

Review of Heavy Meson Spectroscopy

Feng-Kun Guo

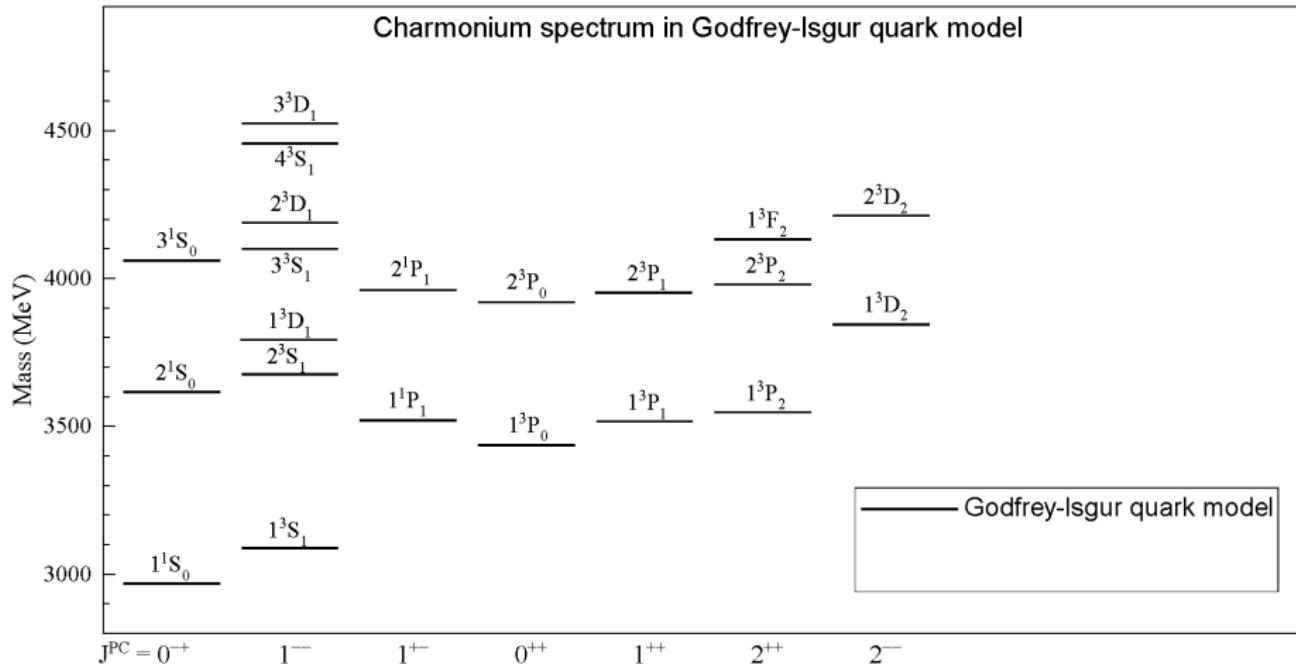
Institute of Theoretical Physics, Chinese Academy of Sciences

The 7th Asia-Pacific Conference on Few-Body Problems in Physics

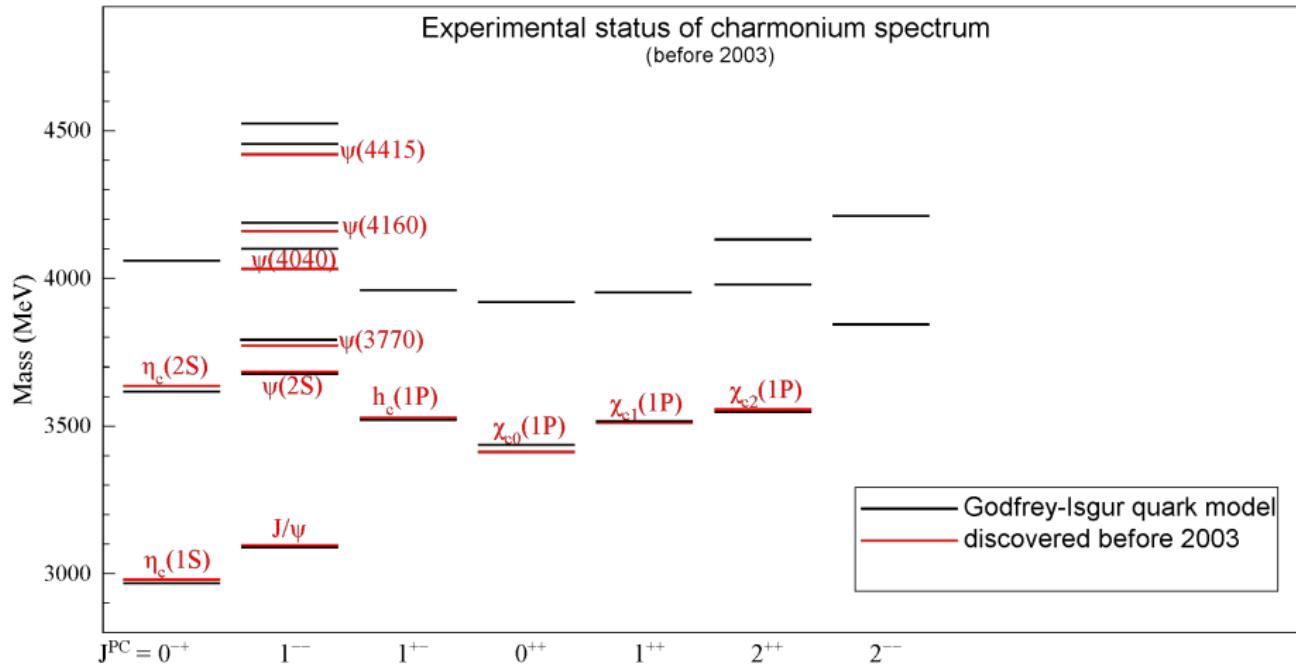
Guilin, August 25 – 30, 2017

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- Charmonium(-like) structures and heavy-strange mesons
 - Kinematic effects versus genuine resonances
 - Examples: $Z_c(3900)$, $X(3872)$, $D_{s0}^*(2317)$

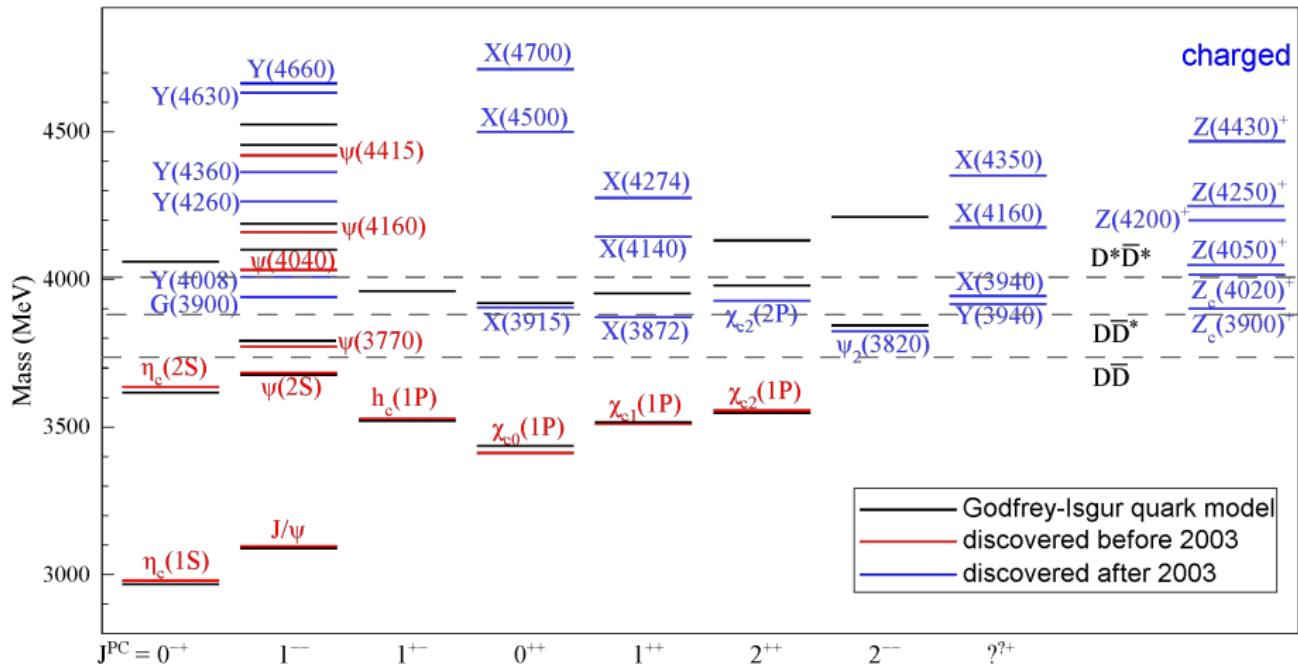
Charmonium spectrum: too many structures



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Heavy-strange mesons

Charm-strange $I = 1$ $c\bar{s}$ mesons $D_{s0}^*(2317)$ and $D_{s1}(2460)$

- $D_{s0}^*(2317)$: 0^+ BaBar (2003)

$M = (2317.7 \pm 0.6) \text{ MeV}$,

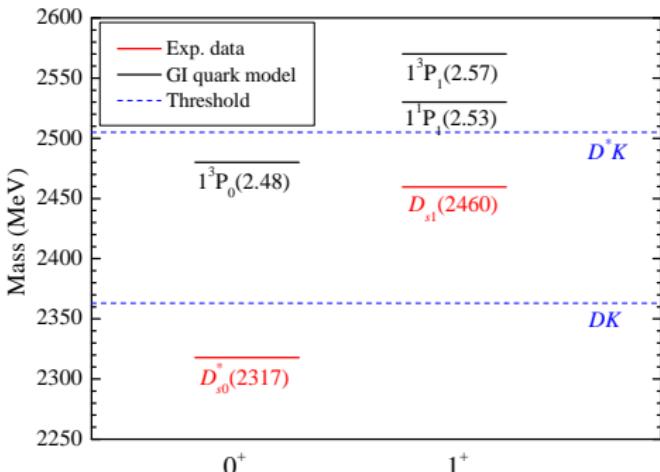
$\Gamma < 3.8 \text{ MeV}$

The only hadronic decay: $D_s\pi$

- $D_{s1}(2460)$: 1^+ CLEO (2003)

$M = (2459.5 \pm 0.6) \text{ MeV}$,

$\Gamma < 3.5 \text{ MeV}$



- Notable features:

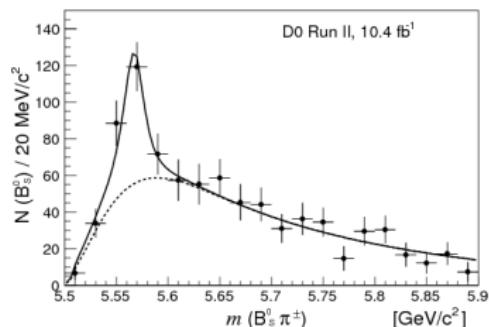
☞ masses are much lower than the quark model predictions for $c\bar{s}$ mesons

☞ $M_{D_{s1}(2460)} - M_{D_{s0}^*(2317)} \simeq M_{D^*} - M_D + 1 \text{ MeV}$

Observation in 2016: an “explicitly exotic” multiquark candidate

- $X(5568)$ by D0 Collaboration ($p\bar{p}$ collisions)

PRL117(2016)022003



$$M = (5567.8 \pm 2.9_{-1.9}^{+0.9}) \text{ MeV}$$

$$\Gamma = (21.9 \pm 6.4_{-2.5}^{+5.0}) \text{ MeV}$$

$B_s^0 \pi^+$: minimal quark contents is $\bar{b} s \bar{d} u$!

☞ difficulties in all possible structure explanations

Burns, Swanson, PLB760(2016)627; FKG, Meißner, Zou, Commun.Theor.Phys. 65 (2016) 593

might be due to kinematic cuts in analysis Yang, Wang, Meißner, PLB767(2017)470

☞ immediately, negative result by LHCb

PRL117(2016)152003

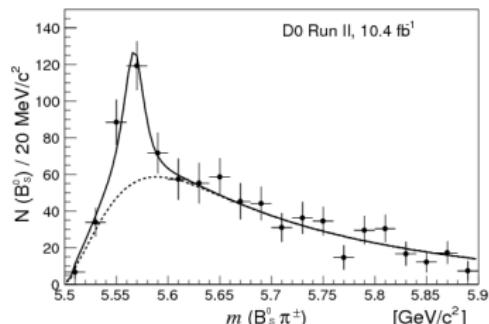
and by CMS

CMS-PAS-BPH-16-002

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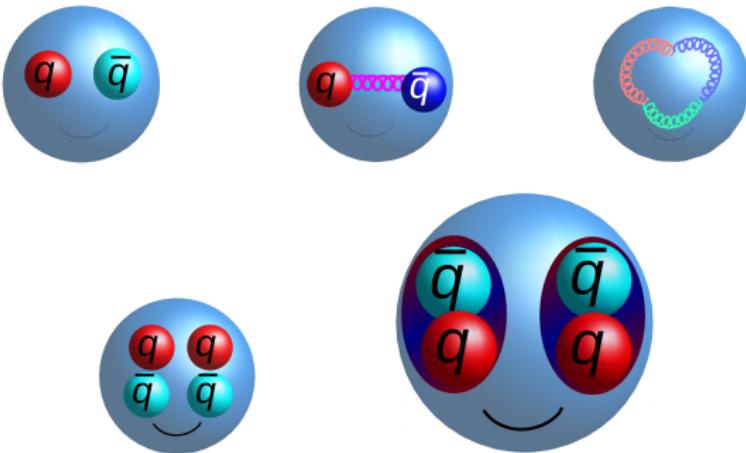
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Interpretations

Always many models for each observed structure:

- **Dynamics** \Rightarrow poles in the S -matrix: genuine physical states. The origins of the poles can be different:

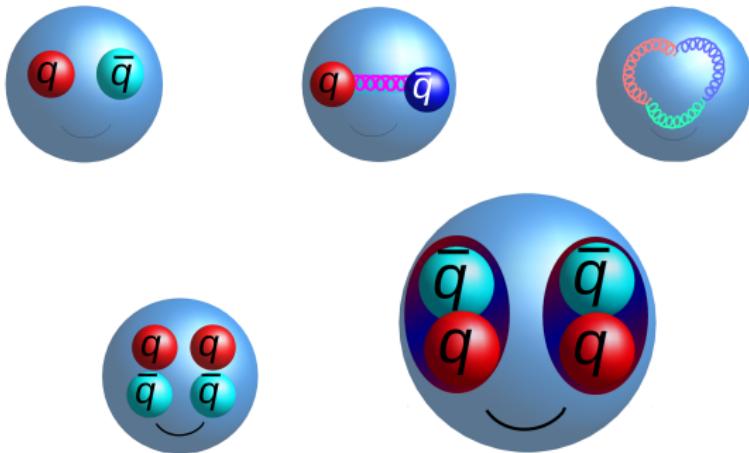


- **Kinematic** effects \Rightarrow branching points of S -matrix
 - normal two-body threshold cusp
 - triangle singularity
 - ...

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 - ☞ triangle singularity
 - ☞ ...

Threshold cusp

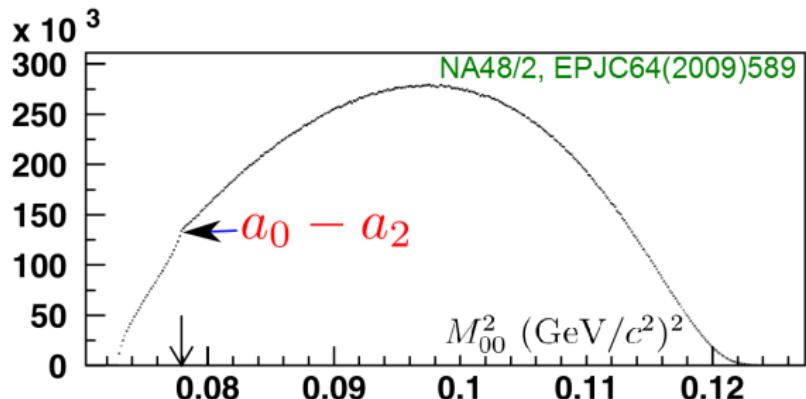
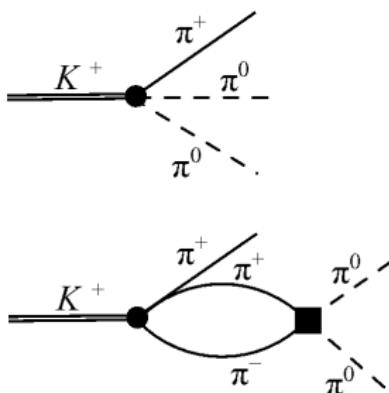
- There is **always** a cusp at an *S*-wave threshold

- Cusp effect has been well-known for a long time:

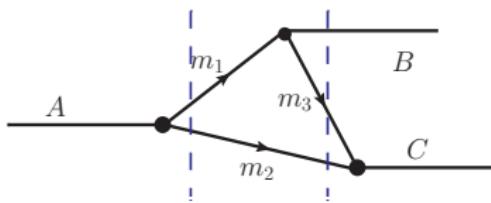
☞ example of the cusp in $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$

☞ strength of the cusp measures the **interaction strength!**

Meißner, Müller, Steininger (1997); Cabibbo (2004); Colangelo, Gasser, Kubis, Rusetsky (2006); ...

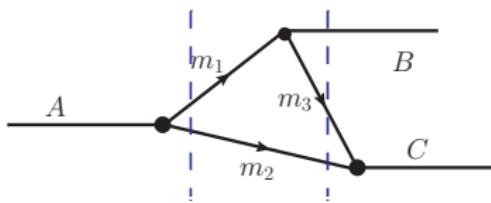


Triangle singularity

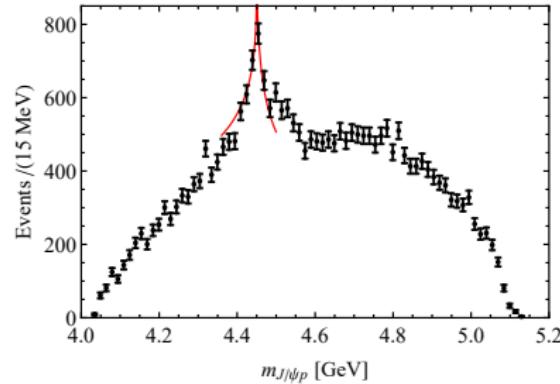


- Conditions (Coleman–Norton theorem (1965)):
 - ☒ all three intermediate particles can go **on shell simultaneously**
 - ☒ $\vec{p}_2 \parallel \vec{p}_3$, particle-3 can catch up with particle-2 (**as a classical process**)
- requires very special kinematics
⇒ **process dependent!**
- **S-wave TS** can produce a **narrow peak**
mimicking a resonance
Bayar et al., PRD94(2016)074039
- many recent applications
reviews: Q.Zhao, JPS Conf.Proc.13(2017)010008;
FKG et al., arXiv:1705.00141

Triangle singularity



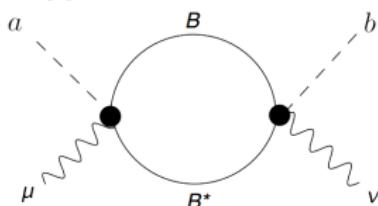
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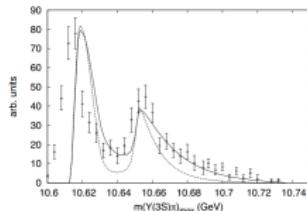
Threshold cusp — example: Z_b/Z_c

- Models of $Z_b(10610, 10650)$, $Z_c(3900, 4020)$ as threshold cusps

☞ Bugg, Swanson:



D. Bugg, EPL96(2011)11002; E. Swanson, PRD91(2015)034009

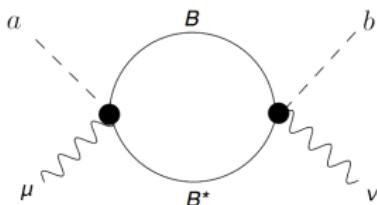


☞ Initial pion radiation: D.-Y.Chen, X.Liu, PRD84(2011)094003; PRD84(2011)034032; Chen, Liu, Matsuki, PRD84(2011)074032; PRL110(2013)232001; ...

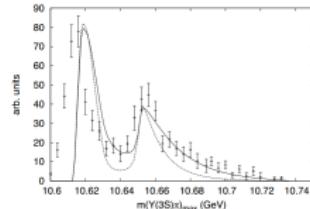
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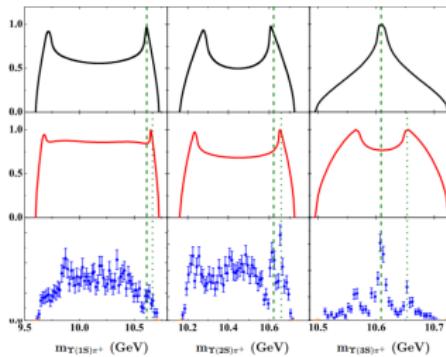
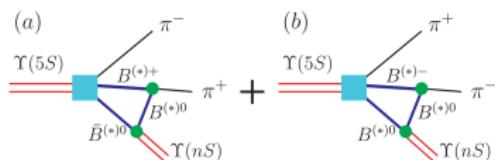
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$\Leftarrow B\bar{B}^* + c.c.$ loops

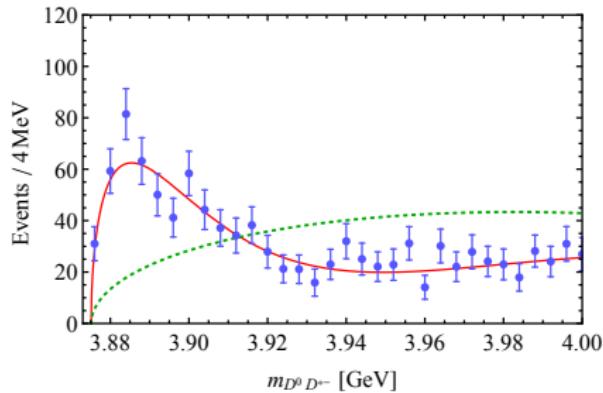
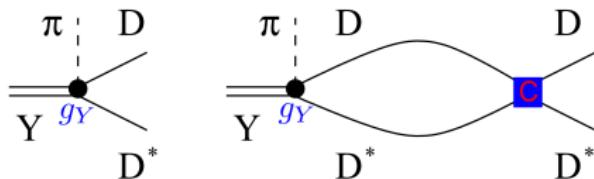
$\Leftarrow B^*\bar{B}^*$ loops

\Leftarrow Belle data

Threshold cusp — example: Z_b/Z_c

- But $Z_c(3900)[Z_b]$ as a narrow peak in $D\bar{D}^*[B\bar{B}^*]$ distribution cannot be only due to cusp: prominent cusp \Rightarrow strong int. \Rightarrow pole!

FKG, Hanhart, Wang, Zhao, PRD91(2015)051504



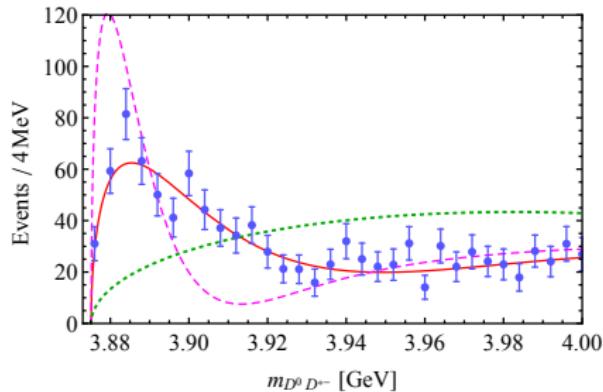
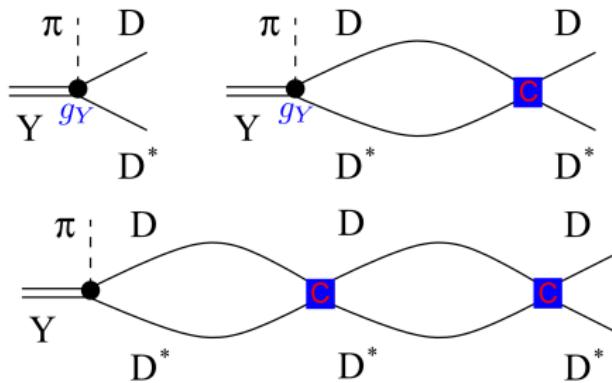
Black curve: up to 1 loop with $C_\Lambda G_\Lambda(E_{\text{th}}) = -1/2$,
no narrow peak any more!

$g_Y [1 + C_\Lambda G_\Lambda(E) + C_\Lambda G_\Lambda(E)C_\Lambda G_\Lambda(E) + \dots]$ produces a pole

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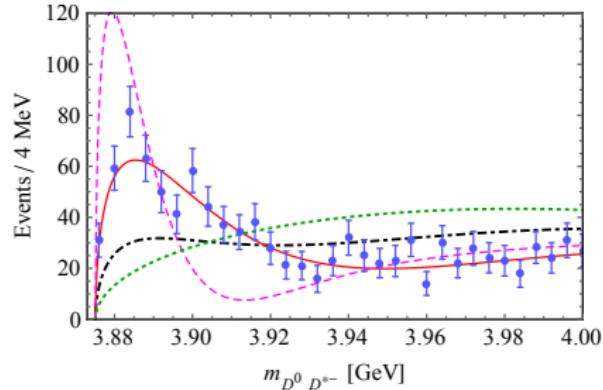
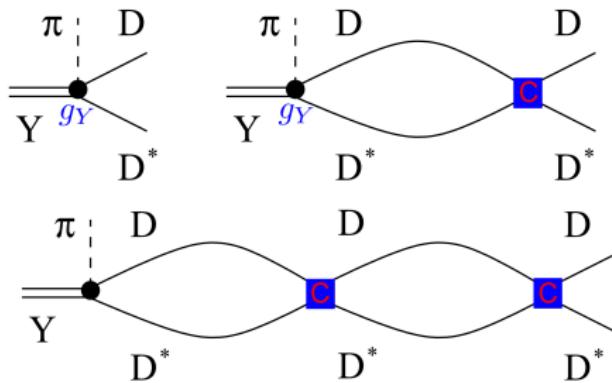
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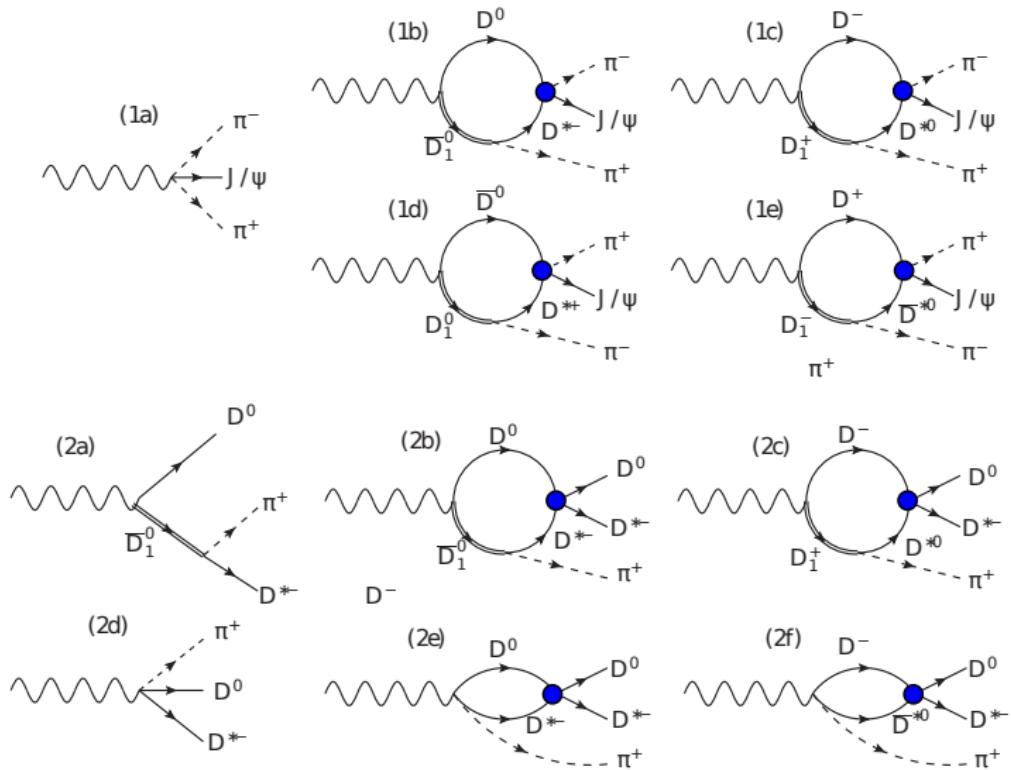


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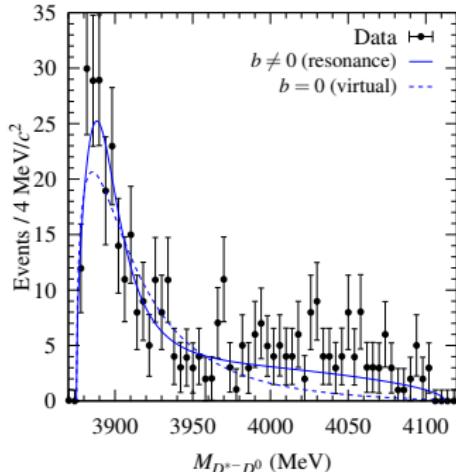
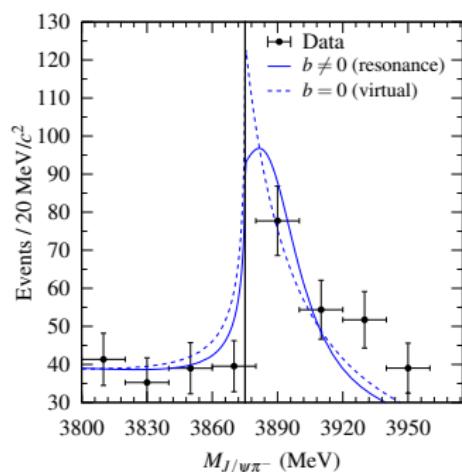
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More about $Z_c(3900)$

Albaladejo, FKG, Hidalgo-Duque, Nieves, PLB755(2016)337



More about $Z_c(3900)$



M_{Z_c} (MeV)	$\Gamma_{Z_c}/2$ (MeV)	Ref.	Final state
3899 ± 6	23 ± 11	[1] (BESIII)	$J/\psi \pi$
3895 ± 8	32 ± 18	[2] (Belle)	$J/\psi \pi$
3886 ± 5	19 ± 5	[3] (CLEO-c)	$J/\psi \pi$
3884 ± 5	12 ± 6	[4] (BESIII)	$\bar{D}^* D$
3882 ± 3	13 ± 5	[5] (BESIII)	$\bar{D}^* D$
$3894 \pm 6 \pm 1$	$30 \pm 12 \pm 6$	$\Lambda_2 = 1.0$ GeV	$J/\psi \pi, \bar{D}^* D$
$3886 \pm 4 \pm 1$	$22 \pm 6 \pm 4$	$\Lambda_2 = 0.5$ GeV	$J/\psi \pi, \bar{D}^* D$
$3831 \pm 26^{+7}_{-28}$	virtual state	$\Lambda_2 = 1.0$ GeV	$J/\psi \pi, \bar{D}^* D$
$3844 \pm 19^{+12}_{-21}$	virtual state	$\Lambda_2 = 0.5$ GeV	$J/\psi \pi, \bar{D}^* D$

resonance pole

or virtual state

new BESIII data on $J/\psi\pi^+\pi^-$ channel published (next slide), needs to be updated

Triangle singularity — example: $Y(4260) \rightarrow Z_c\pi$

- Importance of TS in $Y(4260) \rightarrow Z_c\pi$ already noticed, but Z_c pole still needed

Q.Wang, Hanhart, Q.Zhao, PRL111(2013)132002; PLB725(2013)106

- however, debate continues:

☞ opposite claim: whether Z_c pole is needed is inconclusive

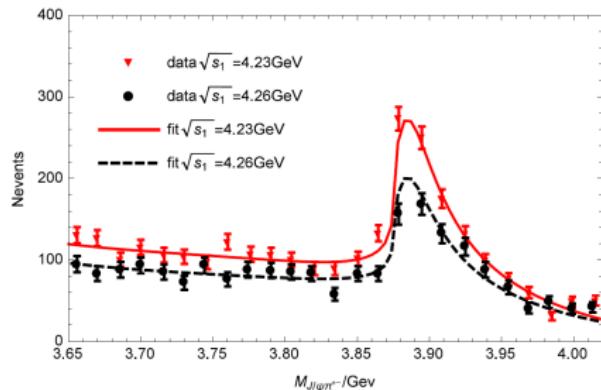
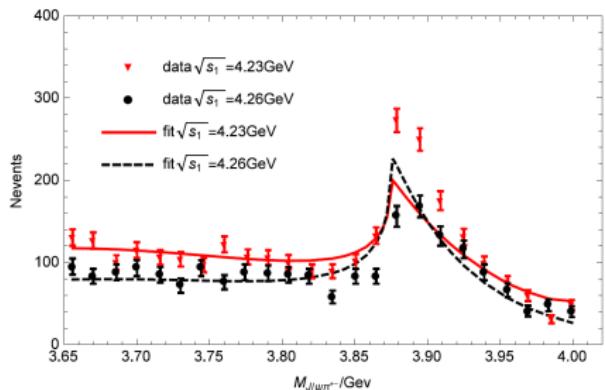
Pilloni et al. (JPAC), PLB772(2017)200

☞ updated combined analysis of $e^+e^- \rightarrow J/\psi\pi\pi$ and $e^+e^- \rightarrow (D\bar{D}^*)^\mp\pi^\pm$

⇒ necessity of Z_c pole (virtual or resonance) Albaladejo et al., PLB755(2016)337

TS only not favored by using the latest data of BESIII, PRL119(2017)072001

Q.-R. Gong et al., arXiv:1612.08159



$Z_c(3900)$: negative lattice results

- So far, no evidence for $Z_c(3900)$ in lattice QCD:

talk by Y. Chen

CLQCD:

PRD89(2014)094506

$I = 1 D\bar{D}^*$ weakly repulsive \Rightarrow no bound state ($M_\pi \geq 300$ MeV)

Prelovsek et al.:

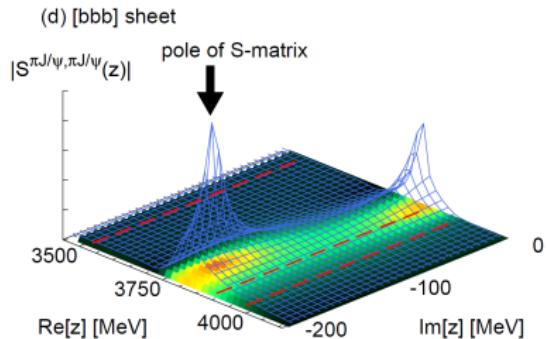
PRD91(2015)014504

"no additional eigenstate" corresponding to $Z_c(3900)$ ($M_\pi = 266$ MeV),

HALQCD: PRL117(2017)242001

virtual state pole with a very low mass and deep in the complex plane ($M_\pi \geq 410$ MeV)

Ikeda for HALQCD, arXiv:1706.07300

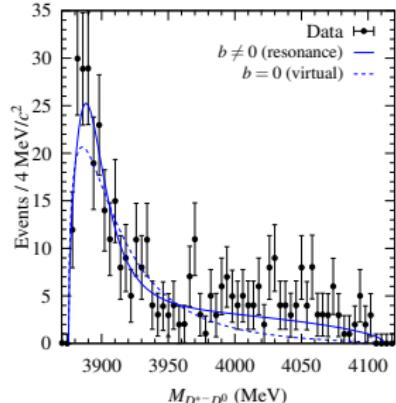
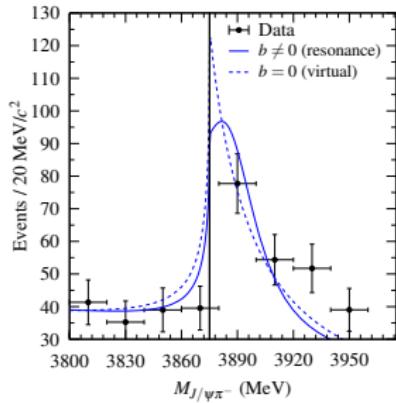


- Are they in conflict with experiments?

For HALQCD: recall the virtual state pole in Albaladejo et al., PLB755(2016)337 is much closer to the threshold

$Z_c(3900)$: Interpreting lattice results by Prelovsek et al.

Albaladejo, Fernandez-Soler, Nieves, EPJC76(2016)573

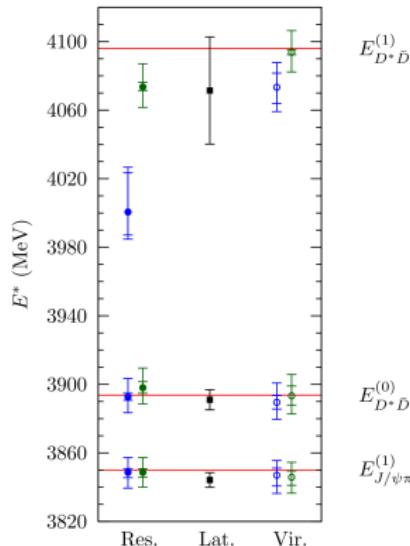
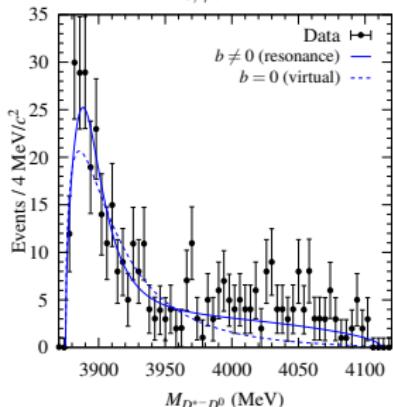
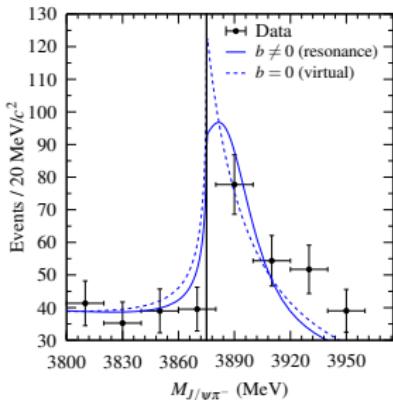


- Model fitted to BESIII data with: (1) resonance, or (2) virtual state
- In finite volume ($L = 2$ fm): consistent with lattice energy levels, but **with a pole in continuum!**

Albaladejo et al., PLB755(2016)337

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Albaladejo et al., PLB755(2016)337

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$X(3872)$ (I)

- Very important features:

$$M_{D^0} + M_{\bar{D}^{*0}} - M_X = (0.00 \pm 0.18) \text{ MeV}; \quad \text{Br}(X \rightarrow D\bar{D}^{*0} + c.c.) > 24\%$$

- Many models:

☞ hadronic molecule

$D\bar{D}^*$ bound state Törnquist (2003); Voloshin (2004); Braaten (2004); Swanson (2004); ...

virtual state Hanhart et al. (2007)

☞ $c\bar{c} + D\bar{D}^*$ coupled-channel effects

Kalashnikova (2005); Meng, Gao, Chao (2005); Zhang, Meng, Zheng (2009); Li, Chao (2009);

Danilkin, Simonov (2010); Zhou, Xiao (2014); ...

if large coupling to $D\bar{D}^*$ \Rightarrow a large $D\bar{D}^*$ component

☞ tetraquark Maiani et al. (2005); ...

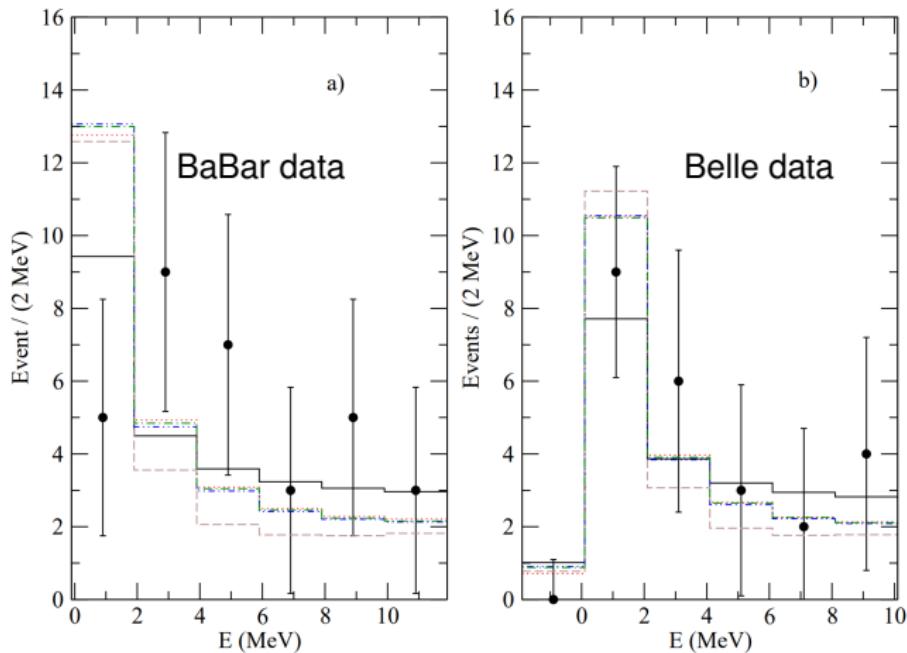
generally predicting too many states

☞ cusp, $c\bar{c}g$, ...: not under active discussion any more

$X(3872)$ (II)

- Very precise measurements in the $X \rightarrow D^0 \bar{D}^0 \pi^0$ channel needed:
current data allow for several possible scenarios with the compositeness ranging from nearly 0 (blue, green, weak coupling to $D\bar{D}^*$) to 1 (black solid, strong coupling to $D\bar{D}^*$)

X.-W. Kang, J. A. Oller, EPJC77(2017)399



$D_{s0}^*(2317)$

- Early studies using only $c\bar{s}$ -type interpolators typically give mass larger than that for $D_{s0}^*(2317)$

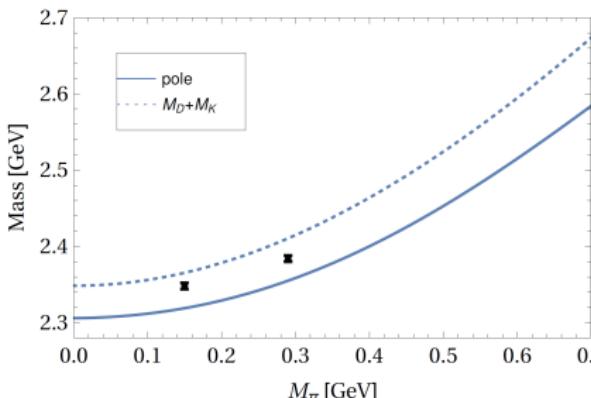
Bali (2003); UKQCD (2003); ...

- $c\bar{s} + DK$ interpolators:

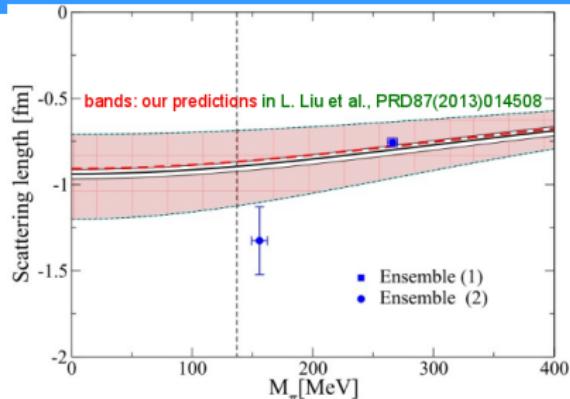
Mohler et al., PRL111(2013)222001; Lang et al., PRD90(2014)034510

- Latest lattice results

Bali et al. (RQCD), arXiv:1706.01247



Exp: to measure the width of $D_{s0}^*(2317) \rightarrow D_s \pi$ precisely

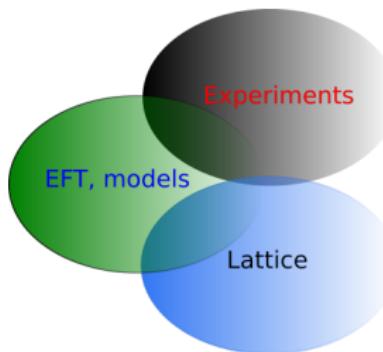


compositeness: = 1.04(0.08)(+0.30)

M_π [MeV]	150	290
$M_{D_{s0}^*(2317)}$ [MeV]	2348 ± 4	2384 ± 3
M_{D_s} [MeV]	1977 ± 1	1980 ± 1

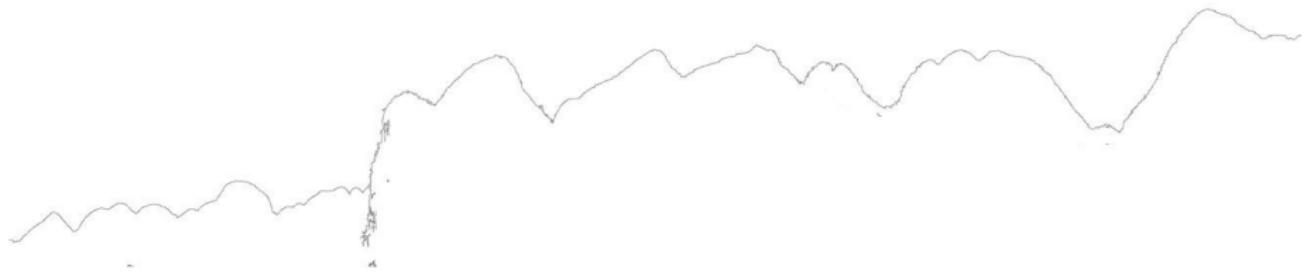
curves: prediction in Du et al., arXiv:1703.10836

Not a summary

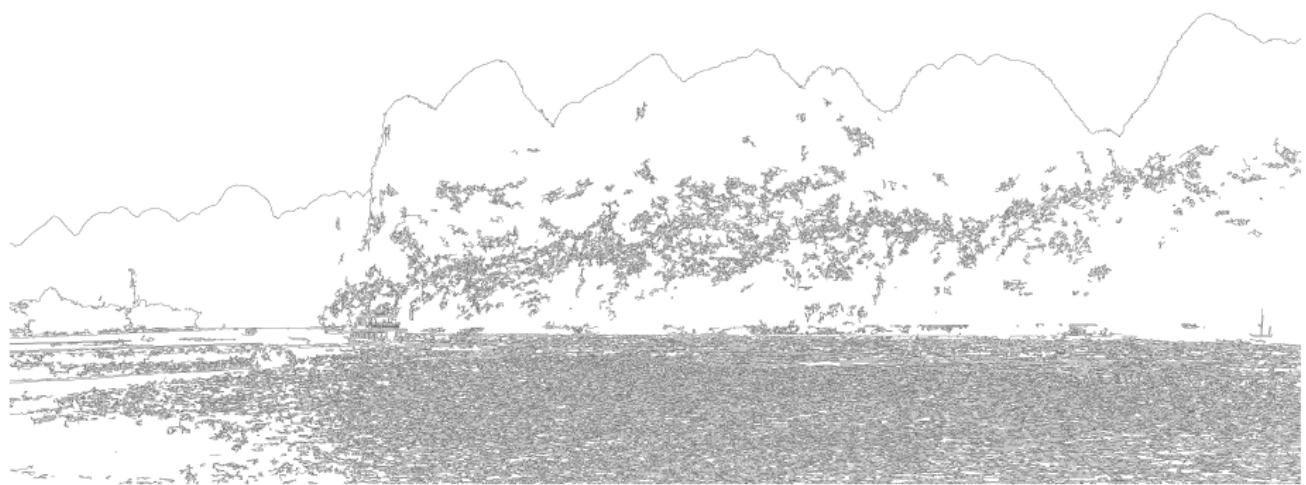


- The study of exotic hadrons is difficult: nonperturbative QCD, confinement
 - ☞ Why are exotic hadrons so scarce?
 - ☞ Searching for and confirming states with exotic quantum numbers
 - ☞ Calculating QCD spectrum using lattice simulations
 - ☞ For the confirmed states: understanding their structures,
why is the spectrum organized as such?
⇒ learning about confinement
- lots of progress in recent years, but still a long way to go
⇒ more joint efforts needed !

Another example of peaking structures



Another example of peaking structures



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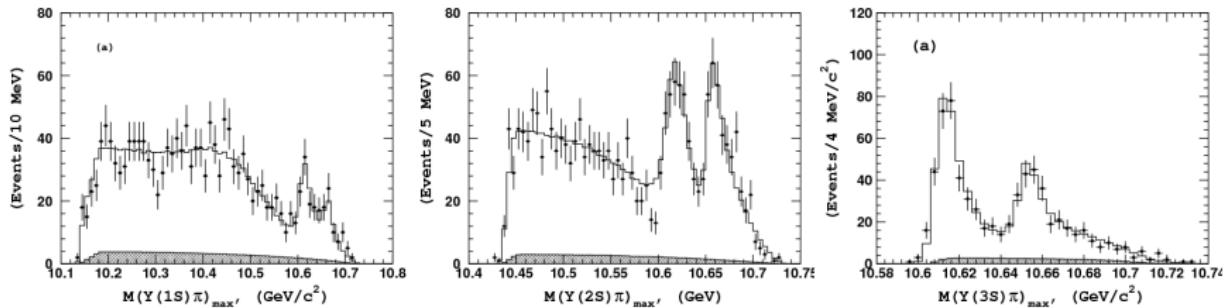


谢 谢 !

Backup slides

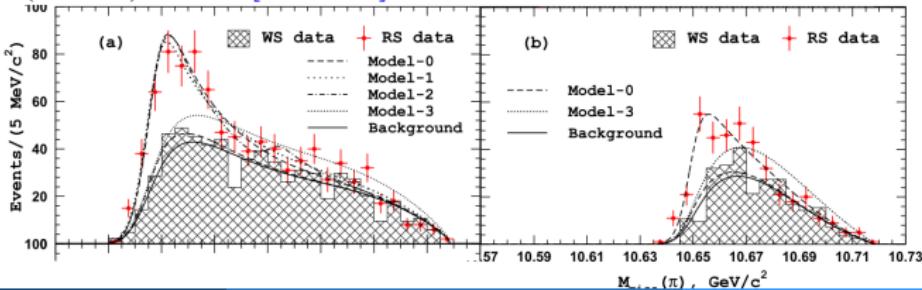
“Explicitly exotic” multiquarks: Z_c^\pm and Z_b^\pm with hidden $Q\bar{Q}$

- Z_c^\pm, Z_b^\pm : charged structures in heavy quarkonium mass region, $Q\bar{Q}\bar{d}u, Q\bar{Q}\bar{u}d$
 $Z_c(3900), Z_c(4020), Z_c(4200), Z_c(4430), \dots$ talks by R.Mitchell, C.-Z.Yuan, L.-M.Zhang
- $Z_b(10610)$ and $Z_b(10650)$: Belle, arXiv:1105.4583; PRL108(2012)122001
 observed in $\Upsilon(10860) \rightarrow \pi^\mp [\pi^\pm \Upsilon(1S, 2S, 3S)/h_b(1P, 2P)]$



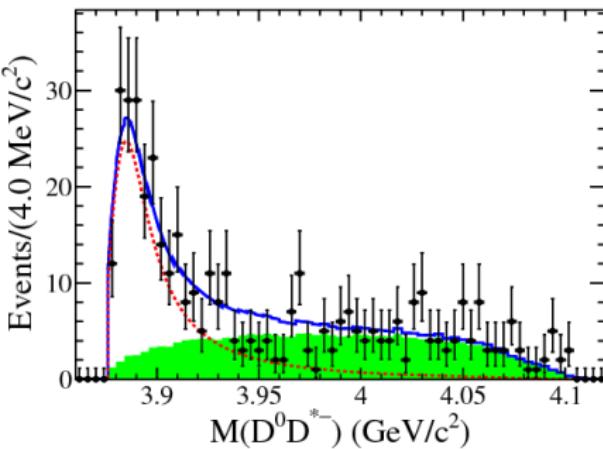
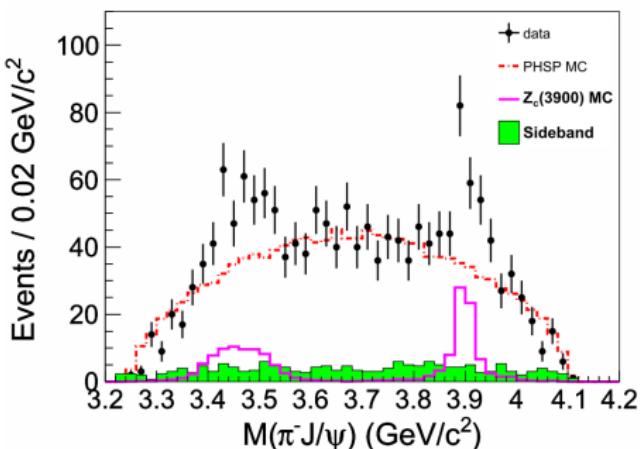
also in $\Upsilon(10860) \rightarrow \pi^\mp [B^{(*)}\bar{B}^*]^\pm$

Belle, arXiv:1209.6450; PRL116(2016)212001



Z_c^\pm and Z_b^\pm with hidden $Q\bar{Q}$ (II)

- $Z_c(3900/3885)^\pm$: structure around 3.9 GeV seen in $J/\psi\pi$ by BESIII and Belle in $Y(4260) \rightarrow J/\psi\pi^+\pi^-$, BESIII, PRL110(2013)252001; Belle, PRL110(2013)252002 and in $D\bar{D}^*$ by BESIII in $Y(4260) \rightarrow \pi^\pm(D\bar{D}^*)^\mp$ BESIII, PRD92(2015)092006



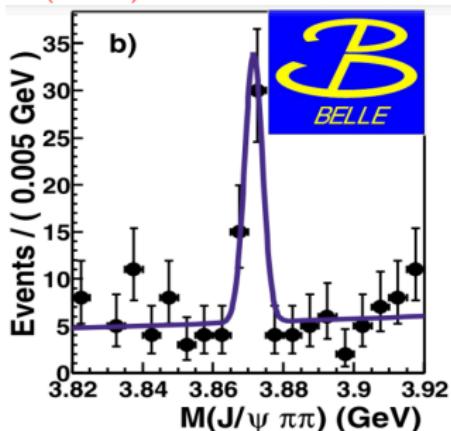
- can be described by the same state

Aldaladejo et al., PLB755(2016)337

$X(3872)$: best established

Belle, BaBar, BESIII, CDF, CMS, D0, LHCb

- $X(3872)$ Belle, PRL91(2003)262001



- Discovered in $B^\pm \rightarrow K^\pm J/\psi \pi\pi$, mass extremely close to the $D^0 \bar{D}^{*0}$ threshold
 $M_X = (3871.69 \pm 0.17) \text{ MeV}$

$$M_{D^0} + M_{D^{*0}} - M_X = (0.00 \pm 0.18) \text{ MeV}$$

- $\Gamma < 1.2 \text{ MeV}$ Belle, PRD84(2011)052004
- $J^{PC} = 1^{++}$ LHCb PRL110(2013)222001
 $\Rightarrow S$ -wave coupling to $D \bar{D}^*$

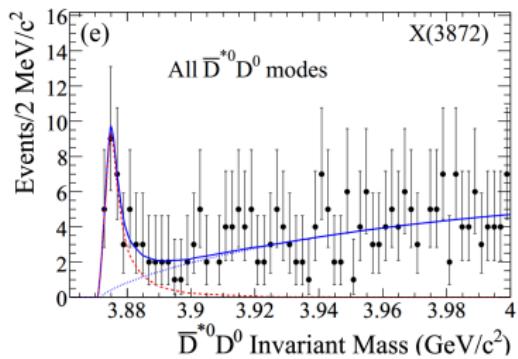
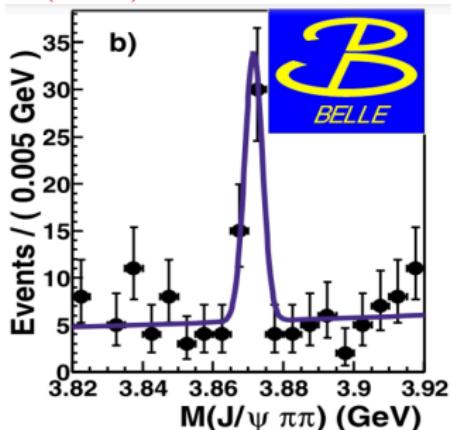
- Observed in the $D^0 \bar{D}^{*0}$ mode as well
BaBar, PRD77(2008)011102
- Large coupling to $D^0 \bar{D}^{*0}$:
 $\mathcal{B}(X \rightarrow D^0 \bar{D}^{*0}) > 24\%$ PDG2016
- Large isospin breaking:

$$\frac{\mathcal{B}(X \rightarrow \omega J/\psi)}{\mathcal{B}(X \rightarrow \pi^+ \pi^- J/\psi)} = 0.8 \pm 0.3$$

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Multiquark states: $X(5568)$?

Commun. Theor. Phys. **65** (2016) 593–595

Vol. 65, No. 5, May 1, 2016

How the $X(5568)$ Challenges Our Understanding of QCD*

Feng-Kun Guo (郭奉坤),^{1,†} Ulf-G. Meißner,^{1,2,3,‡} and Bing-Song Zou (邹冰松)^{1,4,§}

- mass too low for $X(5568)$ to be a $\bar{b}s\bar{u}d$: $M \simeq M_{B_s} + 200$ MeV
 - ☞ $M_\pi \simeq 140$ MeV because pions are pseudo-Goldstone bosons of spontaneous chiral symmetry breaking $SU(2)_L \times SU(2)_R \rightarrow SU(2)_V$
 - ☞ Gell-Mann–Oakes–Renner: $M_\pi^2 \propto m_q$; chiral counting: $M_\pi = \mathcal{O}(p)$
 - ☞ For any matter field: $M_R = \mathcal{O}(p^0) \gg M_\pi$; we expect $M_{\bar{q}q} \sim M_R \gtrsim M_\sigma$

$$M_{\bar{b}s\bar{u}d} \gtrsim M_{B_s} + 500 \text{ MeV} \sim 5.9 \text{ GeV}$$

- heavy quark flavor symmetry predicts an isovector X_c :

$$M_{X_c} = M_{X(5568)} - \bar{M}_{B_s} + \bar{M}_{D_s} + \mathcal{O}\left(\Lambda_{\text{QCD}}^2 \left(\frac{1}{m_c} - \frac{1}{m_b}\right)\right) \simeq (2.24 \pm 0.15) \text{ GeV}$$

but in $D_s\pi$, only isoscalar $D_{s0}^*(2317)$ was observed!

BaBar (2003)

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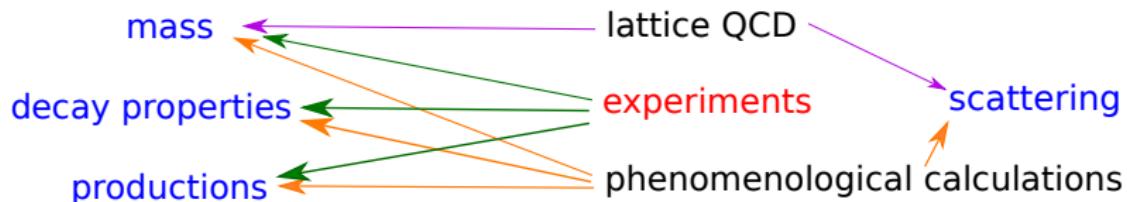
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Difficulties in interpreting experimental observations

If the observed structure are due to a genuine resonance \Rightarrow what is its nature?

Difficult to answer generally!



- Phenomenological calculations:
model dependence often hard to quantify
- Lattice calculations:
energy levels in finite volume, interpretation not straightforward

$X(3872)$ (III)

- Processes driven by short-distance $c\bar{c}$ physics:

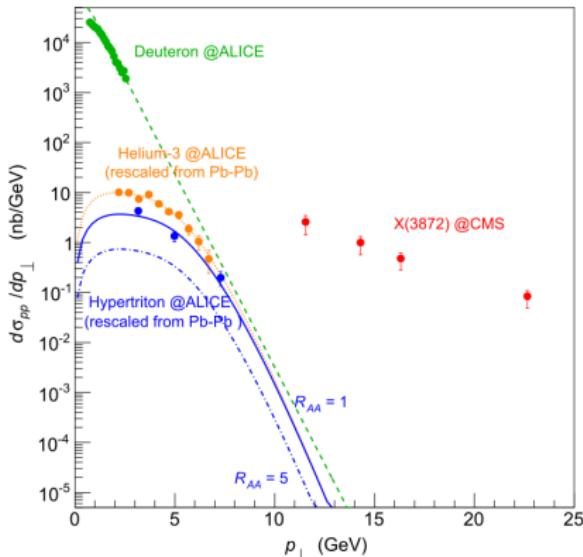
Examples:

- ☞ production of $X(3872)$ in B decays, at hadron colliders with large p_T

Braaten et al. (2004,2005,2006,2009); Meng, Gao, Chao (2005); Bignamini et al. (2009); ...

- Often used to blame the $D\bar{D}^*$ molecular interpretation, e.g.

Esposito et al., PRD92(2015)034028 :



- but deuteron and X are very different at short distances:

☞ deuteron: 6 quarks

☞ X : dominantly produced by $c\bar{c}$ at short distances