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Search for heavy Higgs bosons A/H decaying to a top quark pair in pp collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector

ATLAS Collaboration

A search for heavy pseudoscalar (A) and scalar (H) Higgs bosons decaying into a top quark pair ($t\bar{t}$) has been performed with 20.3 fb^{-1} of proton–proton collision data collected by the ATLAS experiment at the Large Hadron Collider at a center-of-mass energy of $\sqrt{s} = 8$ TeV. Interference effects between the signal process and Standard Model $t\bar{t}$ production, which are expected to distort the signal shape from a single peak to a peak–dip structure, are taken into account. No significant deviation from the Standard Model prediction is observed in the $t\bar{t}$ invariant mass spectrum in final states with an electron or muon, large missing transverse momentum, and at least four jets. The results are interpreted within the context of a type-II two-Higgs-doublet model. Exclusion limits on the signal strength are derived as a function of the mass $m_{A/H}$ and the ratio of the vacuum expectation values of the two Higgs fields, $\tan\beta$, for $m_{A/H} > 500$ GeV.

The production of new particles at the Large Hadron Collider (LHC) with masses close to the TeV scale is predicted by many models of physics beyond the Standard Model (SM). In this Letter, a search for massive pseudoscalar and scalar resonances decaying into a top–antitop quark pair ($t\bar{t}$) is presented. It is the first search in this final state to take into account the significant interference between the signal and the background from SM $t\bar{t}$ production. The search is conducted on a sample of pp collision data corresponding to an integrated luminosity of 20.3 fb^{-1} at a center-of-mass energy of $\sqrt{s} = 8$ TeV, collected with the ATLAS detector [1].

New pseudoscalar (A) and scalar (H) states coupling strongly to $t\bar{t}$ are predicted by a class of models in which the Higgs sector is extended to include a second Higgs doublet, referred to as two-Higgs-doublet models (2HDMs) [2]. These models are motivated by many theories beyond the SM, such as supersymmetry [3–8] and axion models [9]. In 2HDMs of type II [2], such as the Minimal Supersymmetric Standard Model (MSSM), these states decay predominantly into $t\bar{t}$ pairs if $m_{A/H} \geq 500$ GeV and the ratio of the vacuum expectation values of the two Higgs fields, $\tan\beta$, is small ($\tan\beta \lesssim 3$).

To date, this parameter region has not been probed by searches in other final states [10–15] or by previous searches for $t\bar{t}$ resonances [16–20]. The latter, which aim to identify resonant excesses in the $t\bar{t}$ invariant mass ($m_{t\bar{t}}$) spectrum, have a reduced sensitivity to 2HDM signatures as they do not take into account interference effects between the signal and the dominant background from SM $t\bar{t}$ production. These are significant for (pseudo)scalar Higgs bosons with masses above the $t\bar{t}$ production threshold where the interference between the gluon–gluon (gg) initiated loop production and the irreducible background from SM $t\bar{t}$ production yields a non-negligible imaginary term in the amplitude, which at the LHC is dominated by $gg \to t\bar{t}$ production [21–24]. As a result of the interference, the signal shape is distorted from a simple Breit–Wigner peak to a peak–dip structure.

The results of the search are interpreted in a CP-conserving 2HDM of type II with a softly broken Z_2 symmetry [25]. The lighter of the two neutral CP-even states, h, is assumed to be the Higgs boson discovered at a mass of $m_h = 125$ GeV [26, 27] with couplings as predicted by the SM. This corresponds to the condition $\sin(\alpha - \beta) = 1$, referred to as the *alignment limit*, where α denotes the mixing angle between the two CP-even states. The parameter m_{12} of the Z_2 breaking term of the potential is taken to be $m_{12}^2 = m_A^2 \tan\beta/(1 + \tan^2\beta)$. In this model, the production cross-sections and widths of A and H, as well as the signal shape, are determined by $\tan\beta$ and the masses m_A and m_H. Three mass hierarchies are considered. In the first two scenarios, one of the Higgs bosons is significantly heavier than the other so that only the interference pattern of the lighter boson appears in the $m_{t\bar{t}}$ spectrum. In the third scenario, motivated by a number of SM extensions [2], A and H are assumed to be mass degenerate ($m_H = m_A$) such that both processes contribute to the $m_{t\bar{t}}$ spectrum, leading to more significant enhancements of the interference pattern.

This analysis closely follows the resolved-topology analysis in Ref. [17]. Events with signatures compatible with $t\bar{t} \to W^+bW^-\bar{b}$, with one W boson decaying hadronically and the other leptonically, referred to as the lepton-plus-jets channel (ℓ+jets, $\ell = e, \mu$) were collected using single-electron and single-muon triggers. The trigger efficiency is constant in the transverse momentum (p_T) of leptons with $p_{\mathrm{T}} > 25$ GeV [28, 29]. The dominant background arises from SM $t\bar{t}$ production, followed by a contribution from W+jets processes. Data-driven techniques were used to normalize the W+jets background contribution and to estimate the background from multijet events. All other background processes were estimated using Monte Carlo (MC) simulation. The background estimates for all processes are identical to those used in Ref. [17].

The signal process $gg \to A/H \to t\bar{t}$, including the decays of the top quarks and resulting W bosons, was simulated using MadGraph5_aMC@NLO [30] v2.3.3 with the model of Ref. [31], which implements the production of (pseudo)scalar particles through loop-induced gluon–gluon fusion with loop contributions from top and bottom quarks at leading order (LO) in QCD. The CT10 set [32] of parton distribution functions (PDFs) was used and the renormalization and factorization scales were set to $\sqrt{\sum_{\text{decay products}}(p_{\text{T}}^2 + m^2)}$.

For the statistical interpretation, the $t\bar{t}$ invariant mass distributions in the signal regions in data were compared to a combination of the expected distributions from all background processes B, the pure signal process S, and the signal-plus-interference component $S + I$ for a given signal hypothesis, as illustrated in Eq. (1) below. The most reliable description of the $t\bar{t}$ background [33] is obtained at next-to-leading order (NLO) with Powheg-Box [34–37] + Pythia6 [38]. Therefore, the $S + I$ contribution was modeled separately from this background process by modifying the MadGraph5_aMC@NLO software to remove the pure SM $t\bar{t}$ process to yield only the $S + I$ contribution on an event-by-event basis. The $S + I$ events obtained with the modified software can have positive or negative weights. Figure 1 shows the $t\bar{t}$ invariant mass distributions for the S and $S + I$ components in a model with $\tan\beta = 0.68$ and a pseudoscalar of mass $m_A = 500$ GeV ($m_A \ll m_H$). The $S + I$ component exhibits a peak-dip structure with the minimum around $m_{A/H}$ for all signal hypotheses studied in this search. The width of both the S and $S + I$ distribution decreases with increasing $\tan\beta$.

The $S + I$ distributions from the modified MadGraph5_aMC@NLO software were validated against those from the unmodified program. The latter were obtained by generating a large inclusive sample $S + I + B_{t\bar{t}}$ for a given parameter point and a LO SM $t\bar{t}$ background $B_{t\bar{t}}$ sample with the same generator settings. The event selection and reconstruction algorithms were applied separately to each sample. The difference between the resulting two $m_{t\bar{t}}$ distributions corresponds to the $S + I$ component, which agrees with that obtained with the modified software within 0.4% across the whole spectrum. The difference is taken as a systematic uncertainty in $S + I$.

Pythia6 with the Perugia 2011c set of tuned parameters [39] was used to model the parton shower and hadronization for all S and $S + I$ samples and the stable particles obtained after hadronization were passed through the ATLAS fast detector simulation [40]. The effects of additional collisions within the same or nearby bunch crossings were simulated by overlaying additional pp collisions, simulated with Pythia v8.1 [41], on each event. Correction factors were applied to adjust the trigger and selection efficiencies in simulated events to those measured in data. The S and $S + I$ samples with this setup were generated separately for pseudoscalar and scalar Higgs bosons.

Event samples for both the S and $S + I$ components for different values of $(m_{A/H}, \tan\beta)$ were obtained from signal samples S after the detector simulation by applying an event-by-event reweighting. This reweighting substantially reduces the computing time required. The weight is the ratio of the MadGraph5_aMC@NLO matrix elements, calculated from the four-momenta of the incoming gluons and outgoing top quarks of the generated event with the new and the old values of $(m_{A/H}, \tan\beta)$, respectively. All $S + I$ and a small number of S samples were obtained through reweighting. Signal hypotheses with $m_{A/H} < 500$ GeV were not considered as they require an accurate modeling of the Higgs boson decay into virtual top quarks and the implementation of higher-order corrections that are not available in the MadGraph5_aMC@NLO model. The requirement $\tan\beta \geq 0.4$ was imposed to ensure that all amplitudes involving scalars preserve perturbative unitarity [25, 42].

Correction factors K_S were applied to normalize the generated signal (S) cross-section to the value calculated at partial next-to-next-to-leading-order (NNLO) precision in QCD [43–45]. The correction factor

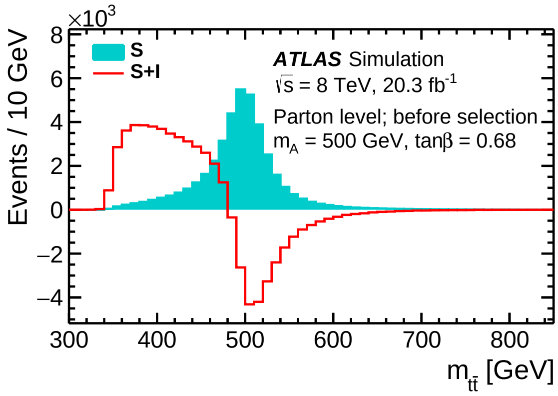

Figure 1: Distributions of the invariant mass of the $t\bar{t}$ pair from the decay of a pseudoscalar A of mass m_A = 500 GeV ($m_A \ll m_H$) at parton level before the emission of final-state radiation and before the parton shower for the pure resonance S (filled) and signal+interference contribution $S + I$ (unfilled). Events from all $t\bar{t}$ decay modes are included and no selection requirements are imposed. The distributions are normalized to an integrated luminosity of 20.3 fb^{-1}.

for the interference component I is $K_I = \sqrt{K_S \times K_B}$, as suggested in Ref. [46], where $K_B = 1.87$ is the correction factor to normalize the total cross-section of the SM $t\bar{t}$ background generated at LO with MadGraph to the cross-section calculated at NNLO accuracy in the strong coupling constant α_S, including resummation of next-to-next-to-leading-logarithmic soft gluon terms. The values of K_S range between two and three for the tested signal hypotheses.

The event selection criteria for the signal regions provide a high selection efficiency for $t\bar{t}$ events in the ℓ+jets channel. Only events with a *resolved topology*, in which the three jets from the hadronically decaying top quark are well separated in the detector, are selected. This is the most efficient selection strategy for signal hypotheses with $m_{A/H} < 800$ GeV. Events with a *merged topology*, in which the hadronically decaying top quark is reconstructed as a single large radius jet, are not considered. The event reconstruction and selection criteria are identical to those for the resolved topology in Ref. [17] except that events that would satisfy the criteria for both topologies are classified as "resolved" instead of "merged".

Events are required to contain exactly one isolated electron or muon that is geometrically matched to the corresponding trigger-level signature. Electron candidates are reconstructed from energy clusters in the electromagnetic calorimeter matched to tracks in the inner detector and required to have $E_{\mathrm{T}} > 25$ GeV and pseudorapidity [47] $|\eta_{\mathrm{cluster}}| < 2.5$ [48]. Muon candidates are reconstructed by combining tracks in the inner detector with tracks in the muon spectrometer [49]. They are required to have $p_{\mathrm{T}} > 25$ GeV and $|\eta| < 2.5$. Events must have large missing transverse momentum, $E_{\mathrm{T}}^{\mathrm{miss}} > 20$ GeV, computed as the magnitude of the negative vector sum of lepton and jet transverse momenta [50]. In addition, $E_{\mathrm{T}}^{\mathrm{miss}} + m_{\mathrm{T}}^{W} >$ 60 GeV, is required to further suppress the contribution from multijet events, where m_{T}^{W} is the lepton–$E_{\mathrm{T}}^{\mathrm{miss}}$ transverse mass [17]. Events must contain at least four hadronic jets with $p_{\mathrm{T}} > 25$ GeV and $|\eta| < 2.5$, reconstructed using the anti-k_t algorithm [51, 52] with radius parameter $R = 0.4$. Jets from additional collisions in the same bunch crossing are rejected using dedicated tracking and vertex requirements [53] applied to jets with $p_{\mathrm{T}} < 50$ GeV and $|\eta| < 2.4$. At least one of the jets must be identified as originating from the decay of a b-hadron (b-jet) using a multivariate tagging algorithm with a 70% efficiency for b-jets [54].

Jets are assigned to the top quarks using a χ^2 algorithm that relies on kinematic constraints and the expected values of the top quark and W boson masses [17]. The invariant mass $m_{t\bar{t}}^{\text{reco}}$ of the candidate $t\bar{t}$ pair is reconstructed from the four selected jets, the lepton, and the $E_{\text{T}}^{\text{miss}}$ vector. The experimental resolution for the $t\bar{t}$ invariant mass is 8% for a resonance mass of 500 GeV. Events in the e+jets and μ+jets channels are further classified into three orthogonal categories, based on whether a b-tagged jet was assigned to either the hadronically or the semileptonically decaying top quark, or to both of them. Each category defines a signal region, hence six independent signal regions are used in the statistical analysis.

The systematic uncertainties can be divided into experimental and modeling uncertainties. The impact on both the normalization and the shape of the $m_{t\bar{t}}^{\text{reco}}$ distributions is taken into account for all systematic uncertainties. The average impact of the dominant uncertainties on the event yields is summarized in Table 1.

Table 1: Average impact of the dominant uncertainties on the estimated yields for the total background and for a pseudoscalar A with m_A = 500 GeV and $\tan\beta$ = 0.68 in percent of the nominal value. The spectra of the e+jets and μ+jets channels are added. Only uncertainties with a yield impact > 0.5% are shown. A bar (–) indicates that an uncertainty is not applicable to a sample.

Systematic uncertainties [%]	Total bkg	S	$S+I$
Luminosity [55]	1.7	1.9	1.9
PDF	2.5	2.1	12
$t\bar{t}$ initial-/final-state radiation	3.2	–	–
$t\bar{t}$ parton shower + fragmentation	4.9	–	–
$t\bar{t}$ normalization	5.7	–	–
$t\bar{t}$ event generator	0.5	–	–
Top quark mass	0.5	2.2	13
Jet energy scale	6.4	4.9	9.3
Jet energy resolution	1.3	1.6	1.7
b-tagging: b-jet efficiency	1.5	1.3	1.1
b-tagging: c-jet efficiency	0.2	0.2	0.8
Electron efficiency	0.3	0.4	0.7
Muon efficiency	0.9	1.0	1.0
Signal MC scales	–	7.3	7.3
Reweighting	–	–	5.0
MC statistical uncertainty	0.5	2.4	11
Total uncertainty	11	10	25

Among the experimental uncertainties, those with the largest impact on the event yields and the shape of the $m_{t\bar{t}}^{\text{reco}}$ distributions are the uncertainties related to the jet energy scale (JES), and the jet energy resolution (JER) [56, 57]. The JER and JES uncertainties are also propagated to the $E_{\text{T}}^{\text{miss}}$ uncertainty. The uncertainties in the b-tagging efficiency and in the misidentification rates for jets from c-quarks, light-flavor quarks, and gluons depend on the jet transverse momentum [54]. The uncertainties related to leptons include those in the reconstruction and isolation efficiency, the single-lepton triggers, and the energy scale and resolution [48, 49].

The uncertainty of 6.5% in the NNLO+NNLL cross-section for SM $t\bar{t}$ production, is the dominant uncertainty in the normalization of the total background estimate. It includes the uncertainties due to the choice of renormalization and factorization scales and PDF set, and uncertainties in α_S [17]. A number of modeling uncertainties affecting the shape of the $m_{t\bar{t}}^{\mathrm{reco}}$ distribution for the SM $t\bar{t}$ background are also taken into account. These include uncertainties related to the choice of NLO event generator, the modeling of the parton shower and fragmentation, the modeling of gluon initial- and final-state radiation, and the value of the top quark mass m_t. A detailed description of the different sources of uncertainty related to the various background components can be found in Ref. [17].

The largest uncertainty in the modeling of the $S + I$ and S components is related to the ±1.0 GeV uncertainty of the value $m_t = 172.5$ GeV. Uncertainties related to the choice of PDF set and renormalization and factorization scales are also considered. The latter is estimated by varying the scales by factors of 0.5 and 2.0, which yields a constant ±7.3% variation across the $m_{t\bar{t}}^{\mathrm{reco}}$ spectrum. An asymmetric variation, for which the bins at the low and high ends of the $m_{t\bar{t}}^{\mathrm{reco}}$ spectrum are taken as anticorrelated [58] is also considered to estimate the impact of the scale variations on the shape of the $m_{t\bar{t}}^{\mathrm{reco}}$ spectrum. For the $S + I$ samples, an additional constant ±5% uncertainty is included to cover the difference between reweighted and generated distributions.

A detailed breakdown of the observed and expected event yields in the e+jets and μ+jets channels and their associated total uncertainties is shown in Table 2. Good agreement is found between the observed number of events in data and the expected total number of background events.

Table 2: Number of events observed in data and expected number of background events after the event selection, before the profile-likelihood fit to the full dataset. The uncertainty in the expected background yields is derived by summing all systematic and statistical uncertainties in quadrature. The "other bkg" component comprises single top quark, $t\bar{t} + W/Z$, Z+jets, diboson and multijet production. Details of the estimation of these backgrounds are given in Ref. [17].

Type	e+jets	μ+jets
$t\bar{t}$	95 000 ± 11 000	93 000 ± 11 000
W+jets	6600 ± 2100	7200 ± 2300
Other bkg.	11 200 ± 1400	6100 ± 600
Total	112 800 ± 13 000	106 300 ± 12 000
Data	115 785	110 218

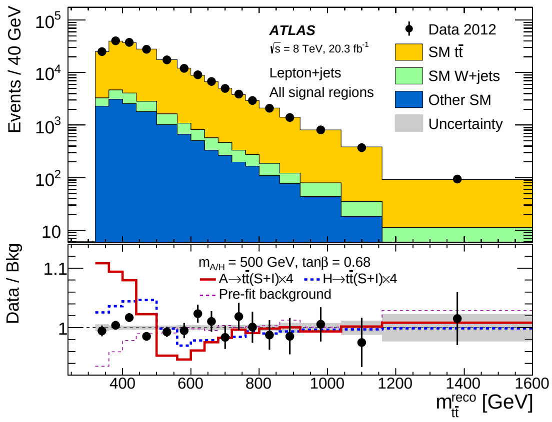

Figure 2: Distribution of $m_{t\bar{t}}^{\mathrm{reco}}$ for the data and the expected background after the profile-likelihood fit under the background-only hypothesis for all signal regions combined. The lines in the bottom panel show the individual $S + I$ distributions (scaled by a factor of four) for a pseudoscalar A (solid) and scalar H (bold dashed) with $m_{A/H} =$ 500 GeV and $\tan\beta = 0.68$ relative to the total background. The total background before the fit is indicated by a dashed line. The last bin includes overflow events.

The exclusion limits are derived separately for each signal hypothesis from a profile-likelihood fit [59] of the expected $m_{t\bar{t}}^{\mathrm{reco}}$ distributions to the observed ones simultaneously in all signal regions, taking both the statistical and systematic uncertainties into account as nuisance parameters [17]. Only bins with $m_{t\bar{t}}^{\mathrm{reco}} > 320$ GeV are considered to avoid threshold effects that are not well described by the simulation. The shape of the binned $m_{t\bar{t}}^{\mathrm{reco}}$ distributions in the presence of interference is parameterized in terms of the signal strength μ as follows [21, 22]:

$$\mu \cdot S + \sqrt{\mu} \cdot I + B = (\mu - \sqrt{\mu}) \cdot S + \sqrt{\mu} \cdot (S + I) + B. \tag{1}$$

The fitted variable is $\sqrt{\mu}$ and the case $\mu = 1$ corresponds to the type-II 2HDM in the alignment limit, while the case $\mu = 0$ corresponds to the background-only hypothesis. This approach relies on the assumption that, for a given signal hypothesis, the shape of the $t\bar{t}$ invariant mass distributions for S and $S + I$ in Eq. (1) does not change if the signal strength is varied. The terms S and $S + I$ on the right-hand side of Eq. (1) correspond to the $m_{t\bar{t}}^{\mathrm{reco}}$ distributions obtained from the S and $S + I$ samples, respectively, while B stands for the expected $m_{t\bar{t}}^{\mathrm{reco}}$ distribution of the total SM background.

The level of agreement between the observed and expected mass spectra is quantified in a fit under the background-only hypothesis with $\mu = 0$ in which only the nuisance parameters are allowed to vary. The observed $m_{t\bar{t}}^{\mathrm{reco}}$ spectra are compatible with the expected spectra after the background-only fit within the (constrained) uncertainty bands, as shown in Figure 2.

The upper limits on μ at 95% confidence level (CL) are obtained with the CL$_s$ method [60] for a number of $(m_{A/H}, \tan\beta)$ values. The upper limits at intermediate points are obtained from a linear interpolation among the three closest points. In Figure 3, the observed and expected exclusion regions for the type-II 2HDM ($\mu = 1$) are shown for the three mass hierarchies discussed in the introduction. The excluded values of $\tan\beta$ for the different mass hypotheses are listed in Table 3.

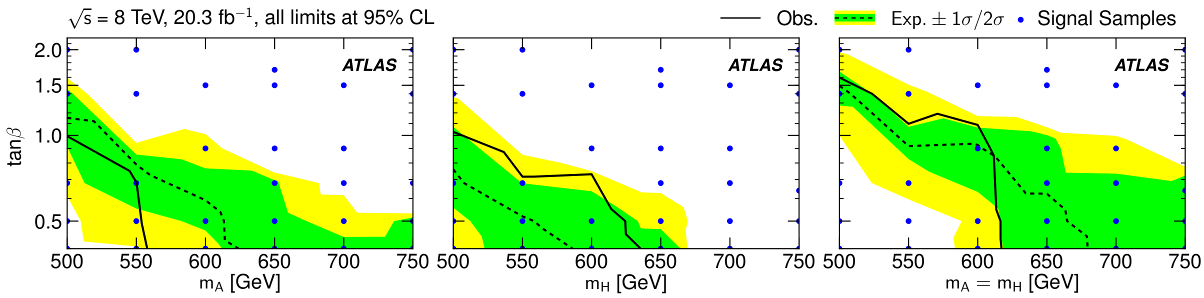

Figure 3: The 95% CL observed and expected exclusion regions for the type-II 2HDM ($\mu = 1$) considering only a pseudoscalar A (left), only a scalar H (middle), and the mass-degenerate scenario $m_A = m_H$ (right). The areas below the black solid (dashed) lines are observed (expected) to be excluded. Blue points indicate parameter values at which signal samples are produced. Upper limits at intermediate points are obtained from a linear triangular interpolation.

Table 3: The 95% CL observed and expected exclusion limits on $\tan\beta$ for a type-II 2HDM in the alignment limit considering only a pseudoscalar A (left), only a scalar H (middle), and the mass-degenerate scenario $m_A = m_H$ (right). A bar (–) indicates that no value of $\tan\beta \geq 0.4$ is excluded.

Mass		m_A		m_H		$m_A = m_H$	
[GeV]	$\tan\beta$:	obs.	exp.	obs.	exp.	obs.	exp.
500		< 1.00	< 1.16	< 1.00	< 0.77	< 1.55	< 1.50
550		< 0.69	< 0.79	< 0.72	< 0.52	< 1.10	< 0.92
600		–	< 0.59	< 0.73	–	< 1.09	< 0.93
650		–	–	–	–	–	< 0.62

In conclusion, the search for massive pseudoscalar and scalar resonances decaying into a top quark pair in 20.3 fb^{-1} of pp collisions at 8 TeV recorded by the ATLAS experiment at the LHC, yields no statistically signifiant deviations from the SM prediction. The results are interpreted in the type-II 2HDM model in the alignment limit, and upper limits are set on the signal strength μ at 95% CL in the $\tan\beta$ versus resonance mass plane. Unlike previous searches for massive resonances in the same final state, this analysis takes into account interference effects between the signal process and the background from SM $t\bar{t}$ production. It tightens significantly the previously published constraints on the 2HDM parameter space in the low $\tan\beta$ and high mass ($m_{A/H} > 500$ GeV) region.

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