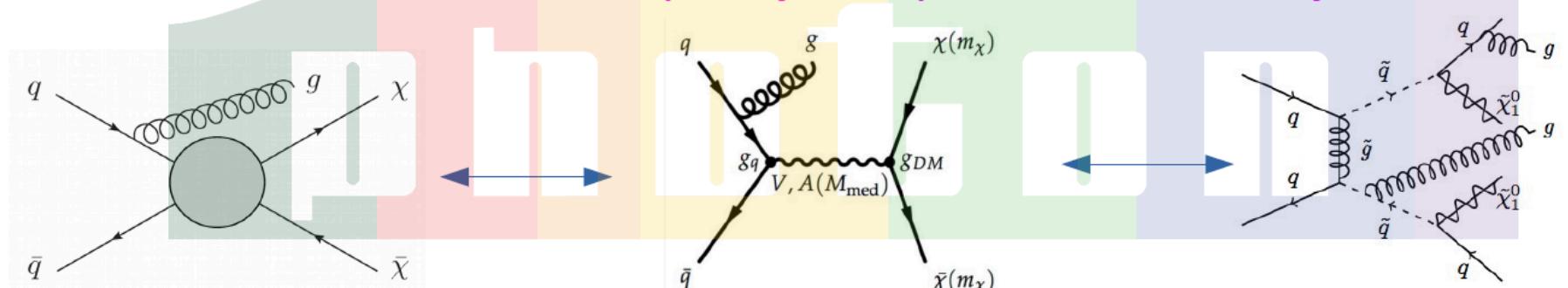


CHARACTERIZATION OF SEARCHES FOR DARK MATTER PRODUCTION AT COLLIDERS



Oliver Buchmueller, Imperial College London

XXVIII INTERNATIONAL SYMPOSIUM ON LEPTON PHOTON INTERACTIONS AT
HIGH ENERGIES AT YAT-SEN UNIVERSITY (SYSU), GUANGZHOU, CHINA



LEPTON PHOTON 2017

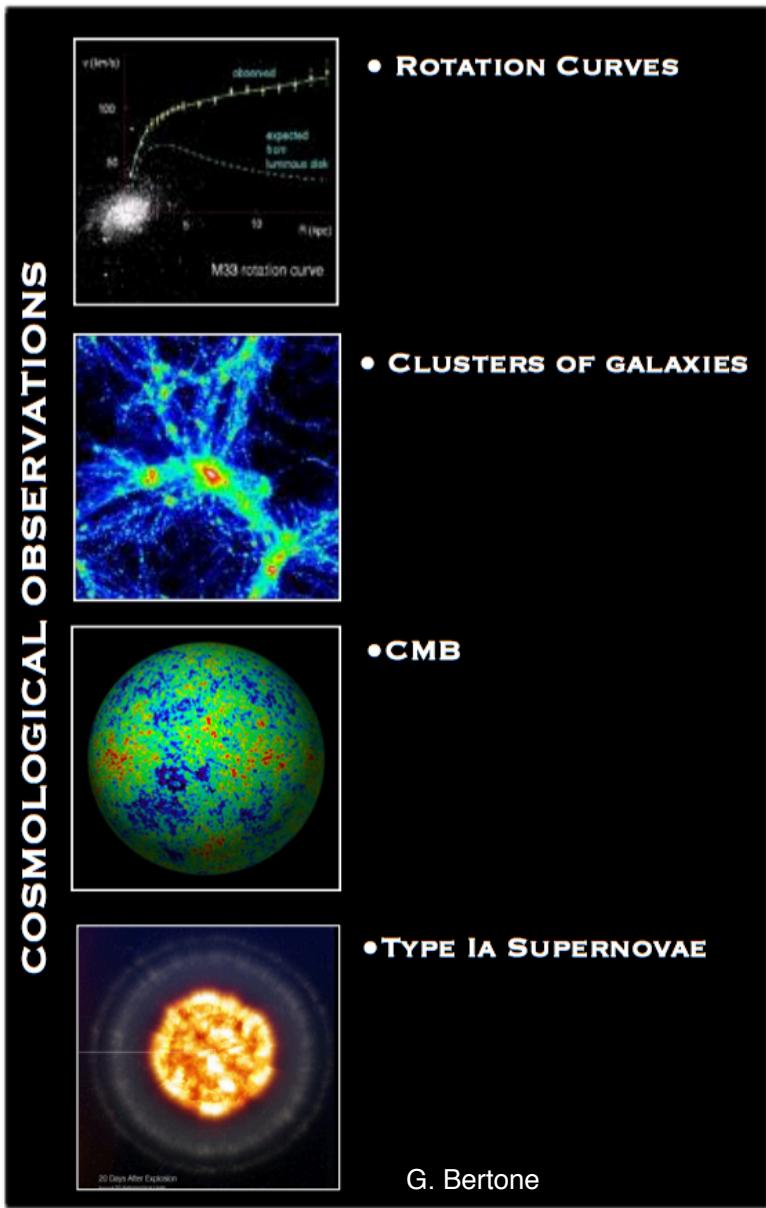
Preface

I will mainly revert to material/searches shown in:

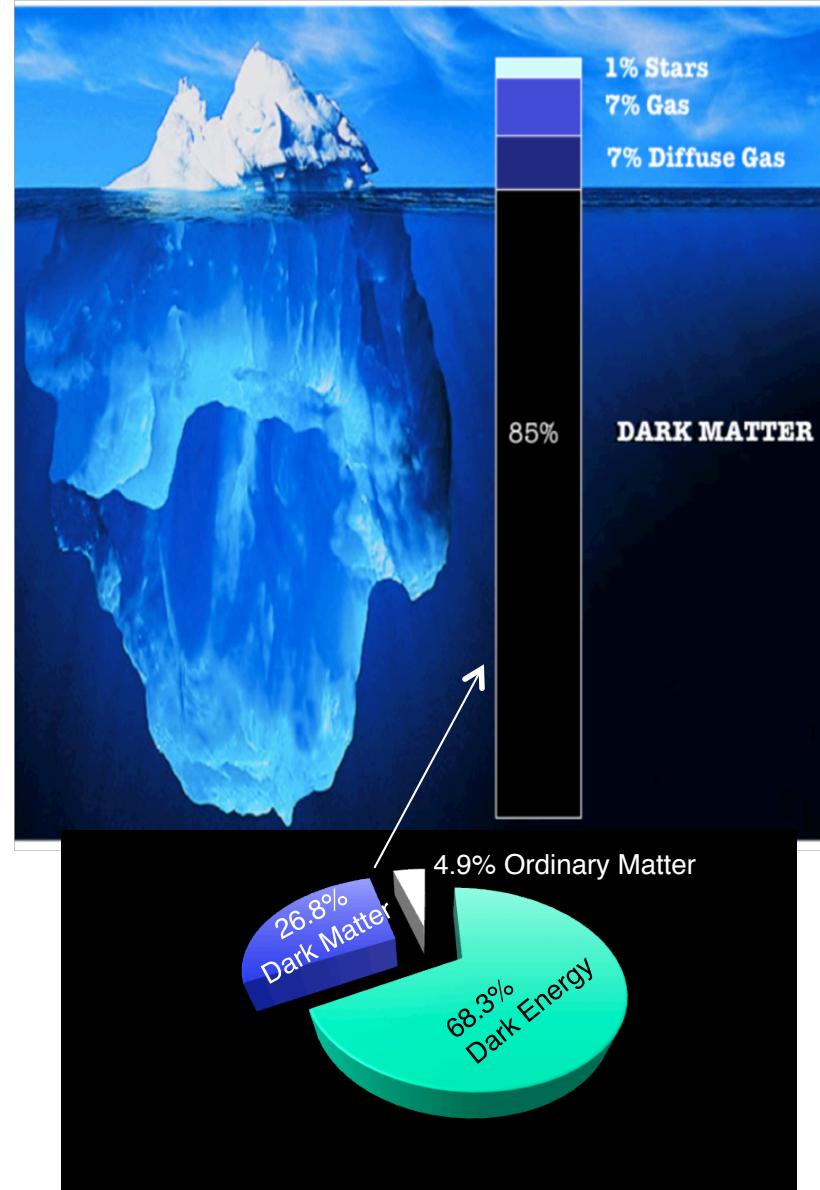
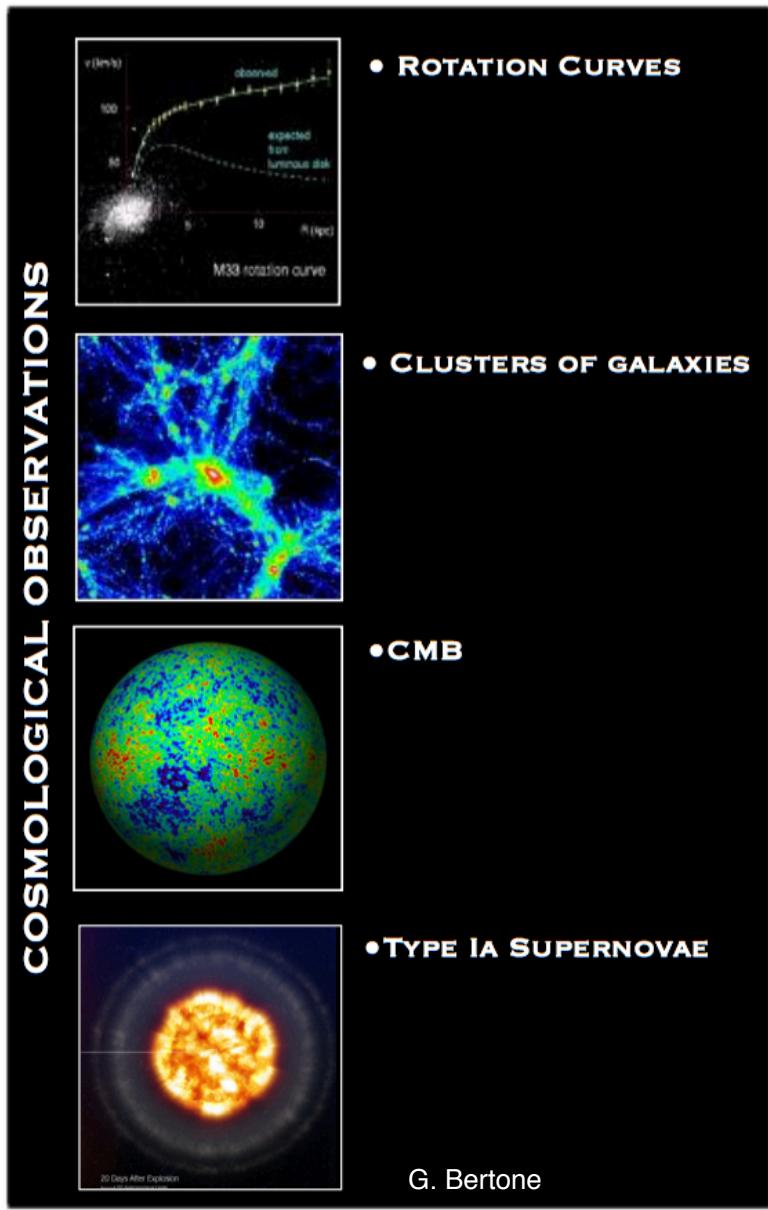
- **Searches for Beyond SM Higgs Bosons, Soshi TSUNO**
- **Searches for SUSY at LHC, Iacopo VIVARELLI**
- **Exotics searches at LHC, Sunil SOMALWAR**

and outline how these results can be (hands-on) interpreted in the context of Dark Matter and in turn be compared with other experiments like Direct Detection or Indirect Detection experiments.

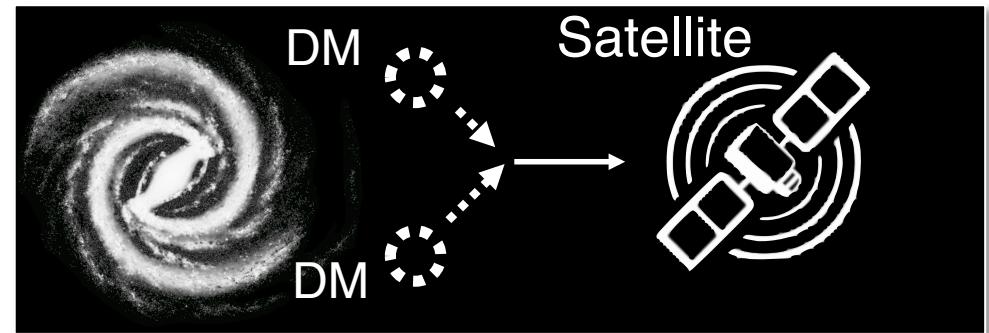
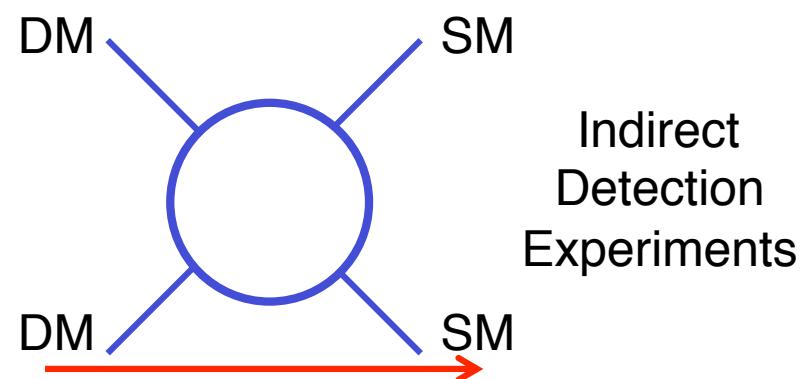
(Very Strong) Evidence for Dark Matter



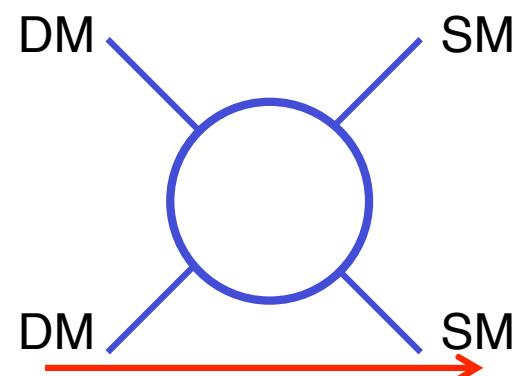
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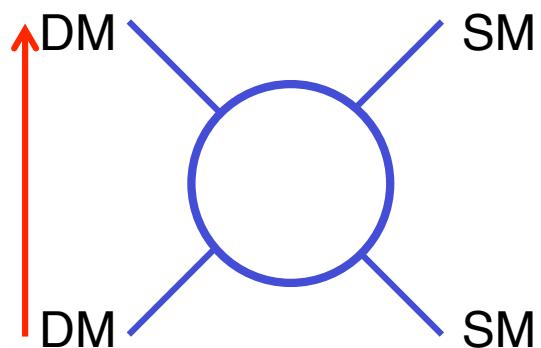
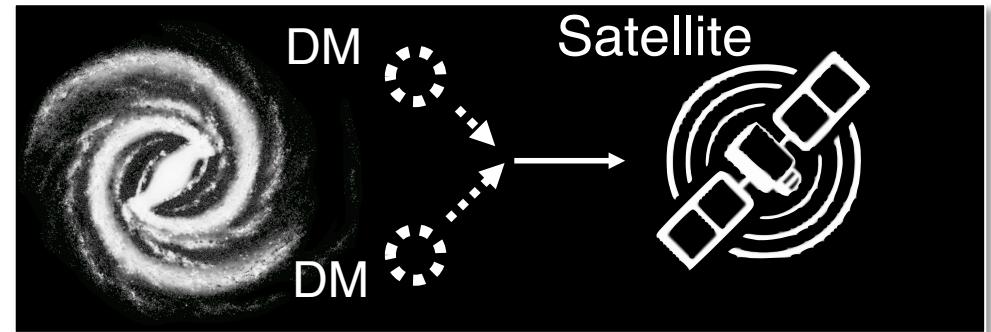
Hunting for Particle DM



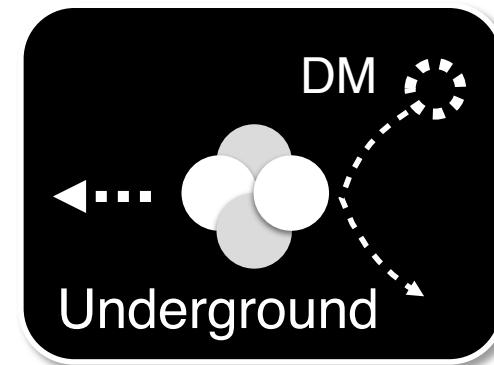
Hunting for Particle DM



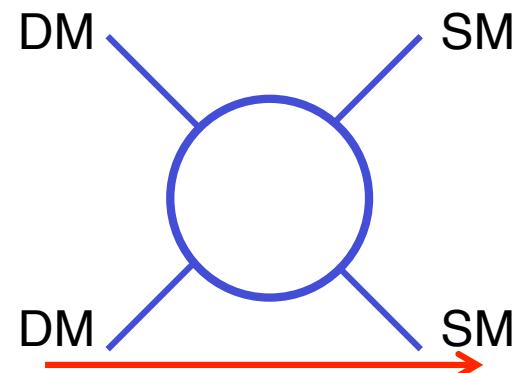
Indirect
Detection
Experiments



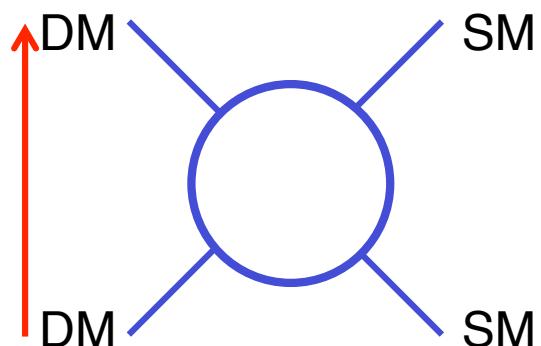
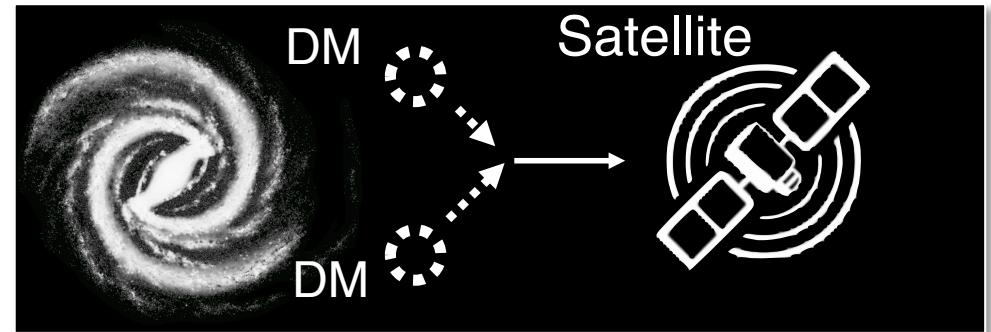
Direct
Detection
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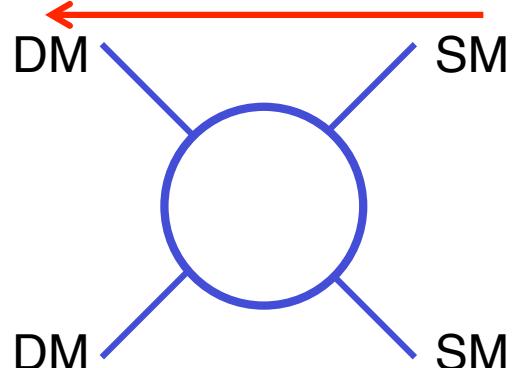
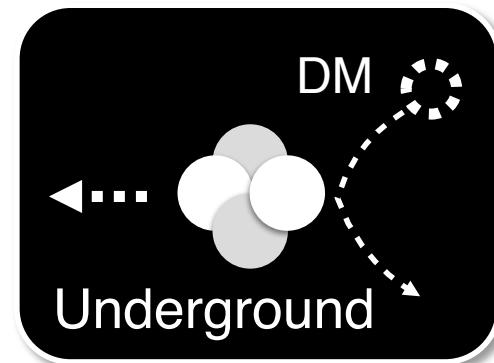
Hunting for Particle DM



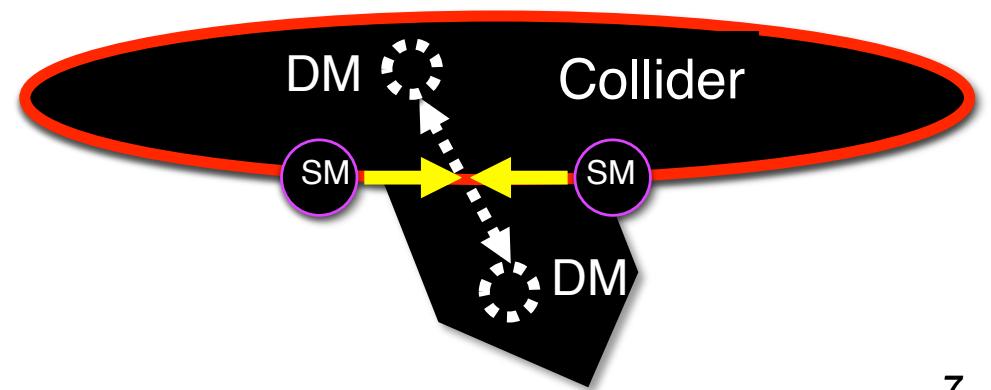
Indirect
Detection
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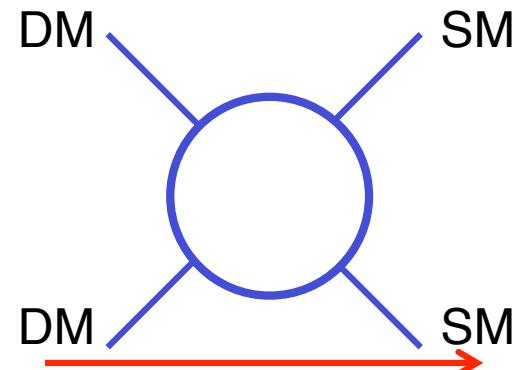
Direct
Detection
Experiments



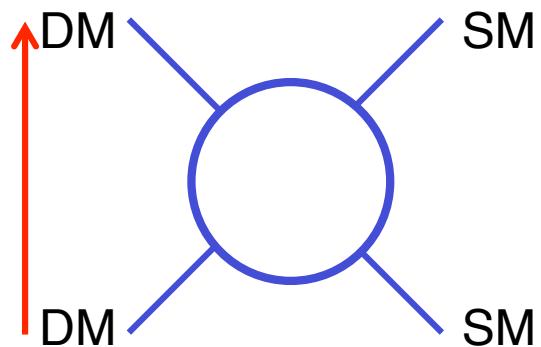
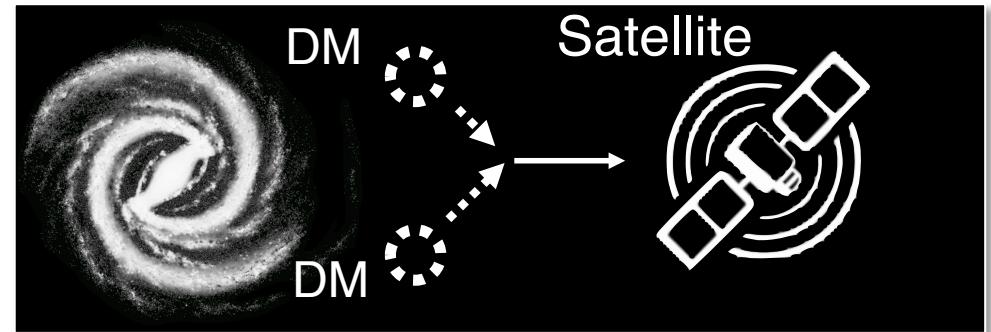
Collider
Experiments



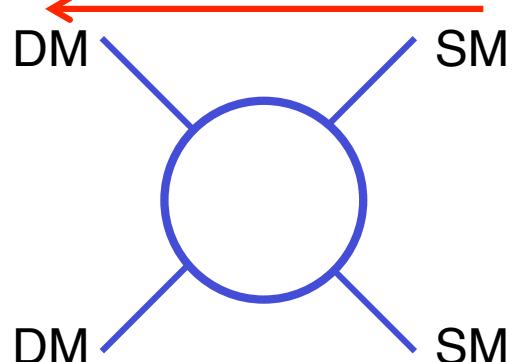
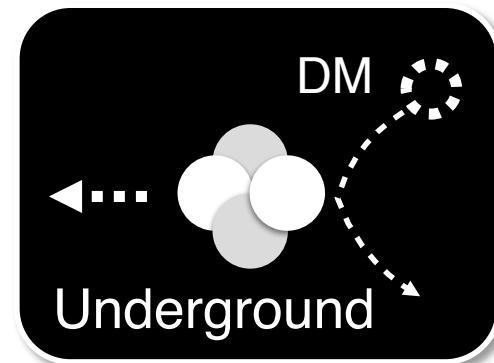
Hunting for Particle DM



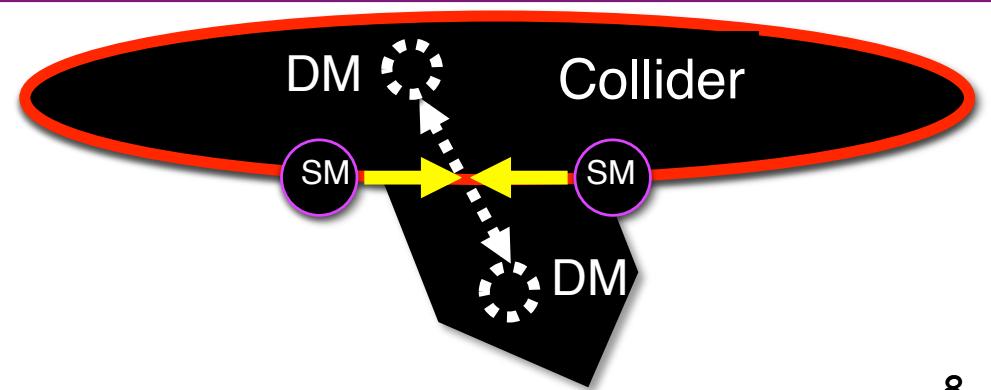
Indirect
Detection
Experiments



Direct
Detection
Experiments



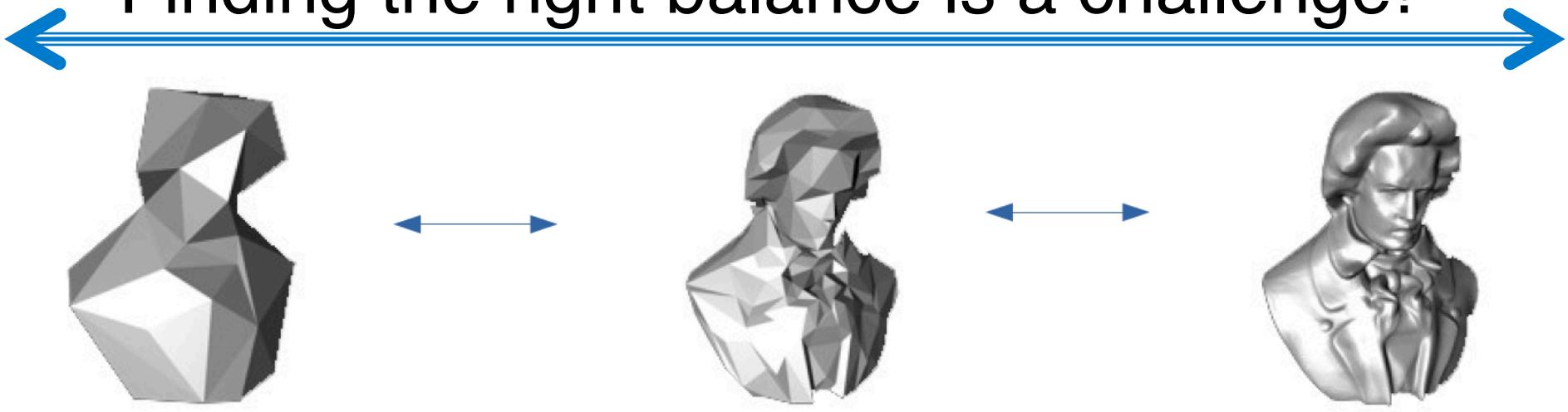
Collider
Experiments



Characterisation of Dark Matter searches at colliders

Simplicity vs. Complexity

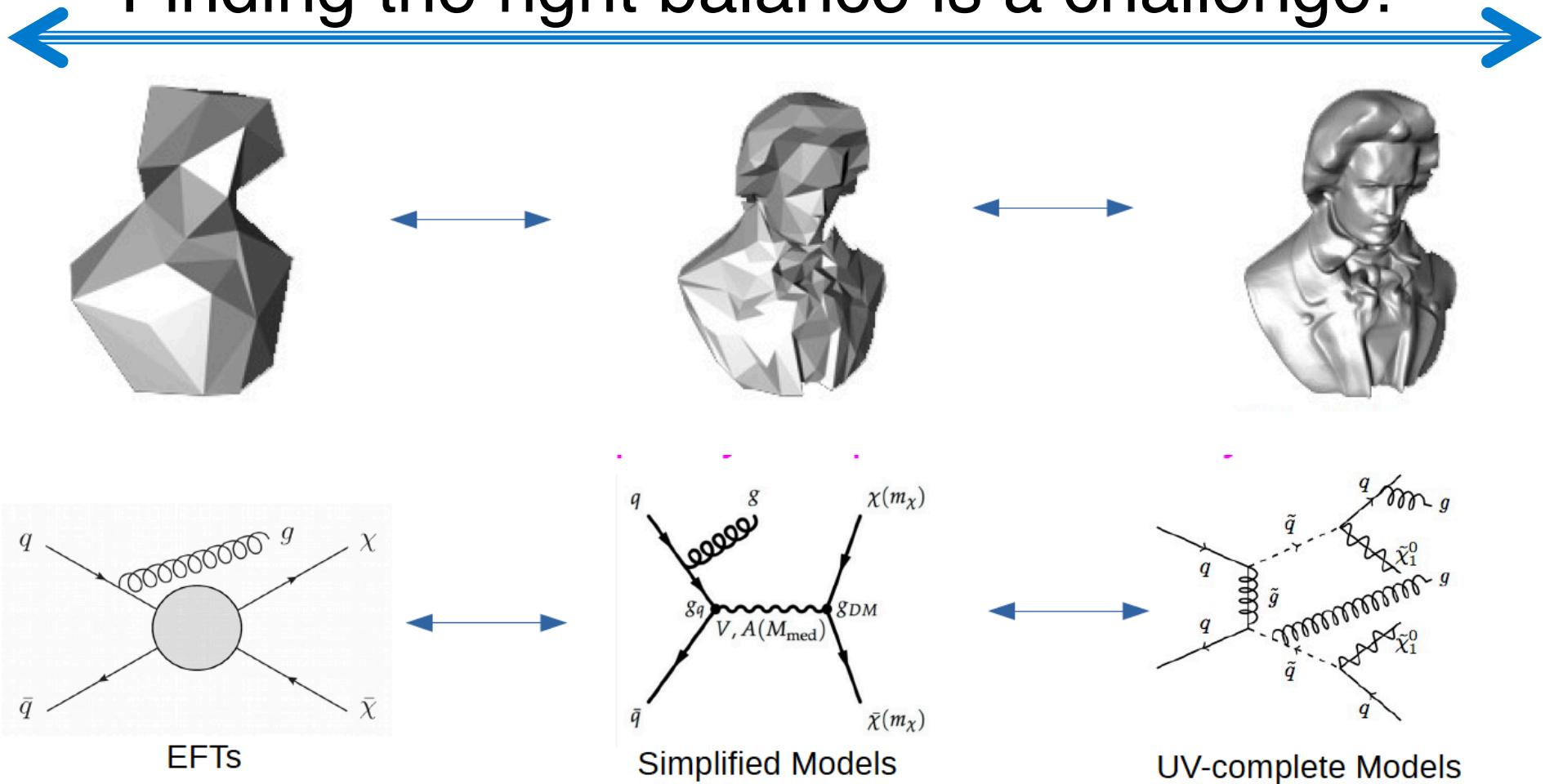
Finding the right balance is a challenge!



Characterisation of Dark Matter searches at colliders

Simplicity vs. Complexity

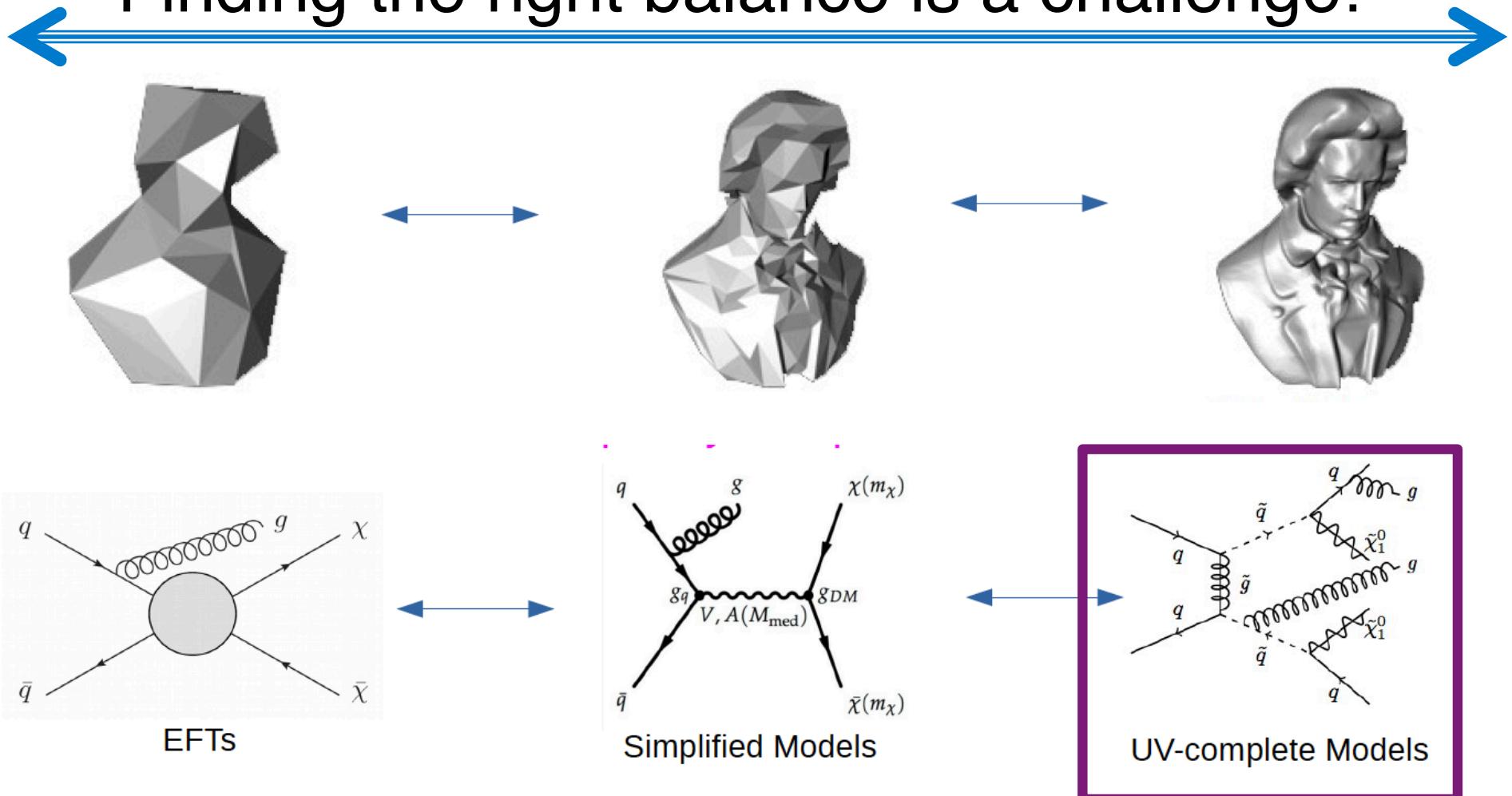
Finding the right balance is a challenge!



Characterisation of Dark Matter searches at colliders

Simplicity vs. Complexity

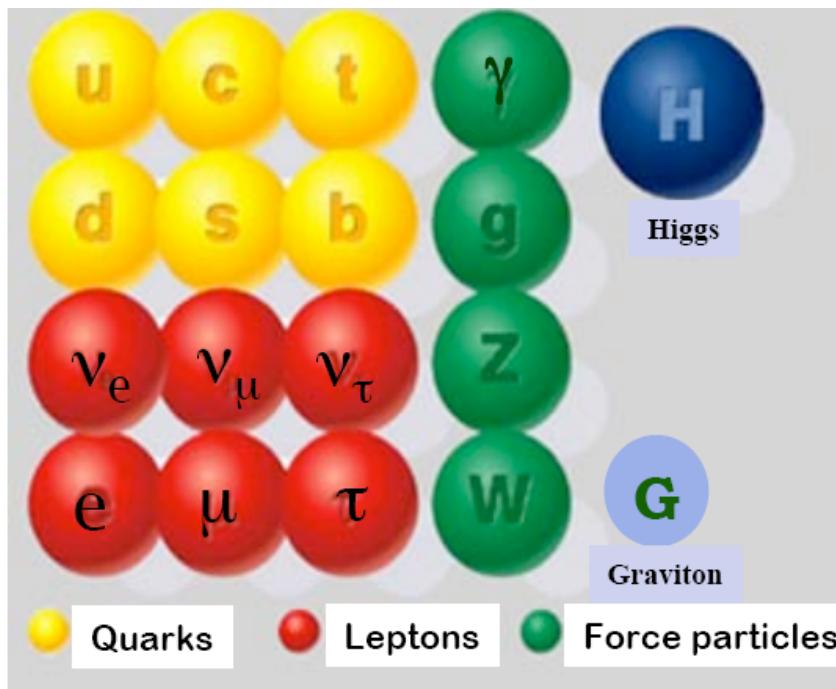
Finding the right balance is a challenge!



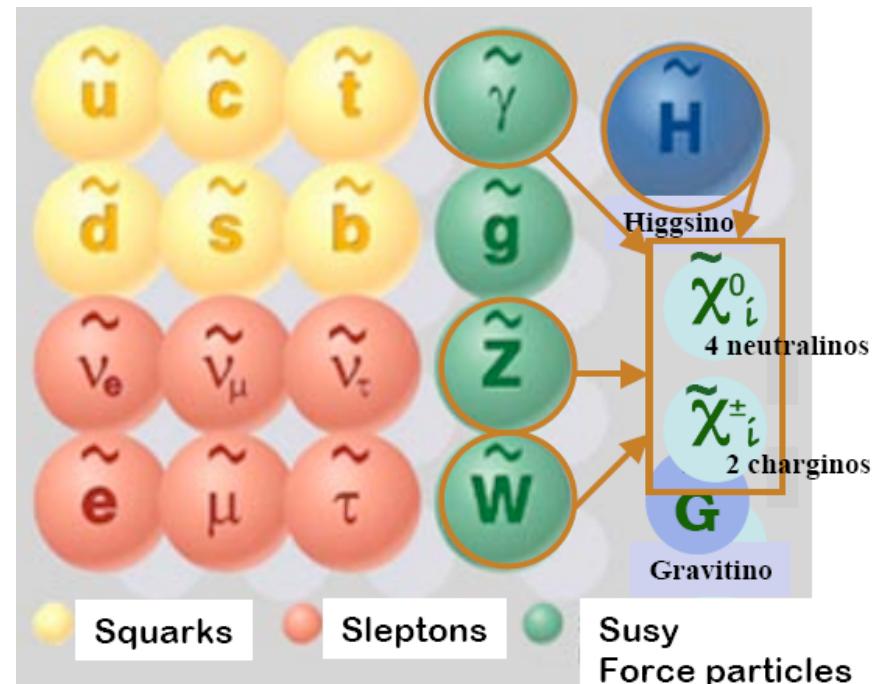
Supersymmetry

Extension of the Standard Model: Introduce a new symmetry
 Spin $\frac{1}{2}$ matter particles (fermions) \Leftrightarrow Spin 1 force carriers (bosons)

Standard Model particles



SUSY particles



New Quantum number: R-parity:

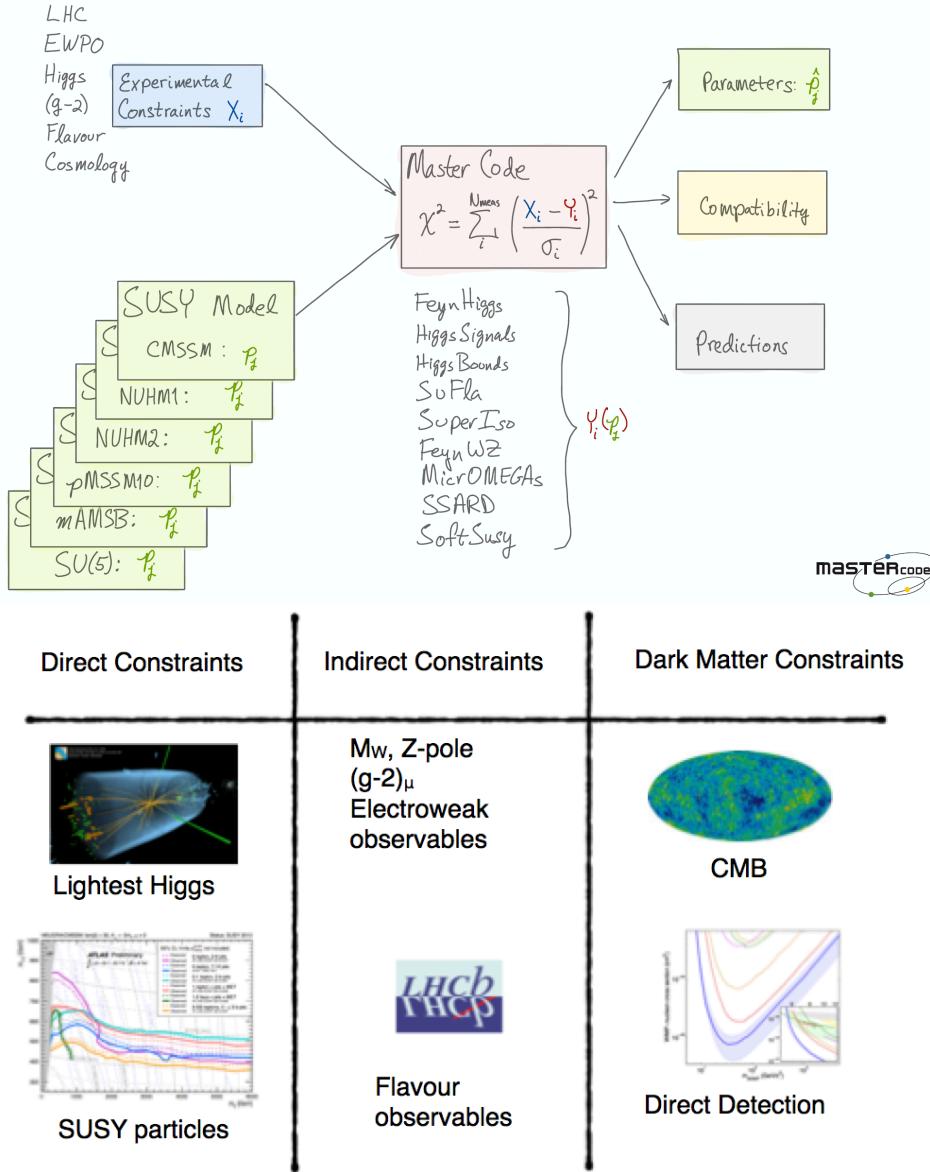
R-parity conservation:

- SUSY particles are produced in pairs
- The lightest SUSY particle (LSP) is stable

$$R_p = (-1)^{B+L+2s} = \begin{cases} +1 & \text{SM particles} \\ -1 & \text{SUSY particles} \end{cases}$$

Dark Matter in Supersymmetry with MasterCode

Global Fit to indirect and direct constraints on SUSY!

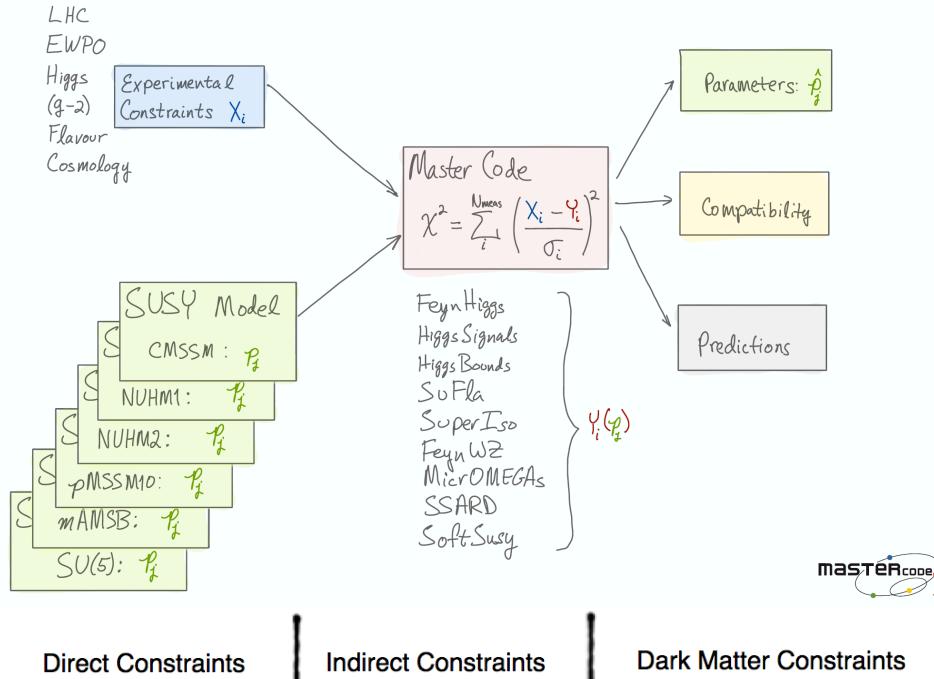


Source:
<http://mastercode.web.cern.ch/mastercode/>

Observable	Source Th./Ex.	Constraint	$\Delta\chi^2$ (CMSSM)	$\Delta\chi^2$ (NUHM1)	$\Delta\chi^2$ ("SM")
m_t [GeV]	[43]	173.2 ± 0.90	0.05	0.06	-
$\Delta\alpha_{\text{had}}^{(0)}(M_Z)$	[42]	0.02749 ± 0.00010	0.009	0.004	-
M_Z [GeV]	[44]	91.1875 ± 0.0021	2.7×10^{-6}	0.26	-
Γ_Z [GeV]	[26] / [44]	$2.4952 \pm 0.0023 \pm 0.001_{\text{SUSY}}$	0.078	0.047	0.14
σ_{had}^0 [nb]	[26] / [44]	41.540 ± 0.037	2.50	2.57	2.54
R_l	[26] / [44]	20.767 ± 0.025	1.05	1.08	1.08
$A_{fb}(l)$	[26] / [44]	0.01714 ± 0.00095	0.72	0.69	0.81
$A_{fb}(P_T)$	[26] / [44]	0.1465 ± 0.0032	0.11	0.13	0.07
R_b	[26] / [44]	0.21629 ± 0.00066	0.26	0.29	0.27
R_c	[26] / [44]	0.1721 ± 0.0030	0.002	0.002	0.002
$A_{fb}(b)$	[26] / [44]	0.0992 ± 0.0016	7.17	7.37	6.63
$A_{fb}(c)$	[26] / [44]	0.0707 ± 0.0035	0.86	0.88	0.80
A_b	[26] / [44]	0.923 ± 0.020	0.36	0.36	0.35
A_c	[26] / [44]	0.670 ± 0.027	0.005	0.005	0.005
A_t (SLD)	[26] / [44]	0.1513 ± 0.0021	3.16	3.03	3.51
$\sin^2 \theta_w(Q_b)$	[26] / [44]	0.2324 ± 0.0012	0.63	0.64	0.59
M_W [GeV]	[26] / [44]	$80.399 \pm 0.023 \pm 0.010_{\text{SUSY}}$	1.77	1.39	2.08
$a_\mu^{\text{EXP}} - a_\mu^{\text{SM}}$	[53] / [42,54]	$(30.2 \pm 8.8 \pm 2.0_{\text{SUSY}}) \times 10^{-10}$	4.35	1.82	11.19 (N/A)
M_h [GeV]	[28] / [55,56]	$> 114.4 [\pm 1.5_{\text{SUSY}}]$	0.0	0.0	0.0
$\text{BR}_{b \rightarrow s \gamma}^{\text{EXP/SM}}$	[45] / [46]	$1.117 \pm 0.076_{\text{EXP}} \pm 0.082_{\text{SM}} \pm 0.050_{\text{SUSY}}$	1.83	1.09	0.94
$\text{BR}(B_s \rightarrow \mu^+ \mu^-)$	[29] / [41]	CMS & LHCb	0.04	0.44	0.01
$\text{BR}_{B \rightarrow \tau \nu}^{\text{EXP/SM}}$	[29] / [46]	$1.43 \pm 0.43_{\text{EXP+TH}}$	1.43	1.59	1.00
$\text{BR}(B_d \rightarrow \mu^+ \mu^-)$	[29] / [46]	$< 4.6 [\pm 0.01_{\text{SUSY}}] \times 10^{-6}$	0.0	0.0	0.0
$\text{BR}_{B \rightarrow X_s \ell \bar{\nu}}^{\text{EXP/SM}}$	[47] / [46]	0.99 ± 0.32	0.02	$\ll 0.01$	$\ll 0.01$
$\text{BR}_{K \rightarrow \ell \nu}^{\text{EXP/SM}}$	[29] / [48]	$1.008 \pm 0.014_{\text{EXP+TH}}$	0.39	0.42	0.33
$\text{BR}_{K \rightarrow \pi \nu}^{\text{EXP/SM}}$	[49] / [50]	< 4.5	0.0	0.0	0.0
ΔM_B^{exp}	[49] / [51,52]	$0.97 \pm 0.01_{\text{EXP}} \pm 0.27_{\text{SM}}$	0.02	0.02	0.01
$\Delta M_{B_s}^{\text{exp/SM}}$	[29] / [46,51,52]	$1.00 \pm 0.01_{\text{EXP}} \pm 0.13_{\text{SM}}$	$\ll 0.01$	0.33	$\ll 0.01$
$\Delta M_{B_d}^{\text{exp/SM}}$	[49] / [51,52]	$1.08 \pm 0.14_{\text{EXP+TH}}$	0.27	0.37	0.33
$\Omega_{\text{CDM}} h^2$	[31] / [13]	$0.1120 \pm 0.0056 \pm 0.012_{\text{SUSY}}$	8.4×10^{-4}	0.1	N/A
σ_8^2	[25]	(m_{χ_1}, σ_8^2) plane	0.13	0.13	N/A
jets + E_T	[18,20]	$(m_0, m_{1/2})$ plane	1.55	2.20	N/A
$H/A, H^\pm$	[21]	$(M_A, \tan \beta)$ plane	0.0	0.0	N/A
Total $\chi^2/\text{d.o.f.}$	All	All	28.8/22	27.3/21	32.7/23 (21.5/22)
p-values			15%	16%	9% (49%)

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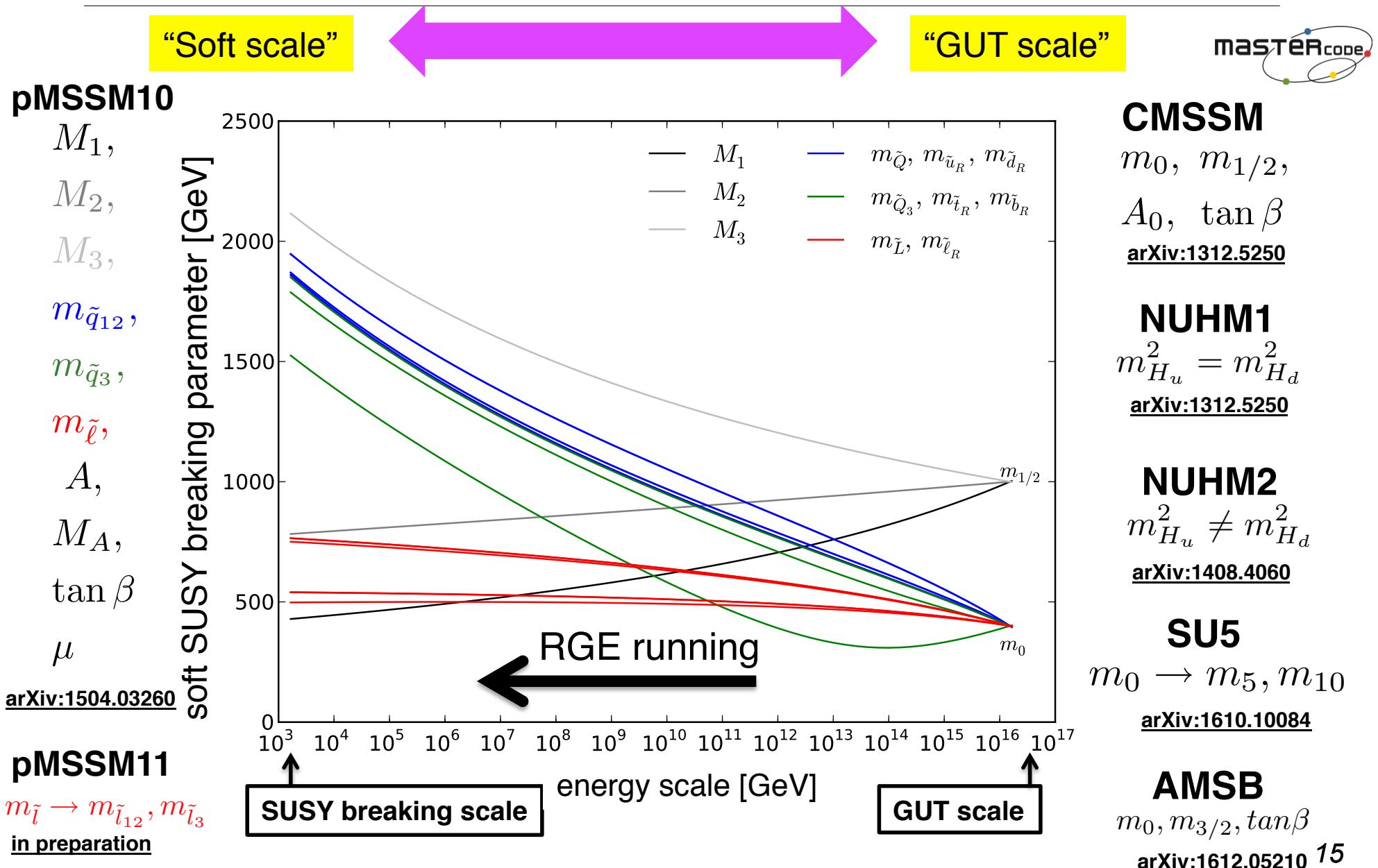
Other “global Fitters” with similar studies are:
 Fittino group: [see e.g. arXiv:1508.05951]
<http://flcwiki.desy.de/Fittino>
 Gambit group: [see e.g. arXiv:1705.07917]
<https://gambit.hepforge.org>
 SuperBayeS: [see e.g. arXiv:1507.07008]



Flavour observables

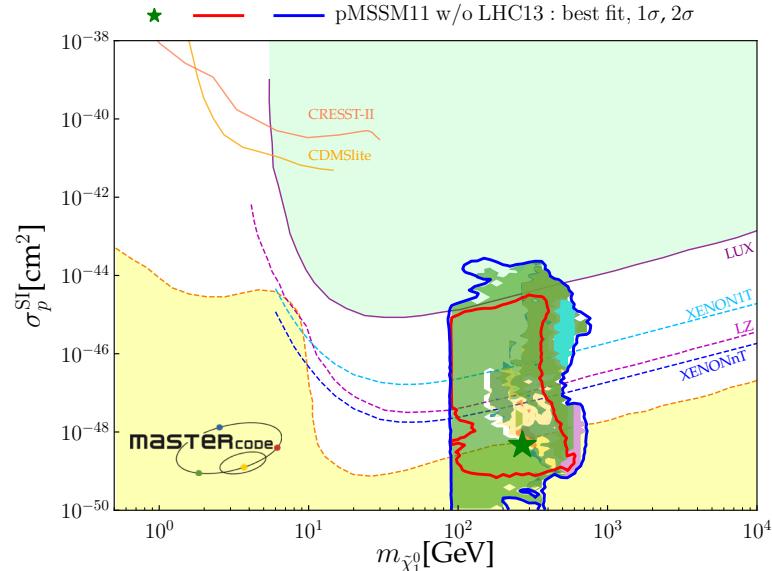
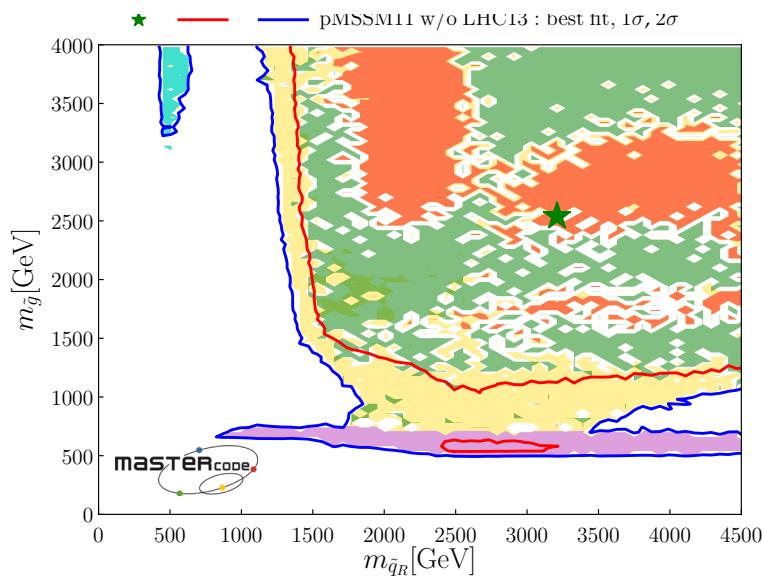
Direct Detection

MasterCode: The two worlds of SUSY models



pMSSM11: Status LHC RUN 1 (pre-LHC 13 TeV)

■ stau coann. ■ $\tilde{\chi}_1^\pm$ coann. ■ slep coann. ■ gluino coann. ■ squark coann.



DM mechanisms:

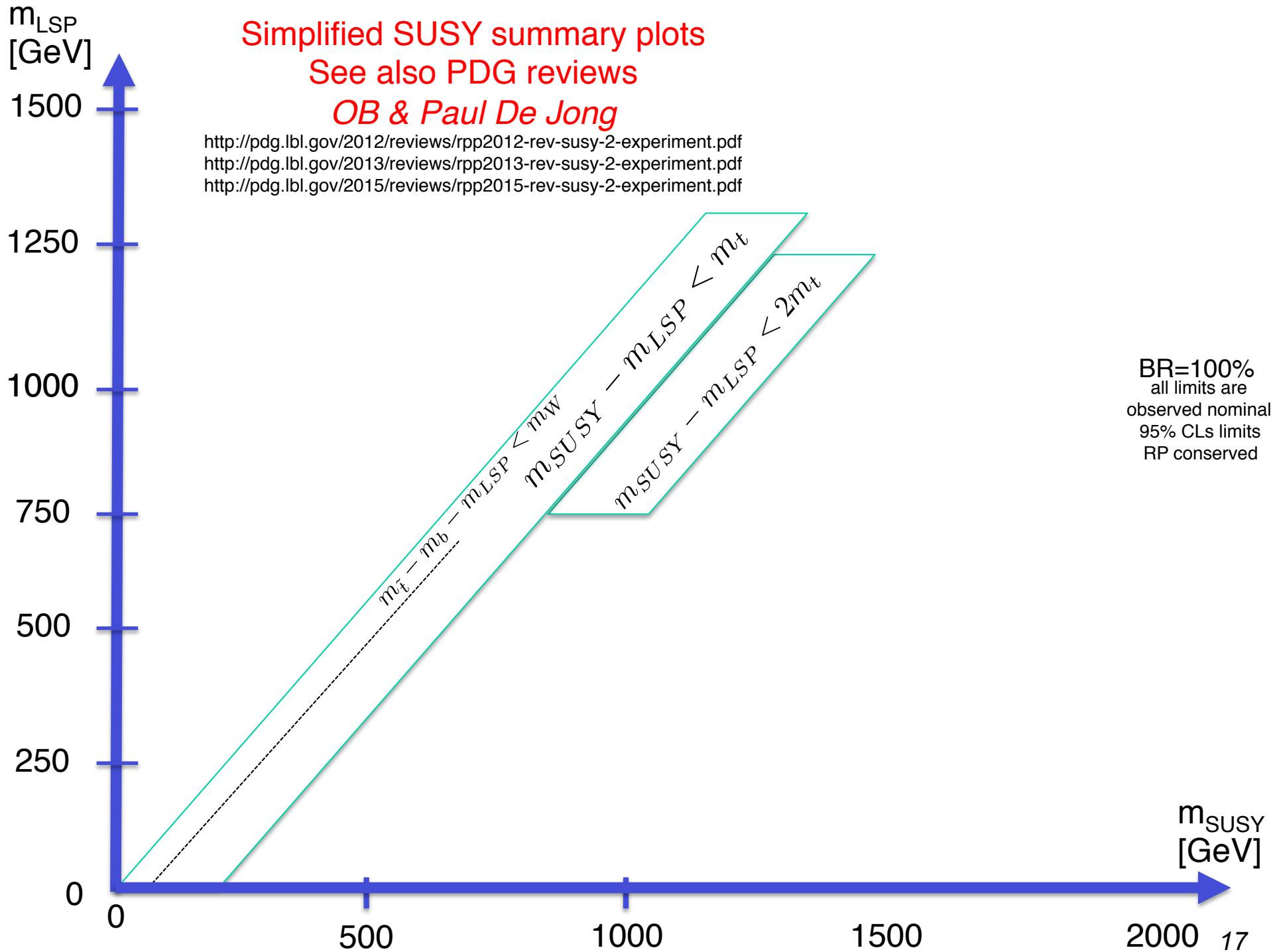
To satisfy cosmological DM density constraint requires, in general, specific relations between sparticle masses that suppress the relic density via coannihilation effects and/or rapid annihilations through direct channel resonances.

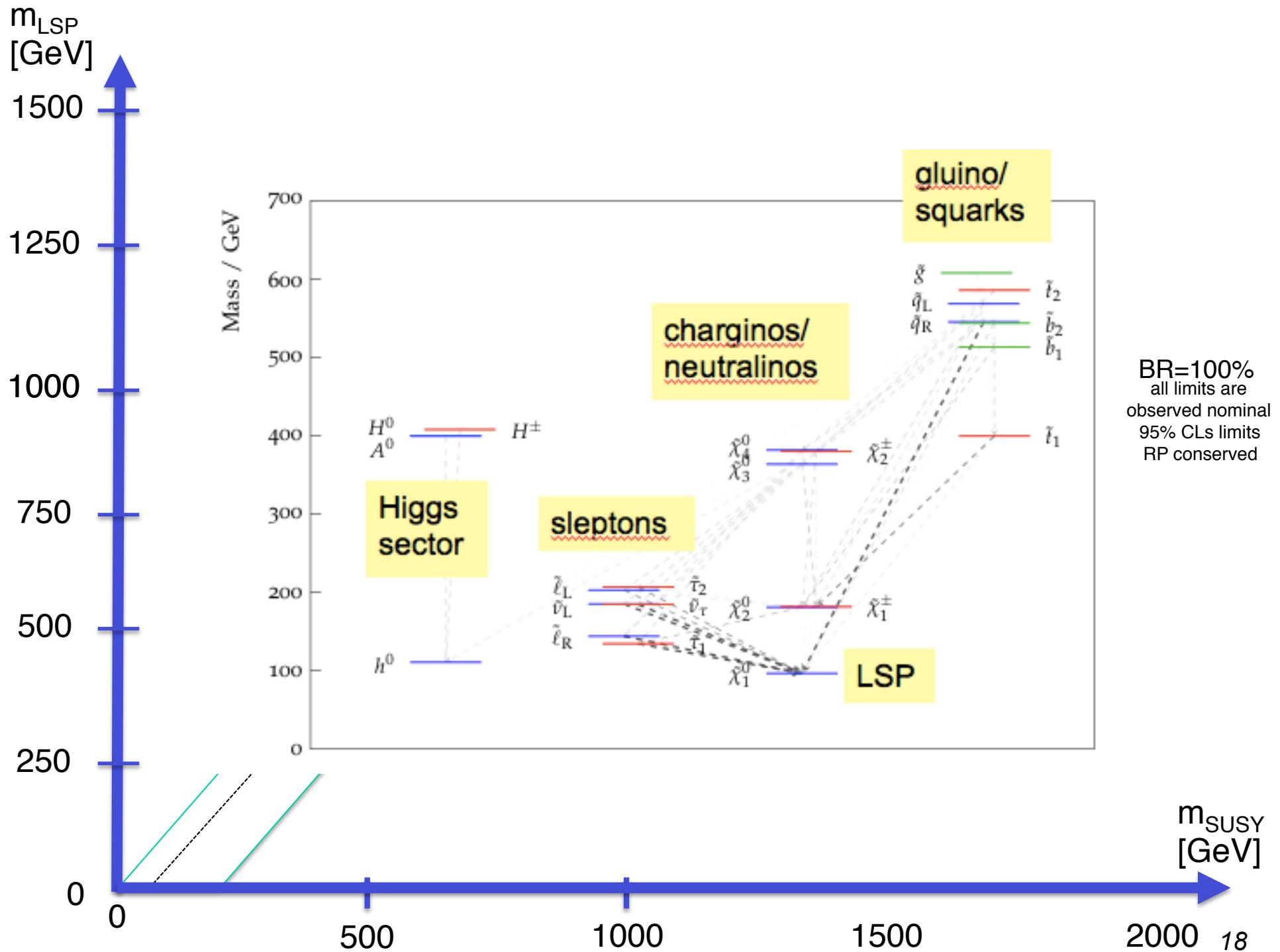
Define indicative measures to highlight different DM mechanisms in the preferred regions of the fit:

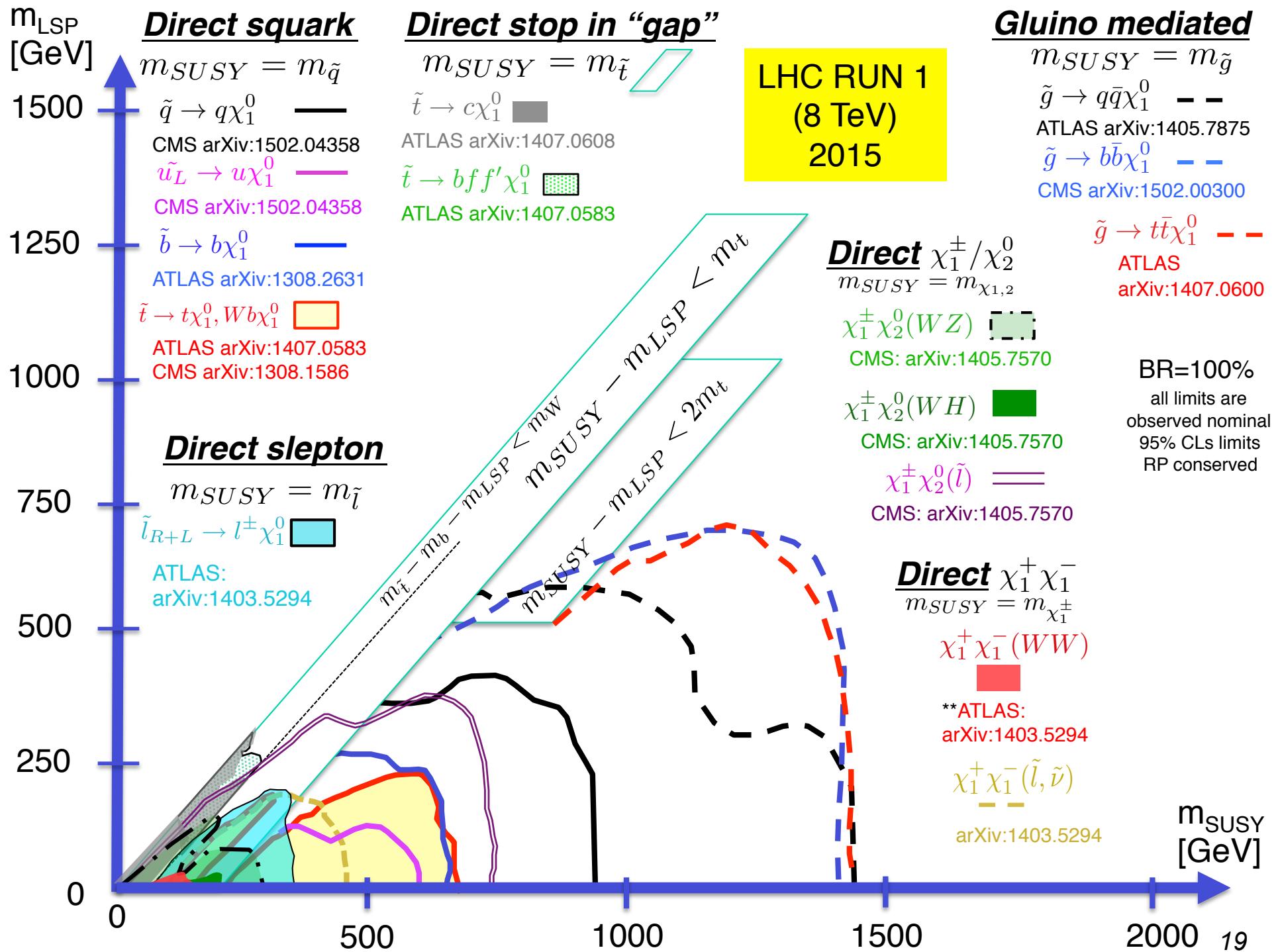
$\left(\frac{M_{\tilde{\tau}}}{m_{\tilde{\chi}_1^0}} - 1 \right) < 0.15$	Stau coannihilation	$\left(\frac{M_{\tilde{l}}}{m_{\tilde{\chi}_1^0}} - 1 \right) < 0.15$	Slepton Co-annihilation
$\left(\frac{M_{\tilde{\chi}_1^\pm}}{m_{\tilde{\chi}_1^0}} - 1 \right) < 0.25$	Chargino Co-annihilation	$\left(\frac{M_{\tilde{g}}}{m_{\tilde{\chi}_1^0}} - 1 \right) < 0.25$	Gluino Co-annihilation
$\left(\frac{M_{\tilde{q}}}{m_{\tilde{\chi}_1^0}} - 1 \right) < 0.20$	Squark Co-annihilation		
$\left \frac{M_B}{m_{\tilde{\chi}_1^0}} - 2 \right < 0.4$	B = h, Z or H/A funnel		
$\left \frac{\mu}{m_{\tilde{\chi}_1^0}} - 1 \right < 0.30$	Higgsino enriched "focus-point" like		

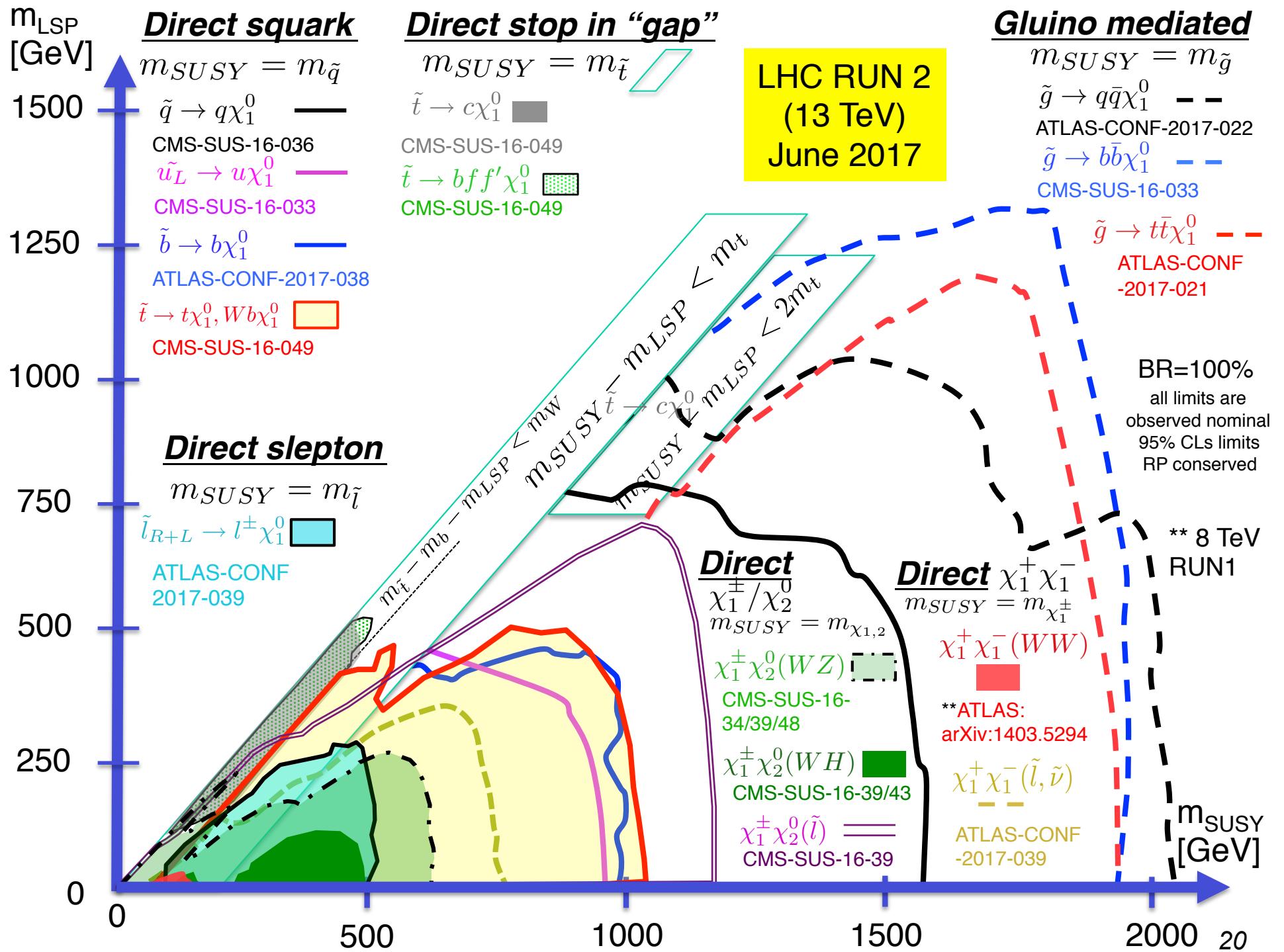
Hybrid regions:
In addition to the 'primary' regions where only one of the conditions is satisfied, there can also be 'hybrid' regions where more than one condition is satisfied. If present, these are indicated using combined colours.

See also arXiv:1508.01173 for further details

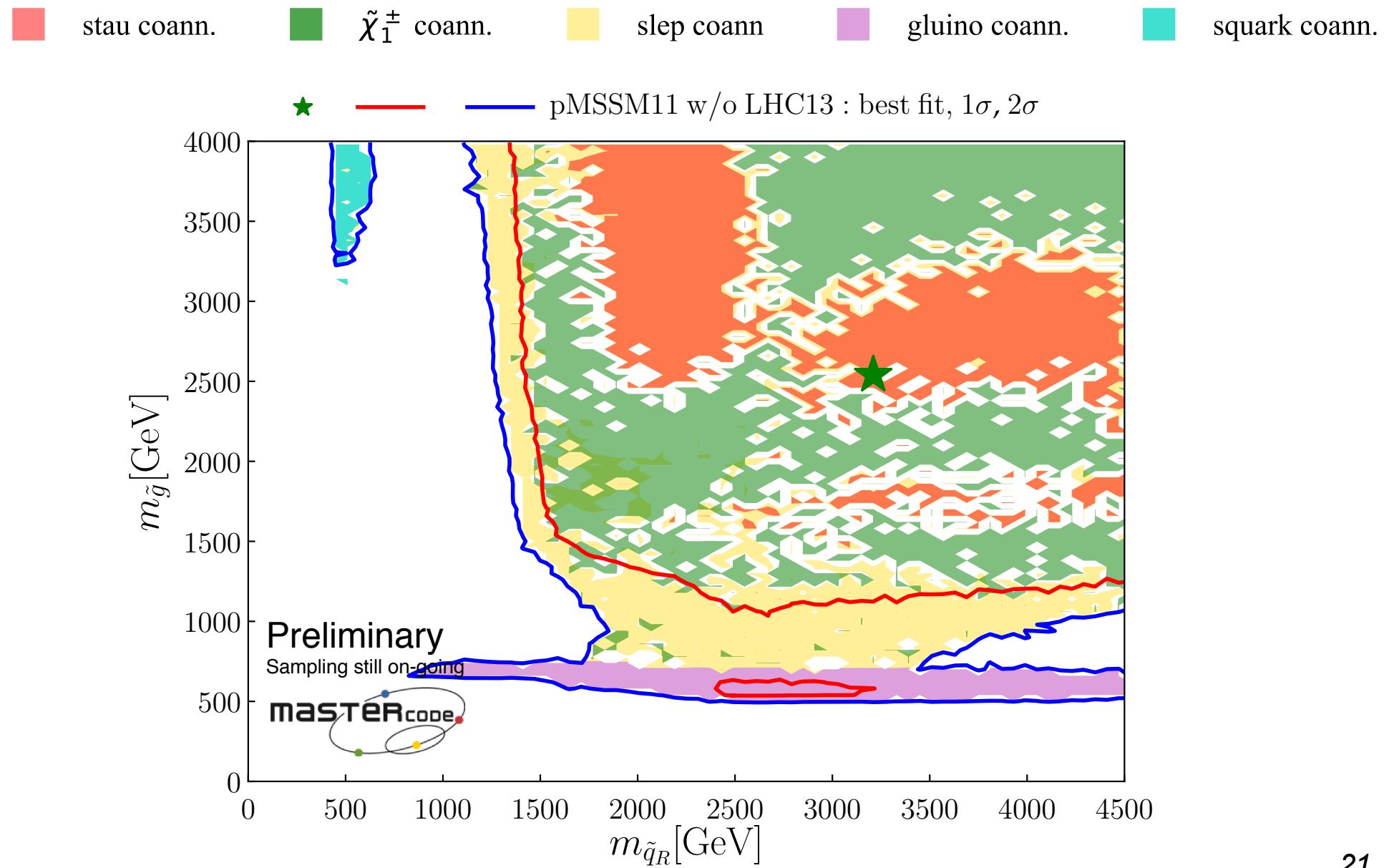




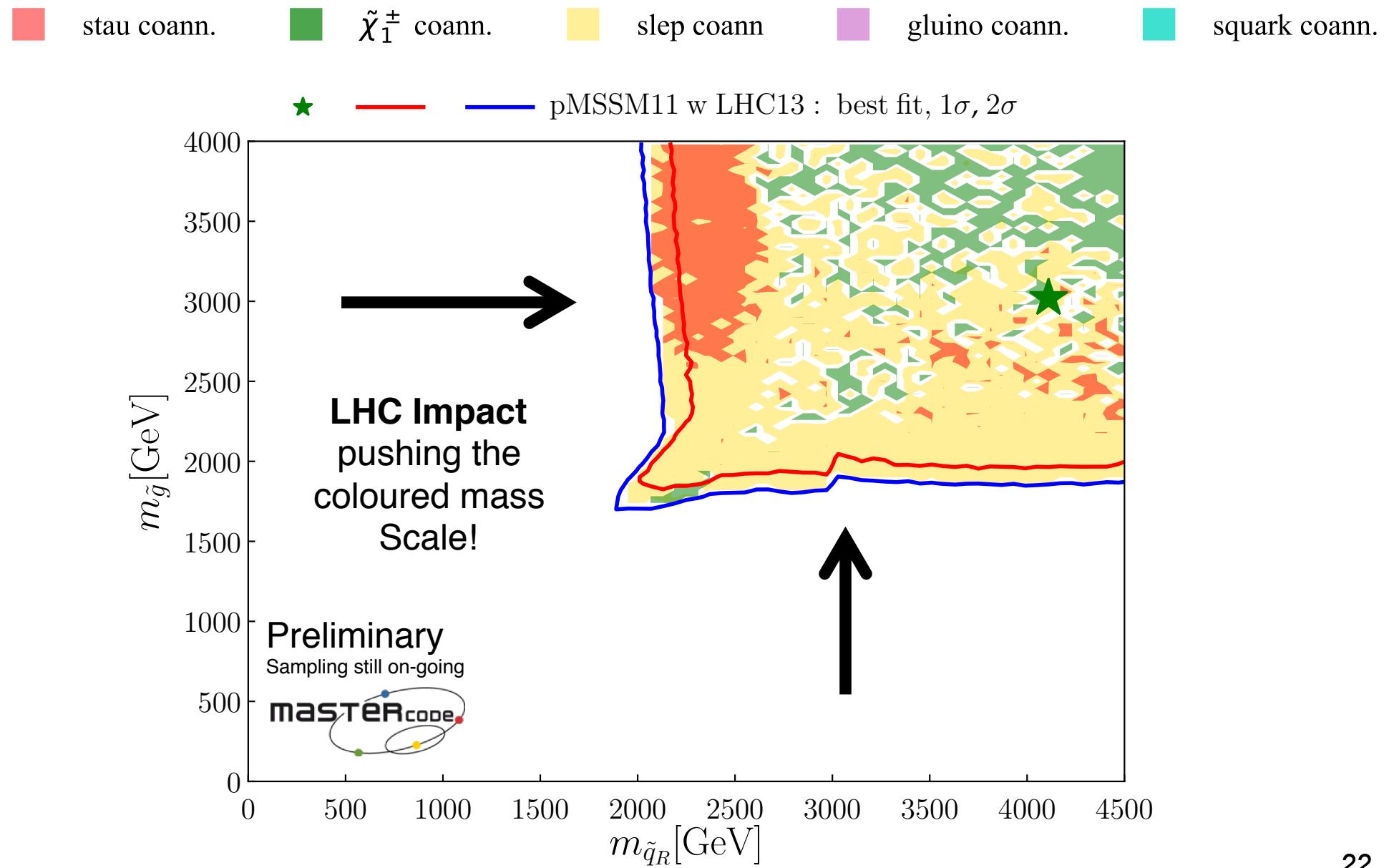




Gluino vs Squark: LHC RUN 1

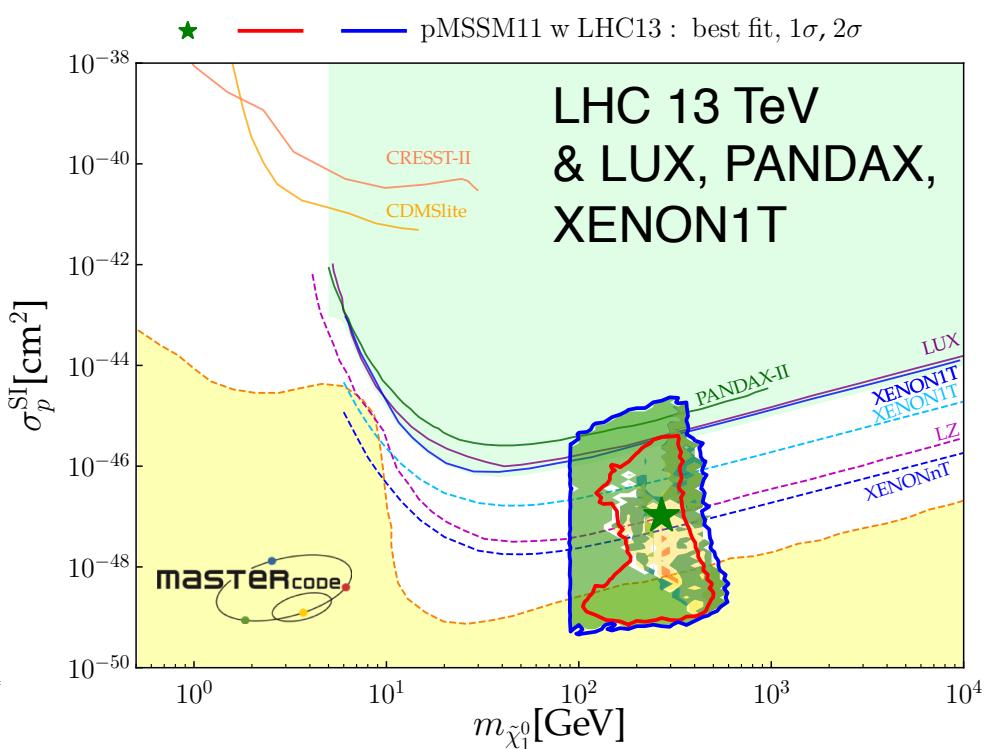
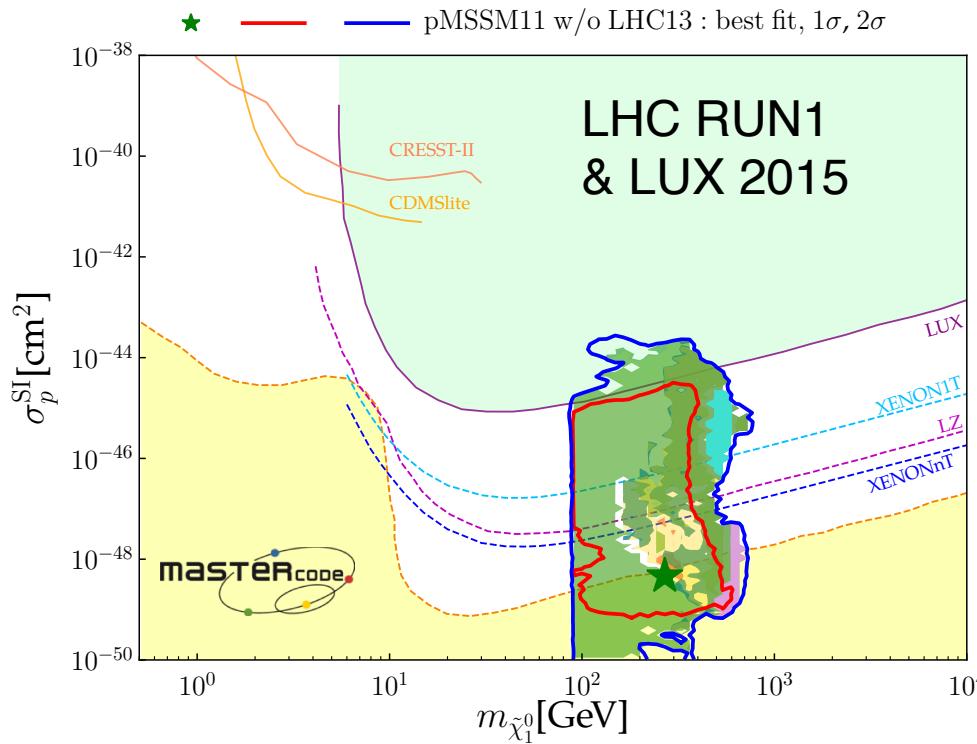


Gluino vs Squark: LHC RUN 2 (2015 + 2016 data)



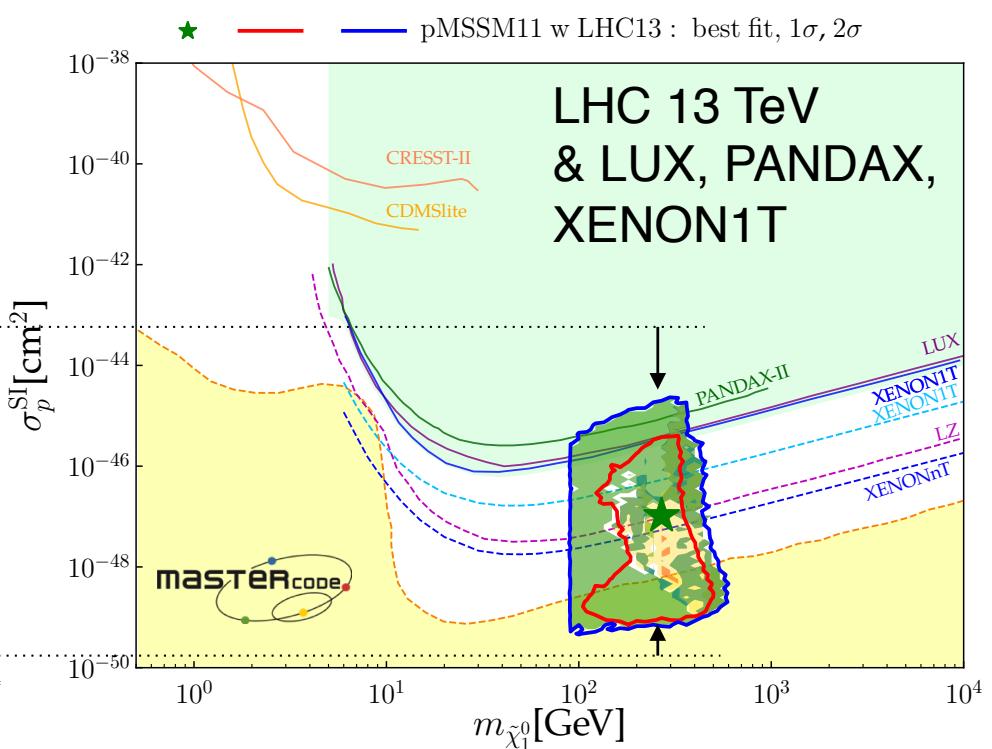
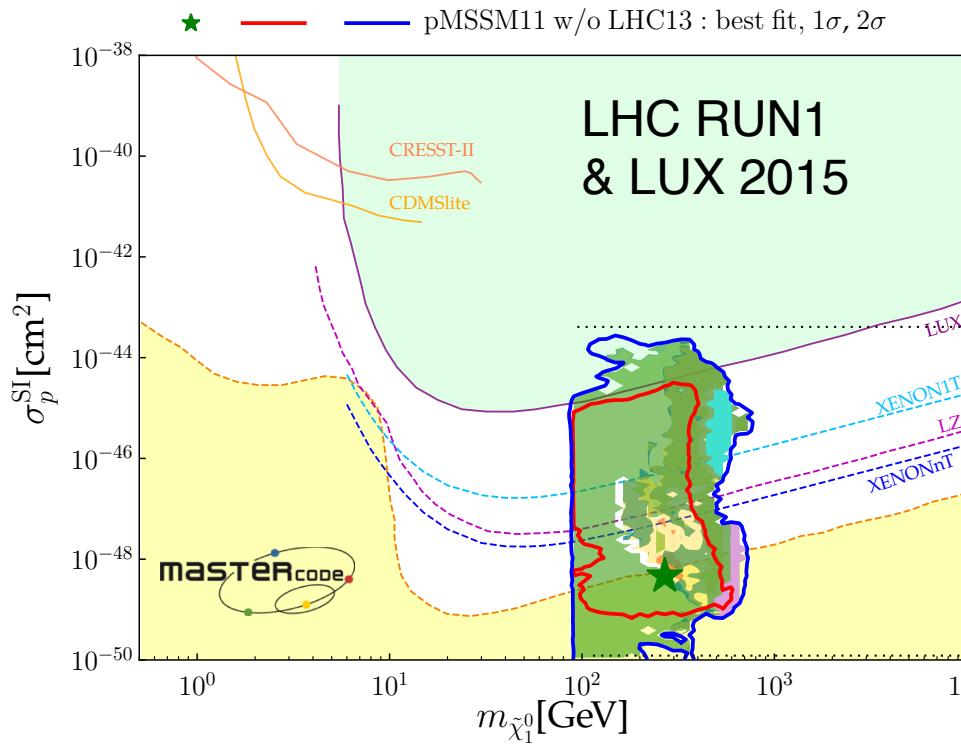
pMSSM11: σ_{SI} vs $m_{\tilde{\chi}_1^0}$

■ stau coann. ■ $\tilde{\chi}_1^\pm$ coann. ■ slep coann. ■ gluino coann. ■ squark coann.



pMSSM11: σ_{SI} vs m_{DM}

■ stau coann. ■ $\tilde{\chi}_1^\pm$ coann. ■ slep coann. ■ gluino coann. ■ squark coann.



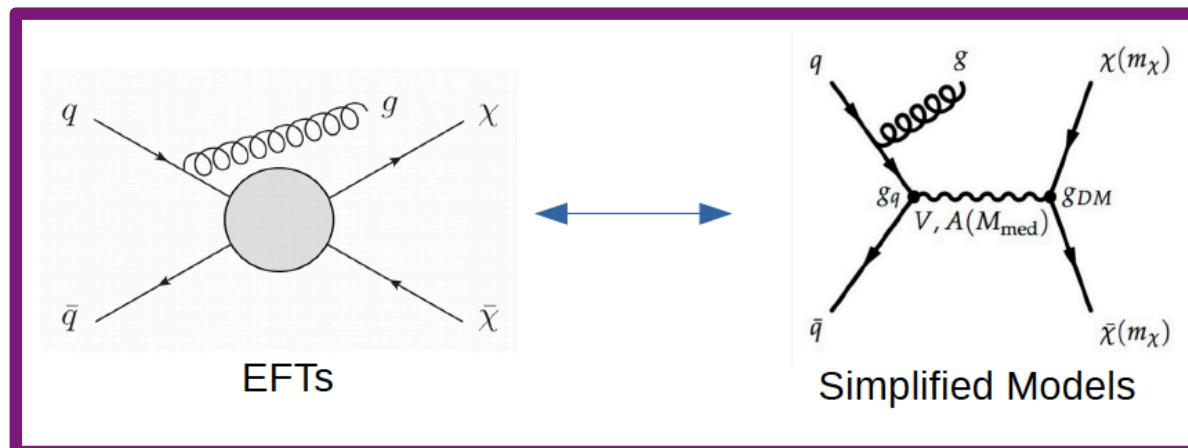
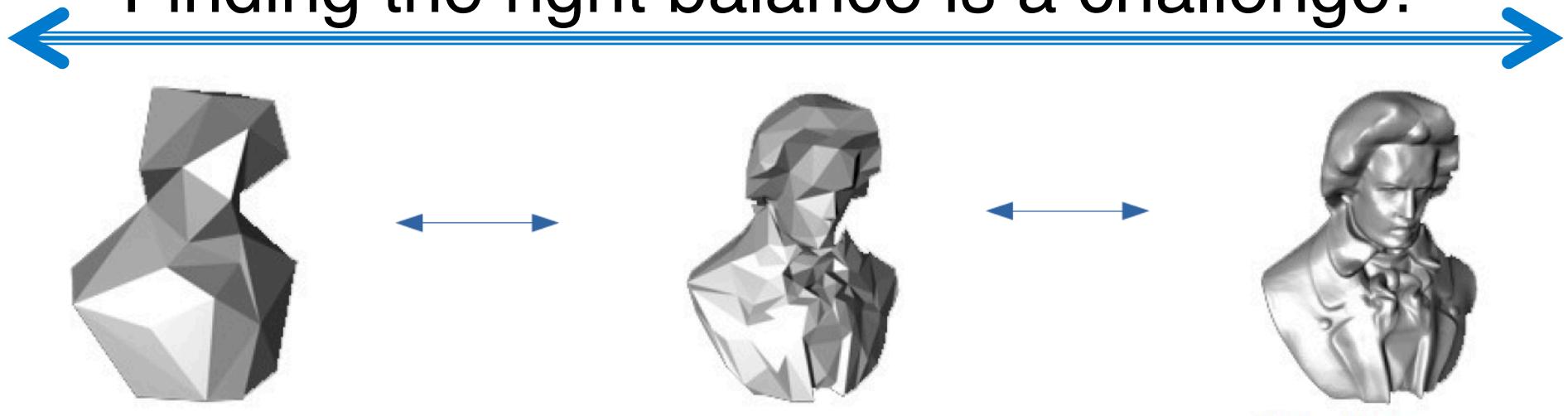
Clear complementarity of collider and DD constraints:

- Collider covers regions not easily or not at all accessible to DD experiments (i.e. low m_{DM} and also very small σ_{SI})
- On the other hand, DD experiments push strongly the preferred region to lower σ_{SI} (and will continue to do so in the future)

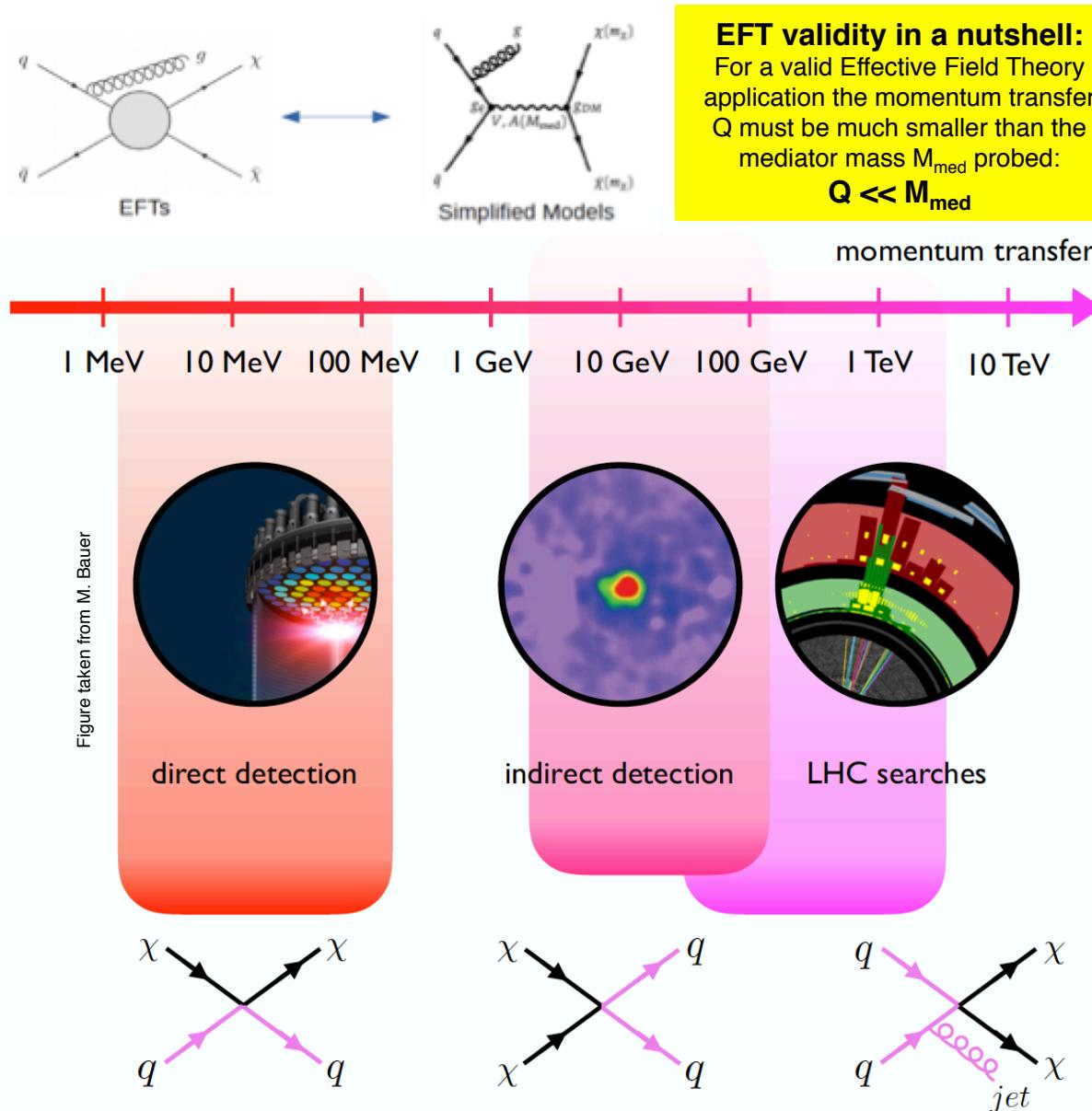
Characterisation of Dark Matter searches at colliders

Simplicity vs. Complexity

Finding the right balance is a challenge!



EFT vs Simplified Model



Therefore, validity requires typically:

DD Experiments:

$M_{\text{med}} >$ few hundred **MeV**

ID Experiments:

$M_{\text{med}} >$ few hundred **GeV**

Collider (LHC):

$M_{\text{med}} >$ few **TeV**

As the LHC probes the TeV scale, a comprehensive application of the EFT for DM searches is not possible.

Therefore, adopt simplified DM models as main vehicle to interpret DM searches for LHC!

For more info about EFT validity for DM collider searches see e.g.:

[arXiv:1307.2253](https://arxiv.org/abs/1307.2253)

[arXiv:1308.6799](https://arxiv.org/abs/1308.6799)

[arXiv:1405.3101](https://arxiv.org/abs/1405.3101)

[arXiv:1402.1275](https://arxiv.org/abs/1402.1275)

LHC Dark Matter Working Group

End of Run-1: Discussion on how to present the Dark Matter search data in the experiments for Run-2, in a DM forum and now DM working group

First collection on
DM models (simplified
and look-alike) for LHC

Dark Matter Benchmark Models for Early LHC Run-2 Searches:
Report of the ATLAS/CMS Dark Matter Forum

arXiv:1507.00966

Guidelines how to compare
data from LHC and non-LHC
search results

Recommendations on presenting LHC
searches for missing transverse energy
signals using simplified s -channel models
of dark matter

arXiv:1603.04156

Guidelines for direct DM production
searches with constraints on the
heavy mediators

Recommendations of the LHC Dark Matter
Working Group: Comparing LHC searches for
heavy mediators of dark matter production in
visible and invisible decay channels

arXiv:1703.05703

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End of Run-1: Discussion on how to present the Dark Matter search data in the experiments for Run-2, in a DM forum and now DM working group

First collection on

Dark Matter Benchmark Models for Early LHC Run-2 Searches:
Report of the ATLAS/CMS-Dark Matter Working Group

This Working Group brings together theorists and experimentalists to define guidelines and recommendations for the benchmark models, interpretation, and characterisation necessary for broad and systematic searches for dark matter at the LHC.

More details can be found at this page:

http://lpcc.web.cern.ch/LPCC/index.php?page=dm_wg

and the mailing list is lhc-dmwg@cern.ch**.

**To join the WG mailing list, go to

<http://simba3.web.cern.ch/simba3/SelfSubscription.aspx?groupName=lhc-dmwg>

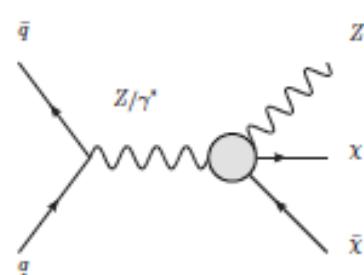
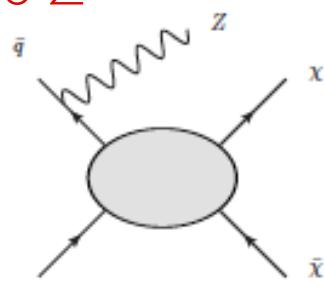
Guidelines for direct DM production searches with constraints on the heavy mediators

Recommendations of the LHC Dark Matter Working Group: Comparing LHC searches for heavy mediators of dark matter production in visible and invisible decay channels

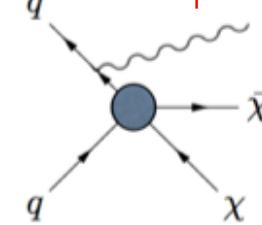
arXiv:1703.05703

Mono-Mania (at the LHC)

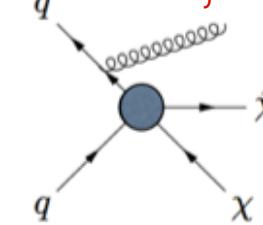
Mono-Z



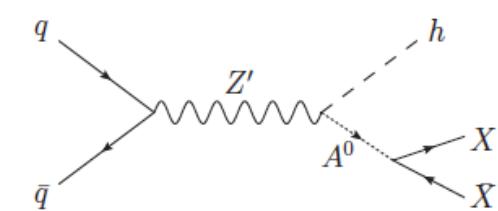
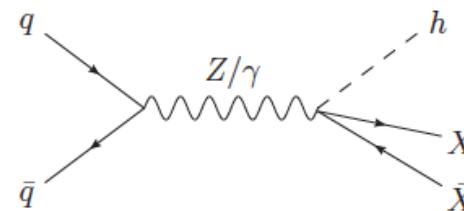
Mono-photon



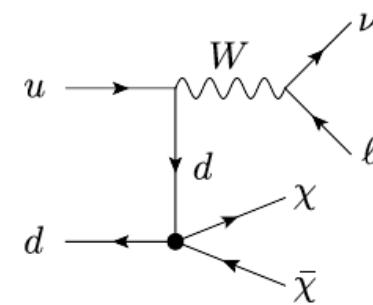
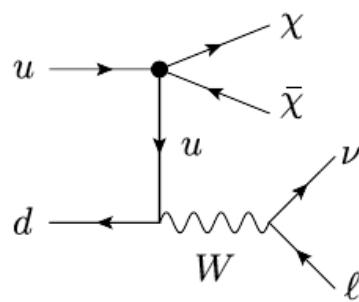
Mono-jet



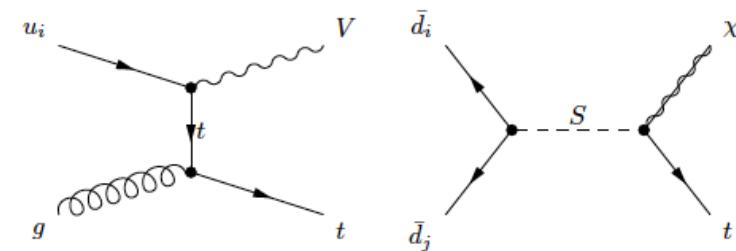
Mono-Higgs



Mono-W

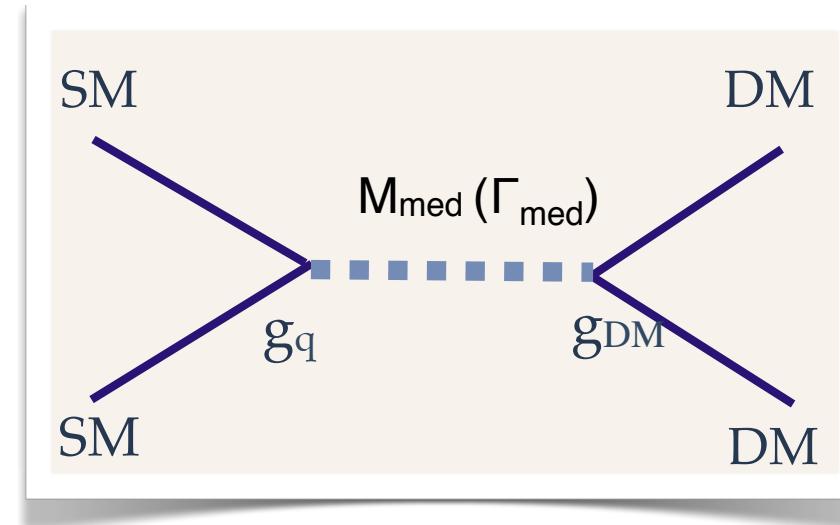


Mono-top



Minimal Simplified Dark Matter Model

See e.g.
 arXiv:1407.8257
 arXiv:1507.00966
 arXiv:1603.04156



s-channel

Define simplified model with
(minimum) 4 parameters

Mediator mass (M_{med})	DM mass (M_{DM})
g_q	g_{DM}

DM	
Dirac fermion	Scalar - real
Majorana fermion	Scalar - complex

Consider comprehensive set
of diagrams for mediator

Vector	Axial-vector
Scalar	Pseudoscalar

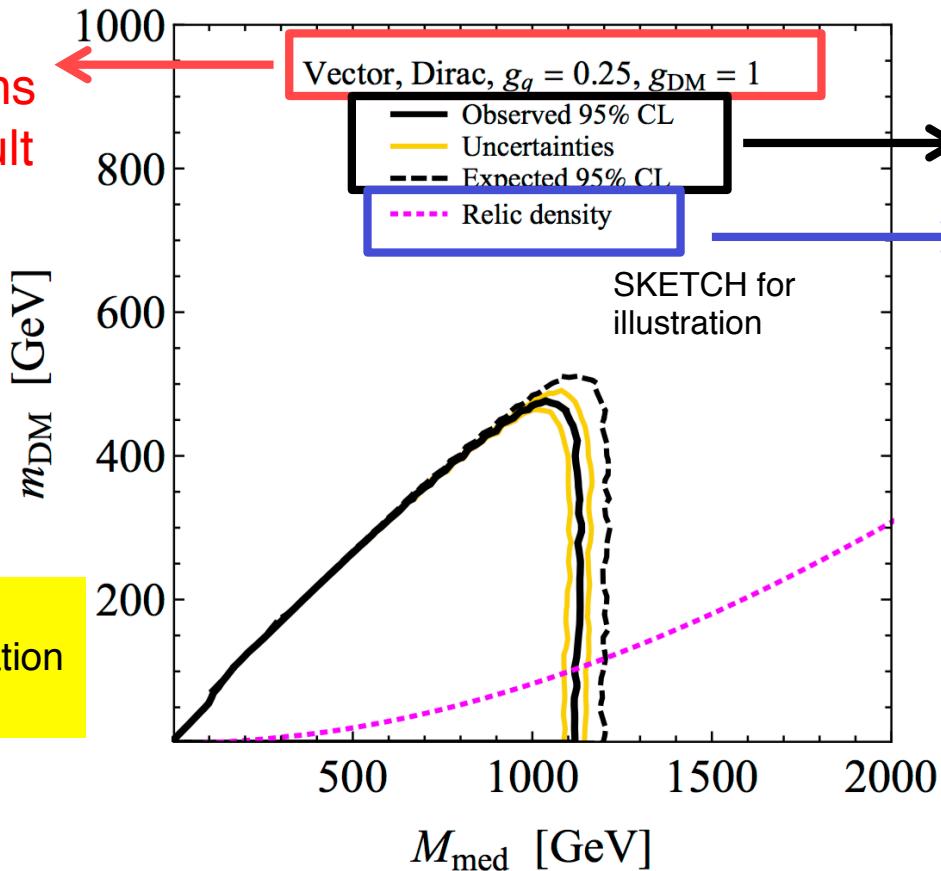
$(\Gamma_{med}$ can also be free as long
As $\Gamma_{med} < M_{med}$)

Mass-Mass plane [$M_{\text{med}} - M_{\text{DM}}$]

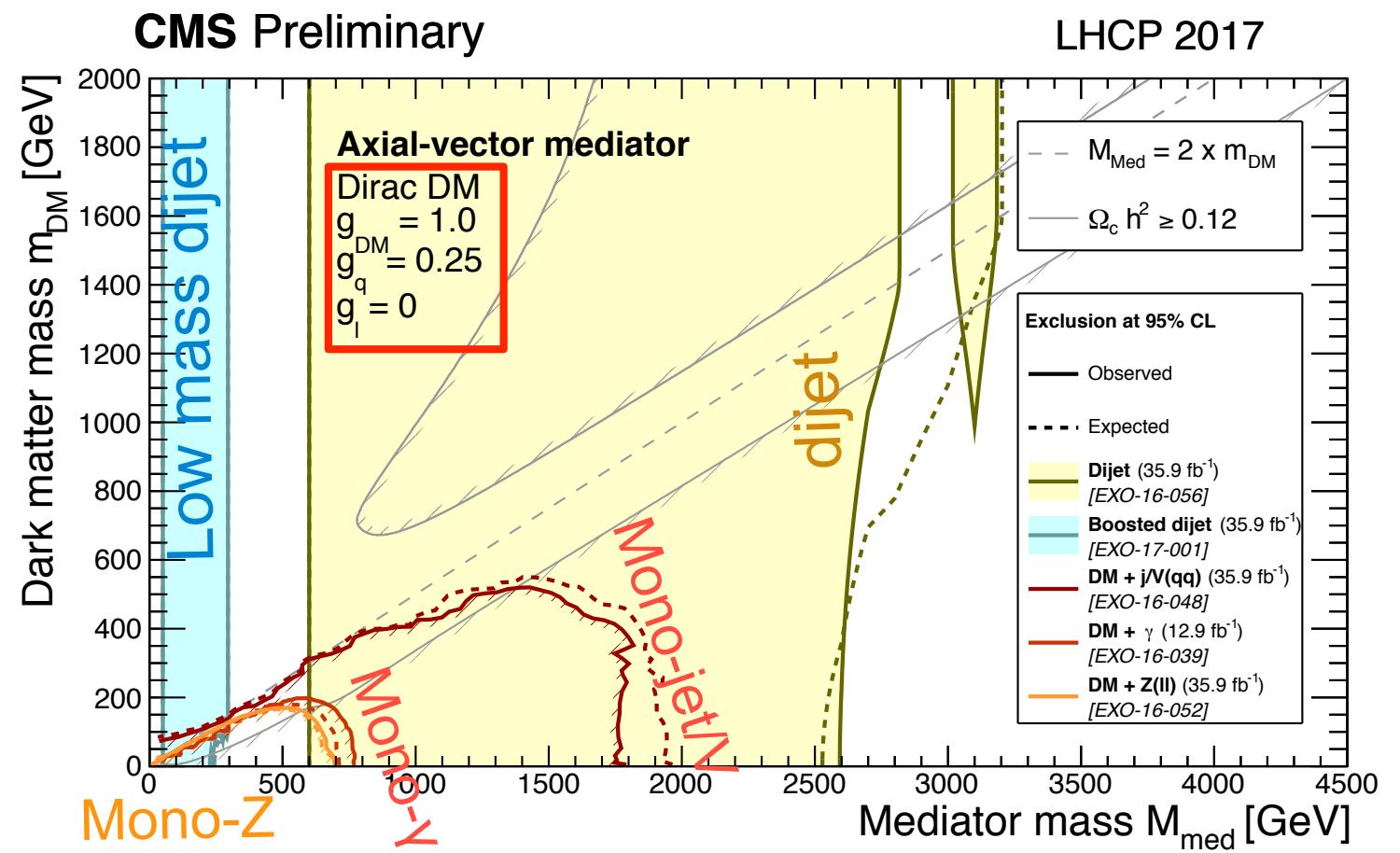
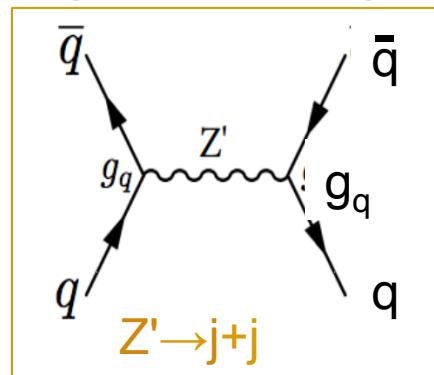
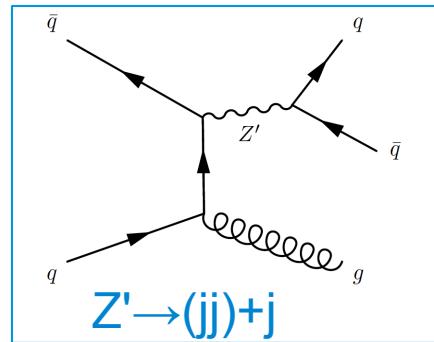
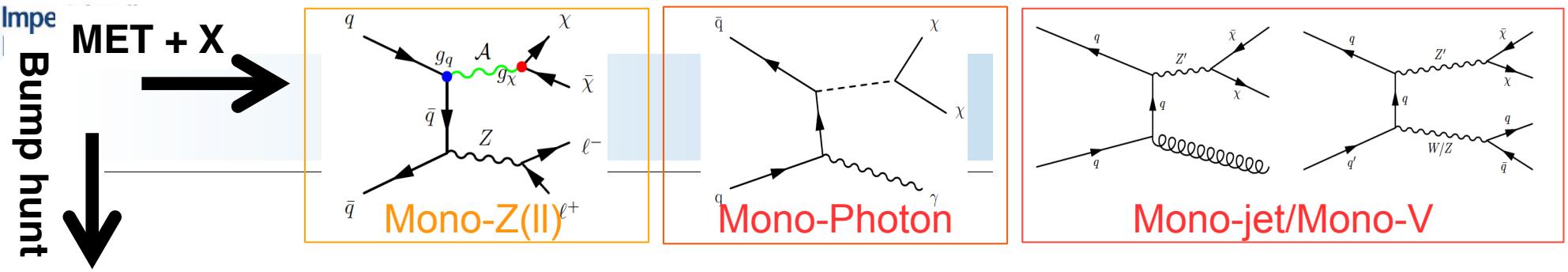
Main result of the interpretation of collider search in simplified model

Clearly state
Main assumptions
entering the result
- Mediator
- DM type
- Couplings

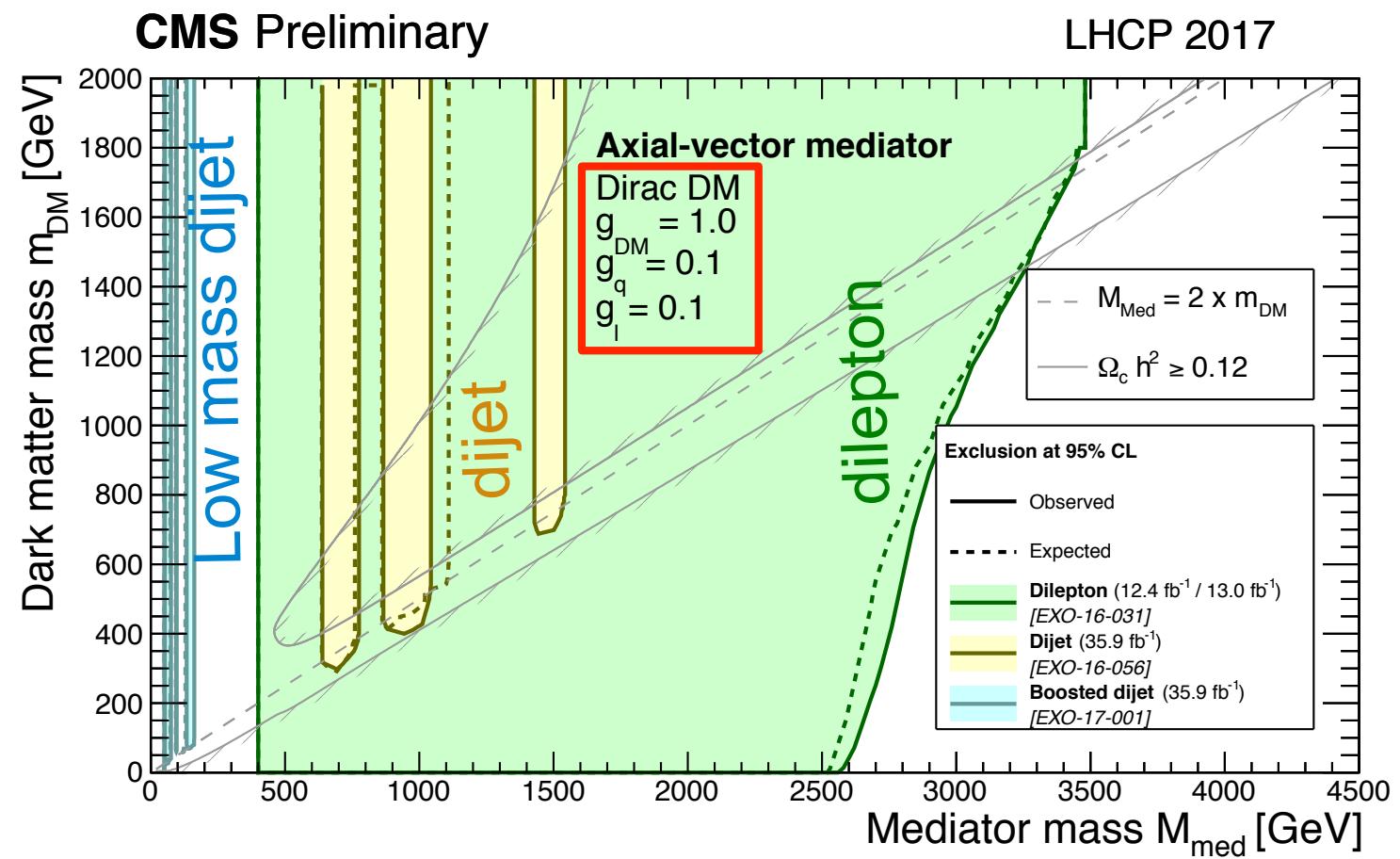
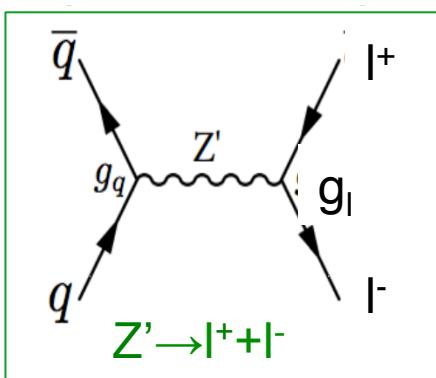
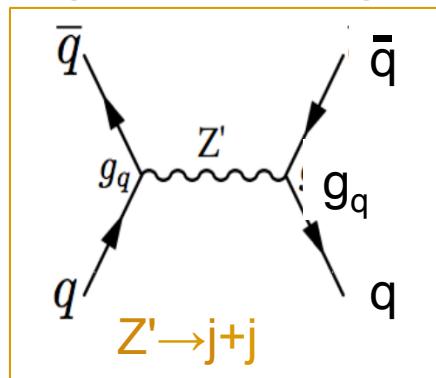
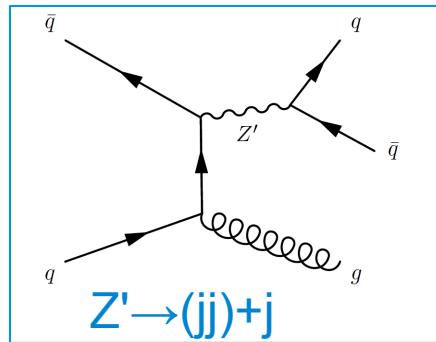
All based on LHC
DM WG recommendation
1603.04156



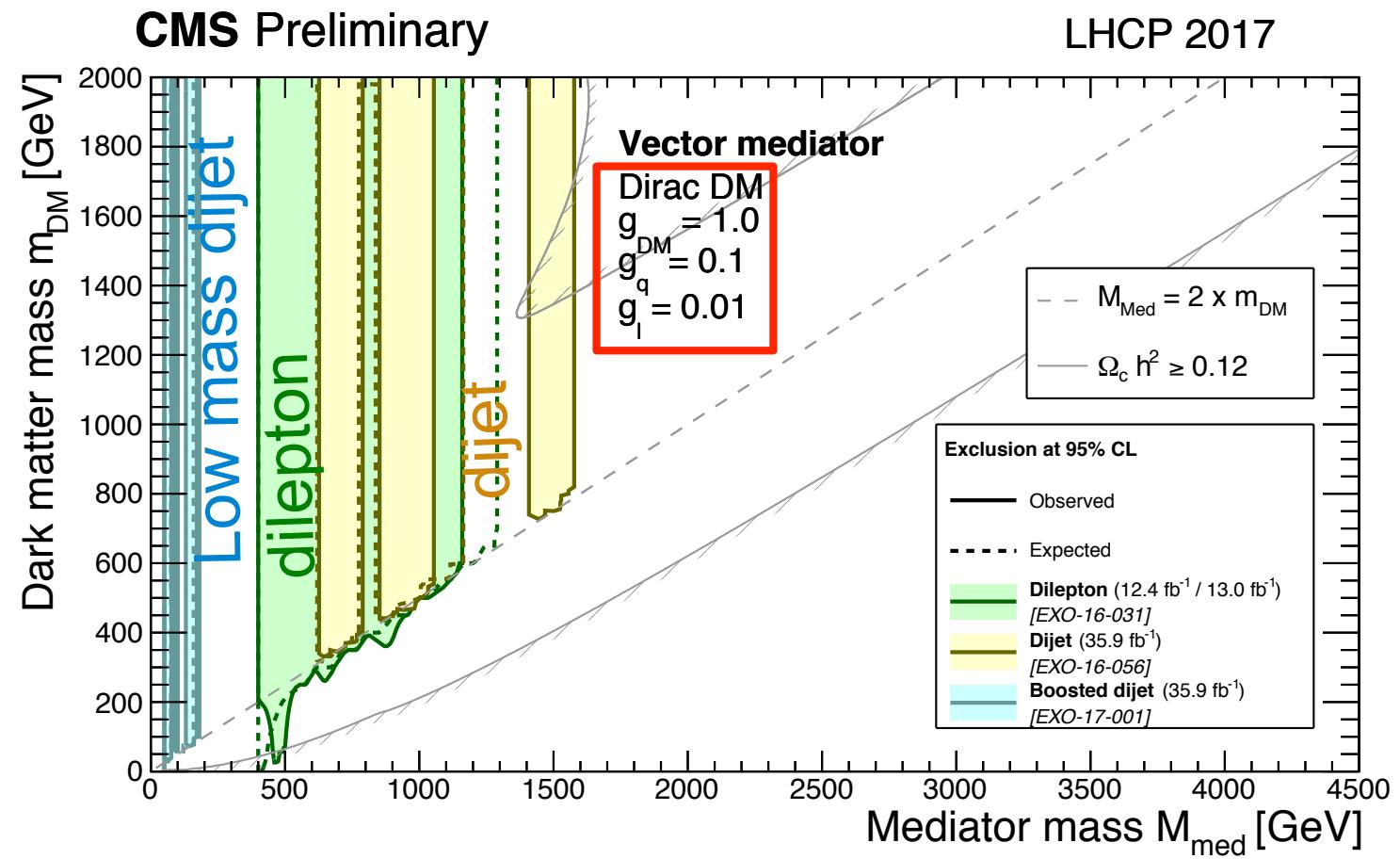
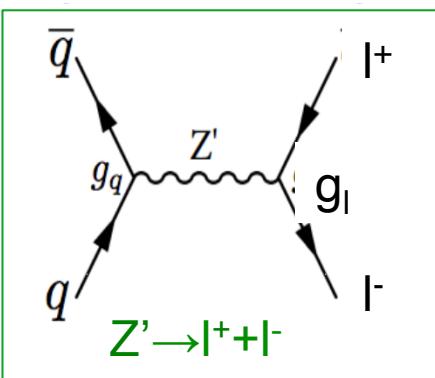
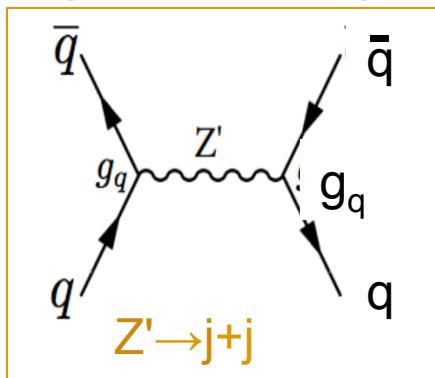
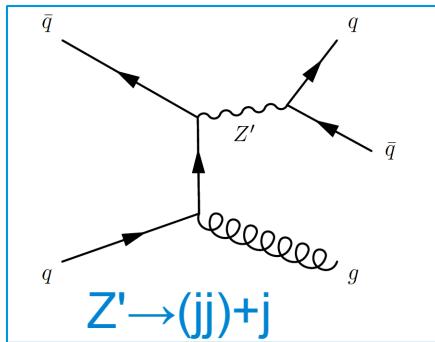
Usual “LHC limits”
For 95% CL [not 90%]
Indicate Relic density
line but do not use it
as “validity” requirement.
Its FI only.
[more caveats and
discussion are provided
in the report]



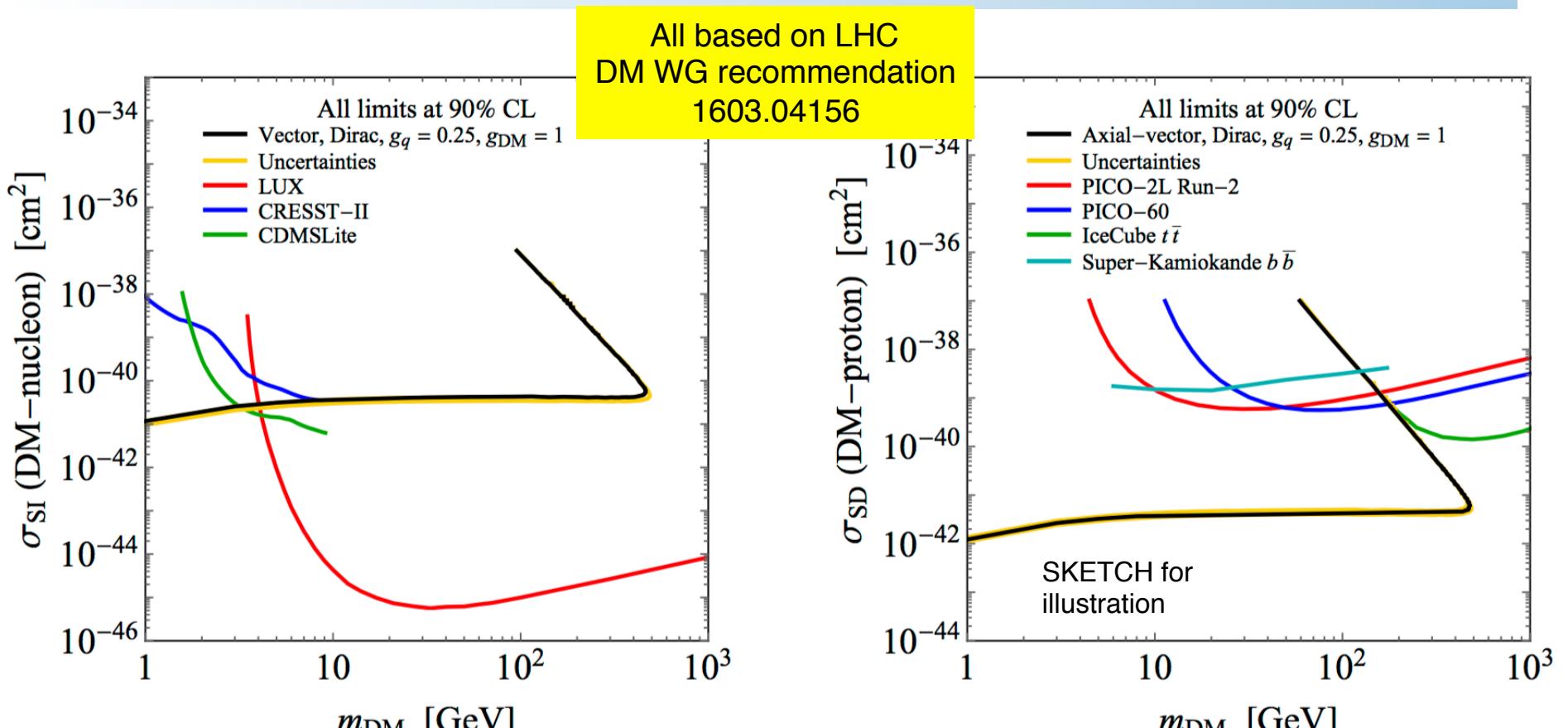
ATLAS very similar



ATLAS very similar



Comparison with Direct Detection



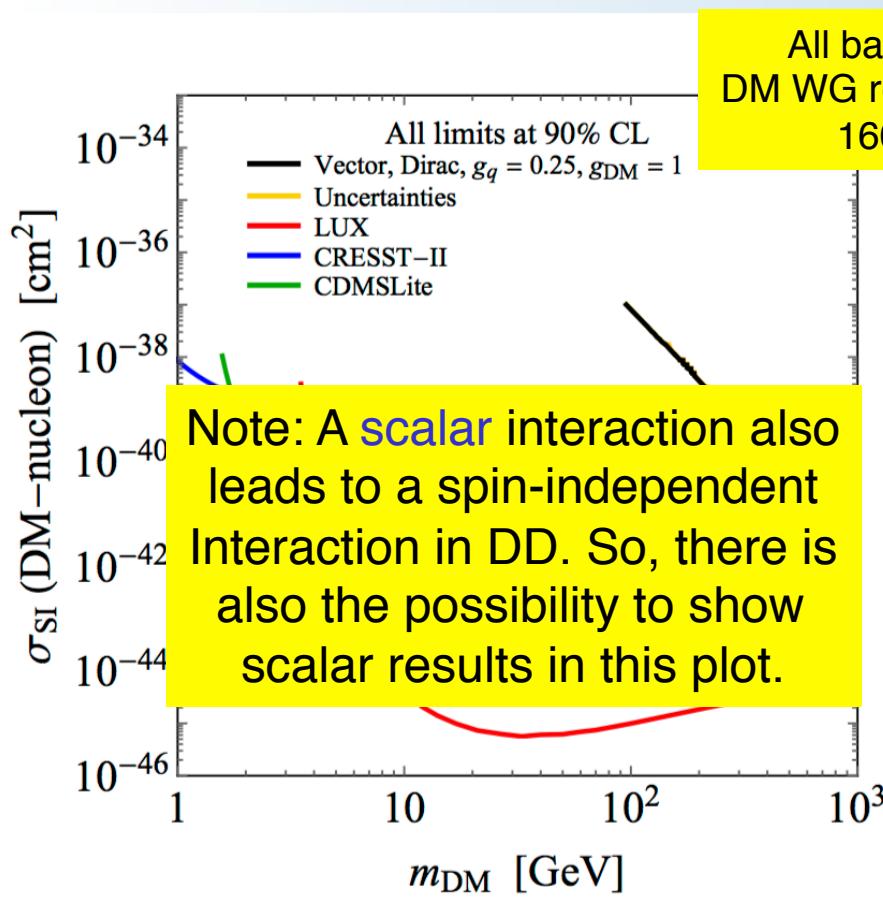
$$\sigma_{\text{SI}} = \frac{9 g_q^2 g_{\text{DM}}^2 \mu_{n\chi}^2}{\pi M_{\text{med}}^4}$$

Vector:

$$\approx 1.1 \times 10^{-39} \text{ cm}^2 \cdot \left(\frac{g_{\text{DM}} g_q}{1} \right)^2 \left(\frac{1 \text{ TeV}}{M_{\text{med}}} \right)^4 \left(\frac{\mu_{n\chi}}{1 \text{ GeV}} \right)^2 \quad \sigma^{\text{SD}} \simeq 2.4 \times 10^{-42} \text{ cm}^2 \cdot \left(\frac{g_q g_{\text{DM}}}{0.25} \right)^2 \left(\frac{1 \text{ TeV}}{M_{\text{med}}} \right)^4 \left(\frac{\mu_{n\chi}}{1 \text{ GeV}} \right)^2$$

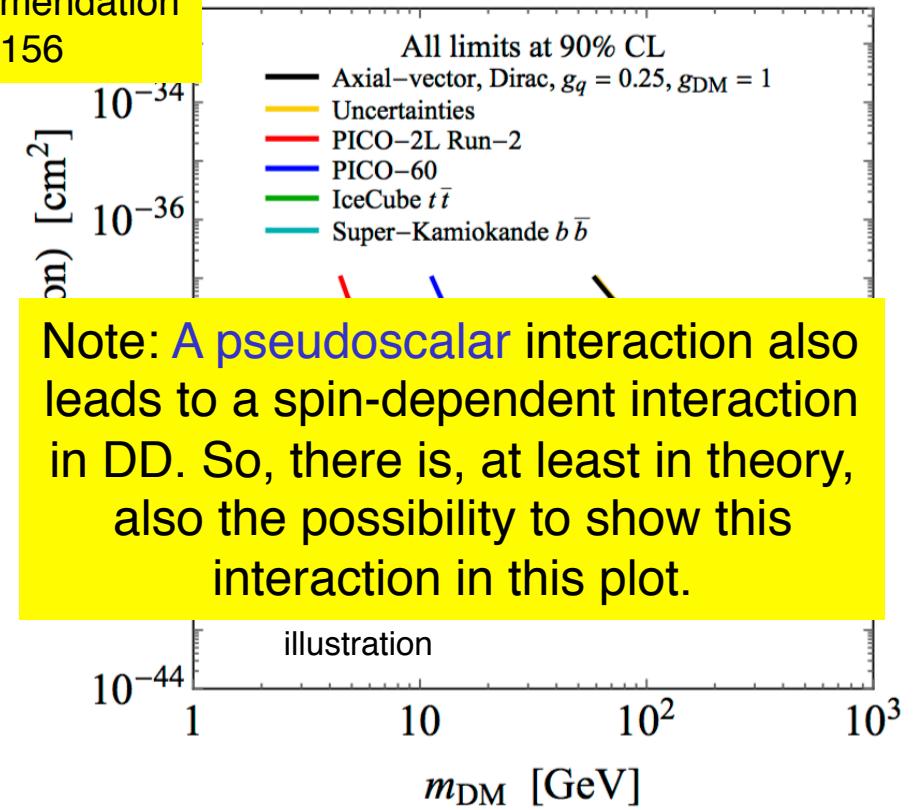
Provide simple formulas to perform the translation of the Mass-mass plane results into these planes. A full derivation of these formulas along with assumptions/caveat discussions is provided in the report.

Comparison with Direct Detection



Scalar:

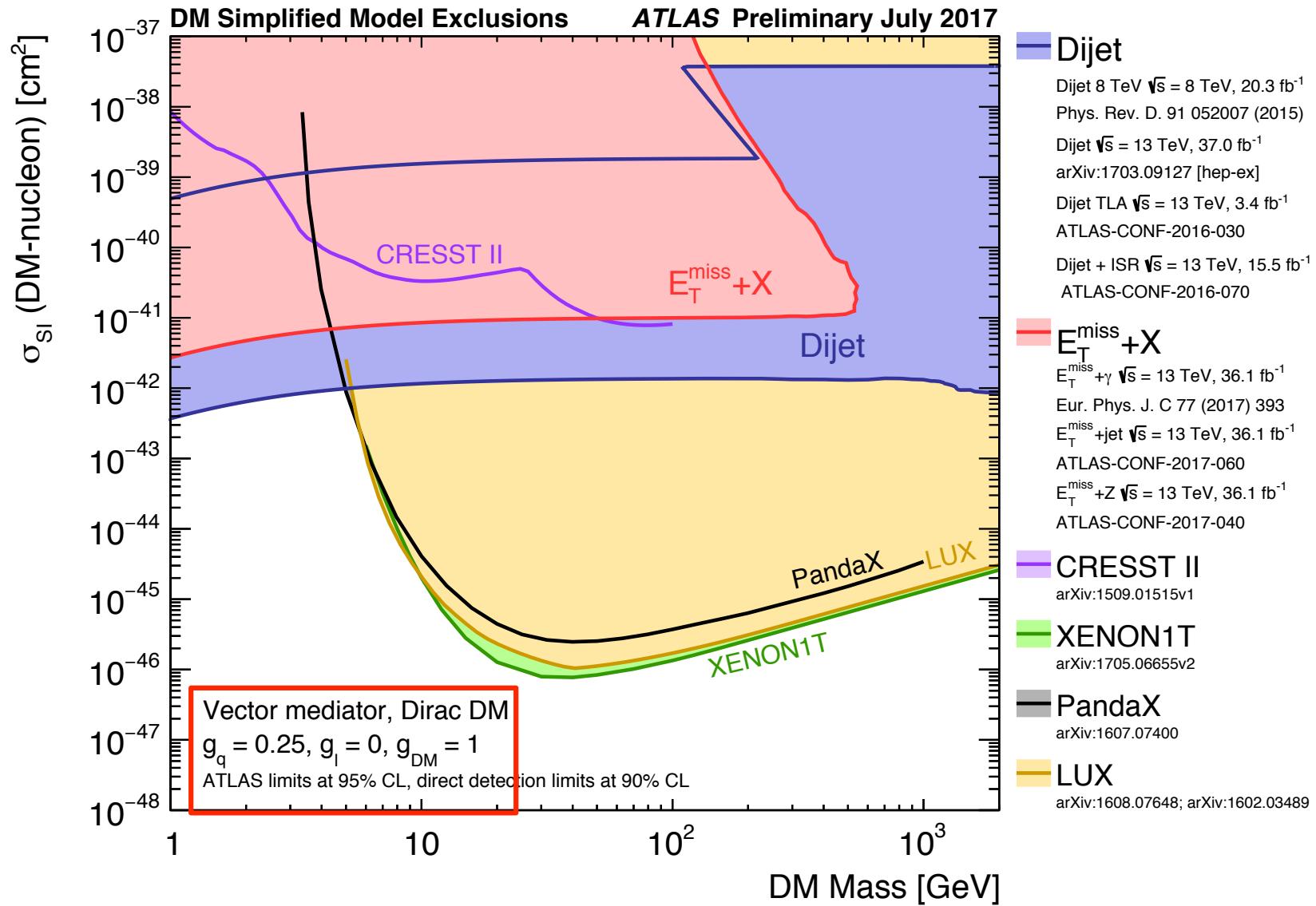
$$\sigma_{\text{SI}} \approx 6.9 \times 10^{-43} \text{ cm}^2 \cdot \left(\frac{g_{\text{DM}} g_q}{1} \right)^2 \left(\frac{125 \text{ GeV}}{M_{\text{med}}} \right)^4 \left(\frac{\mu_{n\chi}}{1 \text{ GeV}} \right)^2$$



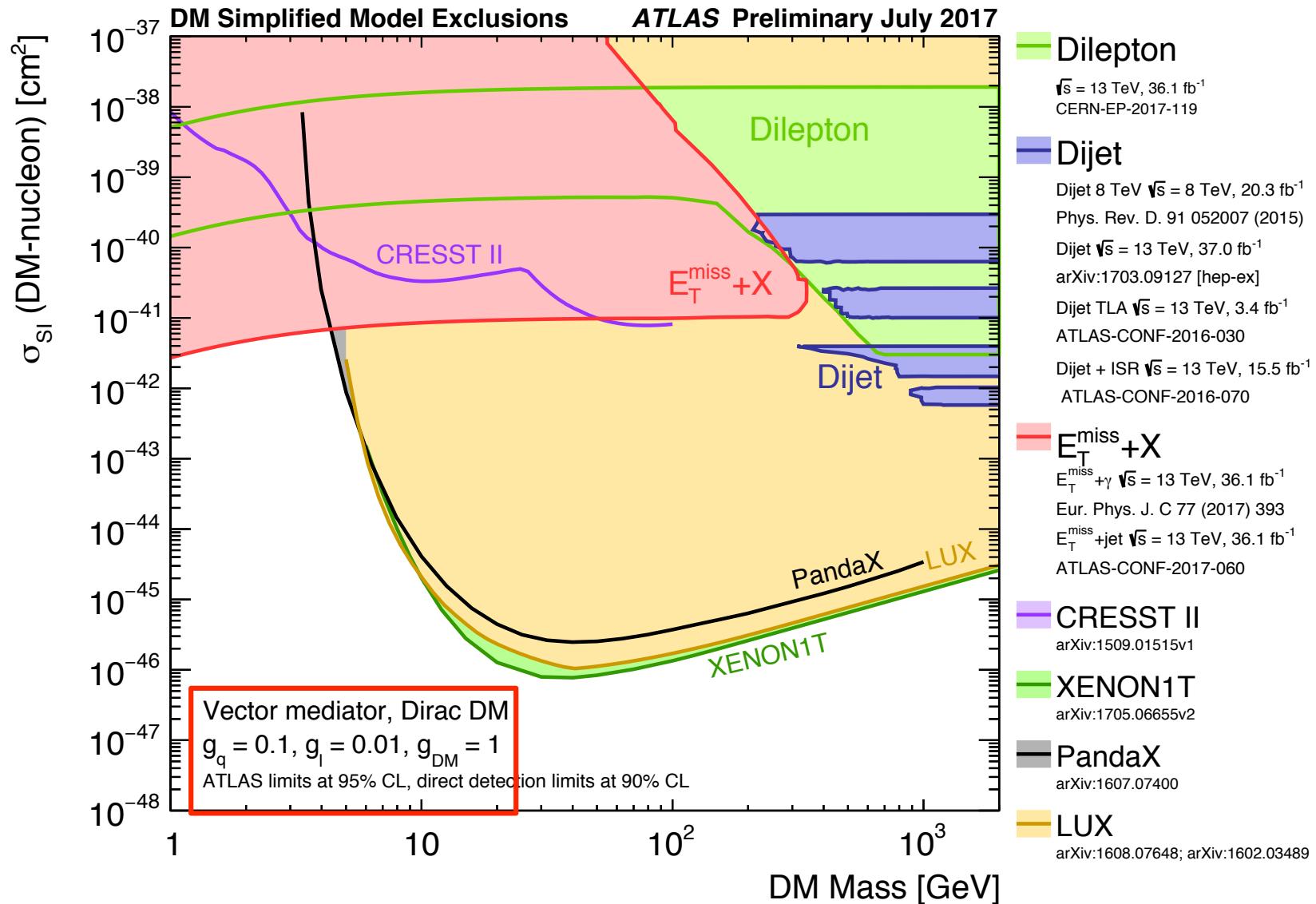
Pseudoscalar
see next slides

Provide simple formulas to perform the translation of the Mass-mass plane results into these planes. A full derivation of these formulas along with assumptions/caveat discussions is provided in the report.

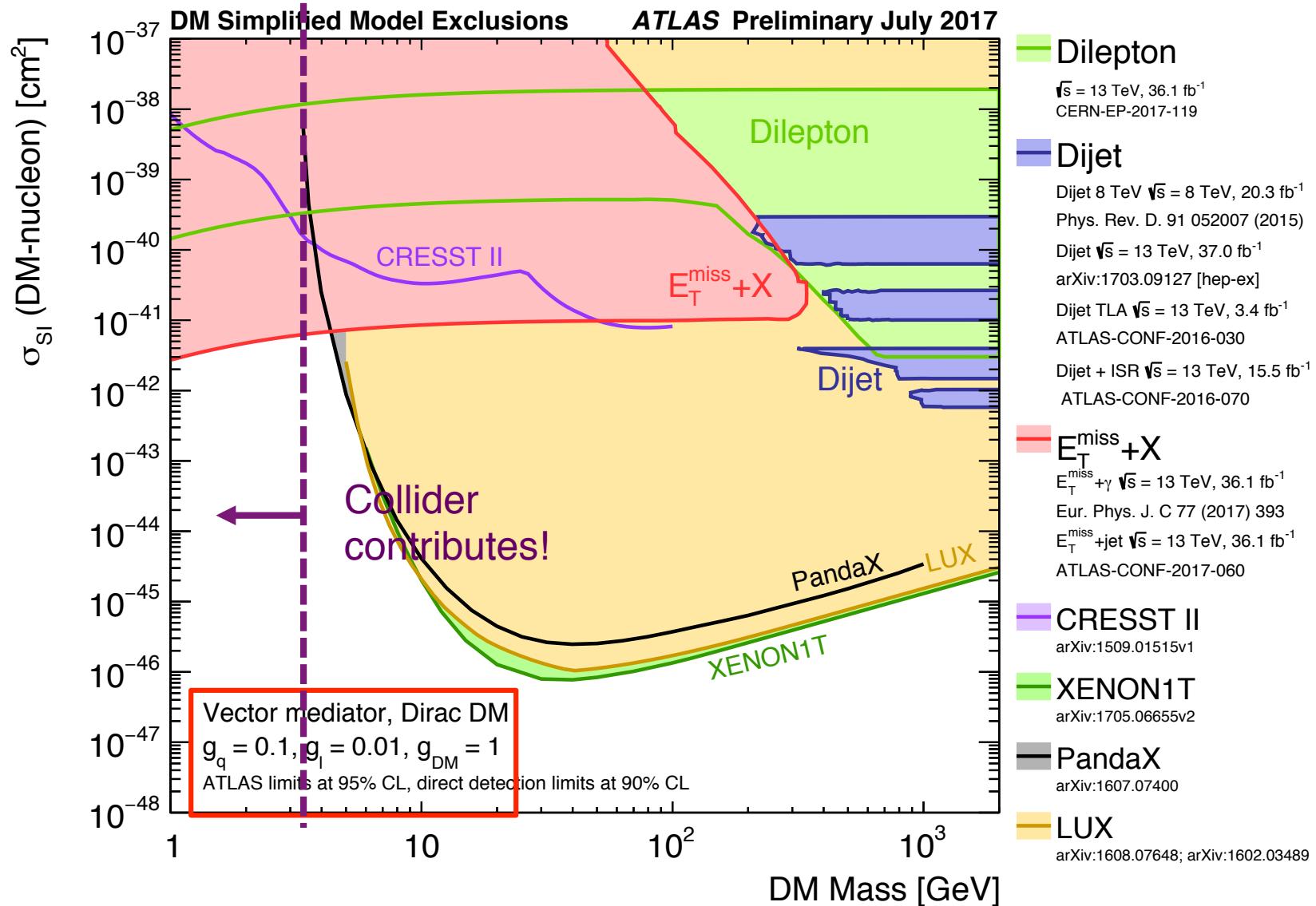
Comparison with Direct Detection: Vector case



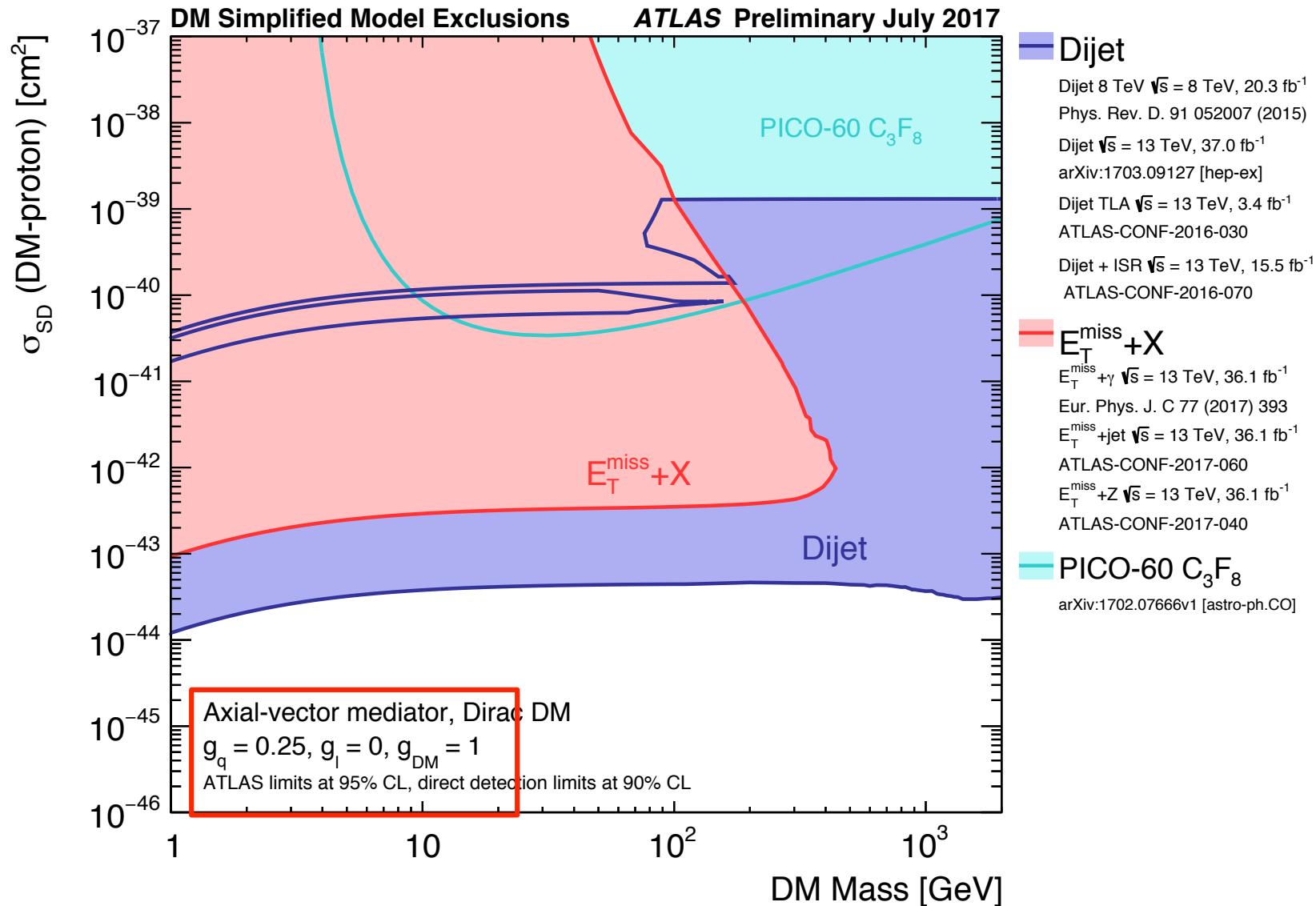
Comparison with Direct Detection: Vector Case



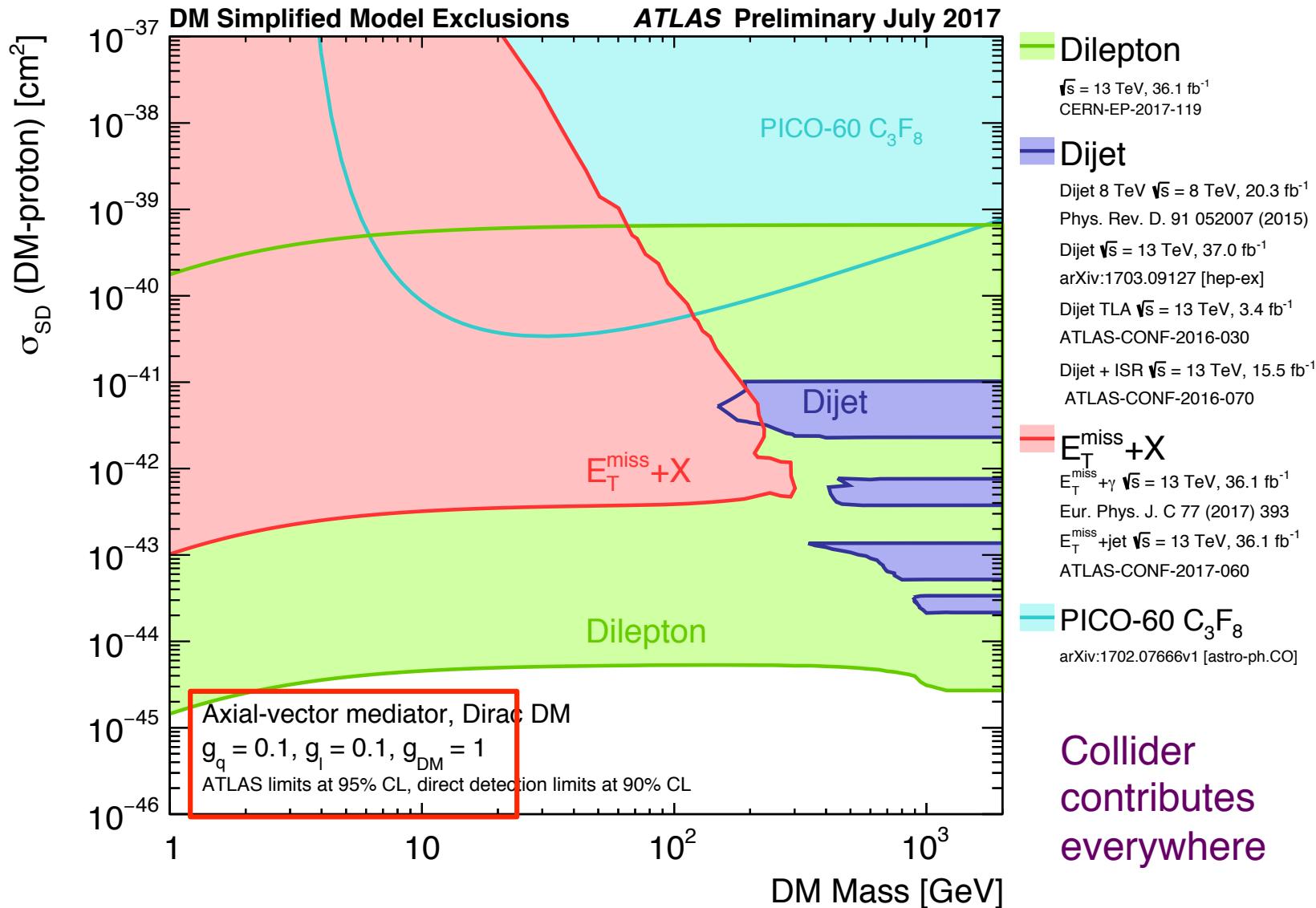
Comparison with Direct Detection: Vector Case



Comparison with Direct Detection: Axial-Vector



Comparison with Direct Detection: Axial-Vector



Scalar and Pseudoscalar

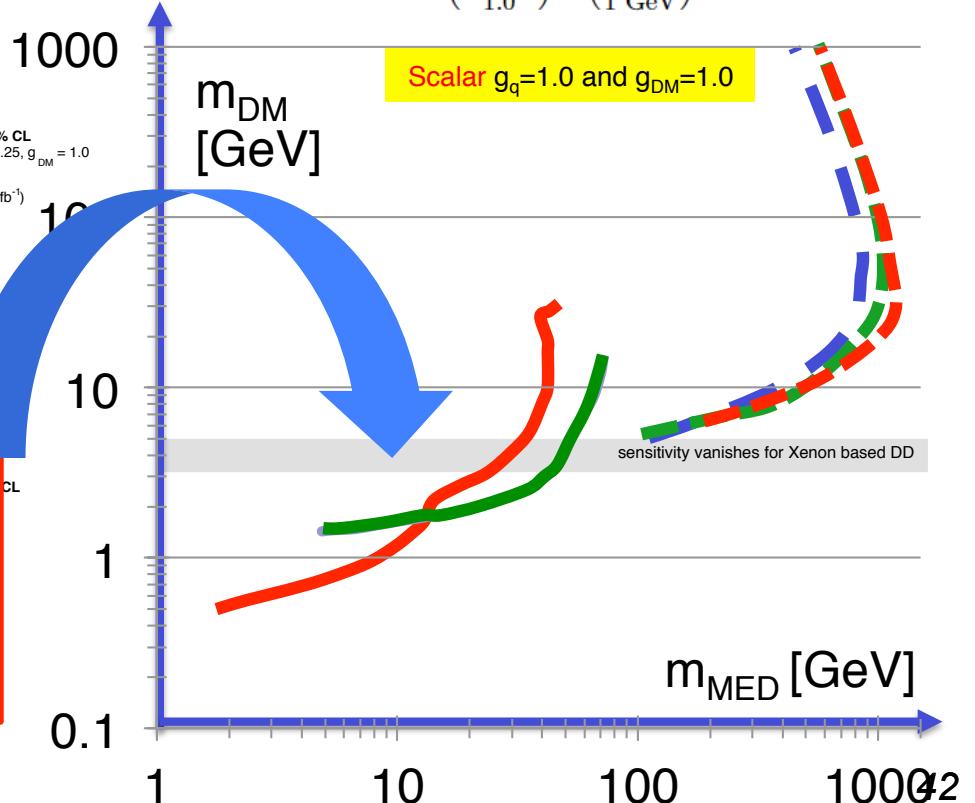
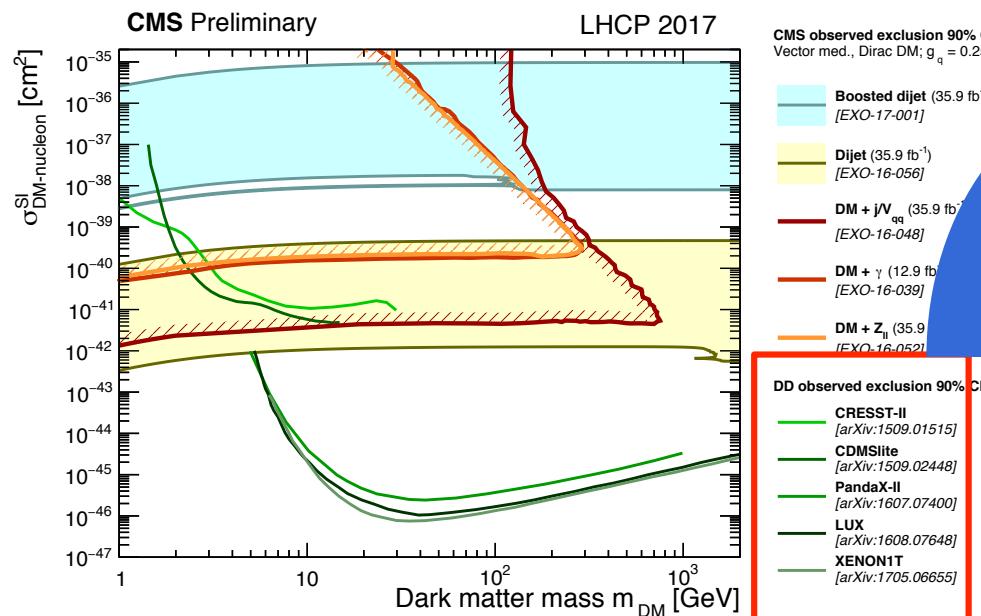
Of course, there is also a well-defined way to compare DD limits with collider results in the collider language: mass-mass plane:

Based on
arXiv:1407.8257

Numerical calculation
provide by. C. McCabe

Scalar simplified model:
 $g_q=1.0$ and $g_{DM}=1.0$

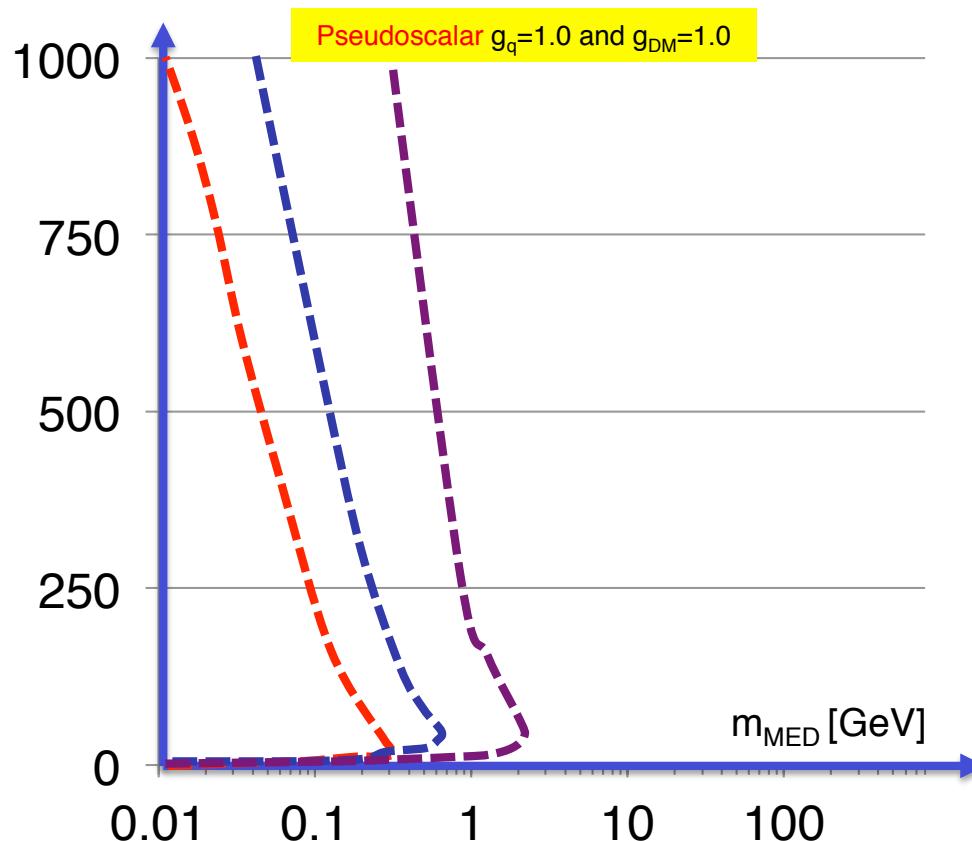
$$M_\phi = 125 \text{ GeV} \left(\frac{6.9 \times 10^{-43} \text{ cm}^2}{\sigma^{\text{SI}}} \right)^{0.25} \cdot \left(\frac{g_{DM} g_q}{1.0} \right)^{0.5} \left(\frac{\mu_n}{1 \text{ GeV}} \right)^{0.5}.$$



Scalar and Pseudoscalar

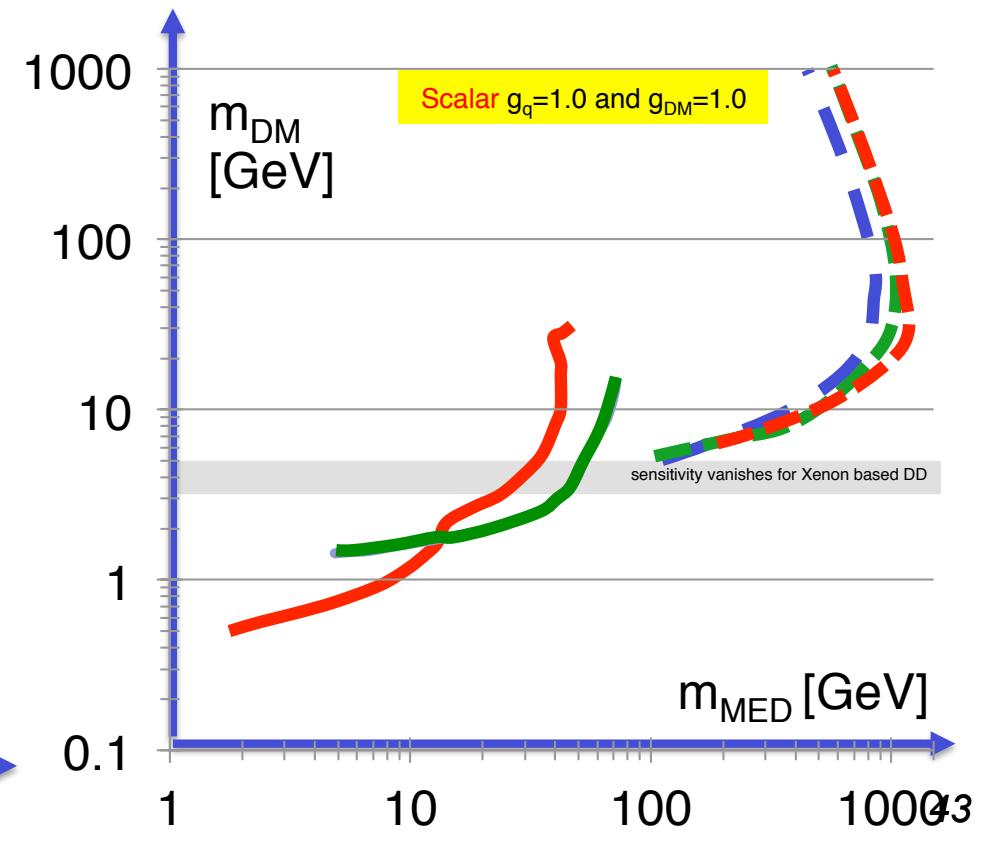
Of course, there is also a well-defined way to compare DD limits with collider results in the collider language: mass-mass plane:

Pseudoscalar (PS)
simplified model:
 $g_q=1.0$ and $g_{DM}=1.0$



Based on
arXiv:1407.8257
Numerical calculation
provide by. C. McCabe

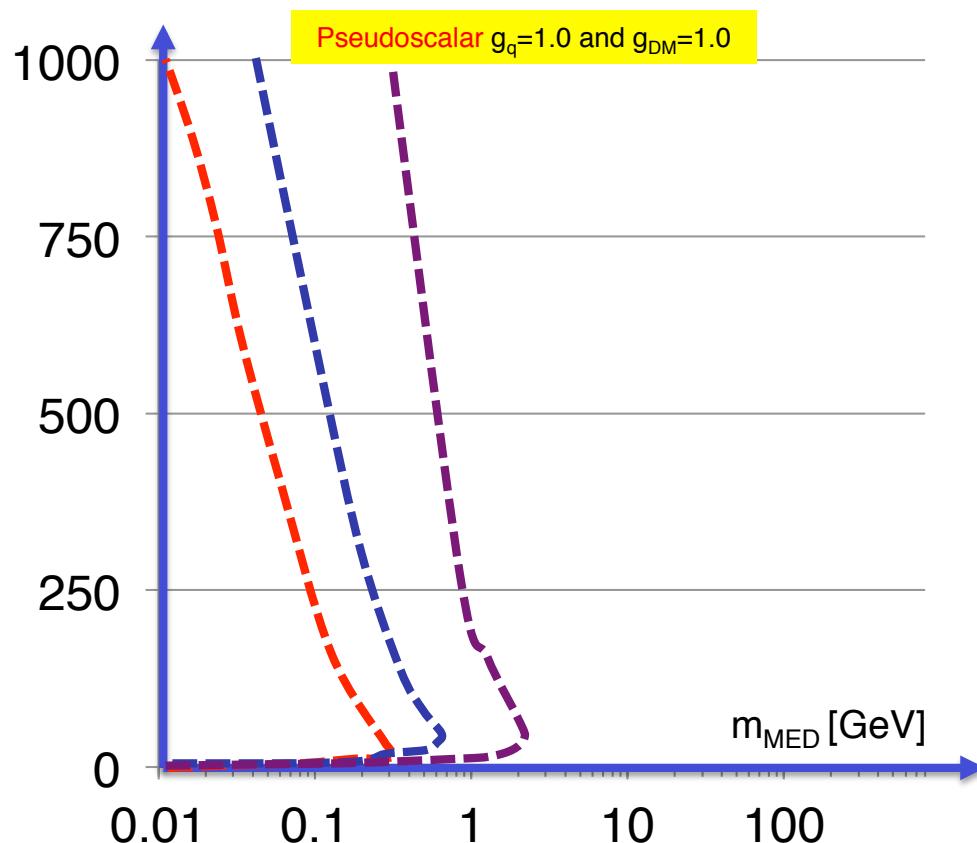
Scalar simplified model:
 $g_q=1.0$ and $g_{DM}=1.0$



DD limits in collider plane: σ_{SI}

Pseudoscalar case

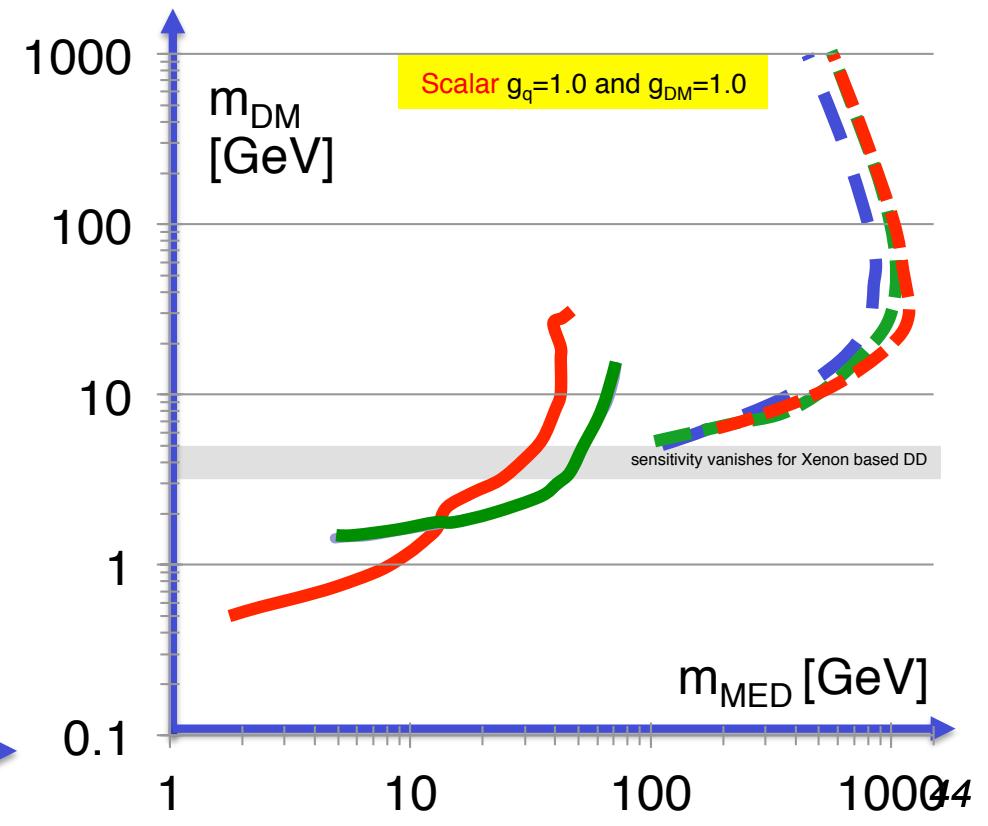
PICO60 ---
 [arXiv:1702.07666]
 XENON1T ----
 [arXiv:1705.06655]
 LZ/XENONnT ----
 15 t-yr projection



Scalar case

CRESST-II —
 [arXiv:1509.01515]
 CDMSlite —
 [arXiv:1509.02448]

PandaX-II ----
 [arXiv:1607.07400]
 LUX ----
 [arXiv:1608.07648]
 XENON1T ----
 [arXiv:1705.06655]

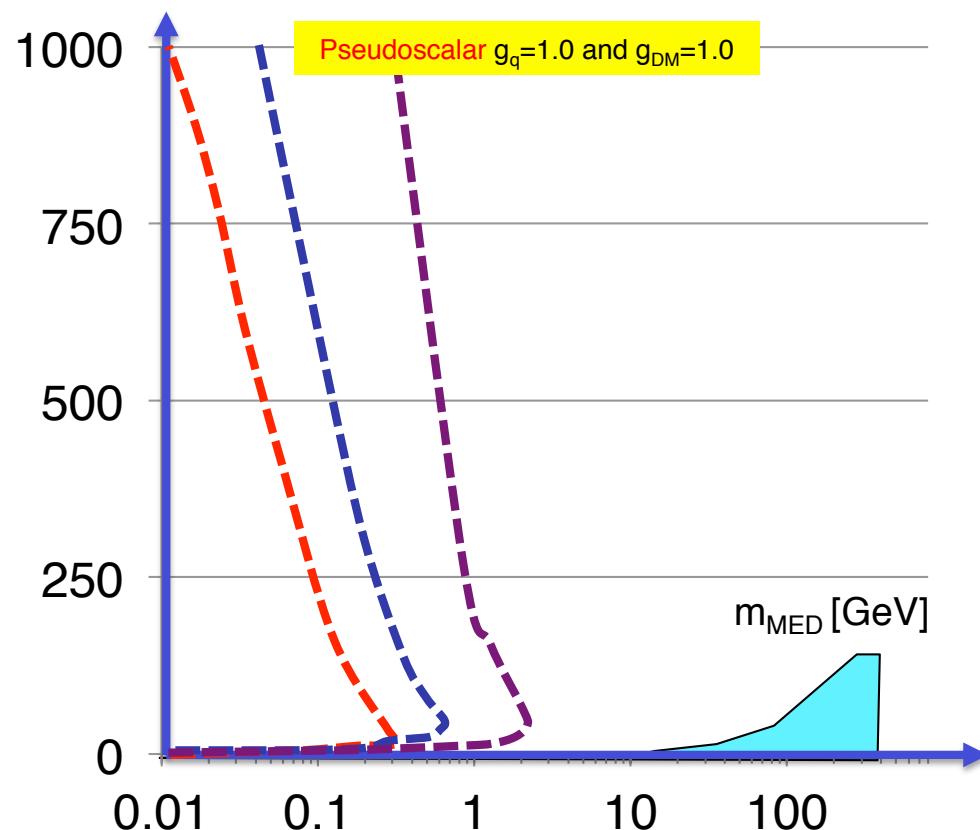


DD limits in collider plane: σ_{SI}

Pseudoscalar case

PICO60 ---
 [arXiv:1702.07666]
 XENON1T ----
 [arXiv:1705.06655]
 LZ/XENONnT ----
 15 t-yr projection

Mono-jet ■
 [EXO-16-048]



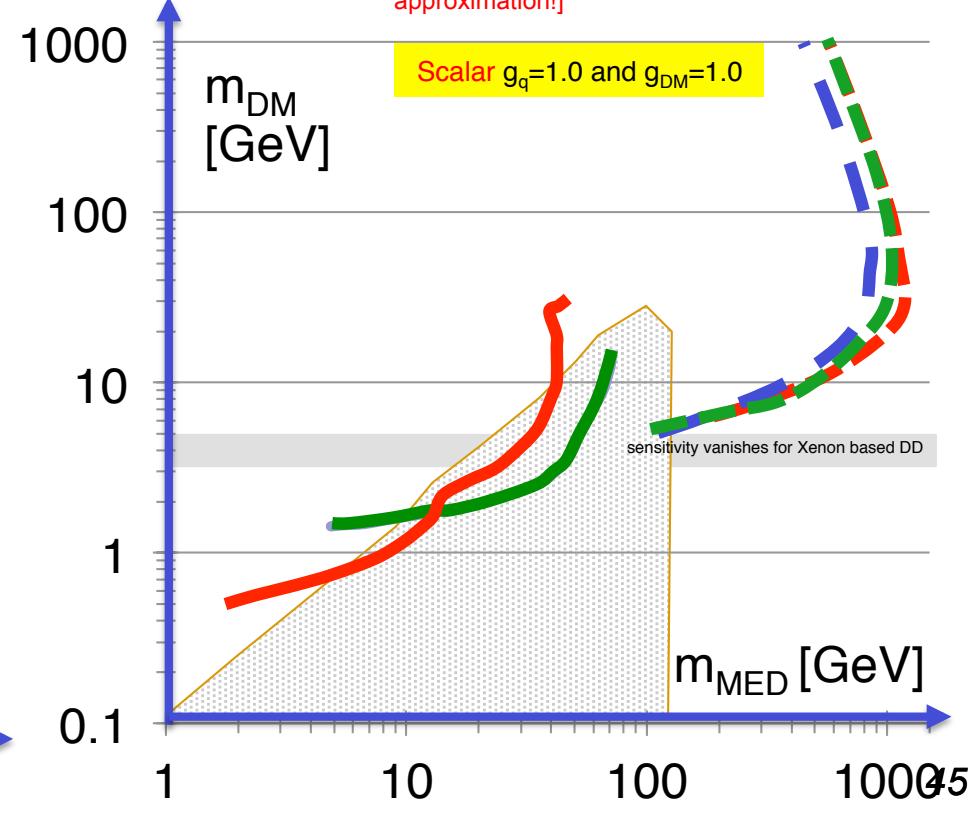
Scalar case

CRESST-II —
 [arXiv:1509.01515]
 CDMSlite —
 [arXiv:1509.02448]

2L + MET □
 [SUS-17-001]

[Note: private
estimate of 2D
exclusion –
only rough
approximation!]

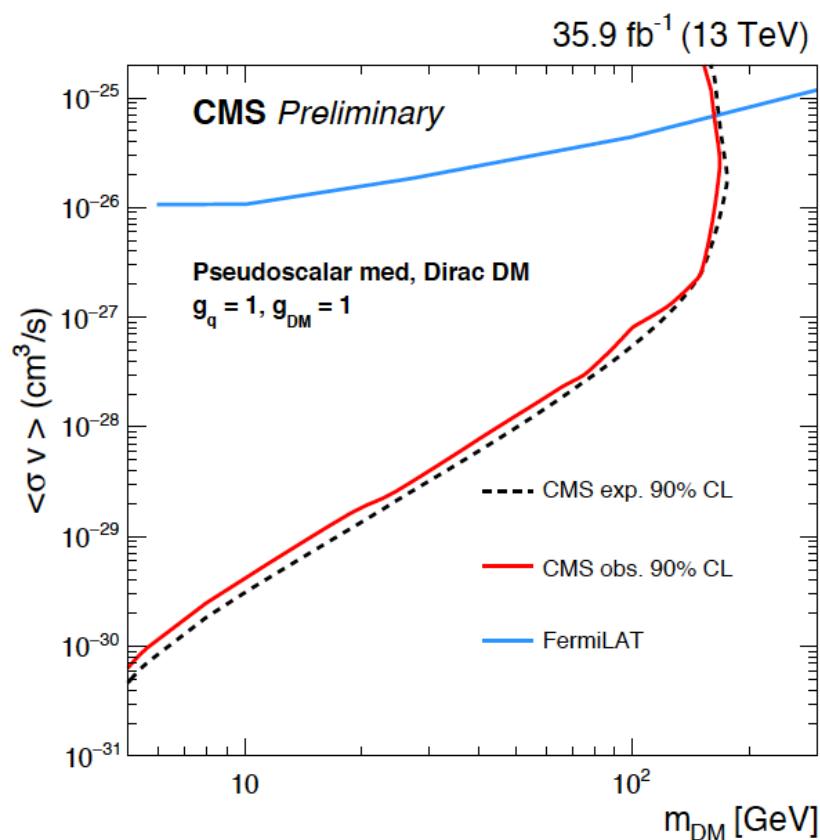
PandaX-II ----
 [arXiv:1607.07400]
 LUX ----
 [arXiv:1608.07648]
 XENON1T ----
 [arXiv:1705.06655]



Comparing with Indirect Detection

$$\langle \sigma v_{\text{rel}} \rangle_q = \frac{3m_q^2}{2\pi v^2} \frac{g_q^2 g_{\text{DM}}^2 m_{\text{DM}}^2}{(M_{\text{med}}^2 - 4m_{\text{DM}}^2)^2 + M_{\text{med}}^2 \Gamma_{\text{med}}^2} \sqrt{1 - \frac{m_q^2}{m_{\text{DM}}^2}}$$

All based on LHC
DM WG recommendation
1603.04156



Due to additional velocity suppression DD experiments, have sensitivity for $M_{\text{med}} >$ few GeV, but indirect detection can provide further constrains on pseudoscalar (PS) interactions.

Hence, compare collider also with ID

Mono-jet — [EXO-16-048]
FermiLAT — arXiv:1503.02641

DD and Collider Complementarity in nutshell ...

Wimp – Nucleon Interaction	
<u>Spin-Independent (SI)</u>	<u>Spin-Dependent (SD)</u>
Basic Mediators	
<u>Vector</u> <i>Besides low DM masses DD provides best sensitivity. Complementarity at low DM masses (<5 GeV)!</i>	<u>Axial-vector</u> <i>DD and collider are equal in overall sensitivity but probe different regions of parameter space! Complementarity in full parameter space!</i>
<u>Scalar</u> <i>Besides low DM masses DD provides best sensitivity. Complementarity at low DM masses (<5 GeV)!</i>	<u>Pseudoscalar</u> <i>Effectively no limits from DD above a few GeV in M_{med}, Collider and ID probe region at larger M_{med}. Complementarity in M_{med}!</i>

Summary

- The LHC experiments have established an impressive variety of very powerful direct searches that can be linked to DM production!
 - The traditional SUSY searches are complemented by mono-X analyses.
- The challenge is now to find a good balance between simplicity and complexity for the DM interpretations of these searches.
 - We have started to outline an interpretation programme that uses simplified DM models.
 - Today, the used simplified models are still very basic and therefore interpretations come along with several assumptions and some caveats – this will evolve with time!
 - Our goal is to establish the “big picture” in order to understand if/where our search strategy might have weak spots or even holes and this also requires appropriate interpretations of the searches and a MEANIGFUL comparison with other experiments.
- We have still almost two decades of data taking in front of us, with a factor 100 increase of statistic still to come!

The story continues ...

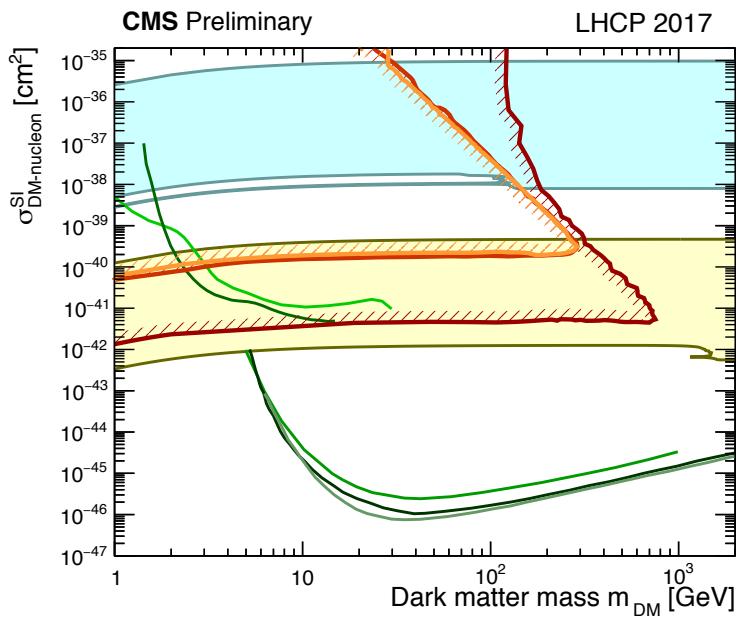
BACKUP

DD limits in collider plane: σ_{SI}

Of course, there is also a well-defined way to compare DD limits with collider results in the collider language: mass-mass plane:

Scalar simplified model:
 $g_q=1.0$ and $g_{\text{DM}}=1.0$

$$M_\phi = 125 \text{ GeV} \left(\frac{6.9 \times 10^{-43} \text{ cm}^2}{\sigma_n^{\text{SI}}} \right)^{0.25} \cdot \left(\frac{g_{\text{DM}} g_q}{1.0} \right)^{0.5} \left(\frac{\mu_n}{1 \text{ GeV}} \right)^{0.5}.$$

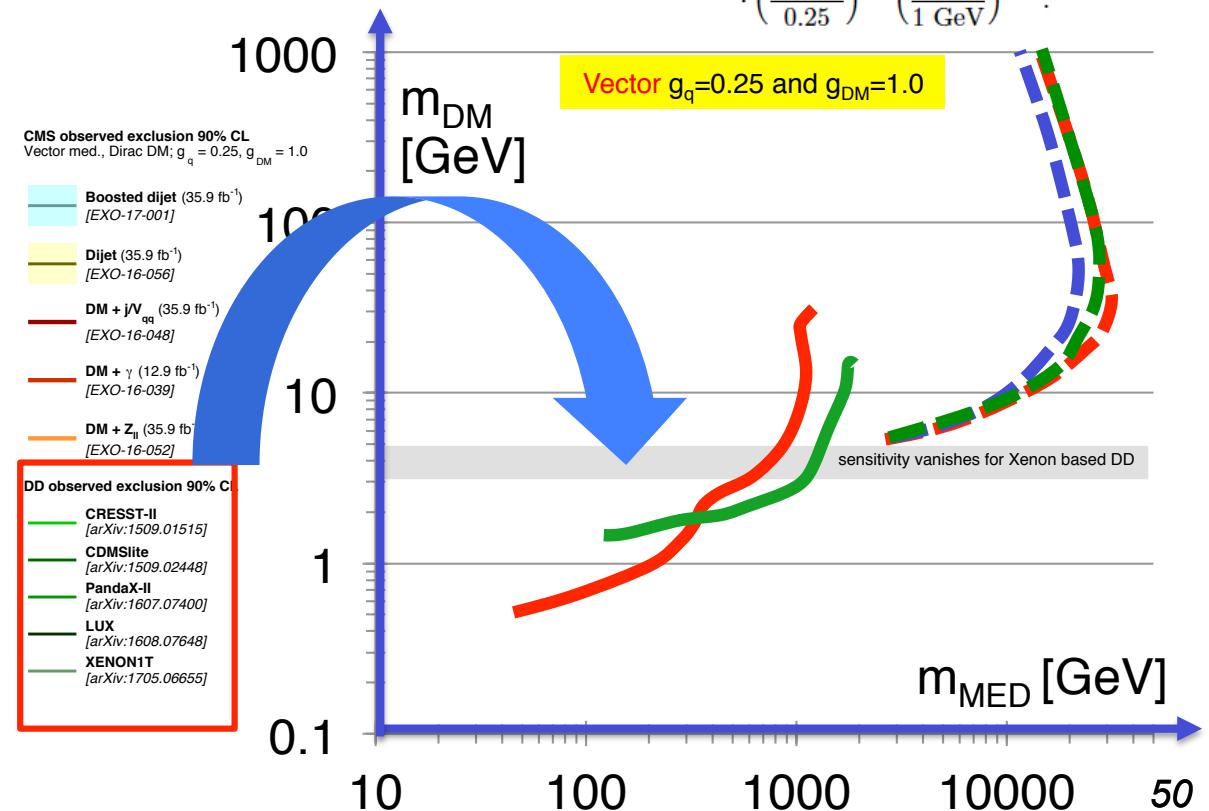


Based on
arXiv:1407.8257

Numerical calculation
provide by C. McCabe

Vector simplified model:
 $g_q=0.25$ and $g_{\text{DM}}=1.0$

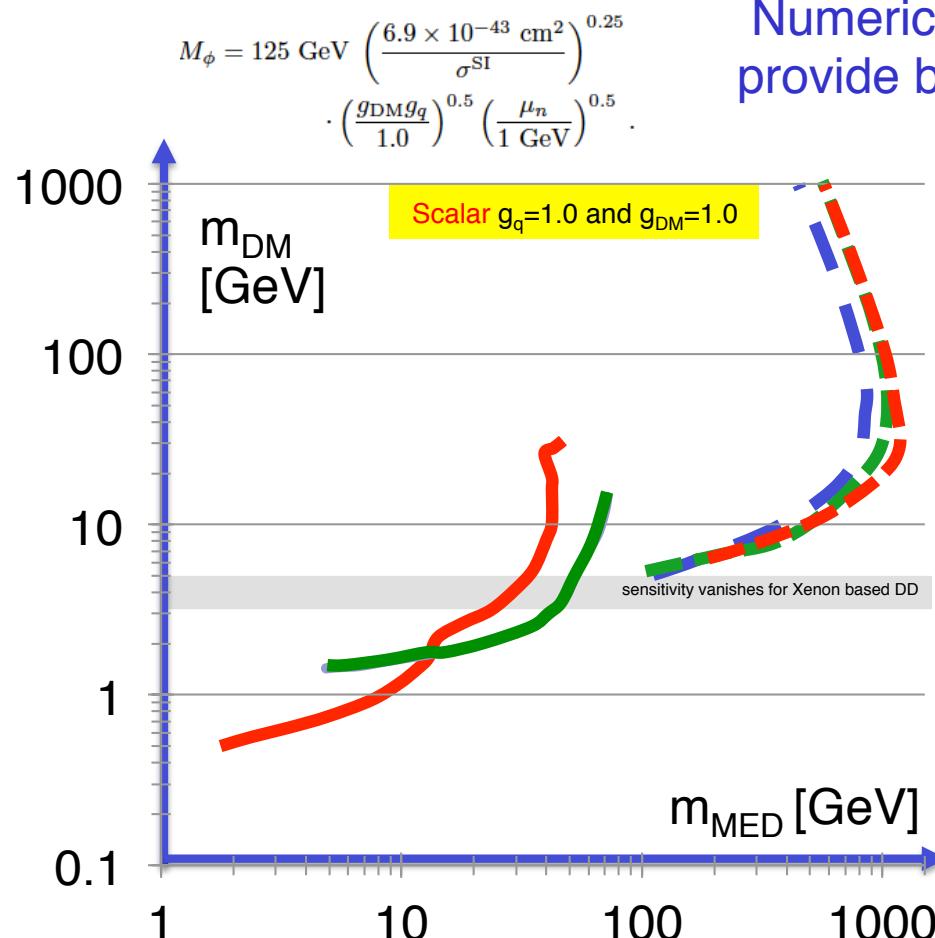
$$M_A = 1 \text{ TeV} \left(\frac{6.971 \times 10^{-41} \text{ cm}^2}{\sigma_n^{\text{SI}}} \right)^{0.25} \cdot \left(\frac{g_{\text{DM}} g_q}{0.25} \right)^{0.5} \left(\frac{\mu_n}{1 \text{ GeV}} \right)^{0.5}.$$



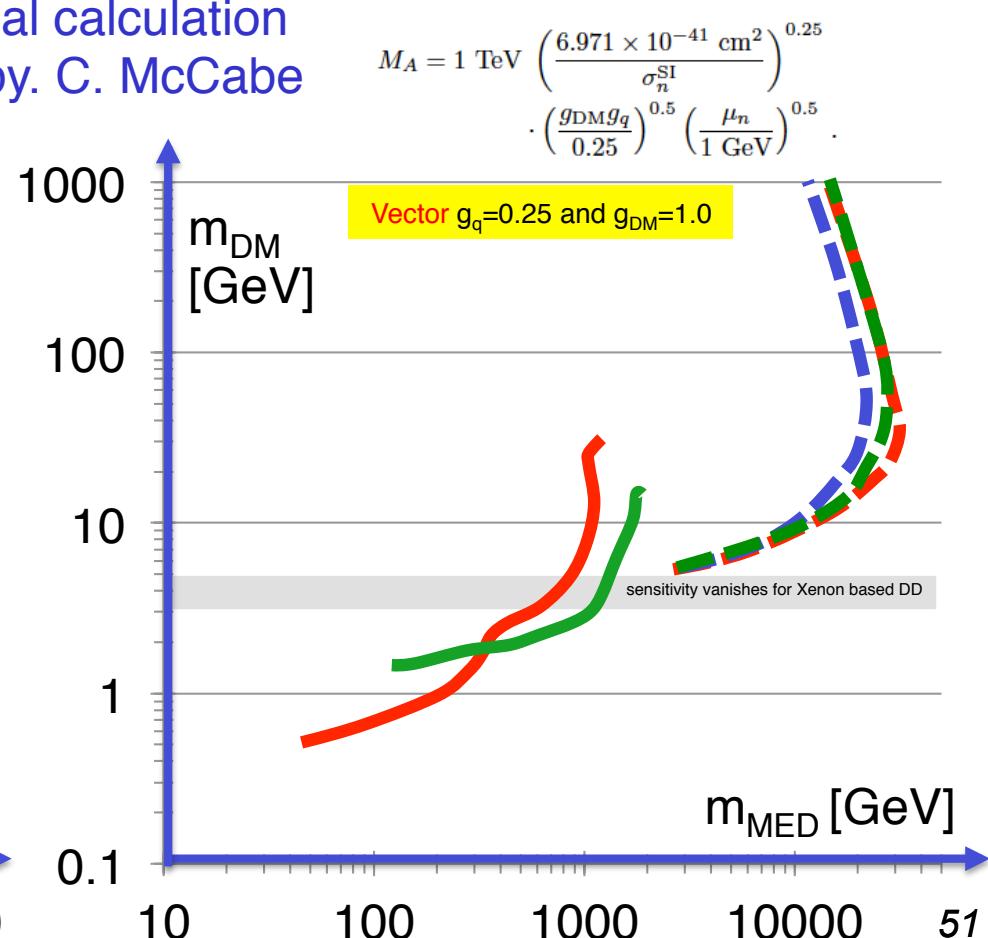
DD limits in collider plane: σ_{SI}

Of course, there is also a well-defined way to compare DD limits with collider results in the collider language: mass-mass plane:

Scalar simplified model:
 $g_q=1.0$ and $g_{\text{DM}}=1.0$



Vector simplified model:
 $g_q=0.25$ and $g_{\text{DM}}=1.0$

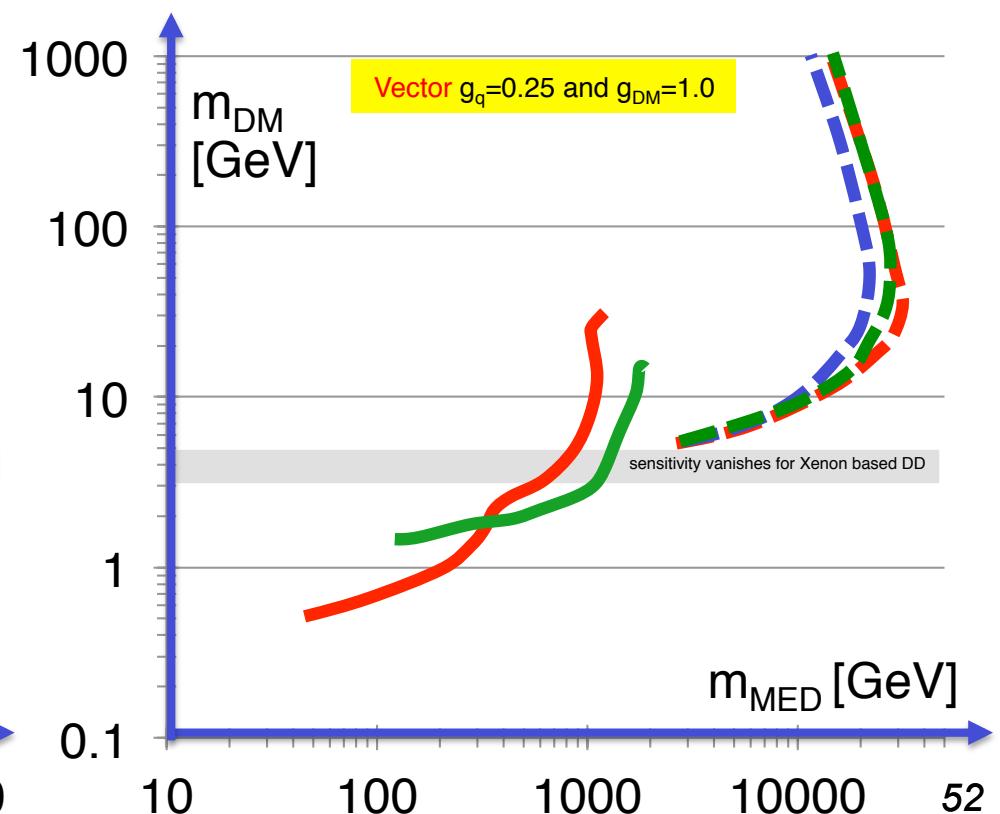
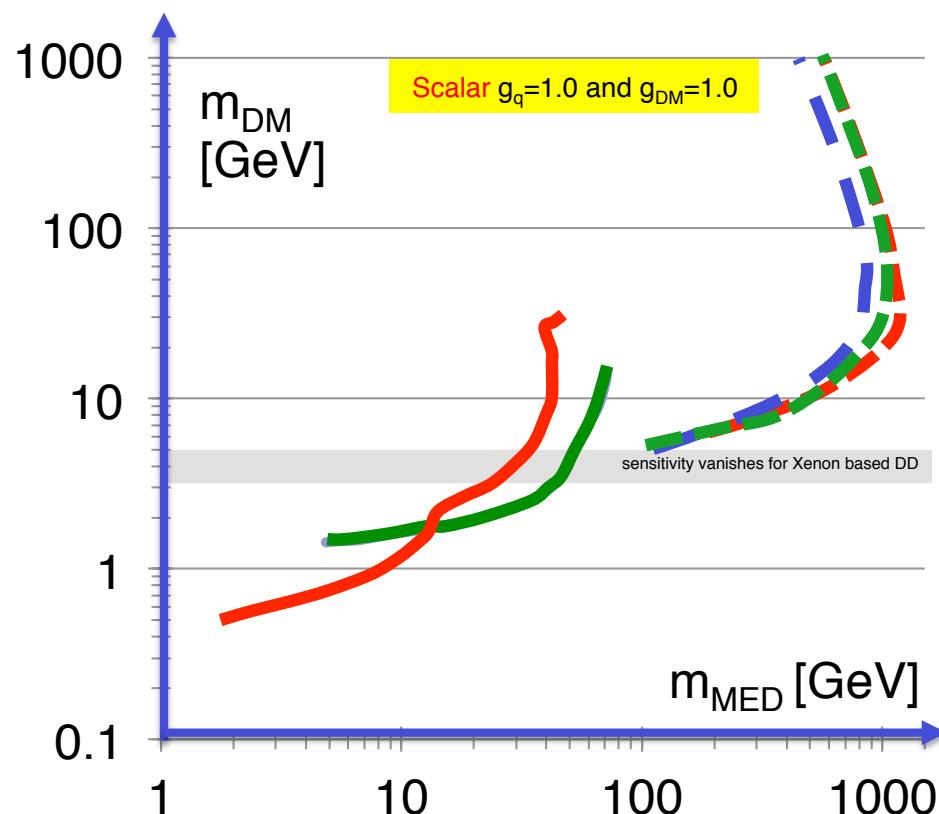


DD limits in collider plane: σ_{SI}

DD Experiments

CRESST-II — [arXiv:1509.01515]
CDMSlite — [arXiv:1509.02448]

PandaX-II ----- [arXiv:1607.07400]
LUX ----- [arXiv:1608.07648]
XENON1T ----- [arXiv:1705.06655]



DD limits in collider plane: σ_{SI}

DD Experiments

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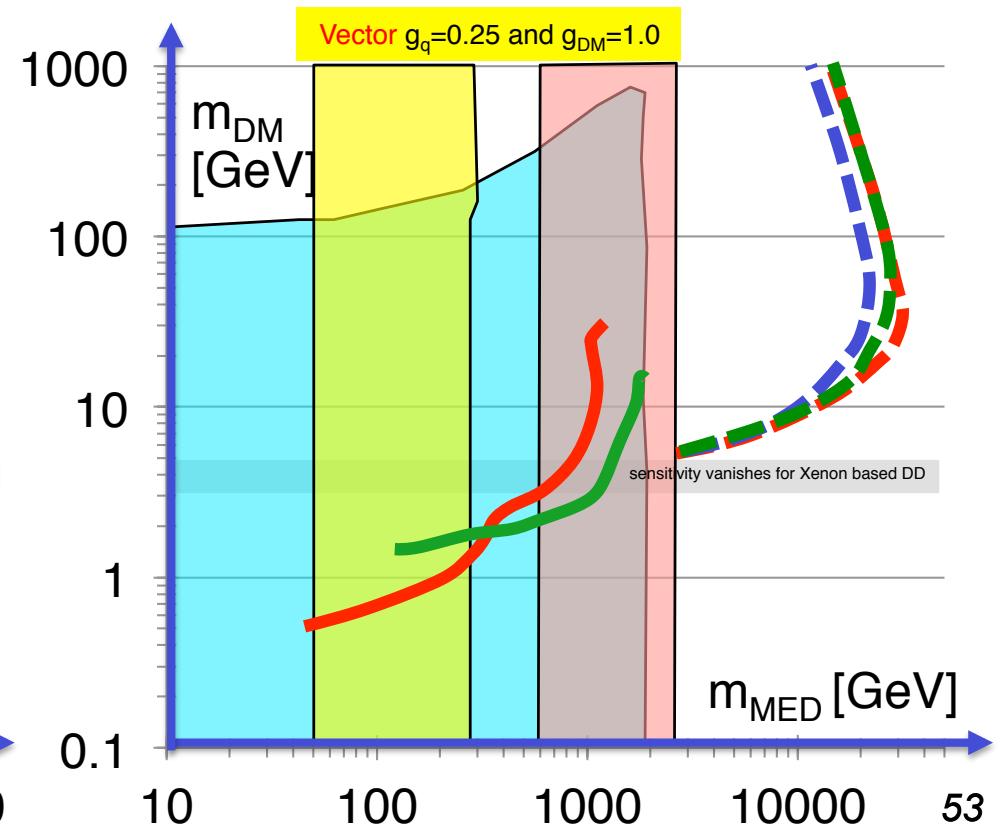
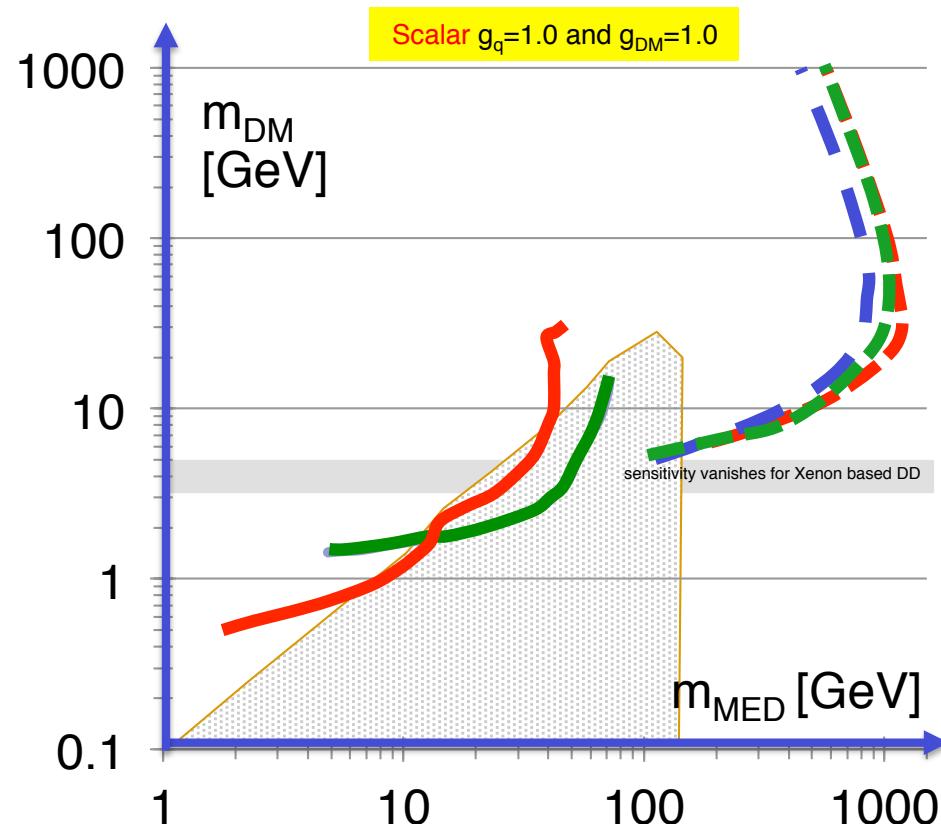
PandaX-II ----- [arXiv:1607.07400]
LUX ----- [arXiv:1608.07648]
XENON1T ----- [arXiv:1705.06655]

Collider

2L + MET [SUS-17-001]

[Note: private estimate of 2D exclusion – only rough approximation!]

Boosted dijet [EXO-17-001]
Dijet [EXO-16-056]
Mono-jet [EXO-16-048]



DD limits in collider plane: σ_{SI}

DD Experiments

PandaX-II

[arXiv:1807.07101]

Collider

Boosted dijet

[FCC-17-001]

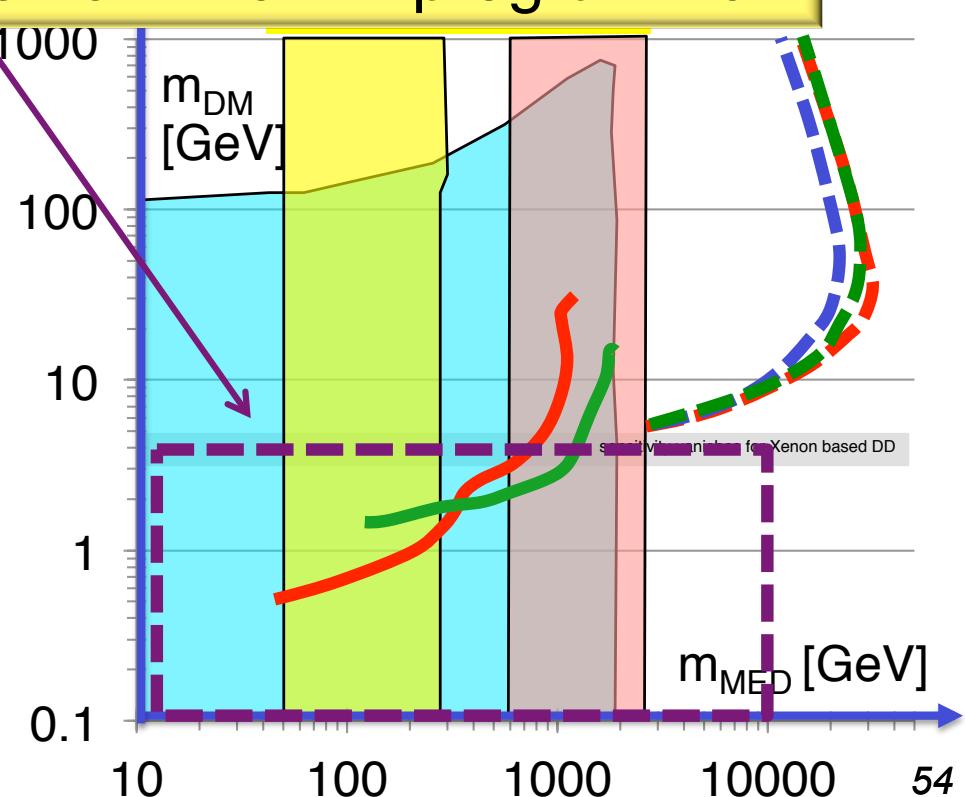
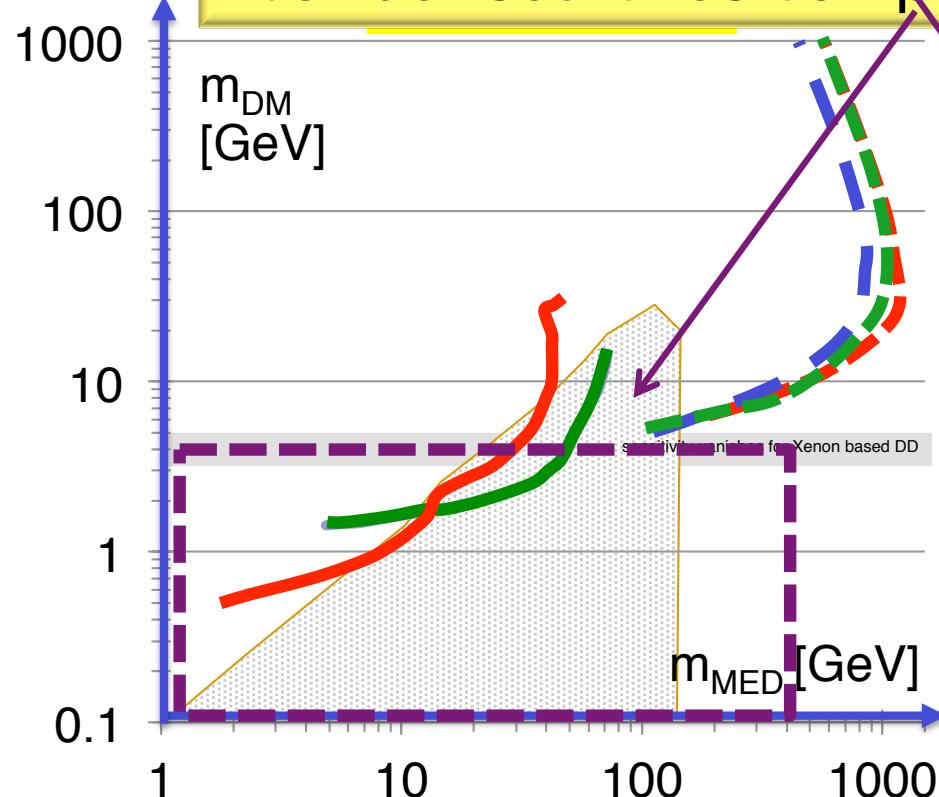
CRE

[arXiv:

CDM

[arXiv:

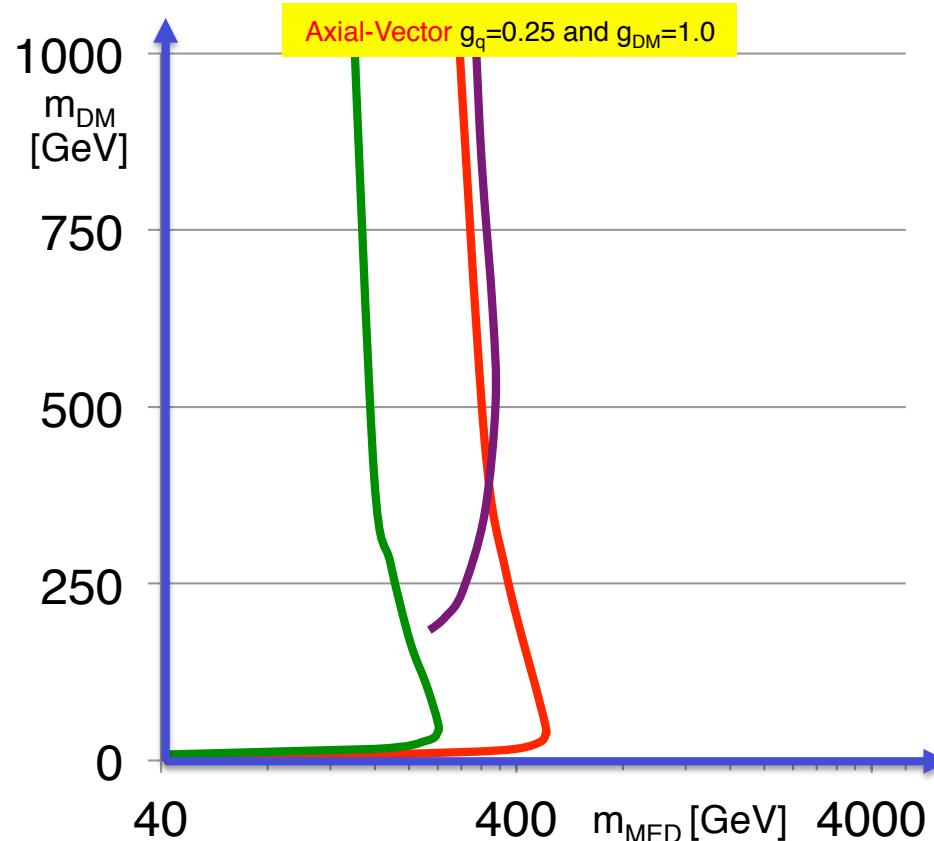
Broad brush bottom line:
For Vector and Scalar interactions, xenon based DD experiments rule, while for low m_{DM} ($<\sim 5$ GeV) collider searches complement the DD programme.



DD limits in collider plane: $\sigma_{\text{SD}(p,n)}$

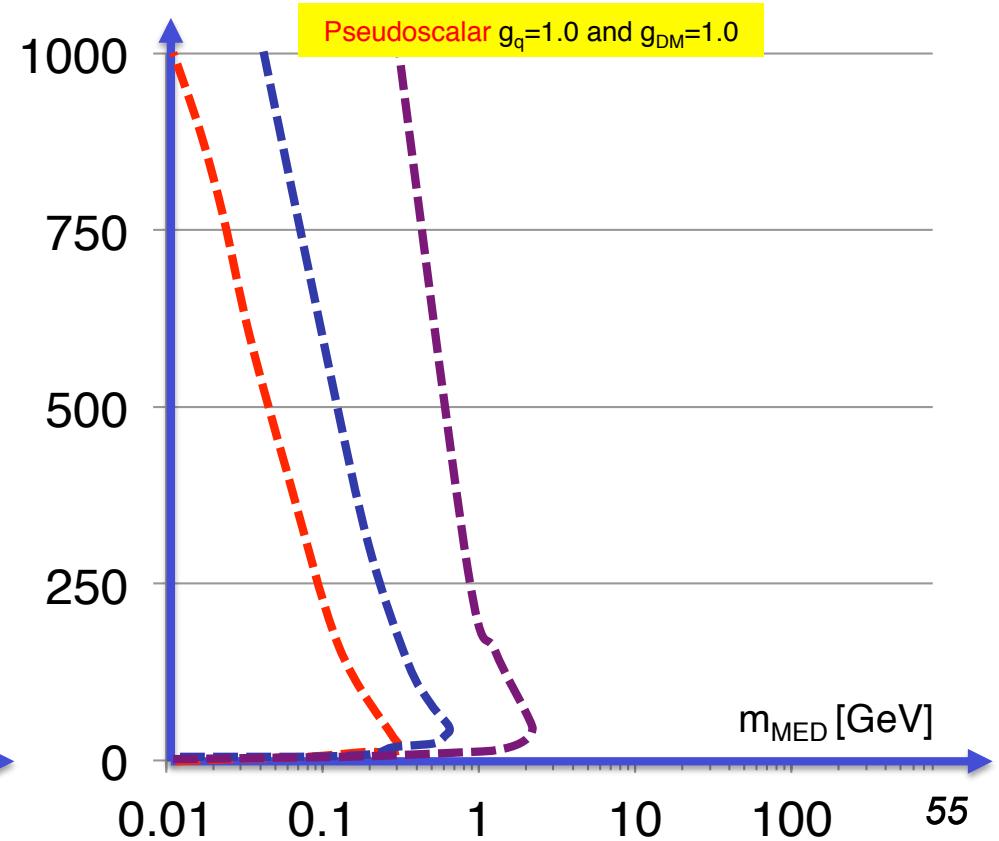
Axial-Vector simplified model:
 $g_q=0.25$ and $g_{\text{DM}}=1.0$

$$M_A = 1 \text{ TeV} \left(\frac{2.38 \times 10^{-42} \text{ cm}^2}{\sigma_{p,n}^{\text{SD}}} \right)^{0.25} \cdot \left(\frac{g_{\text{DM}} g_q}{0.25} \right)^{0.5} \left(\frac{\mu_n}{1 \text{ GeV}} \right)^{0.5}.$$



Based on
arXiv:1407.8257
Numerical calculation
provide by. C. McCabe

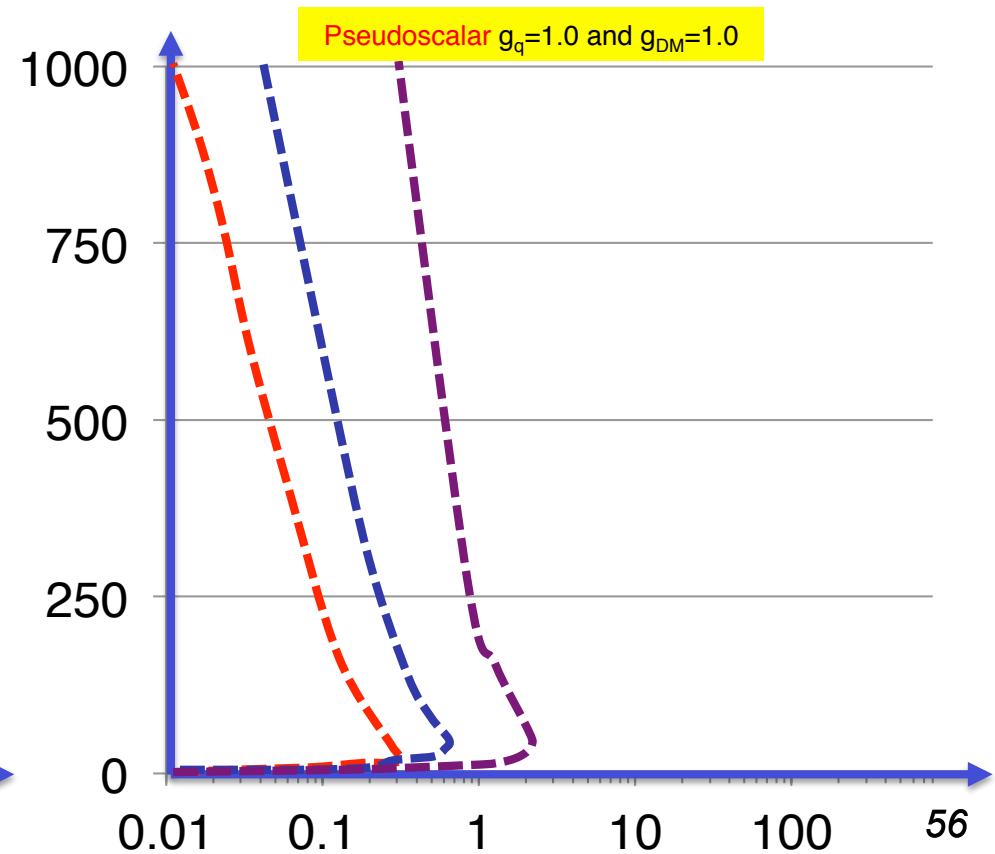
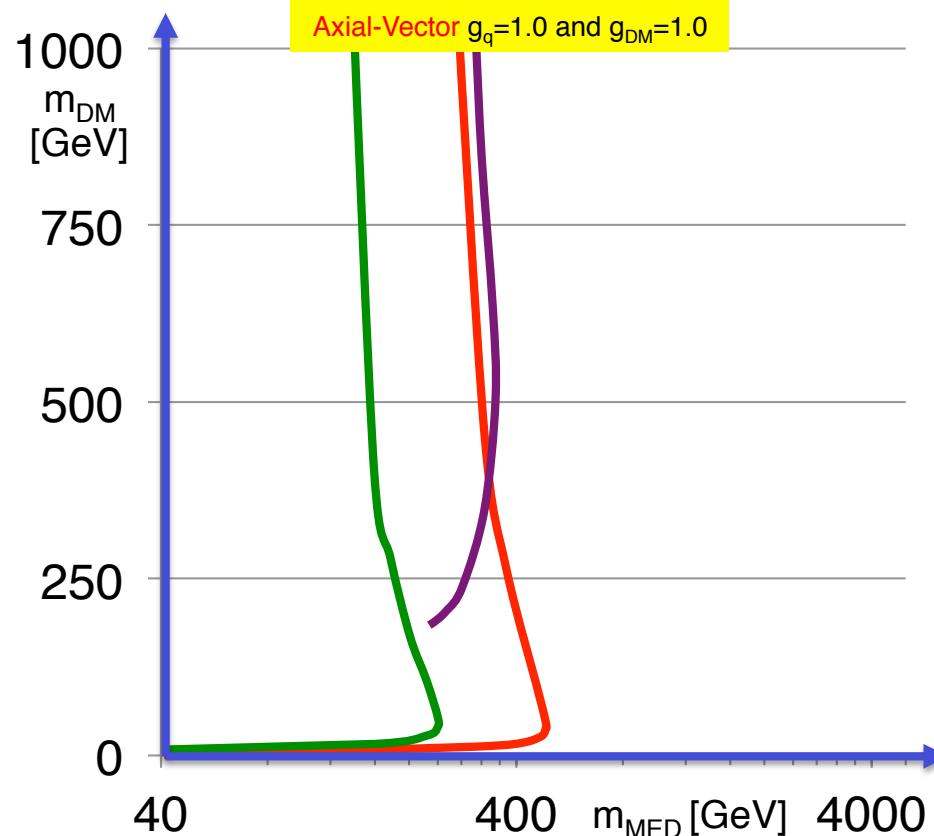
Pseudoscalar (PS)
simplified model:
 $g_q=1.0$ and $g_{\text{DM}}=1.0$



DD limits in collider plane: $\sigma_{\text{SD}(p,n)}$

DD/ID Experiments

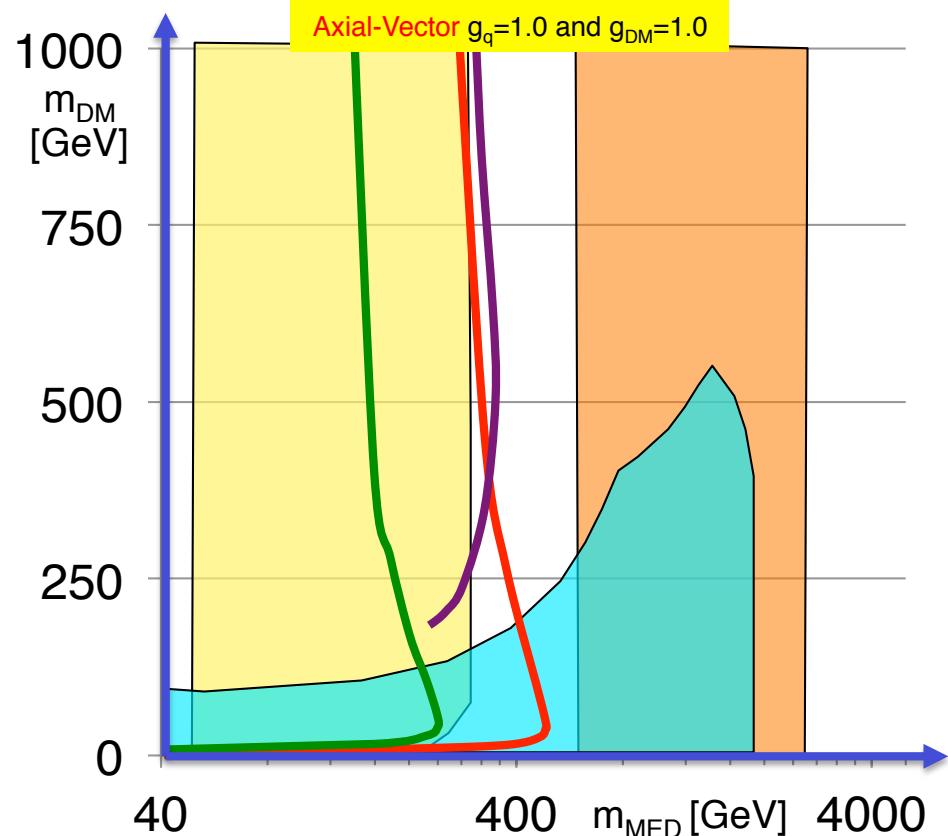
PICO60	—	PICO60	- - -
[arXiv:1702.07666]		[arXiv:1702.07666]	
LUX ($\sigma_{\text{SD},p}$)	—	XENON1T	- - -
[arXiv:1705.03380]		[arXiv:1705.06655]	
ICECUBE	—	LZ/XENONnT	- - -
($t\bar{t}$ bar)			
[arXiv:1601.00653]		15 t-yr projection	



DD limits in collider plane: $\sigma_{SD(p,n)}$

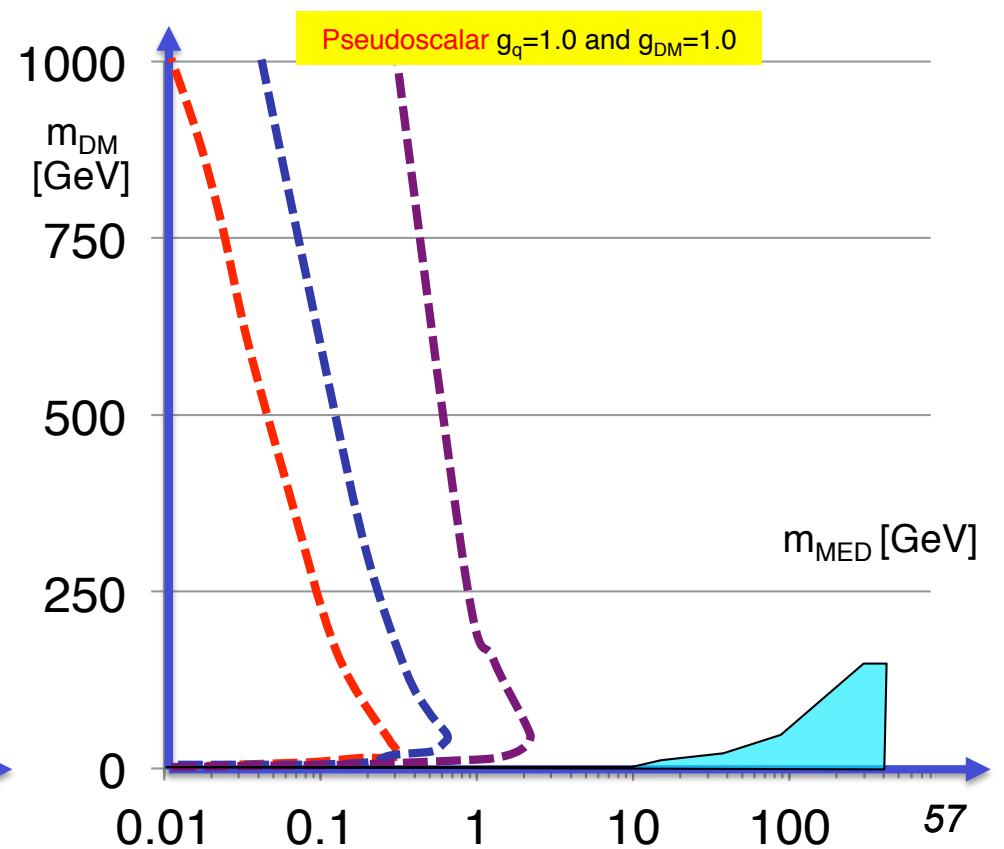
DD/ID Experiments

PICO60	—	PICO60	—
[arXiv:1702.07666]		[arXiv:1702.07666]	
LUX ($\sigma_{SD,p}$)	—	XENON1T	—
[arXiv:1705.03380]		[arXiv:1705.06655]	
ICECUBE (ttbar)	—	LZ/XENONnT	—
[arXiv:1601.00653]		15 t-yr projection	



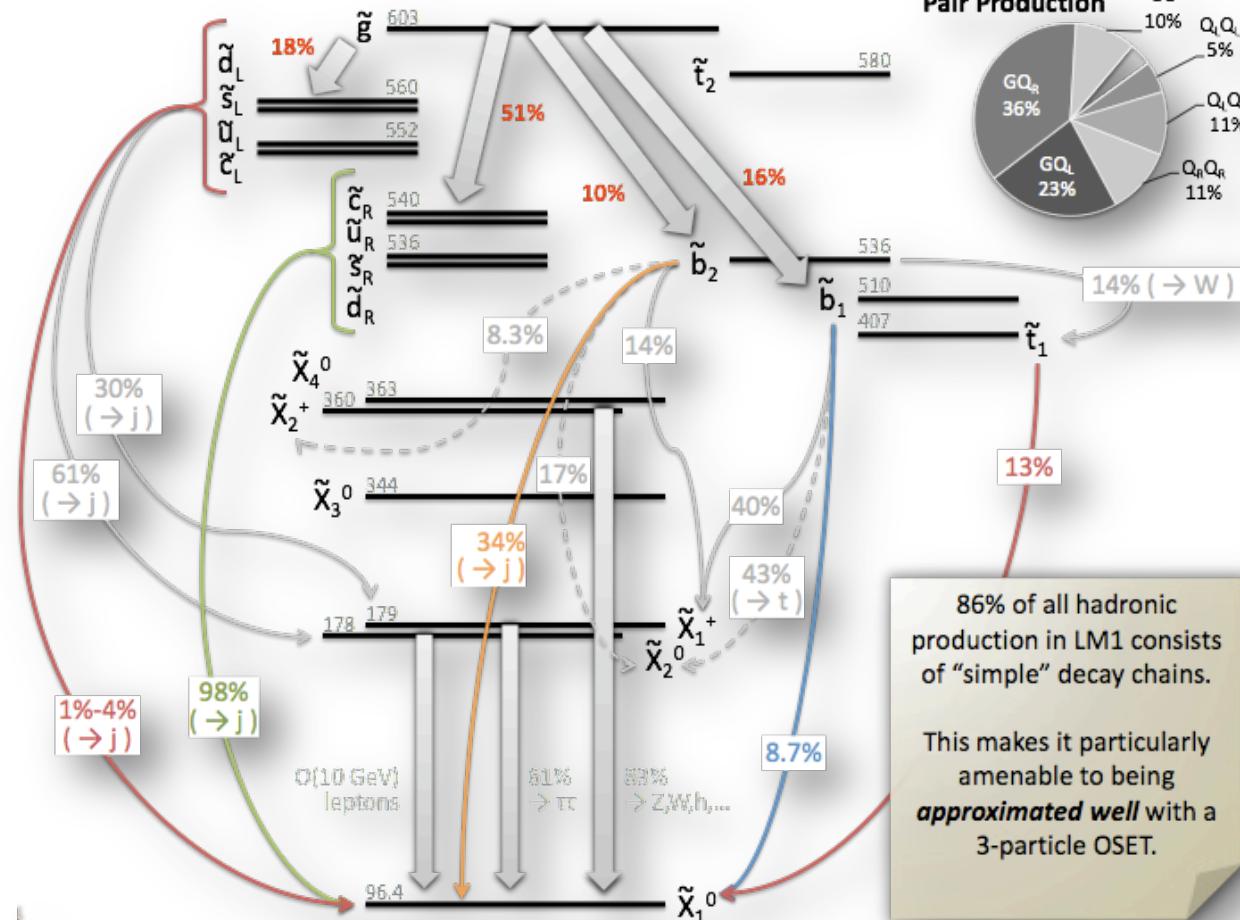
Collider

Boosted dijet	Yellow square
[EXO-17-001]	
Dijet	Red square
[EXO-16-056]	
Mono-jet	Cyan square
[EXO-16-048]	



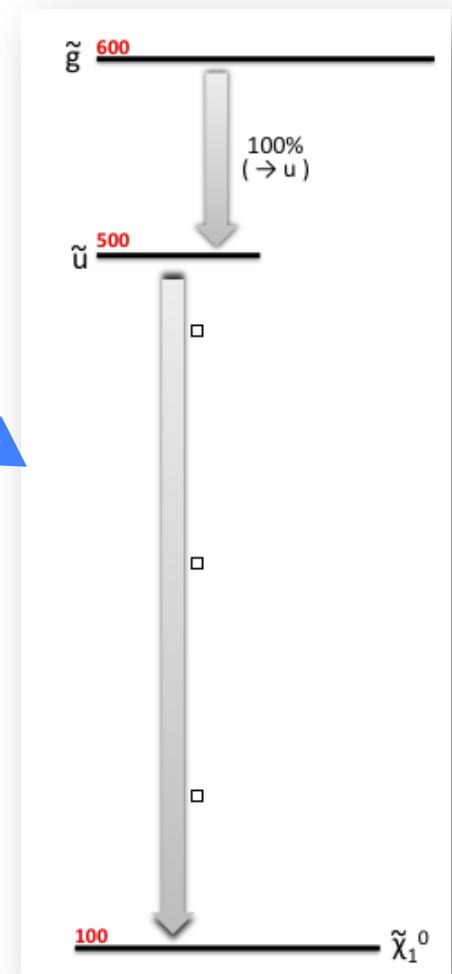
Interpretation of SUSY Searches in Simplified Models

CMSSM



Simplified model spectrum (SMS)
with 3 particles, 2 decay modes

What the individual searches are sensitive to is much more simple...

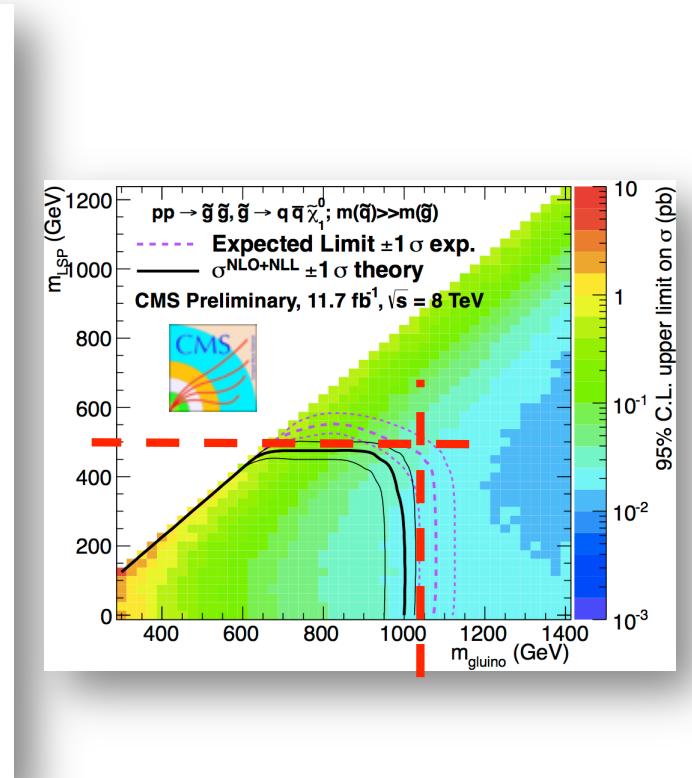


How to summarize SMS limits?

*Approach taken in the 2012 and 2013 Experimental SUSY PDG reviews
[OB & Paul De Jong]:*

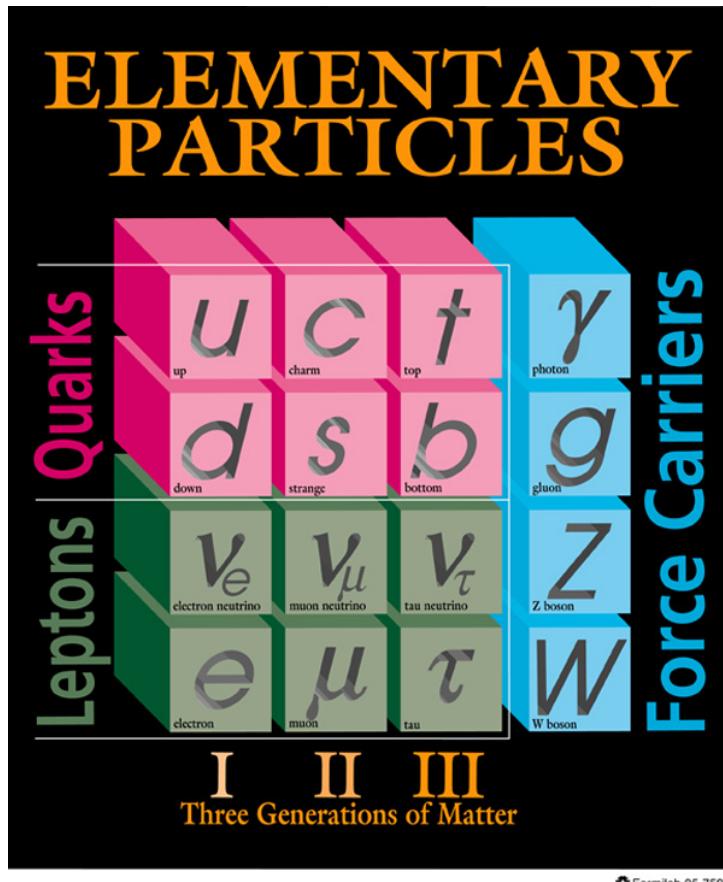
<http://pdg.lbl.gov/2012/reviews/rpp2012-rev-susy-2-experiment.pdf>
<http://pdg.lbl.gov/2013/reviews/rpp2013-rev-susy-2-experiment.pdf>

Model	Assumption	$m_{\tilde{q}}$	$m_{\tilde{g}}$
CMSSM	$m_{\tilde{q}} \approx m_{\tilde{g}}$	1400	1400
	all $m_{\tilde{q}}$	-	800
	all $m_{\tilde{g}}$	1300	-
Simplified model $\tilde{g}\tilde{g}$	$m_{\tilde{\chi}_1^0} = 0$	-	900
	$m_{\tilde{\chi}_1^0} > 300$	-	no limit
Simplified model $\tilde{q}\tilde{q}$	$m_{\tilde{\chi}_1^0} = 0$	750	-
	$m_{\tilde{\chi}_1^0} > 250$	no limit	-
Simplified model $\tilde{g}\tilde{q}, \tilde{g}\bar{\tilde{q}}$	$m_{\tilde{\chi}_1^0} = 0, m_{\tilde{q}} \approx m_{\tilde{g}}$	1500	1500
	$m_{\tilde{\chi}_1^0} = 0, \text{ all } m_{\tilde{g}}$	1400	-
	$m_{\tilde{\chi}_1^0} = 0, \text{ all } m_{\tilde{q}}$	-	900



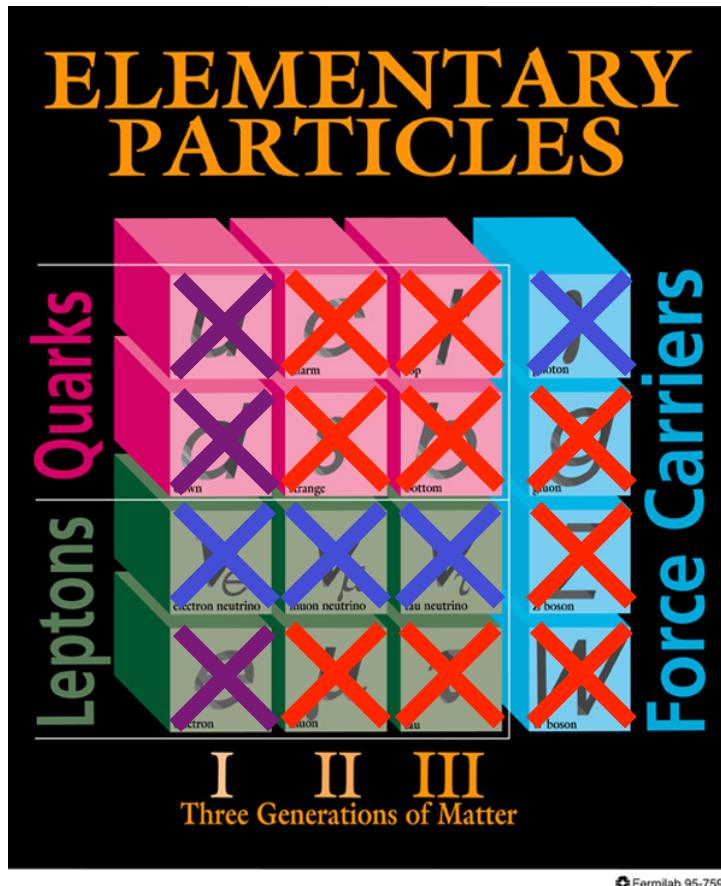
This was an appropriate approach for the rather limited amount of inclusive searches and corresponding SMS interpretations available in 2011 (7 TeV).

Dark Matter: Particle Hypothesis



Known DM properties

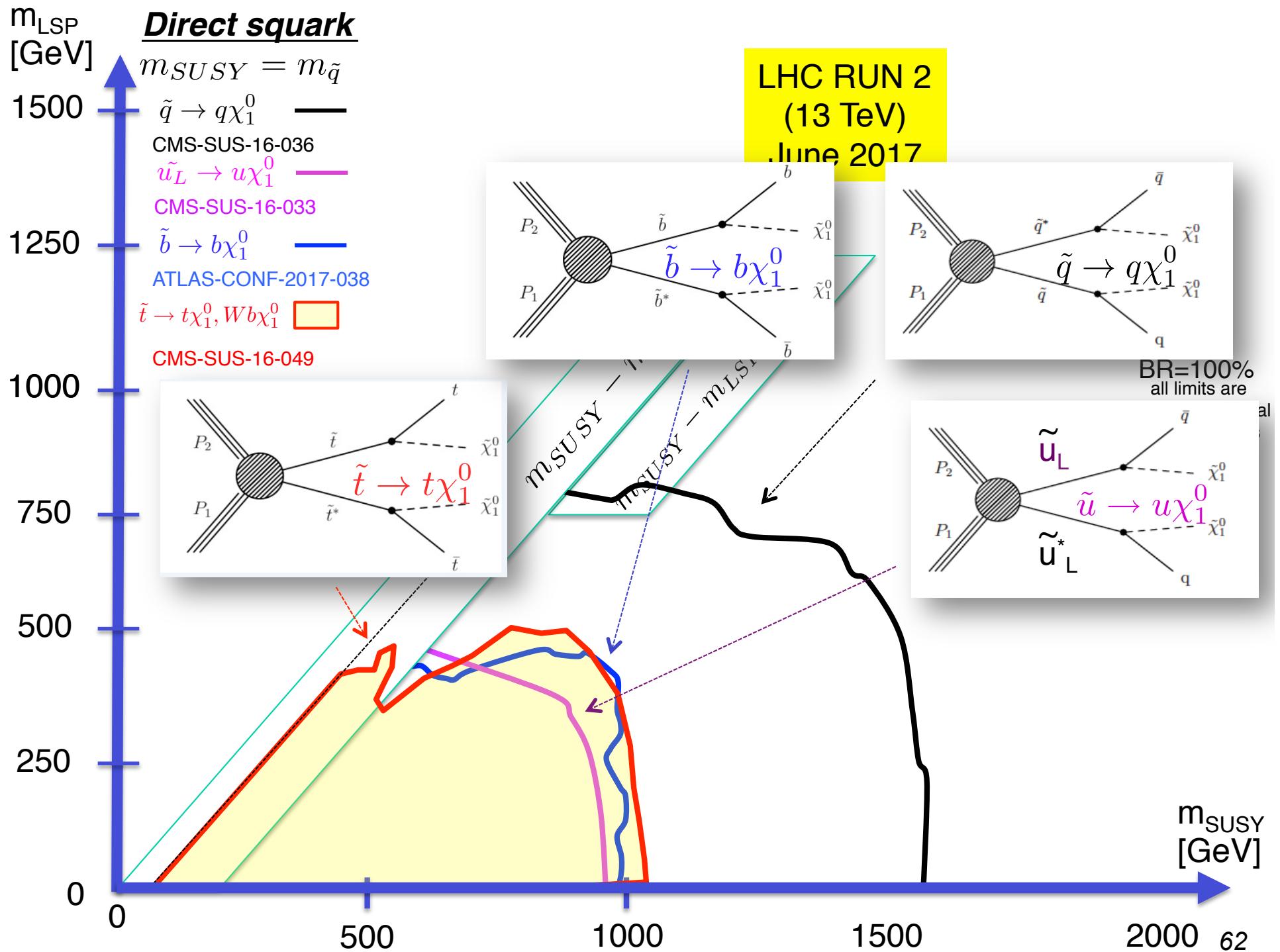
Dark Matter: Particle Hypothesis

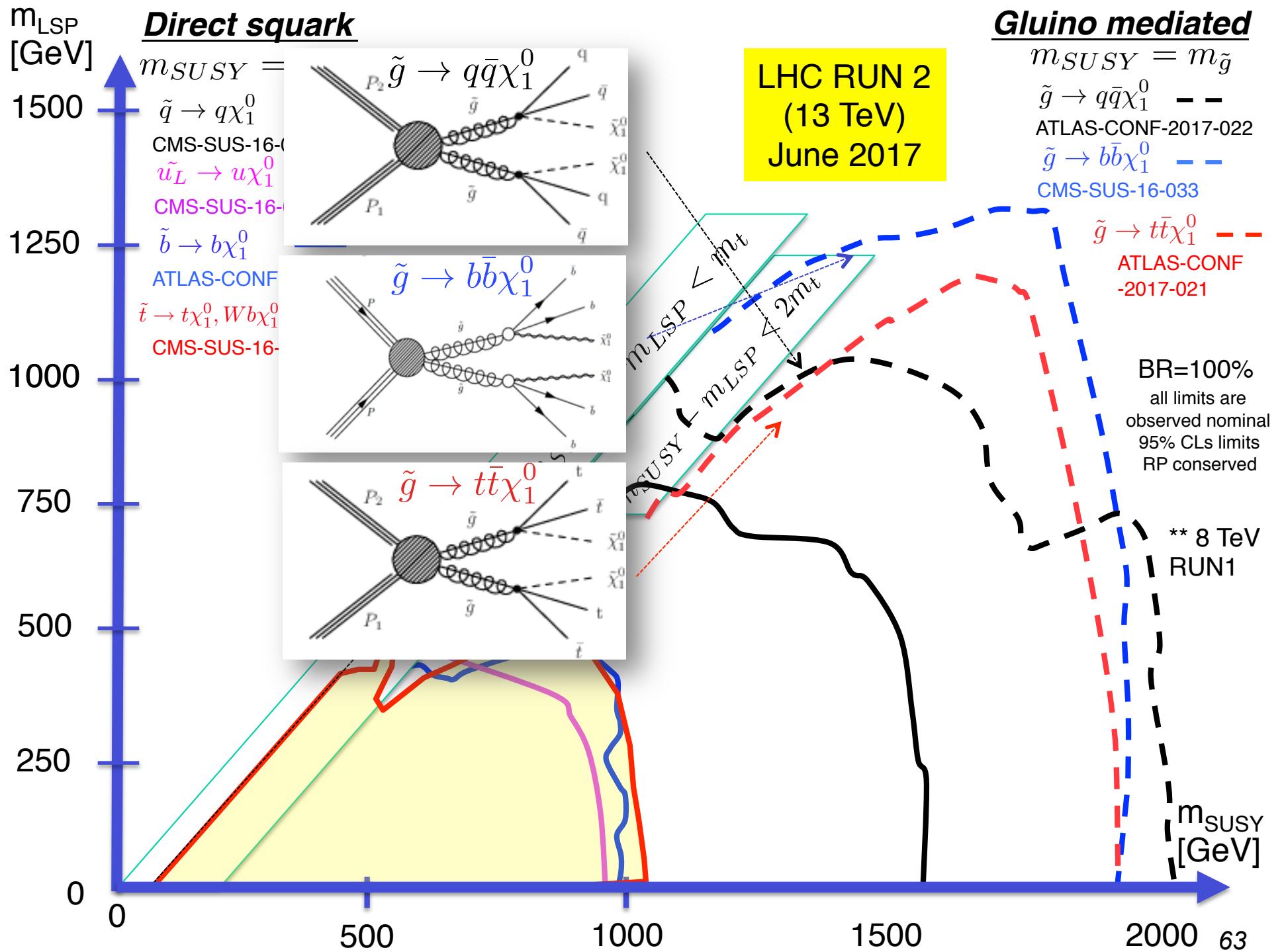


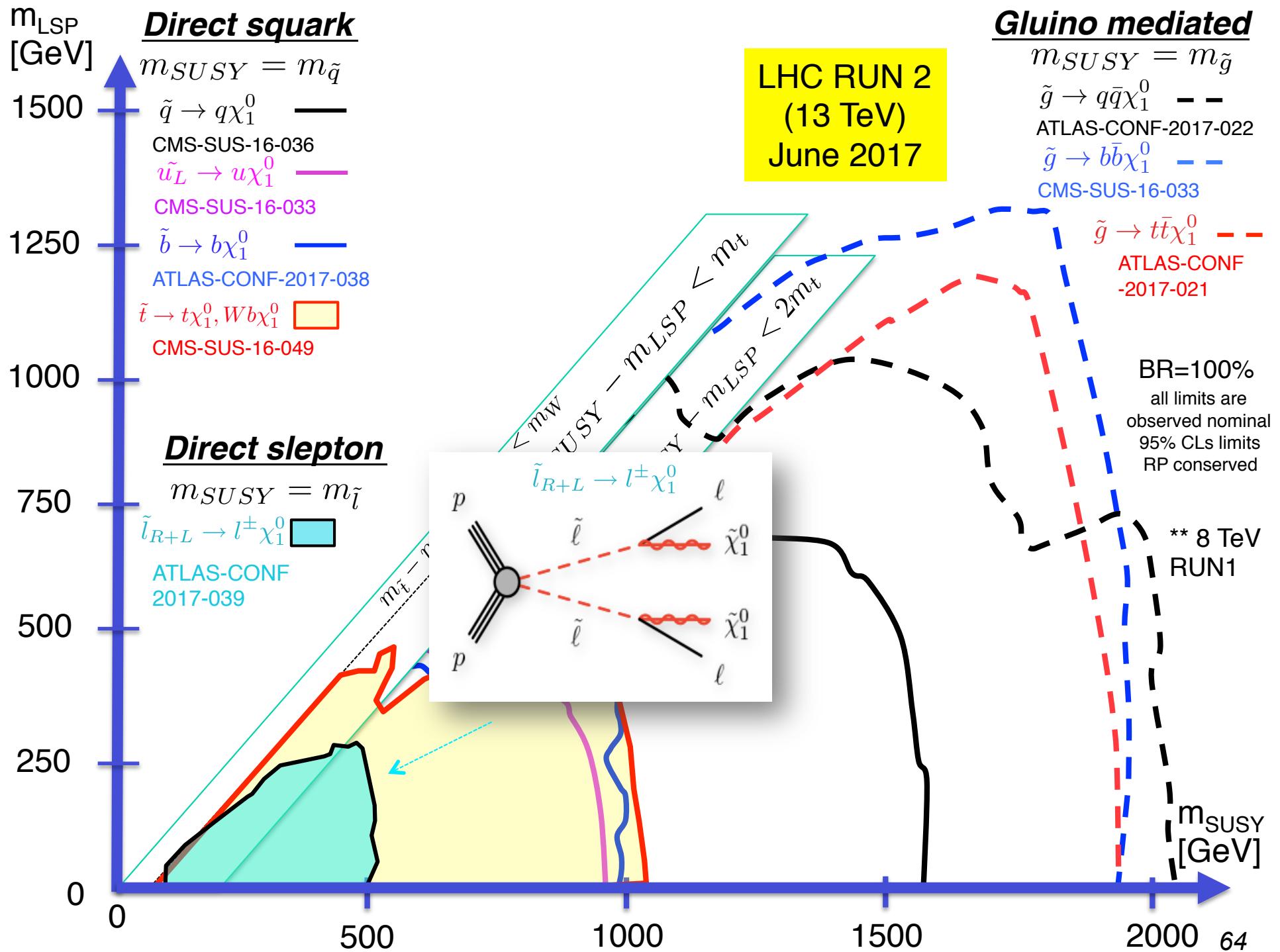
Known DM properties:

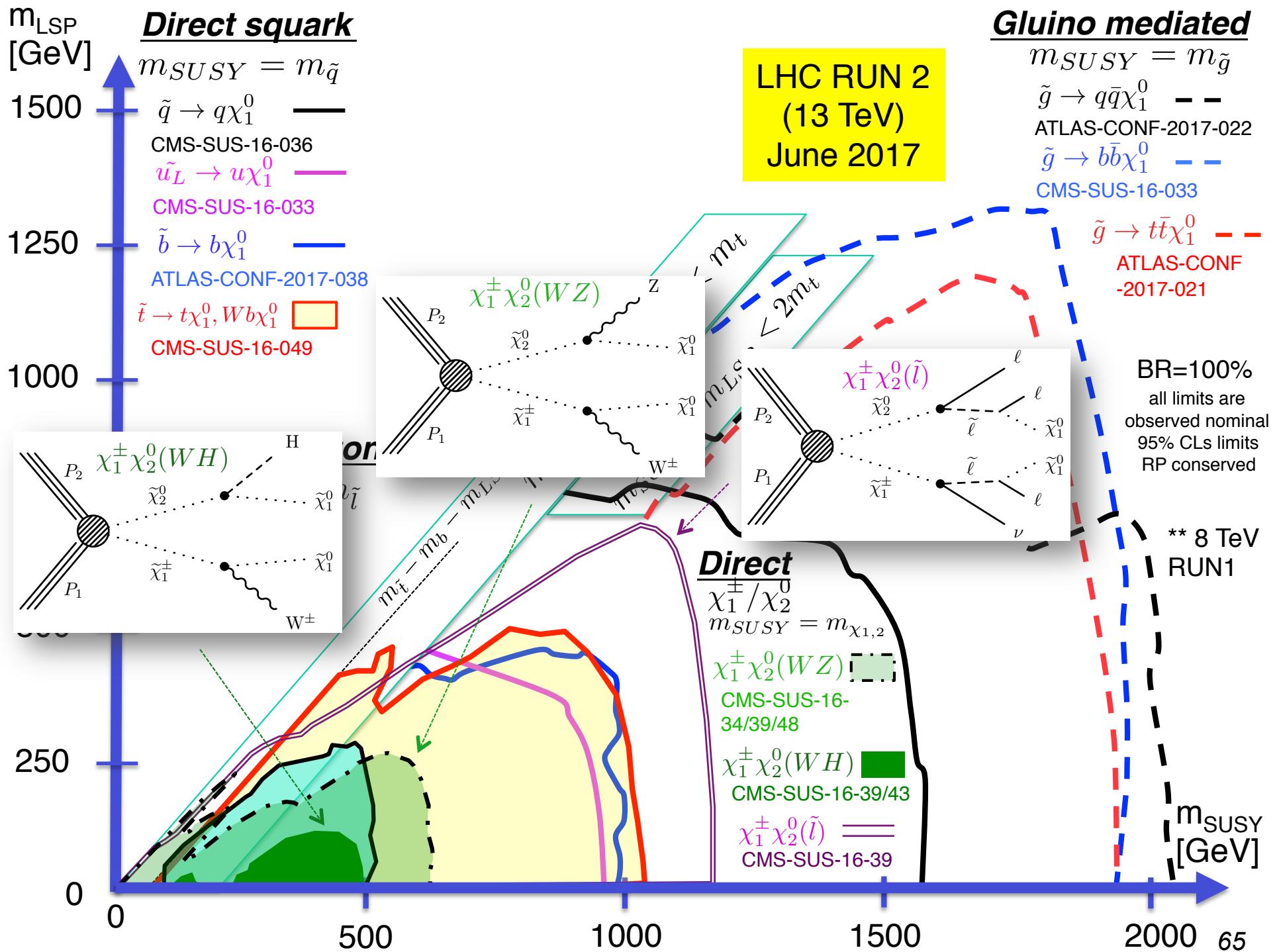
- Gravitationally interacting
- Not short-lived
- Not hot
- Not baryonic

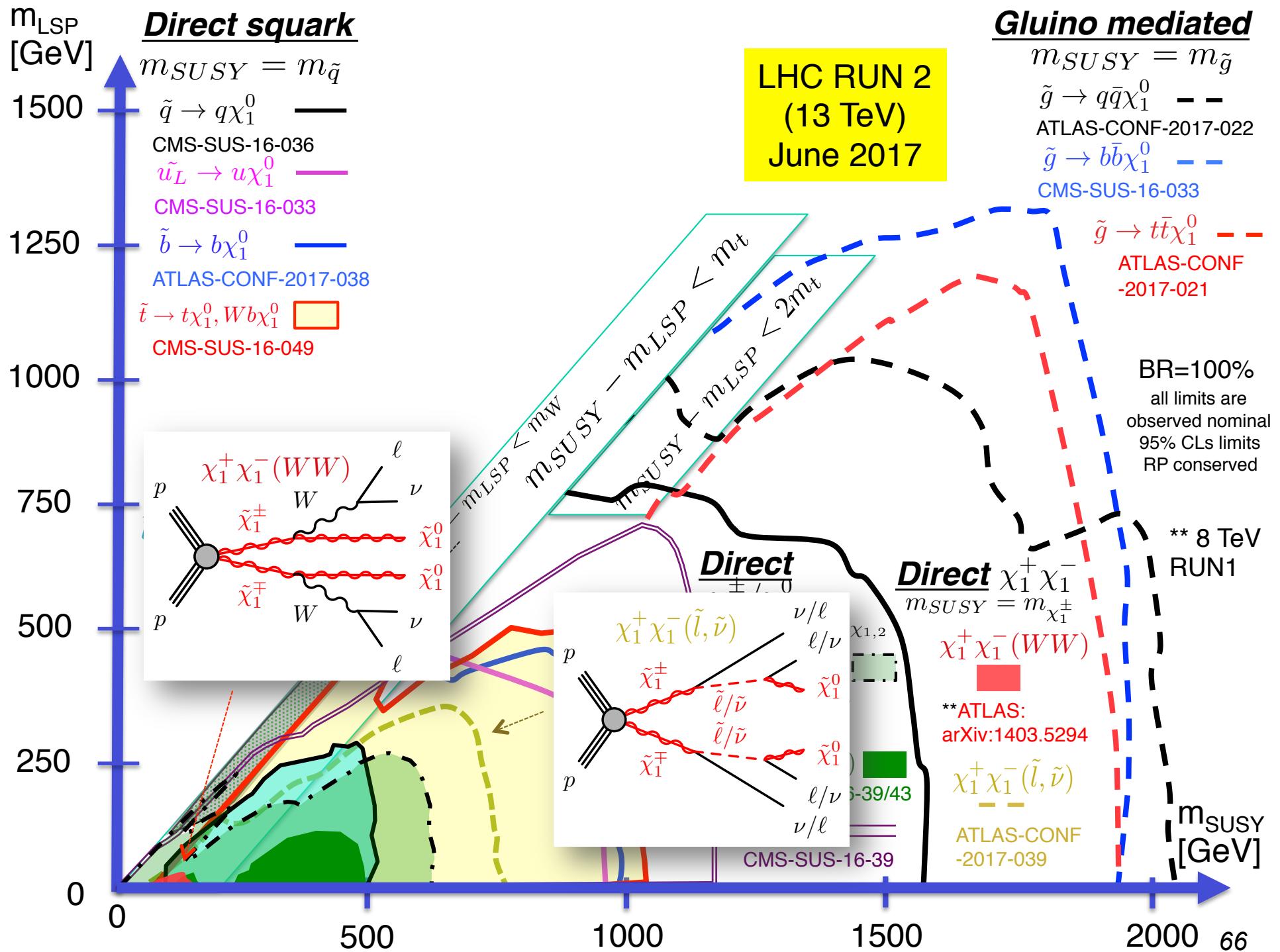
Hypothesis: Dark Matter is a new particle (or particles)

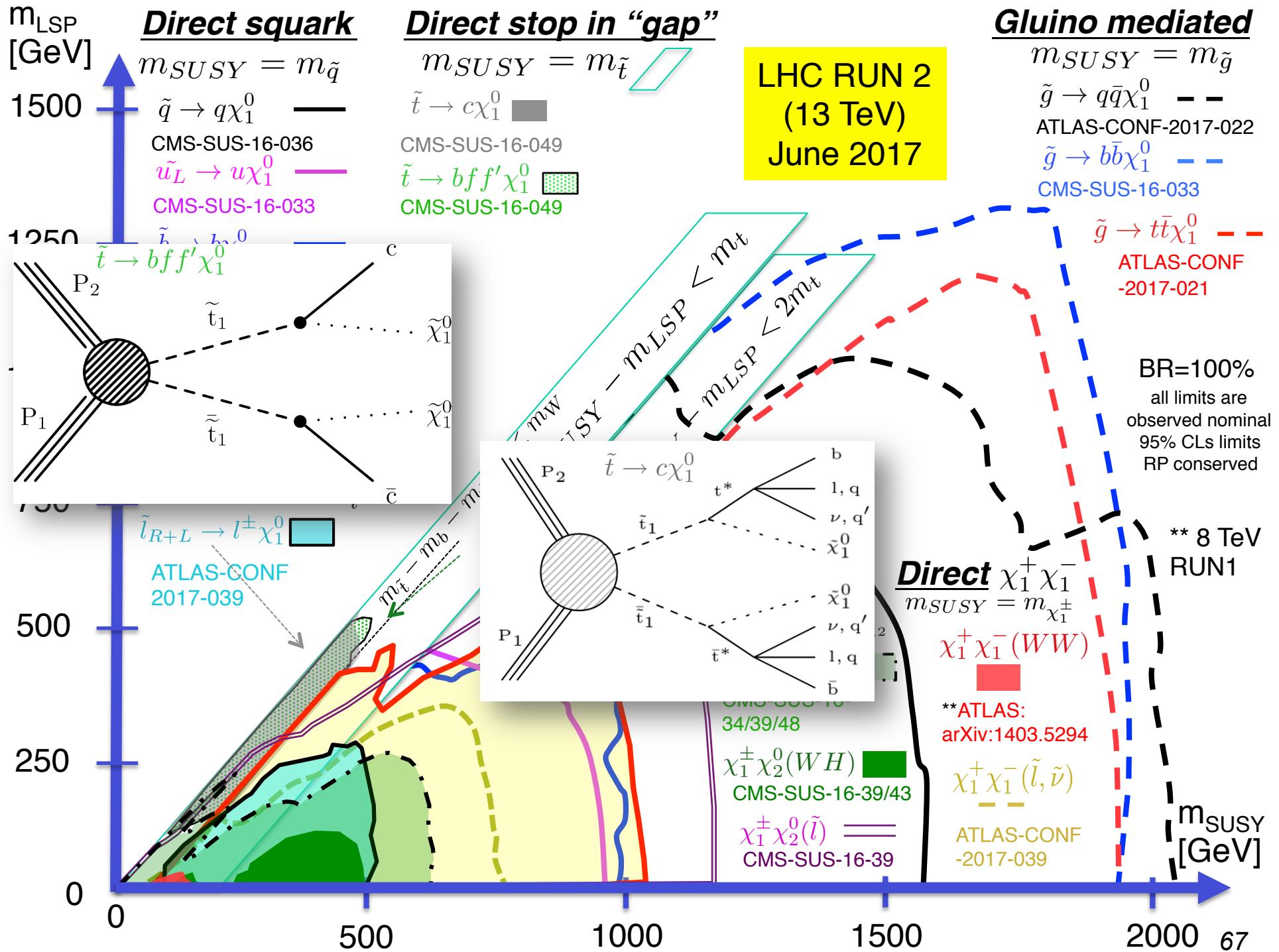












DM Mechanisms

■ stau coann. ■ $\tilde{\chi}_1^\pm$ coann. ■ slep coann.

■ gluino coann.

■ squark coann.

$$\left(\frac{M_{\tilde{\tau}}}{m_{\chi_1^0}} - 1 \right) < 0.15$$

Stau
coannihilation

$$\left| \frac{M_B}{m_{\chi_1^0}} - 2 \right| < 0.4$$

$B = h, Z$ or H/A
funnel

$$\left(\frac{M_{\chi_1^\pm}}{m_{\chi_1^0}} - 1 \right) < 0.25$$

Chargino
Co-annihilation

$$\left| \frac{\mu}{m_{\chi_1^0}} - 1 \right| < 0.30$$

Higgsino enriched
“focus-point” like

$$\left(\frac{M_{\tilde{l}}}{m_{\chi_1^0}} - 1 \right) < 0.15$$

Slepton
Co-annihilation

Hybrid regions:

In addition to the ‘primary’ regions where only one of the conditions is satisfied, there are also ‘hybrid’ regions where more than one condition is satisfied. These are indicated using combined colours.

$$\left(\frac{M_{\tilde{g}}}{m_{\chi_1^0}} - 1 \right) < 0.25$$

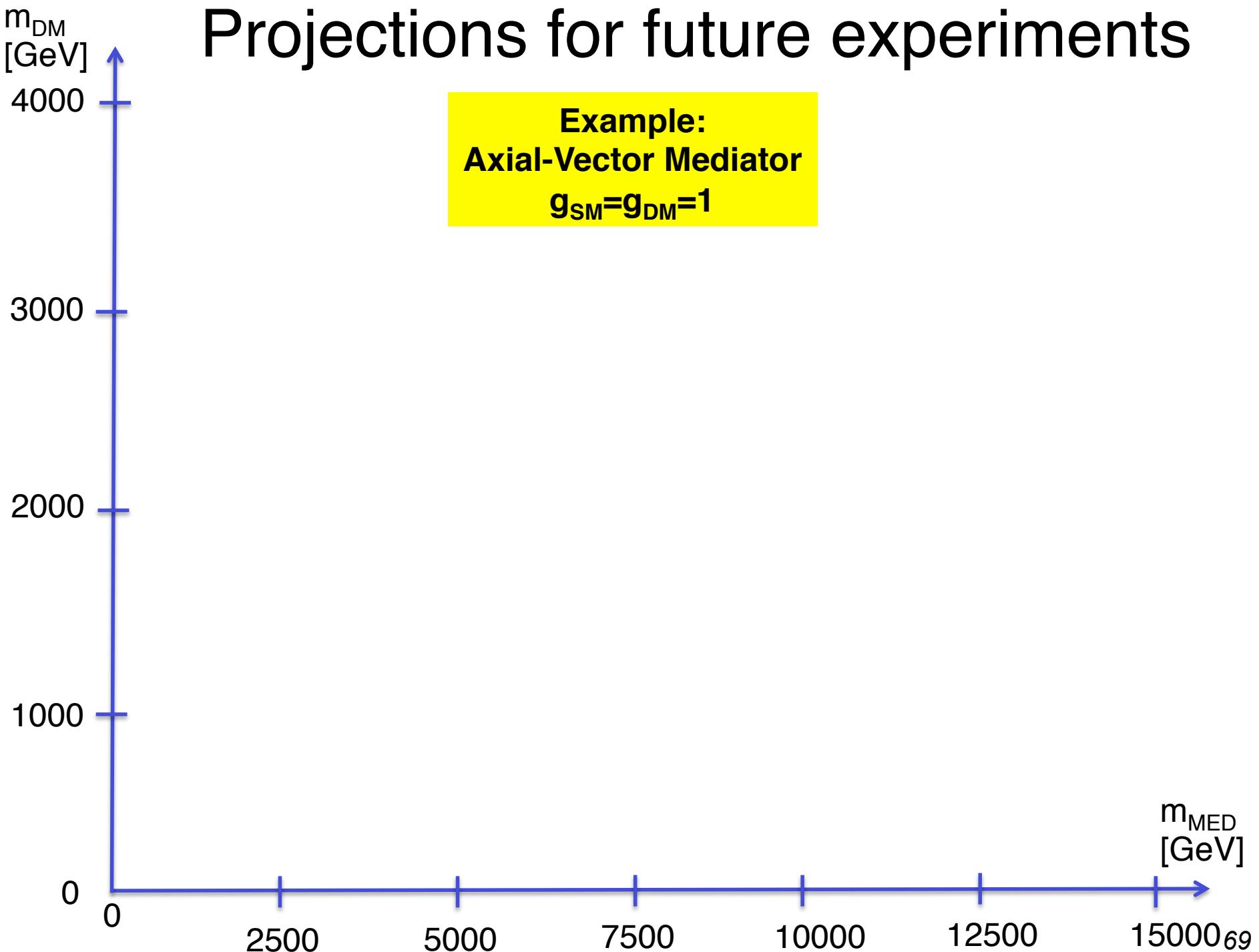
Gluino
Co-annihilation

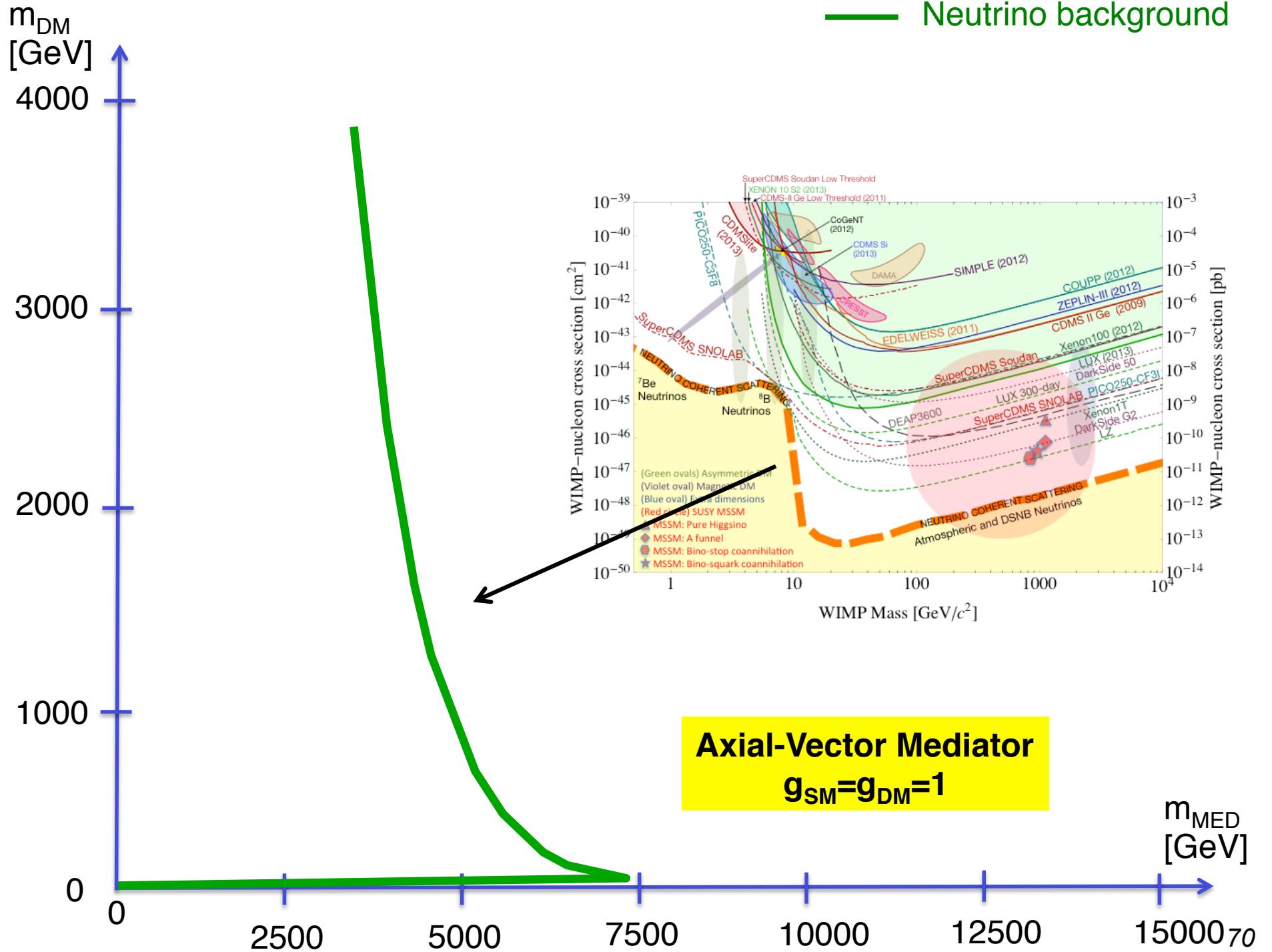
$$\left(\frac{M_{\tilde{q}}}{m_{\chi_1^0}} - 1 \right) < 0.20$$

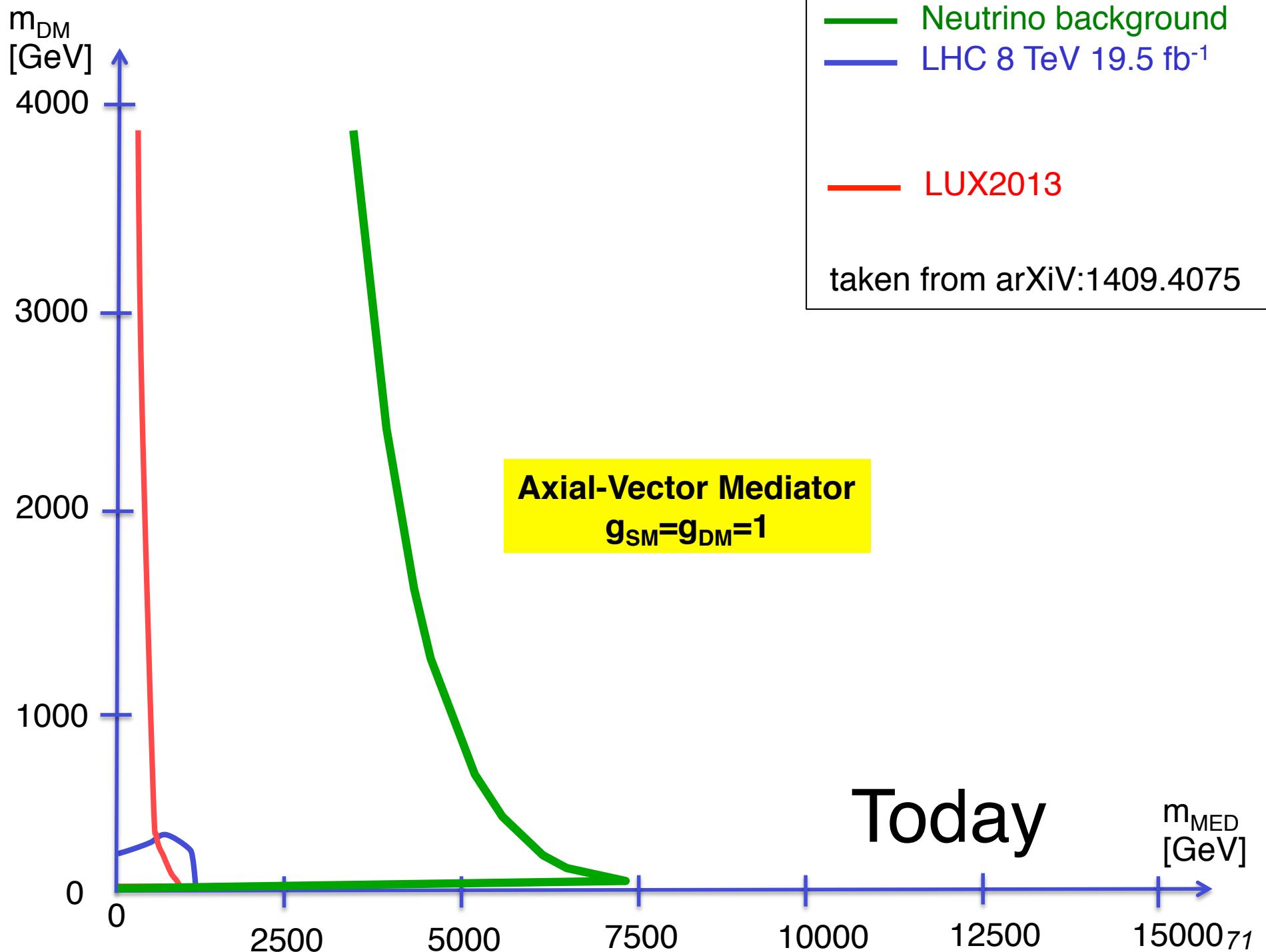
Squark
Co-annihilation

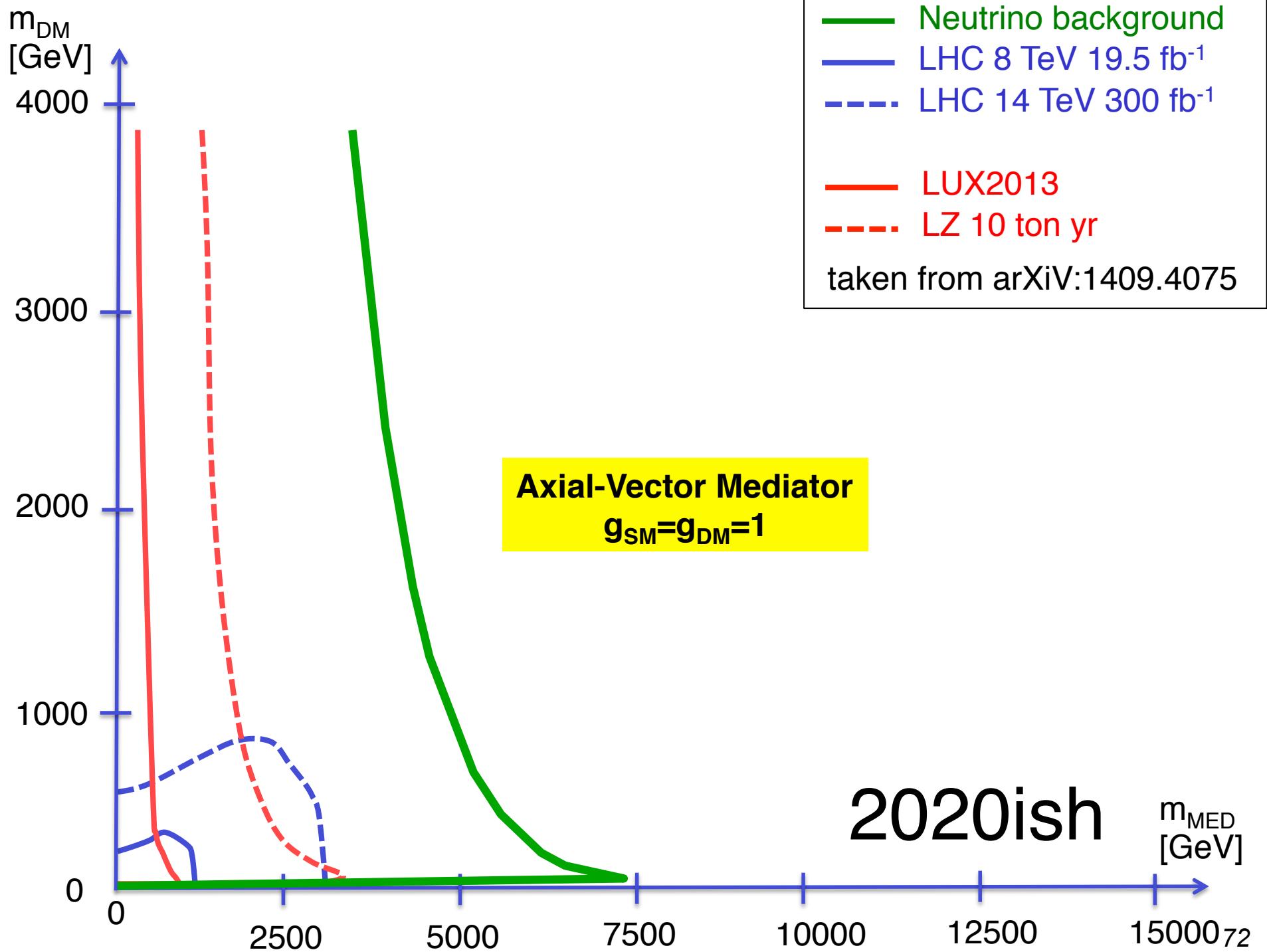
See also arXiv:1508.01173
for further details

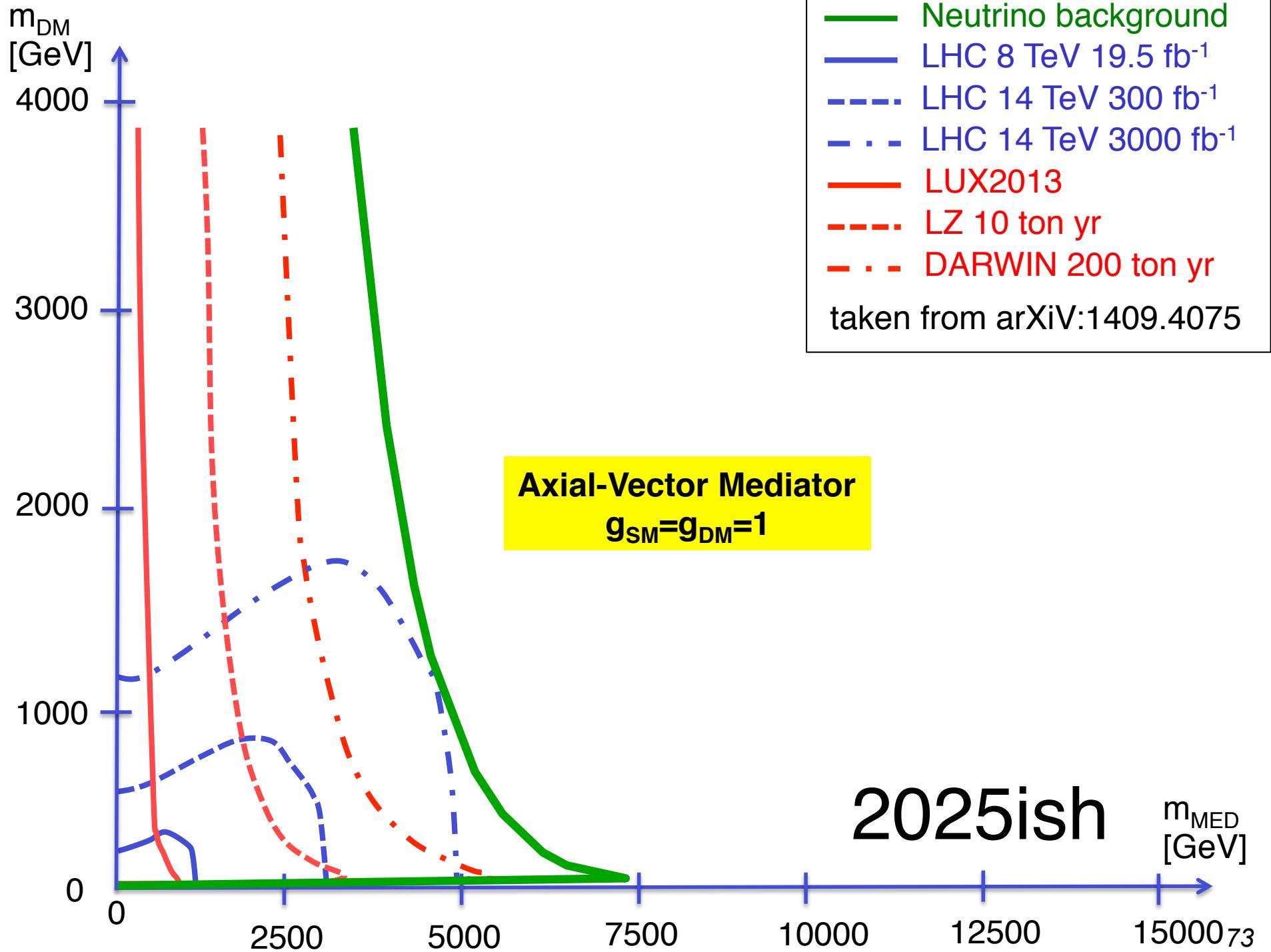
Projections for future experiments

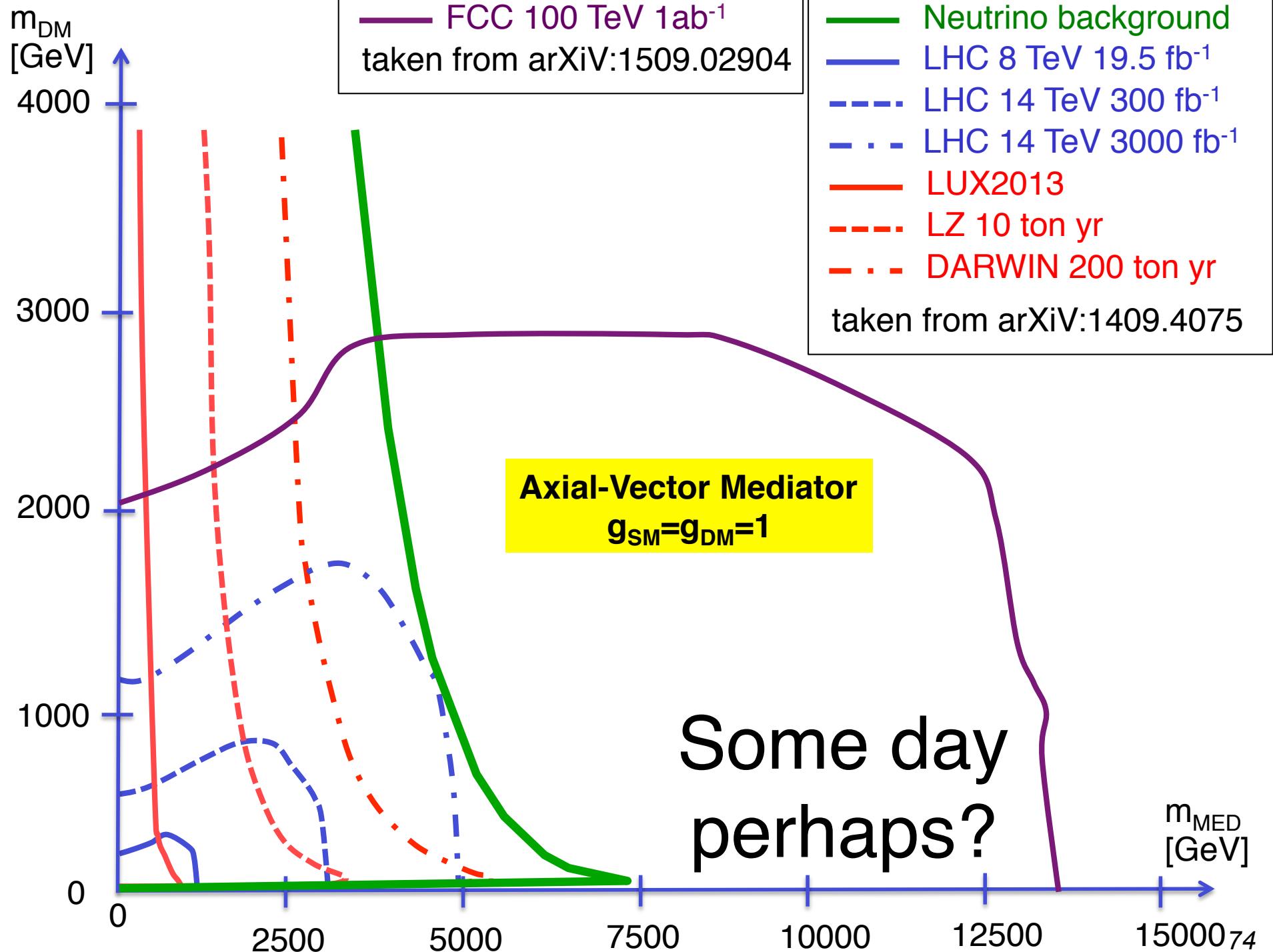


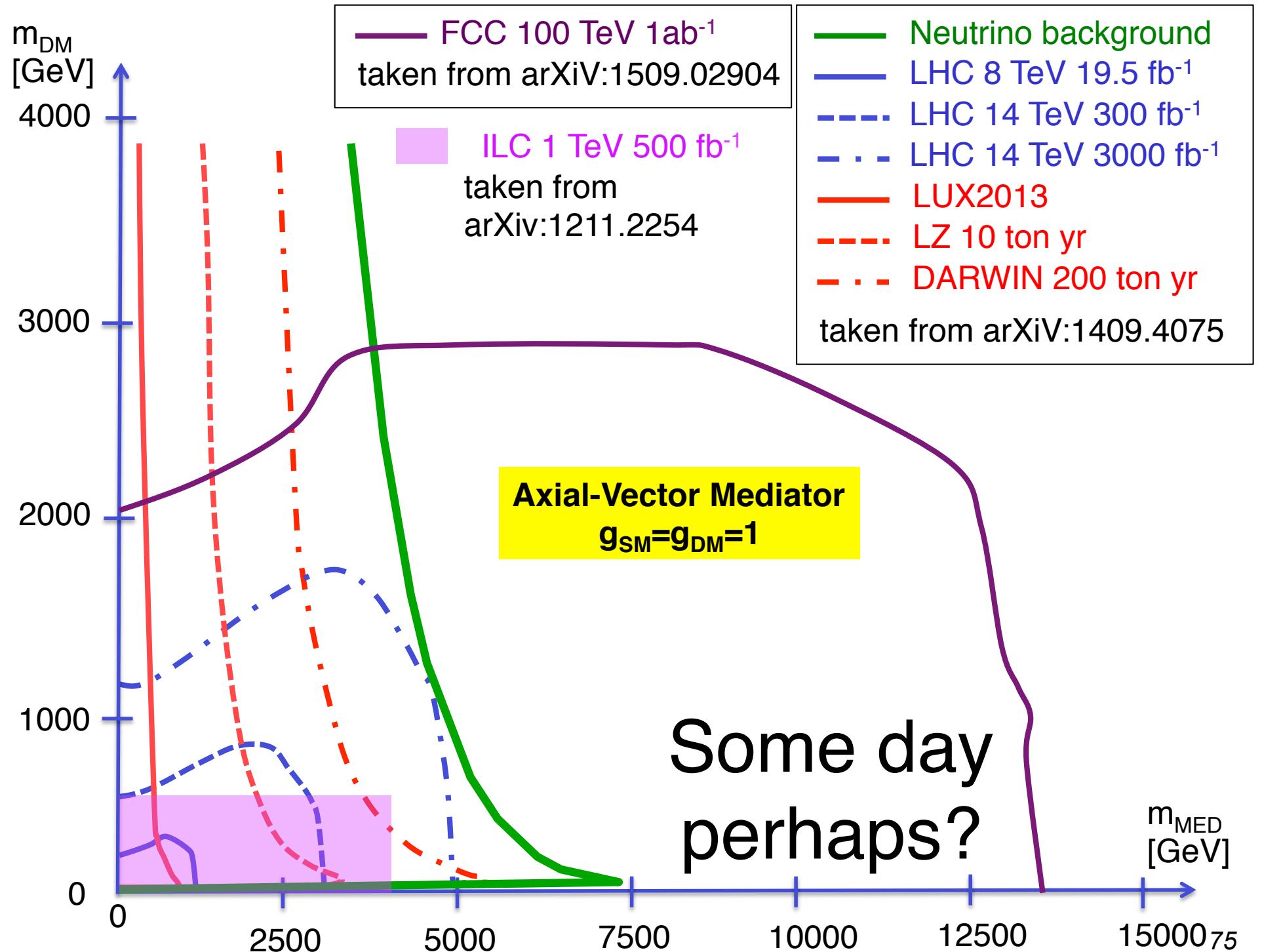




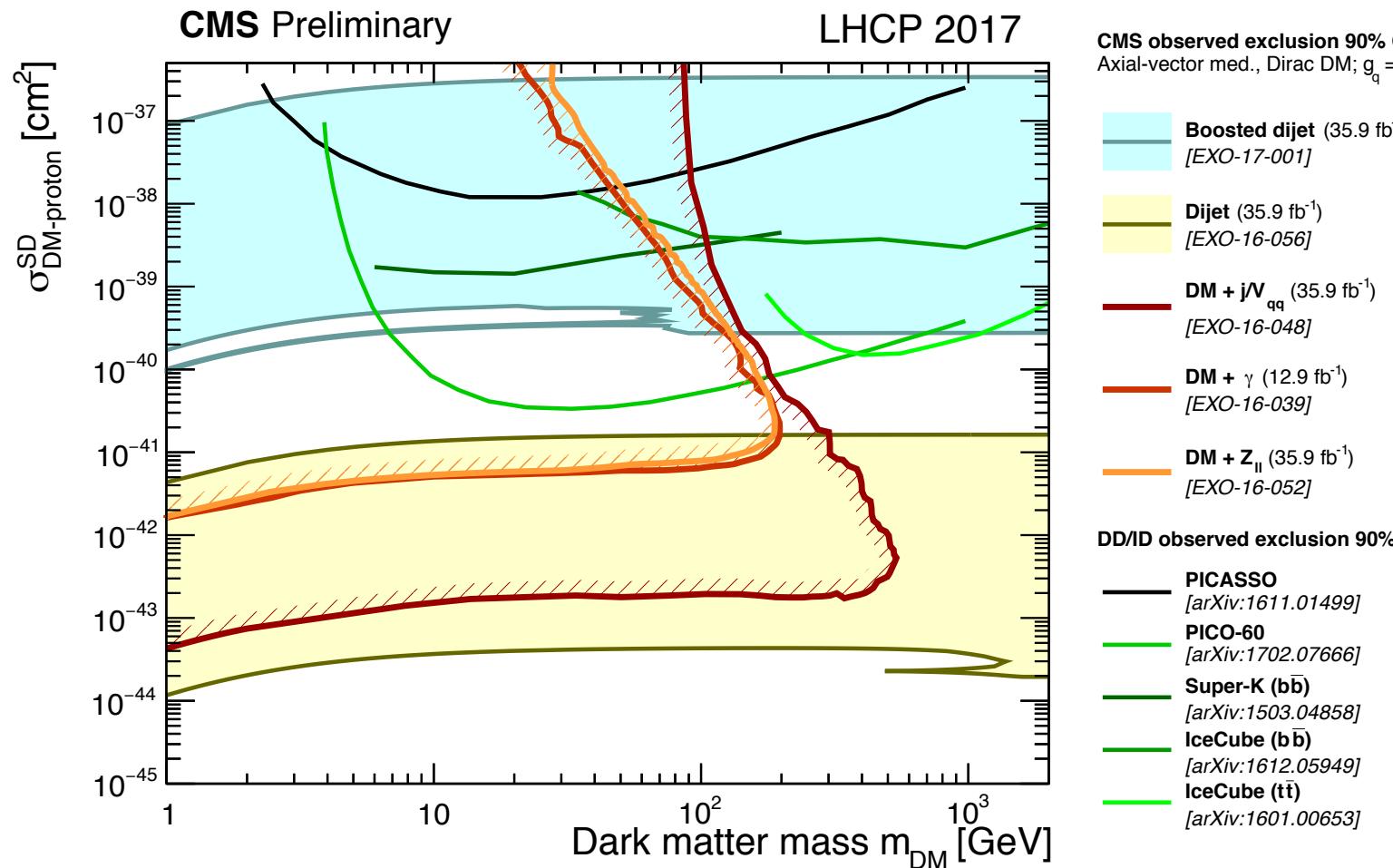




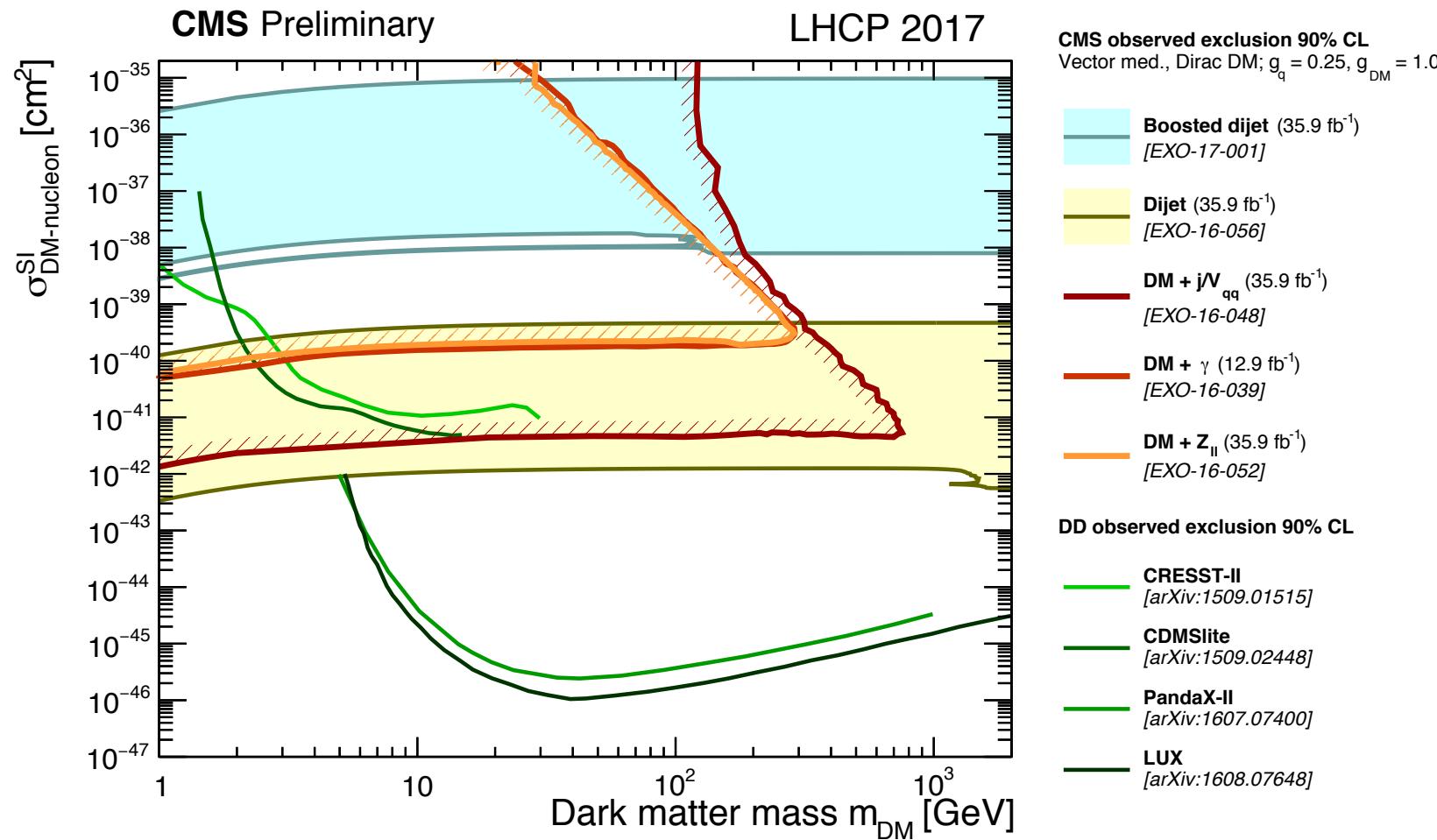




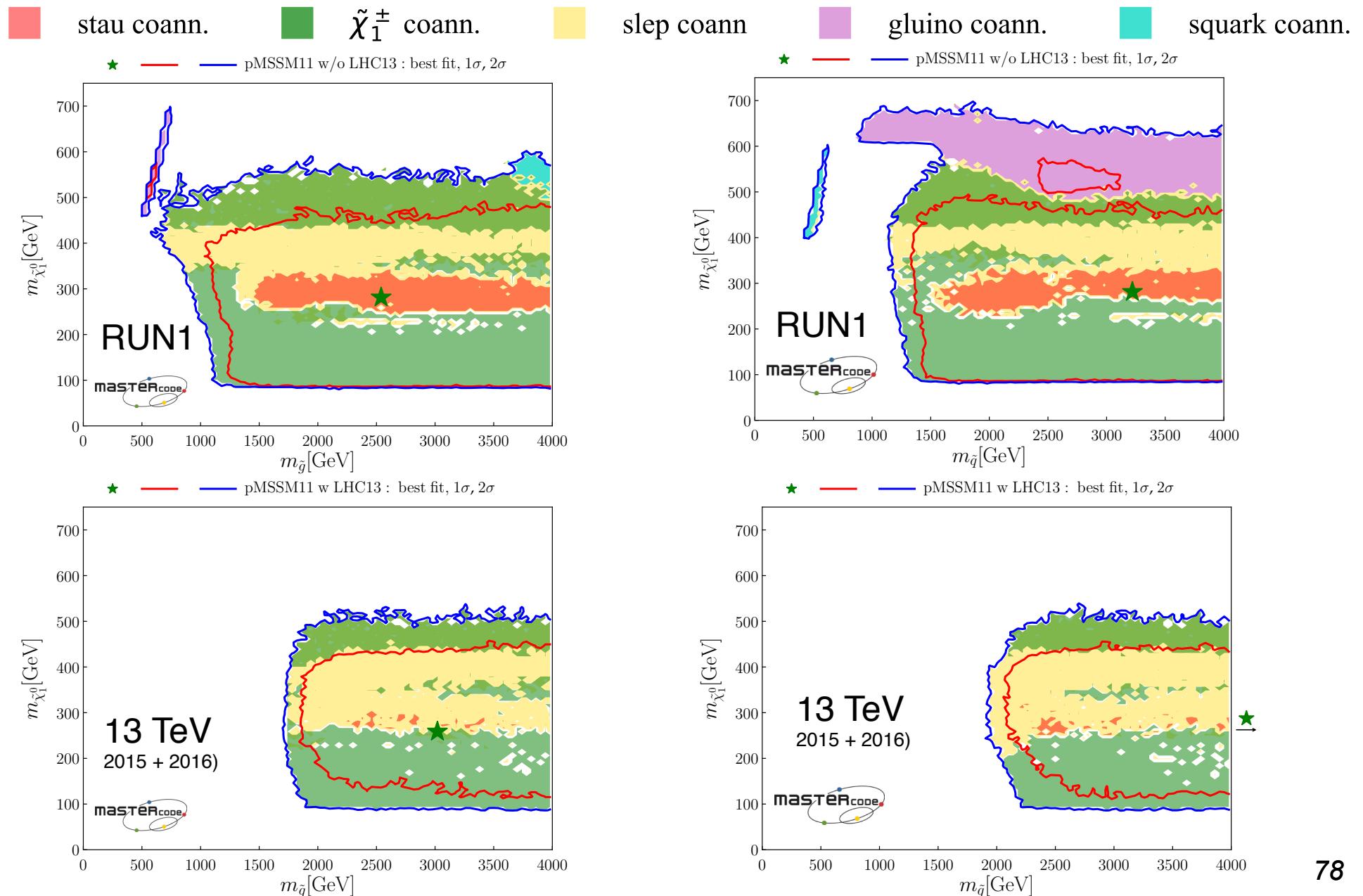
Comparison with Direct Detection Experiments



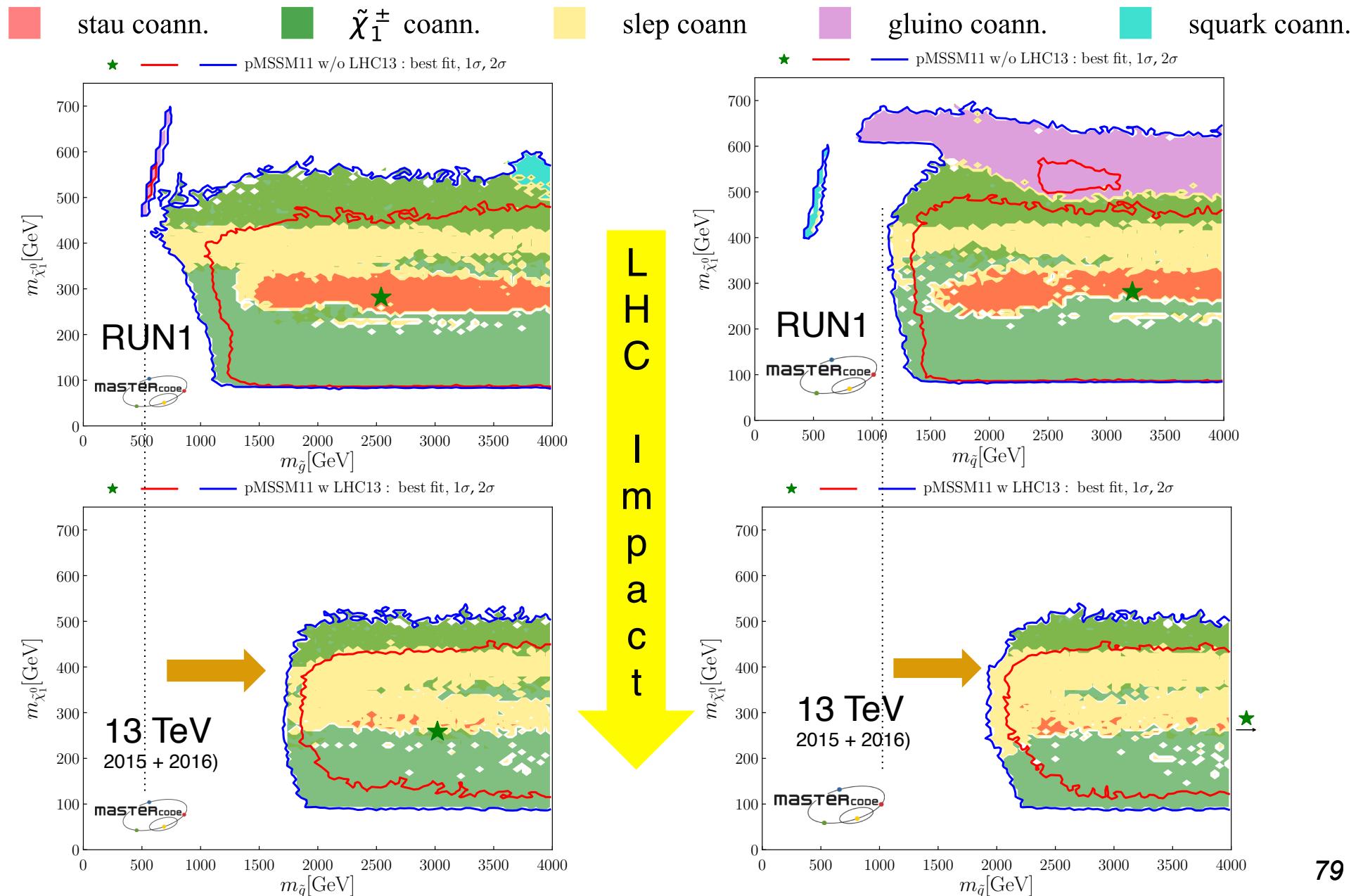
Comparison with Direct Detection Experiments



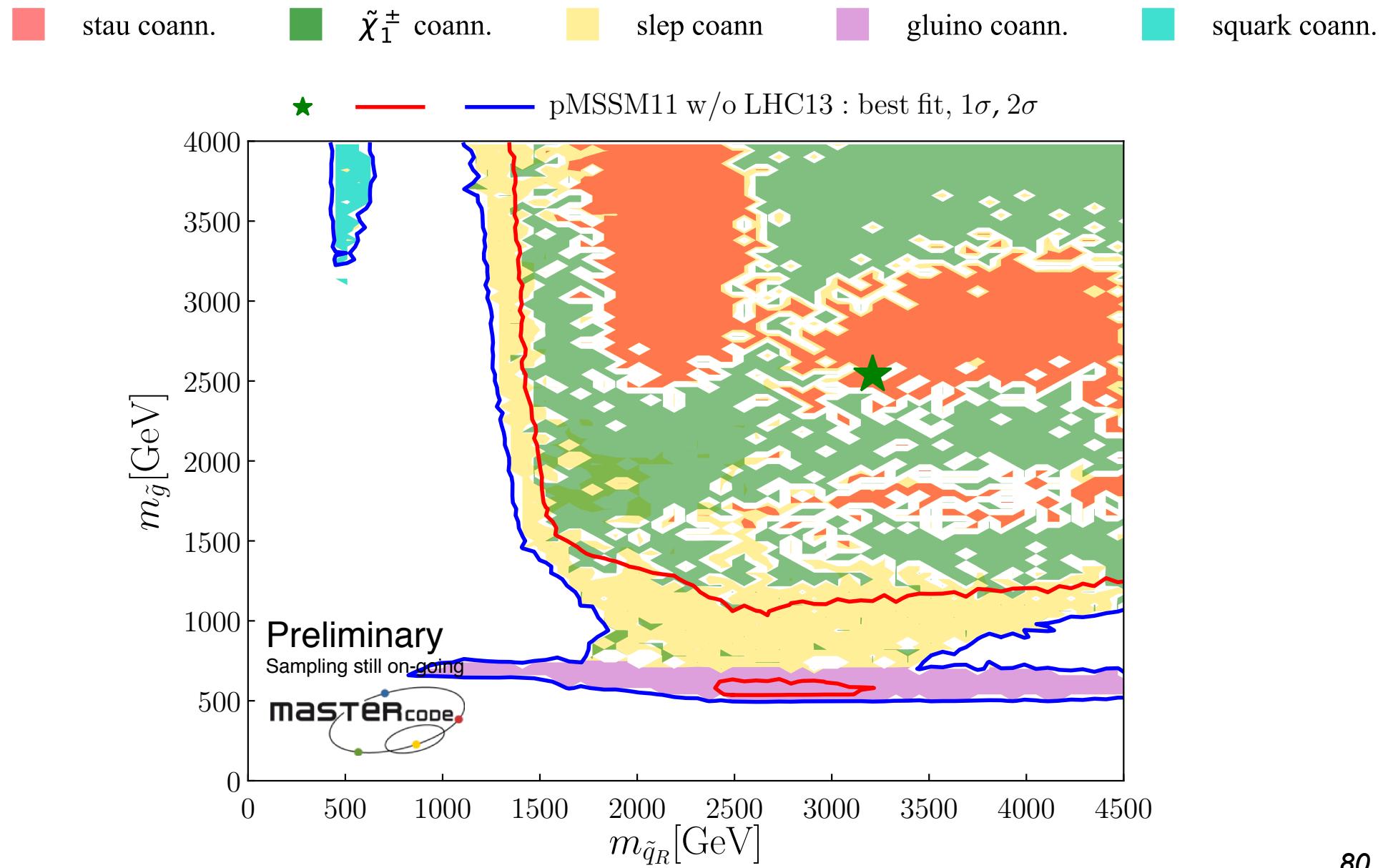
pMSSM11: RUN1 vs 13 TeV (2015 + 2016)



pMSSM11: RUN1 vs 13 TeV (2015 + 2016)



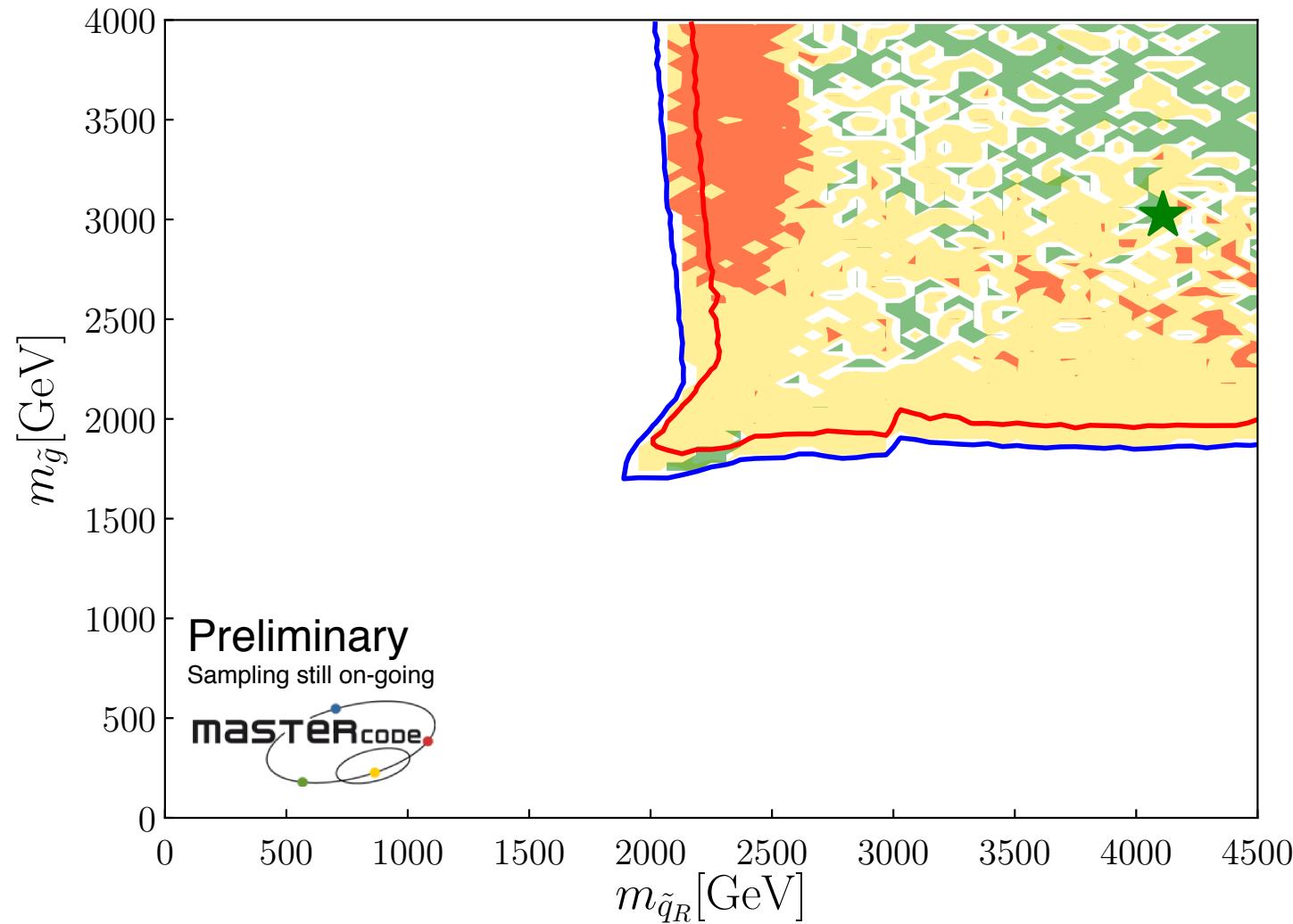
Gluino vs Squark: LHC RUN 1



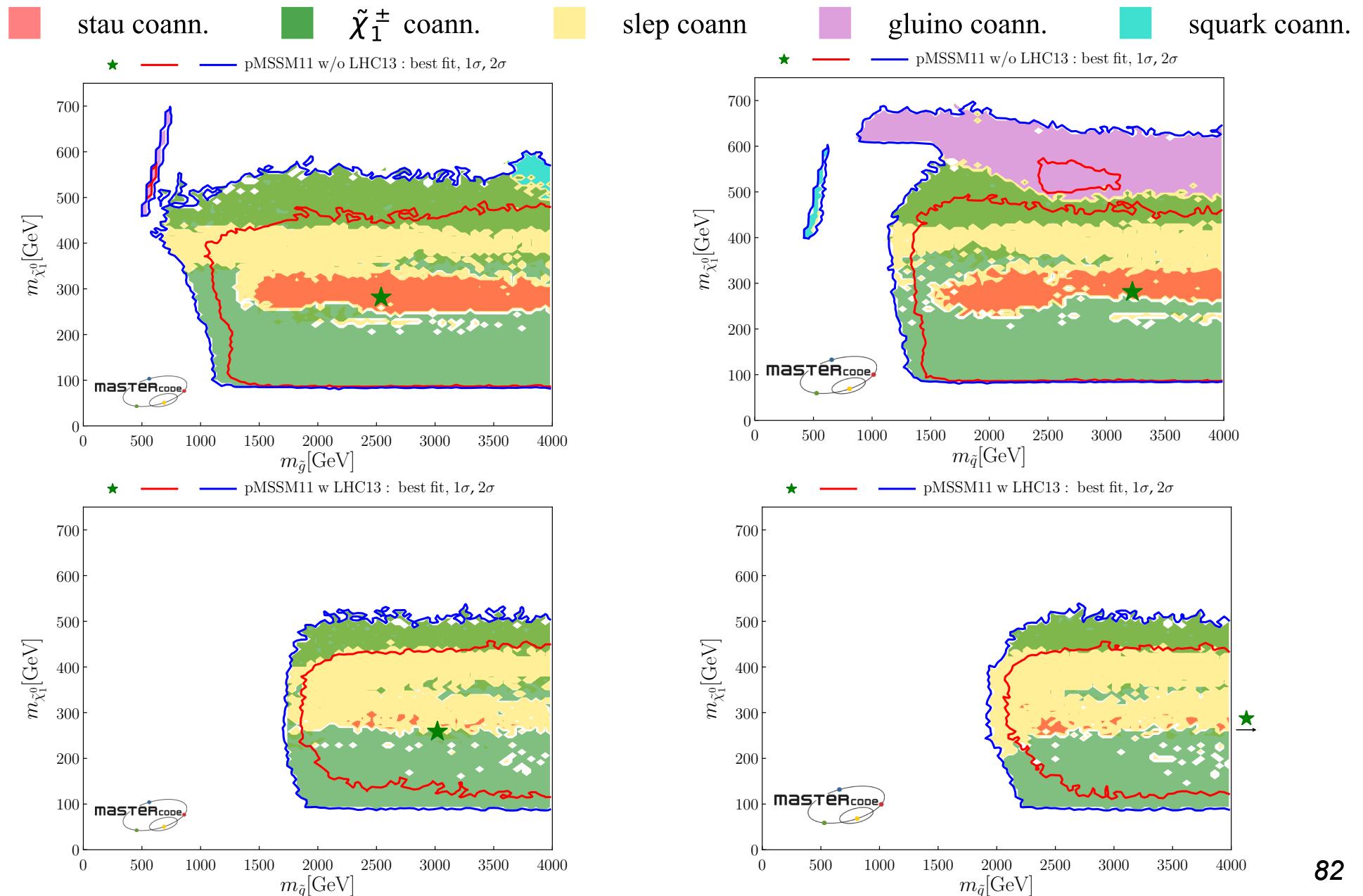
Gluino vs Squark: LHC RUN 2 (2015 + 2016 data)

■ stau coann. ■ $\tilde{\chi}_1^\pm$ coann. ■ slep coann. ■ gluino coann. ■ squark coann.

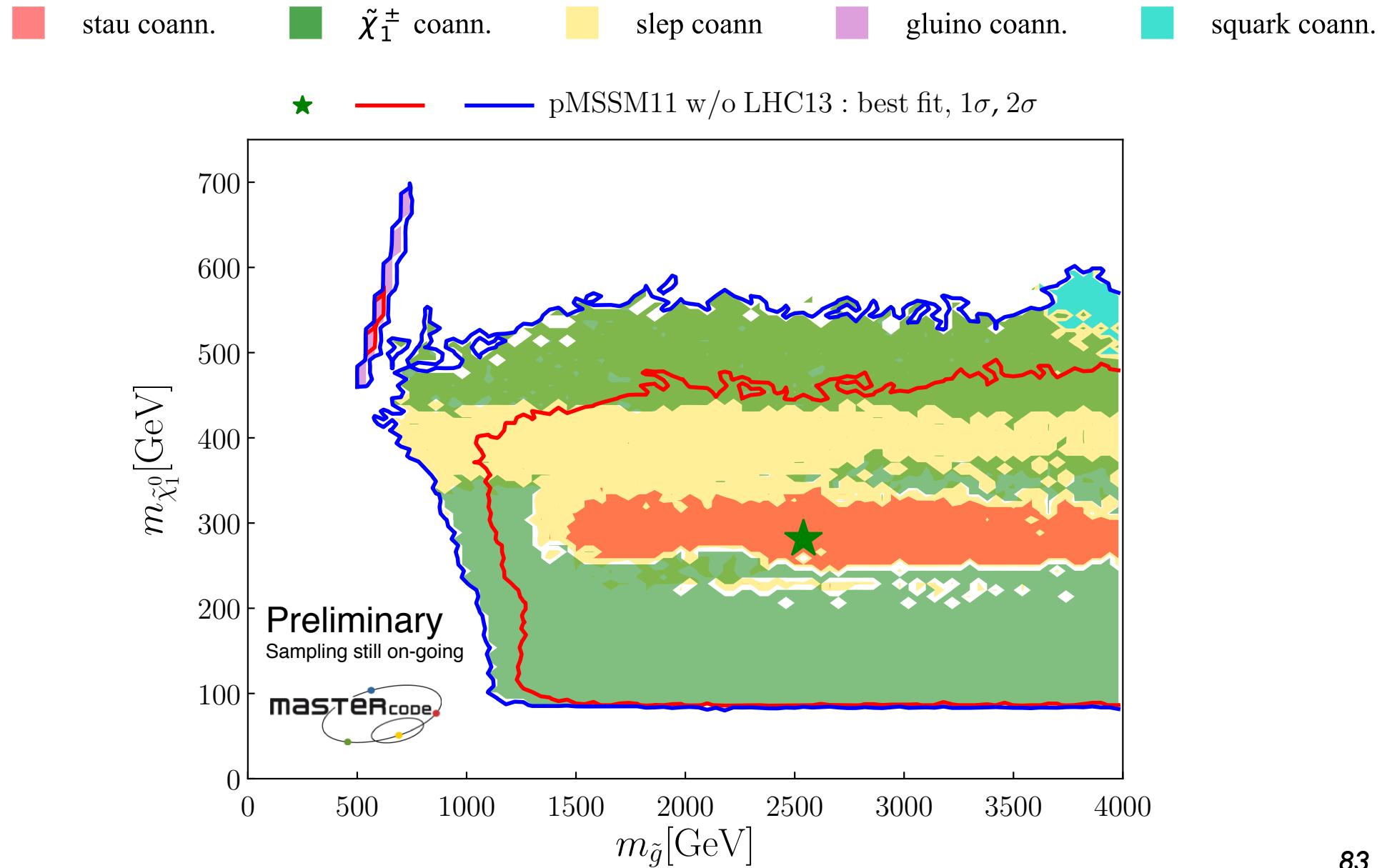
★ —— pMSSM11 w LHC13 : best fit, 1σ , 2σ



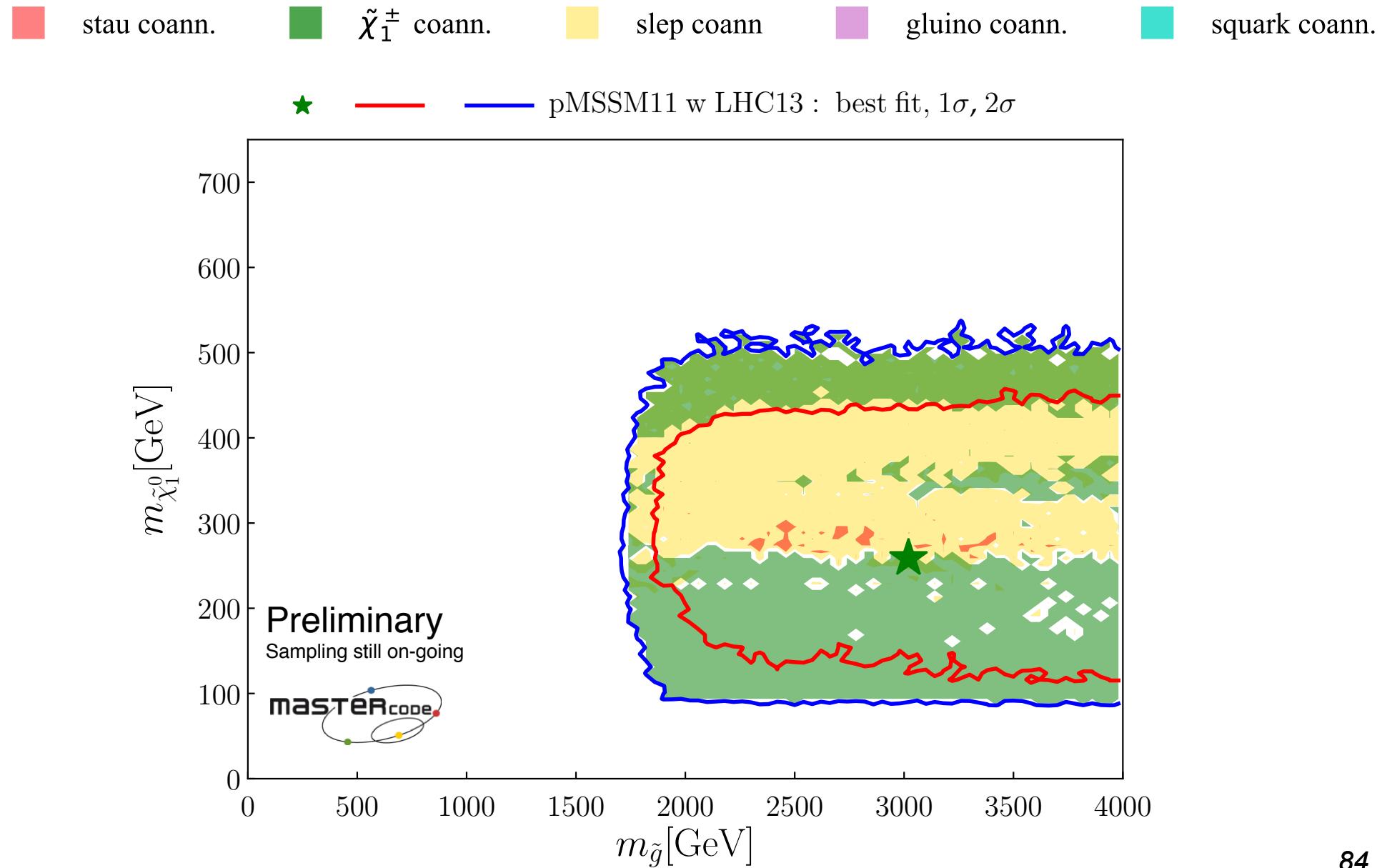
Gluino: LHC RUN 1



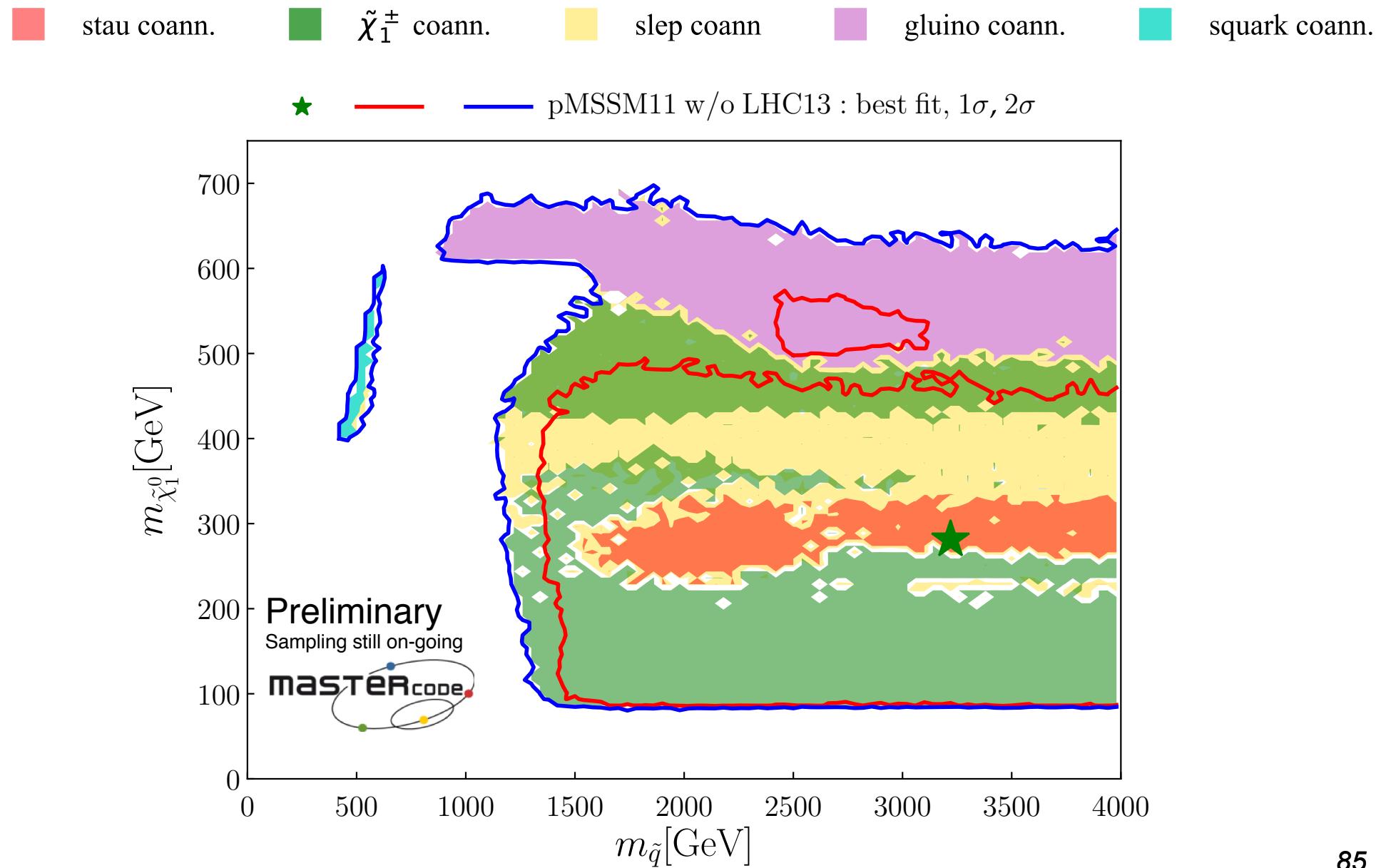
Gluino: LHC RUN 1



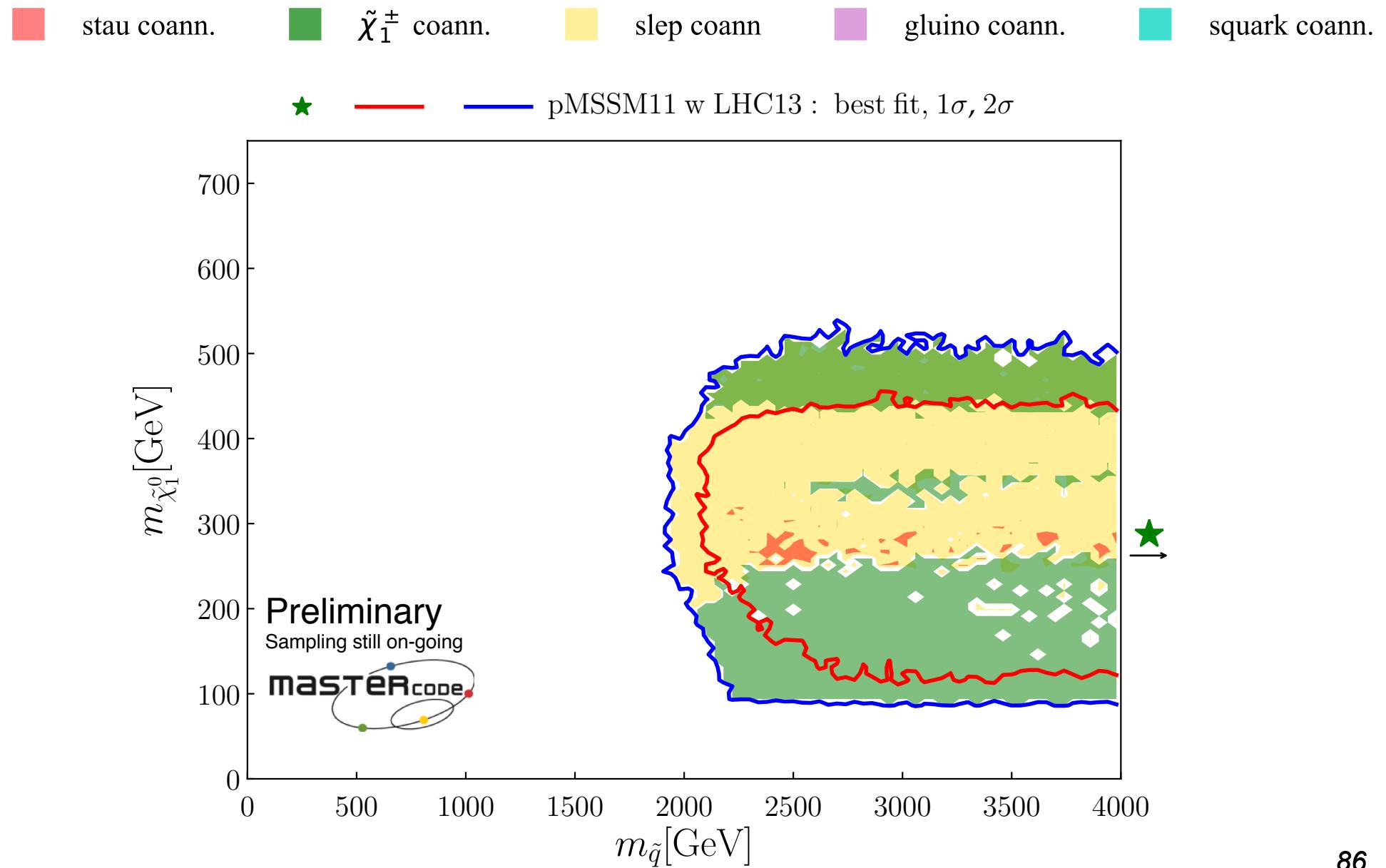
Gluino: LHC RUN 2 (2015 + 2016 data)



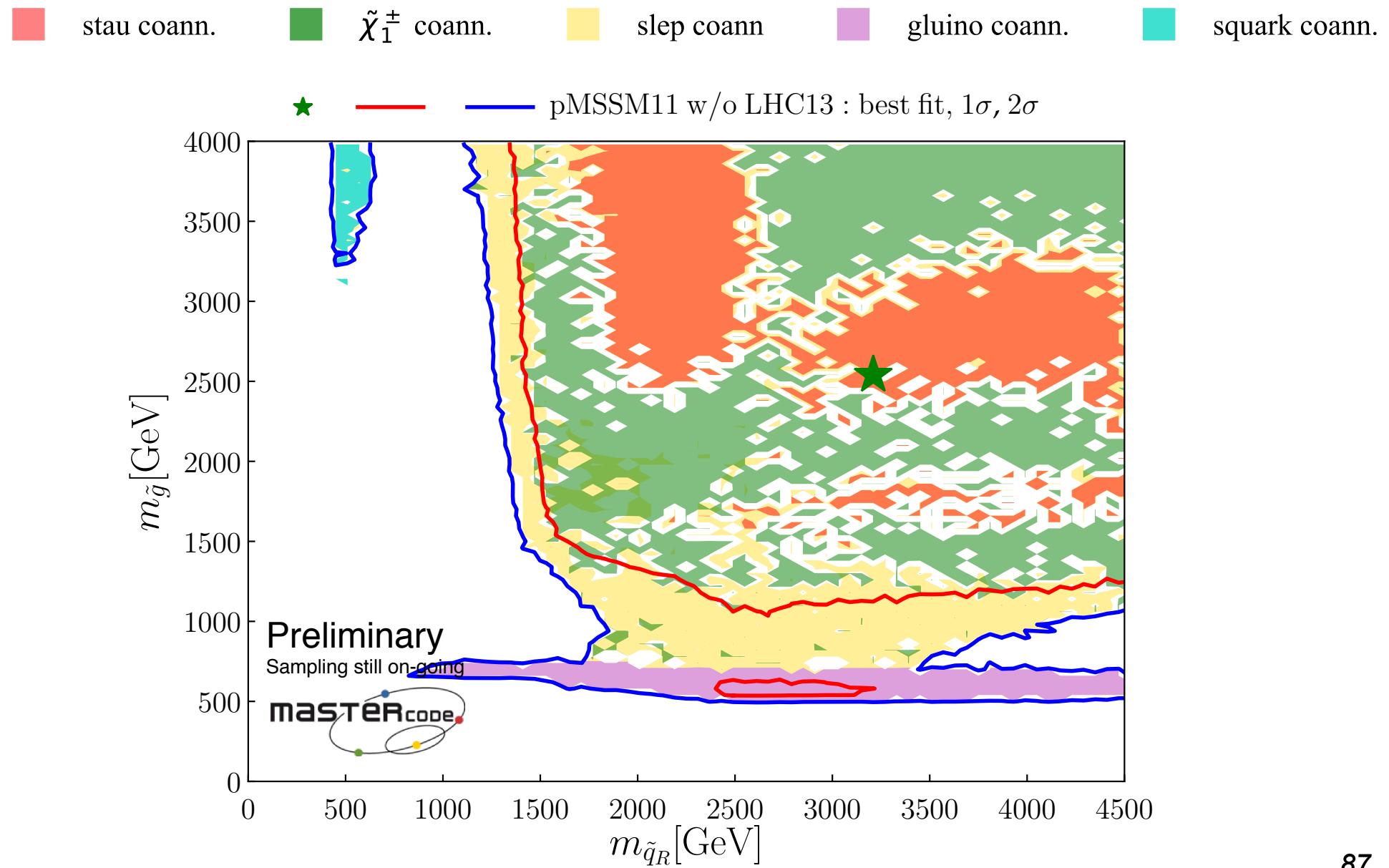
Squark: LHC RUN 1



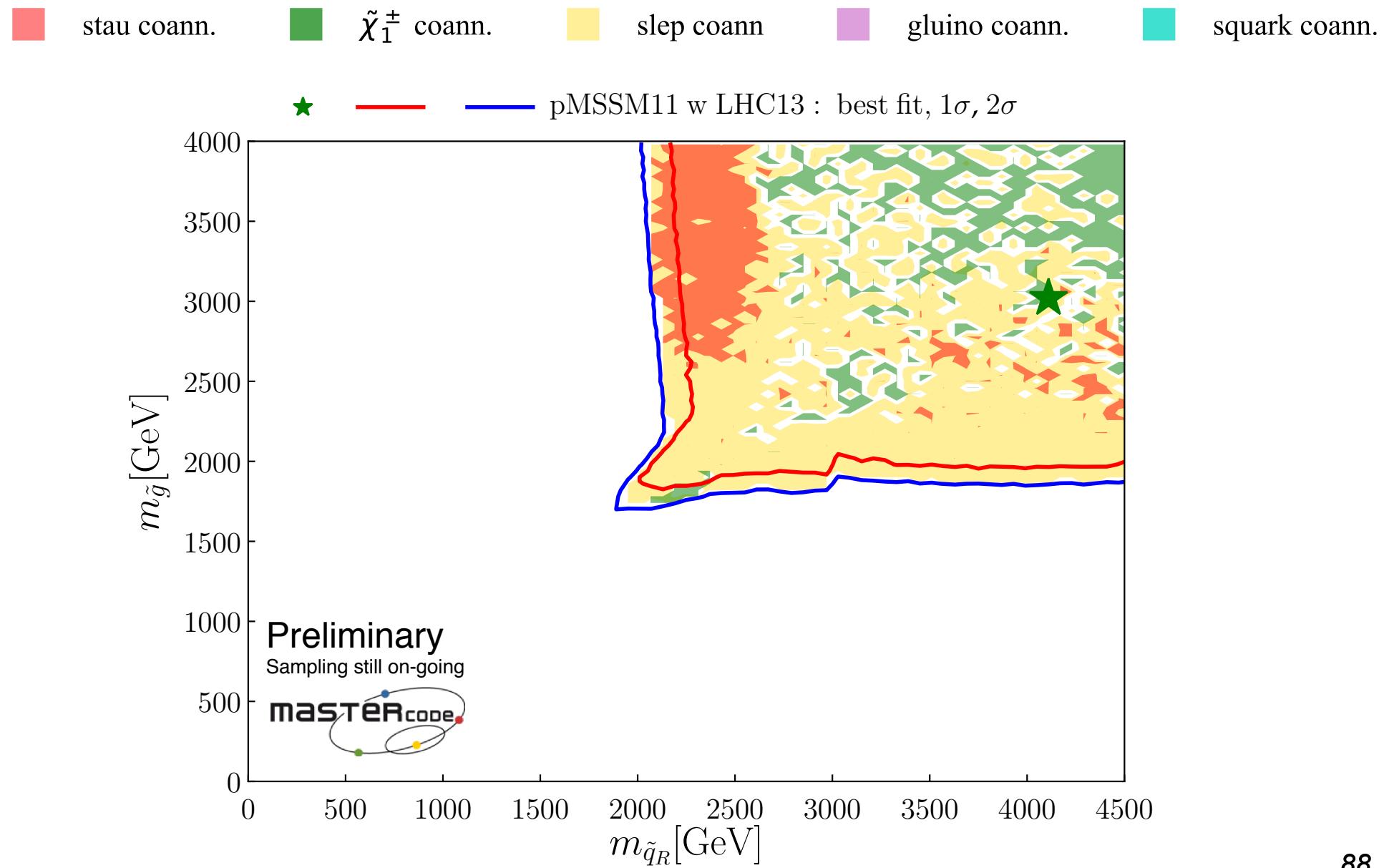
Squark: LHC RUN 2 (2015 + 2016 data)



Gluino vs Squark: LHC RUN 1

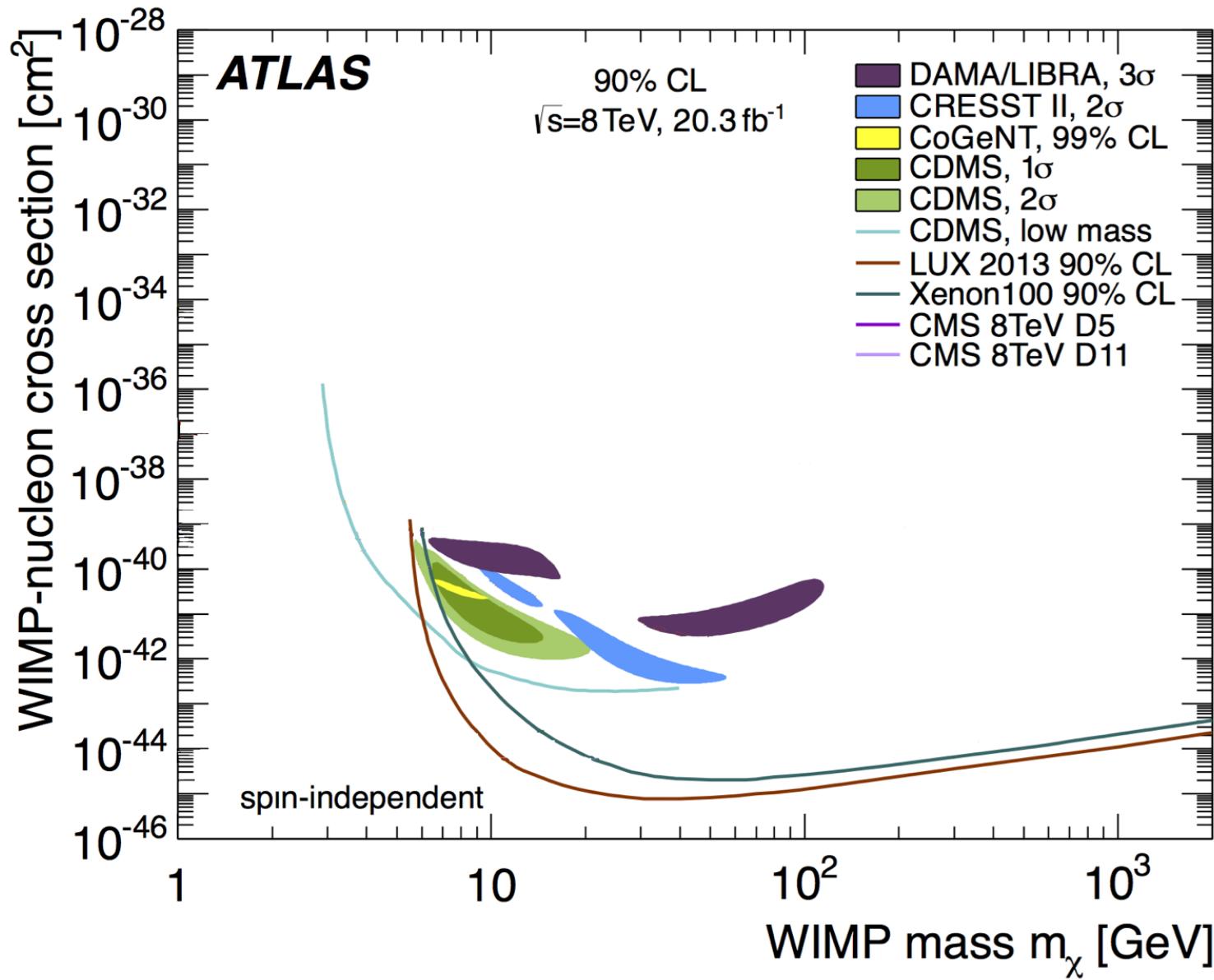


Gluino vs Squark: LHC RUN 2 (2015 + 2016 data)

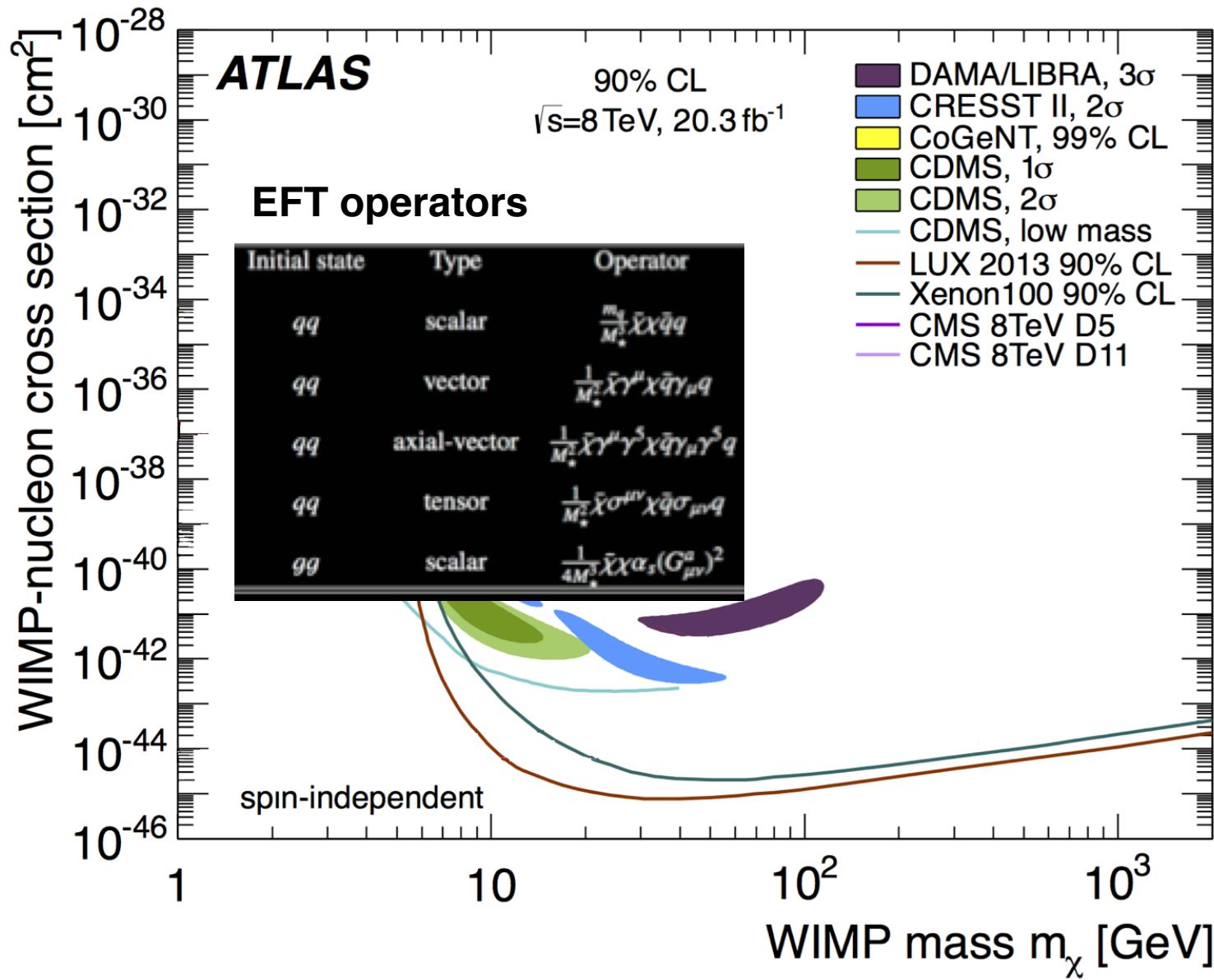


FROM EFT TO SIMPLIFIED MODELS

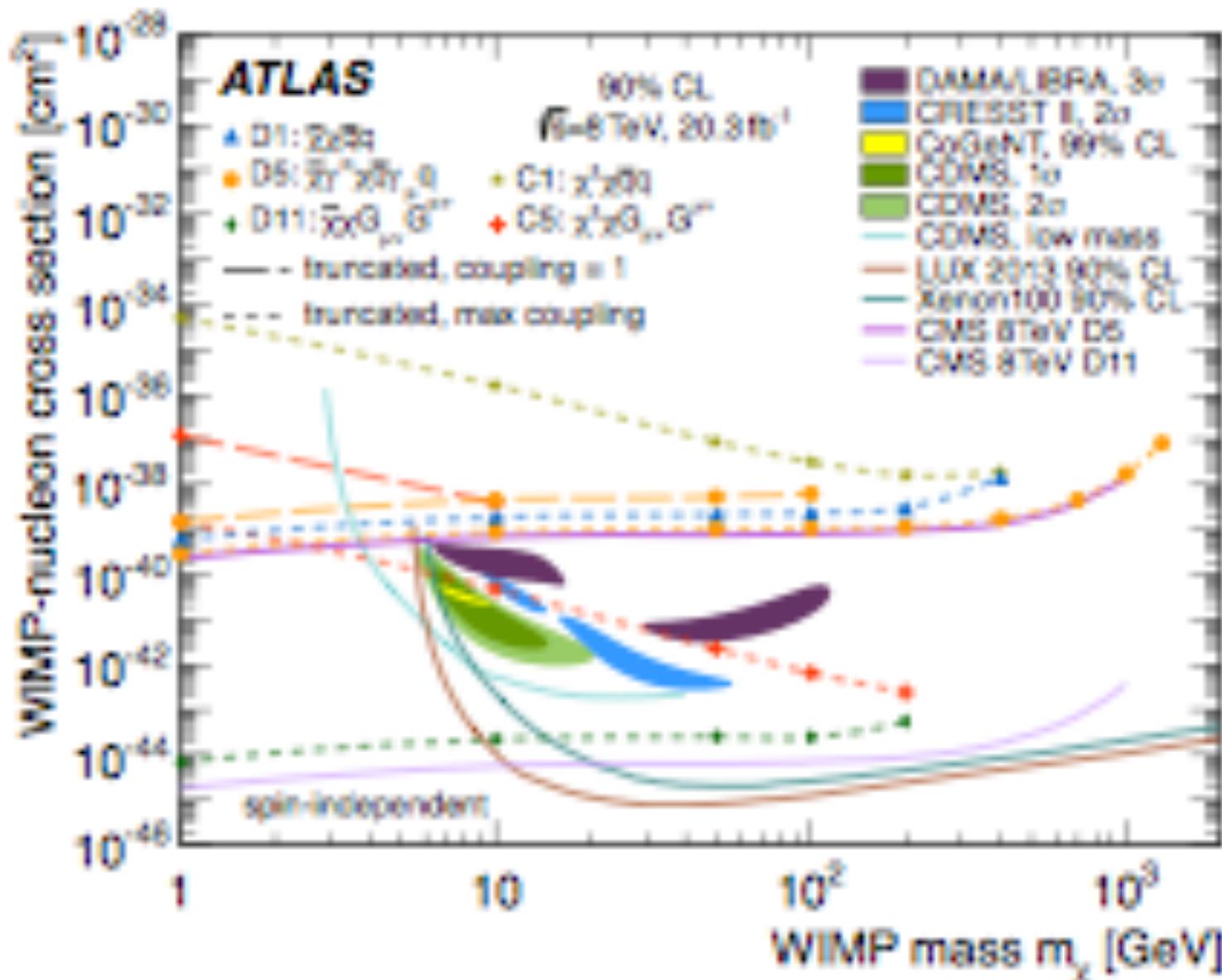
ATLAS Mono-Jet: Comparison with Direct Detection



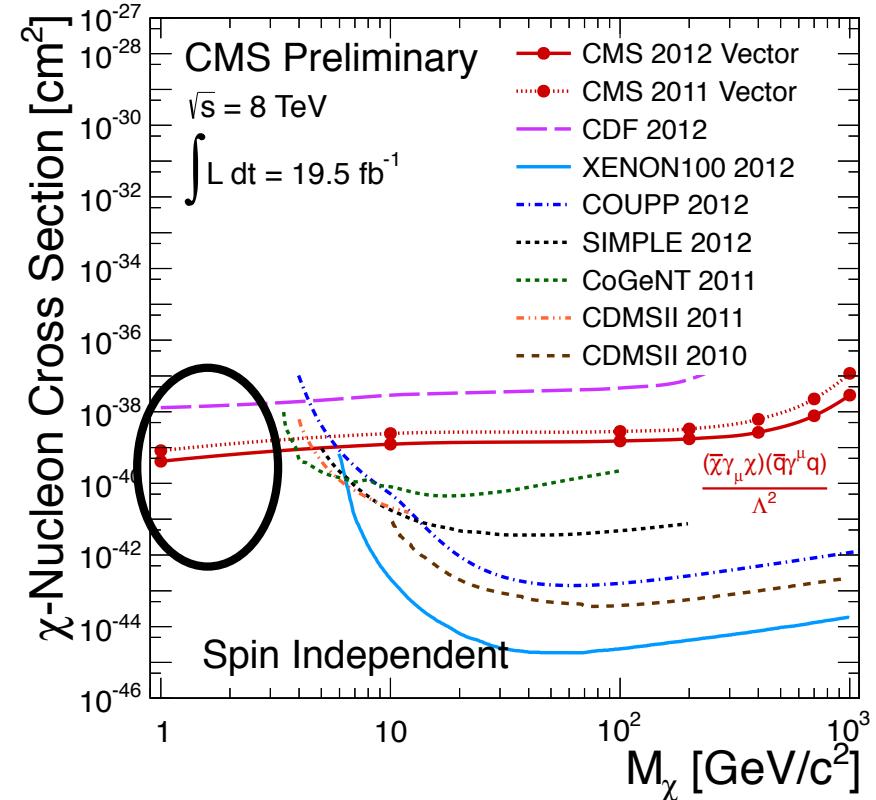
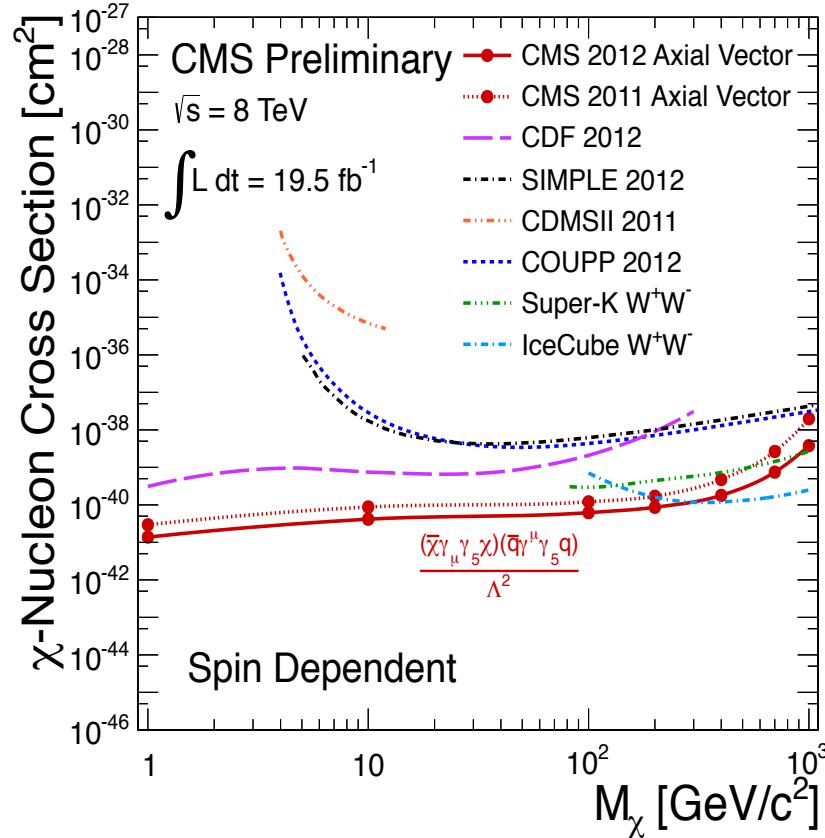
ATLAS Mono-Jet: Comparison with Direct Detection



ATLAS Mono-Jet: Comparison with Direct Detection



Mono-Jet analyses better than direct detection?!



Claim [often made]:

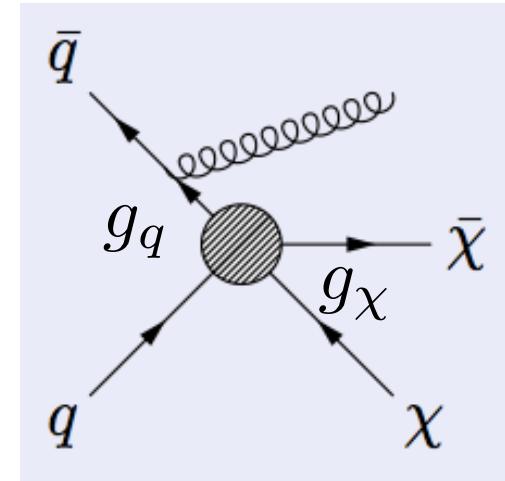
For **low mass** and the entire **spin-dependent** case monojet limits are stronger than direct detection limits!

Effective Field Theory (EFT) Interpretation

Example of considered operators:

$$O_V = \frac{(\bar{\chi} \gamma_\mu \chi)(\bar{q} \gamma_\mu q)}{\Lambda^2} \quad \text{Vector operator, s-channel}$$

$$O_{AV} = \frac{(\bar{\chi} \gamma_\mu \gamma_5 \chi)(\bar{q} \gamma_\mu \gamma_5 q)}{\Lambda^2} \quad \text{Axial vector operator, s-channel}$$

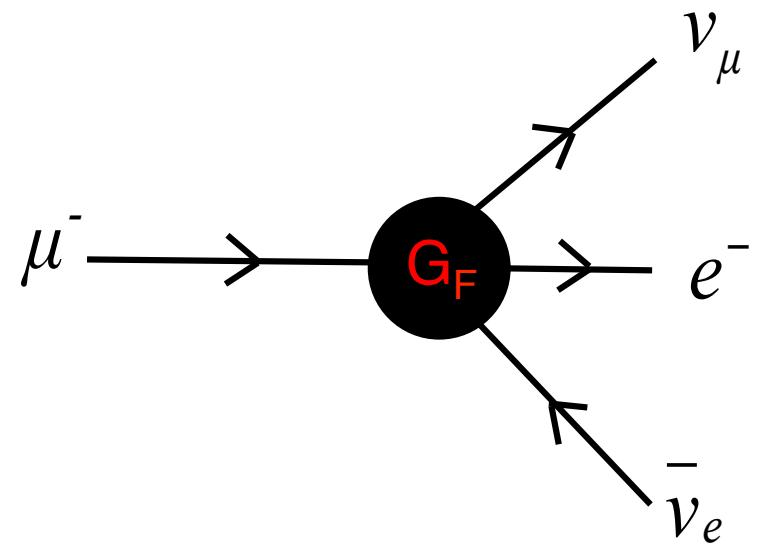


Assumption of EFT

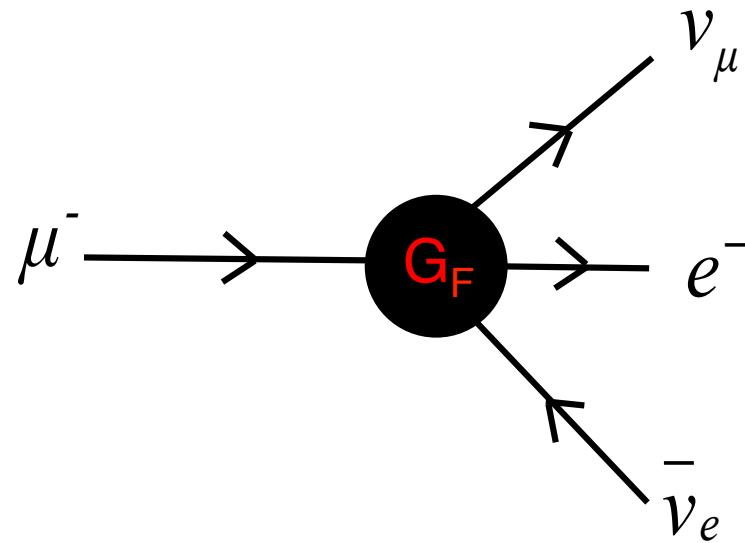
If the operator (e.g. V or AV) mediator is **suitably(!!)** heavy it can be integrated out to obtain the effective V or AV contact operator. **In this case (and only this case)**, the contact interaction scale Λ is related to the parameters entering the Lagrangian:

$$\Lambda = \frac{M_{mediator}}{\sqrt{g_q g_\chi}} \quad (\text{relation in the full theory})$$

Fermi Interaction & Muon Decay



Fermi Interaction & Muon Decay



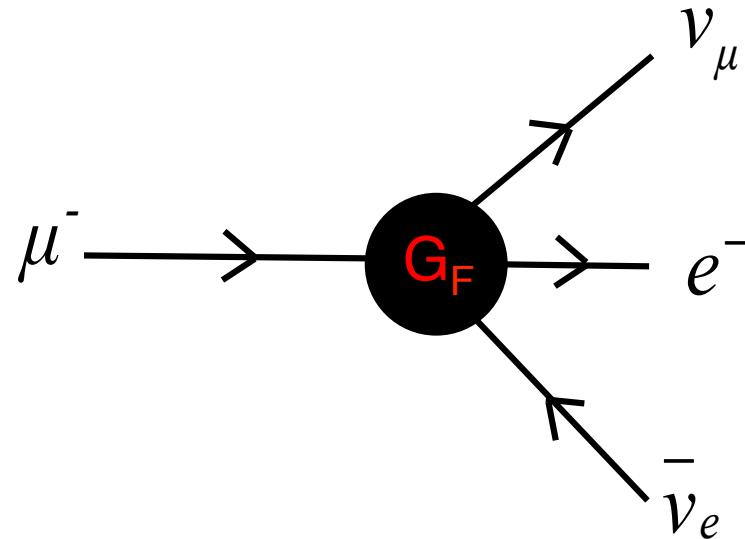
The Fermi 4-point interaction was able to explain well the beta-decay as well as the muon decay with one single interact strengths G_F (Fermi constant)

However, the cross-section grows as the square of the energy:

$$\sigma \propto G_F^2 E^2$$

making it invalid for higher energies!

Fermi Interaction & Muon Decay

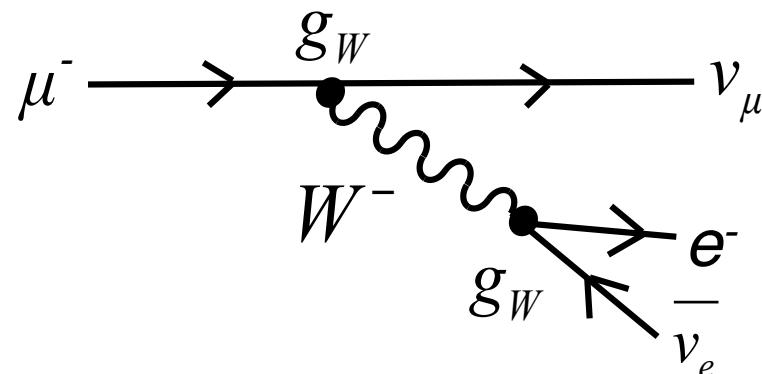


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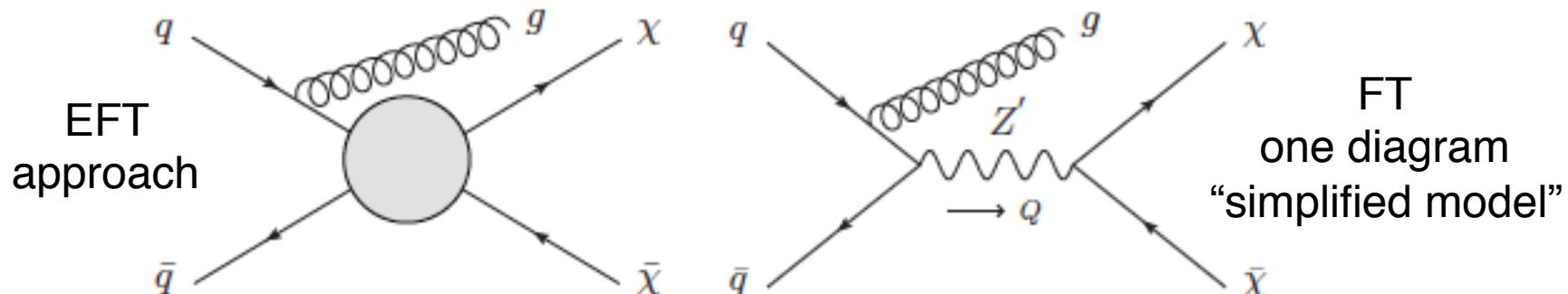
Solution:
Resolve the “blob” and replace the 4-point interaction with an ultraviolet complete theory!

$$\sigma \propto G_F^2 M_W^2$$

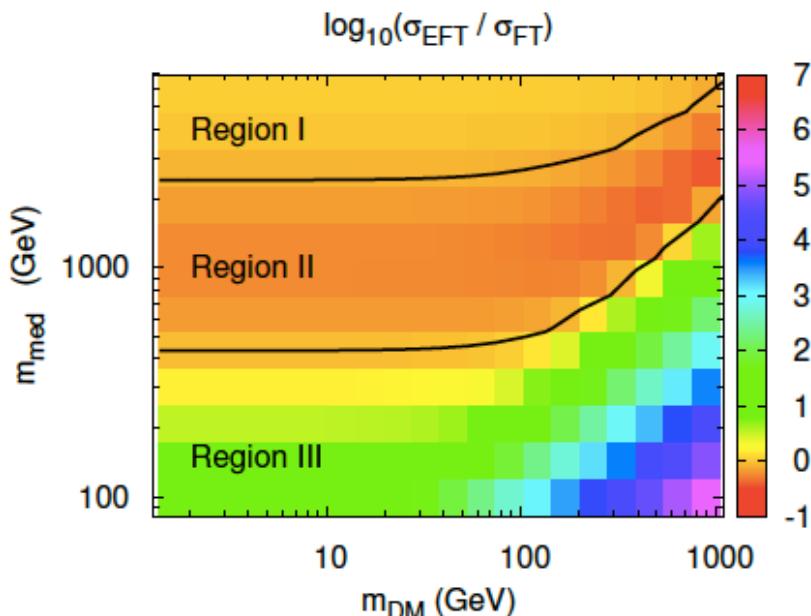
Validity of Effective Field Theory Limits

Recent work from OB, M.Dolan,C.McCabe: arXiv:1308.6799

- Compare Effective Field Theory (EFT) with Full Theory (FT)



Use vector and axial-vector mediators (e.g. Z') as example - scalar are similar in conclusion!



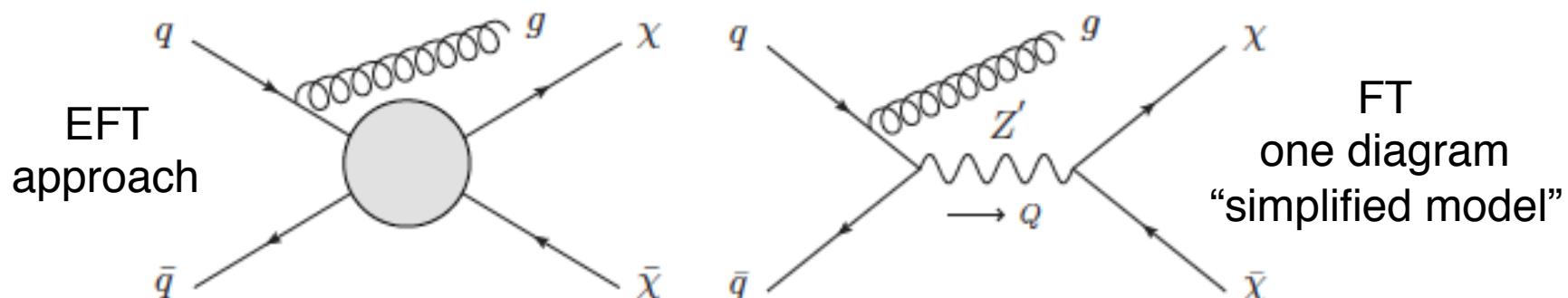
Compare prediction of FT with EFT in $m_{\text{med}} - m_{\text{DM}}$ plane.
Three regions become visible:

- Region I: EFT and FT agree better than 20%
➤ EFT is valid!
- Region II: EFT yields significant weaker limits than FT
➤ EFT limits are too conservative!
- Region III: EFT yields significant stronger limits than FT
➤ EFT limits are too aggressive!

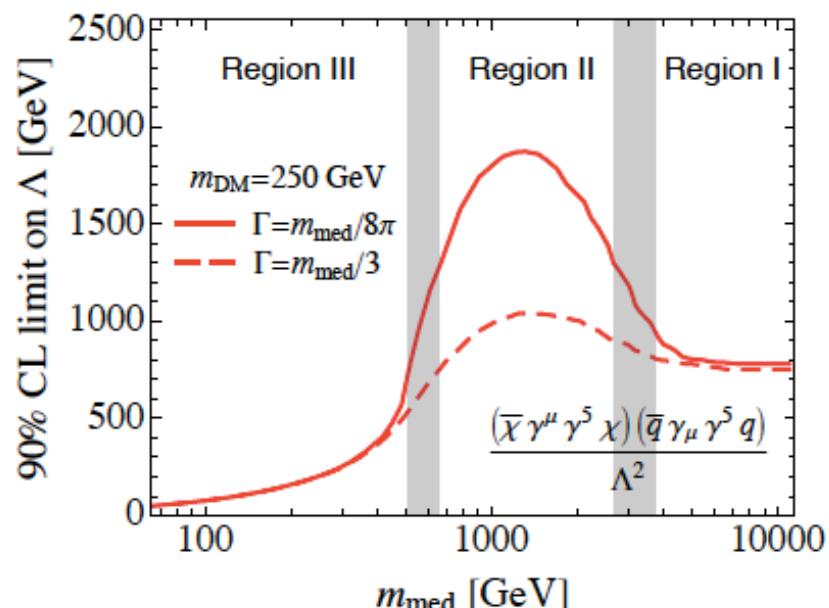
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Use vector and axial-vector mediators (e.g. Z') as example - scalar are similar in conclusion!



Three Regions as function of mediator mass:

Region I: Heavy m_{med}

- EFT is valid!

Region II: Medium m_{med} – Resonant enhancement

- EFT limits are too conservative!

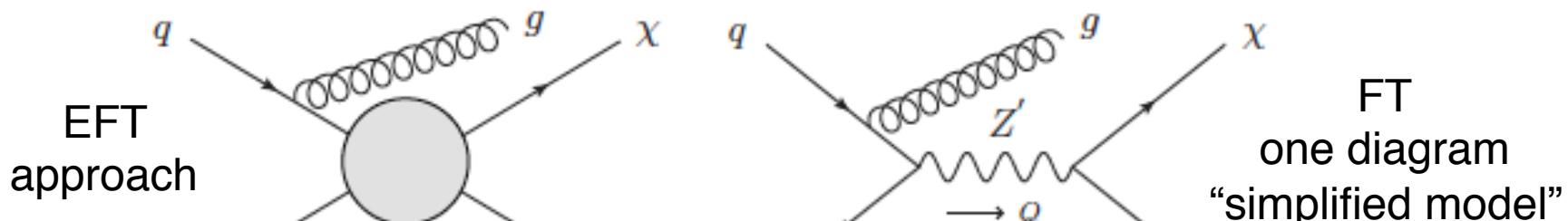
Region III: Low m_{med}

- EFT limits are too aggressive!

Validity of Effective Field Theory Limits

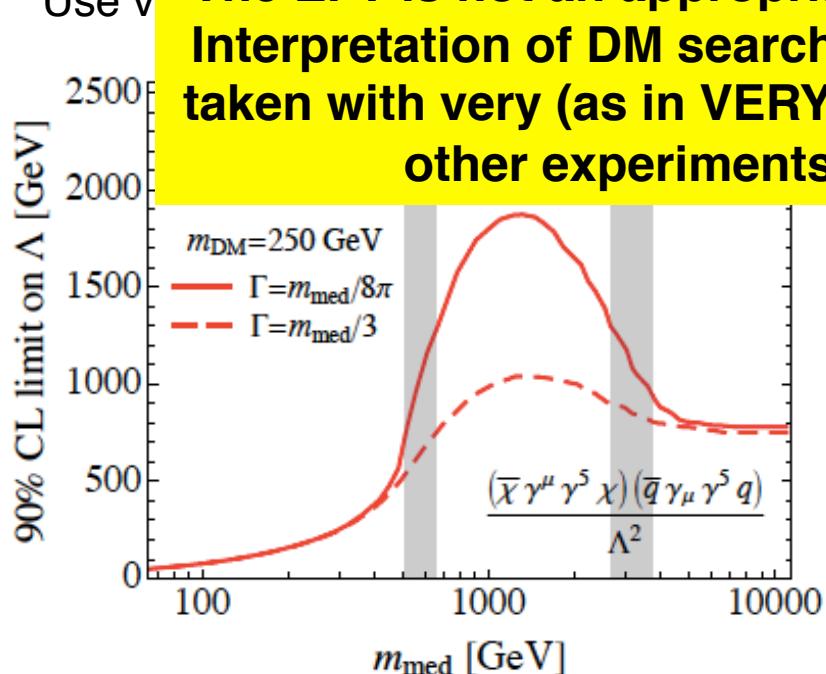
Recent work from OB, M.Dolan,C.McCabe: arXiv:1308.6799

- Compare Effective Field Theory (EFT) with Full Theory (FT)



Conclusion:

The EFT is not an appropriate framework for a comprehensive Interpretation of DM searches at colliders and especially must taken with very (as in VERY) special care when comparing with other experiments such as Direct Detection!



Region I: Heavy m_{med}

- EFT is valid!

Region II: Medium m_{med} – Resonant enhancement

- EFT limits are too conservative!

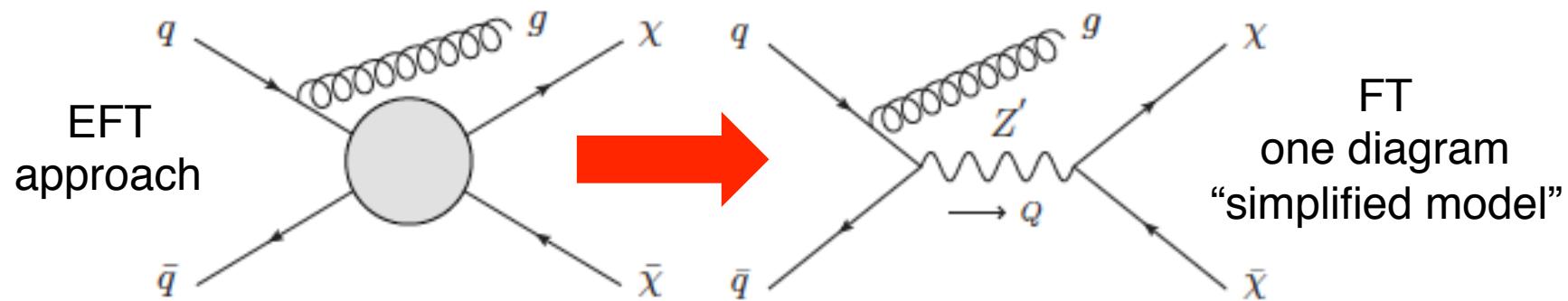
Region III: Low m_{med}

- EFT limits are too aggressive!

Alternative Interpretation Ansatz: Simplified models

Recent work from OB, M.Dolan,C.McCabe: arXiv:1308.6799

- Compare Effective Field Theory (EFT) with Full Theory (FT)

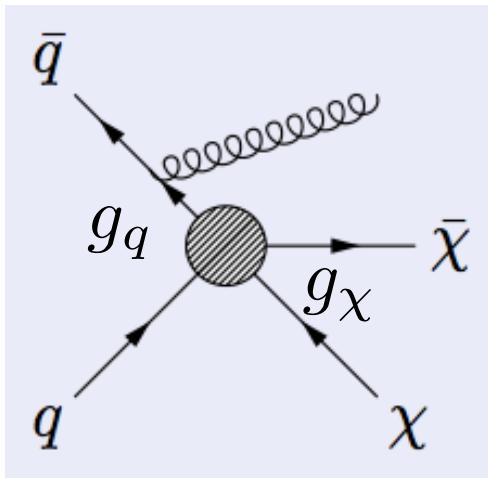
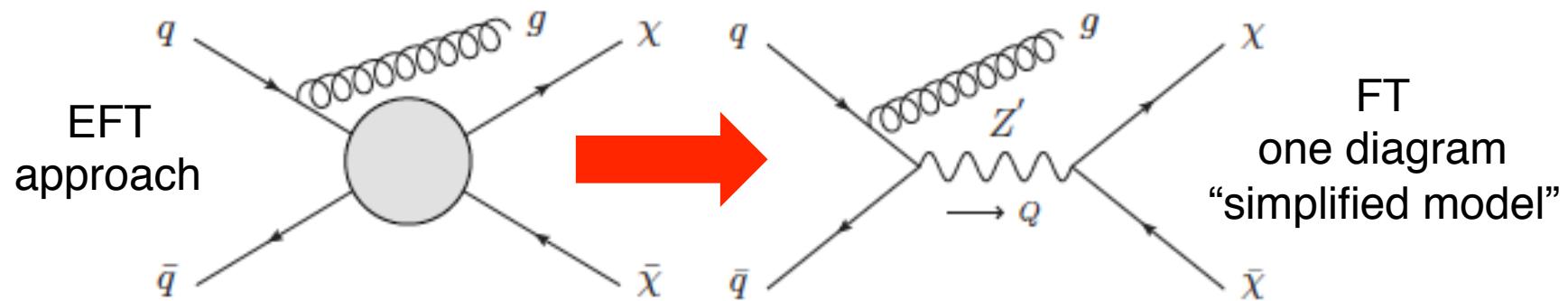


After three years of operation at the LHC the landscape for interpretation of searches has changed dramatically – new superior & modern approaches have replaced in many areas longstanding traditional ones (e.g. SUSY searches)

Alternative Interpretation Ansatz: Simplified models

Recent work from OB, M.Dolan,C.McCabe: arXiv:1308.6799

- Compare Effective Field Theory (EFT) with Full Theory (FT)



The problem is governed by five variables:

- Couplings g_q and g_χ
- Mediator mass m_{med} and mediator width Γ_{med}
- Dark matter candidate mass m_{DM}

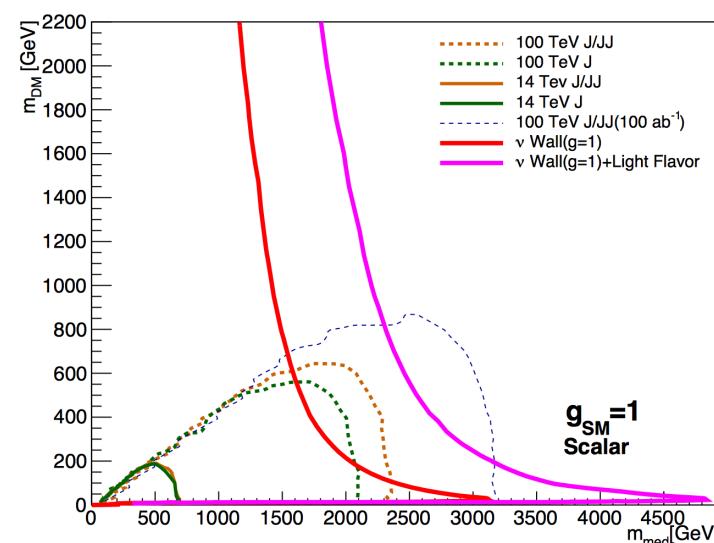
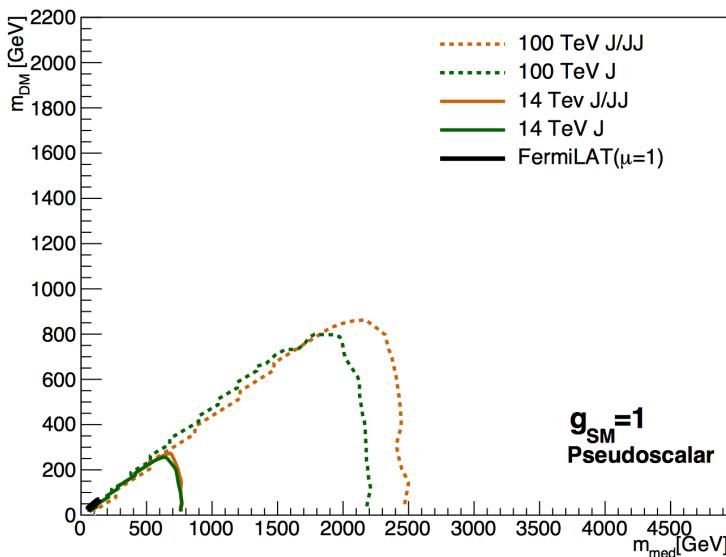
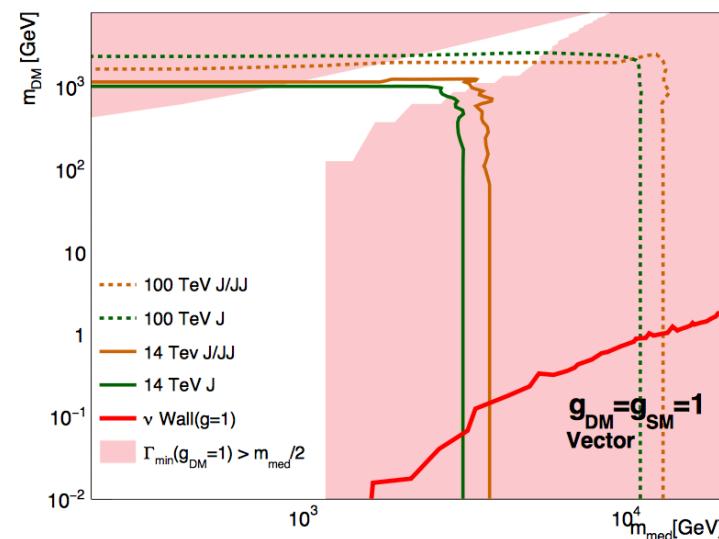
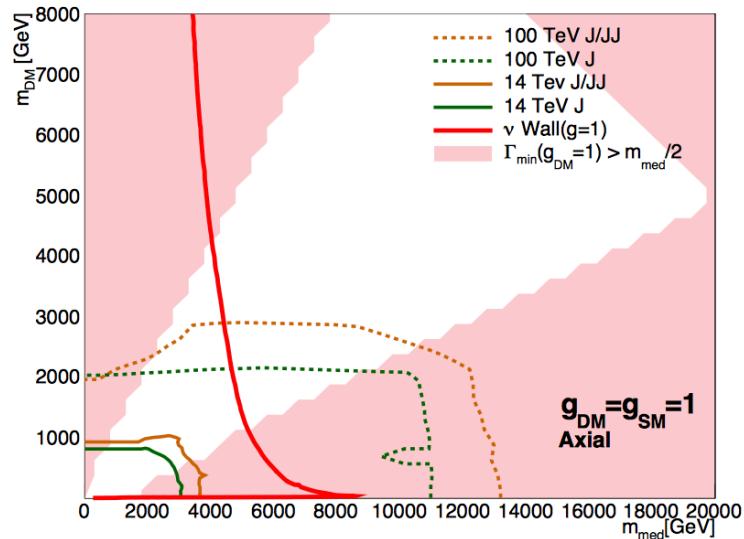
ATLAS & CMS public results

Most results presented in this talk (and many more) can be accessed via the public page of the ATLAS and CMS experiments:

ATLAS SUSY: <https://twiki.cern.ch/twiki/bin/view/AtlasPublic/SupersymmetryPublicResults>

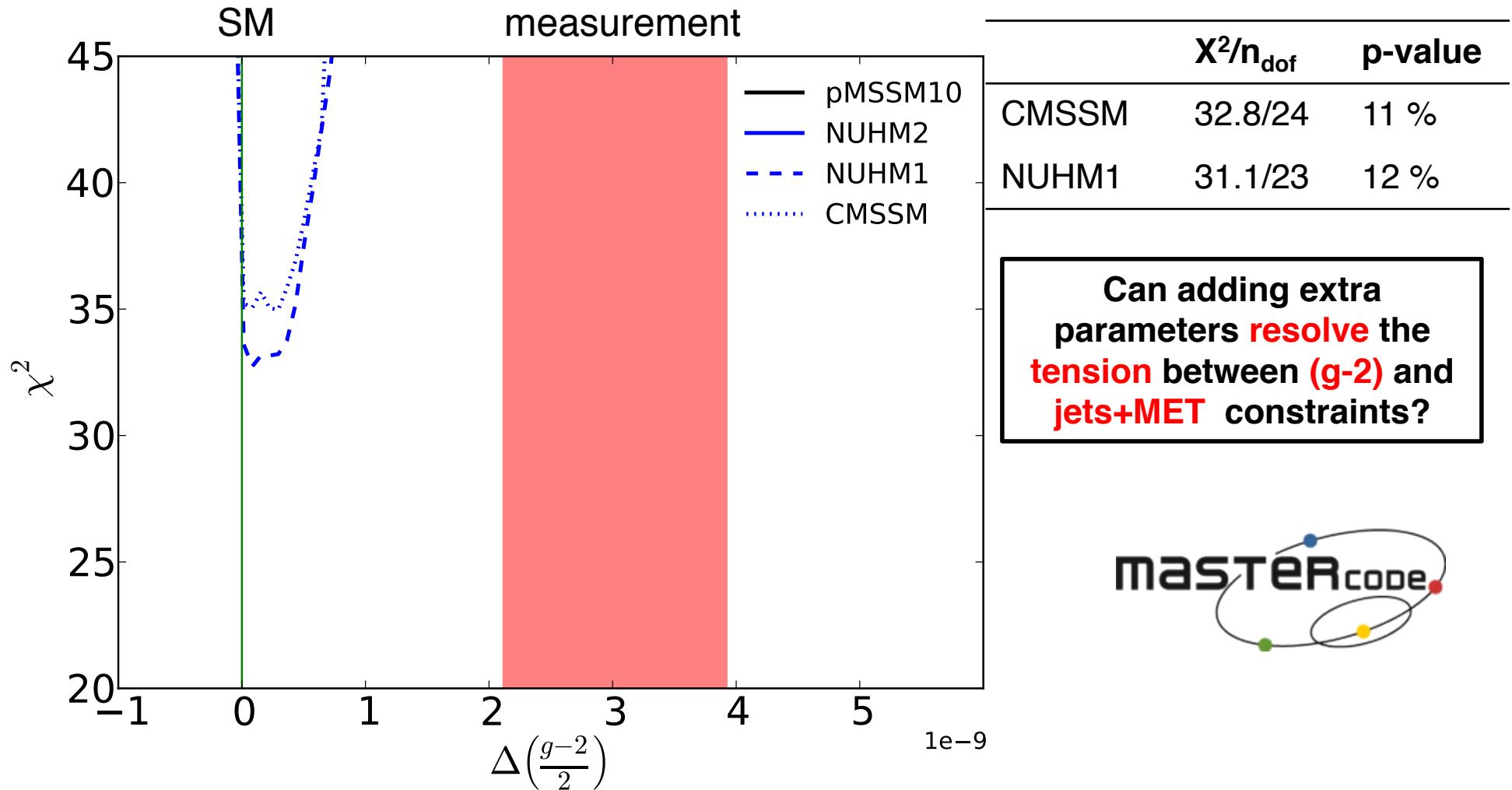
CMS SUSY :<https://twiki.cern.ch/twiki/bin/view/CMSPublic/PhysicsResultsSUS>

100 TeV Prediction from arXiv:1509.02904



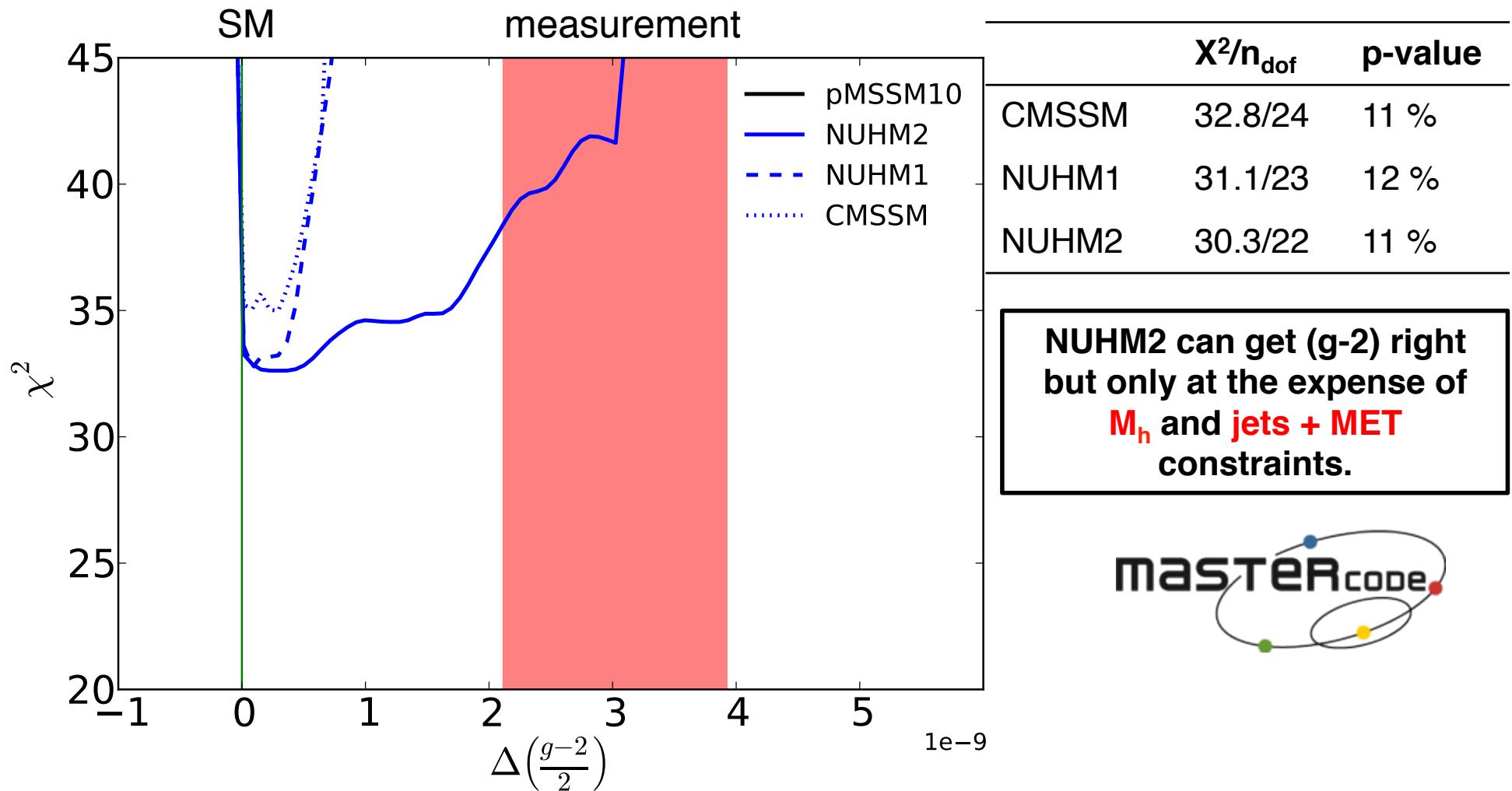
MASTERCODE

Resolving tension ($g-2$) and LHC



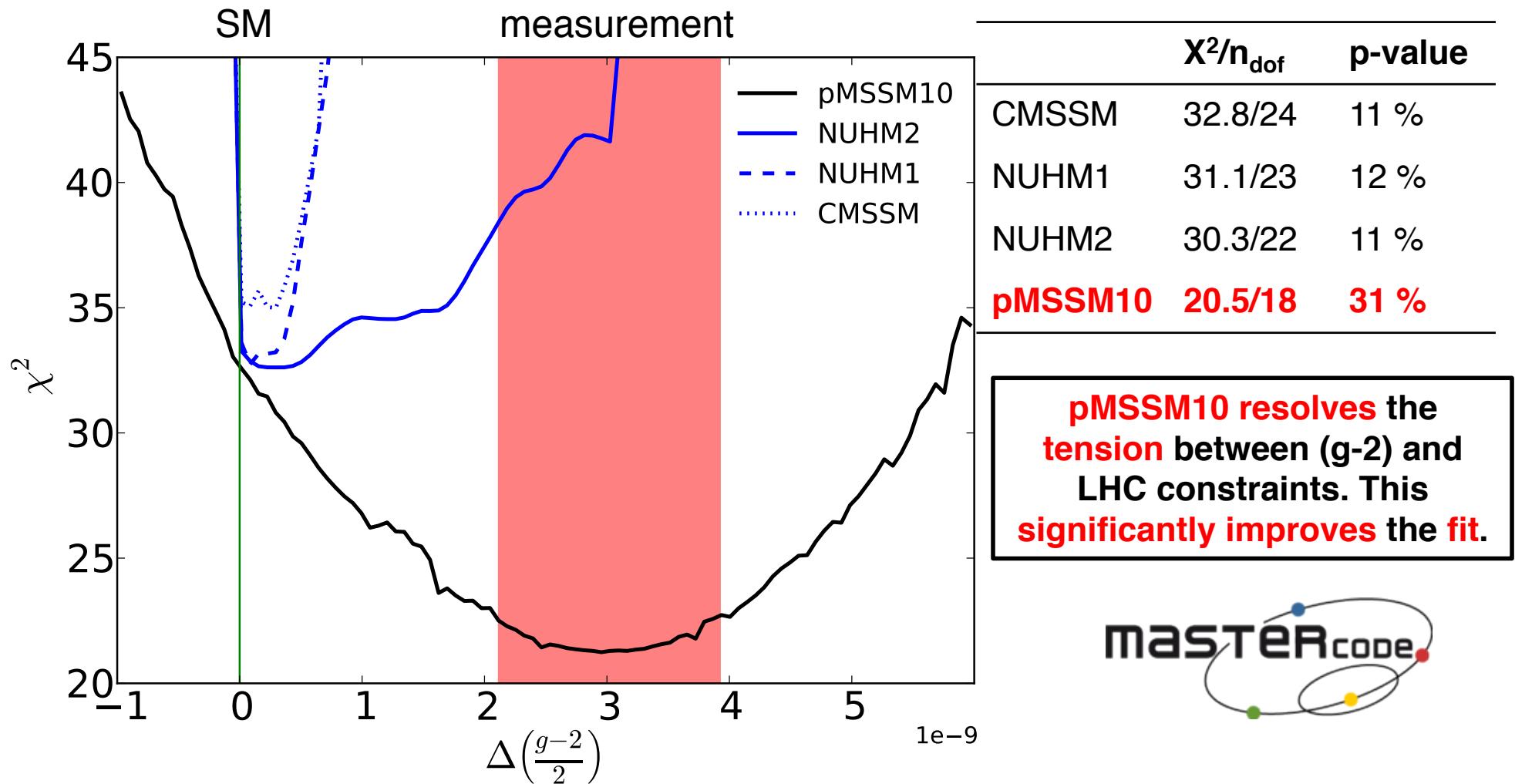
From MasterCode papers:
1312.5250, 1408.4060 and 1504.03260

Resolving tension ($g-2$) and LHC



From MasterCode papers:
1312.5250, 1408.4060 and 1504.03260

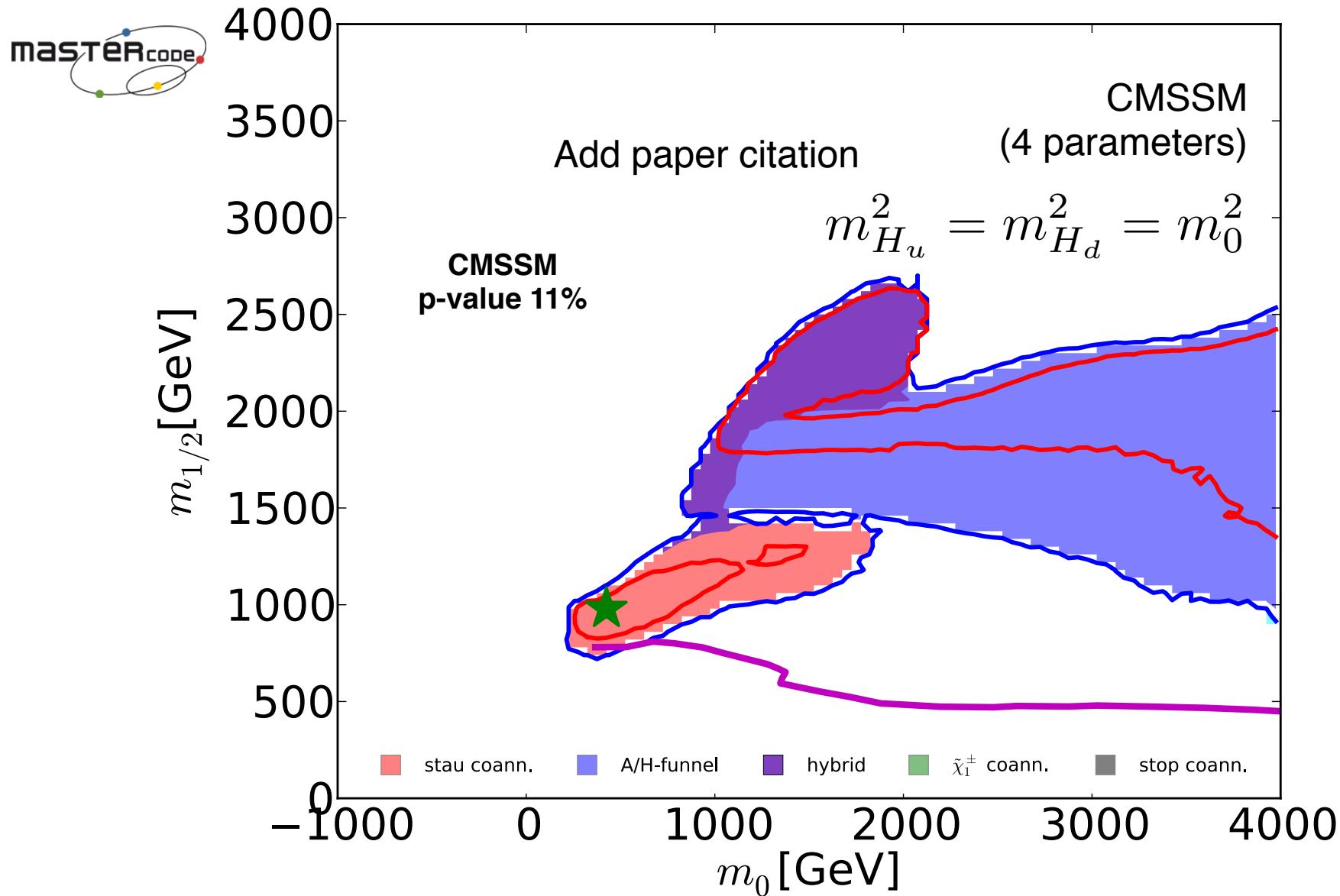
Resolving tension ($g-2$) and LHC



From MasterCode papers:
1312.5250, 1408.4060 and 1504.03260

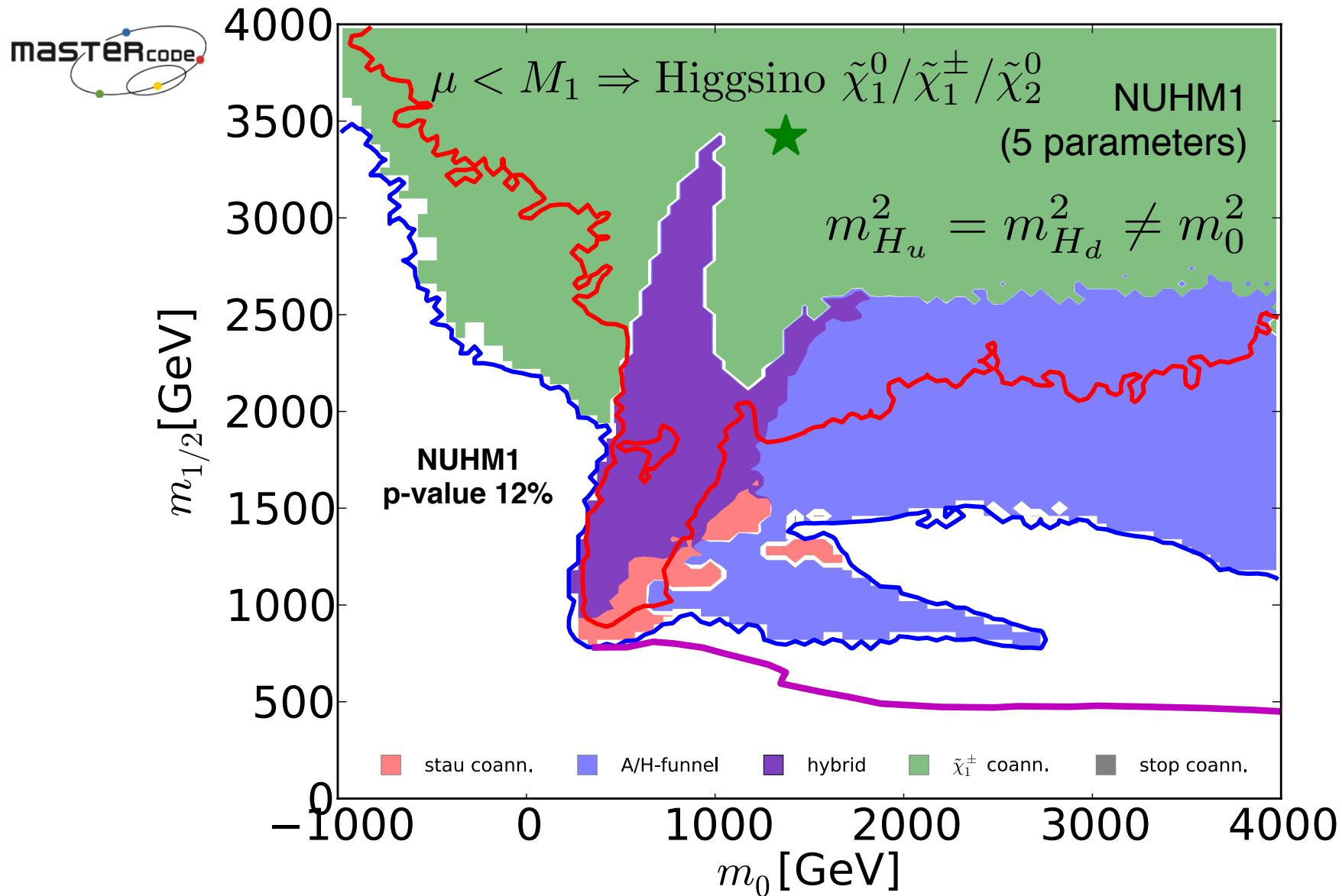
CMSSM

★ — CMSSM: best fit, 1σ , 2σ



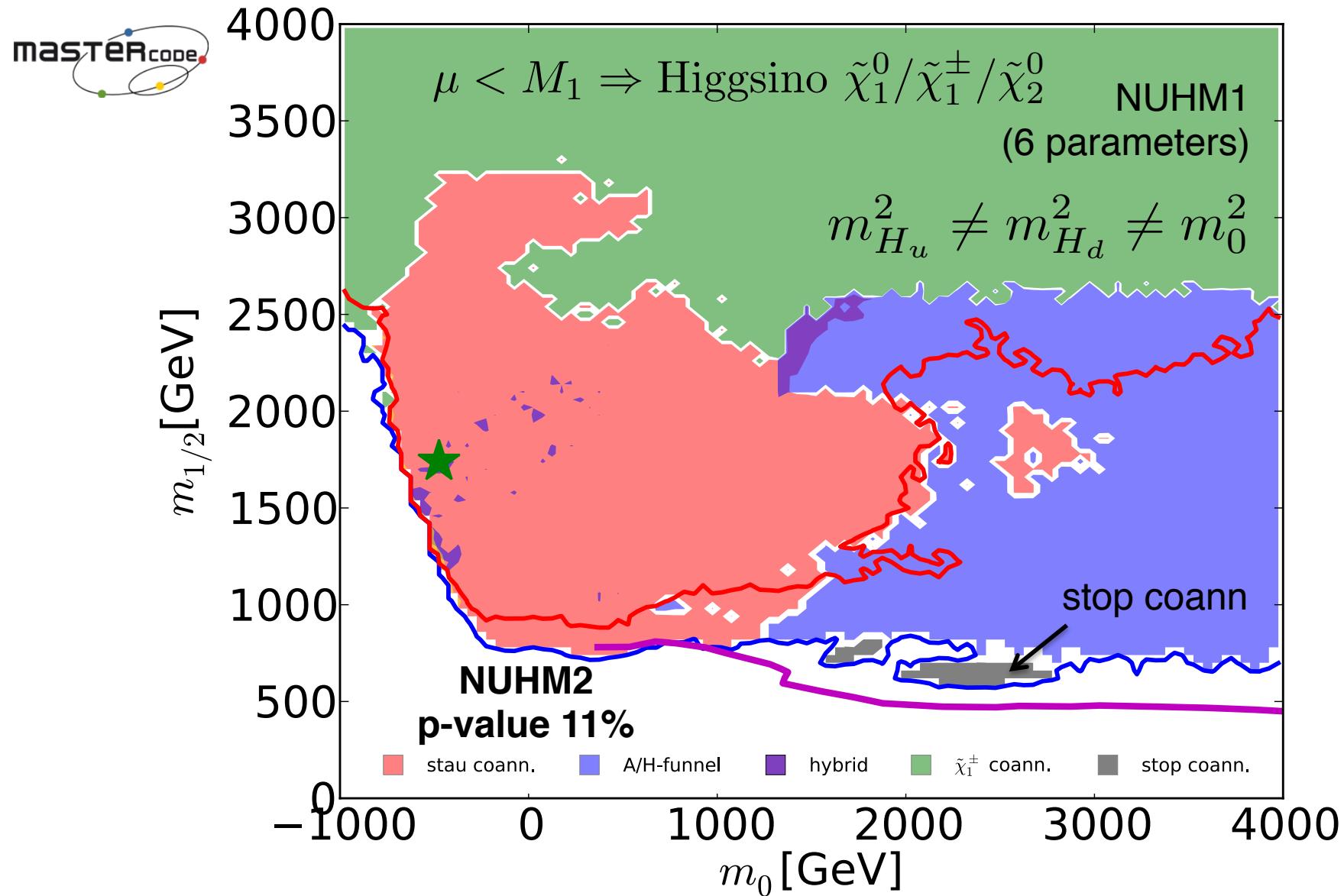
NUHM1

★ — NUHM1: best fit, 1σ , 2σ

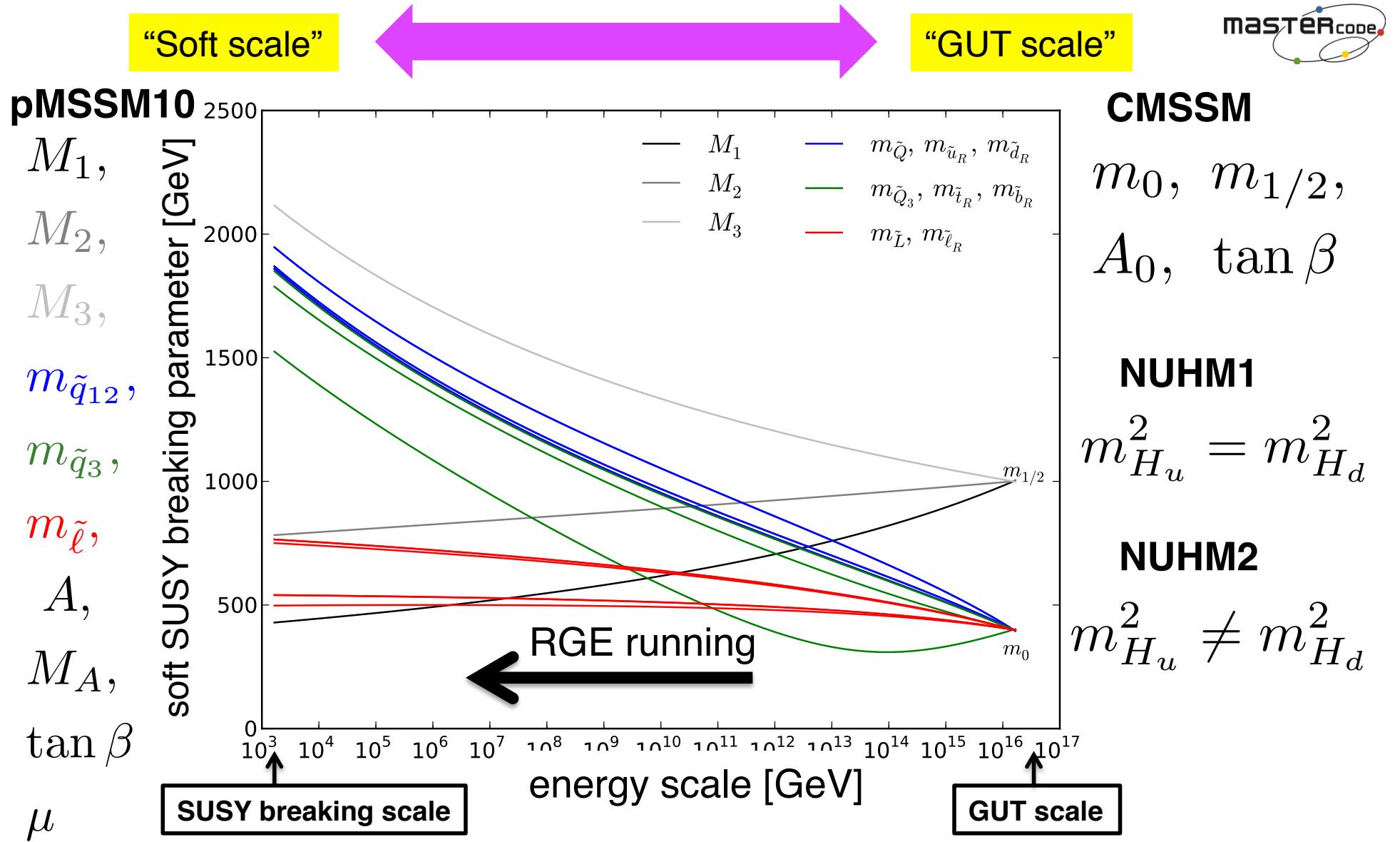


NUHM2

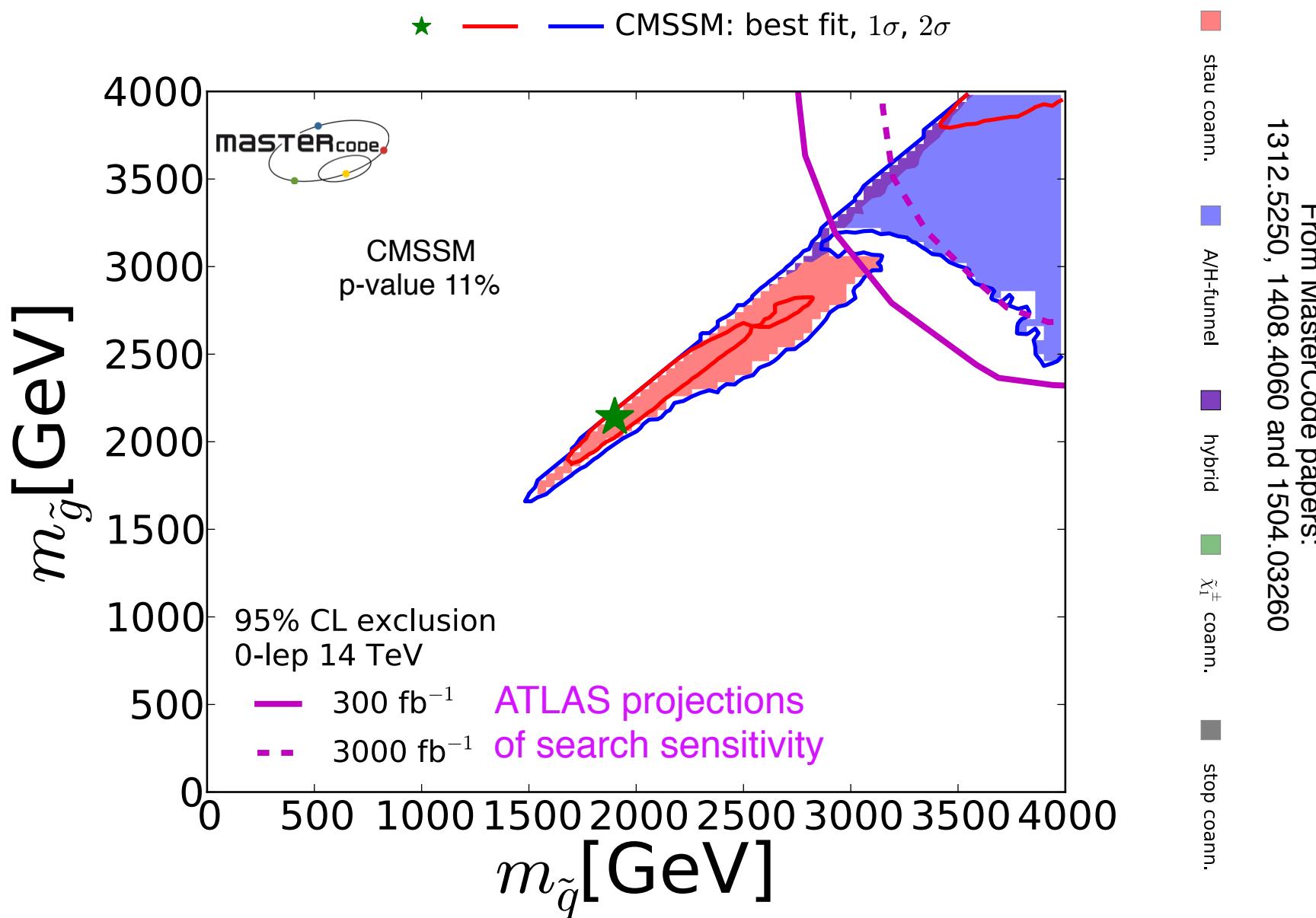
★ — NUHM2: best fit, 1σ , 2σ



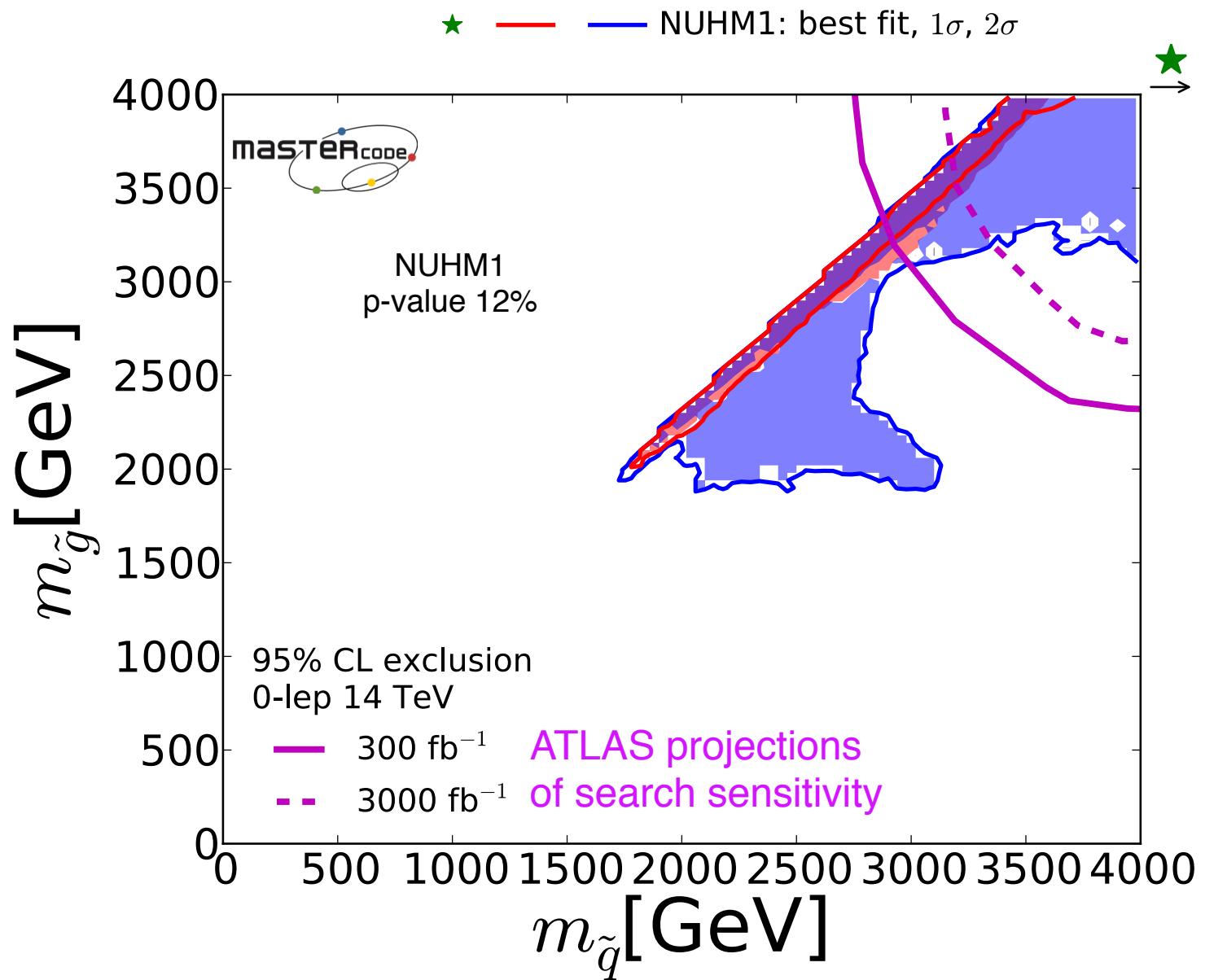
MasterCode: The two worlds of SUSY models



CMSSM Today: Mq-Mg Search plane

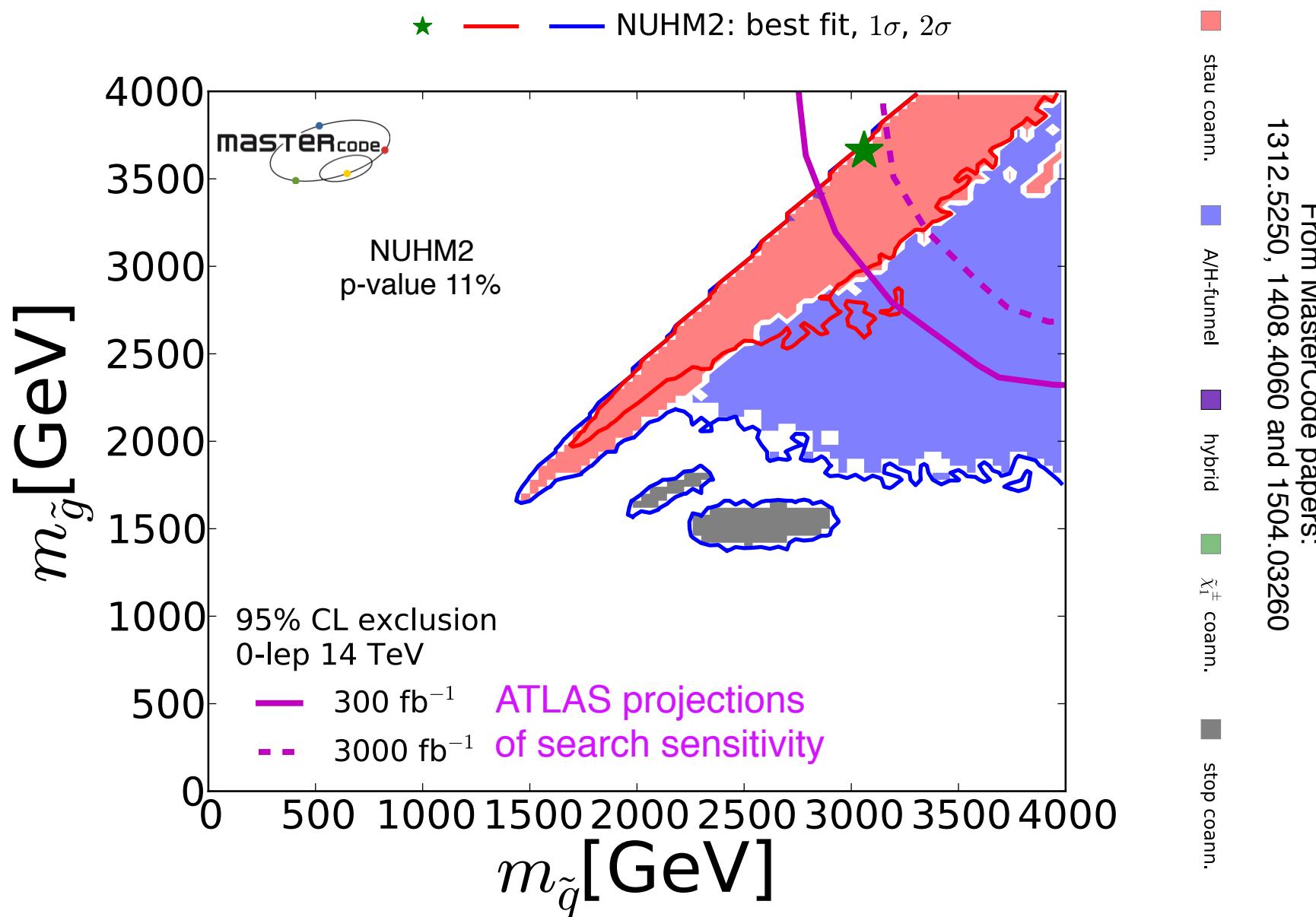


CMSSM Today: Mq-Mg Search plane



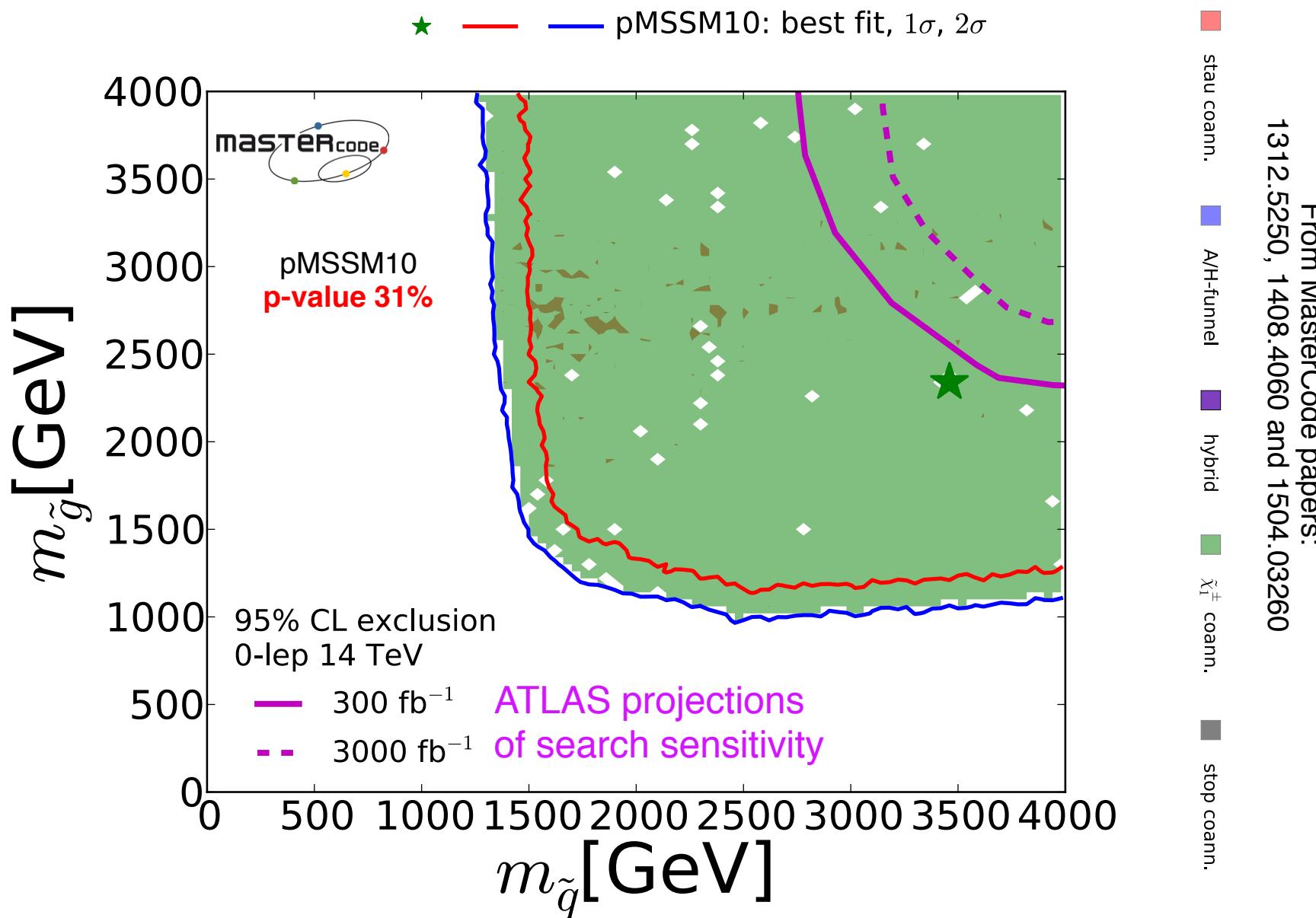
From MasterCode papers:
1312.5250, 1408.4060 and 1504.03260

CMSSM Today: M \tilde{q} -M $\tilde{\mu}$ Search plane

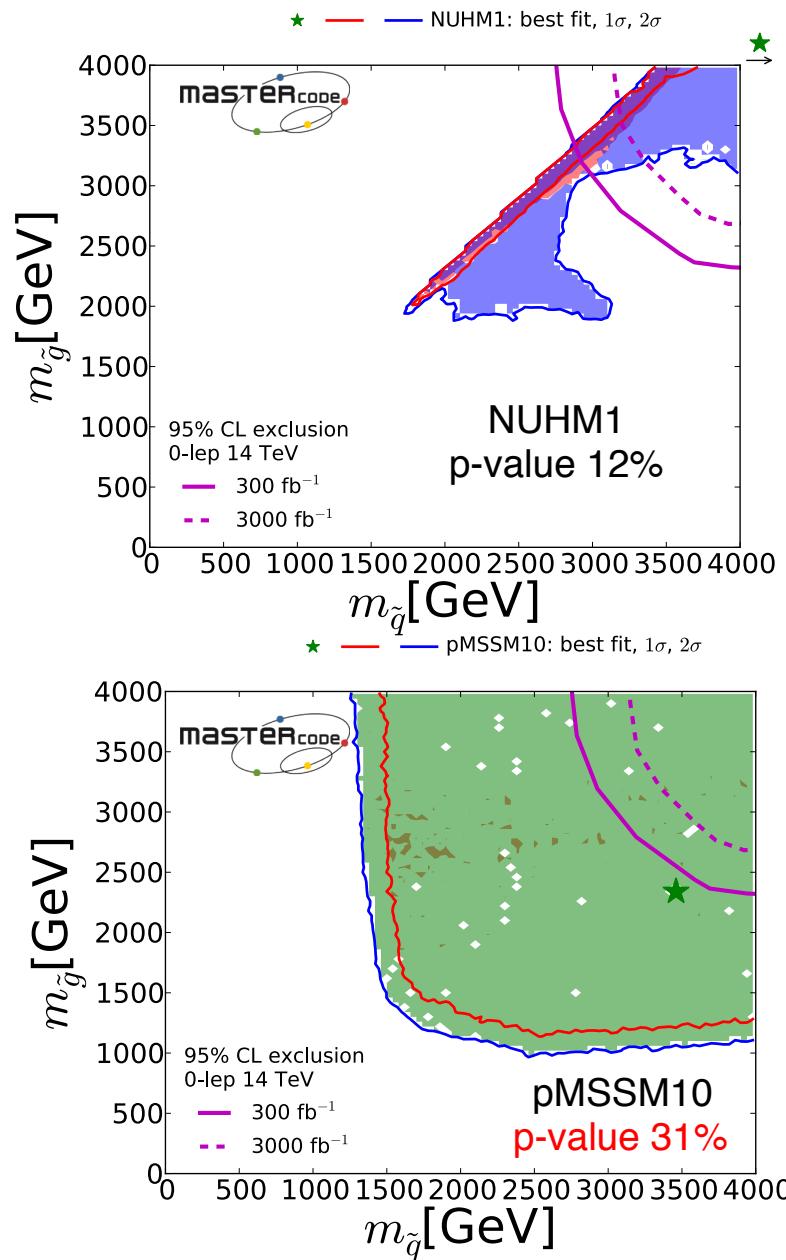
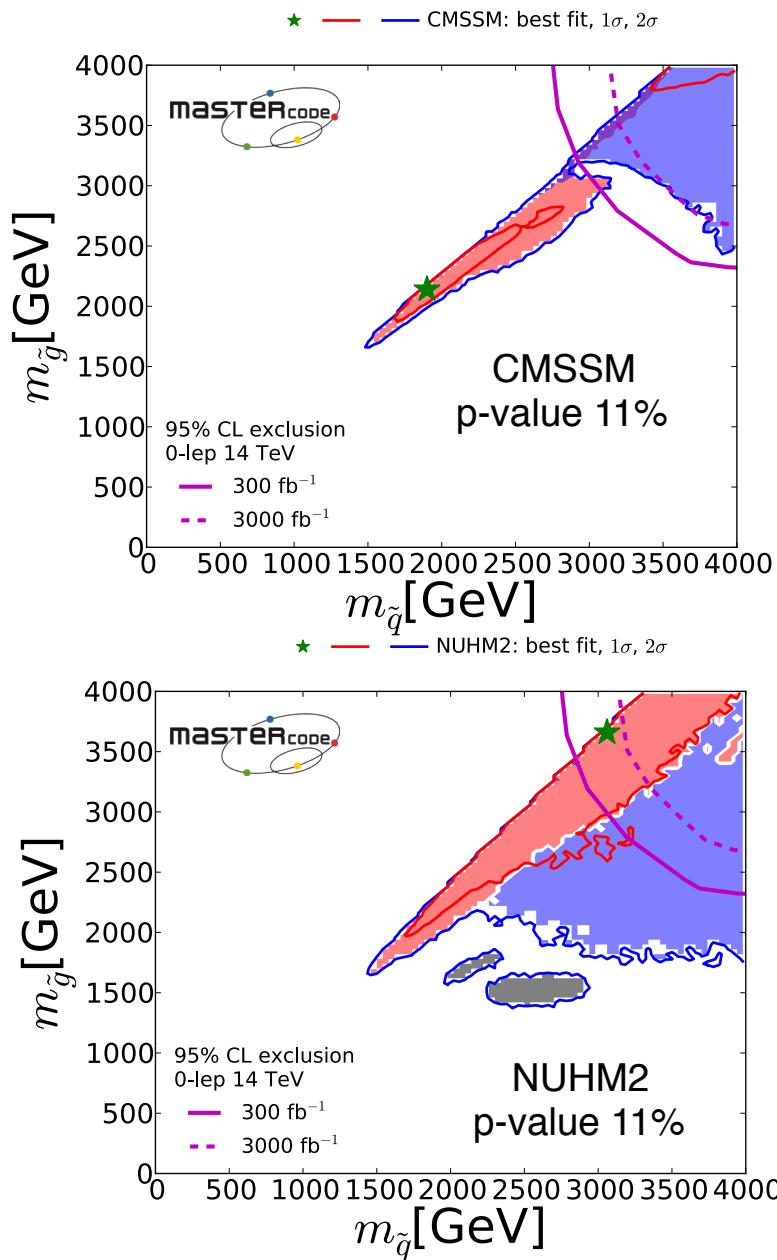


From MasterCode papers:
1312.5250, 1408.4060 and 1504.03260

CMSSM Today: Mq-Mg Search plane

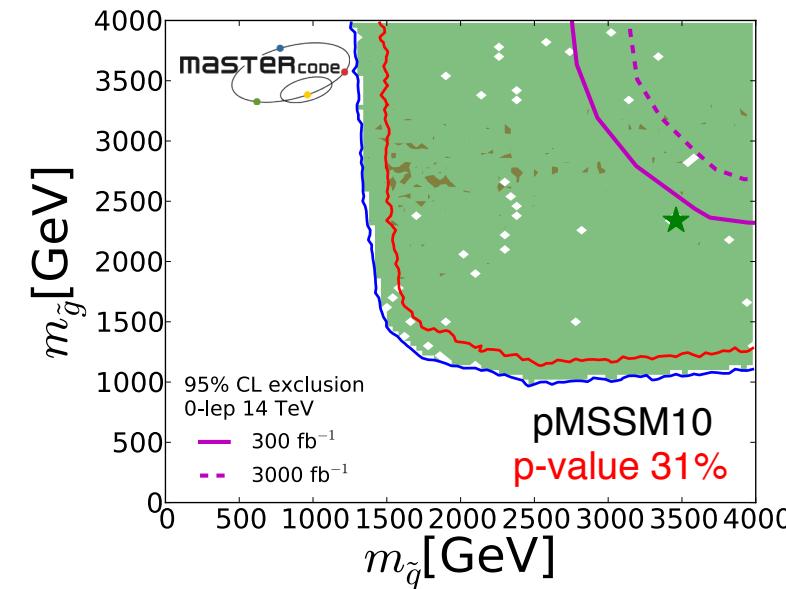
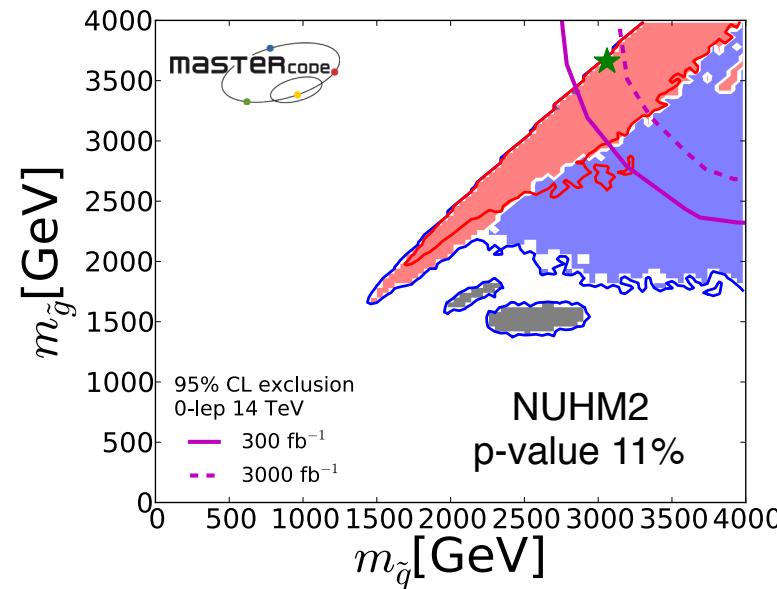


Models in Comparison in “Mq-Mg Search plane”

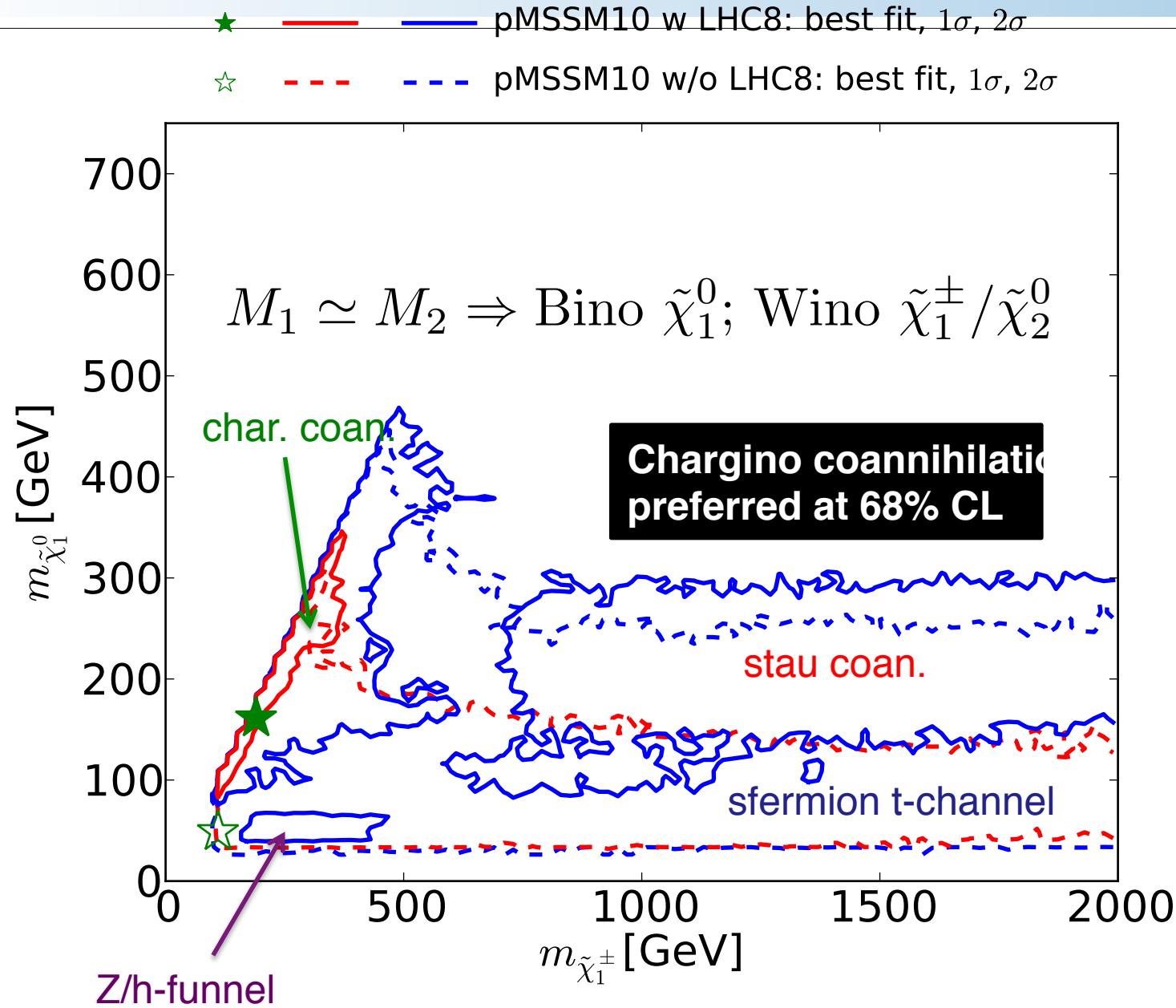


stau coann.
A/H-funnel
hybrid
 $\tilde{\chi}_1^\pm$ coann.
stop coann.

From MasterCode papers:
1312.5250, 1408.4060 and 1504.03260

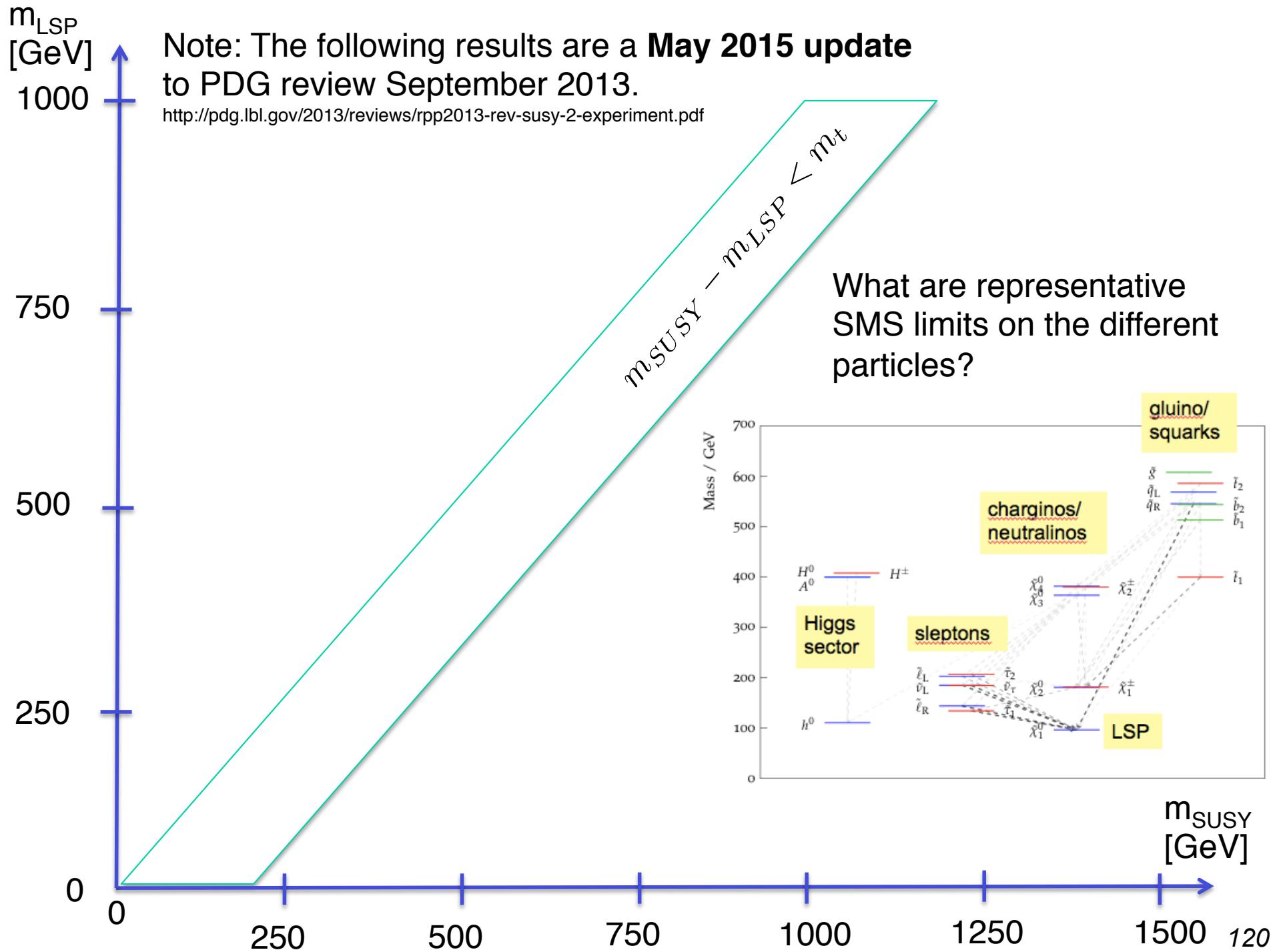


pMSSM10: parameter space

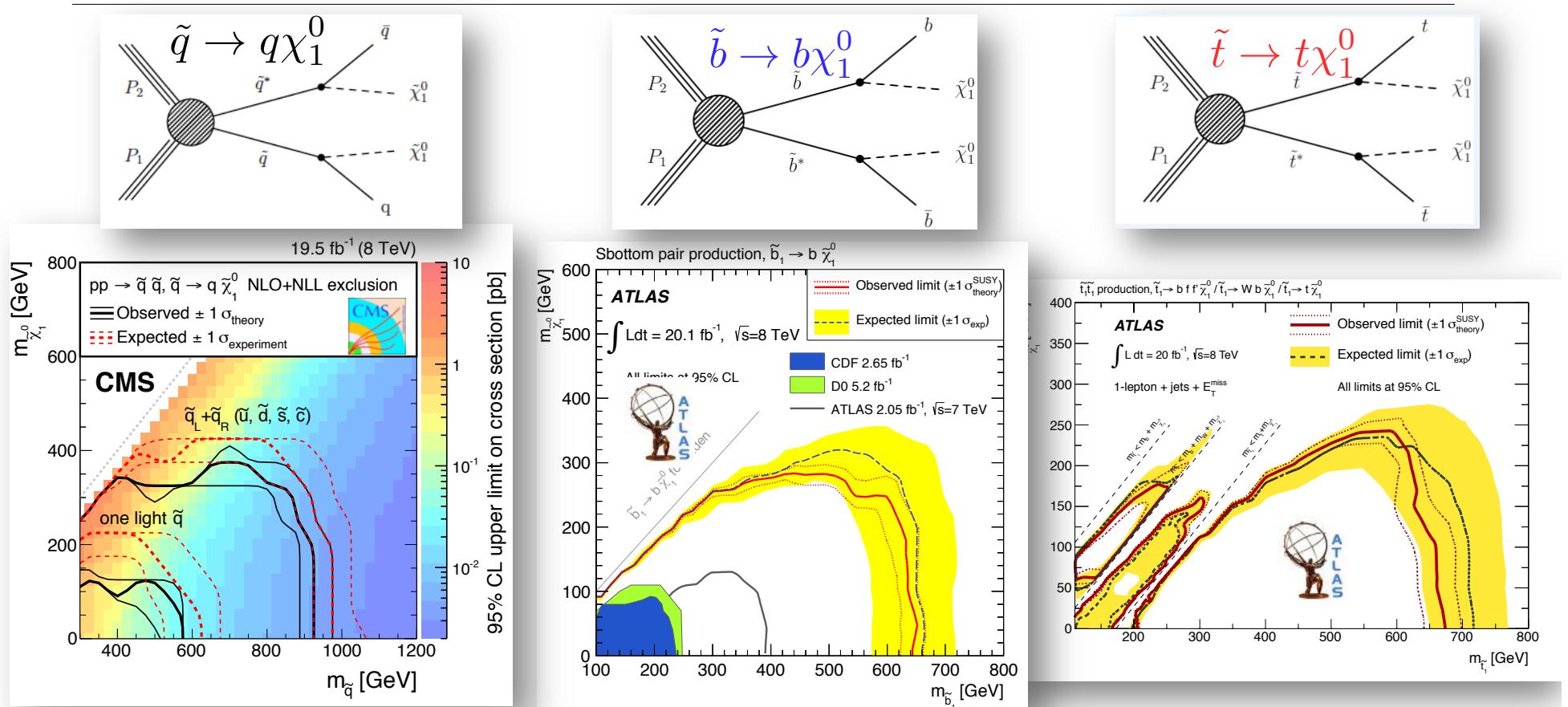


The full story

SUSY SUMMARY PLOT



Direct squark production – chosen limits



CMS arXiv:1502.04358

Signature: Jets + E_T^{miss} with M_{T2}
Limit assumes all 1st & 2nd gen
squarks to be mass degenerate
[or only one light squark]!

ATLAS arXiv:1308.2631

Signature: 2 b-jets + E_T^{miss}

ATLAS arXiv:1407.0583

Signature: 1 Lepton + jets +
 E_T^{miss}

m_{LSP}
[GeV]

Direct squark

$m_{\text{SUSY}} = m_{\tilde{q}}$

$\tilde{q} \rightarrow q\chi_1^0$

CMS arXiv:1502.04358

$\tilde{u}_L \rightarrow u\chi_1^0$

CMS arXiv:1502.04358

$\tilde{b} \rightarrow b\chi_1^0$

ATLAS arXiv:1308.2631

$\tilde{t} \rightarrow t\chi_1^0$

ATLAS arXiv:1407.0583

1000

750

500

250

0

$m_{\text{SUSY}} - m_{\text{LSP}} < m_t$

BR=100%
all limits are
observed nominal
95% CLs limits
RP conserved

m_{SUSY}
[GeV]

0

250

500

750

1000

1250

1500

122

Direct squark

$m_{\text{SUSY}} = m_{\tilde{q}}$

$\tilde{q} \rightarrow q\chi_1^0$

CMS arXiv:1502.04358

$\tilde{u}_L \rightarrow u\chi_1^0$

CMS arXiv:1502.04358

$\tilde{b} \rightarrow b\chi_1^0$

ATLAS arXiv:1308.2631

$\tilde{t} \rightarrow t\chi_1^0$

ATLAS arXiv:1407.0583

$m_{\text{SUSY}} - m_{\text{LSP}} < m_t$

BR=100%
all limits are
observed nominal
95% CLs limits
RP conserved

m_{SUSY}
[GeV]

0

250

500

750

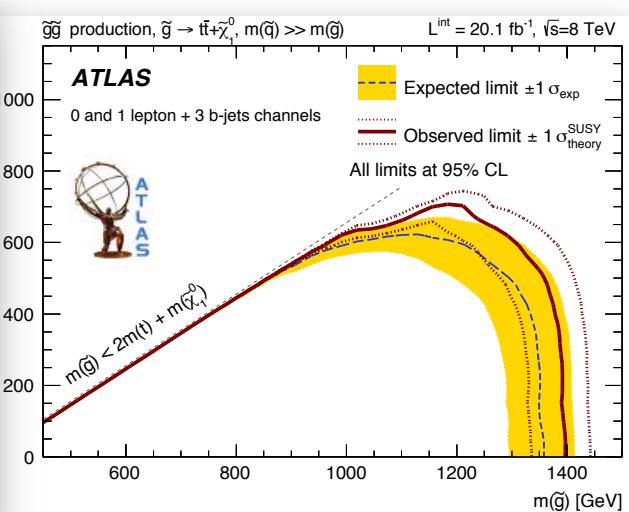
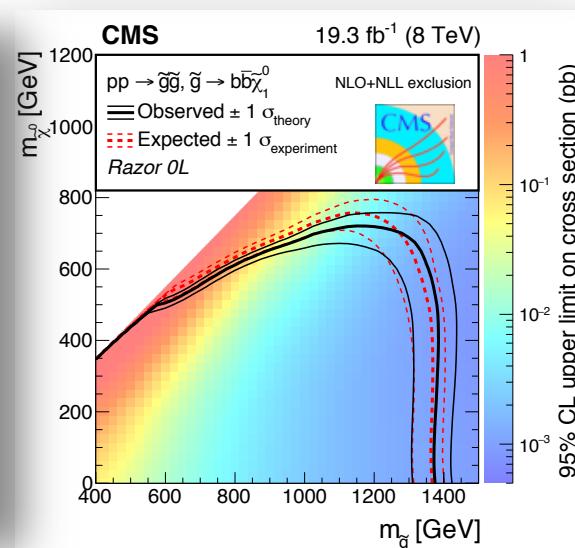
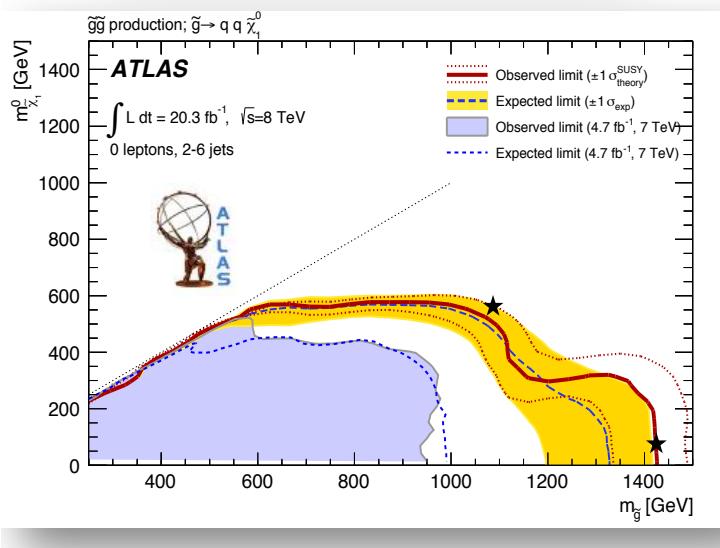
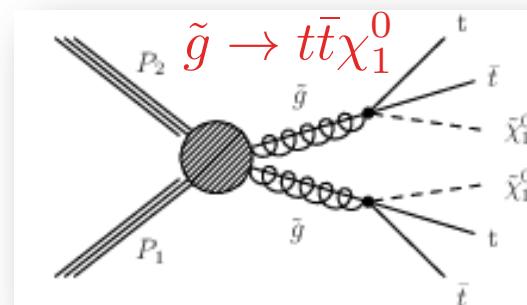
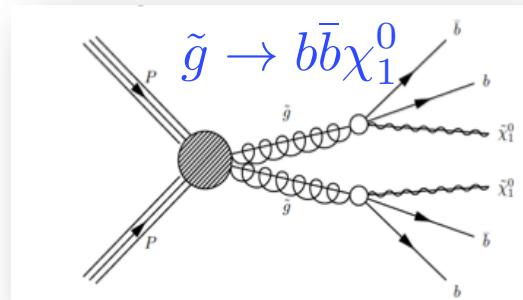
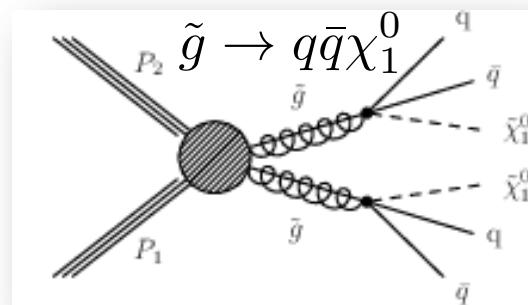
1000

1250

1500

122

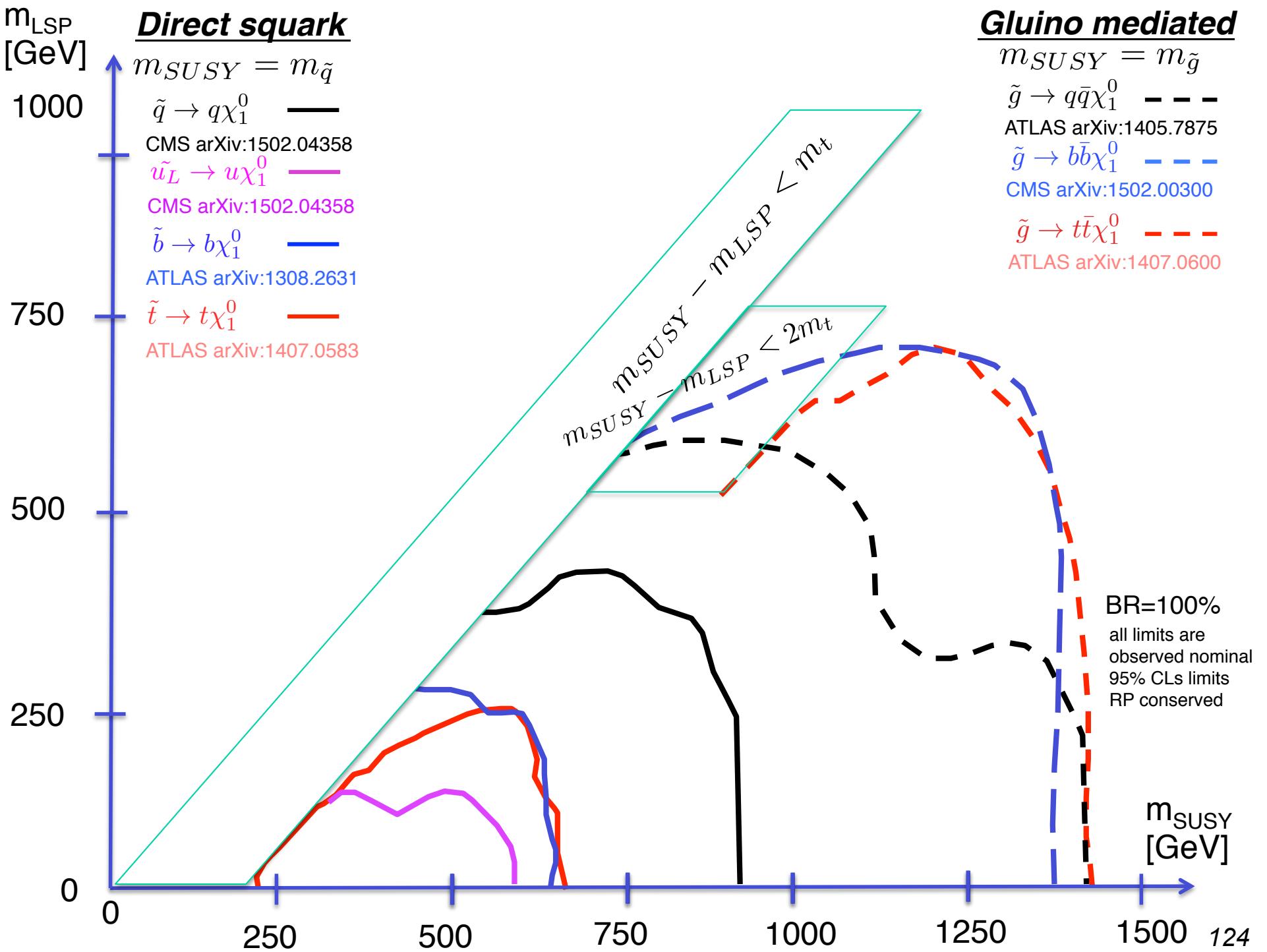
Gluino mediated squark production – limits chosen

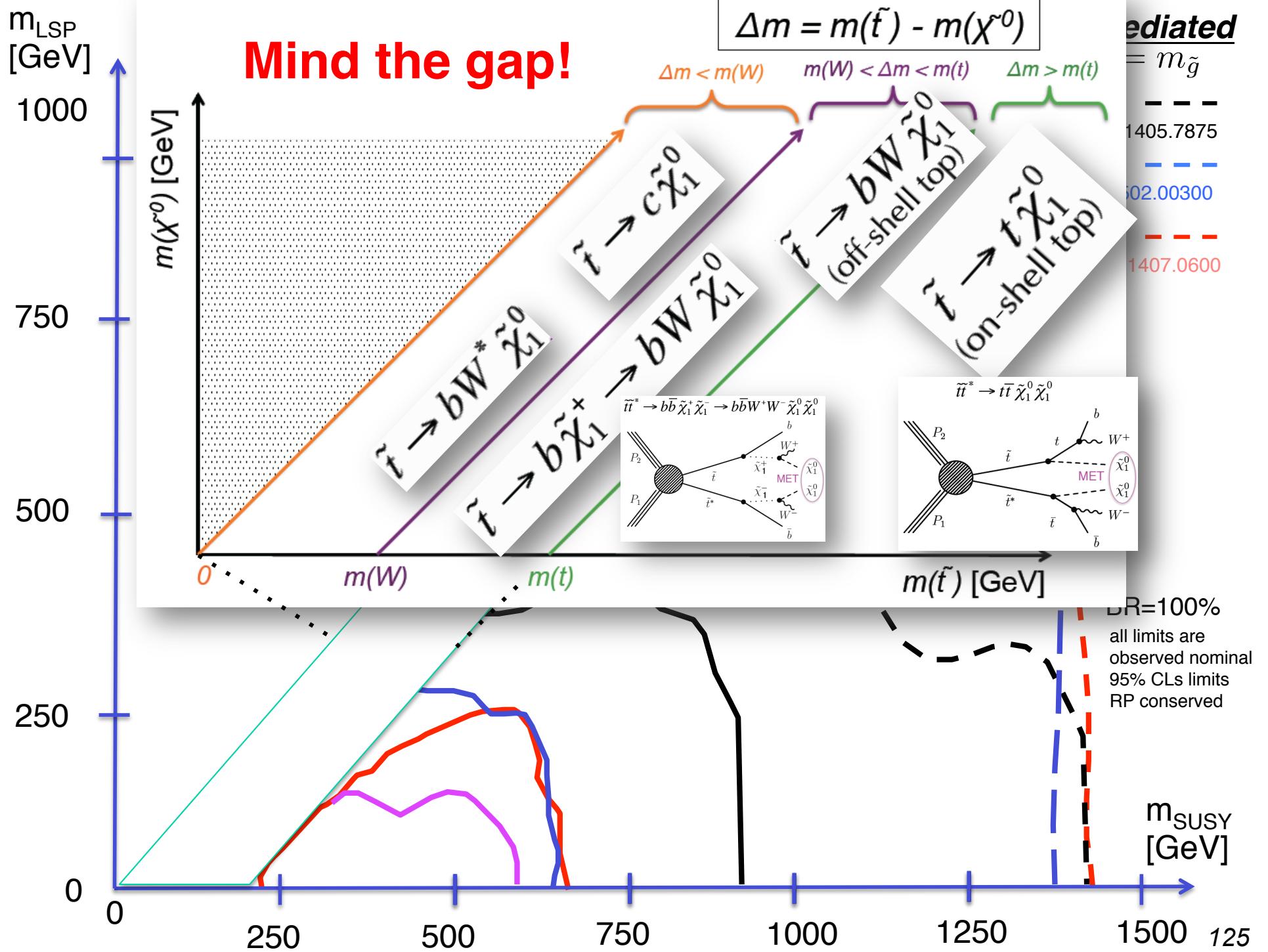


ATLAS arXiv:1405.7875
Signature: 0L + 2-6 Jets
+ E_t^{miss}

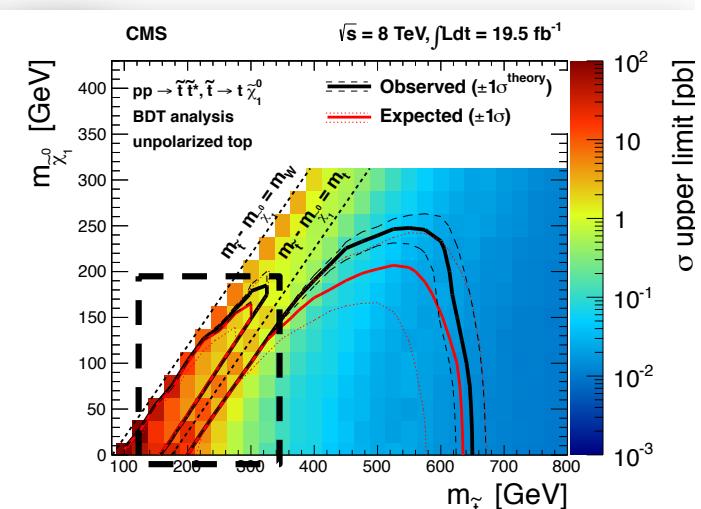
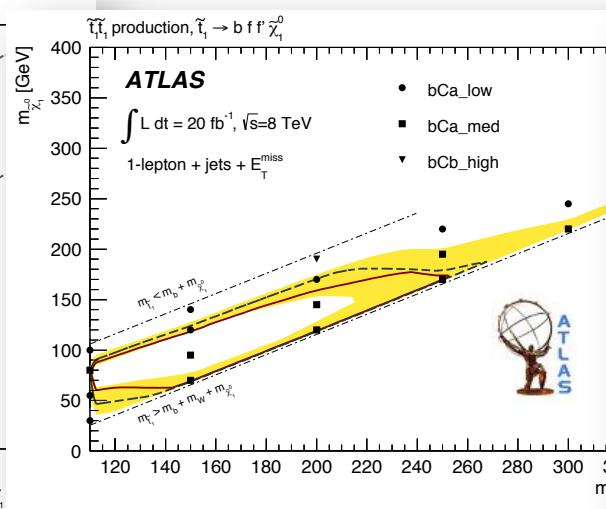
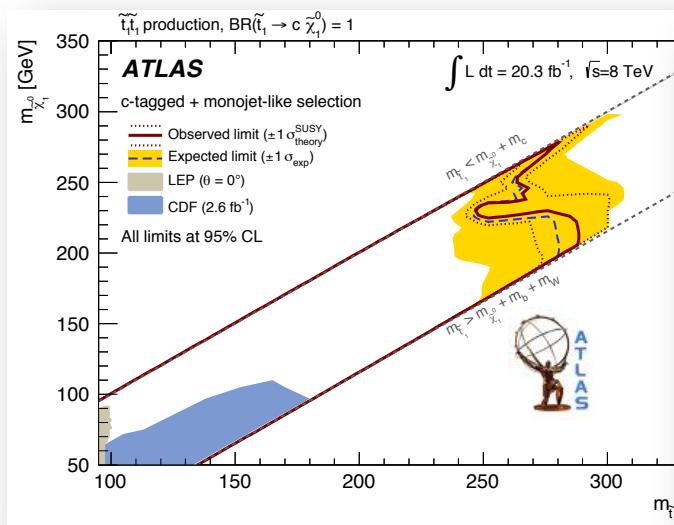
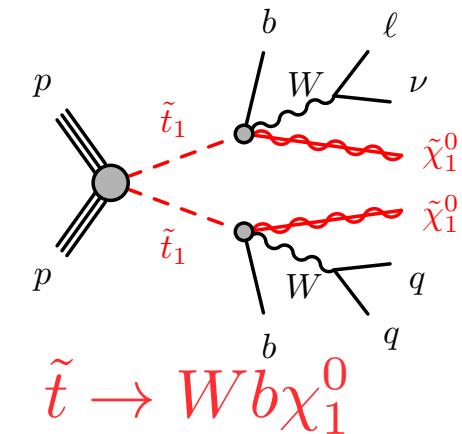
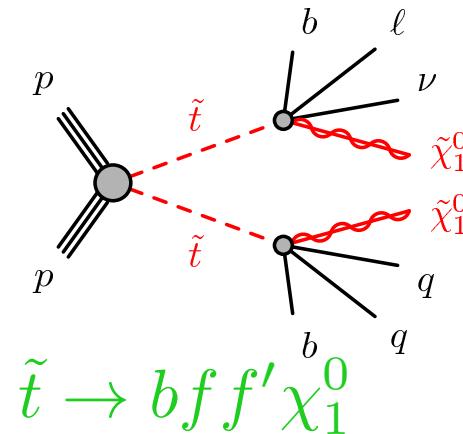
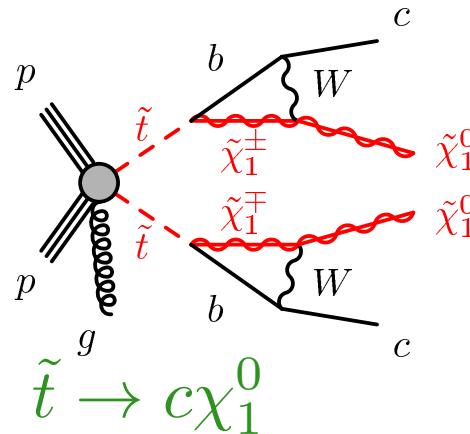
CMS arXiv:1502.00300
Signature: 0L + Razor
+ b-tag

Signature: 0/1 Leptons +
3 b-tag + E_t^{miss}





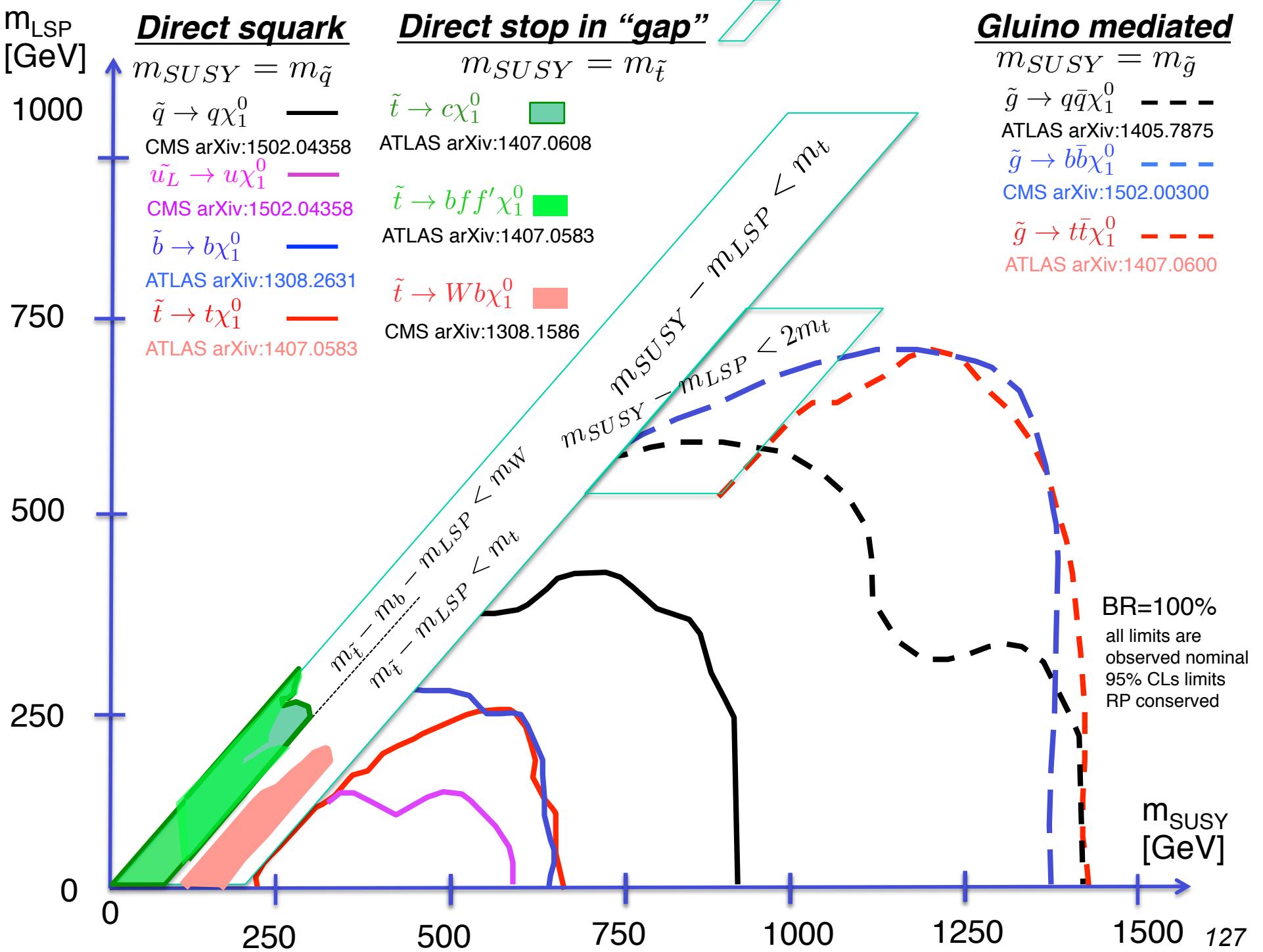
Compressed stop – mind the gap!

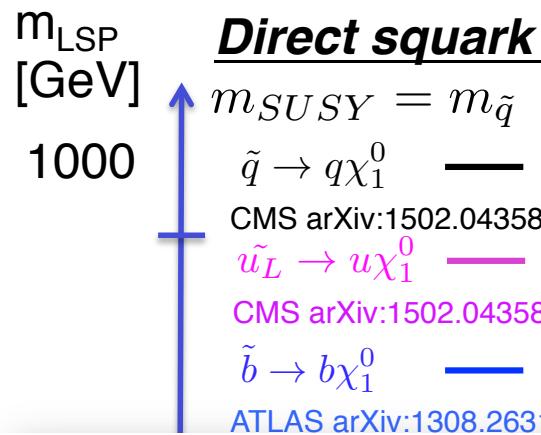


ATLAS arXiv:1407.0608
Mono-jet & c-tag
combined

ATLAS: arXiv:1407.0583
1L + E_T^{mis} & b-tag

CMS arXiv:1308.1586
1L + E_T^{mis} and BDT &
b-tag





Direct stop in “gap”

$m_{\text{SUSY}} = m_{\tilde{t}}$

$\tilde{t} \rightarrow c\chi_1^0$

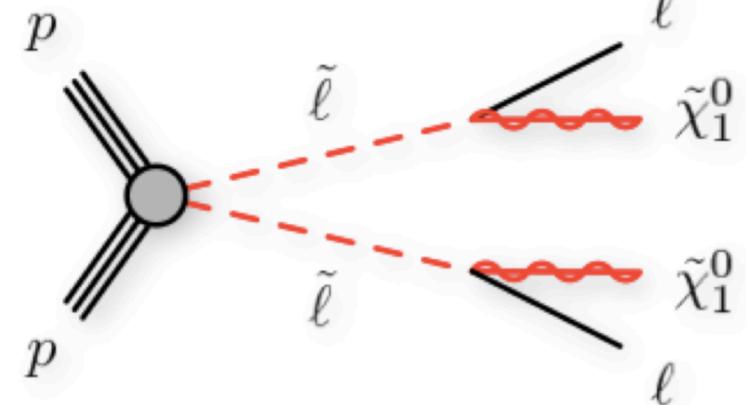
ATLAS arXiv:1407.

$\tilde{t} \rightarrow bff'\chi_1^0$

ATLAS arXiv:1407.

$Wb\chi_1^0$

Xiv:1308.1



Gluino mediated

$m_{\text{SUSY}} = m_{\tilde{g}}$

\tilde{e}_1 —

$\tilde{\chi}_1^0$ - -

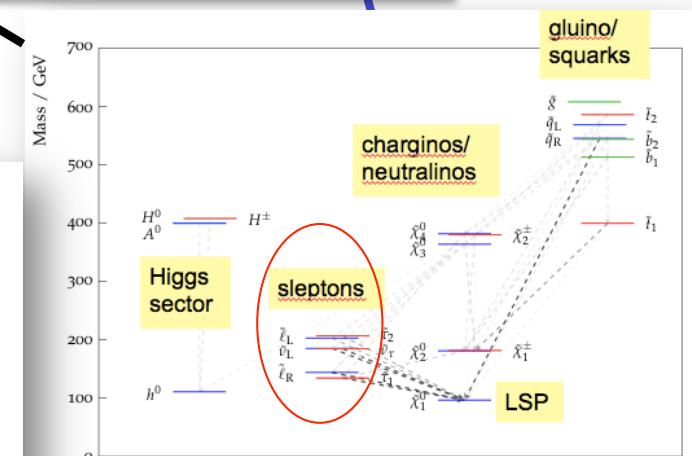
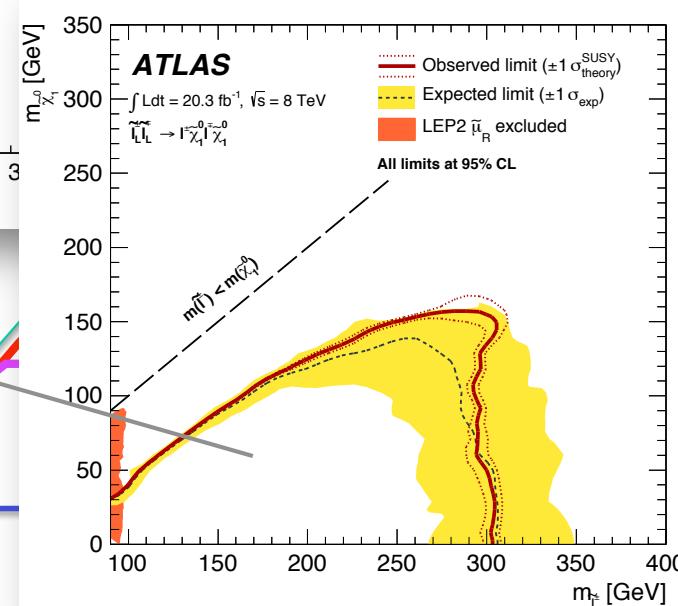
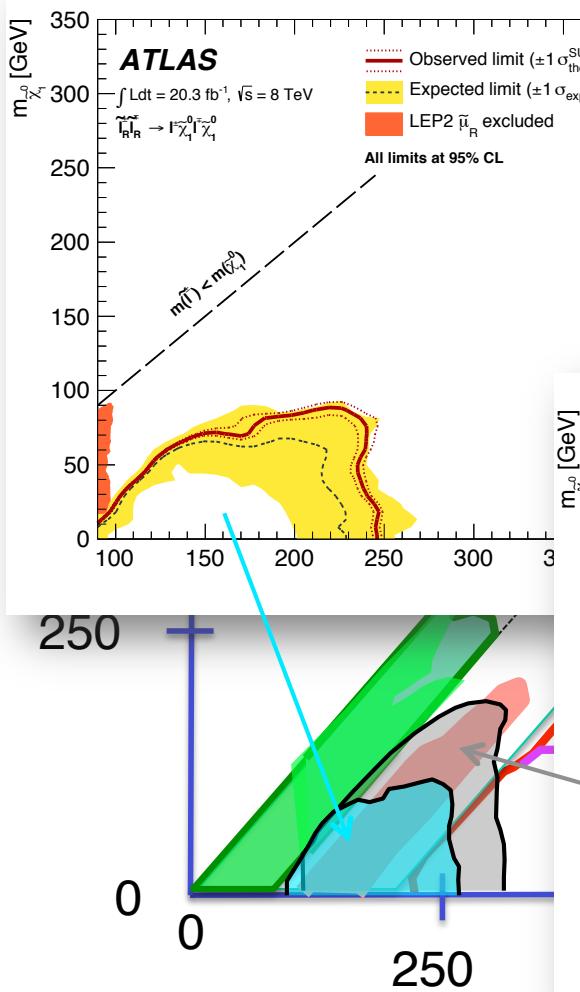
Xiv:1405.7875

\tilde{e}_1 - -

:1502.00300

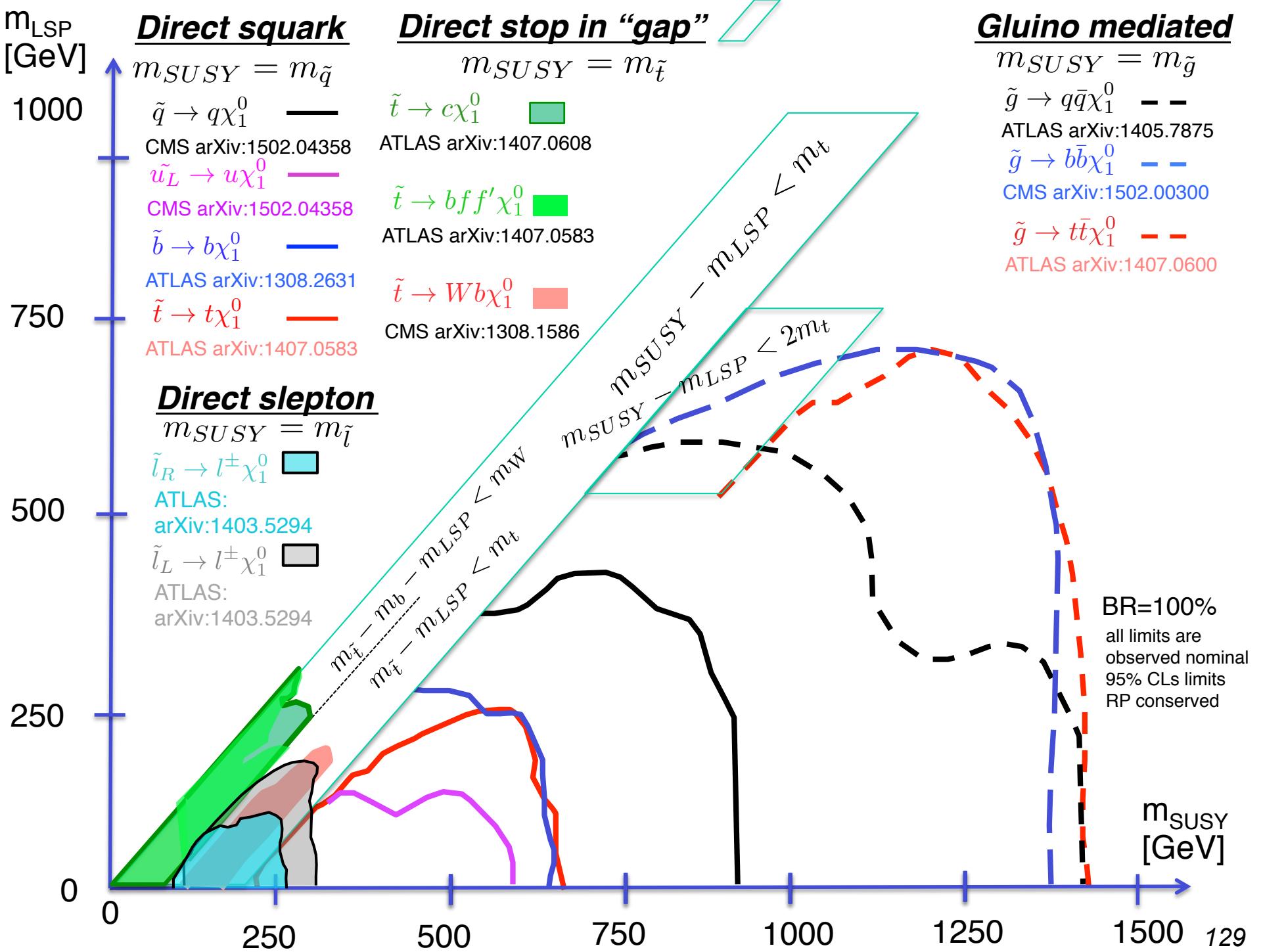
$\tilde{\chi}_1^0$ - -

Xiv:1407.0600



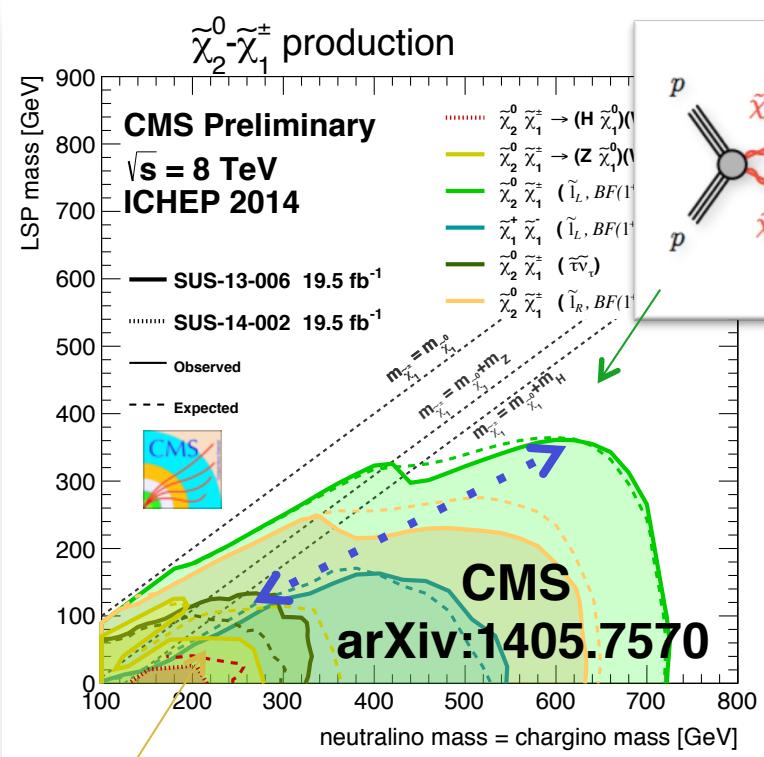
ATLAS arXiv:1403.5294

Signature
2 lepton + E_T^{miss}

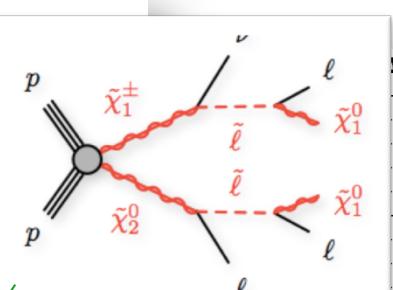


Direct chargino/neutralino production

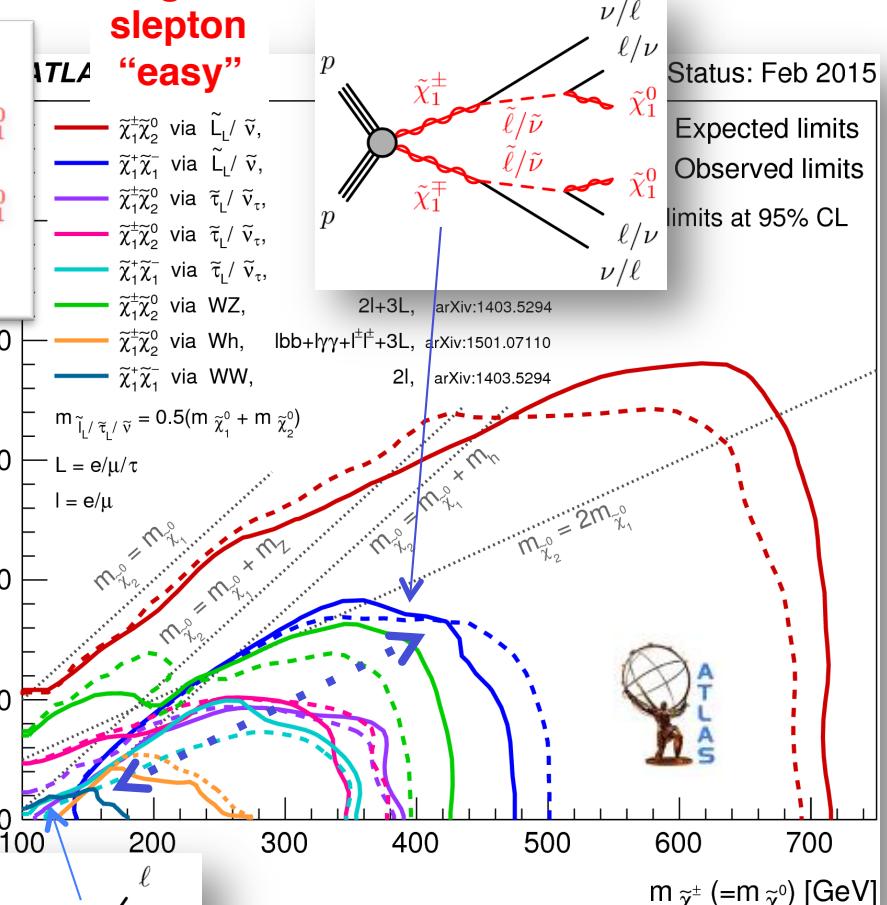
$\tilde{\chi}_2^0 \tilde{\chi}_1^\pm$ production



light slepton
“easy”

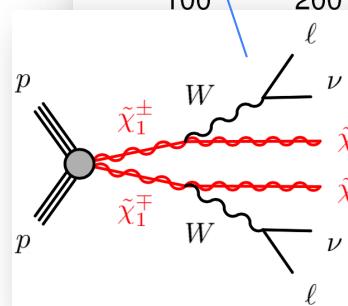
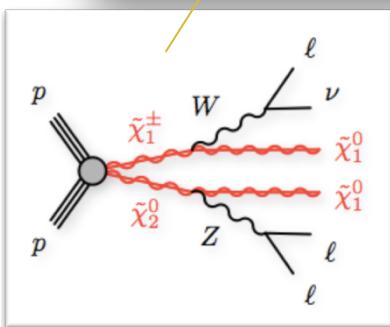


$\tilde{\chi}_1^\pm \tilde{\chi}_1^\pm$ production

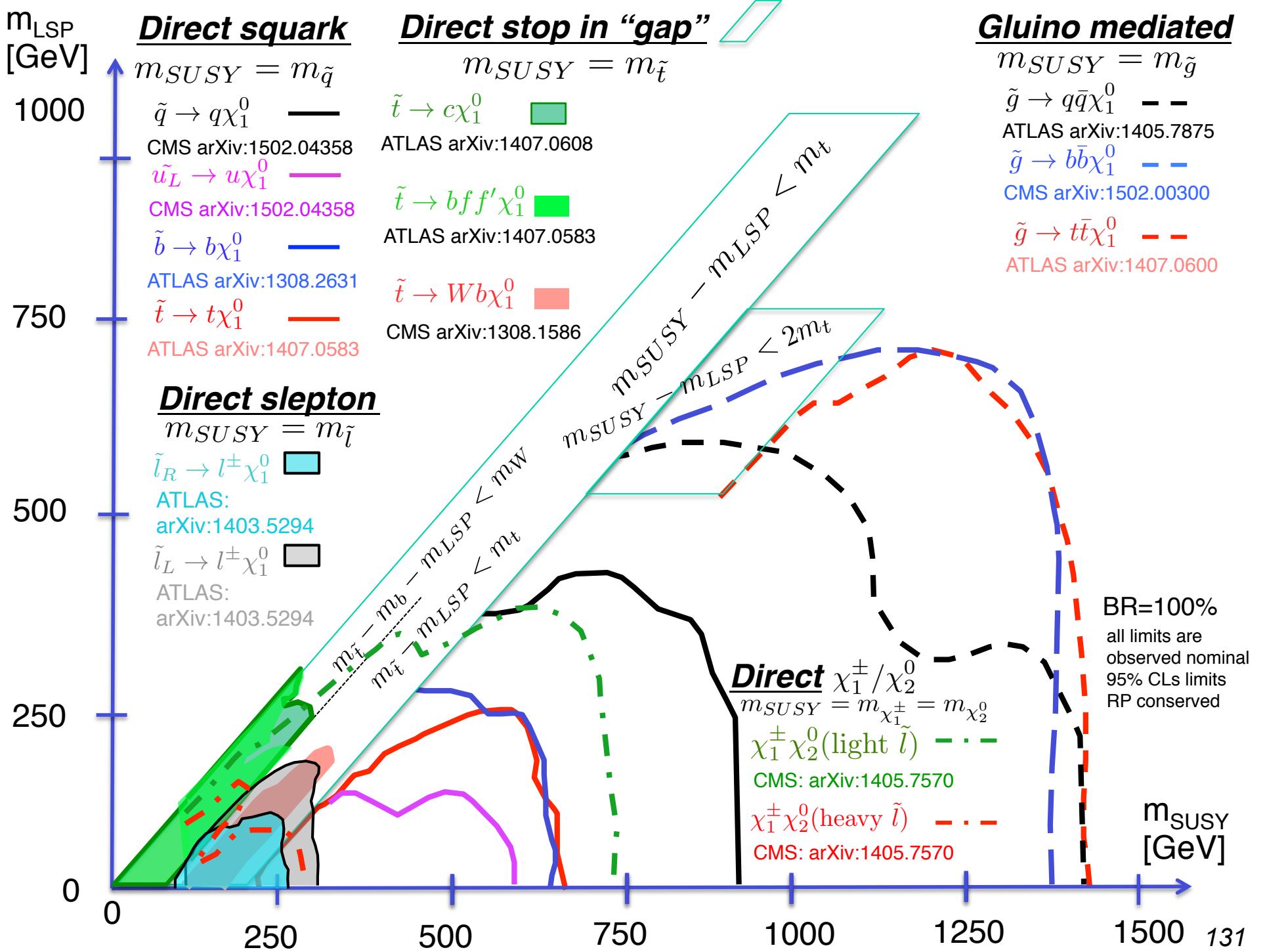


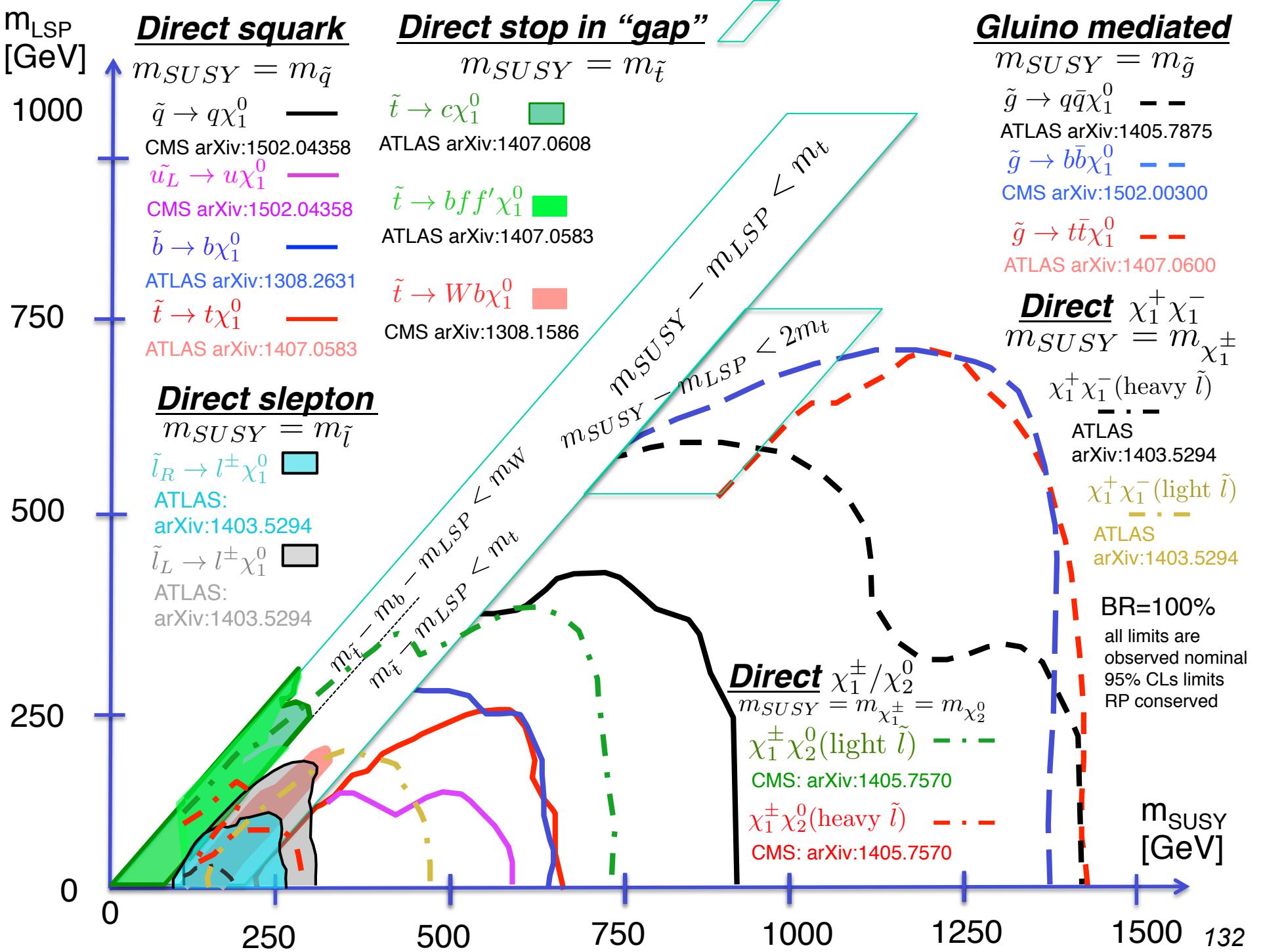
heavy slepton “hard(er)”

Add $Z(l^+l^-)$ +2jets topology in bins of E_t^{miss} to increase sensitivity for “heavy” slepton case

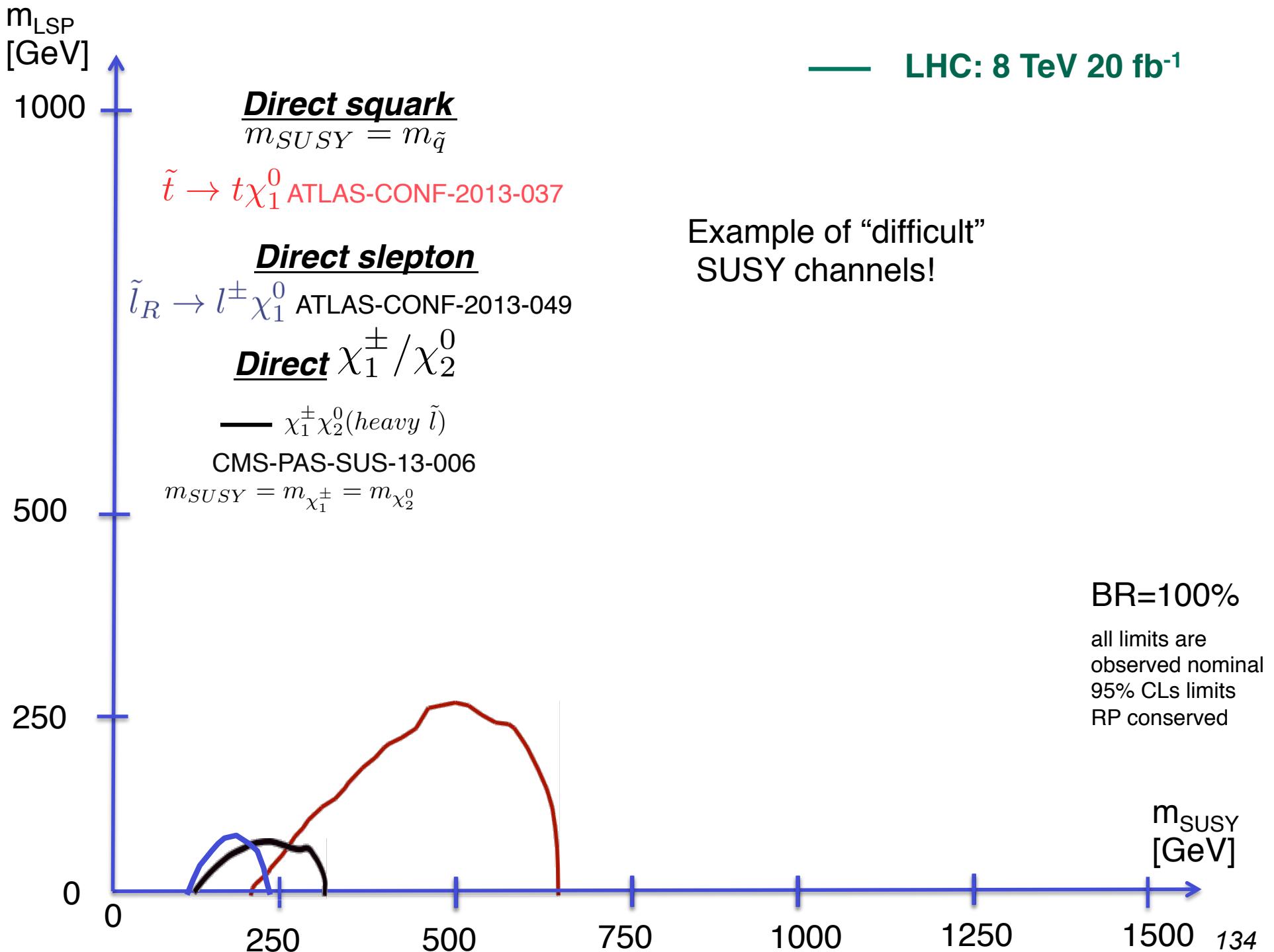


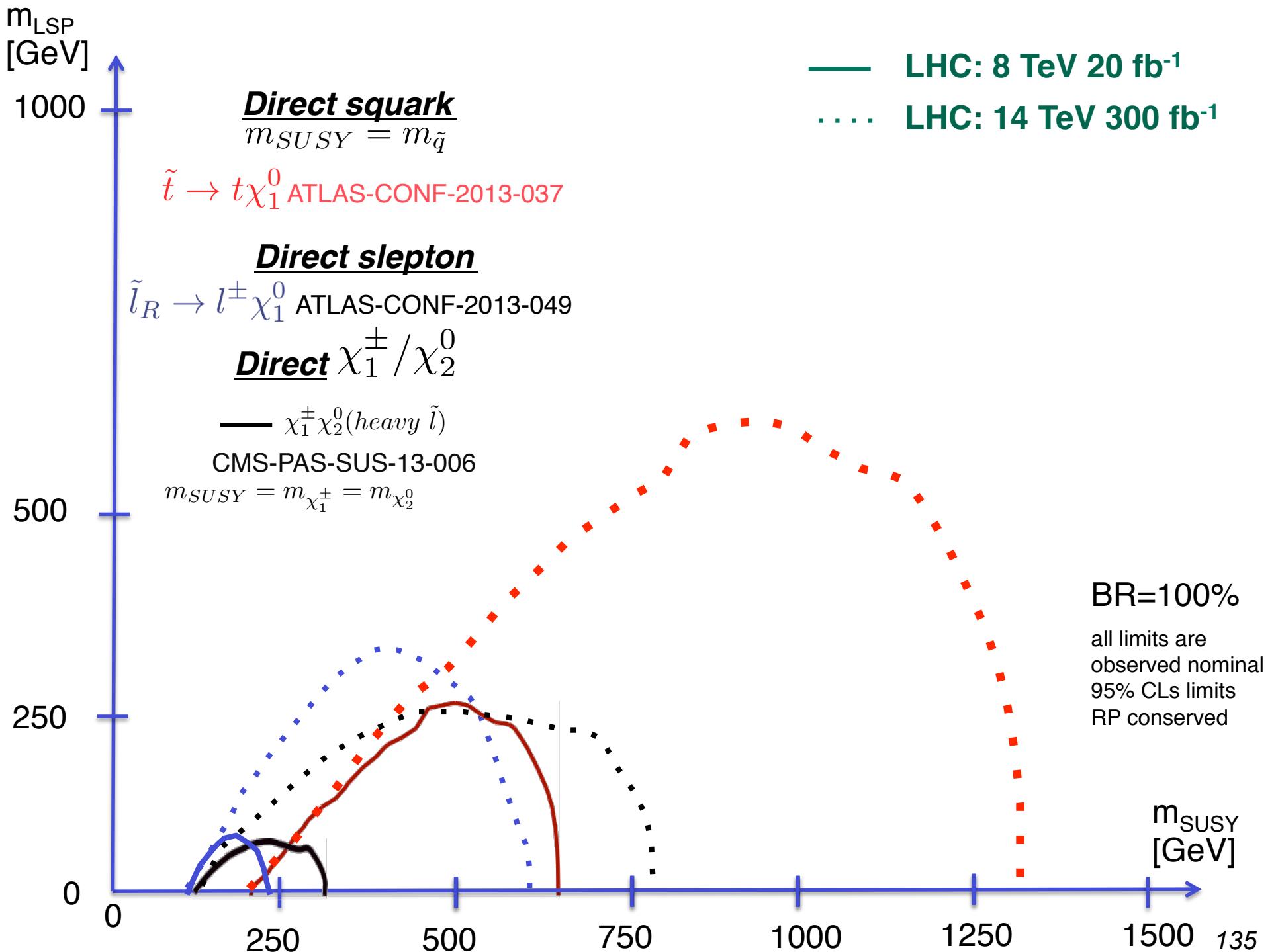
ATLAS arXiv:1403.5294

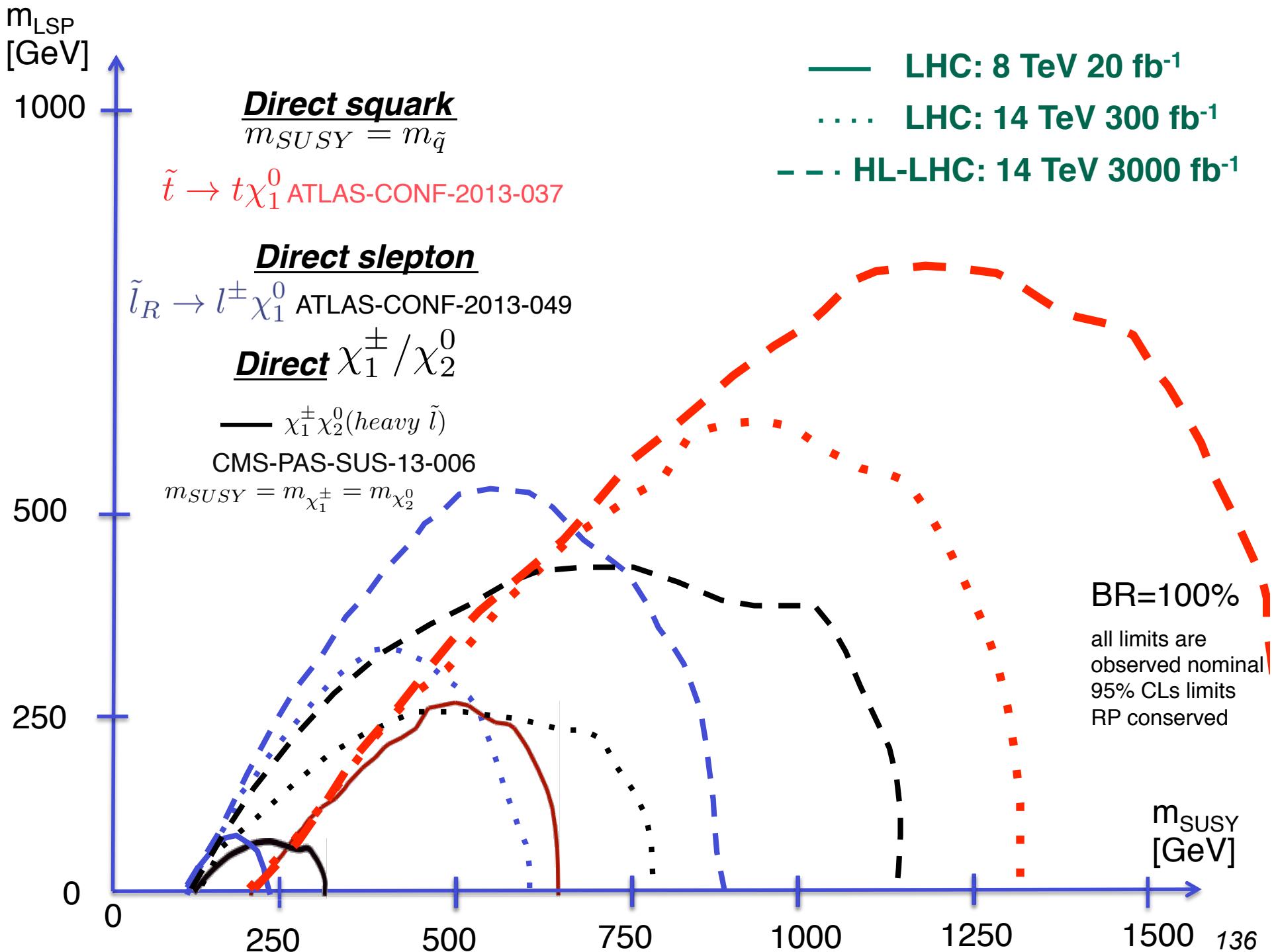


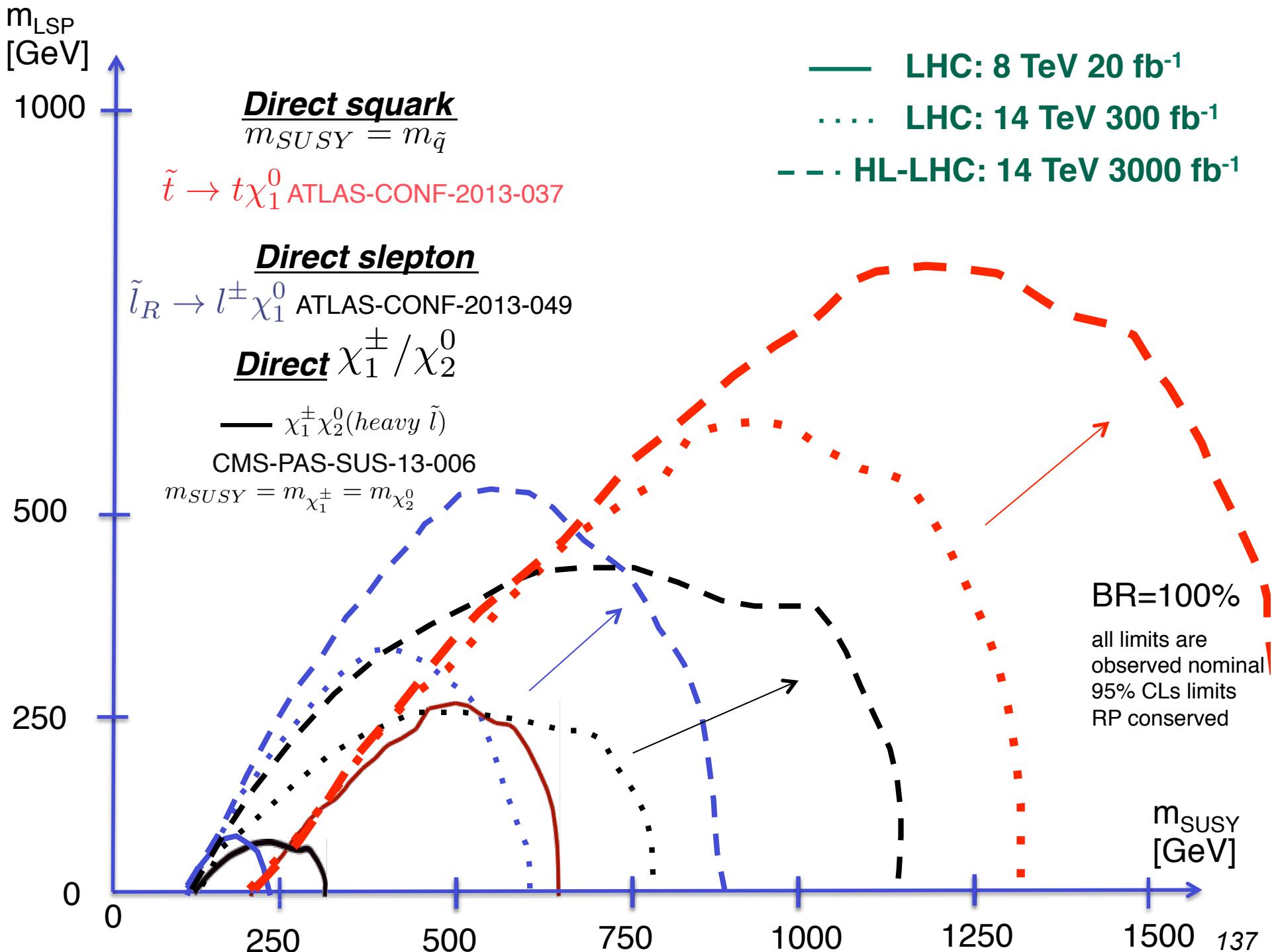


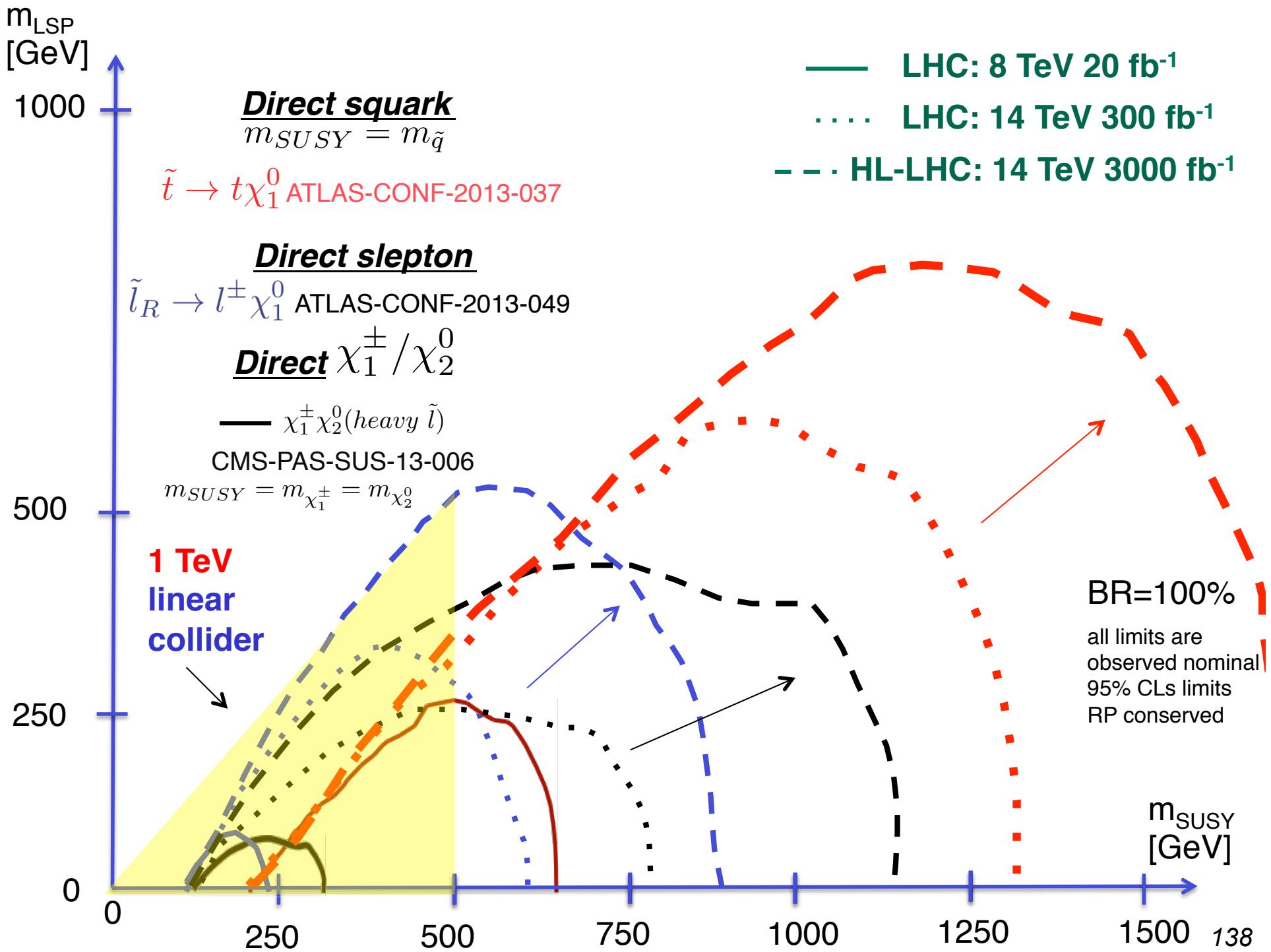
SUSY PROJECTION OF DIFFICULT CHANNELS









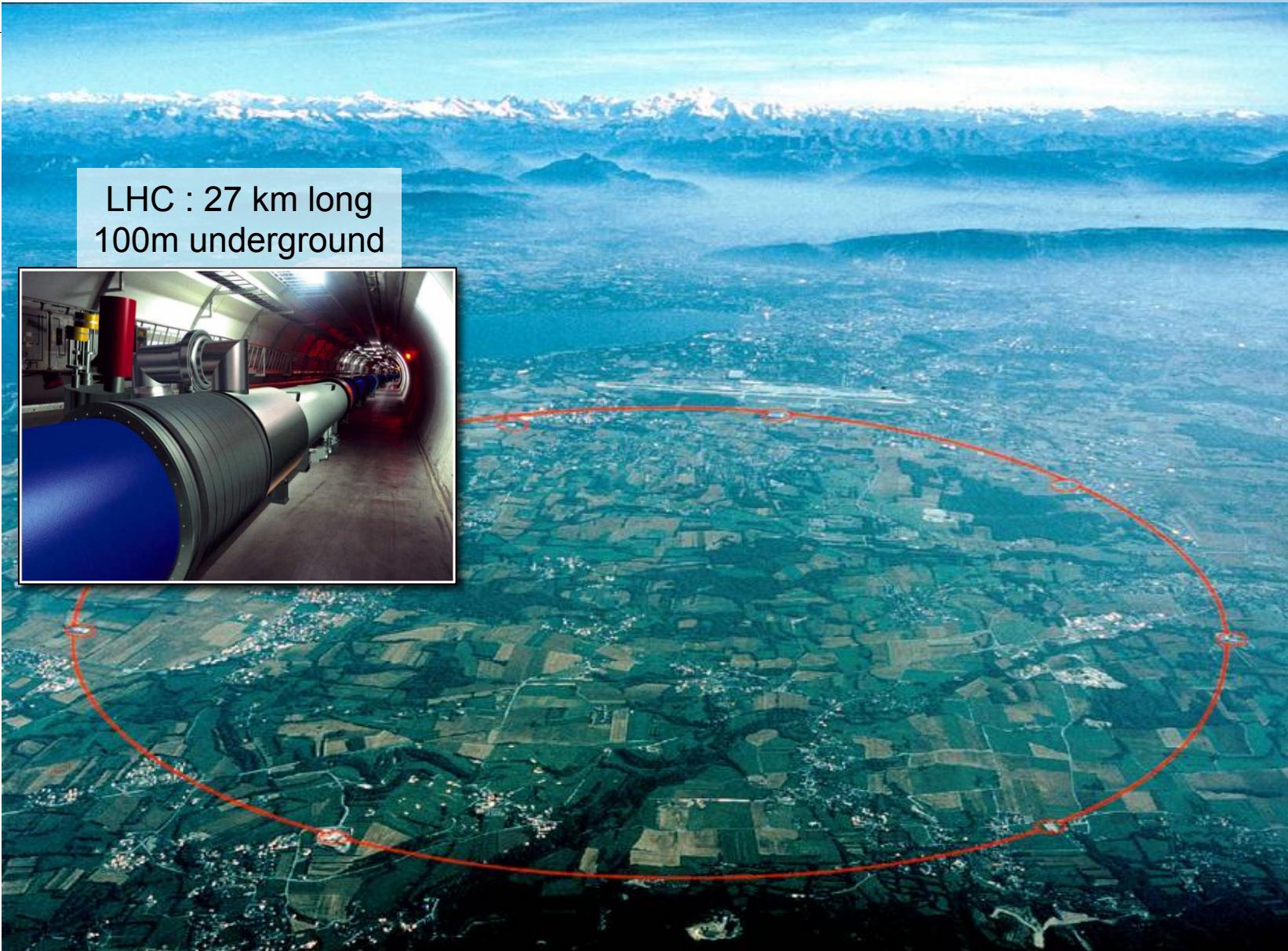


LHC

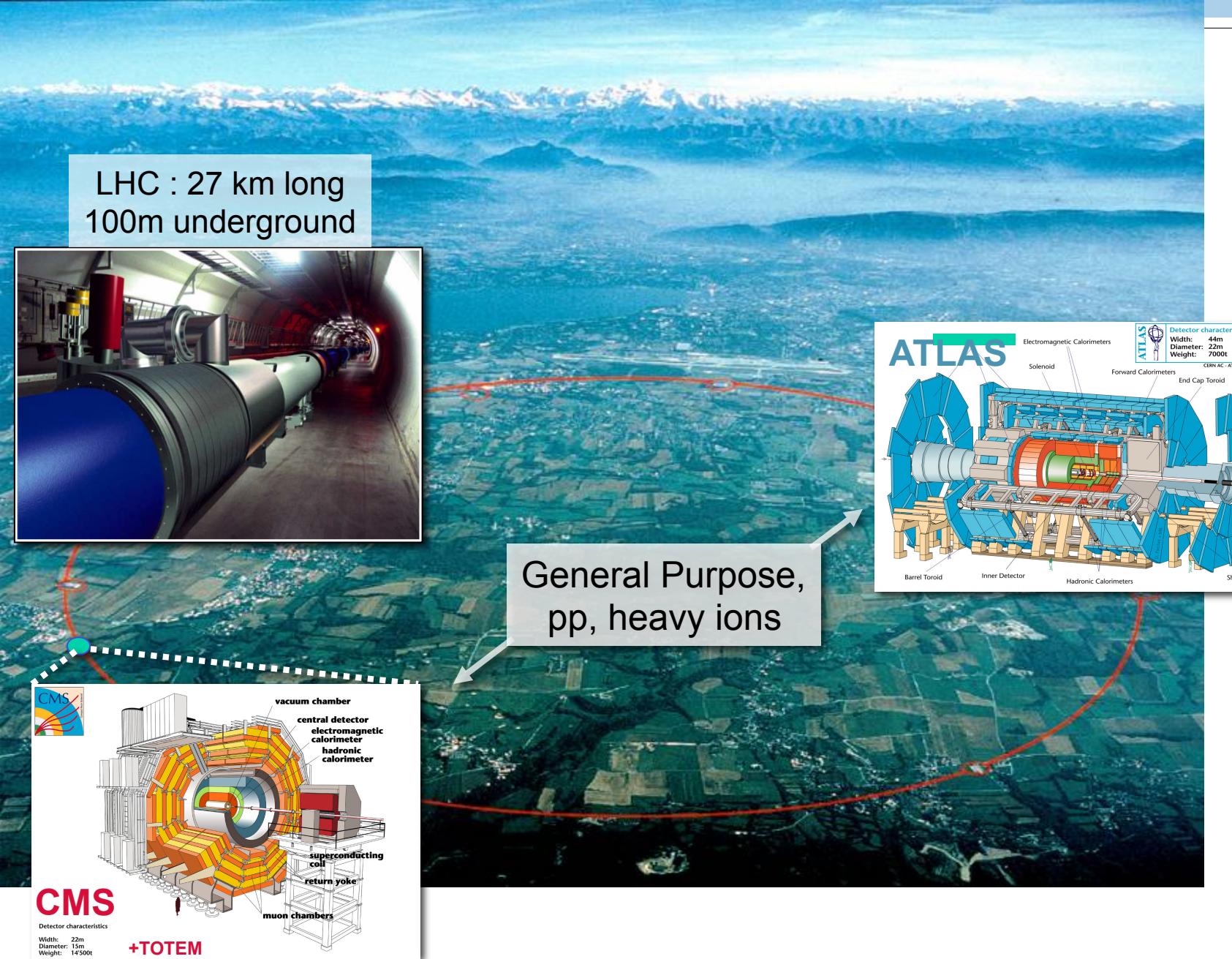
The Large Hadron Collider at CERN



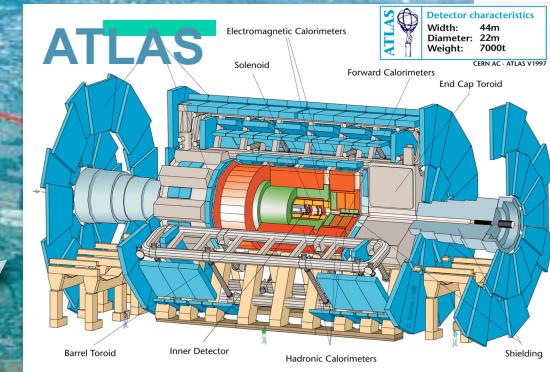
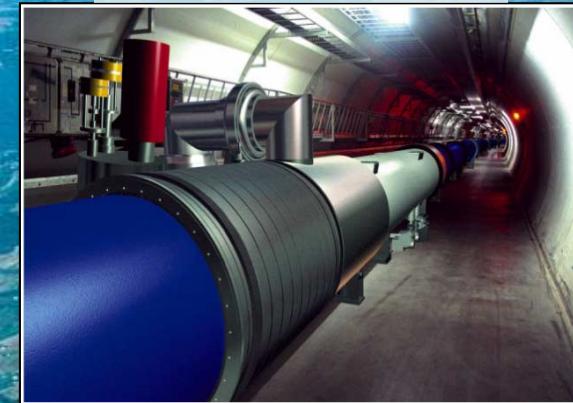
The Large Hadron Collider at CERN



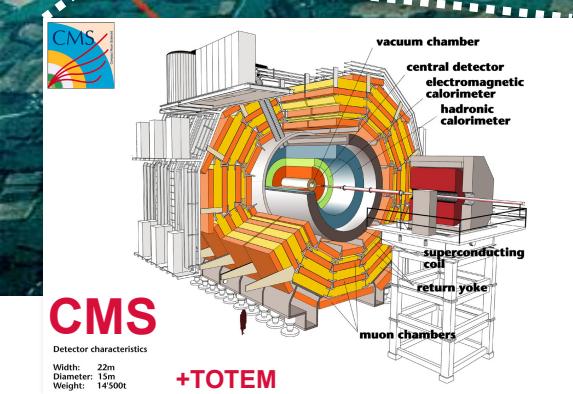
The Large Hadron Collider at CERN



The Large Hadron Collider at CERN



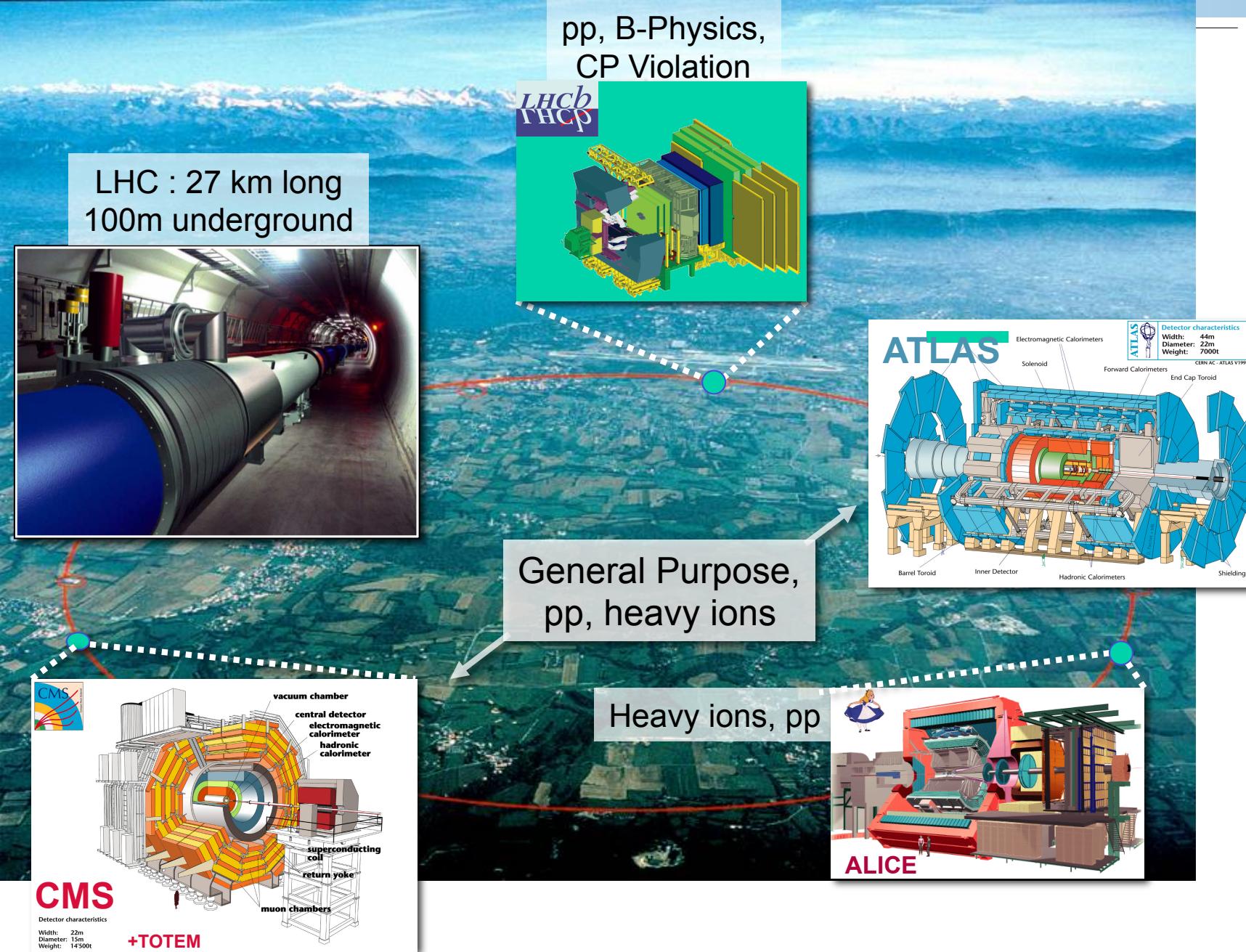
General Purpose,
pp, heavy ions



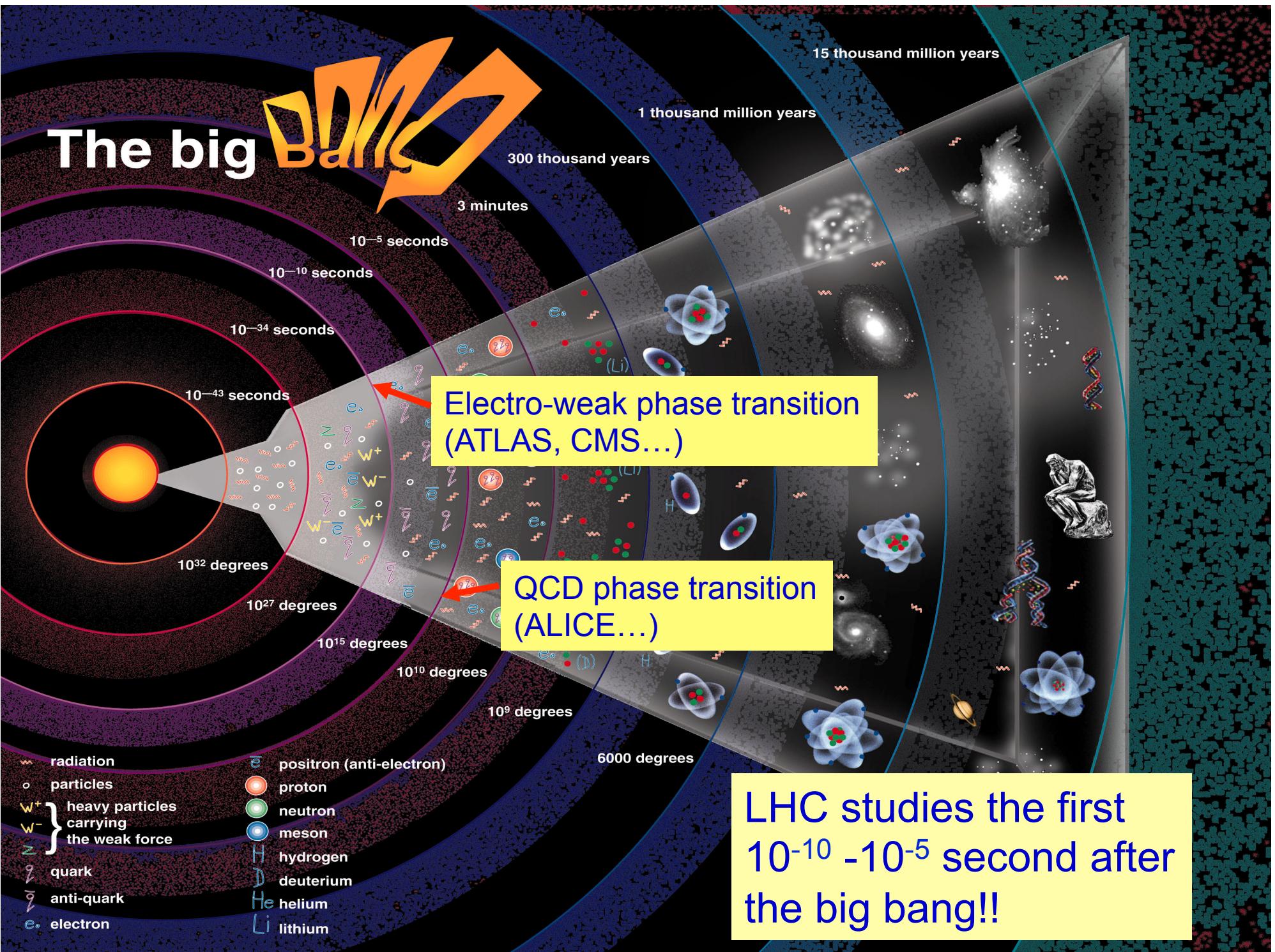
Heavy ions, pp



The Large Hadron Collider at CERN



The big Bang



Comparison with Direct Detection: Vector Case

