STRUCTU A LA HUEY-WEN LIN Founded 855

Nucleon Structure

- § Study nucleon structure since '60s
 e Deep inelastic scattering @ SLAC, more
- § Fundamental QCD property



Exploration of the valence and sea-quark content of the nucleon
 Important for BSM searches

Provides SM cross-section prediction for LHC new-physics search
 § Impressive results but still limited knowledge

Many ongoing/planned experiments: (RHIC, JLab, J-PARC, COMPASS-II, GSI, EIC, LHeC, ...)





How Can LQCD Help?

 § Lattice QCD is an ideal theoretical tool for investigating strong-coupling regime of quantum field theories
 § Physical observables are calculated from the path integral

$$\langle 0 | O(\bar{\psi}, \psi, A) | 0 \rangle = \frac{1}{Z} \int \mathcal{D}A \, \mathcal{D}\bar{\psi} \, \mathcal{D}\psi \, e^{iS(\bar{\psi}, \psi, A)} O(\bar{\psi}, \psi, A)$$

in **Euclidian** space

Quark mass parameter (described by m_{π})
Impose a UV cutoff discretize spacetime
Impose an infrared cutoff finite volume
S Recover physical limit $m_{\pi} \rightarrow m_{\pi}^{\text{phys}}, a \rightarrow 0, L \rightarrow \infty$ x, y, z x, y, z $a \rightarrow 0, L \rightarrow \infty$

How Can LQCD Help?

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in **Euclidi**
Sequark minimum Great to study the quark and gluon
(described degrees of freedom inside hadrons!
Set Impose a ov cutoring
discretize spacetime
Set Impose an infrared cutoff
finite volume
Set Recover physical limit
 $m_{\pi} \rightarrow m_{\pi}^{\text{phys}}, a \rightarrow 0, L \rightarrow \infty$



Are We There Yet?

- § Lattice gauge theory was proposed in the 1970s by Wilson
- > Why haven't we solved QCD yet?
- § Progress is limited by computational resources 1980s Today





§ Greatly assisted by advances in algorithms
> Physical pion-mass ensembles are not uncommon!



Successful Examples

§ Lattice flavor physics provides precise inputs from the SM
 A. El-Khadra, Sep. 2015, INT workshop "QCD for New Physics at the Precision Frontier"
 > Very precise results in many meson systems



errors (in %) (preliminary) FLAG-3 averages

§ We are beginning to do precision calculations in nucleons



Outlíne

§ Lattice Precision Frontier

➢ Nucleon structure with controlled systematics in physical limit (m_π → m_π^{phys}, a → 0, L → ∞)
 ➢ New-physics implications

§ Lattice Pioneer Frontier

> Lattice parton distribution functions





Precision Nucleon Structure

§ Much effort has been devoted to controlling systematics § A state-of-the art calculation (PNDME)

<i>a</i> (fm)	V	$M_{\pi}L$	$oldsymbol{M}_{\pi}$ (MeV)	t _{sep}	# Meas.
0.12	$24^3 \times 64$	4.55	310	8,10,12	64.8k
0.12	$24^3 \times 64$	3.29	220	8,10,12	24k
0.12	$32^3 \times 64$	4.38	220	8,10,12	7.6k
0.12	$40^3 \times 64$	5.49	220	8,10,12,14	64.6k
0.09	$32^3 \times 96$	4.51	310	10,12,14	7.0k
0.09	$48^3 \times 96$	4.79	220	10,12,14	7.1k
0.09	$64^3 \times 96$	3.90	135	10,12,14	56.5k
0.06	$48^3 \times 144$	4.52	310	16,20,22,24	64.0k
0.06	64 ³ × 144	4.41	220	16,20,22,24	41.6k
0.06	96 ³ × 192	3.80	135	16,18,20,22	51.5 k

Tensor/Scalar Charge

$g_T(a, m_\pi, L) =$	$c_1 + c_2 m_\pi^2 + c_3 a + c_4 e^2$	$-m_{\pi}L$
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Tensor/Scalar Charge

§ PNDME's g_{T,S} calculations PNDME, 1506.06411; 1506.04196; in preparation
 ➢ Extrapolate to the physical limit

 $g_T(a, m_\pi, L) = c_1 + c_2 m_\pi^2 + c_3 a + c_4 e^{-m_\pi L}$





Precision Nucleon Couplings

FLAG rating system

PNDME, 1506.06411; 1606.07049



MICHIGAN STATE

Beta Decays & BSM

§ Given precision $g_{S,T}$ and O_{BSM} , predict new-physics scales Low-Energy

Expt
$$\longrightarrow O_{\text{BSM}} = fo(\varepsilon_{s,\tau} g_{s,\tau}) \leftarrow (m_{\pi} \rightarrow 140 \text{ MeV}, a \rightarrow 0)$$



 $\varepsilon_{S,T} \propto \Lambda_{S,T}^{-2}$ Upcoming precision low-energy experiments LANL/ ORNL UCN neutron decay exp't $|B_1 - b|_{\rm BSM} < 10^{-3}$ $|b|_{\rm RSM} < 10^{-3}$ CENPA: ${}^{6}\text{He}(b_{GT})$ at 10^{-3} PNDME, PRD85 054512 (2012); 1306.5435; 1606.07049 $\Lambda_S > 7 \text{ TeV}$ $\Lambda_T > 13 \text{ TeV}$



Electric Dipole Moment

§ Why do we care?

 \sim CP-violating effect \Rightarrow Key ingredient for baryogenesis

 \Rightarrow Why matter exists

➢ Extremely small in SM: ≈ 10⁻³¹ e-cm (expect to probe 10⁻²⁸ soon)
 ➢ Good candidate to constrain BSM models





Electric Dipole Moment

§ Quark EDM (d_q) in nucleon comes from

 $d_N = d_u g_T^{(n,u)} + d_d g_T^{(n,d)} + d_s g_T^{(n,s)}$ \Rightarrow Hadronic contribution: $\langle N | \bar{q} \sigma_{\mu\nu} q | N \rangle$, $q \in \{u, d, s\}$ \Rightarrow Need "disconnected" diagram contributions § Extrapolate to the continuum limit PNDME, 1506.04196; 1506.06411

 $g_T^u = 0.774(66), g_T^d = -0.233(28), g_T^s = 0.008(9)$

 § Implications for new physics?
 ➢ Take split SUSY for example
 ➢ Using our lattice inputs, we can derive an upper limit for the neutron EDM in split SUSY
 Wells, 2003; Arkani-Hamed and Dimopoulos, 2004; Giudice and Romanino, 2004

 $|d_n| < 4 \times 10^{-28} e \cdot \mathrm{cm}$

using $|d_e| < 8.7 \times 10^{-29} e \cdot \text{cm}$ with 90% confidence

ACME Coll., Science Vol. 343 no. 6168 pp. 269-272 (2014)



EM Form Factors





Axíal Form Factors

§ Axial form factor

PNDME, 1705.06834

$$\approx \langle N(\vec{p}_f) | A_{\mu}(\vec{Q}) | N(\vec{p}_i) \rangle = \bar{u}(\vec{p}_f) \left[G_A(Q^2) \gamma_{\mu} + q_{\mu} \frac{\tilde{G}_P(Q^2)}{2M_N} \right] \gamma_5 u(\vec{p}_i)$$

Important input for neutrino experiments, such as DUNE



Axíal Form Factors

§ Axial form factor PNDME, 1705.06834 $\approx \langle N(\vec{p}_f) | A_{\mu}(\vec{Q}) | N(\vec{p}_i) \rangle = \bar{u}(\vec{p}_f) \left[G_A(Q^2) \gamma_{\mu} + q_{\mu} \frac{\tilde{G}_P(Q^2)}{2M_N} \right] \gamma_5 u(\vec{p}_i)$ > Important input for neutrino experiments, such as DUNE $\nu, \overline{\nu}$ Scattering Electron 0.8 Deuteriun MiniBoo G_A/g_A 0.6 a06*m*310 HAH a06m220 ⊷ a06m135 0.4 0.2 0.8 0.2 0.4 0.6 12 1.81.4 1.6

Huey-Wen Lin — 9th Hadron Physics in China & Worldwide @ Nanjing

Lattice Pioneer Frontier

Xiangdong Ji, Phys. Rev. Lett. 111, 039103 (2013)

"Flavor Structure of the Nucleon Sea from Lattice QCD",

PRD 91, 054510 [arXiv:1402.1462]

"Nucleon Helicity and Transversity Parton Distributions from Lattice QCD", to be appeared in Frontier Article in Nuclear Physics B,

[arXiv:1603.06664]



S

PDFs on the Lattice

Long existing obstacles!

§ Lattice calculations rely on operator product expansion, only pro $dx x^{n-1}q(x)$

§ For higher No practi **New Strates** § Calculate quark dist $rightarrow In P_{7} \rightarrow \infty$ ✤ For finite § Feasible with today s resource

Symmetry: You Break it, You Buy It.

> Xiangdong Ji, Phys. Rev. Lett. 111, 039103 (2013)



Huey-Wen Lin — 9th Hadron Physics in China & Worldwide @ Nanjing

imension ops

XI

PDFs on the Lattice

Long existing obstacles!

§ Lattice calculations rely on operator product expansion, only provide moments $\langle x^n \rangle$ $\langle x^n \rangle_q = \int_{-1}^{1} dx \, x^n q(x)$

§ For higher moments, all ops mix with lower-dimension ops \gg No practical proposal to overcome this

- New Strategy (LaMET):
- § Calculate finite-momentum boosted quark distribution

Solution In P_z → ∞ limit, parton distribution is recovered
Solution For finite P_z , corrections are needed
Solution Feasible with today's resources!

Xiangdong Ji, Phys. Rev. Lett. 111, 039103 (2013)



 x_{\perp}

PDFs on the Lattice

- Long existing obstacles!
- § Lattice calculations rely on operator product expansion, only provide moments $\langle x^n \rangle$ $\langle x^n \rangle_q = \int_{-1}^{1} dx \, x^n q(x)$
- § For higher moments, all ops mix with lower-dimension ops
- \sim No practical proposal to overcome this
- New Strategy (LaMET):
- § Calculate finite-momentum boosted quark distribution
- § Feasible with today's resources!

 $\frac{x_{-}}{P_{z}=0}$

Xiangdong Ji, Phys. Rev. Lett. 111, 039103 (2013)



A New Direction

Large-Momentum Effective Theory for PDFs 1) Calculate nucleon matrix elements on the lattice



2) Compute quasi-distribution via

$$\tilde{q}(x,\mu,P_z) = \int \frac{dz}{4\pi} e^{-izk_z} \left\langle P \left| \overline{\psi}(z) \right| \sum \exp\left(-ig \int_0^z dz' A_z(z')\right) \psi(0) \right| P \right\rangle$$

3) Recover true distribution (take $P_z \rightarrow \infty$ limit)

 $\boldsymbol{q(x,\mu)} = \tilde{\boldsymbol{q}(x,\mu,P_z)} + \mathcal{O}(\alpha_s) + \mathcal{O}(M_N^2/P_z^2) + \mathcal{O}(\Lambda_{\rm QCD}^2/P_z^2)$

X. Xiong et al., 1310.7471; J.-W. Chen et al, 1603.06664

Sea Flavor Asymmetry

§ First time in LQCD history to study antiquark distribution! $\gg M_{\pi} \approx 310 \text{ MeV}$



$$\bar{q}(x) = -q(-x)$$

Lost resolution in small-x region Future improvement: larger lattice volume

$$dx\left(\bar{u}(x) - \bar{d}(x)\right) \approx -0.16(7)$$

Experiment	x range	$\int_0^1 [\overline{d(x)} - \overline{u(x)}] dx$
E866	0.015< <i>x</i> <0.35	0.118 ± 0.012
NMC	0.004 < x < 0.80	0.148 ± 0.039
HERMES	0.020 < x < 0.30	0.16 ± 0.03

R. Towell et al. (E866/NuSea), Phys.Rev. D64, 052002 (2001)

Sea Flavor Asymmetry

§ Lattice exploratory study $\approx M_{\pi} \approx 310 \text{ MeV}$



Compared with E866 Too good to be true?

Lost resolution in small-x region

Similar results repeated by ETMC, at $M_{\pi} \approx 373$ MeV ETMC, 1504.07455

Experiment	x range	$\int_0^1 [\overline{d(x)} - \overline{u(x)}] dx$
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R. Towell et al. (E866/NuSea), Phys.Rev. D64, 052002 (2001)

(7)

Helicity Distribution

§ Exploratory study $\gg M_{\pi} \approx 310 \text{ MeV}$





Removing $O(M_N^n/P_z^n)$ errors + $O(\alpha_s)$ + $O(\Lambda_{QCD}^2/P_z^2)$

Solution We see polarized "sea asymmetry" $\int dx \left(\Delta \bar{u}(x) - \Delta \bar{d}(x)\right) \approx 0.14(9)$ So both STAR and PHENIX at RHIC see $\Delta \bar{u} > \Delta \bar{d}$

1404.6880 and 1504.07451

> Other experiments, Fermilab DY exp'ts (E1027/E1039), future EIC

Transversity Distribution



Physical Pion Mass

§ Preliminary results

$M_{\pi} \approx$ **135 MeV**, $a \approx$ **0.09 fm**



Yi-Bo Yang (MSU)

A NEW HOPE

It is a period of war and economic uncertainty.

Turmoil has engulfed the galactic republics.

Basic truths at foundation of the human civilization are disputed by the dark forces of the evil empire.

A small group of QCD Knights from United Federation of Physicists has gathered in a remote location on the third planet of a star called Sol on the inner edge of the Orion-Cygnus arm of the galaxy.

The QCD Knights are the only ones who can tame the power of the Strong Force, responsible for holding atomic nuclei together, for giving mass and shape to matter in the Universe.

They carry secret plans to build the most powerful

Summary & Outlook

Exciting era for studying structure on the lattice

- § We are entering the precision era for charges
- ➢ More nucleon matrix elements with physical pion masses
 ➢ Well-studied systematics → precision inputs
- § Overcoming longstanding obstacle to full x-distribution
 > Most importantly, this can be done with today's computer
 > First lattice approach to study sea asymmetry
- § Systematics, systematics, systematics...

Working on renormalization, statistics, physical pion mass, finer lattice-spacing ensembles, ...



HPCC@MSU

Backup Slides





Approaching the Continuum

§ Worse excited-state contamination

PNDME, 1606.07049



Nucleon Axíal Charge

FLAG-like rating system New: excited-state rating

White paper in progress with representatives from each collaboration



➢ Crucial to have independent/unbiased g_A from LQCD, $\lambda_{exp} = g_A / g_V f_{New Physics}$

Form Factors

§ EM form factors very close to physical pion mass *▶* PNDME (2+1+1f): M_{π} = 130 MeV, a = 0.09fm, $M_{\pi}L$ = 3.9, 2-state, 56K ➢ ETMC (2f): M_{π} = 131 MeV, a = 0.09fm, $M_{\pi}L$ = 3, t_{sep} = 1.3fm, O(10K)? *▶* LHPC (2+1f): M_{π} = 149 MeV, a = 0.12 fm, $M_{\pi}L$ = 4.2, sum., 0(7K) ➢ PACS (2+1f): M_{π} = 145 MeV, *a* = 0.085 fm, $M_{\pi}L$ = 6, t_{sep} = 1.3fm, O(9K) 1.00.8 0.6 0.40.2 Q^2 (GeV²) 0.3 0.2 0.10.20.40.5 0.10.3 0.5§ Expecting more precise results in the next couple years § New way of calculating form factors C. Chang, Tue. Lattice, "Form factors from moments of correlation functions"

Disconnected Diagrams

§ Disconnected diagram

> Multiple ways to calculate this notorious contribution...

Truncated solver, hopping-parameter expansion, hierarchical probing, ...





Strangeness

§ Importance of g_A^s

> Strange-quark intrinsic-spin contribution to proton

✤ Astrophysics application: the CCSN "problem"

3D explosions require $g_A^s \approx -0.2$

§ Global fit: $g_A^s \approx -0.1$ assumptions often used:

$$\Delta \bar{s}(x,Q^2) = \Delta \bar{u}(x,Q^2)$$
$$= \Delta \bar{d}(x,Q^2)$$
$$= \frac{1}{2}\Delta s^+(x,Q^2)$$

§ Lattice status

More players since the last Spin Lighter pion masses Janka, Melson, & Summa (2016)



Strange Form Factors

§ Better determined strange form factors \approx LHPC (2+1f): clover M_{π} = 317 MeV, a = 0.11 fm $\approx \chi$ QCD (2+1f): ov/DWF M_{π} = 207,140 MeV, a = 0.11 fm



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K. Orginos/R. Sufian,

Tue. Lattice

Parton Distribution Functions

§ PDFs are universal quark/gluon distributions inside nucleon Many ongoing/planned experiments (BNL, JLab, J-PARC, COMPASS, GSI, EIC, LHeC, ...)







Electron Ion Collider: The Next QCD Frontier

Imaging of the proton

How are the sea quarks and gluons, and their spins, distributed in space and momentum inside the nucleon? EIC White Paper, 1212.1701





Parton Distribution Functions

- § PDFs are universal quark/gluon distributions inside nucleon
- Many ongoing/planned experiments (BNL, JLab, J-PARC, COMPASS, GSI, EIC, LHeC, ...)
- § Important inputs to discern new physics at LHC Currently dominate errors in Higgs production



(J. Campbell, HCP2012)



Global Analysis

§ Experiments cover diverse kinematics of parton variables

✤ Global analysis takes advantage of all data sets



Choice of data sets and kinematic cuts

Strong coupling constant $\alpha_s(M_Z)$

How to parametrize the distribution

$$xf(x,\mu_0) = a_0 x^{a_1} (1-x)^{a_2} P(x)$$

Assumptions imposed

SU(3) flavor symmetry, charge symmetry, strange and sea distributions

$$s = \bar{s} = \kappa \big(\bar{u} + \bar{d} \big)$$



Global Analysis

§ Discrepancies appear when data is scarce
 § Many groups have tackled the analysis
 > CTEQ, MSTW, ABM, JR, NNPDF, etc. 10⁻⁴

CJ12mid **CT10** 0.8 - MSTW08 ABKM09 0.6 d/u0.4 0.2 = 100 GeV0 0.2 0.4 0.6 0.8 0 x

Jimenez-Delgado, Melnitchouk, Owens, J.Phys. G40 (2013) 09310



New Direction

Large-Momentum Effective Theory (LaMET) X. Ji, PRL. 111, 262002 (2013) Finite-P_z corrections needed Neglect typical lattice corrections for now:



§ Benefit from our pQCD colleagues

New Direction

Large-Momentum Effective Theory (LaMET) X. Ji, PRL. 111, 262002 (2013) Finite-P_z corrections needed ≫ Neglect typical lattice corrections for now:

$$\tilde{q}(x,\mu,P_z) = \int_{-\infty}^{\infty} \frac{dy}{|y|} Z\left(\frac{x}{y},\frac{\mu}{P_z}\right) q(y,\mu) + \mathcal{O}(M_N^2/P_z^2) + \mathcal{O}\left(\Lambda_{\text{QCD}}^2/P_z^2\right)$$

complicated higher-twist operator; smaller P_z correction for nucleon J.-W. Chen et al, 1603.06664 and reference within (extrapolate it away)

§ Some similarity to more broadly studied HQET...

$$O\left(\frac{m_b}{\Lambda}\right) = Z\left(\frac{m_b}{\Lambda}, \frac{\Lambda}{\mu}\right) o(\mu) + O\left(\frac{1}{m_b}\right) + \cdots$$



Systematic Control

§ Much effort has been devoted to controlling systematics
§ A state-of-the art calculation (PNDME)
➢ Statistical effect a = 0.06 fm, 220-MeV pion



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§ A state-of-the art calculation (PNDME)
➢ Statistical effect a = 0.06 fm, 220-MeV pion





Systematic Control

§ Much effort has been devoted to controlling systematics
§ A state-of-the art calculation (PNDME)
➢ Statistical effect (worst case) a = 0.06 fm, 220-MeV pion

