# Six-body calculation for electricdipole response of <sup>6</sup>Li

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Electric-dipole (E1) response reflects the nuclear structure

✓ The peak energy region✓ The shape of peak

- Giant-dipole-resonance (GDR)
  - large and broad peak
  - occurring at high energy region
  - out-of-phase oscillation of all protons and neutrons.

### Pygmy-dipole-resonance (PDR)

- Occurring at low energy region in neutron-proton unbalanced nuclei
- Mechanism is still controversial

The mode is collective, one-particle excitation or ...



 $\alpha$ (<sup>4</sup>He) cluster structure appears in light nuclei (e.g. <sup>6</sup>He, <sup>6</sup>Li, ...)



#### In the case of ${}^{6}$ He, there are two modes of GDR

D. Mikami, W. Horiuchi, and Y. Suzuki, Phys. Rev. C89, 064303 (2014)

Two valence neutrons is bound weakly

#### GDR

occurring at high energy region (~ 30MeV)

#### Soft-dipole mode (SDM)

- out-of-phase oscillation of α and two valence neutron
- occurring at low energy region (~ 3MeV)



α



Its conjecture are not confirmed by fully microscopic (six-body) calculation

The distortion of  $\alpha$  cluster is taken into account naturally

Purpose

- Perform fully microscopic calculation for ground and E1 excited states of <sup>6</sup>Li
- Understand the E1 excitation mechanism of <sup>6</sup>Li

### Problem

#### Difficulty to calculate many-body system

✓ describe the full variety of correlations between particles

Correlated Gaussian (CG) basis expansion
Correlations are described explicitly

 ✓ High calculation cost : the number of bases increases , the cost becomes higher



More important bases are selected

K.Varga and Y.Suzuki Phys. Rev .C52 (1995), 6

Method: Wave function



Wave function  $\rightarrow$  Basis expansion

$$\Psi = \sum_{i} c_i \varphi_i$$

Variational method + Basis expansion

Generalized eigenvalue problem

$$H\mathbf{c} = EB\mathbf{c}$$

 $B_{jk} = \langle \varphi_j | \varphi_k \rangle$  $H_{jk} = \langle \varphi_j | \hat{H} | \varphi_k \rangle$ 

# Method: Basis function

$$\begin{split} \varphi_i &= \mathcal{A}[\exp(-\frac{1}{2}\tilde{x}A_ix)\chi_S^{(i)}\eta] \text{ :basis function} \\ & \text{Correlated Gaussian} \\ \chi_S^{(i)} &= [[\chi^1\chi^2]_{S_{12}^{(i)}}\chi^3]_{S_{123}^{(i)}}...\chi^6]_{S=1} \\ \eta &= |\text{ppnnn}> \end{split}$$



$$oldsymbol{x} = (oldsymbol{x}_1, \cdots oldsymbol{x}_{A-1})$$

:Jacobi coordinates

 $A_i$ :positive-definite matrix $ilde{oldsymbol{x}}Aoldsymbol{x} = \sum_{j,k=1}^N A_{jk}oldsymbol{x}_j\cdotoldsymbol{x}_k$ 

Off diagonal elements describe correlations among particles

Variational parameters  $A_i$  ,  $\chi^{(i)}_S$ 

## Method: Selection of bases

Stochastic variational method (SVM)

K.Varga and Y.Suzuki Phys. Rev .C52 (1995), 6

(1) Candidates of variational parameter set  $\{A^1, A^2, \dots, A^K\}$  are generated randomly

②Energies  $\{E^1, E^2, \dots, E^K\}$  for each set are calculated

(3)Set  $A^i$  giving lowest energy is selected and added in bases

(4) (1~3) are repeated until the energy reaches convergence

## Results: ground state

 $V_{ij}^{NN} = (V_1 + \frac{1}{2}(1 + P_{ij}^{\sigma})V_2 + \frac{1}{2}(1 - P_{ij}^{\sigma})V_3)(\frac{1}{2}u + \frac{1}{2}(2 - u)P_{ij}^{r})$ :Minnesota potential  $V_l = V_{0l} \exp(-\mu_l |\mathbf{r}_i - \mathbf{r}_j|^2)$ D. R. Thompson *et al.* Nucl. Phys. A286 (1977) ,53

	$r_m$ [fm]	$r_p$ [fm]	$r_n$ [fm]	$E_{g.s.}$ [MeV]
<sup>6</sup> He	2.40	1.83	2.64	-30.95
<sup>6</sup> He(*)	2.41	1.83	2.65	-30.98

Experimental  $S_{2N}$  is reproduced by adjusting the u value

(\*)D. Mikami, W. Horiuchi, and Y. Suzuki, Phys. Rev. C89, 064303 (2014)

	$r_m$ [fm]	$r_p$ [fm]	$r_n$ [fm]	$E_{g.s.}$ [MeV]	$S_{2N}$ [MeV]
<sup>6</sup> Li	2.29	2.29	2.29	-33.90	4.0
<sup>6</sup> Li(Exp.)		2.45		-31.99	3.7

F. Ajzenberg-Selove, Nucl. Phys. A413, 1 (1984)

Different u values to reproduce the experimental data of rms radius can be considered

# E1 excited state



#### Model space

W. Horiuchi, Y. Suzuki, K. Arai, Phys. Rev. C 85 (2012) 054002

single-particle (sp) excitation  $\phi_i^{(sp)} = \mathcal{A}[\phi_i^{(^6\mathrm{Li})}\mathcal{Y}_{1\mu}(\boldsymbol{r}_1 - \boldsymbol{x}_{cm})]_J$ 

 $\alpha + \mathbf{p} + \mathbf{n} \text{ decay channel}$  $\phi_i^{(\alpha p n)} = \mathcal{A}[\phi_i^{(\alpha)} \exp(-\frac{1}{2}\tilde{z}B_i z)[[\chi^{(5)}\chi^{(6)}]_{s_1}]\mathcal{Y}_{1\mu}(z_1)]_{J_1}]_J$ 

 $h(^{3}\text{He}) + t(^{3}\text{H})$  decay channel (Not included in this work)

$$\phi_i^{(ht)} = \mathcal{A}[[\phi_i^{(h)}\phi_i^{(t)}]_{J_1} \exp(-\frac{1}{2}a\boldsymbol{y}^2)\mathcal{Y}_{1\mu}(\boldsymbol{y})]_{J_2}]_J$$

 $\alpha + p + n$  $oldsymbol{z}_1$ α Y-type T-type

sp



### Results: E1 strength



# **Results:** Transition density



Like soft-dipole mode

### Results: Transition density





# Summary

- ✓ Six-body calculation for ground and E1 excited state of <sup>6</sup>Li is performed
- ✓ The u parameter of Minnesota potential is adjusted to reproduce the experimental data of two nucleon separation energy for the ground state
- ✓ The E1 strength and transition density show the possibility of three modes of GDR, SDM, GDR of <sup>6</sup>Li and GDR of <sup>4</sup>He in <sup>6</sup>Li

#### Future work

- Add h + t decay channel to the model space
- Detail analyses of the modes of GDR