ALICE results on charmonia and bottomonia

Roberta Arnaldi INFN Torino for the ALICE Collaboration





Quarkonium in AA collisions

the original idea:

quarkonium production suppressed via color screening in the QGP

T.Matsui,H.Satz, PLB178 (1986) 416



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2

Quarkonium in AA collisions

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3

Quarkonium in AA collisions

the original idea:

quarkonium production suppressed via color screening in the QGP (T.Matsui, H.Satz, PLB178 (1986) 416)



Statistical regeneration

cc multiplicity increases with collision energy → enhanced quarkonium production via (re)combination at hadronization or during QGP

P. Braun-Muzinger, J. Stachel, PLB 490(2000)196, R. Thews et al, Phys. Rev. C63:054905(2001)

Sequential melting

differences in the quarkonium binding energies lead to a sequential melting with increasing temperature

Digal, Petrecki, Satz PRD 64(2001) 0940150

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Cold nuclear matter effects

On top of the hot matter mechanisms, other effects, related to cold nuclear matter (CNM), might affect quarkonium production

- nuclear parton shadowing/color glass condensate
- energy loss
- *cc* break-up in nuclear matter

CNM are investigated in pA collisions, addressing:



Role of the various contributions, whose importance depends on kinematic and energy of the collisions

Size of CNM effects, fundamental to interpret quarkonium AA results



Quarkonium measurements: ALICE



Central Barrel $J/\psi \rightarrow e^+e^ |\gamma_{LAB}| < 0.9$ Electrons tracked using ITS and TPC Particle id: ITS, TPC, TOF, TRD

Forward muon arm $J/\psi \rightarrow \mu^+\mu^-$ 2.5< y_{LAB} <4 Muons identified and tracked in the muon spectrometer

acceptance coverage in both yregions down to zero p_T

ALICE measures inclusive J/ψ at mid and forward-y and prompt J/ψ at mid-y

Quarkonium at mid-rapidity

$J/\psi \rightarrow e^+e^-$

- Minimum bias trigger
- Good mass resolution, but low significance especially in Pb-Pb
 → only J/ψ analysis so far

Signal extraction:

- Combinatorial background subtracted via event mixing
- Signal obtained by counting technique



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Quarkonium at forward-rapidity

quarkonium $\rightarrow \mu^+\mu^-$

Dimuon trigger

Good S/B for J/ψ and Y(1S)
 → study of excited resonances still limited in Pb-Pb

Signal extraction:

- Yields extracted from a fit with signal + background shapes
- In Pb-Pb, background subtracted also via mixed-events



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	System	$\sqrt{s_{_{ m NN}}}$ (TeV)	L (MB)	L (dimuon)		
Run 1	рр	0.9, 2.76, 7, 8	1.1 nb ⁻¹ @2.76	19.9 nb ⁻¹ @2.76		
2009-	p-Pb	5.02	51 μb ⁻¹	5-5.8 nb ⁻¹		
	Pb-Pb	2.76	26 μb ⁻¹	69 μb ⁻¹		
Run 2 2015- 2018	рр	5.02, 13	2nb ⁻¹ @5.02	106nb ⁻¹ @5.02		
	p-Pb	5.02, 8.16	0.4 nb ⁻¹ @5.02	8.7-12.9 nb ⁻¹ @8.16		
	Pb-Pb	5.02	19 μb ⁻¹	225 μb ⁻¹		
	Xe-Xe	5.44	-	-		

9

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	Xe-Xe	5.44	-	-	



ALICE talks:

- Hugo Pereira da Costa → Nov 8th Cristiane Janke
 - \rightarrow Nov 9th



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acuum reference for AA - pA, genuine pp physics program

hot matter effects

matter effects

Focus on pA and AA Run 2 results

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Observables

Nuclear modification factor R_{AA}

Medium effects quantified comparing AA particle yield with pp cross section, scaled by a geometrical factor ($\propto N_{coll}$)

$$R_{AA} = \frac{Y_{AA}}{\langle T_{AA} \rangle \sigma_{\rm pp}}$$

- no medium effects $\rightarrow R_{AA} = 1$
- hot/cold matter effects $\rightarrow R_{AA} \neq 1$

Azimuthal anisotropy v_2

Multiple interactions in the medium convert initial geometric anisotropy into particle momenta anisotropy

→ elliptic flow (v₂) is the 2nd coeff. of the Fourier expansion of the azimuthal distributions of the produced particles

$$v_2 = \langle \cos 2(\phi_{\text{particle}} - \Psi_{\text{EP}}) \rangle$$



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quarkonium in AA



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$J/\psi R_{AA}$ at forward-y- Run 1



 Stronger centrality suppression at RHIC, in spite of LHC larger energy densities

• Very different p_{T} dependence

suppression + regeneration mechanisms



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Low $p_T J/\psi$

$J/\psi R_{AA}$ at forward-y- Run 2



 J/ψ suppression in Run2 confirms Run1 observation, with an increased precision



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 $J/\psi R_{AA}$ at mid-y- Run 2



No significant \sqrt{s} -dependence at mid-rapidity, confirming observation at forward-yIncrease at low p_T compared to forward-y



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Comparison with theoretical models



Transport models:

based on thermal rate eq. with continuous J/ψ dissociation and regeneration in QGP and hadronic phase

X. Zhao, R. Rapp NPA 859 (2011) 114, K. Zhou et al, PRC 89 (2011) 05491

Comover model:

J/ψ dissociated via interactions with partons - hadrons + regeneration contribution E. Ferreiro, PLB749 (2015) 98, PLB731 (2014) 57

All models fairly describe the data, as already in Run1

but large uncertainties associated to charm cross section and shadowing (data precision better than the theory one)

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Multi-differential R_{AA} at forward-y





Zhao et al., NPA 859 (2011) 114

 R_{AA} vs p_{T} for different centrality bins (and vice-versa) at $\sqrt{s_{NN}}$ =5.02 TeV

Striking features observed

- \rightarrow no R_{AA} centrality dependence in 0.3< p_T <2 GeV/c
- \rightarrow ~80% suppression for central events at $p_T \sim 10$ GeV/c

Increase in results precision opens up the way to precise model comparisons

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ψ**(2S)** *R*_{AA}

ψ (2S) shows a stronger suppression than the J/ ψ , in semi-central and central collisions

However, the low significance limits the precision of the measurements



Results at $\sqrt{s_{NN}} = 5.02$ TeV compatible with those at $\sqrt{s_{NN}} = 2.76$ TeV

Good agreement also with CMS results at $\sqrt{s_{NN}}$ = 5.02 TeV

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J/ψ elliptic flow



 J/ψ from recombination should inherit the charm flow, leading to a ν_2 signal

ALICE Run 1 result \rightarrow indication of non-zero flow (2.7 σ) **Higher Run2 precision**

 \rightarrow evidence for non-zero flow $(7\sigma \text{ effect in } 4 < p_T < 6 \text{ GeV}/c)$

ALI-DER-139475

ALICE, arXiv:1709.05260

J/ψ elliptic flow: mid and forward-y



 J/ψ from recombination should inherit the charm flow, leading to a v_2 signal

ALICE Run 1 result \rightarrow indication of non-zero flow (2.7 σ)

Higher Run2 precision
→ evidence for non-zero flow
(7σ effect in 4<p_T<6 GeV/c)

First J/ ψ v_2 measurement at mid-y \rightarrow agreement with forward-y result

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ALICE, arXiv:1709.05260

J/ψ elliptic flow: theory models



ALICE, arXiv:1709.05260 Zhou et al., PRC89(2014) 054911 Du et al., NPA943 (2015) 147 J/ψ from recombination should inherit the charm flow, leading to a v_2 signal

- ALICE Run 1 result \rightarrow indication of non-zero flow (2.7 σ)
- Higher Run2 precision \rightarrow evidence for non-zero flow (7 σ effect in 4< p_T <6 GeV/c)
- First J/ ψ v_2 measurement at mid-y \rightarrow agreement with forward-y result

Comparison with models:

- → low p_T : v_2 reproduced including a strong J/ ψ regeneration component
- → high p_T : v_2 underestimated (prompt J/v from CMS also show $v_2 \neq 0$)

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J/ψ elliptic flow: comparison with open charm



ALI-DER-138768

ALICE, arXiv:1707.01005

Similar v_2 observed for open charm

- \rightarrow different kinematic range:
- J/ψ: 2.5 < y< 4, centrality= 20-40% D: |y|< 0.8, centrality= 30-50%
- \rightarrow Low $p_{\rm T} v_2$ larger for D
- Charm quarks strongly interact in the medium
 Comparison between J/ψ and D flow can give insights on flow properties of heavy vs light quarks



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Bottomonia in ALICE



Strong Υ (1S) suppression vs centrality, similar, within uncertainties, to the $\sqrt{s_{NN}}$ = 2.76TeV one

bottomonium states accessible with higher precision in Run 2

Hint for stronger $\Upsilon(2S)$ suppression vs $\Upsilon(1S)$, as observed by CMS

 $R_{AA} (\Upsilon(2S)) = 0.26 \pm 0.12 \pm 0.06 (sys.) <$ $< R_{AA} (\Upsilon(1S)) = 0.40 \pm 0.03 \pm 0.04 (sys.)$

suppression of directly produced $\Upsilon(1S)$? \rightarrow feed-down contribution~30%

Y(1S) in ALICE: theory comparison



Transport and anisotropic hydrodynamical models qualitatively describe the centrality and the p_T evolution

Some tension in the *y* dependence?

No need for significant contribution of regenerated $\Upsilon(1S)$

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quarkonium in pA



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J/ψ in p-Pb collisions

pA collisions are a tool to:

Disentangle among CNM effects Investigate role of CNM effects underlying AA collisions Search for possible hot matter effects?

Two beam configurations: p-Pb and Pb-p $H_{\rm pPb}$ ALICE, inclusive $J/\psi \rightarrow \mu^+\mu^-$ 1.2 0.8 2.03<*y*_{CMS}<3.53 0.6 -4.46<*y*_{CMS}<-2.96 0.4 Clear J/ ψ suppression at forward-y, while R_{pA} p-Pb \sqrt{s_{NN}} = 5.02 TeV (JHEP 02 (2014) 073) 0.2 is compatible with unity at backward-y p-Pb \screwspiperset states and states and states p-Pb \screwspiperset states and states Compatible R_{pPb} at $\sqrt{s_{NN}}$ = 5.02 and 8.16 TeV У_{стs} ALI-PREL-118140 (slightly different $x_{\rm F}$ range) CERN-ALICE-PUBLIC-2017-001

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J/ψ in p-Pb collisions: Run1 vs Run2



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Comparison with theory models

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Good agreement between data and models based on shadowing and/or energy loss, as at $\sqrt{s_{NN}} = 5.02 \text{ TeV}$



Size of theory uncertainties (mainly shadowing) still limits a more quantitative comparison

Ducloue et al, PRD91(2015)114005, Lansberg et al, EPJC77(2017)1, Ma et al, PRD92(2015)071901, Chen et al, PLB765(2017)323, Arleo, Vogt arXiv:1707.09973

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$J/\psi v_2$ in p-Pb

Azimuthal correlations between forward / backward J/ ψ and mid-y charged particles



*p*_T<3 GeV/*c* → *v*₂ compatible with 0 (in line with expectation of no recombination)
 3<*p*_T<6 GeV/*c* → *v*₂>0



- → $v_2 > 0$ → suggests J/ ψ participation to the collective flow of the medium
- → ~5σ total significance (forward + backward, 5.02+8.16 TeV)
- values comparable to J/ψ ν₂ in central Pb-Pb collisions
- → common mechanism at the origin of the J/ ψ v_2 ?



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ψ (2S) in pA collisions

Strong $\psi(2S)$ suppression of at both forward and backward γ



Effect similar to the one observed at $\sqrt{s_{NN}} = 5.02 \text{ TeV}$

No sizeable $\sqrt{s_{NN}}$ dependence, both in y and p_T

J/ψ and ψ (2S) comparison in pA

Strong $\psi(2S)$ suppression of at both forward and backward γ



$\Rightarrow \psi(2S)$ suppression is stronger than the J/ ψ one, in particular at backward-y

unexpected J/ ψ and ψ (2S) different behavior since at LHC energies formation time > crossing time

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J/ψ and ψ (2S) comparison with theory



shadowing/energy loss:

- similar for J/ ψ and ψ (2S)
- not enough to describe the ψ(2S) suppression at backward-y



need final state effects

- soft color exchanges between hadronizing cc and comoving partons (Ma and Venugopalan)
- "classical" comover model, with break-up σ tuned on low energy data (Ferreiro)

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Υ in pA collisions – Run 1



Model predictions describe the measured R_{pPb} at forward y and tend to underestimate the suppression at backward y

Compatible within (large) uncertainties with LHCb results

Run 2 data will be soon available!



Conclusions

New high-precision results on flavor production in pA and AA collisions

- J/ ψ described by interplay of suppression and recombination mechanisms
- *****

AA

- Significant J/ ψ v_2 at intermediate p_T confirms formation by recombination
- Strong $\psi(2S)$ suppression
- Hint for sequential suppression of bottomonium states



- Modification of J/ψ yields, with strong kinematic dependence, understood in terms of "standard" cold nuclear matter effects
- Size of J/ ψ v_2 at intermediate p_T , reminiscent of the Pb-Pb one. Common mechanism at play?
- Strong ψ (2S) suppression due to final state effects?

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Backup slides



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July 12th 2017

Quarkonium sequential melting



the original idea:

quarkonium production suppressed via color screening in the QGP

sequential melting:

differences in the quarkonium binding energies lead to a sequential melting with increasing temperature



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NPQCD17

May 23rd 2017

6

Quarkonium sequential melting

state	J/ψ	χ _c	ψ (2S)	Ƴ (1S)	Υ (2S)	Υ (3S)
Mass(GeV)	3.10	3.51	3.69	9.46	10.0	10.36
∆E (GeV)	0.64	0.22	0.05	1.10	0.54	0.20
r _o (fm)	0.50	0.72	0.90	0.28	0.56	0.78

(Digal, Petrecki, Satz PRD 64(2001) 0940150)

sequential melting:

differences in the quarkonium binding energies lead to a sequential melting with increasing temperature

> Quarkonium as thermometer of the initial QGP temperature



July 12th 2017

Caveat

Even if the "suppression-recombination" approach looks simple, a realistic description of the involved mechanisms is rather complex:

\rightarrow on the theory side:

- Link between suppression and critical temperature requires precise assessment of T_D , $M_{\psi}(T)$, $\Gamma_{\psi}(T)$ from QCD calculations using EFT/LQCD spectral functions
- Short QGP thermalization time at LHC might imply in-medium formation of quarkonia rather than suppression

\rightarrow on the experimental side:

- Precise determination of open charm σ
- Assessment of quarkonium feed-down into lighter states



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Comparison with theoretical models



Transport models: based on thermal rate eq. with continuous J/ ψ dissociation and regeneration in QGP and hadronic phase X. Zhao, R. Rapp NPA 859 (2011) 114, K. Zhou et al, PRC 89 (2011) 05491

Statistical hadronization: J/ψ produced at chemical freeze-out according to their statistical weight A. Andronic et al., NPA 904-905 (2013) 535

Comover model: J/ ψ dissociated via interactions with partons - hadrons + regeneration contribution E. Ferreiro, PLB749 (2015) 98, PLB731 (2014) 57

All models fairly describe the data, as already in Run1

Model	dσ _{J/ψ} /dy [mb] fw-y	shadowing	
Transport, TM1	0.57	EPS09	
Transport, TM2	0.82	EPS09	
Stat. Hadroniz.	0.32	EPS09	
Comovers	0.45-0.7	Glauber-Gribov	

but large uncertainties associated to charm cross section and shadowing

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 $J/\psi R_{AA}$ at mid-y. Run 2



No significant \sqrt{s} -dependence also at mid-rapidity, confirming observation at forward-y

Small R_{AA} increase in most central collisions, wrt forward-y, as expected in a (re)generation scenario (but fluctuations cannot be yet excluded)

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$p_{\rm T}$ dependence of $R_{\rm AA}$



Similar R_{AA} at $\sqrt{s_{NN}} = 2.76$ and 5.02 TeV, with a hint for an increase in the range $2 < p_T < 6$ GeV/c $J/\psi R_{AA}$ is higher at low p_T , where J/ψ from regeneration dominate



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More differential J/ ψ R_{AA} : p_T

Constraints to the theoretical models can be imposed by more differential R_{AA} studies



no centrality dependence in $0.3 < p_T < 2$ GeV/c in central collisions, smaller suppression for low- p_T J/ ψ , as expected by (re)generation



High- p_T J/ ψ : pattern qualitatively similar to the one measured by ATLAS and CMS, reaching R_{AA} ~0.2

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From pA to AA

Once CNM effects are measured in pA, what can we learn on J/ψ production in PbPb?

Hypothesis: $2 \rightarrow 1$ kinematics for J/ ψ production CNM effects (dominated by shadowing) factorize in p-A

CNM obtained as $R_{pA} \times R_{Ap} (R_{pA}^2)$, similar x-coverage as PbPb



Sizeable p_{T} dependent suppression still visible \rightarrow CNM effects not enough to explain AA data at high p_{T}

we get rid of CNM effects with

AA / pA

CNM effects not enough to explain PbPb data at high p_{T}



Evidence for hot matter effects in Pb-Pb!

RAA vs y

Constraints to the theoretical models can be imposed by more differential RAA studies





Hint of enhanced production towards mid-y

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J/ψ elliptic flow

J/ ψ from recombination should inherit the charm flow, leading to a v_2 signal

Effect should be important at LHC energies, in kinematic regions where regeneration plays a role



RHIC results favour $v_2 \sim 0$

 $v_2 \neq 0$ at high p_T \rightarrow possibly due to the energy loss path-length dependence ALICE observes evidence for non-zero flow at intermediate pT (7σ effe 26

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J/ψ elliptic flow: analysis technique

J/ ψ $v_2 = \langle \cos 2(\phi_{\mu\mu} - \Psi_{EP}) \rangle$ is computed using the Event Plane from

SPD ($\Delta\eta$ =1.1) at fw-y TPC ($\Delta\eta$ =0) at mid-y

 $\sim v_2^{J/\psi}$ is obtained modeling <cos 2 ($\phi_{\mu\mu}$ - Ψ_{EP})> vs inv. mass as

 $v_2(m_{\mu\mu}) = v_2^{J/\psi} \alpha(m_{\mu\mu}) + v_2^{bck} (1 - \alpha(m_{\mu\mu}))$

 $\sim v_2 = v_2^{obs} / \sigma_{EP}$





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 α (m_{µµ}) is S/S+B from inv. mass fit

 v_2^{bck} background parametrized by several functions

ψ (2S) in AA collisions

 ψ (2s) is a loosely bound state (binding energy ~60 MeV wrt to ~640 MeV for J/ ψ)

Expected to be more easily dissociated than J/ ψ \rightarrow sequential suppression scenario

► Less clear role played by recombination, taking place
→ at freeze-out, as for J/ψ in the statistical hadronization model

 → in later collision stages, when the system is more diluted (and radial flow is stronger) [sequential regeneration, Rapp, arXiv:1609.04868]

Ratio of charmonium states vs. centrality and vs. $p_{\rm T}$ can give insight on quarkonium behaviour





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ψ**(**2S) *R*_{AA}

At $\sqrt{s_{NN}}$ = 5.02 TeV, results are compatible with CMS, in a similar kinematic range, while some tension exists at lower energy



Results in different kinematic ranges are sensitive to the fraction of primordial and regenerated charmonia, to different medium temperature and flow...

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Low pT J/ ψ at fw-y

Strong R_{AA} enhancement in peripheral collisions for 0<p_T<0.3 GeV/c



if excess is "removed" requiring $p_T^{J/\psi}$ >0.3GeV/c \rightarrow ALICE R_{AA} lowers by 20% at maximum (in the most peripheral bin)

significance of the excess is 5.4 (3.4)σ in 70-90% (50-70%)

behaviour not predicted by transport models

excess might be due to coherent J/ψ photoproduction in PbPb (as measured also in UPC)



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Bottomonia in AA

Three states characterized by very different binding energies:

Y(1S): Eb~1100 MeV Y(2S): Eb~500 MeV Y(3S): Eb~200 MeV





Sensitive in very different ways to the medium

With respect to charmonium:

- Limited recombination effects
 interesting for sequential suppression studies
- More robust theoretical calculations, due to higher b quark mass
- No B hadron feed-down
 → simpler interpretation?

Some drawbacks

- Lower production cross sections
- Non negligible feed-down contributions from higher states



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Υ(2S) in ALICE



Stronger suppression has been observed for the $\Upsilon(2S)$ wrt $\Upsilon(1S)$

Theoretical models describe the R_{AA} ratio (no need for regeneration contribution)

Result is consistent with the centrality-integrated CMS measurement

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Y(1S) in ALICE: theory comparison



Some tension in the R_{AA} evolution vs y with energy, but still large uncertainties



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$p_{\rm T}$ dependence of J/ $\psi R_{\rm pA}$



Slightly different y coverage in ALICE and LHCb, but rather similar p_T dependences

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Shadowing and energy loss models describe R_{pA} vs p_T



Roberta Arnaldi Precision spectroscopy of QGP properties with jets and heavy quarks

ψ (2S) in pA collisions

Being more weakly bound than the J/ ψ , the ψ (2S) is an interesting probe to have further insight on the charmonium behaviour in pA



 ψ (2S) suppression stronger than the J/ ψ one at RHIC and LHC

→ unexpected because time spent by the cc pair in the nucleus (τ_c) is shorter than charmonium formation time (τ_f)

→ shadowing and energy loss, almost identical for J/ ψ and ψ (2S), do not account for the different suppression

QGP+hadron resonance gas or comovers models describe the stronger ψ (2S) suppression \int

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ψ (2S) in p-Pb at $\sqrt{s_{NN}} = 5.02$ TeV

Being more weakly bound than the J/ ψ , the ψ (2S) is an interesting probe to have further insight on the charmonium behaviour in pA

1.5

0.5

 ψ (2S) suppression stronger than the J/ ψ one at RHIC and LHC

- \rightarrow unexpected because time spent by the cc pair in the nucleus (τ_c) is shorter than charmonium $\mathfrak{Q}_{\mathsf{pPb}}$ formation time (τ_f)
- shadowing and energy loss, almost identical for J/ψ and $\psi(2S)$, do not account for the different suppression

√s=200 GeV **PH**^{*}ENIX p+Au preliminary N_{V(2s)}p p+Al N/L N/V ^I d+Au PRL 111 202301 (2013) $\frac{N_{\psi(2s)}}{N_{J/\psi}})^{p_+}$ ±15.6% global uncertainty on forward/backward rapidity points ±16% global uncertainty on midrapidity point rapidity o_{pPb} ALICE, p-Pb $\sqrt{s_{NN}}$ = 5.02 TeV, -4.46 < y_{cms} < -2.96 ALICE, p-Pb $\sqrt{s_{NN}}$ = 5.02 TeV, 2.03 < y_{ome}< 3.53 Inclusive J/ψ , $\psi(2S) \rightarrow \mu^{+}\mu^{-}$ J/ψ: EPS09 LO + comovers (Ferreiro) ELoss (Arleo et al.) — J/ψ: EPS09 LO + comovers (Ferreiro) $\psi(2S)$: EPS09 LO + comovers (Ferreiro) v(2S): EPS09 LO + comovers (Ferreiro) EPS09 NLO (Vogt et al. J/ψ: QGP+HRG (Du et al.) J/ψ: QGP+HRG (Du et al.) EPS09 LO (Ferreiro) 1.5 ψ(2S): QGP+HRG (Du et al.) w(2S); QGP+HRG (Du et al.) • 0.5

• w(2S

QGP+hadron resonance gas or comovers models describe the stronger $\psi(2S)$ suppression

Inclusive J/ ψ , $\psi(2S) \rightarrow \mu^+ \mu$

EPS09 NLO (Vogt et al.)

ELoss (Arleo et al.)

EPS09 LO (Ferreiro)

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12

<N

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12

<N

42

10

$J/\psi v_2$ in pPb

Azimuthal correlations between forward/backward J/ ψ and mid- γ charged particles

Correlations expressed as associated SPD-tracklet yields per dimuon(J/ψ) trigger





additional enhancement at both near and away sides

Low multiplicity



clear away-side correlation
(jets?)



Jet correlations eliminated via subtraction

J/ ψv_2 extracted assuming factorization of J/ ψ and tracklet v_2

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Υ in pA collisions

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Stronger excited states suppression with respect to $\Upsilon(1S)$ Initial state effects similar for the three Υ states \rightarrow Final states effects in p-Pb?

no strong rapidity dependence of $\Upsilon(1S) R_{pA}$

 Υ (1S) R_{pA} described by shadowing and energy loss models



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