



Study of the Lorentz structure of tau decays at Belle

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@China National Convention Center

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- Belle experiment
- Study of Michel parameters at Belle
 - $\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau$
 - $\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau \gamma$
 - $\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau l'^+ l'^-$
- Conclusion & Future Plan

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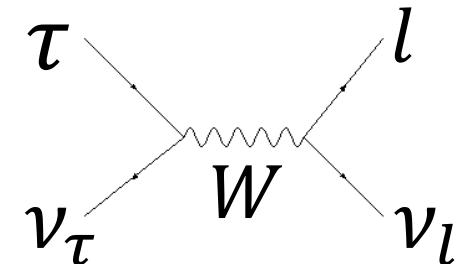
Introduction

Standard Model (SM)'s prediction: A Lorentz structure of charged weak current has a V-A structure.

We study the Lorentz structure through the measurement of **Michel parameters**.

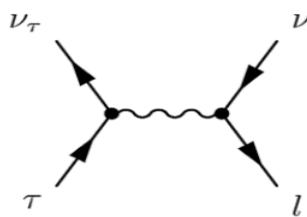
$$\mathcal{L} \propto \sum_{\substack{i=S,V,T \\ \lambda,\rho=L,R}} \{g_{\lambda\rho}^i [\bar{l}'_\lambda \Gamma^i (\nu_{l'})_\xi] [(\bar{\nu}_l)_\kappa \Gamma_i l_\rho]\}$$

In the SM, $g_{LL}^V = 1$, and others are $g_{\lambda\rho}^i = 0$

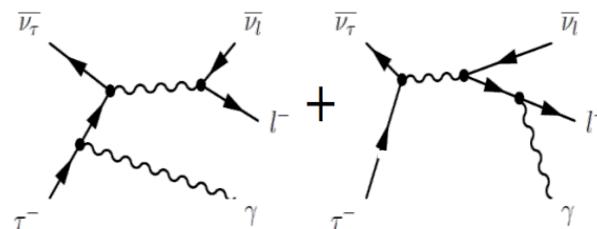


Michel parameters are the bi-linear combination of coupling constants $g_{\lambda\rho}^i$. The different type of Michel parameter appears in three modes shown below.

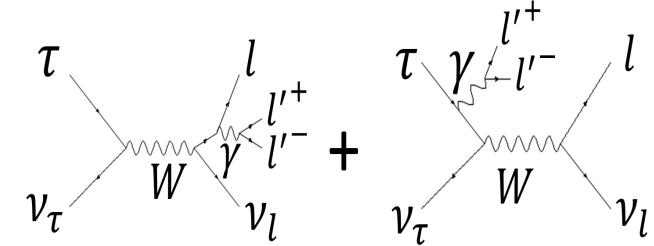
$$\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau$$



$$\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau \gamma$$



$$\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau l'^+ l'^-$$



Introduction

Measurable Michel parameters (MPs) for each tau's leptonic decay and what the coupling constants $g_{\lambda\rho}^i$ are contained, are shown below.

$$\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau,$$

$$\rho = \frac{3}{4} - \frac{3}{4} \left(|g_{LR}^V|^2 + |g_{RL}^V|^2 + 2|g_{LR}^T|^2 + 2|g_{RL}^T|^2 + \Re(g_{LR}^S g_{LR}^{T*} + g_{RL}^S g_{RL}^{T*}) \right)$$

$$\eta = \frac{1}{2} \Re \left(6g_{RL}^V g_{LR}^{T*} + 6g_{LR}^V g_{RL}^{T*} + g_{RR}^S g_{LL}^{V*} + g_{RL}^S g_{LR}^{V*} + g_{LR}^S g_{RL}^{V*} + g_{LL}^S g_{RR}^{V*} \right)$$

$$\xi = 4\Re(g_{LR}^S g_{LR}^{T*}) - 4\Re(g_{RL}^S g_{RL}^{T*}) + |g_{LL}^V|^2 + 3|g_{LR}^V|^2 - 3|g_{RL}^V|^2 - |g_{RR}^V|^2 + 5|g_{LR}^T|^2 - 5|g_{RL}^T|^2 + \frac{1}{4}|g_{LL}^S|^2 - \frac{1}{4}|g_{LR}^S|^2 + \frac{1}{4}|g_{RL}^S|^2 - \frac{1}{4}|g_{RR}^S|^2$$

$$\xi\delta = \frac{3}{16}|g_{LL}^S|^2 - \frac{3}{16}|g_{LR}^S|^2 + \frac{3}{16}|g_{RL}^S|^2 - \frac{3}{16}|g_{RR}^S|^2 - \frac{3}{4}|g_{LR}^T|^2 + \frac{3}{4}|g_{RL}^T|^2 + \frac{3}{4}|g_{LL}^V|^2 - \frac{3}{4}|g_{RR}^V|^2 + \frac{3}{4}\Re(g_{LR}^S g_{LR}^{T*}) - \frac{3}{4}\Re(g_{RL}^S g_{RL}^{T*})$$

$$\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau \gamma$$

$$\bar{\eta} = |g_{RL}^V|^2 + |g_{LR}^V|^2 + \frac{1}{8} \left(|g_{RL}^S + 2g_{RL}^T|^2 + |g_{LR}^S + 2g_{LR}^T|^2 \right) + 2 \left(|g_{RL}^T|^2 + |g_{LR}^T|^2 \right)$$

$$\xi\kappa = |g_{RL}^V|^2 - |g_{LR}^V|^2 + \frac{1}{8} \left(|g_{RL}^S + 2g_{RL}^T|^2 - |g_{LR}^S + 2g_{LR}^T|^2 \right) + 2 \left(|g_{RL}^T|^2 - |g_{LR}^T|^2 \right)$$

Some coupling constants $g_{\lambda\rho}^i$ are not well known as shown in right-side table.

Measurement of MPs is the important clue for studying Lorentz structure.

$$\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau l'^+ l'^-$$

$$Q_{LL} = \frac{1}{4} |g_{LL}^S|^2 + |g_{LL}^V|^2$$

$$Q_{RL} = \frac{1}{4} |g_{RL}^S|^2 + |g_{RL}^V|^2 + |g_{RL}^T|^2$$

$$Q_{RL} = \frac{1}{4} |g_{LR}^S|^2 + |g_{LR}^V|^2 + |g_{LR}^T|^2$$

$$Q_{RR} = \frac{1}{4} |g_{RR}^S|^2 + |g_{RR}^V|^2$$

$$B_{RL} = \frac{1}{16} |g_{RL}^S + 6g_{RL}^T|^2 + |g_{RL}^V|^2$$

$$B_{LR} = \frac{1}{16} |g_{LR}^S + 6g_{LR}^T|^2 + |g_{LR}^V|^2$$

$$I_\alpha = \frac{1}{4} g_{LR}^V (g_{RL}^S + 6g_{RL}^T)^* + \frac{1}{4} g_{RL}^V (g_{LR}^S + 6g_{LR}^T)$$

$$I_\beta = g_{LL}^V g_{RR}^{S*}/2 + g_{RR}^V g_{LL}^S/2$$

Recent Constraints to $g_{\lambda\rho}^i$

$$\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau$$

$ g_{RR}^S < 0.70$	$ g_{LR}^S < 0.99$	$ g_{RL}^S \leq 2$	$ g_{LL}^S \leq 2$
$ g_{RR}^V < 0.17$	$ g_{LR}^V < 0.13$	$ g_{RL}^V < 0.52$	$ g_{LL}^V \leq 1$
$ g_{RR}^T \equiv 0$	$ g_{LR}^T < 0.082$	$ g_{RL}^T < 0.51$	$ g_{LL}^T \equiv 0$

$$\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau$$

$ g_{RR}^S < 0.72$	$ g_{LR}^S < 0.95$	$ g_{RL}^S \leq 2$	$ g_{LL}^S \leq 2$
$ g_{RR}^V < 0.18$	$ g_{LR}^V < 0.12$	$ g_{RL}^V < 0.52$	$ g_{LL}^V \leq 1$
$ g_{RR}^T \equiv 0$	$ g_{LR}^T < 0.079$	$ g_{RL}^T < 0.51$	$ g_{LL}^T \equiv 0$

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Belle Experiment (Tsukuba, Japan)

- Integrated Luminosity: 1000 fb^{-1}
- Energy of Collision e^- (8 GeV) / e^+ (3.5 GeV)
- $\sqrt{s} = 10.58 \text{ GeV}$ (Same energy to generate BB-pair)
- Low background & Energy of center of mass frame is known
- Pair creation ($B\bar{B}$, $\tau^+\tau^-$, etc). One-side can be used for tagging.

$$e^+e^- \rightarrow B\bar{B} \quad \sigma_{BB} \sim 1.05 \text{ nb (neutral+charged)}$$

$$e^+e^- \rightarrow \tau^+\tau^- \quad \sigma_{\tau\tau} = (0.919 \pm 0.003) \text{ nb}$$

Belle is a B factory, and also a **tau-factory**

$$N_{\tau\tau} \sim 9.0 \times 10^8$$

> 1 ab⁻¹

On resonance:

$Y(5S): 121 \text{ fb}^{-1}$

$Y(4S): 711 \text{ fb}^{-1}$

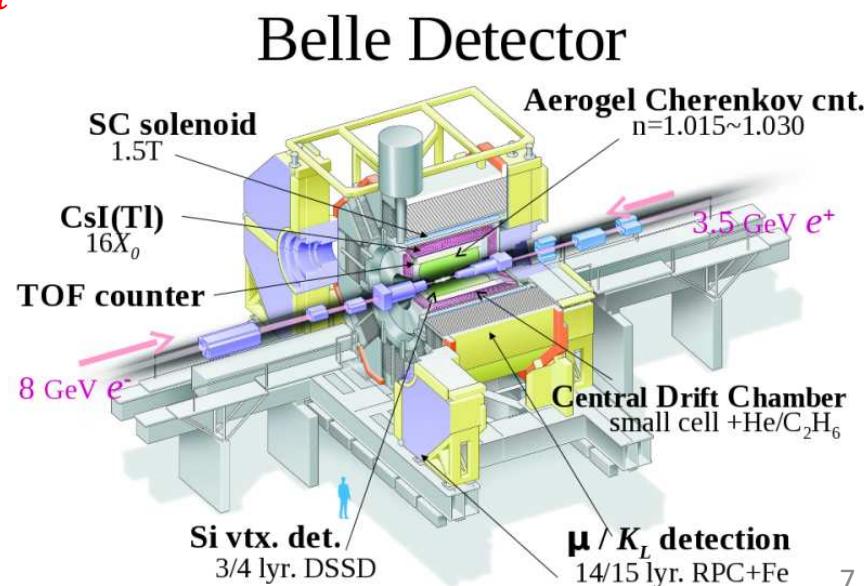
$Y(3S): 3 \text{ fb}^{-1}$

$Y(2S): 25 \text{ fb}^{-1}$

$Y(1S): 6 \text{ fb}^{-1}$

Off reson./scan:

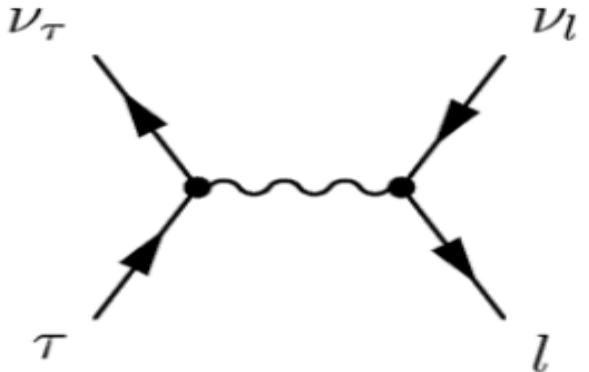
$\sim 100 \text{ fb}^{-1}$



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$$\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau$$



Michel Parameter

Target MPs: ρ, η, ξ, δ

$$\rho = \frac{3}{4} - \frac{3}{4} \left(|g_{LR}^V|^2 + |g_{RL}^V|^2 + 2|g_{LR}^T|^2 + 2|g_{RL}^T|^2 + \Re(g_{LR}^S g_{LR}^{T*} + g_{RL}^S g_{RL}^{T*}) \right)$$

$$\eta = \frac{1}{2} \Re \left(6g_{RL}^V g_{LR}^{T*} + 6g_{LR}^V g_{RL}^{T*} + g_{RR}^S g_{LL}^{V*} + g_{RL}^S g_{LR}^{V*} + g_{LR}^S g_{RL}^{V*} + g_{LL}^S g_{RR}^{V*} \right)$$

$$\xi = 4\Re(g_{LR}^S g_{LR}^{T*}) - 4\Re(g_{RL}^S g_{RL}^{T*}) + |g_{LL}^V|^2 + 3|g_{LR}^V|^2 - 3|g_{RL}^V|^2 - |g_{RR}^V|^2 +$$

$$+ 5|g_{LR}^T|^2 - 5|g_{RL}^T|^2 + \frac{1}{4}|g_{LL}^S|^2 - \frac{1}{4}|g_{LR}^S|^2 + \frac{1}{4}|g_{RL}^S|^2 - \frac{1}{4}|g_{RR}^S|^2$$

$$\xi\delta = \frac{3}{16}|g_{LL}^S|^2 - \frac{3}{16}|g_{LR}^S|^2 + \frac{3}{16}|g_{RL}^S|^2 - \frac{3}{16}|g_{RR}^S|^2 - \frac{3}{4}|g_{LR}^T|^2 + \frac{3}{4}|g_{RL}^T|^2 +$$

$$+ \frac{3}{4}|g_{LL}^V|^2 - \frac{3}{4}|g_{RR}^V|^2 + \frac{3}{4}\Re(g_{LR}^S g_{LR}^{T*}) - \frac{3}{4}\Re(g_{RL}^S g_{RL}^{T*})$$

Formalism of $\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau$'s differential decay width

$$\frac{d\Gamma(\tau^\mp)}{d\Omega dx} = \frac{4G_F^2 M_\tau E_{\max}^4}{(2\pi)^4} \sqrt{x^2 - x_0^2} \left(x(1-x) + \frac{2}{9} \rho (4x^2 - 3x - x_0^2) + \eta x_0(1-x) \right)$$

$$+ \frac{1}{3} P_\tau \cos\theta_\ell \xi \sqrt{x^2 - x_0^2} \left[1 - x + \frac{2}{3} \delta (4x - 4 + \sqrt{1 - x_0^2}) \right] \right), \quad x = \frac{E_\ell}{E_{\max}}, \quad x_0 = \frac{m_\ell}{E_{\max}}$$

In the SM: $\rho = \frac{3}{4}$, $\eta = 0$, $\xi = 1$, $\delta = \frac{3}{4}$

Formalism of $\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau$'s differential decay width

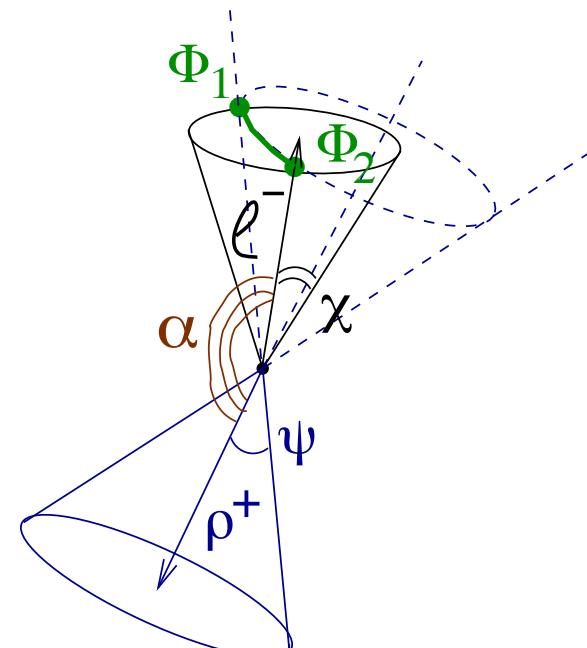
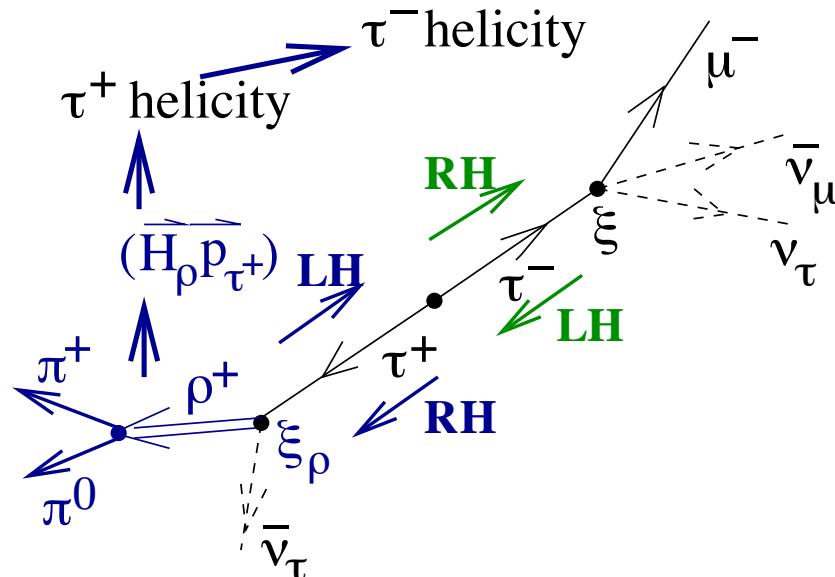
$$\frac{d\Gamma(\tau^\mp)}{d\Omega dx} = \frac{4G_F^2 M_\tau E_{\max}^4}{(2\pi)^4} \sqrt{x^2 - x_0^2} \left(x(1-x) + \frac{2}{9} \rho (4x^2 - 3x - x_0^2) + \eta x_0(1-x) \right)$$

$$\pm \frac{1}{3} P_\tau \cos \theta_\ell \xi \sqrt{x^2 - x_0^2} \left[1 - x + \frac{2}{3} \delta (4x - 4 + \sqrt{1 - x_0^2}) \right], \quad x = \frac{E_\ell}{E_{\max}}, \quad x_0 = \frac{m_\ell}{E_{\max}}$$

In the SM: $\rho = \frac{3}{4}$, $\eta = 0$, $\xi = 1$, $\delta = \frac{3}{4}$

Point of this analysis:

1. Information of τ – spin direction is needed.
2. To extract τ – spin direction, we use a mode $\tau^\pm \rightarrow \rho^\pm \nu$ for tag-side.



Method, study of $(\ell\nu\nu; \rho\nu)$ and $(\rho\nu; \rho\nu)$ events

Reference: D. Epifanov's slide@PhiPsi-2017

Effect of τ spin-spin correlation is used to measure ξ and δ MP.

Events of the $(\tau^\mp \rightarrow \ell^\mp \nu\nu; \tau^\pm \rightarrow \rho^\pm \nu)$ topology are used to measure: ρ , η , $\xi_\rho \xi$ and $\xi_\rho \xi \delta$, while $(\tau^\mp \rightarrow \rho^\mp \nu; \tau^\pm \rightarrow \rho^\pm \nu)$ events are used to extract ξ_ρ^2 .

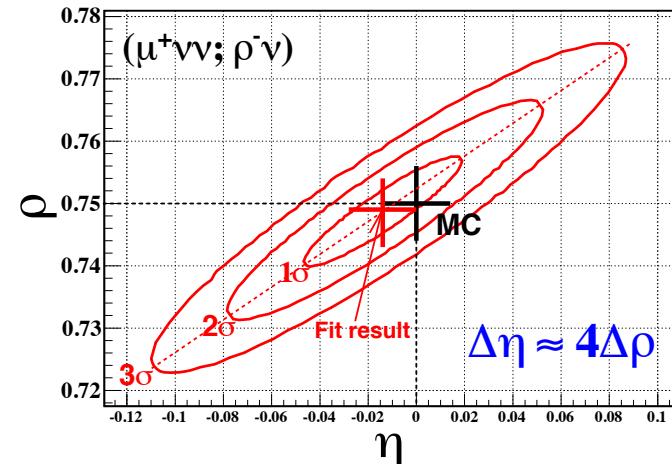
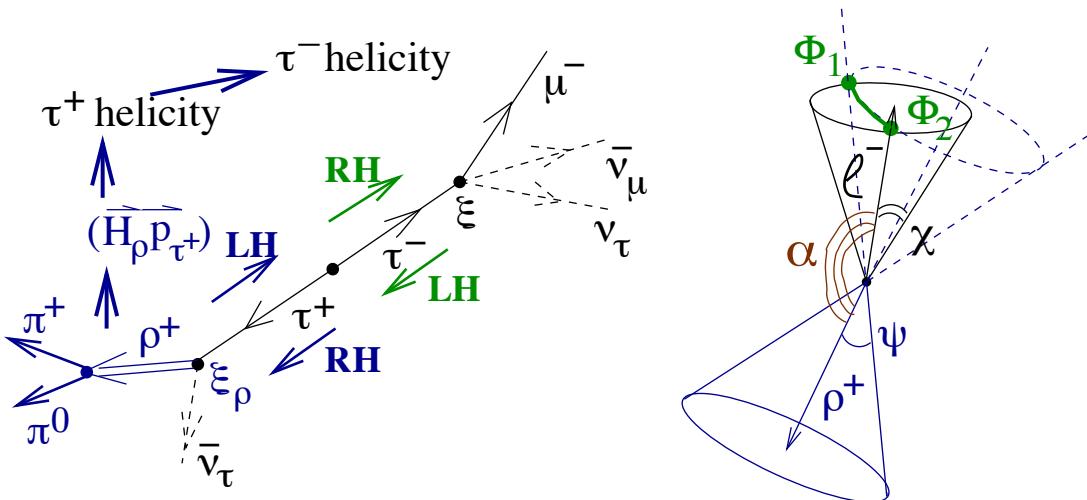
$$\frac{d\sigma(\ell^\mp \nu\nu, \rho^\pm \nu)}{dE_\ell^* d\Omega_\ell^* d\Omega_\rho^* dm_{\pi\pi}^2 d\tilde{\Omega}_\pi d\Omega_\tau} = A_0 + \rho A_1 + \eta A_2 + \xi_\rho \xi A_3 + \xi_\rho \xi \delta A_4 = \sum_{i=0}^4 A_i \Theta_i$$

$$\mathcal{F}(\vec{z}) = \frac{d\sigma(\ell^\mp \nu\nu, \rho^\pm \nu)}{dp_\ell d\Omega_\ell dp_\rho d\Omega_\rho dm_{\pi\pi}^2 d\tilde{\Omega}_\pi} = \int_{\Phi_1}^{\Phi_2} \frac{d\sigma(\ell^\mp \nu\nu, \rho^\pm \nu)}{dE_\ell^* d\Omega_\ell^* d\Omega_\rho^* dm_{\pi\pi}^2 d\tilde{\Omega}_\pi d\Omega_\tau} \left| \frac{\partial(E_\ell^*, \Omega_\ell^*, \Omega_\rho^*, \Omega_\tau)}{\partial(p_\ell, \Omega_\ell, p_\rho, \Omega_\rho, \Phi_\tau)} \right| d\Phi_\tau$$

$$L = \prod_{k=1}^N \mathcal{P}^{(k)}, \quad \mathcal{P}^{(k)} = \mathcal{F}(\vec{z}^{(k)}) / \mathcal{N}(\vec{\Theta}), \quad \mathcal{N}(\vec{\Theta}) = \int \mathcal{F}(\vec{z}) d\vec{z}, \quad \vec{\Theta} = (1, \rho, \eta, \xi_\rho \xi_\ell, \xi_\rho \xi_\ell \delta_\ell)$$

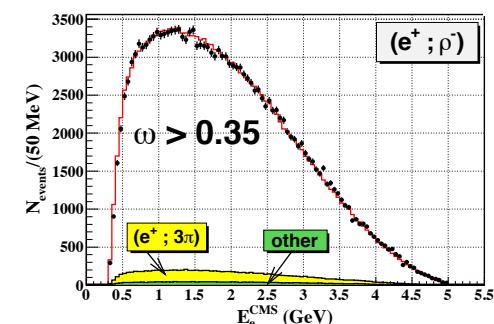
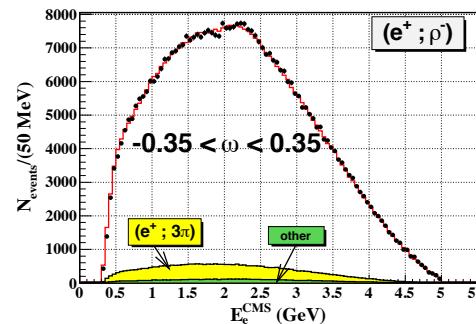
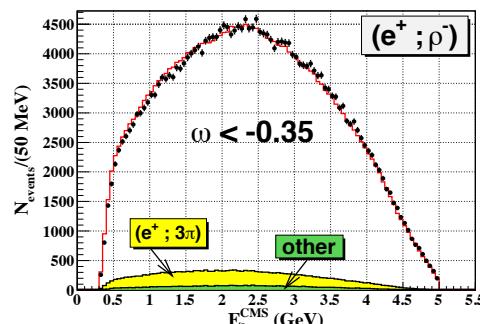
$$\mathcal{P}_{\text{total}} = (1 - \sum_{i=1}^4 \lambda_i) \mathcal{P}_{\text{signal}}^{\ell-\rho} + \lambda_1 \mathcal{P}_{\text{bg}}^{\ell-3\pi} + \lambda_2 \mathcal{P}_{\text{bg}}^{\pi-\rho} + \lambda_3 \mathcal{P}_{\text{bg}}^{\rho-\rho} + \lambda_4 \mathcal{P}_{\text{bg}}^{\text{other}} (\text{MC})$$

MP are extracted in the unbinned maximum likelihood fit of $(\ell\nu\nu; \rho\nu)$ events in the 9D phase space $\vec{z} = (p_\ell, \cos \theta_\ell, \phi_\ell, p_\rho, \cos \theta_\rho, \phi_\rho, m_{\pi\pi}^2, \cos \tilde{\theta}_\pi, \tilde{\phi}_\pi)$ in CMS.



Data fits and systematic uncertainties

$$\text{Helicity sensitive variable } \omega = \frac{1}{\Phi_2 - \Phi_1} \int_{\Phi_1}^{\Phi_2} (\vec{H}_{\rho^\pm}, \vec{n}_{\tau^\pm}) d\Phi = <(\vec{H}_{\rho^\pm}, \vec{n}_{\tau^\pm})>_{\Phi_\tau}$$



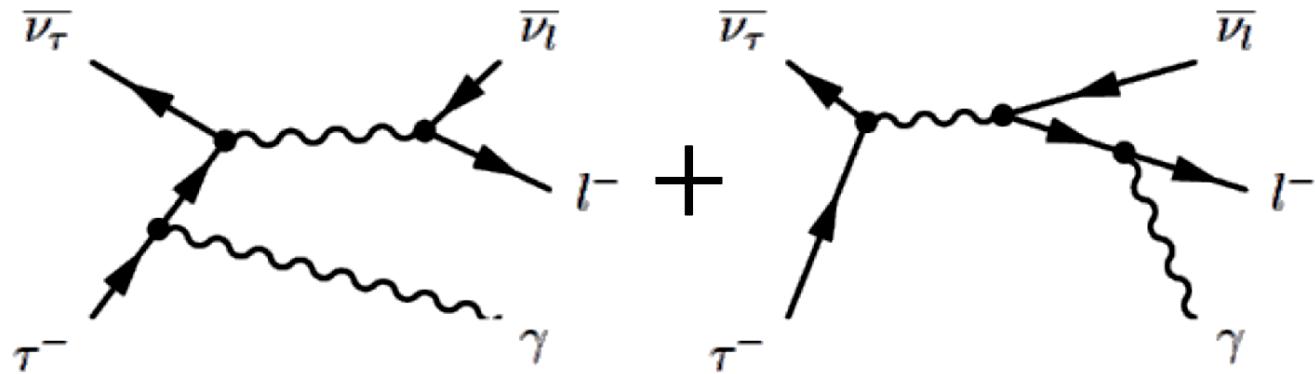
Spin-spin correlation manifests itself through momentum-momentum correlations of final lepton and pions.

Source	$\Delta(\rho)$, %	$\Delta(\eta)$, %	$\Delta(\xi_\rho \xi)$, %	$\Delta(\xi_\rho \xi \delta)$, %
Physical corrections				
ISR+ $\mathcal{O}(\alpha^3)$	0.10	0.30	0.20	0.15
$\tau \rightarrow \ell \nu \nu \gamma$	0.03	0.10	0.09	0.08
$\tau \rightarrow \rho \nu \gamma$	0.06	0.16	0.11	0.02
Background	0.20	0.60	0.20	0.20
Apparatus corrections				
Resolution \oplus brems.	0.10	0.33	0.11	0.19
$\sigma(E_{\text{beam}})$	0.07	0.25	0.03	0.15
Normalization				
$\Delta \mathcal{N}$	0.11	0.50	0.17	0.13
without Data/MC corr.	0.29	0.95	0.38	0.38
trigger eff. corr.	~ 1	~ 2	~ 3	~ 3

We are working on the Data/MC efficiency corrections (trigger, ℓ ID, track rec., π^0 rec.).

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Target MPs: $\bar{\eta}, \xi\kappa$

$$\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau \gamma$$

$$\begin{aligned}\bar{\eta} &= |g_{RL}^V|^2 + |g_{LR}^V|^2 + \frac{1}{8} \left(|g_{RL}^S + 2g_{RL}^T|^2 + |g_{LR}^S + 2g_{LR}^T|^2 \right) + 2 \left(|g_{RL}^T|^2 + |g_{LR}^T|^2 \right) \\ \xi\kappa &= |g_{RL}^V|^2 - |g_{LR}^V|^2 + \frac{1}{8} \left(|g_{RL}^S + 2g_{RL}^T|^2 - |g_{LR}^S + 2g_{LR}^T|^2 \right) + 2 \left(|g_{RL}^T|^2 - |g_{LR}^T|^2 \right)\end{aligned}$$

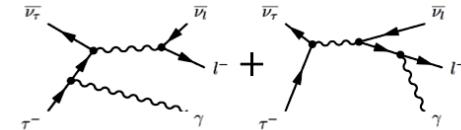
$$\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau \gamma$$

Target MPs:

$$\bar{\eta}, \quad \xi\kappa$$

New Michel parameters ($\bar{\eta}, \xi\kappa$) appear in a radiative decay of tau.

$$\begin{aligned}\bar{\eta} &= |g_{RL}^V|^2 + |g_{LR}^V|^2 + \frac{1}{8} \left(|g_{RL}^S + 2g_{RL}^T|^2 + |g_{LR}^S + 2g_{LR}^T|^2 \right) + 2 \left(|g_{RL}^T|^2 + |g_{LR}^T|^2 \right) \\ \xi\kappa &= |g_{RL}^V|^2 - |g_{LR}^V|^2 + \frac{1}{8} \left(|g_{RL}^S + 2g_{RL}^T|^2 - |g_{LR}^S + 2g_{LR}^T|^2 \right) + 2 \left(|g_{RL}^T|^2 - |g_{LR}^T|^2 \right)\end{aligned}$$



The purpose of this study is to measure the values of $\bar{\eta}, \xi\kappa$

MP	ρ	η	$\xi\delta$	ξ_h	$\bar{\eta}$	$\xi\kappa$
SM	0.75	0	0.75	1	0	0
EX	0.747 ± 0.010	0.013 ± 0.020	0.746 ± 0.021	0.995 ± 0.007	not measured yet	

PDG Chin. Phys. C38, 090001 (2014).

□ Observation of γ is equivalent to the measurement of polarization of the daughter lepton:

- coupling of τ with the right handed daughter lepton $\propto 1 - \xi'$
- $\xi' = -\xi - 4\xi\kappa + 8\xi\delta/3$



Final piece to reveal
the $V - A$ structure!

□ $\bar{\eta}$ gives a constraint on each term

$$\bar{\eta} = |g_{RL}^V|^2 + |g_{LR}^V|^2 + \frac{1}{8} \left(|g_{RL}^S + 2g_{RL}^T|^2 + |g_{LR}^S + 2g_{LR}^T|^2 \right) + 2 \left(|g_{RL}^T|^2 + |g_{LR}^T|^2 \right)$$

$\tau \rightarrow e\nu_e\nu_\tau$		
$ g_{RR}^S < 0.70$	$ g_{RR}^V < 0.17$	$ g_{RR}^T \equiv 0$
$ g_{LR}^S < 0.99$	$ g_{LR}^V < 0.13$	$ g_{LR}^T < 0.082$
$ g_{RL}^S < 2.01$	$ g_{RL}^V < 0.52$	$ g_{RL}^T < 0.51$
$ g_{LL}^S < 2.01$	$ g_{LL}^V < 1.005$	$ g_{LL}^T \equiv 0$
$\tau \rightarrow \mu\nu_\mu\nu_\tau$		
$ g_{RR}^S < 0.72$	$ g_{RR}^V < 0.18$	$ g_{RR}^T \equiv 0$
$ g_{LR}^S < 0.95$	$ g_{LR}^V < 0.12$	$ g_{LR}^T < 0.079$
$ g_{RL}^S < 2.01$	$ g_{RL}^V < 0.52$	$ g_{RL}^T < 0.51$
$ g_{LL}^S < 2.01$	$ g_{LL}^V < 1.005$	$ g_{LL}^T \equiv 0$

PDG Tau decav parameters

$$\left. \begin{array}{l} |g_{RL}^V| < 0.52 \\ |g_{RL}^T| < 0.51 \\ |g_{RL}^S| < 2.01 \\ |g_{LR}^S| < 0.95 \end{array} \right\}$$

These values are not well known yet

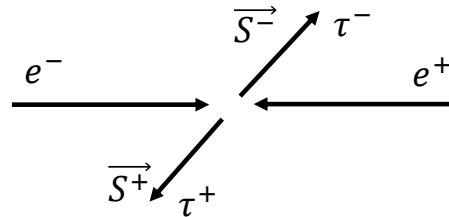


Further constraint on NP models

Reference: N. Shimizu's slide@TAU2016

Method

□ $d\Gamma(\tau^- \rightarrow l^-\bar{\nu}\nu\gamma) \propto A^- - \vec{B}^- \cdot \vec{S}^-$



$$d\Gamma = \frac{n_V + \left(1 - \frac{4}{3}\rho\right)(2n_S + n_V - n_T) + \bar{\eta}(2n_S - 2n_V + n_T)}{-\vec{S}^- \cdot \vec{\xi} \sum_{k=l,\gamma} \vec{e}^k} \left\{ n_V^k - \frac{1}{3} \left(1 - \frac{4}{3}\delta\right)(2n_S^k + 5n_V^k - n_T^k) + \xi_k(2n_S^k - 2n_V^k + n_T^k) \right\} \vec{B}^-$$

where \vec{e}^k is direction of $k = l, \gamma$ and n_*^k are known function, whose arguments are experimentally observable.

□ Measurement of $\xi\kappa$ requires information of τ 's spin

□ We use correlation of the $\tau\tau$ pairs

- The cross section $\frac{d\sigma}{d\Omega_\tau}$ under the definite direction of spin (\vec{S}^-, \vec{S}^+) is expressed as $\frac{d\sigma}{d\Omega_\tau} = D_0 + D_{ij} S_i^- S_j^+$: D_0 spin-independent, D_{ij} spin-dependent coefficients.

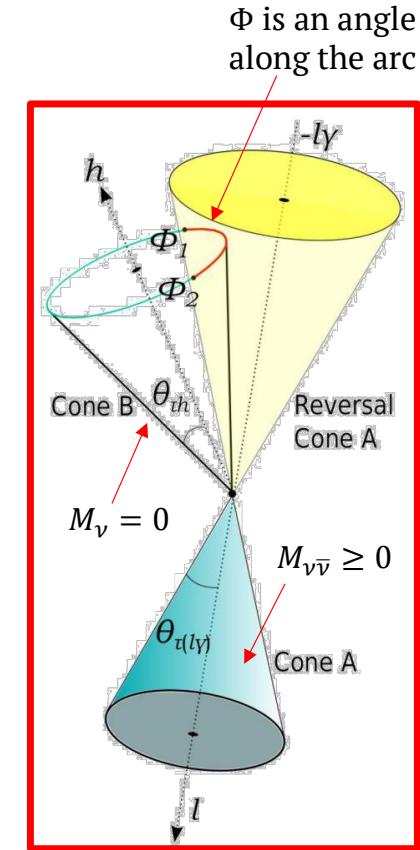
□ We use $\tau^+ \rightarrow \rho^+(\rightarrow \pi^+\pi^0)\bar{\nu}$ decay as a spin analyzer.

- $d\Gamma(\tau^+ \rightarrow \rho^+\nu) = A^+ - \vec{B}^+ \cdot \vec{S}^+$



$$d\sigma(\tau^-\tau^+ \rightarrow (l^-\bar{\nu}\nu\gamma)(\rho^+\bar{\nu})) \propto D_0 A^- A^+ + D_{ij} B_i^- B_j^+$$

PDF of signal



□ The Michel Parameters are fitted using the likelihood function: $\mathcal{L}(\bar{\eta}, \xi\kappa) = \sum \log(PDF)$

- "bare" PDF = $\frac{\frac{d\sigma}{dE_l^* d\Omega_l^* dE_\gamma^* d\Omega_\gamma^* d\Omega_\rho^* dm_\rho^2 d\tilde{\Omega}_\pi d\Omega_\tau}}{\tau^- \text{ side}} \frac{\frac{d\sigma}{dP_l d\Omega_l dP_\gamma d\Omega_\gamma dP_\rho d\Omega_\rho dm_\rho^2 d\tilde{\Omega}_\pi d\Phi}}{\tau^+ \text{ side}}$

Convert to CMS

$$J = \left| \frac{\partial(E_l^*, \Omega_l^*)}{\partial(P_l, \Omega_l)} \right| \cdot \left| \frac{\partial(E_\gamma^*, \Omega_\gamma^*)}{\partial(P_\gamma, \Omega_\gamma)} \right| \cdot \left| \frac{\partial(\Omega_\rho^*, \Omega_\tau)}{\partial(P_\rho, \Omega_\rho, \Phi)} \right|$$

$$\frac{d\sigma}{dP_l d\Omega_l dP_\gamma d\Omega_\gamma dP_\rho d\Omega_\rho dm_\rho^2 d\tilde{\Omega}_\pi d\Phi}$$

integrate out

Method

PDF of signal

$$S(\vec{x}) = \frac{d\sigma}{dP_l d\Omega_l dP_\gamma d\Omega_\gamma dP_\rho d\Omega_\rho dm_\rho^2 d\tilde{\Omega}_\pi} = \int_{\Phi_1}^{\Phi_2} \textcolor{red}{d\Phi} \frac{d\sigma}{dP_l d\Omega_l dP_\gamma d\Omega_\gamma dP_\rho d\Omega_\rho dm_\rho^2 d\tilde{\Omega}_\pi \textcolor{red}{d\Phi}}$$

- $S(\vec{x})$: the visible PDF for observable \vec{x} ($N_{\text{dim.}} = 12$)
- $S(\vec{x})$ has a form: $\vec{x} = \{P_l, \Omega_l, P_\gamma, \Omega_\gamma, P_\rho, \Omega_\rho, m_{\pi\pi}^2, \tilde{\Omega}_\pi\}$

□ Event selection & existence of BG → the visible PDF is formulated as sum of signal and BGs

$$\bullet P_{tot}(\vec{x}) = (1 - \sum \lambda_i) \cdot \frac{S(\vec{x}) \varepsilon(\vec{x})}{\int d\vec{x} S(\vec{x}) \varepsilon(\vec{x})} + \sum \lambda_i \frac{B_i(\vec{x}) \varepsilon(\vec{x})}{\int d\vec{x} B_i(\vec{x}) \varepsilon(\vec{x})}$$

i : index of background

$B_i(\vec{x})$: PDF of i -th background

$S(\vec{x})$: PDF of signal

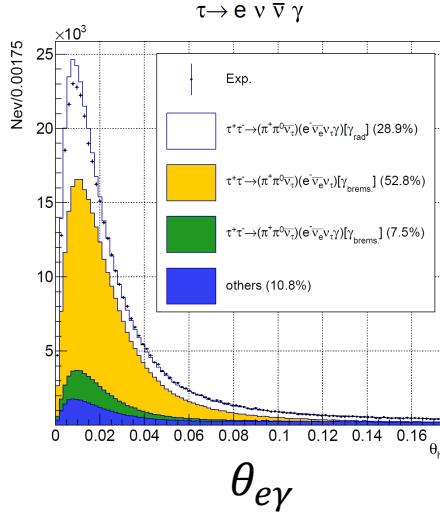
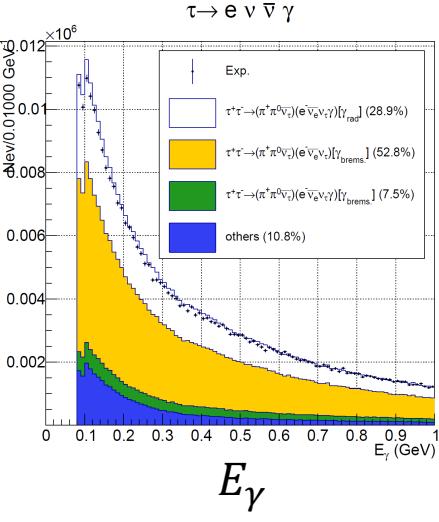
λ_i : fraction of i -th background component

$\varepsilon(\vec{x})$: selection efficiency

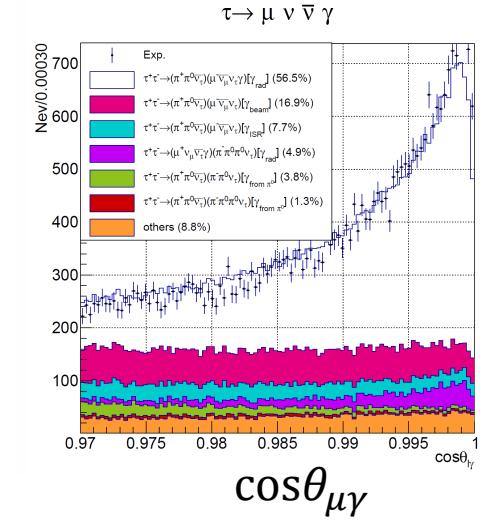
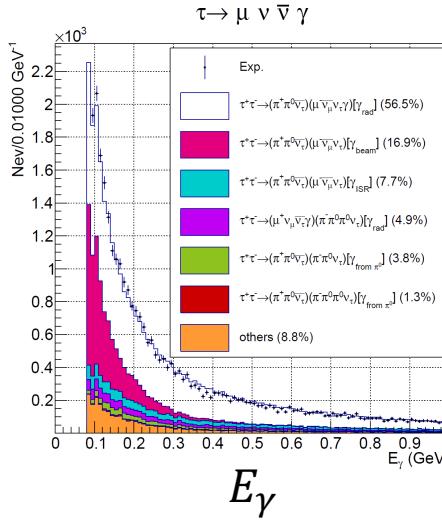
Then Michel parameters are fitted by likelihood function:

$$\mathcal{L}(\bar{\eta}, \xi \kappa) = \sum \log(PDF)$$

$\tau \rightarrow e v \bar{v} \gamma$ candidates



$\tau \rightarrow \mu v \bar{v} \gamma$ candidates



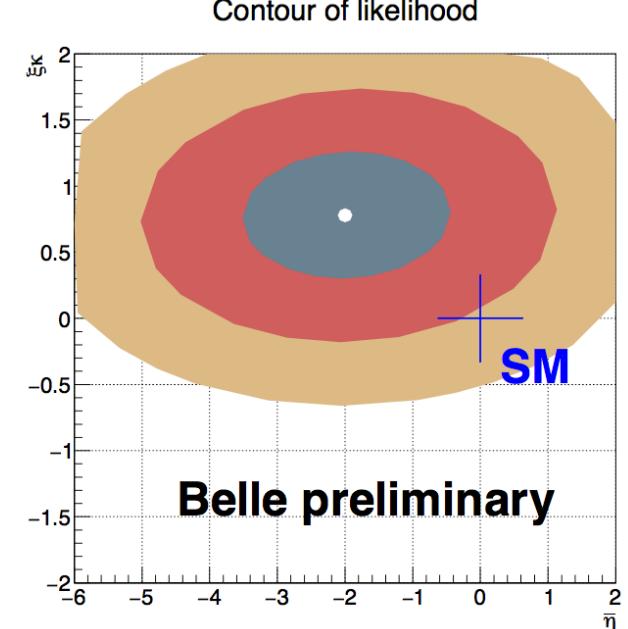
Result

Source	$\sigma_{\bar{\eta}}^e$	$\sigma_{\xi\kappa}^e$	$\sigma_{\bar{\eta}}^\mu$	$\sigma_{\xi\kappa}^\mu$
Normalization	4.3	0.94	0.15	0.04
Background PDF	2.5	0.24	0.67	0.22
Branching ratios	3.8	0.05	0.25	0.01
Cluster merge in ECL	2.2	0.46	0.02	0.06
Detector resolution	0.74	0.20	0.22	0.02
Data/MC eff. corr.	1.9	0.14	0.04	0.04
Total	7.0	1.1	0.76	0.24

Belle preliminary

$$\bar{\eta} = -1.3 \pm 1.5 \pm 0.8$$

$$\xi\kappa = 0.5 \pm 0.4 \pm 0.2$$



Reference: N. Shimizu's slide@TAU2016

Reference: D. Epifanov's slide@PhiPsi-2017

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 - $\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau$
 - $\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau \gamma$
 - $\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau l'^+ l'^-$
- Conclusion & Future Plan

Target MPs:

$$Q_{LL}, Q_{RL}, Q_{LR}, Q_{RR}, B_{RL}, B_{LR}, I_\alpha, I_\beta$$

$$\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau l'^+ l'^-$$

$$Q_{LL} = \frac{1}{4} |g_{LL}^S|^2 + |g_{LL}^V|^2$$

$$Q_{RL} = \frac{1}{4} |g_{RL}^S|^2 + |g_{RL}^V|^2 + |g_{RL}^T|^2$$

$$Q_{LR} = \frac{1}{4} |g_{LR}^S|^2 + |g_{LR}^V|^2 + |g_{LR}^T|^2$$

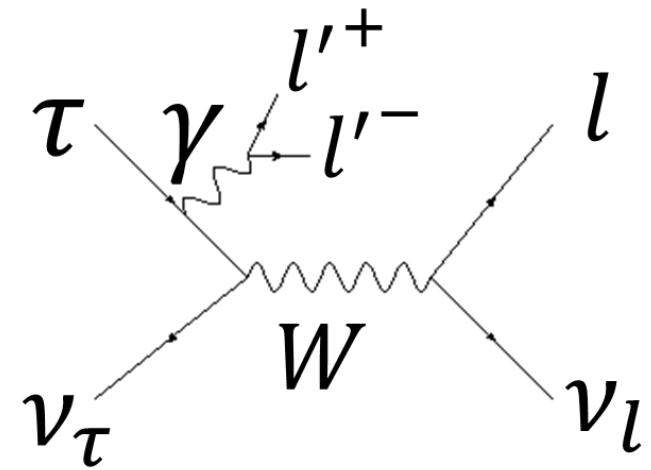
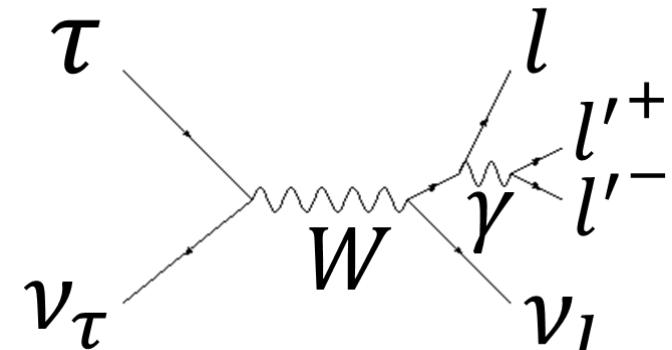
$$Q_{RR} = \frac{1}{4} |g_{RR}^S|^2 + |g_{RR}^V|^2$$

$$B_{RL} = \frac{1}{16} |g_{RL}^S + 6g_{RL}^T|^2 + |g_{RL}^V|^2$$

$$B_{LR} = \frac{1}{16} |g_{LR}^S + 6g_{LR}^T|^2 + |g_{LR}^V|^2$$

$$I_\alpha = \frac{1}{4} g_{LR}^V (g_{RL}^S + 6g_{RL}^T)^* + \frac{1}{4} g_{RL}^{V*} (g_{LR}^S + 6g_{LR}^T)$$

$$I_\beta = g_{LL}^V g_{RR}^{S*}/2 + g_{RR}^{V*} g_{LL}^S/2$$

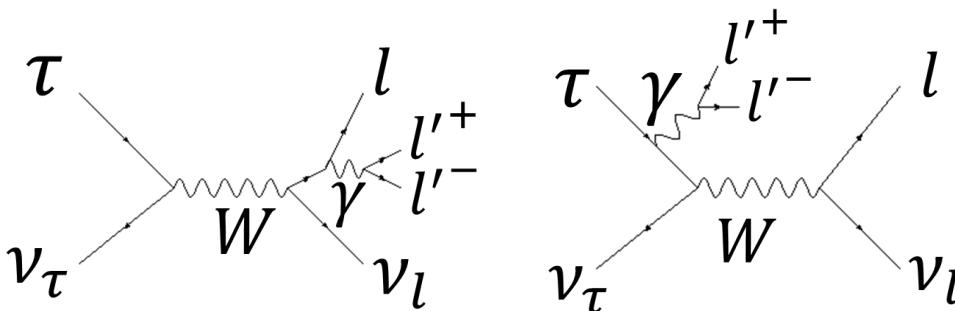


$$\tau^\pm \rightarrow l^\pm l'^+ l'^- \nu_\tau \nu_l \quad (l, l' = e, \mu)$$

Target MPs:

$$Q_{LL}, Q_{RL}, Q_{LR}, Q_{RR}, B_{RL}, B_{LR}, I_\alpha, I_\beta$$

We measure the Michel-like parameters through the measurement of branching fraction of five-body leptonic decays of tau



Prediction from theory	
Channel	
$\text{BR}(\tau^- \rightarrow e^- e^+ e^- \bar{\nu}_e \nu_e) \times 10^5$	4.21 ± 0.01
$\text{BR}(\tau^- \rightarrow e^- \mu^+ \mu^- \bar{\nu}_e \nu_e) \times 10^7$	1.247 ± 0.001
$\text{BR}(\tau^- \rightarrow \mu^- e^+ e^- \bar{\nu}_\mu \nu_\tau) \times 10^5$	1.984 ± 0.004
$\text{BR}(\tau^- \rightarrow \mu^- \mu^+ \mu^- \bar{\nu}_\mu \nu_\tau) \times 10^7$	1.183 ± 0.001

Ref. [JHEP 1604, 185 (2016)]

$$Q_{LL} = \frac{1}{4} |g_{LL}^S|^2 + |g_{LL}^V|^2$$

$$Q_{RL} = \frac{1}{4} |g_{RL}^S|^2 + |g_{RL}^V|^2 + |g_{RL}^T|^2$$

$$Q_{LR} = \frac{1}{4} |g_{LR}^S|^2 + |g_{LR}^V|^2 + |g_{LR}^T|^2$$

$$Q_{RR} = \frac{1}{4} |g_{RR}^S|^2 + |g_{RR}^V|^2$$

$$B_{RL} = \frac{1}{16} |g_{RL}^S + 6g_{RL}^T|^2 + |g_{RL}^V|^2$$

$$B_{LR} = \frac{1}{16} |g_{LR}^S + 6g_{LR}^T|^2 + |g_{LR}^V|^2$$

$$I_\alpha = \frac{1}{4} g_{LR}^V (g_{RL}^S + 6g_{RL}^T)^* + \frac{1}{4} g_{RL}^{V*} (g_{LR}^S + 6g_{LR}^T)$$

$$I_\beta = g_{LL}^V g_{RR}^{S*}/2 + g_{RR}^{V*} g_{LL}^S/2$$

Reference:

W. Fetscher, H. J. Gerber and K. F. Johnson,
Phys. Lett. B 173, 102 (1986)

$$\begin{array}{ccc} l^\pm & l'^+ & l'^- \\ x_1 = \frac{2E_1}{M_\tau} & x_2 = \frac{2E_2}{M_\tau} & x_3 = \frac{2E_3}{M_\tau} \end{array}$$

$$\frac{d\Gamma_5}{dx_1 d\Omega_1 dx_2 d\Omega_2 dx_3 d\Omega_3} = \frac{M^2 |\vec{p}_1| |\vec{p}_2| |\vec{p}_3|}{3 \cdot 2^{21} \pi^{10}} \underline{\mathcal{T}_{\alpha\beta}^s I^{\alpha\beta}(P)}$$

$$\mathcal{T}_{\alpha\beta}^s I^{\alpha\beta}(P) = e^4 |G_{\ell\ell'}|^2 \left[\left(Q_{LL} T_{LL}^Q + Q_{RL} T_{RL}^Q + B_{RL} T_{RL}^B + L \leftrightarrow R \right) + \Re e \left(I_\alpha T_\alpha^I + I_\beta T_\beta^I \right) \right]$$

($T_{LL}^Q, T_{RL}^Q, T_{LR}^Q, T_{RR}^Q, T_{RL}^B, T_{LR}^B, T_\alpha^I$, and T_β^I are the function of kinematic variables)

Strategy of Measurement of MPs in $\tau^\pm \rightarrow l^\pm l'^+ l'^- \nu_\tau \nu_l$

- We do not use a kinematical fitting due to a lack of statistics
- We only do the measurement of branching fraction of $\tau^\pm \rightarrow l^\pm l'^+ l'^- \nu_\tau \nu_l$
- We use a discrepancy of branching fraction between MC and data

$$\frac{d\Gamma_5}{dx_1 d\Omega_1 dx_2 d\Omega_2 dx_3 d\Omega_3} = \frac{M^2 |\vec{\mathbf{p}}_1| |\vec{\mathbf{p}}_2| |\vec{\mathbf{p}}_3|}{3 \cdot 2^{21} \pi^{10}} \underline{\mathcal{T}_{\alpha\beta}^s I^{\alpha\beta}(P)}$$

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Strategy of Measurement of MPs in $\tau^\pm \rightarrow l^\pm l'^+ l'^- \nu_\tau \nu_l$

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Integrate a differential decay width

$$BR_{\text{exp}} = BR_{\text{SM}} [Q_{LL} + bQ_{LR} + cB_{LR} + Q_{RR} + dQ_{RL} + eB_{RL} + \Re(fI_\alpha + gI_\beta)]$$

$\Delta b \sim g$ are the integrated value of the terms related to T_{RL}^Q , T_{LR}^Q , T_{RR}^Q , T_{RL}^B , T_{LR}^B , T_α^I , and T_β^I

Strategy of Measurement of MPs in $\tau^\pm \rightarrow l^\pm l'^+ l'^- \nu_\tau \nu_l$

- We do not use a kinematical fitting due to a lack of statistics
- We only do the measurement of branching fraction of $\tau^\pm \rightarrow l^\pm l'^+ l'^- \nu_\tau \nu_l$
- We use a discrepancy of branching fraction between MC and data

$$\frac{d\Gamma_5}{dx_1 d\Omega_1 dx_2 d\Omega_2 dx_3 d\Omega_3} = \frac{M^2 |\vec{p}_1| |\vec{p}_2| |\vec{p}_3|}{3 \cdot 2^{21} \pi^{10}} \overline{\mathcal{T}_{\alpha\beta}^s I^{\alpha\beta}(P)}$$

\downarrow

$$\mathcal{T}_{\alpha\beta}^s I^{\alpha\beta}(P) = e^4 |G_{\ell\ell'}|^2 \left[\left(Q_{LL} T_{LL}^Q + Q_{RL} T_{RL}^Q + B_{RL} T_{RL}^B + L \leftrightarrow R \right) + \Re e \left(I_\alpha T_\alpha^I + I_\beta T_\beta^I \right) \right]$$

Integrate a differential decay width

$$BR_{\text{exp}} = BR_{\text{SM}} [Q_{LL} + bQ_{LR} + cB_{LR} + Q_{RR} + dQ_{RL} + eB_{RL} + \Re(fI_\alpha + gI_\beta)]$$

$\Delta b \sim g$ are the integrated value of the terms related to $T_{RL}^Q, T_{LR}^Q, T_{RR}^Q, T_{RL}^B, T_{LR}^B, T_\alpha^I$, and T_β^I

Assume the discrepancy Δ

$$BR_{\tau^\pm \rightarrow l^\pm l'^+ l'^- \nu_\tau \nu_l}^{\text{Measured}} = BR_{\tau^\pm \rightarrow l^\pm l'^+ l'^- \nu_\tau \nu_l}^{\text{SM predicted}} + \Delta$$

We assume that, the discrepancy Δ is brought by only one term in BR_{exp}

We give constraints to Michel parameters by the discrepancy Δ :

$$Q_{RL} < \Delta / (BR_{\text{SM}} b), \quad B_{RL} < \Delta / (BR_{\text{SM}} c), \quad Q_{RR} < \Delta / (BR_{\text{SM}}), \quad Q_{LR} < \Delta / (BR_{\text{SM}} d),$$

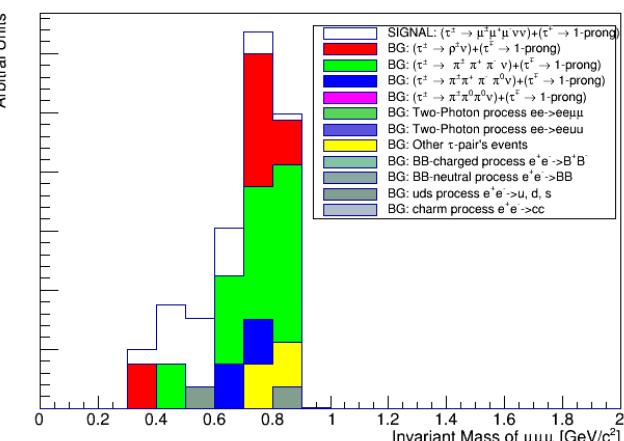
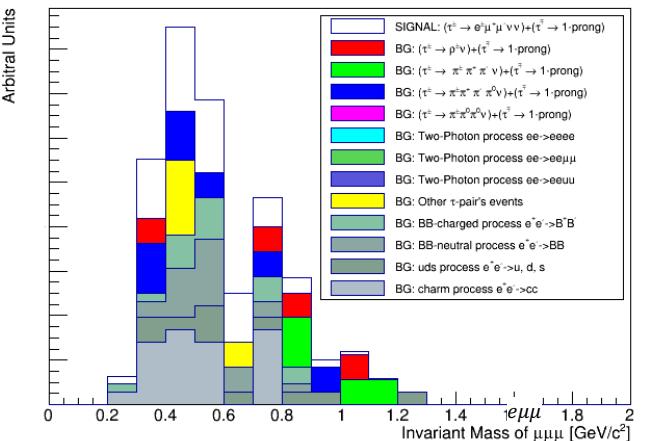
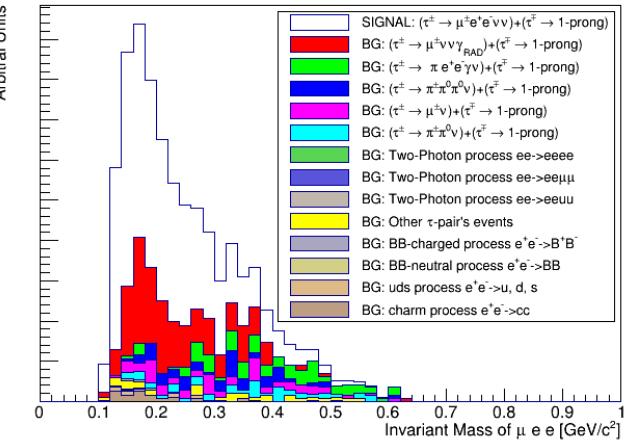
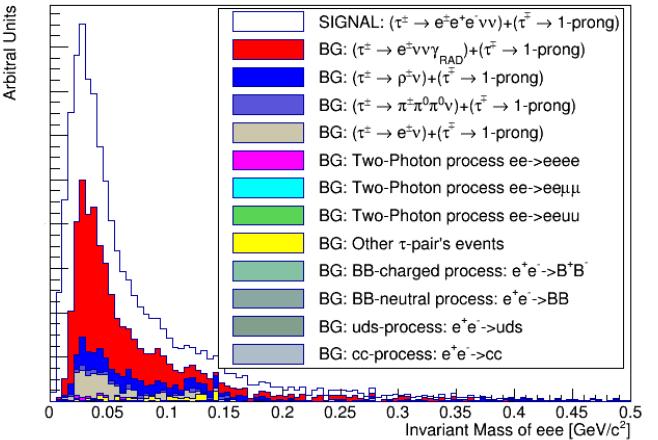
$$B_{LR} < \Delta / (BR_{\text{SM}} e), \quad I_\alpha < \Delta / (BR_{\text{SM}} f), \quad I_\beta < \Delta / (BR_{\text{SM}} g),$$

Result of MC study:

Monte Carlo (MC) sample of 4 million signal decays was used.

Histograms shown below are plotted with assuming the SM-predicted branching ratio.

τ^- decay mode	$e^-e^+\bar{\nu}_e\nu_\tau$	$\mu^-\bar{\nu}_\mu\nu_\tau$	$e^-\mu^+\bar{\nu}_e\nu_\tau$	$\mu^-\mu^+\bar{\nu}_\mu\nu_\tau$
Detection efficiency	(1.769±0.004)%	(1.204±0.003)%	(3.561±0.006)%	(1.674±0.004)%
Main backgrounds	$e^-\bar{\nu}_e\nu_\tau\gamma$ $\rightarrow e^-\bar{\nu}_e\nu_\tau(e^+e^-)$ $\pi^-\pi^0\nu_\tau$ $\rightarrow \pi^-(\gamma\gamma)\nu_\tau$ $\rightarrow \pi^-((e^+e^-)\gamma)\nu_\tau$ (mis-ID π as e)	$\mu^-\bar{\nu}_\mu\nu_\tau\gamma$ $\rightarrow \mu^-\bar{\nu}_\mu\nu_\tau(e^+e^-)$ $\pi^-\pi^0\nu_\tau$ $\rightarrow \pi^-(e^+e^-\gamma)\nu_\tau$ $\pi^-\pi^0\nu_\tau$ $\rightarrow \pi^-(\gamma\gamma)(\gamma\gamma)\nu_\tau$ $\rightarrow \pi^-((e^+e^-)\gamma)(\gamma\gamma)\nu_\tau$ (mis-ID π as μ)	$\pi^-\pi^0\nu_\tau$ $\rightarrow \pi^-(\gamma\gamma)\nu_\tau$ $\rightarrow \pi^-((e^+e^-)\gamma)\nu_\tau$ $\pi^-\pi^+\pi^-\nu_\tau$ (mis-ID π as μ, e)	$\pi^-\pi^0\nu_\tau$ $\rightarrow \pi^-(\gamma\gamma)\nu_\tau$ $\rightarrow \pi^-((e^+e^-)\gamma)\nu_\tau$ $\pi^-\pi^+\pi^-\nu_\tau$ (mis-ID π as μ)
Expected number of signal events	1300	430	8	4
Fraction of the signal	47%	50%	22%	16%



$$\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau l'^+ l'^-$$

Preliminary Estimation of systematic uncertainties (for branching fraction)

$$\tau^\pm \rightarrow e^\pm e^+ e^- \nu_\tau \nu_e$$

about 10%

Dominant errors:

- Statistical error
- Particle Identification's error

$$\tau^\pm \rightarrow \mu^\pm e^+ e^- \nu_\tau \nu_\mu$$

about 15%

Dominant errors:

- Statistical error
- Particle Identification's error

$$\tau^\pm \rightarrow e^\pm \mu^+ \mu^- \nu_\tau \nu_e$$

about 70%

Dominant errors:

- Statistical error

$$\tau^\pm \rightarrow \mu^\pm \mu^+ \mu^- \nu_\tau \nu_\mu$$

about 70%

Dominant errors:

- Statistical error

The data will be opened soon and the preliminary result of branching fractions will be finalized

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Conclusion & Future Plan

- To verify the SM's prediction that the Lorentz structure of the charged weak current has a V-A coupling, we study the Lorentz structure.
- For the study of Lorentz structure, a measurement of Michel parameters is ongoing for three modes below at Belle.
 - $\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau$ (Target MPs: ρ, η, ξ, δ)
 - $\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau \gamma$ (Target MPs: $\bar{\eta}, \xi\kappa$)
 - $\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau l'^+ l'^-$ (Target MPs: $Q_{LL}, Q_{RL}, Q_{LR}, Q_{RR}, B_{RL}, B_{LR}, I_\alpha, I_\beta$)
- Belle has a good environment for the measurement of MPs of tau-lepton.
 - Low background & Energy of center of mass frame is known
 - Pair creation ($B\bar{B}$, $\tau^+\tau^-$, etc). One-side can be used for tagging.
 - Large number of tau-pairs ($N_{\tau\tau} \sim 9.0 \times 10^8$)
- For $\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau$ and $\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau \gamma$, the result of $\rho, \eta, \xi, \delta, \bar{\eta}, \xi\kappa$ will be published.
- For $\tau^- \rightarrow l^- \bar{\nu}_l \nu_\tau l'^+ l'^-$, we try to give the constraints to MPs ($Q_{LL}, Q_{RL}, Q_{LR}, Q_{RR}, B_{RL}, B_{LR}, I_\alpha, I_\beta$) from branching fractions and the estimation of its systematic errors is ongoing and finalized soon then open data.

Thank you!

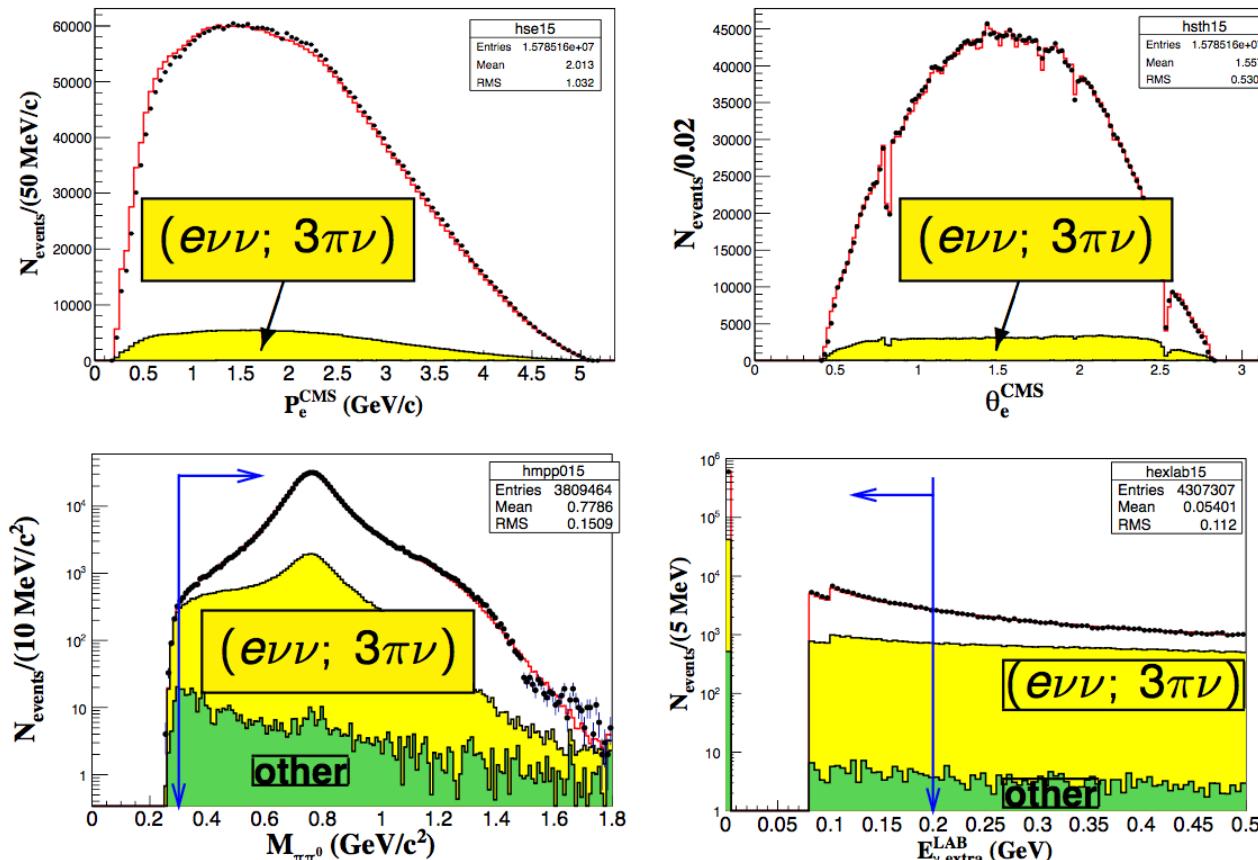
Backup

Backup $(\tau^- \rightarrow l^-\bar{\nu}_l\nu_\tau)$

Selection criteria

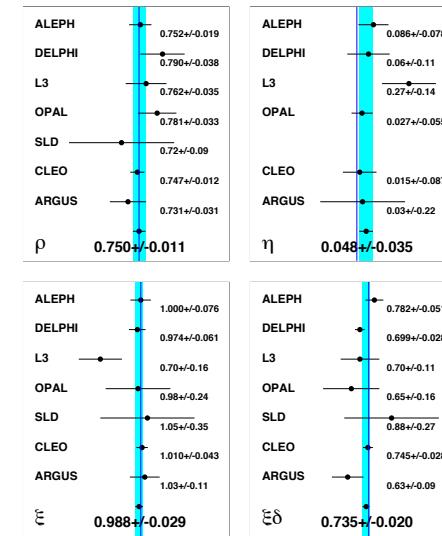
- After the standard preselections we take events with two oppositely charged tracks, one of them is identified as lepton ($eID, \mu ID > 0.9$) and the other one as pion ($PID(\pi/K) > 0.4$).
- π^0 candidate is reconstructed from the pair of gammas ($E_{\gamma}^{LAB} > 80$ MeV) satisfying $115 \text{ MeV}/c^2 < M_{\gamma\gamma} < 150 \text{ MeV}/c^2, P_{\pi^0}^{CMS} > 0.3 \text{ GeV}/c$.
- $\cos(\vec{P}_{\text{lep}}, \vec{P}_{\pi}) < 0, \cos(\vec{P}_{\text{lep}}, \vec{P}_{\pi^0}) < 0, 0.3 \text{ GeV}/c^2 < M_{\pi\pi^0} < 1.8 \text{ GeV}/c^2$.
- $E_{\text{rest}\gamma}^{LAB} < 0.2 \text{ GeV}$

Detection efficiency $\varepsilon_{\text{det}} \simeq 12\%$



Status of Michel parameters in τ decays

Michel par.	Measured value	Experiment	SM value
ρ	$0.747 \pm 0.010 \pm 0.006$	CLEO-97	$3/4$
(e or μ)	1.2%		
η	$0.012 \pm 0.026 \pm 0.004$	ALEPH-01	0
(e or μ)	2.6%		
ξ	$1.007 \pm 0.040 \pm 0.015$	CLEO-97	1
(e or μ)	4.3%		
$\xi\delta$	$0.745 \pm 0.026 \pm 0.009$	CLEO-97	$3/4$
(e or μ)	2.8%		
ξ_h	$0.992 \pm 0.007 \pm 0.008$	ALEPH-01	1
(all hadr.)	1.1%		



With $\times 300$ Belle statistics we can improve MP uncertainties by one order of magnitude
In BSM models the couplings to τ are expected to be enhanced in comparison with μ .

Also contribution from New Physics in τ decays can be amplified by $(\frac{m_\tau}{m_\mu})^n$.

- **Type II 2HDM:** $\eta_\mu(\tau) = \frac{m_\mu M_\tau}{2} \left(\frac{\tan^2 \beta}{M_{H^\pm}^2} \right)^2$; $\frac{\eta_\mu(\tau)}{\eta_e(\mu)} = \frac{M_\tau}{m_e} \approx 3500$
- **Tensor interaction:** $\mathcal{L} = \frac{g}{2\sqrt{2}} W^\mu \left\{ \bar{\nu} \gamma_\mu (1 - \gamma^5) \tau + \frac{\kappa_\tau^W}{2m_\tau} \partial^\nu \left(\bar{\nu} \sigma_\mu n u (1 - \gamma^5) \tau \right) \right\}$,
 $-0.096 < \kappa_\tau^W < 0.037$: DELPHI Abreu EPJ C16 (2000) 229.
- **Unparticles:** Moyotl PRD 84 (2011) 073010, Choudhury PLB 658 (2008) 148.
- **Lorentz and CPTV:** Hollenberg PLB 701 (2011) 89
- **Heavy Majorana neutrino:** M. Doi *et al.*, Prog. Theor. Phys. 118 (2007) 1069.

Corrections, detector effects, background

Physical corrections:

- All $\mathcal{O}(\alpha^3)$ QED and electroweak higher order corrections to $e^+ e^- \rightarrow \tau^+ \tau^- (\gamma)$ are included
- Radiative leptonic decays $\tau^- \rightarrow \ell^- \bar{\nu}_\ell \nu_\tau \gamma$
- Radiative decay $\tau^- \rightarrow \pi^- \pi^0 \nu_\tau \gamma$

Detector effects:

- Track momentum resolution
- γ energy and angular resolution
- Effect of external bremsstrahlung for $e - \rho$ events
- Beam energy spread
- EXP/MC efficiency corrections (trigger, track rec., π^0 rec., ℓ ID, π ID)

Background:

The main background comes from $(\ell \nu \nu; \pi 2\pi^0 \nu)(\sim 10\%)$, $(\pi \nu; \pi \pi^0 \nu)(\sim 1.5\%)$ and $(\rho^+ \nu; \rho^- \nu)(\sim 0.5\%)$ events, it is included in PDF analytically. The remaining background($\sim 2.0\%$) is taken into account using MC-based approach.

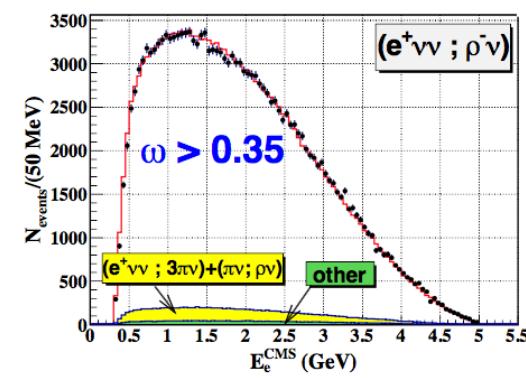
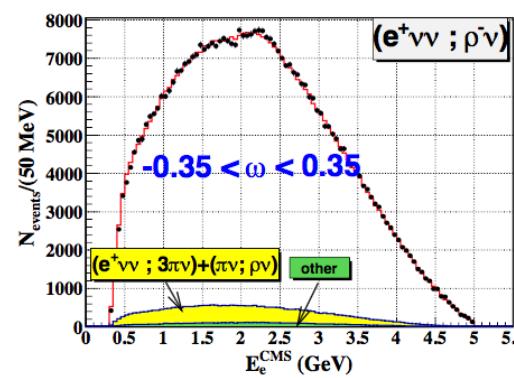
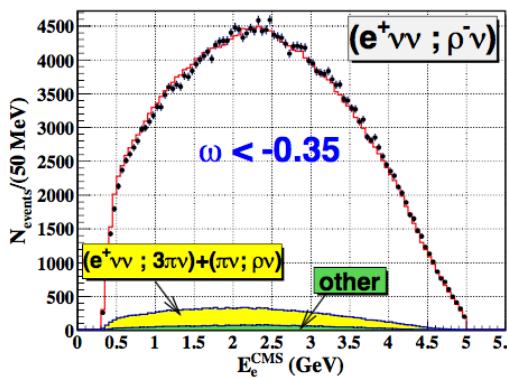
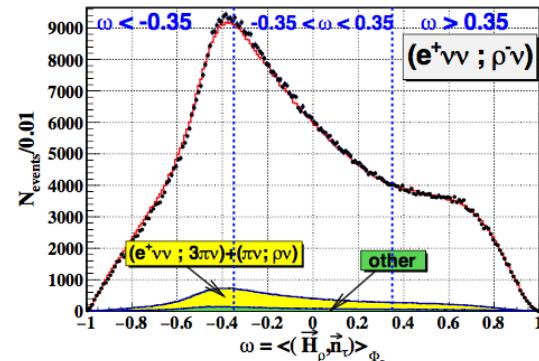
Background from the non- $\tau\tau$ events is $\lesssim 0.1\%$.

Helicity sensitive variable ω

M. Davier et. al Phys. Lett. B 306 (1993) 411.

Helicity sensitive variable ω is introduced as:

$$\omega = \frac{1}{\Phi_2 - \Phi_1} \int_{\Phi_1}^{\Phi_2} (\vec{H}_{\rho^\pm}, \vec{n}_{\tau^\pm}) d\Phi = <(\vec{H}_{\rho^\pm}, \vec{n}_{\tau^\pm})>_{\Phi_\tau}$$



Spin-spin correlation manifests itself through momentum-momentum correlations of final lepton and pions.

Backup $(\tau^- \rightarrow l^-\bar{\nu}_l\nu_\tau\gamma)$

Method

Reference: N. Shimizu's slide@TAU2016

40

PDF of signal

$$S(\vec{x}) = \frac{d\sigma}{dP_l d\Omega_l dP_\gamma d\Omega_\gamma dP_\rho d\Omega_\rho dm_\rho^2 d\tilde{\Omega}_\pi} = \int_{\Phi_1}^{\Phi_2} \frac{d\Phi}{dP_l d\Omega_l dP_\gamma d\Omega_\gamma dP_\rho d\Omega_\rho dm_\rho^2 d\tilde{\Omega}_\pi} \frac{d\sigma}{d\Phi}$$

- $S(\vec{x})$: the visible PDF for observable \vec{x} ($N_{\text{dim.}} = 12$)
- $S(\vec{x})$ has a form: $\vec{x} = \{P_l, \Omega_l, P_\gamma, \Omega_\gamma, P_\rho, \Omega_\rho, m_{\pi\pi}^2, \tilde{\Omega}_\pi\}$

□ Event selection & existence of BG → the visible PDF is formulated as sum of signal and BGs

- $P_{tot}(\vec{x}) = (1 - \sum \lambda_i) \cdot \frac{S(\vec{x})\varepsilon(\vec{x})}{\int d\vec{x} S(\vec{x})\varepsilon(\vec{x})} + \sum \lambda_i \frac{B_i(\vec{x})\varepsilon(\vec{x})}{\int d\vec{x} B_i(\vec{x})\varepsilon(\vec{x})}$

i : index of background

$B_i(\vec{x})$: PDF of i -th background

$S(\vec{x})$: PDF of signal

λ_i : fraction of i -th background component

$\varepsilon(\vec{x})$: selection efficiency

□ Normalization is evaluated by MC

- $S(\vec{x}) = A_0(\vec{x}) + A_{\bar{\eta}}(\vec{x}) \cdot \bar{\eta} + A_{\xi\kappa}(\vec{x}) \cdot \xi\kappa$

$$\sigma_0 = \int d\vec{x} A_0(\vec{x})$$

- $\int d\vec{x} S(\vec{x})\varepsilon(\vec{x}) = \int d\vec{x} \varepsilon(\vec{x}) A_0(\vec{x}) \frac{A_0 + A_{\bar{\eta}}\bar{\eta} + A_{\xi\kappa}\xi\kappa}{A_0} = \frac{\sigma_0 \bar{\varepsilon}}{N_{\text{sel}}} \sum_k \frac{A_0 + A_{\bar{\eta}}\bar{\eta} + A_{\xi\kappa}\xi\kappa}{A_0}$

$$= \frac{\sigma_0 \bar{\varepsilon}}{N_{\text{sel}}} [1 + N_{\bar{\eta}} \cdot \bar{\eta} + N_{\xi\kappa} \cdot \xi\kappa]$$

σ_0 : absolute normalization

$N_{\bar{\eta}}, N_{\xi\kappa}$: relative normalization

Event selection

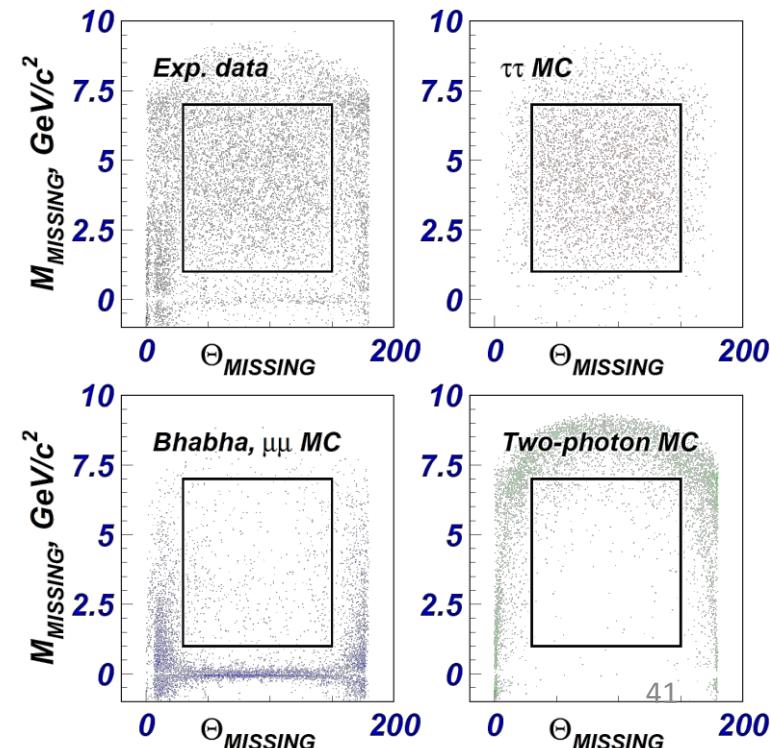
□ We use all data taken with $\Upsilon(4S)$ energy: 703 fb^{-1}

□ Preselection of $\tau\tau$

1. exactly two oppositely charged tracks
 - $dr < 0.5 \text{ cm}$, $|dz| < 2.5 \text{ cm}$, one $P_t > 0.5 \text{ GeV}/c$, the other $P_t > 0.1 \text{ GeV}/c$
2. ECL cluster energy $< 9 \text{ GeV}$
3. opening angle of two tracks $20^\circ < \psi < 175^\circ$
4. $N_\gamma < 5$ for $E_\gamma > 80 \text{ MeV}$
5. $1 \text{ GeV}/c^2 < M_{\text{missing}} < 7 \text{ GeV}/c^2$
6. $30^\circ < \theta_{\text{missing}} < 175^\circ$

In particular, the last two requirements well discriminate other physics processes like

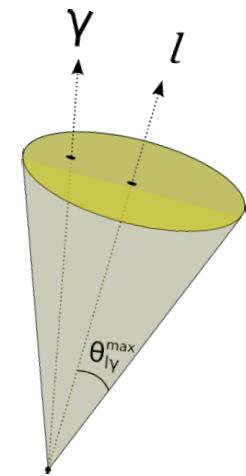
- Bhabha $e^+e^- \rightarrow e^+e^-$
- $e^+e^- \rightarrow \mu^+\mu^-$
- Two photon processes



Event selection

□ Final selection

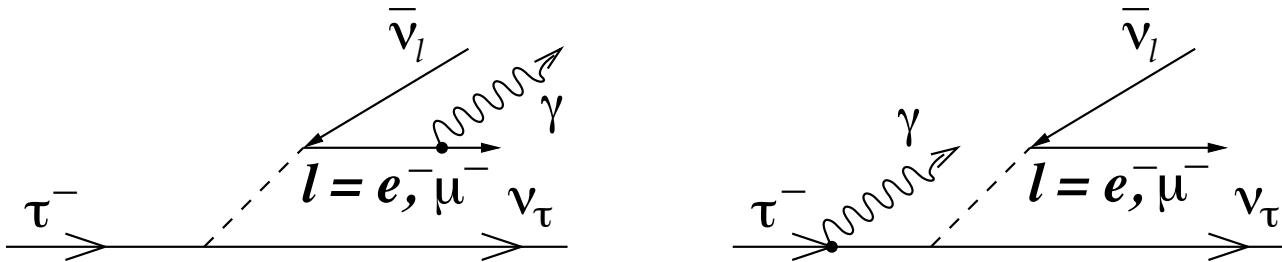
- Electron: $\frac{P_e}{P_e + P_x} > 0.9$, Muon: $\frac{P_\mu}{P_\mu + P_\pi + P_K} > 0.9$
- Pion: $\frac{P_\pi}{P_\pi + P_K} > 0.4$
- π^0 candidate: $115 \text{ MeV}/c^2 < m_{\gamma\gamma} < 150 \text{ MeV}/c^2$ for $E_\gamma > 80 \text{ MeV}$
- ρ candidate: $0.5 \text{ GeV}/c^2 < m_{\pi\pi} < 1.5 \text{ GeV}/c^2$
- $\theta_{\rho(l\gamma)} > 90^\circ$
- $\cos\theta_{e\gamma} > 0.9848$, $\cos\theta_{\mu\gamma} > 0.9700$
- Energy of photons which are not associated with any tracks
 $E_{\text{extray}} < 0.2 \text{ GeV}$ for $\tau \rightarrow e\nu\bar{\nu}\gamma$ and
 $E_{\text{extray}} < 0.3 \text{ GeV}$ for $\tau \rightarrow \mu\nu\bar{\nu}\gamma$



Michel parameters in $\tau \rightarrow \ell \nu \nu \gamma$, ($\ell = e, \mu$) (II)

C. Fronsdal and H. Uberall, Phys. Rev. **113** (1959) 654. ($m_\ell = 0$)

A. B. Arbuzov and T. V. Kopylova, JHEP **1609** (2016) 109. ($m_\ell \neq 0$)



Photon carries information about spin state of outgoing lepton, as a result two additional parameters, $\bar{\eta}$ and $\xi\kappa$, can be extracted.

These parameters were measured in τ decays at Belle for the first time.

$$\frac{d\Gamma(\tau^\mp \rightarrow \ell^\mp \nu_\ell \nu_\tau \gamma)}{dx dy d\Omega_\ell d\Omega_\gamma} = \Gamma_0 \frac{\alpha}{64\pi^3} \frac{\beta_\ell}{y} \left[F(x, y, d) \pm P_\tau (\beta_\ell \cos \theta_\ell G(x, y, d) + \cos \theta_\gamma H(x, y, d)) \right],$$

$$\Gamma_0 = G_F^2 m_\tau^5 / 192\pi^3, \quad \beta_\ell = \sqrt{1 - m_\ell^2/E_\ell^2}, \quad x = 2E_\ell/m_\tau, \quad y = 2E_\gamma/m_\tau, \quad d = 1 - \beta_\ell \cos \theta_{\ell\gamma}$$

$$F = F_0 + \bar{\eta} F_1, \quad G = G_0 + \xi\kappa G_1, \quad H = H_0 + \xi\kappa H_1, \quad \frac{d\sigma(\ell^\mp \nu_\nu \gamma, \rho^\pm \nu)}{dE_\ell^* d\Omega_\ell^* dE_\gamma^* d\Omega_\gamma^* d\Omega_\rho^* dm_{\pi\pi}^2 d\tilde{\Omega}_\pi d\Omega_\tau} = A_0 + \bar{\eta} A_1 + \xi\kappa A_2$$

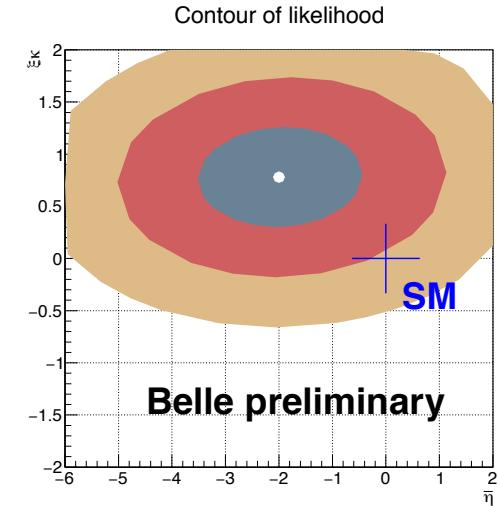
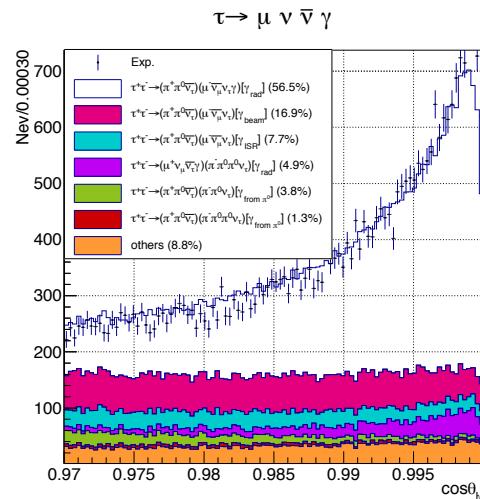
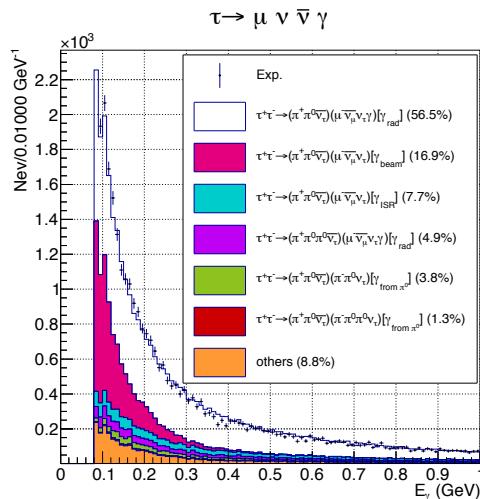
$$\mathcal{F}(\vec{z}) = \frac{d\sigma(\ell^\mp \nu_\nu \gamma, \rho^\pm \nu)}{dp_\ell d\Omega_\ell dp_\gamma d\Omega_\gamma dp_\rho d\Omega_\rho dm_{\pi\pi}^2 d\tilde{\Omega}_\pi} = \int_{\Phi_1}^{\Phi_2} \frac{d\sigma(\ell^\mp \nu_\nu \gamma, \rho^\pm \nu)}{dE_\ell^* d\Omega_\ell^* dE_\gamma^* d\Omega_\gamma^* d\Omega_\rho^* dm_{\pi\pi}^2 d\tilde{\Omega}_\pi d\Omega_\tau} |\text{JACOBIAN}| d\Phi_\tau$$

$$\mathcal{L} = \prod_{k=1}^N \mathcal{P}^{(k)}, \quad \mathcal{P}^{(k)} = \frac{\mathcal{F}(\vec{z}^{(k)})}{\mathcal{N}(\vec{\Theta})} = \frac{\mathcal{F}_0 + \mathcal{F}_1 \bar{\eta} + \mathcal{F}_2 \xi\delta}{\mathcal{N}_0 + \mathcal{N}_1 \bar{\eta} + \mathcal{N}_2 \xi\delta}, \quad \mathcal{N}_k = \int \mathcal{F}_k(\vec{z}) d\vec{z}, \quad (k = 0, 1, 2)$$

$\bar{\eta}$ and $\xi\delta$ are extracted in the unbinned maximum likelihood fit of $(\ell \nu \nu \gamma; \rho \nu)$ events in the 12D phase space in CMS.

Michel parameters in $\tau \rightarrow \ell \nu \nu \gamma$, ($\ell = e, \mu$) (II)

$N_{\tau\tau} = 646 \times 10^6$, selected: 71171 ($\mu \nu \nu \gamma$; $\rho \nu$) and 776834 ($e \nu \nu \gamma$; $\rho \nu$) events



Source	$\sigma_{\bar{\eta}}^e$	$\sigma_{\xi \kappa}^e$	$\sigma_{\bar{\eta}}^\mu$	$\sigma_{\xi \kappa}^\mu$
Normalization	4.3	0.94	0.15	0.04
Background PDF	2.5	0.24	0.67	0.22
Branching ratios	3.8	0.05	0.25	0.01
Cluster merge in ECL	2.2	0.46	0.02	0.06
Detector resolution	0.74	0.20	0.22	0.02
Data/MC eff. corr.	1.9	0.14	0.04	0.04
Total	7.0	1.1	0.76	0.24

Belle preliminary

$$\bar{\eta} = -1.3 \pm 1.5 \pm 0.8$$

$$\xi \kappa = 0.5 \pm 0.4 \pm 0.2$$

Backup ($\tau^\pm \rightarrow l^\pm l'^+ l'^- \nu_\tau \nu_l$)

Preliminary Results

$$\tau^\pm \rightarrow l^\pm l'^+ l'^- \nu_\tau \nu_l \quad (l, l' = e, \mu)$$

Preliminary : Estimation of systematic uncertainties

$$\tau^\pm \rightarrow e^\pm e^+ e^- \nu_\tau \nu_e$$

contents	syst. error (SVD1)	syst. error (SVD2)
PID correction	7.3%	5.0%
Tracking efficiency	1.1%	1.1%
Trigger correction	0.1%	0.1%
Tag-side	0.35%	0.35%
Luminosity	1.4%	1.4%
Background	5.5%	2.8%
Selection Cut	—	—
Total	9.3%	6.0%

$$\tau^\pm \rightarrow \mu^\pm e^+ e^- \nu_\tau \nu_\mu$$

contents	syst. error (SVD1)	syst. error (SVD2)
PID correction	6.9%	6.4%
Tracking efficiency	1.1%	1.1%
Trigger correction	0.1%	0.1%
Tag-side	0.35%	0.35%
Luminosity	1.4%	1.4%
Background	9.6%	4.8%
Selection Cut	—	—
Total	12.0%	8.2%

~ 10%

~ 15%

$$\tau^\pm \rightarrow e^\pm \mu^+ \mu^- \nu_\tau \nu_e$$

contents	syst. error (SVD1)	syst. error (SVD2)
PID correction	8.7%	7.4%
Tracking efficiency	1.1%	1.1%
Trigger correction	0.1%	0.1%
Tag-side	0.35%	0.35%
Luminosity	1.4%	1.4%
Background	71%	35%
Selection Cut	—	—
Total	72%	36%

$$\tau^\pm \rightarrow \mu^\pm \mu^+ \mu^- \nu_\tau \nu_\mu$$

contents	syst. error (SVD1)	syst. error (SVD2)
PID correction	6.2%	8.4%
Tracking efficiency	1.1%	1.1%
Trigger correction	0.1%	0.1%
Tag-side	0.35%	0.35%
Luminosity	1.4%	1.4%
Background	71%	35%
Selection Cut	—	—
Total	72%	36%

~ 70%

~ 70%

We are working to finalize the result and the estimation of systematic errors will be finished soon.

Theoretical Formula of BRs depends on Michel parameters

- Method of Monte Carlo integral is used
- Take the ratio of BR_{SM} to avoid considering the complicated common factor appears in the theoretical formula

$$b \sim g = \frac{1}{N_{gen}} \sum_{\mathbf{x} \in \Omega} \frac{d\Gamma_{NP}(\mathbf{x})}{d\Gamma_{SM}(\mathbf{x})}$$

$$\begin{aligned} b \sim g &= \frac{1}{BR_{SM}} BR_{NP} = \frac{1}{\Gamma_{SM}} \int d\Gamma_{NP} d(PS) = \frac{1}{\Gamma_{SM}} \int \frac{d\Gamma_{NP}}{d\Gamma_{SM}/\Gamma_{SM}} [(d\Gamma_{SM}/\Gamma_{SM}) d(PS)] \\ &= \frac{1}{\Gamma_{SM}} \int \frac{d\Gamma_{NP}}{d\tilde{\Gamma}_{SM}} [(d\tilde{\Gamma}_{SM}) d(PS)] \approx \frac{1}{\Gamma_{SM}} \frac{1}{N_{gen}} \sum_{\mathbf{x} \in \Omega} \frac{d\Gamma_{NP}(\mathbf{x})}{d\tilde{\Gamma}_{SM}(\mathbf{x})} = \frac{1}{N_{gen}} \sum_{\mathbf{x} \in \Omega} \frac{d\Gamma_{NP}(\mathbf{x})}{d\Gamma_{SM}(\mathbf{x})}, \quad (5.2) \end{aligned}$$

where, $d\tilde{\Gamma}_{SM} = d\Gamma_{SM}/\Gamma_{SM}$ is a normalized differential decay width of the SM, Ω is an allowed phase space (PS), \mathbf{x} follows the distribution of $d\Gamma_{SM}$, and N_{gen} is the number of generated events.

From next page, we show the result of calculated coefficients for our four target modes.

$$\tau^\pm \rightarrow e^\pm e^+ e^- \nu_\tau \nu_e$$

$$BR_{\text{exp}} = BR_{\text{SM}} \{ Q_{LL} + (1.051 \pm 0.036) Q_{LR} + (-0.2053 \pm 0.1431) B_{LR} + L \leftrightarrow R \\ + \Re[(0.2416 \pm 0.0002) I_\alpha + (0.8606 \pm 0.0001) I_\beta] \} + BR_{\text{NLO}}. \quad (7.19)$$

The formulation of coupling constant g_{jk}^i is written by,

$$BR_{\text{exp}} = BR_{\text{SM}} \{ |g_{LL}^V|^2 (1 + \frac{|g_{LL}^S|^2}{4|g_{LL}^V|^2}) + (0.2501 \pm 0.0001) |g_{RL}^S|^2 + (0.8465 \pm 0.1073) |g_{RL}^V|^2 \\ + (2.693 \pm 0.215) |g_{RL}^T|^2 + \Re[-(0.1540 \pm 0.1073) g_{RL}^S g_{RL}^{T*} + (0.4303 \pm 0.0001) g_{LL}^S g_{RR}^{V*} \\ + (0.06039 \pm 0.00004) g_{LR}^S g_{RL}^{V*} + (0.3623 \pm 0.0002) g_{LR}^V g_{RL}^{T*}] + L \leftrightarrow R \} + BR_{\text{NLO}}.$$

$$\tau^\pm \rightarrow \mu^\pm e^+ e^- \nu_\tau \nu_\mu$$

$$BR_{\text{exp}} = BR_{\text{SM}} \{ Q_{LL} + (1.220 \pm 0.049) Q_{LR} + (-0.8717 \pm 0.1957) B_{LR} + L \leftrightarrow R \\ + \Re[(181.3 \pm 0.1) I_\alpha + (104.4 \pm 0.1) I_\beta] \} + BR_{\text{NLO}}. \quad (7.22)$$

The formulation of coupling constant g_{jk}^i is written by,

$$BR_{\text{exp}} = BR_{\text{SM}} \{ |g_{LL}^V|^2 (1 + \frac{|g_{LL}^S|^2}{4|g_{LL}^V|^2}) + (0.2506 \pm 0.0001) |g_{RL}^S|^2 + (0.3484 \pm 0.1468) |g_{RL}^V|^2 \\ + (1.699 \pm 0.294) |g_{RL}^T|^2 + \Re[-(0.6538 \pm 0.1468) g_{RL}^S g_{RL}^{T*} + (52.20 \pm 0.01) g_{LL}^S g_{RR}^{V*} \\ + (45.33 \pm 0.01) g_{LR}^S g_{RL}^{V*} + (272.0 \pm 0.1) g_{LR}^V g_{RL}^{T*}] + L \leftrightarrow R \} + BR_{\text{NLO}}.$$

Because of pseudo peculiarity which appears in some terms, some result include large error. This pseudo peculiarity is caused mainly by the factor of virtual gamma conversion ($\gamma \rightarrow ee$)
 $|1/q_{ee}|^2 \sim |1/O(1\text{MeV})^2|^2$
in the matrix element.

$$\tau^\pm \rightarrow e^\pm \mu^+ \mu^- \nu_\tau \nu_e$$

↓ from my note

$$BR_{\text{exp}} = BR_{\text{SM}} \{ Q_{LL} + (1.226 \pm 0.001)Q_{LR} + (-0.8456 \pm 0.0001)B_{LR} + L \leftrightarrow R \\ + \Re[(0.2253 \pm 0.0001)I_\alpha + (0.5231 \pm 0.0001)I_\beta] \} + BR_{\text{NLO}}. \quad (7.24)$$

The formulation of coupling constant g_{jk}^i is written by,

$$BR_{\text{exp}} = BR_{\text{SM}} \{ |g_{LL}^V|^2 (1 + \frac{|g_{LL}^S|^2}{4|g_{LL}^V|^2}) + (0.2536 \pm 0.0001)|g_{RL}^S|^2 + (0.3802 \pm 0.0001)|g_{RL}^V|^2 \\ + (1.775 \pm 0.001)|g_{RL}^T|^2 + \Re[-(0.6342 \pm 0.0001)g_{RL}^S g_{RL}^{T*} + (0.2616 \pm 0.0001)g_{LL}^S g_{RR}^{V*} \\ + (0.05633 \pm 0.00001)g_{LR}^S g_{RL}^{V*} + (0.3380 \pm 0.0001)g_{LR}^V g_{RL}^{T*}] + L \leftrightarrow R \} + BR_{\text{NLO}}.$$

$$\tau^\pm \rightarrow \mu^\pm \mu^+ \mu^- \nu_\tau \nu_\mu$$

$$BR_{\text{exp}} = BR_{\text{SM}} \{ Q_{LL} + (1.216 \pm 0.005)Q_{LR} + (-0.8459 \pm 0.0005)B_{LR} + L \leftrightarrow R \\ + \Re[-(18.00 \pm 0.01)I_\alpha + (197.3 \pm 0.1)I_\beta] \} + BR_{\text{NLO}}. \quad (7.26)$$

The formulation of coupling constant g_{jk}^i is written by,

$$BR_{\text{exp}} = BR_{\text{SM}} \{ |g_{LL}^V|^2 (1 + \frac{1}{4}|g_{LL}^S|^2) + (0.2512 \pm 0.0001)|g_{RL}^S|^2 + (0.3704 \pm 0.0001)|g_{RL}^V|^2 \\ + (1.745 \pm 0.015)|g_{RL}^T|^2 + \Re[-(0.6344 \pm 0.0004)g_{RL}^S g_{RL}^{T*} + (98.67 \pm 0.01)g_{LL}^S g_{RR}^{V*} \\ - (4.510 \pm 0.001)g_{LR}^S g_{RL}^{V*} - (27.060 \pm 0.006)g_{LR}^V g_{RL}^{T*}] + L \leftrightarrow R \} + BR_{\text{NLO}}.$$

POSSIBLE MEASUREMENT

Because of expected statistics, we concentrate on the measurement through two modes:

$$\tau^\pm \rightarrow e^\pm e^+ e^- \nu_\tau \nu_e, \tau^\pm \rightarrow \mu^\pm e^+ e^- \nu_\tau \nu_\mu$$

The statistics of other two modes ($\tau^\pm \rightarrow e^\pm \mu^+ \mu^- \nu_\tau \nu_e, \tau^\pm \rightarrow \mu^\pm \mu^+ \mu^- \nu_\tau \nu_\mu$) expected to be small because of its expected branching fraction is small. And the measurement of Michel parameters is difficult from these two modes.

Expected BRs from the Standard Model

Prediction from theory	
Channel	
$\text{BR}(\tau^- \rightarrow e^- e^+ e^- \bar{\nu}_e \nu_\tau) \times 10^5$	4.21 ± 0.01
$\text{BR}(\tau^- \rightarrow e^- \mu^+ \mu^- \bar{\nu}_e \nu_\tau) \times 10^7$	1.247 ± 0.001
$\text{BR}(\tau^- \rightarrow \mu^- e^+ e^- \bar{\nu}_\mu \nu_\tau) \times 10^5$	1.984 ± 0.004
$\text{BR}(\tau^- \rightarrow \mu^- \mu^+ \mu^- \bar{\nu}_\mu \nu_\tau) \times 10^7$	1.183 ± 0.001

Ref. [JHEP 1604, 185 (2016)]

Result of Monte Carlo simulation

	$e^\pm e^+ e^- \nu_\tau \nu_e$	$\mu^\pm e^+ e^- \nu_\tau \nu_\mu$	$e^\pm \mu^+ \mu^- \nu_\tau \nu_e$	$\mu^\pm \mu^+ \mu^- \nu_\tau \nu_\mu$
Detection Efficiency	1.76 %	1.20%	3.56%	1.67%
Main Background(s)	$e\nu_\tau \nu_e \gamma, \pi\pi^0 \nu_\tau$	$\mu\nu_\tau \nu_\mu \gamma, \pi\pi^0 \pi^0 \nu_\tau$ $, \pi\pi^0 (\rightarrow e^+ e^- \gamma) \nu_\tau$	$\pi\pi^0 \nu_\tau$	$\pi\pi^+ \pi^- \nu_\tau$
Expected number of signals at Belle	1300	430	8	4
Purity of signal	47%	50%	37%	16%

Previous Experiment

CLEOII measured the branching fraction of $\tau^\pm \rightarrow (e/\mu^\pm) e^+ e^- \nu_\tau \nu_{e/\mu}$

Result of CLEOII

$$Br(\tau \rightarrow ee^+ e^- \nu_\tau \nu_e) = (2.7^{+1.5+0.4+0.1}_{-1.1-0.4-0.3}) \times 10^{-5}$$

$$Br(\tau \rightarrow \mu e^+ e^- \nu_\tau \nu_\mu) < 3.2 \times 10^{-5} \text{ (at 90% C.L.)}$$

CLEO-II experiment

Decay mode	Number of events
$\tau^\pm \rightarrow e^\pm e^+ e^- \nu_\tau \nu_e$	5
$\tau^\pm \rightarrow \mu^\pm e^+ e^- \nu_\tau \nu_\mu$	1

Integrated luminosity $3.6 fb^{-1}$
 $N_{\tau\tau} = (3.28 \pm 0.05) \times 10^6$

- Main source of systematic error
 - Uncertainties of lepton identification efficiency
 - Uncertainties of reconstruction efficiency of slow tracks

Reference: Phys. Rev. Lett. 76, 2637 (1996)

Detection efficiency:

$$\tau \rightarrow ee^+ e^- \nu_\tau \nu_e \quad (2.7 \pm 0.1)\%$$

$$\tau \rightarrow \mu e^+ e^- \nu_\tau \nu_\mu \quad (1.9 \pm 0.1)\%$$

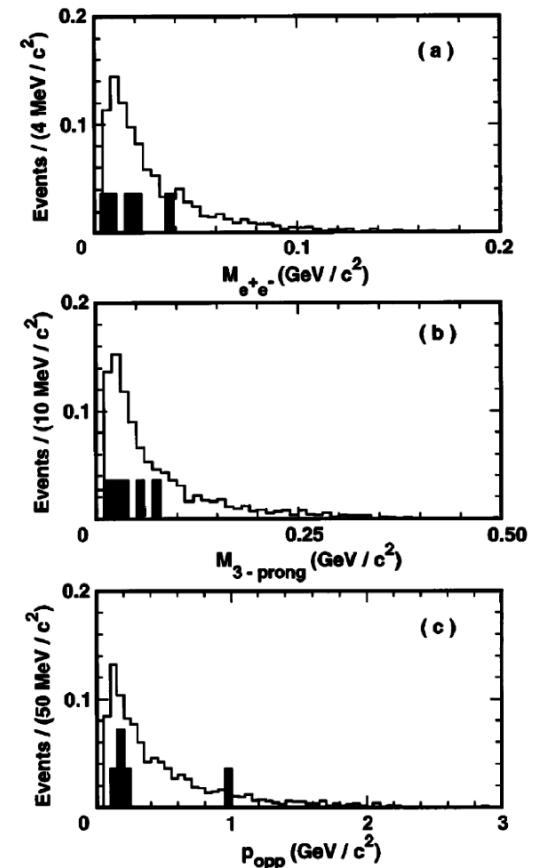


FIG. 3. Comparison of the kinematical distributions of the $\tau \rightarrow ee^+ e^- \nu_\tau \nu_e$ Monte Carlo (solid line) and the data (shaded histogram) for events passing all selection requirements: (a) the e^+e^- invariant mass averaged over two possible combinations, $M_{e^+e^-}$, (b) the 3-prong invariant mass, $M_{3\text{-prong}}$, and (c) the momentum of the electron on the 3-prong side with the charge opposite to that of the parent tau, p_{opp} . The normalization of the plots is arbitrary.

Selection

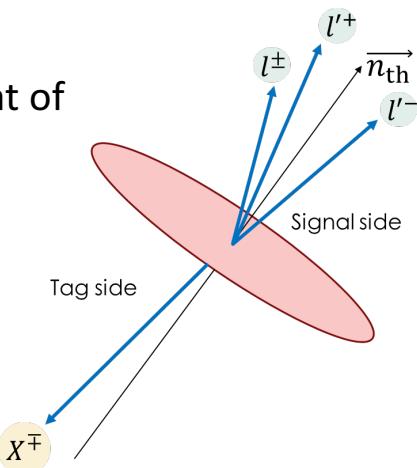
Pre-selection of tau-pair and thrust selection are applied at the first stage

Thrust selection is applied as following method (in the CM-frame).

- Define thrust vector by $\overrightarrow{n_{\text{th}}} = \max\left\{\frac{\sum_{\text{sig,tag}} |\overrightarrow{n_{\text{th}}} \cdot \overrightarrow{p_i}|}{\sum_{\text{sig,tag}} |\overrightarrow{p_i}|}\right\}$
- Separate signal and tag side by:
 $[\{\overrightarrow{n_{\text{th}}} \cdot \overrightarrow{p_{i,\text{sig}}} < 0\} \&\& \{\overrightarrow{n_{\text{th}}} \cdot \overrightarrow{p_{\text{tag}}} > 0\}] \text{ or } [\{\overrightarrow{n_{\text{th}}} \cdot \overrightarrow{p_{i,\text{sig}}} > 0\} \&\& \{\overrightarrow{n_{\text{th}}} \cdot \overrightarrow{p_{\text{tag}}} < 0\}].$
- Require the number of charged tracks in signal side to be three and that of tag side to be one
- Require $\sum_{\text{sig}} Q = \pm 1$ and $\sum_{\text{tag}} Q = \mp 1$

$\overrightarrow{p_{i,\text{sig}}}$: momentum vector of electrons in signal-side

$\overrightarrow{p_{\text{tag}}}$: momentum vector of charged particle in tag-side



※ Pre-selection of tau-pair and second stage selection are written in a backup

We use 1-prong decay of tau as tag

Pre-Selection of Tau-pair

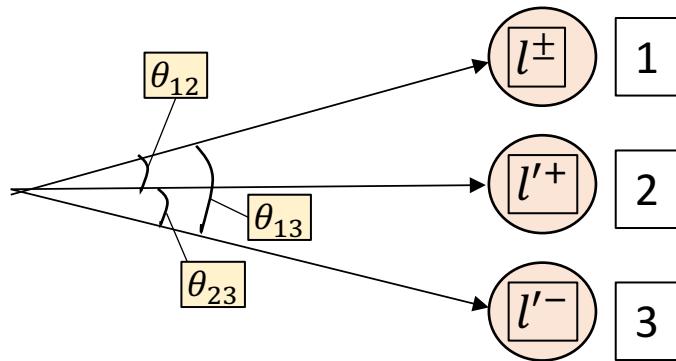
Table 3: Pre-selection criteria of tau-pair

Index	Selection Criteria
1	$2 < \text{Number of charged tracks} < 8$
2	$ \text{Sum of charge} \leq 2$
3	Sum of momenta of charged tracks in the CM frame (P^{CM}) $< 10 \text{ GeV}/c$
4	Sum of energy deposit in the ECL $E(\text{ECL}) < 10 \text{ GeV}$
5	Maximum Pt of charged track (Pt_{max}) $> 0.5 \text{ GeV}/c$
6	Event vertex $ r < 0.5 \text{ cm}$, $ z < 3.0 \text{ cm}$
7	For 2 track events, 7-1,7-2, and 7-3 must be satisfied:
7-1	Sum of $P^{\text{CM}} < 9 \text{ GeV}/c$
7-2	Sum of $E(\text{ECL}) < 9 \text{ GeV}$
7-3	$5 \text{ deg} < \theta_{\text{missing momentum}} < 175 \text{ deg}$
8	$E_{\text{rec}} = [\text{Sum of } P^{\text{CM}} + \text{Sum of } E_{\gamma}^{\text{CM}} \text{ (energy of } \gamma \text{ in the CM frame)}] > 3 \text{ GeV}$.or. $Pt_{\text{max}} > 1.0 \text{ GeV}/c$
9	For 2-4 track events, 9-1 and 9-2 must be satisfied:
9-1	$E_{\text{tot}} = [E_{\text{rec}} + P_{\text{miss}}^{\text{CM}}] < 9 \text{ GeV}$.or. maximum opening angle $< 175 \text{ deg}$
9-2	[Number of tracks within $30 < \theta < 130 \text{ deg}$] ≥ 2 .or. [Sum of $E(\text{ECL}) - \text{Sum of } E_{\gamma}^{\text{CM}}$] $< 5.3 \text{ GeV}$
10	Maximum opening angle $> 20 \text{ deg}$

Second Stage Selection

Explanation of defined word

Sum of $\cos \theta_{ij}$: $(\sum_{i < j} \cos \theta_{ij})$



Second Stage Selection

$$\tau^\pm \rightarrow e^\pm e^+ e^- \nu_\tau \nu_e$$

1. Number of charged track = 4
2. Total charge (sum of Q_{sig} + sum of Q_{tag}) = 0
3. Number of photons (with $E(\gamma)_{CM} > 0.06$ GeV) ≤ 8
4. Total ECL energy deposition < 9 GeV
5. $1.5 \text{ GeV}/c^2 < M_{\text{missing}} < 7 \text{ GeV}/c^2$
6. Number of tracks in signal side = 3 $\&\&$ Number of track in tag side = 1
7. Max transverse momentum of electron in signal side $|\vec{p}_{t_i}| > 0.15 \text{ GeV} / c$ (CM-frame)

$$\tau^\pm \rightarrow e^\pm e^+ e^- \nu_\tau \nu_e$$

8. Reconstructed vertex position of $\gamma (\rightarrow e^+ e^-)$ should be $r(xy - \text{plane}) < 1.5\text{cm}$
9. Reconstructed vertex position of $\gamma (\rightarrow e^+ e^-)$ should be $r(xyz - \text{space}) < 3.0\text{cm}$
10. The number of γ in signal-side ≤ 1 $\&\&$ The sum of energy of γ in signal-side < 0.5 GeV
11. Sum of $\cos \theta_{ij}$: $(\sum_{i < j} \cos \theta_{ij}) > 2.90$
12. $|dz_i|$ of electrons in signal-side should be $|dz_i| < 1\text{cm}$
13. The momentum of virtual gamma in lab frame $< 3 \text{ GeV}/c$
14. Polar angle of the missing momentum $30 \text{ deg} < \theta < 150 \text{ deg}$

Second Stage Selection

$$\tau^\pm \rightarrow \mu^\pm e^+ e^- \nu_\tau \nu_\mu$$

1. Number of charged track = 4
2. Total charge (sum of Q_{sig} + sum of Q_{tag}) = 0
3. Number of photons (with $E(\gamma)_{CM} > 0.06 \text{ GeV}$) ≤ 8
4. Total number of gamma in sig & tag -side ≤ 4
5. Total ECL energy deposition $< 9 \text{ GeV}$
6. $1.5 \text{ GeV}/c^2 < M_{\text{missing}} < 7 \text{ GeV}/c^2$
7. Number of charged tracks in signal side = 3 && Number of charged track in tag side = 1
8. Reconstructed vertex position of $\gamma (\rightarrow e^+ e^-)$ should be $r(xy - \text{plane}) < 1.5 \text{ cm}$
9. The sum of energy of γ in signal-side $< 0.5 \text{ GeV}$
10. Sum of $\cos \theta_{ij}$: $(\sum_{i < j} \cos \theta_{ij}) > 2.93$
11. Polar angle of the missing momentum $30 \text{ deg} < \theta < 150 \text{ deg}$

Second Stage Selection

$$\tau^\pm \rightarrow e^\pm \mu^+ \mu^- \nu_\tau \nu_e$$

1. Number of charged track = 4
2. Total charge (sum of Q_{sig} + sum of Q_{tag}) = 0
3. Number of photons (with $E(\gamma)_{CM} > 0.06 \text{ GeV}$) ≤ 8
4. Total ECL energy deposition $< 9 \text{ GeV}$
5. $1.5 \text{ GeV}/c^2 < M_{\text{missing}} < 7 \text{ GeV}/c^2$
6. Number of tracks in signal side = 3 $\&\&$ Number of track in tag side = 1
8. The number of γ in signal-side ≤ 1 $\&\&$ The sum of energy of γ in signal-side $< 0.3 \text{ GeV}$
9. Sum of $\cos \theta_{ij}$: $(\sum_{i < j} \cos \theta_{ij}) > 2.70$
10. $(E_c/p)_\mu$ of muons in signal-side should be < 0.5
11. Polar angle of the missing momentum $30 \text{ deg} < \theta < 150 \text{ deg}$
12. Invariant mass of $\mu^+ \mu^-$ should be $< 0.4 \text{ GeV}/c^2$

Second Stage Selection

$$\tau^\pm \rightarrow \mu^\pm \mu^+ \mu^- \nu_\tau \nu_\mu$$

1. Number of charged track = 4
2. Total charge (sum of Q_{sig} + sum of Q_{tag}) = 0
3. Number of photons (with $E(\gamma)_{CM} > 0.06 \text{ GeV}$) ≤ 8
4. Total ECL energy deposition $< 9 \text{ GeV}$
5. $1.5 \text{ GeV}/c^2 < M_{\text{missing}} < 7 \text{ GeV}/c^2$
6. Number of tracks in signal side = 3 && Number of track in tag side = 1
8. The number of γ in signal-side = 0
9. Sum of $\cos \theta_{ij}$: $(\sum_{i < j} \cos \theta_{ij}) > 2.85$
10. $(E_c)_\mu$ of muons in signal-side should be $< 0.4 \text{ GeV}/c^2$
11. Polar angle of the missing momentum $30 \text{ deg} < \theta < 150 \text{ deg}$
12. $|dz_i|$ of muons in signal-side should be $|dz_i| < 0.1 \text{ cm}$