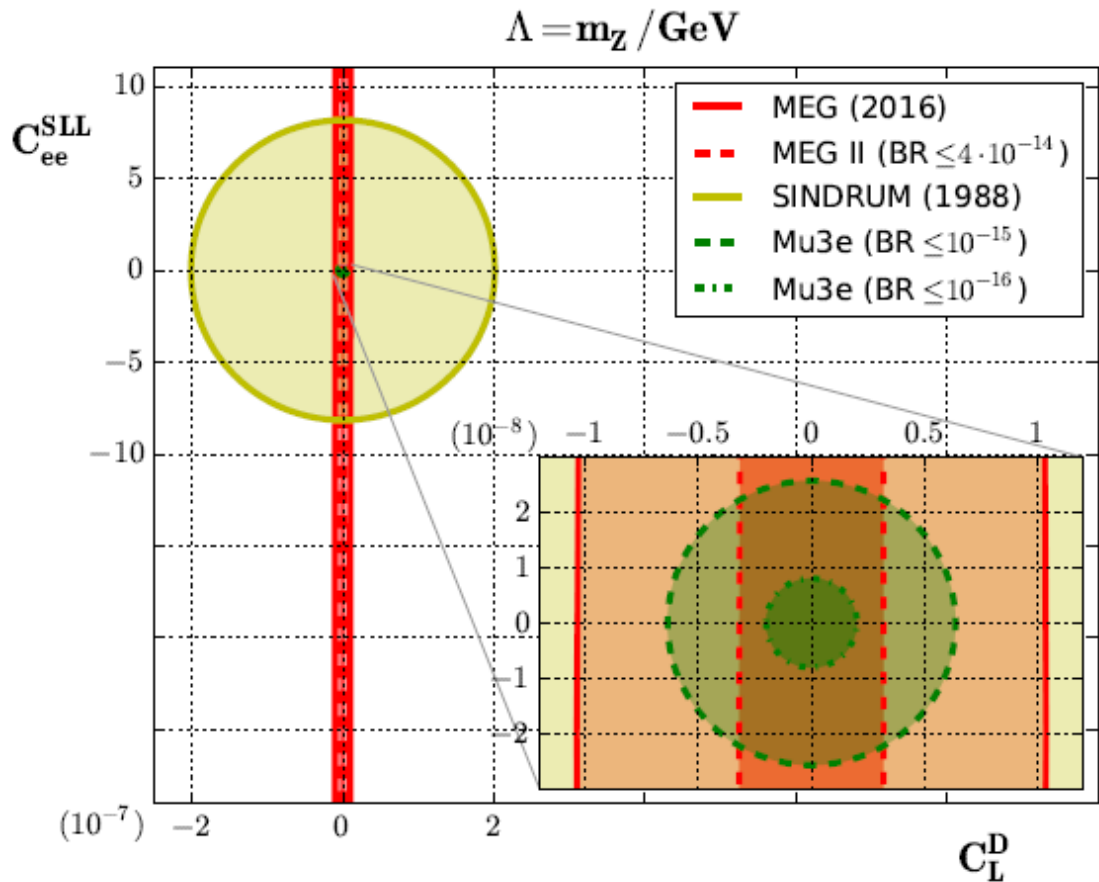
Efficiencies: 5.1% for $\mu^- \gamma$ and 3.0% for $e^- \gamma$

K. Hayasaka et al., Phys. Lett. B666 (2008) 16

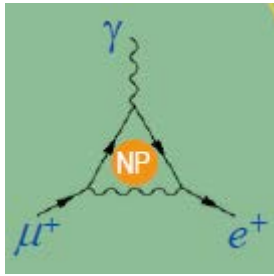
$$\mathcal{L}_{eff} = \mathcal{L}_{SM} + \sum_{d>4} \frac{c_n^{(d)}}{\Lambda^{d-4}} \mathcal{O}^{(d)}$$

$$O_L^D = e m_\mu (\bar{e} \sigma^{\mu\nu} P_L \mu) F_{\mu\nu}$$

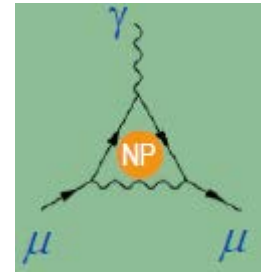
$$O_{ff}^{SLL} = (\bar{e} P_L \mu) (\bar{f} P_L f)$$



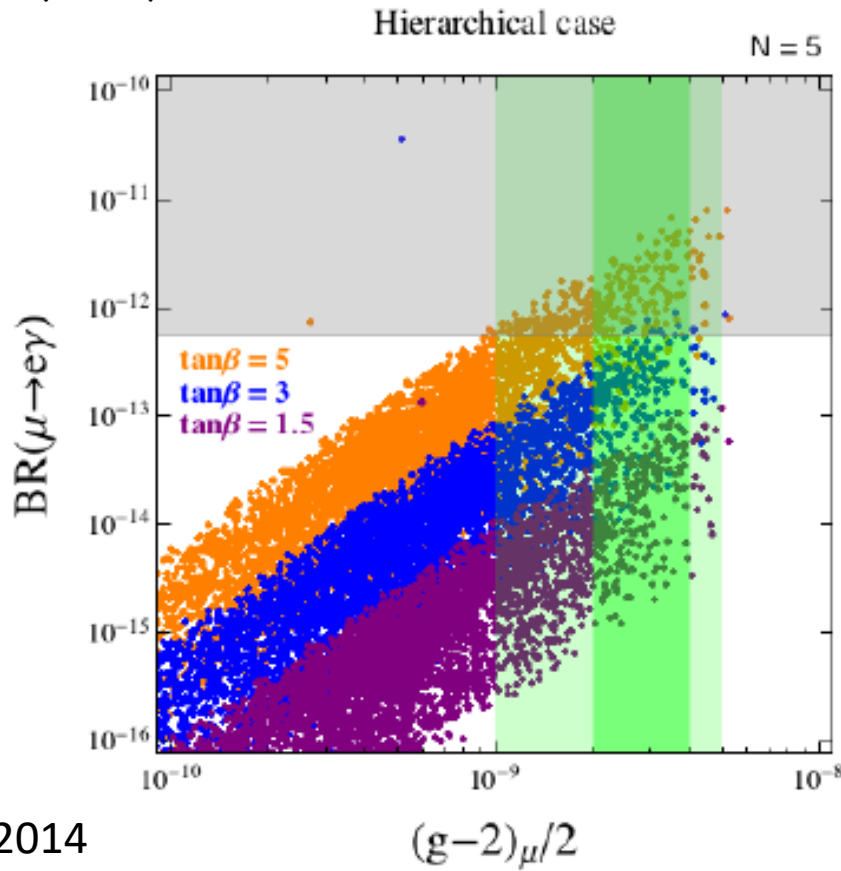
Crivellin et al., 2016



$\mu \rightarrow e \gamma$



$g-2$



Calibbi et al., 2014

The sensitivity is limited by the accidental background

$$n_{\text{sig}} \propto R_{\mu}, n_{\text{phys.b.}} \propto R_{\mu}, n_{\text{acc.b.}} \propto R_{\mu}^2$$

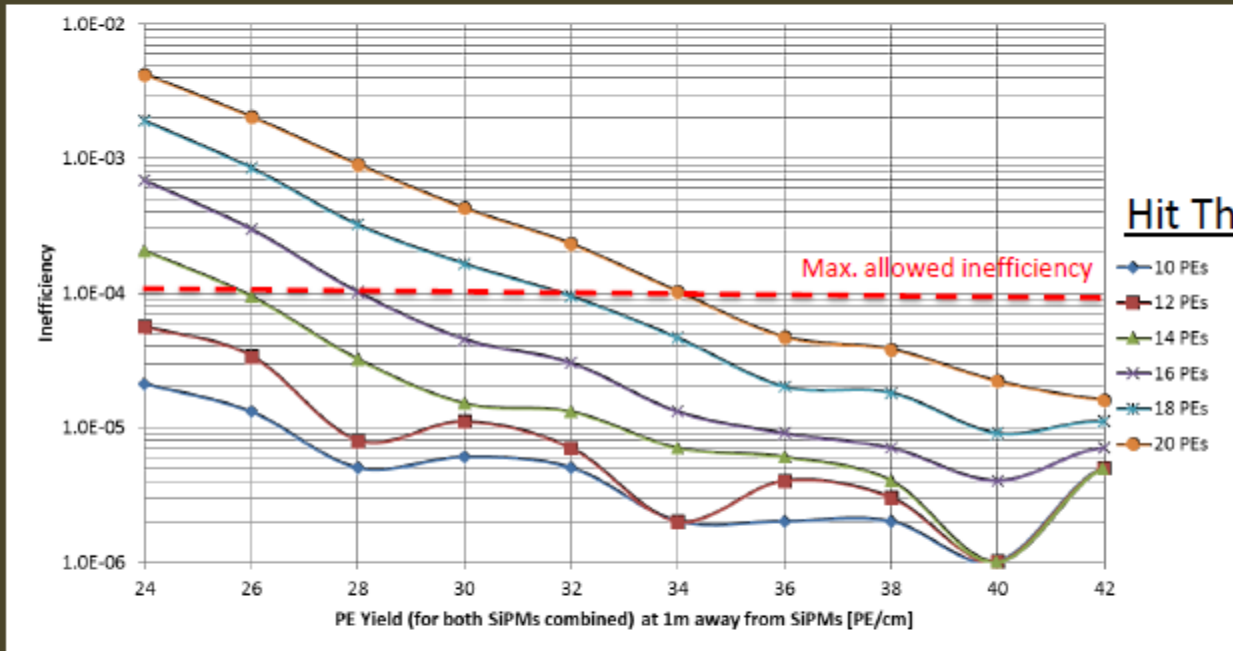
The n. of acc. backg events ($n_{\text{acc.b.}}$) depends quadratically on the muon rate and on how well we measure the experimental quantities: e- γ relative timing and angle, positron and photon energy

Effective BRback ($n_{\text{back}}/R_{\mu} T$)

$$BR_{\text{acc}} \propto R_{\mu} \times \Delta t_{e\gamma} \times \Delta \theta_{e\gamma}^2 \times \Delta E_e \times \Delta E_{\gamma}^2$$

Integral on the detector resolutions of the Michel and radiative decay spectra

Cosmic Ray Veto



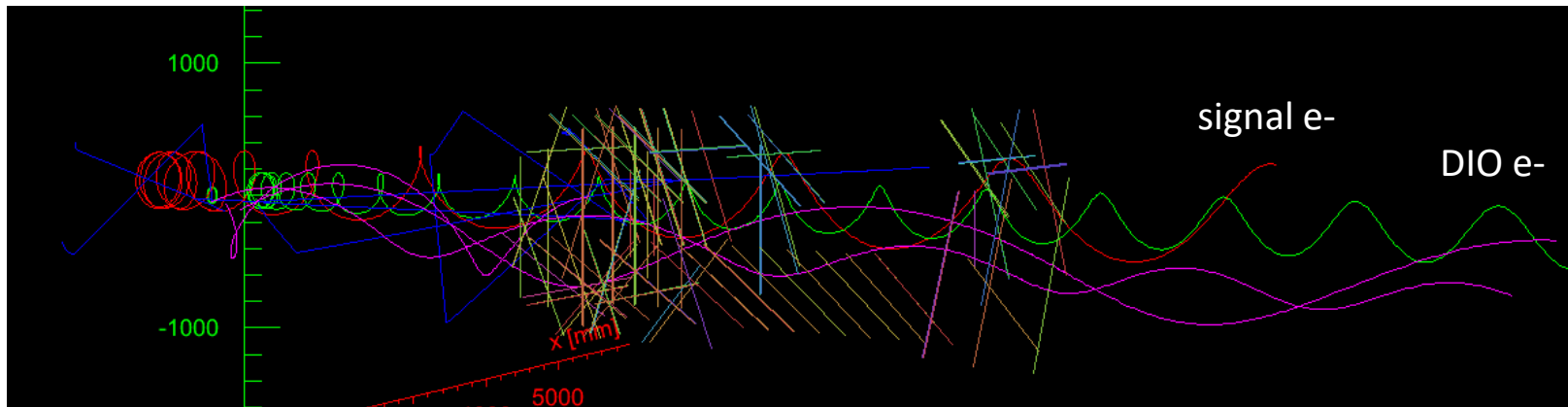
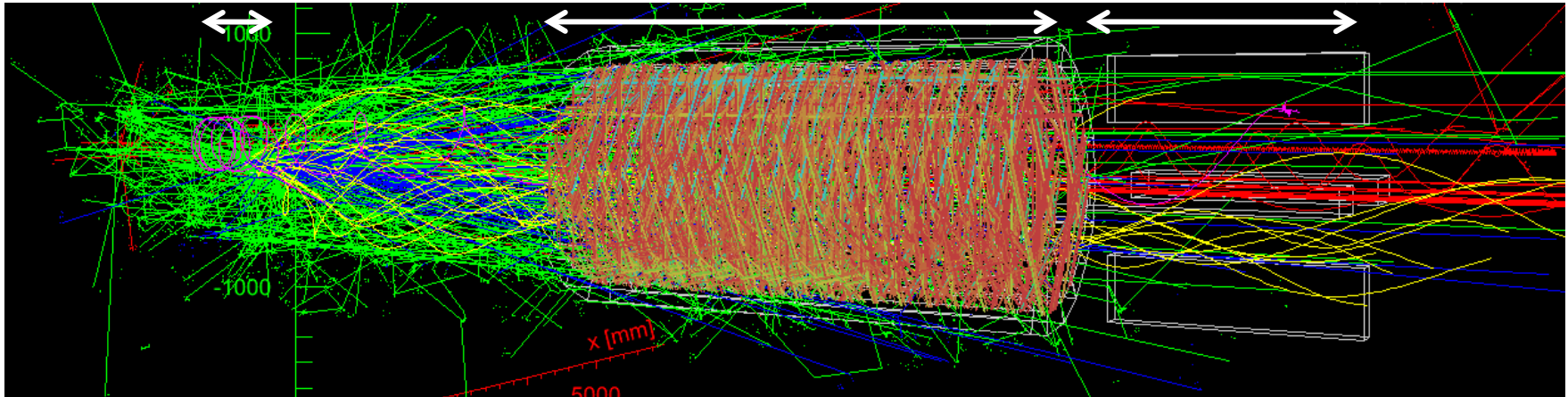
- Use detailed simulation and reconstruction to convince ourselves we can reasonably achieve the required veto efficiency

Mu2e Pattern Recognition

Stopping Target

Straw Tracker

Crystal Calorimeter

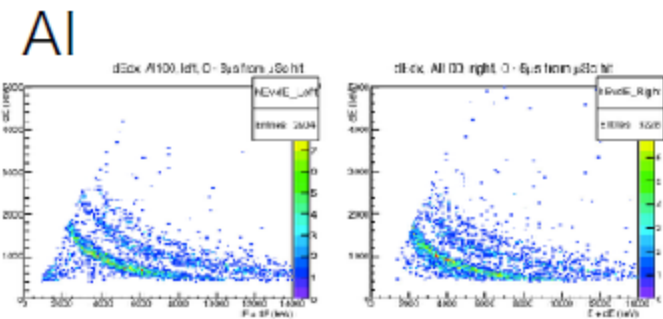
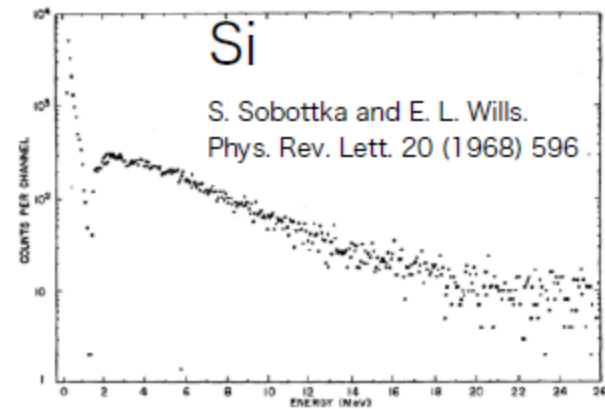


(particles with hits within ± 40 ns of signal electron t_{mean})



Background Assessment

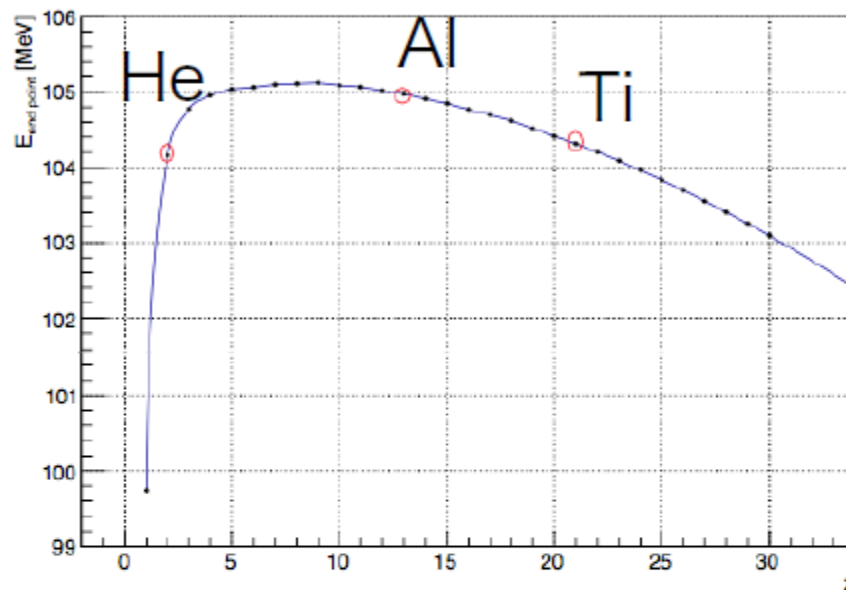
- Proton emission from muon captures
- No data available for Al
 - Si data only (near Al)
 - 15% proton emission / muon capture on ^{28}Si
 - $50\text{MeV}/c < p < 200\text{MeV}/c$
- Significant contribution to the detector hit rate
- AlCap experiment at PSI
 - Joint effort between COMET/Mu2e



· 0.05 protons/ muon capture
· Tran Hoai Nam (Osaka Univ.) PhD

Selection of the Target Material

- DIO E_{endpoing} extends to the $E_{\mu\text{-e}}$
 - Recoil energy
 - Muon binding energy
- Select the target material with high $E_{\mu\text{-e}}$ and avoid using the material with larger E_{endpoint} around the target
 - When the target is made of aluminium, we should avoid using materials from $Z=5$ to $Z=12$.
 - He ($Z=2$) is OK to use around the target
- Lifetime of muon in muonic atoms
 - Shorter in larger Z because of the larger nuclear muon capture rate



	Al	Ti
lifetime	864 ns	330 ns
time window	0.3	0.2
signal	1	1.5
net	0.3	0.3

