

# Higher twist effects in $e^+e^-$ annihilation at high energies

Shu-yi Wei (魏树一)

Shandong University

WSY, Yu-kun Song, Zuo-tang Liang, **Phys.Rev.D89,(2014)014024**

WSY, Kai-bao Chen, Yu-kun Song, Zuo-tang Liang, **arXiv:1410.4314**

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## Process Involving Hadron States in Initial or Final State

$$\text{Cross Section} = \text{Hard Part} \otimes \text{PDFs/FFs} + \dots$$

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### PDFs and FFs

- Transverse Momentum Dependence
- Spin Dependence



### Experiments

- Azimuthal Asymmetries
- Spin Asymmetries

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**Collinear Expansion**



Higher Twist Terms

## PDFs

R.K. Ellis, W. Furmanski and R. Petronzio

NPB 1982, 1983

J.W. Qiu, G. Sterman NPB, 1991

## TMD PDFs

Z.T. Liang and X.N. Wang

PRD, 2007

## Collinear Expansion



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## Fragmentation Functions (Leading Order Twist-3 Contributions)

$$e^+e^- \rightarrow h + X$$

Tree Level

FFs (PRD 89, (2014) 014024)

$$e^+e^- \rightarrow h + \bar{q} + X$$

$\Rightarrow$

TMD FFs (arXiv:1410.4314)



$$d\sigma = \frac{g_Z^4}{32s} L_{\mu'\nu'}(l_1, l_2) D_F^{\mu'\mu}(q) D_F^{\nu'\nu*}(q) W_{\mu\nu}(q, p, S) \frac{d^3 p}{(2\pi)^2 2E_p}$$



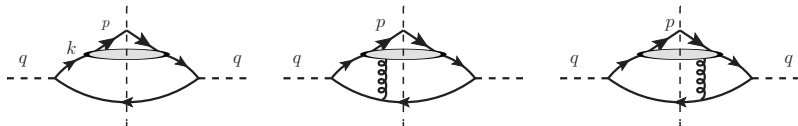
$$d\sigma = \frac{g_Z^4}{32s} L_{\mu'\nu'}(l_1, l_2) D_F^{\mu'\mu}(q) D_F^{\nu'\nu^*}(q) W_{\mu\nu}(q, p, S) \frac{d^3 p}{(2\pi)^2 2E_p}$$

$$W_{\mu\nu} = \left| \begin{array}{c} \text{Diagram 1: A vertex with incoming momentum } q \text{ and outgoing momentum } p. \text{ A loop with momentum } k \text{ is shown.} \\ \text{Diagram 2: A vertex with incoming momentum } q \text{ and outgoing momentum } q. \text{ A loop with momentum } k \text{ is shown, with a vertical dashed line through it.} \end{array} \right|^2$$

$$W_{\mu\nu}^{(0)}(q, p, S) = \int \frac{d^4 k}{(2\pi)^4} \text{Tr} \left[ \hat{H}_{\mu\nu}^{(0)}(k, q) \hat{\Pi}^{(0)}(k, p, S) \right]$$

$$\hat{H}_{\mu\nu}^{(0)}(k) = \Gamma_{\mu}^q (\not{q} - \not{k}) \Gamma_{\nu}^q (2\pi) \delta_+((q-k)^2)$$

$$\hat{\Pi}^{(0)}(k) = \frac{1}{2\pi} \sum_X \int d^4 \xi e^{-ik\xi} \langle 0 | \psi(0) | hX \rangle \langle hX | \bar{\psi}(\xi) | 0 \rangle$$



$$W_{\mu\nu} = W_{\mu\nu}^{(0)} + W_{\mu\nu}^{(1)} + W_{\mu\nu}^{(2)} + \dots$$

Where,

$$W_{\mu\nu}^{(1,L)}(q, p, S) = \int \frac{d^4 k_1}{(2\pi)^4} \frac{d^4 k_2}{(2\pi)^4} \text{Tr}[\hat{H}_{\mu\nu}^{(1,L)\rho}(k_1, k_2, q) \hat{\Pi}_{\rho}^{(1,L)}(k_1, k_2, p, S)]$$

$$\hat{H}_{\mu\nu}^{(1,L)\rho}(k_1, k_2) = \Gamma_{\mu}^q (\not{q} - \not{k}_1) \gamma^{\rho} \frac{\not{k}_2 - \not{q}}{(k_2 - q)^2 - i\epsilon} \Gamma_{\nu}^q (2\pi) \delta_{+}((q - k_1)^2)$$

$$\hat{\Pi}_{\rho}^{(1,L)}(k_1, k_2) = \frac{1}{2\pi} \sum_X \int d^4 \xi d^4 \eta e^{-ik_1 \xi} e^{-i(k_2 - k_1) \eta} \langle 0 | g A_{\rho}(\eta) \psi(0) | hX \rangle \langle hX | \bar{\psi}(\xi) | 0 \rangle$$

## Taylor Expansion

$$\begin{aligned}\hat{H}_{\mu\nu}^{(0)}(k) &= \hat{H}_{\mu\nu}^{(0)}(z) + \left. \frac{\partial \hat{H}_{\mu\nu}^{(0)}(k)}{\partial k_\rho} \right|_{k=p^+/z} \omega_\rho^{\rho'} k_{\rho'} + \dots \\ &= \hat{H}_{\mu\nu}^{(0)}(z) - (\hat{H}_{\mu\nu}^{(1L)\rho}(z, z) + \hat{H}_{\mu\nu}^{(1R)\rho}(z, z)) \omega_\rho^{\rho'} k_{\rho'} + \dots \\ \hat{H}_{\mu\nu}^{(1L)\rho}(k_1, k_2) &= \hat{H}_{\mu\nu}^{(1L)\rho}(z_1, z_2) + \dots\end{aligned}$$

## Gluon Fields

$$A_\rho = A^+ \bar{n}_\rho + \omega_\rho^{\rho'} A_{\rho'}$$

## Ward Identities

$$p^+ \bar{n}_\rho \hat{H}_{\mu\nu}^{(1L)\rho}(z_1, z_2) = -\frac{z_1 z_2}{z_2 - z_1 - i\epsilon} \hat{H}_{\mu\nu}^{(0)}(z_1)$$

$$W_{\mu\nu}^{(0)} = H^{(0)}(k) \Phi_0 = H_{\mu\nu}^{(0)}(z) \times \phi_0 + H_{\mu\nu}^{(1)}(z) \times \phi_1 + \dots$$

$$W_{\mu\nu}^{(1)} = H^{(1)}(k) \Phi_1 = H_{\mu\nu}^{(0)}(z) \times \phi_2 + H_{\mu\nu}^{(1)}(z) \times \phi_3 + \dots$$

$$W_{\mu\nu} = W_{\mu\nu}^{(0)} + W_{\mu\nu}^{(1)} + W_{\mu\nu}^{(2)} + \dots \quad \Rightarrow \quad \tilde{W}_{\mu\nu} = \tilde{W}_{\mu\nu}^{(0)} + \tilde{W}_{\mu\nu}^{(1)} + \tilde{W}_{\mu\nu}^{(2)} + \dots$$

## Inclusive Process

$$\tilde{W}_{\mu\nu}^{(0)} = \frac{1}{2} \text{Tr} \left[ \hat{h}_{\mu\nu}^{(0)} \hat{\Xi}^{(0)}(z_B, p, S) \right] \quad \Rightarrow \quad \text{twist 2, twist 3, } \dots$$

$$\tilde{W}_{\mu\nu}^{(1L)} = -\frac{1}{4p \cdot q} \text{Tr} \left[ \hat{h}_{\mu\nu}^{(1)\rho} \omega_{\rho}^{\rho'} \hat{\Xi}^{(1)}(z_B, p, S) \right] \quad \Rightarrow \quad \text{twist 3, } \dots$$

## Semi-inclusive Process

$$\tilde{W}_{\mu\nu}^{(0,\text{si})} = \frac{1}{2} \text{Tr} \left[ \hat{h}_{\mu\nu}^{(0)} \hat{\Xi}^{(0)}(z_B, k'_{\perp}, p, S) \right] \quad \Rightarrow \quad \text{twist 2, twist 3, } \dots$$

$$\tilde{W}_{\mu\nu}^{(1,L,\text{si})} = -\frac{1}{4p \cdot q} \text{Tr} \left[ \hat{h}_{\mu\nu}^{(1)\rho} \omega_{\rho}^{\rho'} \hat{\Xi}^{(1,L)}(z_B, k'_{\perp}, p, S) \right] \quad \Rightarrow \quad \text{twist 3, } \dots$$

## Inclusive Process

$$\hat{\Xi}^{(0)}(z, p, S) = \sum_X \int \frac{p^+ d\xi^-}{2\pi} e^{-ik^+\xi^-} \langle 0 | \mathcal{L}^\dagger(0^-, \infty) \psi(0) | \text{hX} \rangle \langle \text{hX} | \bar{\psi}(\xi^-) \mathcal{L}(\xi^-, \infty) | 0 \rangle$$

$$\hat{\Xi}_\rho^{(1)}(z, p, S) = \sum_X \int \frac{p^+ d\xi^-}{2\pi} e^{-ik^+\xi^-} \langle 0 | \mathcal{L}^\dagger(0, \infty) [D_\rho(0) \psi(0)] | \text{hX} \rangle \langle \text{hX} | \bar{\psi}(\xi^-) \mathcal{L}(\xi^-, \infty) | 0 \rangle$$

## Semi-inclusive Process

$$\hat{\Xi}^{(0)}(z, k_\perp, p, S) = \sum_X \int \frac{p^+ d\xi^- d^2\xi_\perp}{2\pi} e^{-ik^+\xi^- + ik_\perp \cdot \xi_\perp} \langle 0 | \mathcal{L}^\dagger(0, \infty) \psi(0) | \text{hX} \rangle \langle \text{hX} | \bar{\psi}(\xi) \mathcal{L}(\xi, \infty) | 0 \rangle$$

$$\hat{\Xi}_\rho^{(1)}(z, k_\perp, p, S) = \sum_X \int \frac{p^+ d\xi^- d^2\xi_\perp}{2\pi} e^{-ik^+\xi^- + ik_\perp \cdot \xi_\perp} \langle 0 | \mathcal{L}^\dagger(0, \infty) \hat{D}_\rho(0) \psi(0) | \text{hX} \rangle \langle \text{hX} | \bar{\psi}(\xi) \mathcal{L}(\xi, \infty) | 0 \rangle$$

$$W_0^{\mu\nu} \propto \text{Tr}[\gamma^\mu \not{p} \gamma^\nu \hat{\Xi}^{(0)}]$$

## Gamma Matrices Expansion

$$\begin{aligned}\hat{\Xi}^{(0)}(z, p, S, k_\perp) &= \Xi_\alpha^{(0)}(z, p, S, k_\perp) \gamma^\alpha + \tilde{\Xi}_\alpha^{(0)}(z, p, S, k_\perp) \gamma_5 \gamma^\alpha + \dots \\ \hat{\Xi}_\rho^{(1)}(z, p, S, k_\perp) &= \Xi_{\rho\alpha}^{(1)}(z, p, S, k_\perp) \gamma^\alpha + \tilde{\Xi}_{\rho\alpha}^{(1)}(z, p, S, k_\perp) \gamma_5 \gamma^\alpha + \dots\end{aligned}$$

## Parity Conservation

$$\Xi_\alpha^{(0)}(V, A) = \Xi^{(0)\alpha}(\tilde{V}, -\tilde{A})$$

**V** : Vector

$$\tilde{\Xi}_\alpha^{(0)}(V, A) = -\tilde{\Xi}^{(0)\alpha}(\tilde{V}, -\tilde{A})$$

**A** : Axis Vector

---


$$\tilde{V}^\mu = V_\mu$$

## Spin Density Matrix (Particle Rest Frame)

$$\text{Spin } \frac{1}{2} \quad \rho = \begin{pmatrix} \rho_{++} & \rho_{+-} \\ \rho_{-+} & \rho_{--} \end{pmatrix} = \frac{1}{2}(1 + S^i \sigma^i)$$

$$\text{Spin } 1 \quad \rho = \begin{pmatrix} \rho_{++} & \rho_{+0} & \rho_{+-} \\ \rho_{0+} & \rho_{00} & \rho_{0-} \\ \rho_{-+} & \rho_{-0} & \rho_{--} \end{pmatrix} = \frac{1}{3} + \frac{1}{2} S^i \sigma^i + T^{ij} \Sigma^{ij}$$

Relativistic Case  $p^\rho = (p^+, M^2/2p^+, \vec{0}_\perp)$

$$S^\rho = \lambda_h \left( \frac{p^+}{M} \bar{n}^\rho - \frac{M}{2p^+} n^\rho \right) + S_\perp^\rho$$

$$T^{\mu\nu} = \frac{1}{2} \left\{ \frac{4}{3} S_{LL} \left[ \left( \frac{p^+}{M} \right)^2 \bar{n}^\mu \bar{n}^\nu + \left( \frac{M}{2p^+} \right)^2 n^\mu n^\nu - \frac{1}{2} (\bar{n}^{\{\mu} n^{\nu\}} - g_\perp^{\mu\nu}) \right] \right. \\ \left. + \left[ \frac{p^+}{M} \bar{n}^{\{\mu} - \frac{M}{2p^+} n^{\{\mu} \right] S_{LT}^{\nu\}} + S_{TT}^{\mu\nu} \right] \right\}$$

## Inclusive Process

$$\tilde{W}_{\mu\nu}^{(0)} = \frac{1}{2} \text{Tr} \left[ \hat{h}_{\mu\nu}^{(0)} \gamma_\alpha \right] \Xi^{(0)\alpha}(z_B, p, S) + \frac{1}{2} \text{Tr} \left[ \hat{h}_{\mu\nu}^{(0)} \gamma_5 \gamma_\alpha \right] \tilde{\Xi}^{(0)\alpha}(z_B, p, S)$$

$$\tilde{W}_{\mu\nu}^{(1L)} = -\frac{1}{4p \cdot q} \text{Tr} \left[ \hat{h}_{\mu\nu}^{(1)\rho} \gamma_\alpha \right] \omega_\rho^{\rho'} \Xi^{(1)\alpha}(z_B, p, S) - \frac{1}{4p \cdot q} \text{Tr} \left[ \hat{h}_{\mu\nu}^{(1)\rho} \gamma_5 \gamma_\alpha \right] \omega_\rho^{\rho'} \tilde{\Xi}^{(1)\alpha}(z_B, p, S)$$

$$\Xi^{(0)\alpha}(z, p, S) = \sum_X \int \frac{p^+ d\xi^-}{2\pi} e^{-ik^+\xi^-} \text{Tr}[\gamma^\alpha \langle 0 | \mathcal{L}^\dagger(0^-, \infty) \psi(0) | hX \rangle \langle hX | \bar{\psi}(\xi^-) \mathcal{L}(\xi^-, \infty) | 0 \rangle]$$

$$\Xi_\rho^{(1)\alpha}(z, p, S) = \sum_X \int \frac{p^+ d\xi^-}{2\pi} e^{-ik^+\xi^-} \text{Tr}[\gamma^\alpha \langle 0 | \mathcal{L}^\dagger(0, \infty) [D_\rho(0) \psi(0)] | hX \rangle \langle hX | \bar{\psi}(\xi^-) \mathcal{L}(\xi^-, \infty) | 0 \rangle]$$

...



## Lorentz Structures Up to Twist-3

Spin Independent

$$z\Xi^{(0)\alpha} = p^\alpha D_1(z)$$

Spin Vector

$$\begin{aligned} z\Xi^{(0)\alpha} &= M \epsilon_{\perp}^{\alpha\gamma} S_{\perp\gamma} D_T(z) \\ z\tilde{\Xi}^{(0)\alpha} &= \lambda_h p^\alpha \Delta D_{1L}(z) + M S_{\perp}^{\alpha} \Delta D_T(z) \\ z\Xi^{(1)\rho\alpha} &= M \epsilon_{\perp}^{\rho\gamma} S_{\perp\gamma} p^\alpha \xi_{\perp S}^{(1)}(z) \\ z\tilde{\Xi}^{(1)\rho\alpha} &= i M S_{\perp}^{\rho} p^\alpha \tilde{\xi}_{\perp S}^{(1)}(z) \end{aligned}$$

Spin Tensor

$$\begin{aligned} z\Xi^{(0)\alpha} &= S_{LL} p^\alpha D_{1LL}(z) + M S_{LT}^{\alpha} D_{LT}(z) \\ z\tilde{\Xi}^{(0)\alpha} &= M \epsilon_{\perp}^{\alpha\gamma} S_{LT,\gamma} \Delta D_{LT}(z) \\ z\Xi^{(1)\rho\alpha} &= M S_{LT}^{\rho} \xi_{LTS}^{(1)}(z) p^\alpha \\ z\tilde{\Xi}^{(1)\rho\alpha} &= i M \epsilon_{\perp}^{\rho\gamma} S_{LT,\gamma} \tilde{\xi}_{LTS}^{(1)}(z) p^\alpha \end{aligned}$$



## Leading twist fragmentation functions (possibility density)

$D_1(z)$       quark  $\rightarrow \sum_S$  hadron

$\Delta D_{1L}(z)$       longitudinally polarized quark  $\rightarrow$  longitudinally polarized hadron

$D_{1LL}(z)$       quark  $\rightarrow$  LL polarized hadron

## Twist-3 fragmentation functions

Vector Polarized Hadron

$D_T(z), \Delta D_T(z)$

Tensor Polarized Hadron

$D_{LT}(z), \Delta D_{LT}(z)$

## Gauge Invariant Hadronic Tensor

Spin-0  $\tilde{W}_{\mu\nu} = -\frac{2}{z} (c_1^q d_{\mu\nu} + i c_3^q \epsilon_{\perp\mu\nu}) D_1(z)$

---

Spin- $\frac{1}{2}$   $\tilde{W}_{\mu\nu} = \frac{2}{z} \left\{ - (c_1^q d_{\mu\nu} + i c_3^q \epsilon_{\perp\mu\nu}) D_1(z) + \lambda_h (c_3^q d_{\mu\nu} + i c_1^q \epsilon_{\perp\mu\nu}) \Delta D_{1L}(z) \right.$   
 $\left. + \frac{M}{p \cdot q} [c_1^q (q - 2p/z)_{\{\mu} \epsilon_{\perp\nu\}\gamma} S_{\perp}^{\gamma} + i c_3^q (q - 2p/z)_{[\mu} S_{\perp\nu]}] D_T(z) \right.$   
 $\left. - \frac{M}{p \cdot q} [c_3^q (q - 2p/z)_{\{\mu} S_{\perp\nu\}} - i c_1^q (q - 2p/z)_{[\mu} \epsilon_{\perp\nu]\gamma} S_{\perp}^{\gamma}] \Delta D_T(z) \right\}$

---

Spin-1  $\tilde{W}_{\mu\nu} = \frac{2}{z} \left\{ - (c_1^q d_{\mu\nu} + i c_3^q \epsilon_{\perp\mu\nu}) [D_1 + S_{LL} D_{1LL}] + \lambda_h (c_3^q d_{\mu\nu} + i c_1^q \epsilon_{\perp\mu\nu}) \Delta D_{1L} \right.$   
 $\left. + \frac{M}{p \cdot q} [c_1^q (q - 2p/z)_{\{\mu} \epsilon_{\perp\nu\}\gamma} S_{\perp}^{\gamma} + i c_3^q (q - 2p/z)_{[\mu} S_{\perp\nu]}] D_T(z) \right.$   
 $\left. + \frac{M}{p \cdot q} [c_1^q (q - 2p/z)_{\{\mu} S_{LT,\nu\}} - i c_3^q (q - 2p/z)_{[\mu} \epsilon_{\perp\nu]\gamma} S_{LT}^{\gamma}] D_{LT}(z) \right.$   
 $\left. - \frac{M}{p \cdot q} [c_3^q (q - 2p/z)_{\{\mu} S_{\perp\nu\}} - i c_1^q (q - 2p/z)_{[\mu} \epsilon_{\perp\nu]\gamma} S_{\perp}^{\gamma}] \Delta D_T(z) \right.$   
 $\left. - \frac{M}{p \cdot q} [c_3^q (q - 2p/z)_{\{\mu} \epsilon_{\perp\nu\}\gamma} S_{LT}^{\gamma} + i c_1^q (q - 2p/z)_{[\mu} S_{LT,\nu]}] \Delta D_{LT}(z) \right\}$

## Spin-0 Hadrons - $e^+e^- \rightarrow Z_0 \rightarrow hX$

$$\frac{d^2\sigma}{dzdy} = \sum_q \frac{2\pi\alpha^2}{Q^2} \chi T_0^q(y) D_1^{q \rightarrow h}(z).$$

$$D_1^{q \rightarrow h}(z) = \frac{z}{4} \sum_X \int \frac{d\xi^-}{2\pi} e^{-ip^+\xi^-/z} \text{Tr} \left[ \gamma^+ \langle 0 | \mathcal{L}^\dagger(0, \infty) \psi(0) | hX \rangle \langle hX | \bar{\psi}(\xi^-) \mathcal{L}(\xi^-, \infty) | 0 \rangle \right]$$

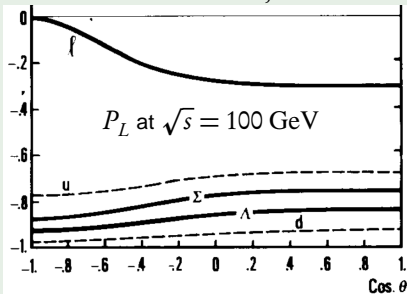
For  $e^+e^- \rightarrow \gamma^* \rightarrow q\bar{q} \rightarrow hX$ ,

$$\frac{d\sigma^{\text{em}}}{dz} = \sum_q \frac{4\pi\alpha^2}{3Q^2} e_q^2 D_1^{q \rightarrow h}(z)$$

## Spin-1/2 Hadrons - $e^+e^- \rightarrow Z_0 \rightarrow hX$

$$\frac{d^2\sigma}{dzdy} = \chi \frac{2\pi\alpha^2}{Q^2} \left\{ [T_0(y)D_1(z) + \lambda_h T_1(y)\Delta D_{1L}(z)] + \frac{4M}{zQ^2} [T_2(y)D_T(z)\epsilon_{\perp}^{l_1 S_{\perp}} + T_3(y)\Delta D_T(z)l_{\perp} \cdot S_{\perp}] \right\}$$

$$P_{Lh}(z, y) = \frac{\sum_q P_q(y) T_0^q(y) \Delta D_{1L}^{q \rightarrow h}(z)}{\sum_q T_0^q(y) D_1^{q \rightarrow h}(z)}$$



NPB, 1980, J.E. Augustin, F.M. Renard

$P_q$ : Longitudinal polarization of quark

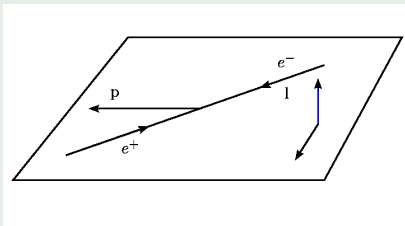
$\Delta D_{1L}^{q \rightarrow h}(z)$ : Longitudinally polarized quark  $\rightarrow$  Longitudinally polarized hadron

## Spin-1/2 Hadrons - $e^+e^- \rightarrow Z_0 \rightarrow hX$

$$\frac{d^2\sigma}{dzdy} = \chi \frac{2\pi\alpha^2}{Q^2} \left\{ [T_0(y)D_1(z) + \lambda_h T_1(y)\Delta D_{1L}(z)] \right. \\ \left. + \frac{4M}{zQ^2} [T_2(y)D_T(z)\epsilon_{\perp}^{l_{\perp}S_{\perp}} + T_3(y)\Delta D_T(z)l_{\perp} \cdot S_{\perp}] \right\}$$

$$P_{hx}(z, y) = -\frac{4M}{zQ} \frac{\sum_q \tilde{T}_3^q(y) \Delta D_T^{q \rightarrow h}(z)}{\sum_q T_0^q(y) D_1^{q \rightarrow h}(z)}$$

$$P_{hy}(z, y) = \frac{4M}{zQ} \frac{\sum_q \tilde{T}_2^q(y) D_T^{q \rightarrow h}(z)}{\sum_q T_0^q(y) D_1^{q \rightarrow h}(z)}$$



$P_q$ : Longitudinal polarization of quark

$\Delta D_{1L}^{q \rightarrow h}(z)$ : Longitudinally polarized quark  $\rightarrow$  Longitudinally polarized hadron

## Spin-1/2 Hadrons - $e^+e^- \rightarrow \gamma^* \rightarrow hX$

$$\frac{d\sigma^{\text{em}}}{dzdy} = \frac{2\pi\alpha^2 e_q^2}{Q^2} \left\{ A(y)D_1(z) + \frac{4M}{zQ^2} B(y)D_T(z)\epsilon_{\perp}^{l_1 s_{\perp}} \right\}.$$

$$P_{Lh}(z, y) = 0$$

$$P_{hx}^{\text{em}}(z, y) = 0$$

$$P_{hy}^{\text{em}}(z, y) = \frac{4M}{zQ} \frac{\sqrt{y(1-y)}(1-2y)}{(1-y)^2 + y^2} \frac{\sum_q e_q^2 D_T^{q \rightarrow h}(z)}{\sum_q e_q^2 D_1^{q \rightarrow h}(z)}$$

$$MS_{\perp}^2 D_T(z) = \frac{z}{4} \sum_X \int \frac{p^+ d\xi^-}{2\pi} e^{-ip^+ \xi^- / z} \epsilon_{\perp}^{\alpha\gamma} S_{\perp\gamma} \text{Tr}[\gamma_{\alpha} \langle 0 | \mathcal{L}^{\dagger}(0, \infty) \psi(0) | hX \rangle \langle hX | \bar{\psi}(\xi^-) \mathcal{L}(\xi^-, \infty) | 0 \rangle]$$

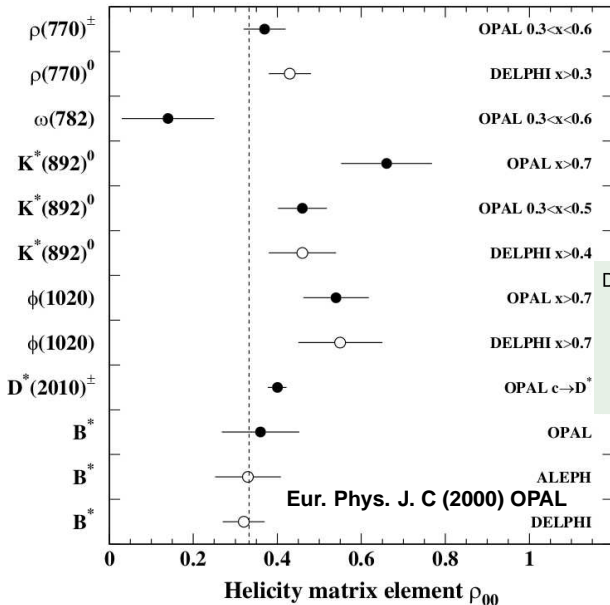
## Spin-1 Hadrons - $e^+e^- \rightarrow Z_0 \rightarrow hX$

$$\begin{aligned} \frac{d\sigma}{dzdy} = & \chi \frac{2\pi\alpha^2}{Q^2} \left\{ [T_0(y)D_1(z) + T_0(y)S_{LL}D_{1LL}(z) + \lambda_h T_1(y)\Delta D_{1L}(z)] \right. \\ & + \frac{4M}{zQ^2} [T_2(y)\epsilon_{\perp}^{l_1 S_{\perp}} D_T(z) + T_3(y)l_{\perp} \cdot S_{\perp} \Delta D_T(z)] \\ & \left. + \frac{4M}{zQ^2} [T_2(y)l_{\perp} \cdot S_{LT} D_{LT}(z) + T_3(y)\epsilon_{\perp}^{l_1 S_{LT}} \Delta D_{LT}(z)] \right\} \end{aligned}$$

Spin alignment

$$\rho_{00} = \frac{1}{3} - \frac{1}{3} \frac{\sum_q t_0^q D_{1LL}^{q \rightarrow h}(z)}{\sum_q t_0^q D_1^{q \rightarrow h}(z)}$$





$)]\Delta D_{1L}(z)]$

DELPHI Data for  $k^{*0}$

$0.1 \sim 0.3$	$0.33 \pm 0.05$
$z \geq 0.3$	$0.41 \pm 0.07$
$z \geq 0.4$	$0.46 \pm 0.08$
$z \geq 0.5$	$0.47 \pm 0.10$

PLB (1997) DELPHI

## Spin-1 Hadrons - $e^+e^- \rightarrow \gamma^* \rightarrow hX$

$$\frac{d\sigma^{\text{em}}}{dzdy} = \frac{2\pi\alpha^2 e_q^2}{Q^2} \left\{ A(y) [D_1(z) + S_{LL} D_{1LL}(z)] \right. \\ \left. + \frac{M}{zQ} \sqrt{y(1-y)} B(y) [|\vec{S}_\perp| \sin \phi_s D_T(z) - |\vec{S}_{LT}| \cos \phi_{LT} D_{LT}(z)] \right\}.$$

Spin alignment

$$\rho_{00}^{\text{em}} = \frac{1}{3} - \frac{1}{3} \frac{\sum_q e_q^2 D_{1LL}^{q \rightarrow h}(z)}{\sum_q e_q^2 D_1^{q \rightarrow h}(z)}.$$

## Semi-inclusive Process

$$\tilde{W}_{\mu\nu}^{(0)} = \frac{1}{2} \text{Tr} \left[ \hat{h}_{\mu\nu}^{(0)} \gamma_\alpha \right] \Xi^{(0)\alpha}(z_B, p, S, k'_\perp) + \frac{1}{2} \text{Tr} \left[ \hat{h}_{\mu\nu}^{(0)} \gamma_5 \gamma_\alpha \right] \tilde{\Xi}^{(0)\alpha}(z_B, p, S, k'_\perp)$$

$$\tilde{W}_{\mu\nu}^{(1L)} = -\frac{1}{4p \cdot q} \text{Tr} \left[ \hat{h}_{\mu\nu}^{(1)\rho} \gamma_\alpha \right] \omega_{\rho'}^{\rho'} \Xi_{\rho'}^{(1)\alpha}(z_B, p, S, k'_\perp) - \frac{1}{4p \cdot q} \text{Tr} \left[ \hat{h}_{\mu\nu}^{(1)\rho} \gamma_5 \gamma_\alpha \right] \omega_{\rho'}^{\rho'} \tilde{\Xi}_{\rho'}^{(1)\alpha}(z_B, p, S, k'_\perp)$$

$$\Xi^{(0)\alpha}(z, k_\perp, p, S) = \sum_X \int \frac{p^+ d\xi^- d^2\xi_\perp}{2\pi} e^{-ik^+\xi^- + ik_\perp \cdot \xi_\perp}$$

$$\text{Tr}[\gamma^\alpha \langle 0 | \mathcal{L}^\dagger(0, \infty) \psi(0) | hX \rangle \langle hX | \bar{\psi}(\xi) \mathcal{L}(\xi, \infty) | 0 \rangle]$$

$$\Xi_\rho^{(1)\alpha}(z, k_\perp, p, S) = \sum_X \int \frac{p^+ d\xi^- d^2\xi_\perp}{2\pi} e^{-ik^+\xi^- + ik_\perp \cdot \xi_\perp}$$

$$\text{Tr}[\gamma^\alpha \langle 0 | \mathcal{L}^\dagger(0, \infty) \hat{D}_\rho(0) \psi(0) | hX \rangle \langle hX | \bar{\psi}(\xi) \mathcal{L}(\xi, \infty) | 0 \rangle]$$

...

## New Lorentz Structures up to twist-3

Spin Independent

$$\begin{aligned}
 z\Xi_{\alpha}^{(0)}(z, k_{\perp}, p) &= k_{\perp\alpha} \hat{D}^{\perp}(z, k_{\perp}) \\
 z\tilde{\Xi}_{\alpha}^{(0)}(z, k_{\perp}, p) &= \epsilon_{\perp\alpha k_{\perp}} \Delta \hat{D}^{\perp}(z, k_{\perp}) \\
 z\Xi_{\rho\alpha}^{(1)}(z, k_{\perp}, p) &= p_{\alpha} k_{\perp\rho} \zeta_{\perp}^{(1)}(z, k_{\perp}) \\
 z\tilde{\Xi}_{\rho\alpha}^{(1)}(z, k_{\perp}, p) &= i p_{\alpha} \epsilon_{\perp\rho k_{\perp}} \tilde{\zeta}_{\perp}^{(1)}(z, k_{\perp})
 \end{aligned}$$

Spin Vector

$$\begin{aligned}
 z\Xi_{\alpha}^{(0)} &= p_{\alpha} \frac{\epsilon_{\perp}^{k_{\perp} S_{\perp}}}{M} \hat{D}_{1T}^{\perp} + k_{\perp\alpha} \frac{\epsilon_{\perp}^{k_{\perp} S_{\perp}}}{M} \hat{D}_T^{\perp} + \lambda_b \epsilon_{\perp\alpha k_{\perp}} \hat{D}_L^{\perp} \\
 z\tilde{\Xi}_{\alpha}^{(0)} &= p_{\alpha} \frac{k_{\perp} \cdot S_{\perp}}{M} \Delta \hat{D}_{1T}^{\perp} + \frac{\epsilon_{\perp}^{k_{\perp} S_{\perp}}}{M} \epsilon_{\perp\alpha k_{\perp}} \Delta \hat{D}_T^{\perp} + \lambda_b k_{\perp\alpha} \Delta \hat{D}_L^{\perp} \\
 z\Xi_{\rho\alpha}^{(1)} &= p_{\alpha} \left[ k_{\perp\rho} \frac{\epsilon_{\perp}^{k_{\perp} S_{\perp}}}{M} \zeta_T^{(1)\perp} + \lambda_b \epsilon_{\perp\rho k_{\perp}} \zeta_L^{(1)\perp} \right] \\
 z\tilde{\Xi}_{\rho\alpha}^{(1)} &= i p_{\alpha} \left[ \frac{\epsilon_{\perp}^{k_{\perp} S_{\perp}}}{M} \epsilon_{\perp\rho k_{\perp}} \tilde{\zeta}_T^{(1)\perp} + \lambda_b k_{\perp\rho} \tilde{\zeta}_L^{(1)\perp} \right]
 \end{aligned}$$

## New Lorentz Structures (Spin Tensor)

$$\begin{aligned}
z\Xi_{\alpha}^{(0)} &= p_{\alpha} \frac{S_{LT} \cdot k_{\perp}}{M} \hat{D}_{1LT}^{\perp} + p_{\alpha} \frac{k_{\perp\gamma} k_{\perp\delta} S_{TT}^{\gamma\delta}}{M^2} \hat{D}_{1TT}^{\perp} \\
&\quad + k_{\perp\alpha} S_{LL} \hat{D}_{LL}^{\perp} + k_{\perp\alpha} \frac{k_{\perp} \cdot S_{LT}}{M} \hat{D}_{LT}^{\perp} + S_{TT\alpha\beta} k_{\perp}^{\beta} \hat{D}_{TT}^{\perp A} + k_{\perp\alpha} \frac{k_{\perp\gamma} k_{\perp\delta} S_{TT}^{\gamma\delta}}{M^2} \hat{D}_{TT}^{\perp C} \\
z\check{\Xi}_{\alpha}^{(0)} &= p_{\alpha} \frac{\epsilon_{\perp}^{k_{\perp} S_{LT}}}{M} \Delta \hat{D}_{1LT}^{\perp} + p_{\alpha} \frac{\epsilon_{\perp k_{\perp\gamma} k_{\perp\delta} S_{TT}^{\gamma\delta}}}{M^2} \Delta \hat{D}_{1TT}^{\perp} \\
&\quad + \epsilon_{\perp\alpha k_{\perp}} \left[ S_{LL} \Delta \hat{D}_{LL}^{\perp} + \frac{k_{\perp} \cdot S_{LT}}{M} \Delta \hat{D}_{LT}^{\perp} + \frac{k_{\perp\gamma} k_{\perp\delta} S_{TT}^{\gamma\delta}}{M^2} \Delta \hat{D}_{TT}^{\perp C} \right] + \epsilon_{\perp\alpha\beta} S_{TT}^{\beta\gamma} k_{\perp\gamma} \Delta \hat{D}_{TT}^{\perp A} \\
z\Xi_{\rho\alpha}^{(1)} &= p_{\alpha} \left[ k_{\perp\rho} S_{LL} \xi_{LL}^{\perp} + k_{\perp\rho} \frac{k_{\perp} \cdot S_{LT}}{M} \xi_{LT}^{\perp} + S_{TT\rho\beta} k_{\perp}^{\beta} \xi_{TT}^{\perp A} + k_{\perp\rho} \frac{k_{\perp\gamma} k_{\perp\delta} S_{TT}^{\gamma\delta}}{M^2} \xi_{TT}^{\perp C} \right] \\
z\check{\Xi}_{\rho\alpha}^{(1)} &= i p^{\alpha} \left[ \epsilon_{\perp\rho k_{\perp}} S_{LL} \check{\xi}_{LL}^{\perp} + \epsilon_{\perp\rho k_{\perp}} k_{\perp} \cdot S_{LT} \check{\xi}_{LT}^{\perp} + \epsilon_{\perp\rho\beta} S_{TT}^{\beta\gamma} k_{\perp\gamma} \check{\xi}_{TT}^{\perp A} + \epsilon_{\perp\rho k_{\perp}} \frac{k_{\perp\gamma} k_{\perp\delta} S_{TT}^{\gamma\delta}}{M^2} \check{\xi}_{TT}^{\perp C} \right]
\end{aligned}$$

### New leading twist fragmentation functions (possibility density)

$\hat{D}_{1T}^\perp(z, k_\perp)$  quark  $\rightarrow$  transversely polarized hadron

$\Delta\hat{D}_{1T}^\perp(z, k_\perp)$  longitudinally polarized quark  $\rightarrow$  transversely polarized hadron

$\hat{D}_{1LT}^\perp(z, k_\perp)$  quark  $\rightarrow$  LT polarized hadron

$\Delta\hat{D}_{1LT}^\perp(z, k_\perp)$  longitudinally polarized quark  $\rightarrow$  LT polarized hadron

$\hat{D}_{1TT}^\perp(z, k_\perp)$  quark  $\rightarrow$  TT polarized hadron

$\Delta\hat{D}_{1TT}^\perp(z, k_\perp)$  longitudinally polarized quark  $\rightarrow$  TT polarized hadron

### New Twist-3 fragmentation functions

unpolarized hadron: 2;      polarization vector: 4;      polarization tensor: 8;

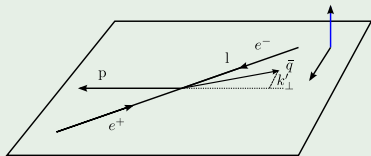
Spin-0 Hadrons -  $e^+e^- \rightarrow Z_0 \rightarrow h + \bar{q} + X$ 

$$\frac{d\sigma^{(\text{si,unp})}}{dydzd^2k'_\perp} = \frac{\alpha^2\chi}{2\pi Q^2} \left\{ T_0^q(y)\hat{D}_1(z,k'_\perp) + \frac{4}{zQ^2} \right. \\ \left. [T_2^q(y)l_\perp \cdot k'_\perp \hat{D}^\perp(z,k'_\perp) + T_3^q(y)\epsilon_\perp^{l_1 k'_\perp} \Delta \hat{D}^\perp(z,k'_\perp)] \right\}$$

Azimuthal Asymmetries,

$$A_{\text{unp}}^{\cos\varphi} = -\frac{2|\vec{k}'_\perp|}{zQ} \frac{\sum_q \tilde{T}_2^q(y)\hat{D}^{\perp q \rightarrow b}}{\sum_q T_0^q(y)\hat{D}_1^{q \rightarrow b}}$$

$$A_{\text{unp}}^{\sin\varphi} = \frac{2|\vec{k}'_\perp|}{zQ} \frac{\sum_q \tilde{T}_3^q(y)\Delta \hat{D}^{\perp q \rightarrow b}}{\sum_q T_0^q(y)\hat{D}_1^{q \rightarrow b}}$$



Spin-0 Hadrons -  $e^+e^- \rightarrow \gamma^* \rightarrow h + \bar{q} + X$ 

$$\frac{d\sigma^{(\text{si,unp,em})}}{dydzd^2k'_\perp} = \frac{\alpha^2 e_q^2}{2\pi Q^2} \left\{ A(y) \hat{D}_1(z, k'_\perp) + \frac{4l_\perp \cdot k'_\perp}{zQ^2} B(y) \hat{D}^\perp(z, k'_\perp) \right\}$$

Azimuthal Asymmetries,

$$A_{\text{unp,em}}^{\cos\varphi}(z, y, k'_\perp) = -\frac{2|\vec{k}'_\perp| \tilde{B}(y) \sum_q e_q^2 D^{\perp q \rightarrow h}(z, k'_\perp)}{zQ A(y) \sum_q e_q^2 D_1^{q \rightarrow h}(z, k'_\perp)}$$

$$A_{\text{unp,em}}^{\sin\varphi}(z, y, k'_\perp) = 0$$



Spin-1/2 Hadrons -  $e^+e^- \rightarrow Z_0 \rightarrow h + \bar{q} + X$ 

$$\frac{d\sigma^{(si,1/2)}}{dydzd^2k'_{\perp}} = \frac{d\sigma^{(si,unp)}}{dydzd^2k'_{\perp}} + \frac{d\sigma^{(si,Vpol)}}{dydzd^2k'_{\perp}}$$

$$\begin{aligned} \frac{d\sigma^{(si,Vpol)}}{dydzd^2k'_{\perp}} = & \frac{\alpha^2 \chi}{2\pi Q^2} \left\{ T_0^q(y) \frac{\epsilon_{\perp}^{k'_{\perp} S_{\perp}}}{M} \hat{D}_{1T}^{\perp}(z, k'_{\perp}) + T_1^q(y) [\lambda_b \Delta \hat{D}_{1L}^{\perp}(z, k'_{\perp}) + \frac{k'_{\perp} \cdot S_{\perp}}{M} \Delta \hat{D}_{1T}^{\perp}(z, k'_{\perp})] \right. \\ & + \frac{4\lambda_b}{zQ^2} [T_2^q(y) \epsilon_{\perp}^{l_{\perp} k'_{\perp}} \hat{D}_L^{\perp}(z, k'_{\perp}) + T_3^q(y) l_{\perp} \cdot k'_{\perp} \Delta \hat{D}_L^{\perp}(z, k'_{\perp})] \\ & + \frac{4\epsilon_{\perp}^{k'_{\perp} S_{\perp}}}{zMQ^2} [T_2^q(y) l_{\perp} \cdot k'_{\perp} \hat{D}_T^{\perp}(z, k'_{\perp}) + T_3^q(y) \epsilon_{\perp}^{l_{\perp} k'_{\perp}} \Delta \hat{D}_T^{\perp}(z, k'_{\perp})] \\ & \left. + \frac{4M}{zQ^2} [T_2^q(y) \epsilon_{\perp}^{l_{\perp} S_{\perp}} \hat{D}_T(z, k'_{\perp}) + T_3^q(y) l_{\perp} \cdot S_{\perp} \Delta \hat{D}_T(z, k'_{\perp})] \right\}. \end{aligned}$$

## Polarization

$$P_{Lh}(y, z, k'_\perp) = \frac{\sum_q T_0^q(y) P_q(y) \Delta \hat{D}_{1L}(z, k'_\perp)}{\sum_q T_0^q(y) \hat{D}_1(z, k'_\perp)}$$

$$P_{hn}(y, z, k'_\perp) = - \frac{|\vec{k}'_\perp|}{M} \frac{\sum_q T_0^q(y) \hat{D}_{1T}^\perp(z, k'_\perp)}{\sum_q T_0^q(y) \hat{D}_1(z, k'_\perp)}$$

$$P_{ht}(y, z, k'_\perp) = - \frac{|\vec{k}'_\perp|}{M} \frac{\sum_q P_q(y) T_0^q(y) \Delta \hat{D}_{1T}^\perp(z, k'_\perp)}{\sum_q T_0^q(y) \hat{D}_1(z, k'_\perp)}$$

## Polarization

$$P_{Lh}(y, z, k'_\perp) = \frac{\sum_q T_0^q(y) P_q(y) \Delta \hat{D}_{1L}(z, k'_\perp)}{\sum_q T_0^q(y) \hat{D}_1(z, k'_\perp)} \left[ 1 + \frac{M}{Q} \hat{\Delta}(y, z, k'_\perp) \right] \\ + \frac{4}{zQ} \frac{\sum_q \left[ \tilde{T}_2^q(y) k'_y \hat{D}_L^\perp(z, k'_\perp) - \tilde{T}_3^q(y) k'_x \Delta \hat{D}_L^\perp(z, k'_\perp) \right]}{\sum_q T_0^q(y) \hat{D}_1(z, k'_\perp)}$$

$$P_{hn}(y, z, k'_\perp) = - \frac{|\vec{k}'_\perp|}{M} \frac{\sum_q T_0^q(y) \hat{D}_{1T}^\perp(z, k'_\perp)}{\sum_q T_0^q(y) \hat{D}_1(z, k'_\perp)} + \dots$$

$$P_{ht}(y, z, k'_\perp) = - \frac{|\vec{k}'_\perp|}{M} \frac{\sum_q P_q(y) T_0^q(y) \Delta \hat{D}_{1T}^\perp(z, k'_\perp)}{\sum_q T_0^q(y) \hat{D}_1(z, k'_\perp)} + \dots$$

Spin-1/2 Hadrons -  $e^+e^- \rightarrow \gamma^* \rightarrow h + \bar{q} + X$ 

$$\frac{d\sigma^{(\text{si, Vpol, em})}}{dydzd^2k'_\perp} = \frac{\alpha^2 e_q^2}{2\pi Q^2} \left\{ A(y) \frac{\epsilon_\perp^{k'_\perp S_\perp}}{M} \hat{D}_{1T}^\perp(z, k'_\perp) \right. \\ \left. + \frac{4B(y)}{zMQ^2} \left[ \lambda_b M \epsilon_\perp^{l_\perp k'_\perp} \hat{D}_L^\perp(z, k'_\perp) + \epsilon_\perp^{k'_\perp S_\perp} l_\perp \cdot k'_\perp \hat{D}_T^\perp(z, k'_\perp) + M^2 \epsilon_\perp^{l_\perp S_\perp} \hat{D}_T(z, k'_\perp) \right] \right\}$$

Polarization (Leading Twist)

$$P_{\text{Lh}}^{(0)(\text{em})}(y, z, k'_\perp) = P_{\text{ht}}^{(0)(\text{em})}(y, z, k'_\perp) = 0, \quad P_{\text{hn}}^{(0)(\text{em})}(y, z, k'_\perp) = -\frac{|\vec{k}'_\perp|}{M} \frac{\sum_q e_q^2 \hat{D}_{1T}^\perp(z, k'_\perp)}{\sum_q e_q^2 \hat{D}_1(z, k'_\perp)}$$

Twist-3

$$P_{\text{Lh}}^{(\text{em})}(y, z, k'_\perp) = \frac{4k'_y}{zQ} \frac{\tilde{B}(y)}{A(y)} \frac{\sum_q e_q^2 \hat{D}_L^\perp(z, k'_\perp)}{\sum_q e_q^2 \hat{D}_1(z, k'_\perp)}$$

Spin-1 Hadrons -  $e^+e^- \rightarrow Z_0 \rightarrow h + \bar{q} + X$ 

$$\frac{d\sigma^{(si,1)}}{dydzd^2k'_\perp} = \frac{d\sigma^{(si,unp)}}{dydzd^2k'_\perp} + \frac{d\sigma^{(si,Vpol)}}{dydzd^2k'_\perp} + \frac{d\sigma^{(si,Tpol)}}{dydzd^2k'_\perp}$$

$$\frac{d\sigma^{(si,Tpol)}}{dydzd^2k'_\perp} = \frac{d\sigma^{(si,LL)}}{dydzd^2k'_\perp} + \frac{d\sigma^{(si,LT)}}{dydzd^2k'_\perp} + \frac{d\sigma^{(si,TT)}}{dydzd^2k'_\perp}$$

$$[\mathbf{LL}] = \frac{\alpha^2 \chi}{2\pi Q^2} S_{LL} \left\{ T_0^q(y) \hat{D}_{1LL}(z, k'_\perp) + \frac{4}{zQ^2} [T_2^q(y)(l_\perp \cdot k'_\perp) \hat{D}_{LL}^\perp(z, k'_\perp) + T_3^q(y) \epsilon_\perp^{l_\perp k'_\perp} \Delta \hat{D}_{LL}^\perp(z, k'_\perp)] \right\}$$

## Spin Alignment

$$\rho_{00} = \frac{1}{3} - \frac{1}{3} \frac{\sum_q T_0^q(y) \hat{D}_{1LL}(z, k'_\perp)}{\sum_q T_0^q(y) \hat{D}_1(z, k'_\perp)} \left[ 1 + \frac{M}{Q} \hat{\Delta}(y, z, k'_\perp) \right]$$

$$- \frac{4}{3} \frac{\sum_q \left[ \tilde{T}_2^q(y) k'_x \hat{D}_{LL}^\perp(z, k'_\perp) - \tilde{T}_3^q(y) k'_y \Delta \hat{D}_{LL}^\perp(z, k'_\perp) \right]}{zQ \sum_q T_0^q(y) \hat{D}_1(z, k'_\perp)}$$

$$\begin{aligned}
 [\text{LT Terms}] = & \frac{\alpha^2 \chi}{2\pi Q^2} S_{LT}^\alpha \left\{ T_0^q(y) \frac{k'_{\perp\alpha}}{M} \hat{D}_{1LT}^\perp(z, k'_\perp) + T_1^q(y) \frac{\epsilon_{\perp k'\alpha}}{M} \Delta \hat{D}_{1LT}^\perp(z, k'_\perp) \right. \\
 & + \frac{4}{zQ^2} T_2^q(y) [(l_\perp \cdot k'_\perp) \frac{k'_{\perp\alpha}}{M} \hat{D}_{LT}^\perp(z, k'_\perp) + M l_{\perp\alpha} \hat{D}_{LT}(z, k'_\perp)] \\
 & \left. + \frac{4}{zQ^2} T_3^q(y) [\epsilon_{\perp}^{l_1 k'_\perp} \frac{k'_{\perp\alpha}}{M} \Delta \hat{D}_{LT}^\perp(z, k'_\perp) + M \epsilon_{\perp l\alpha} \Delta \hat{D}_{LT}(z, k'_\perp)] \right\}
 \end{aligned}$$

$$\begin{aligned}
 [\text{TT Terms}] = & \frac{\alpha^2 \chi}{2\pi Q^2} S_{TT}^{\alpha\beta} \left\{ T_0^q(y) \frac{k'_{\perp\alpha} k'_{\perp\beta}}{M^2} \hat{D}_{1TT}^\perp(z, k'_\perp) + T_1^q(y) \frac{\epsilon_{\perp k'\alpha} k'_{\perp\beta}}{M^2} \Delta \hat{D}_{1TT}^\perp(z, k'_\perp) \right. \\
 & + \frac{4}{zQ^2} T_2^q(y) [(l_\perp \cdot k'_\perp) \frac{k'_{\perp\alpha} k'_{\perp\beta}}{M^2} \hat{D}_{TT}^{\perp C}(z, k'_\perp) + l_{\perp\alpha} k'_{\perp\beta} \hat{D}_{TT}^{\perp A}(z, k'_\perp)] \\
 & \left. + \frac{4}{zQ^2} T_3^q(y) [\epsilon_{\perp}^{l_1 k'_\perp} \frac{k'_{\perp\alpha} k'_{\perp\beta}}{M^2} \Delta \hat{D}_{TT}^{\perp C}(z, k'_\perp) + \epsilon_{\perp l\alpha} k'_{\perp\beta} \Delta \hat{D}_{TT}^{\perp A}(z, k'_\perp)] \right\}
 \end{aligned}$$



**Inclusive**  $e^+e^- \rightarrow Z_0 \rightarrow h + X$

- There is a leading twist longitudinal polarization for spin-1/2 hadrons and also spin alignment ( $\rho_{00} \neq 1/3$ ) for vector mesons.
- On the twist-3 level, there are transverse polarizations for spin-1/2 hadrons that in and perpendicular to the leptonic plane.

**Semi-inclusive**  $e^+e^- \rightarrow Z_0 \rightarrow h + \bar{q}(jet) + X$

- For spin-0 hadrons, there are two azimuthal asymmetries on twist-3 level.
- For spin-1/2 hadrons, there is a longitudinal polarization and also transverse polarizations that in and transverse to the production plane on the leading twist.
- For vector mesons, all five tensor polarizations have leading twist contributions.

$$S_{LL}, S_{LT}^n, S_{LT}^t, S_{TT}^{nn} \text{ and } S_{TT}^{nt}$$

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$$S_{LL}, S_{LT}^n, S_{LT}^t, S_{TT}^{nn} \text{ and } S_{TT}^{nt}$$

**Thanks for your attention!**