



Exotic states from LHCb

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On behalf of LHCb Collaboration

4th workshop on the XYZ particles
Beihang University, 2016.11.23

Standard and Exotic Hadrons

- Mesons beyond $q\bar{q}$ and baryons beyond three-quark configurations are called as exotic hadrons
- Their existence is not forbidden by QCD as long as they maintain color-singlet configurations
- Multiquark objects were proposed separately by Gell-Mann and Zweig in 1964
- Recent discoveries in heavy quark states have revived hopes for conclusive proofs for existence of exotic hadrons

STANDARD



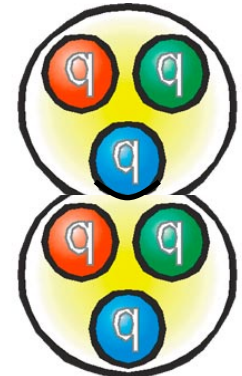
meson



baryon



mesonic molecule ?

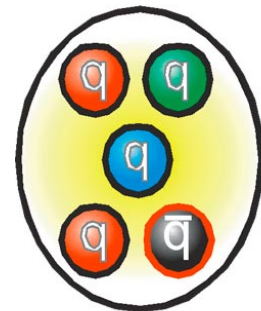


e.g. deuteron

EXOTIC



tetraquark ?

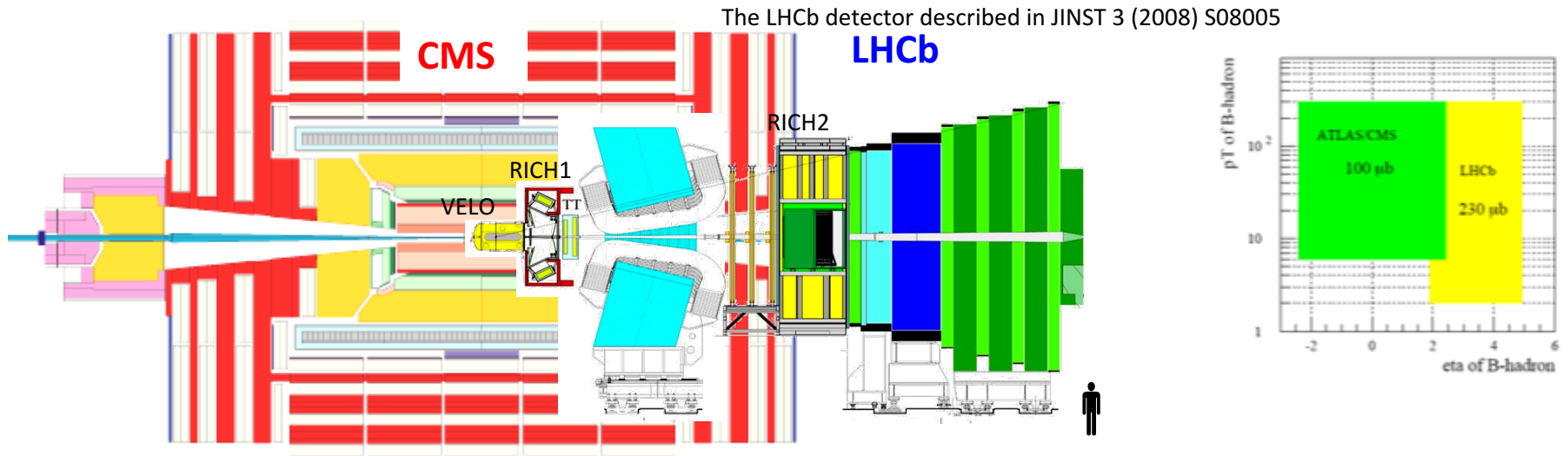


pentaquark ?



hybrid ?

Detectors and data sets



- LHCb detector advantages:

Momentum:

$$\Delta p/p = 0.4 \sim 0.6\% \quad (5 - 100 \text{ GeV}/c)$$

Mass :

$$\sigma_m = 8 \text{ MeV}/c^2 \text{ for } B \rightarrow J/\psi X \text{ (constrained } m_{J/\psi})$$

RICH $K - \pi$ separation:

$$\epsilon(K \rightarrow K) \sim 95\% \quad \text{mis-ID } \epsilon(\pi \rightarrow K) \sim 5\%$$

All LHCb results use total RUN-I data (3.0 fb^{-1})

$$Z_c(4430)^- \rightarrow \psi' \pi^- \text{ and} \\ P_c(4380)^+ P_c(4450)^+ \rightarrow J/\psi p$$

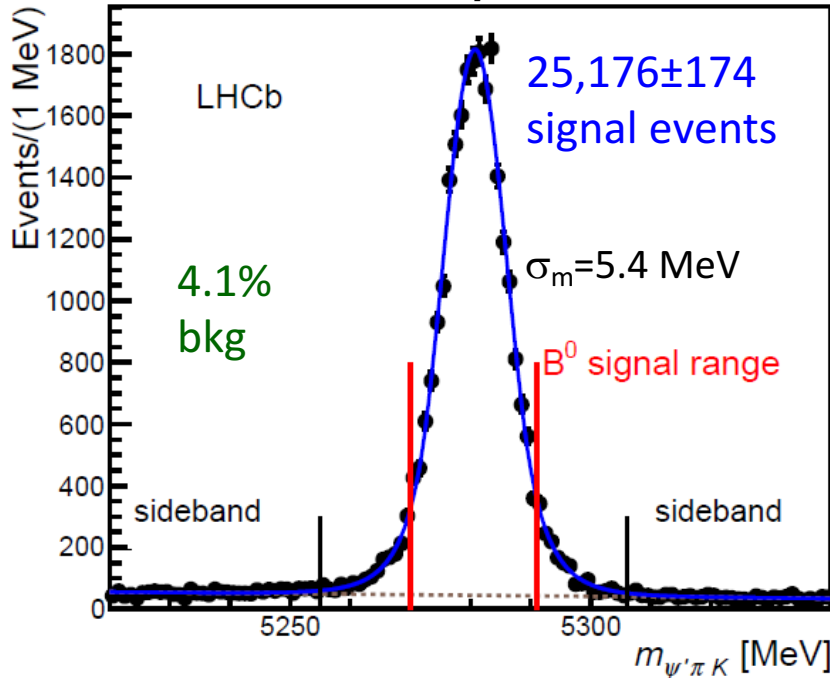
Full amplitude fits and
model-independent approaches

LHCb data samples (3 fb^{-1})



LHCb-PAPER-2014-014

PRL **112**, 222002 (2014)



vs. bkg in
Belle: 7.8%

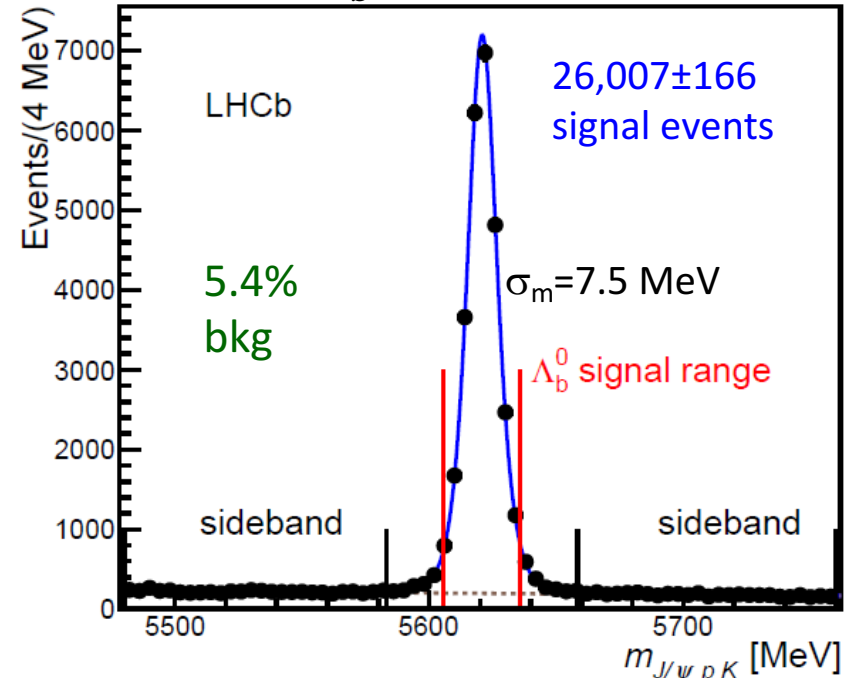
vs

Belle: $2,010 \pm 50$ [PRL **100**, 142001 (2008); PRD **88**, 074026 (2013)]

BaBar: $2,021 \pm 53$ [PRD **79**, 112001 (2009)]

LHCb-PAPER-2015-029

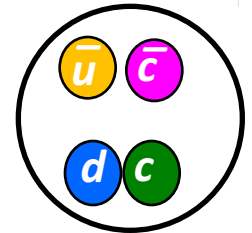
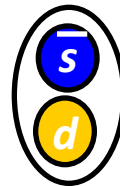
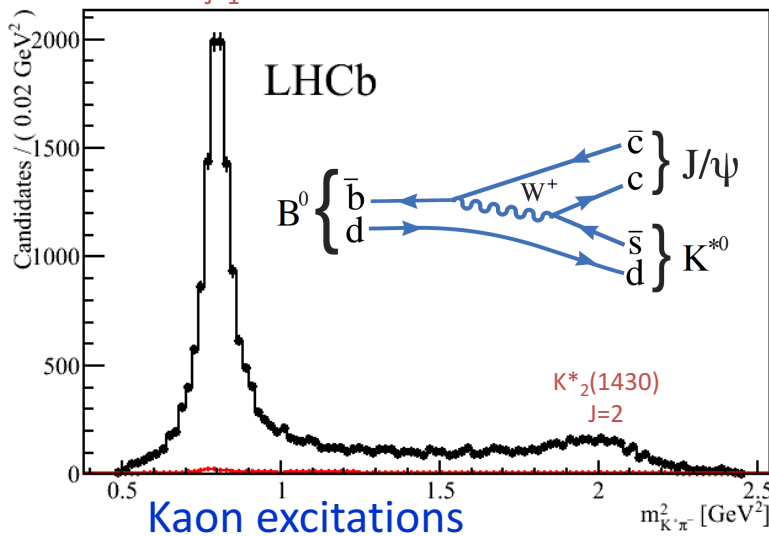
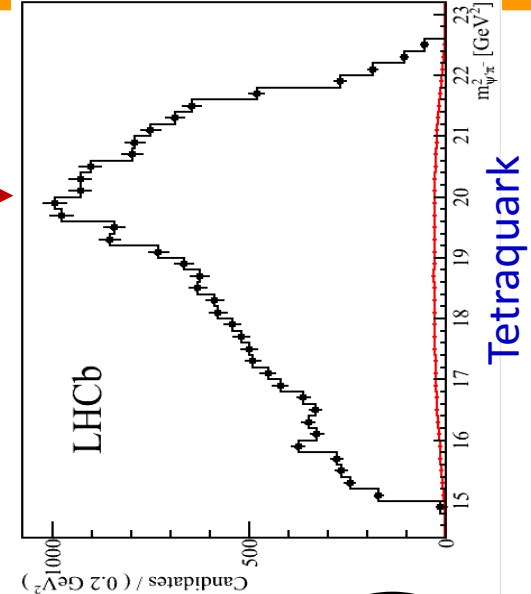
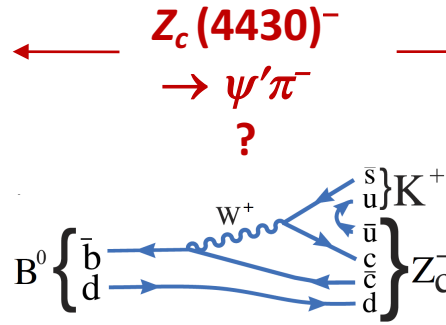
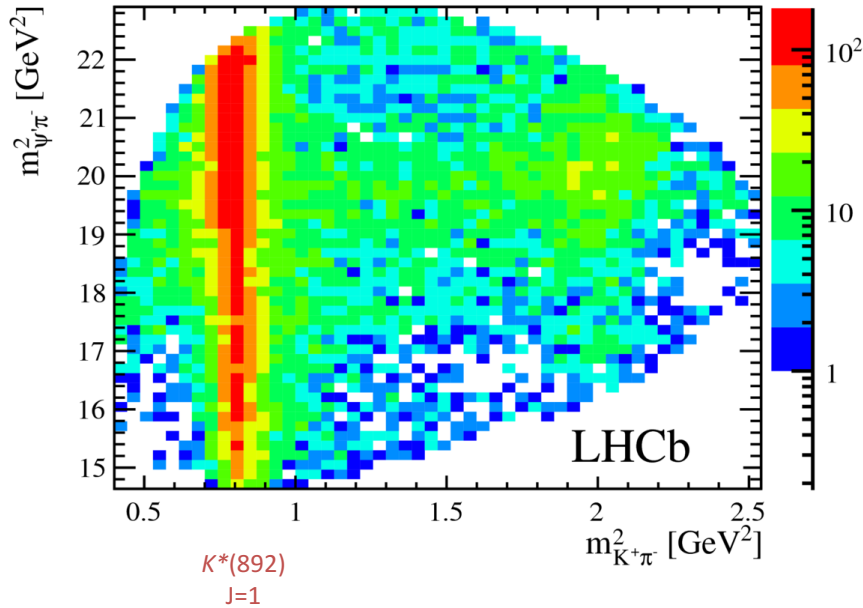
PRL **115**, 07201 (2015)



- > factor of 10 better statistics than at the B factories, at smaller bkg
- Comparable signal statistics and bkg levels between the B^0 and Λ_b data samples

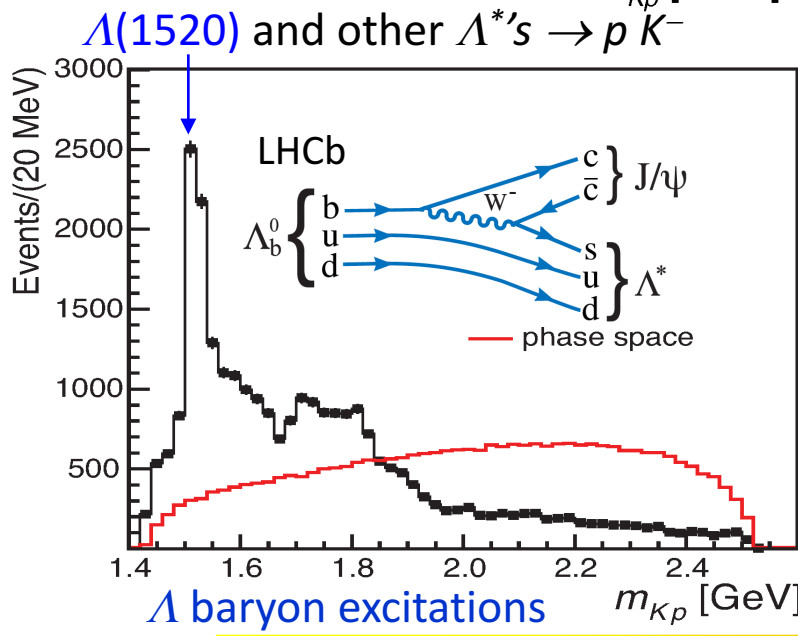
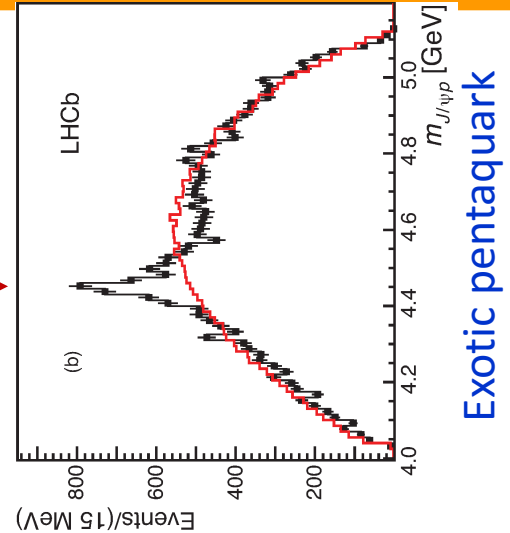
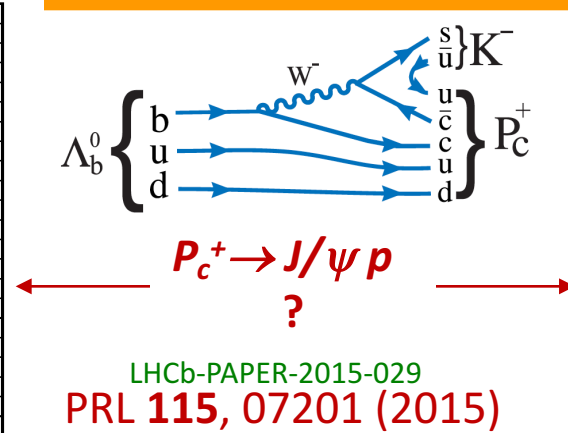
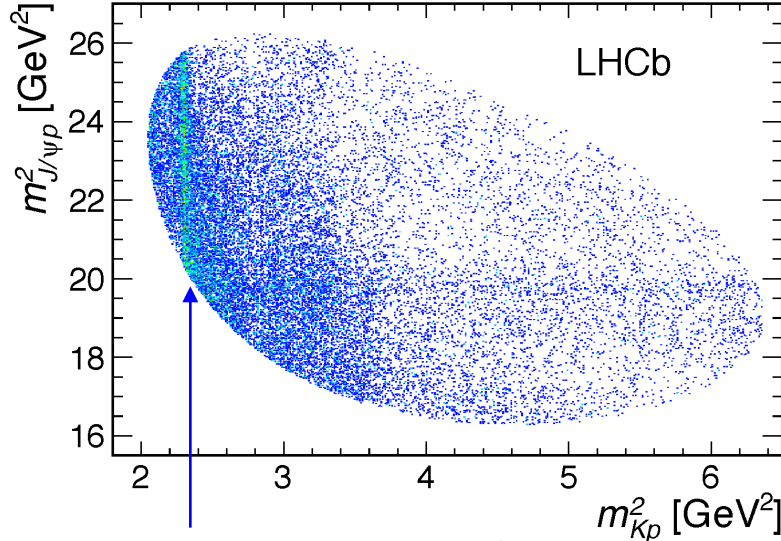
$B^0 \rightarrow J/\psi \pi^- K^+$

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PRL 112, 222002 (2014)

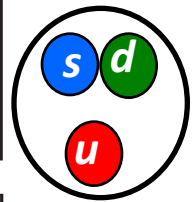
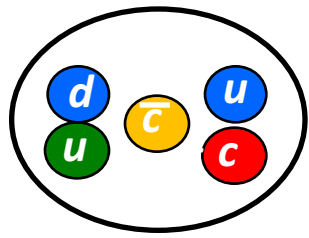


Is it a reflection of interfering K^* 's $\rightarrow \pi K^+$?
Proper amplitude analysis necessary to check

$\Lambda_b^0 \rightarrow J/\psi p K^-$: unexpected structure in $m_{J/\psi p}$



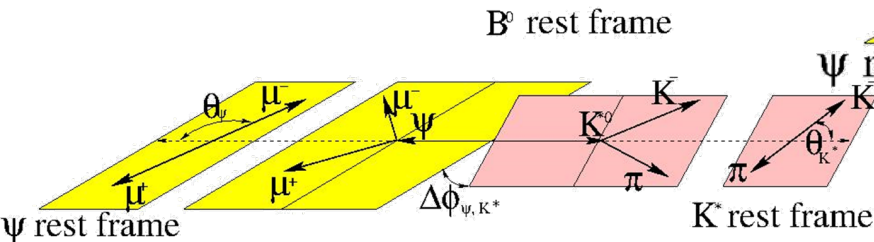
- Unexpected, narrow peak in $m_{J/\psi p}$



Is it a reflection of interfering Λ^* 's $\rightarrow p K^-$?
Proper amplitude analysis necessary to check

Full amplitude analyses

- Helicity frame



4D maximum likelihood fit

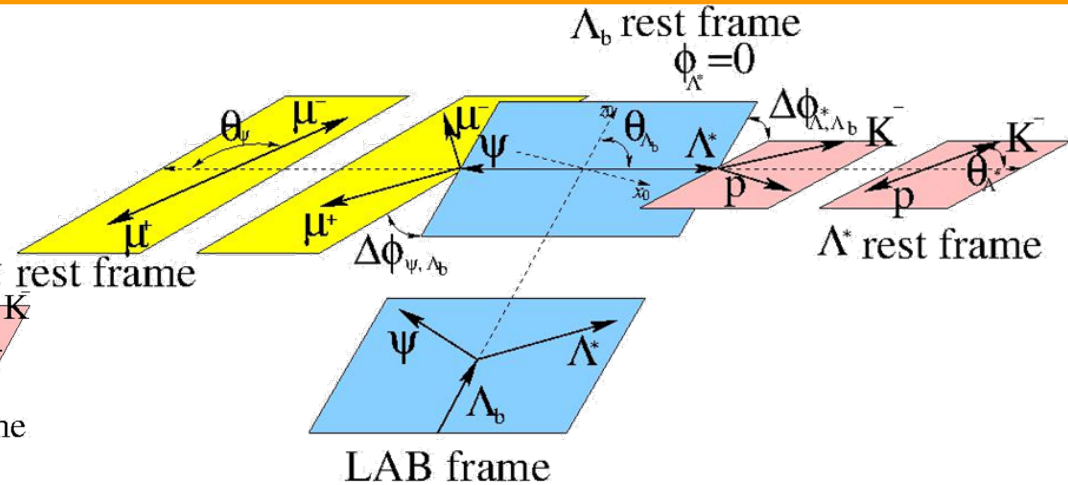
$$\Omega \equiv (\theta_{K^*}, \theta_{\psi}, \Delta\phi_{\psi, K^*})$$

$$\text{PDF}(\mathbf{m}_{K\pi/Kp}, \Omega) = \left| \text{MatrixEle}(\mathbf{m}_{K\pi/Kp}, \Omega \mid \underbrace{J_R^P, M_R, \Gamma_R, H_R}_{\text{Fitting parameters}}) \right|^2 \times \text{eff}(\mathbf{m}_{K\pi/Kp}, \Omega) + \text{PDF}_{\text{bkg}}(\mathbf{m}_{K\pi/Kp}, \Omega)$$

M_R, Γ_R fixed to known values, and varied within errors for systematics.

1-3 independent **complex** helicity couplings

H_R per K^* resonance



6D maximum likelihood fit

$$\Omega \equiv (\theta_{\Lambda_b}, \theta_{\Lambda^*}, \Delta\phi_{\Lambda^*, \Lambda_b}, \theta_{\psi}, \Delta\phi_{\psi, \Lambda_b})$$

4-6 independent **complex** helicity couplings

H_R per Λ^* resonance

The decay angles greatly increase discrimination power between resonances of different J^P
 Matrix elements for exotic decays are coherently added later to these for conventional resonance

Model of conventional resonances

Well established states from PDG

State	J^P	M_0 (MeV)	Γ_0 (MeV)	No high- M_0 high- J^P	
				# of complex couplings	
Only state with $P=(-1)^J$ decays to $K\pi$				Red.	Ext.
NR	0^+	—	—	1	1
$K^*(800)^0$	0^+	682	547	1	1
$K^*(892)^0$	0^+	896	49	3	3
$K^*(1410)^0$	1^-	1414	232	3	3
$K^*(1430)^0$	0^+	1425	270	1	1
$K_2^*(1430)^0$	2^+	1432	109	3	3
$K^*(1680)^0$	1^-	1717	322	3	3
$K_3^*(1780)^0$	3^-	1776	159	0	3
Total # of free parameters				28	34

Using Extend model to set significance of exotic states

State	J^P	M_0 (MeV)	Γ_0 (MeV)	No high- M_0 high- J^P & limit L		All states all L
				# of complex couplings		
				Red.	Ext.	
$\Lambda(1405)$	$1/2^-$	1405	50	3	4	
$\Lambda(1520)$	$3/2^-$	1520	16	5	6	
$\Lambda(1600)$	$1/2^+$	1600	150	3	4	
$\Lambda(1670)$	$1/2^-$	1670	35	3	4	
$\Lambda(1690)$	$3/2^-$	1690	60	5	6	
$\Lambda(1800)$	$1/2^-$	1800	300	4	4	
$\Lambda(1810)$	$1/2^+$	1810	150	3	4	
$\Lambda(1820)$	$5/2^+$	1820	80	1	6	
$\Lambda(1830)$	$5/2^-$	1830	95	1	6	
$\Lambda(1890)$	$3/2^+$	1890	100	3	6	
$\Lambda(2100)$	$7/2^-$	2100	200	1	6	
$\Lambda(2110)$	$5/2^+$	2110	200	1	6	
$\Lambda(2350)$	$9/2^+$	2350	150	0	6	
$\Lambda(2585)$	$5/2^-?$	2585	200	0	6	
Total # of free parameters				64	146	

- Baryons have more than doubling of known states to include

Fits with conventional resonances only

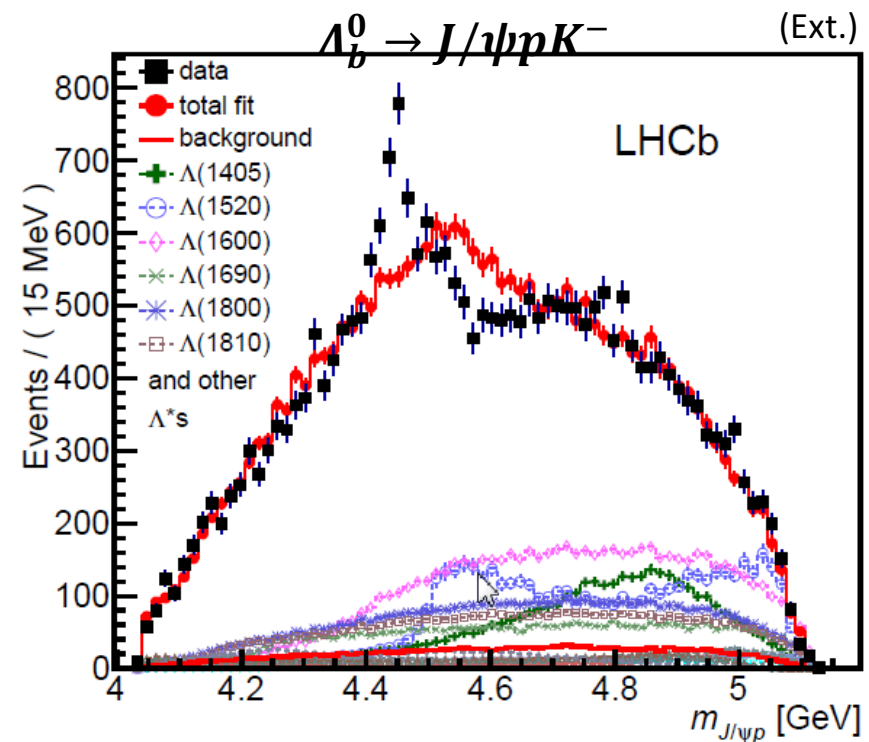
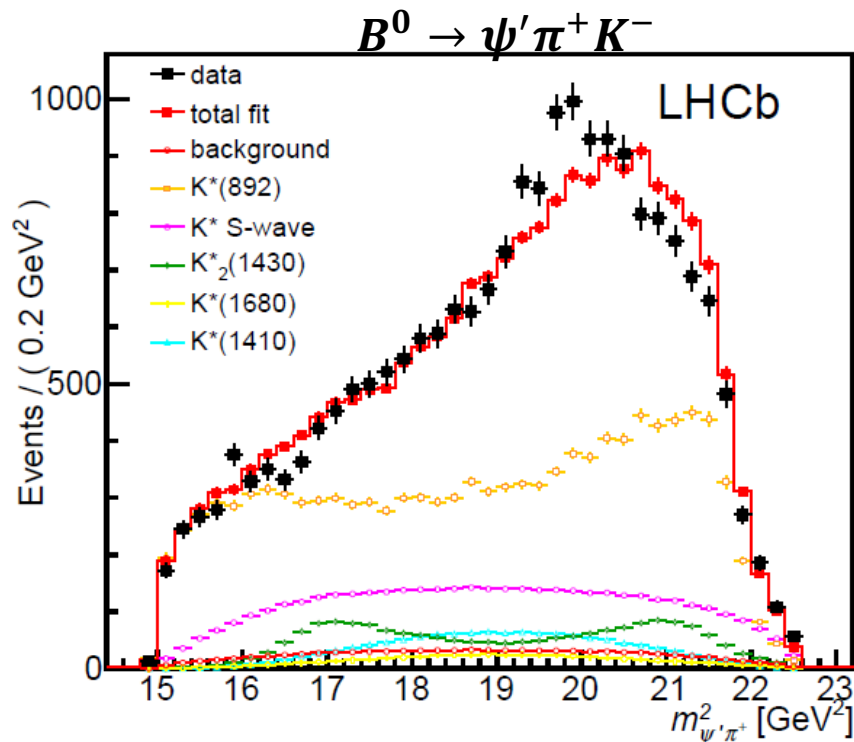


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PRL 112, 222002 (2014)

LHCb-PAPER-2015-029

PRL 115, 07201 (2015)



- Two models for conventional resonances: **Extended and Reduced**, cross-check with each other
- Extended uses all established resonances, for significances and syst.
- Reduced keeps only significant contributions

Cannot describe the data with the conventional resonances (Extended) alone

Fits including exotic hadrons

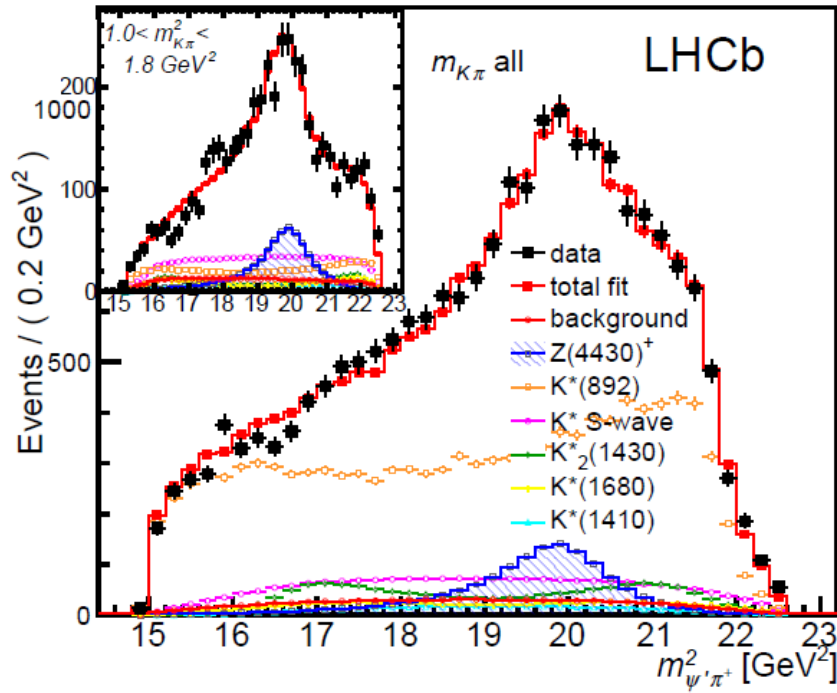
LHCb-PAPER-2014-014

PRL 112, 222002 (2014)

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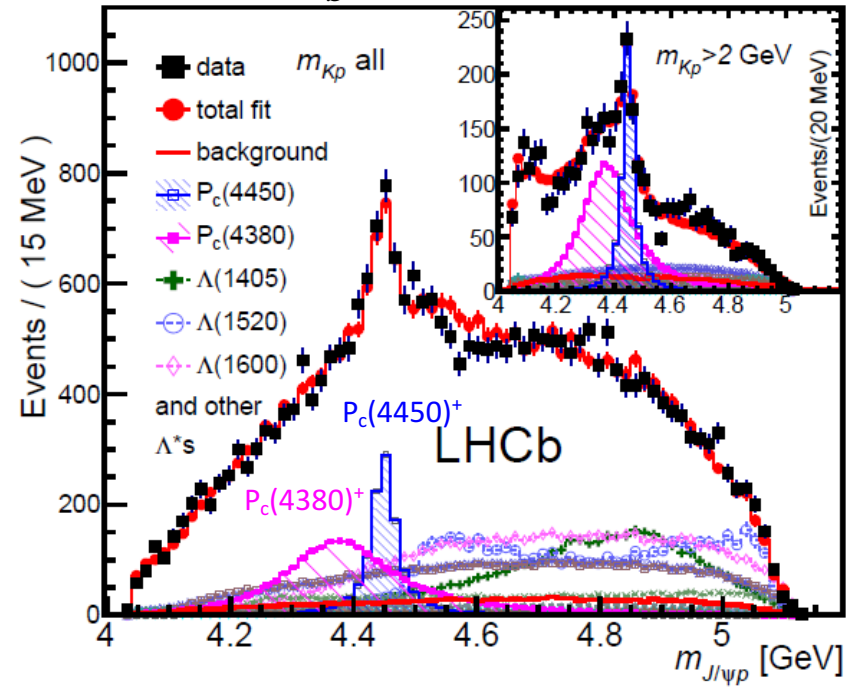
PRL 115, 07201 (2015)

$$B^0 \rightarrow \psi' \pi^+ K^-$$



$$\Lambda_b^0 \rightarrow J/\psi p K^-$$

(Red.)



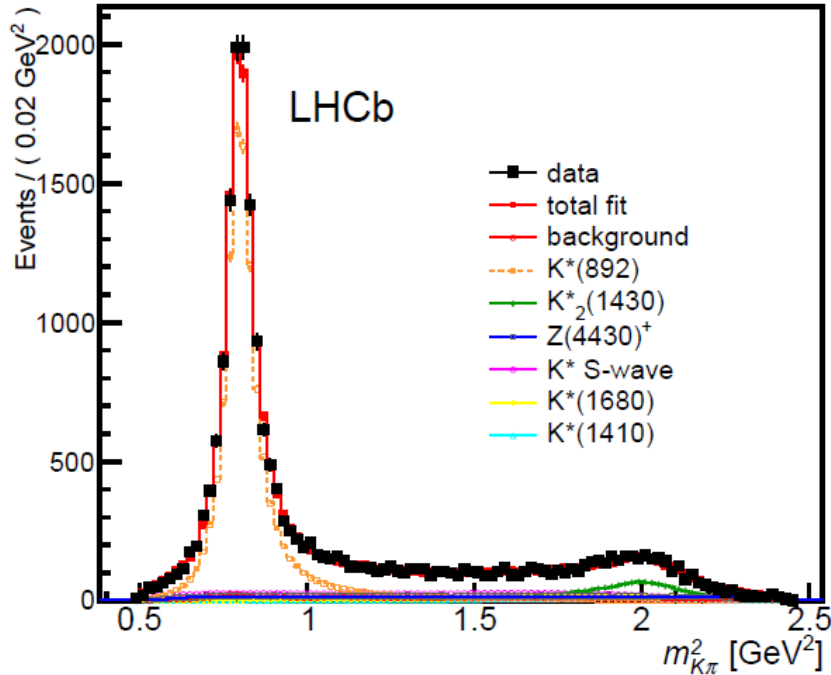
State	Mass (MeV)	Width (MeV)	Fit frac. (%)	Sig.	State	Mass (MeV)	Width (MeV)	Fit frac. (%)	Sig.
$Z_c(4430)^+$	$4475 \pm 7_{-25}^{+15}$	$172 \pm 13_{-34}^{+37}$	$5.9 \pm 0.9_{-3.3}^{+1.5}$	14σ	$P_c(4450)^+$	$4449.8 \pm 1.7 \pm 2.5$	$39 \pm 5 \pm 19$	$4.1 \pm 0.5 \pm 1.1$	12σ
Belle	$4485 \pm 22_{-11}^{+28}$	$200 \pm 46_{-35}^{+26}$	$10.3 \pm 3.5_{-2.3}^{+4.3}$	5σ	$P_c(4380)^+$	$4380 \pm 8 \pm 29$	$205 \pm 18 \pm 86$	$8.4 \pm 0.7 \pm 4.2$	9σ

• $J^P=1^+$ at 9.7σ incl. syst. (in Belle at 3.4σ)

• Best fit has $J^P=(3/2^-, 5/2^+)$, also $(3/2^+, 5/2^-)$ & $(5/2^+, 3/2^-)$ cannot be ruled out

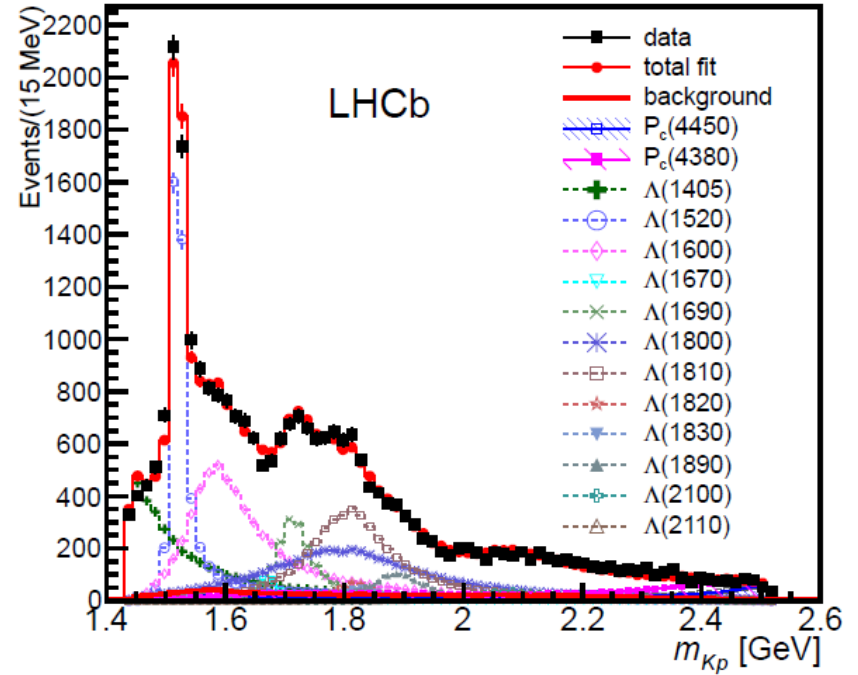
Fits including exotic hadrons

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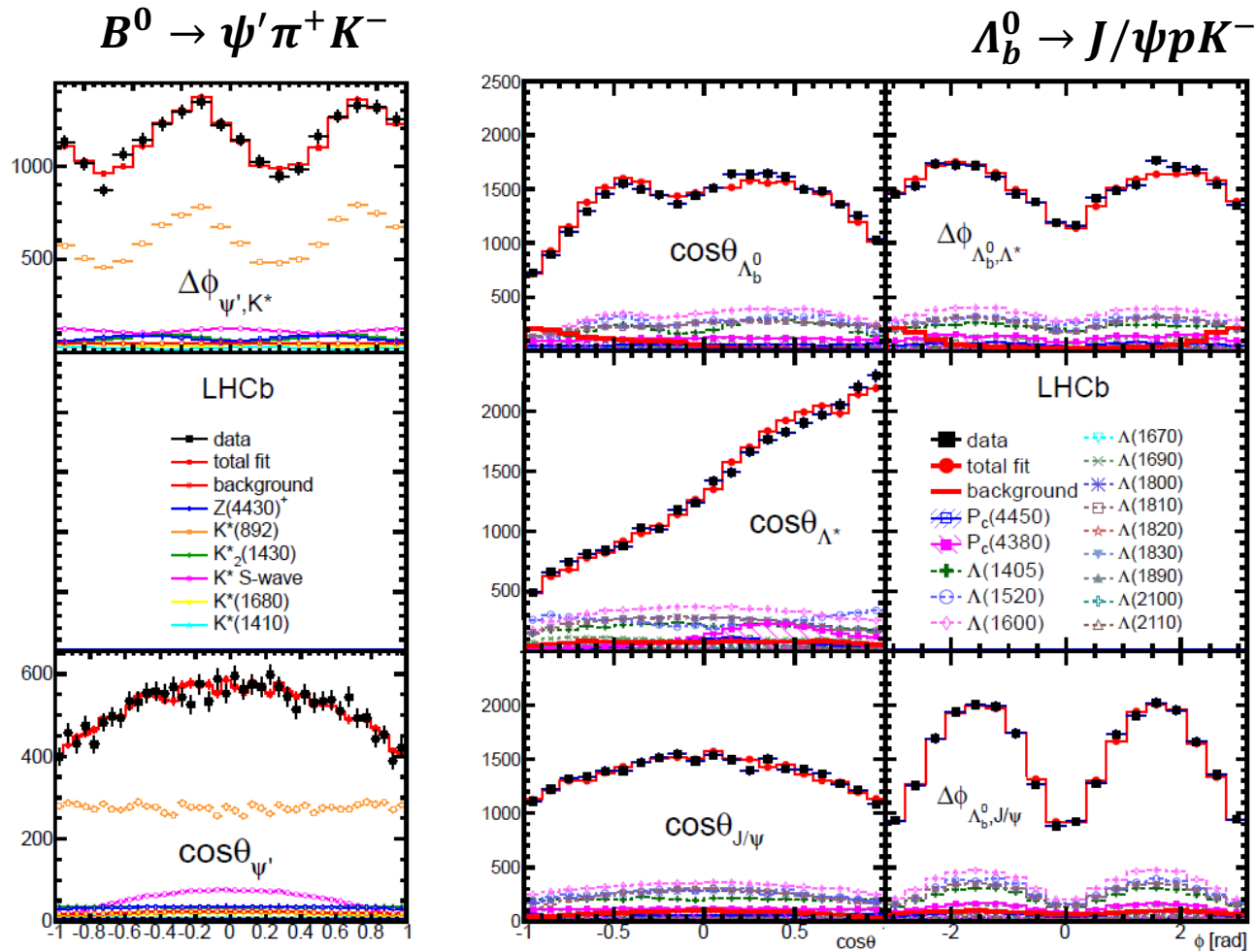
$$\Lambda_b^0 \rightarrow J/\psi p K^-$$

(Red.)



- Data are well described by the fits
- Conventional resonances dominate the rate
- Exotics spread across wide range of these masses

Decay angular projections

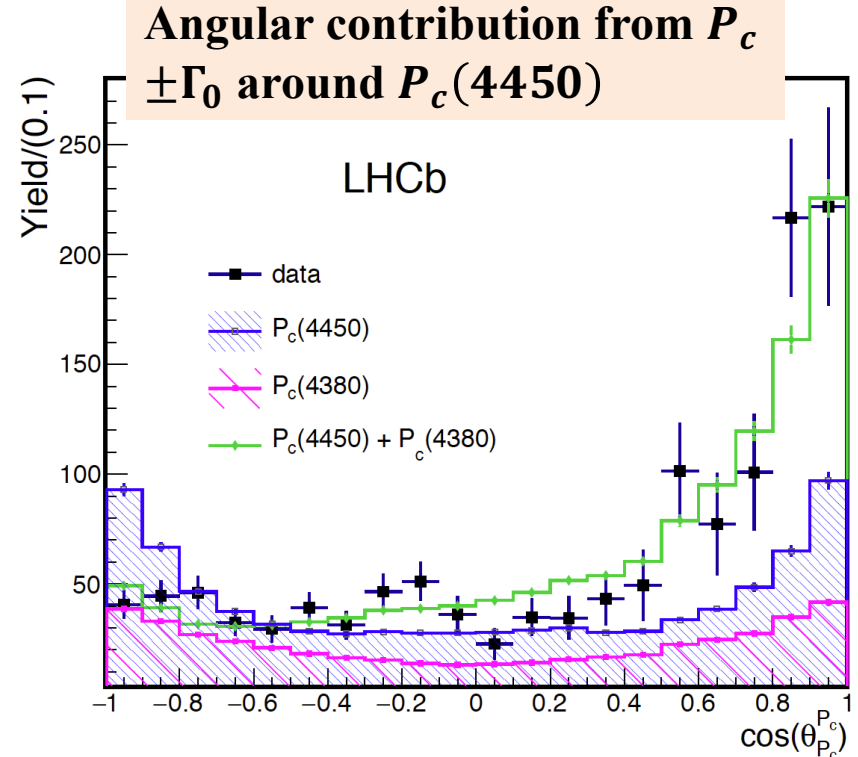
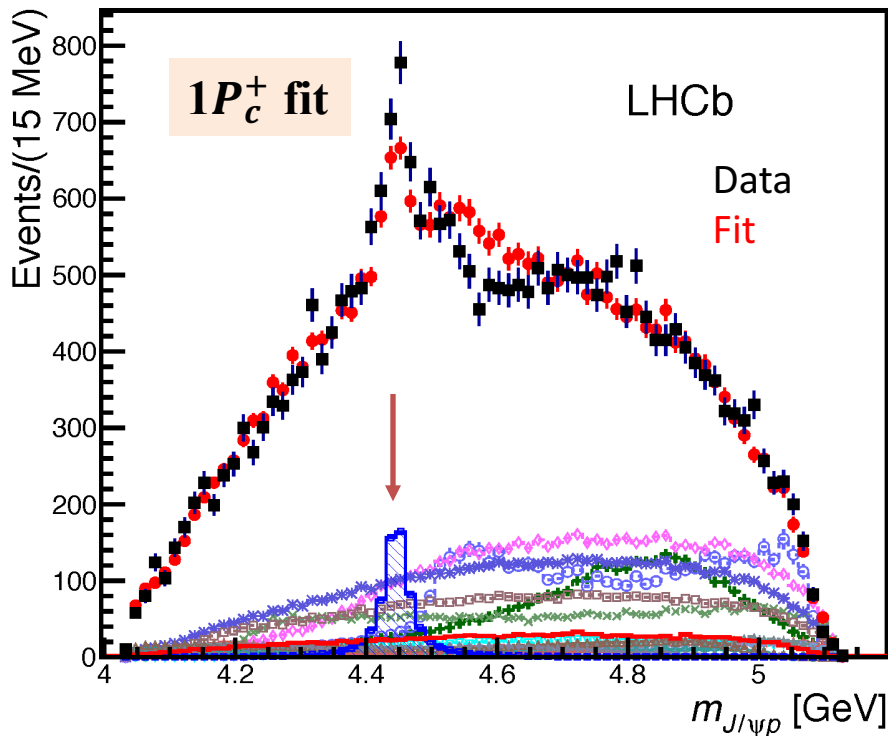


- They greatly increase discrimination power between resonances of different J^P
- Correctly describe interference

Why two P_c^+ with opposite parity



- 1 P_c^+ cannot fit well the peak
- Strong asymmetry angular distribution from P_c^+
- Can only be explained by interference two opposite parity P_c^+



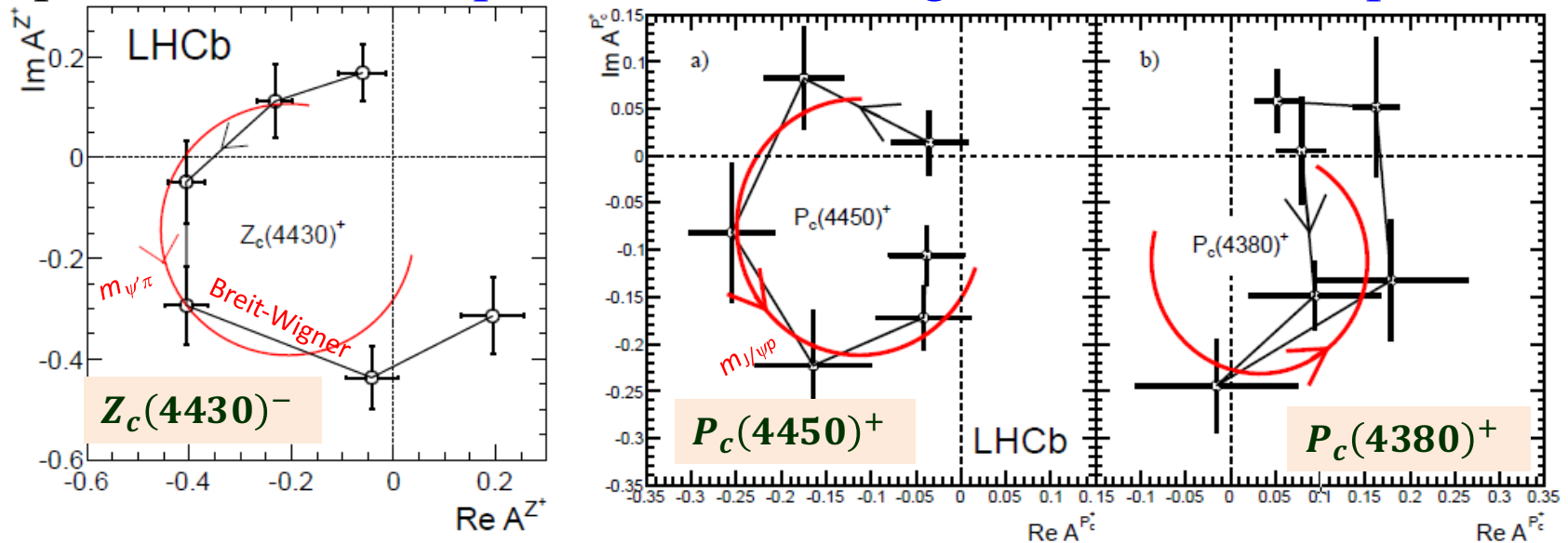
CERN-THESIS-2016-086

Argand diagrams for resonant behavior

LHCb-PAPER-2015-029

PRL 115, 07201 (2015)

Exotic hadron amplitudes modelled by 6 complex coefficients near the peak mass, to compare with Breit-Wigner resonance amplitude



Good evidence for resonant character

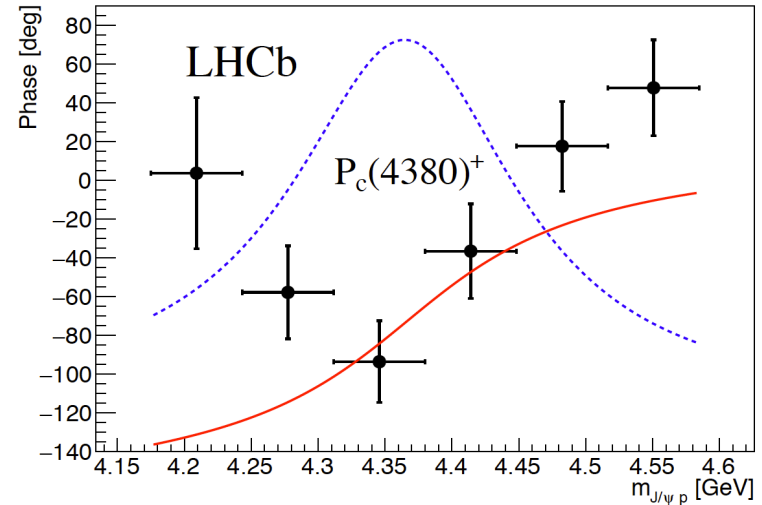
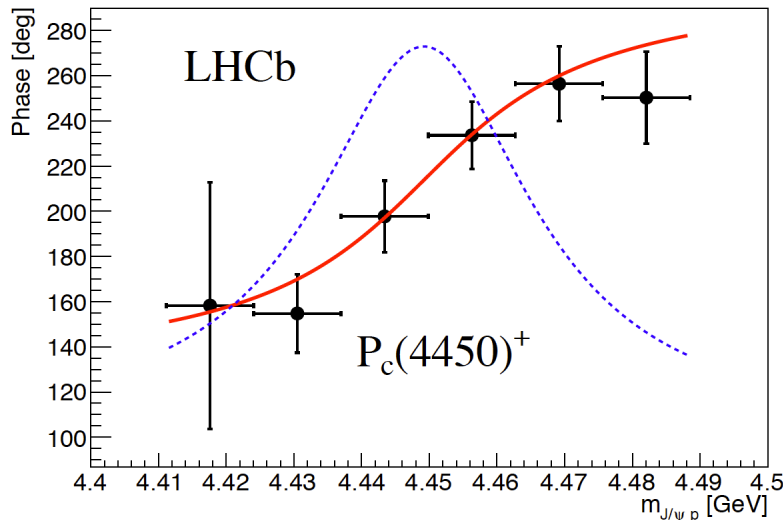
Large errors: not conclusive

- The results are still dependent on the model of conventional hadrons
- Simultaneous PWA of the latter is not possible since exotics reflect into variables characterizing conventional hadrons
- However, we can assume exotics are not present and test for their presence in model-independent way - next 5 slides

Argand diagrams for resonant behavior

CERN-THESIS-2016-086

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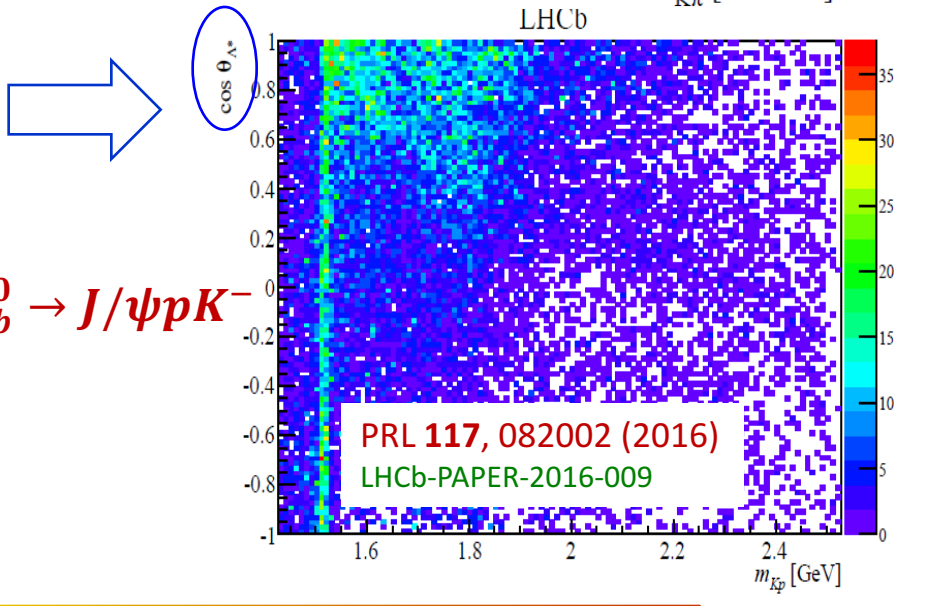
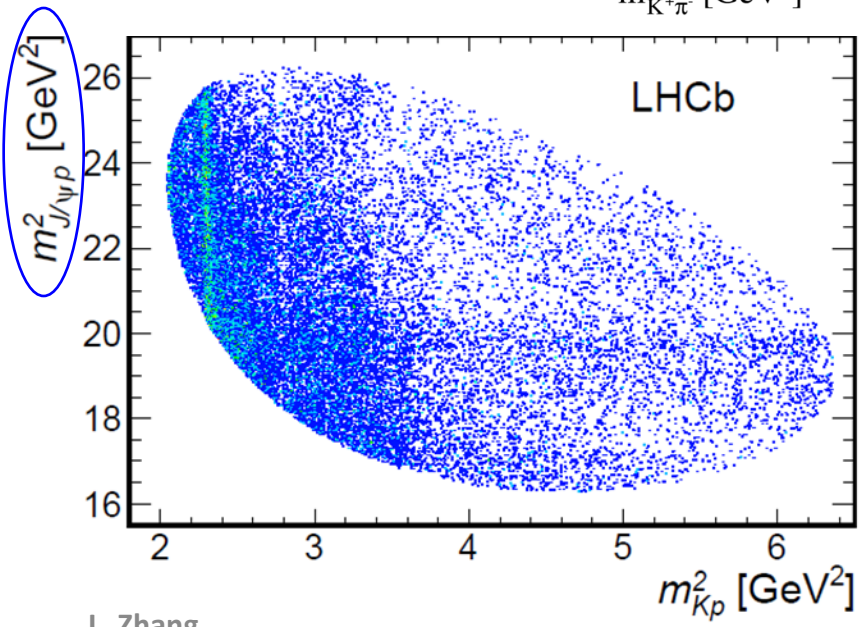
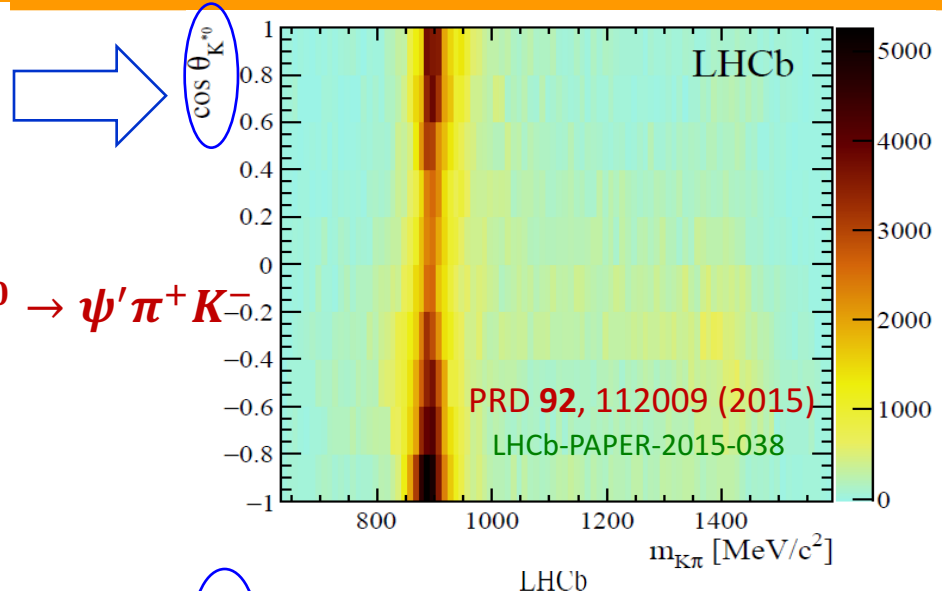
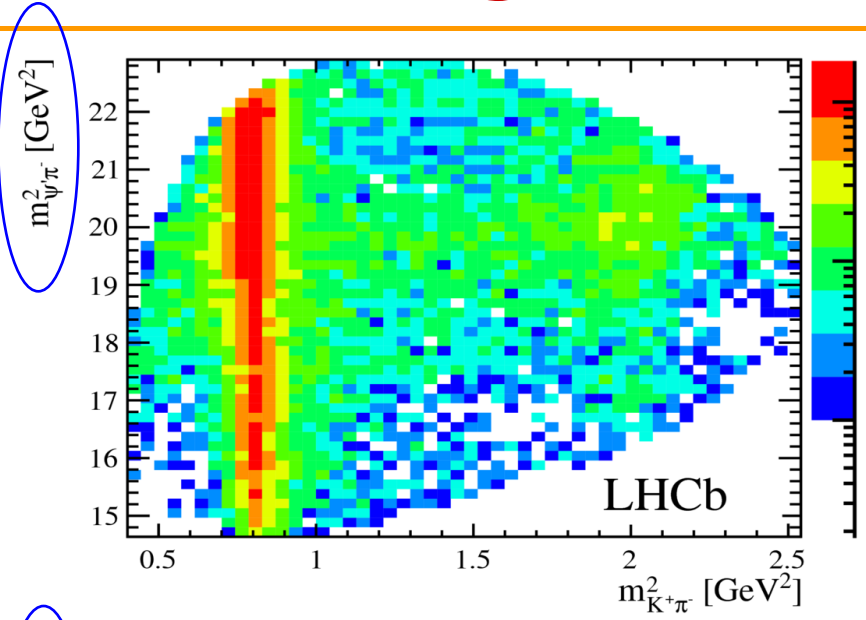


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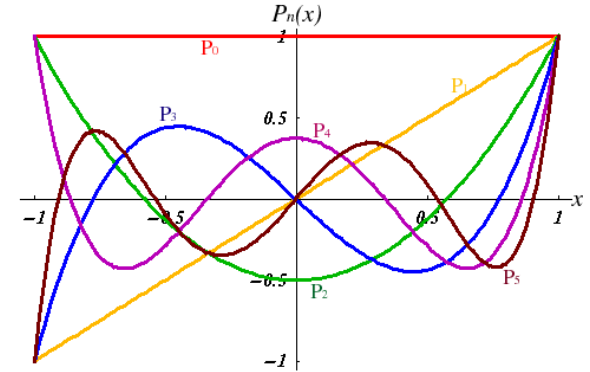
Rectangular Dalitz plane



Legendre moments

$$\frac{dN}{d \cos \theta} = \sum_{l=0}^{l_{\max}} \langle P_l^U \rangle P_l(\cos \theta) \quad \theta = \theta_{K^*} \quad \text{or} \quad \theta_{\Lambda^*}$$

$$\langle P_l^U \rangle = \int_{-1}^{+1} \frac{dN}{d \cos \theta} P_l(\cos \theta) d \cos \theta \propto \sum_{i=1}^{n_{\text{events}}} \frac{1}{\mathcal{E}_i} P_l(\cos \theta_i)$$



K^*/Λ^* can contribute only to low-order moments

Reflections of exotic hadrons can contribute to low **and high** order moments:

K^*/Λ^* -only hypothesis called H_0

$$l_{\max} = 2J_{\max}$$

- Detecting non-zero moments above $2J_{\max}$ signals presence of exotics

J_{\max} is the highest spin of K^*/Λ^* resonance possible

Setting l_{\max} as function of $m_{K\pi}/m_{K\rho}$

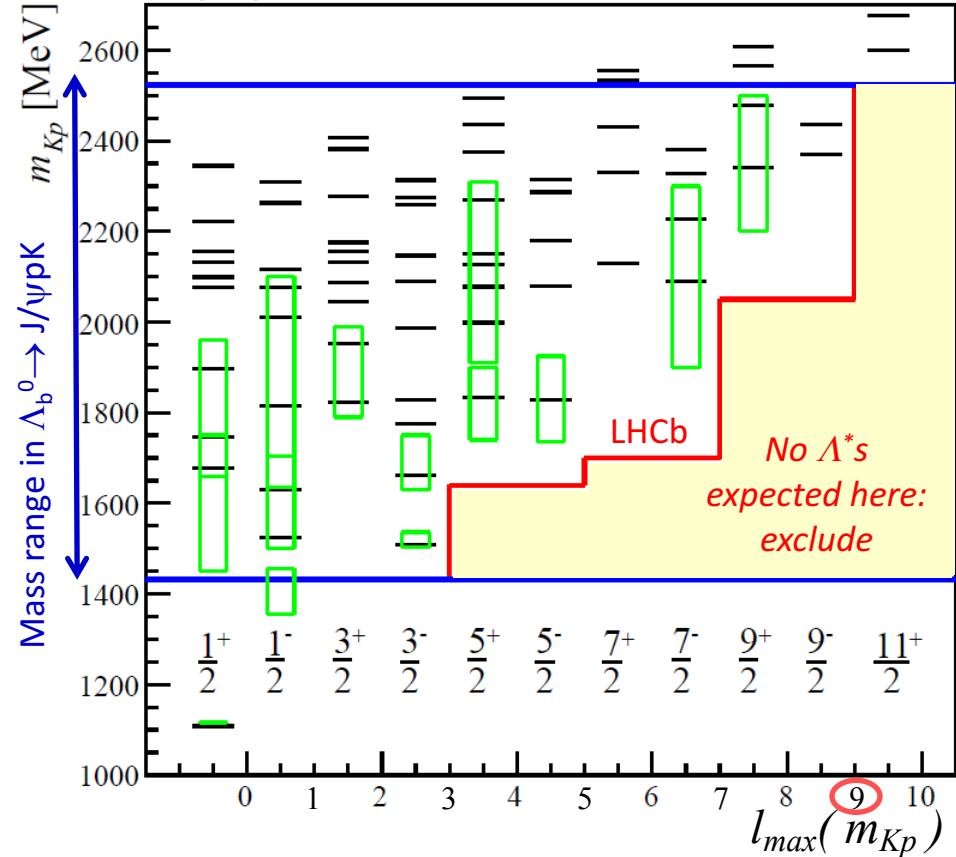
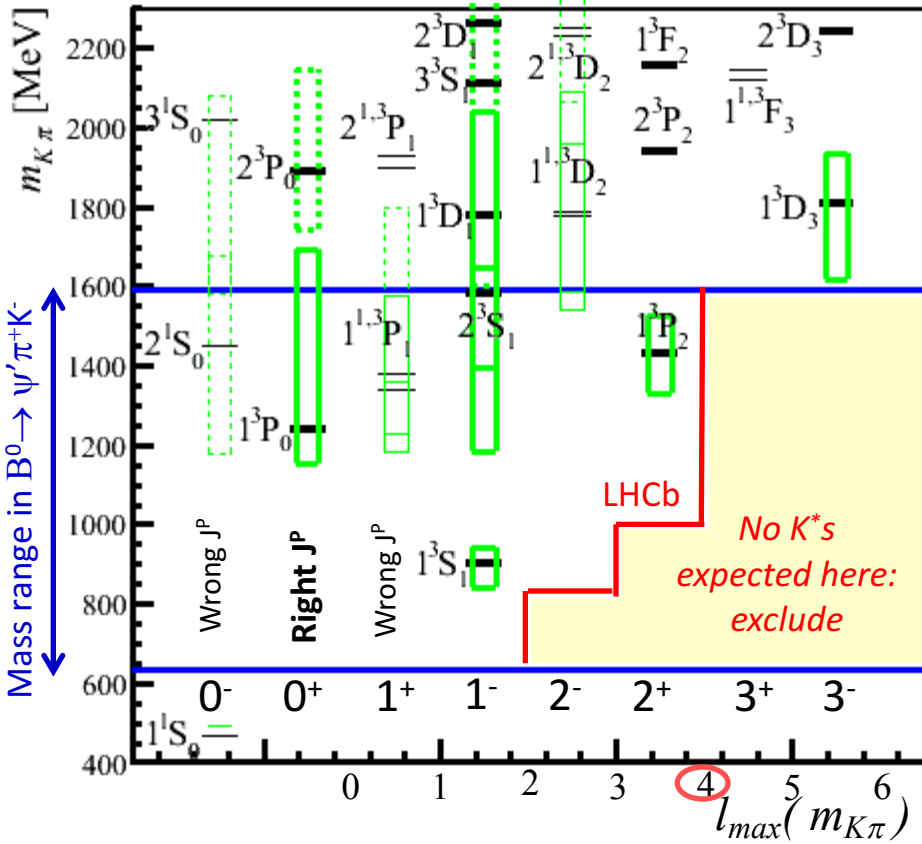
From **known K^*/Λ^* resonances**, quark model predictions as a guide

Much fewer known states than predicted!

K^* mass predictions by Godfrey-Isgur, PRD 32, 189 (1985)

Known K^*/Λ^* states: boxes $M_0 \pm \Gamma_0$

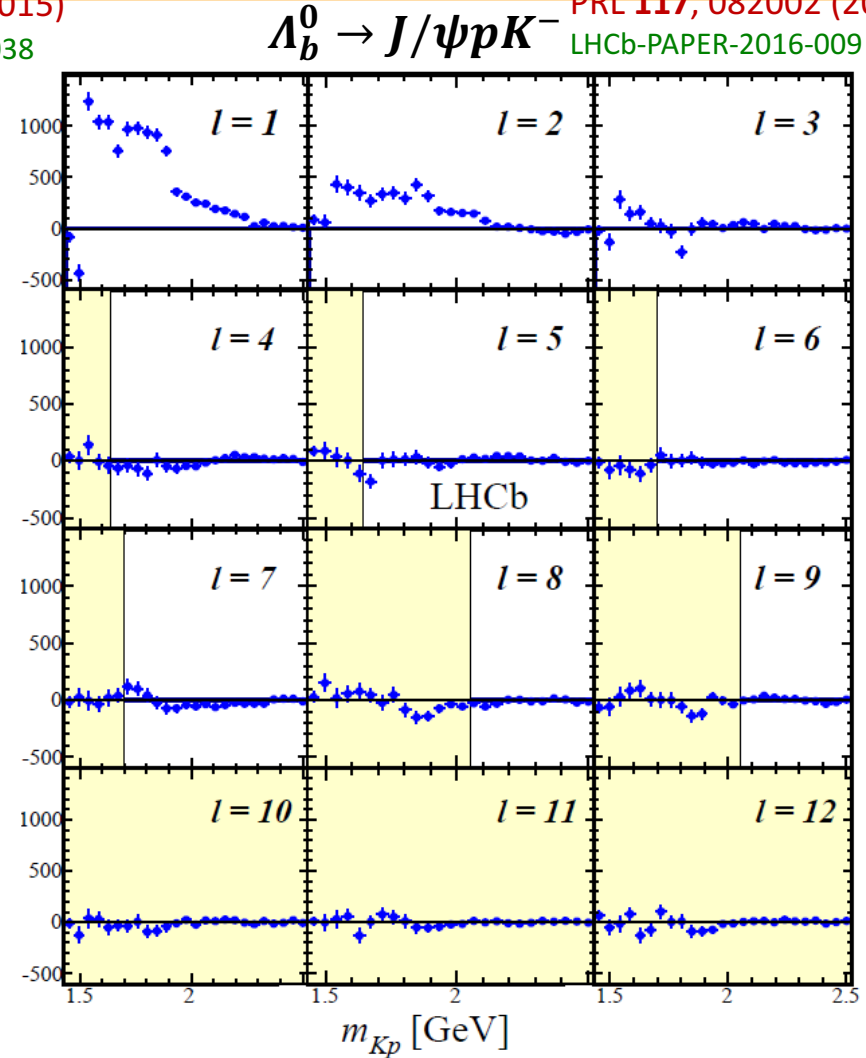
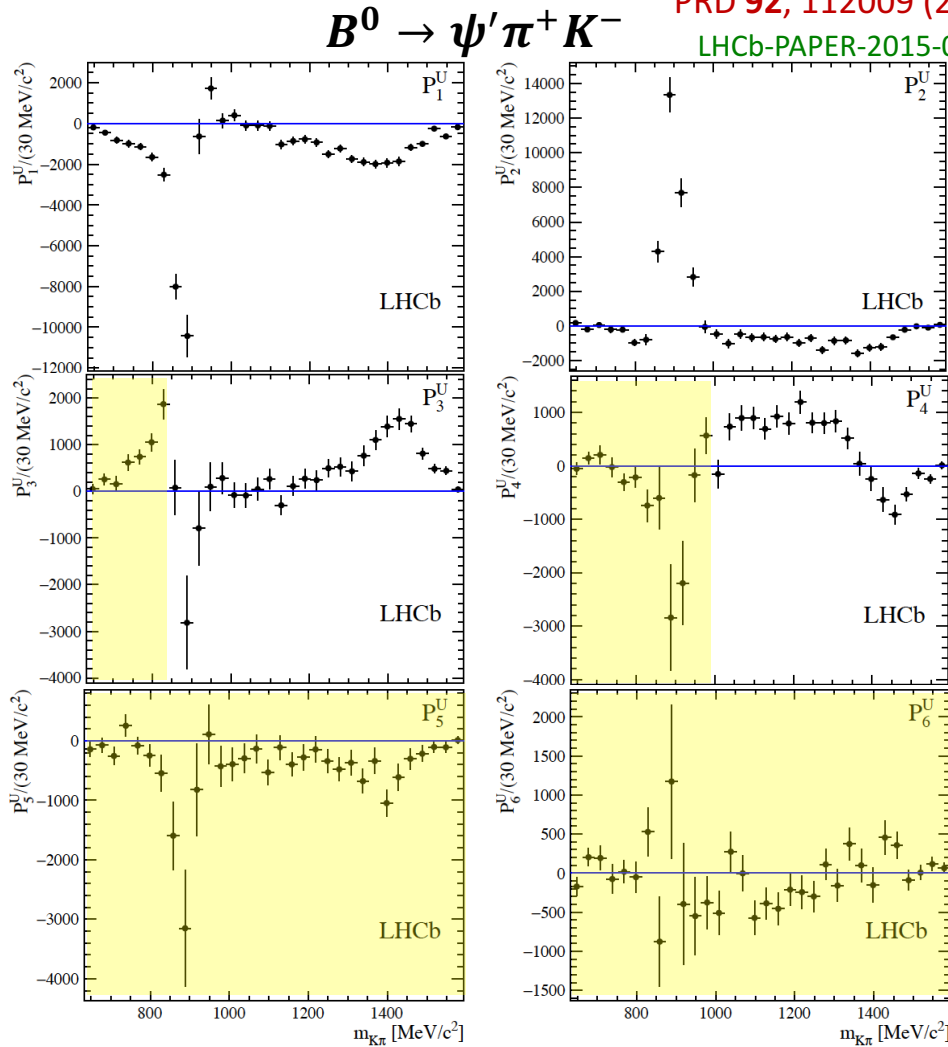
Λ^* mass predictions by Loring-Metsch-Petry EPJ, A10, 447 (2001)



Moments coefficients

PRD 92, 112009 (2015)
LHCb-PAPER-2015-038

PRL 117, 082002 (2016)
LHCb-PAPER-2016-009



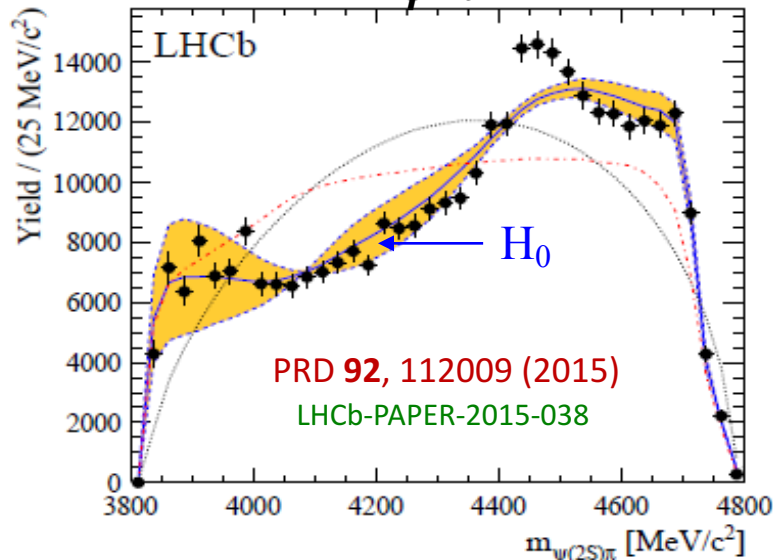
Non-zero values in the hatched area imply exotic contribution

Null exotic hypothesis (H_0) vs data

- Toy simulates the reflection of mass and angular structure of $K\pi/K\rho$ onto $m_{\psi'\pi/\psi\rho}$
- Limits of l_{\max} used (i.e. zero moments for $> l_{\max}$ order)
- $m_{\psi'\pi/\psi\rho}$ can not be explained by the reflections of conventional (non)resonances alone

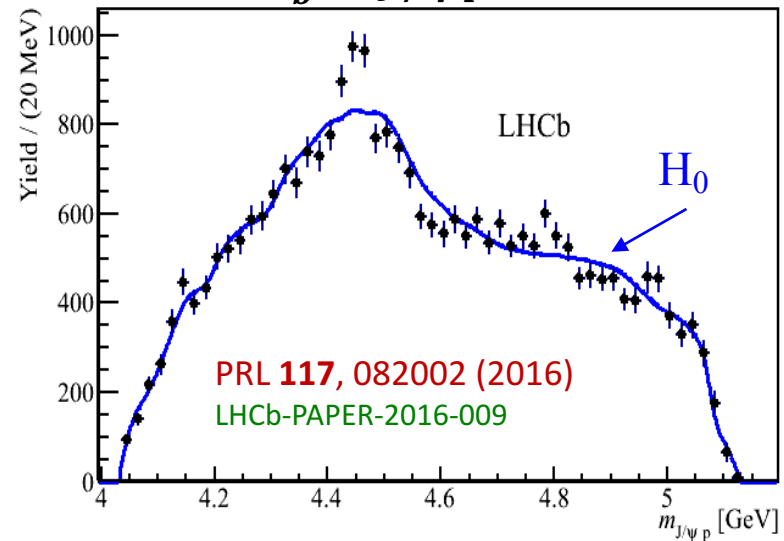
H_0 rejected by 15.2σ

$B^0 \rightarrow \psi'\pi^+ K^-$



H_0 rejected by 9.0σ

$\Lambda_b^0 \rightarrow J/\psi p K^-$





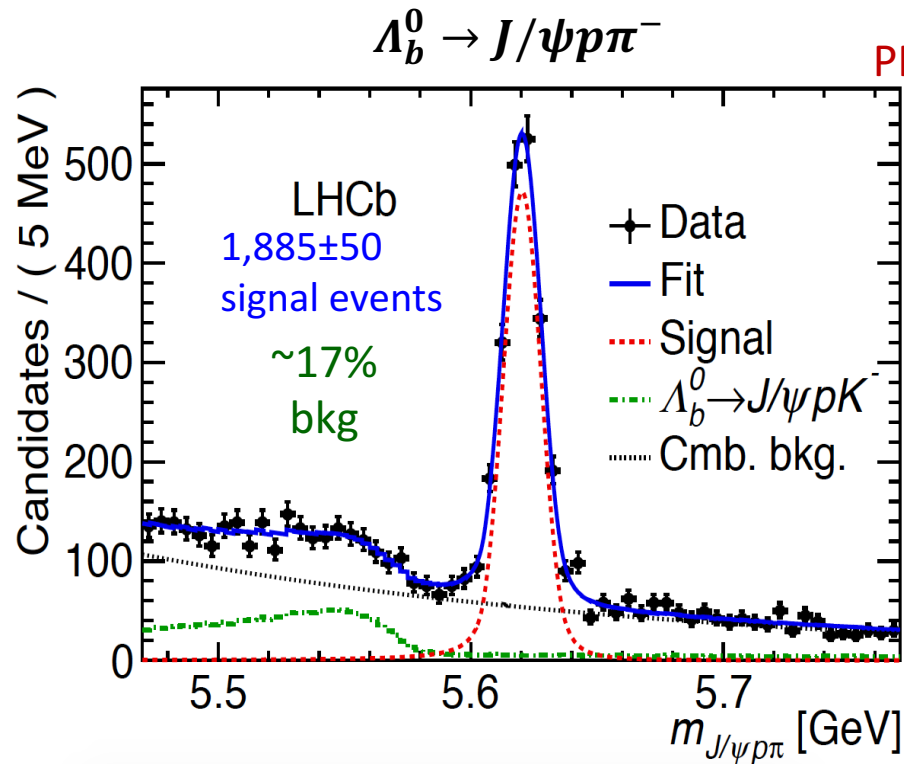
Exotic contributions in Cabibbo-suppressed decays $\Lambda_b^0 \rightarrow J/\psi p \pi^-$

Cabibbo suppressed decays $\Lambda_b^0 \rightarrow J/\psi p \pi^-$



LHCb-PAPER-2016-015

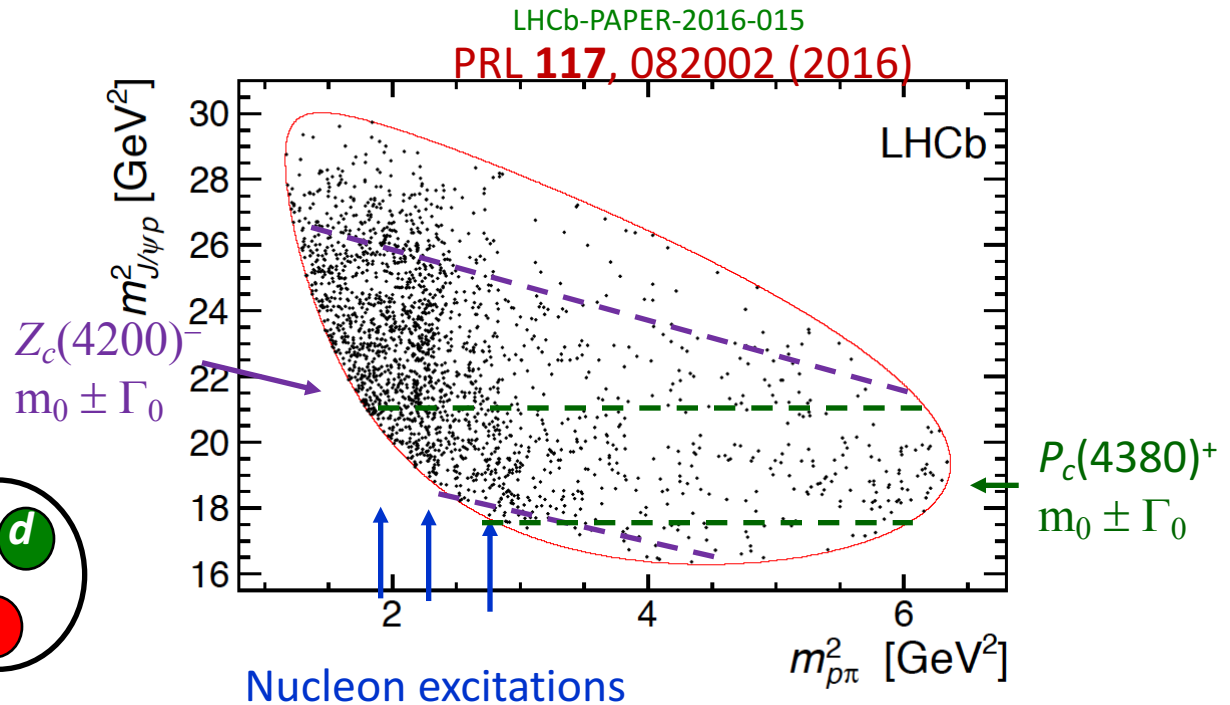
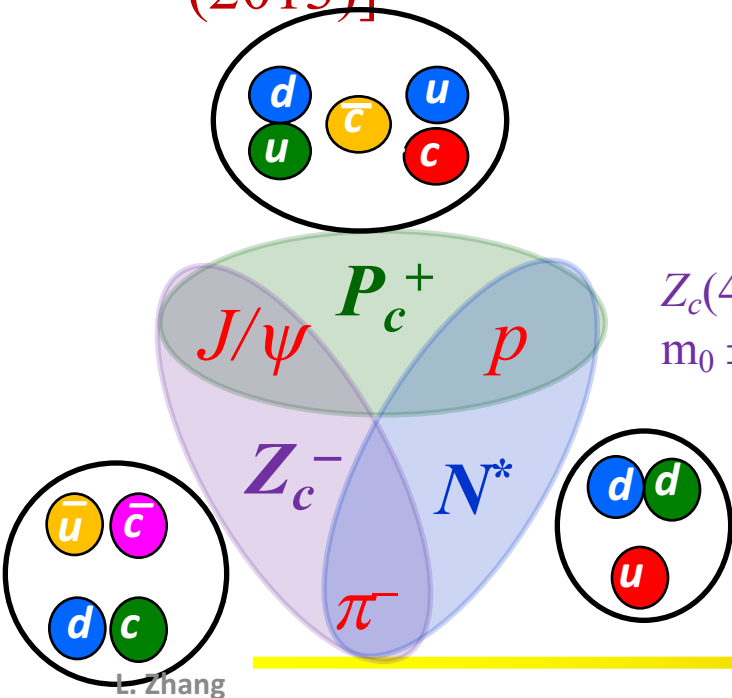
PRL **117**, 082003 (2016)



- More than a factor of 10 lower signal statistics in $\Lambda_b^0 \rightarrow J/\psi p \pi^-$ analysis than in $\Lambda_b^0 \rightarrow J/\psi p K^-$

$\Lambda_b^0 \rightarrow J/\psi p \pi^-$: Cabibbo suppressed

- Less statistics, more complex because of Z_c^-
- Here the exotic hadron contributions are examined for $P_c(4380)^+$, $P_c(4450)^+ \rightarrow J/\psi p$ and $Z_c(4200)^- \rightarrow J/\psi \pi^-$
- $Z_c(4200)^- m_0 = 4196_{-32}^{+35}$ MeV, $\Gamma = 370_{-149}^{+99}$ MeV $J^P=1^+$ by Belle (6.2σ) in $B^0 \rightarrow J/\psi \pi^- K^+$ decays [PRD 88, 074026 (2013)]



Model of N^* and exotic states

Better established states from PDG

Only significant states limit L All states limit L

$\Lambda_b^0 \rightarrow J/\psi p \pi^-$

State	J^P	M_0 (MeV)	Γ_0 (MeV)	# of complex couplings	
				Red.	Ext.
NR $p\pi$	$1/2^-$	-	-	4	4
$N(1440)$	$1/2^+$	1430	350	3	4
$N(1520)$	$3/2^-$	1515	115	3	3
$N(1535)$	$1/2^-$	1535	150	4	4
$N(1650)$	$1/2^-$	1655	140	1	4
$N(1675)$	$5/2^-$	1675	150	3	5
$N(1680)$	$5/2^+$	1685	130	0	3
$N(1700)$	$3/2^-$	1700	150	0	3
$N(1710)$	$1/2^+$	1710	100	0	4
$N(1720)$	$3/2^+$	1720	250	3	5
$N(1875)$	$3/2^-$	1875	250	0	3
$N(1900)$	$3/2^+$	1900	200	0	3
$N(2190)$	$7/2^-$	2190	500	0	3
$N(2300)$	$1/2^+$	2300	340	0	3
$N(2570)$	$5/2^-$	2570	250	0	3
Total # of free parameters				40	106

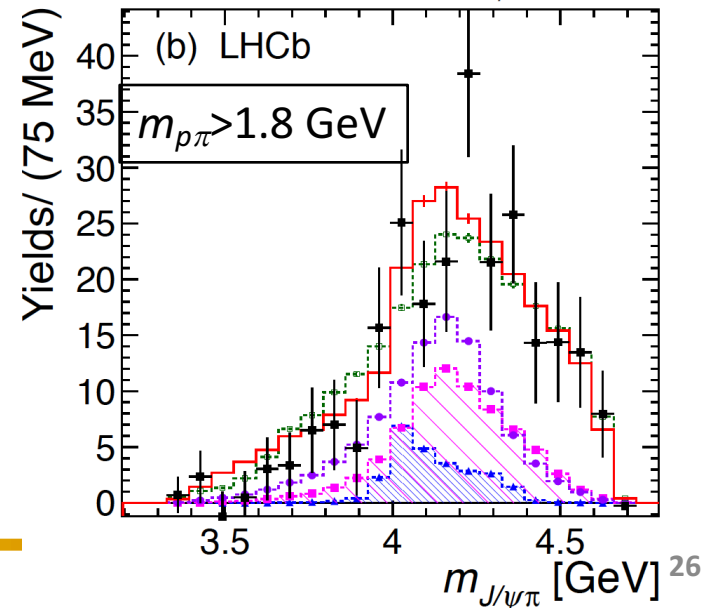
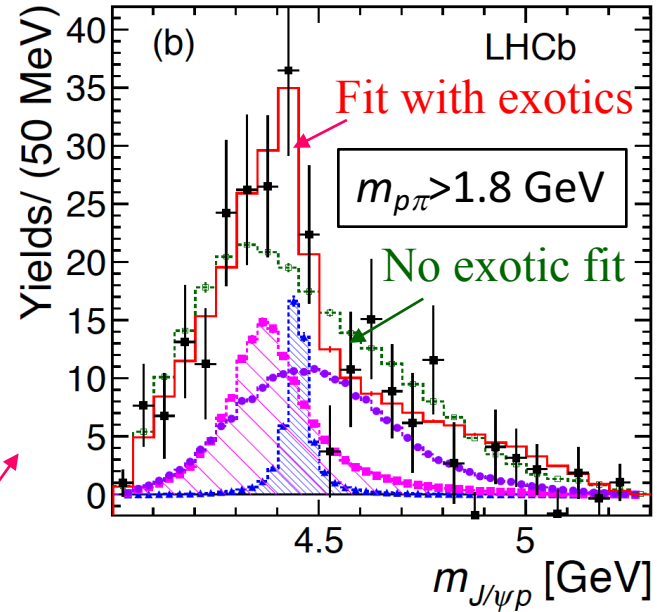
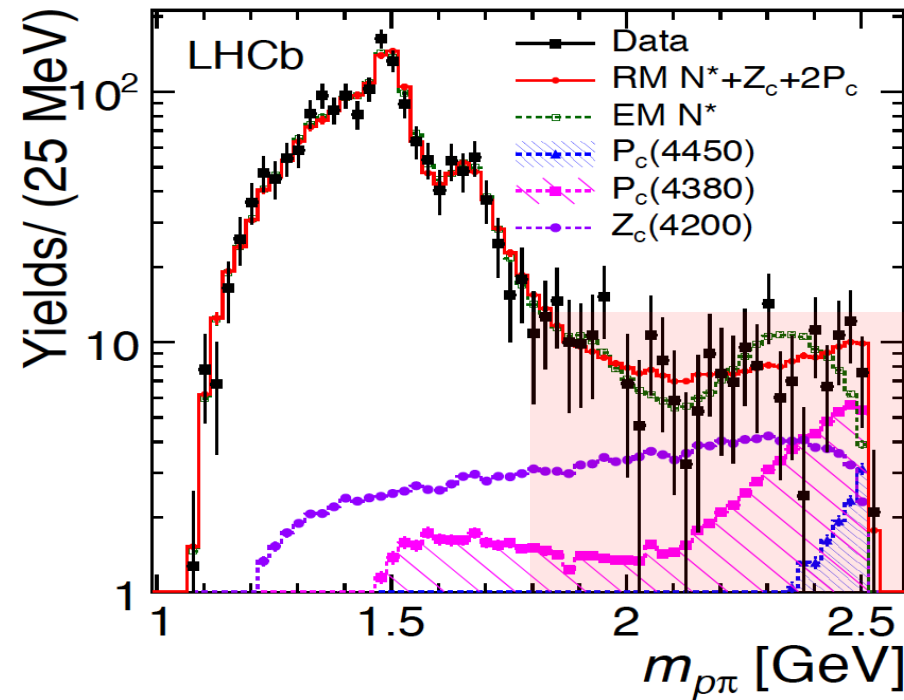
- Reduced model for central values
- Extended for significance and systematics
- Almost as many free parameters as in $\Lambda_b^0 \rightarrow J/\psi p K^-$
- Fixed m_0 and Γ_0 for the N^* and exotic states
- $Z_c(4200)$: 10 free parameters
- Each P_c : 4 free parameters + 8 fixed to that from $\Lambda_b^0 \rightarrow J/\psi p K^-$

Full amplitude fits to $\Lambda_b^0 \rightarrow J/\psi p \pi^-$

- Significance of $P_c(4380)^+$, $P_c(4450)^+$, $Z_c(4200)^-$ taken together is 3.1σ including syst.
- Evidence for exotic hadron contributions to $\Lambda_b^0 \rightarrow J/\psi p \pi^-$!

LHCb-PAPER-2016-015

PRL 117, 082002 (2016)



Further results for $\Lambda_b^0 \rightarrow J/\psi p \pi^-$



- Individual exotic hadron contributions are not significant if others are present. **More data are needed**
- Significance of two P_c is 3.3σ , **if assume production of $Z_c(4200)^-$ is negligible**. **No independent confirmation of the P_c^+ states**
- The data are consistent with the P_c states production rates expected from the $J/\psi p K^-$ measurement and Cabibbo suppression

LHCb-PAPER-2016-015

PRL **117**, 082002 (2016)

State	Fit fraction (%)	BR($\Lambda_b \rightarrow P_c^+ \pi^-$)/BR($\Lambda_b \rightarrow P_c^+ K^-$)
$Z_c(4200)^-$	$7.7 \pm 2.8_{-4.0}^{+3.4}$	---
$P_c(4380)^+$	$5.1 \pm 1.5_{-1.6}^{+2.6}$	$0.050 \pm 0.016_{-0.016}^{+0.026} \pm 0.025$
$P_c(4450)^+$	$1.6 \quad \begin{matrix} +0.8 & +0.6 \\ -0.6 & -0.5 \end{matrix}$	$0.033 \quad \begin{matrix} +0.016 & +0.011 \\ -0.014 & -0.010 \end{matrix} \pm 0.009$

0.07~0.08

Expectation assuming the second diagram negligible:

[Cheng and Chua, PRD **92**, 096009 (2015) arXiv:1509.03708]

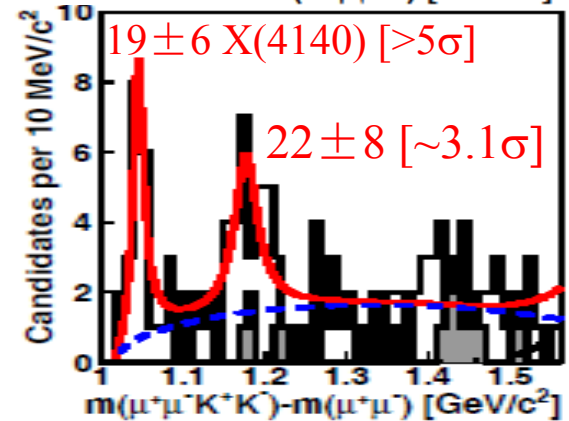
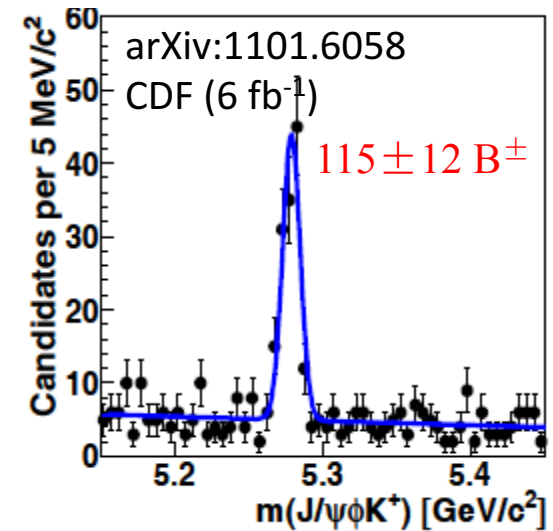


**$X \rightarrow J/\psi\phi$ in $B^+ \rightarrow J/\psi\phi K^+$
decays**

$X(4140)$ and $X(4274)$



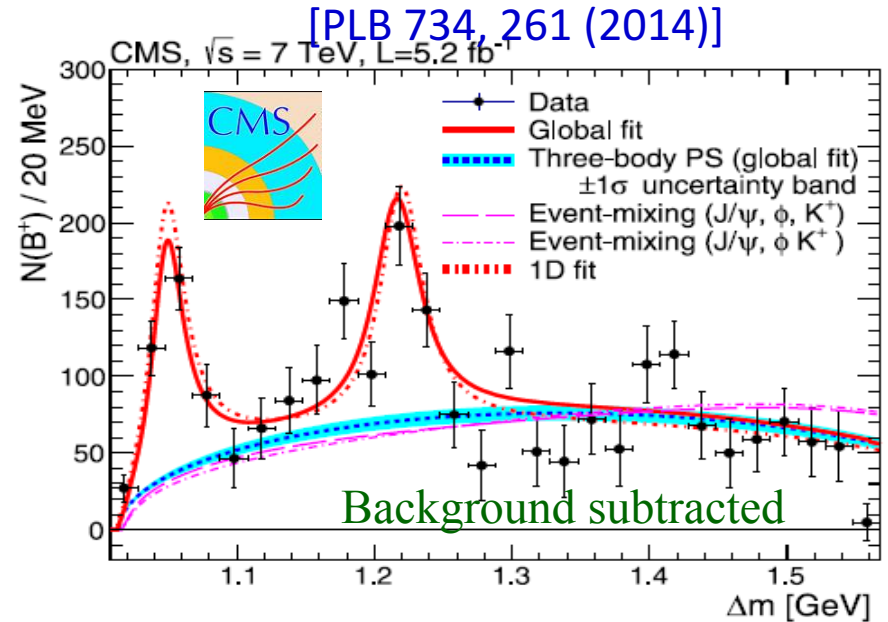
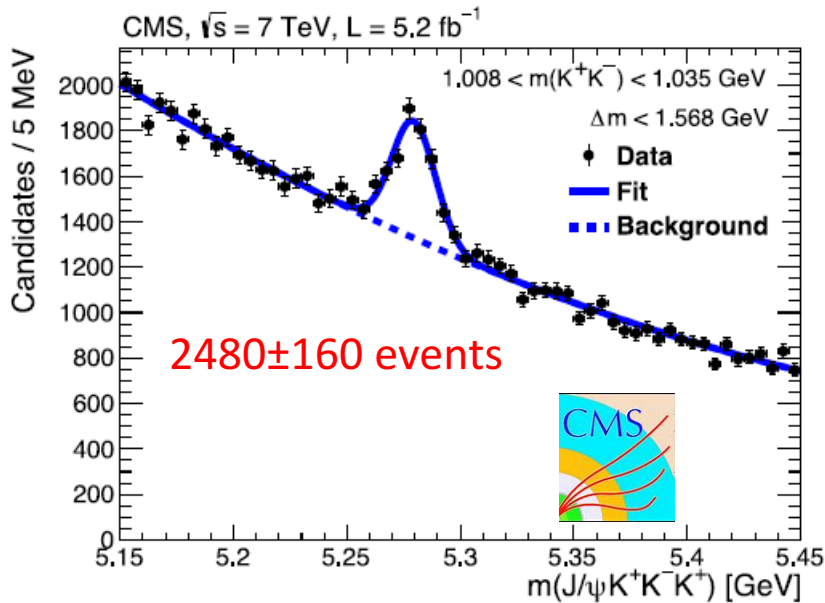
- CDF observed a narrow ($J/\psi\phi$) structure in $B^+ \rightarrow J/\psi\phi K^+$ decays [Initial publication on 2.7 fb^{-1} PRL102, 242002 (2009)]
 - $M = 4143.4 \pm 3.0 \pm 0.6 \text{ MeV}$
 - $\Gamma = 15.3_{-6.1}^{+10.4} \pm 2.5 \text{ MeV}$
 - Necessarily exotic since it is narrow and above the $D_s^+ D_s^-$ threshold
 - $[c\bar{s}c\bar{s}]$ tetraquark ?
 - Hint of a second structure: $X(4274)$
- Not confirmed by B-factories and LHCb with 0.37 fb^{-1} data



X(4140) and X(4274) from CMS



- Crucial to check by different experiments with larger statistics.

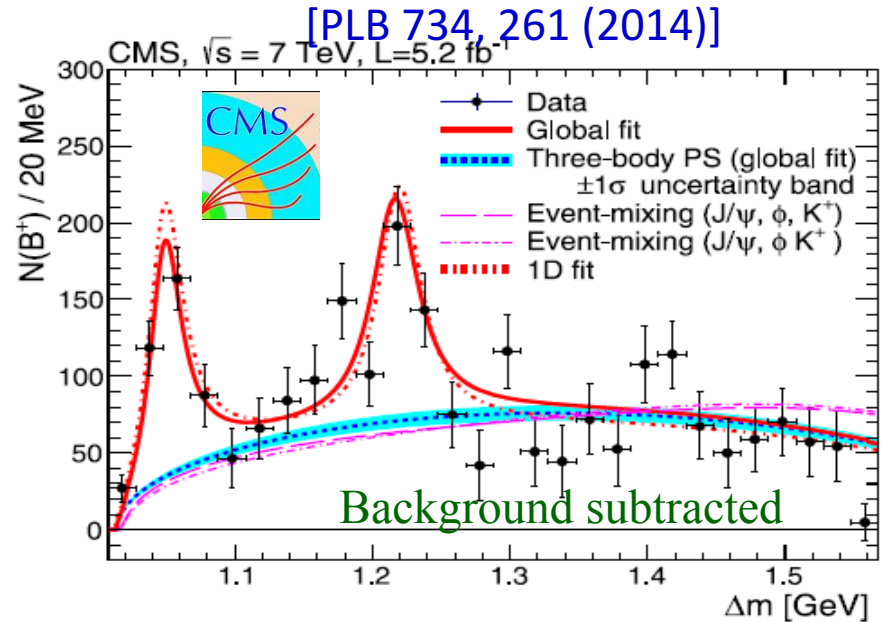
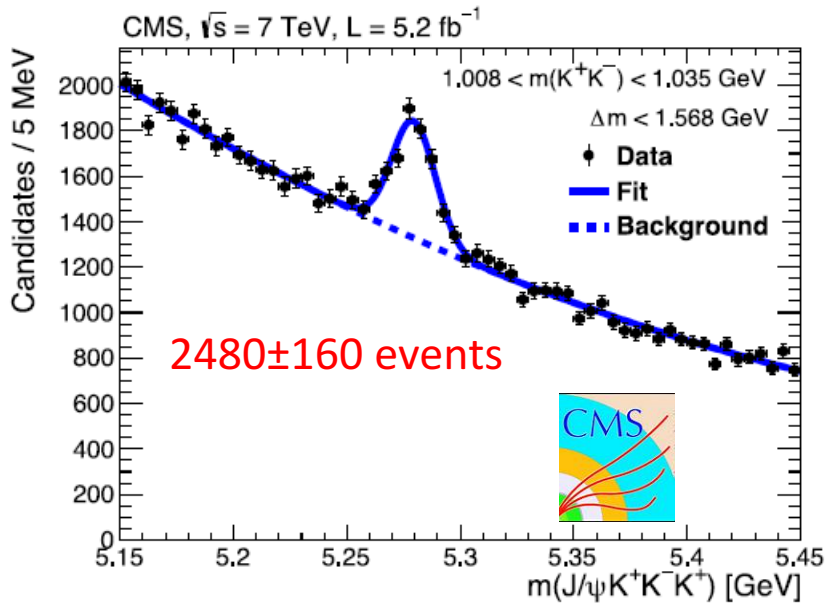


X(4140)	CDF [arXiv:1101.6058]	CMS [PLB 734, 261 (2014)]	DØ [PRD 89, 012004 (2014)]
Significance	$>5\sigma$	$>5\sigma$	3.1σ
M_0 (MeV)	$4143.4 \pm 3.0 \pm 0.6$	$4148.0 \pm 2.4 \pm 6.3$	$4159.0 \pm 4.3 \pm 6.6$
Γ_0 (MeV)	$15.3^{+10.4}_{-6.1} \pm 2.5$	$28^{+15}_{-11} \pm 19$	$19.9 \pm 12.6^{+1}_{-8}$

X(4140) and X(4274) from CMS



- Crucial to check by different experiments with larger statistics.

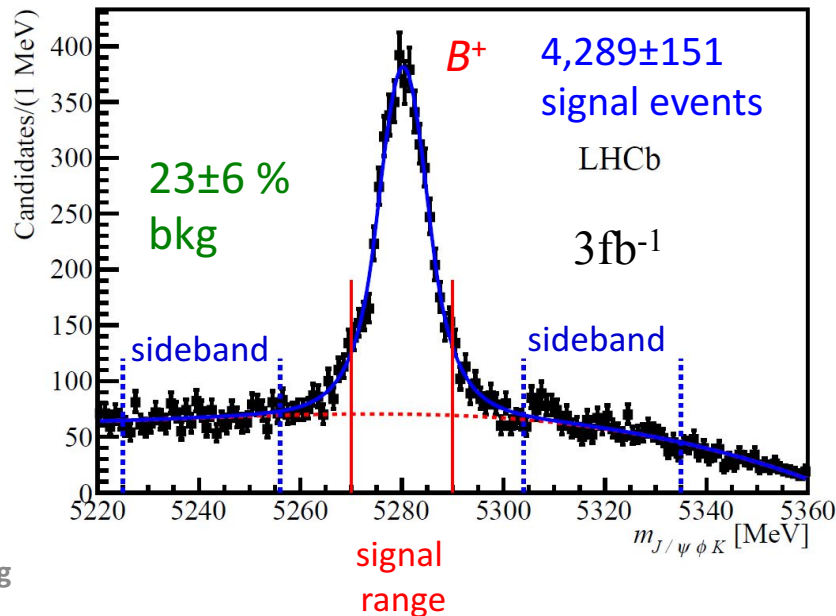
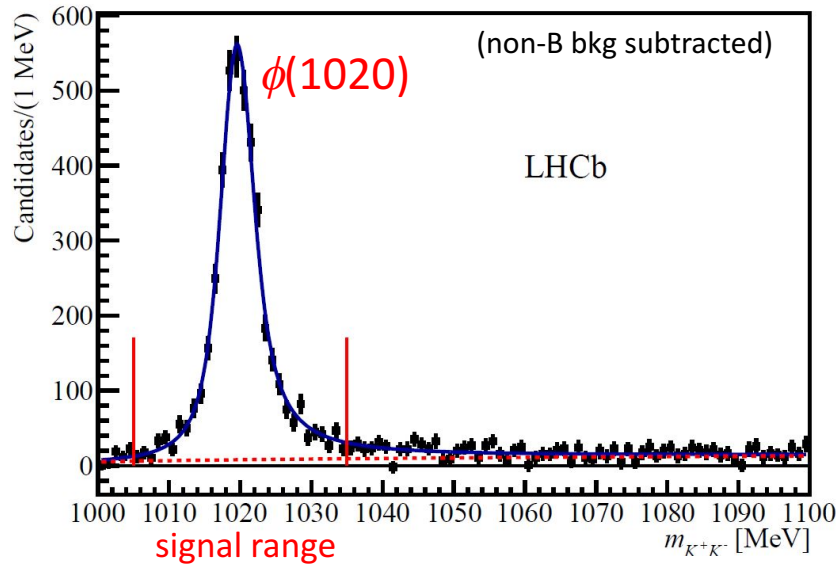


X(4274-4351)?	CDF [arXiv:1101.6058]	CMS [PLB 734, 261 (2014)]	DØ [PRD 89, 012004 (2014)]
Significance	3.1σ	>3σ	
M_0 (MeV)	$4274.4^{+8.4}_{-6.7} \pm 1.9$	$4313.8 \pm 5.3 \pm 7.3$	4328.5 ± 12.0
Γ_0 (MeV)	$32.3^{+21.9}_{-15.3} \pm 7.6$	$28^{+15}_{-11} \pm 19$	

LHCb $B^+ \rightarrow J/\psi\phi K^+$ data samples

LHCb-PAPER-2016-018 arXiv: 1606.07895

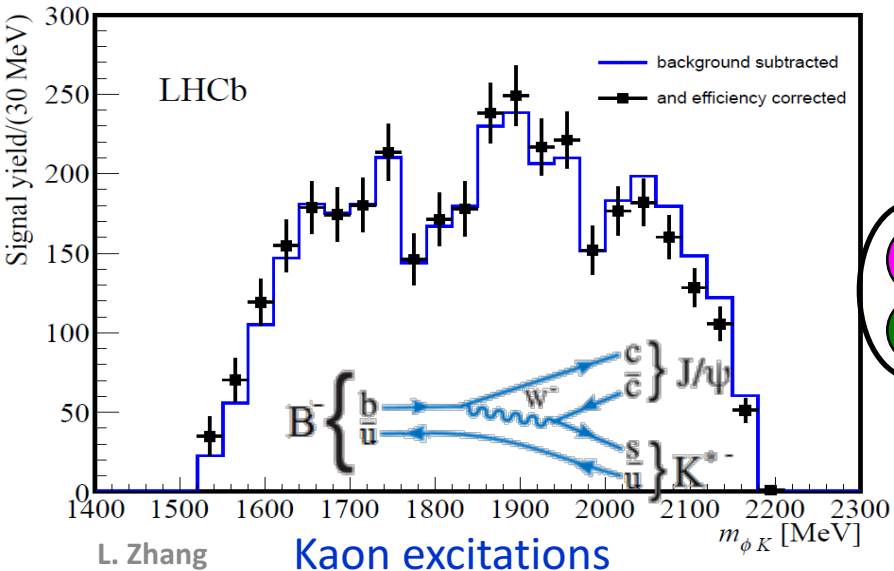
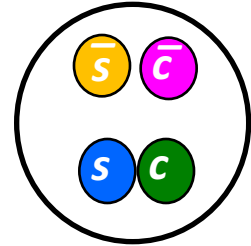
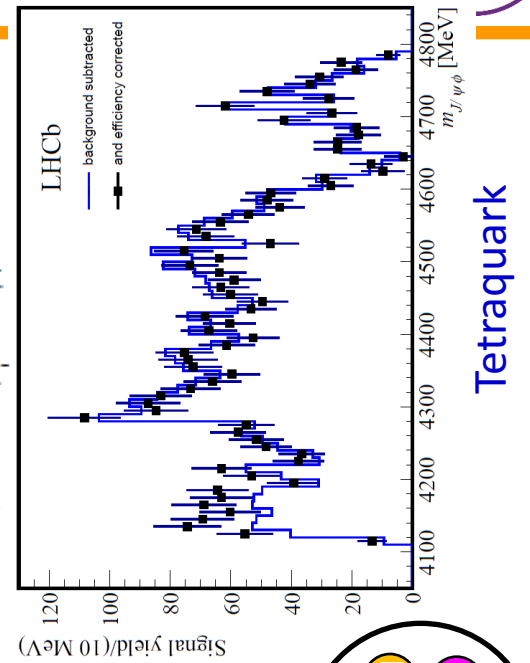
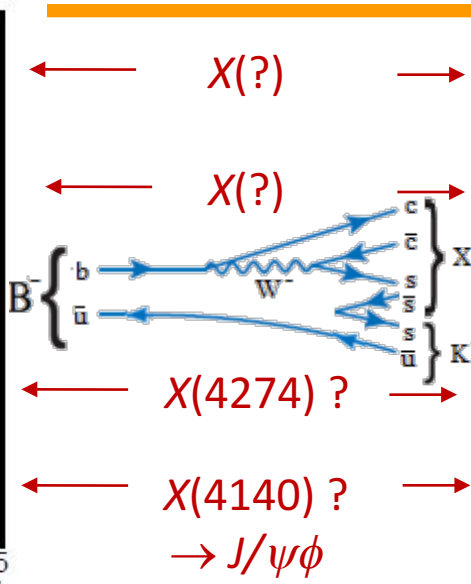
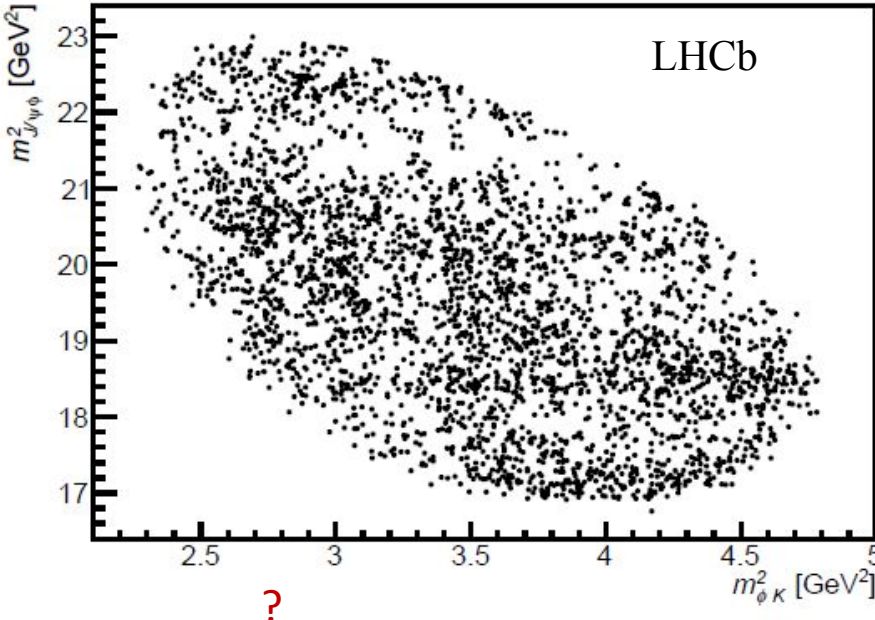
LHCb-PAPER-2016-019 arXiv: 1606.07898



- Statistically, the most powerful $B^+ \rightarrow J/\psi\phi K^+$ sample analyzed so far

Use sidebands to subtract background

$B^+ \rightarrow J/\psi \phi K^+$



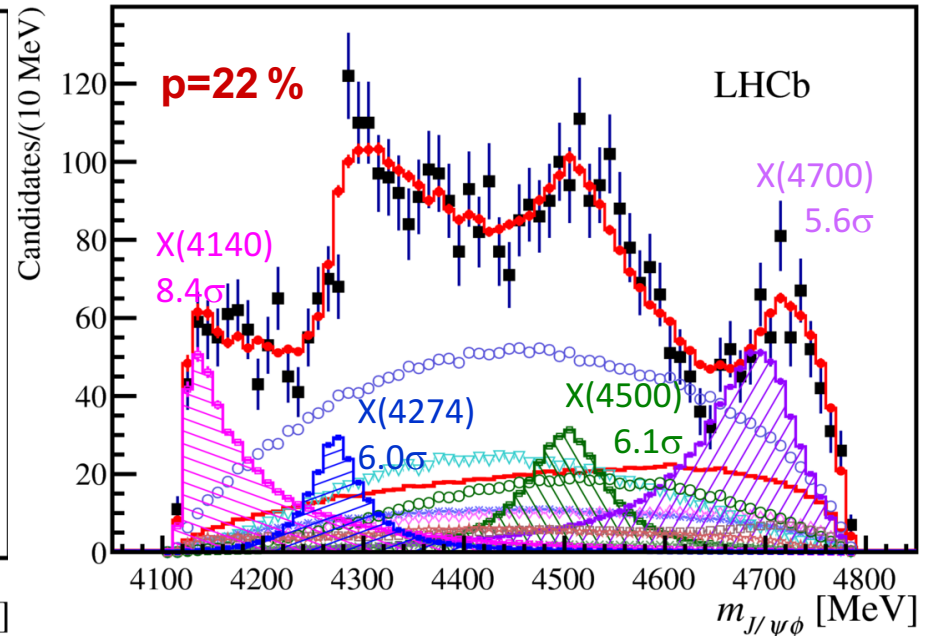
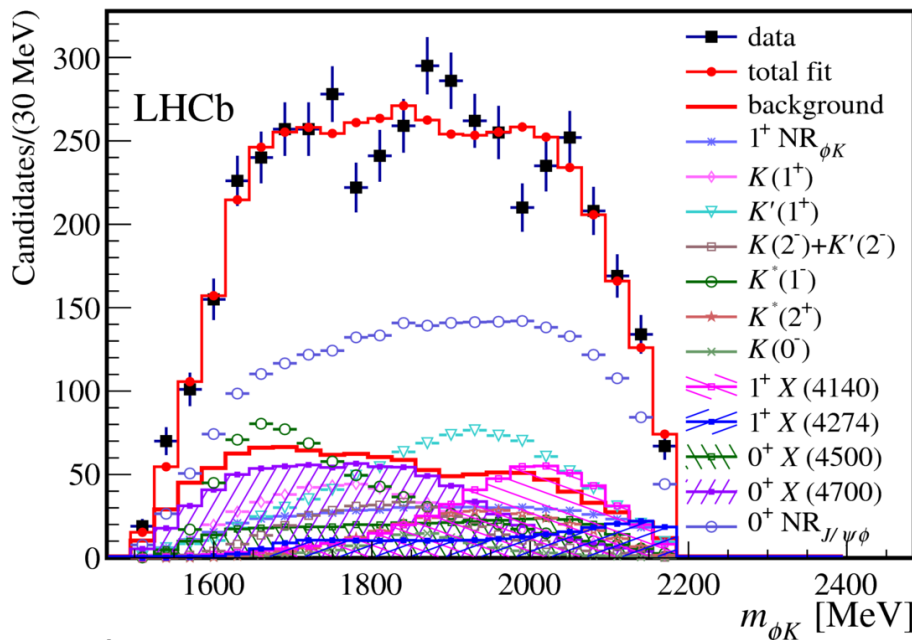
Is it a reflection of interfering $K^{*-} \rightarrow \phi K^-$?

Proper amplitude analysis necessary to check

Amplitude fit including 4 exotic X

LHCb-PAPER-2016-019 arXiv: 1606.07898

- Four X states + NR $J/\psi\phi$ give very significant improvements over the models with K^* s alone
- Default model also includes NR $\phi K + 7 K^*$ (float M_0 and Γ_0) that are significant
- These results add significantly to the knowledge of K spectroscopy (results in the paper and backup)





- J^{PC} are useful for interpretations of the states
- $X(4140)$ & $X(4274)$: identified as $J^{PC} = 1^{++}$ at $> 5\sigma$
- $X(4500)$ & $X(4700)$: $J^{PC} = 0^{++}$ at $> 4\sigma$

Contribution	sign. or Ref.	M_0 [MeV]	Γ_0 [MeV]	Fit results	
				FF %	
All $X(1^+)$				16 ± 3	${}^{+6}_{-2}$
$X(4140)$	8.4σ	4146.5 ± 4.5	83 ± 21	13 ± 3.2	${}^{+4.8}_{-2.0}$
Average other experiments		4143.4 ± 1.9	15.7 ± 6.3		
$X(4274)$	6.0σ	4273.3 ± 8.3	56 ± 11	7.1 ± 2.5	${}^{+3.5}_{-2.4}$
CDF	[28]	4274.4	32		
CMS	[25]	4313.8 ± 5.3	38		
All $X(0^+)$				28 ± 5	${}^{+7}_{-7}$
$\text{NR}_{J/\psi\phi}$	6.4σ			46 ± 11	${}^{+11}_{-21}$
$X(4500)$	6.1σ	4506 ± 11	92 ± 21	6.6 ± 2.4	${}^{+3.5}_{-2.3}$
$X(4700)$	5.6σ	4704 ± 10	120 ± 31	12 ± 5	${}^{+9}_{-5}$

substantially larger



Search for $B_S^0 \pi^\pm$ state

A new $B_S^0 \pi^\pm$ state claimed by DØ



[DØ: arXiv:1602.07588, PRL 117, 022003 (2016)]

Claimed evidence of an exotic state

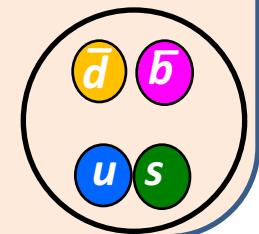
◆ $X(5568)^\pm \rightarrow B_S^0 \pi^\pm, B_S^0 \rightarrow J/\psi \phi, J/\psi \rightarrow \mu^+ \mu^-, \phi \rightarrow K^+ K^-$

$$M = 5567.8 \pm 2.9_{-1.9}^{+0.9} \text{ MeV}$$

$$\Gamma = 21.9 \pm 6.4_{-2.5}^{+5.0} \text{ MeV}$$

◆ Fraction of B_S^0 from X^\pm decay: $\rho_X^{\text{DØ}} = (8.6 \pm 1.9 \pm 1.4) \%$

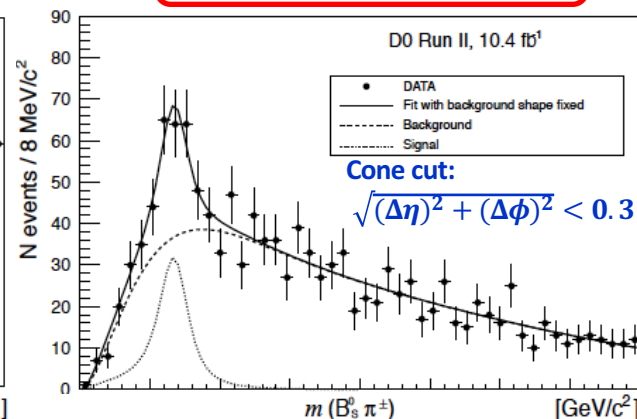
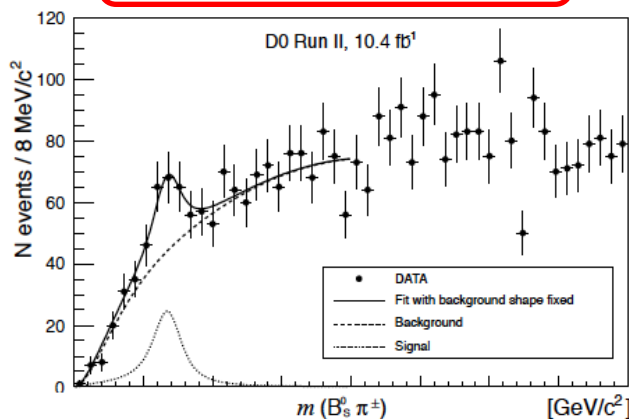
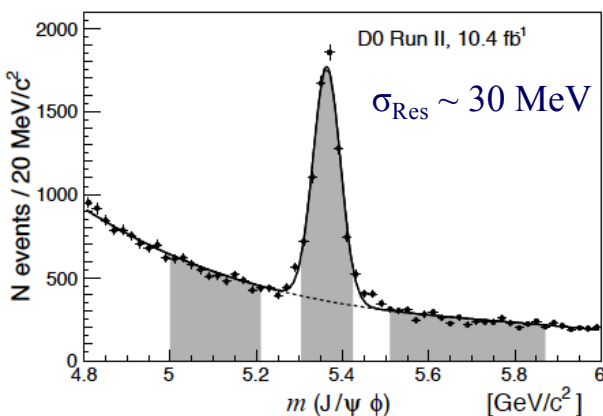
◆ If confirmed, would be unique with 4 different flavours



$$N(B_S) \sim 5500$$

$$N(X) = 106 \pm 23$$

$$N(X) = 133 \pm 31$$



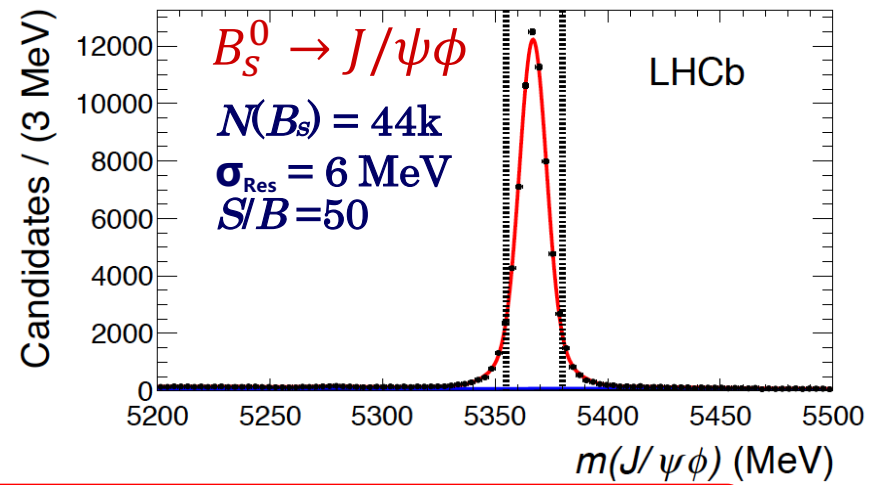
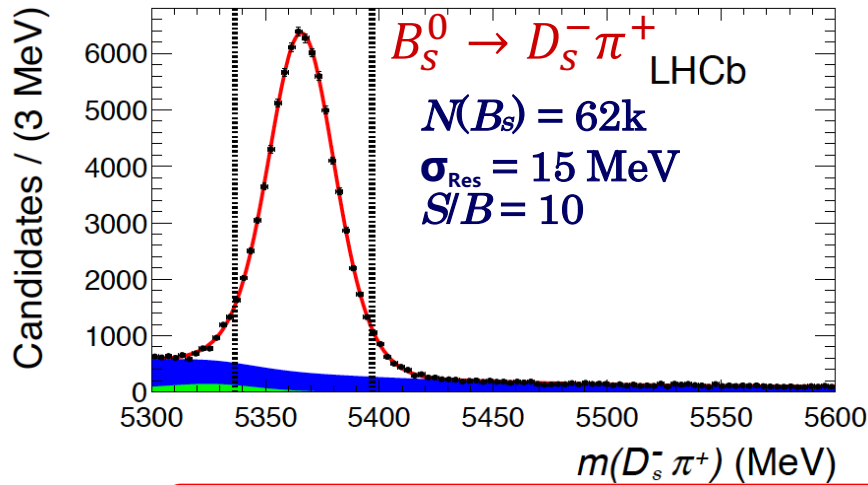
3.9σ

5.1σ

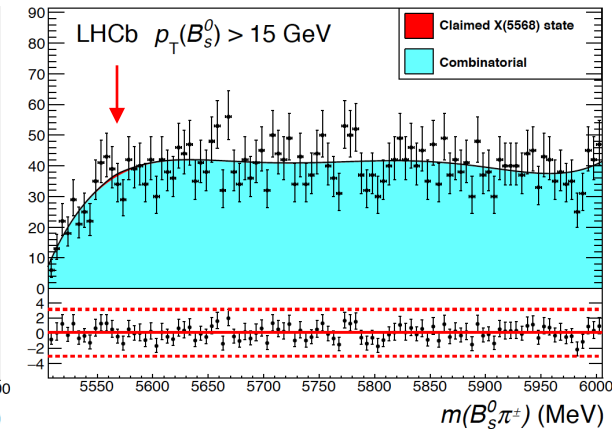
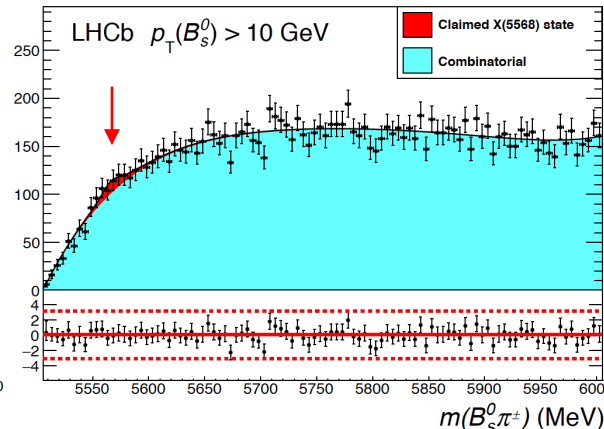
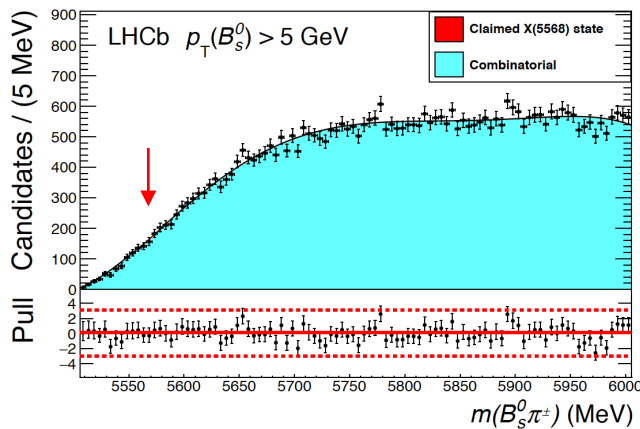
Signal significance

Samples from LHCb

LHCb-PAPER-2016-029
PRL 117, 152003 (2016)



B_S^0 sample 20x larger and much cleaner than DØ



No evident X(5568) in $B_S^0 \pi^\pm$ sample for 3 different $p_T(B_S^0)$ cuts

Upper limits



- At 90% (95%) CL

$$\rho_X^{\text{LHCb}} = \frac{\sigma(pp \rightarrow X + \text{anything}; X \rightarrow B_s^0 \pi^{\pm})}{\sigma(pp \rightarrow B_s^0 + \text{anything})}$$

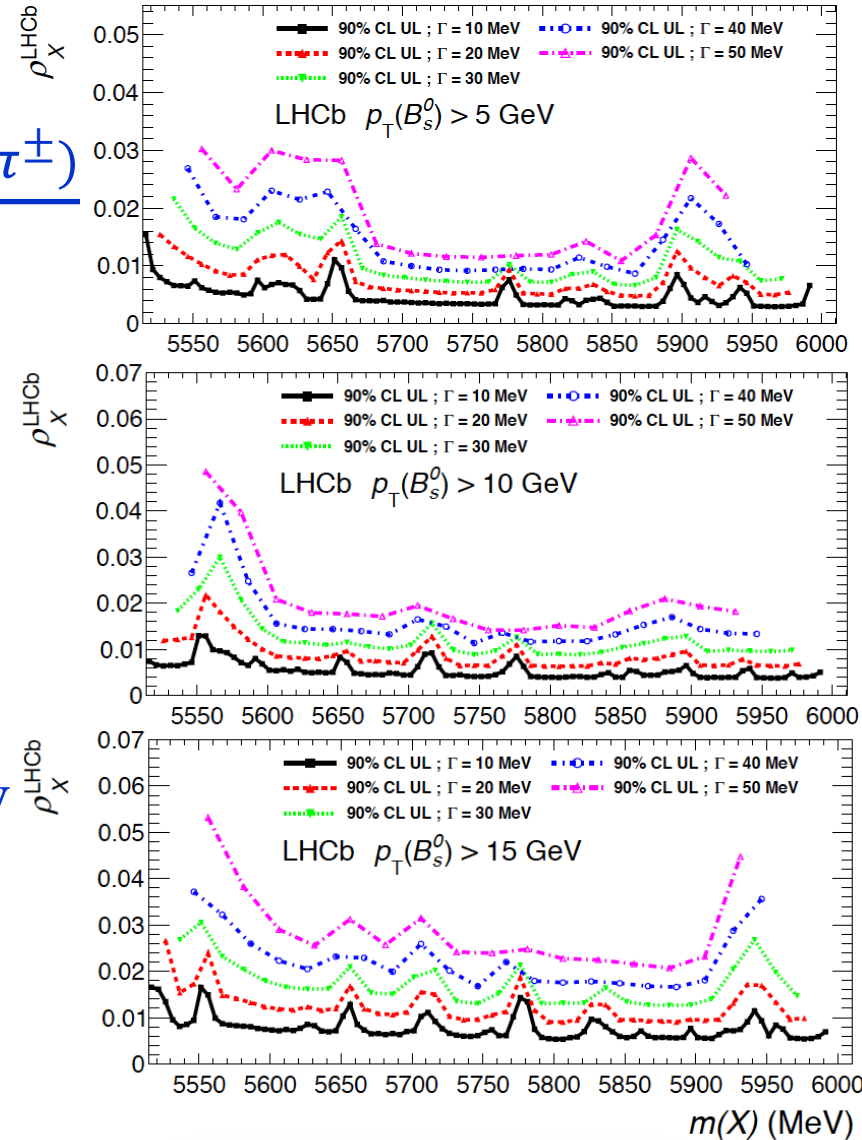
in LHCb acceptance

$$\rho_X^{\text{LHCb}}(p_T(B_s^0) > 5 \text{ GeV}) < 0.011 \text{ (0.012)}$$

$$\rho_X^{\text{LHCb}}(p_T(B_s^0) > 10 \text{ GeV}) < 0.021 \text{ (0.024)}$$

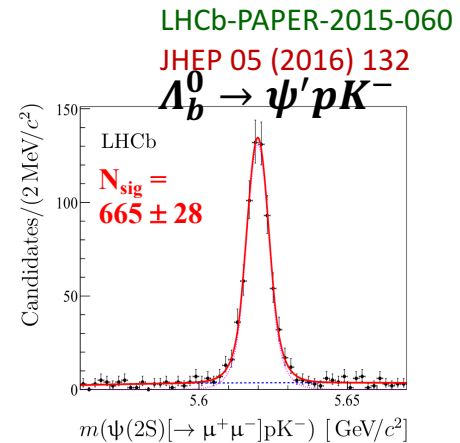
$$\rho_X^{\text{LHCb}}(p_T(B_s^0) > 15 \text{ GeV}) < 0.018 \text{ (0.020)}$$

- No significant $B_s^0 \pi^{\pm}$ states for any mass and width below 6 GeV
- Upper limit is set as a function of $m(X)$ and $\Gamma(X)$



Summary

- No significant signal for a tetraquark of $B_s^0 \pi^\pm$ for $m < 6$ GeV
- $Z_c(4430)$ state is confirmed, and two pentaquark states $P_c(4380)^+$ and $P_c(4450)^+$ are observed
- Using amplitude analysis, 3.1σ evidence is seen for the exotic hadron contributions in $\Lambda_b^0 \rightarrow J/\psi p \pi^-$ (more data needed)
- Four $X \rightarrow J/\psi \phi$ observed in the 1st full amplitude analysis of $B^+ \rightarrow J/\psi \phi K^+$
- More studied can be made with Run-II data, in similar channels where J/ψ replaced by ψ', χ_{c1}, η_c etc.

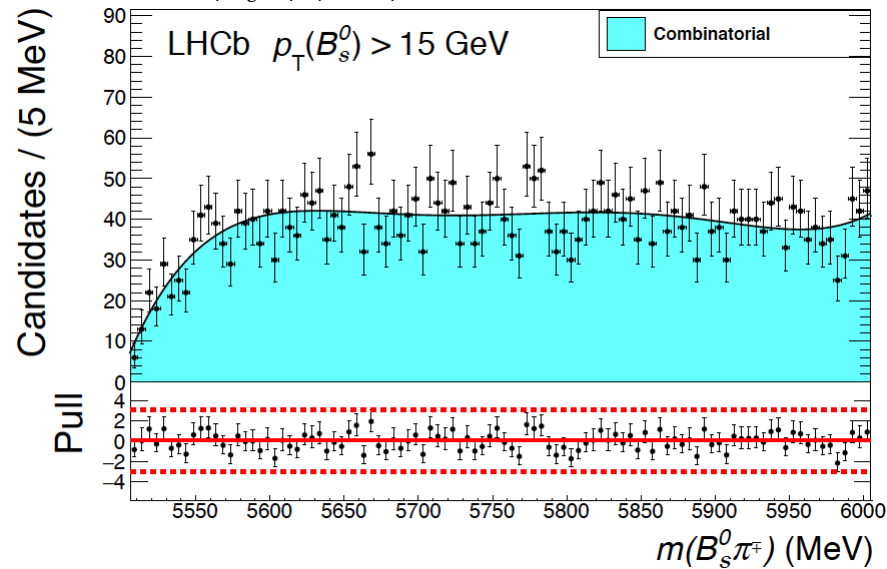
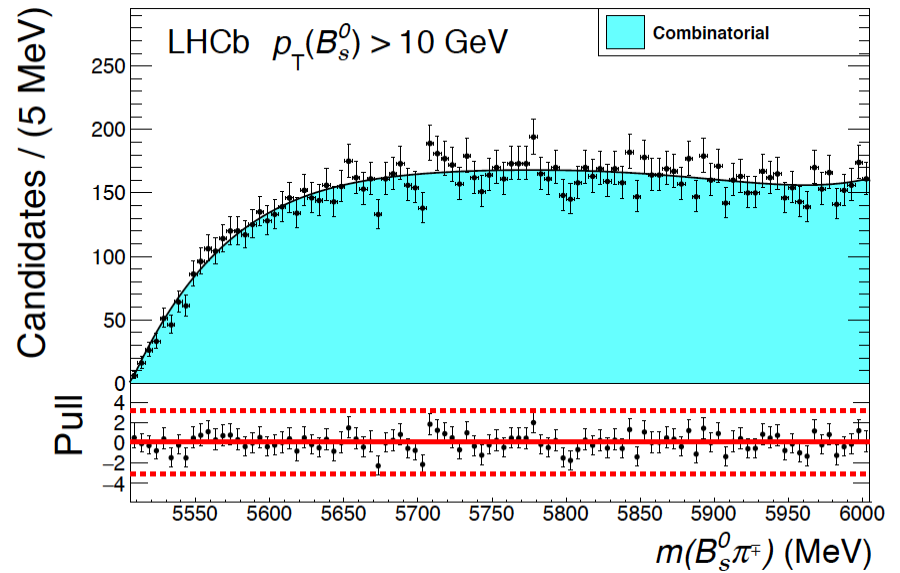
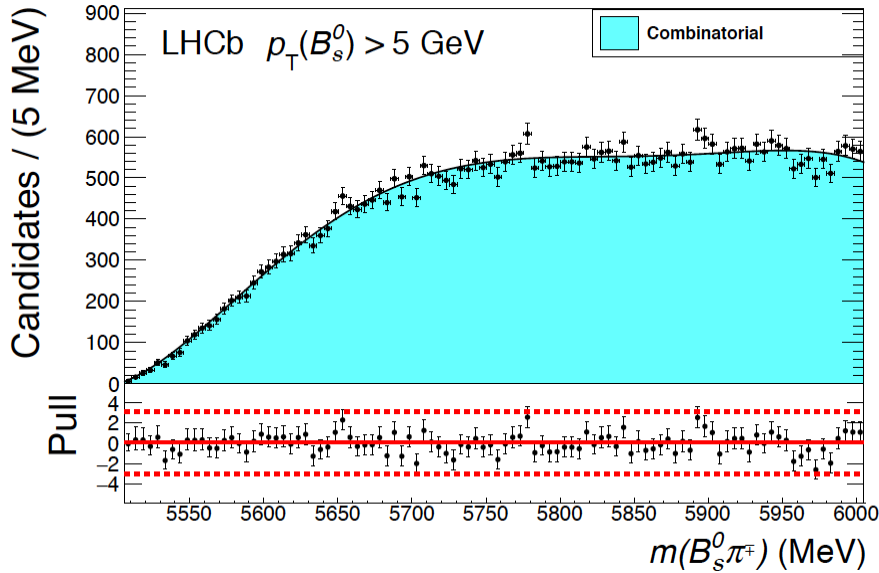




BACKUP



Search for $B_S^0\pi^\pm$ state: fits w/o signal



LHCb-PAPER-2016-029
arXiv:1608.00435

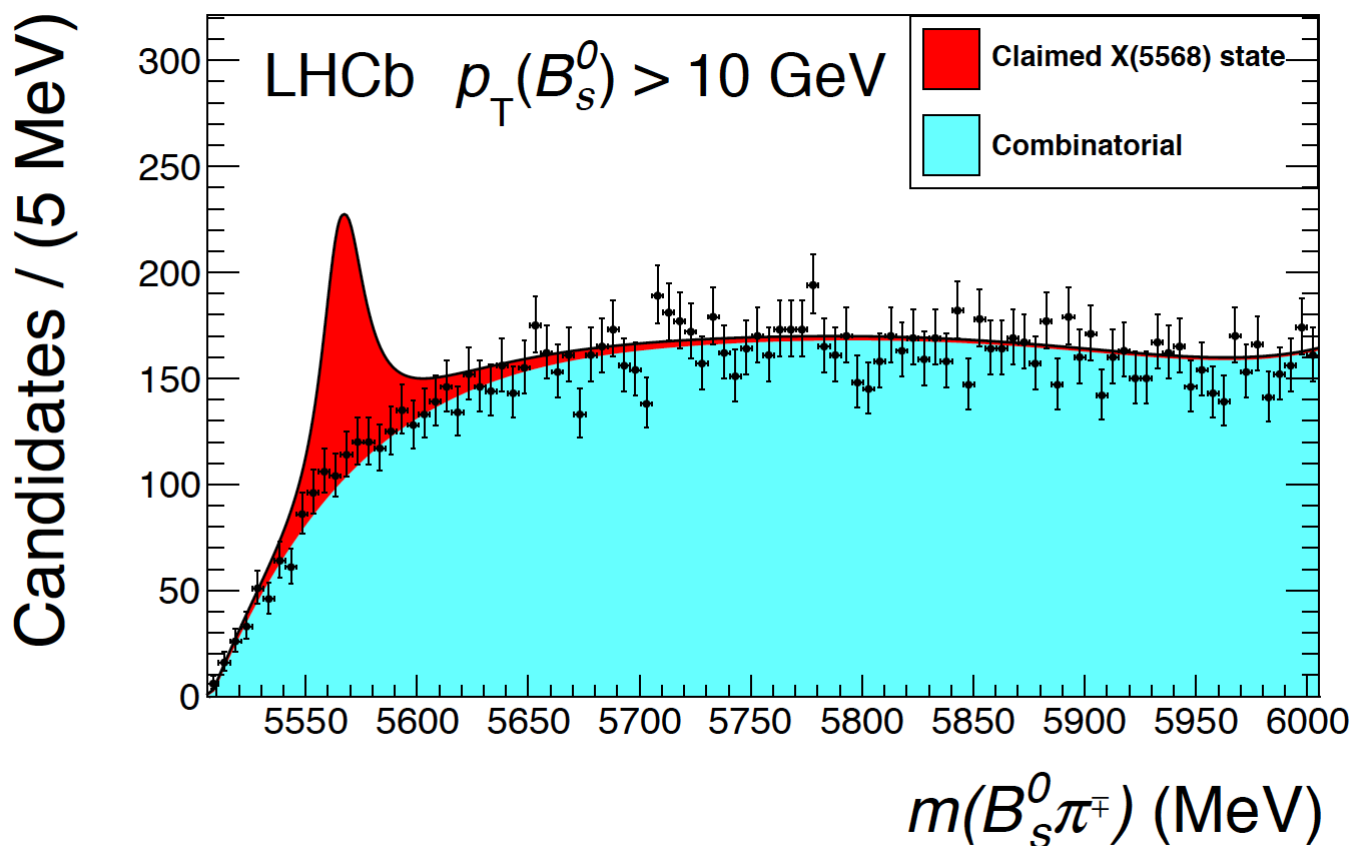
Search for $B_S^0 \pi^\pm$ state

LHCb-PAPER-2016-029

arXiv:1608.00435



- Superimpose signal component assuming $\rho_X^{\text{LHCb}} = \rho_X^{\text{D}\phi} = 8.6\%$

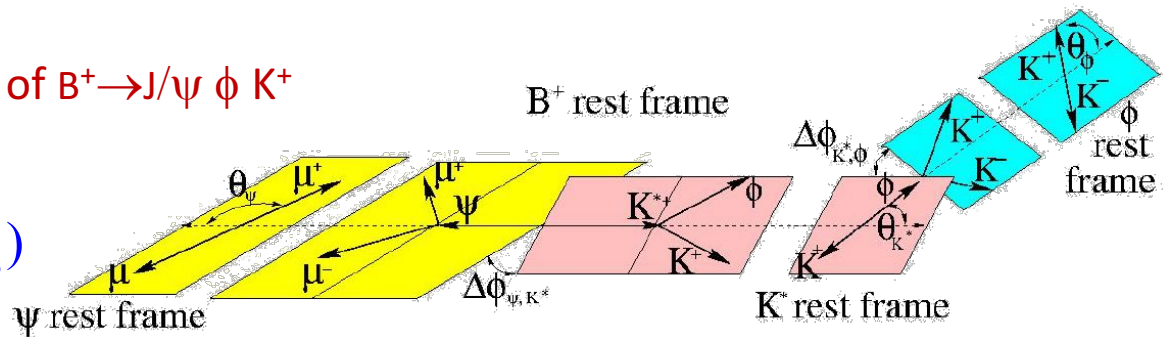


Amplitude analysis needed

- All previous analyses performed naïve 1D mass fits to $m_{J/\psi\phi}$
 - Ad hoc assumptions about kaon contributions (e.g. 3-body phase-space distribution, incoherent)
 - No sensitivity to J^{PC} of X structures

Perform first amplitude analysis of $B^+ \rightarrow J/\psi \phi K^+$

$$\Omega \equiv (\theta_{K^*}, \theta_{\psi}, \Delta\phi_{\psi, K^*}, \theta_{\phi}, \Delta\phi_{K^*, \phi})$$



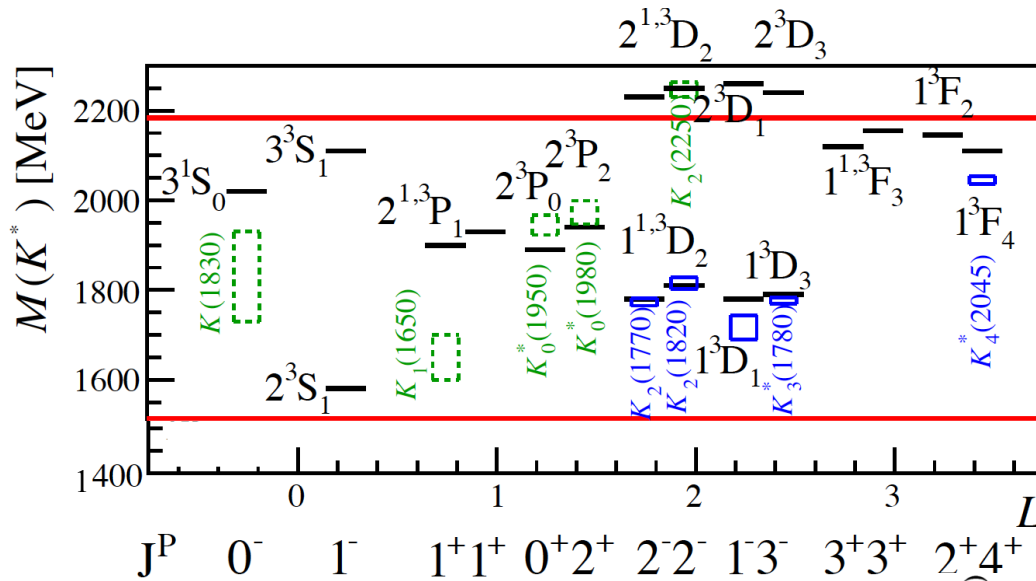
6D maximum likelihood fit

1-4 independent **complex** helicity couplings

H_R per K^* resonance

$$\left| \text{MatrixEle}(m_{K\phi}, \Omega | J_R^P, M_R, \Gamma_R, H_R) \right|^2$$

Amplitude fits with kaon excitations only



Mass range in $B^+ \rightarrow J/\psi \phi K^+$

Guidance from quark model prediction

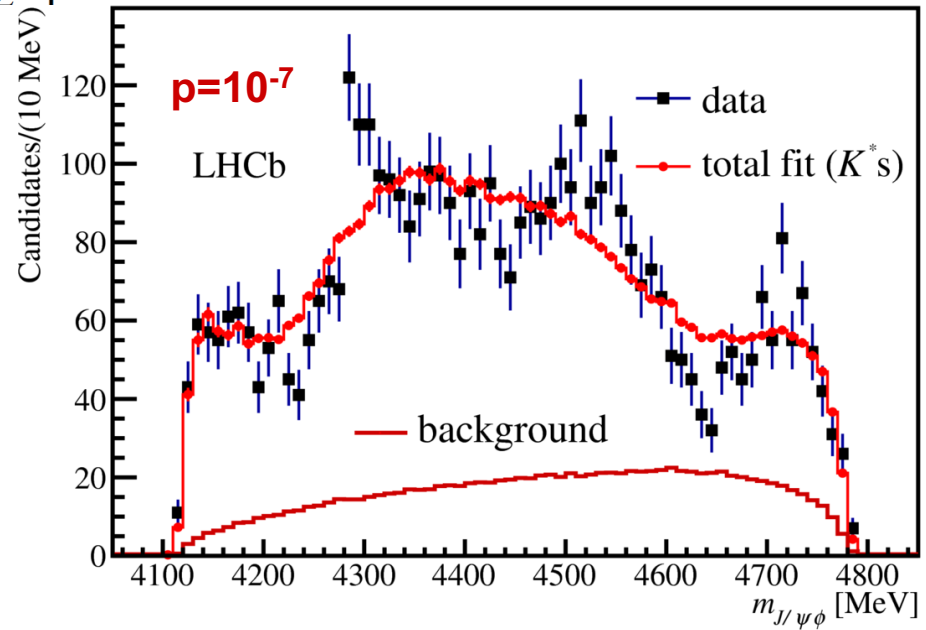
Godfrey-Isgur, PRD 32, 189 (1985)

Established

Unconfirmed

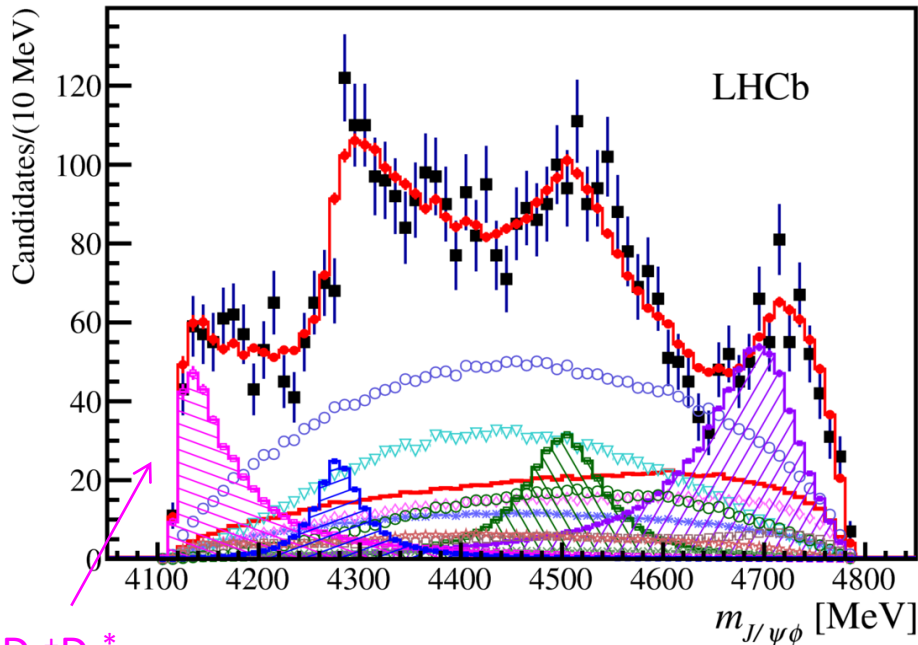
LHCb-PAPER-2016-019 arXiv: 1606.07898

- M_0 and Γ_0 of K^* s are kept float in the fits
- Data can not be described by K^* only
 - Example fit: $12 K^* + NR \phi K$



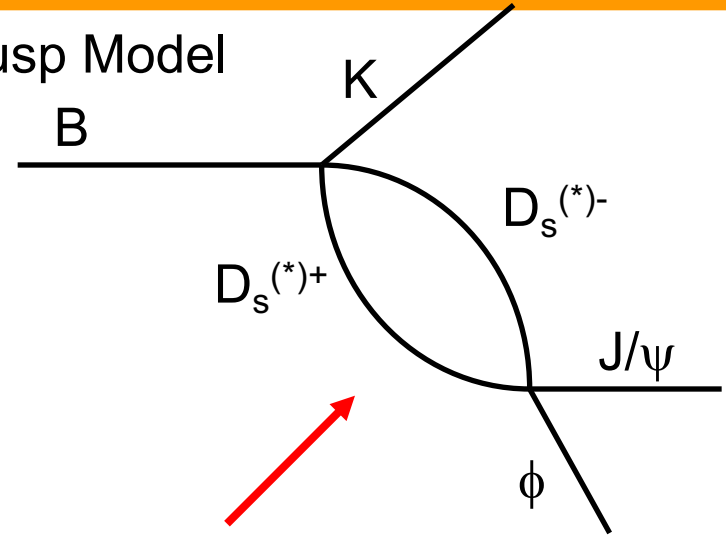
Is X(4140) a $D_s^+ D_s^{*-}$ cusp ?

LHCb-PAPER-2016-019 arXiv: 1606.07898



$D_s^+ D_s^{*-}$ cusp

Cusp Model



Cusp model by
E. S. Swanson, arXiv:1504.07952
(see also PRD91, 034009 (2015))

- The cusp is preferred by 1.6σ over the Breit-Wigner amplitude for X(4140) from the fit likelihood ratio
- No success in describing any other J/ $\psi\phi$ mass structures as a cusp

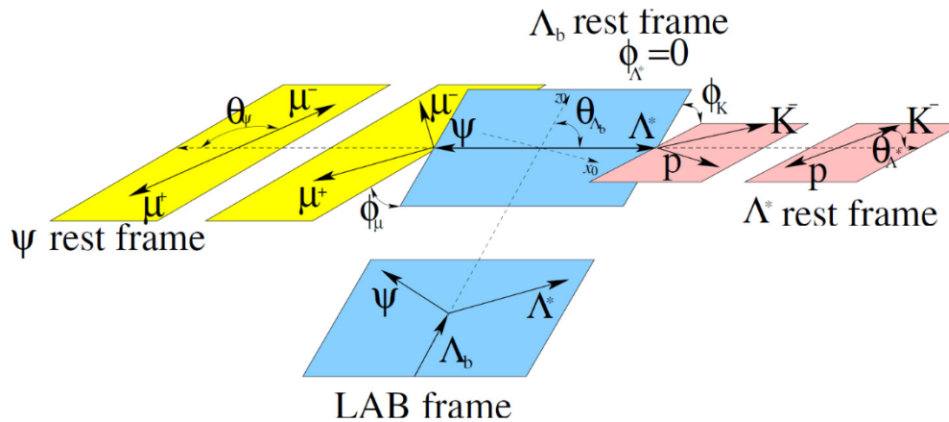
Amplitude Analysis Formalism

- Helicity formalism
 - Allows for the conventional $\Lambda^* \rightarrow pK$ resonances to interfere with pentaquark states $P_c^+ \rightarrow J/\psi p$
 - Use $m(K^-p)$ & 5 decay angles as fit parameters.

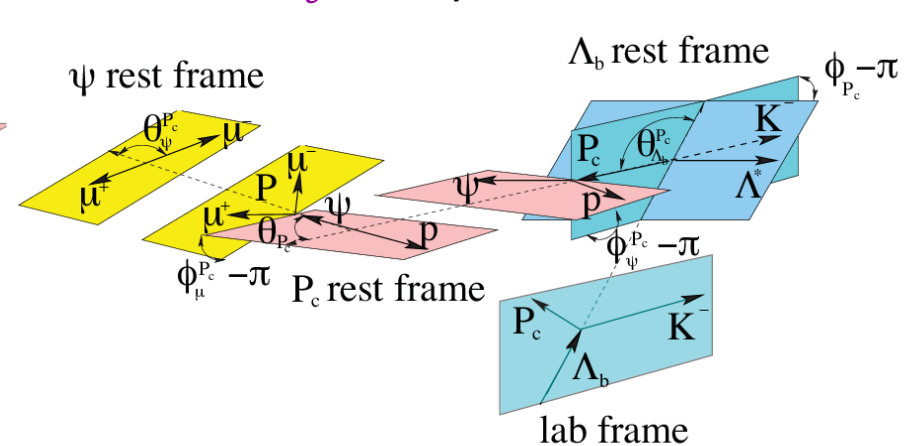
So 6D fit

$$|\mathcal{M}|^2 = \sum_{\lambda_{\Lambda_b^0}} \sum_{\lambda_p} \sum_{\Delta\lambda_\mu} \left| \mathcal{M}_{\lambda_{\Lambda_b^0}, \lambda_p, \Delta\lambda_\mu}^{\Lambda^*} + e^{i\Delta\lambda_\mu\alpha_\mu} \sum_{\lambda_p^{P_c}} d_{\lambda_p^{P_c}, \lambda_p}^{\frac{1}{2}}(\theta_p) \mathcal{M}_{\lambda_{\Lambda_b^0}, \lambda_p^{P_c}, \Delta\lambda_\mu}^{P_c} \right|^2$$

Λ^* Decay Chain



P_c^+ Decay Chain



Systematic uncertainty

LHCb-PAPER-2015-029
PRL 115, 07201 (2015)



Source	$\Lambda_b^0 \rightarrow J/\psi p K^-$				Fit fractions (%)			
	M_0 (MeV)		Γ_0 (MeV)		low	high	$\Lambda(1405)$	$\Lambda(1520)$
	low	high	low	high				
Extended vs. reduced	21	0.2	54	10	3.14	0.32	1.37	0.15
Λ^* masses & widths	7	0.7	20	4	0.58	0.37	2.49	2.45
Proton ID	2	0.3	1	2	0.27	0.14	0.20	0.05
$10 < p_p < 100$ GeV	0	1.2	1	1	0.09	0.03	0.31	0.01
Nonresonant	3	0.3	34	2	2.35	0.13	3.28	0.39
Separate sidebands	0	0	5	0	0.24	0.14	0.02	0.03
J^P ($3/2^+$, $5/2^-$) or ($5/2^+$, $3/2^-$)	10	1.2	34	10	0.76	0.44		
$d = 1.5 - 4.5$ GeV $^{-1}$	9	0.6	19	3	0.29	0.42	0.36	1.91
$L_{\Lambda_b^0}^{P_c} \Lambda_b^0 \rightarrow P_c^+ (\text{low/high}) K^-$	6	0.7	4	8	0.37	0.16		
$L_{P_c} P_c^+ (\text{low/high}) \rightarrow J/\psi p$	4	0.4	31	7	0.63	0.37		
$L_{\Lambda_b^0}^{\Lambda^*} \Lambda_b^0 \rightarrow J/\psi \Lambda^*$	11	0.3	20	2	0.81	0.53	3.34	2.31
Efficiencies	1	0.4	4	0	0.13	0.02	0.26	0.23
Change $\Lambda(1405)$ coupling	0	0	0	0	0	0	1.90	0
Overall	29	2.5	86	19	4.21	1.05	5.82	3.89
sFit/cFit cross check	5	1.0	11	3	0.46	0.01	0.45	0.13

Systematic uncertainty

LHCb-PAPER-2016-015

arXiv: 1606.06999,
to appear in PRL

$$\Lambda_b^0 \rightarrow J/\psi p \pi^-$$

Source	$P_c(4450)^+$	$P_c(4380)^+$	$Z_c(4200)^-$
N^* masses and widths	± 0.05	± 0.23	± 0.31
P_c^+ , Z_c^- masses and widths	± 0.32	± 1.27	± 1.56
Additional N^*	$+0.08$ -0.23	$+0.59$ -0.55	$+0.71$ -2.92
Inclusion of $Z_c(4430)^-$	$+0.01$	$+0.97$	$+2.87$
Exclusion of $Z_c(4200)^-$	-0.15	$+1.61$	-
Other J^P	$+0.38$ -0.00	$+0.92$ -0.28	$+0.00$ -2.16
Blatt–Weisskopf radius	± 0.11	± 0.17	± 0.21
$L_{\Lambda_b^0}^{N^*}$ in $\Lambda_b^0 \rightarrow J/\psi N^*$	± 0.07	± 0.46	± 0.04
$L_{\Lambda_b^0}^{P_c}$ in $\Lambda_b^0 \rightarrow P_c^+ \pi^-$	-0.05	-0.17	$+0.09$
$L_{\Lambda_b^0}^{Z_c}$ in $\Lambda_b^0 \rightarrow Z_c^- p$	± 0.07	± 0.22	± 0.53
K-matrix model	-0.03	$+0.11$	-0.02
P_c^+ couplings	± 0.14	± 0.31	± 0.36
Background subtraction	-0.07	-0.13	-0.39
Total	$+0.55$ -0.48	$+2.61$ -1.58	$+3.43$ -4.04

Systematic uncertainty

LHCb-PAPER-2016-019

arXiv: 1606.07898



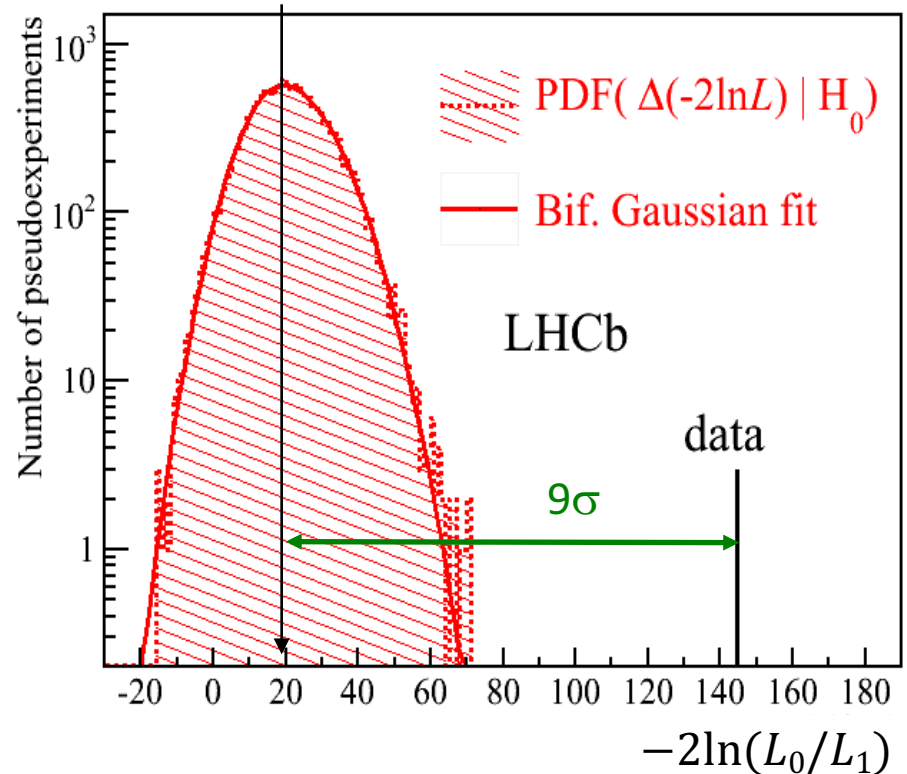
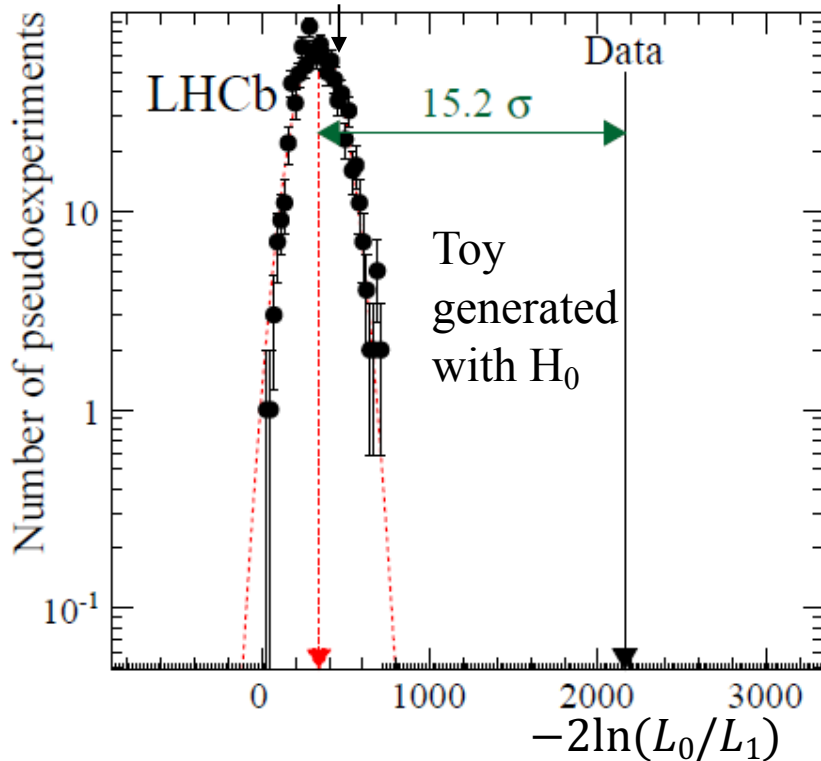
sys. var.	1 ⁺			X(4140)				X(4274)				0 ⁺			X(4500)				X(4700)			NR	
	FF	M_0	Γ_0	FF	M_0	Γ_0	FF	FF	M_0	Γ_0	FF	M_0	Γ_0	FF	M_0	Γ_0	FF	FF					
K^*	+2.0	+3.6	+17.1	+2.2	+11.2	+7.9	+1.4	+1.8	+9.3	+13.8	+2.0	+7.5	+38.6	+6.7	+8.0	-11.0	-8.6	-16.6	-1.7	-18.9	-13.5	-4.8	-16.6
model	-1.7	-2.6	-11.7	-1.9	-2.5	-8.5	-1.5	+0.3	+1.3	+10.8	+1.7	+9.0	+12.4	+1.5	+1.2	-4.7	-9.6	-11.2	-1.6	-6.8	-24.9	-0.8	-8.5
L	+3.2	+2.2	+7.3	+2.1	+10.6	+1.4	+1.0	+0.3	+1.3	+10.8	+1.7	+9.0	+12.4	+1.5	+1.2	+0.0	-1.2	-6.2	-0.5	-0.8	-4.6	-1.2	-8.5
var.	+0.0	-1.2	-6.2	-0.5	-0.8	-4.6	-1.2	-1.7	+6.3	+0.3	+0.2	+7.1	-15.7	-1.7	-9.1	-1.2	+0.0	+1.2	+0.2	+1.9	-2.5	0.5	-1.6
NR exp.	+0.4	-0.2	-0.1	+0.4	-0.2	+0.6	+0.8	+0.1	+0.8	-0.1	-0.3	+0.9	-5.8	-0.9	-1.1	+0.5	+1.7	+3.2	+0.1	-0.1	+1.7	+0.0	+1.1
X cusp	+2.2			+0.9	+6.4	-5.4	-1.4	-0.5	-2.4	-2.6	-0.2	-1.5	-3.1	-0.7	-2.5	+0.5	+3.7	+3.4	+0.4	+1.2	+7.0	+0.8	+1.6
Γ_{tot}	-0.6	+0.2	+1.5	-0.4	+3.2	+0.2	-0.3	-2.5	-4.6	-11.1	-0.5	-3.9	-6.1	-1.4	-1.4	+1.7	+0.0	+0.2	+0.2	+0.1	+0.0	+1.2	+2.7
d=1.5	-0.9	+1.1	+5.3	-0.5	+2.2	+0.8	-0.4	+4.2	-4.3	+7.1	+1.2	-9.3	+5.8	+0.7	+4.7	+1.7	+0.0	+0.2	+0.2	+0.1	+0.0	+1.2	+2.7
d=5.0	+1.1	-0.2	-2.0	+0.6	+0.2	-0.8	+0.3	+4.2	-4.3	+7.1	+1.2	-9.3	+5.8	+0.7	+4.7	+1.7	+0.0	+0.2	+0.2	+0.1	+0.0	+1.2	+2.7
Left s.	+0.1	-0.4	-2.0	+0.1	+0.4	-0.8	+0.1	+4.2	-4.3	+7.1	+1.2	-9.3	+5.8	+0.7	+4.7	+1.7	+0.0	+0.2	+0.2	+0.1	+0.0	+1.2	+2.7
Right s.	-0.3	+0.3	+2.6	-0.2	-0.6	+1.0	+0.0	+4.2	-4.3	+7.1	+1.2	-9.3	+5.8	+0.7	+4.7	+1.7	+0.0	+0.2	+0.2	+0.1	+0.0	+1.2	+2.7
β	+1.2	-0.6	-3.6	+1.2	+1.7	-0.7	+0.9	+4.2	-4.3	+7.1	+1.2	-9.3	+5.8	+0.7	+4.7	+1.7	+0.0	+0.2	+0.2	+0.1	+0.0	+1.2	+2.7
No w^{MC}	+1.6	+0.0	+0.0	+0.1	+0.0	+0.0	+1.4	+4.2	-4.3	+7.1	+1.2	-9.3	+5.8	+0.7	+4.7	+1.7	+0.0	+0.2	+0.2	+0.1	+0.0	+1.2	+2.7
ϕ window	+2.5	+1.1	+4.7	+2.4	-1.6	+1.4	+1.8	+4.2	-4.3	+7.1	+1.2	-9.3	+5.8	+0.7	+4.7	+1.7	+0.0	+0.2	+0.2	+0.1	+0.0	+1.2	+2.7
Total	+5.9	+4.6	+20.7	+4.7	+17.2	+8.4	+3.5	+6.5	+12.0	+20.8	+3.2	+13.9	+42.0	+7.2	+11.0	+6.5	+12.0	+20.8	+3.2	+13.9	+42.0	+7.2	+11.0
sys.	-2.1	-2.8	-13.5	-2.0	-3.6	-11.1	-2.4	-6.7	-14.5	-20.4	-2.3	-24.1	-33.3	-5.3	-21.0	-6.7	-14.5	-20.4	-2.3	-24.1	-33.3	-5.3	-21.0
Stat.	2.8	4.5	20.7	3.2	8.3	10.9	2.5	5.1	11.1	21.2	2.4	10.1	30.7	4.9	10.7	5.1	11.1	21.2	2.4	10.1	30.7	4.9	10.7
$p_{\text{T}}^{K^*} > 500$	-1.3	+1.6	+1.7	-2.7	+7.8	+12.2	+0.2	-9.6	-10.9	-18.6	-3.2	-4.7	-12.7	-6.6	-17.1	-9.6	-10.9	-18.6	-3.2	-4.7	-12.7	-6.6	-17.1

Log-likelihood ratio method

- Likelihood ratio for $\text{PDF}(m_{\psi'\pi/\psi\rho}|H_0)$ and $\text{PDF}(m_{\psi'\pi/\psi\rho}|H_1)$
- H_1 can well present $m_{\psi'\pi/\psi\rho}$ spectrum in data using $l_{\max}=30/31$

$B^0 \rightarrow \psi'\pi^-K^+$

$\Lambda_b \rightarrow J/\psi p K^-$

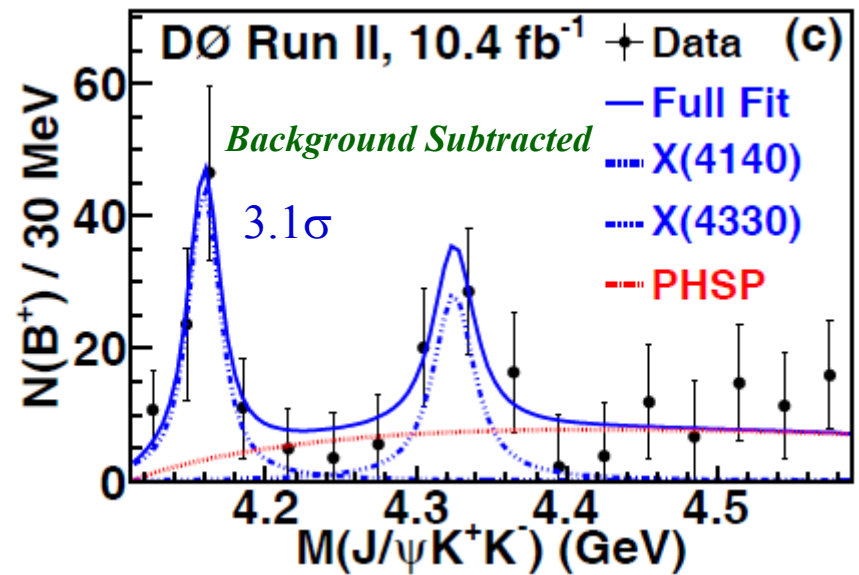
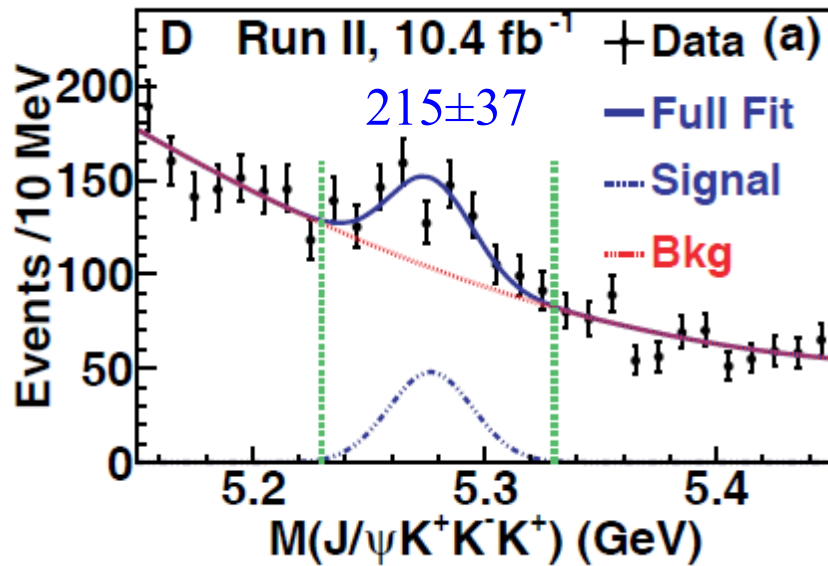


However, this approach cannot characterize exotics – amplitude analysis is still necessary

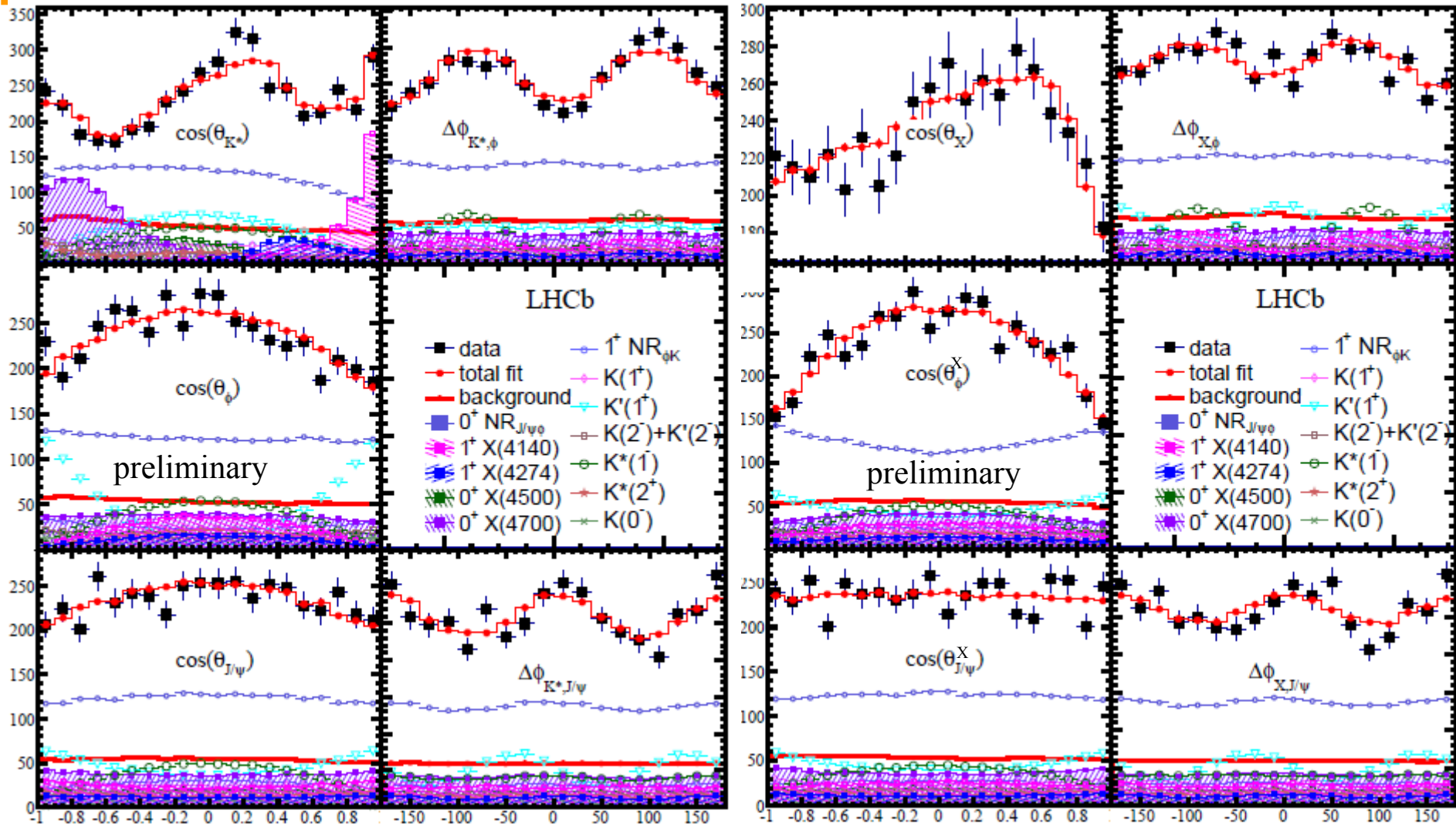
X(4140) from



[Phys. Rev. D 89, 012004 (2014)]



Fitted angles in $J/\psi\phi K$



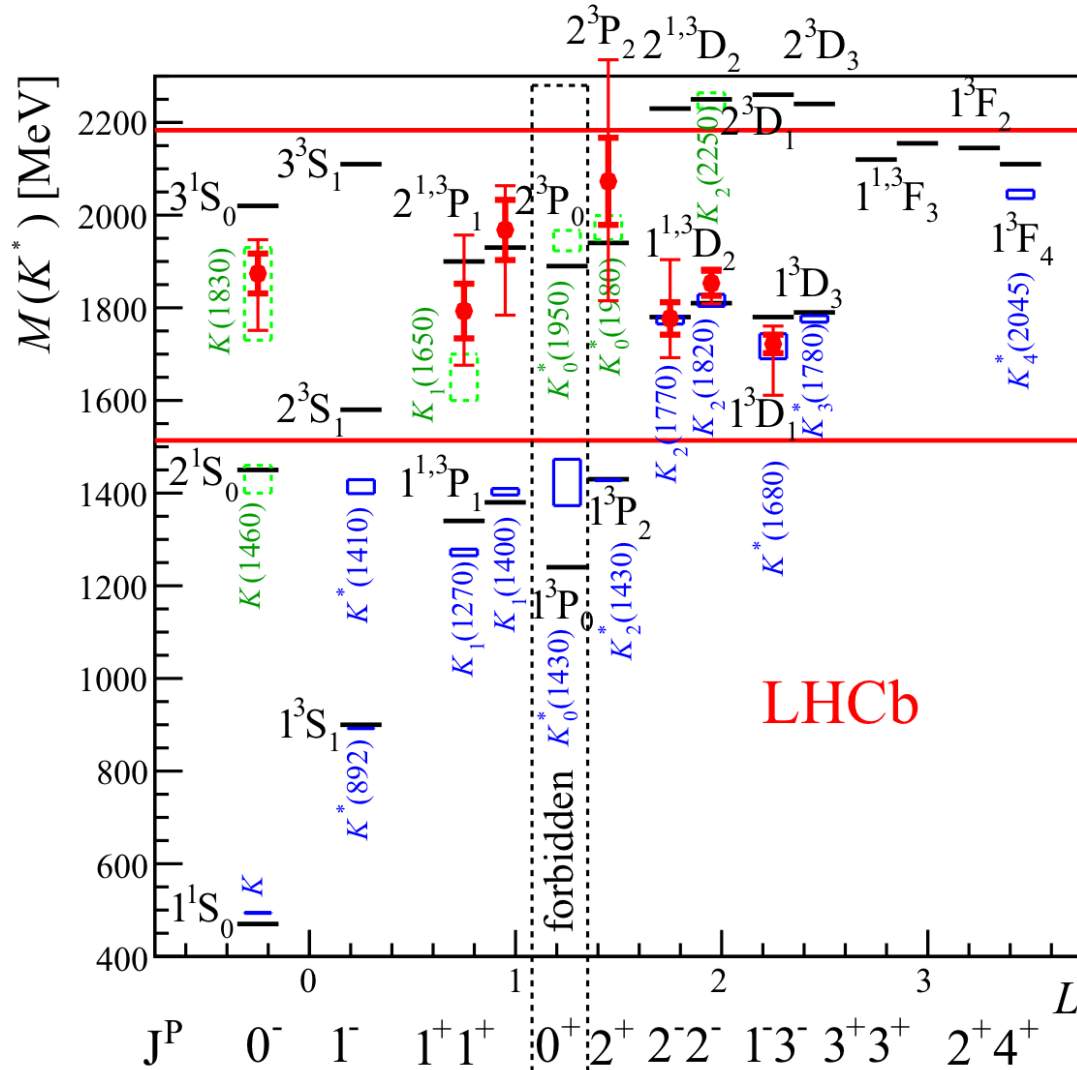
• Fit quality is good in all fitted variables

K* results

Default model: 6 K* resonances (of 4 different J^P) + 1 NR φK

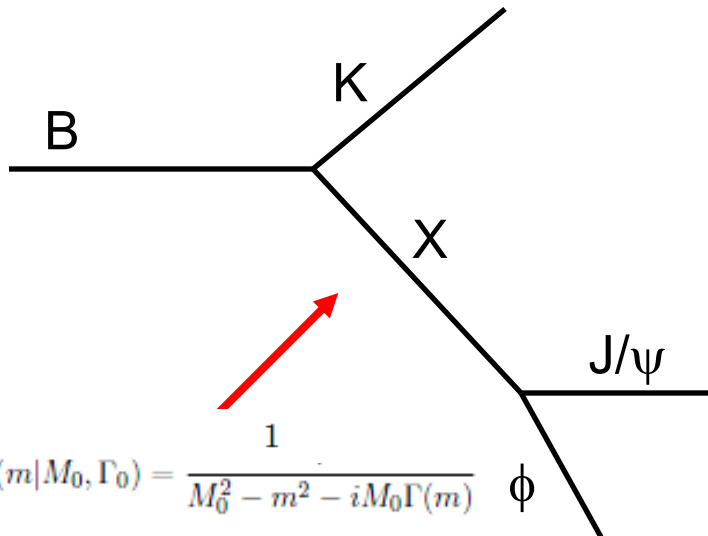
Our results are given by the red points with error bars

Excellent agreement between our results and both theory and previous experiments

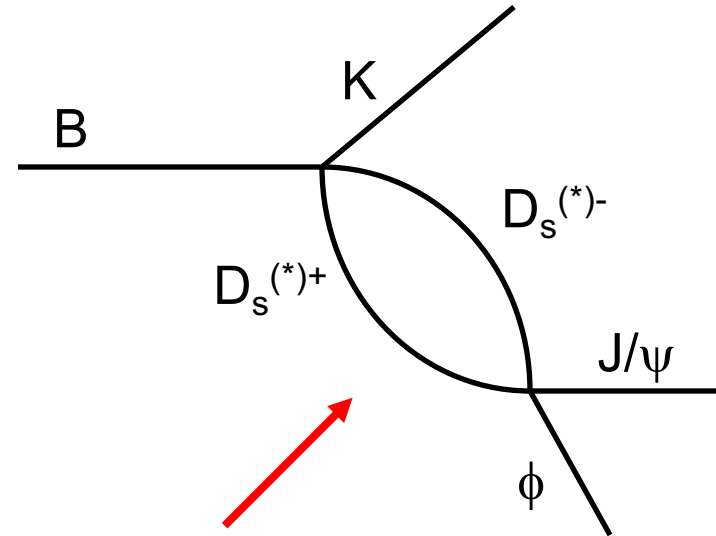


X exotics as $D_s^{(*)}D_s^{(*)\bar{}}$ cusps?

Breit-Wigner Model



Cusps Model



$$BW(m|M_0, \Gamma_0) = \frac{1}{M_0^2 - m^2 - iM_0\Gamma(m)} \phi$$

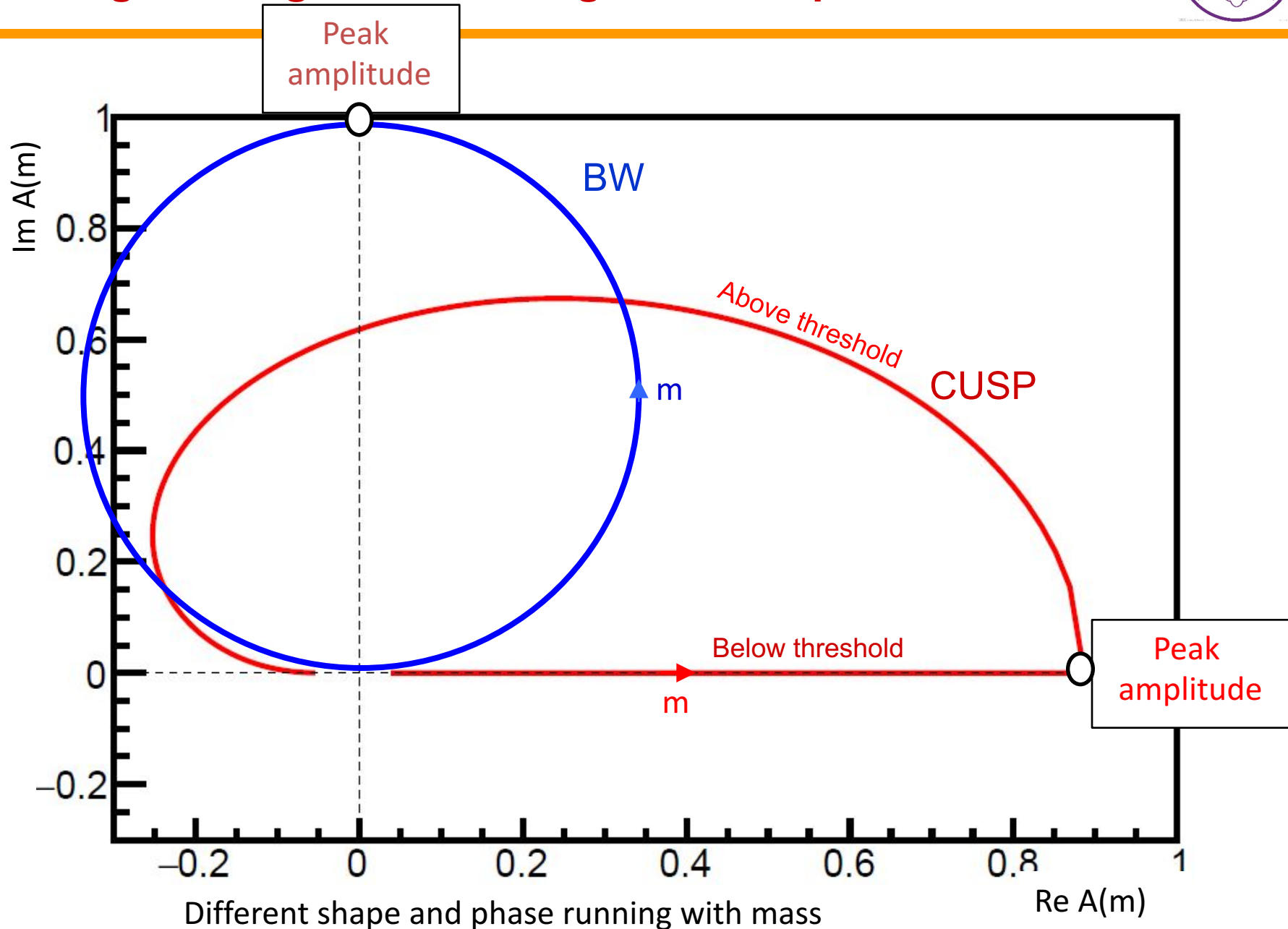
$$\Pi(m|\beta_0) = \int \frac{d^3q}{(2\pi)^3} \frac{e^{-2q^2/\beta_0}}{m - M_A - M_B - \frac{q^2}{2\mu_{AB}} + i\epsilon}$$

Cusp model by

E. S. Swanson, arXiv:1504.07952

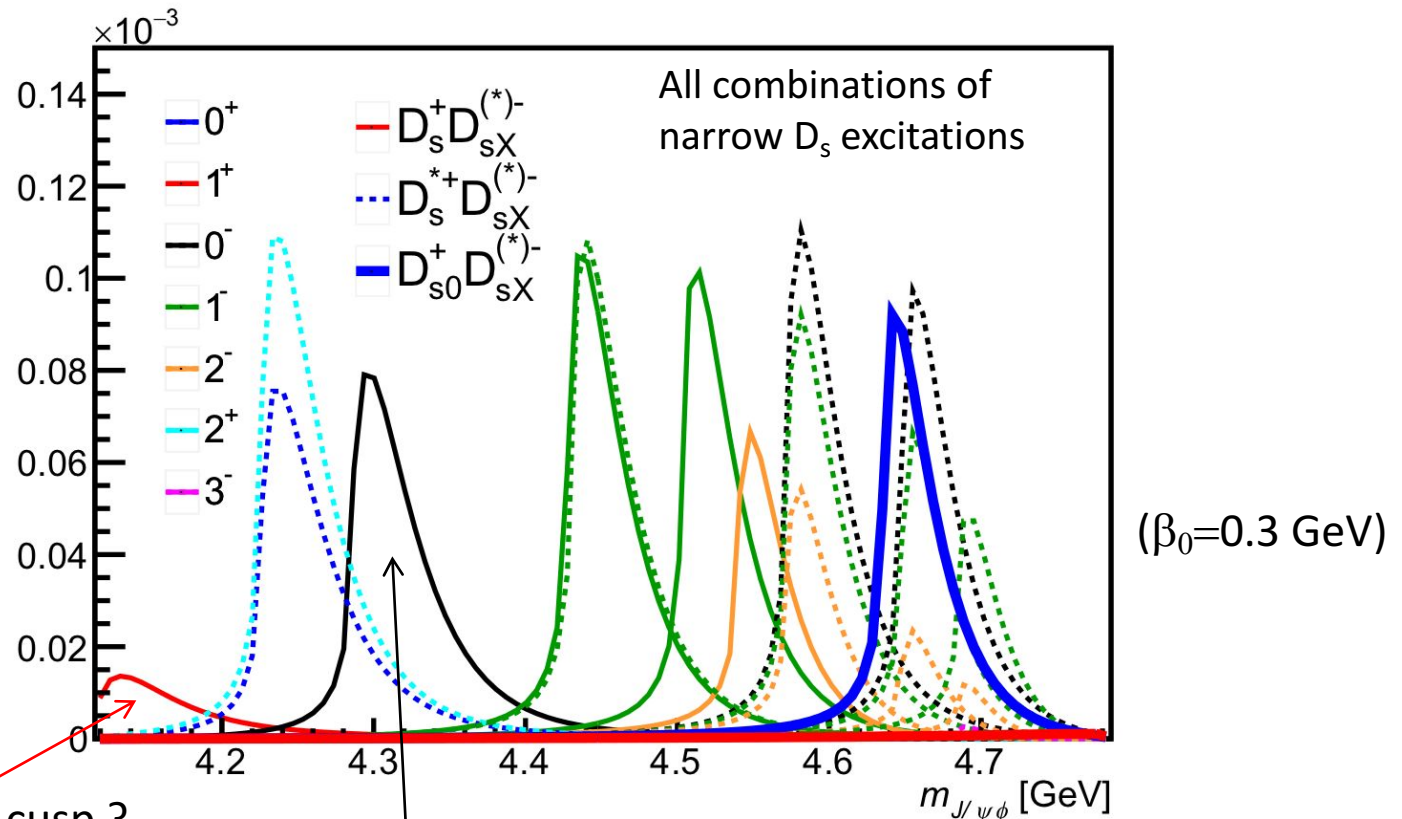
(see also PRD91, 034009 (2015))

Argand diagram: Breit-Wigner vs cusp



Cusps

- Cusp peaks at the sum of masses of the virtual narrow- $D_{sX}^{(*)}$ pairs.
- Width of cusp in Swanson model is controlled with a free parameter (β_0)
- J^P of cusp determined by J^P s of virtual D_s pairs (cusps occur in S-wave)



Is X(4140) a $D_s^+ D_s^{*-}$ cusp ?

Right $J^P=1^+$

Is X(4274) a $D_s^+ D_{s0}^{*-}$ cusp ?

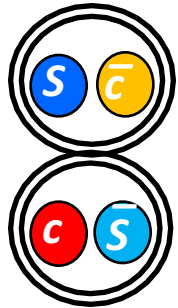
Wrong $J^P=0^-$

Theoretical interpretations of X(4140), X(4274)



Molecular models

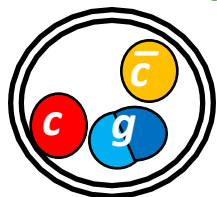
- The determination of the quantum numbers of X(4140) as $J^{PC}=1^{++}$ rules out many interpretations. Namely, 0^{++} or 2^{++} $D_s^* \bar{D}_s^*$ molecules. The large width is also not expected for true molecular bound states.



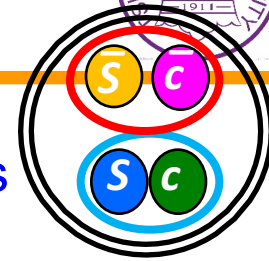
- However, X(4140) may be a 1^{++} $D_s \bar{D}_s^*$ cusp (form of rescattering)

Hybrid models

- Hybrid charmonium states proposed for X(4140) would have $J^{PC}=1^{-+}$. Thus they are also ruled out.



Tightly-bound tetraquark models

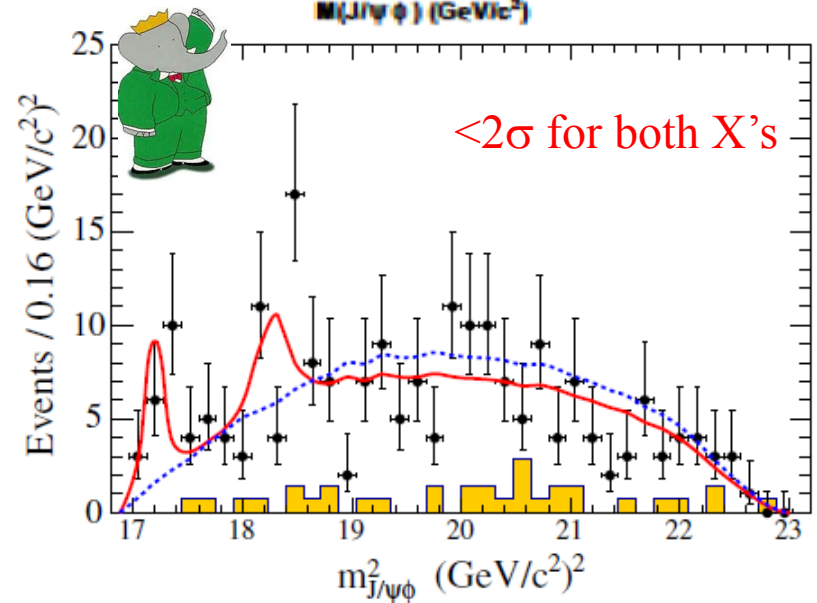
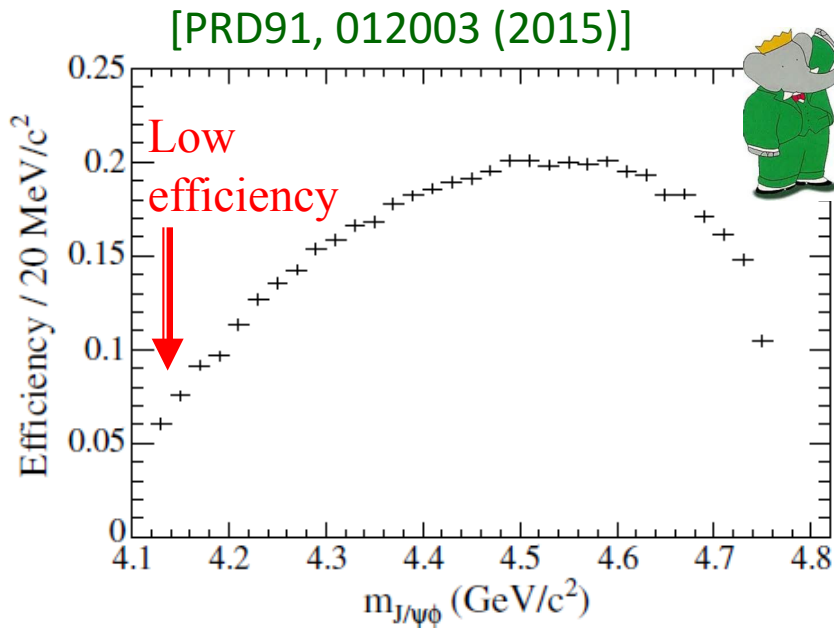
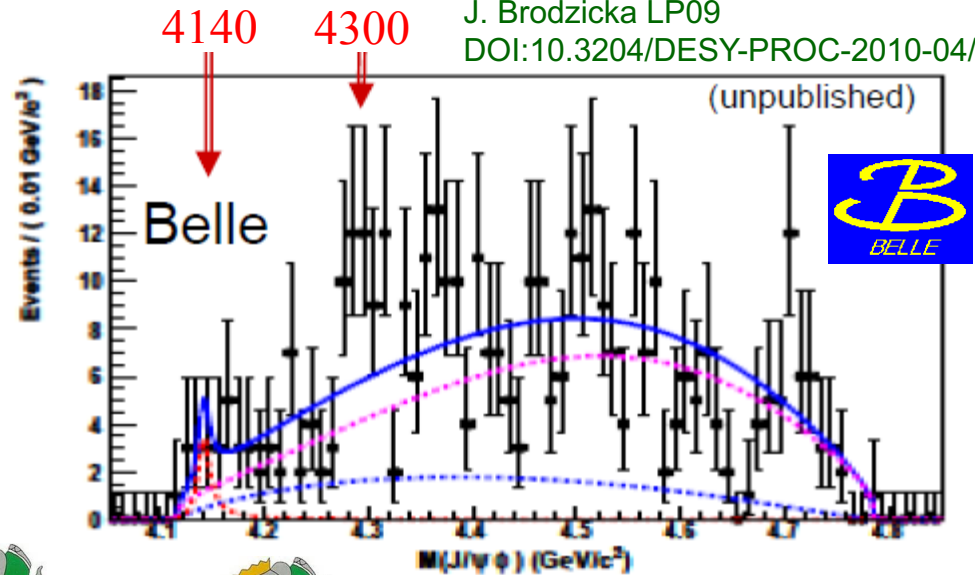


- There are tetraquark models which predict states with $J^{PC}=0^{-+}$, 1^{-+} or 0^{++} , 2^{++} near X(4140); these can be ruled out.
- A tetraquark model implemented by Stancu [JP G37, 075017 (2010), arXiv:0906.2485] correctly assigns 1^{++} to X(4140) and predicts a second 1^{++} state at a mass not much higher than X(4274)
- A Lattice calculation by Padmanth et al [PRD92, 034501 (2015)], based on a diquark tetraquark model, found no evidence for a 1^{++} tetraquark below 4.2 GeV

X(4140) from B-factories

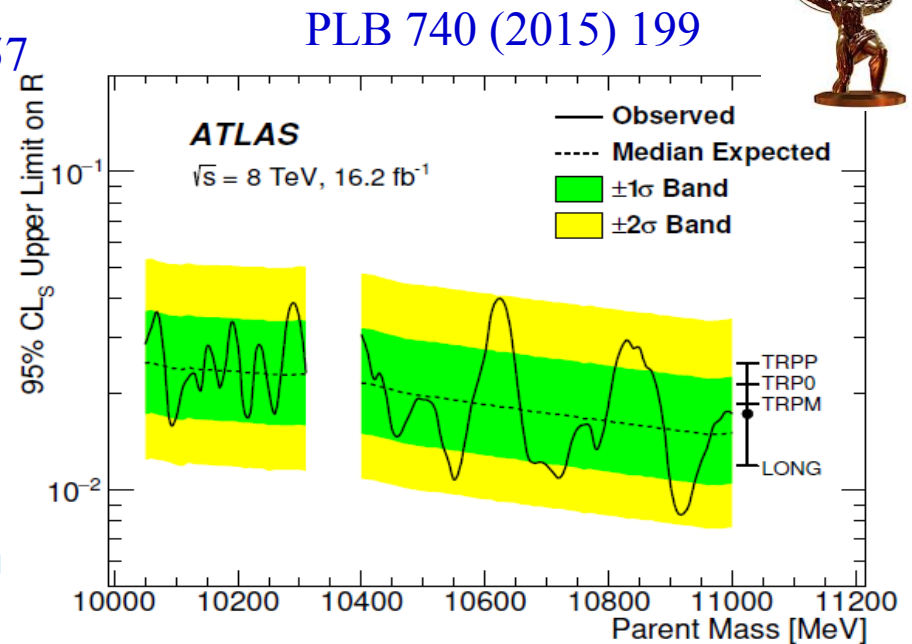
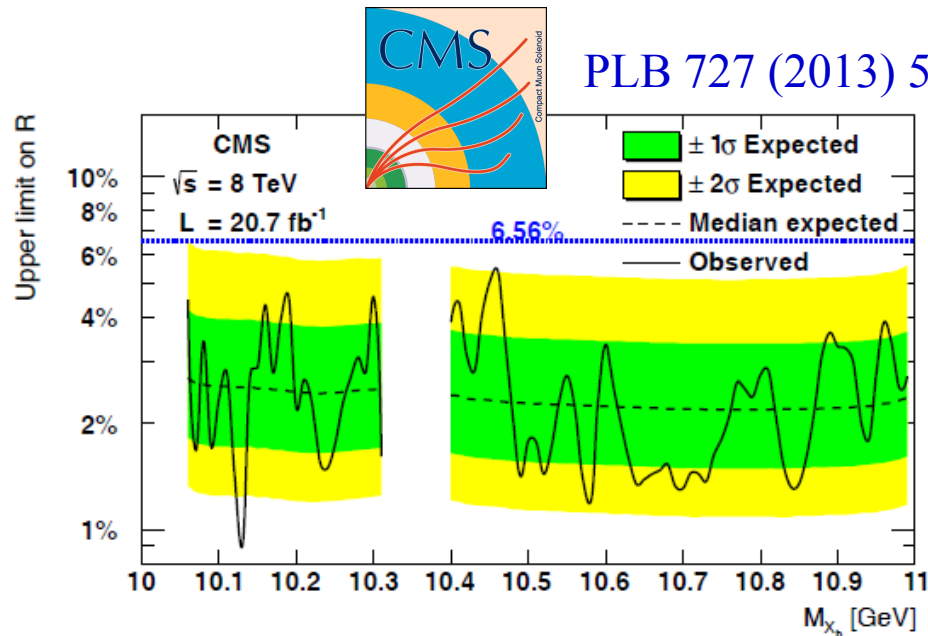
- No evidences reported by B-factories
 - Low efficiency at threshold region

J. Brodzicka LP09
DOI:10.3204/DESY-PROC-2010-04/38



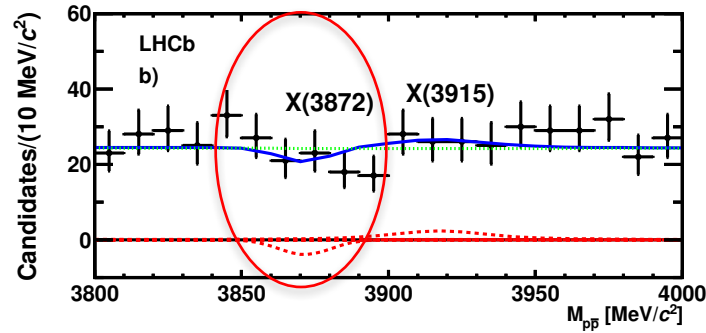
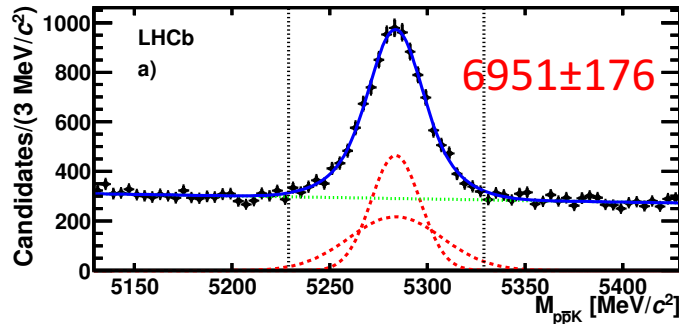
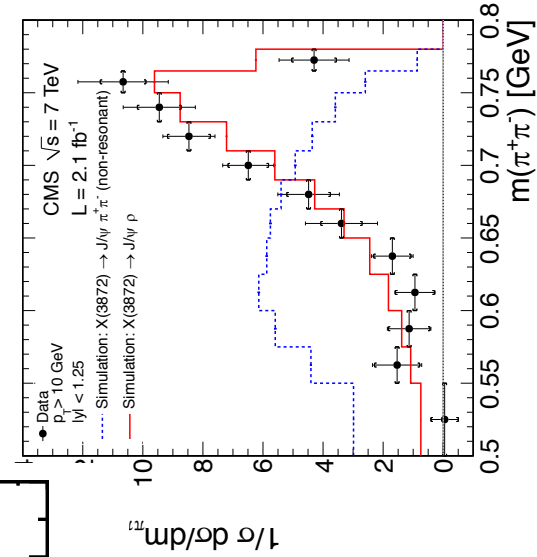
Search for X_b

- $X_b \rightarrow \Upsilon(1S)\pi^+\pi^-$: bottomonium counterpart of X(3872)
- Set 95% CL upper limit on $R = \frac{(\sigma\mathcal{B})_{X_b}}{(\sigma\mathcal{B})_{\Upsilon(2S)}}$ as a function of M_{X_b}



X(3872) decay

- CMS: $\pi^-\pi^+$ spectrum in $X(3872) \rightarrow J/\psi \pi^+ \pi^-$ decays, consistent with ρ^0 [JHEP 1304, 154 (2013)]
- LHCb: search for $X(3872) \rightarrow \bar{p}p$ decays, in $B^+ \rightarrow K^+ \bar{p}p$ with 1 fb^{-1} [EPJC73, 2642 (2013)]

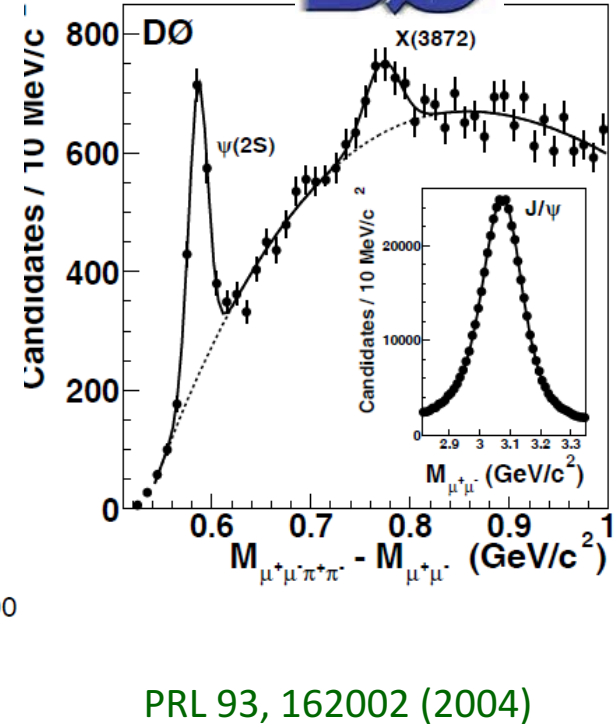
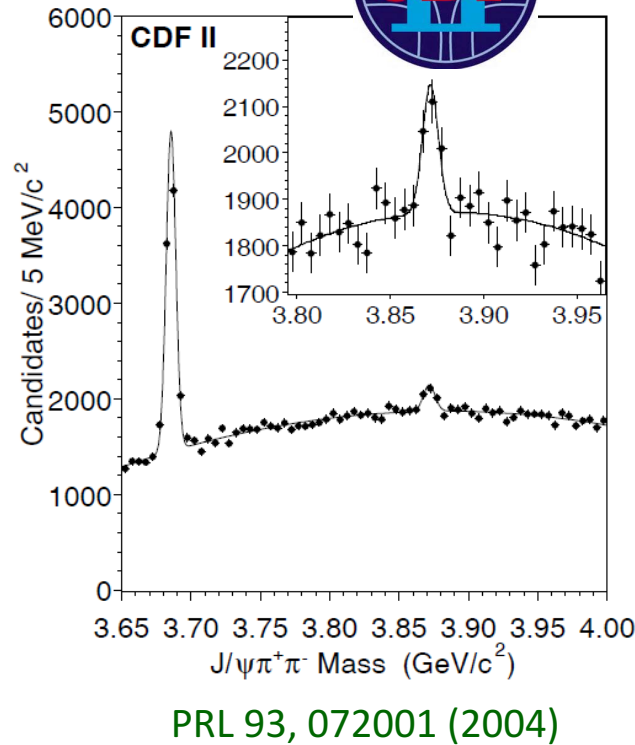
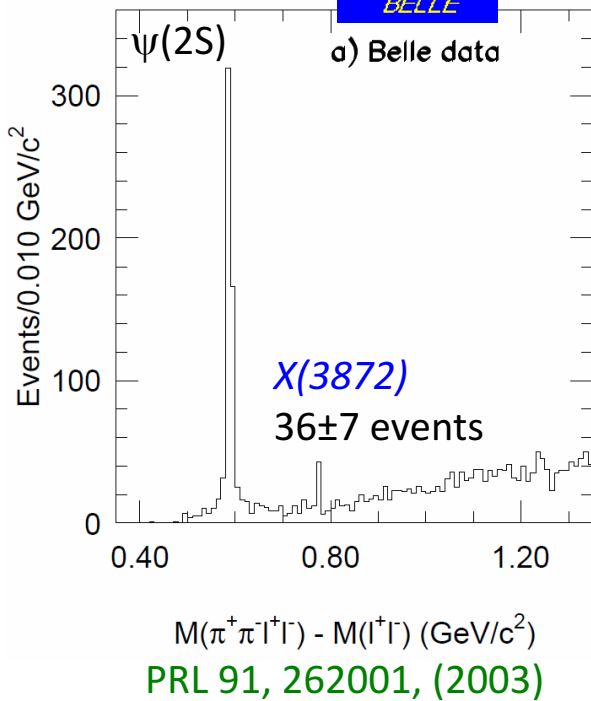


- Upper limit:

$$\frac{\text{BR}(X(3872) \rightarrow p\bar{p})}{\text{BR}(X(3872) \rightarrow J/\psi \pi^+ \pi^-)} < 2.0 \times 10^{-3}$$

X(3872)

- First exotic observed by Belle in 2003, and then confirmed by CDF and other experiments

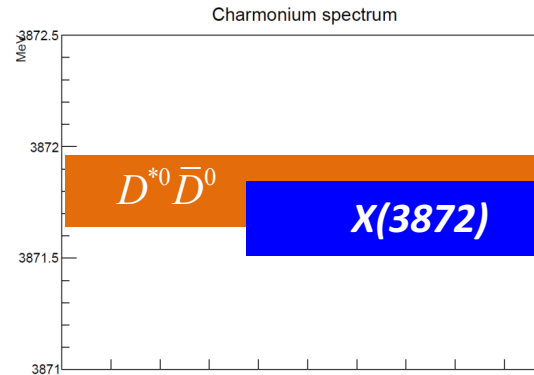


X(3872)

- Studied now in details:

- Mass = 3871.69 ± 0.17 MeV [PDG], Width < 1.2 MeV

$$m_{X(3872)} - (m_{D^{*0}} + m_{D^0}) = -0.11 \pm 0.22 \text{ MeV} \quad \text{[PDG]}$$



- Decays: open charm, charmonium ($J/\psi \pi^+ \pi^-$, ...)
- Production both in B decays and hadron collisions

Like conventional charmonium state

- It also decays with strong isospin violation

$$\frac{\mathcal{B}(X \rightarrow \omega J/\psi)}{\mathcal{B}(X \rightarrow \rho J/\psi)} \approx 1$$

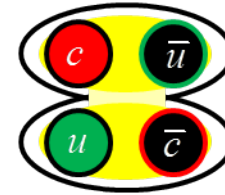
Comparable with a $D^{*0} \bar{D}^0$ molecular interpretation

X(3872) J^{PC} : Motivation

X(3872) quantum numbers will help to better understand the state.

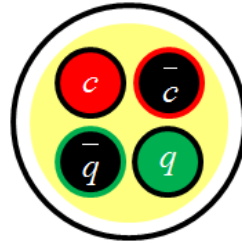
➤ $D^{*0}\bar{D}^0$ molecule, i.e. a $((u\bar{c})(c\bar{u}))$ system?

○ Requires $J^{PC} = 1^{++}$



➤ Tetraquark?

○ Allows with $J^{PC} = 1^{++}$

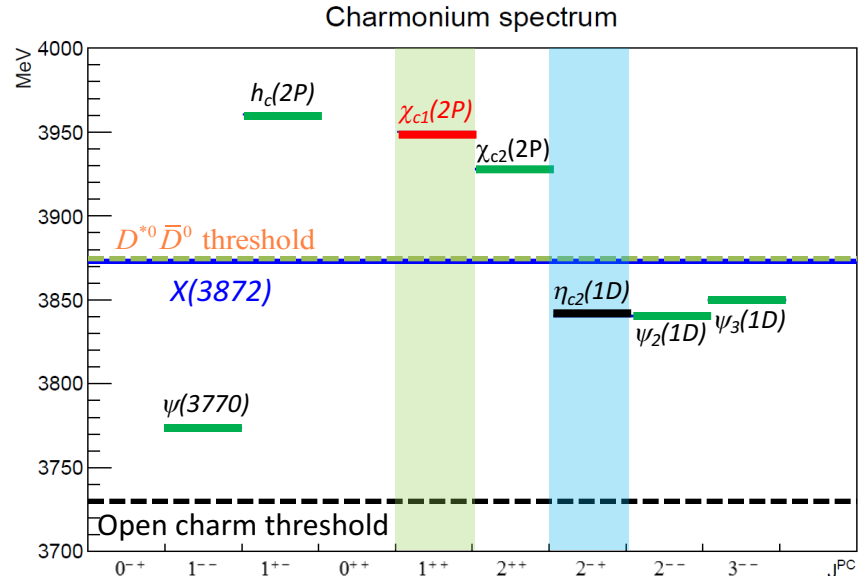
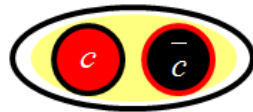


➤ Conventional charmonium states?

○ $\chi_{c1}(2^3P_1)$ requires $J^{PC} = 1^{++}$

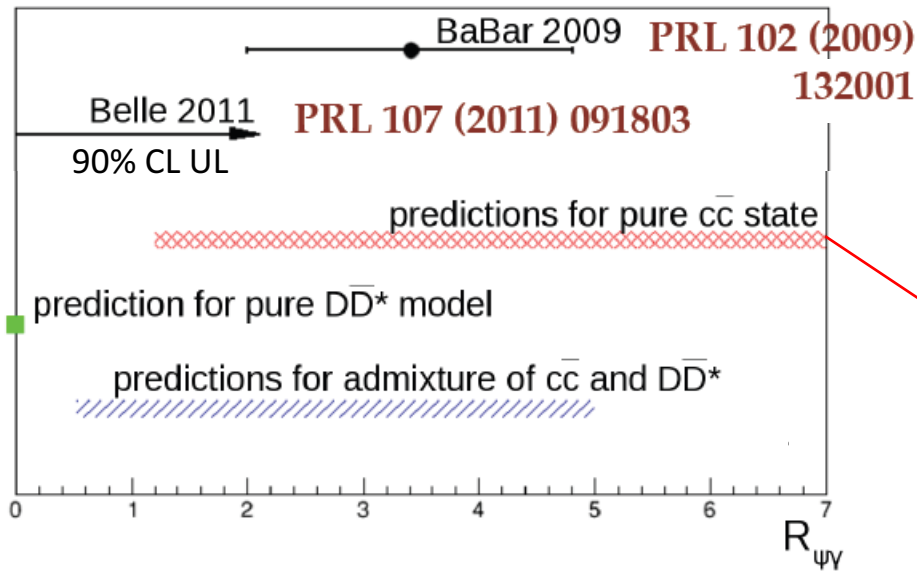
○ $\eta_{c2}(1^1D_2)$ requires $J^{PC} = 2^{-+}$

$(n^{2s+1}L_J)$



Radiative decays of X(3872)

- Measurement of $R_{\psi\gamma} = \frac{\mathcal{B}(X(3872) \rightarrow \psi(2S)\gamma)}{\mathcal{B}(X(3872) \rightarrow J/\psi\gamma)}$
a good probe for internal structure of X(3872)



efficiency($\psi(2S)\gamma$) / efficiency($J/\psi\gamma$) ~ 1

$\psi(2S)\gamma$ disfavored by phase-space
but enhanced by E1 matrix element:
 $|\langle 2S|r|2P\rangle|^2 \gg |\langle 1S|r|2P\rangle|^2$

For theory predictions, see references in Nucl.Phys. B886, 665 (2014)

- Previous measurements by BaBar and Belle barely consistent and favoring the opposite conclusions

Radiative decays of X(3872)@LHCb



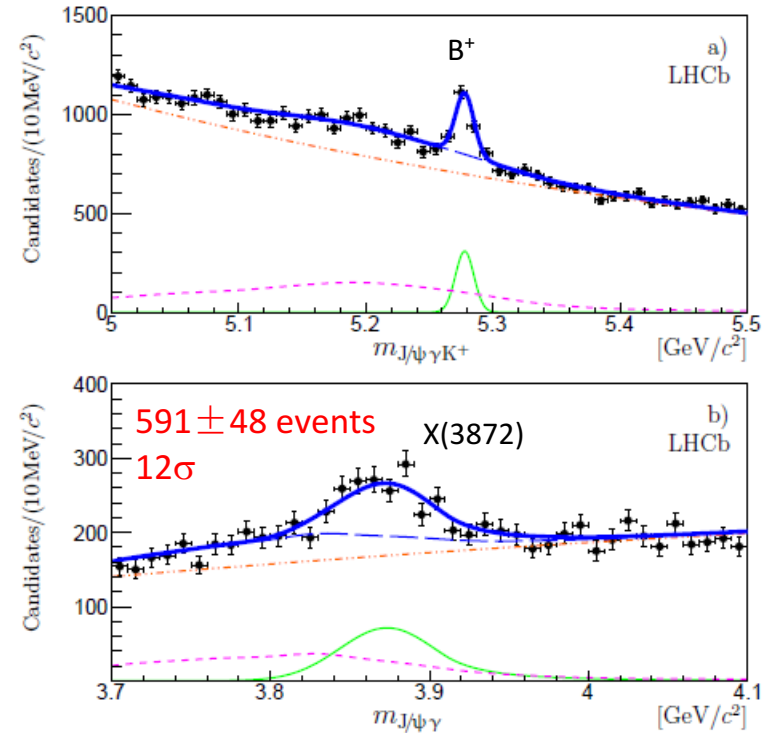
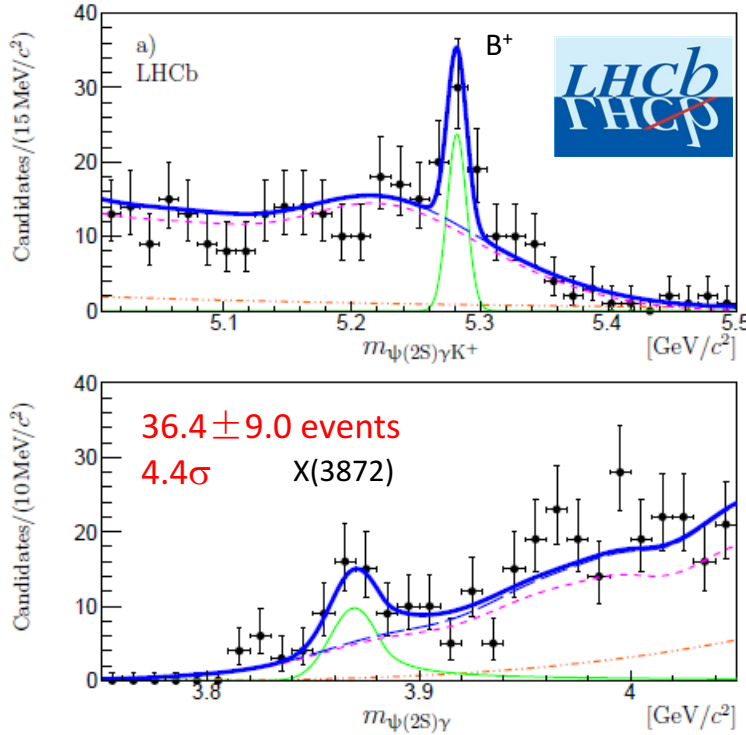
Nucl.Phys. B886, 665 (2014)

$B^+ \rightarrow X(3872)K^+$,
 $X(3872) \rightarrow \psi(2S)\gamma$

3fb⁻¹

$B^+ \rightarrow X(3872)K^+$,
 $X(3872) \rightarrow J/\psi\gamma$

Projections of 2D fit to $m_{\psi\gamma K^+}$ vs $m_{\psi\gamma}$



The most significant evidence for $X(3872) \rightarrow \psi(2S)\gamma$ to date!

efficiency($\psi(2S)\gamma$) / efficiency($J/\psi\gamma$) \sim 0.2

Detecting soft photons at hadron collider is hard.

Radiative decays of X(3872)@LHCb



Signal events:

Signal significance:

$B^+ \rightarrow X(3872)K^+$,

$X(3872) \rightarrow \psi(2S)\gamma, J/\psi\gamma$

$\psi(2S)\gamma, J/\psi\gamma$

$25.4 \pm 7.3,$

23.0 ± 6.4

$3.6\sigma, 3.5\sigma$

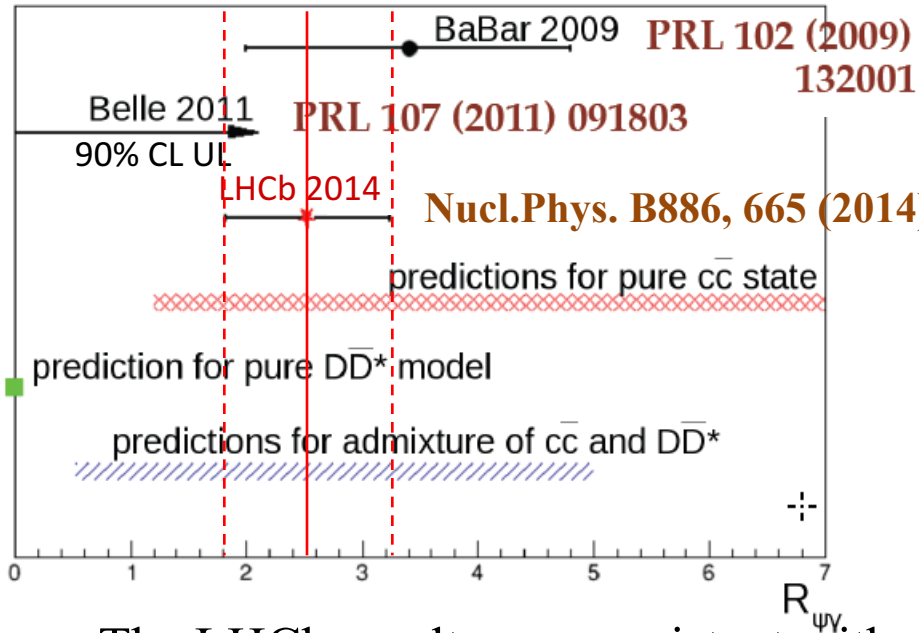
$5.0^{+11.9}_{-11.0},$

$30.0^{+8.2}_{-7.4}$

$0.4\sigma, 4.9\sigma$

$36.4 \pm 9.0, 591.0 \pm 48.0$

$4.4\sigma, 12\sigma$

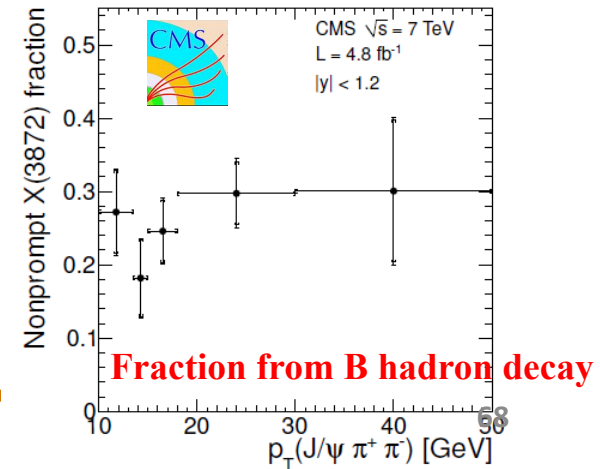
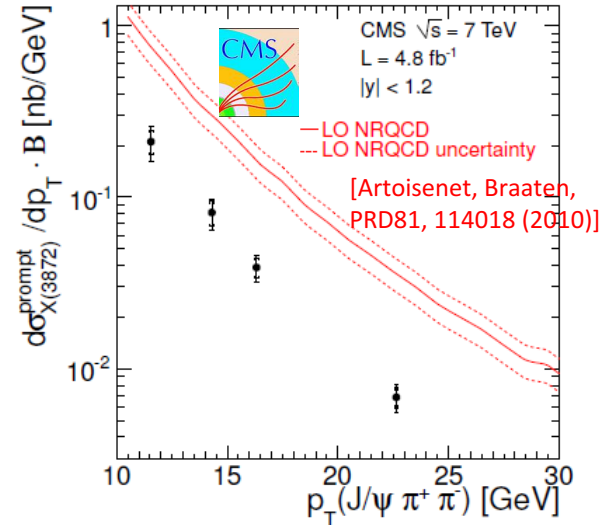
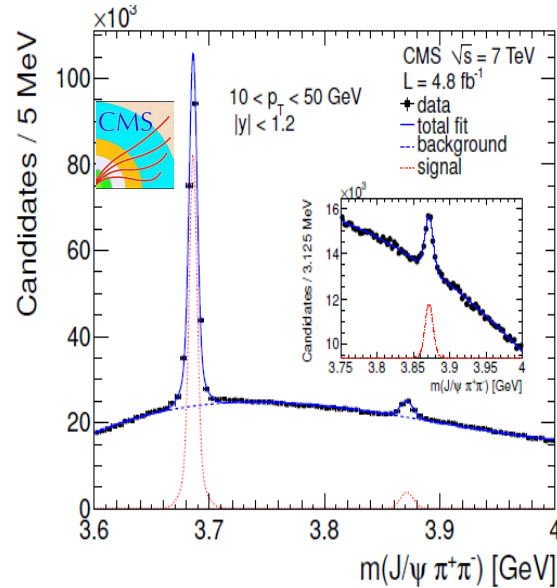
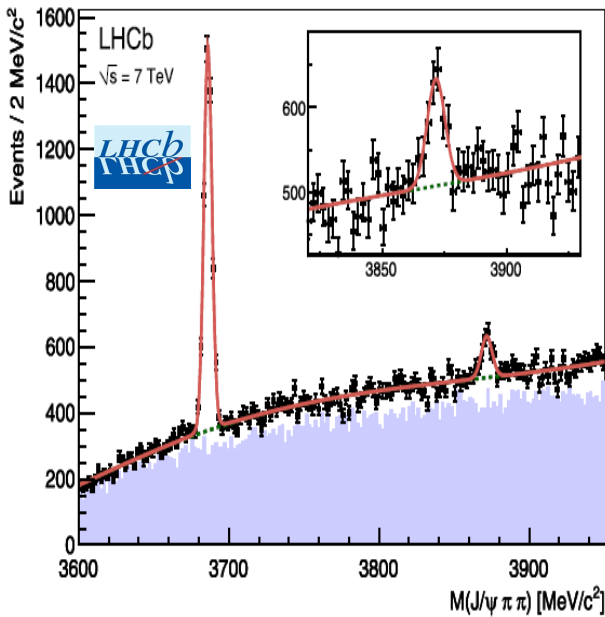


$$R_{\psi\gamma} = 2.48 \pm 0.64 \pm 0.29$$

- The LHCb results are consistent with, but more precise than, the BaBar and Belle results
- The results are not consistent with the expectations for pure molecular X(3872)
- The favorite X(3872) interpretation is a mixture of a $\chi_{c1}(2^3P_1)$ charmonium state and of $D^{*0}\bar{D}^0$ molecule or cusp [arXiv:1404.3723]

X(3872) production

- All hadron collisions reported X(3872)
 - CDF [PRL93, 072001 (2004)], D0 [PRL93, 162002 (2004)], [LHCb](#) [EPJC72, 1972 (2012)] and [CMS](#) [JHEP 1304, 154 (2013)].
- Study regions: central ([CMS](#), $|y| < 1.2$) or forward ([LHCb](#), $2.5 < y < 4.5$)



LHCb: $\sigma_{inclusive} \mathcal{B}(X(3872) \rightarrow J/\psi \pi^+ \pi^-)$ [$2.5 < y < 4.5$, $p_T > 5$ GeV] = $5.4 \pm 1.3 \pm 0.8$ nb

CMS: $\sigma_{prompt} \mathcal{B}(X(3872) \rightarrow J/\psi \pi^+ \pi^-)$ [$|y| < 1.2$, $p_T > 10$ GeV] = $1.06 \pm 0.11 \pm 0.15$ nb

CMS: Fraction of X(3872) from B hadron = $(26.3 \pm 2.3 \pm 1.6)\%$